

The Muon g-2 Experiment at Fermilab: measuring the magnetic field

Anna Driutti

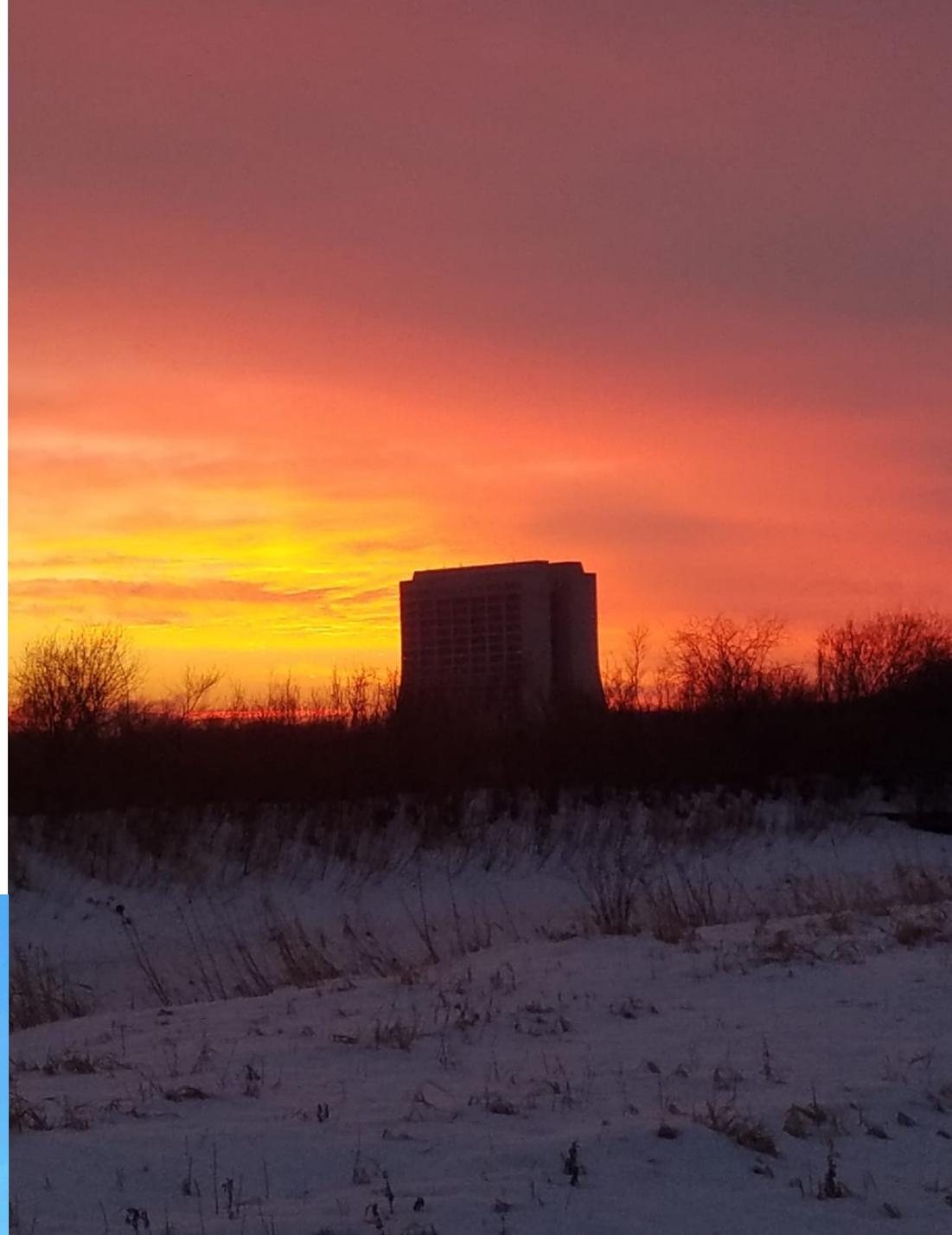
University and INFN Pisa



UNIVERSITÀ DI PISA



Fermilab 2022 Summer School
Pisa, 19 July 2022



Fermilab Muon g-2 Experiment

Muon anomaly is determined with:

$$a_\mu = \underbrace{\frac{\omega_a}{\tilde{\omega}'_p(T_r)}}_{\text{ratio of frequencies (} R_\mu \text{) measured by us}} \underbrace{\frac{\mu'_p(T_r)}{\mu_e(H)} \frac{\mu_e(H)}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}}_{\text{fundamental factors (combined uncertainty 25 ppb)}}$$

ratio of frequencies (R_μ)
measured by us

fundamental factors
(combined uncertainty 25 ppb):

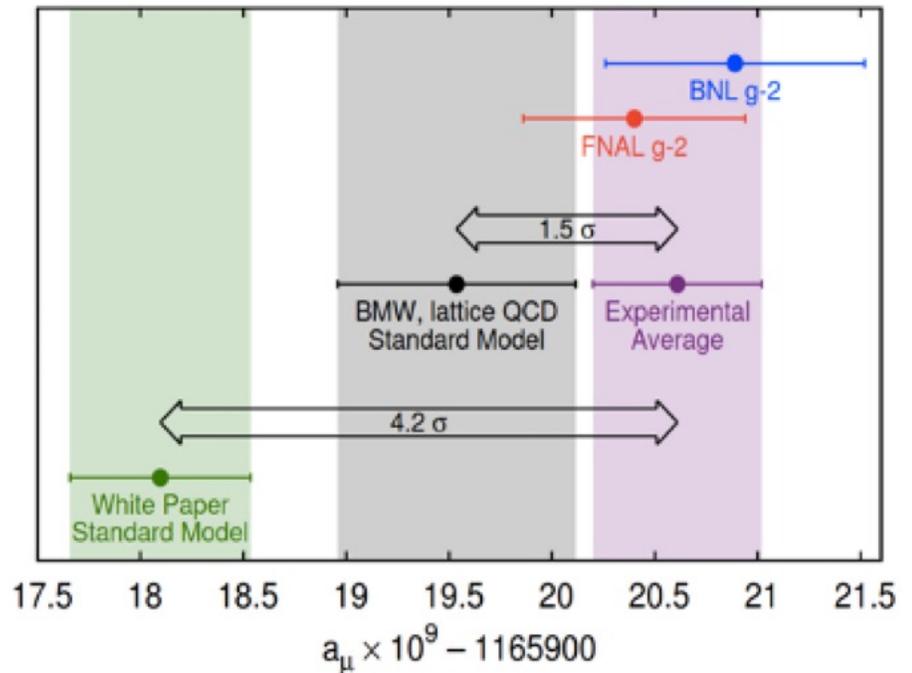
- $\mu'_p(T_r)/\mu_e(H)$ from [Metrologia **13**, 179 (1977)]
- $\mu_e(H)/\mu_e$ from [Rev. Mod. Phys. **88** 035009 (2016)]
- m_μ/m_e from [Phys. Rev. Lett. **82**, 711 (1999)]
- $g_e/2$ from [Phys. Rev. A **83** 052122 (2011)]

