

Weak Interaction in Nuclear Astrophysics

-Main actor: Gamow-Teller transitions-

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Euro. SS on Exp. Nucl. Astrophysics

Santa Tecla, Italy, Sep. 18 – Sep.27, 2011

GT : Important weak response, simple $\sigma\tau$ operator

- ✧ Representing “Spin Isospin” response of nuclei,
that are unique quantum numbers in Atomic Nuclei.
- ✧ Good Probe to study the Key Part of the Nuclear Structure.
- ✧ Astrophysical Interest.
- ✧ Studied by β decay and Charge-Exchange reactions

Crucial Weak Processes during the Core Collapse

mainly by τ & $\sigma\tau$



**Gamow-Teller (GT)
transitions**

Langanke & Martinez-Pinedo
Rev.Mod.Phys.75('04)819

Balantekin & Fuller
J.Phys.G 29('03)2513

$$p + e^- \rightleftharpoons n + \nu_e,$$

$$n + e^+ \rightleftharpoons p + \bar{\nu}_e.$$

(A,Z)=nuclei in the
Fe, Ni region

$$(A, Z) + e^- \rightleftharpoons (A, Z-1) + \nu_e,$$

$$(A, Z) + e^+ \rightleftharpoons (A, Z+1) + \bar{\nu}_e.$$

$$\nu + N \rightleftharpoons \nu + N$$

$$N +$$

**β -decay, e-capture,
 ν -induced reactions**

$$\nu + (A, Z) \rightleftharpoons \nu + (A, Z),$$

$$\nu + e^\pm \rightleftharpoons \nu + e^\pm,$$

$$\nu + (A, Z) \rightleftharpoons \nu + (A, Z)^*,$$

$$e^+ + e^- \rightleftharpoons \nu + \bar{\nu},$$

$$(A, Z)^* \rightleftharpoons (A, Z) + \nu + \bar{\nu}.$$

Properties of GT transitions

Caused by the $\sigma\tau$ operator : a simple operator !

1) $|i\rangle$ and $|f\rangle$ states should have similar spatial shapes.

- there is no space-type operator -

2) σ operator: states with $j_>$ and $j_<$ configurations are connected.

Selection Rules

$$|\Delta J| = 0, 1$$

3) τ operator: isospin quantum number T plays an important role (isospin selection rule)

$$|\Delta T| = 0, 1$$

→ GT transitions are sensitive to Nuclear Structure !

→ GT transitions in each nucleus are UNIQUE !

***Nucleus = Bell

Operators = Hammers ??

* $B(Op)$ ex. $B(E1)$, $B(M1)$, $B(GT)$...

Reduced Transition Strength

$$\propto |\langle f | Op | i \rangle|^2$$

*various kinds of Operators (Op)

Various Operators / Various Hammers!



wooden hammers



metal hammers

hammers
=operators

The sound from the bell is different depending on hammers!

The mode of nuclear excitation is decided by an operator!

Various Reaction Mechanism / How and Where you hit the bell!



how and where you hit
=reaction mechanism

The sound from the bell is different how and where you hit!

The strength of nuclear excitation is dependent on them!

Vibration Modes in Nuclei (Operators)

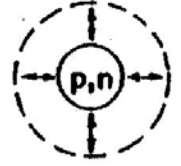
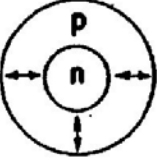
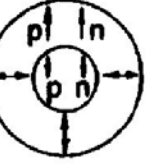
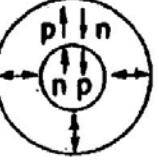
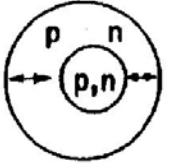
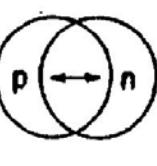
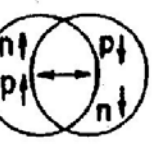
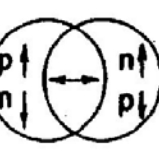
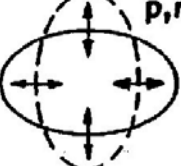


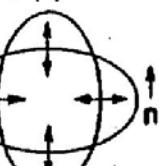


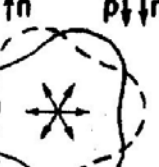

Microscopic classification of giant resonances

	$\Delta S = 0$ $\Delta T = 0$	$\Delta S = 0$ $\Delta T = 1$	$\Delta S = 1$ $\Delta T = 0$	$\Delta S = 1$ $\Delta T = 1$
L = 0		$\sum \tau_i$ IAS		$\sum \vec{\sigma}_i \tau_i$ GTR
2 nd order	$\sum r_i^2$ ISGMR	$\sum r_i^2 \tau_i$ IVGMR	$\sum r_i^2 \vec{\sigma}_i$ ISSMR	$\sum r_i^2 \vec{\sigma}_i \tau_i$ IVSMR
L = 1		$\sum r_i Y_m^1 \tau_i$ IVGDR	$\sum r_i Y_m^1 \vec{\sigma}_i$ ISSDR	$\sum r_i Y_m^1 \vec{\sigma}_i \tau_i$ IVSDR
2 nd order	$\sum r_i^3 Y_m^1$ ISGDR			
L = 2	$\sum r_i^2 Y_m^2$ ISGQR	$\sum r_i^2 Y_m^2 \tau_i$ IVGQR	$\sum r_i^2 Y_m^2 \vec{\sigma}_i$ ISSQR	$\sum r_i^2 Y_m^2 \vec{\sigma}_i \tau_i$ IVSQR
L = 3	$\sum r_i^3 Y_m^3$ ISGOR	$\sum r_i^3 Y_m^3 \tau_i$ IVGOR	$\sum r_i^3 Y_m^3 \vec{\sigma}_i$ ISSOR	$\sum r_i^3 Y_m^3 \vec{\sigma}_i \tau_i$ IVSOR

$\Delta S=1$:
spin excitation

$\Delta T=1$:
IV excitation
(isospin related!)

Vibration Modes in Nuclei (Schematic)

	Electric Mode ($\Delta S=0$)		Magnetic Mode ($\Delta S=1$)	
	IS ($\Delta T=0$)	IV ($\Delta T=1$)	IS ($\Delta T=0$)	IV ($\Delta T=1$)
L=0				
L=1				
L=2				
L=3				

IS-Giant
Monopole
Resonance
(GMR)

Vibration Modes in Nuclei (Schematic)

	Electric Mode ($\Delta S=0$)		Magnetic Mode ($\Delta S=1$)	
	IS ($\Delta T=0$)	IV ($\Delta T=1$)	IS ($\Delta T=0$)	IV ($\Delta T=1$)
L=0				
L=1				
L=2				
L=3				

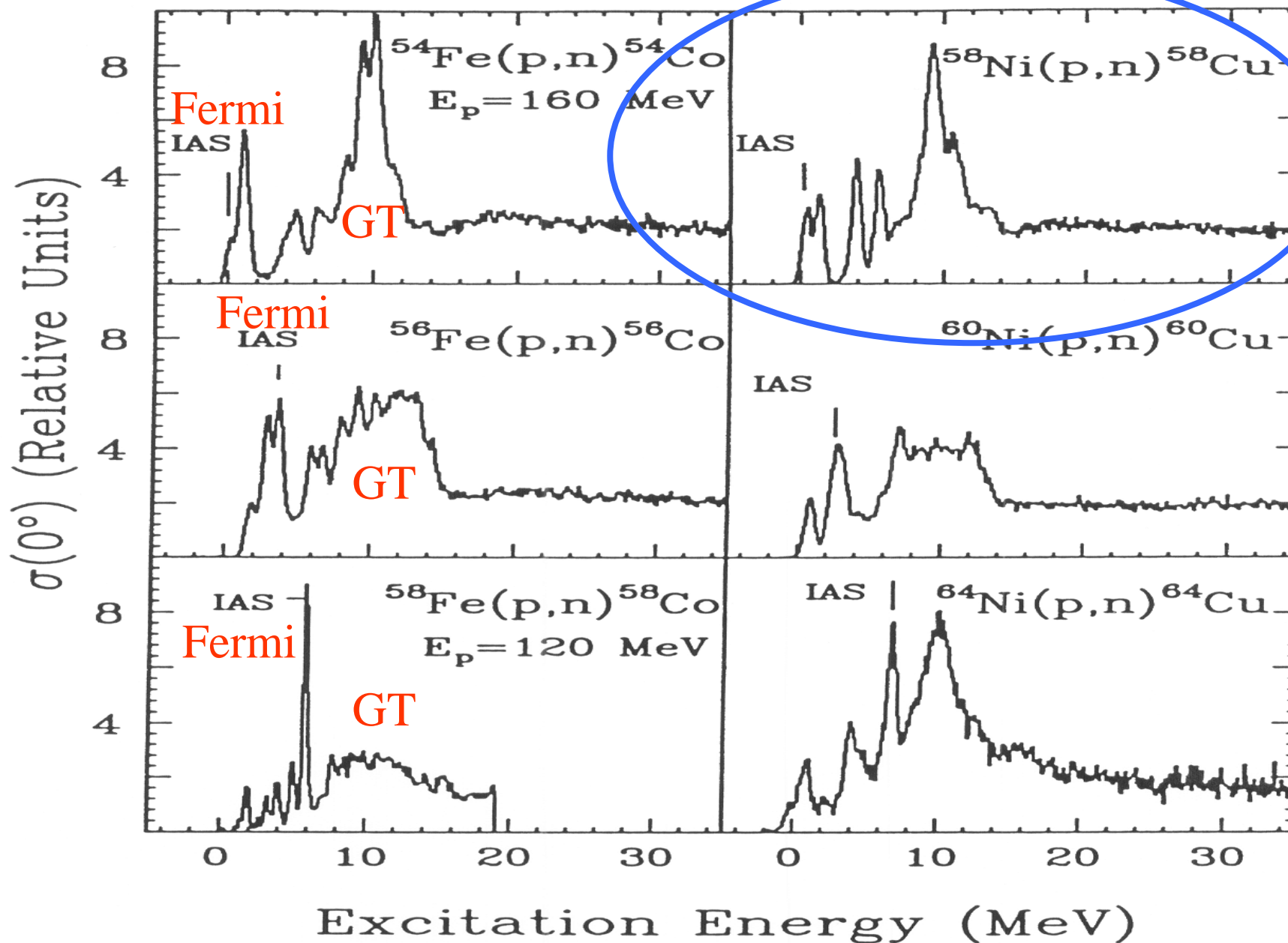
IV-Giant
Monopole
Resonance
(IVGMR)

Vibration Modes in Nuclei (Schematic)

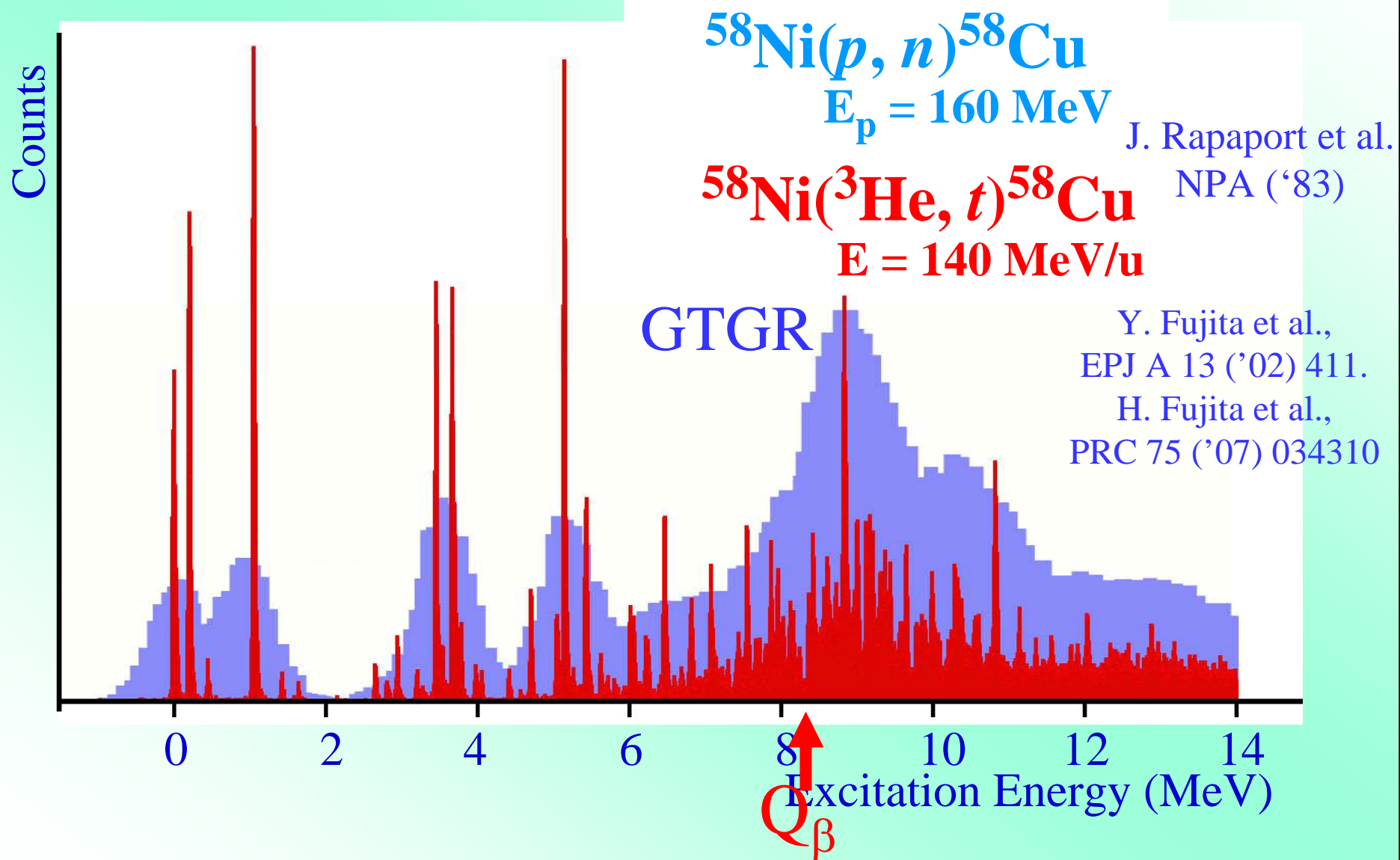
	Electric Mode ($\Delta S=0$)		Magnetic Mode ($\Delta S=1$)	
	IS ($\Delta T=0$)	IV ($\Delta T=1$)	IS ($\Delta T=0$)	IV ($\Delta T=1$)
L=0				
L=1				
L=2				
L=3				

Gamow-Teller mode
($\sigma\tau$)

(p, n) spectra for Fe and Ni Isotopes



Comparison of (p, n) and ($^3\text{He}, t$) 0° spectra



**Basic common understanding of β -decay and Charge-Exchange reaction

β decays :

Absolute $B(\text{GT})$ values,

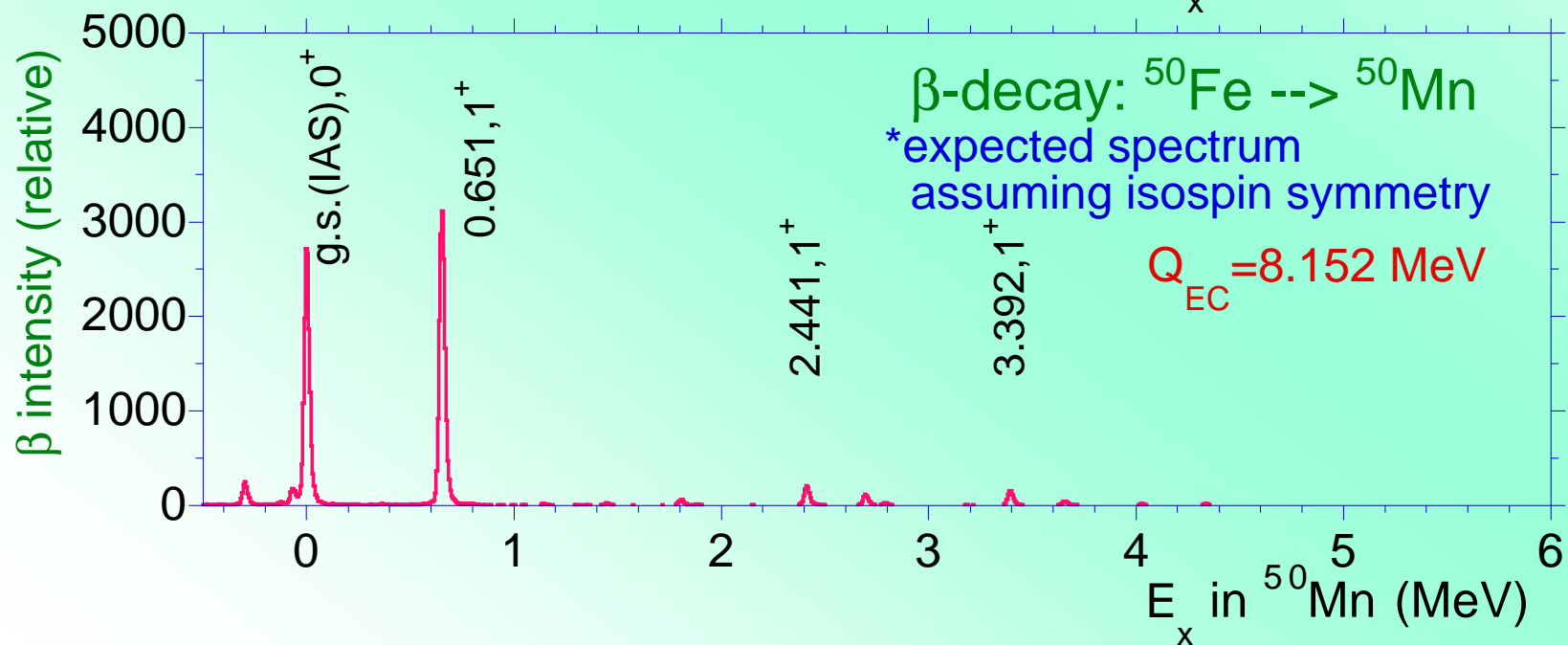
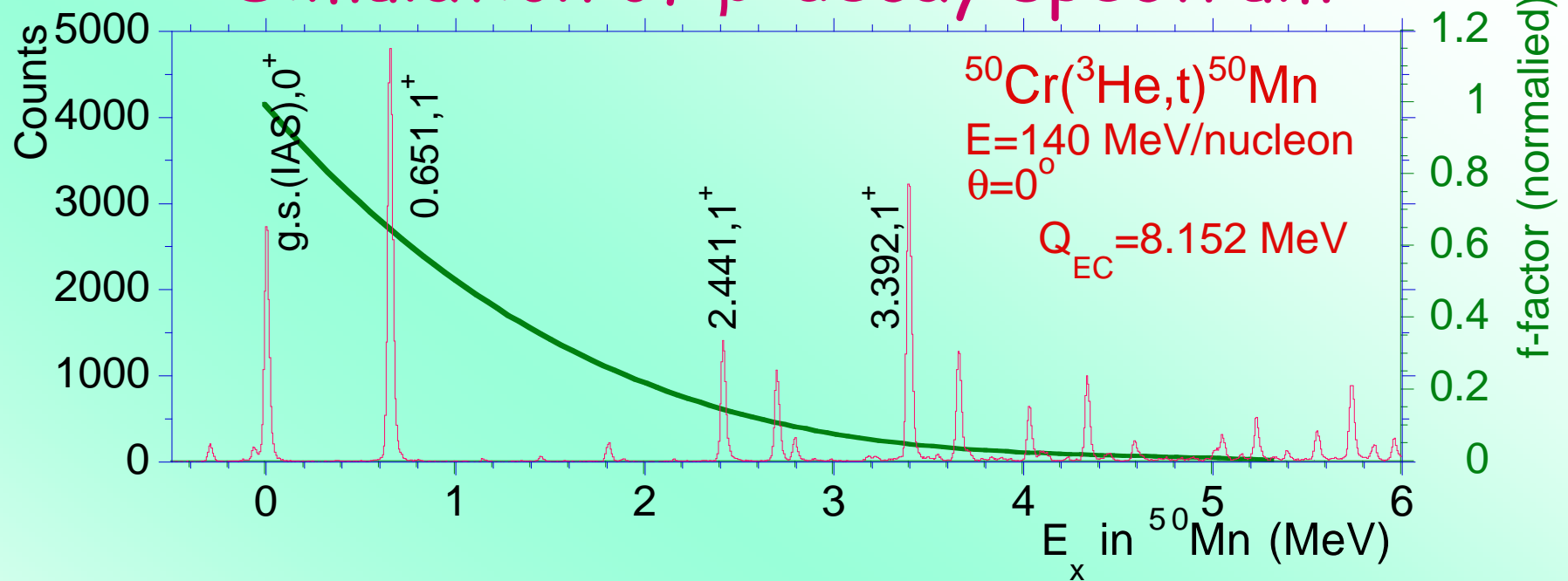
but usually the study is limited to low-lying state

