

Gravitational Wave Observations

Archisman Ghosh

archisman.ghosh@ugent.be



Lecture 1: 2023 Jun 06

Discovery and Detectors

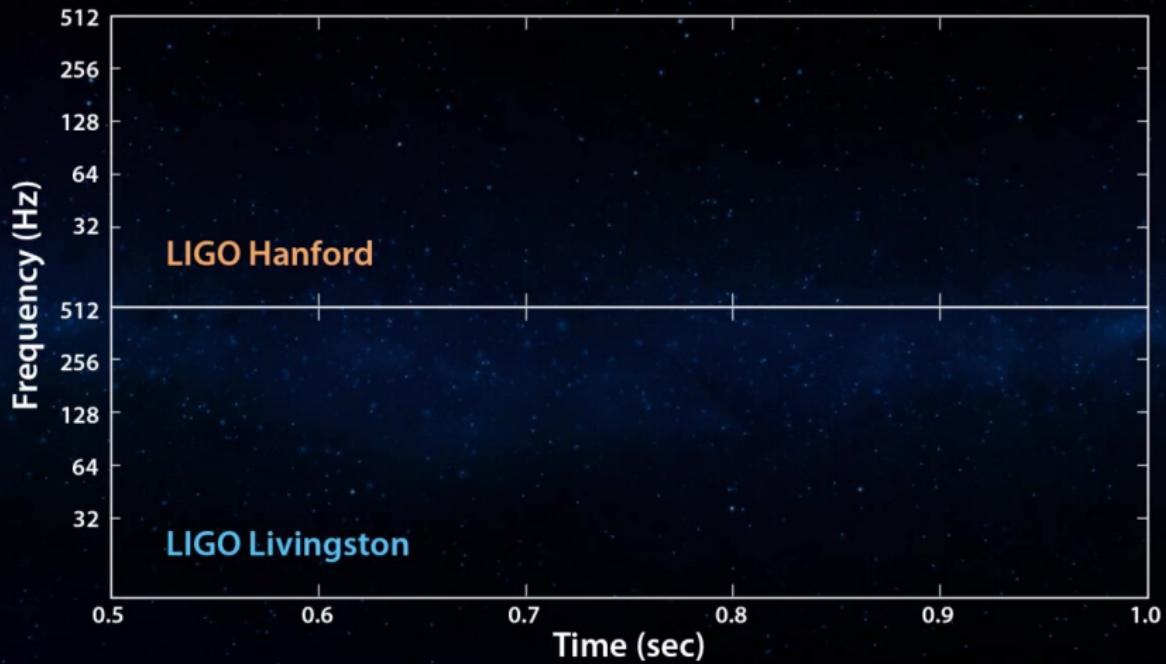


Gravitational wave observations: course plan

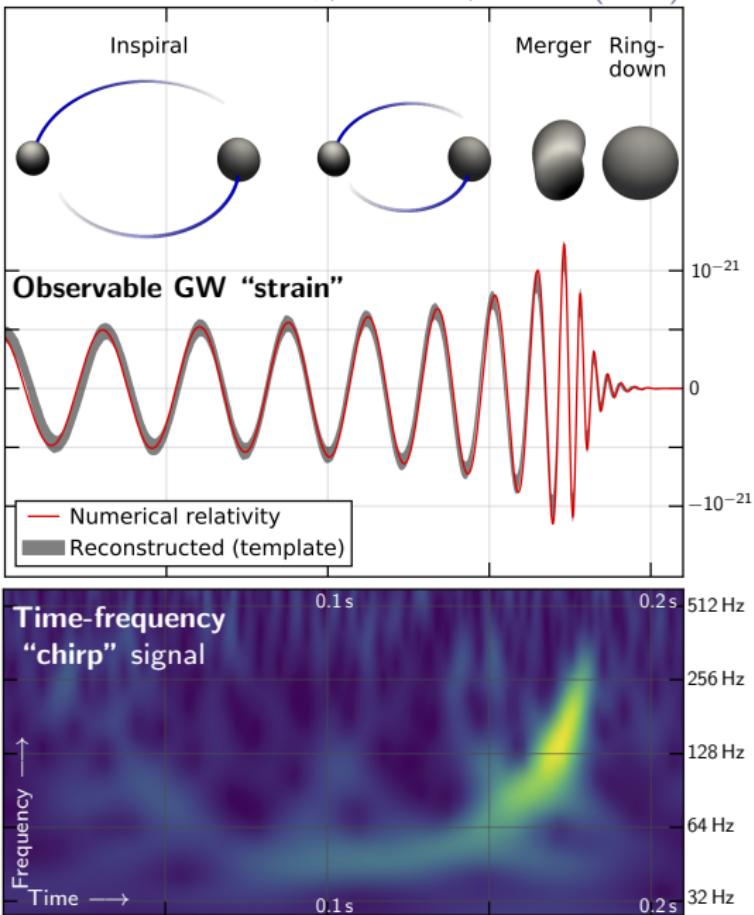
Lecture 1: Discovery and the detectors that made it possible

