

Giornata nazionale EIC_NET 2022

Transversity 2022 highlights and the Physics of SIDIS

Marco Radici



Transversity 2022 in numbers



65 registered
49 in person, **16** remote

57 talks, **8** from remote
16 short talks

+

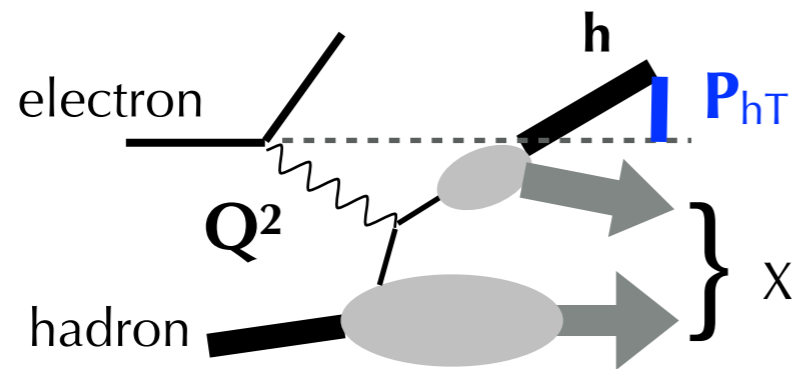
Piet Mulders' fest for
his 70th birthday



Transversity 2022 outline

Mostly focus on

Semi-Inclusive Deep-Inelastic Scattering (SIDIS)

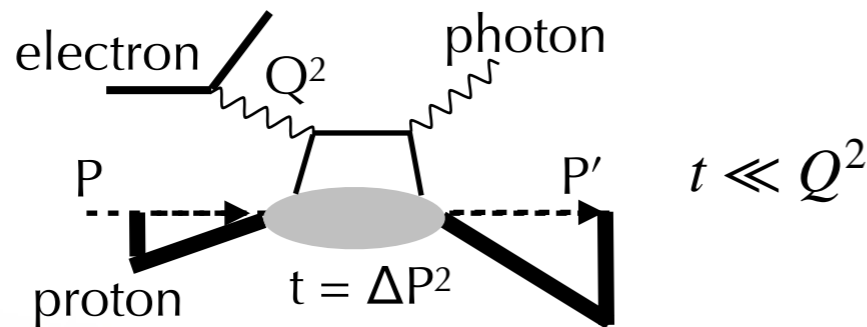


with various final states "h" :
light- / heavy-quark hadrons
jets, hadron-in-jet, etc..

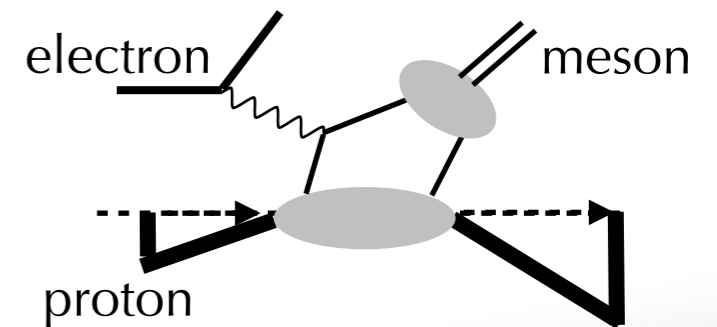
Factorization th.'s available (not everywhere!) for $P_{hT}^2/z^2 \ll Q^2$

but also exclusive processes...

DVCS



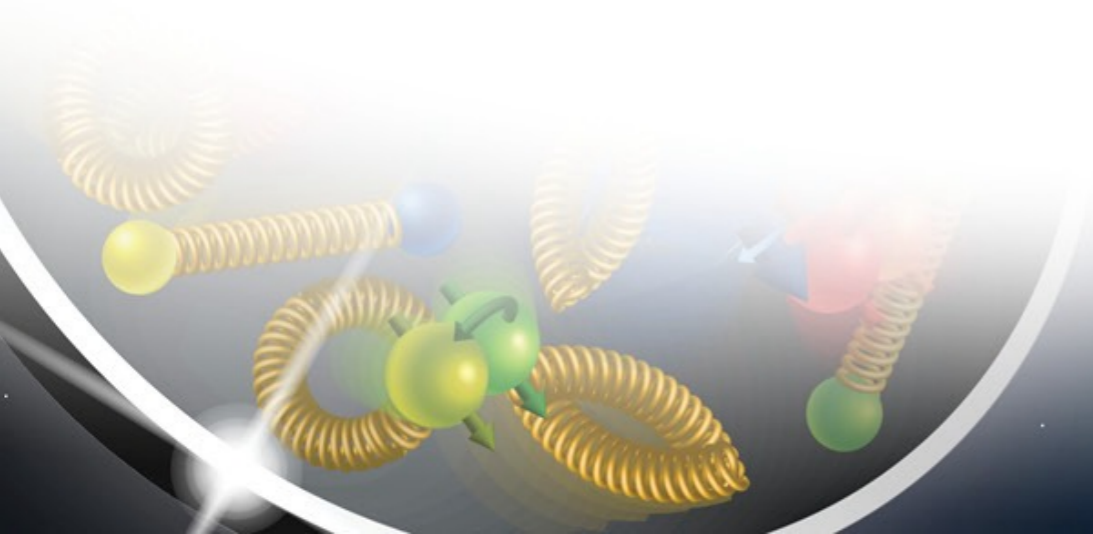
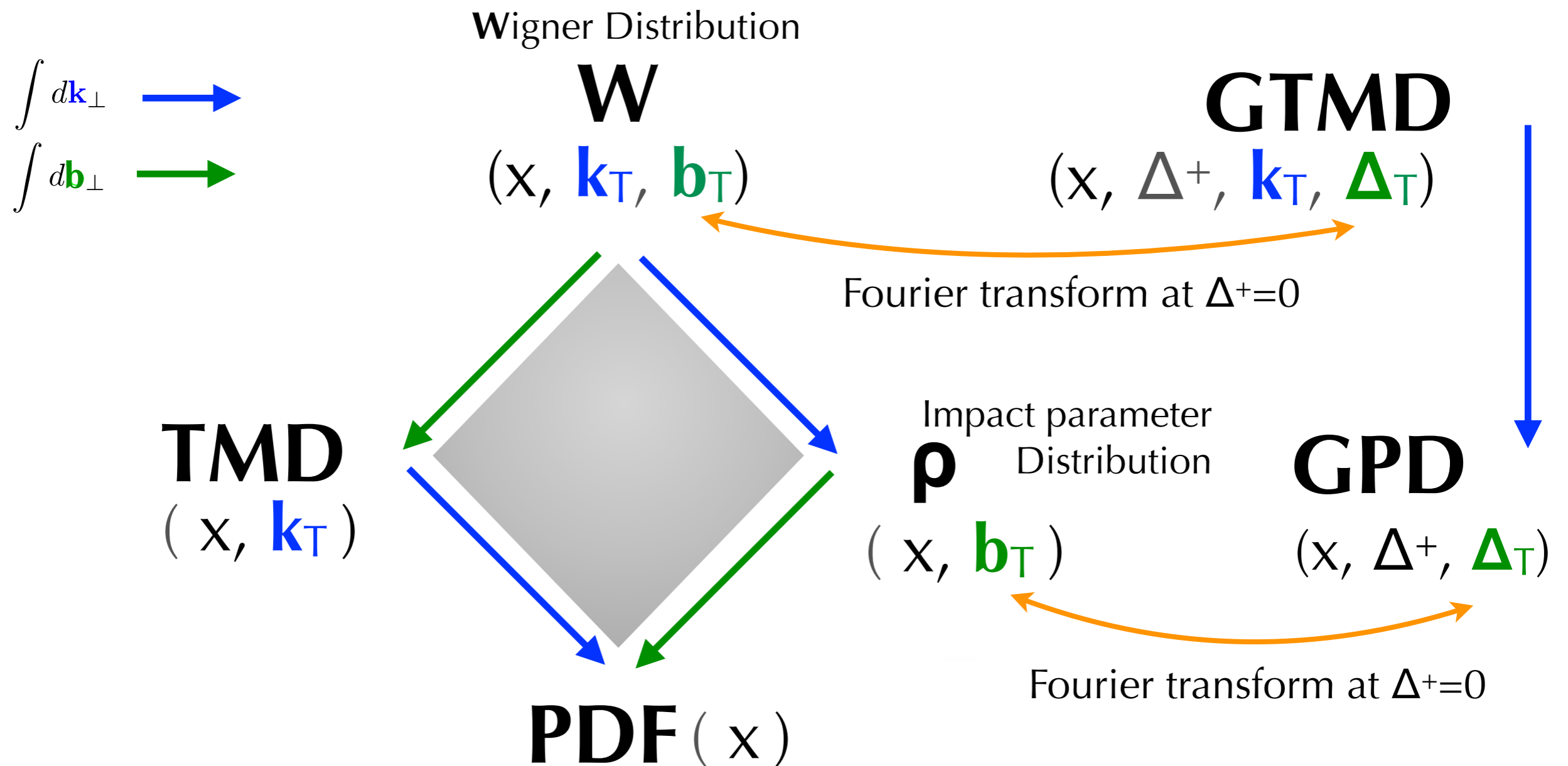
DVMP



and specific diffractive channels

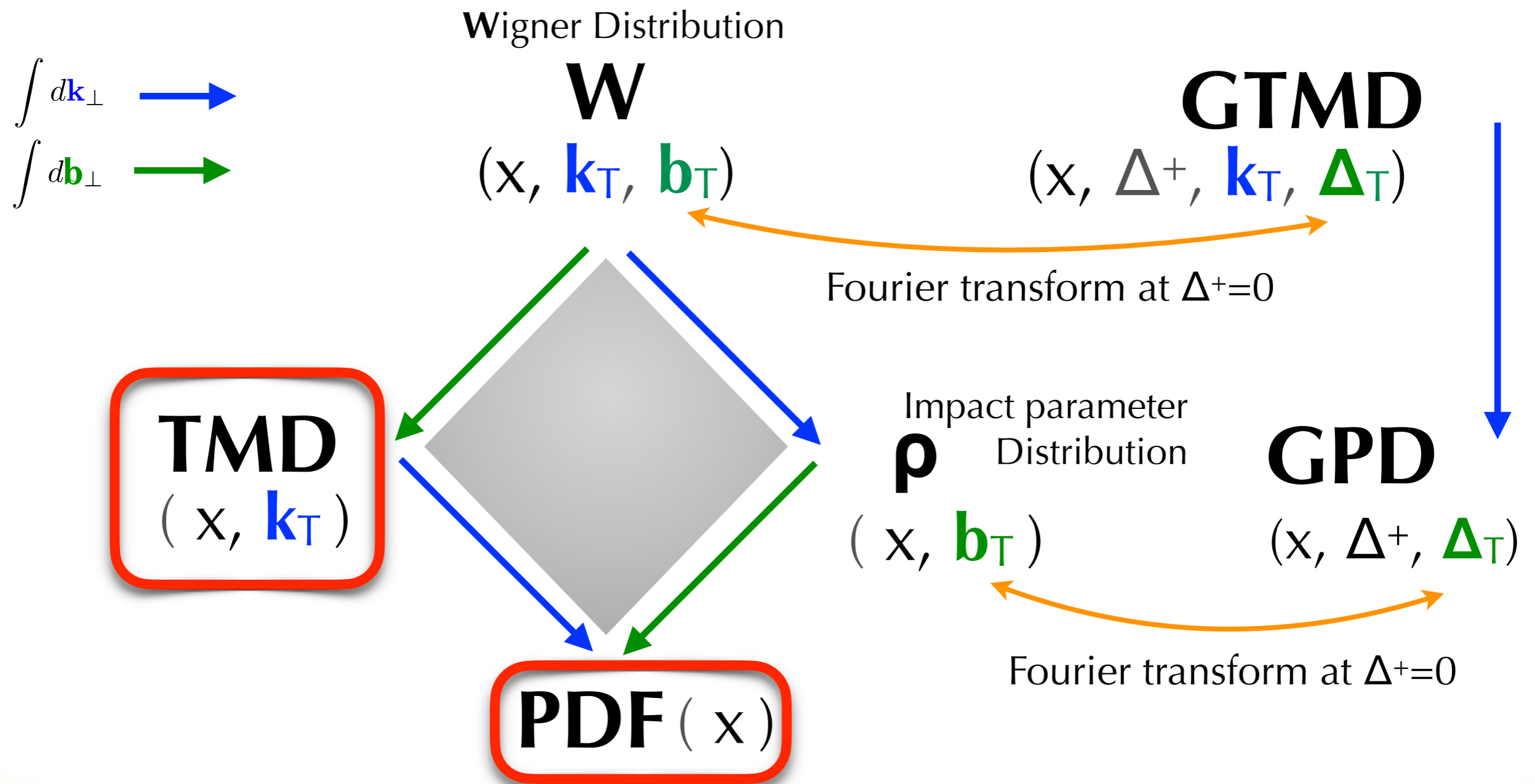
From the point of view of a theoretician...

Lorcé et al., JHEP 1105 (11) 041



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Lorcé et al., JHEP 1105 (11) 041

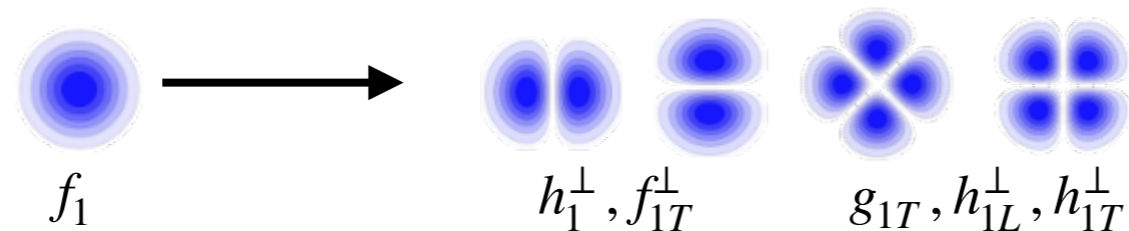


What do we know about them ?
Where do we learn more ?

The TMD “zoo” at leading twist

		Quark polarization		
		Unpolarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 = \odot$	\times	$h_1^\perp = \odot \uparrow - \odot \downarrow$
	L	\times	$g_1 = \odot \rightarrow - \odot \leftarrow$	$h_{1L}^\perp = \odot \nearrow - \odot \nwarrow$
	T	$f_{1T}^\perp = \odot \uparrow - \odot \downarrow$	$g_{1T} = \odot \rightarrow - \odot \leftarrow$	$h_1 = \odot \uparrow - \odot \downarrow$ $h_{1T}^\perp = \odot \nearrow - \odot \nwarrow$

deformations induced by spin-momentum correlations



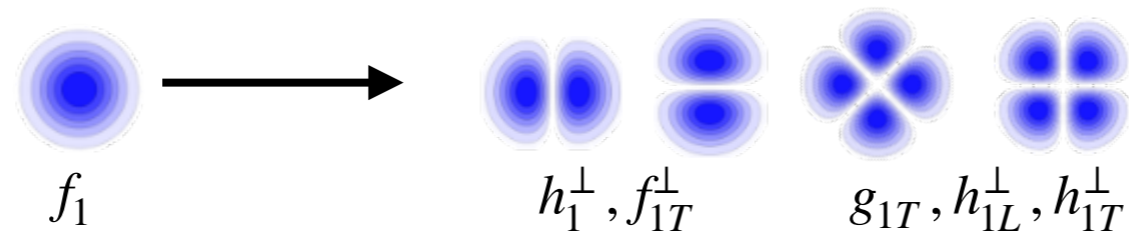
each TMD is connected to a specific measurable SIDIS spin asymmetry

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exist as collinear PDFs

deformations induced by spin-momentum correlations



each TMD is connected to a specific measurable SIDIS spin asymmetry

Transversity

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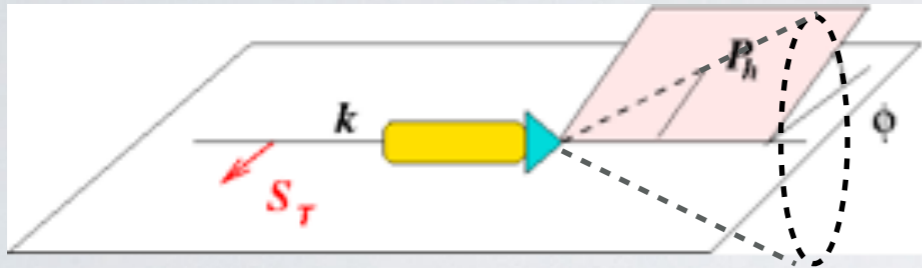
transversity

- chiral-odd structure **also in collinear kin.**
- only way to determine the tensor charge $\delta^q(Q^2) = \int_0^1 dx h_1^{q-\bar{q}}(x, Q^2)$
- no chiral-odd structures in SM Lagrangian; potential doorway to BSM

Example: in SMEFT's, neutron EDM d_n is source of strong CP violation

bounds from exp. $\longrightarrow d_n = \delta u d_u + \delta d d_d + \delta s d_s$ tensor charges

Mechanisms for transversity



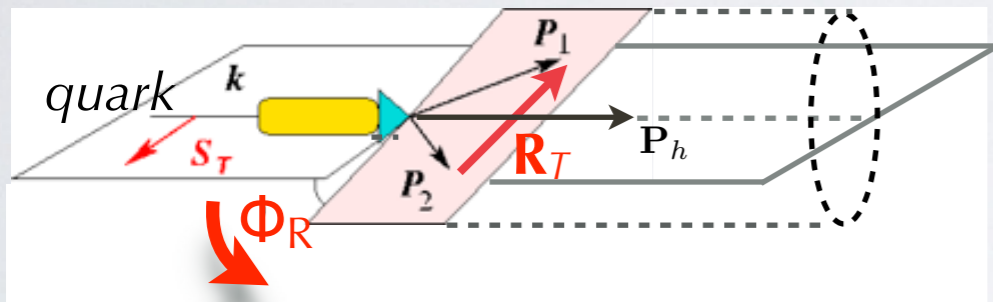
Collins effect

$$\mathbf{S}_T \cdot \mathbf{k} \times \mathbf{P}_{hT}$$

Collins, N.P. **B396** (93) 161

$$h_1(x, k_\perp) \otimes H_1^\perp(z, P_\perp)$$

transversity as TMD



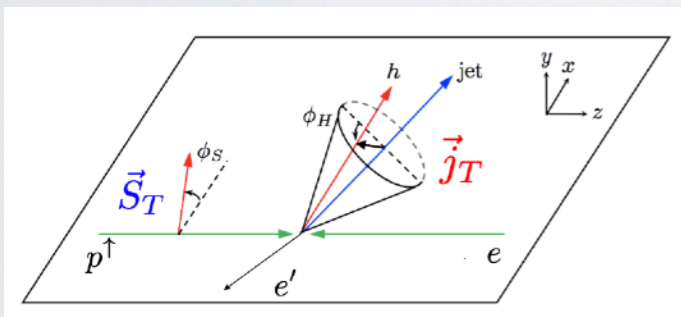
di-hadron mechanism

$$\mathbf{S}_T \cdot \mathbf{P}_2 \times \mathbf{P}_1 = \mathbf{S}_T \cdot \mathbf{P}_h \times \mathbf{R}_T$$

