

# The rates of neutron-releasing reactions in He-burning phases and their astrophysical consequences

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1. Slow neutron captures in stars: a general reminder
2. The neutron sources and their activation
3. s-Processing from the  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  reaction: status & needs for new measures
4. s-Processing from the  $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$  reaction: status & needs for new measures

# The distribution of s-elements: two components

830

KÄPPELER *ET AL.* (1989)

Vol. 257

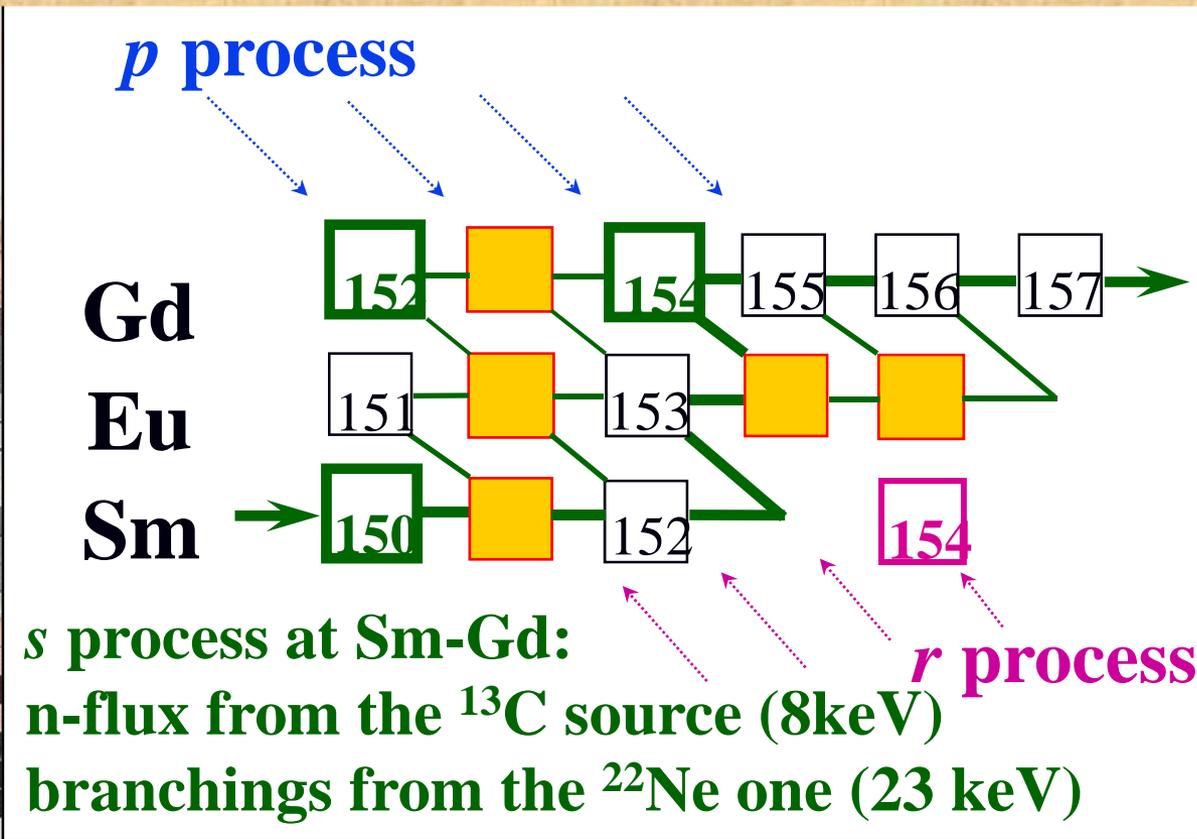
**Main component: TP-AGB Stars of Low & Intermediate Mass (Iben & Truran 1977; Gallino et al. 1998; Busso et al. 1999): repeated neutron fluences (originally imagined as exponentially distributed, they are actually more complex in nature).**

**Weak component: Massive stars during He and C burning (Raiteri et al. 1991 a,b; Kaeppler et al. 1994; Pignatari et al. 2010)**

