

From Flavor to Top, Higgs, and beyond

Through the language of Effective Field Theories

Electroweak, Strong, and New Interactions:
A symposium to celebrate Guido Martinelli's 70th birthday

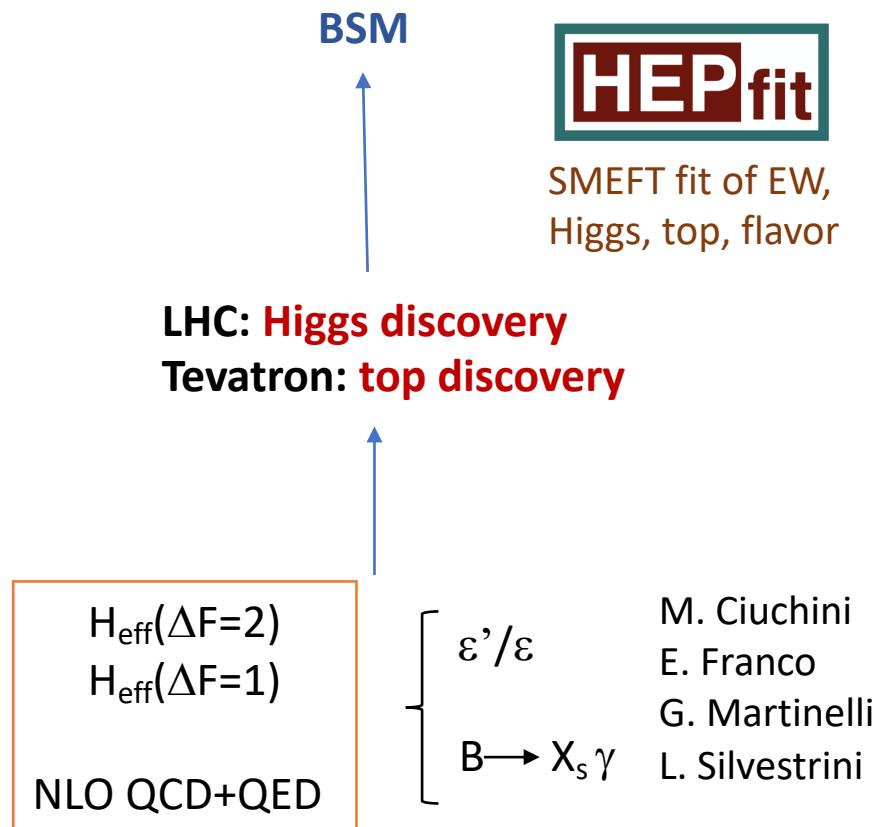
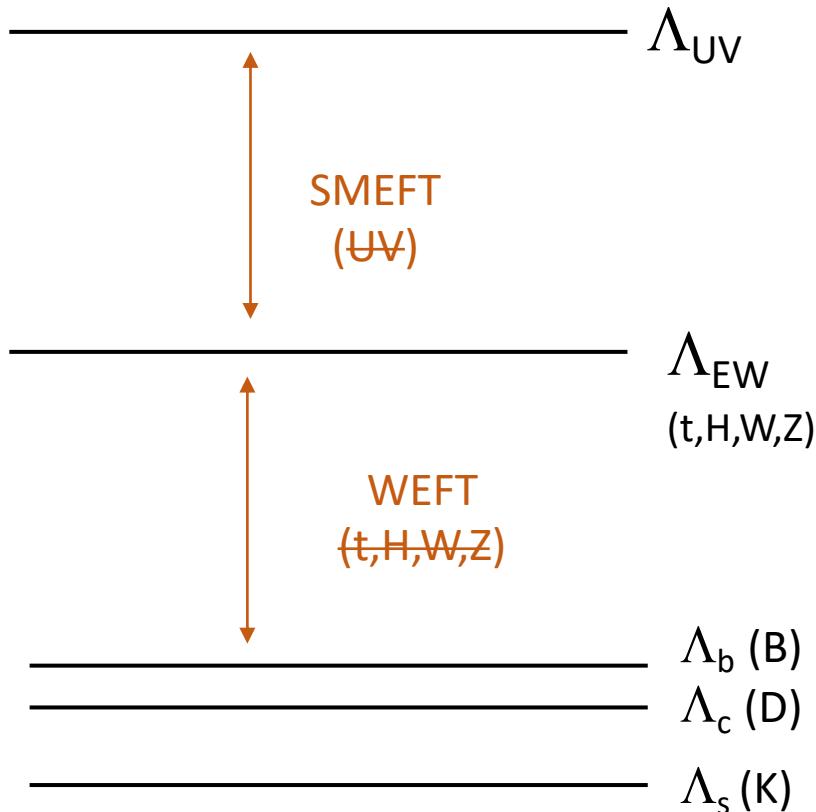
Roma – Accademia Nazionale dei Lincei
26 Settembre 2022



Laura Reina
(Florida State University)

