

Cavity magnonics in strong coupling regime – from magnon-polariton hybrid states to perspectives for quantum sensing

Angelo Leo

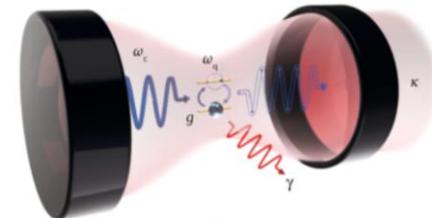
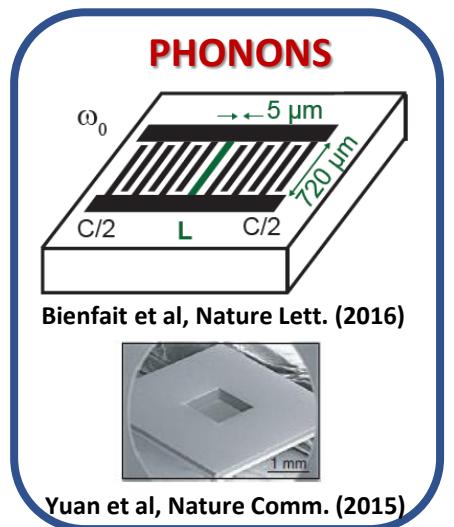
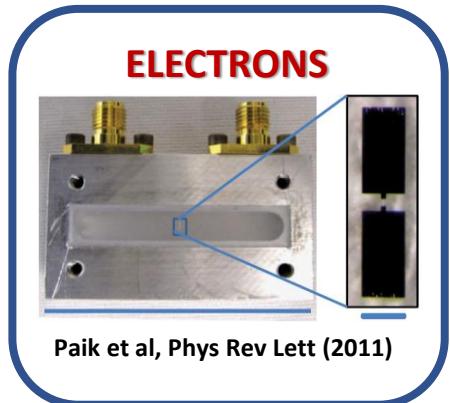
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Istituto Nazionale di Fisica Nucleare, Lecce, Italy



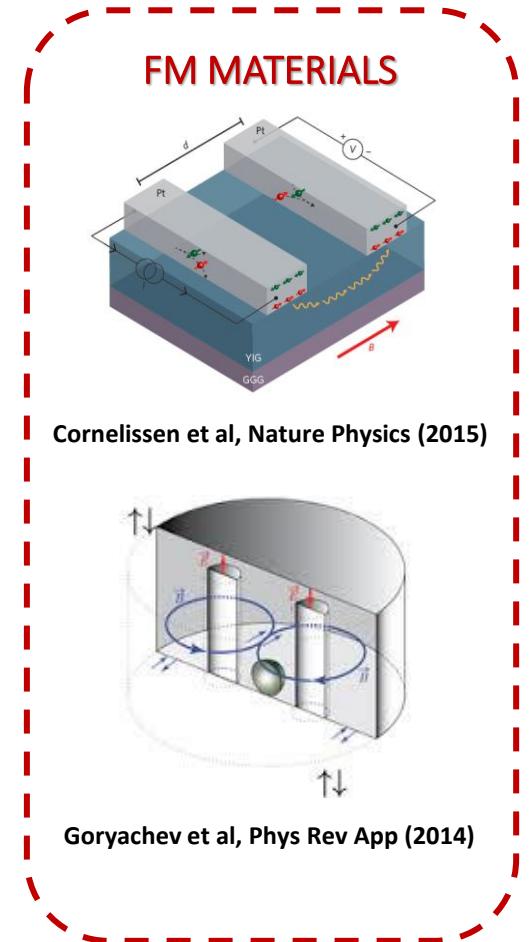
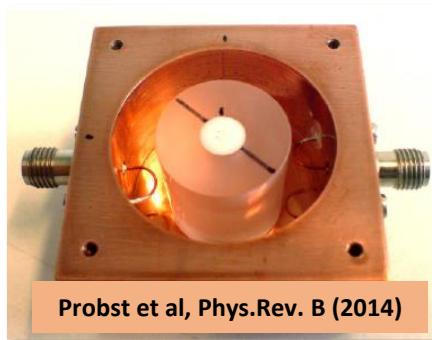
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Cavity QED and hybrid systems for QC and QS



3D MW Cavities

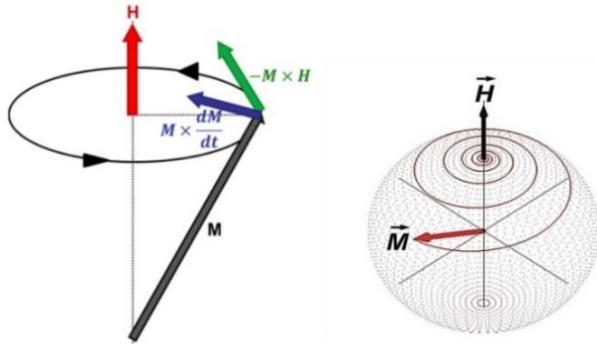


Classical approach

Dynamics of Magnetization

Landau – Lifshitz – Gilbert eq.

$$\frac{d\vec{M}}{dt} = \gamma_e \mu_0 (\vec{M} \times \vec{H}_{eff}) + \frac{\alpha}{M_S} \vec{M} \times \frac{d\vec{M}}{dt}$$



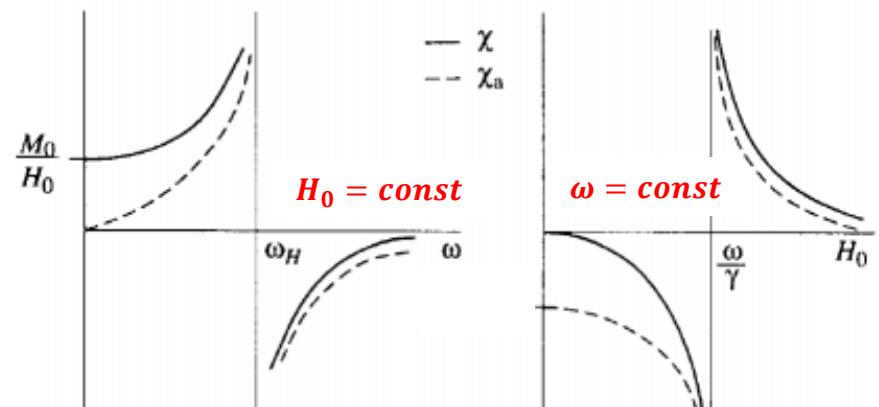
Magnetic susceptibility

$$\chi_a = \frac{\gamma_e M_0 \omega}{\omega_H^2 - \omega^2}$$

$$\vec{\chi} = \begin{pmatrix} \chi & i\chi_a & 0 \\ -i\chi_a & \chi & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\chi = \frac{\gamma_e M_0 \omega_H}{\omega_H^2 - \omega^2}$$

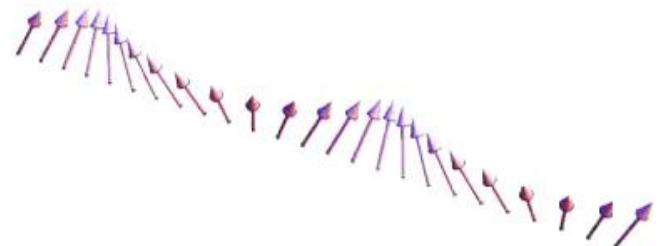
when $\omega = \omega_H = \gamma_e H_0$ **FMR**



Spin waves and magnons – overview

$$\hat{\mathcal{H}} = -2 \frac{J}{\hbar^2} \sum_{j,\delta} \hat{\mathbf{S}}_j \cdot \hat{\mathbf{S}}_{j+\delta} - \frac{g\mu_B H_0}{\hbar} \sum_j \hat{S}_{jz}$$

Exchange among
spins of NN sites with ext. field



Holstein-Primakoff transformations

$$S_j^+ |n_j\rangle = \hbar\sqrt{2s} \left(1 - \frac{(n_j - 1)}{2s}\right)^{1/2} \sqrt{n_j} |n_j - 1\rangle$$

$$S_j^+ = \hbar\sqrt{2s} \left(1 - \frac{m_j^+ m_j^-}{2s}\right)^{1/2} m_j^-$$

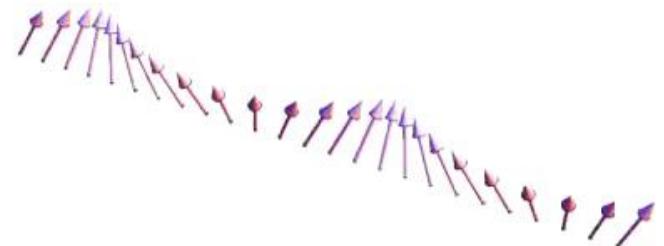
$$S_j^- |n_j\rangle = \hbar\sqrt{2s} \left(1 - \frac{(n_j - 1)}{2s}\right)^{1/2} \sqrt{n_j + 1} |n_j + 1\rangle$$

$$S_j^- = \hbar\sqrt{2s} \left(1 - \frac{m_j^+ m_j^-}{2s}\right)^{1/2} m_j^+$$

Spin waves and magnons – overview

$$\hat{\mathcal{H}} = -2 \frac{J}{\hbar^2} \sum_{j,\delta} \hat{\mathbf{S}}_j \cdot \hat{\mathbf{S}}_{j+\delta} - \frac{g\mu_B H_0}{\hbar} \sum_j \hat{S}_{jz}$$

Exchange among
spins of NN sites with ext. field



At low temperature...

$$m_k^+ = N^{-1/2} \sum_j e^{ik \cdot r_j} m_j^+$$

$$m_k^- = N^{-1/2} \sum_j e^{-ik \cdot r_j} m_j^-$$

$$\hat{\mathcal{H}}_0 = \sum_k \hat{n}_k \hbar \omega_k \quad \hbar \omega_k = 4J s Z (1 - \gamma_k) + g \mu_B H_0$$

If the j^{th} site has Z nearest neighbours

$$\hat{\mathcal{H}} = -2JNZs^2 - g\mu_B H_0 Ns + \hat{\mathcal{H}}_0$$

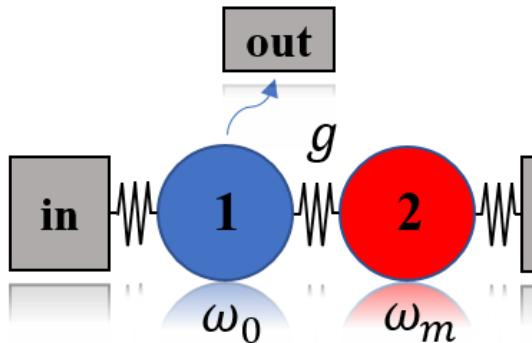
$$\hat{\mathcal{H}}_0 = \sum_k (4JaZ(1 - \gamma_k) + g\mu_B H_0) m_k^+ m_k^-$$

n of magnons @ k

$$\hat{n}_k = m_k^+ m_k^-$$

$$\epsilon_k = (n_k + 1/2) \hbar \omega_k$$

Analogies with quantum systems – cavity magnon polariton systems



Confined EM field

$$\mathcal{H}_c = \hbar\omega_0 \left(a^\dagger a + \frac{1}{2} \right)$$

Magnons

$$\mathcal{H}_m = \hbar\omega_m \left(m^\dagger m + \frac{1}{2} \right)$$

only with resonant interactions:

$$\mathcal{H} = \mathcal{H}_c + \mathcal{H}_m + \mathcal{H}_{int}$$

$$\frac{\mathcal{H}}{\hbar} = \omega_0 a^\dagger a + \omega_m m^\dagger m + \underline{g(a^\dagger m + m^\dagger a)}$$

