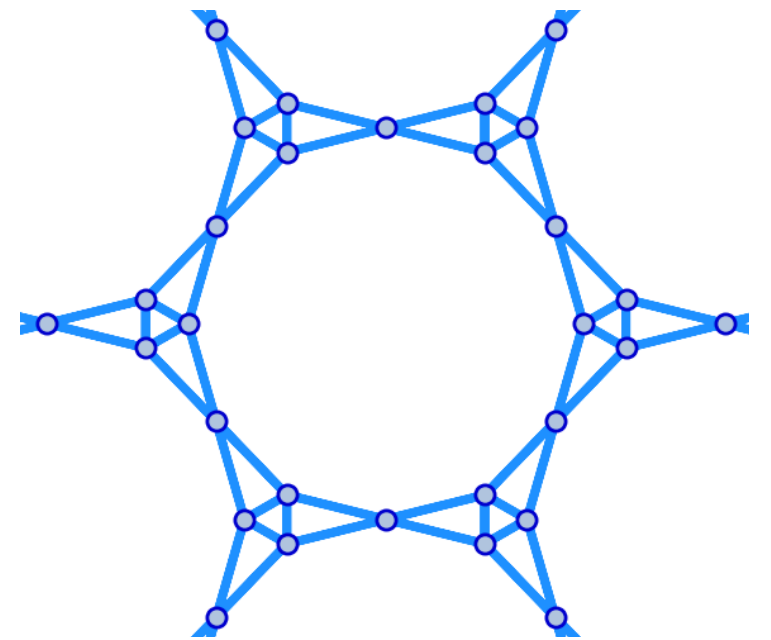
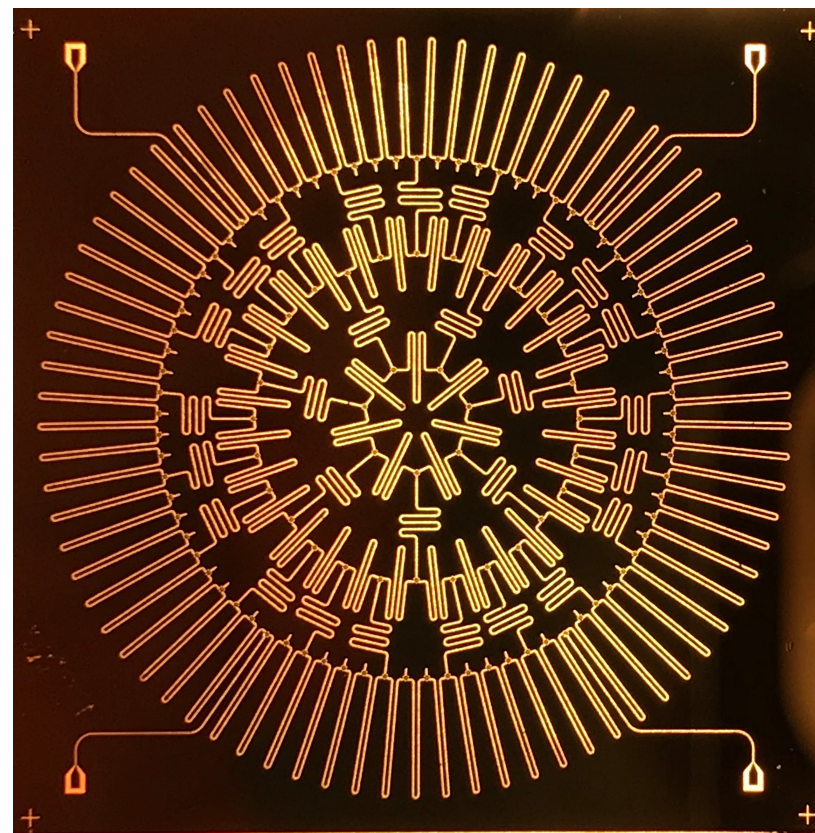
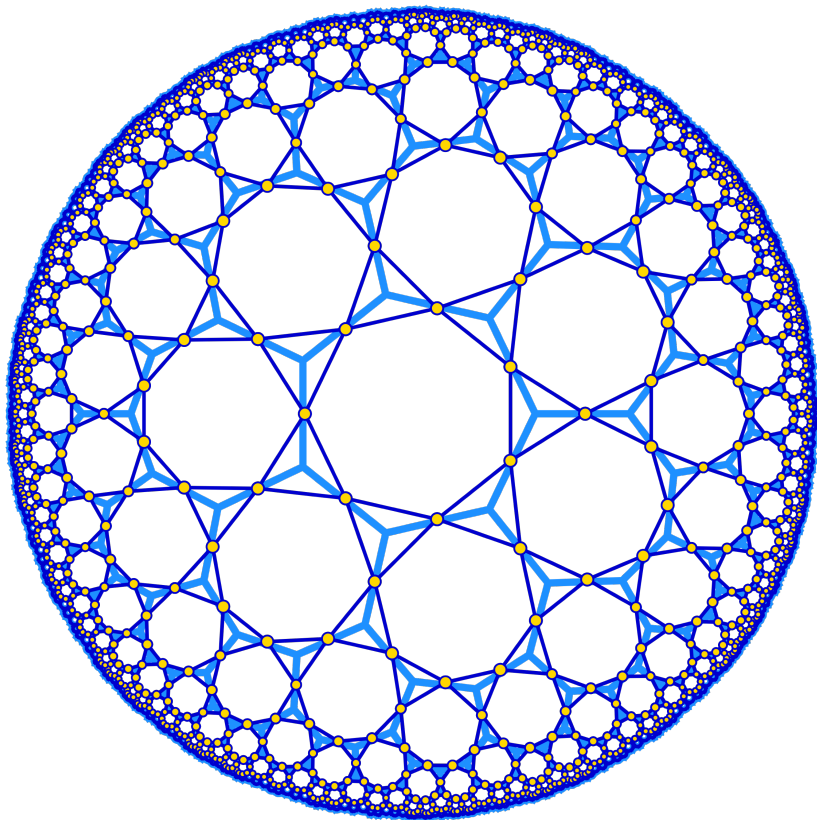


Circuit QED Lattices

Alicia Kollár

Department of Physics and JQI, University of Maryland

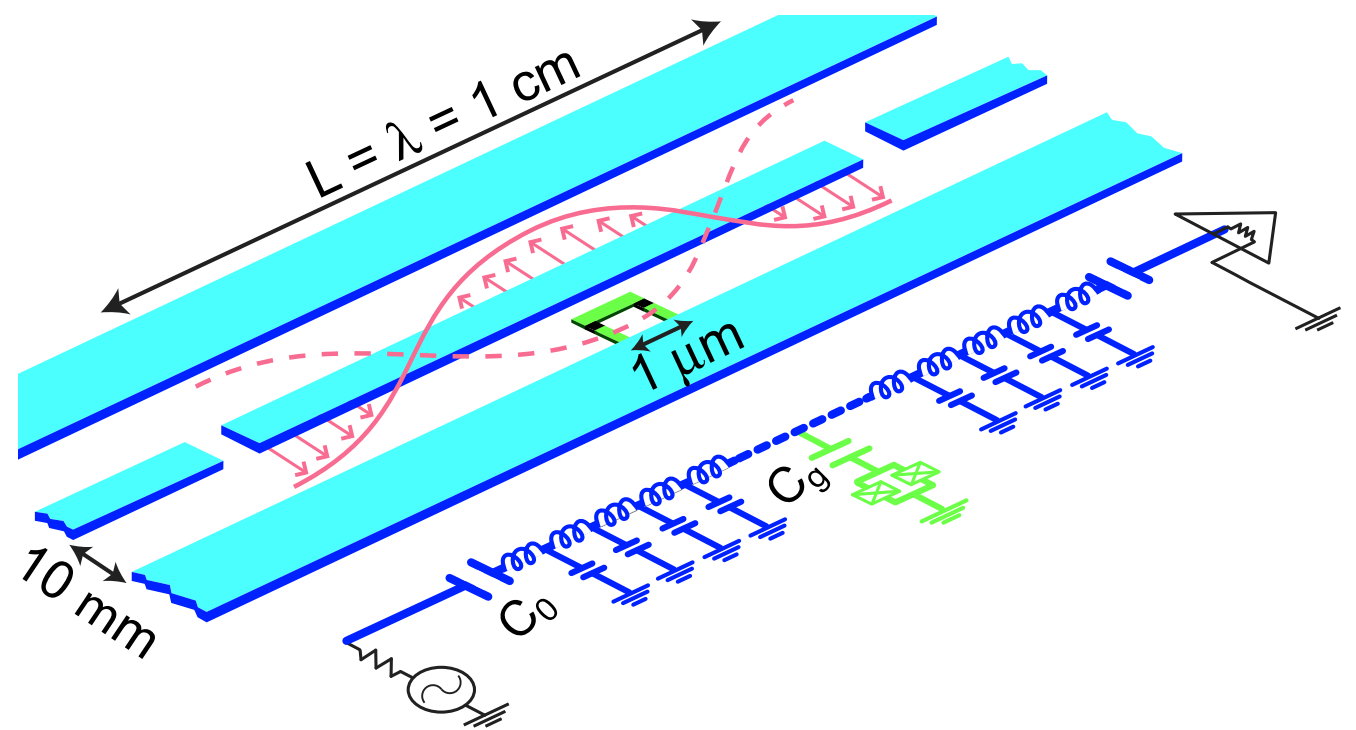


Outline

- Coplanar Waveguide (CPW) Lattices
 - Deformable lattice sites
 - Line-graph lattices
 - Interacting photons
- Band Engineering
 - Hyperbolic lattice
 - Gapped flat bands
- Mathematical Connections
 - Planar Gaps
 - Maximal Gaps
 - Quantum Error Correcting Codes
- Experimental Developments

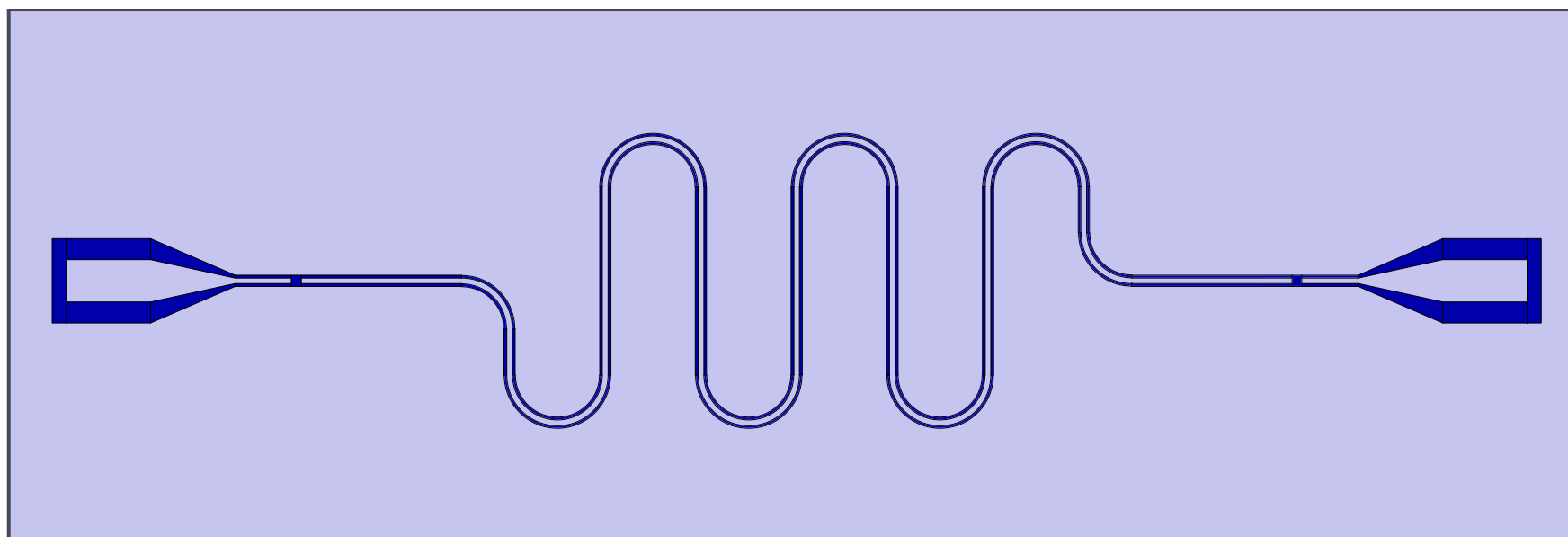
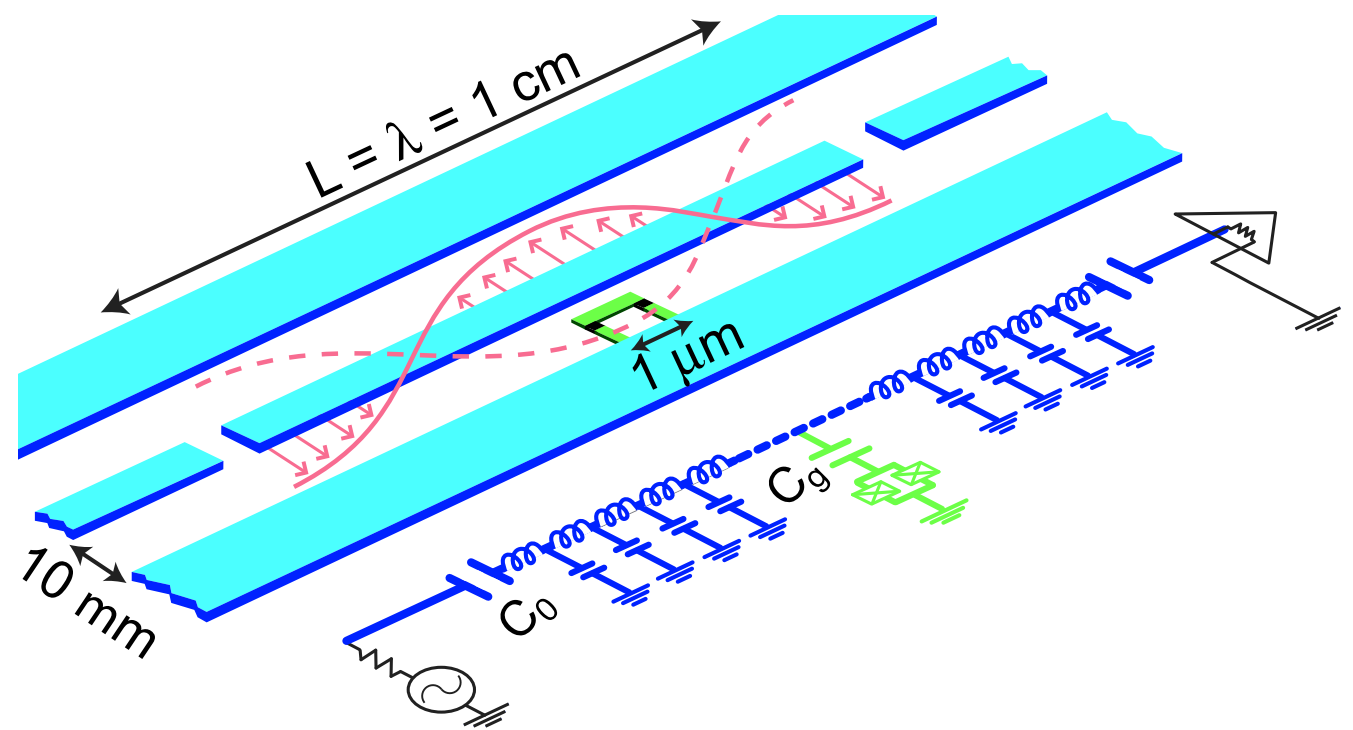
Microwave Coplanar Waveguide Resonators

- 2D analog of coaxial cable
- Cavity defined by cutting center pin
- Voltage antinode at “mirror”



Microwave Coplanar Waveguide Resonators

- 2D analog of coaxial cable
- Cavity defined by cutting center pin
- Voltage antinode at “mirror”

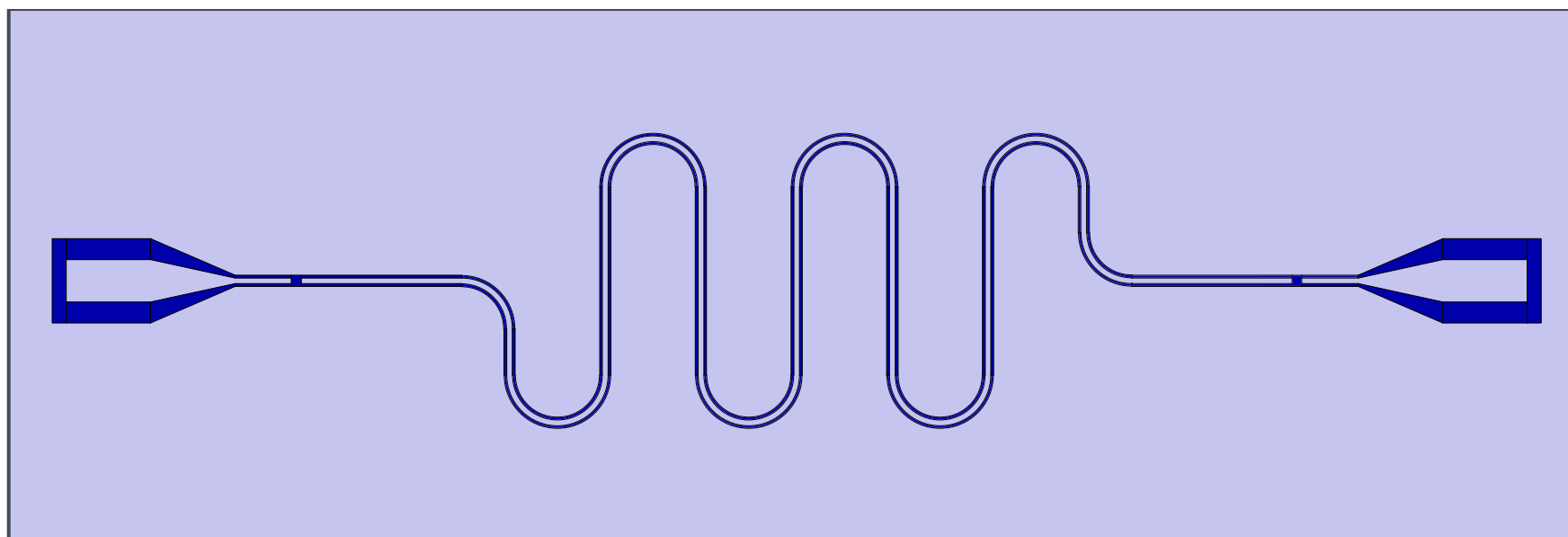
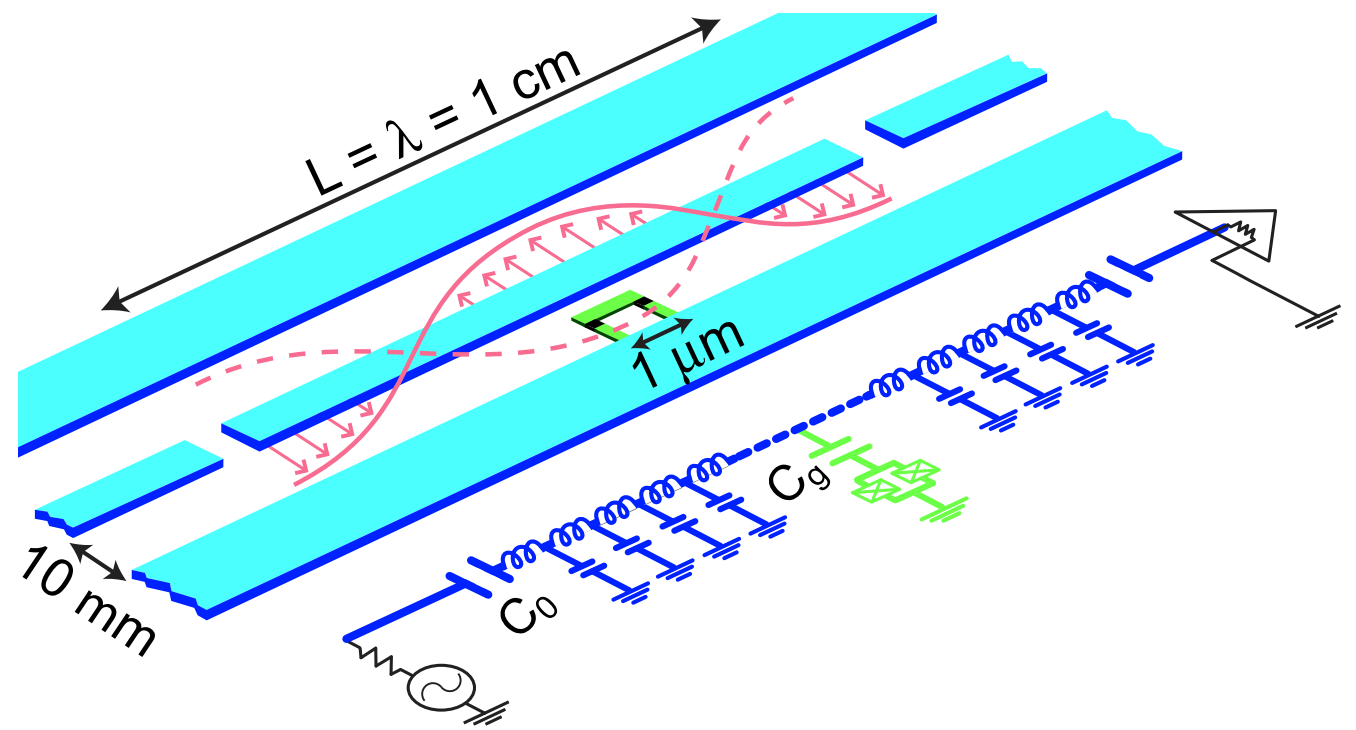


Microwave Coplanar Waveguide Resonators

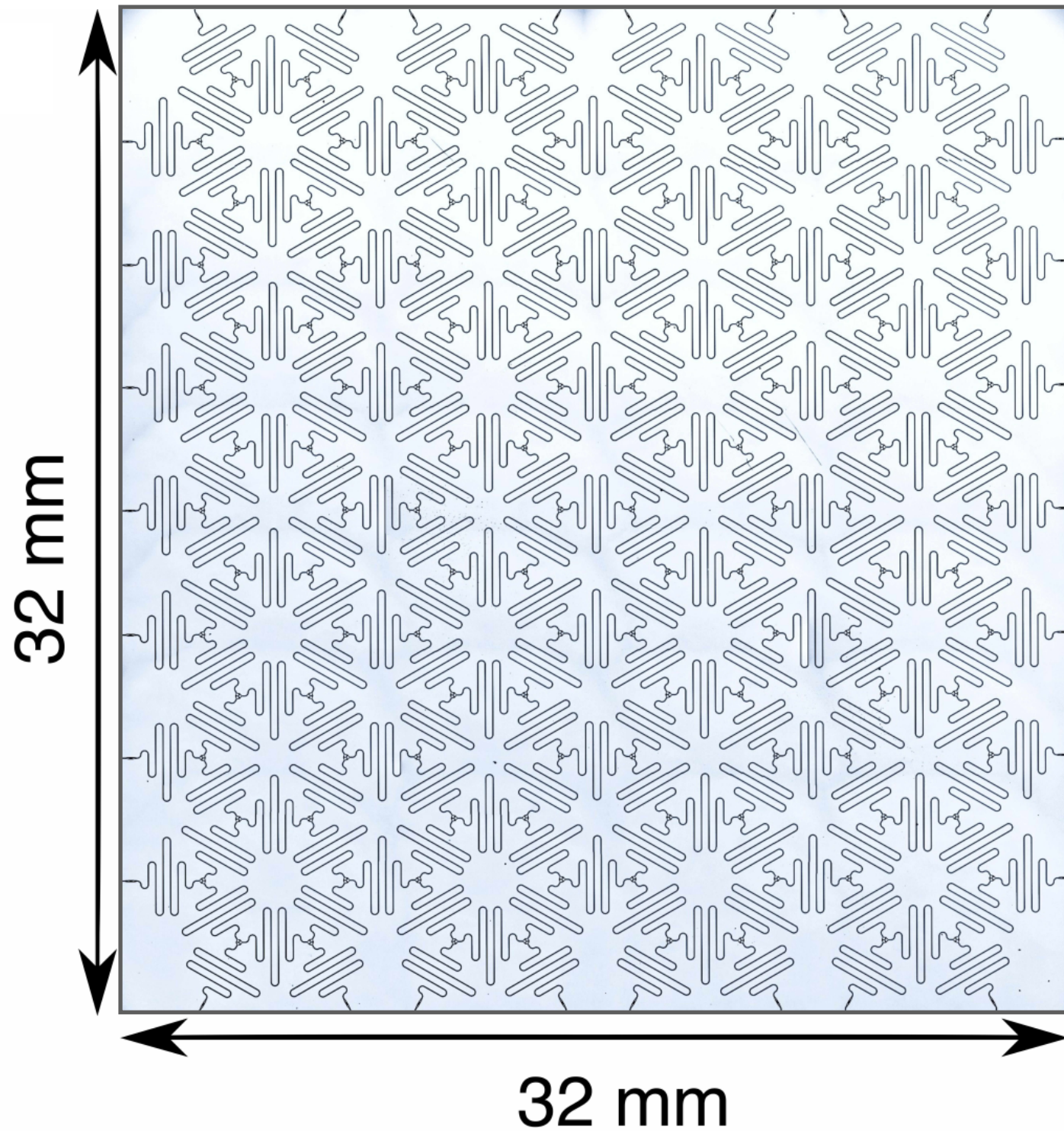
- 2D analog of coaxial cable
- Cavity defined by cutting center pin
- Voltage antinode at “mirror”

Harmonic oscillator

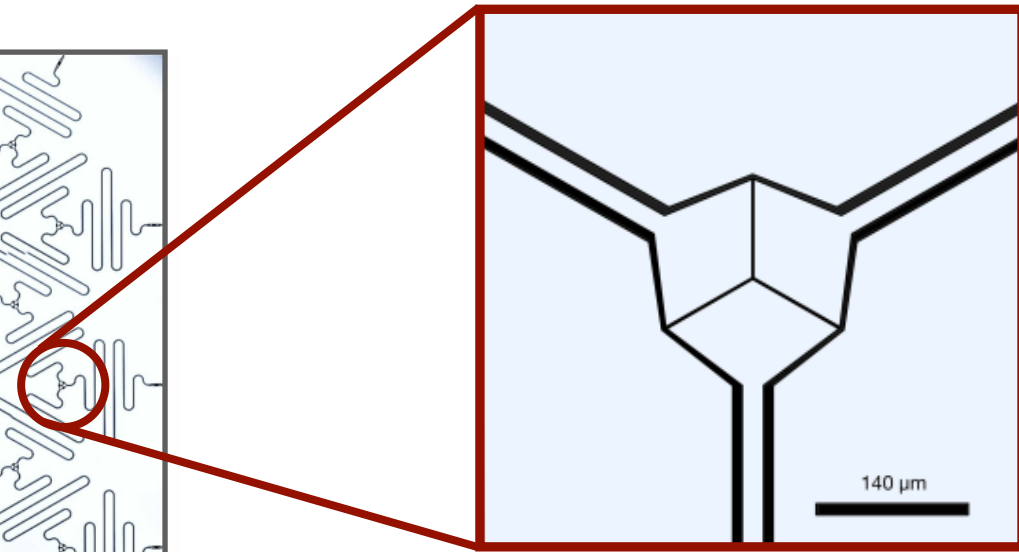
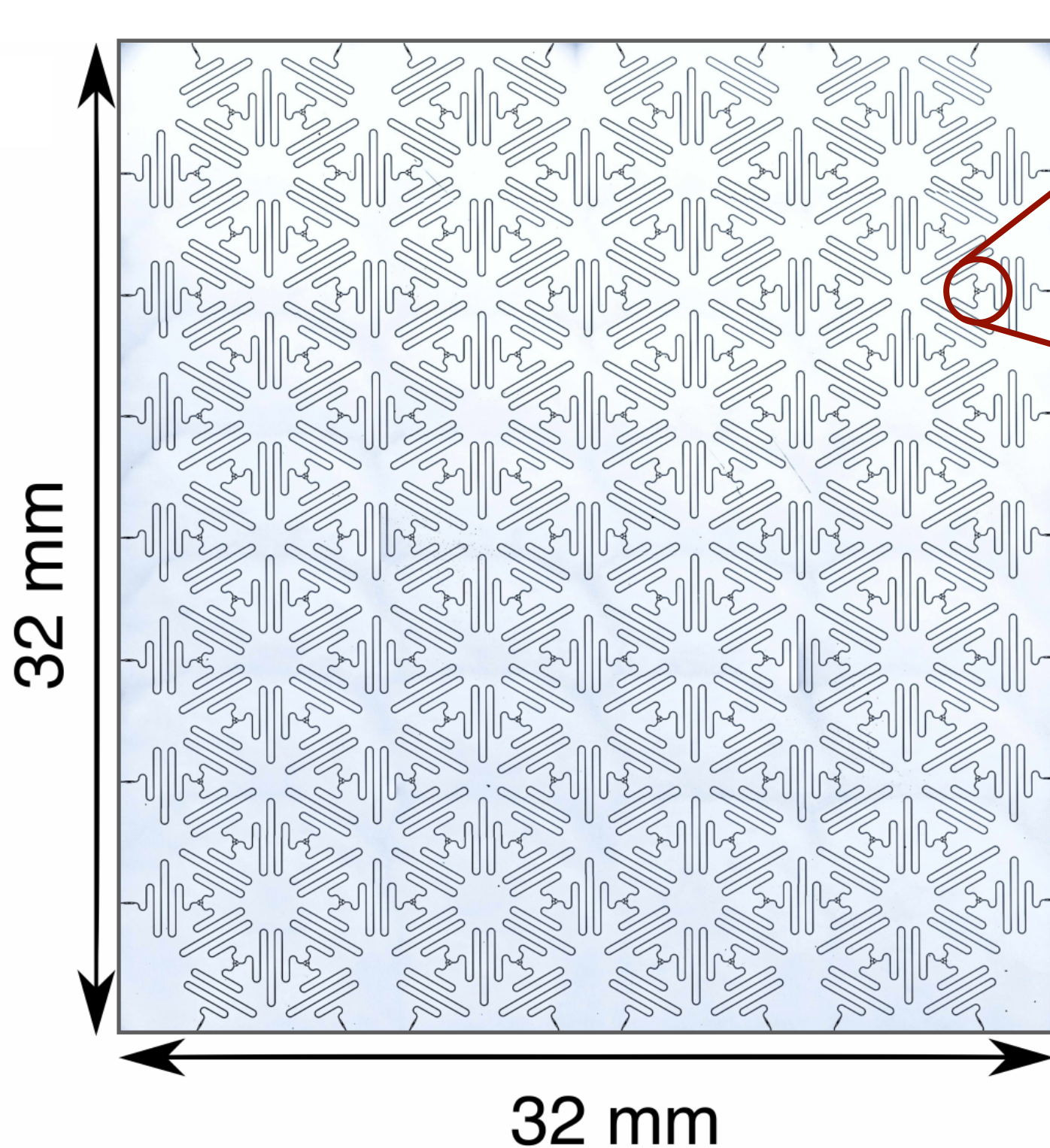
$$\hat{H} = \frac{1}{2C} \hat{n}^2 + \frac{1}{2L} \hat{\phi}^2$$



CPW Lattices

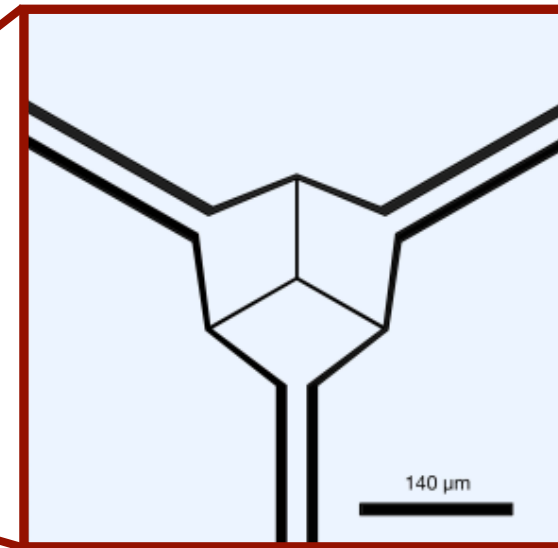
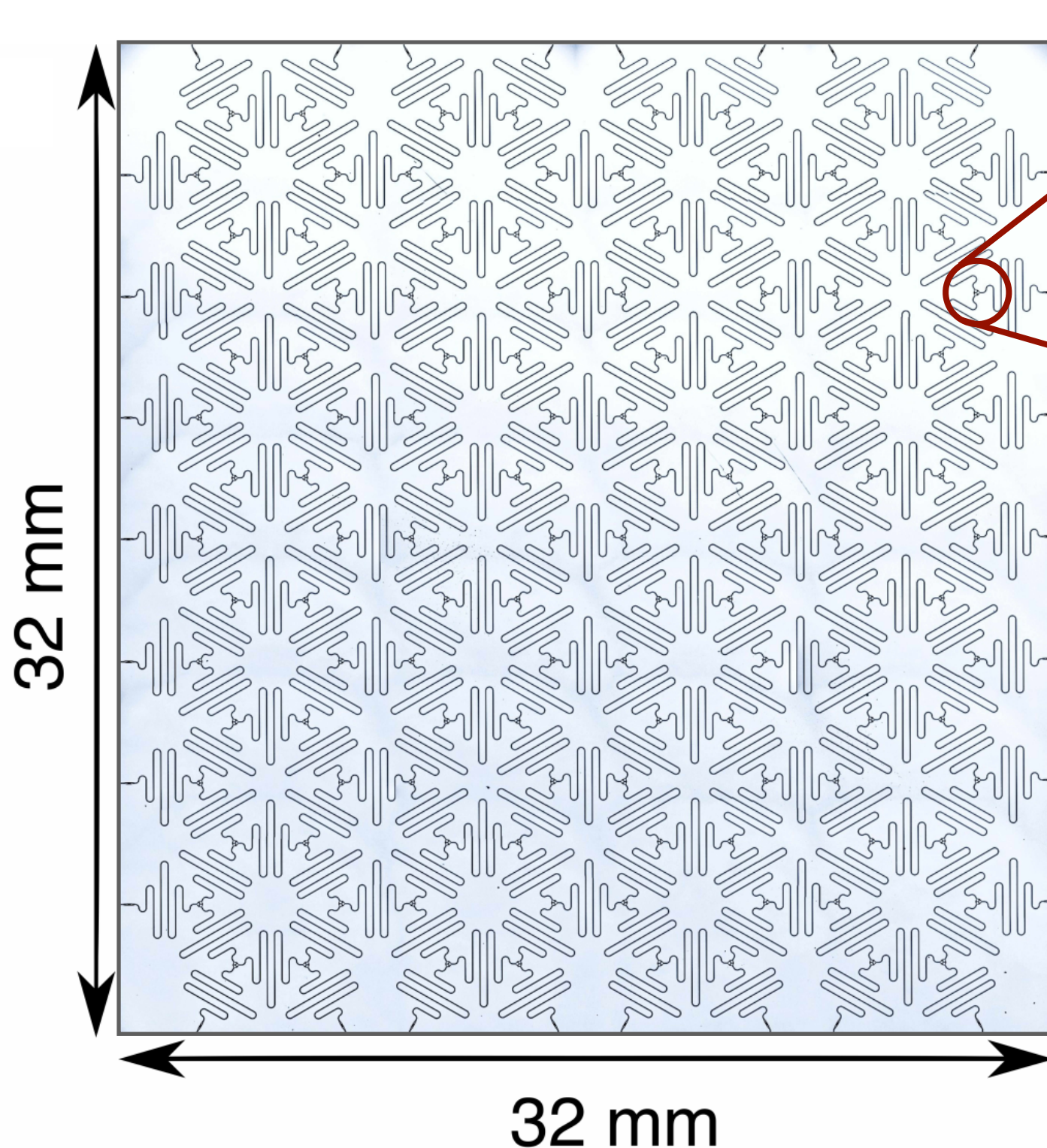


CPW Lattices



- Capacitive coupling of resonators

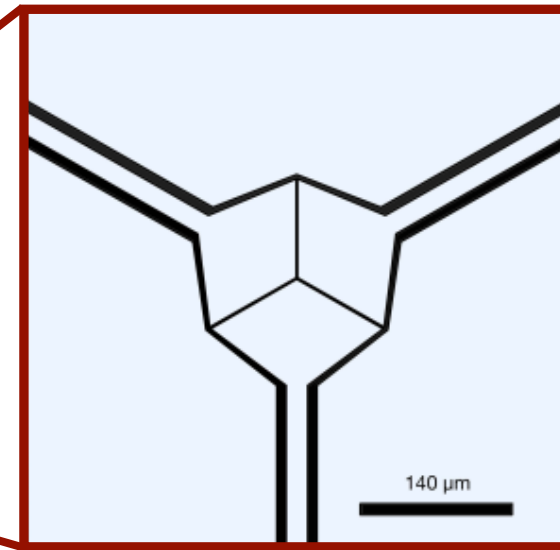
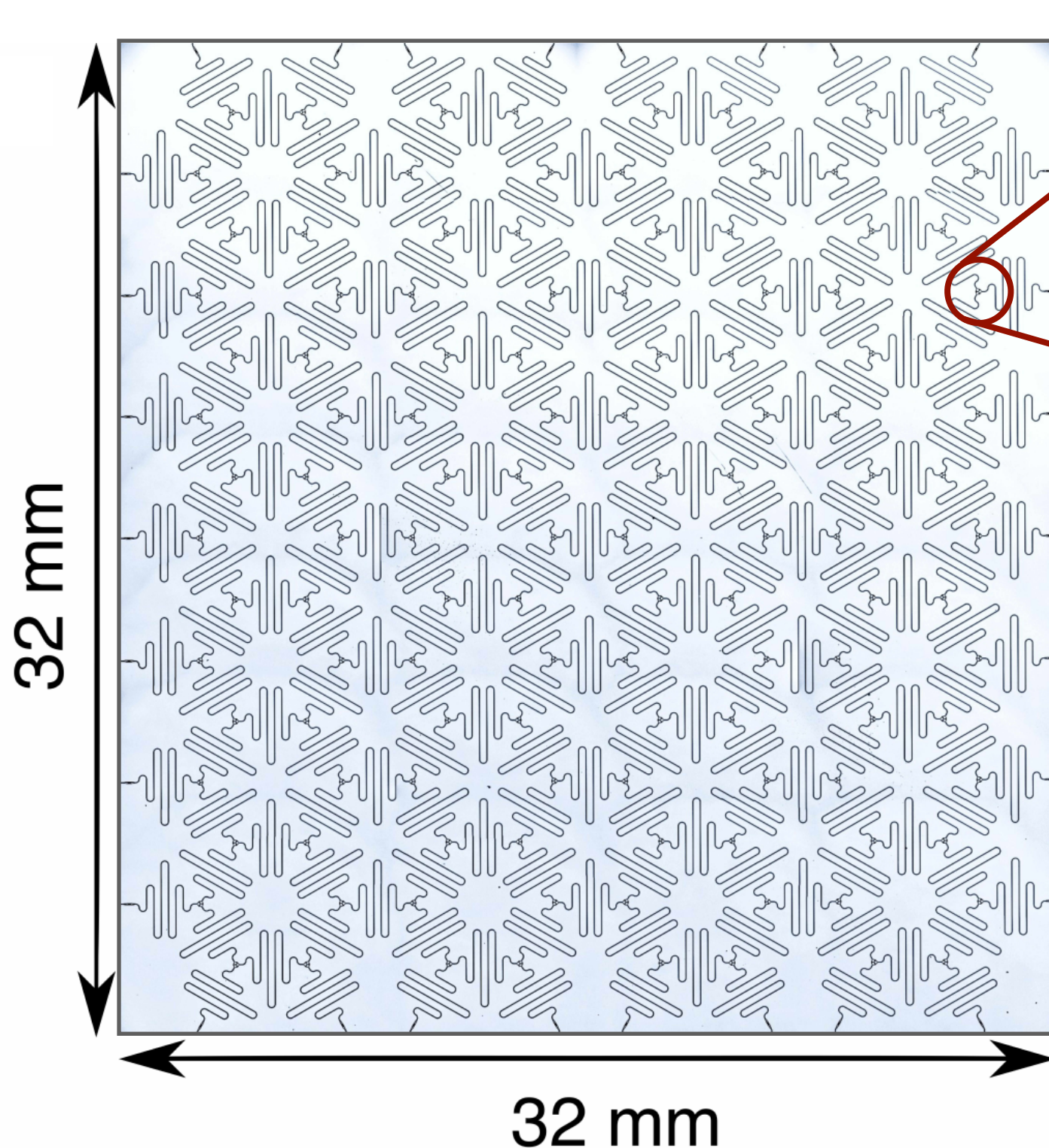
CPW Lattices



- Capacitive coupling of resonators
- Tight-binding solid

$$\mathbf{H}_{\text{TB}} = \omega_0 \sum_{\mathbf{i}} \mathbf{a}_{\mathbf{i}}^{\dagger} \mathbf{a}_{\mathbf{i}} - t \sum_{\langle \mathbf{i}, \mathbf{j} \rangle} (\mathbf{a}_{\mathbf{i}}^{\dagger} \mathbf{a}_{\mathbf{j}} + \mathbf{a}_{\mathbf{j}}^{\dagger} \mathbf{a}_{\mathbf{i}})$$

CPW Lattices



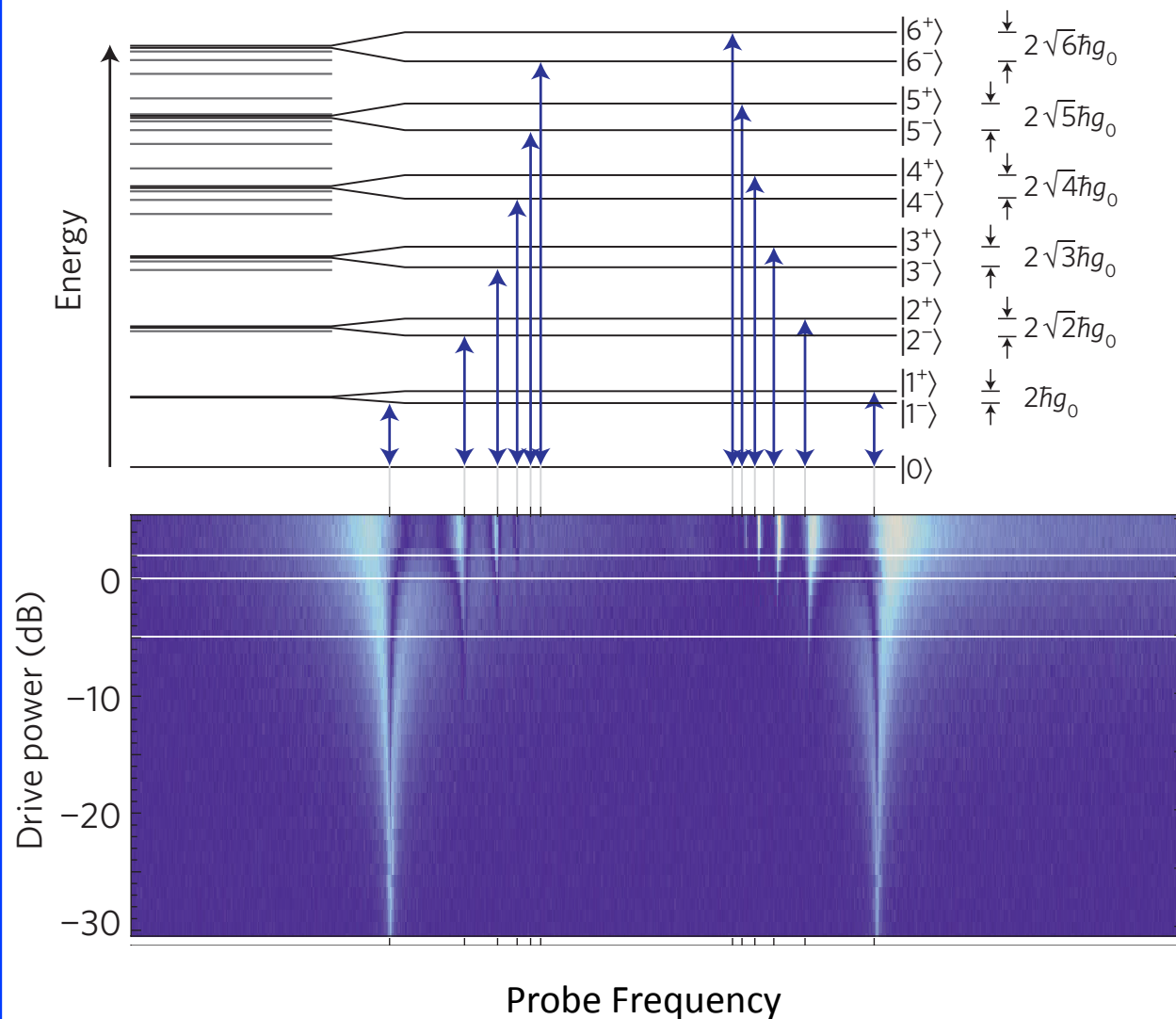
- Capacitive coupling of resonators
- Tight-binding solid
- $t < 0$

$$\mathbf{H}_{\text{TB}} = \omega_0 \sum_{\mathbf{i}} \mathbf{a}_{\mathbf{i}}^{\dagger} \mathbf{a}_{\mathbf{i}} - t \sum_{\langle \mathbf{i}, \mathbf{j} \rangle} (\mathbf{a}_{\mathbf{i}}^{\dagger} \mathbf{a}_{\mathbf{j}} + \mathbf{a}_{\mathbf{j}}^{\dagger} \mathbf{a}_{\mathbf{i}})$$

Combining Lattices and Qubits

Number-Resolved Transitions

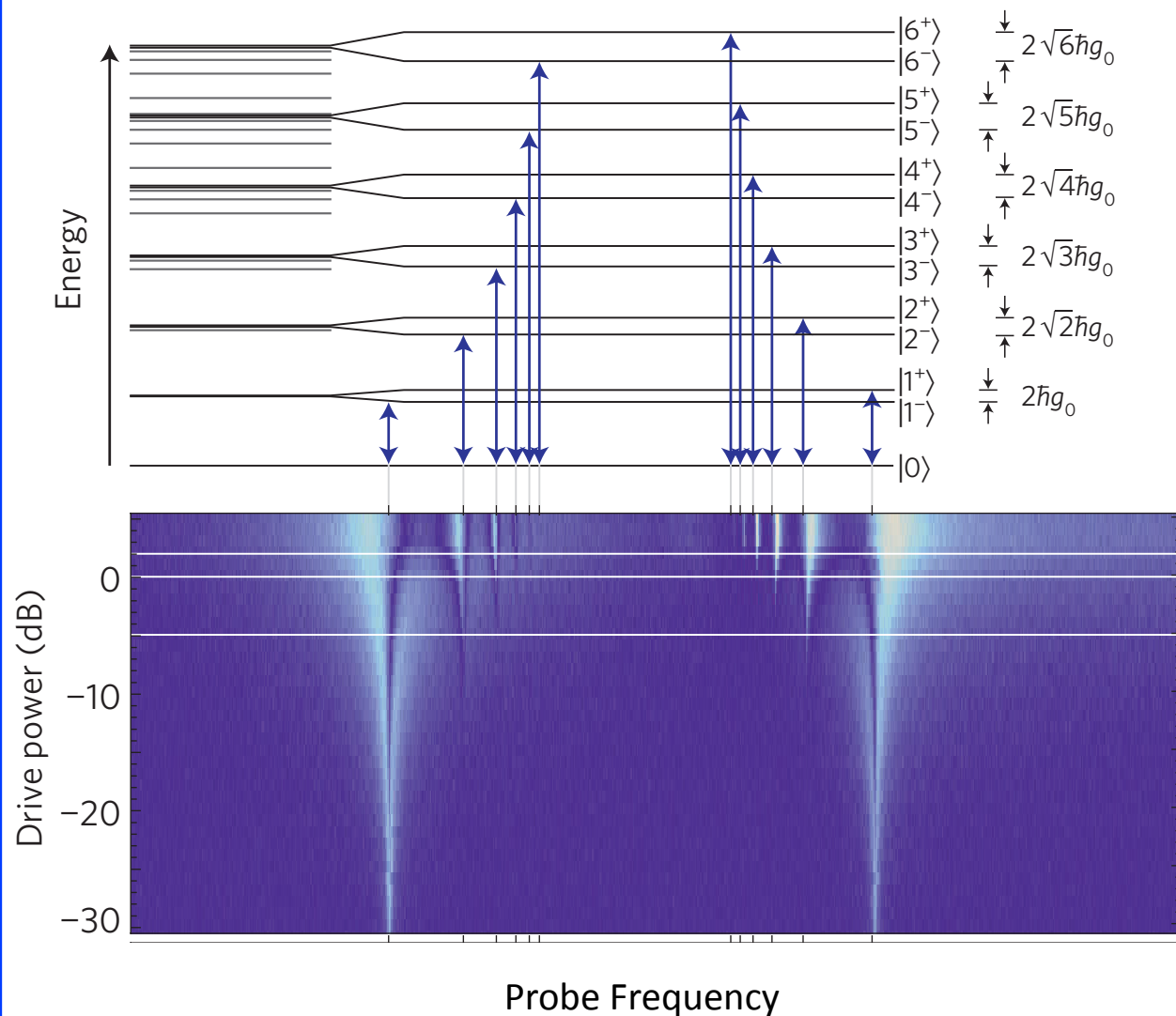
$$H_{JC} = \omega_c a^\dagger a + \frac{1}{2} \omega_q \sigma_z + g_0 (a^\dagger \sigma^- + a \sigma^+)$$



Combining Lattices and Qubits

Number-Resolved Transitions

$$H_{JC} = \omega_c a^\dagger a + \frac{1}{2} \omega_q \sigma_z + g_0 (a^\dagger \sigma^- + a \sigma^+)$$



Qubits in Photonic Crystals

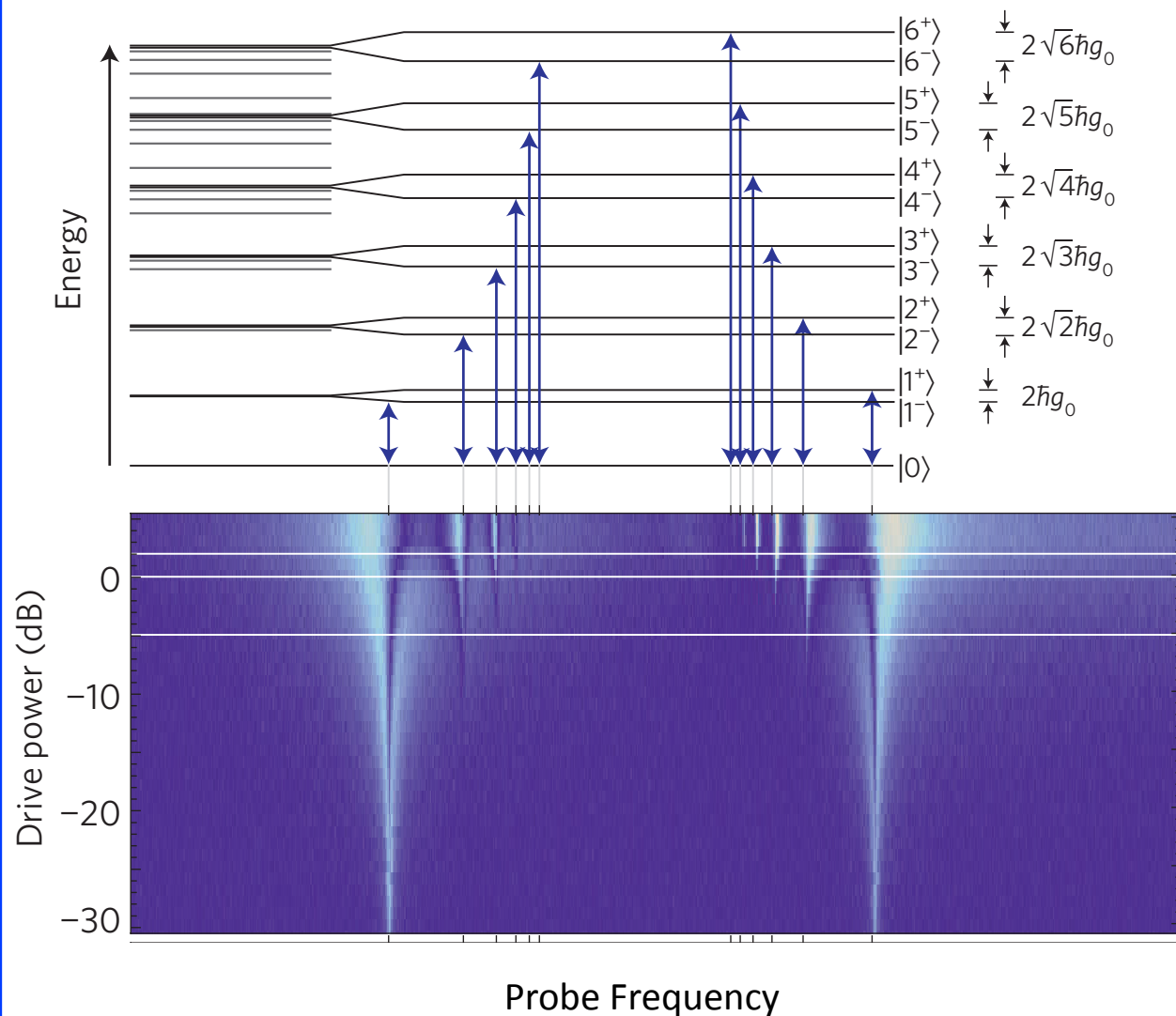
- Effective swap interaction between qubits
- All modes in parallel

$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

Combining Lattices and Qubits

Number-Resolved Transitions

$$H_{JC} = \omega_c a^\dagger a + \frac{1}{2} \omega_q \sigma_z + g_0 (a^\dagger \sigma^- + a \sigma^+)$$



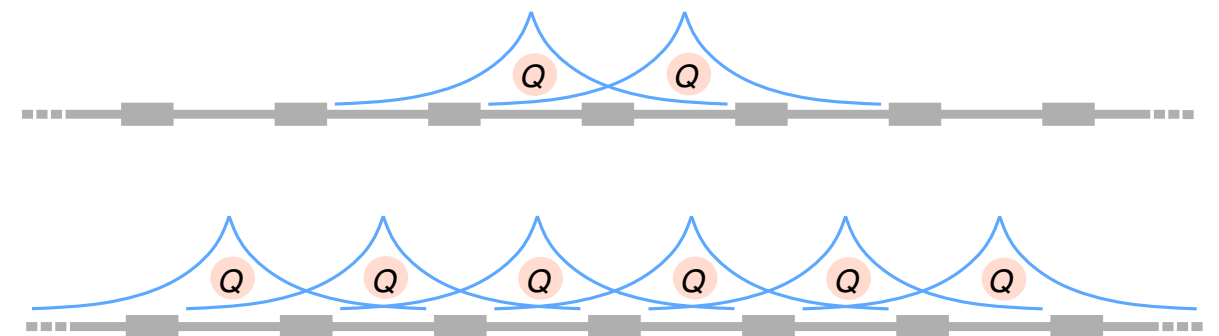
Qubits in Photonic Crystals

- Effective swap interaction between qubits
- All modes in parallel

$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

1D-Photonic Crystal

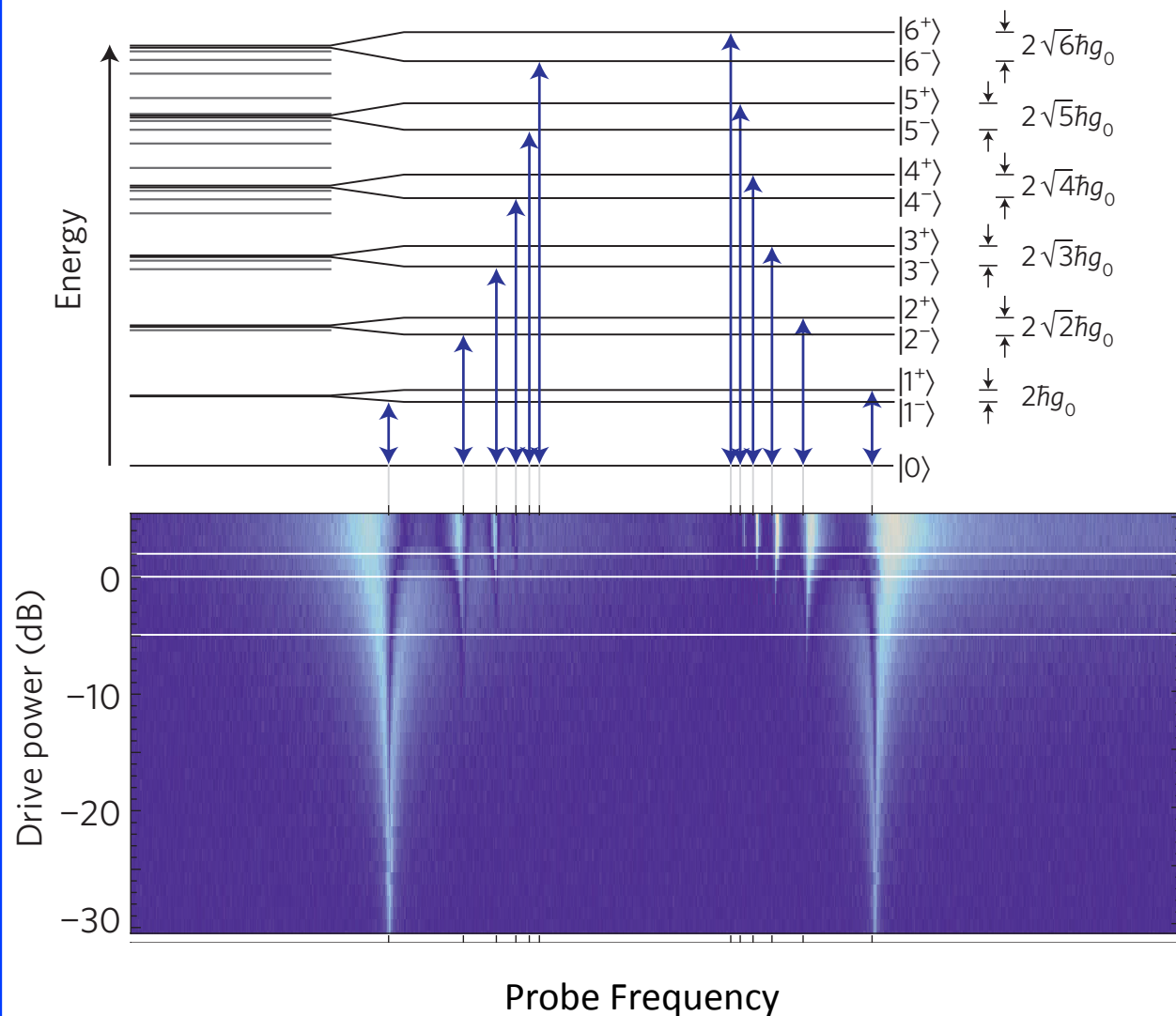
- Exponentially localized bound state



Combining Lattices and Qubits

Number-Resolved Transitions

$$H_{JC} = \omega_c a^\dagger a + \frac{1}{2} \omega_q \sigma_z + g_0 (a^\dagger \sigma^- + a \sigma^+)$$



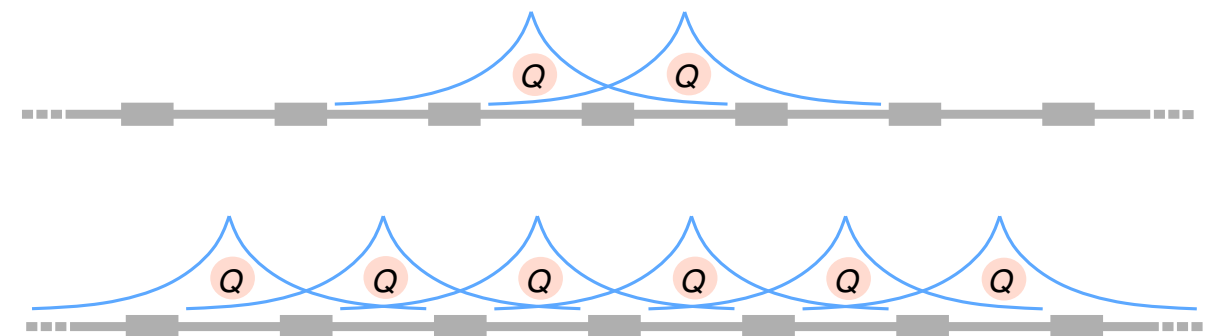
Qubits in Photonic Crystals

- Effective swap interaction between qubits
- All modes in parallel

$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

1D-Photonic Crystal

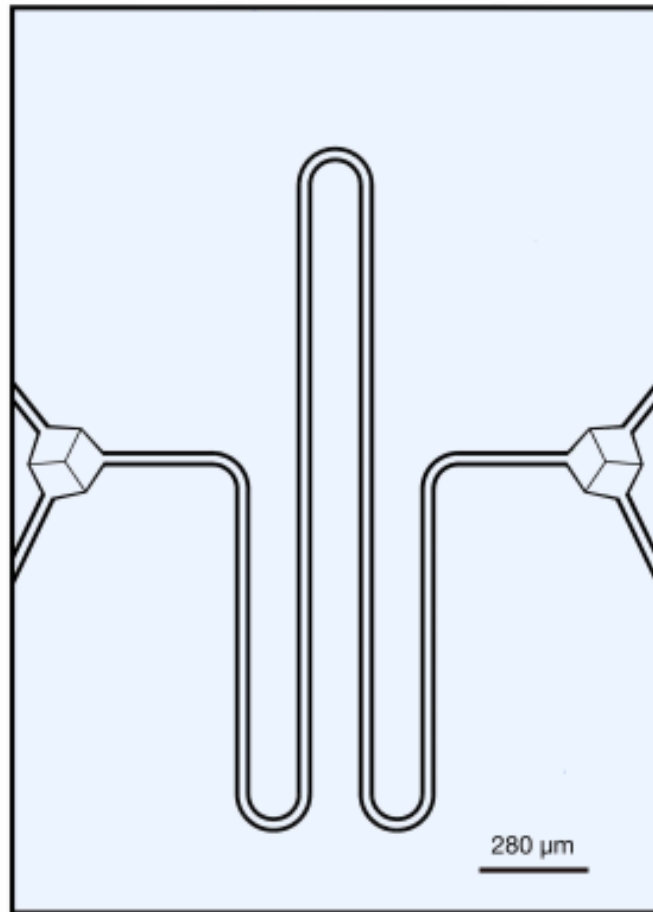
- Exponentially localized bound state



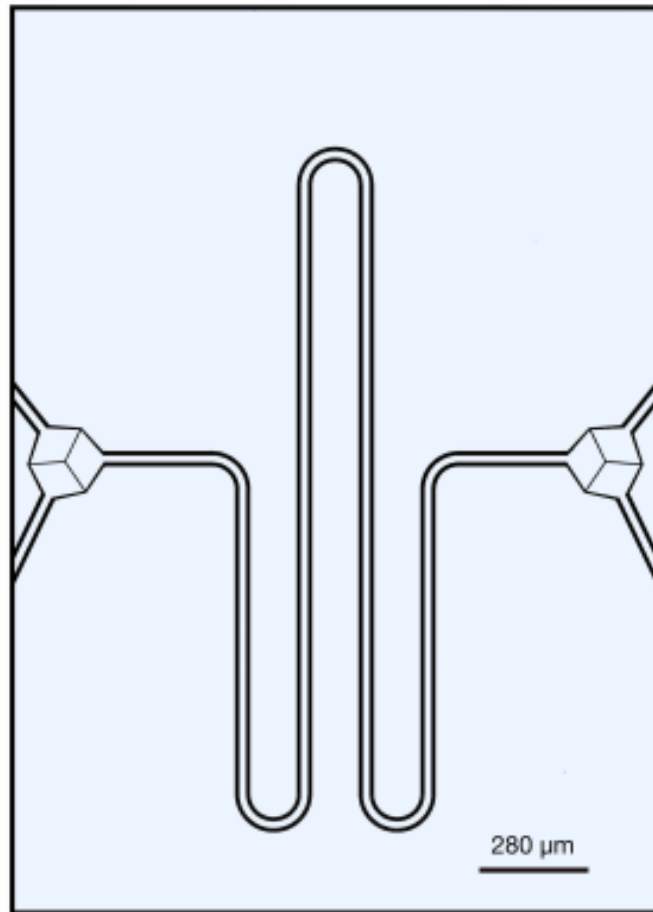
New Regimes:

- New lattices

Deformable Resonators

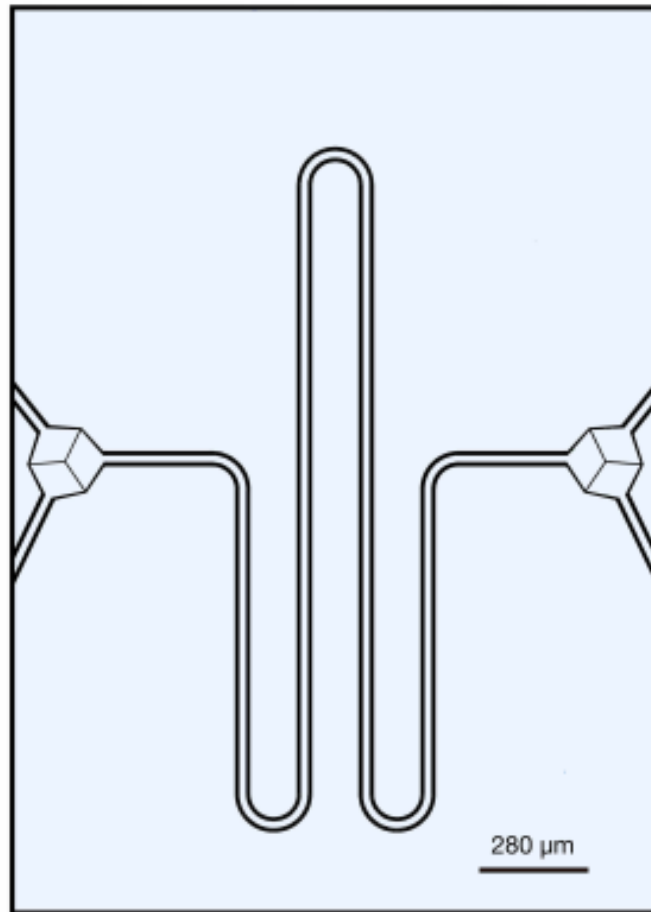


Deformable Resonators



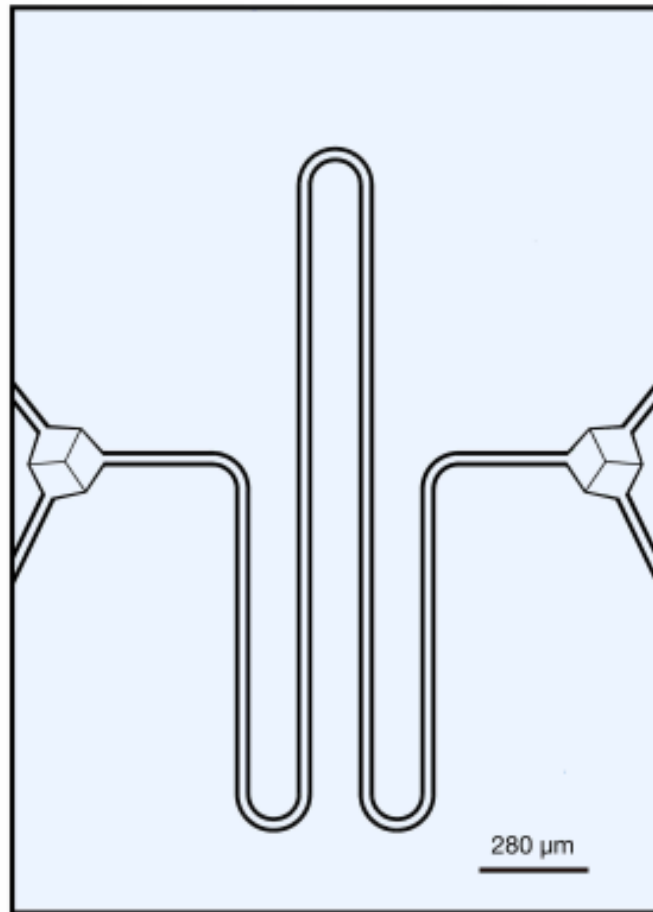
- Frequency depends only on length

Deformable Resonators



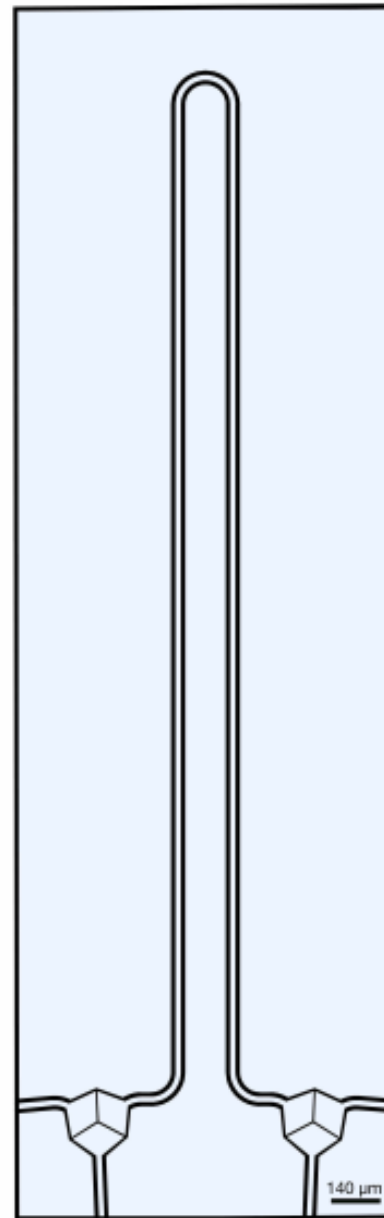
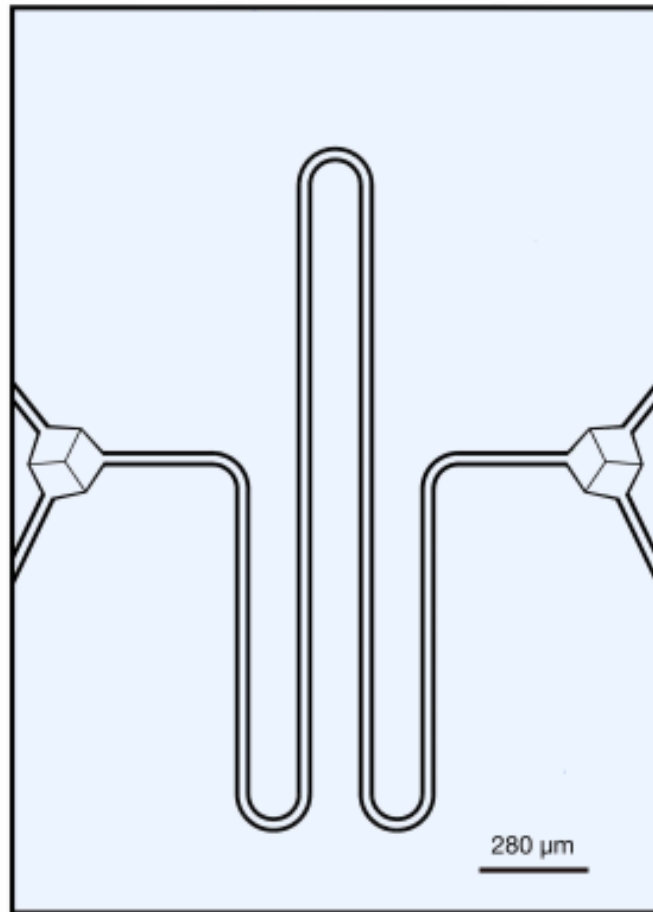
- Frequency depends only on length
- Coupling depends on ends

Deformable Resonators



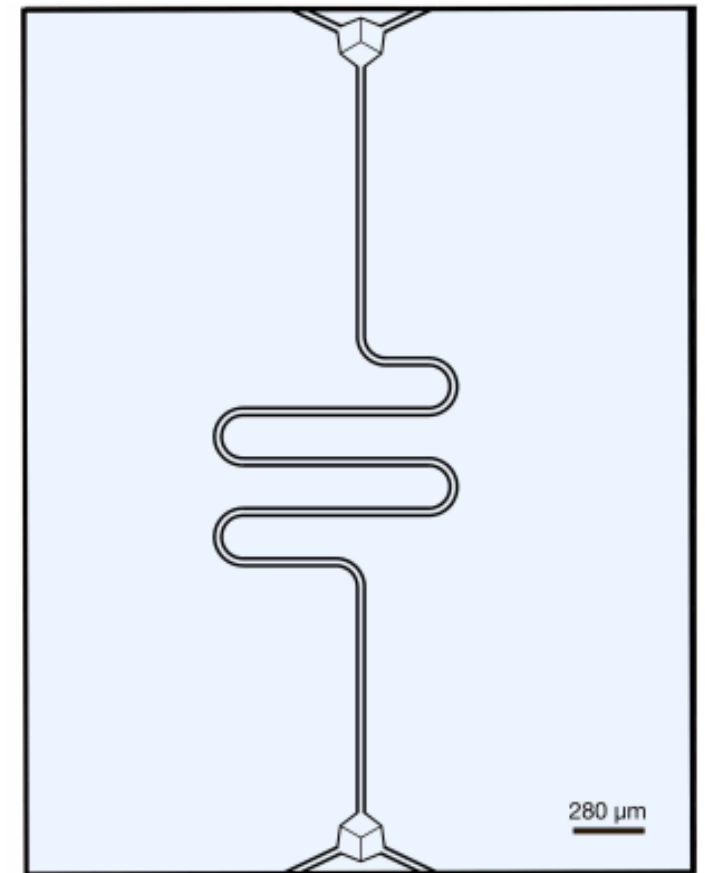
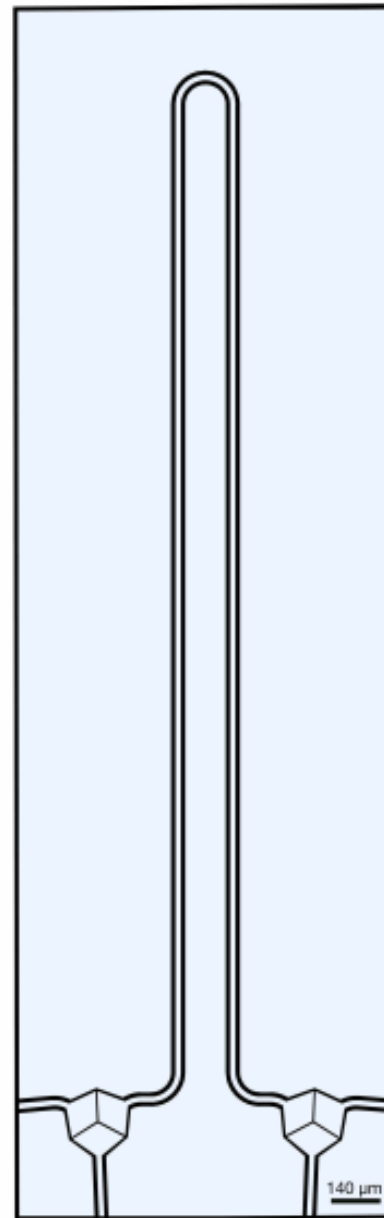
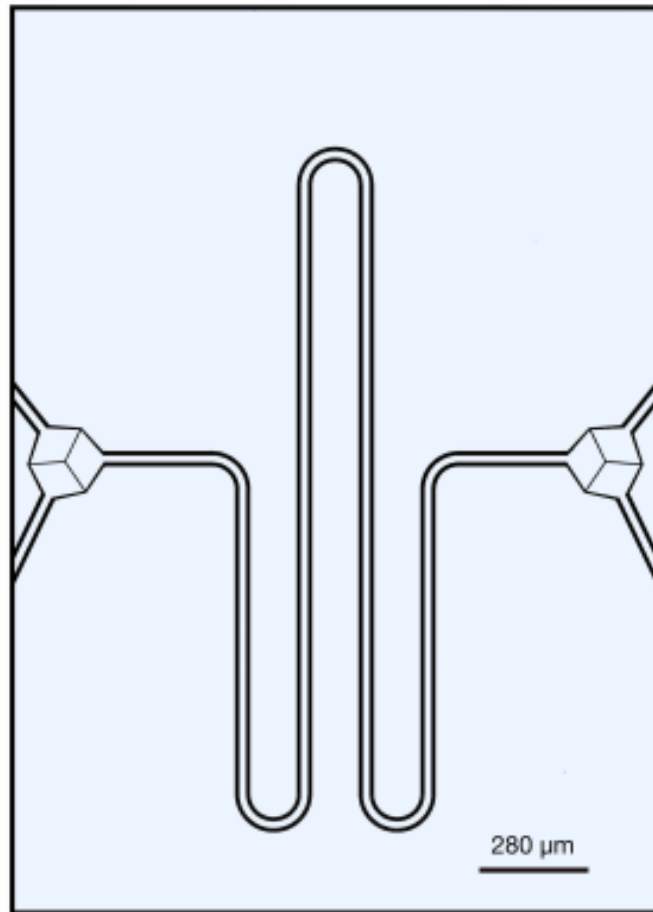
- Frequency depends only on length
- Coupling depends on ends
- “Bendable”

Deformable Resonators



- Frequency depends only on length
- Coupling depends on ends
- “Bendable”

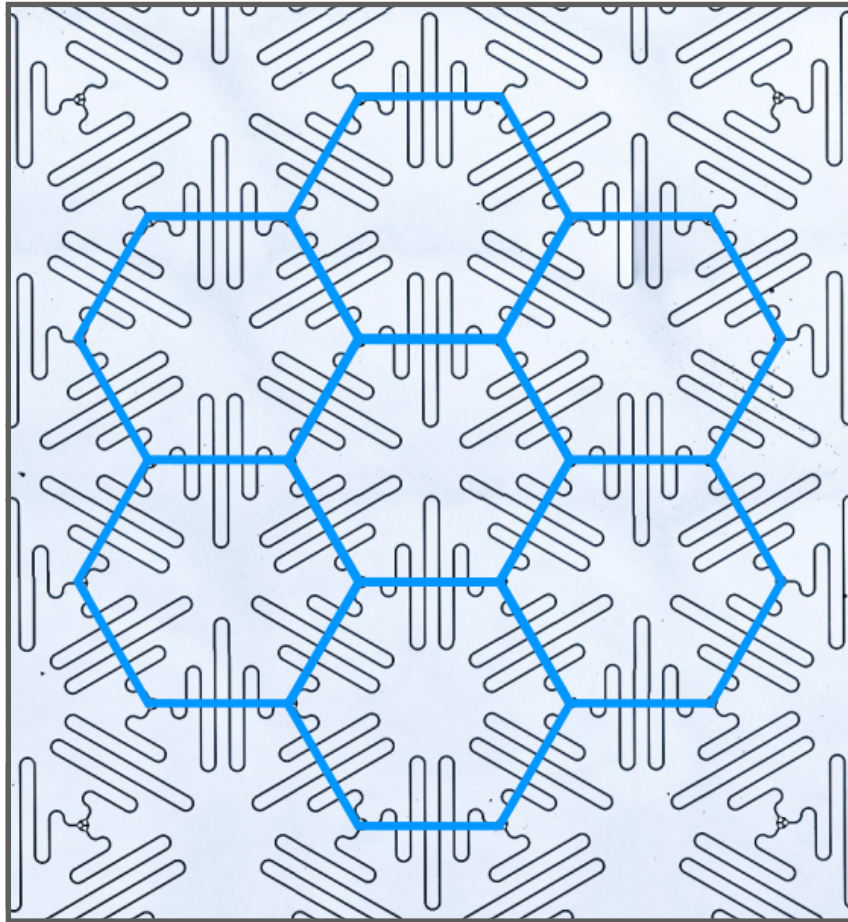
Deformable Resonators



- Frequency depends only on length
- Coupling depends on ends
- “Bendable”

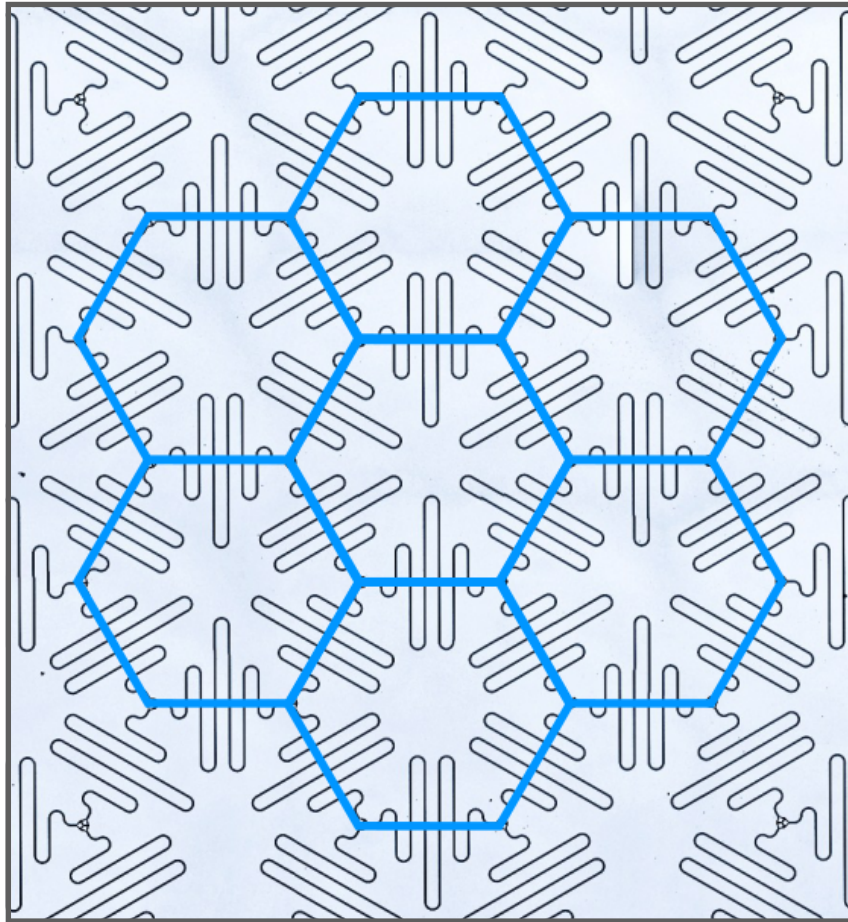
Layout and Effective Lattices

Resonator Lattice



Layout and Effective Lattices

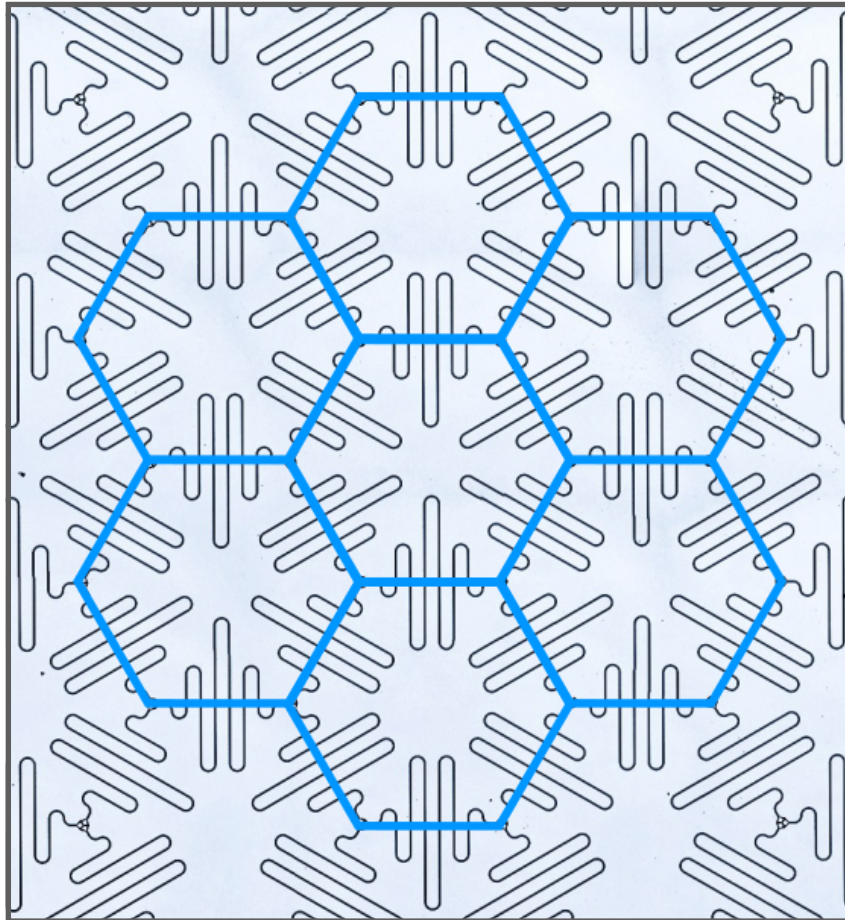
Resonator Lattice



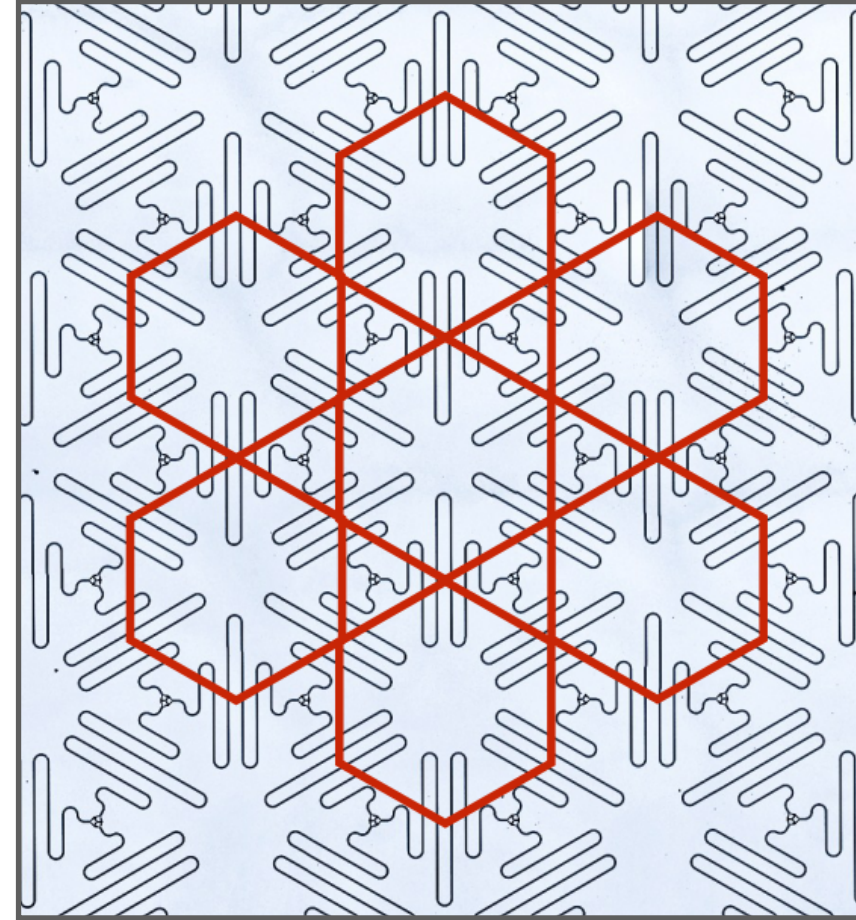
- An *edge* on each resonator

Layout and Effective Lattices

Resonator Lattice



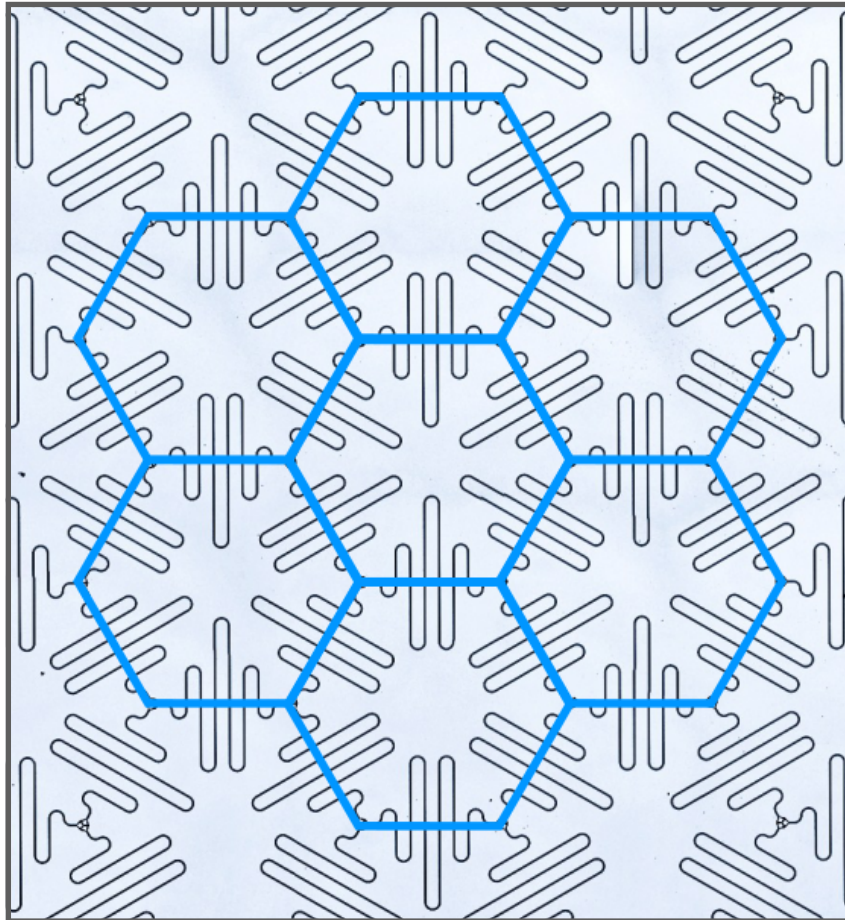
Effective Photonic Lattice



- An *edge* on each resonator

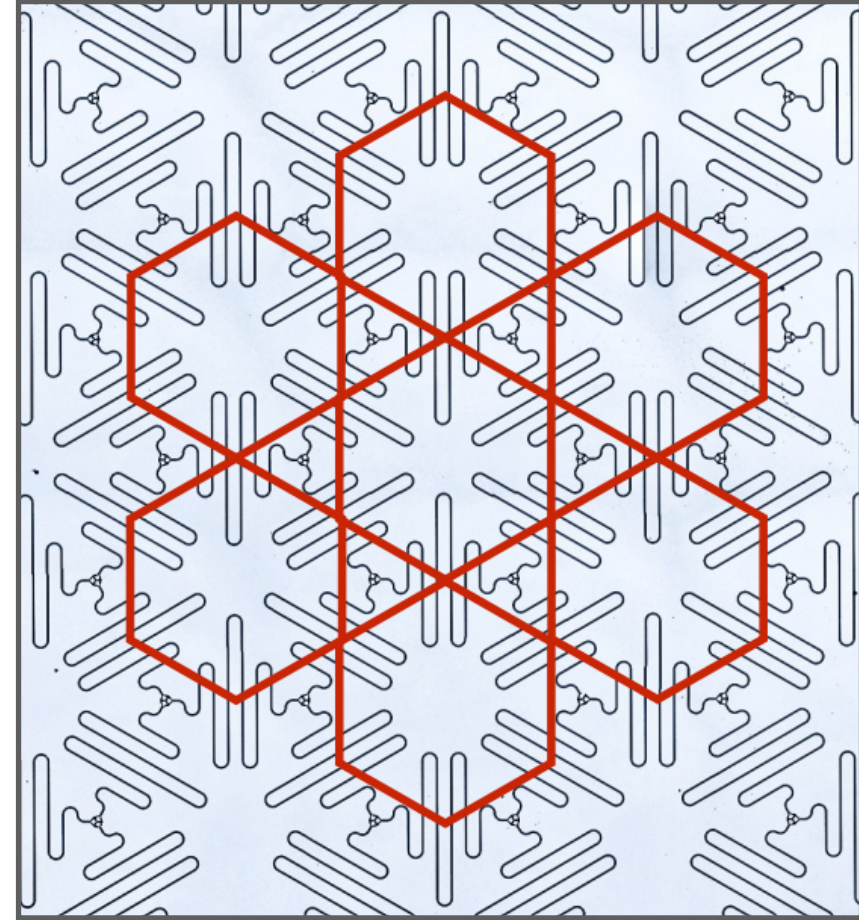
Layout and Effective Lattices

Resonator Lattice



- An *edge* on each resonator

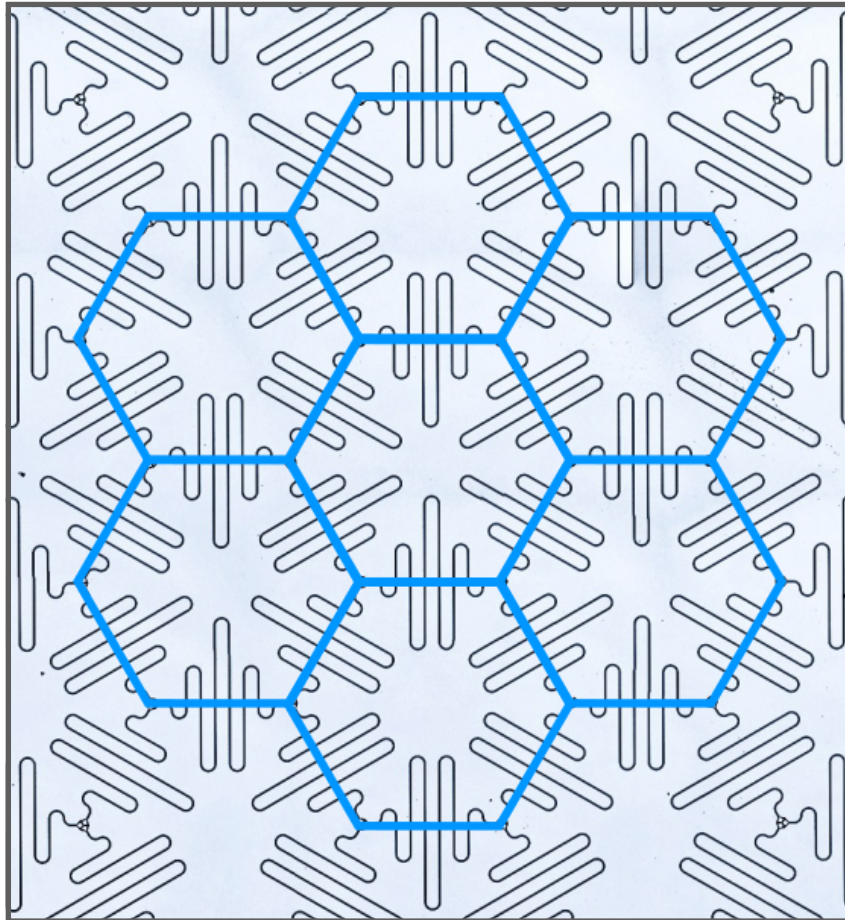
Effective Photonic Lattice



- A *vertex* on each resonator

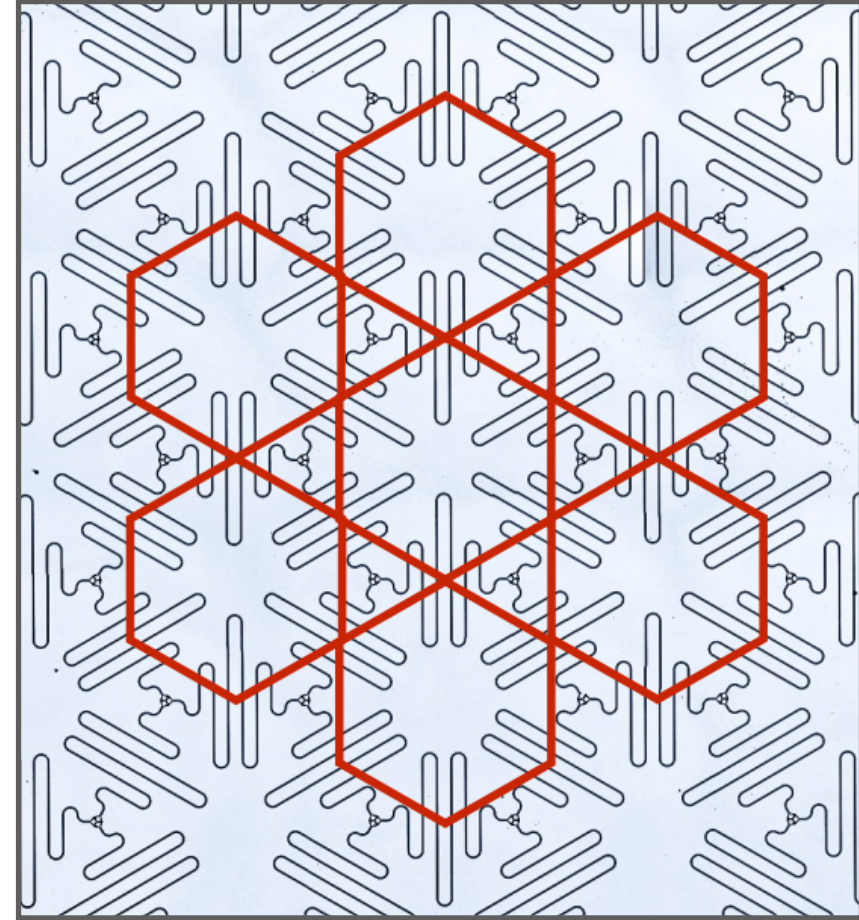
Layout and Effective Lattices

Resonator Lattice



- An *edge* on each resonator

Effective Photonic Lattice

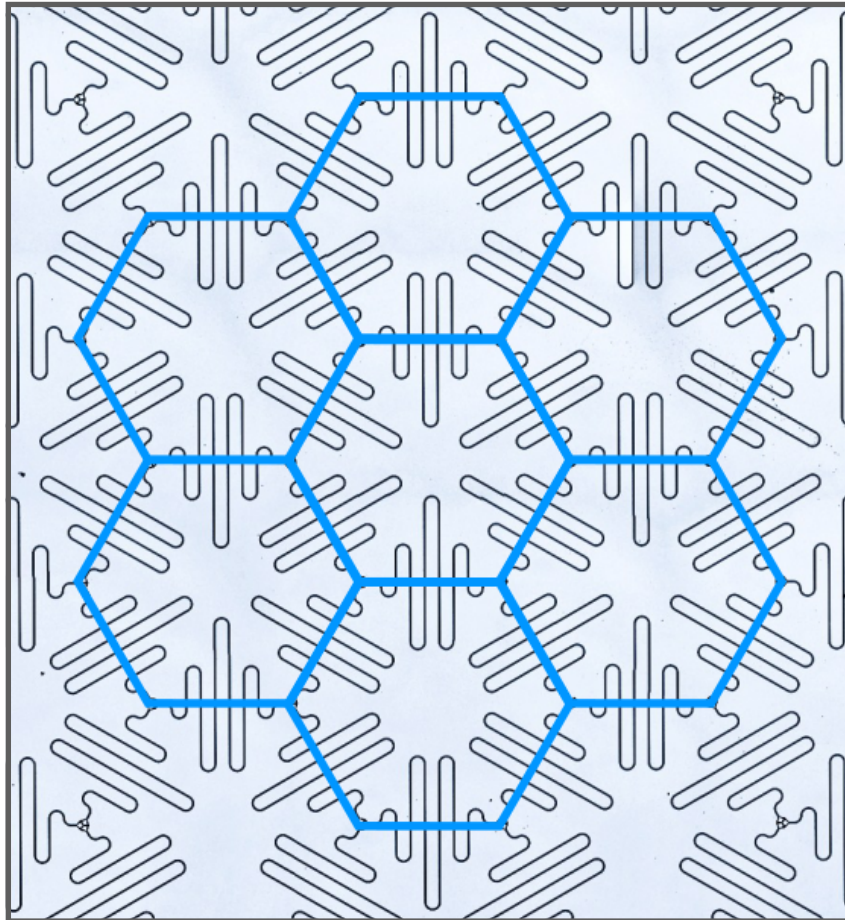


- A *vertex* on each resonator

Layout X

Layout and Effective Lattices

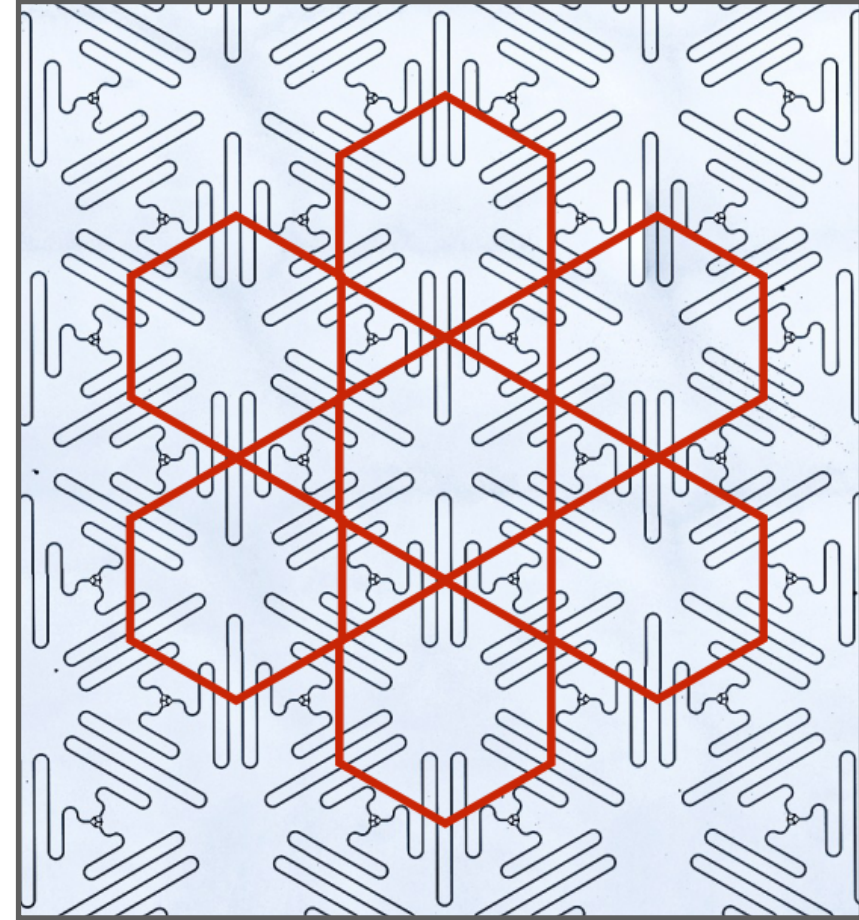
Resonator Lattice



- An *edge* on each resonator

Layout X

Effective Photonic Lattice



- A *vertex* on each resonator

Line Graph $L(X)$

Outline

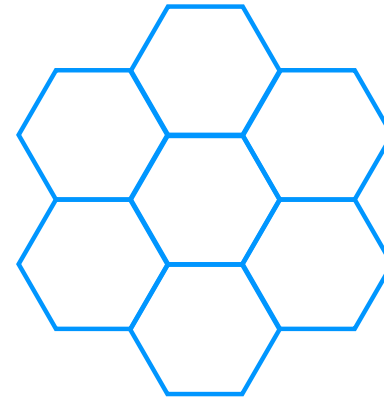
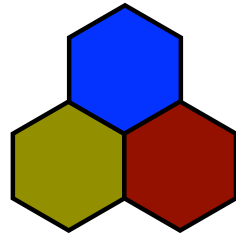
- Coplanar Waveguide (CPW) Lattices
 - Deformable lattice sites
 - Line-graph lattices
 - Interacting photons
- Band Engineering
 - Hyperbolic lattice
 - Gapped flat bands
- Mathematical Connections
 - Planar Gaps
 - Maximal Gaps
 - Quantum Error Correcting Codes
- Experimental Developments

Outline

- Coplanar Waveguide (CPW) Lattices
 - Deformable lattice sites
 - Line-graph lattices
 - Interacting photons
- Band Engineering
 - Hyperbolic lattice
 - Gapped flat bands
- Mathematical Connections
 - Planar Gaps
 - Maximal Gaps
 - Quantum Error Correcting Codes
- Experimental Developments

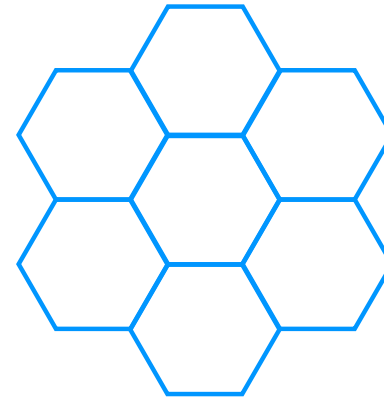
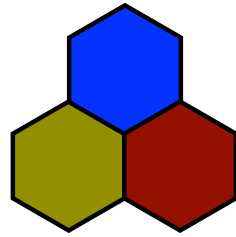
Projecting to Flat 2D

$n = 6$
flat

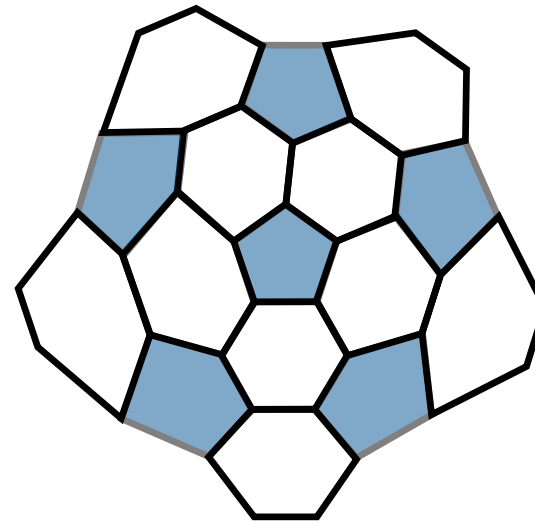
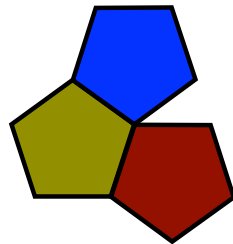


Projecting to Flat 2D

$n = 6$
flat

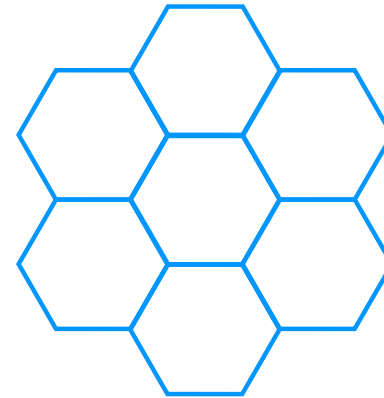
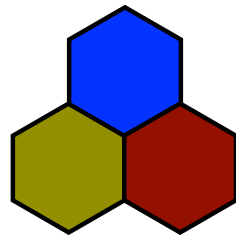


$n = 5$
spherical

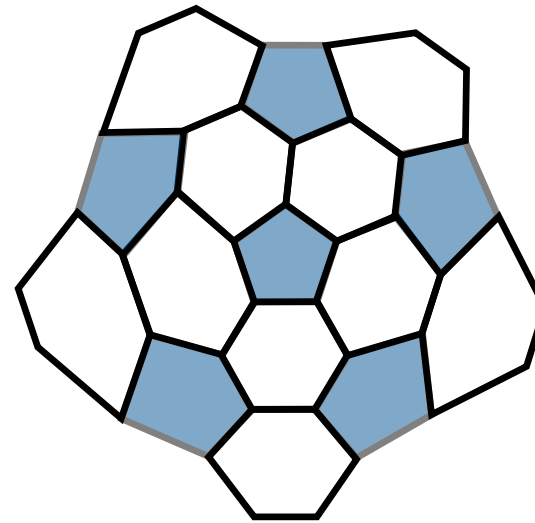
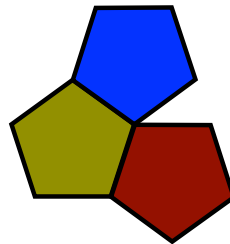


Projecting to Flat 2D

$n = 6$
flat



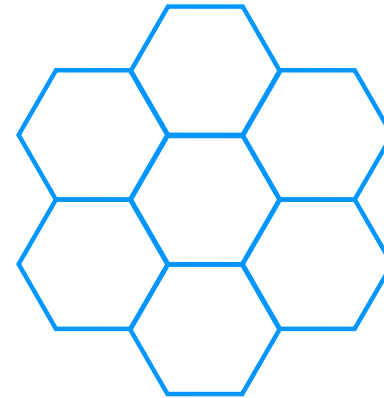
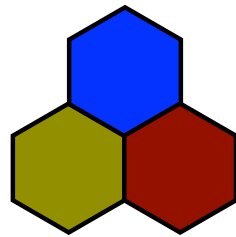
$n = 5$
spherical



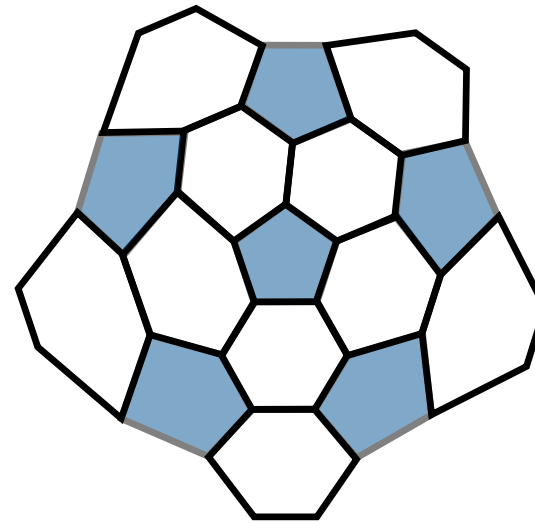
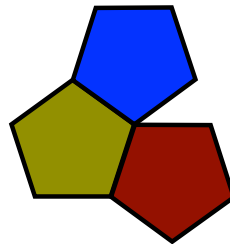
Distance is not
preserved.

Projecting to Flat 2D

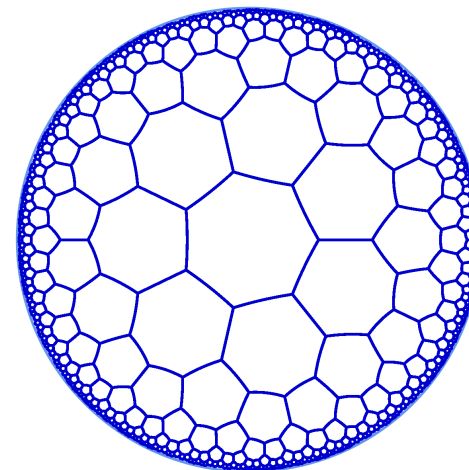
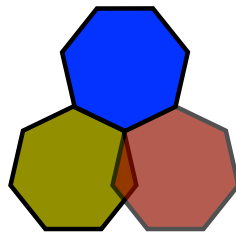
$n = 6$
flat



$n = 5$
spherical



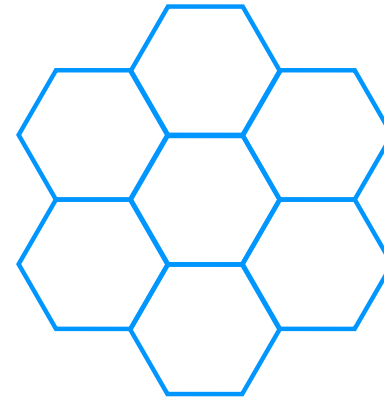
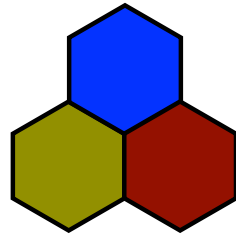
$n = 7$
hyperbolic



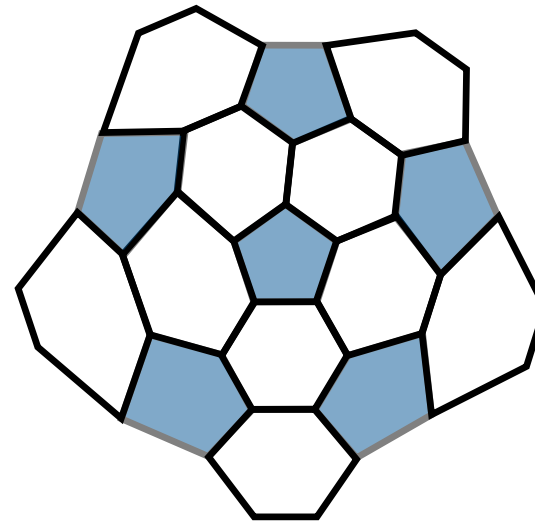
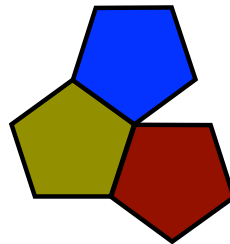
Distance is not
preserved.

Projecting to Flat 2D

$n = 6$
flat

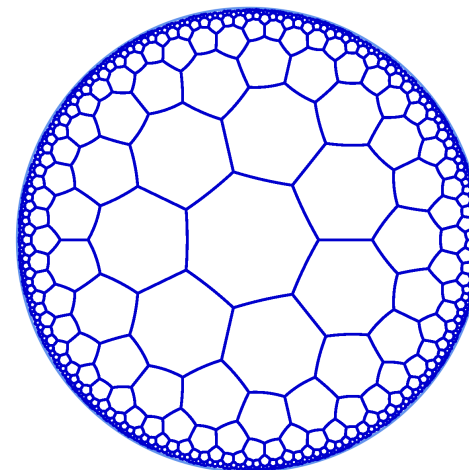
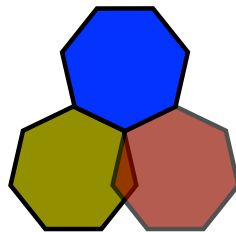


$n = 5$
spherical



Distance is not preserved.

$n = 7$
hyperbolic

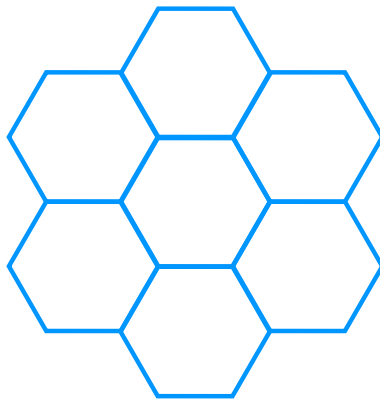
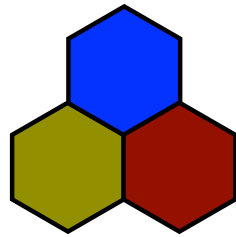


Distance is not preserved.

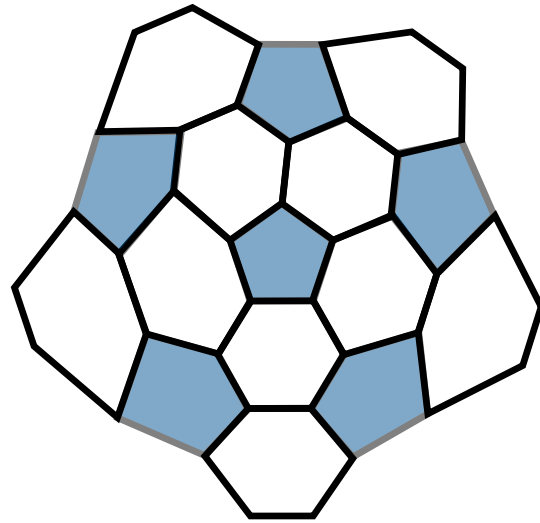
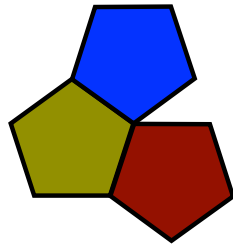
Projecting to Flat 2D

$n = 6$

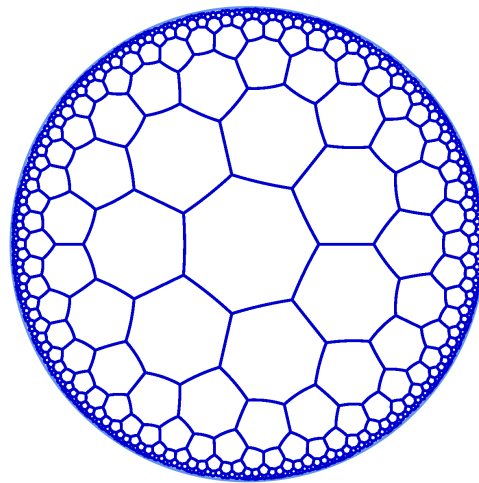
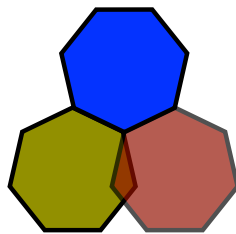
flat



$n = 5$
spherical



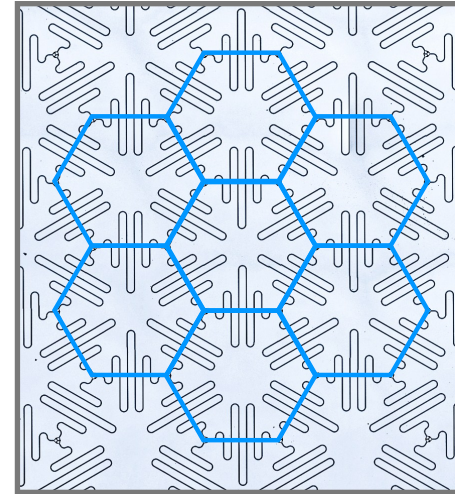
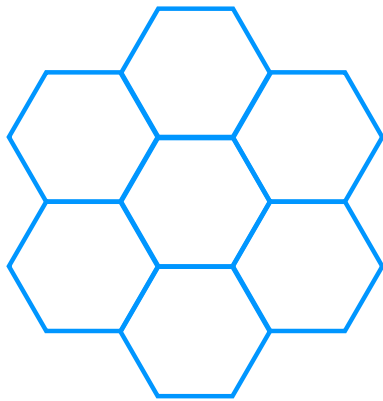
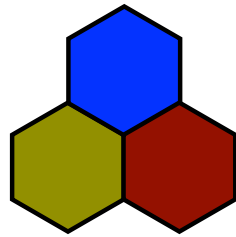
$n = 7$
hyperbolic



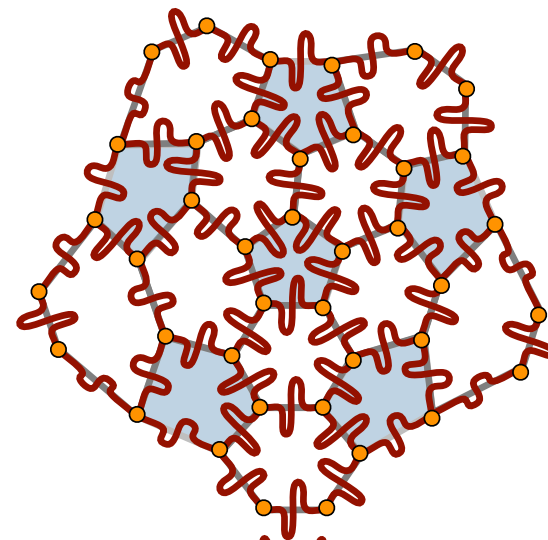
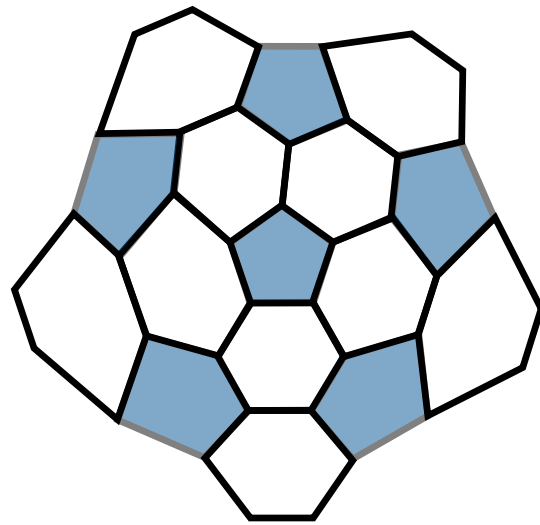
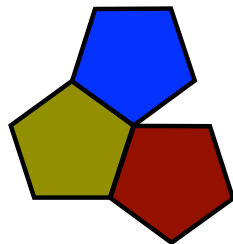
Projecting to Flat 2D

$n = 6$

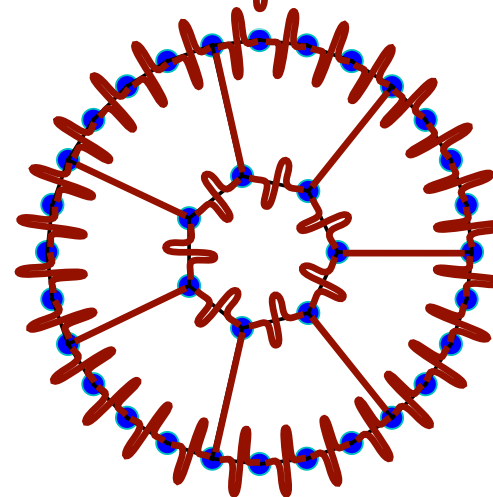
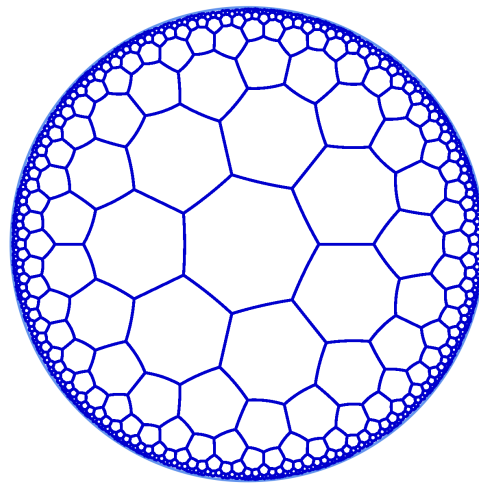
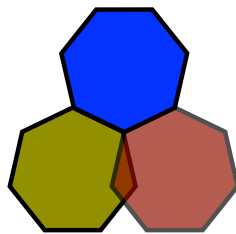
flat



$n = 5$
spherical

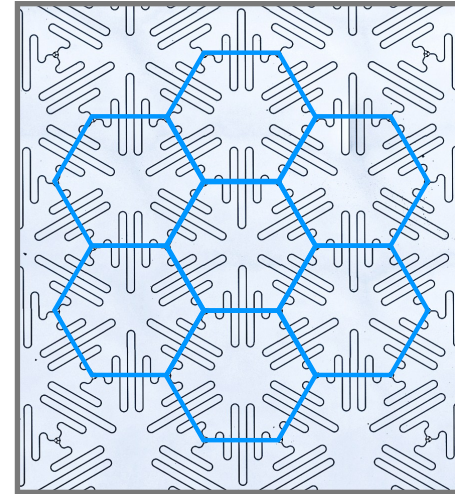
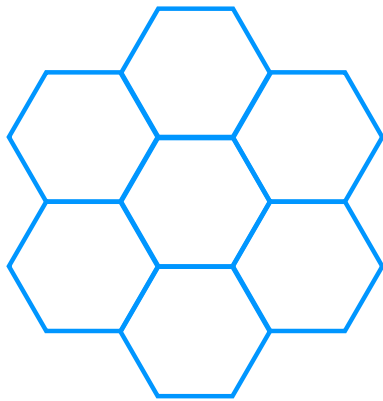
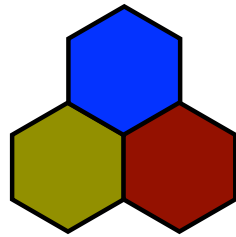


$n = 7$
hyperbolic

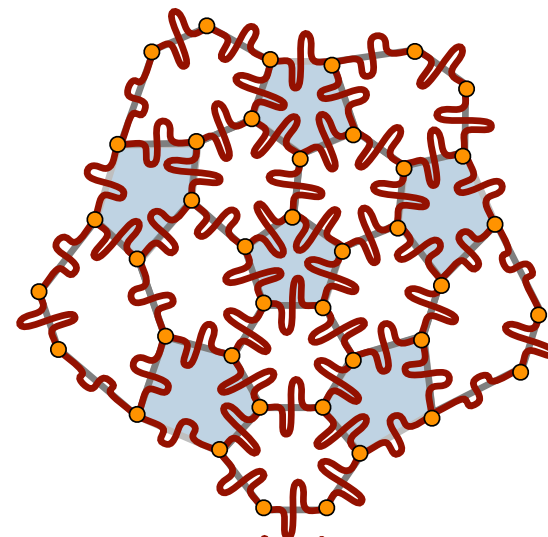
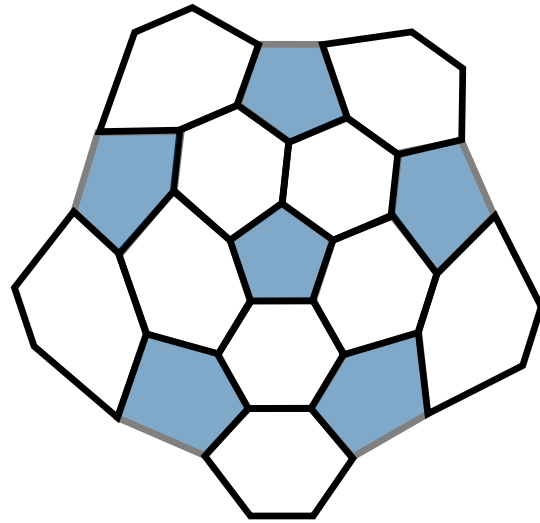
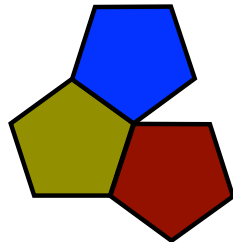


Projecting to Flat 2D

$n = 6$
flat

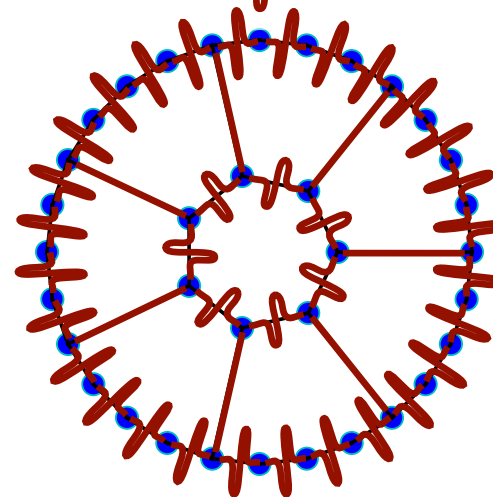
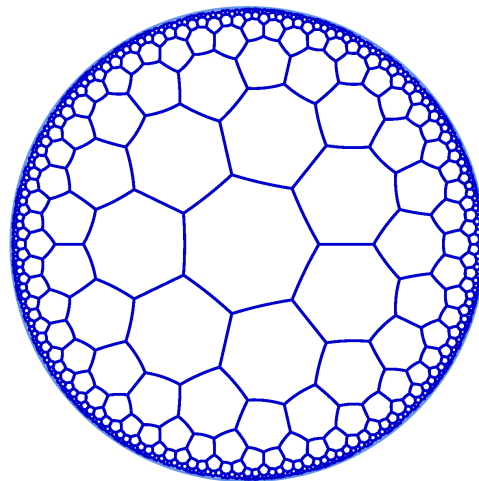
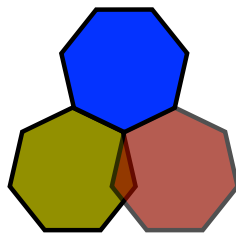


$n = 5$
spherical



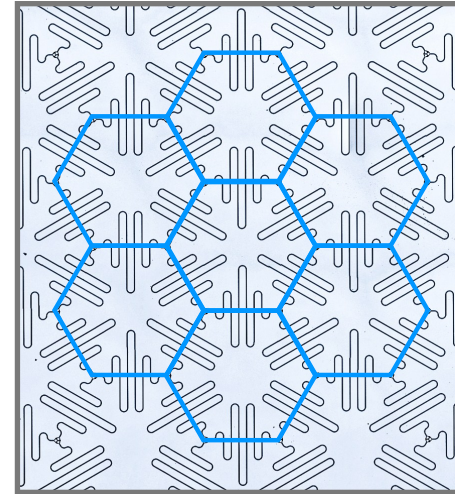
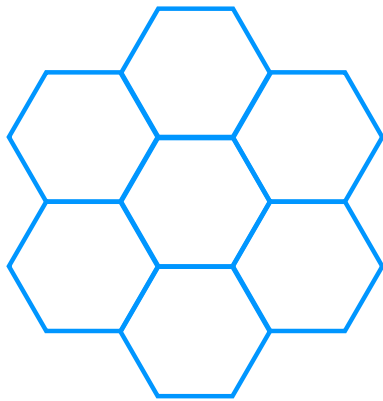
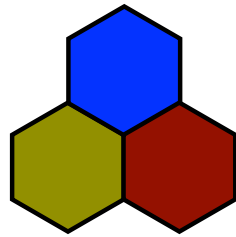
Distance is not preserved.

