

Sexten school "Hands-on the Extreme Universe with High Energy Gamma-ray data" 18-24 July 2022

PARTICLE ACCELERATION IN

PULSACS

AND

PULSAR WIND NEBULAE

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SUMMARY

PULSARS AND THEIR MAGNETOSPHERES THE GOLDREICH AND JULIAN MAGNETOSPHERE GAPS, PLASMA SUPPLY, MULTIPLICITY 0 THE PULSAR WIND

PULSAR WIND NEBULAE

- Ø DYNAMICS
- PARTICLE ACCELERATION 0
- **RECENTS FROM GAMMA-RAYS** 0
- PARTICLE ESCAPE

Pulsar -> pulsar wind -> pulsar wind nebulae

IN CRAVITATIONAL COLLAPSE

CONSERVATION OF ANGULAR MOMENTUM CONSERVATION OF MAGNETIC FLUX



 $M_f R_f^2 \ \Omega_f = M_i R_i^2 \Omega_i$

 $B_i R_i^2 = B_f R_f^2$

$P_i \sim 10^5 \text{s}$ $P_f < 0.1 \text{ms}$ $B_i \sim 100 \text{G}$ $P_f < 0.1 \text{ms}$ $B_f \sim 10^{12} \text{G}$





RELATIVISTIC EFFECTS $v_{\rm rot} = \omega R_{\rm NS} \simeq 0.2c \left(\frac{P}{10^{-3}{\rm s}}\right)^{-1} \left(\frac{R_{\rm NS}}{10\,{\rm km}}\right)$

2GM $R_{\rm GR} = \frac{2000}{c^2}$

EXTREME ENVIRONMENT

QED EFFECTS

MAGNETARS....





NEUTRON STARS

•1933: DISCOVERY OF THE NEUTRON (Chadwick 1933)

•1934: NEUTRON STARS PROPOSED (Baade & Zwicky 1934) AS EXPLANATION FOR SUPERNOVA ENERGY RELEASE

•1939: FIRST MEANINGFUL NEUTRON STAR EQUATION OF STATE (Tolman, Oppenheimer & Volkoff 1939): MAXIMUM MASS SMALLER THAN FOR WHITE DWARFS

•1939–1959: NO HOPE OF OBSERVATION -> FORGOTTEN PROBLEM

•1959: REVISED MASS ESTIMATE (Cameron 1959)

•1967: A BRILLIANT IDEA...



"ENERGY EMISSION FROM A NEUTRON STAR" Pacini 1967

Scargle 1967





 $\dot{E} = \frac{B_{\star}^2 R_{\star}^6 \Omega^4 \sin^2 \chi}{6c^3}$





Franco Pacini [1939-2012]



OBLIQUE ROTATING DIPOLE



ENERGY LOSSES:

 $\overrightarrow{\mu} = \mu_0 \left[\sin \alpha \left(\cos \Omega t \ \underline{\mathbf{e}}_1 + \sin \Omega t \ \underline{\mathbf{e}}_2 \right) + \cos \alpha \ \underline{\mathbf{e}}_3 \right]$

Photons

$\overrightarrow{B} = \frac{3\underline{\mathbf{e}}_{\mathbf{R}} (\overrightarrow{\mu} \cdot \underline{\mathbf{e}}_{\mathbf{R}}) - \overrightarrow{\mu}}{R^{3}} \qquad \mu_{0} = \frac{B_{\star} R_{\star}^{3}}{2}$

LARMOR FORMULA

 $\frac{dE}{dt} = -\frac{2}{3c^3}\dot{\mu}^2$

 $\dot{E} = \frac{2}{3c^3} \frac{B_{\star}^2 R_{\star}^6}{\Lambda} \Omega^4 \sin^2 \alpha$

 $\dot{E} = 10^{40} \frac{B_{12}^2}{D_4} \text{ erg/s}$



Hewish & Bell 1968



P=1.33s AND KEEPING SIDEREAL TIME

2GM

 R^2

IF A ROTATING OBJECT:

 $< v_{\rm esc} \rightarrow \Omega R$

Vrot

FOR A WHITE DWARF: P>10 S





$\dot{E} = \frac{B_{\star}^2 R_{\star}^6 \Omega^4 \sin^2 \chi}{6c^3}$

DIPOLE

 $K_{\rm sd} = \frac{6 I_{\star} c^3}{B_{\star}^2 R_{\star}^6 \sin^2 \gamma}$

 $\frac{B_{\star}^2 R_{\star}^6 \Omega^4 \sin^2 \chi}{6c^3} = I_{\star} \Omega \dot{\Omega}$