Collins et al., N.P. **B420** (94)

$$h_1(x) H_1^{\leftarrow}(z, R_T^2)$$

transversity as PDF

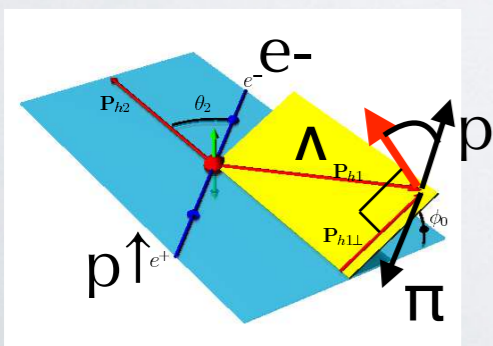


hadron-in-jet Collins effect

$j_T^2 \ll Q^2 = (P_T^{\text{jet}})^2$ hybrid factorisation:

$$h_1(x) [C(z, \mu) \otimes H_1^\perp(z_h, j_T, P_T^{\text{jet}} R)]$$

transversity as PDF



Λ spin transfer

$$h_1(x) H_1(z)$$

transversity as PDF

and also in πp^\uparrow Drell-Yan

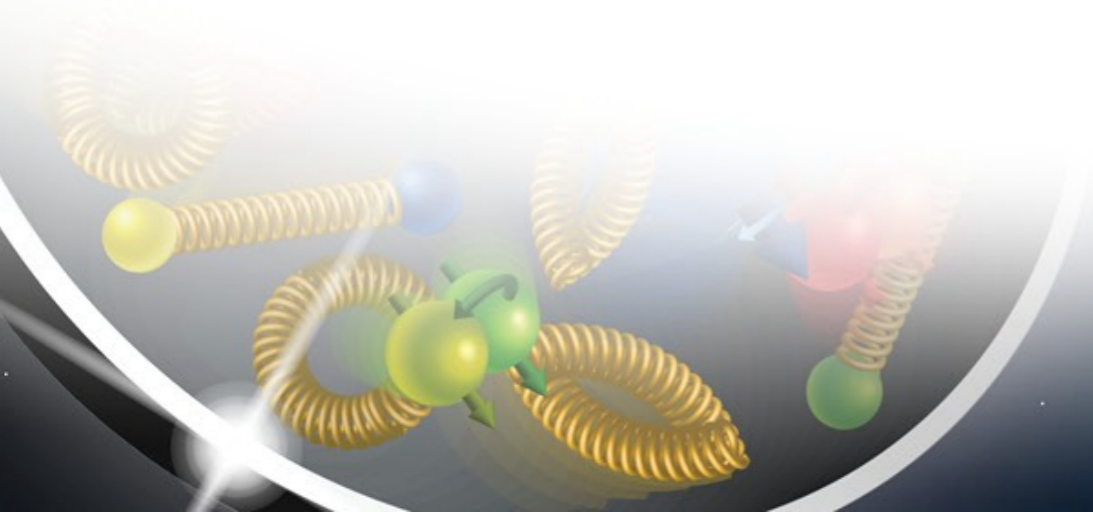
$$h_1^\perp(x_1, k_{1\perp}) \otimes h_1(x_2, k_{2\perp})$$

transversity as TMD

Phenomenology of Transversity

most recent extractions

	Mechanism	Framework	SIDIS	e+e-	p-p collisions	N pts
PV 2018 arXiv:1802.05212	collinear DiFF	LO	✓	✓	✓	78
JAM 2020 arXiv:2002.08384	Collins effect	generalized parton model	✓	✓	✓	517
MEX 2019 arXiv:1912.03289	collinear DiFF	LO	✓	✓	✗	68
CA 2020 arXiv:2001.01573	Collins effect	generalized parton model	✓	✓	✗	76
JAM 2022 arXiv:2205.00999	Collins effect	generalized parton model	✓	✓	✓	634



Phenomenology of Transversity

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other works with STAR data

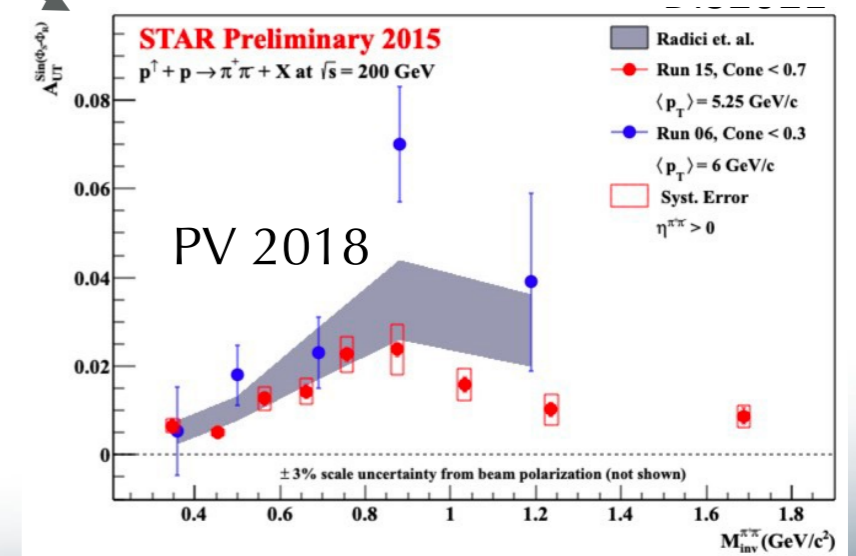
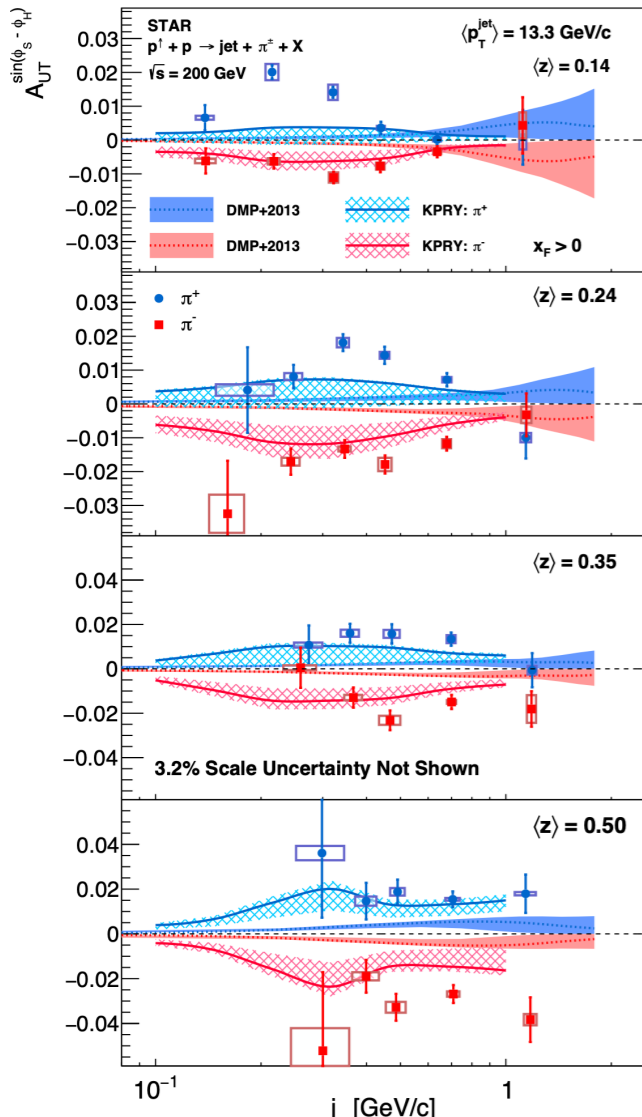
hadron-in-jet
Collins effect

di-hadron mechanism

talks by M. Grosse-Perdekamp
W.W. Jacobs

KPRY *Kang et al.,
P.L. B774 (17) 635*

DMP *D'Alesio et al.,
P.L. B773 (17) 300*



Tensor charge

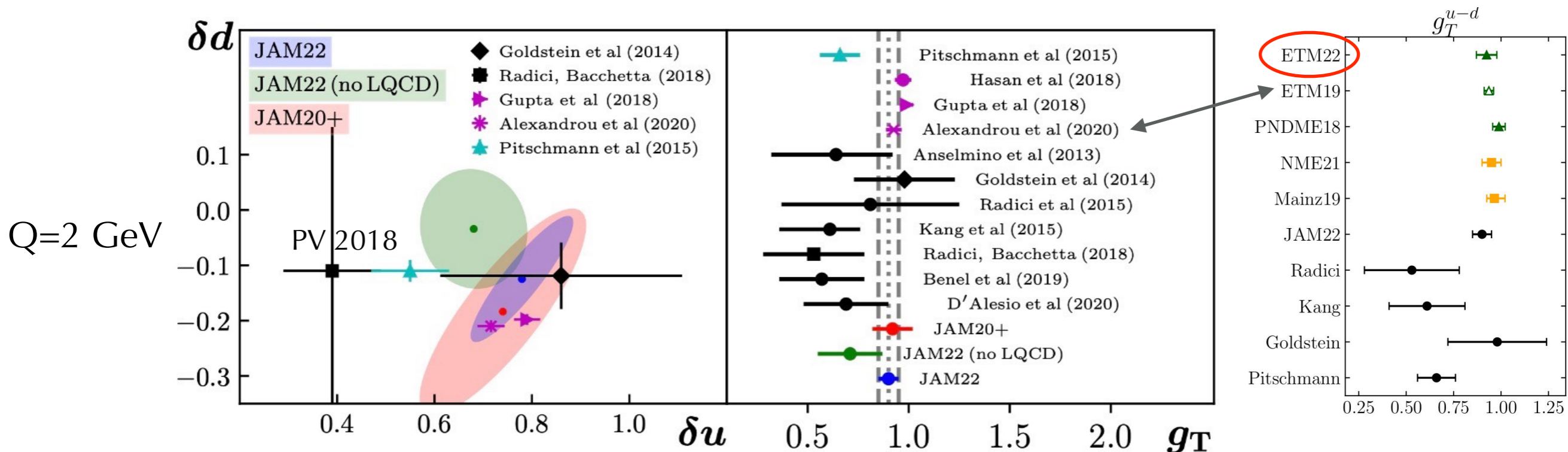
talk by D. Pitonyak

$$\delta^q(Q^2) = \int_0^1 dx h_1^{q-\bar{q}}(x, Q^2)$$

$$g_T = \delta u - \delta d$$

talk by S. Bacchio

arXiv:2202.09871

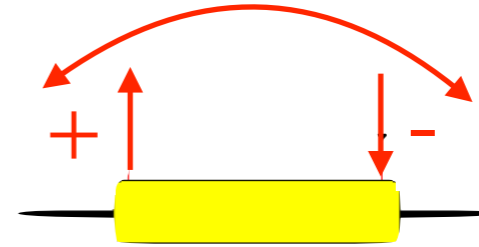


- JAM22 includes Soffer bound $\Rightarrow \delta^d$ similar to others, δ^u still larger (effect of A_N data?)
- JAM22 includes lattice g_T results in the fit \Rightarrow statistically compatible by construction
- JAM22 and PV 2018 do not \Rightarrow tension with lattice **why??**

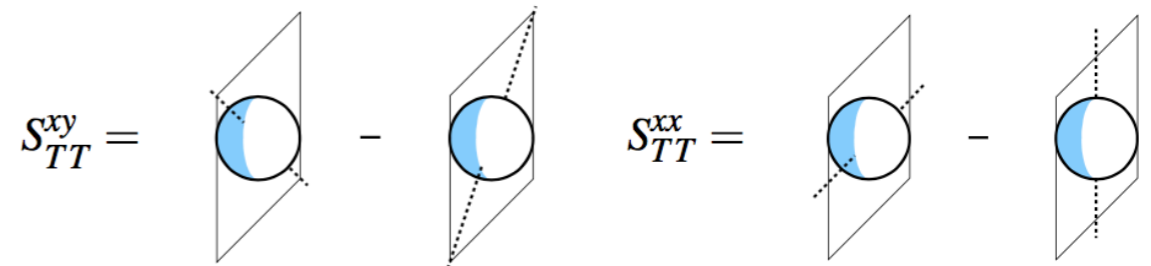
What about gluons ?

in spin-1/2 proton \rightarrow no gluon transversity

$$\max \Delta s_L = |S'_L - S_L|$$



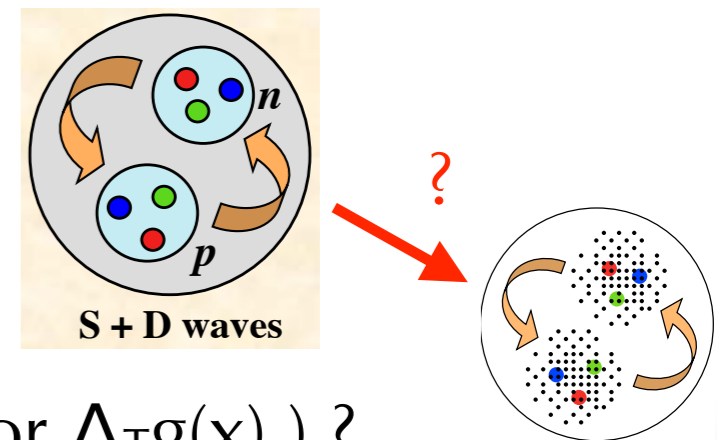
in spin-1 deuteron \rightarrow gluon "transversity"
because for transverse tensor polarization
it can be $\Delta s_L = 2$



Jaffe & Manohar (1989), Artru & Mekhfi (1990), Bacchetta & Mulders (2000)

talks by D. Boer and S. Kumano

since standard convolution model for deuteron
does not reproduce data for the tensor struct. fnct. $b_1(x)$



what is the mechanism for the gluon transversity $h_{1TT^g}(x)$ (or $\Delta_{Tg}(x)$) ?

an objective of the EIC



New and future data for transversity studies

- new 3-D analysis of Collins effect from Hermes, with final $h = \pi, K, p, pbar$

talk by G. Schnell

Airapetian et al., JHEP12 (2020) 010

- Collins effect for ρ^0 measured by Compass
- transversity induced by Λ polarization

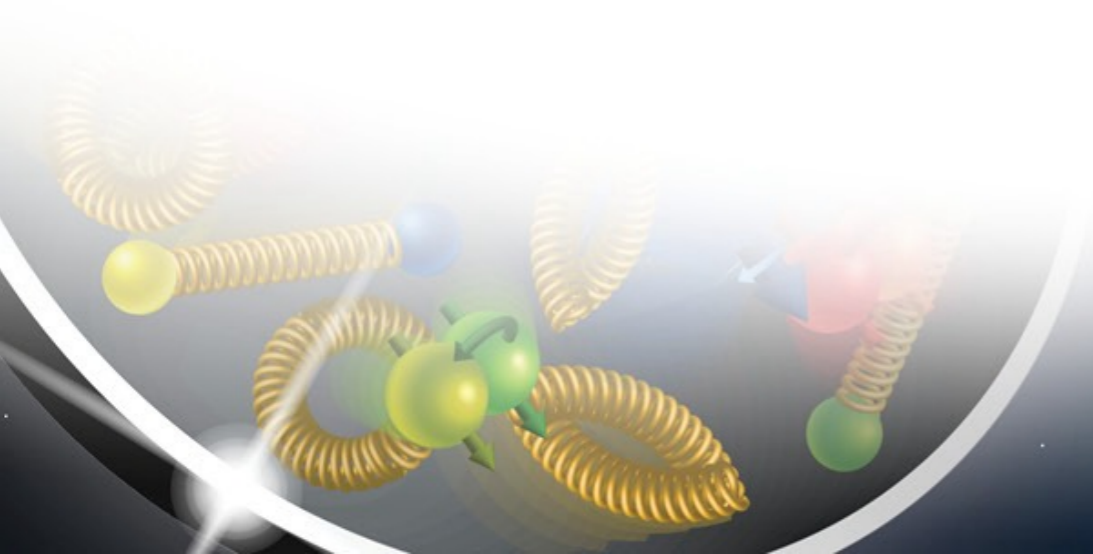
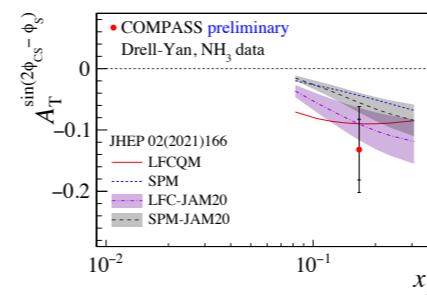
P.L. B824 (22) 136834

talks by A. Bressan

F. Bradamante

- $\pi p \uparrow$ DY by Compass: $h_{1,\pi}^\perp \otimes h_{1,p}$

talk by R. Longo



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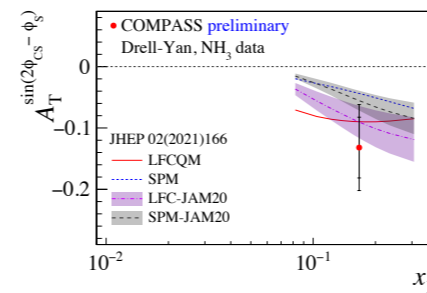
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talk by R. Longo



- Compass run with transversely polarized $^6\text{LiD} \Rightarrow$ will improve h_1^d

- JLab12 Hall-A TSSA with “neutron target” (SoLID) \Rightarrow improve $h_1^{u,d}$

- LHCspin $\Rightarrow p\text{-}p \uparrow$ DY $\Rightarrow h_{1,p}^\perp \otimes h_{1,p}$

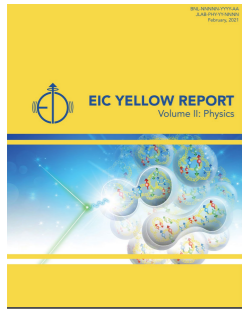
talk by P. Di Nezza

- Amber $\Rightarrow \pi, K$ DY $\Rightarrow h_{1,\pi,K}^\perp \otimes h_{1,p}$

talk by N. Wuerfel

- FermiLab “LongQuest” spin-1 $\Rightarrow pD \uparrow \Rightarrow h_{1,p}^\perp \otimes h_{1,D}$

The EIC impact

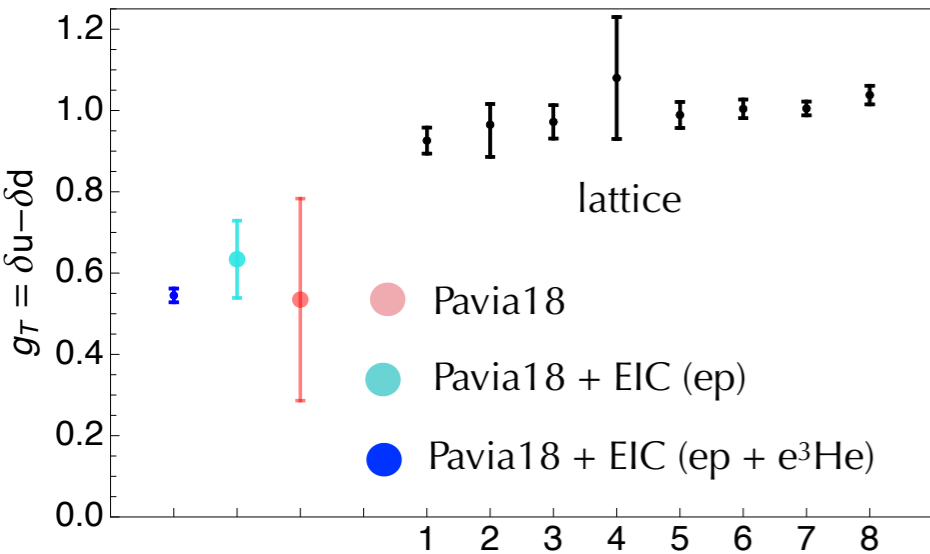
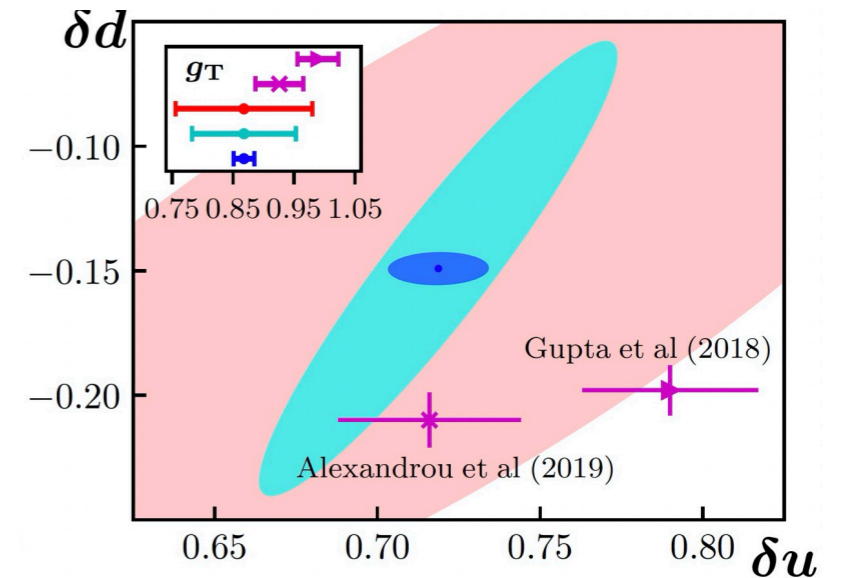


arXiv:2103.05419,
N.P.A in press

Collins effect

- JAM20
- JAM20 + EIC(ep)
- JAM20 + EIC(ep+e³He)