$(^3\text{He},t)$ reaction at 0° :

Relative $B(\text{GT})$ values, but **Highly Excited States**

** Both are important for the study of GT transitions!

Simulation of β -decay spectrum



β -decay transition strength

* β -decay transition strength

$$\frac{K}{ft_{1/2}} = B(\text{GT}) \lambda^2 + B(\text{F}) \quad \text{where} \quad \lambda = \frac{g_A}{g_V}$$

then

$$\frac{1}{t_{1/2}} = \frac{f}{K g_V^2} [g_A^2 B(\text{GT}) + g_V^2 B(\text{F})]$$

↑
transition
strength

↑ ↑
coupling
constants

where

$$B(\text{GT}) \propto |\langle f | \sigma \tau | i \rangle|^2$$

$$B(\text{F}) \propto |\langle f | \tau | i \rangle|^2$$

β -decay & Nuclear Reaction

* β -decay GT tra. rate = $\frac{1}{t_{1/2}} = f \frac{\lambda^2}{K} B(\text{GT})$

$B(\text{GT})$: reduced GT transition strength
 $\propto (\text{matrix element})^2 = |\langle f | \sigma \tau | i \rangle|^2$

*Nuclear (CE) reaction rate (cross-section)
= reaction mechanism

⊗ operator

⊗ structure

$= (\text{matrix element})^2$

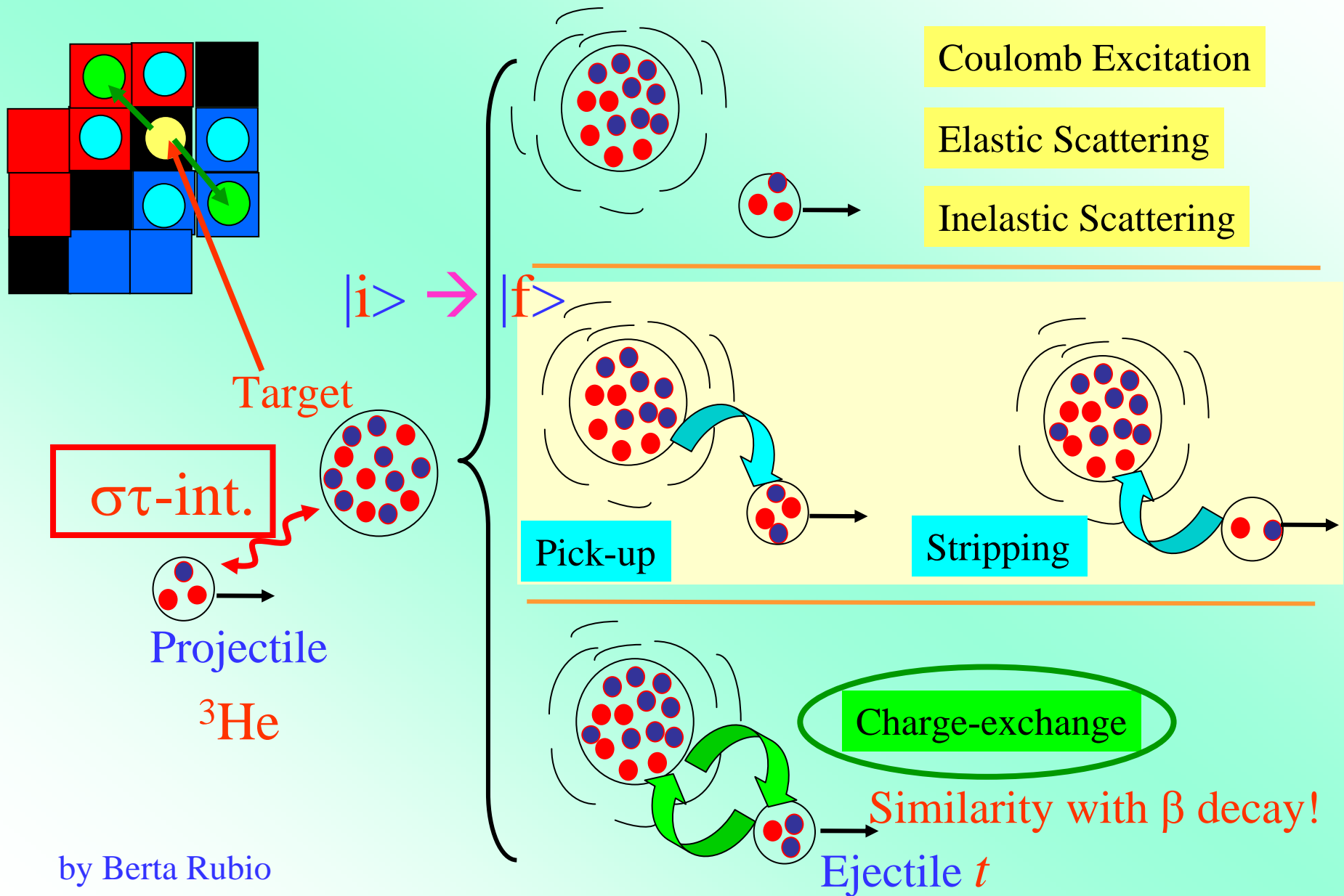
*At intermediate energies ($100 < E_{\text{in}} < 500$ MeV)

→ $d\sigma/d\omega(q=0)$: proportional to $B(\text{GT})$

****Nuclear Excitations****
by
Charge Exchange Reaction
and β -Decay

Study of Weak Response of Nuclei
by means of
Strong Interaction !?

Direct Reactions with Light Projectiles



by Berta Rubio

****Connection between
 β -decay and ($^3\text{He},t$) reaction****

**by means of
Isospin Symmetry**

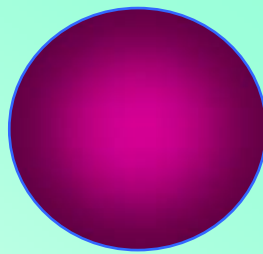
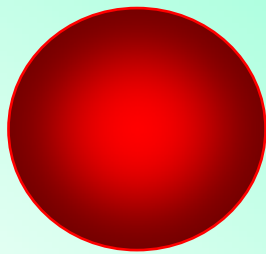
Nuclei & Coin



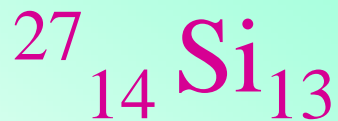
back

front

= Coin



= Isomer



$$T_z = (1/2)N + (-1/2)Z$$

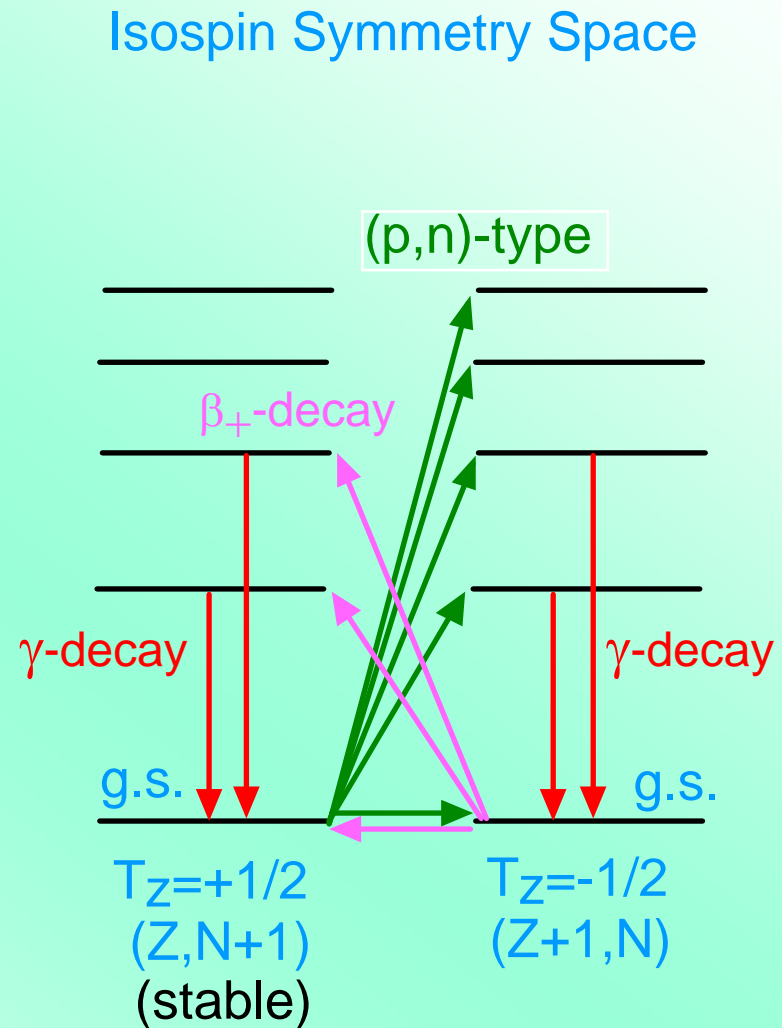
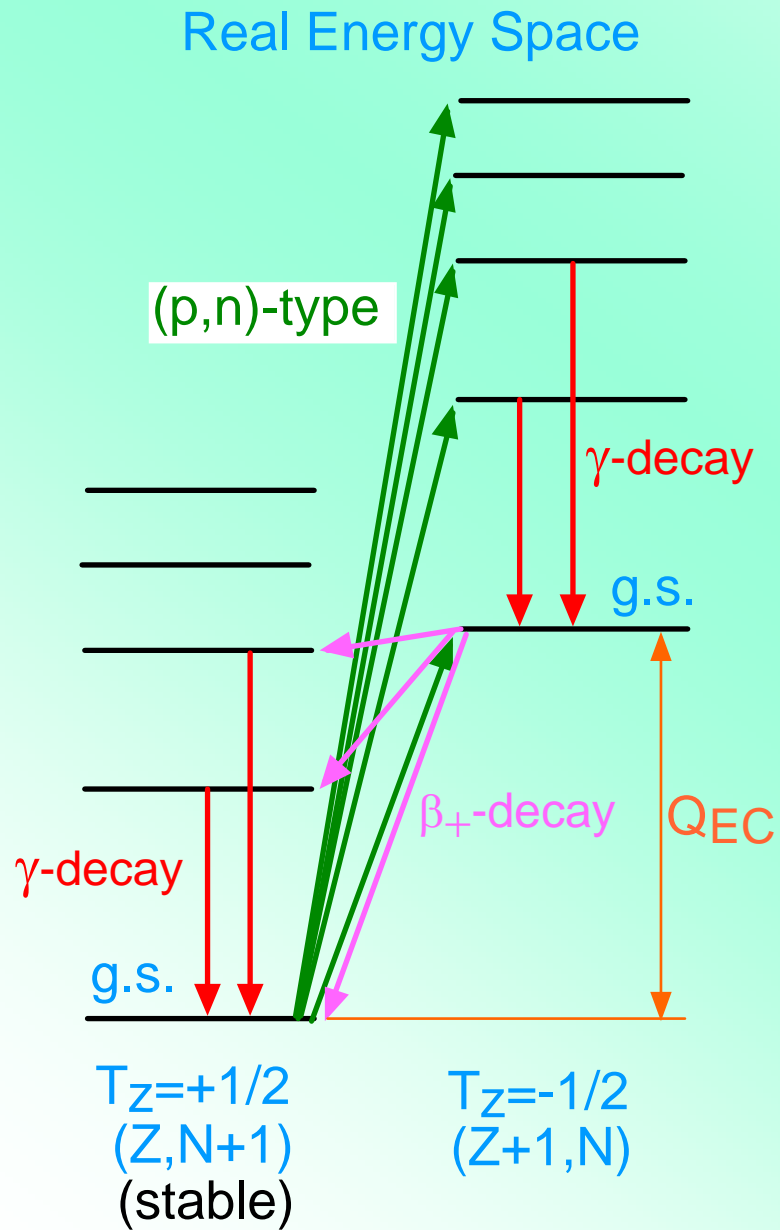
*z-component: conserved

$$T_z = 1/2$$

$$T_z = -1/2$$

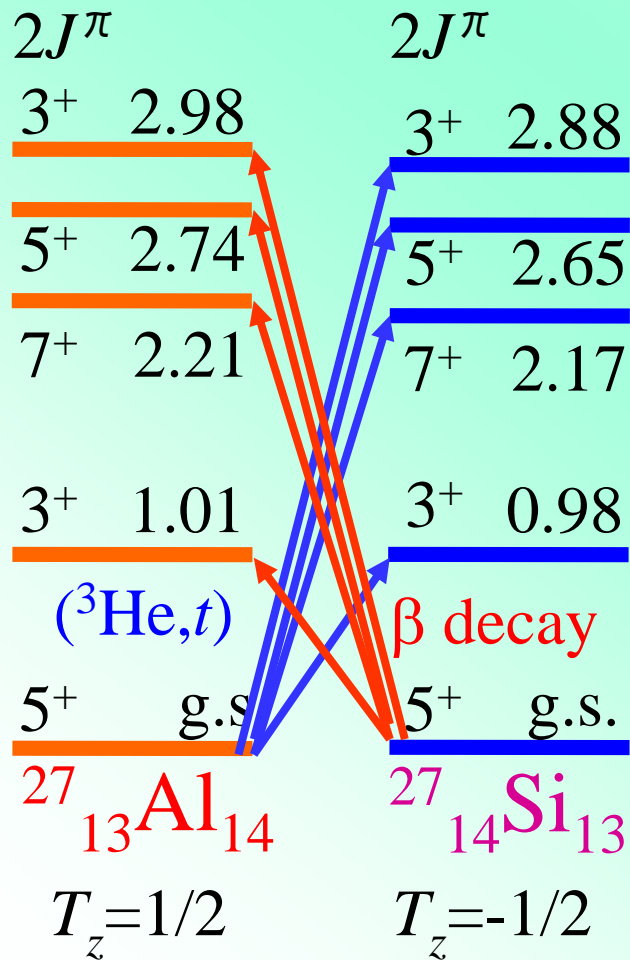
isospin $T = 1/2, 3/2, \dots$

Analogous Structures and Transitions in $T=1/2$ System

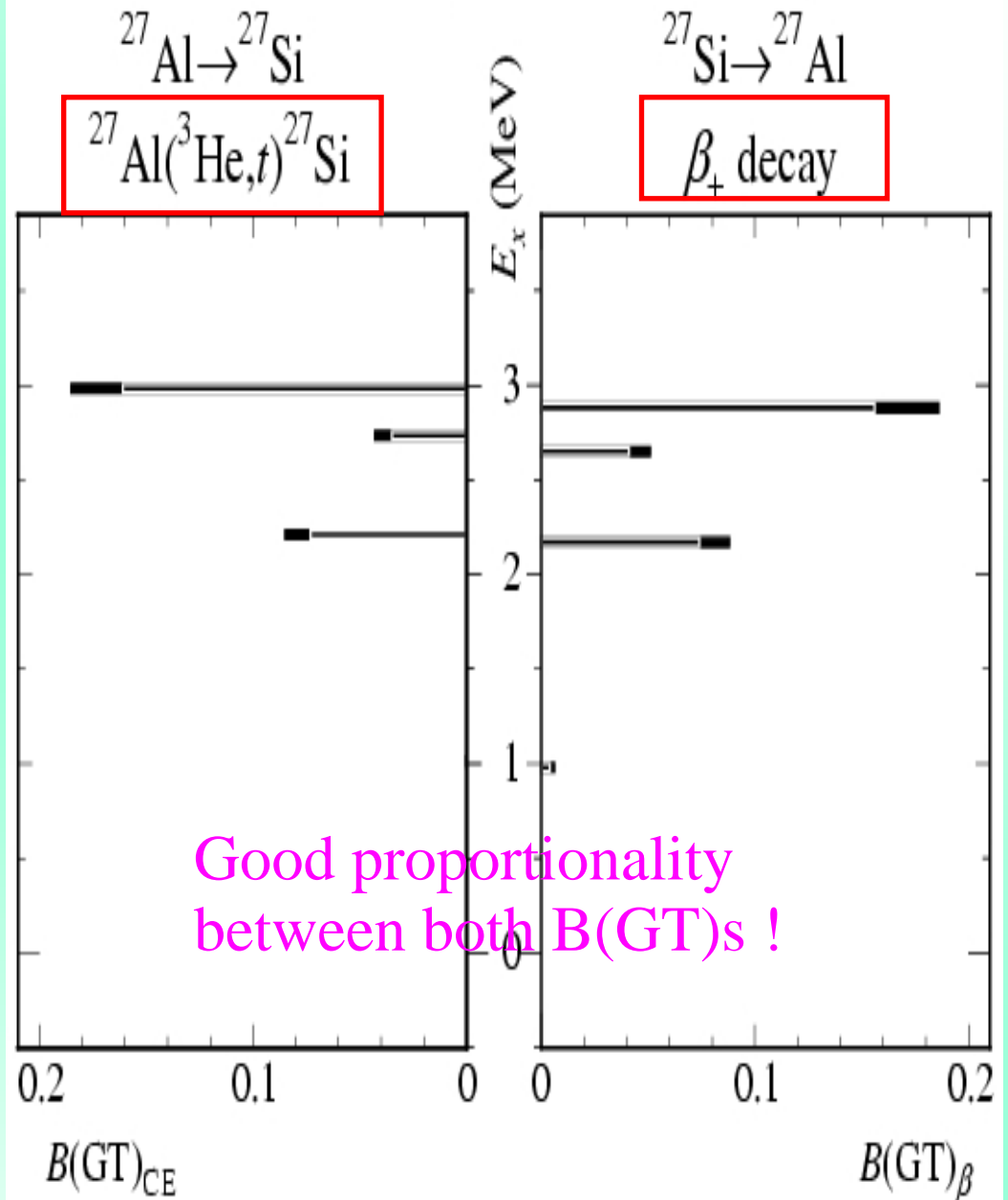


Symmetry in A=27 System

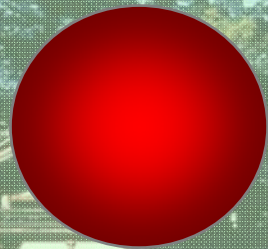
$$\frac{d\sigma}{d\Omega}(q \rightarrow 0) = KM |J|^2 B(\text{GT})$$