ω_a : muon anomalous precession frequency

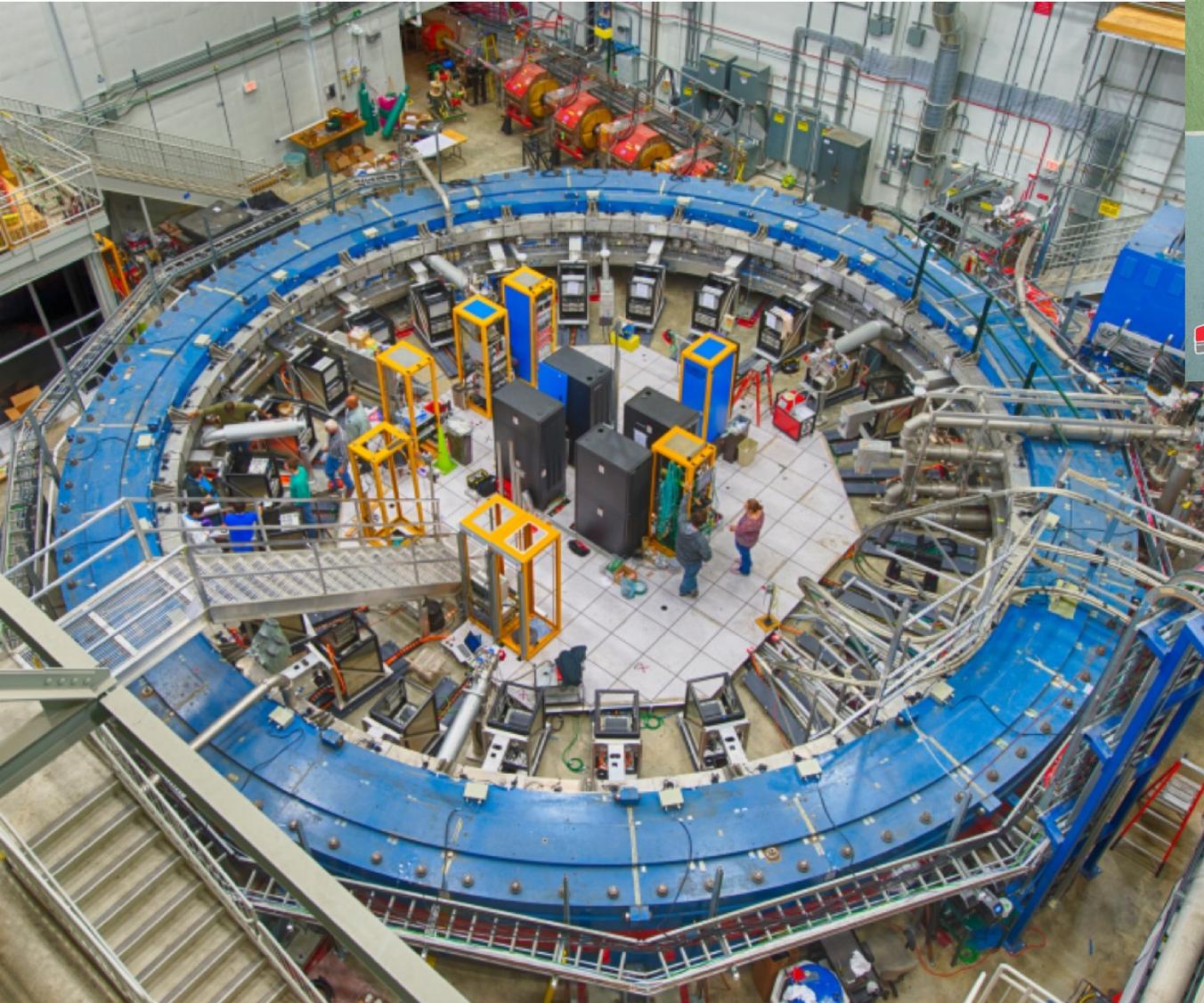
Lorenzo ed Elia

$\tilde{\omega}'_p(T_r)$: magnetic field B in terms of (shielded) proton precession frequency **and** weighted by the muon distribution
(shielded = measured in a spherical water sample at the reference temperature $T_r = 34.7^\circ\text{C}$)

In questo talk



The Muon g-2 Storage Ring



The Muon g-2 Superconducting C shaped magnet

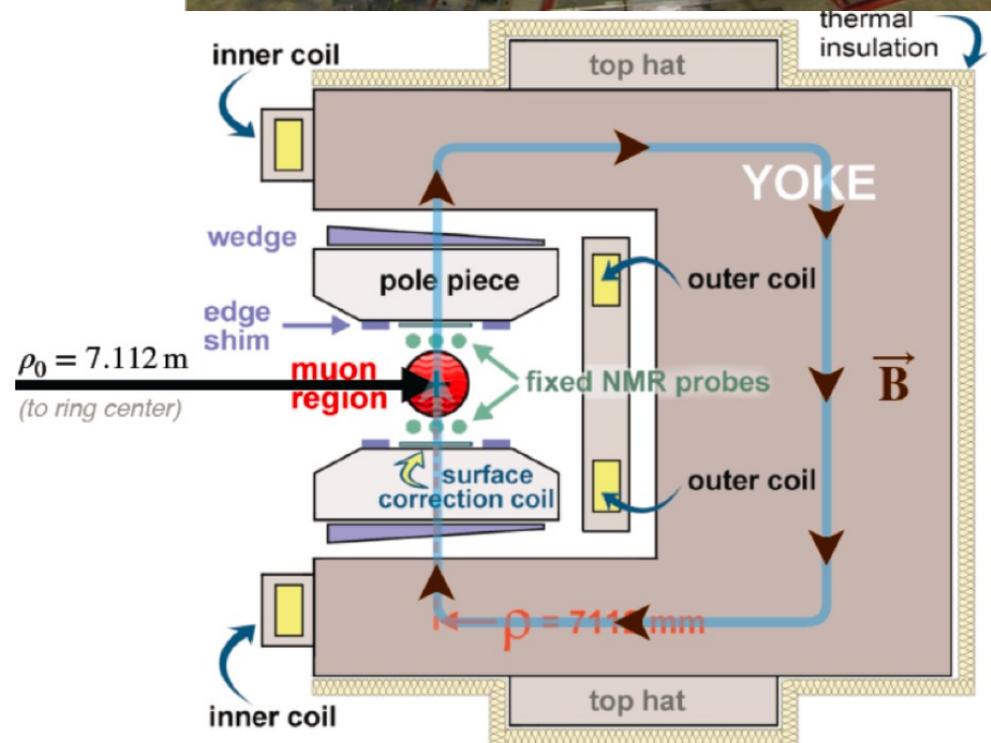
Provides 1.45T vertical and uniform B field:

- **12 iron Yokes:** open on the inside to allow positrons from muon decay to reach the detectors and excited by superconducting coils made of Niobium-Titanium in a Copper matrix
- **72 low-carbon steel poles:** to minimize impurities
- **144 Edge shims:** local sextupole field minimized by changing edge shim thickness
- **864 Steel wedges:** allowed for angle adjustment (compensate quadrupole component) and radial adjustment (shim local dipole field)
- **Surface correction coil:** help to reduce non-uniformities on field higher moment

Fermilab: ~15 ppm RMS (~75 ppm peak-to-peak)
BNL E821: ~35 ppm RMS (~200 ppm peak-to-peak)



Achieved ~3x better field uniformity than at BNL



We measure the ratio of two frequency!

