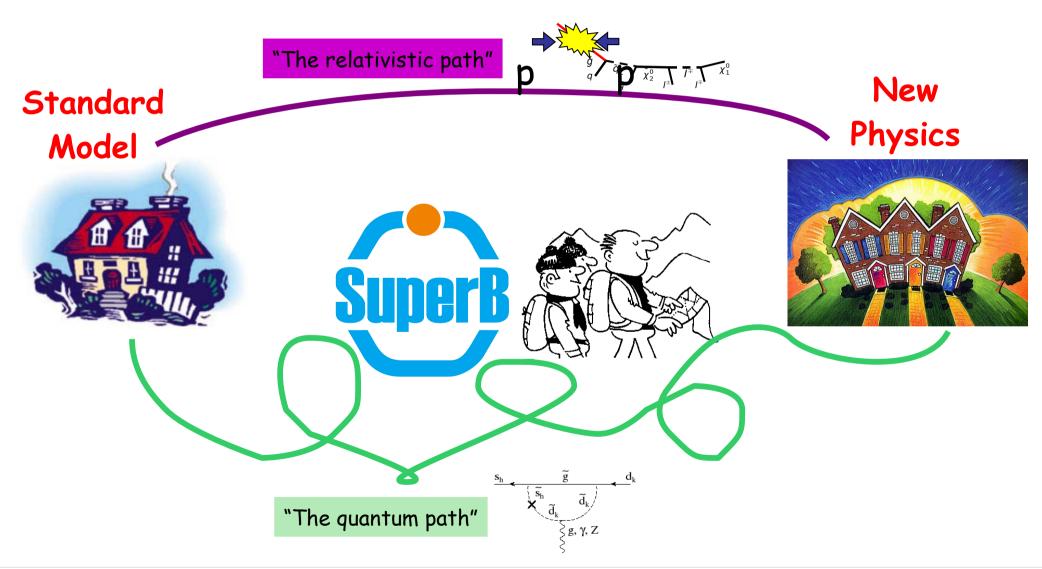
SuperB New Physics

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B physics observables for New

