

# Recent results from Coulomb Dissociation (CD) experiments

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Sep. 2011

Santa Tecla



# Past and future of CD experiments; Recent results from RIBF



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Some (random) tips

solar abundance, nuclear stability, “human” abundance  
thermonuclear process  
relevant nuclear properties –  $\sigma$ , lifetime, or ....

Coulomb dissociation – experimental aspects

advantages

examples

SAMURAI at RIBF

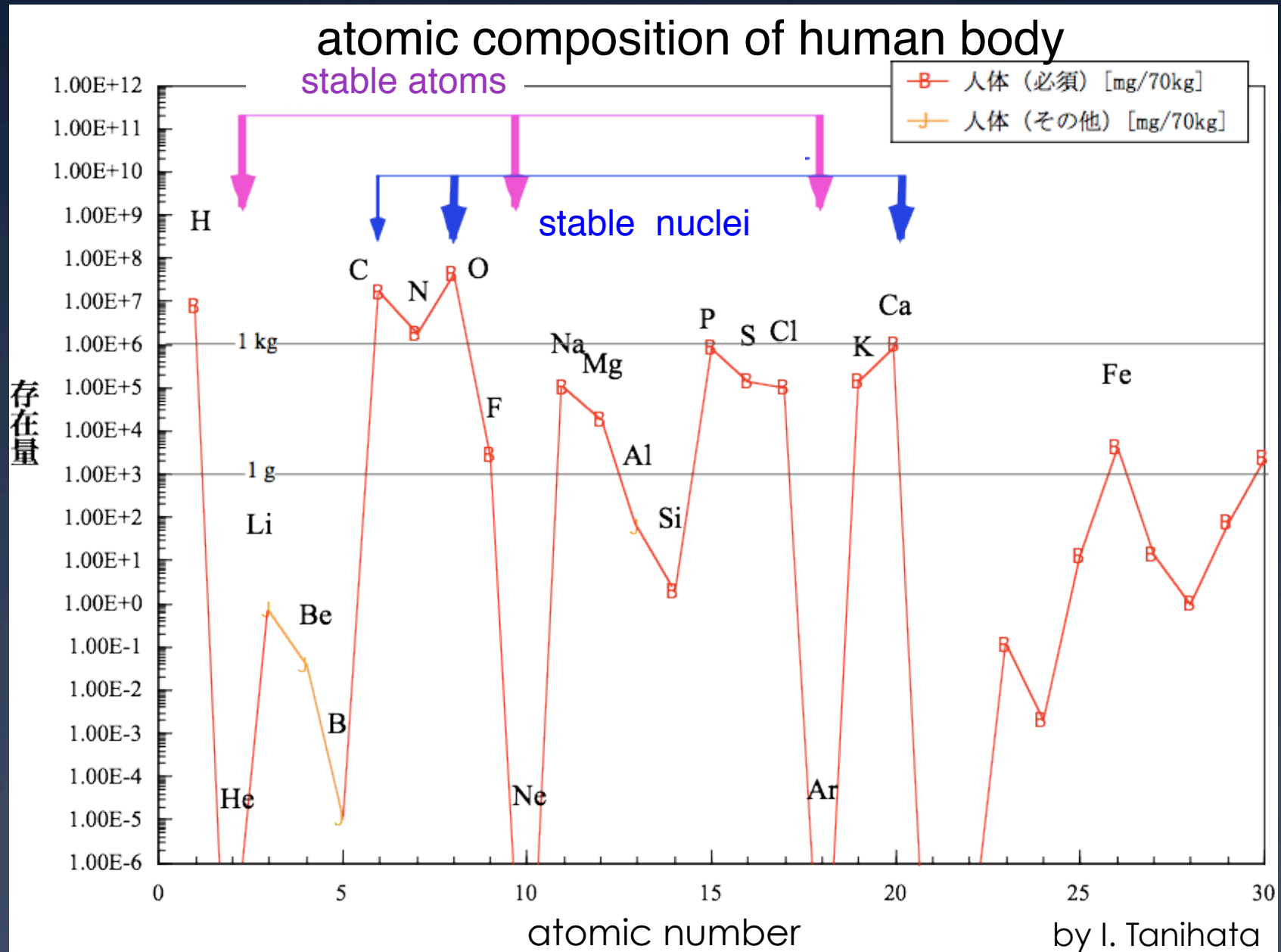
Recent astrophysics studies at RIKEN RI Beam Factory (RIBF)

lifetime measurements for nuclei near and in the r-process path



By the way... Atomic and nuclear stabilities are different. (lucky)

If the magic numbers are common to atom and nucleus, .....



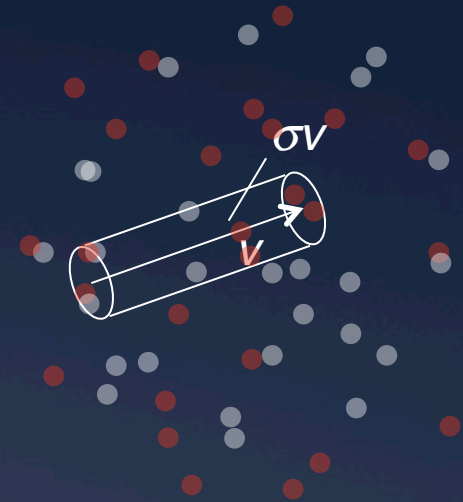


Reaction rate is obtained by the Maxwellian average of  $\sigma$ .

Number of reaction per unit time and unit volume

$$P_{12} = \rho_1 \rho_2 \langle \sigma v \rangle$$

$$\langle \sigma v \rangle = \left( \frac{8}{\pi \mu_{12} (kT)^3} \right)^{1/2} \int dE \sigma(E) E \exp\left[ -\frac{E}{kT} \right]$$



Hot and dense gas

For charged particles:  $\sigma \leftarrow$  Coulomb penetration (tunnel effect)

*e.g.*

$T = 1.5 \times 10^7$  K (sun)  $\rightarrow kT = 1.3$  keV ( $E_G = 20$  keV)  
much lower than the Coulomb barrier

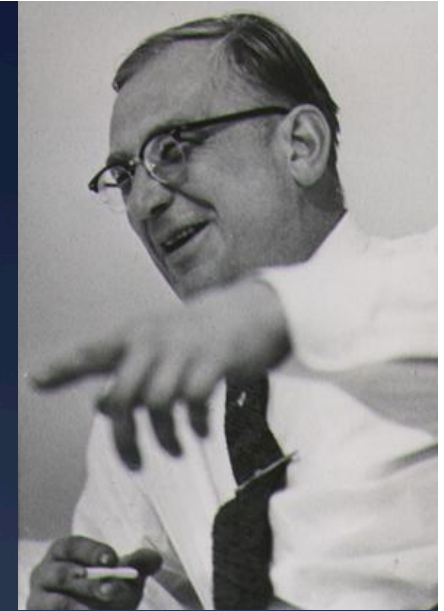
astrophysical S-factor  $\sim$  constant *v.s.*  $E$

$$S = \sigma E \exp[2\pi\eta]$$

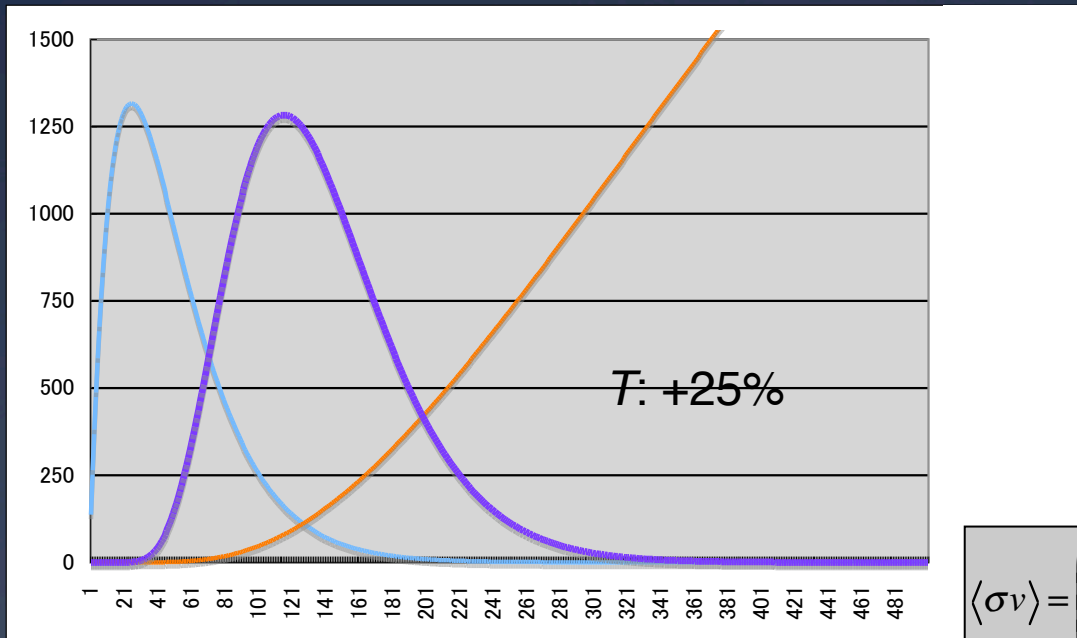
$$\eta = e^2 Z_1 Z_2 / \hbar v$$

The most probable energy is higher than  $kT$ .

“Gamow peak”



Reaction rate  $\langle\sigma v\rangle$  depends strongly on  $T$ .

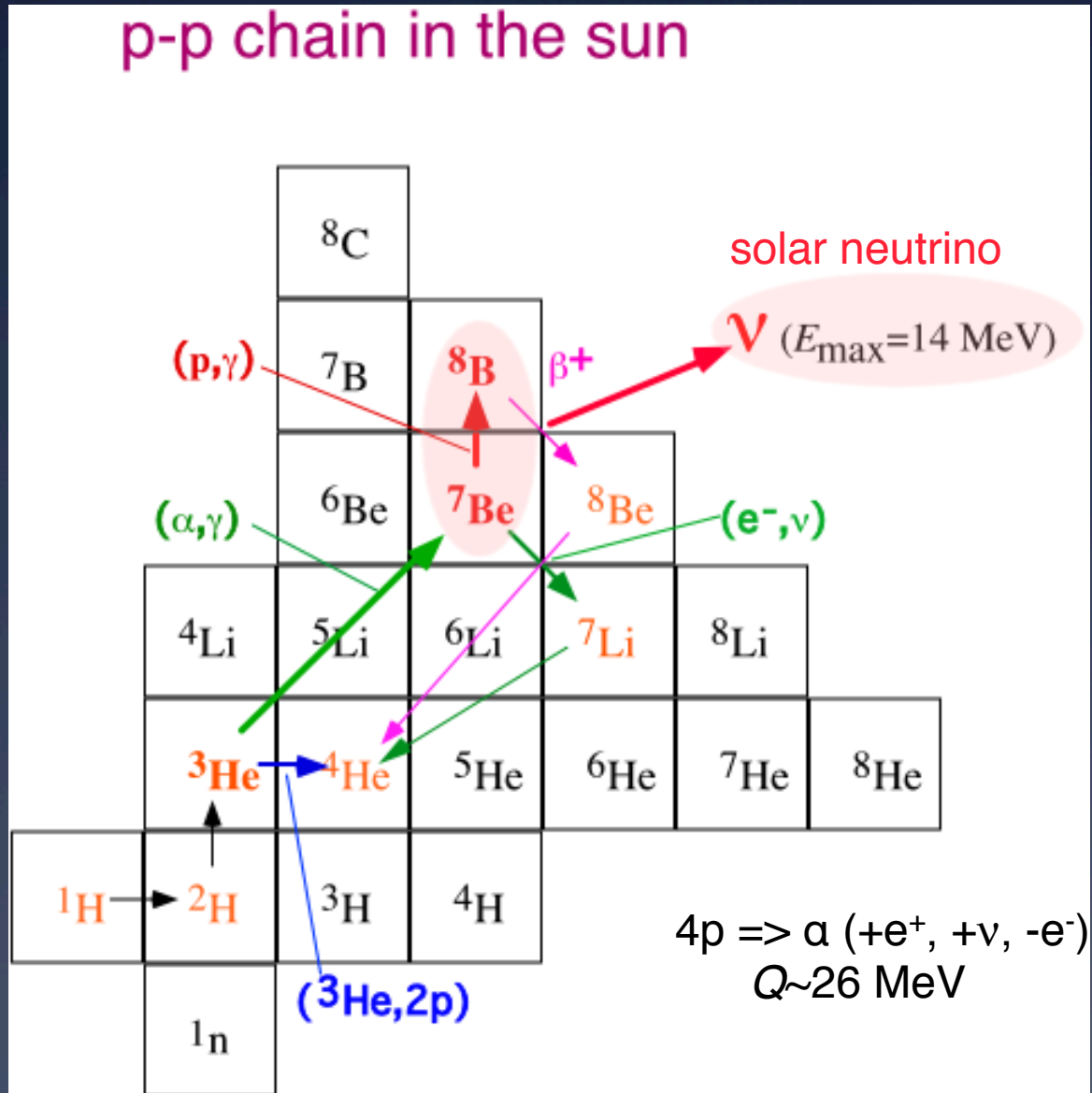


$$\langle\sigma v\rangle = \left(\frac{8}{\pi\mu_{12}(kT)^3}\right)^{1/2} \int dE\sigma(E)E \exp\left[-\frac{E}{kT}\right]$$