GENERIC MULTIPOLE



NOTE



PULSAR SPIN DOWN

ROTATIONAL ENERGY LOSS

 $\dot{E} = \frac{E_0}{(1 + t/\tau)^{\frac{n+1}{n-1}}}$ $n = \frac{\ddot{\Omega}\Omega}{\dot{\Omega}^2}$

 $n = 3 + \frac{\Omega}{\dot{\Omega}} \left(2\frac{\dot{B}_{\star}}{B_{\star}} + 2\dot{\chi}CoTg\chi + 6\frac{\dot{R}_{el}}{R_{el}} - \frac{\dot{I}}{I} \right)$





Abdo+ 2013



FOR n = 3(dipole)

 $\dot{E} = 4\pi^2 I_{\star} \dot{P} P^{-3}$

 $\dot{E} \approx 5 \times 10^{31} \text{erg/s} P^{-3} \dot{P}_{-15}$

 $B_{\star} = \frac{(6I_{\star}c^3)^{1/2}}{2\pi R_{\star}^3} \sqrt{P \dot{P}}$ $B_{\star} \approx 10^{12} G_{\rm V} / P \dot{P}_{-15}$

$\tau = \frac{1}{2\dot{P}} \approx 3 \times 10^7 \text{yr } P \dot{P}_{-15}^{-1}$



PULSARS AS UNIPOLAR INDUCTORS



NO VACUUM PHYSICS

EVEN AN ALIGNED DIPOLE SPINS DOWN

 $\dot{E} = \frac{\mu^2 \Omega^4}{c^3} \left(1 + \sin^2 \chi\right)$ firstly obtained by means of numerical simulations [Spitkovski 2006]

while aligned rotator in vacuum:



PULSAR: VACUUM SOLUTION OF ALIGNED ROTATOR



 $\overrightarrow{\Omega} = \Omega \underline{\mathbf{e}}_{z}$ $\overrightarrow{\mu} = \mu \underline{\mathbf{e}}_{z}$ $\mu = \frac{B_{\star} R_{\star}^{3}}{2}$

 $\overrightarrow{E} + \overrightarrow{v} \wedge \overrightarrow{B} = 0$

IF VACUUM OUTSIDE $\nabla^2 \Phi = 0 \qquad \overrightarrow{E}^{out} = -\overrightarrow{\nabla} \Phi$

 $\Phi \rightarrow \overrightarrow{E}^{out}$

 $\left(E_r^{out} - E_r^{in}\right)_{R_*} = 4\pi\sigma_e \neq 0$

 $\overrightarrow{F}\cdot\overrightarrow{R}$ $P_{R_{\star}}$ $eE_{\parallel} = e \frac{B_{\star} \Omega R_{\star}}{C} \cos^2 \theta \approx \frac{0.2}{D} erg/cm$

FREE CHARGES IN THE STAR ROTATE WITH $\overrightarrow{v} = \Omega R \underline{\mathbf{e}}_{\phi}$

$\overrightarrow{E}^{in} = -\left(\frac{\overrightarrow{\Omega} \wedge \overrightarrow{R}}{c}\right) \wedge \overrightarrow{B}$

BOUNDARY CONDITION AT STAR SURFACE $E_{\theta}^{in}(R_{\star}) = E_{\theta}^{out}(R_{\star})$

FINITE CHARGE SURFACE DENSITY AT STAR SURFACE

 $\frac{eE_{\parallel}}{GM_{\star}m_{e}/R_{\star}^{2}} \approx 8 \times 10^{12} \frac{B_{12}}{P_{100}} \quad \text{AND A LARGE ENOUGH}$ FIELD TO EXTRACT IT!





THE GOLDREICH AND JULIAN MAGNETOSPHERE



PARTICLES FLOWING ALONG OPEN FIELD LINES MAY REACH INFINITY

STAR DEVELOPS COROTATING MAGNETOSPHERE $\overrightarrow{E} = - \frac{\Omega R \sin \theta}{\Theta}$ $\mathbf{e}_{\phi} \wedge \overrightarrow{B}$ THROUGHOUT COROTATION REGION

COROTATION ONLY POSSIBLE UNTIL v < c

LIGHT CYLINDER

C







THE PULSAR POTENTIAL DROP

• Φ_{pc} is the actual potential available

• Φ_{pc} is a "measured" quantity, once you MEASURE \dot{E}

•PULSARS CAN EASILY REACH THE KNEE: -FOR CRAB PERIOD OF 33ms $E_{max} \approx 60 \ PeV$

