

XX Frascati Summer School "Bruno Touschek" in Nuclear, Subnuclear and Astroparticle Physics July 2022, Frascati, Italy

QCD and Jets at Colliders

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Lecture I: QCD and perturbative calculations and tools

Lecture 2: Jet algorithms and substructure

[Includes material from Gavin Salam and Grégory Soyez]







Outline

- Jet algorithms
 - ▶ How are jets made
- ▶ Jet substructure
 - What's inside them, and how to use it

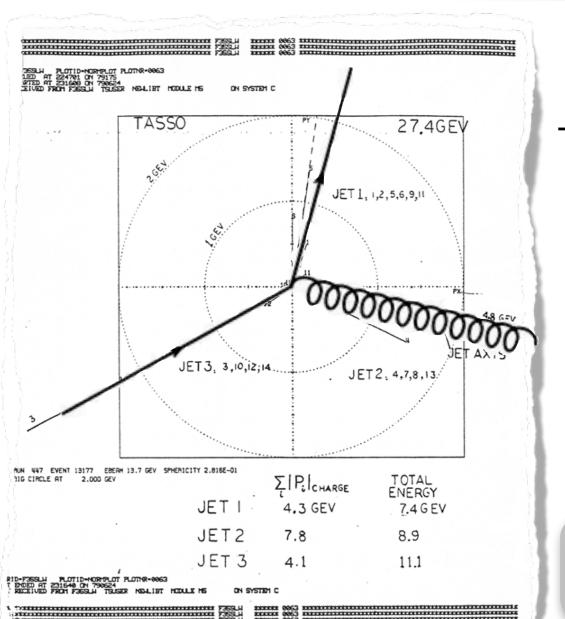
What is a jet?



No, not this....

A jet is something that happens in high energy events: a collimated bunch of hadrons flying roughly in the same direction

Gluon 'discovery'



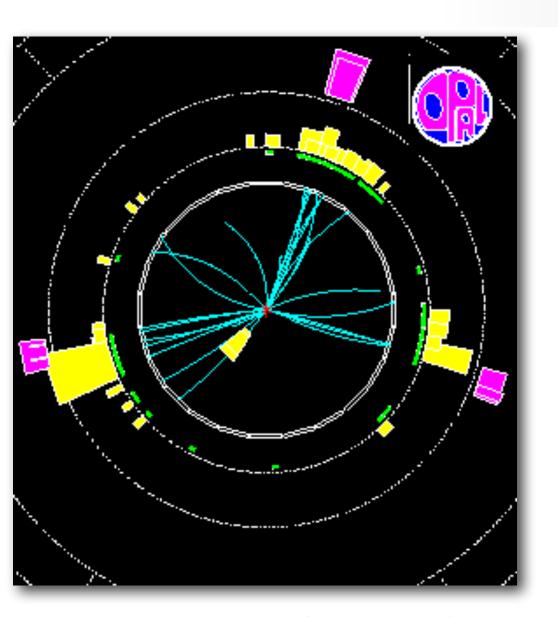
1979:

Three-jet events observed by TASSO, JADE, MARK J and PLUTO at PETRA in e⁺e⁻ collisions at 27.4 GeV

Interpretation: large angle emission of a hard gluon

Jets viewed as a proxy to the initial partons

Why jets



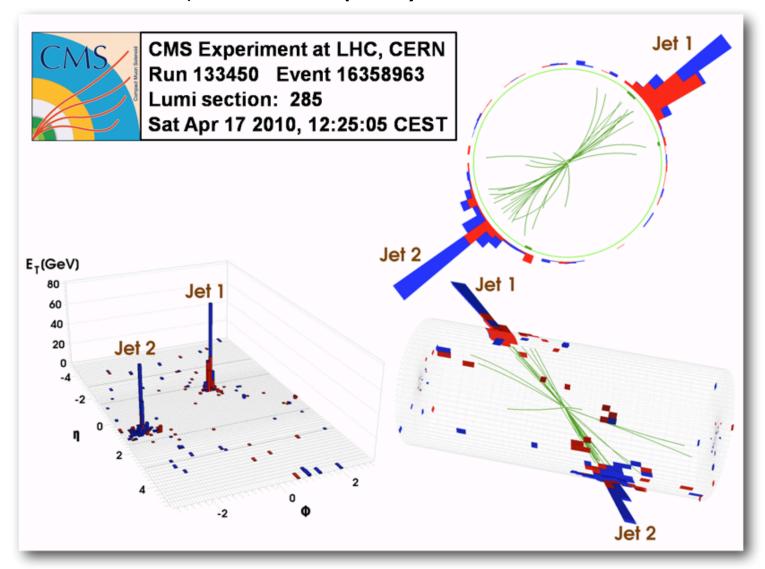
From PETRA to LEP

We could eyeball the collimated bunches, but it becomes impractical with millions of events