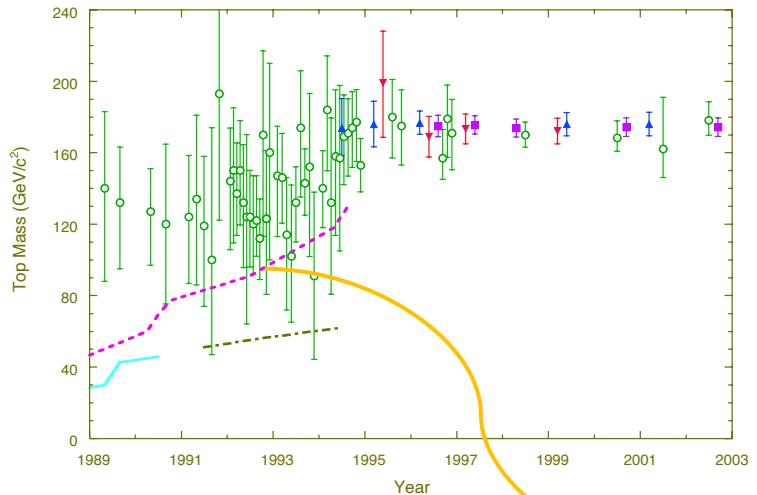


From Flavor to Top, Higgs, and beyond



Top

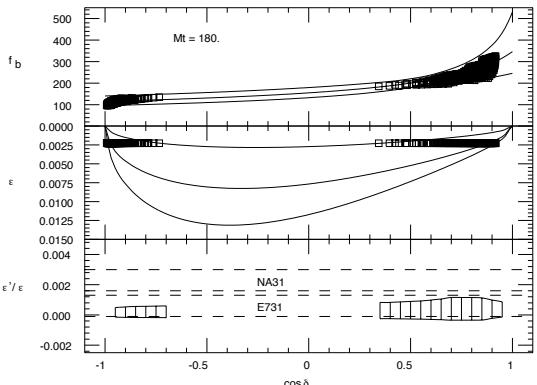
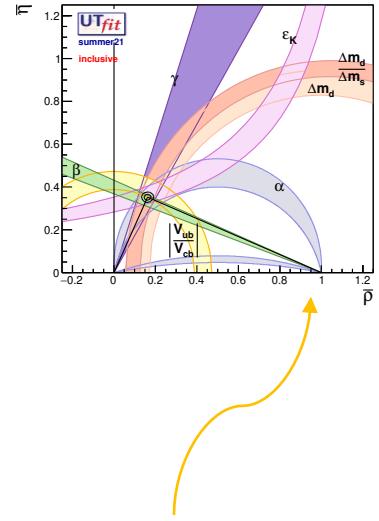
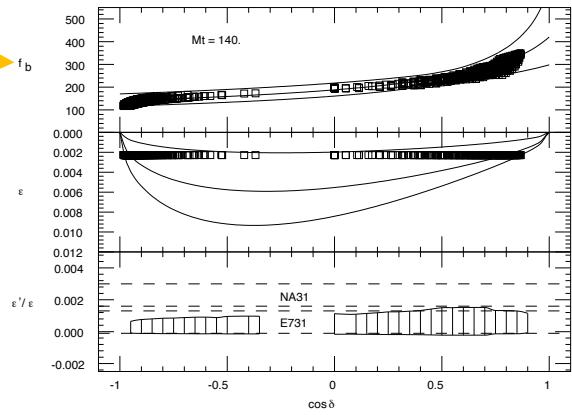
from prediction to discovery



C. Quigg [hep-ph/0404228]

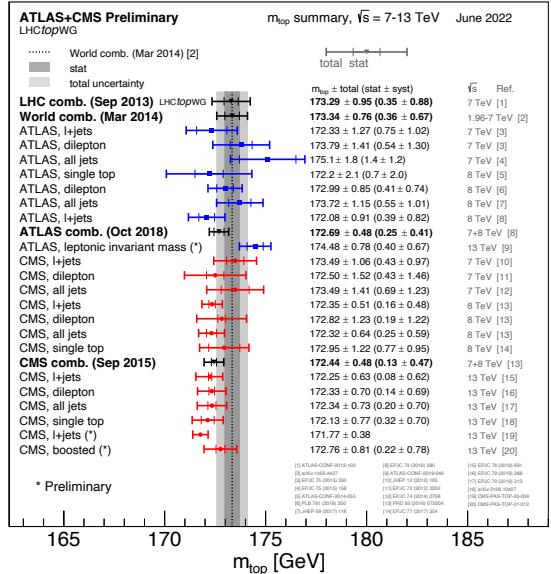
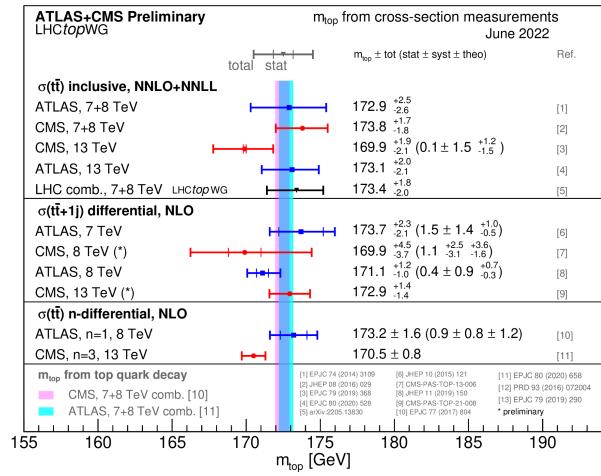
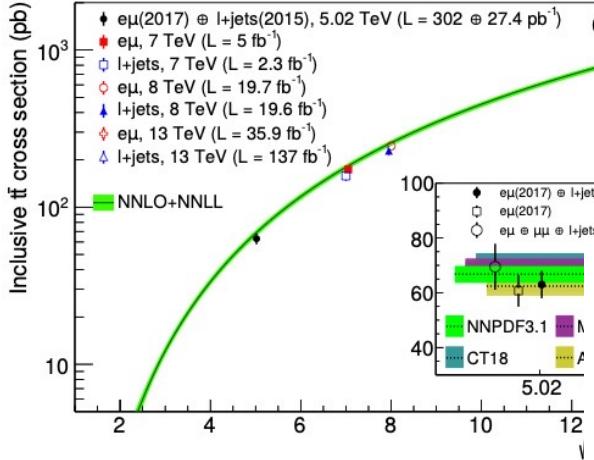
*Effective Hamiltonians
beyond the Leading Order
and their applications*

[L.R. – PhD thesis]



Top

from discovery to precision physics

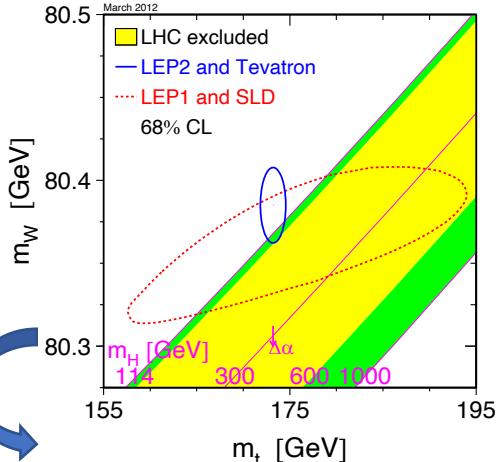


M_t becomes a crucial input in precision fits of the SM (EW and flavor)

Anomalies in Top-quark EW couplings (W,Z,H) possible hint of BSM physics

Higgs

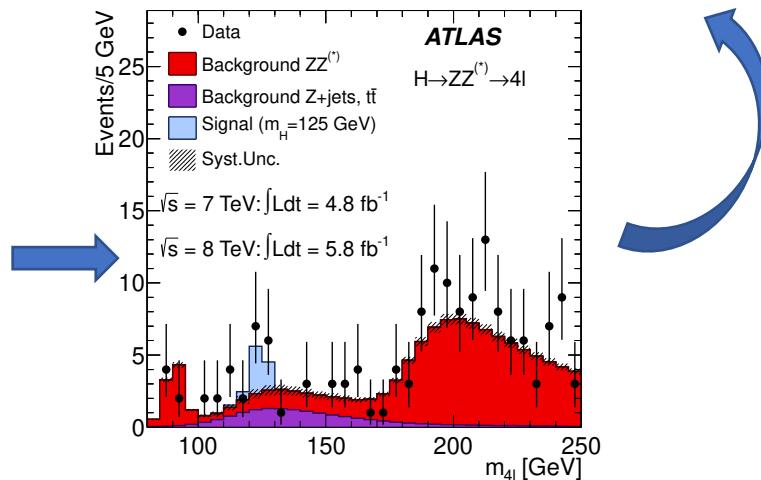
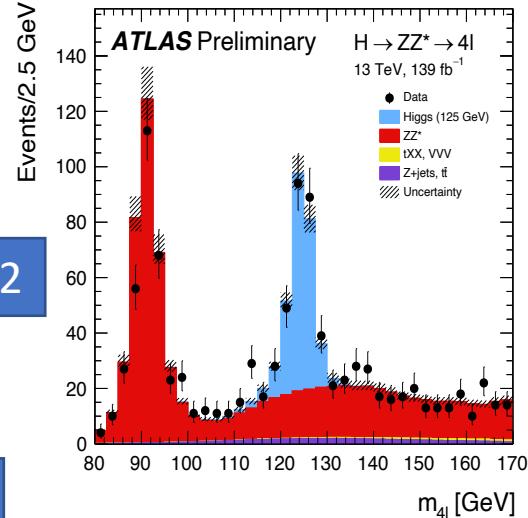
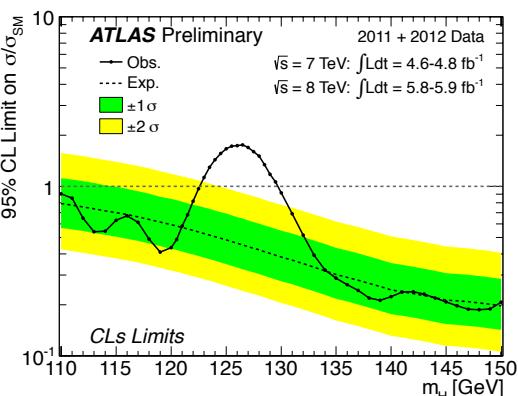
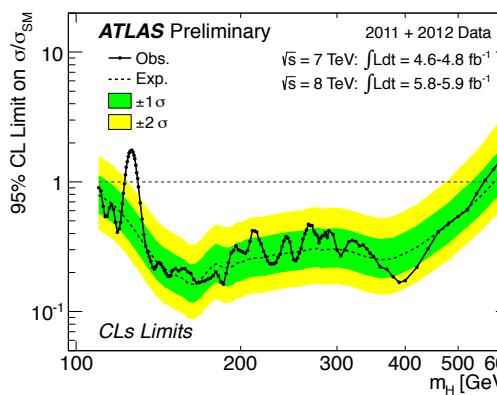
from prediction to discovery



