

# Probing New Physics with Pulsar Timing Arrays

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**Caltech**

based on: [2205.06817](#), [2104.13930](#), and work in progress with D. E. Kaplan, T. Trickle, and NANOGrav collab.



# PULSARS AS COSMIC CLOCKS

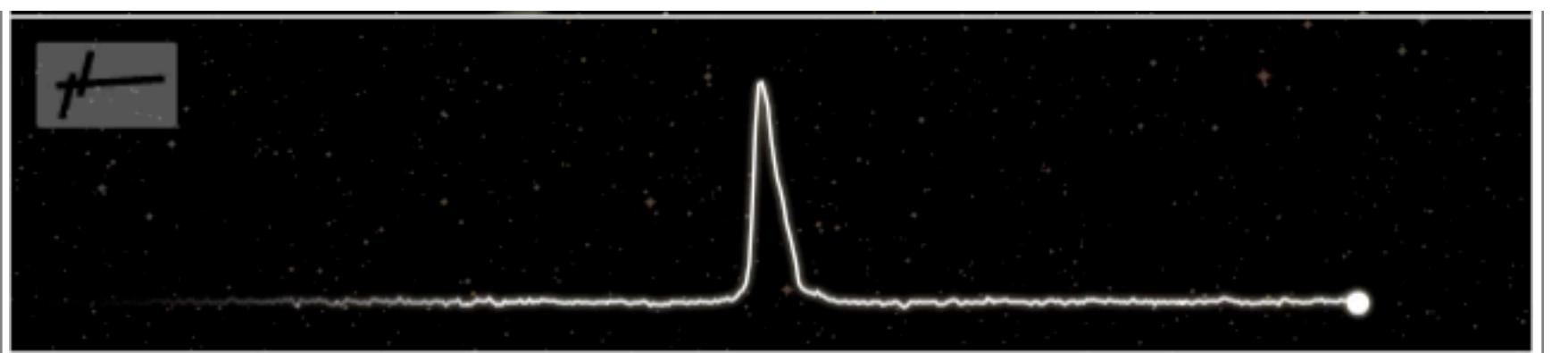
we see a pulse of radiation each time a full rotation is completed

$$\phi \equiv \int \nu(t) dt = \text{integer}$$

for millisecond pulsar the rotation frequency is extremely stable

$$\nu(t) = \nu_0 + \dot{\nu}_0 t \quad \nu_0 / \dot{\nu}_0 \sim 10^{-23} - 10^{-20} \text{ Hz}$$

the pulse arrival times,  $\vec{t}_{\text{TOA},p}$ , can be accurately determined

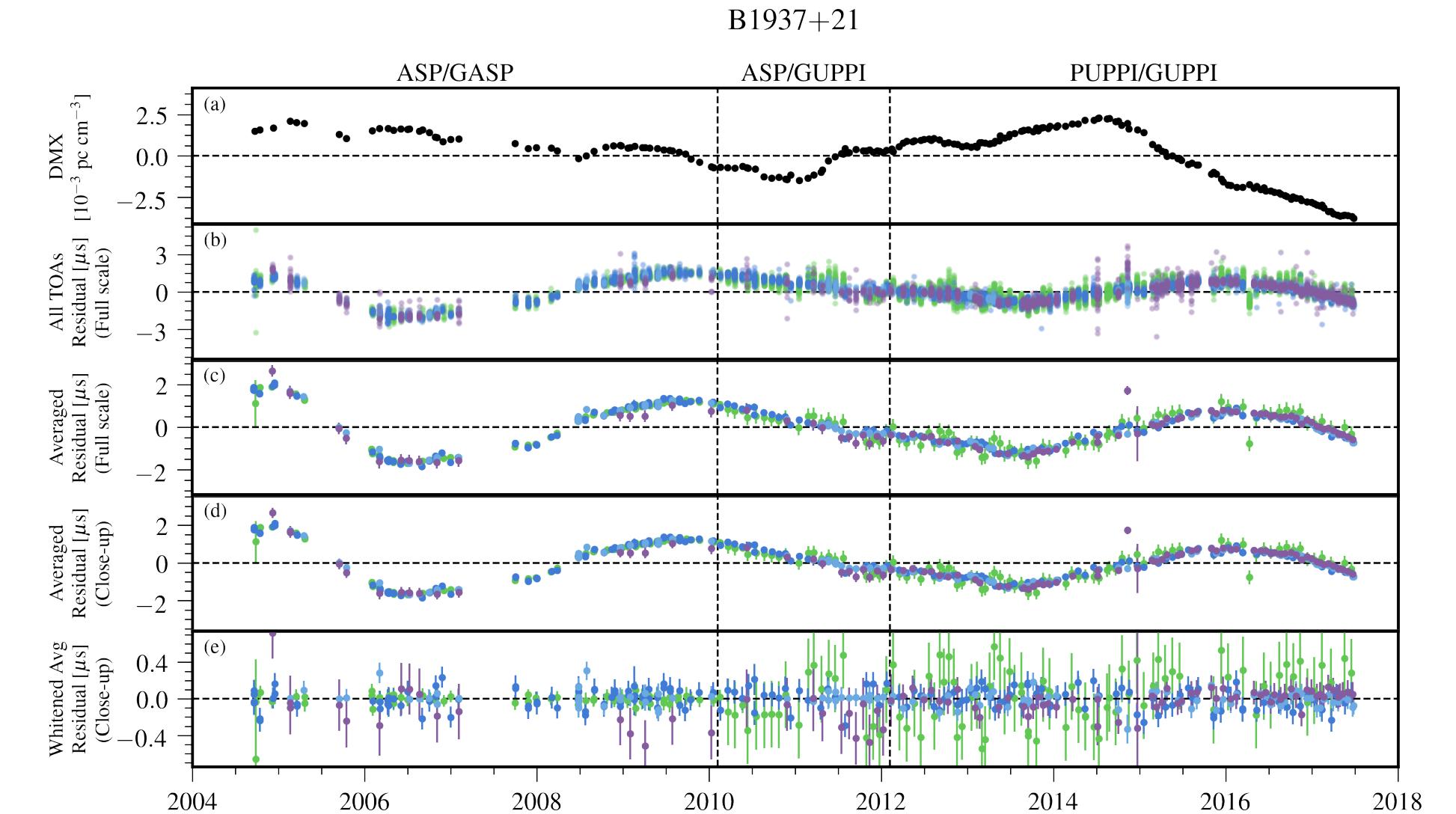


timing residuals

$$\vec{\delta t} \equiv \vec{t}_{\text{TOA}} - \vec{t}_{\text{TOA},p}$$

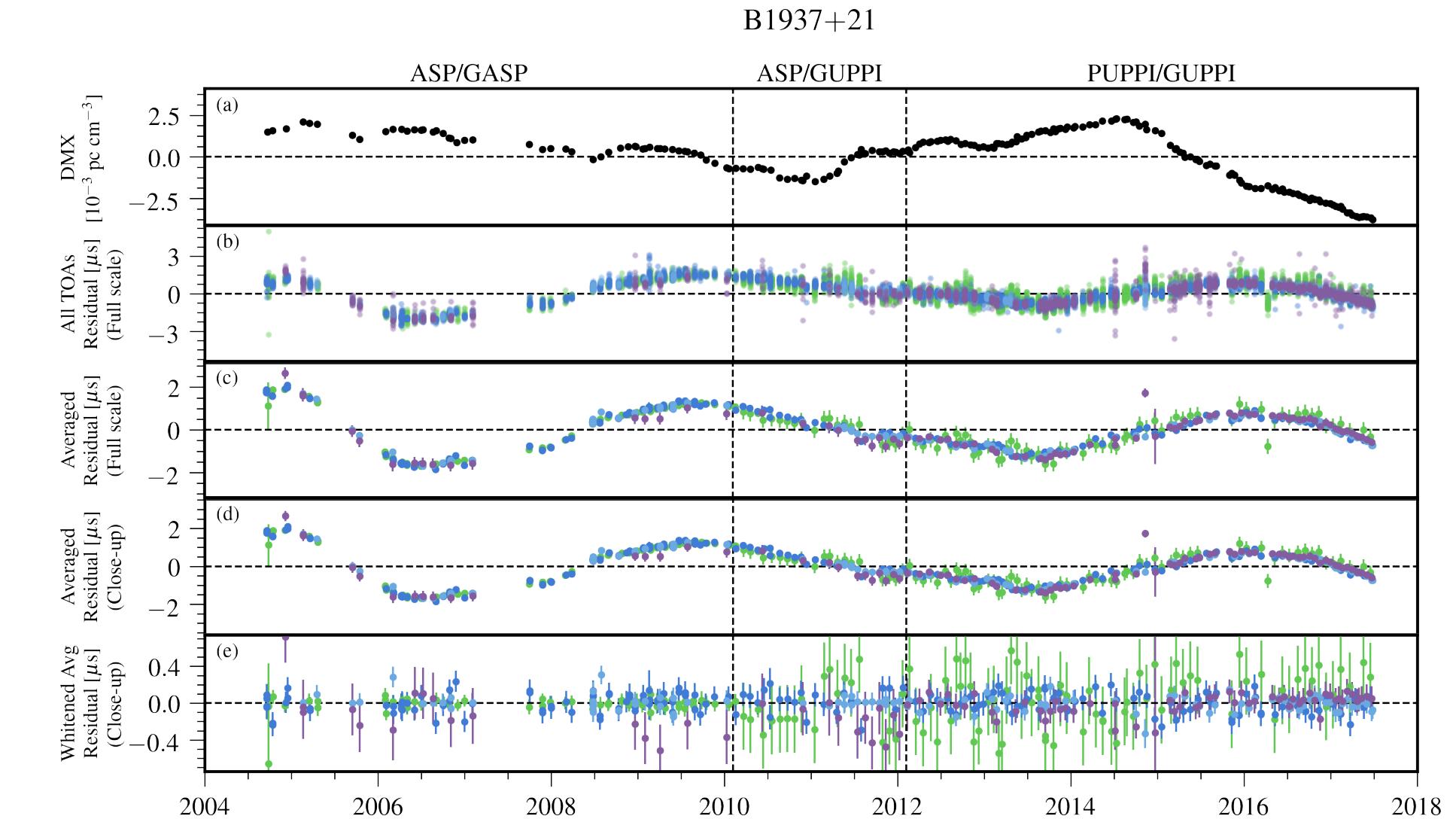
# THE TIMING RESIDUALS

$$\vec{\delta t} = \vec{n} + \mathbf{F}\vec{a} + \mathbf{M}\vec{\epsilon} + \vec{h}(\vec{\theta})$$



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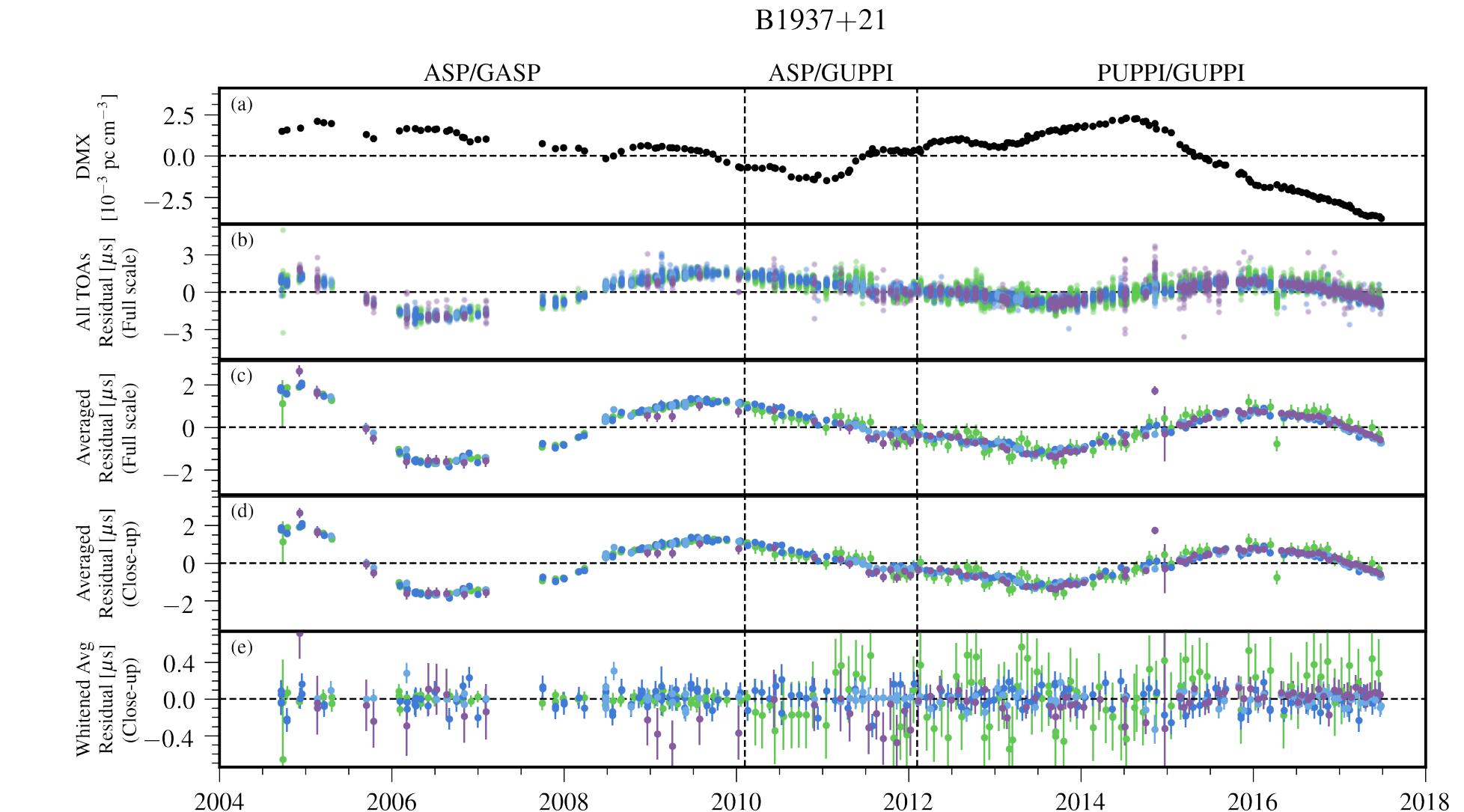
**white noise**

$$\langle n_{i,\mu} n_{j,\nu} \rangle = E_\mu^2 \sigma_i \delta_{ij} \delta_{\mu\nu} + Q_\mu^2 \delta_{ij} \delta_{\mu\nu}$$

EFAC                    EQUAD

# THE TIMING RESIDUALS

$$\vec{\delta t} = \vec{n} + \boxed{\mathbf{F}\vec{a}} + \mathbf{M}\vec{\epsilon} + \vec{h}(\vec{\theta})$$

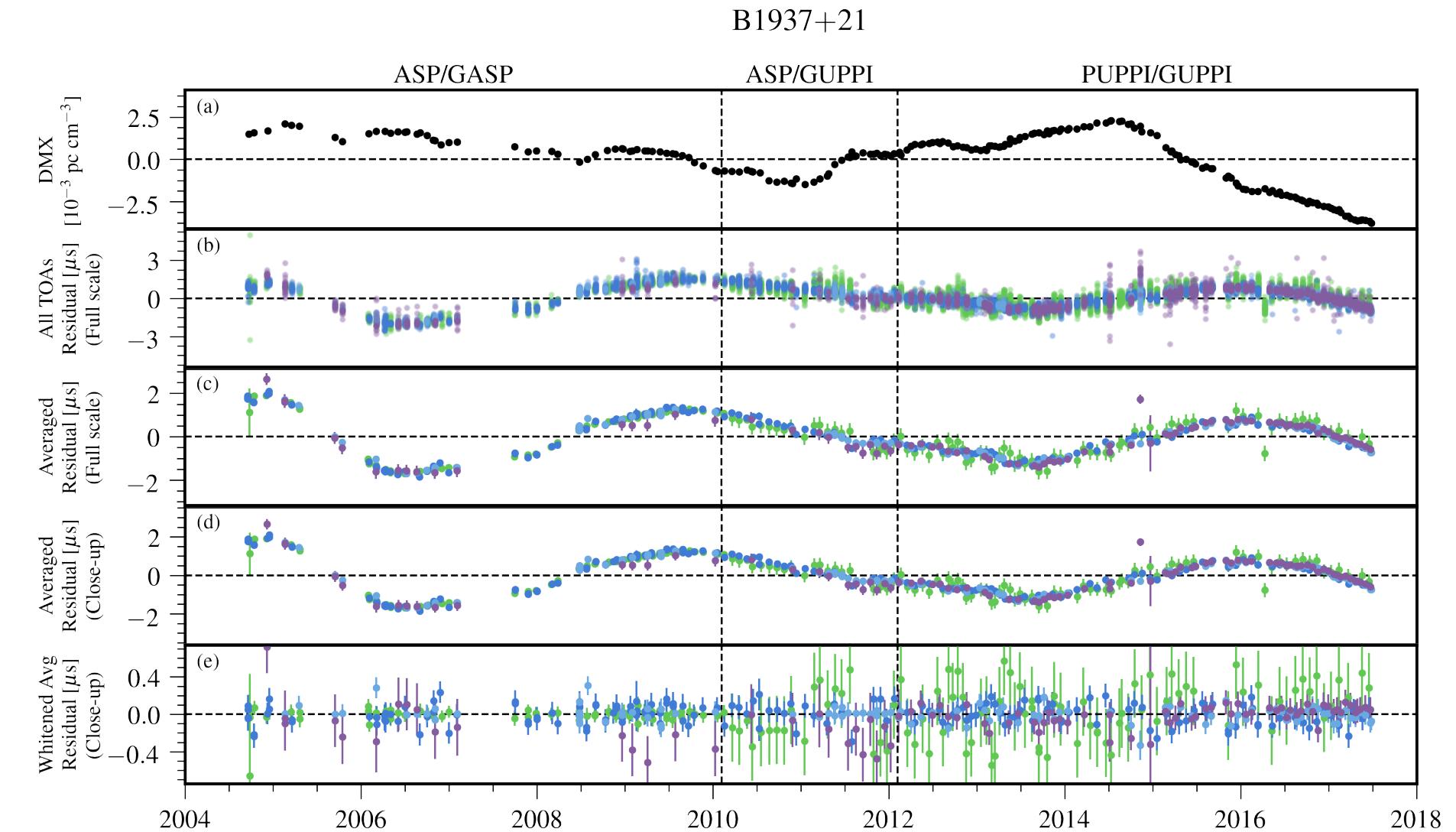


**red noise**

$$\mathbf{F} = \begin{pmatrix} \sin(2\pi t_1/T) & \cos(2\pi t_1/T) & \dots & \sin(2\pi N_f t_1/T) & \cos(2\pi N_f t_1/T) \\ \sin(2\pi t_2/T) & \cos(2\pi t_2/T) & \dots & \sin(2\pi N_f t_2/T) & \cos(2\pi N_f t_2/T) \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ \sin(2\pi t_N/T) & \cos(2\pi t_N/T) & \dots & \sin(2\pi N_f t_N/T) & \cos(2\pi N_f t_N/T) \end{pmatrix}$$