The classification of particles into jets is best done using a clustering algorithm

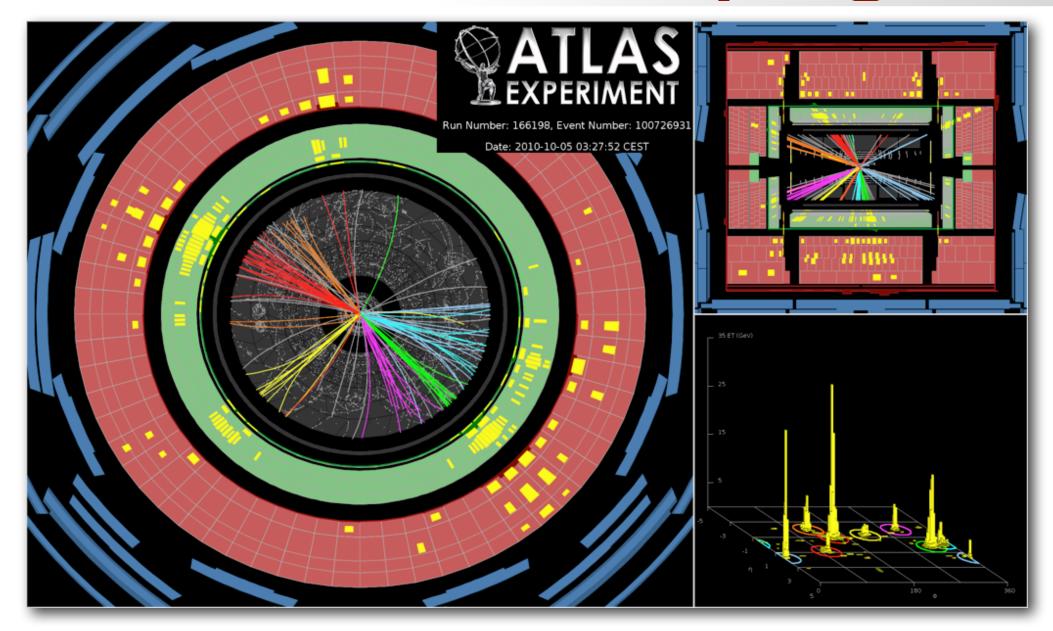


A few decades after PETRA and LEP, the event displays got prettier, but jets are still pretty much the same