$n = 7$
hyperbolic

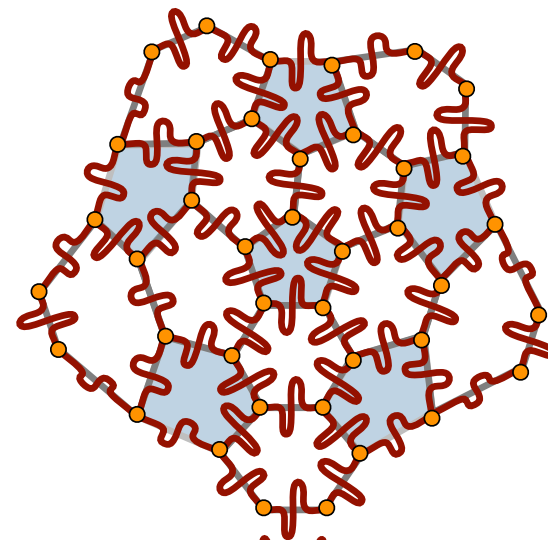
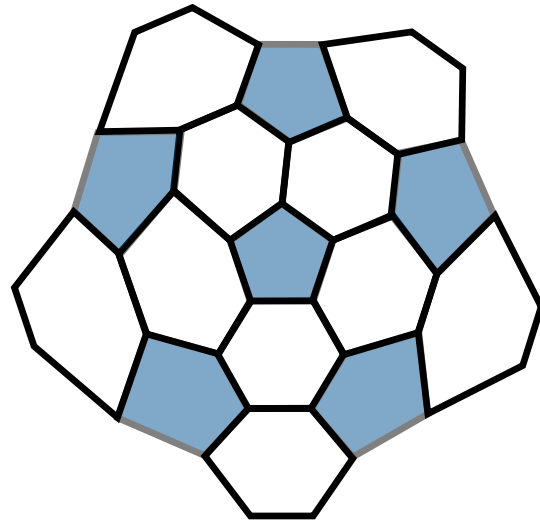
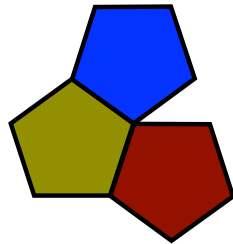


Projecting to Flat 2D

$n = 6$
flat

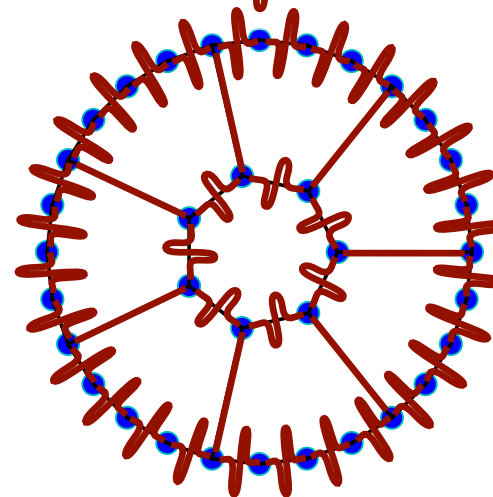
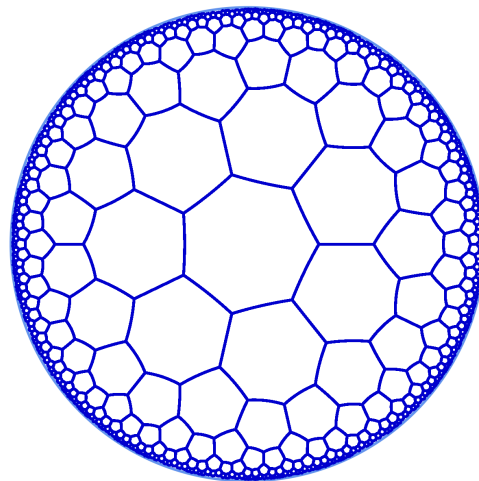
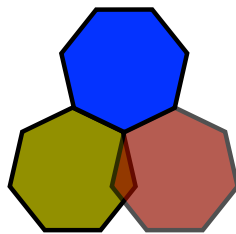


$n = 5$
spherical



Distance is not
preserved.

$n = 7$
hyperbolic

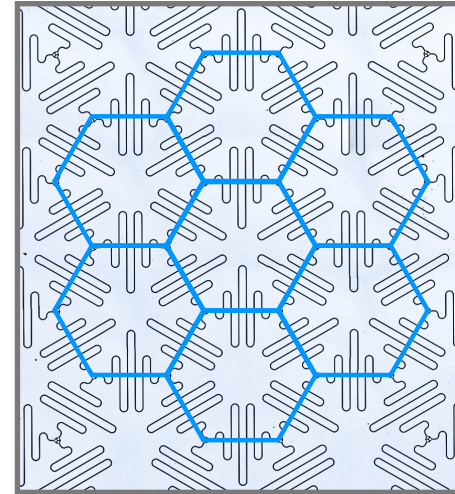
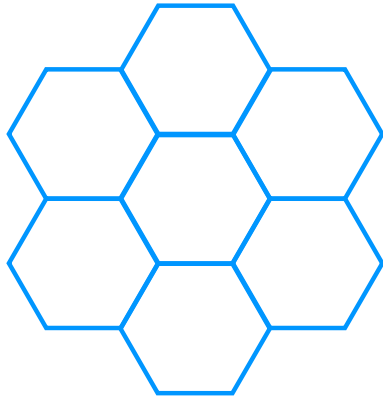
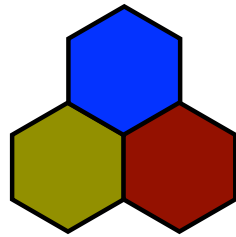


t is preserved.

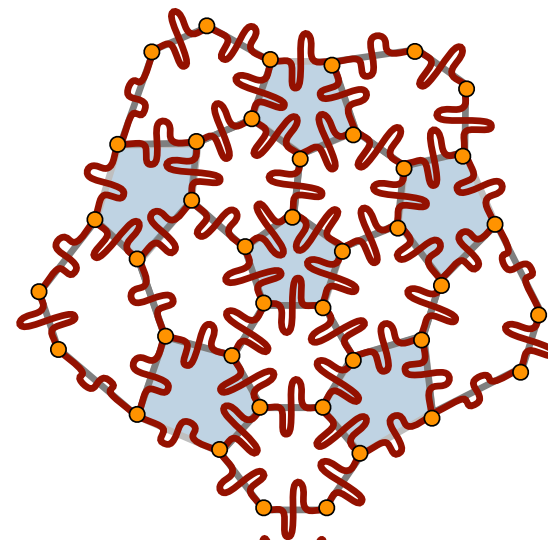
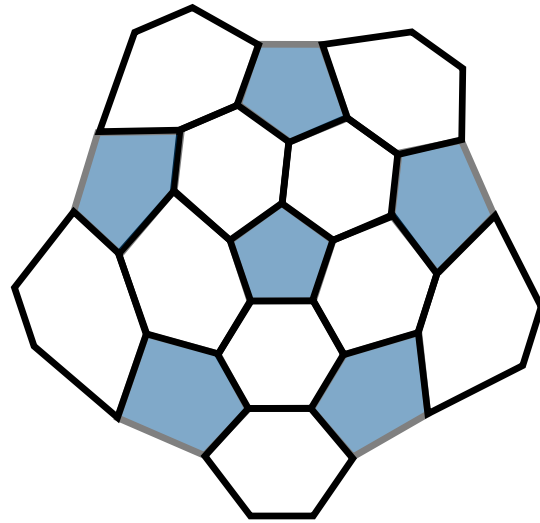
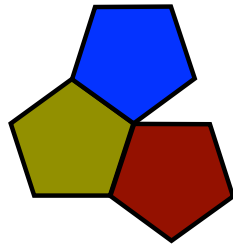
Projecting to Flat 2D

$n = 6$

flat



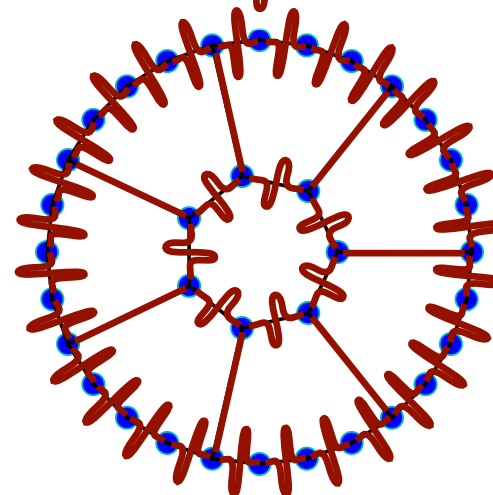
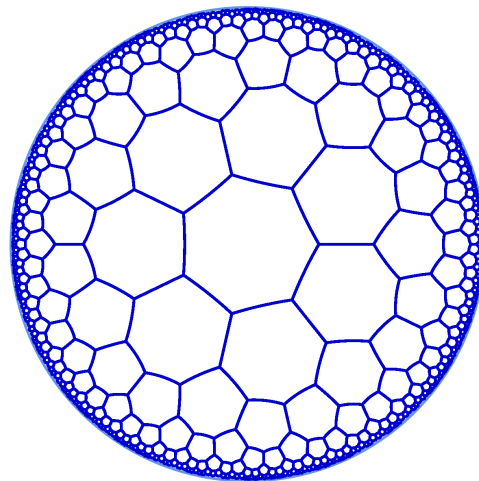
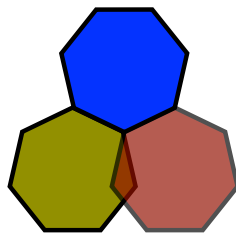
$n = 5$
spherical



Distance is not
preserved.

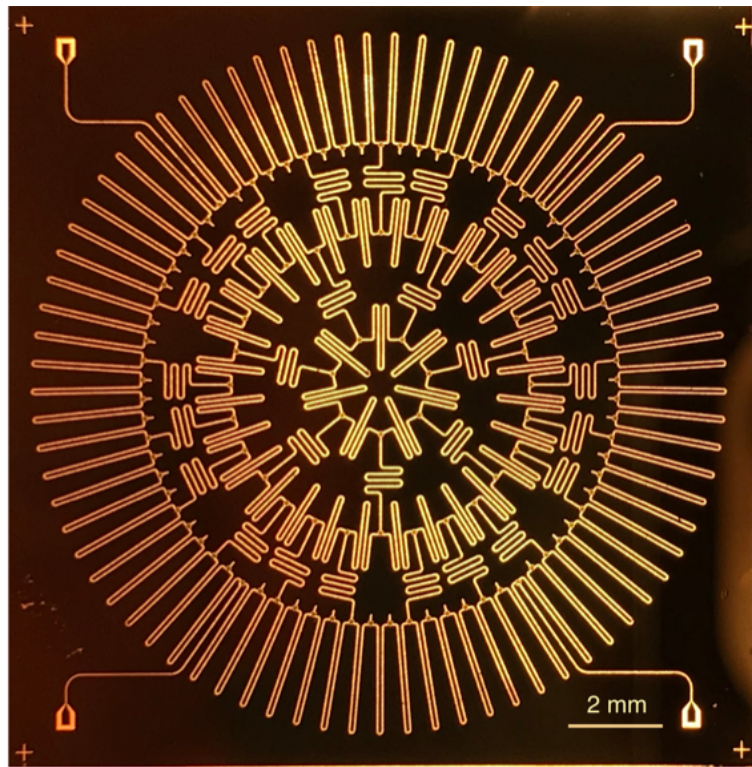
t *is* preserved.

$n = 7$
hyperbolic

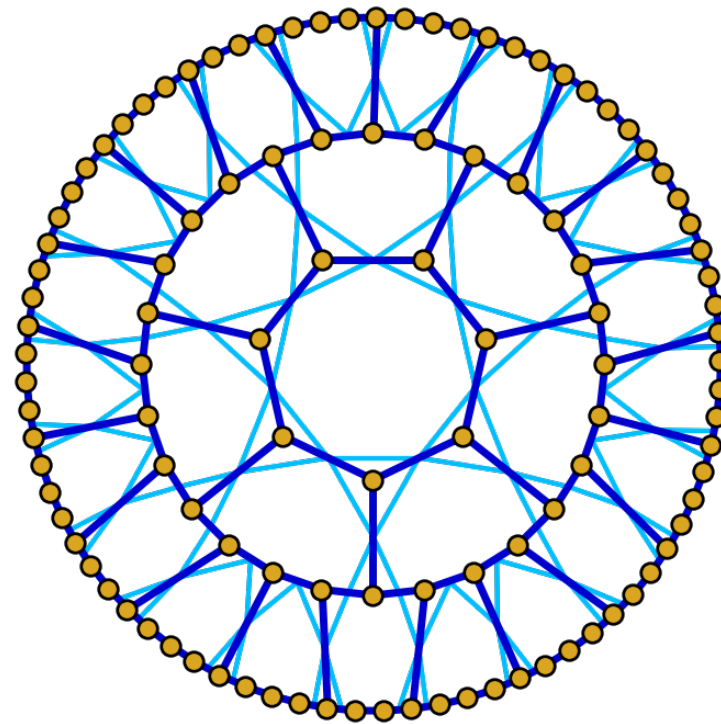
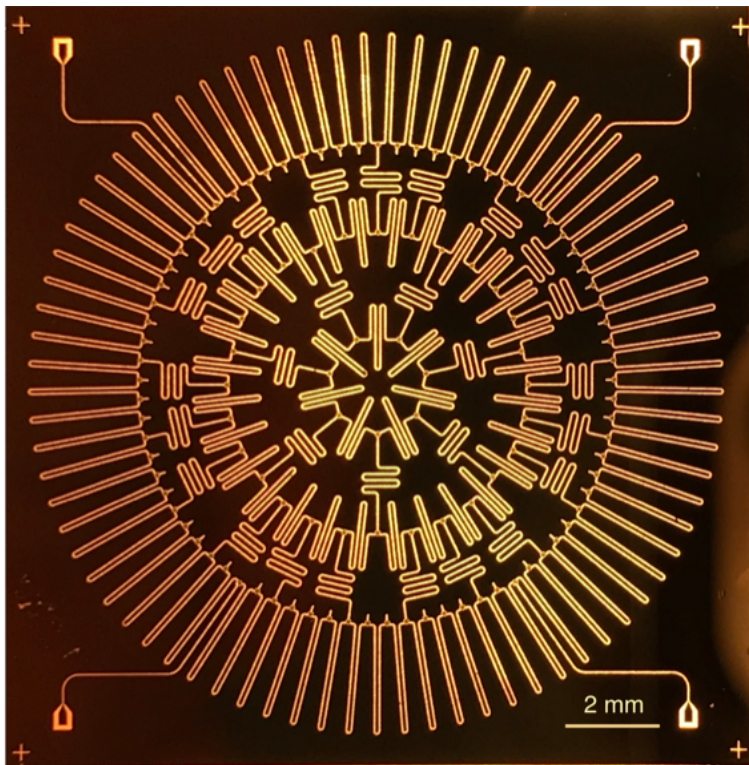


Graph *is* preserved.

Heptagon-Kagome Device

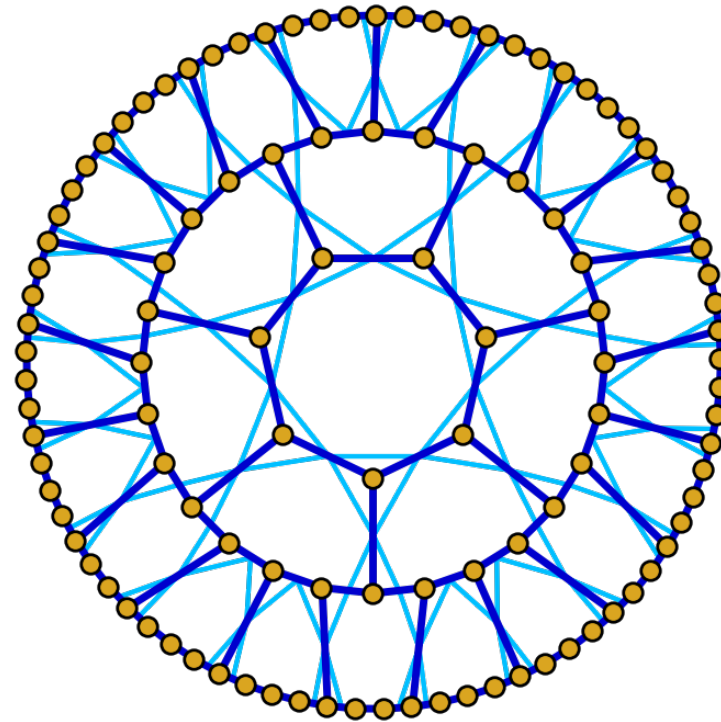
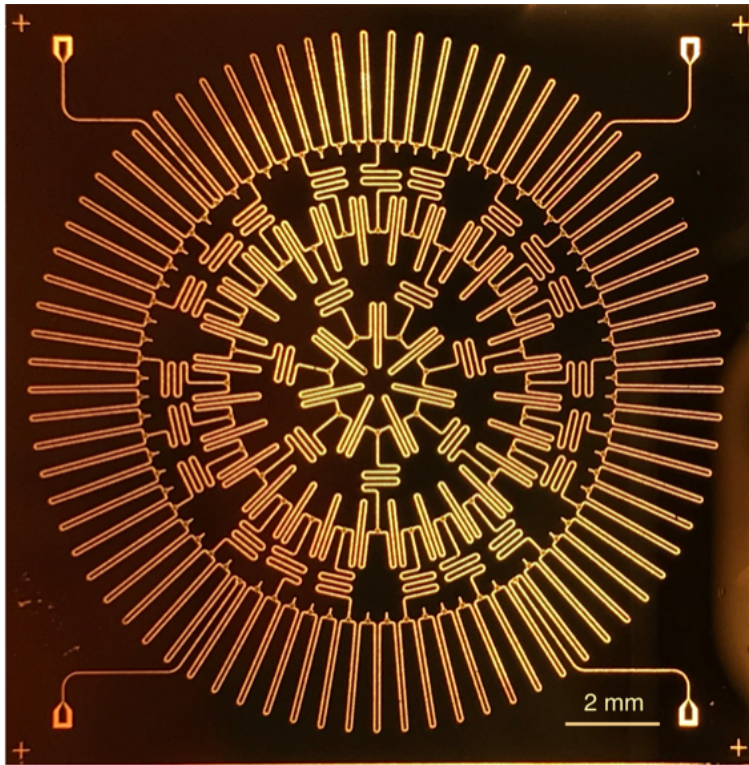


Heptagon-Kagome Device



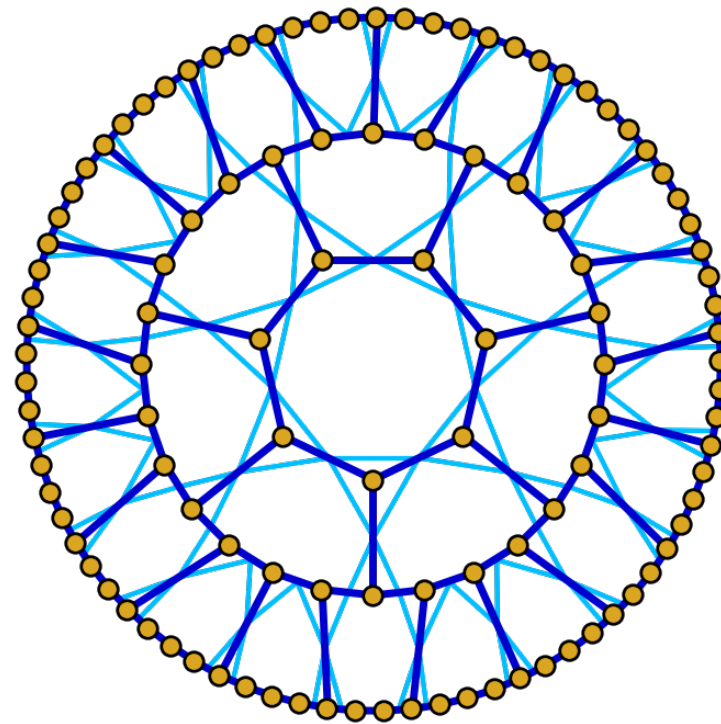
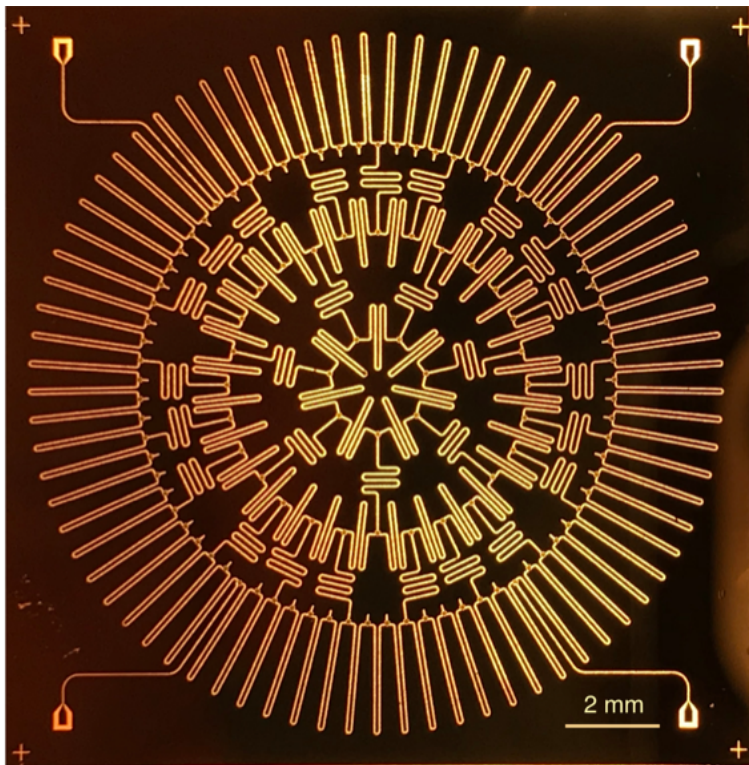
- 2 shells

Heptagon-Kagome Device



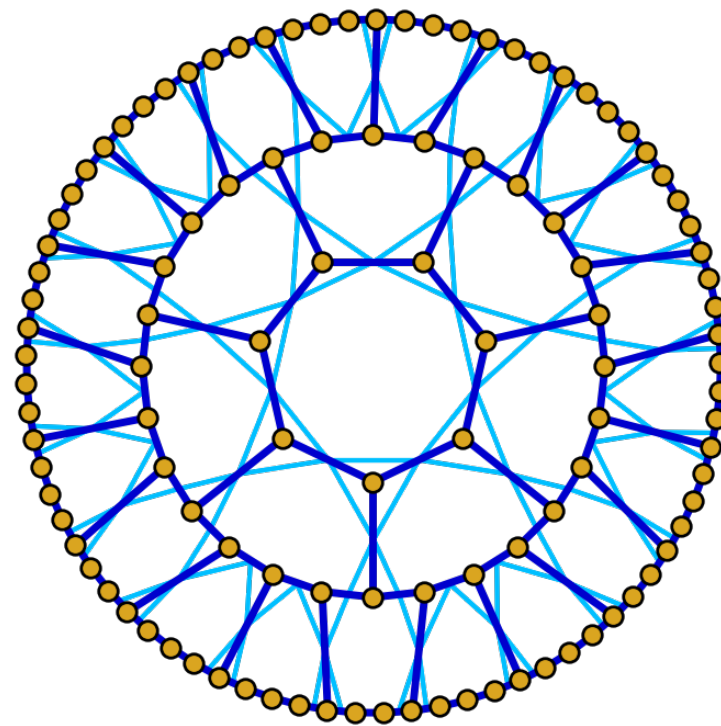
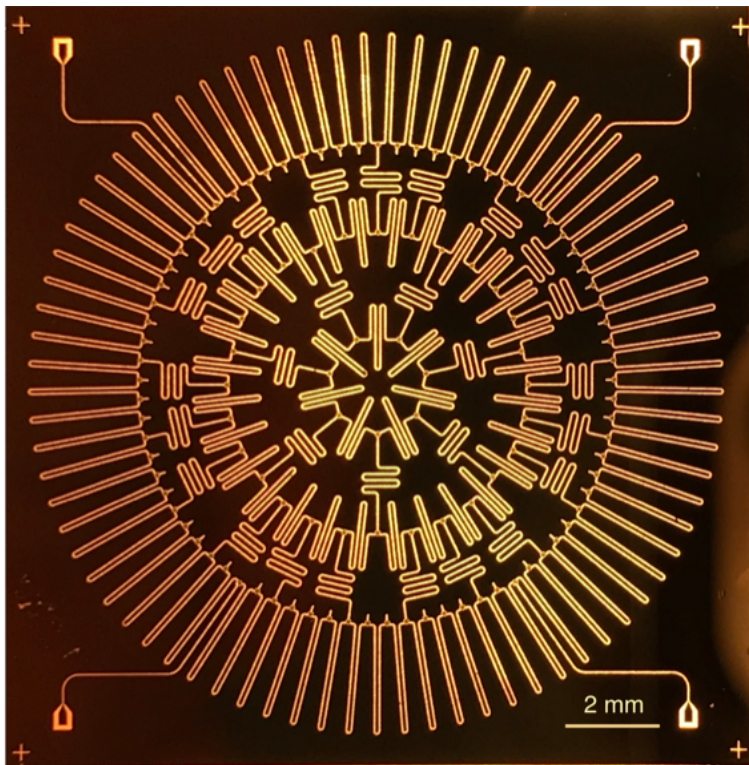
- 2 shells
- Operating frequency: 16 GHz

Heptagon-Kagome Device

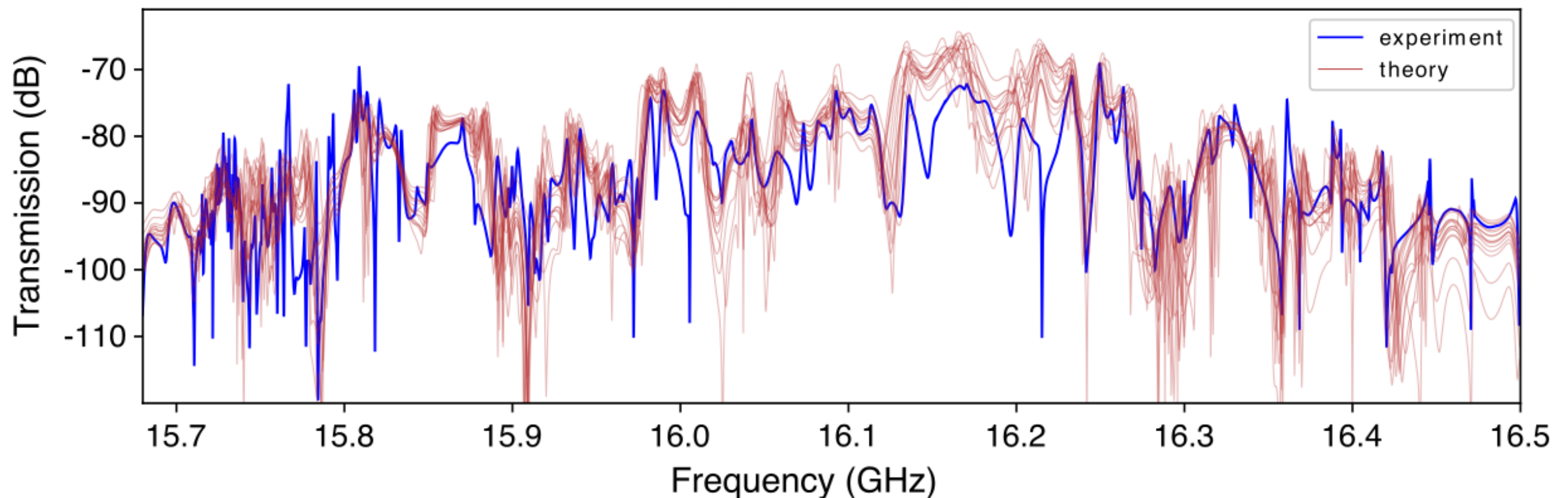


- 2 shells
- Operating frequency: 16 GHz
- 4 input-output ports

Heptagon-Kagome Device

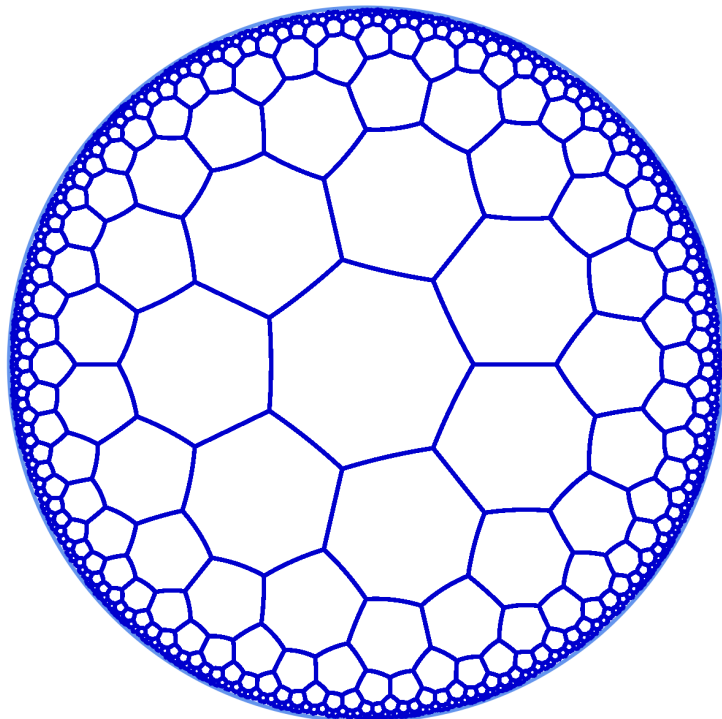
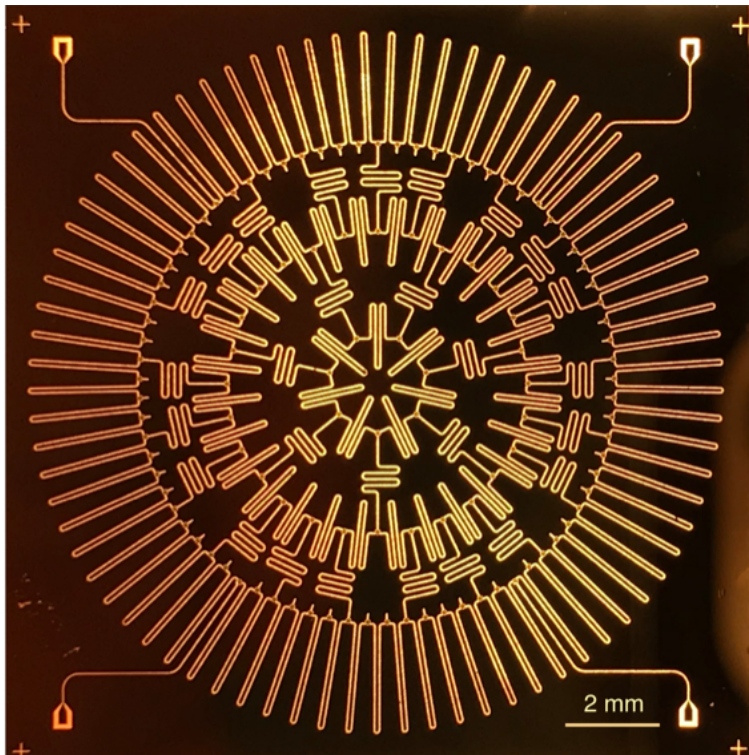


- 2 shells
- Operating frequency: 16 GHz
- 4 input-output ports



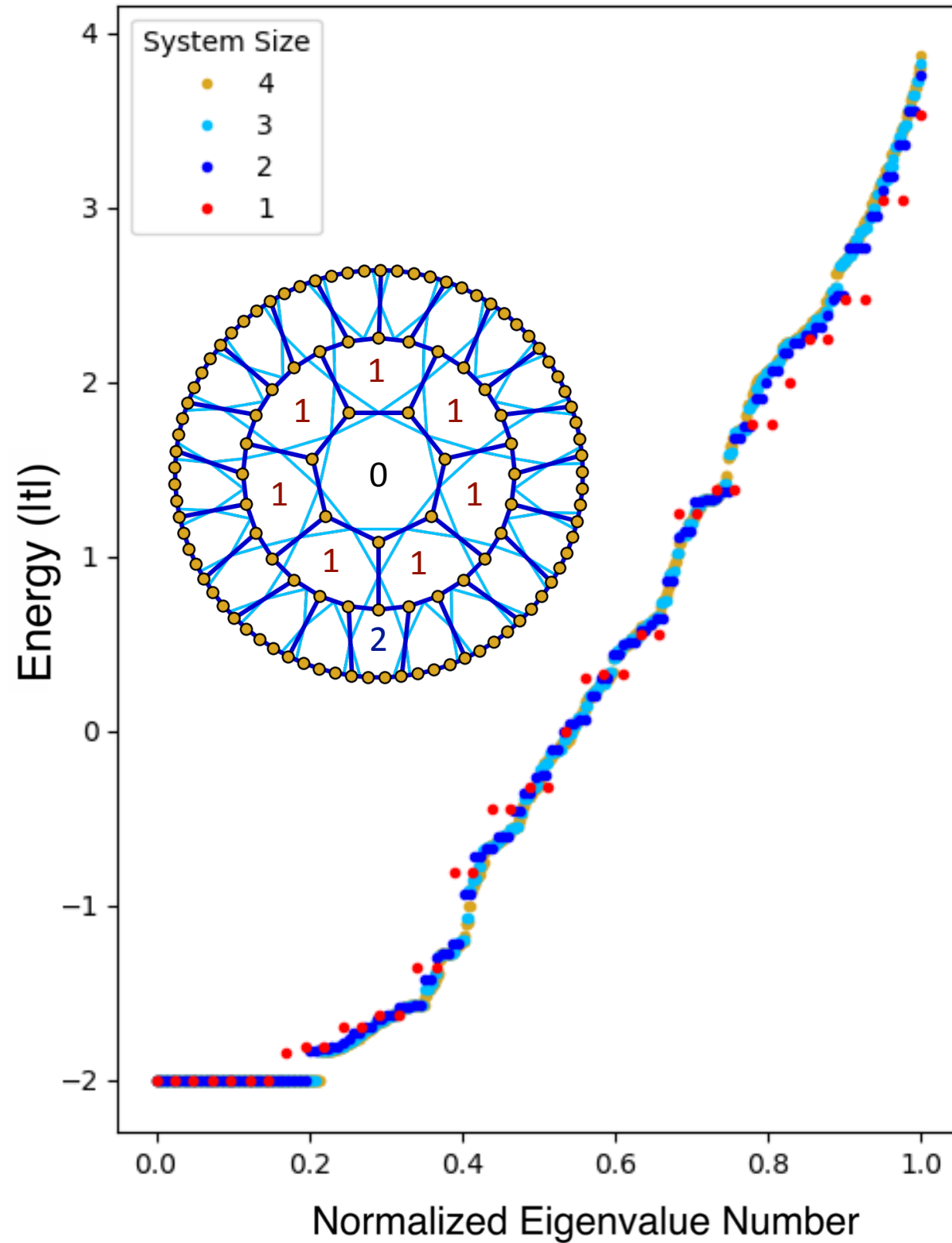
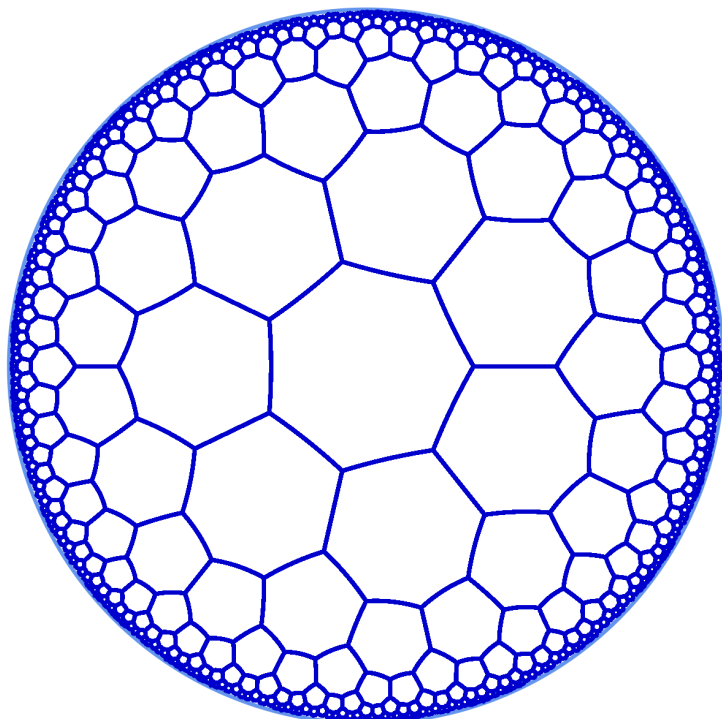
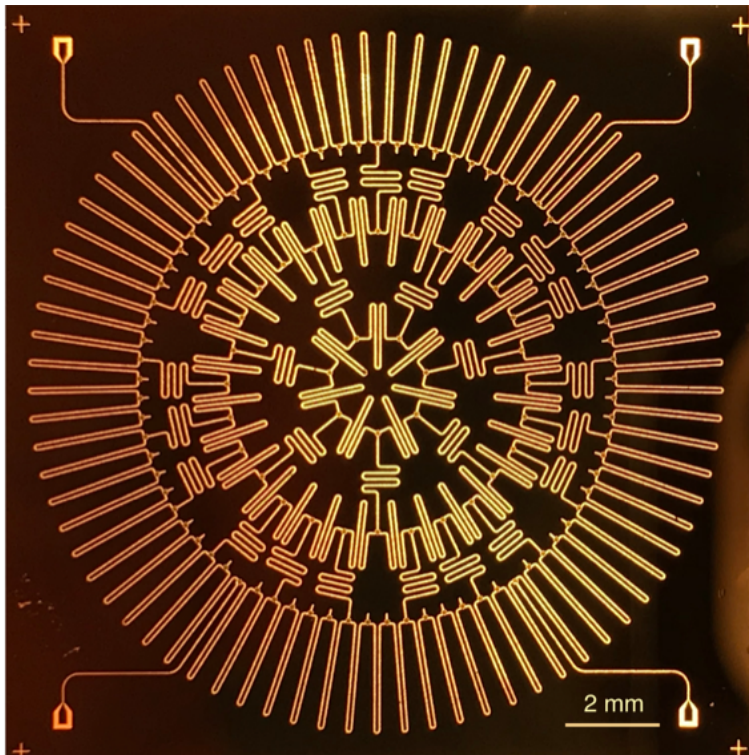
Spectrum Calculations

What is the spectrum of this?



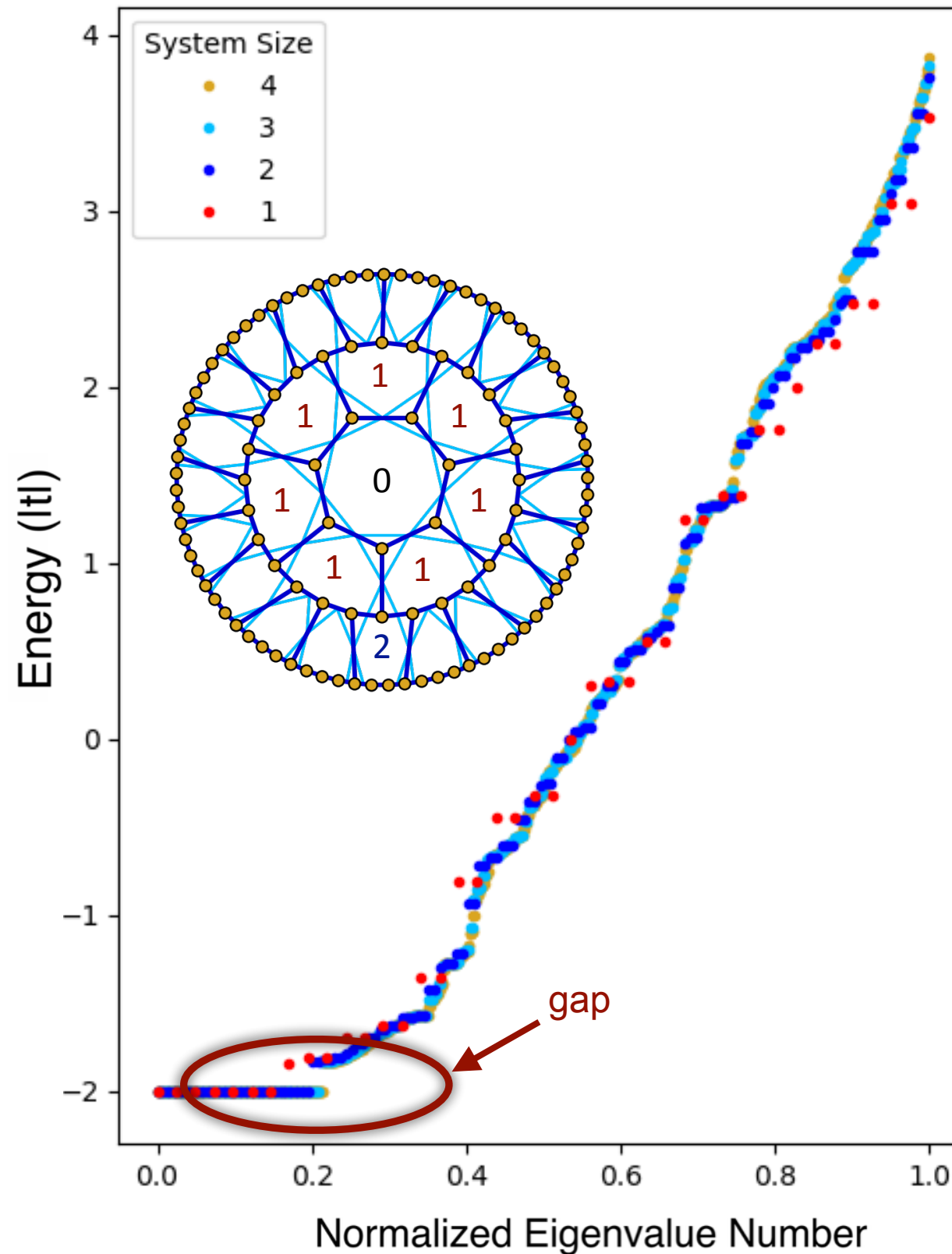
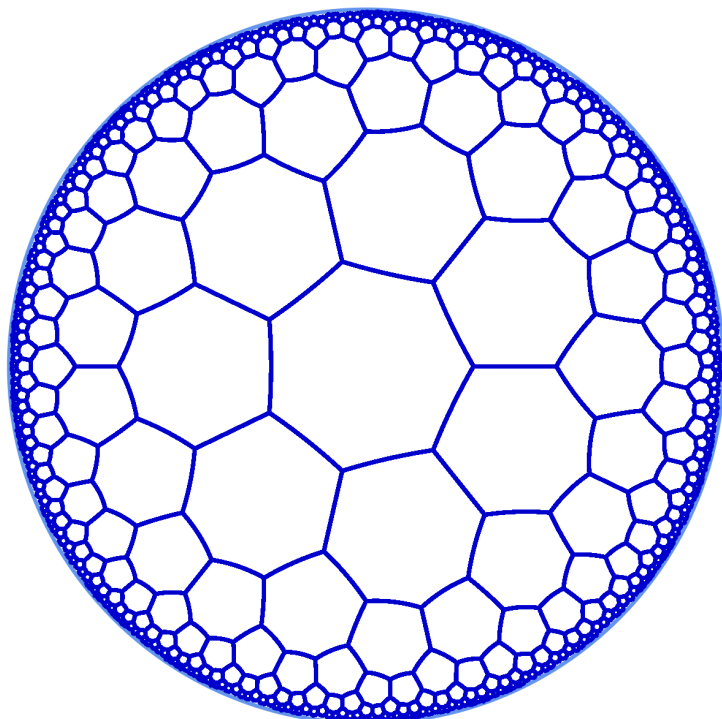
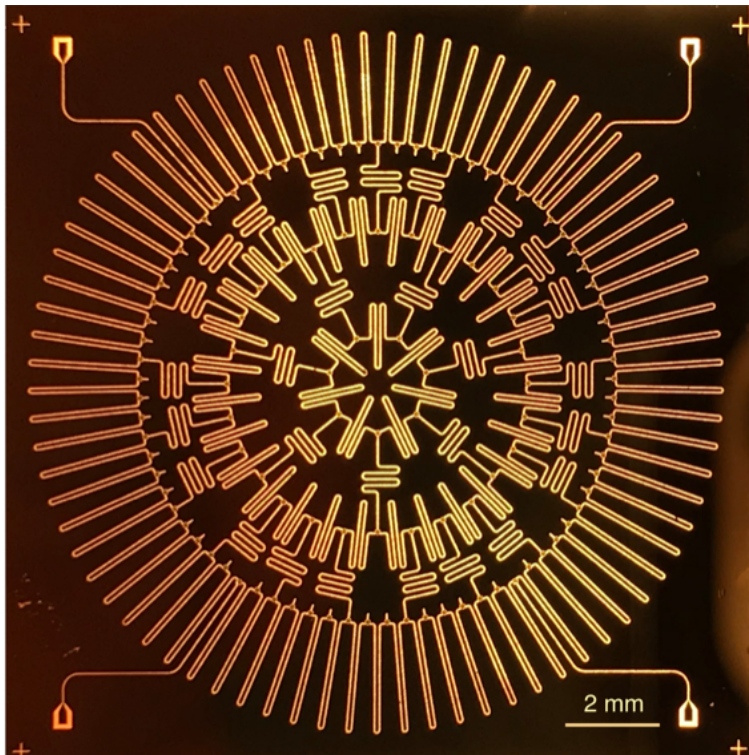
Spectrum Calculations

What is the spectrum of this?



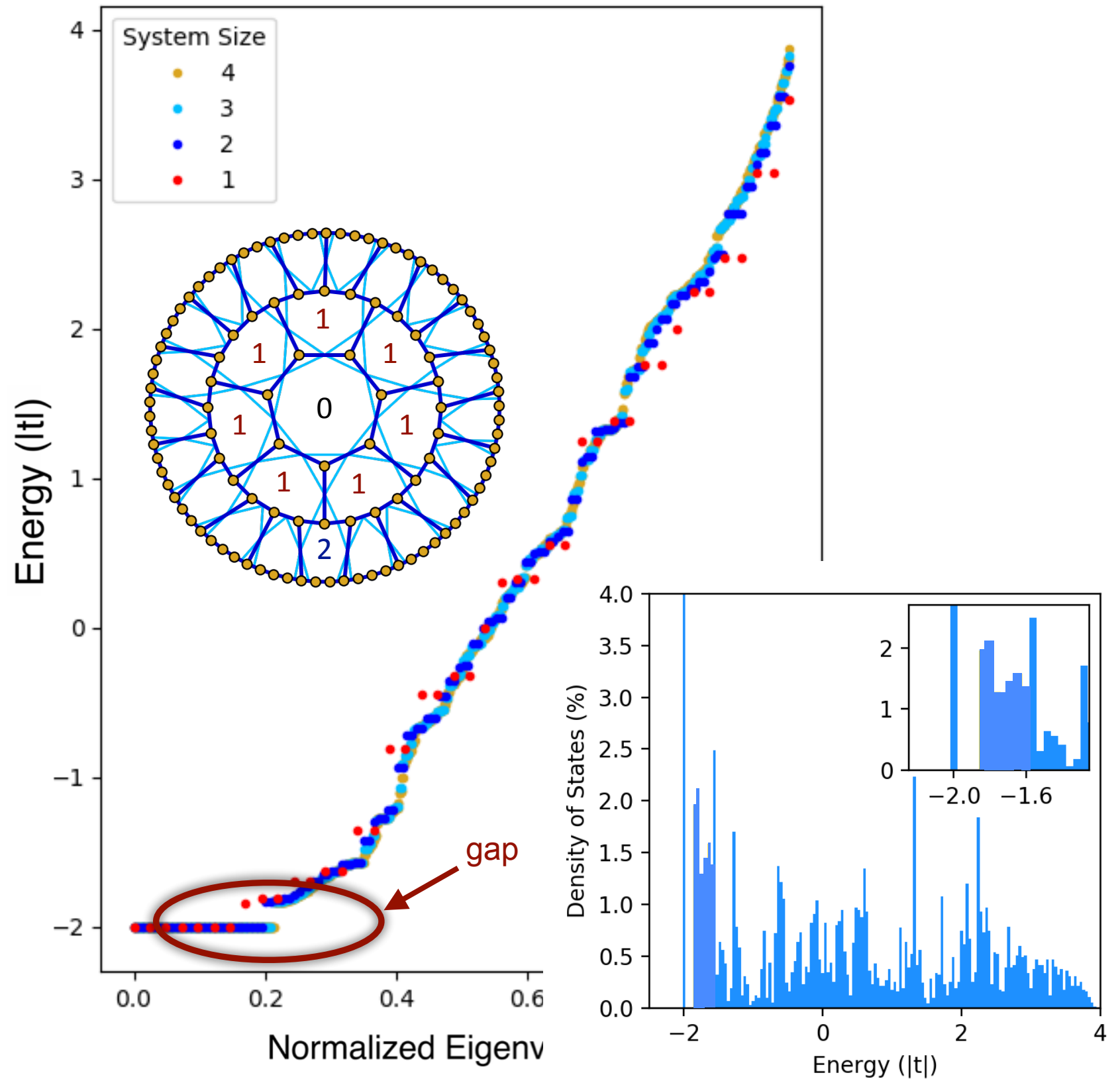
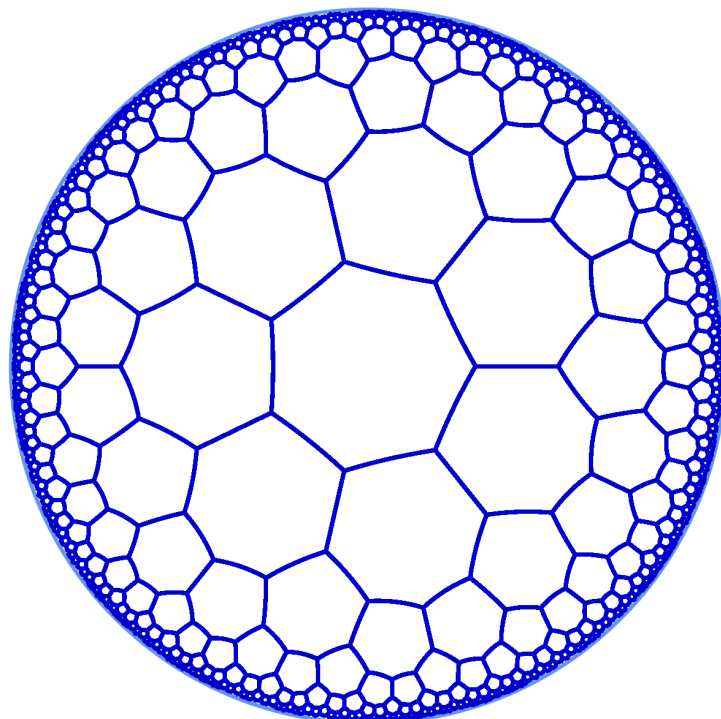
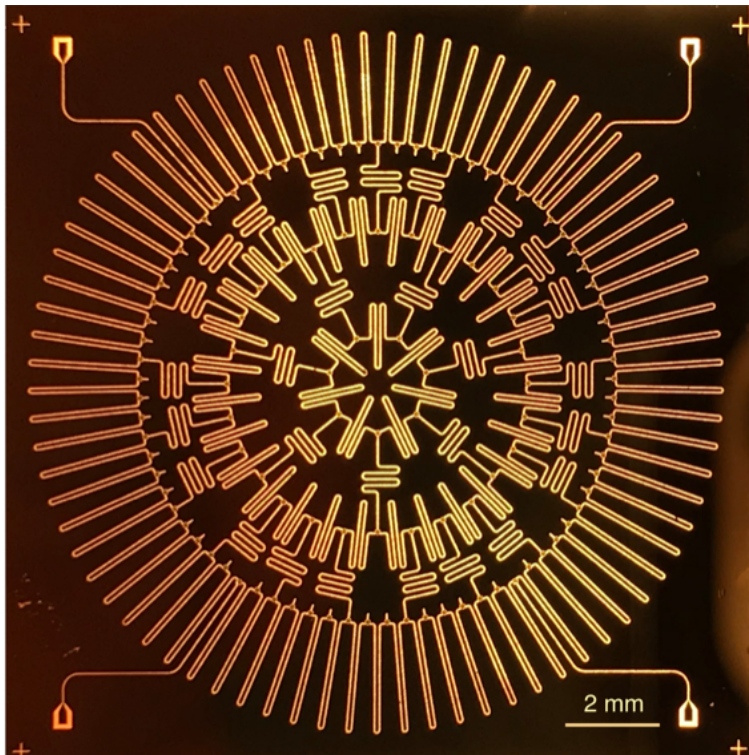
Spectrum Calculations

What is the spectrum of this?



Spectrum Calculations

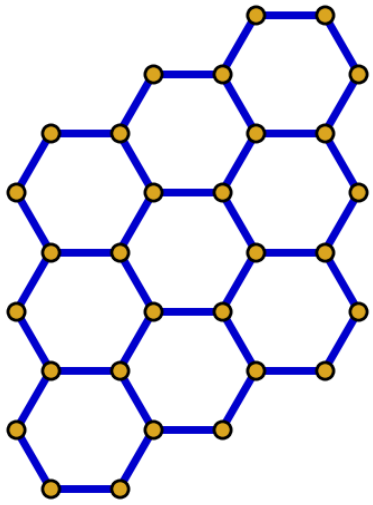
What is the spectrum of this?



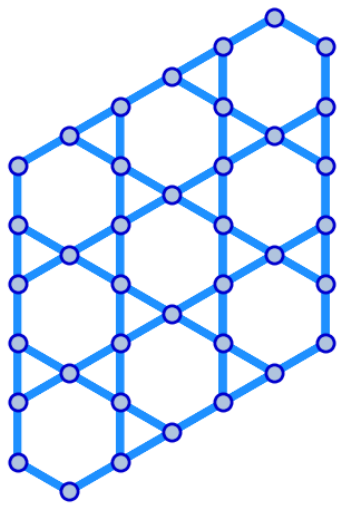
Line-Graph Lattices

Graphene

Layout X



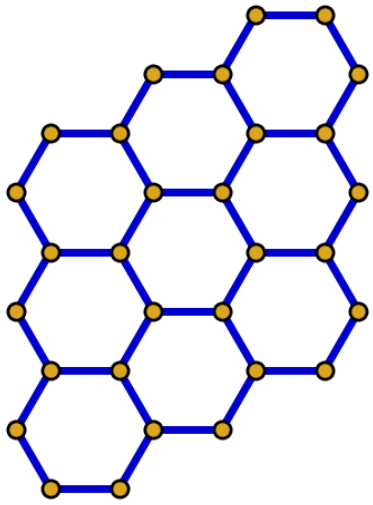
Line Graph $L(X)$



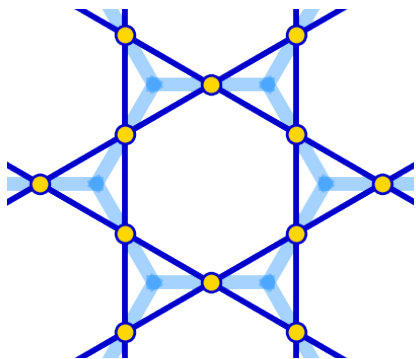
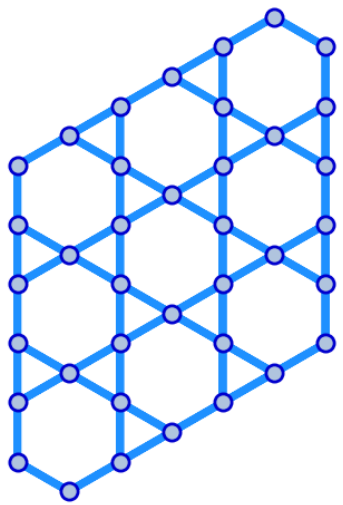
Line-Graph Lattices

Graphene

Layout X



Line Graph $L(X)$

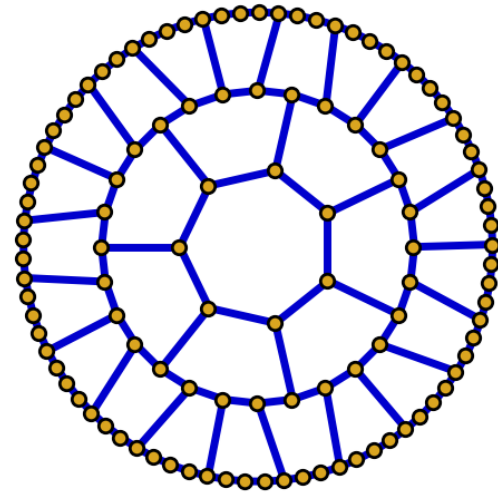
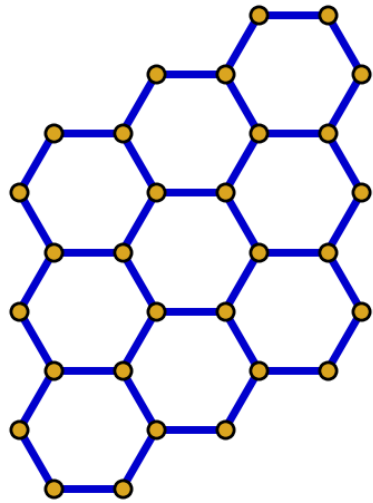


Line-Graph Lattices

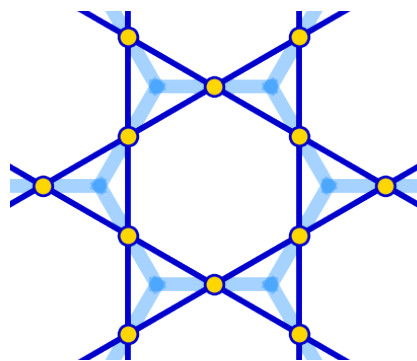
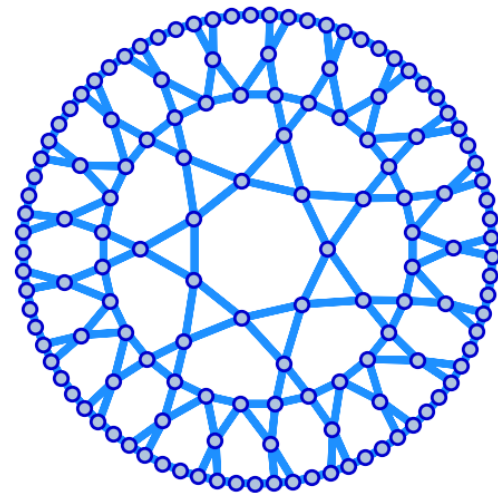
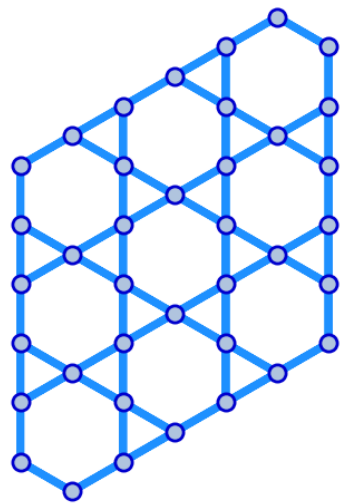
Graphene

Heptagon-Graphene

Layout X



Line Graph $L(X)$



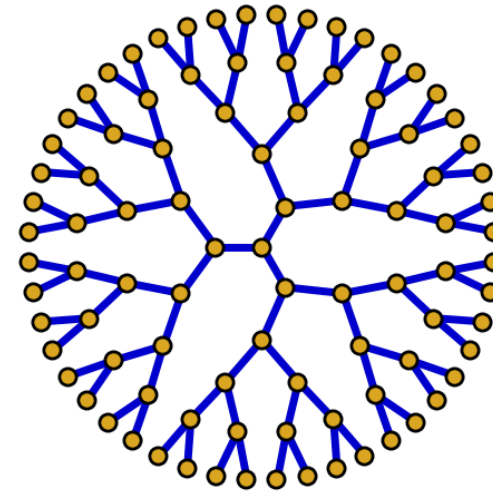
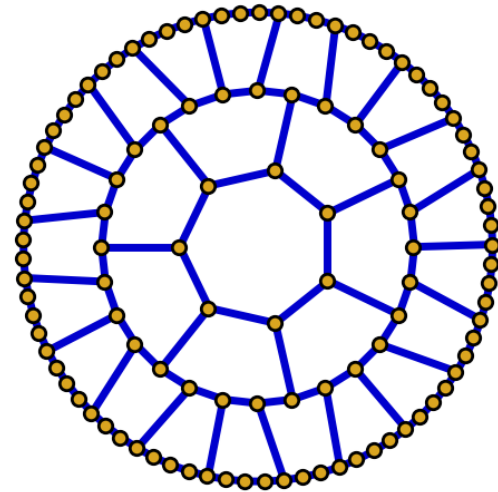
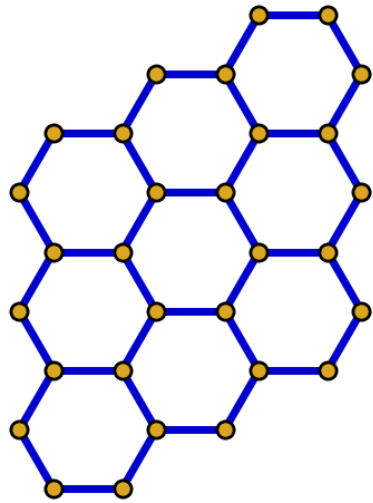
Line-Graph Lattices

Graphene

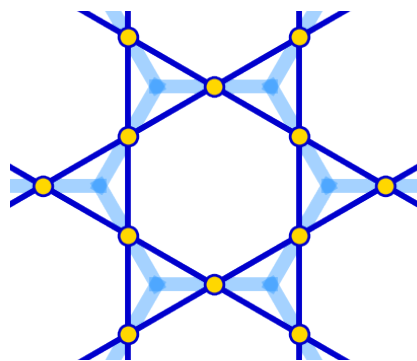
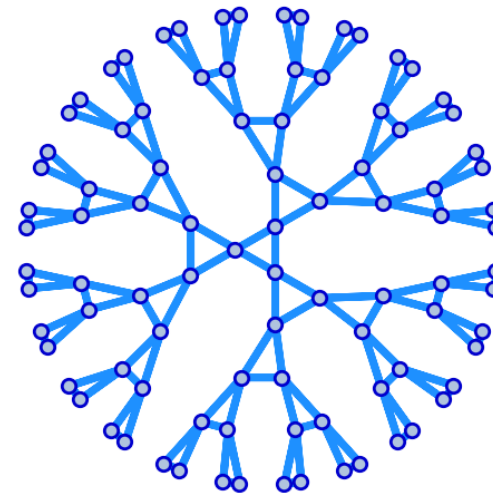
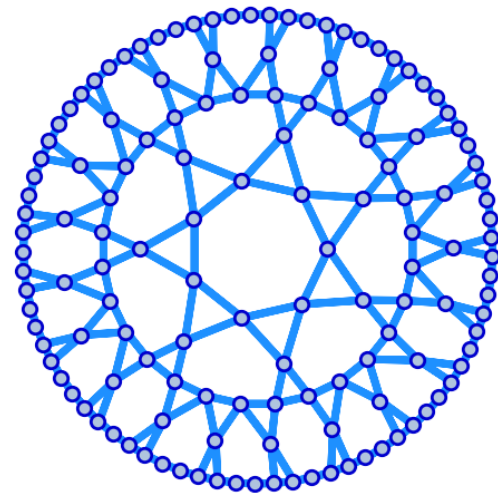
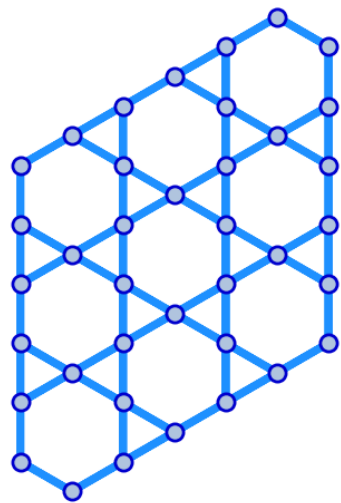
Heptagon-Graphene

Tree

Layout X



Line Graph $L(X)$



Line-Graph Lattices

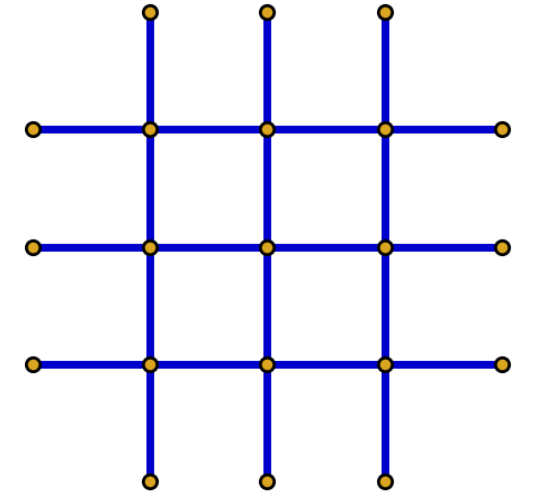
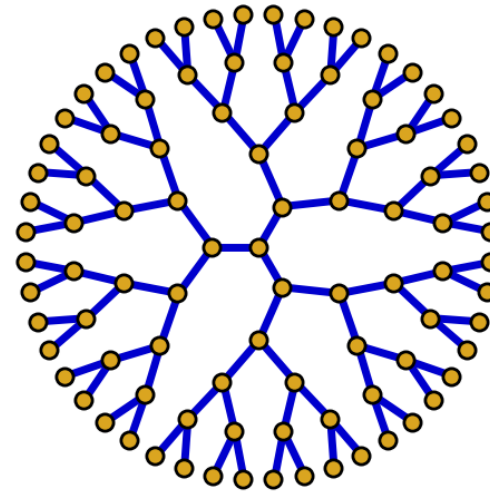
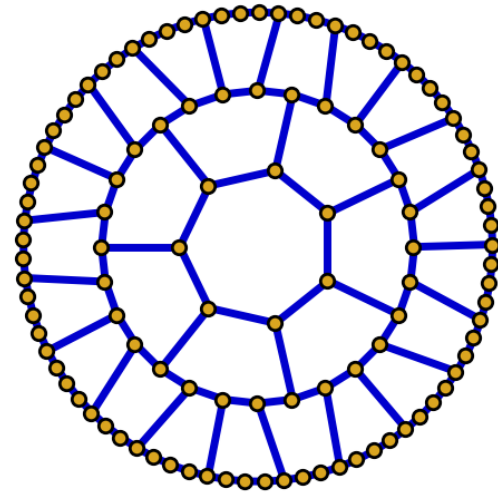
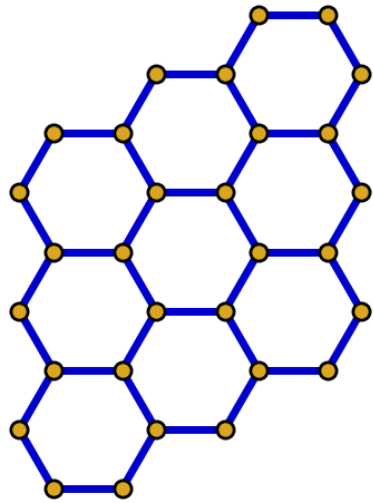
Graphene

Heptagon-Graphene

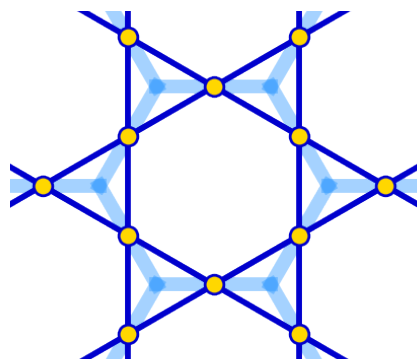
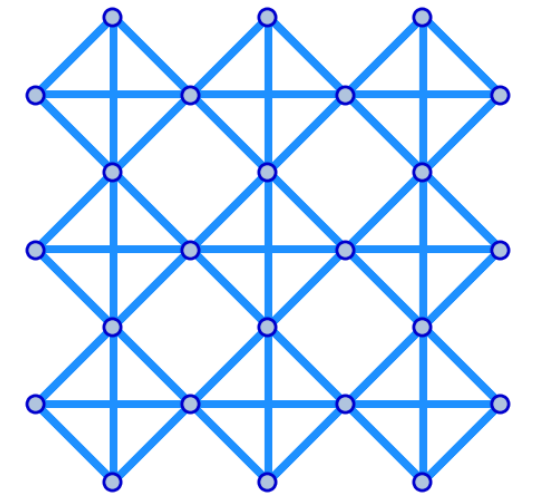
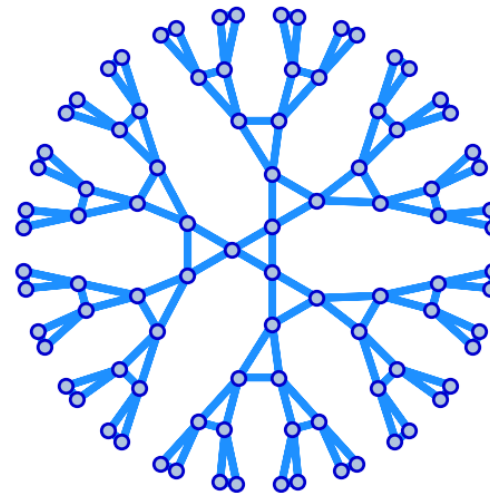
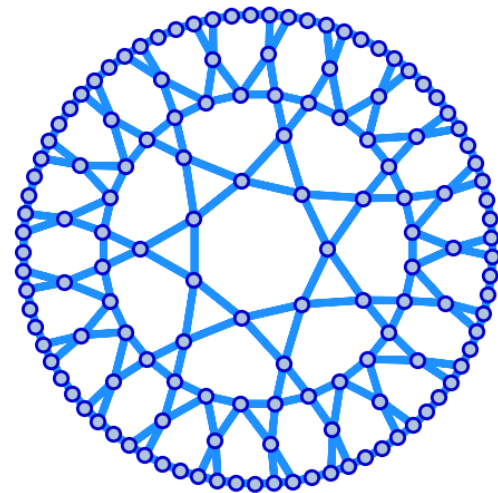
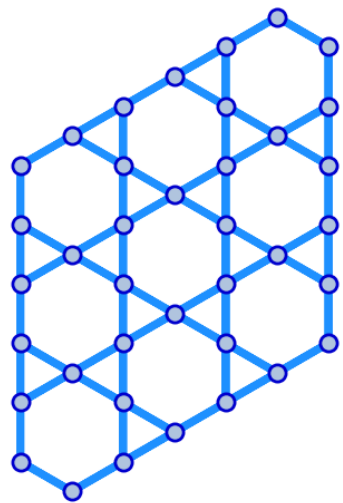
Tree

Square

Layout X

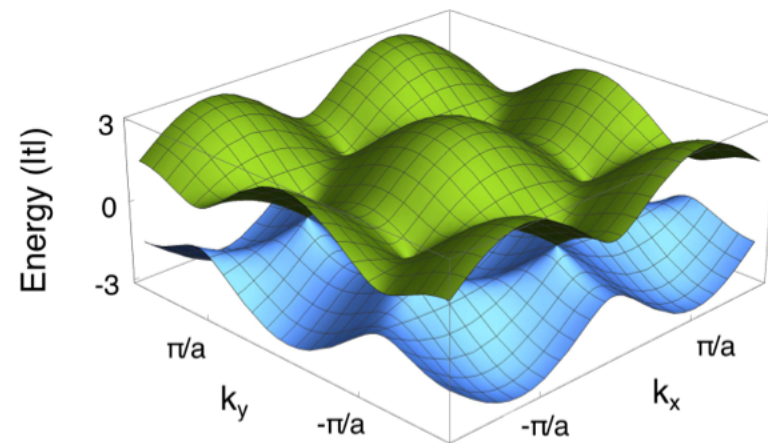
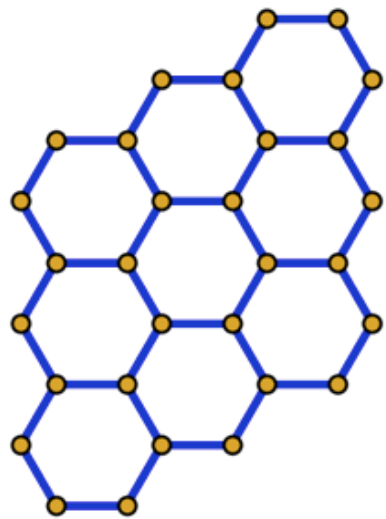


Line Graph $L(X)$



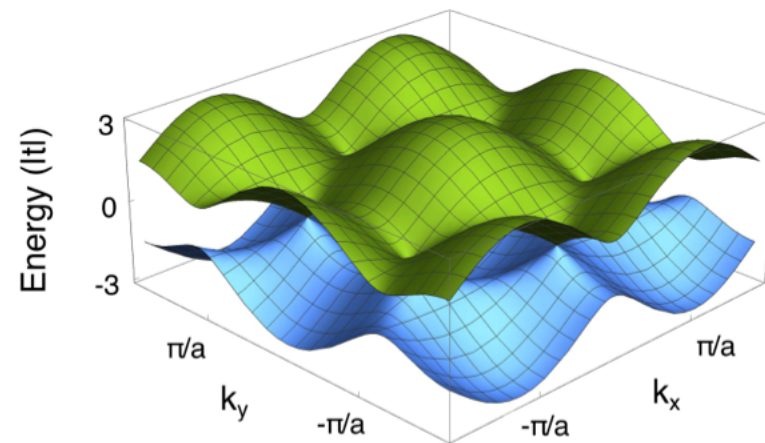
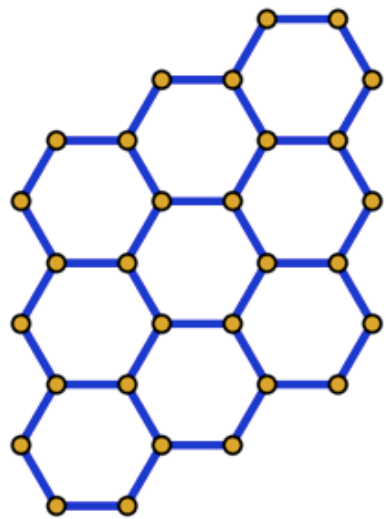
Band Structure Correspondence

Layout X

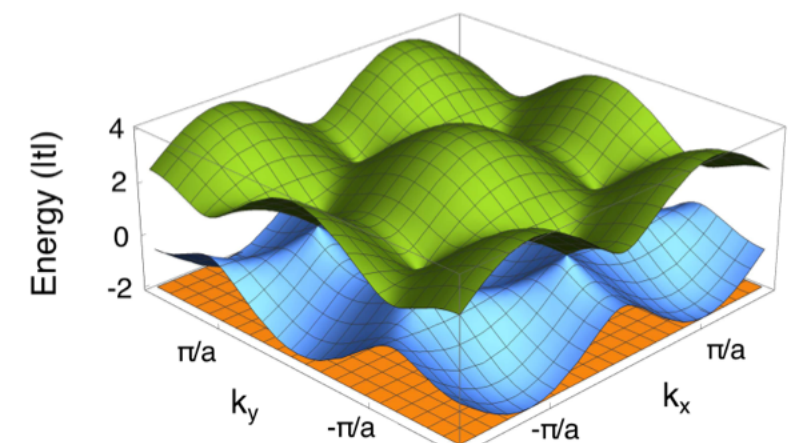
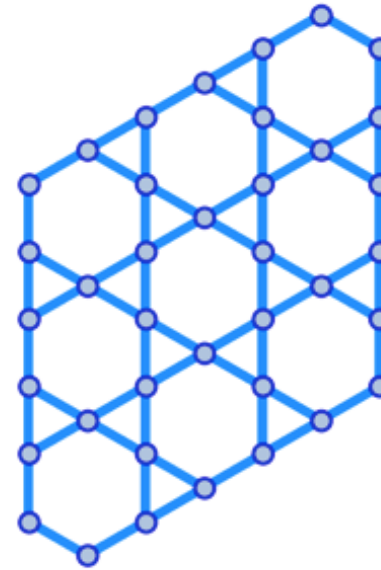


Band Structure Correspondence

Layout X

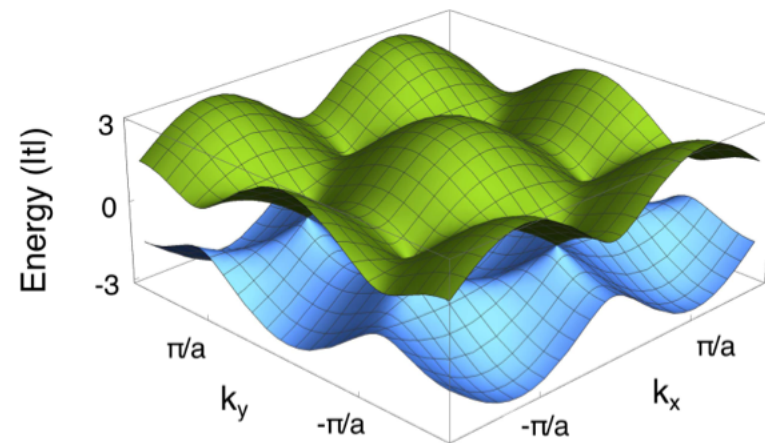
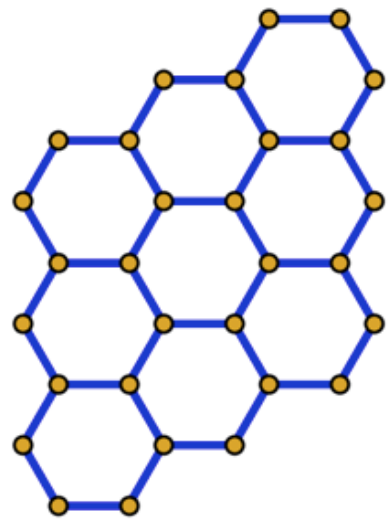


Line Graph $L(X)$

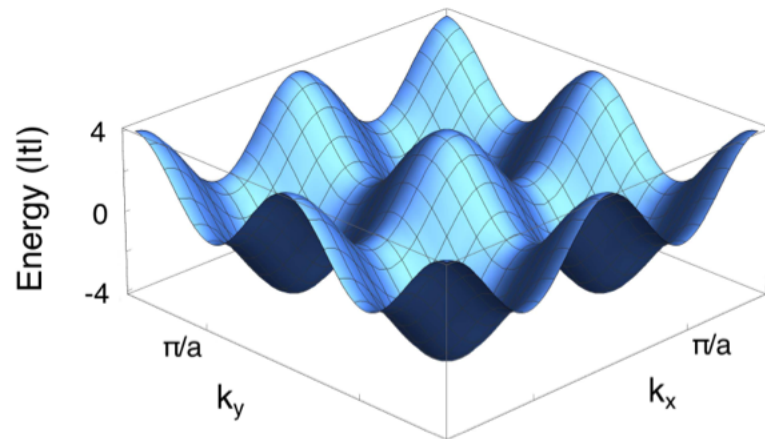
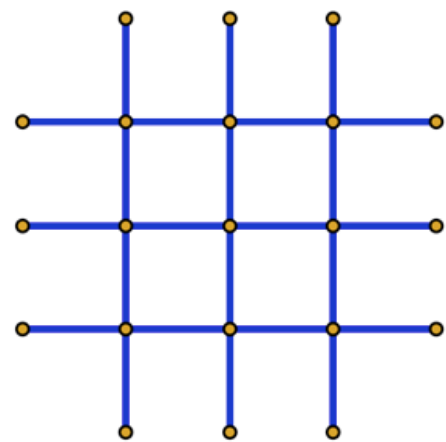
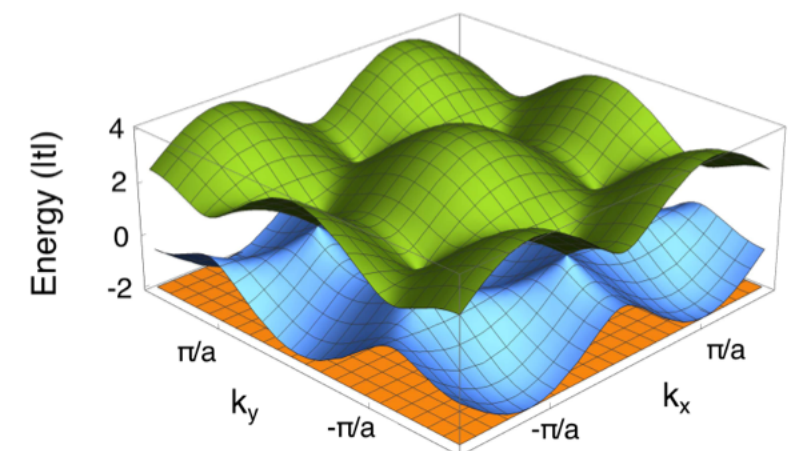
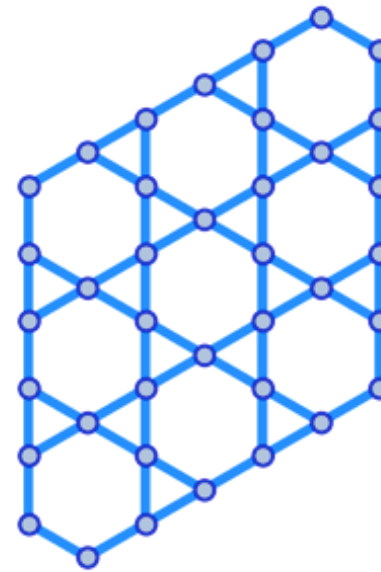


Band Structure Correspondence

Layout X

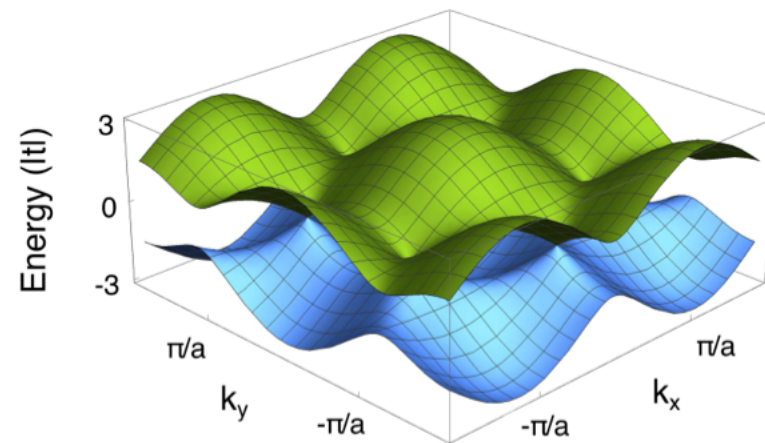
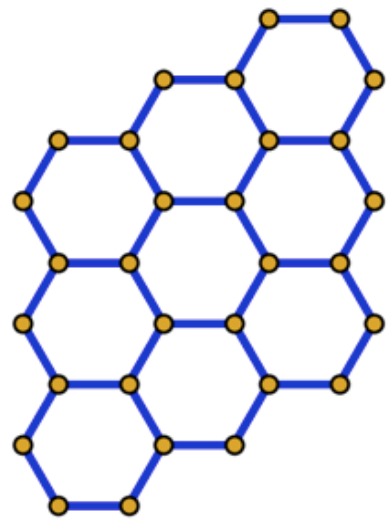


Line Graph $L(X)$

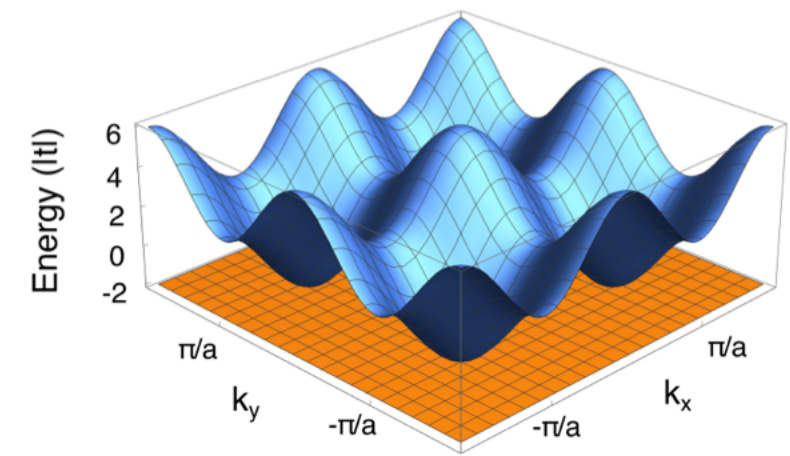
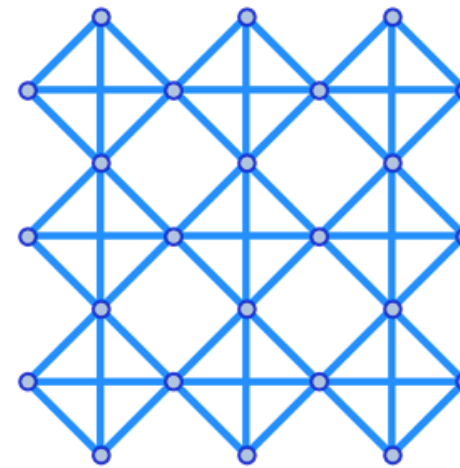
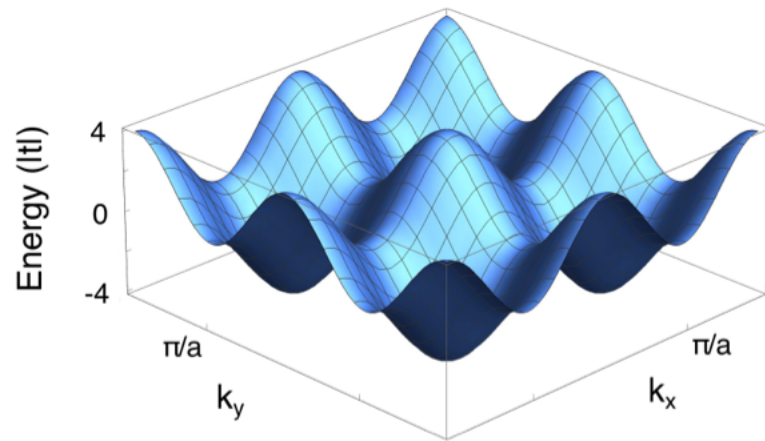
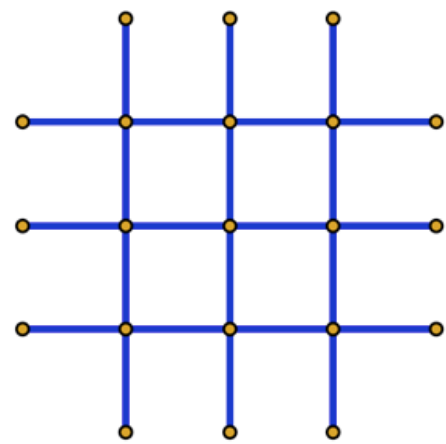
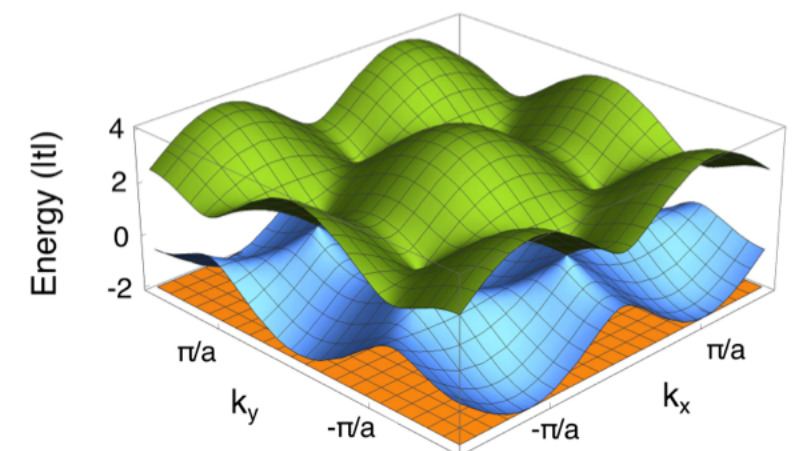
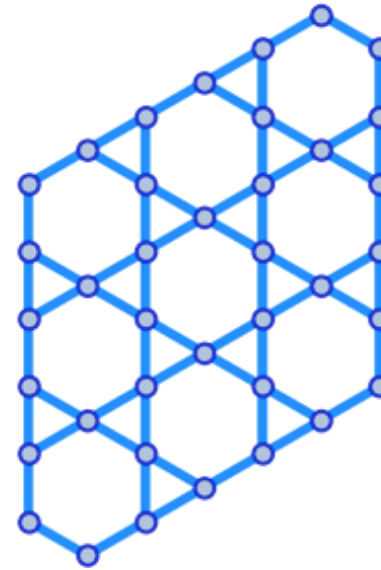


Band Structure Correspondence

Layout X

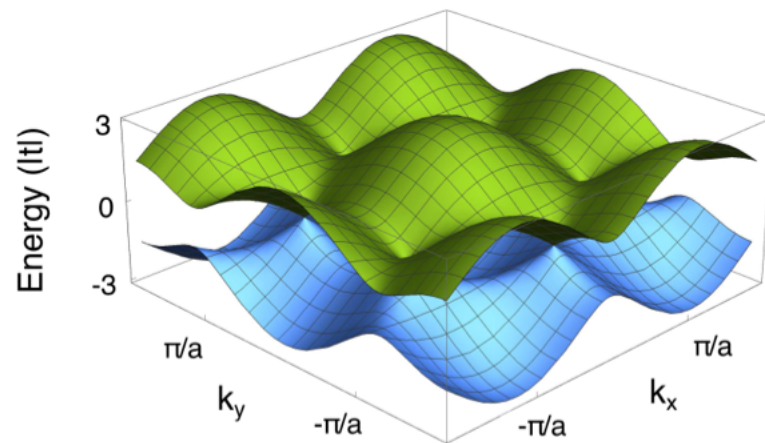
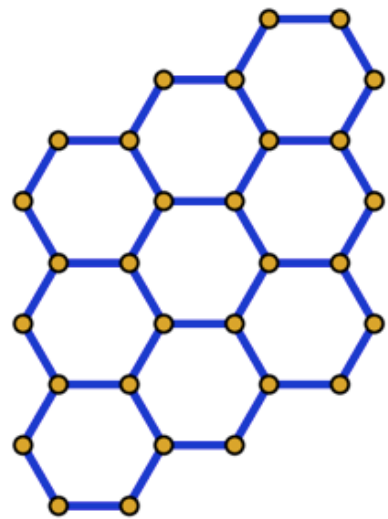


Line Graph $L(X)$

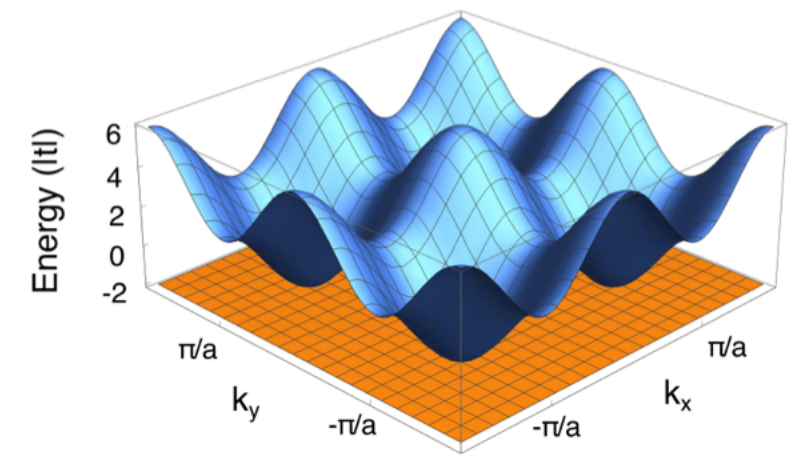
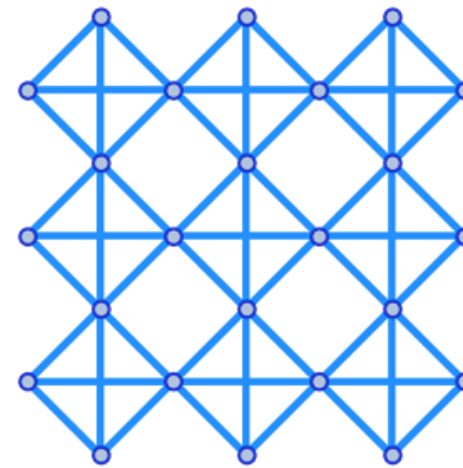
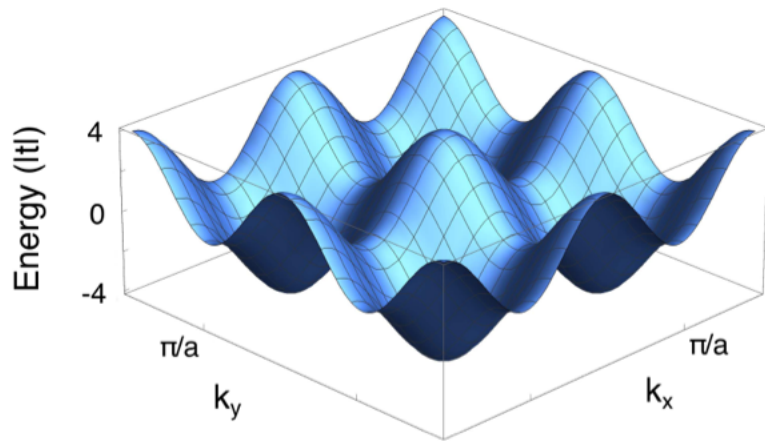
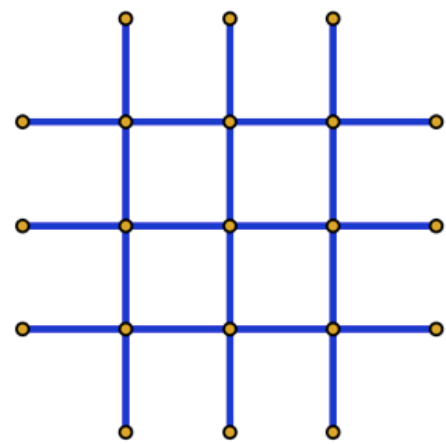
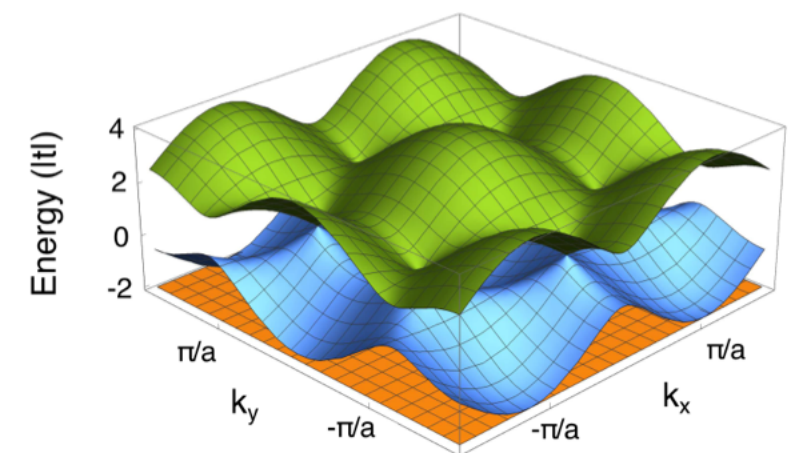
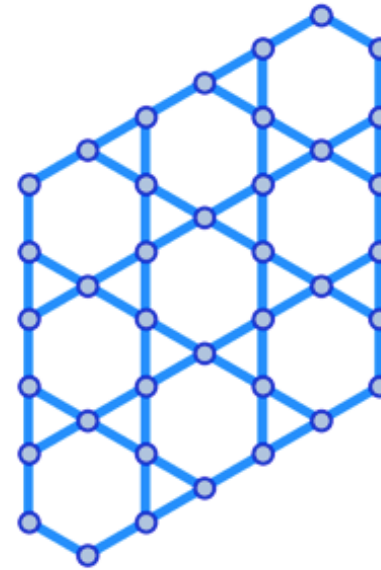


Band Structure Correspondence

Layout X



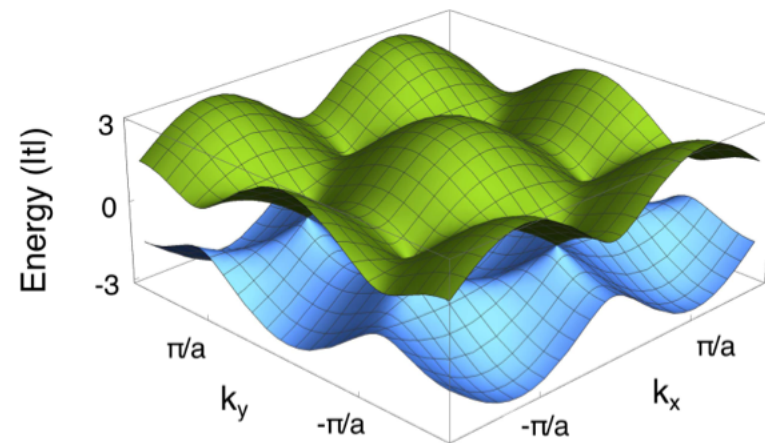
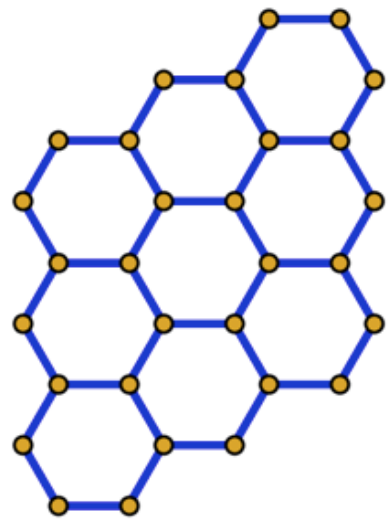
Line Graph $L(X)$



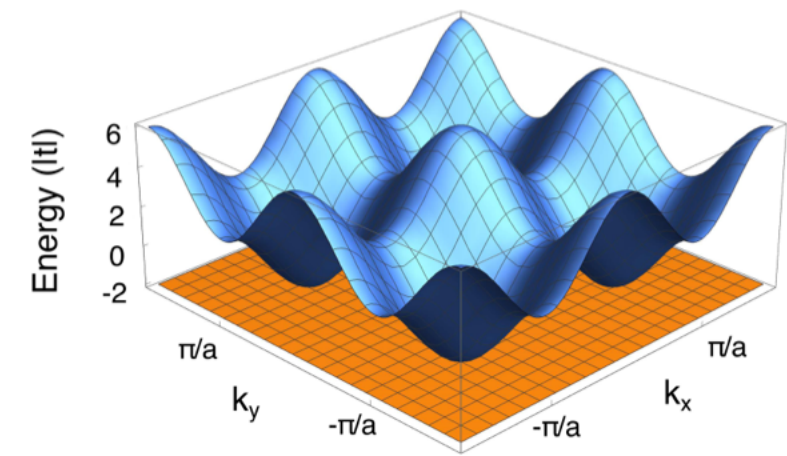
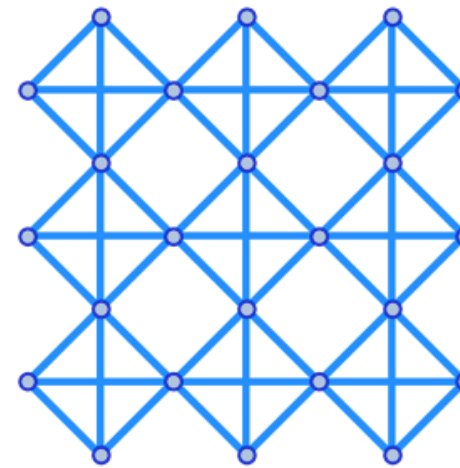
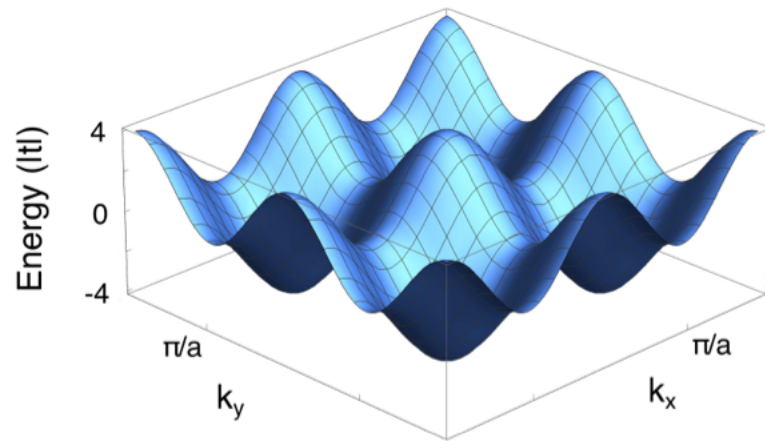
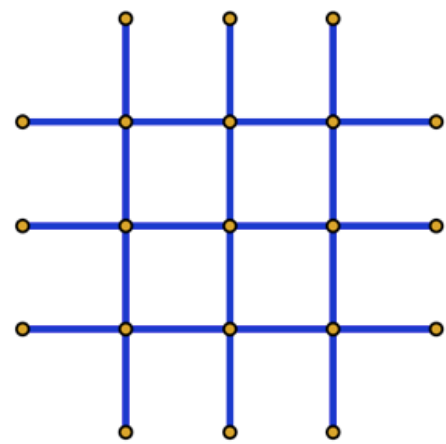
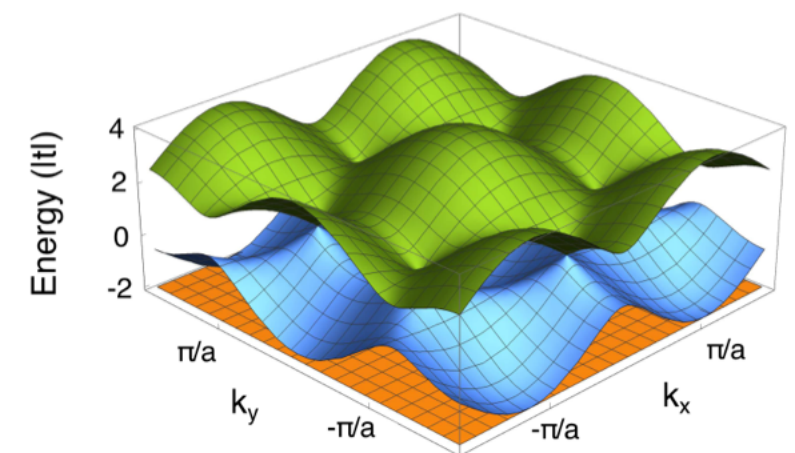
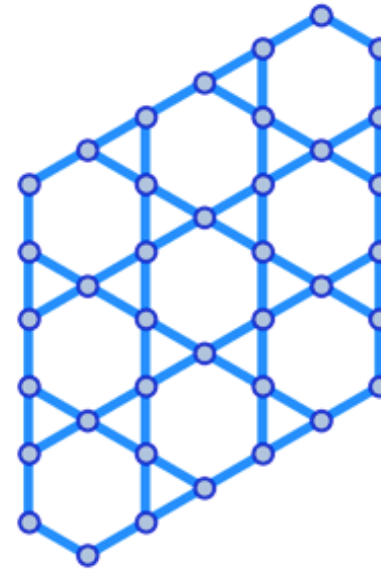
$$E_{\bar{H}_s} = \begin{cases} d - 2 + E_{H_X} \\ -2 \end{cases}$$

Band Structure Correspondence

Layout X



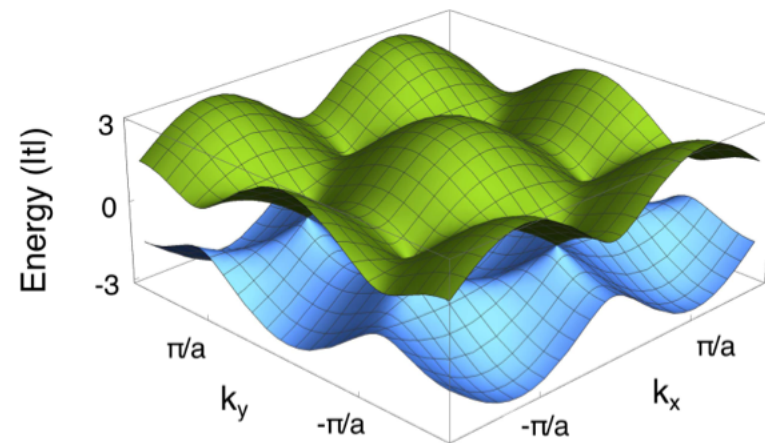
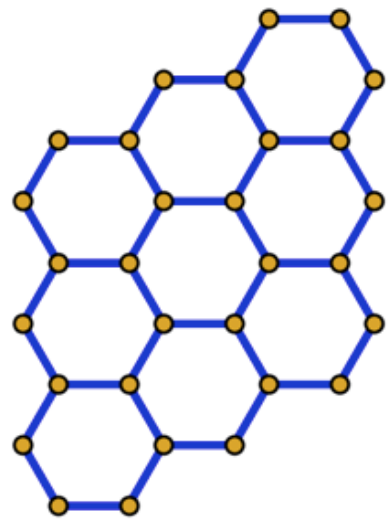
Line Graph $L(X)$



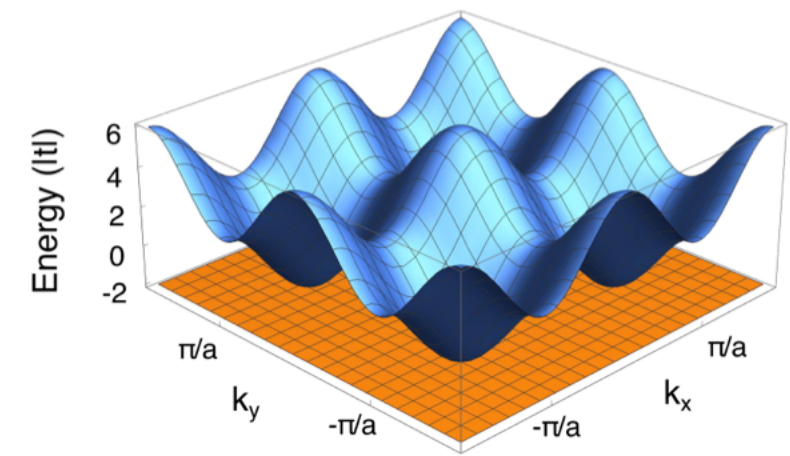
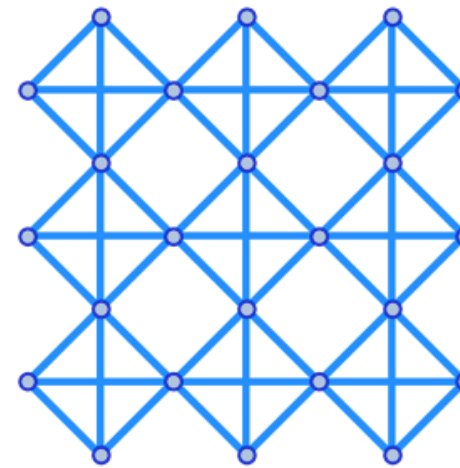
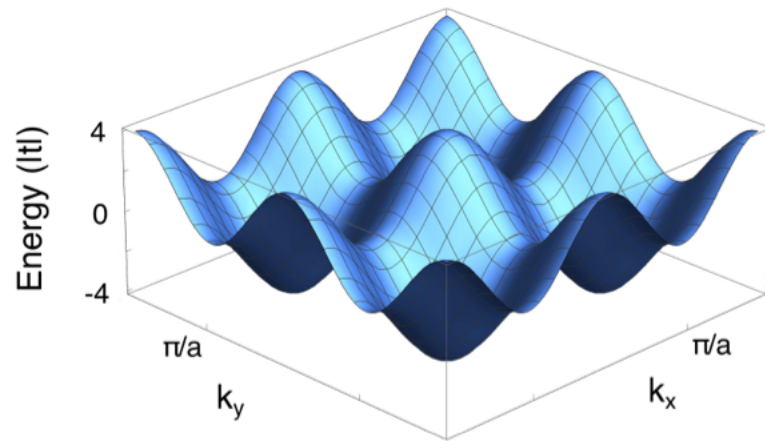
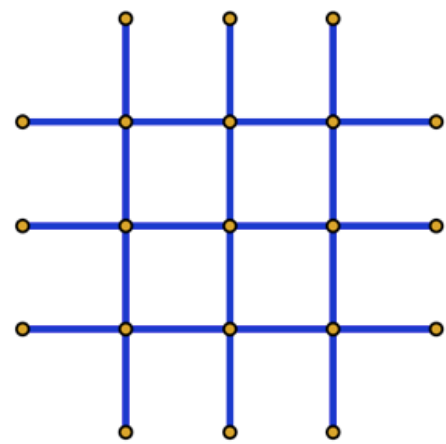
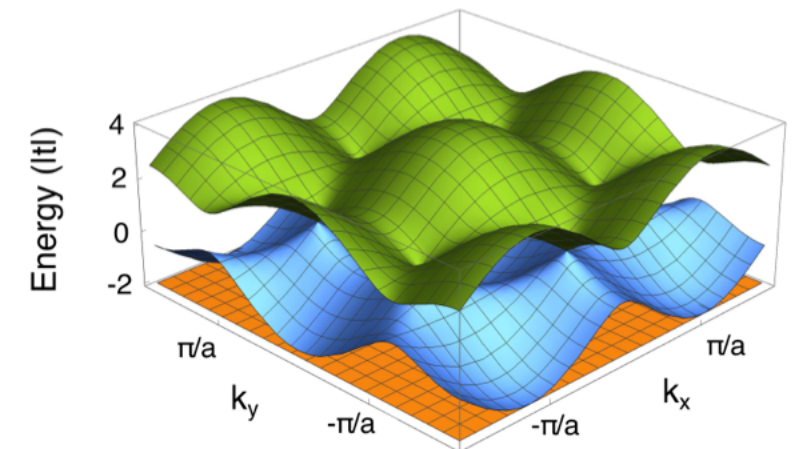
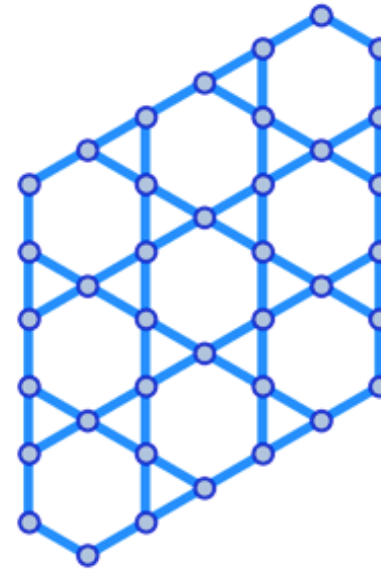
$$E_{\bar{H}_s} = \begin{cases} d - 2 + E_{H_X} & \leftarrow \text{Shifted Bands} \\ -2 \end{cases}$$

Band Structure Correspondence

Layout X



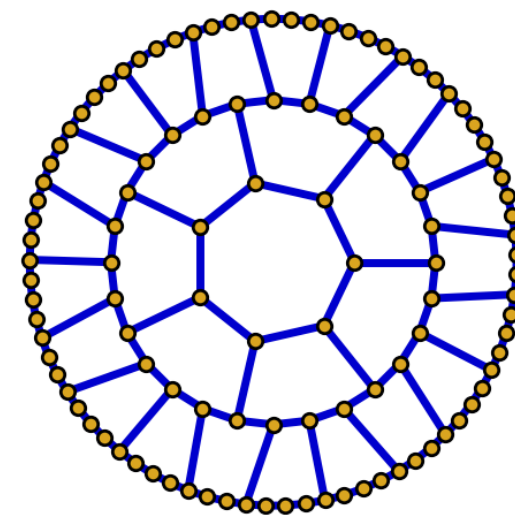
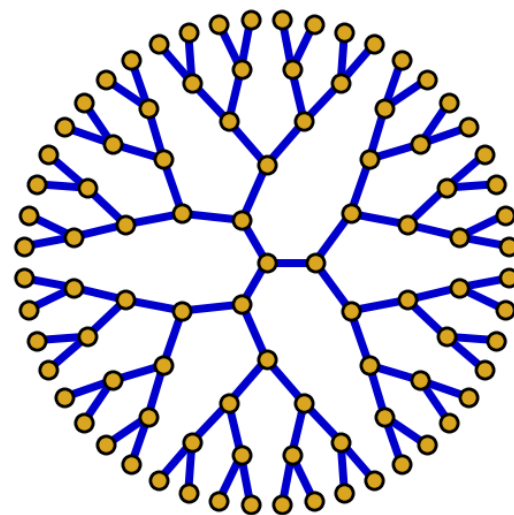
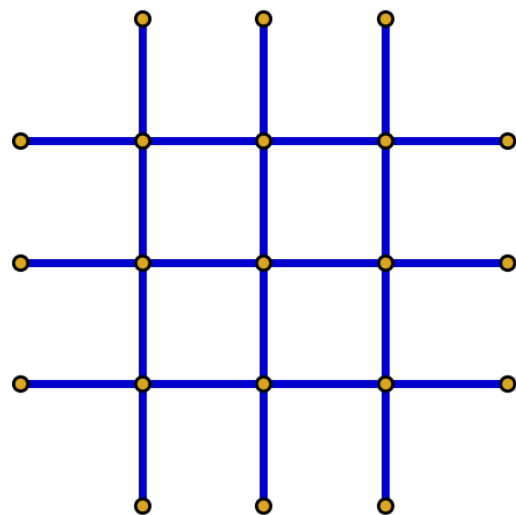
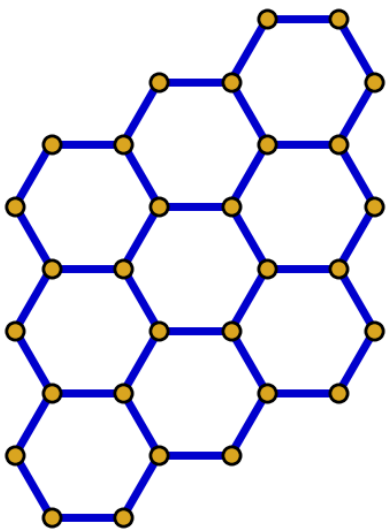
Line Graph $L(X)$