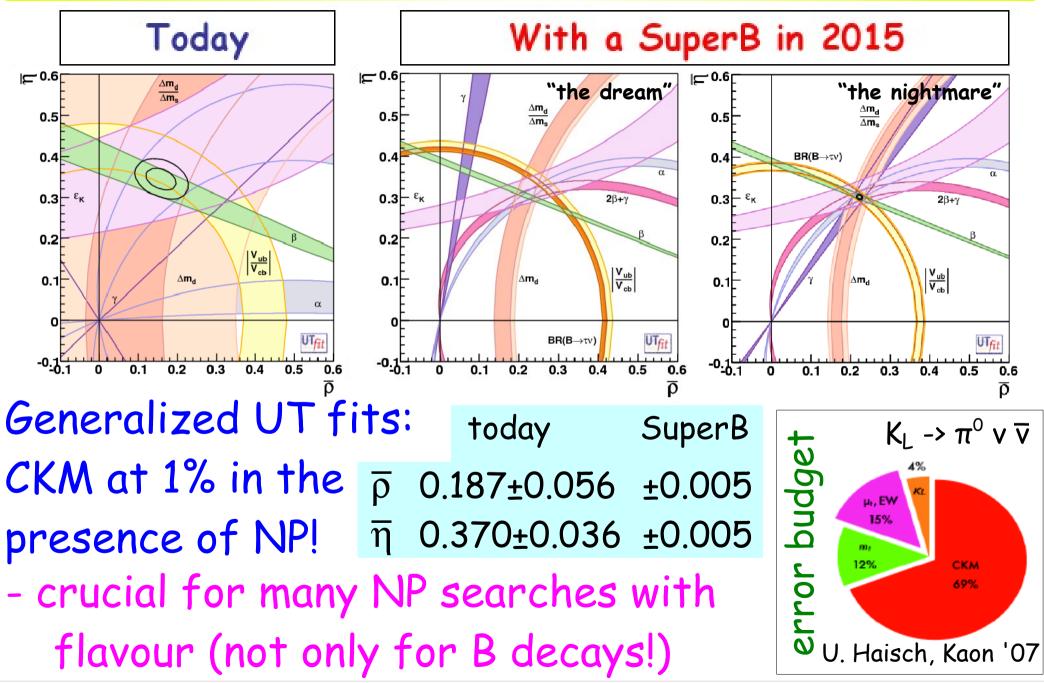
Observable	B factories (2 ab^{-1})	$\operatorname{Super} B$ (75 ab^{-1})
$sin(2\beta) (J/\psi K^0)$	0.018	0.005 (†)
$cos(2\beta) (J/\psi K^{*0})$	0.30	0.05
$sin(2\beta)$ (Dh ⁰)	0.10	0.02
$\cos(2\beta)$ (Dh ⁰)	0.20	0.04
$S(J/\psi \pi^0)$ $S(D^+D^-)$	0.10	0.02
$S(D^+D^-)$ $S(\phi K^0)$	0.20 0.13	0.03 0.02 (*)
$S(\phi K^{-})$ $S(\eta' K^{0})$	0.05	0.01 (*)
$S(K_{0}^{0}K_{0}^{0}K_{0}^{0})$	0.15	0.02 (*)
$S(K_{S}^{0}K_{S}^{0}K_{S}^{0})$ $S(K_{S}^{0}\pi^{0})$	0.15	0.02 (*)
$S(\omega K_s^0)$	0.17	0.03 (*)
$S(f_0 K_S^0)$	0.12	0.02 (*)
$\gamma (B \rightarrow DK, D \rightarrow CP \text{ eigensta})$	ates) $\sim 15^{\circ}$	2.5°
$\gamma (B \rightarrow DK, D \rightarrow \text{suppressed})$		2.0°
$\gamma (B \rightarrow DK, D \rightarrow \text{multibody s})$		1.5°
$\gamma (B \rightarrow DK, \text{ combined})$	$\sim 6^{\circ}$	1-2°
$\alpha (B \rightarrow \pi \pi)$	$\sim 16^{\circ}$	3°
$\alpha (B \rightarrow \rho \rho)$	$\sim 7^{\circ}$	$1-2^{\circ}(*)$
$\alpha (B \rightarrow \rho \pi)$	$\sim 12^{\circ}$	20
α (combined)	$\sim 6^{\circ}$	$1-2^{\circ}$ (*)
$2\beta + \gamma \left(D^{(*)\pm}\pi^{\mp}, D^{\pm}K^{0}_{S}\pi^{\mp} \right)$	20°	50
$ V_{cb} $ (exclusive)	4% (*)	1.0% (*)
$ V_{cb} $ (inclusive)	1% (+)	0.5% (*)
$ V_{ub} $ (exclusive)	8% (*)	3.0% (*)
$ V_{ub} $ (inclusive)	8% (*)	2.0% (*)
$BR(B \rightarrow \tau \nu)$	20%	4% (†)
$BR(B \rightarrow \mu\nu)$ $BR(B \rightarrow D\tau\nu)$	visible	5%
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$BR(B \rightarrow \rho \gamma)$ $BR(B \rightarrow \nu \gamma)$	Observable	Error with 1 ab^{-1}
$BR(B \rightarrow \omega \gamma)$ $A_{CP}(B \rightarrow K^* \gamma)$	$\Delta\Gamma$	0.16 ps^{-1}
$A_{CP}(B \rightarrow \rho\gamma)$	Г	$0.07 \ {\rm ps}^{-1}$
$A_{CP}(b \rightarrow s\gamma)$	β_s from angular analysis	20°
$A_{CP}(b \rightarrow (s + d)\gamma)$ $S(K_{S}^{0}\pi^{0}\gamma)$	$A^s_{ m SL}$	0.006
$S(R_S^{(R-\gamma)})$ $S(\rho^0\gamma)$	$A_{ m CH}$	0.004
$A_{CP}(B \rightarrow K^*\ell\ell)$	$\mathcal{B}(B_s \to \mu^+ \mu^-)$	121
$A^{FB}(B \rightarrow K^*\ell\ell)s_0$	$ V_{td}/V_{ts} $	0.08
$A^{FB}(B \rightarrow X_s \ell \ell) s_0$	$\mathcal{B}(B_s \to \gamma \gamma)$	38%
$BR(B \rightarrow K \nu \overline{\nu})$ $BR(B \rightarrow \pi \nu \overline{\nu})$	$\beta_s \text{ from } J/\psi\phi$	10°
	$\rho_s \operatorname{nom} J/\psi \phi$	10