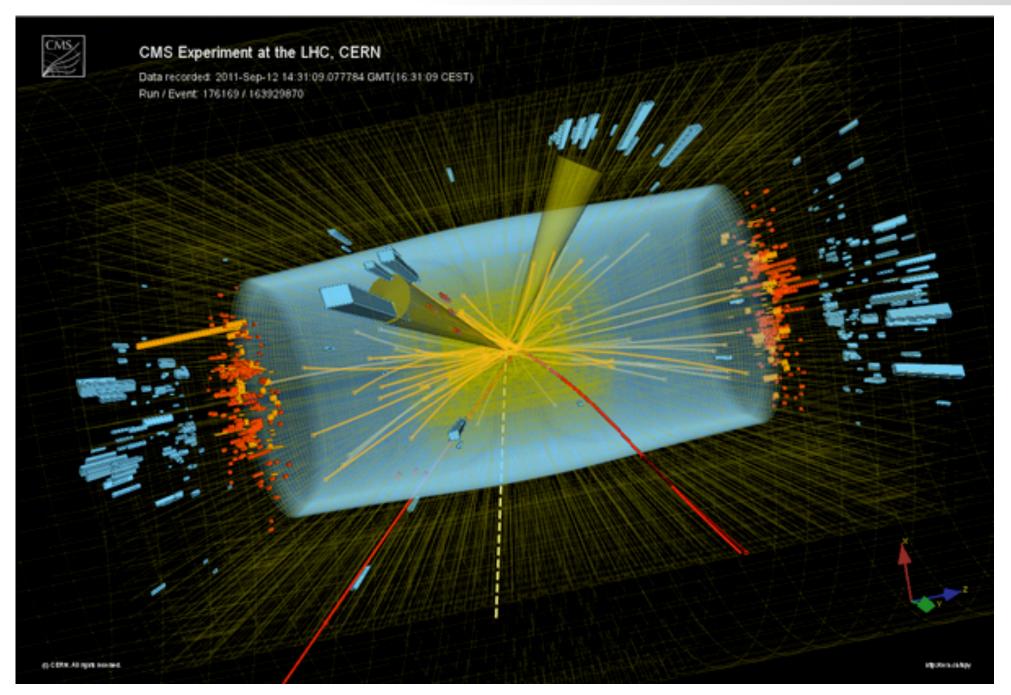
Dijet event from CMS

Jets @ LHC

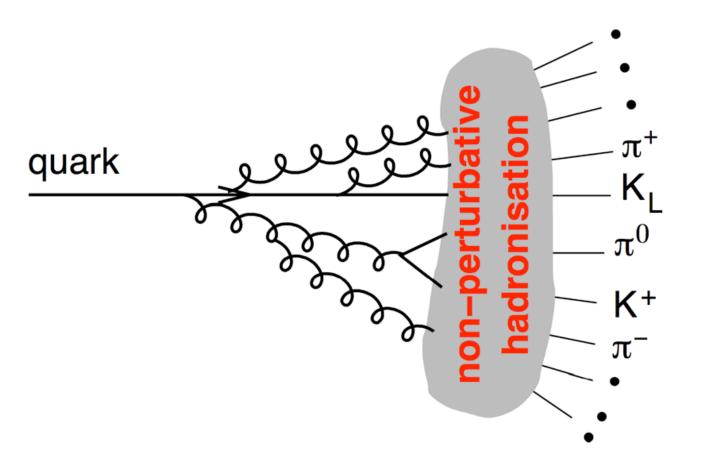


8(!) jets event from ATLAS

Jets @ LHC



Why do jets happen?