•NEW BORN (FAST SPINNING) MAGNETARS CAN BE ZEVATRONS IN PRINCIPLE

•POTENTIAL DROPS LARGER THAN Φ_{pc} and up to Some fraction of Φ_{tot} can be achieved in the MAGNETOSPHERE IF FOR SOME REASONS MORE FIELD LINES ARE OPEN



POTENTIAL DROP BETWEEN POLE AND INFINITY $\Delta \Phi_{\text{tot}} = -\frac{B_{\star} \Omega R_{\star}^2}{3c} = \frac{R_L}{R_{\star}} \Delta \Phi_{pc}$ $E_{max} = e\Phi = 2 \times 10^{17} \frac{B_{12}}{P_{100}^2} \text{ eV}$





Goldreich & Julian 1969



VACUUM FIELD NOT SELF-CONSISTENT:

- NEGLECT OF DISPLACEMENT CURRENT (Deutsch 1955)
- NEGLECT PARTICLE INERTIA
- NEGLECT OF MONOPOLE TERM (Michel 1969)

• CHARGE SUPPLY UNCLEAR

CHARGE SEPARATED FLOW OR QUASI NEUTRAL PLASMA?

 $G M_{\star} m_e \approx 10^9 \text{ K}$ $R_{\star}k_{R}$



WHY DOES AN ALIGNED ROTATOR SPIN DOWN?



 $\dot{N}_{GJ} = \frac{\Omega B}{2\pi ce} \left(\pi R_{\star}^2 \frac{R_{\star}}{R_{\star}} \right) c \approx \frac{\sqrt{c} \dot{E}}{\frac{\rho}{2}}$ $E_{drop} = e\Delta\Phi_{pc} = e\sqrt{\dot{E}/c}$

FORCE-FREE SOLUTION: accounting for particles in the magnetosphere $\dot{E} = \frac{\mu^2 \Omega^4}{c^3} \left(1 + \sin^2 \chi\right)$

ENERGY LOSS BY A GJ FLUX OF PARTICLES LEAVING THE STAR POLAR CAP AT PULSAR DROP

 $\dot{E}_{\text{part}} = \dot{N}_{GJ} E_{drop}$

 $\dot{E}_{\text{part}} = \dot{E}$





POLAR CAPS



$J = J_{\star} \frac{B}{B_{\star}}$

CHARGE DENSITY LOCALLY NEEDED CAN BE SUPPLIED FROM THE STAR ONLY FOR FIELD LINES WITH DECREASING B_{\parallel}/B a fraction of the total for oblique rotators

• UNSCREENED ELECTRIC FIELD IS LEFT IN ALL CASES IF PARTICLES FROM STAR SURFACE ONLY







SPACE CHARGE GAPS AND VACUUM GAPS





 $\boldsymbol{\Omega} \bullet \boldsymbol{B} > 0$





Usov & Melrose 1995 Harding 2007

Charged particles are bounded to the star due to lattice structure in strong magnetic field

> Thermoionic emission temperatures:

 $T_i \approx 3.5 \times 10^5 \text{ K} \left(\frac{B_{\star}}{10^{12} \text{G}}\right)^{0.73}$

Particles below those temperature are bounded to the star



MAGNETIC PAIR PRODUCTION



IN VACUUM PAIR LIVES ONLY $\Delta t \sim \hbar/\Delta E$

 $B_{\rm OED}$

В

PAIR BECOMES REAL IF $D > \lambda_C = \frac{h}{m_e c}$ $D \approx \frac{\lambda_C}{4\pi} \left(\frac{h\nu}{m_e c^2}\right) \left(\frac{B}{B_{OED}}\right) = \chi \frac{\lambda_C}{4\pi}$

$$\chi = \frac{\epsilon_{\gamma}}{m_e c^2} \frac{B}{B_{\text{QED}}} = 0.4 \left(\frac{\epsilon_{\gamma}}{10 \text{MeV}}\right) \left(\frac{B}{10^{12} \text{G}}\right)$$

PAIR CREATION LENGTH:

 $h\nu > 2m_ec^2 \Rightarrow \frac{\text{PHOTON CAN}}{\text{CREATE PAIR}} \begin{cases} h\nu \to 2m_ec^2\gamma \\ h\nu/c \to 2m_ec^2\gamma\beta \end{cases} \xrightarrow{\Delta E} \sim \frac{2(m_ec^2)^2}{h\nu}$

IF B-FIELD $\Delta \phi = \frac{c \Delta t}{---} \Rightarrow D \sim c \Delta t \Delta \Phi$