assuming constant  $S$

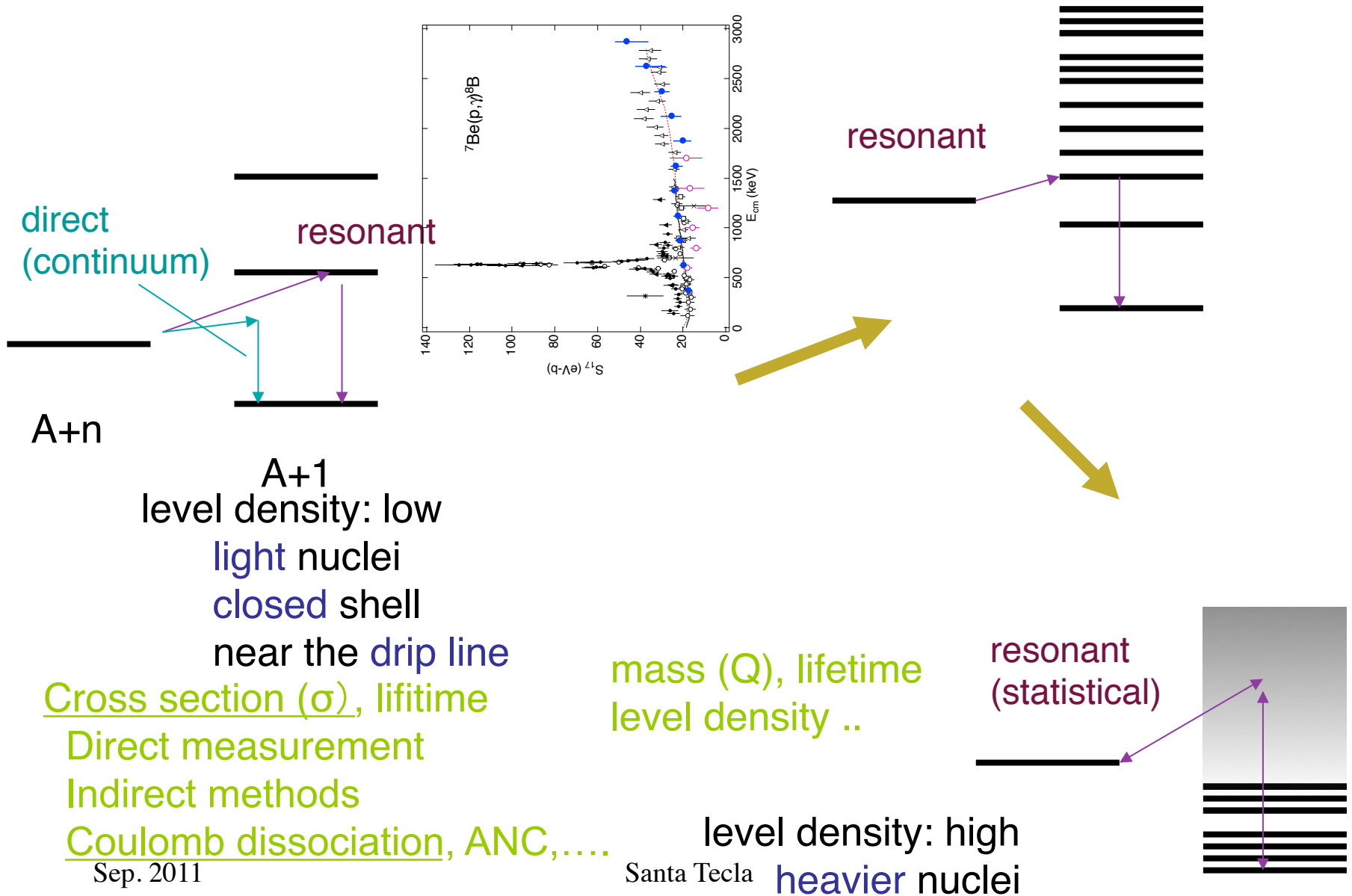


# Competition between reaction (p-capture) and $\beta^+$ decay → measurements of $\sigma$ / lifetime



Situation depends much on the region in the nuclear chart.

$(n,\gamma)$  (or  $(p,\gamma)$ ) – radiative capture





Process can be (mostly) controlled by the reaction  $Q$ -values (not by  $\sigma$ ). *e.g.* r-process

equilibrium in explosive conditions: between the inverse processes  
*e.g.* radiative capture  $\leftrightarrow$  photo-disintegration

Saha equation for  $A+n \leftrightarrow (A+1)+\gamma$

$$\frac{N(Z, A + 1)}{N(Z, A)} = N_n \left( \frac{h^2}{2\pi m_{An} kT} \right)^{3/2} \frac{2j_{Z,A+1} + 1}{(2j_{Z,A} + 1)(2j_n + 1)} \frac{G_{Z,A+1}^{norm}}{G_{Z,A}^{norm}} e^{Q_{m\gamma}/kT}$$

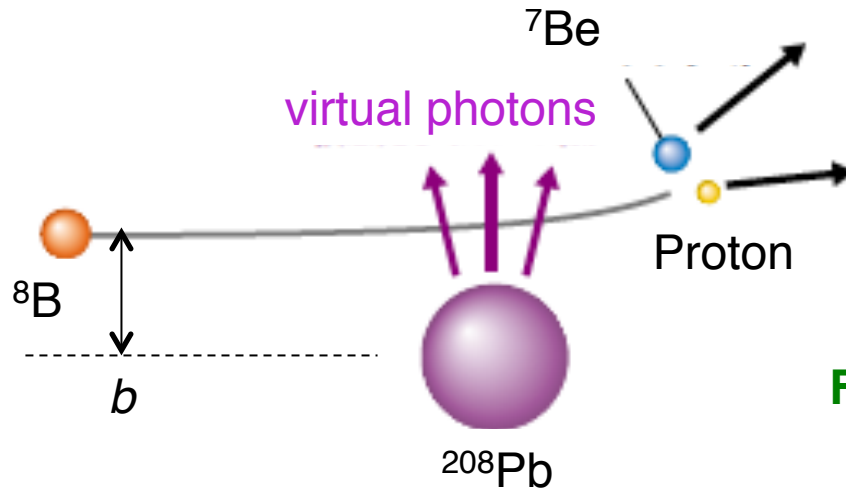
level (density)  
temperature

nuclear mass

$$Q_{ng} = m(A)c^2 + m(n)c^2 - m(A+1)c^2$$

Coulomb dissociation an indirect method for  $\sigma$  measurements  
 = photodisintegration by virtual photon

( $\gamma$  Trojan Horse)

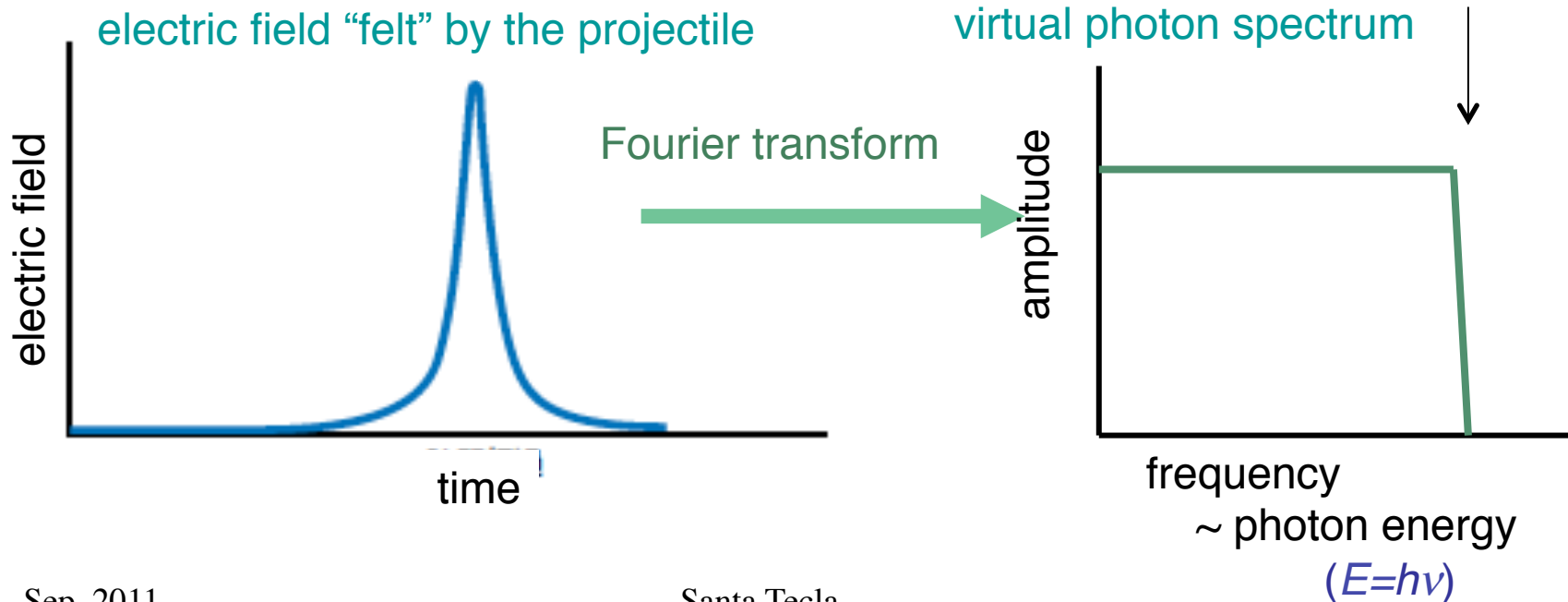


$$E_{in} \nearrow : E_{max} \nearrow$$

$$b \nearrow : E_{max} \searrow$$

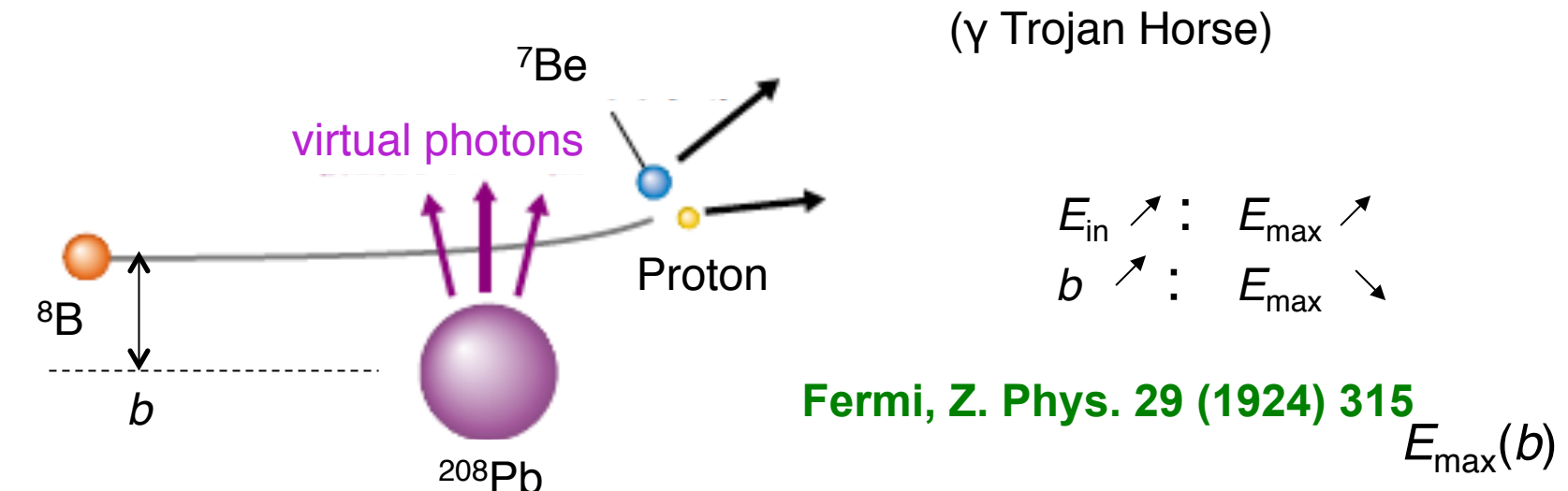
Fermi, Z. Phys. 29 (1924) 315

$E_{max}(b)$

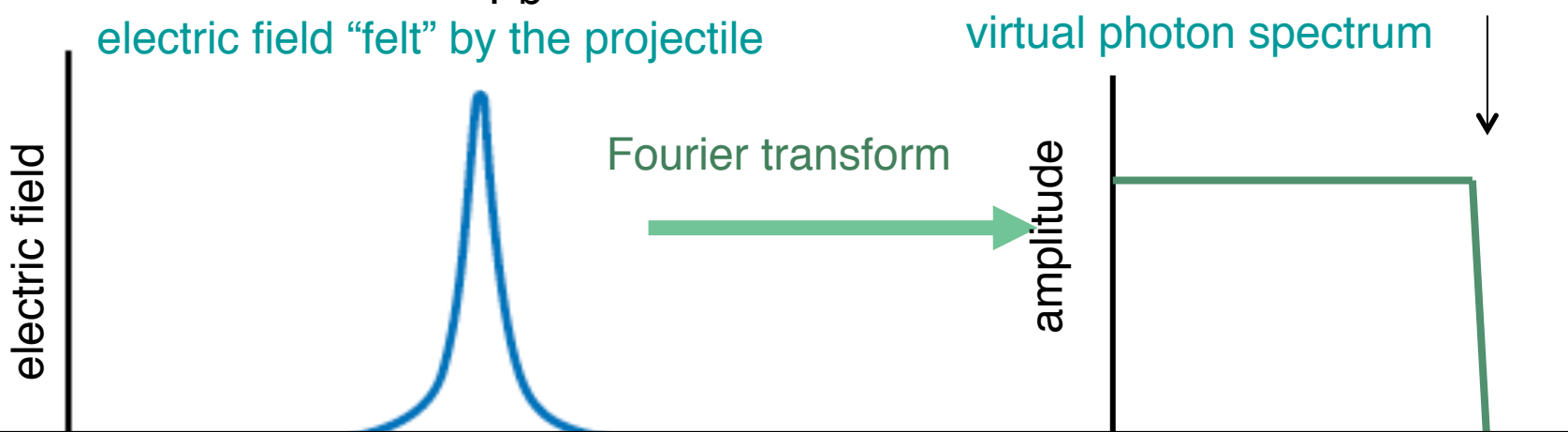




Coulomb dissociation an indirect method for  $\sigma$  measurements  
 = photodisintegration by virtual photon



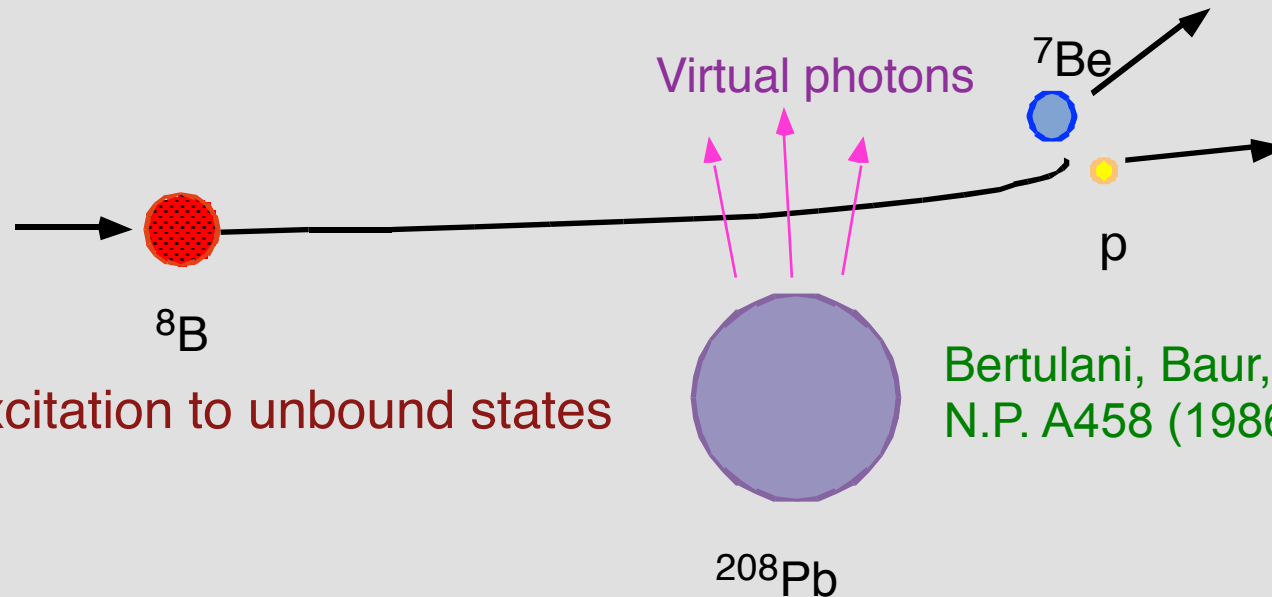
Fermi, Z. Phys. 29 (1924) 315



"Fast beam" can cover the energy range of nuclear excitation.