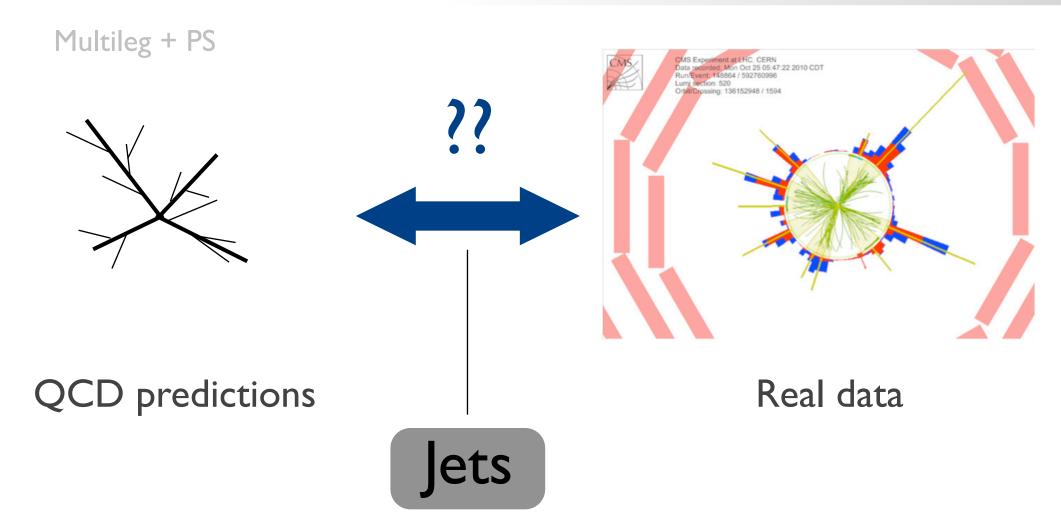
Gluon emission

$$\int \alpha_s \frac{dE}{E} \frac{d\theta}{\theta} \gg 1$$

Non-perturbative physics

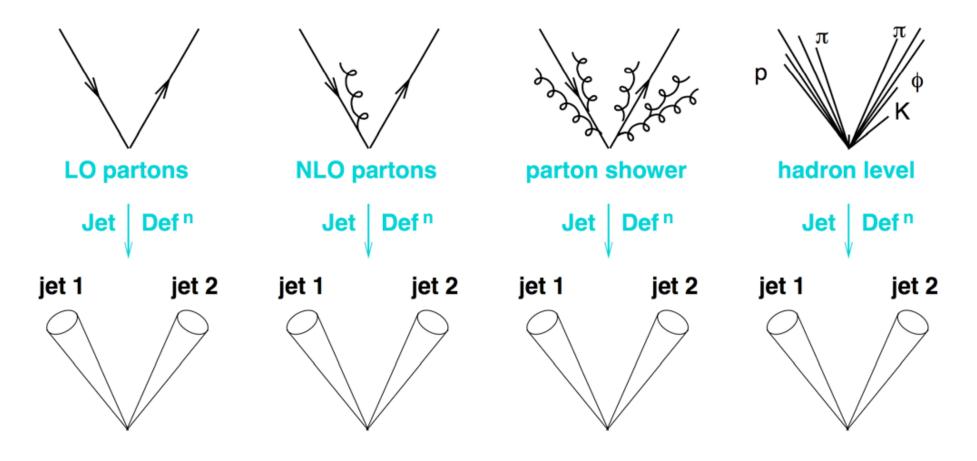
$$\alpha_s \sim 1$$

Taming reality



One purpose of a 'jet clustering' algorithm is to reduce the complexity of the final state, simplifying many hadrons to simpler objects that one can hope to calculate

