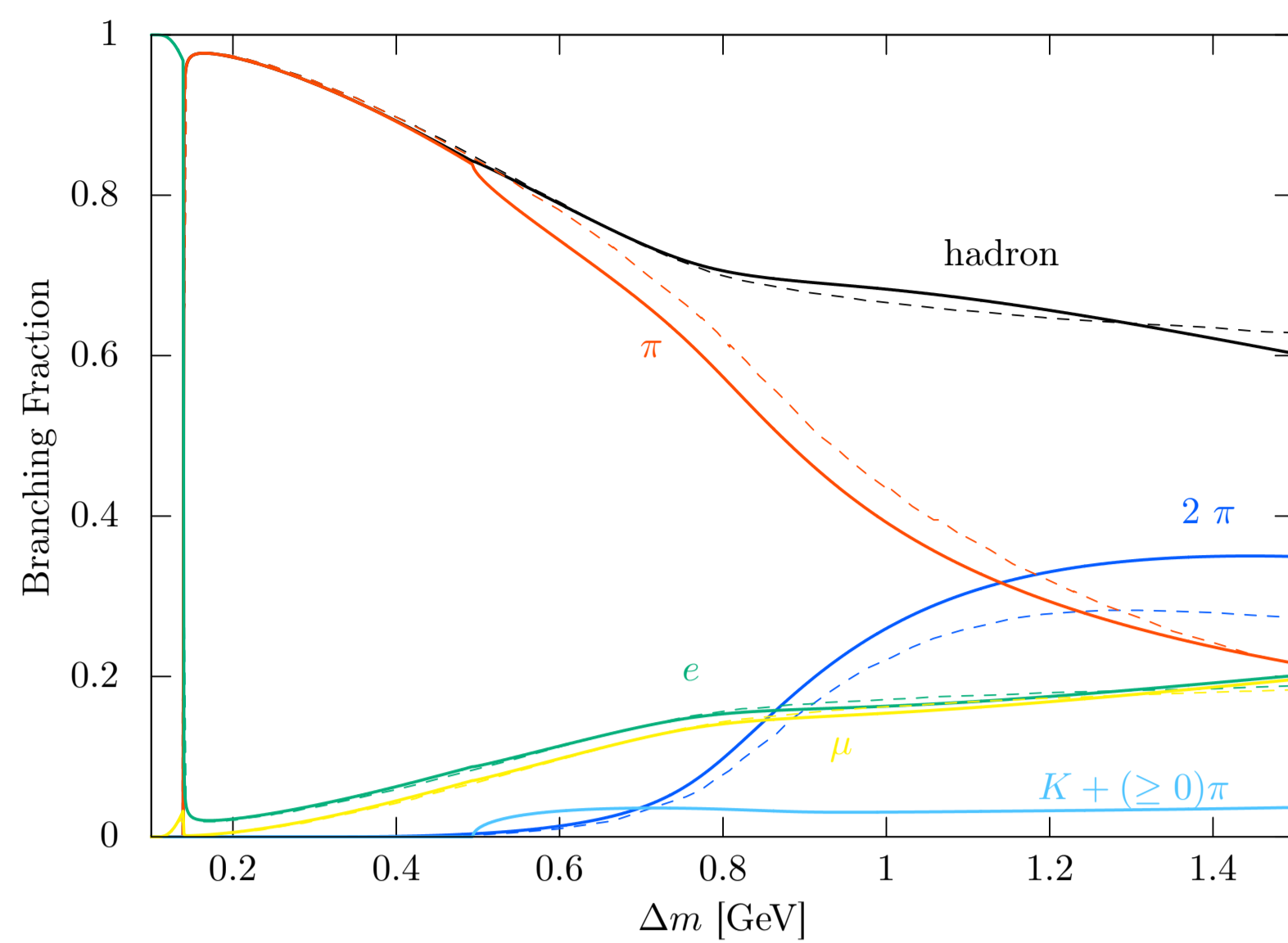