Coupling strength of each spin:

$$\frac{g_{0,i}}{2\pi} = \frac{\left(\eta \gamma_e \sqrt{\frac{(\mu_0 \hbar \omega_0)}{V_c}} \right)}{4\pi}$$

As a spin ensemble:

$$g_i = g_{0,i} \sqrt{Ns}$$

$$g_i \propto \sqrt{\omega_{eff}} = \sqrt{\omega_0 V_m / V_a}$$

Overlap among subsystems:

$$\eta = \int_{V_m}^0 \frac{\vec{h} \cdot \vec{M}}{|\vec{h}_{max}| |\vec{M}_{max}|} dV \leq 1$$



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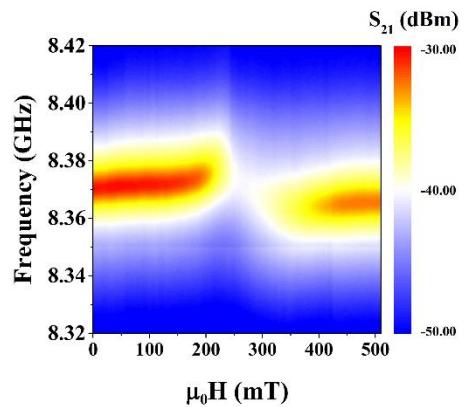
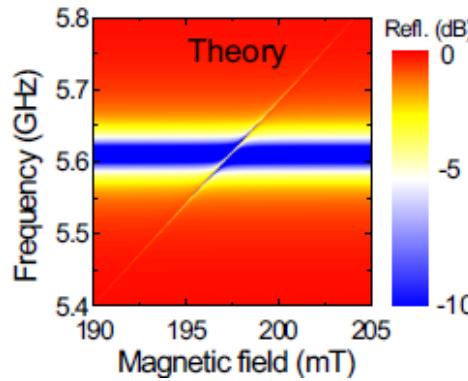


INFN

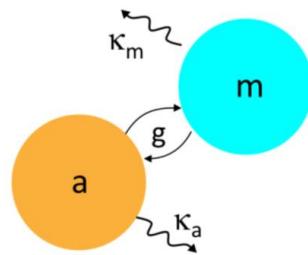


Analogies with quantum systems – cavity magnon polariton systems

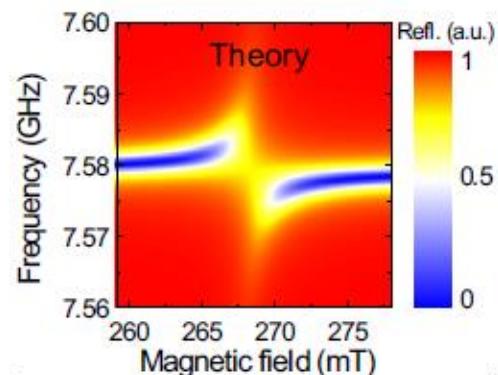
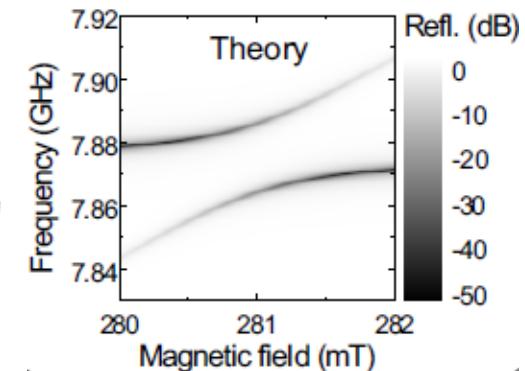
Zhang et al, Phys Rev Lett (2014)



Mn_2FeO_4 NPs in a low Q cavity (Unisalento)

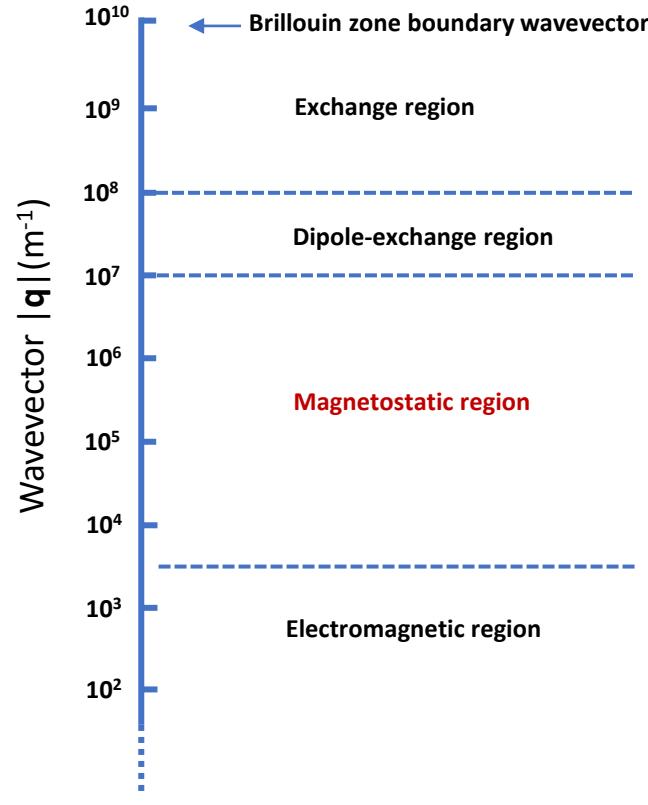


Zhang et al, Phys Rev Lett (2014)



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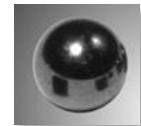
Magnetostatic modes of a magnetic sphere



Kittel mode ($m = n = 1$)
 $f = \gamma H_0$

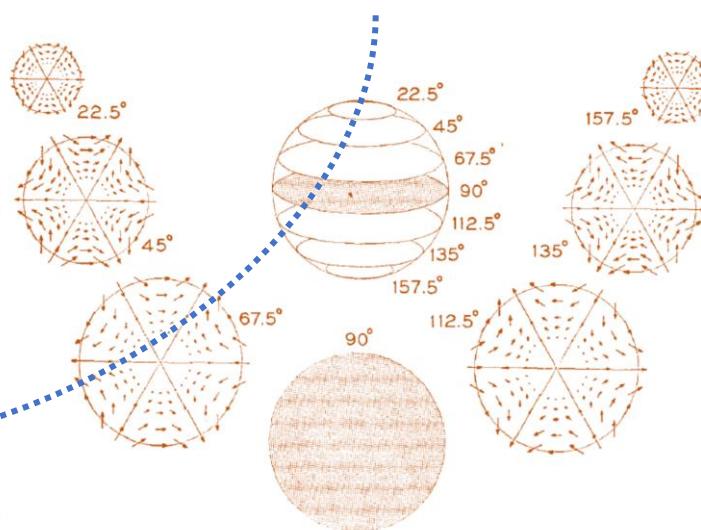
Magnetic dipolar interaction

$$\frac{d}{dt} \vec{M} = -\gamma \mu_0 \left(\vec{M} \times \vec{H}_{\text{eff}} - \eta \vec{M} \times \frac{d}{dt} \vec{M} \right)$$



Indexes of MSMs

$$n + 1 + \xi_0 \frac{P_n^{m'}(f, H_0)}{P_n^m(f, H_0)} \pm m \frac{\gamma f M_S}{\gamma^2 H_i^2 - f^2} = 0$$

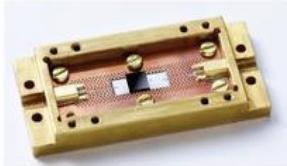


Walker, J.Appl.Phys. (1958)

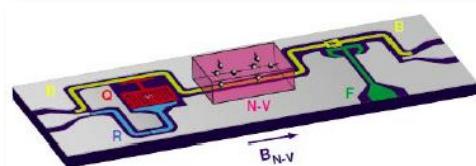
Yttrium Iron Garnet for cavity magnonics

Paramagnetic spin ensembles

NV- centers in diamond

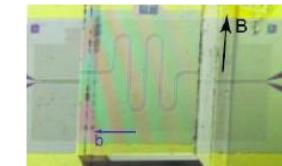


Putz et al., Nature Physics (2014)



Phys Rev Lett 107, 220501 (2011)

Rare-earth doped crystal

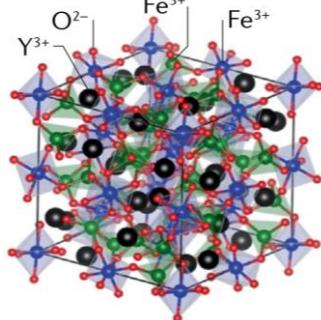
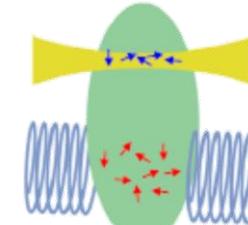


Bushev et al., Phys Rev B (2011)

Low spin density: 10^{12} - 10^{18} cm^{-3}



Optical Light
Microwave



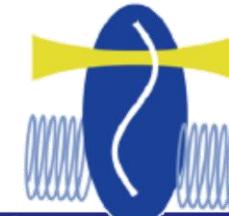
- High Curie Temperature
- Insulator
- Low magnetic damping
- Absorption coefficient
- Large Verdet constant
- large size of the lattice cell

$T_c = 560 \text{ K}$
 $BG = 2.68 \text{ eV}$
 $\alpha_0 \sim 10^{-5}$
 $\sim 10/\text{cm} @ 0.8 \text{ THz}$
 $\sim 100 \text{ }^\circ/\text{mm/T}$
 1.2376 nm