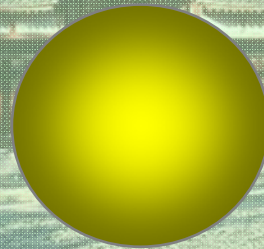
Experiments



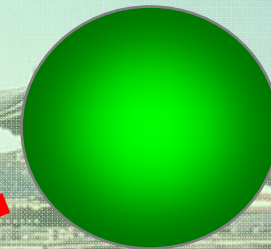
T=1 Isospin Symmetry



GT



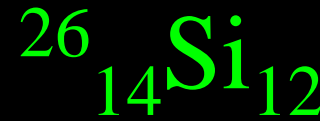
GT



$$T_z = +1$$



$$T_z = 0$$



$$T_z = -1$$

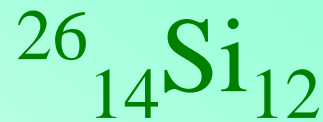
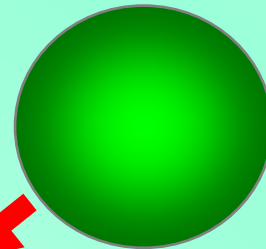
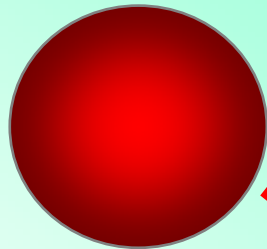
Nuclei & Coin



= Coin

back

front



$$T_z = +1$$

$$T_z = 0$$

$$T_z = -1$$

= Nuclei

Symmetry in

- 1) structure
- 2) transitions

isospin $T=1$ triplet

Transitions in real & isospin space (T=1)

Symmetry Transitions from T=1 Nuclei

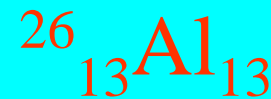
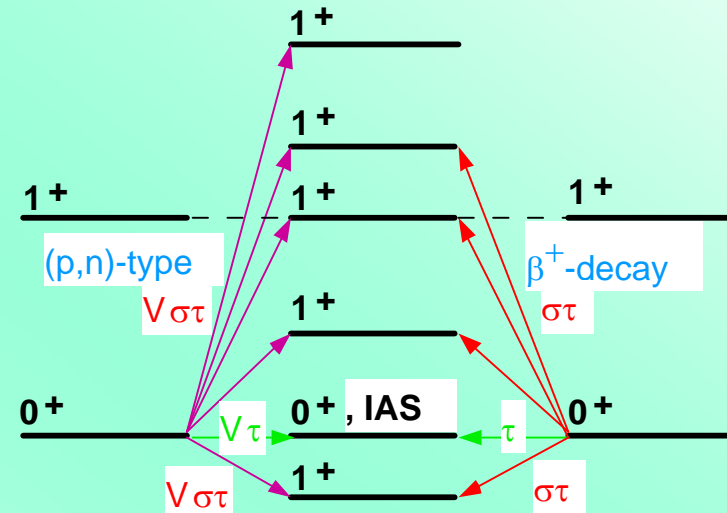
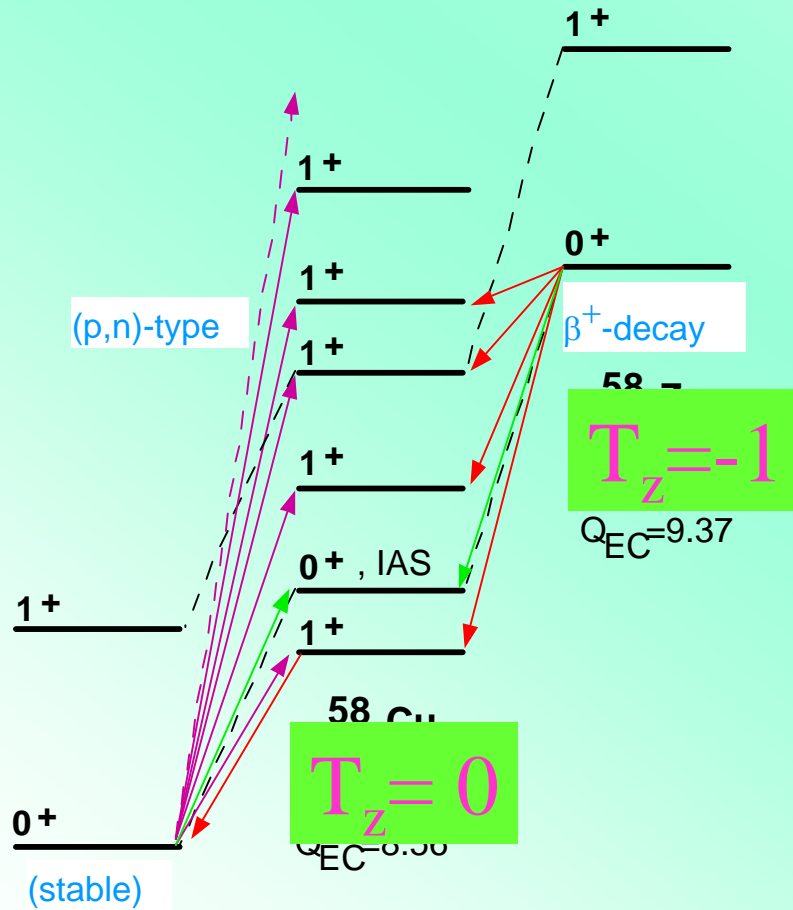
$$T_z=+1 \rightarrow T_z=0 \leftarrow T_z=-1$$

(in real energy space)

Symmetry Transitions from T=1 Nuclei

$$T_z=+1 \rightarrow T_z=0 \leftarrow T_z=-1$$

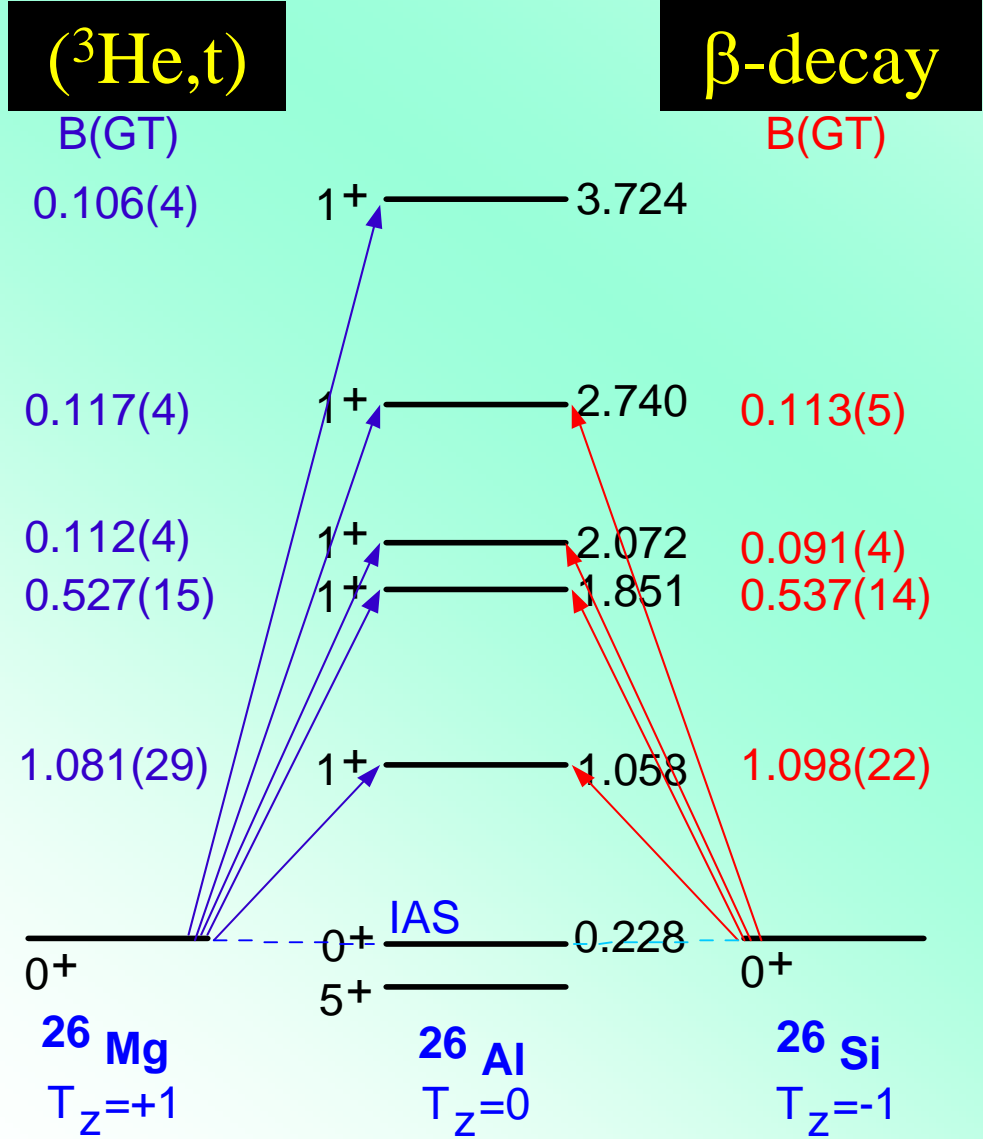
(in isospin symmetry space*)



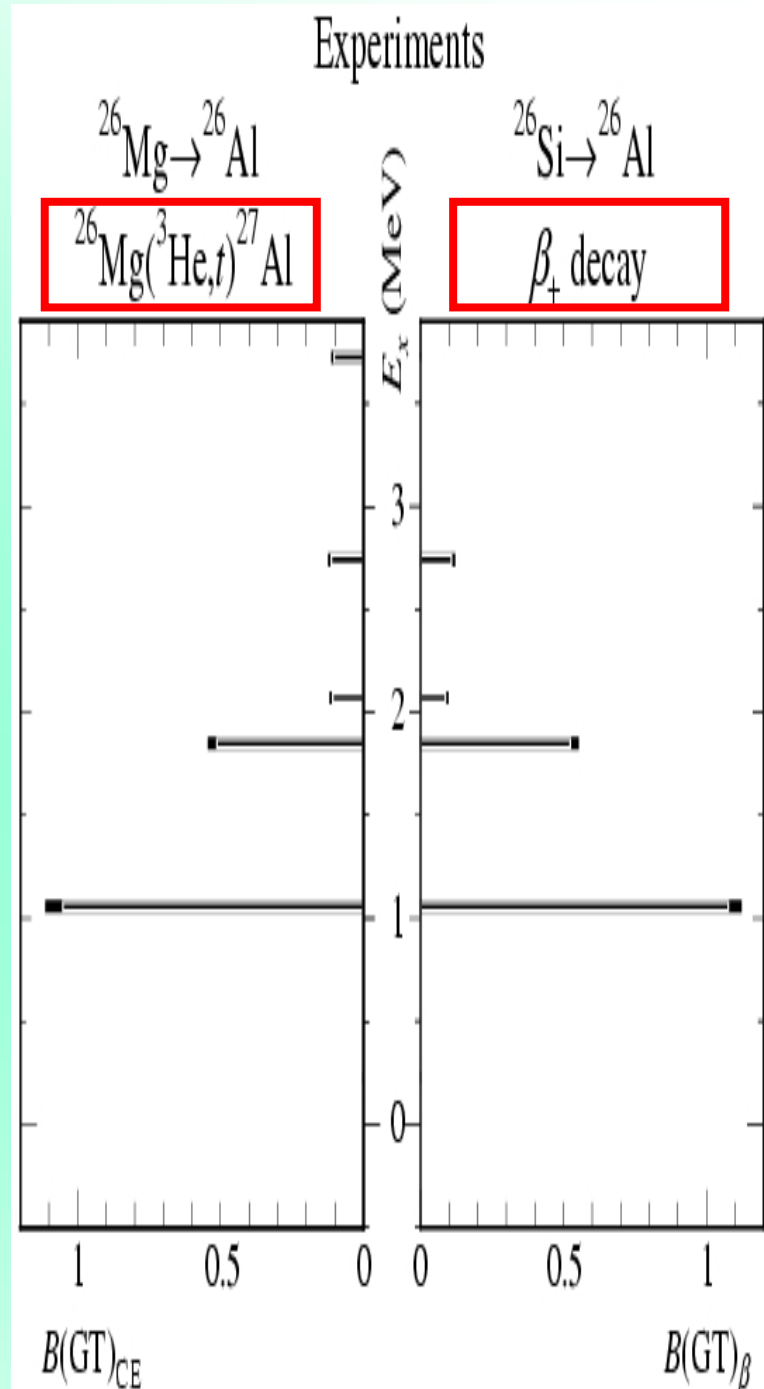
$T_z = +1$

*after the correction of Coulomb displacement energy

B(GT) values from Symmetry Transitions (A=26)

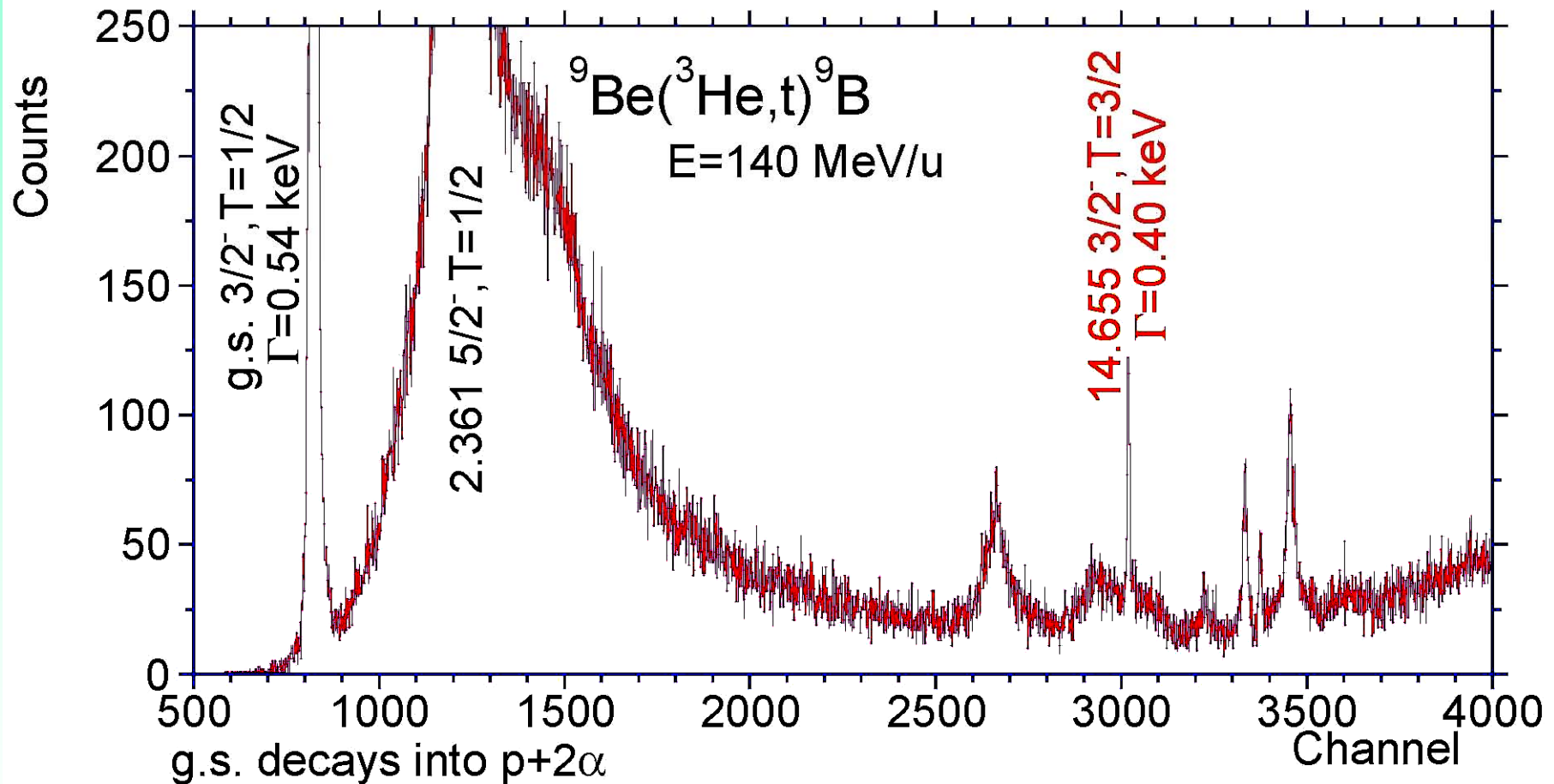


Y. Fujita et al., PRC 67 ('03) 064312



****(³He,t): high resolution and sensitivity !**

${}^9\text{Be}({}^3\text{He},t){}^9\text{B}$ spectrum (II)



**Isospin selection rule prohibits
proton decay of $T=3/2$ state!**

C. Scholl, Koeln

Relationship: Decay and Width

Heisenberg's Uncertainty Principle

$$\Delta x \cdot \Delta p \approx \hbar$$

$$\Delta t \cdot \Delta E \approx \hbar$$

Width $\Gamma = \Delta E$

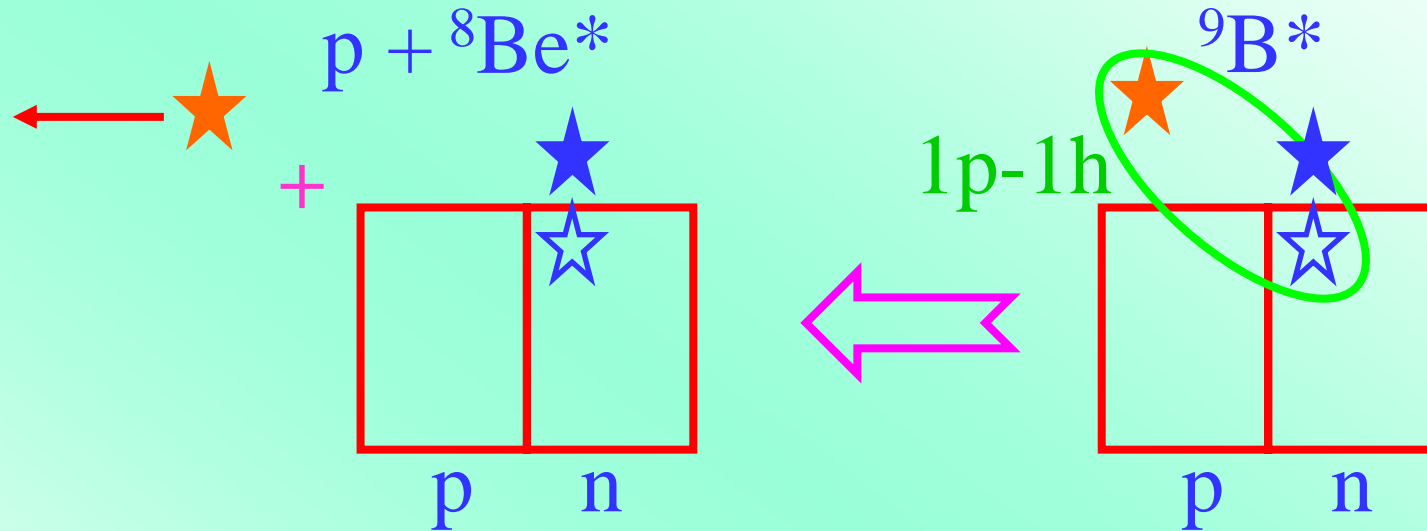
*if: Decay is Fast,

then: Width of a State is Wider !