Physics discovery all these are good enough, even those NP independent any of them can be the best one if the appropriate NP scenario is realized other considerations

(SuperB only, no lattice, ...) are required to select few flagship measurements

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Overture: CKM matrix at 1%



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evaluating NP discovery potential

using EFT's

- fully NP model free encompass any NP
- refined dim. analysis
 up to "O(1)" factors
- not fully specified
 may miss correlations
- virtual NP contrib. only new particle content unspecified

using explicit models

- theoretically biased may miss the point partially cured using several models
- fully specified
 can exploit correlations
- real processes calculable known particle content interplay with direct searches possible

EFT approach to New Flavour Physics a game of scale and couplings

 $\mathcal{L} = \mathcal{L}_{SM} + \Sigma_{k} \left(\Sigma_{i} C_{i}^{k} Q_{i}^{(k+4)} \right) / \Lambda^{k}$

NP flavour effects are governed by two players:
i) the value of the new physics scale Λ
ii) the effective flavour-violating couplings C's

In explict models:

Λ ~ mass of virtual particles (SM: M_W)
 C ~ loop coupling x flavour coupling
 (SM/MFV: a_W x CKM)

B physics: why?

SM FCNC and CPV processes occur at the loop level and thus could receive O(1) NP corrections but effects >20% are excluded

common misconception: this result points to MFV (or even establishes MFV)

if NP < 1 TeV

- * suppression of flavourviolating couplings required in all sectors *possibly* pointing to MFV
- * SUSY can stabilize the Fermi scale with "mild" fine-tuning

if 1 < NP < 10-100 TeV

* suppression of flavourviolating couplings needed in sector 1-2 only. No indication of MFV

* SUSY can still stabilize the Fermi scale with "moderate" fine-tuning

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SuperB physics goals NP found at LHC * determine the flavour- and CP-violating couplings the NP Lagrangian * look for the effects of heavier states beyond the LHC discovery reach NP not found at LHC * look for indirect NP signals coming from the 1-100 TeV energy range * exclude regions of the NP parameter space

Reconstructing the MSSM Lagrangian

Parameters	MSSM		SM
gauge+Higgs	14		6
masses	30	(36)	9 (12)
mixing angles	39	(54)	3 (6)
phases	41	(56)	1 (2)
Total	124	(160)	19 (26)

MSSM parameters match: FC vs FV&CPV 50-110

* flavour and high-p_T physics are complementary
 * the region of "small" masses AND large FV&CPV couplings is largely ruled out already

from "LHC ready" to "full LHC"



Are the FV couplings of new particles below the TeV scale measurable at a SuperB even in the worst case (MFV)?

Likely, but not guaranteed (yet)!

The EFT analysis of Δ F=2 processes finds SuperB sensitivity of ~0.6 TeV for strict MFV couplings in the low tanB regime