High net spin density $2.1 \times 10^{22} \text{ cm}^{-3}$



Optical Light
Microwave



Standard use

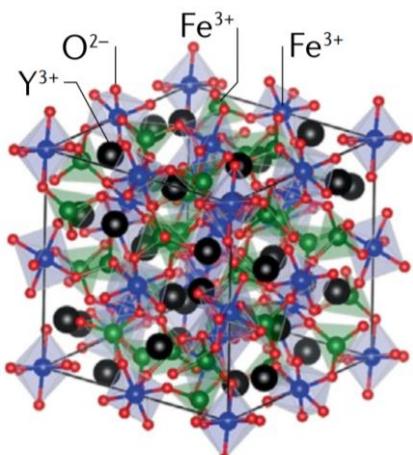
- MW oscillator



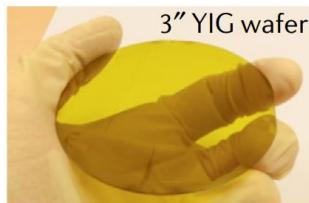
- Optical isolator



Yttrium Iron Garnet for coherent cavity magnonics

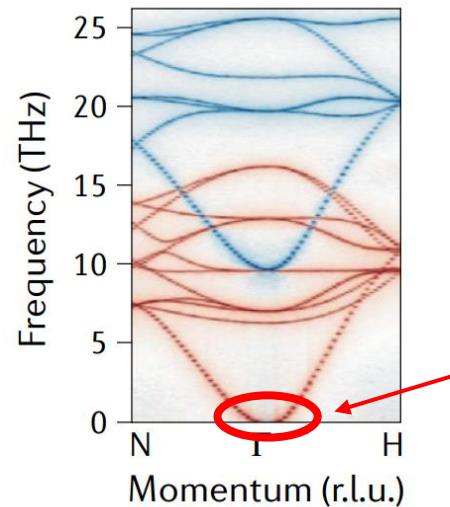


Fabrication



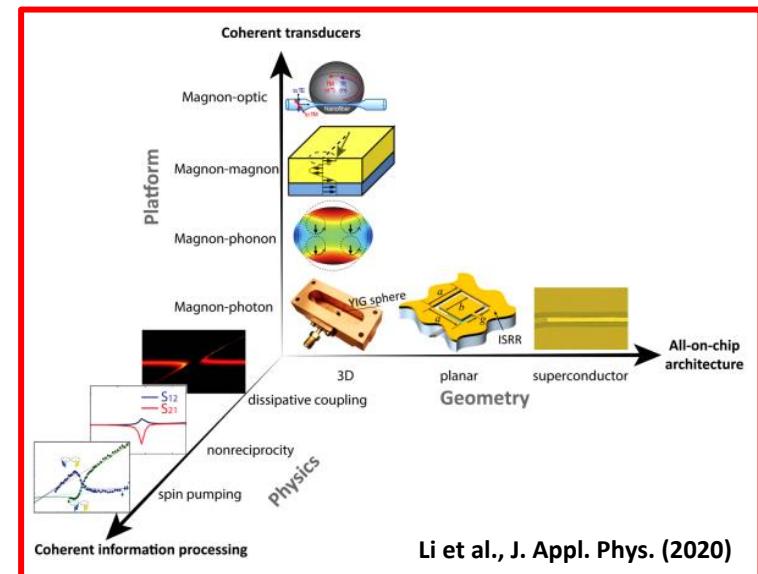
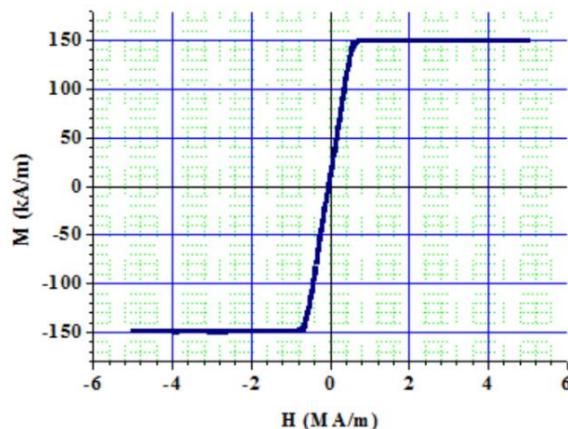
- LPE
- PLD
- Magnetron sputtering

P. Pirro et al., Nature Reviews. (2021)



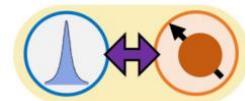
There are 20 distinct exchange magnon branches but, for cavity (quantum) magnonics, the interest is focused on **Kittel mode**

Vibrating sample magnetometry @ UniSalento

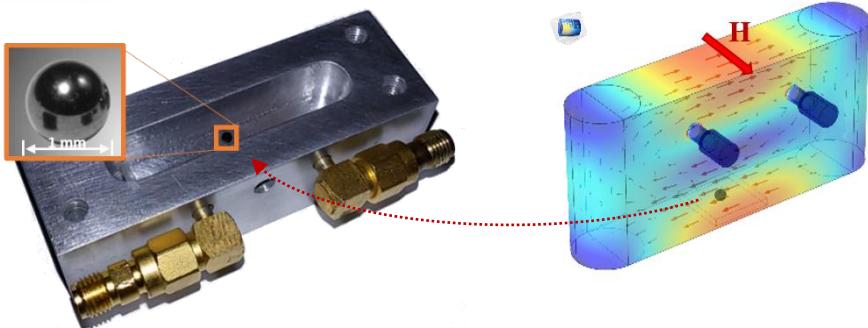


Li et al., J. Appl. Phys. (2020)

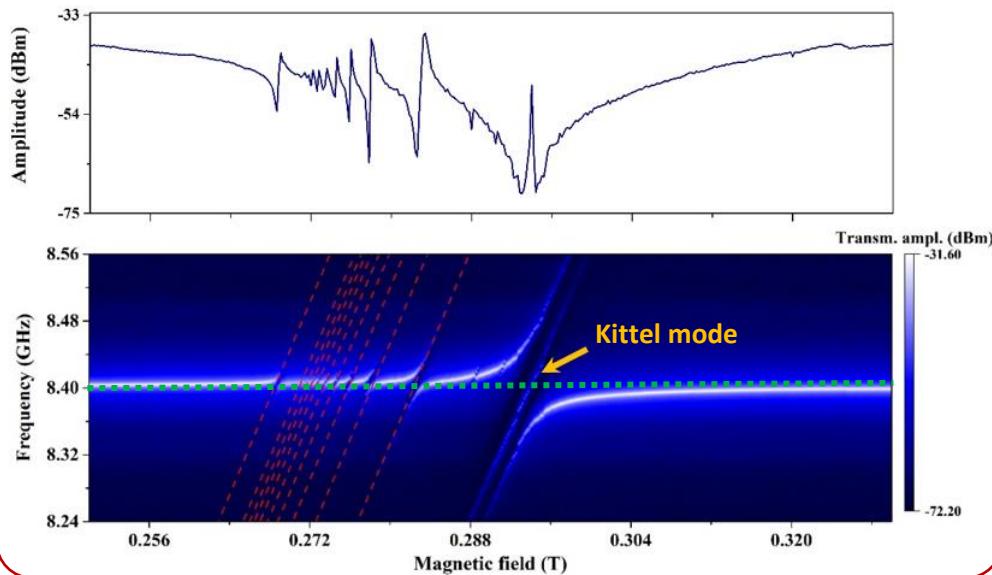
Strong coupling regime



setup



Spectroscopic investigation



Input-output formalism:

$$T(\omega) = \frac{\kappa_c}{i(\omega - \omega_c) - \kappa_c + \sum_m \frac{1}{2} \kappa_m + i(\omega - \omega_m)} |g_m|^2$$

Cavity mode at:

$$\omega_c$$

Total cavity decay rate:

$$\kappa_c = \frac{1}{2} (2\kappa_{i,o} + \kappa_{int})$$

Magnon mode damping rate:

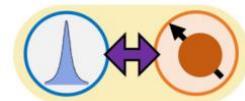
$$\kappa_m \quad (m = \text{FMR, MSMs})$$

Magnon mode at
Coupling strength

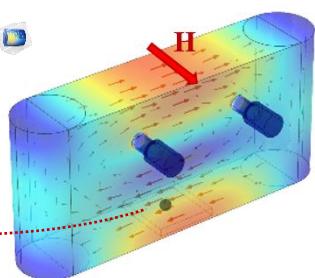
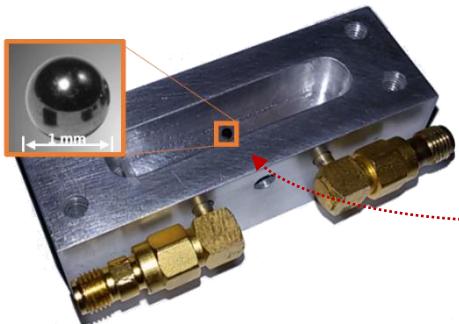
$$\omega_m \quad (m = \text{FMR, MSMs})$$

$$g_m \quad (m = \text{FMR, MSMs})$$

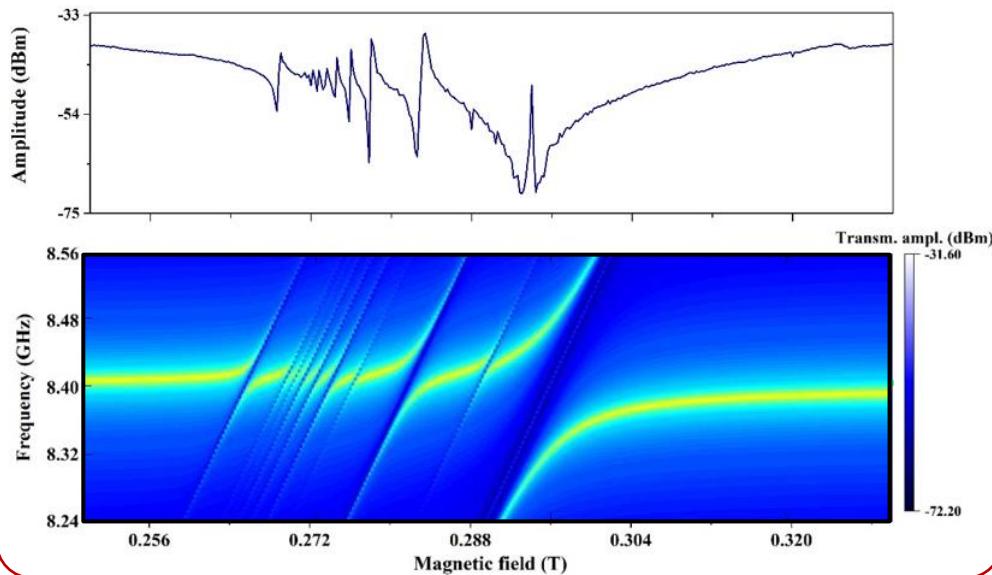
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Magnon mode at

$$\omega_m \quad (m = \text{FMR, MSMs})$$

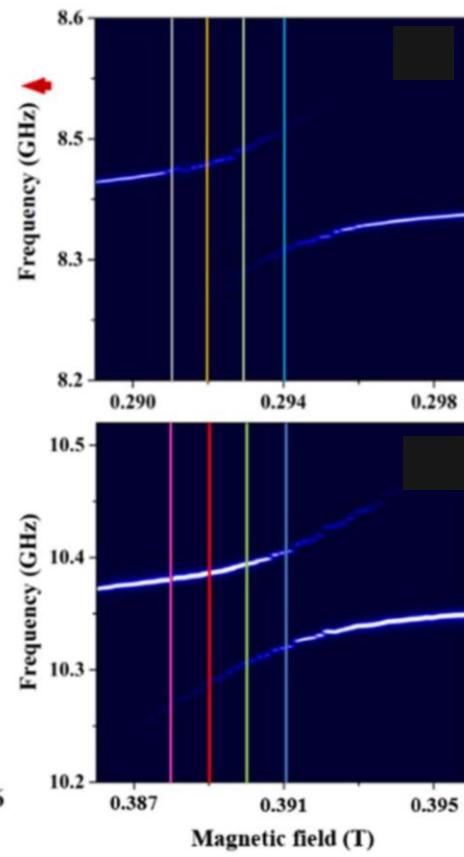
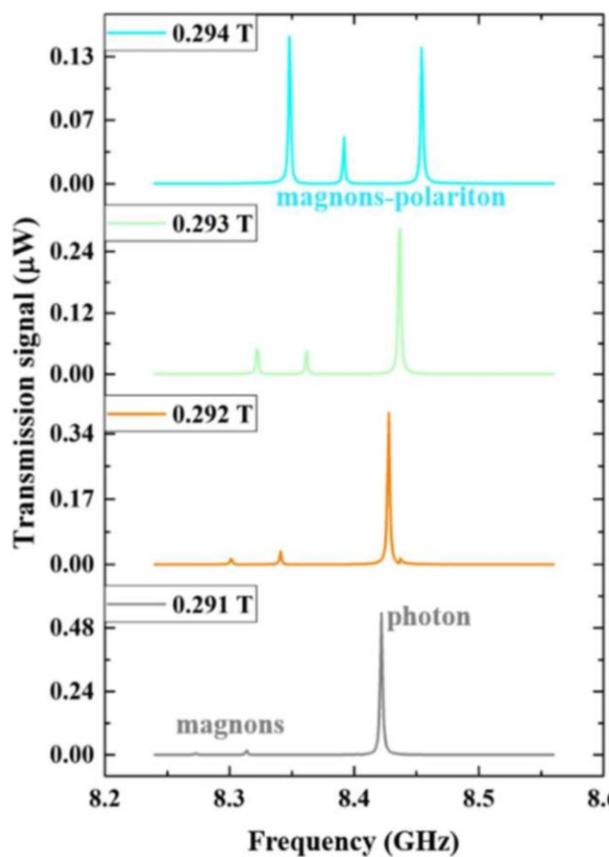
Coupling strength

$$g_m \quad (m = \text{FMR, MSMs})$$

Strong coupling regime

Transmitted signal

$$T(\omega) = \frac{\kappa_c}{i(\omega - \omega_c) - \kappa_c + \sum_m \frac{|g_m|^2}{-\frac{1}{2}\kappa_m + i(\omega - \omega_m)}}$$