# THE TIMING RESIDUALS

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**red noise**

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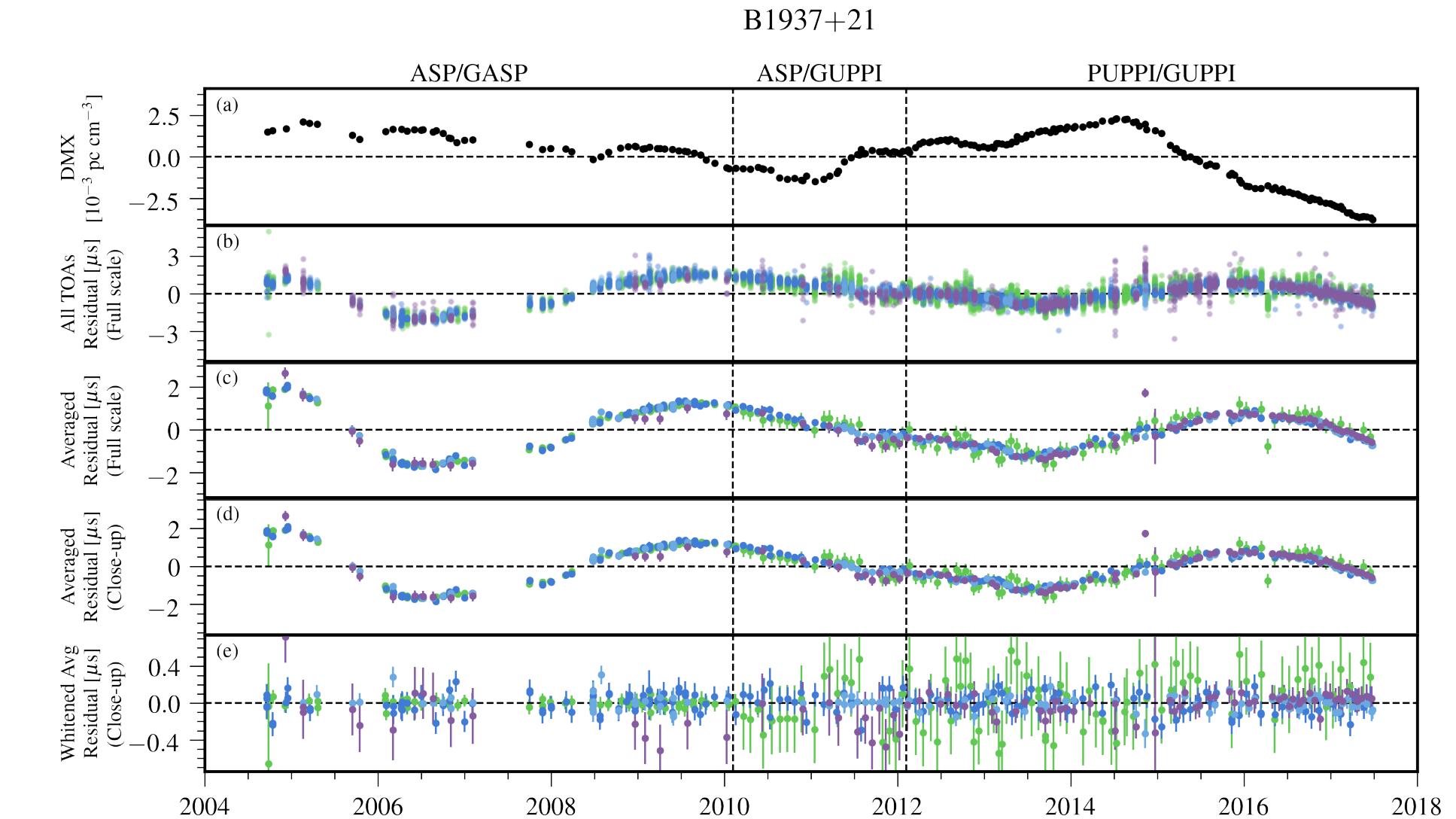
$$p(\vec{a}|\eta) = \frac{\exp\left(-\frac{1}{2}\vec{a}^T \boldsymbol{\phi}^{-1} \vec{a}\right)}{\sqrt{\det(2\pi\boldsymbol{\phi})}}$$

$$[\phi]_{(ak)(bj)} = \boxed{\Gamma_{ab}\rho_k\delta_{kj}} + \boxed{\kappa_{ak}\delta_{kj}\delta_{ab}}$$

common process      intrinsic noise

# THE TIMING RESIDUALS

$$\vec{\delta t} = \vec{n} + \mathbf{F}\vec{a} + \mathbf{M}\vec{\epsilon} + \vec{h}(\vec{\theta})$$



**timing-ephemeris model variations**

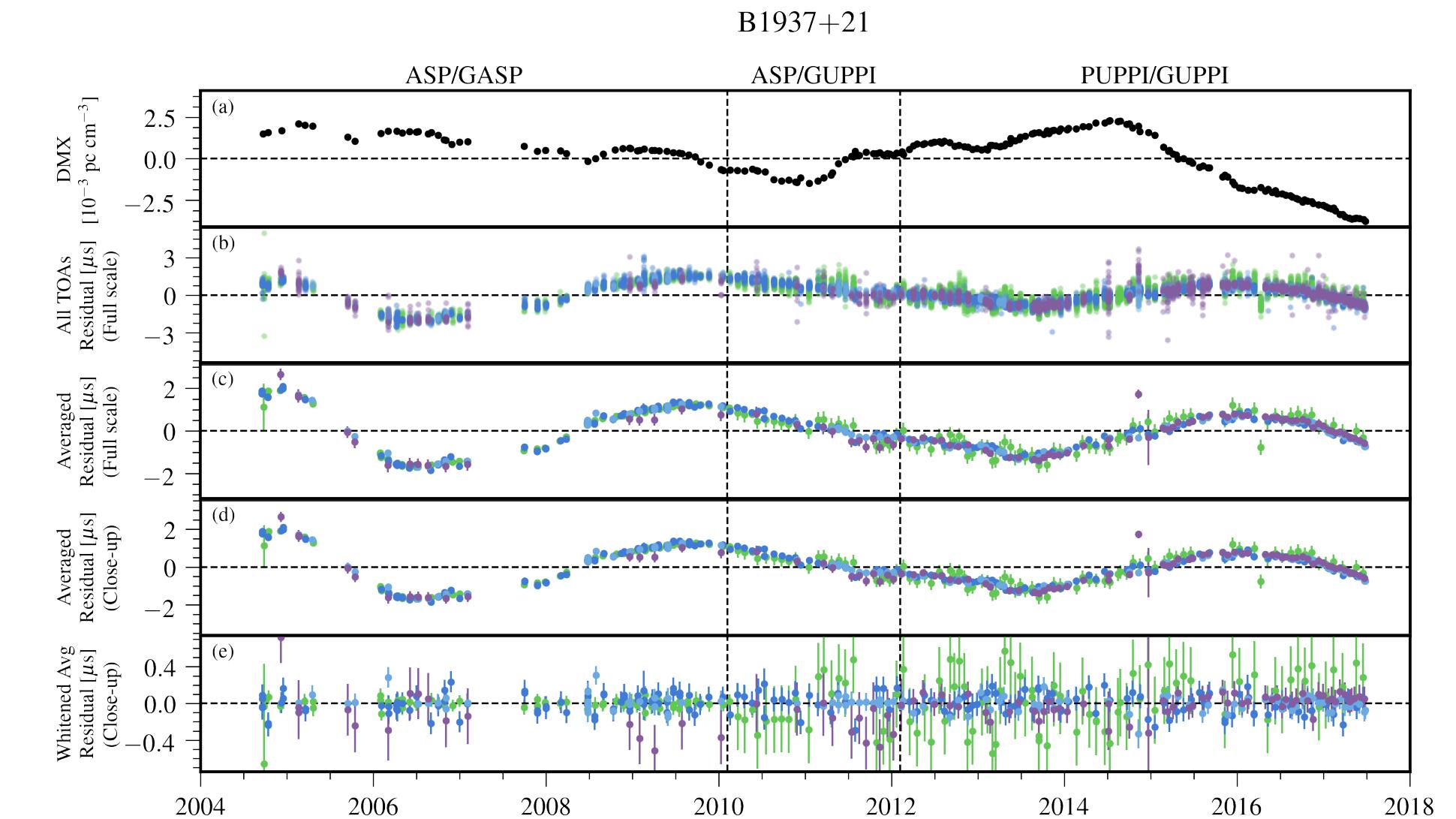
$$M_{ij} = \frac{\partial t_i}{\partial \beta_j}$$

$$\epsilon_j = \beta_j - \beta_{0,j}$$

# THE TIMING RESIDUALS

$$\vec{\delta t} = \vec{n} + \mathbf{F}\vec{a} + \mathbf{M}\vec{\epsilon} + \vec{h}(\vec{\theta})$$

**deterministic signal**



# THE PTA LIKELIHOOD

$$p(\vec{\delta t} | \vec{b}, \vec{\eta}) = \frac{\exp\left(-\frac{1}{2}(\vec{\delta t} - \mathbf{T}\vec{b})^T \mathbf{N}^{-1}(\vec{\delta t} - \mathbf{T}\vec{b})\right)}{\sqrt{\det(2\pi\mathbf{N})}} \times \frac{\exp\left(-\frac{1}{2}\vec{b}^T \mathbf{B}^{-1}\vec{b}\right)}{\sqrt{\det(2\pi\mathbf{B})}} \times p(\vec{\eta})$$

Annotations:

- A green arrow points from  $\vec{\eta} = \{\text{EFAC, EQUAD, } \phi\}$  to the term  $p(\vec{\eta})$ .
- A green arrow points from  $\mathbf{T} = [\mathbf{M}, \mathbf{F}]$  to the term  $(\vec{\delta t} - \mathbf{T}\vec{b})^T \mathbf{N}^{-1}(\vec{\delta t} - \mathbf{T}\vec{b})$ .
- A green arrow points from  $\vec{b} = \begin{bmatrix} \vec{\epsilon} \\ \vec{a} \end{bmatrix}$  to the term  $\vec{b}^T \mathbf{B}^{-1}\vec{b}$ .
- A green arrow points from  $\mathbf{B} = \begin{pmatrix} \infty & 0 \\ 0 & \phi \end{pmatrix}$  to the term  $\vec{b}^T \mathbf{B}^{-1}\vec{b}$ .
- A green arrow points from the label "priors" to the term  $p(\vec{\eta})$ .

# THE PTA LIKELIHOOD

$$p(\vec{\delta t}|\vec{b}, \vec{\eta}) = \frac{\exp\left(-\frac{1}{2}(\vec{\delta t} - \mathbf{T}\vec{b})^T \mathbf{N}^{-1}(\vec{\delta t} - \mathbf{T}\vec{b})\right)}{\sqrt{\det(2\pi\mathbf{N})}} \times \frac{\exp\left(-\frac{1}{2}\vec{b}^T \mathbf{B}^{-1}\vec{b}\right)}{\sqrt{\det(2\pi\mathbf{B})}} \times p(\vec{\eta})$$

$\vec{\eta} = \{\text{EFAC, EQUAD, } \phi\}$  ↗  
 $\mathbf{T} = [\mathbf{M}, \mathbf{F}]$  ↗  
 $\vec{b} = \begin{bmatrix} \vec{\epsilon} \\ \vec{a} \end{bmatrix}$  ↗  
 $\mathbf{B} = \begin{pmatrix} \infty & 0 \\ 0 & \phi \end{pmatrix}$   
 priors ↘

we do not care about the specific noise realization, all the physics is in the parameters controlling the noise statistical properties (i.e.  $\eta$ )

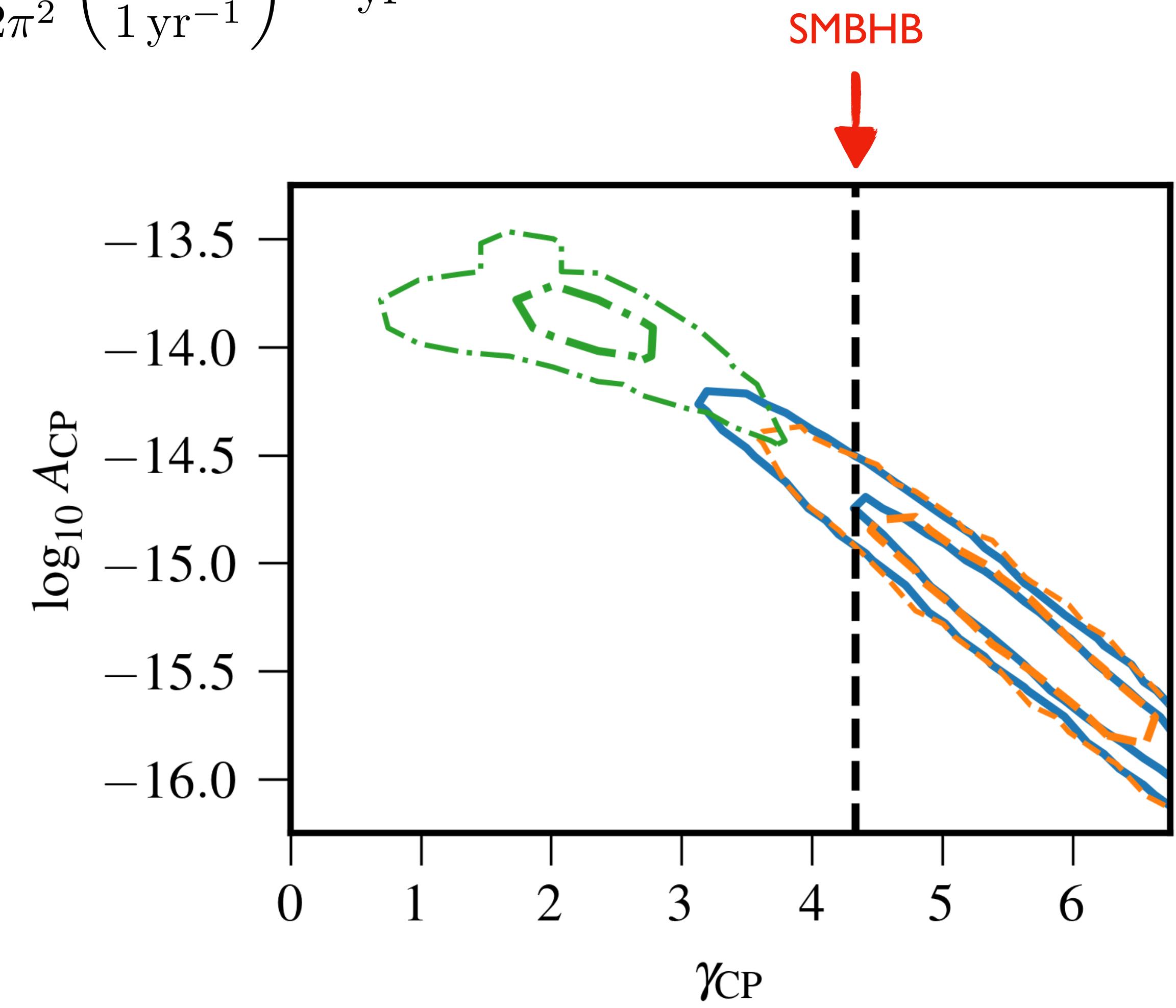
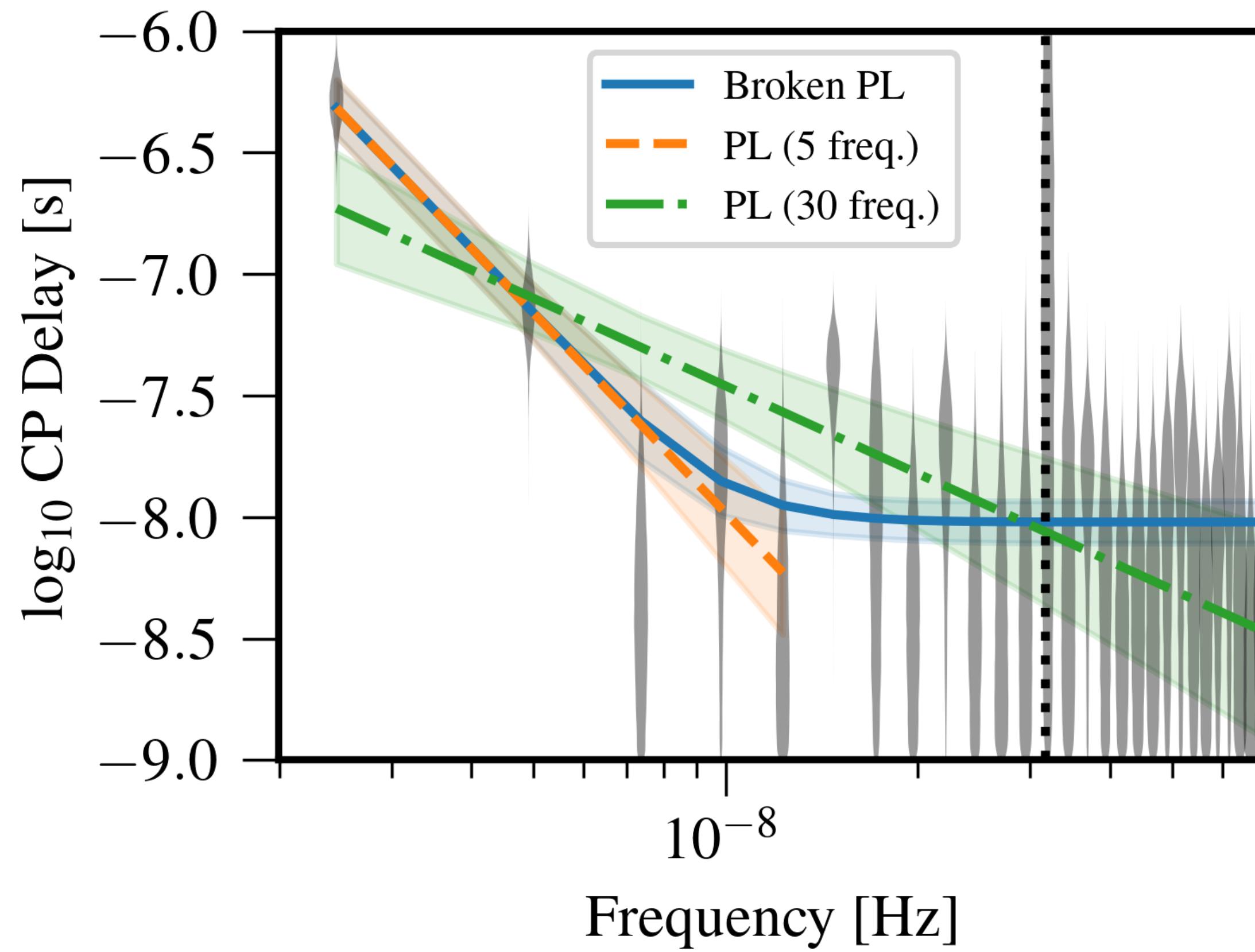
$$p(\{\vec{\delta t}\}|\vec{\eta}) = \int p(\{\vec{\delta t}\}|\{\vec{b}\}, \vec{\eta}) d^{N_p} \vec{b} = \frac{\exp\left(-\frac{1}{2}\vec{\delta t}^T \mathbf{C}^{-1}\vec{\delta t}\right)}{\sqrt{\det(2\pi\mathbf{C})}} \times p(\vec{\eta})$$

$$\mathbf{C} = \mathbf{N} + \mathbf{T}\mathbf{B}\mathbf{T}^T$$

# A FIRST GLIMPSE?

[NANOGrav 12.5 yr \[2009.04496\]](#)