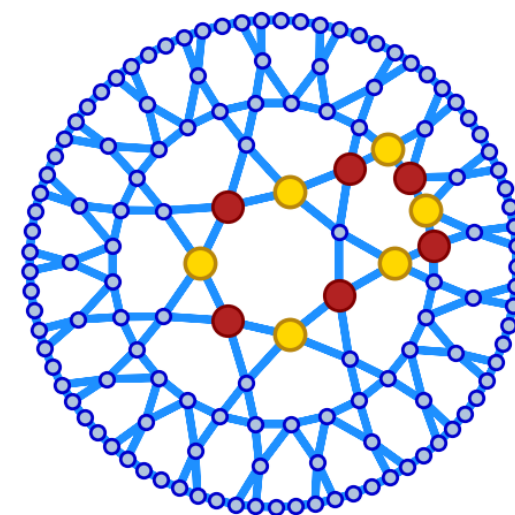
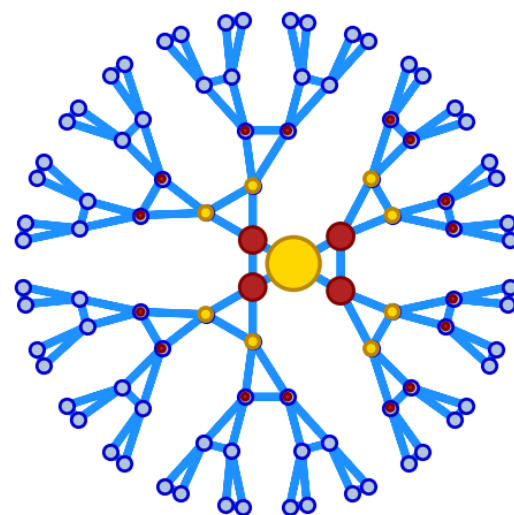
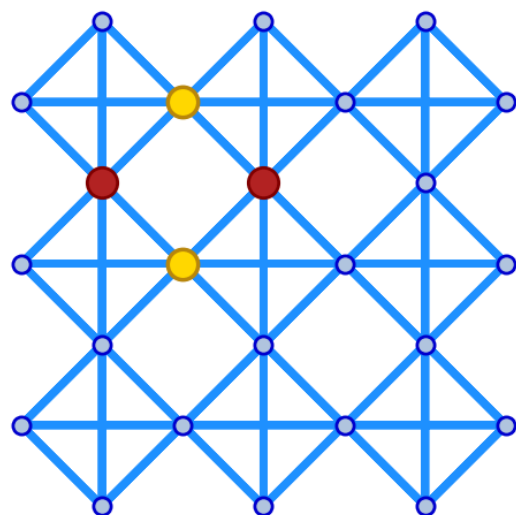
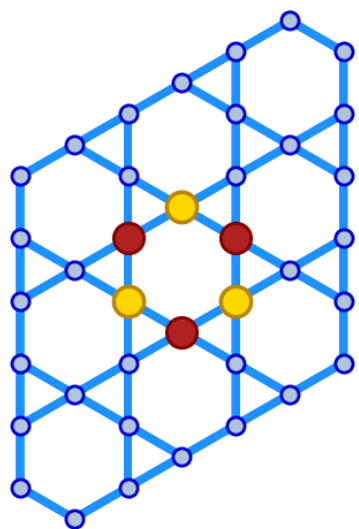
$$E_{\bar{H}_s} = \begin{cases} d - 2 + E_{H_X} & \leftarrow \text{Shifted Bands} \\ -2 & \leftarrow \text{Flat Band(s)} \end{cases}$$

Density of States and Flat-Band States

Layout X



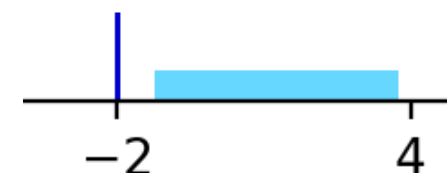
Line Graph $L(X)$



DOS X

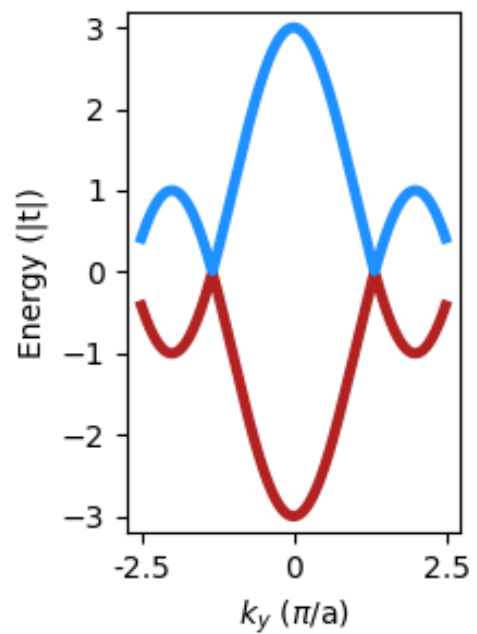
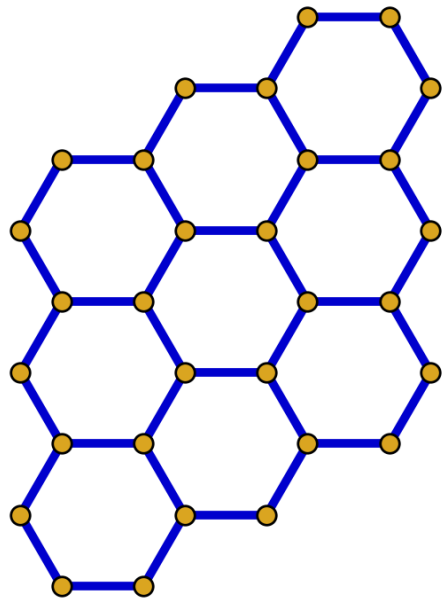


DOS $L(X)$



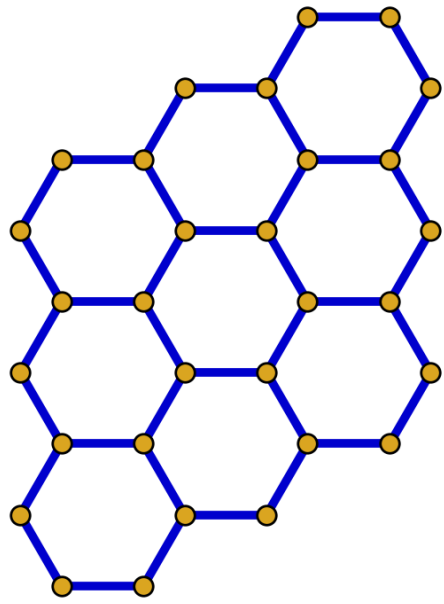
Subdivision Graphs and Optimally Gapped Flat Bands

X



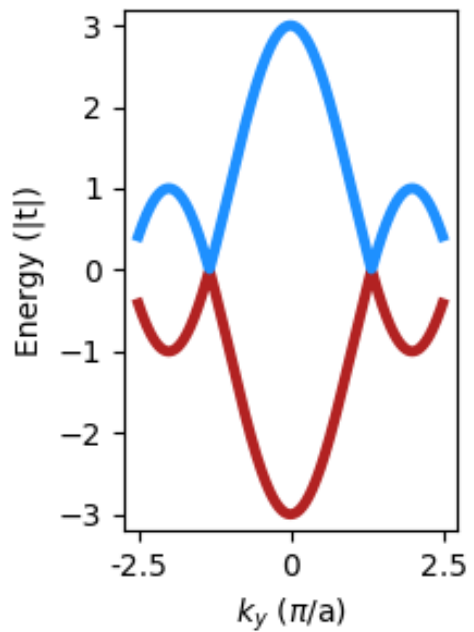
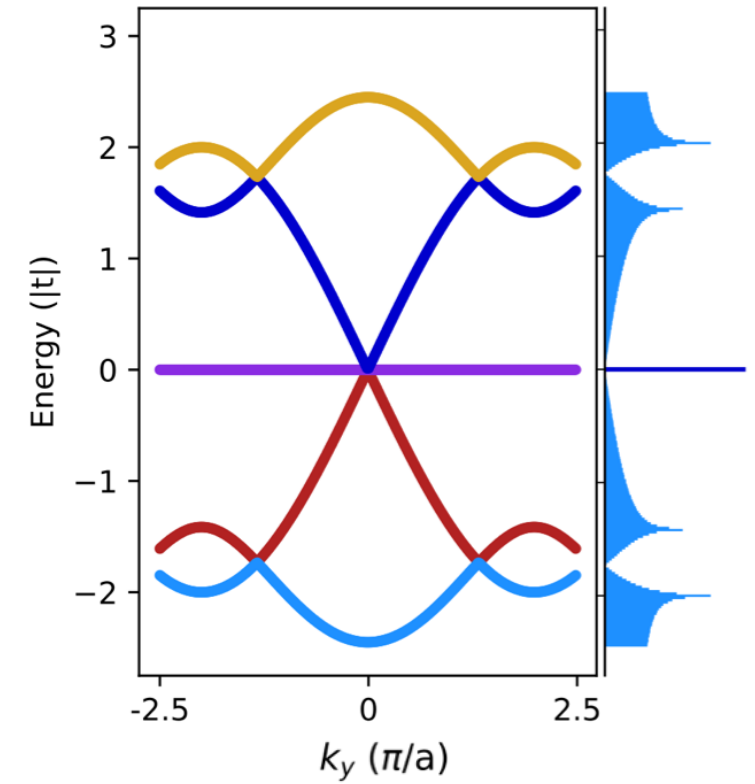
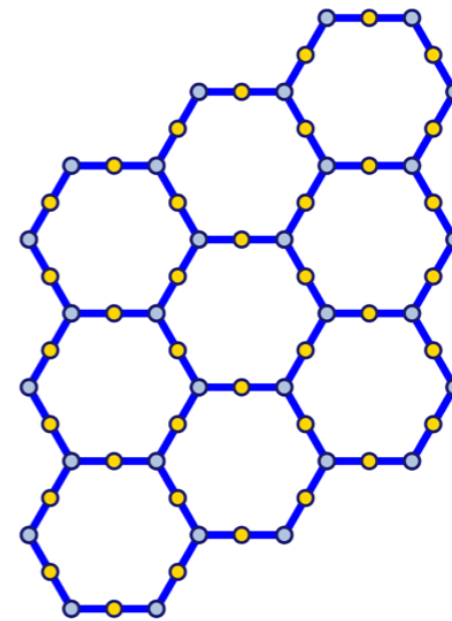
Subdivision Graphs and Optimally Gapped Flat Bands

X



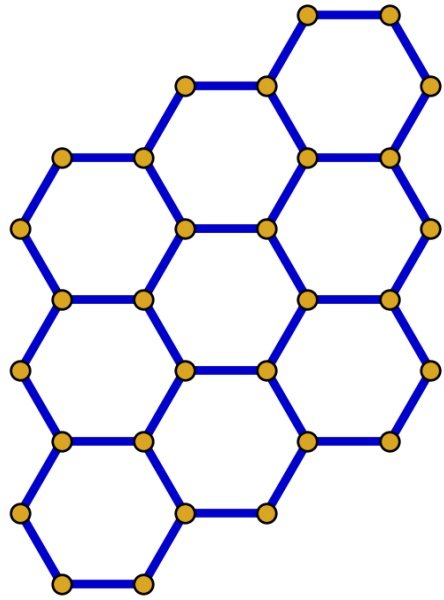
$$E_{\mathcal{S}(X)} = \begin{cases} \pm\sqrt{E_X + 3} \\ 0 \end{cases}$$

$\mathcal{S}(X)$



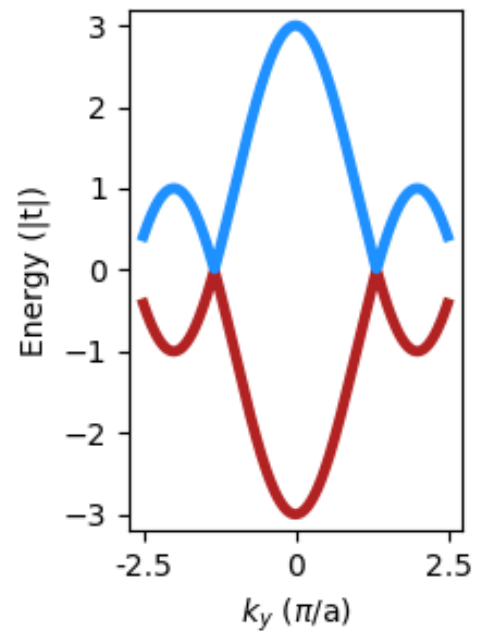
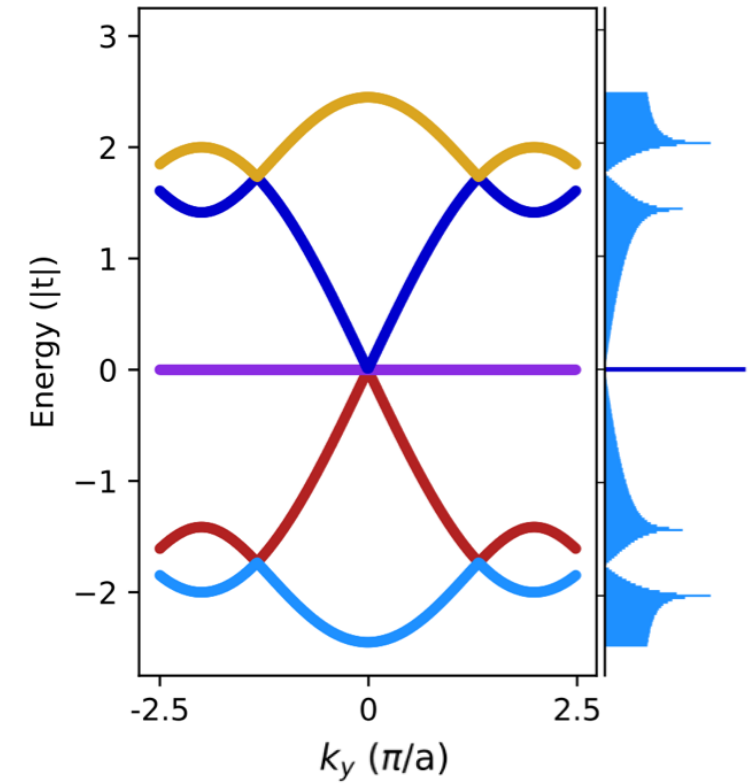
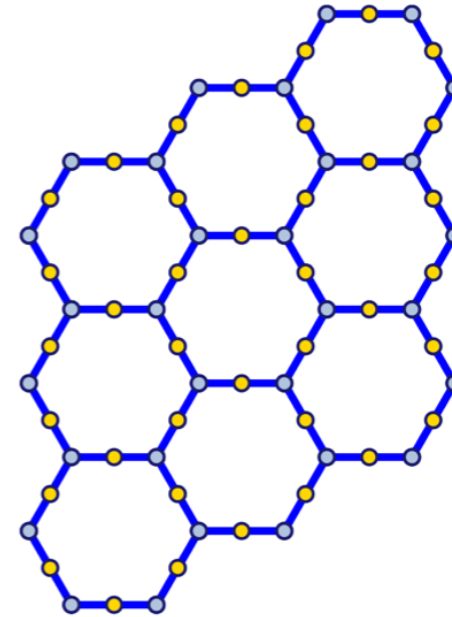
Subdivision Graphs and Optimally Gapped Flat Bands

X



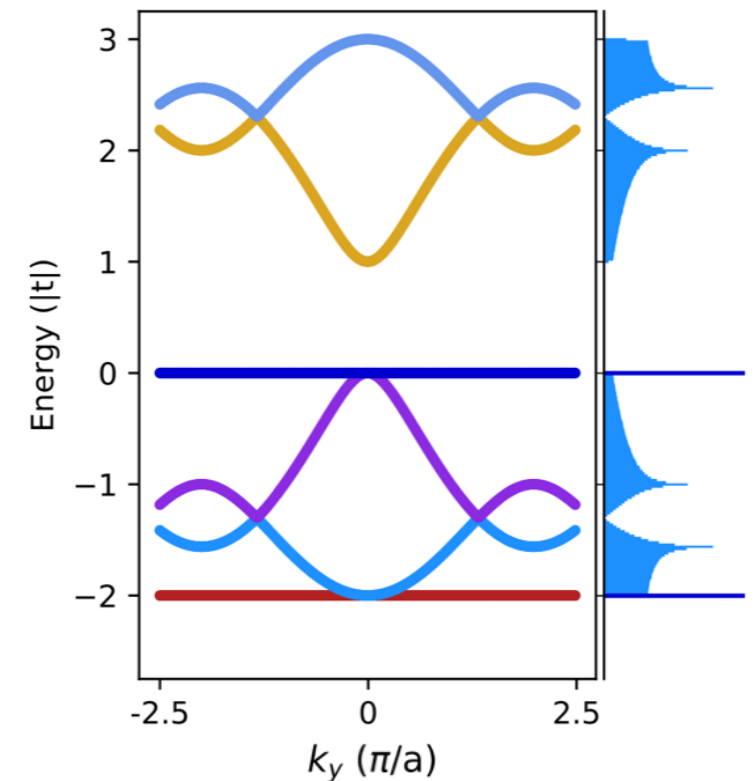
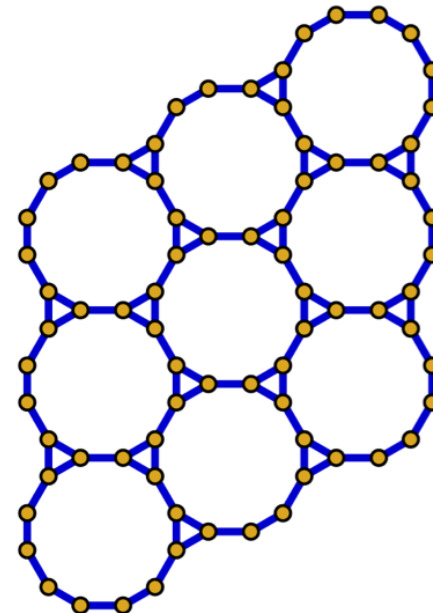
$$E_{\mathcal{S}(X)} = \begin{cases} \pm\sqrt{E_X + 3} \\ 0 \end{cases}$$

$\mathcal{S}(X)$

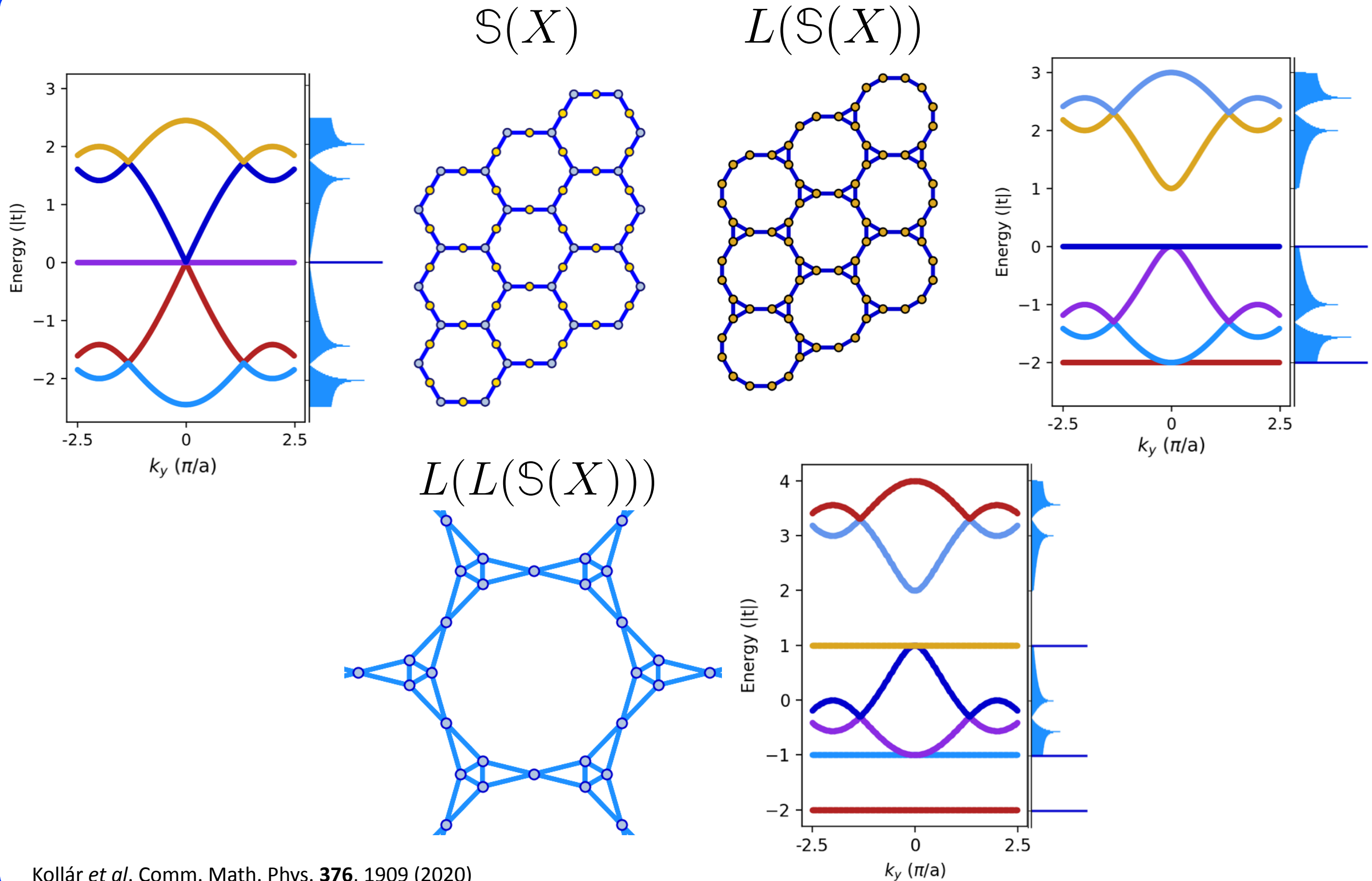


$$E_{L(\mathcal{S}(X))} = \begin{cases} \frac{1 \pm \sqrt{1 + 4(E_X + 3)}}{2} \\ 0 \\ -2 \end{cases}$$

$L(\mathcal{S}(X))$



Subdivision Graphs and Optimally Gapped Flat Bands



New Lattices for Photon-Mediated Interactions

Photonic Crystal + qubits

- Effective swap interaction
- All modes in parallel

$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

New Lattices for Photon-Mediated Interactions

Photonic Crystal + qubits

- Effective swap interaction
- All modes in parallel

$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

- **Hyperbolic Lattice**

- Follows hyperbolic metric

New Lattices for Photon-Mediated Interactions

Photonic Crystal + qubits

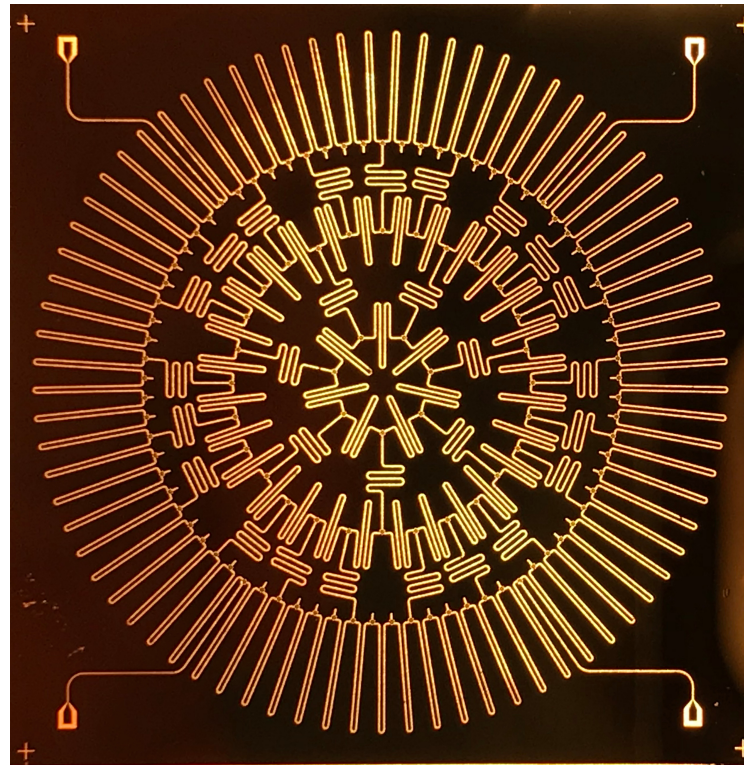
- Effective swap interaction
- All modes in parallel

$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

- **Hyperbolic Lattice**

- Follows hyperbolic metric

AJK *et al.* Nature **571** (2019)



New Lattices for Photon-Mediated Interactions

Photonic Crystal + qubits

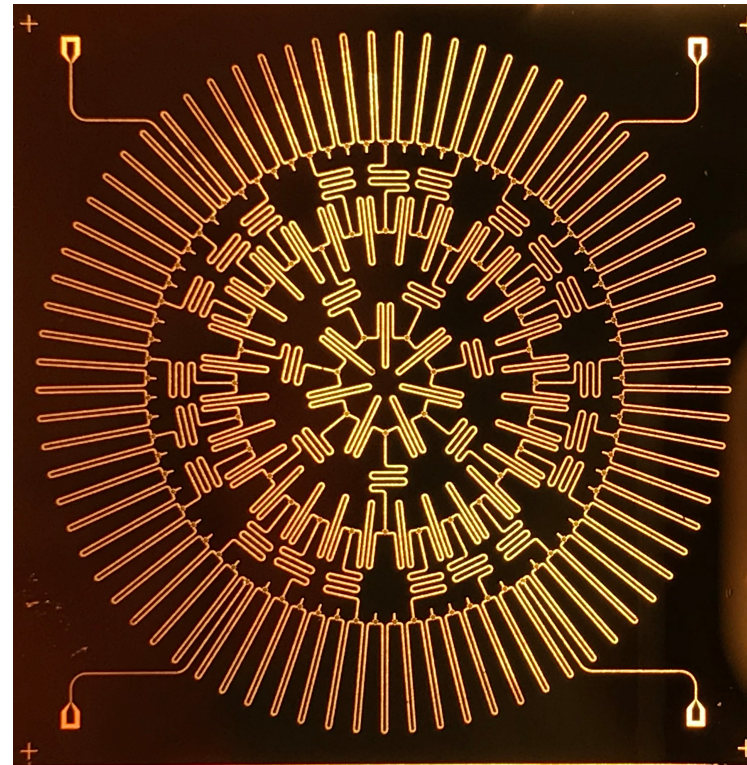
- Effective swap interaction
- All modes in parallel

$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

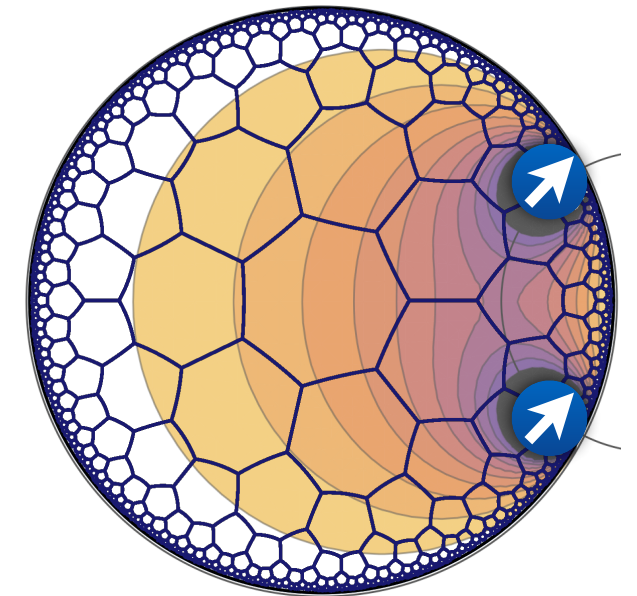
- **Hyperbolic Lattice**

- Follows hyperbolic metric

AJK *et al.* Nature **571** (2019)



Bienias, AJK *et al.* Phys. Rev. Lett. **128**, 013601 (2022)



New Lattices for Photon-Mediated Interactions

Photonic Crystal + qubits

- Effective swap interaction
- All modes in parallel

$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

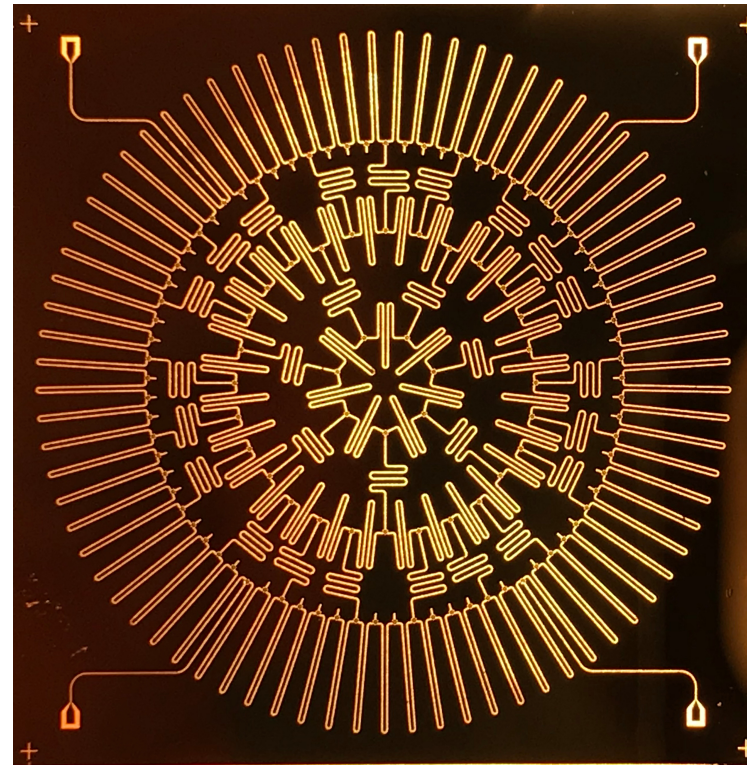
- **Hyperbolic Lattice**

- Follows hyperbolic metric

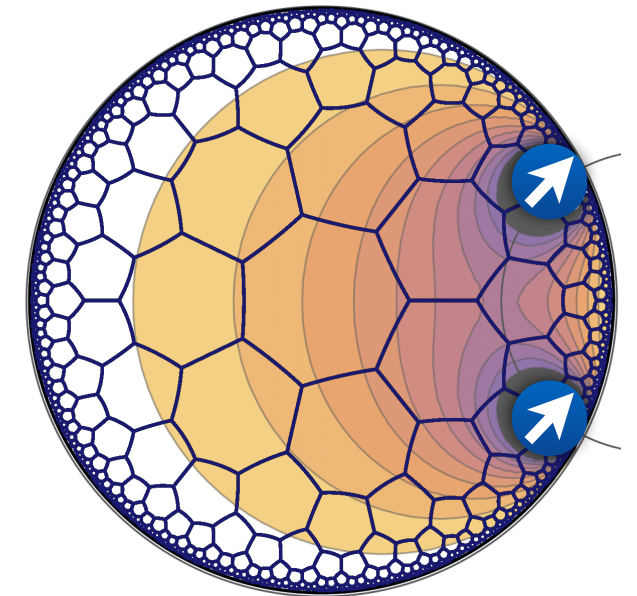
- **Flat-Band Lattice**

- Frustrated Magnet

AJK *et al.* Nature **571** (2019)



Bienias, AJK *et al.* Phys. Rev. Lett. **128**, 013601 (2022)



New Lattices for Photon-Mediated Interactions

Photonic Crystal + qubits

- Effective swap interaction
- All modes in parallel

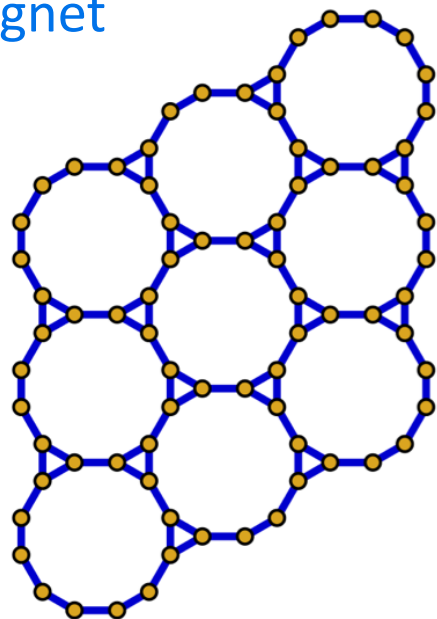
$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

- **Hyperbolic Lattice**

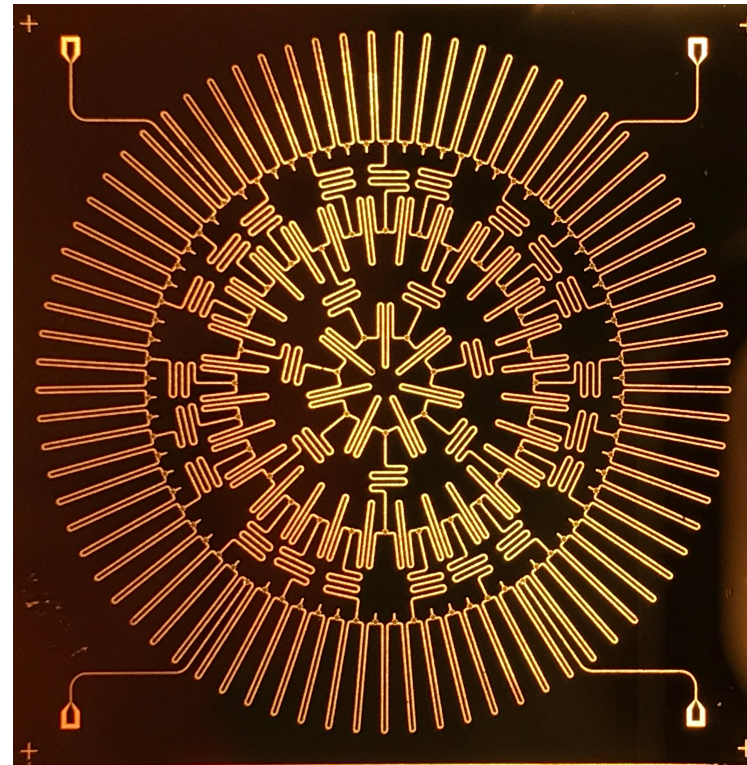
- Follows hyperbolic metric

- **Flat-Band Lattice**

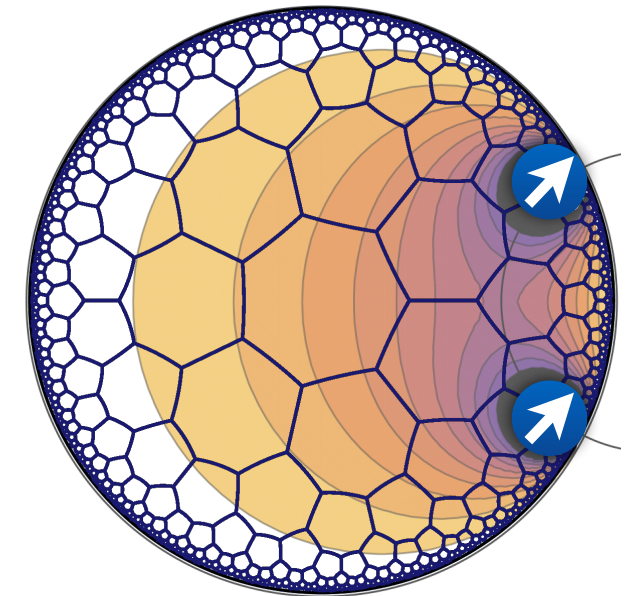
- Frustrated Magnet



AJK *et al.* Nature **571** (2019)



Bienias, AJK *et al.* Phys. Rev. Lett. **128**, 013601 (2022)



New Lattices for Photon-Mediated Interactions

Photonic Crystal + qubits

- Effective swap interaction
- All modes in parallel

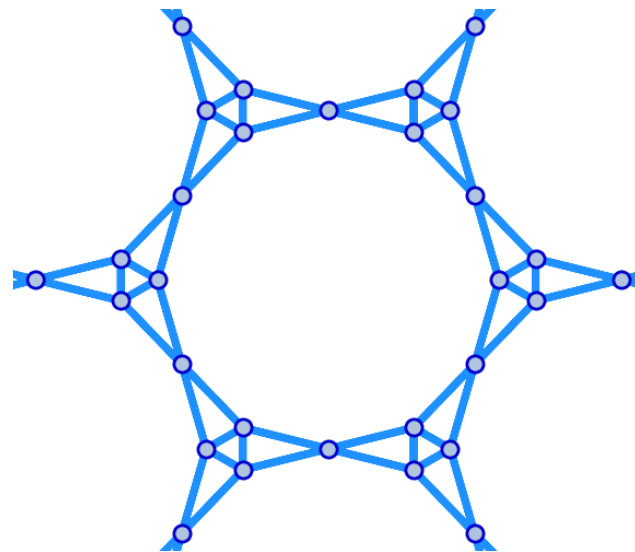
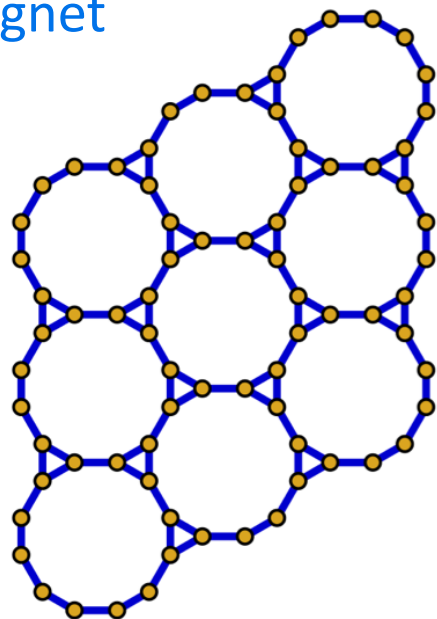
$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

- **Hyperbolic Lattice**

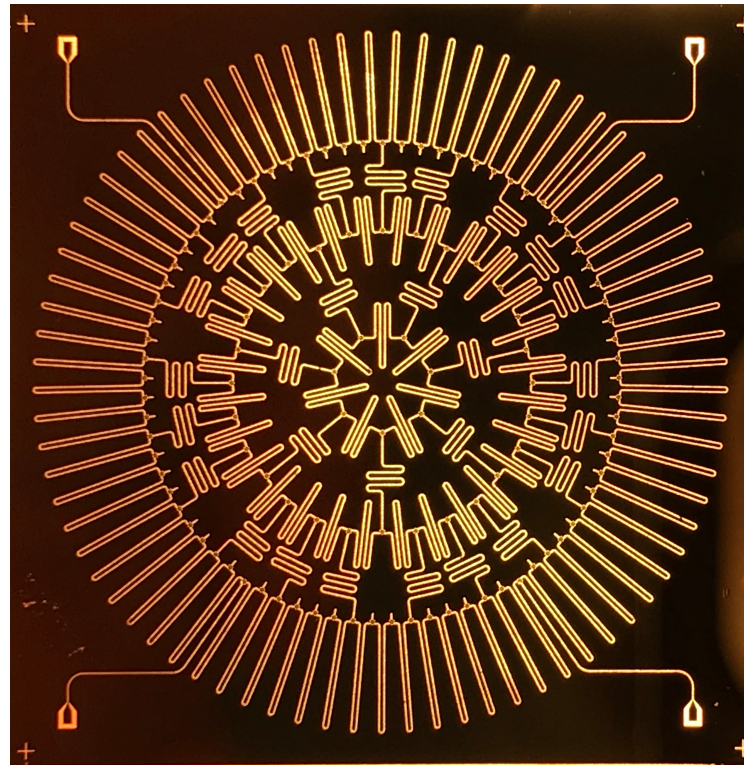
- Follows hyperbolic metric

- **Flat-Band Lattice**

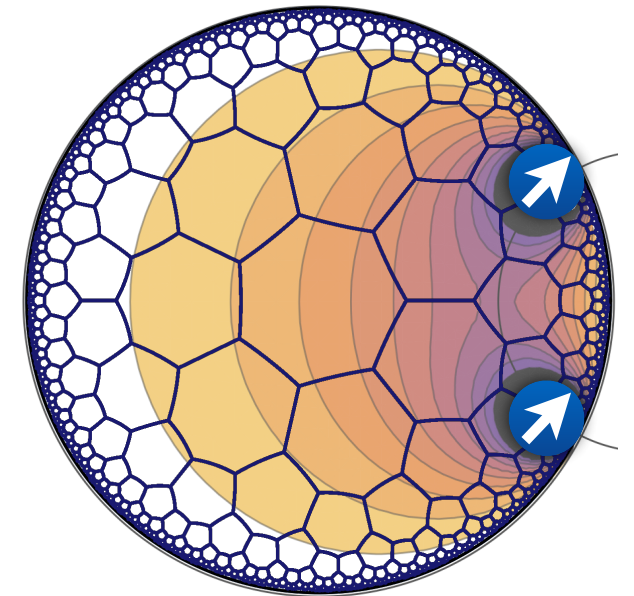
- Frustrated Magnet



AJK *et al.* Nature **571** (2019)



Bienias, AJK *et al.* Phys. Rev. Lett. **128**, 013601 (2022)



New Lattices for Photon-Mediated Interactions

Photonic Crystal + qubits

- Effective swap interaction
- All modes in parallel

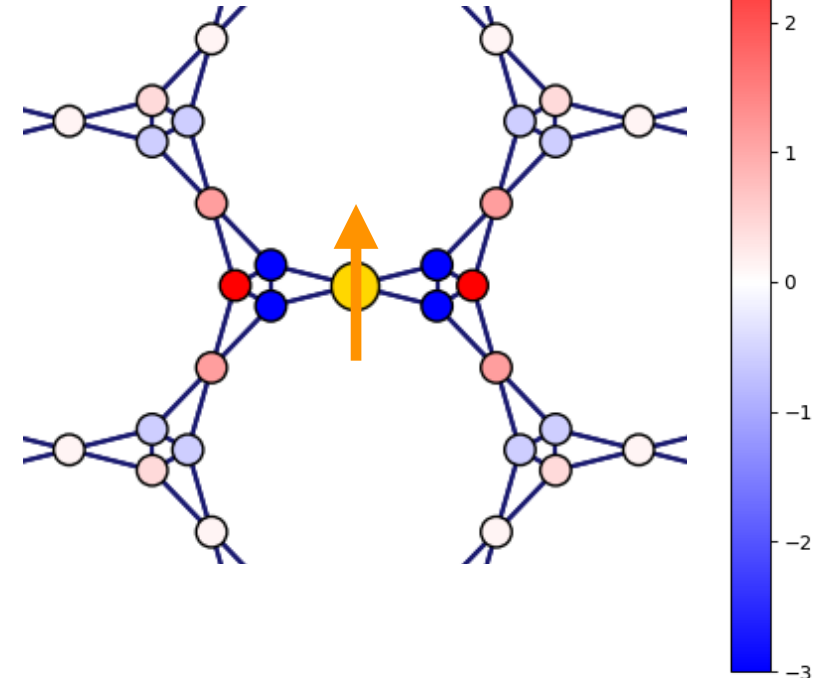
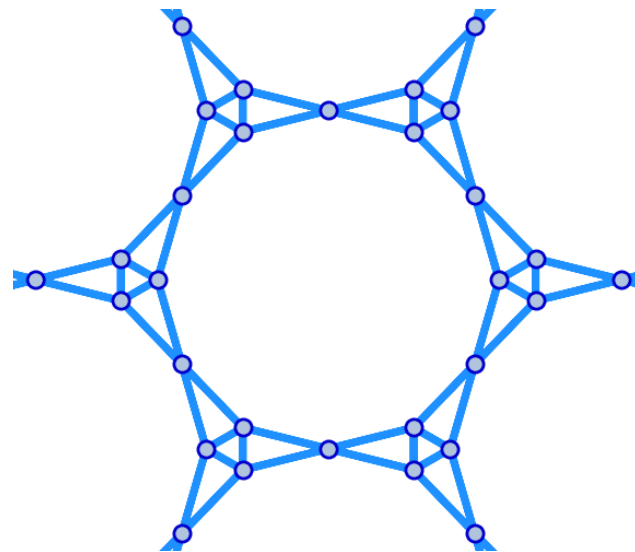
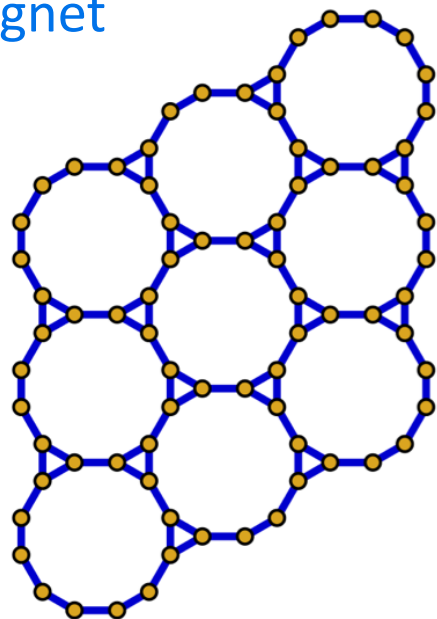
$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

• Hyperbolic Lattice

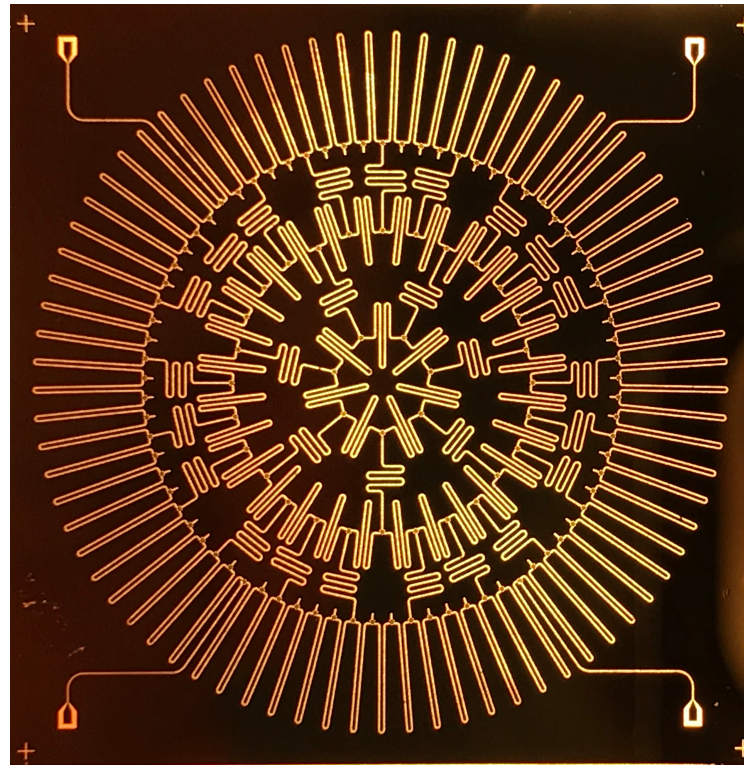
- Follows hyperbolic metric

• Flat-Band Lattice

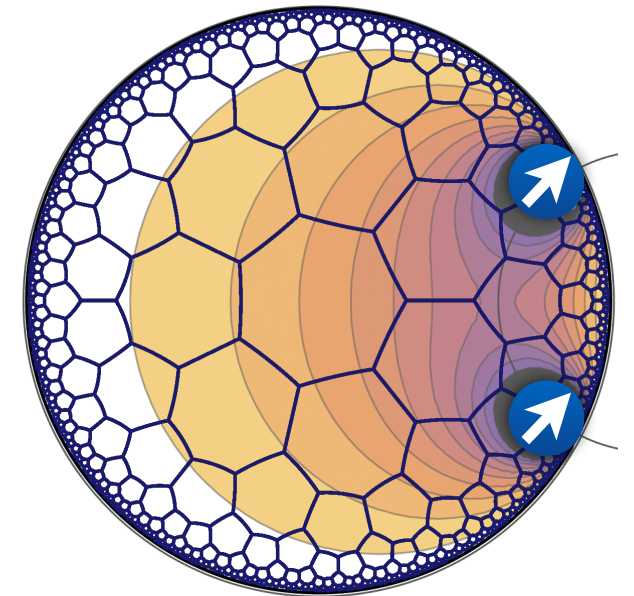
- Frustrated Magnet



AJK *et al.* Nature **571** (2019)



Bienias, AJK *et al.* Phys. Rev. Lett. **128**, 013601 (2022)



Outline

- Coplanar Waveguide (CPW) Lattices
 - Deformable lattice sites
 - Line-graph lattices
 - Interacting photons
- Band Engineering
 - Hyperbolic lattice
 - Gapped flat bands
- Mathematical Connections
 - Planar Gaps
 - Maximal Gaps
 - Quantum Error Correcting Codes
- Experimental Developments

Outline

- Coplanar Waveguide (CPW) Lattices
 - Deformable lattice sites
 - Line-graph lattices
 - Interacting photons
- Band Engineering
 - Hyperbolic lattice
 - Gapped flat bands
- Mathematical Connections
 - Planar Gaps
 - Maximal Gaps
 - Quantum Error Correcting Codes
- Experimental Developments

Other Maximal Gaps?

Two Driving Questions

- Even larger gaps possible at other energies?
- Where can planar graphs have gaps?

Other Maximal Gaps?

Two Driving Questions

- Even larger gaps possible at other energies?
- Where can planar graphs have gaps?

Thm:

No large 3-regular graph can have a gap larger than 2.

Other Maximal Gaps?

Two Driving Questions

- Even larger gaps possible at other energies?
- Where can planar graphs have gaps?

Thm:

No large 3-regular graph can have a gap larger than 2.

- Have found 2 such gaps.
- Conjecture that these are the only ones.

Other Maximal Gaps?

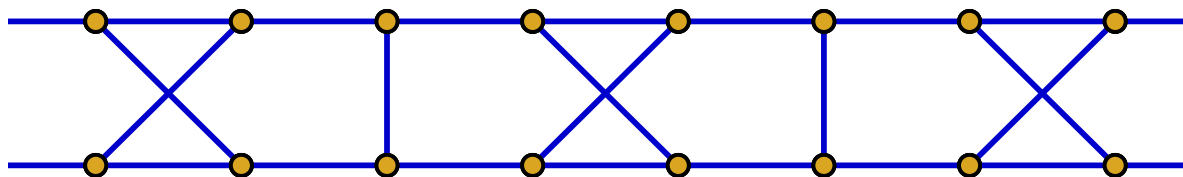
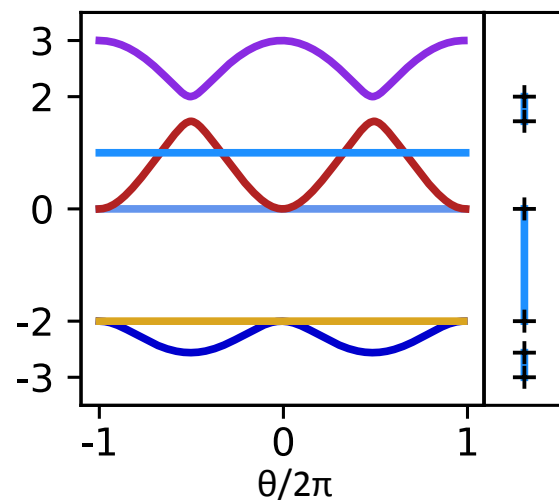
Two Driving Questions

- Even larger gaps possible at other energies?
- Where can planar graphs have gaps?

Thm:

No large 3-regular graph can have a gap larger than 2.

- Have found 2 such gaps.
- Conjecture that these are the only ones.



Other Maximal Gaps?

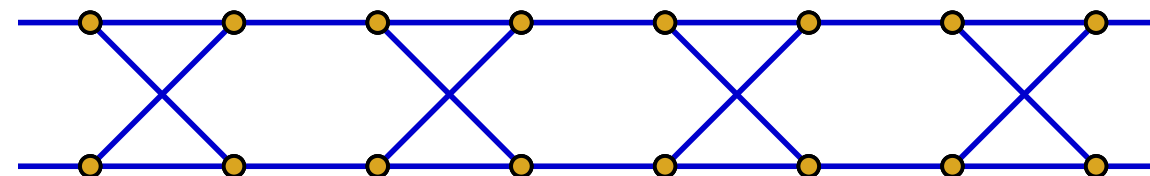
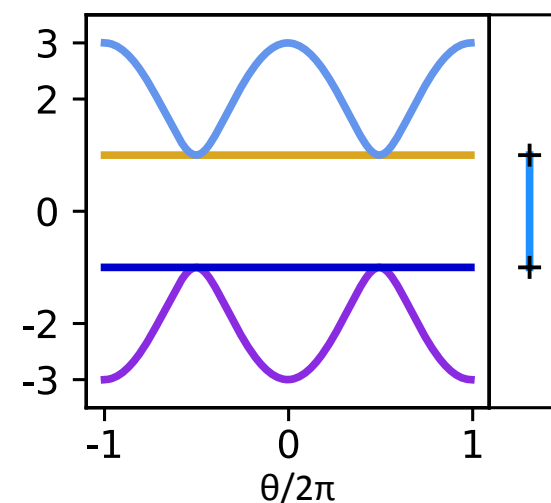
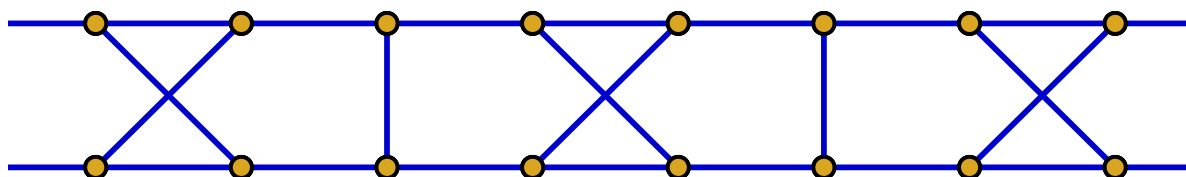
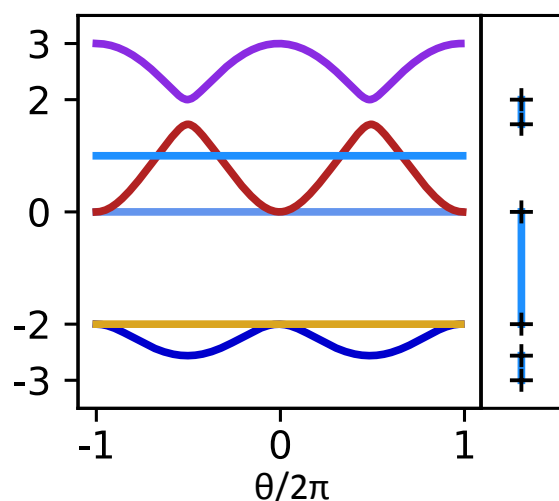
Two Driving Questions

- Even larger gaps possible at other energies?
- Where can planar graphs have gaps?

Thm:

No large 3-regular graph can have a gap larger than 2.

- Have found 2 such gaps.
- Conjecture that these are the only ones.



Other Maximal Gaps?

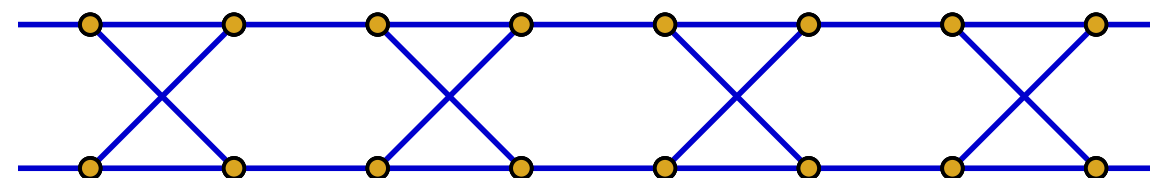
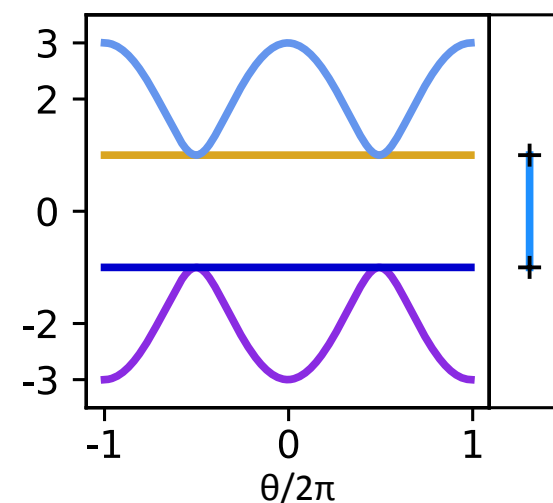
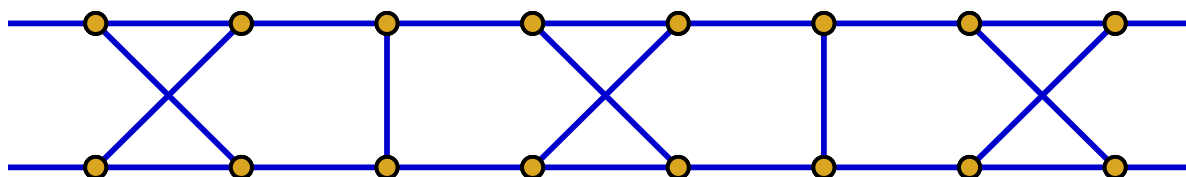
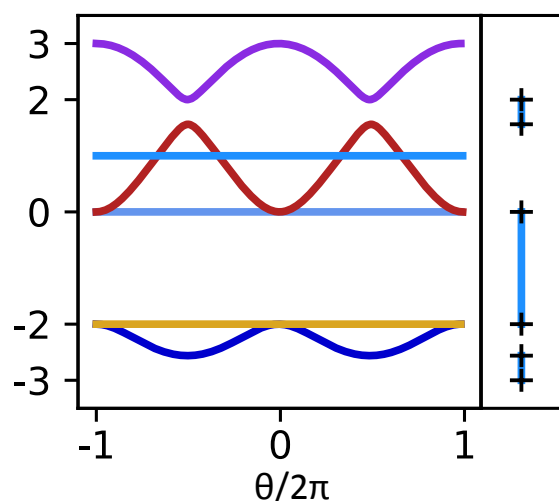
Two Driving Questions

- Even larger gaps possible at other energies?
- Where can planar graphs have gaps?

Thm:

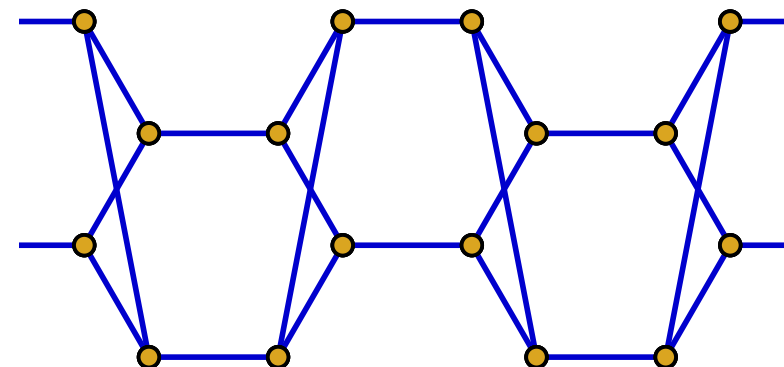
No large 3-regular graph can have a gap larger than 2.

- Have found 2 such gaps.
- Conjecture that these are the only ones.



A.K.A.

$n=2, m=0$ carbon nanotube



Kollár *et al.* Comm. AMS 1,1 (2021)

Guo, Mohar Lin. Alg. and Appl. 449, 68-75 (2014)

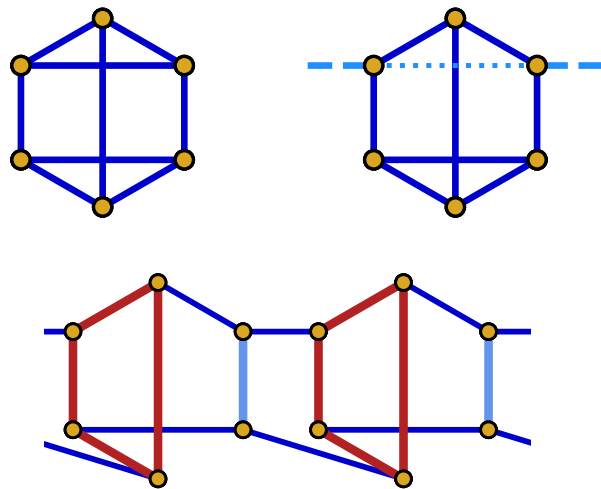
Abelian Covers and Planar Gaps

New Lattice Viewpoint

Abelian Covers and Planar Gaps

New Lattice Viewpoint

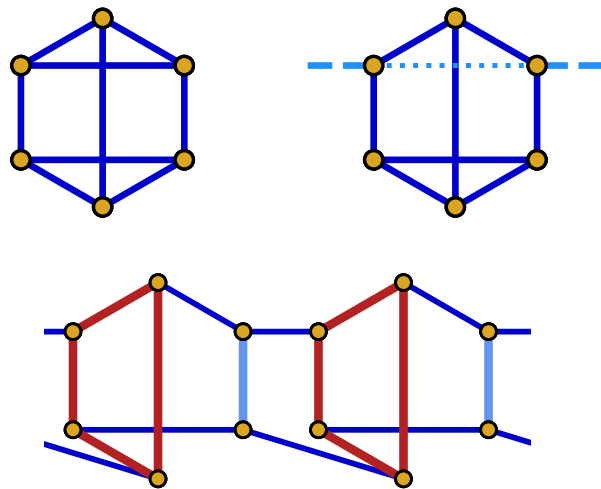
- Use method of Abelian covers to construct examples.
 - “Unwrap” small graph to form lattice



Abelian Covers and Planar Gaps

New Lattice Viewpoint

- Use method of Abelian covers to construct examples.
 - “Unwrap” small graph to form lattice

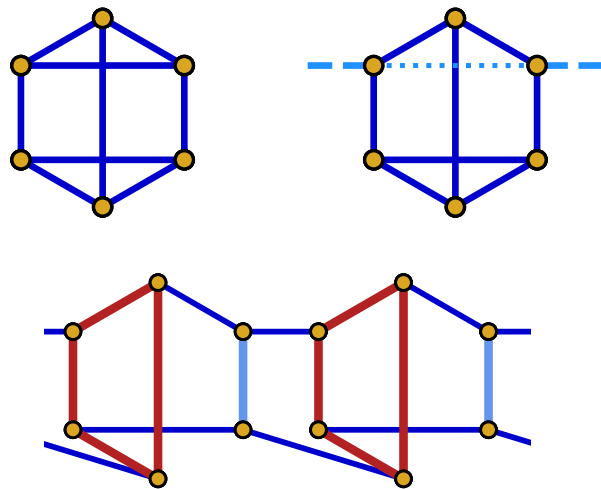


- Initial energies are $k=0$ energies of the lattice
 - Small graphs and their spectra tabulated.
 - “Periodic table” of unit cells to start from.

Abelian Covers and Planar Gaps

New Lattice Viewpoint

- Use method of Abelian covers to construct examples.
 - “Unwrap” small graph to form lattice



- Initial energies are $k=0$ energies of the lattice
 - Small graphs and their spectra tabulated.
 - “Periodic table” of unit cells to start from.

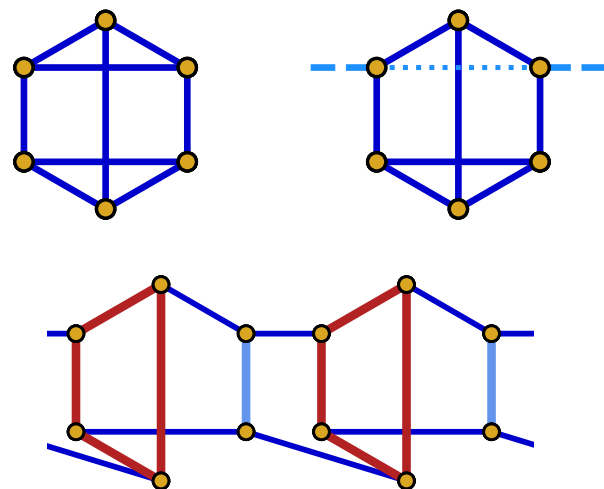
Thm:

All points in $[-3,3)$ can be gapped by large 3-regular planar graphs.

Abelian Covers and Planar Gaps

New Lattice Viewpoint

- Use method of Abelian covers to construct examples.
 - “Unwrap” small graph to form lattice

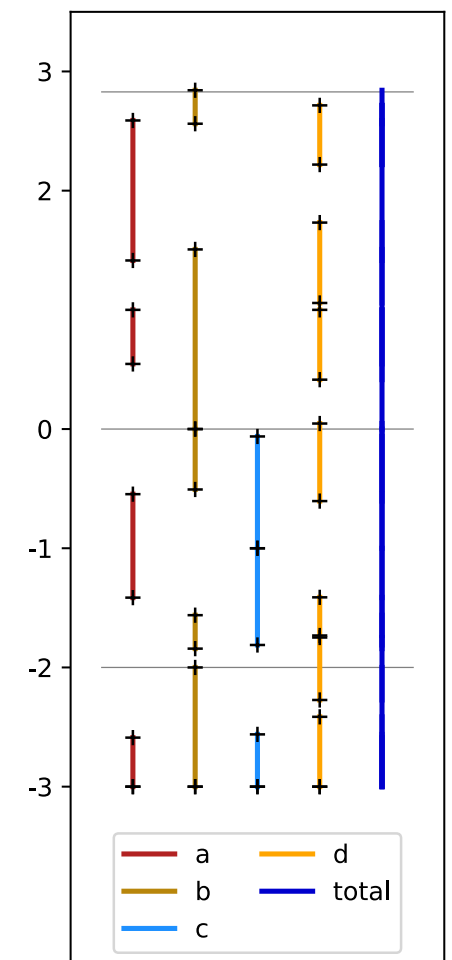
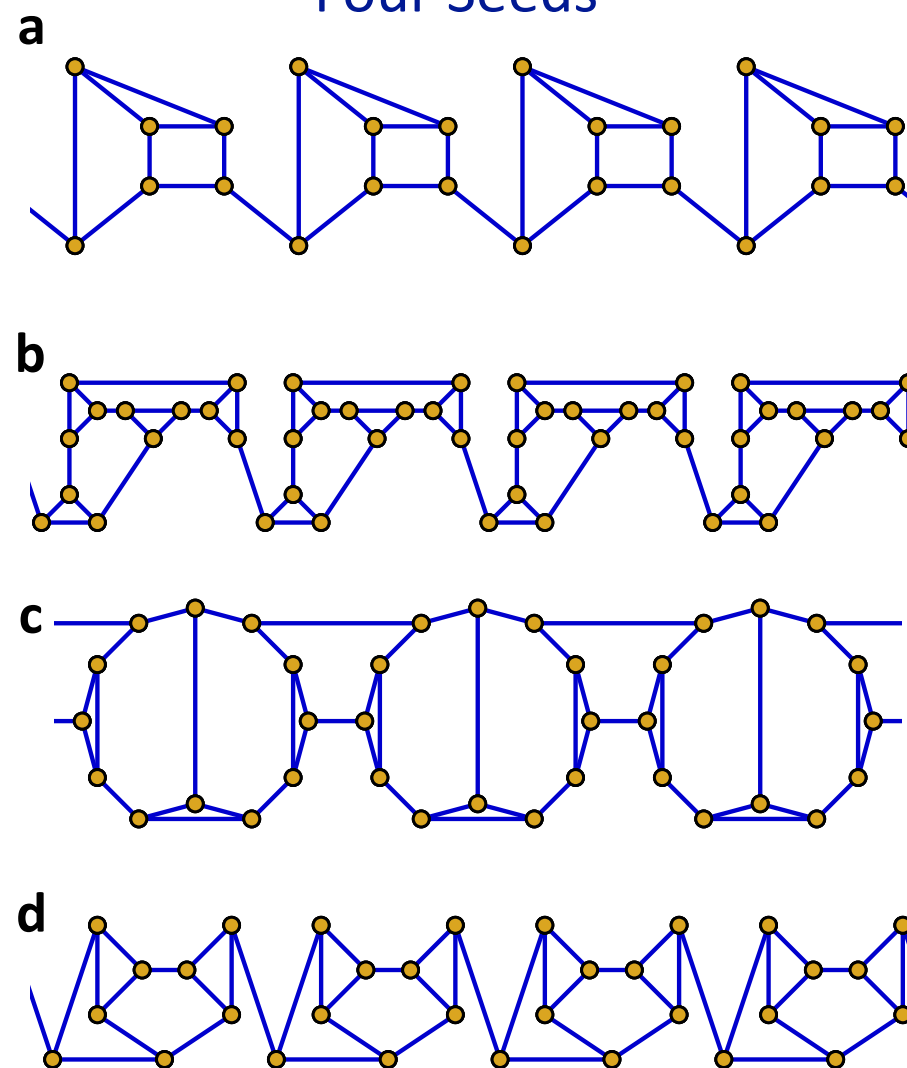


- Initial energies are $k=0$ energies of the lattice
 - Small graphs and their spectra tabulated.
 - “Periodic table” of unit cells to start from.

Thm:

All points in $[-3,3)$ can be gapped by large 3-regular planar graphs.

Four Seeds



- Combined gaps cover $[-3, 2\sqrt{2}]$.
- Iteration of $L(S(X))$ covers the rest.

Line-Graph Subsystem Codes

Thm: (Chapman and Flammia)

A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

Chapman *et al.* Quantum **4**, 278 (2020)

Line-Graph Subsystem Codes

Thm: (Chapman and Flammia)

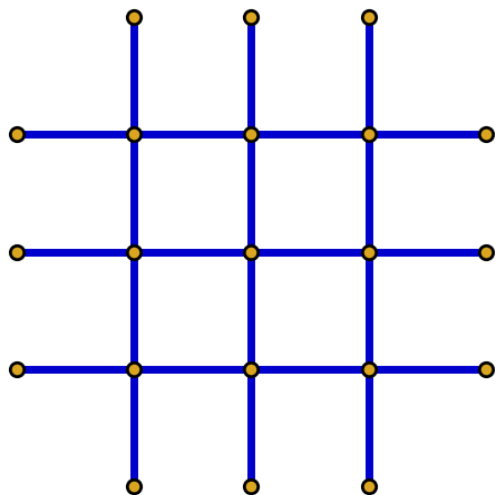
A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

Chapman *et al.* Quantum **4**, 278 (2020)

The Checkerboard-Lattice Code

- Built on the square lattice

Fermion Lattice



Line-Graph Subsystem Codes

Thm: (Chapman and Flammia)

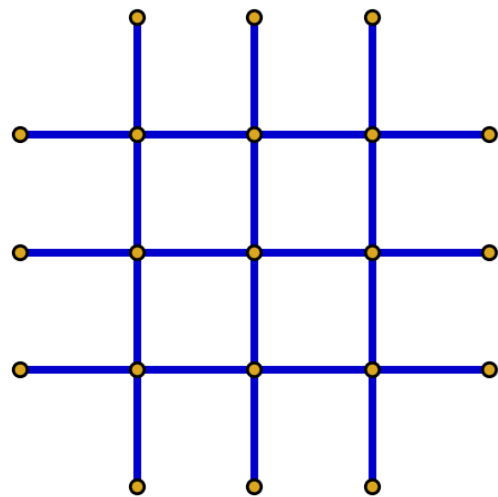
A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

Chapman *et al.* Quantum **4**, 278 (2020)

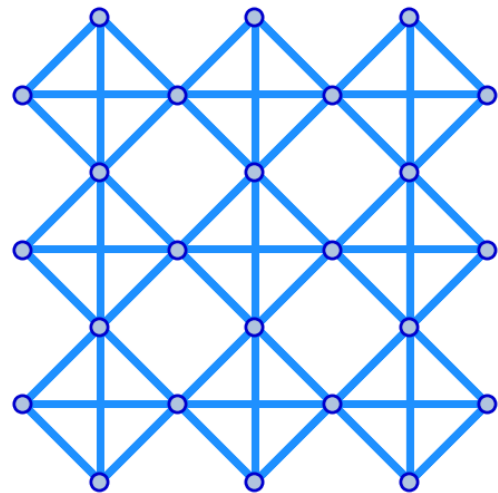
The Checkerboard-Lattice Code

- Built on the square lattice

Fermion Lattice



Anticommutation Relations



Line-Graph Subsystem Codes

Thm: (Chapman and Flammia)

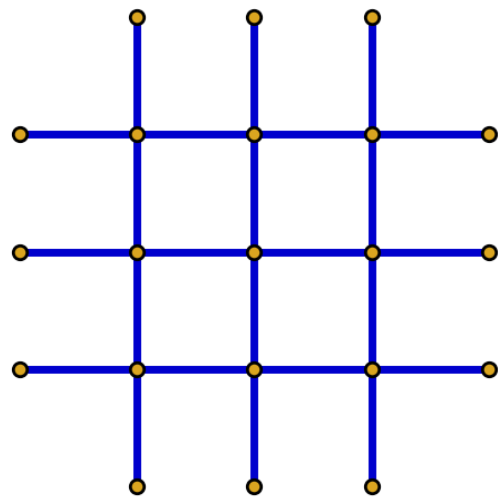
A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

Chapman *et al.* Quantum **4**, 278 (2020)

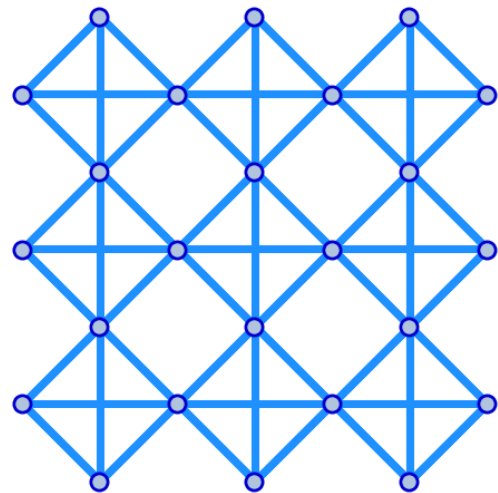
The Checkerboard-Lattice Code

- Built on the square lattice

Fermion Lattice



Anticommutation Relations



- Three Ingredients
 - Two commuting free-fermion models on the square lattice
 - Set of stabilizers

Line-Graph Subsystem Codes

Thm: (Chapman and Flammia)

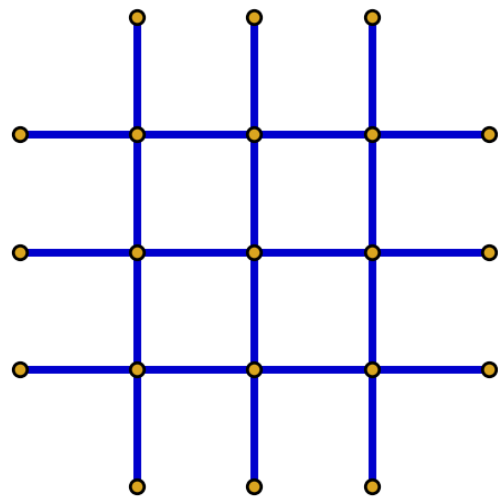
A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

Chapman *et al.* Quantum **4**, 278 (2020)

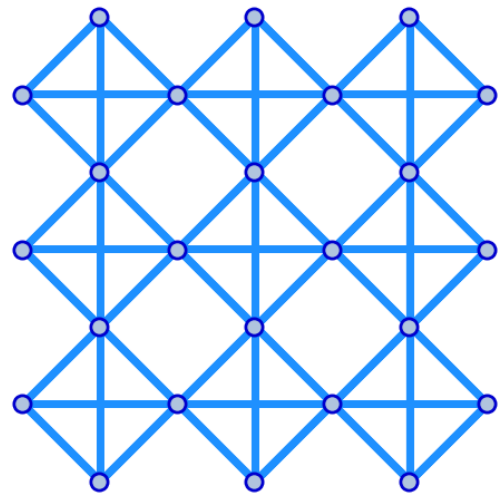
The Checkerboard-Lattice Code

- Built on the square lattice

Fermion Lattice



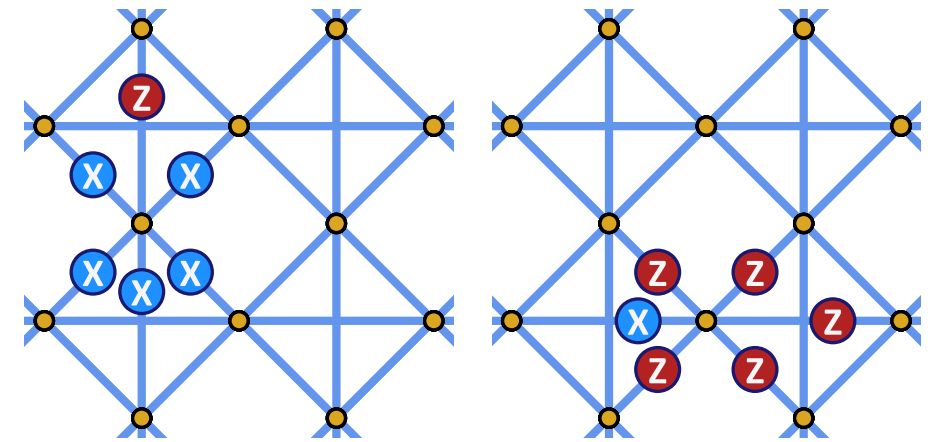
Anticommutation Relations



- Three Ingredients

- Two commuting free-fermion models on the square lattice
- Set of stabilizers

Free-Fermion 1



Line-Graph Subsystem Codes

Thm: (Chapman and Flammia)

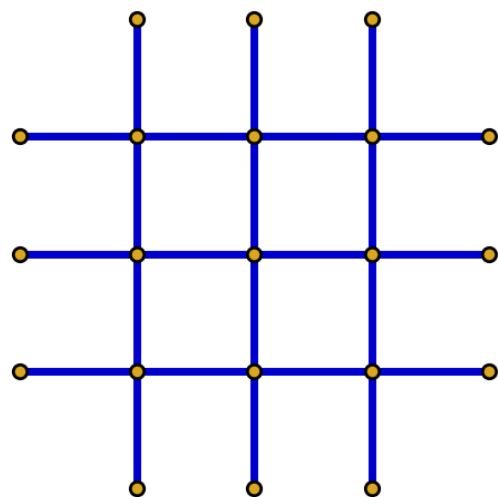
A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

Chapman *et al.* Quantum **4**, 278 (2020)

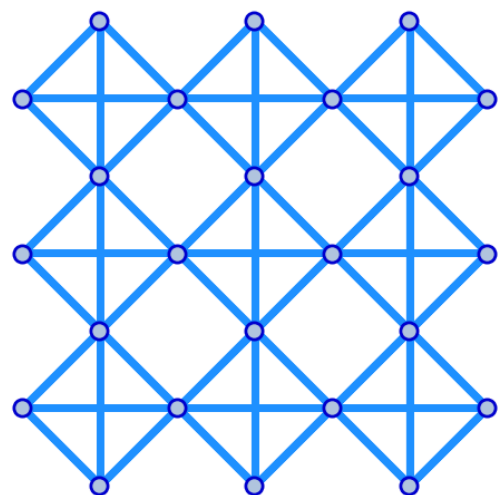
The Checkerboard-Lattice Code

- Built on the square lattice

Fermion Lattice



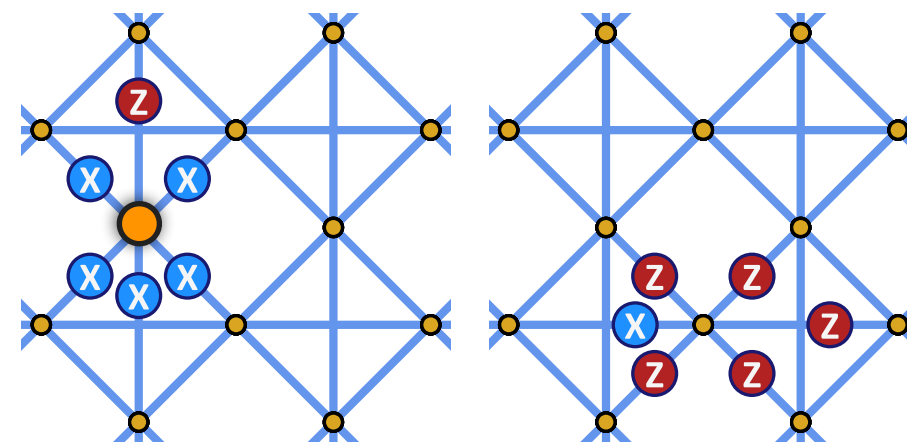
Anticommutation Relations



- Three Ingredients

- Two commuting free-fermion models on the square lattice
- Set of stabilizers

Free-Fermion 1



Line-Graph Subsystem Codes

Thm: (Chapman and Flammia)

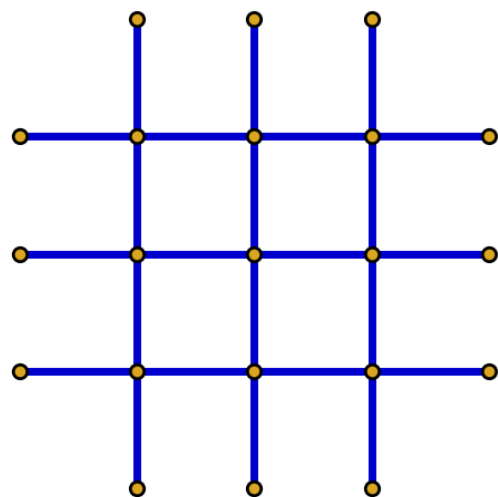
A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

Chapman *et al.* Quantum **4**, 278 (2020)

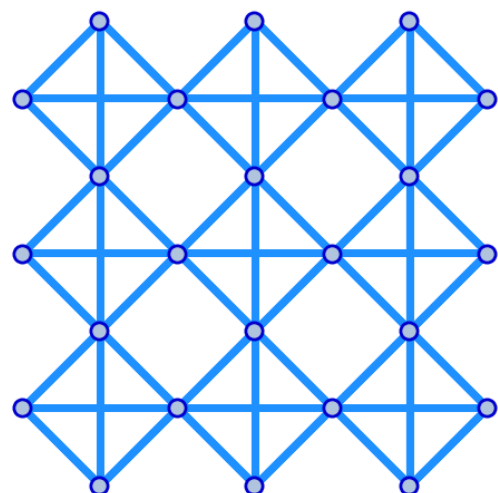
The Checkerboard-Lattice Code

- Built on the square lattice

Fermion Lattice



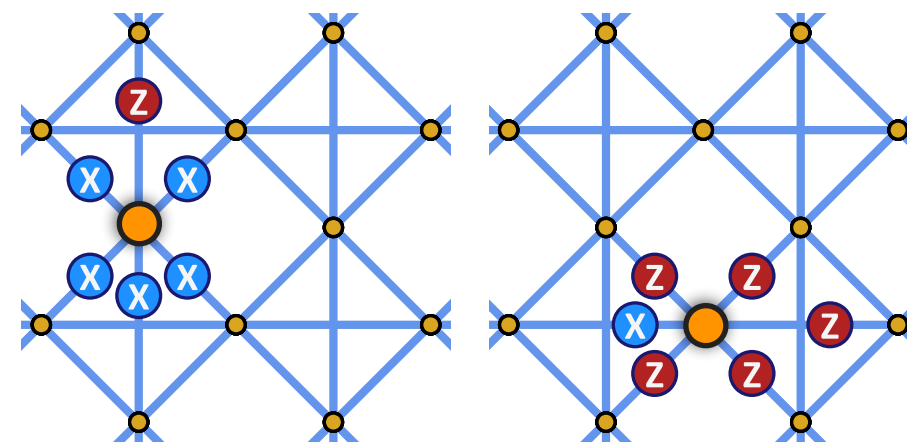
Anticommutation Relations



- Three Ingredients

- Two commuting free-fermion models on the square lattice
- Set of stabilizers

Free-Fermion 1



Line-Graph Subsystem Codes

Thm: (Chapman and Flammia)

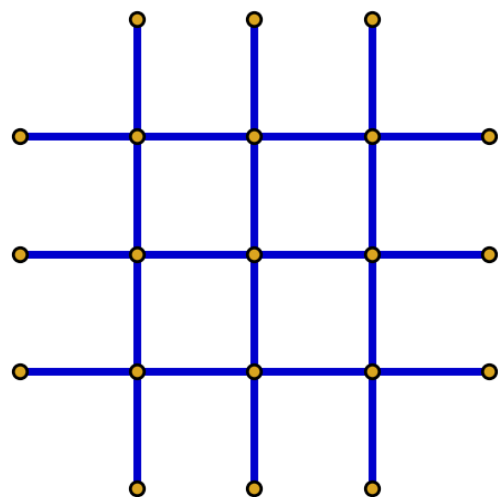
A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

Chapman *et al.* Quantum **4**, 278 (2020)

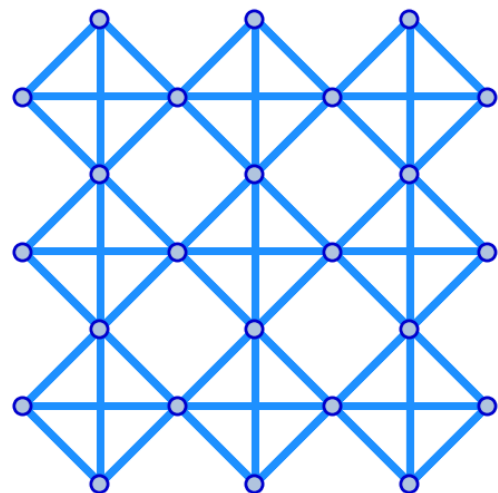
The Checkerboard-Lattice Code

- Built on the square lattice

Fermion Lattice



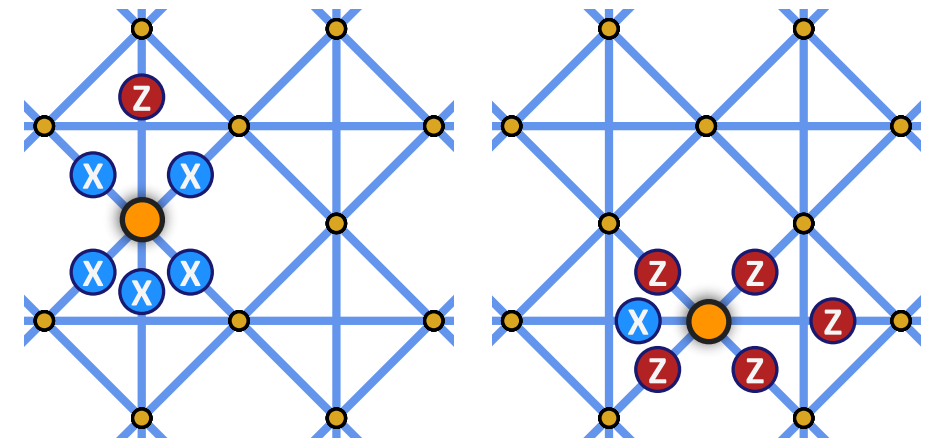
Anticommutation Relations



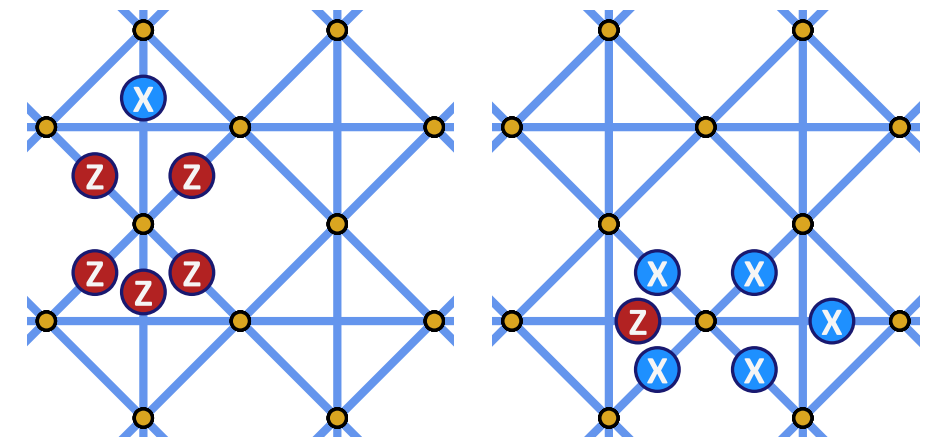
- Three Ingredients

- Two commuting free-fermion models on the square lattice
- Set of stabilizers

Free-Fermion 1



Free-Fermion 2



Line-Graph Subsystem Codes

Thm: (Chapman and Flammia)

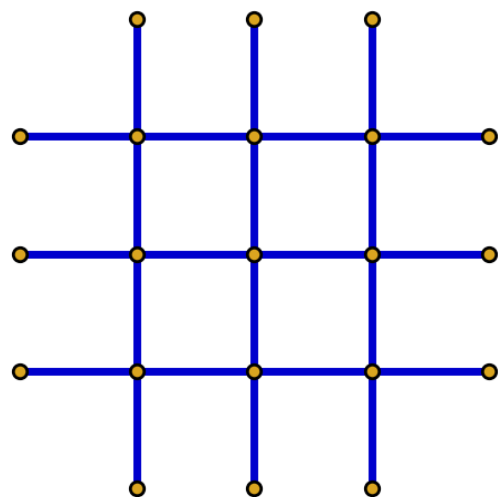
A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

Chapman *et al.* Quantum **4**, 278 (2020)

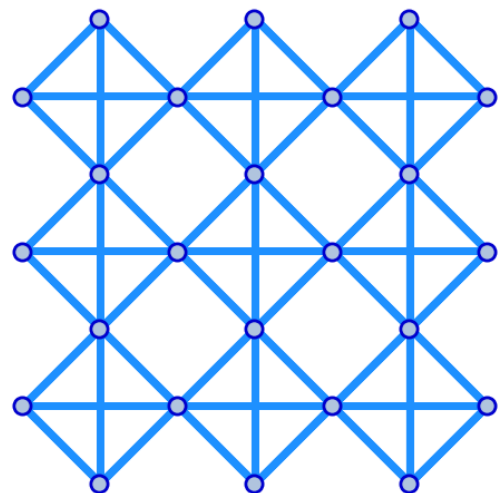
The Checkerboard-Lattice Code

- Built on the square lattice

Fermion Lattice



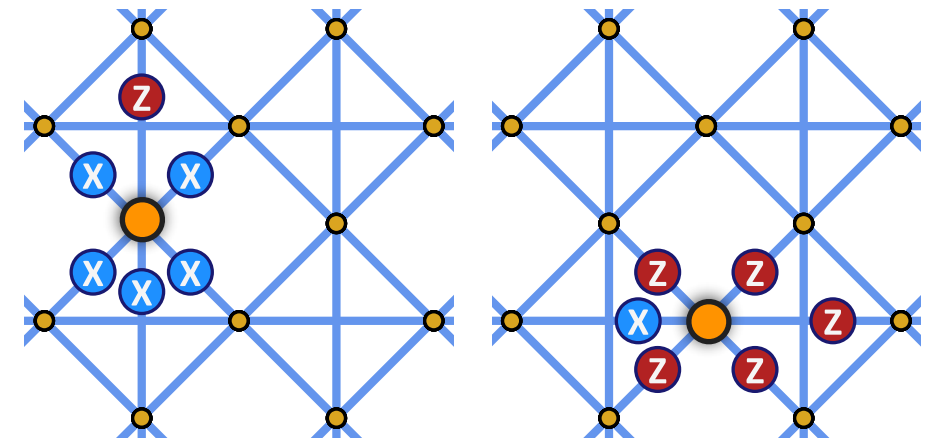
Anticommutation Relations



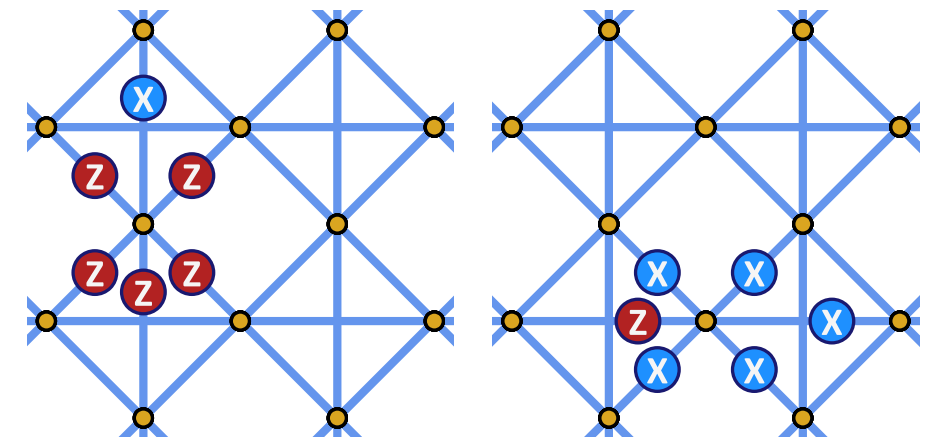
- Three Ingredients

- Two commuting free-fermion models on the square lattice
- Set of stabilizers

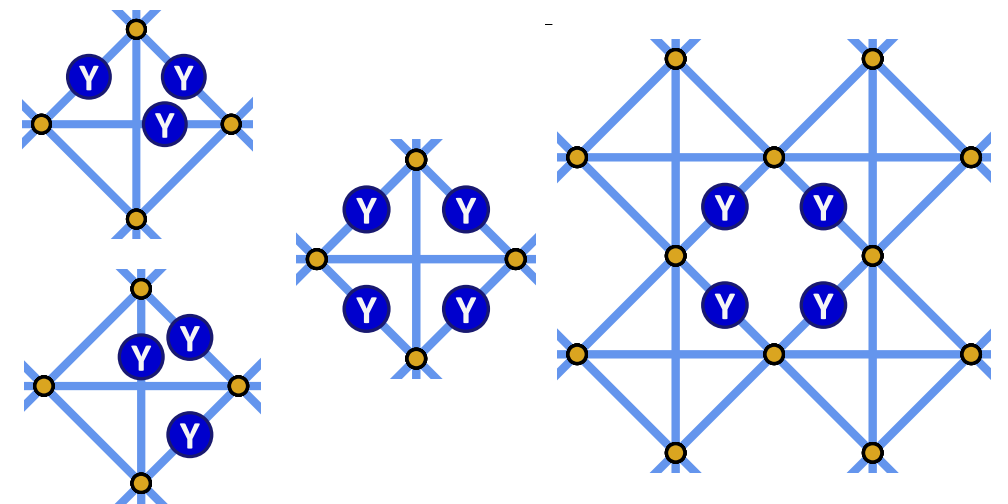
Free-Fermion 1



Free-Fermion 2

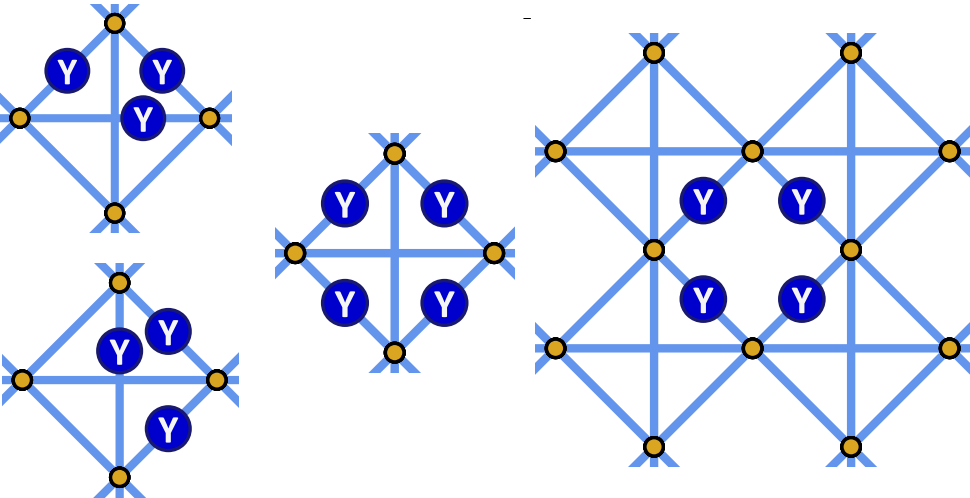
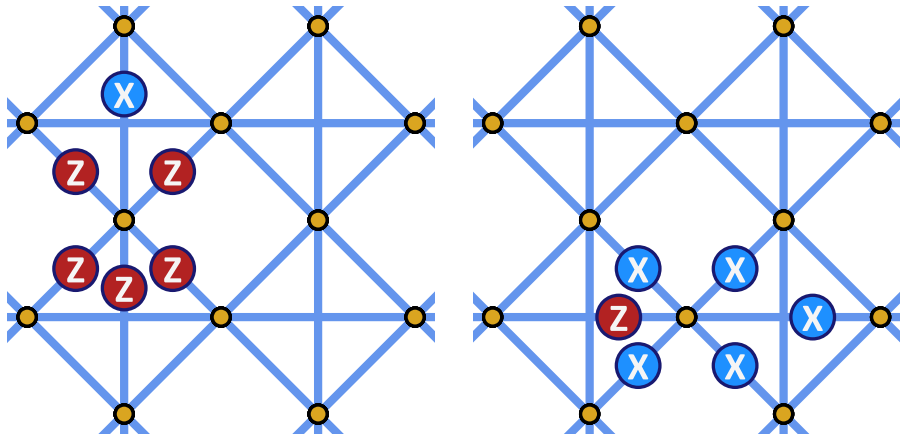
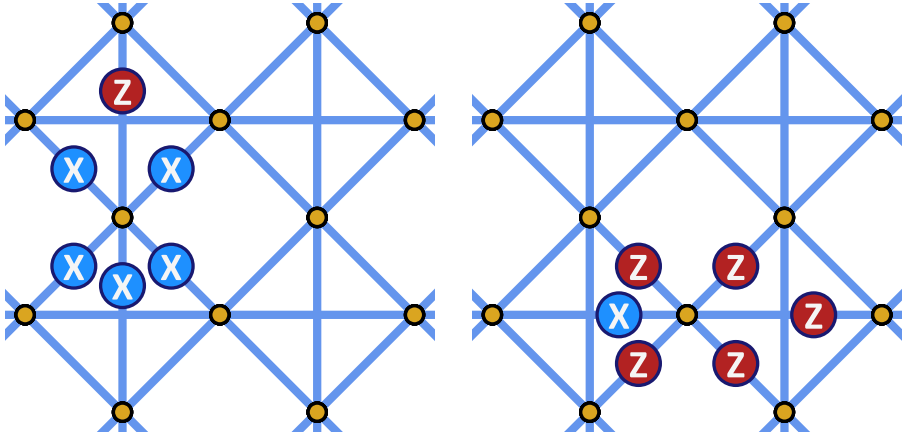


Stabilizers



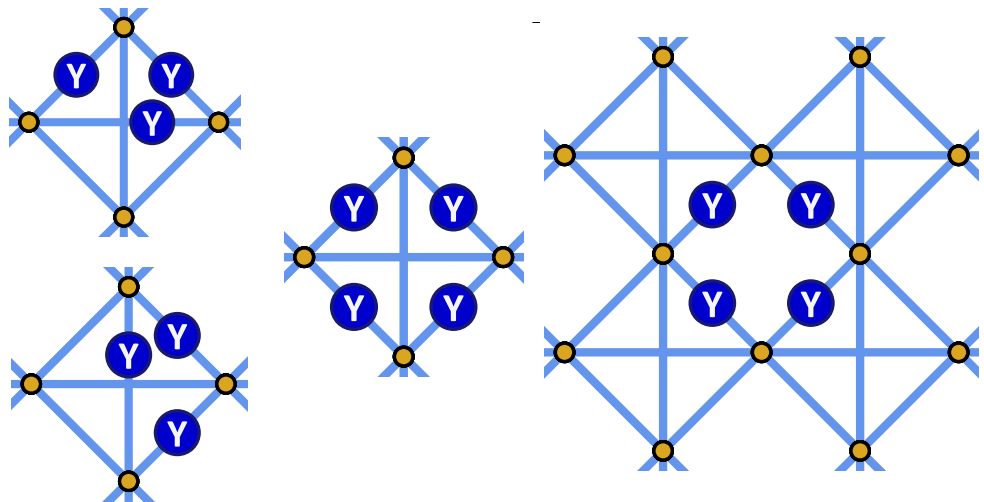
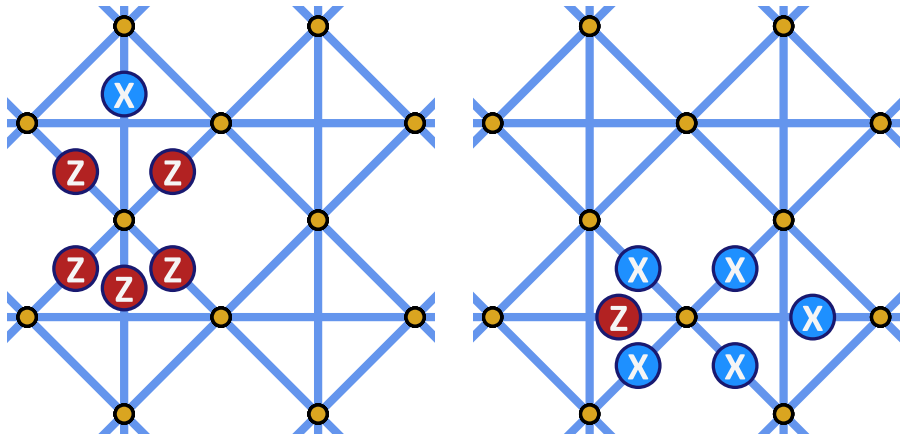
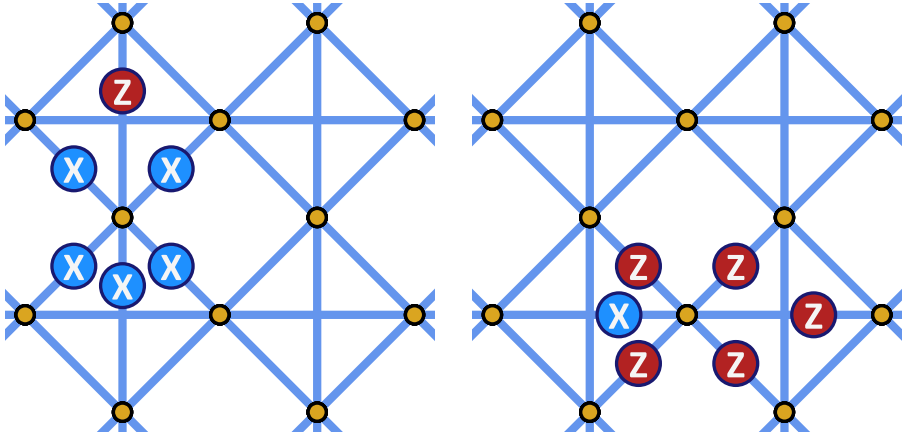
Checkerboard Lattice Code

Hamiltonian Terms/Gauge Generators

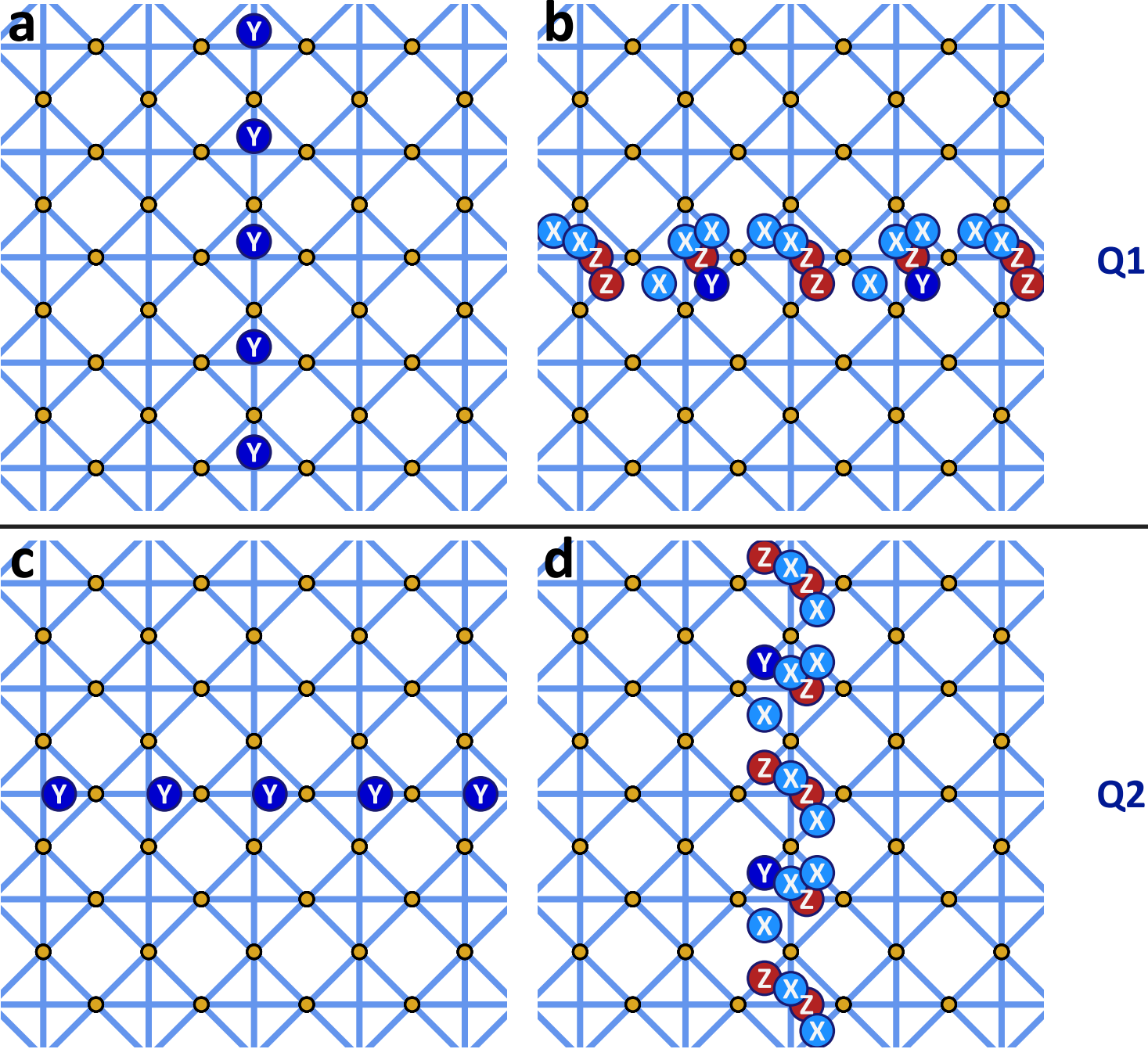


Checkerboard Lattice Code

Hamiltonian Terms/Gauge Generators



Exact Logical Qubits



Outline

- Coplanar Waveguide (CPW) Lattices
 - Deformable lattice sites
 - Line-graph lattices
 - Interacting photons
- Band Engineering
 - Hyperbolic lattice
 - Gapped flat bands
- Mathematical Connections
 - Planar Gaps
 - Maximal Gaps
 - Quantum Error Correcting Codes
- Experimental Developments

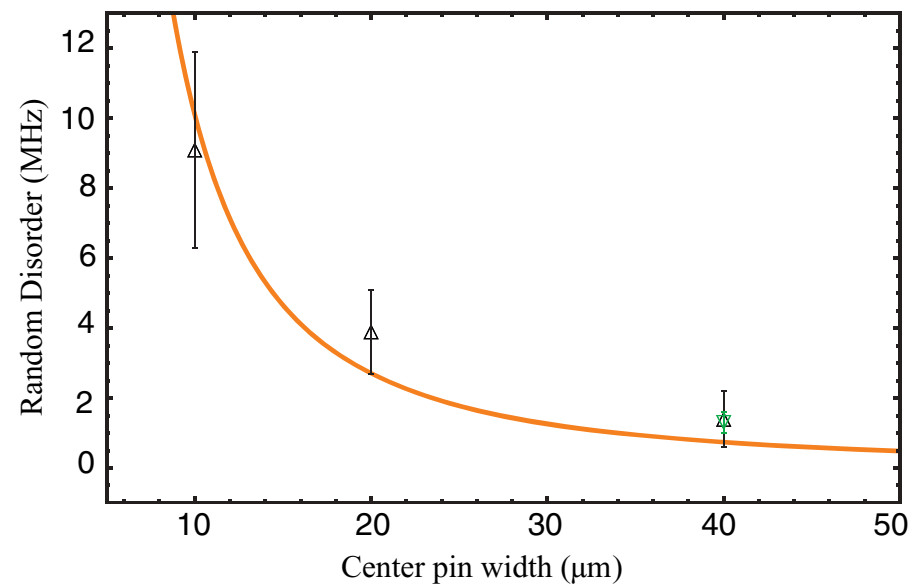
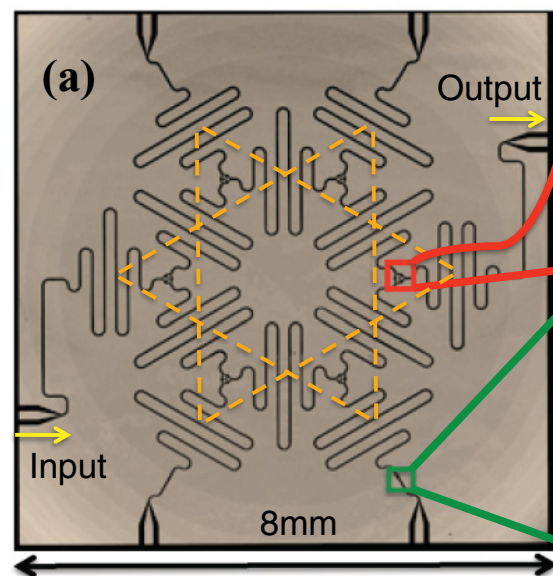
Outline

- Coplanar Waveguide (CPW) Lattices
 - Deformable lattice sites
 - Line-graph lattices
 - Interacting photons
- Band Engineering
 - Hyperbolic lattice
 - Gapped flat bands
- Mathematical Connections
 - Planar Gaps
 - Maximal Gaps
 - Quantum Error Correcting Codes
- Experimental Developments

Intrinsic Fabrication Disorder

Previous Benchmarks

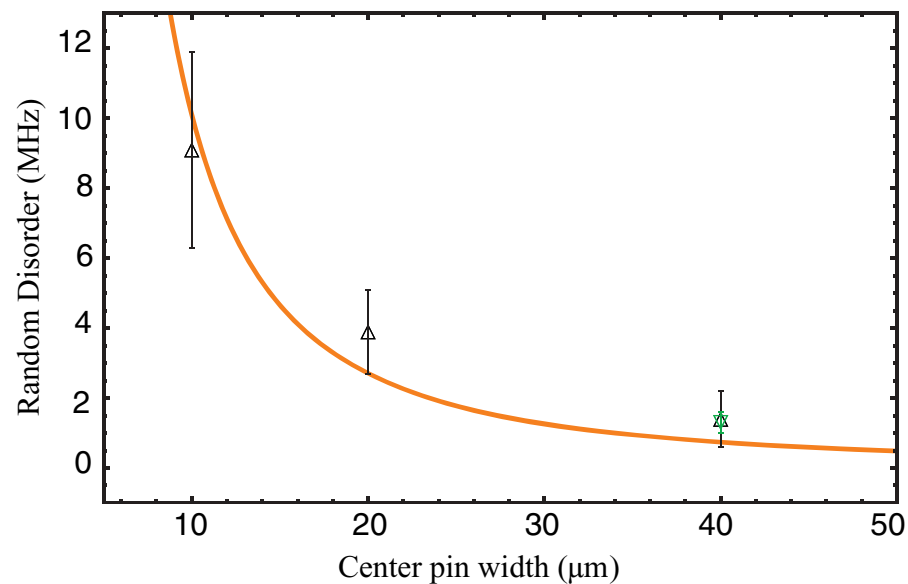
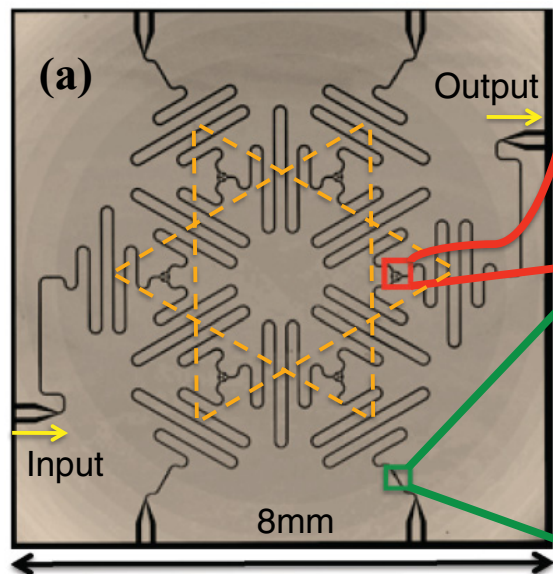
- Kagome star normal modes
- Fabricated at Princeton
- Fabrication disorder $\sim 3e-4$



Intrinsic Fabrication Disorder

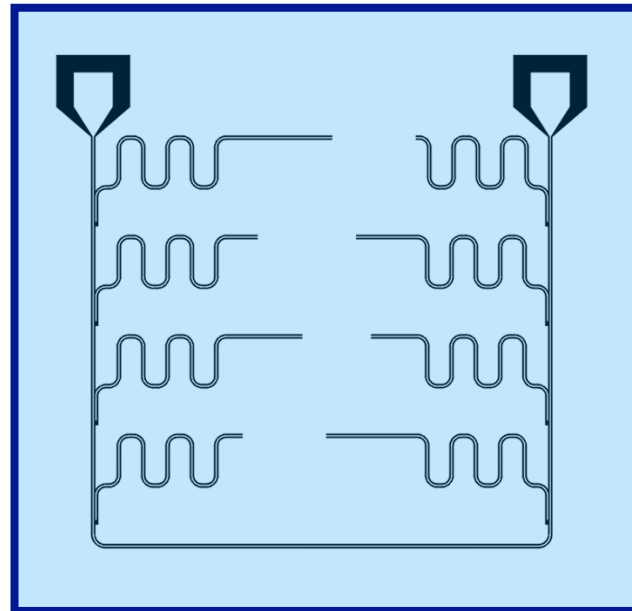
Previous Benchmarks

- Kagome star normal modes
- Fabricated at Princeton
- Fabrication disorder $\sim 3e-4$



Current Devices

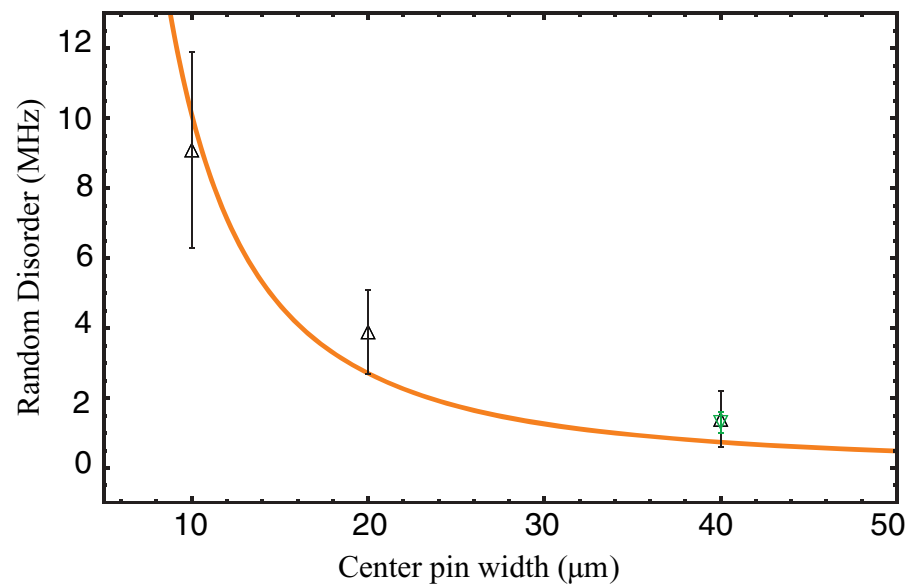
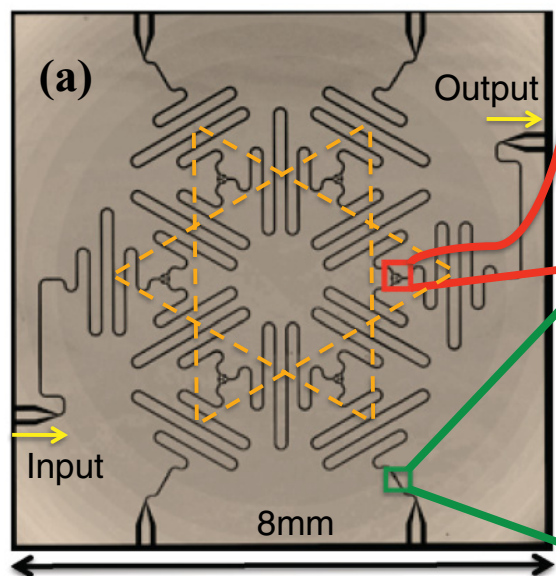
- Fabricated at UMD
- Fabrication disorder $\sim 3e-4$



Intrinsic Fabrication Disorder

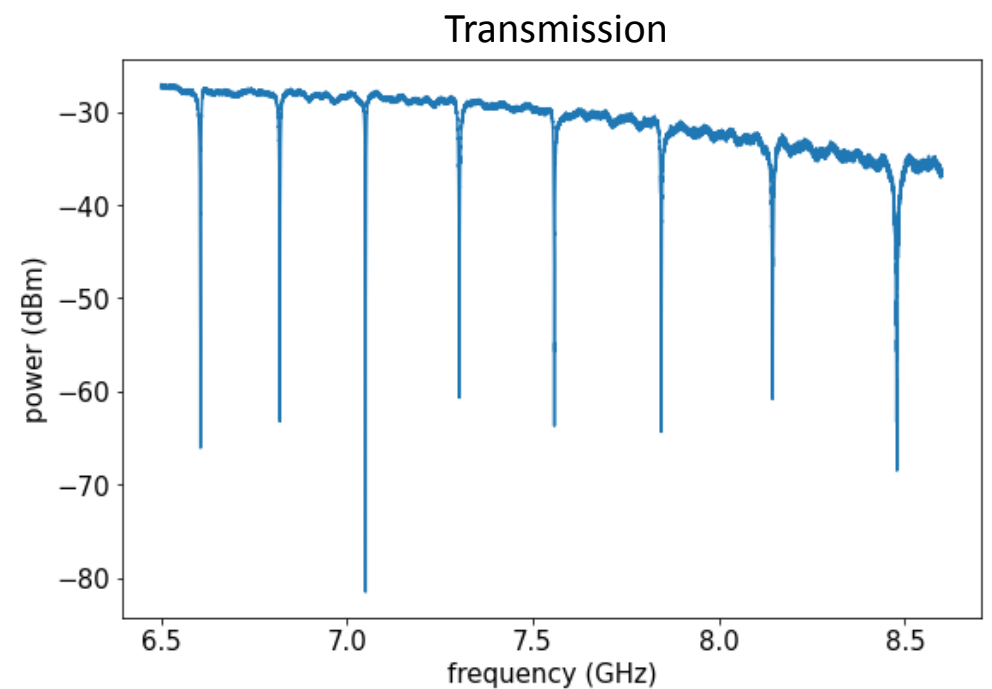
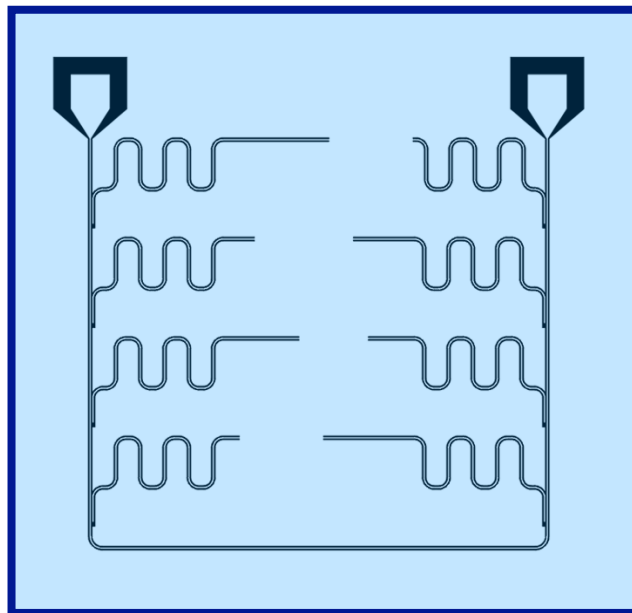
Previous Benchmarks

- Kagome star normal modes
- Fabricated at Princeton
- Fabrication disorder $\sim 3e-4$



Current Devices

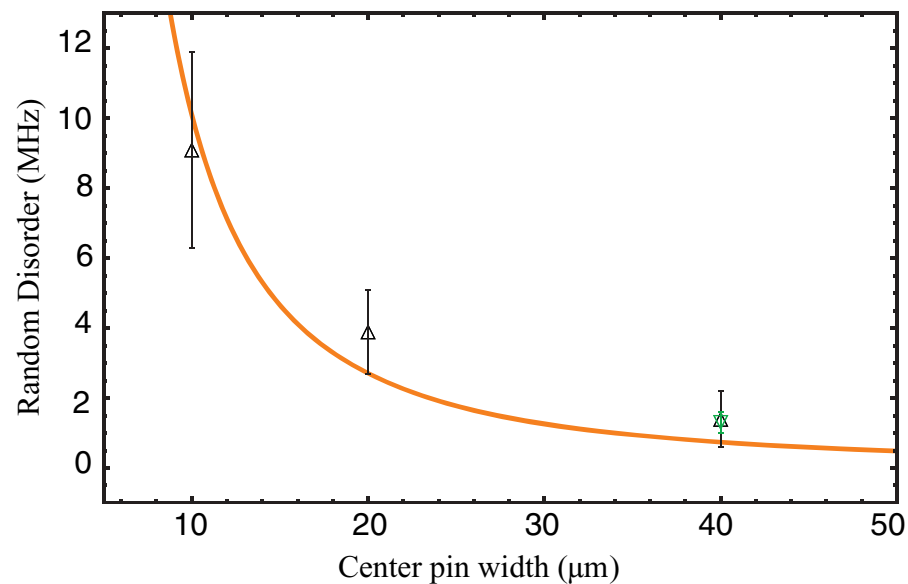
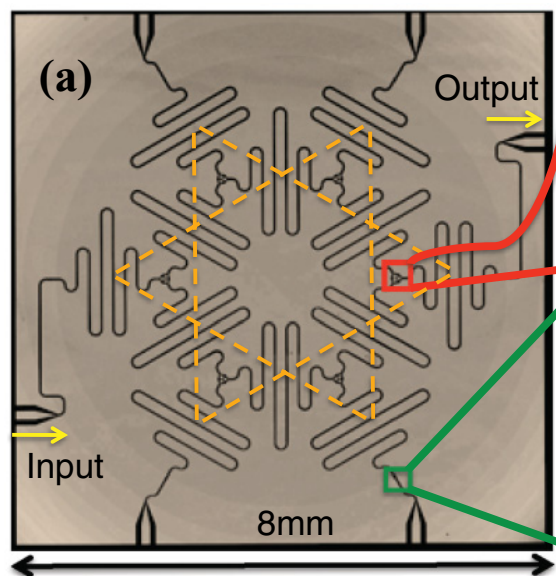
- Fabricated at UMD
- Fabrication disorder $\sim 3e-4$



