

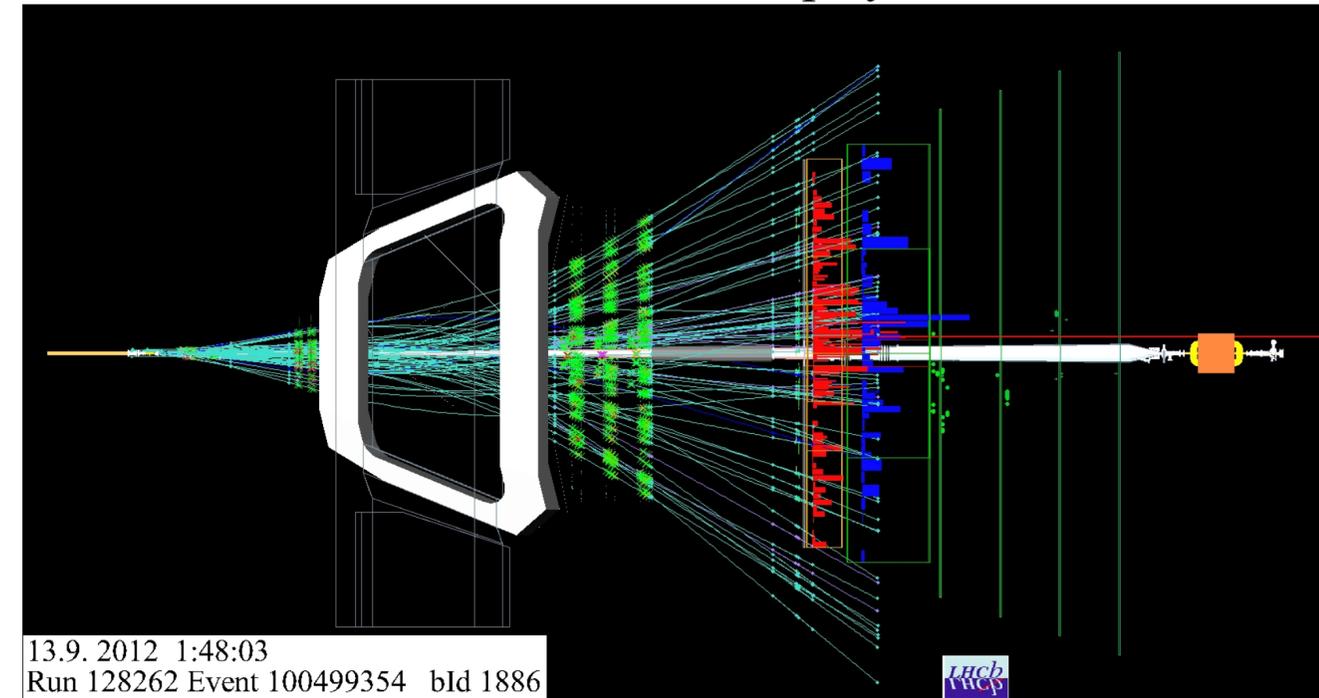
# W mass measurement at LHCb

Lorenzo Sestini - INFN Padova

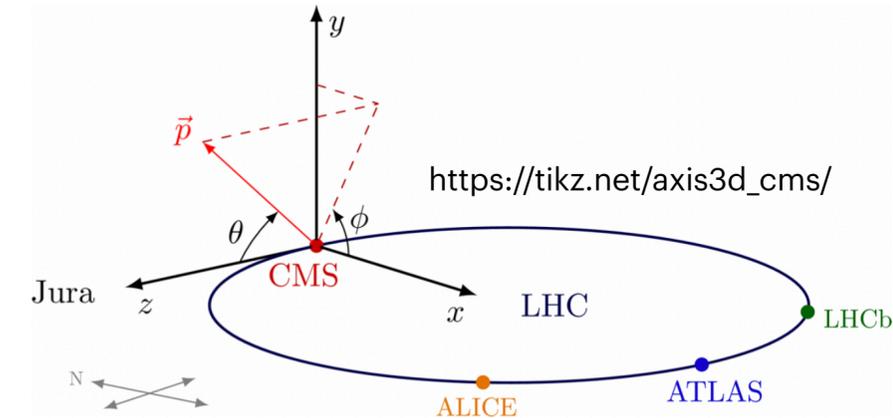
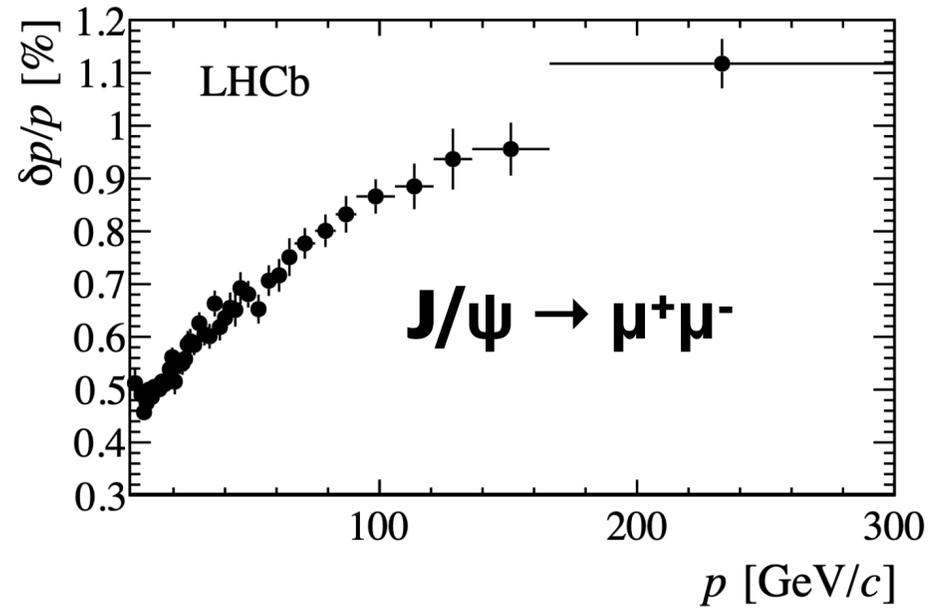
# Introduction

- The environment of proton-proton colliders is probably the **most challenging** for the W boson mass measurement
- Experimental conditions are not as clean as at lepton colliders
- The W boson production modeling is more under control at proton-anti-proton colliders
- **LHCb** has been designed as a flavour physics experiment, but it has developed a wide and rich **electroweak program**
- At LHCb millions of W bosons can be analyzed to extract the W boson mass measurements, **the challenge has been accepted!**