f_{clock} : blinded clock
 ω_a^{meas} : measured precession frequency

ω_a

Lorenzo

Elia

beam dynamics corrections

C_e : electric field correction
 C_p : pitch correction
 C_{ml} : muon loss correction
 C_{pa} : phase-acceptance correction

$$R_\mu = \frac{f_{clock} \cdot \omega_a^{meas} \cdot (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{calib} \cdot \omega'_p(x, y, \phi) \otimes M(x, y, \phi) \cdot (1 + B_k + B_q)}$$

Paolo

$\tilde{\omega}'_p(T_r)$

In questo talk

field corrections

f_{calib} : absolute magnetic field calibration
 $\omega'_p(x, y, \phi)$: field maps
 $M(x, y, \phi)$: muon beam distribution

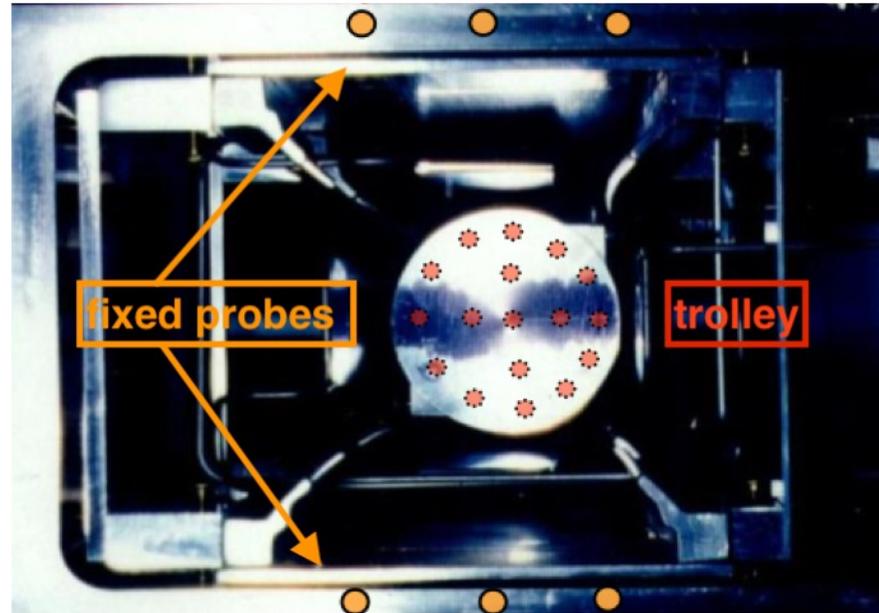
B_k : transient field from eddy current in kicker
 B_q : transient field from quad vibration

The Field Maps

$$R_{\mu} = \left(\frac{f_{clock} \cdot \omega_a^{meas} \cdot (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{calib} \cdot \omega'_p(x, y, \phi) \otimes M(x, y, \phi) \cdot (1 + B_k + B_q)} \right)$$

2 types of field probes

- 378 fixed NMR probes above and below storage region
→ measure B-field 24/7

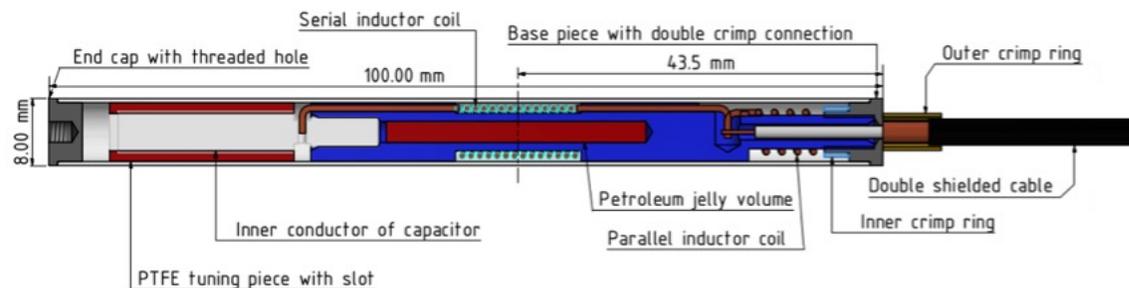


- Trolley with 17-probe NMR
→ 2D profile of B over the entire azimuth when beam is OFF



Electronics

NMR probes

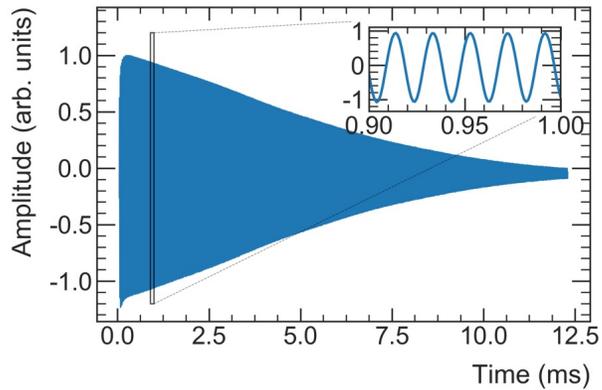


Pulsed proton Nuclear Magnetic Resonance (NMR) probes filled with petroleum jelly

The Field Maps

$$R_{\mu} = \left(\frac{f_{clock} \cdot \omega_a^{meas} \cdot (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{calib} \cdot \omega'_p(x, y, \phi) \otimes M(x, y, \phi) \cdot (1 + B_k + B_q)} \right)$$

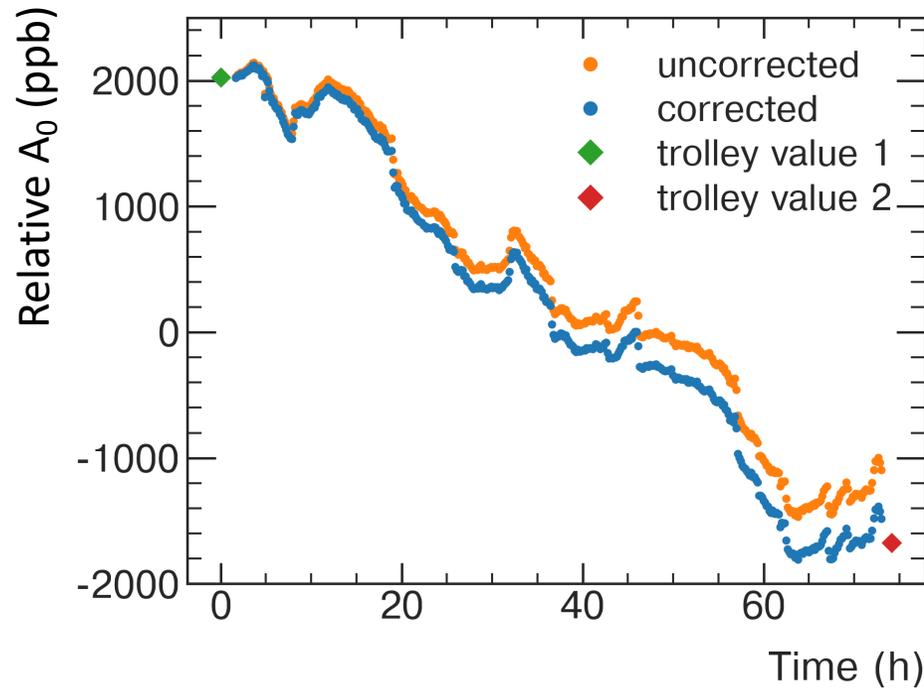
- 378 fixed NMR probes above and below storage region
→ measure B-field 24/7