55 70 90 110 130 150 170 190 210  
MASS NUMBER

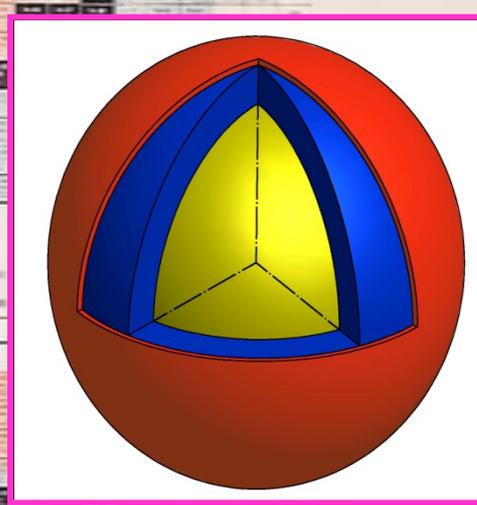
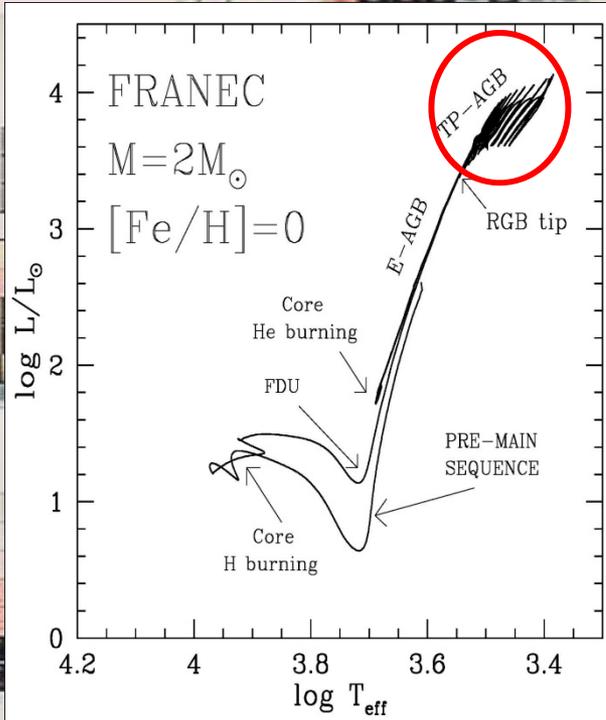
FIG. 2.—The product of s-process abundance times cross section as a function of mass number. The symbols correspond to empirical values for s-only isotopes (squares) or to neutron magic isotopes which are predominantly produced by the s-process (circles). The respective abundances are taken from the solar abundance table of Cameron (1981). Error bars include the cross section uncertainties only. The calculated solid lines correspond to the strong and weak component in the exponential neutron fluence distribution.

# Neutron density and neutron source from the main component



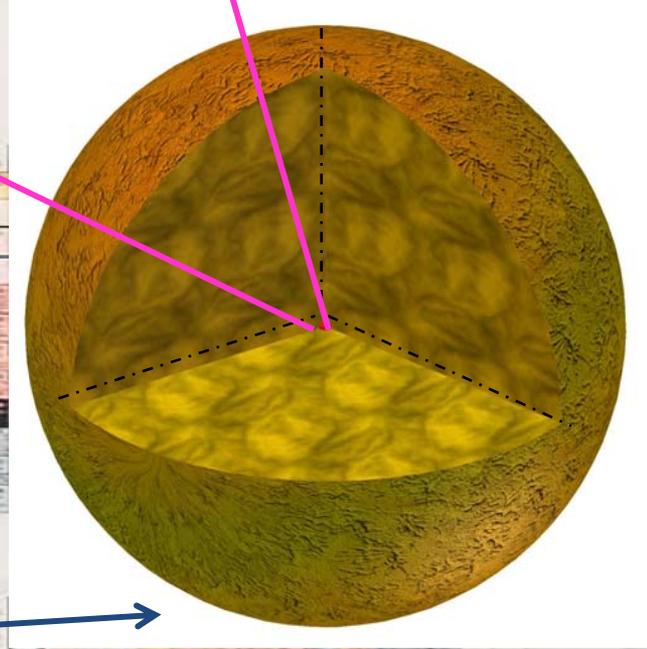
Precise cross sections, decay rate estimates near branching points and comparison with stellar models made clear that the main component is dominated by a regime with low  $n$ -density ( $^{13}\text{C}(a,n)^{16}\text{O}$  reaction, at about 8 keV,  $n_n = 10^7 \text{ cm}^{-3}$ ), with small "corrections" from the marginal activation of the  $^{22}\text{Ne}(a,n)^{25}\text{Mg}$  source at about 22-23 keV, with a higher neutron density, up to  $10^{10} \text{ cm}^{-3}$ .

# AGB stars



Earth-Sun  
( $\sim 200 R_{\text{SUN}}$ )

Detailed description: A blue double-headed arrow pointing left and right, indicating the scale of the Earth-Sun distance, which is approximately 200 solar radii ( $R_{\text{SUN}}$ ).