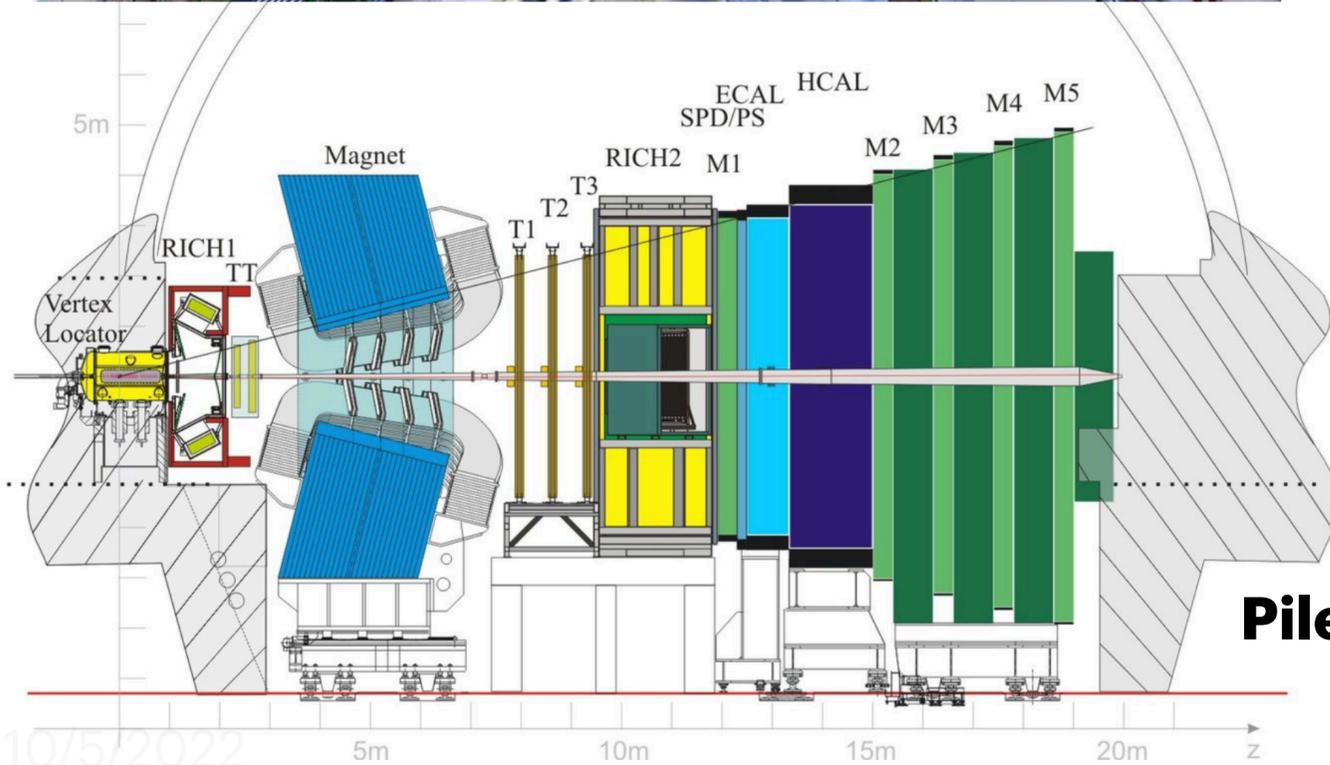
LHCb Event Display



## General purpose detector in the forward region

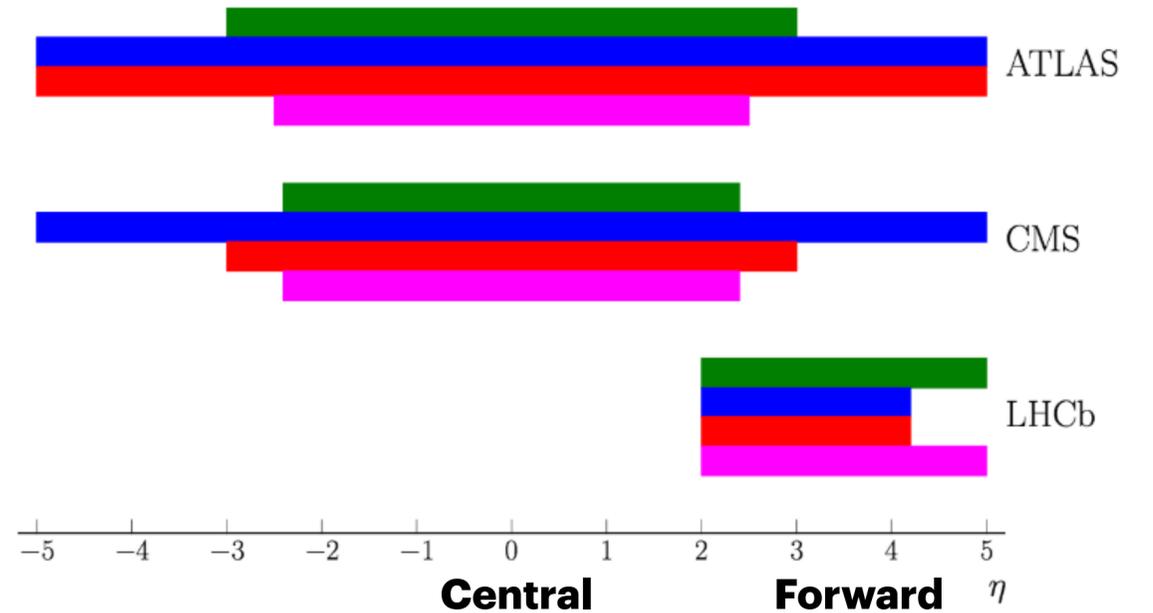


$$\eta \equiv -\ln \left[ \tan \left( \frac{\theta}{2} \right) \right]$$

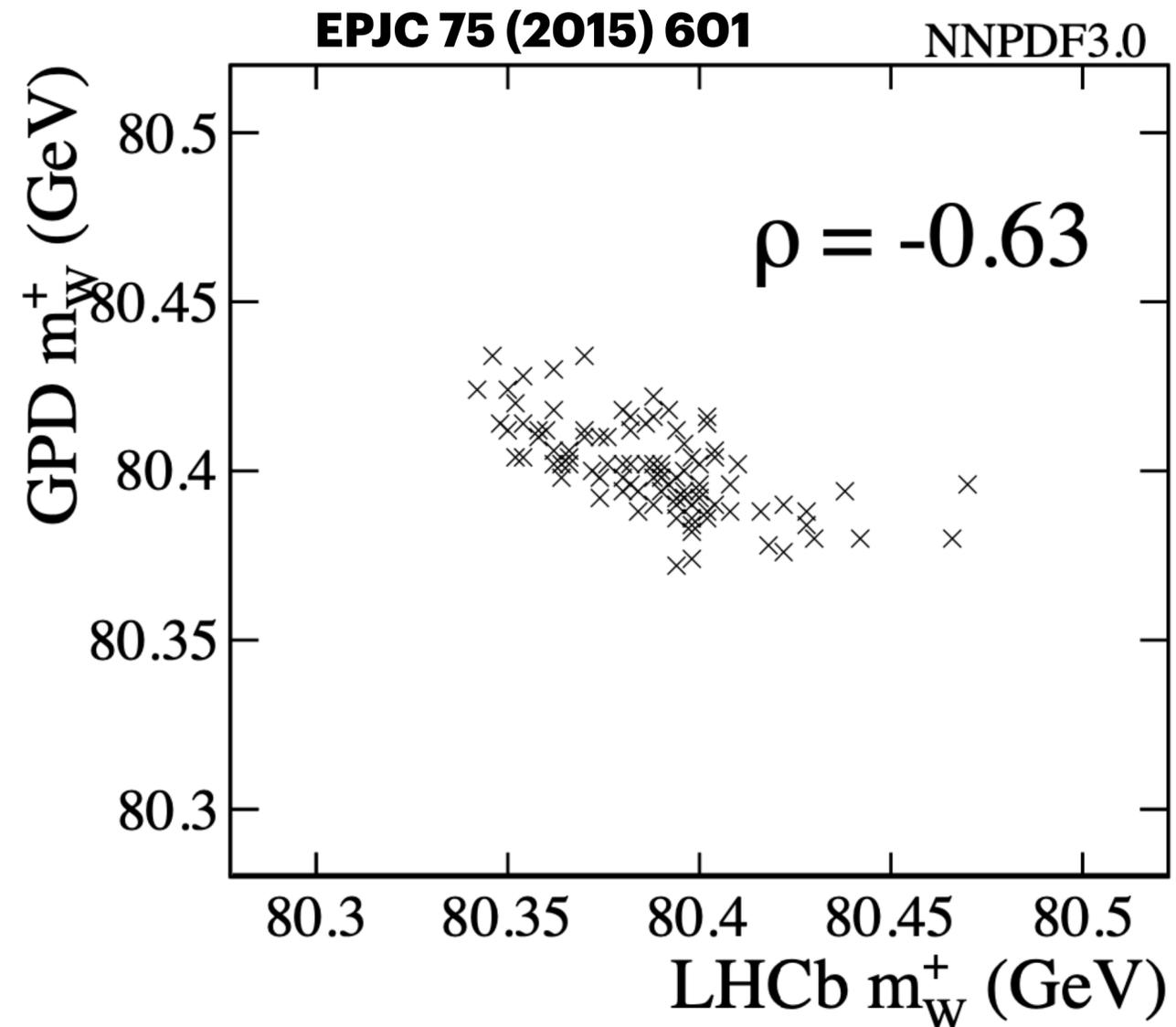


**Pile-up ~1 in Run 1-2**

Lorenzo Sestini



# Complementarity

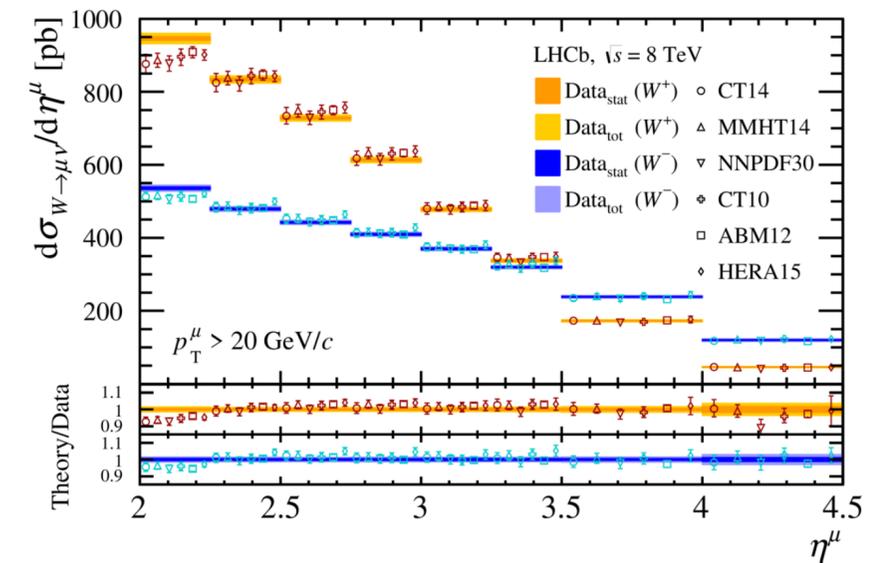
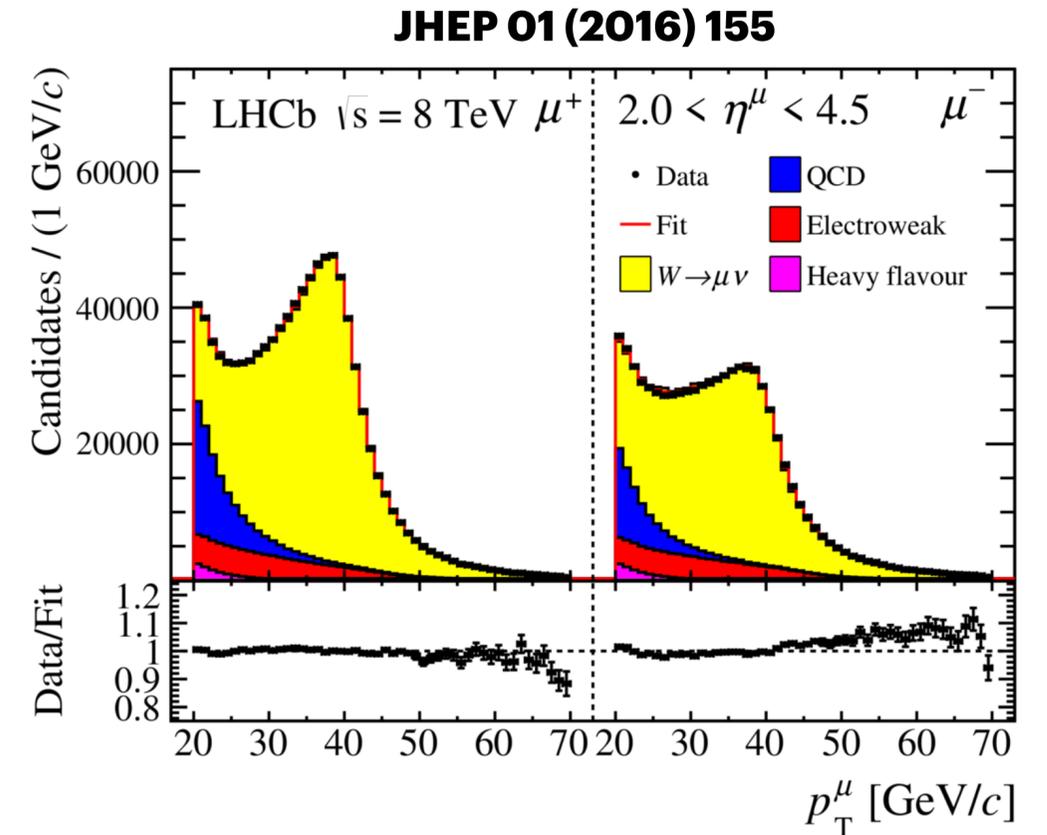
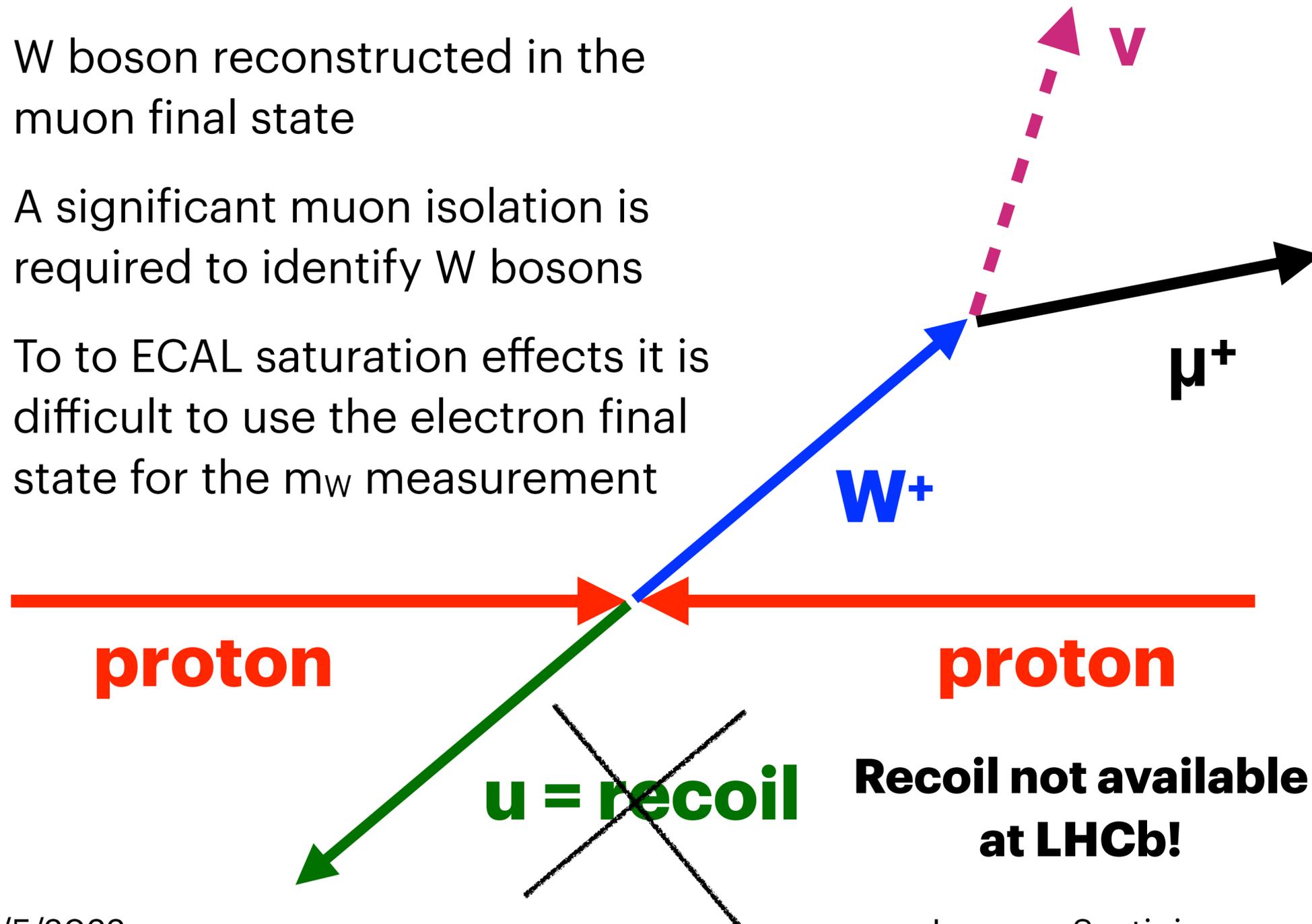


- PDFs uncertainties in the W mass measurement are anti-correlated between the central and forward region
- Combining ATLAS/CMS+LHCb can reduce the PDFs uncertainty
- All the three experiments can significantly contribute in a LHC-wide average
- **The overall average is ultimately the quantity that matters**

**GPD = General Purpose Detectors = ATLAS/CMS**

# W boson identification at LHCb

- W boson reconstructed in the muon final state
- A significant muon isolation is required to identify W bosons
- To to ECAL saturation effects it is difficult to use the electron final state for the  $m_W$  measurement