Coulomb dissociation

successfully applied to astrophysical capture reactions



= Coulomb excitation to unbound states

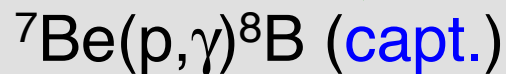
Bertulani, Baur, Rebel  
N.P. A458 (1986) 188



↓ virtual photon theory or DWBA



↓ detailed balance



large  $\sigma$ , thick target (intermediate energy)

experiments with weak RI beams

## detailed balance

$$\sigma_{(\gamma,p)} = \frac{(2j_7 + 1)(2j_1 + 1)}{2(2j_8 + 1)} \frac{k_{17}^2}{k_\gamma^2} \sigma_{(p,\gamma)} \quad 100 \sim 1000$$

## virtual photon number (intermediate energy)

$$\left( \frac{d\sigma}{dE_\gamma} \right)_{\text{C.D.}} = \frac{n}{E_\gamma} \sigma_{(\gamma,p)} \quad 100 \sim 1000$$

## thick target

## charged particle detection

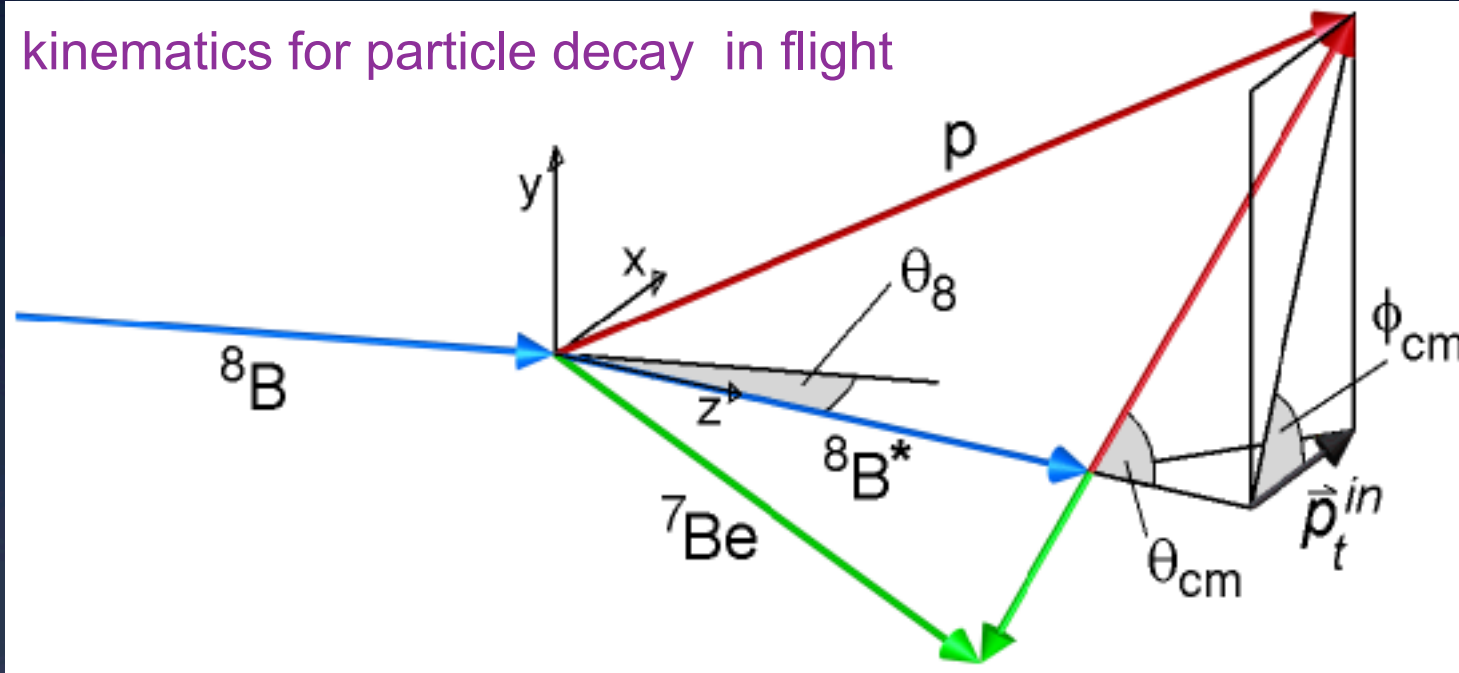
but

Indirect --- nucl. force / higher order / E2 / relativistic effects...

← reaction theory

# Invariant mass (relative energy) measurement

kinematics for particle decay in flight



$\Delta E_{\text{rel}}$  : Independent of  $\Delta E_{\text{in}}$

good for experiments with poor-quality RI beams

$$\Delta E_{\text{rel}} \approx 2 \sqrt{\frac{A_1 A_2}{A_1 + A_2}} \sqrt{T_0 E_{\text{rel}}} \Delta \chi$$

$$\Delta \chi = \Delta \theta, \Delta v / v$$

p+X,  $T_0=100$  AMeV,  $E_{\text{rel}}=1$  MeV,  
 $\Delta q=0.5$  deg.  $\Delta v=1\%$

↓

$$\Delta E_{\text{rel}}=200 \text{ keV}$$



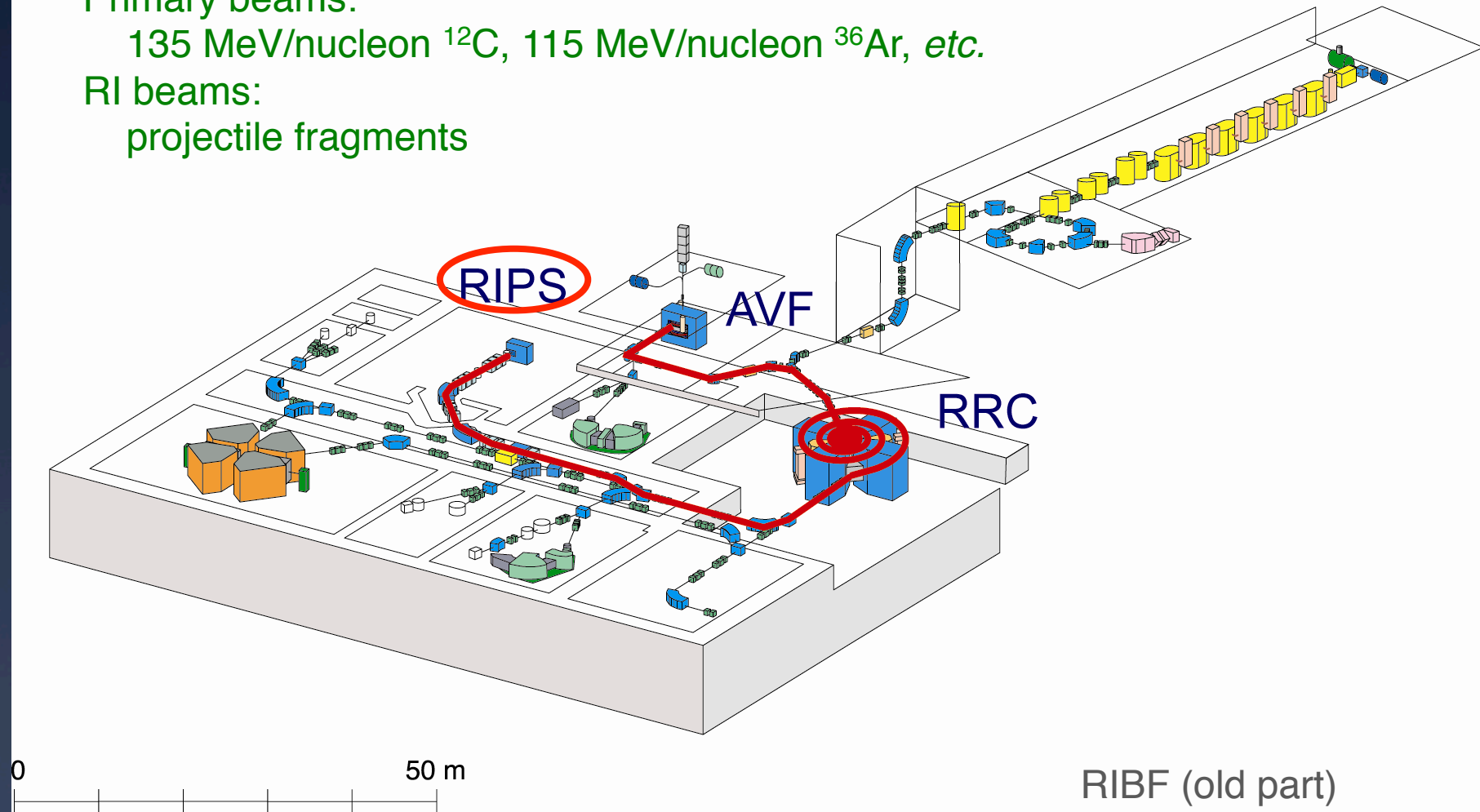
# Experiments at RIPS, RIKEN RIBF

Primary beams:

135 MeV/nucleon  $^{12}\text{C}$ , 115 MeV/nucleon  $^{36}\text{Ar}$ , etc.

RI beams:

projectile fragments



# Experiments at RIKEN (CD) and Louvain la Neuve (direct) on $^{13}\text{N}(p,\gamma)^{14}\text{O}$

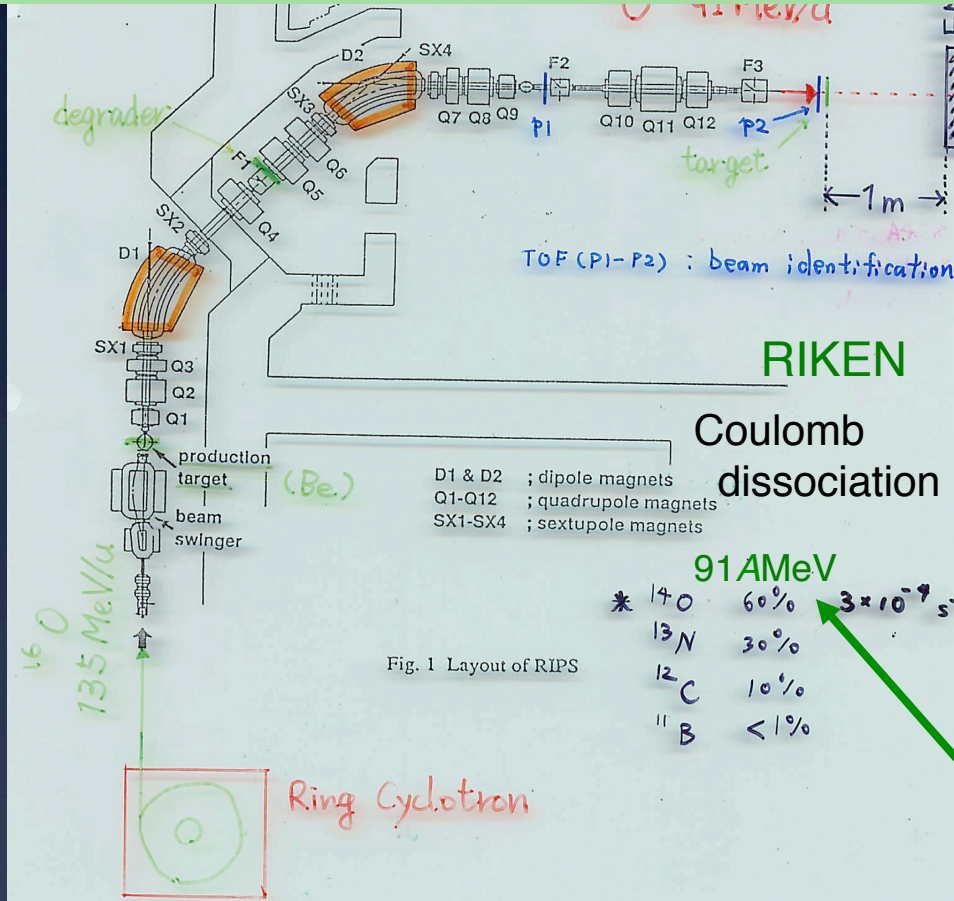
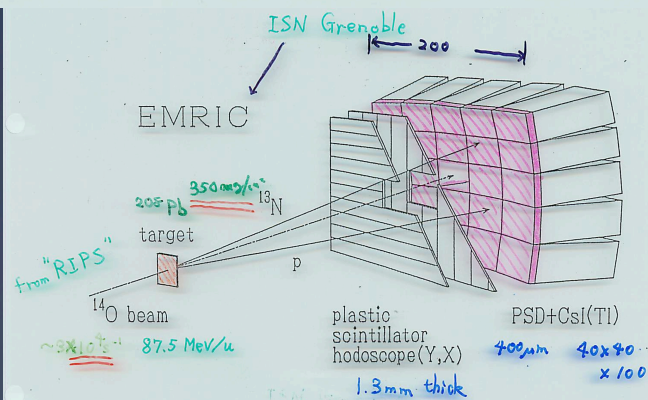


Fig. 1 Layout of RIPS

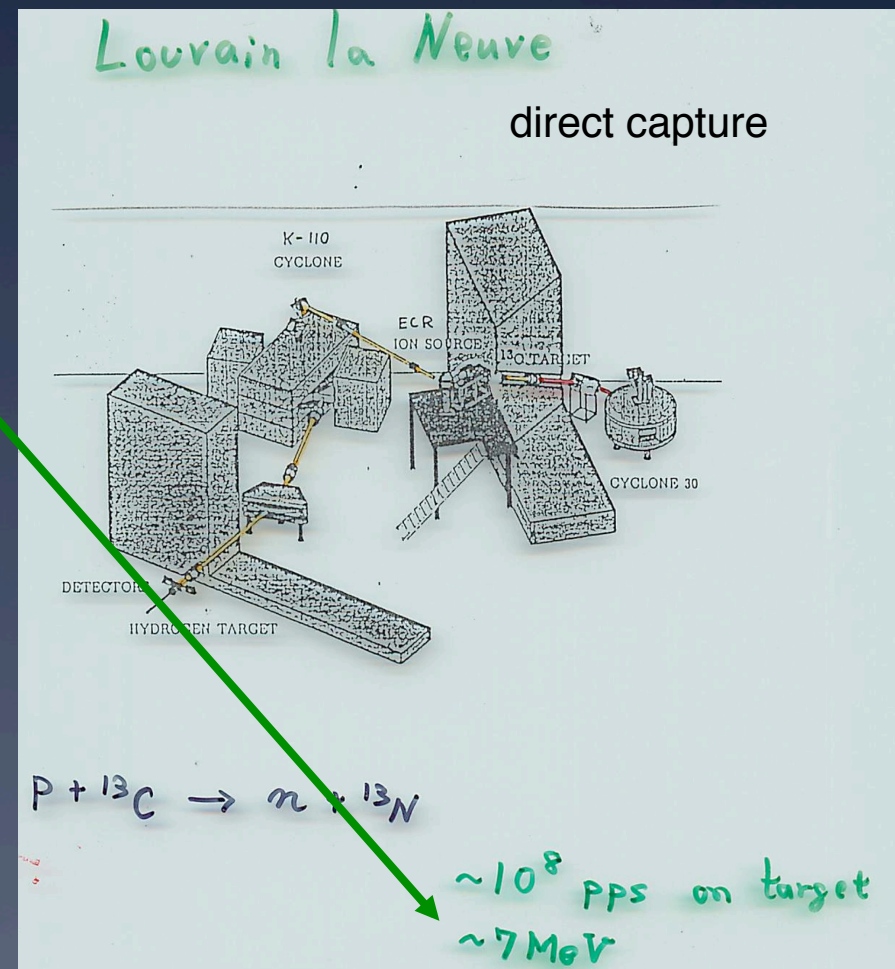
91 AMeV

$^{14}\text{O}$	60%	$3 \times 10^{-9} \text{ s}^{-1}$
$^{13}\text{N}$	30%	
$^{12}\text{C}$	10%	
$^{11}\text{B}$	< 1%	