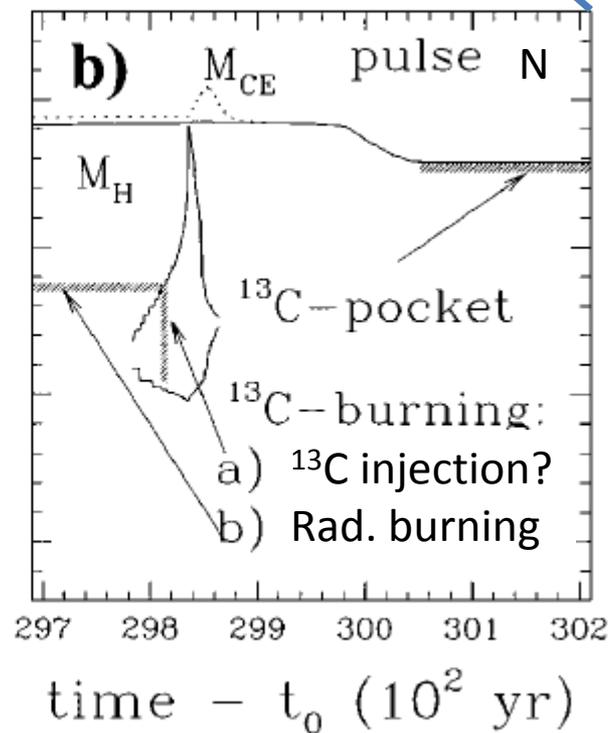
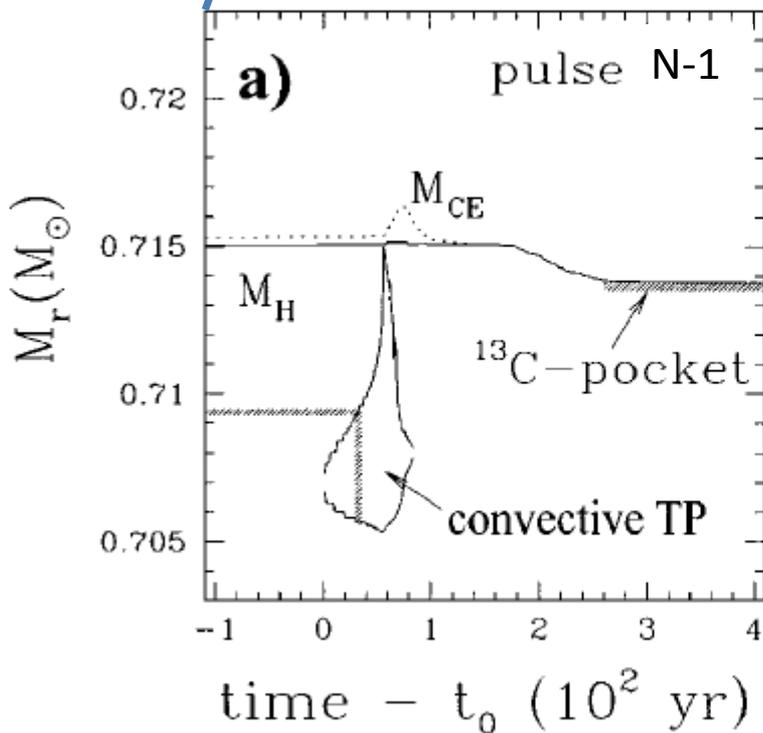
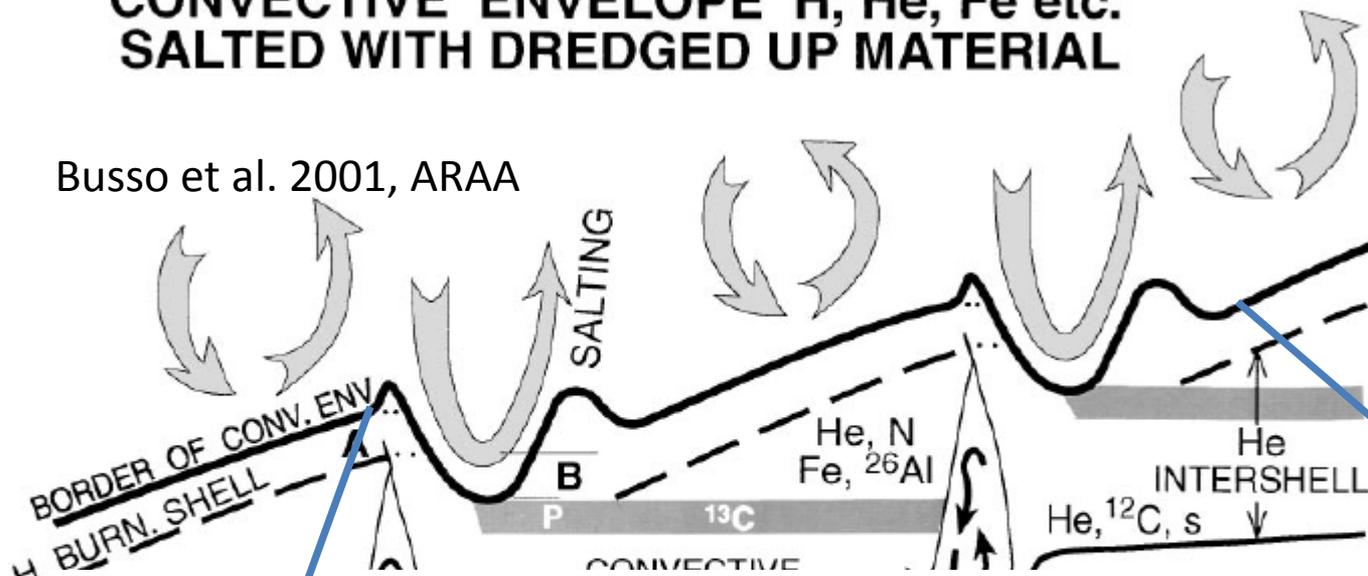
## MAIN REFERENCES:

- Chieffi et al. 1998
- Straniero et al. 2005
- Cristallo et al 2009

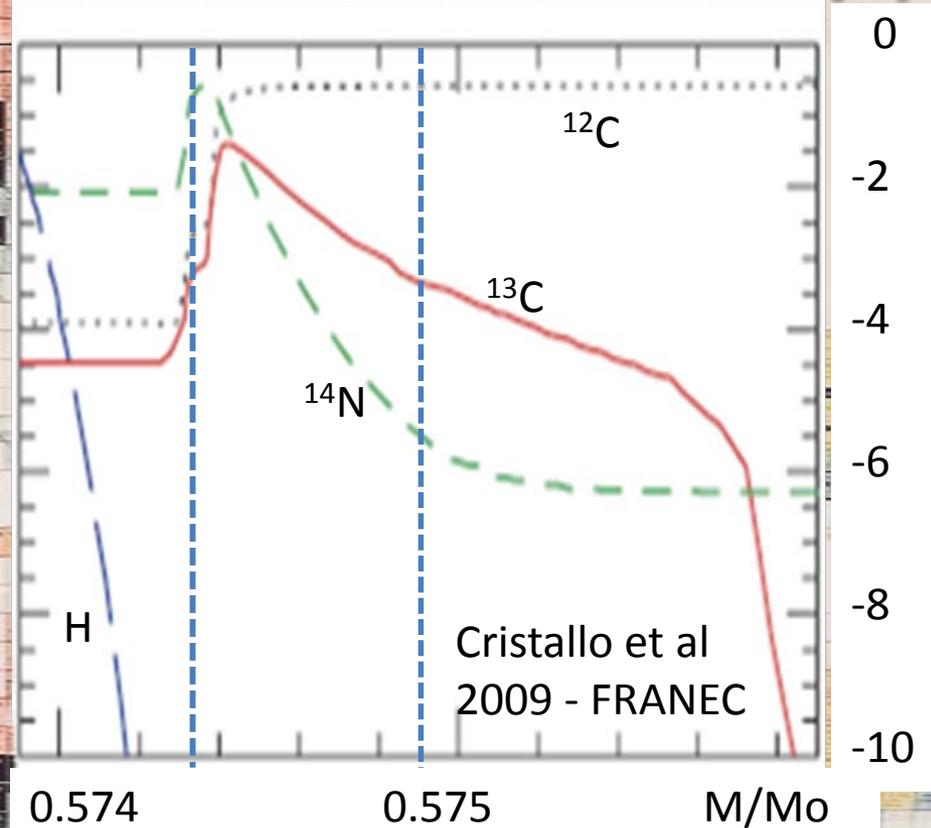
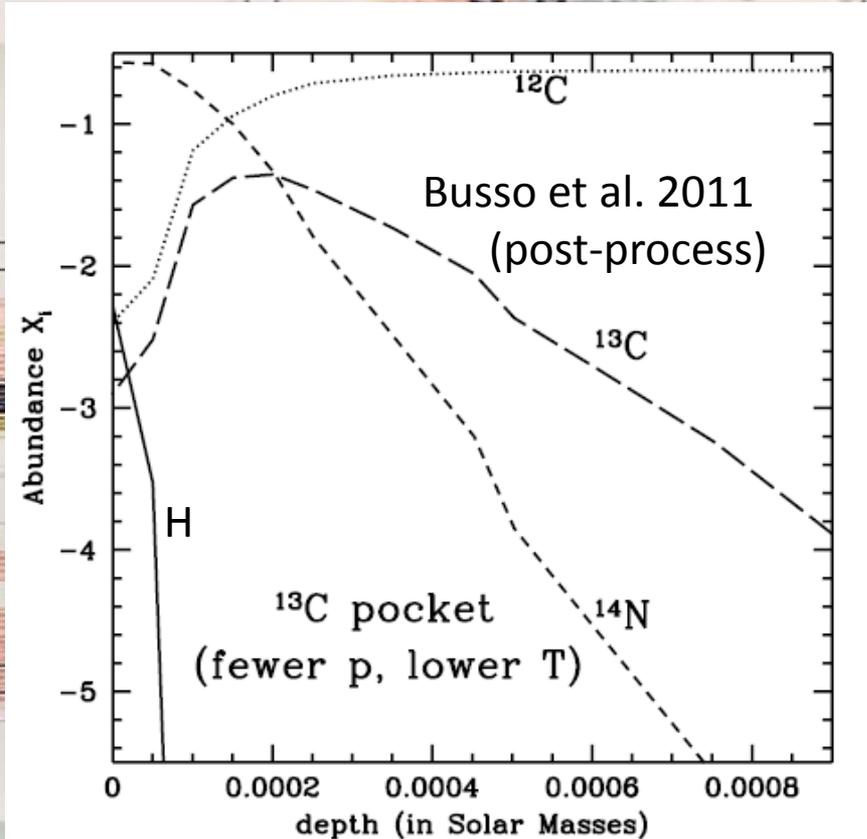
# CONVECTIVE ENVELOPE H, He, Fe etc. SALTED WITH DREDGED UP MATERIAL

Busso et al. 2001, ARAA

MASS COORDINATE



# Current ways of modelling the $^{13}\text{C}$ neutron source formation



Obtained through  $n(\text{H}) = n(\text{H}_{\text{bce}})\exp(-k\Delta M)$

(bc) In any case, the profile of protons is in general exponential-like, as expected (more or less) in a diffusive phenomenon

FIRST PROBLEM: if the subthreshold-resonance contribution to the rate of the  $^{13}\text{C}+\alpha$  is much smaller than assumed (Kato et al. 2005; Kubono et al. 2003) then we might need more time to burn  $^{13}\text{C}$  than available in the interpulse

→  $^{13}\text{C}$  would burn partly convectively, with all the inherent complications (extra energy, shell splitting, possibly higher neutron densities.....).

THIS POINT SHOULD BE VERIFIED CLEARLY  
(Heil et al. 2008 go toward opposite directions).  
REAL MEASURES, NOT EXTRAPOLATIONS, NEEDED

In the rare, but important cases in which Proton Ingestion Episodes (PIE: Critsallo et al. 2009b) occur, the structure is heavily affected, the convective shell is splitted,  $^{13}\text{C}$  forms over a wide range of temperatures: better knowledge of the rate is needed from few keV up to about 30 keV.