From EW fits
 $M_H = 94^{+29}_{-24}$ GeV
 $M_H < 152\text{--}171$ GeV

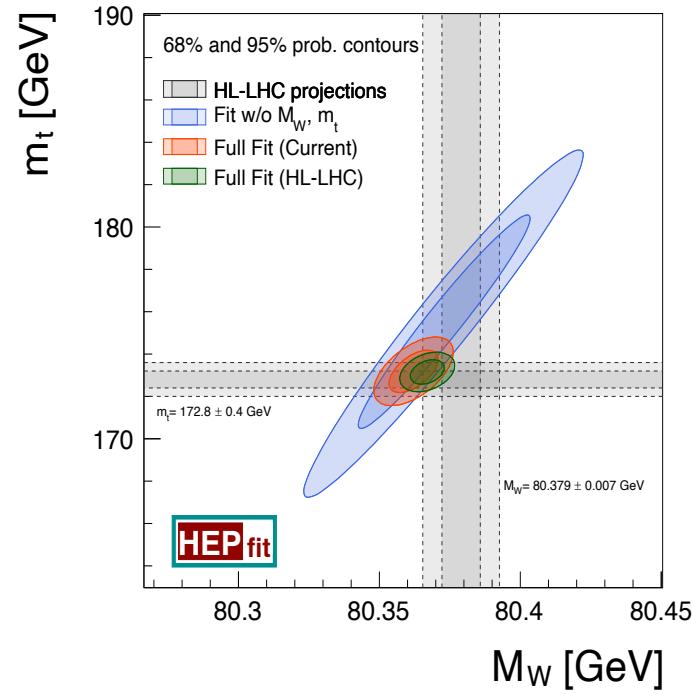
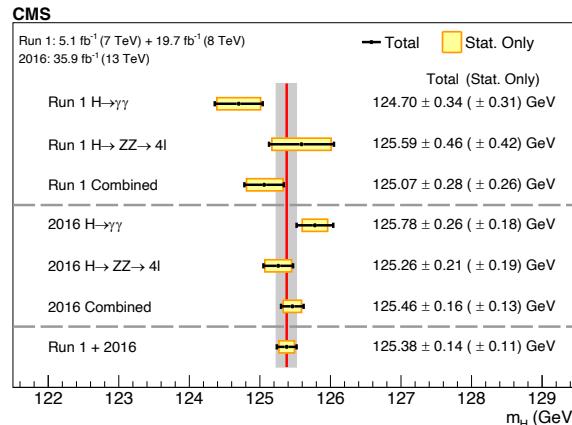
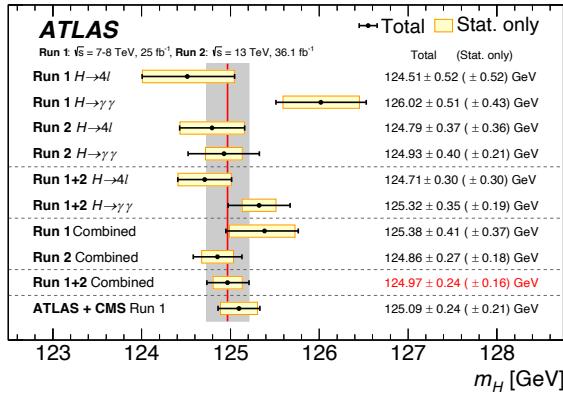
LHC@Run1+2