# W mass measurement at LHCb

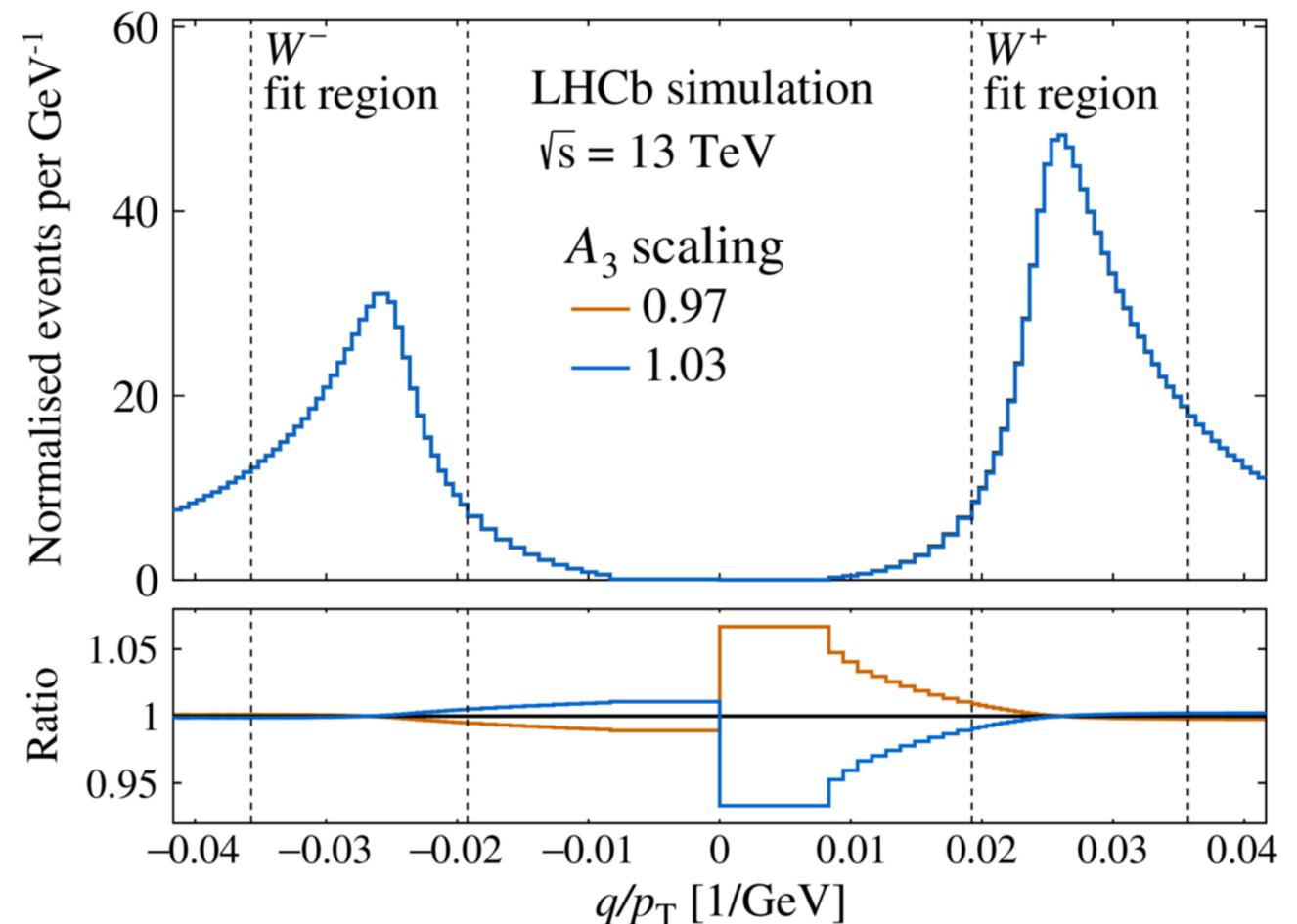
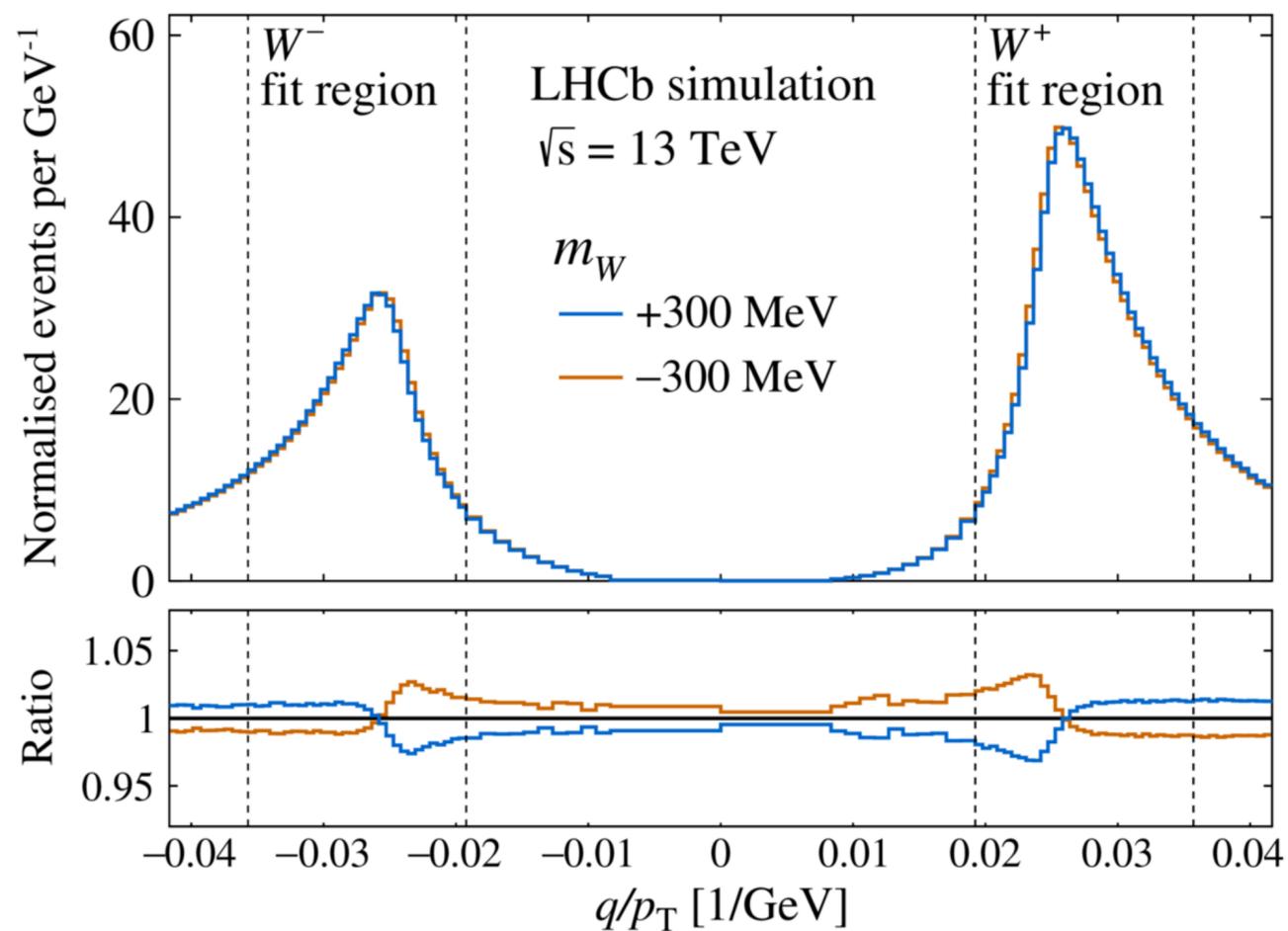
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- Measurement with muon final state, just a part of the Run 2 dataset has been used (1.7 fb<sup>-1</sup>)

- **Simultaneous fit to W boson q/p<sub>T</sub> and Z<sup>0</sup> boson ϕ\***

$$\phi^* = \frac{\tan((\pi - \Delta\phi)/2)}{\cosh(\Delta\eta/2)} \sim \frac{p_T^Z}{M}$$

- 28 < p<sub>T</sub>(μ) < 52 GeV is the optimal range for the fit: **2.4M W candidates**



# Modeling

$$\frac{d\sigma}{dp_1 dp_2} = \left[ \frac{d\sigma(m)}{dm} \right] \left[ \frac{d\sigma(y)}{dy} \right] \left[ \frac{d\sigma(p_T, y)}{dp_T dy} \left( \frac{d\sigma(y)}{dy} \right)^{-1} \right] \left[ (1 + \cos^2 \theta) + \sum_{i=0}^7 A_i(p_T, y) P_i(\cos \theta, \phi) \right]$$

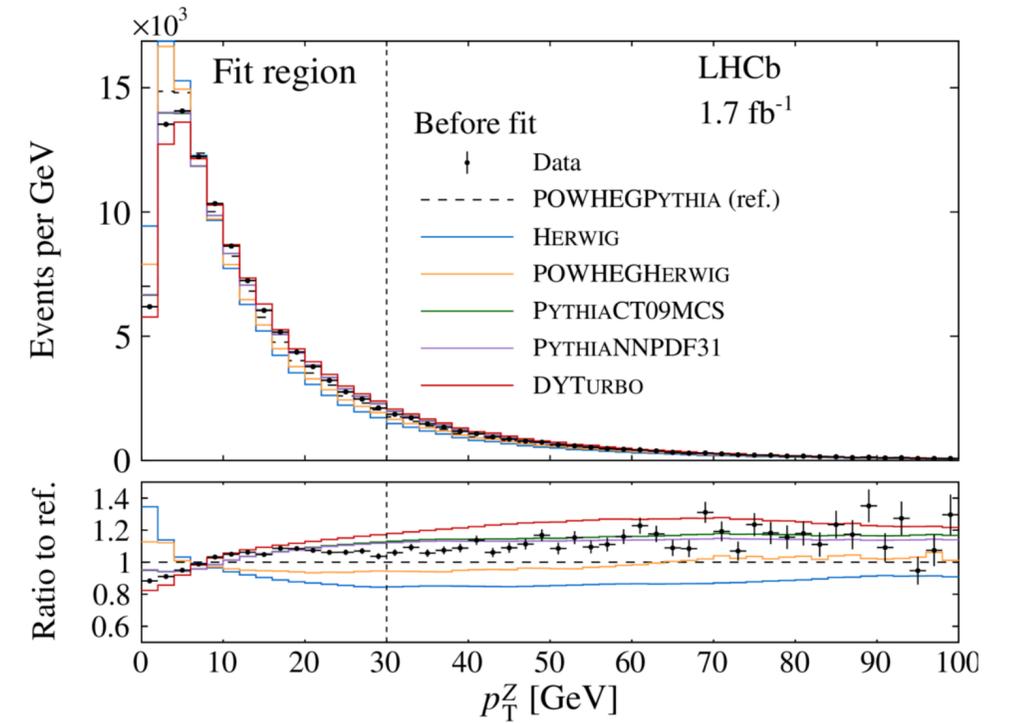
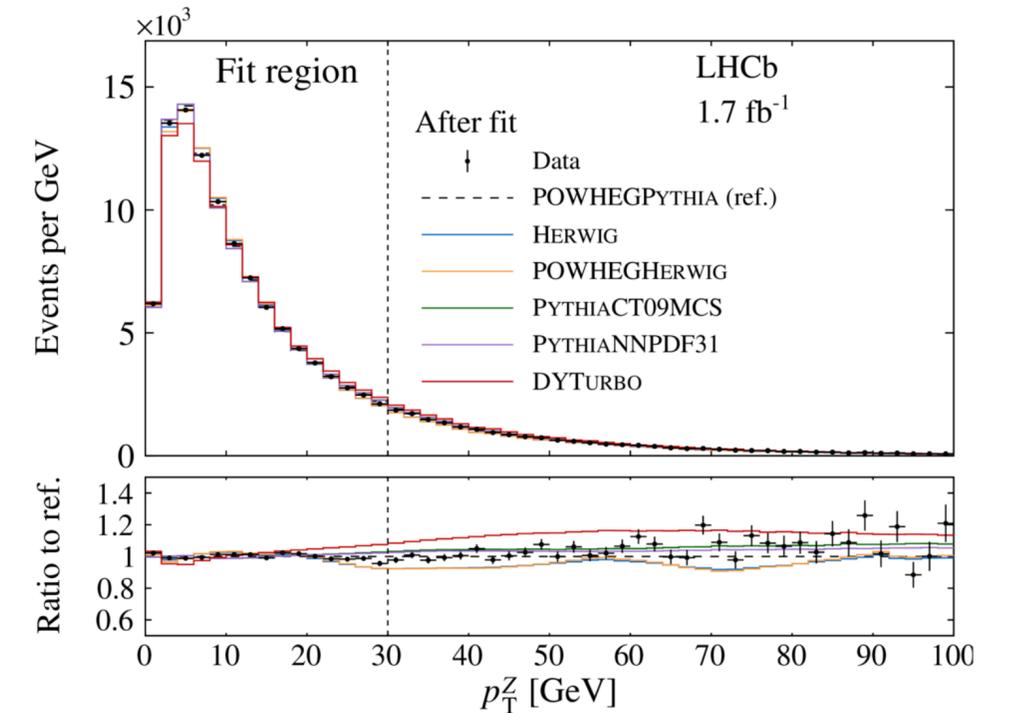
**Inv. mass:**  
**Breit Wigner**

