

Charged dark matter in supersymmetric Twin Higgs models

based on [2202.10488] by Marcin Badziak, Giovanni Grilli di Cortona, Keisuke Harigaya and MŁ

Michał Łukawski

Faculty of Physics
University of Warsaw

14.07.2022

Summary

- 1 Hierarchy problem
- 2 Twin Higgs
- 3 Twin stau as DM candidate

Hierarchy problem

In SM $m_h^2 = (m_h^0)^2 + \delta m_h^2$:

$$\delta m_h^2 = \frac{3}{4\pi^2} \left(-y_t^2 + \frac{g^2}{4} + \frac{g'^2}{8 \cos^2 \theta_W} + \lambda \right) \Lambda^2$$

Higgs mass in Standard Model (SM) is not natural, $m_h \ll \Lambda_{\text{cut-off}}$.

Higgs mass in Standard Model (SM) is not even technically natural - no reason to believe that any hierarchy between m_h and new physics should hold.

Hierarchy problem

In SM $m_h^2 = (m_h^0)^2 + \delta m_h^2$:

$$\delta m_h^2 = \frac{3}{4\pi^2} \left(-y_t^2 + \frac{g^2}{4} + \frac{g'^2}{8 \cos^2 \theta_W} + \lambda \right) \Lambda^2$$

Higgs mass in Standard Model (SM) is not natural, $m_h \ll \Lambda_{\text{cut-off}}$.

Higgs mass in Standard Model (SM) is not even technically natural - no reason to believe that any hierarchy between m_h and new physics should hold.

Hierarchy problem

In SM $m_h^2 = (m_h^0)^2 + \delta m_h^2$:

$$\delta m_h^2 = \frac{3}{4\pi^2} \left(-y_t^2 + \frac{g^2}{4} + \frac{g'^2}{8 \cos^2 \theta_W} + \lambda \right) \Lambda^2$$

Higgs mass in Standard Model (SM) is not natural, $m_h \ll \Lambda_{\text{cut-off}}$.

Higgs mass in Standard Model (SM) is not even technically natural - no reason to believe that any hierarchy between m_h and new physics should hold.

Twin Higgs

1. double the particle content adding twin sector (in particular second higgs H')
2. impose Z_2 symmetry interchanging particles between sectors
3. the scalar potential is $SU(4)$ invariant due to Z_2 symmetry

$$V(\mathcal{H}) = -m_{\mathcal{H}}^2(H^2 + H'^2) + \lambda(H^2 + H'^2)^2 = -m_{\mathcal{H}}^2\mathcal{H}^\dagger\mathcal{H} + \lambda(\mathcal{H}^\dagger\mathcal{H})^2$$

4. spontaneous symmetry breaking of $SU(4) \rightarrow SU(3)$ generates SM Higgs as one of Nambu-Goldstone bosons
5. Quadratically divergent gauge contributions to the potential

$$\delta V = \frac{9\Lambda^2 g^2}{64\pi^2}(H^\dagger H + H'^\dagger H') = \frac{9g^2\Lambda^2}{64\pi^2}\mathcal{H}^\dagger\mathcal{H}$$

$SU(4)$ symmetric!

Twin Higgs

1. double the particle content adding twin sector (in particular second higgs H')
2. impose Z_2 symmetry interchanging particles between sectors
3. the scalar potential is $SU(4)$ invariant due to Z_2 symmetry

$$V(\mathcal{H}) = -m_{\mathcal{H}}^2(H^2 + H'^2) + \lambda(H^2 + H'^2)^2 = -m_{\mathcal{H}}^2\mathcal{H}^\dagger\mathcal{H} + \lambda(\mathcal{H}^\dagger\mathcal{H})^2$$

4. spontaneous symmetry breaking of $SU(4) \rightarrow SU(3)$ generates SM Higgs as one of Nambu-Goldstone bosons
5. Quadratically divergent gauge contributions to the potential

$$\delta V = \frac{9\Lambda^2 g^2}{64\pi^2}(H^\dagger H + H'^\dagger H') = \frac{9g^2\Lambda^2}{64\pi^2}\mathcal{H}^\dagger\mathcal{H}$$

$SU(4)$ symmetric!

