

### The AMPIS project

Status and Outlook

R. Caravita INFN - TIFPA



H2020 MSCA COFUND G.A. 754496

### Testing the universality of the free fall with cold antiatoms



#### **One needs**

- an intense beam of cold antiatoms **—** the **real** bottleneck
- a very high-resolution imaging detector for antiparticles
- a classical deflectometer/light interferometer

# AMPIS

### **Positronium** (Ps)

- electron + positron
- short-lived (only 142 ns!)
- table-top experiments

### AEgIS

₽

#### Antihydrogen $(\overline{H})$

- antiproton + positron
- only «stable» atom
- small amounts only @ CERN

Hydrogen



Two joint ventures with one single goal – developing the first source of cold antiatoms for a test of the Universality of the Free-Fall



Antihydrogen

### **A Metastable Positronium Inertial Sensor**



The crucial step: a laser cooled beam of Ps to enhance its collimation



Consider now an extra counter-propagating laser with negative photon momenta i.e., symmetric w.r.t. null velocity along its propagation axis. Transitions with negative momentum exchange are now allowed, leading to a symmetrical effect on the velocity distribution as the sign of the velocity detuning  $m\delta/k$  is also inverted by the  $k \rightarrow -k$  transformation.



Secondment, 2021



PHYSICAL REVIEW A

VOLUME 48, NUMBER 1

**JULY 1993** 

#### Measurement of the positronium $1 {}^{3}S_{1} - 2 {}^{3}S_{1}$ interval by continuous-wave two-photon excitation

M. S. Fee\* and S. Chu Physics Department, Stanford University, Stanford, California 94305

A. P. Mills, Jr., R. J. Chichester, and D. M. Zuckerman AT&T Bell Laboratories, Murray Hill, New Jersey 07974

The ultimate reduction in the positronium temperature would, of course, be achieved by laser cooling the atoms. Subkelvin positronium temperatures may be possible with laser cooling techniques. The laser cooling of positronium presents an unusual challenge since (i) the groundstate atom annihilates in 140 ns, (ii) the Doppler width at room temperature is ~ 500 GHz, and (iii) the single photon recoil velocity corresponds to a Doppler shift of 6.1 GHz. Thus, the "chirped slowing" commonly used for heavier atoms is not feasible. Even if it were possible to

#### Strong impact on the community

- Cold beam for interferometric studies
- Higher resolution spectroscopy limited by T
- Opens to the possibility of antimatter BECs
- Achievement awaited for 30 years!





### Cooling transition

- $-1^{3}S \rightarrow 2^{3}P$
- Wavelength = 243.024 nm
- Lifetime = 3.2ns



**Effective two-level system** with an extra loss mechanism (annihilation from the ground state)





### Cooling transition

- $-1^{3}S \rightarrow 2^{3}P$
- Wavelength = 243.024 nm
- Lifetime = 3.1ns
- Probing transition
  - Intermediate state excitation
    - $1^{3}S \rightarrow 3^{3}P$
    - Wavelength = 205.045nm
  - Photoionization
    - $3^{3}P \rightarrow$  continuum
    - Wavelength = 1064nm

**Two-step incoherent photoionization** sensitive to the Ps velocity distribution through the probe laser bandwidth







### Cooling transition

- $-1^{3}S \rightarrow 2^{3}P$
- Wavelength = 243.024 nm
- Lifetime = 3.1ns
- Probing transition
  - Intermediate state excitation
    - $1^{3}S \rightarrow 3^{3}P$
    - Wavelength = 205.045nm
  - Photoionization
    - $3^{3}P \rightarrow$  continuum
    - Wavelength = 1064nm

**Magnetic quenching** is our primary competing effect altering natural spontaneous emission cycles



### **Detecting Ps via SSPALS**

### Single-Shot Positron Annihilation Lifetime Spectroscopy (SSPALS)





#### Measurement classes in a laser cooling run







Amount of 1<sup>3</sup>S Ps in ground state

 $A_{205}$ 

Amount of 1<sup>3</sup>S Ps surviving photoionization by probe laser



 $A_{243}$ 

Amount of 1<sup>3</sup>S Ps

surviving quenching

and cooling cycles



 $A_{205+243}$ 

Amount of 1<sup>3</sup>S Ps surviving quenching, cooling cycles and subsequent photoionization



