



**Paolo Franzini and
Juliet Lee-Franzini**

**MEMORIAL
SYMPOSIUM**

19 Aprile 2022

dalle 10:00 alle 17:00

Auditorium Bruno Touschek

19 April 2022

INFN - Laboratori Nazionali di Frascati

The Muon Magnetic Anomaly

Graziano Venanzoni – INFN Pisa

Paolo and Juliet

- As many other colleagues I had the privilege to be in Frascati during the KLOE years, first as a student (in the 90s) and after as a researcher (in the 2000s).
- I was deeply influenced by Paolo and Juliet: their passion for physics, enthusiasm, informal approach, scientific clearness and rigor were always an example and a guide.
- As a co-spoke in g-2 I often think to them when I'm in front of a difficult decision: "What Paolo and Juliet would have done in this case?"
- I miss Paolo and Juliet very much and I'm sure I would have been different without having meet them!



Paolo e Juliet “legacy”

(from Gino’s contribution to DIF06 proceedings)



Paolo, Juliet and g-2

- Paolo and Juliet were particularly interesting to g-2 and sigma(had).
- During the years I had almost continuous interactions with them. Often in my office (sorry Erika!) but sometimes in theirs.
- Their contribution/criticism/advice was always very constructive and helpful.
- Juliet was a referee of all the ISR memo's and Paolo strongly edited the physics (PLB) papers.
- When in 2009 I decided to embark in the g-2 adventure they gave me a strong encourage and support.
- In the last years of my activity at LNF (~2015/16) I met often Paolo and discuss with him the g-2 physics and experiment.

Paolo and Juliet and g-2



$$\sigma(e^+e^- \rightarrow \text{hadrons} + \gamma)$$

Juliet Lee-Franzini

Laboratori Nazionali di Frascati

Cape Cod, June 2003

$$g - 2$$

Paolo Franzini

Rome and Karlsruhe

Karlsruhe - Fall 2001



Measurement of $\sigma(e^+e^- \rightarrow \pi^+\pi^-)$ at e^+e^- Colliders

Juliet Lee-Franzini
LNF, INFN, I-00044, Frascati, Italy

Abstract

At the DAΦNE-II workshop a session was devoted to the prospects of measuring the hadronic cross section at the new DAΦNE. The session included six papers, two theoretical and four experimental ones. The theory treatises, one on the muon anomaly and the other on measuring the hadronic cross section using initial state radiation at e^+e^- colliders, set the background for the four experimental discussions. I summarize in the following the salient points of the session.

INTRODUCTION

The electroweak contribution is evaluated to the percent level: $(15.4 \pm 0.2) \times 10^{-10}$. Coming to more uncertain terms, he showed that the sign of the light by light scattering contribution is unambiguously positive but assigned to its value a 50% error: $\delta_{\text{had}}^{\text{bl}} a_\mu = (8 \pm 4) \times 10^{-10}$. He devoted most of his time to what still remains the most uncertain term, the hadronic vacuum polarization effects. He advocates using the analysis by HMNT which gives, including the corrected CMD2 results, $\delta_{\text{had}}^{\text{VP}} a_\mu = [(692.4 \pm 6.4)_{\text{lo}} - (9.79 \pm 0.095)_{\text{hlo}}] \times 10^{-10}$. Adding everything: $a_\mu = (11659176.3 \pm 7.4) \times 10^{-10}$ [2]. In conclusion the discrepancy between the experimental measurement and the SM evaluation is now $(32 \pm 10) \times 10^{-10}$, about 3σ .

?????

file=...g-2.tex

KLOE Memo

Paolo Franzini
13 November 2010

Magnetic moment, magneton, gyromagnetic or Landé g-factor

From the Dirac equation the electron magnetic moment $\mu_{e,D}$ is:

$$\mu_{e,D} = g \frac{e\hbar}{2m_e c} \frac{J}{\hbar}$$

with the gyromagnetic, or Landé g -factor, $g=2$, to lowest order in QED, using Gaussian units. g is dimension-less. J dimensionally is mvr and is measured in $\text{gram} \times \text{centimeter}^2 \times \text{second}^{-1}$. J/\hbar is the spin in the usual \hbar units. Introducing the Bohr magneton, making explicit the fact that it applies to the electron,

$$\mu_{B,e} = \frac{e\hbar}{2m_e c}$$

Please avoid saying that the anomaly, a dimensionless quantity, is a magnetic moment (anomalous or not) which is a dimensionfull quantity! That is a gross conceptual mistake.

Paolo and Juliet and the born of ISR...

- The g-2 physics through the measurement of sigma(had) was part of the KLOE program
- However DAFNE was not optimized for an energy scan...



KLOE Workshop at Karlsruhe 1996

Tesi di Laurea

R Measurements at KLOE

a_μ and $\alpha_{\text{QED}}(M_Z)$

PAOLO FRANZINI



Studio di fattibilità di una misura di sezione
d'urto adronica con il rivelatore KLOE,
finalizzata alla determinazione delle correzioni
adroniche al momento magnetico anomalo
del muone

Anno Accademico 1995-1996

Candidato:

Giorgio Cabibbo

Relatore interno:

Prof. Luciano Paoluzzi

Relatore esterno:

Prof. Paolo Franzini

Hadronic Contributions to the Muon g-2

Riccardo Barbieri

Dip. di Fisica, Univ. di Pisa and INFN, Sez. di Pisa
and

Ettore Remiddi

Dip. di Fisica, Univ. di Bologna and INFN, Sez. di Bologna

1991

The obvious target in the domain of the precision measurements of the muon g-2 is the detection of the electroweak corrections due to vector and Higgs boson exchange. The present experimental value of the muon anomaly is [1]



THE DAFNE STUDY GROUP

Then (~1998/99) came the idea of ISR...



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Available online at www.sciencedirect.com



Nuclear Physics B (Proc. Suppl.) 181–182 (2008) 280–285

NUCLEAR PHYSICS B
PROCEEDINGS
SUPPLEMENTS

www.elsevierphysics.com

Initial State Radiation: A success story

W. Kluge^{a*}

^aInstitut für Experimentelle Kernphysik, Universität Karlsruhe,
Postfach 3640, D 76021 Germany

The investigation of events with *Initial State Radiation (ISR)* and subsequent *Radiative Return* has become an impressively successful and guiding tool in low and intermediate energy hadron physics with electron positron colliders: it allows to measure hadronic cross sections and the ratio R from threshold up to the maximum energy of the colliders running at fixed energy, to clarify reaction mechanisms and reveal substructures (intermediate states and their decay mechanisms) and to search for new highly excited mesonic states with $J^{PC} = 1^{--}$. While being discussed since the sixties-seventies *ISR* became a powerful tool for experimentalists only with the development of *EVA-PHOKHARA*, a Monte Carlo generator developed over almost 10 years, while increasing its complexity, which is user friendly, flexible and easy to implement into the software of existing detectors.

Radiative return at NLO and the measurement of the hadronic cross-section in electron-positron annihilation

Germán Rodrigo^{1a}, Henryk Czyz^{2,3b}, Johann H. Kühn^{1,4c}, and Marcin Szopa²

¹ TH-Division, CERN, CH-1211 Geneva 23, Switzerland.

² Institute of Physics, University of Silesia, PL-40007 Katowice, Poland.

³ Institute of Advanced Study, University of Bologna, I-40138 Bologna, Italy

⁴ Institut für Theoretische Teilchenphysik, Universität Karlsruhe, D-76128 Karlsruhe, Germany.

Received: February 1, 2008

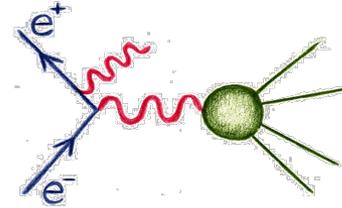
From Hadronic Cross Section to the measurement of the Vacuum Polarization at KLOE: a fascinating endeavour

Graziano Venanzoni^{1,*}

¹INFN Sezione di Pisa, Pisa, Italy

EPJ Web of Conferences **166**, 00021 (2018)

Abstract. The KLOE experiment at the ϕ -factory DAΦNE in Frascati is the first to have employed Initial State Radiation (ISR) to precisely determine the $e^+e^- \rightarrow \pi^+\pi^-(\gamma)$ cross section below 1 GeV. Such a measurement is particularly important to test the Standard Model (SM) calculation for the $(g-2)$ of the muon, where a long standing 3σ discrepancy is observed. I will review the ISR activity in KLOE in the last 20 years from the measurement of the hadronic cross section to the first direct determination of the time-like complex running $\alpha(s)$ in the region below 1 GeV.



Physics Letters B 459 (1999) 279–287

Measuring $\sigma(e^+e^- \rightarrow \text{hadrons})$ using tagged photons

S. Binner, J.H. Kühn, K. Melnikov¹

Institut für Theoretische Teilchenphysik, Universität Karlsruhe, D-76128 Karlsruhe, Germany

Eur. Phys. J. C 6, 637–645 (1999)

DOI 10.1007/s100529800953

THE EUROPEAN
PHYSICAL JOURNAL C
© Springer-Verlag 1999

The hadronic contribution to the muon $g - 2$ from hadron production in initial state radiation events at the e^+e^- collider DAΦNE

S. Spagnolo^a

Dipartimento di Fisica dell'Università di Lecce and INFN, Sezione di Lecce, via Arnesano I-73100 Lecce, Italy

KLOE MEMO n°195

August 13, 1999

Measurement of the hadronic cross section $\sigma(e^+e^- \rightarrow \pi^+\pi^-)$ from initial state radiative events $\pi^+\pi^-\gamma$ with the KLOE detector

Gabriella Cataldi, Achim Denig, Wolfgang Kluge, Graziano Venanzoni

Institut für Experimentelle Kernphysik, Universität Karlsruhe

Measurement of the hadronic cross-section with KLOE using a radiated photon in the initial state

G. Cataldi (Karlsruhe U.), A. Denig (Karlsruhe U.), W. Kluge (Karlsruhe U.), S. Muller (Karlsruhe U.), G. Venanzoni (Karlsruhe U.) (Nov, 1999)

Published in: *Frascati Phys.Ser.* 16 (1999) 569–578 • Contribution to: *3rd Workshop on Physics and Detectors for DAPHNE (DAPHNE 99)*, 569–578

ISR@Karlsruhe (1999-2000s)

- I was lucky to be in Karlsruhe from 1999-2000 (as EURODAFNEpostdoc) working on ISR at KLOE



Wolfgang Kluge



Achim Denig



Stefan Mueller



Debora Leone



Hans Kuhn



Fred Jegerlehner
(Berlin)

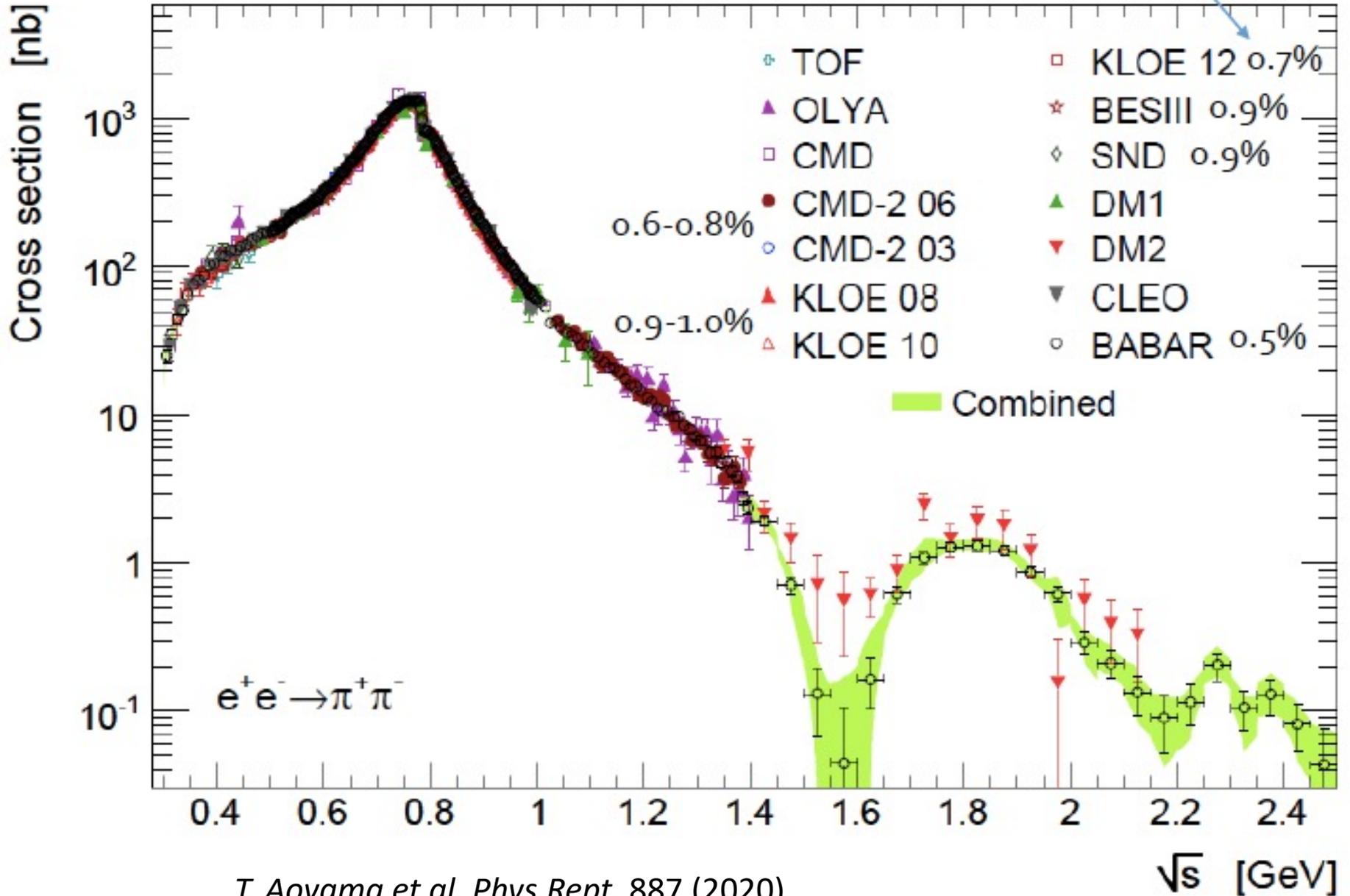


Barbara Valeriani



Paolo Beltrame
(Jesuit now)

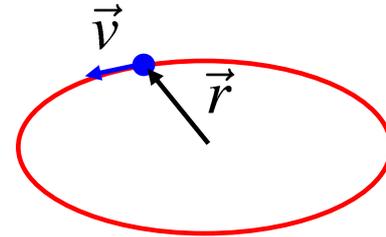
- Also important for ISR were the contributions of S. Spagnolo (Lecce), F. Nguyen (Roma 3), M. Incagli, S. Di Falco (Pisa), and H. Czyz and G. Rodrigo on theory...and many other (experimental and theory) colleagues joined later



The Muon g-2

- A charge particle in a plane orbit has **angular momentum** L and **magnetic moment** μ

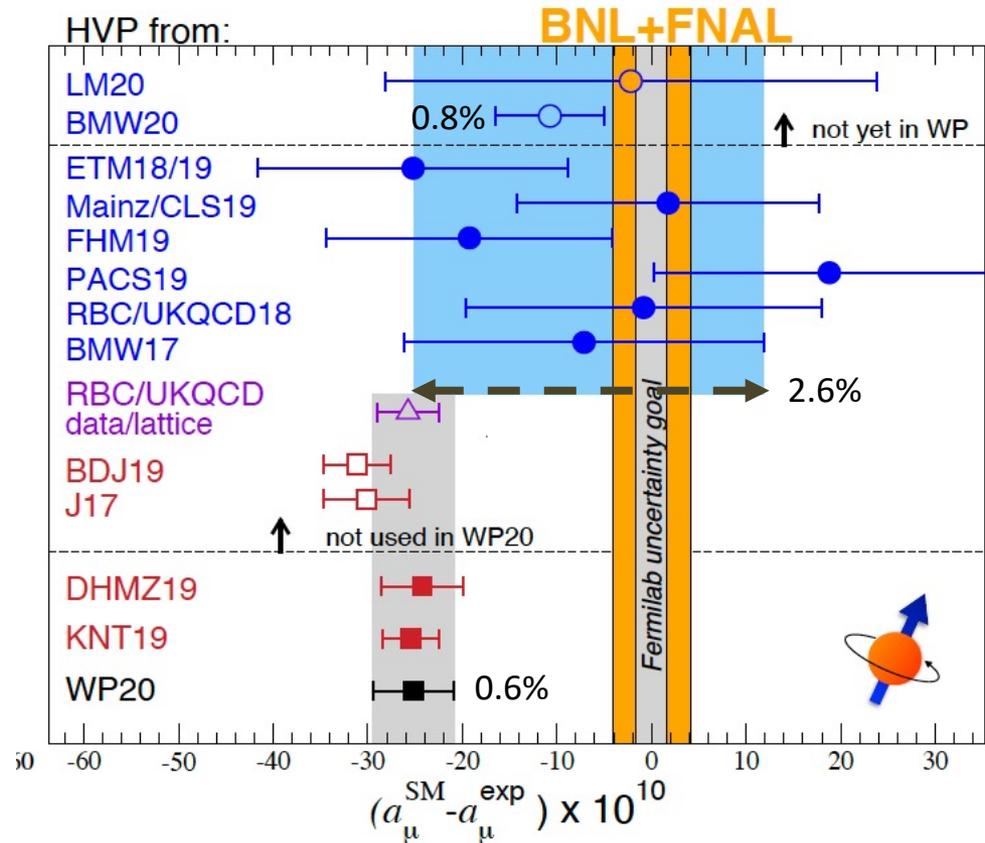
$$\mu = \frac{q}{2m} \vec{L}$$



- The ratio $\mu/(q/2m)L$ is called giromagnetic ratio g . Classically **$g=1$**
- For an elementary particle of Spin = 1/2 (e^- , μ) the eq. Dirac's predicts **$g = 2$**
$$\vec{\mu} = \frac{e}{2m} \vec{\sigma} \equiv g\mu_B \vec{S}; \quad \vec{S} = \vec{\sigma}/2, \quad g = 2$$
- The magnetic anomaly is defined as **$a = (g-2)/2$** . **$g=2 \rightarrow a=0$** according to Dirac

However: a_{μ}^{HLO}

- Uncertainty SM dominated by HVP (HLO) contribution
- Two approaches:
 - Dispersive (e^+e^-)
 - Lattice
- FNAL measurement is 4.2σ from SM prediction based on e^+e^-
- However using lattice: **BMW20**
 - 1.5σ tra exp e SM
- Hot topic and intensive work going on!
- MUonE experimental proposal at CERN!

