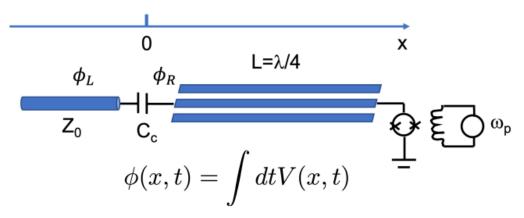
Modeling of JPA

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Aim: modelling of a JPA without simplifications of the nonlinear components

Model by Claudio Gatti



The Lagrangian density of the circuit in figure 1 is obtained by the sum of the capacitive energies minus the inductive energies:

$$\mathcal{L} = \left(\frac{1}{2}C_L\dot{\phi_L}^2 - \frac{1}{2L_L}\phi_L'^2\right)\sigma(-x) + \left(\frac{1}{2}C_R\dot{\phi_R}^2 - \frac{1}{2L_R}\phi_R'^2\right)[\sigma(x) - \sigma(x - L)](1) + \left(\frac{1}{2}C_c(\dot{\phi}_L - \dot{\phi}_R)\right)\delta(x) + \left[\frac{C_J}{2}\dot{\phi}_J^2 + E_J(\phi_{ext})\cos\left(2\pi\frac{\phi_J}{\phi_0}\right)\right]\delta(x - L)$$

where $\phi_{L,R}(x)$ is the flux variable defined as $\phi(x) = \int_{-\infty}^{t} dt' V(t')$ in the left (right) transmission line and $C_{L,R}$ and $L_{L,R}$ are their capacitances and inductances per unit length, and $\phi_J = \phi_R(x = L)$. $\sigma(x)$ is the Heaviside (step) function. The (symmetric) SQUID Josephson energy is:

$$E_J(\phi_{ext}) = 2E_{0J}\cos\left(\pi\frac{\phi_{ext}}{\phi_0}\right) \tag{2}$$

where E_{0J} is the Josephson energy of the single junction.

Wave equations for ϕ_L and ϕ_R

$$\begin{split} \sigma(-x)\left[C_L\ddot{\phi}_L - \frac{1}{L_L}\phi_L''\right] + \delta(x)\left[C_c(\ddot{\phi}_L - \ddot{\phi}_R) + \frac{1}{L_L}\phi_L'\right] &= 0 (4) \\ \left[\sigma(x) - \sigma(x-L)\right]\left[C_R\ddot{\phi}_R - \frac{1}{L_R}\phi_R''\right] + \delta(x)\left[-C_c(\ddot{\phi}_L - \ddot{\phi}_R) - \frac{1}{L_R}\phi_R'\right] + \\ \delta(x-L)\left[C_J\ddot{\phi}_R + \frac{1}{L_R}\phi_R' + \frac{2\pi}{\phi_0}E_J(\phi_{ext})\sin\left(2\pi\frac{\phi_R}{\phi_0}\right)\right] &= 0 \end{split}$$

The solution is the sum of an outgoing and an incoming wave

and $v_p = 1/\sqrt{LC}$. The other equations give the boundary conditions: $\phi(x,t) = \phi^{in}(t - x/v_n) + \phi^{out}(t + x/v_n)$

$$\phi' = \frac{1}{v_p} \left(\dot{\phi}^{out} - \dot{\phi}^{in} \right) \qquad (x = 0)$$

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$$\dot{\phi} = \dot{\phi}^{in} + \dot{\phi}^{out} \qquad C_J \ddot{\phi}_R + \frac{1}{L_R} \phi'_R + \frac{2\pi}{\phi_0} E_J(\phi_{ext}) \sin\left(2\pi \frac{\phi_R}{\phi_0}\right) = 0 \qquad (x = L)$$

$$2\ddot{\phi}_R^{in} + \frac{1}{CcZ_0}\dot{\phi}_R^{in} = \ddot{\phi}_L^{in} + \frac{1}{CcZ_0}\dot{\phi}_R^{out} \qquad (x=0)$$