Ring Cyclotron



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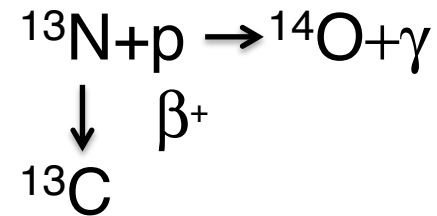
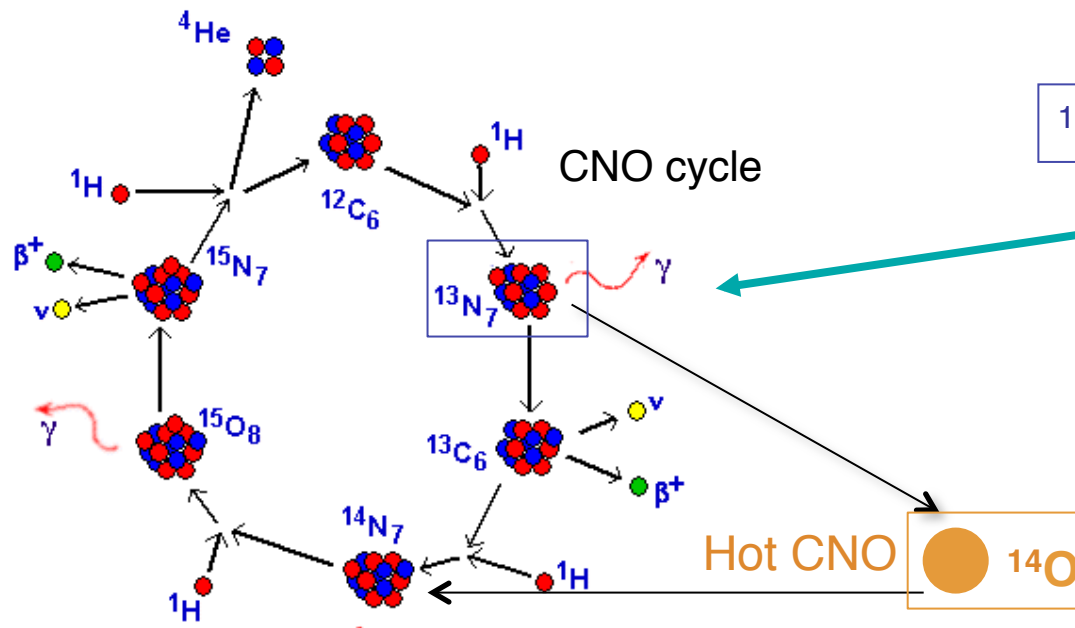
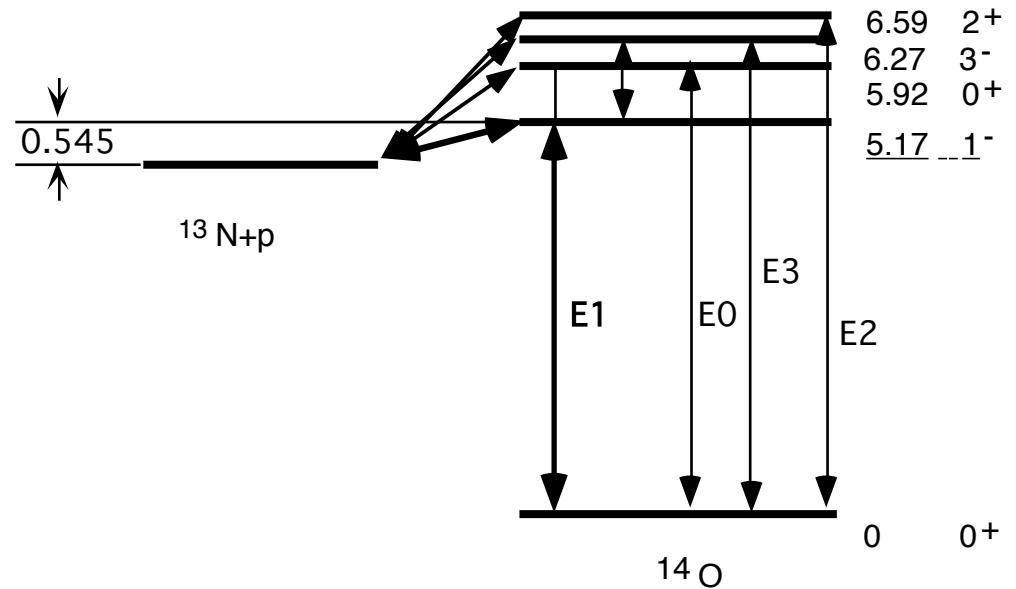


# $^{13}\text{N}(p,\gamma)^{14}\text{O}$ – CNO v.s. hot CNO

In hot CNO cycle,  
 $^{13}\text{C}$  is bypassed.  
 faster cycle.

$^{14}\text{O}: t_{1/2} = 1\text{min.}$

$^{13}\text{N}: t_{1/2} = 10\text{min.}$

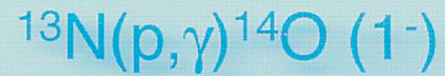


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Coulomb dissociation is efficient.



	C.D. RIKEN ('91)	Direct Louvain ('91)	ratio
beam(s <sup>-1</sup> )	3×10 <sup>4</sup>	3×10 <sup>8</sup>	10 <sup>-4</sup>
σ	10 mb	100 μb	10 <sup>2</sup>
target	350 mg/cm <sup>2</sup>	200 μg/cm <sup>2</sup>	60 ← n. of atoms
efficiency	0.5	2×10 <sup>-3</sup>	350
data taking	36 h	30 h	1.2
total counts	1.5×10 <sup>4</sup>	85	180



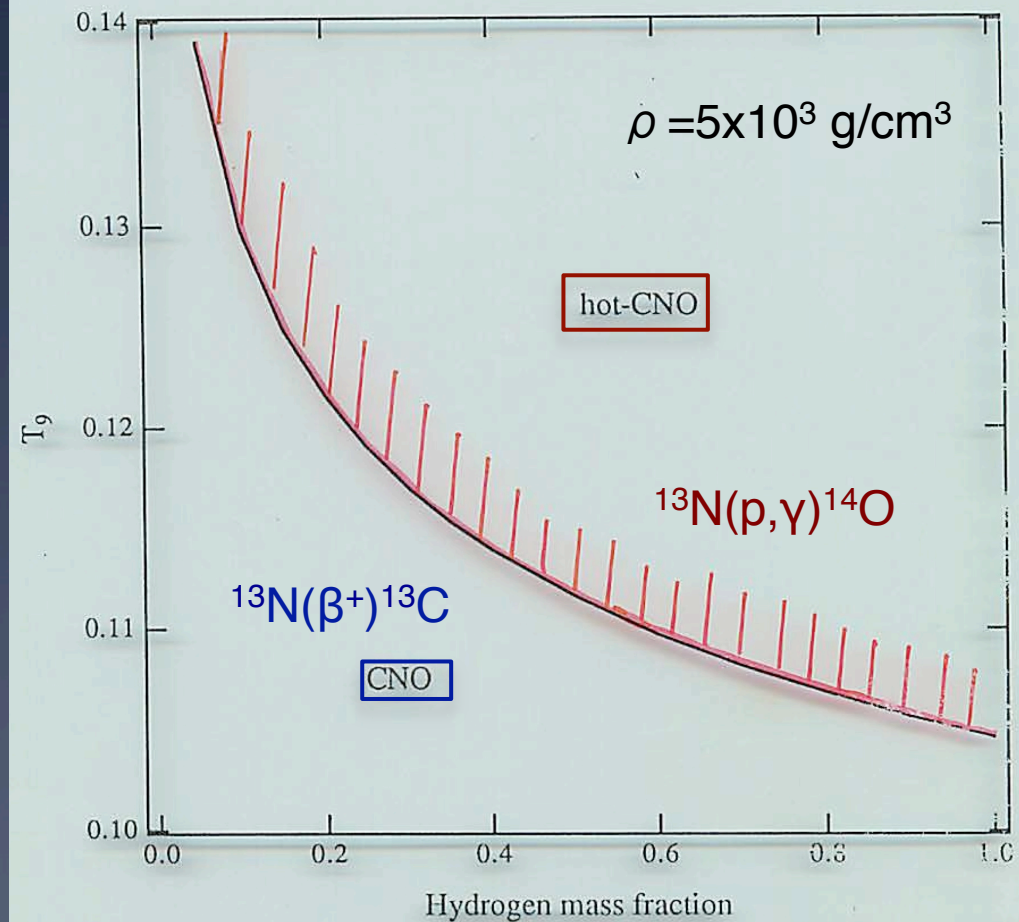
# results -- measured radiative width / CNO hot-CNO boundaries

$\Gamma_\gamma$	present	(P, $\rho$ )
$^{14}\text{O} (1^-)$	$3.1 \pm 0.6$ eV	$(3.8 \pm 1.2) \times 10^8$ (3, 2) $3.4 \pm 0.9$
$^{13}\text{N} (1/2^+)$	$0.59 \pm 0.18$ eV	$0.50 \pm 0.04^*$ eV

reaction rate  $\leq \omega \Gamma_\gamma \cdot E_0$

$$\langle \sigma v \rangle \propto \omega \Gamma_\gamma (kT)^{-3/2} \exp\left[-\frac{E_0}{kT}\right]$$

$$P_{12} = \rho_1 \rho_2 \langle \sigma v \rangle$$



## “typical” numbers for (p,γ) experiments

	direct	C. D.	ANC
$\sigma(\text{mb})$	$< 0.1$	10	1
Target ( $\text{mg}/\text{cm}^2$ )	0.1	100	1
$\sigma \cdot t$	$< 0.01$	100	1
Beam ( $\text{s}^{-1}$ )	$10^8$	$10^4$	$10^5$
eff.	$10^{-3}$	$10^{-1}$	1
relative yield	$< 10^3$	$10^5$	$10^5$
	$^{13}\text{N}$ -LLN	$^{14}\text{O}$ -RIKEN	$^8\text{B}$ -TAMU

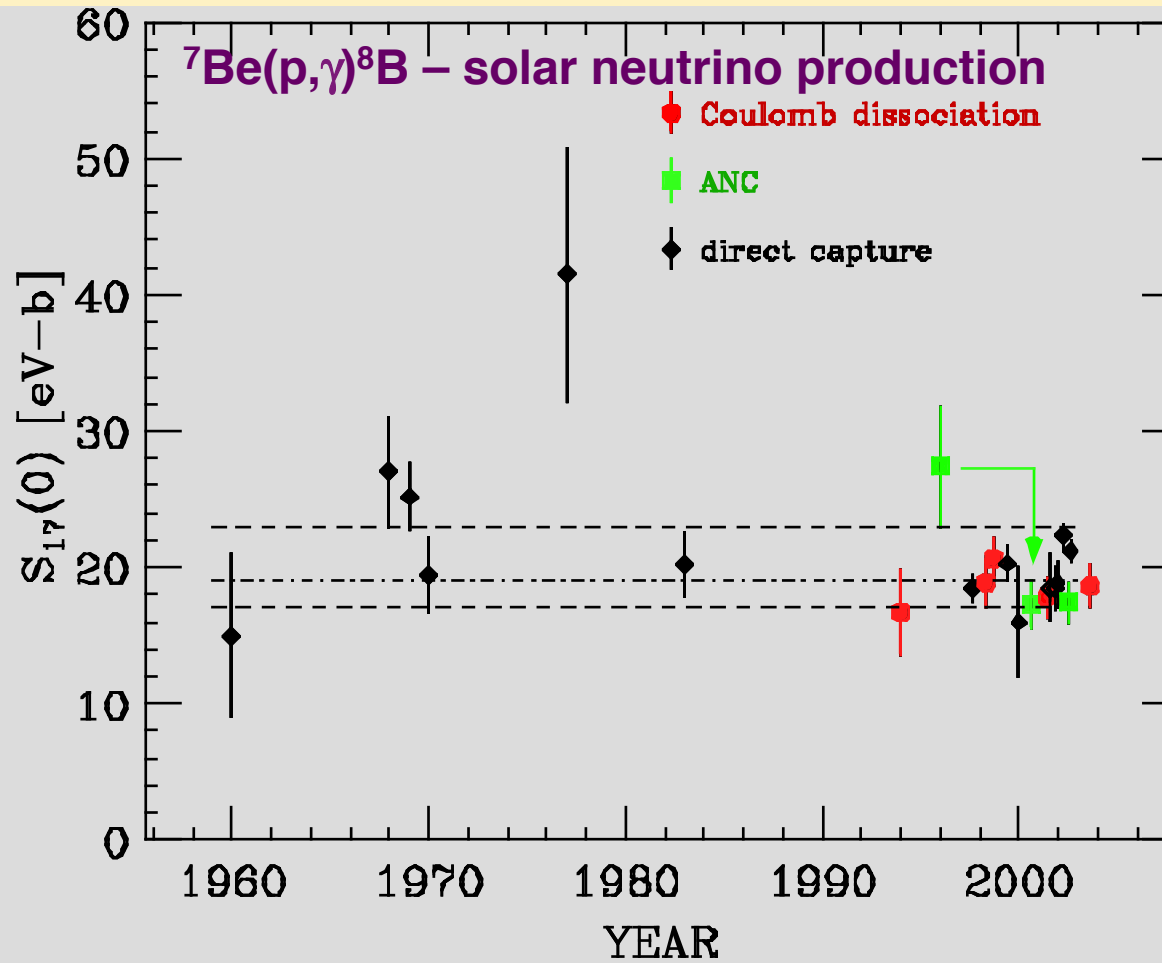
ANC:

Asymptotic Normalization Coefficient  
by transfer (breakup)

Indirect measurements appears to be better ...