Lecture 2: Modelling and data analysis

Lecture 3: Observational science and future prospects



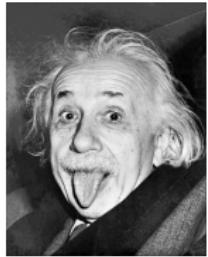
GW150914: BBH



What are gravitational waves?

Ripples in the curvature of spacetime!

General relativity



Einstein equations:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^4}T_{\mu\nu}$$

spacetime matter
 ↑
 coupling

metric $g_{\mu\nu}$

curvature $R_{\mu\nu\lambda\sigma}$ “Riemann tensor”

$$\eta_{\mu\nu} = \begin{pmatrix} -1 & & & \\ & 1 & & \\ & & 1 & \\ & & & 1 \end{pmatrix}$$

$$T^\mu_\nu = \begin{pmatrix} -\rho & & & \\ & \mathcal{P} & & \\ & & \mathcal{P} & \\ & & & \mathcal{P} \end{pmatrix}$$

ρ : energy density

non-linear

\mathcal{P} : pressure

Linearized GR

Small perturbations h about a flat background:

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$

flat background metric, $\eta_{\mu\nu}$

small deformations, $h_{\mu\nu} \ll 1$

“trace reversed” $\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}g_{\mu\nu}h$ (convenient variable)

In Lorenz gauge, where $\partial_\mu \bar{h}^{\mu\alpha} = 0$,

$$\square \bar{h}_{\mu\nu} = 0 \quad \text{wave equation!} \quad \text{for relative deformation or strain } h \equiv \frac{\delta l}{l}$$

GW polarizations

Two physical degrees of freedom

rest are “gauge” d.o.f

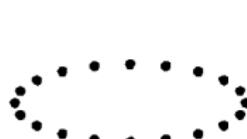
Gauge transformation possible to **TT gauge**

where d.o.f are transverse and traceless

Two polarizations: + and \times (“plus” and “cross”)

$$\begin{pmatrix} 0 & & & \\ h_+ & h_\times & & \\ h_\times & -h_+ & & \\ & & & 0 \end{pmatrix}$$

transverse, traceless in **TT gauge**



+ polarization



\times polarization

Are gravitational waves physical?

1916–1918: **Einstein**'s calculations (incl. flux, quadrupole formula, . . .)

1922: **Eddington** brought up the importance of gauge artefacts

1936: **Einstein** and **Rosen** claimed that GWs could not exist!

Robertson (reviewer) was convinced otherwise → **Infeld** → **Einstein**

1956: **Pirani** rephrased in terms of co-ordinate independent observables

1957: **Feynman** demonstrated GWs could transmit energy (Chapel Hill)

Bondi | **Weber** (resonant bar detectors)

1975: **Hulse-Taylor** binary **pulsar** ⇒ GWs exist! (Nobel Prize 1993)

2015: Direct detection by **LIGO-Virgo** (Nobel Prize 2017)

Pulsars

Spinning neutron stars

strong magnetic fields $\mathcal{O}(10^8\text{-}10^{15}) \text{ G}$

radio emission along magnetic axis

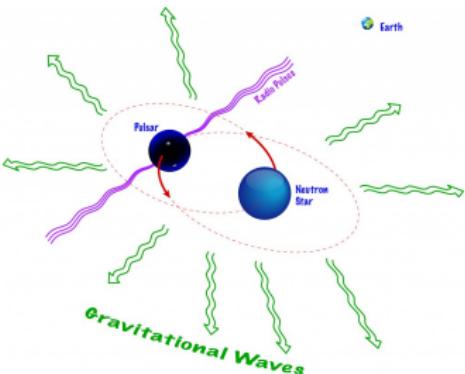
not aligned with spin axis

\Rightarrow pulses of radio emission

most precise clocks!