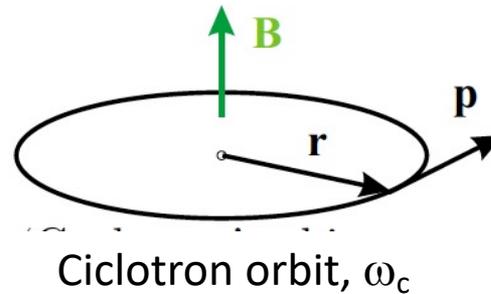


The Muon g-2 Theory Initiative has recently submitted a contribution to 2022 Snowmass Summer Study (<https://arxiv.org/pdf/2203.15810.pdf>) where they expect a main update in 2023. New results from the two different approaches (lattice and data-driven a_{μ}^{HLO}) are expected by 2025, possibly with reduced error

$a_\mu = (g-2)/2$ can also be measured very precisely...

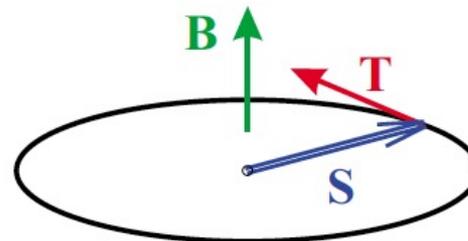
- A charged particle with spin put in a magnetic field (uniform) rotates in a circular orbit with angular frequency (called cyclotron):

$$\omega_c = \frac{qB}{m}$$



- The presence of the magnetic field acts on the spin by rotating it around the field direction (precession frequency of the spin)

$$\omega_s = g \frac{qB}{2m}$$

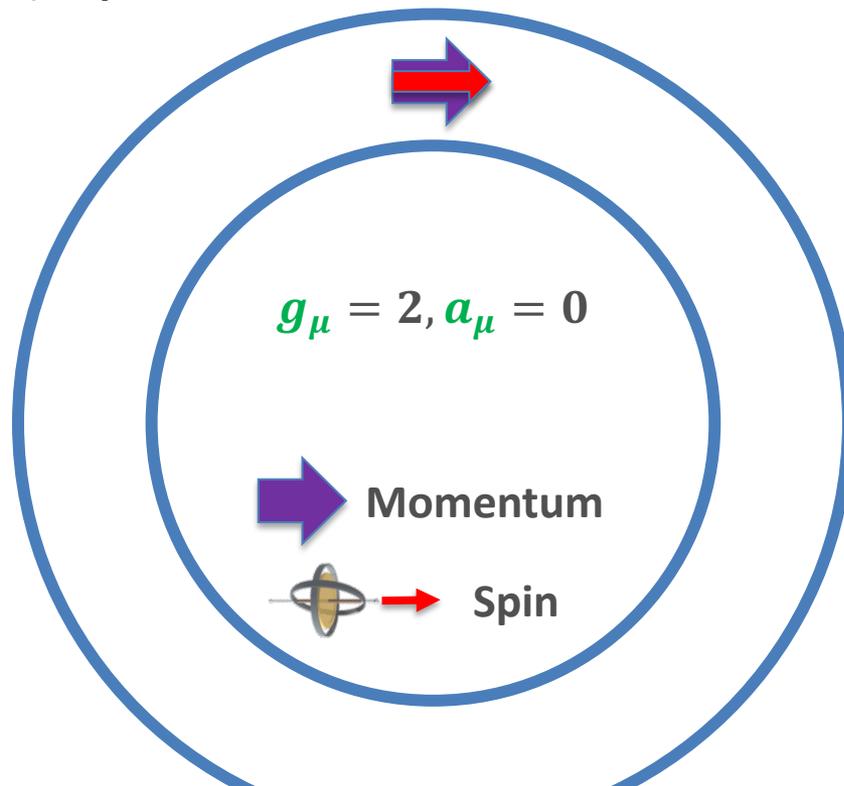


How to measure the muon anomaly?

- The frequency with which the spin moves ahead of the momentum in a magnetic field B (anomalous precession frequency ω_a) is:

$$\omega_a = \omega_s - \omega_c = a \frac{eB}{m}$$

- If $g=2$ ($a=0$) spin remains locked to momentum

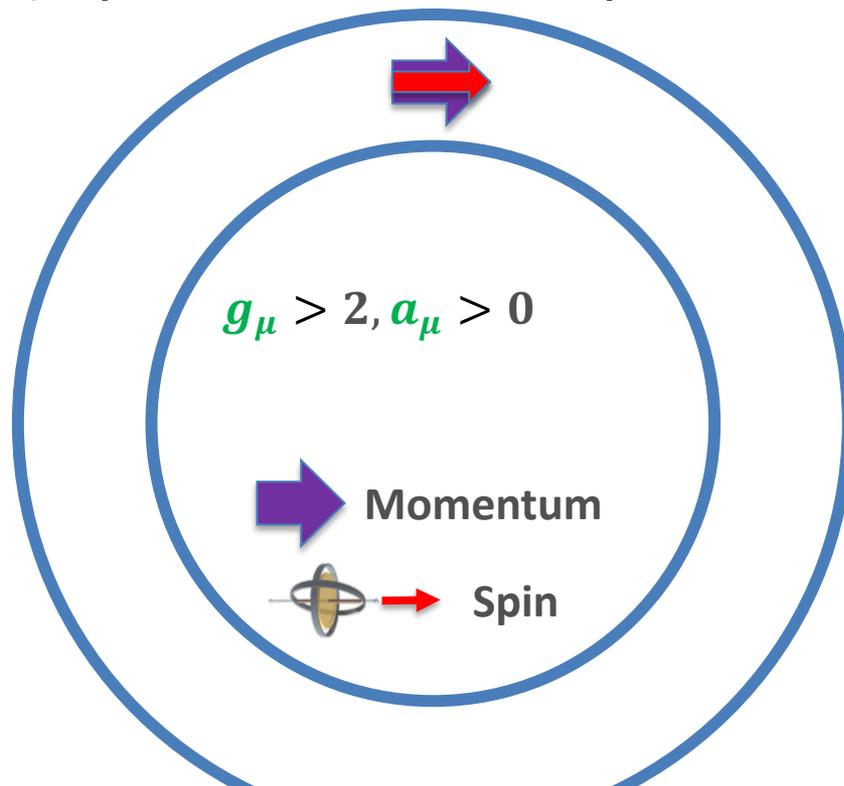


How to measure the muon anomaly?

- The frequency with which the spin moves ahead of the momentum in a magnetic field B (anomalous precession frequency ω_a) is:

$$\omega_a = \omega_s - \omega_c = a \frac{eB}{m}$$

- If $g > 2$ ($a > 0$) spin advances respect to the momentum



Current experiments
 $\delta a_\mu < 1\text{ppm}$

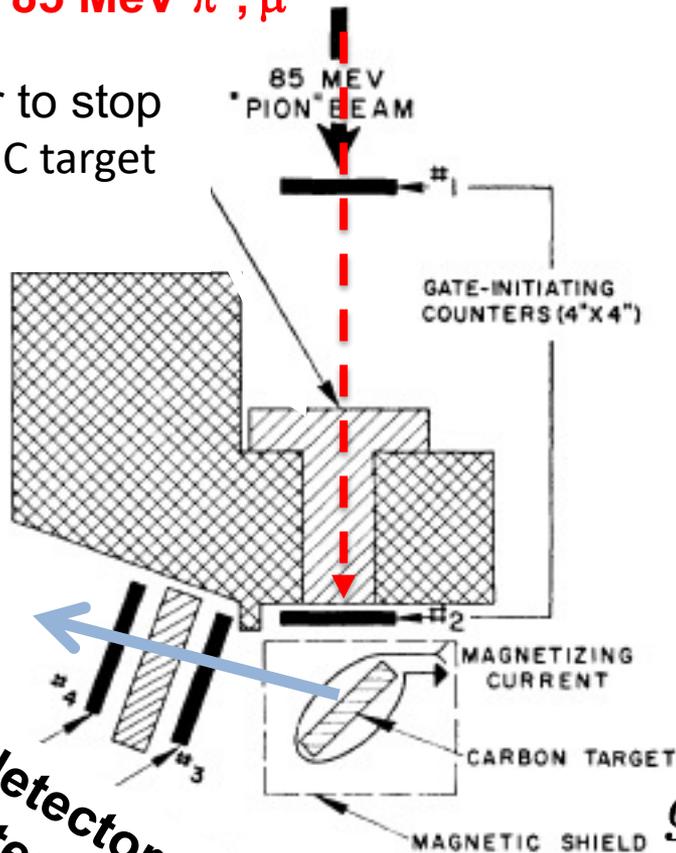
A (brief) history of the muon $g-2$ experiments

History: the first measurement of g_μ

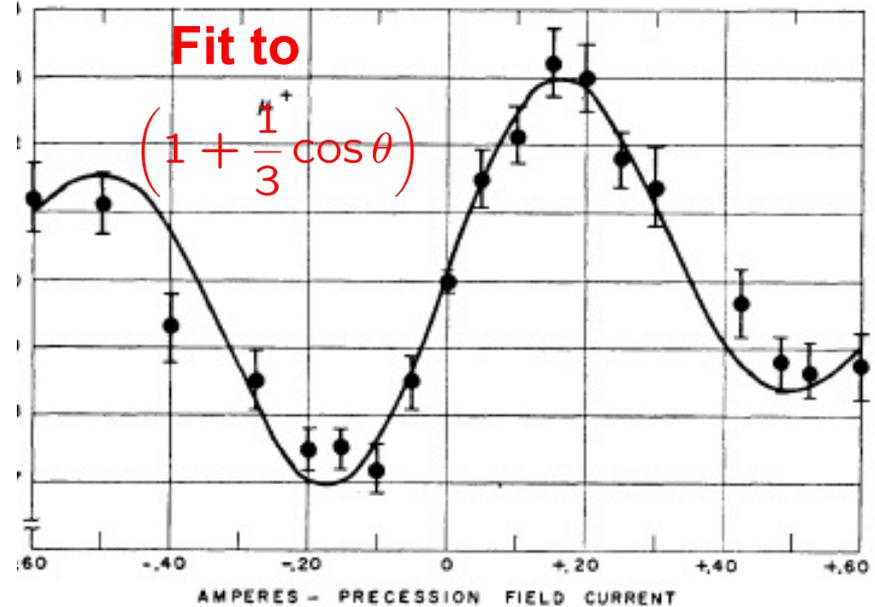
- 1957: Garwin, Lederman, Weinrich at Nevis (Just after Yang and Lee parity violation paper - confirmation)

85 MeV π^+, μ^+

degrader to stop π^+ before C target



Direct measurement of g -- asym vs field



$$g_\mu = 2.00 \pm 0.10$$

5% uncertainty

muons behave like electrons

F. Farley, E. Picasso The Muon ($g-2$) Experiments at CERN
Ann.Rev.Nucl.Part.Sci. 29 (1979) 243-282

The CERN muon $g-2$ experiments (1960-1979)

F.J.M. Farley, Y.K. Semertzidis / Progress in Particle and Nuclear Physics 52 (2004) 1-83



Fig. 10. The first experimental magnet in which muons were stored at CERN for up to 30 turns. Left to right: Georges Charpak, Francis Farley, Bruno Nicolai, Hans Sens, Antonio Zichichi, Carl York and Richard Garwin.

The 47 years of muon $g - 2$

F.J.M. Farley^{a,*}, Y.K. Semertzidis^b

^aYale University, New Haven, CT 06520, USA

^bBrookhaven National Laboratory, Upton, NY 11973, USA

Received 30 October 2003

The history of the muon ($g - 2$) experiments

B. Lee Roberts*

21

SciPost Phys. Proc. 1, 032 (2019)

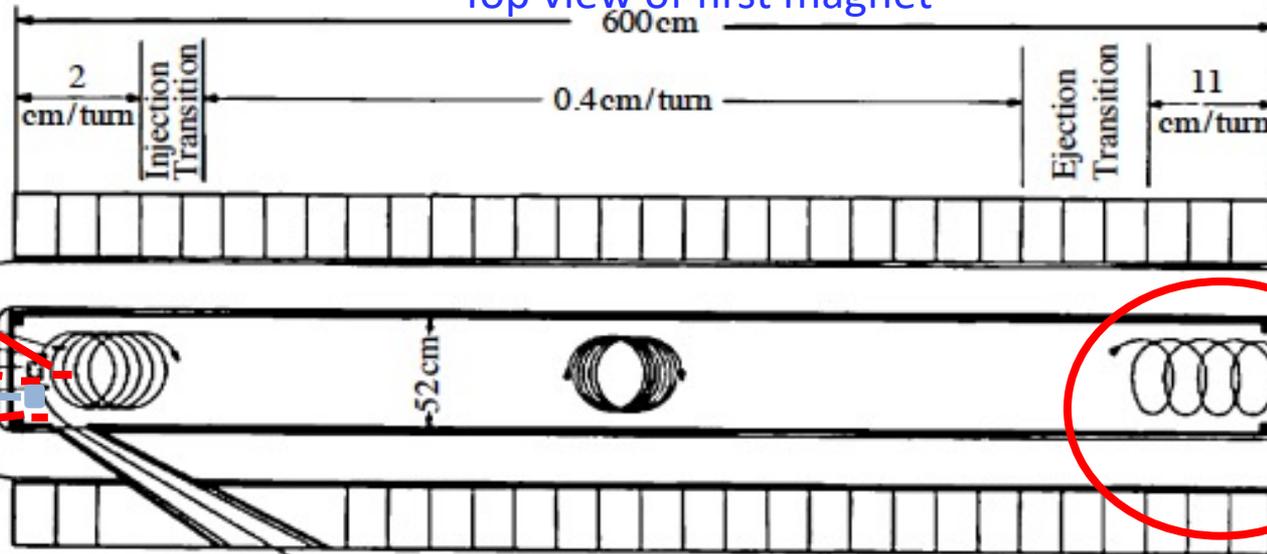
They measure a_μ since
 they measure the spin
 relative to the
 momentum

$$\begin{aligned}\vec{\omega}_a &= \omega_S - \omega_C = \\ &= -\frac{Qe}{m} a_\mu \vec{B}\end{aligned}$$

CERN I, 1958-1962

Top view of first magnet

With 100 MeV/c muons

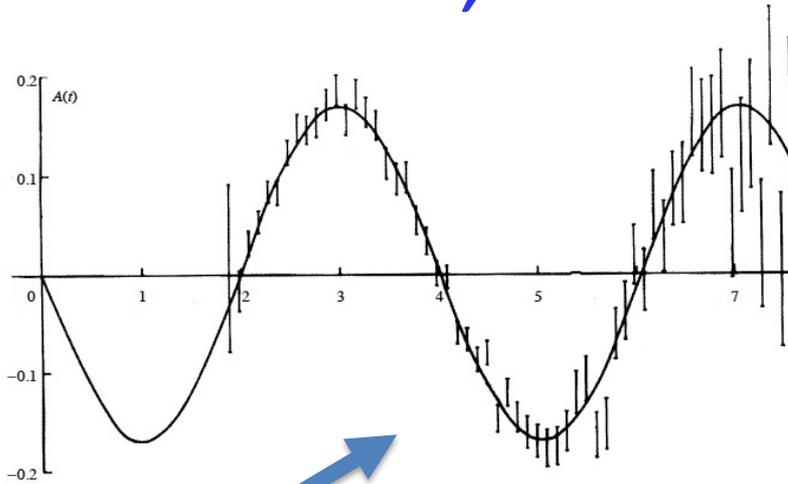


Incident μ 1 · 2 · 3

Exiting μ 4 · 6 · 6' · 5 · 7 · 7'
followed by 7 · 7' or 6' · 6

- Inject polarized muon into a long magnet ($B \approx 1.5$ T) with a small gradient – particles drift in circular orbits to the other end: **7.5 μ s = 1600 turns**
- Extract muons with a large gradient into a polarization monitor where they stopped
- Time in the magnetic field was measured by counters
- Measure the time dependent forward-backward decay asymmetry