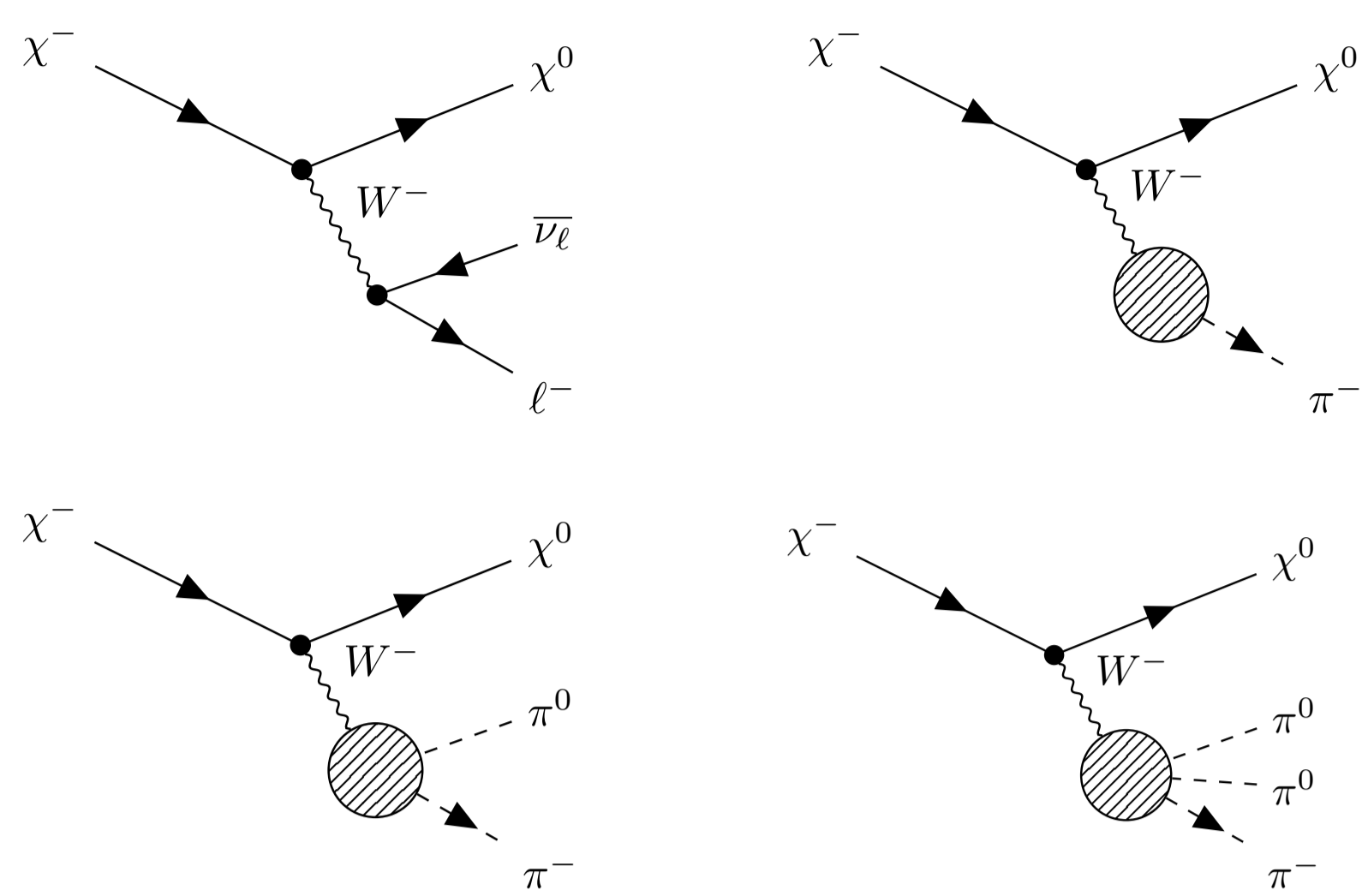