Twin Higgs

1. double the particle content adding twin sector (in particular second higgs H')
2. impose Z_2 symmetry interchanging particles between sectors
3. the scalar potential is $SU(4)$ invariant due to Z_2 symmetry

$$V(\mathcal{H}) = -m_{\mathcal{H}}^2(H^2 + H'^2) + \lambda(H^2 + H'^2)^2 = -m_{\mathcal{H}}^2\mathcal{H}^\dagger\mathcal{H} + \lambda(\mathcal{H}^\dagger\mathcal{H})^2$$

4. spontaneous symmetry breaking of $SU(4) \rightarrow SU(3)$ generates SM Higgs as one of Nambu-Goldstone bosons
5. Quadratically divergent gauge contributions to the potential

$$\delta V = \frac{9\Lambda^2 g^2}{64\pi^2}(H^\dagger H + H'^\dagger H') = \frac{9g^2\Lambda^2}{64\pi^2}\mathcal{H}^\dagger\mathcal{H}$$

$SU(4)$ symmetric!

Twin Higgs

1. double the particle content adding twin sector (in particular second higgs H')
2. impose Z_2 symmetry interchanging particles between sectors
3. the scalar potential is $SU(4)$ invariant due to Z_2 symmetry

$$V(\mathcal{H}) = -m_{\mathcal{H}}^2(H^2 + H'^2) + \lambda(H^2 + H'^2)^2 = -m_{\mathcal{H}}^2\mathcal{H}^\dagger\mathcal{H} + \lambda(\mathcal{H}^\dagger\mathcal{H})^2$$

4. spontaneous symmetry breaking of $SU(4) \rightarrow SU(3)$ generates SM Higgs as one of Nambu-Goldstone bosons
5. Quadratically divergent gauge contributions to the potential

$$\delta V = \frac{9\Lambda^2 g^2}{64\pi^2}(H^\dagger H + H'^\dagger H') = \frac{9g^2\Lambda^2}{64\pi^2}\mathcal{H}^\dagger\mathcal{H}$$

$SU(4)$ symmetric!

Twin Higgs

1. double the particle content adding twin sector (in particular second higgs H')
2. impose Z_2 symmetry interchanging particles between sectors
3. the scalar potential is $SU(4)$ invariant due to Z_2 symmetry

$$V(\mathcal{H}) = -m_{\mathcal{H}}^2(H^2 + H'^2) + \lambda(H^2 + H'^2)^2 = -m_{\mathcal{H}}^2\mathcal{H}^\dagger\mathcal{H} + \lambda(\mathcal{H}^\dagger\mathcal{H})^2$$

4. spontaneous symmetry breaking of $SU(4) \rightarrow SU(3)$ generates SM Higgs as one of Nambu-Goldstone bosons
5. Quadratically divergent gauge contributions to the potential

$$\delta V = \frac{9\Lambda^2 g^2}{64\pi^2}(H^\dagger H + H'^\dagger H') = \frac{9g^2\Lambda^2}{64\pi^2}\mathcal{H}^\dagger\mathcal{H}$$

$SU(4)$ symmetric!

Twin Higgs models

General Twin Higgs potential could be written

$$V(H, H') = \lambda(H^2 + H'^2)^2 - m_{\mathcal{H}}^2(H^2 + H'^2) + \Delta\lambda(H^4 + H'^4) + \Delta m^2 H^2$$

In minimal setting 4 parameters, but we know mass of Higgs m_h and EW vev v

We have only two parameters v'/v and mass of the heavy higgs $m_{\mathcal{H}}$.

Twin Higgs models

General Twin Higgs potential could be written

$$V(H, H') = \lambda(H^2 + H'^2)^2 - m_{\mathcal{H}}^2(H^2 + H'^2) + \Delta\lambda(H^4 + H'^4) + \Delta m^2 H^2$$

In minimal setting 4 parameters, but we know mass of Higgs m_h and EW vev v

We have only two parameters v'/v and mass of the heavy higgs $m_{\mathcal{H}}$.