CERN I, 1958-1962



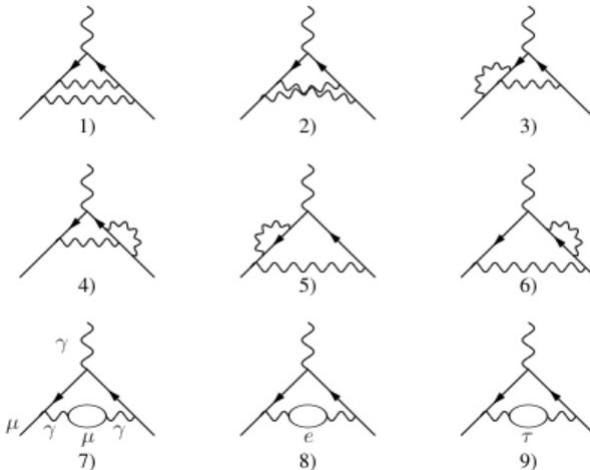
$$A(t) = A_0 \sin(\omega_s t + \varphi)$$

$$\omega_s = a_\mu (e/mc) \bar{B}$$

$\simeq 10^3 \mu^+$ recorded

$$a_\mu(\text{expt}) = 0.001162(5) \quad (4300 \text{ ppm})$$

$$a_\mu(\text{theory}) = 0.001165$$



QED, second order

$$a_\mu \approx 0.5 \left(\frac{\alpha}{\pi}\right) + 0.766 \left(\frac{\alpha}{\pi}\right)^2$$

CERN II, 1962-1968: The First Storage Ring (proton beam) π production target inside

Top view of the second magnet

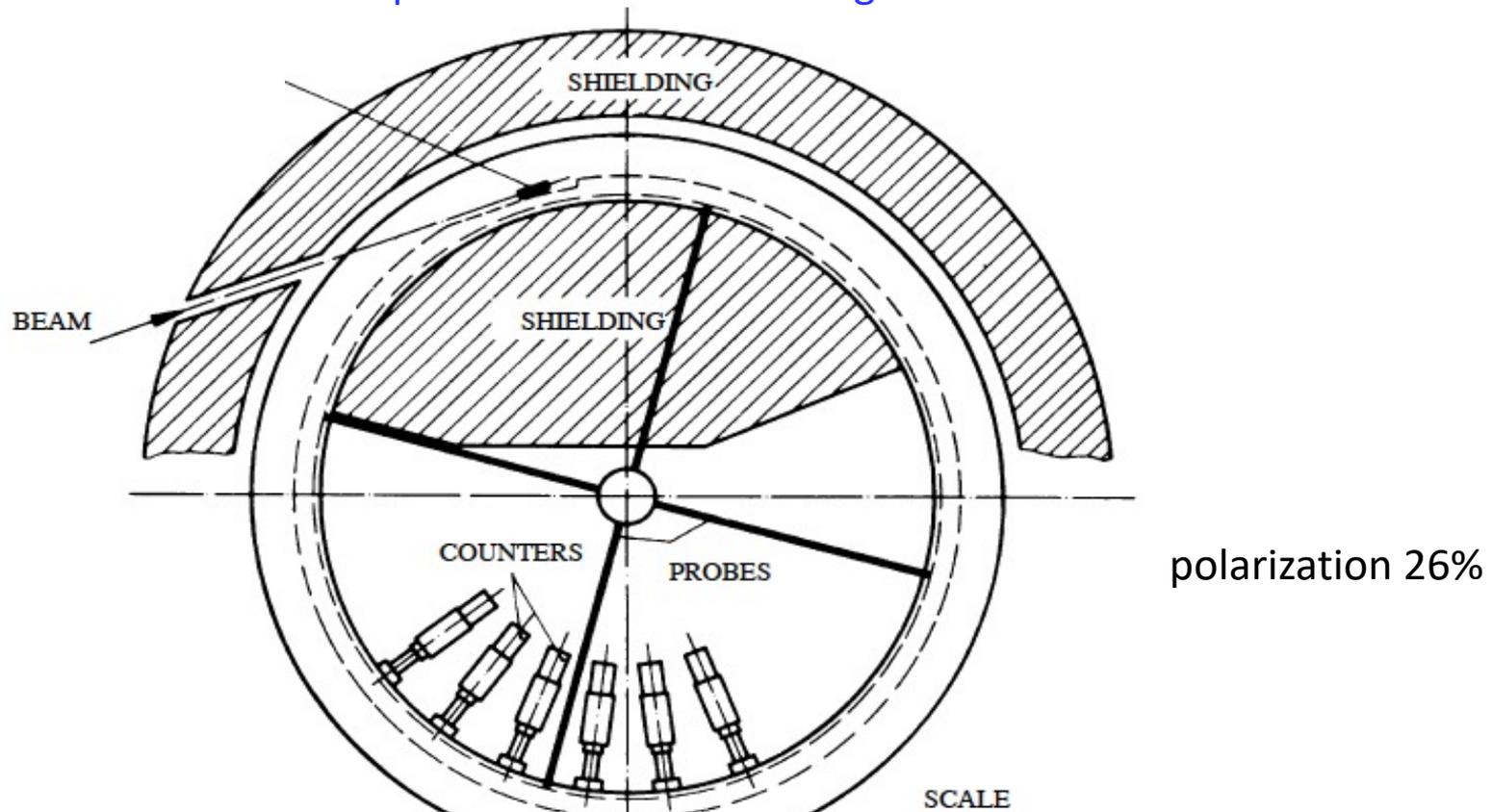
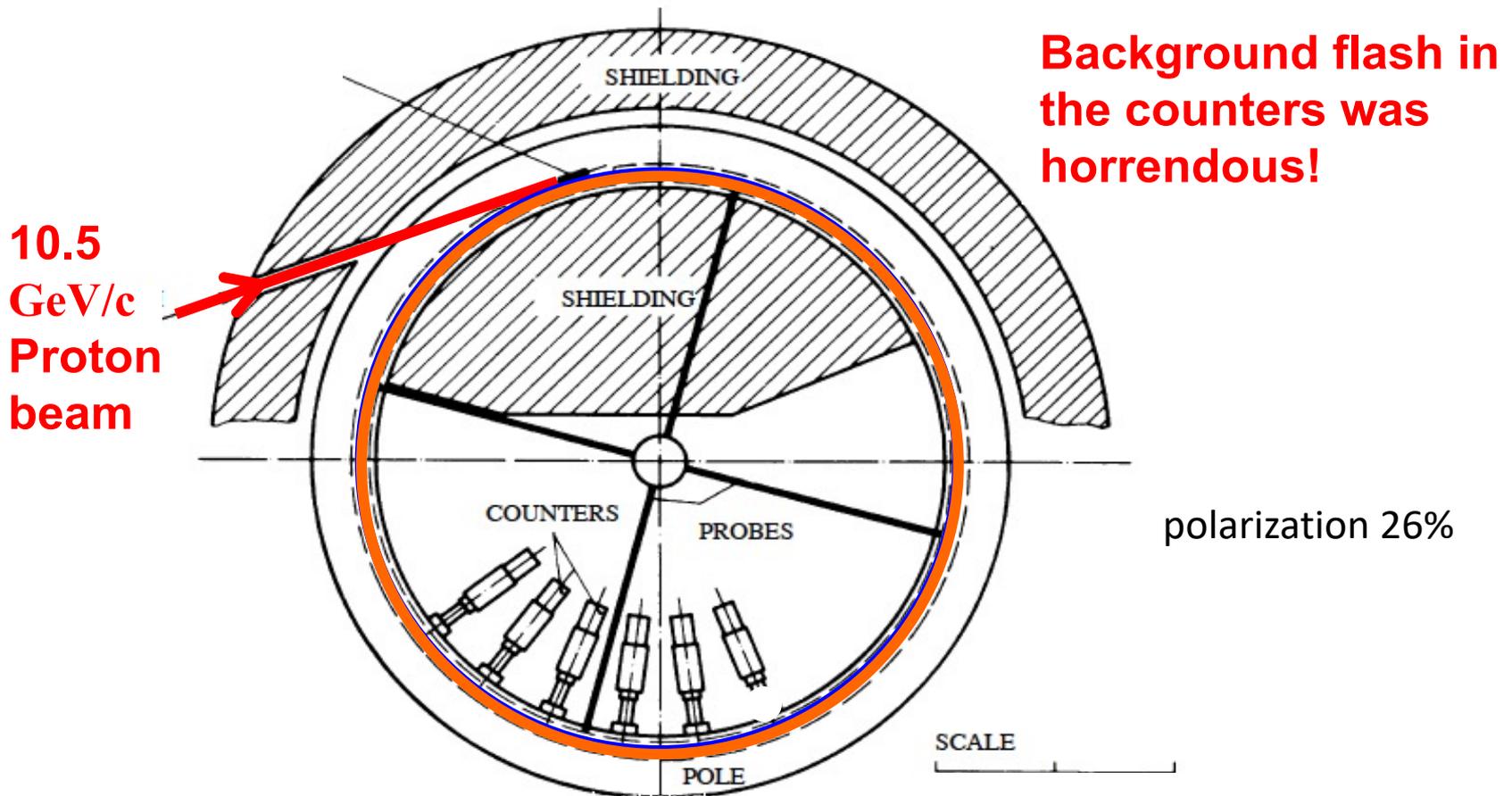


Fig. 17. The first muon storage ring: diameter 5 m, muon momentum 1.3 GeV/c, time dilation factor 12. The injected pulse of 10.5 GeV protons produces pions at the target, which decay in flight to give muons.

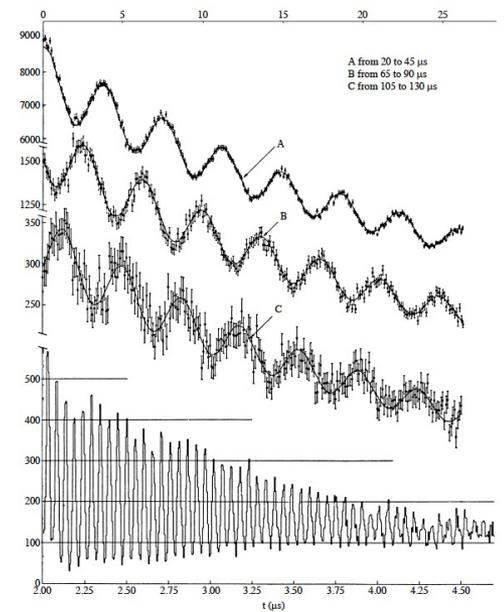
- Go to $p_{\mu} = 1.27 \text{ GeV}/c$, $\gamma_{\mu} = 12$; $\gamma\tau = 27 \mu\text{s}$;
- Used a weak-focusing magnetic storage ring; $B_z = 1.71 \text{ T}$
- $p + N \rightarrow \pi \rightarrow \mu$ which are stored



CERN II, 1962-1968

$$f(t) \simeq N_0 e^{-\lambda t} [1 + A \cos(\omega_a t + \phi)]$$

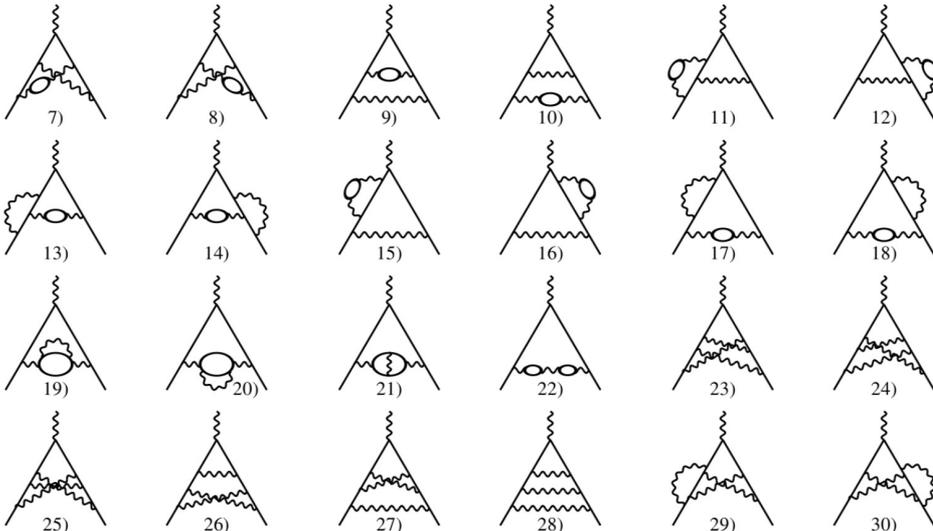
$$\frac{\delta\omega_a}{\omega_a} = \frac{\sqrt{2}}{\omega_a A \gamma \tau \sqrt{N}}$$



$$a_\mu(\text{expt}) = 0.00116616(31) \quad (266 \text{ ppm})$$

$$a_\mu(\text{theory}) = 0.00116587(3)$$

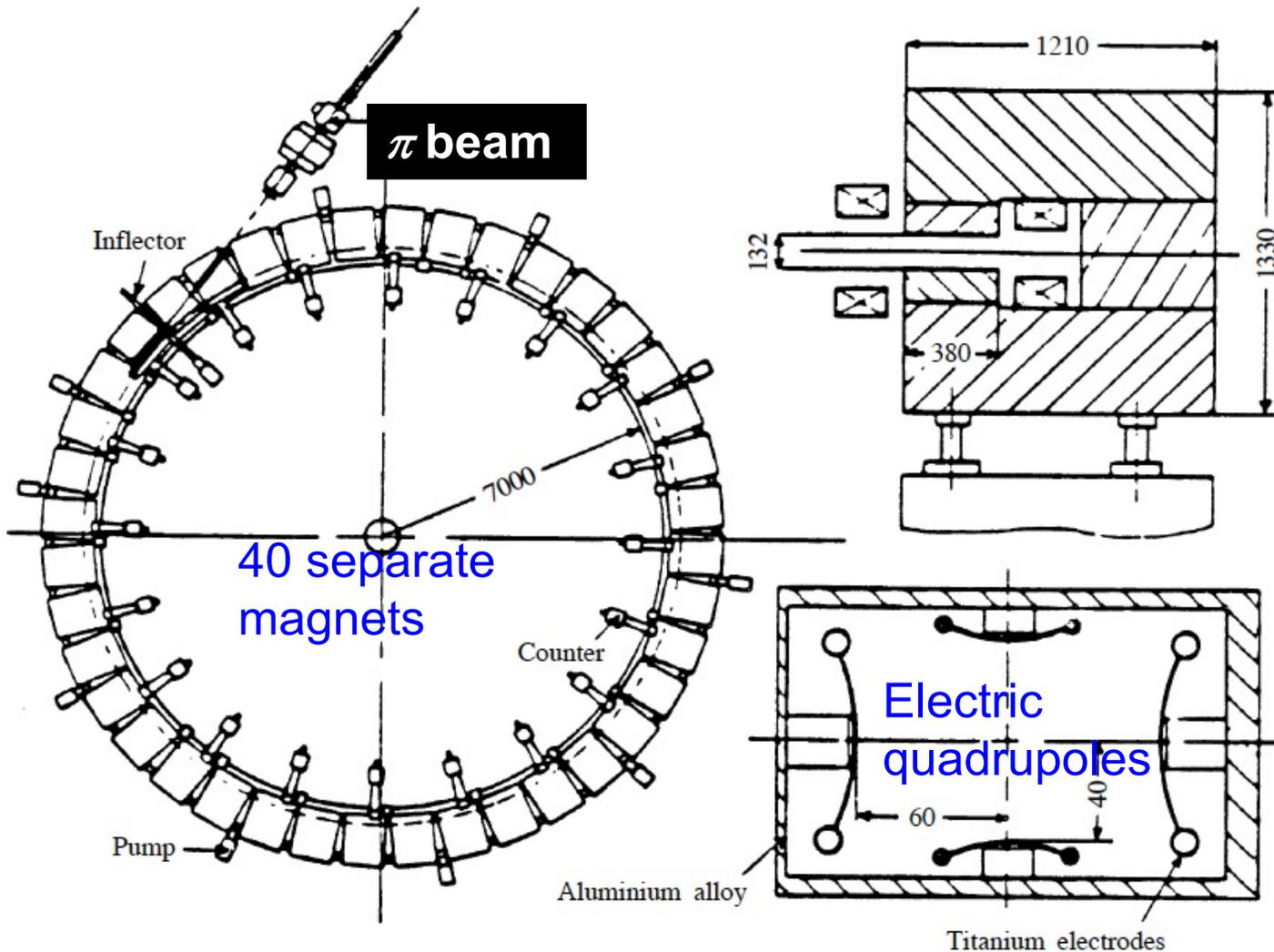
QED, third order



$$a_\mu \approx 0.5 \left(\frac{\alpha}{\pi}\right) + 0.766 \left(\frac{\alpha}{\pi}\right)^2 + 24.050 \left(\frac{\alpha}{\pi}\right)^3$$

CERN III, 1969-1976

The third magnet, second storage ring. Pion injection, E-field focusing, Magic momentum

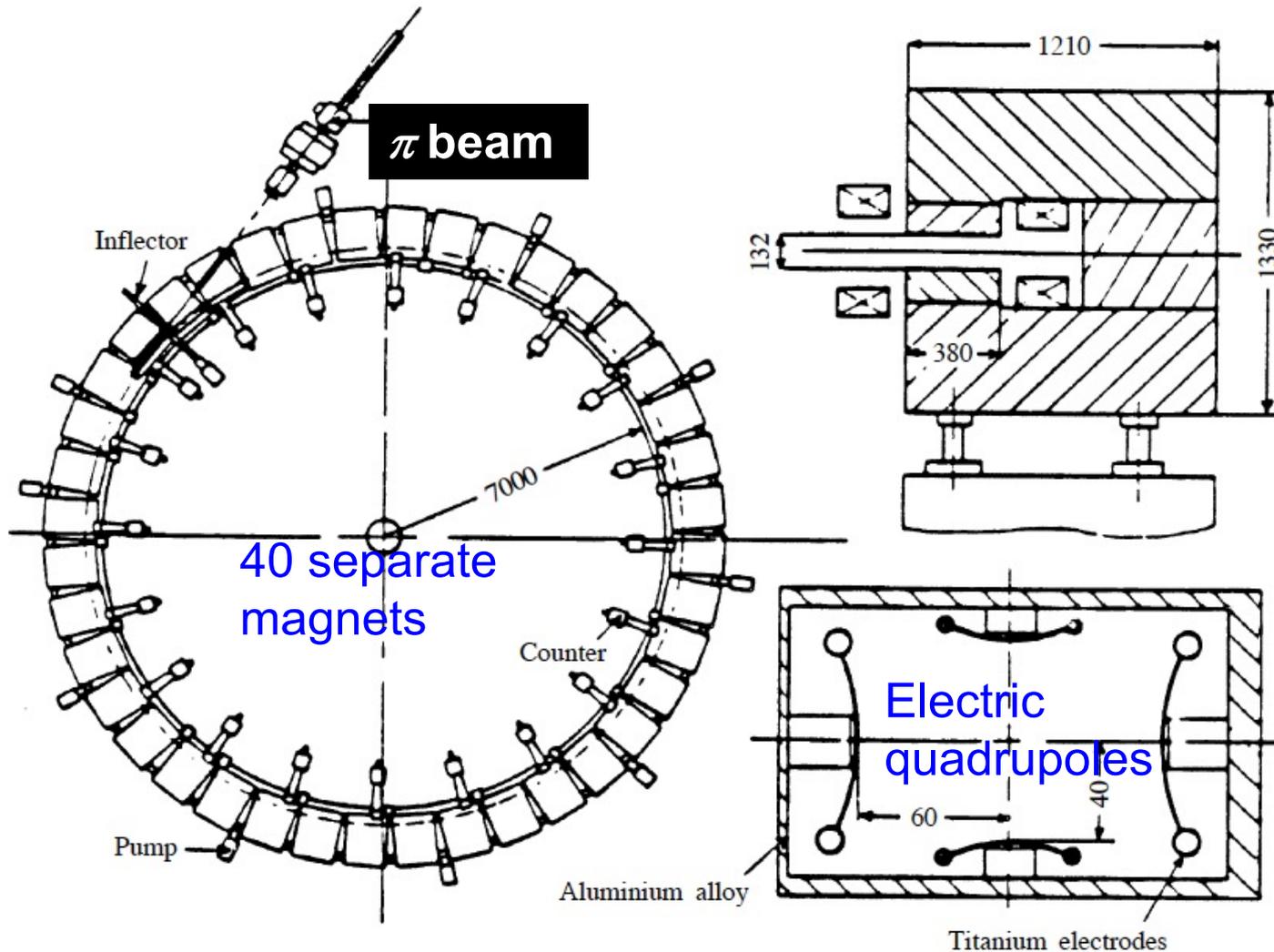


Still have pion flash at injection!