Cavity mode at:

ω_c

Total cavity decay rate:

$$\kappa_c = \frac{1}{2}(2\kappa_{i,o} + \kappa_{int})$$

Magnon mode damping rate:

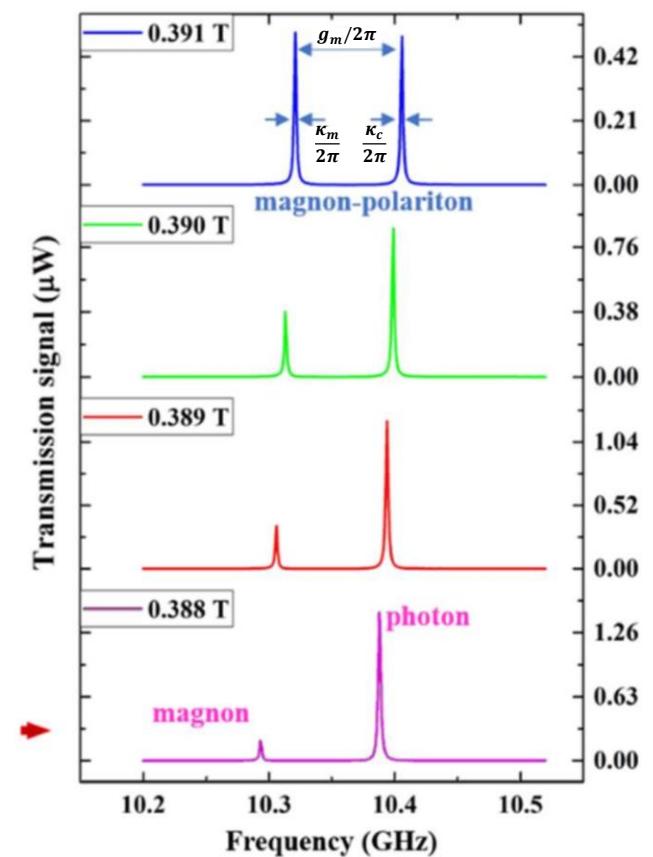
$$\kappa_m (m = \text{FMR, MSMs})$$

Magnon mode at:

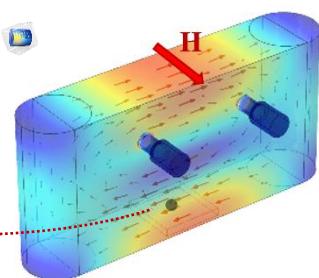
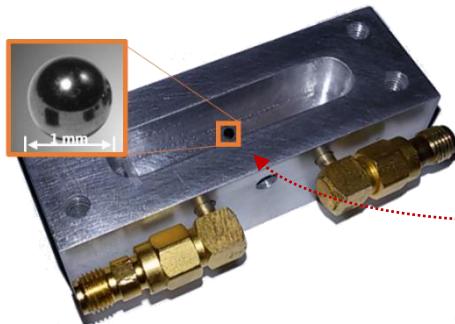
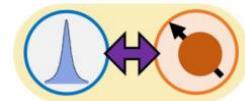
$$\omega_m (m = \text{FMR, MSMs})$$

Coupling strength

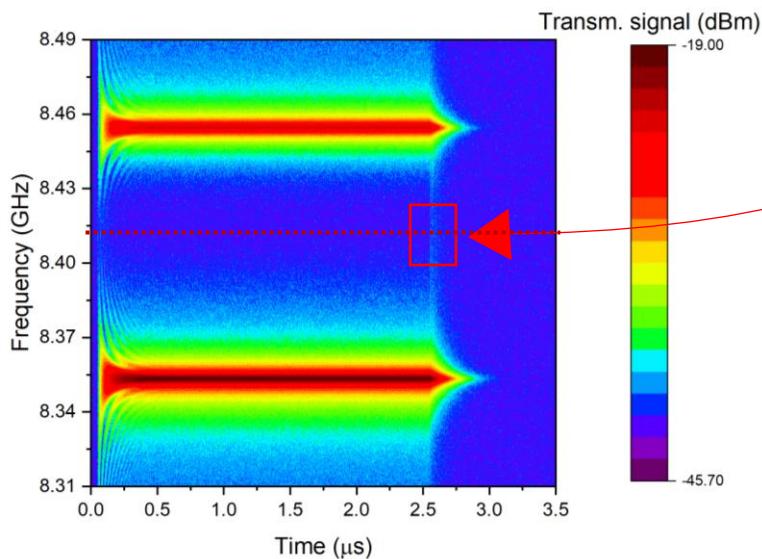
$$g_m (m = \text{FMR, MSMs})$$



Rabi oscillations



Time-resolved measurements

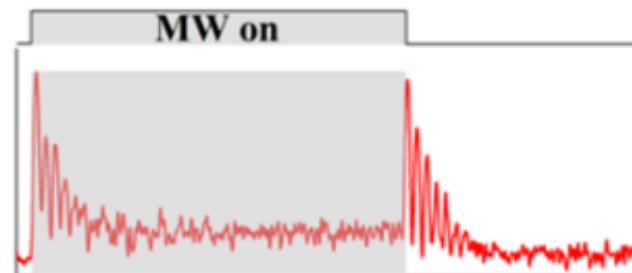


Rabi oscillations:

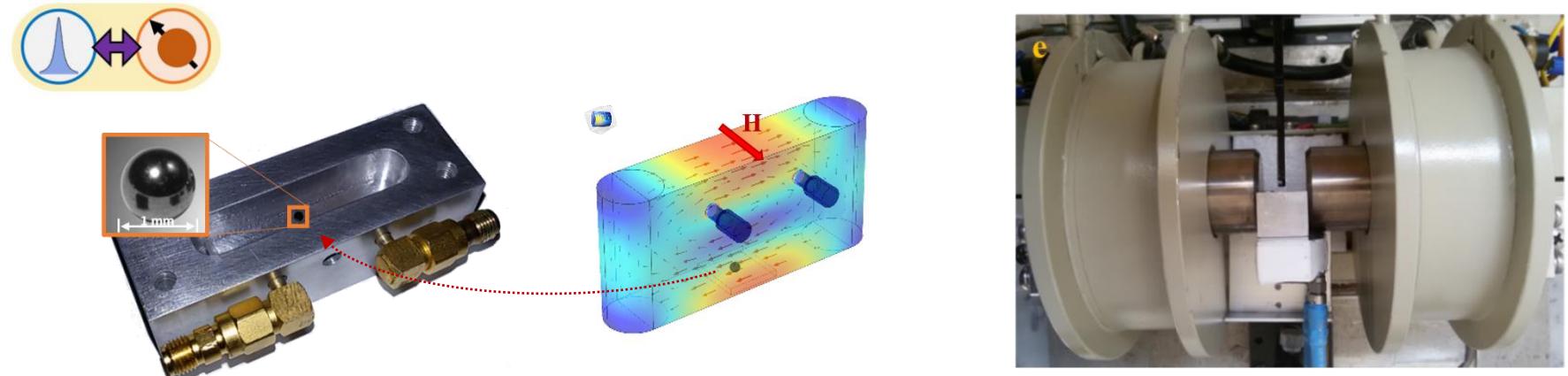
$$P\left(\frac{\omega_i}{2\pi}\right) \propto A_i \cdot \log \left[\sin^2\left(\frac{g_i}{2\pi}t\right) \cdot \exp\left(-\frac{2\kappa_c - \kappa_{int}}{2\pi}t\right) \right]$$

Coupling among resonators

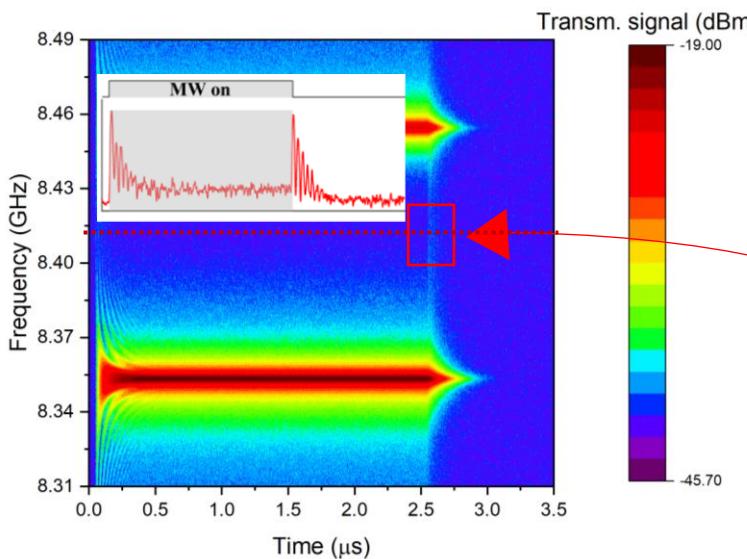
MW cavity relaxation



Rabi oscillations



Time-resolved measurements



Volterra equation

$$a(t) = \int_0^t d\tau \mathcal{K}(t-\tau) a(\tau) + \mathcal{F}(t)$$

$$\mathcal{F}(t) = - \int_0^t d\tau \cdot \eta(\tau) e^{-\kappa(t-\tau)}$$

$$\mathcal{K}(t-\tau) = g_m^2 \int d\omega \rho(\omega) \frac{\left[e^{-i(\omega - \omega_c - i(\kappa_m - \kappa_c))(t-\tau)} - 1 \right]}{-i(\omega - \omega_c - i(\kappa_m - \kappa_c))} \cdot e^{-\kappa(t-\tau)}$$

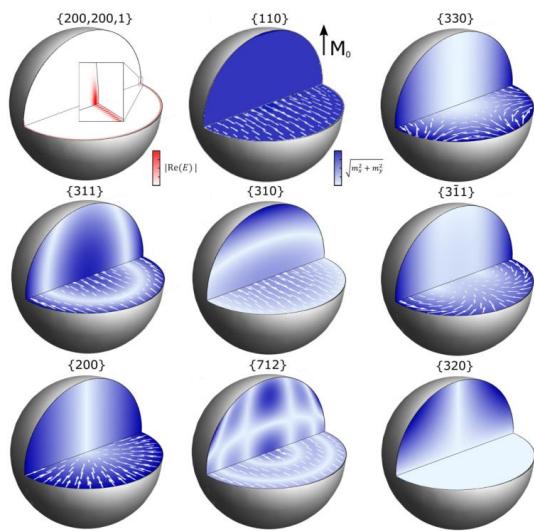
Rabi oscillations:

$$P\left(\frac{\omega_i}{2\pi}\right) \propto A_i \cdot \log \left[\sin^2\left(\frac{g_i}{2\pi}t\right) \cdot \exp\left(-\frac{2\kappa_c - \kappa_{int}}{2\pi}t\right) \right]$$

