

The Science and T.T. behind Sibylla Biotech SRL

Pietro Faccioli

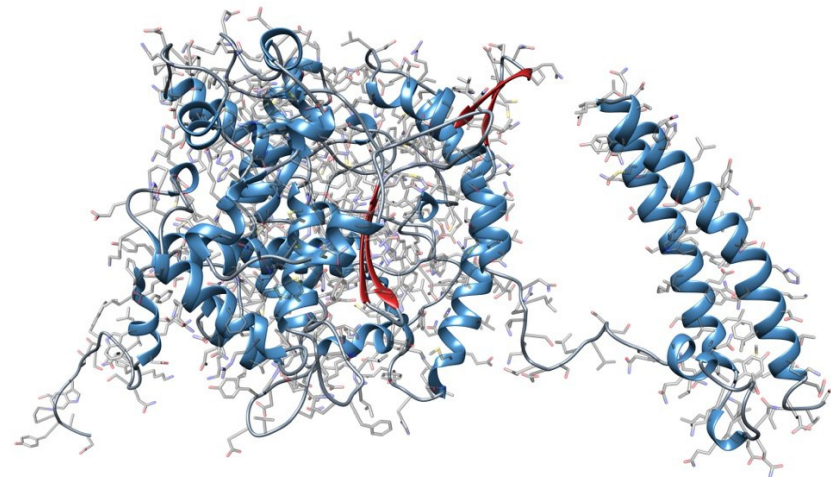


UNIVERSITÀ DEGLI STUDI
DI TRENTO

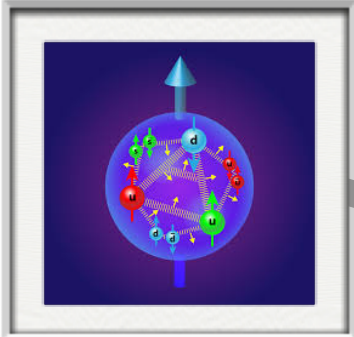
Dipartimento di Fisica



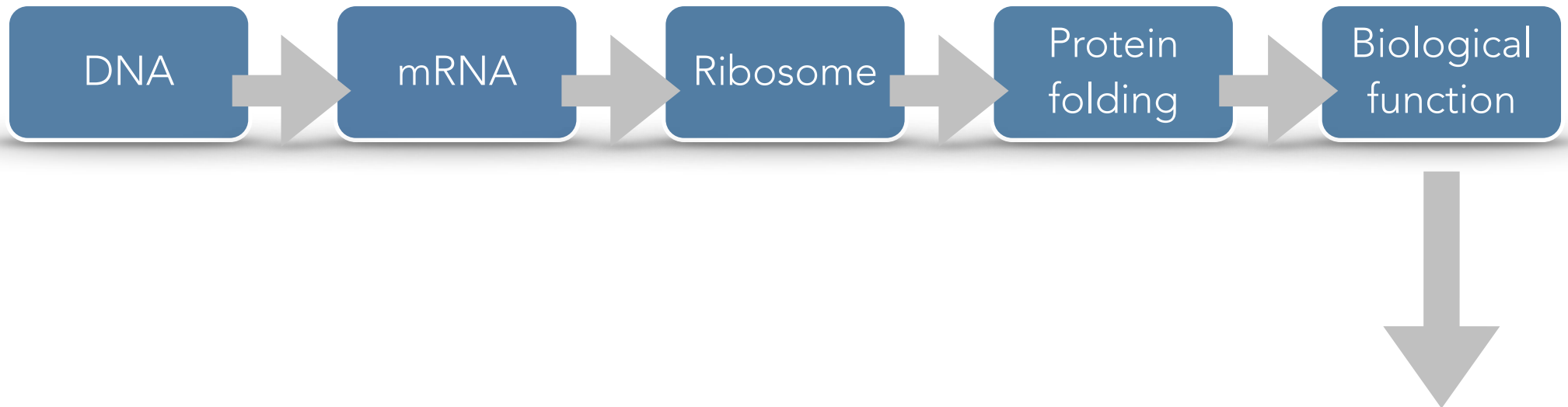
Trento Institute for
Fundamental Physics
and Applications



A SCIENTIFIC JOURNEY



FUNDAMENTAL DOGMA OF MOLECULAR BIOLOGY

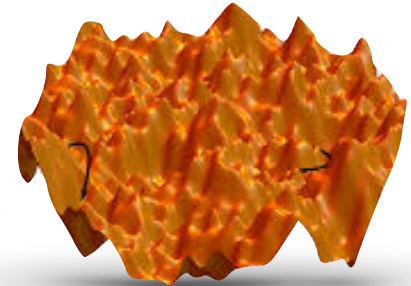
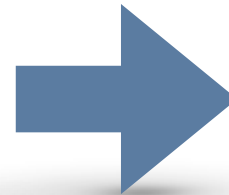


...proteins are Life's nanomachines!

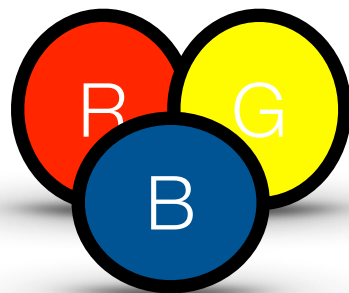
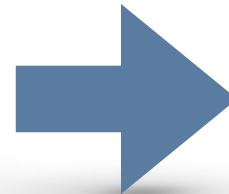
PROTEINS AND HADRONS ARE VERY SPECIAL PHYSICAL SYSTEMS



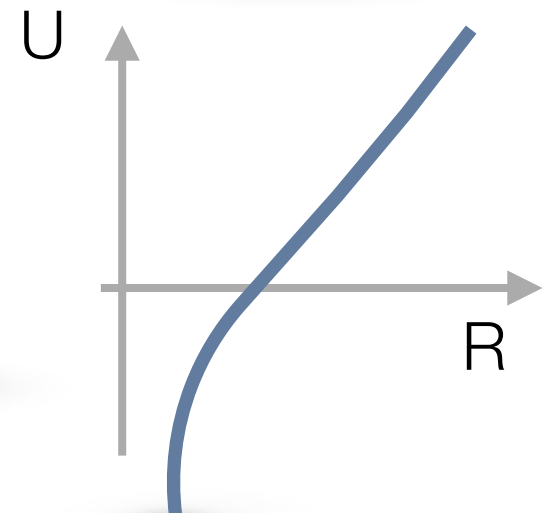
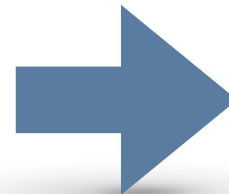
Random polypeptide



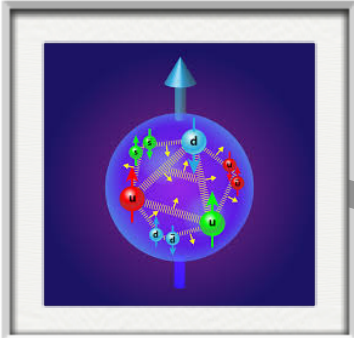
Protein



Baryon

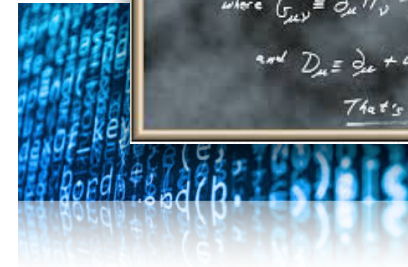


PHASE 1: MATHEMATICAL FORMALISM & HIGH PERFORMANCE COMPUTING

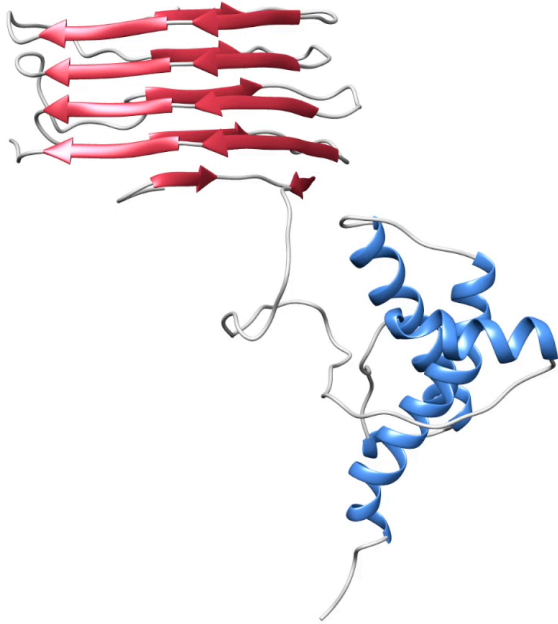


$$\mathcal{L} = \frac{1}{4g^2} G_{\mu\nu}^a G_{\mu\nu}^a + \sum_f \bar{\psi}_f (i \not{\partial} \not{D}_\mu + m_f) \psi_f$$

where $G_{\mu\nu}^a \equiv \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + gf_{abc} A_\mu^b A_\nu^c$
and $D_\mu \equiv \partial_\mu + i t^a A_\mu^a$
That's it!



