

# Implementation of mass-resolution effect in amplitude analysis

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Mengzhen Wang

# Introduction

$$\text{PDF}_{\text{obs}} = \left[ \text{PDF}_{\text{phys}} \otimes G \right] \times \epsilon \quad \text{Efficiency}$$

Data distribution

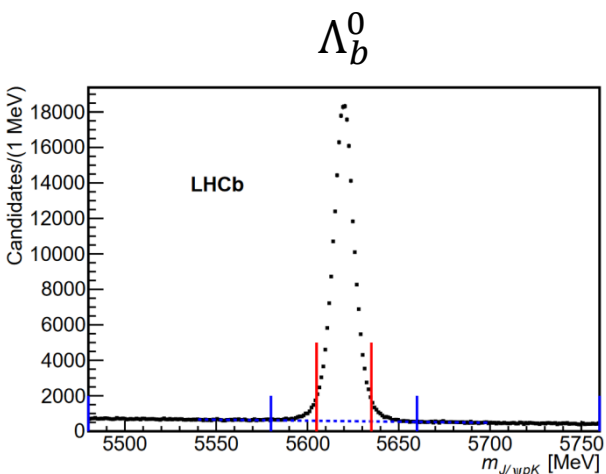
Detector Resolution

Kinematic properties  
of physics progress

- Typical resolution ( $\sigma$ ) @ LHCb  $\sim$  **MeV**

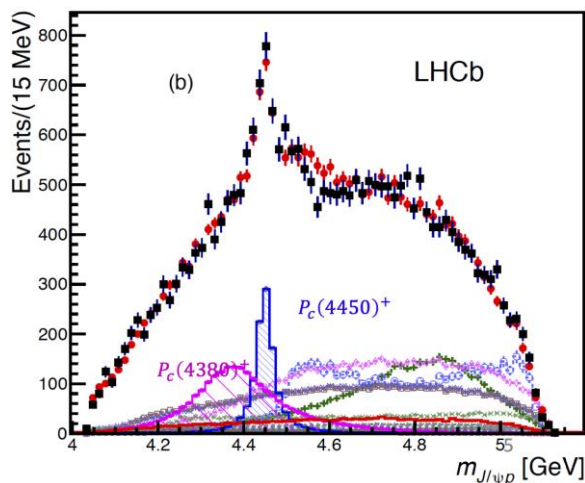
$P_c(4312)^+, P_c(4440)^+, P_c(4457)^+$

$P_c(4380)^+, P_c(4450)^+$



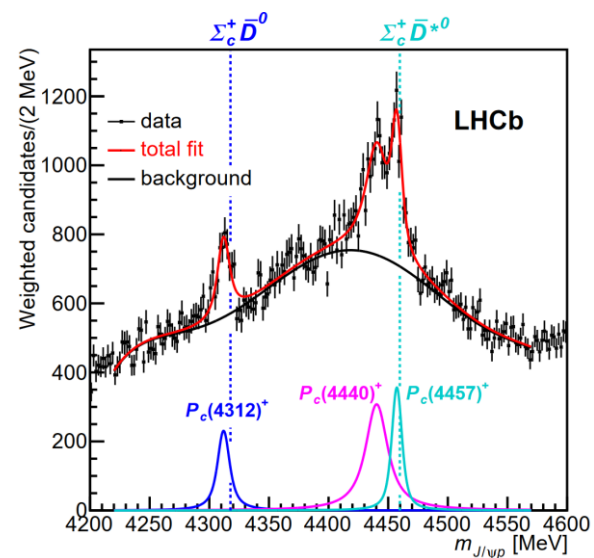
$\Gamma \lll \sigma$

$\text{PDF}_{\text{phys}} \otimes G \approx G$



$\Gamma \sim 100 \text{ MeV} \gg \sigma$

$\text{PDF}_{\text{phys}} \otimes G \approx \text{PDF}_{\text{phys}}$



$\Gamma \sim 10 \text{ MeV} \sim \sigma$

$\text{PDF}_{\text{phys}} \otimes G$

# Introduction

$$\text{PDF}_{\text{obs}} = \underbrace{\text{PDF}_{\text{phys}}}_{\text{Data distribution}} \otimes \underbrace{G}_{\text{Kinematic properties of physics progress}} \times \underbrace{\epsilon}_{\text{Detector Resolution}} \times \text{Efficiency}$$