Intrinsic Fabrication Disorder

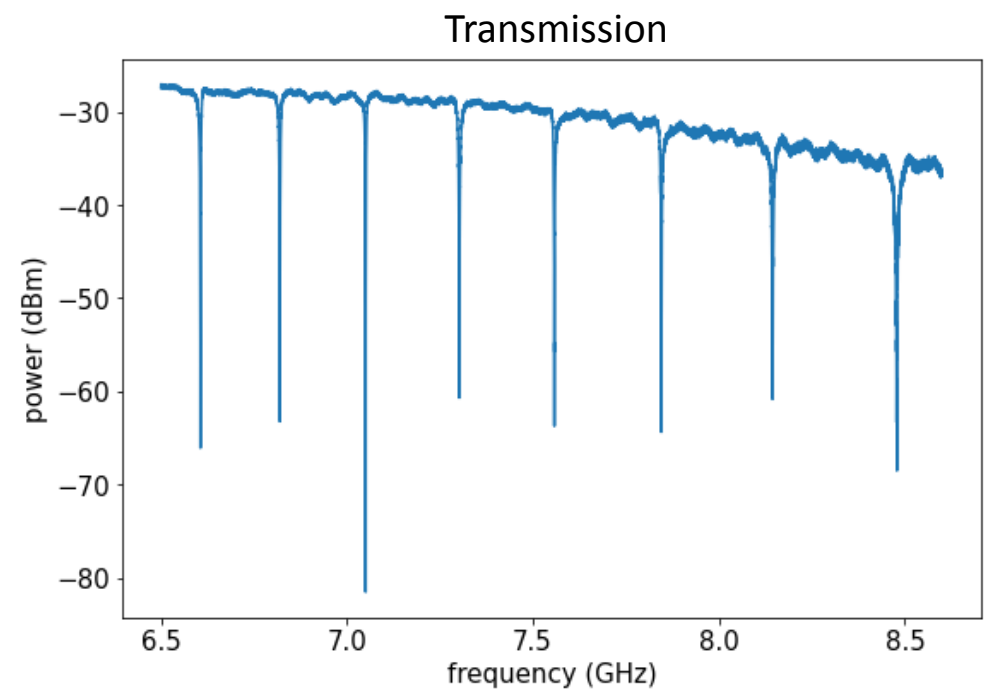
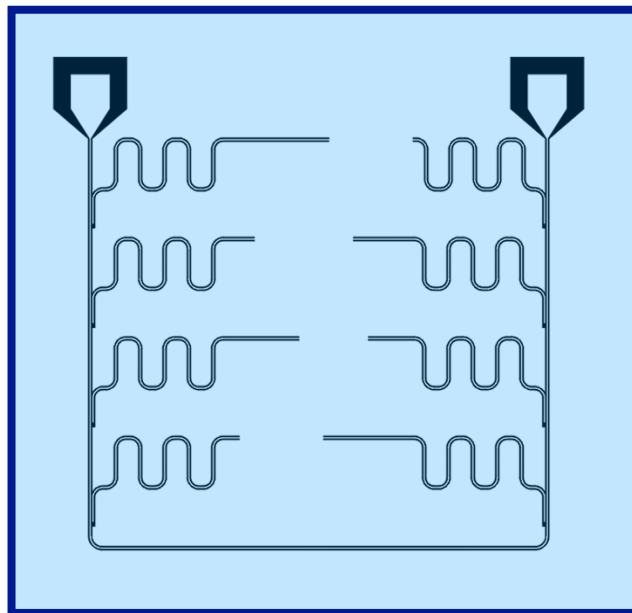
Previous Benchmarks

- Kagome star normal modes
- Fabricated at Princeton
- Fabrication disorder $\sim 3e-4$



Current Devices

- Fabricated at UMD
- Fabrication disorder $\sim 3e-4$

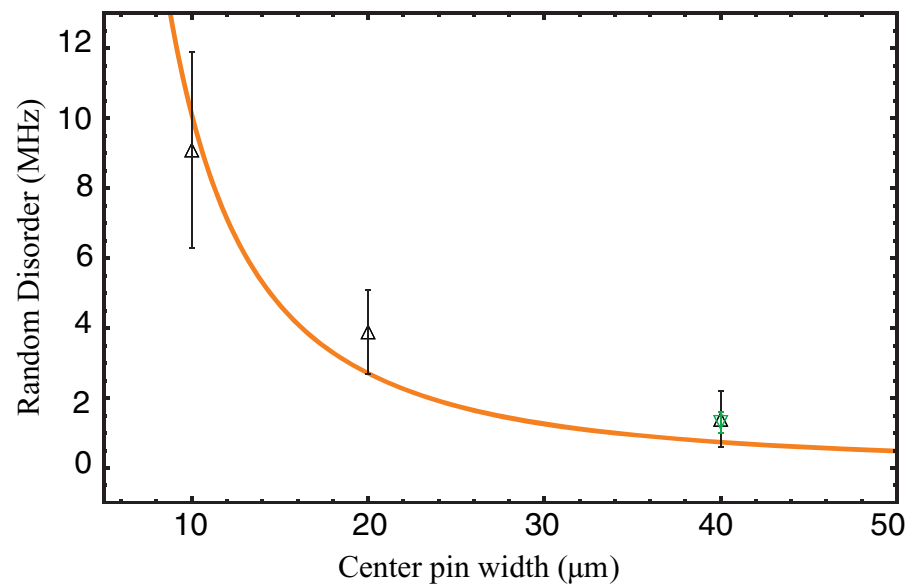
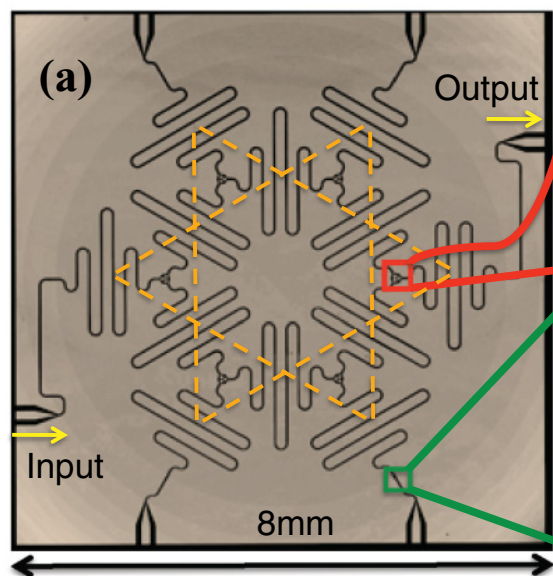


- Parallel measurement

Intrinsic Fabrication Disorder

Previous Benchmarks

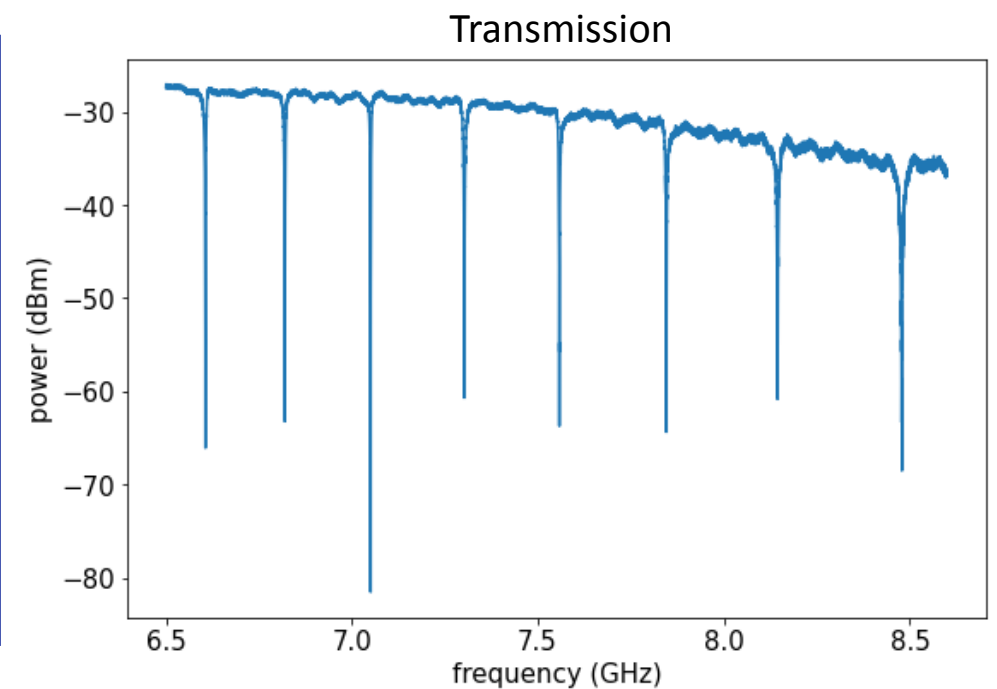
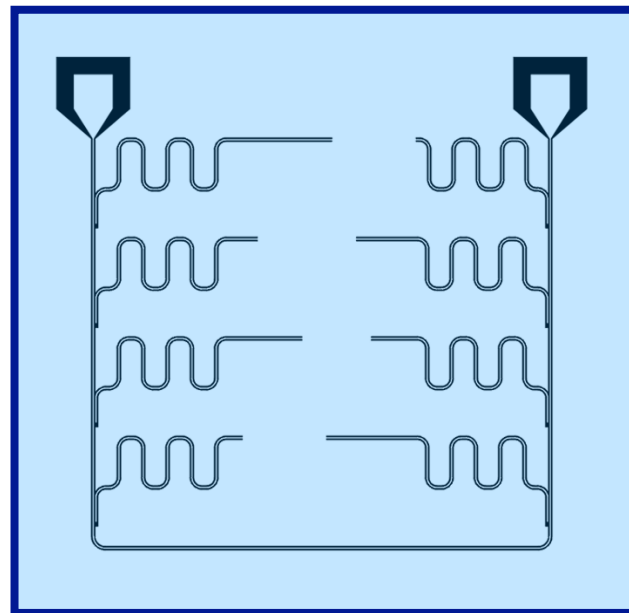
- Kagome star normal modes
- Fabricated at Princeton
- Fabrication disorder $\sim 3e-4$



Underwood *et al.* PRA **86**, 023837 (2012)

Current Devices

- Fabricated at UMD
- Fabrication disorder $\sim 3e-4$

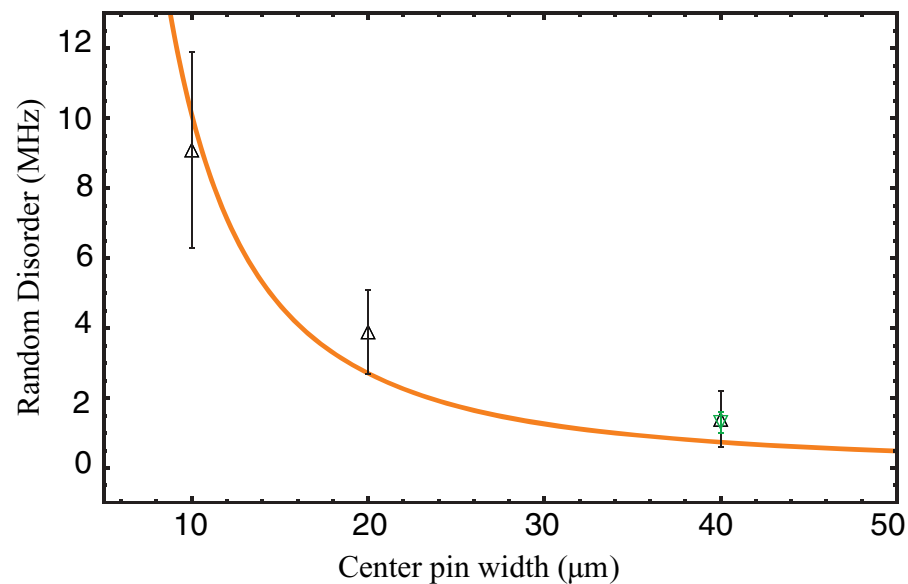
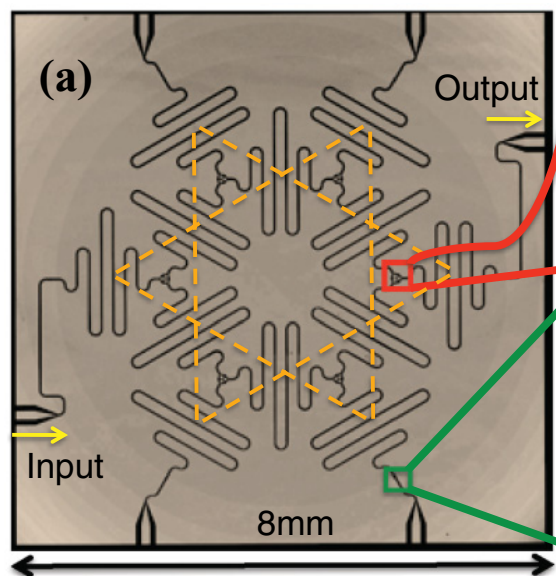


- Parallel measurement
- Disorder extracted from comb spacing

Intrinsic Fabrication Disorder

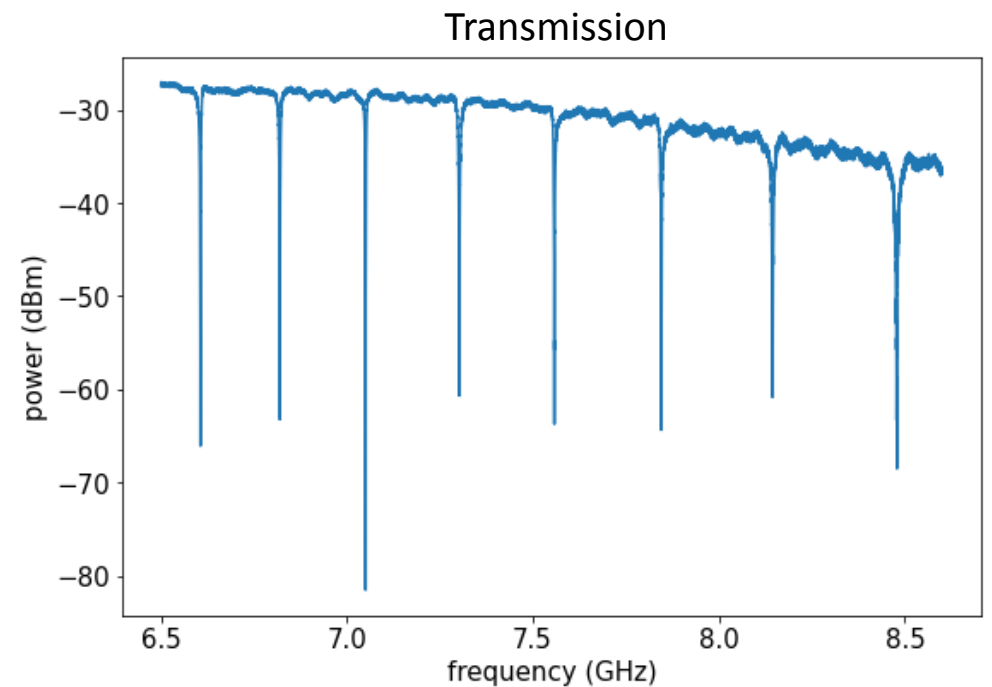
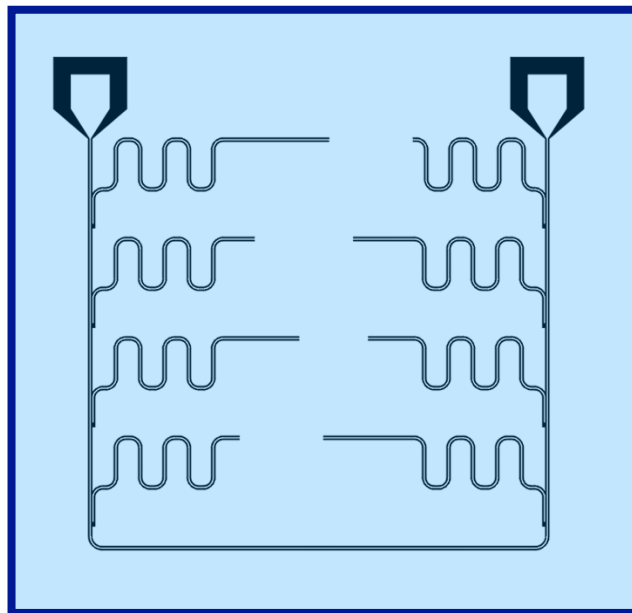
Previous Benchmarks

- Kagome star normal modes
- Fabricated at Princeton
- Fabrication disorder $\sim 3e-4$

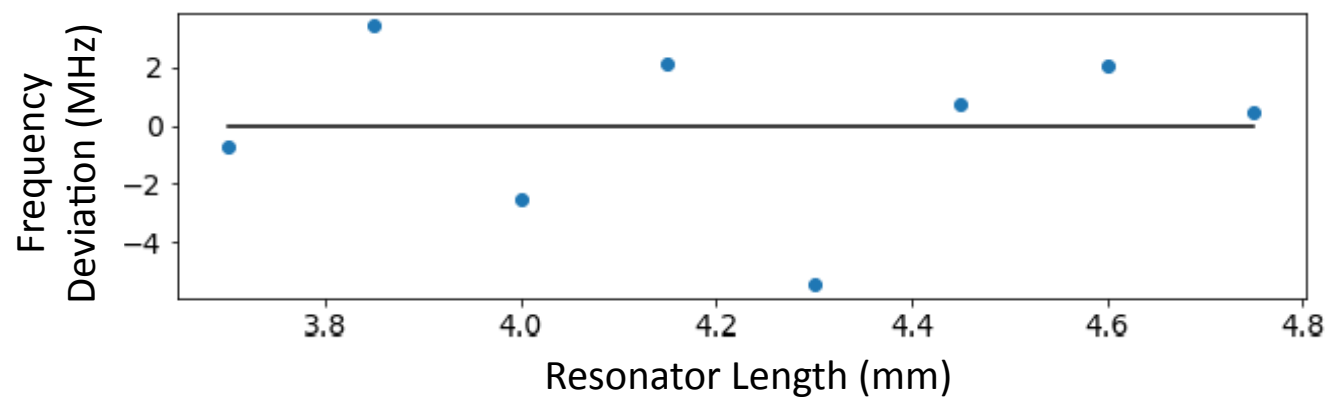


Current Devices

- Fabricated at UMD
- Fabrication disorder $\sim 3e-4$



- Parallel measurement
- Disorder extracted from comb spacing



Disorder Mitigation

Systematic v. Random Disorder

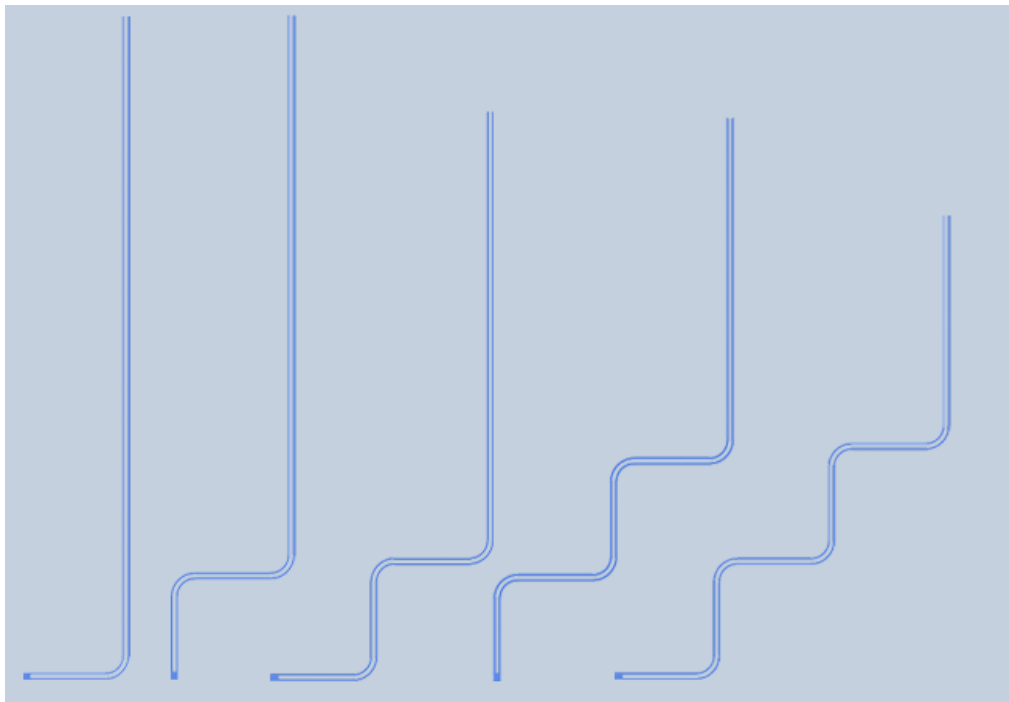
- Fabrication disorder $\sim 3e-4$
- Shape-dependent disorder $\sim 2-3e-3$

Disorder Mitigation

Systematic v. Random Disorder

- Fabrication disorder $\sim 3e-4$
- Shape-dependent disorder $\sim 2-3e-3$

Numerical Test Geometries

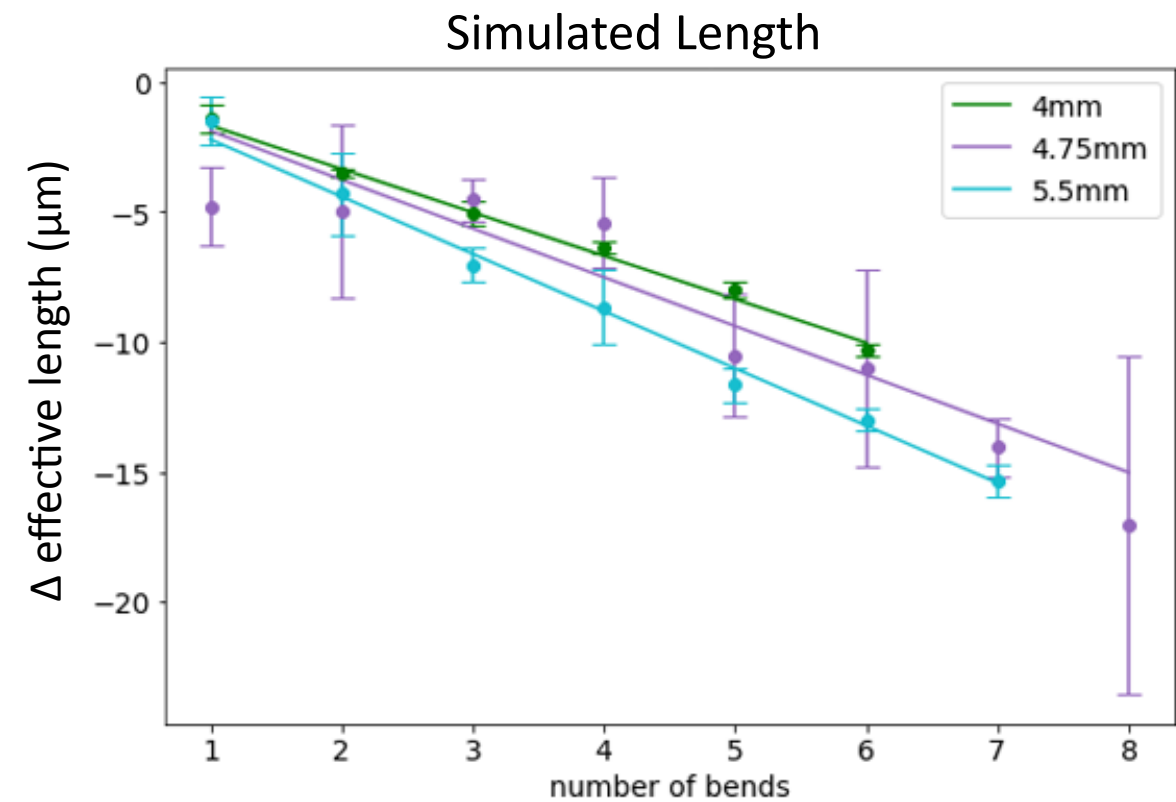
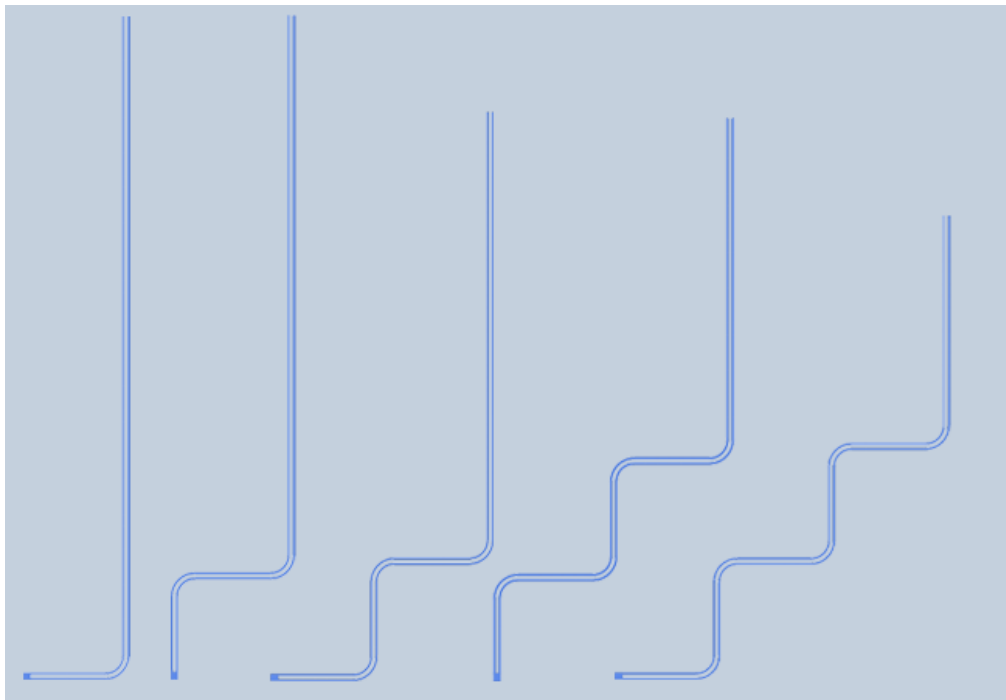


Disorder Mitigation

Systematic v. Random Disorder

- Fabrication disorder $\sim 3e-4$
- Shape-dependent disorder $\sim 2-3e-3$

Numerical Test Geometries

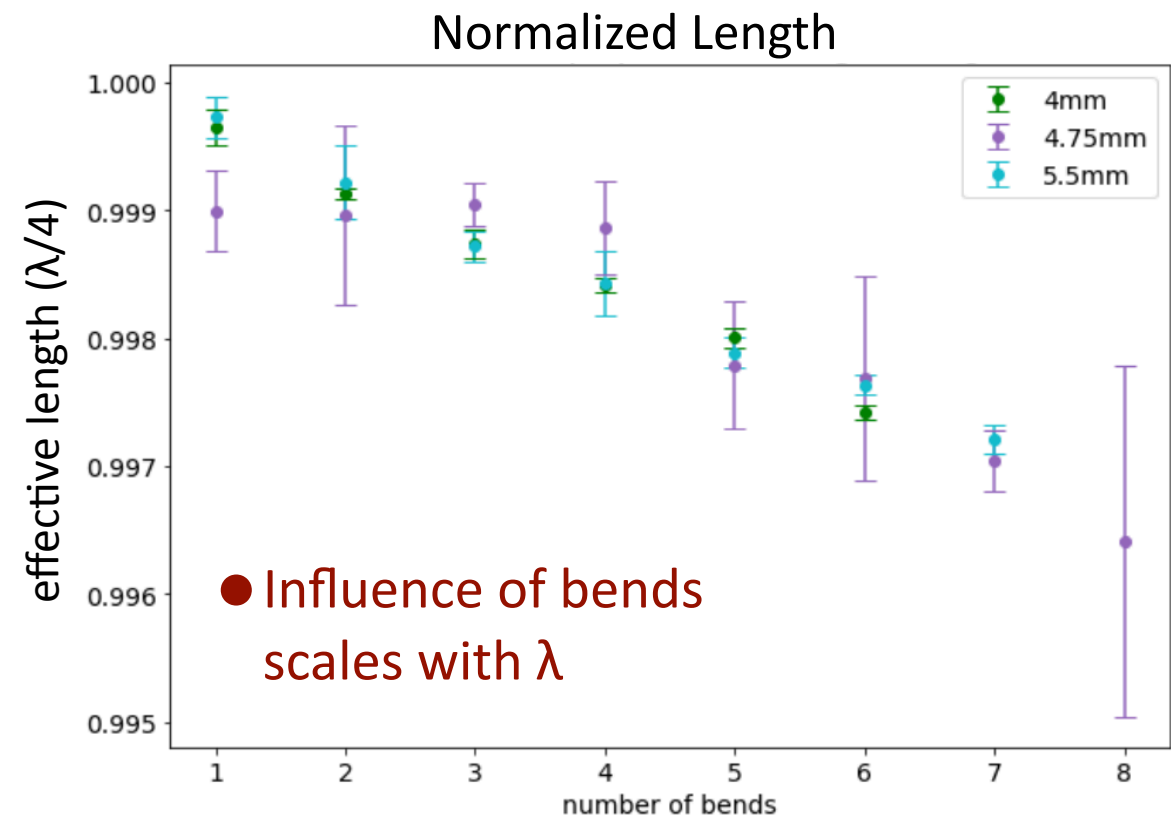
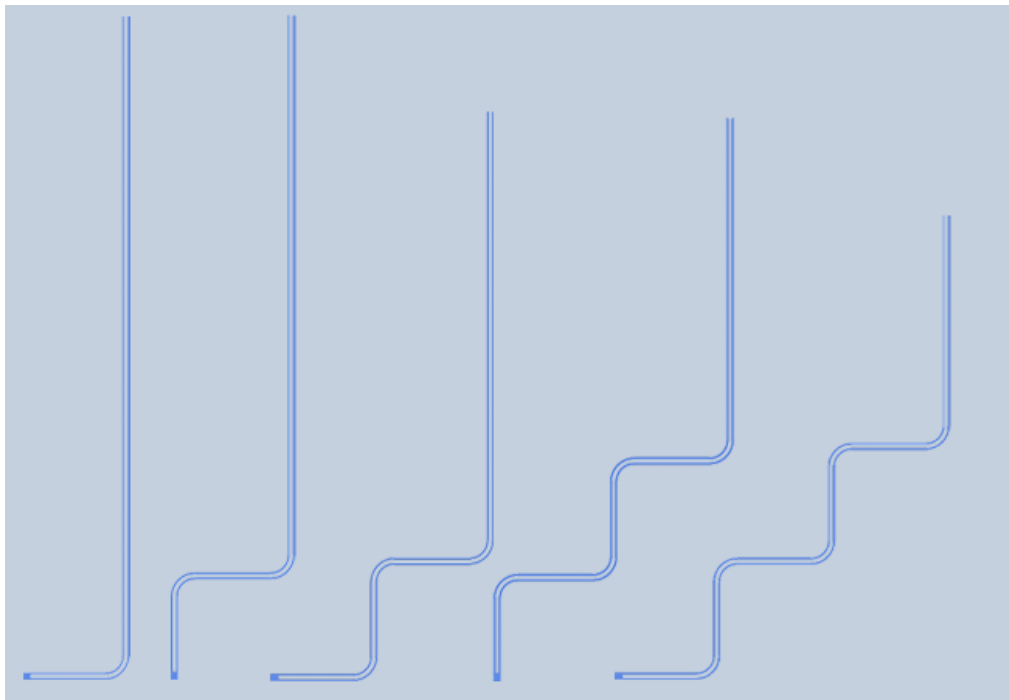


Disorder Mitigation

Systematic v. Random Disorder

- Fabrication disorder $\sim 3e-4$
- Shape-dependent disorder $\sim 2-3e-3$

Numerical Test Geometries

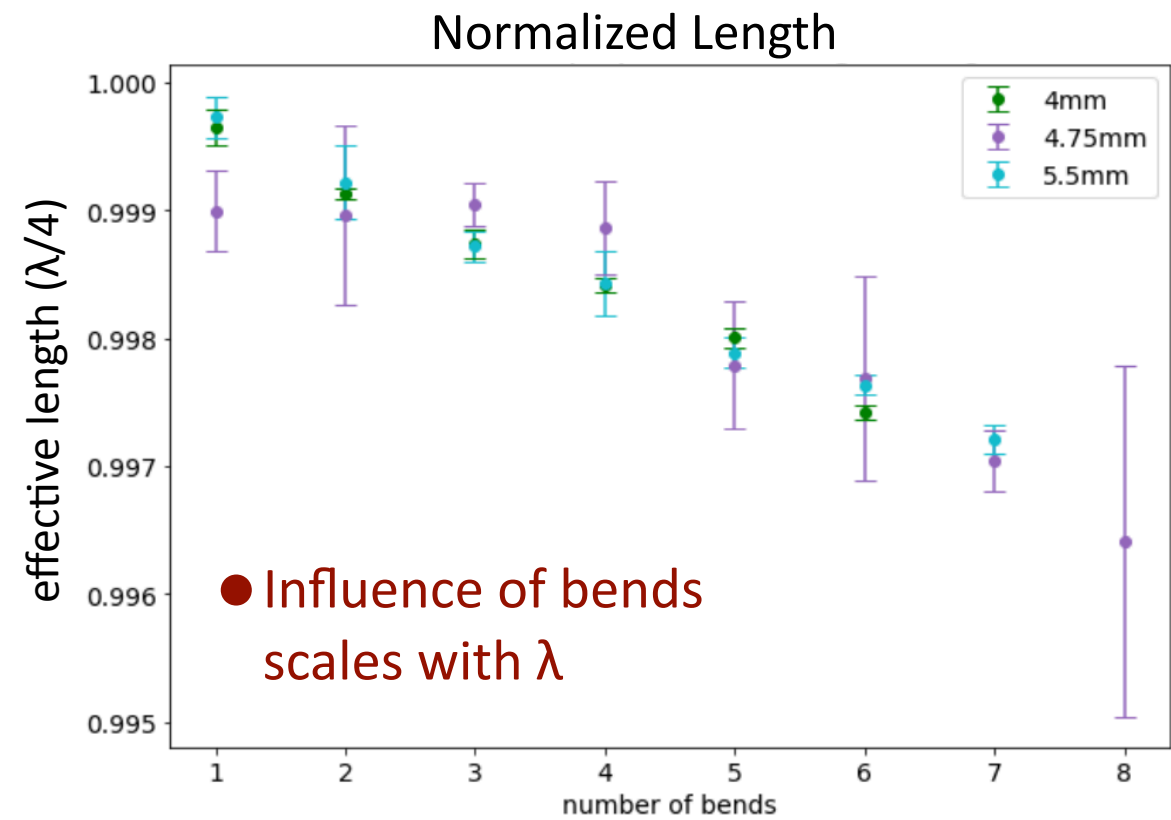
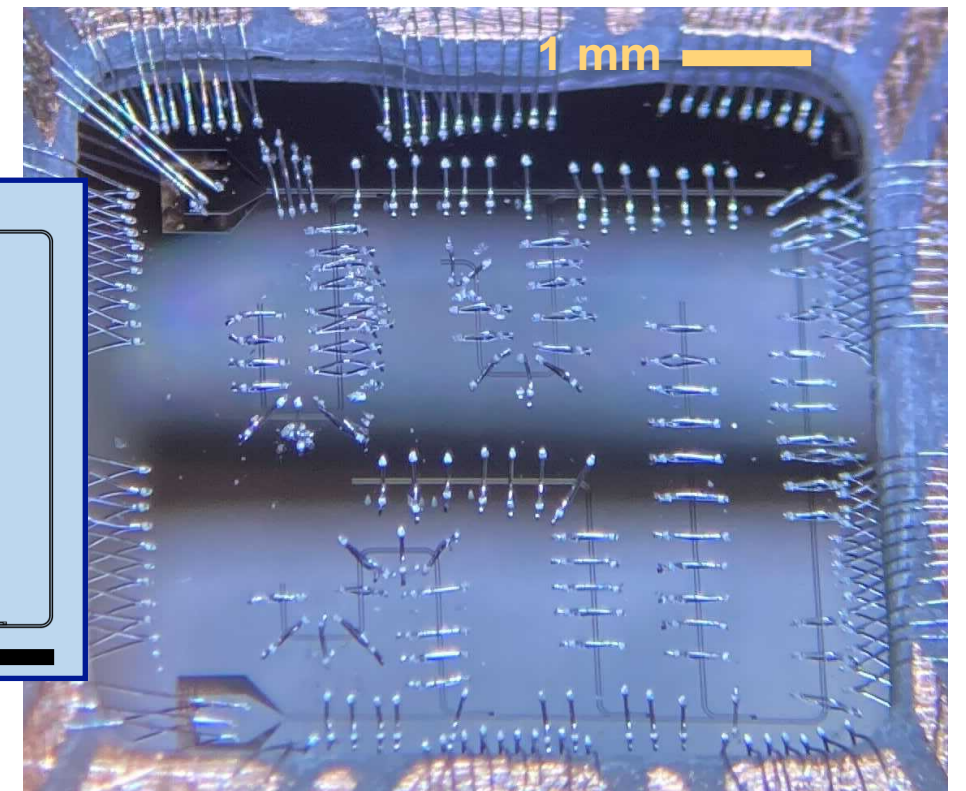
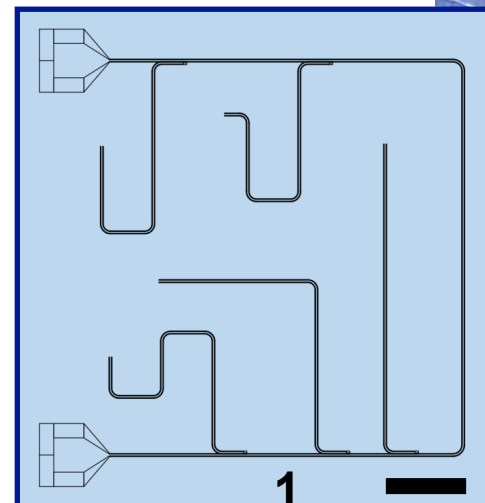
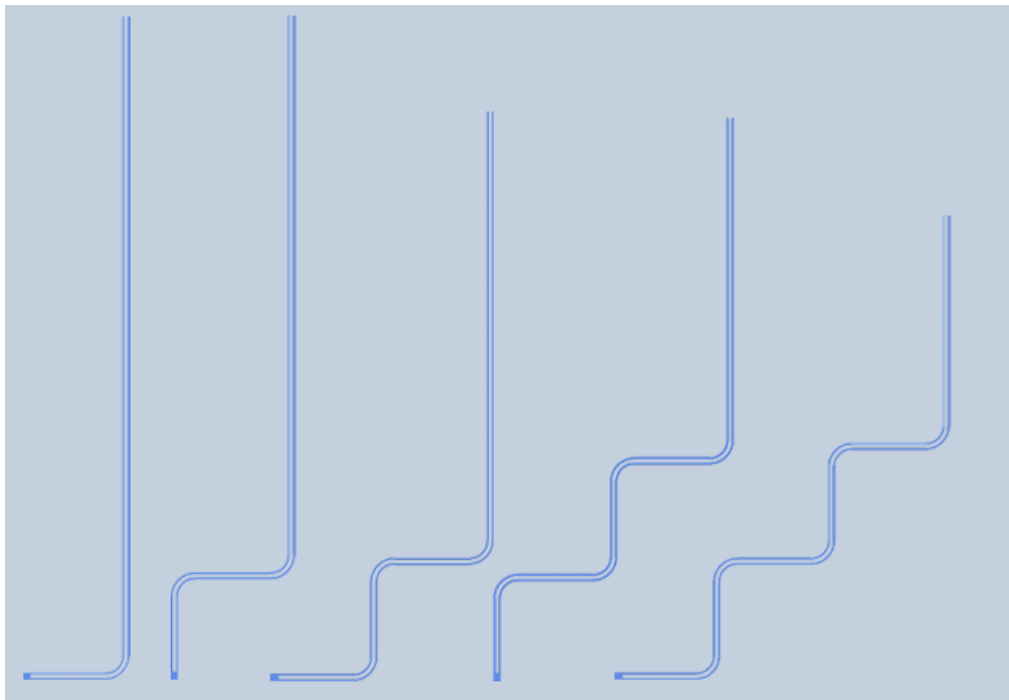


Disorder Mitigation

Systematic v. Random Disorder

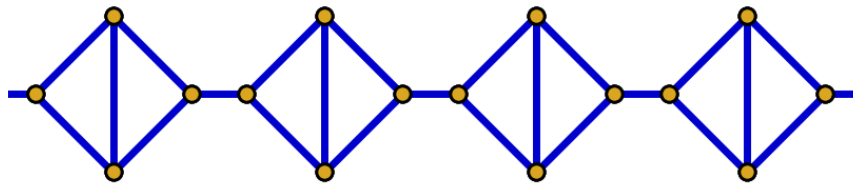
- Fabrication disorder $\sim 3e-4$
- Shape-dependent disorder $\sim 2-3e-3$

Numerical Test Geometries

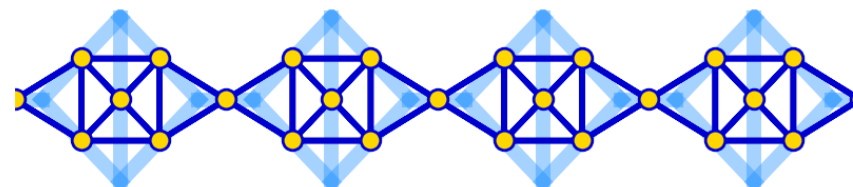


Quasi-1D Lattice Device

Hardware Layout

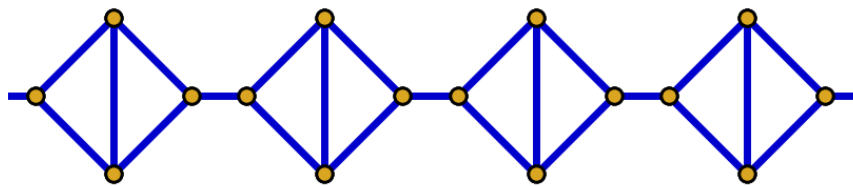


Effective Lattice

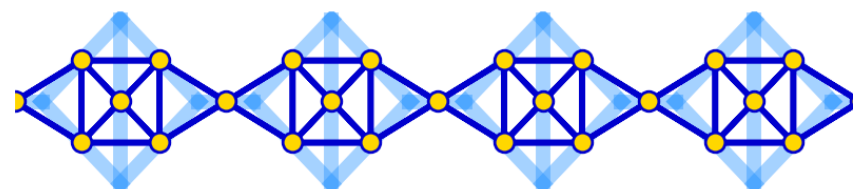


Quasi-1D Lattice Device

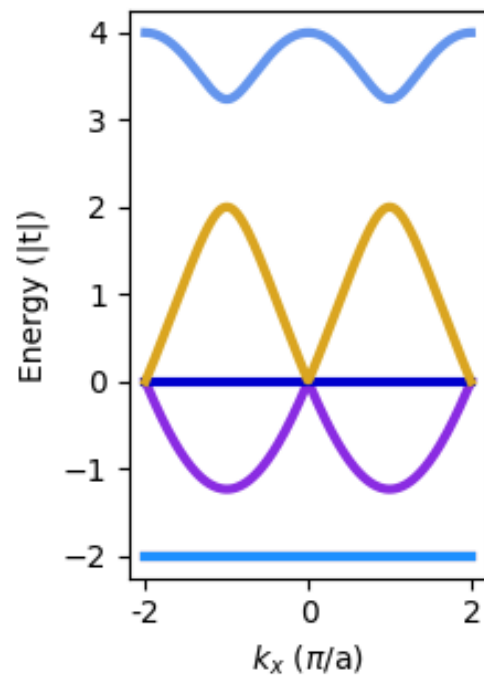
Hardware Layout



Effective Lattice



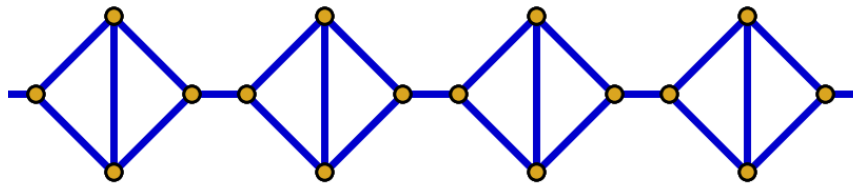
Band Structure



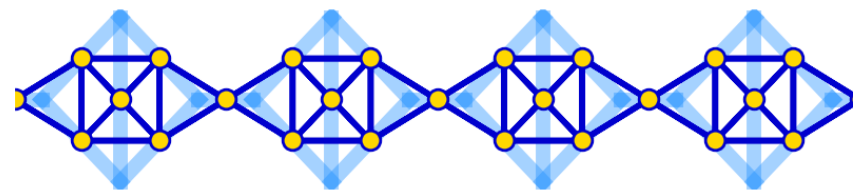
- Flat bands
 - Gapped
 - Ungapped
- Linear bands
- Quadratic bands

Quasi-1D Lattice Device

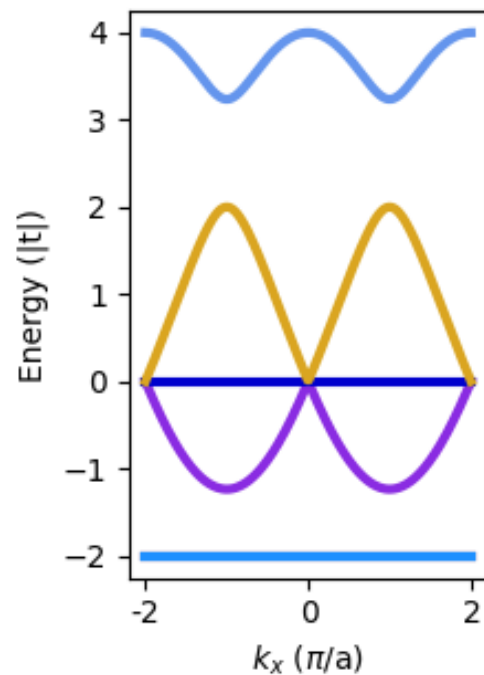
Hardware Layout



Effective Lattice



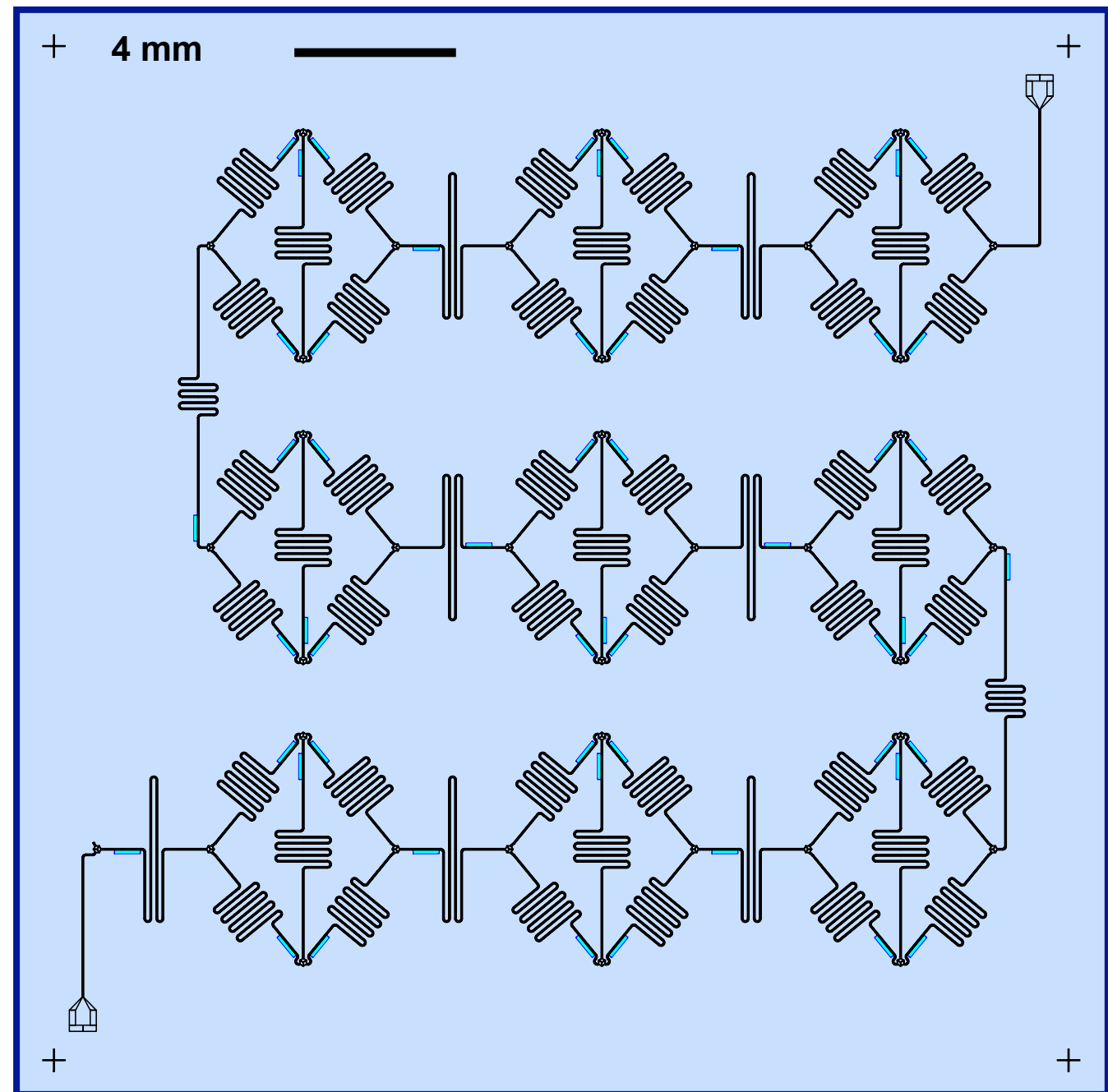
Band Structure



- Flat bands
- Gapped
- Ungapped
- Linear bands
- Quadratic bands

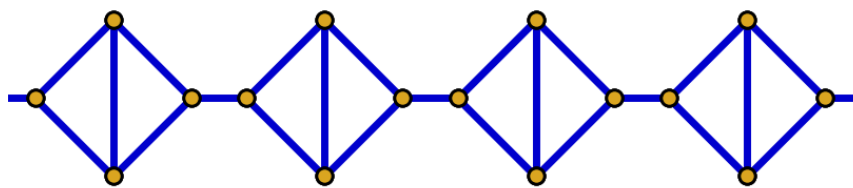
Device Design

(preliminary)

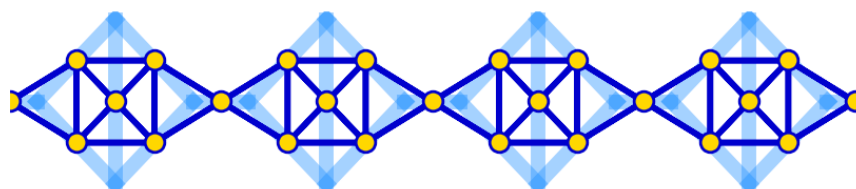


Quasi-1D Lattice Device

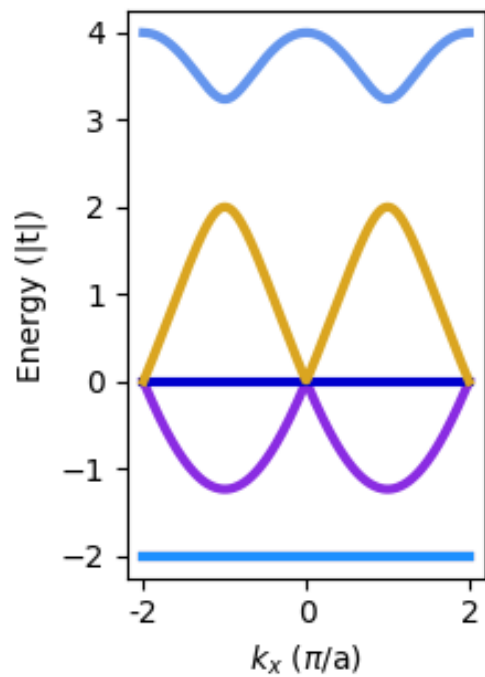
Hardware Layout



Effective Lattice



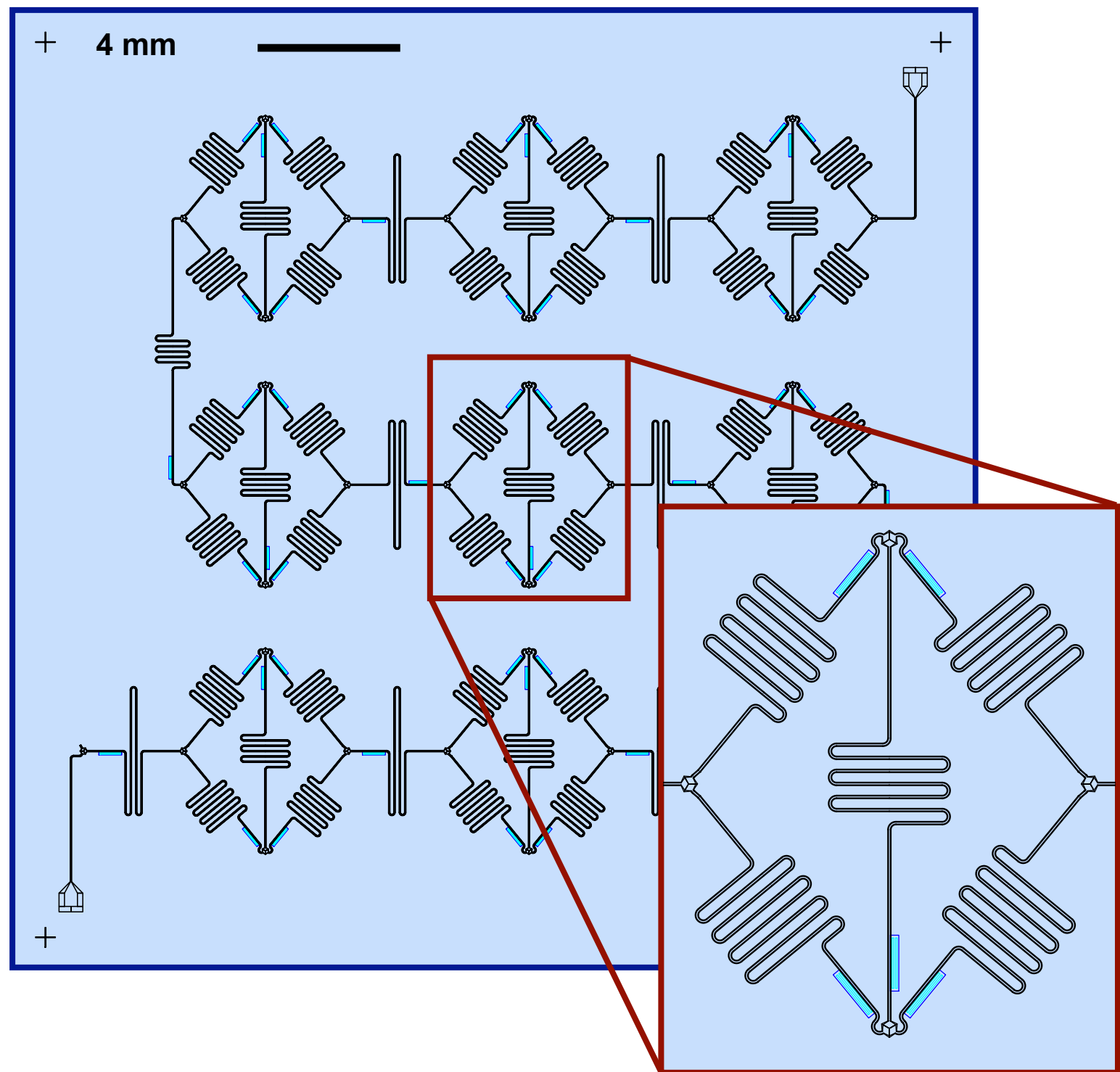
Band Structure



- Flat bands
- Gapped
- Ungapped
- Linear bands
- Quadratic bands

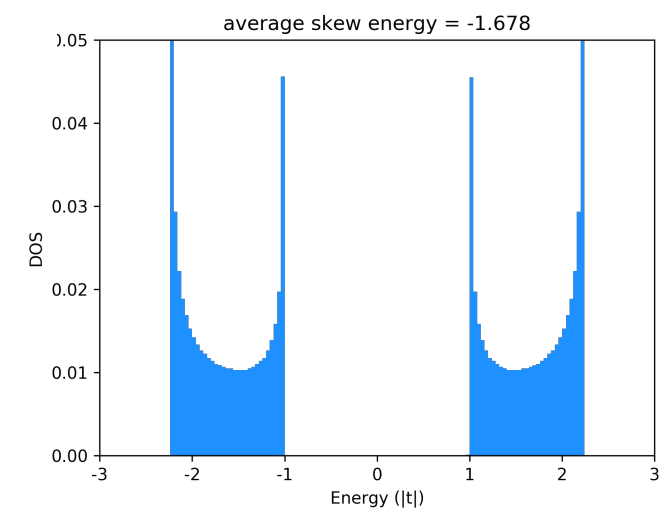
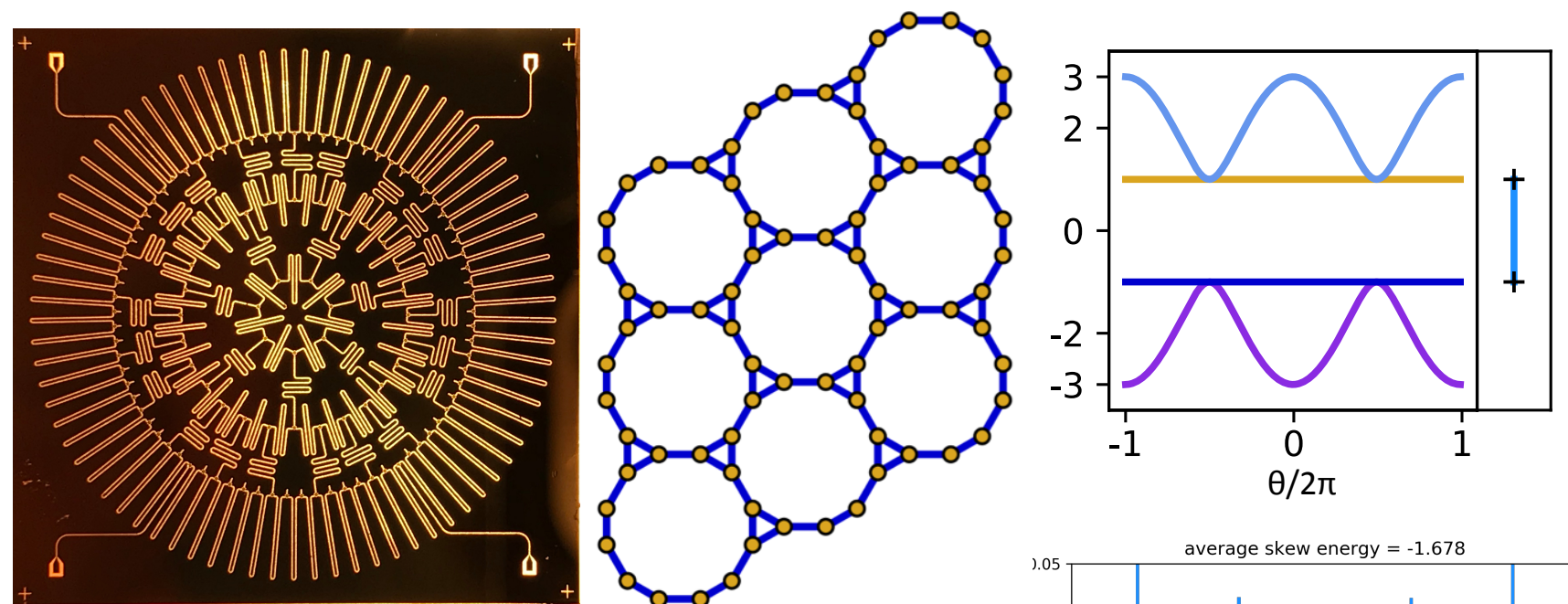
Device Design

(preliminary)



Conclusion and Outlook

- Circuit QED lattices
 - Artificial photonic materials
 - Interacting photons
- Hyperbolic lattices
 - On-chip fabrication
- Flat-band lattices
 - Optimal gaps
- Mathematics
 - Graph Spectra
 - Gap Sets
 - Abelian Covers



Kollár *et al.* Nature **571** (2019)
 Kollár *et al.* Comm. Math. Phys. **376**, 1909 (2020)
 Boettcher *et al.* Phys. Rev. A **102**, 032208 (2020)
 Kollár *et al.* Comm. AMS **1,1** (2021)
 Boettcher *et al.* arXiv:2105.0187 (2021)
 Bienias *et al.* Phys. Rev. Lett. **128**, 013601 (2022)
 Chapman, Flammia, AJK, PRX Quantum **3**, 030321 (2022)

Conclusion and Outlook

- Circuit QED lattices

- Artificial photonic materials
- Interacting photons

- Hyperbolic lattices

- On-chip fabrication

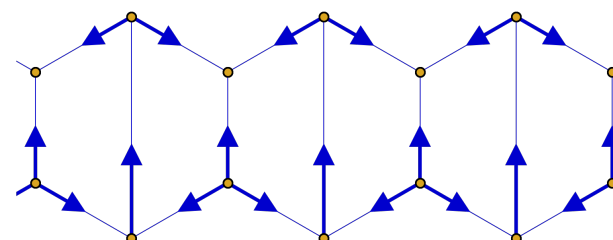
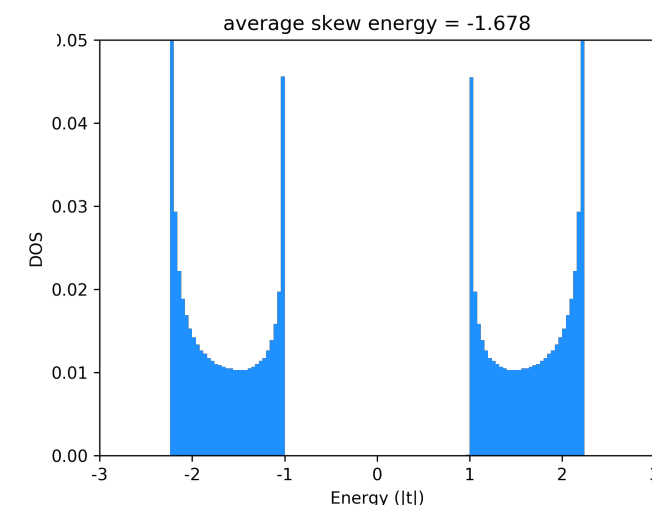
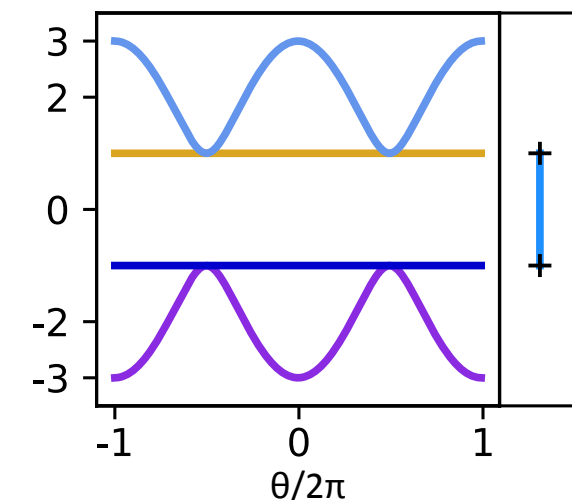
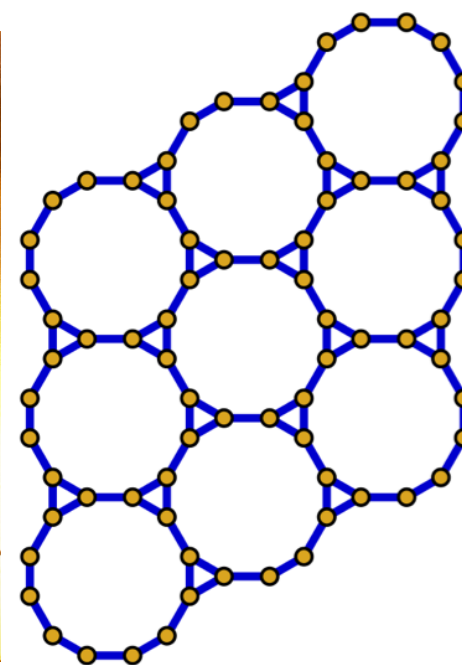
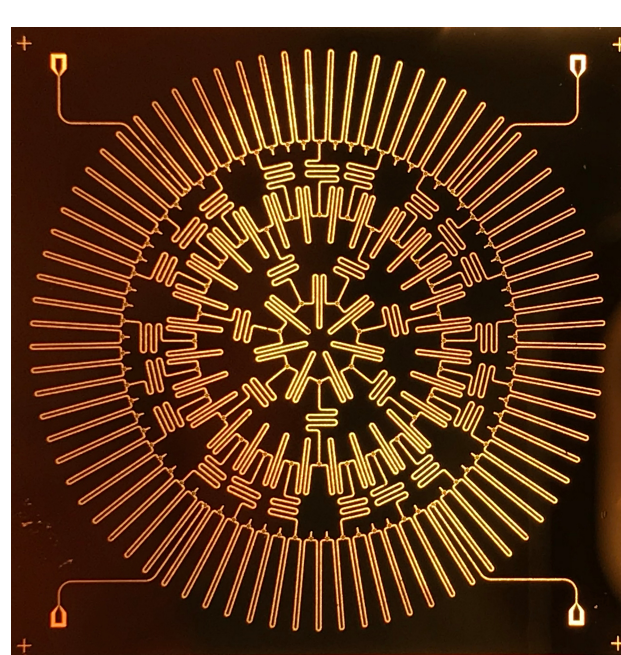
- Flat-band lattices

- Optimal gaps

- Mathematics

- Graph Spectra
- Gap Sets
- Abelian Covers

- Outlook



Kollár *et al.* Nature **571** (2019)

Kollár *et al.* Comm. Math. Phys. **376**, 1909 (2020)

Boettcher *et al.* Phys. Rev. A **102**, 032208 (2020)

Kollár *et al.* Comm. AMS **1,1** (2021)

Boettcher *et al.* arXiv:2105.0187 (2021)

Bienias *et al.* Phys. Rev. Lett. **128**, 013601 (2022)

Chapman, Flammia, AJK, PRX Quantum **3**, 030321 (2022)

Conclusion and Outlook

- Circuit QED lattices

- Artificial photonic materials
- Interacting photons

- Hyperbolic lattices

- On-chip fabrication

- Flat-band lattices

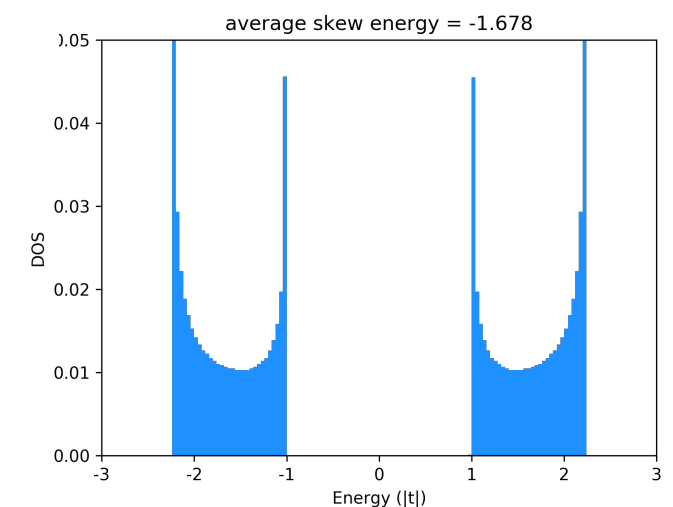
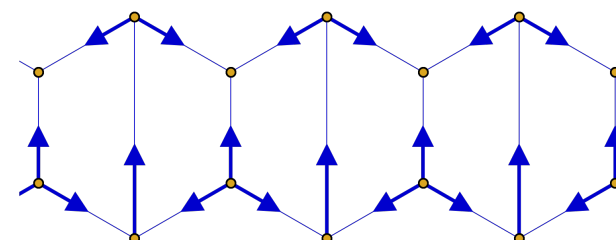
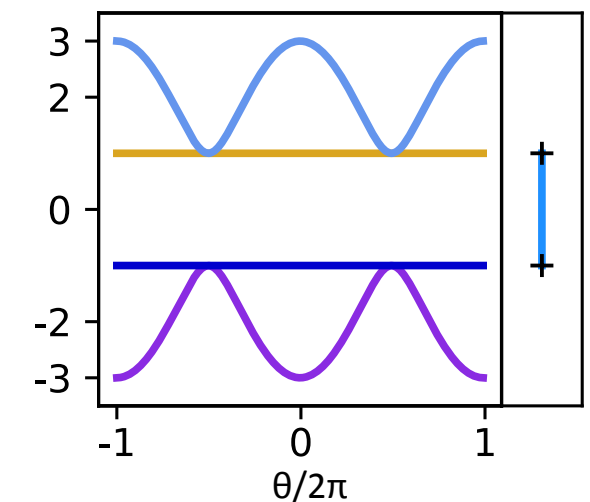
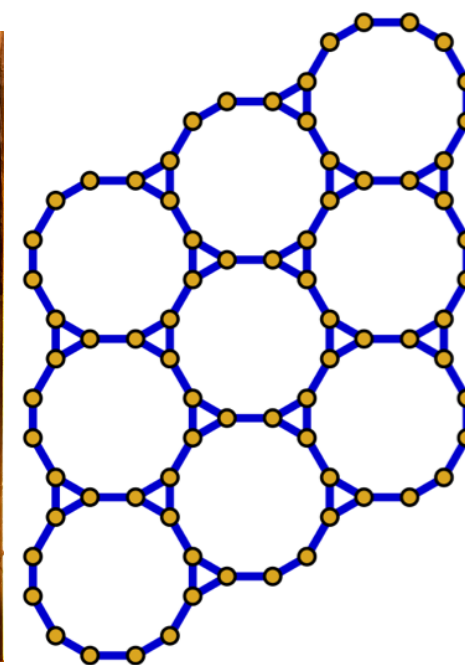
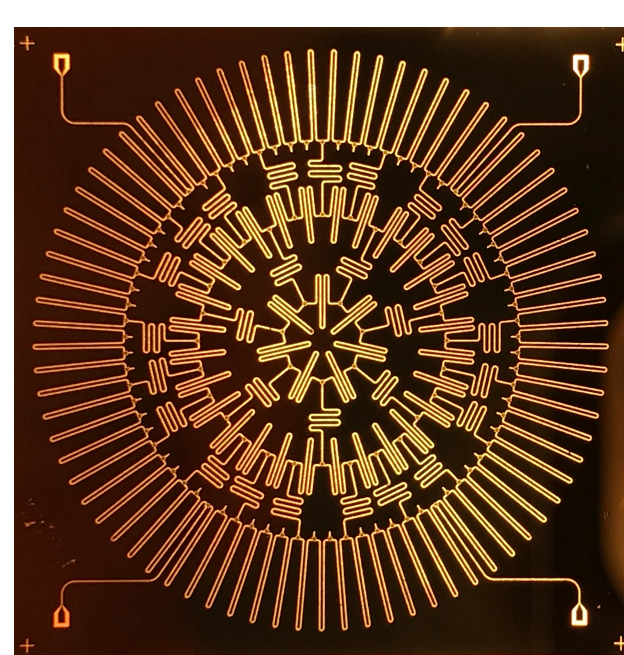
- Optimal gaps

- Mathematics

- Graph Spectra
- Gap Sets
- Abelian Covers

- Outlook

- Frustrated and hyperbolic interactions



Kollár *et al.* Nature **571** (2019)

Kollár *et al.* Comm. Math. Phys. **376**, 1909 (2020)

Boettcher *et al.* Phys. Rev. A **102**, 032208 (2020)

Kollár *et al.* Comm. AMS **1,1** (2021)

Boettcher *et al.* arXiv:2105.0187 (2021)

Bienias *et al.* Phys. Rev. Lett. **128**, 013601 (2022)

Chapman, Flammia, AJK, PRX Quantum **3**, 030321 (2022)

Conclusion and Outlook

- Circuit QED lattices

- Artificial photonic materials
- Interacting photons

- Hyperbolic lattices

- On-chip fabrication

- Flat-band lattices

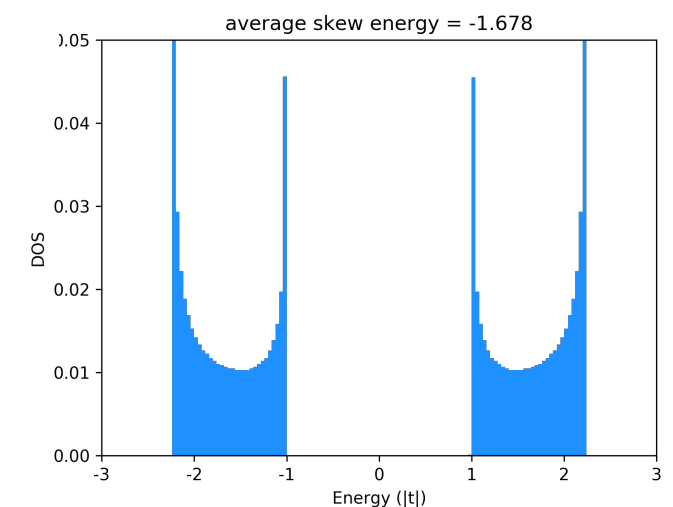
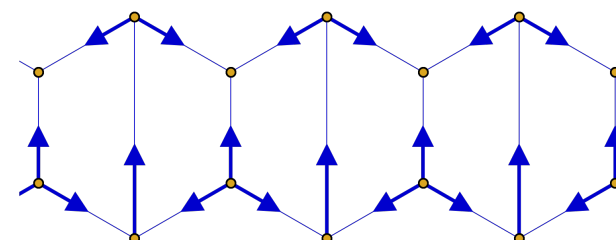
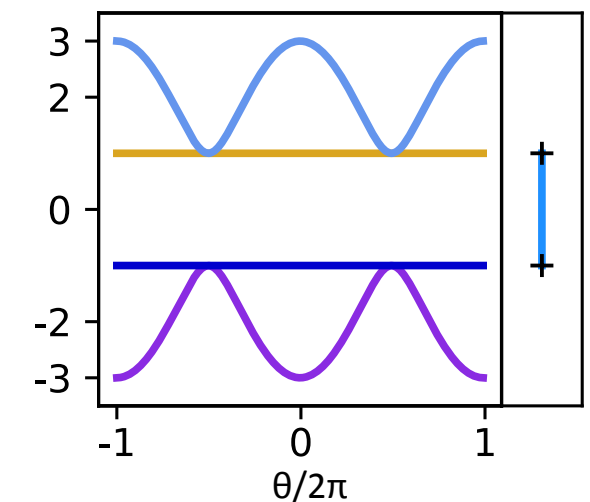
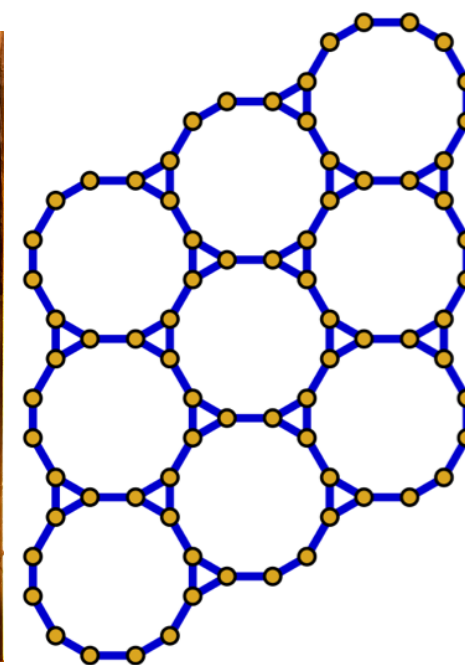
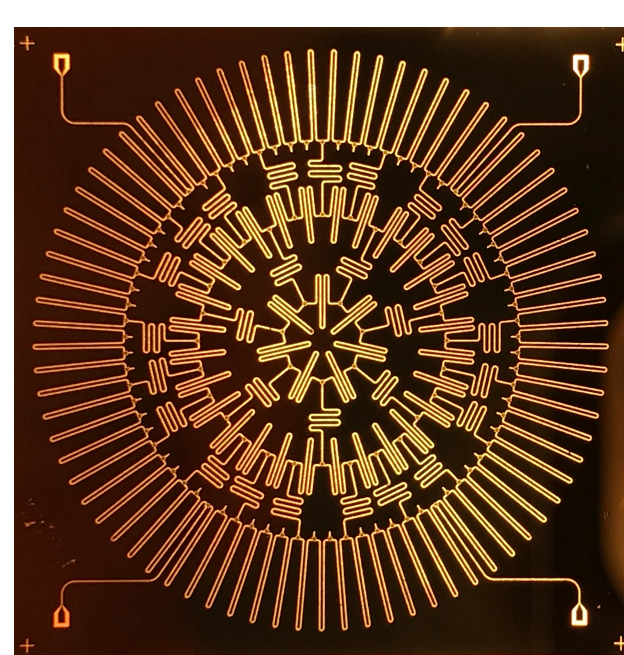
- Optimal gaps

- Mathematics

- Graph Spectra
- Gap Sets
- Abelian Covers

- Outlook

- Frustrated and hyperbolic interactions
- Many-body physics in flat bands



Kollár *et al.* Nature **571** (2019)
 Kollár *et al.* Comm. Math. Phys. **376**, 1909 (2020)
 Boettcher *et al.* Phys. Rev. A **102**, 032208 (2020)
 Kollár *et al.* Comm. AMS **1,1** (2021)
 Boettcher *et al.* arXiv:2105.0187 (2021)
 Bienias *et al.* Phys. Rev. Lett. **128**, 013601 (2022)
 Chapman, Flammia, AJK, PRX Quantum **3**, 030321 (2022)

Conclusion and Outlook

- Circuit QED lattices

- Artificial photonic materials
- Interacting photons

- Hyperbolic lattices

- On-chip fabrication

- Flat-band lattices

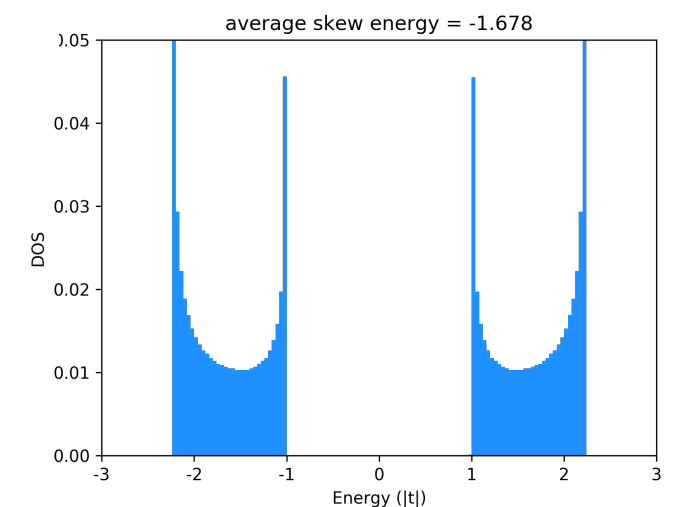
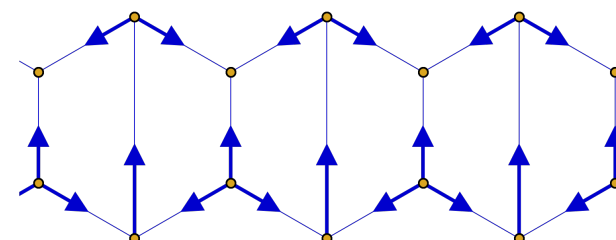
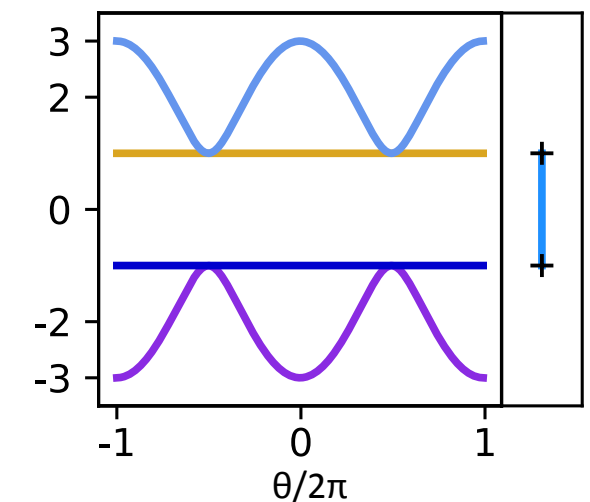
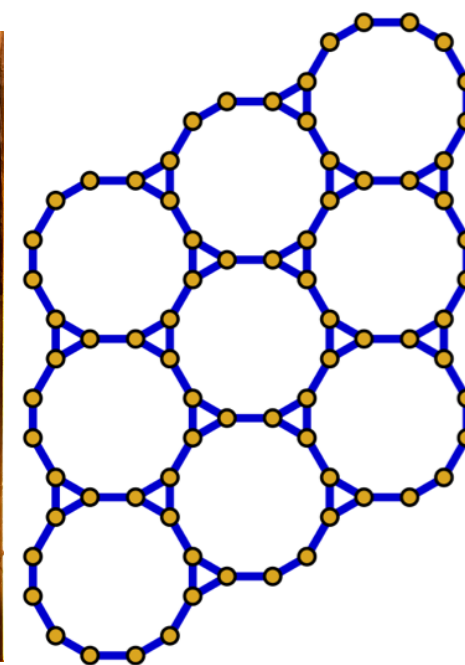
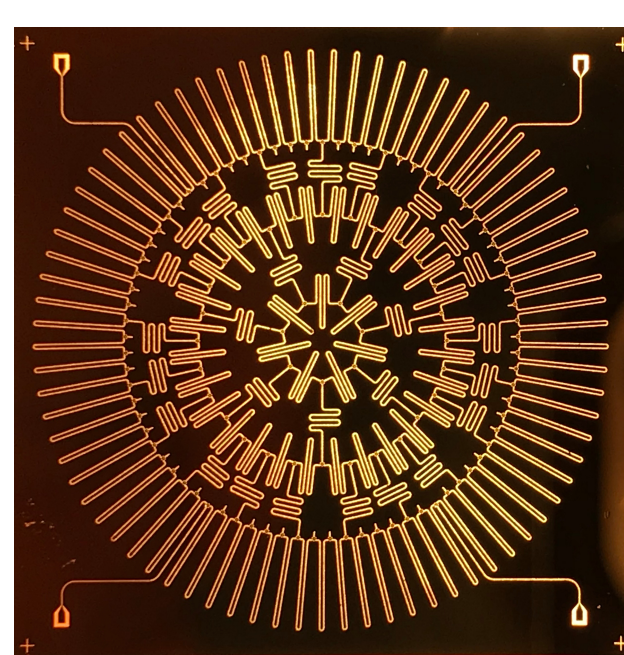
- Optimal gaps

- Mathematics

- Graph Spectra
- Gap Sets
- Abelian Covers

- Outlook

- Frustrated and hyperbolic interactions
- Many-body physics in flat bands
- Exactly-solvable 3D line-graph codes



Kollár *et al.* Nature **571** (2019)
 Kollár *et al.* Comm. Math. Phys. **376**, 1909 (2020)
 Boettcher *et al.* Phys. Rev. A **102**, 032208 (2020)
 Kollár *et al.* Comm. AMS **1,1** (2021)
 Boettcher *et al.* arXiv:2105.0187 (2021)
 Bienias *et al.* Phys. Rev. Lett. **128**, 013601 (2022)
 Chapman, Flammia, AJK, PRX Quantum **3**, 030321 (2022)

Conclusion and Outlook

- Circuit QED lattices

- Artificial photonic materials
- Interacting photons

- Hyperbolic lattices

- On-chip fabrication

- Flat-band lattices

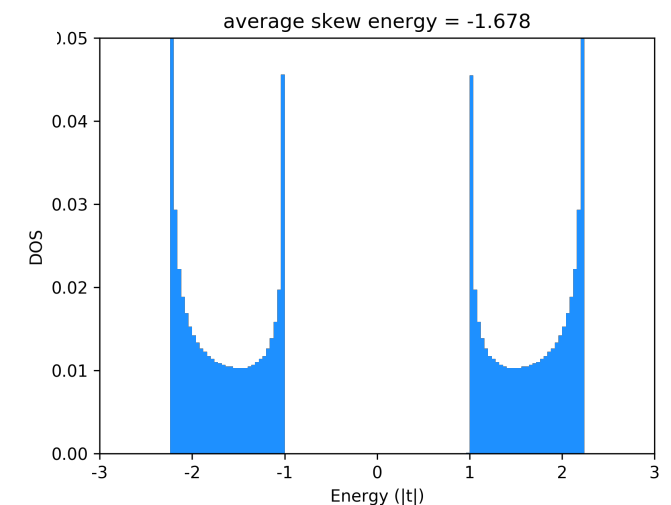
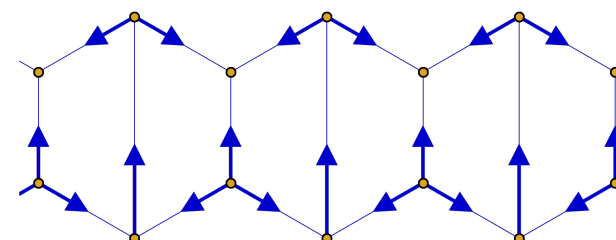
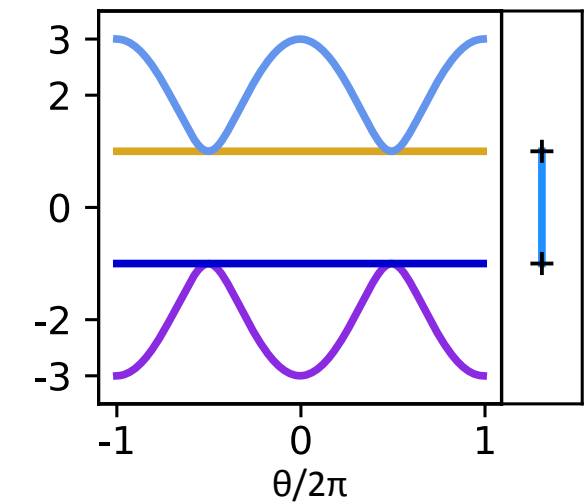
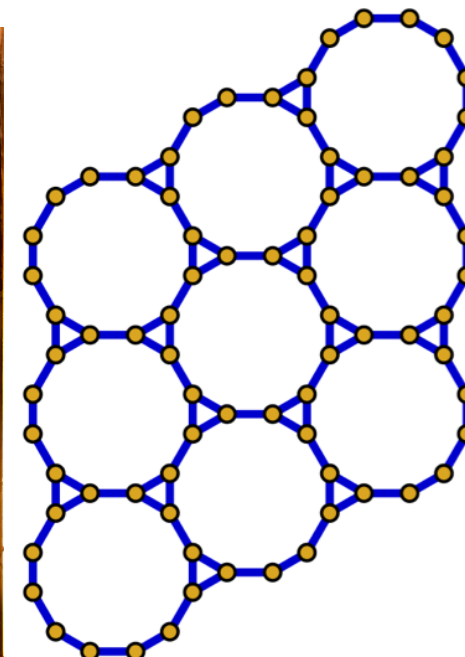
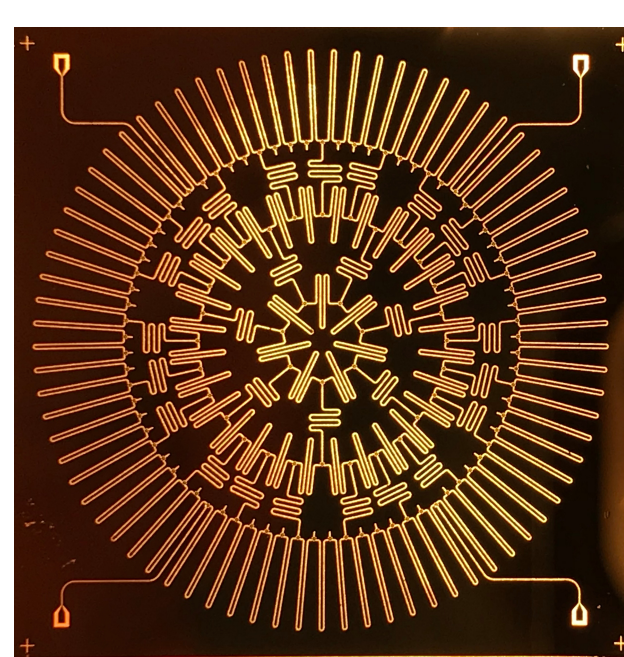
- Optimal gaps

- Mathematics

- Graph Spectra
- Gap Sets
- Abelian Covers

- Outlook

- Frustrated and hyperbolic interactions
- Many-body physics in flat bands
- Exactly-solvable 3D line-graph codes
- Leapfrog Fullerenes



Kollár *et al.* Nature **571** (2019)
 Kollár *et al.* Comm. Math. Phys. **376**, 1909 (2020)
 Boettcher *et al.* Phys. Rev. A **102**, 032208 (2020)
 Kollár *et al.* Comm. AMS **1,1** (2021)
 Boettcher *et al.* arXiv:2105.0187 (2021)
 Bienias *et al.* Phys. Rev. Lett. **128**, 013601 (2022)
 Chapman, Flammia, AJK, PRX Quantum **3**, 030321 (2022)

Circuit QED Lattices

Andrew Houck
EE, Princeton

Peter Sarnak
Math, Princeton



Alexey Gorshkov
NIST, JQI

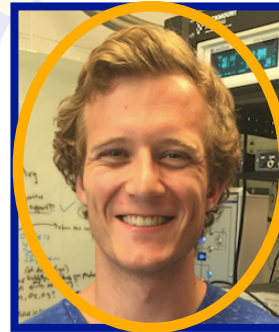
Steve Flammia
AWS



Alicia Kollár

Department of Physics and JQI, University of Maryland

Mattias
Fitzpatrick



Maya
Amouzegar



Martin
Ritter



Jeffrey
Wack



Adrian
Chapman



Przemislav
Bienias



Igor
Boettcher

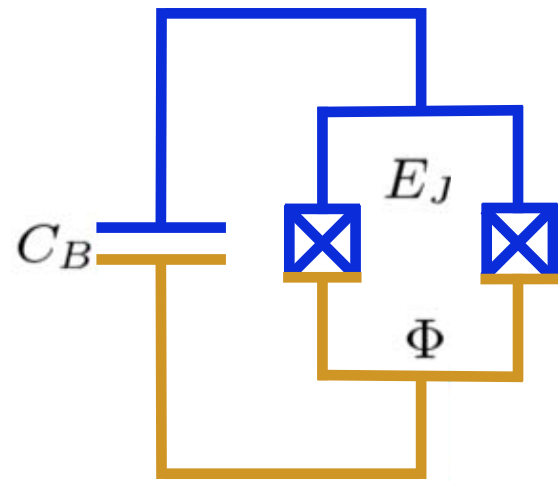


David
Long

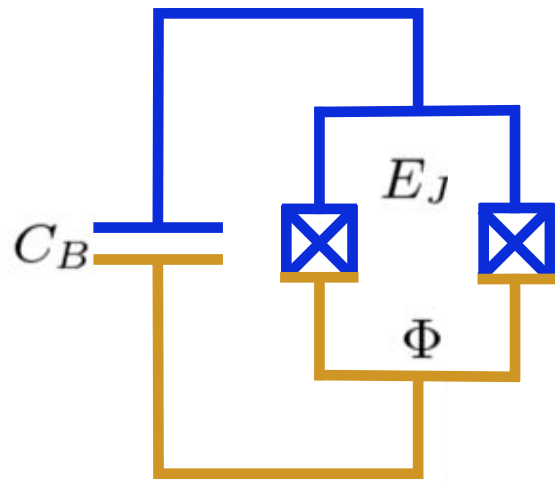




Transmon Qubit



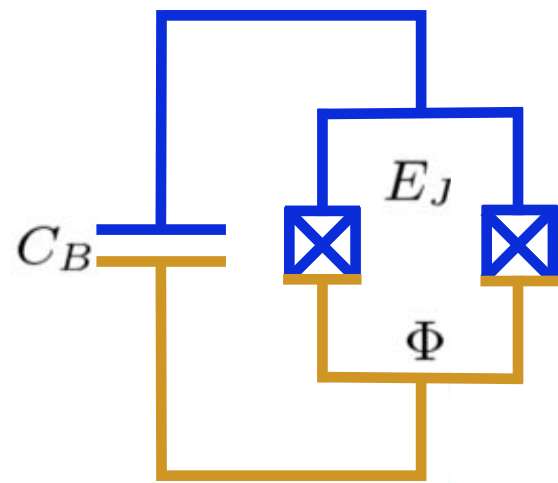
Transmon Qubit



Anharmonic oscillator

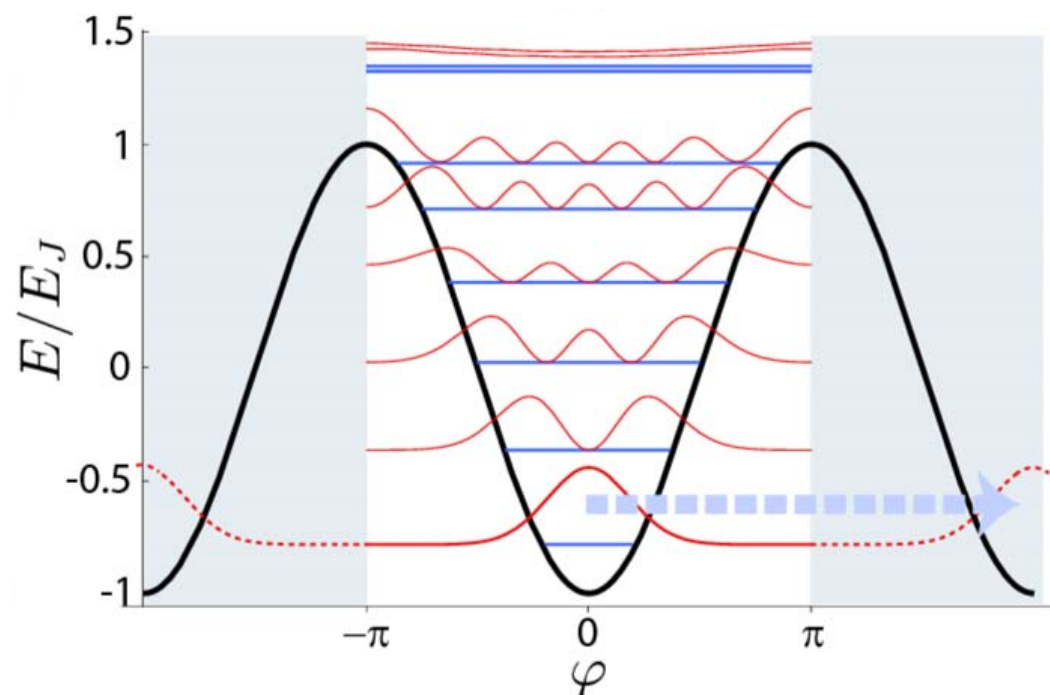
$$\hat{H} = 4E_C \hat{n}^2 - E_J \cos \hat{\varphi}$$

Transmon Qubit

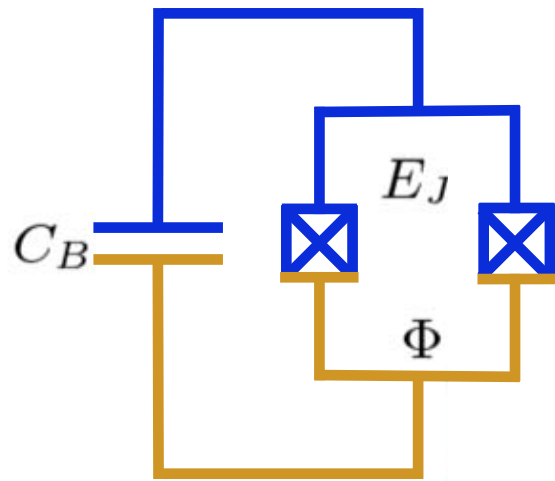


Anharmonic oscillator

$$\hat{H} = 4E_C \hat{n}^2 - E_J \cos \hat{\varphi}$$



Transmon Qubit



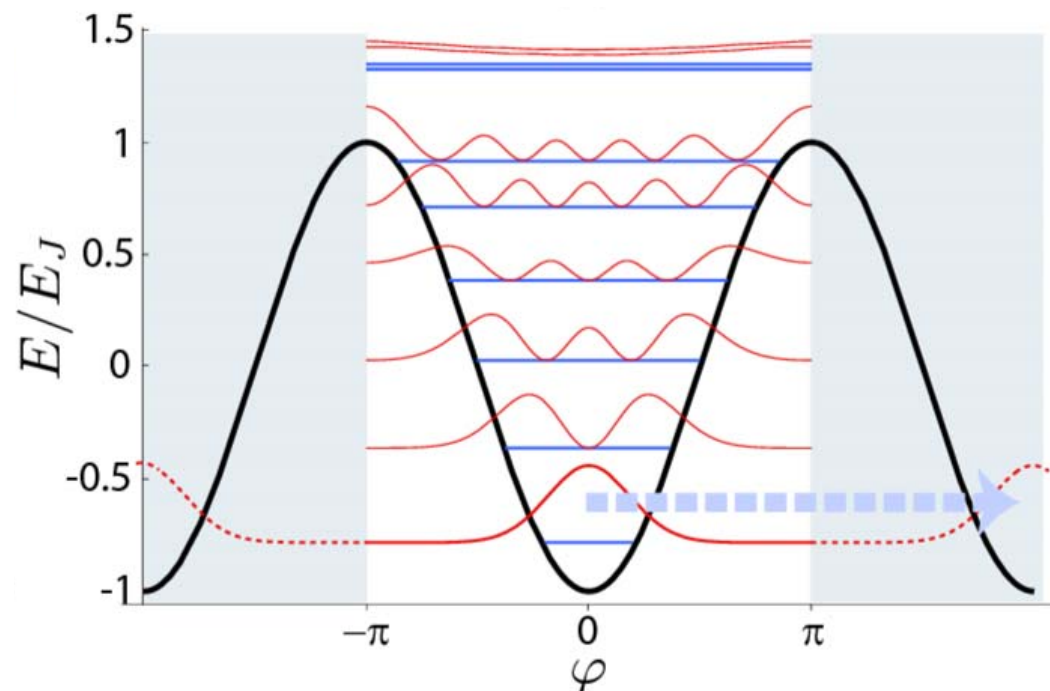
Qubit-Cavity

(Jaynes-Cummings Model)

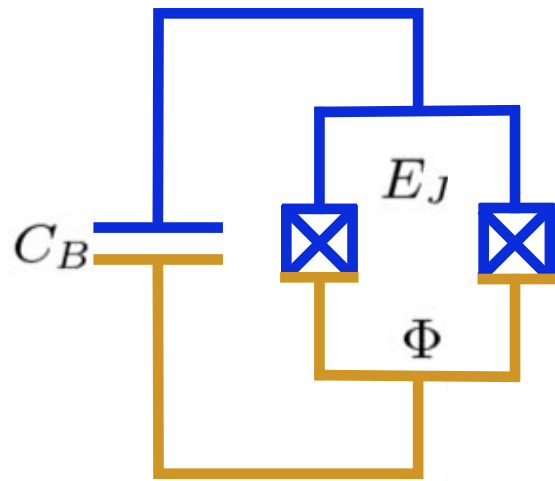
$$H_{JC} = \omega_c a^\dagger a + \frac{1}{2} \omega_q \sigma_z + g_0 (a^\dagger \sigma^- + a \sigma^+)$$

Anharmonic oscillator

$$\hat{H} = 4E_C \hat{n}^2 - E_J \cos \hat{\varphi}$$

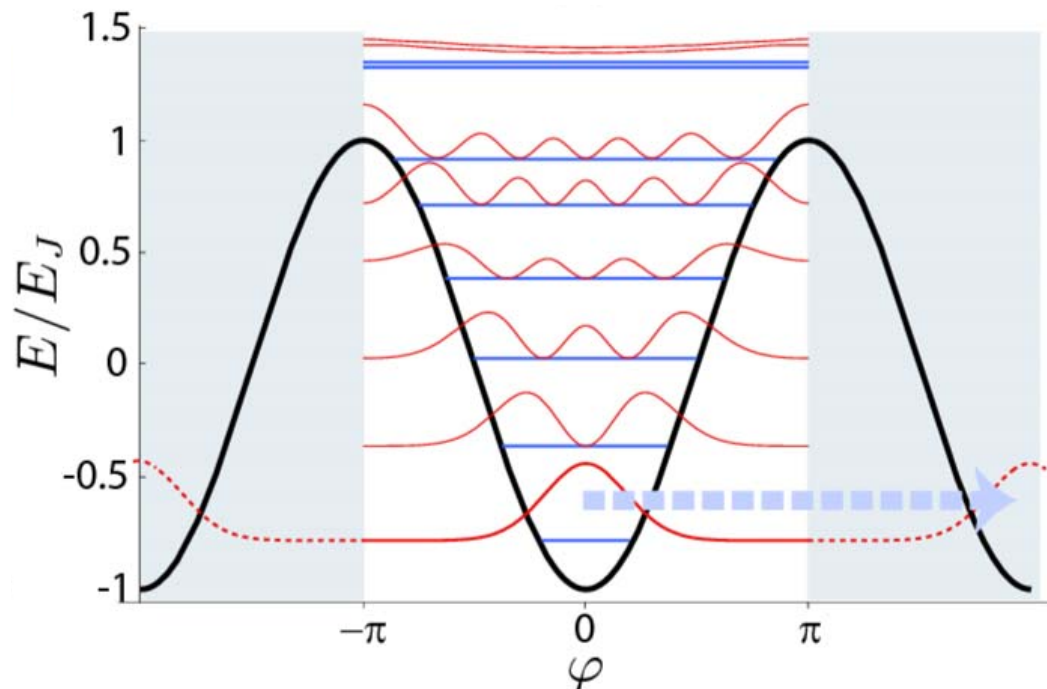


Transmon Qubit



Anharmonic oscillator

$$\hat{H} = 4E_C \hat{n}^2 - E_J \cos \hat{\varphi}$$

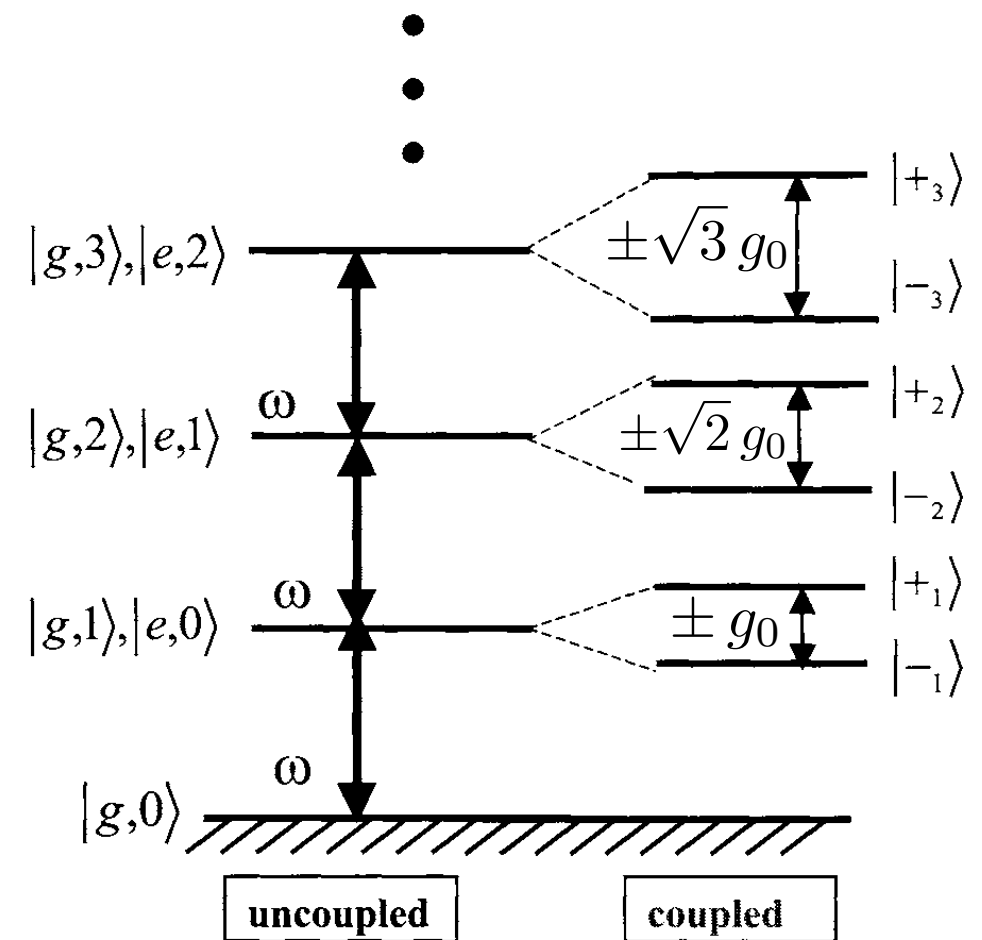


Qubit-Cavity

(Jaynes-Cummings Model)

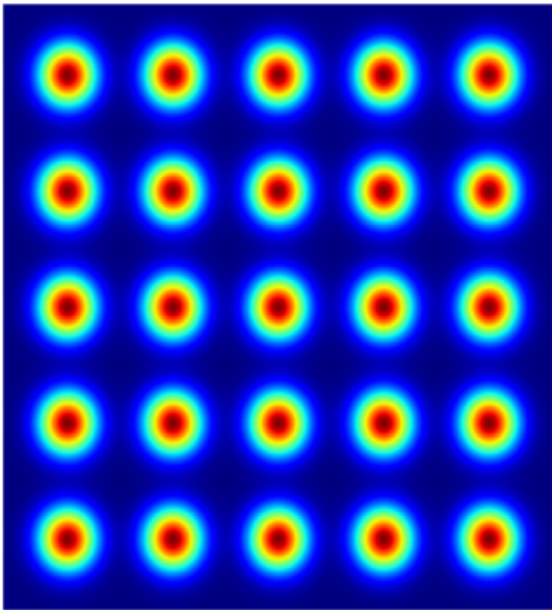
$$H_{JC} = \omega_c a^\dagger a + \frac{1}{2} \omega_q \sigma_z + g_0 (a^\dagger \sigma^- + a \sigma^+)$$

$$|\pm_n\rangle = \frac{1}{\sqrt{2}} (|g, n\rangle \pm |e, n-1\rangle),$$



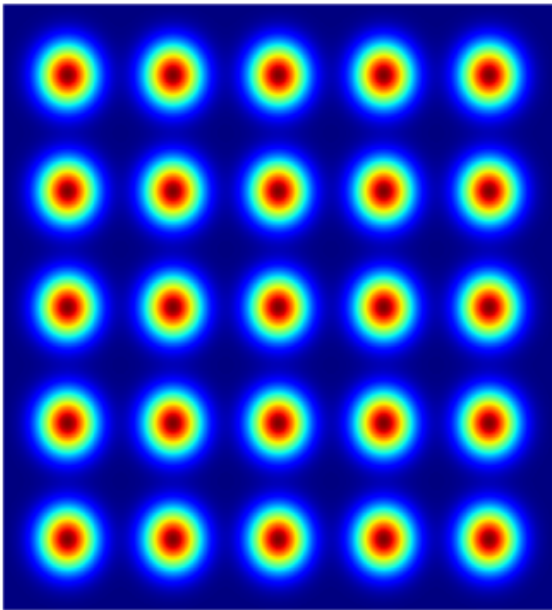
The Graph is Everything

Regular Lattice

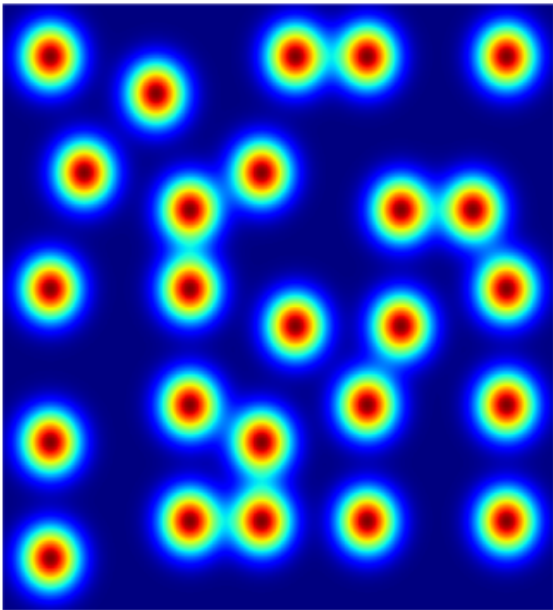


The Graph is Everything

Regular Lattice

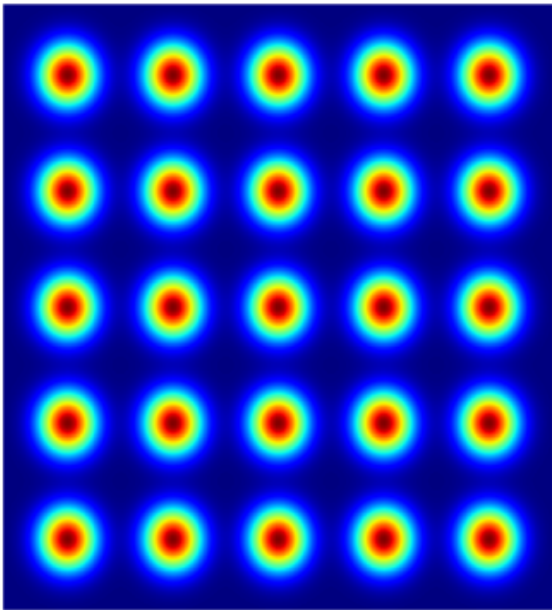


Disordered Lattice

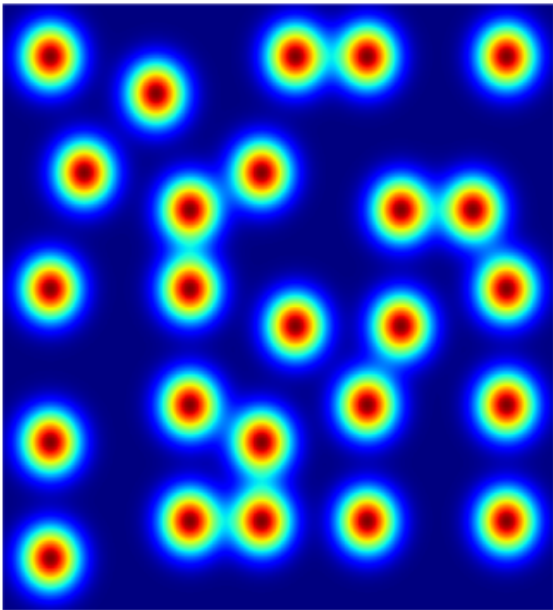


The Graph is Everything

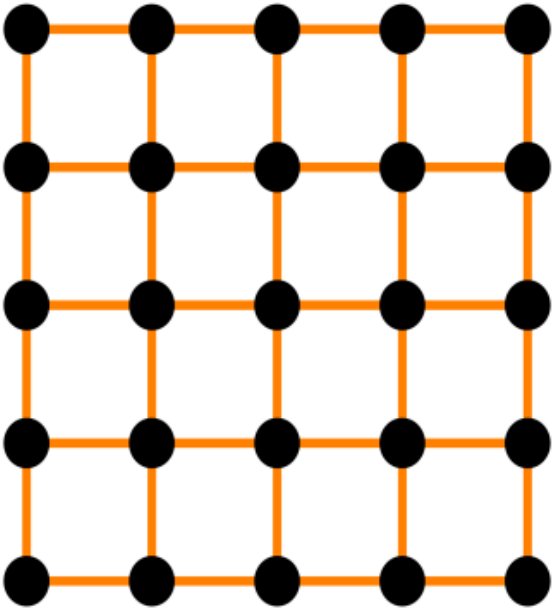
Regular Lattice



Disordered Lattice

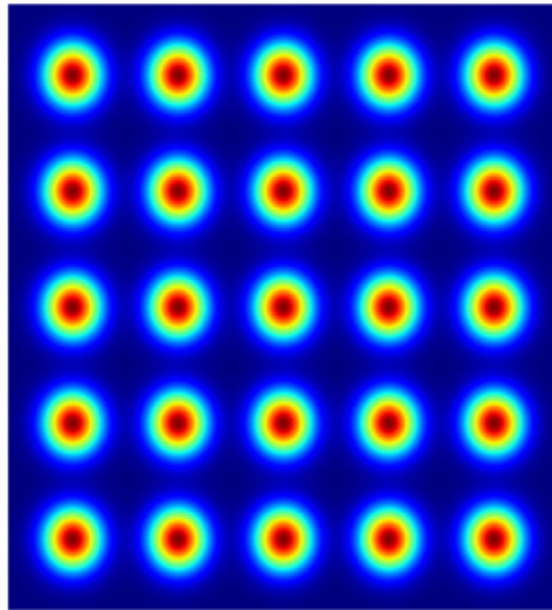


Regular Tight-Binding Graph

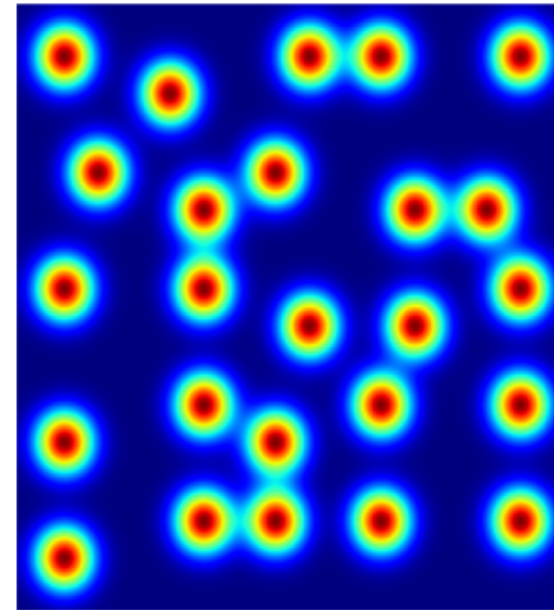


The Graph is Everything

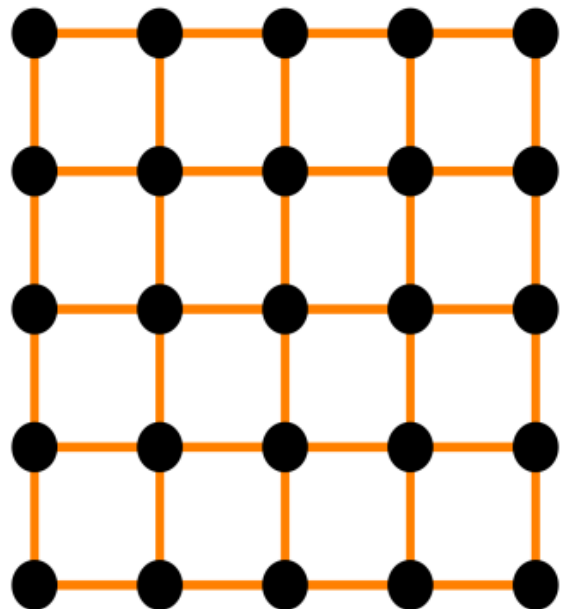
Regular Lattice



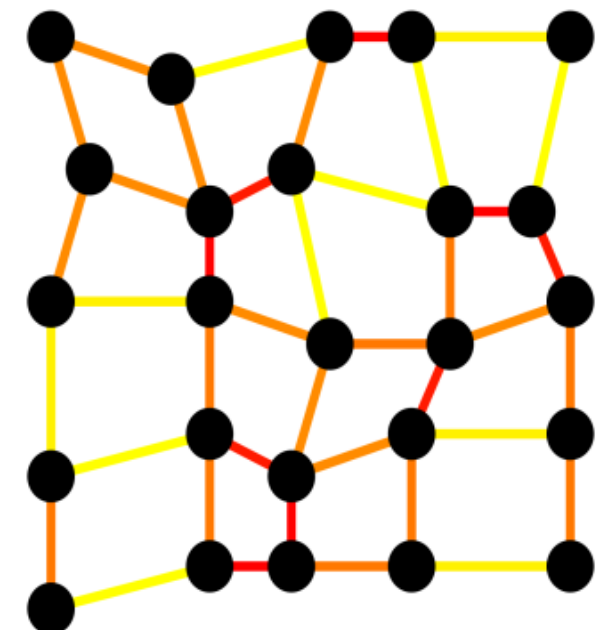
Disordered Lattice



Regular Tight-Binding Graph

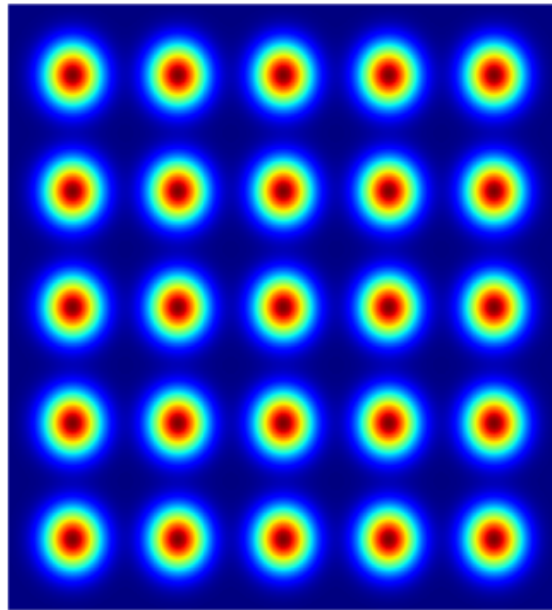


Disordered TB Graph

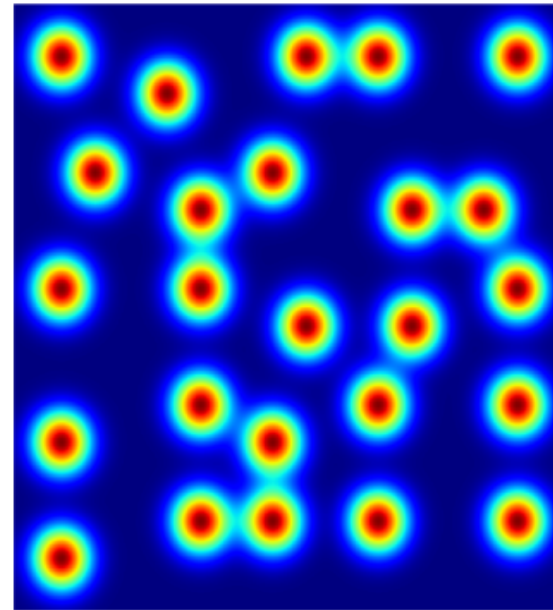


The Graph is Everything

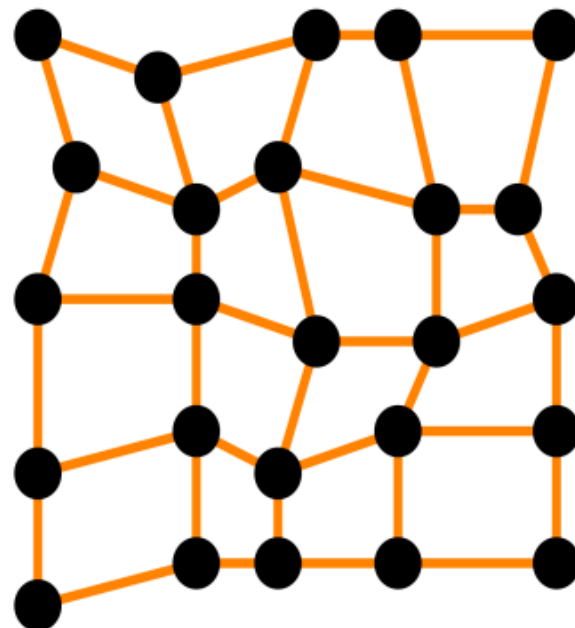
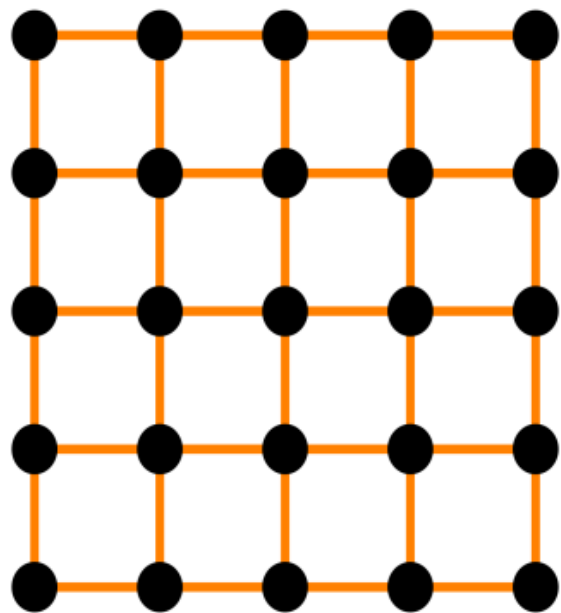
Regular Lattice



