

Collapse models make particles jiggle... and emit photons

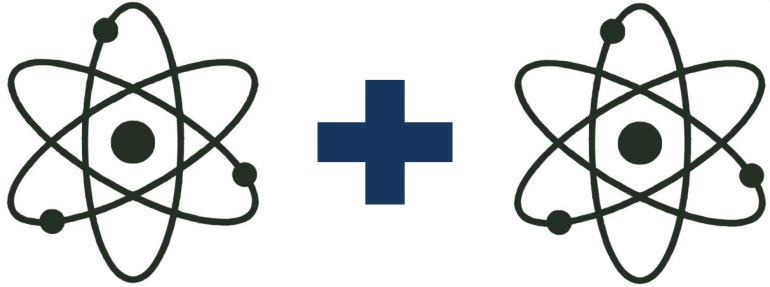
The Hitchhiker's Advanced Guide to Quantum Collapse Models and their impact in science, philosophy, technology and biology

31st October – 4th November 2022

Angelo Bassi

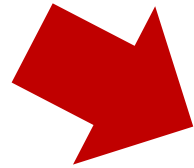
University of Trieste & INFN - Italy

Quantum superpositions

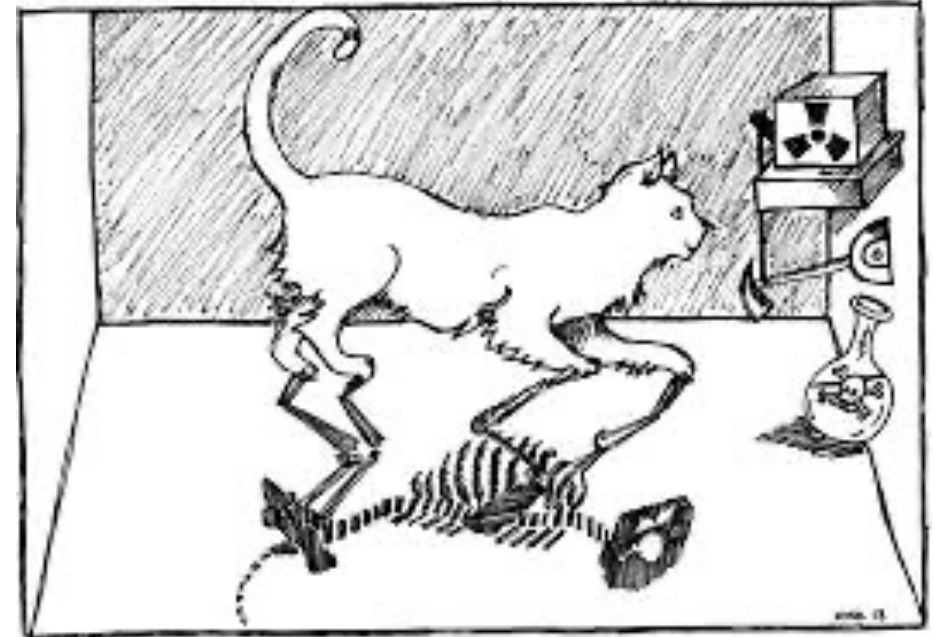


Microscopic superpositions
Experimentally verified

Cats are made of atoms + linearity of the theory

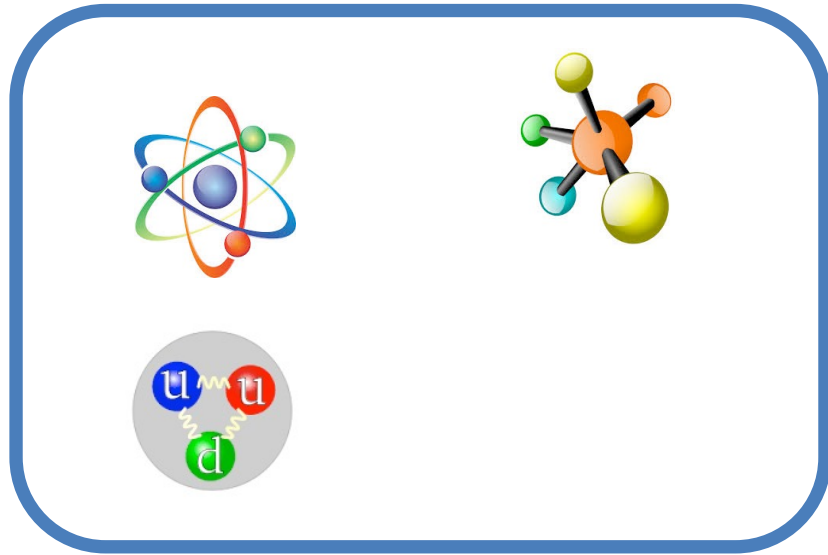


Macroscopic
superpositions
Never seen

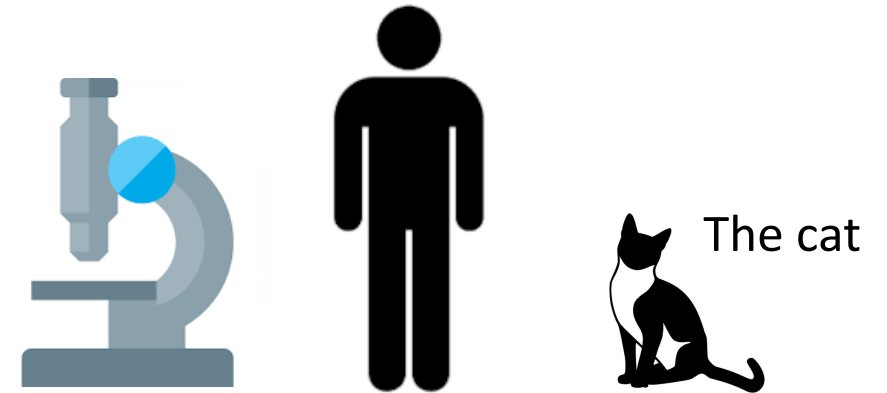


Standard Quantum Mechanics

Quantum world



Classical world

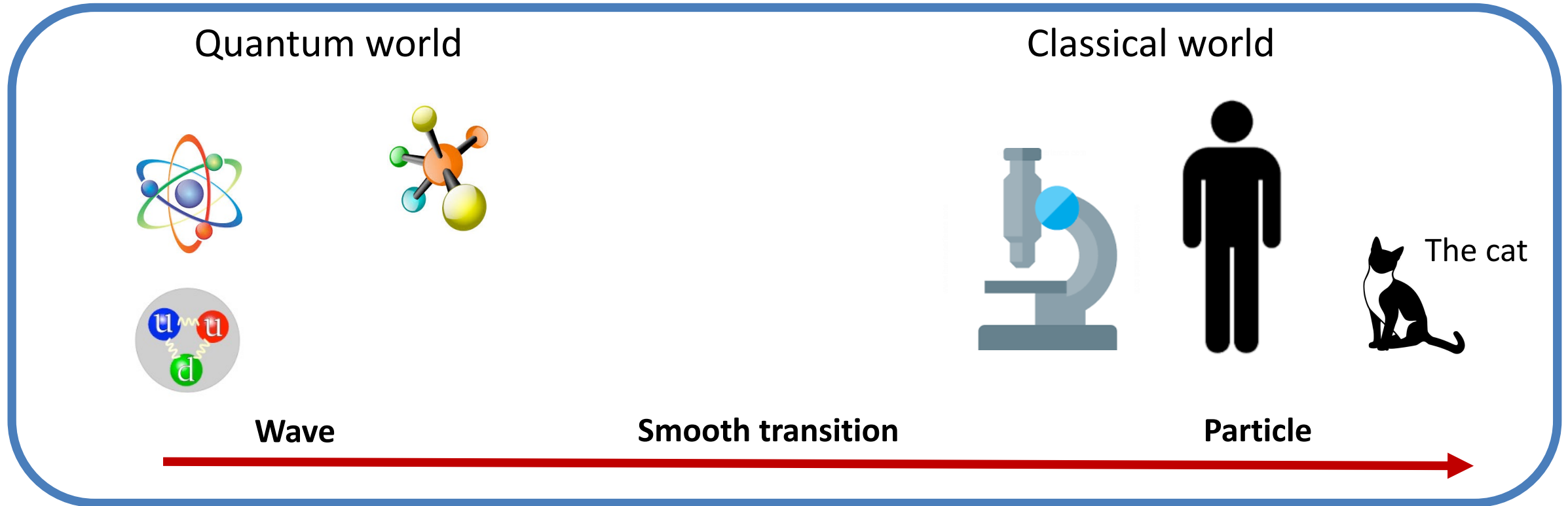


Quantum - Classical
divide

The wave function gives the probabilities
of outcomes of measurements

The Copenhagen interpretation assumes a **mysterious division** between the microscopic world governed by quantum mechanics and a macroscopic world of apparatus and observers that obeys classical physics [...]

Spontaneous wave function collapse



The Schrödinger equation is **modified**. The new dynamics is **nonlinear** in such a way to describe the quantum micro-world, the classical macro-world, as well as the transition from one to the other.

The dynamics of collapse models

A. Bassi and G.C. Ghirardi, *Phys. Rept.* 379, 257 (2003), A. Bassi, K. Lochan, S. Satin, T.P. Singh and H. Ulbricht, *Rev. Mod. Phys.* 85, 471 (2013)

$$d|\psi_t\rangle = \left[-\frac{i}{\hbar} \hat{H} dt + \int d^3\mathbf{x} \left(\hat{M}(\mathbf{x}) - \langle \hat{M}(\mathbf{x}) \rangle_t \right) dW_t(\mathbf{x}) - \frac{1}{2} \iint d^3\mathbf{x} d^3\mathbf{y} \mathcal{G}(\mathbf{x} - \mathbf{y}) \left(\hat{M}(\mathbf{x}) - \langle \hat{M}(\mathbf{x}) \rangle_t \right) \left(\hat{M}(\mathbf{y}) - \langle \hat{M}(\mathbf{y}) \rangle_t \right) dt \right] |\psi_t\rangle$$