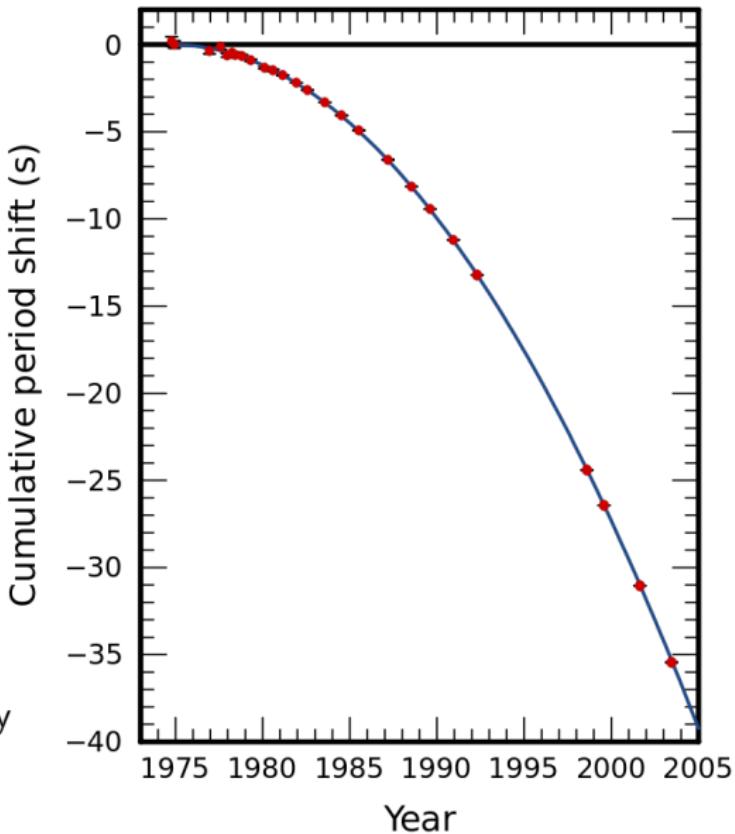
Jocelyn Bell 1967



$$\dot{P} = -\frac{192\pi}{5c^5} \frac{\mu}{M} \left(\frac{2\pi G M}{P} \right)^{5/3} \frac{1 + \frac{73}{24}e^2 + \frac{37}{96}e^4}{(1 - e^2)^{7/2}}$$

PSR B1913+16 orbital decay

Hulse-Taylor (1975)



How can we detect GWs?

rigid rulers	not practicable
resonant mass	bar detectors lunar GW antenna!
proper length between “freely falling” masses	LIGO ET LISA
modulation of time dilation	pulsar timing array
CMB polarization	Planck BICEP

GW strain

Radiation zone:

$$R \ll \lambda \ll d_L$$

$$h_{ij}^{\text{TT}} \simeq \frac{2G}{c^4 d_L} \ddot{\tilde{l}}_{ij}(t - \frac{r}{c})$$

↑ retarded time
changing quadrupole moment

Order of magnitude estimates —

$$h \sim \frac{G}{c^4} \frac{\ddot{I}}{d_L} \quad I \sim M R^2 \quad \ddot{I} \sim \omega^2 I$$

$$M \sim 1 \text{kg}, \quad R \sim 1 \text{m}, \quad \omega \sim 1 \text{s}^{-1} \quad d_L \gg c/\omega \quad \Rightarrow \quad h \ll \frac{G}{c^5} M R^2 \omega^3$$

$$G = 6.64 \times 10^{-11} \text{ Nm}^2 \text{kg}^{-2}$$

$$c = 3 \times 10^8 \text{ms}^{-1} \quad h \sim \frac{G}{c^4} \frac{M R^2 \omega^2}{d_L} = \frac{10^{-44}}{d_L/\text{m}} \quad \text{or} \quad h \ll 10^{-53}$$

spacetime is **stiff**

GW energy

$$L_{\text{GW}} = -\frac{dE}{dt} = \frac{1}{5} \frac{G}{c^5} \langle \ddot{\bar{T}}_{ij} \ddot{\bar{T}}^{ij} \rangle$$

gravitational **Larmor formula**

Order of magnitude estimates —

$$L_{\text{GW}} \sim \frac{G}{c^5} \dot{I}^2 \quad I \sim M R^2 \quad \ddot{I} \sim \omega^3 I$$