Precise Estimate of Chargino Decay

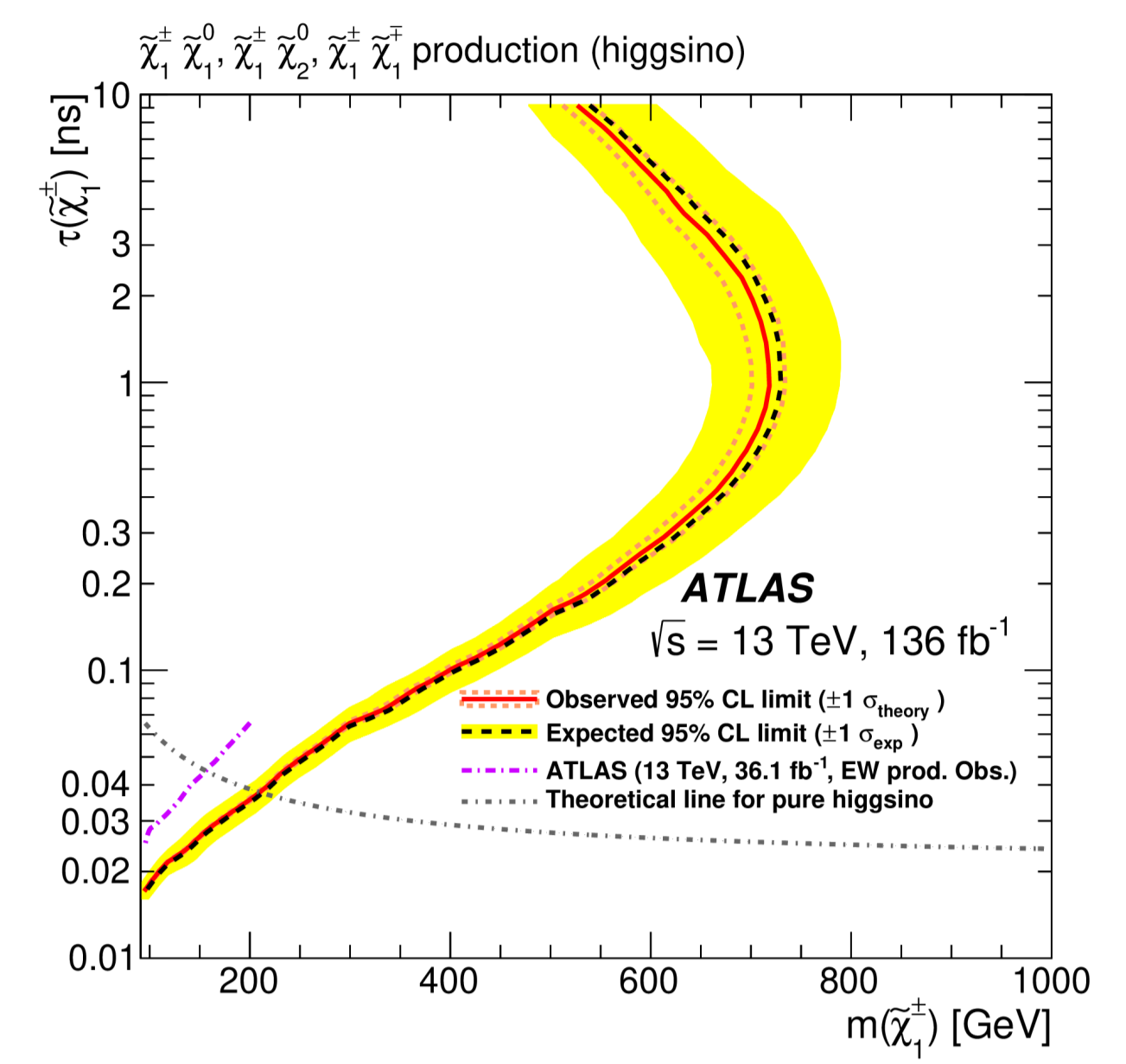
Masahiro Ibe, Masataka Mishima, Yuhei Nakayama, Satoshi Shirai, JHEP 07 (2021) 098

Chargino Decay and Collider Constraint

Chargino decays through weak interaction ...



Collider constraints are sensitive to decay rates and BFs!



Universal Short-Distance Correction

W-boson self energy

$$G_F^0 \rightarrow G_F^0 \times \left\{ 1 + \frac{\sum_{W,W'}^{1PI}(0)}{m_W^2} + \frac{\alpha}{4\pi} \left[3 \log \frac{m_W}{\mu_{IR}} + F_V^{\text{Virtual(EW)}} \left(\frac{m_W}{m_{\chi^\pm}} \right) \right] \right\} \quad G_F^0 = \frac{e^2}{4\sqrt{2}s_W^2 m_W^2}$$

Replacement of G_F^0 with $G_F \simeq 1.166 \times 10^{-5} \text{ GeV}^{-2}$ removes W-boson self-energy, because