Not as bad as for CERN2

CERN III, 1969-1976

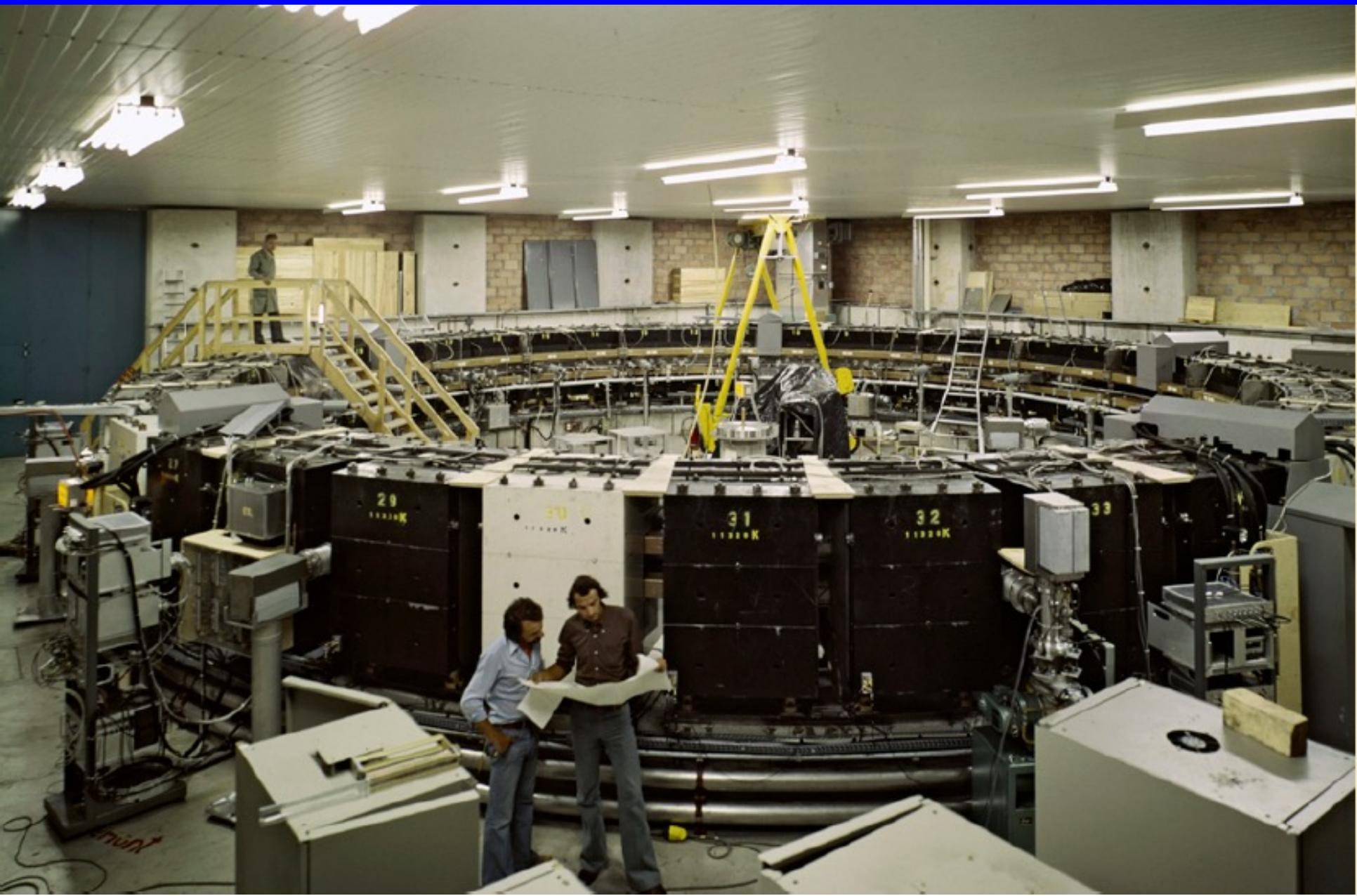
- Inject pions at 3.2 GeV Muon lifetime dilates to 64 μ s
- Use $\pi \rightarrow \mu$ decay to kick muons onto stable orbits



Still have pion flash at injection!

Not as bad as for CERN2

3rd Muon g-2 experiment at Cern



CERN III, 1969-1976

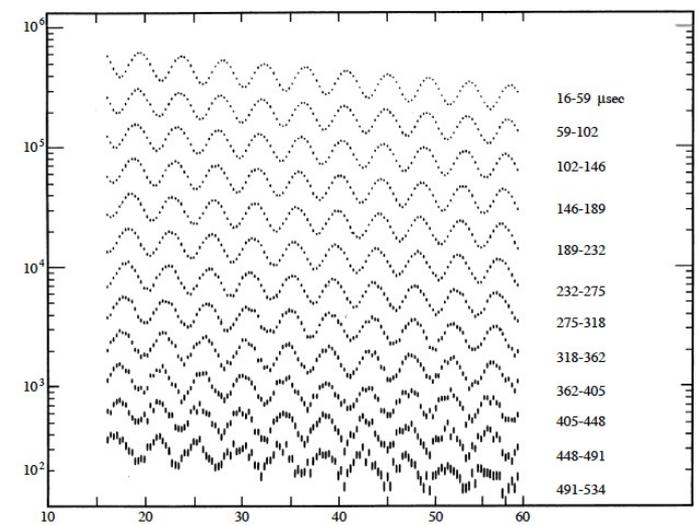
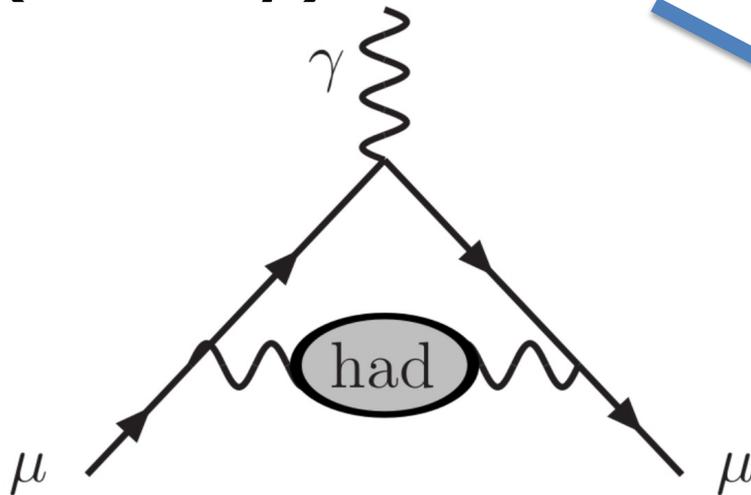


Fig. 25. The second muon storage ring: decay electron counts versus time (in microseconds) after injection. The range of time for each line is shown on the right (in microseconds).

$$a_{\mu}(\text{expt}) = 0.001165924(9) \quad (7.3 \text{ ppm})$$

$$a_{\mu}(\text{theory}) = 0.001165921(13)$$



HVP (hadronic vacuum polarization)

$$a_{\mu}^{\text{HAD}} \sim 700 \cdot 10^{-10} \quad (\sim 60 \text{ ppm})$$

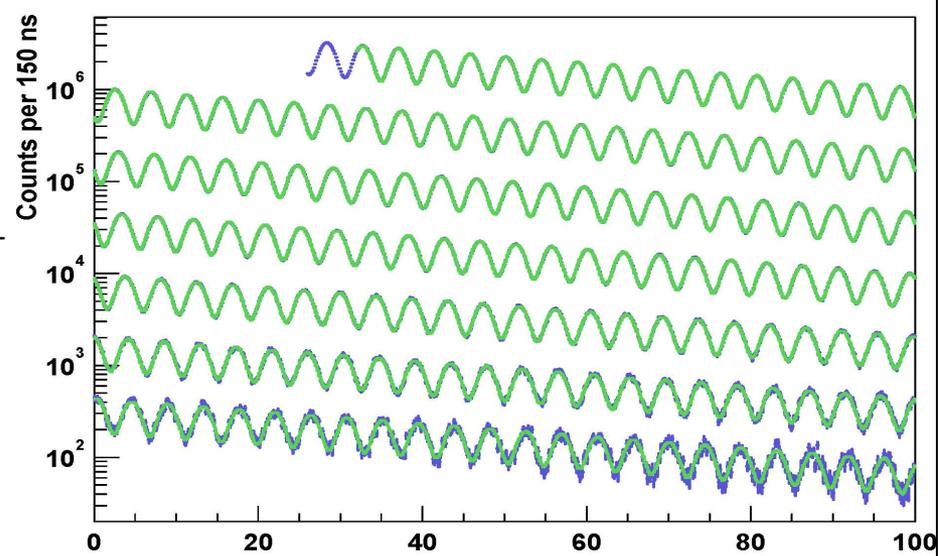
1984-2001: Measurement of a_μ at BNL (E821)

The measurement of the g-2 of the muon has been repeated with x15 better accuracy at Brookhaven National Laboratory (USA)



E821, 1984-2001

$$3.6 \times 10^9 e^-$$

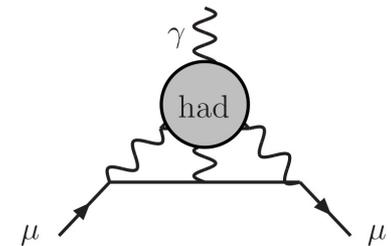
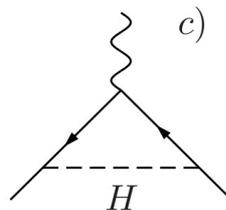
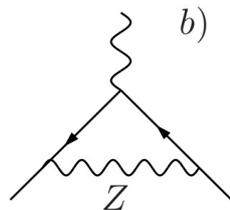
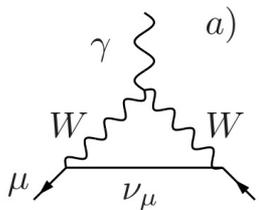


(2006)

$$a_\mu(\text{expt}) = 0.00116592089(63) \quad (0.54 \text{ ppm})$$

$$a_\mu(\text{theory})_{\text{(HMNT 03)}} = 0.00116591820(73) \quad \sim 3\sigma \text{ btw TH and EXP!}$$

EW ↙ ↘ Hadronic light-by-light



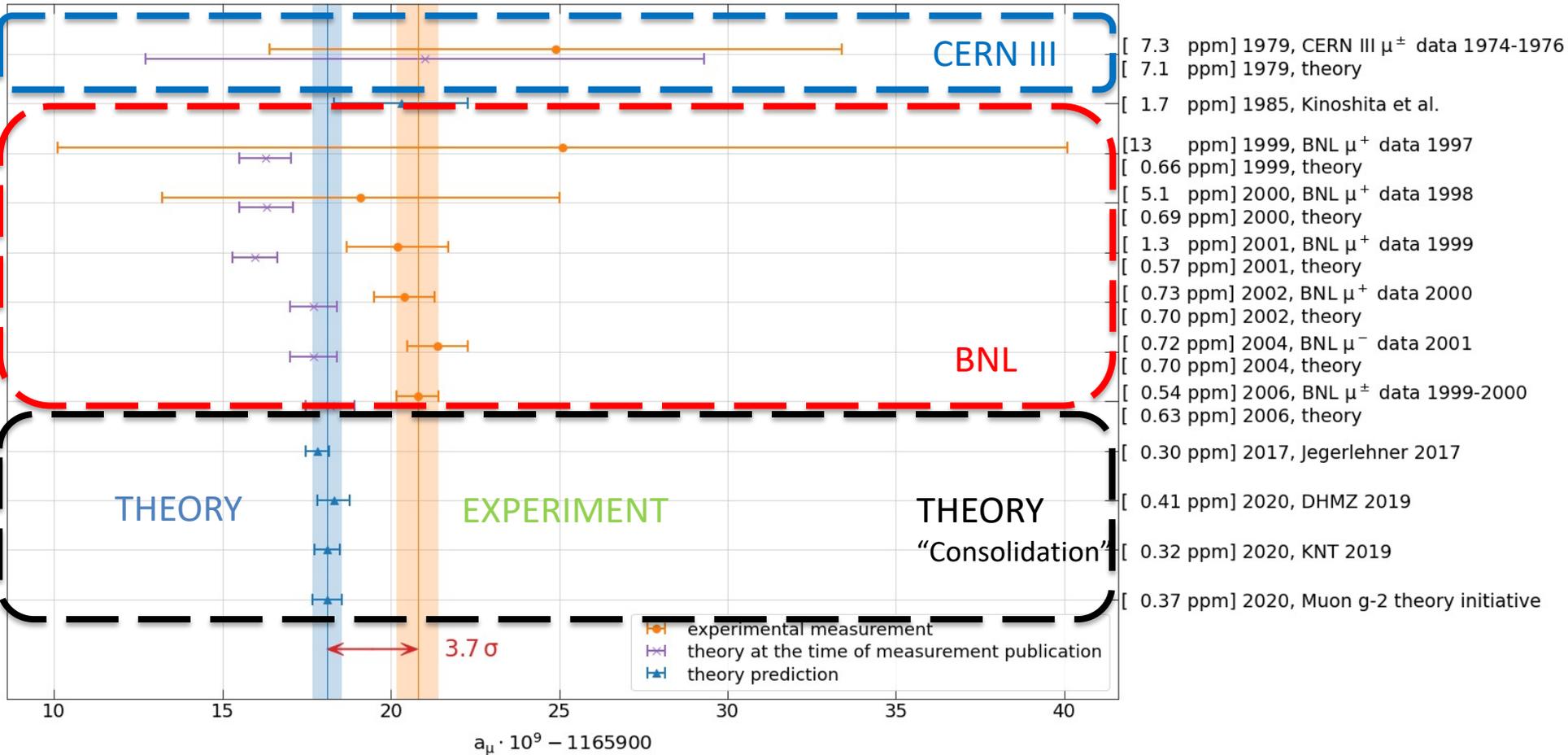
$$a_\mu^{\text{EW}} \sim 15.3 \cdot 10^{-10} \quad (\sim 1.3 \text{ ppm})$$