DM is TH models

- twin tau ($m_{\tau'} \approx 65 - 130 \text{ GeV}$, [1505.07109])
- twin electrons ($m_{e'} \approx 2 - 5 \text{ MeV}$ [1908.03559])
- twin baryons ($m_{baryon} \approx 5 \text{ GeV}$, [1506.03520])

Twin electromagnetism necessarily broken!

Self-interactions of DM are constrained and for self-coupling $g = g_{em}$ we have $m_{DM} \gtrsim 200 \text{ GeV}$. [1610.04611]

Observation:

SUSY partners obtain large soft masses and can escape that bound, while preserving unbroken $U'_{em}(1)$

DM is TH models

- twin tau ($m_{\tau'} \approx 65 - 130 \text{ GeV}$, [1505.07109])
- twin electrons ($m_{e'} \approx 2 - 5 \text{ MeV}$ [1908.03559])
- twin baryons ($m_{baryon} \approx 5 \text{ GeV}$, [1506.03520])

Twin electromagnetism necessarily broken!

Self-interactions of DM are constrained and for self-coupling $g = g_{em}$ we have $m_{DM} \gtrsim 200 \text{ GeV}$. [1610.04611]

Observation:

SUSY partners obtain large soft masses and can escape that bound, while preserving unbroken $U'_{em}(1)$

DM is TH models

- twin tau ($m_{\tau'} \approx 65 - 130 \text{ GeV}$, [1505.07109])
- twin electrons ($m_{e'} \approx 2 - 5 \text{ MeV}$ [1908.03559])
- twin baryons ($m_{baryon} \approx 5 \text{ GeV}$, [1506.03520])

Twin electromagnetism necessarily broken!

Self-interactions of DM are constrained and for self-coupling $g = g_{em}$ we have $m_{DM} \gtrsim 200 \text{ GeV}$. [1610.04611]

Observation:

SUSY partners obtain large soft masses and can escape that bound, while preserving unbroken $U'_{em}(1)$

Few remarks on SUSY TH

- In SUSY the potential is fixed by particle content, F-term, and gauge interactions, D-term.
- The $SU(4)$ invariant potential may be generated in two way, F-term SUSY TH [1611.08615] and D-term SUSY TH D-term [1703.02122]
- The main difference is captured by preferable values of $\tan \beta$, which are small (F-term) or large (D-term).
- D-term SUSY TH allows for fine-tuning $\mathcal{O}(10\%)$ while heavy stop (2 TeV) [1703.02122]

Few remarks on SUSY TH

- In SUSY the potential is fixed by particle content, F-term, and gauge interactions, D-term.
- The $SU(4)$ invariant potential may be generated in two way, F-term SUSY TH [1611.08615] and D-term SUSY TH D-term [1703.02122]
- The main difference is captured by preferable values of $\tan \beta$, which are small (F-term) or large (D-term).
- D-term SUSY TH allows for fine-tuning $\mathcal{O}(10\%)$ while heavy stop (2 TeV) [1703.02122]

Few remarks on SUSY TH

- In SUSY the potential is fixed by particle content, F-term, and gauge interactions, D-term.
- The $SU(4)$ invariant potential may be generated in two way, F-term SUSY TH [1611.08615] and D-term SUSY TH D-term [1703.02122]
- The main difference is captured by preferable values of $\tan \beta$, which are small (F-term) or large (D-term).
- D-term SUSY TH allows for fine-tuning $\mathcal{O}(10\%)$ while heavy stop (2 TeV) [1703.02122]

Few remarks on SUSY TH

- In SUSY the potential is fixed by particle content, F-term, and gauge interactions, D-term.
- The $SU(4)$ invariant potential may be generated in two way, F-term SUSY TH [1611.08615] and D-term SUSY TH D-term [1703.02122]
- The main difference is captured by preferable values of $\tan \beta$, which are small (F-term) or large (D-term).
- D-term SUSY TH allows for fine-tuning $\mathcal{O}(10\%)$ while heavy stop (2 TeV) [1703.02122]