**Rapidity: fixed**  
**order pQCD**

**p<sub>T</sub>: parton**  
**shower**

**Angular**  
**coefficients: fixed**  
**order pQCD**

- As for ATLAS, Powheg+Pythia is used as baseline simulation
- Pythia QCD parameters are fitted to match p<sub>T</sub>(Z<sup>0</sup> → μ<sup>+</sup>μ<sup>-</sup>) distribution
- Templates reweighted also to match DYTurbo
- Pythia, Photos, Herwig for QED description
- Three different PDFs sets: NNPDF3.1, CT18, MSHT20

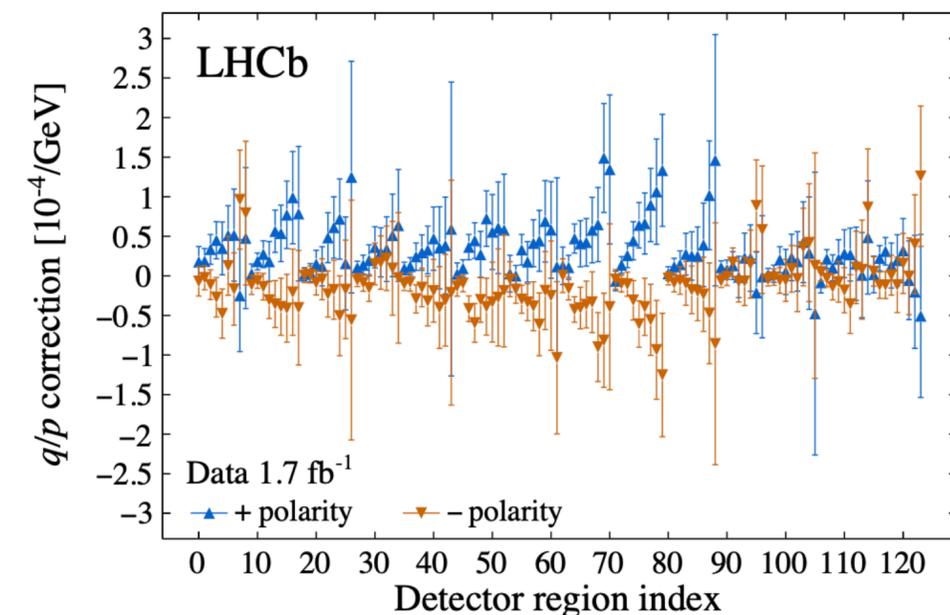
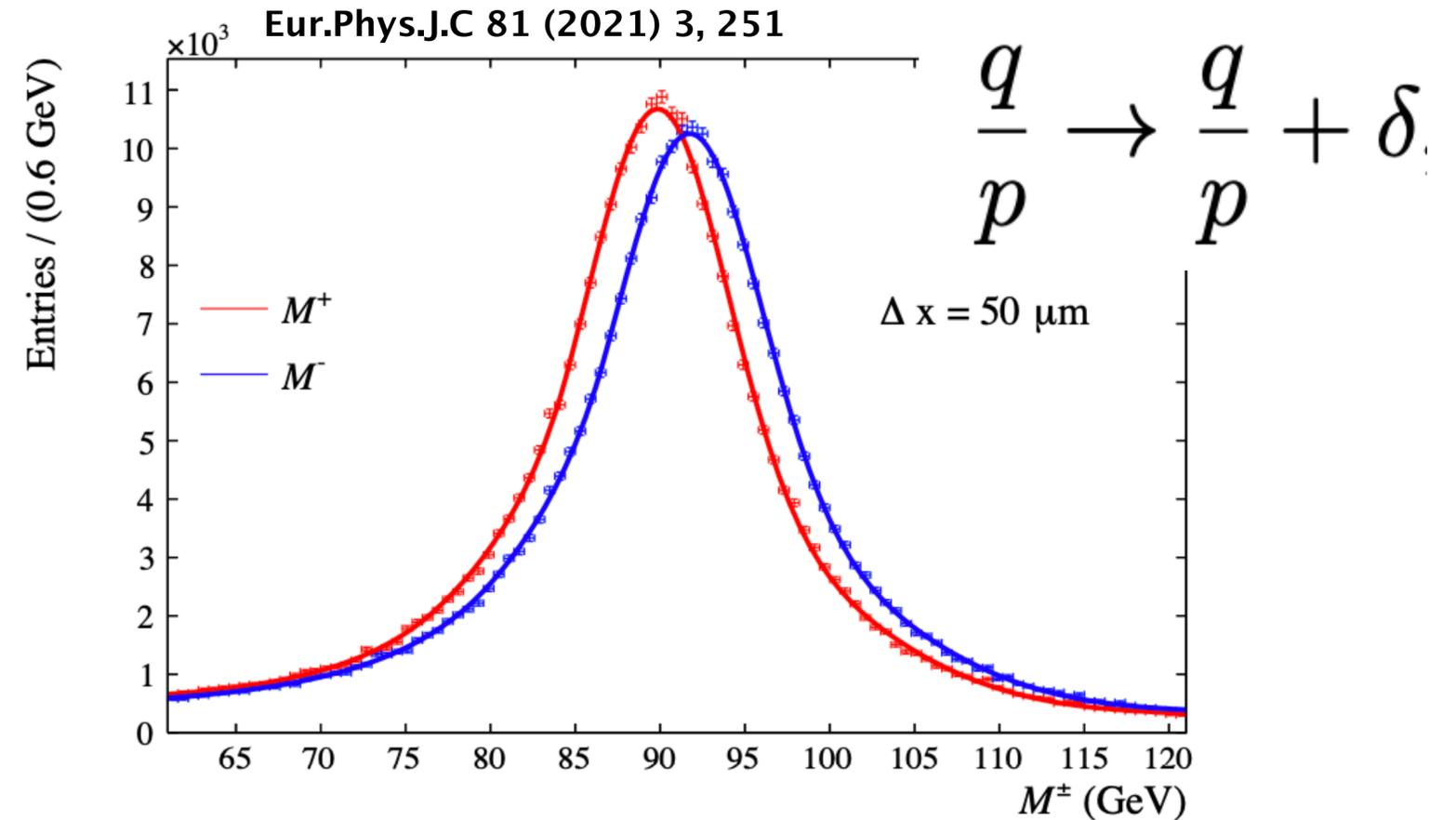


# Muon curvature biases

- ATLAS determined the curvature bias ( $\delta$ ) in E/p calibration for electrons: usable only if muon and electron reconstruction has a comparable performance
- Due to saturation effects in ECAL, at LHCb electrons are not usable for this purpose
- **Pseudo-mass method applied to  $Z^0 \rightarrow \mu^+\mu^-$ :** does not depend from the magnitude of the momentum

$$\mathcal{M}^\pm = \sqrt{2p^\pm p_T^\pm \frac{p^\mp}{p_T^\mp} (1 - \cos \theta)}$$

$$\delta \approx A \frac{\langle \frac{1}{p^+} \rangle + \langle \frac{1}{p^-} \rangle}{2}$$

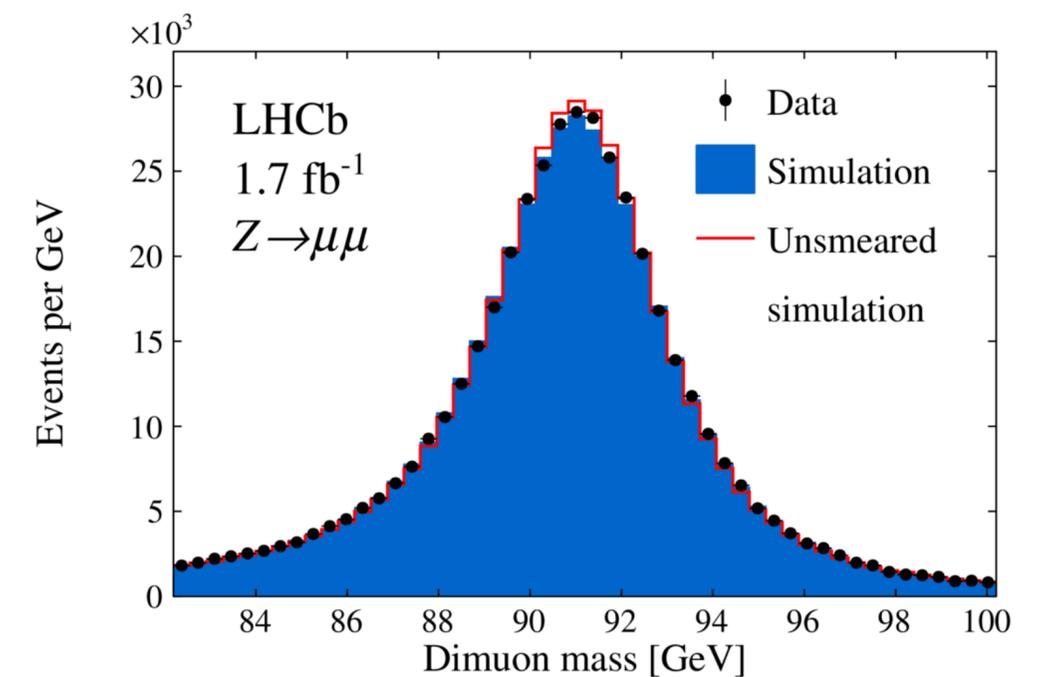
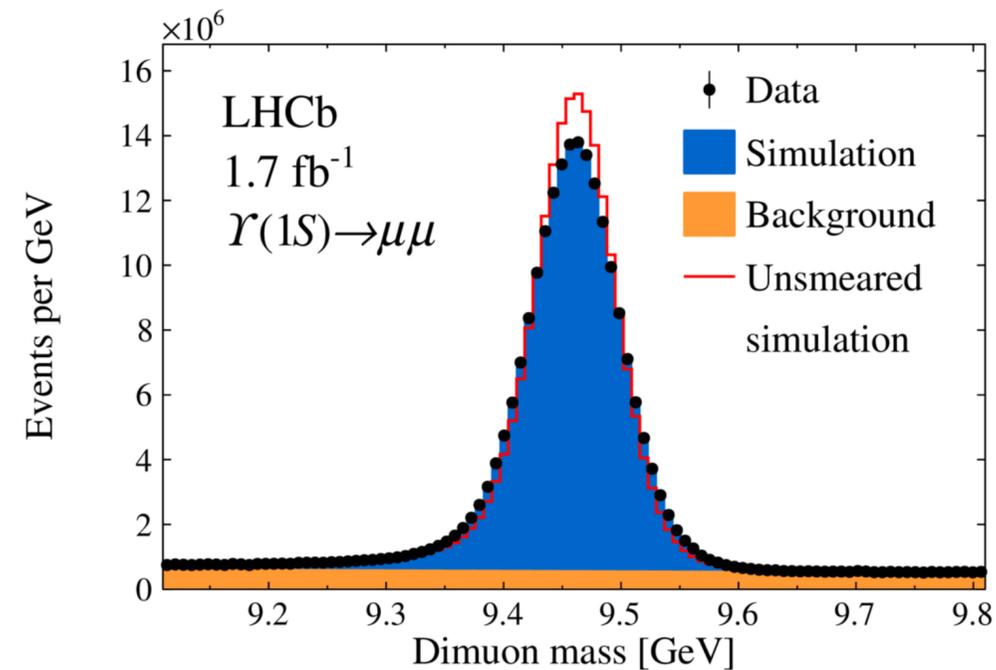
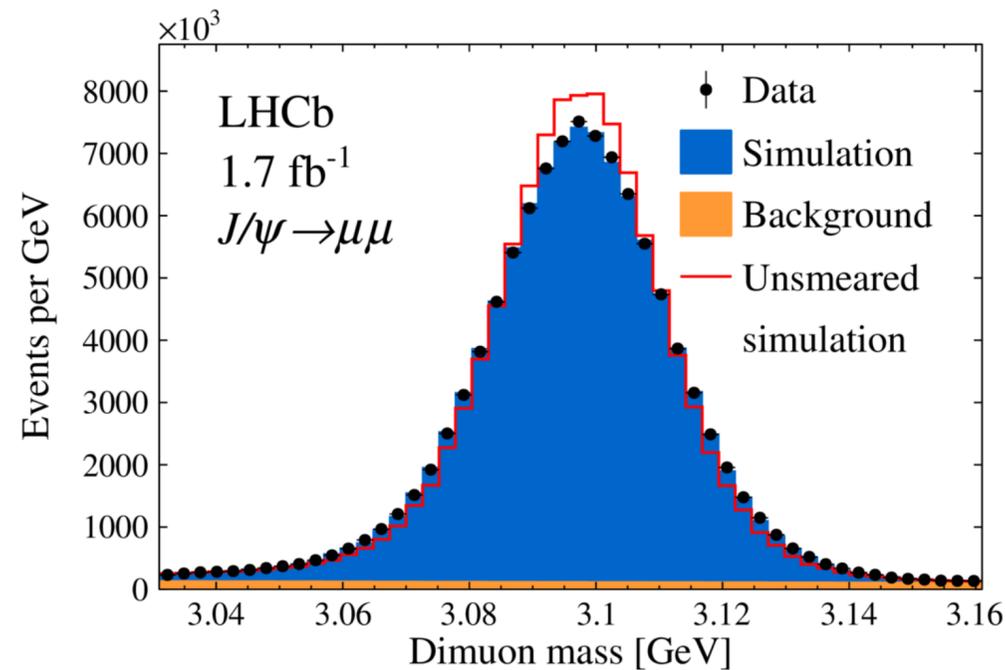