$$a_\mu^{\text{HLbL}} \sim 9.2 \cdot 10^{-10} \quad (\sim 0.8 \text{ ppm})$$

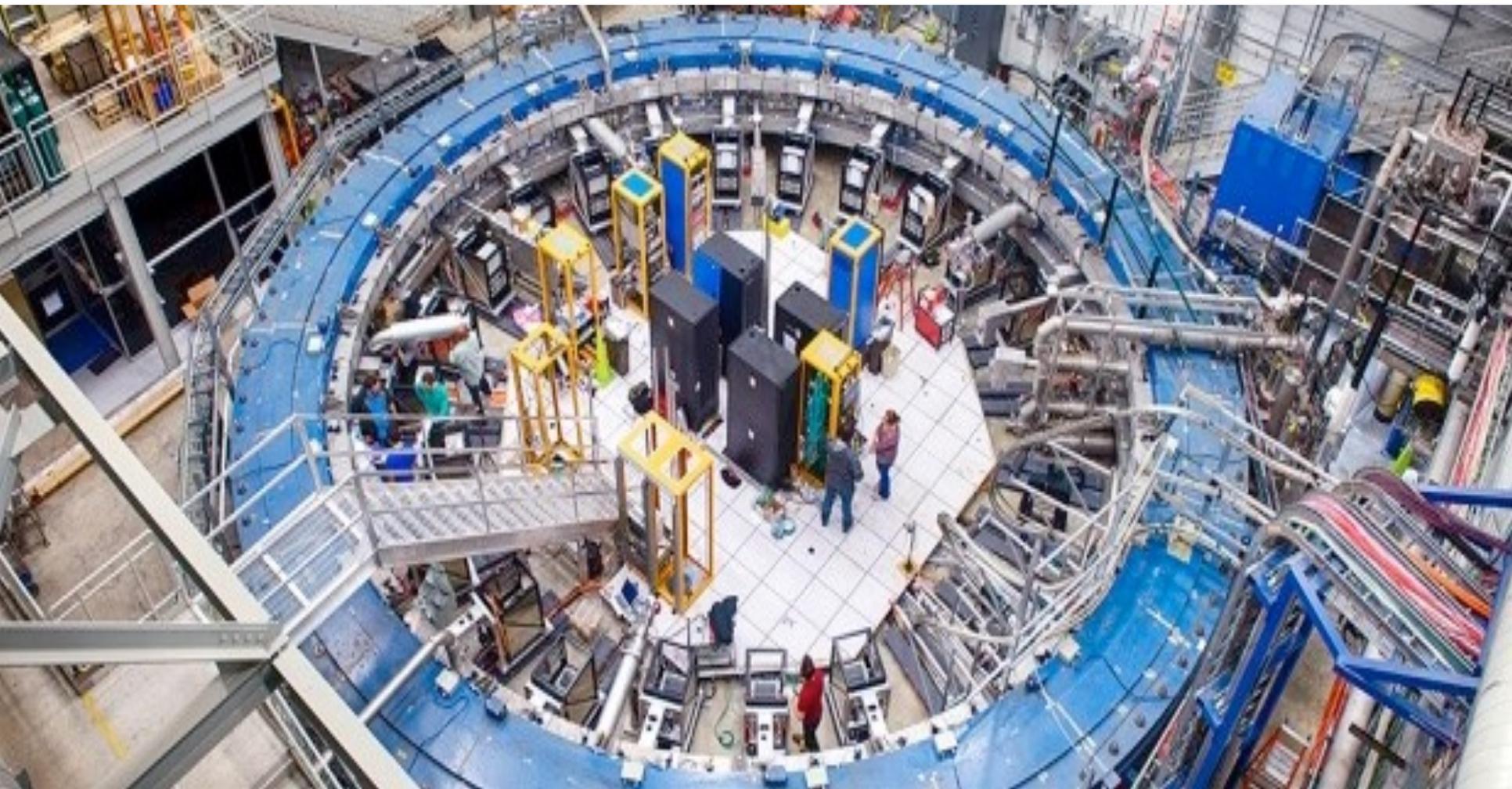
Summary of g-2 theory and measurements (2020)

History of muon anomaly measurements and predictions

$a_\mu = (g-2)/2 = \text{Muon (magnetic) anomaly}$



Fermilab Muon g-2 Storage Ring (2009- present)



4 key elements for E989 at FNAL

- Consolidated method (same ring of the BNL experiment)
- More muons (x20)
- improved beam and detector → Reduced systematics
- New crew → new ideas

• E821 at Brookhaven

$$\left. \begin{array}{l} \sigma_{\text{stat}} = \pm 0.46 \text{ ppm} \\ \sigma_{\text{syst}} = \pm 0.28 \text{ ppm} \end{array} \right\} \sigma = \pm 0.54 \text{ ppm}$$

• E989 at Fermilab $\hookrightarrow 0.2\omega_a \oplus 0.17\omega_p$

$$\left. \begin{array}{l} \sigma_{\text{stat}} = \pm 0.1 \text{ ppm} \\ \sigma_{\text{syst}} = \pm 0.1 \text{ ppm} \end{array} \right\} \sigma = \pm 0.14 \text{ ppm}$$

$0.07\omega_a \oplus 0.07\omega_p$

Key ingredients

1) Polarized muons

~97% polarized for forward decay



2) Precession proportional to $(g-2)$

$$\omega_a = \omega_{spin} - \omega_{cyclotron} = \left(\frac{g-2}{2} \right) \frac{eB}{mc}$$

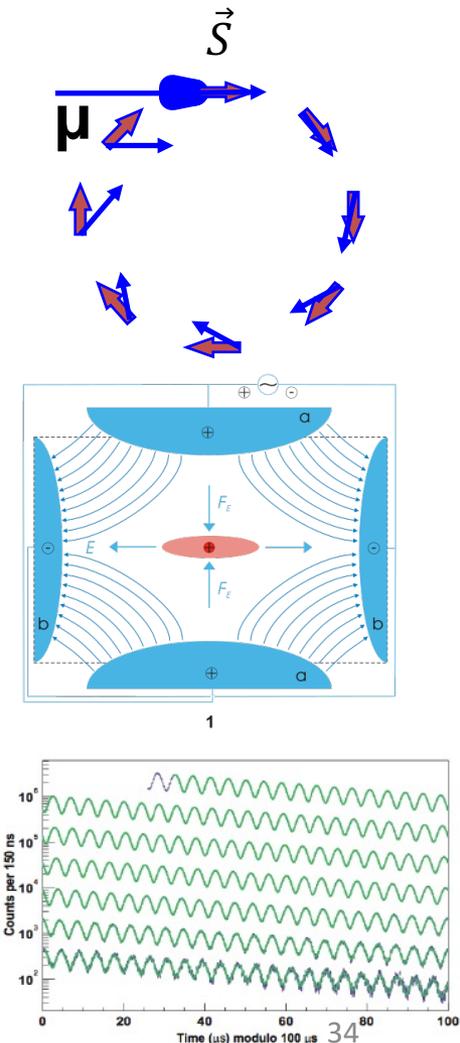
Measure 2 quantities $a_\mu = (g-2)/2$

3) P_μ magic momentum = 3.09 GeV/c

$$\vec{\omega}_a = \frac{e}{mc} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} \right]$$

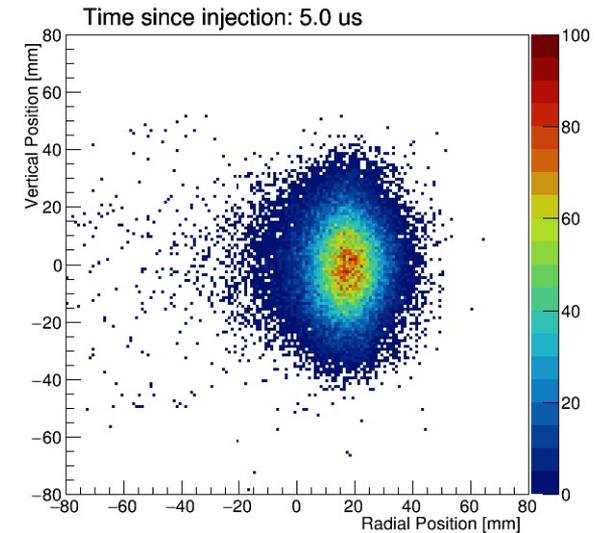
E field doesn't affect muon spin when $\gamma = 29.3$

4) Decay e^+ emitted preferably in spin direction of the muon



However there are beam dynamics effects

- The muon beam oscillates and breathes as a whole
- The full equation is more complex and corrections due to radial (x) and vertical (y) beam motion are needed



$$\vec{\omega}_a = \vec{\omega}_s - \vec{\omega}_c =$$

$$= -\frac{e}{mc} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} \right]$$

- Running at $\gamma_{\text{magic}}=29.3$ ($p=3.094$ GeV/c) this coefficient is null
- Because of momentum spread ($<0.2\%$) \rightarrow **E-field Correction**

- Vertical beam oscillation \rightarrow **Pitch correction**

Extracting a_μ (simplified)

By expressing B in terms of the precession frequency ω_p' of a proton shielded in a spherical water sample:

$$a_\mu = \frac{\omega_a}{\tilde{\omega}_p'} \frac{\mu_p' m_\mu g_e}{\mu_e m_e 2}$$

External (precise) data

$$B = \frac{\hbar \omega_p'}{2\mu_p'}$$

$$e = \frac{4m_e \mu_e}{\hbar g_e}$$

$$R' = \frac{\omega_a}{\tilde{\omega}_p'}$$

ratio of muon to proton precession
in the same magnetic dipole field

$\tilde{\omega}_p'$ = Proton Larmor precession frequency **weighted for the muon distribution**

Muon g-2 collaboration



USA

- Boston
- Cornell
- Illinois
- James Madison
- Kentucky
- Massachusetts
- Michigan
- Michigan State
- Mississippi
- North Central
- Northern Illinois
- Regis
- Virginia
- Washington