Few remarks

1. assume Z_2 symmetric soft breaking terms and $\tan\beta$
2. lightest supersymmetric particle (LSP) is stable
3. **twin stau** is Z_2 partner of stau, supersymmetric partner of tau
4. The mass matrix of stau is given by

$$m_{\tilde{\tau}}^2 = \begin{pmatrix} m_{L3}^2 + \Delta_{\tilde{\tau}_L} + m_{\tau}^2 & -\mu v y_{\tau} \sin(\beta) \\ -\mu v y_{\tau} \sin(\beta) & m_{E3}^2 + \Delta_{\tilde{\tau}_R} + m_{\tau}^2 \end{pmatrix}$$

5. for pure $\tilde{\tau}'_L$ and $\tilde{\tau}'_R$ visible stau is LSP
6. off-diagonal term is larger in twin sector, for mixed state twin stau may be LSP

Few remarks

1. assume Z_2 symmetric soft breaking terms and $\tan\beta$
2. lightest supersymmetric particle (LSP) is stable
3. **twin stau** is Z_2 partner of stau, supersymmetric partner of tau
4. The mass matrix of stau is given by

$$m_{\tilde{\tau}}^2 = \begin{pmatrix} m_{L_3}^2 + \Delta_{\tilde{\tau}_L} + m_{\tau}^2 & -\mu\nu y_{\tau} \sin(\beta) \\ -\mu\nu y_{\tau} \sin(\beta) & m_{\tilde{e}_3}^2 + \Delta_{\tilde{\tau}_R} + m_{\tau}^2 \end{pmatrix}$$

5. for pure $\tilde{\tau}'_L$ and $\tilde{\tau}'_R$ visible stau is LSP
6. off-diagonal term is larger in twin sector, for mixed state twin stau may be LSP

Few remarks

1. assume Z_2 symmetric soft breaking terms and $\tan \beta$
2. lightest supersymmetric particle (LSP) is stable
3. **twin stau** is Z_2 partner of stau, supersymmetric partner of tau
4. The mass matrix of stau is given by

$$m_{\tilde{\tau}}^2 = \begin{pmatrix} m_{L_3}^2 + \Delta_{\tilde{\tau}_L} + m_{\tau}^2 & -\mu v y_{\tau} \sin(\beta) \\ -\mu v y_{\tau} \sin(\beta) & m_{\tilde{e}_3}^2 + \Delta_{\tilde{\tau}_R} + m_{\tau}^2 \end{pmatrix}$$

5. for pure $\tilde{\tau}'_L$ and $\tilde{\tau}'_R$ visible stau is LSP
6. off-diagonal term is larger in twin sector, for mixed state twin stau may be LSP

Few remarks

1. assume Z_2 symmetric soft breaking terms and $\tan \beta$
2. lightest supersymmetric particle (LSP) is stable
3. **twin stau** is Z_2 partner of stau, supersymmetric partner of tau
4. The mass matrix of stau is given by

$$m_{\tilde{\tau}}^2 = \begin{pmatrix} m_{L_3}^2 + \Delta_{\tilde{\tau}_L} + m_{\tau}^2 & -\mu\nu y_{\tau} \sin(\beta) \\ -\mu\nu y_{\tau} \sin(\beta) & m_{e_3}^2 + \Delta_{\tilde{\tau}_R} + m_{\tau}^2 \end{pmatrix}$$

5. for pure $\tilde{\tau}'_L$ and $\tilde{\tau}'_R$ visible stau is LSP
6. off-diagonal term is larger in twin sector, for mixed state twin stau may be LSP

Few remarks

1. assume Z_2 symmetric soft breaking terms and $\tan \beta$
2. lightest supersymmetric particle (LSP) is stable
3. **twin stau** is Z_2 partner of stau, supersymmetric partner of tau
4. The mass matrix of stau is given by

$$m_{\tilde{\tau}}^2 = \begin{pmatrix} m_{L_3}^2 + \Delta_{\tilde{\tau}_L} + m_{\tau}^2 & -\mu\nu y_{\tau} \sin(\beta) \\ -\mu\nu y_{\tau} \sin(\beta) & m_{\tilde{e}_3}^2 + \Delta_{\tilde{\tau}_R} + m_{\tau}^2 \end{pmatrix}$$