Quantum mechanics + collapse in space

Nonlinear

Stochastic

$$M(\mathbf{x}) = ma^\dagger(\mathbf{x})a(\mathbf{x}) \quad \langle M(\mathbf{x}) \rangle_t = \langle \psi_t | M(\mathbf{x}) | \psi_t \rangle$$

Collapse operator \sim position

$$\mathbb{E}[dW_t(\mathbf{x})] = 0 \quad \mathbb{E}[dW_t(\mathbf{x})dW_t(\mathbf{y})] = \mathcal{G}(\mathbf{x} - \mathbf{y})dt$$

Noise driving the collapse

$$\mathcal{G}(\mathbf{x}) = \frac{\lambda}{m_0^2} e^{-\mathbf{x}^2/4r_C^2}$$

$$\mathcal{G}(\mathbf{x}) = \frac{G}{\hbar} \frac{1}{|\mathbf{x}|}$$

CSL model

P. Pearle, *Phys. Rev. A* 39, 2277 (1989).

G.C. Ghirardi et al., *Phys. Rev. A* 42, 78 (1990)

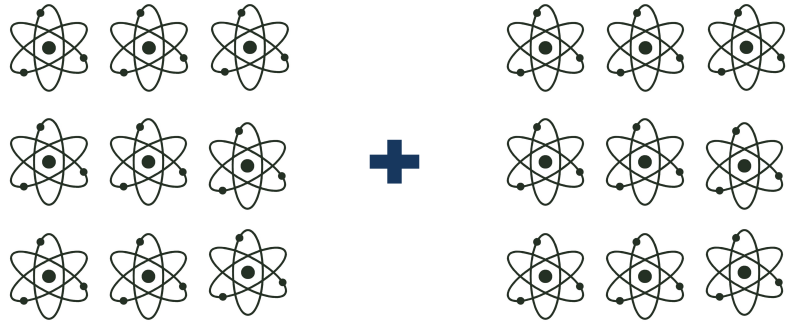
DP model

L. Diosi, *Phys. Rev. A* 40, 1165 (1989)

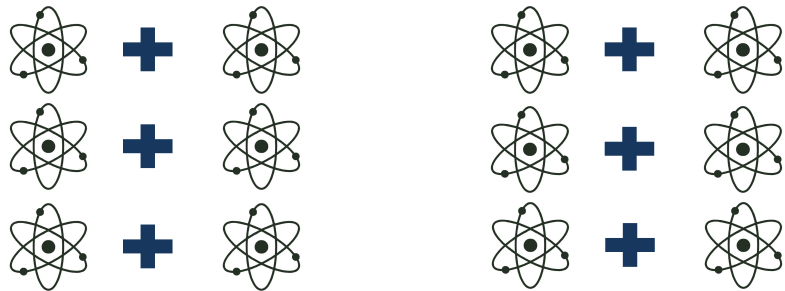
Collapse dynamics in a nutshell



Microscopic superposition in space. Collapse very weak, modulo tiny deviations



Macroscopic superposition in space. Collapse very strong. The larger the delocalization in space and the number of particles, the faster the collapse



Many-body single-particle superpositions in space. Collapse very weak, modulo tiny deviations

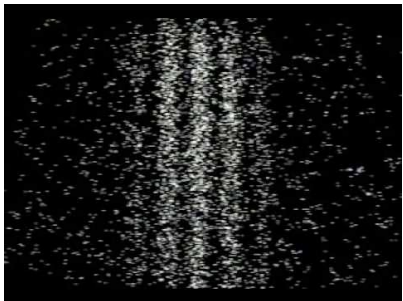


Superpositions in other d.o.f. very weak if they do not imply delocalization in space

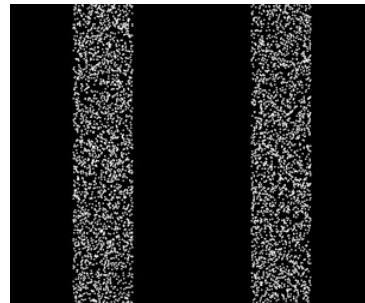
How to test collapse models

Interferometric experiments

Create a large superposition, in terms of mass, distance and duration, and perform a “double slit” experiment