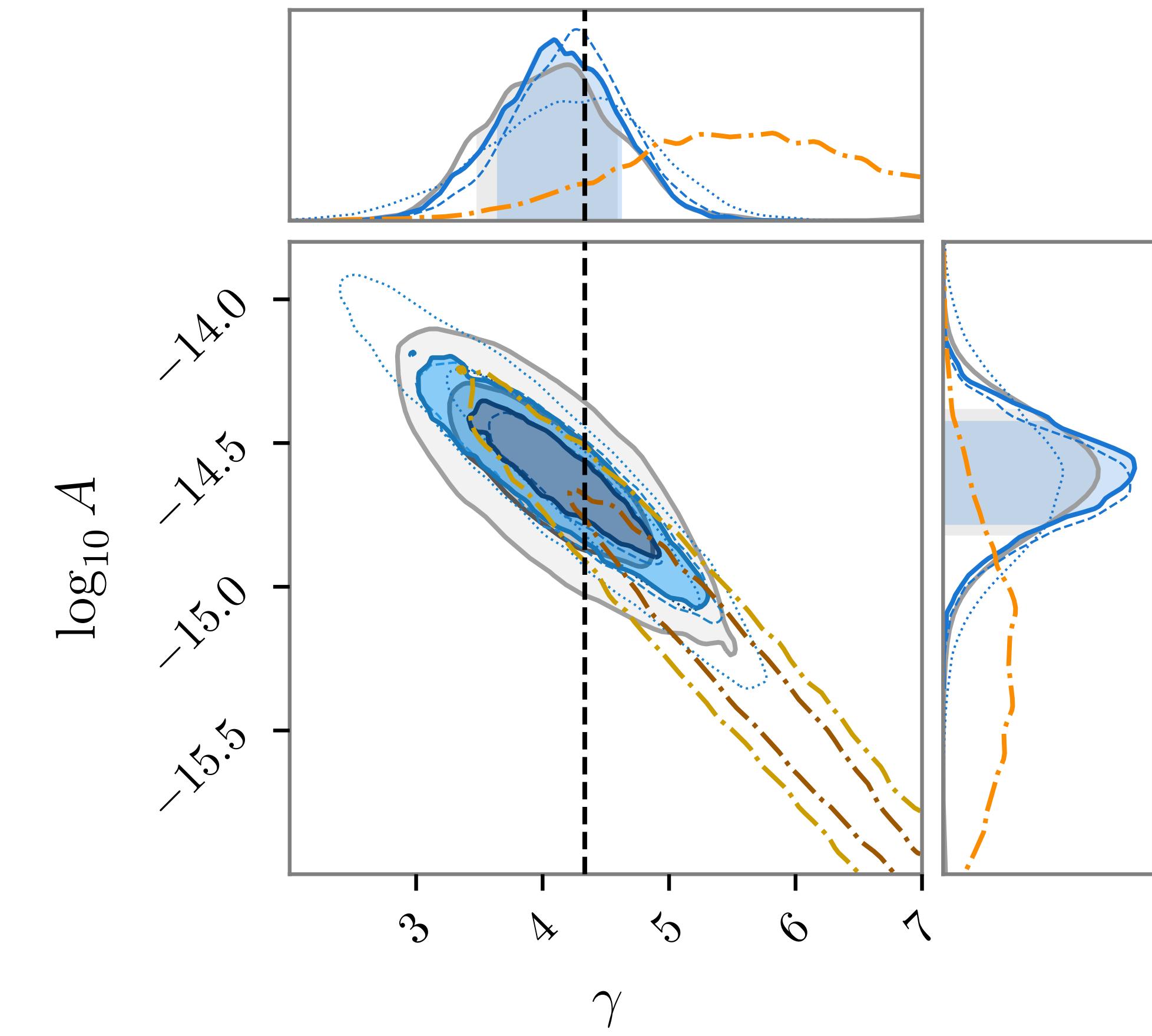
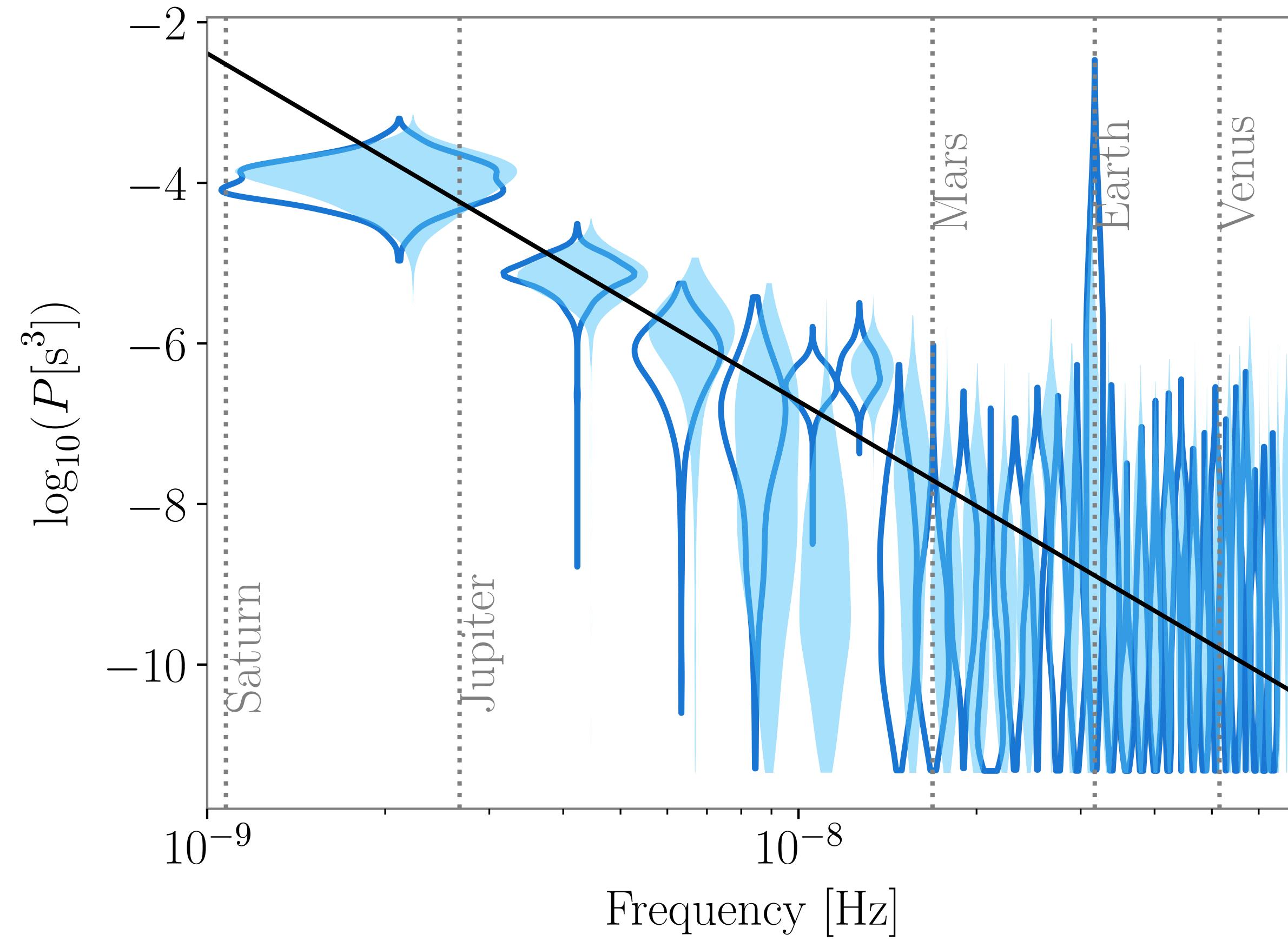
$$\rho(f) = \frac{A^2}{12\pi^2} \left( \frac{f}{1 \text{ yr}^{-1}} \right)^{-\gamma} \text{ yr}^2$$



# A FIRST GLIMPSE?

PPTA DR2 [2107.12112]

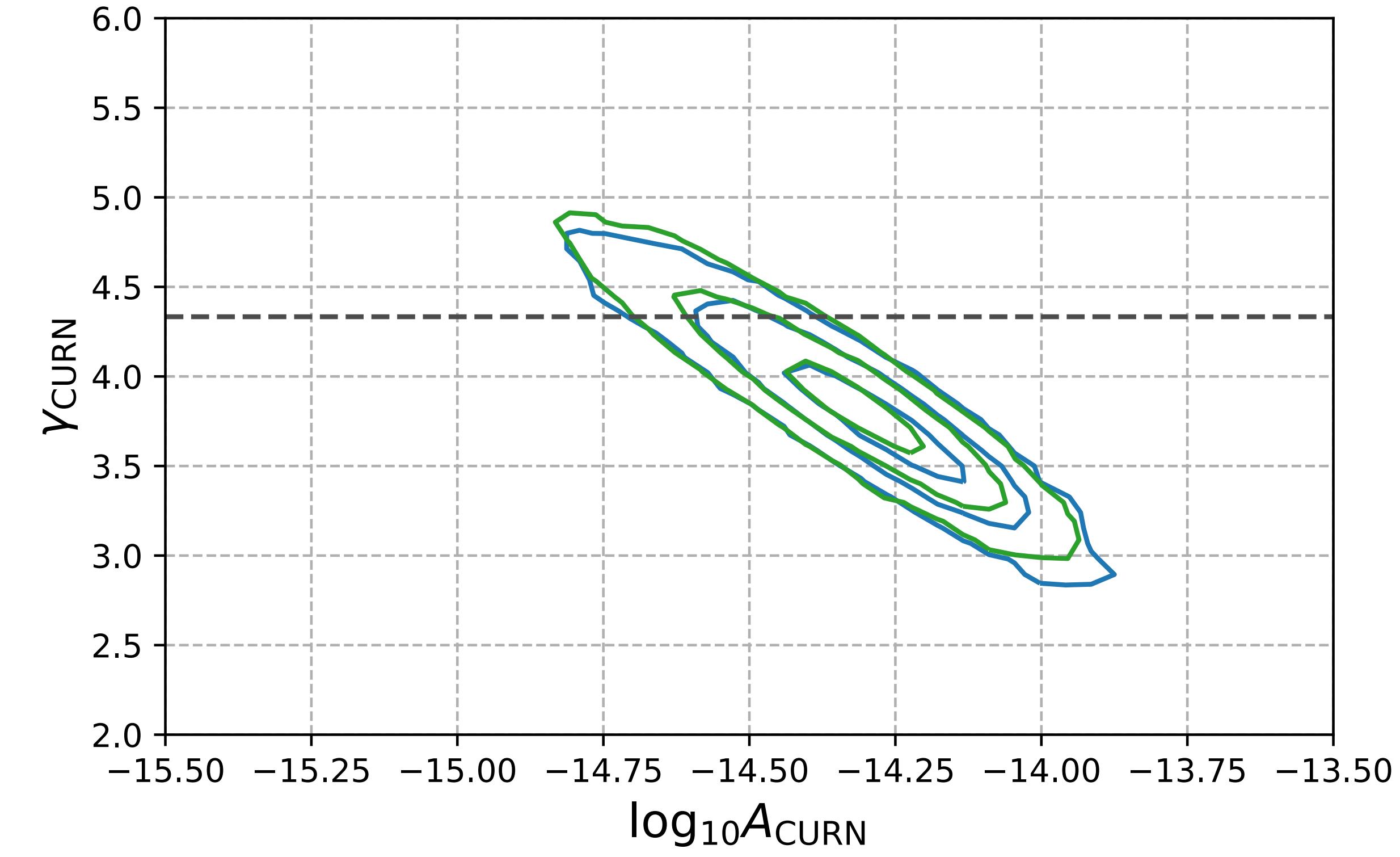
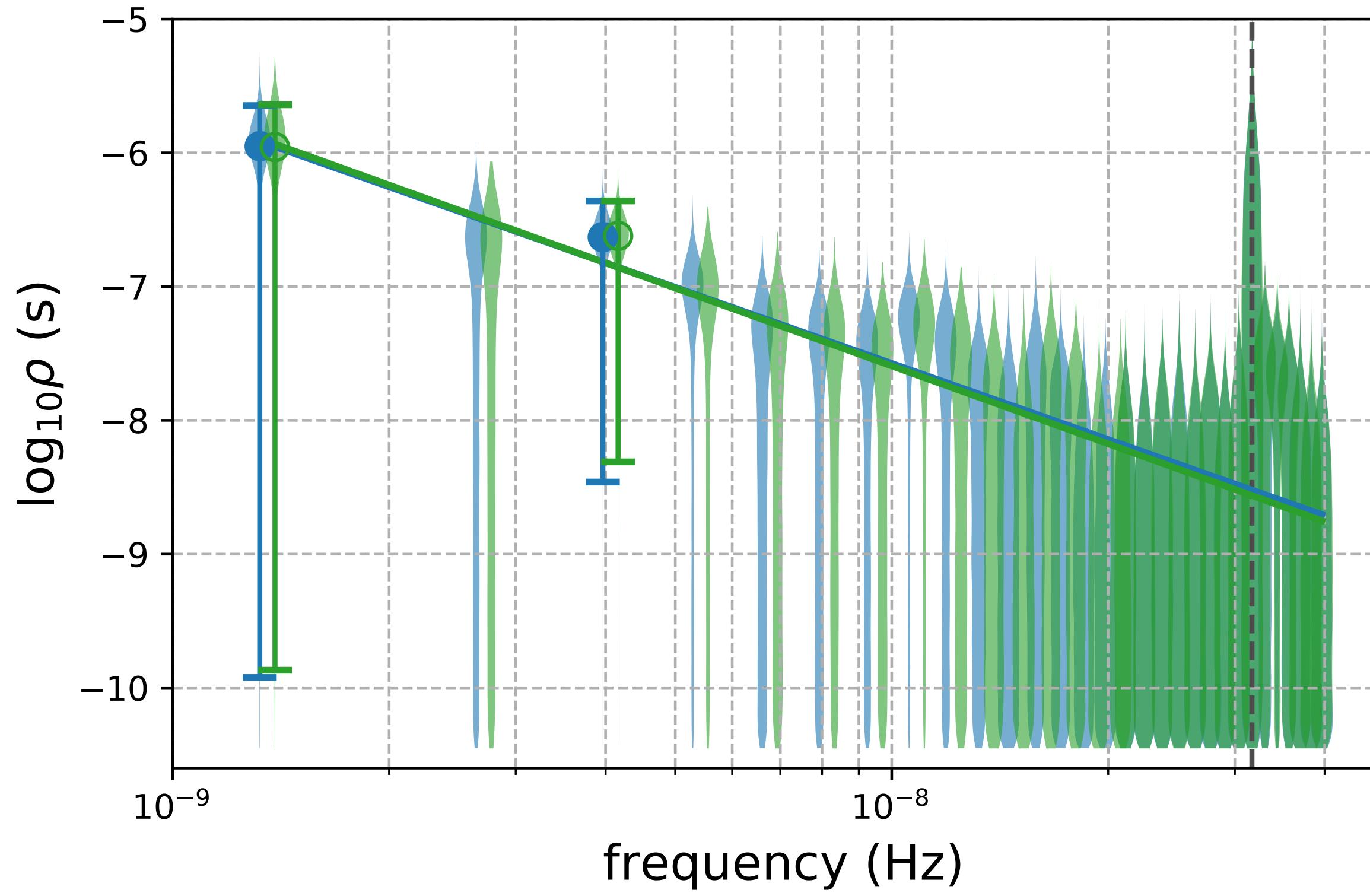
$$\rho(f) = \frac{A^2}{12\pi^2} \left( \frac{f}{1 \text{ yr}^{-1}} \right)^{-\gamma} \text{yr}^2$$



# A FIRST GLIMPSE?

[EPTA DR2 \[2110.13184\]](#)

$$\rho(f) = \frac{A^2}{12\pi^2} \left( \frac{f}{1 \text{ yr}^{-1}} \right)^{-\gamma} \text{ yr}^2$$