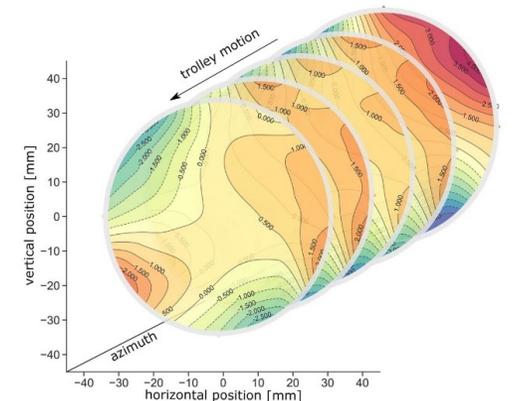
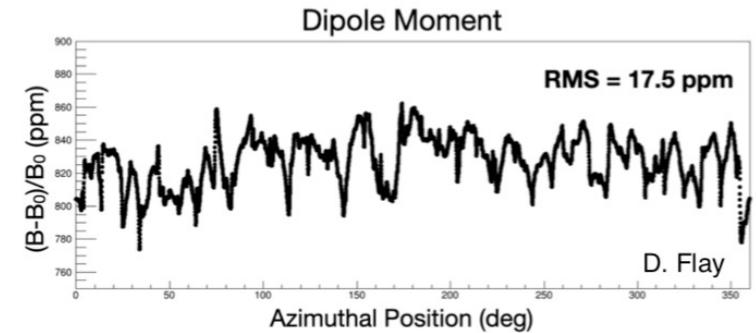
2 types of field probes



Final field from interpolation



- Trolley with 17-probe NMR
→ 2D profile of B over the entire azimuth when beam is OFF

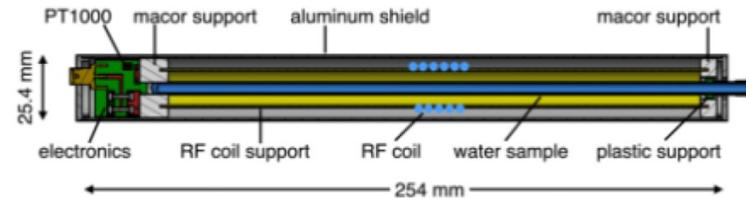
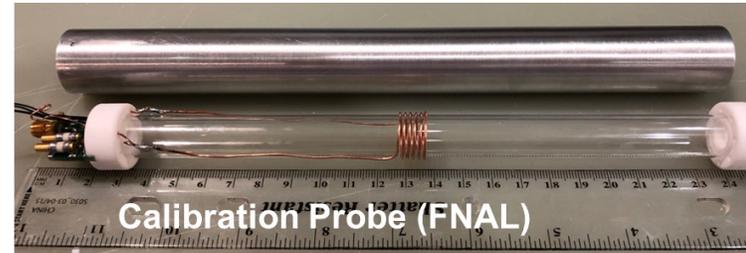


Calibration of the field probes

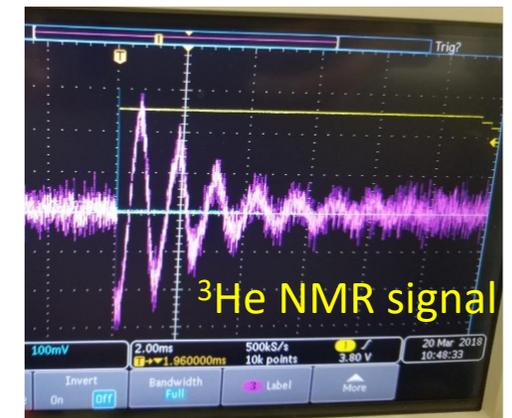
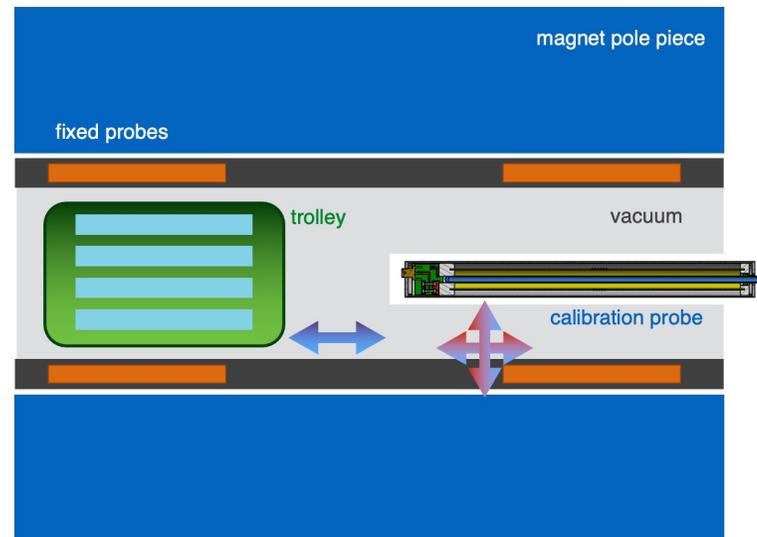
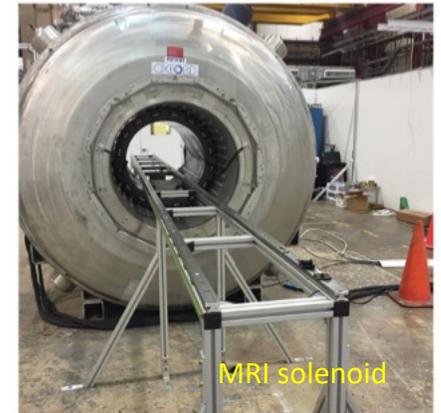
$$R_{\mu} = \left(\frac{f_{clock} \cdot \omega_a^{meas} \cdot (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{calib} \cdot \omega'_p(x, y, \phi) \otimes M(x, y, \phi) \cdot (1 + B_k + B_q)} \right)$$

Absolute field calibration:

- NMR probes calibrated using absolute probes
- Absolute probes provide field in terms of proton shielded in a spherical sample of water at an exact temperature: $\omega_p(Tr)$
- Corrections due to sample shape, temp and magnetization from studies in an MRI solenoid
- Trolley probes calibrated inside the ring
- Cross-check provided by a novel ^3He NMR probe was developed with different systematics



ω_p : free proton precession frequency
Using proton NMR: $\hbar\omega_p = 2\mu_p B$



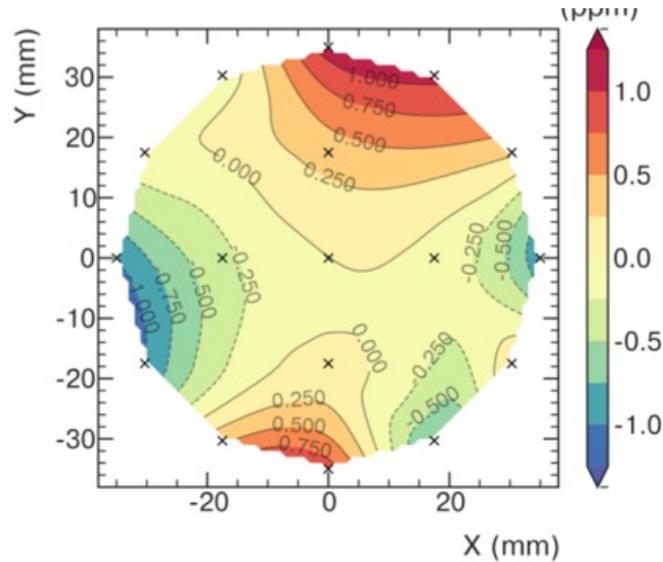
Convolution

$$R_\mu = \left(\frac{f_{clock} \cdot \omega_a^{meas} \cdot (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{calib} \cdot \omega'_p(x, y, \phi) \otimes M(x, y, \phi) \cdot (1 + B_k + B_q)} \right)$$