$$M \sim 1 \text{kg}, \quad R \sim 1 \text{m}, \quad \omega \sim 1 \text{s}^{-1} \quad d_L \gg c/\omega$$

$$G = 6.64 \times 10^{-11} \text{ Nm}^2 \text{kg}^{-2}$$

$$c = 3 \times 10^8 \text{ ms}^{-1}$$

$$L_{\text{GW}} \sim \frac{(1W)^2}{c^5/G} \sim 10^{-53} \text{ W}$$

GW energy

$$L_{\text{GW}} = -\frac{dE}{dt} = \frac{1}{5} \frac{G}{c^5} \langle \ddot{\bar{I}}_{ij} \ddot{\bar{I}}^{ij} \rangle$$

gravitational **Larmor formula**

Order of magnitude estimates —

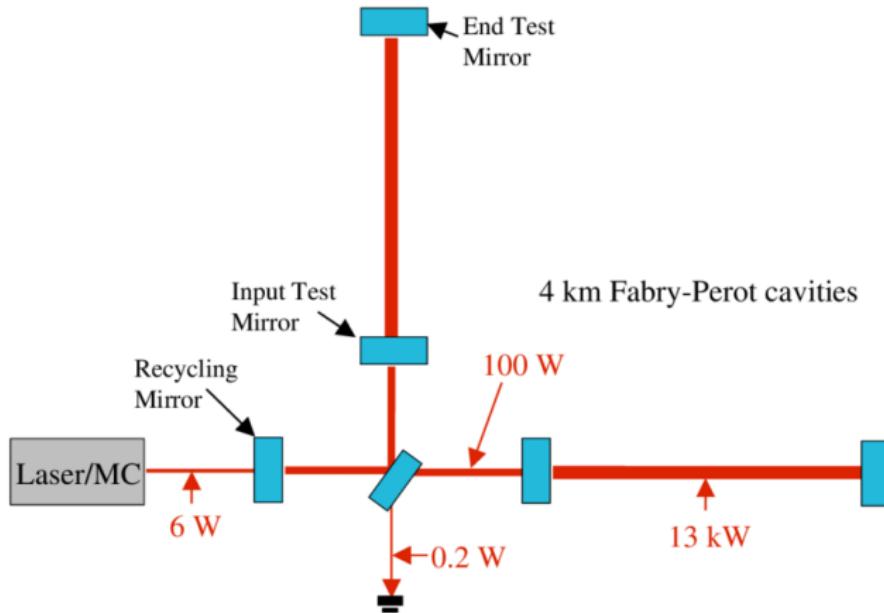
$$L_{\text{GW}} \sim \frac{G}{c^5} \dot{I}^2 \quad I \sim M R^2 \quad \ddot{I} \sim \omega^3 I \sim I \frac{v^3}{R^3}$$

$$\text{virial} \Rightarrow \frac{GM}{R} \approx v^2$$

$$L_{\text{GW}} \sim \frac{G}{c^5} \left(\frac{Mv^3}{R} \right)^2 = \frac{c^5}{G} \left(\frac{GM}{Rc^2} \right)^2 \left(\frac{v}{c} \right)^6 \approx \frac{c^5}{G} \left(\frac{v}{c} \right)^{10}$$

$$\text{“Planck luminosity” } \frac{c^5}{G} = 3.63 \times 10^{52} \text{ W}$$

Interferometric GW detectors



Can we measure a strain of 10^{-21} ?

Over 4 km detectors: *distances* of 10^{-18} m!

nucleus of an atom: 10^{-15} m
molecules (mirror surface): 10^{-10} m

Can we measure a strain of 10^{-21} ?

Michelson interferometer:

$$h = \frac{\Delta I}{I} \approx \frac{\lambda_{\text{laser}}}{L_{\text{ifo}}} \approx \frac{10^{-6} \text{ m}}{10^3 \text{ m}} \approx 10^{-9}$$

Fabry-Pérot cavity (wave-optics): light bounces back and forth

$$L \rightarrow L_{\text{eff}} = \sim 140 \times 4 \text{ km} \approx 600 \text{ km}$$

Note: $\lambda_{\text{GW}} = \frac{c}{f_{\text{GW}}} \approx \frac{3 \times 10^8 \text{ ms}^{-1}}{300 \text{ s}^{-1}} \approx 1000 \text{ km}$

$$h = \frac{\Delta I}{I} \approx \frac{\lambda_{\text{laser}}}{L_{\text{eff}}} \approx \frac{10^{-6} \text{ m}}{10^6 \text{ m}} \approx 10^{-12}$$

Can we measure a strain of 10^{-21} ?