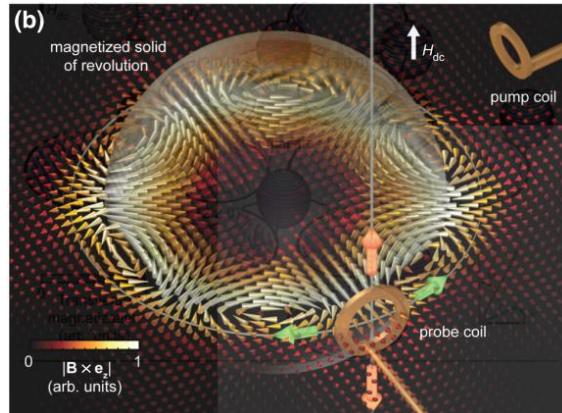
Coupling among resonators

MW cavity relaxation

Effect of reservoir on MSMs visualization



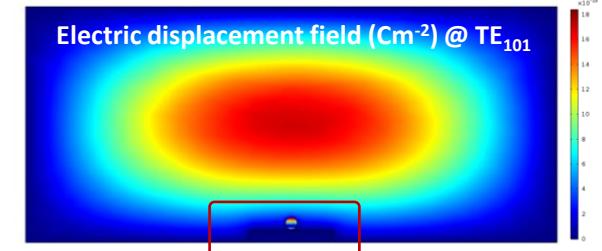
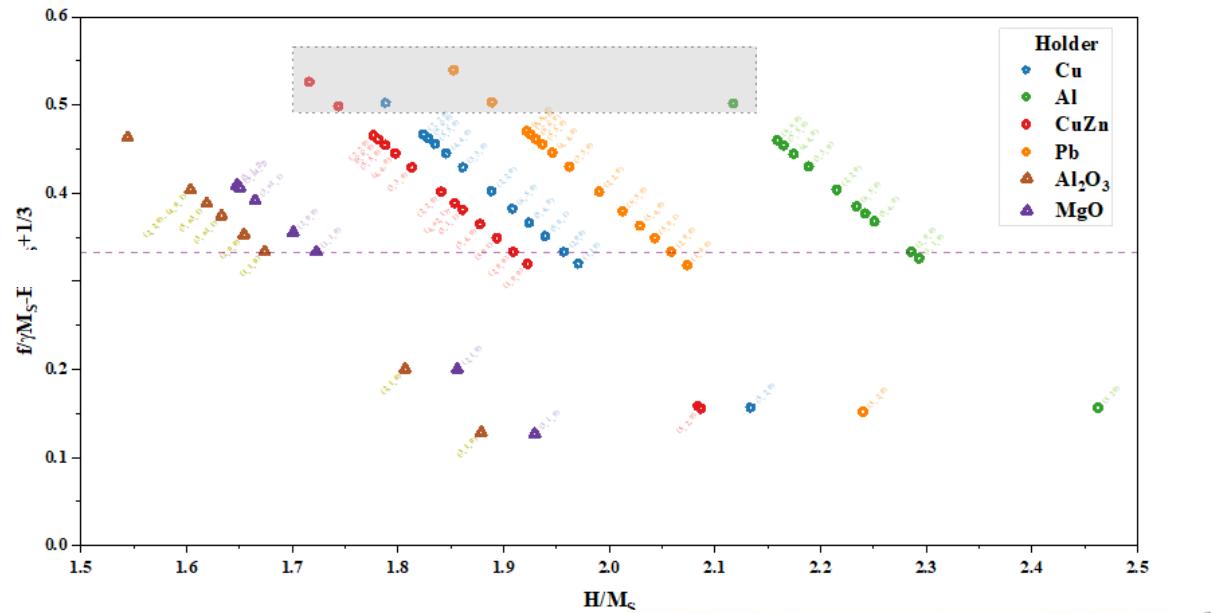
J.A. Haigh Phys Rev B (2018)



A. Gloppe, Phys Rev Appl (2018)

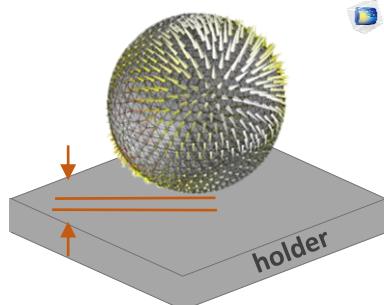
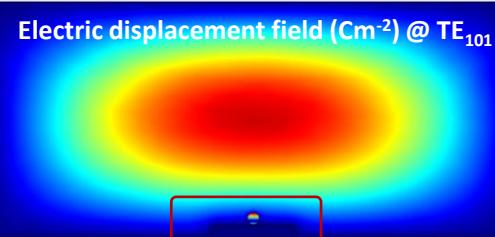


$$\text{MSMs identification based on: } n + 1 + \xi_0 \frac{P_n^{m'}(f, H_0)}{P_n^m(f, H_0)} \pm m \frac{\gamma f M_S}{\gamma^2 H_i^2 - f^2} = 0$$



Effect of reservoir on MSMs visualization

RF electric field @ TE₁₀₁

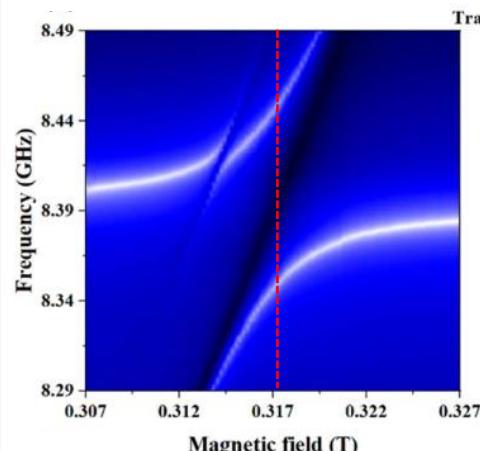


Surf. electric currents on YIG sphere at cavity working frequency

multimode hybridization mediated by cavity photons

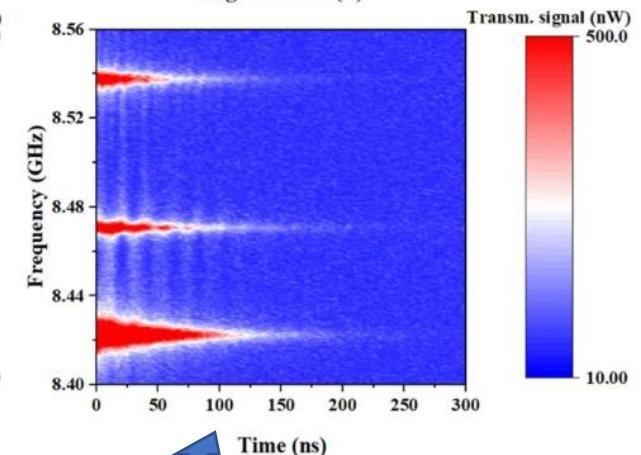
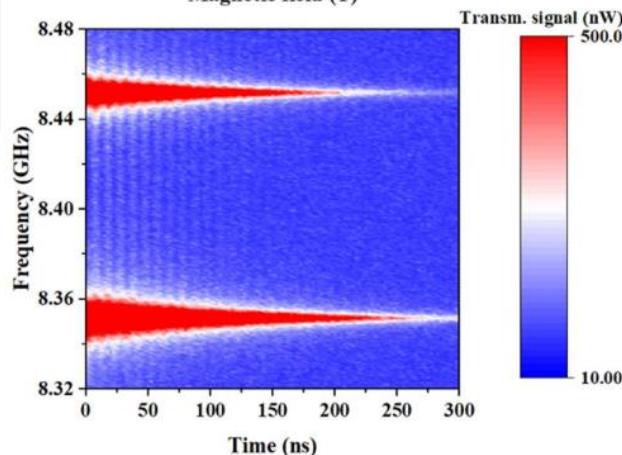
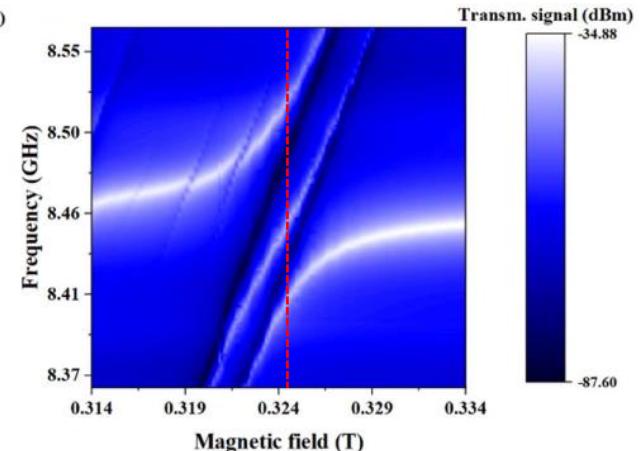


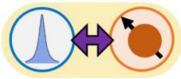
MgO



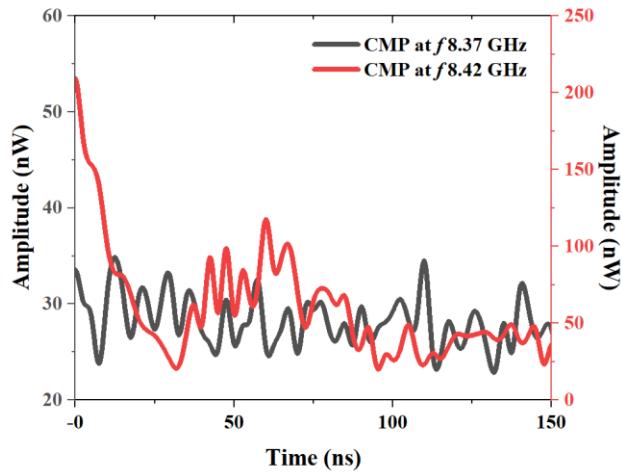
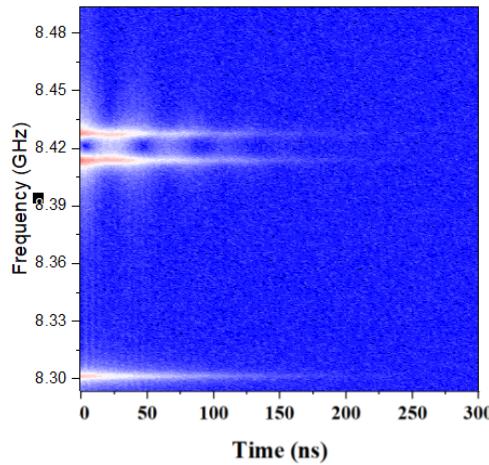
Holder electrical conductivity