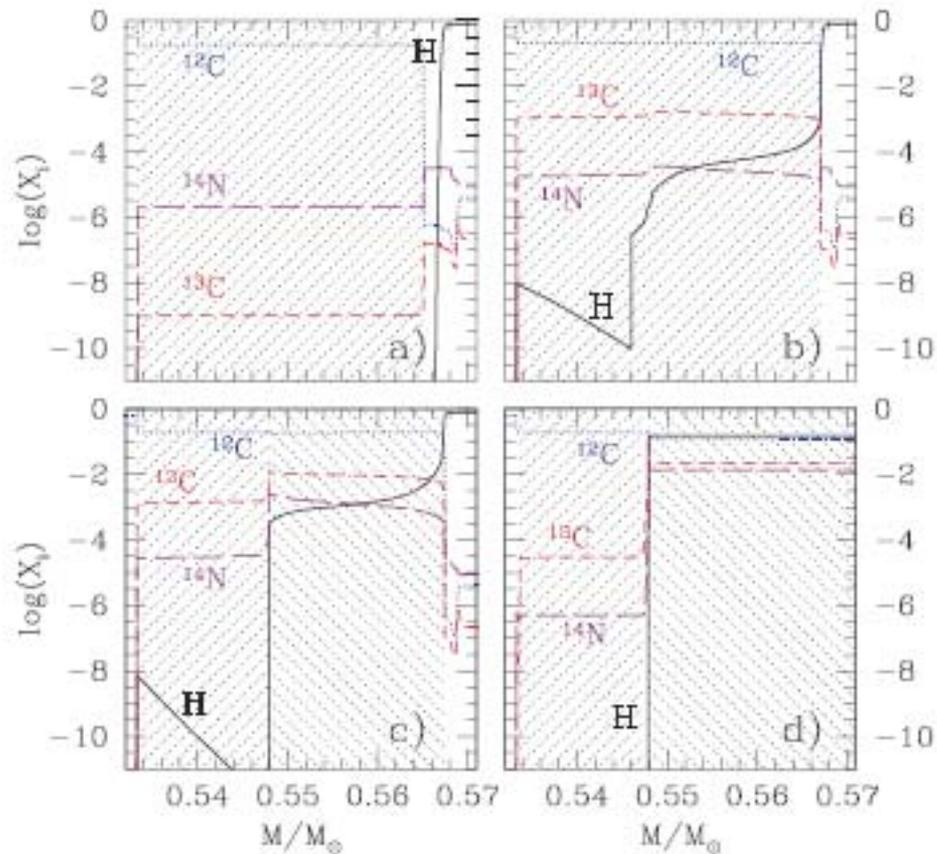
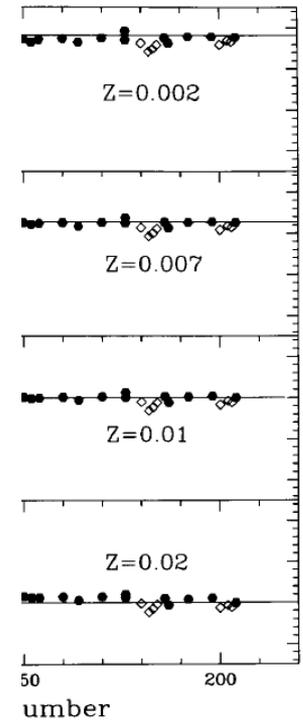
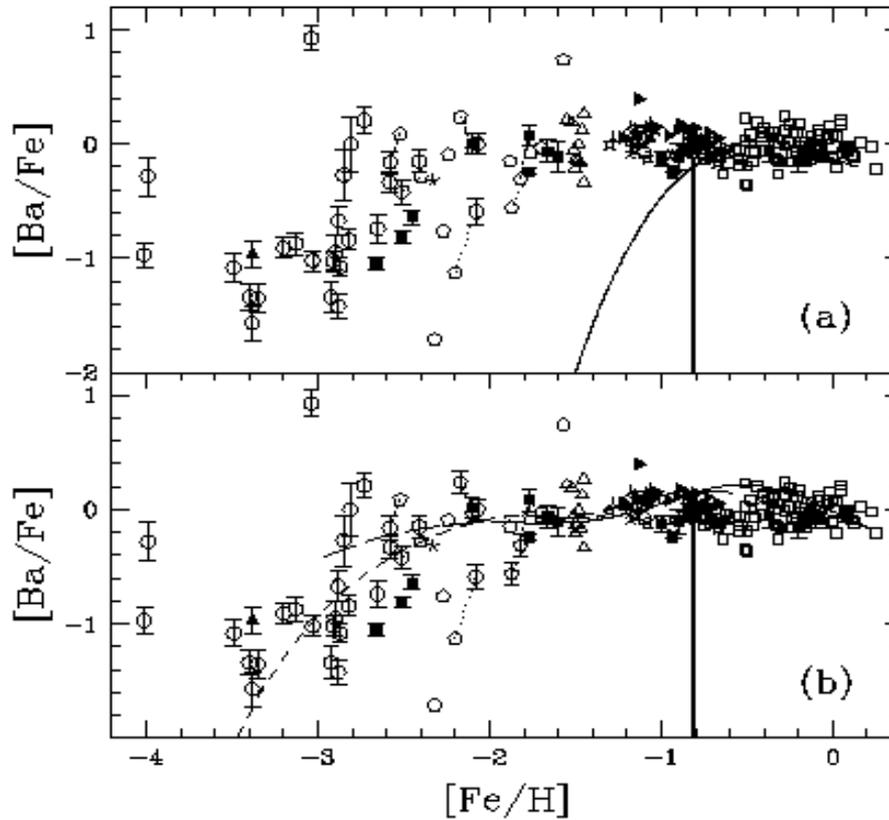
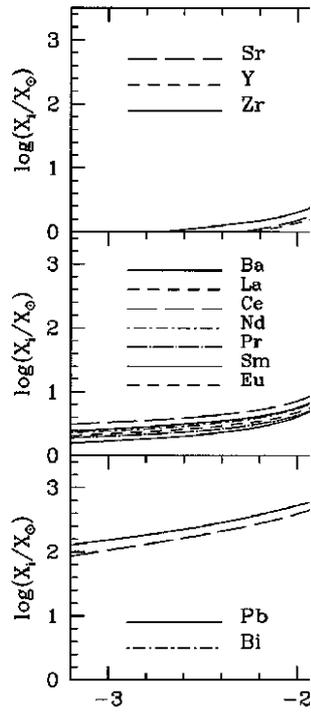