In a typical SSPALS experiment, one constructs the S parameters by comparing the 1<sup>3</sup>S Ps amount in a background reference measurement (**bck**) to the amount in the signal measurement (**sig**)

$$S_{\rm bck}^{\rm sig} := \frac{A_{\rm sig} - A_{\rm bck}}{A_{\rm bck}}$$



Oservable for laser cooling:

$$S_{243}^{205+243} := \frac{A_{205+243} - A_{243}}{A_{243}}$$

Secondment, 2021



31-May-22



### **Expected signal**











Supervision of one PhD student (CZ)



### **Overall experimental setup**









Best 1S-3P transition measurement so far



### **Comparison with expected signal**



Get rid of the 150 G magnetic field, and try again with higher statistics (~ one integrated month of data taking with the single PbWO<sub>4</sub>)

Secondment, 2021

### **Improved cooling laser**

#### Supervision of one PhD student (NG)



Tripled Alexandrite laser Broaband long pulses at 243 nm (an idea I've been proudly pushing for years <sup>(C)</sup>)

FELLINI support for line broadening with RF generator





31-May-22



### PHYSICAL REVIEW B 105, 115422 (2022)

#### Forward emission of positronium from nanochanneled silicon membranes

S. Mariazzi<sup>®</sup>,<sup>1,2,\*</sup> B. Rienäcker<sup>®</sup>,<sup>3,†</sup> R. Magrin Maffei<sup>®</sup>,<sup>4,5</sup> L. Povolo<sup>®</sup>,<sup>1,2</sup> S. Sharma,<sup>2</sup> R. Caravita<sup>®</sup>,<sup>2</sup> L. Penasa<sup>®</sup>,<sup>1,2</sup> P. Bettotti,<sup>1</sup> M. Doser,<sup>3</sup> and R. S. Brusa<sup>®</sup>,<sup>1,2</sup>



Published with FELLINI acknowledgment



### Improved Position- and timing- sensitive detector

Supervision of one PhD student (LG)

### MCP-TimePix3 ASIC hybrid

- Event-by-event reconstruction with positional and timing informations
- 10 bit charge readout of TimePix3: 2D centroiding reconstruction and subpixel resolution
- Benefits from the high detection efficiency of chevron MCPs

### Preliminary results (with e<sup>+</sup>)

- 15 ns timing resolution
- 12 um spatial resolution
- 41 % total detection efficiency

Resolution down to 5 um with UV light

In press (with FELLINI ack.)





- Many technological advances towards the first beam of cold, long-lived positronium atoms for inertial measurements
- Development of efficient positron-positronium targets suitable for transmission geometries
- Development of a novel **detector** for **high-resolution** timing and imaging antimatter events
- Solid perspective to establish soon laser cooling of positronium
- Constructing the first inertial sensor employing antimatter atoms is within reach!







### Thank you for your attention





#### H2020 MSCA COFUND G.A. 754496

The Antimatter Experiment: Gravity, Interferometry, Spectroscopy (**AEgIS**) collaboration aims at performing direct experimental tests of the Weak Equivalence Principle (**WEP**) using **anti-atoms**.

The chosen method is the **direct detection of the free-fall trajectory** of antihydrogen atoms, produced in a **pulsed** fashion

#### The CERN accelerator complex Complexe des accélérateurs du CERN







### **Conceptual experimental scheme**



#### Inertial sensing with a moiré deflectometer

- 1) Anti-atoms from an incoherent, uncollimated source
- 2) A set of two gratings selects the trajectories
- 3) Probed by a position- and timing-sensitive detector

#### Pulsed antihydrogen source

- 1) Cold antiproton  $(\bar{p})$  plasma preparation in a Penning trap
- 2) Pulsed positronium (Ps) from positron (e<sup>+</sup>) conversion
- 3) Two-step laser excitation of Ps to Rydberg levels (Ps\*)
- 4) Rydberg antihydrogen ( $\overline{H}^*$ ) via charge-exchange,  $\sigma \propto n_{Ps}^4$ 5) Beam formation







**AEgIS** concluded its **Phase 1** in 2018, achieving its goal of demonstrating the first pulsed H source



### Main experimental results obtained during the AEgIS Phase 1

- Validation of the inertial sensing methodology with antiprotons
- Ps excitation to Rydberg levels in strong magnetic fields
- First pulsed antihydrogen source demonstration