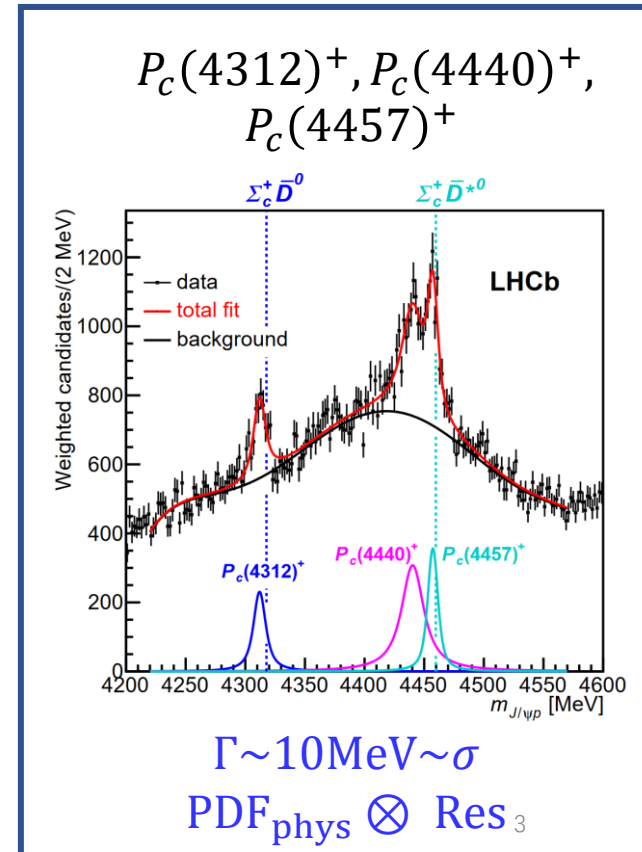
- Typical resolution ( $\sigma$ ) @ LHCb  $\sim$  **MeV**

Today's menu:

Challenges of implementing mass resolution effect in AmAn on huge dataset

Several tricks as potential solutions

Use  $\Lambda_b^0 \rightarrow J/\psi p K^-$  Run1+2 AmAn as example (CERN-THESIS-2021-314, LHCb unofficial)



# Amplitude fit on huge dataset

arXiv:0905.0724

- Unbinned maximum likelihood fit (sFit)

$$\begin{aligned}
 -\ln L(\vec{\omega}) &= -\alpha \sum_i W_i \ln \mathcal{P}_{sig}(m_{Kp,i}, \Omega_i | \vec{\omega}) \\
 &= -\alpha \sum_i W_i \ln \frac{\epsilon(m_{Kp}, \Omega) \times [(|\mathcal{M}(m_{Kp,i}, \Omega_i | \vec{\omega})|^2 \times \Phi(m_{Kp,i})) \otimes G(m_{Kp,i}, \Omega_i)]}{I(\vec{\omega})} \\
 &\approx -\alpha \sum_i W_i \ln \frac{\epsilon(m_{Kp}, \Omega) \times \Phi(m_{Kp,i}) \times [|\mathcal{M}(m_{Kp,i}, \Omega_i | \vec{\omega})|^2 \otimes G(m_{Kp,i}, \Omega_i)]}{I(\vec{\omega})} \\
 &= -\alpha \sum_i W_i \ln \epsilon(m_{Kp,i}, \Omega_i) \Phi(m_{Kp}) - \alpha \sum_i W_i \ln \frac{|\mathcal{M}(m_{Kp,i}, \Omega_i | \vec{\omega})|^2 \otimes G(m_{Kp,i}, \Omega_i)}{I(\vec{\omega})},
 \end{aligned}$$

$W_i$ : single-event weight for BKG subtraction;  $M$ : matrix element;  $G$ : resolution function  
 $\epsilon$ : PHSP-dependent efficiency;  $\Phi$ : PHSP density;  $I$ : normalization factor

$$\begin{aligned}
 I(\vec{\omega}) &= \iint \epsilon(m_{Kp}, \Omega) \times [(|\mathcal{M}(m_{Kp}, \Omega | \vec{\omega})|^2 \times \Phi(m_{Kp})) \otimes G(m_{Kp}, \Omega)] dm_{Kp} d\Omega \\
 &\approx \iint \epsilon(m_{Kp}, \Omega) \times \Phi(m_{Kp}) \times [|\mathcal{M}(m_{Kp}, \Omega | \vec{\omega})|^2 \otimes G(m_{Kp}, \Omega)] dm_{Kp} d\Omega. \\
 &= C \sum_j (|\mathcal{M}(m_{Kp,j}, \Omega_j | \vec{\omega})|^2 \times \Phi(m_{Kp})) \otimes G(m_{Kp,j}, \Omega_j),
 \end{aligned}$$