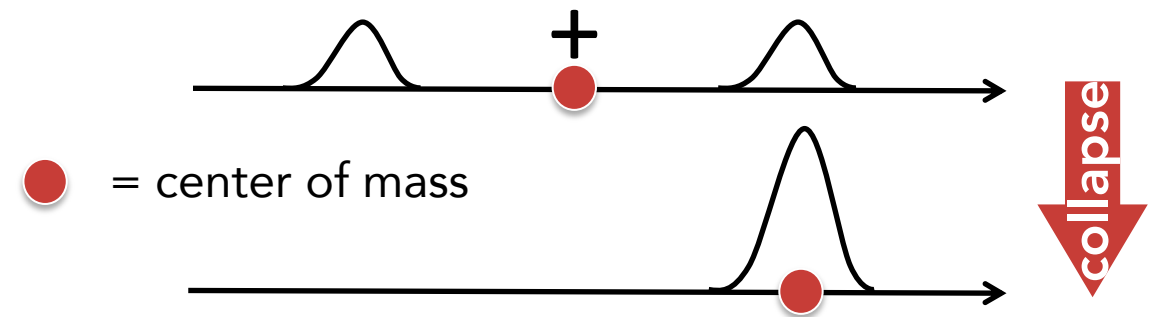
Prediction of quantum mechanics
(no environmental noise)



Prediction of collapse models
(no environmental noise)

Non interferometric experiments

S. Donadi, L. Ferialdi & A. Bassi, “Collapse dynamics are diffusive”, arXiv:2209.09697



A collapse of the wave function changes the position of the center of mass → **Collapse-induced Brownian motion**



Quantum prediction
(no environmental noise)

Collapse prediction
(no environmental noise)

Advantages and disadvantages

Interferometric experiments



These are a **direct test** of the quantum superposition principle and of collapse models.



They are **difficult**. The whole field of quantum optomechanics boomed also with the aim of creating macroscopic quantum states.

Non interferometric experiments



They are a **direct test** of collapse models and an **indirect test** of the quantum superposition principle.



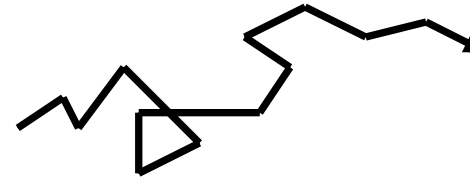
They are **easier** because **no quantum superposition** is needed to test the collapse-induced Brownian motion.

How to test the collapse noise

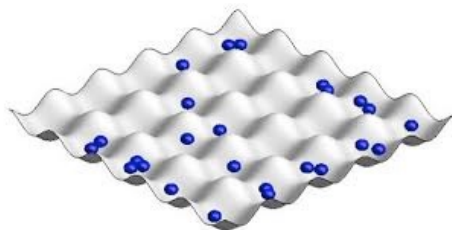
Quantum Mechanics



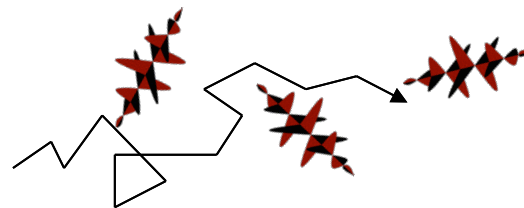
Collapse models



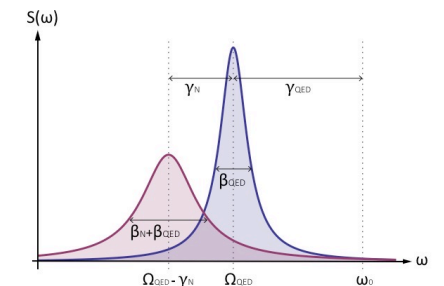
A **gas** will **expand** (heat up) faster than what predicted by QM



Charged particles will **emit** radiation, whereas QM predicts no emission



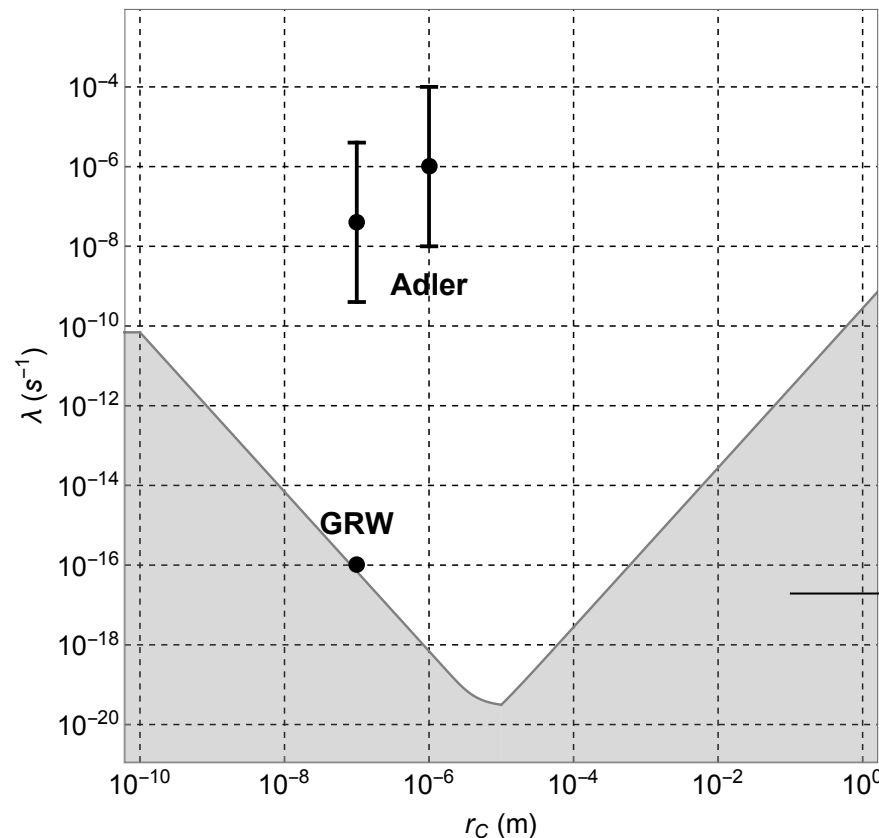
A **cantilever's** motion cannot be **cooled down** below a given limit



Tests of the CSL model

$$\mathcal{G}(\mathbf{x}) = \frac{\lambda}{m_0^2} e^{-\mathbf{x}^2/4r_C^2}$$

Two phenomenological parameters. λ measures the strength of the collapse, r_C the space resolution of the collapse. m_0 is a reference mass, equal to that of a nucleon



• = Theoretical guesses

Lower bound: for such values of the parameters, the collapse is too weak and ineffective at the “macroscopic” level.
Working assumption: a graphene disk with $N = 10^{11}$ amu, delocalized over $d = 10^{-5} \text{ m}$, should collapse in $T = 10^{-2} \text{ s}$