Inertial sensing with antimatter

Received 5 Nov 2013 | Accepted 27 Jun 2014 | Published 28 Jul 2014

DOI: 10.1038/ncomms5538 OPEN

A moiré deflectometer for antimatter



### **Key findings**

- Talbot-Lau interferometry is a effective tool for gratings alignment in 3D
- Near-field interference provides an undeflected reference
- aN sensitivity with 100 keV antiprotons and emulsion detectors

$$^{\mathbf{SD}}_{\mathbf{A}y} = \frac{F_{\parallel}}{m} \tau^2 \longrightarrow F_{min} \approx 5 \cdot 10^{-16} \,\mathrm{N}$$



### Ps excitation to Rydberg levels in strong magnetic fields

#### PHYSICAL REVIEW A 102, 013101 (2020)

#### Rydberg-positronium velocity and self-ionization studies in a 1T magnetic field and cryogenic environment





### **Key findings**

- Positronium excited to n = 15 17 in a 1T magnetic field
- Rydberg Ps self-ionizes due to the motional Stark electric field
- Limiting factor: Ps cannot be excited at higher levels than n = 17

$$\vec{E}_{MS} = \vec{v} \times \vec{B} \approx 0.5 \div 1 \, kV \, cm^{-1}$$

$$n_{\max} = \sqrt[4]{\frac{1.3 \cdot 10^{11}}{9vB\sin\theta}}$$





https://doi.org/10.1038/s42005-020-00494-z

OPEN

## Pulsed production of antihydrogen



- 0.05  $\overline{H}^*$  produced per cycle (110 s) with 10<sup>6</sup> antiprotons and n = 17 Ps<sup>\*</sup>
- Initial temperature of the antiproton plasma in the asymmetric trap of 400 K
- Unexpectedly long 25 us bunch: currently no explanation for this effect

Published with FELLINI





Main goal of AEgIS Phase 2: a first proof-of-concept inertial measurement with pulsed antihydrogen

### Take-home messages from the AEgIS Phase 1

- The antihydrogen source intensity must be increased by 2/3 orders of magnitude
- The temperature of the produced atoms must be reduced by 1 order of magnitude
- The first gravitational measurement has to be designed to use Rydberg antihydrogens
- The free-fall should take place in the most homogeneous volume of the AEgIS magnet

### **New AEgIS Phase 2 configuration**

- Positronium conversion target on-axis
- Laser excitation in a Doppler-free scheme
- Positrons passing through resting antiprotons





### A closer look to the new experimental scheme



#### The way to a higher production cross-section

- Motional Stark effect is cancelled at first order
- Rydberg levels as high as n = 35 are reachable
- No azimuthal asymmetry in the antiproton trap
- Bigger electrodes with smaller imperfections

**Caveat**: some heat transferred to the pbar plasma from the passage of the positron burst (small)



Limit from motional Stark ionization

[1] S. Mariazzi et al., J. Phys. B (2021) 085004



### **AEgIS-2 trap redesign**

#### AEgIS-2 coaxial trap design

- Cryoactuators for B field alignment
- Double laser passage with fiber monitoring
- Movable target holder and ionization grid
- Cryofilters for minimal Joule heating

Optimized Nanochanneled Si Converter





### First antiprotons from ELENA in AEgIS

Supervision of two PhD students (SH and MV)



#### First antiprotons from ELENA observed on 20 October, 2021



Fellini annual meeting

### **Capture of antiprotons**



First antiproton from ELENA caught from AEgIS-2 on 11th of November, 2021



Fellini annual meeting

### AEgIS Phase 2 timeline (COVID-corrected)

	2021	2022	2023	2024	2025
WEP test on a pulsed Rydberg Hbar beam					
Installation and test of new trap electrodes					
Connection to the ELENA beamline					
Improvement of the Hbar source flux					
Development of pulsed Hbar beam via trapped antiprotons					
Interaction of Hbar* with gratings					
Proof-of-concept inertial sensing with pulsed Hbar beam					

#### Summary

- Demonstration of a first pulsed source of H
  <sup>\*</sup>
- O(10<sup>3</sup>) higher  $\overline{H}^*$  production rate
- Pulsed beam of low energy  $\overline{H}^*$
- Gratings design validation with a beam of Rydberg atoms ( $\overline{H}^*$ , Ps<sup>\*</sup>,  $\overline{p}$ ion ...)
- First attempt of the measurement!