5. for pure $\tilde{\tau}'_L$ and $\tilde{\tau}'_R$ visible stau is LSP
6. off-diagonal term is larger in twin sector, for mixed state twin stau may be LSP

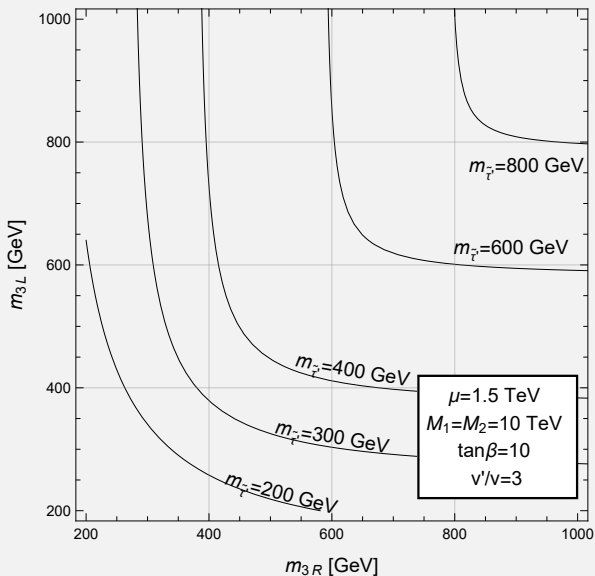
Few remarks

1. assume Z_2 symmetric soft breaking terms and $\tan \beta$
2. lightest supersymmetric particle (LSP) is stable
3. **twin stau** is Z_2 partner of stau, supersymmetric partner of tau
4. The mass matrix of stau is given by

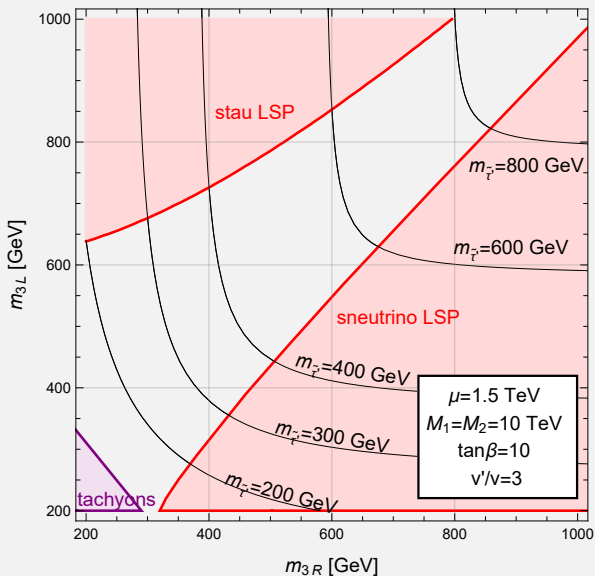
$$m_{\tilde{\tau}}^2 = \begin{pmatrix} m_{L_3}^2 + \Delta_{\tilde{\tau}_L} + m_{\tau}^2 & -\mu v y_{\tau} \sin(\beta) \\ -\mu v y_{\tau} \sin(\beta) & m_{\tilde{e}_3}^2 + \Delta_{\tilde{\tau}_R} + m_{\tau}^2 \end{pmatrix}$$

5. for pure $\tilde{\tau}'_L$ and $\tilde{\tau}'_R$ visible stau is LSP
6. off-diagonal term is larger in twin sector, for mixed state twin stau may be LSP