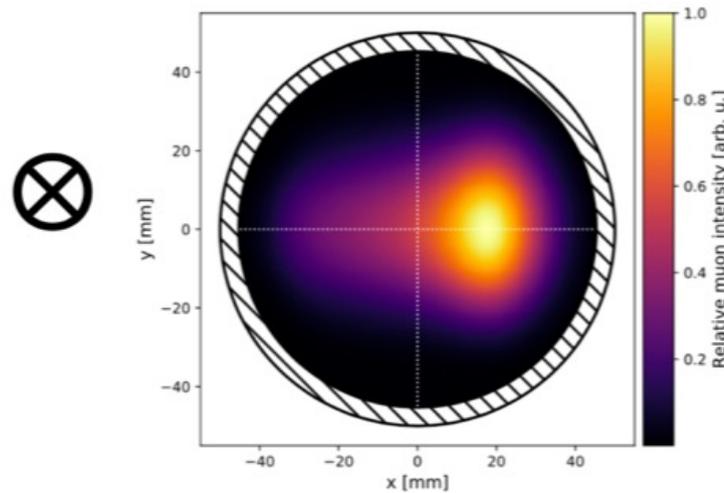
Field maps are weighted by beam distribution

$$\tilde{\omega}_p = \left\langle \frac{\int \omega_p(x, y, \phi) M(x, y, \phi) dx dy}{\int M(x, y, \phi) dx dy} \right\rangle$$

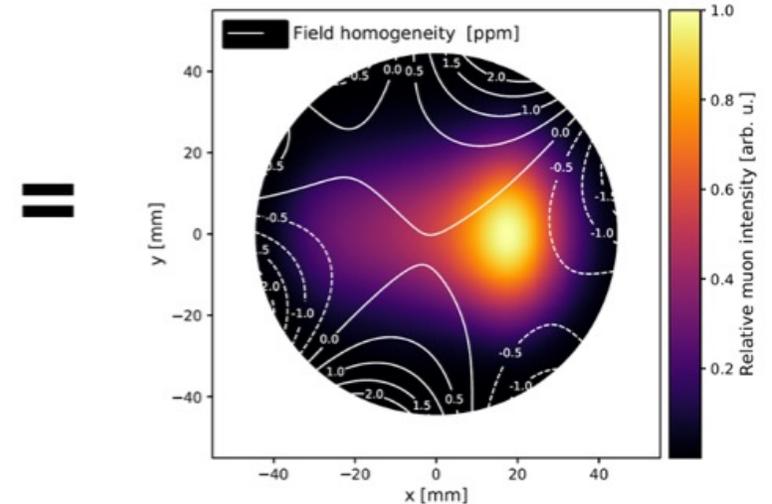
Field from Trolley and fixed probe



Beam distribution extrapolated from trackers data propagated around the ring with simulations



Field seen by the muons (tot Run1 unc. 0.56 ppm)



Quads transient field Correction

$$R_{\mu} = \left(\frac{f_{clock} \cdot \omega_a^{meas} \cdot (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{calib} \cdot \omega'_p(x, y, \phi) \otimes M(x, y, \phi) \cdot (1 + B_k + B_q)} \right)$$

- due to mechanical vibrations from pulsing the quads
- mapped using special NMR probes

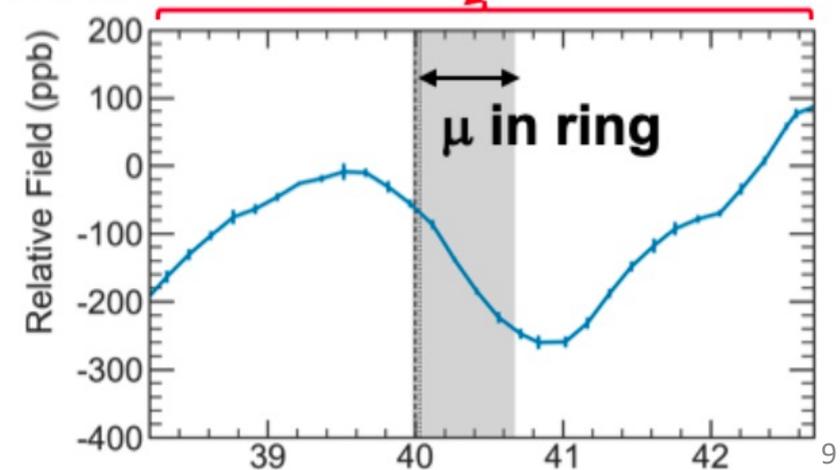
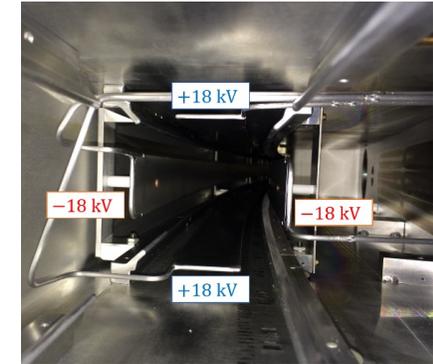
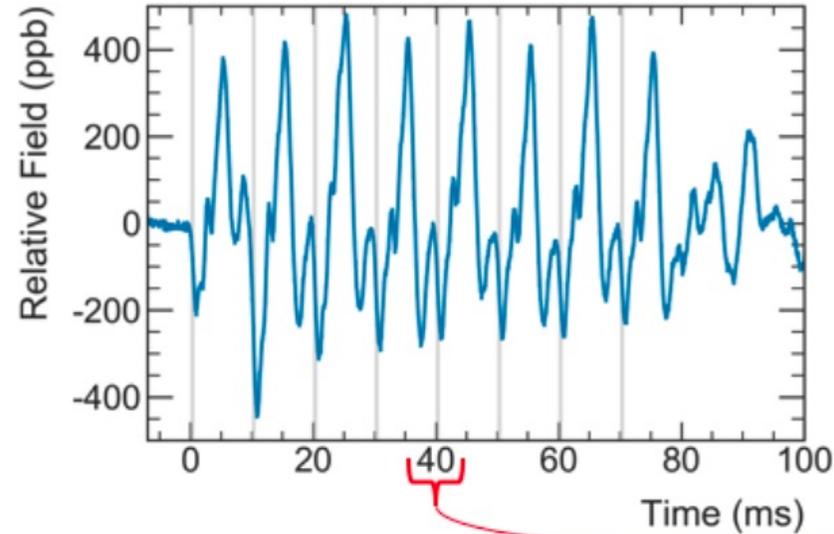


Results for Run1:

$$B_q \sim 17 \text{ ppb} \quad \delta_{B_q} \sim 92 \text{ ppb}$$

ppb = parti per miliardo (10^9)