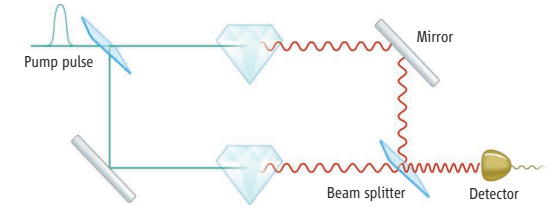
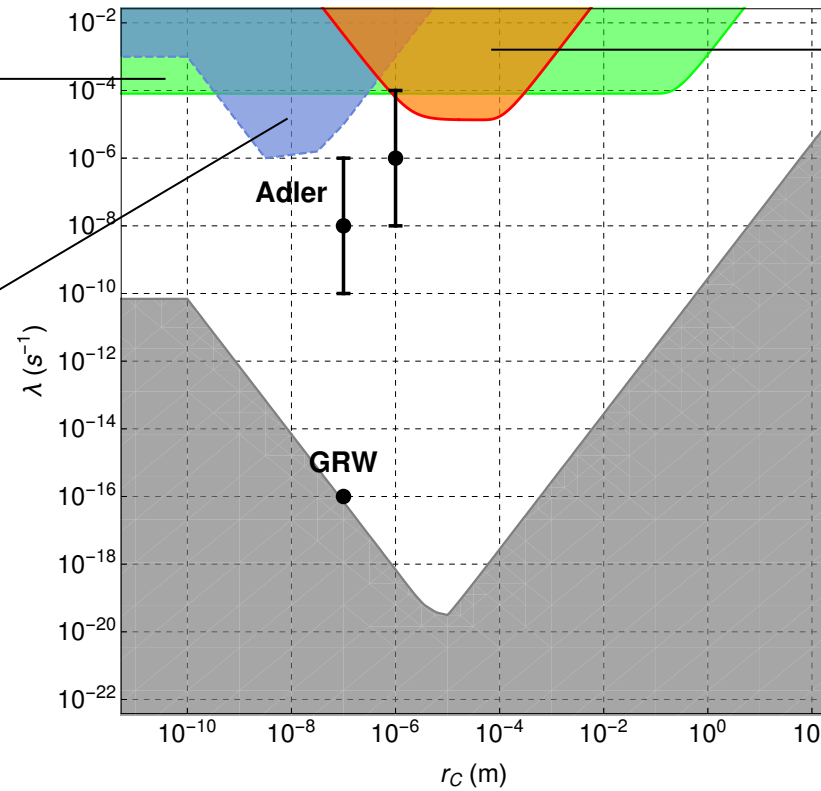
Interferometric Experiments



Atom Interferometry

T. Kovachy *et al.*, Nature 528, 530 (2015)

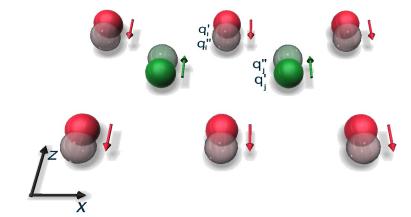
$M = 87$ amu
 $d = 0.54$ m
 $T = 1$ s



Entangling Diamonds

K. C. Lee *et al.*, Science. 334, 1253 (2011).
 S. Belli *et al.*, PRA 94, 012108 (2016).
 B. Schirnski *et al.*, ArXiv:2209.06635.

$M = 10^{16}$ amu
 $d = 10^{-11}$ m \rightarrow in reality much smaller
 $T = 10^{-12}$ s



Molecular Interferometry

Y.Y Fein *et al.*, Nature Physics 15, 1242 (2019)
 M. Toros *et al.*, PLA 381, 3921 (2017)

$M = 10^5$ amu
 $d = 10^{-7}$ m
 $T = 10^{-3}$ s

To improve interferometric tests, it will likely be necessary to go to micro-gravity environment in outer space \rightarrow MAQRO

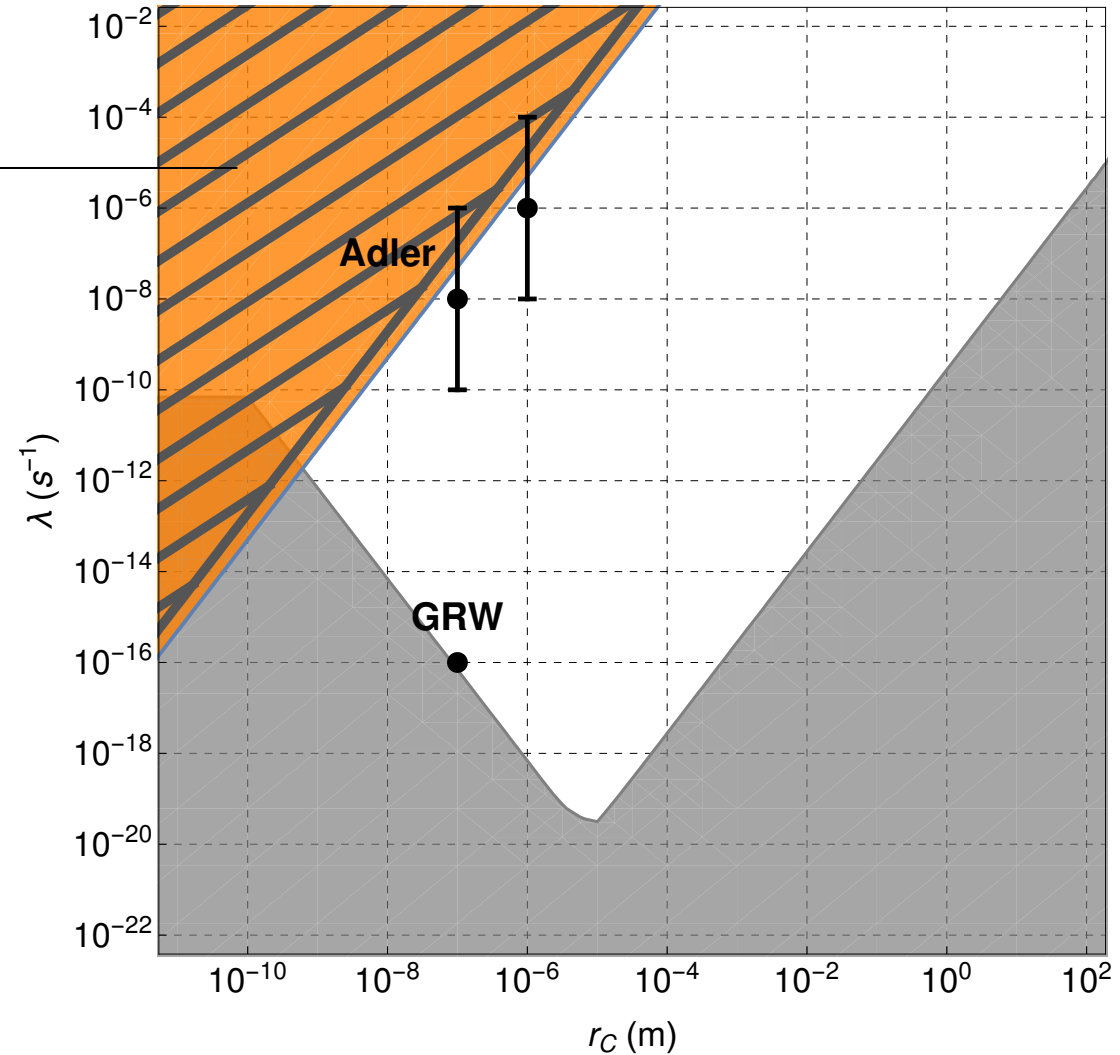
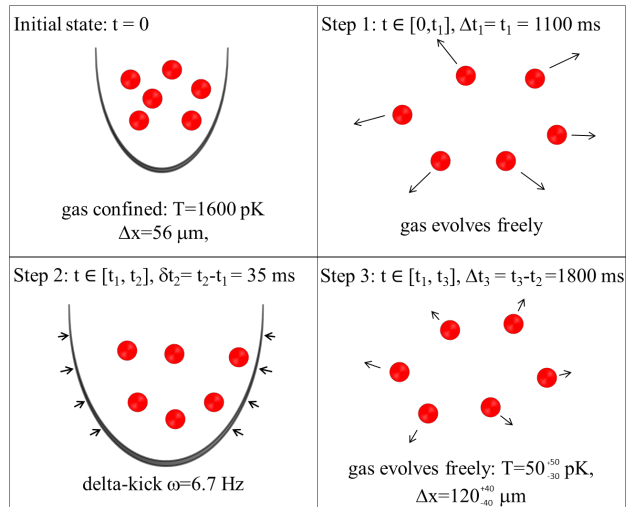
Non - Interferometric Experiments