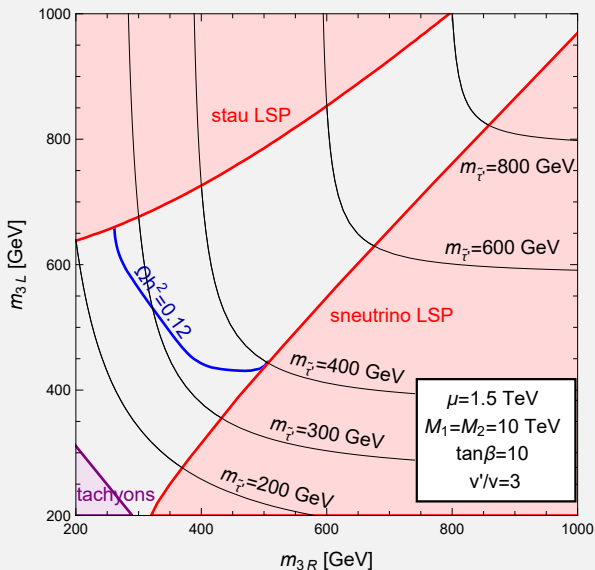
Decoupling limit



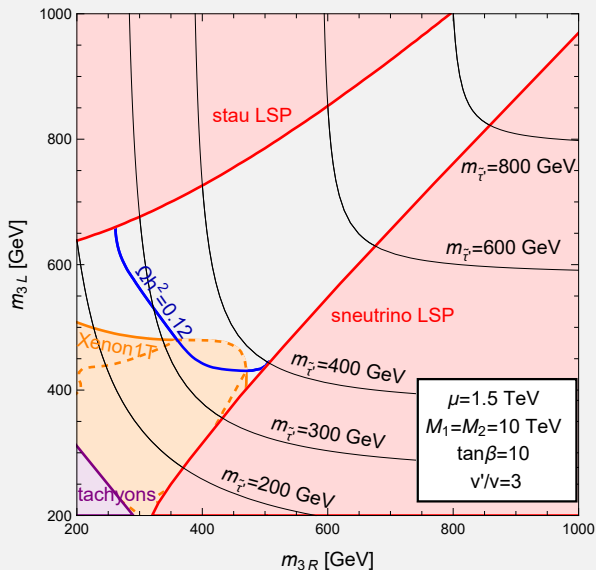
Decoupling limit



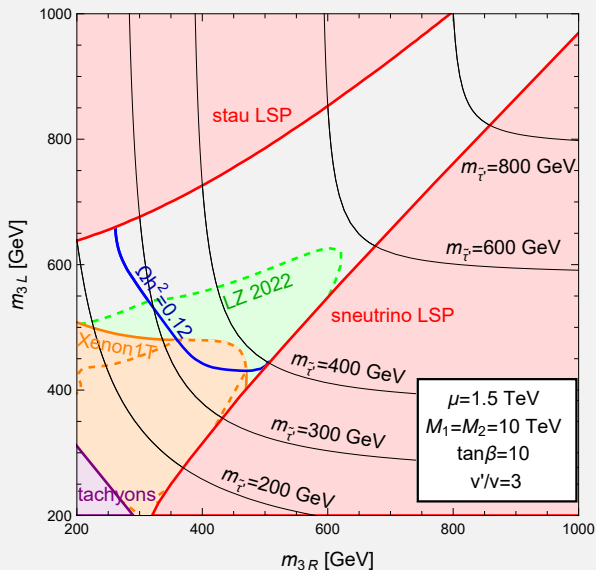
Decoupling limit



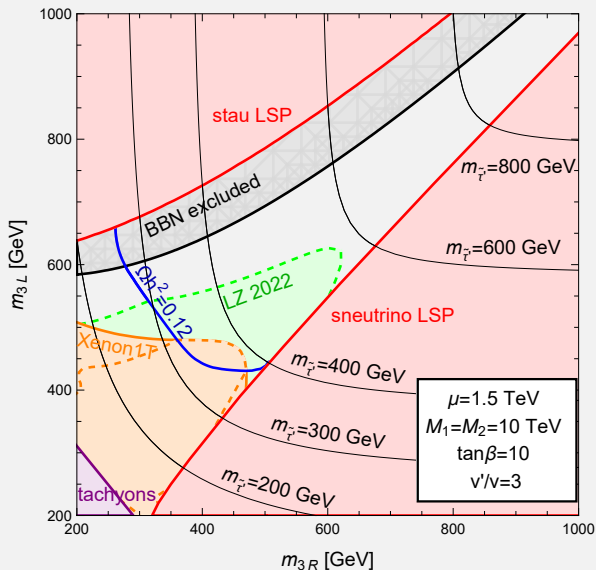
Decoupling limit



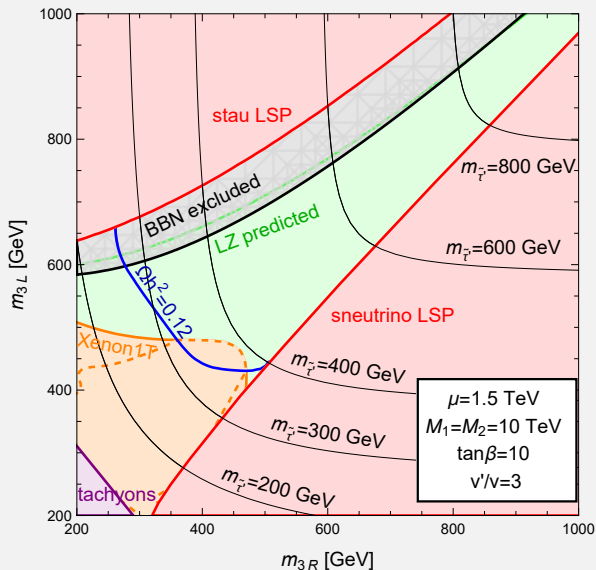
Decoupling limit



Decoupling limit



Decoupling limit



Conclusions

- Supersymmetric Twin Higgs models allow for naturally light Higgs boson, stable under quantum corrections
- in TH models usually one needs to break twin electromagnetism to obtain DM
- in SUSY completions large soft masses allow for LSP charged under twin EM
- twin stau DM will be probed by LZ experiment

Conclusions

- Supersymmetric Twin Higgs models allow for naturally light Higgs boson, stable under quantum corrections
- in TH models usually one needs to break twin electromagnetism to obtain DM
- in SUSY completions large soft masses allow for LSP charged under twin EM
- twin stau DM will be probed by LZ experiment

Conclusions

- Supersymmetric Twin Higgs models allow for naturally light Higgs boson, stable under quantum corrections