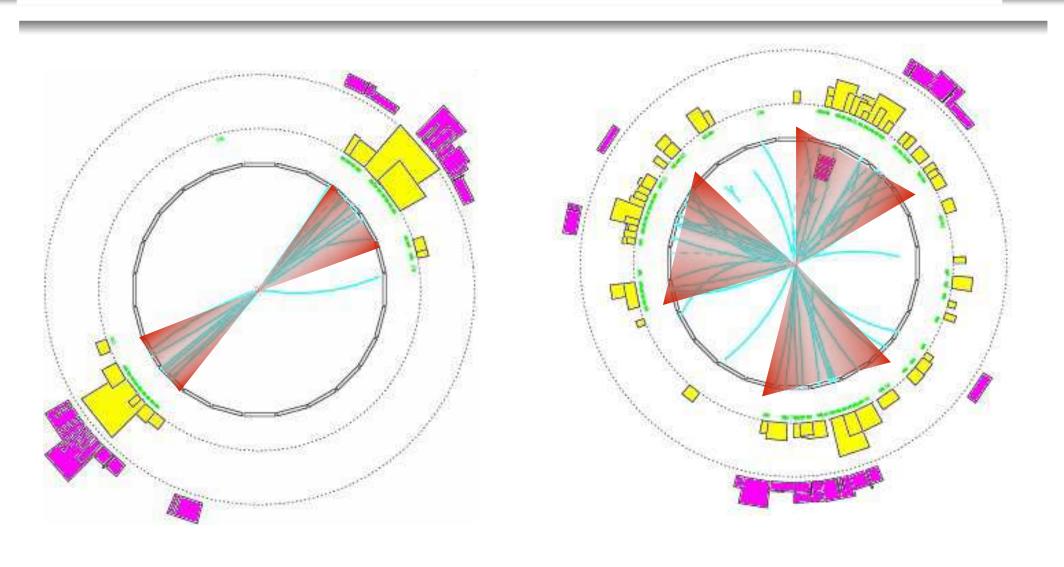
Jet definitions as projections



Projection to jets should be resilient to QCD effects

NB: projections are NOT unique: a jet is NOT EQUIVALENT to a parton

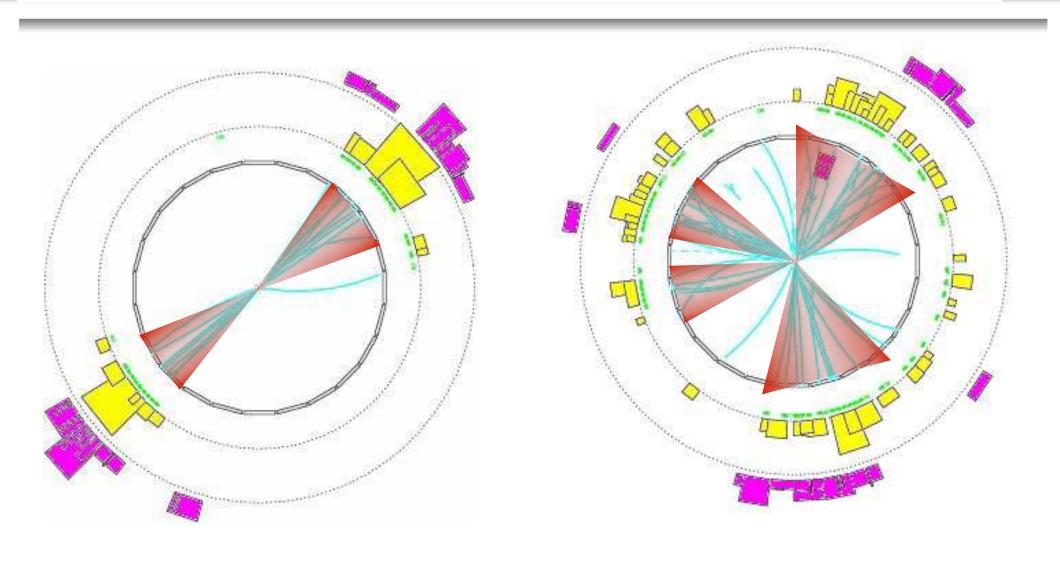




2 clear jets

3 jets?



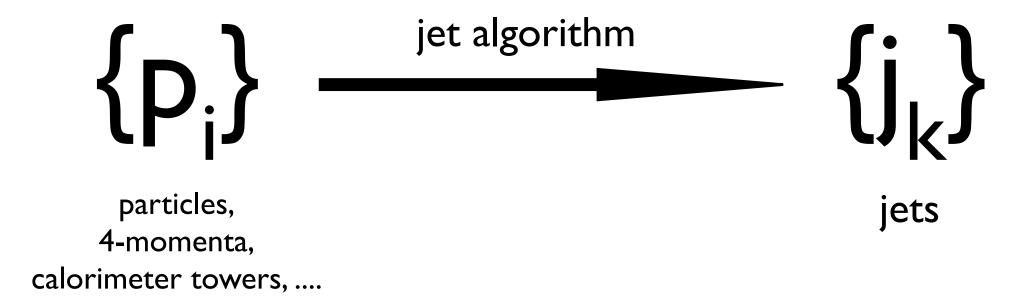


2 clear jets

3 jets? or 4 jets?

Jet clustering algorithm

A **jet algorithm** maps the momenta of the final state particles into the momenta of a certain number of jets:



Most algorithms contain a resolution parameter, **R**, which controls the extension of the jet

"Jet [definitions] are legal contracts between theorists and experimentalists"
-- M| Tannenbaum

Jets can serve two purposes

- They can be observables, that one can measure and calculate
- They can be **tools**, that one can employ to extract specific properties of the final state

Different clustering algorithms have different properties and characteristics that can make them more or less appropriate for each of these tasks

IRC safety

An observable is **infrared and collinear safe** if, in the limit of a **collinear splitting**, or the **emission of an infinitely soft** particle, the observable remains **unchanged**:

$$O(X; p_1, \dots, p_n, p_{n+1} \to 0) \to O(X; p_1, \dots, p_n)$$

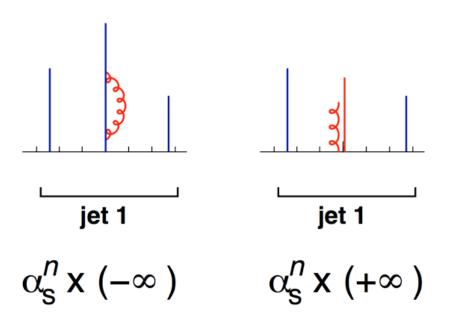
 $O(X; p_1, \dots, p_n \parallel p_{n+1}) \to O(X; p_1, \dots, p_n + p_{n+1})$

This property ensures cancellation of **real** and **virtual** divergences in higher order calculations

If we wish to be able to calculate a jet rate in perturbative QCD the jet algorithm that we use must be IRC safe: soft emissions and collinear splittings must not change the hard jets

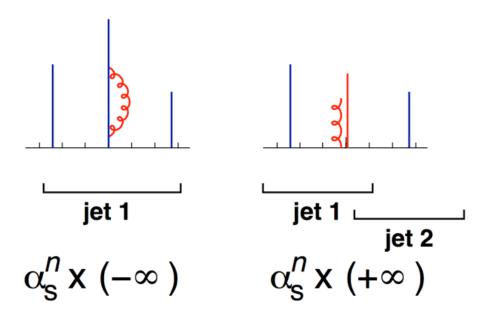
Reconstructing jets must respect rules





Infinities cancel

Collinear Unsafe



Infinities do not cancel

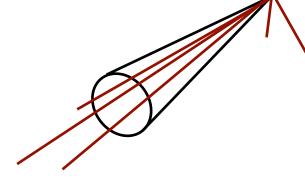
Perturbative calculations of jet observable will only be possible with collinear (and infrared) safe jet definitions