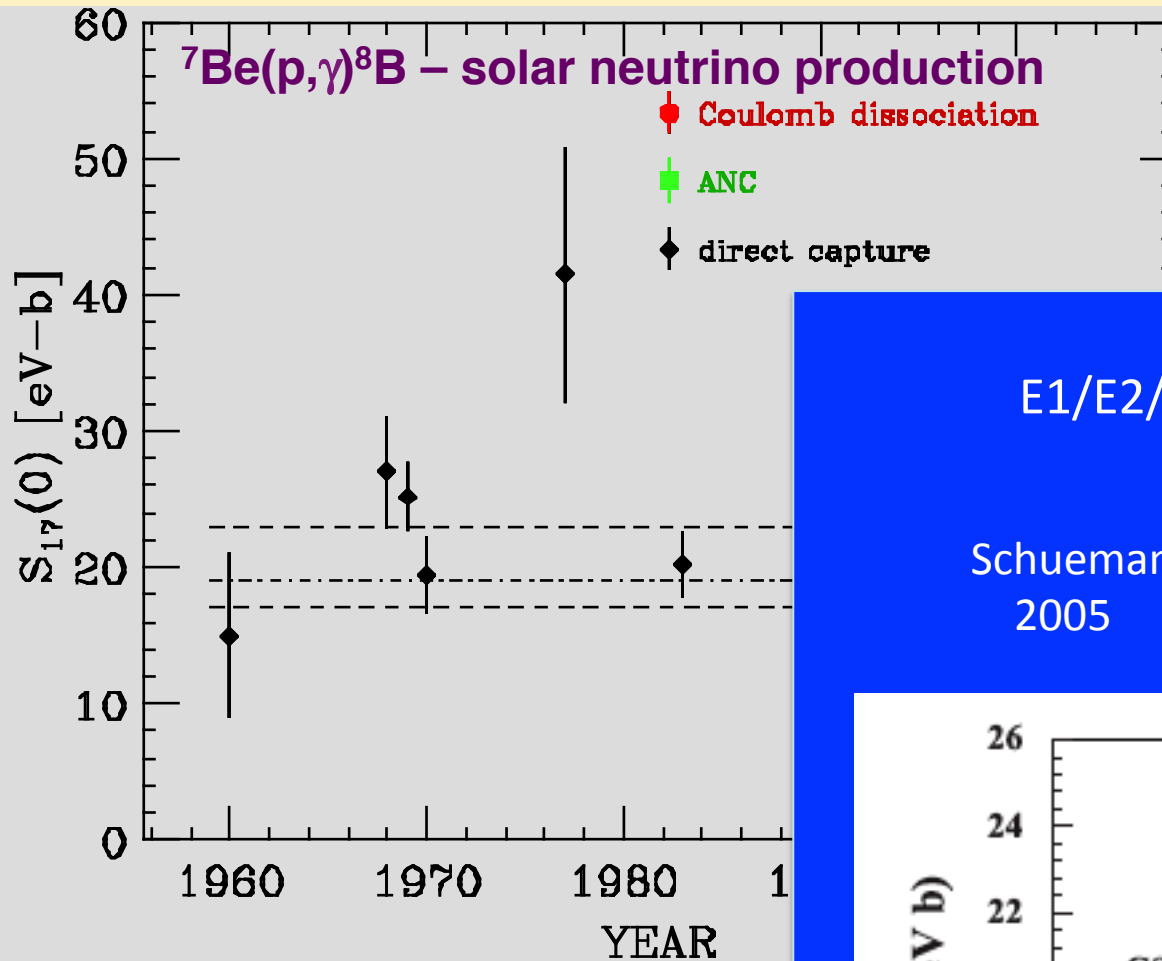
... but direct measurements have several advantages.

# Recommended values change as experimental and theoretical development



$S_{17}$  at  $E=0$

Recommended values change as experimental and theoretical development

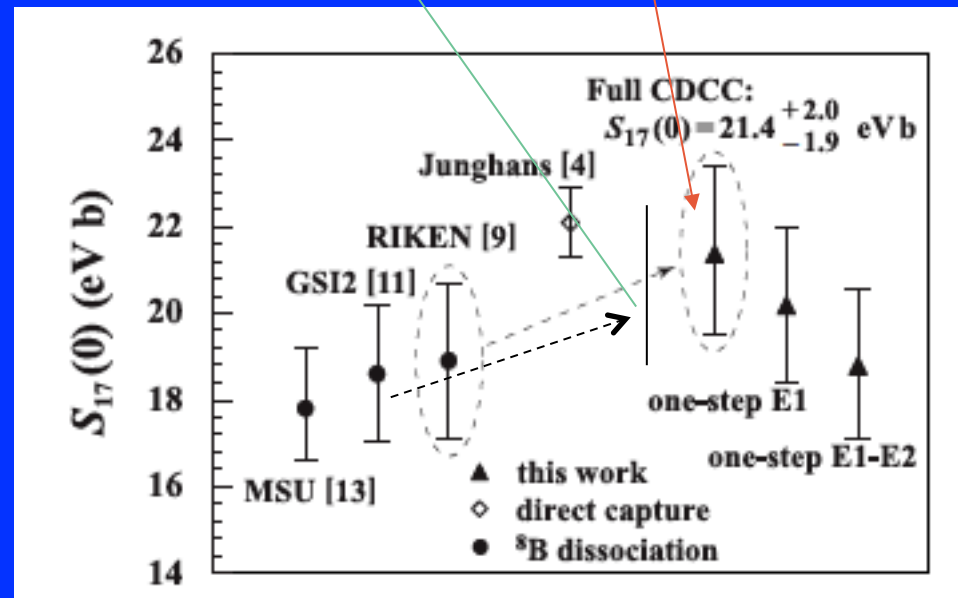


$S_{17}$  at  $E=0$

E1/E2/nucl. interference

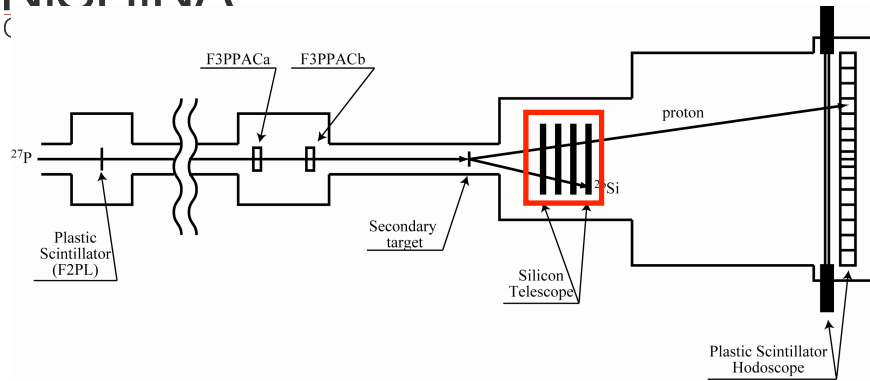
Ogata *et al.* (CDCC)

Schuemann *et al.*  
2005

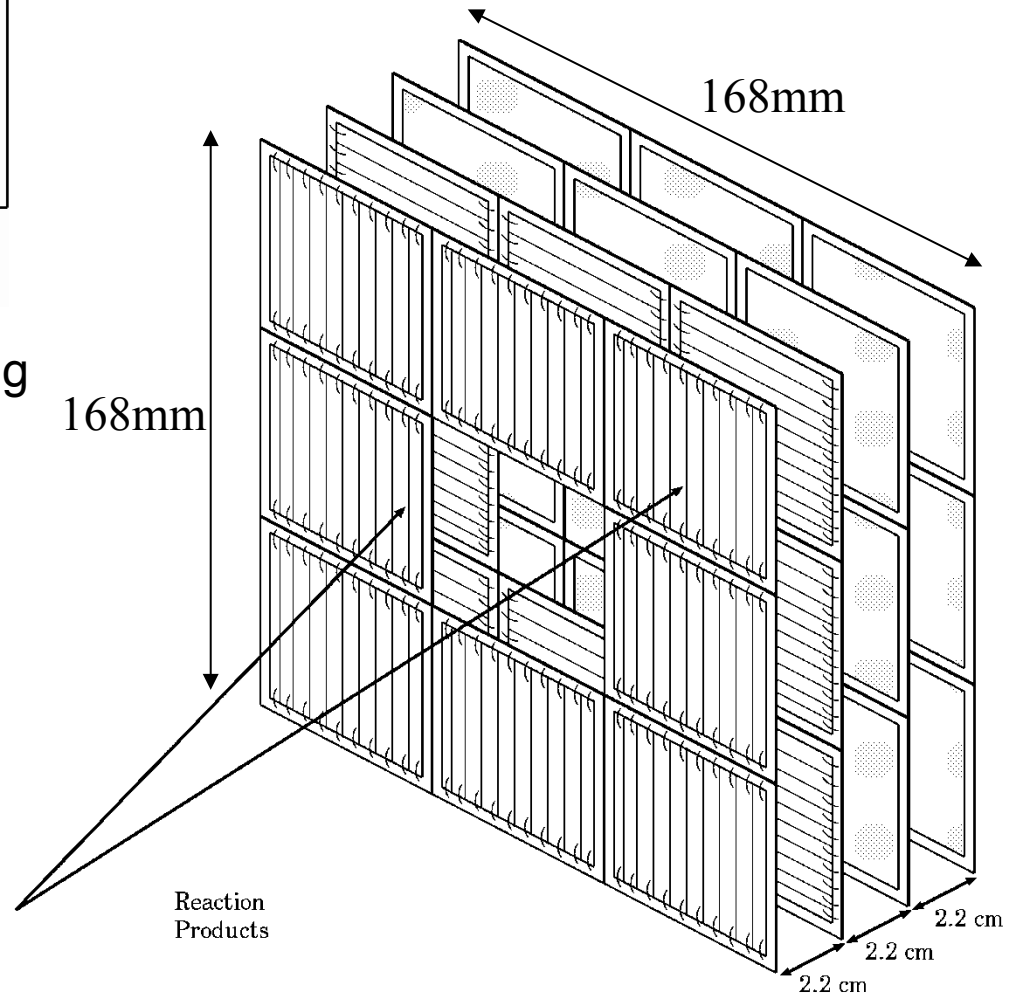




## Silicon telescope



- Positions (angles) of proton and  $^{22}\text{Mg}$
- Kinetic energy of  $^{22}\text{Mg}$ , isotope ID
  - $50 \times 50 \times 0.5 \text{ mm}^3 \text{ Si} \times 8$
  - Hole at the center
  - 4 layers to stop  $^{22}\text{Mg}$
  - 2 layers at upstream
    - 5mm strips for position
  - Other 2 layers: Single element

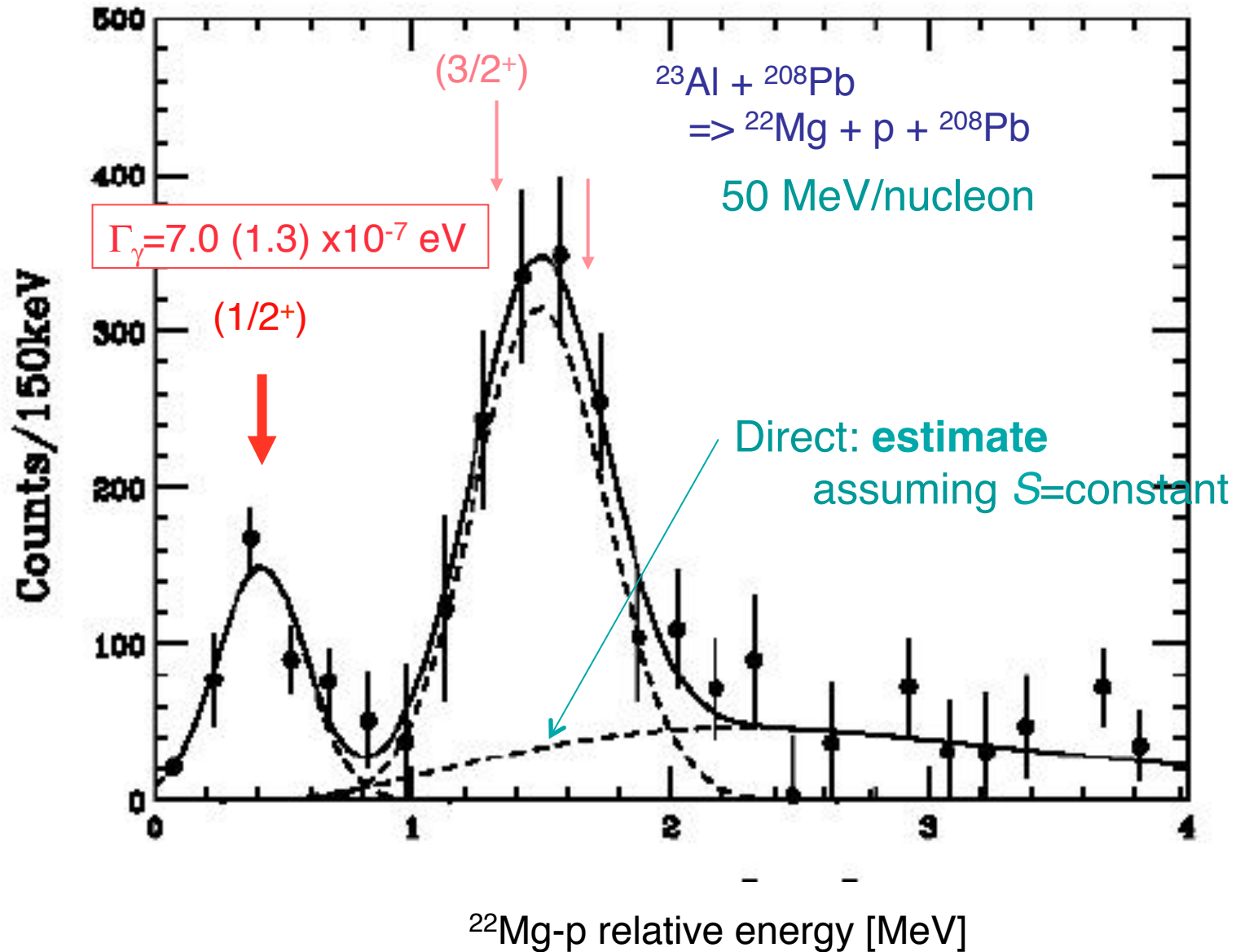


Proton energy by TOF

Coulomb dissociation -  $^{22}\text{Mg}(p,\gamma)^{23}\text{Al}$  result – demonstration of its efficiency