LHC@Run1: $M_H = 125.09 \pm 0.24$ GeV



Higgs

from discovery to precision physics

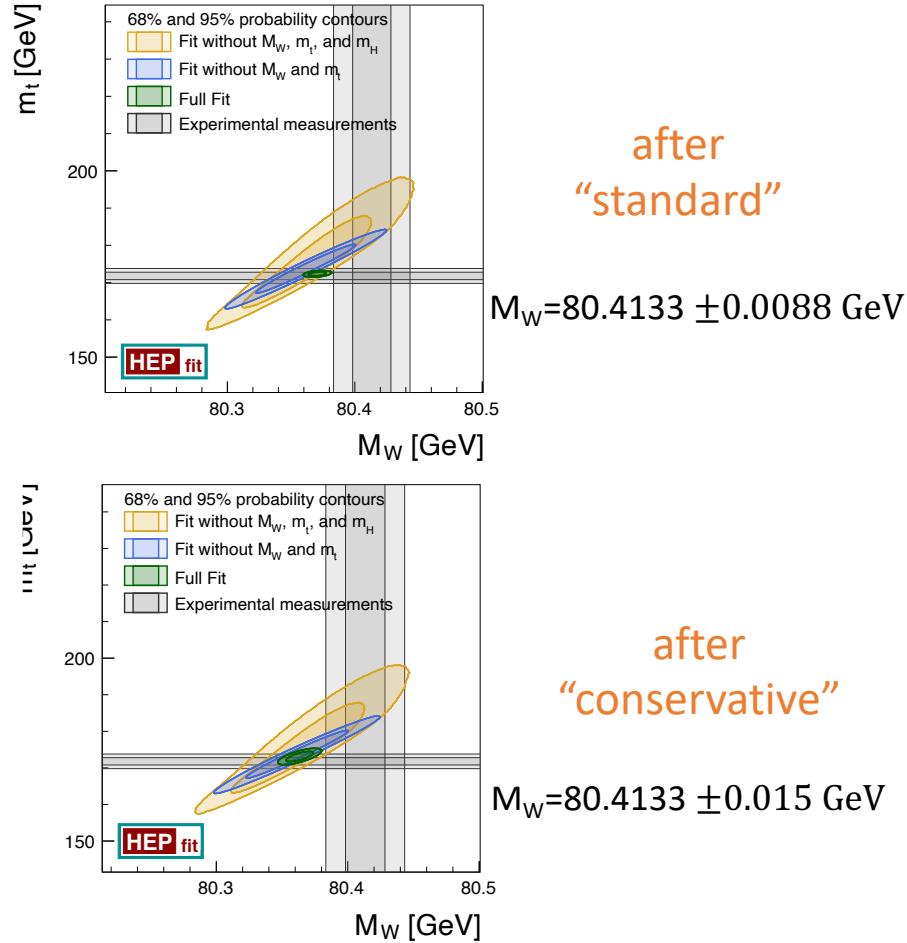
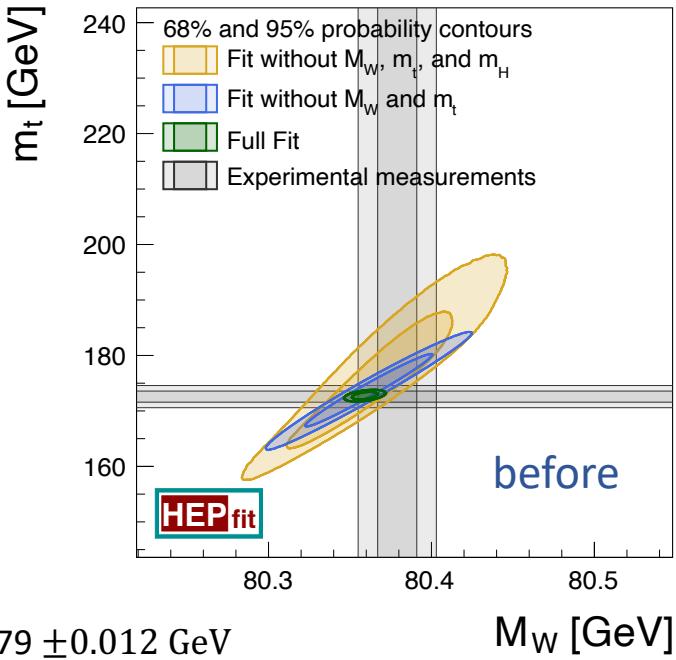


M_H promoted to EW precision observable

Stress-testing the SM

A recent challenge: CDF new M_W measurement

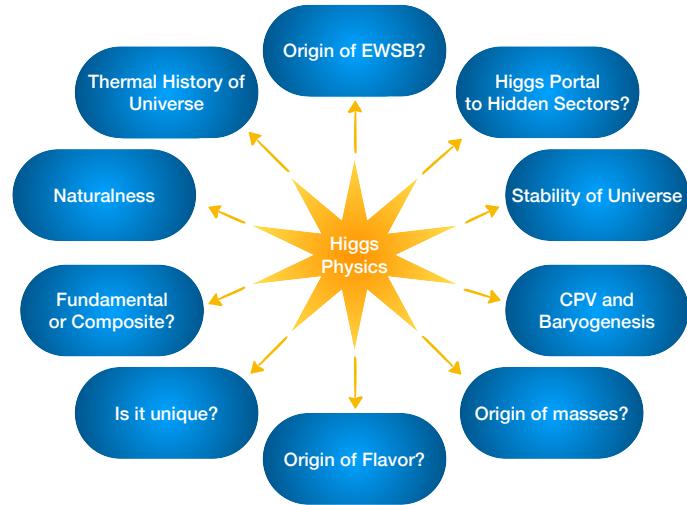
De Blas et al.
[2204.04204]



The big open questions

What is the origin of the EW scale?

The discovery of the Higgs boson has sharpened the big open questions and given us a unique handle on BSM physics.



- Why the $M_H \ll M_{\text{planck}}$ **hierarchy problem?**
- What are the implications for **Naturalness**?
- Can we uncover the origin of BSM physics from precision measurement of Higgs properties (couplings, width, ...). **Elementary vs composit? One Higgs? More?**
- Can we measure the shape of the **Higgs potential** → **Higgs self coupling(s)**
- Can Higgs properties give us **insights on flavor** and vice versa?
 - Couplings to heavy flavors (bottom, **top**, ..)
 - Couplings to light quarks and leptons

The LHC era: exploring the TeV scale



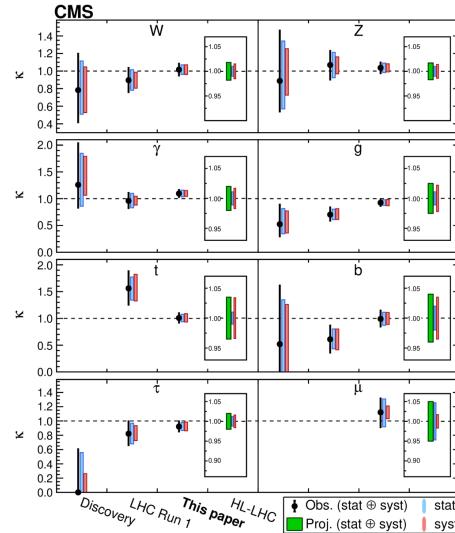
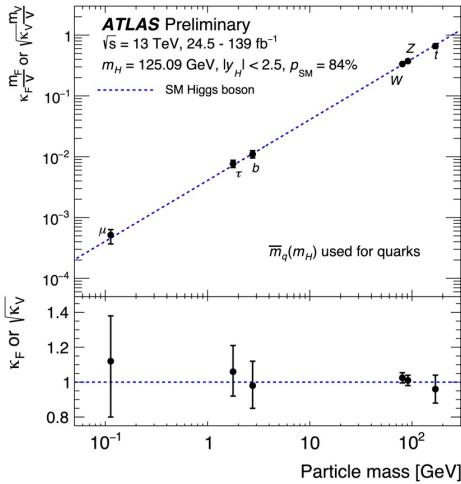
Indirectly via Higgs and Top:

- Run 2 delivery for Higgs couplings outperformed expectations
- LHC will define top physics till the next high-energy collider
 - $e^+e^- > 500 \text{ GeV}$
 - $pp @ 100 \text{ TeV}$
 - $\mu^+\mu^- > 10 \text{ TeV}$

- ➔ 2-fold increase in statistics by the end of Run 3
- ➔ 20-fold increase in statistics by the end of HL-LHC!