Cold atom gas

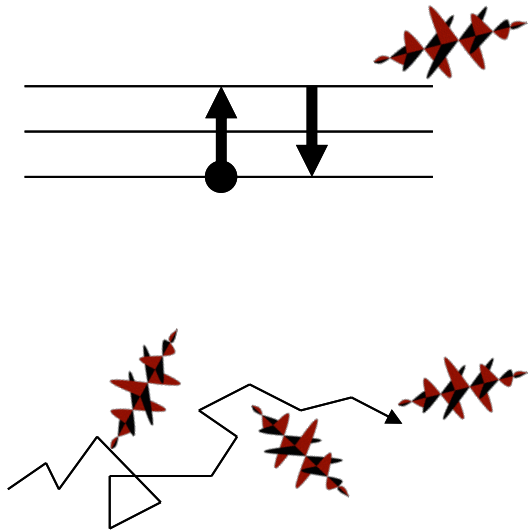
F. Laloë *et al.* Phys. Rev. A 90, 052119 (2014)

T. Kovachy *et al.*, Phys. Rev. Lett. 114, 143004 (2015)

M. Bilardello *et al.*, Physica A 462, 764 (2016)

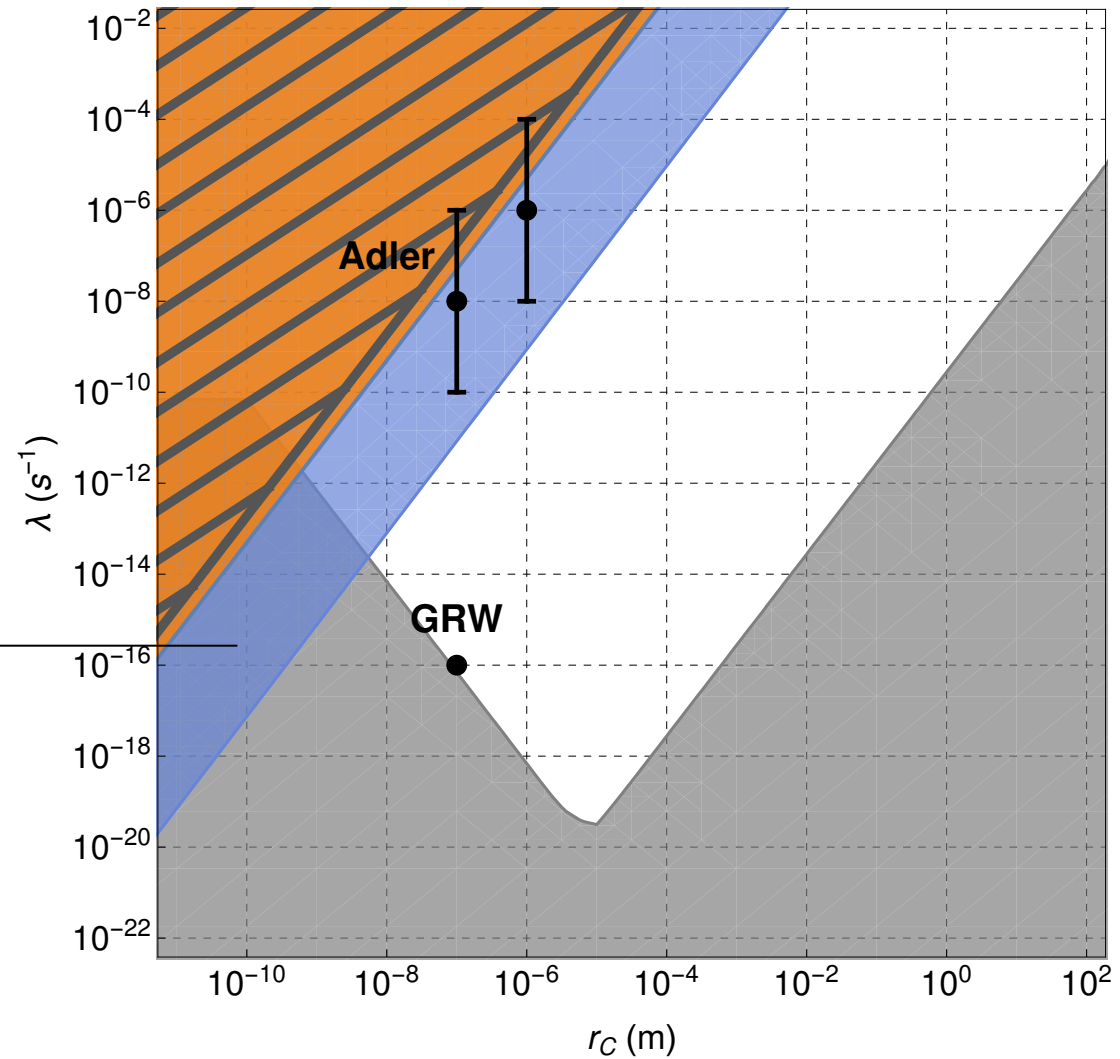


Non - Interferometric Experiments

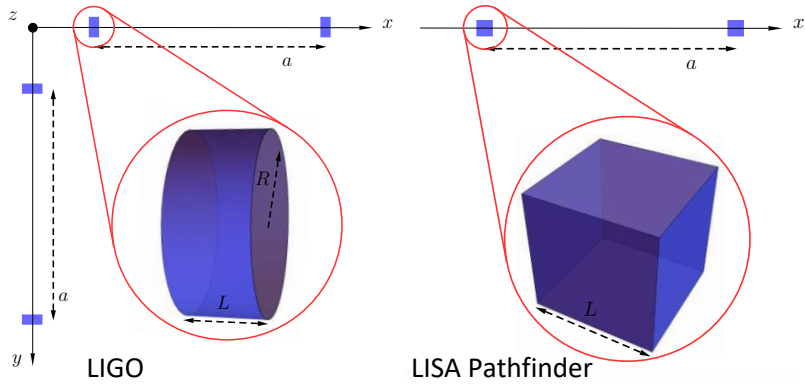


X rays

S.L. Adler *et al.*, Jour. Phys. A 40, 13395 (2009)
S.L. Adler *et al.*, Journ. Phys. A 46, 245304 (2013)
A. Bassi & S. Donadi, Annals of Phys. 340, 70 (2014)
S. Donadi & A. Bassi, Journ. Phys. A 48, 035305 (2015)
C. Curceanu *et al.*, J. Adv. Phys. 4, 263 (2015)
+ several more