USA National Labs

- Argonne
- Brookhaven
- Fermilab



China

- Shanghai Jiao Tong



Germany

- Dresden
- Mainz



Italy

- Frascati
- Molise
- Naples
- Pisa
- Roma Tor Vergata
- Trieste
- Udine



Korea

- CAPP/IBS
- KAIST



Russia

- Budker/Novosibirsk
- JINR Dubna



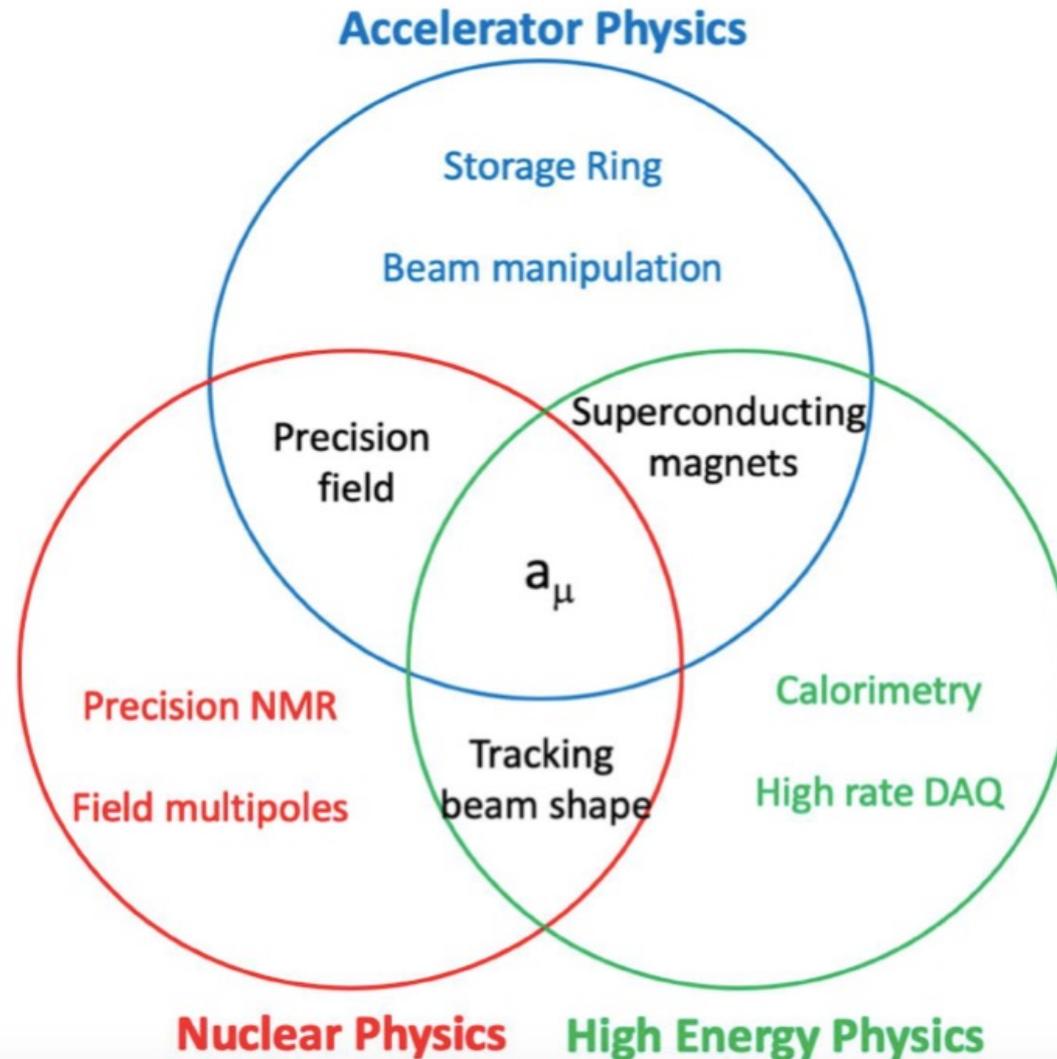
United Kingdom

- Lancaster/Cockcroft
- Liverpool
- Manchester
- University College London



>200 collaborators
35 Institutions
7 countries

Three different communities to measure a_μ

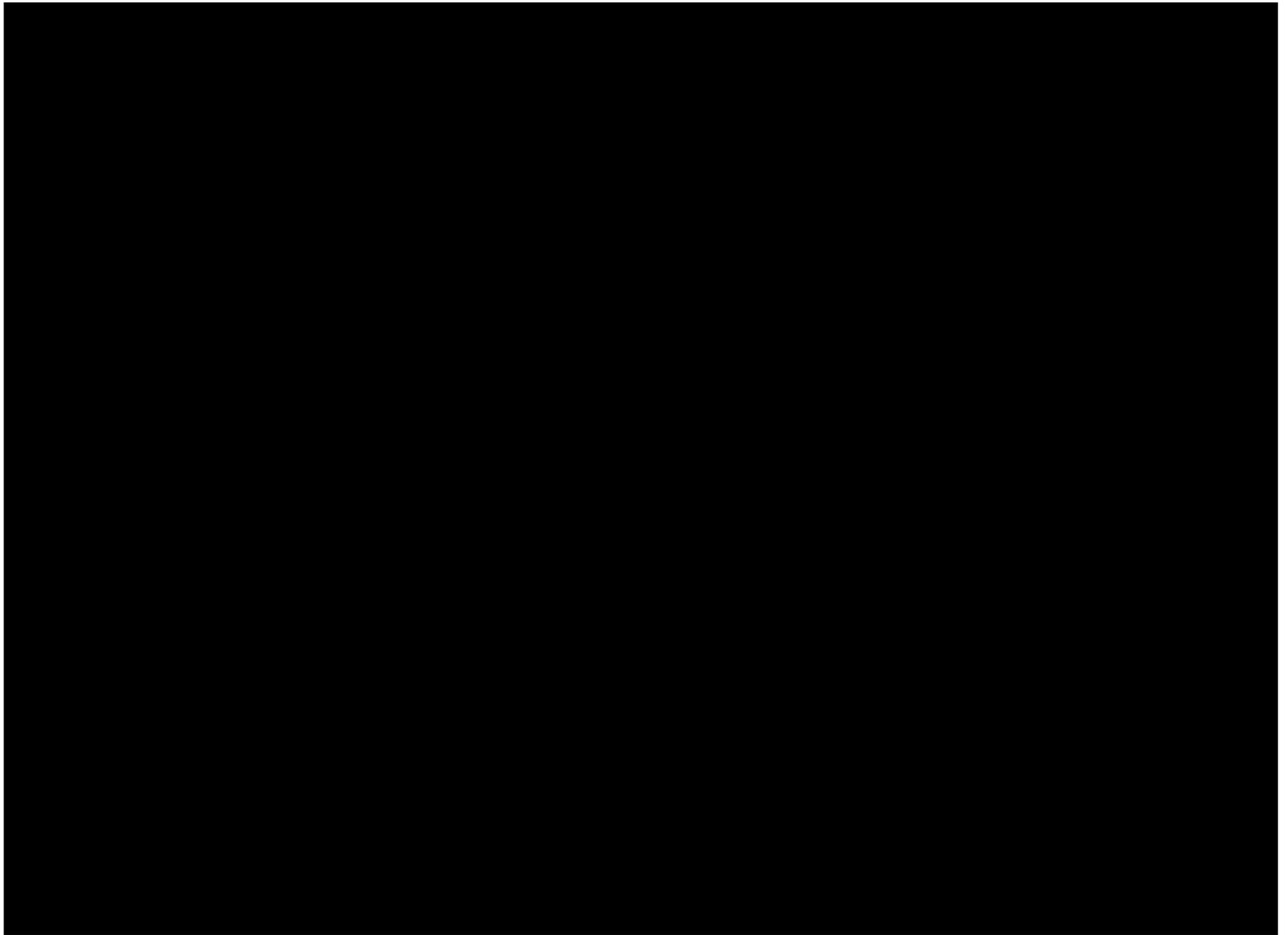


June 2013: The ring leaves from BNL

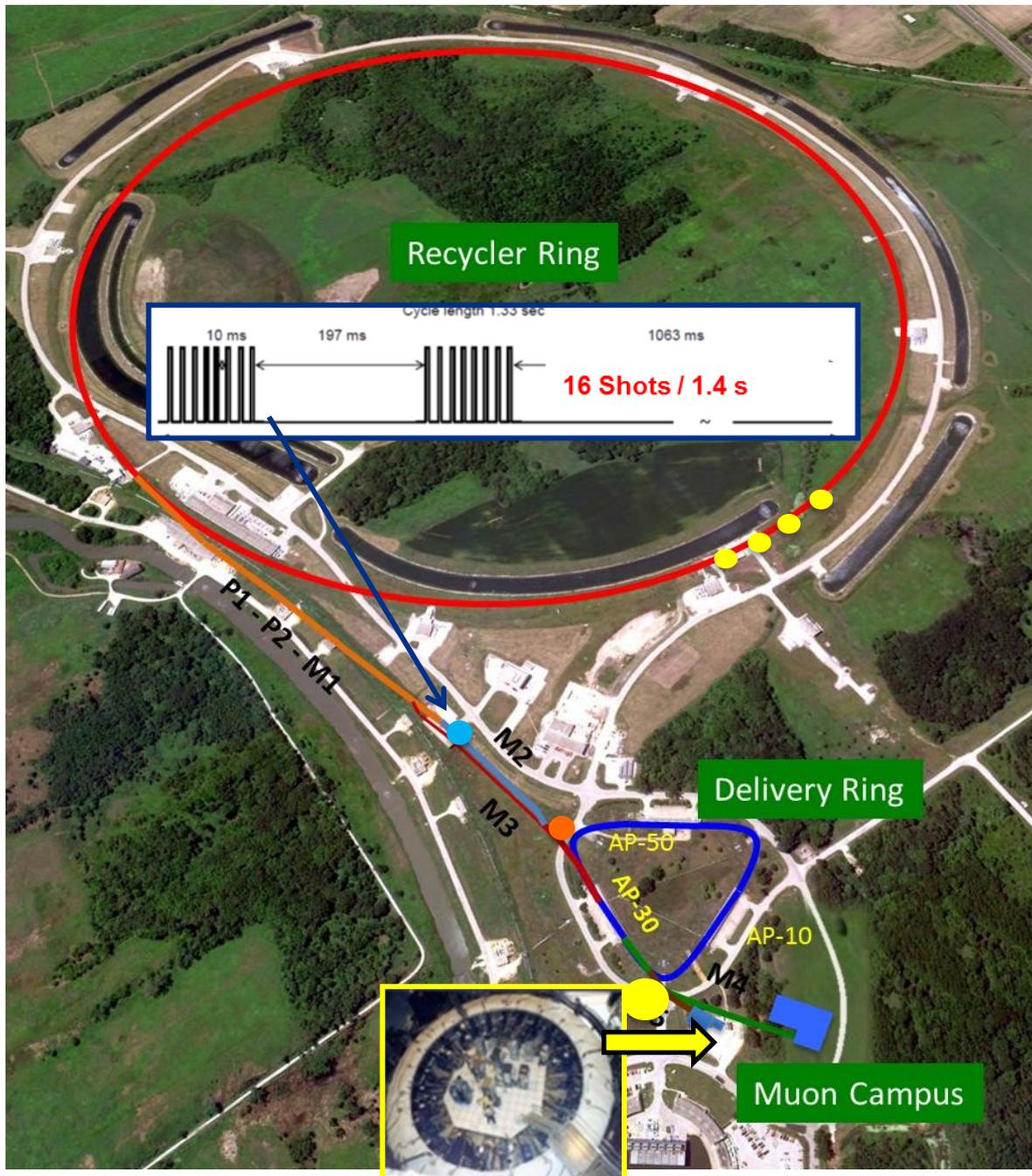


26 July 2013:...the ring arrives to FNAL





Creating the Muon Beam for g-2



- 8 GeV p batch into Recycler
- Split into 4 bunches
- Extract 1 by 1 to strike target
- Long FODO channel to collect $\pi \rightarrow \mu\nu$
- $\rho/\pi/\mu$ beam enters DR; protons kicked out; π decay away
- μ enter storage ring

- APRIL 2017
- RING
- FIELD
- PRECESSION

muons

Inflector

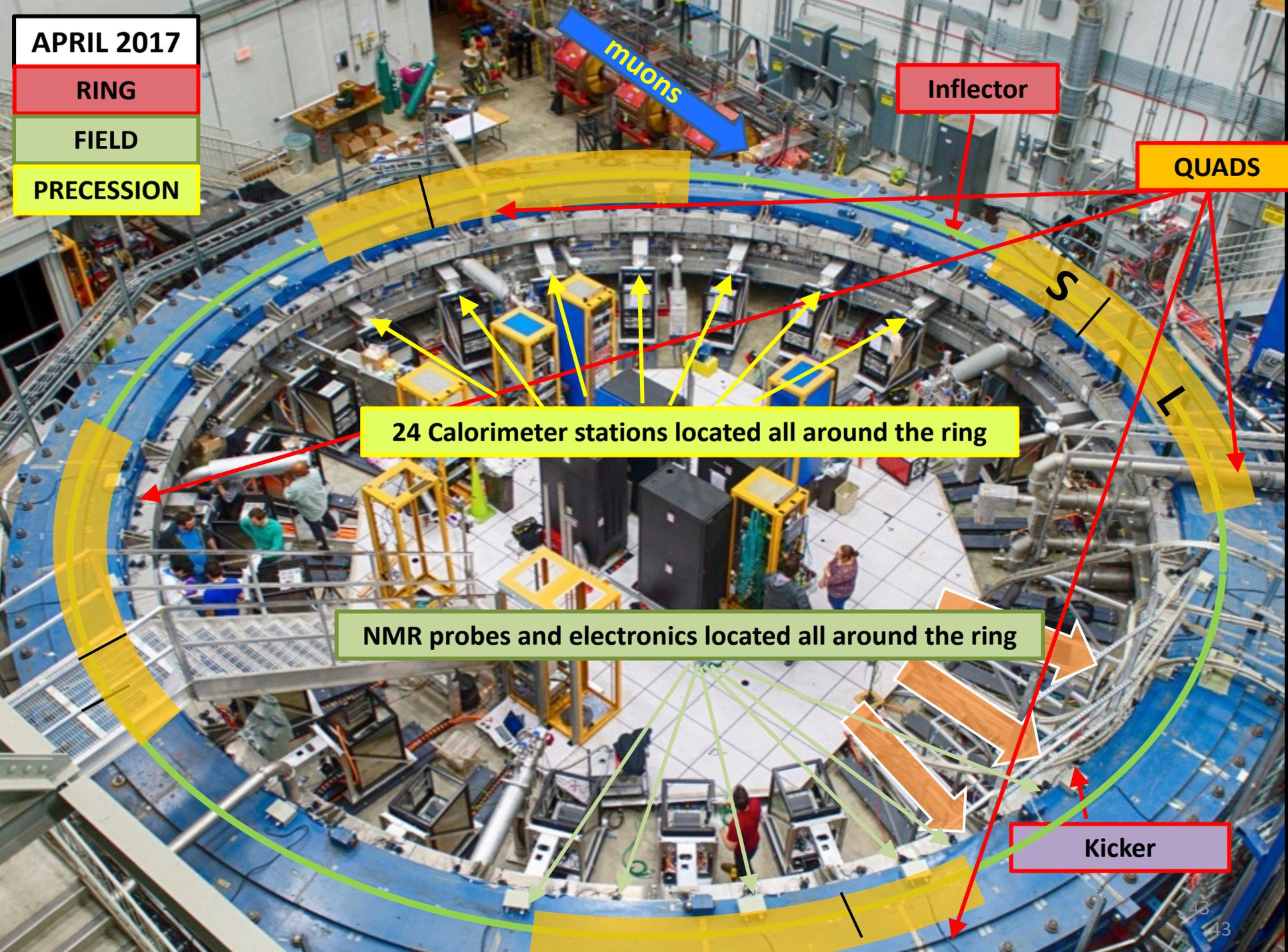
QUADS

S
L

24 Calorimeter stations located all around the ring

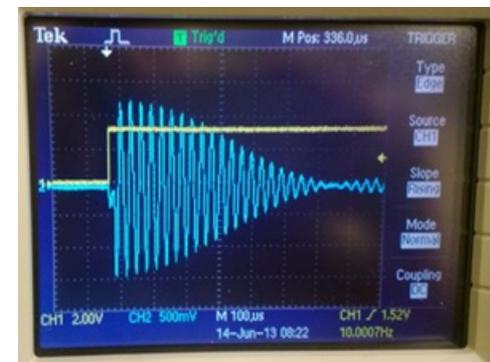
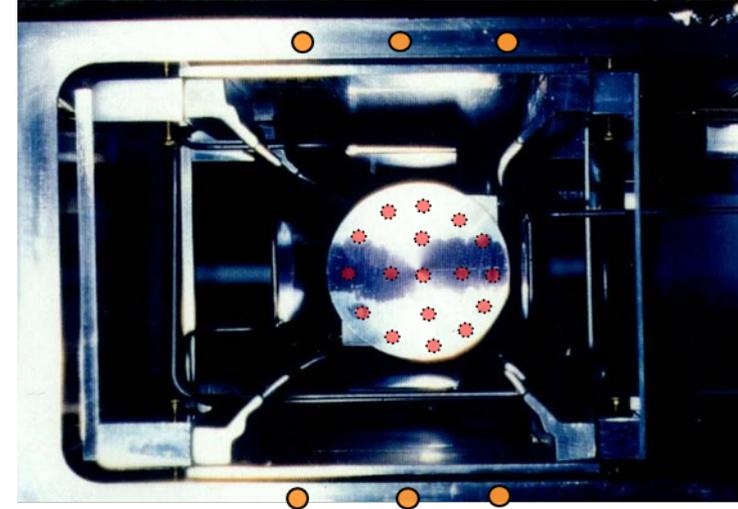
NMR probes and electronics located all around the ring

Kicker



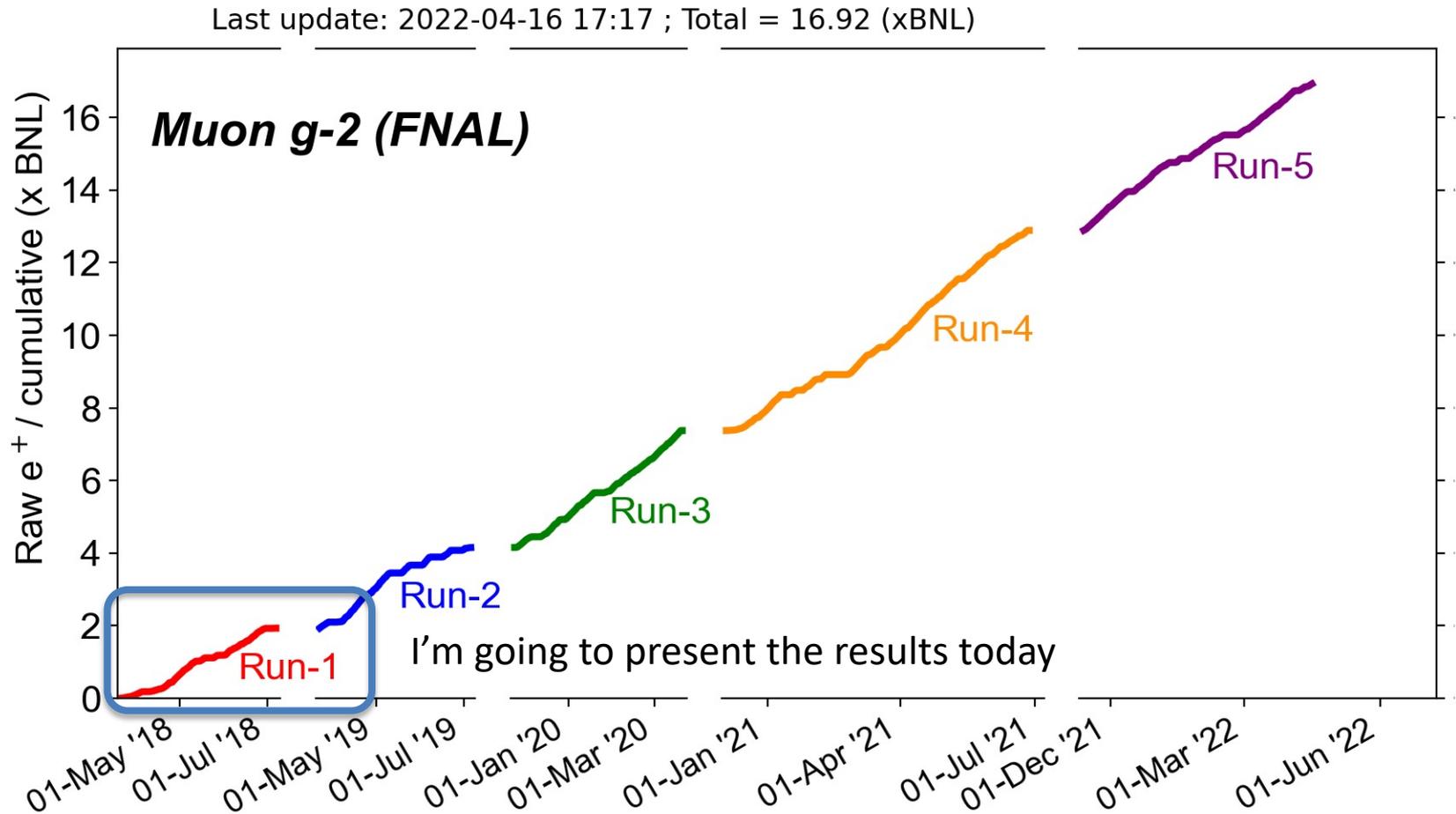
Measuring the magnetic field

- 378 Fixed probes monitor field 24/7
- 17-probe NMR trolley maps the magnetic field over the muon storage region
 - Trolley runs every 2-3 days
- Free induction decay signal of the probes digitized and analyzed to extract a precession frequency



E989 collected data

We have collected ~16 x BNL over the last 4 years:



A blinded analysis

- The analysis is twofold blinded:
 - Clock frequency blinding (HW)
 - Unknow offset in the analysis of ω_a (Software)
- The HW blinding factor is known only to two people outside the collaboration and revealed at the completion of the analysis



blinding the clock in 2018

Locked Clock Panel



RUN1: analysis structure

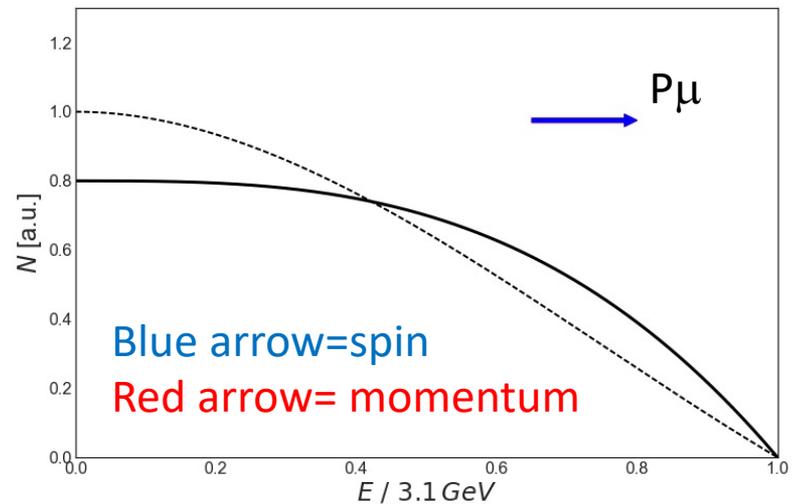
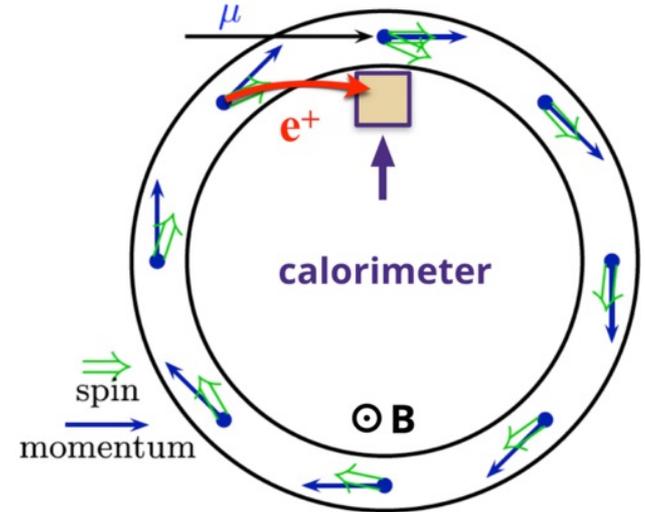
- Multiple analysis groups with different methodologies:
 - Six groups analyse ω_a with 2 different energy and time reconstructions and 4 different analysis methods
 - Two groups for the analysis of ω_p + one group for calibration
 - Different groups for beam dynamics corrections

ω_a Measurement

- The number of positrons is modulated by the anomalous precession frequency

$$N_0 e^{-t/\tau} [1 - A \cos(\omega_a t + \phi)]$$

- 4 different analysis methods:
 - T: simple energy threshold >1.7 GeV
 - A: asymmetry weighted with threshold >1.1 GeV
 - R: ratio method
 - Q: No clustering: total energy above minimal threshold
- A-method used to provide ω_a



E and t are the measured observables.