Figure 2: Evolution of selected key isotopes within the He-intershell during the PIE.

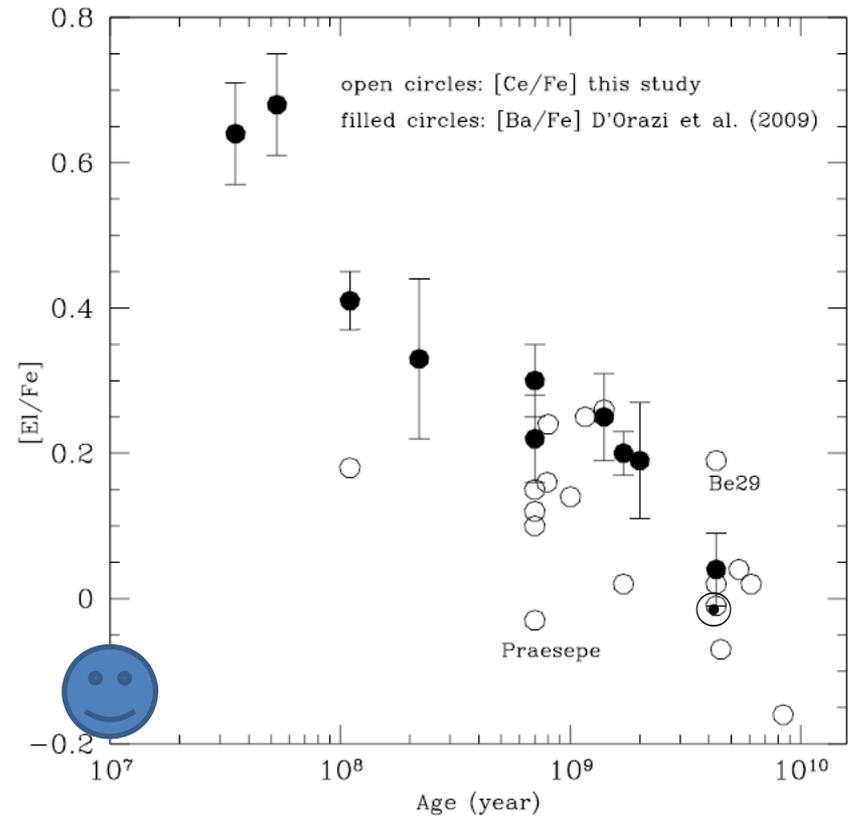
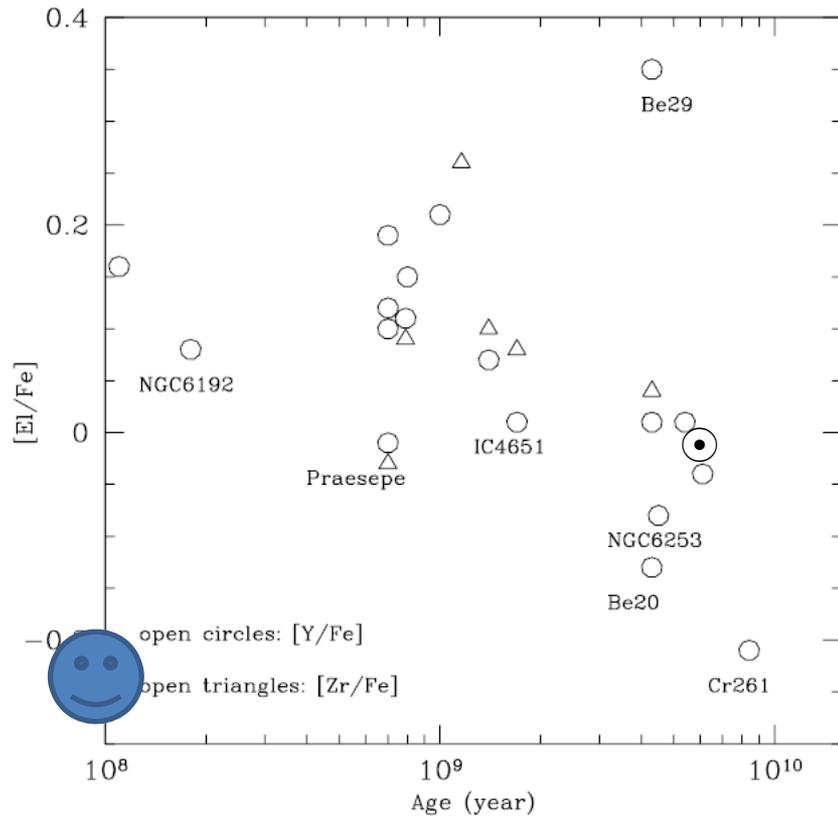
# Average s-process increases in the Galaxy