Higgs

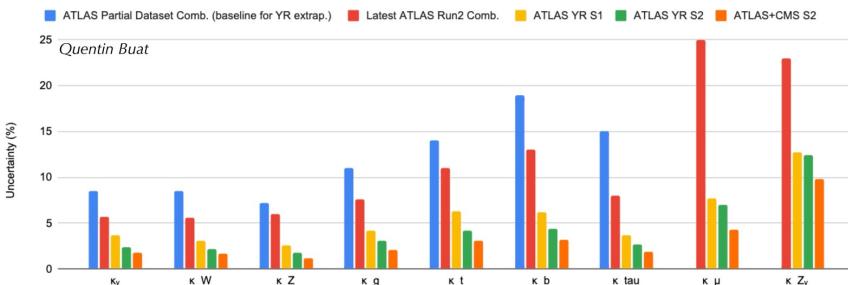
zooming on couplings



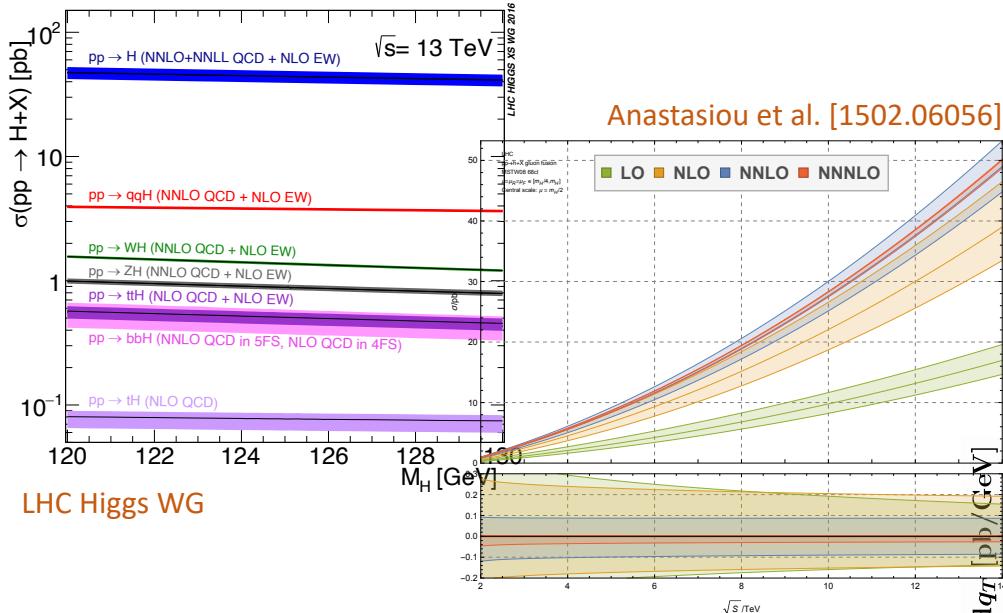
$$\kappa = g_x / g_x^{\text{SM}}$$

- Couplings to W/Z at 5-10 %, couplings to 3rd generation fermions to 10-20%
- First measurements of couplings to 2nd generation fermions
- HL-LHC projections from YR: 2-5 % on most couplings and <50% on Higgs self-coupling

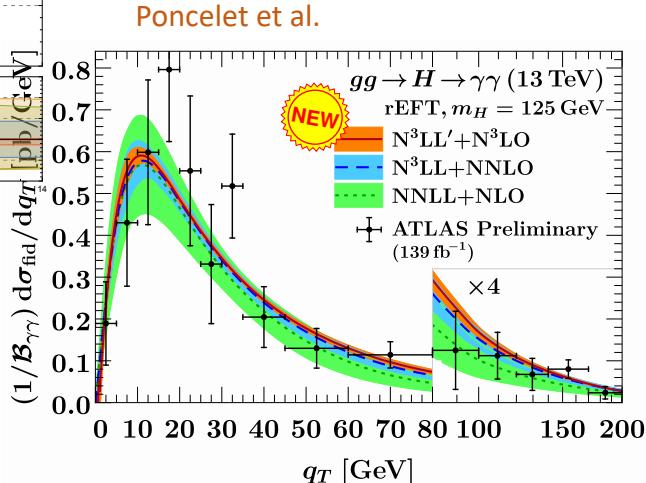
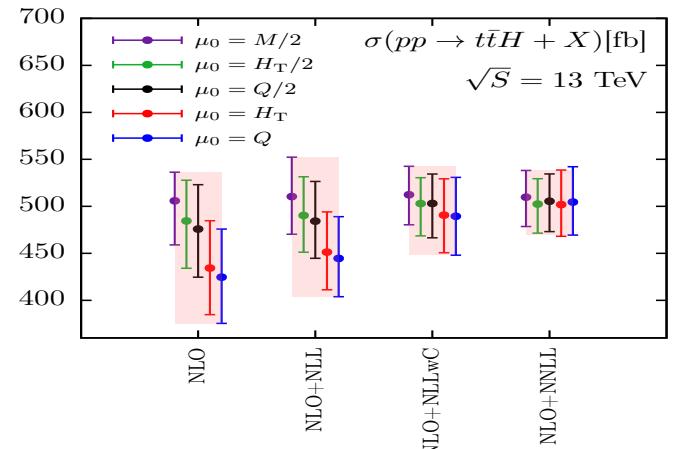
Full Run2 results drastically improve partial Run 2 results (baseline for YR projections)



Theory has come a long way

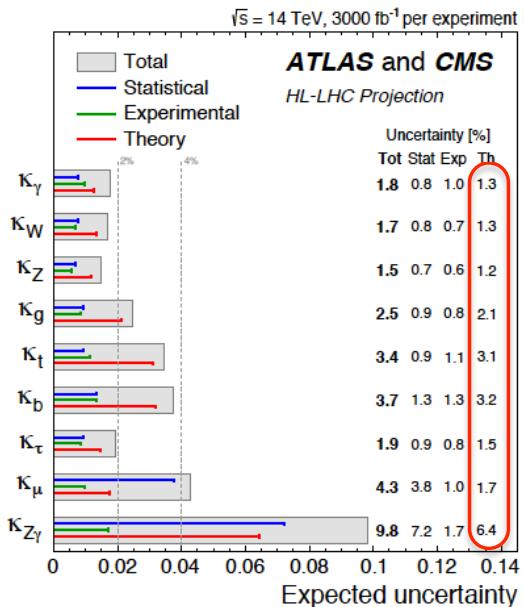
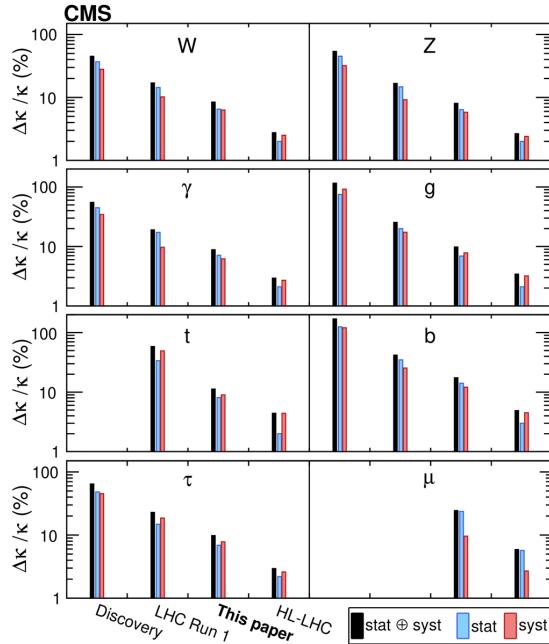


Several backgrounds also known at
NLO QCD+EW or improved NLO (+NNLL)
(e.g. W/Z+j, ttbb, ttW, ttZ, tt γ , ...)