The fit equation

$$N_0 e^{-\frac{t}{\tau}} (1 + A \cdot A_{BO}(t) \cos(\omega_a t + \phi \cdot \phi_{BO}(t))) \cdot N_{CBO}(t) \cdot N_{VW}(t) \cdot N_y(t) \cdot N_{2CBO}(t) \cdot J(t)$$

$$A_{BO}(t) = 1 + A_A \cos(\omega_{CBO}(t) + \phi_A) e^{-\frac{t}{\tau_{CBO}}}$$

$$\phi_{BO}(t) = 1 + A_\phi \cos(\omega_{CBO}(t) + \phi_\phi) e^{-\frac{t}{\tau_{CBO}}}$$

$$N_{CBO}(t) = 1 + A_{CBO} \cos(\omega_{CBO}(t) + \phi_{CBO}) e^{-\frac{t}{\tau_{CBO}}}$$

$$N_{2CBO}(t) = 1 + A_{2CBO} \cos(2\omega_{CBO}(t) + \phi_{2CBO}) e^{-\frac{t}{2\tau_{CBO}}}$$

$$N_{VW}(t) = 1 + A_{VW} \cos(\omega_{VW}(t)t + \phi_{VW}) e^{-\frac{t}{\tau_{VW}}}$$

$$N_y(t) = 1 + A_y \cos(\omega_y(t)t + \phi_y) e^{-\frac{t}{\tau_y}}$$

$$J(t) = 1 - k_{LM} \int_{t_0}^t \Lambda(t) dt \quad \text{Muon Loss term}$$

$$\omega_{CBO}(t) = \omega_0 t + A e^{-\frac{t}{\tau_A}} + B e^{-\frac{t}{\tau_B}}$$

$$\omega_y(t) = F \omega_{CBO}(t) \sqrt{2\omega_c / F \omega_{CBO}(t) - 1}$$

$$\omega_{VW}(t) = \omega_c - 2\omega_y(t)$$

ω_y, ω_{VW} vertical oscillations

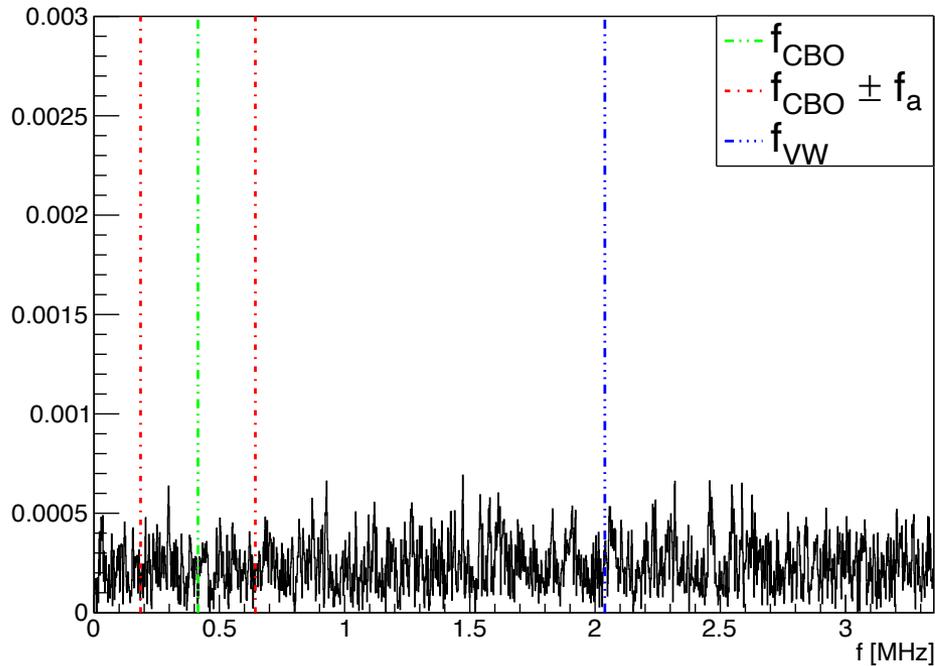
$\omega_{CBO}, \omega_{2CBO}$, radial oscillation

Red = free parameters
Blue = fixed parameters

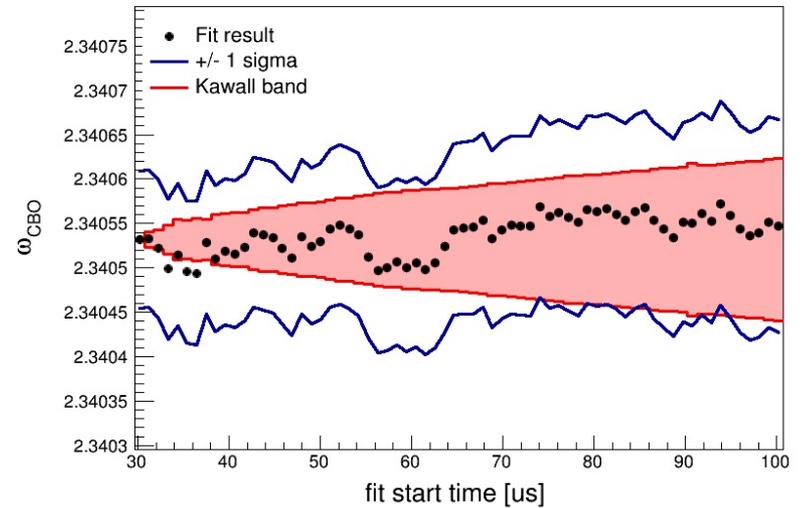
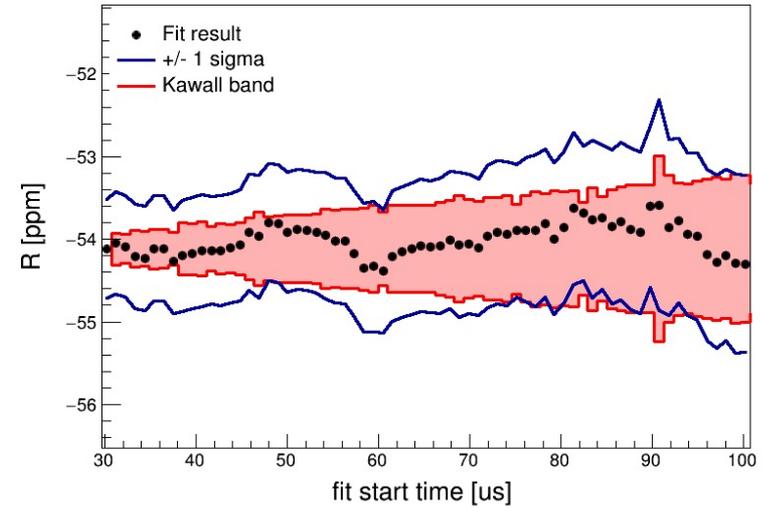
Final Fit

$$R_{(\text{blinded})} = (1 + \omega_{\text{blind}} / \omega_{\text{ref}}) [\text{ppm}]$$

Fourier transform of residuals

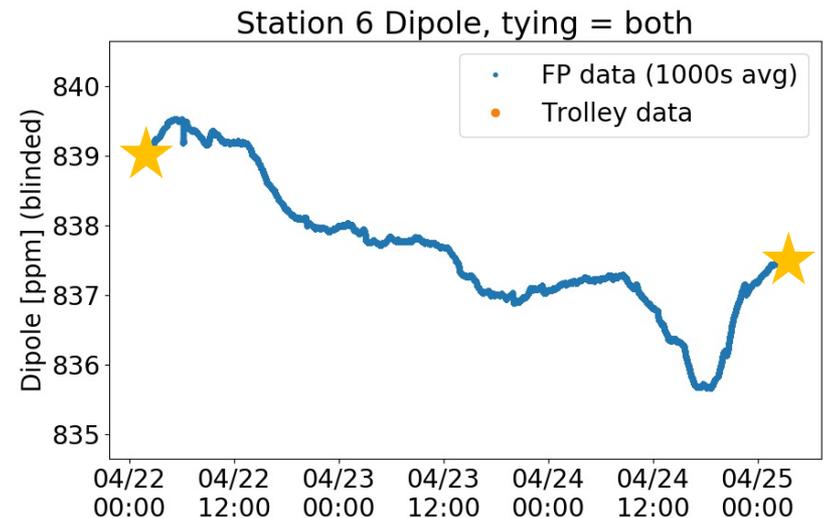
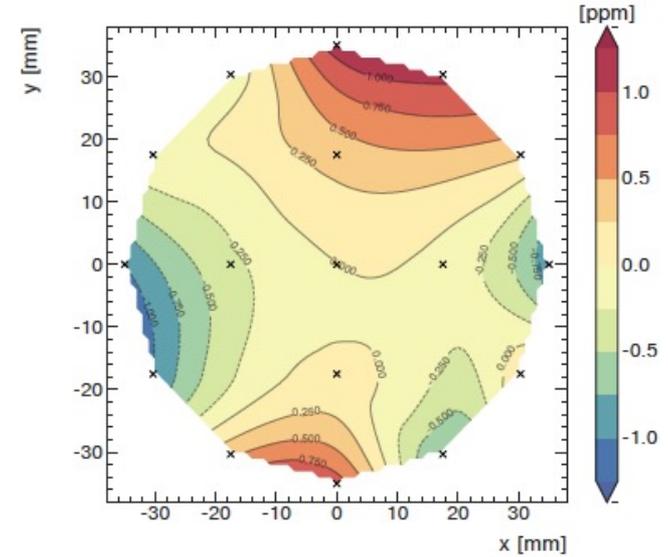


No unaccounted frequencies



Measuring the magnetic field

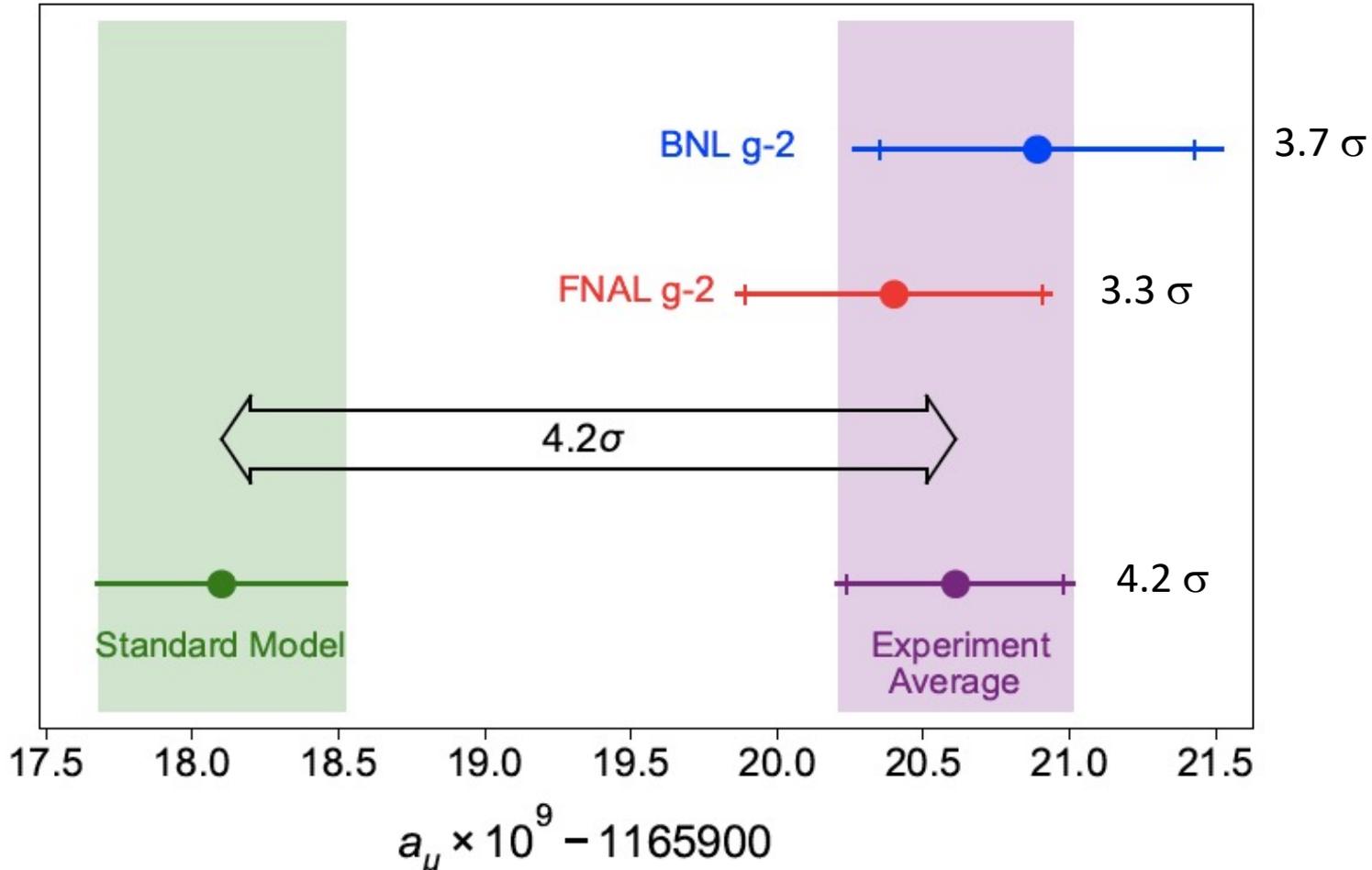
- 378 Fixed probes monitor field 24/7
- 17-probe NMR trolley maps the magnetic field over the muon storage region
 - Trolley runs every 2-3 days
- Free induction decay signal of the probes digitized and analyzed to extract a precession frequency



Result

$$a_\mu(\text{Exp}) = 116\,592\,061(41) \times 10^{-11} \quad (0.35 \text{ ppm}) \quad \text{FNAL+BNL}$$

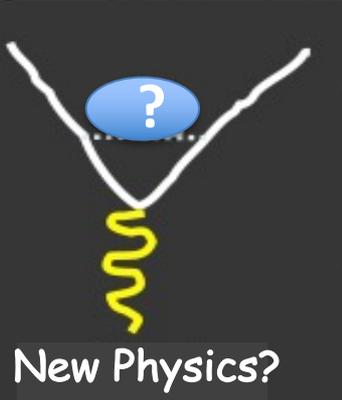
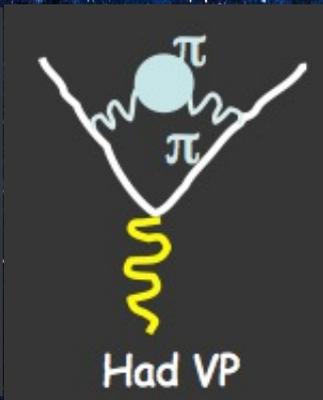
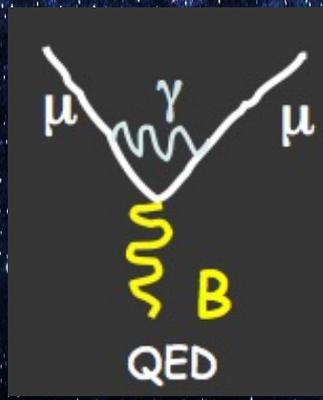
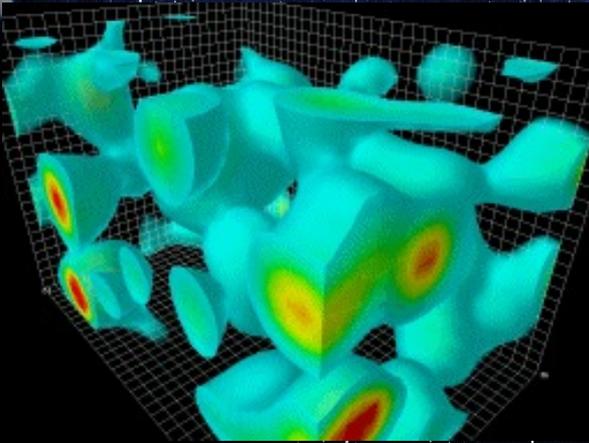
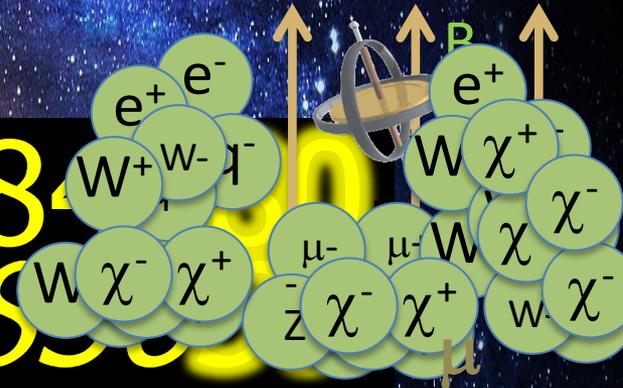
$$a_\mu(\text{SM}) = 116\,591\,810(43) \times 10^{-11} \quad (0.37 \text{ ppm}) \quad \text{Th. Initiative WP20}$$



$$a_\mu(\text{Exp}) - a_\mu(\text{SM}) = (251 \pm 59) \times 10^{-11}$$

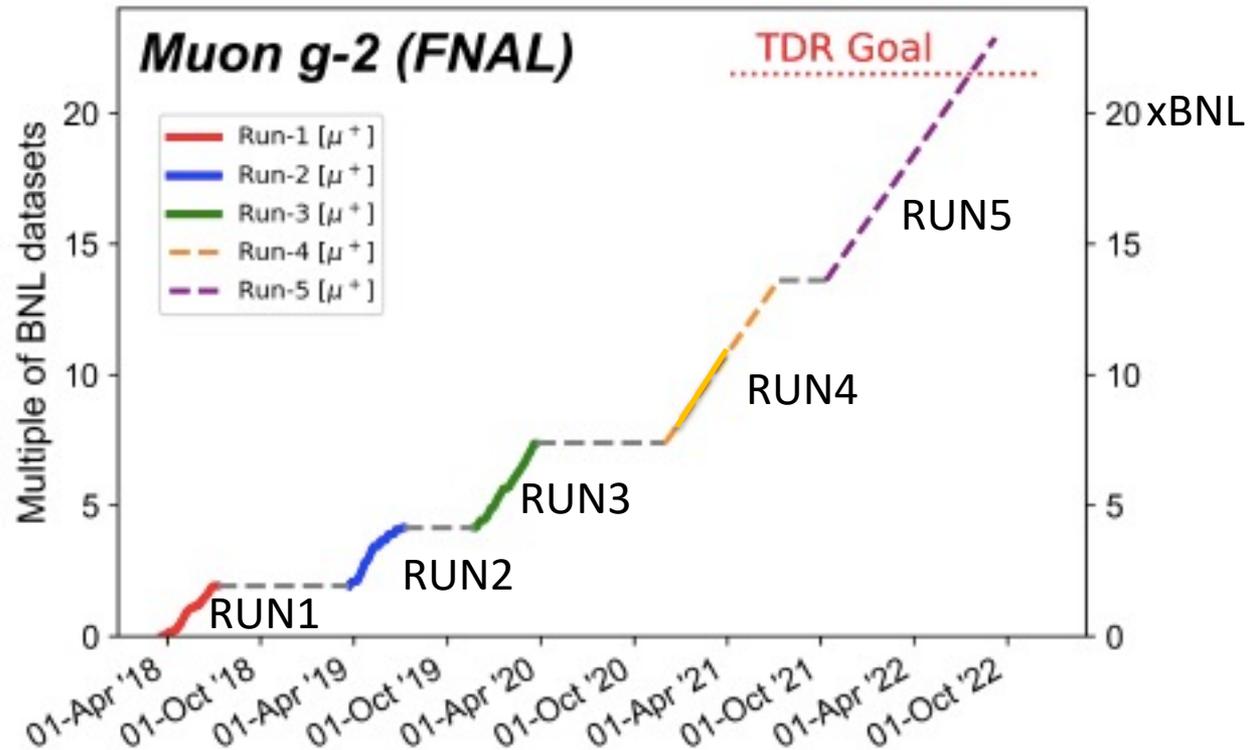
Are we seeing something new ?