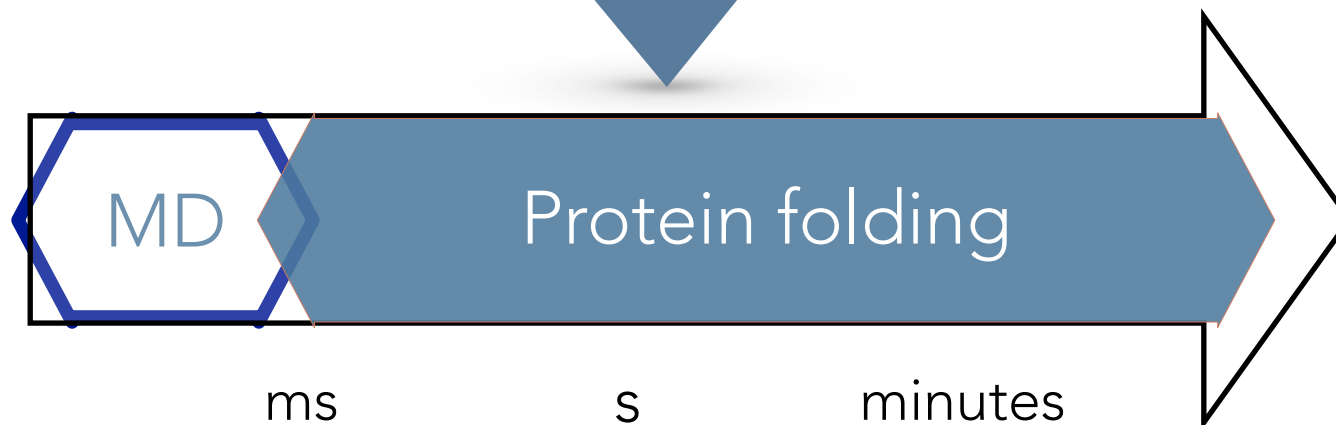
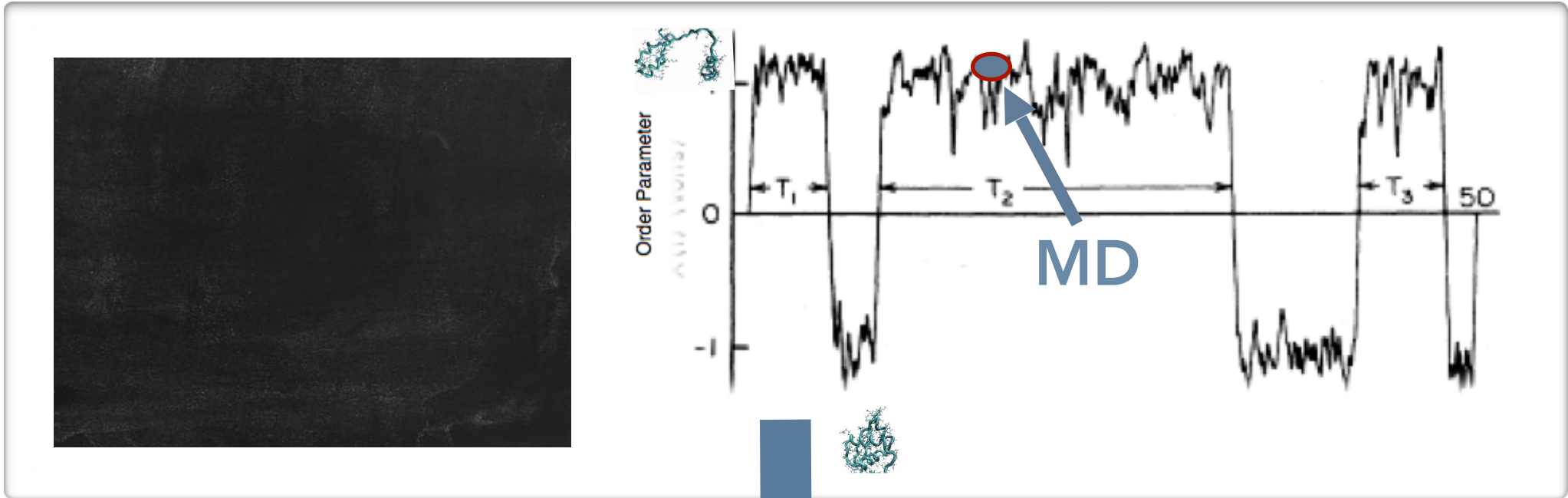
REDUCTIONIST'S APPROACH TO MOLECULAR BIOLOGY



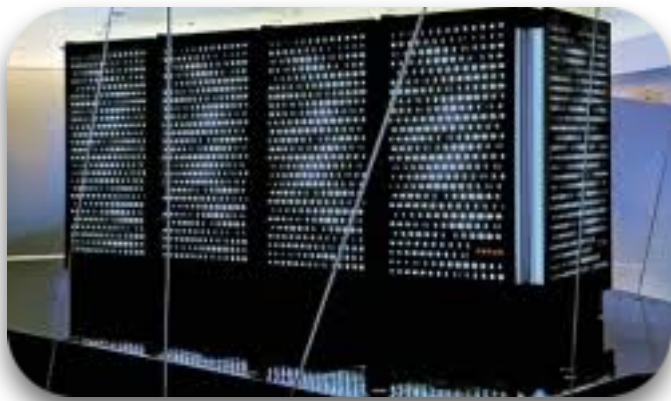
Challenge:

Integrate $\sim 10^6$ coupled
Newton-type equations
looking for **extremely
rare events**

PROTEIN DYNAMICS IS FULL OF RARE EVENT PROBLEMS



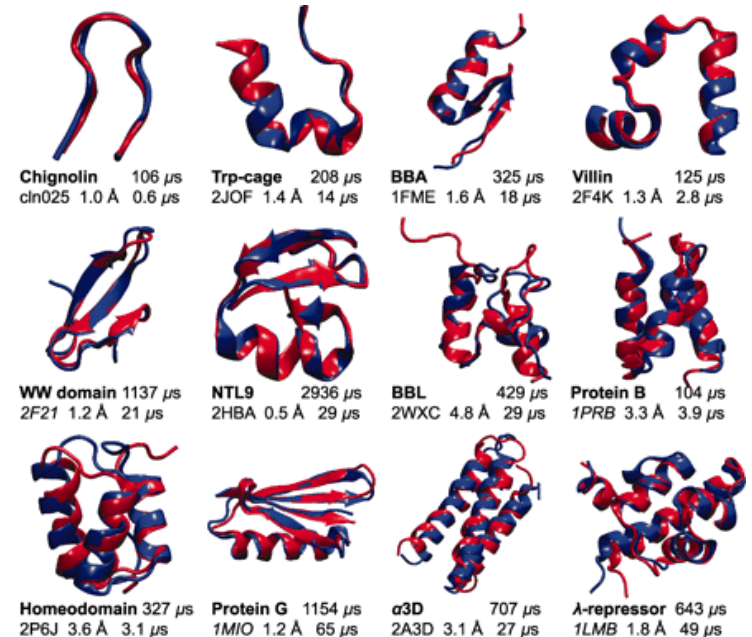
MD YIELDS CORRECT PROTEIN NATIVE STATES



Anton supercomputer
(DES Research)



MD



Atomic-Level Characterization of the Structural Dynamics of Proteins
David E. Shaw, *et al.*
Science **330**, 341 (2010);
DOI: 10.1126/science.1187409

How Fast-Folding Proteins Fold

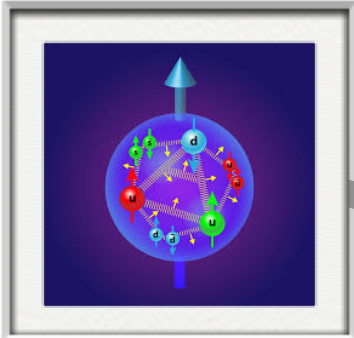
Kresten Lindorff-Larsen,^{1*}† Stefano Piana,^{1*}† Ron O. Dror,¹ David E. Shaw^{1,2†}

ZOOLOGY OF ENHANCED SAMPLING METHODS

Markov State Models, Milestoning, Transition Path Sampling, Transition Interface Sampling, Forward Flux Sampling, Temperature Accelerated Molecular Dynamics, Metadynamics, Umbrella Sampling, Blue Moon Sampling, String Method, Stochastic Difference, ... [and counting]

They are **all too computationally demanding** for many biologically relevant problems.

PHASE 1: MATHEMATICAL FORMALISM & HIGH PERFORMANCE COMPUTING



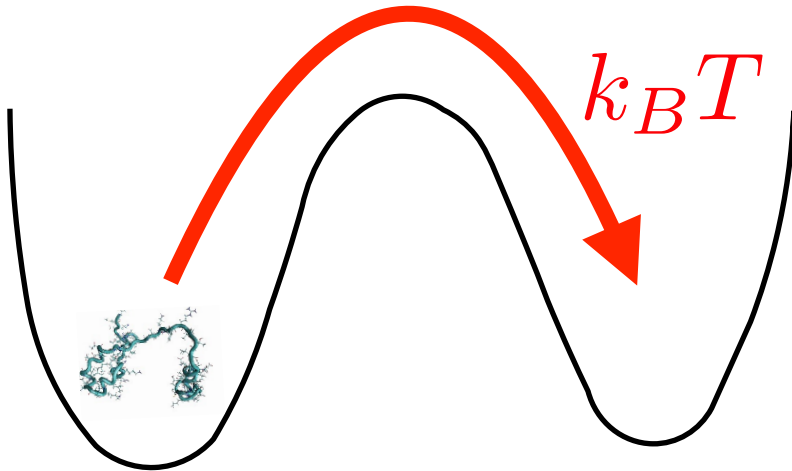
$$\mathcal{L} = \frac{1}{4g^2} G_{\mu\nu}^a G_{\mu\nu}^a + \sum_f \bar{\psi}_f (i \not{\partial} + m_f) \psi_f$$

where $G_{\mu\nu}^a \equiv \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + gf_{abc} A_\mu^b A_\nu^c$
and $D_\mu \equiv \partial_\mu + i t^a A_\mu^a$
That's it!