$\mathcal{L}=10 \text{ fb}^{-1}$, 8223 data pts.
proton [GeV]: 5x41, 5x100, 10x100, 18x275
³He [GeV]: 5x41, 5x100, 18x100

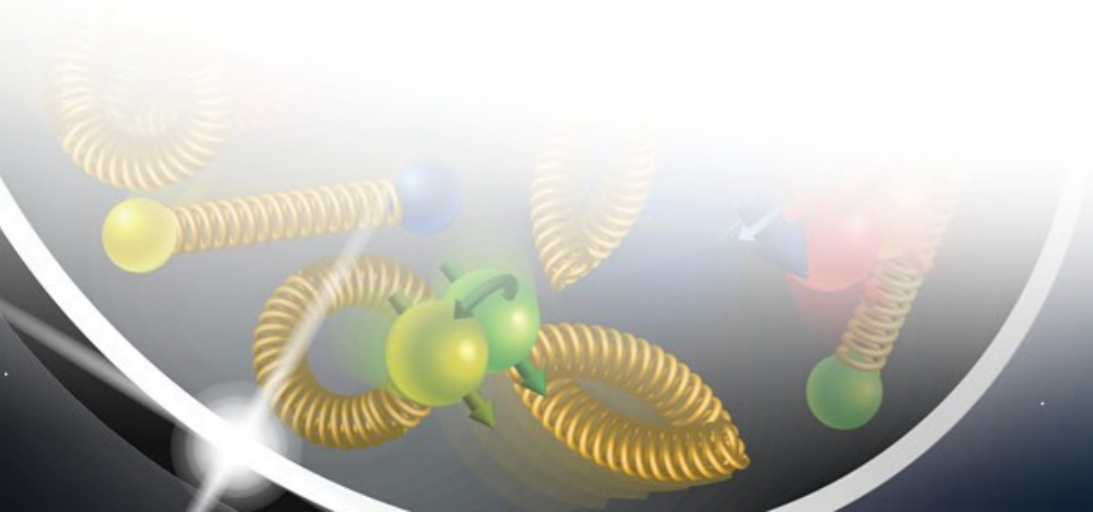
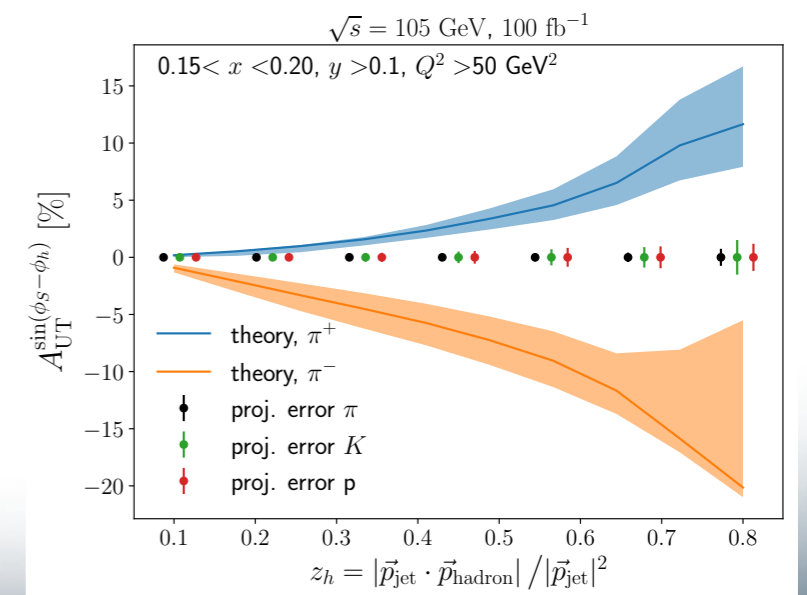


di-hadron mechanism

$\mathcal{L}=10 \text{ fb}^{-1}$, 3852 data pts, proton & ³He [GeV]: 10x100

hadron-in-jet Collins effect

Arratia et al., arXiv:2007.07281



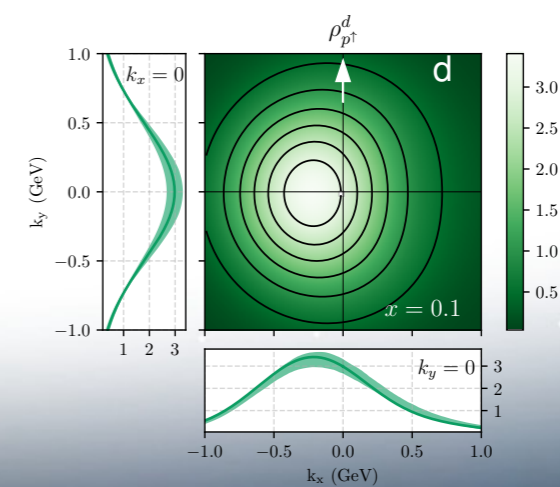
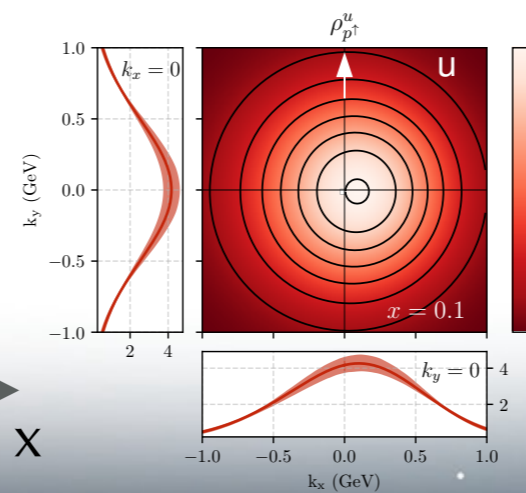
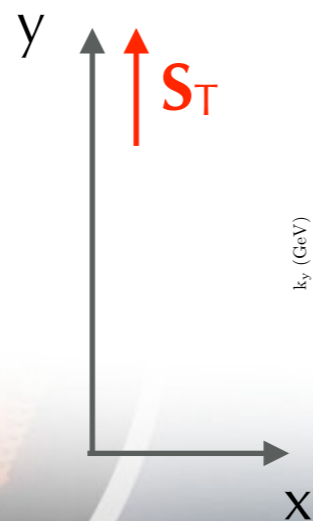
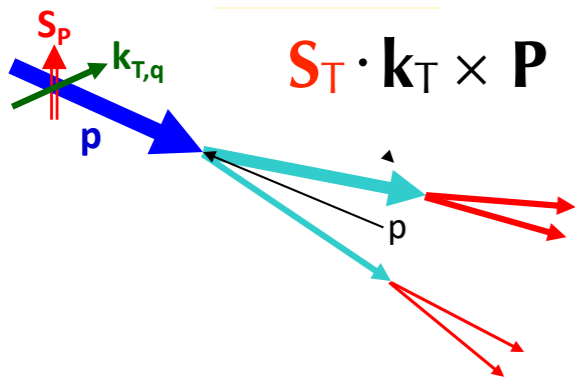
Sivers effect

Sivers

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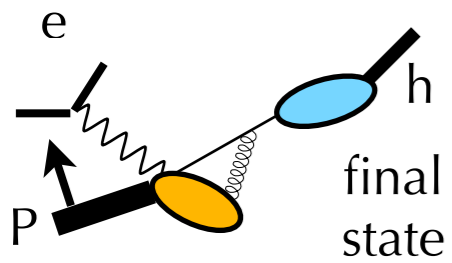
distortion of quark momentum distribution by nucleon spin

Bacchetta et al., P.L. **B827** (22) 136961, arXiv:2004.14278



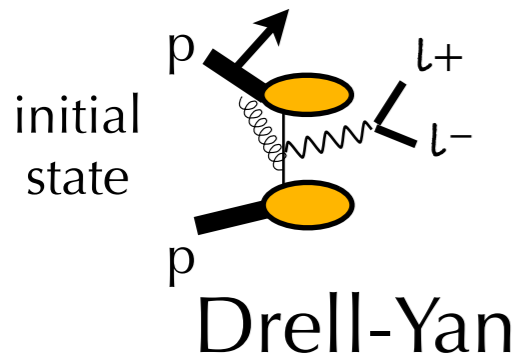
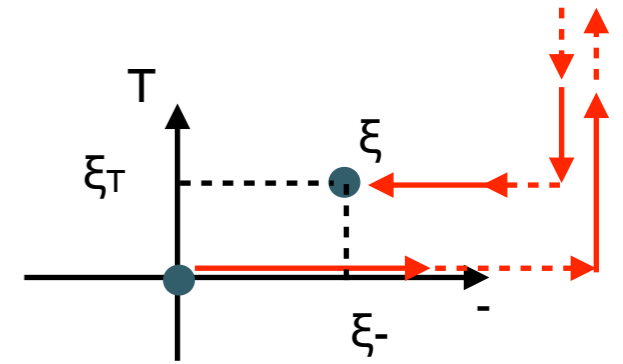
Sivers

the quark Sivers TMD is not universal !



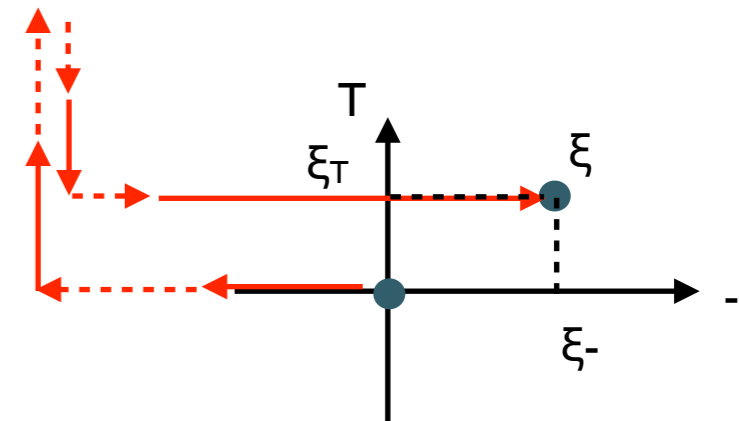
SIDIS

in SIDIS, gauge link structure is "future pointing" \rightarrow describes residual color final-state interactions



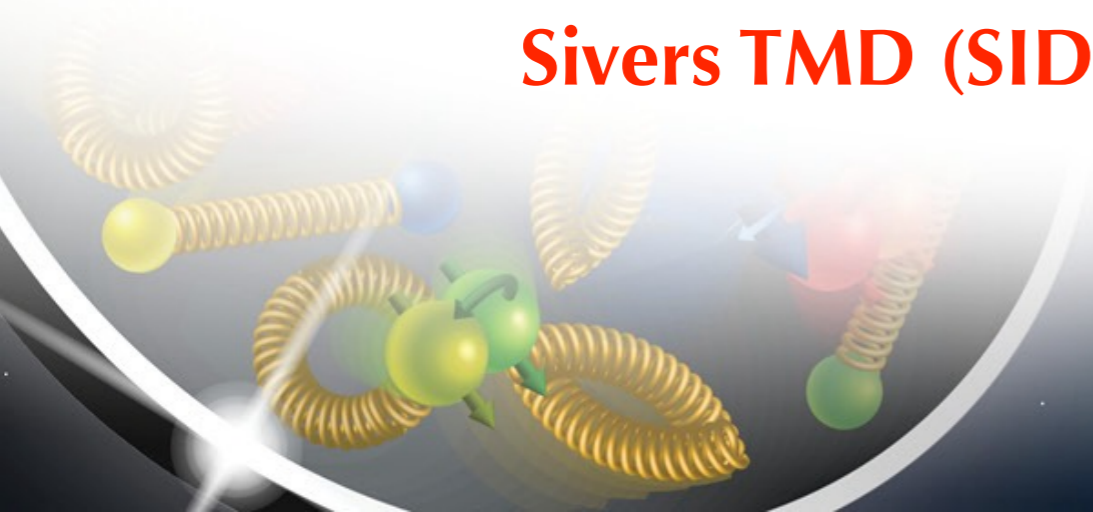
Drell-Yan

in Drell-Yan, gauge link structure is "past pointing" \rightarrow describes color initial-state interactions



Prediction of QCD:

Sivers TMD (SIDIS) = - Sivers TMD (Drell-Yan)



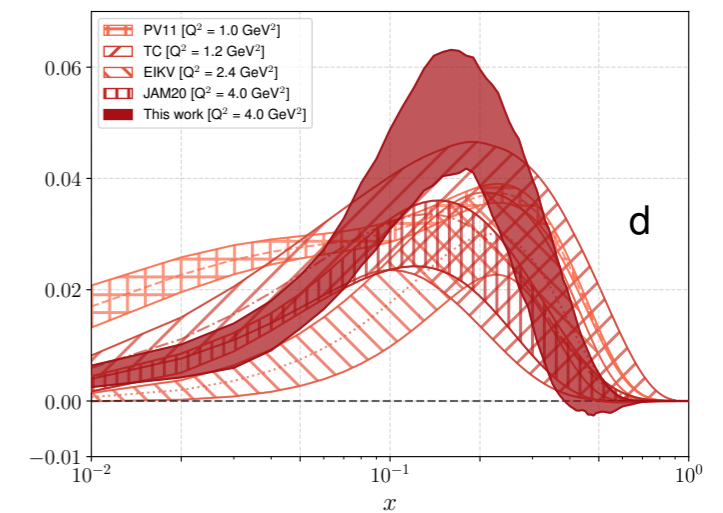
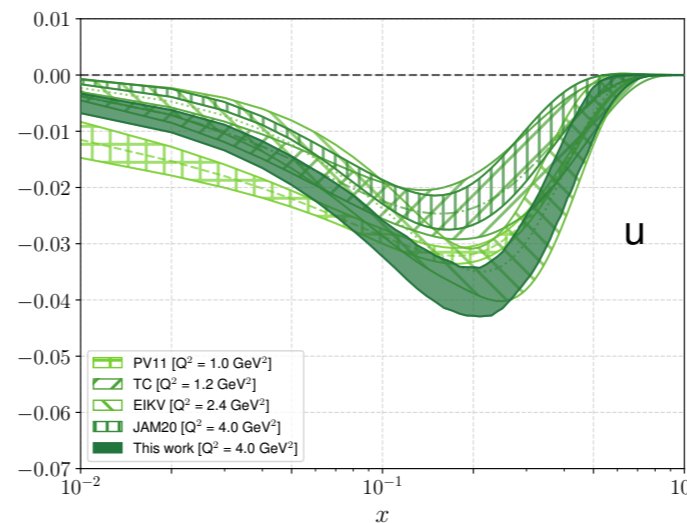
Sivers Phenomenology

most recent extractions of quark Sivers

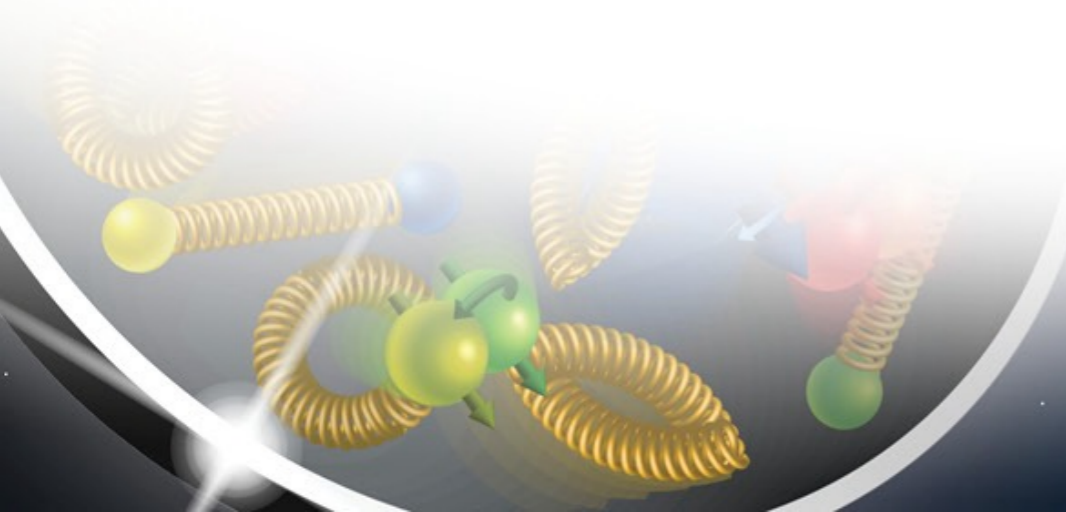
	Framework	SIDIS	DY	W/Z production	e+e-	N of points
JAM 2020 arXiv:2002.08384	extended parton model	✓	✓	✓	✓	517
Pavia 2020 arXiv:2004.14278	LO+NLL	✓	✓	✓	✗	150
EKT 2020 arXiv:2009.10710	NLO+N ² LL	✓	✓	✓	✗	243
BPV 2020 arXiv:2012.05135 arXiv:2103.03270	ζ prescription	✓	✓	✓	✗	76

all parametrizations are in fair agreement for valence flavors

sea-quarks $\sim O(10^{-3})$ smaller



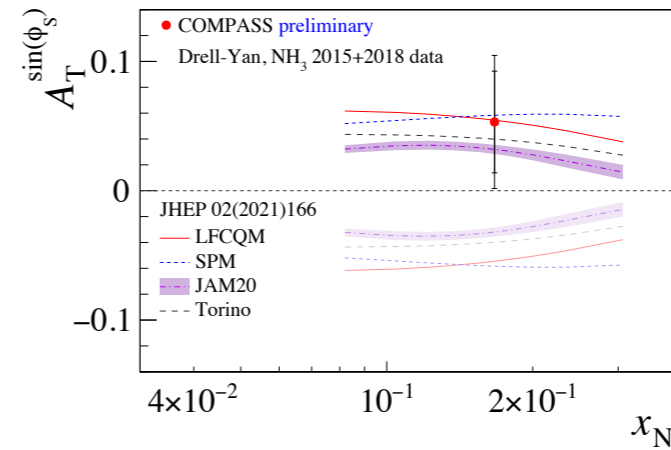
Bacchetta et al., arXiv:2004.14278



The Sign Change Puzzle

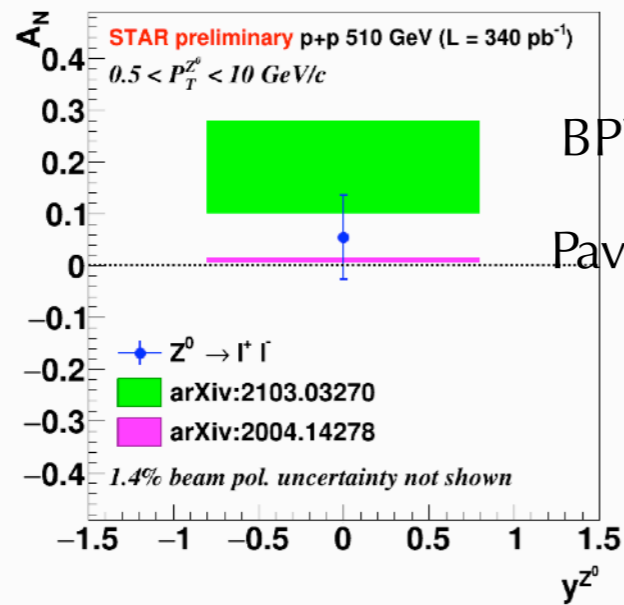
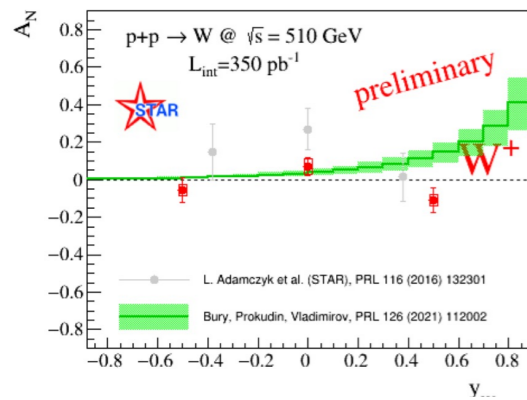
- π p^\uparrow DY by Compass: $f_{1,\pi} \otimes f_{1T,p}^\perp$
compatible with sign change