- in TH models usually one needs to break twin electromagnetism to obtain DM
- in SUSY completions large soft masses allow for LSP charged under twin EM
- twin stau DM will be probed by LZ experiment

Conclusions

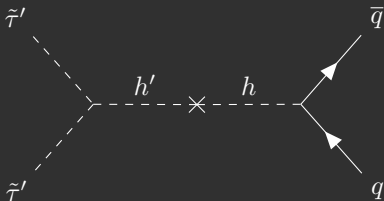
- Supersymmetric Twin Higgs models allow for naturally light Higgs boson, stable under quantum corrections
- in TH models usually one needs to break twin electromagnetism to obtain DM
- in SUSY completions large soft masses allow for LSP charged under twin EM
- twin stau DM will be probed by LZ experiment

Thank you

Direct detection

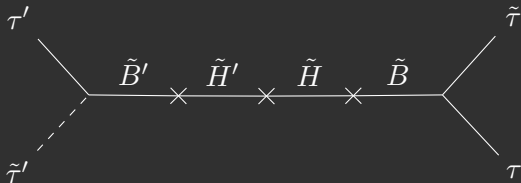
Twin stau can interact with nucleons from visible sector via Higgs portal.
The relevant coupling in decoupling limit is

$$\lambda_{h\tilde{\tau}'\tilde{\tau}'} = \frac{g}{m_{W'}} \left[\left(\frac{1}{2} c_{\theta_{\tilde{\tau}'}}^2 - s_{2\theta_{\tilde{\tau}'}}^2 \right) m_{Z'}^2 c_{2\theta_{\tilde{\tau}'}} - m_{\tau'}^2 + \frac{m_{\tau'}}{2} \mu \tan \beta s_{2\theta_{\tilde{\tau}'}} \right] \frac{v}{v'}$$