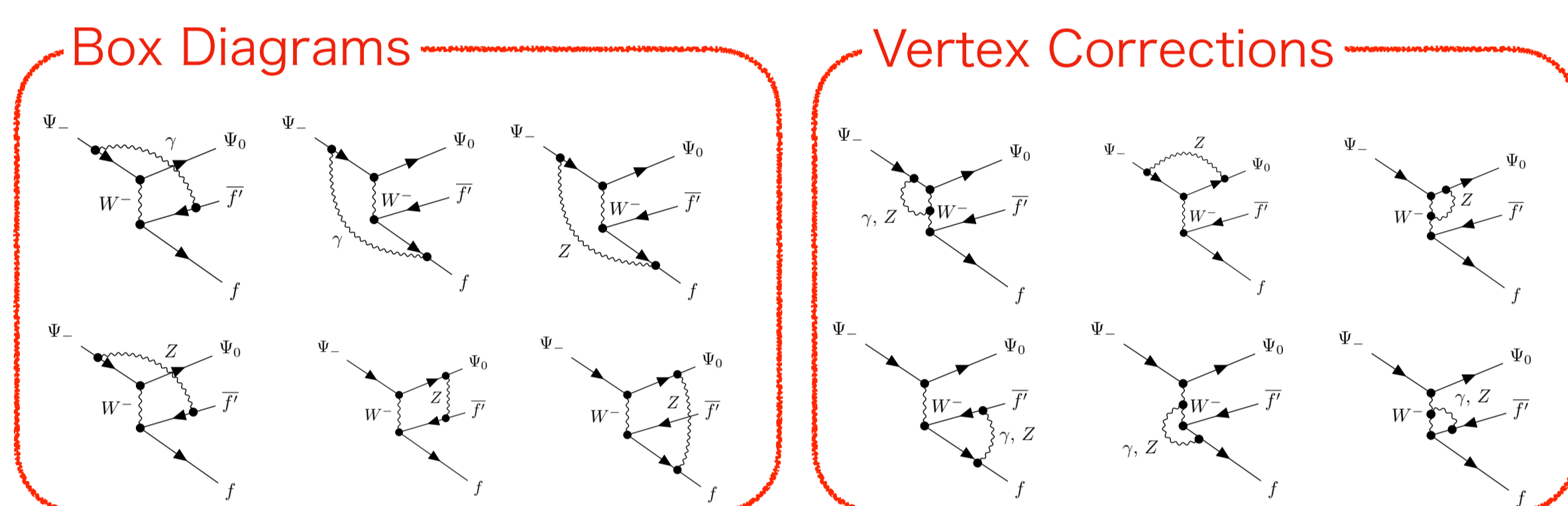
$$G_F^0 = G_F \times (1 - \Delta r), \quad \Delta r|_{1\text{-loop}} = \frac{\sum_{W,W'}^{1PI}(0)}{m_W^2} + \frac{\alpha}{4\pi} \left(\frac{6}{s_W^2} + \frac{7 - 4s_W^2}{s_W^4} \log c_W \right)$$

Encapsulation of short-distance corrections:

$$G_F^0 \rightarrow G_F \times \sqrt{S_{EW}^X(\mu_{IR})}; \quad S_{EW}^X(\mu_{IR}) = 1 + \frac{\alpha}{4\pi} \left[6 \log \frac{m_W}{\mu_{IR}} + 2F_V^{\text{Virtual(EW)}} \left(\frac{m_W}{m_{\chi^\pm}} \right) - 2 \left(\frac{6}{s_W^2} + \frac{7 - 4s_W^2}{s_W^4} \log c_W \right) \right] + O(\alpha^2).$$

On-shell matching defines the effective Fermi theory:

$$\mathcal{L}_{FF} = -2\sqrt{2}G_F (\bar{\Psi}_0 \gamma^\mu \Psi_-) (\bar{\Psi}_f \gamma_\mu P_L \Psi_{f'}) + (\text{CT})$$