Coulomb dissociation -  $10^4$  pps  $^{23}\text{Al}$

$\leftrightarrow 10^{12}$  pps  $^{22}\text{Mg} ! + ^1\text{H}$

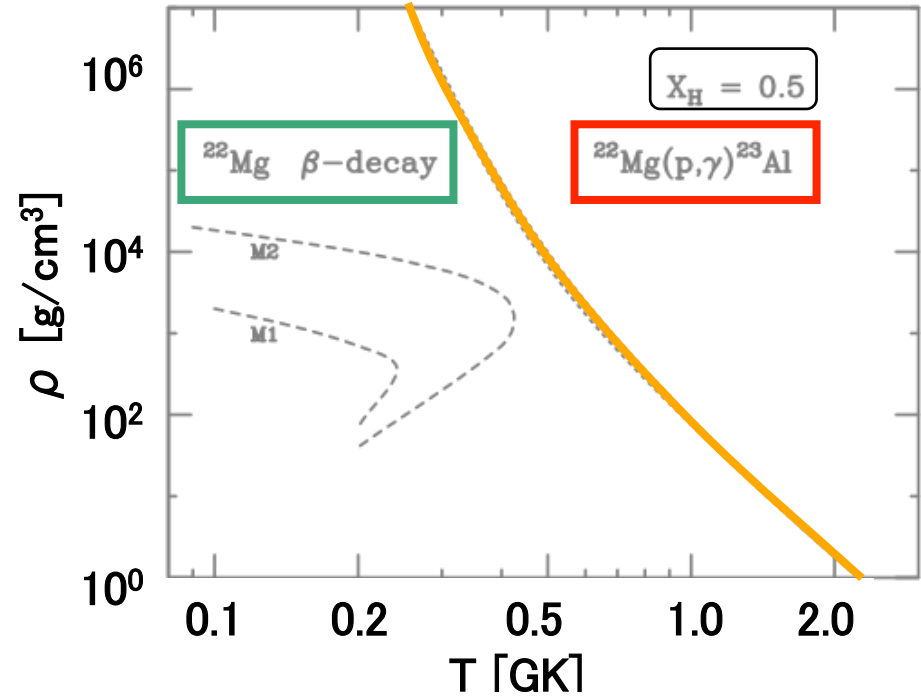
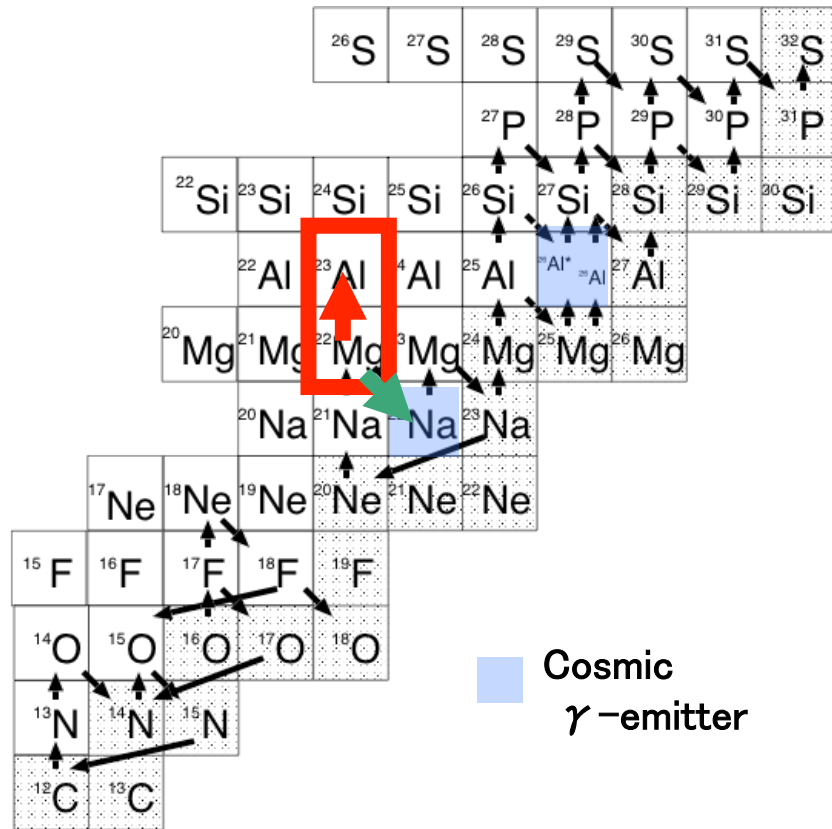


$^{22}\text{Mg}(p,\gamma)^{23}\text{Al}$  may not contribute in novae.  $\Leftarrow$  our result

our resonance data w. direct component

- $\rightarrow$  reaction rate
- $\rightarrow$  competition with  $\beta$  decay

c.f. Lectures by Trache



Nova Model

M1 : J. Jose et al. *Astrophys. J.* 520 347 (1999)

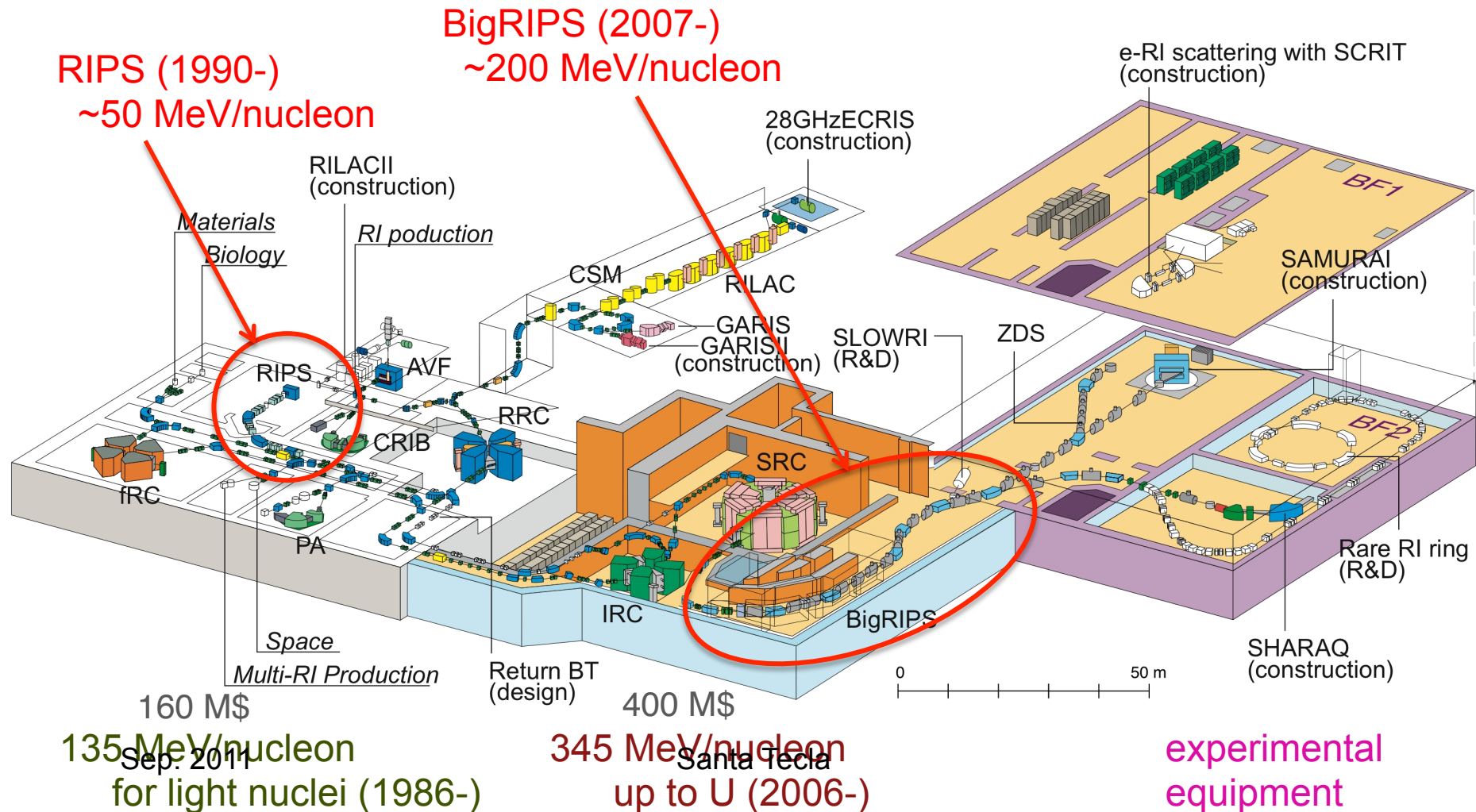
M2 : C. Iliadis et al. *Astrophys. J. Supp.* 142 105 (2002)

*In novae ( $\leftrightarrow$  X-ray bursts)  
 $\beta$  decay may be always favored.*

Similar work at RIKEN on  $^{26}\text{Si}(p,\gamma)^{27}\text{P}$ : Togano et al., *Phys. Rev. C* 84. 035808 (2011)

# RIBF new facility

**RIBF** – a new generation RIB facility in operation  
with world highest capability of producing exotic nuclei in coming years!

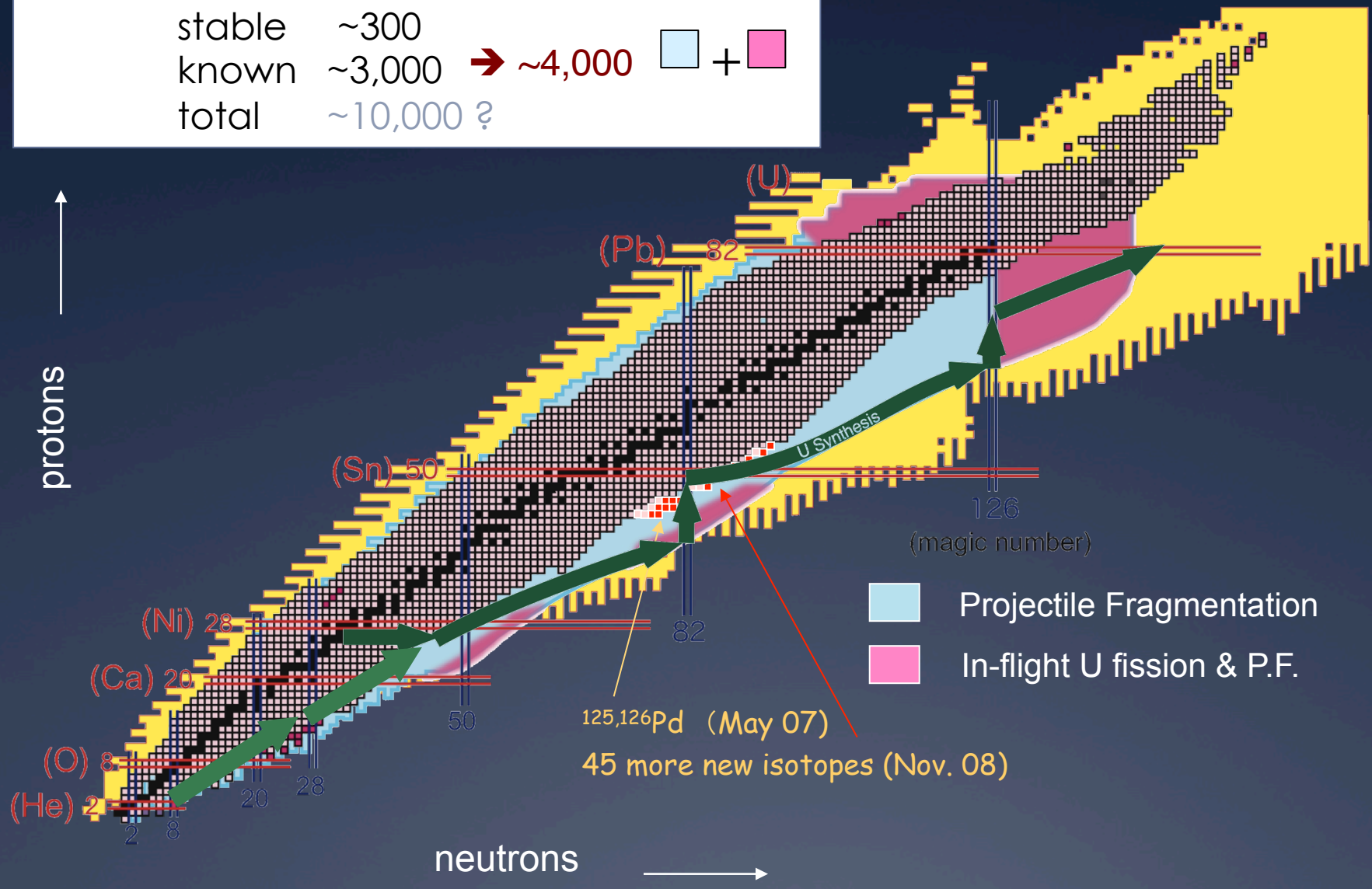




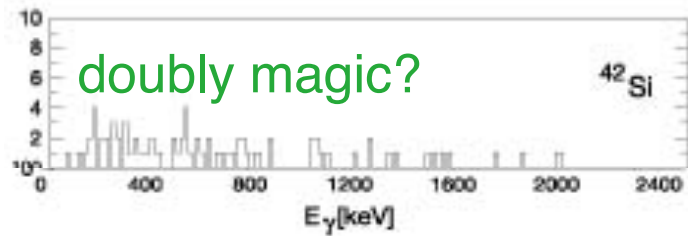
# Atomic nuclei (isotopes)

stable ~300  
 known ~3,000 → ~4,000  
 total ~10,000 ?