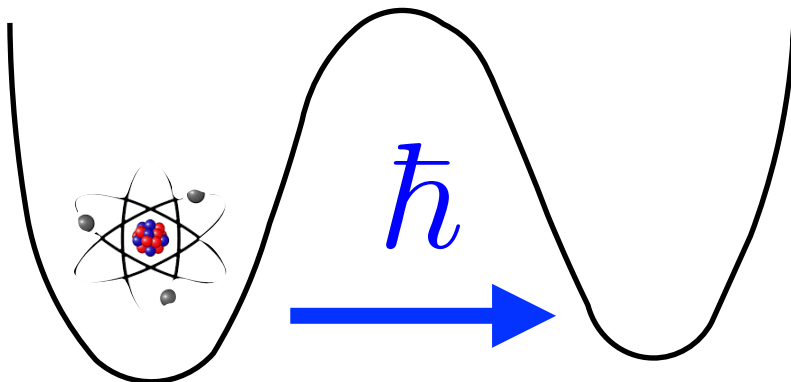
A USEFUL ANALOGY

Thermal activation



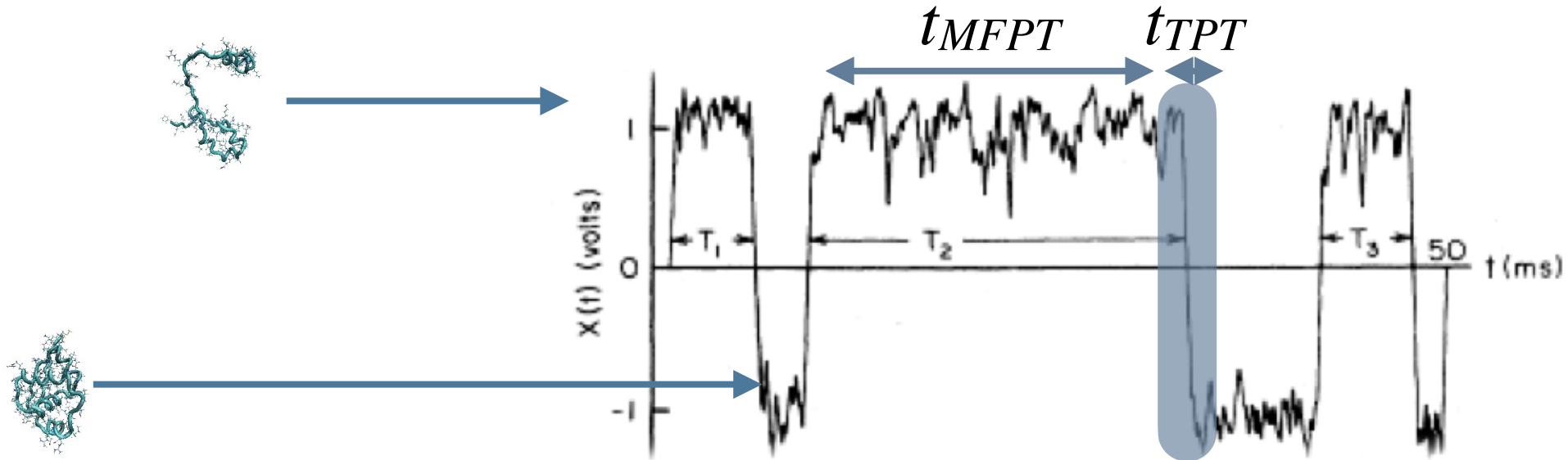
$$P(x_f, t|x_i) = \int_{x_i}^{x_f} \mathcal{D}q e^{-\frac{\beta}{4M\gamma} \int_0^t d\tau (M\ddot{q} + M\gamma\dot{q} + \nabla U(q))^2}$$

Quantum tunneling



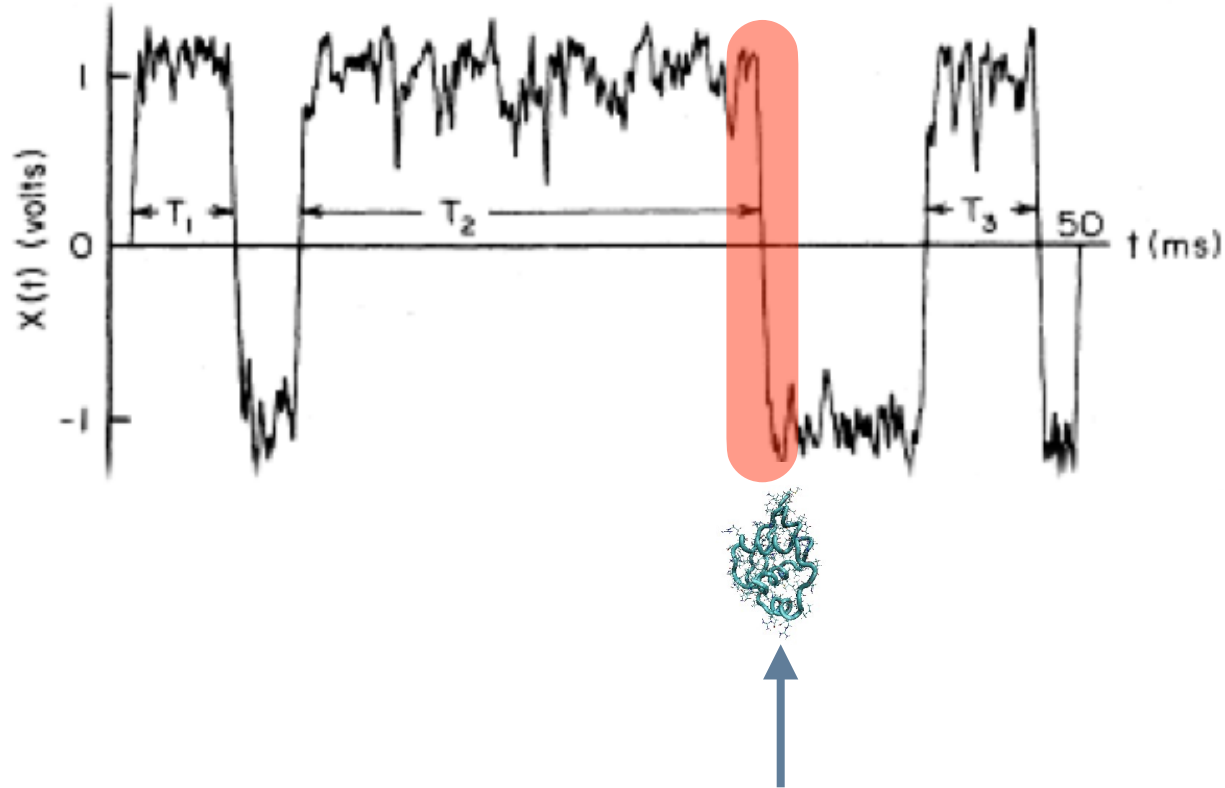
$$K_E(x_f, t|x_i) = \int_{x_i}^{x_f} \mathcal{D}q e^{-\frac{1}{\hbar} \int_0^t d\tau (\frac{M}{2} \dot{q}^2 + U(q))}$$

ADVANTAGES OF PATH INTEGRALS



$$t_{TPT} \sim \tau_0 \log \left[\log \left(\frac{t_{MFPT}}{\tau_0} \right) \right]$$

IS THIS A "FREE LUNCH"?



All atom 3D structure of the native state **are given in input**, not predicted

FULLY EXPLOITING THEORETICAL PHYSICS TOOLS

072336-4 Bartolucci, Orioli, and Faccioli

between the Gibbs distribution and the SCR estimate forward- and backward-committors, as in Eq. (A3). Introducing the distribution

$$P^{(P)}(x, t) \equiv \int dx_i P^{(P)}(x, t | x_i, 0) \rho_0(x_i), \quad (22)$$

the density in Eq. (22) reads

$$m_{SCR}(x) = \frac{1}{t_f - \tau_0} \int_{\tau_0}^{t_f} dt Q^{(R)}(x, t_f - t) P^{(P)}(x, t).$$

Using the detailed balance condition, we find $P^{(P)} = e^{-\beta U(x)} \frac{1}{Z_R} Q^{(P)}(x, t)$. Then, inserting this result into Eq. we find

$$m_{SCR}(x) = \frac{e^{-\beta U(x)}}{Z_R (t_f - \tau_0)} \int_{\tau_0}^{t_f} dt Q^{(R)}(x, t_f - t) Q^{(P)}(x, t).$$

Finally, recalling that $Q^{(R)}(x, t)$ and $Q^{(P)}(x, t)$ are time-independent in the SCR and using Eqs. (17) and we recover a fundamental result of TPT [cf. Eq. (A. Appendix A)],

$$m_{SCR}(x) \propto e^{-\beta U(x)} q_{SCR}^+(x) (1 - q_{SCR}^+(x)).$$

Within the same framework, it is possible to do the reactive current in the SCR in complete analogy Eq. (22),

$$J_{SCR}^i(x) = \frac{-D}{t_f - \tau_0} \int_{\tau_0}^{t_f} dt Q^{(R)}(x, t_f - t) \times (\vec{\nabla} - \overleftarrow{\nabla} + \beta \nabla U(x)) P^{(P)}(x, t).$$

$$\begin{aligned} V_{eff}^R(\mathbf{X}) &\simeq \frac{D_0(1-b)}{\pi b \Omega} \nabla^2 V_{eff}(\mathbf{X}) \\ &+ \frac{1}{2} \left(\frac{D_0(1-b)}{\pi b \Omega} \right)^2 \nabla^4 V_{eff}(\mathbf{X}) \\ &+ \frac{1}{6} \left(\frac{D_0(1-b)}{\pi b \Omega} \right)^3 \nabla^6 V_{eff}(\mathbf{X}) - \frac{D_0^2(1-b^3)}{3\pi(b\Omega)^3} \left(\partial_i \partial_j V_{eff}(\mathbf{X}) \right)^2. \end{aligned} \quad (24)$$

Note that the first line is the leading order term (i.e. $L = 1$), while the second and third lines display the order $L = 2$ and $L = 3$ corrections, respectively.