# Challenges: speed of the fit (1)

- Complicated amplitude formula

$$\begin{aligned}
 -\ln L(\vec{\omega}) &= -\alpha \sum_i W_i \ln \mathcal{P}_{sig}(m_{Kp,i}, \Omega_i | \vec{\omega}) \\
 &= -\alpha \sum_i W_i \ln \frac{\epsilon(m_{Kp}, \Omega) \times [(|\mathcal{M}(m_{Kp,i}, \Omega_i | \vec{\omega})|^2 \times \Phi(m_{Kp,i})) \otimes G(m_{Kp,i}, \Omega_i)]}{I(\vec{\omega})} \\
 &\approx -\alpha \sum_i W_i \ln \frac{\epsilon(m_{Kp}, \Omega) \times \Phi(m_{Kp,i}) \times [|\mathcal{M}(m_{Kp,i}, \Omega_i | \vec{\omega})|^2 \otimes G(m_{Kp,i}, \Omega_i)]}{I(\vec{\omega})} \\
 &= -\alpha \sum_i W_i \ln \epsilon(m_{Kp,i}, \Omega_i) \Phi(m_{Kp}) - \alpha \sum_i W_i \ln \frac{|\mathcal{M}(m_{Kp,i}, \Omega_i | \vec{\omega})|^2 \otimes G(m_{Kp,i}, \Omega_i)}{I(\vec{\omega})},
 \end{aligned}$$

$$\begin{aligned}
 |\mathcal{M}|^2 &= \sum_{\lambda_{\Lambda_b^0} = \pm \frac{1}{2}, \lambda_p = \pm \frac{1}{2}, \Delta\lambda_\mu = \pm 1} |\mathcal{M}_{\lambda_{\Lambda_b^0}, \lambda_p, \Delta\lambda_\mu}^{\Lambda^*} + e^{i\Delta\lambda_\mu \alpha_\mu} \sum_{\lambda_p^{P_c}} d_{\lambda_p^{P_c}, \lambda_p}^{\frac{1}{2}} (\theta_p) \mathcal{M}_{\lambda_{\Lambda_b^0}, \lambda_p^{P_c}, \Delta\lambda_\mu}^{P_c} \\
 &+ e^{i\Delta\lambda_\mu \alpha_\mu^{Z_{cs}}} \sum_{\lambda_p^{Z_{cs}}} d_{\lambda_p^{Z_{cs}}, \lambda_p}^{\frac{1}{2}} (\theta_p^{Z_{cs}}) \mathcal{M}_{\lambda_{\Lambda_b^0}, \lambda_p^{Z_{cs}}, \Delta\lambda_\mu^{Z_{cs}}}^{Z_{cs}}|^2.
 \end{aligned}$$

A complicated function to describe decay properties @ 6D PHSP

# Challenges: speed of the fit (2)

- Huge number of data & MC events

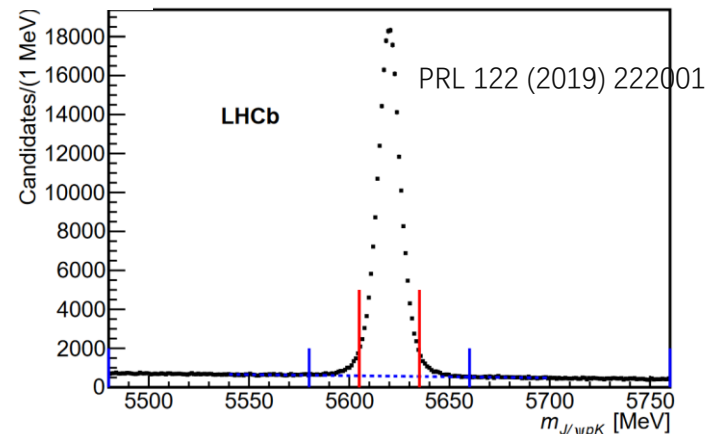
$$-\ln L(\vec{\omega}) = -\alpha \sum_i W_i \ln \epsilon(m_{Kp,i}, \Omega_i) \Phi(m_{Kp}) - \alpha \sum_i W_i \ln \frac{|\mathcal{M}(m_{Kp,i}, \Omega_i | \vec{\omega})|^2 \otimes G(m_{Kp,i}, \Omega_i)}{I(\vec{\omega})},$$

$$I(\vec{\omega}) = C \sum_j (|\mathcal{M}(m_{Kp,j}, \Omega_j | \vec{\omega})|^2 \times \Phi(m_{Kp})) \otimes G(m_{Kp,j}, \Omega_j),$$