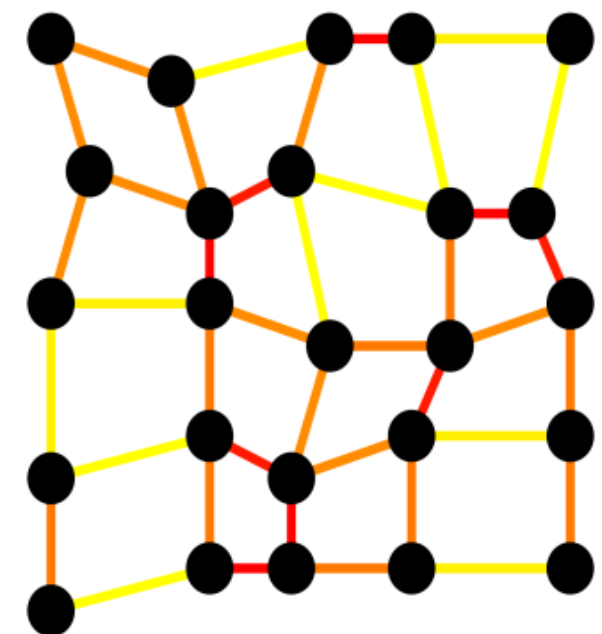
Disordered Lattice



Regular Tight-Binding Graph

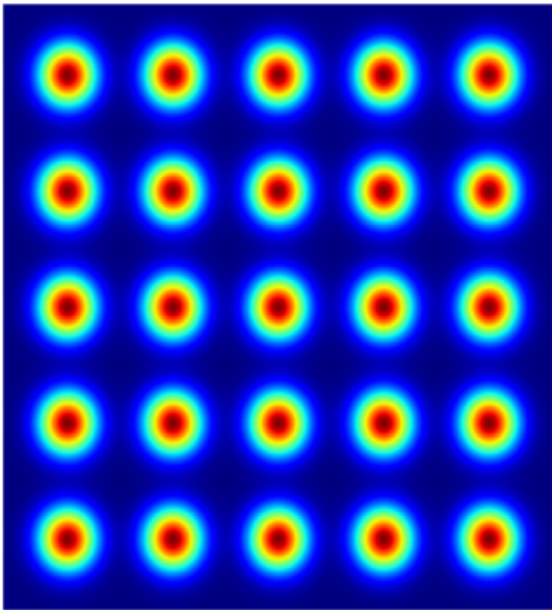


Disordered TB Graph

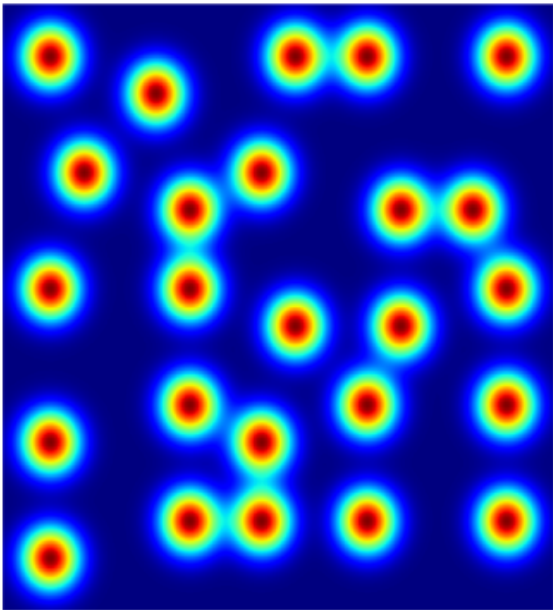


The Graph is Everything

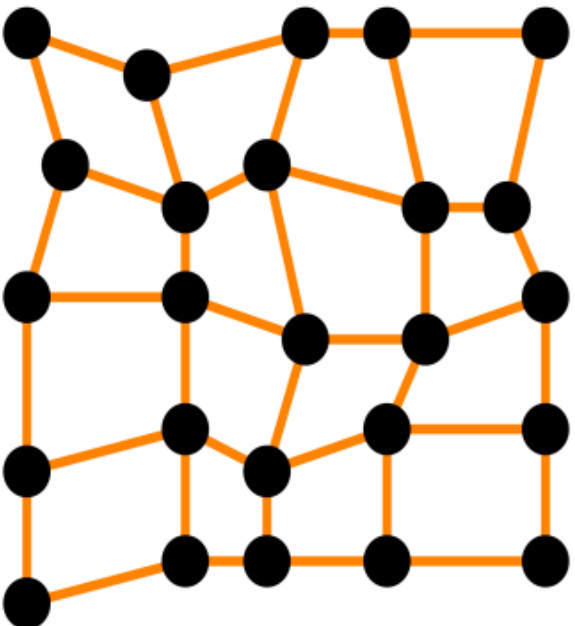
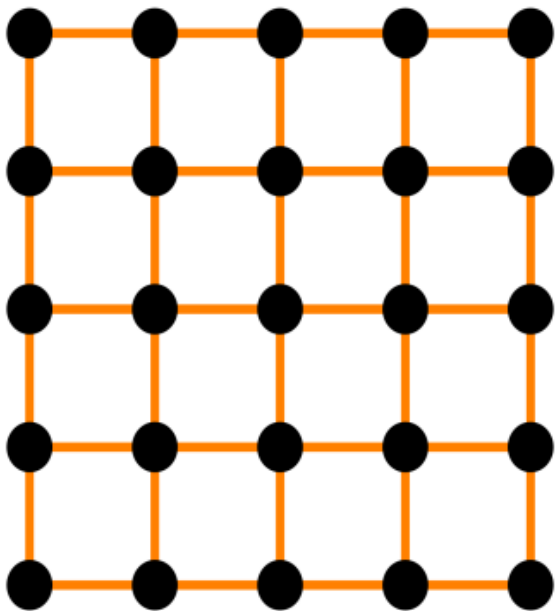
Regular Lattice



Disordered Lattice

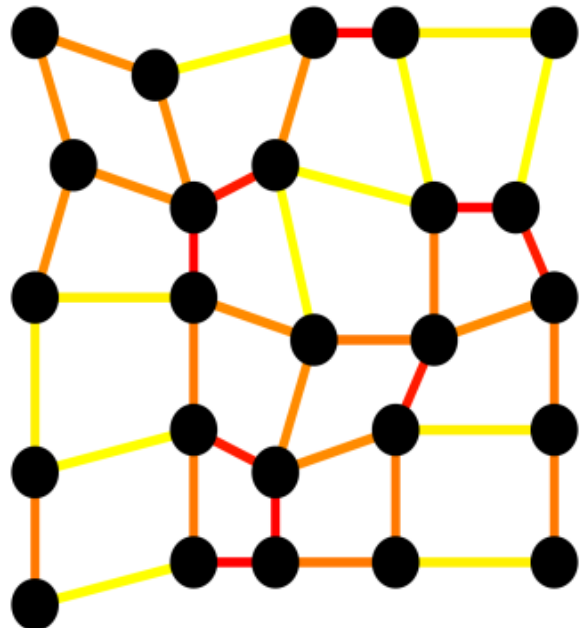


Regular Tight-Binding Graph



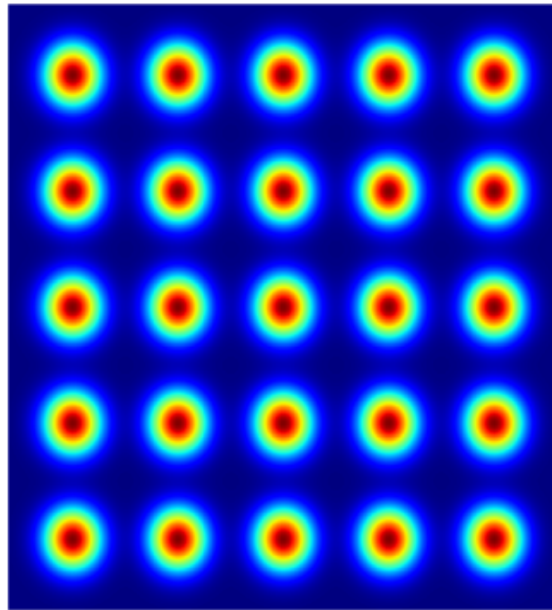
\neq

Disordered TB Graph

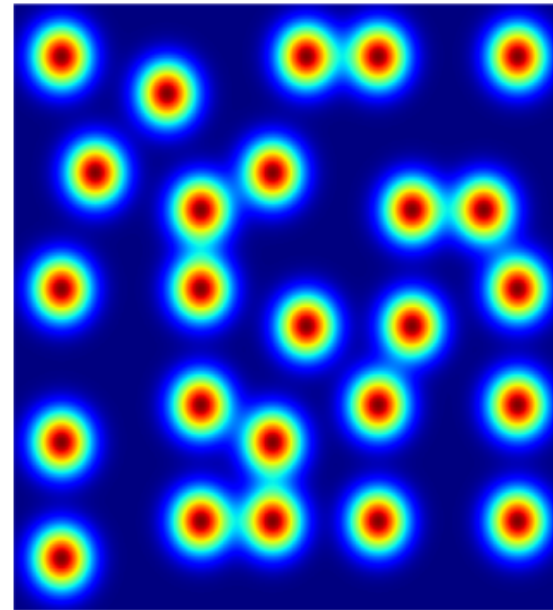


The Graph is Everything

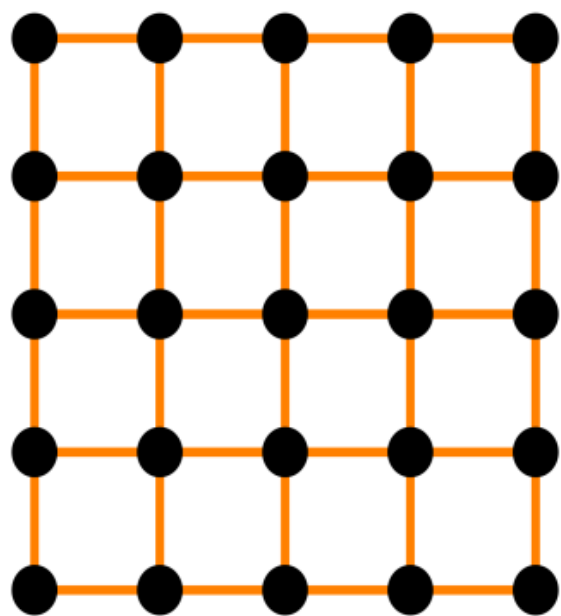
Regular Lattice



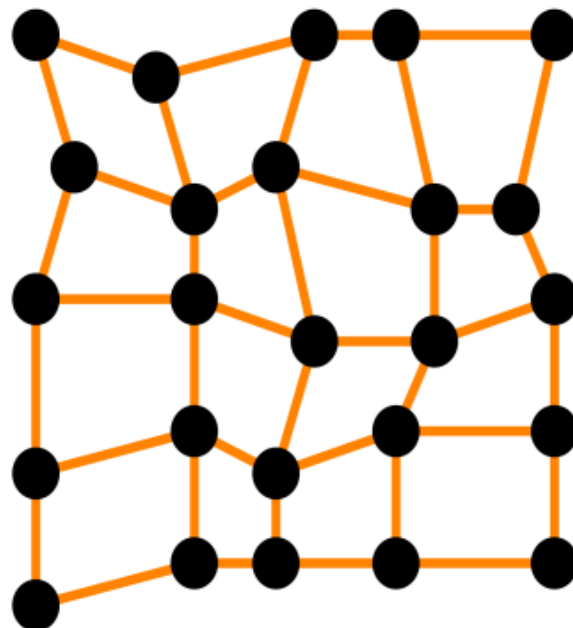
Disordered Lattice



Regular Tight-Binding Graph

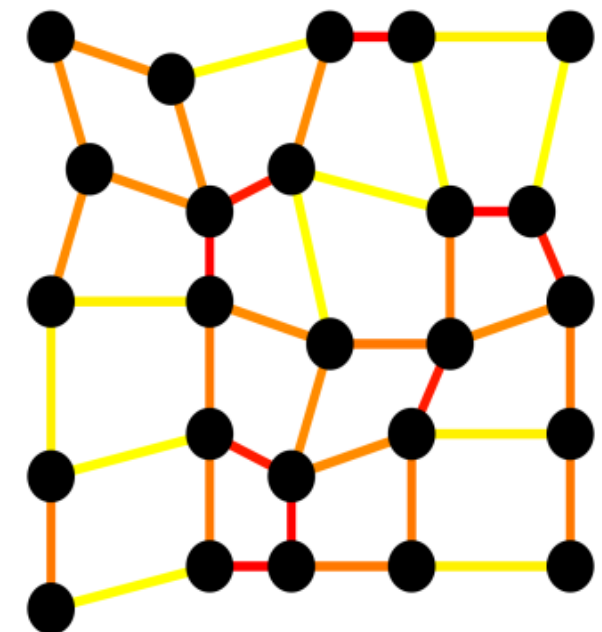


=



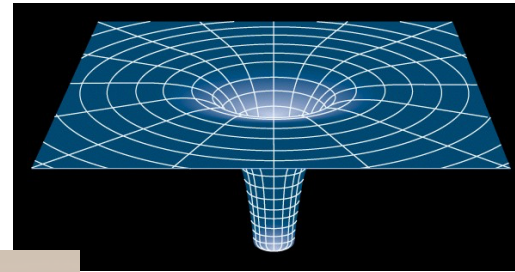
≠

Disordered TB Graph



Applications of Hyperbolic Systems

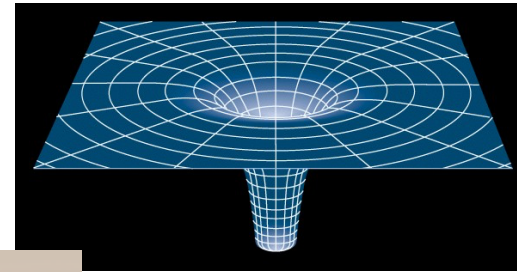
- General relativity
 - Curved space-time
- 2D materials
 - graphene, fullerenes



Applications of Hyperbolic Systems

- General relativity

- Curved space-time



- 2D materials

- graphene, fullerenes



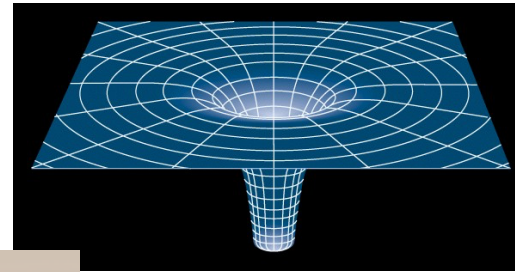
- Mathematics

- Trees 
- Cayley graphs of non-commutative groups
- Automorphic forms

Applications of Hyperbolic Systems

- General relativity

- Curved space-time



- 2D materials

- graphene, fullerenes



- Mathematics

- Trees 
- Cayley graphs of non-commutative groups
- Automorphic forms

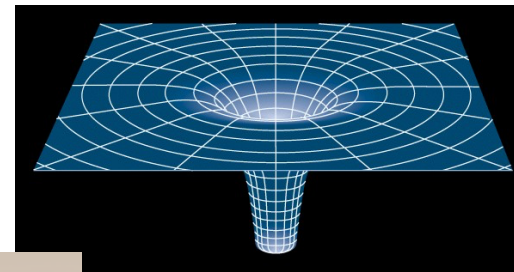
- Computer Science

- Trees 
- Efficient communication networks
- Tamper-resistant networks

Applications of Hyperbolic Systems

- General relativity

- Curved space-time



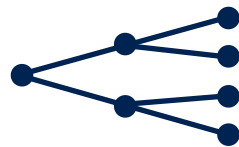
- 2D materials

- graphene, fullerenes



- Mathematics

- Trees

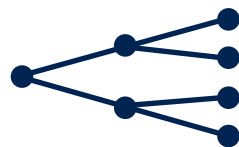


- Cayley graphs of non-commutative groups

- Automorphic forms

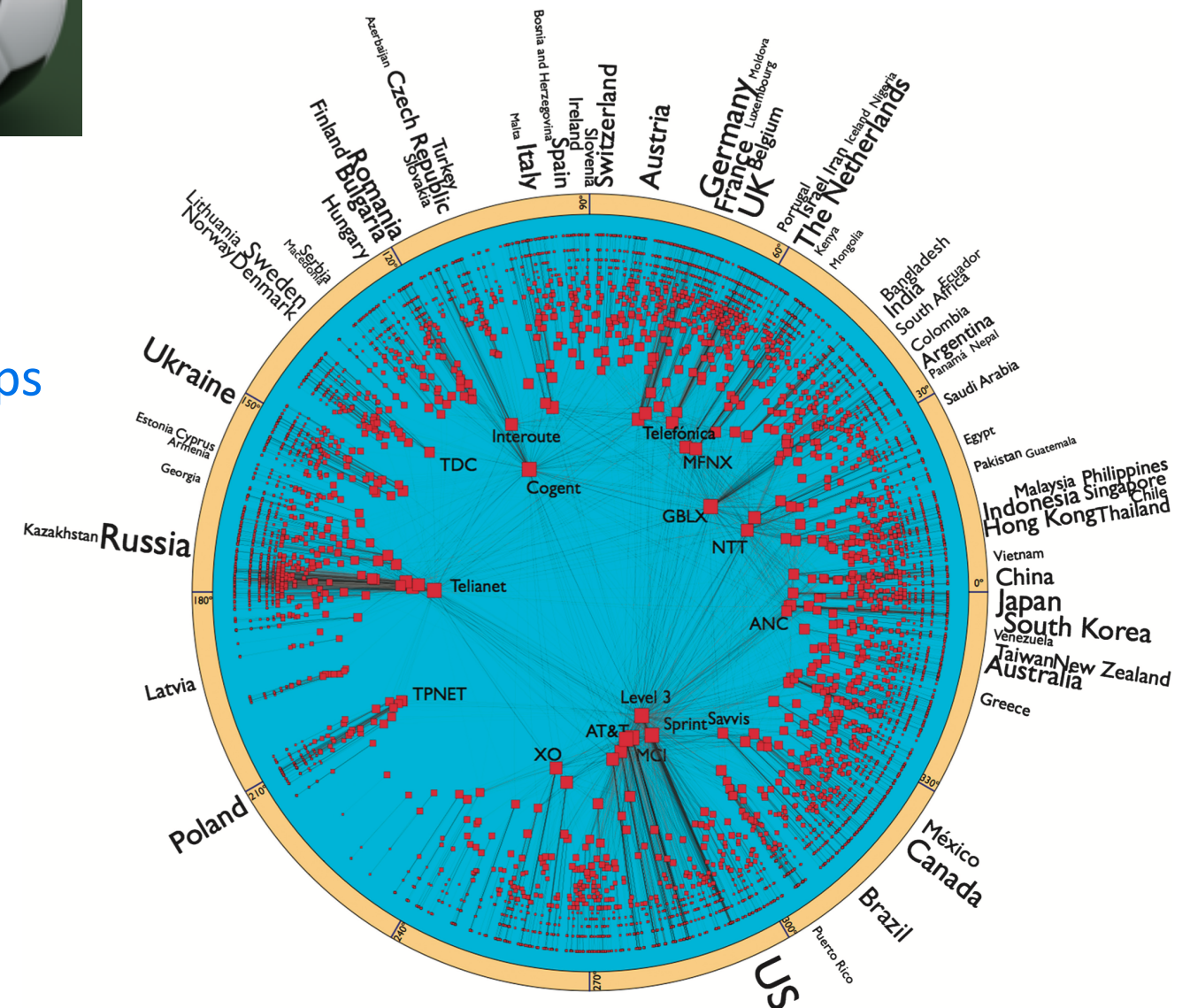
- Computer Science

- Trees



- Efficient communication networks

- Tamper-resistant networks



Continuum Limit and Green's Function

High Energy Limit of The Spectrum

Continuum Limit and Green's Function

High Energy Limit of The Spectrum

- Long-wavelength modes

Continuum Limit and Green's Function

High Energy Limit of The Spectrum

- Long-wavelength modes
- Lattice should course-grain out

Continuum Limit and Green's Function

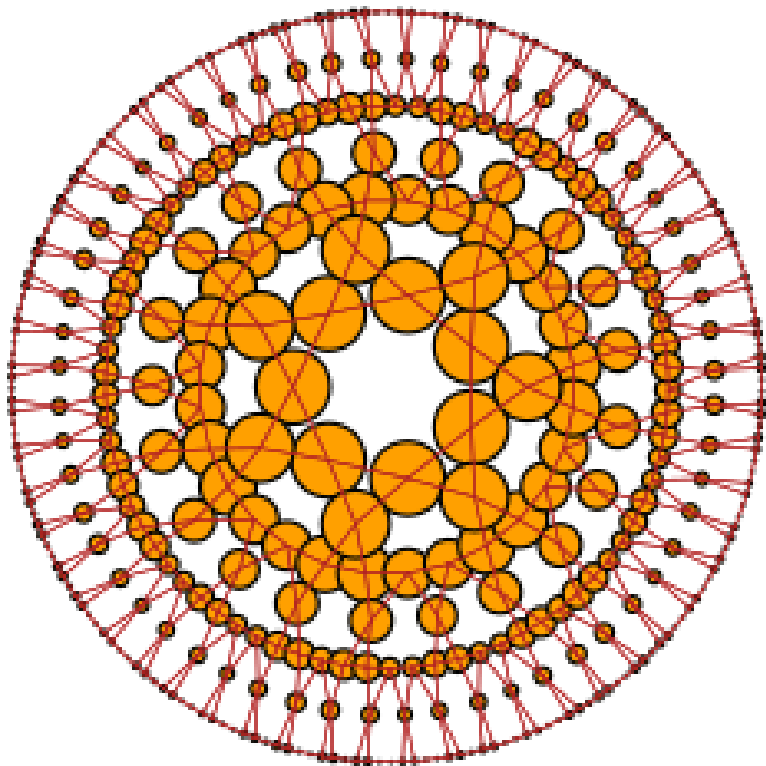
High Energy Limit of The Spectrum

- Long-wavelength modes
- Lattice should coarse-grain out
- Hyperbolic particle in a box

Continuum Limit and Green's Function

High Energy Limit of The Spectrum

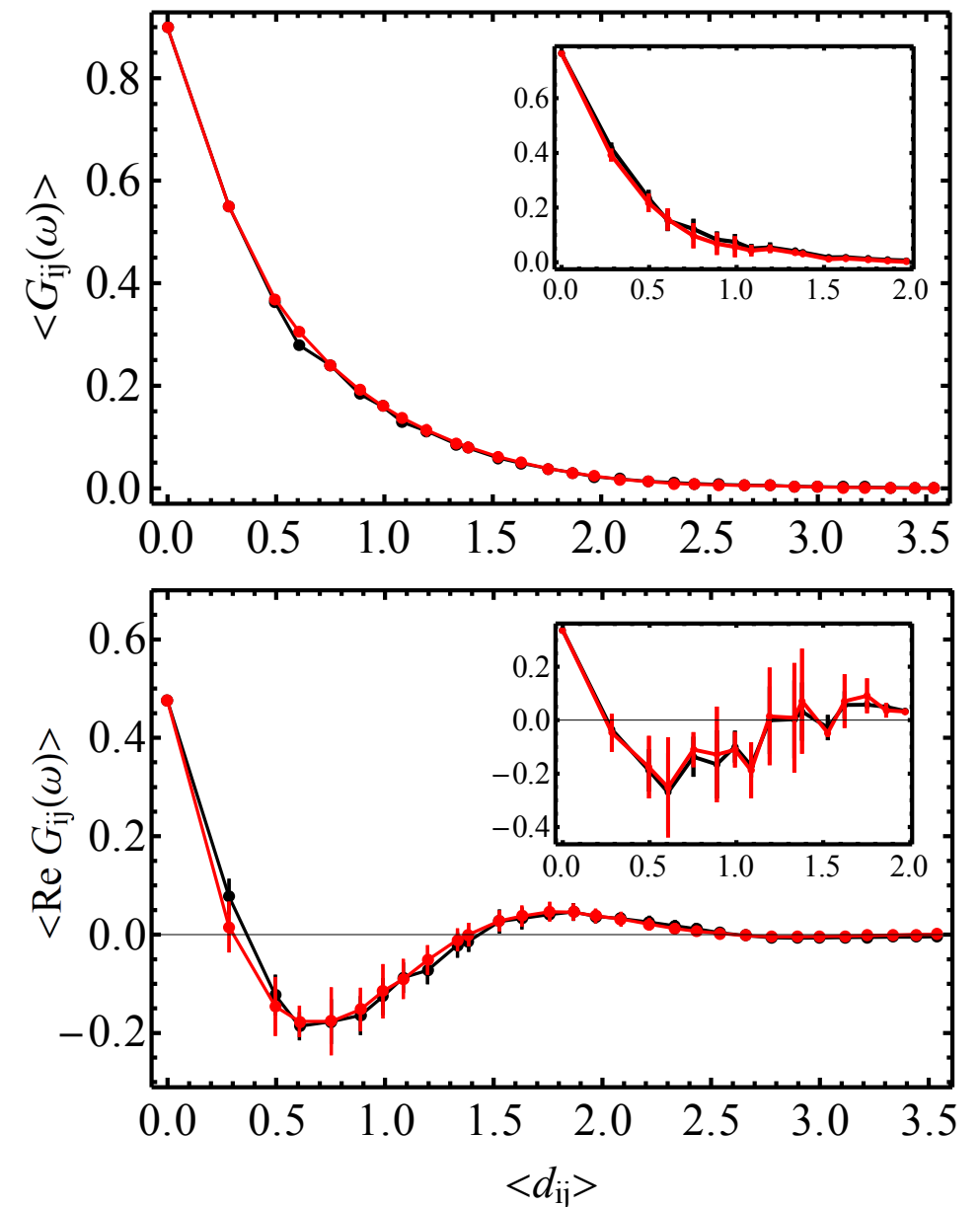
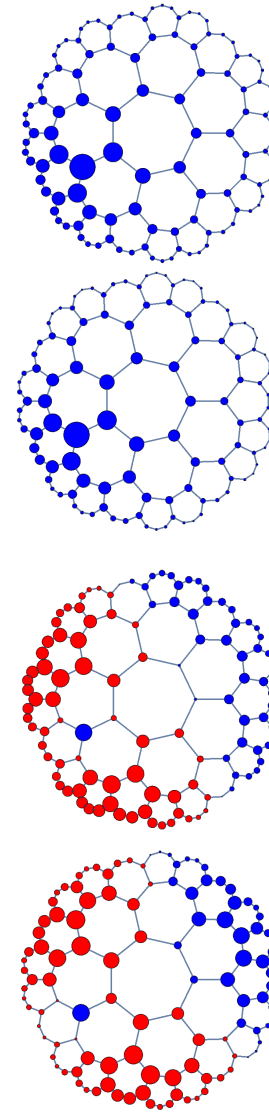
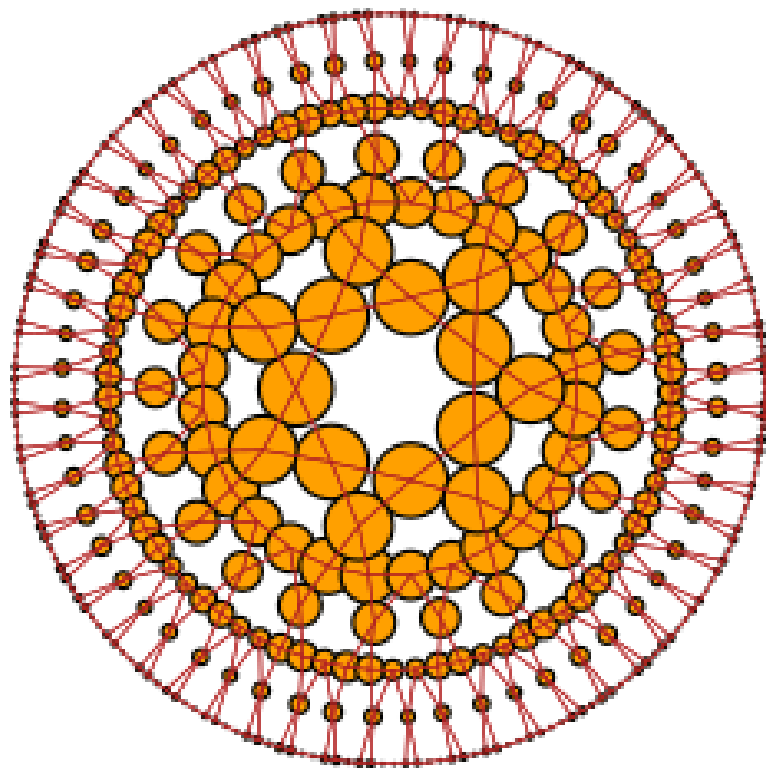
- Long-wavelength modes
- Lattice should course-grain out
- Hyperbolic particle in a box



Continuum Limit and Green's Function

High Energy Limit of The Spectrum

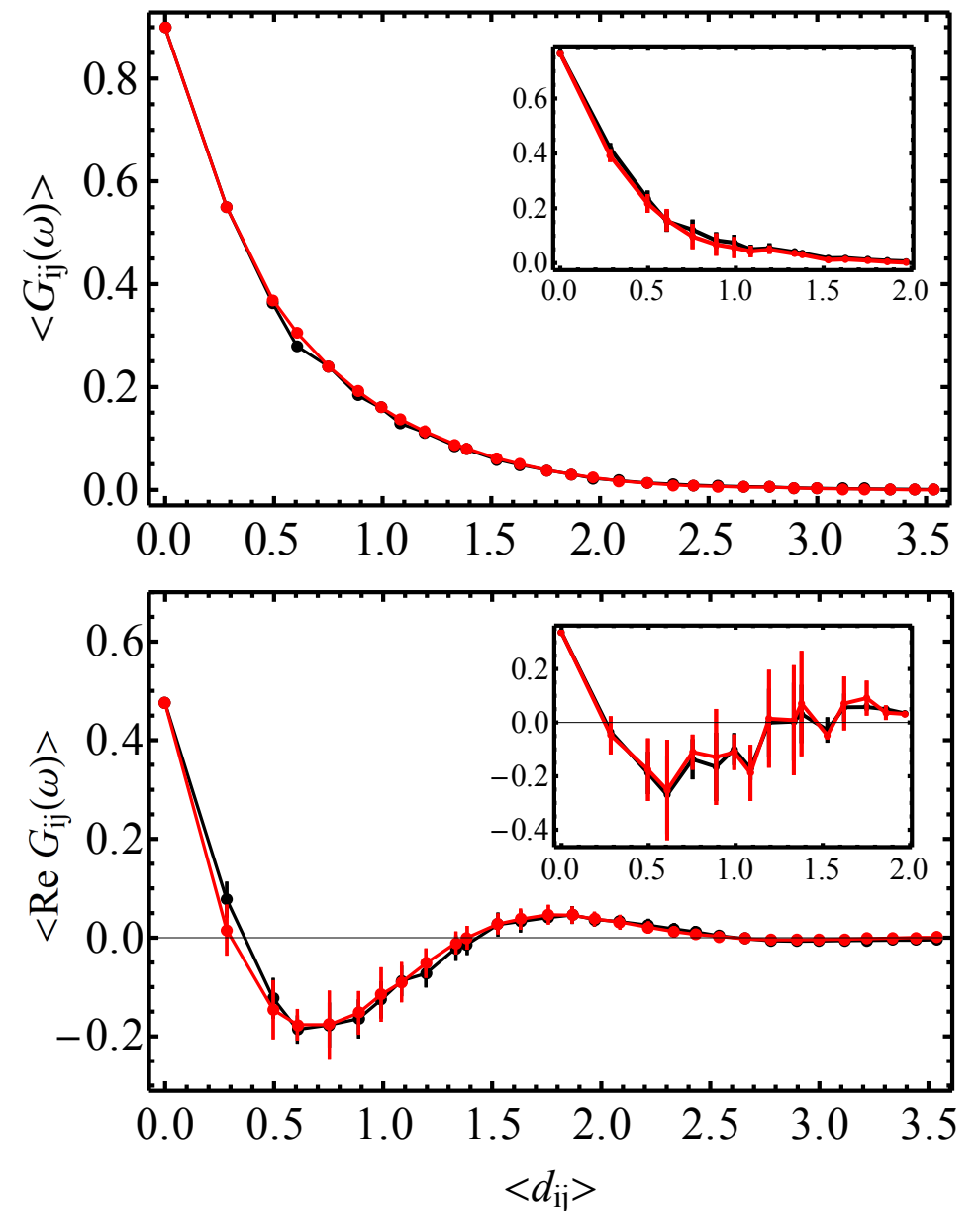
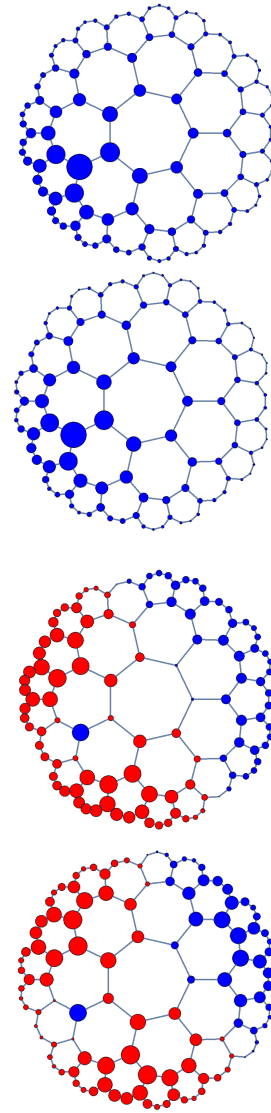
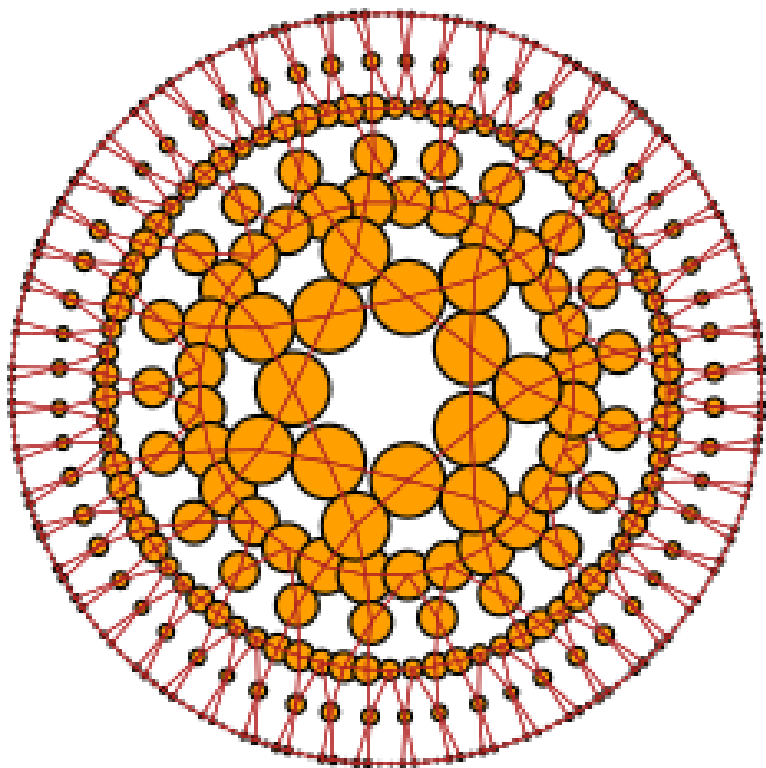
- Long-wavelength modes
- Lattice should course-grain out
- Hyperbolic particle in a box



Continuum Limit and Green's Function

High Energy Limit of The Spectrum

- Long-wavelength modes
- Lattice should course-grain out
- Hyperbolic particle in a box



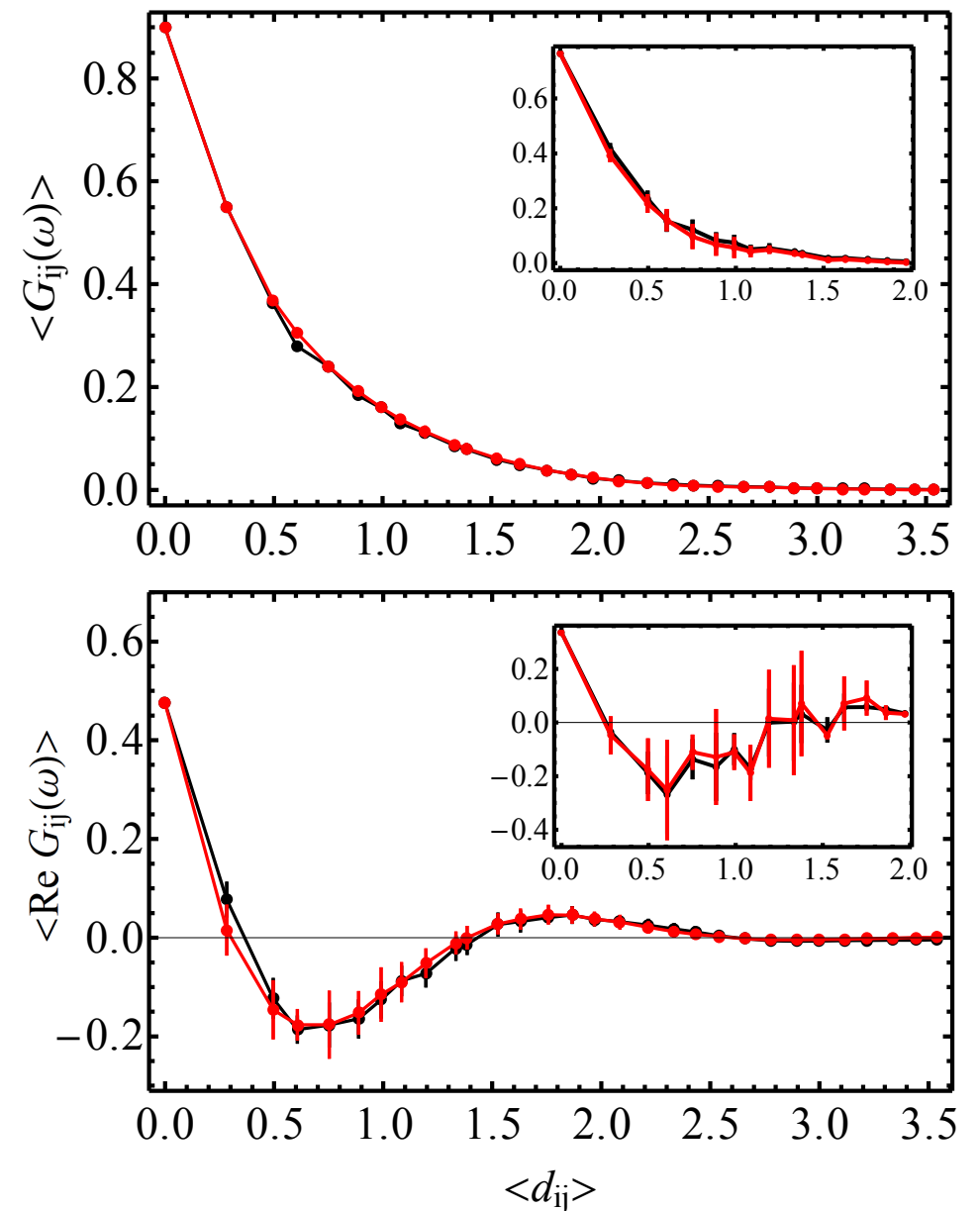
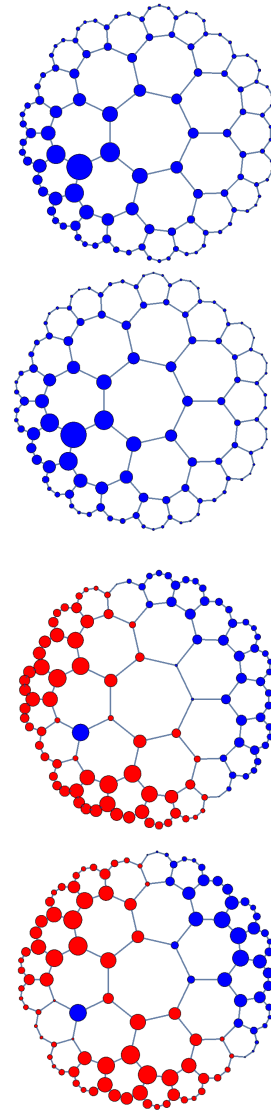
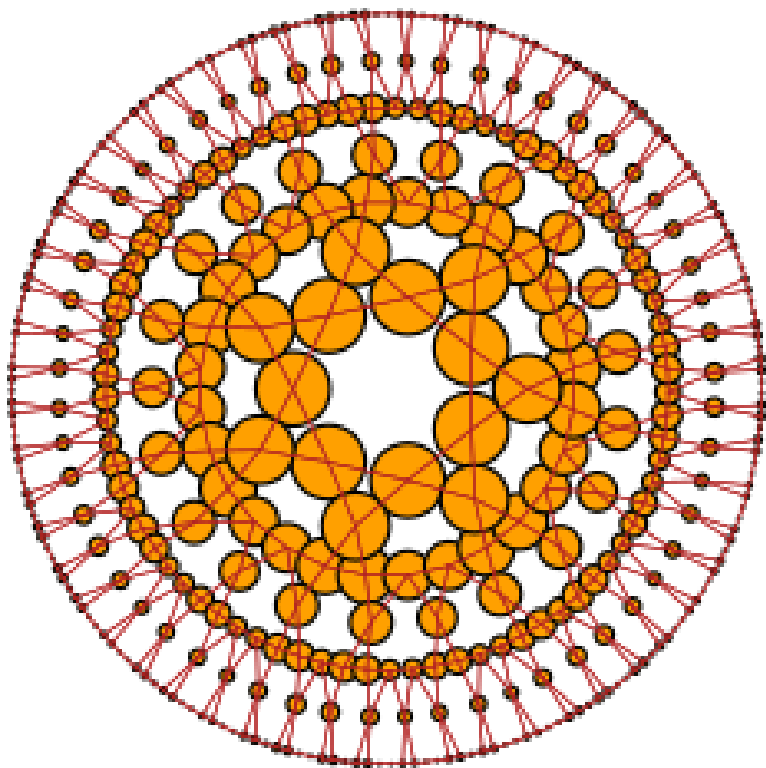
Quantitative Match for Large System Sizes

- Green's function

Continuum Limit and Green's Function

High Energy Limit of The Spectrum

- Long-wavelength modes
- Lattice should course-grain out
- Hyperbolic particle in a box



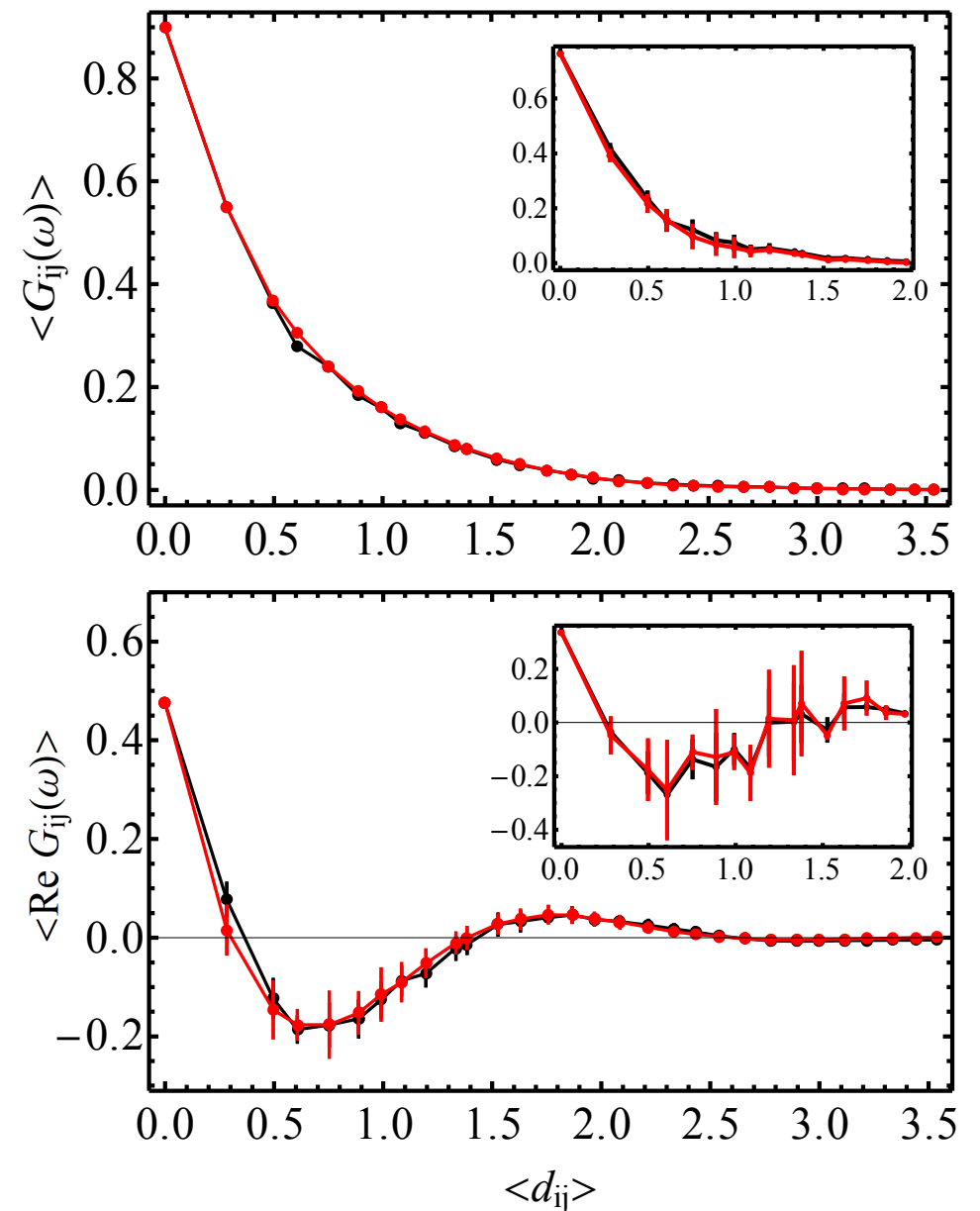
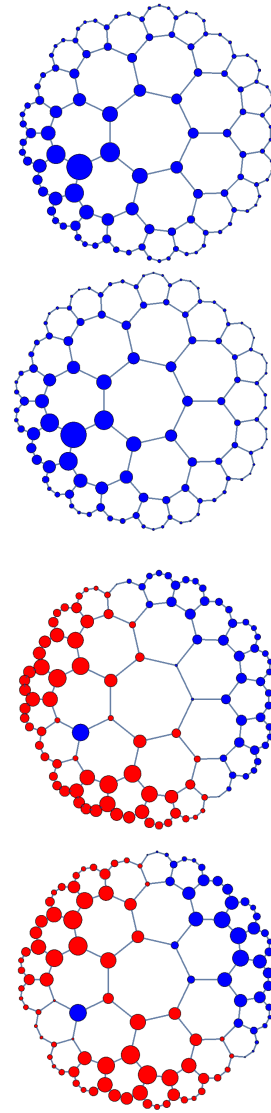
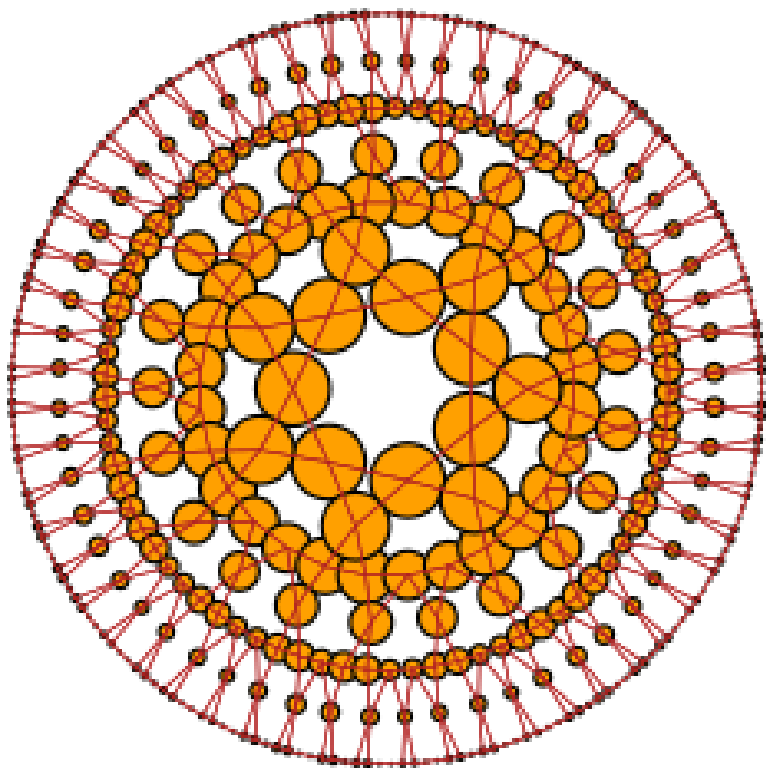
Quantitative Match for Large System Sizes

- Green's function
- "Ground" state energy

Continuum Limit and Green's Function

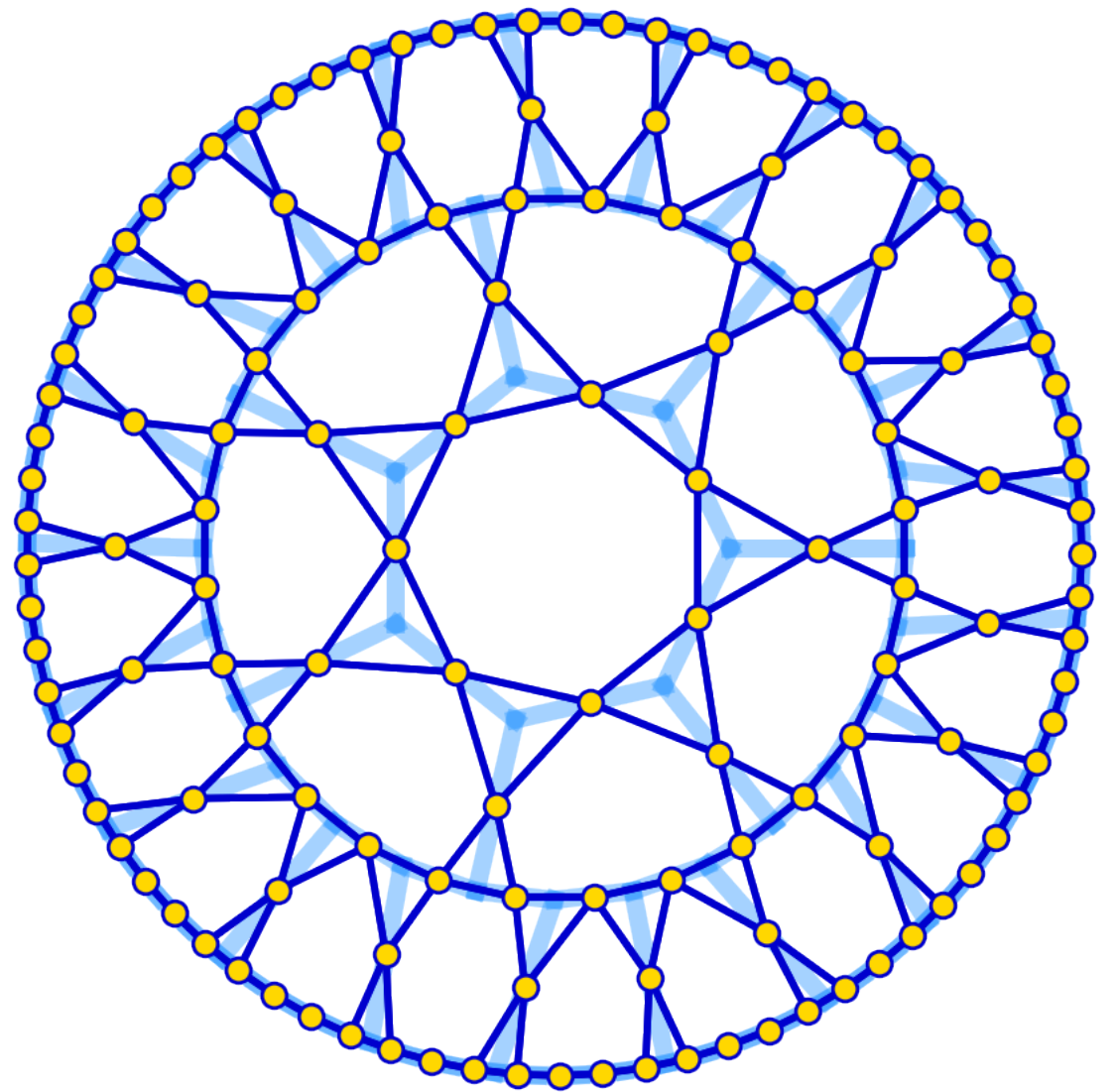
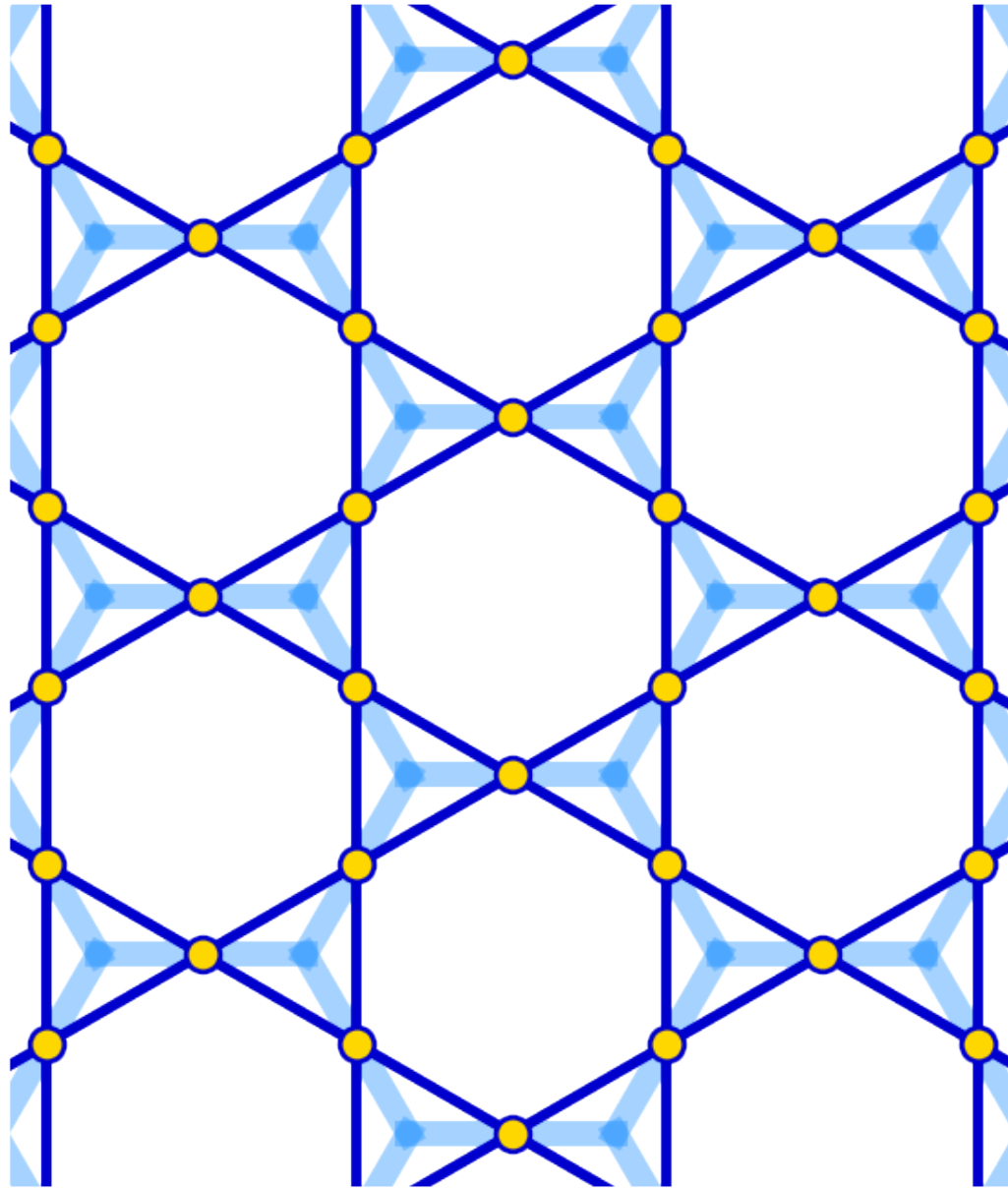
High Energy Limit of The Spectrum

- Long-wavelength modes
- Lattice should course-grain out
- Hyperbolic particle in a box



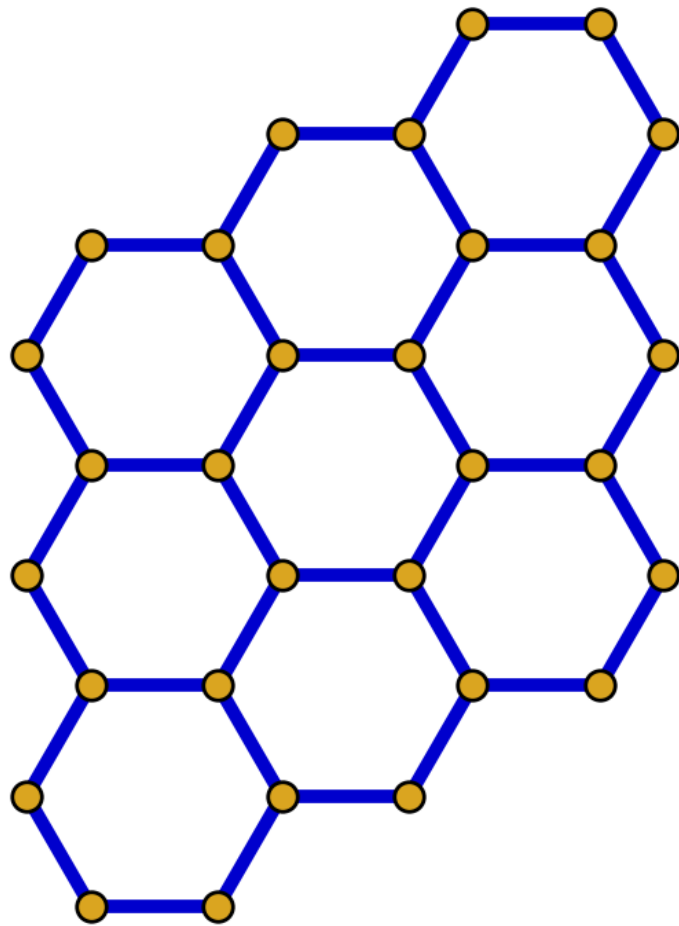
Quantitative Match for Large System Sizes

- Green's function
- "Ground" state energy
- "First" excited state energies.



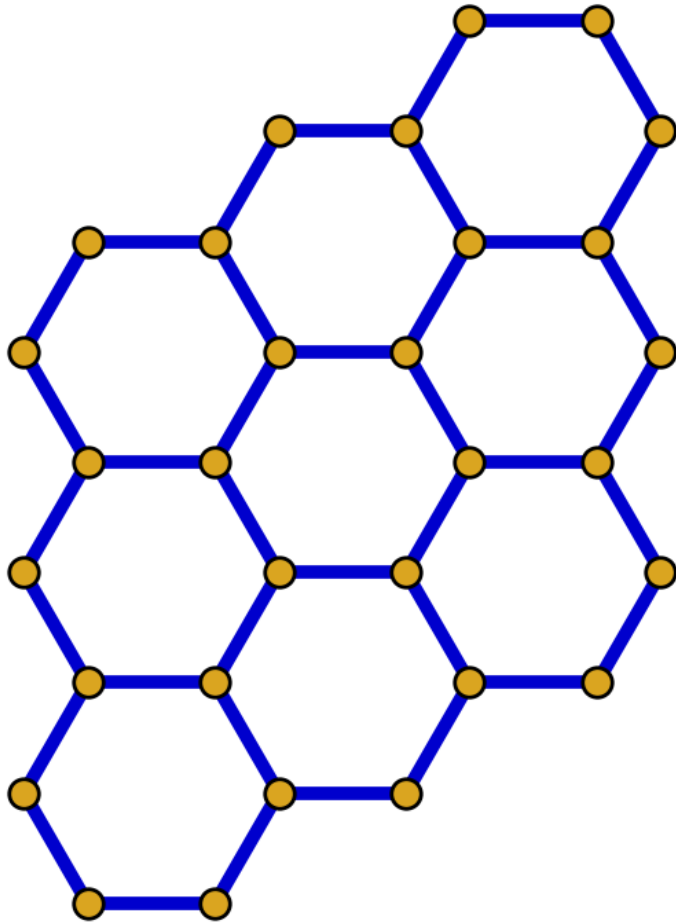
Bipartite and Non-Bipartite Graphs

Bipartite

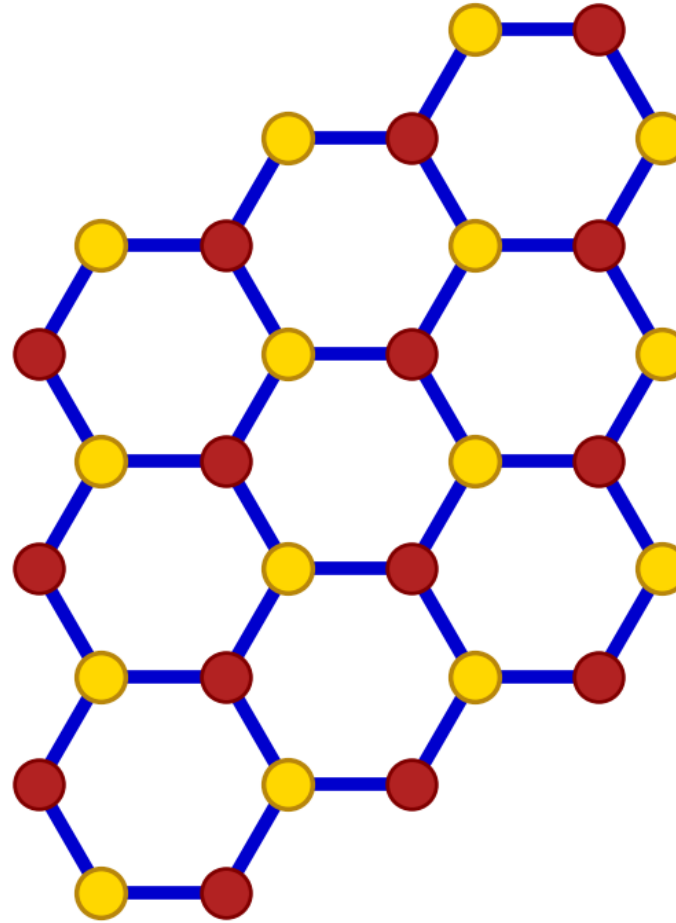


Bipartite and Non-Bipartite Graphs

Bipartite

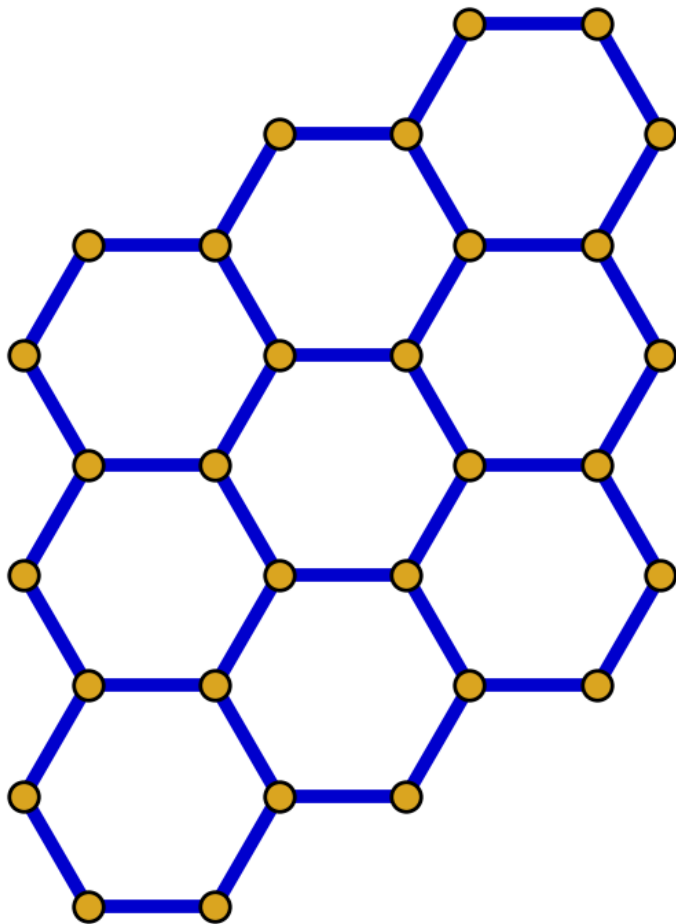


$E = -3$

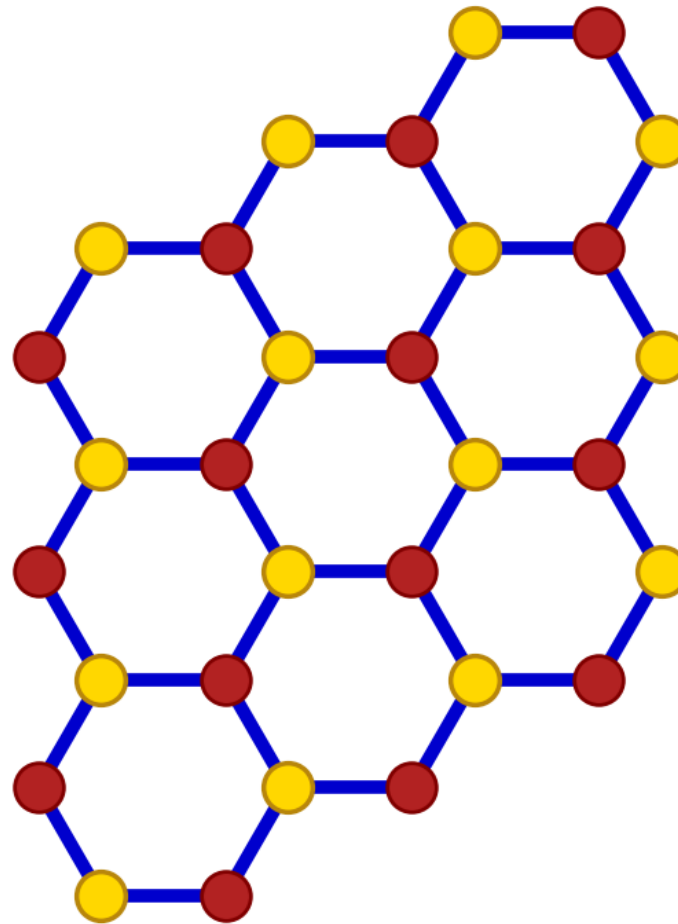


Bipartite and Non-Bipartite Graphs

Bipartite



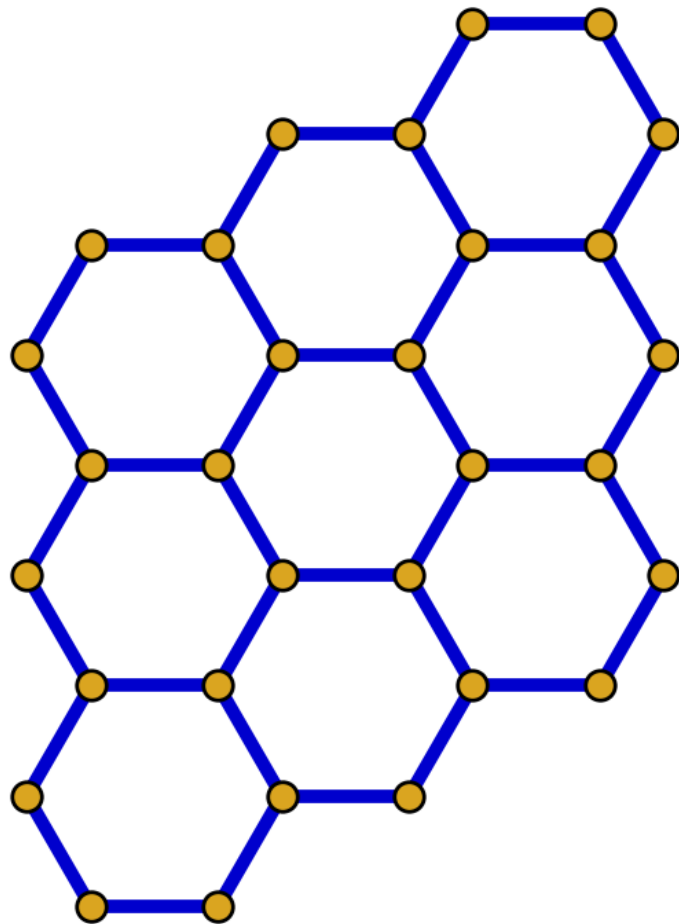
$E = -3$



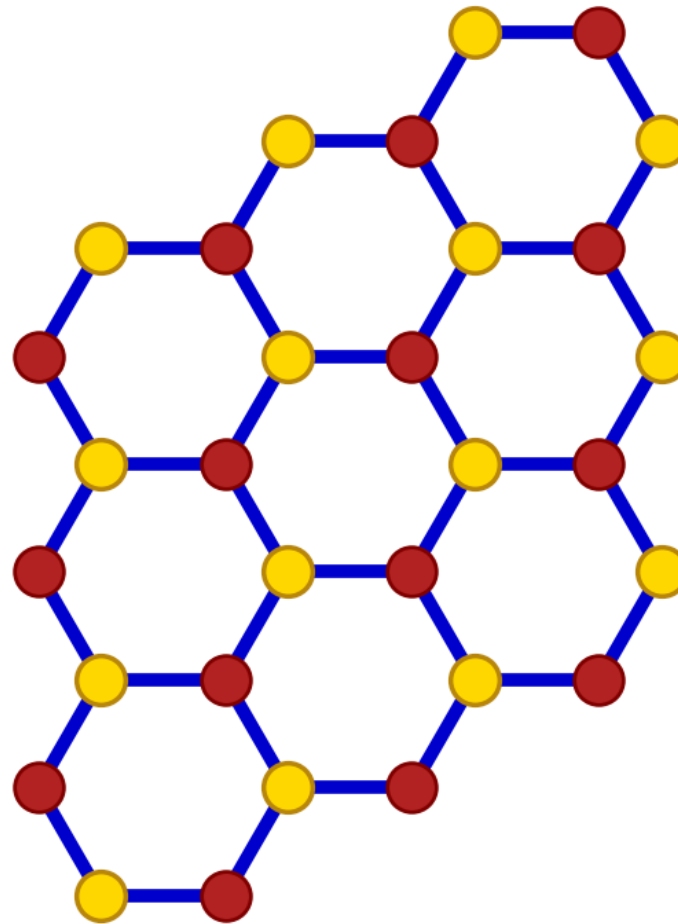
● All neighbors opposite sign

Bipartite and Non-Bipartite Graphs

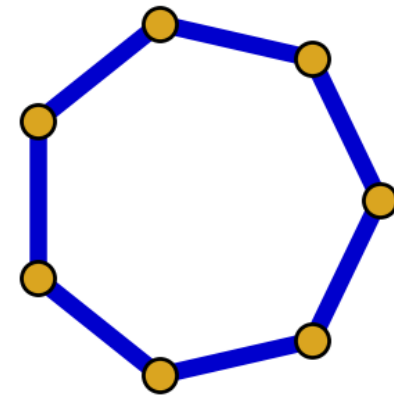
Bipartite



$$E = -3$$



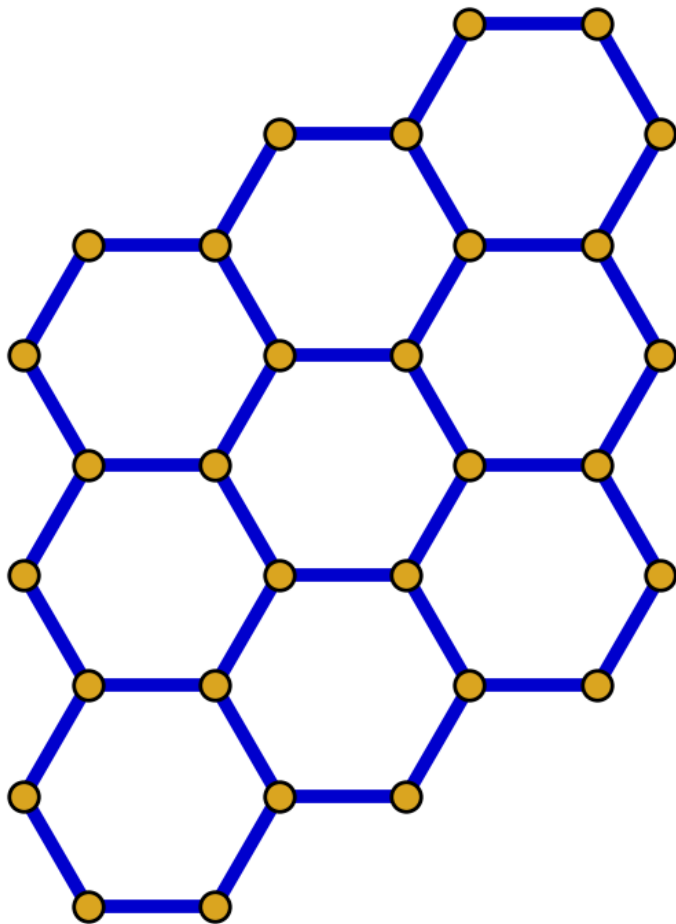
Non-Bipartite



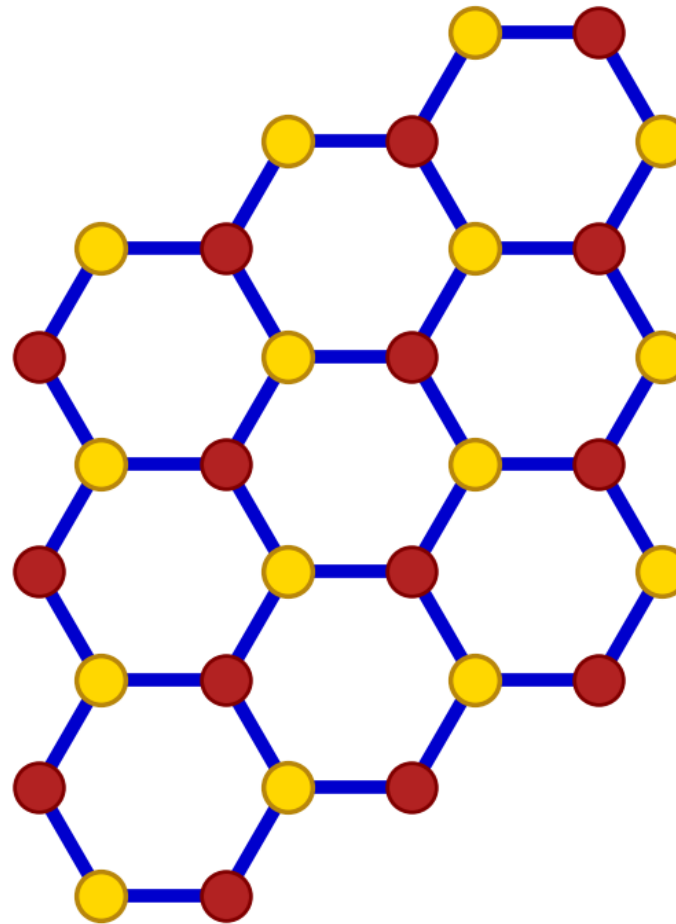
● All neighbors opposite sign

Bipartite and Non-Bipartite Graphs

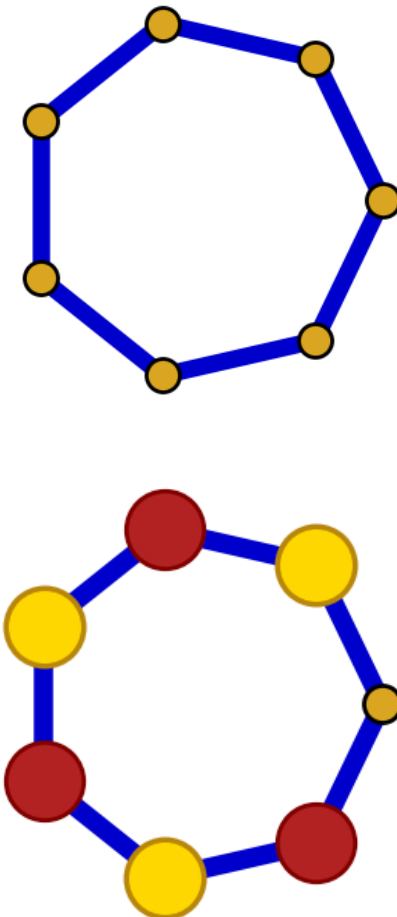
Bipartite



$$E = -3$$



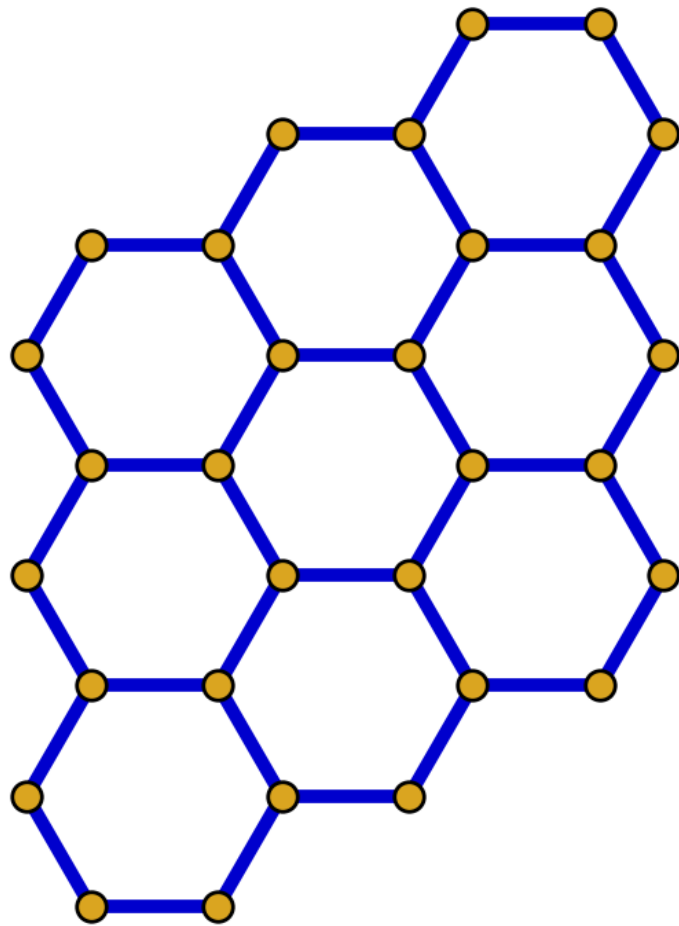
Non-Bipartite



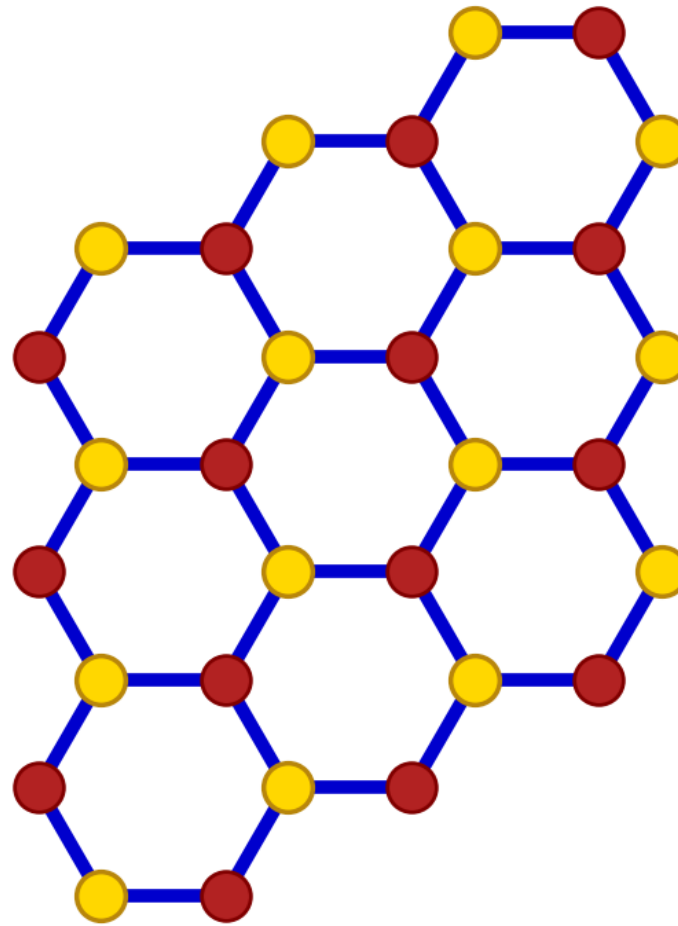
● All neighbors opposite sign

Bipartite and Non-Bipartite Graphs

Bipartite

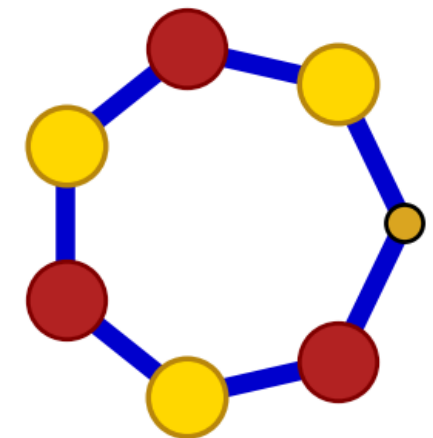
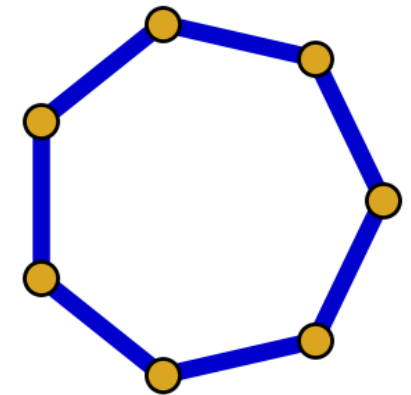


$$E = -3$$



● All neighbors opposite sign

Non-Bipartite



● Not all neighbors can be opposite sign

Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$

$$\bar{H}_s(X) = H_{L(X)}$$

Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Incidence Operator

- From X to $L(X)$

$$M : \ell^2(X) \rightarrow \ell^2(L(X))$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$

$$\bar{H}_s(X) = H_{L(X)}$$

Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$

$$\bar{H}_s(X) = H_{L(X)}$$

Incidence Operator

- From X to $L(X)$

$$M : \ell^2(X) \rightarrow \ell^2(L(X))$$

$$M(v, e) = \begin{cases} 1, & \text{if } e \text{ and } v \text{ are incident,} \\ 0 & \text{otherwise.} \end{cases}$$

Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$

$$\bar{H}_s(X) = H_{L(X)}$$

Incidence Operator

- From X to $L(X)$

$$M : \ell^2(X) \rightarrow \ell^2(L(X))$$

$$M(v, e) = \begin{cases} 1, & \text{if } e \text{ and } v \text{ are incident,} \\ 0 & \text{otherwise.} \end{cases}$$

$$M^t M = D_X + H_X$$

$$M M^t = 2I + \bar{H}_s(X)$$

Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$

$$\bar{H}_s(X) = H_{L(X)}$$

Incidence Operator

- From X to $L(X)$

$$M : \ell^2(X) \rightarrow \ell^2(L(X))$$

$$M(v, e) = \begin{cases} 1, & \text{if } e \text{ and } v \text{ are incident,} \\ 0 & \text{otherwise.} \end{cases}$$

$$M^t M = D_X + H_X$$

$$M M^t = 2I + \bar{H}_s(X)$$

$$D_X + H_X \simeq 2I + \bar{H}_s(X)$$

Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$

$$\bar{H}_s(X) = H_{L(X)}$$

Incidence Operator

- From X to $L(X)$

$$M : \ell^2(X) \rightarrow \ell^2(L(X))$$

$$M(v, e) = \begin{cases} 1, & \text{if } e \text{ and } v \text{ are incident,} \\ 0 & \text{otherwise.} \end{cases}$$

$$M^t M = D_X + H_X$$

$$M M^t = 2I + \bar{H}_s(X)$$

$$D_X + H_X \simeq 2I + \bar{H}_s(X)$$

$$E_{\bar{H}_s} = \begin{cases} d - 2 + E_{H_X} \\ -2 \end{cases}$$

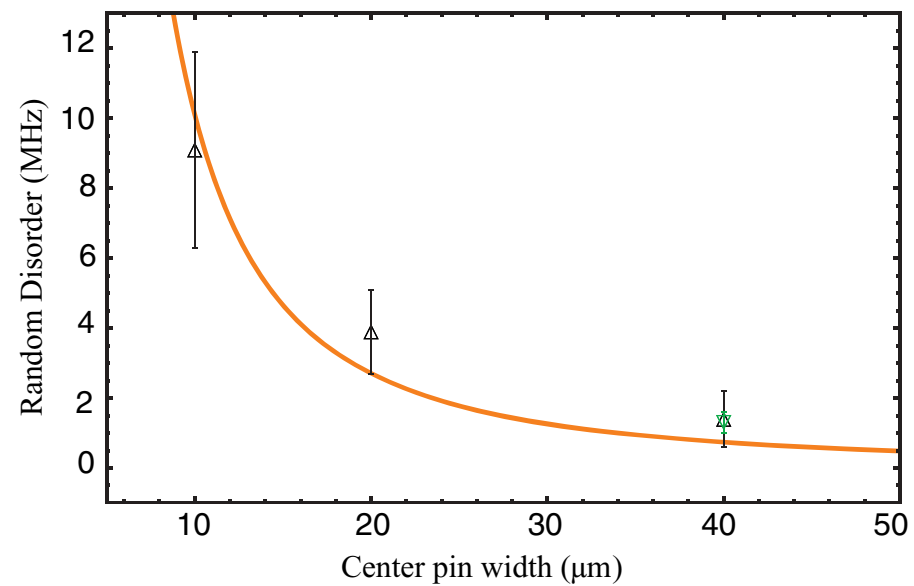
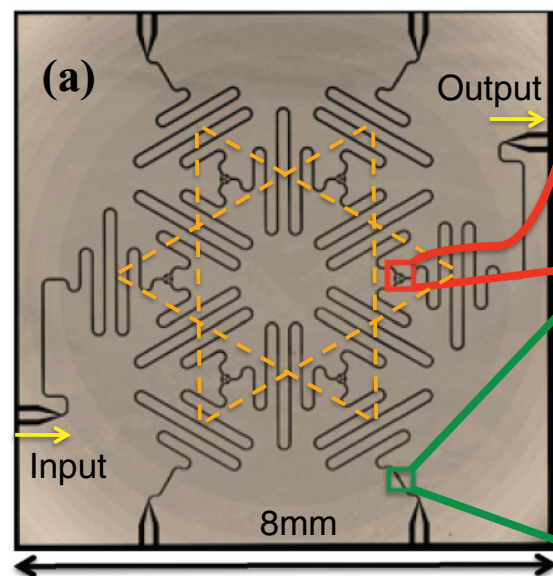
Empty rounded rectangular box at the top of the page.

Large empty rounded rectangular box occupying the majority of the page.

Intrinsic Fabrication Disorder

Previous Benchmarks

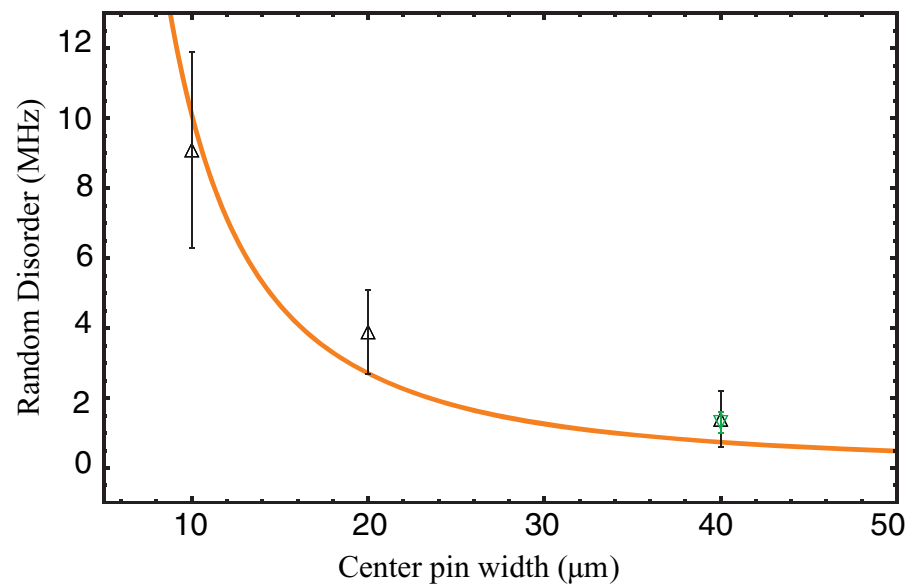
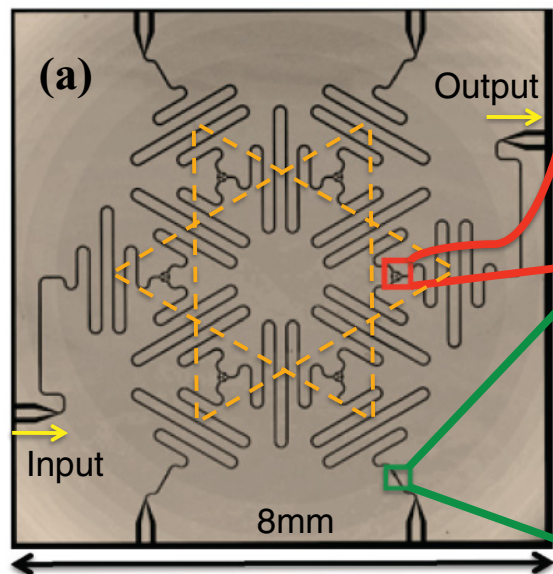
- Kagome star normal modes
- Fabricated at Princeton
- Fabrication disorder $\sim 3e-4$



Intrinsic Fabrication Disorder

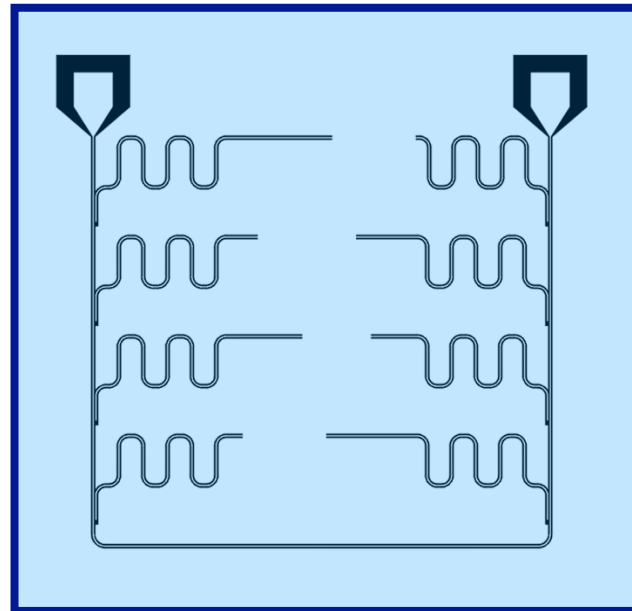
Previous Benchmarks

- Kagome star normal modes
- Fabricated at Princeton
- Fabrication disorder $\sim 3e-4$



Current Devices

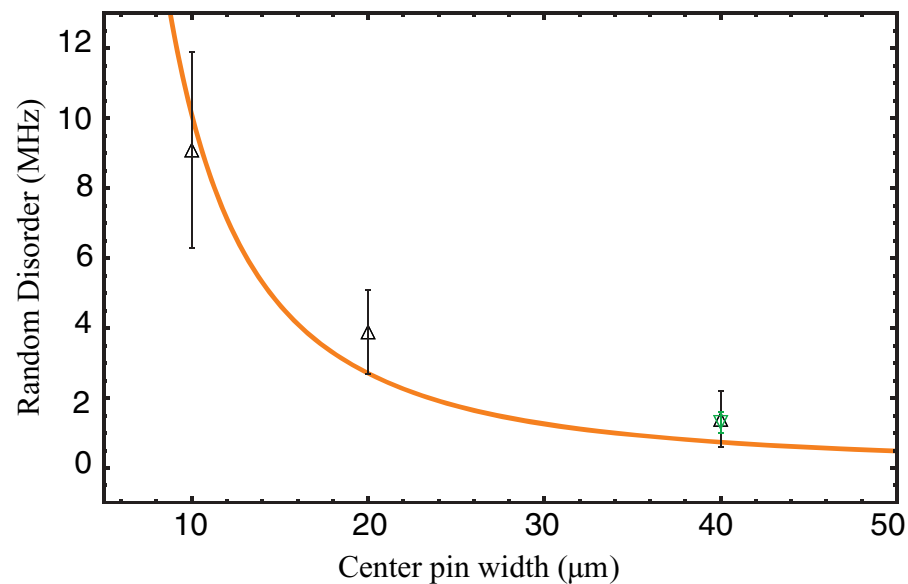
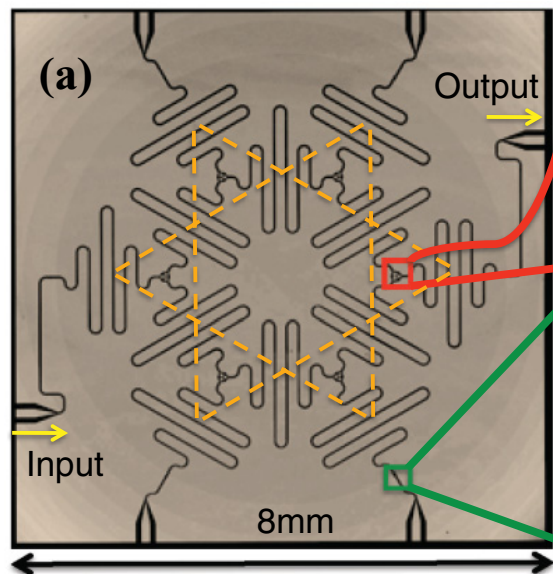
- Fabricated at UMD
- Fabrication disorder $\sim 3e-4$



Intrinsic Fabrication Disorder

Previous Benchmarks

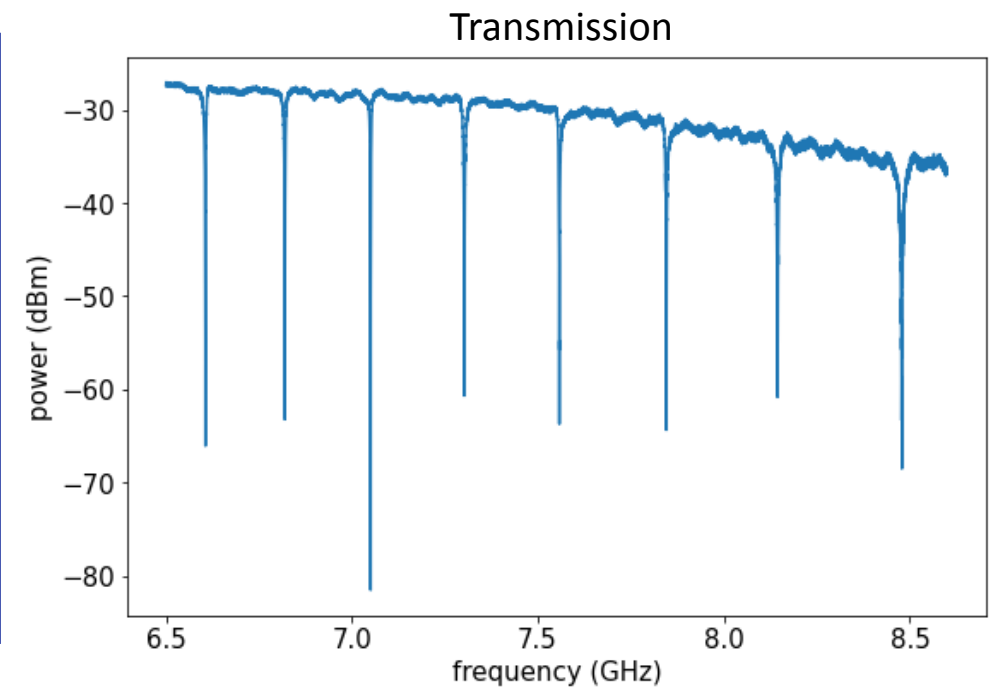
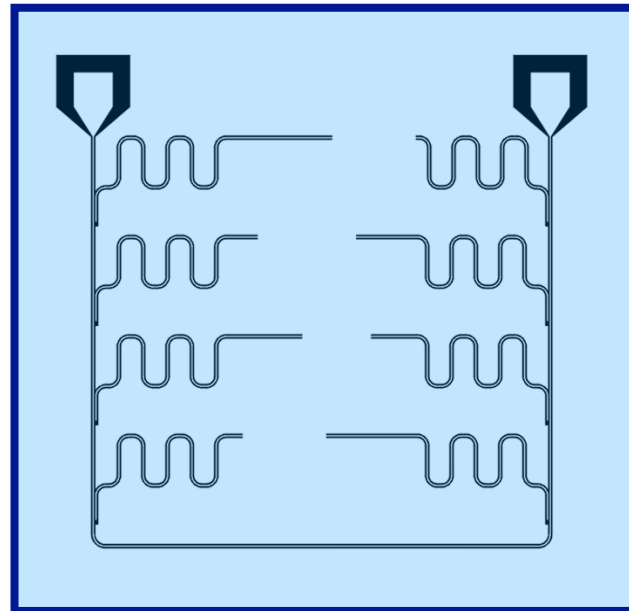
- Kagome star normal modes
- Fabricated at Princeton
- Fabrication disorder $\sim 3e-4$



Underwood *et al.* PRA **86**, 023837 (2012)

Current Devices

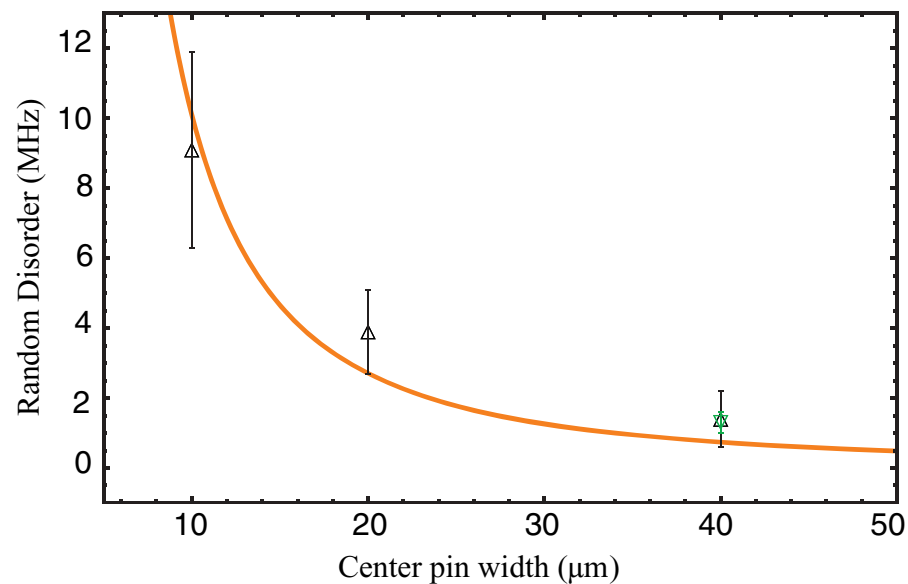
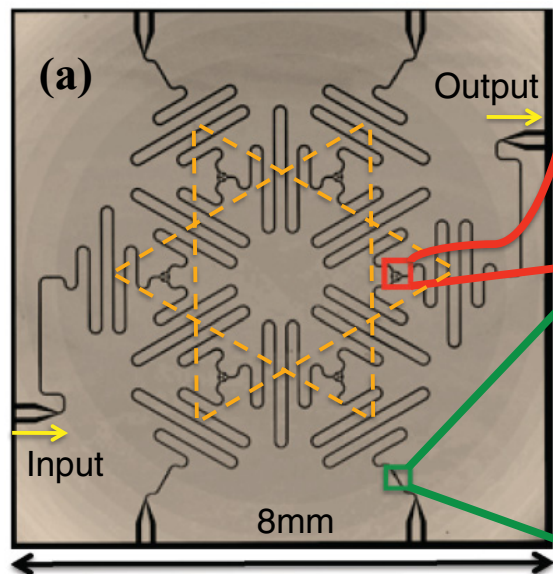
- Fabricated at UMD
- Fabrication disorder $\sim 3e-4$



Intrinsic Fabrication Disorder

Previous Benchmarks

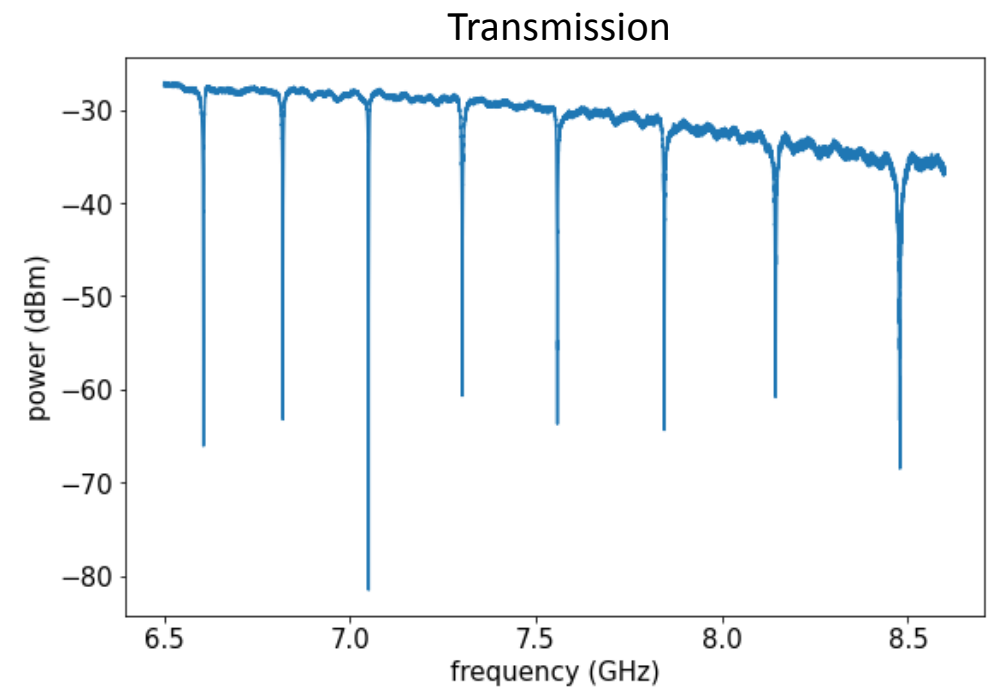
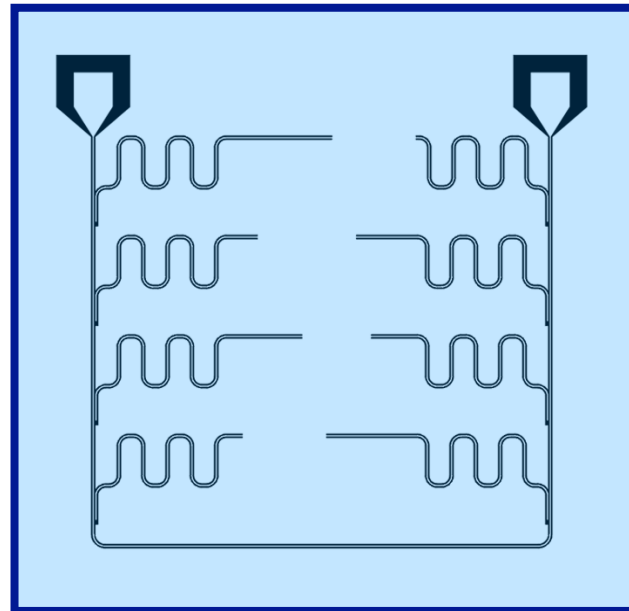
- Kagome star normal modes
- Fabricated at Princeton
- Fabrication disorder $\sim 3e-4$



Underwood *et al.* PRA **86**, 023837 (2012)

Current Devices

- Fabricated at UMD
- Fabrication disorder $\sim 3e-4$

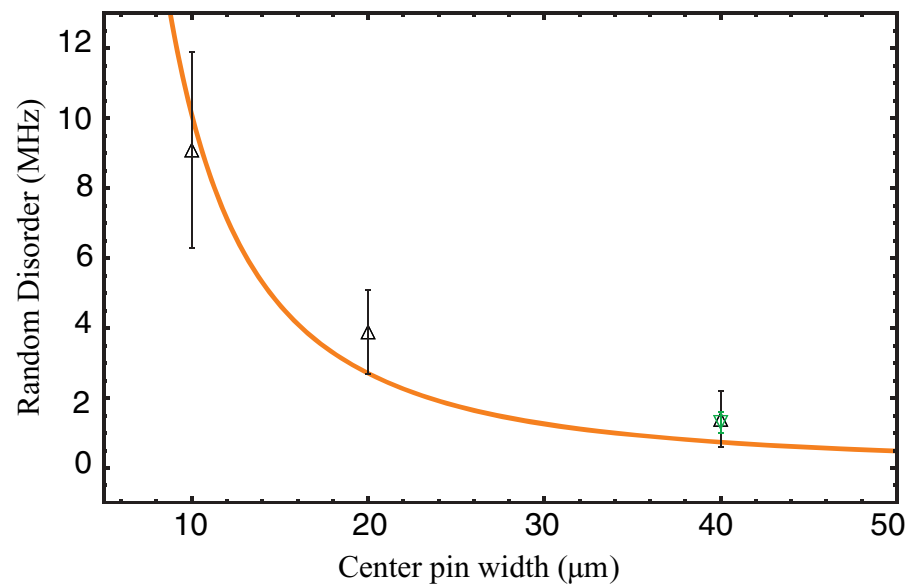
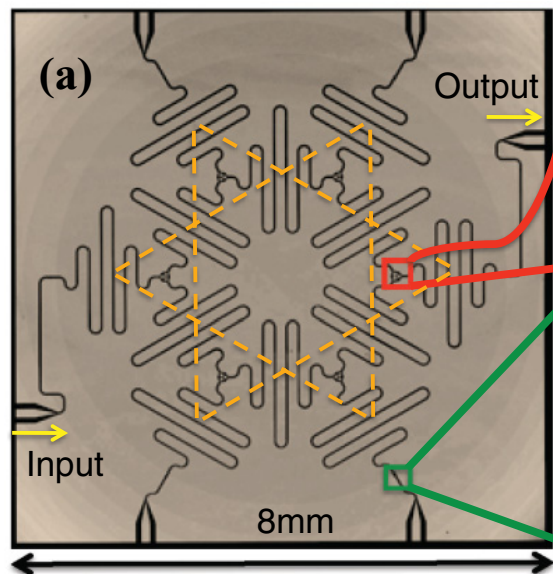


- Parallel measurement

Intrinsic Fabrication Disorder

Previous Benchmarks

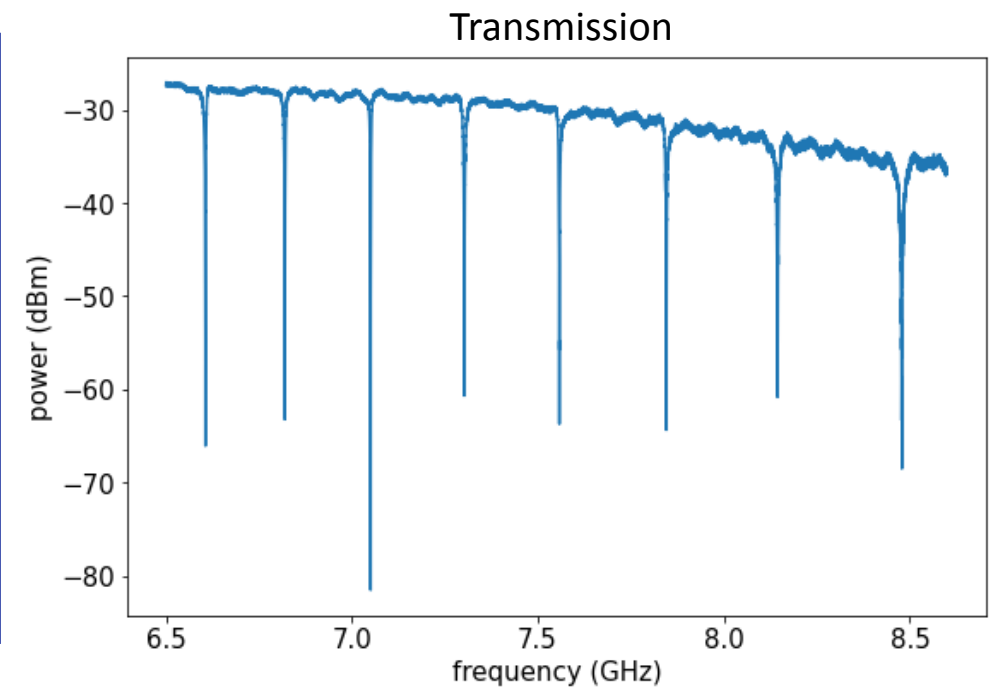
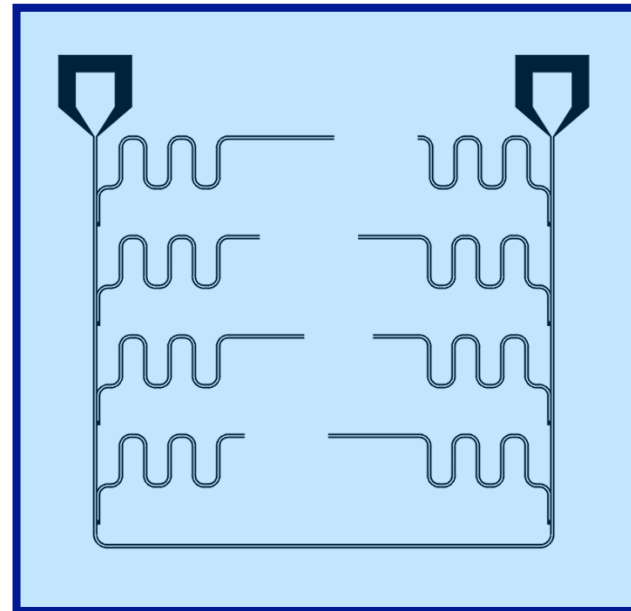
- Kagome star normal modes
- Fabricated at Princeton
- Fabrication disorder $\sim 3e-4$



Underwood *et al.* PRA **86**, 023837 (2012)

Current Devices

- Fabricated at UMD
- Fabrication disorder $\sim 3e-4$

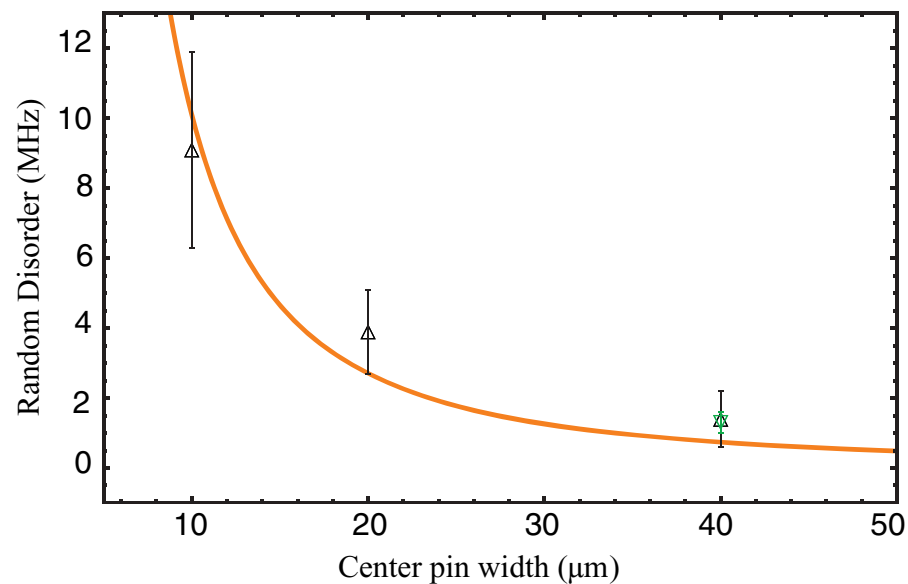
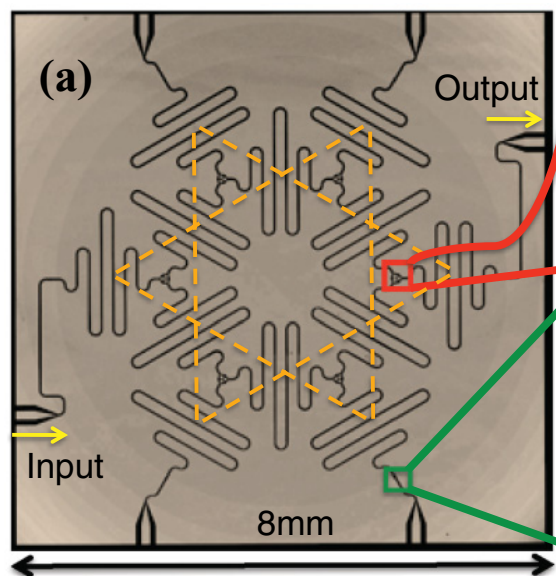


- Parallel measurement
- Disorder extracted from comb spacing

Intrinsic Fabrication Disorder

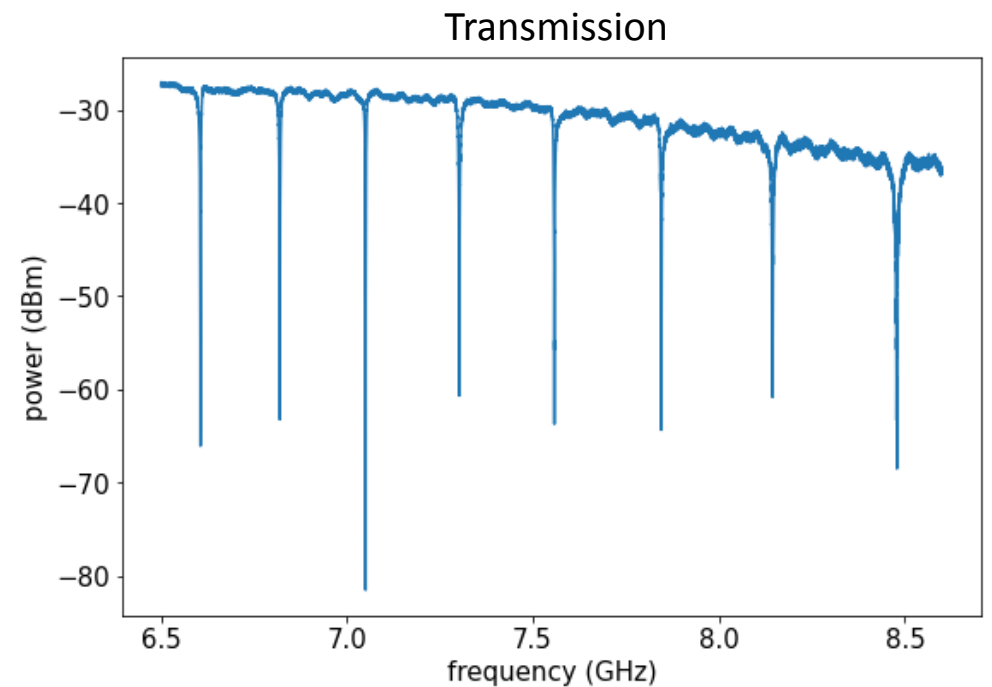
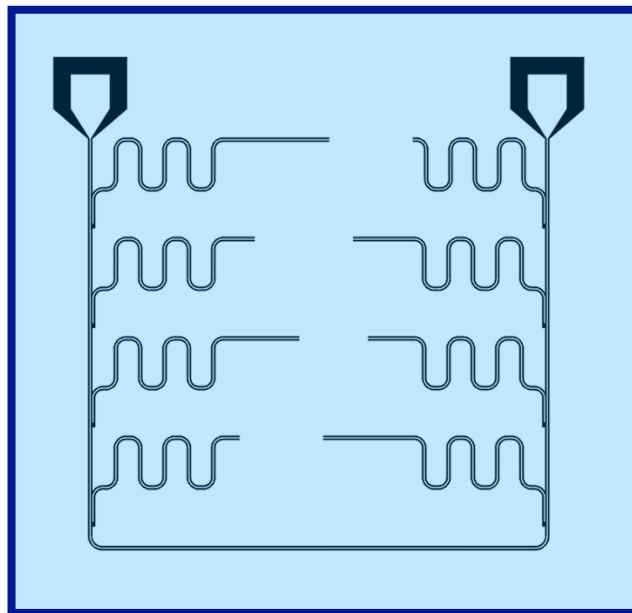
Previous Benchmarks

- Kagome star normal modes
- Fabricated at Princeton
- Fabrication disorder $\sim 3e-4$

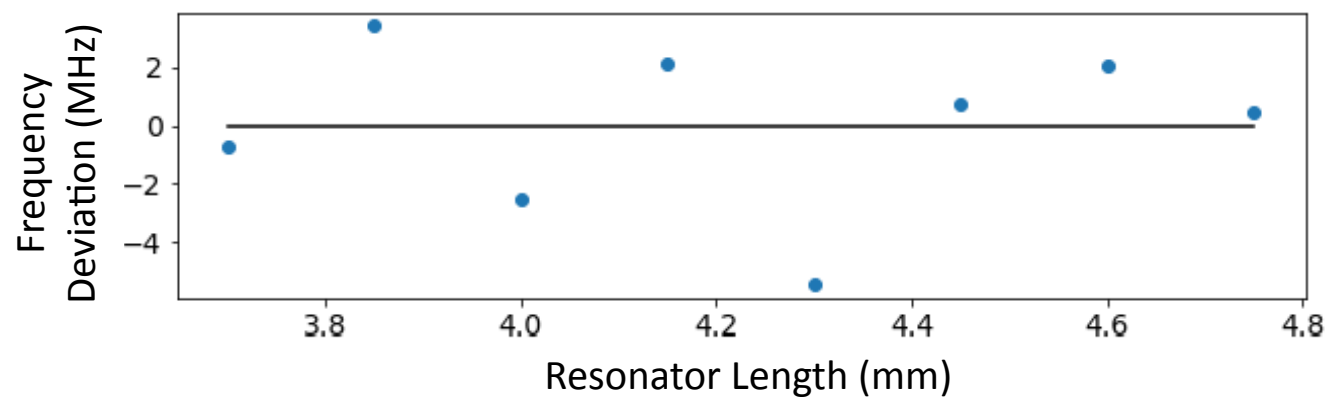


Current Devices

- Fabricated at UMD
- Fabrication disorder $\sim 3e-4$



- Parallel measurement
- Disorder extracted from comb spacing



Disorder Mitigation

Systematic v. Random Disorder

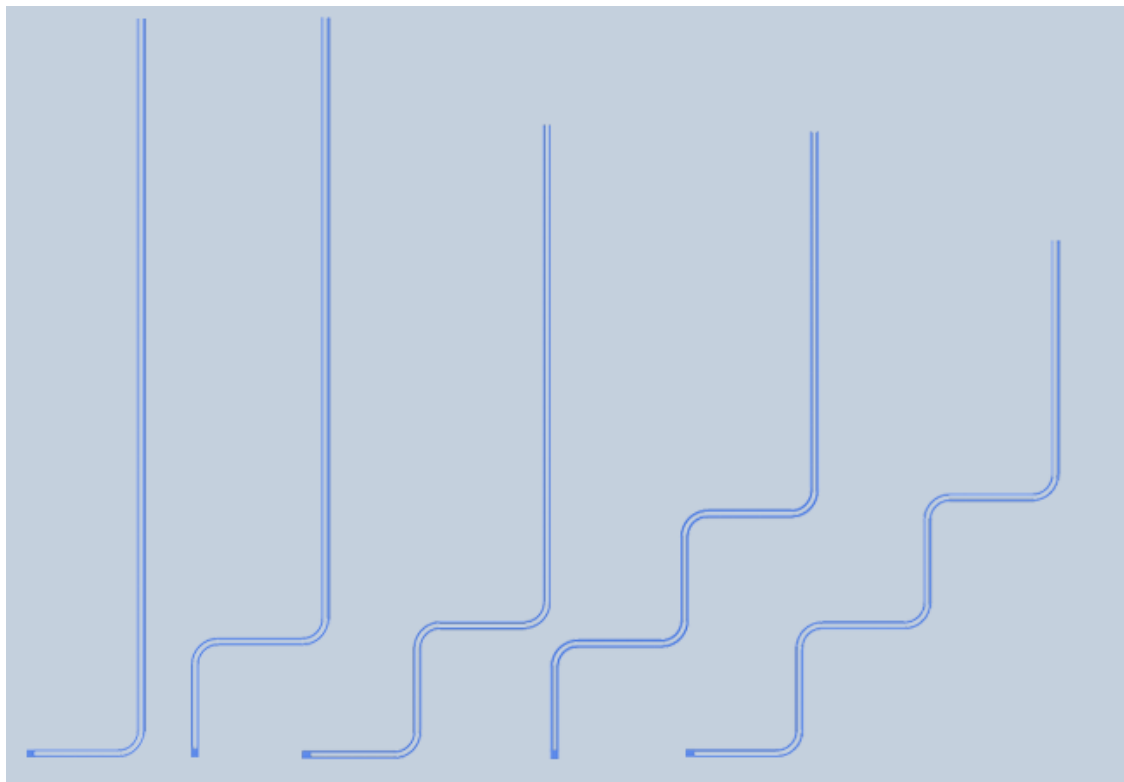
- Fabrication disorder $\sim 3e-4$
- Shape-dependent disorder $\sim 2-3e-3$