We can measure *a fraction* of a wavelength not just dark and bright spots

Photodetector (counts photons):

shot noise \Leftrightarrow Poisson statistics

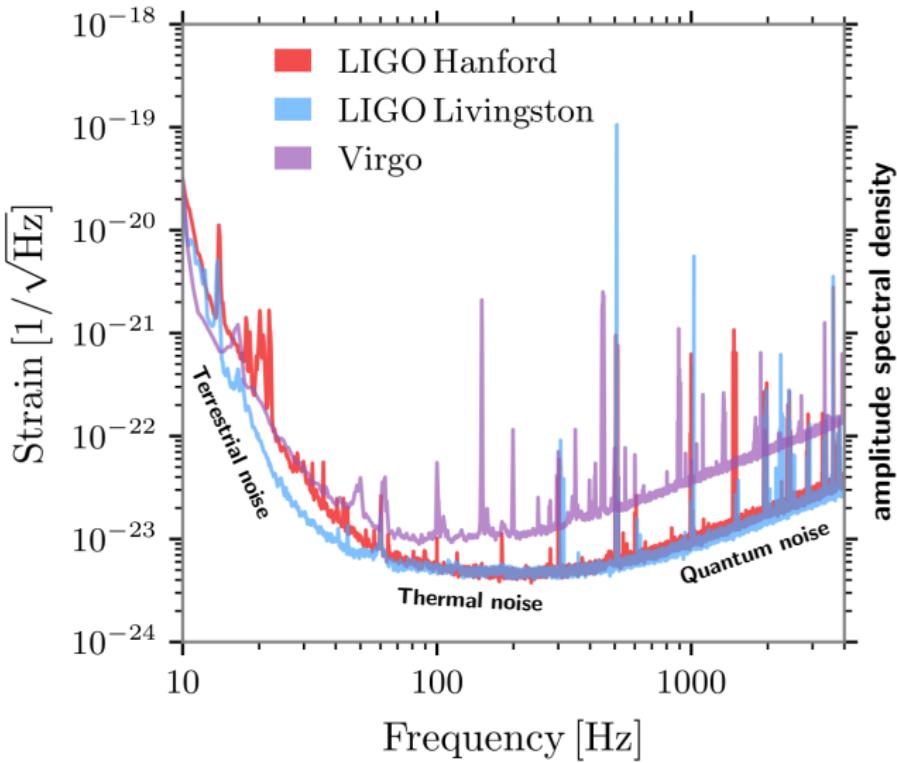
$$\frac{\delta I}{\lambda_{\text{laser}}} \approx \frac{\Delta N_{\text{photon}}}{N_{\text{photon}}} \approx \frac{N_{\text{photon}}^{1/2}}{N_{\text{photon}}} \approx N_{\text{photon}}^{-1/2}$$

$$N_{\text{photon}} = \frac{P_{\text{laser}} \tau}{hc/\lambda_{\text{laser}}} \qquad \qquad \qquad \tau \lesssim \frac{1}{f_{\text{GW}}} , \qquad 1 \text{ W laser}$$

$$\lesssim \frac{P_{\text{laser}} \lambda_{\text{laser}}}{hc f_{\text{GW}}} = \frac{10^{-6}}{6 \times 10^{-34} \times 3 \times 10^8 \times 300} \approx 10^{16} \text{ photons}$$

$$h = \frac{\Delta I}{I} \approx \frac{\lambda_{\text{laser}}}{L_{\text{eff}}} \times N_{\text{photon}}^{-1/2} \approx \frac{10^{-6} \text{ m}}{10^6 \text{ m}} \times 10^{-8} \approx 10^{-20}$$