We emphasize that the result of the EST construction is a new expression for the *same* path integral (15), in which the UV cutoff has been lowered from Ω to $b\Omega$. Equivalently, the path integral is discretized according to a larger elementary time step, $\Delta t \rightarrow \Delta t/b$:

$$Z^{\Delta t}(t) \equiv \int_{\Delta t} \mathcal{D}\mathbf{X} e^{-S_{eff}[\mathbf{X}]} \propto \int_{\Delta t/b} \mathcal{D}\mathbf{X} e^{-S_{eff}[\mathbf{X}] - \int_0^t d\tau V_{eff}^R[\mathbf{X}(\tau)]} \equiv Z_{EST}^{\Delta t/b}(t) \quad (25)$$

In these expressions, the symbol $\int_{\Delta t}$ denotes the fact that the path integral is discretized according to an elementary time step Δt and we have suppressed the subscript " $<$ ", in the paths. It can be shown that the proportionality factor between $Z^{\Delta t}(t)$ and $Z_{EST}^{\Delta t/b}(t)$

PRL 114, 098103 (2015) PHYSICAL REVIEW LETTERS

$$\mathcal{P}_{\text{bias}}[X] = \int \mathcal{D}Y e^{-S_{\text{bias}}[X, Y] - U(X, Y_i) / k_B T}, \quad (3)$$

where the functional $S_{\text{bias}}[X, Y]$ is defined as

$$\begin{aligned} S_{\text{bias}} &\equiv \frac{1}{4k_B T} \int_0^t d\tau \left[\sum_{i=1}^N \frac{1}{\gamma_i m_i} (m_i \dot{x}_i + m_i \gamma_i \dot{y}_i + \nabla_i U - \mathbf{F}_i^{\text{bias}})^2 \right. \\ &\left. + \sum_{j=1}^N \frac{1}{\gamma_j m_j} (m_j \dot{y}_j + m_j \gamma_j \dot{x}_j + \nabla_j U)^2 \right]. \end{aligned} \quad (4)$$

The Onsager-Machlup functional $S_{\text{OM}}[X, Y]$ entering Eq. (2) is recovered, setting $\mathbf{F}_i^{\text{bias}} = 0$ in Eq. (4).

Let us now return to the problem of computing the reaction pathways in the *unbiased* Langevin dynamics [Eq. (1)]. Using the standard reweighting trick we can write the variational condition $(\delta/\delta X) \mathcal{P}[X] = 0$ as

$$\frac{\delta}{\delta X} [\mathcal{P}_{\text{bias}}[X] (e^{-S_{\text{OM}}[X, Y]} - S_{\text{bias}}[X, Y; t])_{\text{bias}}] = 0. \quad (5)$$

We now introduce our main approximation, by restricting the search for the optimum path $X(\tau)$ within an ensemble of trajectories generated by integrating the *biased* Langevin equation. By definition, these paths have a large statistical weight in the biased dynamics, i.e., they lie in the functional vicinity of some path $\tilde{X}(\tau)$ which satisfies $(\delta/\delta X) \mathcal{P}[\tilde{X}] = 0$. Thus, the typical biased paths approximately satisfy the stationary condition

$$\frac{\delta}{\delta X} \mathcal{P}[X]$$

This equation states that for which the force is least. In derived in the context of solvent-induced transitions. Let us now emphasize that the history-dependent ratchet-and-rhythm developed in Ref.

$$-\frac{k_B T}{2} \nabla$$

$z_m(t)$ decays to zero as time t (we obey the condition). Let us

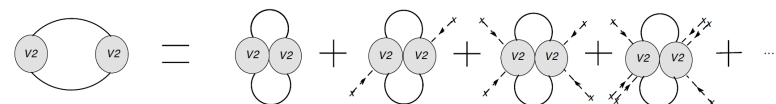


FIG. 3: Diagrammatic representation of the local time-derivative expansion of a non-local diagram —Eq. (49)—. Solid lines are fast-mode propagators, while dashed lines represent a single time derivative acting on the corresponding vertex function.

Notice that each term in the perturbative expansion (35) generates a new vertex, with an increasing power of the $x_{>}(\tau)$ field. The couplings to the fast modes depend implicitly on the time τ , through the slow modes $x_{<}(\tau)$.

By Wick theorem, each term in the series (34) can be related to a Feynman graph with vertices given by (36) and propagators given by —see appendix A —:

$$\langle x_{>}^i(\tau_1) x_{>}^j(\tau_2) \rangle_0 = \sum_{\{|\omega_n|, |\omega_m| \in S_b\}} G_{>}^{ij}(\omega_n, \omega_m) e^{i(\omega_n \tau_1 + \omega_m \tau_2)} = \sum_{\{|\omega_n| \in S_b\}} \delta_{ij} \frac{2}{\beta \gamma t \omega_n^2} e^{i\omega_n(\tau_2 - \tau_1)}. \quad (37)$$

The expansion (34) can be re-organized as the exponent of the sum performed over only connected diagrams:

$$e^{-\beta S_{>}[x_{<}(\tau)]} = e^{\Sigma(\text{all connected diagrams})}. \quad (38)$$

Hence, the path integral (26) for the slow modes can be given the following exact diagrammatic representation

$$Z(t) \equiv \int \mathcal{D}x_{<} e^{-\beta S_{eff}[x_{<}(t)] + \Sigma(\text{all connected diagrams})}. \quad (39)$$

Below we give a classification of all the connected diagrams that may give a contribution to the expansion above.

064108-3 Orioli, a Beccara, and Faccioli

J. Chem. Phys. 147, 064108 (2017)

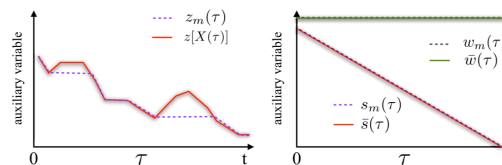


FIG. 1: Illustrative representation of the dynamics of the auxiliary variables introduced in the path integral representation of rMD (left panel) and in the derivation self-consistent path sampling algorithm (right panel).

of such a variable is frozen any time z_m becomes smaller than $z(X)$ and any time the collective coordinate $z(X)$ is increasing. Its time derivative is otherwise set equal to $\dot{z}(X)$. Therefore, by choosing the initial conditions $z_m(0) = z(X(0))$, $z_m(\tau)$ is identically set equal to the minimum value attained by the collective coordinate z until time τ (see left panel of Fig. 1).

The functional $S_{MD}[X, z_m]$ in the exponent of Eq. (8) coincides with an OM action with the addition of the unphysical biasing force \mathbf{F}_i ,

$$S_{MD} = \sum_{i=1}^N \Gamma_i \int_0^t d\tau [m_i \dot{x}_i + m_i \gamma_i \dot{y}_i + \nabla_i U - \mathbf{F}_i]^2. \quad (9)$$

In Eq. (8), $\Phi[z_m, X]$ denotes a Jacobian factor that needs to be introduced in order to ensure that the statistical weight of the paths is not affected by the measure of the $\int \mathcal{D}z_m$ integral, i.e.,

$$\int \mathcal{D}z_m \Phi[z_m, X] \delta \left[z_m(\tau) - \int_0^\tau d\tau' z[X(\tau')] \theta(-\dot{z}[X(\tau')]) \right]$$

III. SELF-CONSISTENT PATH SAMPLING

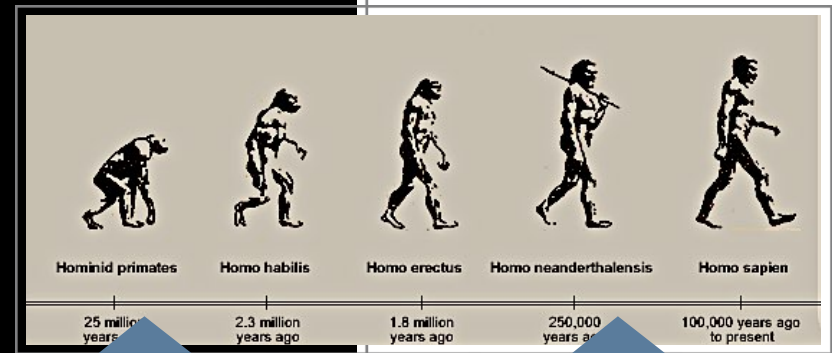
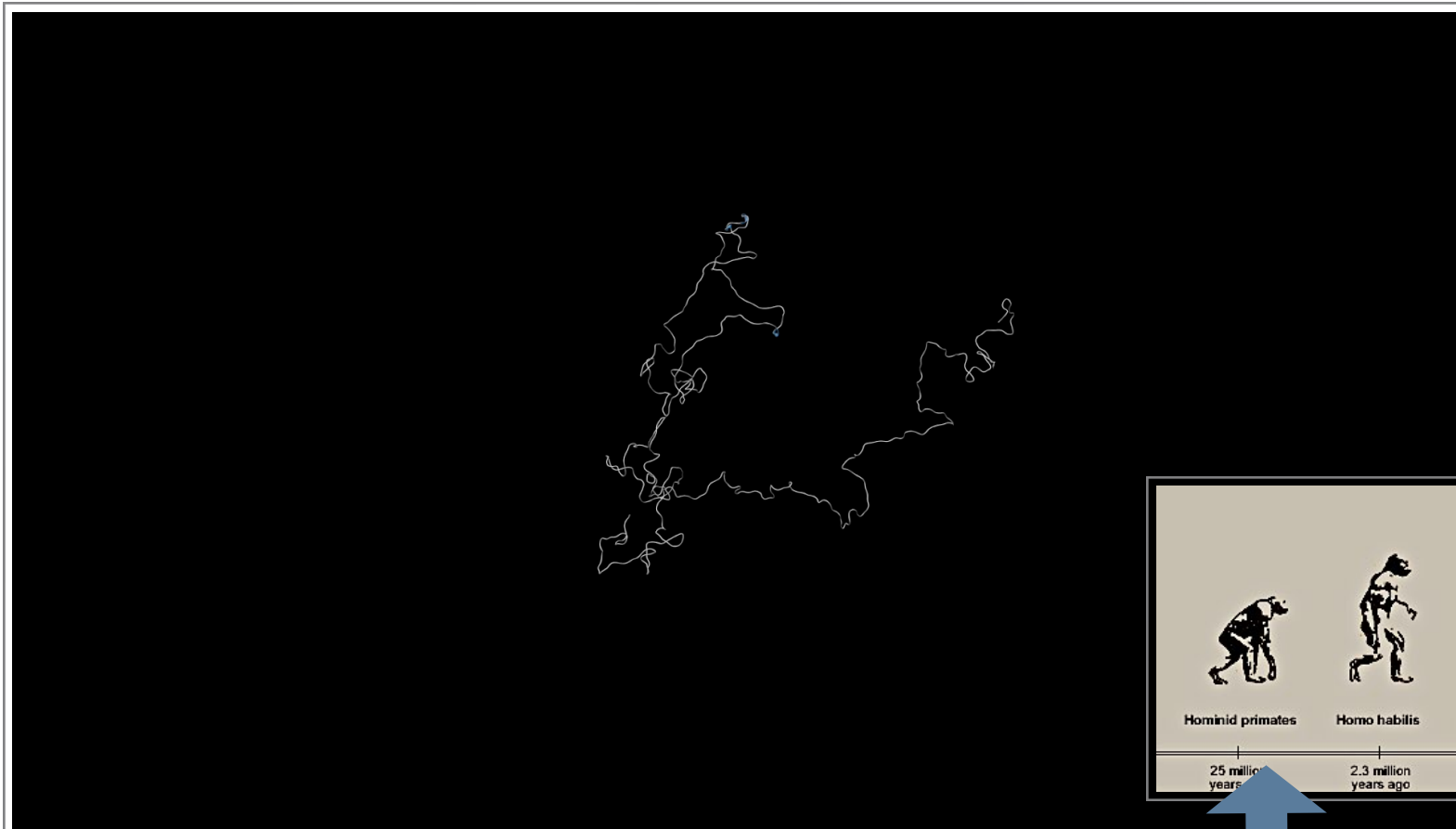
Let us now introduce our new algorithm, which provides major improvement with respect to the rMD and BF schemes discussed in Sec. II A. Indeed, it follows directly from the unbiased Langevin equation and allows us to remove the systematic errors associated to the choice of biasing coordinate.