# Muon calibration

- Muon momentum scale and resolution obtained with several dimuon resonances samples
- Reconstruction efficiencies with tag & probe  $Z^0 \rightarrow \mu^+\mu^-$

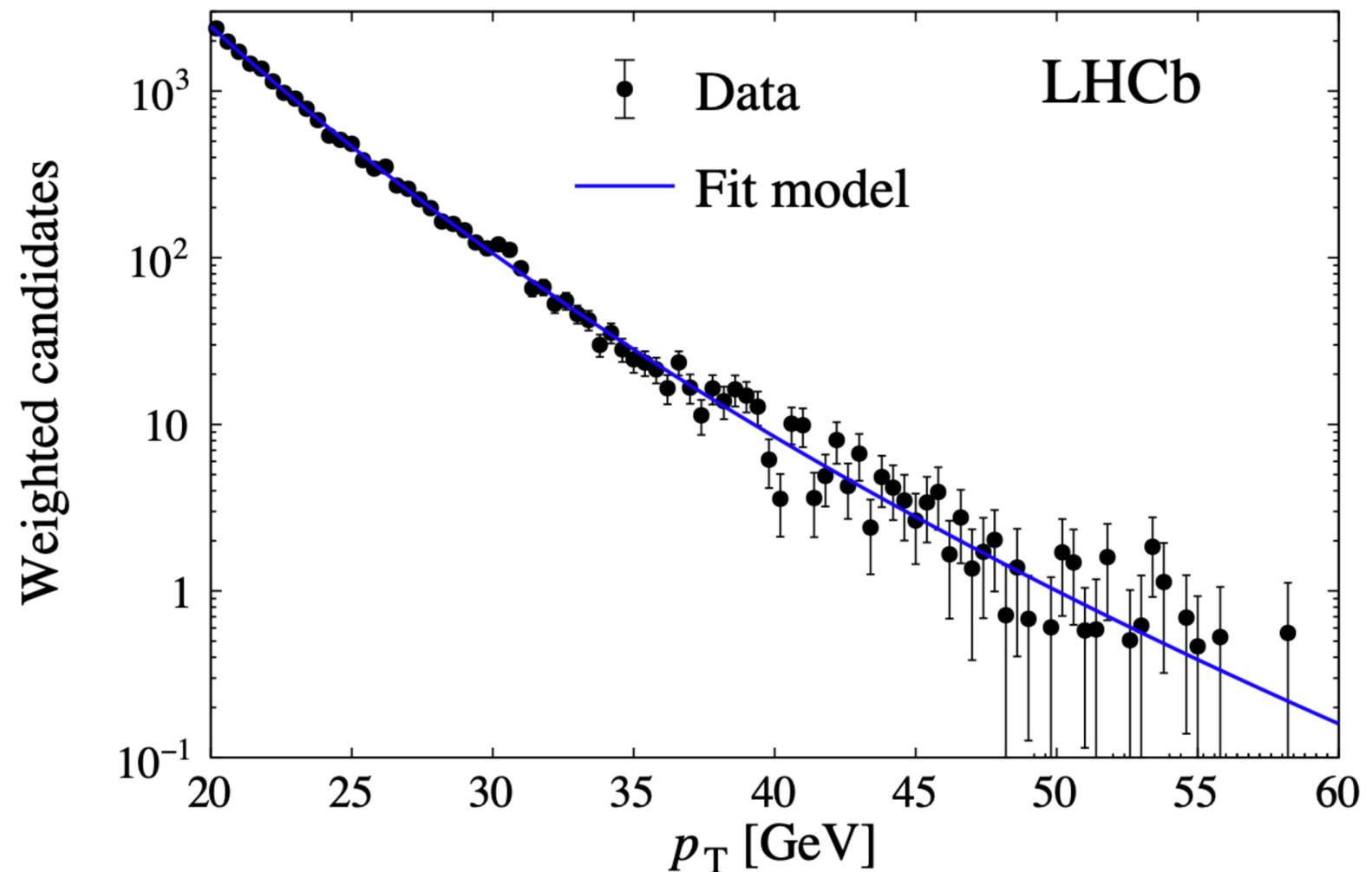
$$\frac{q}{p} \rightarrow \frac{q}{p \cdot \mathcal{N}(1 + \alpha \sigma_{MS})} + \mathcal{N}\left(\delta, \frac{\sigma_\delta}{\cosh \eta}\right)$$

scale    smearing    bias



# Background

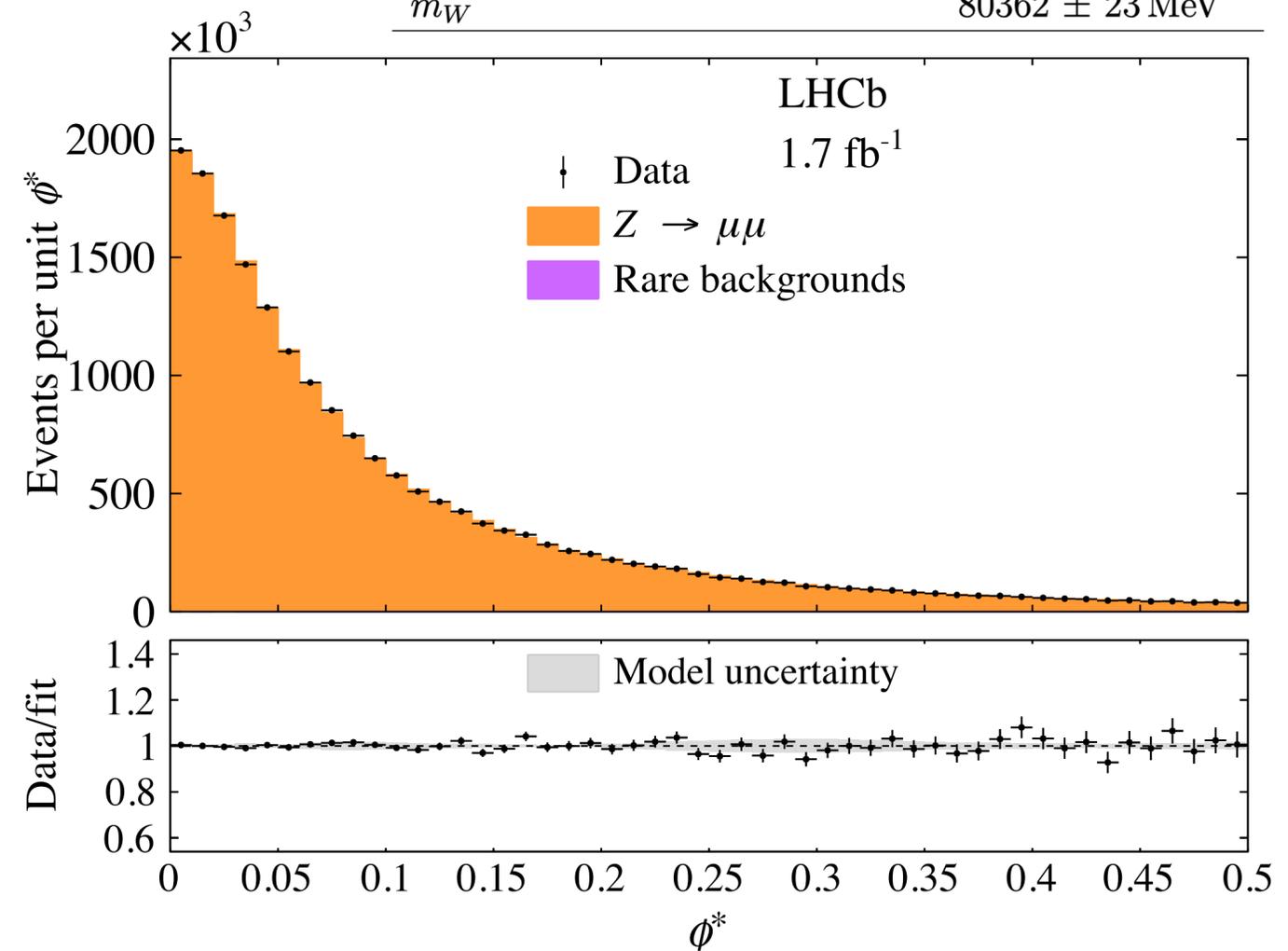
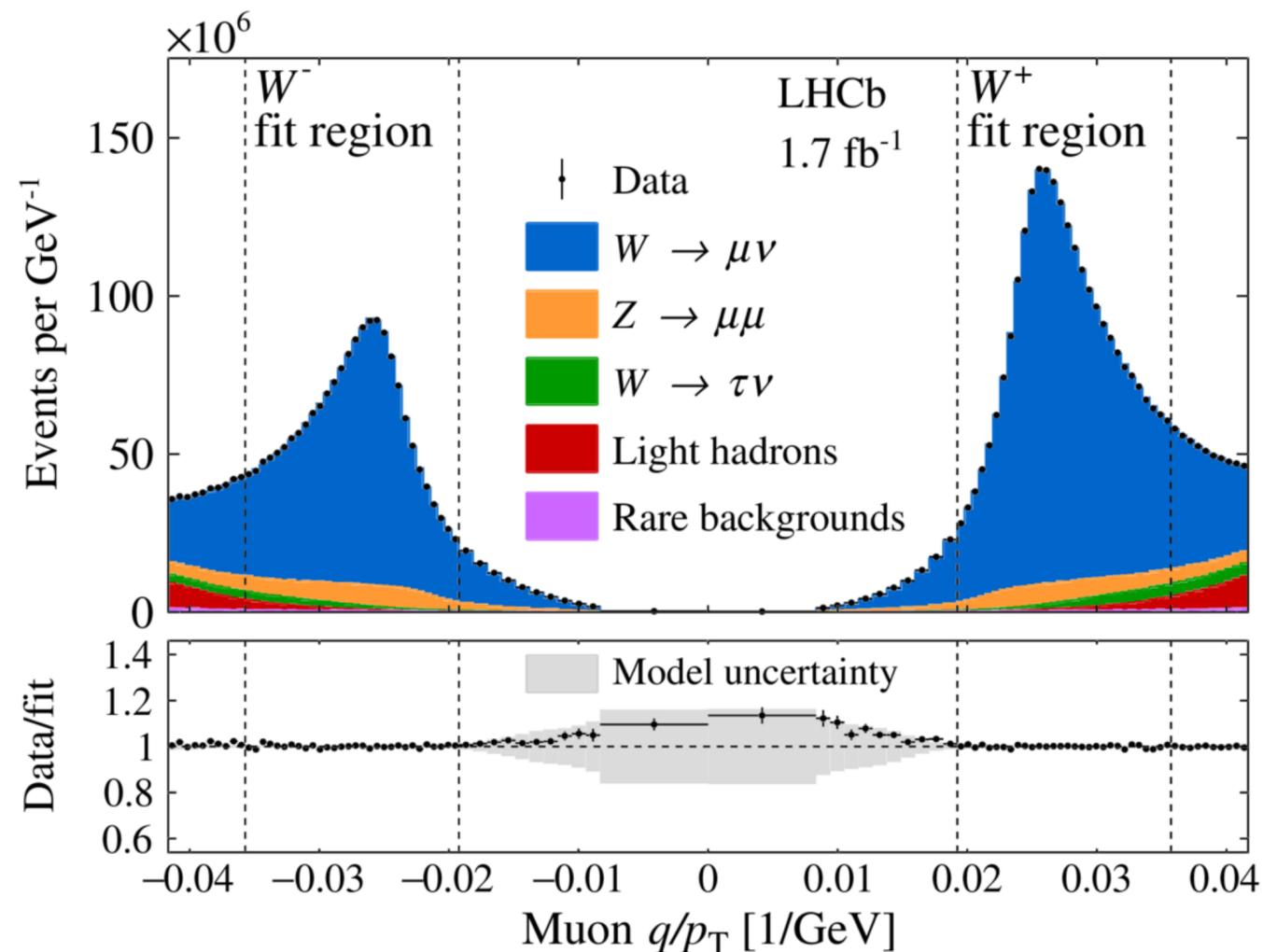
- Most of backgrounds are modeled with simulated samples: single-top, quark/anti-quark (t, b, c), Z/W decays, Drell-Yan
- QCD background (decays-in-flight) has been obtained with a data-driven technique, by inverting the muon identification cuts (i.e. impact parameter/isolation)
- This model (Hagedorn distribution) accurately describes the region of the Jacobian peak