The $\Delta F=2$ effective Hamiltonian

$$H_{eff}^{\Delta B=2} = \sum_{i=1}^{5} C_{i}(\mu) Q_{i}(\mu) + \sum_{i=1}^{3} \widetilde{C}_{i}(\mu) \widetilde{Q}_{i}(\mu)$$

$$Q_{1} = \overline{q}_{L}^{\alpha} \gamma_{\mu} b_{L}^{\alpha} \overline{q}_{L}^{\beta} \gamma^{\mu} b_{L}^{\beta} \quad (SM/MFV)$$

$$Q_{2} = \overline{q}_{R}^{\alpha} b_{L}^{\alpha} \overline{q}_{R}^{\beta} b_{L}^{\beta} \qquad Q_{3} = \overline{q}_{R}^{\alpha} b_{L}^{\beta} \overline{q}_{R}^{\beta} b_{L}^{\beta}$$

$$Q_{4} = \overline{q}_{R}^{\alpha} b_{L}^{\alpha} \overline{q}_{L}^{\beta} b_{R}^{\beta} \qquad Q_{5} = \overline{q}_{R}^{\alpha} b_{L}^{\beta} \overline{q}_{L}^{\beta} b_{R}^{\beta}$$

$$\widetilde{Q}_{1} = \overline{q}_{R}^{\alpha} \gamma_{\mu} b_{R}^{\alpha} \overline{q}_{R}^{\beta} \gamma^{\mu} b_{R}^{\beta} \qquad \widetilde{Q}_{3} = \overline{q}_{L}^{\alpha} b_{R}^{\beta} \overline{q}_{L}^{\beta} b_{R}^{\beta}$$

7 new operators beyond MFV involving quarks with different chiralities

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 H_{eff} can be recast in terms of the high-scale $C_i(\Lambda)$

- $C_i(\Lambda)$ can be extracted from the data (one by one) - the associated NP scale Λ can be defined as

$$\Lambda = \sqrt{\frac{LF_i}{C_i(\Lambda)}}$$

tree/strong interact. NP: L ~ 1 perturbative NP: L ~ α_s^2 , α_W^2

Flavour structures:

MFVnext-to-MFVgeneric $- F_1 = F_{SM} \sim (V_{tq}V_{tb}^*)^2$ $- |F_i| \sim F_{SM}$ $- |F_i| \sim 1$ $- F_{i\neq 1} = 0$ - arbitrary- arbitrary- arbitraryphasesphasesphases

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present lower bound on the NP scale (TeV) B + K B only

Scenario	strong/tree	α_s loop	α_W loop
MFV	5.5	0.5	0.2
NMFV	62	6.2	2
General	24000	2400	800

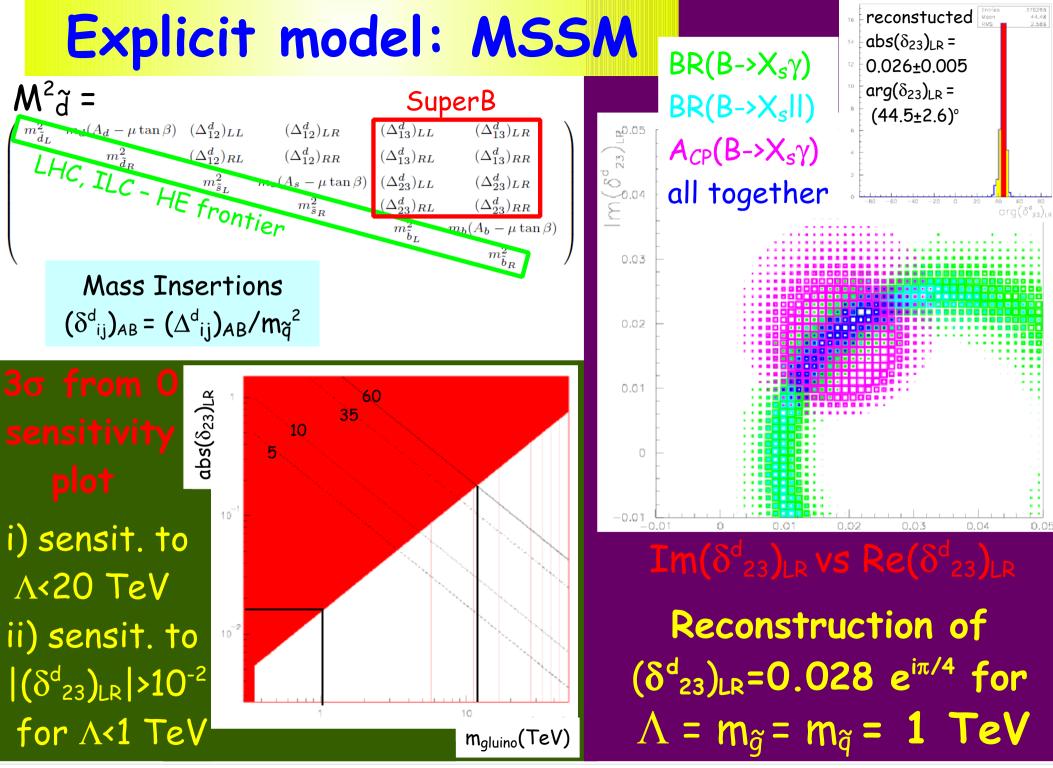
$\rm strong/tree$	α_s loop	α_W loop		
_	_	_		
14	1.4	0.4		
2200	220	66		