Our starting point is path integral representation of the *unbiased* Langevin dynamics (2). We introduce two dumb auxiliary variables $w_m(\tau)$ and $s_m(\tau)$ into this path integral by means of appropriate functional Dirac deltas,

$$\begin{aligned} p(X_N, t | X_0) &= \int_{X_0}^{X_N} \mathcal{D}X \cdot e^{-S[X]} \int_{S(0)} \mathcal{D}S_m \int_{\hat{w}(0)} \mathcal{D}w_m \\ &\cdot \delta \left[w_m(\tau) - \int_0^\tau d\tau' \hat{w}(\tau') \theta(-\dot{\hat{w}}(\tau')) \theta(w_m(\tau') - \hat{w}(\tau')) \right] \\ &\cdot \delta \left[s_m(\tau) - \int_0^\tau d\tau' \tilde{s}(\tau') \theta(-\dot{\tilde{s}}(\tau')) \theta(s_m(\tau') - \tilde{s}(\tau')) \right], \end{aligned} \quad (12)$$

where $\tilde{s}(\tau)$ and $\hat{w}(\tau)$ are two external time-dependent functions to be defined below. In analogy with the path integral repre-

HUGE COMPUTATIONAL GAIN

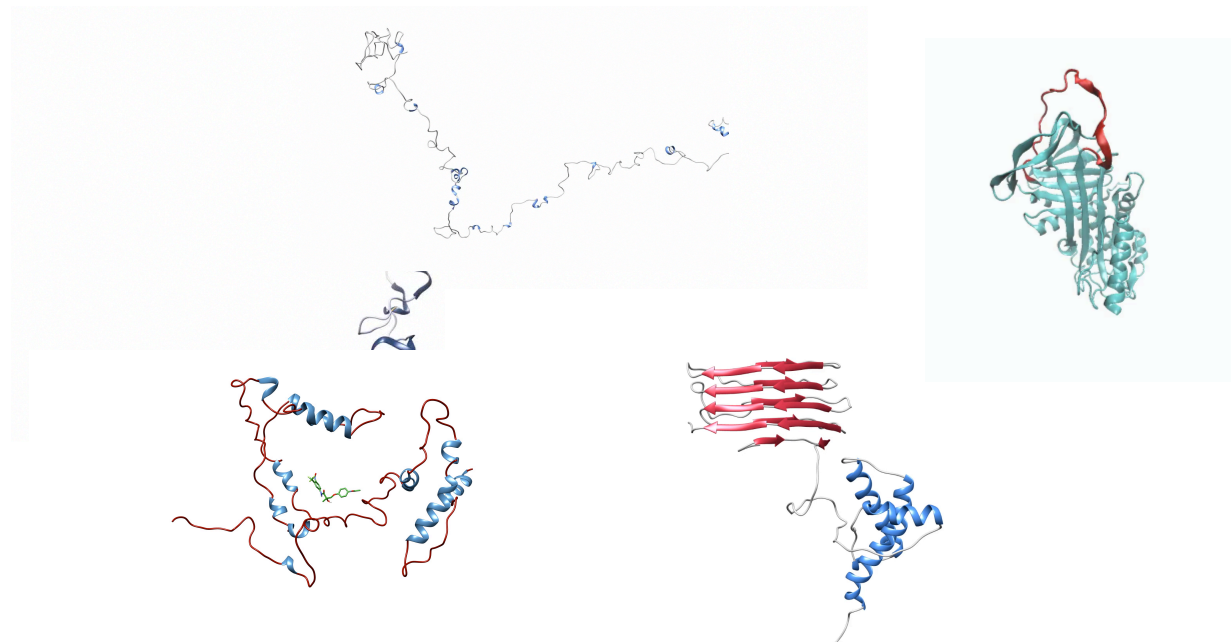
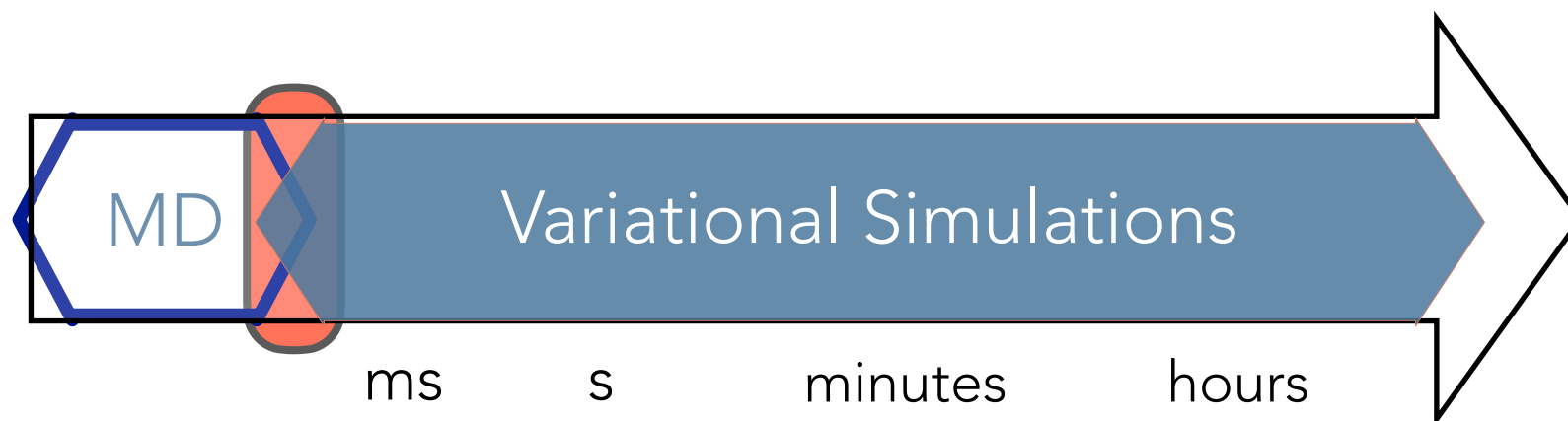