 $\chi \gg 1 \Rightarrow$ EXTREMELY EFFECTIVE PAIR CREATION $\chi < 1 \Rightarrow$ PAIR CREATION SUPPRESSED AS $exp(-4/3\chi)$

> RAPIDLY CHANGING QUANTITY IN THE MAGNETOSPHERE









- closed magnetosphere
 - base of the wind (open field lines)

Petri 2017

DIAGNOSTICS OF THE CASCADE GAMMA-RAYS MULTIPLICITY DEATHLINE



γ**-RAYS**

$L_{\rm radio} \leq 10^{-10} \dot{E}$

$$L_{\gamma} \sim 10^{-2} E$$



DEATHLINE AND MULTIPLICITY



Harding 2007

Peak of curvature radiation photon energy $\epsilon_{\gamma,\text{CR}} = 3\lambda_C \gamma^3 / (2\rho_c)$

Photon energy of Inverse Compton in Klein-Nishina limit:

 $\epsilon_{\gamma,\mathrm{IC}} \sim \gamma$

 $\Rightarrow \epsilon_{\gamma, CR} \ll \epsilon_{\gamma, IC}$

Pair production of CR photons require larger Lorents factors than IC

HIGHEST MULTIPLICITY FOR SYNCHROTRON/ CURVATURE POWERED GAPS: K< few x 10⁵

BUT.... PWNe



THE PULSAR WIND

THE ALTONE



Contopoulos, Kazanas, Fendt 2001



ALIGNED ROTATING DIPOLE VERY SIMILAR TO SPLIT MONOPOLE $\overrightarrow{B} = B_L \left[\left(\frac{R_L}{R} \right)^2 \underline{\mathbf{e}}_{\mathbf{R}} - \left(\frac{R_L}{R} \right) \sin \theta \underline{\mathbf{e}}_{\phi} \right]$ $\vec{E} = -B_L\left(\frac{R_L}{R}\right)\sin\theta \,\underline{\mathbf{e}}_{\theta}$

YET ANOTHER WAY TO LOOK AT THE STAR SPIN-DOWN: POYNTING FLUX THROUGH THE LIGHT CYLINDER SURFACE:

$$\dot{E} = \pi S_L R_L^2 \qquad S_L = \frac{B_L^2}{4\pi} c \pi R_L^2 \qquad B_L = B_\star \left(\frac{R_\star}{R_L}\right)$$
$$\dot{E} \approx \frac{B_\star^2 \Omega^4 R_\star^6}{c^3}$$





IONS

ENERGY FLUX THAT LEAVES THE PSR

 $\dot{E} = \kappa \dot{N}_{GJ} m_e \Gamma c^2 \left(1 + \frac{m_i}{\kappa m_e} + \sigma \right)$

ELECTRON-POSITRON PAIRS

CAN WE CONSTRAIN κ AND σ ? WHAT ABOUT IONS ?

ENERCY BUDGET

$$\sigma = \frac{B^2}{4\pi n_+ m_\rho c^2 \Gamma^2}$$

$\kappa =$ multiplicity due to pair creation

MAGNETIC FIELD



TIME DEPENDENT FORCE FREE SOLUTION



Spitkovsky 2006

Current





HE CURRENT SHEET



 $\overrightarrow{\mu} \cdot \underline{\mathbf{e}}_{\mathbf{R}} = \mu \left\{ \sin \theta \sin \chi \cos \phi \cos \Omega t + \sin \theta \sin \chi \sin \phi \sin \Omega t + \cos \theta \cos \chi \right\}$ $\vec{\mu} \cdot \underline{\mathbf{e}}_{\mathbf{R}}$

PATTERN PROPAGATES AT SPEED V

 $\Psi(t,\theta,\phi) = \sin\theta \sin\chi \cos(\phi - \Omega t) + \cos\theta \ \cos\chi = 0$



$\chi = 10^{\circ}$





 $\chi = 30^{\circ}$

$\Psi(t, R, \theta, \phi) = \sin \theta \sin \chi \cos \left| \phi - \Omega \left(t - \frac{R}{V} \right) \right| + \cos \theta \ \cos \chi = 0$





BASIC PICTURE FOR YOUNG SYSTEMS



Adapted from Kennel & Coroniti 1984



unshocked ISM

> log R [pc] SHOCKED EJECTA

UNSHOCKED EJECTA



RADIO (VLA) UV (Astro-1)

THE CRAB NEBULA AT DIFFERENT FREGUENCIES

IR (Spitzer)



X-Ray (Chandra)

Pixel Size

Hard X-Ray (HEFT)

Visible (Hubble)

THE CRAB NEBULA SPECTRUM

BROAD BAND NON-THERMAL SPECTRUM



synchrotron radiation by relativistic particles in the nebular B field Inverse Compton scattering with local photon field





FOR ICS ON CMB $\epsilon_{\gamma} \approx 0.37 \ (E_{\rm e}/{\rm PeV})^{1.3} \ {\rm PeV}$

THE CRAB NEBULA IN CAMMA-RAYS

THE ONLY ESTABLISHED GALACTIC **PEVATRON!!!**

 $E_{\rm e} \approx 2.4 ~{\rm PeV}$

Amato & Olmi 2021







BUT....