Detector noise



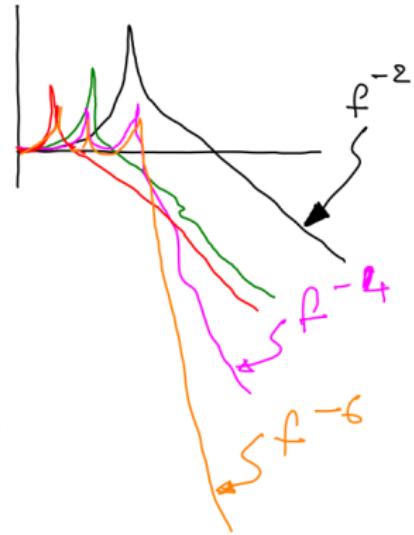
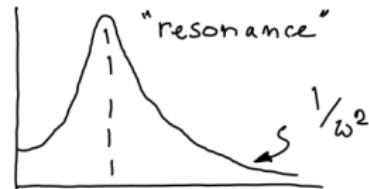
Terrestrial noise: vibration isolation

Seismic noise | anthropogenic noise

$$\ddot{x} + \gamma \dot{x} + \omega_0^2 x = f$$

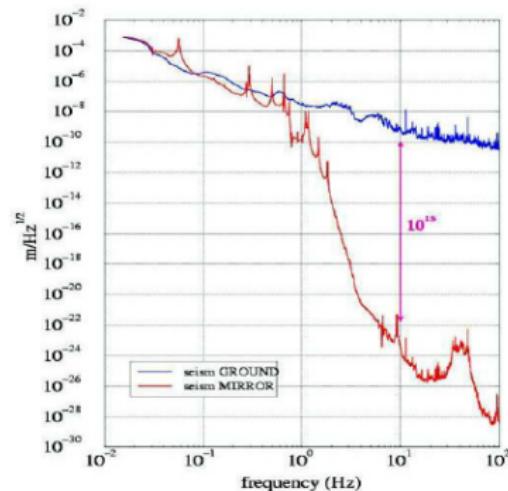
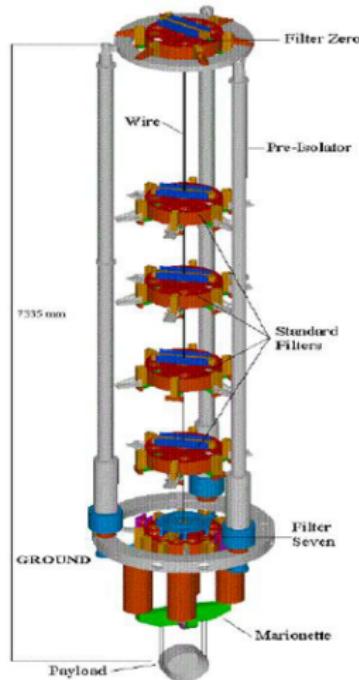
$$x = \frac{f}{\omega_0^2 - \omega^2 + i\gamma\omega}$$

$$x \propto \omega^{-2} \quad (\omega \gg \omega_0)$$



multi-stage pendulum | inverted pendulum
low natural frequency

Multi-stage suspension system



Virgo super-attenuator

Thermal noise

Mirror and its coating get heated up \Rightarrow Brownian motion!

$$S_n(\omega) \approx \frac{4k_B T}{M\omega^2} \operatorname{Re} \left[\frac{i\omega}{\omega_0^2 - \omega^2 + i\omega_0^2 \phi(\omega)} \right]$$
$$\Rightarrow \sqrt{S_n(f)} \sim f^{-5/2}$$

Thermal motion of atoms is the limiting factor! CRYOGENICS

Quantum noise

Photon shot noise | radiation pressure noise

⇒ standard quantum limit

Quantum squeezing to surpass quantum limit?

Frequency-dependent squeezing!

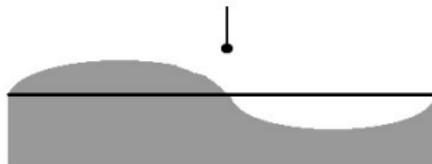
Gravity gradient noise / Newtonian noise

Newtonian gravity fluctuates, e.g. due to surface waves!

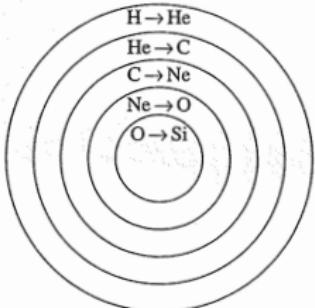
important for future detectors

Cannot be filtered

active subtraction?