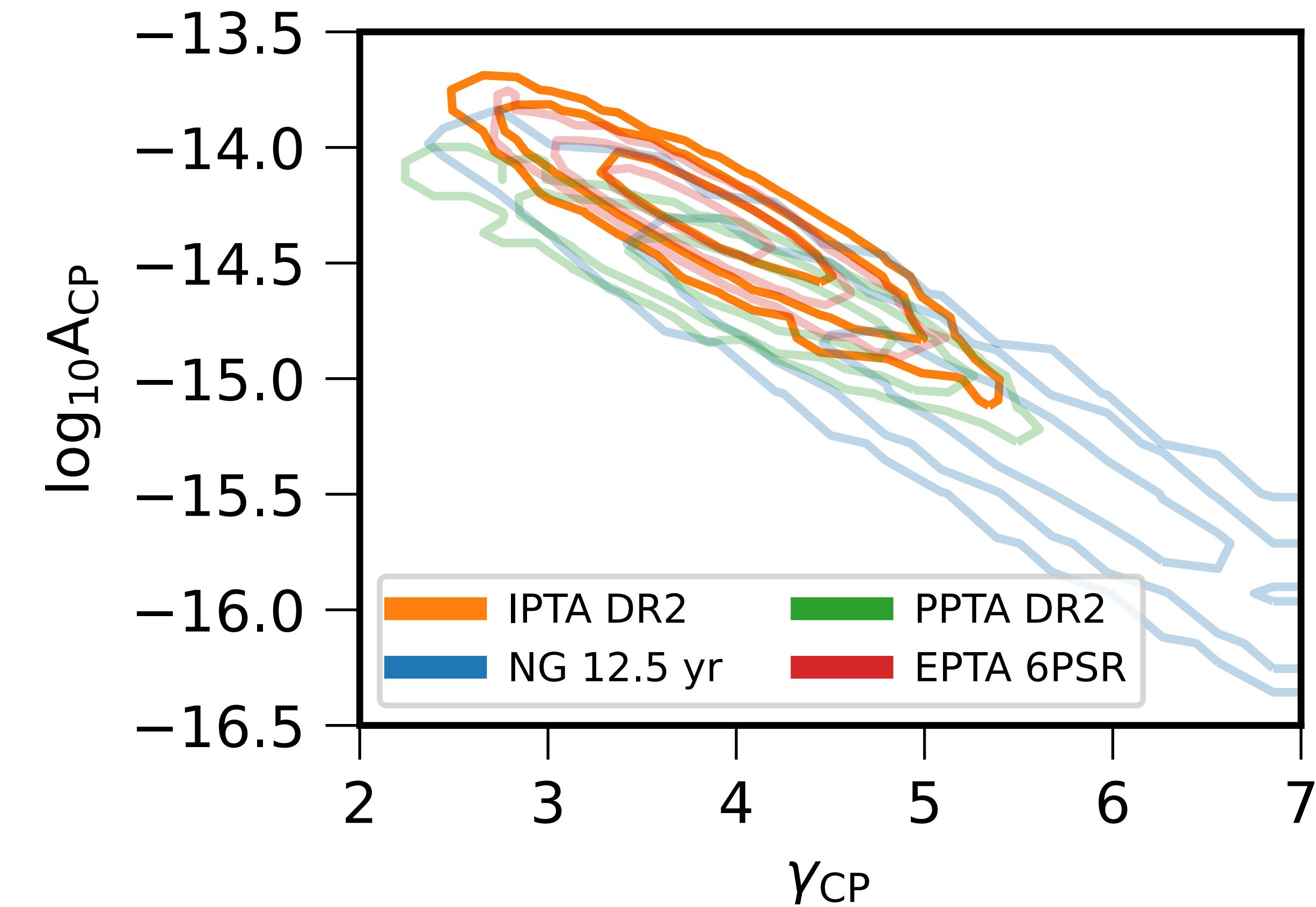
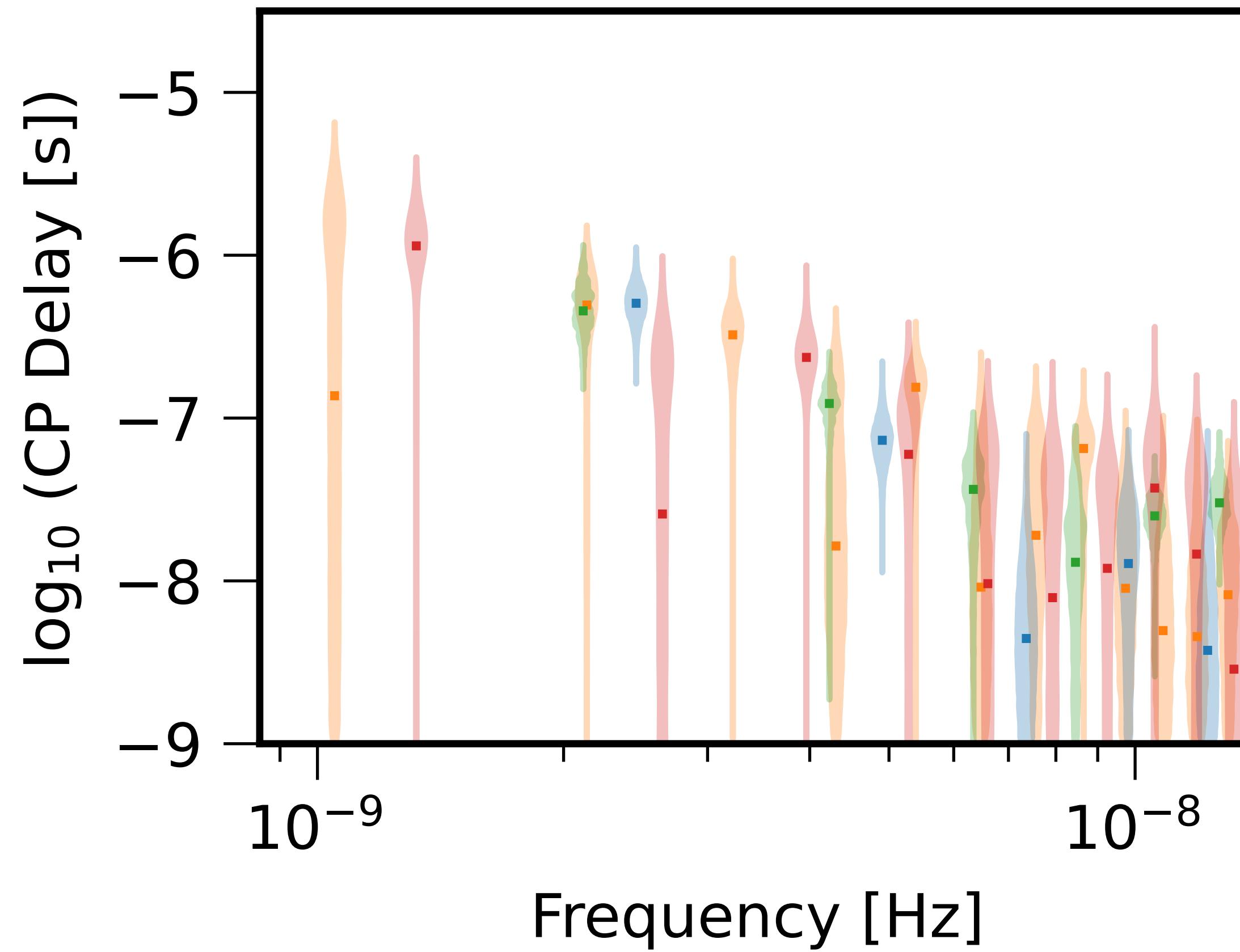
● DR2 EP Free Spectrum    — DR2 EP Power Law

○ DR2 42 Free Spectrum    — DR2 42 Power Law

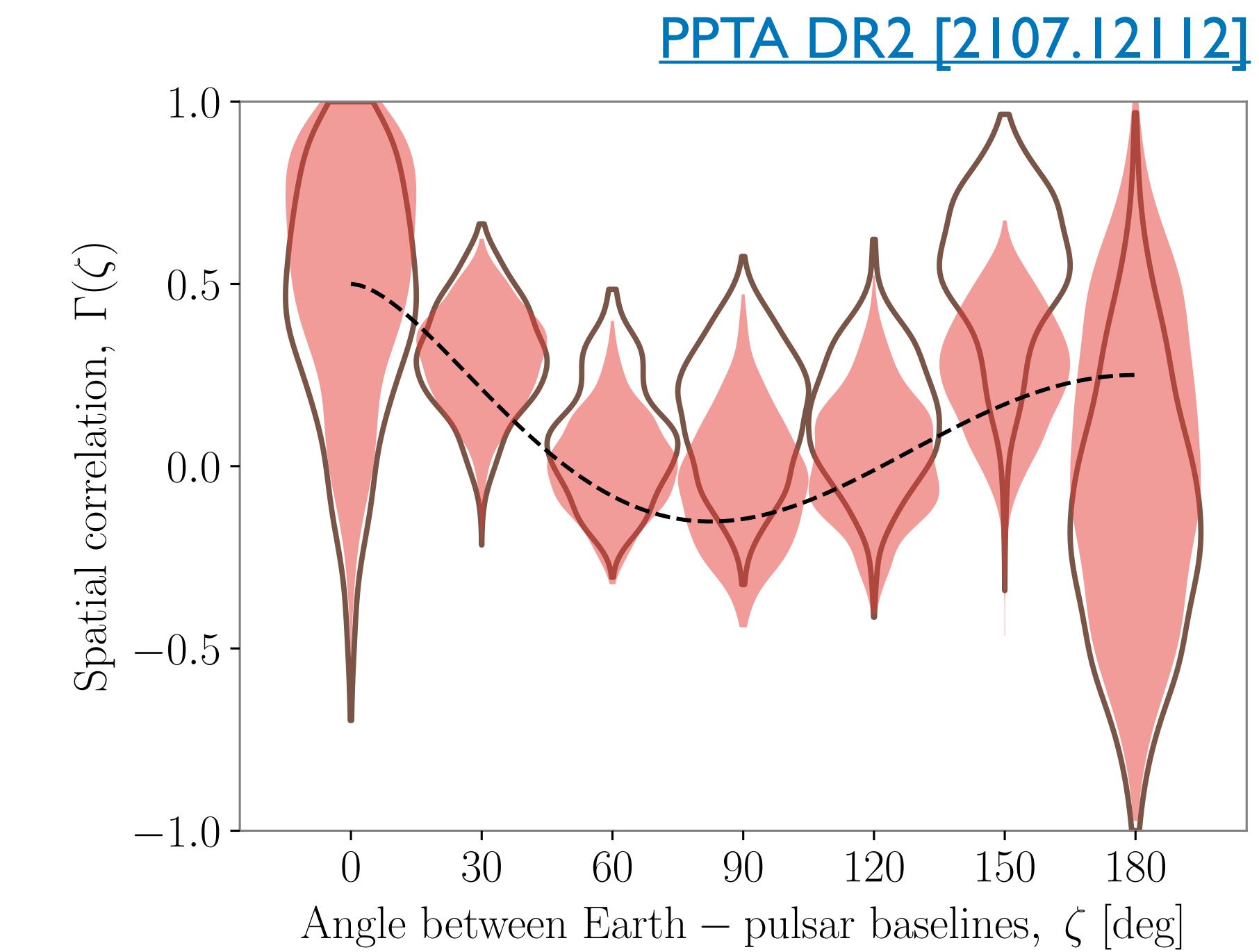
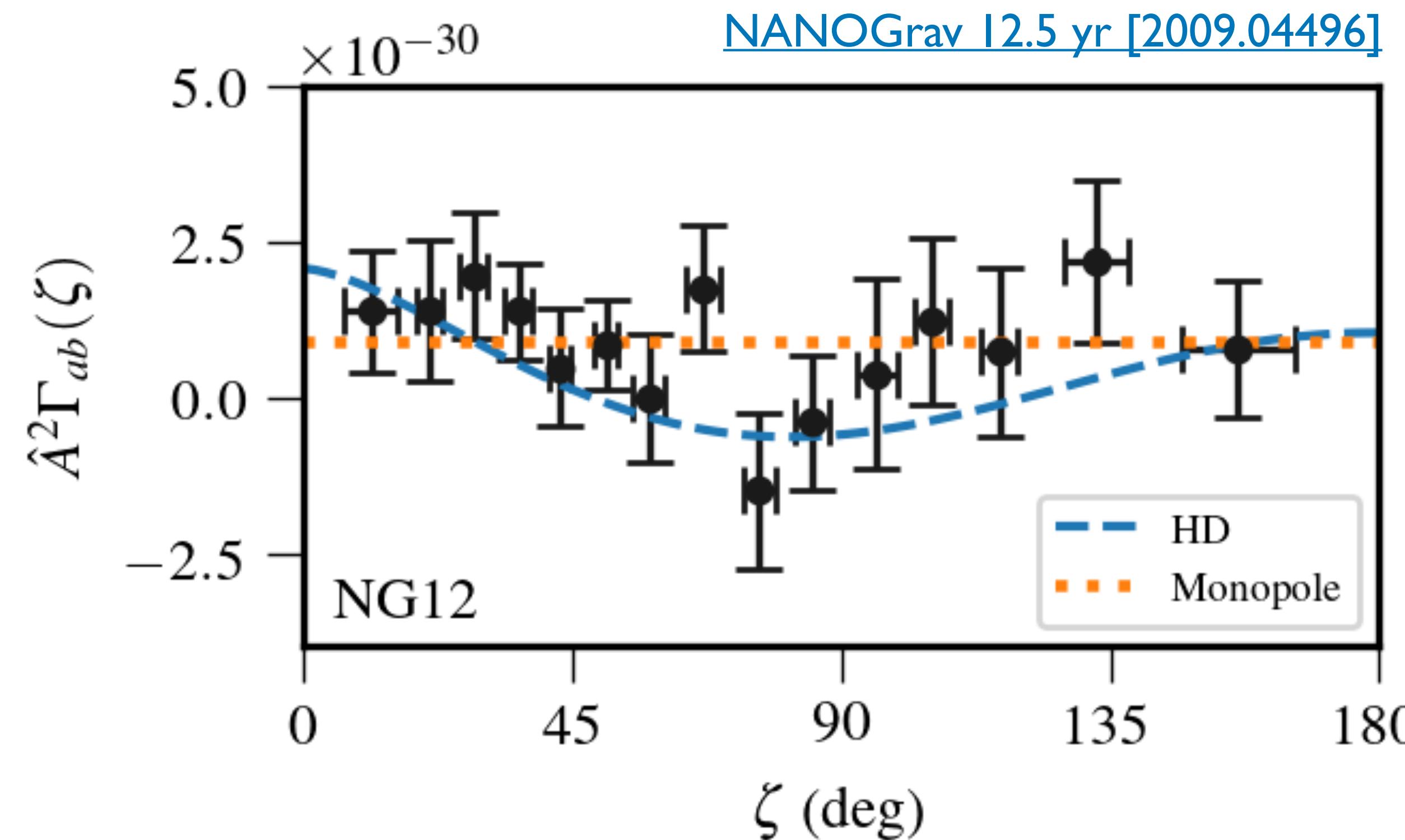
# A FIRST GLIMPSE?

[IPTA DR2 \[2201.03980\]](#)

$$\rho(f) = \frac{A^2}{12\pi^2} \left( \frac{f}{1 \text{ yr}^{-1}} \right)^{-\gamma} \text{ yr}^2$$