Sybilla & INFN joint project on ACE2

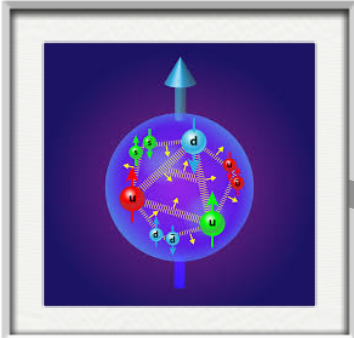
Using top all-purpose supercomputers

Using top special-purpose supercomputer

VENTURING INTO THE BIO-ZONE

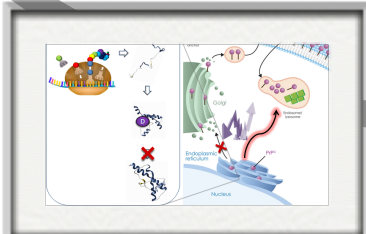


THE WAY TO PHARMACOLOGICAL RESEARCH



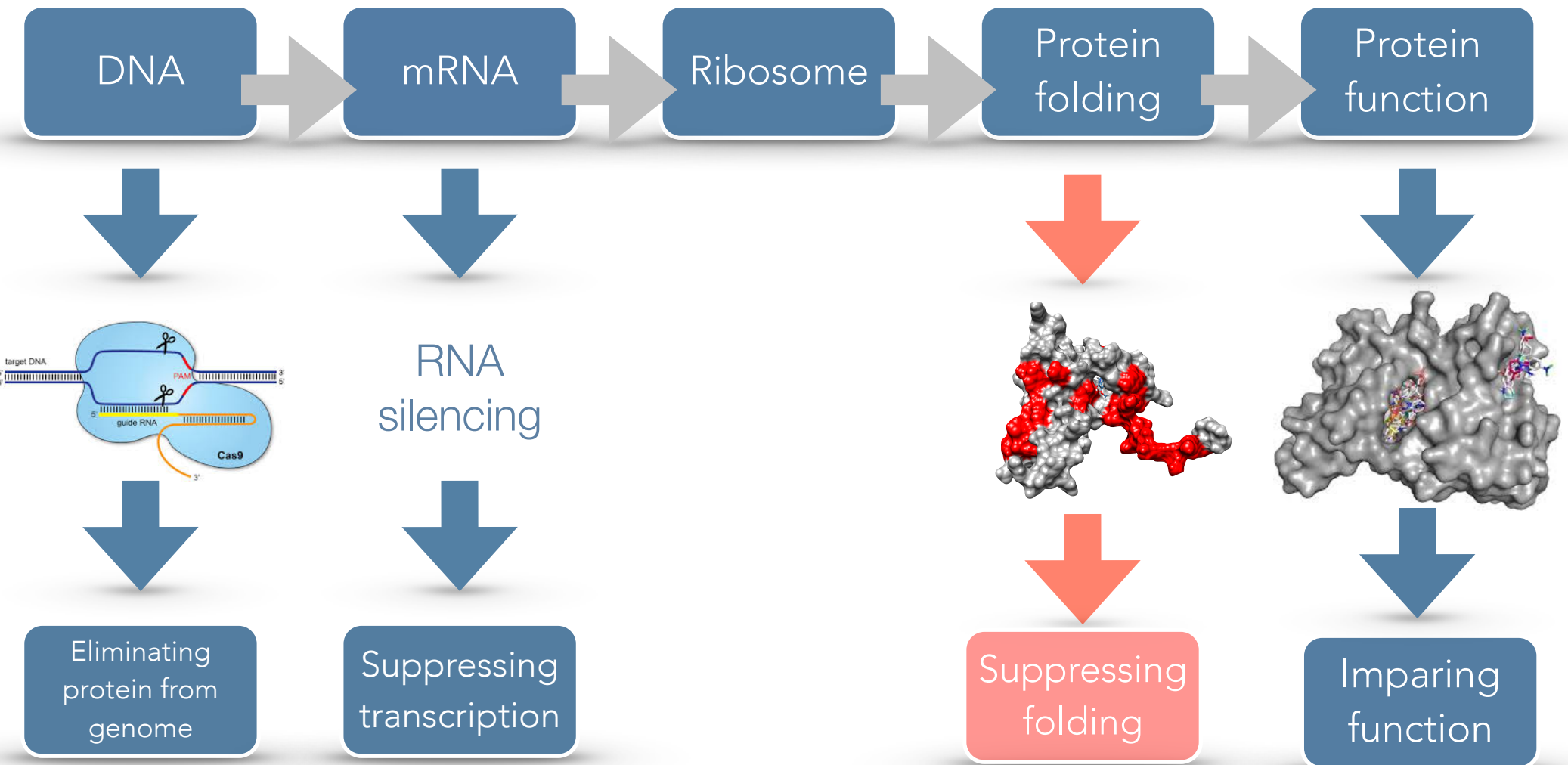
$$\mathcal{L} = \frac{1}{4g^2} G_{\mu\nu}^a G_{\mu\nu}^a + \sum_f \bar{\psi}_f (i\gamma^\mu \partial_\mu + m_f) \psi_f$$

where $G_{\mu\nu}^a \equiv \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + gf_{abc} A_\mu^b A_\nu^c$
and $D_\mu \equiv \partial_\mu + i t^a A_\mu^a$
That's it!

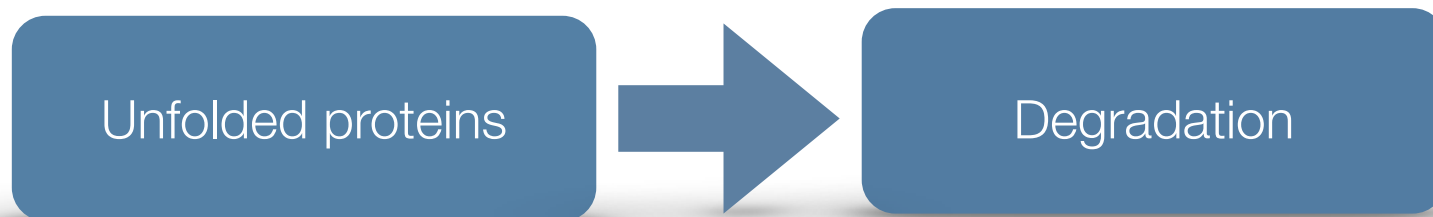
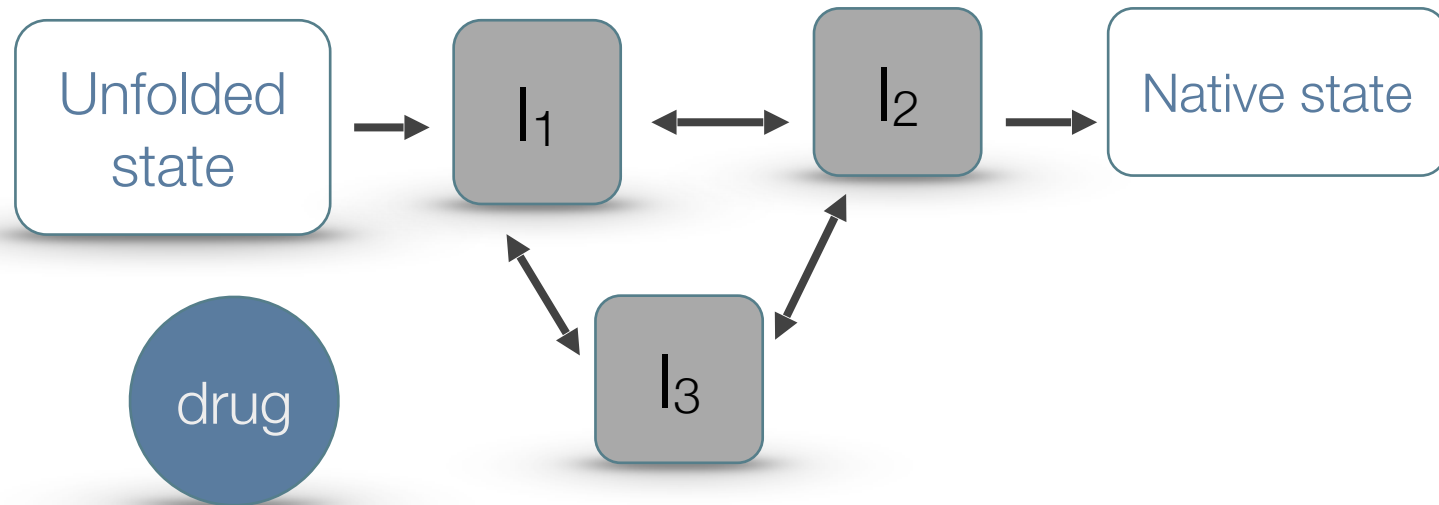


PHARMACOLOGICAL PROTEIN INACTIVATION BY FOLDING INTERMEDIATE TARGETING

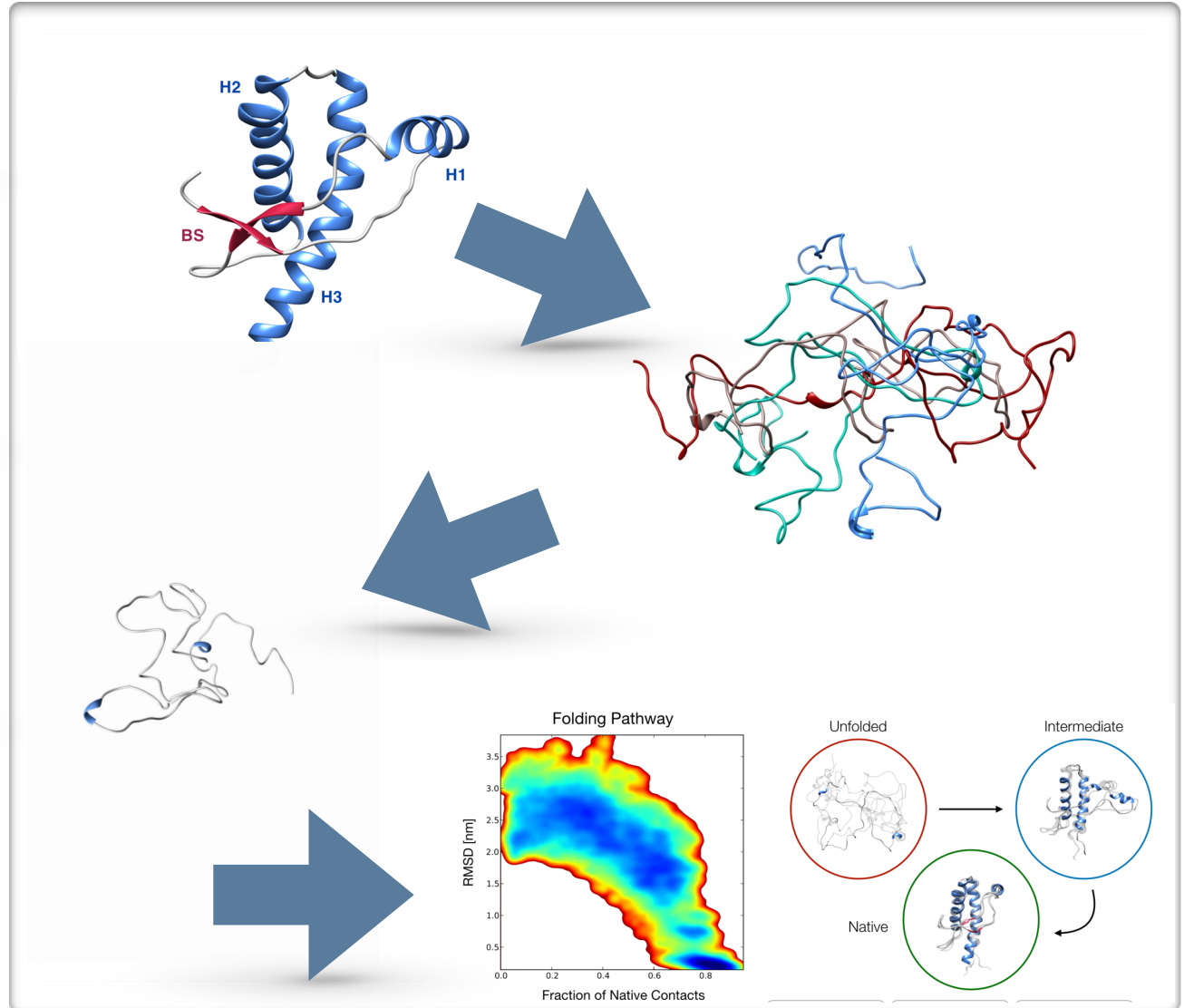
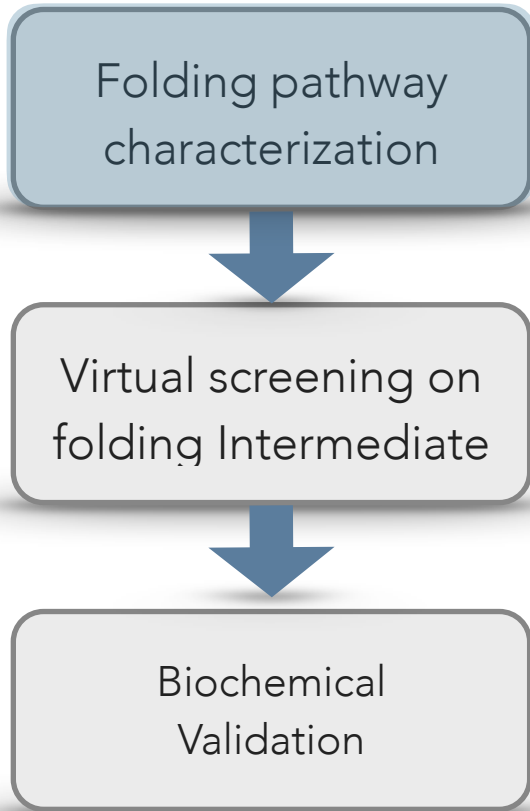
patent file # 102018000007535 (with E. Biasini)