Repeat PDF calculation for all data events

> 0.3 Million  $\Lambda_b^0 \rightarrow J/\psi p K^-$  candidates in data

> 2.5 Million simulated  $\Lambda_b^0 \rightarrow J/\psi p K^-$  events



- To calculate  $L$ , the amplitude is to be computed for **~3 Million** times

# Implementation of $\otimes G$

- An ideal solution:

- Find an analytical expression of  $|M|^2 \otimes G$ , but usually not practical
- **Interference term**
- **Complicated line-shape functions** (for example K-Matrix with coupled-channel effects considered)

- A naïve solution:

- Do a standard numerical convolution

$$|M|^2 \otimes G = \int_{-\infty}^{\infty} |\mathcal{M}(m_{J\psi p} = m_a + x)|^2 G(x) dx$$

$$\approx \int_{-\alpha}^{\alpha} |\mathcal{M}(m_{J\psi p} = m_a + x)|^2 G(x) dx$$

$$\approx \sum_{i=-n}^{i=n} |\mathcal{M}(m_{J\psi p} = m_a + i\Delta x)|^2 G(i\Delta x) \Delta x.$$

Fit Speed  $\propto \frac{1}{2n}$

Not good for complicated analysis

Need further tricks to speed up the process

# No convolution for MC events

- MC events just used for normalization factor

$$I(\vec{\omega}) = C \sum_j (|\mathcal{M}(m_{Kp,j}, \Omega_j | \vec{\omega})|^2 \times \Phi(m_{Kp})) \otimes G(m_{Kp,j}, \Omega_j),$$

- One can safely (?) remove the convolution operator

$$\begin{aligned} I(\vec{\omega}) &\approx \iint [\epsilon(m_{Kp}, \Omega) \times |\mathcal{M}(m_{Kp}, \Omega | \vec{\omega})|^2] \otimes G(m_{Kp}, \Omega) dm_{Kp} d\Omega \\ &= \iint [\epsilon(\Phi + \Phi') \times |\mathcal{M}(\Phi + \Phi' | \vec{\omega})|^2] \times G(\Phi') d\Phi d\Phi' \\ &= \iint [\epsilon(\Phi) \times |\mathcal{M}(\Phi | \vec{\omega})|^2] \times G(\Phi') d\Phi d\Phi' \\ &= \int \epsilon(\Phi) \times |\mathcal{M}(\Phi | \vec{\omega})|^2 d\Phi \times \int G(\Phi') d\Phi' \\ &= \int \epsilon(\Phi) \times |\mathcal{M}(\Phi | \vec{\omega})|^2 d\Phi, \end{aligned}$$

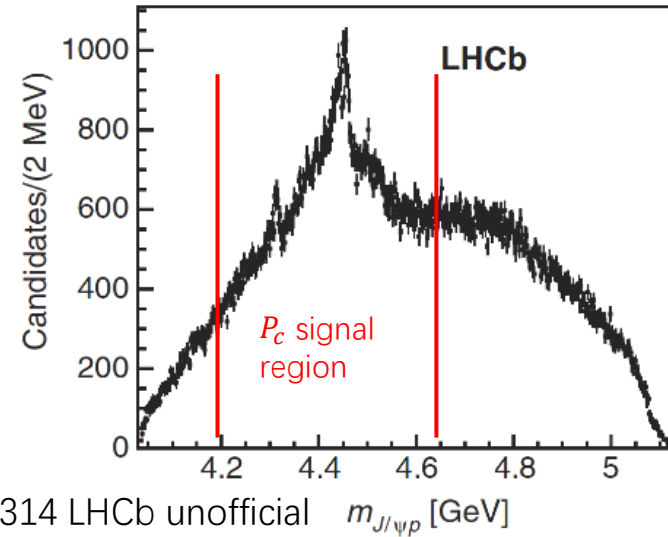
- Effect on the fit speed

- $\frac{1}{2n} \rightarrow \frac{N_{data} + N_{MC}}{2n \times N_{data} + N_{MC}} \gg \frac{1}{2n}$  considering that  $N_{MC} \gg N_{data}$