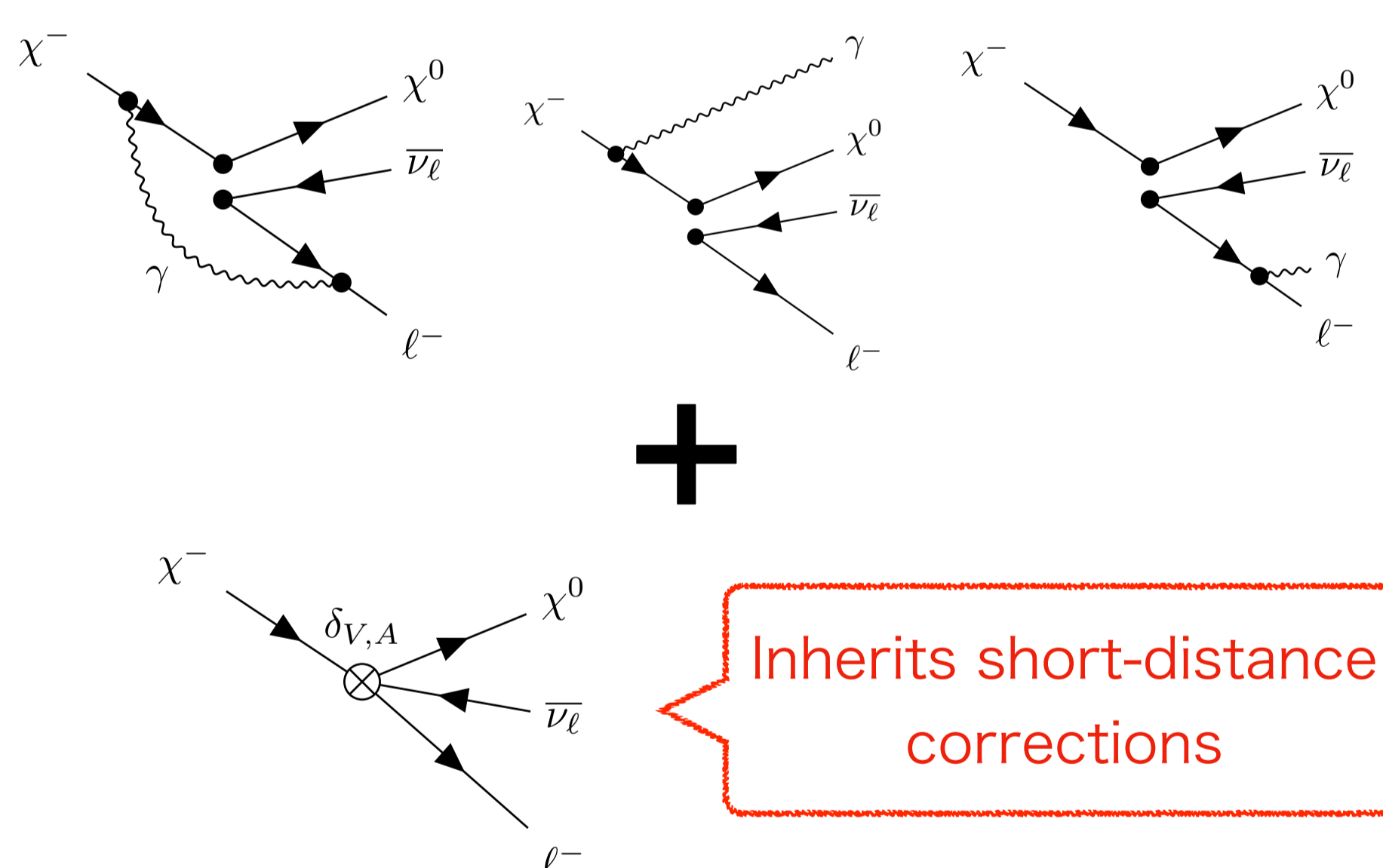
Leptonic Mode

Effective Fermi Theory

$$\mathcal{L}_{FF}^{\text{lepton}} = -2\sqrt{2}G_F (\bar{\Psi}_{\chi^0} \gamma^\mu \Psi_{\chi^-}) (\bar{\ell} \gamma_\mu P_L \nu_\ell)$$

$$\delta \mathcal{L}_{FF}^{\text{lepton}} = -2\sqrt{2} \delta_V G_F (\bar{\Psi}_{\chi^0} \gamma^\mu \Psi_{\chi^-}) (\bar{\ell} \gamma_\mu P_L \nu_\ell) - 2\sqrt{2} \delta_A G_F (\bar{\Psi}_{\chi^0} \gamma^\mu \gamma_5 \Psi_{\chi^-}) (\bar{\ell} \gamma_\mu P_L \nu_\ell)$$

Determined by EW theory



$$\Gamma_{1\text{-loop}}(\chi^- \rightarrow \chi^0 \ell^- \bar{\nu}_\ell(\gamma)) = \Gamma_{\text{tree}}^{\ell, m_\ell=0} f \left(\frac{m_\ell}{\Delta m} \right) S_{EW}^{\chi, \text{lep}} \left[1 + \frac{\alpha}{4\pi} g \left(\frac{m_\ell}{\Delta m} \right) \right]$$