Higgs

zooming on couplings, a little more ...



Generically:
 $\Delta\kappa/\kappa \sim O(v^2/\Lambda^2)$

For new physics at 1 TeV
expect deviations of $O(6\%)$

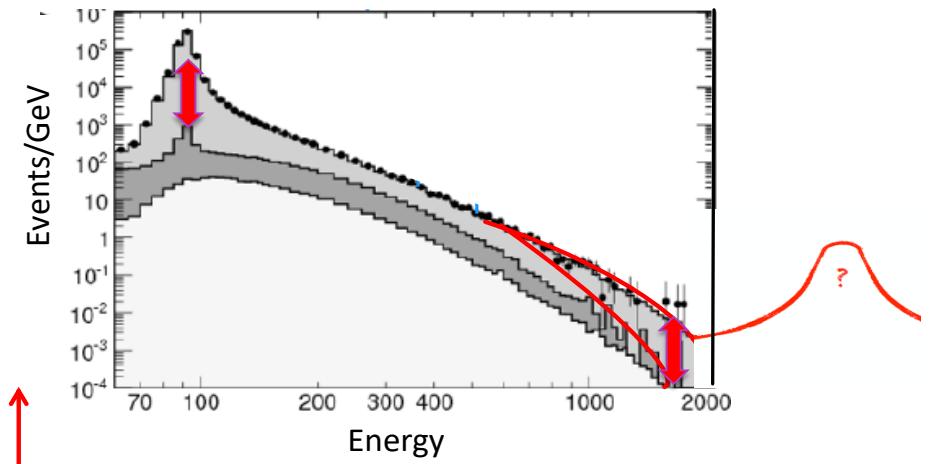
Improved systematics
probes higher scales



Theory could become main
limitation

Theory need to improve modeling and interpretation of LHC events, in particular when new physics may not be a simple rescaling of SM interactions

Beyond total rates

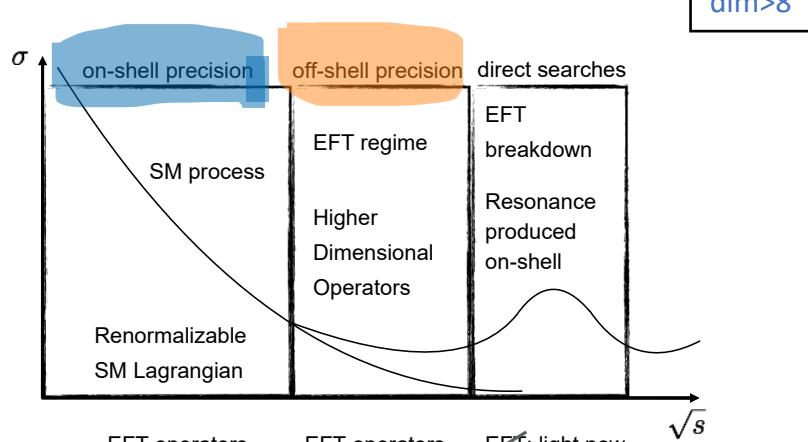


Need SM precision calculations at differential level both at **lower energy**, where rates are large and at **higher energy** where rates are small but effects of new physics may be more visible.

Extending the SM via effective interactions above the EW scale → SMEFT

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \left(\frac{1}{\Lambda^2} \sum_i C_i O_i + \text{h.c.} \right) + O(\Lambda^{-4})$$

dim=6



Examples: EFT operators with Higgses

EFT operators
with derivatives

~~EFT~~: light new
physics

Crucial to control EFT sensitive regions

Ex: Interpreting $t\bar{t}Z$ measurements

Anomalous top couplings

$$\mathcal{L} = e\bar{u}(p_t) \left[\gamma^\mu (C_{1,V} + \gamma_5 C_{1,A}) + \frac{i\sigma^{\mu\nu} q_\nu}{M_Z} (C_{2,V} + i\gamma_5 C_{2,A}) \right] v(p_t) Z_\mu$$

Effective operators

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \left(\frac{1}{\Lambda^2} \sum_i C_i O_i + \text{h.c.} \right) + O(\Lambda^{-4})$$

$$O_{uZ} = -s_W O_{uB} + c_W O_{uW}$$

$$O_{uB} = (\bar{q}\sigma^{\mu\nu} u)(\epsilon\varphi^* B_{\mu\nu})$$

$$O_{uW} = (\bar{q}\tau^I \sigma^{\mu\nu} u)(\epsilon\varphi^* W_{\mu\nu}^I)$$

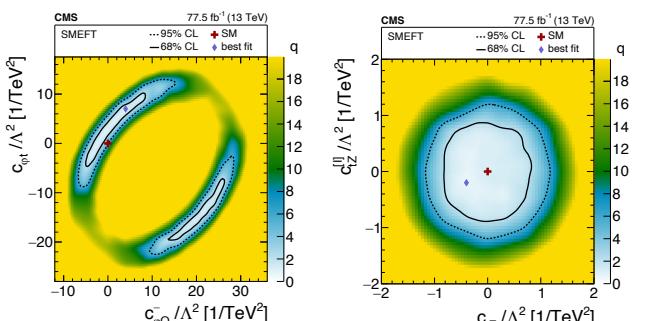
$$O_{\varphi u} = (\bar{u}\gamma^\mu u)(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)$$

$$O_{\varphi q}^- = O_{\varphi q}^1 - O_{\varphi q}^3$$

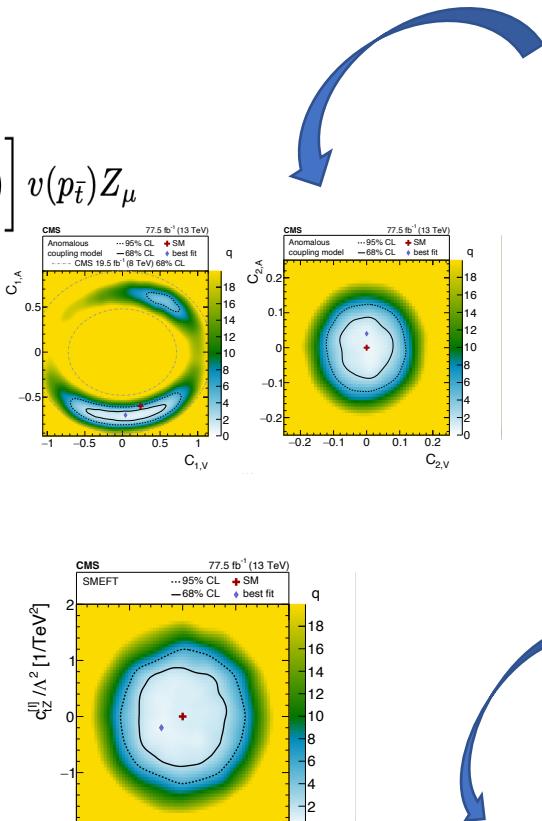
$$O_{\varphi q}^1 = (\bar{q}\gamma^\mu q)(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)$$

$$O_{\varphi q}^3 = (\bar{q}\tau^I \gamma^\mu q)(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)$$