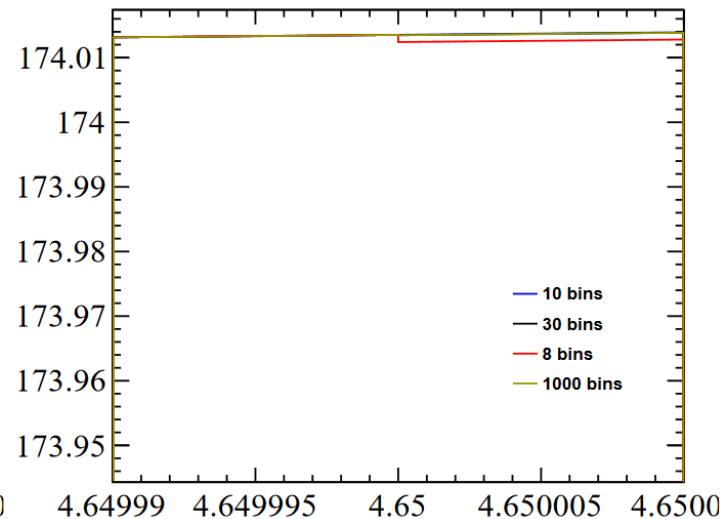
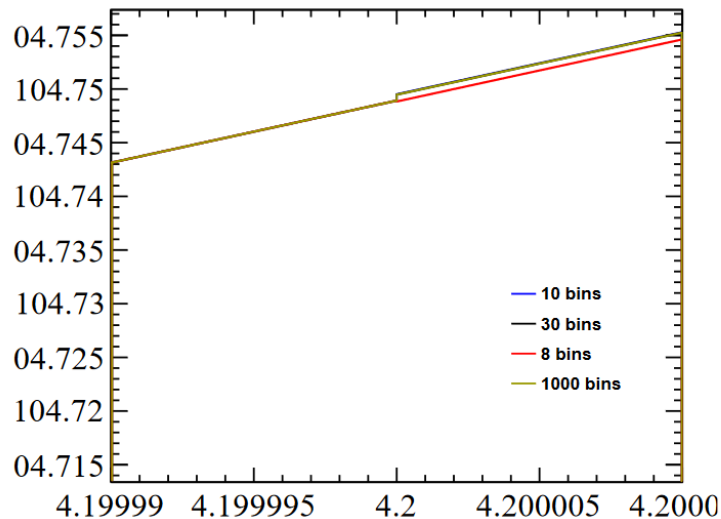


# Only convolution for near-peak events

- Convolution not necessary for events far away from  $P_c$  peaks
  - Mass shape varies slowly
  - $PDF_{phys} \otimes Res \approx PDF_{phys}$  outside of  $P_c$  signal region
- Need to check continuity of PDF



CERN-THESIS-2021-314 LHCb unofficial  $m_{J/\psi p}$  [GeV]