Non - Interferometric Experiments

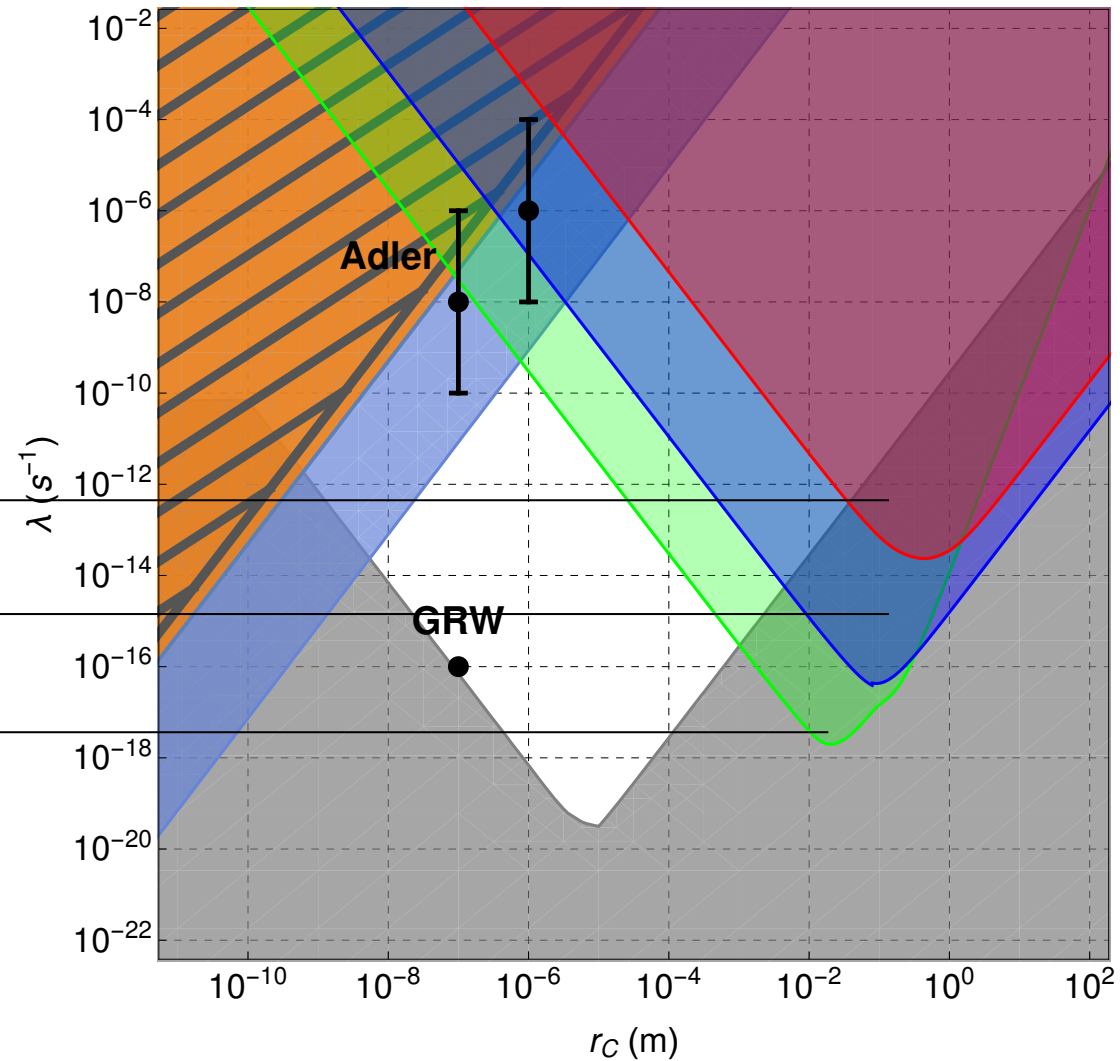
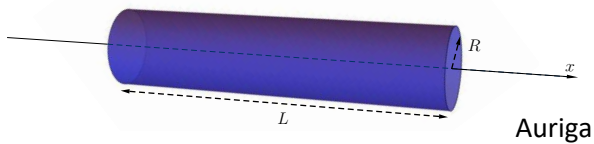


Auriga

Ligo

Lisa Pathfinder

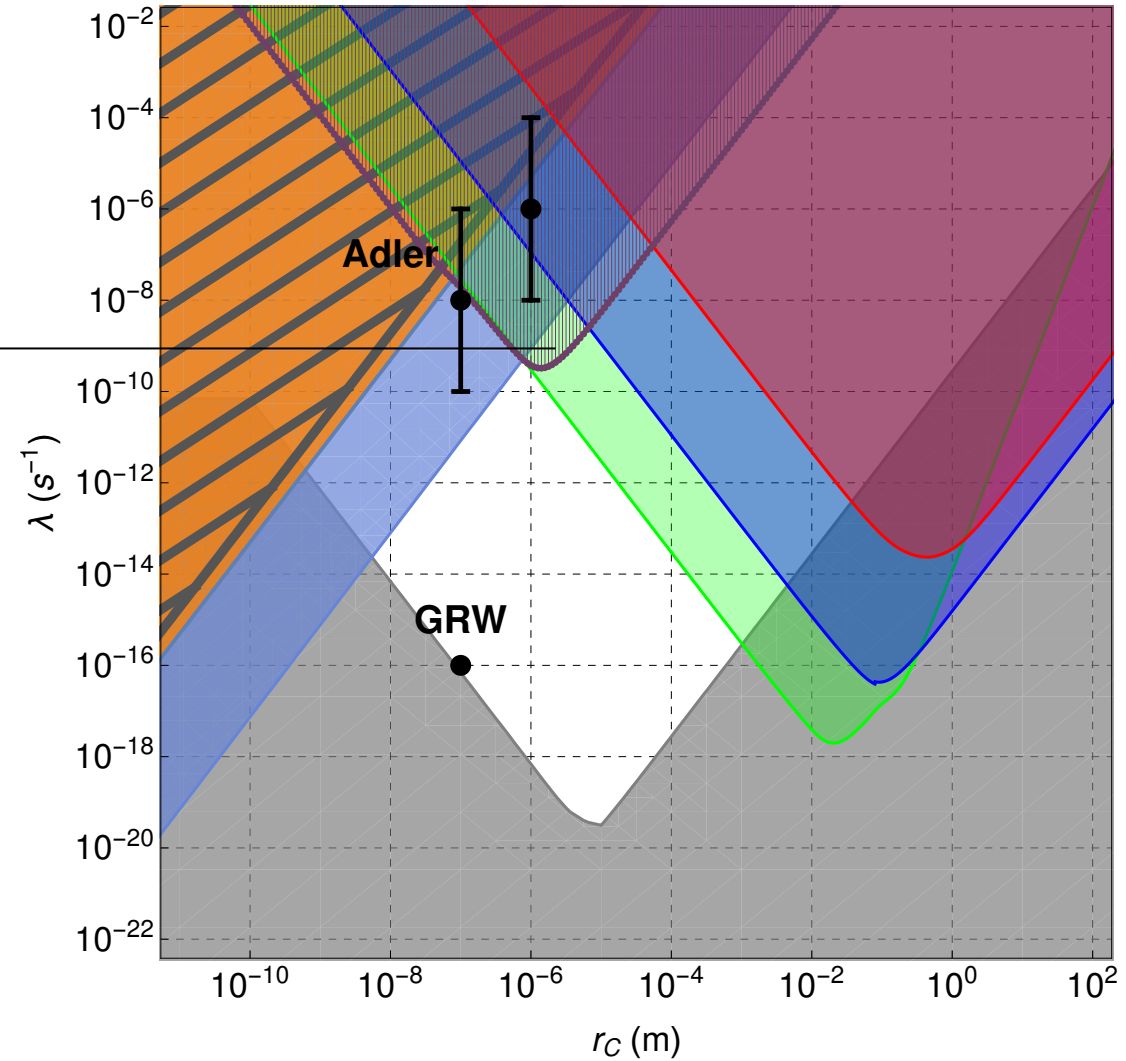
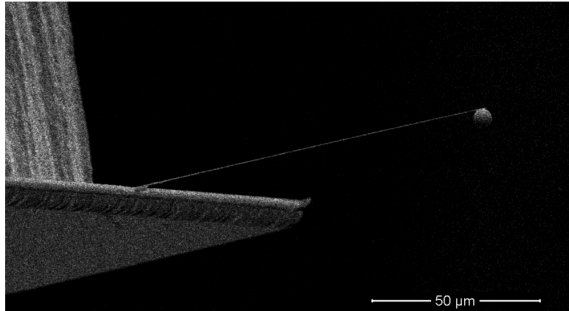
M. Carlesso *et al.* Phys. Rev. D 94, 124036 (2016)