Cu



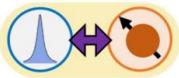


Nonreciprocity

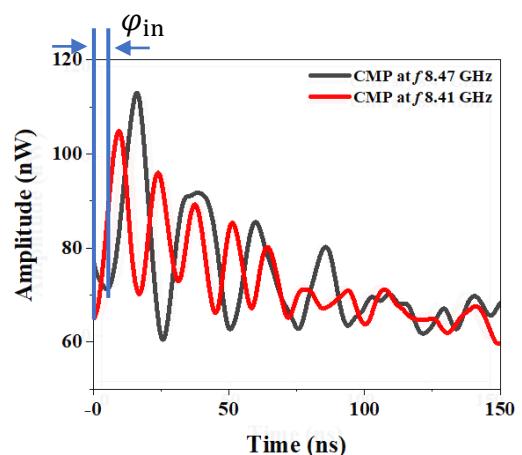
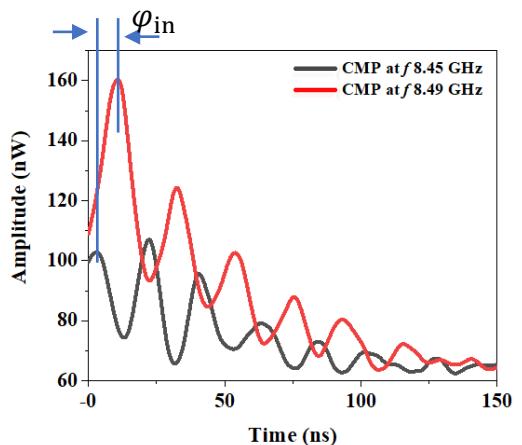
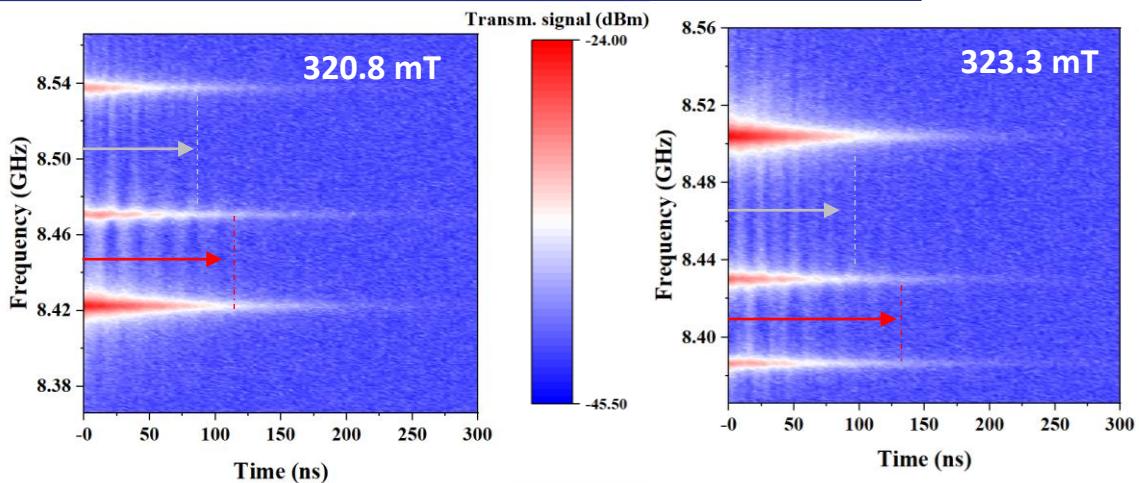
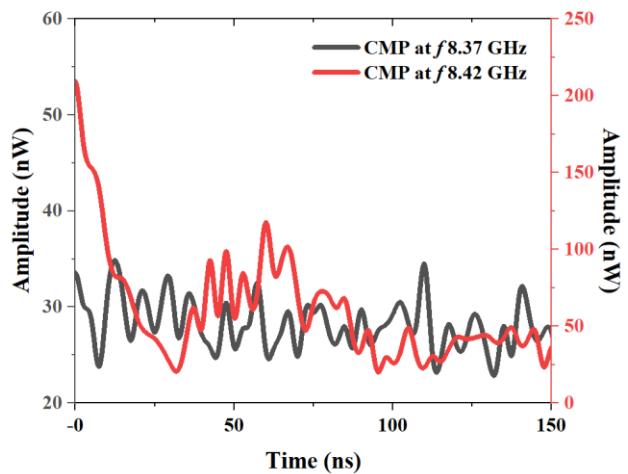
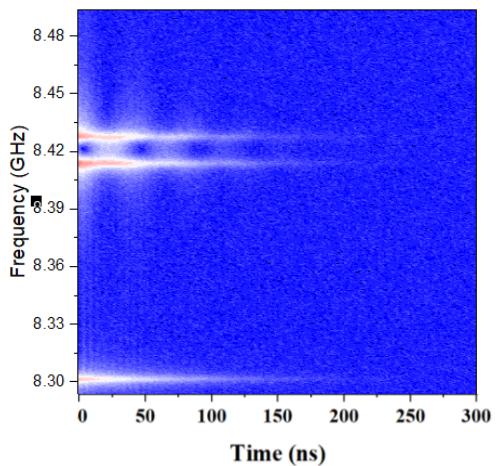
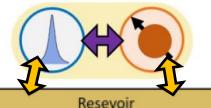


multimode hybridization mediated by cavity photons

$$\frac{1}{4\pi^2} \begin{pmatrix} \omega^2 - \omega_{FMR}^2 + i2\alpha\omega_{FMR}\omega & g_1^2\omega_{FMR}^2 & 0 \\ g_1^2\omega_{FMR}^2 & \omega^2 - \omega_c^2 + i2\alpha\omega_c\omega & g_2^2\omega_{(2,0,0)}^2 \\ 0 & g_2^2\omega_{(2,0,0)}^2 & \omega^2 - \omega_{(2,0,0)}^2 + i2\alpha\omega_{(2,0,0)}\omega \end{pmatrix} \begin{pmatrix} m_{FMR} \\ h_\omega \\ m_{(2,0,0)} \end{pmatrix} = \omega^2 \begin{pmatrix} 0 \\ \Gamma \\ 0 \end{pmatrix} h_0$$



Nonreciprocity



multimode hybridization mediated by cavity photons

$$\frac{1}{4\pi^2} \begin{pmatrix} \omega^2 - \omega_{FMR}^2 + i2\alpha\omega_{FMR}\omega & g_1^2\omega_{FMR}^2 \\ g_1^2\omega_{FMR}^2 & \omega^2 - \omega_c^2 + i2\alpha\omega_c\omega \\ 0 & g_2^2\omega_{(2,0,0)}^2 \end{pmatrix} \begin{pmatrix} m_{FMR} \\ h_\omega \\ m_{(2,0,0)} \end{pmatrix} = \omega^2 \begin{pmatrix} 0 \\ \Gamma \\ 0 \end{pmatrix} h_0$$



“quantum unidirectional processing of information”

C. Zhang et al., PRB (2021)

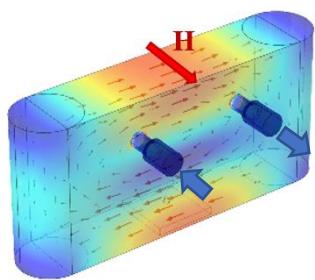
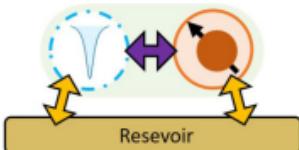


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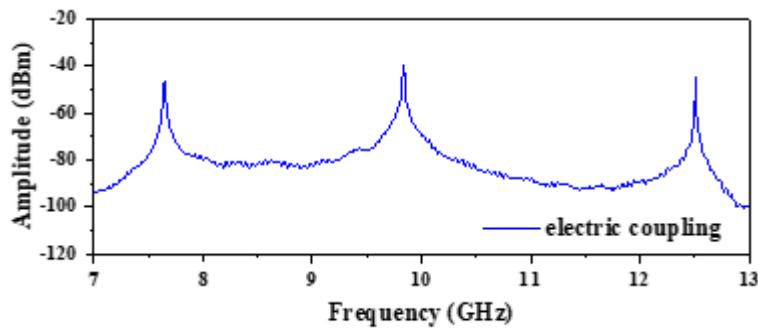
Leo et al. - submitted

Investigation in dissipative regime of CMP hybrid system

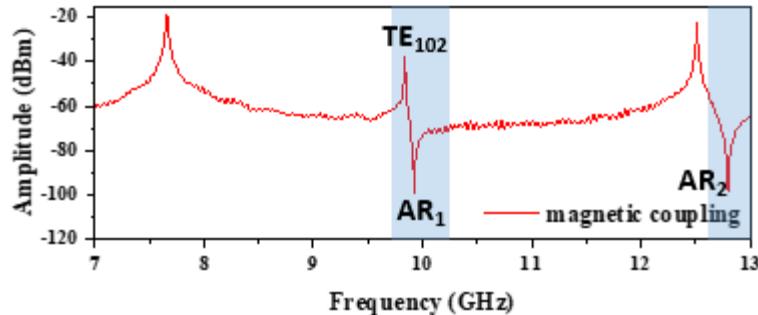


Cavity Lenz effect

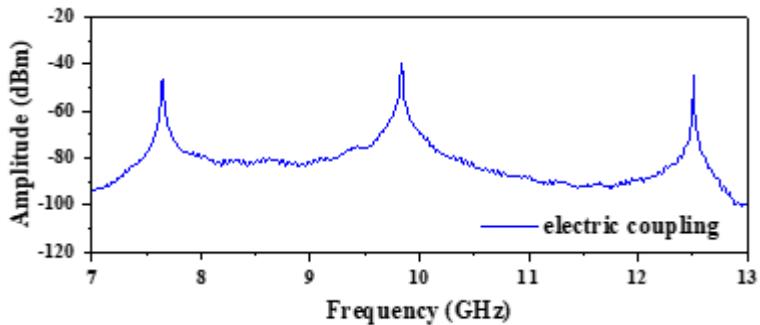
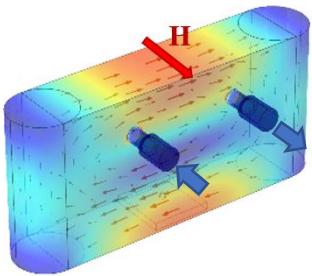
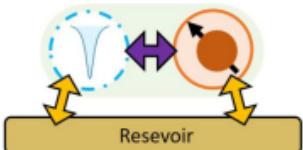
$$\begin{pmatrix} \omega^2 - \omega_c^2 + i2\kappa_a\omega_c\omega & i\omega^2 K_F \\ -i\omega_m(K_A - K_L) & \omega - \omega_r + i\kappa_m\omega \end{pmatrix} \begin{pmatrix} I \\ M \end{pmatrix} = 0$$



Switch of antennas/probes

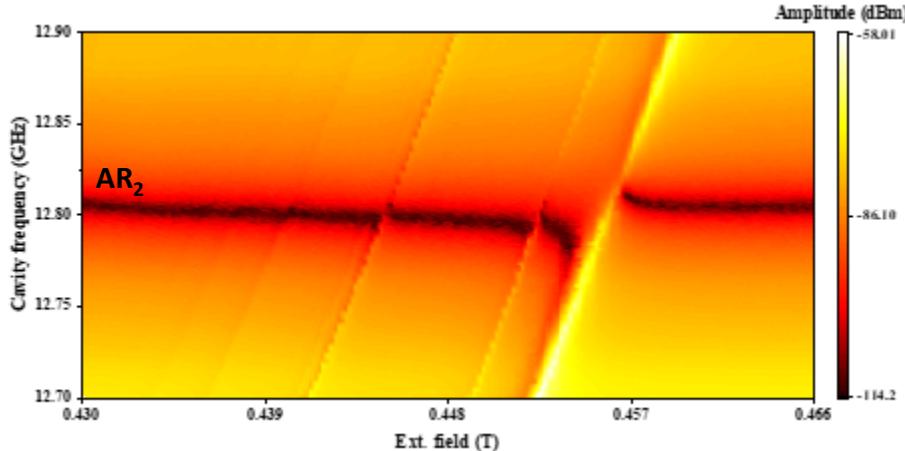
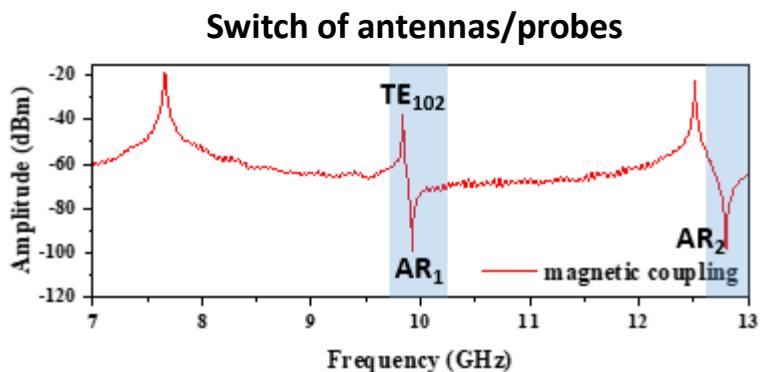
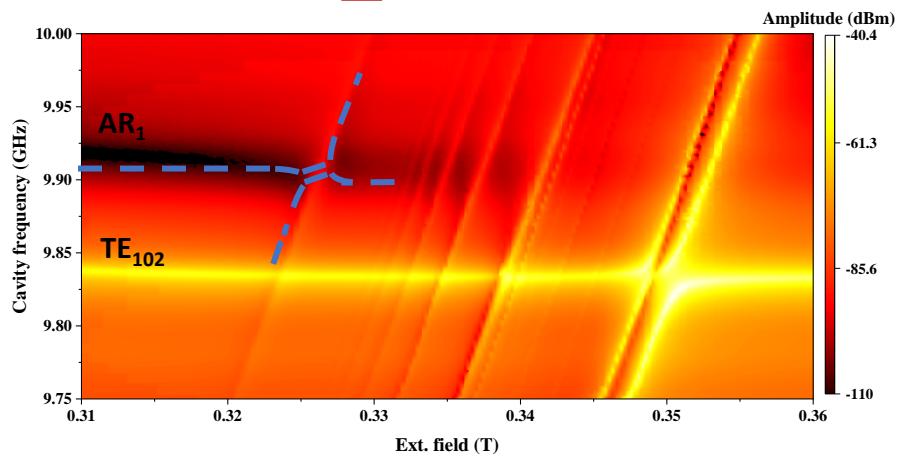