# PDF constructor in fit framework

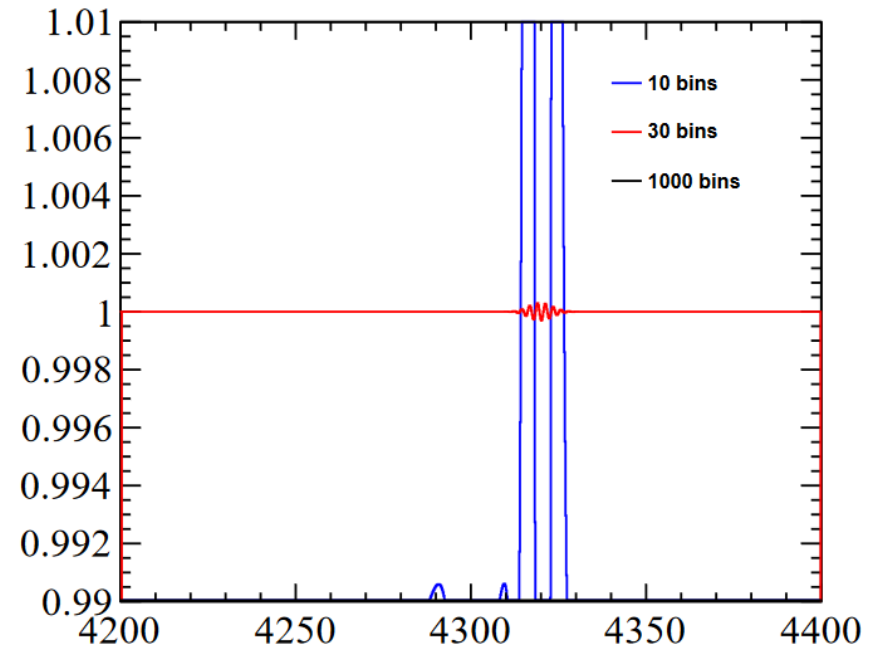
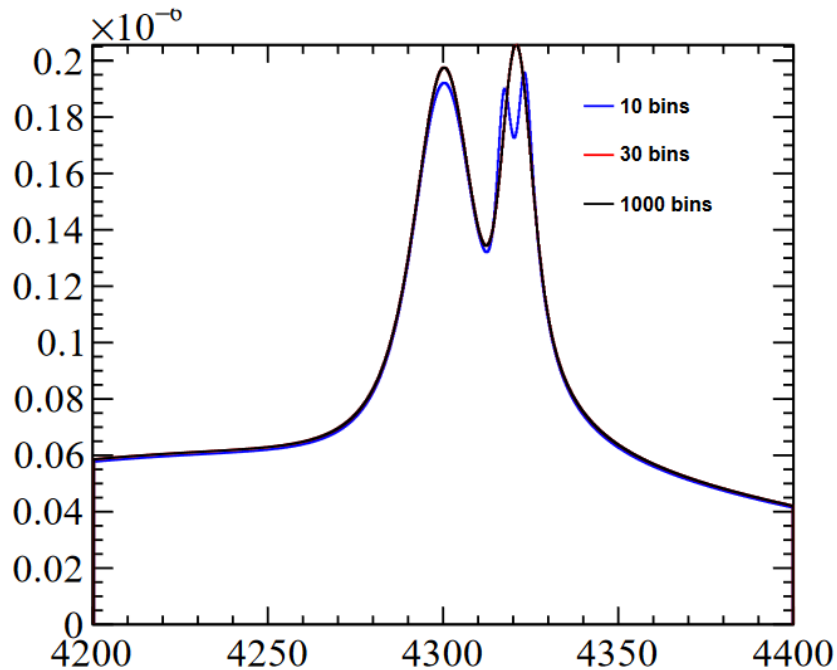
- Rewrite the PDF constructor, only recalculate the **mass-dependent part of amplitude** when doing convolution
  - No need to recalculate  **$\Lambda^*$  chain amplitude** &  **$P_c$  angular terms**

$$\begin{aligned}
 |\mathcal{M}|^2 &= \sum_q |M_{\Lambda^*, Z_{cs}, q}| + \sum_j f_{P_{c,j}, q} R_{P_{c,j}}|^2 \\
 &= \sum_q (|M_{\Lambda^*, Z_{cs}, q}|^2 + M_{\Lambda^*, Z_{cs}, q}^* \times \sum_j f_{P_{c,j}, q} R_{P_{c,j}} + \\
 &\quad M_{\Lambda^*, Z_{cs}, q} \times \sum_j f_{P_{c,j}, q}^* R_{P_{c,j}}^* + \sum_{j,k} f_{P_{c,j}, q} f_{P_{c,k}, q}^* R_{P_{c,j}} R_{P_{c,k}}^*),
 \end{aligned}$$

$$\begin{aligned}
 |\mathcal{M}|^2 \otimes G &= \sum_q (|M_{\Lambda^*, Z_{cs}, q}|^2 \\
 &\quad + M_{\Lambda^*, Z_{cs}, q}^* \times \sum_j f_{P_{c,j}, q} (R_{P_{c,j}} \otimes G) \\
 &\quad + M_{\Lambda^*, Z_{cs}, q} \times \sum_j f_{P_{c,j}, q}^* (R_{P_{c,j}}^* \otimes G) \\
 &\quad + \sum_{j,k} f_{P_{c,j}, q} f_{P_{c,k}, q}^* (R_{P_{c,j}} R_{P_{c,k}}^* \otimes G))
 \end{aligned}$$