talk by R. Longo



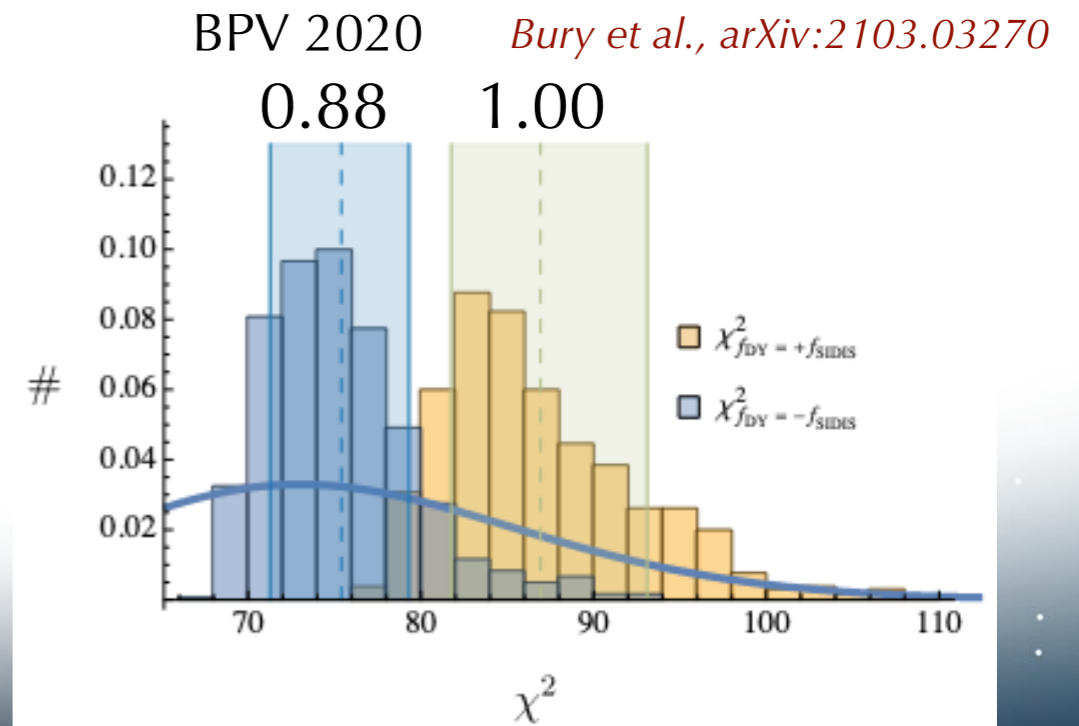
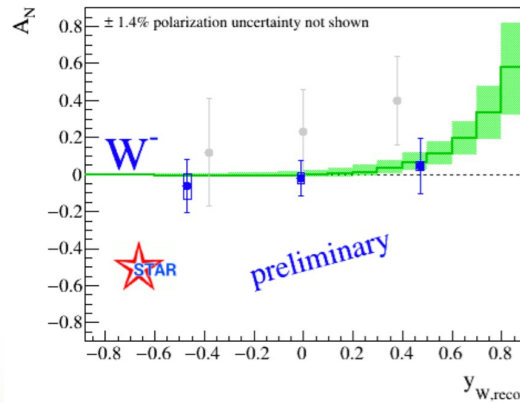
sign change

no sign change



predictions on recent STAR DY data

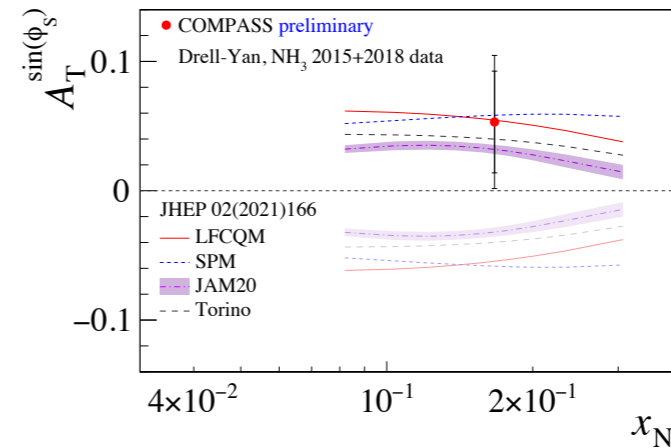
talk by W.W. Jacobs



The Sign Change Puzzle

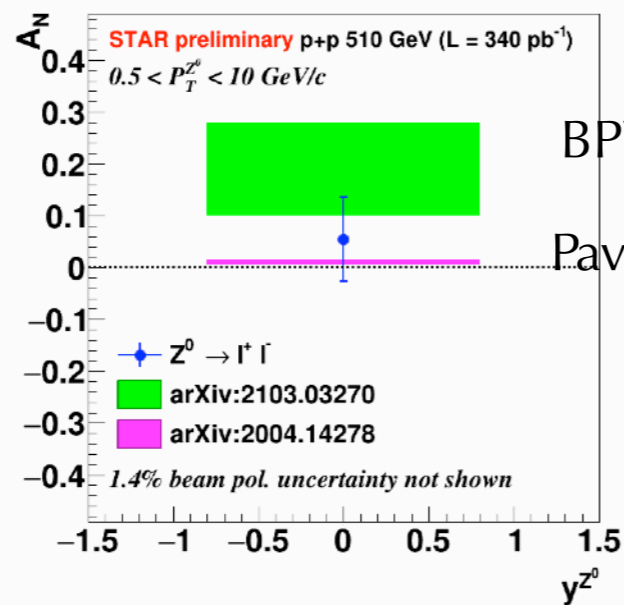
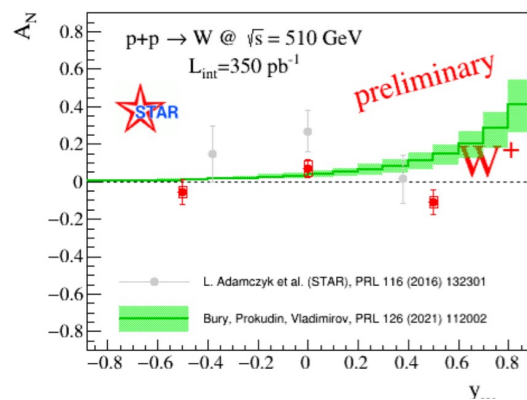
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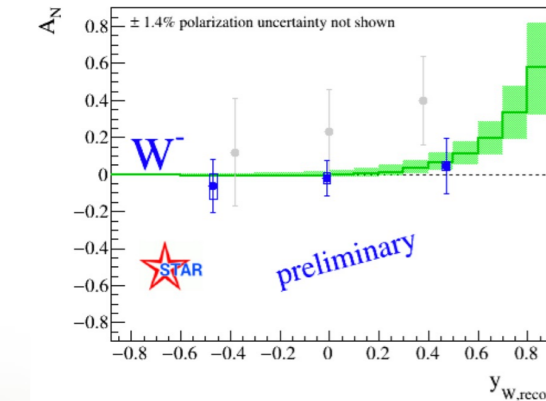
sign change

no sign change

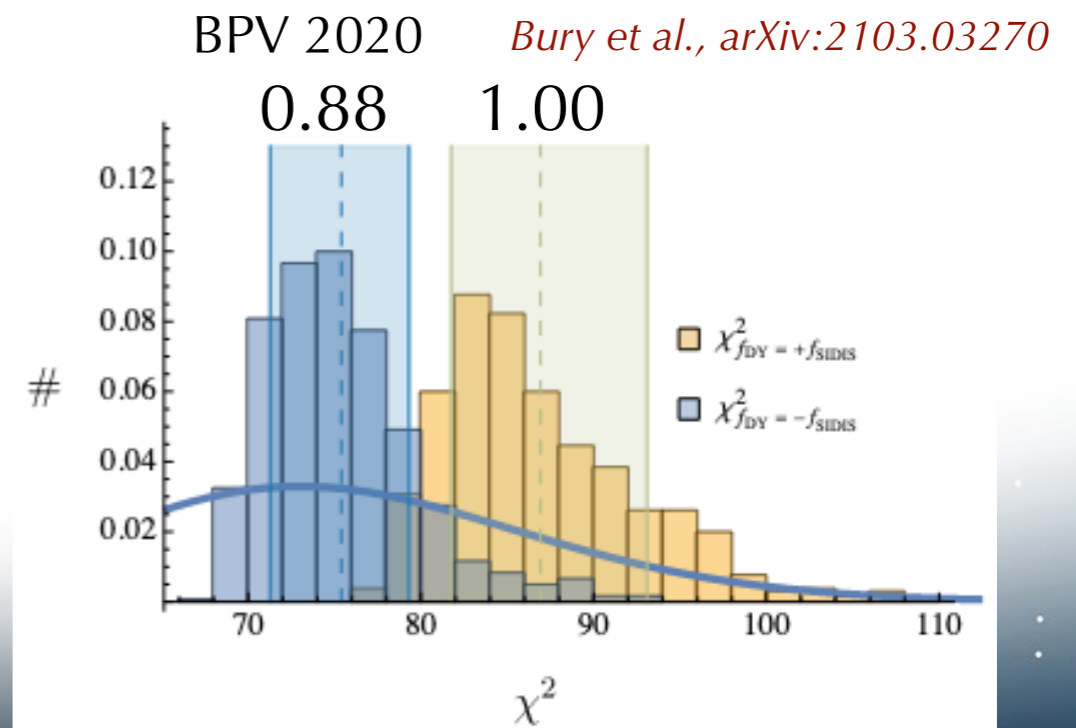


predictions on recent STAR DY data

talk by W.W. Jacobs



still not enough to confirm
sign change:



What about gluons ?

- TMDs are related to hadronic matrix elements of bilocal operators; color gauge links must connect the two points to restore color gauge invariance; **gluons have a more complicated structure than quarks:**

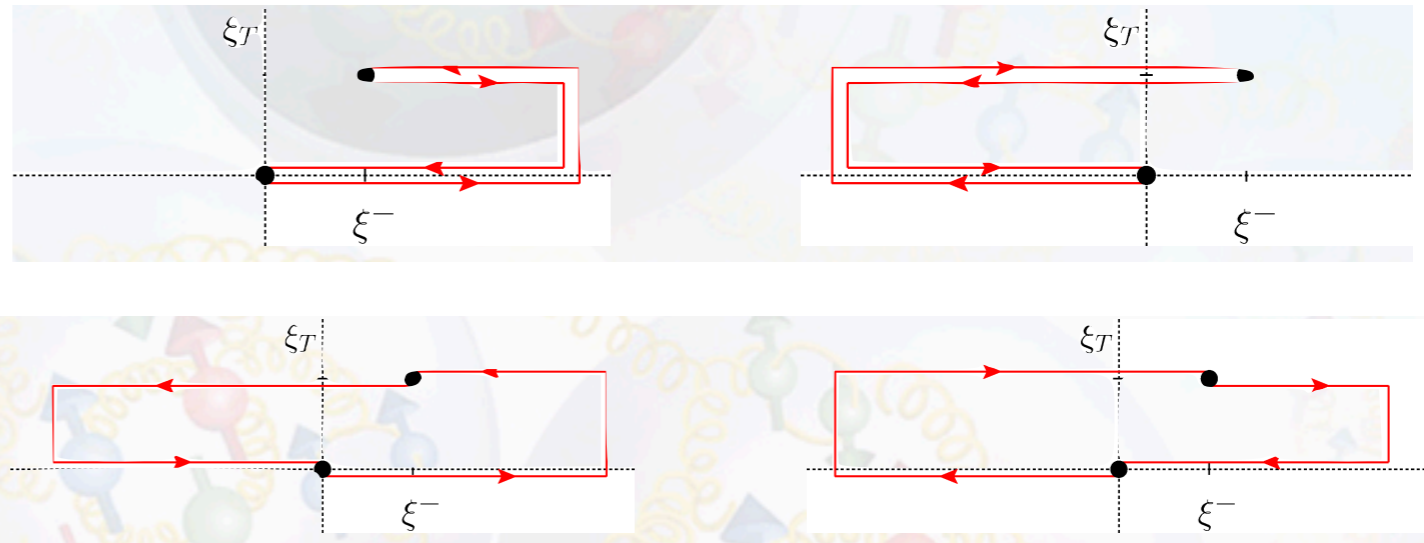
different processes



different TMDs !

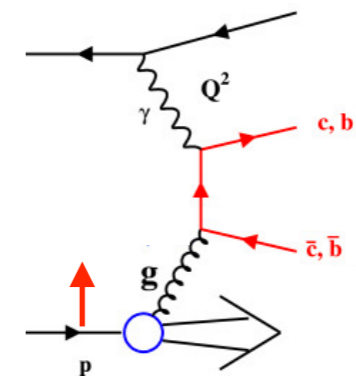
f-type (WW)
[+,+] , [-,-]

d-type (dipole)
[+,-] , [-,+]



- $f_{1T}^{\perp,g[+,+]}$ can be extracted in $ep^{\uparrow} \rightarrow e'Q\bar{Q}X$ at the EIC

- $f_{1T}^{\perp,g[+,-]}$ at small x related to the spin-dep. Odderon only contribution to $pp^{\uparrow} \rightarrow h^{\pm}X$ at $x_F < 0$; RHIC / NICA ?



talk by D. Boer

$f_{1T}^{\perp,g[+,+]}$	$ep^{\uparrow} \rightarrow e'Q\bar{Q}X$	EIC
	$ep^{\uparrow} \rightarrow e' \text{jet jet } X$	EIC
$f_{1T}^{\perp,g[-,-]}$	$p^{\uparrow}p \rightarrow \gamma\gamma X$	RHIC
$f_{1T}^{\perp,g[+,-]}$	$p^{\uparrow}A \rightarrow \gamma^{(*)} \text{jet } X$	RHIC
	$p^{\uparrow}A \rightarrow h X (x_F < 0)$	RHIC & NICA

New and future data for Sivers studies

- new 3-D analysis of Collins effect from Hermes, with final $h = \pi, K, p, pbar$

talk by G. Schnell

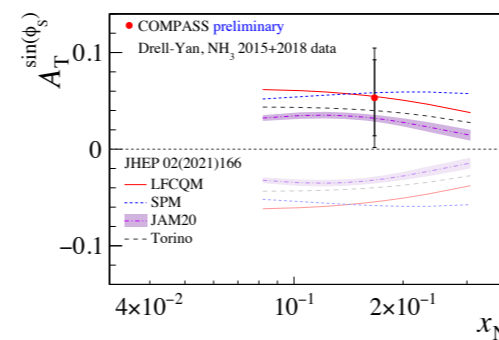
Airapetian et al., JHEP12 (2020) 010

- Sivers effect for ρ^0 measured by Compass

talks by A. Bressan
F. Bradamante

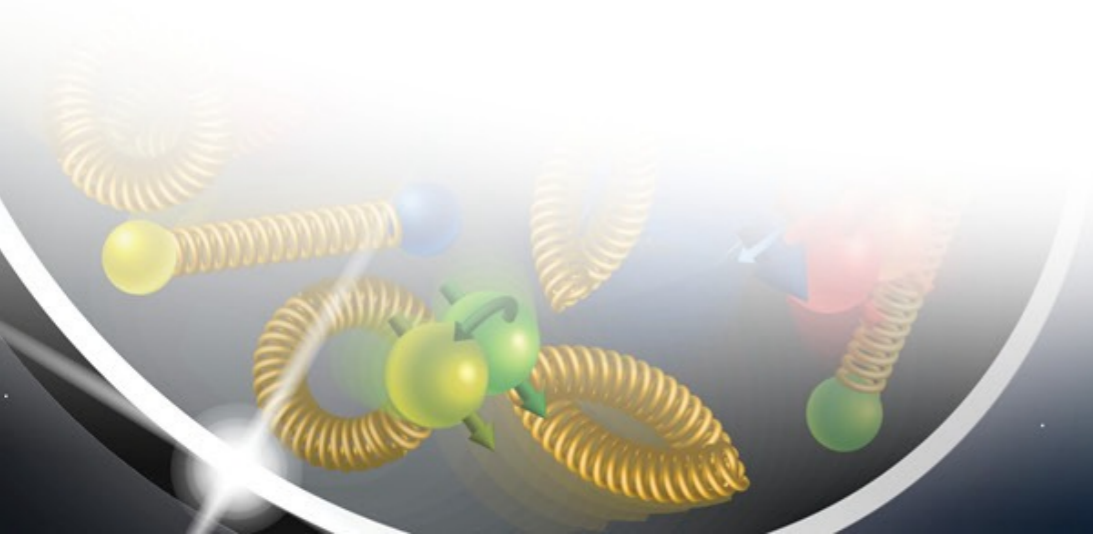
- $\pi p \uparrow$ DY by Compass: $f_{1,\pi} \otimes f_{1T,p}^\perp$

talk by R. Longo



sign change

no sign change



New and future data for Sivers studies

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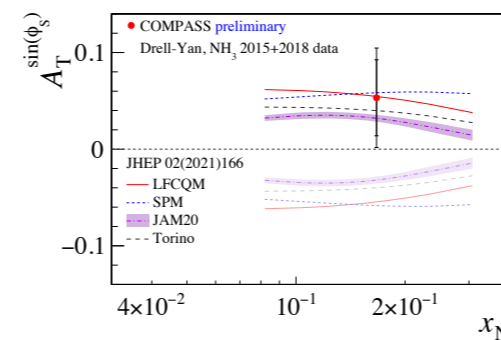
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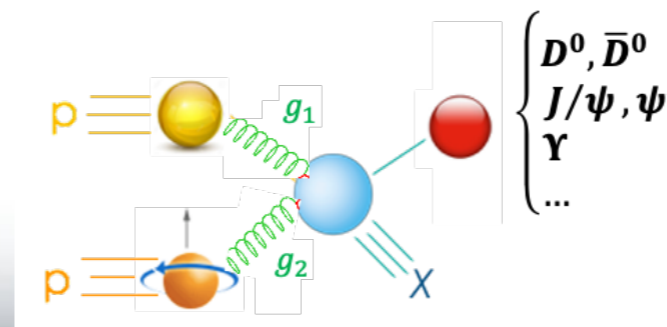
sign change

no sign change

- FermiLab E1039 "SpinQuest" $\Rightarrow pp \uparrow$ & $pD \uparrow \Rightarrow f_{1,p} \otimes f_{1T,p,D}^\perp$

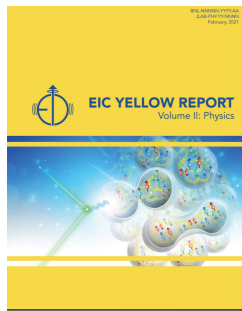
talk by N. Wuerfel

- LHCspin $\Rightarrow p-p \uparrow$ DY $\Rightarrow f_{1,p} \otimes f_{1T,p}^\perp$