Investigation in dissipative regime of CMP hybrid system

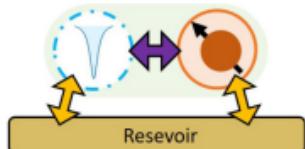


Cavity Lenz effect

$$\begin{pmatrix} \omega^2 - \omega_c^2 + i2\kappa_a\omega_c\omega & i\omega^2 K_F \\ -i\omega_m(K_A - K_L) & \omega - \omega_r + i\kappa_m\omega \end{pmatrix} \begin{pmatrix} I \\ M \end{pmatrix} = 0$$



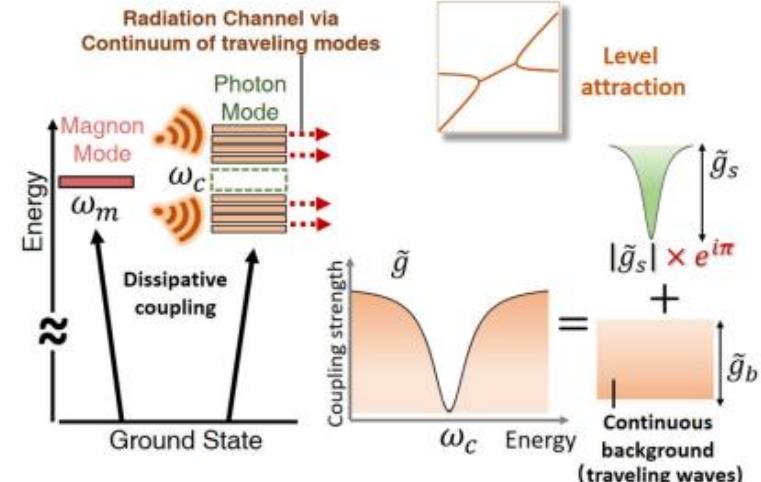
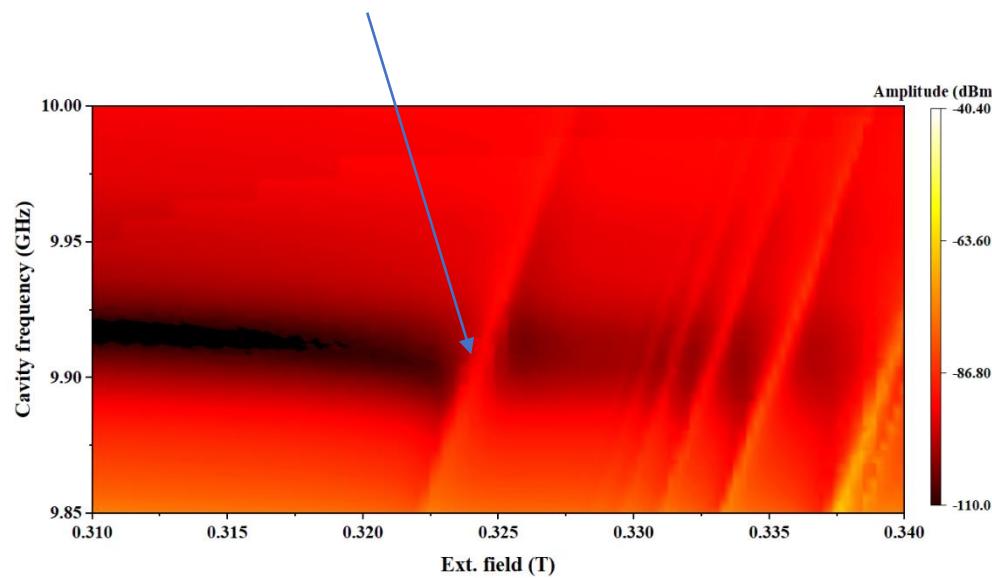
Investigation in dissipative regime of CMP hybrid system



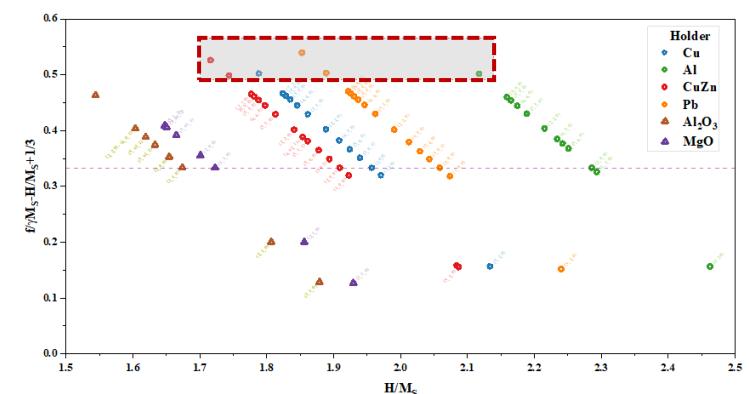
Cavity Lenz effect

$$\begin{pmatrix} \omega^2 - \omega_c^2 + i2\kappa_a\omega_c\omega & i\omega^2 K_F \\ -i\omega_m(K_A - K_L) & \omega - \omega_r + i\kappa_m\omega \end{pmatrix} \begin{pmatrix} I \\ M \end{pmatrix} = 0$$

«Exceptional point»



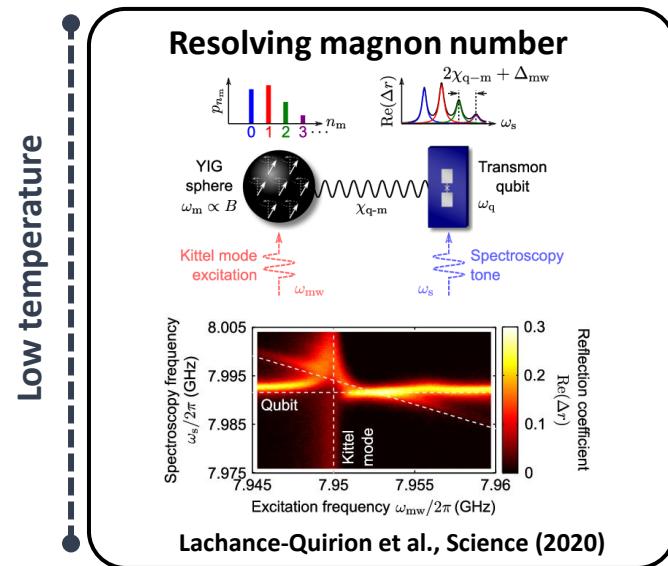
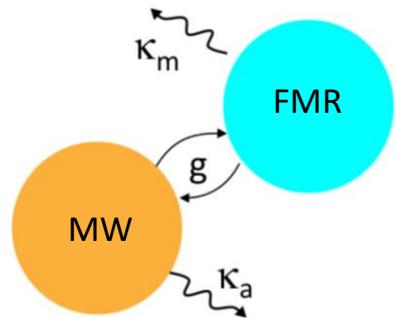
B. Yao et al., PRB (2019)



«Topological energy transfer»

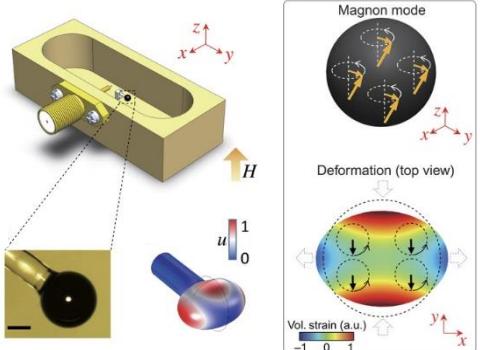
H. Xu et al., Nature 2016

FMR-based information transducer and perspectives for QS

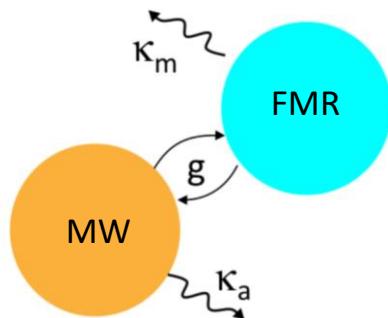


FMR-based information transducer and perspectives for QS

Phononic coupling with magnons

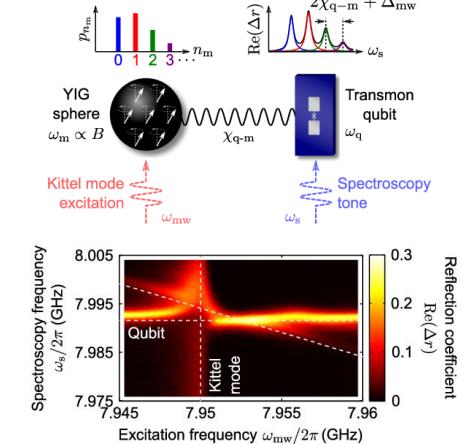


X. Zhang et al., Science Advances (2016)



Low temperature

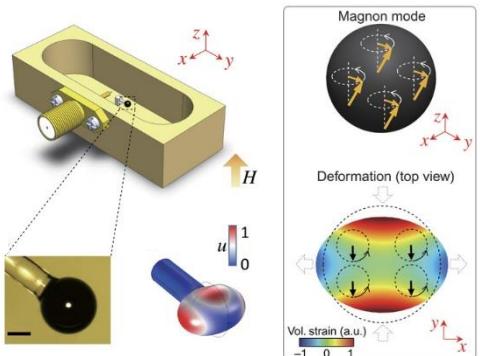
Resolving magnon number



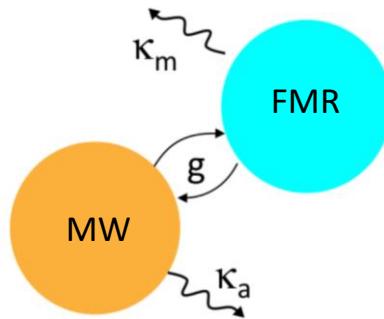
Lachance-Quirion et al., Science (2020)