> 1 particle/day (Goal Intensity)  
 (EPAX prediction)

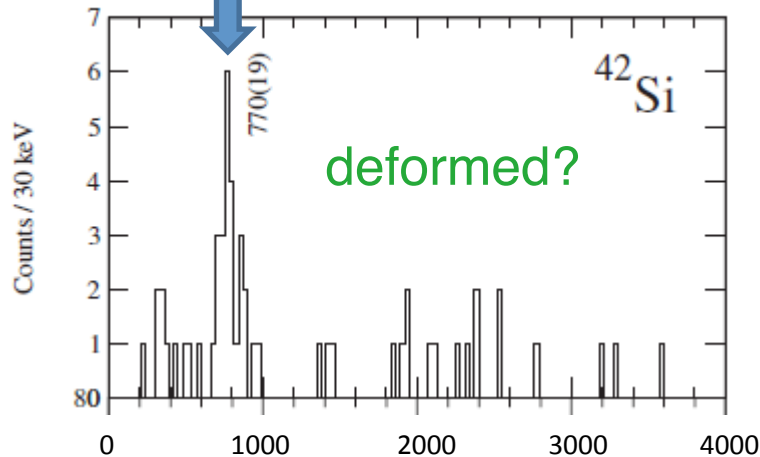


“doubly magic”  $^{42}\text{Si}$  ( $Z=14$ ,  $N=28$ ) - confirmation of the low-lying  $2^+$  / higher states



NSCL → Fridman et al. Nature, PRC74(2006)  
 ( No low-energy peak, small 2p removal  $\sigma$ )  
 Energy ~100A MeV  
 Intensity Integral  $1 \times 10^8$  Particles

2006 GANIL data ( ? days )



GANIL → Bastin et al. PRL 99(2007)022503  
 Energy ~40A MeV  
 Intensity ~125 cps

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Takeuchi *et al.*

2010 RIKEN data (~90 hours)



200 A MeV, 40,000 cps

High RI beam intensity  
 x300 of GANIL  
 x100 of MSU – integral  
 High DALI2 efficiency

high statistics

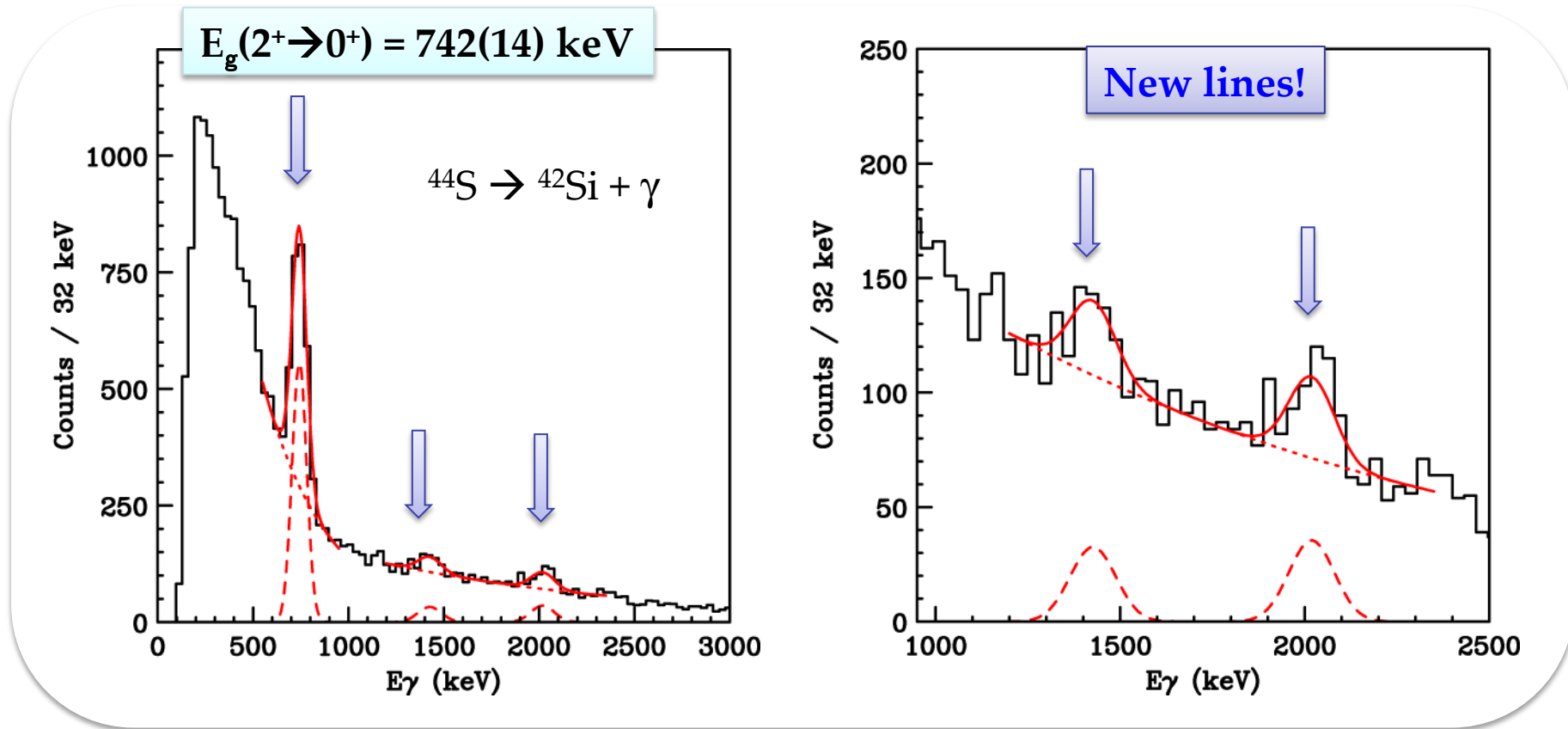
→  $\gamma$ - $\gamma$  coincidence

→ states above  $2^+$

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# Results - $\gamma$ single spectrum -

\*Widths are fixed to simulated values.



GANIL exp. : 770(19) keV

B.Bastin et al., Phys. Rev. Lett. 99, 022503 (2007).

We confirmed the  $\gamma$ line observed at  
GANIL.



“Seven Samurai”, movie by Akira Kurosawa (1954)





Superconducting coils have been successfully excited.

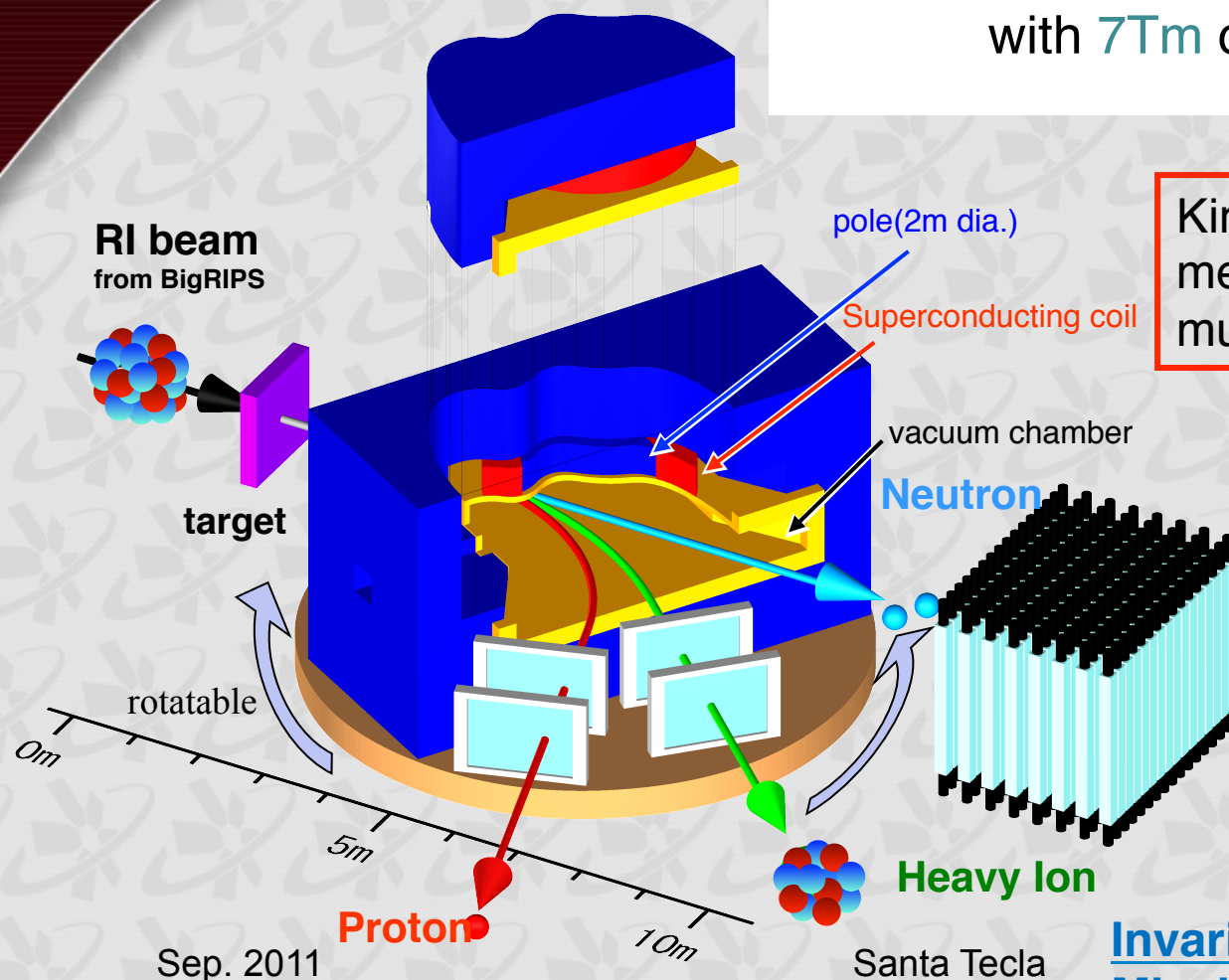
Apr. 2011



# SAMURAI

-- new spectrometer in RIBF --

Superconducting Analyzer for MUlti-particle  
from RAdio Isotope Beam  
with 7Tm of bending power



Kinematically complete  
measurements by detecting  
multiple particles in coincidence

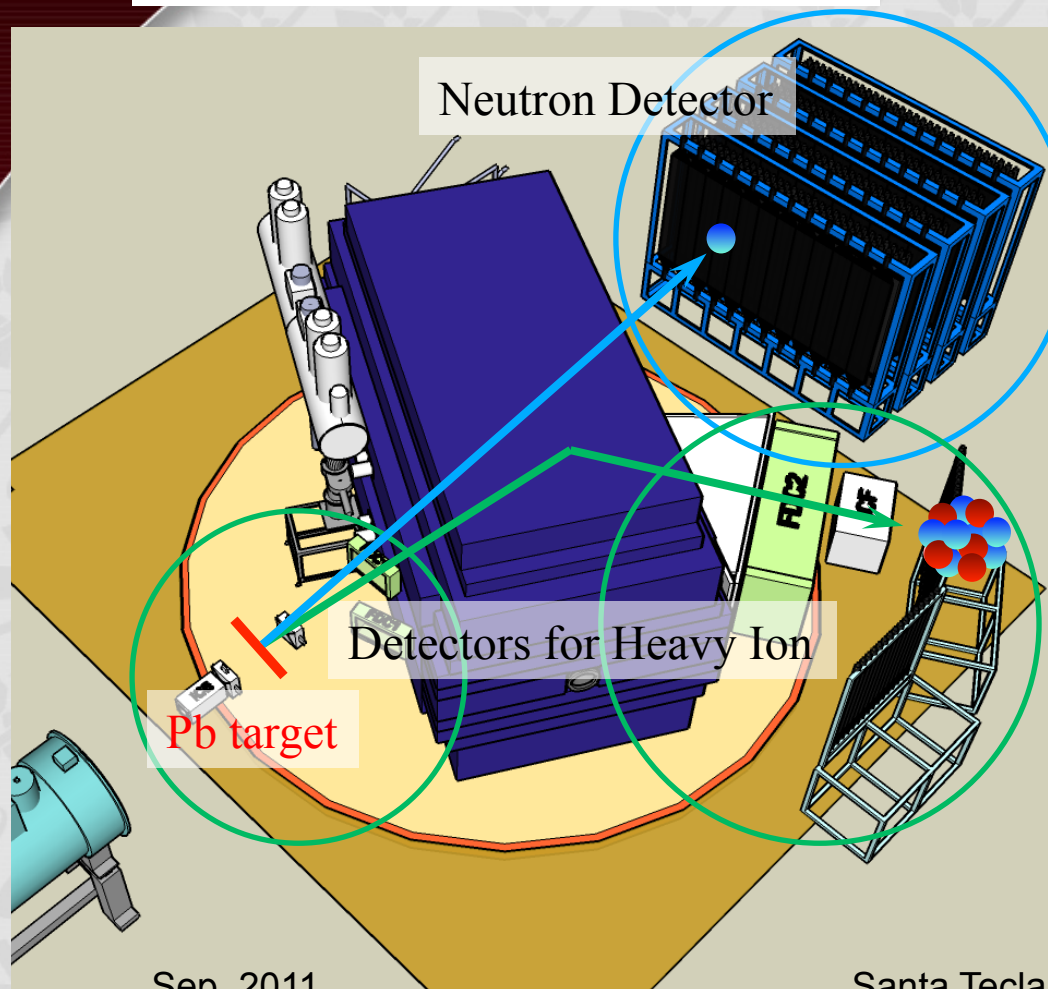
- Superconducting Magnet  
3T with 2m dia. pole  
(designed resolution 1/700)  
80cm gap (vertical)
- Heavy Ion Detectors
- Proton Detectors
- Neutron Detectors
- Large Vacuum Chamber
- Rotational Stage

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Invariant Mass Measurement  
Missing Mass Measurement

# Detector System - ( $\gamma, n$ ) measurement (CD) mode

( $\gamma, n$ ) reaction: neutron-rich side



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## • Detectors for Heavy Ion

- Position measurement
  - Drift Chambers
  - Beam
  - Fragments before/after
- Charge measurement
  - Ion Chambers
  - Beam / Fragments
- Velocity measurement
  - Plastic hodoscope
  - Cherenkov counter
- Total E measurement
  - Pure CsI detector

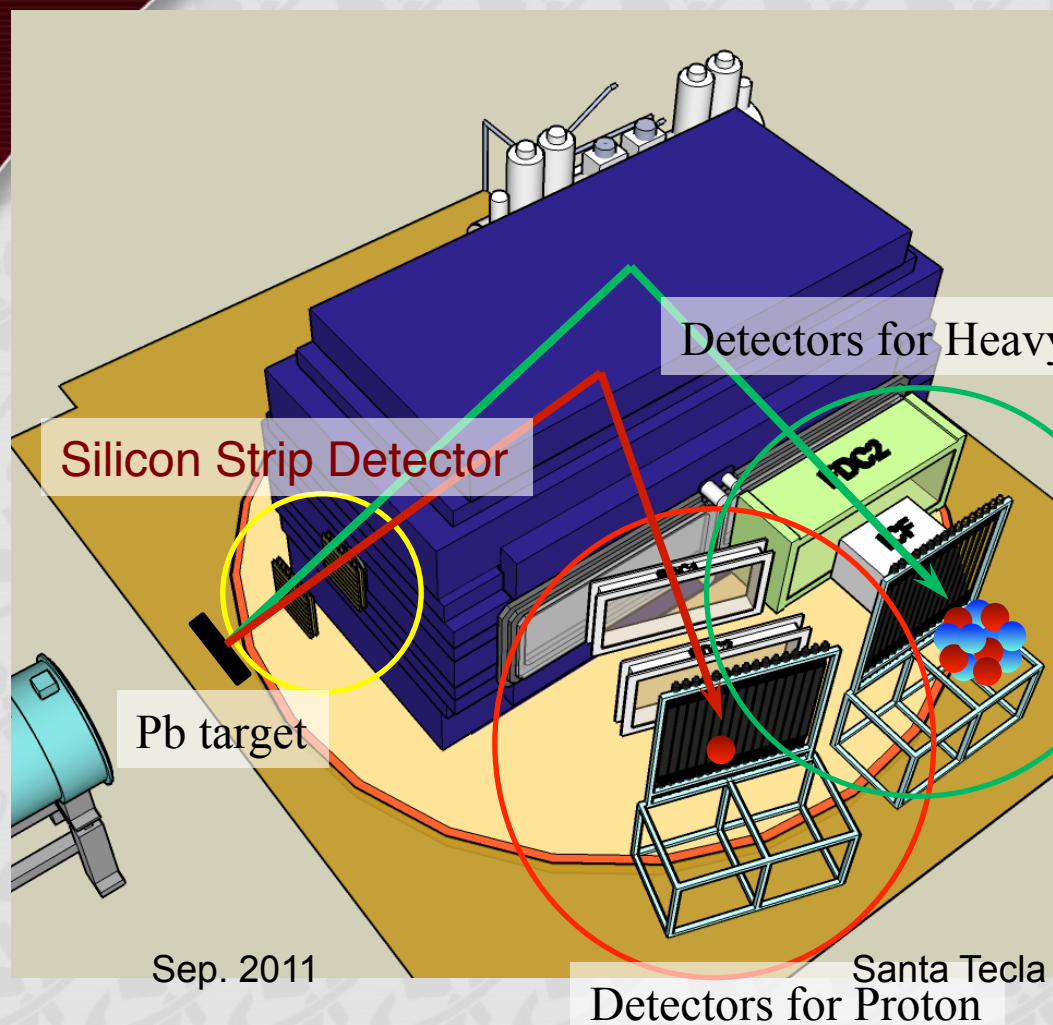
## • Neutron Detector

- Plastic scintillator
  - 240 modules
  - (120mm x 120mm x 1.8m / module)
- Effective Area: 3.6 m (H) x 1.8 m (V)
  - ~ 100 % coverage @  $E_{rel} < 3$  MeV
  - ~ 40 % coverage @  $E_{rel} \sim 10$  MeV
- Efficiency ~ 66 % (Harf: ~40 %)