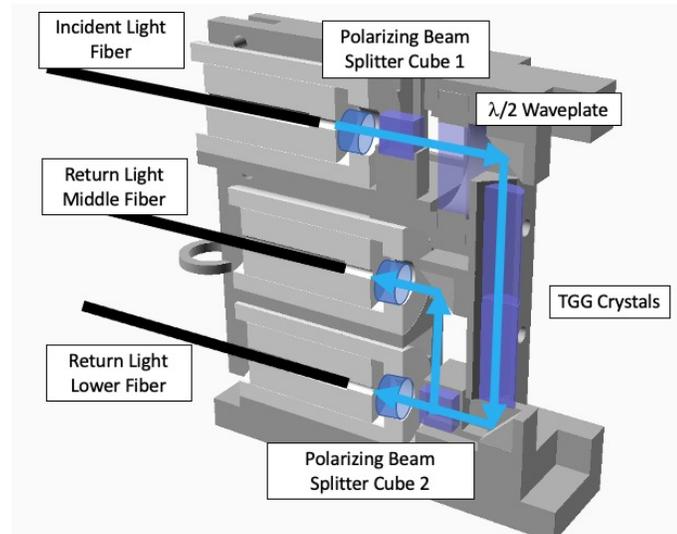
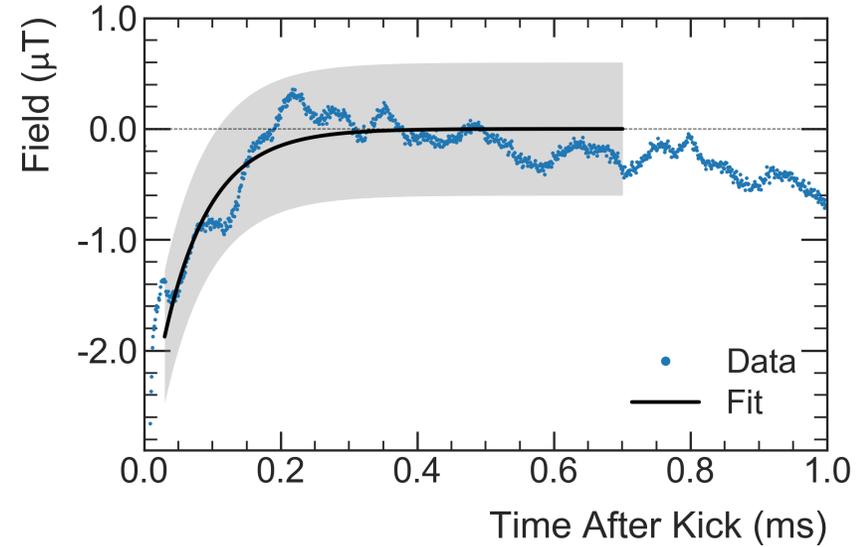
- δ_{B_q} dominated by incomplete map
- expected to be reduced by factor 2 for Run 2 and after



Kicker transient field Correction

$$R_{\mu} = \left(\frac{f_{clock} \cdot \omega_a^{meas} \cdot (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{calib} \cdot \omega'_p(x, y, \phi) \otimes M(x, y, \phi) \cdot (1 + B_k + B_q)} \right)$$

- Due to eddy currents induced by the fast magnetic pulses in surrounding materials
- Measured using a magnetometer and cross-checked with second magnetometer



Results for Run1:

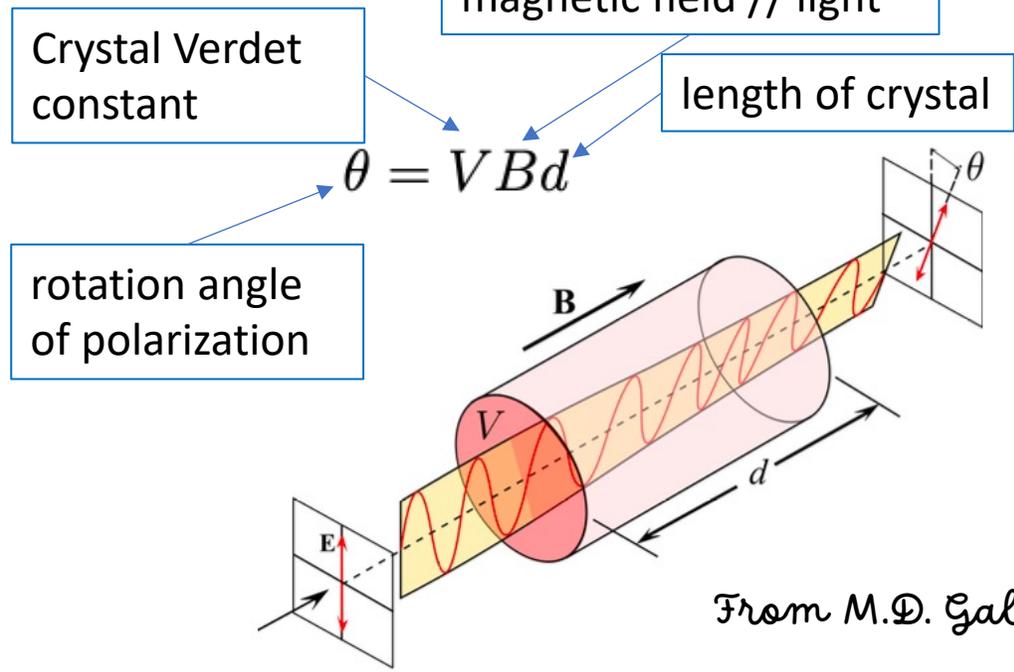
$$B_k \sim 30 \text{ ppb} \quad \delta_{B_k} \sim 40 \text{ ppb}$$

ppb = parti per miliardo (10^9)