# WHAT ABOUT CORRELATIONS?

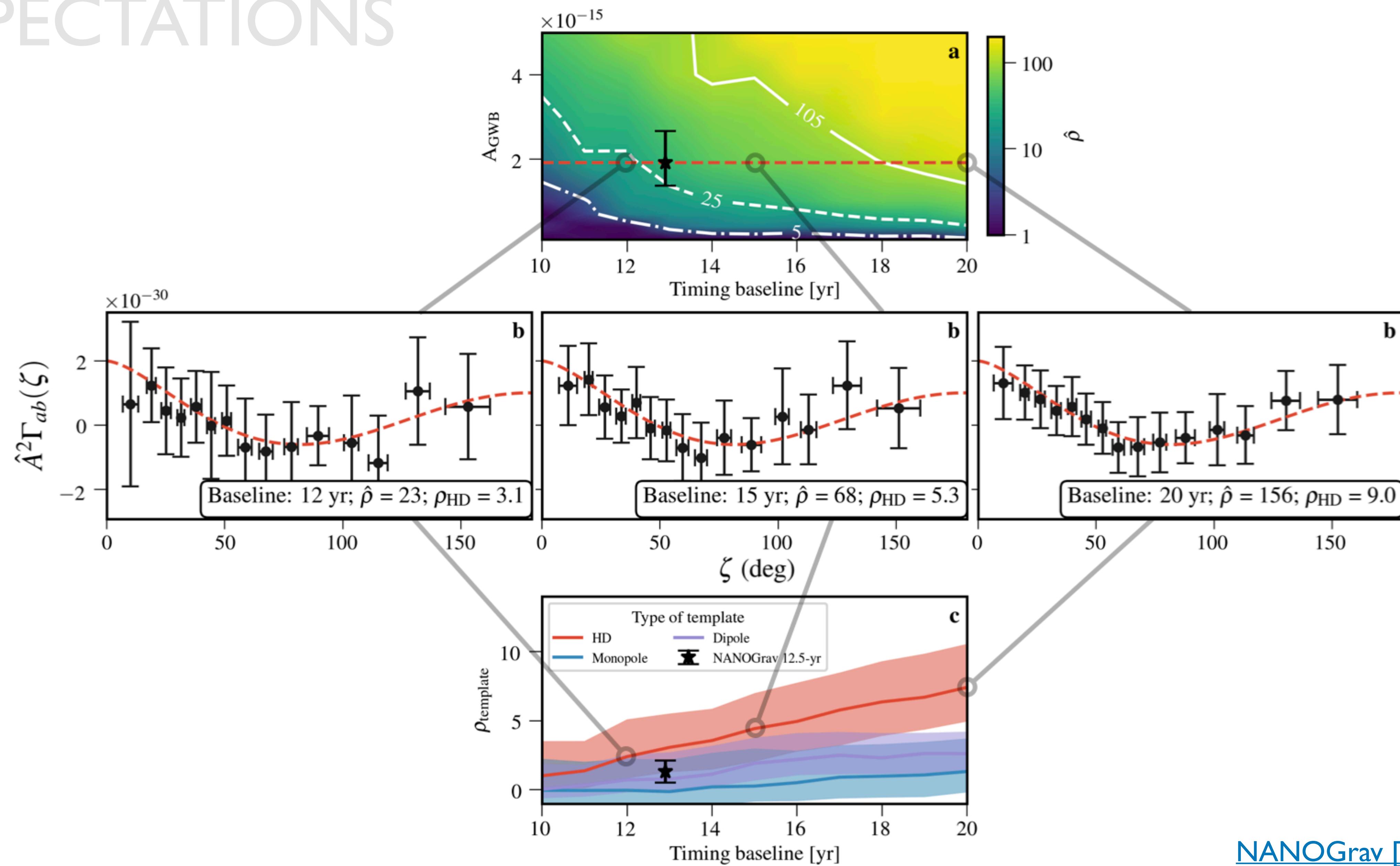


**no** evidence for **monopole** correlation (error in reference clocks)

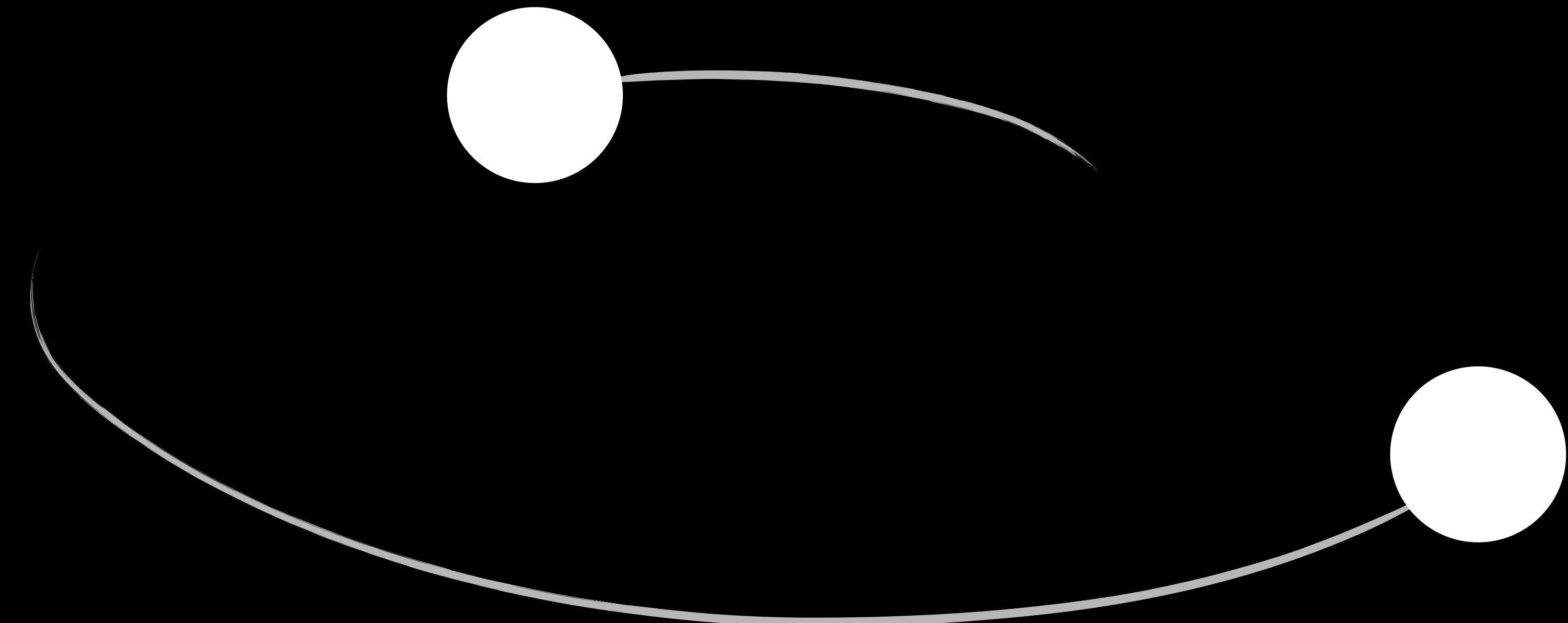
**no** evidence for **dipole** correlation (error in position of SSB)

evidences for Hellings and Downs correlation not yet conclusive

# EXPECTATIONS



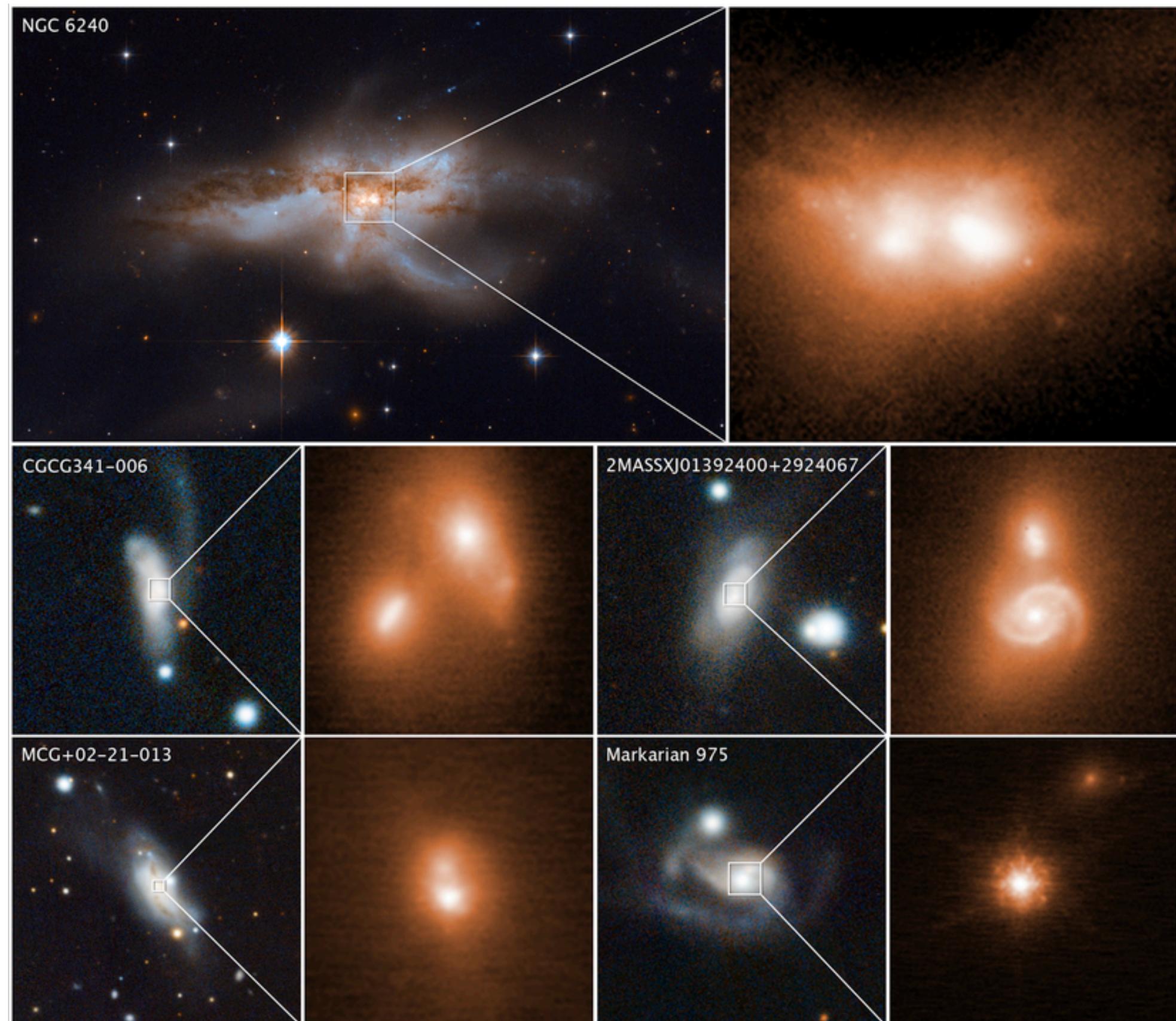
NANOGrav [2010.11950]



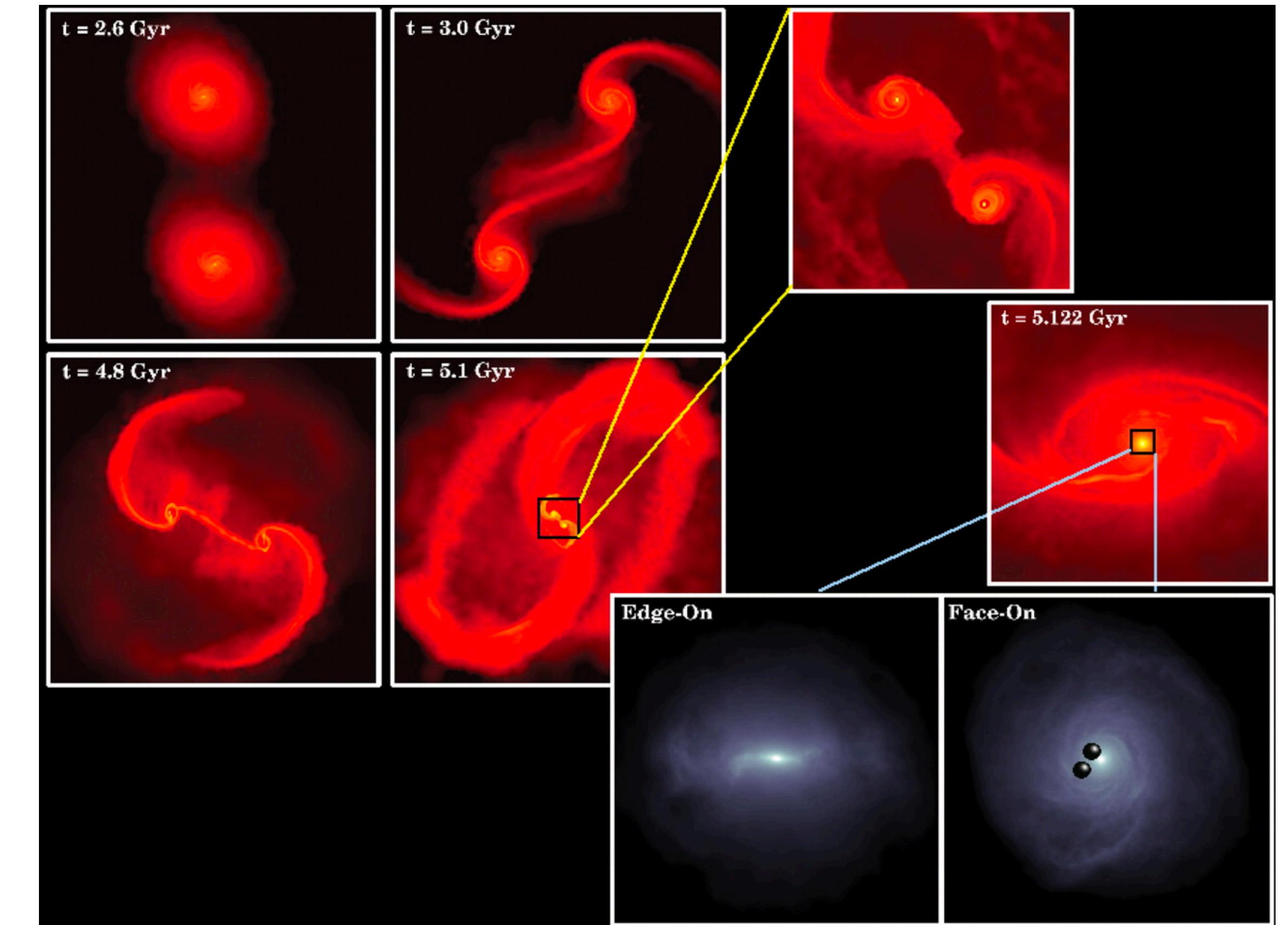
GWB from Supermassive Black Holes Binaries

# THE PRIMARY SUSPECT

[Aaron Evans \[HST Proposal 10592\]](#)



[Mayer et al. \[0706.1562\]](#)



Super Massive Black Holes Binaries (SMBHBs) formed in galaxy mergers

# BHBMs GW SPECTRUM

$$h_c^2(f) \sim f^{-4/3} \int dz d\mathcal{M} \frac{dN^2}{dz d\mathcal{M}} \frac{\mathcal{M}^{5/3}}{(1+z)^{1/3}}$$

merger rate      chirp mass  
↓                  ↓  
spectral shape       $A_{\text{GWB}}^2$

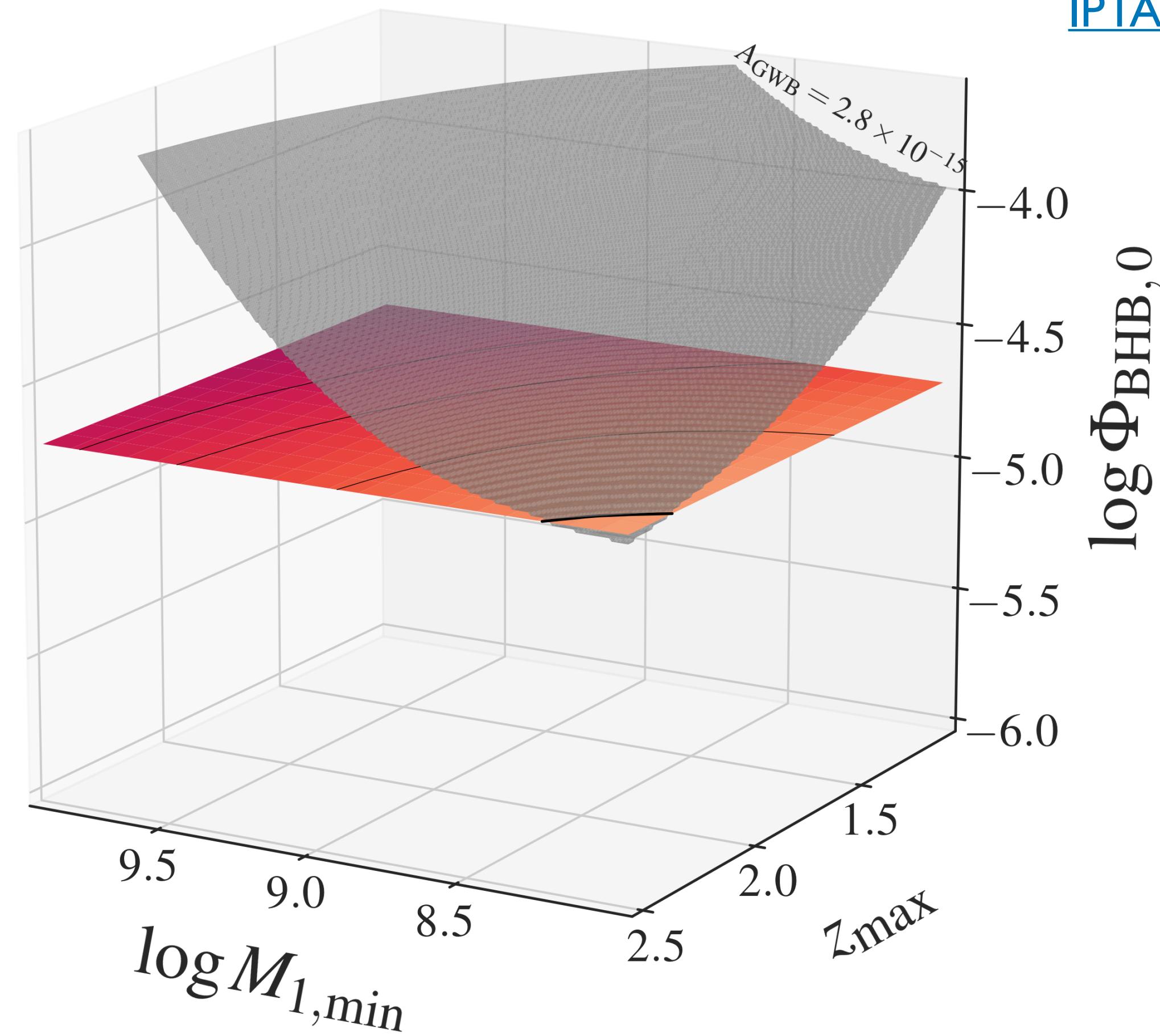
[Phinney \[astro-ph/0108028\]](#)

$$h_c(f) = A_{\text{GWB}} \left( \frac{f}{\text{yr}^{-1}} \right)^{-2/3} \quad \Leftrightarrow \quad S_{ab}(f) = \Gamma_{ab} \frac{A_{\text{GWB}}^2}{12\pi^2} \left( \frac{f}{f_{\text{yr}}} \right)^{-13/3}$$

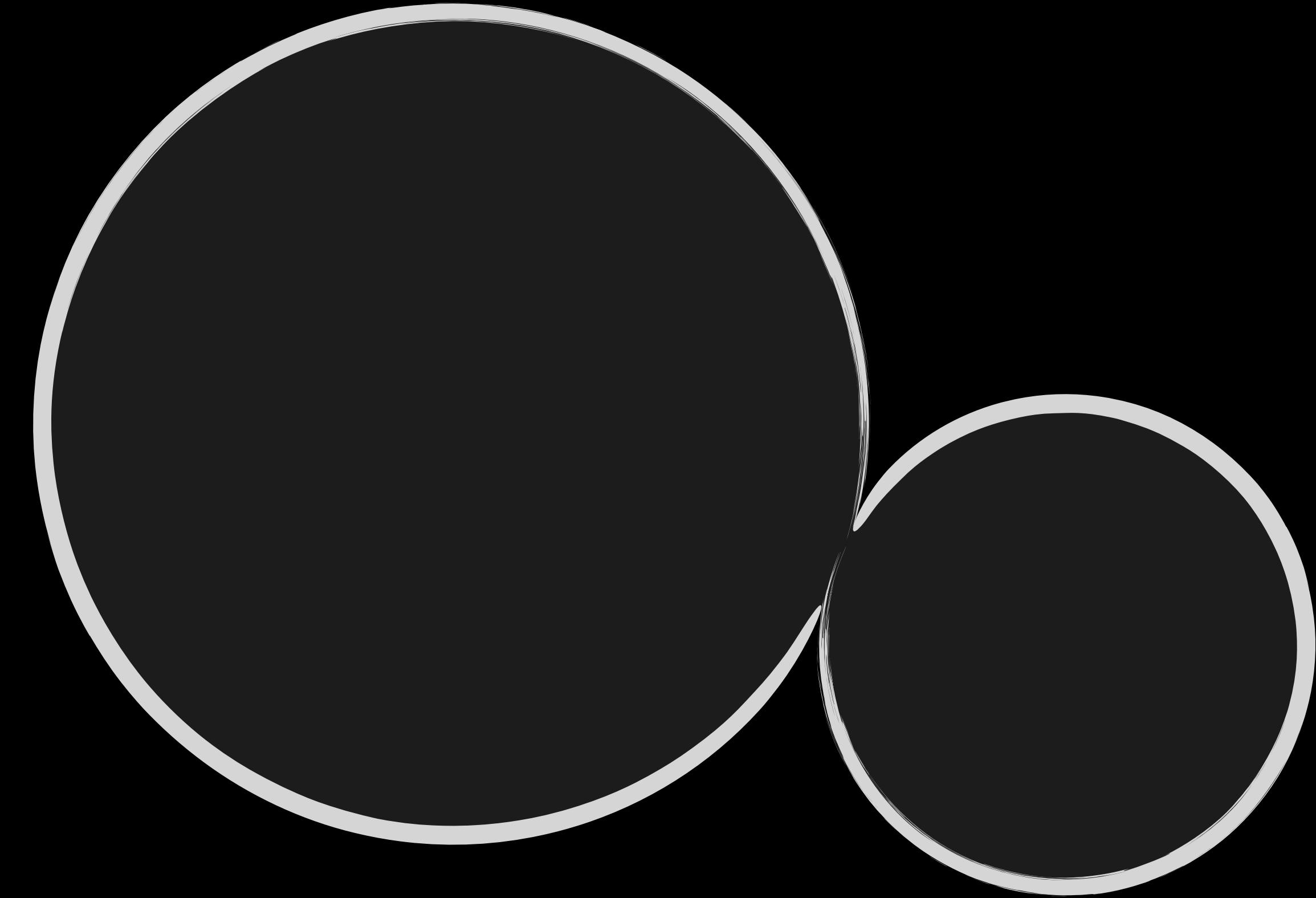
Many open questions: Origin? Delay after galaxy merger? Final-parsec problem? ...