*if $\Delta t = 10^{-20}$ sec $\rightarrow \Delta E \sim 100$ keV (particle decay)

$\Delta t = 10^{-15}$ sec $\rightarrow \Delta E \sim 1$ eV (fast γ decay)

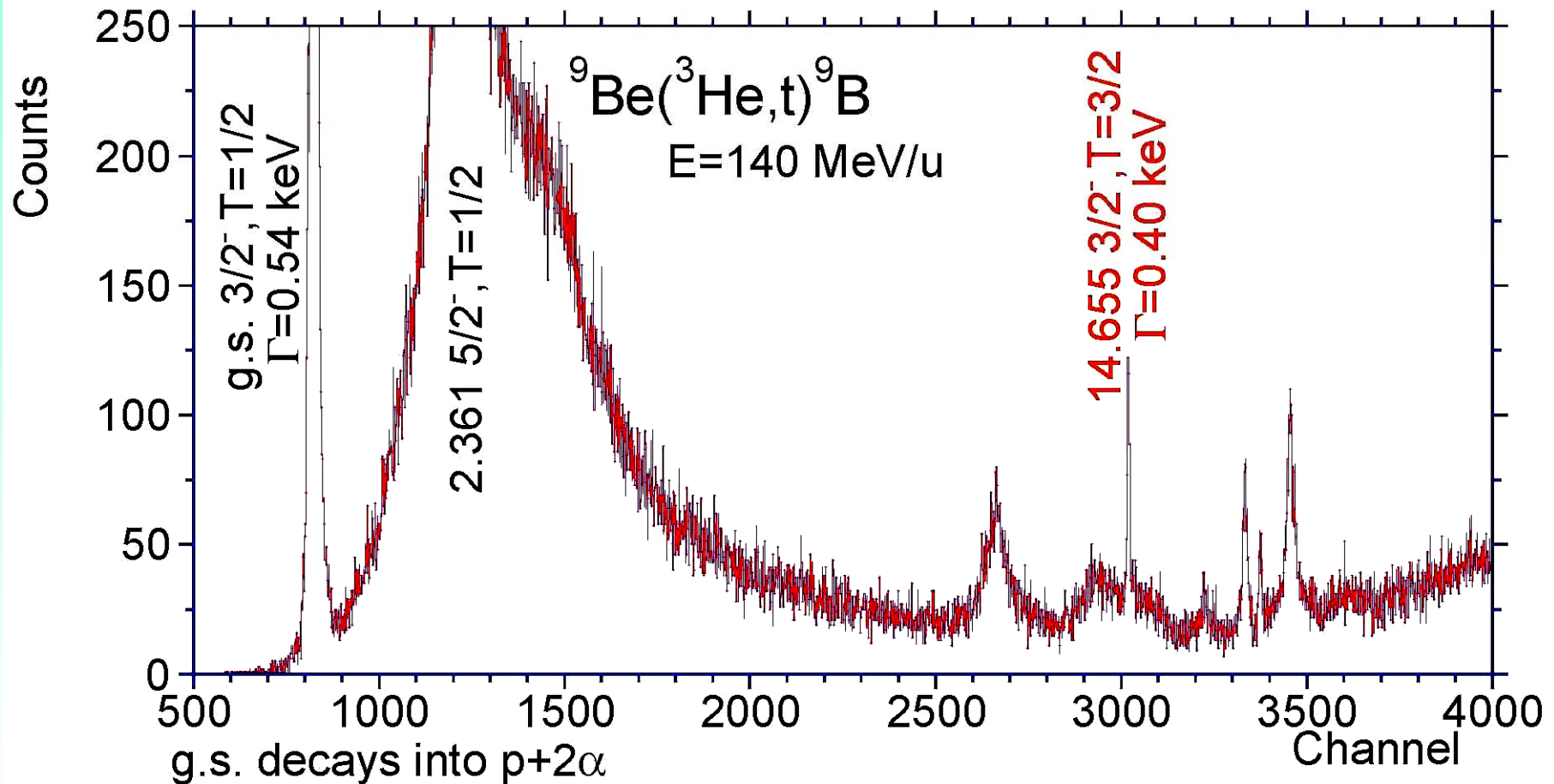
Isospin Selection Rule : in p -decay of ${}^9\text{B}$



$$\begin{array}{rcl}
 \mathbf{T}_z : & -1/2 & + \mathbf{0} & = & -1/2 \\
 \mathbf{T} : & 1/2 & + \mathbf{0} \text{ (low lying)} & = & 1/2 \\
 \mathbf{T} : & 1/2 & + \mathbf{1} \text{ (higher Ex)} & = & 1/2 \ \& \ 3/2
 \end{array}$$

* $T=1$ state in ${}^8\text{Be}$ is
only above
 $E_x=16.6$ MeV

${}^9\text{Be}({}^3\text{He},t){}^9\text{B}$ spectrum (III)

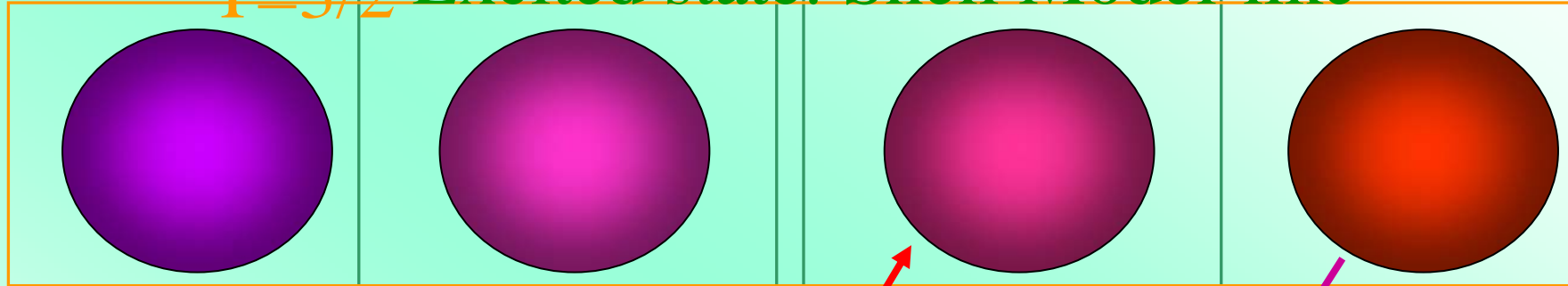


14.7 MeV $T=3/2$ state is very weak!

Strength ratio of g.s. & 14.7 MeV $3/2^-$ states: **140:1**

Shell Structure and Cluster Structure

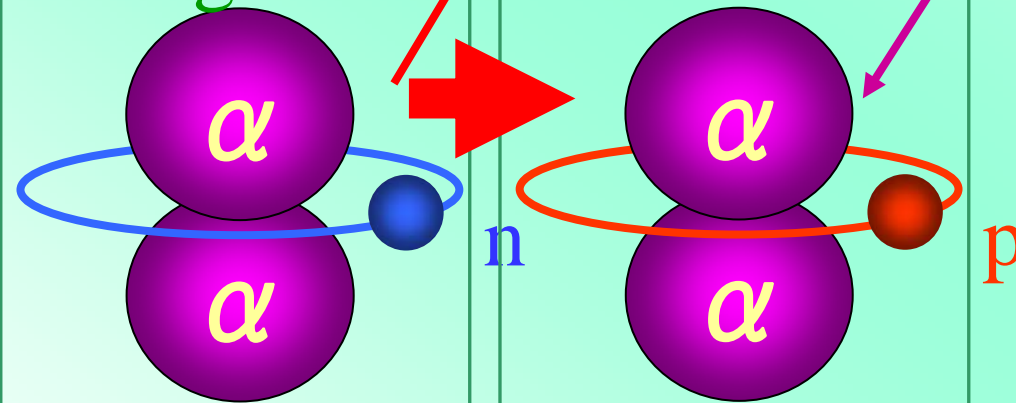
$T=3/2$ Excited state: Shell Model-like



${}^9\text{Li}$
 $T_z=3/2$
 neutron: $p_{3/2}$ closed

${}^9\text{C}$
 $T_z=-3/2$
 proton: $p_{3/2}$ closed

g.s.: Cluster-like



${}^9\text{Be}$
 $T_z=1/2$

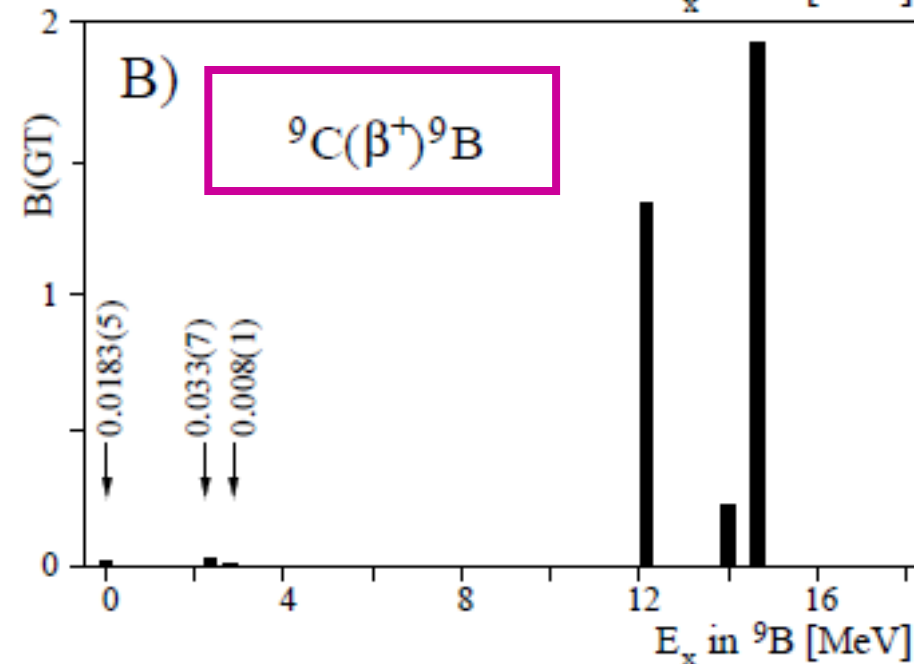
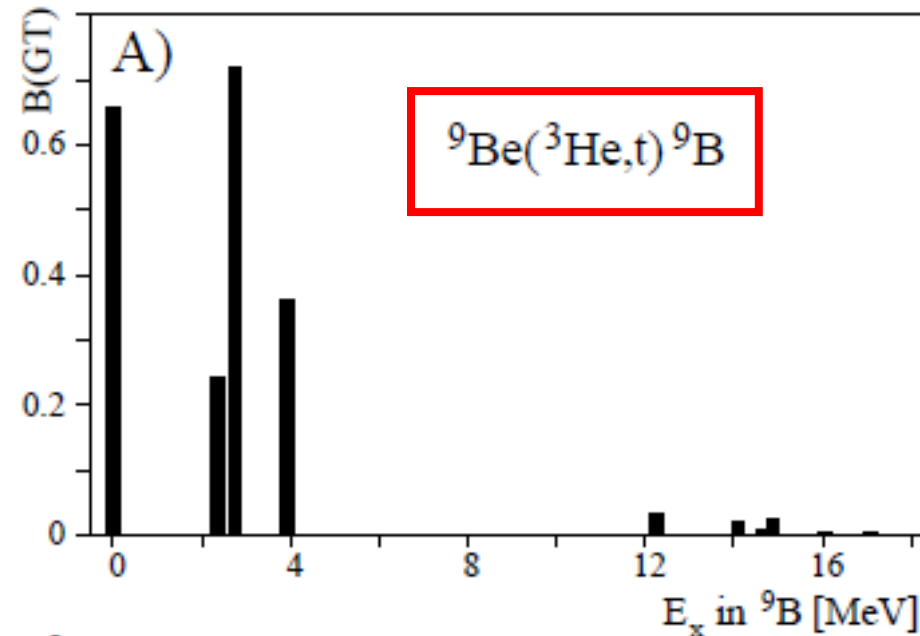
${}^9\text{B}$
 $T_z=-1/2$

suggestion by
 Y. Kanada-En'yo

β -decay and ($^3\text{He}, t$) results

C. Scholl et al,
PRC 84, 014308 (2011)

L. Buchmann et al.,
PRC 63 (2001) 034303.
U.C. Bergmann et al.,
Nucl. Phys. A 692 (2001) 427.



High-resolution Experiment

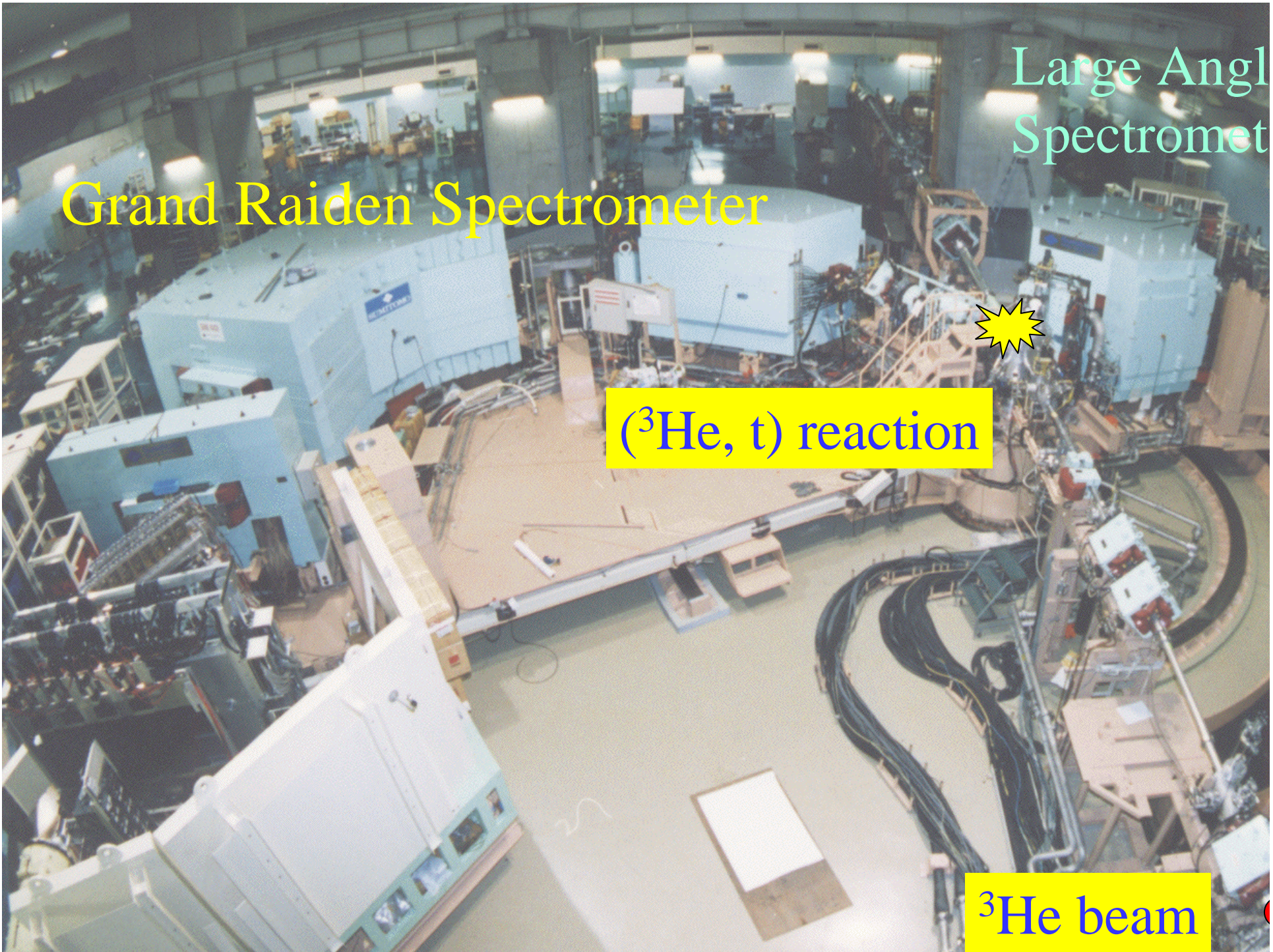
-beam matching techniques-
(dispersion matching techniques)

Large Angle
Spectrometer

Grand Raiden Spectrometer

$(^3\text{He}, t)$ reaction

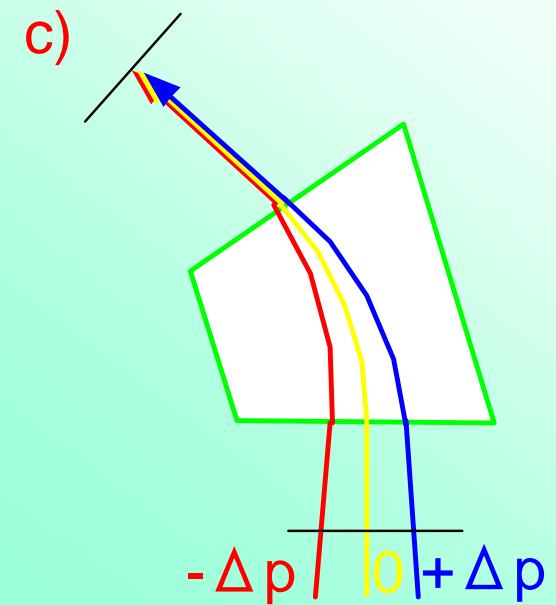
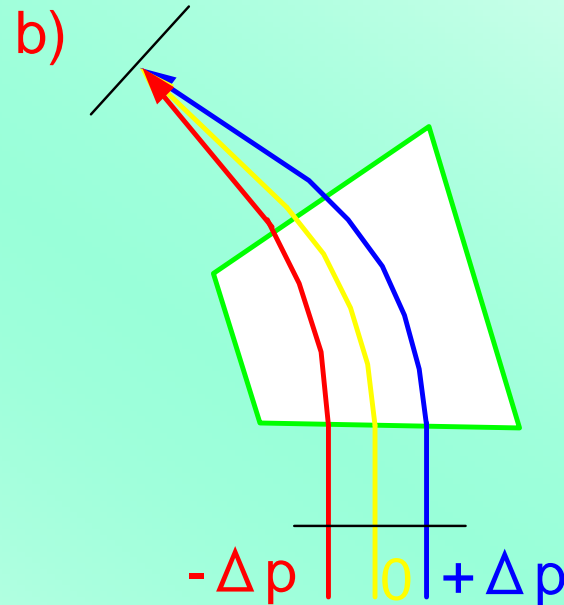
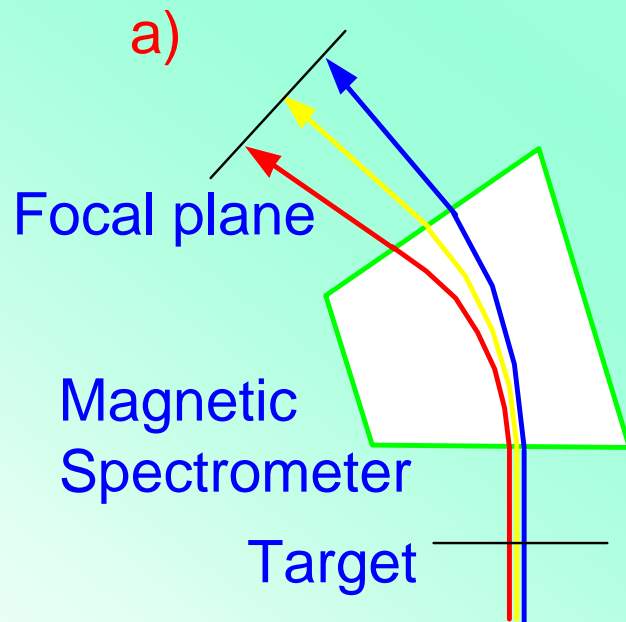
^3He beam



Matching Techniques

Y. Fujita et al., N.I.M. B 126 (1997) 274.

H. Fujita et al., N.I.M. A 484 (2002) 17.



*Achromatic beam
transportation*

$\Delta E \sim 200$ keV
for 140MeV/u³He beam

*Lateral dispersion
matching*

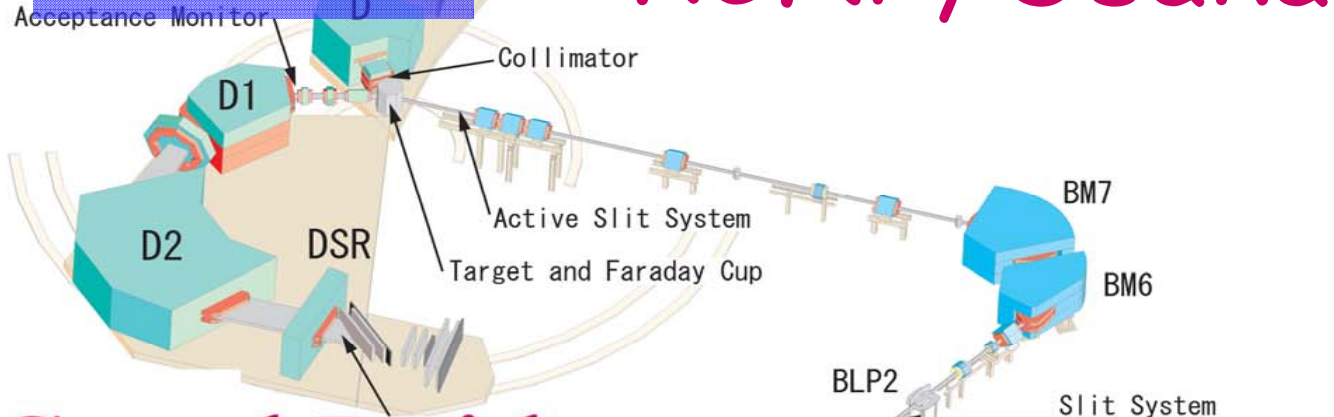
$\Delta E \sim 35$ keV
Horiz. angle resolution
 $\Delta\theta_{sc} > 15$ mrad

*Angular dispersion
matching*

$\Delta\theta_{sc} \sim 5$ mrad

$\Delta E = 30 \text{ keV}$

RCNP, Osaka Univ.