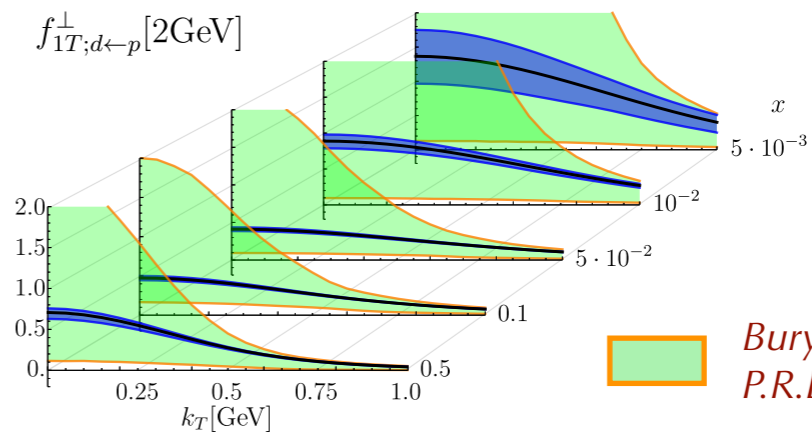
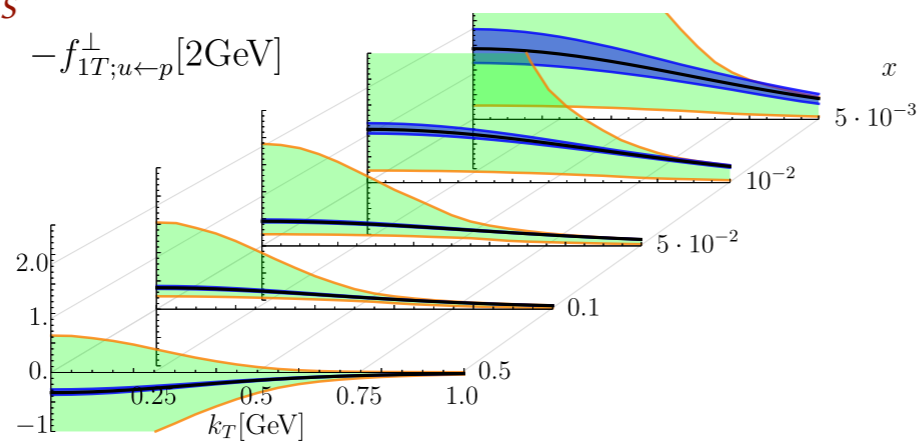



talk by P. Di Nezza


The EIC Impact



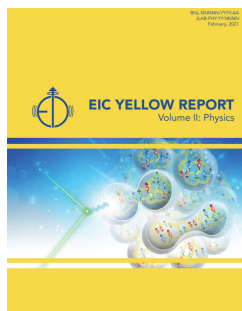
arXiv:2103.05419,
N.P.A in press



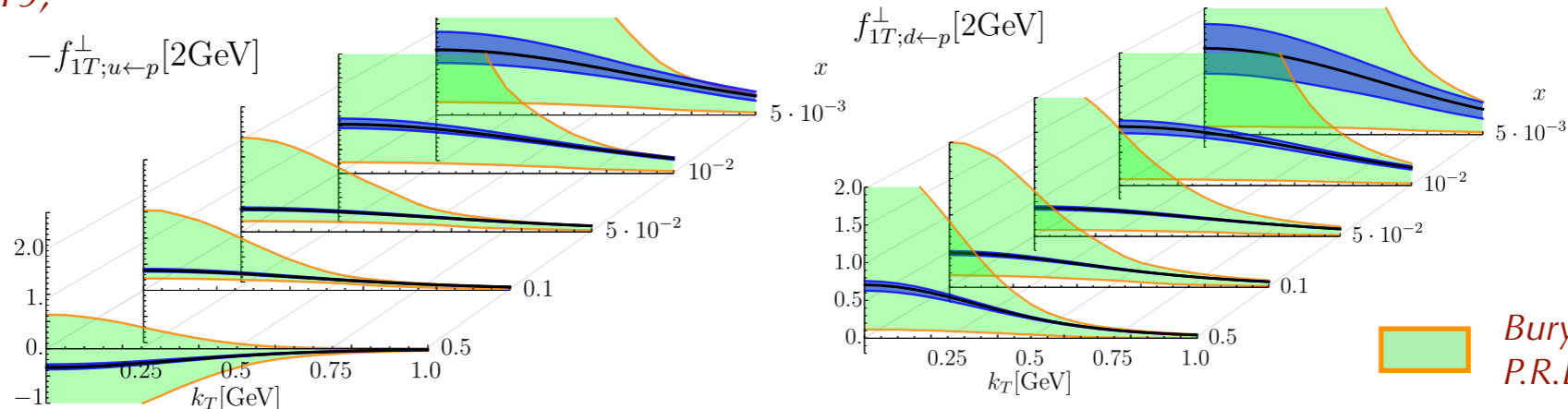
 *Bury et al.,*
P.R.L. 126 (21) 112002

 + EIC 5×41 GeV
 18×275

The EIC Impact



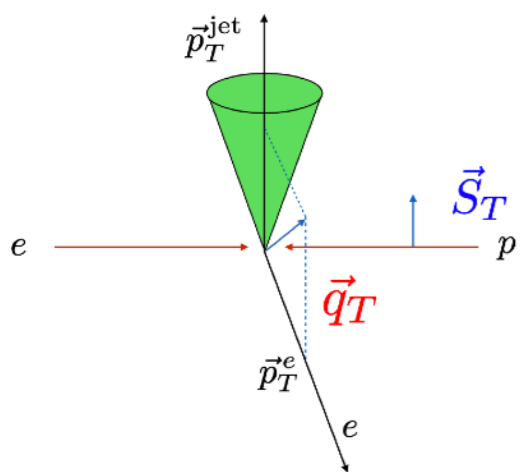
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Bury et al.,
P.R.L. **126** (21) 112002
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opportunities with jets and Heavy Flavors

talk by F. Ringer



electron-jet azimuthal correlations

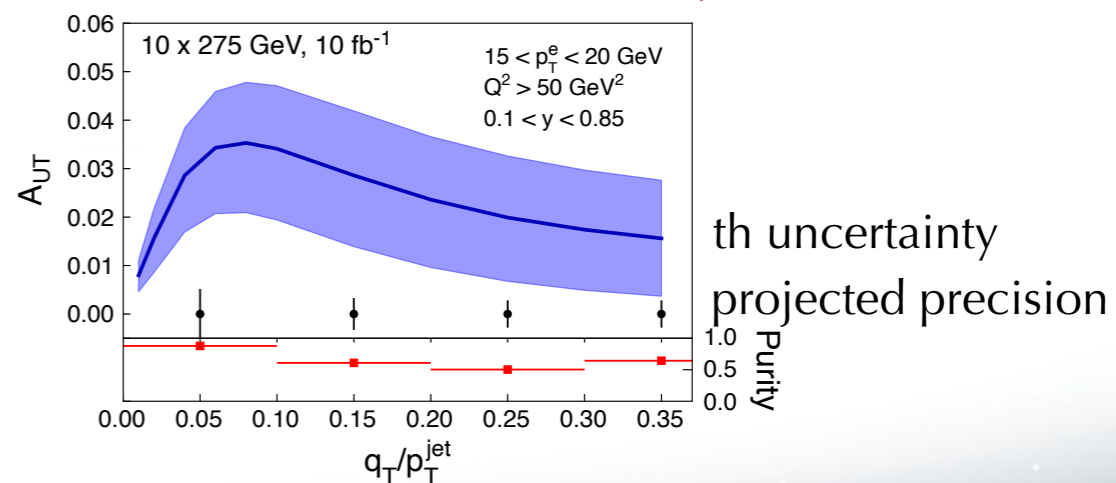
$$|\vec{q}_T| = |\vec{p}_T^e + \vec{p}_T^{\text{jet}}| \ll |\vec{p}_T^{\text{jet}}|$$



$$A_{\text{UT}} \sim d\sigma(S_T) - d\sigma(-S_T)$$

Sivers effect free from TMD FF

Arratia et al., arXiv:2007.07281



also access to gluon Sivers TMD from $D^0\bar{D}^0$, charm di-jets and J/ψ production

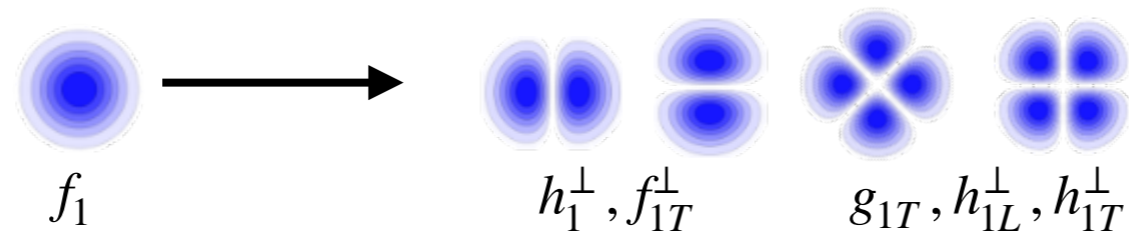
Zheng et al., arXiv:1805.05290

Rajesh et al., arXiv:2108.04866

The TMD “zoo” at leading twist

		Quark polarization		
		Unpolarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 = \odot$	\times	$h_1^\perp = \odot \uparrow - \odot \downarrow$
	L	\times	$g_1 = \odot \rightarrow - \odot \leftarrow$	$h_{1L}^\perp = \odot \nearrow - \odot \nwarrow$
	T	$f_{1T}^\perp = \odot \uparrow - \odot \downarrow$	$g_{1T} = \odot \rightarrow - \odot \leftarrow$	$h_1 = \odot \uparrow - \odot \downarrow$ $h_{1T}^\perp = \odot \nearrow - \odot \nwarrow$

deformations induced by spin-momentum correlations



each TMD is connected to a specific measurable SIDIS spin asymmetry

The unpolarized quark TMD f_{1q}

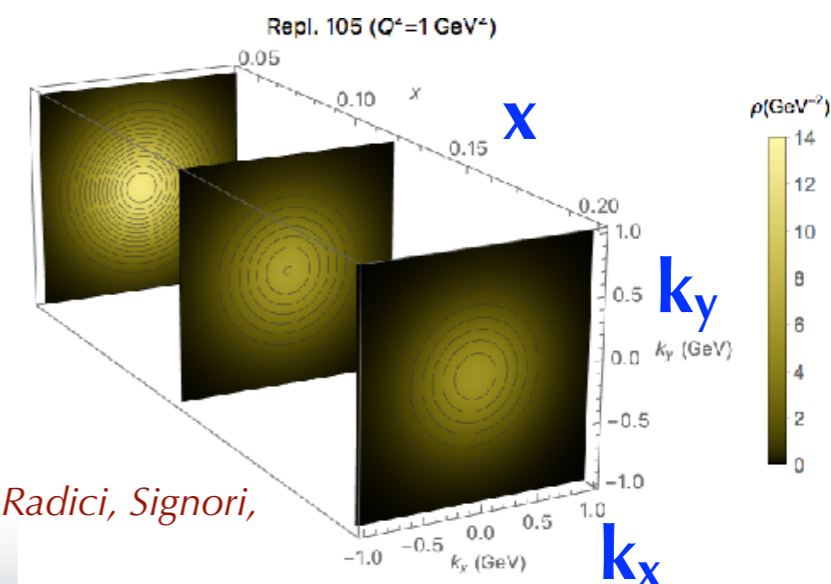
the best known TMD (most recent fits)

Lessons to be learnt :

- non-perturbative k_T dependence is not a simple Gaussian
- average $\langle k_T^2 \rangle$ strongly depends on x , and might depend on flavor (in particular for fragmentation; recent attempt on SV19)
- Gaussian non perturbative evolution seems preferred
- modern fits can reach $N^3LL+NNLO$ perturbative accuracy with reduced $\chi^2 \sim 1$ on thousands data points

	Framework	HERMES	COMPASS	DY	Z production	N of points	χ^2/N_{points}
PV 2017 arXiv:1703.10157	NLL	✓	✓	✓	✓	8059	1.5
SV 2017 arXiv:1706.01473	NNLL'	✗	✗	✓	✓	309	1.23
BSV 2019 arXiv:1902.08474	NNLL'	✗	✗	✓	✓	457	1.17
SV 2019 arXiv:1912.06532	N^3LL	✓	✓	✓	✓	1039	1.06
PV 2019 arXiv:1912.07550	N^3LL	✗	✗	✓	✓	353	1.07
SV19 + flavor dep. arXiv:2201.07114	N^3LL	✗	✗	✓	✓	309	<1.08>
MAPTMD 2022 arXiv:2206.07598	N^3LL	✓	✓	✓	✓	2031	1.06

tomography in momentum space



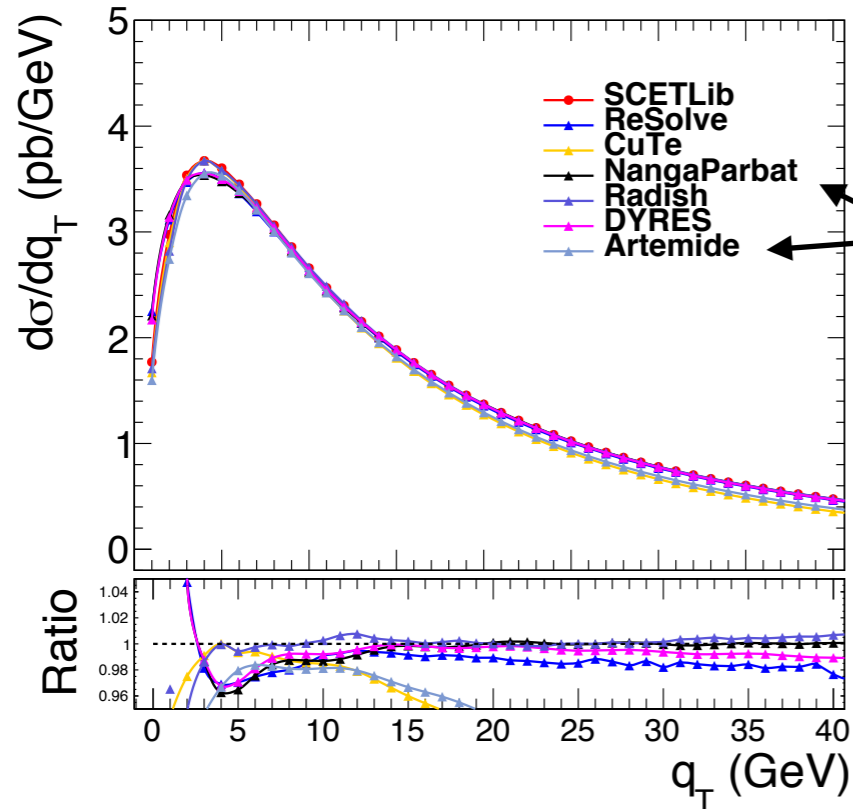
PV 2017
Bacchetta, Delcarro, Pisano, Radici, Signori,
JHEP **06** (17) 081

The unpolarized quark TMD f_{1q}

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same accuracy as PDF
benchmarking codes @LHC

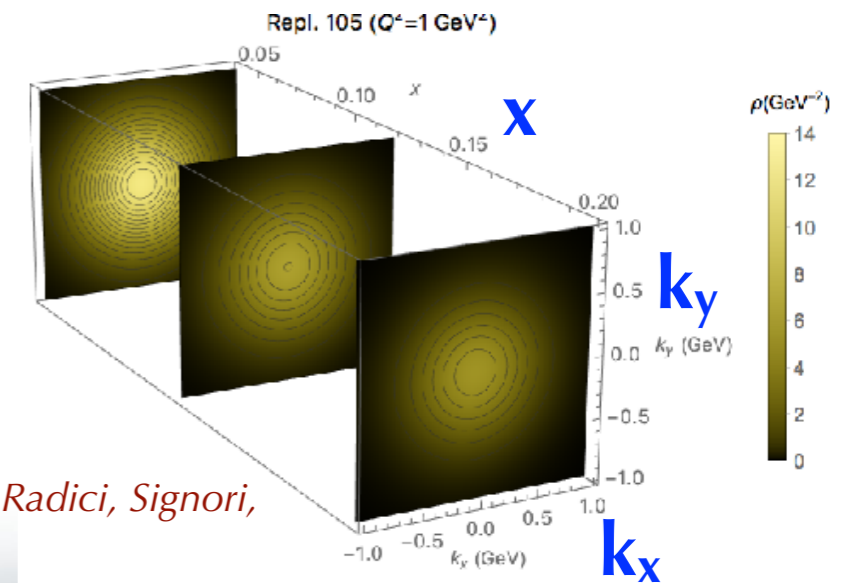


Z production at $\eta=0$ (ATLAS kin)

G. Bozzi, I. Scimemi (eds.) et al.,
Yellow Report of CERN EW WG, in preparation

tomography in
momentum space

PV 2017
Bacchetta, Delcarro, Pisano, Radici, Signori,
JHEP **06** (17) 081



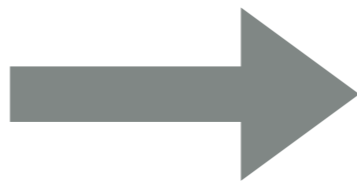
The MAPTMD22 fit

the best known TMD (most recent fits)

Bacchetta et al., arXiv:2206.07598

the new MAPTMD22 fit

talk by V. Bertone



	Framework	HERMES	COMPASS	DY	Z production	N of points	χ^2/N_{points}
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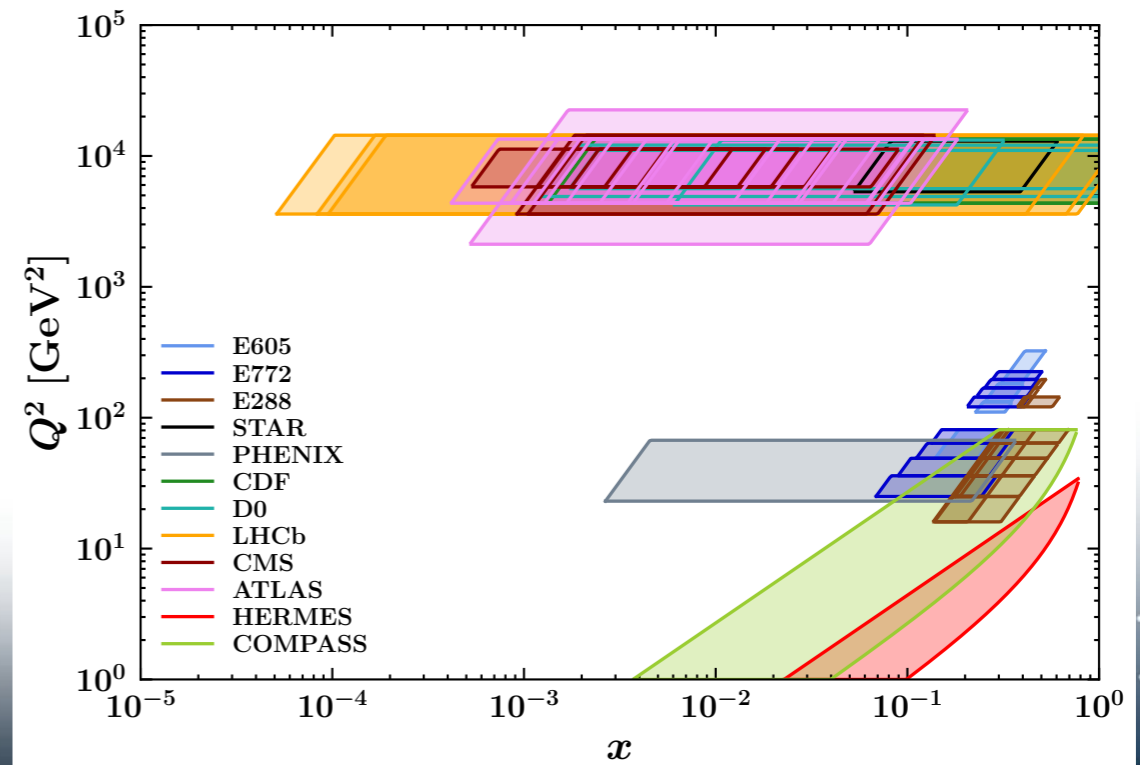
kin. cuts and coverage