# IS THE AMPLITUDE RIGHT?

[IPTA DR2 \[2201.03980\]](#)



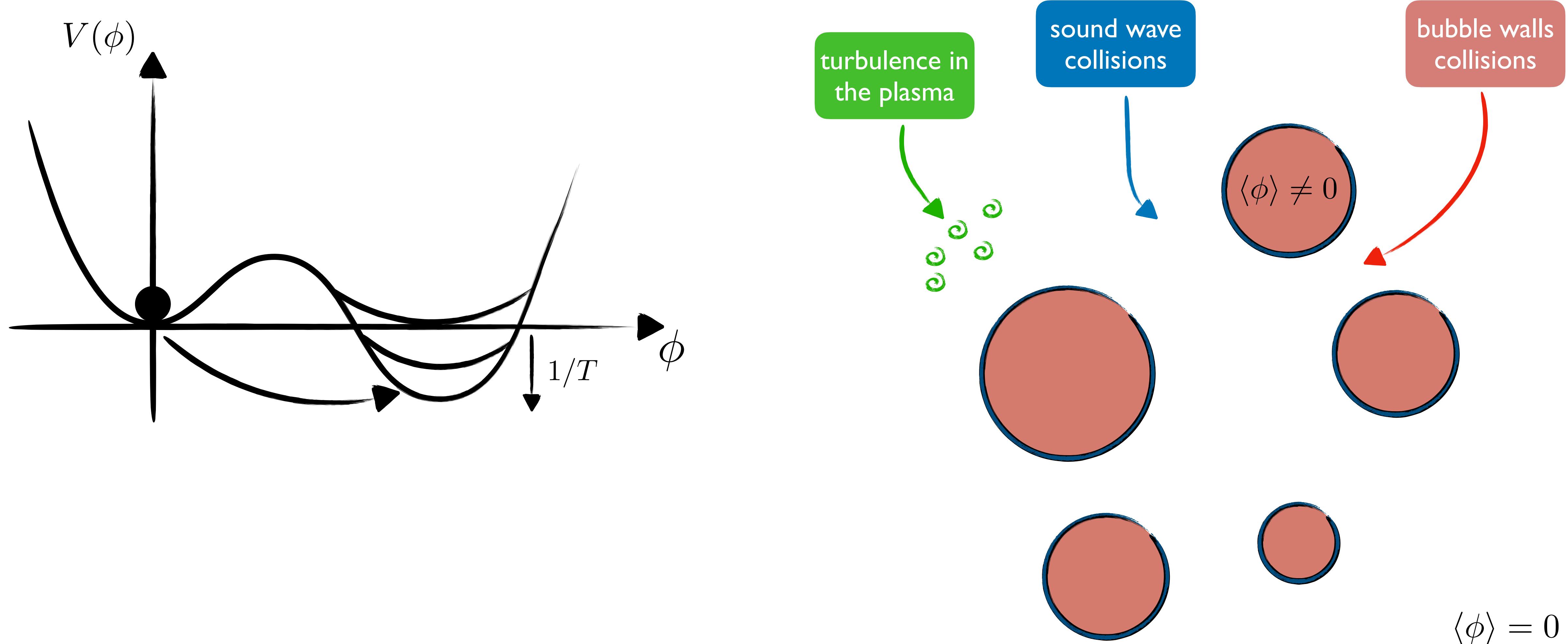
local SMBHB density,  $\Phi_{\text{BHB},0}$ , needs to be larger than previous estimates, e.g., [I708.03491](#), by a factor of  $\mathcal{O}(20)$  in order to explain the observed signal.



GWB from cosmological phase transitions

**three sources of GW**

$$\Omega_{\text{GW}}(f) = \Omega_\phi(f) + \Omega_{\text{sw}}(f) + \Omega_{\text{turb}}(f)$$



# SPECTRUM PARAMETRIZATION

$$\Omega(f) = \mathcal{R} \Delta(v_w) \left( \frac{\kappa \alpha_*}{1 + \alpha_*} \right)^p \left( \frac{H_*}{R_*} \right)^q \mathcal{S}(f/f_*^0)$$

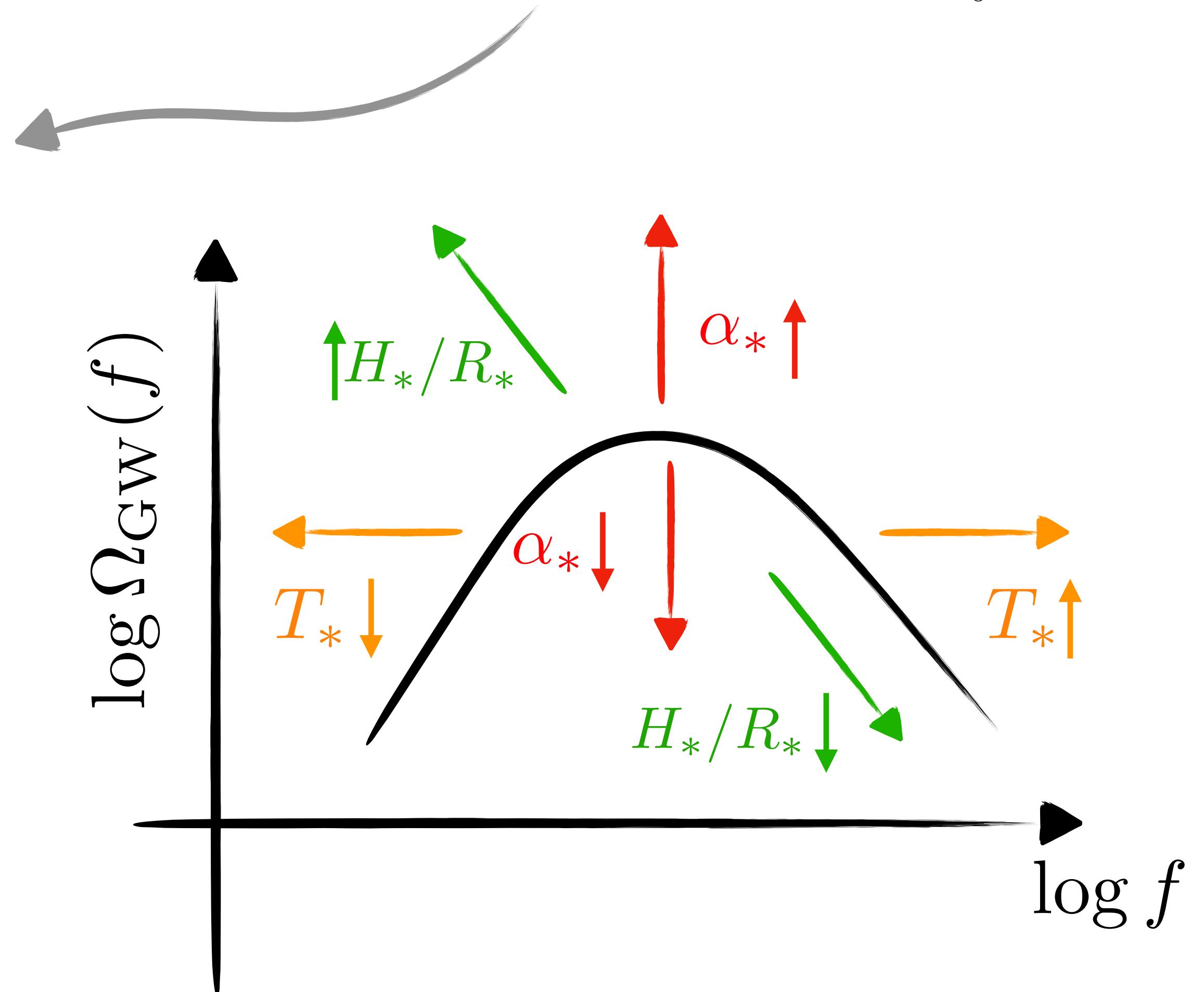
$$f_*^0 \sim 10^{-10} \text{ Hz} \left( \frac{\beta}{R_*} \right) \left( \frac{T_*}{\text{MeV}} \right)$$

$\alpha_*$  : transition strength

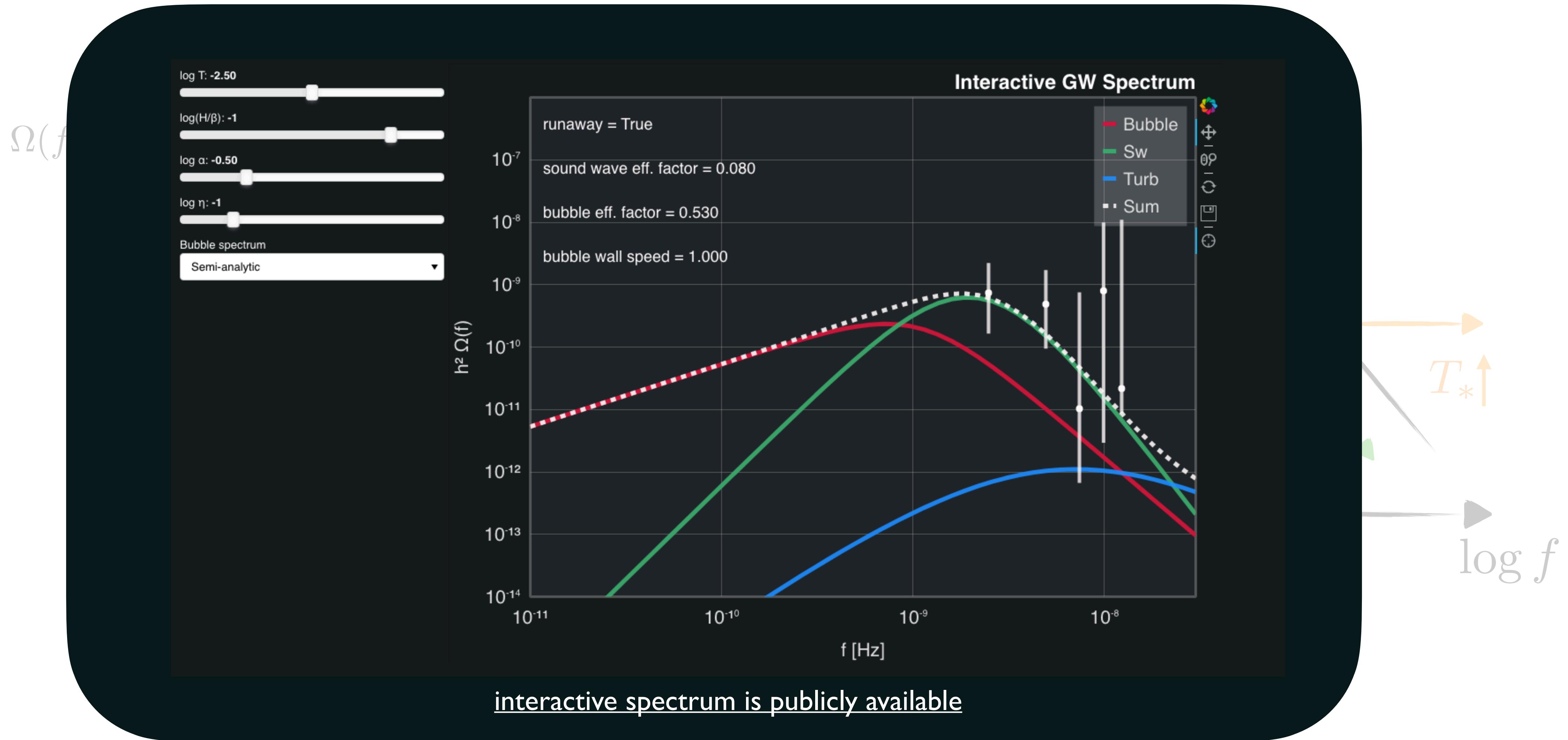
$H_*/R_*$  : transition timescale

$T_*$  : transition temperature

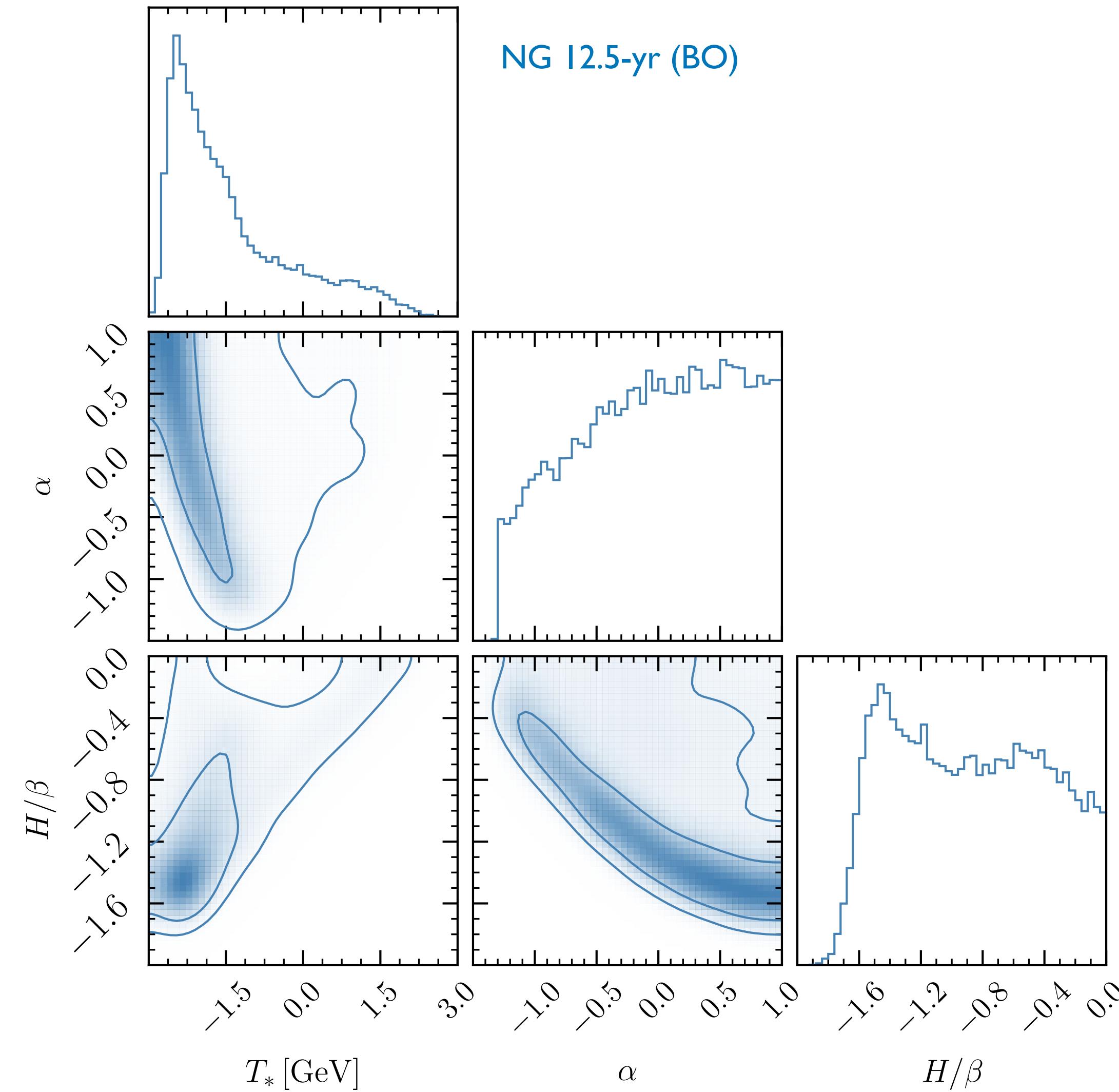
$$\Omega_{\text{gw}}(f) \equiv \frac{1}{\rho_c} \frac{d\rho_{\text{gw}}}{d \log f} = \frac{4\pi^2}{3H_0^2} f^3 S_h(f)$$