WHAT WE KNOW: • MOST EFFICIENT ACCELERATORS IN NATURE $\epsilon_{\rm acc} \lesssim 30\%$

• ENERGY FLUX THAT LEAVES THE PSR

WE DO NOT KNOW:

- WHAT THE ACCELERATION MECHANISM(S) IS (ARE) POSSIBILITIES DEPEND ON

WIND COMPOSITION (IONS? K?) WIND MAGNETIZATION (σ ?)

BIG OPEN GUESTION

 $\dot{E} = \kappa \dot{N}_{GJ} m_e \Gamma c^2 \left(1 + \frac{m_i}{\kappa m_e} + \sigma \right) \qquad \sigma = \frac{B^2}{4\pi n_{\pm} m_e c^2 \Gamma^2}$





Adapted from Kennel & Coroniti 1984 [Del Zanna & Olmi 2017]

BASIC FICTURE OF A PWN



 R_{TS} = termination shock radius R_N = nebula radius

 $\dot{E} t$ Ė $-=P_{PWN}=$ $4\pi cR_{TS}^2$ $4\pi R_N^3$ $R_{TS} = \left(\frac{\nu_N}{c}\right)^{1/2} R_N$



THE TERMINATION SHOCK



Adapted from Kennel & Coroniti 1984 [Del Zanna & Olmi 2017]



1/2 $\left(\frac{v_N}{c}\right)$ $R_{TS} = ($ R_N

DISSIPATION AND PARTICLE ACCELERATION AT TS



10/20 STATIC MODELS OF PWNE

[Rees & Gunn 1974, Kennel & Coroniti 1984, Emmering & Chevalier 1987, Begelman & Li 1992]

FROM DYNAMICS AND RADIATION MODELING OF OPTICAL /X-RAY EMISSION OF THE NEBULA

- particle spectral index(es) $\rightarrow \gamma = 2.3$
- wind Lorentz factor $\rightarrow \Gamma = 3 \times 10^6$
- wind magnetization $\rightarrow \sigma = v_N/c \approx 3 \times 10^{-3}$
- particle injection rate $\rightarrow \dot{N} \approx 10^{38} s^{-1}$



ACCORE LESSE

POWER-LAW DEVELOPS BUT SLOW PROCESS! scattering on small-scale turbulence: $E_{\rm MAX} \propto t^{1/2}$



ERMI ACCELERCATION RELATIVISTIC UNMAGNETIZED!)



ACCELERATION COMPLETELY SUPPRESSED FOR $\sigma > 10^{-3}$ $E_{\rm MAX} \approx \sigma^{-1/4}$

ERMI ACCELERATION (RELATIVISTIC MACHNETIZED)



IN PRINCIPLE VERY FLAT SPECTRA AT LOW ENERGY

FERMI ACCELERATION IN

RESULTS DEPEND ON













FORCED MAGNETIC RECONNECTION



SUCH LARGE K DIFFICULT TO ACCOUNT FOR

IF REALIZED, RECONNECTION BEFORE THE SHOCK

FORCED MAGNETIC RECONNECTION

INTERACTION WITH X-POINT

DC ACCELERATION

THEN ADVECTION INTO MAJOR ISLANDS

BROAD SPECTRUM

BUT

 $r_L \sigma$

 $\sigma > 30$

 $> few \times 10$ $r_L \sigma$



 $\kappa > few \times 10^7$





RESONANT CYCLOTRON ABSORPTION IN ION DOPED PLASMA shock front



the transition



Configuration at the leading edge ~ cold ring in momentum space

PARTICLE ACCELERATION MECHANISMS: SUMMARY OF REQUIREMENTS

FERMI MECHANISM

DRIVEN MAGNETIC RECONNECTION

ION CYCLOTRON ABSORPTION IN ION DOPED PLASMA





MAGNETIZATION: REQUIRES HIGH

PLASMA MULTIPLICITY: REQUIRES HIGH

PLASMA MULTIPLICITY: REQUIRES LOW

HIGH MAGNETIZATION SPEEDS UP PROCESS