Half volume is ready

# Detector System - ( $\gamma$ , $p$ ) measurement mode

( $\gamma$ ,  $p$ ) reaction: proton-rich side



- Detectors for Heavy Ion

- same as ( $\gamma$ , n)

- Detectors for Proton

- Proton Drift Chamber
- Plastic Hodoscope

- Silicon Strip Detector

- **Broad dynamic range (~10,000)**

Both proton & heavy ion ( $Z < 50$ ) hit the detector

- **High density signal processing**  
Signals of about 2500ch in total

**Under development**

Based on ASIC technology  
in collaboration with

Texas A&M Univ. and

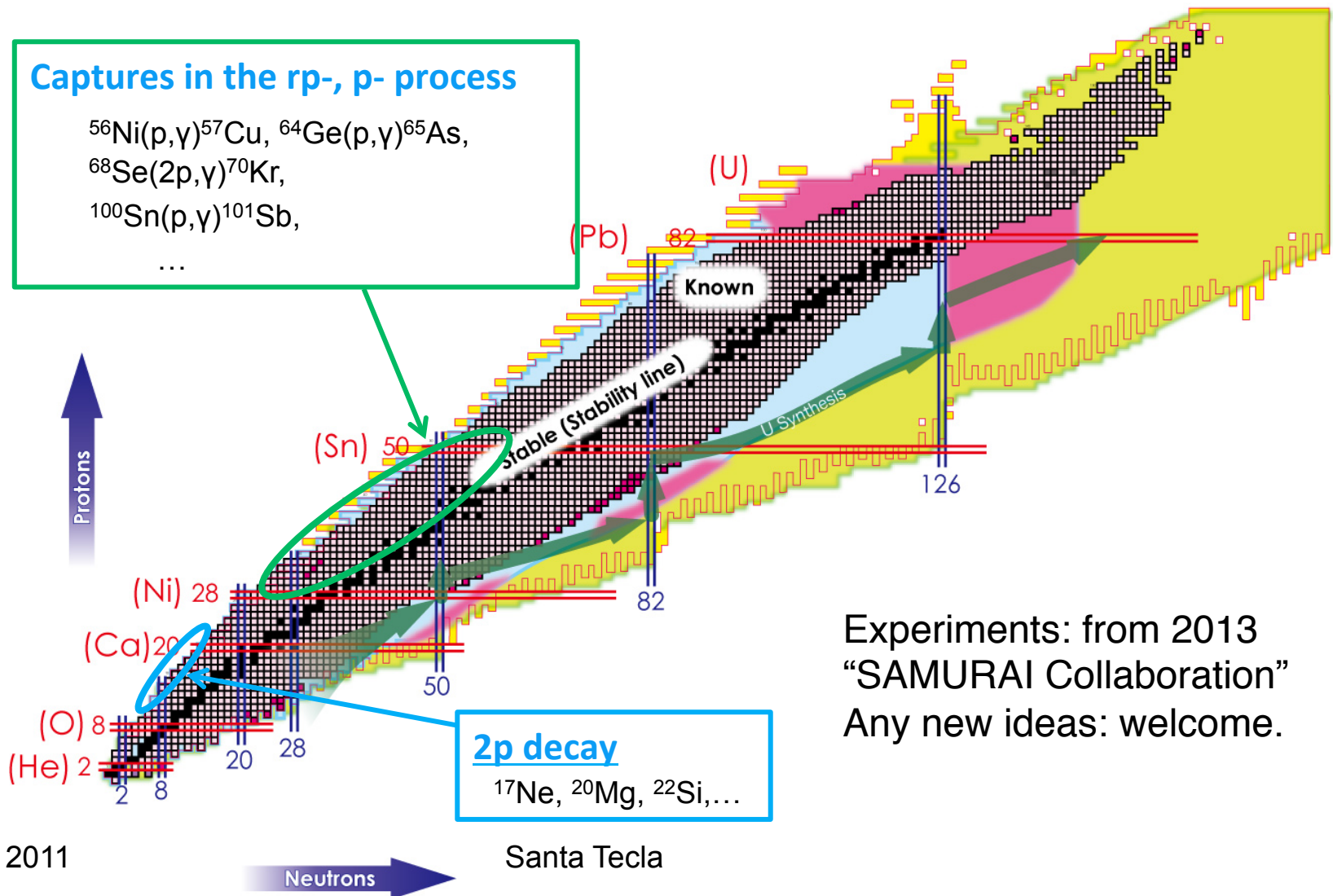
Washington Univ. in St. Louis

HINP16C --- 16ch processing in 1 chip  
two output for energy and timing

**Capability for 2p capture cases**



# Possible experiments in the first stage



Experiments: from 2013  
 “SAMURAI Collaboration”  
 Any new ideas: welcome.

# Toward the r-process path at the RIBF new facility: 1<sup>st</sup> $\beta$ decay measurements

2009 Dec.

U beam to access  $A \sim 110$  region

Intensity 0.8 p nA max.

0.1-0.2 p nA on average

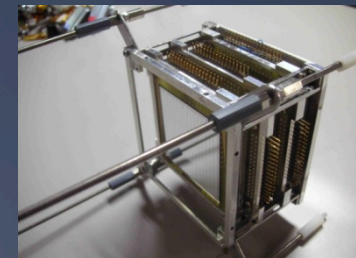
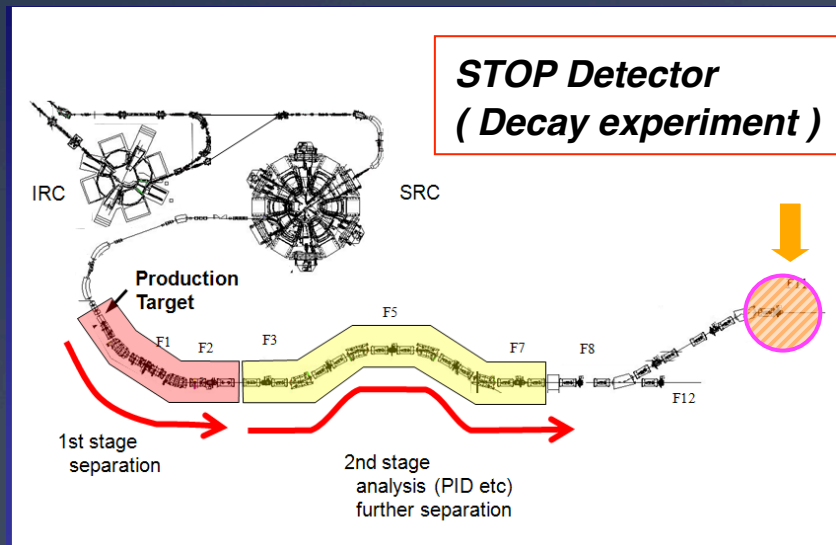
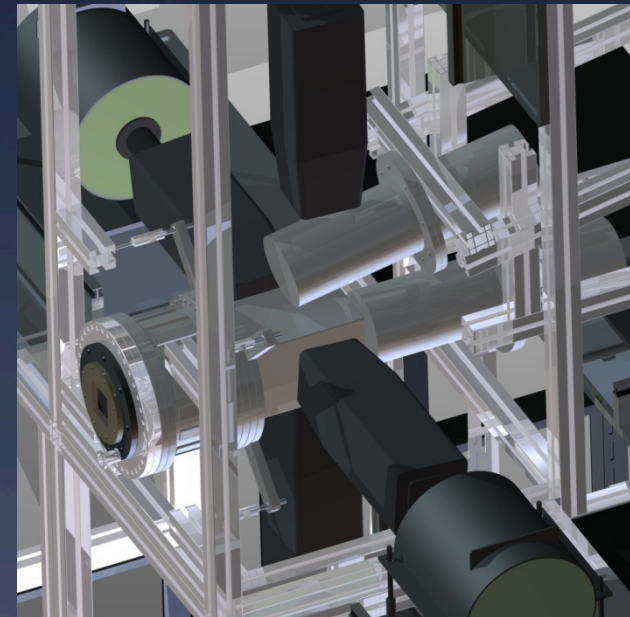
(1/10000 – 1/1000 of the RIBF goal)

Half life measurements for r-process nuclei

Beta-gamma spectroscopy

Delayed gamma spectroscopy for isomers

Nishimura, Sumikama, *et al.*



Clovers (RIKEN)

LaBr<sub>3</sub> (Milano)

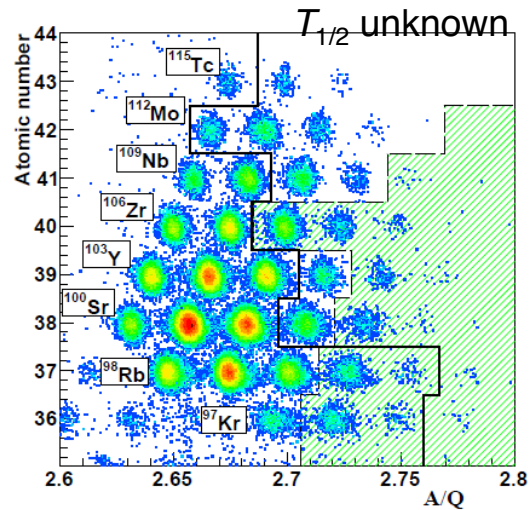
9 layers of DSSD (RIKEN, TUS)



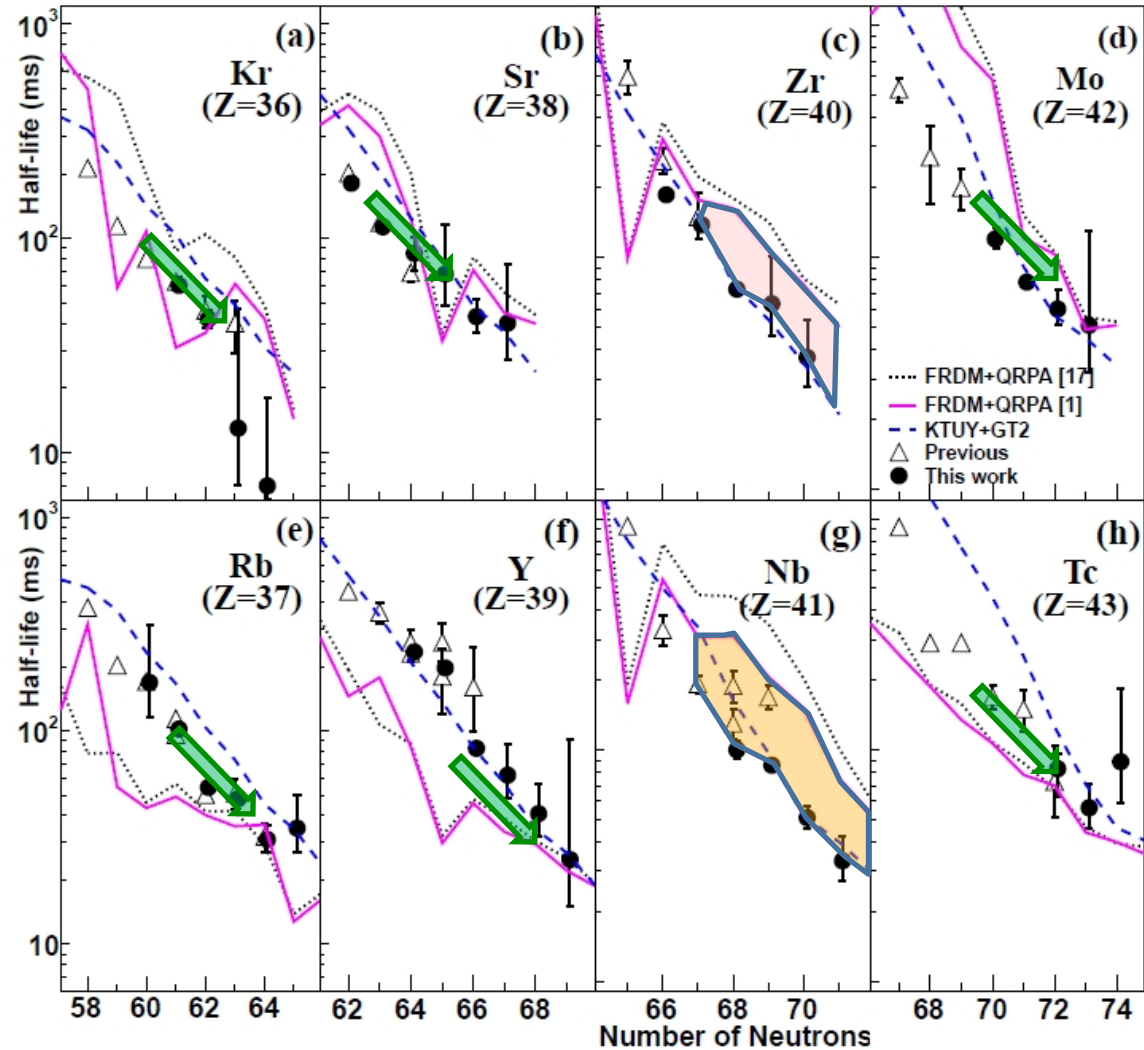
# Toward the r-process path at the RIBF new facility - $\beta$ decay half-life

Systematic studies of  $T_{1/2} \leftrightarrow \text{Mass}, Q_{\beta}, S_n$

Nishimura *et al.* PRL106, 052502 (2011)



- 8 hours data acquisition
- $T_{1/2}$  data of 38 isotopes including first data for 18 isotopes



$T_{1/2}$  are shorter for Zr/Nb compared with FRDM theories!

## Summary

Cross section determination is important in light nuclei (low-level density).

Coulomb dissociation, an indirect tool for astrophysical capture processes, is powerful for the cases involving unstable nuclei.

SAMURAI at RIKEN RIBF will provide new opportunities of astrophysical  $(p,\gamma)$  and  $(n,\gamma)$  studies (Coulomb dissociation).

Studies of near and in the r-process path are to be made at RIKEN RIBF, which has currently the world highest capability of RI beam production.

Visit <http://www.rarf.riken.jp/Eng/facilities/RIBF.html>.  
Join “collaborations”.





Thank you!

Grazzii

ありがとうございました