DY $q_T/Q < 0.2$

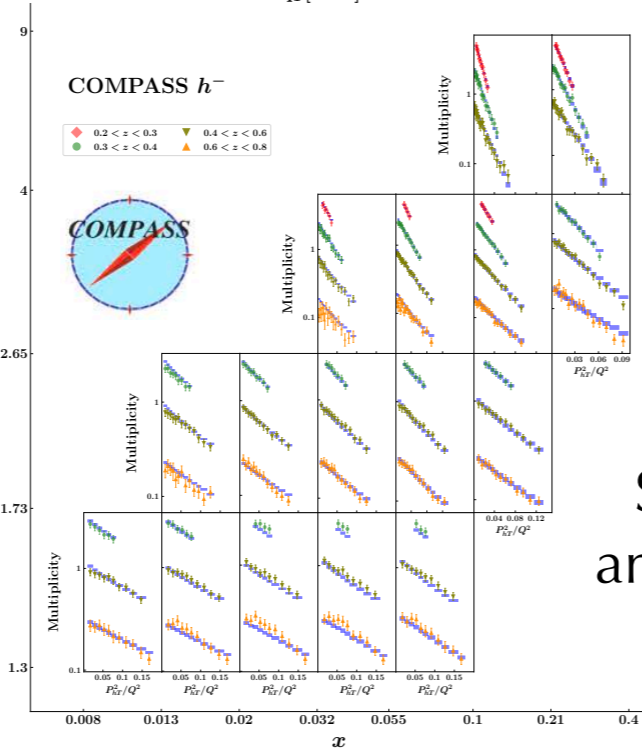
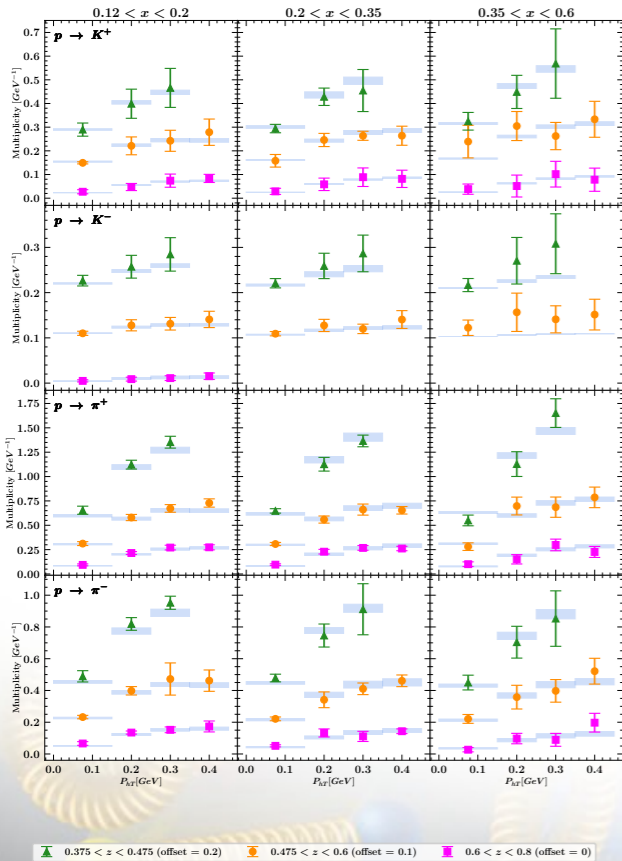
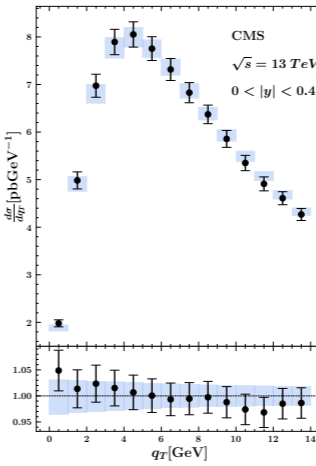
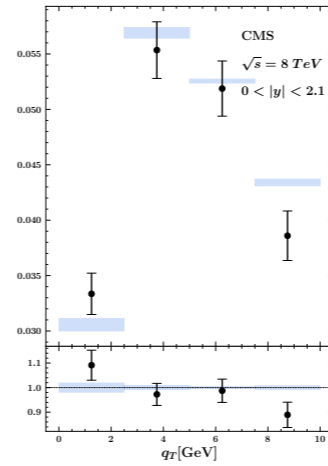
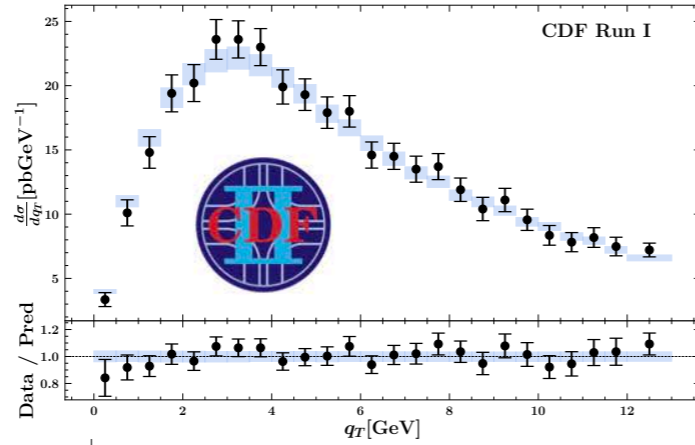
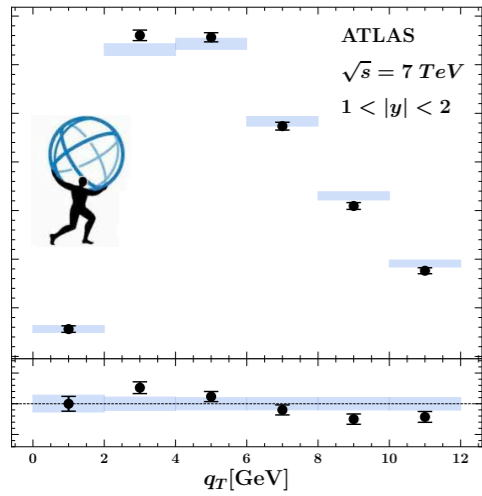
SIDIS $P_{hT} = \min[\min[0.2 Q, 0.5 zQ] + 0.3, zQ]$

$Q > 1.4 \text{ GeV}, 0.2 < z < 0.7$

2031 exp. pts., 21 parameters

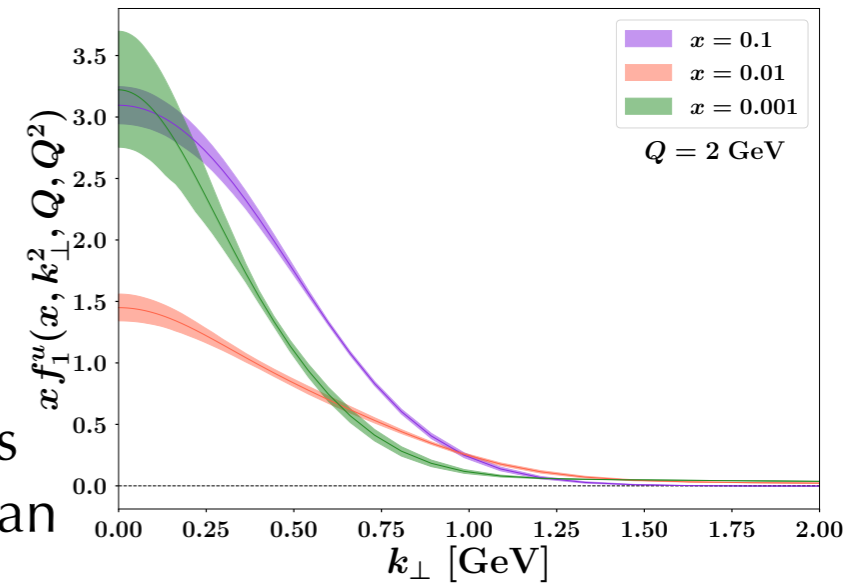


The MAPTMD22 fit



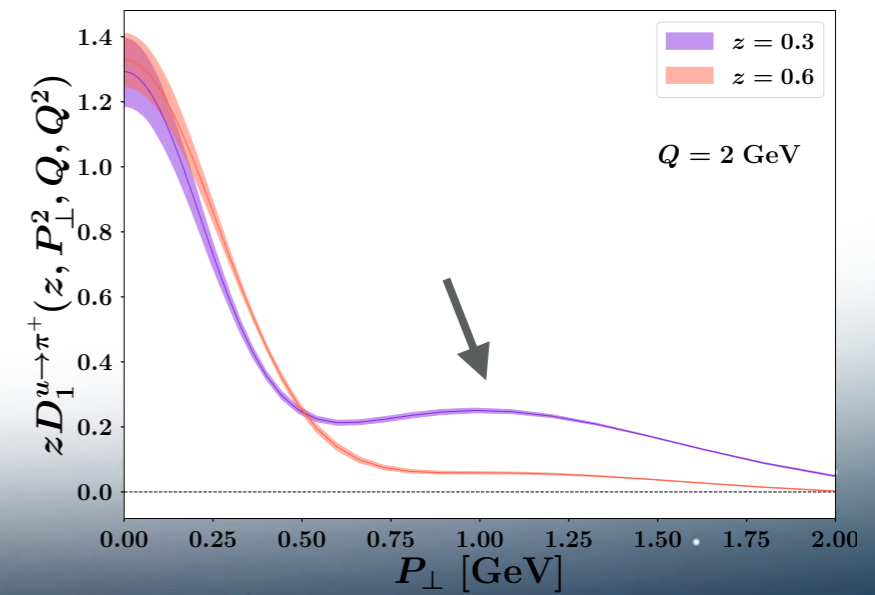
TMD PDF

Sum of two Gaussians
 and a weighted Gaussian
 x-dep. widths



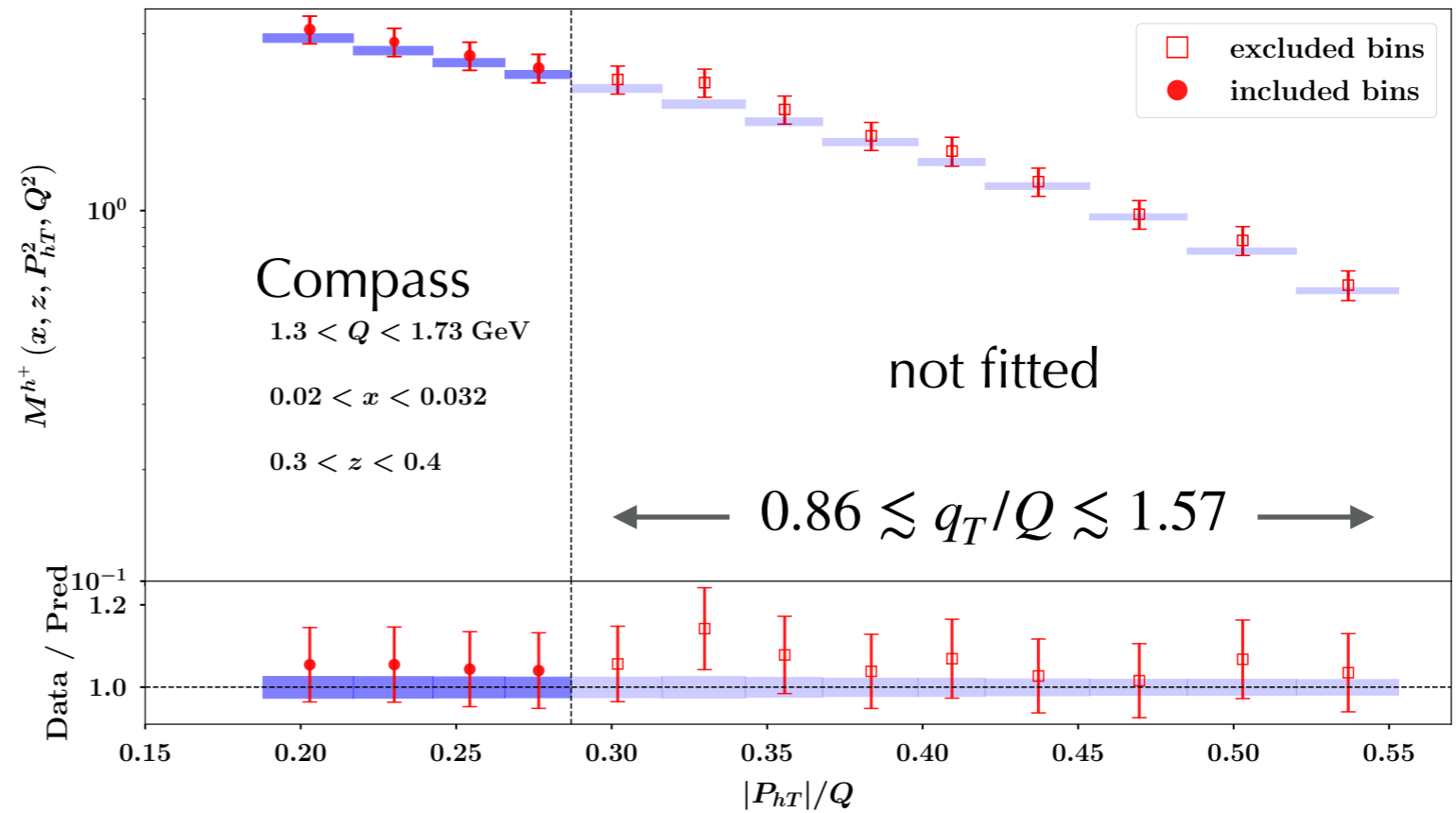
TMD FF

Sum of a Gaussians
 and a weighted Gaussian
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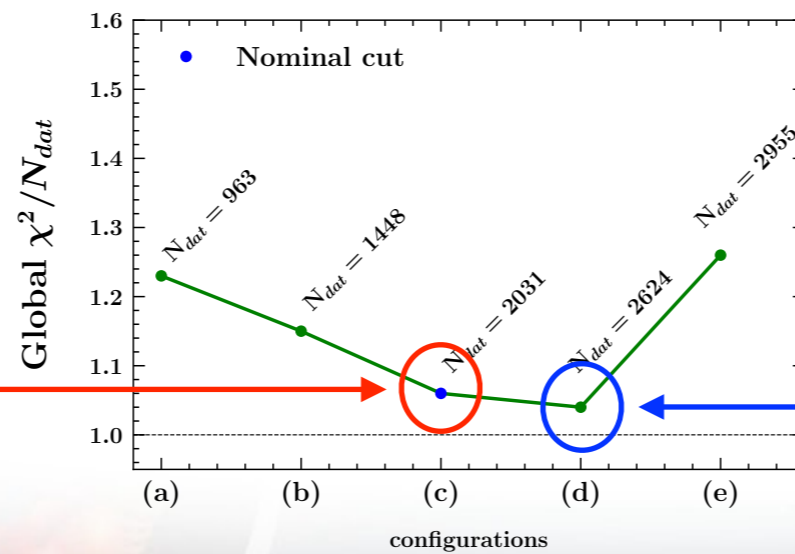


The MAPTMD22 fit

it happens also in other bins
and for Hermes data as well



baseline fit

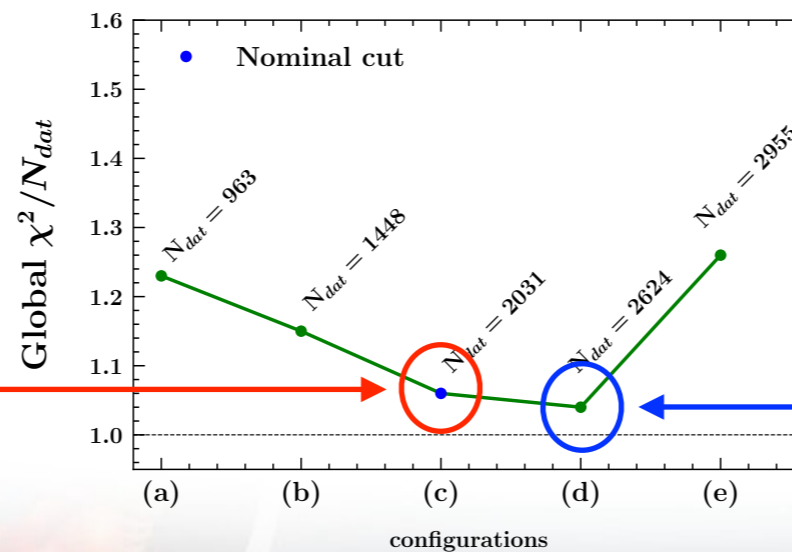
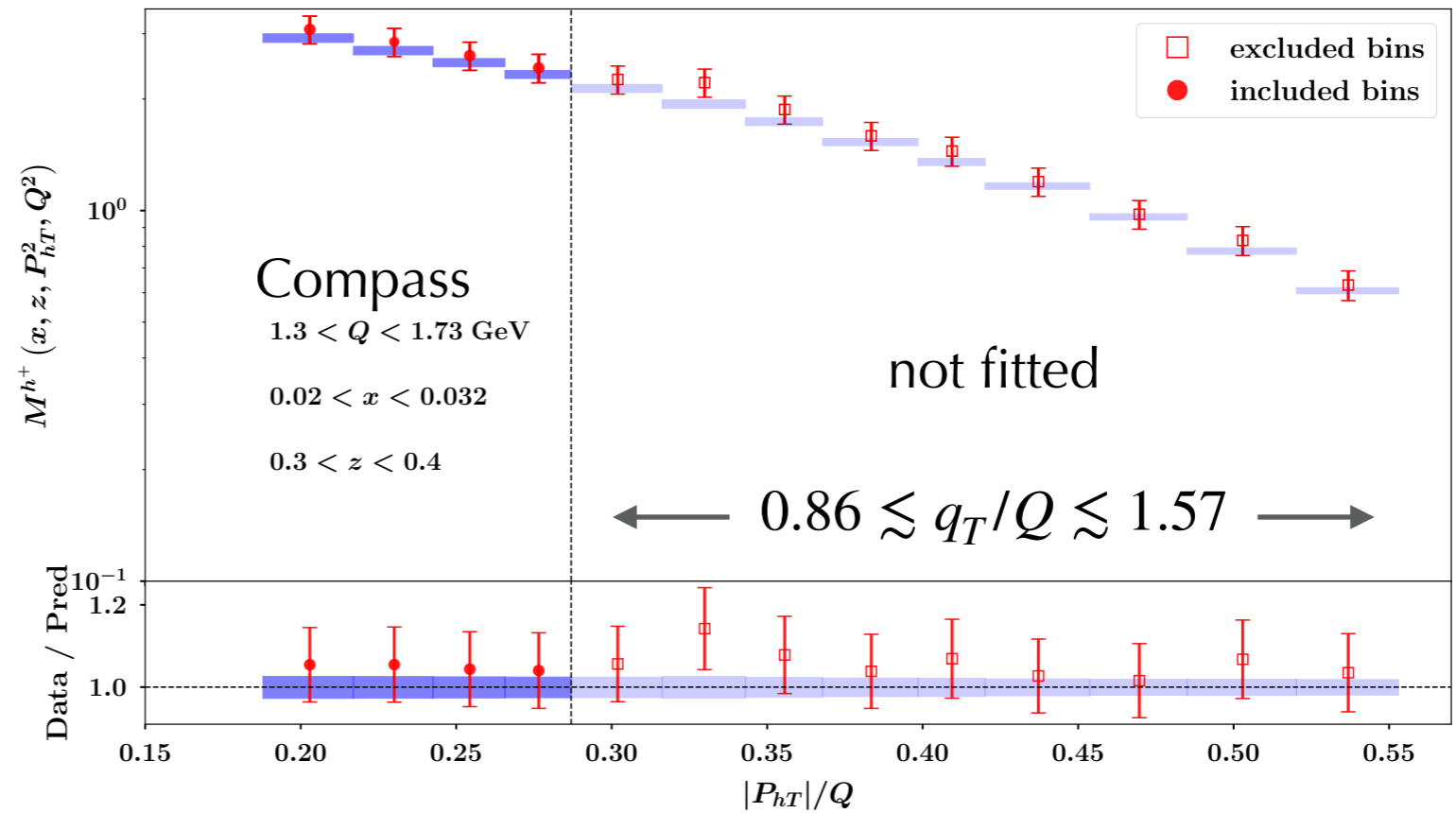


without $q_T < Q$

The MAPTMD22 fit

Where is the limit of applicability of the TMD formalism ?!