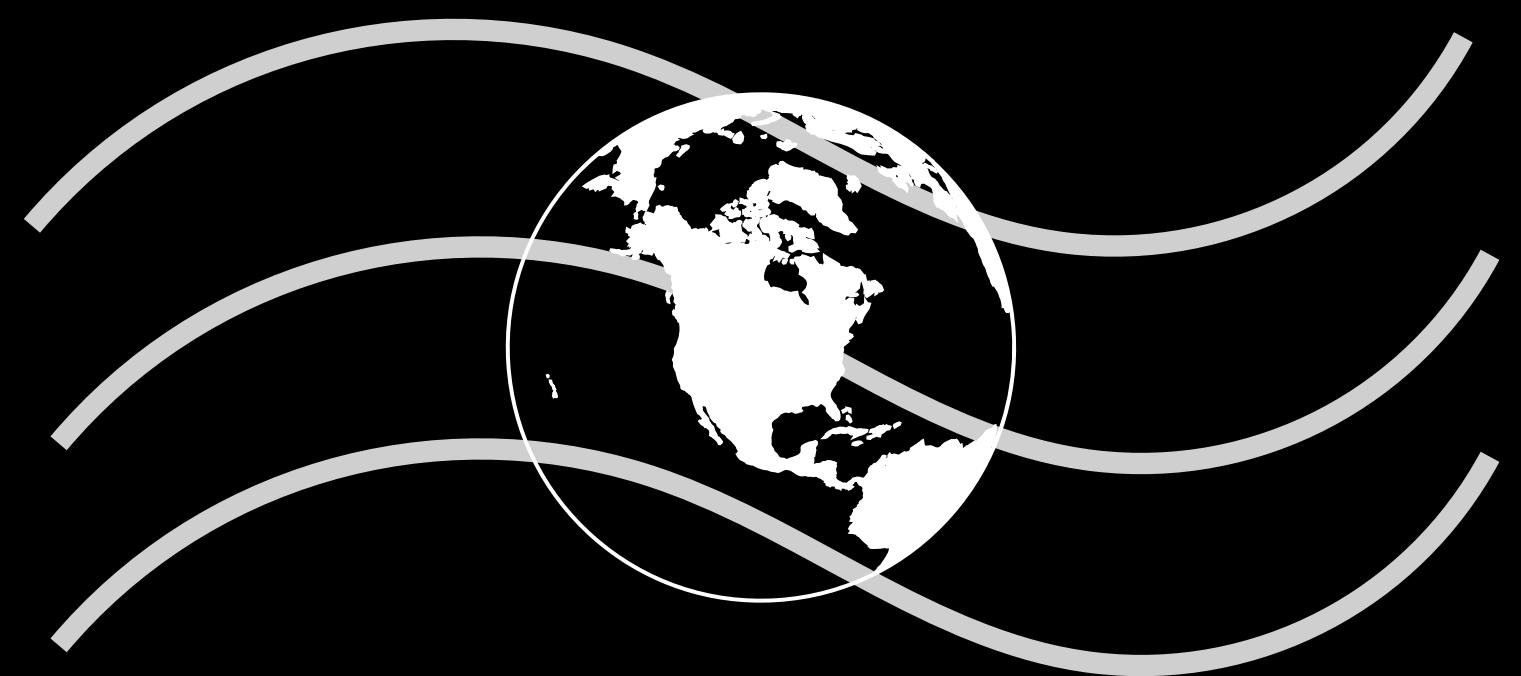
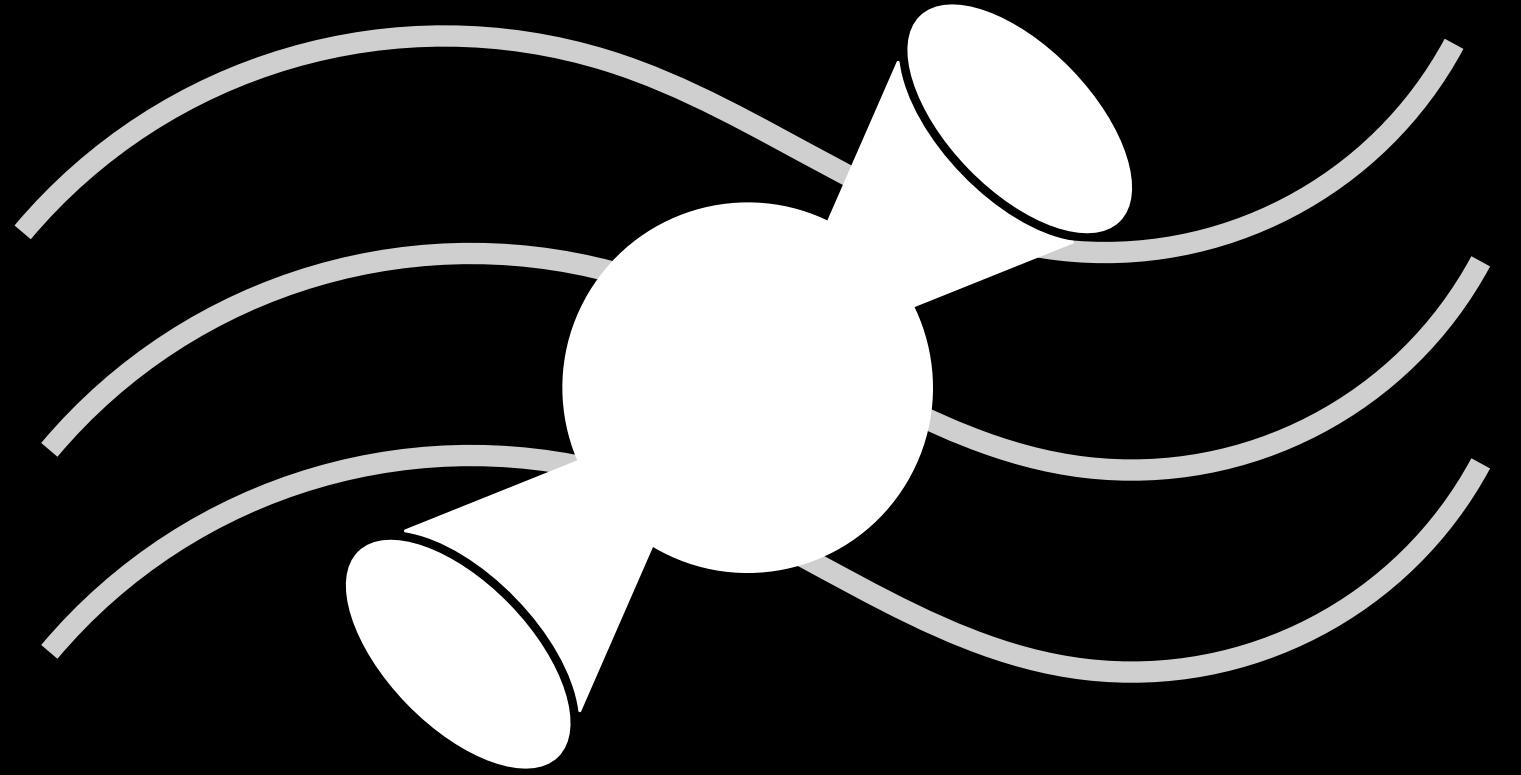


# SPECTRUM PARAMETRIZATION



# SOME RESULTS





ULDM signals in PTAs

# ULDM FIELD

$$10^{-24} \text{ eV} \lesssim m_\phi \ll \text{eV}$$

$$\rho_\phi \sim 0.4 \text{ GeV cm}^{-3}$$

lower masses are already ruled out by CMB observables

local ULDM density

$$\phi(\vec{x}, t) = \frac{\sqrt{2\rho_\phi}}{m_\phi} \hat{\phi}(\vec{x}) \cos(m_\phi t + \gamma(\vec{x}))$$

$$\hat{\phi}(\vec{x})$$

$$\gamma(\vec{x}) = \vec{k} \cdot \vec{x} + \tilde{\gamma}$$

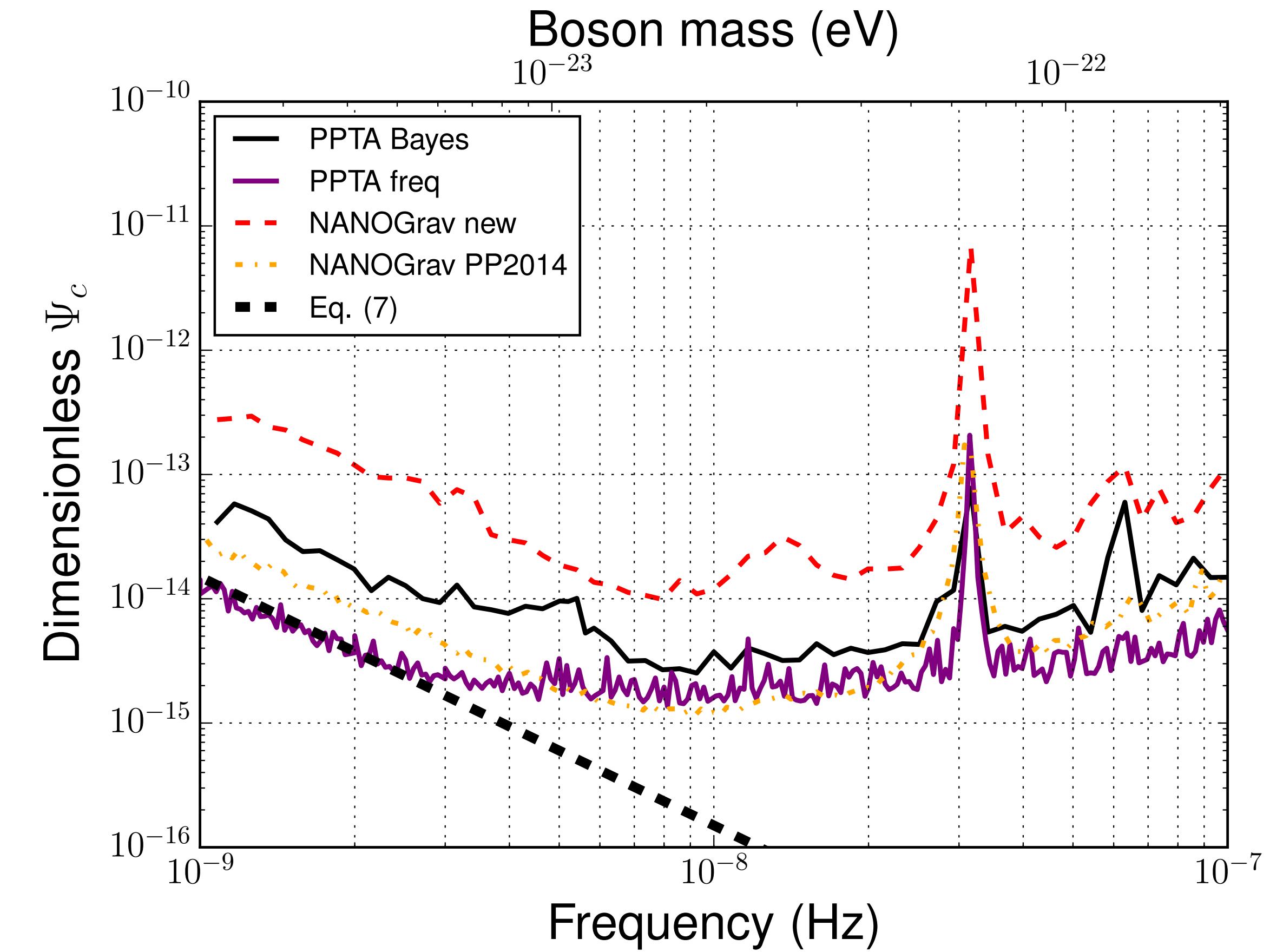
random amplitude with zero mean and unit variance

phase of the signal

# PREVIOUS SEARCHES - gravitational

$$\rho_\phi(t) \Rightarrow \delta g_{\mu\nu}(t) \Rightarrow \frac{\delta\omega}{\omega} \sim \int \dot{\Phi} dz$$

$$s(t) \sim \frac{G\rho_\phi}{m_\phi^3} \cos(2m_\phi t)$$



- [Pulsar timing signal from ultralight scalar dark matter](#)
- [Search for ultralight scalar dark matter with NANOGrav pulsar timing arrays](#)
- [Constraints on ultralight scalar dark matter from pulsar timing](#)
- [Parkes Pulsar Timing Array constraints on ultralight scalar-field dark matter](#)

[PPTA \[1810.03227\]](#)

# PREVIOUS SEARCHES - direct couplings

## • Doppler shift

- Scalar DM       $\mathbf{a} \sim \nabla\phi$

velocity suppressed

- Vector DM       $\mathbf{E} \sim \partial_{\mathbf{t}}\phi$

[High-precision search for dark photon dark matter with the Parkes Pulsar Timing Array -  
Phys. Rev. Research 4, L012022 \(2022\)](#)

## • Fifth force

- [Test of the Equivalence Principle Using a Rotating Torsion Balance - Phys. Rev. Lett. 100, 041101 \(2008\)](#)

- [MICROSCOPE Mission: First Constraints on the Violation of the Weak Equivalence Principle by a Light Scalar Dilaton - Phys. Rev. Lett. 120, 141101 \(2018\)](#)

## • Atomic clock shifts

- [Dark Matter Direct Detection with Accelerometers - \[1512.06165\]](#)

- [Precision Metrology Meets Cosmology: Improved Constraints on Ultralight Dark Matter from Atom-Cavity Frequency Comparisons - PRL 125, 201302 \(2020\)](#)

## • Orbital period in binary pulsars

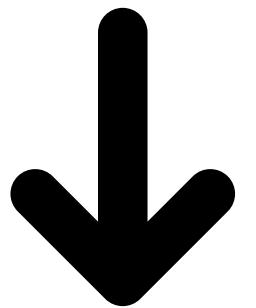
- [Ultralight Dark Matter Resonates with Binary Pulsars - Phys. Rev. Lett. 118, 261102 \(2017\)](#)

# NEW PTA SIGNALS - spin fluctuations

$$\mathcal{L}_{\phi,\text{QCD}} \supset \frac{\phi}{\Lambda} \left( \frac{d_g \beta_3}{2g_3} G_{\mu\nu}^A G_A^{\mu\nu} - \sum_{q=u,d} (d_q + \gamma_q d_g) m_q \bar{q} q \right)$$

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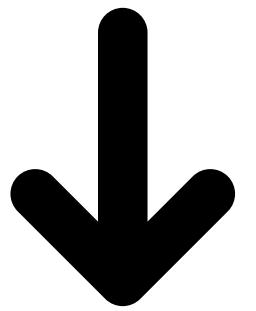
$$\frac{\delta m_{p,n}}{m_{p,n}} \simeq \frac{1}{\Lambda} \left( d_g + C_n d_{\hat{m}} \right) \phi$$

$$C_n \sim 0.048$$

[T. Damour, J. F. Donoghue \[1007.2792\]](#)

# NEW PTA SIGNALS - spin fluctuations

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$$\frac{\delta m_{p,n}}{m_{p,n}} \simeq \frac{1}{\Lambda} \left[ \begin{array}{c|c} \hline d_g & C_n d_{\hat{m}} \\ \hline \end{array} \right] \phi$$

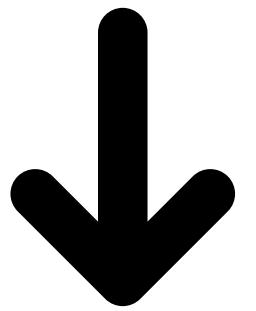
QCD scale quark masses  
fluctuations fluctuations

$$C_n \sim 0.048$$

[T. Damour, J. F. Donoghue \[1007.2792\]](#)

# NEW PTA SIGNALS - spin fluctuations

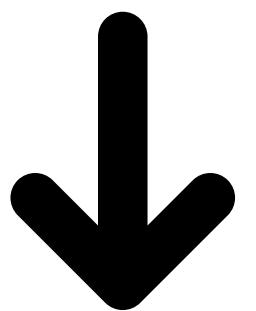
$$\mathcal{L}_{\phi, \text{QCD}} \supset \frac{\phi}{\Lambda} \left( \frac{d_g \beta_3}{2g_3} G_{\mu\nu}^A G_A^{\mu\nu} - \sum_{q=u,d} (d_q + \gamma_q d_g) m_q \bar{q} q \right)$$



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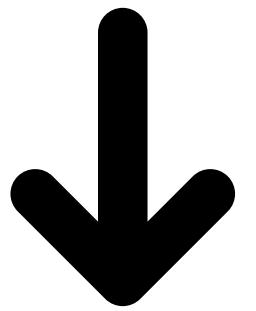
[T. Damour, J. F. Donoghue \[1007.2792\]](#)



$$\frac{\delta I}{I_0} = \eta \frac{\delta M}{M_0} + \delta \eta \frac{\delta m_n}{m_n}$$

# NEW PTA SIGNALS - spin fluctuations

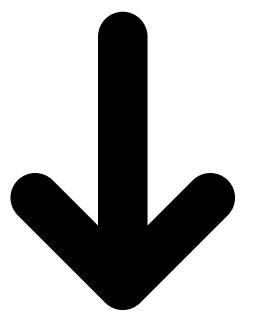
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$$\frac{\delta m_{p,n}}{m_{p,n}} \simeq \frac{1}{\Lambda} \left( d_g + C_n d_{\hat{m}} \right) \phi$$

$$C_n \sim 0.048$$

[T. Damour, J. F. Donoghue \[1007.2792\]](#)



in the simplest model of a spherical non-rotating neutron star

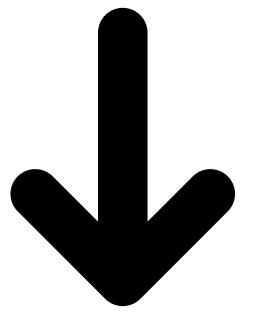
$$\eta = 1/3$$

$$\delta\eta = -16/3$$

$$\frac{\delta I}{I_0} = \boxed{\eta} \frac{\delta M}{M_0} + \boxed{\delta\eta} \frac{\delta m_n}{m_n}$$

# NEW PTA SIGNALS - spin fluctuations

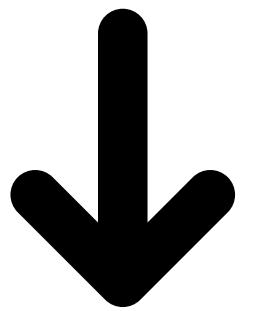
$$\mathcal{L}_{\phi, \text{QCD}} \supset \frac{\phi}{\Lambda} \left( \frac{d_g \beta_3}{2g_3} G_{\mu\nu}^A G_A^{\mu\nu} - \sum_{q=u,d} (d_q + \gamma_q d_g) m_q \bar{q} q \right)$$



$$\frac{\delta m_{p,n}}{m_{p,n}} \simeq \frac{1}{\Lambda} \left( d_g + C_n d_{\hat{m}} \right) \phi$$

$$C_n \sim 0.048$$

[T. Damour, J. F. Donoghue \[1007.2792\]](#)



in the simplest model of a spherical neutron star

$$\eta = 1/3$$

$$\delta\eta = -16/3$$

$$\frac{\delta I}{I_0} = n \frac{\delta M}{M_0} + \delta\eta \frac{\delta m_n}{m_n}$$

$$\frac{\delta M}{M_0} = \sum_{f \in \{e, \mu, p, n\}} Y_f \frac{m_f}{m_n} \frac{\delta m_f}{m_f}$$