Death of main sequence stars



(gravito-thermal instability) negative specific heat \Rightarrow
red giant 'core-halo' structure (hot core, cool outer layers)

progenitor up to $8 M_{\odot}$: planetary nebula + **white dwarf** (e^- degeneracy)

maximum mass of white dwarf $\approx 1.3 M_{\odot}$ (Chandrasekhar limit)
above that \Rightarrow **neutron star** (n degeneracy)

progenitor $8 - 25 M_{\odot}$: **core collapse supernova** (SNe Ib, Ic, II)

most of heavier elements up to ^{56}Fe produced in SNe explosions
first stars: very low **metallicity**

Neutron stars

$$M_{\text{NS}} \gtrsim 1.3 M_{\odot}$$

$$R_{\text{NS}} \lesssim 12 \text{ km}$$

$$\omega \gtrsim \text{kHz}$$

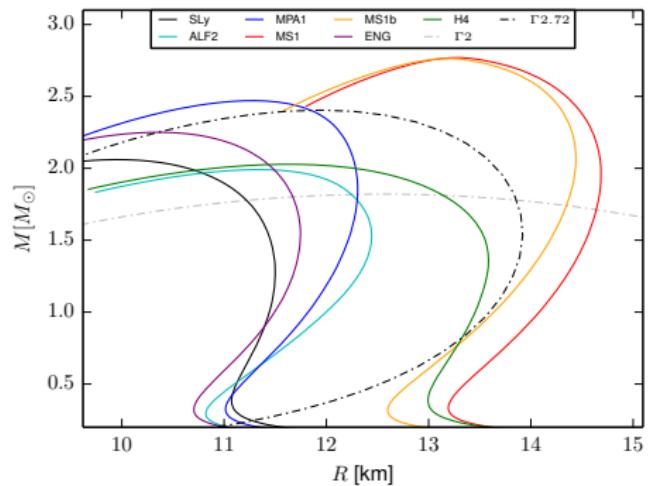
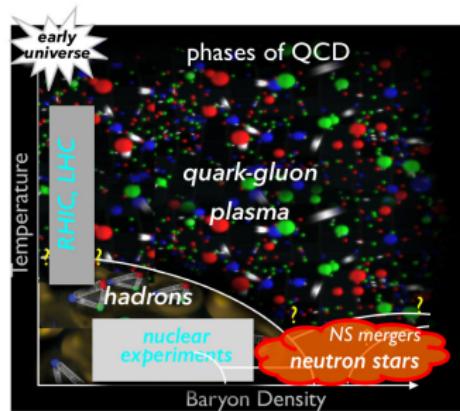
solar mass in a city!

giant nucleus; cold ball of strong interactions

equation-of-state

$$\rho = \mathcal{P}(\rho) \quad \Leftrightarrow \quad M(R)$$

maximum mass?



Pulsars

Spinning neutron stars

strong magnetic fields $\mathcal{O}(10^8\text{-}10^{15}) \text{ G}$

radio emission along magnetic axis

not aligned with spin axis

\Rightarrow pulses of radio emission

most precise clocks!



Jocelyn Bell 1967

Black holes

Schwarzschild solution (2016)

$$\text{Schwarzschild radius} = \frac{2GM}{c^2}$$

$$ds^2 = - \left(1 - \frac{2GM}{rc^2}\right) dt^2 + \frac{dr^2}{\left(1 - \frac{2GM}{rc^2}\right)} + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

Astrophysical BH rotating

\Rightarrow **Kerr solution:** mass m ; "spin" \vec{a}

No restriction on mass | astrophysically expect few $M_\odot - 60 M_\odot$

lower mass gap, PISN mass gap | primordial BHs?

fixed shape (in a given field): no horizon deformability

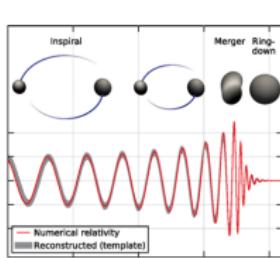
no-hair theorem \Rightarrow quasinormal mode $\omega, \tau = \text{function}(m, a)$

area theorem

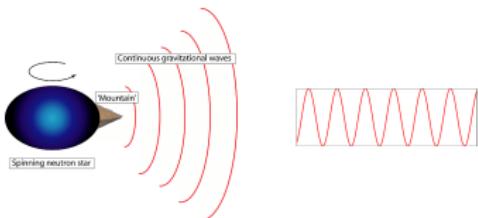
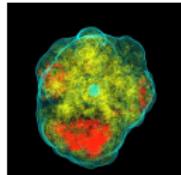
Hawking temperature $T_H \sim \frac{1}{M}$