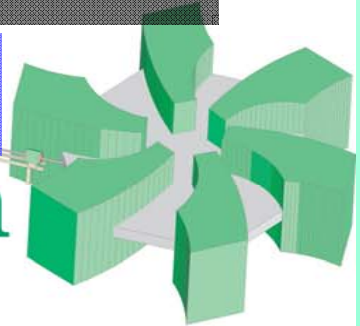
Grand Raiden

WS Beam Line

Dispersion Matching Techniques were applied!

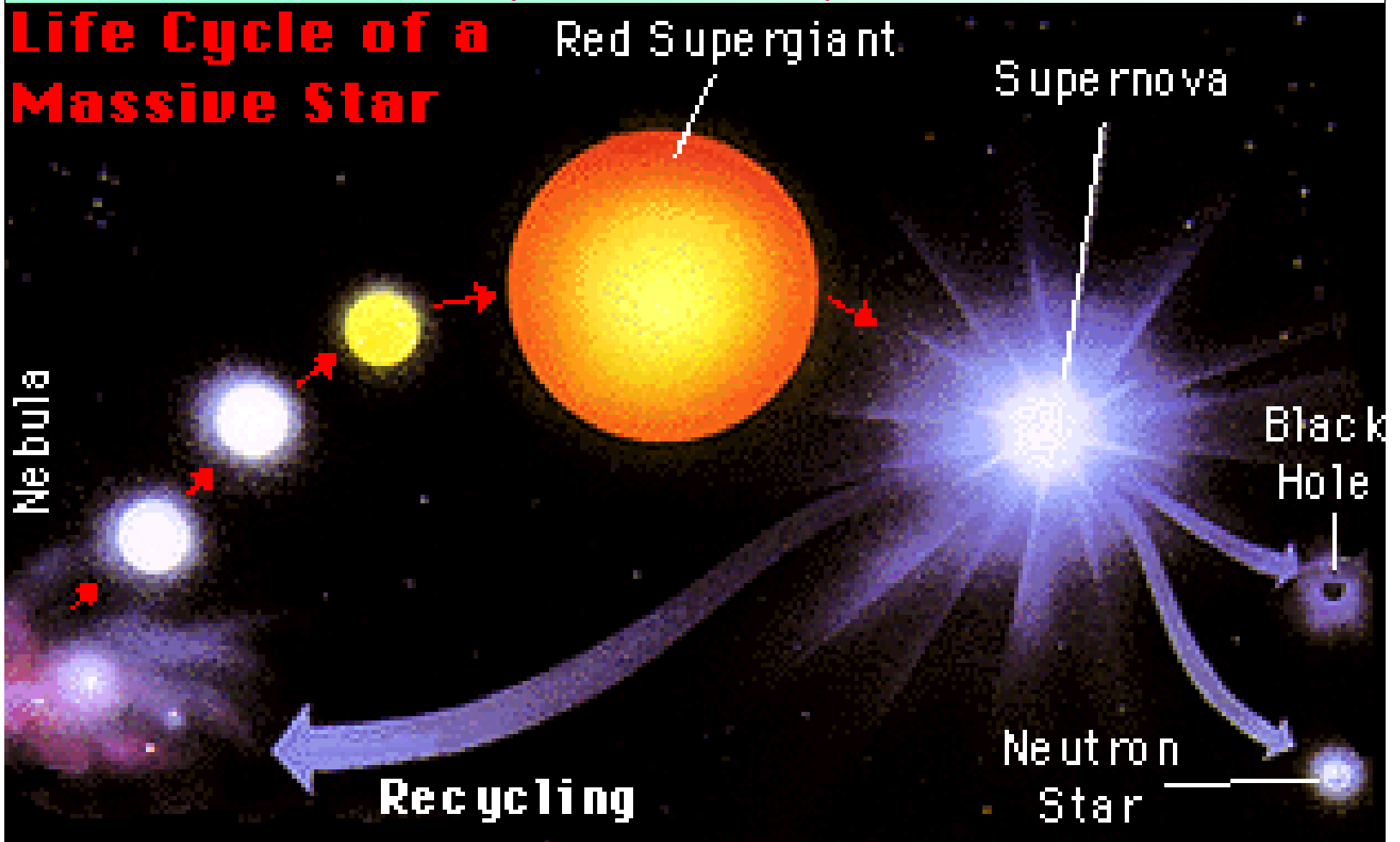
$\Delta E = 150 \text{ keV}$

Ring Cyclotron



Supernova Cycle

Life Cycle of a Massive Star



Crucial Weak Processes during the Core Collapse

$\sigma\tau$: important

(A,Z) =nuclei in the Cr, Mn, Fe, Co, Ni region
 pf -shell Nuclei !

Langanke & Martinez-Pinedo
 Rev.Mod.Phys.75('04)819

Balantekin & Fuller
 J.Phys.G 29('03)2513

$$p + e^- \rightleftharpoons n + \nu_e,$$

$$n + e^+ \rightleftharpoons p + \bar{\nu}_e,$$

→ $(A, Z) + e^- \rightleftharpoons (A, Z-1) + \nu_e,$

→ $(A, Z) + e^+ \rightleftharpoons (A, Z+1) + \bar{\nu}_e,$

$$\nu + N \rightleftharpoons \nu + N,$$

$$N + N \rightleftharpoons N + N + \nu + \bar{\nu},$$

$$\nu + (A, Z) \rightleftharpoons \nu + (A, Z),$$

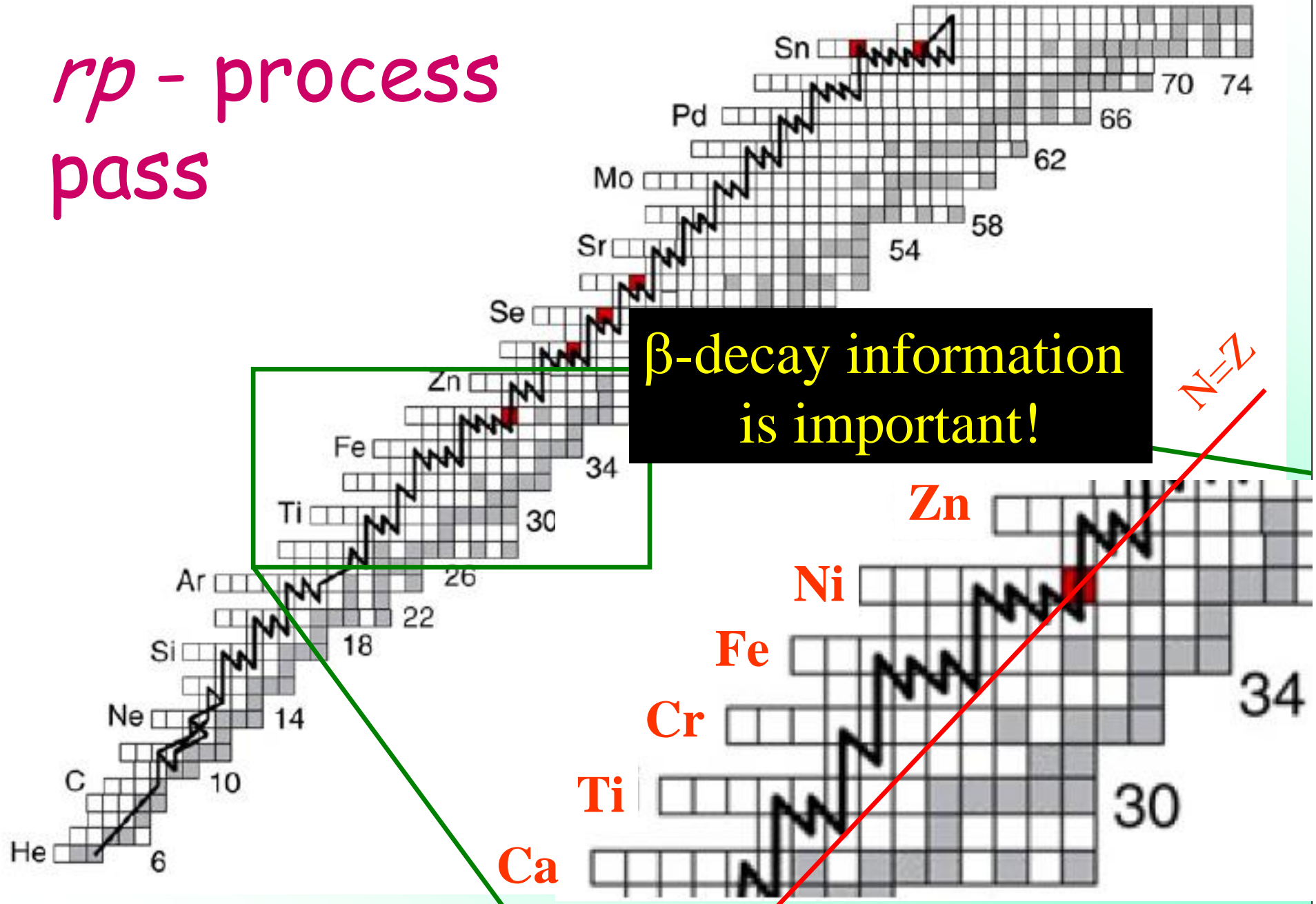
$$\nu + e^\pm \rightleftharpoons \nu + e^\pm,$$

$$\nu + (A, Z) \rightleftharpoons \nu + (A, Z)^*,$$

$$e^+ + e^- \rightleftharpoons \nu + \bar{\nu},$$

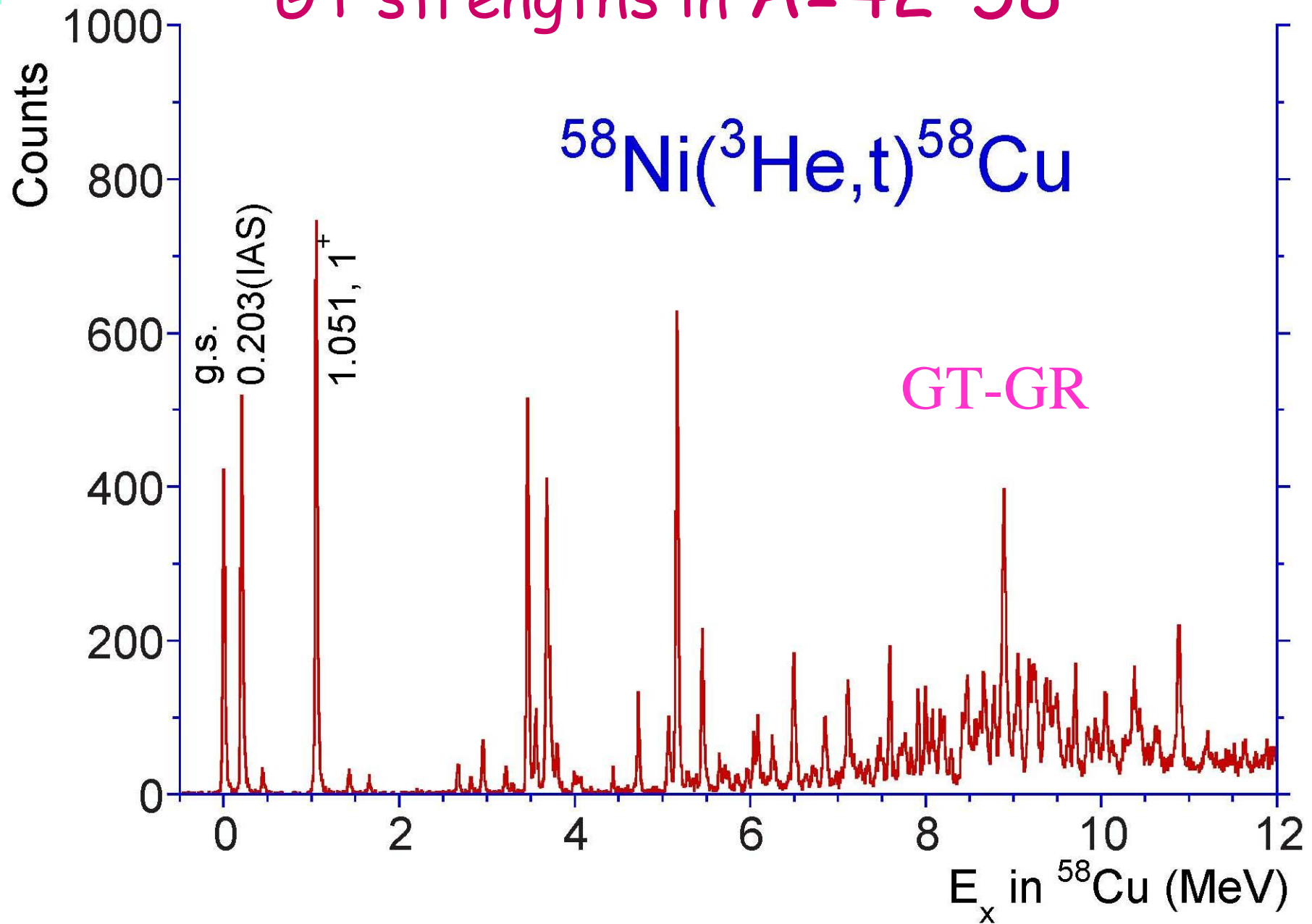
$$(A, Z)^* \rightleftharpoons (A, Z) + \nu + \bar{\nu}.$$

rp - process pass



GT strengths in $A=42-58$

$^{58}\text{Ni}({}^3\text{He},t){}^{58}\text{Cu}$



***Exotic GT transitions from Unstable Nuclei

- Combined ($^3\text{He},t$) and β -decay Study -

**Derivation of "absolute" $B(GT)$ values

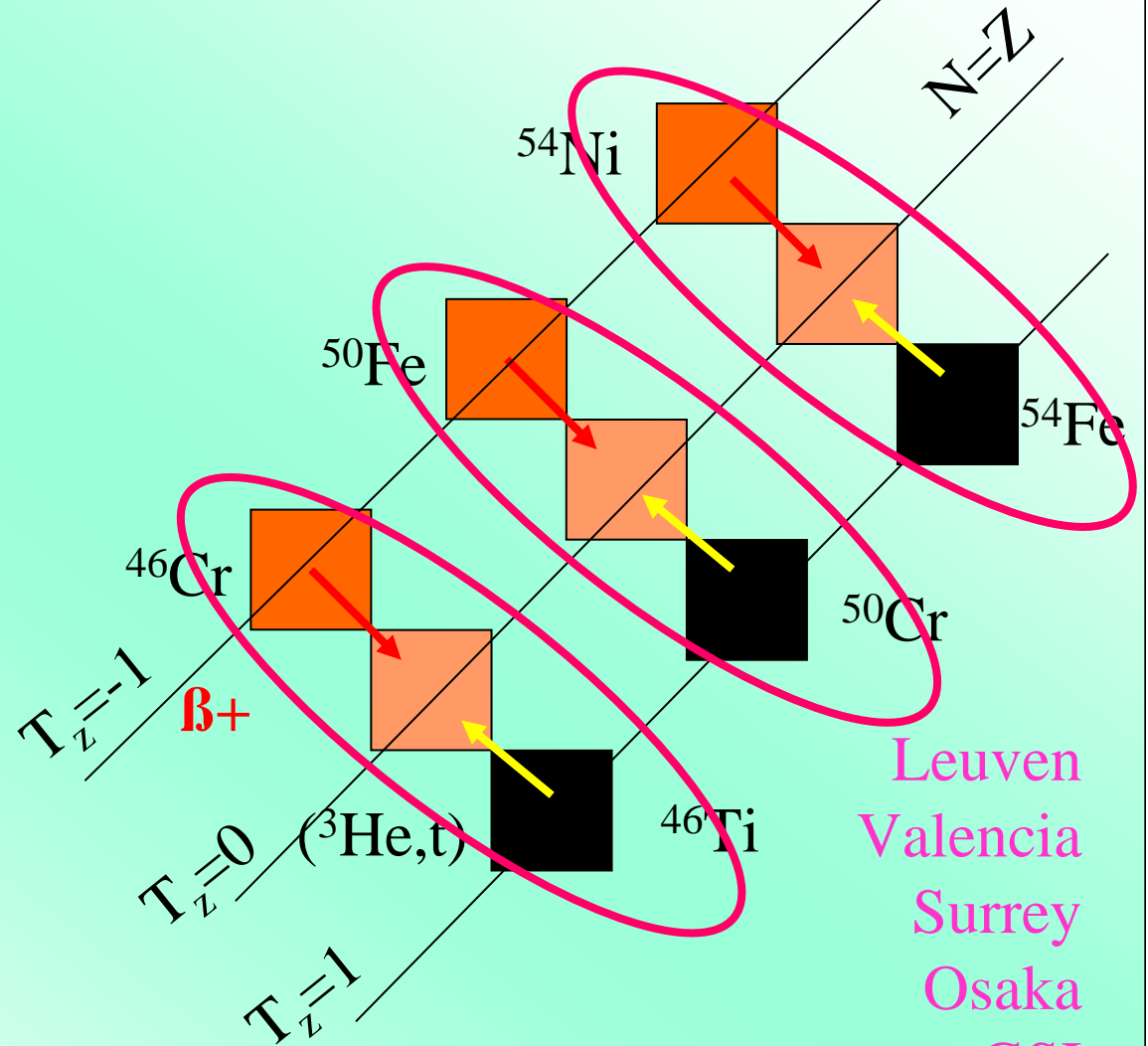
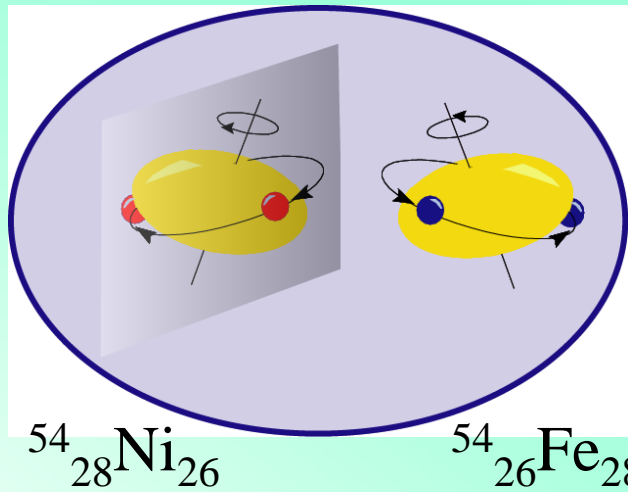
* β -decay: $T_{1/2}$ and absolute $B(GT)$ values
but only for the low-lying states

* $(^3\text{He},t)$ reaction: highly-excited states can be
accessed but only the relative $B(GT)$ values

Let's combine CE reaction and β -decay data
starting from mirror nuclei !