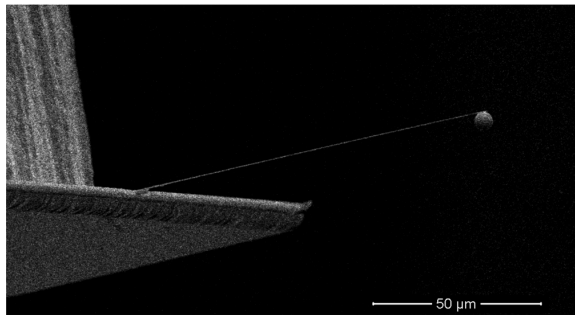
Non - Interferometric Experiments

Cantilever

A. Vinante *et al.*, Phys. Rev. Lett. 116, 090402 (2016)

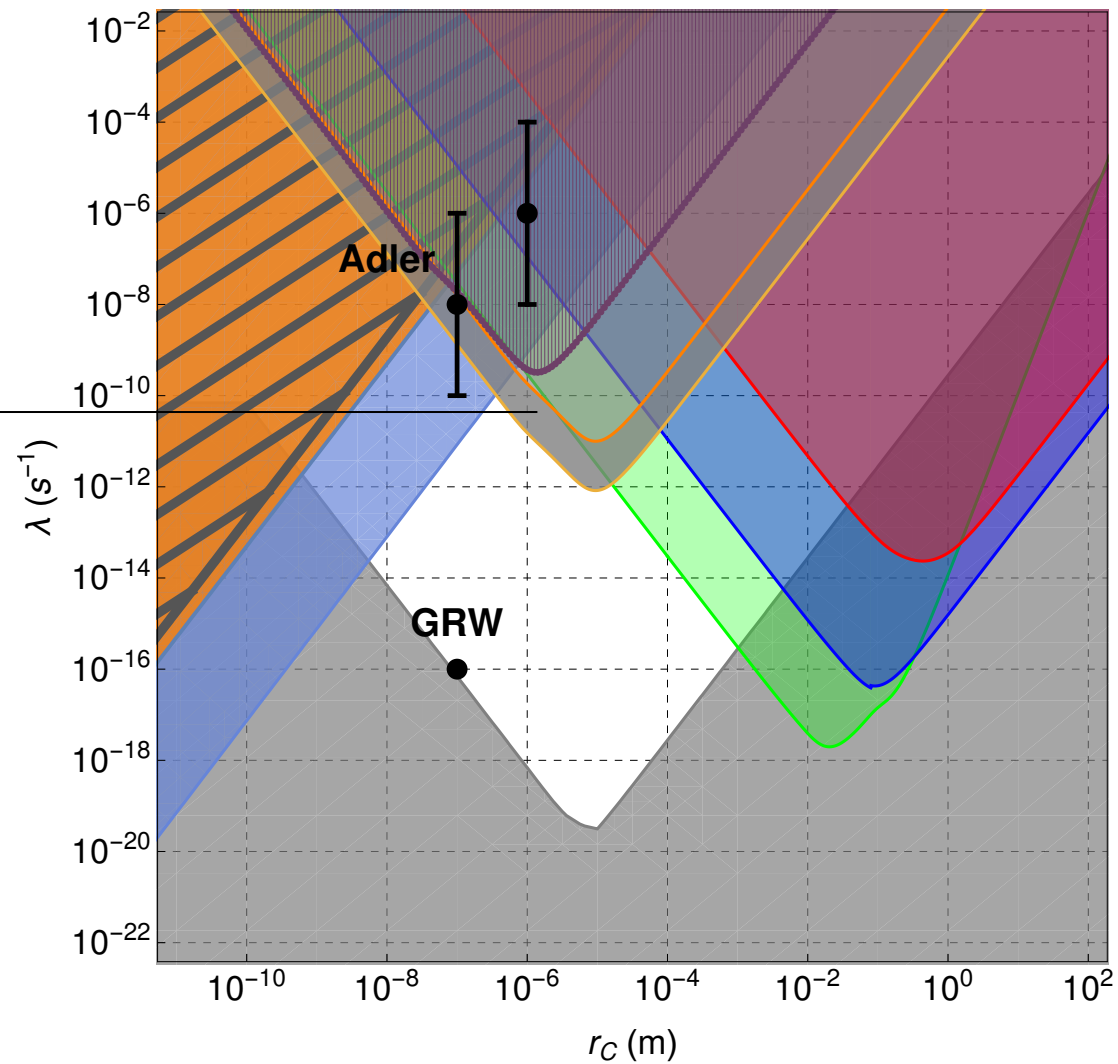


Non - Interferometric Experiments



Cantilever – update 1

A. Vinante *et al.*, *Phys. Rev. Lett.* 119, 110401 (2017).



Non - Interferometric Experiments

Cantilever - Update 2

A. Vinante *et al.*, *Phys. Rev. Lett.* 125, 100404 (2020)

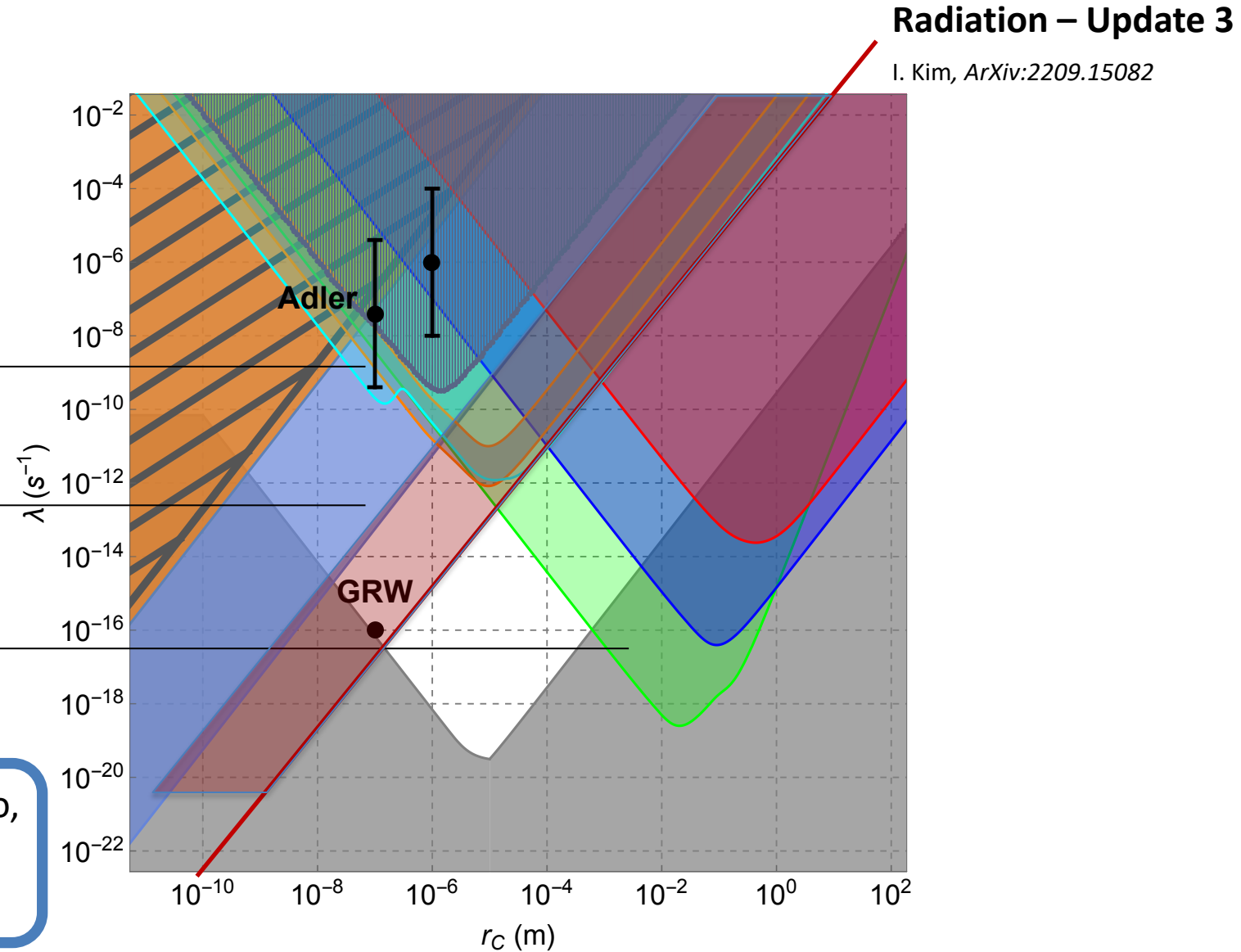
Radiation – Update 1

K. Pispicchia *et al.*, *Entropy* 19, 319 (2017)

Gravitational Wave detectors – Update 1

M. Carlesso *et al.*, *N. Journ. Phys* 20, 083022 (2018)

M. Carlesso, S. Donadi, L. Ferialdi, M. Paternostro,