Lifetime of stau

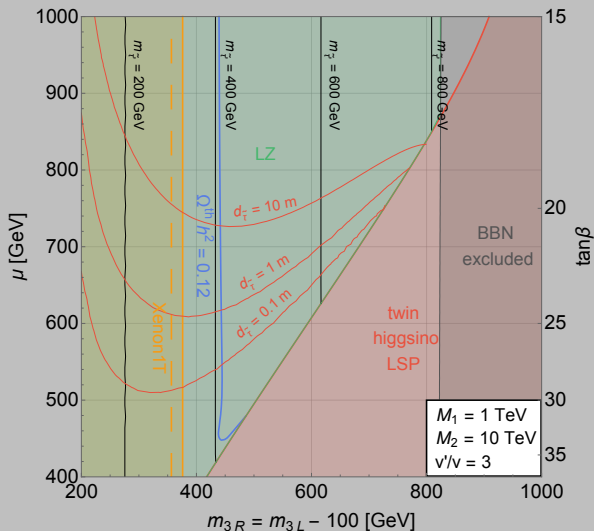
Effective $\tilde{\tau}\tilde{\tau}'^\dagger\tau\tau'$ operator from diagram:



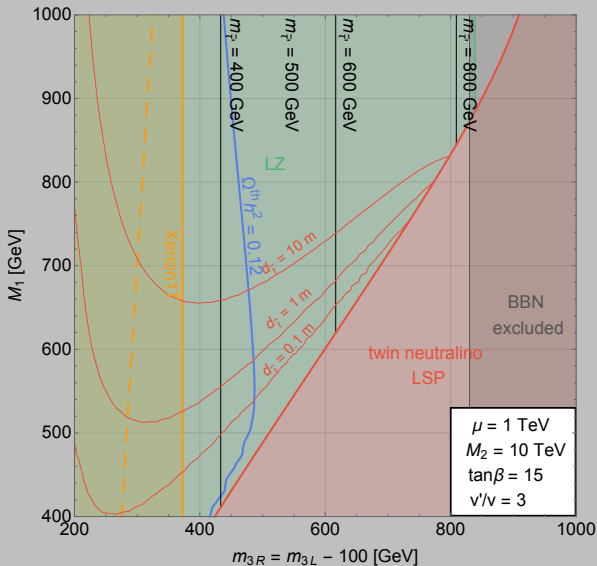
$$\frac{1}{M} \tilde{\tau}\tilde{\tau}'^\dagger\tau\tau' = \frac{g'^4 v v' \varepsilon_{\tilde{H}} m_{\tilde{\tau}}^2 (M_1^2 + m_{\tilde{\tau}}^2)}{(M_1^2 - m_{\tilde{\tau}}^2)^2 (\mu^2 - m_{\tilde{\tau}}^2)^2} \tilde{\tau}\tilde{\tau}'^\dagger\tau\tau'$$

$$d_{\tilde{\tau}} \simeq 2.7 \text{ m} \left(\frac{m_{\tilde{\tau}}}{300 \text{ GeV}} \right)^2 \left(\frac{M}{10^6 \text{ GeV}} \right)^2 \left(\frac{10 \text{ GeV}}{m_{\tilde{\tau}} - m_{\tilde{\tau}'}} \right)^5 \quad (1)$$

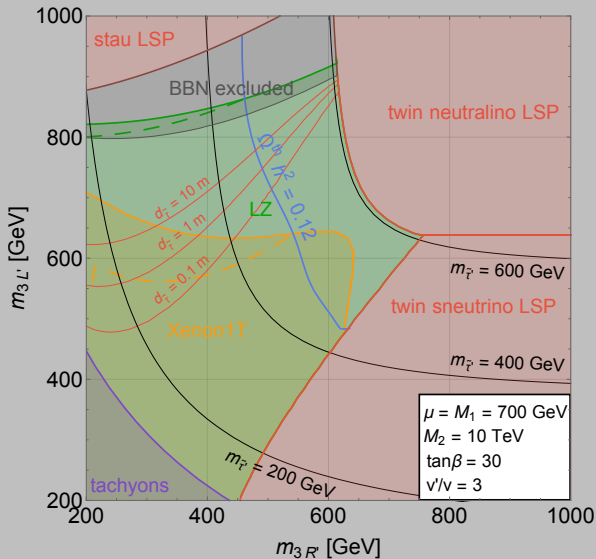
Light Higgsino



Light bino



Light higgsino and bino



Breaking Z_2 in Yukawa