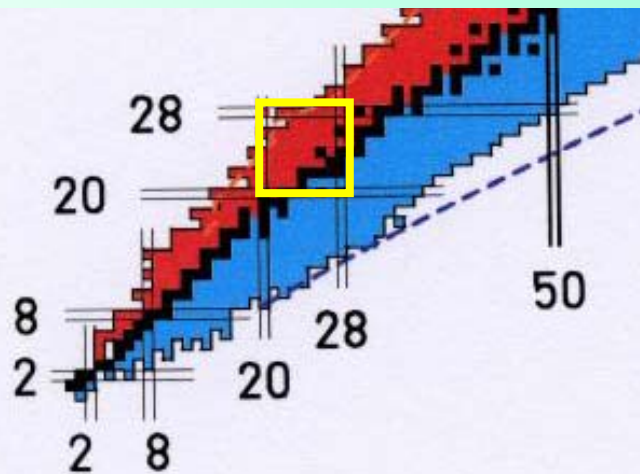
T=1 Isospin Symmetry in *pf*-shell Nuclei

Mirror nuclei



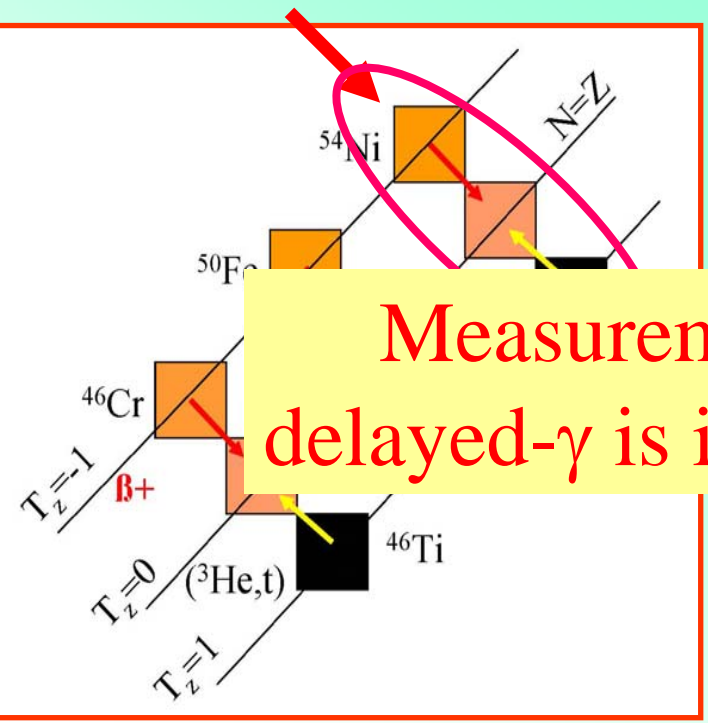
Leuven
Valencia
Surrey
Osaka
GSI
CNS

by B. Rubio

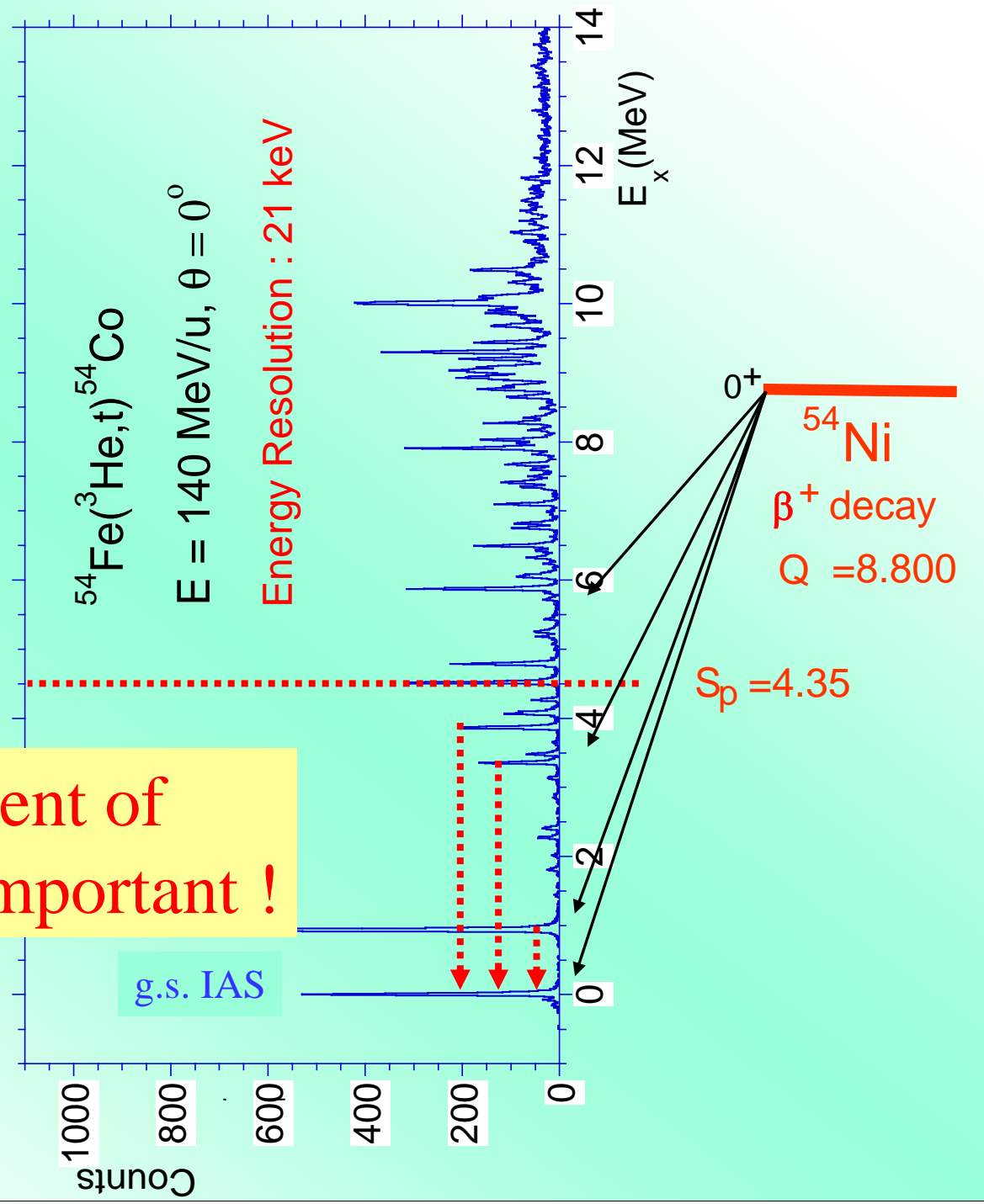


^{54}Ni β -decay measurement

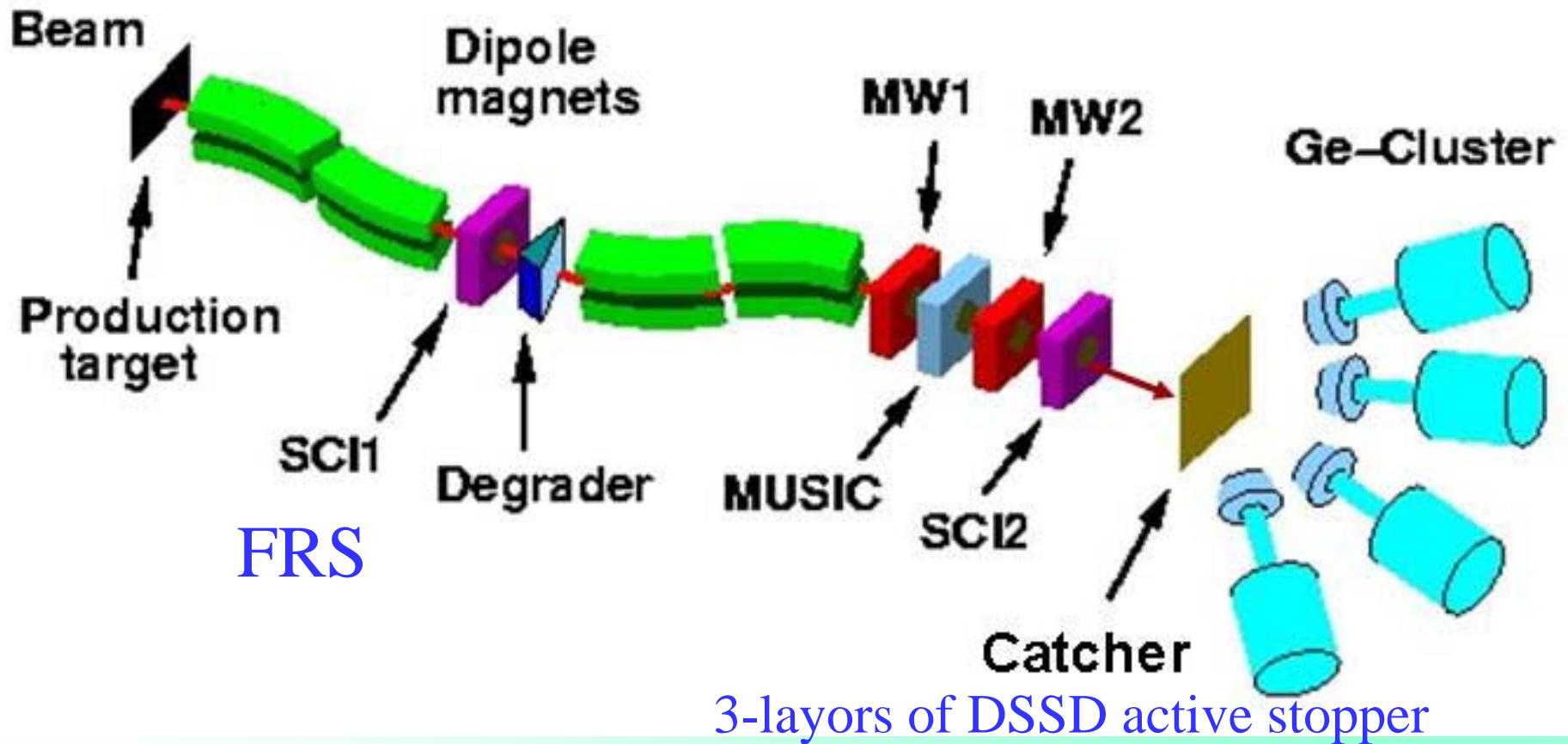
- at GSI (FRS facility)
- RISING (stopped beam campaign)



Measurement of delayed- γ is important !



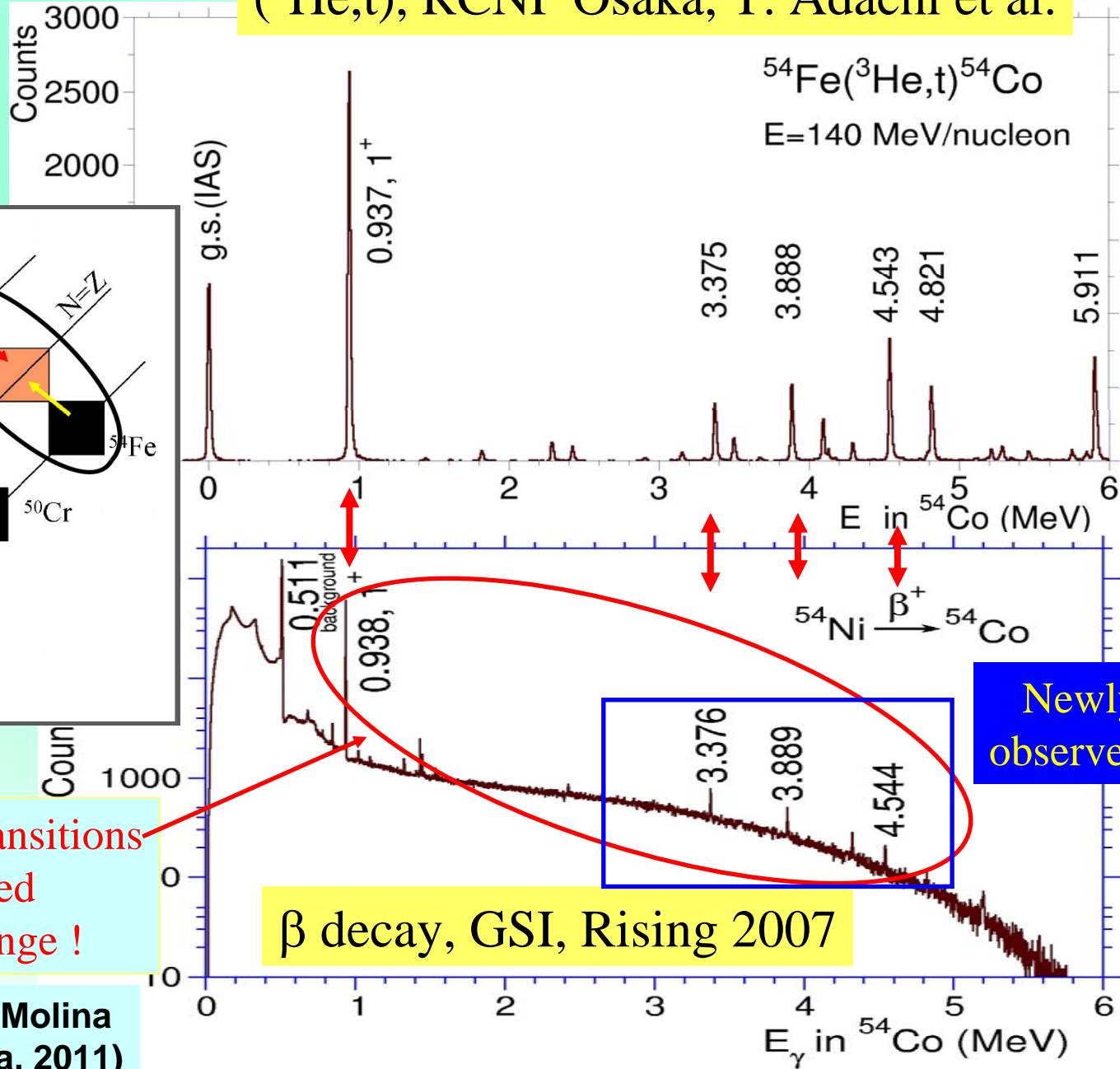
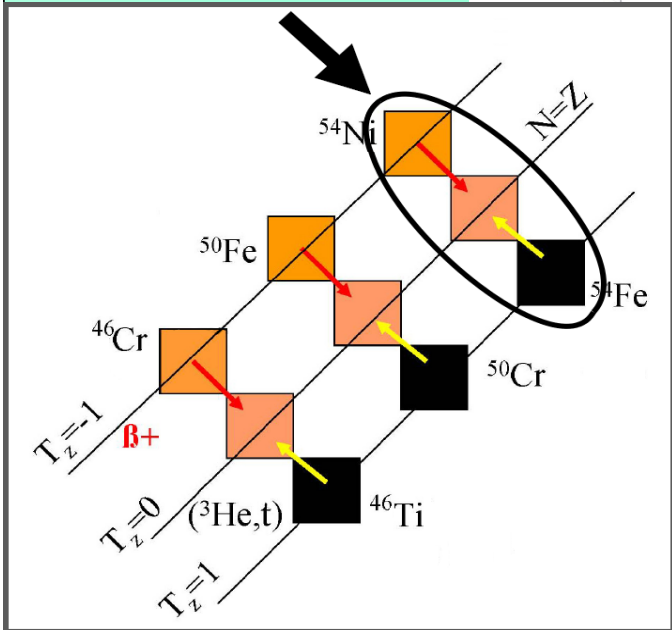
GSI: RISING set up - active stopper campaign -



FRS

3-layers of DSSD active stopper

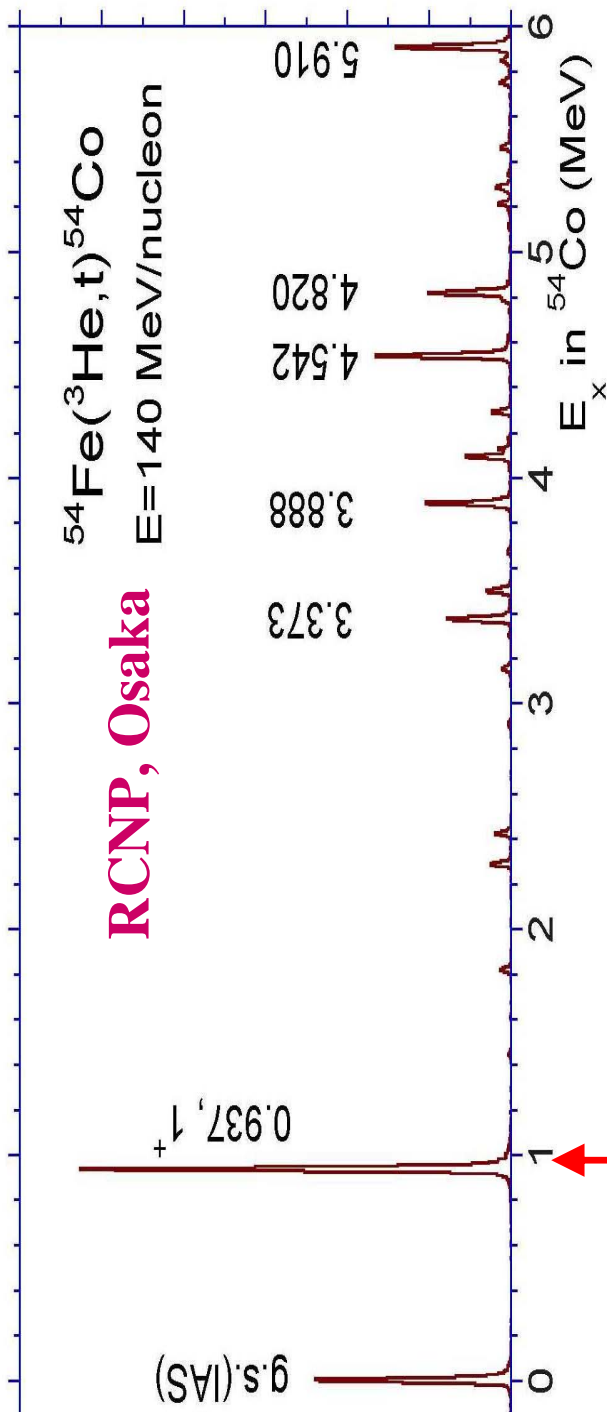
$(^3\text{He},t)$, RCNP Osaka, T. Adachi et al.



Corresponding Transitions were observed in a wide E_x range !