Lepton mass effect @tree-level Long-distance correction @NLO

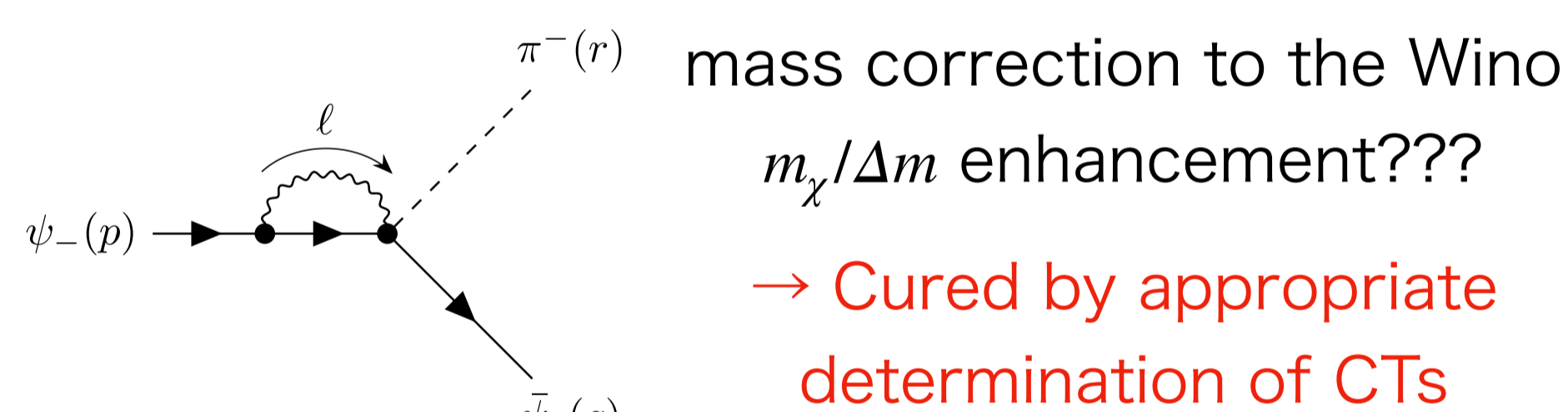
Single Pion Mode

Pion's shift symmetry ensures that the decay amplitude is suppressed by Δm

$$\mathcal{L}_{CC}^{\text{pion}} = -\sqrt{2} G_F V_{ud}^* F_\pi (\partial_\mu \pi^-)^* \times \bar{\Psi}_{\chi^0} \gamma^\mu \Psi_{\chi^-} + \text{h.c.} \quad p \sim \Delta m$$

But QED corrections break the symmetry:

$$\partial_\mu \rightarrow D_\mu = \partial_\mu - ieA_\mu$$



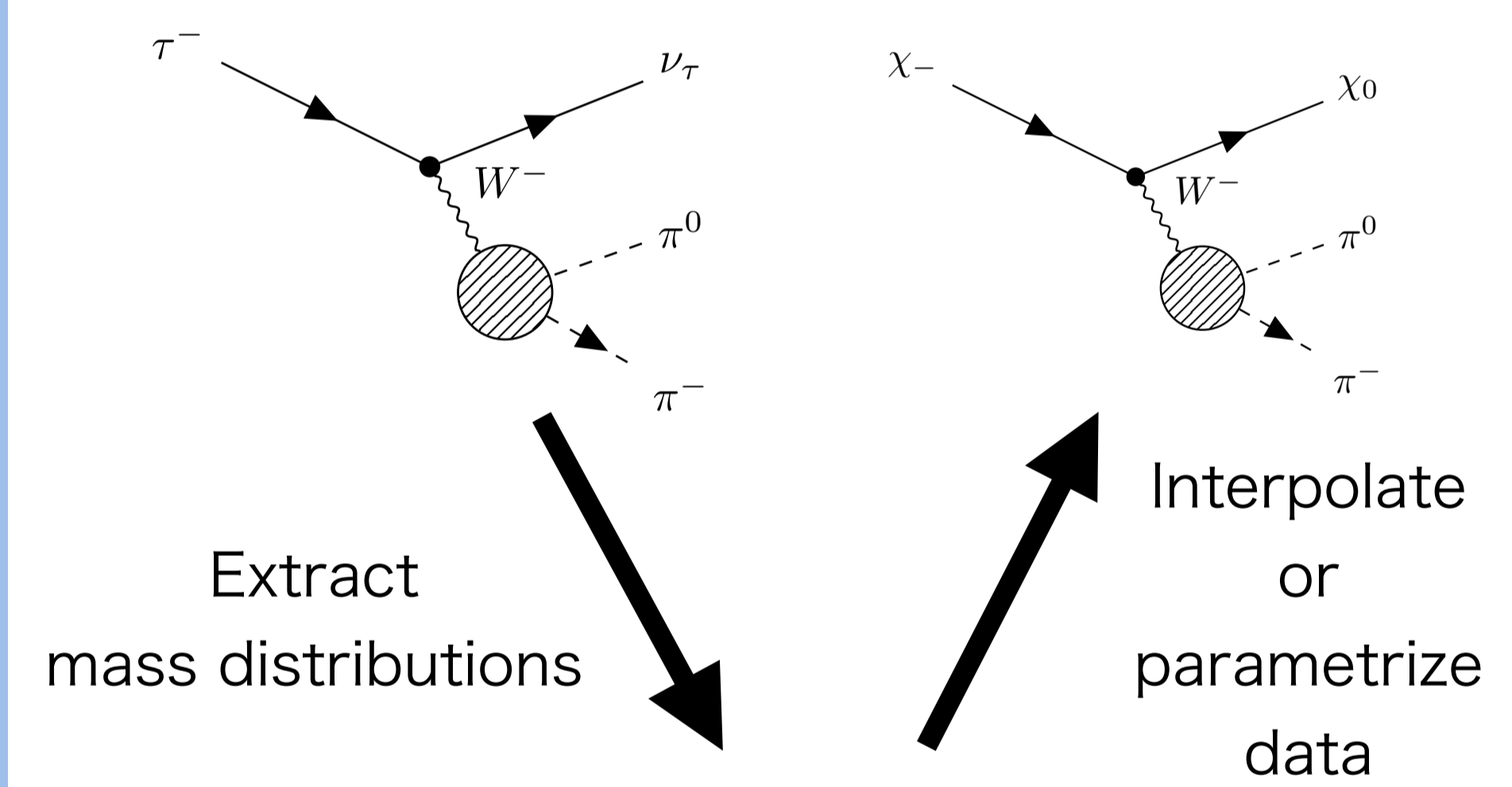
$$\mathcal{L}_Y = e^2 \left\{ \frac{1}{\sqrt{2}} F_0^2 G_F \left[Y_1 \bar{\Psi}_{\chi^-} \gamma_\mu \Psi_{\chi^0} \langle u^\mu \{ Q_R, Q_W \} \rangle + \hat{Y}_1 \bar{\Psi}_{\chi^-} \gamma_\mu \Psi_{\chi^0} \langle u^\mu \{ Q_L, Q_W \} \rangle + Y_2 \bar{\Psi}_{\chi^-} \gamma_\mu \Psi_{\chi^0} \langle u^\mu \{ Q_R, Q_W \} \rangle + Y_3 m_\chi \bar{\Psi}_{\chi^-} \Psi_{\chi^0} \langle Q_R Q_W \rangle + iY_4 \bar{\Psi}_{\chi^-} \gamma_\mu \Psi_{\chi^0} \langle Q_L^c Q_W \rangle + iY_5 \bar{\Psi}_{\chi^-} \gamma_\mu \Psi_{\chi^0} \langle Q_R^c Q_W \rangle + \text{h.c.} \right] + Y_6 \bar{\Psi}_{\chi^-} (i\cancel{\partial} - eA) \Psi_{\chi^-} + \hat{Y}_7 m_\chi \bar{\Psi}_{\chi^-} \Psi_{\chi^-} \right\}$$

We have to determine the following combination of the LECs by matching this with the Fermi theory:

$$\frac{\delta \Gamma_\pi}{\Gamma_\pi} - \frac{\delta \Gamma_\pi}{\Gamma_\pi} \supset e^2 \left[-\hat{Y}_6 - \frac{4}{3} (Y_1 + \hat{Y}_1) - 4 \left(Y_2 + \hat{Y}_2 - \frac{m_{\chi^\pm}}{\Delta m_\pm} Y_3 \right) \right]$$