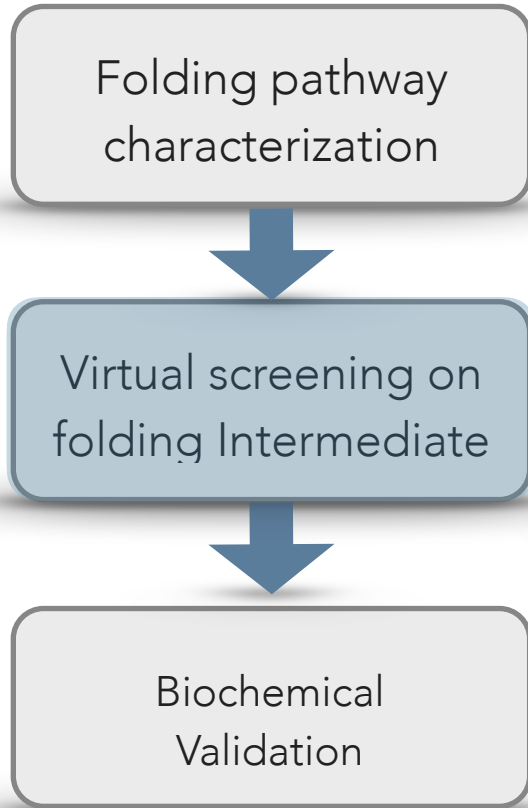
PHARMACOLOGICAL PROTEIN INACTIVATION BY FOLDING INTERMEDIATE TARGETING



PPI-FIT PIPELINE



PPI-FIT PIPELINE



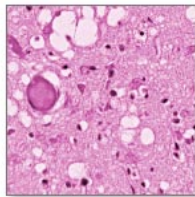
with
L. Barreca's Lab

The image shows a 3D ribbon diagram of a protein structure, colored in gold and blue, set against a black background. A blue mesh overlay is visible on the protein structure, highlighting a specific region. The protein is shown in a ribbon representation, with the main chain in gold and a specific region highlighted in blue. A blue arrow points from a black box above to this 3D structure.

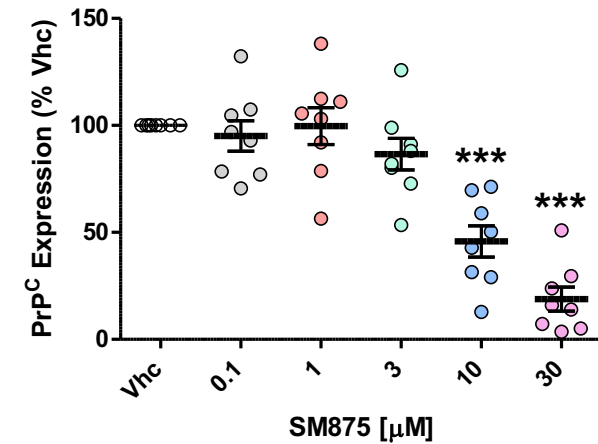
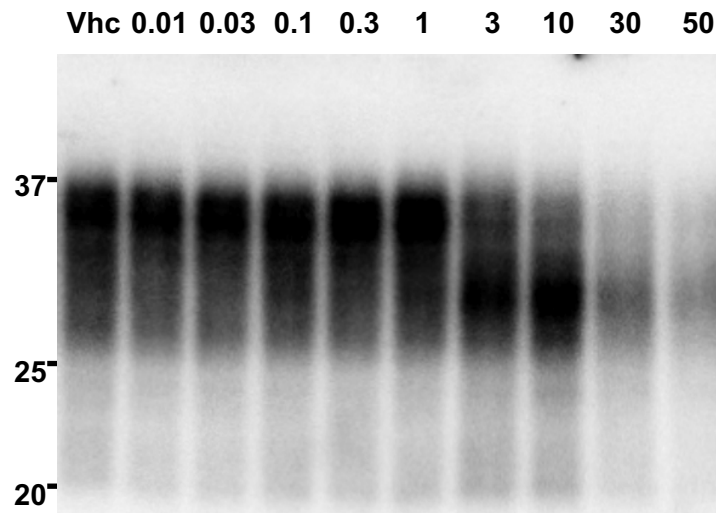
A FIRST VALIDATION

Inactivation of Cellular Prion protein

Brain shrinkage and deterioration occurs rapidly



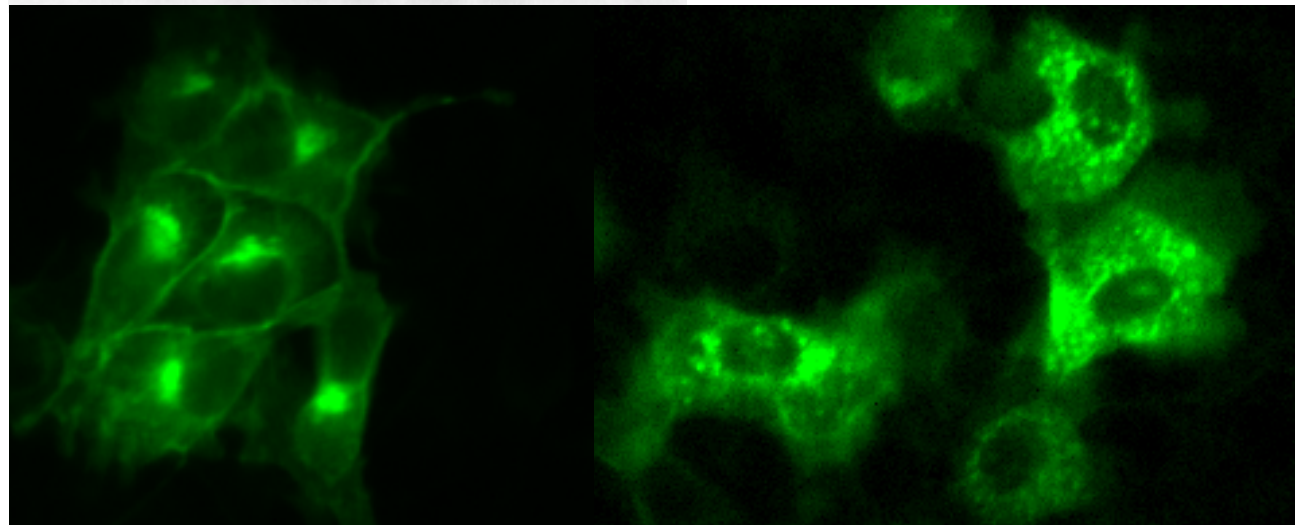
Brain section showing spongiform pathology characteristic of Creutzfeldt-Jakob



PHARMACOLOGICAL PROTEIN INACTIVATION BY TARGETING FOLDING INTERMEDIATES

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doi: <https://doi.org/10.1101/2020.03.31.018069>