β decay, GSI, Rising 2007

Ph.D. F. Molina (Valencia, 2011)



Comparison

B(GT) $({}^3\text{He}, t)$

B(GT) β -decay

0.46

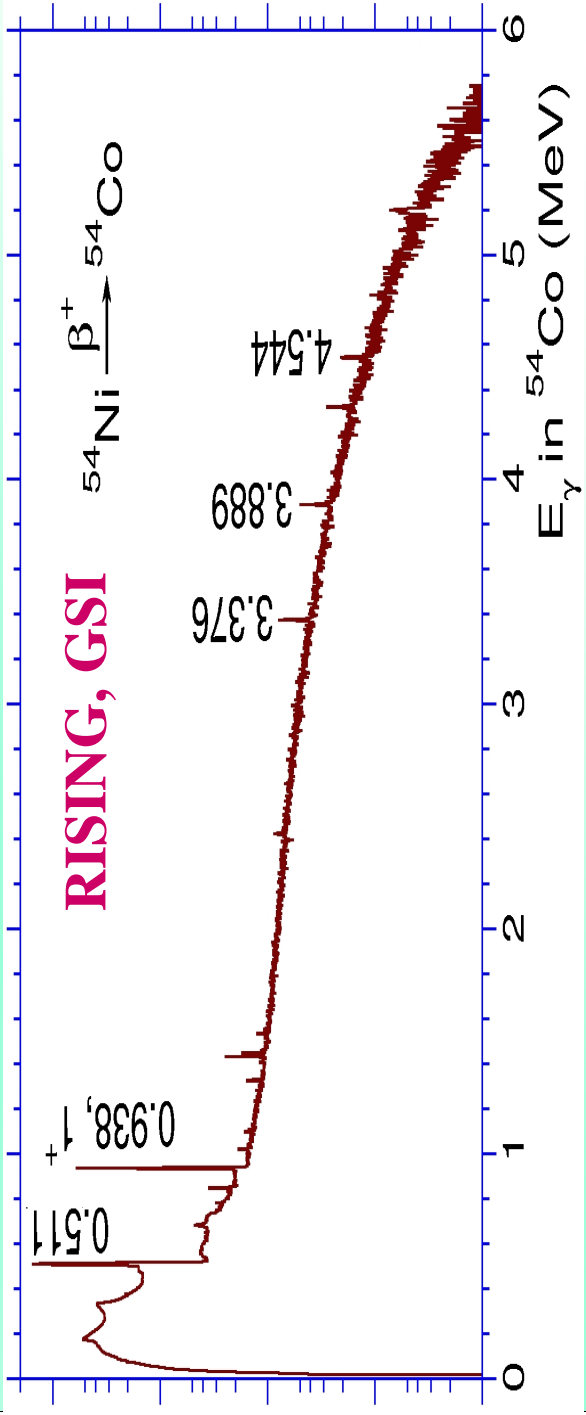
0.48

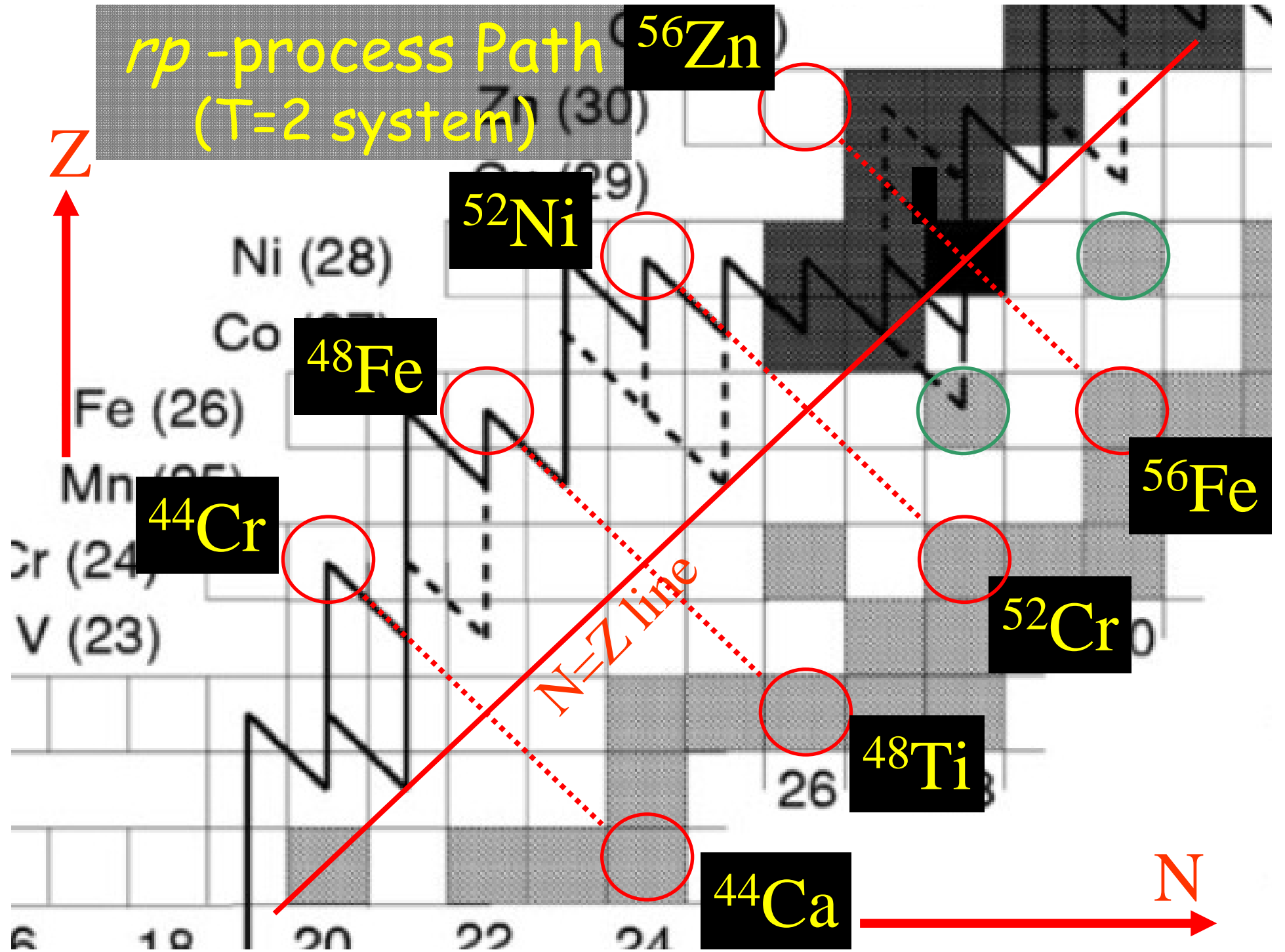
70.0

70.0

60.0

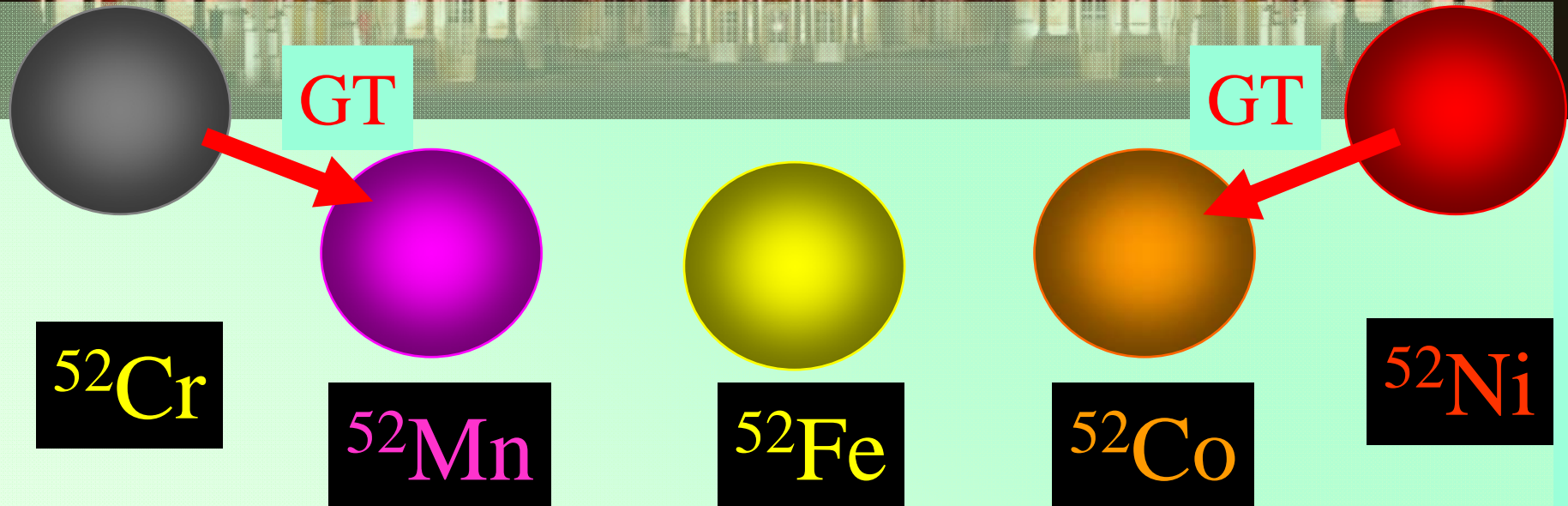
70.0



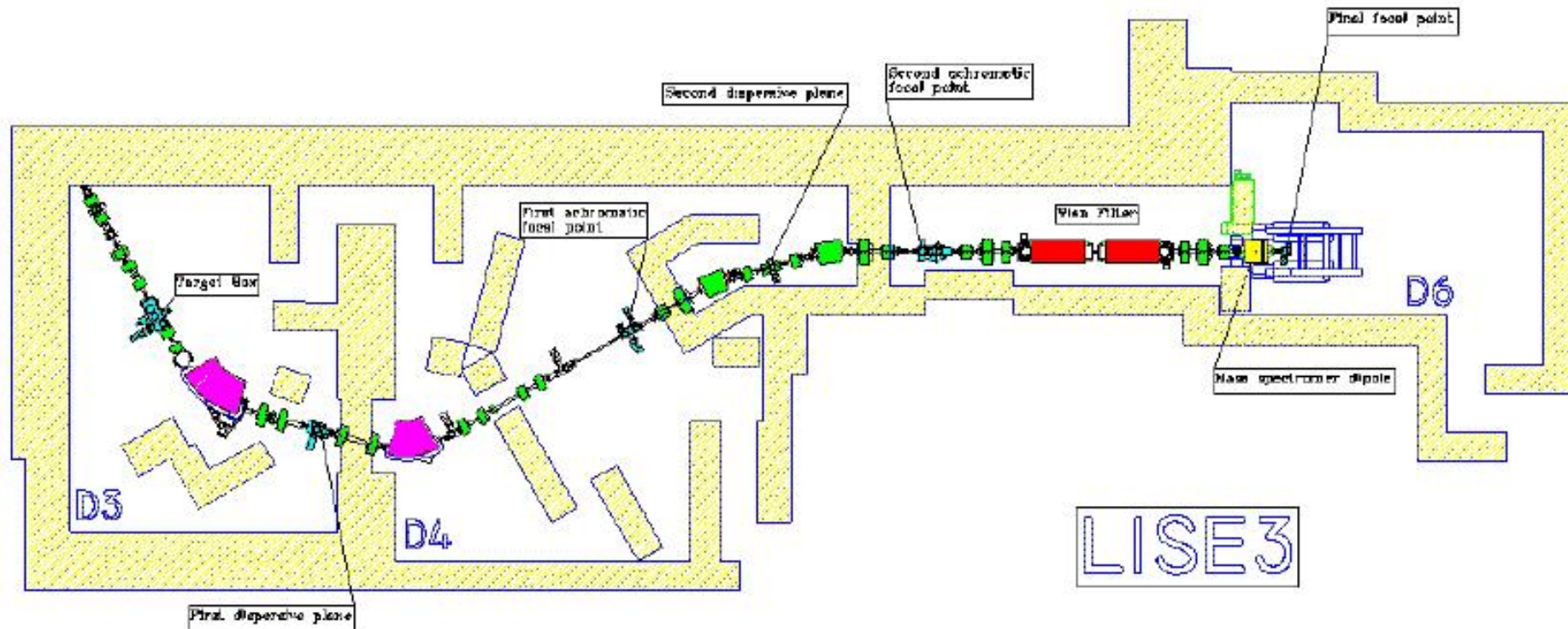


Chateau de Versailles

T=2 Isospin Symmetry



GANIL LISE3 fragment separator

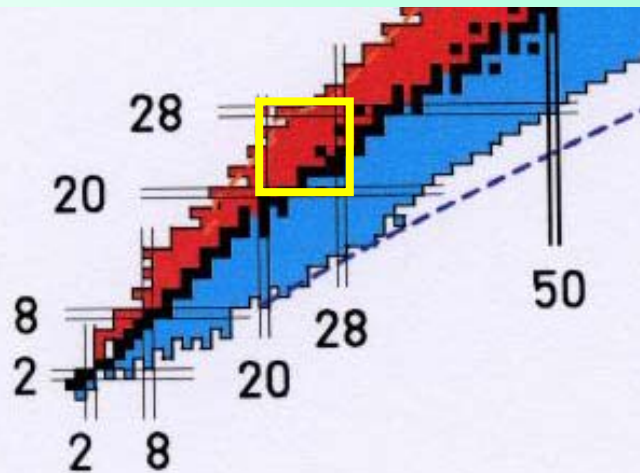
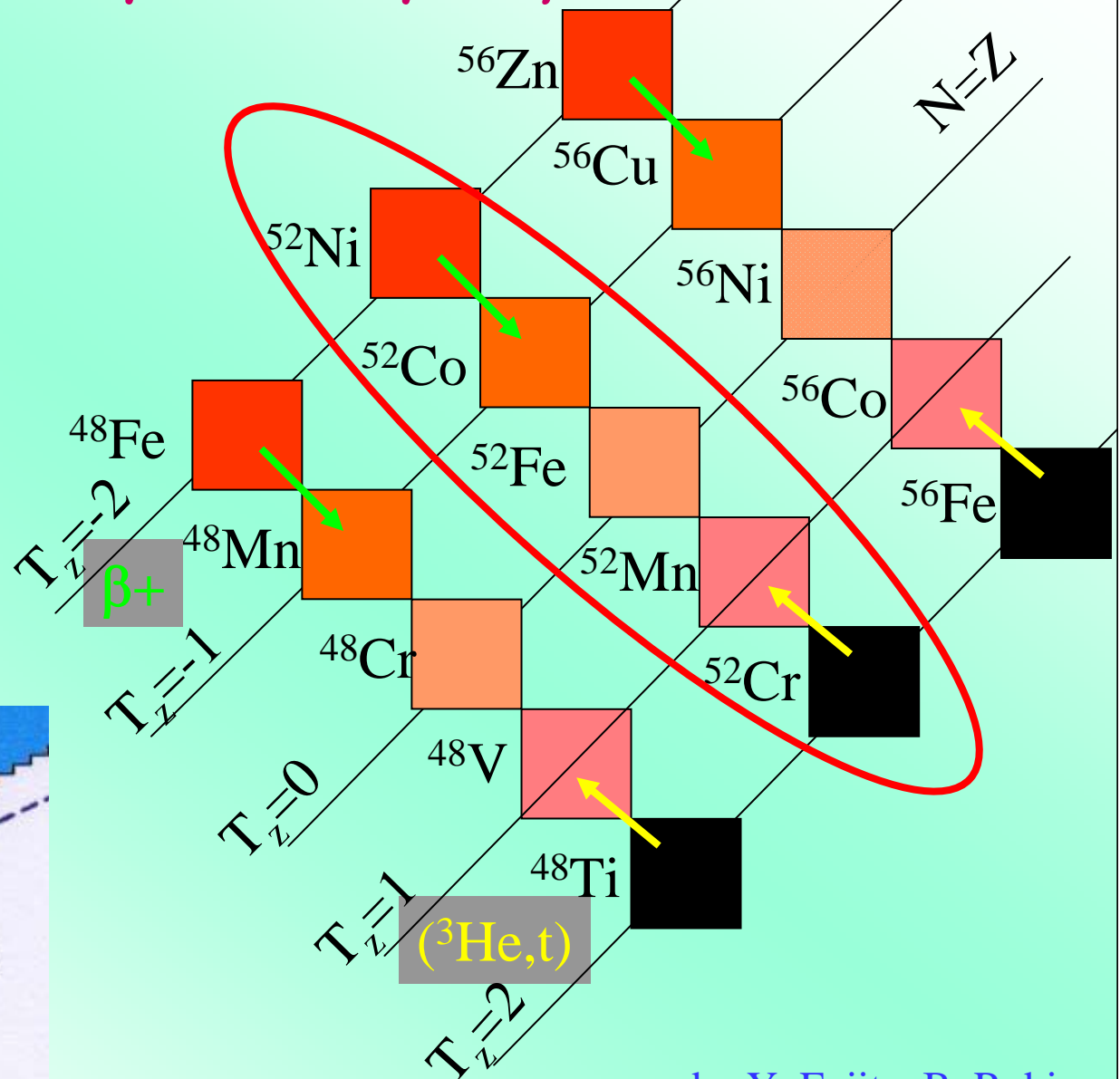
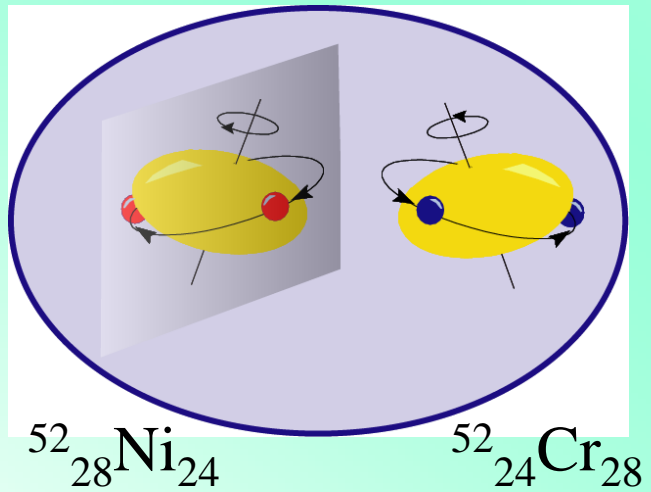


^{58}Ni beam: $\sim 79\text{MeV/u}$, $3.5\text{ e}\mu\text{A}$, production target: Ni

p-decay: by DSSD, γ -decay: by Ge detectors

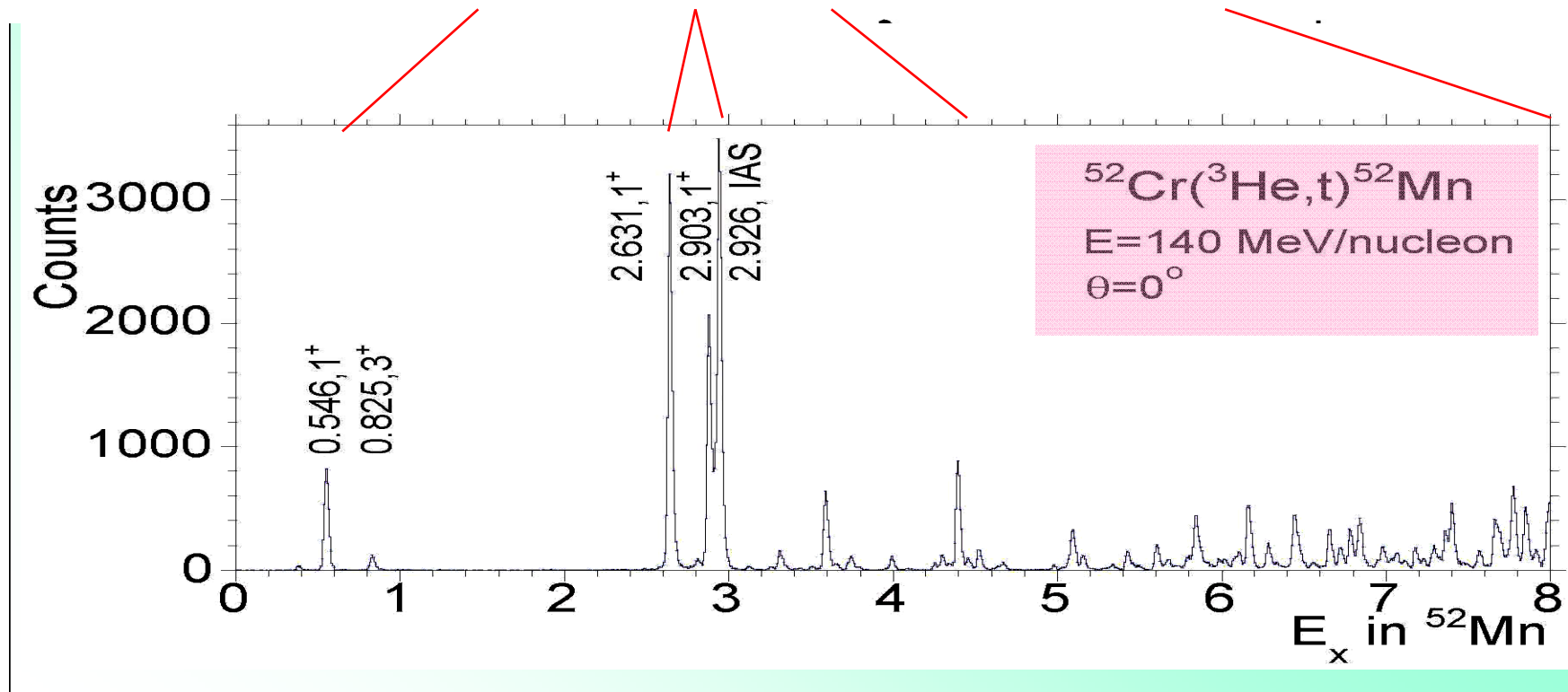
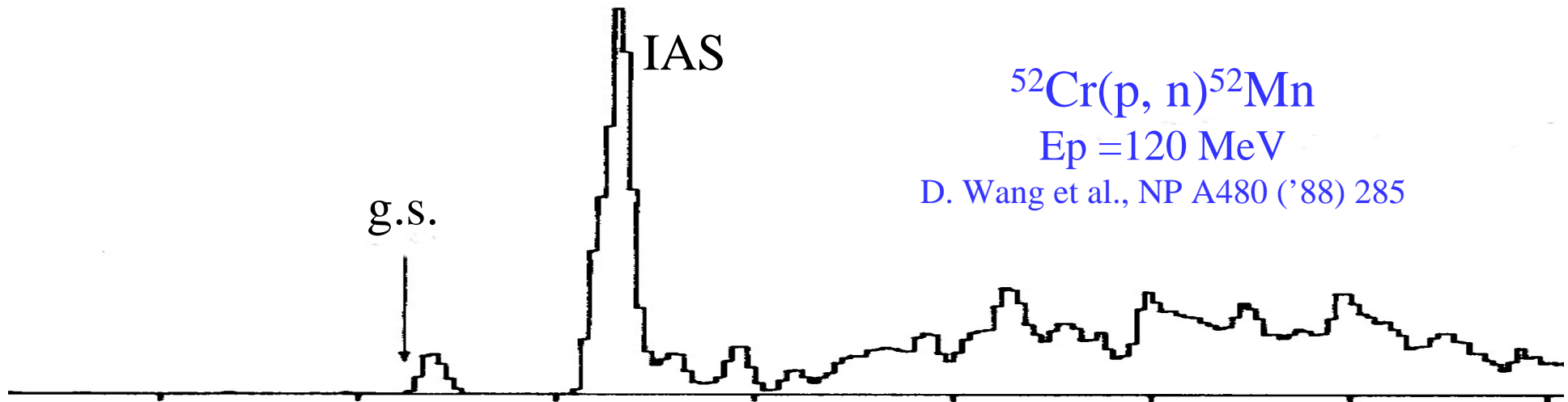
T = 2 Isospin Symmetry in *pf*-shell Nuclei