it happens also in other bins and for Hermes data as well



baseline fit

without $q_T < Q$

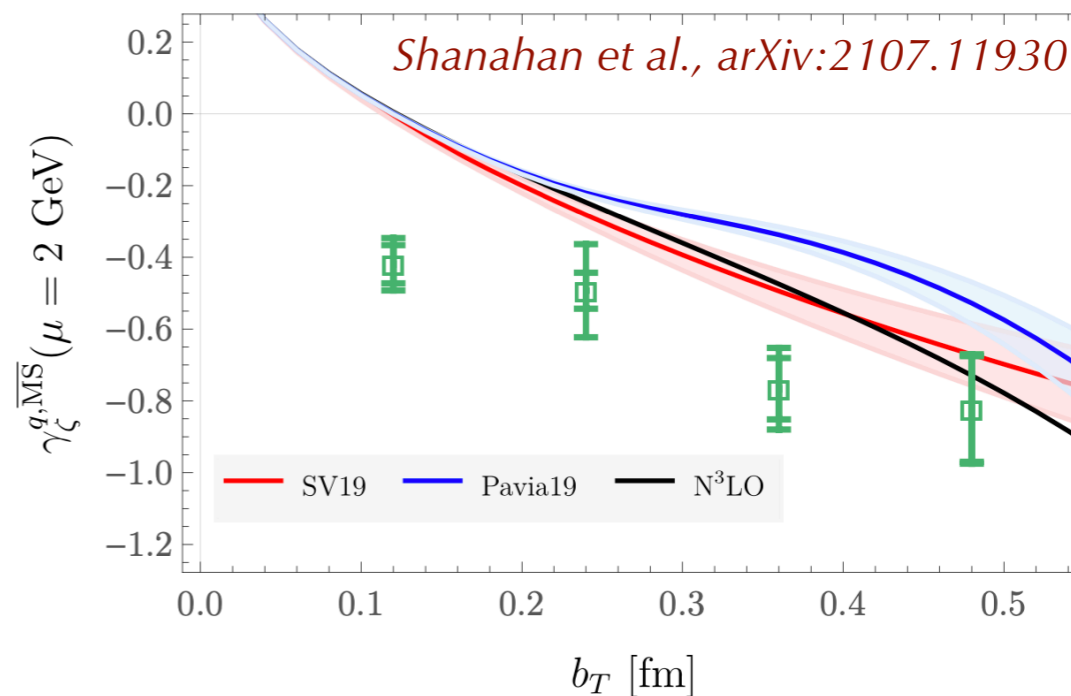
The TMD evolution of f_{1q}

$$\text{TMD}(x, b_T; \mu_f, \zeta_f) = \text{Evo}(\mu_f, \zeta_f; \mu_i, \zeta_i) \text{TMD}(x, b_T; \mu_i, \zeta_i)$$

$$\text{Evo}(\mu_f, \zeta_f; \mu_i, \zeta_i) = \exp\left[S_{\text{pert}}(\mu_f, \mu_i; \zeta_f)\right] \exp\left[\frac{1}{2}K(b_T, \mu_b) \ln(\zeta_f/\zeta_i)\right]$$

Collins-Soper kernel $K = \gamma_\zeta = -2\mathcal{D}$ drives the evolution in rapidity ζ (including the unknown non perturbative part); can be computed on lattice

talk by M. Wagman



The TMD evolution of f_{1q}

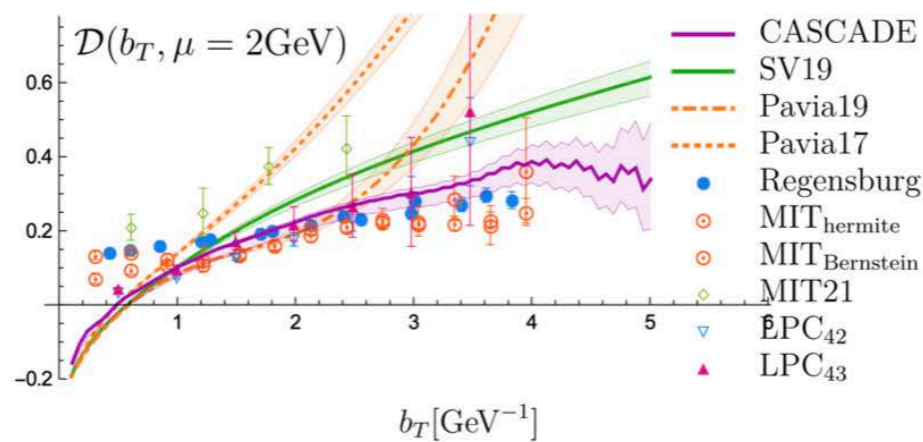
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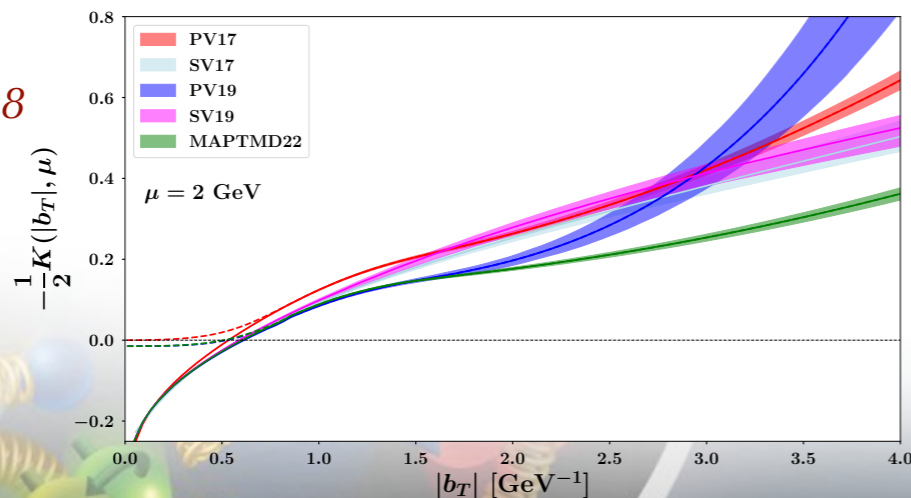
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Martinez & Vladimirov,
arXiv:2206.01105

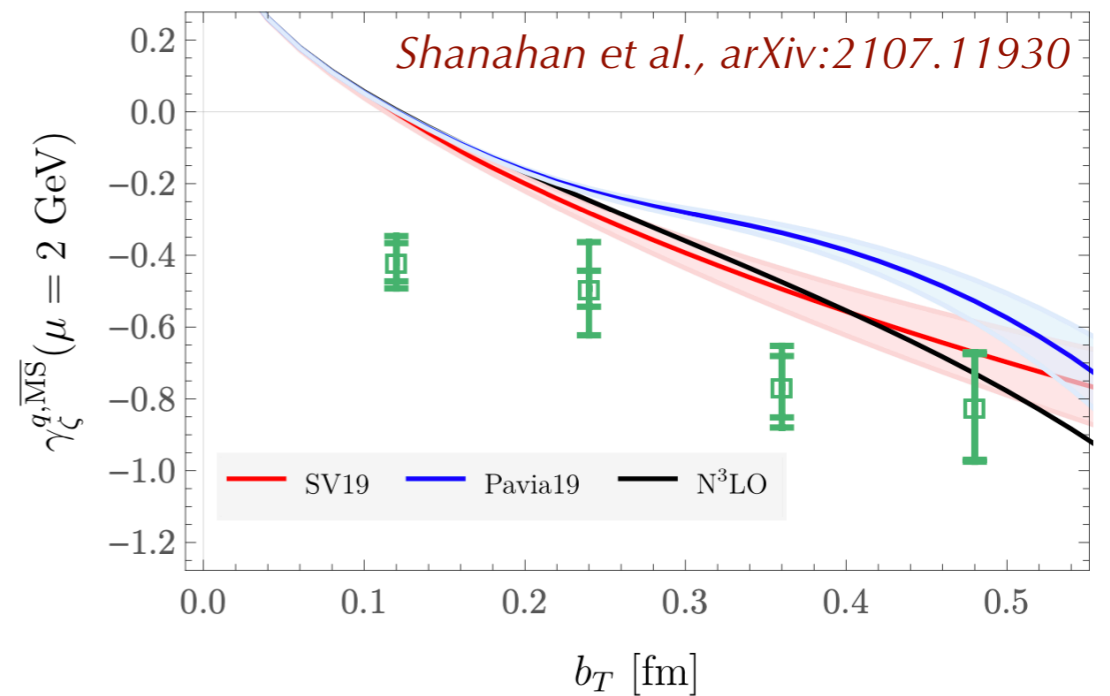
A. Vladimirov talk (CPHI 2022)



Bacchetta et al.,
arXiv:2206.07598



talk by M. Wagman



What about gluons ?

- useful channels: heavy-quarkonium production

talk by M. Echevarria

$$p + p \rightarrow \eta_{c,b} + X$$

$$p + p \rightarrow \chi_{c,b} + X$$

$$p + p \rightarrow H^0 + X$$

$$p + p \rightarrow \gamma + \gamma + X$$

$$p + p \rightarrow J/\psi + \gamma^* + X$$

$$p + p \rightarrow J/\psi + Z + X$$

$$p + p \rightarrow J/\psi + J/\psi + X$$

$$p + p \rightarrow \eta_c + \eta_c + X$$

$$e + p \rightarrow e + c + \bar{c} + X$$

$$e + p \rightarrow e + J/\psi + jet + X$$

$$e + p \rightarrow e + J/\psi + \pi + X$$

$$e + p \rightarrow e + J/\psi + X$$

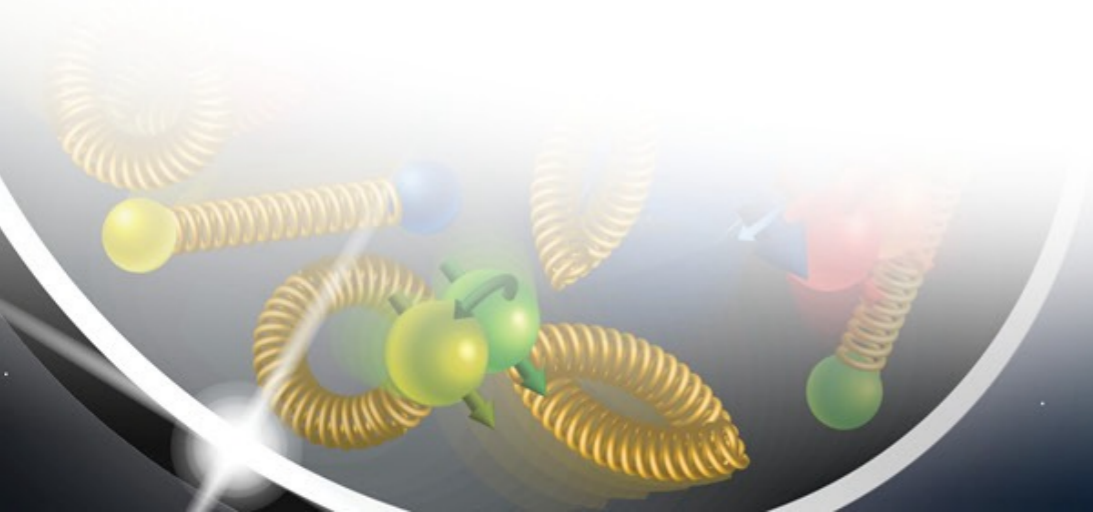
$$e^+ + e^- \rightarrow J/\psi + \pi + X$$

factorization proven

ansatz

2 soft mechanisms:

- soft gluon resum.
- formation of bound state



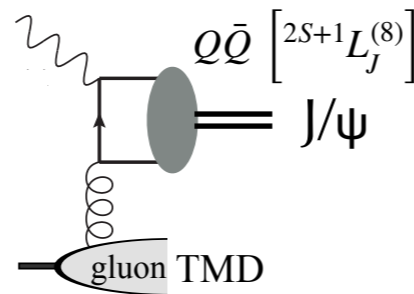
What about gluons ?

- useful channels: heavy-quarkonium production

talk by M. Echevarria

- example: J/ψ production

Bacchetta et al., arXiv:1809.02056
D'Alesio et al., arXiv:1908.00446



$$p + p \rightarrow \eta_{c,b} + X$$

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factorization proven

ansatz

2 soft mechanisms:

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- formation of bound state

- cross section has same structure for quarks: $d\sigma^0 \rightarrow f_1^g \otimes A [\gamma^* g \rightarrow J/\psi] + \cos 2\phi_{J/\psi} h_1^{\perp g} \otimes B [\gamma^* g \rightarrow J/\psi]$

Boer et al., arXiv:2004.06740

Boer et al., arXiv:2102.00003

D'Alesio et al., arXiv:2110.07529

talks by C. Pisano
 L. Maxia
 R. Kishore

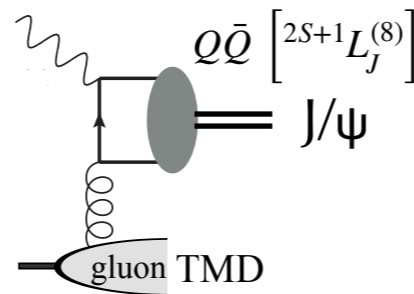
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talks by C. Pisano
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factorization proven

ansatz

- 2 soft mechanisms:
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unknown "Shape Function"

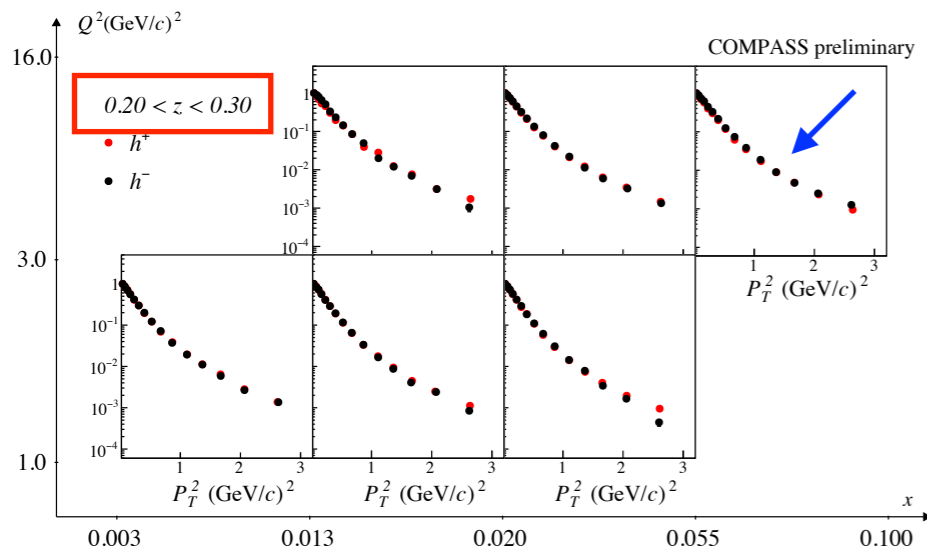
depend on process and (color) structure of $QQ\bar{Q}$
 must be extracted from experiment

opportunity at the EIC

Echevarria, arXiv:1907.06494
Fleming et al., arXiv:1910.03586

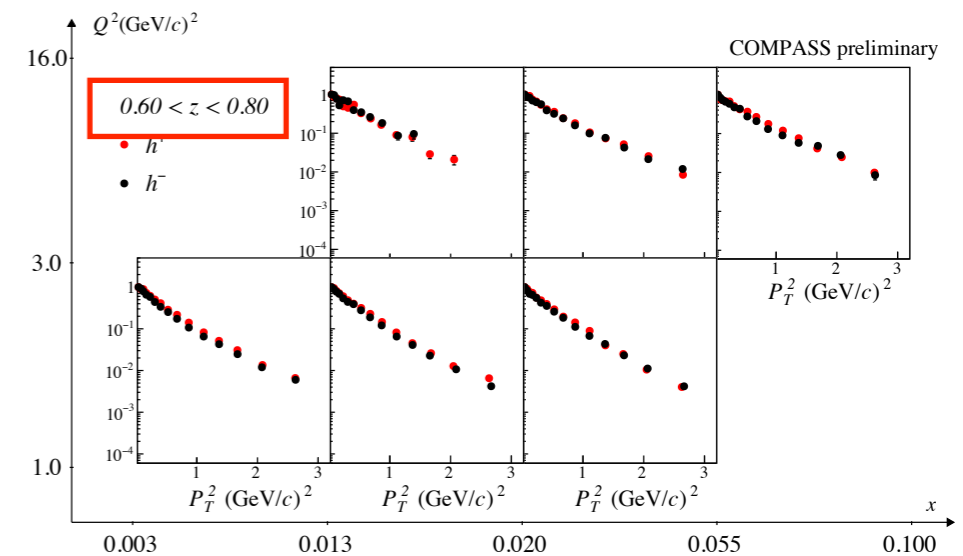
New SIDIS data for unpol. proton target

Compass: 2016-17 run on unpol. LH₂ target; only 11% of data analyzed
 4-D analysis (x , Q^2 , z , P_{hT}) bins; unidentified charged hadrons h^\pm
 QED radiative corrections included
 contamination from exclusive VM decay subtracted bin by bin



P_{hT} -distributions

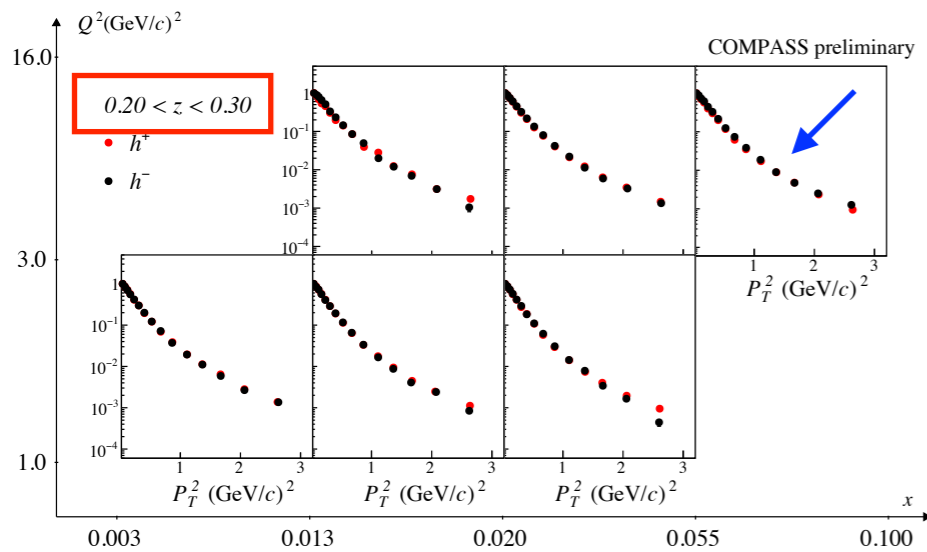
- no evidence of flavor dep.
- clear z , Q^2 , x dep.
- deviation from Gaussian at $P_{hT} > 1$ GeV



talks by A. Bressan
 J. Matousek

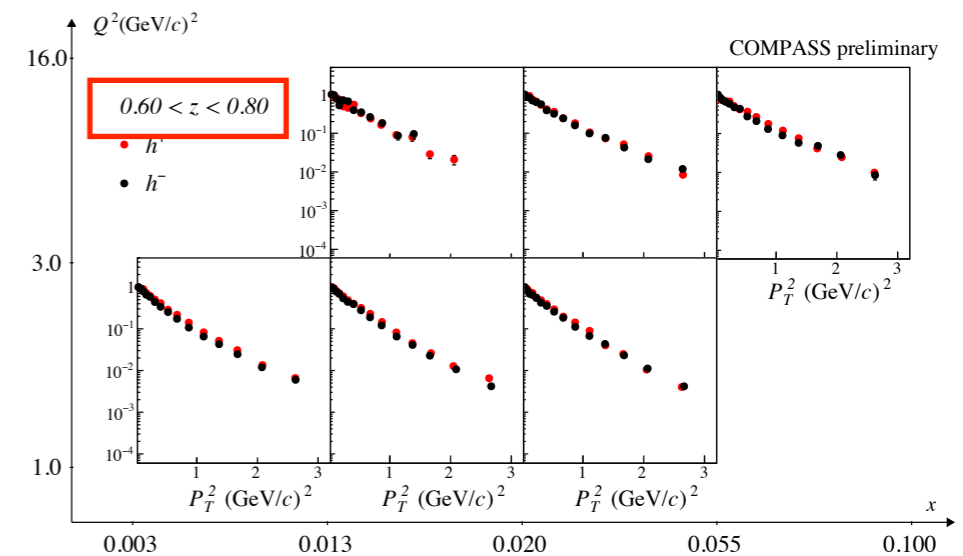
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 contamination from exclusive VM decay subtracted bin by bin