Disorder Mitigation

Systematic v. Random Disorder

- Fabrication disorder $\sim 3e-4$
- Shape-dependent disorder $\sim 2-3e-3$

Numerical Test Geometries

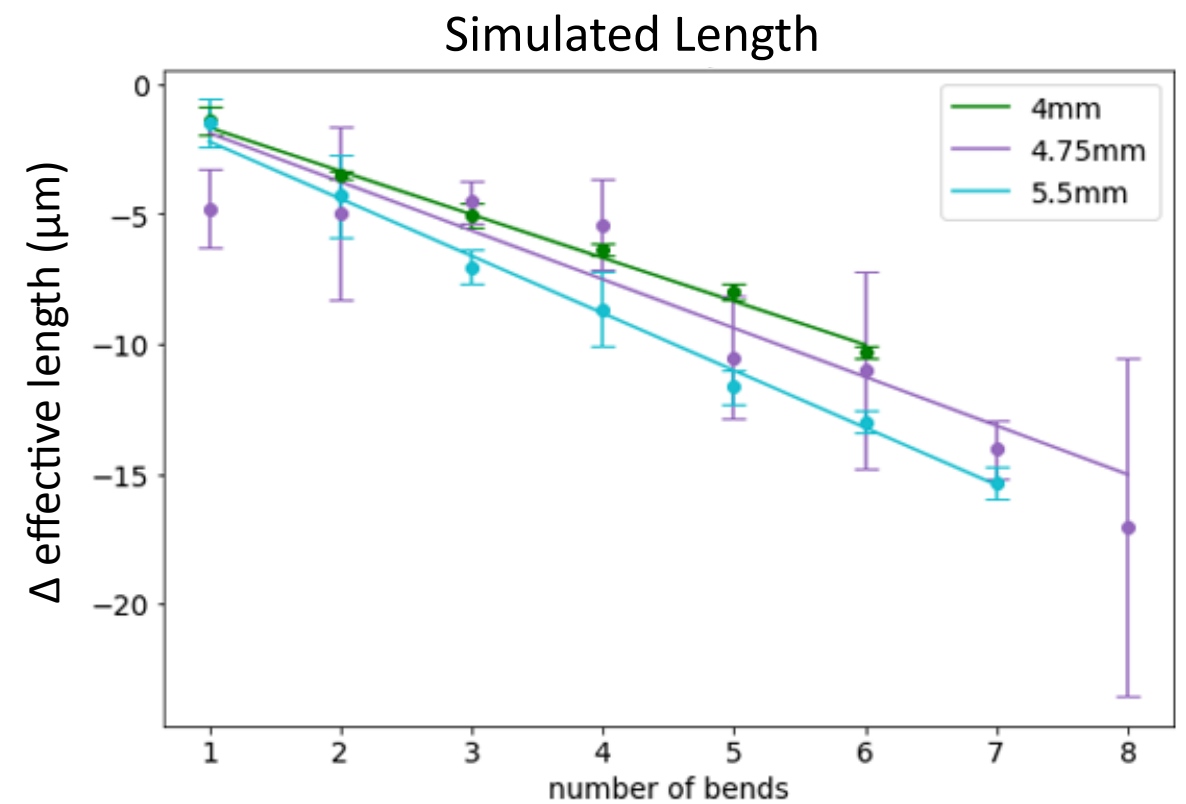
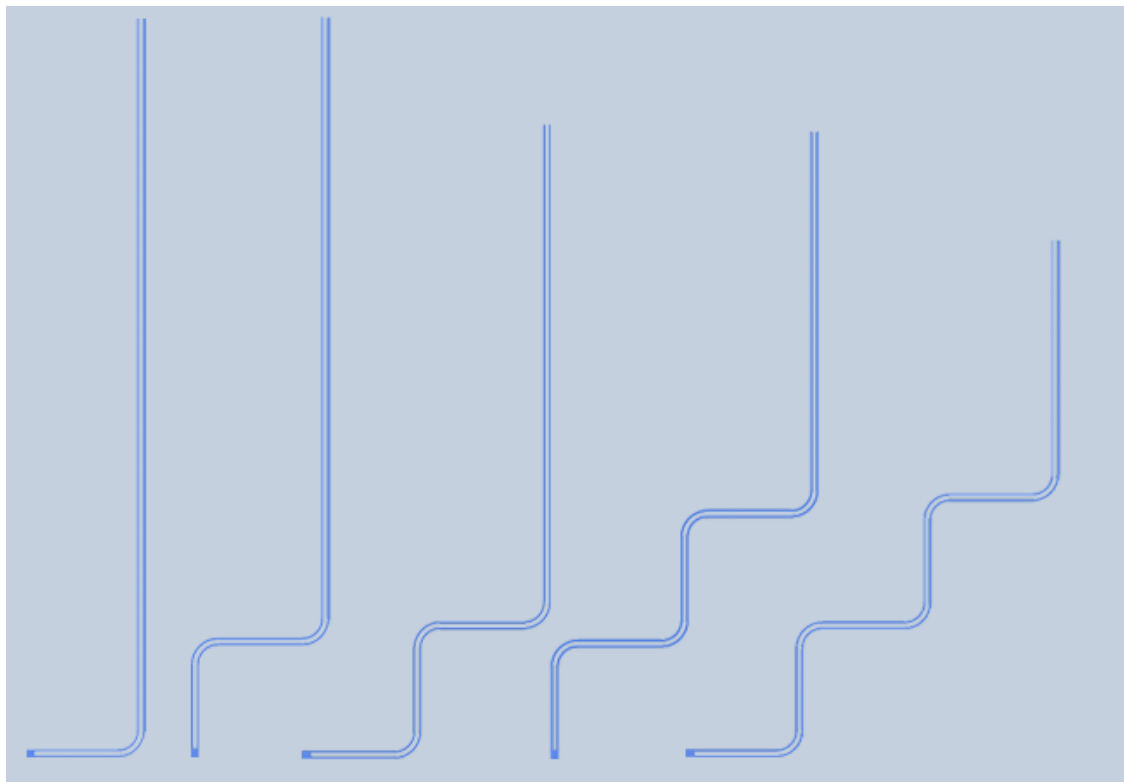


Disorder Mitigation

Systematic v. Random Disorder

- Fabrication disorder $\sim 3e-4$
- Shape-dependent disorder $\sim 2-3e-3$

Numerical Test Geometries

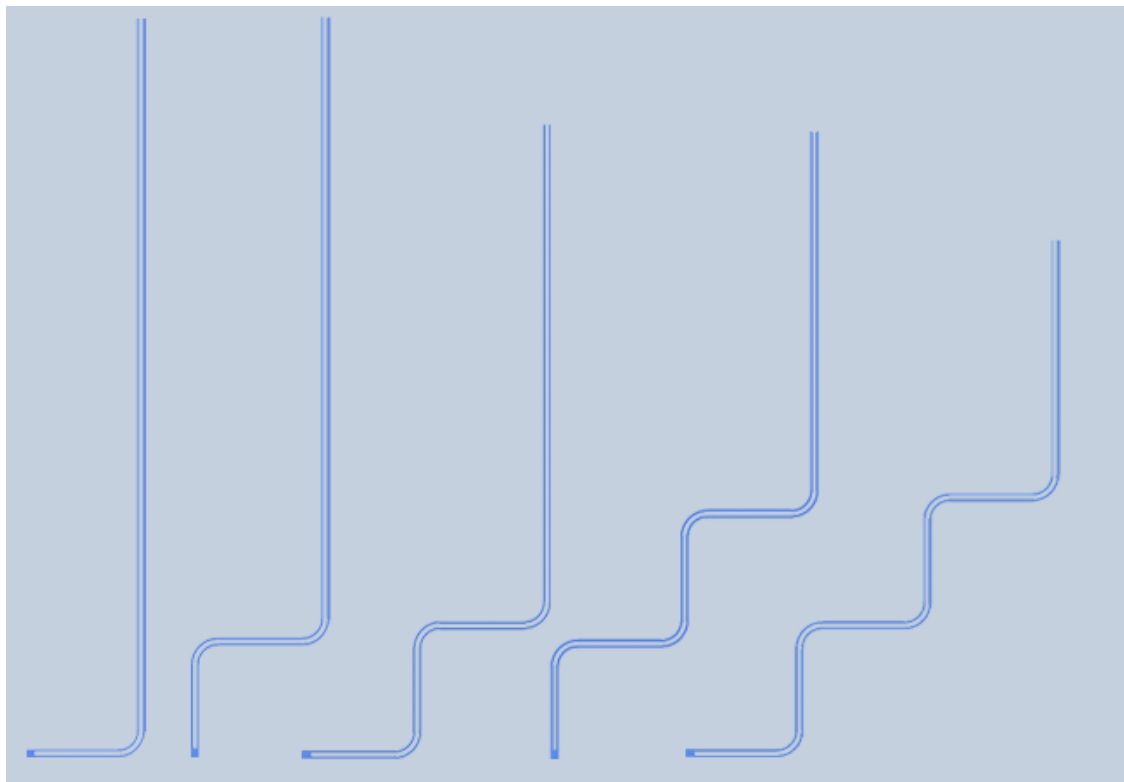


Disorder Mitigation

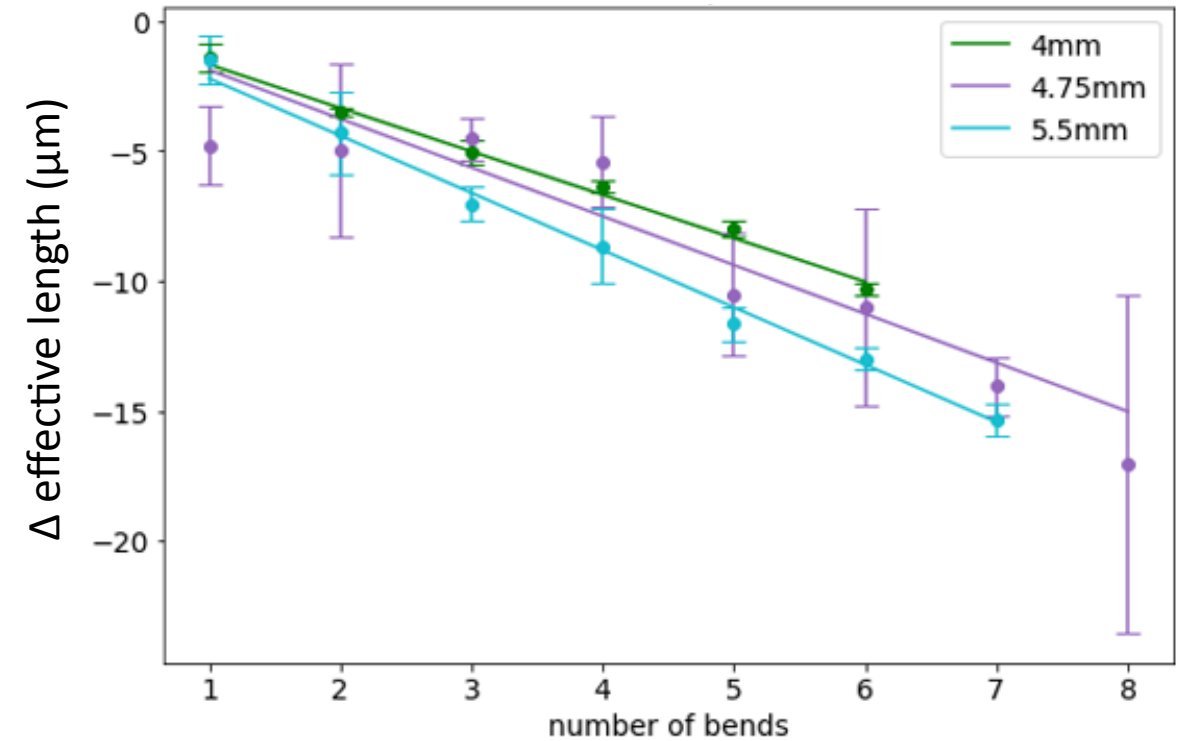
Systematic v. Random Disorder

- Fabrication disorder $\sim 3e-4$
- Shape-dependent disorder $\sim 2-3e-3$

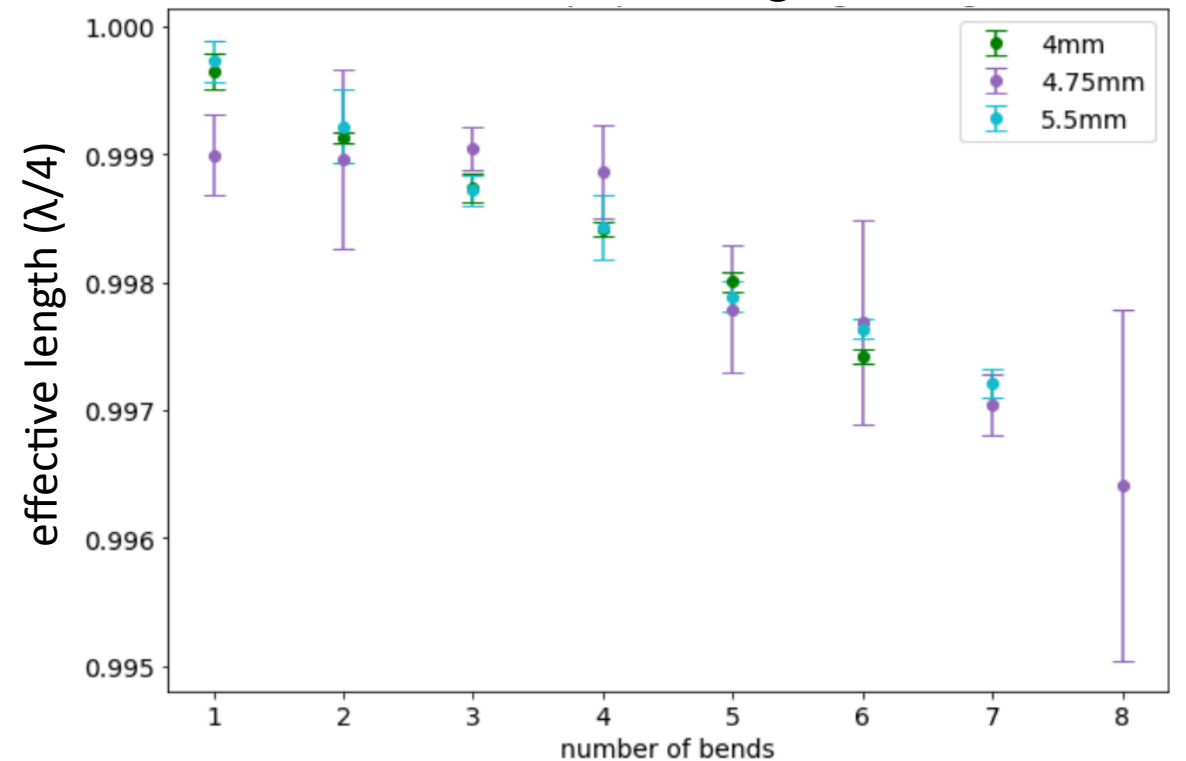
Numerical Test Geometries



Simulated Length



Normalized Length

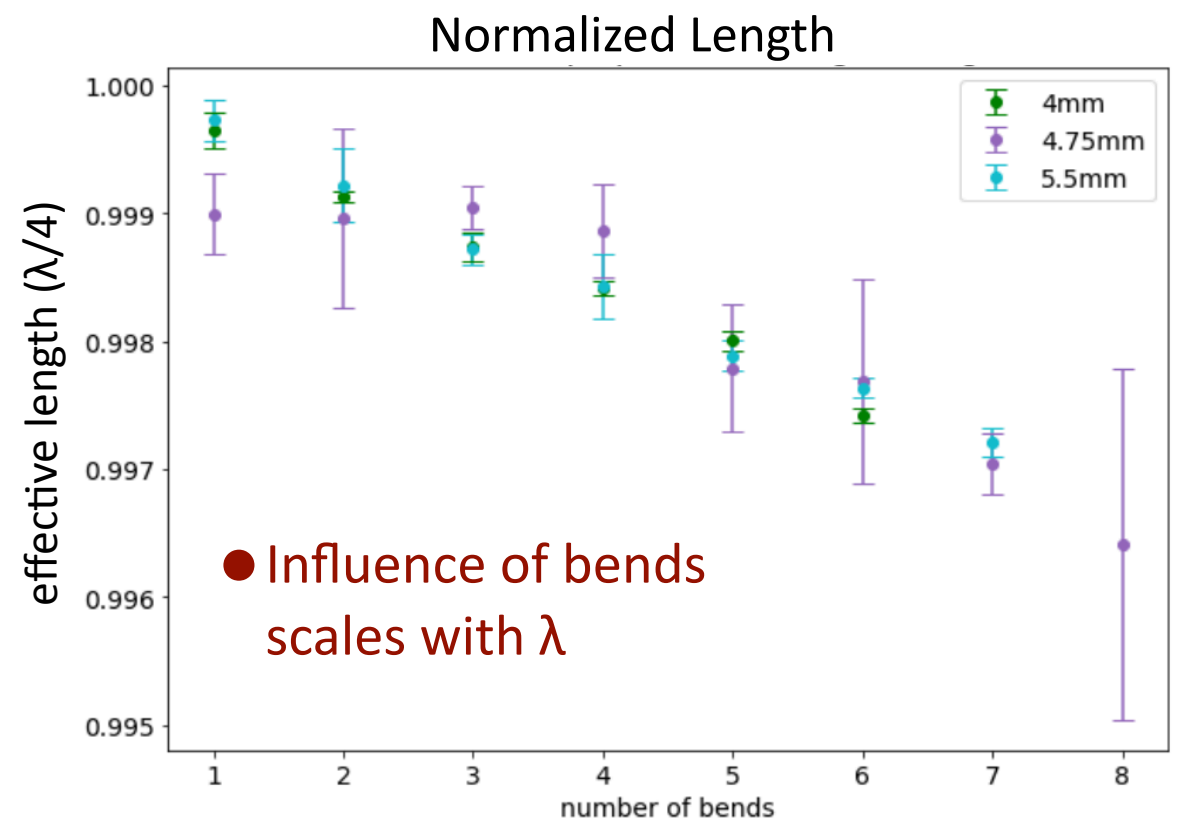
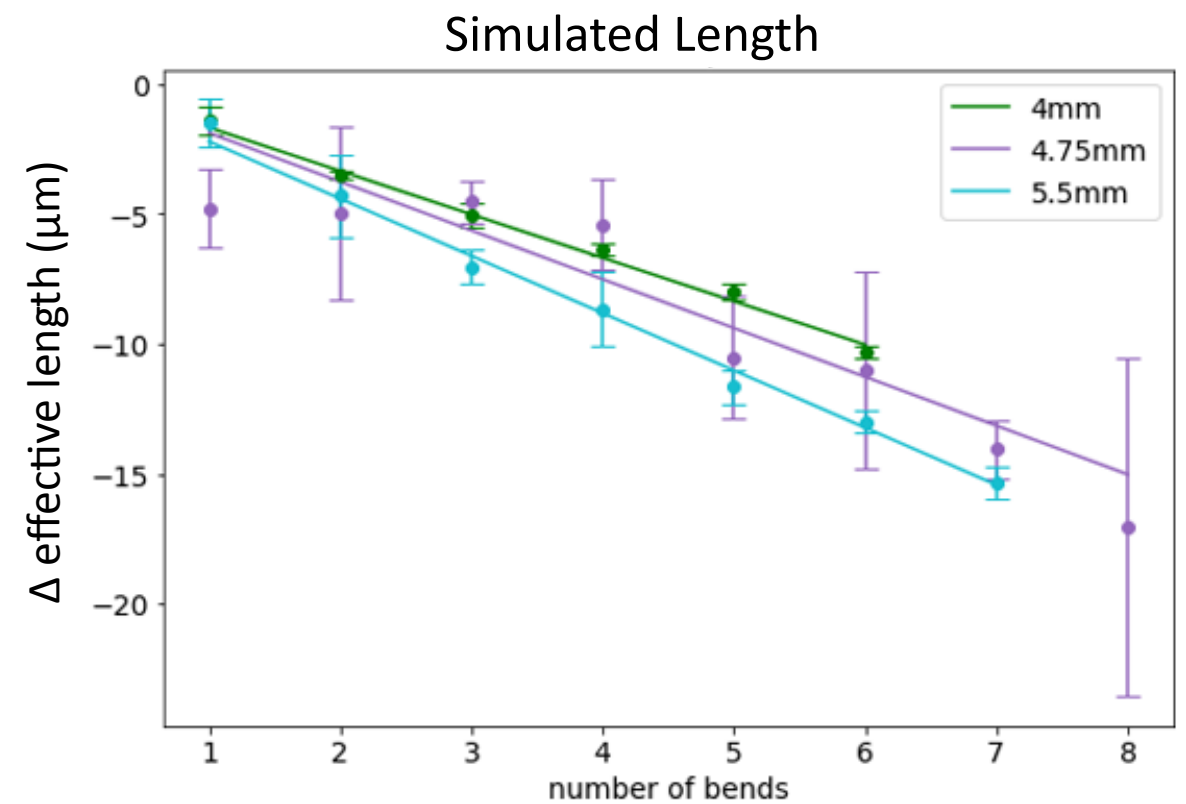
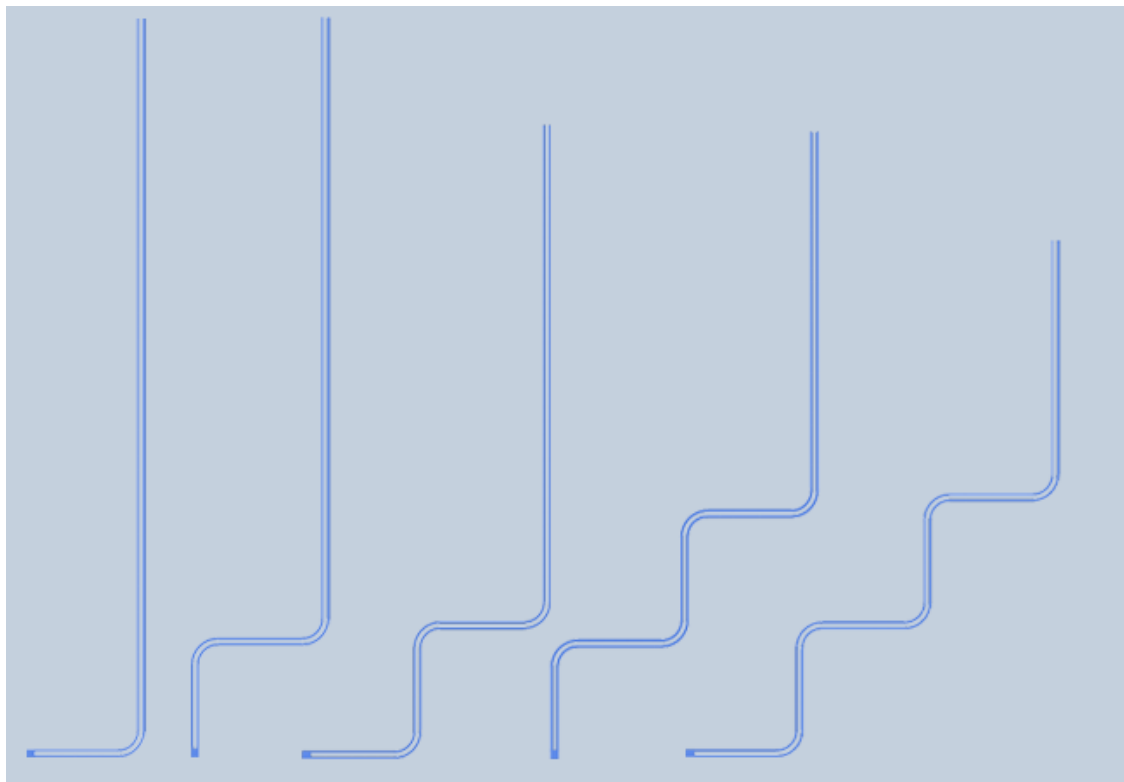


Disorder Mitigation

Systematic v. Random Disorder

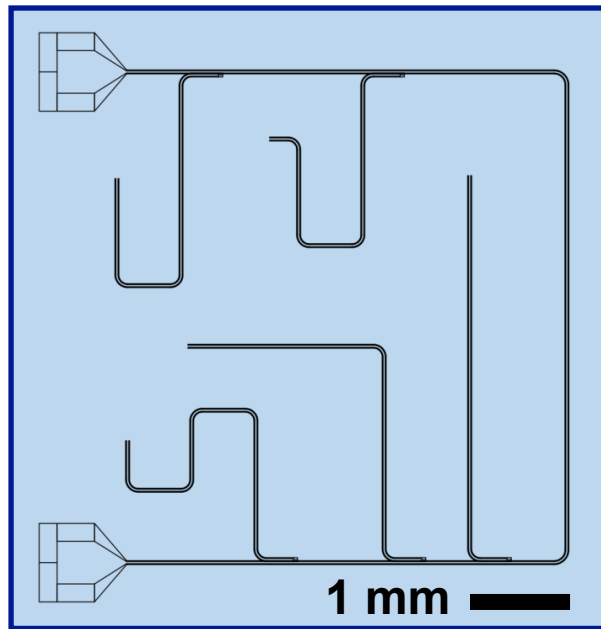
- Fabrication disorder $\sim 3e-4$
- Shape-dependent disorder $\sim 2-3e-3$

Numerical Test Geometries



Disorder Mitigation

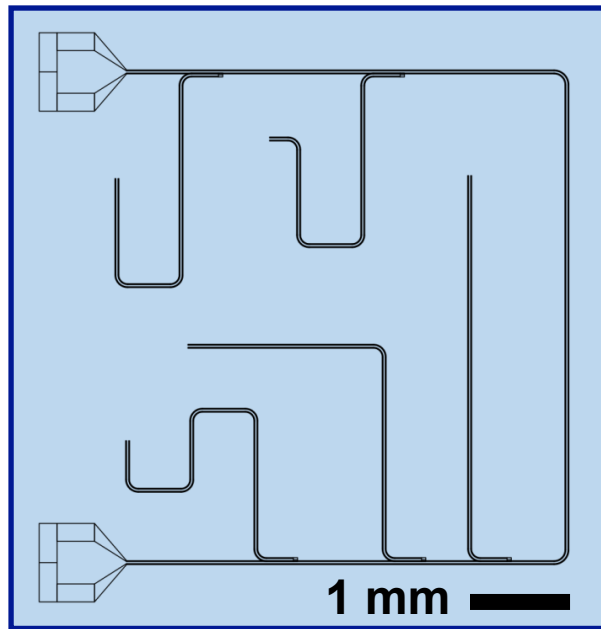
First Generation Test Device



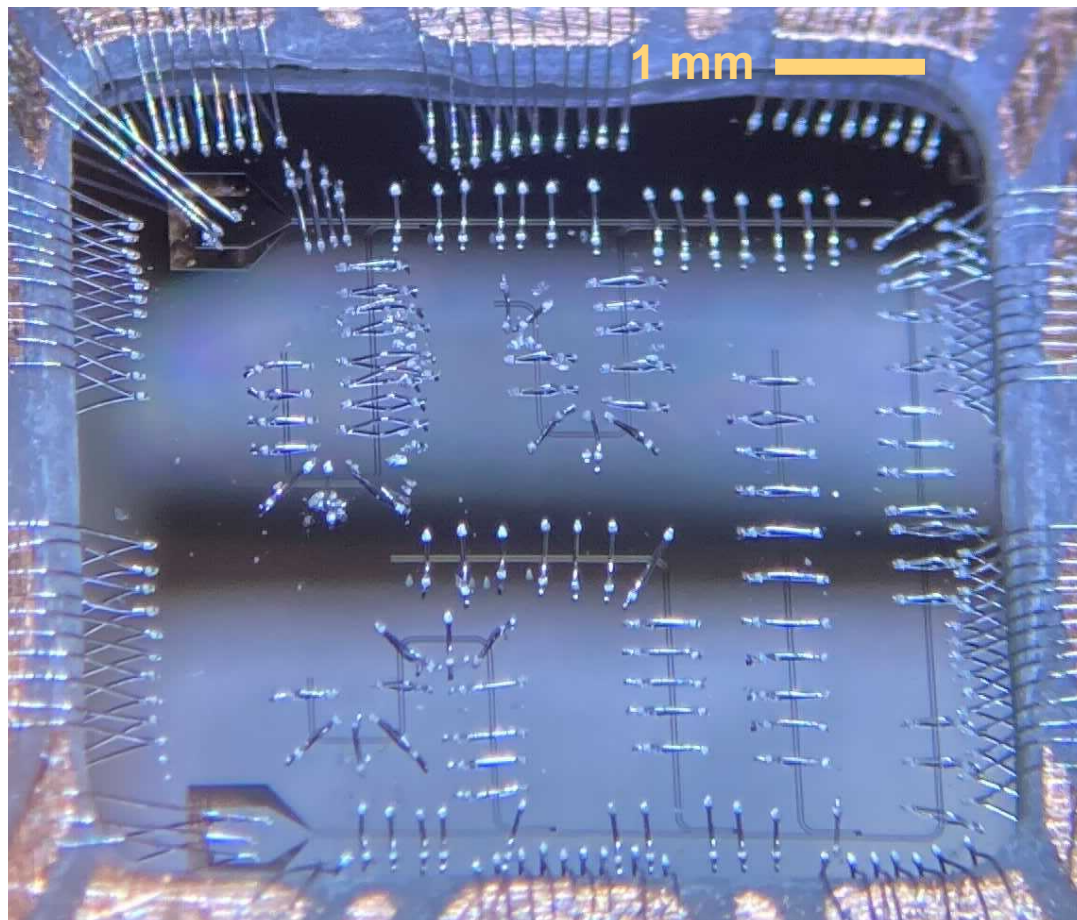
- Varied number of bends

Disorder Mitigation

First Generation Test Device

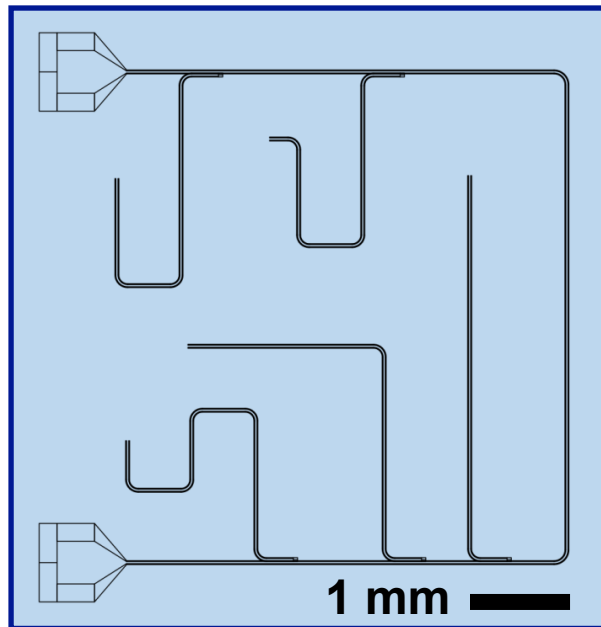


- Varied number of bends



Disorder Mitigation

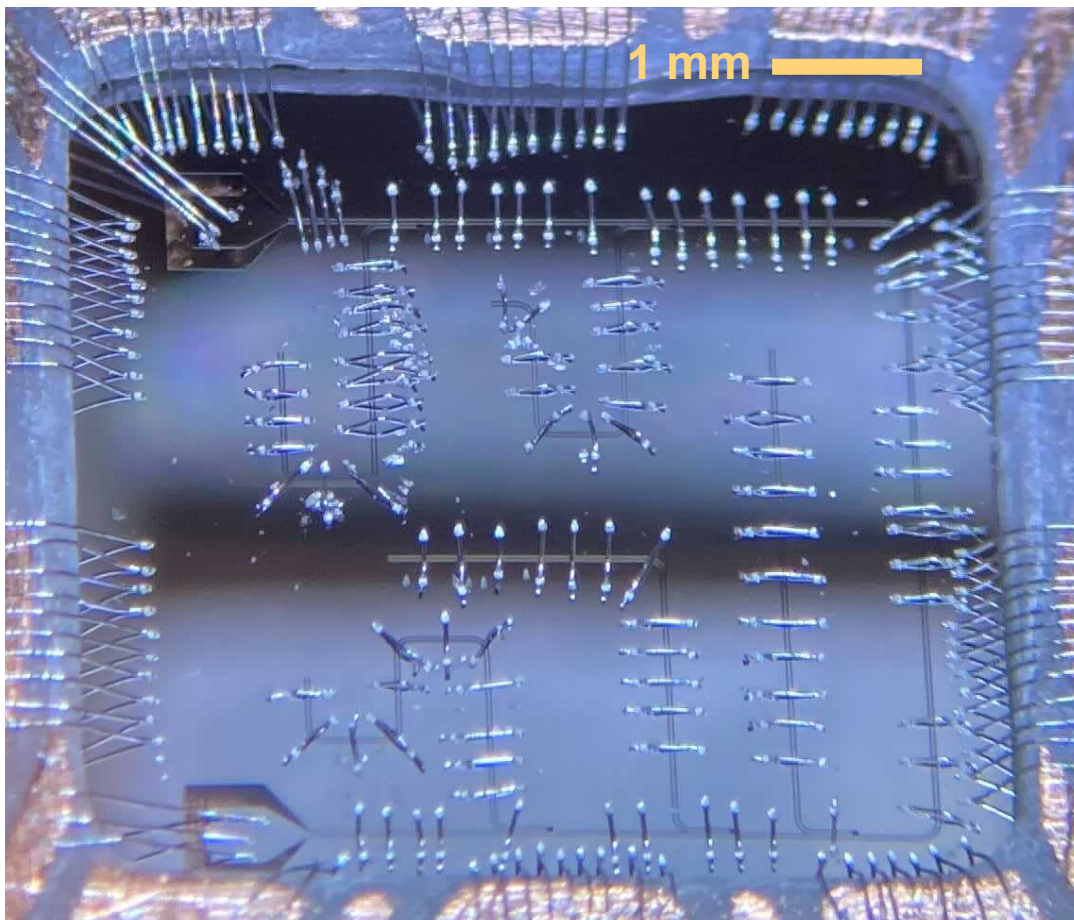
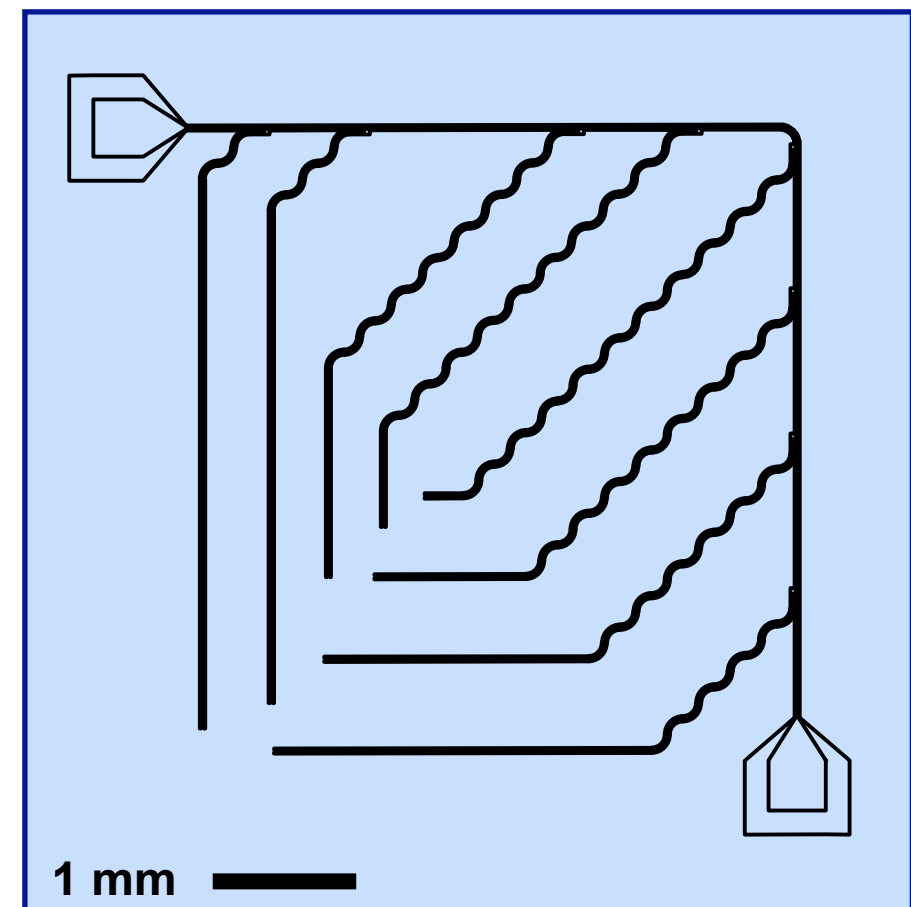
First Generation Test Device



- Varied number of bends

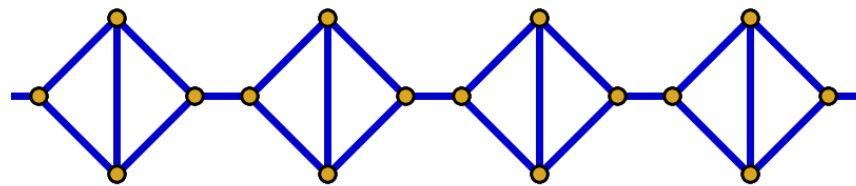
Second Generation Device

- Higher dynamic range (in progress)

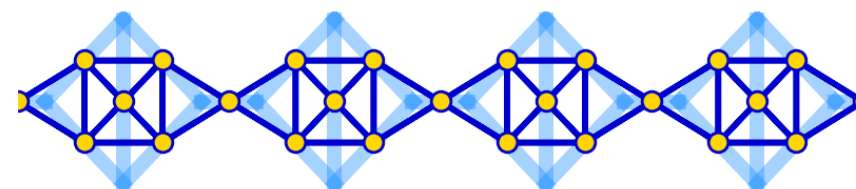


Quasi-1D Lattice Device

Hardware Layout

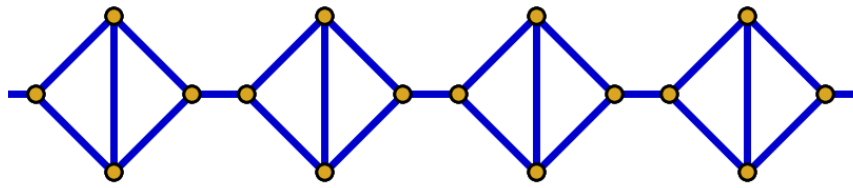


Effective Lattice

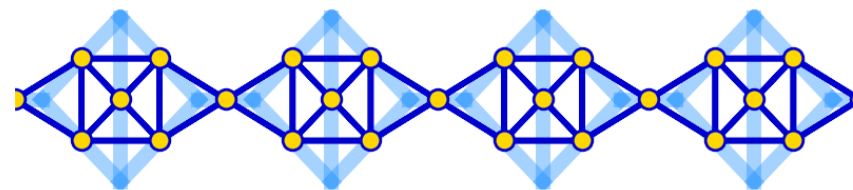


Quasi-1D Lattice Device

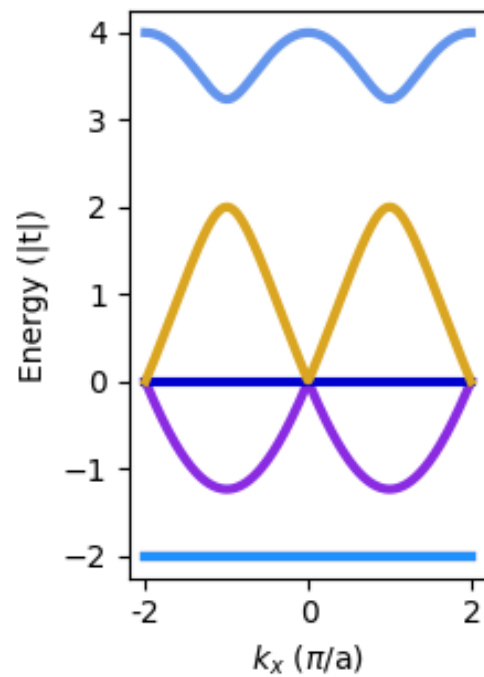
Hardware Layout



Effective Lattice



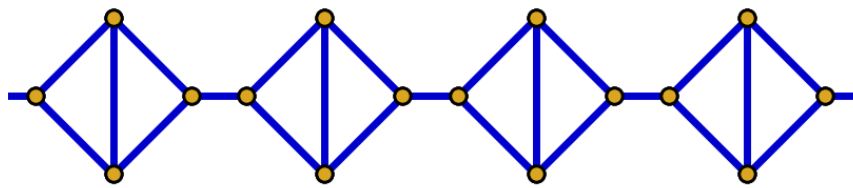
Band Structure



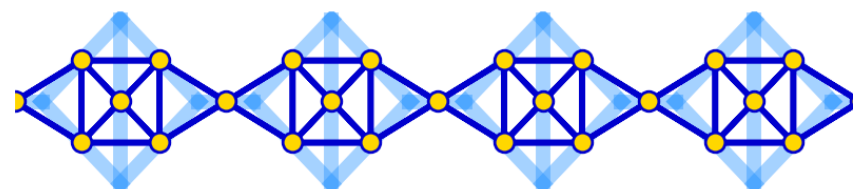
- Flat bands
 - Gapped
 - Ungapped
- Linear bands
- Quadratic bands

Quasi-1D Lattice Device

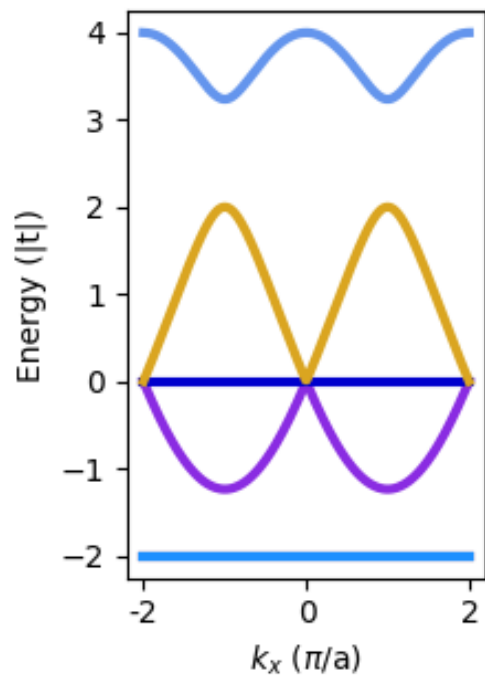
Hardware Layout



Effective Lattice



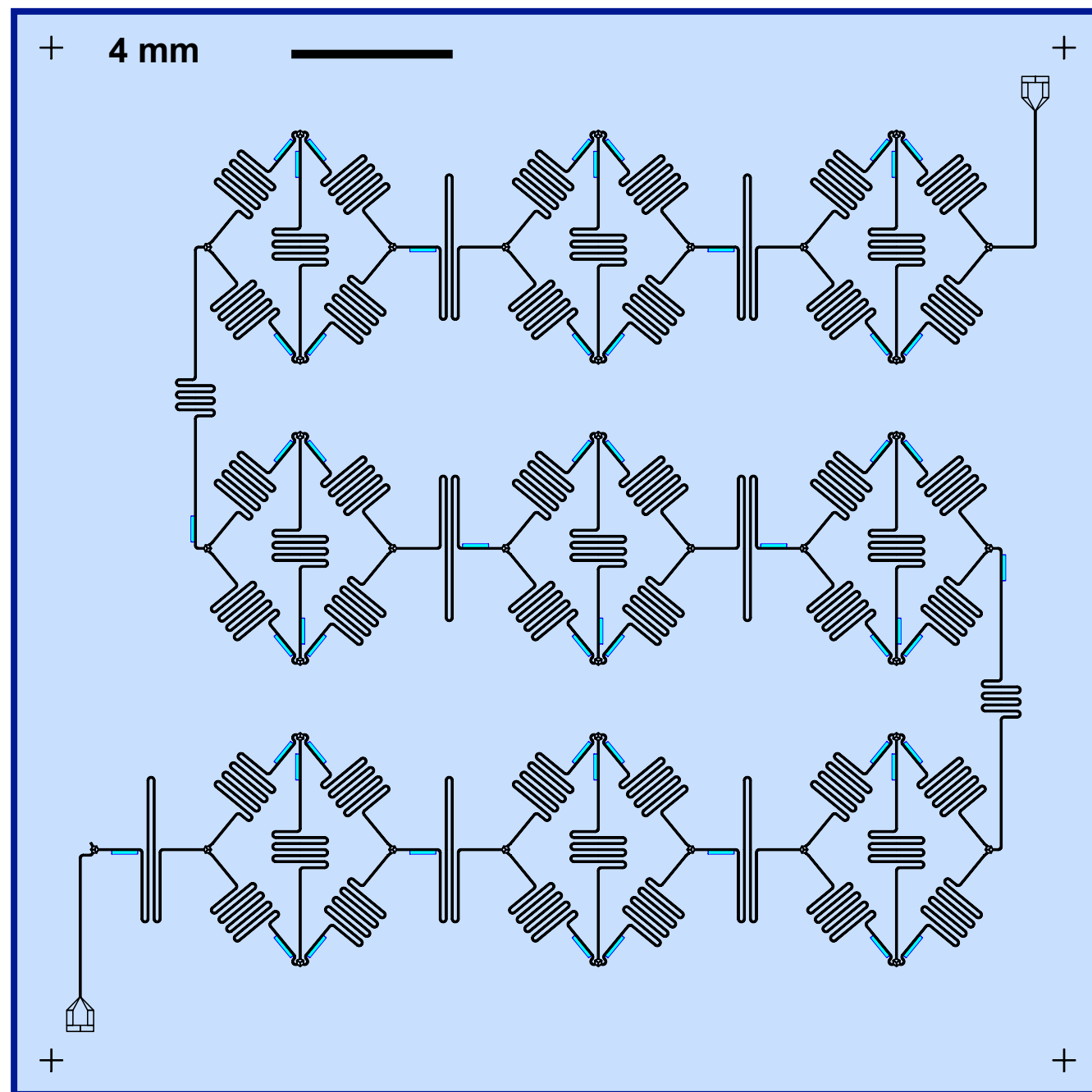
Band Structure



- Flat bands
- Gapped
- Ungapped
- Linear bands
- Quadratic bands

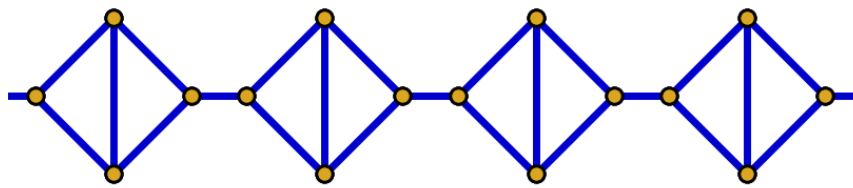
Device Design

(preliminary)

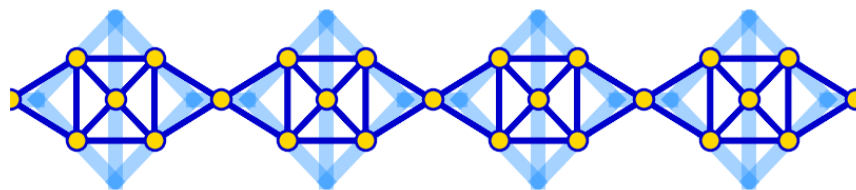


Quasi-1D Lattice Device

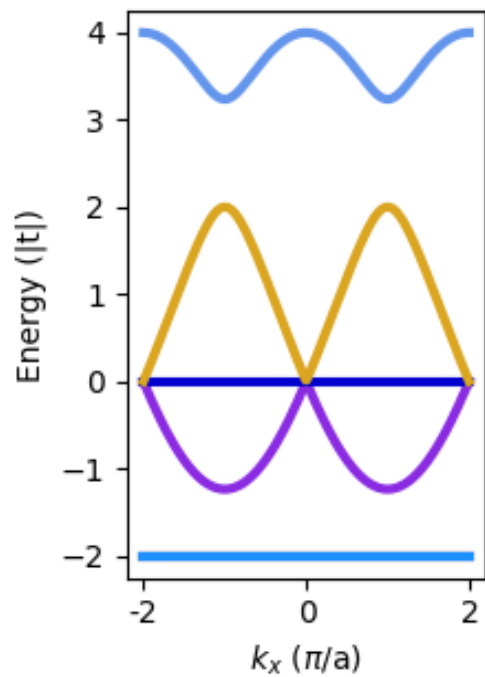
Hardware Layout



Effective Lattice



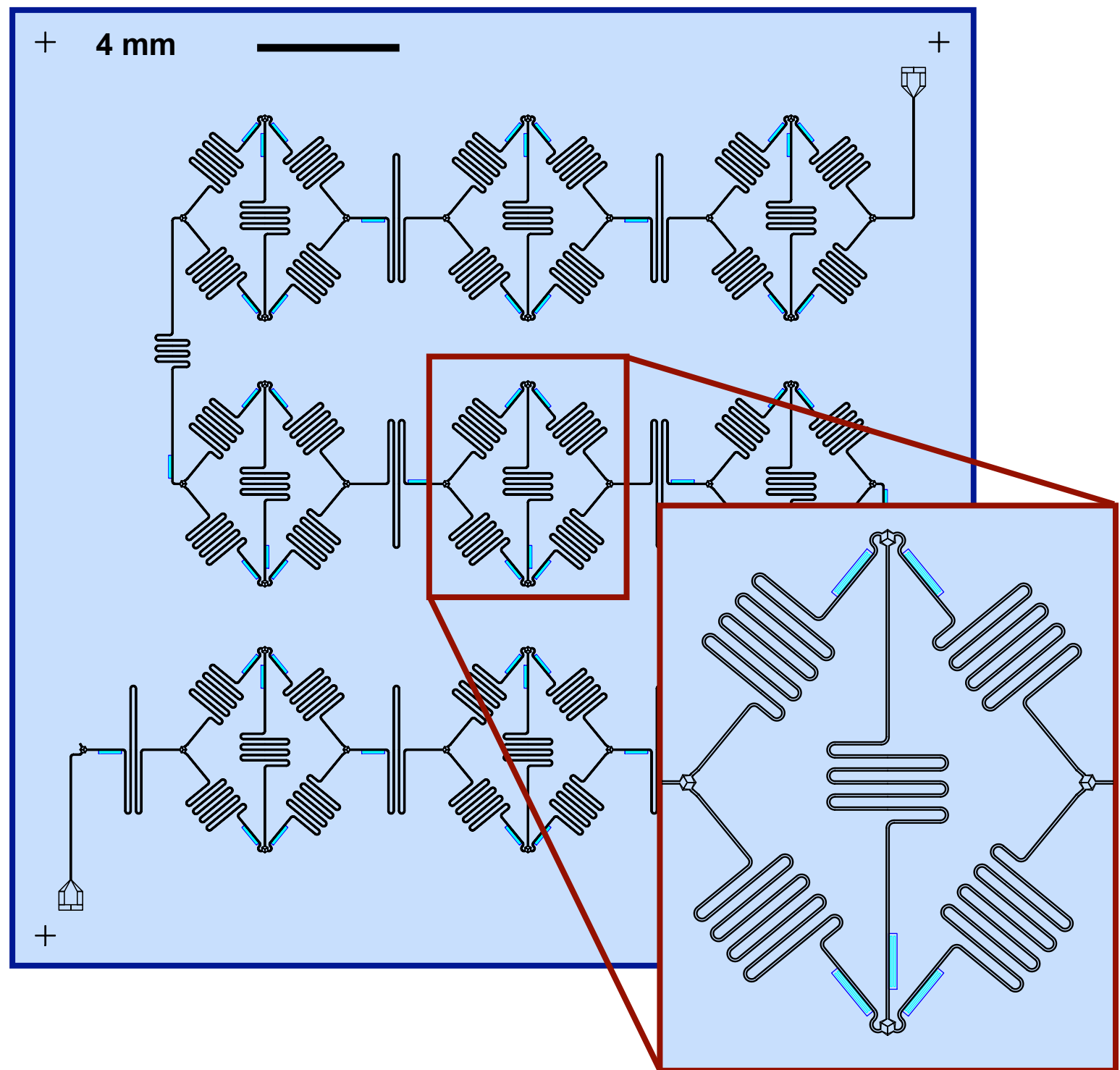
Band Structure



- Flat bands
- Gapped
- Ungapped
- Linear bands
- Quadratic bands

Device Design

(preliminary)



Empty rounded rectangular box at the top of the page.

Large empty rounded rectangular box occupying the majority of the page.

Empty rounded rectangular box at the top of the page.

Large empty rounded rectangular box occupying the majority of the page.

Line Graphs and Quantum Error Correction

Thm: (Chapman and Flammia)

A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

- Spin-model energies found from half-filling of magnetic models on the root graph.
- Gaps in and between these spectra dictate robustness of the code.

Numerical Phenomenology



Error suppression is limited by energy differences between orientations, not single-particle gaps

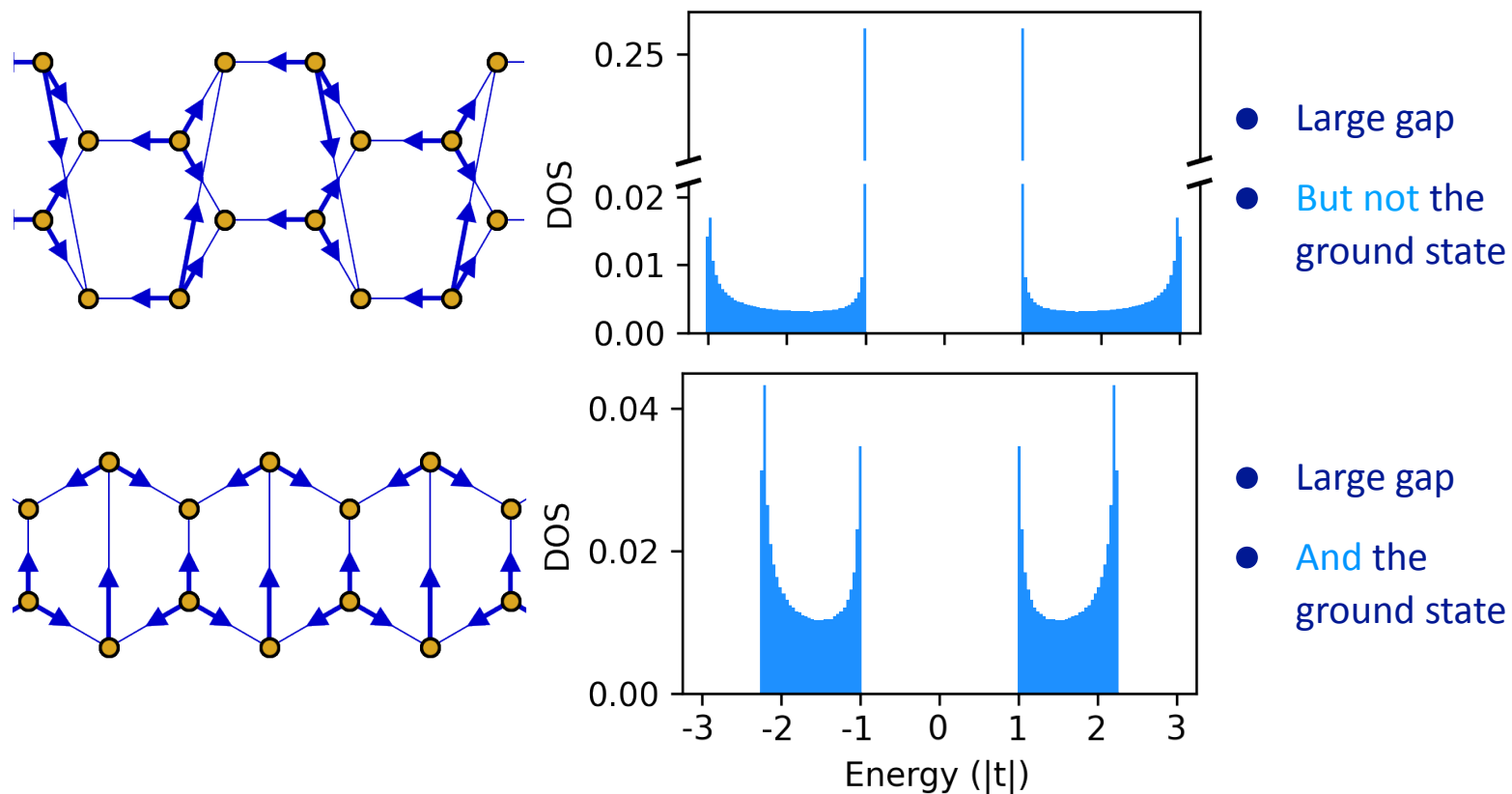
Line Graphs and Quantum Error Correction

Thm: (Chapman and Flammia)

A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

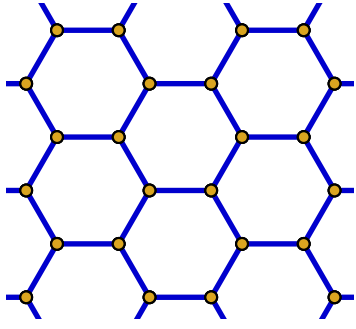
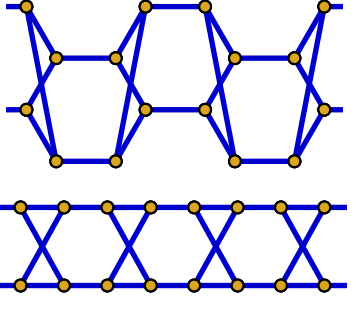
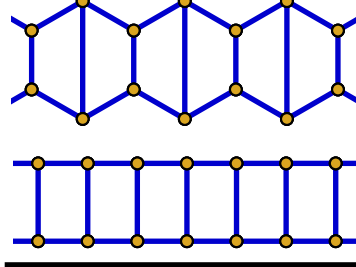
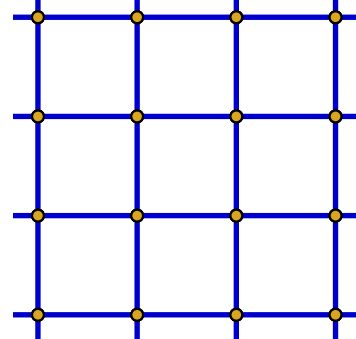
- Spin-model energies found from half-filling of magnetic models on the root graph.
- Gaps in and between these spectra dictate robustness of the code.

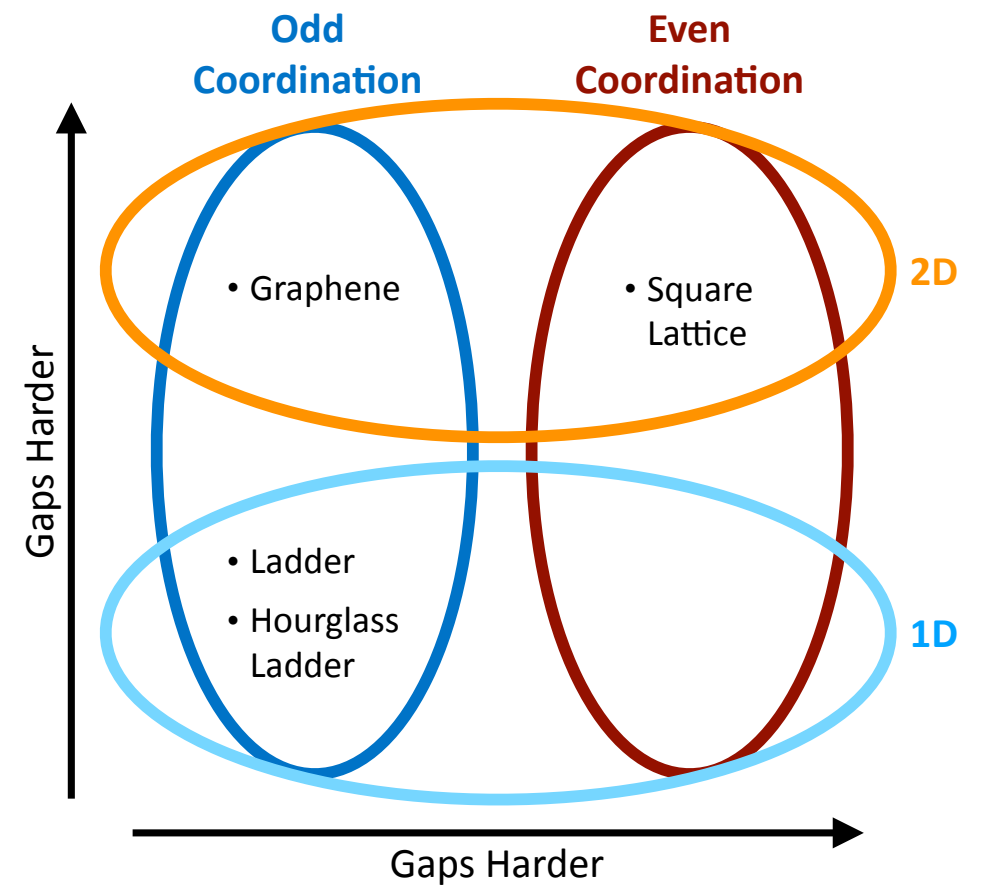
Numerical Phenomenology



→ Error suppression is limited by energy differences between orientations, not single-particle gaps

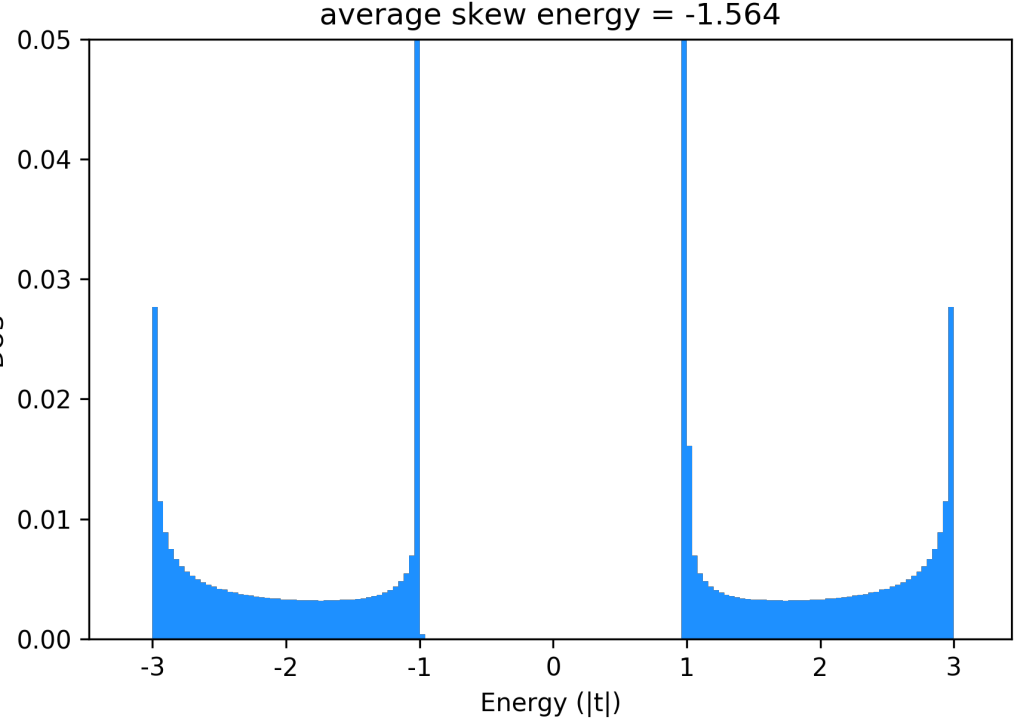
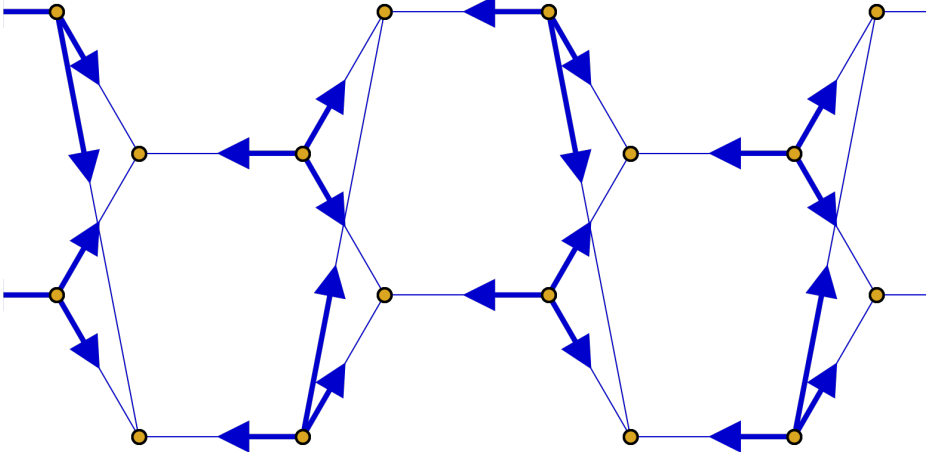
Lattice Gap Examples

	Orientation	Known	Odd	Fluxes	Single-particle gap	Distinct
a 	Elementary		No	0	No	No
	Ground State	Yes	No	0	No	
b 	Elementary		No	0	Yes, Large	Yes
	Ground State	No	—	—	—	
c 	Elementary		No	0	No	Yes
	Ground State	Yes	Yes	π	Yes, Large	
d 	Elementary		No	0	No	Yes
	Ground State	Yes	No	π	No	



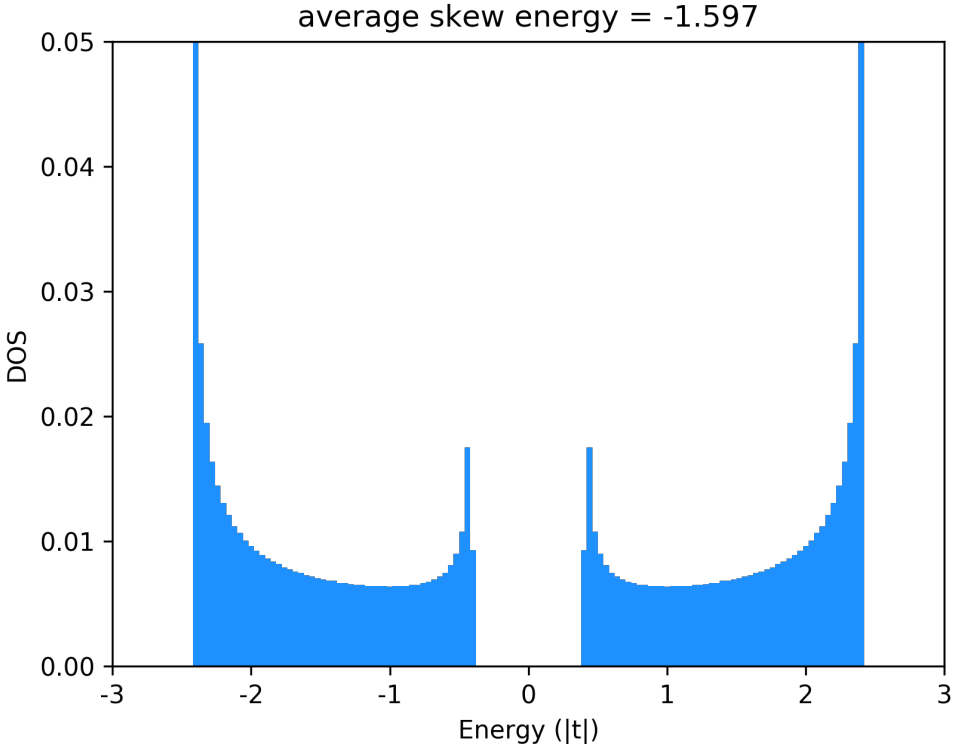
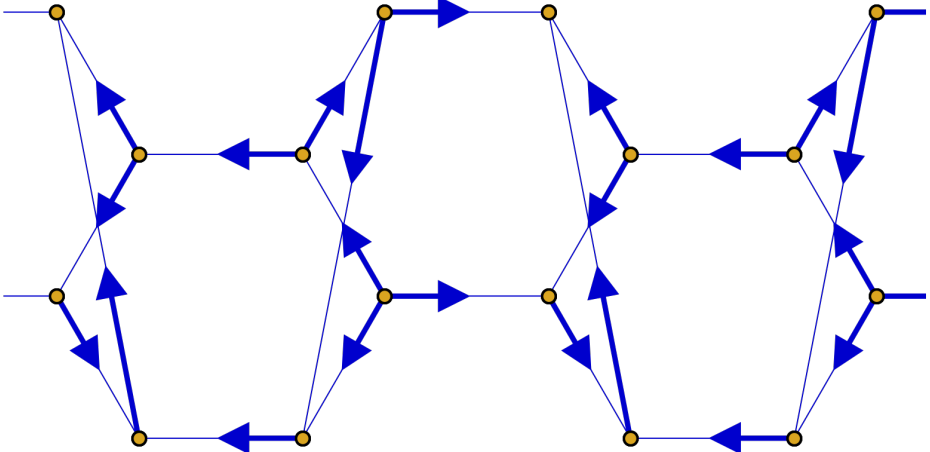
Free-Fermion Solutions

2,0 nanotube : elementary, max SP gap



- Large gap
- But not the ground state

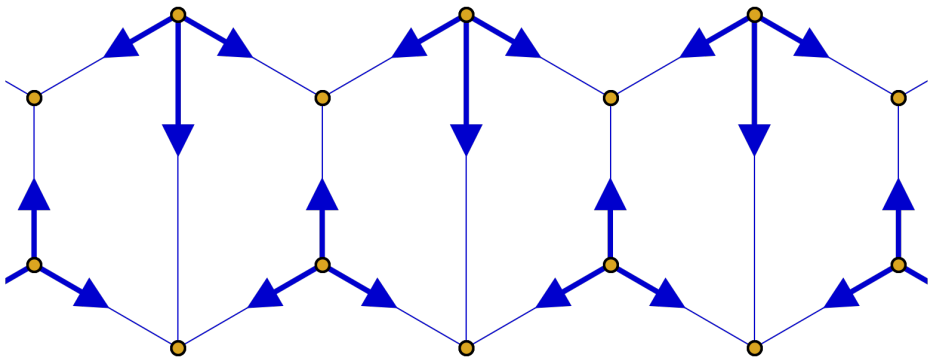
2,0 nanotube : min skew E



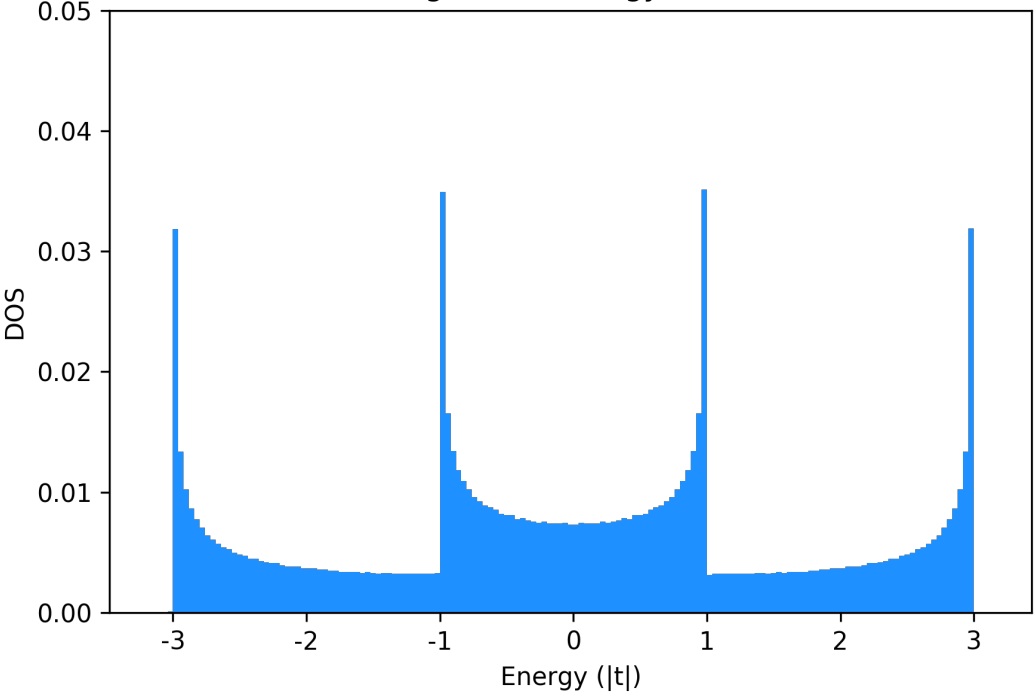
- Modest gap
- Ground state

Free-Fermion Solutions

1,1 nanotube : elementary orientation

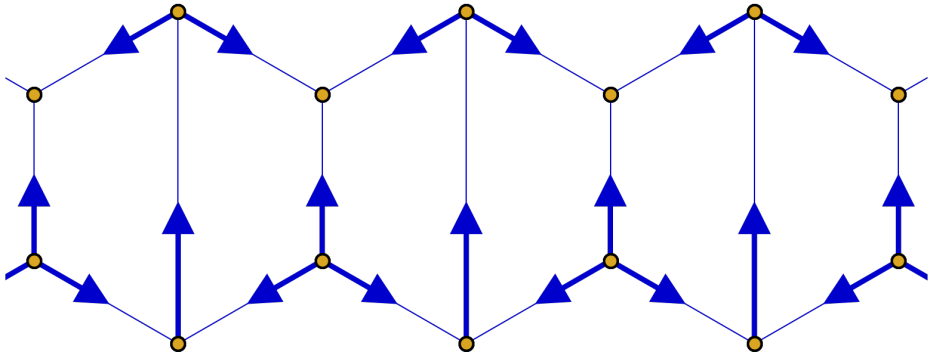


average skew energy = -1.436

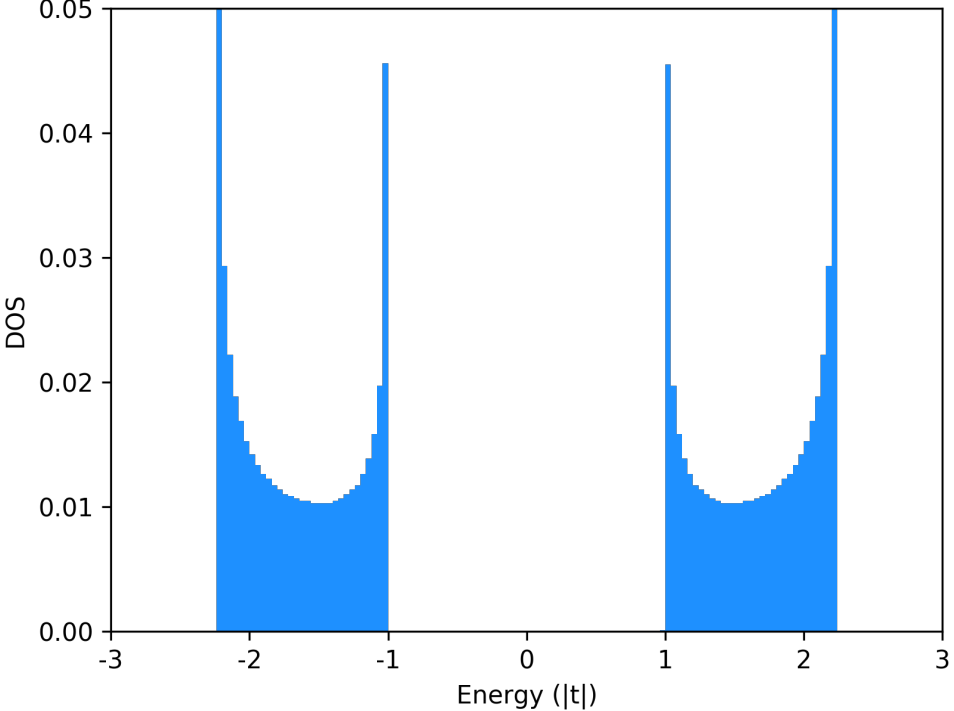


- Non-magnetic orientation
- No gap

1,1 nanotube : max SP gap, min skew E



average skew energy = -1.678

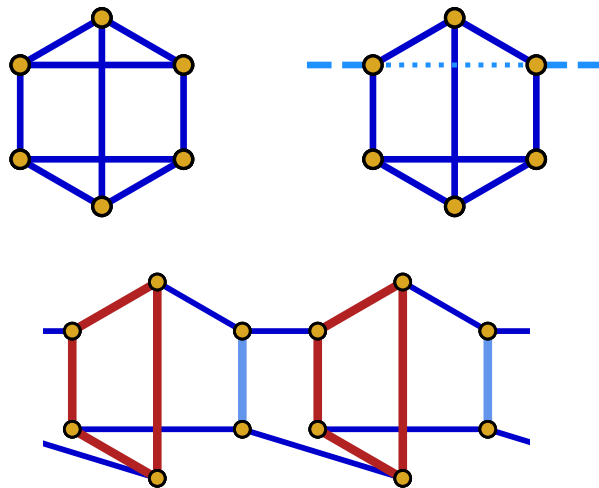


- Large gap
- Ground State

Mathematical Outlook: Abelian Covers and Error Correction

New Lattice Viewpoint

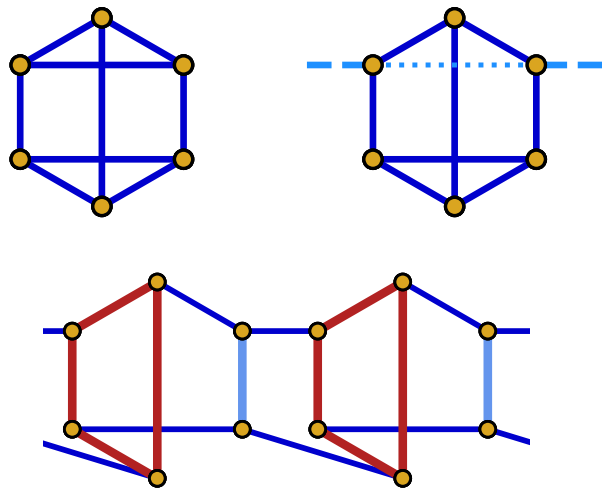
- Use method of Abelian covers to construct examples.
 - “Unwrap” small graph to form lattice



Mathematical Outlook: Abelian Covers and Error Correction

New Lattice Viewpoint

- Use method of Abelian covers to construct examples.
 - “Unwrap” small graph to form lattice

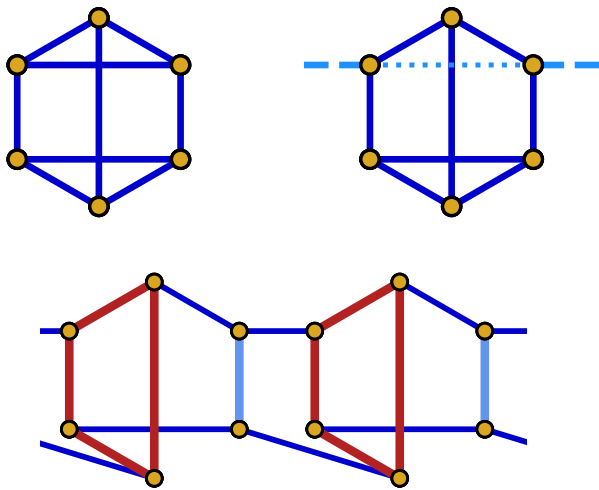


- Initial energies are $k=0$ energies of the lattice
 - Small graphs and their spectra tabulated.
 - “Periodic table” of unit cells to start from.

Mathematical Outlook: Abelian Covers and Error Correction

New Lattice Viewpoint

- Use method of Abelian covers to construct examples.
 - “Unwrap” small graph to form lattice



- Initial energies are $k=0$ energies of the lattice
 - Small graphs and their spectra tabulated.
 - “Periodic table” of unit cells to start from.

Connections to Error-Correcting Codes

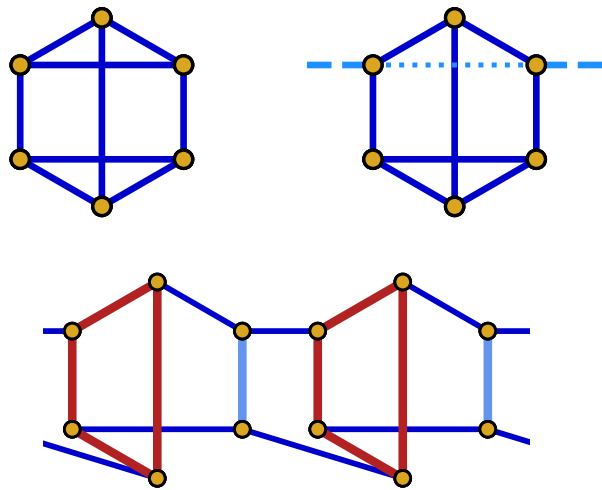
Thm: (Chapman and Flammia)
A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

- Spin-model energies found from half-filling of magnetic models on the root graph.
- Gaps in and between these spectra dictate robustness of the code.
- In progress: using Abelian cover method to categorize large gaps in this sense.

Mathematical Outlook: Abelian Covers and Error Correction

New Lattice Viewpoint

- Use method of Abelian covers to construct examples.
 - “Unwrap” small graph to form lattice

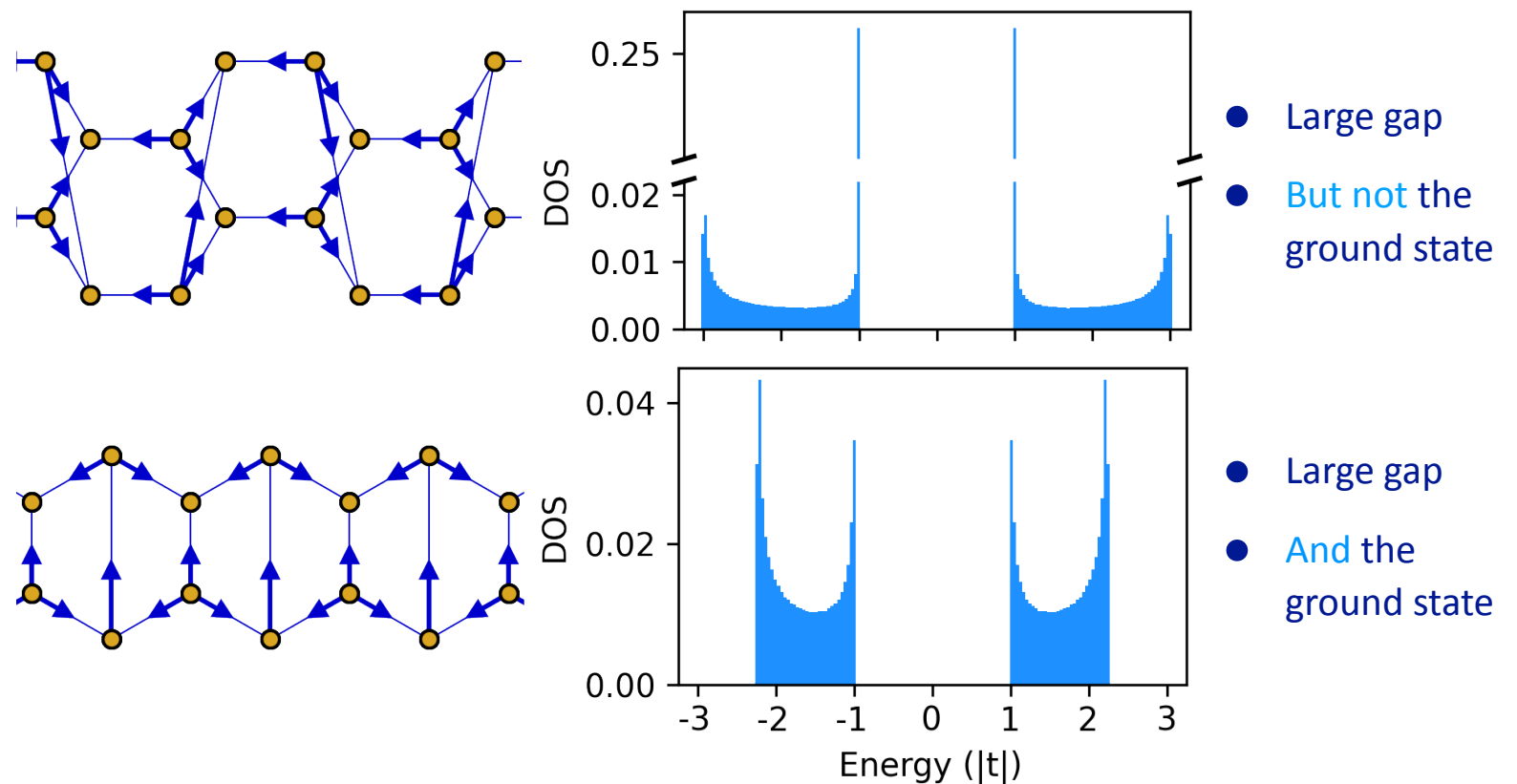


- Initial energies are $k=0$ energies of the lattice
 - Small graphs and their spectra tabulated.
 - “Periodic table” of unit cells to start from.

Connections to Error-Correcting Codes

Thm: (Chapman and Flammia)
 A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

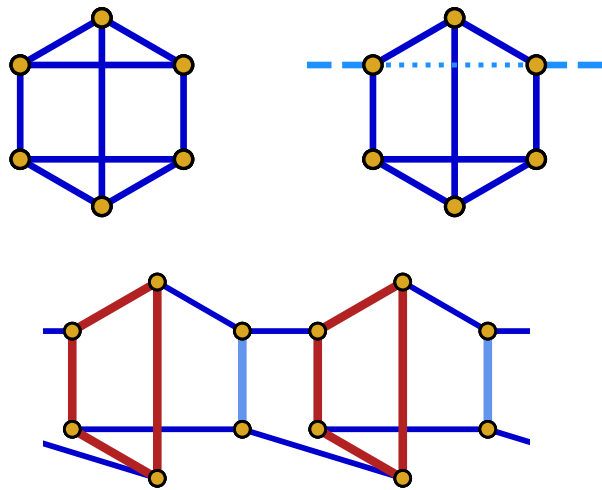
- Spin-model energies found from half-filling of magnetic models on the root graph.
- Gaps in and between these spectra dictate robustness of the code.
- In progress: using Abelian cover method to categorize large gaps in this sense.



Mathematical Outlook: Abelian Covers and Error Correction

New Lattice Viewpoint

- Use method of Abelian covers to construct examples.
 - “Unwrap” small graph to form lattice

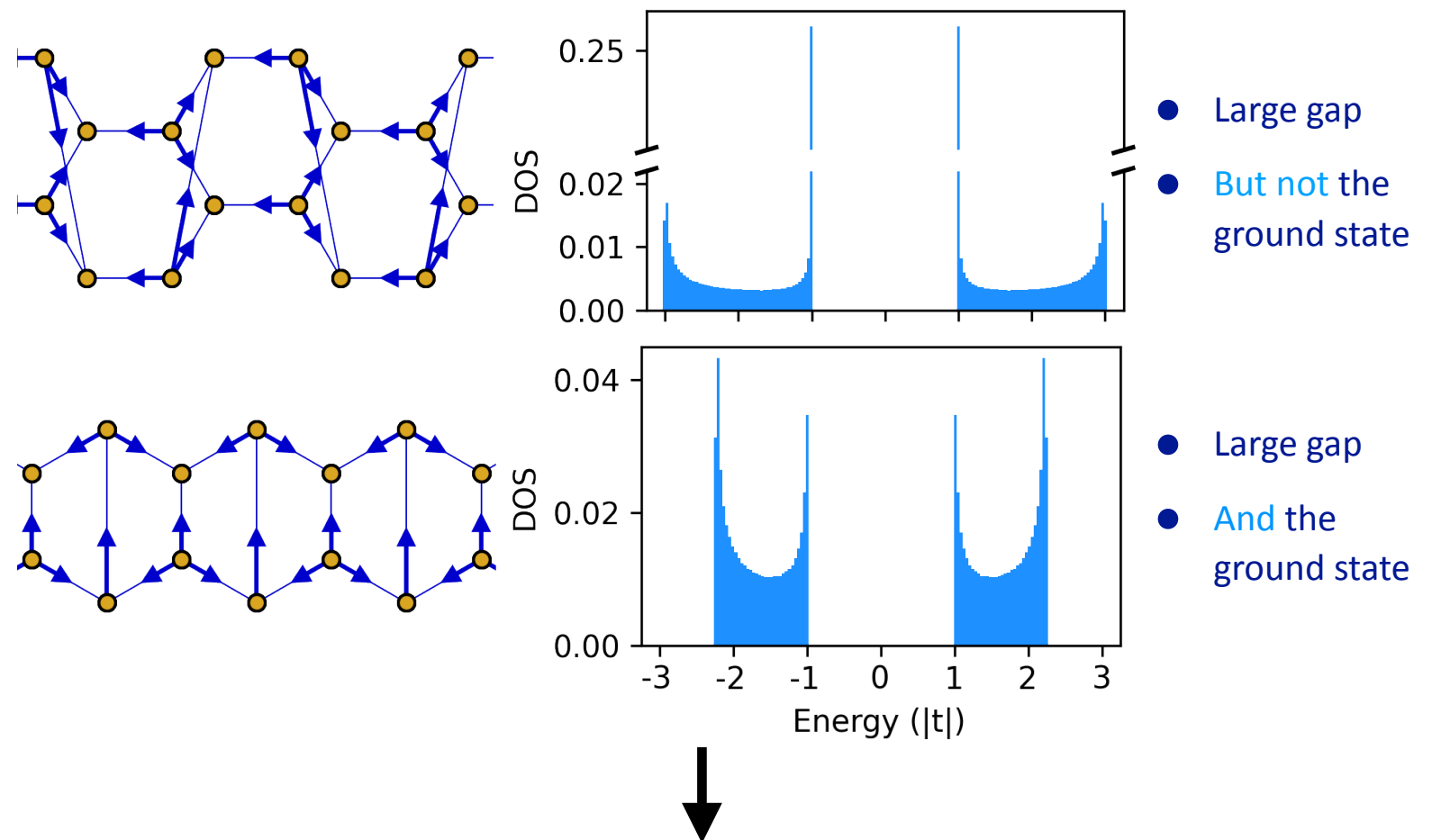


- Initial energies are $k=0$ energies of the lattice
 - Small graphs and their spectra tabulated.
 - “Periodic table” of unit cells to start from.

Connections to Error-Correcting Codes

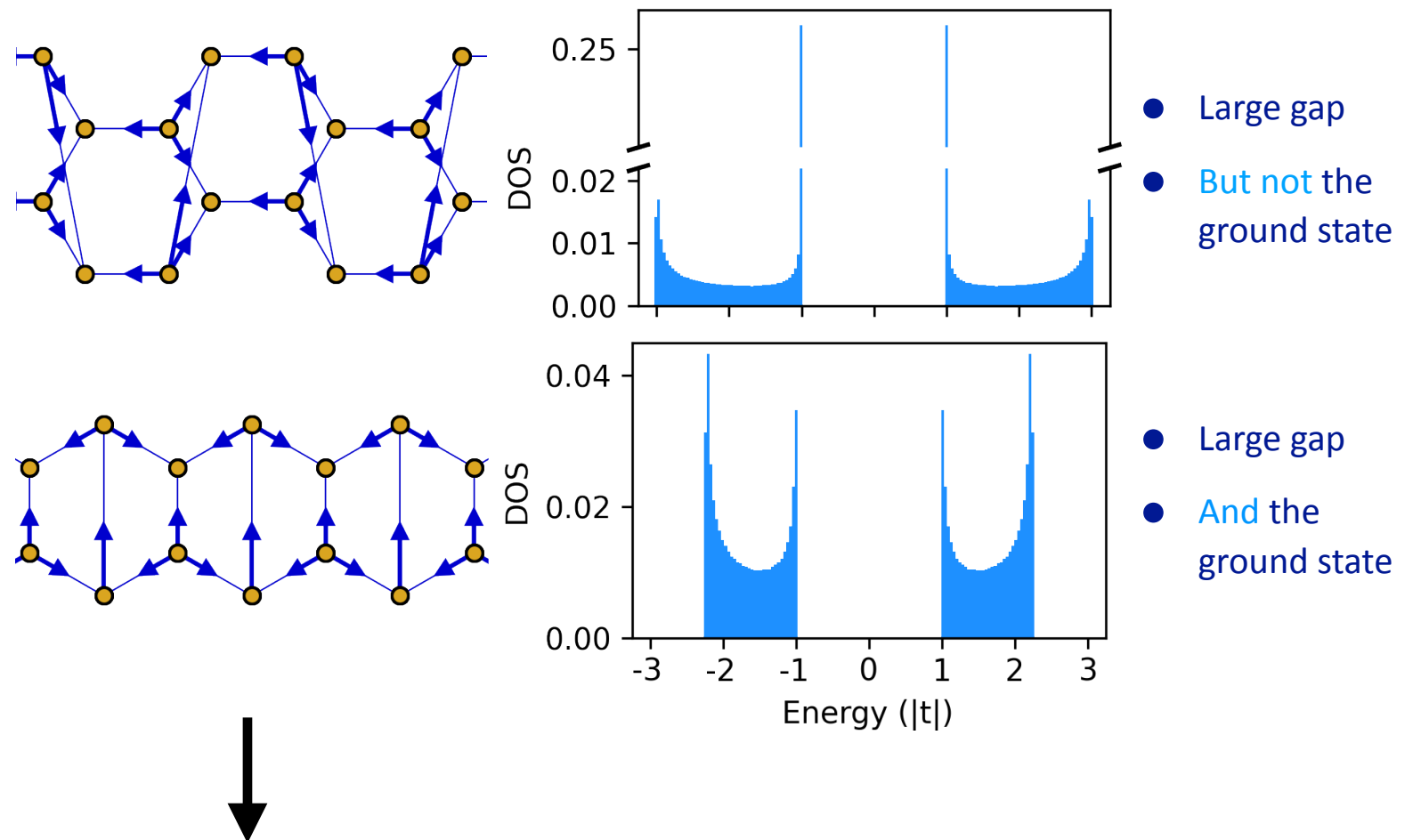
Thm: (Chapman and Flammia)
 A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

- Spin-model energies found from half-filling of magnetic models on the root graph.
- Gaps in and between these spectra dictate robustness of the code.
- In progress: using Abelian cover method to categorize large gaps in this sense.



So far, error suppression is limited by energy differences between orientations, not single-particle gaps

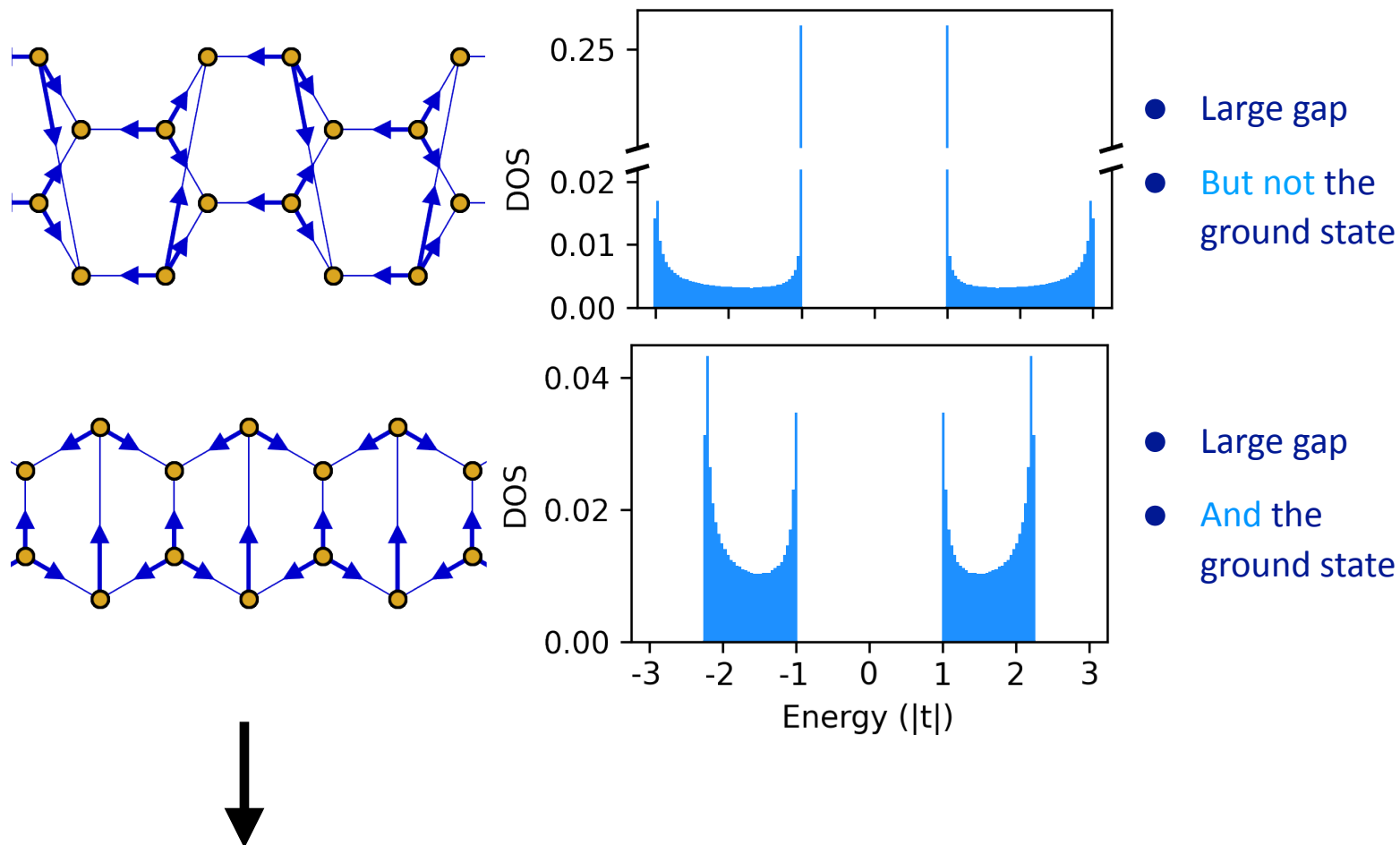
Connections to Error Correction



Connections to Error Correction

Thm: (Chapman and Flammia)

A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

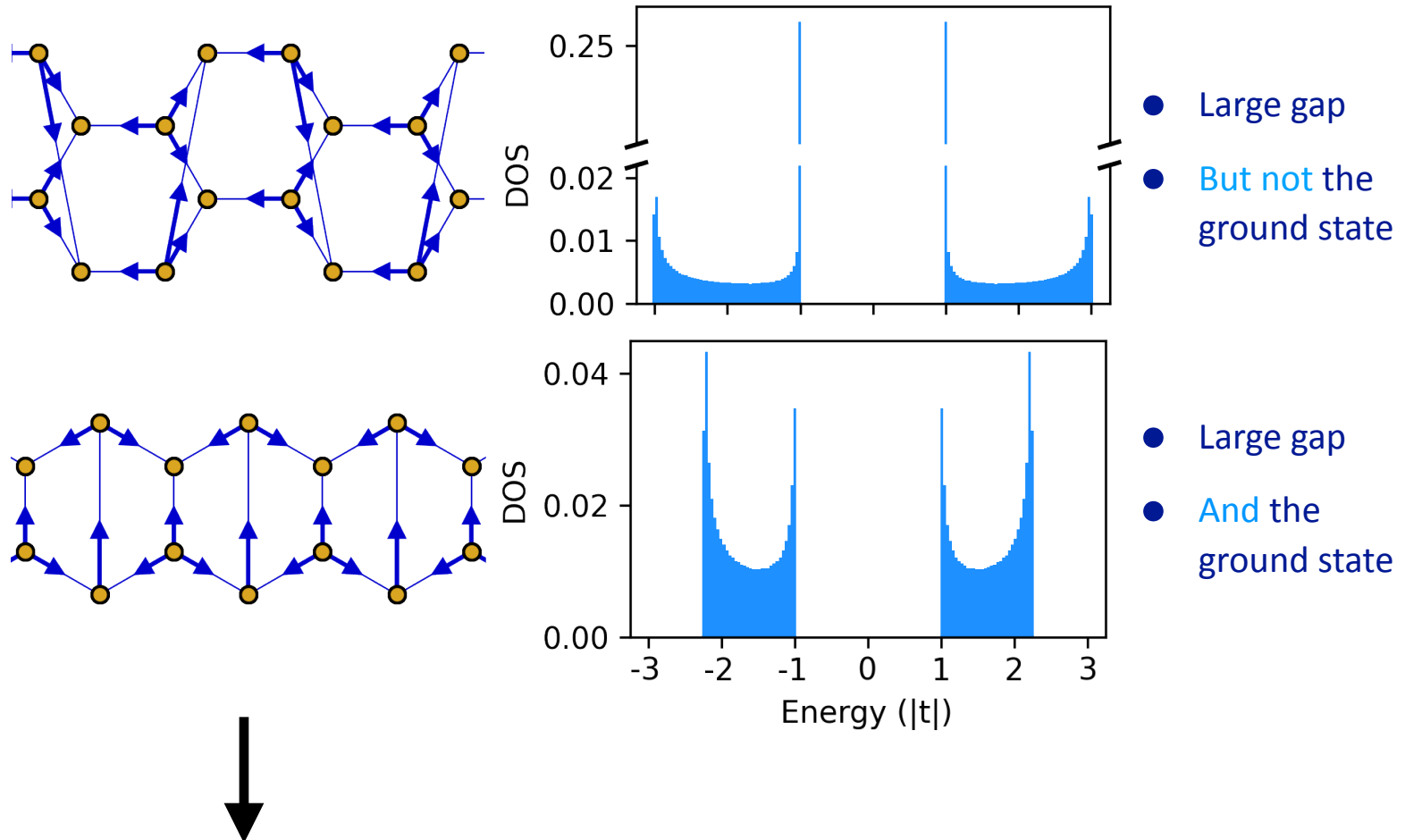


Connections to Error Correction

Thm: (Chapman and Flammia)

A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

- Spin-model energies found from the skew energy of the oriented root graph.

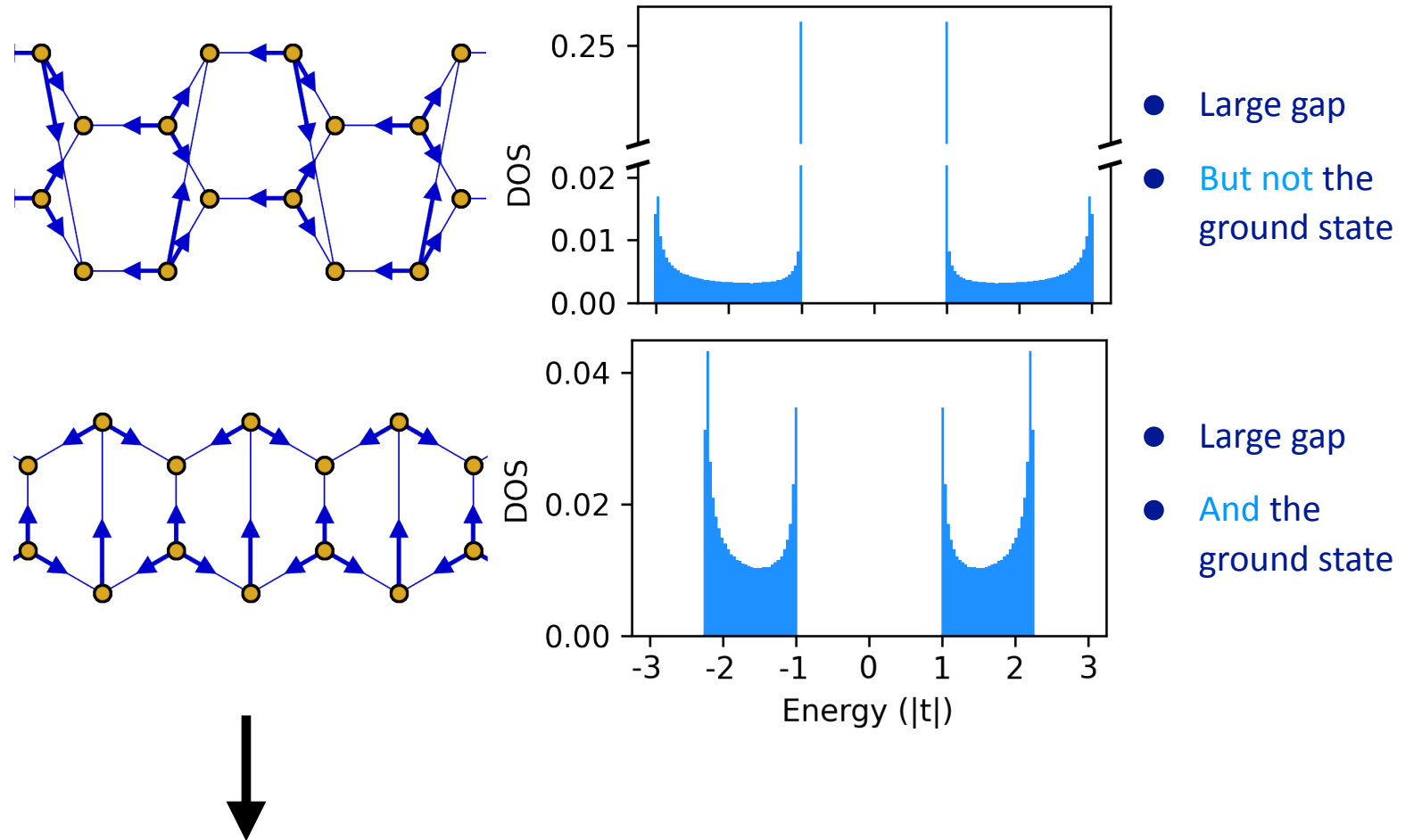


Connections to Error Correction

Thm: (Chapman and Flammia)

A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

- Spin-model energies found from the skew energy of the oriented root graph.
- Gaps in and between these spectra dictate robustness of the code.

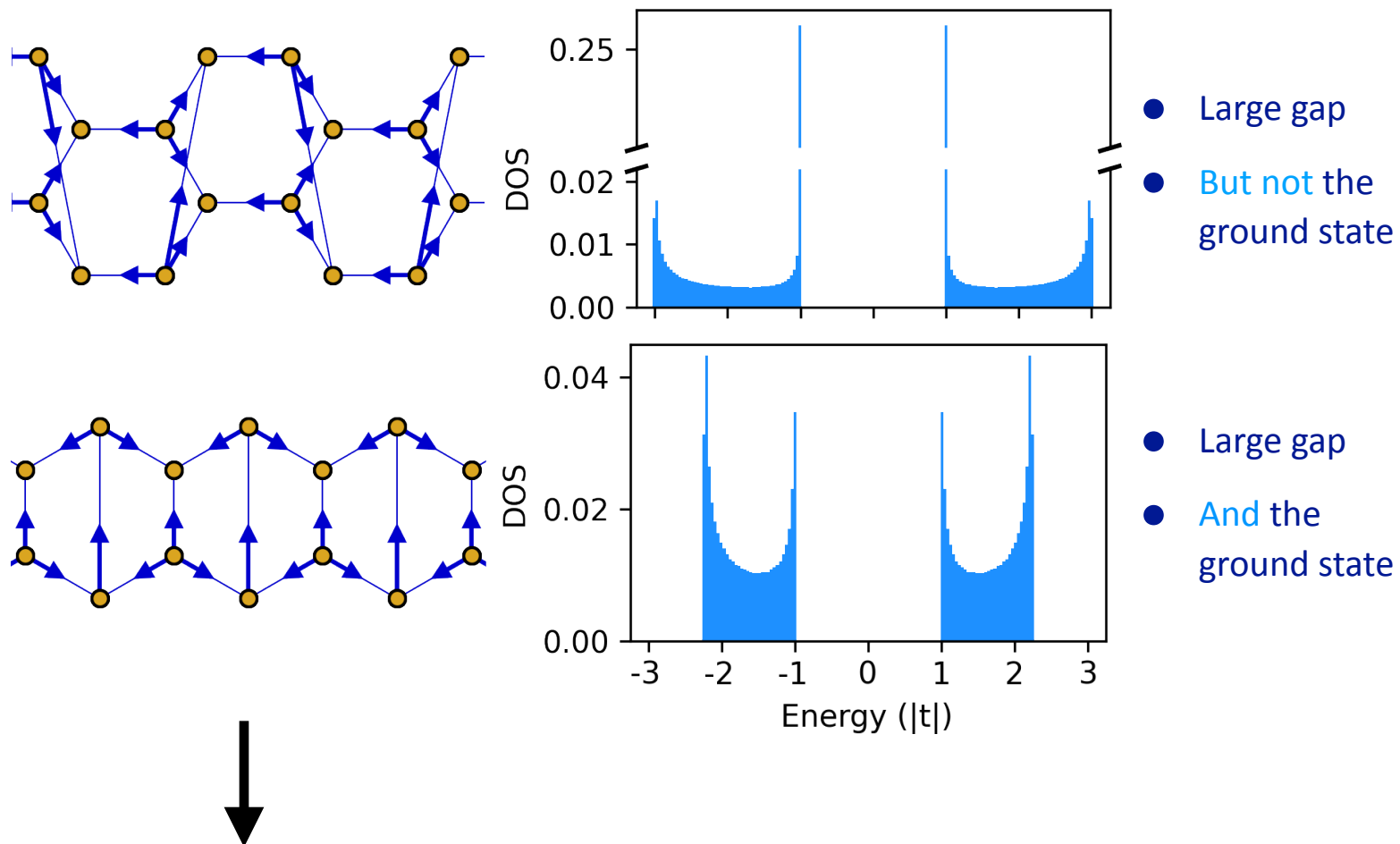


Connections to Error Correction

Thm: (Chapman and Flammia)

A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

- Spin-model energies found from the skew energy of the oriented root graph.
- Gaps in and between these spectra dictate robustness of the code.
- In progress: using Abelian cover search to categorize large gaps in this sense.

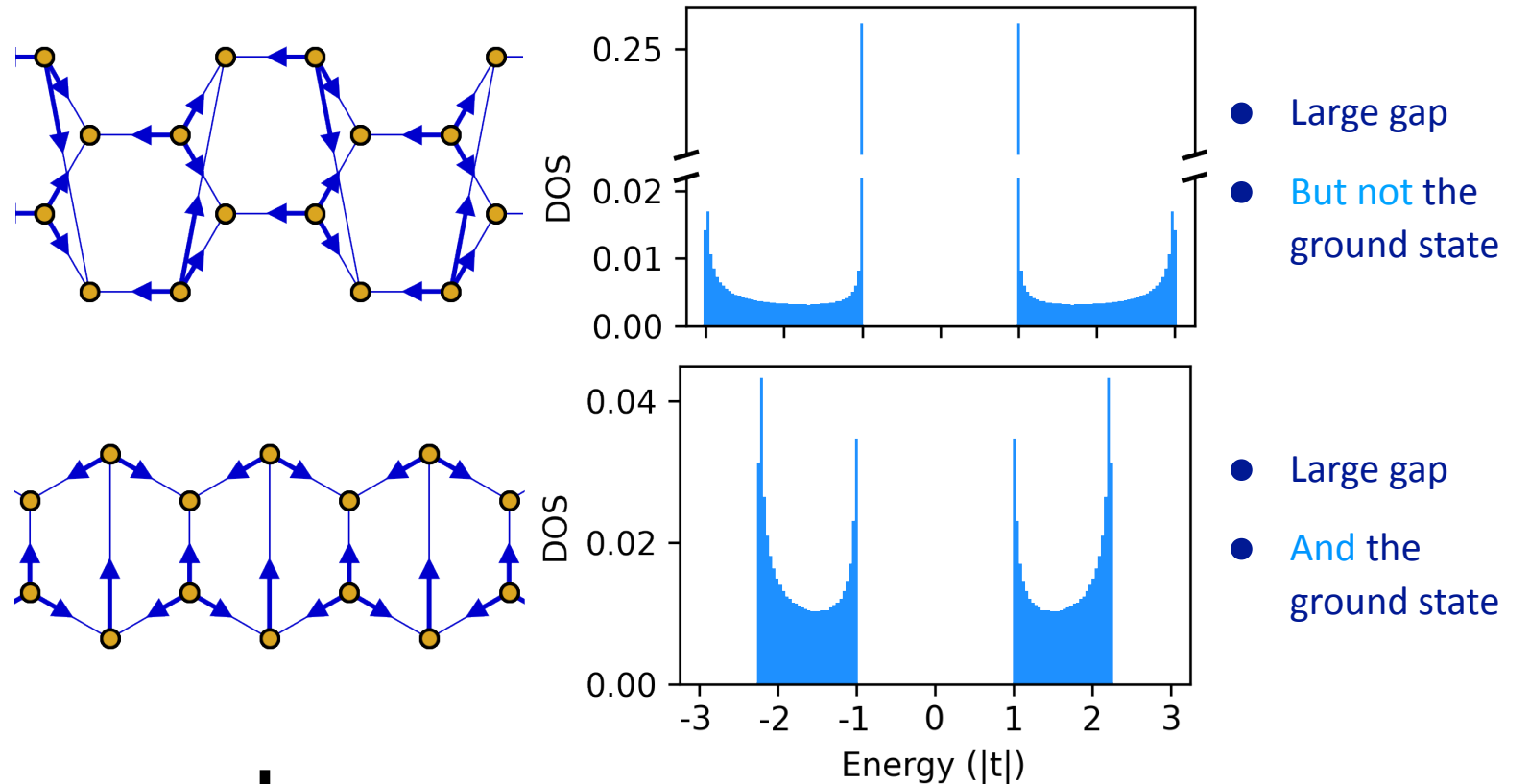


Connections to Error Correction

Thm: (Chapman and Flammia)

A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

- Spin-model energies found from the skew energy of the oriented root graph.
- Gaps in and between these spectra dictate robustness of the code.
- In progress: using Abelian cover search to categorize large gaps in this sense.



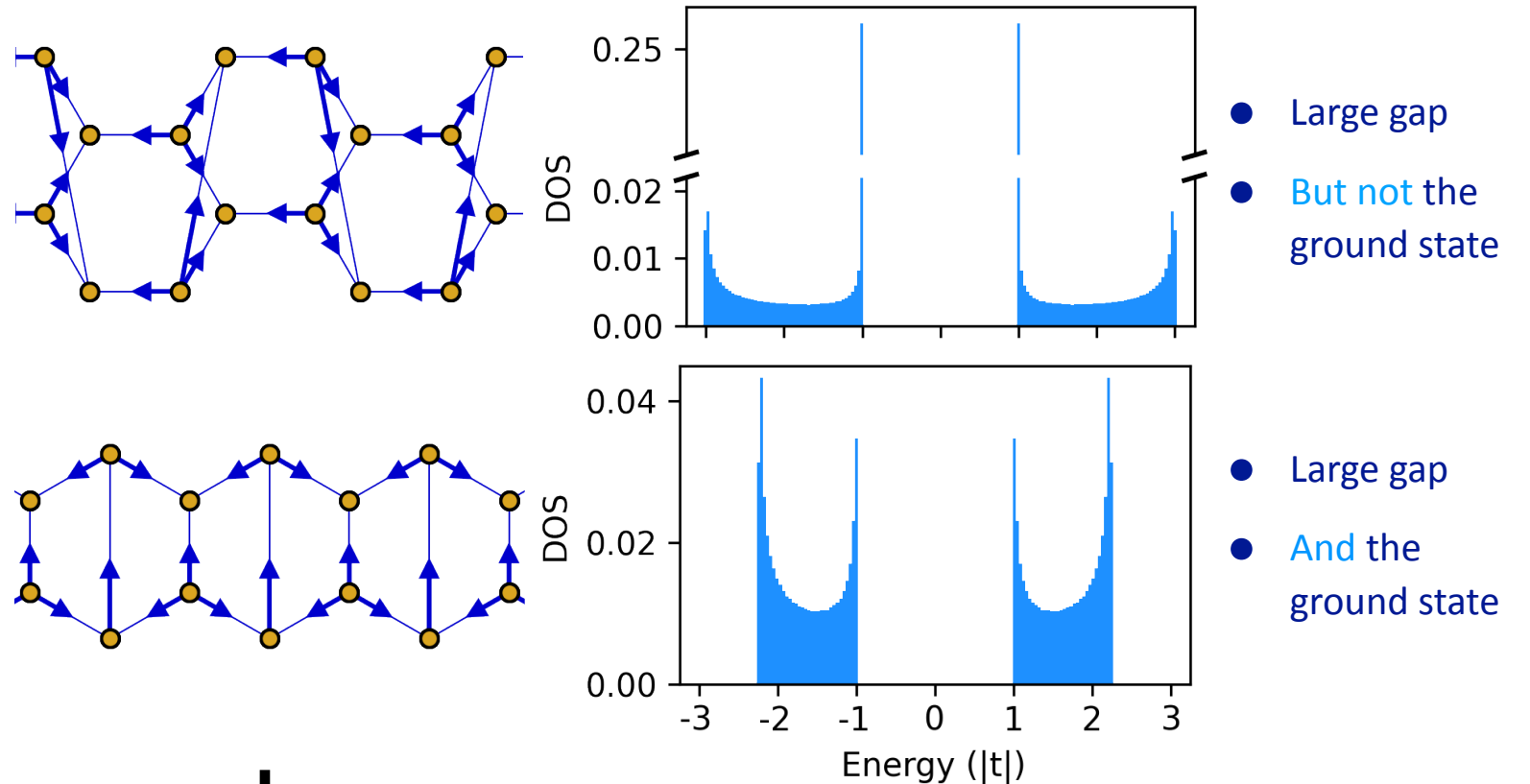
- Abelian covers of small regular graphs yield examples with a large gap within the ground state orientation

Connections to Error Correction

Thm: (Chapman and Flammia)

A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

- Spin-model energies found from the skew energy of the oriented root graph.
- Gaps in and between these spectra dictate robustness of the code.
- In progress: using Abelian cover search to categorize large gaps in this sense.



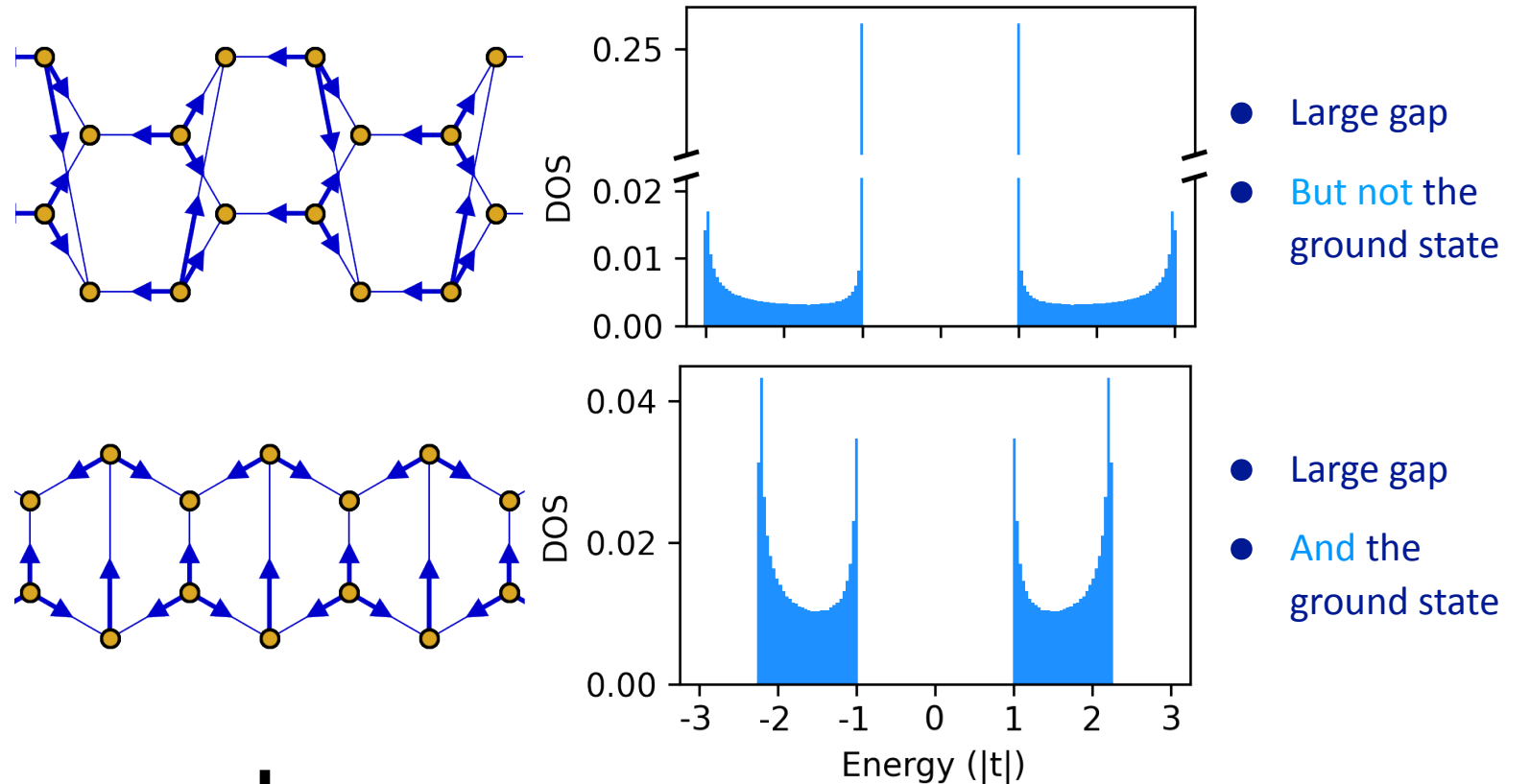
- Abelian covers of small regular graphs yield examples with a large gap within the ground state orientation
- But, skew energy gaps between orientations remain small

Connections to Error Correction

Thm: (Chapman and Flammia)

A spin model can be solved exactly by mapping to free fermions if and only if the anticommutation relations of its terms have the structure of a line graph.

- Spin-model energies found from the skew energy of the oriented root graph.
- Gaps in and between these spectra dictate robustness of the code.
- In progress: using Abelian cover search to categorize large gaps in this sense.