FMR-based information transducer and perspectives for QS

Phononic coupling with magnons

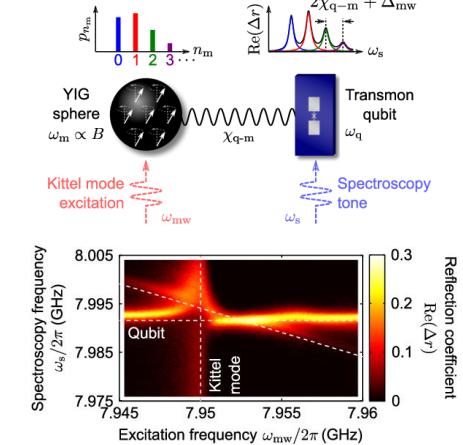


X. Zhang et al., Science Advances (2016)



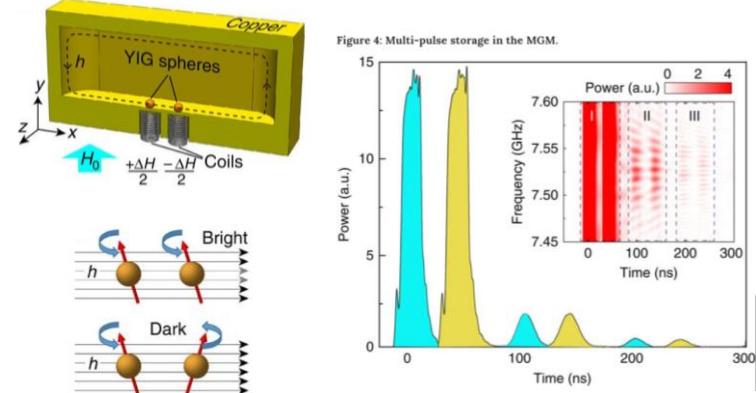
Low temperature

Resolving magnon number



Lachance-Quirion et al., Science (2020)

Magnon dark modes and gradient memories



X. Zhang et al., Nature Comms (2015)



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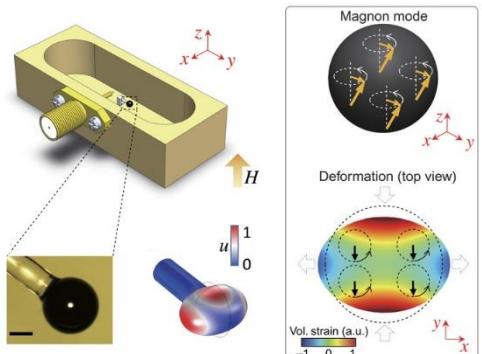
INFN



Three-dimensional Magnonics

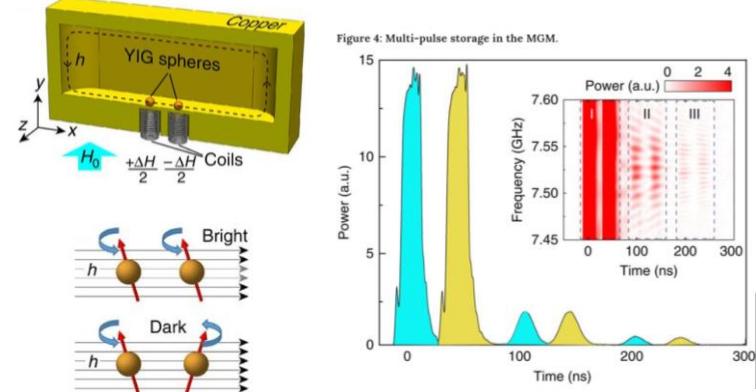
FMR-based information transducer and perspectives for QS

Phononic coupling with magnons

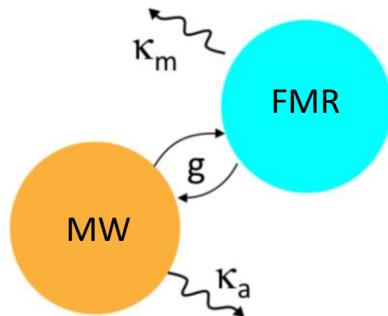


X. Zhang et al., Science Advances (2016)

Magnon dark modes and gradient memories

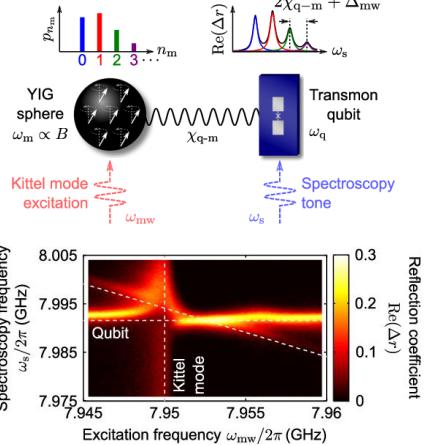


X. Zhang et al., Nature Comms (2015)



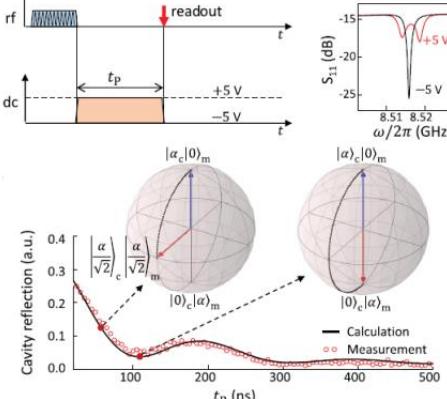
Low temperature

Resolving magnon number



Lachance-Quirion et al., Science (2020)

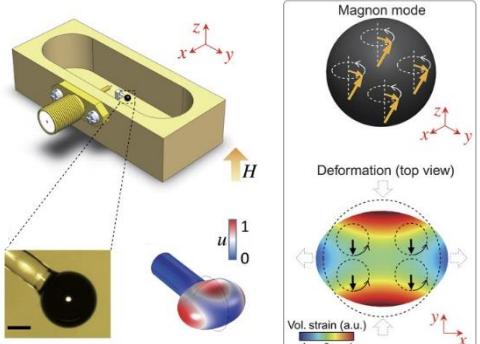
Gate operations



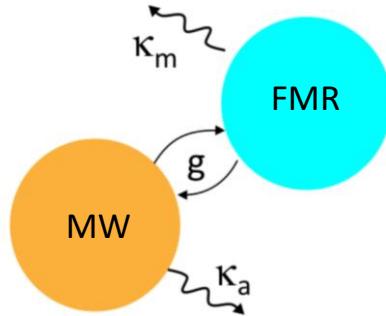
Xu et al, Phys. Rev. Lett. (2021)

FMR-based information transducer and perspectives for QS

Phononic coupling with magnons

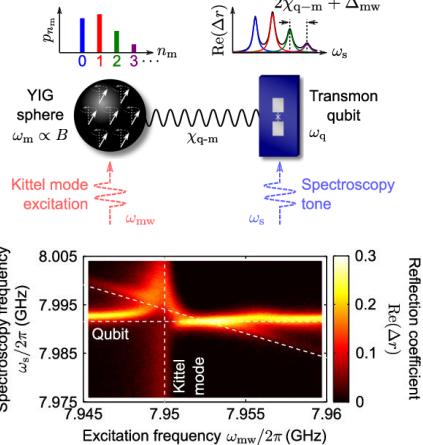


X. Zhang et al., Science Advances (2016)



Low temperature

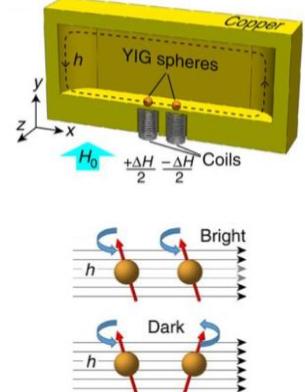
Resolving magnon number



Lachance-Quirion et al., Science (2020)

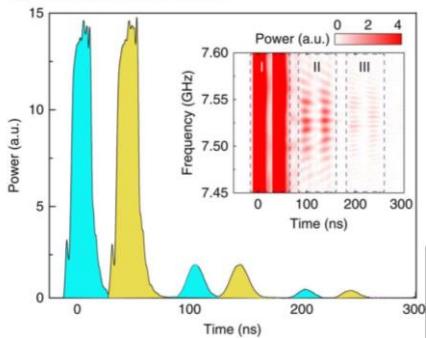
Room temperature

Magnon dark modes and gradient memories

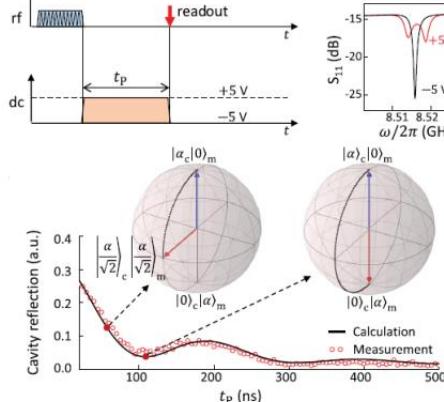


X. Zhang et al., Nature Comms (2015)

Figure 4: Multi-pulse storage in the MGM.

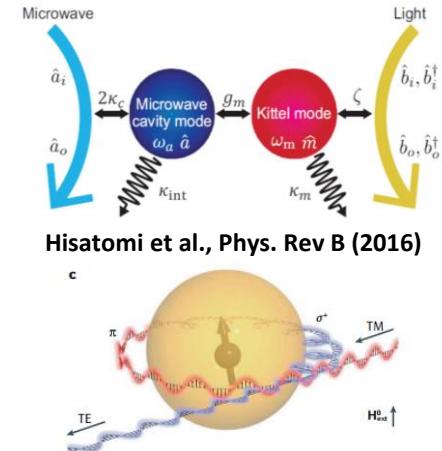


Gate operations



Xu et al, Phys. Rev. Lett. (2021)

MW – opt. conversion



Zhang et al., Phys. Rev Lett (2016)

Three-dimensional Magnonics

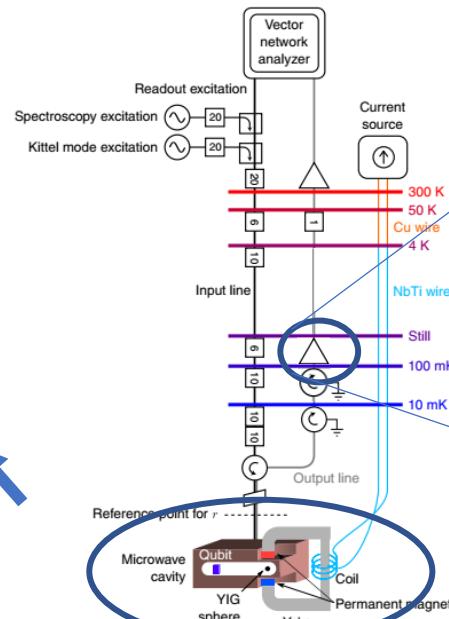
DARTWARS: Detector Array Readout with Traveling Wave AmplifieRS



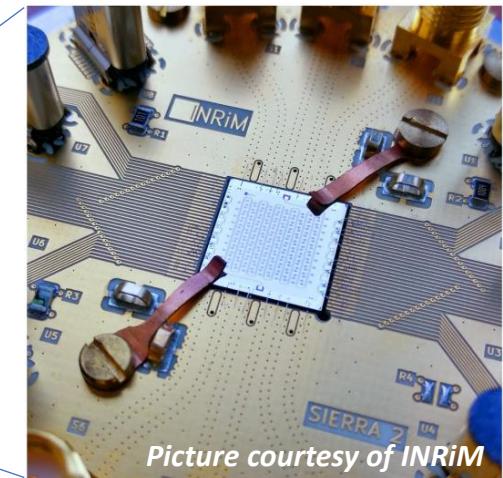
DARTWARS project [INFN]:

1. development of high-performance amplifiers – both KIT and TWJPA – optimizing design, new materials and fabrication processes
 2. demonstration of readout of various detectors/devices with improved performances thanks to the amplification with added noise at the quantum level
- [LE] higher-level investigation of cavity-magnon polaritons using DARTWARS

KIT and TWJPA amplifiers



D. Lachance-Quirion et al., Science (2020)



Picture courtesy of INRIM

From D. Lachance-Quirion et al., SciAdv (2018)

Dilution Fridge – 10 mK