# W mass fit result

- Several QCD parameters and  $A_3$  scaling are also extracted from the  $m_W$  template fit

Parameter	Value
Fraction of $W^+ \rightarrow \mu^+ \nu$	$0.5288 \pm 0.0006$
Fraction of $W^- \rightarrow \mu^- \nu$	$0.3508 \pm 0.0005$
Fraction of hadron background	$0.0146 \pm 0.0007$
$\alpha_s^Z$	$0.1243 \pm 0.0004$
$\alpha_s^W$	$0.1263 \pm 0.0003$
$k_T^{\text{intr}}$	$1.57 \pm 0.14 \text{ GeV}$
$A_3$ scaling	$0.975 \pm 0.026$
$m_W$	$80362 \pm 23 \text{ MeV}$



# Systematics and cross-checks

Source	Size [ MeV]
Parton distribution functions	9
Theory (excl. PDFs) total	17
Transverse momentum model	11
Angular coefficients	10
QED FSR model	7
Additional electroweak corrections	5
Experimental total	10
Momentum scale and resolution modelling	7
Muon ID, trigger and tracking efficiency	6
Isolation efficiency	4
QCD background	2
Statistical	23
Total	32

**Statistical uncertainty still large: with the full Run 2 dataset a total uncertainty < 20 MeV is already possible**

## Cross checks:

- W-like measurement of  $Z^0$  boson mass
- Consistency of orthogonal subsets: muon charge, magnet polarities,  $\phi$ ,  $\eta$
- Fit  $p_T$  range
- Fit model freedom
- NNLO vs NLO PDFs

$$m_W = 80362 \pm 23_{\text{stat}} \pm 10_{\text{exp}} \pm 17_{\text{theory}} \pm 9_{\text{PDF}} \text{ MeV, } \mathbf{NNPDF3.1}$$

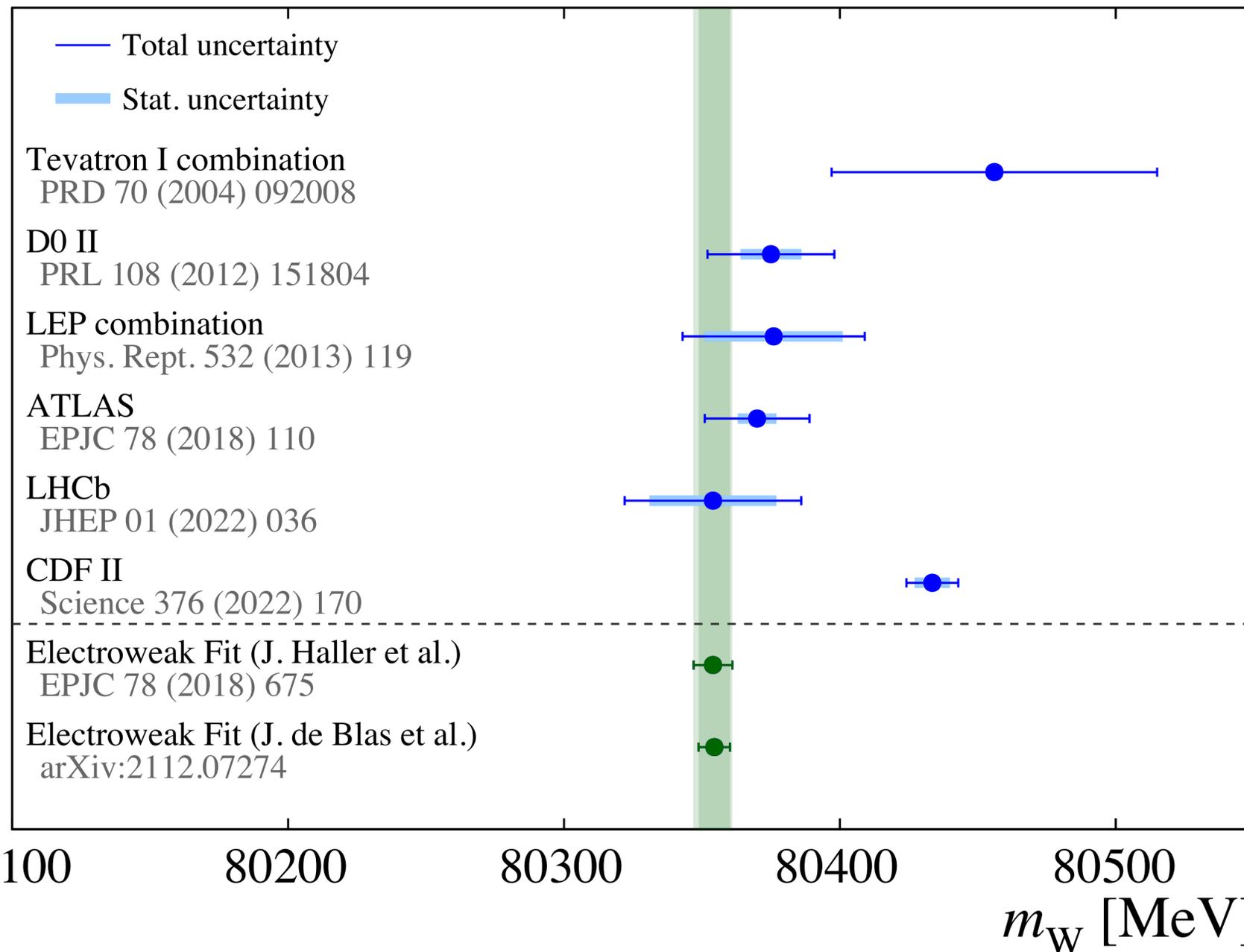
$$m_W = 80350 \pm 23_{\text{stat}} \pm 10_{\text{exp}} \pm 17_{\text{theory}} \pm 12_{\text{PDF}} \text{ MeV, } \mathbf{CT18}$$

$$m_W = 80351 \pm 23_{\text{stat}} \pm 10_{\text{exp}} \pm 17_{\text{theory}} \pm 7_{\text{PDF}} \text{ MeV, } \mathbf{MSHT20}$$

## Final result:

$$m_W = 80354 \pm 23_{\text{stat}} \pm 10_{\text{exp}} \pm 17_{\text{theory}} \pm 9_{\text{PDF}} \text{ MeV}$$

# Comparison with CDF



- Significant displacement between new CDF II measurement and other most precise measurements
- LHC measurements are closer to the Electroweak Fit prediction with respect to CDF II
- However precision of CDF II measurement is much better

# Comparison with CDF

## Uncertainties (in MeV)

	CDF	ATLAS	LHCb
<b>Statistical</b>	6.4	6.8	23
Lepton energy/ momentum scale	2 ( $\mu$ ) + 6 (e)	7* ( $\mu$ ) + 7* (e)	7 ( $\mu$ )
<b>PDFs</b>	4	7*	9
Model (excl. PDFs)	3.5	8*	17
<b>Total</b>	9.4	18.5	31.4

## Modeling

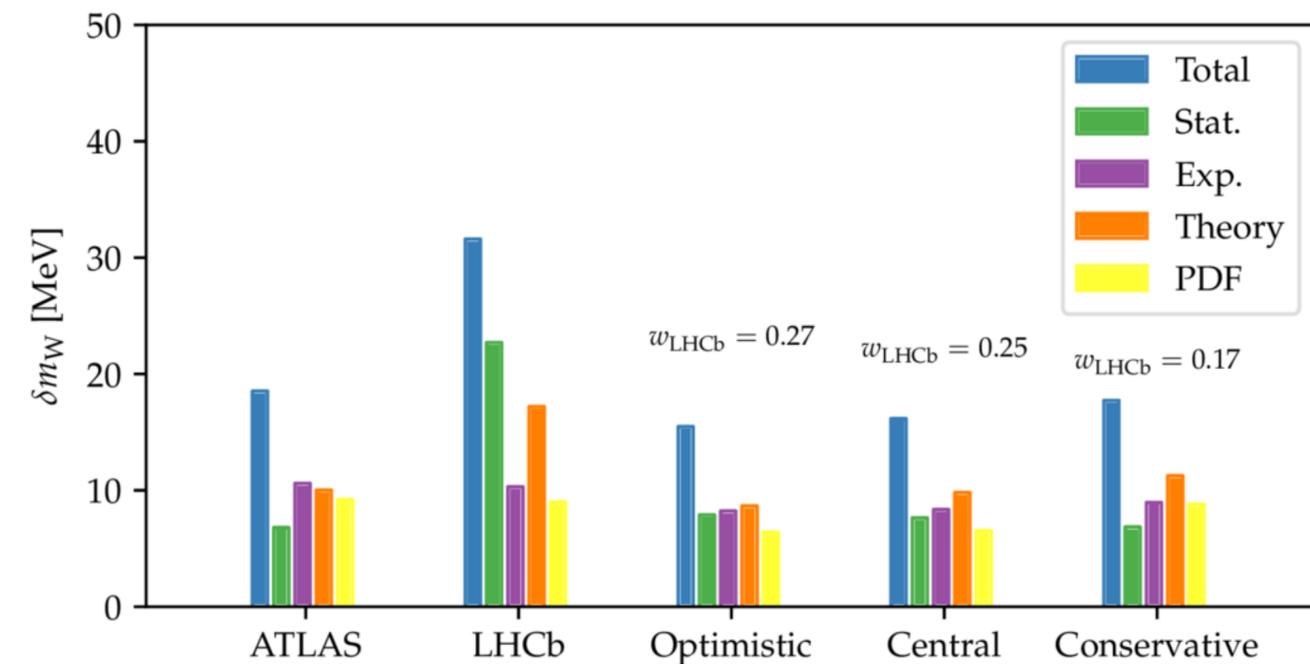
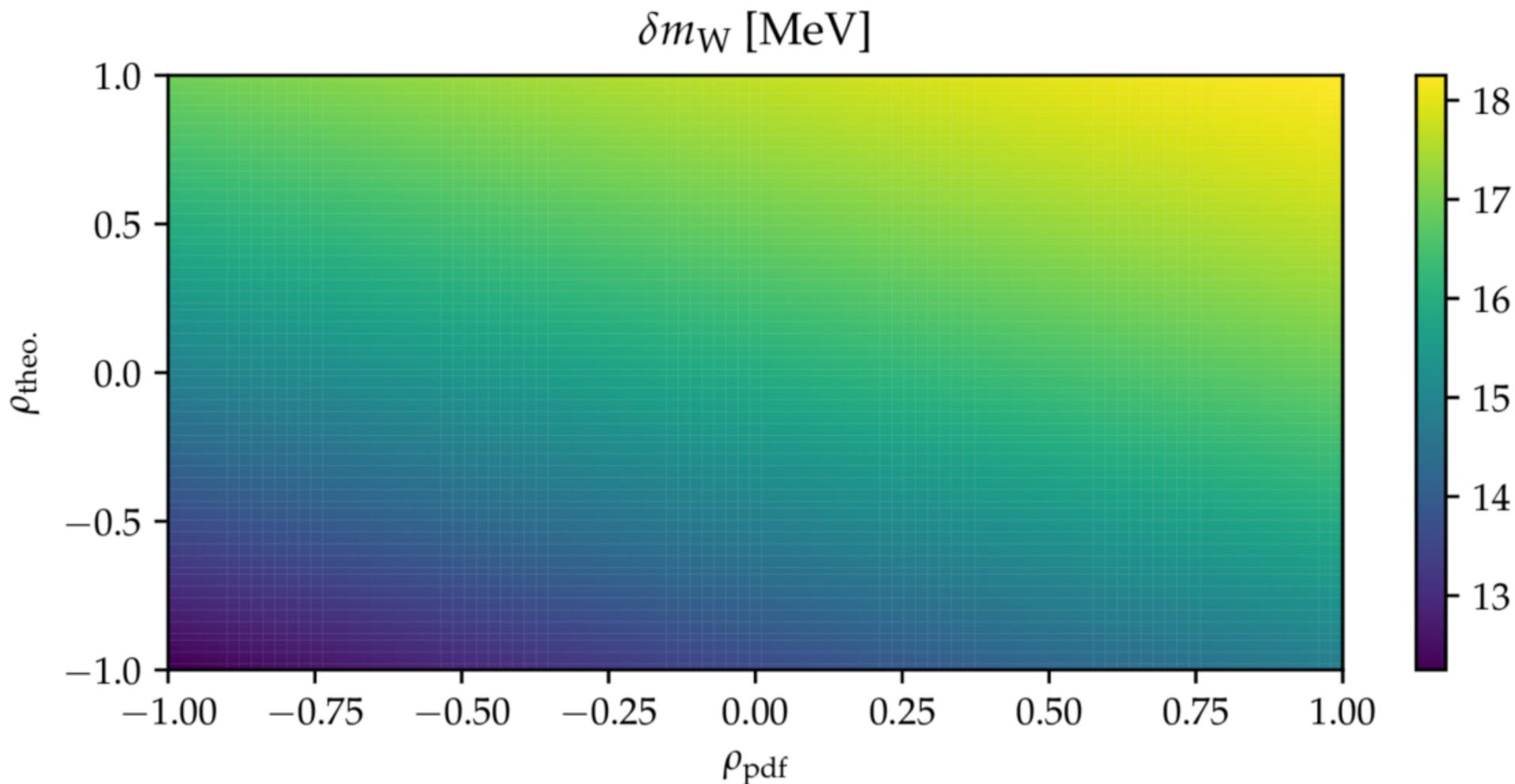
	CDF	ATLAS	LHCb
<b>Baseline</b>	RESBOS	Powheg+Pythia	Powheg+Pythia
<b>Reweight</b>	-	DYNNLO	DYTURBO
<b>Parton shower</b>	data-driven	data-driven	data-driven
<b>QED</b>	PHOTOS+HORACE	PHOTOS	Pythia+PHOTOS+Herwig