- Abelian covers of small regular graphs yield examples with a large gap within the ground state orientation
- But, skew energy gaps between orientations remain small
- Error suppression limited by skew energy, so far

The Triangle Models

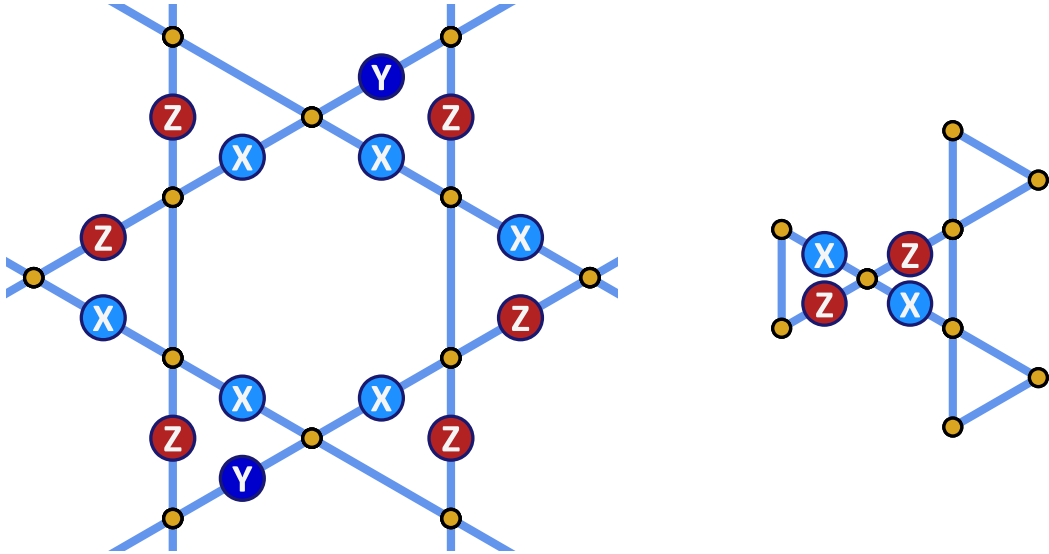
Three Combined Models

- Free-fermion model : Kitaev Honeycomb
- Stabilizer code : Wen Plaquette
- Paramagnet to couple the two

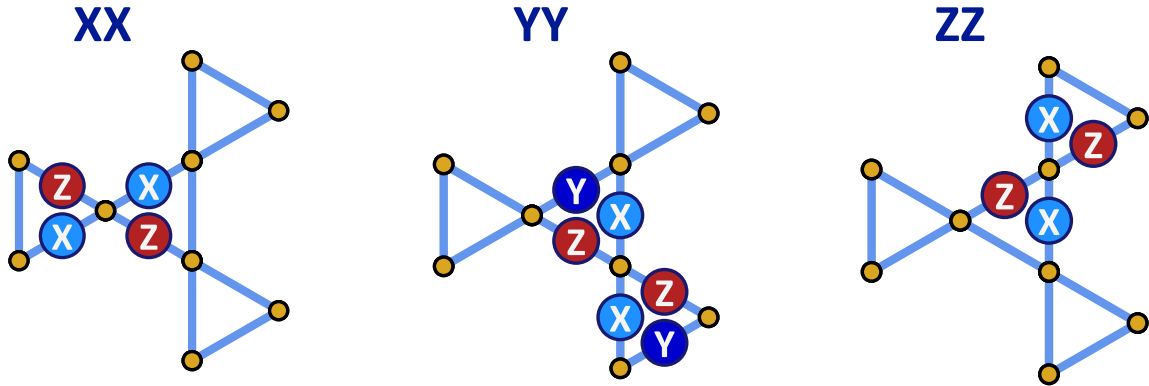
Effective Qubits

	X	Y	Z
Q_P			
Q_S			
Q_F			

Wen Plaquette Model



Kitaev Honeycomb Model



Exact logicals **without** fermion participation

Empty rounded rectangular box at the top of the page.

Large empty rounded rectangular box occupying the majority of the page.

Photon-Mediated Interactions

Photonic Crystal + qubits

- Effective swap interaction
- All modes in parallel

$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

Douglas *et al.* Nat. Photon. (2015)

Calajó *et al.* PRA (2016)

Liu *et al.* Nature Physics (2016)

Sundaresan *et al.* PRX (2019)

Ferreira *et al.* arXiv 2001.0324 (2020)

Photon-Mediated Interactions

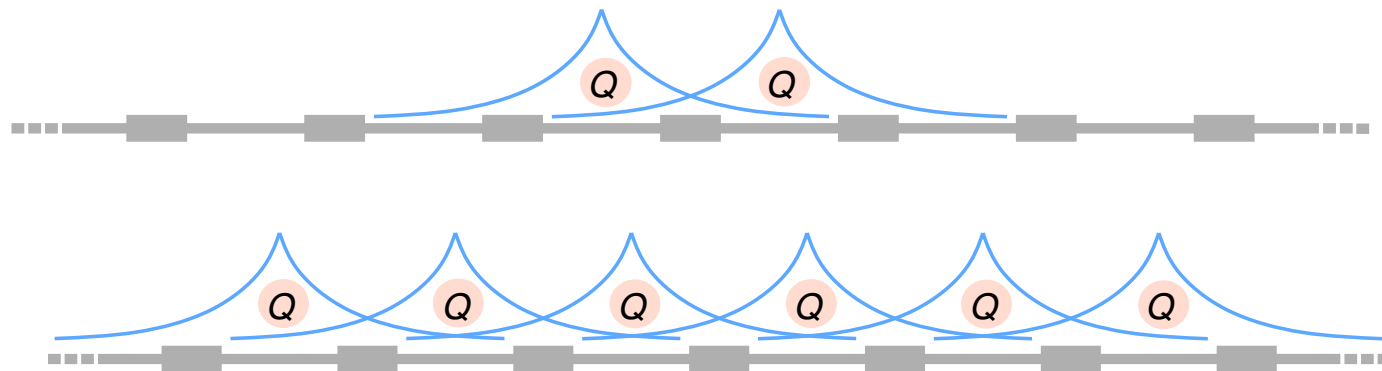
Photonic Crystal + qubits

- Effective swap interaction
- All modes in parallel

$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

1D-Photonic Crystal

- Exponentially localized bound state



Douglas *et al.* Nat. Photon. (2015)
Calajó *et al.* PRA (2016)
Liu *et al.* Nature Physics (2016)
Sundaresan *et al.* PRX (2019)
Ferreira *et al.* arXiv 2001.0324 (2020)

Photon-Mediated Interactions

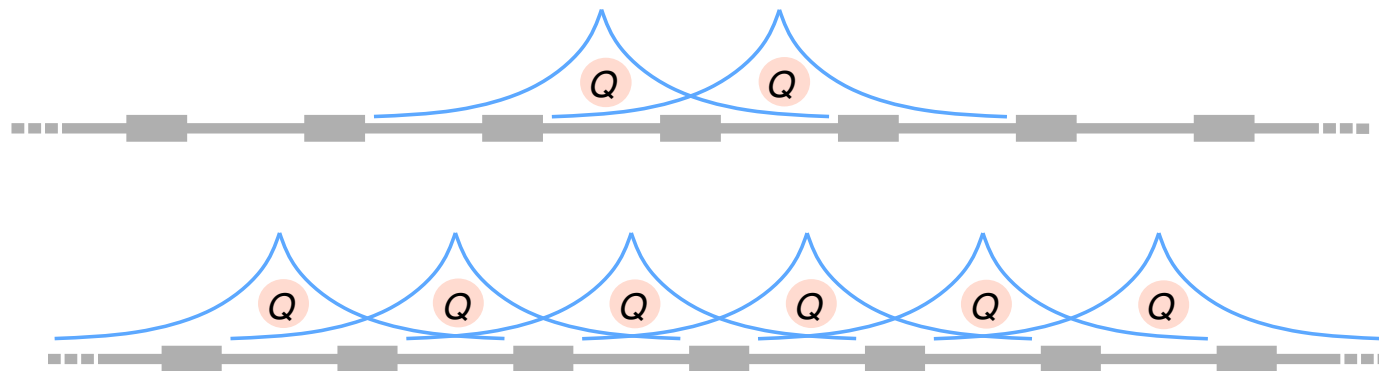
Photonic Crystal + qubits

- Effective swap interaction
- All modes in parallel

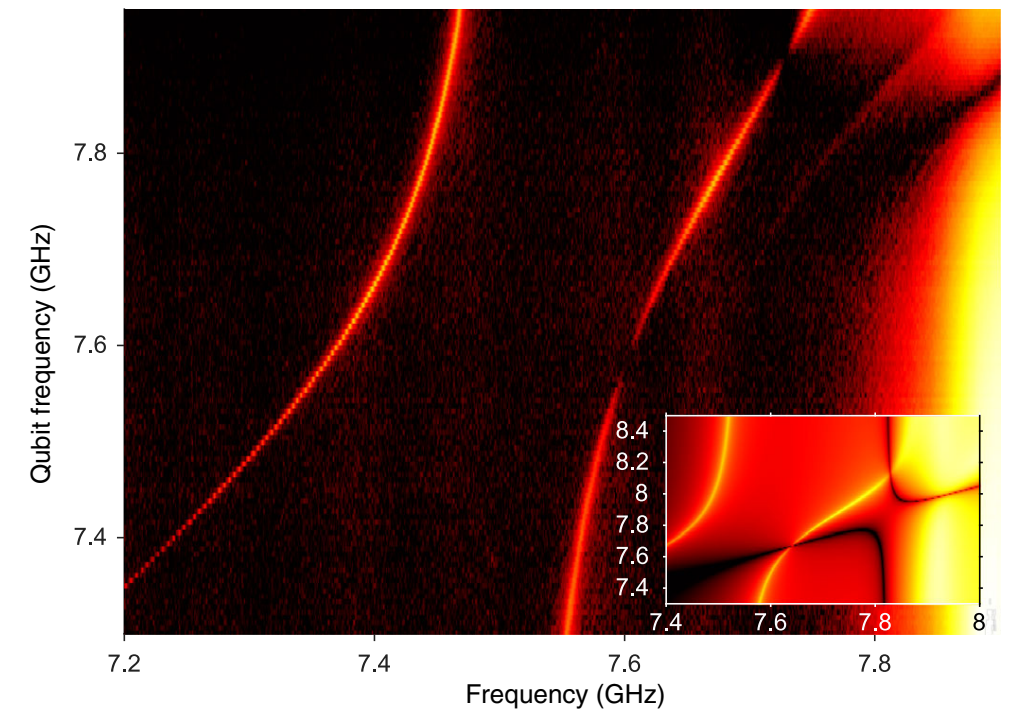
$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

1D-Photonic Crystal

- Exponentially localized bound state



Photon-Mediated Avoided Crossing



- Douglas *et al.* Nat. Photon. (2015)
Calajó *et al.* PRA (2016)
Liu *et al.* Nature Physics (2016)
Sundaresan *et al.* PRX (2019)
Ferreira *et al.* arXiv 2001.0324 (2020)

Photon-Mediated Interactions

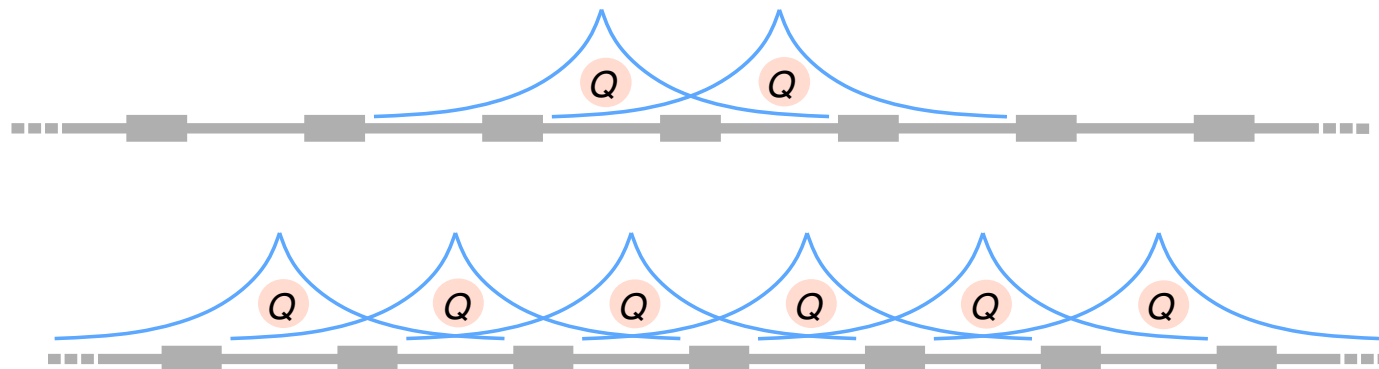
Photonic Crystal + qubits

- Effective swap interaction
- All modes in parallel

$$H = \hbar \sigma_1^+ \sigma_2^- \sum_m \frac{g_m^2}{\Delta(m)} \psi_m(x_1) \psi_m^*(x_2) + h.c.$$

1D-Photonic Crystal

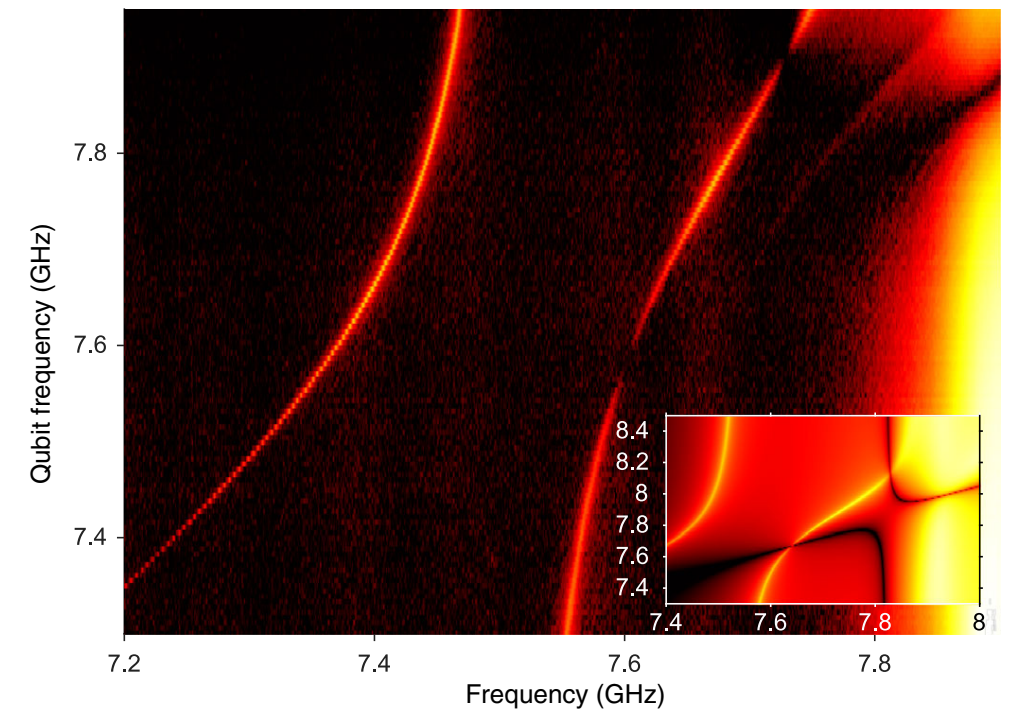
- Exponentially localized bound state



New Regimes:

- New lattices
- Different coupling scheme

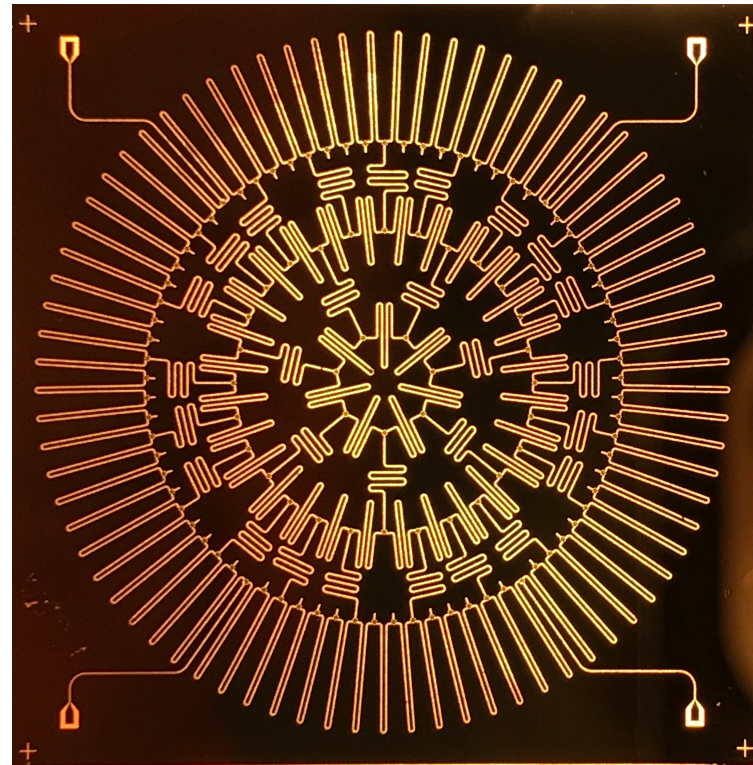
Photon-Mediated Avoided Crossing



- Douglas *et al.* Nat. Photon. (2015)
Calajó *et al.* PRA (2016)
Liu *et al.* Nature Physics (2016)
Sundaresan *et al.* PRX (2019)
Ferreira *et al.* arXiv 2001.0324 (2020)

New Lattices for Photon-Mediated Interactions

AJK *et al.* Nature **571** (2019)



Hyperbolic Lattice

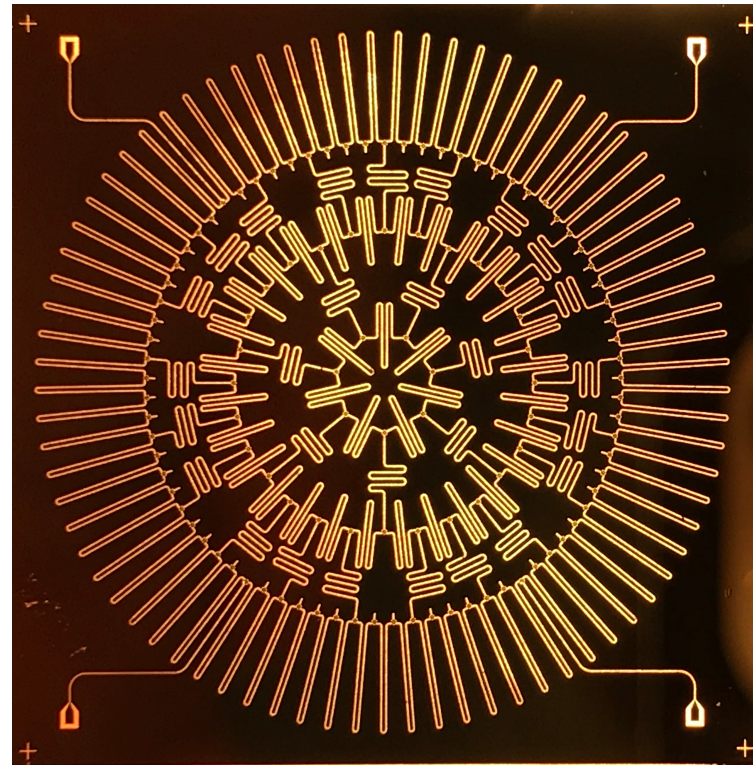
- Follows hyperbolic metric

New Lattices for Photon-Mediated Interactions

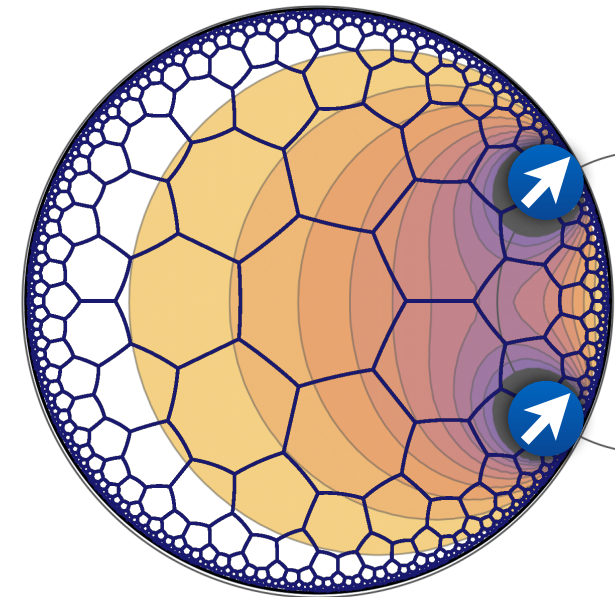
Hyperbolic Lattice

- Follows hyperbolic metric

AJK *et al.* Nature **571** (2019)



Bienias, AJK *et al.* arXiv:2105.06490 (2021)

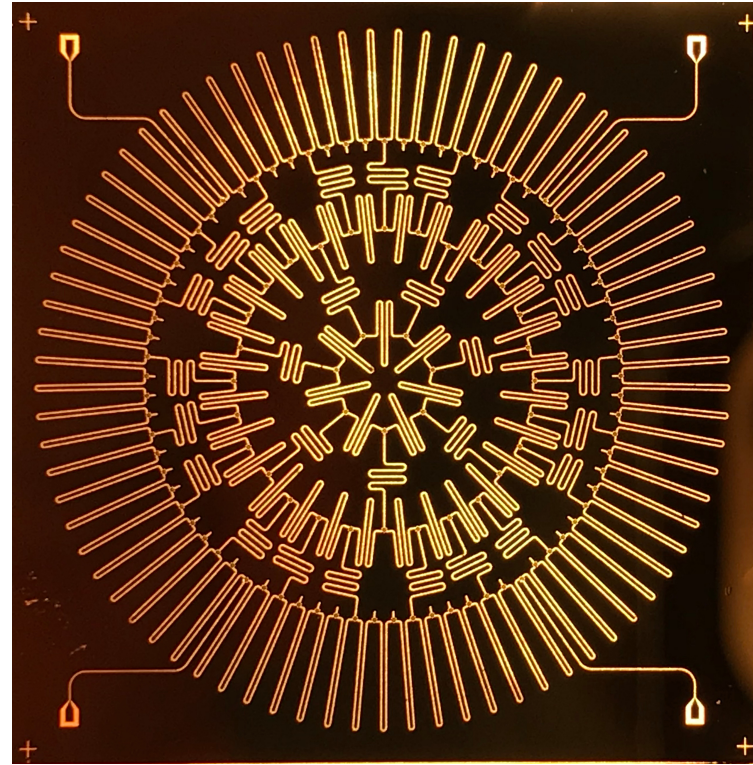


New Lattices for Photon-Mediated Interactions

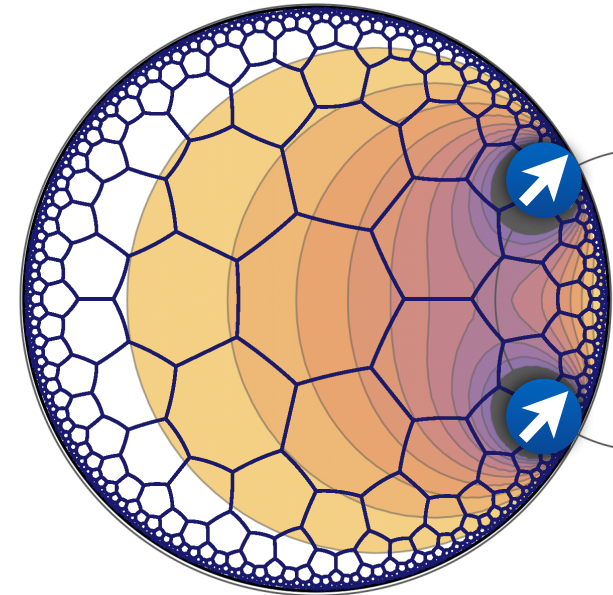
Hyperbolic Lattice

- Follows hyperbolic metric

AJK *et al.* Nature **571** (2019)



Bienias, AJK *et al.* arXiv:2105.06490 (2021)



Flat-Band Lattice

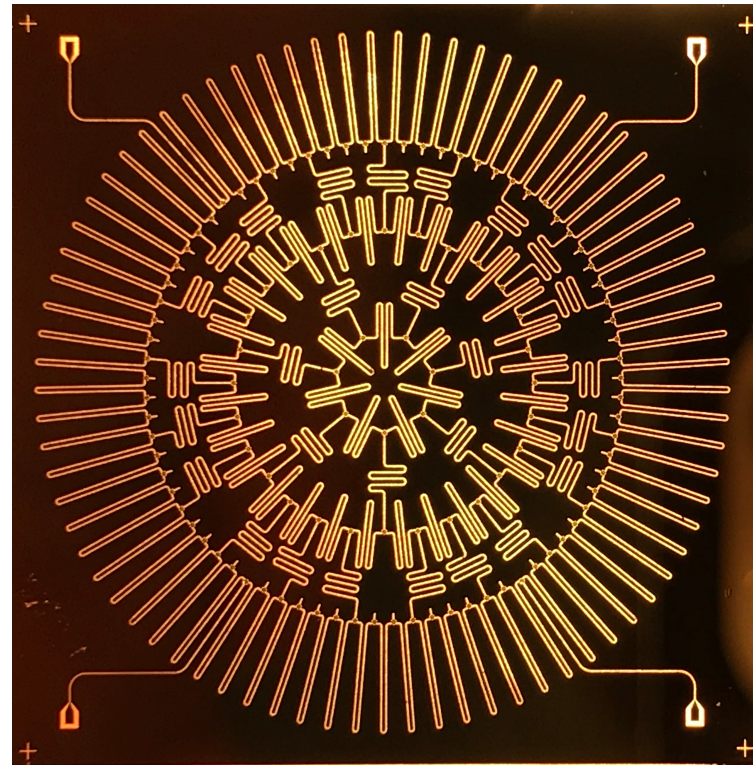
- Frustrated Magnet

New Lattices for Photon-Mediated Interactions

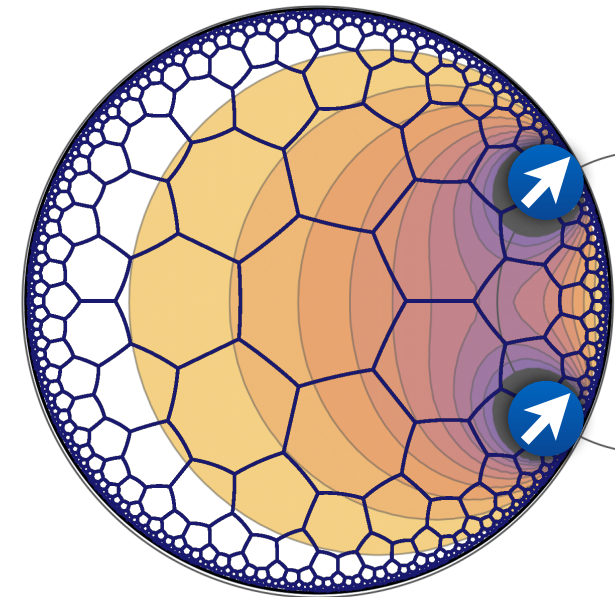
Hyperbolic Lattice

- Follows hyperbolic metric

AJK *et al.* Nature **571** (2019)

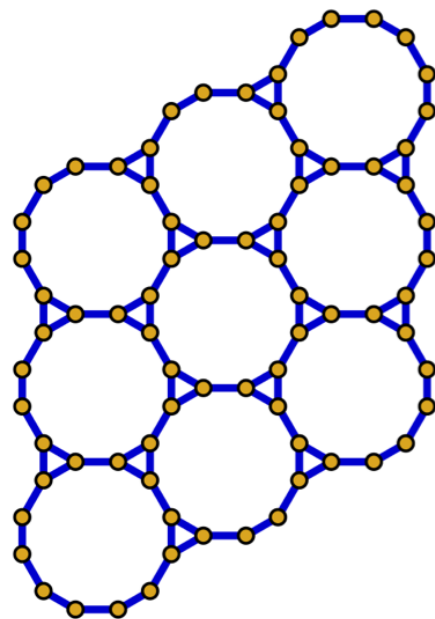


Bienias, AJK *et al.* arXiv:2105.06490 (2021)



Flat-Band Lattice

- Frustrated Magnet

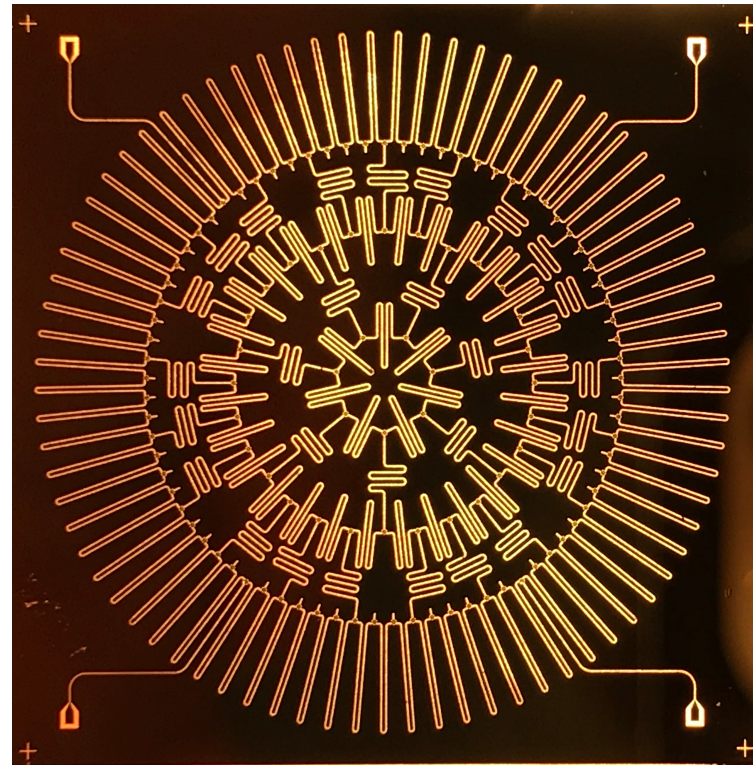


New Lattices for Photon-Mediated Interactions

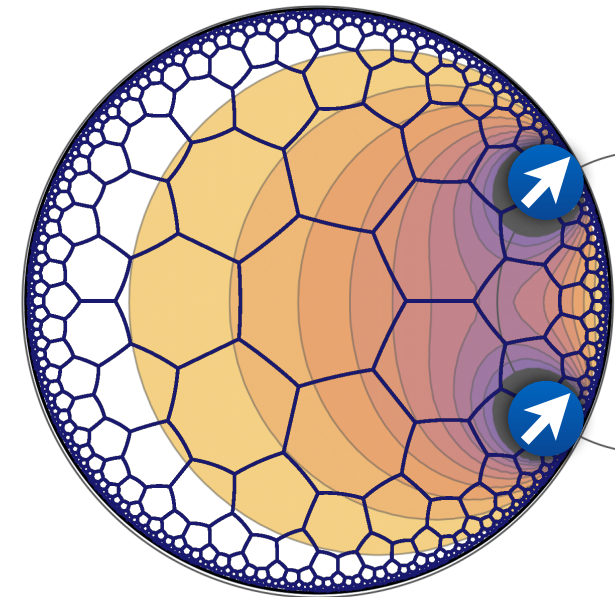
Hyperbolic Lattice

- Follows hyperbolic metric

AJK *et al.* Nature **571** (2019)

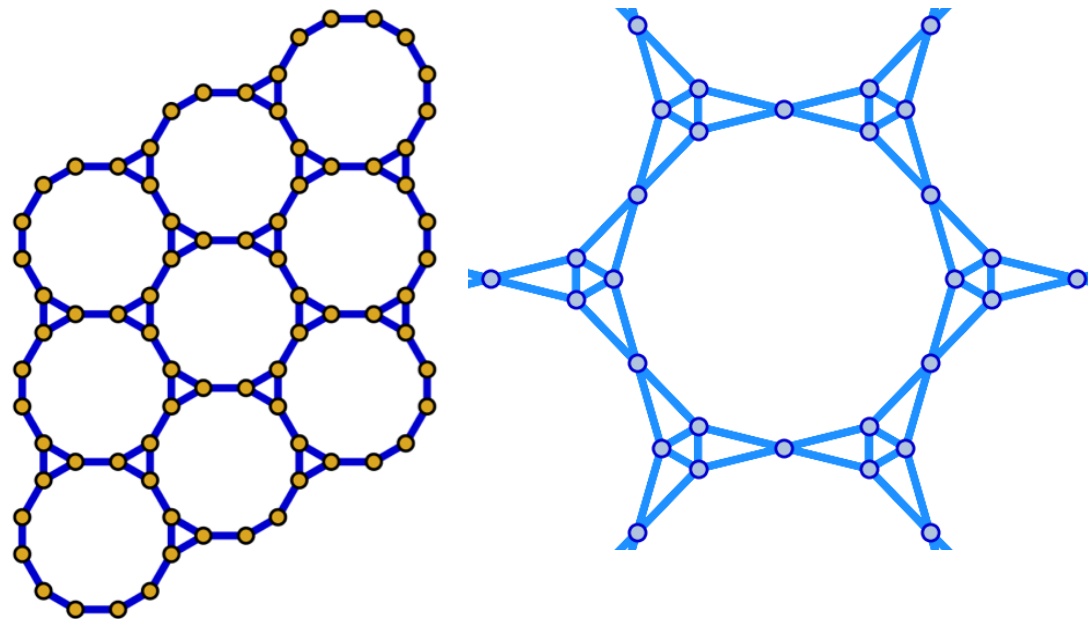


Bienias, AJK *et al.* arXiv:2105.06490 (2021)



Flat-Band Lattice

- Frustrated Magnet

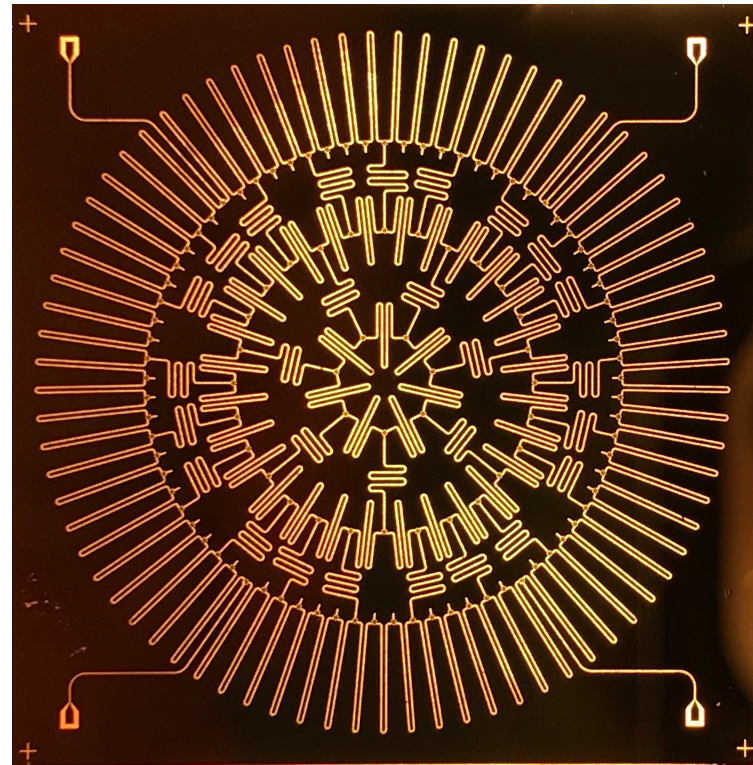


New Lattices for Photon-Mediated Interactions

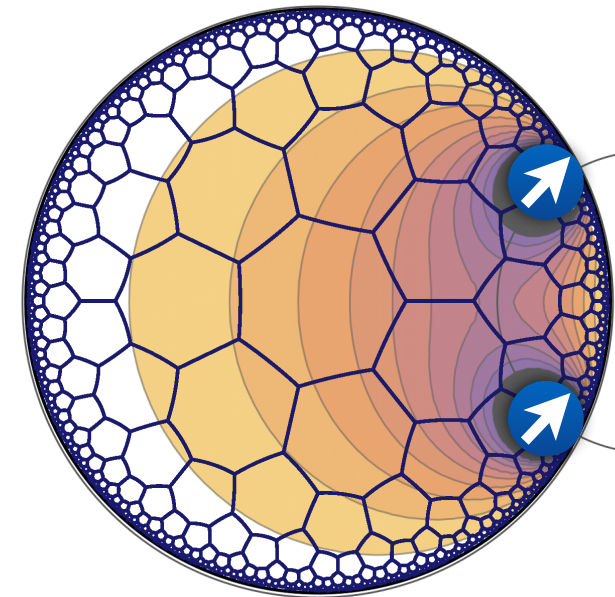
Hyperbolic Lattice

- Follows hyperbolic metric

AJK *et al.* Nature **571** (2019)

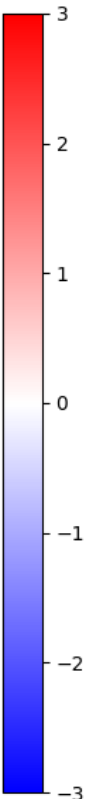
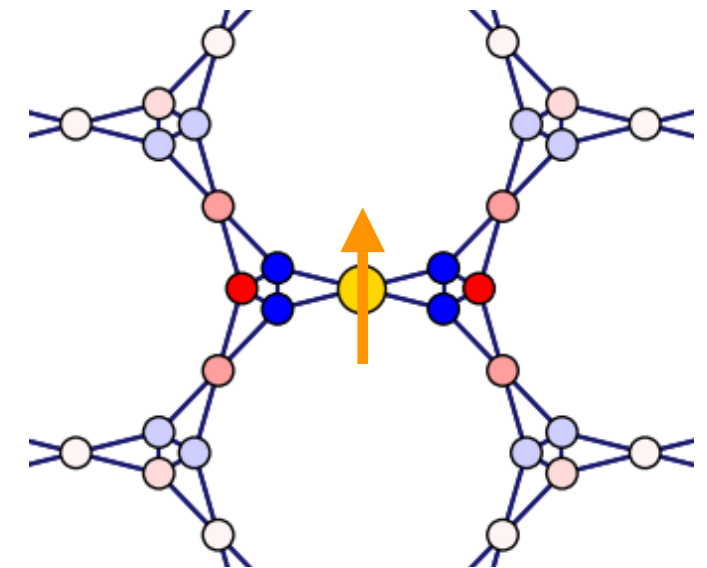
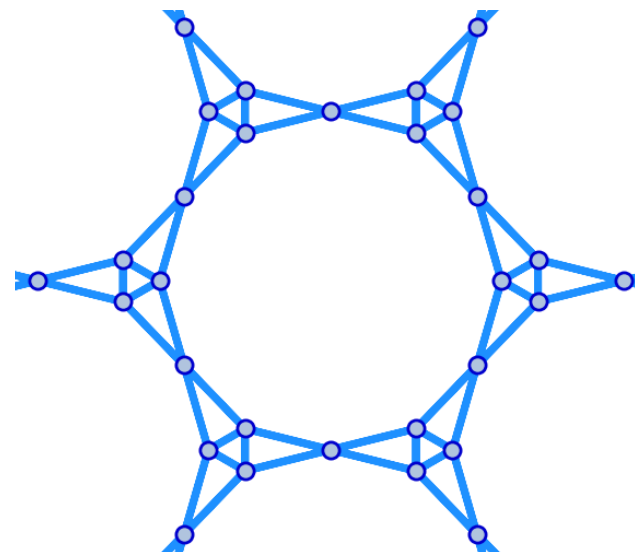
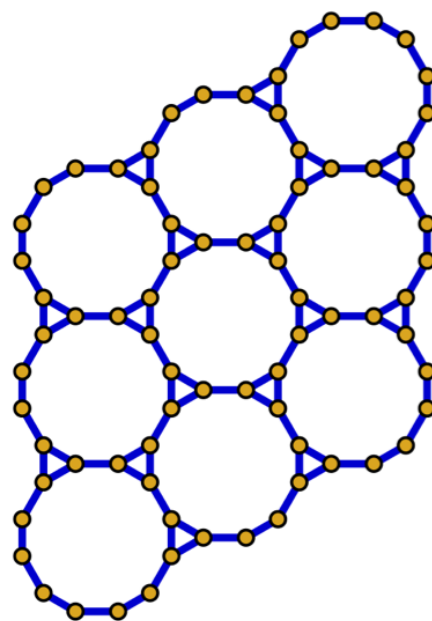


Bienias, AJK *et al.* arXiv:2105.06490 (2021)



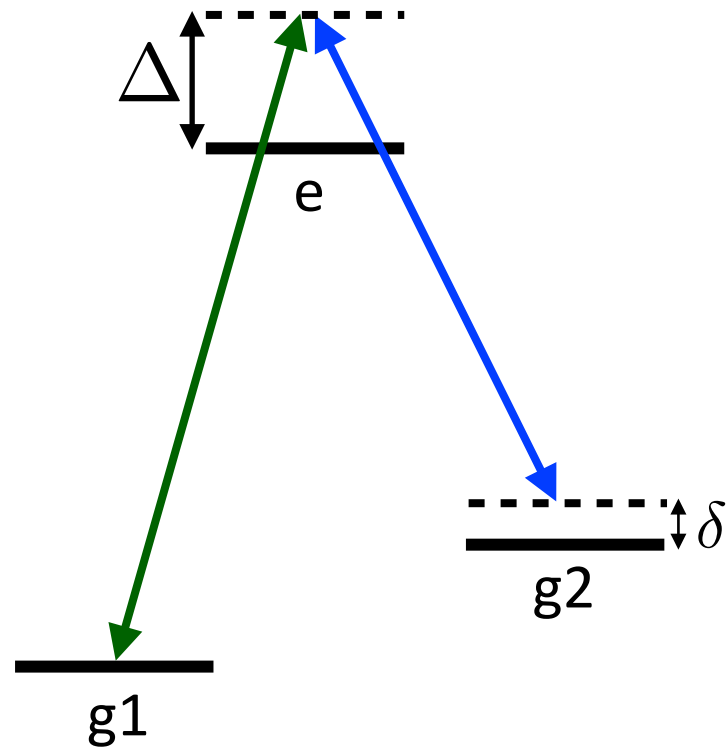
Flat-Band Lattice

- Frustrated Magnet

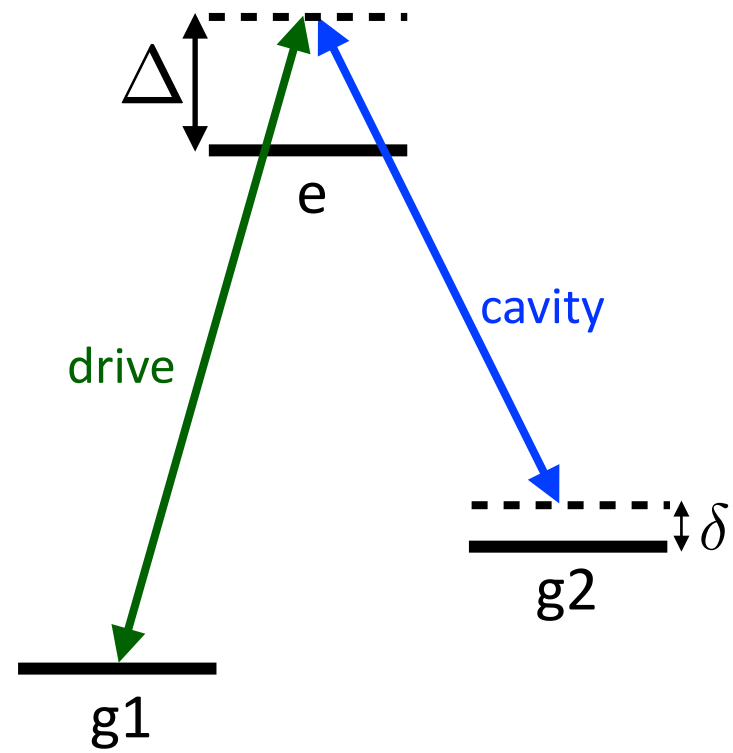


AJK *et al.* Comm. Math. Phys. **376**, 1909 (2020)

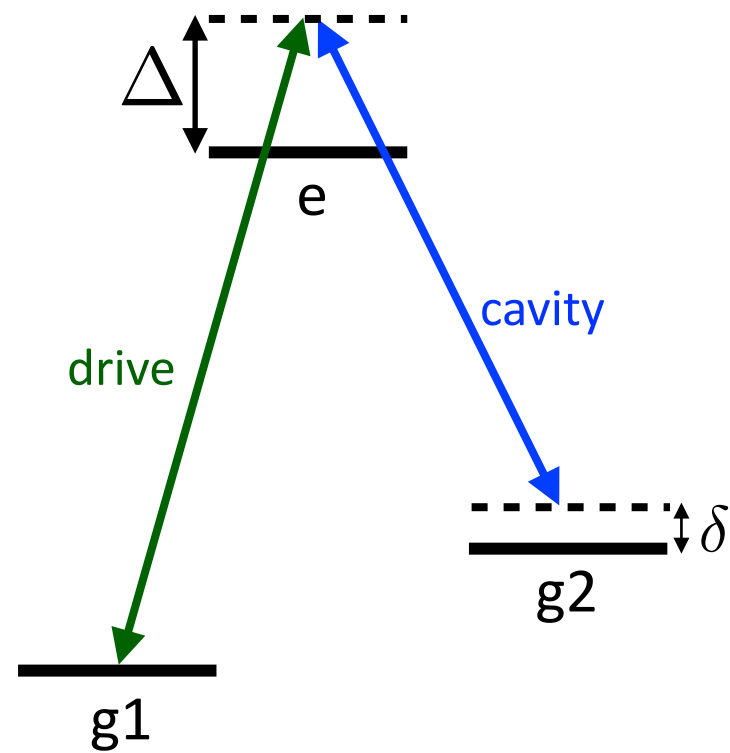
Raman-Coupled Spin Models



Raman-Coupled Spin Models

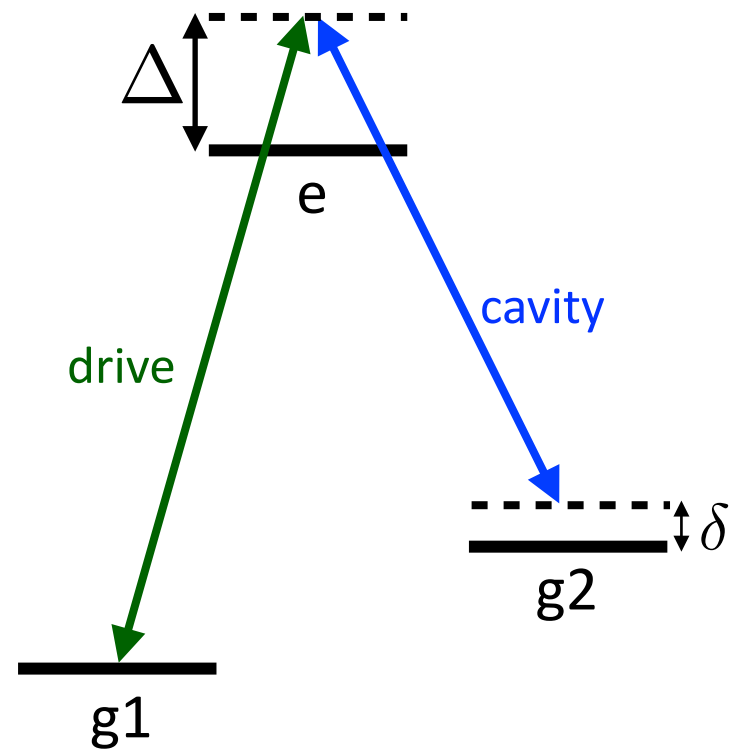


Raman-Coupled Spin Models



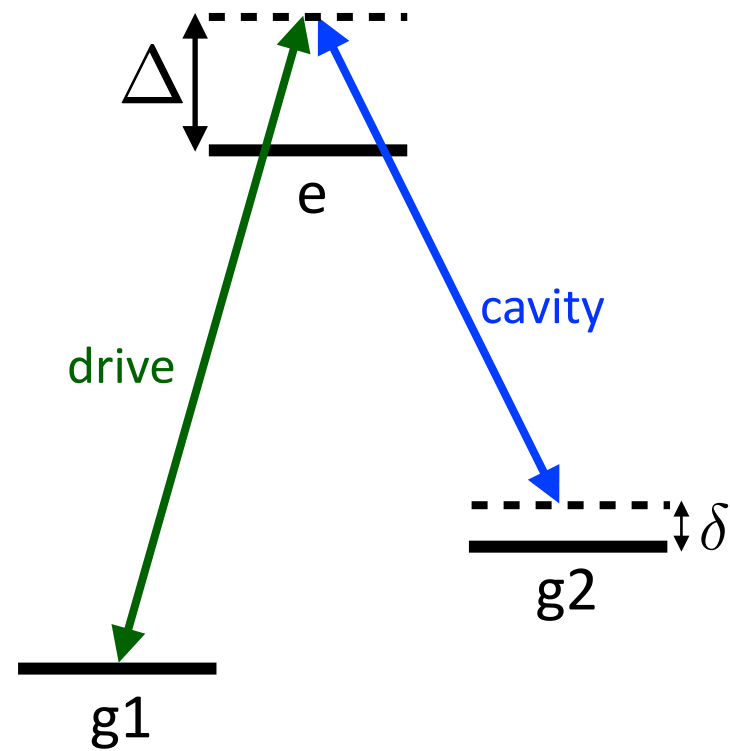
- Microwave-activated coupling

Raman-Coupled Spin Models



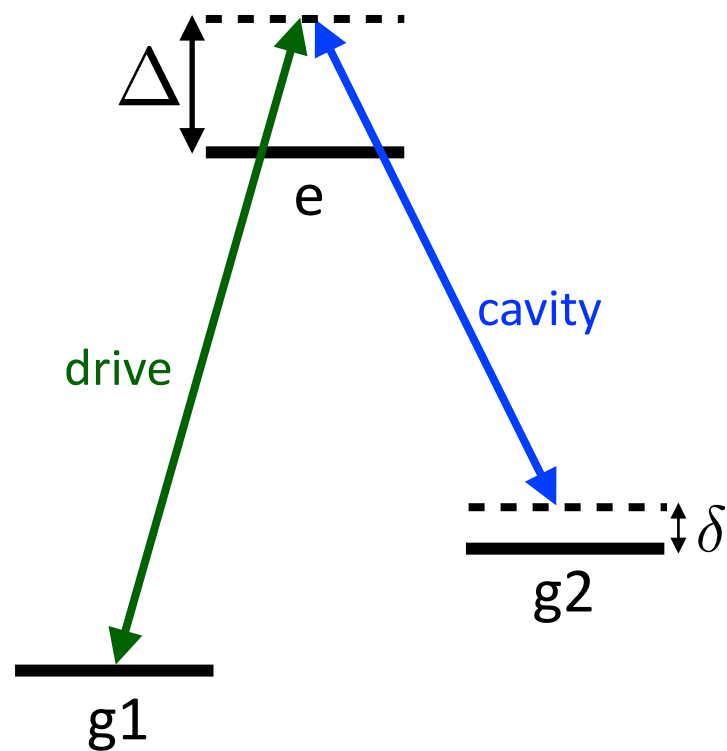
- Microwave-activated coupling
- Two relevant detunings

Raman-Coupled Spin Models



- Microwave-activated coupling
- Two relevant detunings
- Effective swap interaction

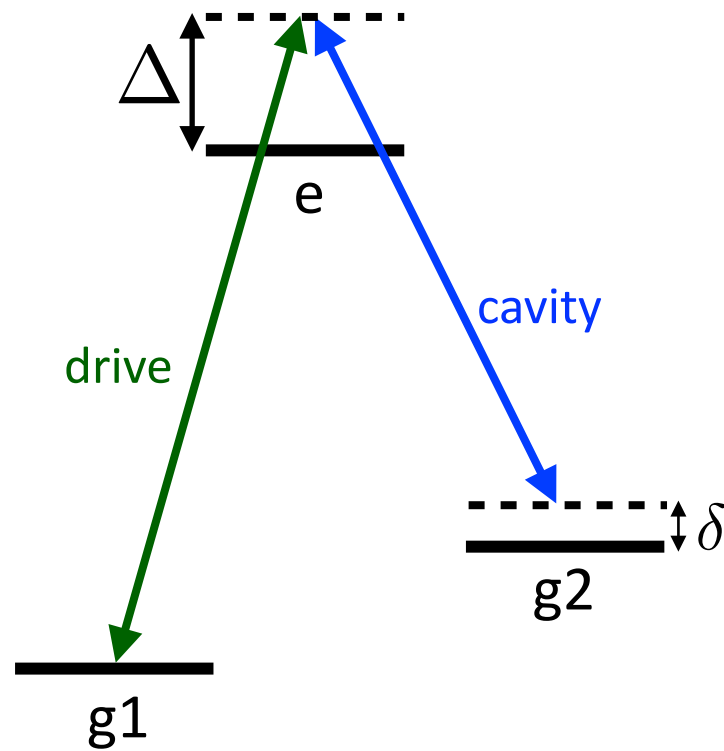
Raman-Coupled Spin Models



- Microwave-activated coupling
- Two relevant detunings
- Effective swap interaction

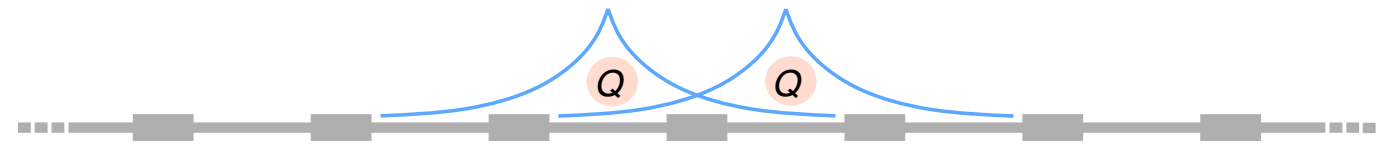
$$H_{Raman} = \hbar \frac{g^2 \Omega^2}{\Delta^2 \delta} \sigma_1^+ \sigma_2^- + h.c.$$

Raman-Coupled Spin Models



1D-Photonic Crystal + Single Drive

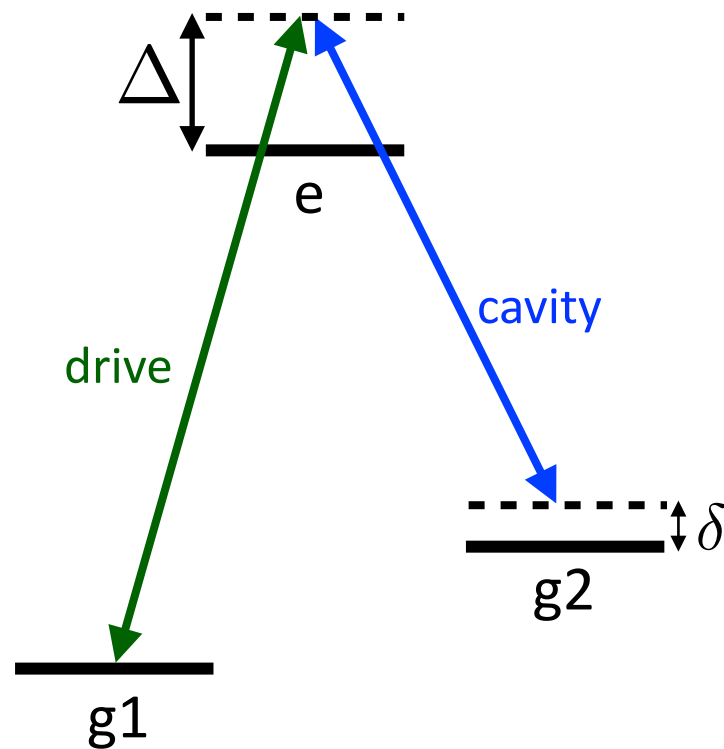
- Exponentially localized interaction



- Microwave-activated coupling
- Two relevant detunings
- Effective swap interaction

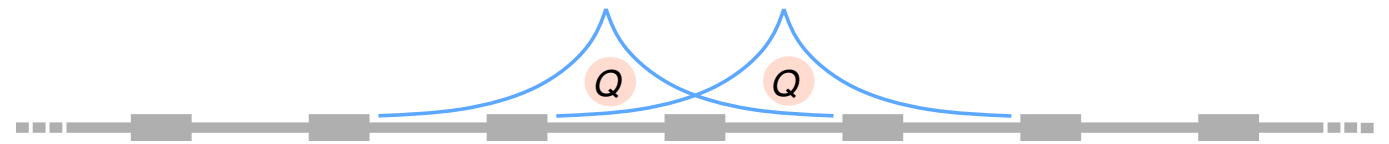
$$H_{Raman} = \hbar \frac{g^2 \Omega^2}{\Delta^2 \delta} \sigma_1^+ \sigma_2^- + h.c.$$

Raman-Coupled Spin Models



1D-Photonic Crystal + Single Drive

- Exponentially localized interaction

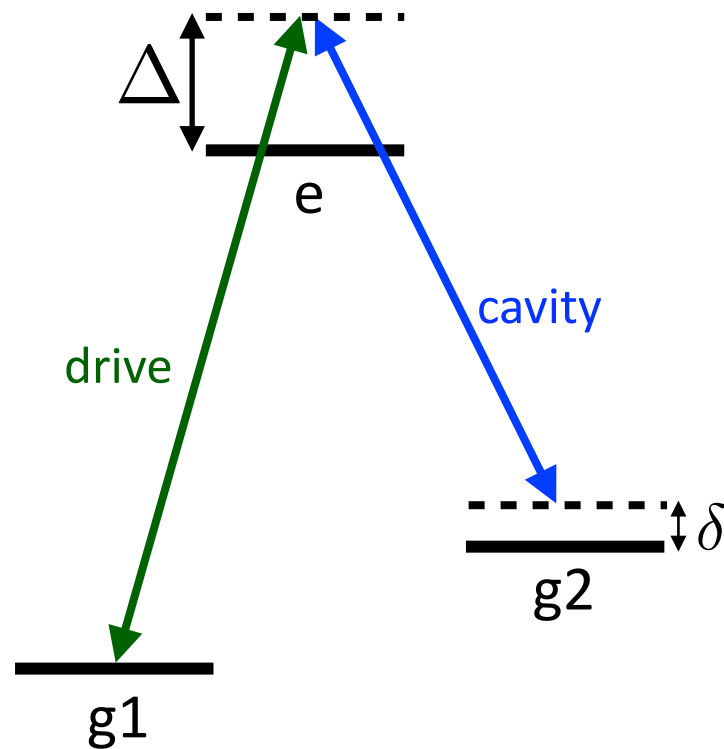


1D-Photonic Crystal + Multiple Drives

- Microwave-activated coupling
- Two relevant detunings
- Effective swap interaction

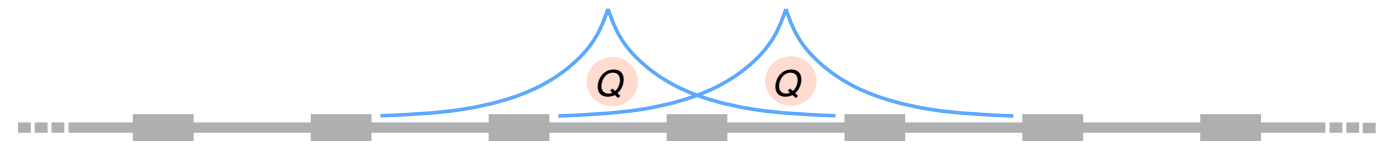
$$H_{Raman} = \hbar \frac{g^2 \Omega^2}{\Delta^2 \delta} \sigma_1^+ \sigma_2^- + h.c.$$

Raman-Coupled Spin Models



1D-Photonic Crystal + Single Drive

- Exponentially localized interaction



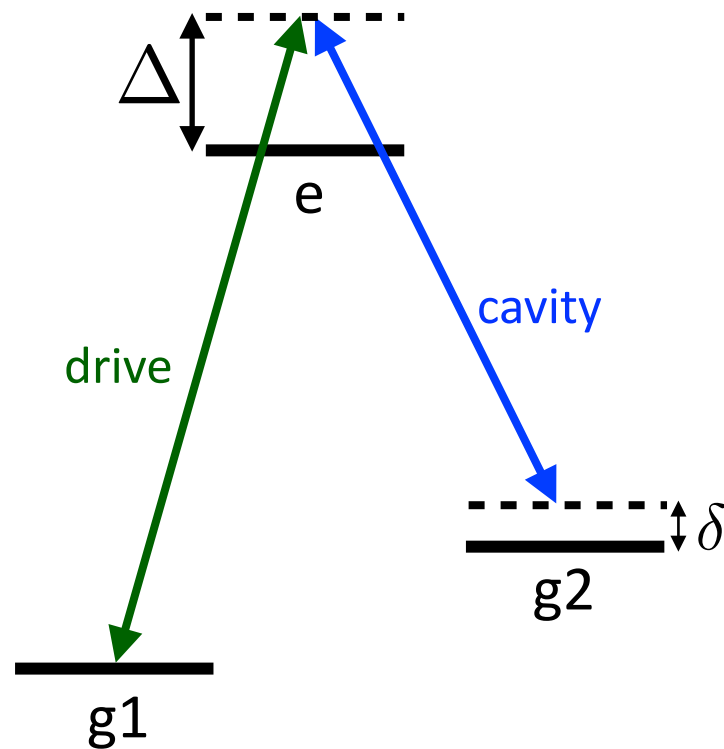
1D-Photonic Crystal + Multiple Drives

- Superposition of exponentials

- Microwave-activated coupling
- Two relevant detunings
- Effective swap interaction

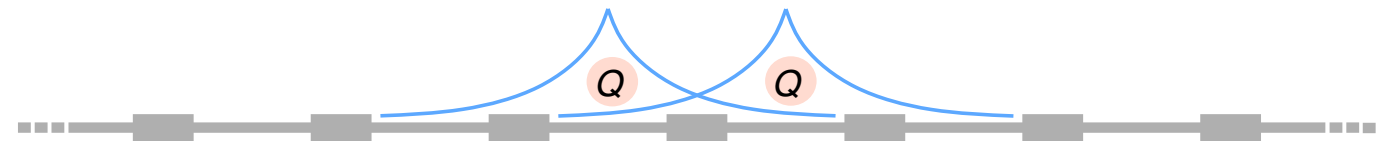
$$H_{Raman} = \hbar \frac{g^2 \Omega^2}{\Delta^2 \delta} \sigma_1^+ \sigma_2^- + h.c.$$

Raman-Coupled Spin Models



1D-Photonic Crystal + Single Drive

- Exponentially localized interaction



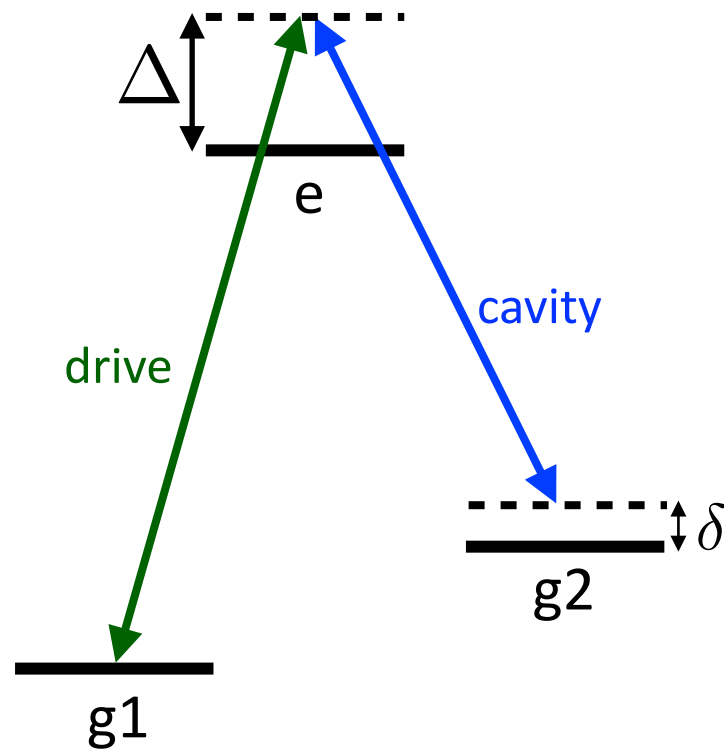
1D-Photonic Crystal + Multiple Drives

- Superposition of exponentials
- Approximate power-law interaction

- Microwave-activated coupling
- Two relevant detunings
- Effective swap interaction

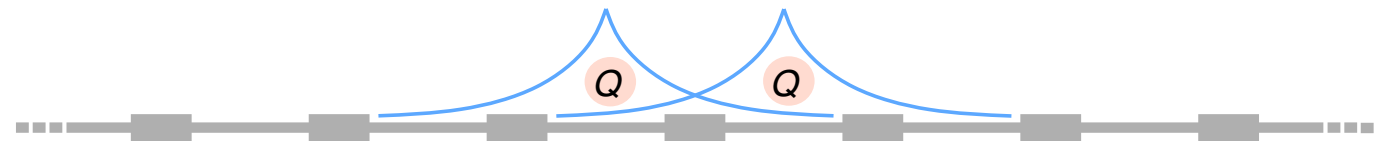
$$H_{Raman} = \hbar \frac{g^2 \Omega^2}{\Delta^2 \delta} \sigma_1^+ \sigma_2^- + h.c.$$

Raman-Coupled Spin Models



1D-Photonic Crystal + Single Drive

- Exponentially localized interaction



1D-Photonic Crystal + Multiple Drives

- Superposition of exponentials
- Approximate power-law interaction



Need 3-level qubit

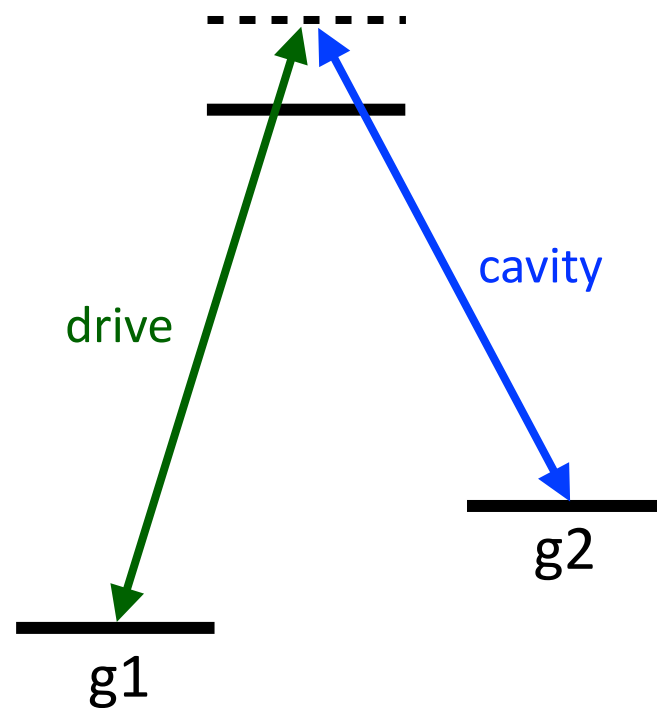
- Microwave-activated coupling
- Two relevant detunings
- Effective swap interaction

$$H_{Raman} = \hbar \frac{g^2 \Omega^2}{\Delta^2 \delta} \sigma_1^+ \sigma_2^- + h.c.$$

Raman Transitions in Fluxonium

Rabi oscillation

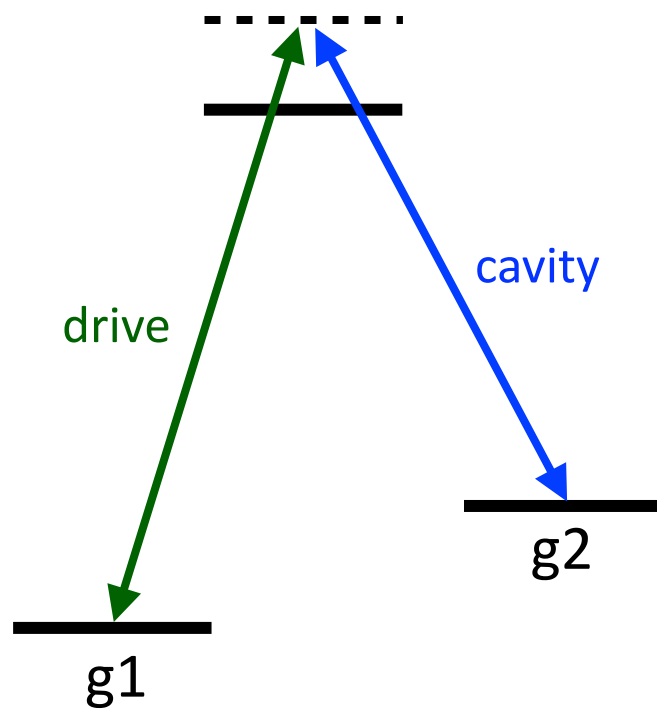
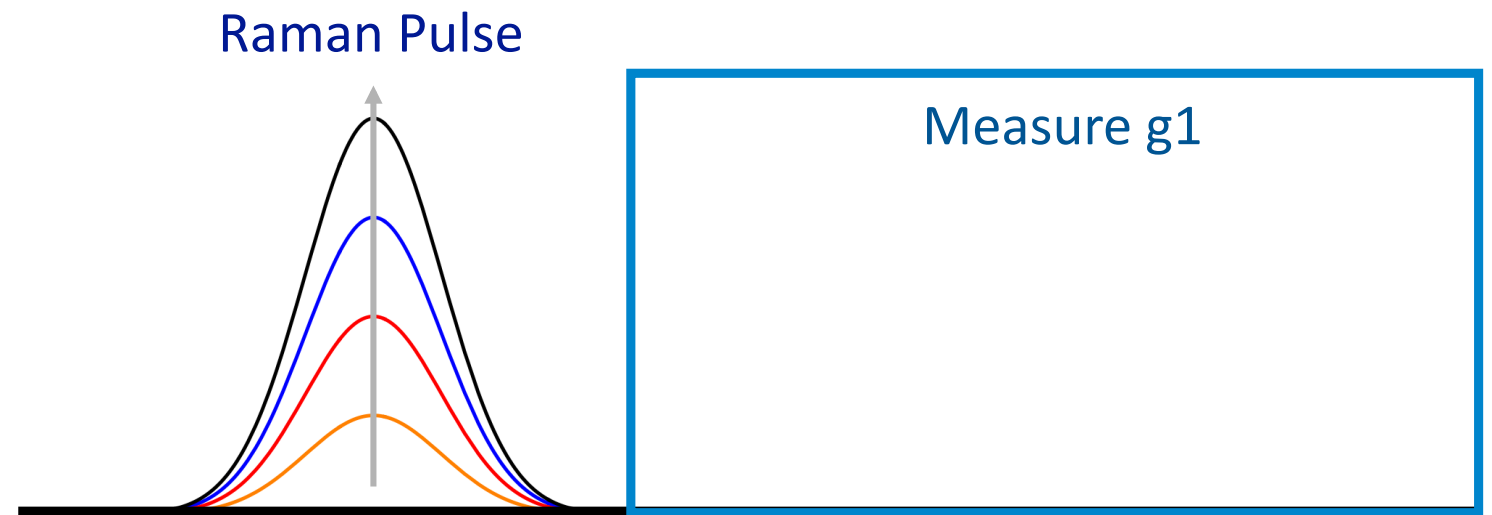
- Gaussian pulse off-resonant of plasmon
- Vacuum Rabi rate of fluxon



Raman Transitions in Fluxonium

Rabi oscillation

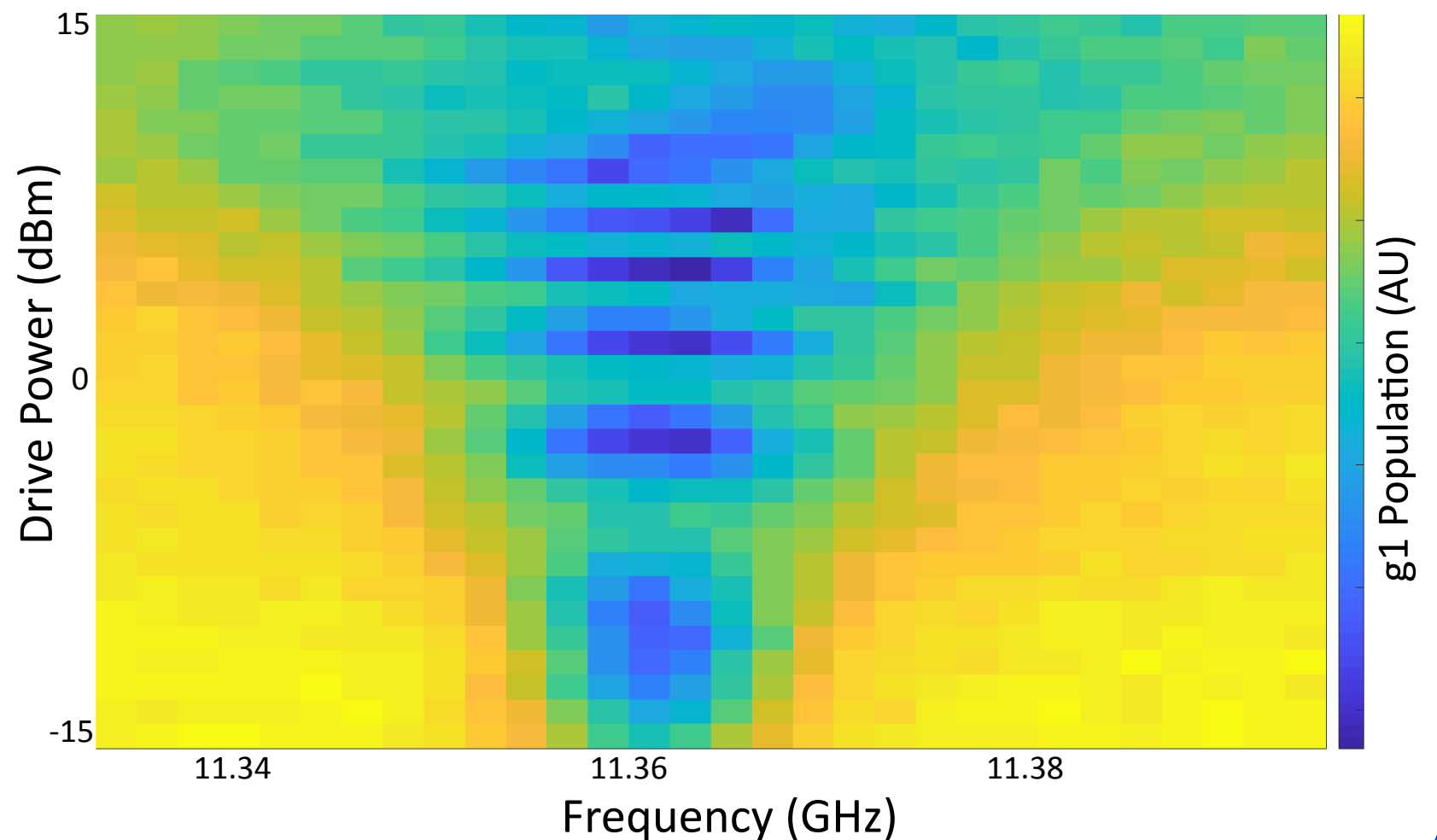
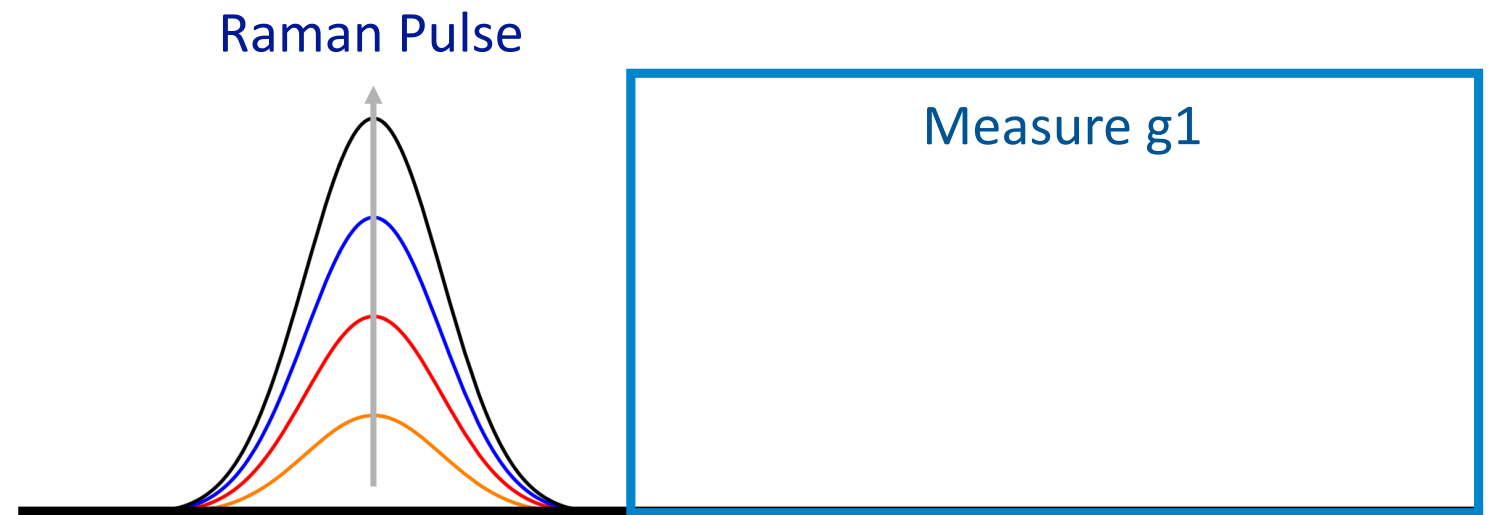
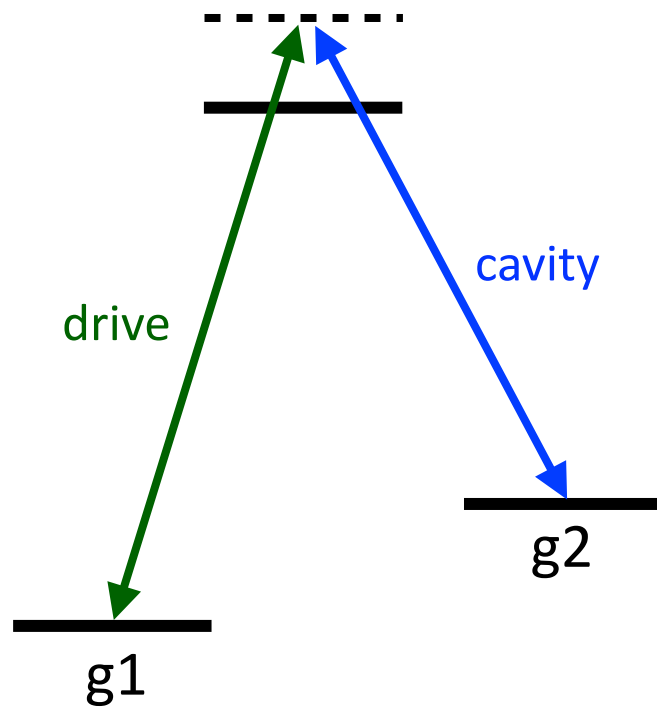
- Gaussian pulse off-resonant of plasmon
- Vacuum Rabi rate of fluxon



Raman Transitions in Fluxonium

Rabi oscillation

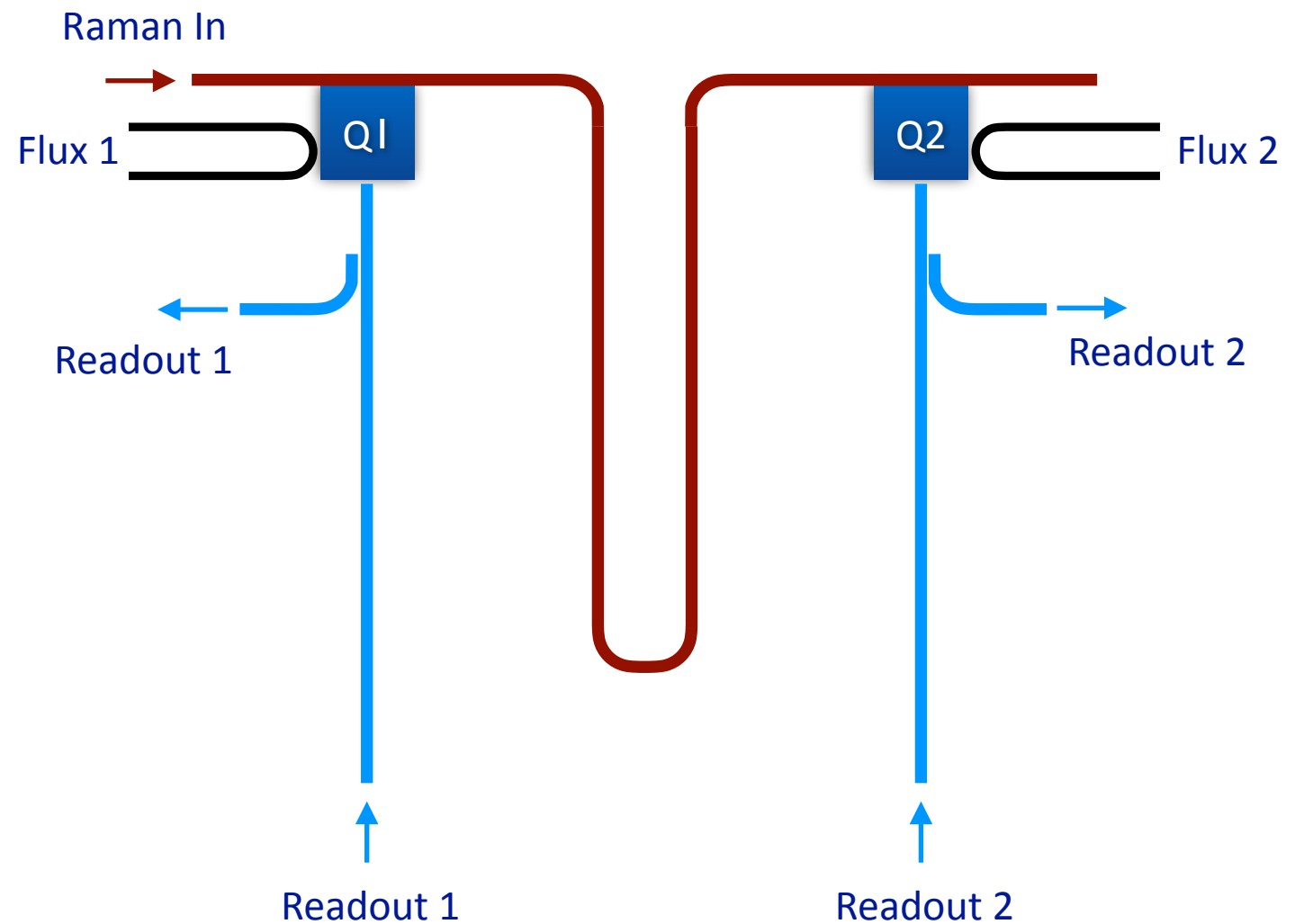
- Gaussian pulse off-resonant of plasmon
- Vacuum Rabi rate of fluxon



Second-Generation Raman Device

Redesigned Device

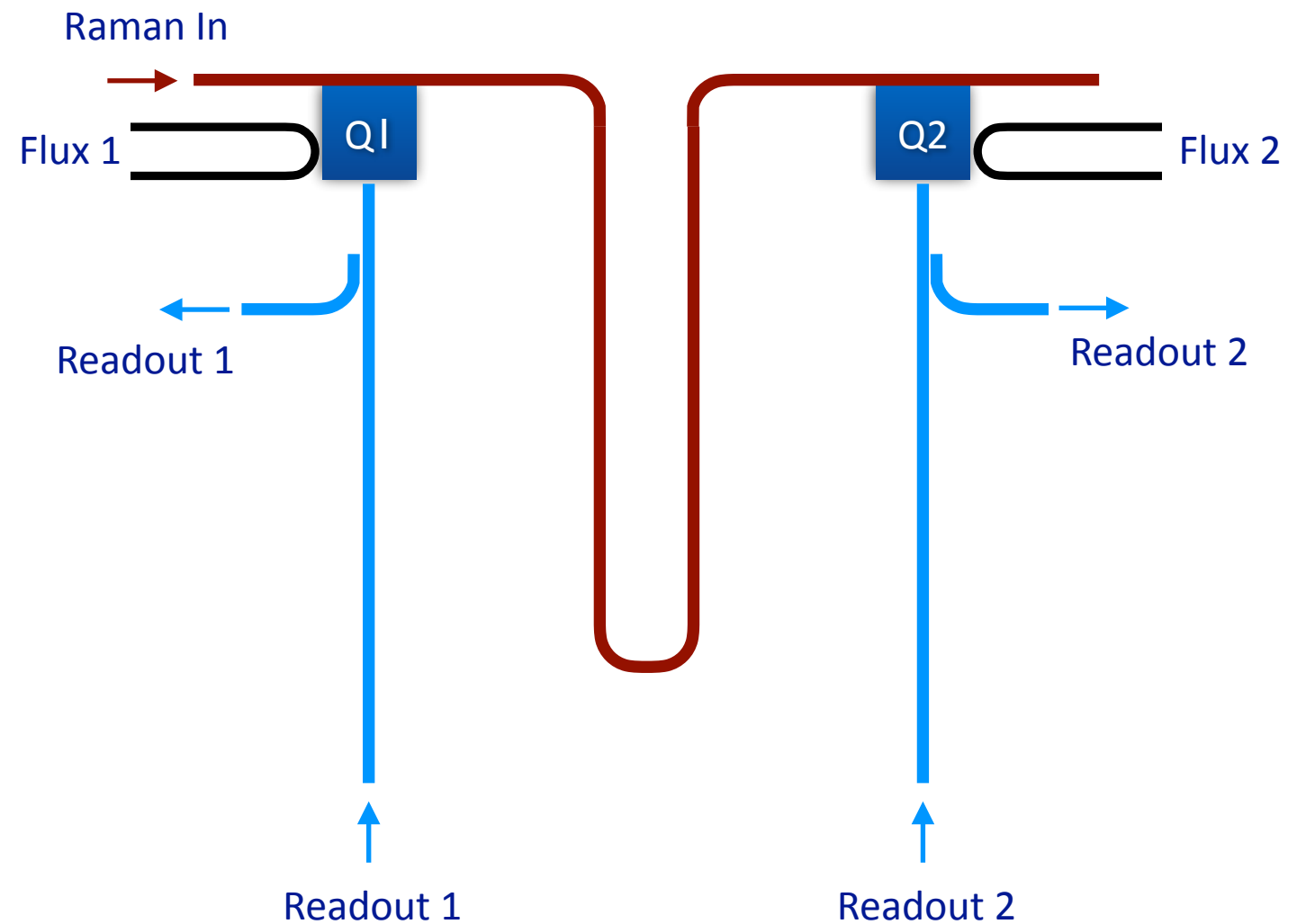
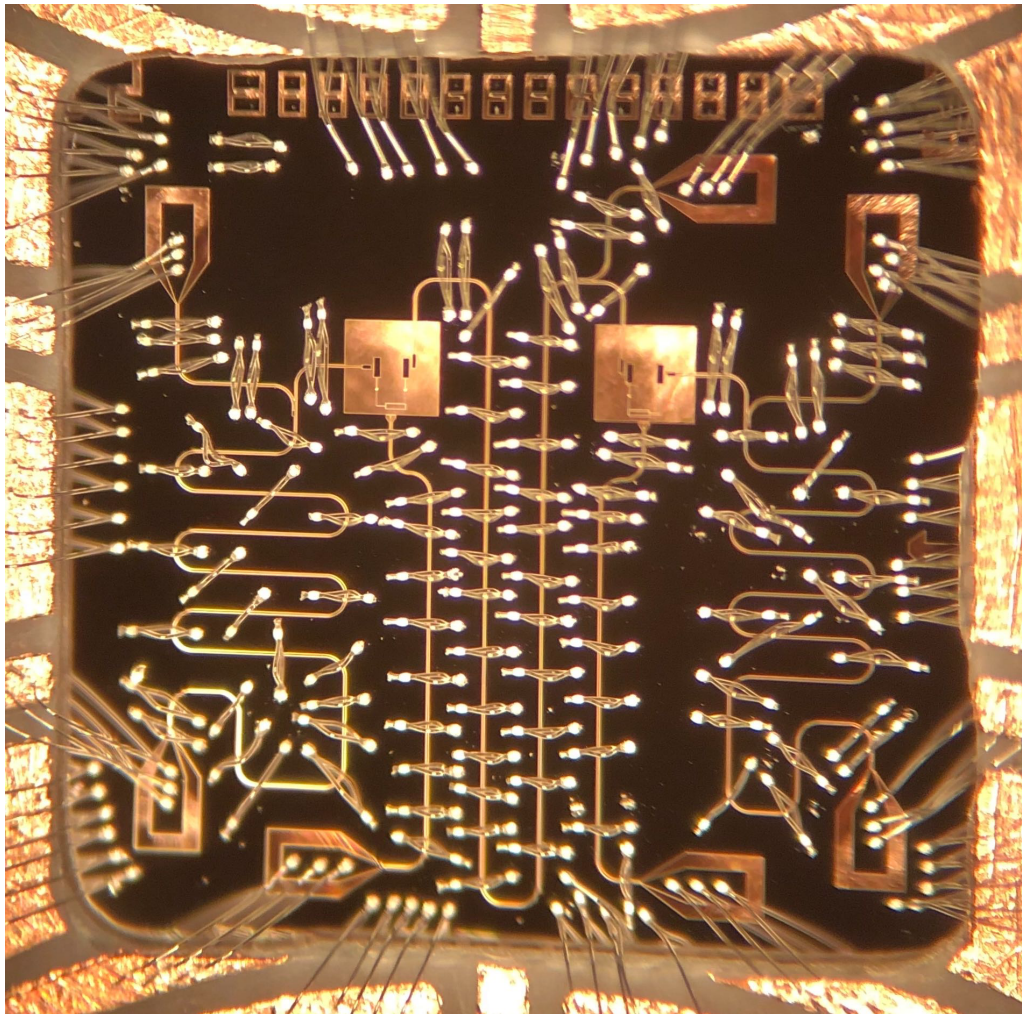
- 3-cavities
- Separate resonators allow
 - Optimized readout
 - Parallel readout and coupling



Second-Generation Raman Device

Redesigned Device

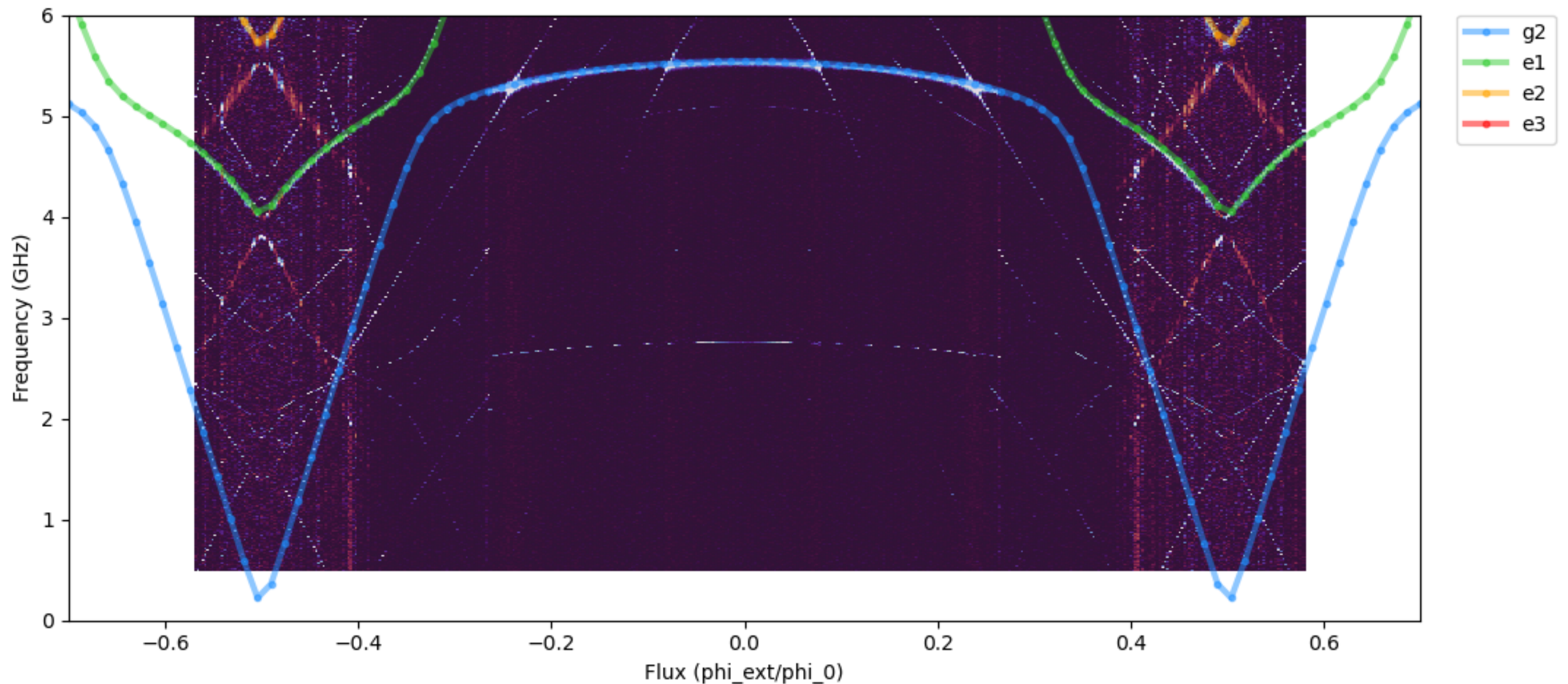
- 3-cavities
- Separate resonators allow
 - Optimized readout
 - Parallel readout and coupling



Second-Generation Raman Device

Redesigned Device

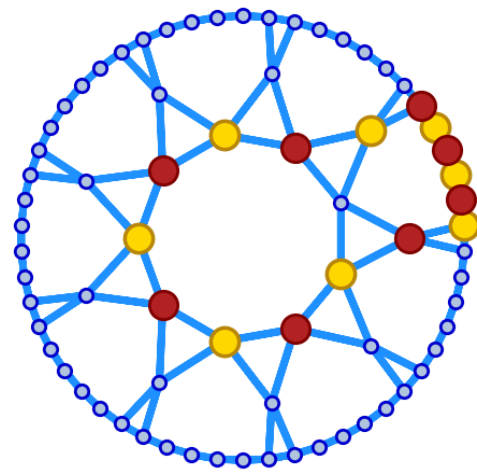
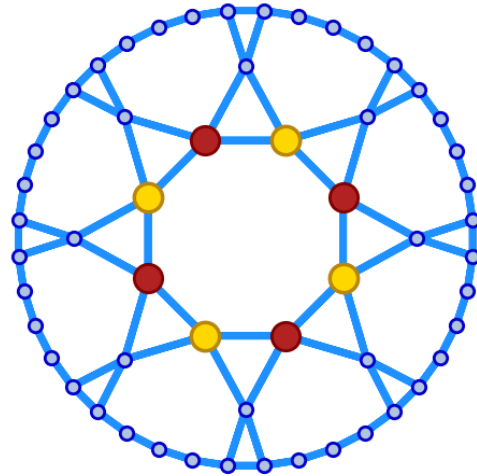
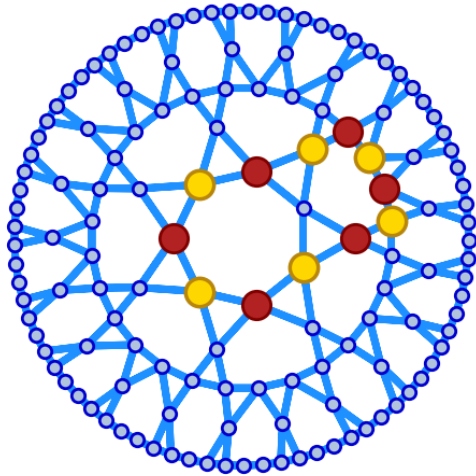
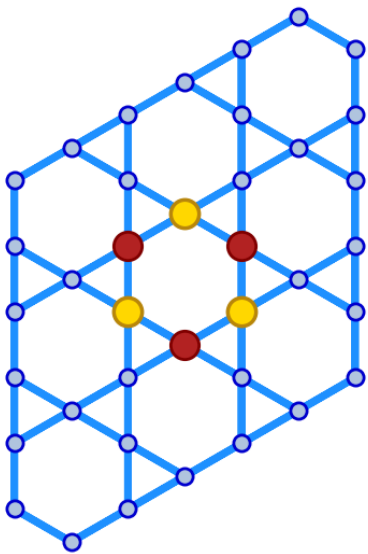
- 3-cavities
- Separate resonators allow
 - Optimized readout
 - Parallel readout and coupling



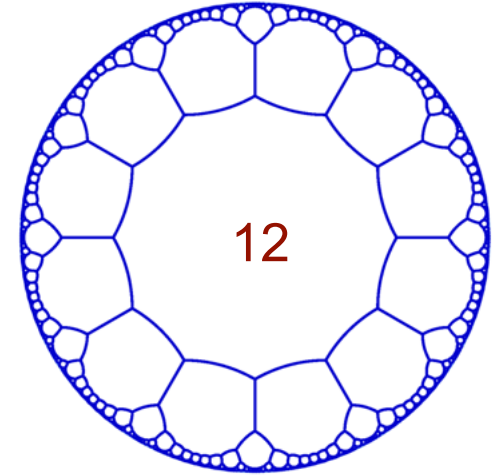
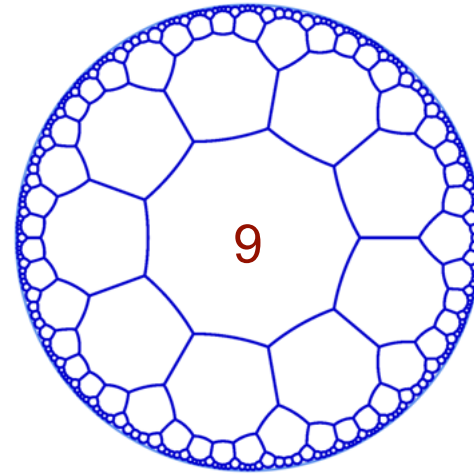
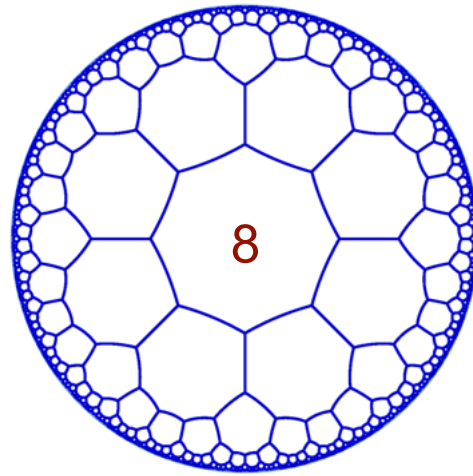
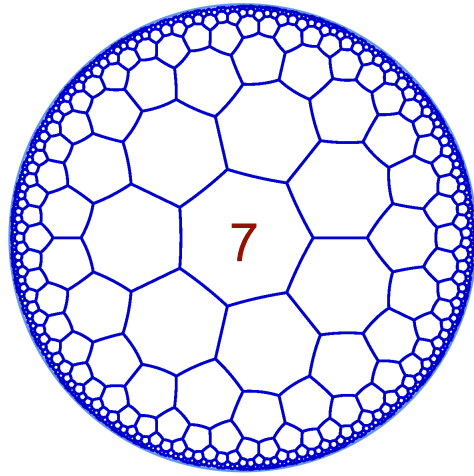
Empty rounded rectangular box at the top of the page.

Large empty rounded rectangular box occupying the majority of the page.

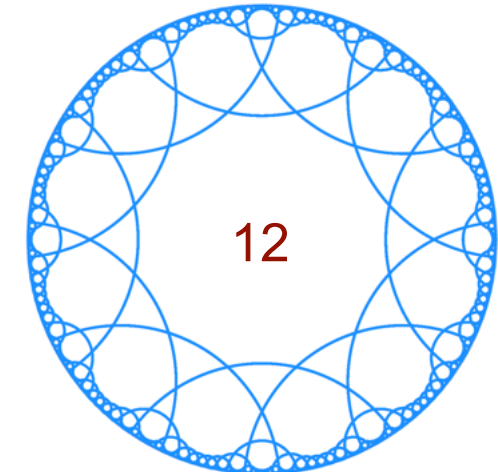
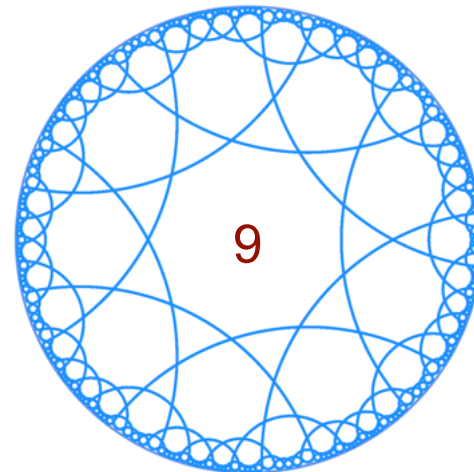
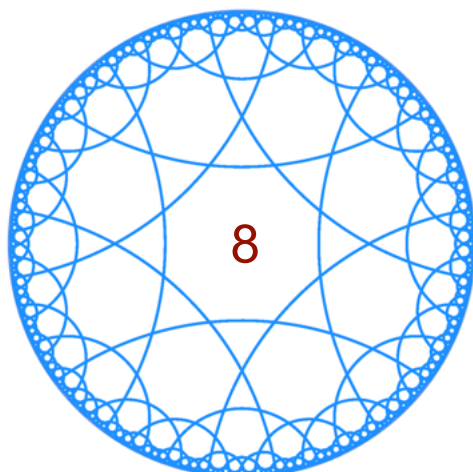
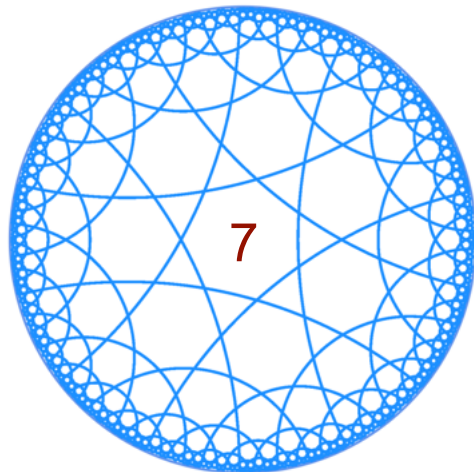
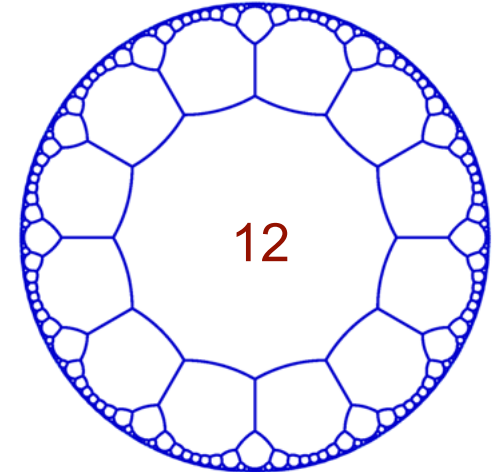
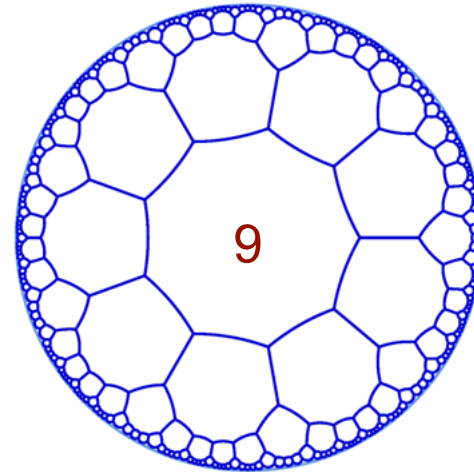
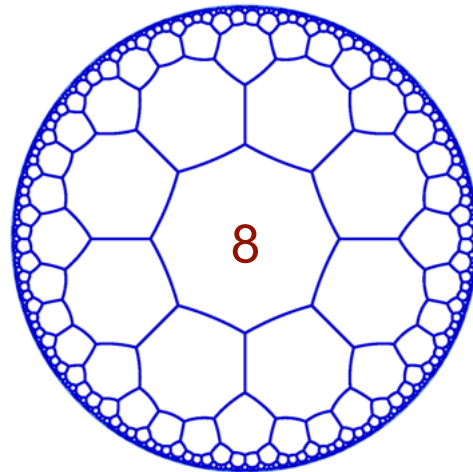
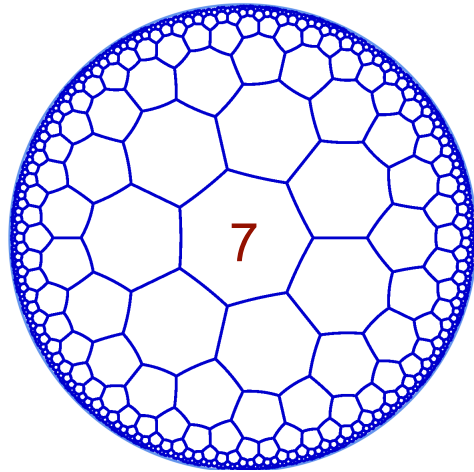
Full-Wave Flat-Band States



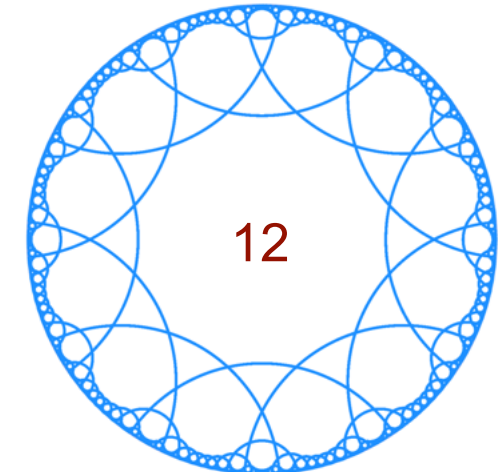
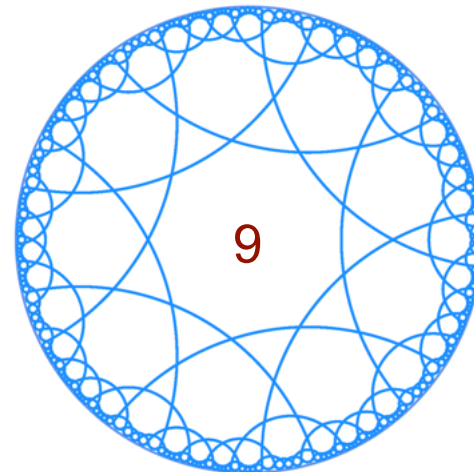
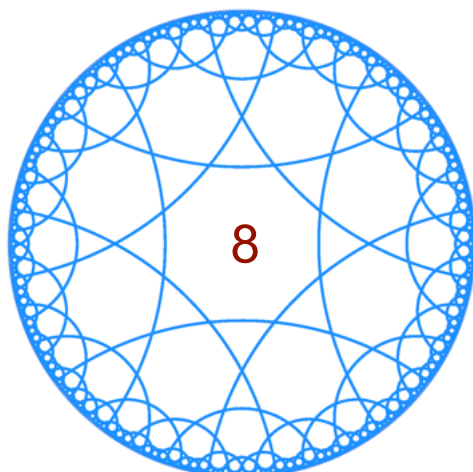
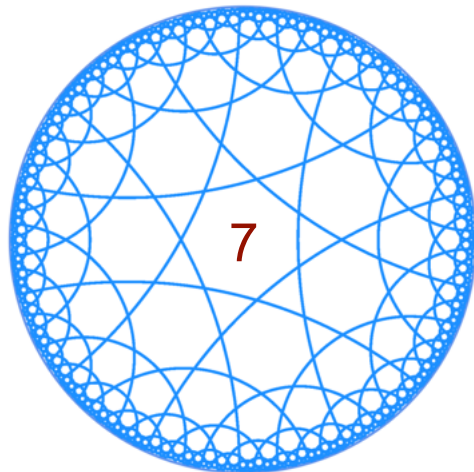
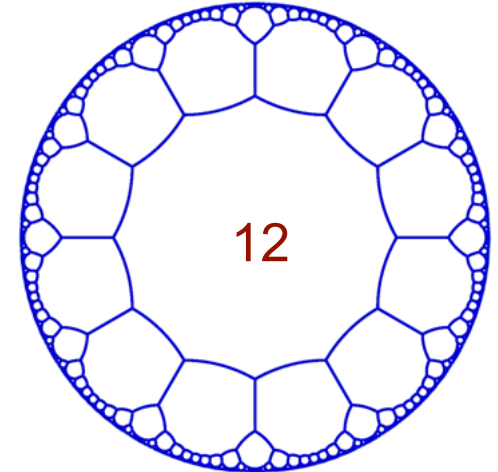
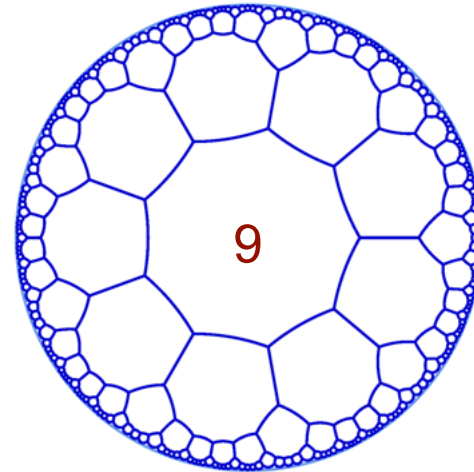
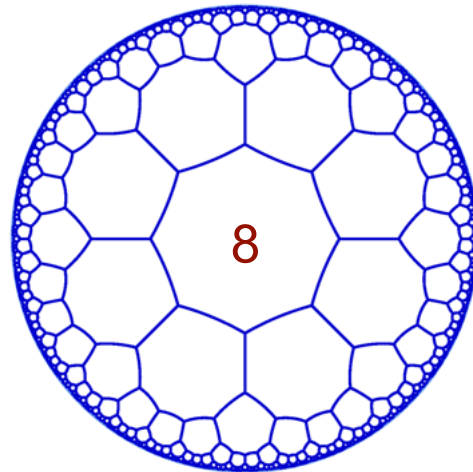
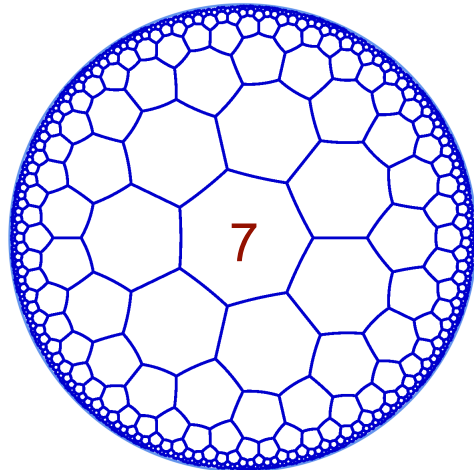
Hyperbolic Lattices and Curvature



Hyperbolic Lattices and Curvature



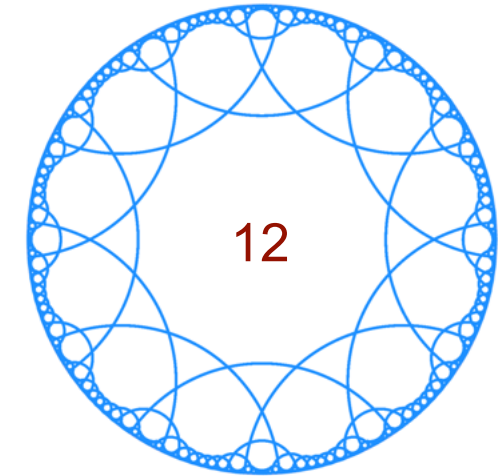
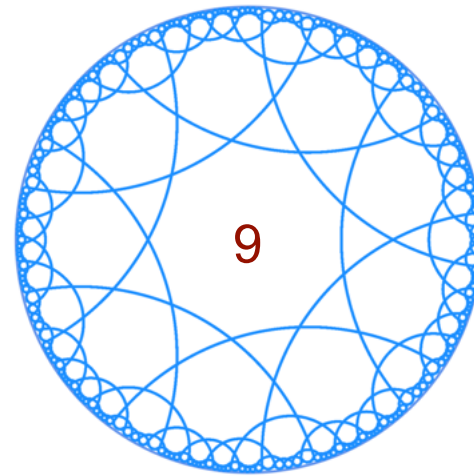
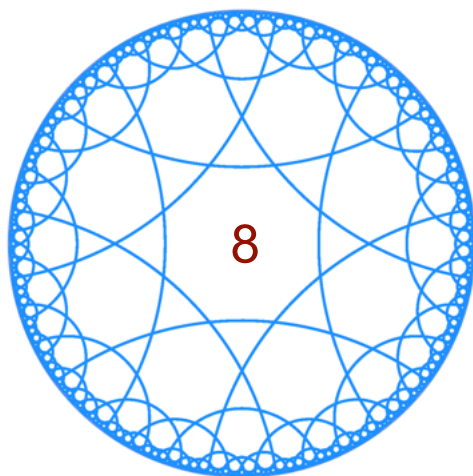
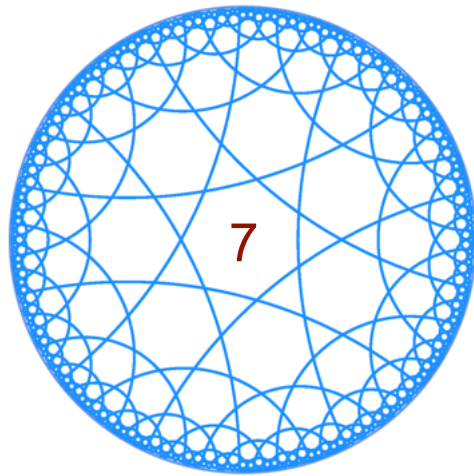
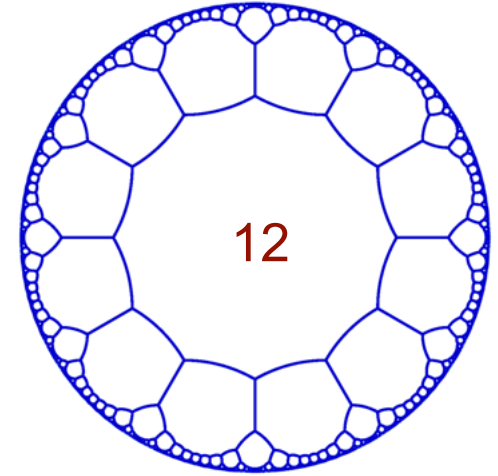
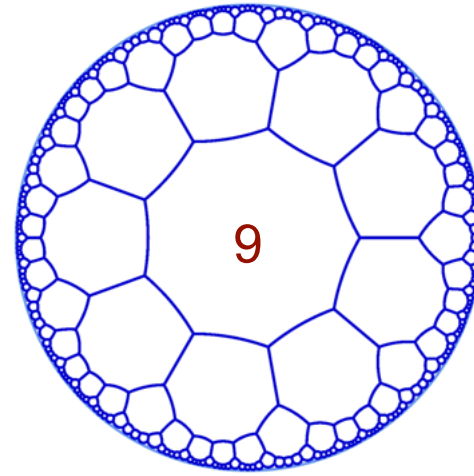
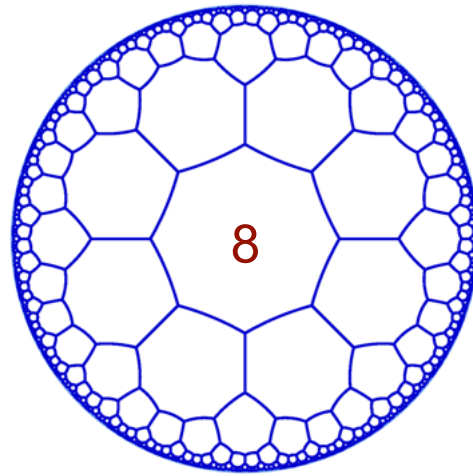
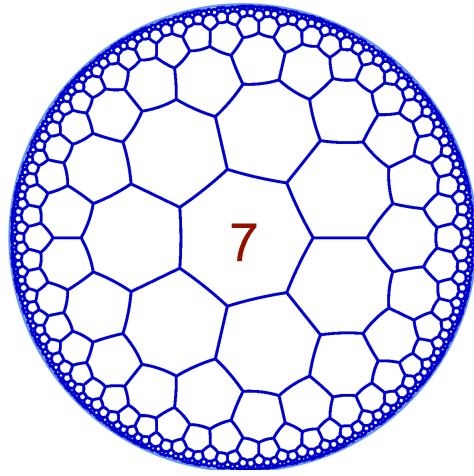
Hyperbolic Lattices and Curvature