# Performance

- Numerical convolution of 30 bins result in a **precision of < 0.1% level**



- Extra time consumption: **10% level**

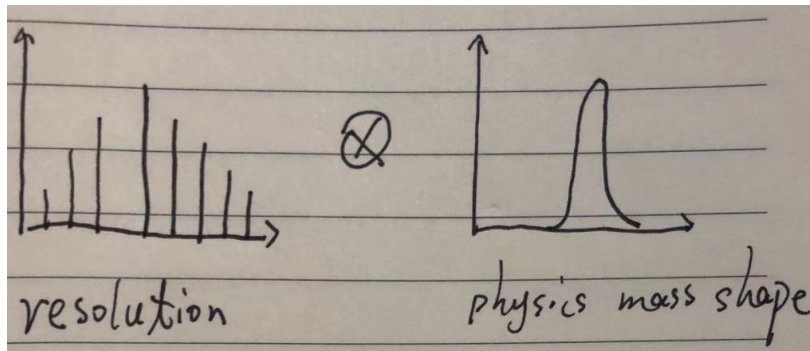
# Other tips

# Set a lower limit on natural widths

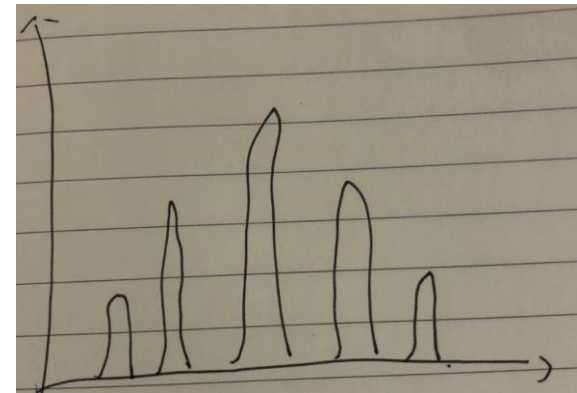
- Numerical convolution only valid when  $\Gamma \gg$  bin width

$$\begin{aligned} |\mathcal{M}|^2 \otimes G &= \int_{-\infty}^{\infty} |\mathcal{M}(m_{J\psi p} = m_a + x)|^2 G(x) dx \\ &\approx \int_{-\alpha}^{\alpha} |\mathcal{M}(m_{J\psi p} = m_a + x)|^2 G(x) dx \\ &\approx \sum_{i=-n}^{i=n} |\mathcal{M}(m_{J\psi p} = m_a + i\Delta x)|^2 G(i\Delta x) \Delta x. \end{aligned}$$

“Sampling” the PDF in each Gaussian bins



$\Gamma <$  bin width



# Set a lower limit on natural widths

- A super narrow natural width may also break the statement:
  - One can safely remove the convolution operator for MC events

$$\begin{aligned} I(\vec{\omega}) &\approx \int \int [\epsilon(m_{Kp}, \Omega) \times |\mathcal{M}(m_{Kp}, \Omega | \vec{\omega})|^2] \otimes G(m_{Kp}, \Omega) dm_{Kp} d\Omega \\ &= \int \int [\epsilon(\Phi + \Phi') \times |\mathcal{M}(\Phi + \Phi' | \vec{\omega})|^2] \times G(\Phi') d\Phi d\Phi' \\ &= \int \int [\epsilon(\Phi) \times |\mathcal{M}(\Phi | \vec{\omega})|^2] \times G(\Phi') d\Phi d\Phi' \\ &= \int \epsilon(\Phi) \times |\mathcal{M}(\Phi | \vec{\omega})|^2 d\Phi \times \int G(\Phi') d\Phi' \\ &= \int \epsilon(\Phi) \times |\mathcal{M}(\Phi | \vec{\omega})|^2 d\Phi, \end{aligned}$$

Correct in math

But in real life it suffers from  
stat. fluctuation of MC sample

- If super narrow peak with reasonable fit fraction
  - Huge (crazy) weights for few events under the narrow peak
  - Hurt the effective stat. power of MC

$$n_{eff} = \frac{(\sum Weights)^2}{(\sum Weight^2)}$$

# Conclusion

- When you find it necessary to implement detector resolution effect in an amplitude analysis, maybe you find it not friendly on the fit speed, which is already very slow
- Several tricks might be helpful:
  - No convolution for MC events
  - No convolution for data events far away from peaking region
  - Try to reorganize the PDF constructor to avoid repeating calculate the same things
- Would be good to set a reasonable lower limit on the natural width of states of interest

Thank you for your attention !

Any questions or comments ?