\*given separately for  $p_T$  and  $m_T$  fits, combined assuming 50% correlation

**Notice: CDF measurement took profit of the PDFs determination at LHC**

# LHC combination

- LHC measurements combination is not trivial, it depends on several correlations
- A naive expectation on **ATLAS+LHCb combination** is given



**PDFs uncertainty correlation is expected to be negative**

# Future prospects at LHC

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Scenario	Experiments	$\delta m_W$ (MeV)		
		Tot	Exp	PDF
Default	2×GPD + LHCb	9.0	4.7	7.7
Default	1×GPD + LHCb	10.1	6.5	7.7
Default	2×GPD	12.0	5.8	10.5
PDF4LHC(3-sets)	2×GPD + LHCb	13.6	4.8	12.7
PDF4LHC(3-sets)	1×GPD + LHCb	14.6	7.3	12.7
PDF4LHC(3-sets)	2×GPD	17.7	5.5	16.9
$\delta_{\text{exp}}^{\text{LHCb}} = 0$	2×GPD + LHCb	8.7	4.0	7.7
$\delta_{\text{exp}}^{\text{LHCb}} = 0$	1×GPD + LHCb	9.8	5.9	7.9
$\delta_{\text{exp}}^{\text{LHCb}} = 0$	2×GPD	12.0	5.8	10.5
$\delta_{\text{exp}}^{\text{GPD}} = 0$	2×GPD + LHCb	7.9	1.9	7.7
$\delta_{\text{exp}}^{\text{GPD}} = 0$	1×GPD + LHCb	7.9	1.9	7.7
$\delta_{\text{exp}}^{\text{GPD}} = 0$	2×GPD	10.5	0.1	10.5
$\delta_{\text{PDF}} = 0$	2×GPD + LHCb	4.6	4.6	0.0
$\delta_{\text{PDF}} = 0$	1×GPD + LHCb	5.8	5.8	0.0
$\delta_{\text{PDF}} = 0$	2×GPD	5.5	5.5	0.0

**GPD = General Purpose Detector = ATLAS/CMS**

- Not a precise extrapolation, just a way to visualize the contribution of the three experiments to the  $m_W$  combination
- Only the PDF uncertainty is considered for the model
- Statistical uncertainty not included

# Conclusions

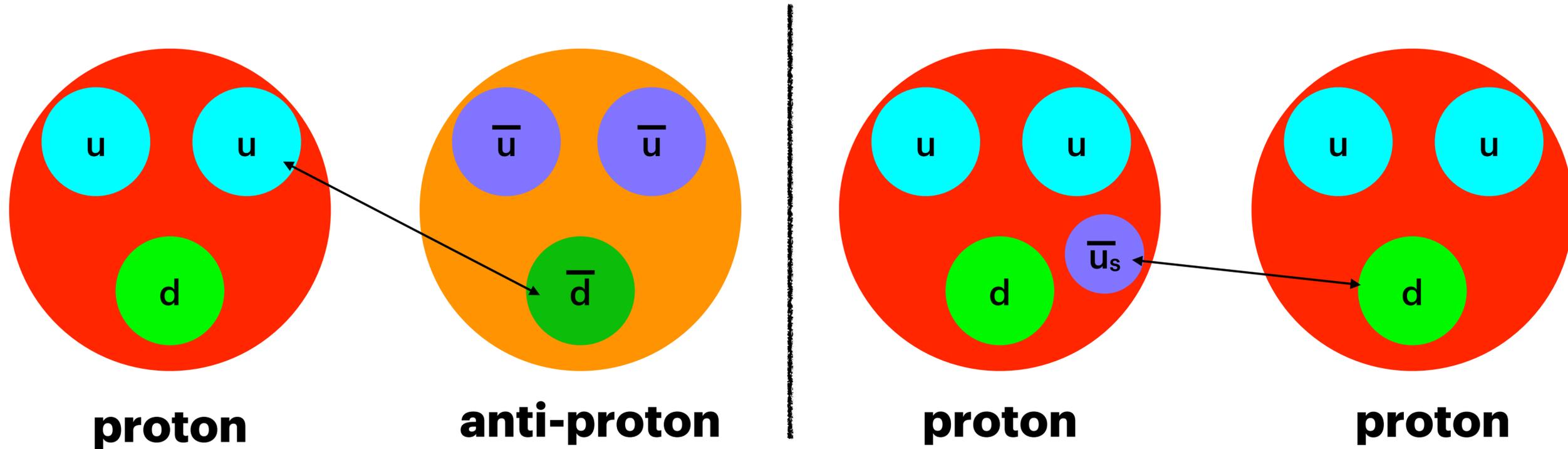
- LHCb have already performed a W mass boson measurement, just with a part of Run 2 data
- A measurement with a precision  $<20$  MeV is possible with the available dataset, the work is currently on-going
- LHCb Upgrade-II for the HL-LHC (<https://arxiv.org/abs/1808.08865>) will allow 100x higher statistics (x50 from luminosity and x2 from having an ECAL with sufficient dynamic range for  $Z \rightarrow ee$ )
- There are several ideas to improve the modeling systematic uncertainty
- **The combination of the measurements from the three experiments is fundamental to obtain the final precision at LHC**
- **We have many years before the next lepton collider, LHC could be the the only way to confirm CDF result in the short period**

**Thanks for your attention!**



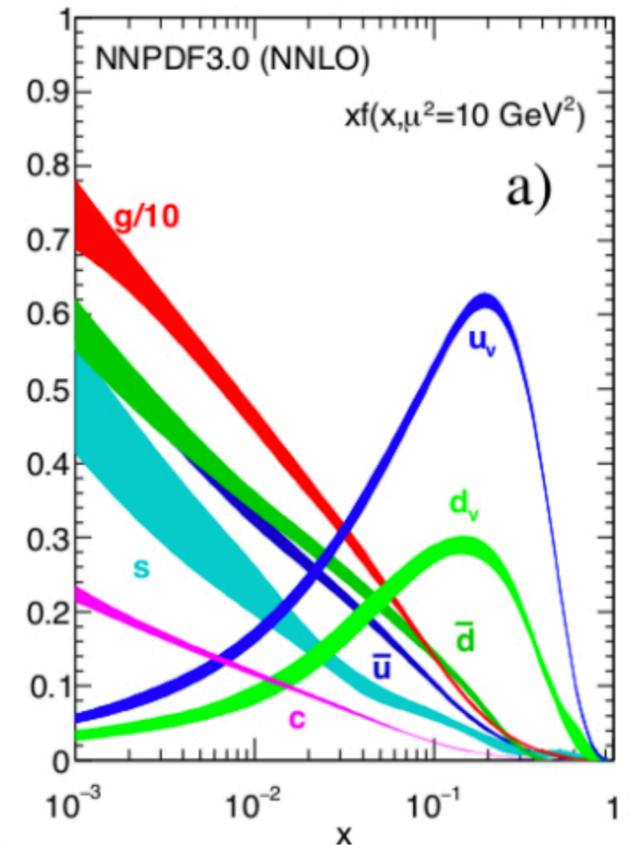
**Backup**

# Proton-anti-proton vs proton-proton



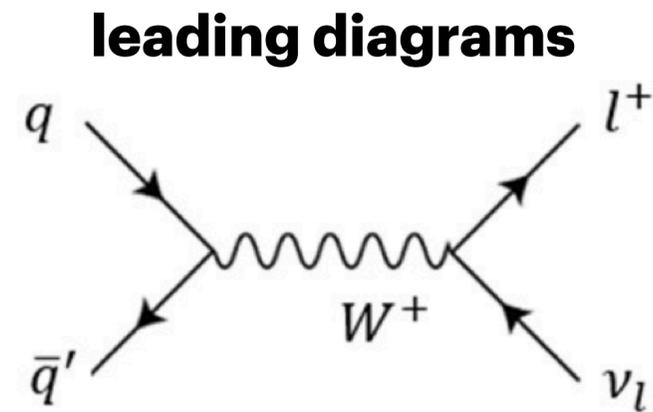
- At Tevatron W boson is mainly produced via valence quark interactions
- At LHC mainly through valence-sea quarks interaction

**Knowledge of Parton Distribution Functions is fundamental for modeling the W boson production**

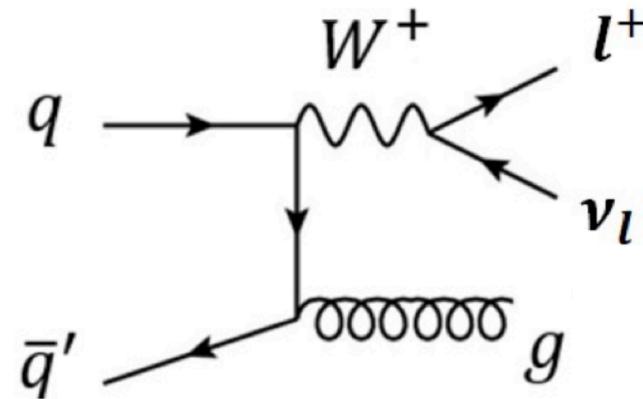


$x$  = fraction of proton momentum taken by the parton

# W boson cross section



$\mathbf{P_T(W)=0}$



$\mathbf{P_T(W)>0}$

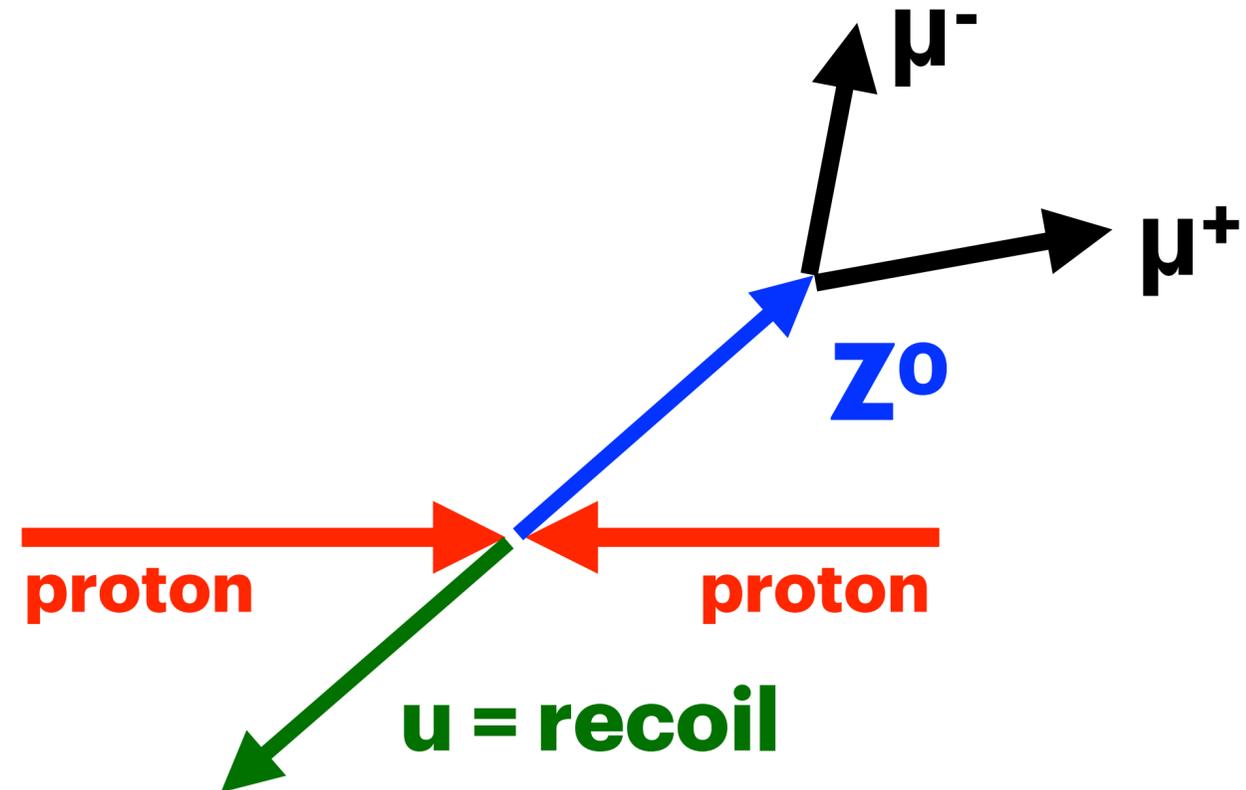
$A_i$  = angular coefficients: ratio between helicity dependent and unpolarized cross-sections