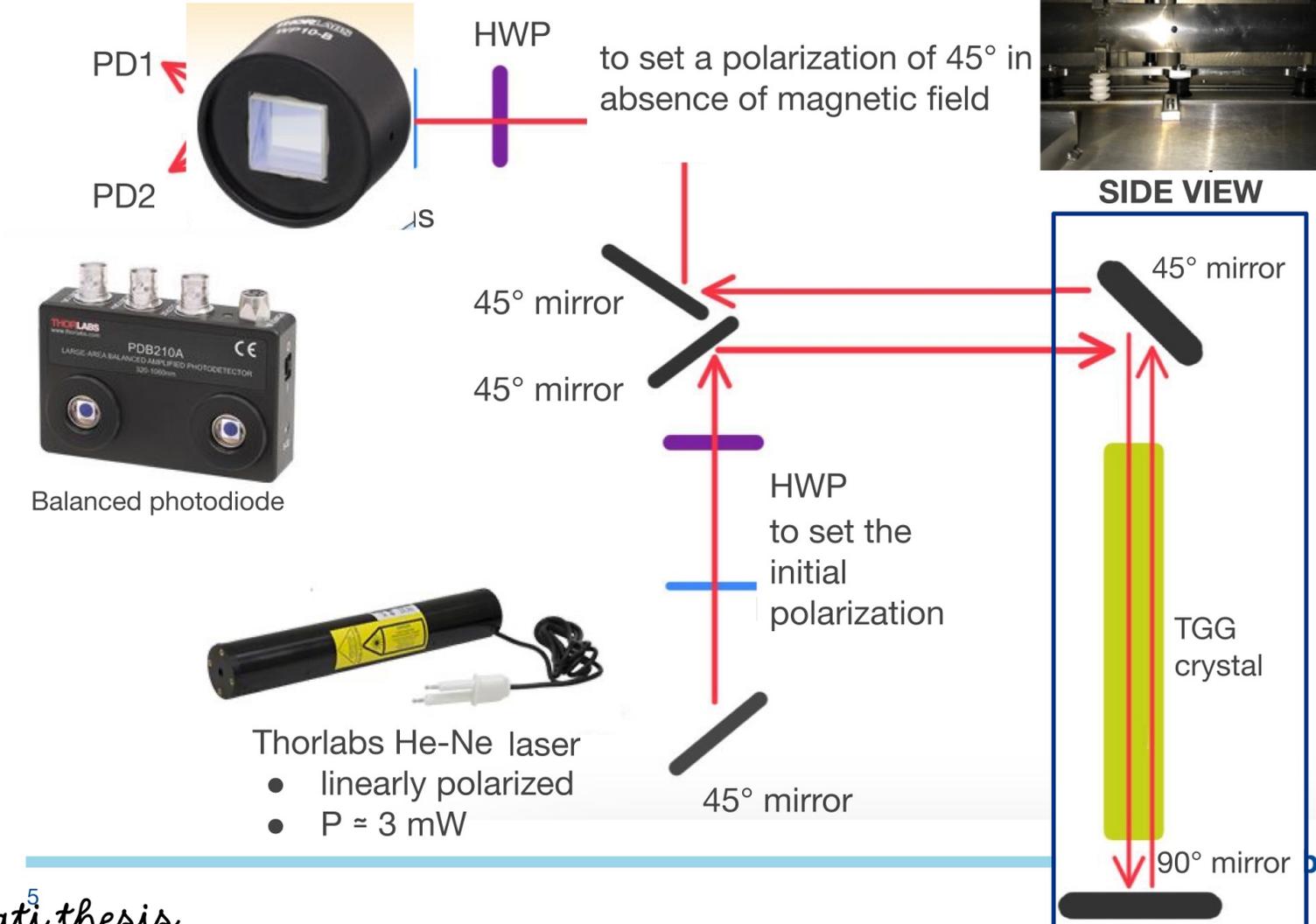
The INFN breadboard magnetometer

- Aiming at few mG sensitivity :
 - Run2/3: **30ppb (5mG)**
 - Run4/5: **10ppb (2mG)**

- Based on Faraday observation that an external magnetic field influences light propagating in materials:



Setup scheme (upper view)



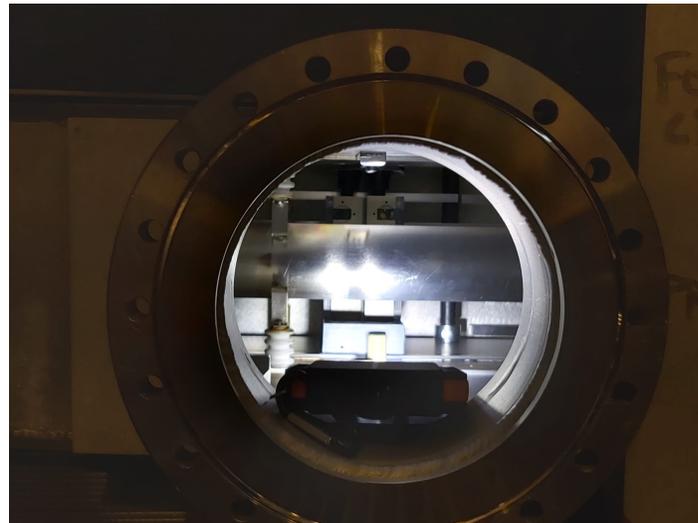
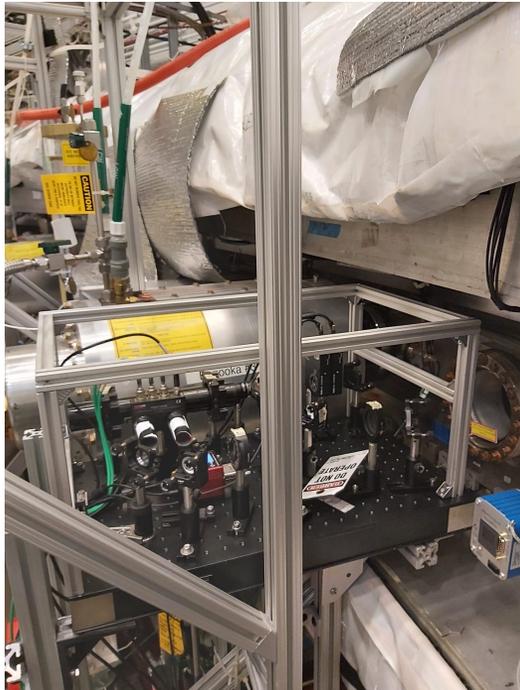
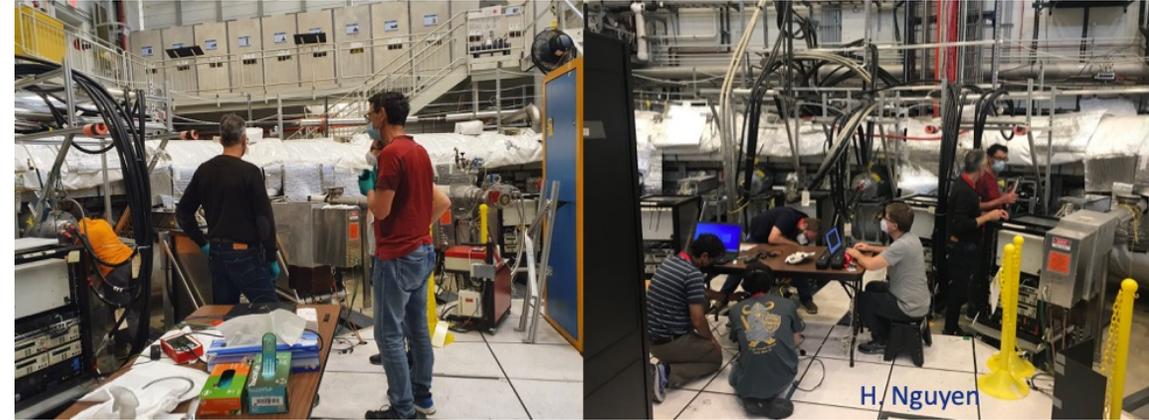
The INFN breadboard magnetometer in Pisa