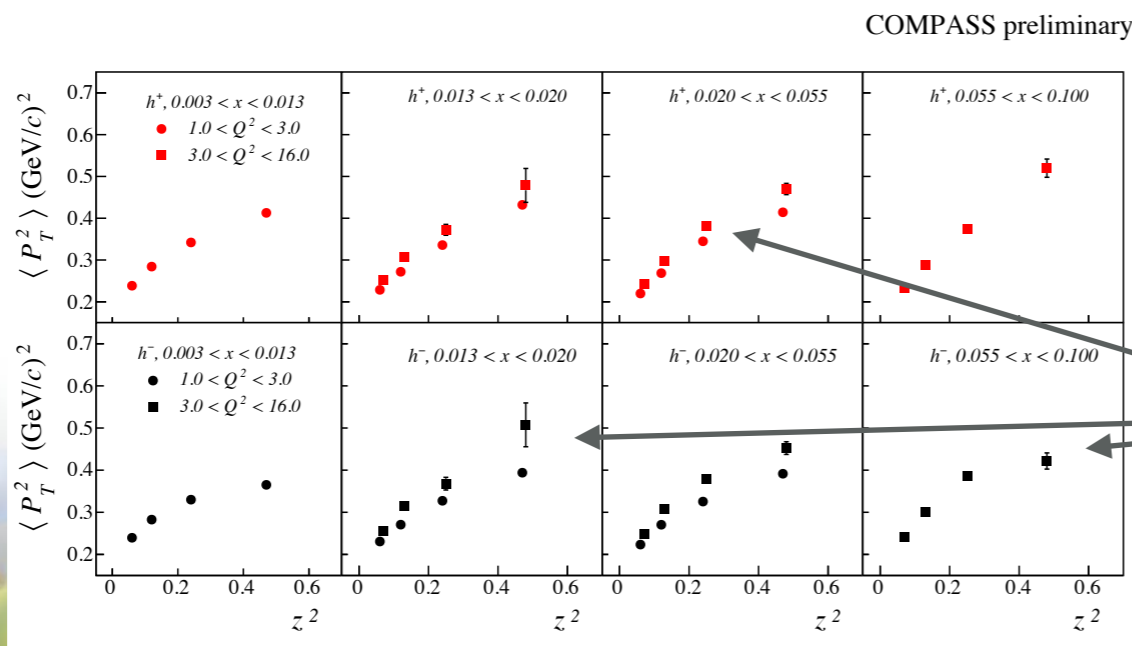


P_{hT} -distributions

- no evidence of flavor dep.
- clear z, Q^2, x dep.
- deviation from Gaussian at $P_{hT} > 1$ GeV



$$\langle P_{hT}^2 \rangle(z^2)$$



talks by A. Bressan
 J. Matousek

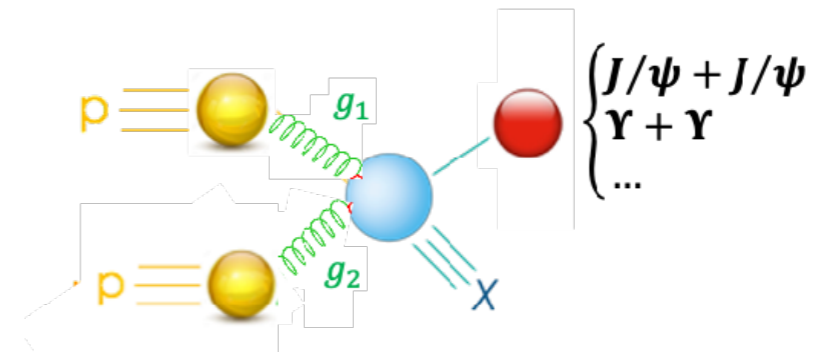
parton model: $\langle P_{hT}^2 \rangle = z^2 \langle k_\perp^2 \rangle + \langle P_\perp^2 \rangle$

deviations : $\langle k_\perp^2 \rangle(x) ?$
 $\langle P_\perp^2 \rangle(z) ?$

Future data for unpol. gluon TMDs

- LHCspin => ex. $pp^{(\uparrow)} \rightarrow J/\psi + J/\psi + X$

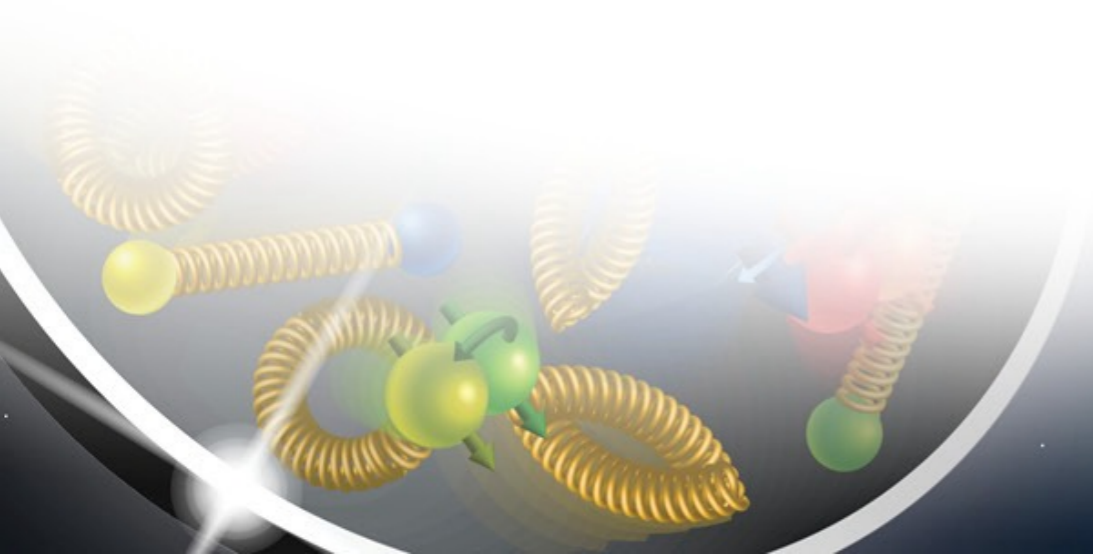
talk by P. Di Nezza



- also complementarity of colliders:

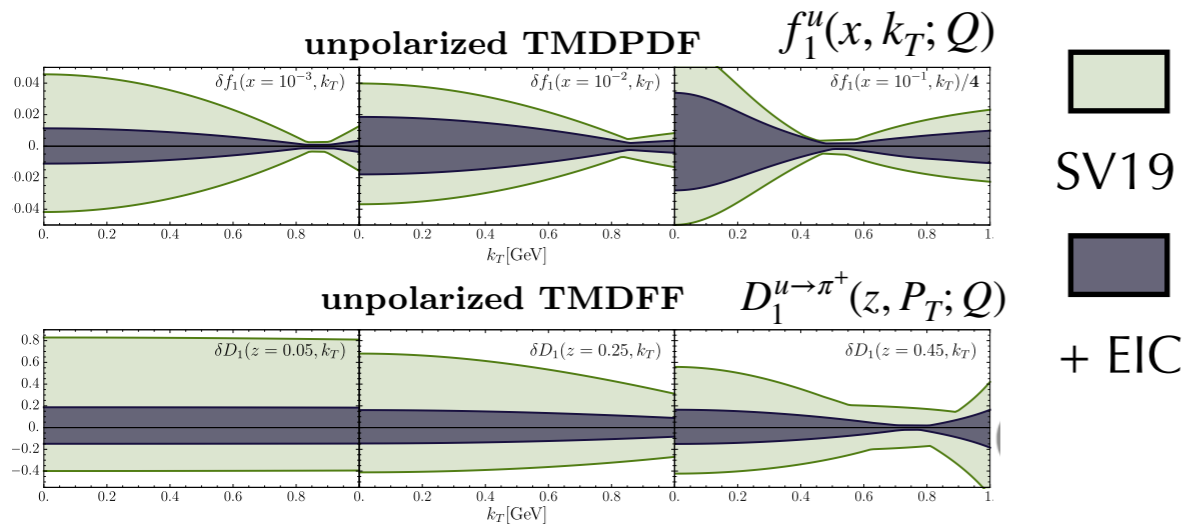
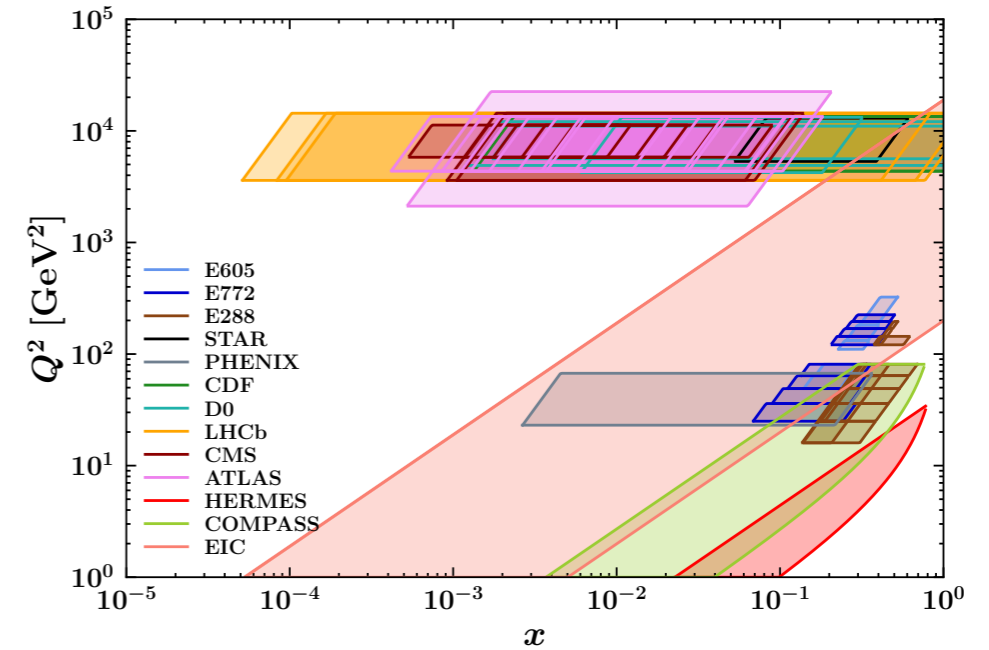
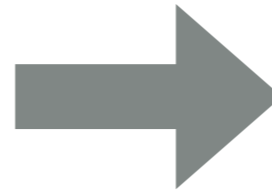
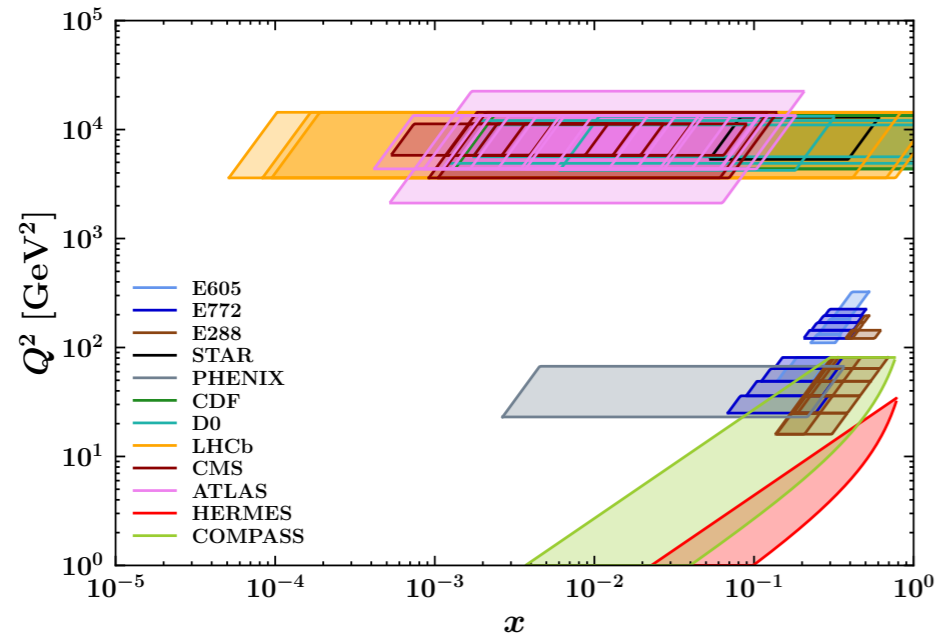
$f_1^g^{[+,+]}$	$pp \rightarrow \gamma J/\psi X$ $pp \rightarrow \gamma \Upsilon X$	LHC LHC
$f_1^g^{[+,-]}$	$pp \rightarrow \gamma \text{jet} X$	LHC & RHIC
$h_1^{\perp g^{[+,+]}}$	$ep \rightarrow e' Q \bar{Q} X$ $ep \rightarrow e' \text{jet jet} X$ $pp \rightarrow \eta_{c,b} X$ $pp \rightarrow H X$	EIC EIC LHC & NICA LHC
$h_1^{\perp g^{[+,-]}}$	$pp \rightarrow \gamma^* \text{jet} X$	LHC & RHIC

Boer, talk at IWHSS2020



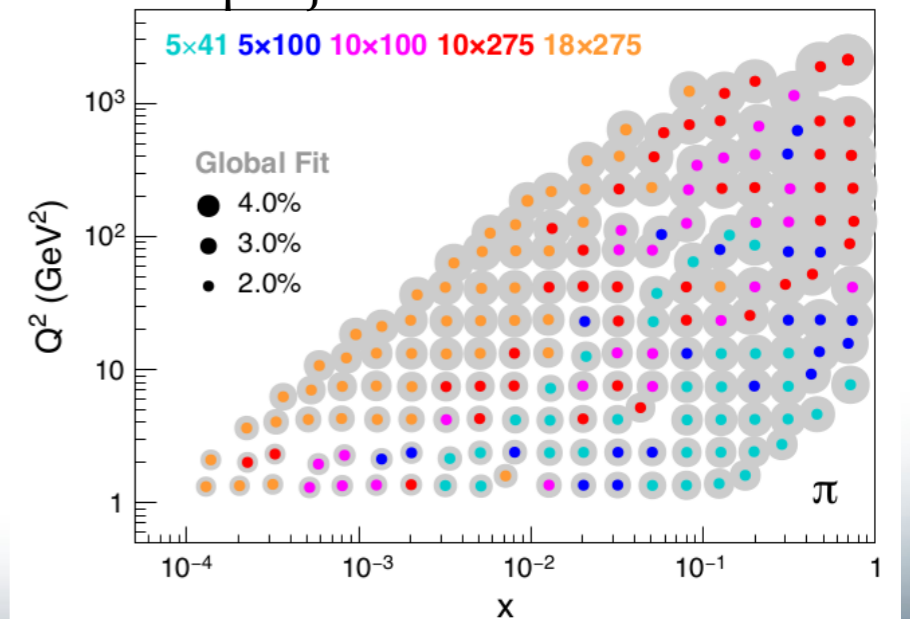
The EIC impact

MAPTMD22 coverage



Abdul Khalek et al., arXiv:2103.05419, N.P.A in press

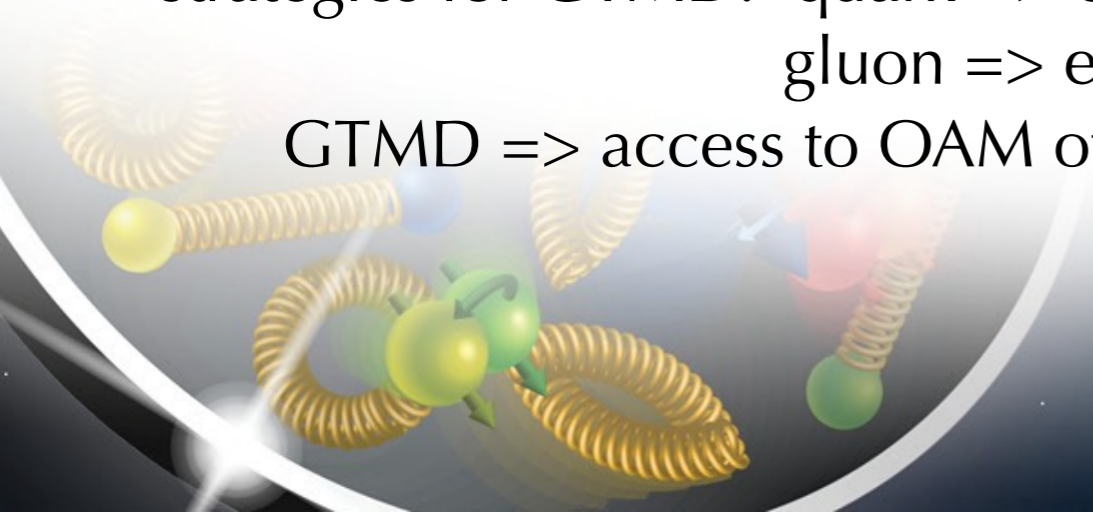
projected errors for $d\sigma^0$



Adam et al., ATHENA Coll.

More stuff ...

- unpolarized azimuthal asymm.: 3-D analysis of $A_{UU}^{\cos \phi}$, $A_{UU}^{\cos 2\phi}$ from Compass
also for di-hadron final state talk by A. Moretti
- twist-3 beam spin asymm. (BSA): $A_{LU}^{\sin \phi}$ from Compass and Hermes
contains $e(x, k_{\perp}) \otimes H_1^{\perp}(z, P_{\perp})$ talks by A. Moretti
G. Schnell
- twist-3 BSA: $A_{LU}^{\sin \phi}$ from CLAS(6+12) with di-hadron final state
contains $e(x) H_1^{\perp}(z, M_{\pi\pi})$
+ decomposition of di-hadron FF in partial waves talks by C. Dilks
A. Courtoy
- JLab BSA with 2 back-to-back hadrons: first evidence of Fracture Funct. talks by T. Hayward
F. Benmokhtar
- exclusive processes for GPD extraction talks by Dupre', d'Hose, Hobart, Kumericki, Sznajder
- strategies for GTMD: quark => exclusive double DY
gluon => exclusive di-jet in (pol.) e-p at the EIC
GTMD => access to OAM of quarks and gluons talks by S. Bhattacharya
F. Yuan



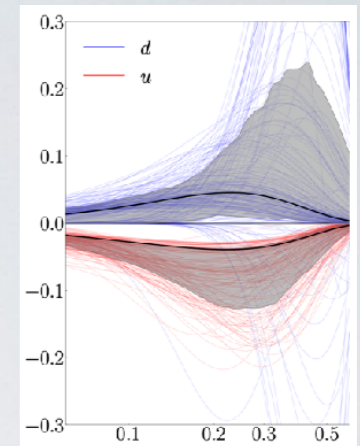


THANK YOU
for your
ATTENTION!

Backup

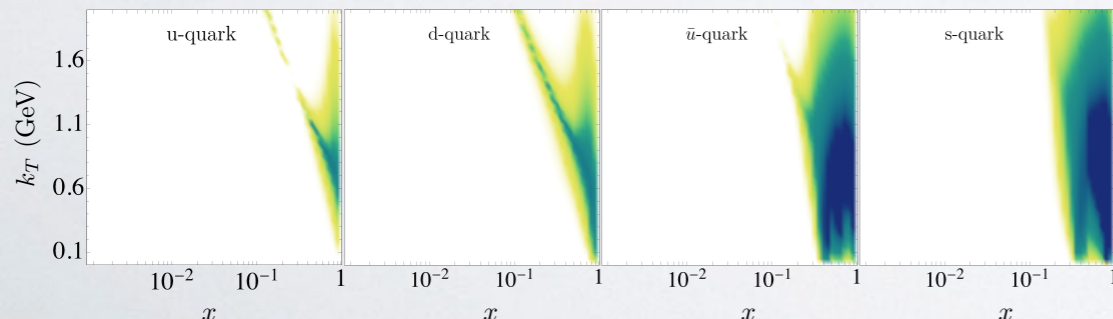
Remarks on Sivers extractions

- Most fits use all correlated projections of same data set; moreover, ETK20 artificially enhance weight of STAR data by factor 13, still getting tension between STAR and SIDIS data ($\chi^2/N_{pts} = 1.44$)



- JAM20 and TO-CA use Generalized Parton Model (GPM) with no TMD evolution and incompatible with Sivers sign change SIDIS-DY; TO-CA use GPM and version CGI-GPM (compatible with sign change), but they get better χ^2 with GPM

- Hard to compare BPV20 with rest of works in CSS formalism; in any case, there are violations of positivity bound for sea quarks at large x



coloured areas

Bury et al., arXiv:2103.03270