Multi-Meson Mode

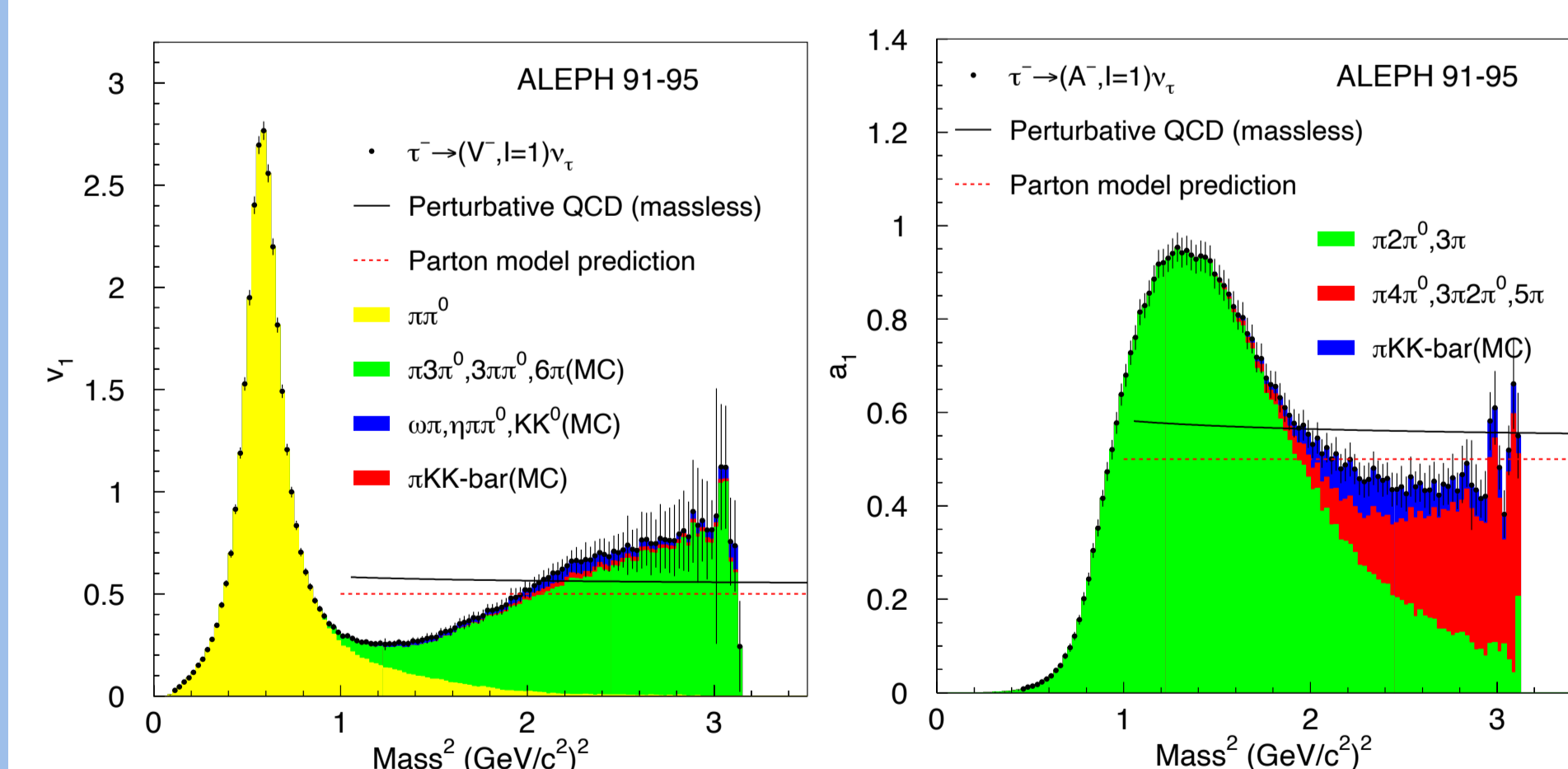
We can obtain info of hadronization from tau lepton decay [Chen et al., hep-ph/9607421]



$$\langle \pi^- \pi^0 | \bar{d} \gamma^\mu u | 0 \rangle$$

Nonperturbative QCD encoded

Mass distributions [ALEPH Collaboration]



Numerical Results (Preliminary)