Travaglio et al. 1999

# A new problem

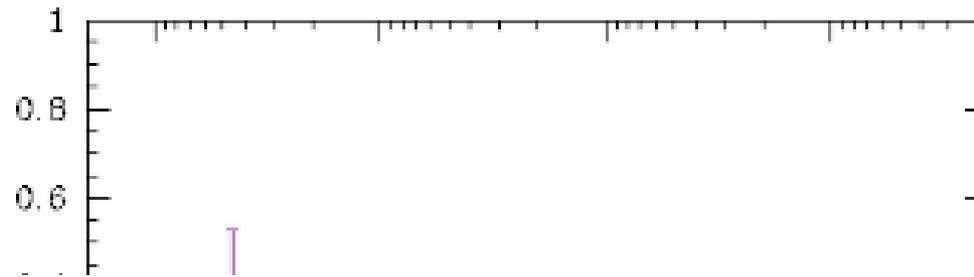
We have just completed a survey of s-process observations in the young Galaxy



First hints in D'Orazi et al. 2009 from Ba. But now we have 5 elements at the two s-process abundance peaks, IN 23 CLUSTERS!!

For young clusters the traditional picture of s-processing fails

If we expand the  
processing in old



much s-  
abundances

## **SECOND PROBLEM, FOR INTERPRETING THE YOUNG CLUSTERS:**

If we need a more extended  $^{13}\text{C}$  pocket in very low mass stars ( $1.2-1.4 M_{\odot}$ ), with a relatively low  $^{13}\text{C}$  abundance, but extending to higher T ( $10-11 \text{ keV}$ ) would this  $^{13}\text{C}$  burn in the radiative region or remain partially unburnt up to the pulse?

Lower  $^{13}\text{C}$  abundance at higher T (easily consumed): here we might have less problems if the rate is smaller. But what if it is larger? Then higher neutron density, change of the branching analysis!?

You need very large pockets, with low  $^{13}\text{C}$  abundances, extending to where T is higher (more than the usual  $8 \text{ KeV}$ ). A precise rate for  $^{13}\text{C}+\alpha$  is crucial: we would burn it in the presence of a T gradient.

[Such pockets would need non-diffusive mixing, i.e. by wave instabilities]

# The $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction rate

Addressed theoretically and experimentally by numberless authors since the Caughlan and Fowler (1988) compilation



**FIRST MOTIVATION:** verify that there are no increases w. respect to K94.

This would put on solid grounds calculations that were so far based on speculations

On the contrary, a strong increase would need a revision of the s-process in AGB stars, mainly going toward lower masses, with lower  $T$  ( $< 23$  keV) in the thermal pulses.

But for these stars, at solar metallicity, nobody can find dredge-up. How can we reconstitute the s-elements to the Galaxy??

Changes in the stellar models??

Temperature (GK)

Fig. A.22. Previous reaction rates: Ref. [22]