UTfit collaboration, arXiv:0707.0636

SuperB: typically 3x present bounds $\Lambda_{MFV} > 0.6$, $\Lambda_{NMFV} > 1.2$, $\Lambda_{GFV} > 220 \text{ TeV}$ ΔF =2 processes alone do not allow the detection at SuperB of particles with M > 0.6 TeV in the MFV scenario $\Delta B=1$ processes, in particular radiative decays, have not been fully exploited yet Present bounds D'Ambrosio et al., hep-ph/0207036 $b \rightarrow s \gamma : \Lambda_{MFV} > 9 \text{ TeV}$

 $b \rightarrow s II: \Lambda_{MFV} > 3 TeV$

Dedicated experimental & theoretical studies needed to assess the SuperB potential for these processes (starting here and now with Hurth, Renga, Walsh, ...)



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- assessing the SuperB potential on flavour-nasty SUSY models (such as mSUGRA) would be highly desirable. Furthermore the IRC asked to evaluate SuperB performances on LHC benchmarks (starting here with help from Shindou and from our high- p_T friends Heinemeyer and Ronga)

 consider other models besides SUSY (some LHT already in the CDR, here extra dimensions with Kou)

An issue to debate: 10^{36} vs 10^{35}

we plan to complete the exercise for B physics during this week. Yet O(1) differences on accessible scales and couplings can be anticipated although any such difference could be crucial, they are likely not very impressive to present and difficult to defend, given the intrinsic uncertainty of the EFT approach which was heavily used we should look for gualitative differences in the physics that can be done with 75/ab vs the physics possible with say 10-15/ab

This means that we have to look at qualifying points which are difficult to achieve even at 10³⁶! A couple comes to my mind: i) the "full LHC" label, i.e. the possibility to to measure flavour effects in the whole LHC discovery energy range, fully playing the complementarity game with high-pT searches ii) the possibility of measuring theoretically interesting values of LFV BRs and of being complementary with the MEG measurement These goals are not attainable with 10³⁵ but we have to clearly show that they are at SuperB

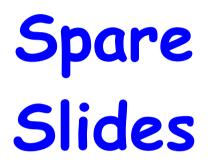
Working Group on B physics

- * totally informal working group setup
- * program subject to last-minute changes
- * few presentations (some of them, notably mine, rather working plans). For theory, we have:
 - Paolo Gambino (by phone) on Vub
 - Emi Kou on B -> tvy and extra dimensions
 - Tetsuo Shindou on SUSY breaking scenarios
 - Tobias Hurth on b -> s y and b -> s Il
 - Florian Domingo on Next-to-MSSM
- * a lot of spare time for discussions, work, new ideas, coffee, local attractions, ...

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Let's make it a SuperB workshop!!!