**Angular-integrated cross-section**

$$\frac{d\sigma}{dp_T^2 dy dm d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma}{dp_T^2 dy dm} \times \left[ (1 + \cos^2\theta) + A_0 \frac{1}{2}(1 - 3\cos^2\theta) \right. \\ \left. + A_1 \sin 2\theta \cos\phi + A_2 \frac{1}{2} \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta \right. \\ \left. + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right].$$

# Experimental techniques

- Large sample of  $Z^0 \rightarrow \mu^+\mu^-$  for tuning and validation
- $Z^0$  fully reconstructed
- energy scale and resolution can be determined by comparing  $Z^0$  data and simulation
- Tag & Probe technique to measure lepton efficiencies in data



- Templates depend from the W boson production model, should be corrected for data/simulation differences

$$\frac{d\sigma}{dp_1 dp_2} = \left[ \frac{d\sigma(m)}{dm} \right] \left[ \frac{d\sigma(y)}{dy} \right] \left[ \frac{d\sigma(p_T, y)}{dp_T dy} \left( \frac{d\sigma(y)}{dy} \right)^{-1} \right] \left[ (1 + \cos^2 \theta) + \sum_{i=0}^7 A_i(p_T, y) P_i(\cos \theta, \phi) \right]$$

**Inv. mass: Breit Wigner** (red arrow pointing to  $\frac{d\sigma(m)}{dm}$ )  
**Rapidity: fixed order pQCD** (green arrow pointing to  $\frac{d\sigma(y)}{dy}$ )  
 **$p_T$ : parton shower** (blue arrow pointing to  $\frac{d\sigma(p_T, y)}{dp_T dy}$ )  
**Angular coefficients: fixed order pQCD** (purple arrow pointing to the bracketed term)

# Uncertainties

## Experimental uncertainties

- Muon momentum calibration and scale
- Background processes
- Differences between data and simulation for lepton efficiencies

## Theoretical uncertainties

- Parton Distribution Functions
- Modeling of  $p_T(W)$
- Modeling of angular coefficients  $A_i$
- Modeling of QED radiation

# LHCb performance

Int. J. Mod. Phys. A 30, 1530022 (2015)

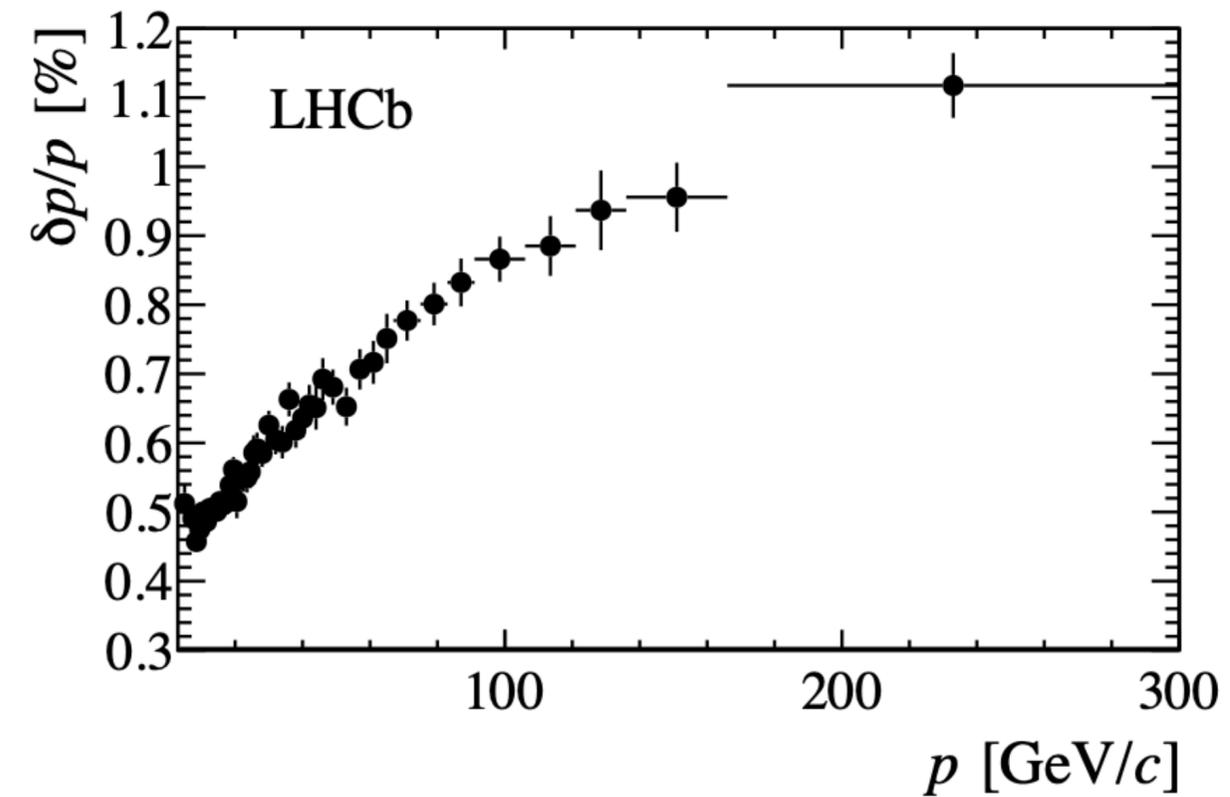
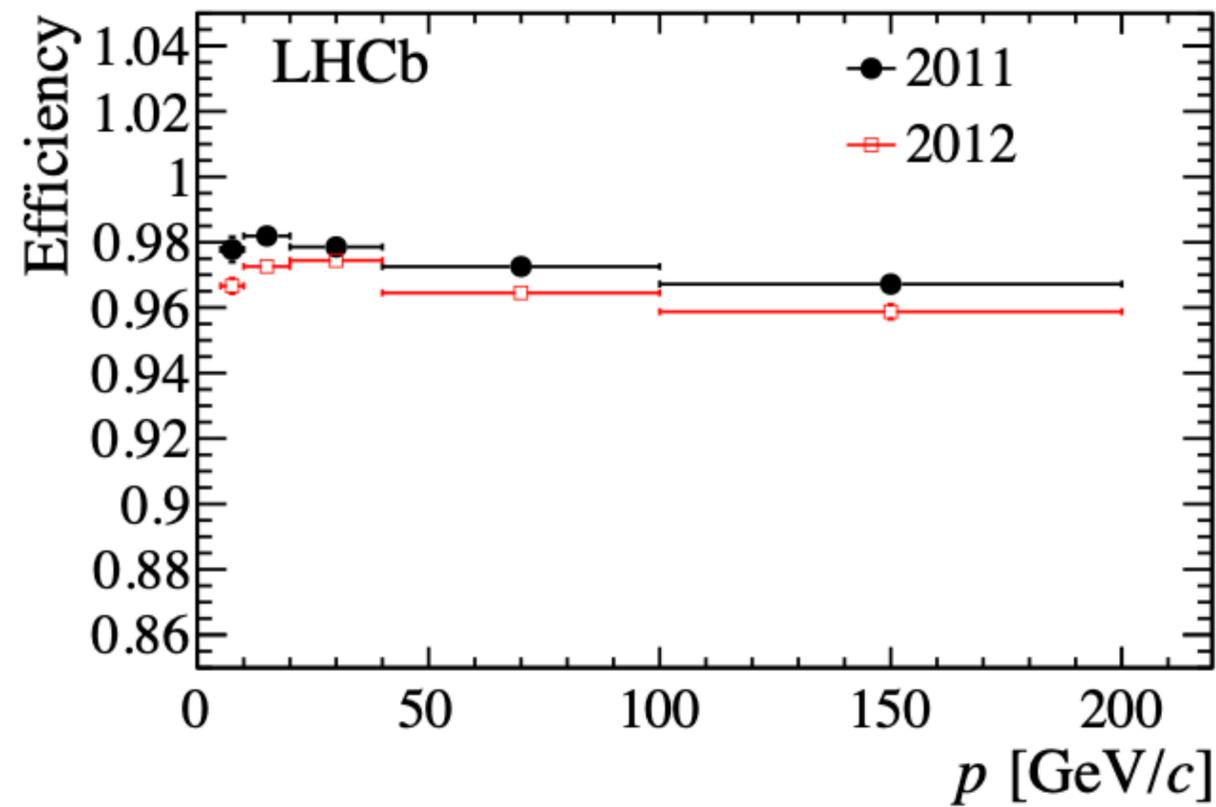
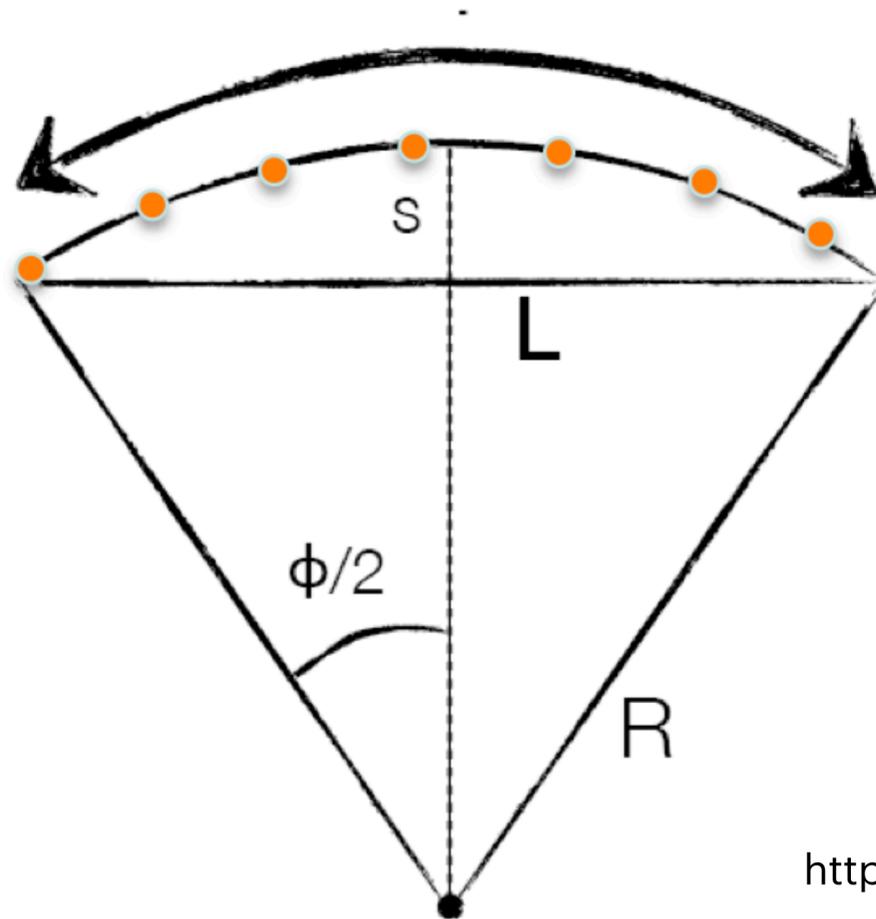


Figure 17: Relative momentum resolution versus momentum for long tracks in data obtained using  $J/\psi$  decays.

# Tracking



$$p \left[ \frac{GeV}{c} \right] = 0.3 B[T] R[m]$$

$$R = \frac{L^2}{8s}$$

$$\frac{q}{p} \rightarrow \frac{q}{p \cdot \mathcal{N}(1 + \alpha \sigma_{MS})} + \mathcal{N} \left( \delta, \frac{\sigma_\delta}{\cosh \eta} \right)$$

scale    smearing    bias

[https://www.desy.de/~garutti/LECTURES/ParticleDetectorSS12/L9\\_Tracking.pdf](https://www.desy.de/~garutti/LECTURES/ParticleDetectorSS12/L9_Tracking.pdf)