cosmic censorship \Rightarrow maximum a given m

Gravitational-wave sources

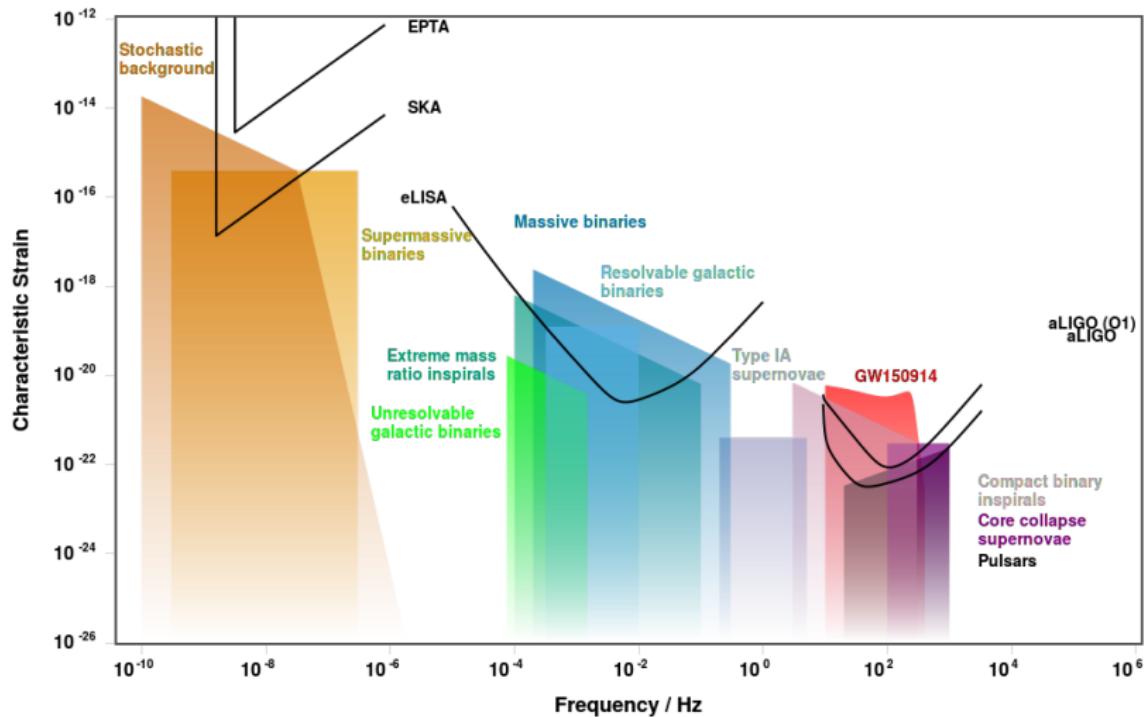


	Modelled	Unmodelled
Transient	Compact binary coalescences NS-NS, NS-BH, BBH	Bursts Supernova explosions
Persistent	Continuous waves Spinning deformed NS	Stochastic background Astrophysical + Cosmological



Other gravitational-wave detectors

Moore, Cole, & Berry, <http://rhcole.com/apps/GWplotter/>



electromagnetic waves

detect **intensity**

$$N_{\text{obs}} \sim \text{sensitivity}^{3/2}$$

incoherent superposition

strongly interacting

affected by gas and dust

deep imaging on small area

high angular resolution

wavelength $\ll / <$ size of source

\Rightarrow **image**

gravitational waves

detect **amplitude**

$$N_{\text{obs}} \sim \text{sensitivity}^3$$

coherent \Rightarrow sensitivity to **phase**

weakly interacting

affected minimally by medium

all sky sensitivity

poor angular resolution

wavelength $\sim / >$ size of source

analogous to **sound**

Geometrized units

$$c = 1, G = 1$$

$$[M] = [L] = [t]$$

Measure in **seconds!**

distances in seconds ✓
masses in seconds?

Schwarzschild radius of the Sun = 3 km

$$\begin{aligned}\Rightarrow \quad 2M_{\odot} &\approx 10^{-5} \text{ s} = 0.01 \text{ ms} \\ \Rightarrow \quad 60M_{\odot} &\approx 0.3 \text{ ms}\end{aligned}$$