$g(\text{expt})$ 2.0023318420
 $g(\text{theory})$ 2.002331836

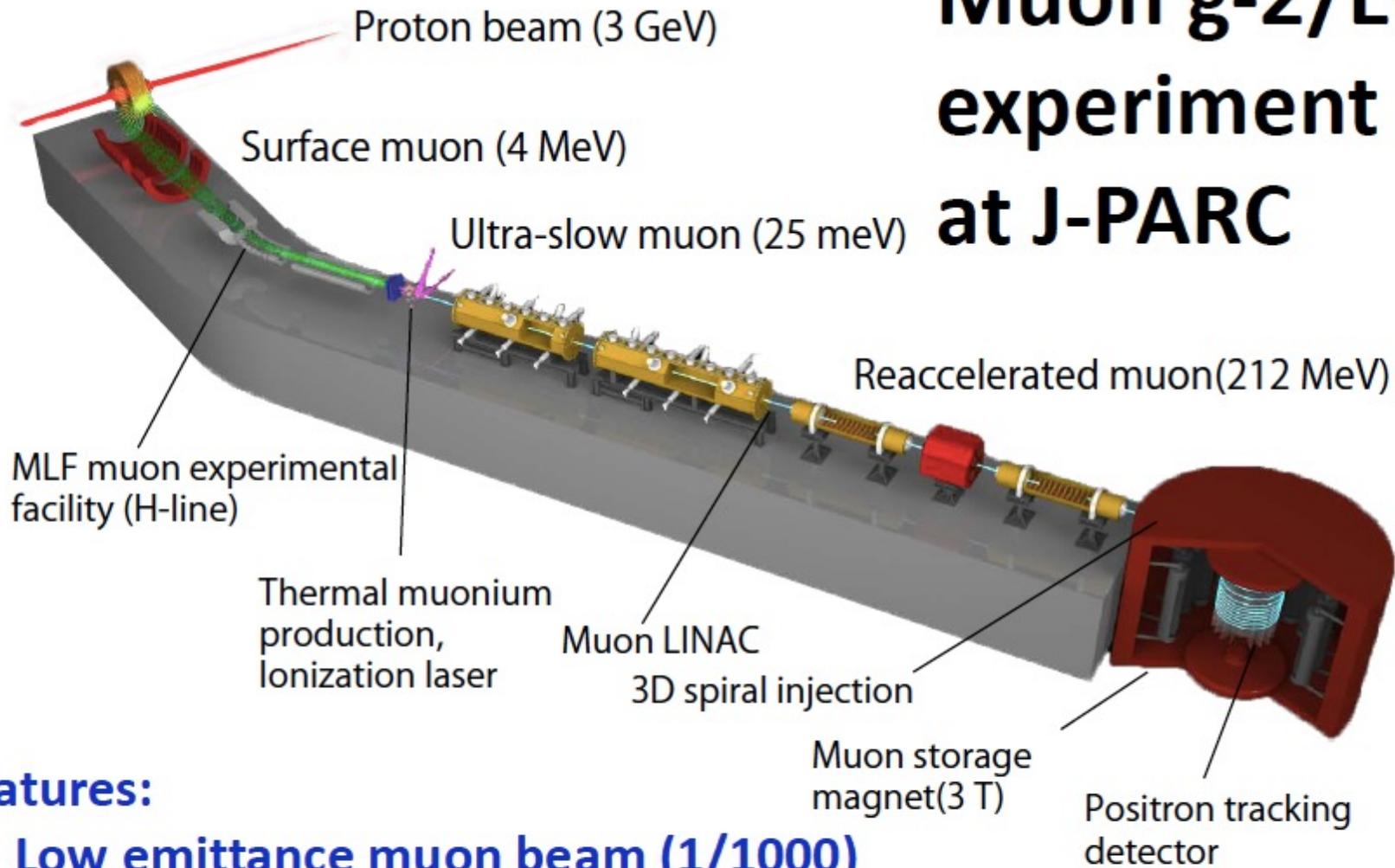


What awaits us

- RUN₁ is only 6% of the final dataset
- Analysis of RUN_{2/3} ongoing. Results expected by early 2023 (~230 ppb uncertainty).
- RUN₅ ongoing should allow to achieve the x20 BNL project goal
- RUN₆ in 2022-2023 planned with μ^-



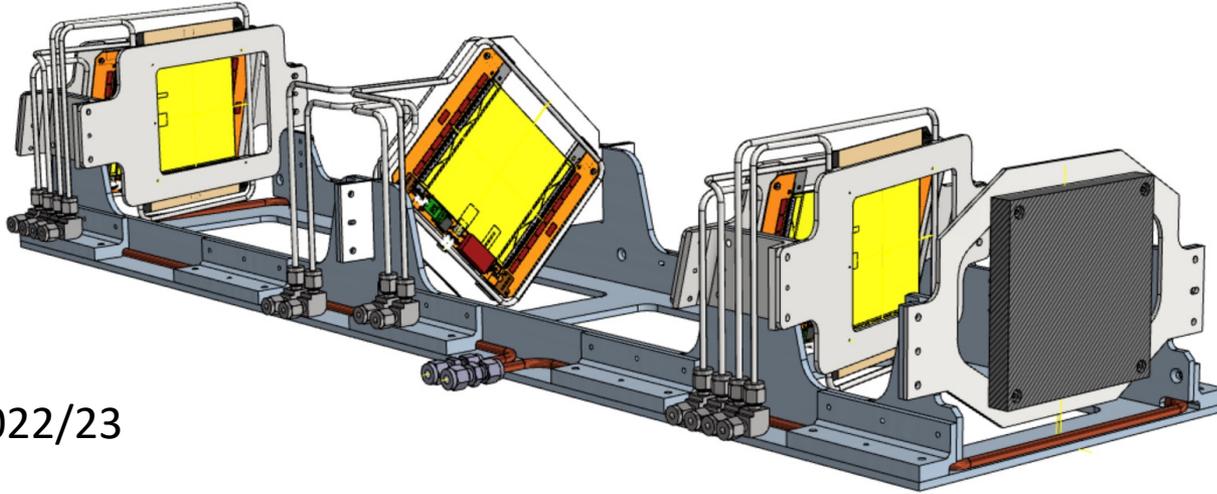
Muon g-2/EDM experiment at J-PARC



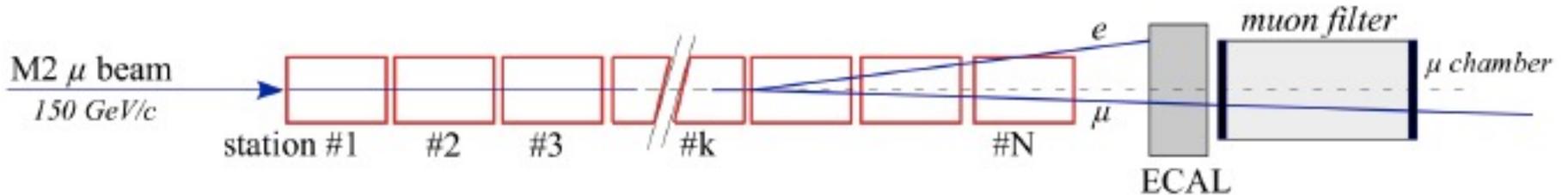
Features:

- **Low emittance muon beam (1/1000)**
- **No strong focusing (1/1000) & good injection eff. (x10)**
- **Compact storage ring (1/20)**
- **Tracking detector with large acceptance**
- **Completely different from BNL/FNAL method**

A third way for HVP...MUonE at CERN



Test RUN 2022/23



Alternative measurement of HVP for a_μ

-C. M. Carloni Calame et al *PLB* 746 (2015) 325

-G. Abbiendi et al *Eur.Phys.J.C* 77 (2017) 3, 139

-Lol <https://cds.cern.ch/record/2677471/files/SPSC-I-252.pdf>

Conclusions (with Paolo's words)

Unsatisfactory points:

Paolo's presentation in 2001

(4.2 σ (e+e-))

1. ~~2.6~~ σ is not very compelling and is also author dependent.
2. M-C is $\sim 3^{\sim 2} \times$ EW contribution. What about LEP, $b \rightarrow s\gamma$, M_W , M_{top} , $\Re(\epsilon'/\epsilon)$, $\sin 2\beta$...
3. Hadronic corrections difficult, e.g. light-by-light (also HVP!)
4. SUSY as a theory is not very precise at the moment. It has too many unknown, free parameters. There is no exp. evidence for it nor a prediction follows from the possible effect in the muon anomaly.

Soon better statistics and both signs muons.

Still very exciting at present.

2nd KLOE Physics Workshop
Otranto, 10th-12th June 2002



Thanks!

The Muon g-2 Collaboration (Elba 2019)



ONE THING IS FOR SURE: THE HUNT IS ON, AND
NEW DISCOVERIES ARE ON THE HORIZON.

Muon g-2 @ J-Park

Muon g-2 @ Fermilab

Theory Initiative

MUonE @ CERN

Stay tuned!



a_μ : Unblinding

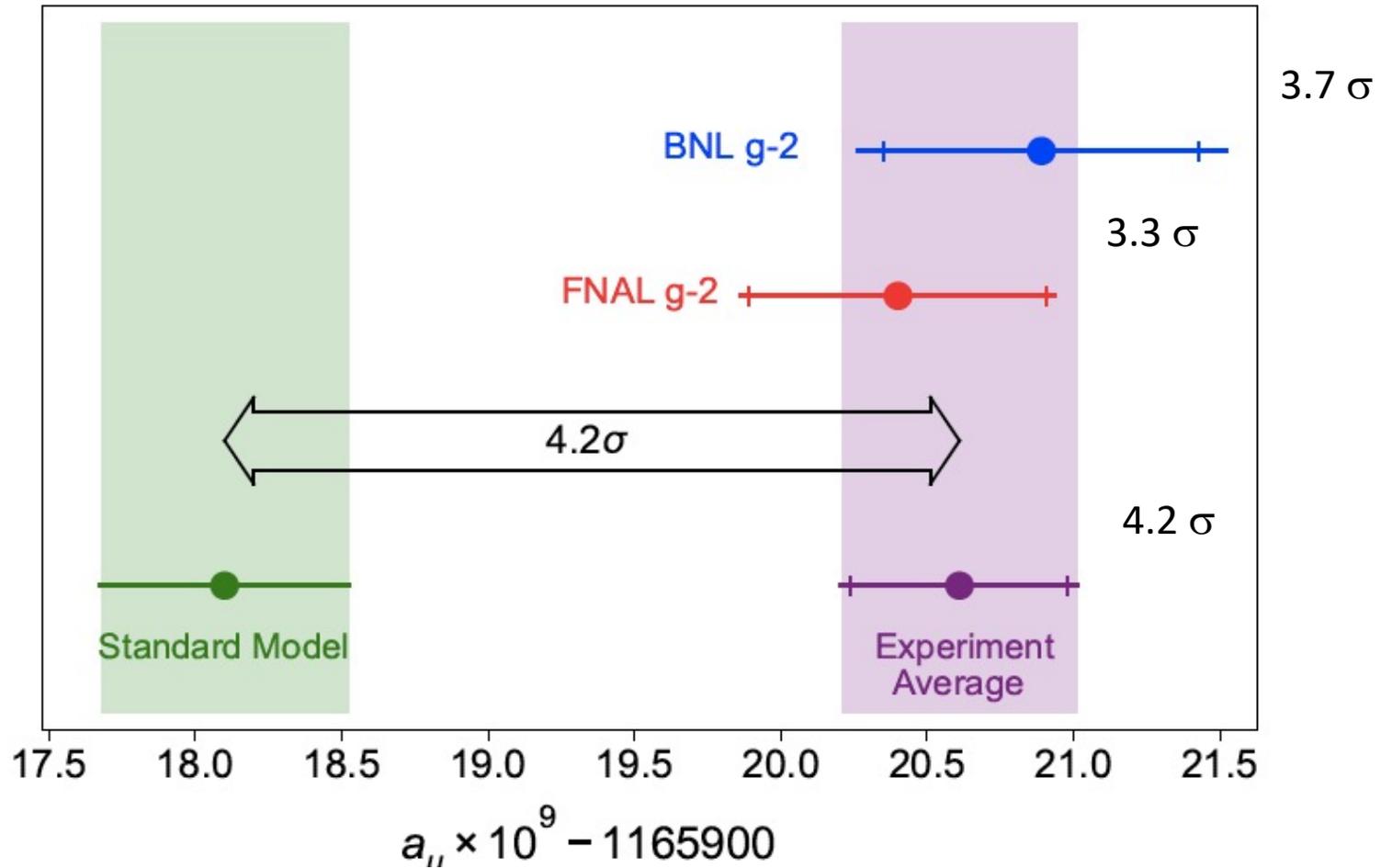
Quantity	Correction Terms (ppb)	Uncertainty (ppb)
ω_a (statistical)	–	434
ω_a (systematic)	–	56
C_e	489	53
C_p	180	13
C_{ml}	-11	5
C_{pa}	-158	75
$f_{calib} \langle \omega'_p(x, y, \phi) \times M(x, y, \phi) \rangle$	–	56
B_q	-17	92
B_k	-27	37
$\mu'_p(34.7^\circ)/\mu_e$	–	10
m_μ/m_e	–	22
$g_e/2$	–	0
Total	–	462

434 ppb stat \oplus 157 ppb syst error

$$a_\mu(\text{FNAL}) = 116\,592\,040(54) \times 10^{-11} \quad (0.46 \text{ ppm})$$

Result

1:40000 chance that the SM is correct!

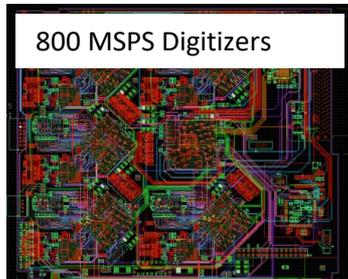
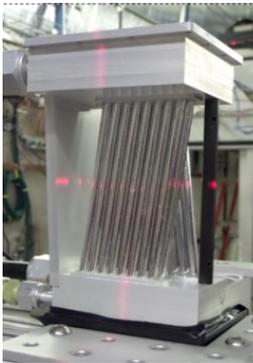


$$a_{\mu}(\text{FNAL}) = 116\,592\,040(54) \times 10^{-11} \quad (0.46 \text{ ppm})$$

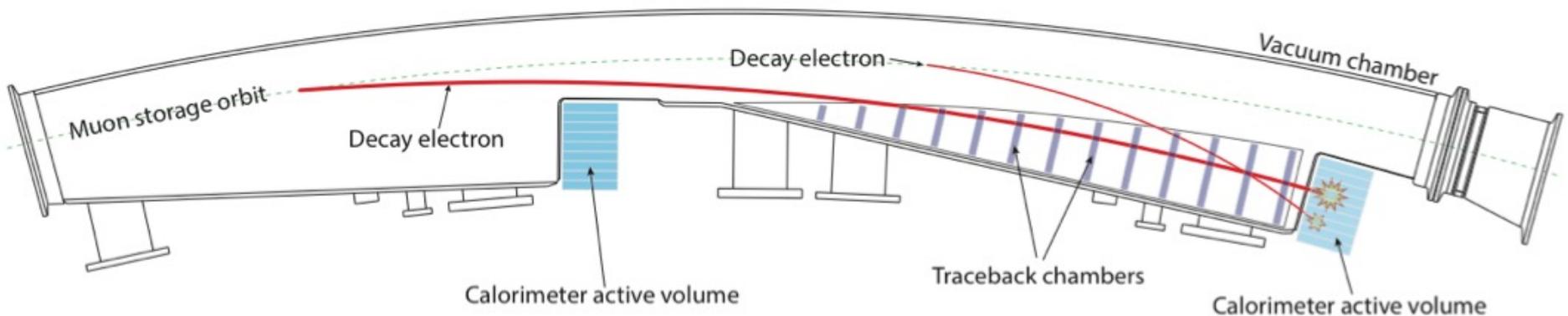
Detector systems



- Calorimeters: fast PbF_2 crystal arrays with SiPM readout \rightarrow greatly reduce pileup
- State of the art laser calibration system
- WFD electronics \rightarrow greatly reduced energy threshold
- Two straw tube trackers to precisely monitor properties of stored muons



Top view of 1 of 12 vacuum chambers



Extracting a_μ (more realistic)

Corrections due to beam dynamics

$$R'_\mu = \left(\frac{f_{clock} \cdot \omega_a^{meas} \cdot (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{calib} \cdot \omega'_p(x, y, \phi) \otimes M(x, y, \phi) \cdot (1 + B_k + B_q)} \right)$$

Corrections due to transient magnetic fields

$$a_\mu = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} \frac{\mu'_p(T_r)}{\mu_e(H)} \frac{\mu_e(H)}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$

What we measure

External inputs (known at 24 ppb)