- 2020: Built, characterized and tested in Pisa

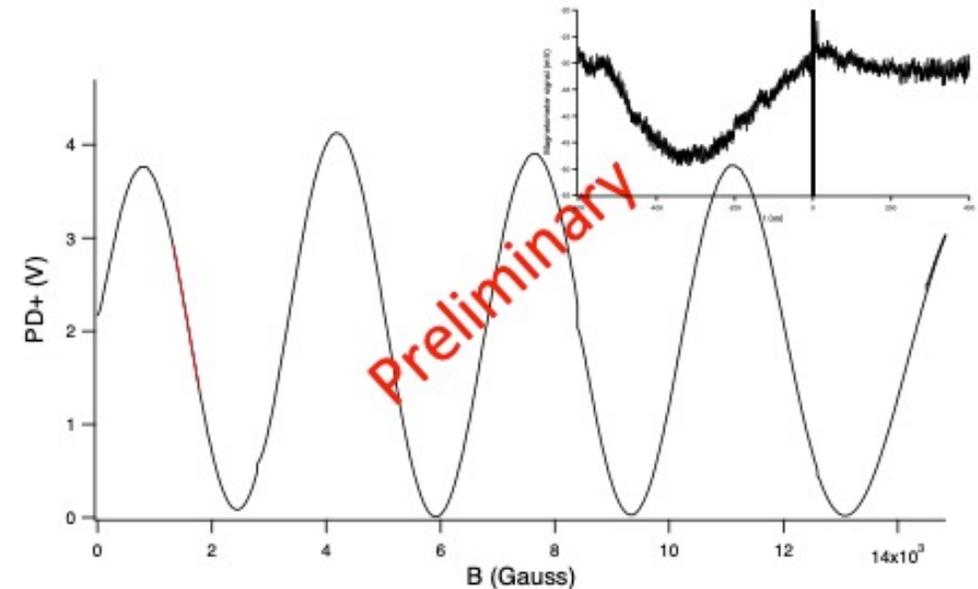


The INFN breadboard magnetometer at FNAL

- Summer 2021: Installed and tested with magnet off
- Summer 2022: installed and tested with magnet on (last week) – planned more measurements during the summer shutdown



A. Gioiosa

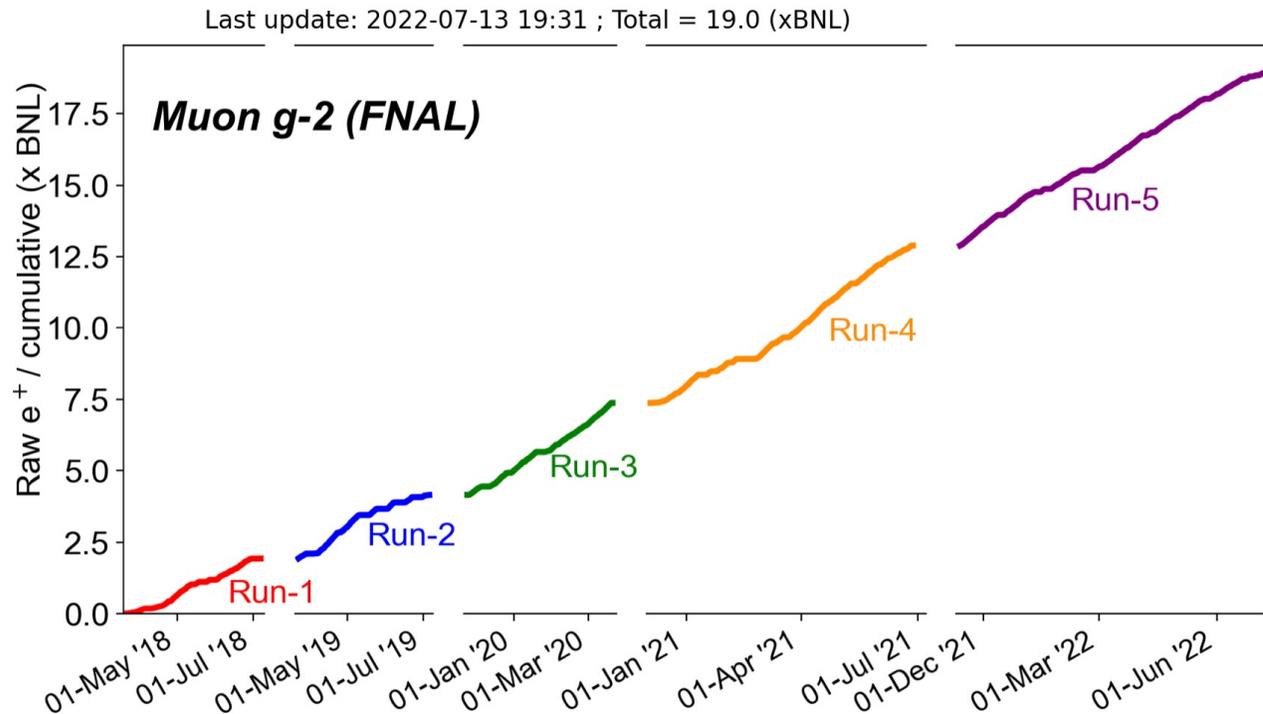


M. Sorbara (OPS meeting)

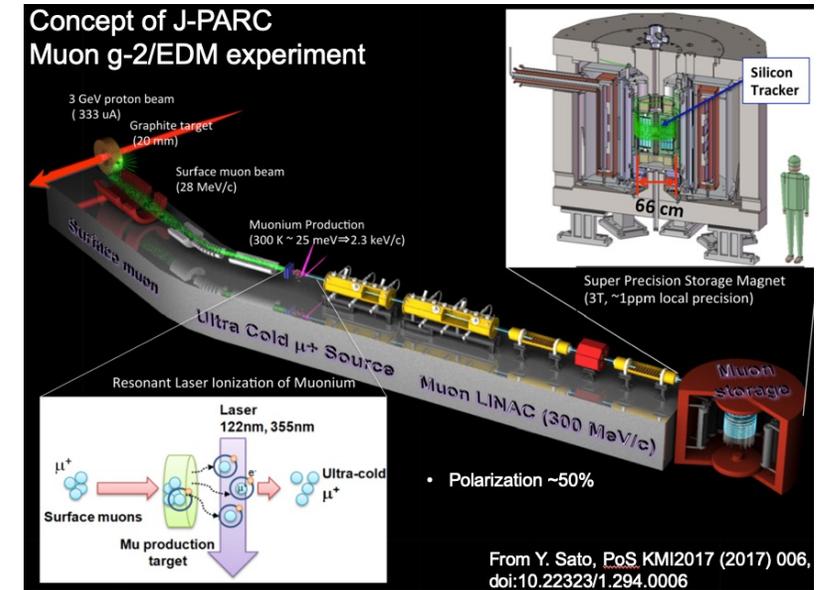
Muon g-2 experiment at J-PARC

What's next?

Analysis of data from Run 2-5

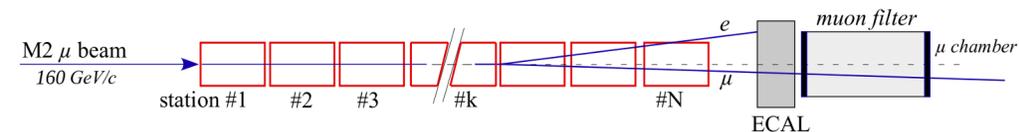
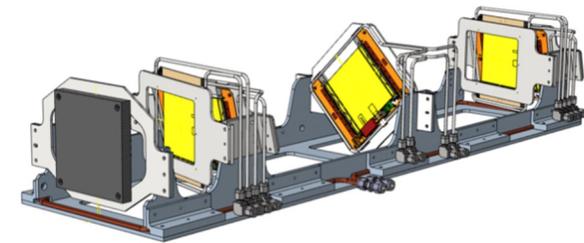


- Run-2** and **Run-3** analysis are ongoing:
 - > reduce combined exp. error by 2 times
- Run-4** + **Run-5** (ended beginning of July)
 - > stat. uncert. down to 100 ppb in total



MUonE experiment at CERN

$$a_{\mu}^{HLO} = \frac{\alpha_0}{\pi} \int_0^1 dx (1-x) \Delta\alpha_{had}[t(x)] \quad (\text{spacelike approach})$$





Thank you!