Technology Transfer Initiative



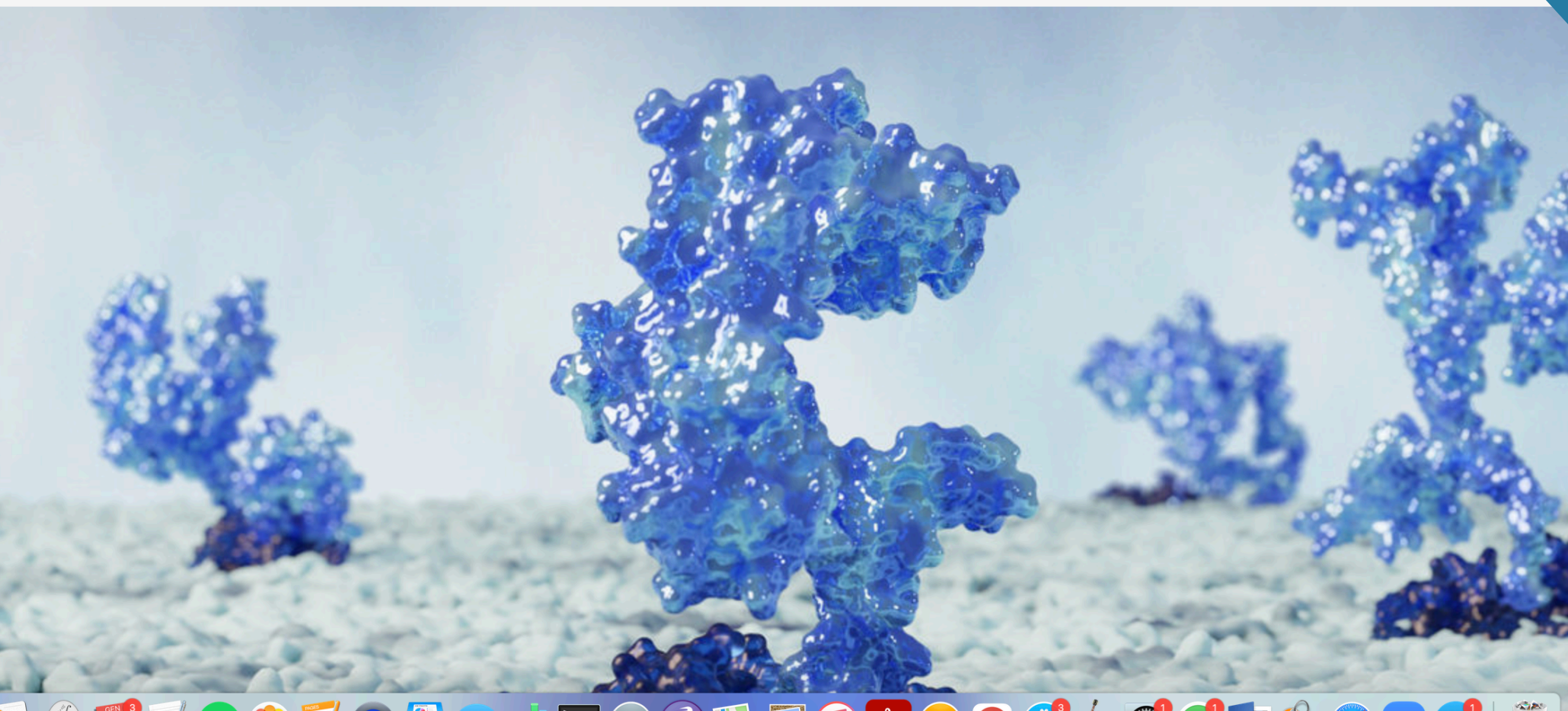
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A few facts about Siybilla Biotech



THE SPINOFF PRIZE

A NATURE RESEARCH AWARD IN PARTNERSHIP WITH:

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OUTLOOK | 24 June 2021 | Correction [08 July 2021](#)

Turning transient structures into drug targets

Start-up Sibylla Biotech has developed a drug-discovery platform to look for protein folding intermediates to target therapeutically.

Future

Sibylla is closing its Series-A round!
(formal announcing expected in the Summer)

Sibylla is **currently hiring** several computational
physicists!

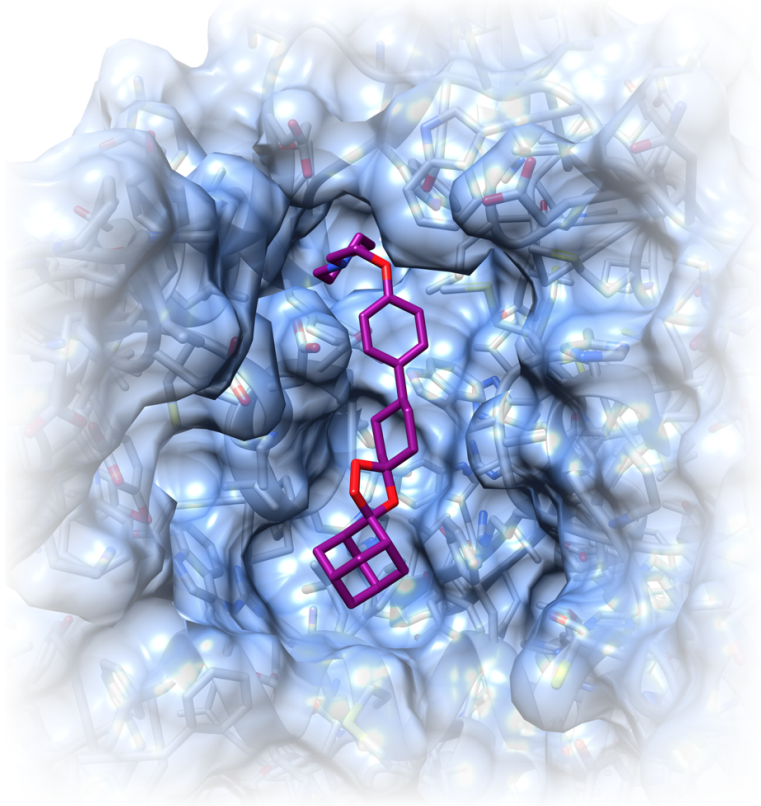
WHAT DO NEW/EMERGING TECHNOLOGIES HAVE TO OFFER?



Partners:

U. Trento, Space Pharma, CJD Foundation (Israel), U. Tel Aviv, U. Santiago de Compostela, INFN

A MAIN LIMITING FACTOR



Impossible to crystallize
folding intermediates
on Earth



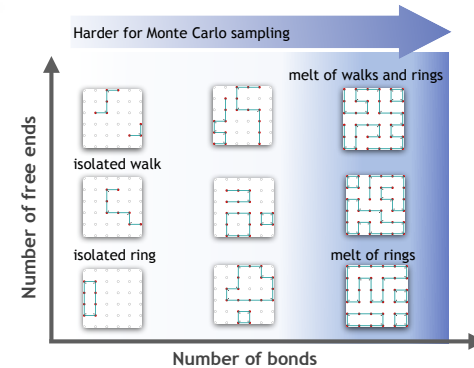
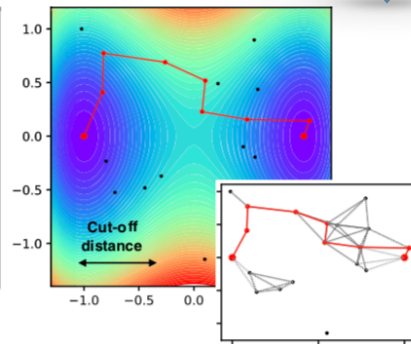
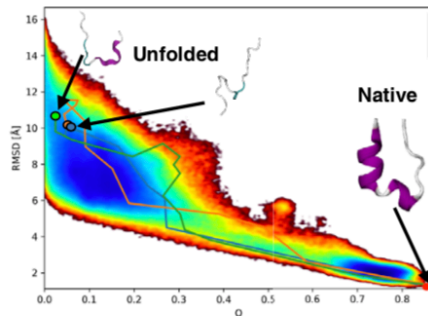
Microgravity
conditions may
provide the solution!

WHAT DO NEW/EMERGING TECHNOLOGIES HAVE TO OFFER?

Quantum Computing

Theoretical Physics

Artificial Intelligence

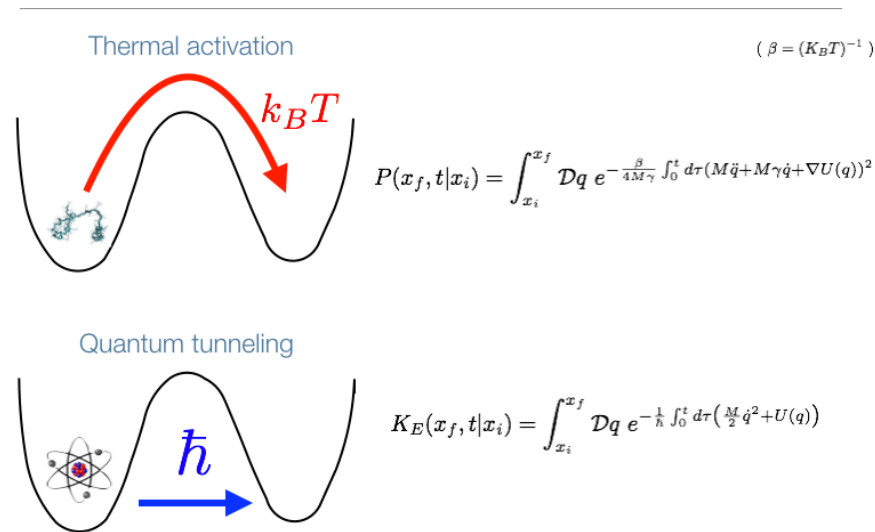


Partners:

U. Trento, Q@Trento, INFN, SISSA, BEC-CNR

Final considerations (very subjective!)

Fundamental science can breed new ideas



Cross-disciplinarity is key to tackle **complexity**.

Seek for colleagues with different background...

.....and **learn to talk a lot!**

Final considerations (very subjective!)

If you have a good idea... **”money is not the limiting factor issue”!**

Patenting is not the Enemy of Science! ...

..but just don't wait too much to look out !!!

Acknowledgments



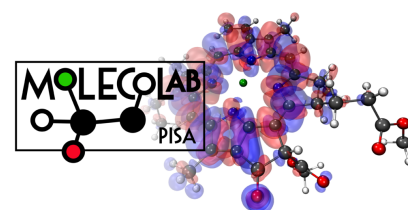
Trento Institute for
Fundamental Physics
and Applications



Italy:

Trento: E. Biasini, A. Ianeselli (2014-2017), G. Spagnoli S. A Beccara (2009-2017), S. Orioli (2014-2018), E. Schneider (2012-2015), M. Carli (2017), M. Turelli (2018), F. Mascherpa (2014), *G. Garberoglio, F. Pederiva, M. Sega*, R. Covino (2012-2015),

Pisa: B. Mennucci, L. Cupellini, S. Jurinovich



Perugia: L. Barreca

SISSA: C. Micheletti, A. Laio

Europe:

U. Zurich: B. Schuler

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CEA-Saclay: H. Orland

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U. Mass.: A. Gershenson