$$C_{J}\ddot{\phi}_{J} + \frac{1}{Z_{0}}\dot{\phi}_{J} + \frac{2\pi}{\phi_{0}}E_{J}(\phi_{ext})\sin\left(2\pi\frac{\phi_{J}}{\phi_{0}}\right) = \frac{2}{Z_{0}}\dot{\phi}_{R}^{in} \qquad (x = L) \qquad \phi_{R}^{in}(t, x = L) = \phi_{R}^{in}(t - L/v_{p}, x = 0)$$

$$\phi_{R}^{out}(t, x = L) = \phi_{R}^{out}(t + L/v_{p}, x = 0) = \phi_{J}(t) - \phi_{R}^{in}(t, x = L)$$

$$= \phi_{J}(t) - \phi_{R}^{in}(t - L/v_{p}, x = 0)$$

$$\phi_{L}^{out} = \phi_{R}^{out} - \phi_{R}^{in} + \phi_{L}^{in} \qquad (x = 0)$$

$$\phi_R^{out}(t, x = L) = \phi_R^{out}(t + L/v_p, x = 0) = \phi_J(t) - \phi_R^{in}(t, x = L)$$
$$= \phi_J(t) - \phi_R^{in}(t - L/v_p, x = 0)$$

$$\phi_L^{out} = \phi_R^{out} - \phi_R^{in} + \phi_L^{in} \qquad (x = 0)$$

Transient Analysis of Lossless Transmission Lines

Manuscript received July 7, 1967; revised September 8, 1967.

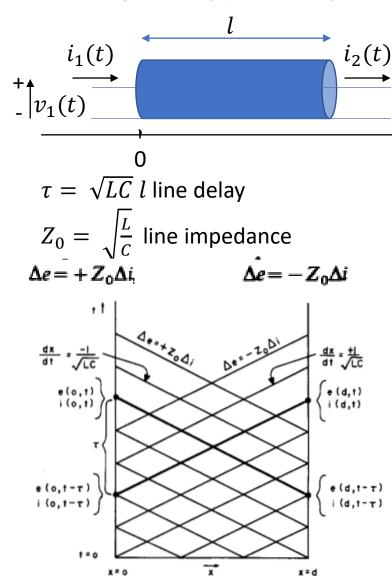


Fig. 1. Characteristic curves for a uniform transmission line.

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 $v_2(t)$

$$v = v(x,t) \ i = i(x,t)$$

$$L \frac{\partial i}{\partial t} + \frac{\partial v}{\partial x} = 0 \qquad i_{tt} - \frac{1}{LC} i_{xx} = 0$$

$$C \frac{\partial v}{\partial t} + \frac{\partial i}{\partial x} = 0 \qquad v_{tt} - \frac{1}{LC} v_{xx} = 0$$

The method of characteristics for solving these equations is based on a transformation in the x-t plane which accomplishes the conversion of (1) and (2) into a pair of ordinary differential equations. Each of these two ordinary differential equations holds true along a different family of characteristic curves in the x-t plane, one family corresponding to the forward or incident wave and the other to the backward or reflected wave. These two families of characteristic curves are, in effect, the basis for a new coordinate system instead of the canonical coordinates x = constant and t = constant. $\frac{dx}{dt} = 1/\sqrt{LC} \text{ and } \frac{dx}{dt} = -1/\sqrt{LC}$

$$\sqrt{\frac{L}{C}}di + \left(Ri + \sqrt{\frac{L}{C}}Ge\right)dx + de = 0$$
 (3)

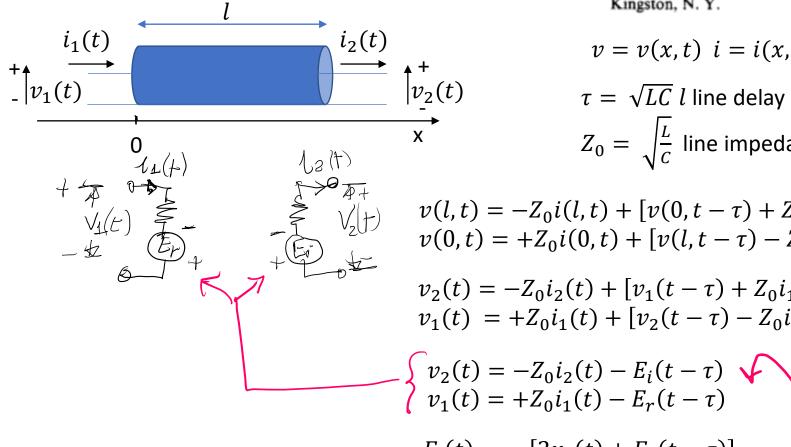
which holds true only along the forward characteristics, defined as $dx/dt = 1/\sqrt{L/C}$, and