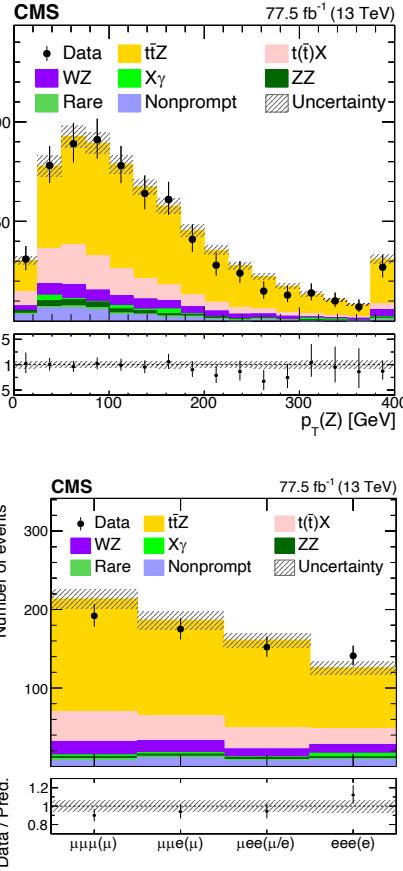
...



CMS [arXiv:1907.11270]



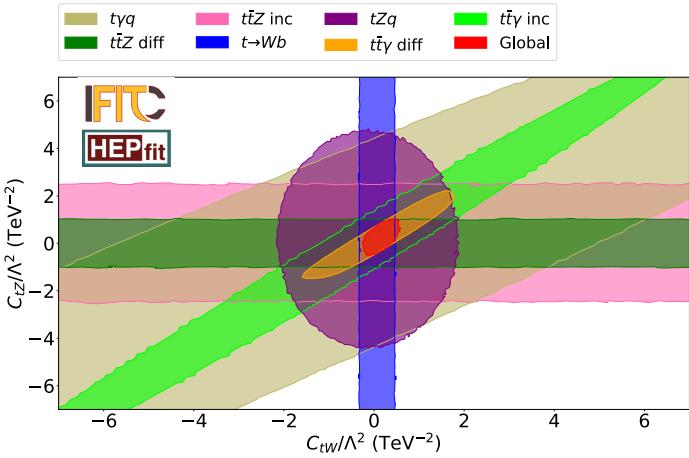
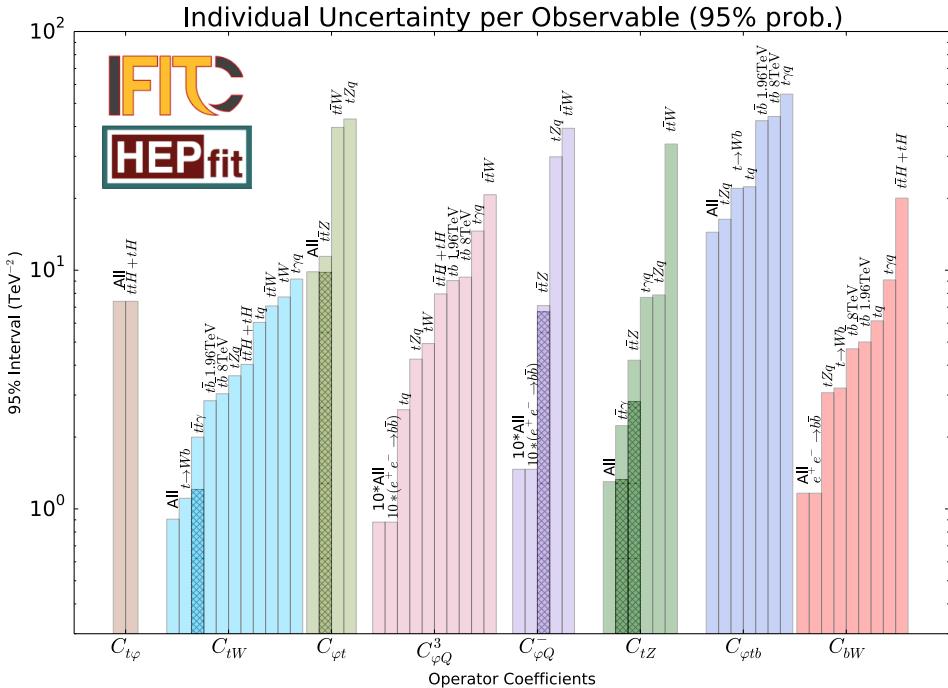
ttZ searches in 3l and 4l signatures



EFT allows multiple probes

Global fits of top observables

V. Miralles, et al. [arXiv:2107.13917]

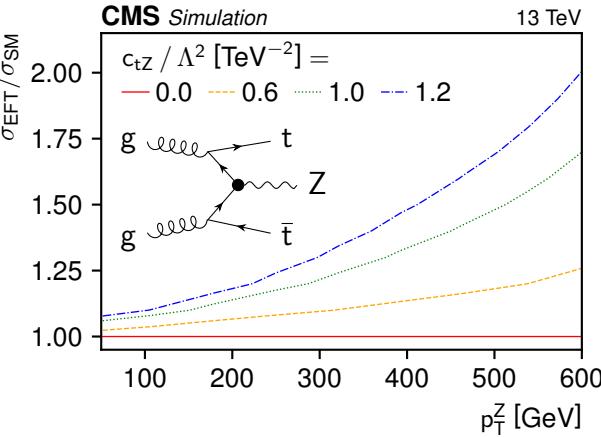


Kinematic distributions add substantial constraining power

Accurate modelling of $\bar{t}tX$ differential cross sections and signatures becomes crucial

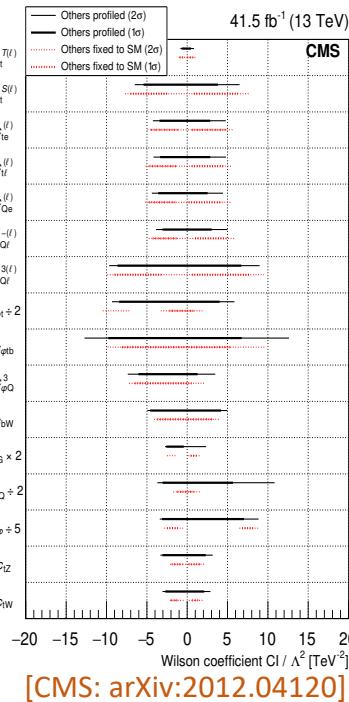
... exploring boosted kinematics and off-shell signatures

Top pair + boosted Z/H

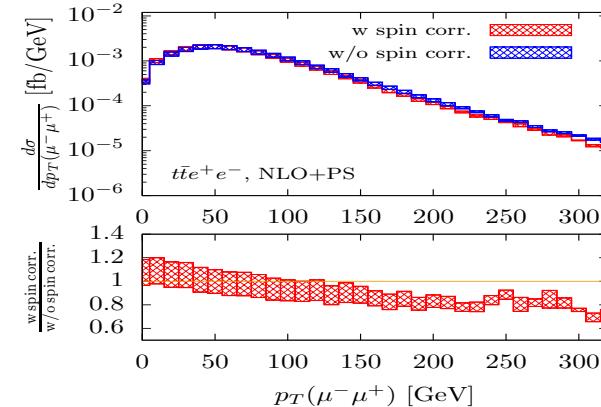


$\delta\eta_{\text{SM}} \sim g_{\text{BSM}}^2 \frac{E^2}{M^2}$ Effects in tails of distributions but also anomalous shapes