# NEW PTA SIGNALS - spin fluctuations

$$\mathcal{L}_{\phi, \text{QCD}} \supset \frac{\phi}{\Lambda} \left( \frac{d_g \beta_3}{2 a_3} G_{\mu\nu}^A G_A^{\mu\nu} - \sum (d_q + \gamma_q d_g) m_q \bar{q} q \right)$$

$$h(t) = \frac{\sqrt{2\rho_\phi}}{m_\phi^2 \Lambda} \left( \vec{y} \cdot \vec{d} \right) \hat{\phi}_P \sin \left( m_\phi t + \gamma(\vec{x}_P) \right)$$

where  $\vec{d}$  contains the ULDM couplings and  $\vec{y}$  contains the sensitivity parameters:

$$\{y_g, y_{\hat{m}}, y_\mu, y_e\} = \eta \{1, C_n, 6 \times 10^{-3}, 5 \times 10^{-5}\} + \delta\eta \{1, C_n, 0, 0\}$$

in the

$$\delta\eta = -16/3$$



# NEW PTA SIGNALS - clock shift

pulses arrival time are measured with atomic clocks (most of them based on the hyperfine transition of the ground state of the Cesium atom)

$$f \propto \left(m_e \alpha^2\right) \left[\alpha^2 F_{\text{rel}}(Z\alpha)\right] \left(\mu \frac{m_e}{m_p}\right)^\zeta$$

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atomic      relativistic      hyperfine  
energy scale    correction    splitting

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ULDM fluctuations induce frequency fluctuations

$$\frac{\delta f_A}{f_A} \simeq \left[ \frac{\delta m_e}{m_e} + (4 + K_A) \frac{\delta \alpha}{\alpha} + \zeta \left( \frac{\delta m_e}{m_e} + C_A \sum_{q=u,d} \frac{\delta m_q}{m_q} - \frac{\delta m_p}{m_p} \right) \right]$$

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$$\frac{\delta F_{\text{rel}}}{F_{\text{rel}}} = K_A \frac{\delta \alpha}{\alpha} \quad \frac{\delta \mu}{\mu} = C_A \frac{\delta m_q}{m_q}$$

# NEW PTA SIGNALS - clock shift

pulses arrival time are measured with atomic clocks (most of them based on the hyperfine transition of the ground state of the Cesium atom)

$$h(t) = \frac{\sqrt{2\rho_\phi}}{m_\phi^2 \Lambda} \left( \vec{y} \cdot \vec{d} \right) \hat{\phi}_E \sin(m_\phi t + \gamma(\vec{x}_E))$$

where  $\vec{d}$  contains the ULDM couplings an  $\vec{y}$  contains the sensitivity parameters:

$$\{y_g, y_\gamma, y_{\hat{m}}, y_e\} \simeq \left\{ \zeta, \xi_A, \zeta \left( C_n + \hat{C}_A \right), 1 + \zeta \right\}$$

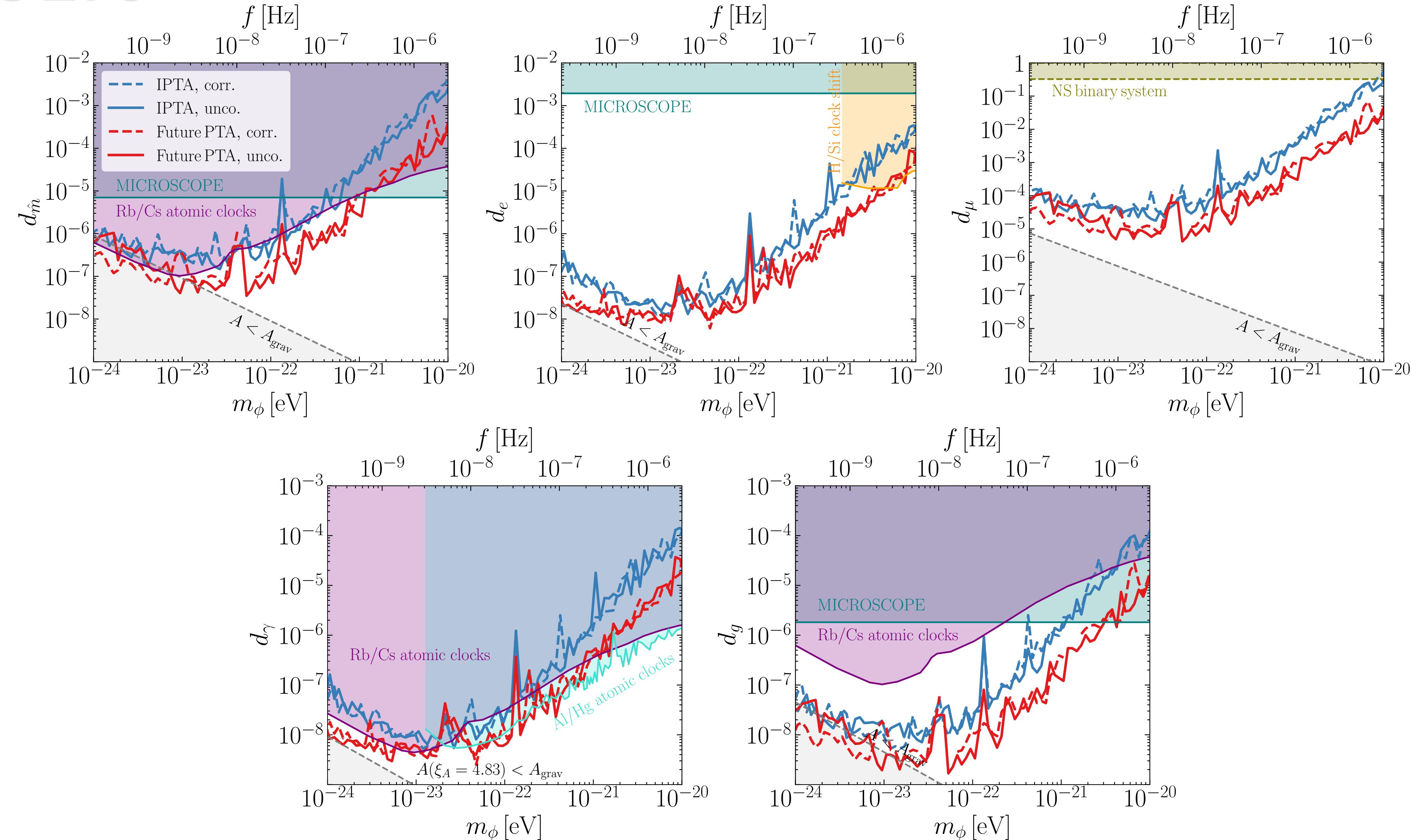
$F_{\text{rel}}$

$\alpha$

$\mu$

$m_q$

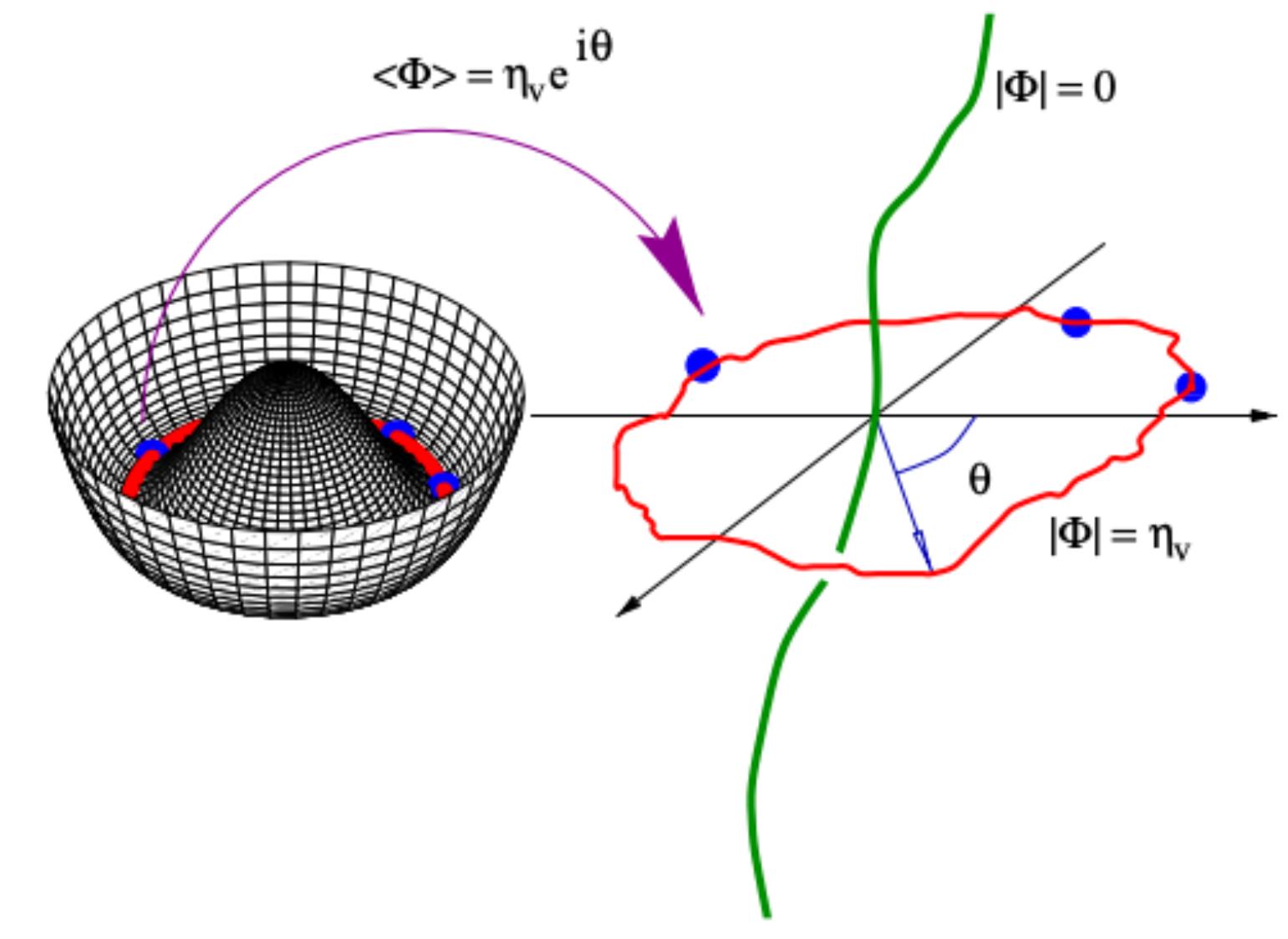
# RESULTS



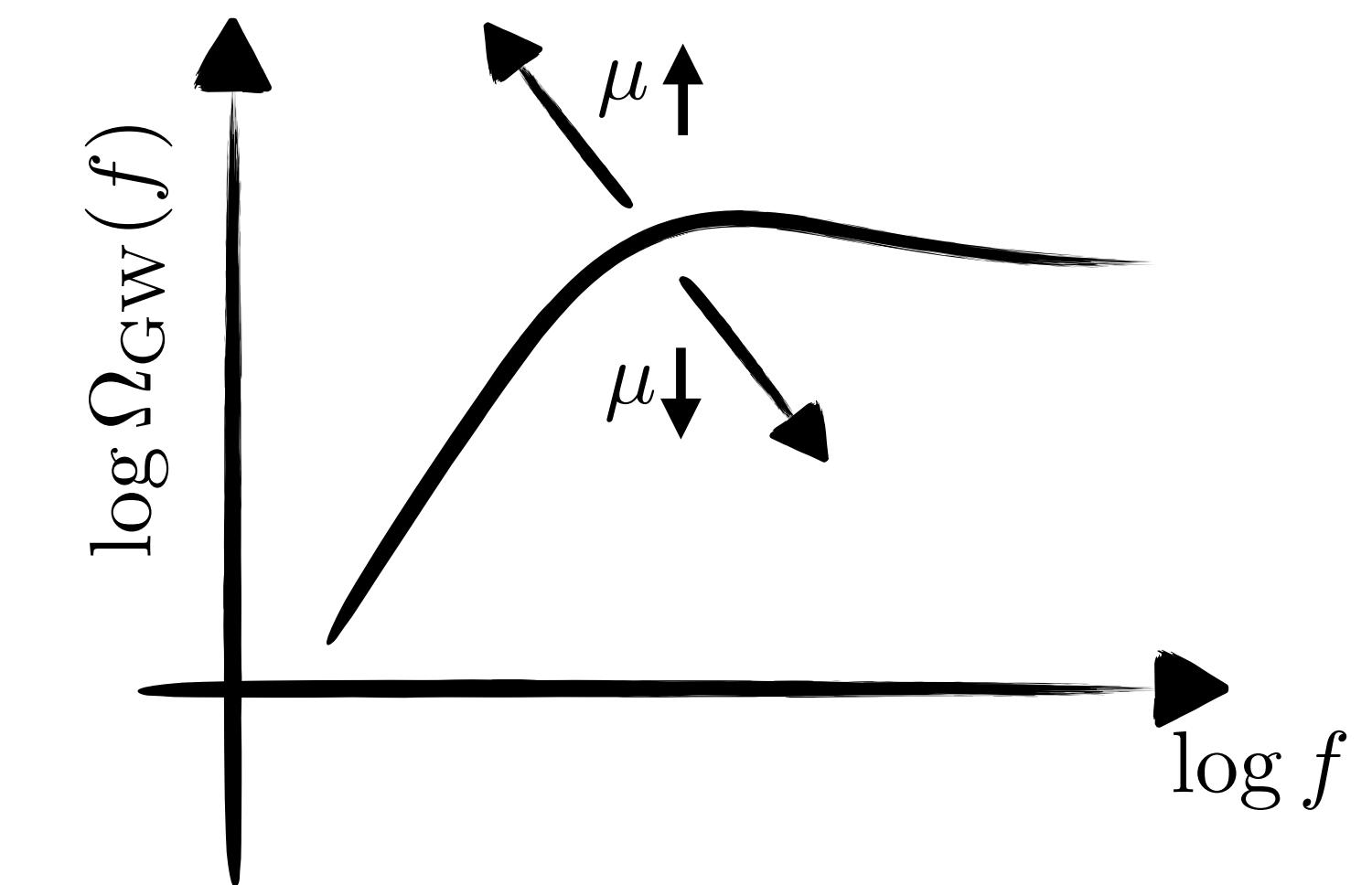
# COSMIC STRINGS



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$\mu$  : string tension



# DM SUBSTRUCTURES & PBH

single transiting objects give deterministic signal that can be easily parametrized

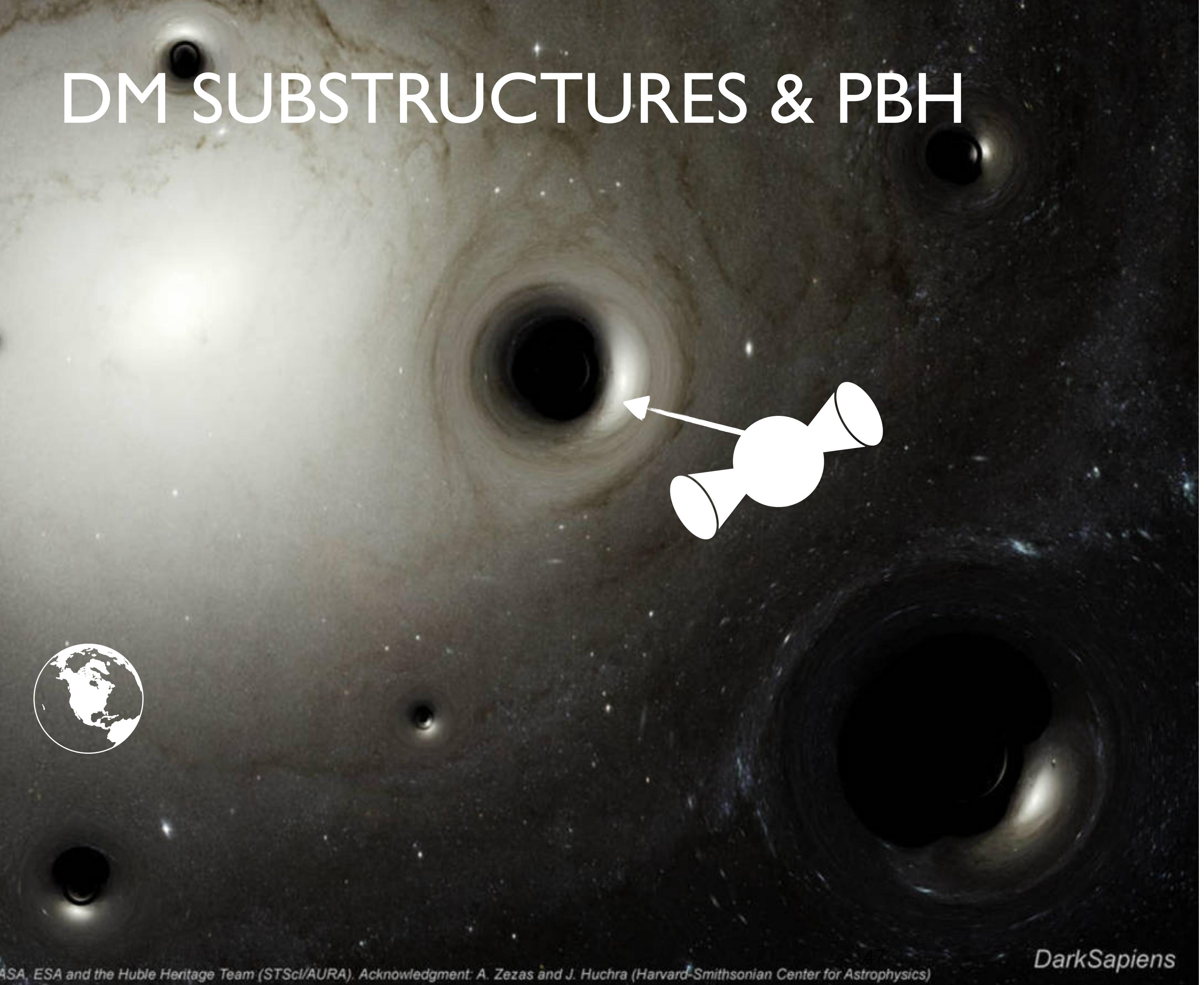
dynamic events ( $\tau \ll T_{\text{obs}}$ )

$$h(t) \sim (t - t_0)\theta(t - t_0)$$

static events ( $\tau \gg T_{\text{obs}}$ )

$$h(t) \sim t^3$$

at low masses the single event is not detectable but the “noise” induce by the entire population could. Still an open problem!



# CONCLUSIONS

Are we close to the detection of a GWB (?)

Astrophysical or new-physics?

PTAs can be a powerful discovery tool (see Phase Transitions, Cosmic Strings, DM substructures ...)

PTAs can set stringent constraint on new-physics model

work needs to be done to fully exploit PTAs as anomalies detectors