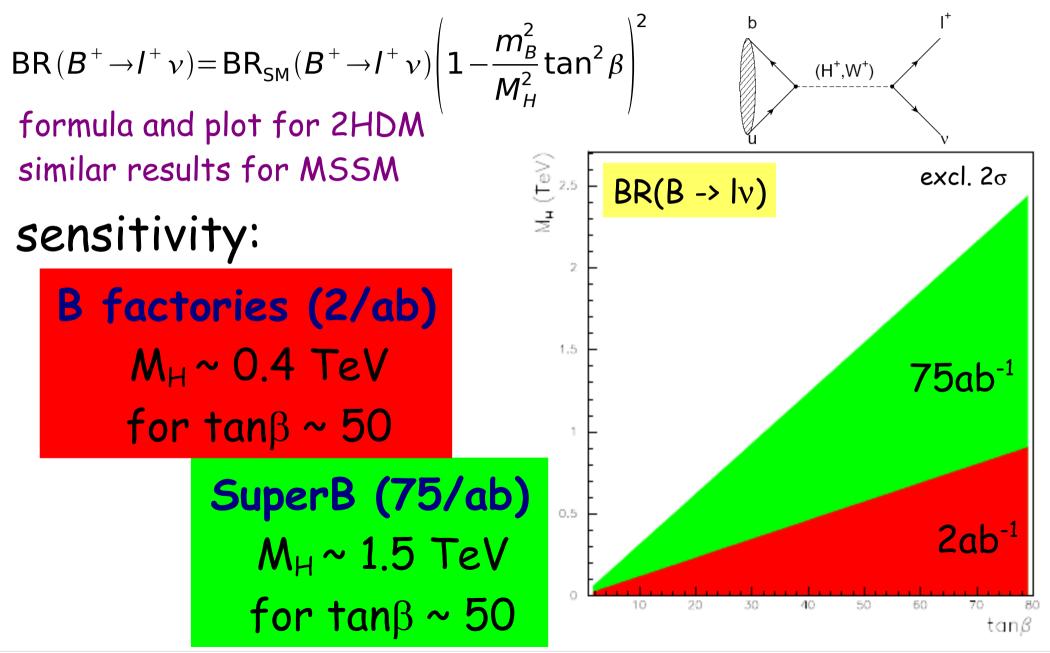


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Conclusions

The SFF physics program complements and extends the NP searches with high p_T at LHC (i) NP is behind the corner: once the LHC finds it, (only) a SFF can measure systematically the new FV & CPV couplings, i.e. the flavour structure of NP (ii) NP is a few corners ahead: NP at scales beyond the LHC reach could give measurable effects at a SFF: unique opportunity to unveil the 1-100 TeV range

Higgs-mediated NP in MFV at large tanß



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Minimal Flavour Violation

Gabrielli, Giudice, NPB433 Buras et al., NPB500 D'Ambrosio et al., NPB645

No new sources of flavour D'Am and CP violation beyond the SM

 NP contributions governed by SM Yukawa couplings ex.: Constrained MSSM (MSUGRA), Universal Extra Dim.
 NP only modifies SM top contribution to FCNC & CPV unless other Yukawa couplings are enhanced; for example large tanβ enhances bottom contributions

1HDM/2HDM at small tanß same operators as in H_{eff}^{SM} NP in K and B correlated 2HDM at large tanß new operators wrt H_{eff} SM NP in K and B uncorrelated



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Crucial questions for NP searches with flavour

can NP be flavour blind? "no",
 NP couples to SM which violates flavour

 2. can a "worst case" be defined? "yes", through the class of models with Minimal Flavour Violation
 NP follows the SM pattern of flavour and CP symmetry breaking Gabrielli, Giudice, NPB433 Buras et al. NPB500

Buras et al, NPB500 D'Ambrosio et al., NPB645

Conclusions (ii)

Only part of the SuperB physics program relies on theory upgrades. For this part, theoretical errors of O(1-2%) are needed: feasible for LQCD; challenging but possibly reachable in inclusive measurements; factorization needs checking on channel basis SuperB and S-LHCb physics programs are largely complementary. As for the part in common, they are competitive but SuperB can measure more and th. cleaner channels