Gaussian Curvature

$$K = -\frac{1}{R^2}$$

Hyperbolic Lattices and Curvature

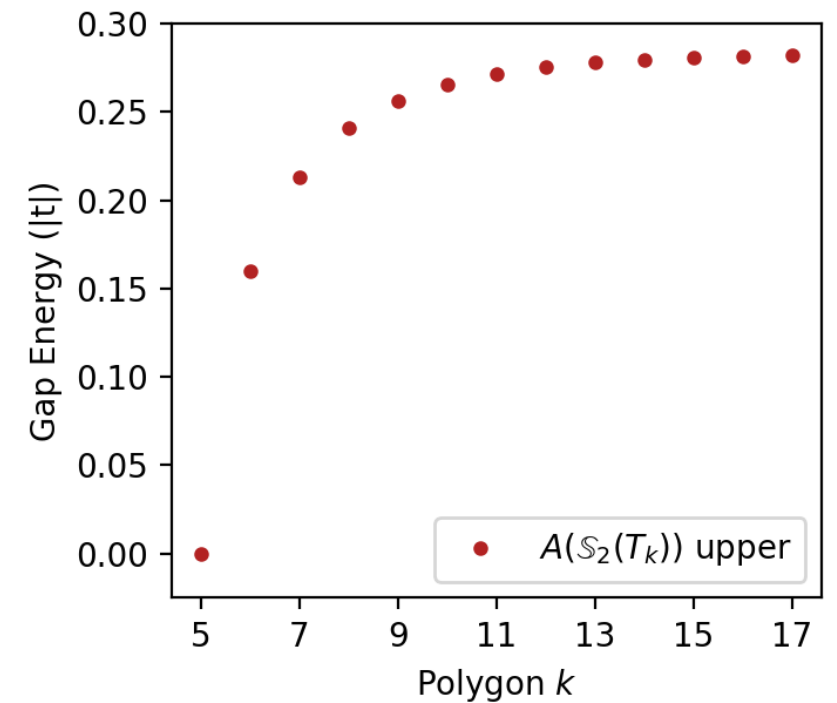
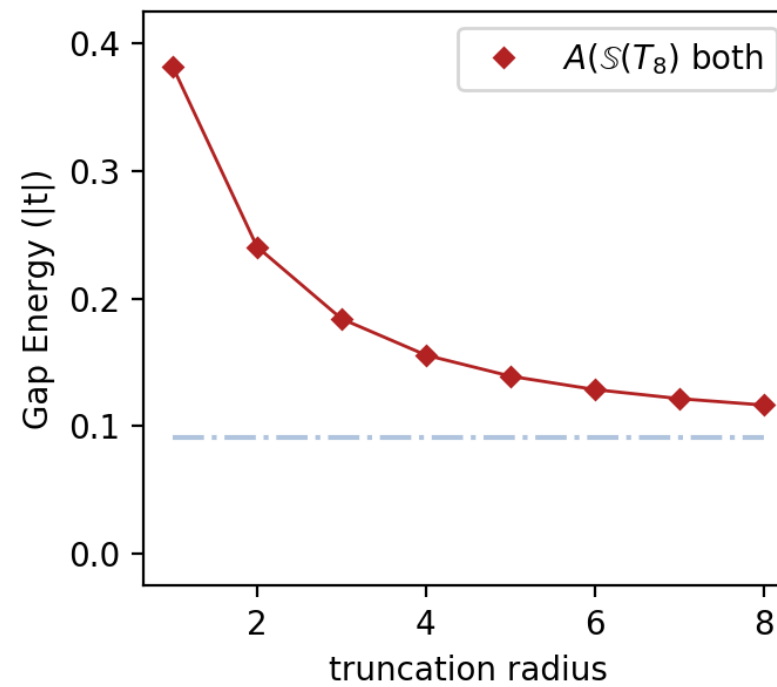
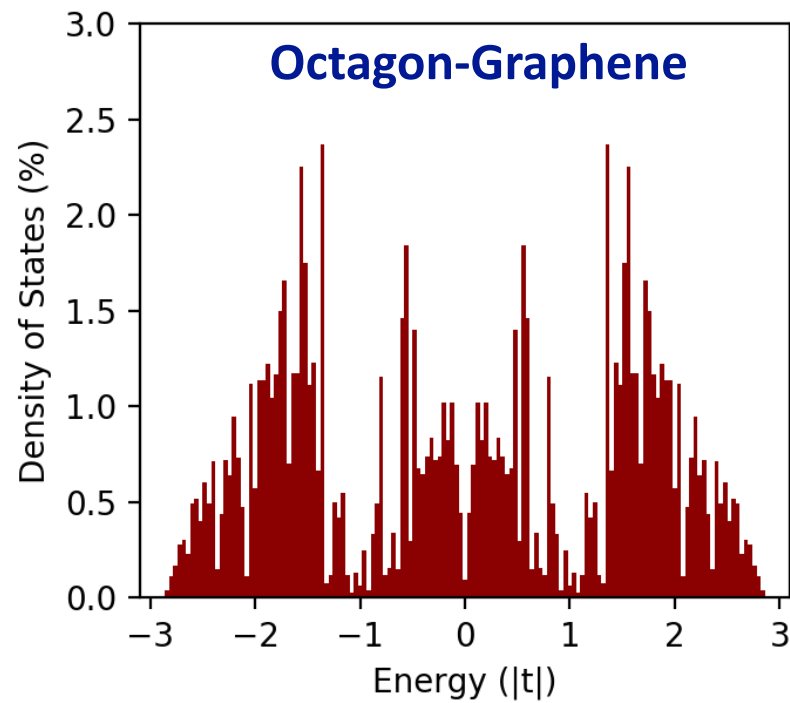
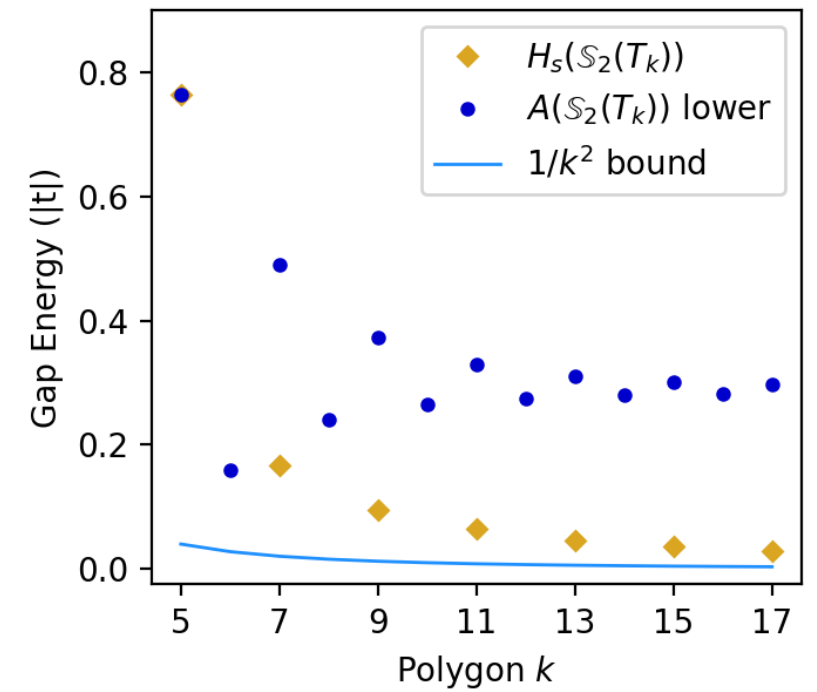
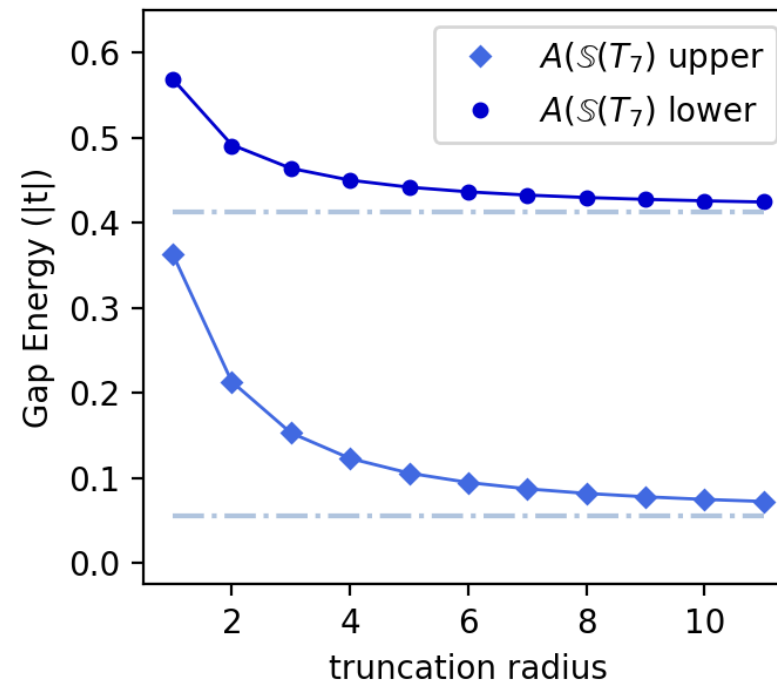
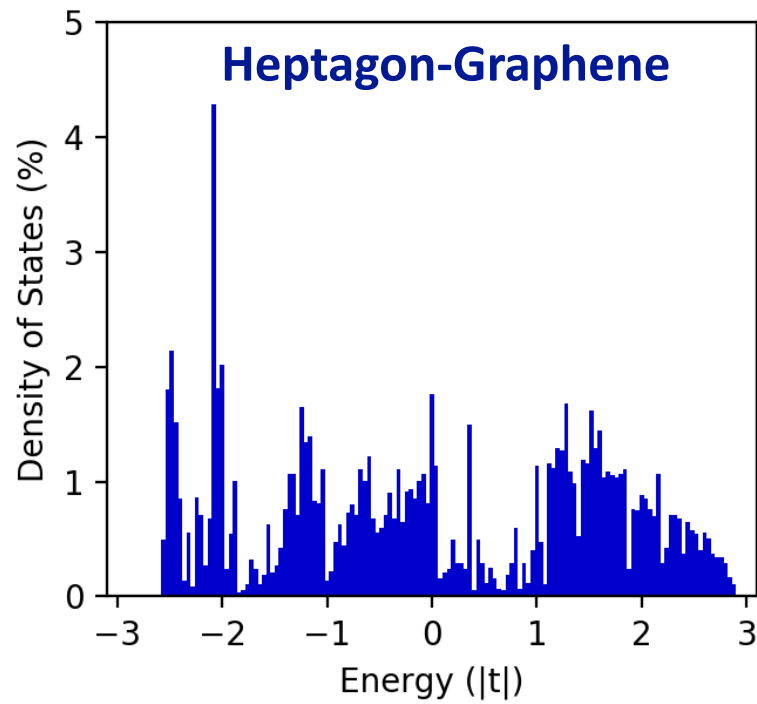


Gaussian Curvature

$$K = -\frac{1}{R^2}$$

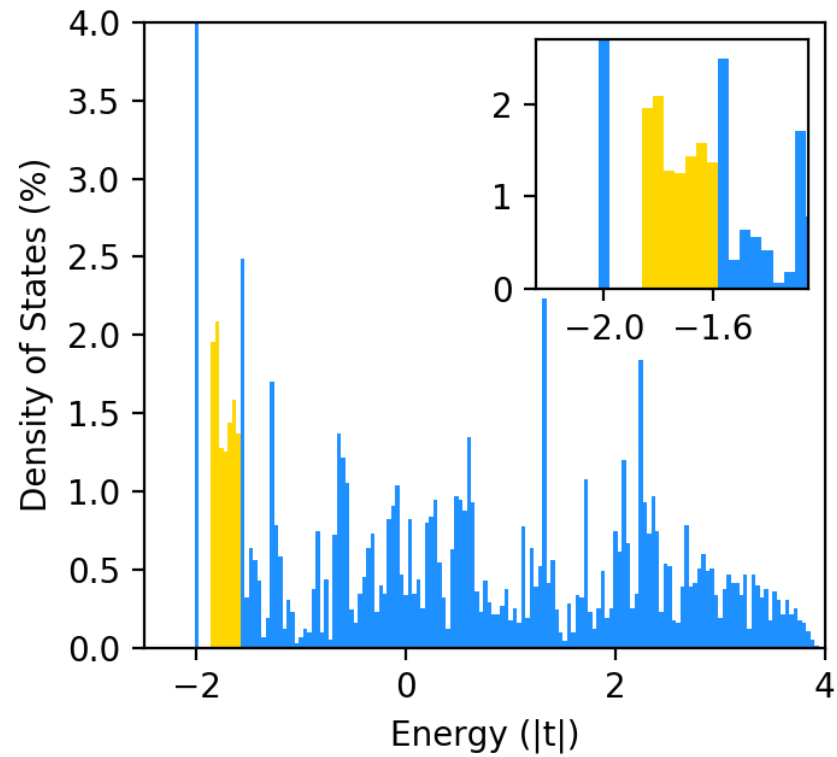
Tiling Polygon (n)	Lattice Constant	Medial Lattice Constant
7	0.566	0.492
8	0.727	0.633
9	0.819	0.714
10	0.879	0.767
11	0.921	0.804
12	0.952	0.831

Hyperbolic Numerics

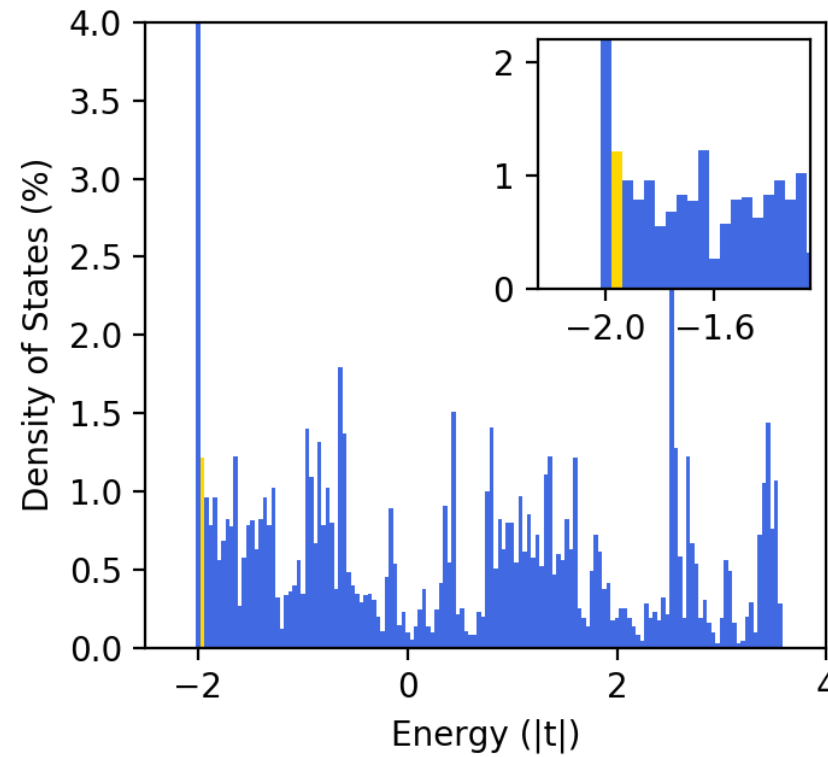


Hyperbolic Numerics

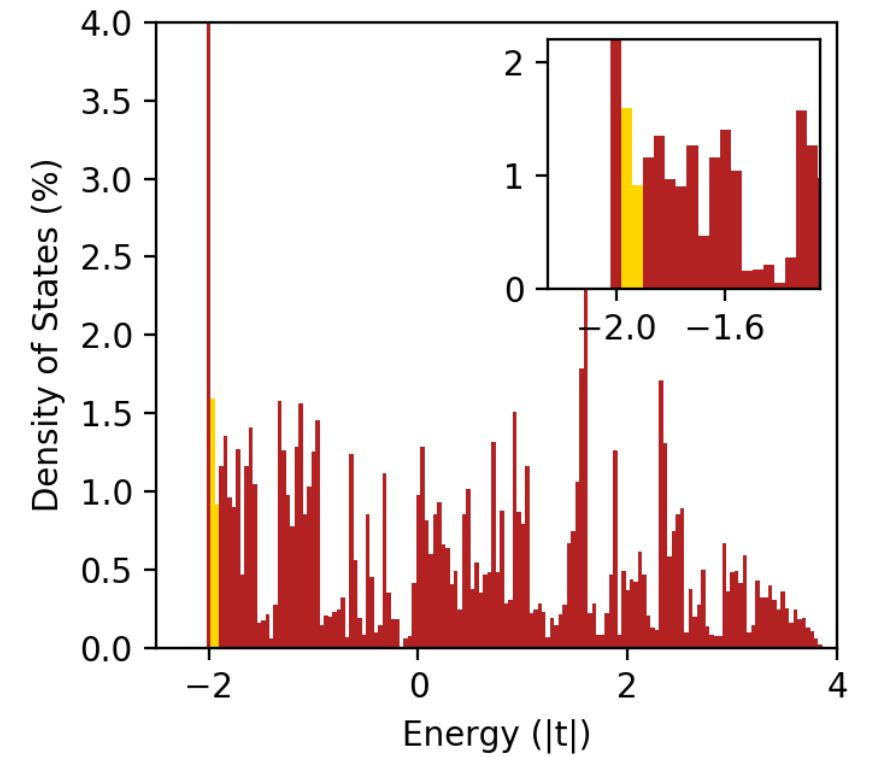
Heptagon-Kagome



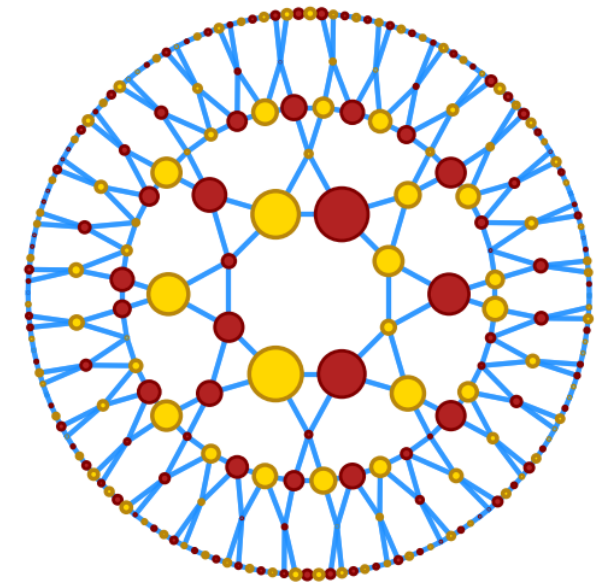
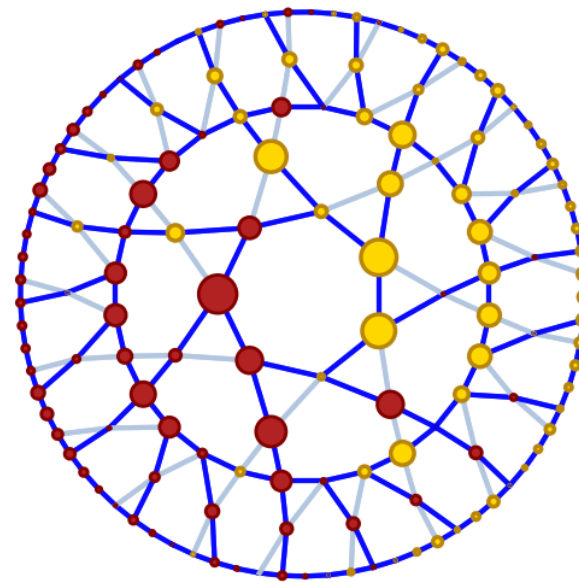
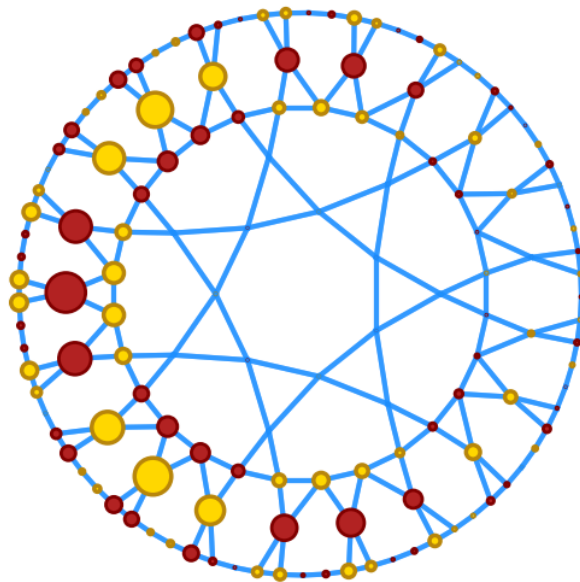
Heptagon-Kagome (HW)



Octagon-Kagome



First "mid-gap" state



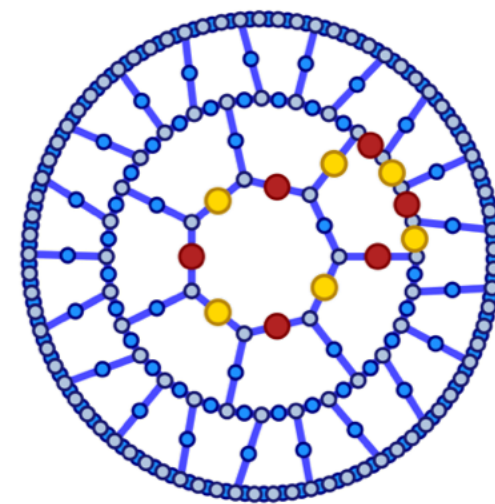
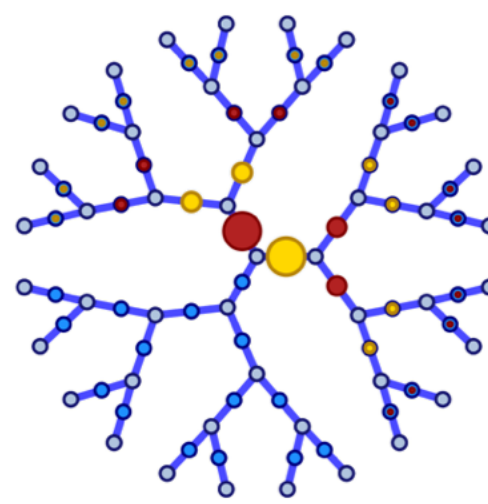
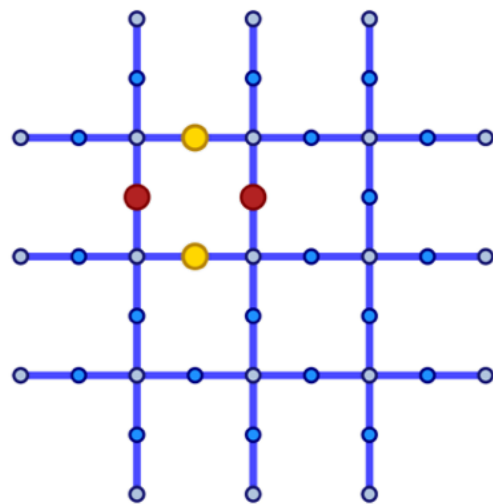
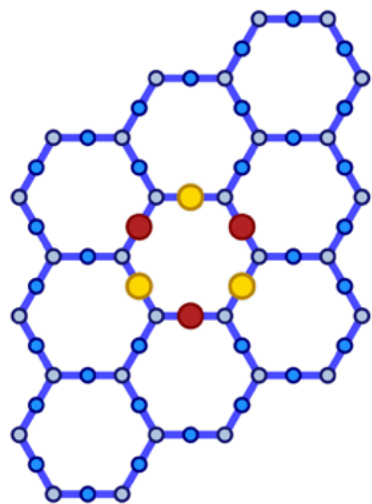
Empty rounded rectangular box at the top of the page.

Large empty rounded rectangular box occupying the majority of the page.

Subdivision Graphs: Flat Bands at 0

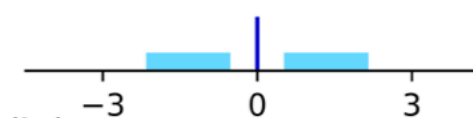
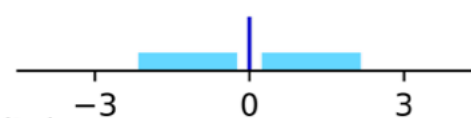
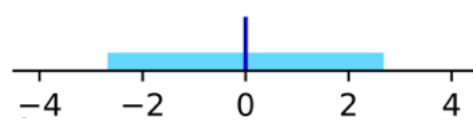
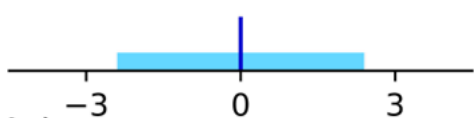
Subdivision

$S(X)$



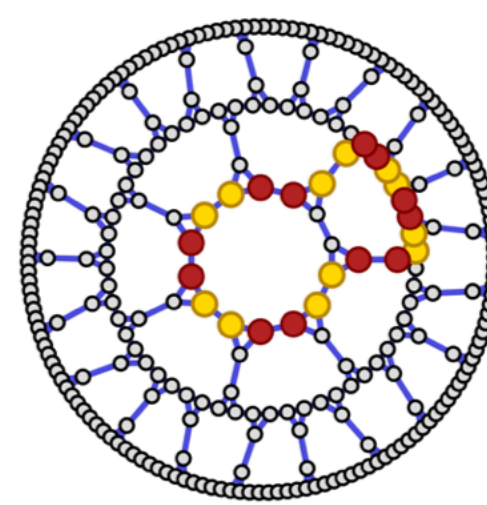
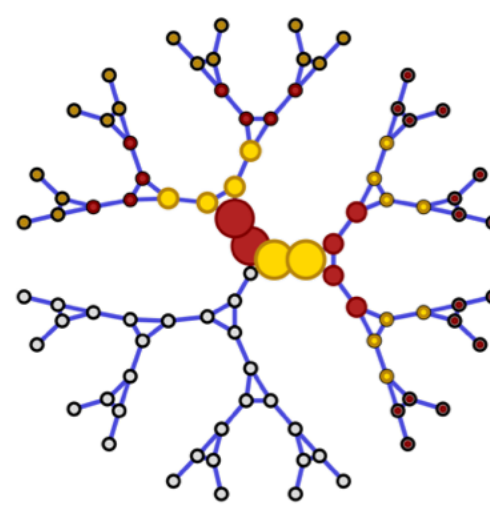
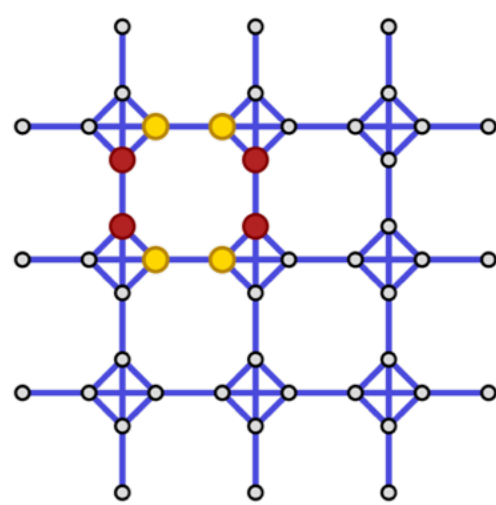
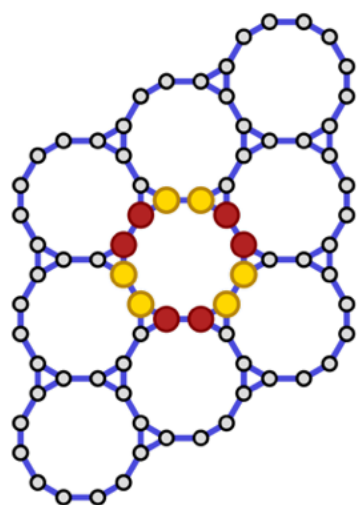
DOS

$S(X)$



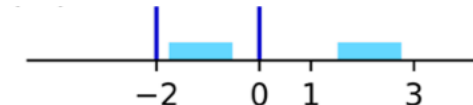
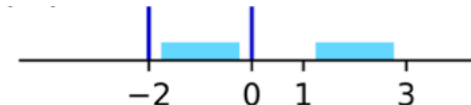
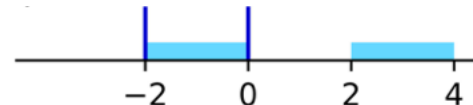
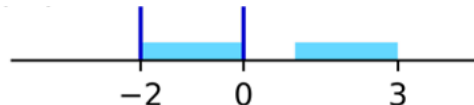
Line Graph

$L(S(X))$



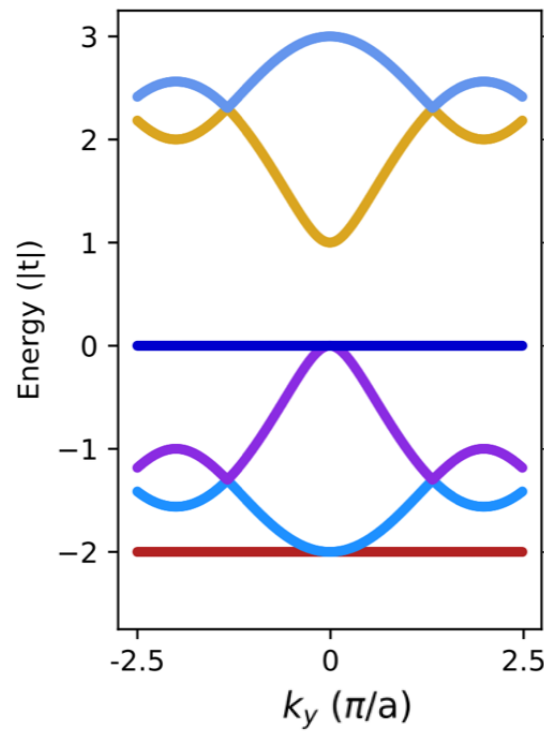
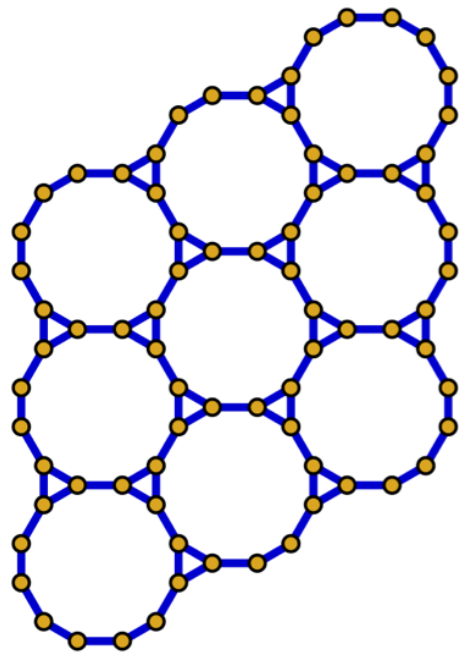
DOS

$L(S(X))$



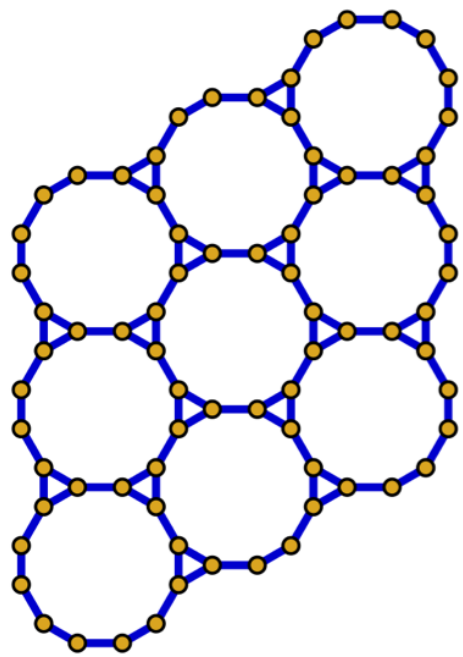
Subdivision Graphs and Optimally Gapped Flat Bands

X



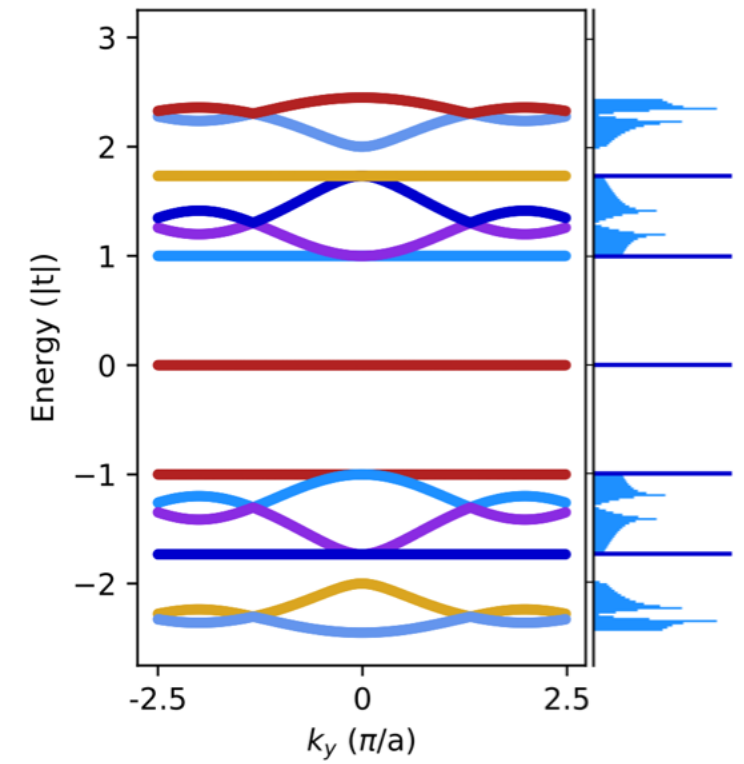
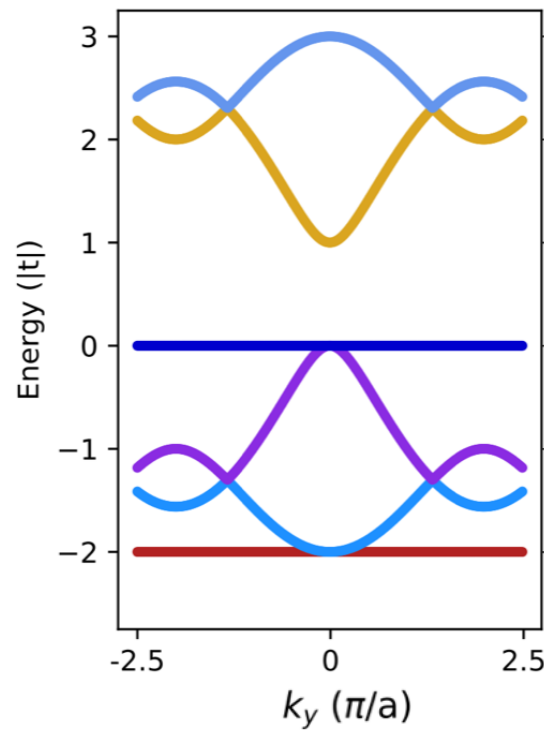
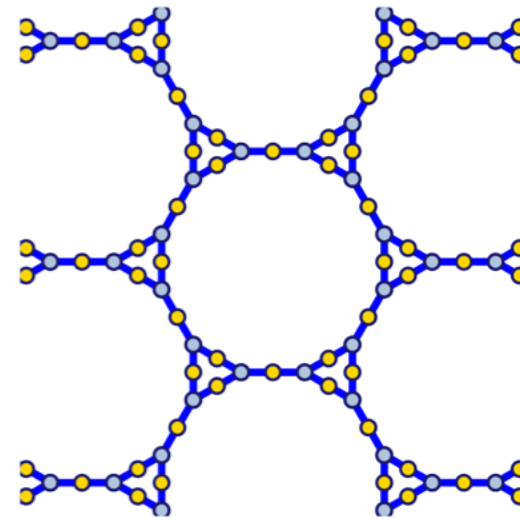
Subdivision Graphs and Optimally Gapped Flat Bands

X



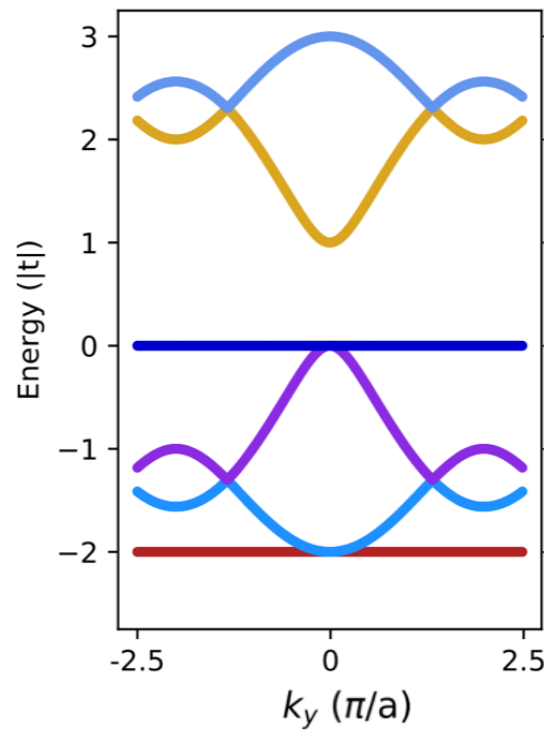
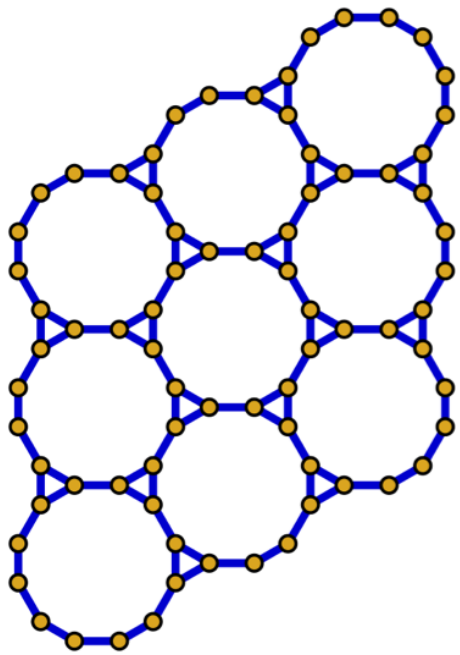
$$E_{\mathbb{S}(X)} = \begin{cases} \pm\sqrt{E_X + 3} \\ 0 \end{cases}$$

$\mathbb{S}(X)$



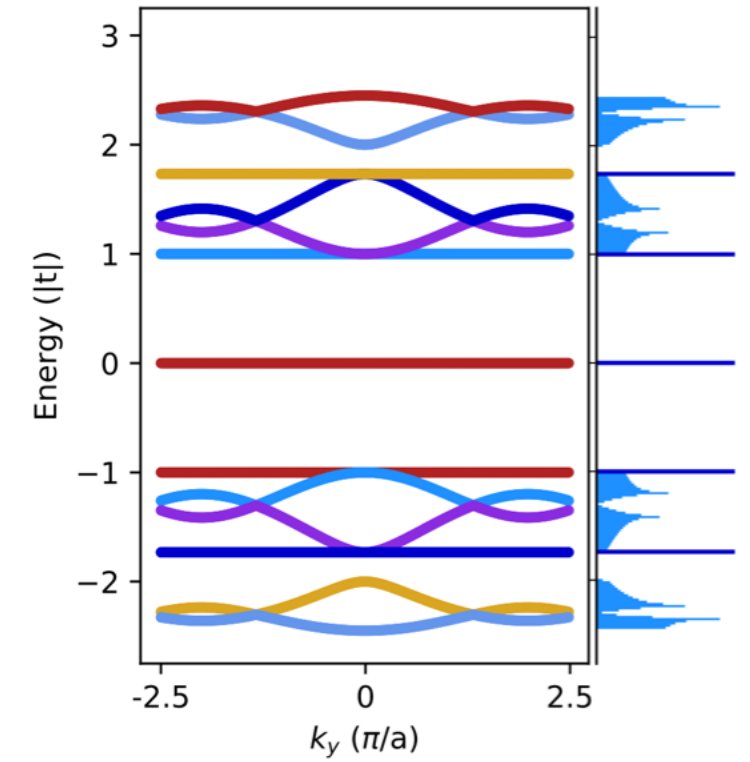
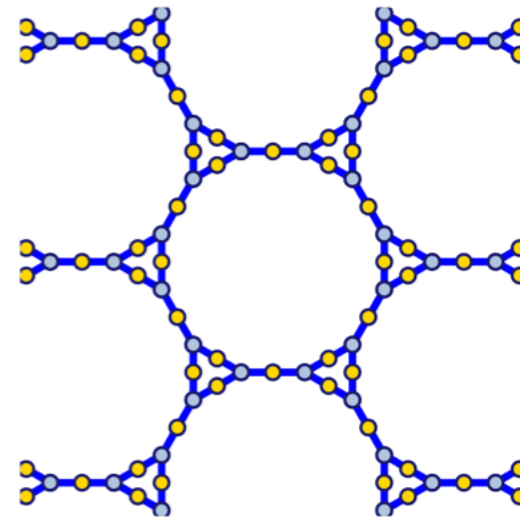
Subdivision Graphs and Optimally Gapped Flat Bands

X

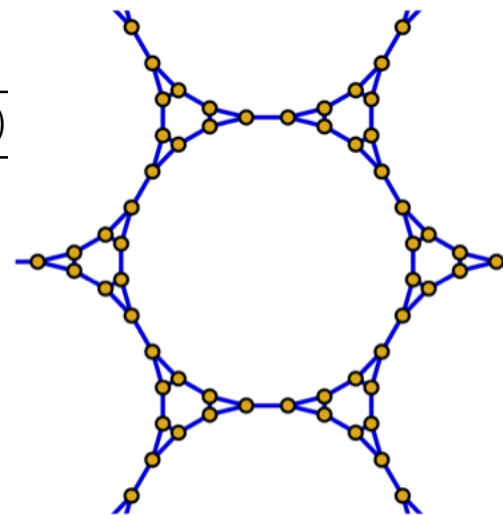


$$E_{\mathcal{S}(X)} = \begin{cases} \pm\sqrt{E_X + 3} \\ 0 \end{cases}$$

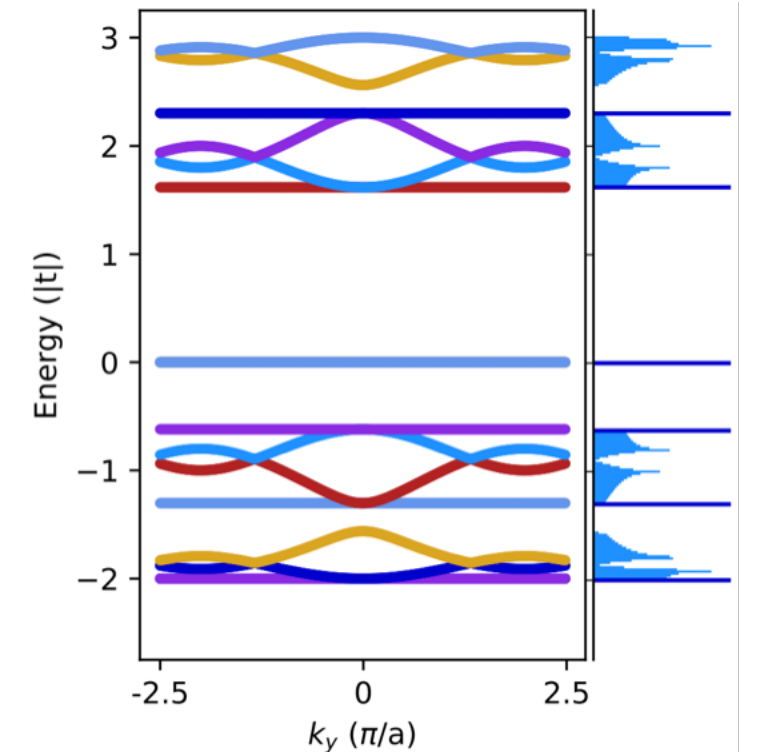
$\mathcal{S}(X)$



$L(\mathcal{S}(X))$

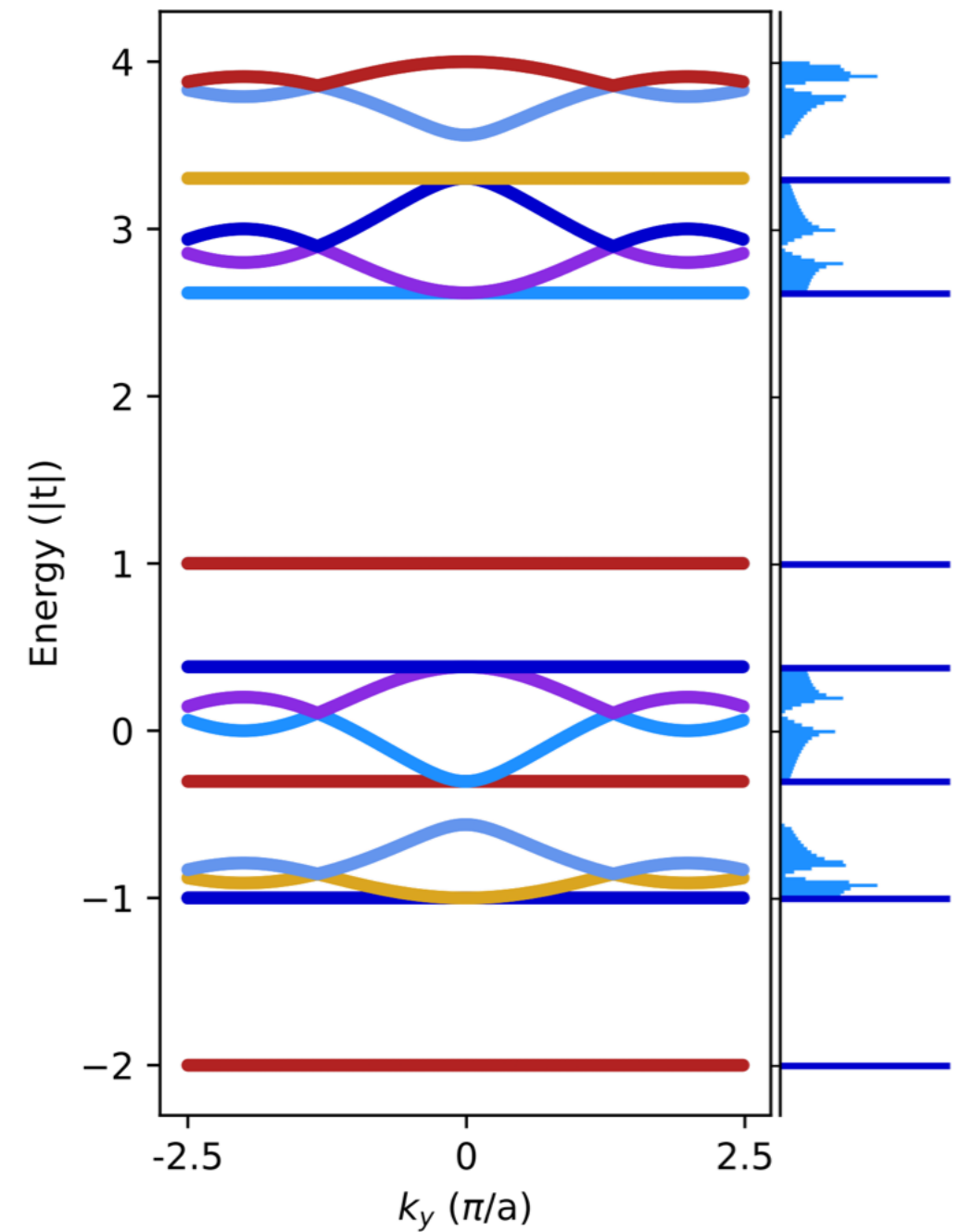
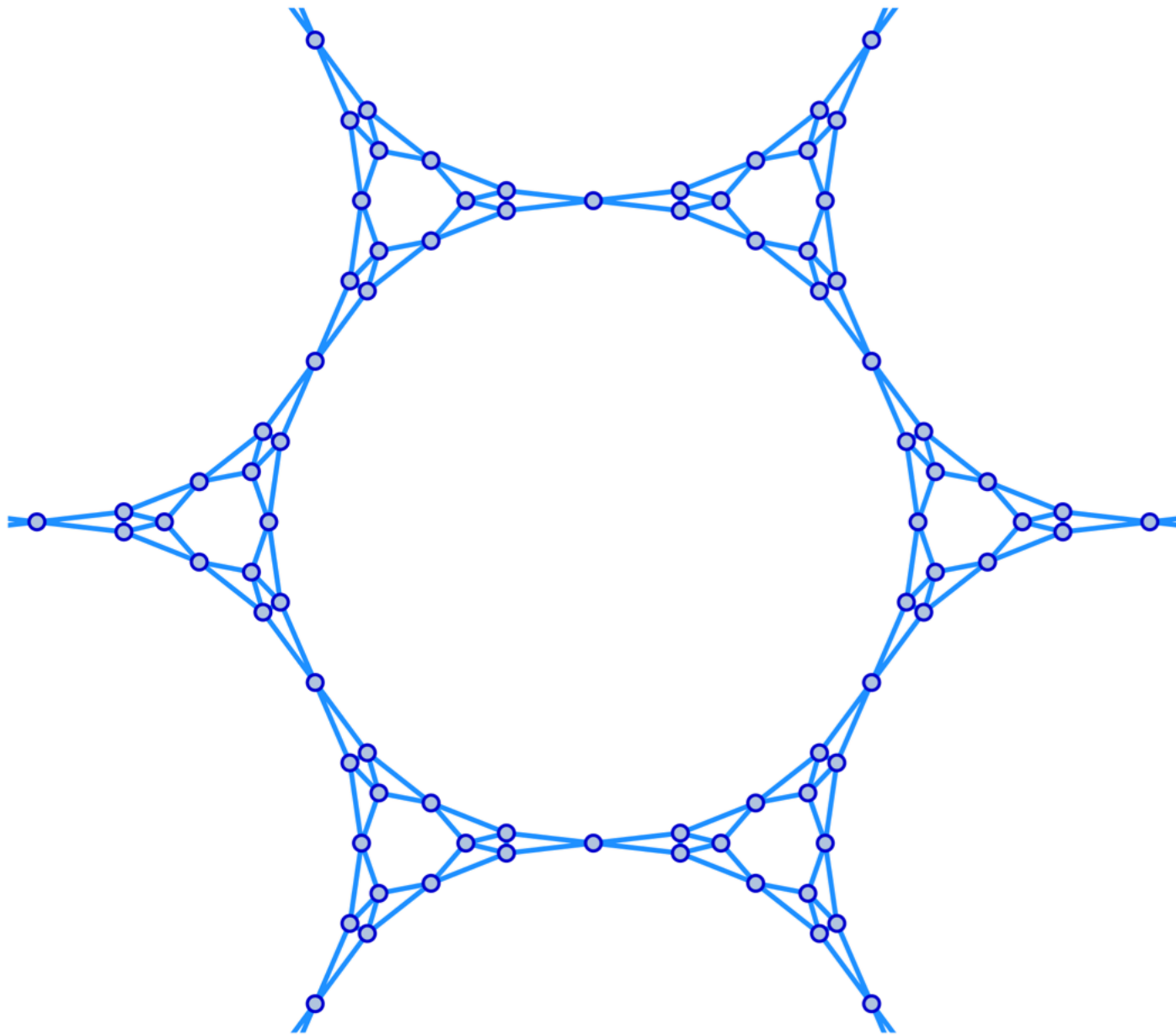


$$E_{L(\mathcal{S}(X))} = \begin{cases} \frac{1 \pm \sqrt{1 + 4(E_X + 3)}}{2} \\ 0 \\ -2 \end{cases}$$



Subdivision Graphs and Optimally Gapped Flat Bands

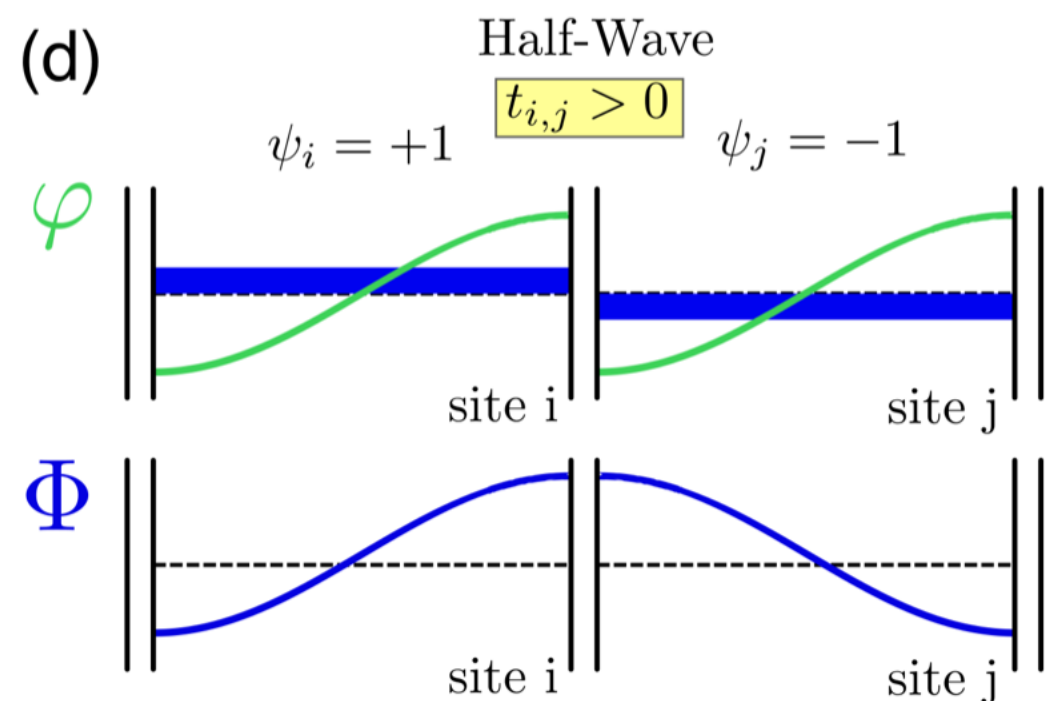
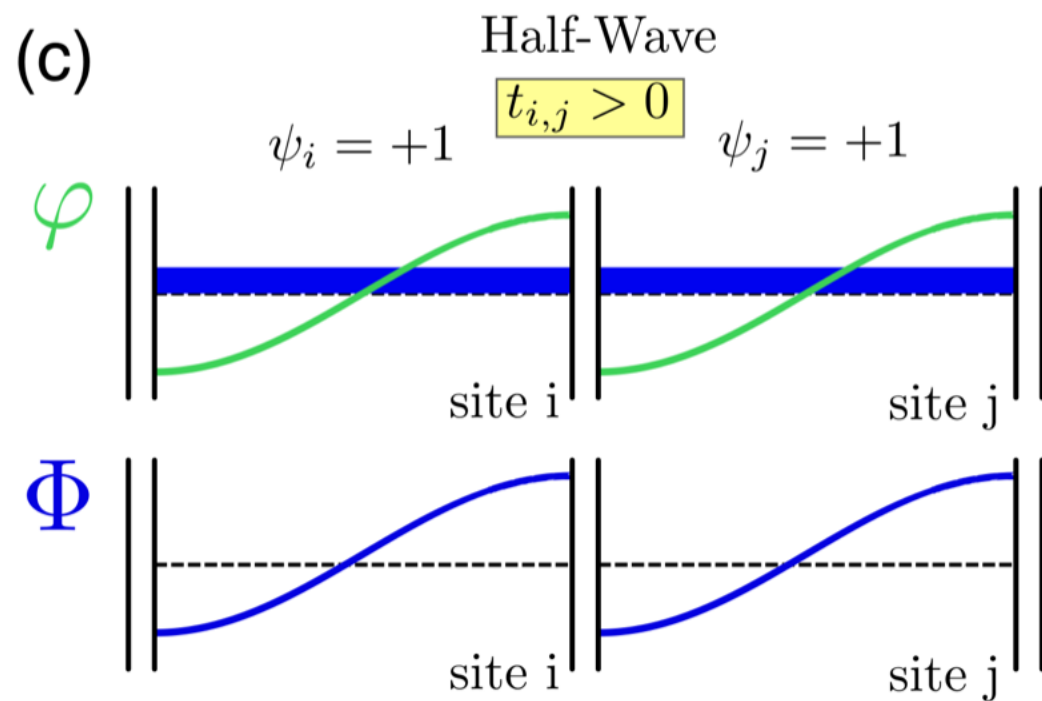
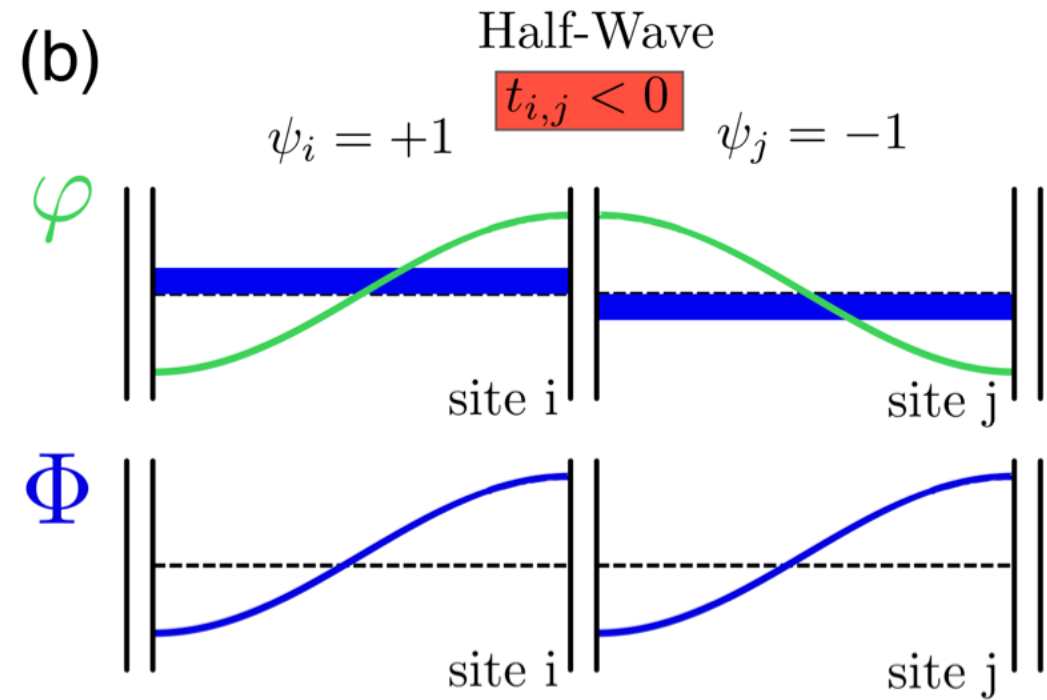
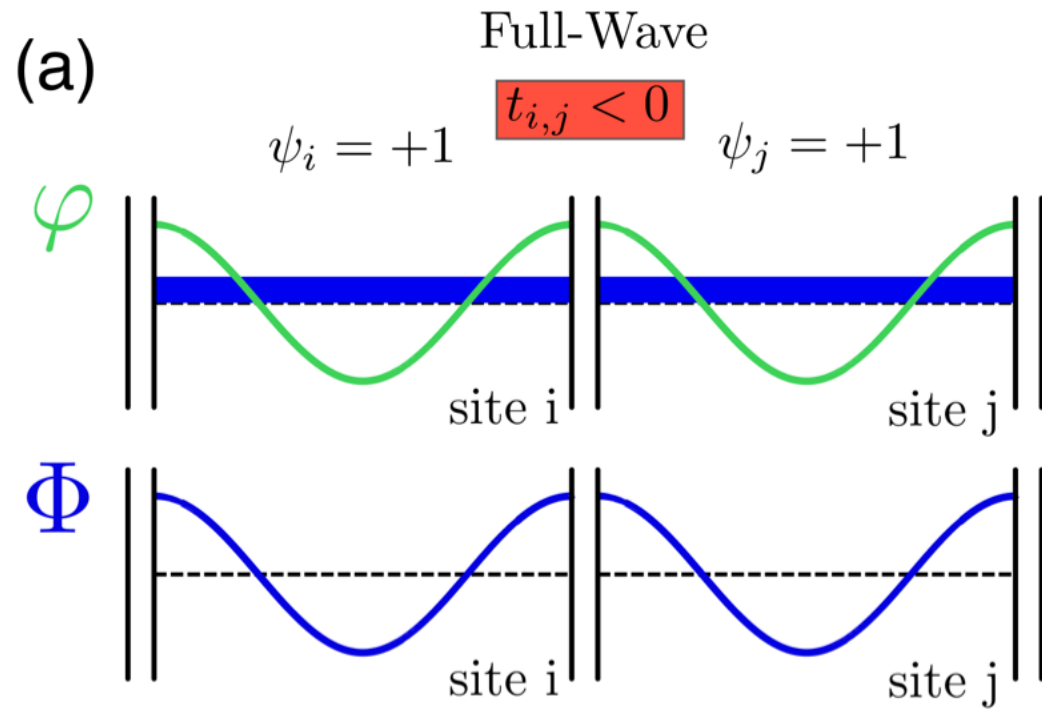
$$L(L(\mathcal{S}(X)))$$



Empty rounded rectangular box at the top of the page.

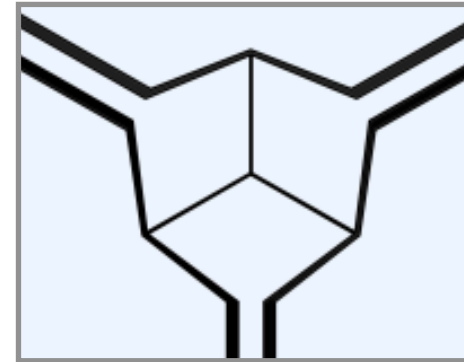
Large empty rounded rectangular box occupying the majority of the page.

Tight Binding



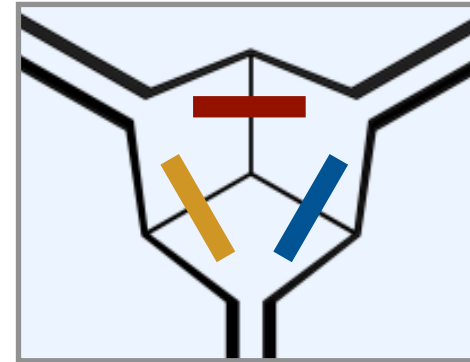
S-Wave and P-Wave On-Site Wave Functions

$$\mathcal{H} = \sum_{\text{coupling capacitors}} \omega C_c \Phi^+ \Phi^-$$



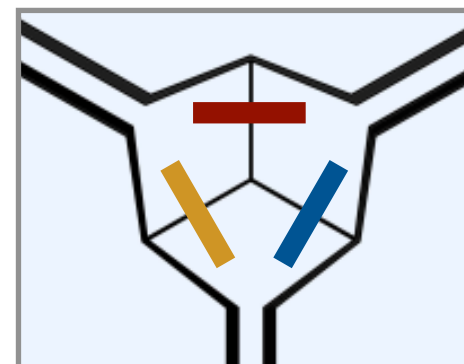
S-Wave and P-Wave On-Site Wave Functions

$$\mathcal{H} = \sum_{\text{coupling capacitors}} \omega C_c \Phi^+ \Phi^-$$

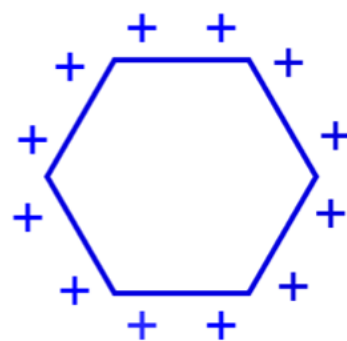
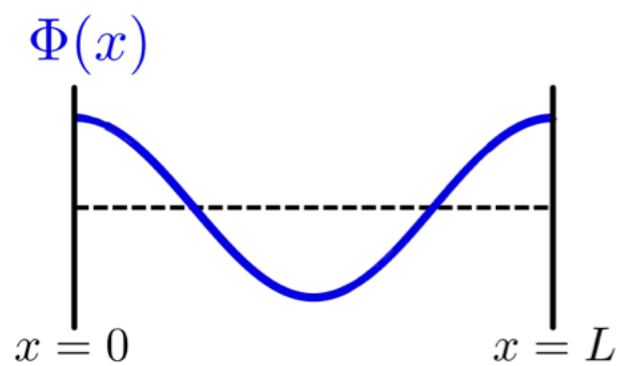


S-Wave and P-Wave On-Site Wave Functions

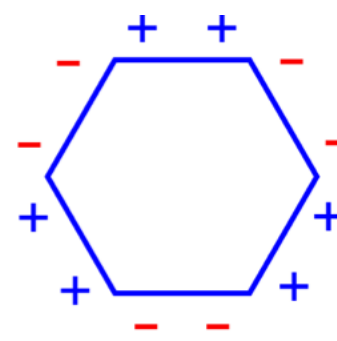
$$\mathcal{H} = \sum_{\text{coupling capacitors}} \omega C_c \Phi^+ \Phi^-$$



Full-wave



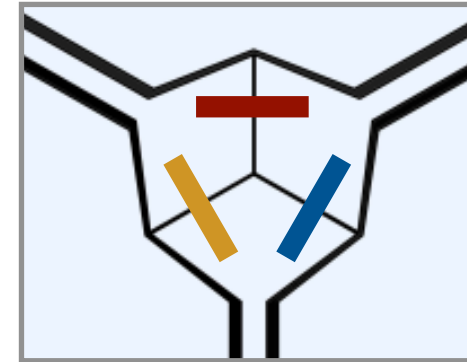
$$E = +2$$



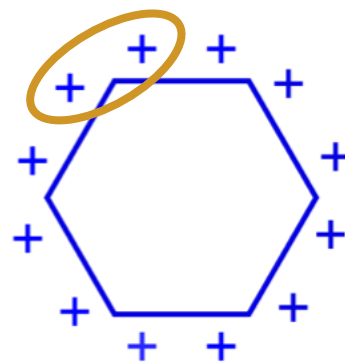
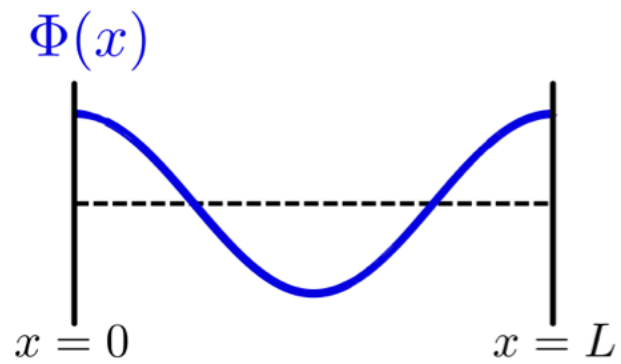
$$E = -2$$

S-Wave and P-Wave On-Site Wave Functions

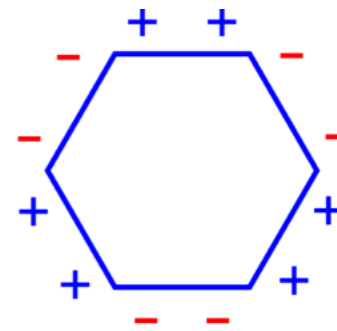
$$\mathcal{H} = \sum_{\text{coupling capacitors}} \omega C_c \Phi^+ \Phi^-$$



Full-wave



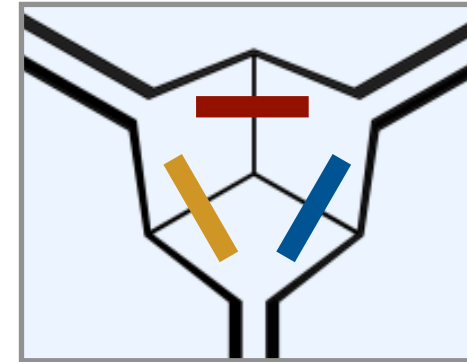
$$E = +2$$



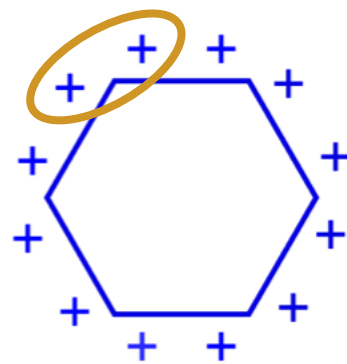
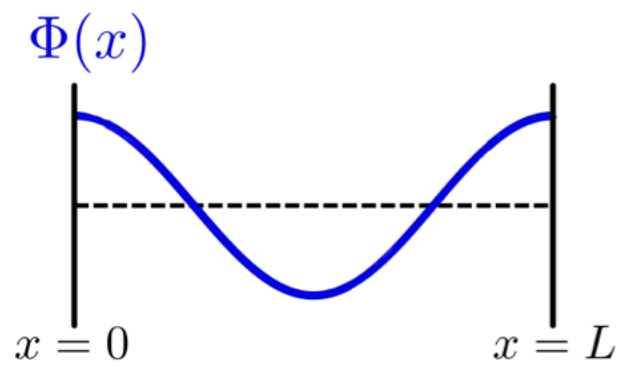
$$E = -2$$

S-Wave and P-Wave On-Site Wave Functions

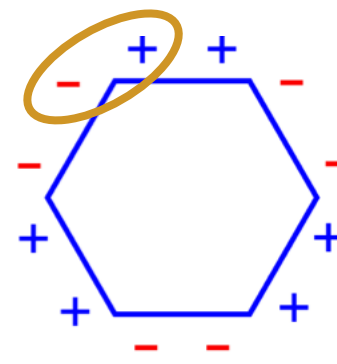
$$\mathcal{H} = \sum_{\text{coupling capacitors}} \omega C_c \Phi^+ \Phi^-$$



Full-wave



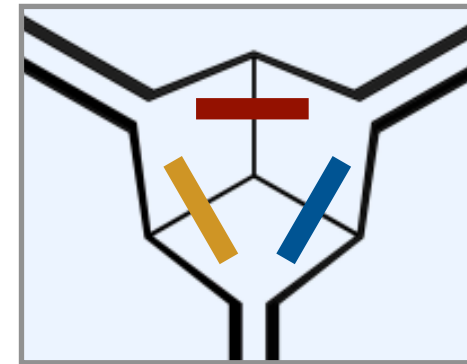
$$E = +2$$



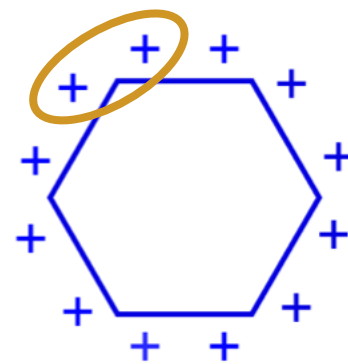
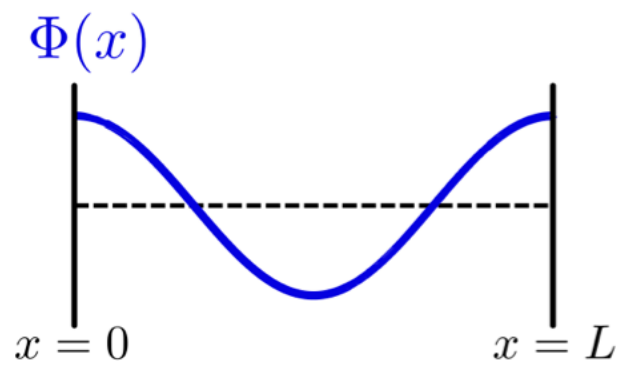
$$E = -2$$

S-Wave and P-Wave On-Site Wave Functions

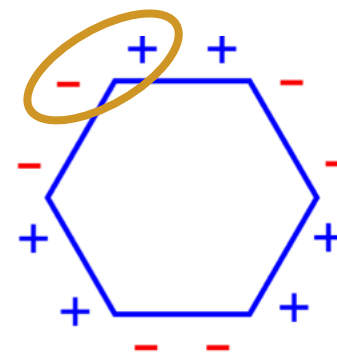
$$\mathcal{H} = \sum_{\text{coupling capacitors}} \omega C_c \Phi^+ \Phi^-$$



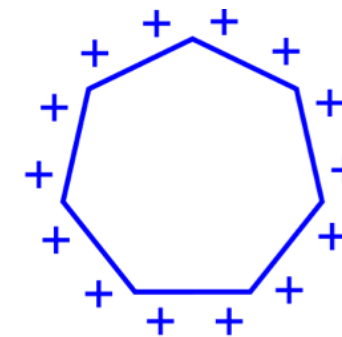
Full-wave



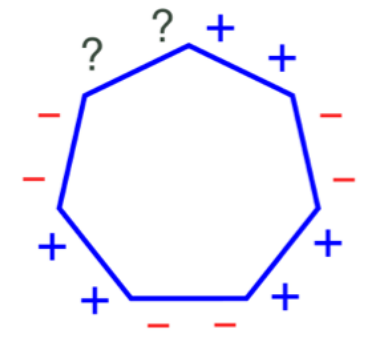
$$E = +2$$



$$E = -2$$



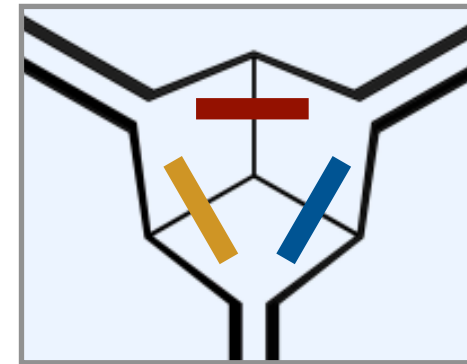
$$E = +2$$



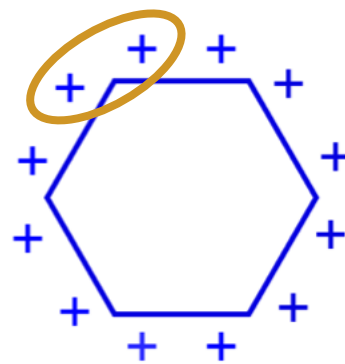
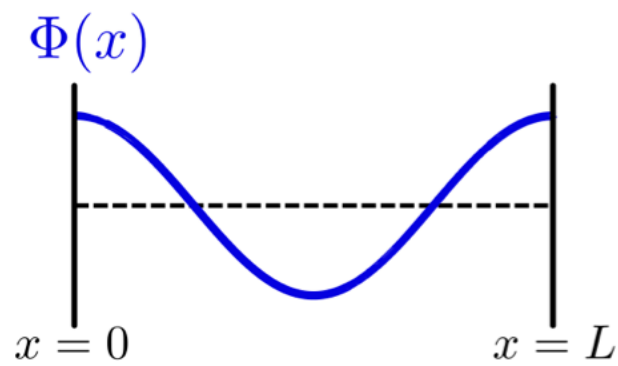
$$E = -2$$

S-Wave and P-Wave On-Site Wave Functions

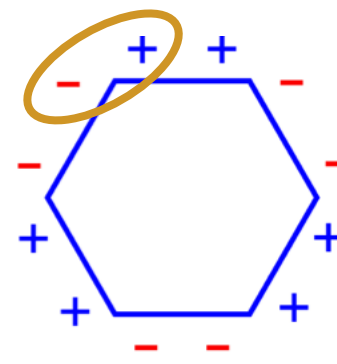
$$\mathcal{H} = \sum_{\text{coupling capacitors}} \omega C_c \Phi^+ \Phi^-$$



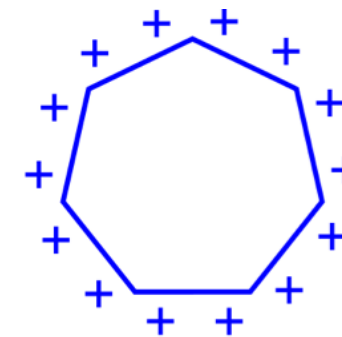
Full-wave



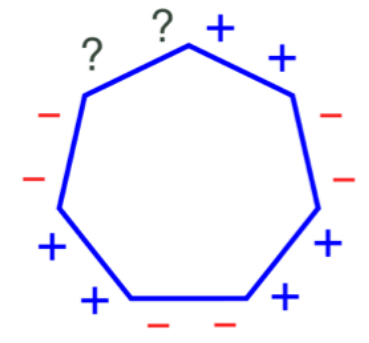
$$E = +2$$



$$E = -2$$

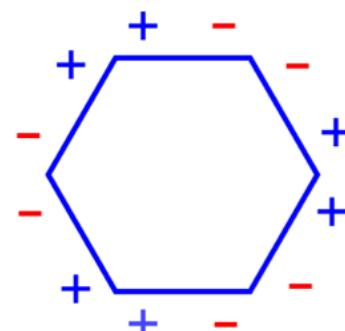
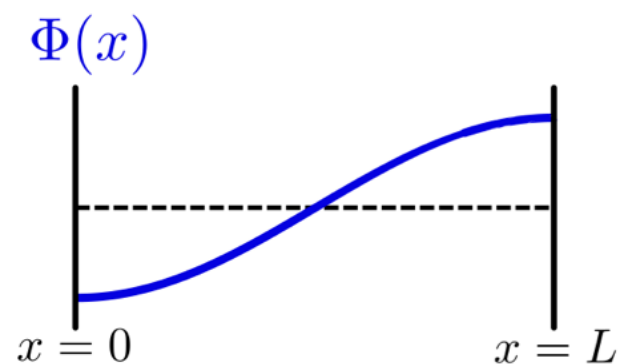


$$E = +2$$

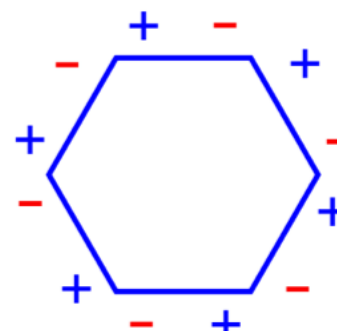


$$E = -2$$

Half-wave



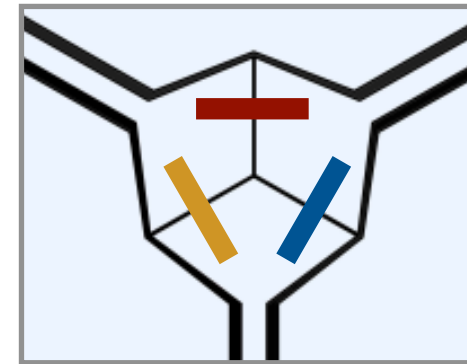
$$E = +2$$



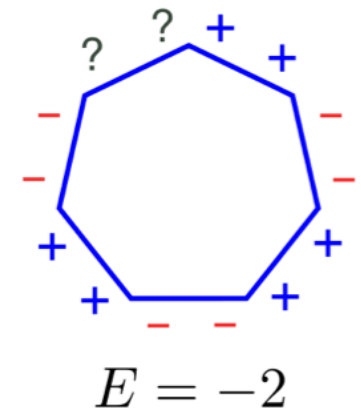
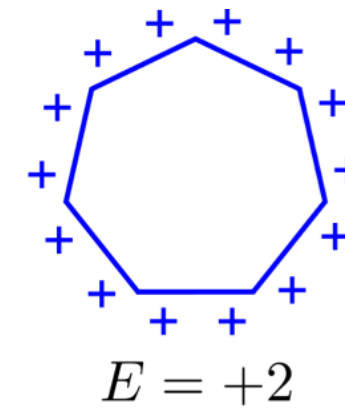
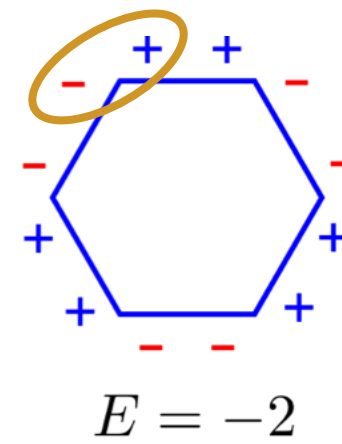
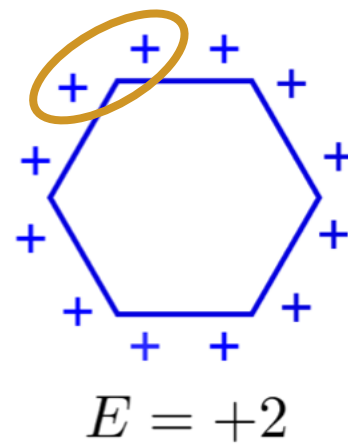
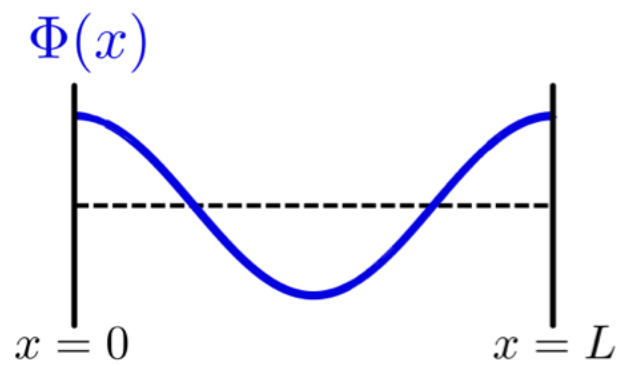
$$E = -2$$

S-Wave and P-Wave On-Site Wave Functions

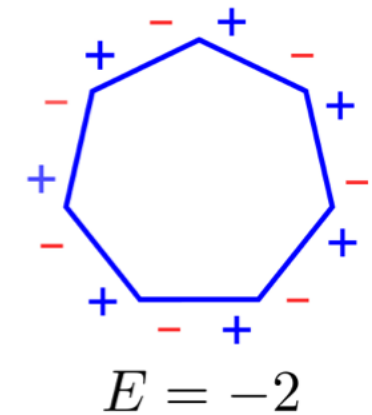
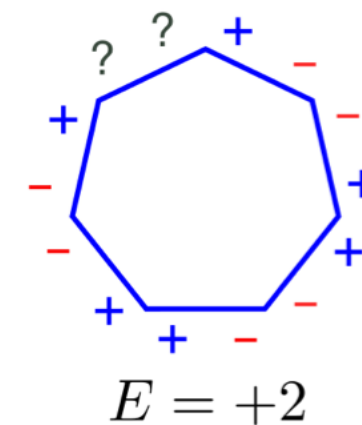
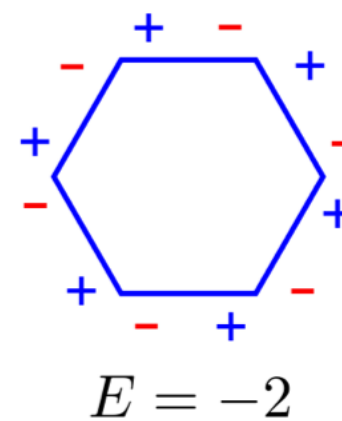
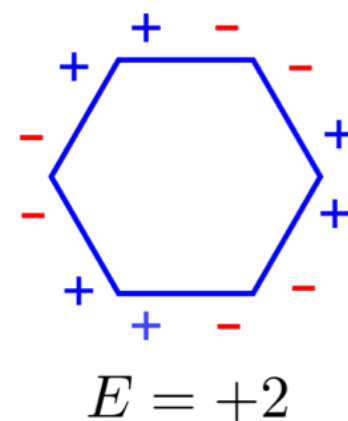
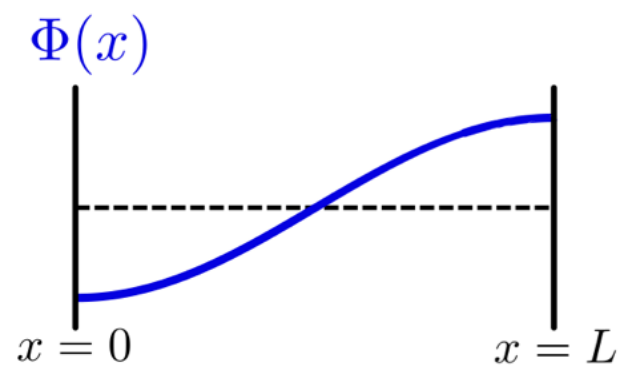
$$\mathcal{H} = \sum_{\text{coupling capacitors}} \omega C_c \Phi^+ \Phi^-$$



Full-wave



Half-wave



Half-Wave Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Half-Wave Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$
- Mixed positive and negative hopping

$$\bar{H}_a(X) \neq H_{L(X)}$$

Half-Wave Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$
- Mixed positive and negative hopping

$$\bar{H}_a(X) \neq H_{L(X)}$$

Incidence Operator

- From X to $L(X)$

$$N : \ell^2(X) \rightarrow \ell^2(L(X))$$

Half-Wave Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$
- Mixed positive and negative hopping

$$\bar{H}_a(X) \neq H_{L(X)}$$

Incidence Operator

- From X to $L(X)$

$$N : \ell^2(X) \rightarrow \ell^2(L(X))$$

$$N(v, e) = \begin{cases} 1, & \text{if } e^+ = v, \\ -1 & \text{if } e^- = v, \\ 0 & \text{otherwise.} \end{cases}$$

Half-Wave Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$
- Mixed positive and negative hopping

$$\bar{H}_a(X) \neq H_{L(X)}$$

Incidence Operator

- From X to $L(X)$

$$N : \ell^2(X) \rightarrow \ell^2(L(X))$$

$$N(v, e) = \begin{cases} 1, & \text{if } e^+ = v, \\ -1 & \text{if } e^- = v, \\ 0 & \text{otherwise.} \end{cases}$$

$$N^t N = D_X - H_X$$

$$N N^t = 2I + \bar{H}_a(X)$$

Half-Wave Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$
- Mixed positive and negative hopping

$$\bar{H}_a(X) \neq H_{L(X)}$$

Incidence Operator

- From X to $L(X)$

$$N : \ell^2(X) \rightarrow \ell^2(L(X))$$

$$N(v, e) = \begin{cases} 1, & \text{if } e^+ = v, \\ -1 & \text{if } e^- = v, \\ 0 & \text{otherwise.} \end{cases}$$

$$N^t N = D_X - H_X$$

$$N N^t = 2I + \bar{H}_a(X)$$

$$D_X - H_X \simeq 2I + \bar{H}_a(X)$$

Half-Wave Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$
- Mixed positive and negative hopping

$$\bar{H}_a(X) \neq H_{L(X)}$$

Incidence Operator

- From X to $L(X)$

$$N : \ell^2(X) \rightarrow \ell^2(L(X))$$

$$N(v, e) = \begin{cases} 1, & \text{if } e^+ = v, \\ -1 & \text{if } e^- = v, \\ 0 & \text{otherwise.} \end{cases}$$

$$N^t N = D_X - H_X$$

$$N N^t = 2I + \bar{H}_a(X)$$

$$D_X - H_X \simeq 2I + \bar{H}_a(X)$$

$$E_{\bar{H}_a} = \begin{cases} d - 2 - E_{H_X} \\ -2 \end{cases}$$

Half-Wave Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$
- Mixed positive and negative hopping

$$\bar{H}_a(X) \neq H_{L(X)}$$

Incidence Operator

- From X to $L(X)$

$$N : \ell^2(X) \rightarrow \ell^2(L(X))$$

$$N(v, e) = \begin{cases} 1, & \text{if } e^+ = v, \\ -1 & \text{if } e^- = v, \\ 0 & \text{otherwise.} \end{cases}$$

$$N^t N = D_X - H_X$$

$$N N^t = 2I + \bar{H}_a(X)$$

$$D_X - H_X \simeq 2I + \bar{H}_a(X)$$

$$E_{\bar{H}_a} = \begin{cases} d - 2 - E_{H_X} \\ -2 \end{cases}$$

- Identical on bipartite graphs

Half-Wave Band Structure Correspondence

Layout Tight-Binding Hamiltonian

- Bounded self-adjoint operator on X

$$H_X$$

Effective Hamiltonian

- Bounded self-adjoint operator on $L(X)$
- Mixed positive and negative hopping

$$\bar{H}_a(X) \neq H_{L(X)}$$

Incidence Operator

- From X to $L(X)$

$$N : \ell^2(X) \rightarrow \ell^2(L(X))$$

$$N(v, e) = \begin{cases} 1, & \text{if } e^+ = v, \\ -1 & \text{if } e^- = v, \\ 0 & \text{otherwise.} \end{cases}$$

$$N^t N = D_X - H_X$$

$$N N^t = 2I + \bar{H}_a(X)$$

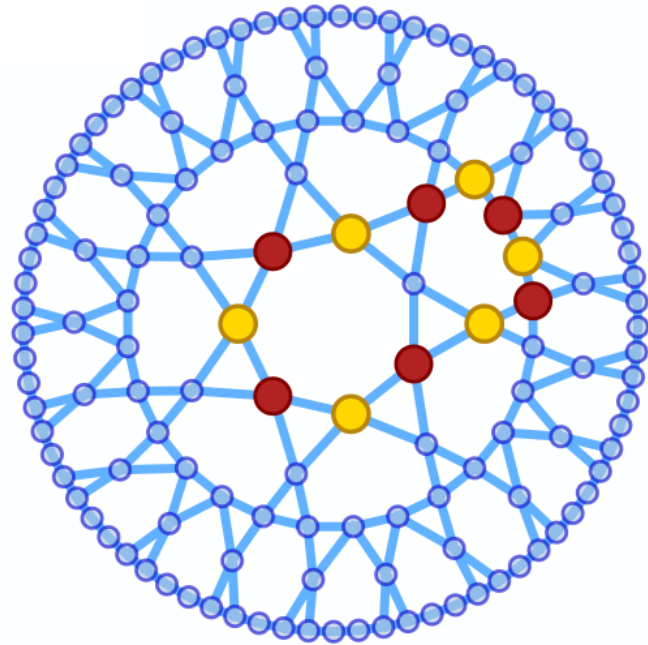
$$D_X - H_X \simeq 2I + \bar{H}_a(X)$$

$$E_{\bar{H}_a} = \begin{cases} d - 2 - E_{H_X} \\ -2 \end{cases}$$

- Identical on bipartite graphs
- Inverted otherwise

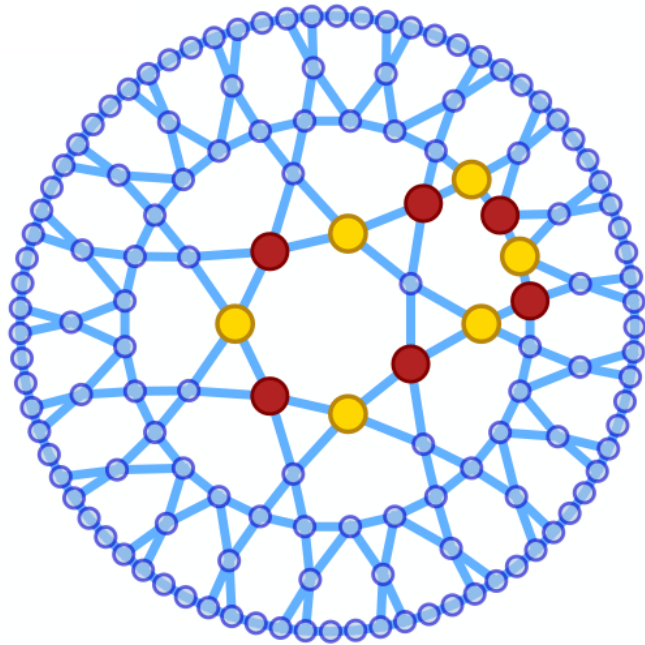
Full-Wave v Half-Wave Flat Band States

FW

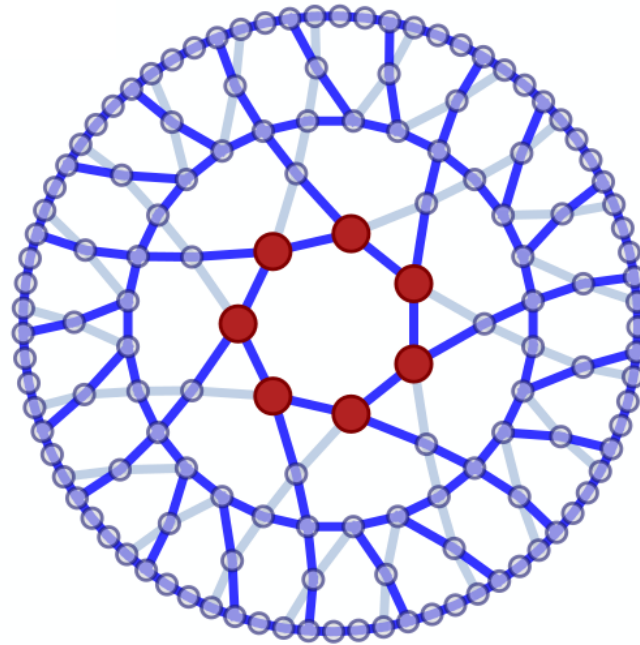


Full-Wave v Half-Wave Flat Band States

FW

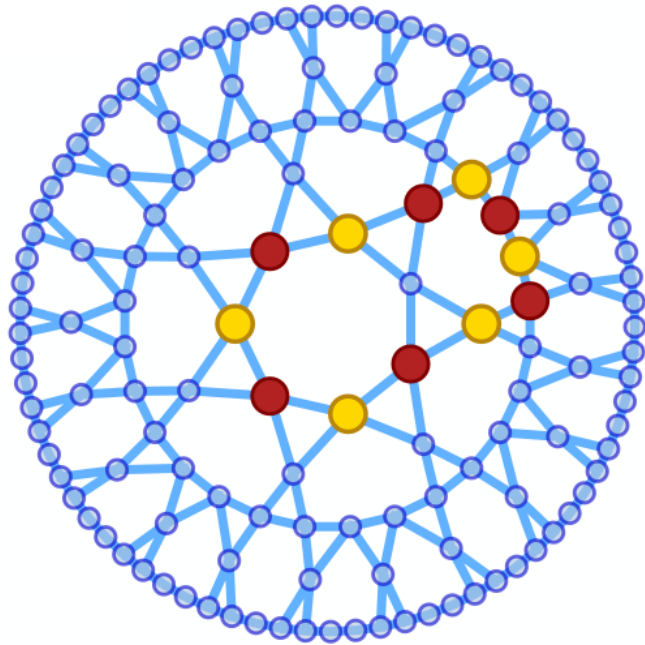


HW

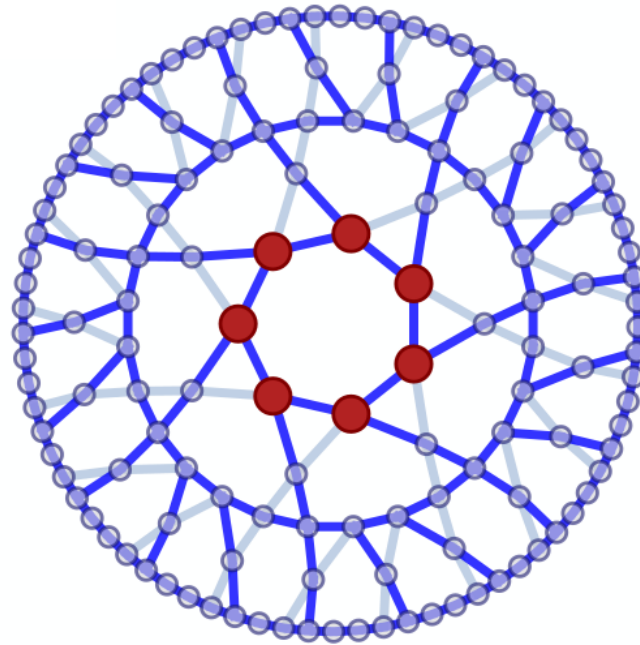


Full-Wave v Half-Wave Flat Band States

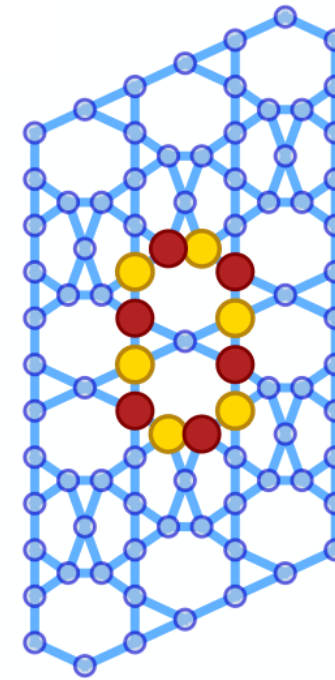
FW



HW

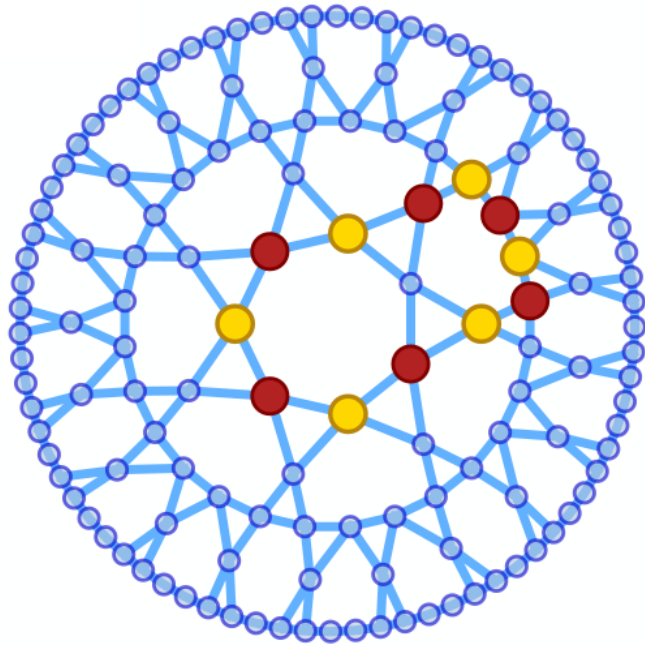


FW

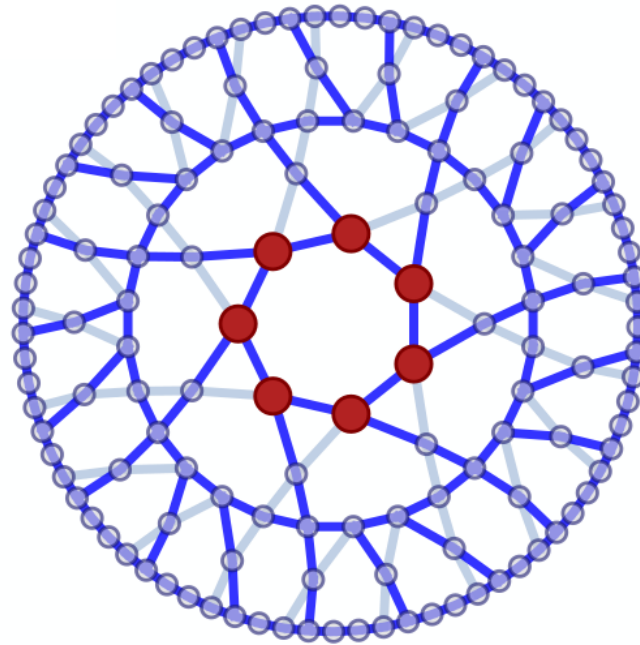


Full-Wave v Half-Wave Flat Band States

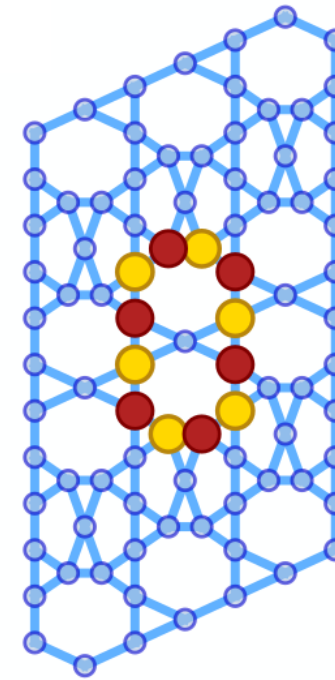
FW



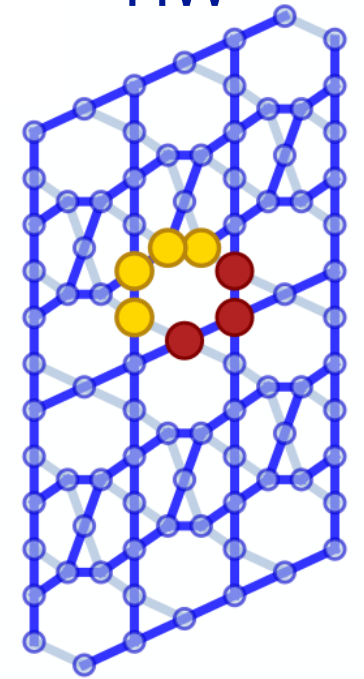
HW



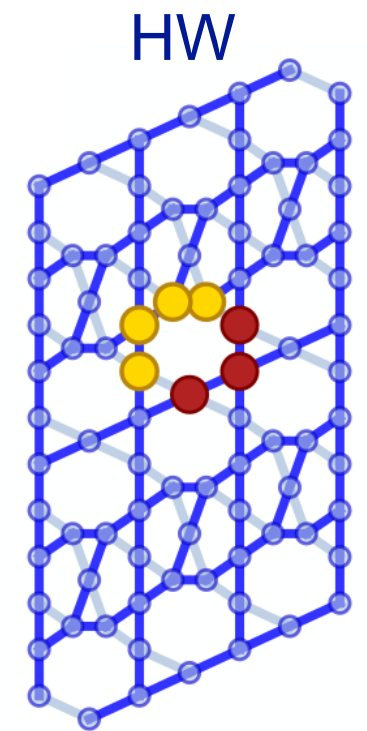
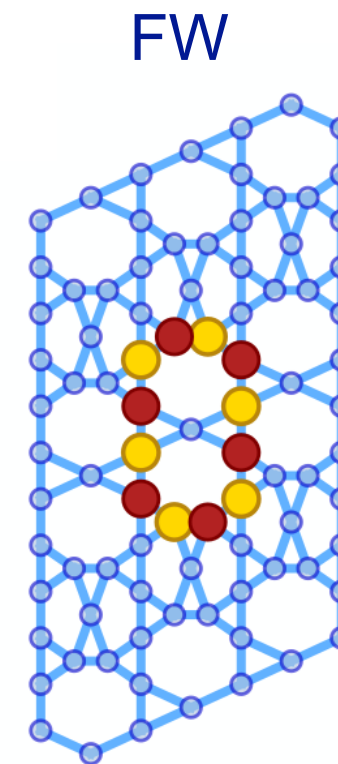
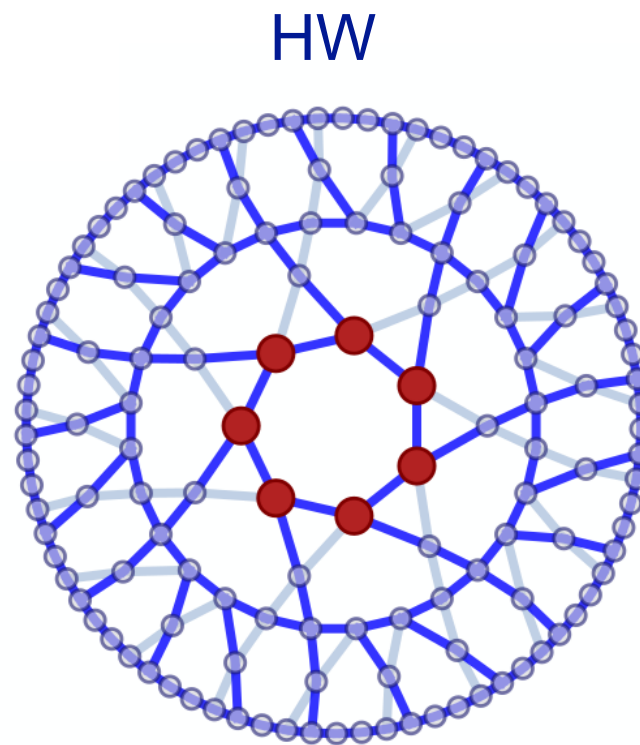
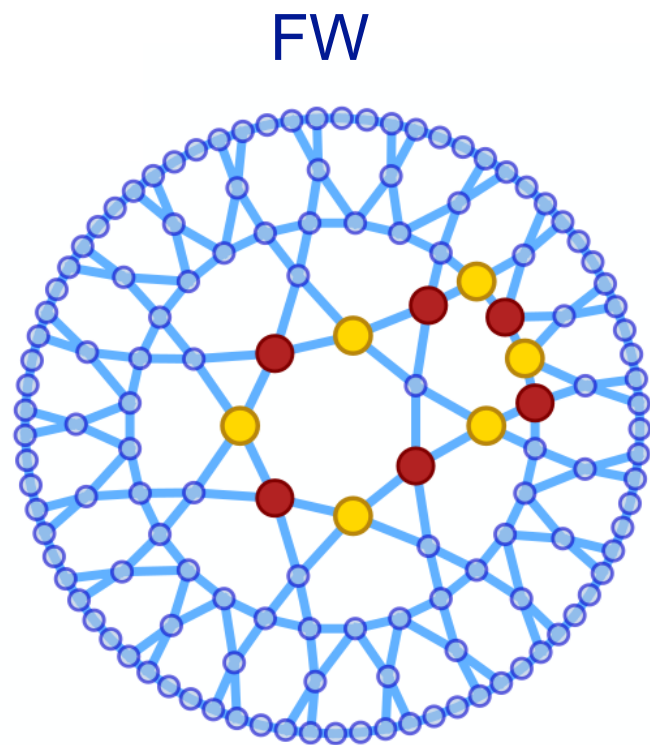
FW



HW

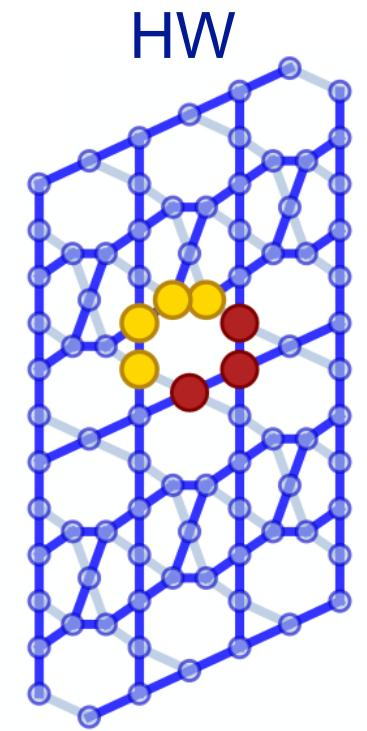
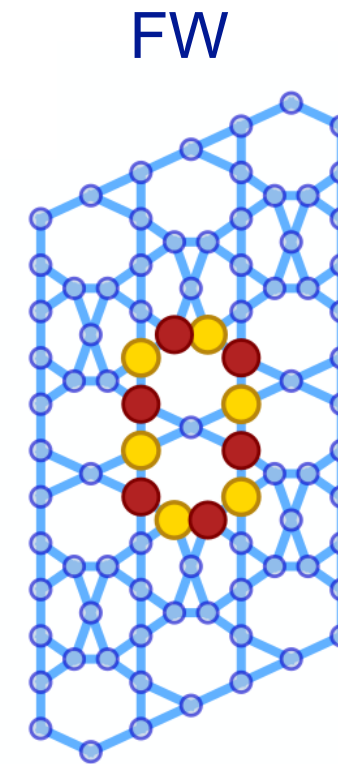
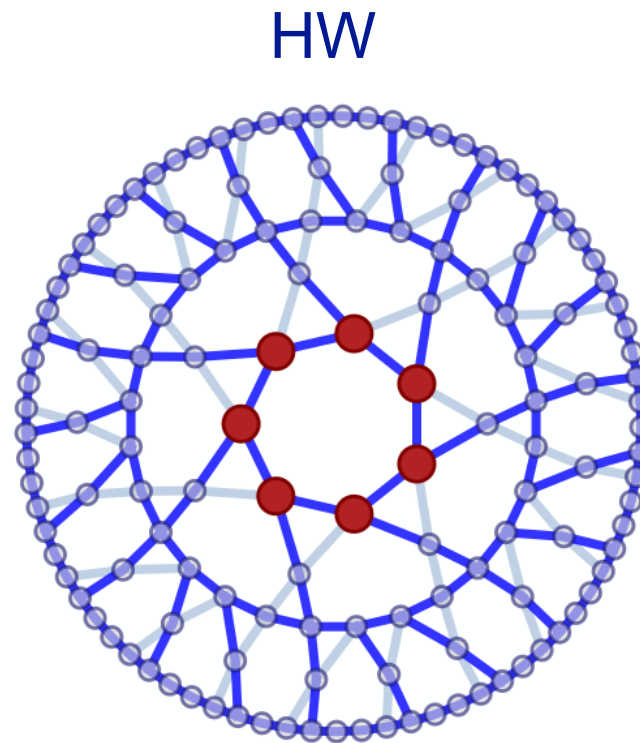
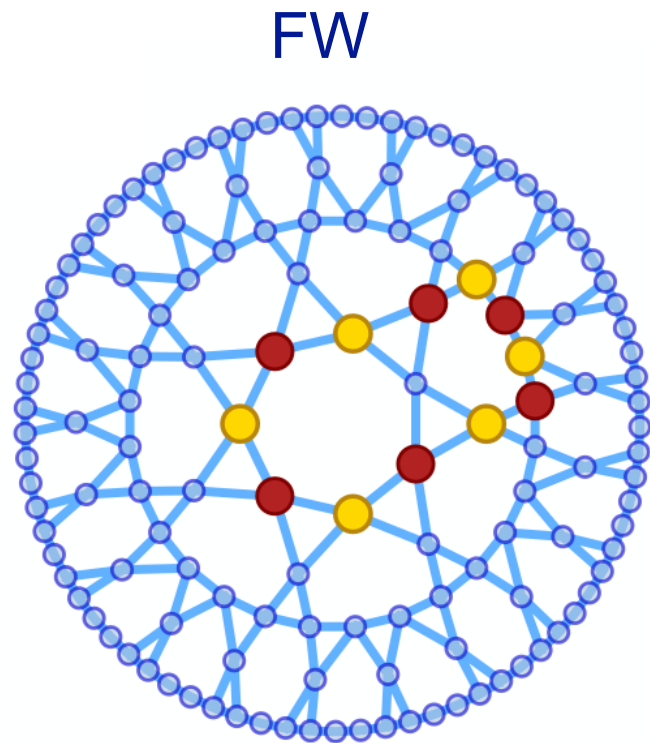


Full-Wave v Half-Wave Flat Band States



- Full-wave has localized states on only **even** cycles of the layout.

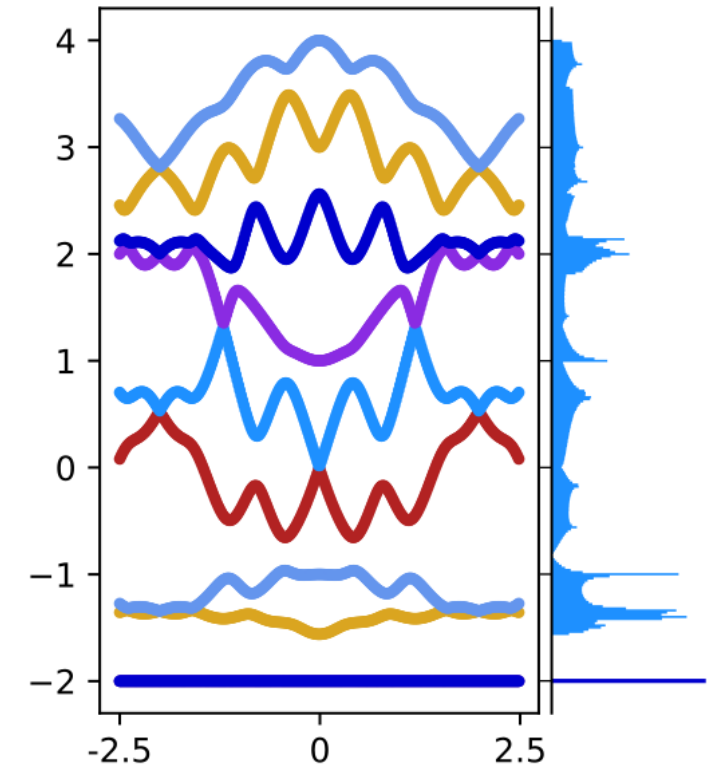
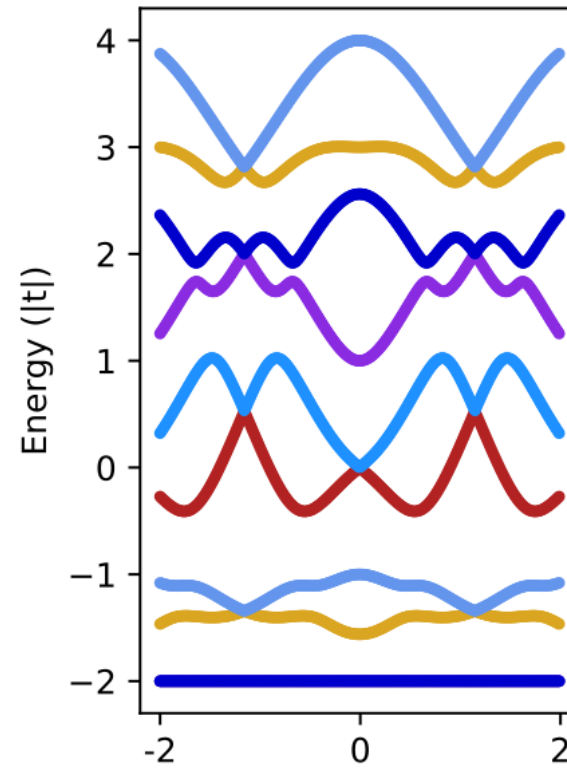
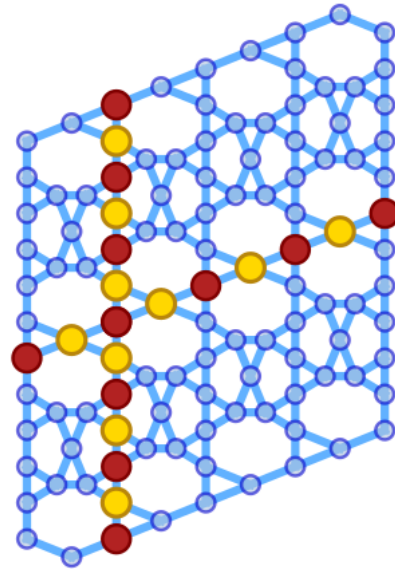
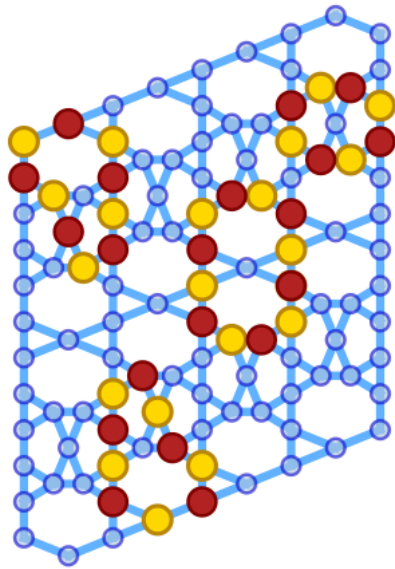
Full-Wave v Half-Wave Flat Band States



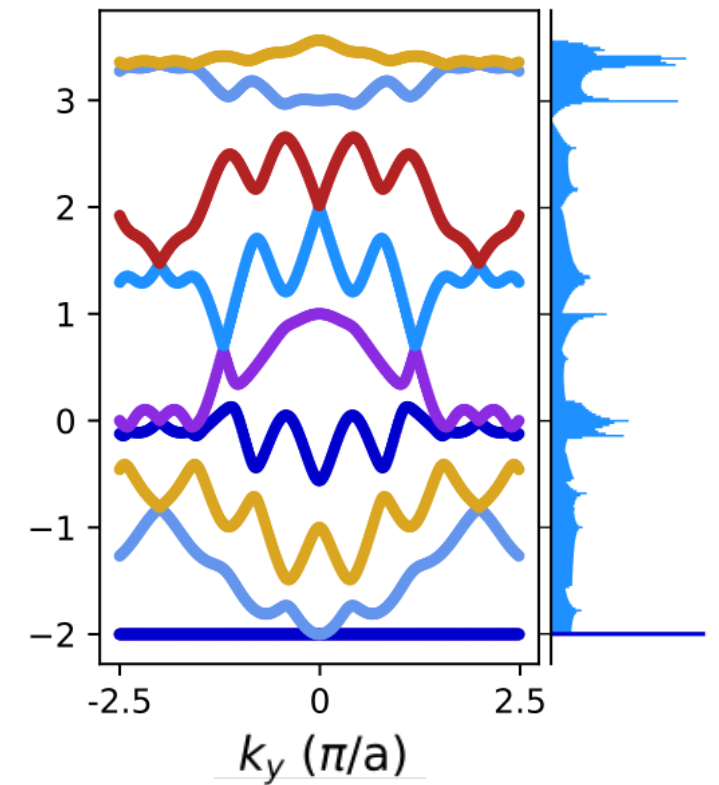
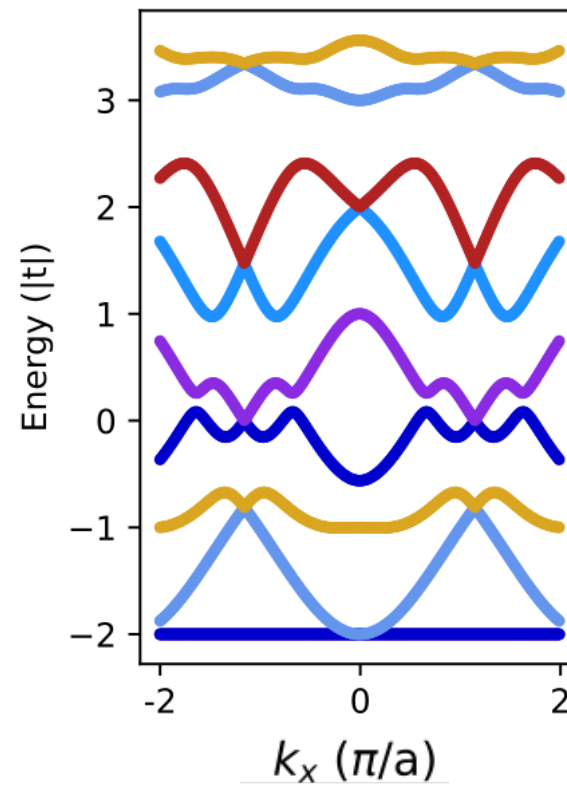
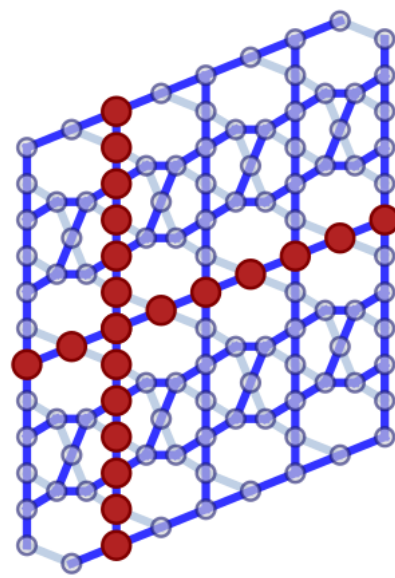
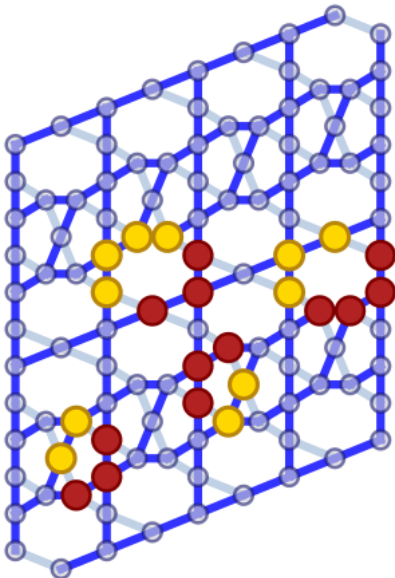
- Full-wave has localized states on only **even** cycles of the layout.
- Half-wave has localized states on **any** cycle of the layout.

Full-Wave Half-Wave Correspondence

FW

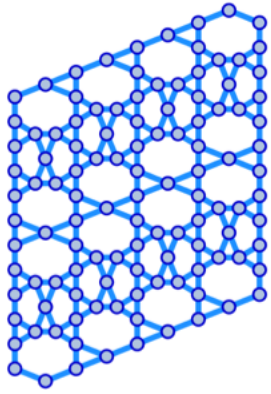


HW

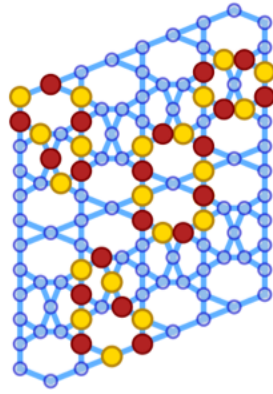


Real-Space Topology

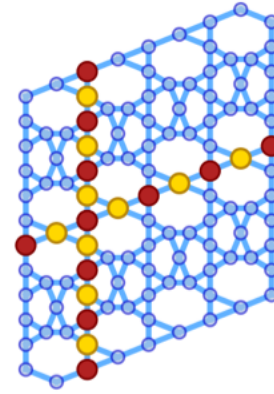
(i)



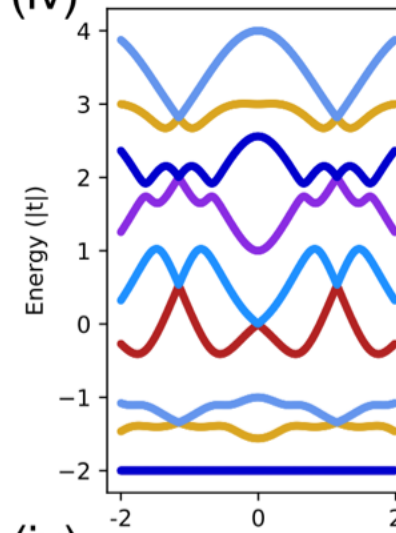
(ii)



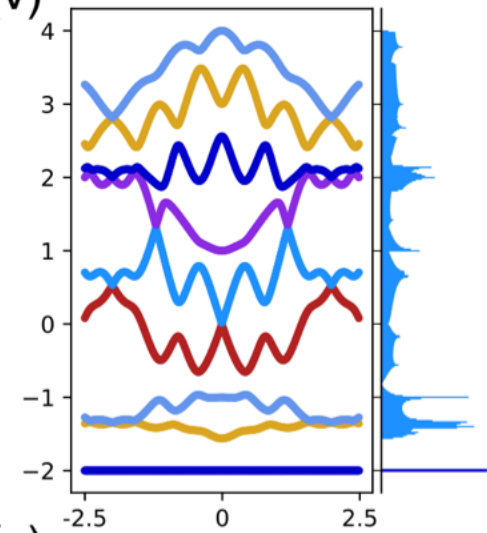
(iii)



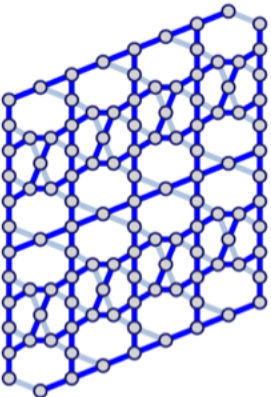
(iv)



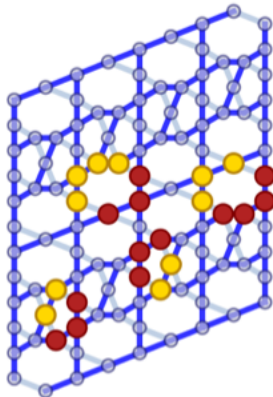
(v)



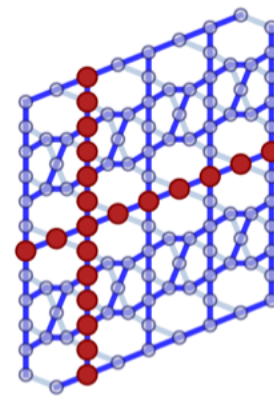
(i)



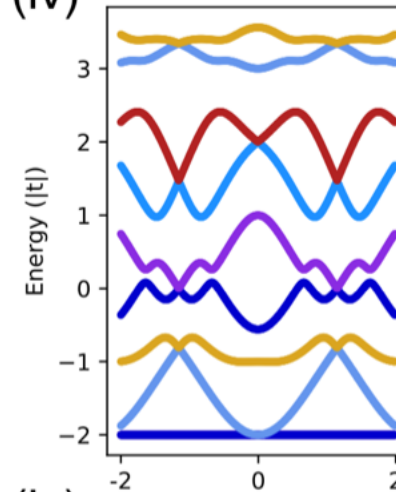
(ii)



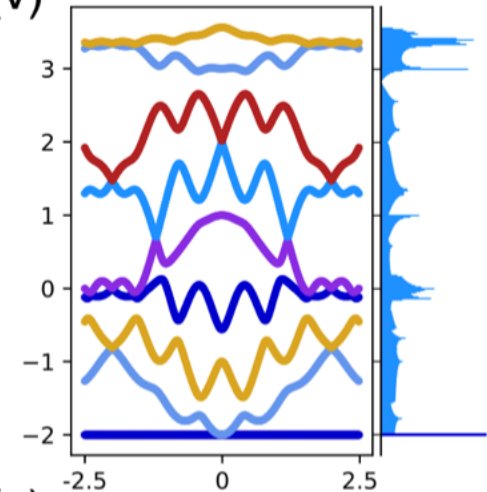
(iii)



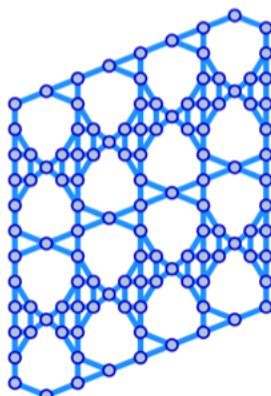
(iv)



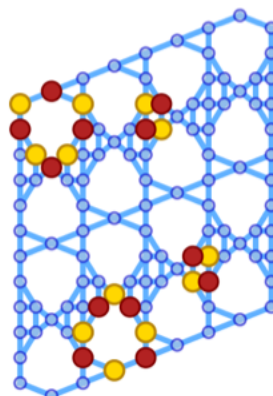
(v)



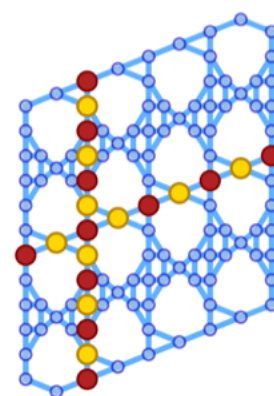
(i)



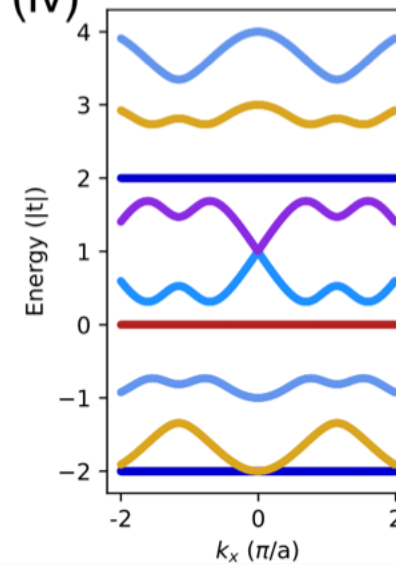
(ii)



(iii)



(iv)



(v)