H. Ulbricht, A. Bassi,
Nature Physics 18, 243-250 (2022)



Non - Interferometric Experiments

Cantilever - Update 2

A. Vinante *et al.*, *Phys. Rev. Lett.* 125, 100404 (2020)

Radiation – Update 1

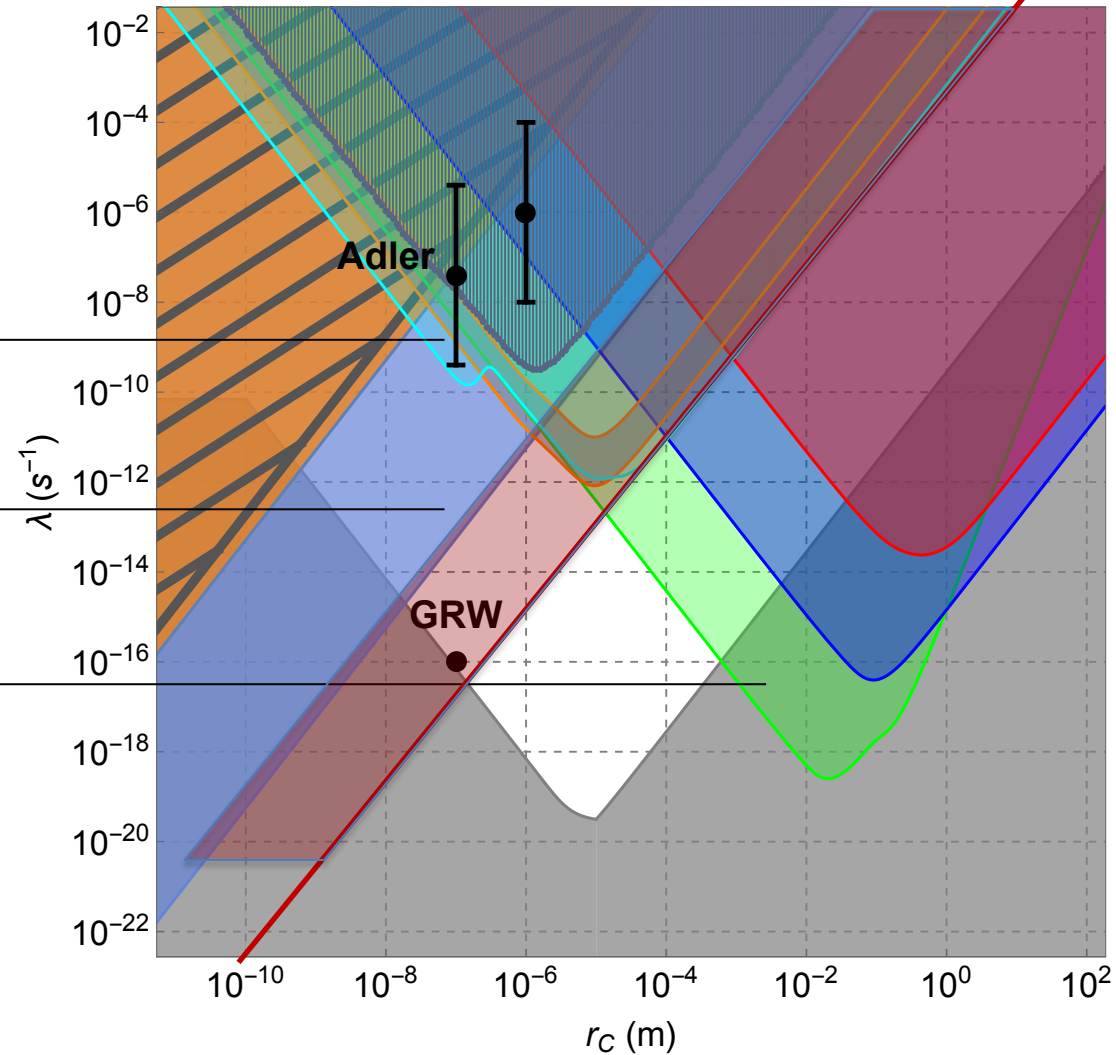
K. Pispicchia *et al.*, *Entropy* 19, 319 (2017)

Gravitational Wave detectors – Update 1

M. Carlesso *et al.*, *N. Journ. Phys* 20, 083022 (2018)

Radiation – Update 3

I. Kim, *ArXiv:2209.15082*



Non-white noises will change
the picture significantly

Acknowledgments

The Group (www.qmts.it)

- Postdocs: S. Donadi, J.L. Gaona Reyes
- Ph.D. students: F. Cesa, L. Figurato, A. Ghundi, M. Vischi, G. Di Bartolomeo



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