Top+additional leptons

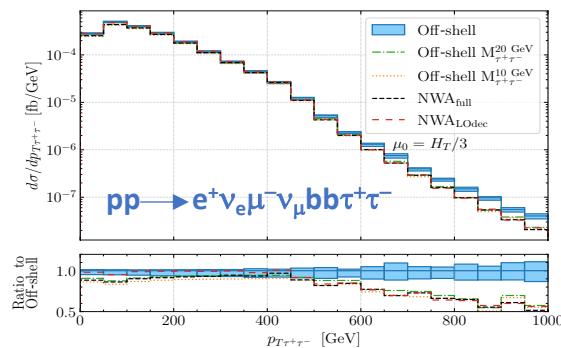


[CMS: arXiv:2012.04120]



M. Ghezzi et al.
[2112.08892]

Off-shell studies

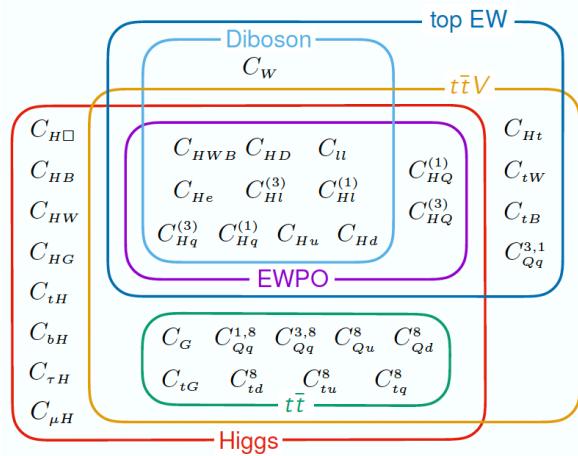


G. Bevilacqua et al. [2203.15688]

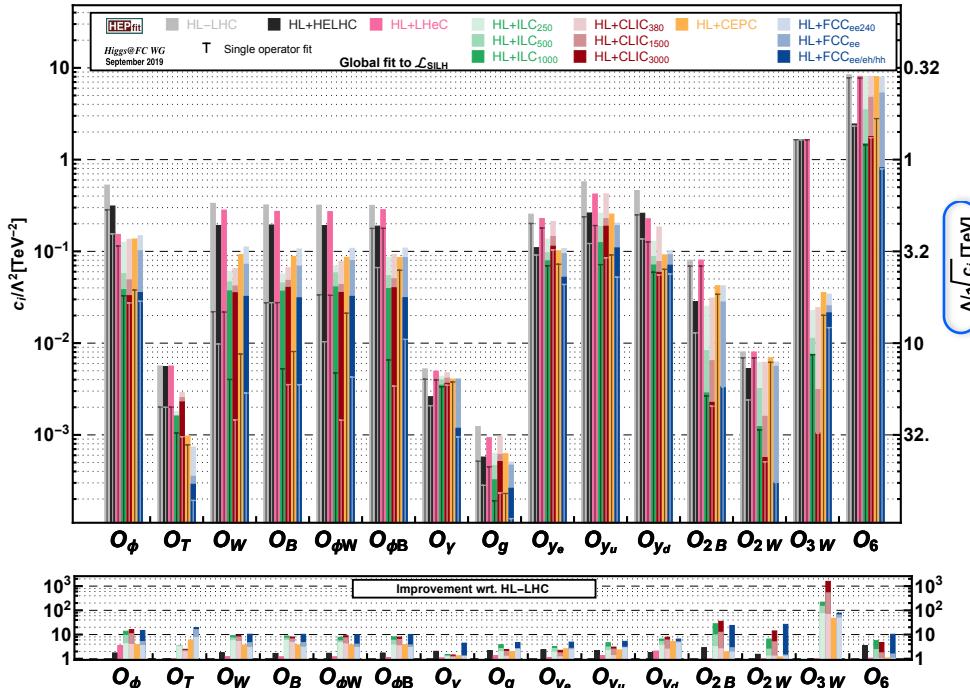
Pointing to the need for precision in modelling signatures from $t\bar{t}+X$ processes in regions where on-shell calculations may not be accurate enough

EFT global fits

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \left(\frac{1}{\Lambda^2} \sum_i C_i O_i + \text{h.c.} \right) + O(\Lambda^{-4})$$



EW + Higgs



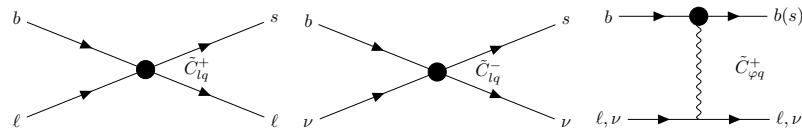
EFT connects different processes with large correlations: pattern of coefficients give insights on underlying BSM model

... adding EW + Higgs + top and flavor!

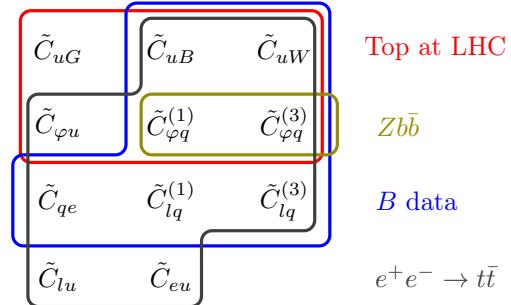
$$\mathcal{L}_{\text{SM}}^{\text{EFT}} \xrightarrow{\Lambda \ll \Lambda_{EW}} \mathcal{L}_{\text{Weak}}^{\text{EFT}} = \sum_{i=1}^{10} C_i^{\text{WEFT}} \mathcal{O}_i^{\text{WEFT}}$$

where

$\mathcal{O}_i^{\text{WEFT}} \rightarrow$ 4-fermion operators of quarks(except t) and leptons
 $C_i^{\text{WEFT}} \rightarrow$ depend on C_i^{SMEFT}



Strong constraint from B-meson semileptonic decays
and intriguing relation with flavor anomalies

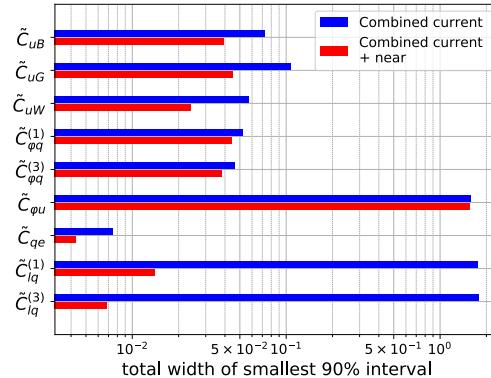


Top at LHC

$Z b\bar{b}$

B data

$e^+ e^- \rightarrow t\bar{t}$



near:
including Belle II
and HL-LHC

Bissman et al. [arXiv:2-12.10456]