Mirror nuclei

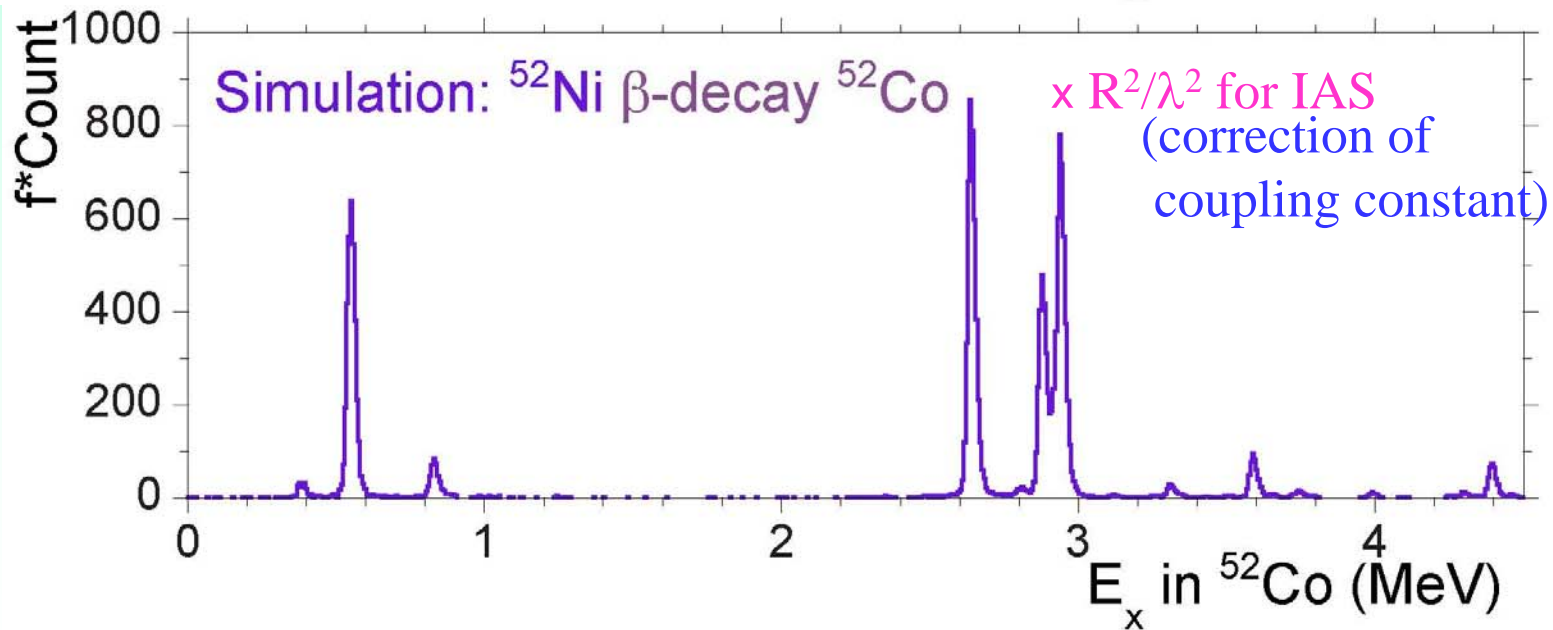
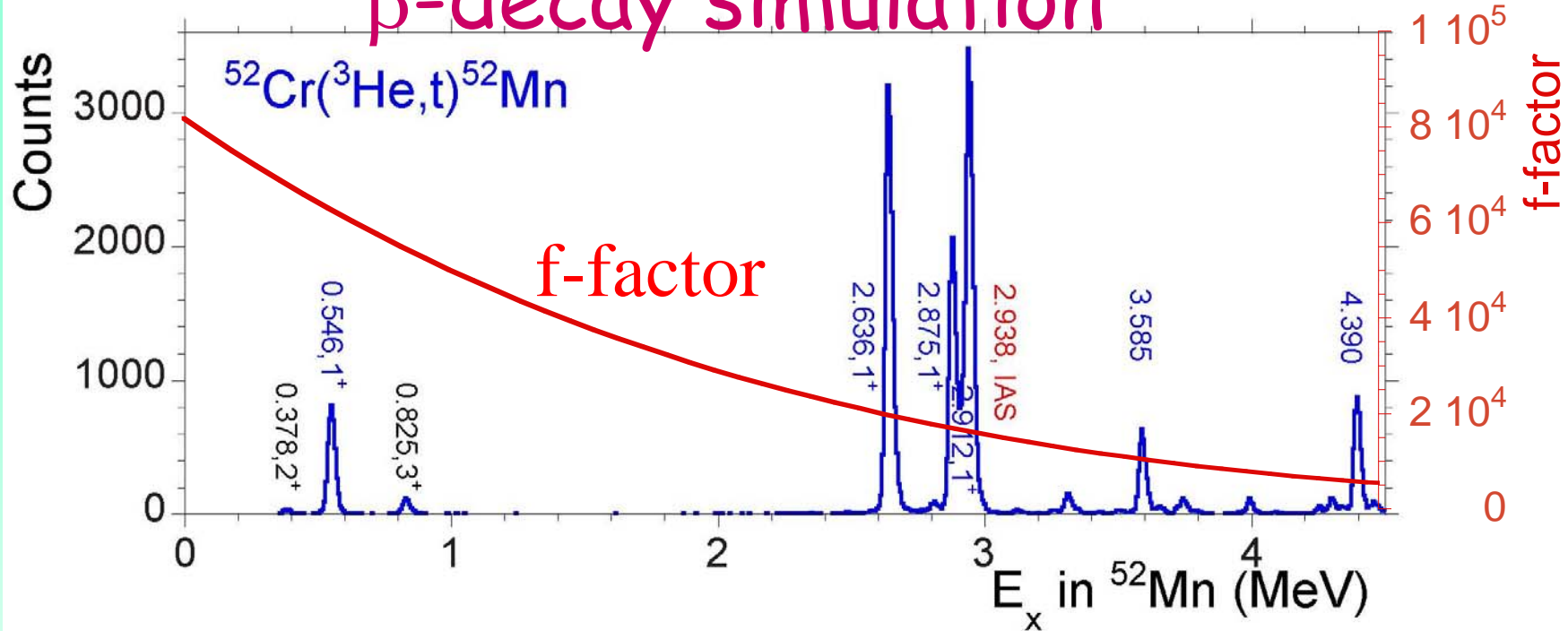


by Y. Fujita, B. Rubio

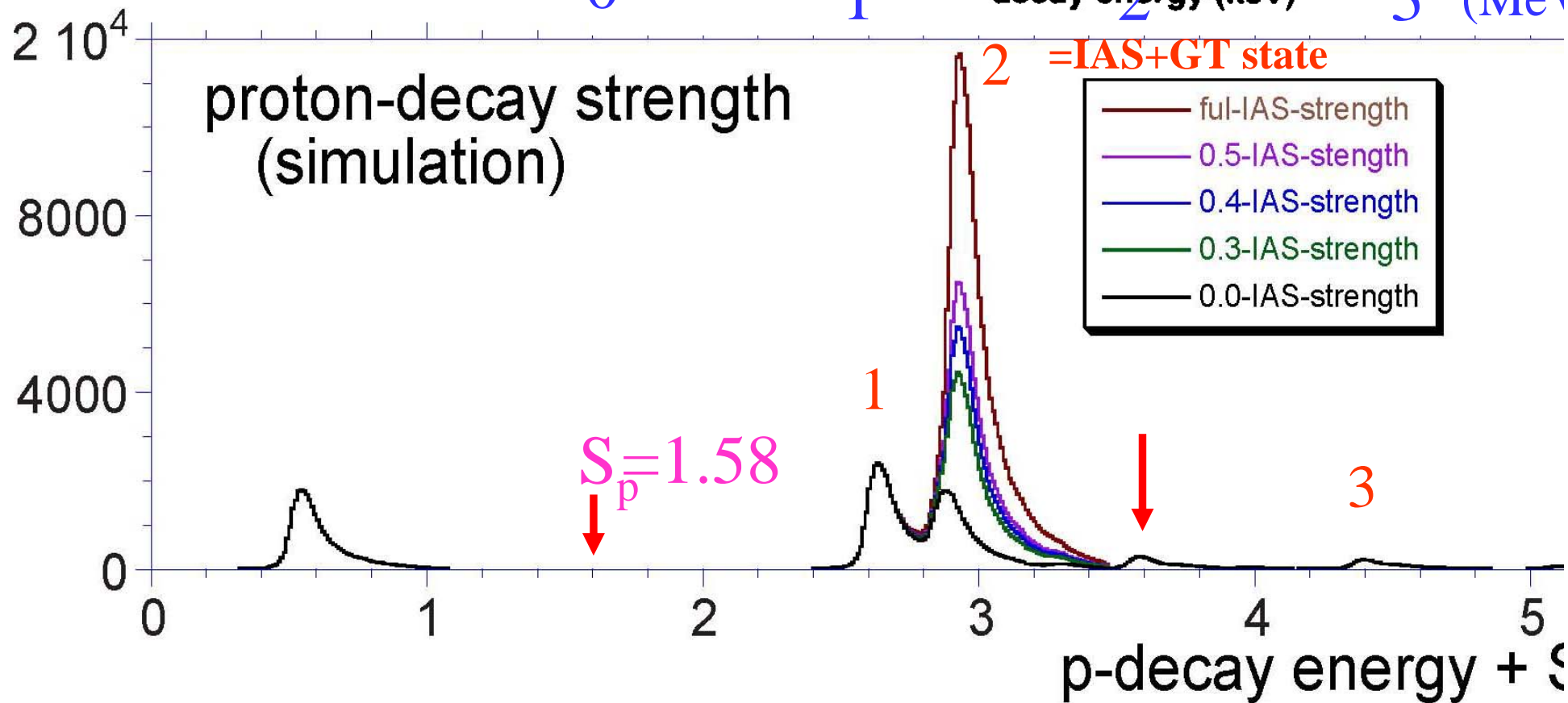
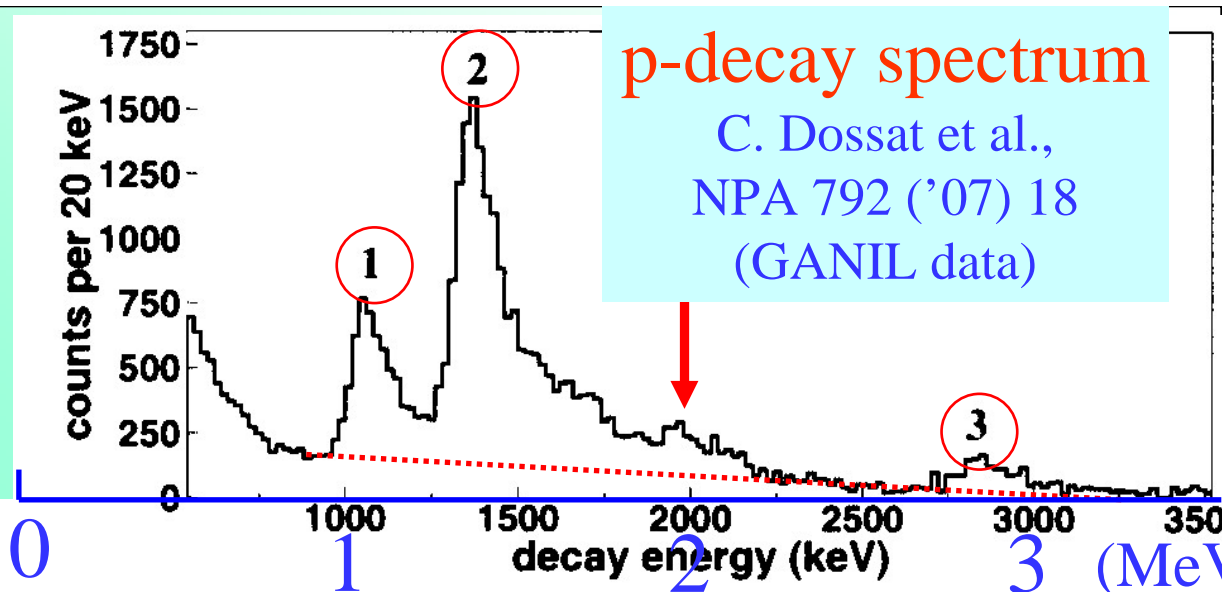
Comparison: (p, n) and ($^3\text{He}, t$)



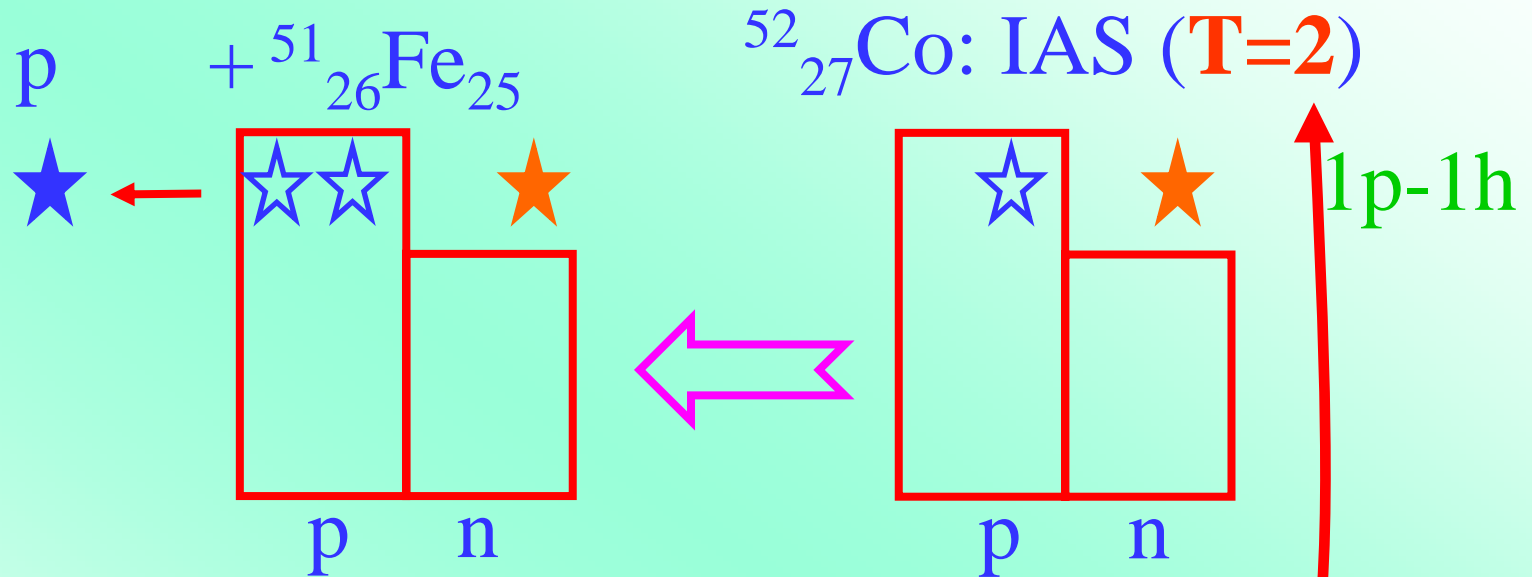
β -decay simulation



p-decay
of ^{52}Co



Isospin Selection Rule : *p*-decay of IAS in ^{52}Co



$$T_z : -1/2 + (-1/2) = -1$$

$$T : 1/2 + 1/2 \text{ (at low lying)} = 1$$

Therefore, delayed *p*-decay of the IAS is forbidden!

Summary

GT ($\sigma\tau$) operator : a simple operator !

- * GT transitions: sensitive to the structure of $|i\rangle$ and $|f\rangle$
- * Isospin quantum number T plays an important role

High resolution of the ($^3\text{He},t$) reaction

- * Width & fine structures of GT transitions
- * Precise comparison with mirror β -decay results

- GT transitions in each nucleus are UNIQUE !
- Vital part of the Nuclear Structure is revealed !

**GT transitions:
- transitions with full of personality -**

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Review

Spin–isospin excitations probed by strong, weak and electro-magnetic interactions

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(2011) 549-606

GT-study Collaborations

Bordeaux (France) : β decay

GANIL (France) : β decay

Gent (Belgium) : (^3He , t), (d, ^2He), (γ , γ'), theory

GSI, Darmstadt (Germany) : β decay, theory

ISOLDE, CERN (Switzerland) : β decay

iThemba LABS. (South Africa) : (p, p'), (^3He , t)

Istanbul (Turkey): (^3He , t), β decay

Jyvaskyla (Finland) : β decay

Koeln (Germany) : γ decay, (^3He , t), theory

KVI, Groningen (The Netherlands) : (d, ^2He)

Leuven (Belgium) : β decay

LTH, Lund (Sweden) : theory

Osaka University (Japan) : (p, p'), (^3He , t), theory

Surrey (GB) : β decay

TU Darmstadt (Germany) : (e, e'), (^3He , t)

Valencia (Spain) : β decay

Michigan State University (USA) : theory, (t, ^3He)

Muenster (Germany) : (d, ^2He), (^3He ,t)

Univ. Tokyo and CNS (Japan) : theory, β decay