# Production of $^{25}\text{Mg}$ in the envelope of a 5 $M_{\odot}$ Star for different rates

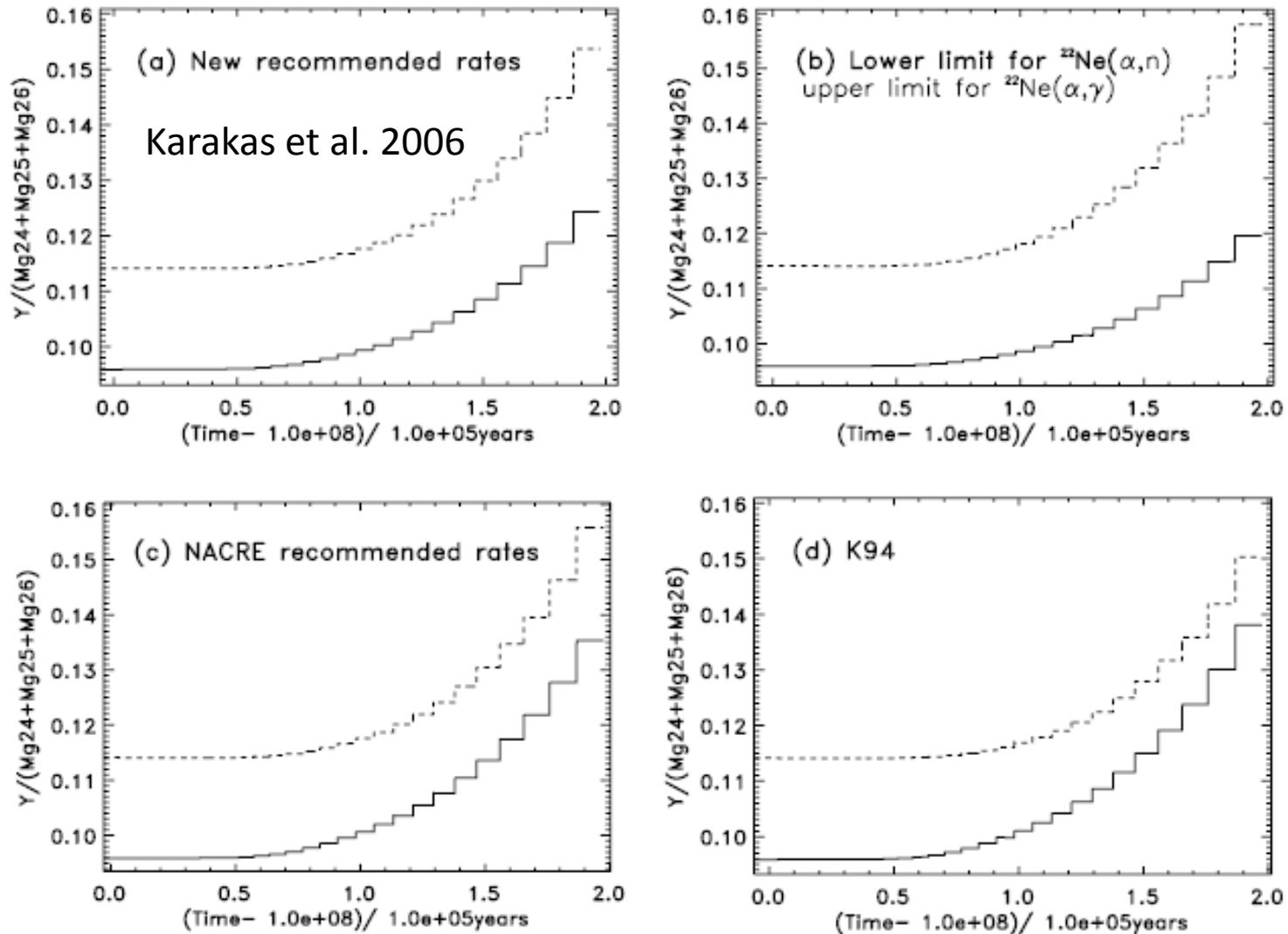
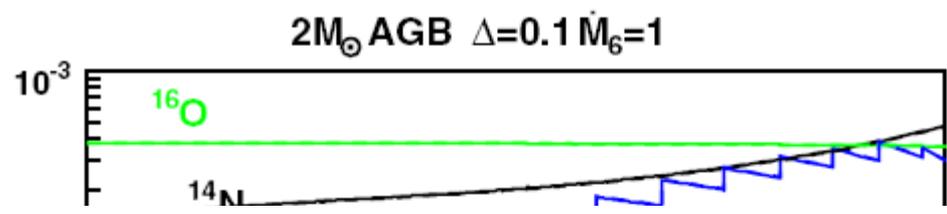


Fig. 5.—Evolution of the  $^{25}\text{Mg}$  (solid lines) and  $^{26}\text{Mg}$  (dashed lines) abundances at the surface of the  $5 M_{\odot}$ ,  $Z = 0.02$  VW93 model for four different choices of the  $^{22}\text{Ne} + \alpha$  reaction rates.

# Another problem: O-Na anticorrelations in GC

In the nearly 100 recent years, (



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AND I DID NOT SPEAK OF MASSIVE STARS, I.E. OF THE WEAK COMPONENT OF THE S-PROCESS, PRODUCED BY THE  $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$  REACTION. (PIGNATARI ET AL. 2010)

The of t

THERE He- & C BURNING CONTRIBUTE: RATE CHANGES WOULD IMPLY A DIFFERENT BALANCE BETWEEN THE TWO AND A DIFFERENT COMPETITION WITH  $^{12}\text{C}+\alpha$ .

field

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THE WEAK COMPONENT WOULD CHANGE, HENCE YOU WOULD NEED CHANGES ON THE MAIN ONE ( $^{13}\text{C}$ ).....

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Ne.

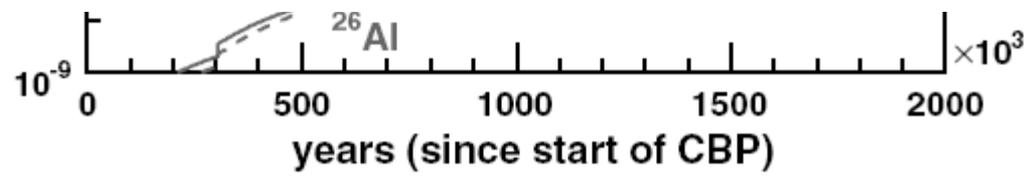
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THIS IS A VERY IMPORTANT ISSUE, BUT WOULD REQUIRE ANOTHER TALK!

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But The

& Lagarde 2011 second-generation deep mixing phenomena.



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# Conclusions

$^{13}\text{C}+\alpha$

1. Changes in the rate might on one side specify whether the situations in which  $^{13}\text{C}$  burns partly convectively are rare (as now believed) or more common.  
This risk would be real for a remarkable REDUCTION of the rate (by more than

## HOW TO SAVE THE CONSERVATIVE SCENARIO:

1. NO CHANGE, OR MODERATE DECREASE, IN THE  $^{22}\text{Ne}+\alpha$  W. RESPECT TO THE PRESENT RATE.
2. ONLY MODERATE CHANGES FOR  $^{13}\text{C}+\alpha$ , NO REDUCTION BY MORE THAN A FACTOR OF THREE AS COMPARED TO NACRE.

GCs.

3. In Massive Stars, changes in the rate  $\rightarrow$  change in the balance between He- and C-burning for the weak component, change in the competition with  $^{12}\text{C}+\alpha$ .
4. More  $^{22}\text{Ne}$  saved? Maybe explanation of the Ne excess of planetary nebulae!!