$$-\sqrt{\frac{L}{C}}di + \left(Ri - \sqrt{\frac{L}{C}}Ge\right)dx + de = 0$$
 (4)

which holds true only along the backward characteristics, defined as $dx/dt = -1/\sqrt{LC}$.

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$$v = v(x,t)$$
 $i = i(x,t)$ $\tau = \sqrt{LC} \ l$ line delay $Z_0 = \sqrt{\frac{L}{C}}$ line impedance

$$v(l,t) = -Z_0 i(l,t) + [v(0,t-\tau) + Z_0 i(0,t-\tau)]$$

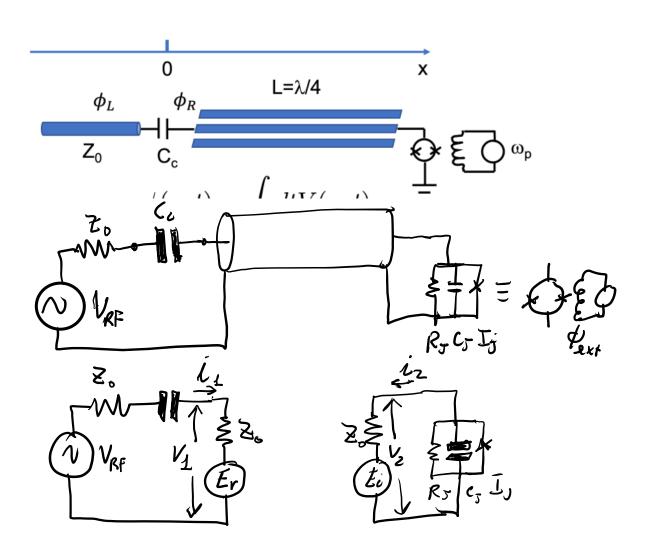
$$v(0,t) = +Z_0 i(0,t) + [v(l,t-\tau) - Z_0 i(l,t-\tau)]$$

$$v_2(t) = -Z_0 i_2(t) + [v_1(t-\tau) + Z_0 i_1(t-\tau)]$$

$$v_1(t) = +Z_0 i_1(t) + [v_2(t-\tau) - Z_0 i_2(t-\tau)]$$

$$E_i(t) = -[2v_1(t) + E_r(t - \tau)]$$

$$E_r(t) = -[2v_2(t) + E_i(t - \tau)]$$



In our case we have:

$$\begin{split} \dot{v}_1 &= -\frac{1}{2 C_c Z_0} v_1 + \frac{1}{2} \dot{v}_{rf}(t) - \frac{1}{2} \dot{E}_r(t - \tau) - \frac{1}{2 C_c Z_0} E_r(t - \tau) \\ \dot{v}_2 &= -\left(\frac{1}{C_j Z_0} + \frac{1}{C_j R_j}\right) v_2 - \frac{2 I_j}{C_j} \cos\left(2\pi \frac{\phi_{ext}}{\phi_0}\right) \sin \varphi - \frac{1}{C_j Z_0} E_i(t - \tau) \\ \dot{\varphi} &= \frac{2 e}{\hbar} v_2 \end{split}$$

where

$$\dot{v}_{rf}(t) = v_{rf} \omega \cos \omega t$$

$$\phi_{ext}(t) = \phi_{ext} \cos \omega_p t$$

$$Z_0 = 50 \Omega$$

$$\tau = \frac{T}{4} = \frac{\pi}{2 \,\omega}$$

We start with Ei and Er set to zero

The values of Ei and Er are updated during the integration according to:

$$E_i(t) = -[2v_1(t) + E_r(t - \tau)]$$

$$E_r(t) = -[2v_2(t) + E_i(t - \tau)]$$

Next step is to integrate the system and do some checks with published data