Cone algorithms

The first rigorous definition of cone jets in QCD is due to Sterman and Weinberg

Phys. Rev. Lett. **39**, 1436 (1977)

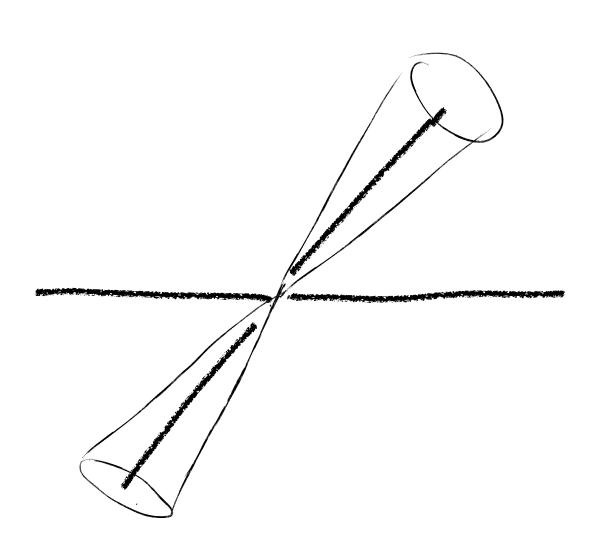
To study jets, we consider the partial cross section $\sigma(E,\theta,\Omega,\epsilon,\delta) \text{ for } e^+e^- \text{ hadron production events, in which all but}$ a fraction $\epsilon <<1$ of the total e^+e^- energy E is emitted within some pair of oppositely directed cones of half-angle $\delta <<1$, lying within two fixed cones of solid angle Ω (with $\pi\delta^2 <<\Omega <<1$) at an angle θ to the e^+e^- beam line. We expect this to be measur-



Two-jet rate:

$$\sigma(E,\theta,\Omega,\varepsilon,\delta) = \left(\frac{d\sigma}{d\Omega}\right)_0 \Omega \left[1 - \left(\frac{g_E^2}{3\pi^2}\right) \left\{3\ln\delta + 4\ln\delta \ln 2\varepsilon + \frac{\pi^3}{3} - \frac{5}{2}\right\}\right]$$

2 particles = 2 jets



3 particles =) collinear large angle

Jet algorithms

The Sterman-Weinberg definition is "inclusive enough" for IRC safety

Good for 2 jets and e+e- collisions

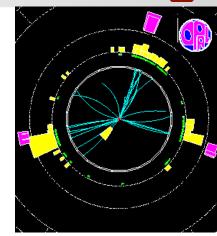
What happens in a more general case, where more than two jets are likely to exist?

Where do we place the cones? How many?

Iterative jet algorithms

Two main approaches to jet clustering

I. Find regions where a lot of energy flows



2. Decide which particles are "close", aggregate them

In HEP these are usually called **cone** and **sequential recombination** algorithms respectively

(in other fields they are often called partitional-type clustering and agglomerative hierarchical clustering)

Two main classes of jet algorithms

Sequential recombination algorithms

Bottom-up approach: combine particles starting from closest ones

How? Choose a **distance measure**, iterate recombination until few objects left, call them jets

Works because of mapping closeness \Leftrightarrow QCD divergence Examples: Jade, k_t , Cambridge/Aachen, anti- k_t ,

Usually trivially made IRC safe, but their algorithmic complexity scales like N³

Cone algorithms

Top-down approach: find coarse regions of energy flow.

How? Find stable cones (i.e. their axis coincides with sum of momenta of particles in it)

Works because QCD only modifies energy flow on small scales Examples: JetClu, MidPoint, ATLAS cone, CMS cone, SISCone......

Can be programmed to be fairly fast, at the price of being complex and IRC unsafe

Recombination algorithms

- ▶ First introduced in e⁺e⁻ collisions in the '80s
- ▶ Typically they work by calculating a 'distance' between particles, and then recombine them pairwise according to a given order, until some condition is met (e.g. no particles are left, or the distance crosses a given threshold)

IRC safety can usually be seen to be trivially guaranteed

JADE algorithm

Distance:

$$y_{ij} = \frac{2E_i E_j (1 - \cos \theta_{ij})}{Q^2}$$

- Find the minimum y_{min} of all y_{ij}
- If y_{min} is below some jet resolution threshold y_{cut}, recombine i and j into a single new particle ('pseudojet'), and repeat
- ▶ If no y_{min} < y_{cut} are left, all remaining particles are jets

Problem of this particular algorithm:

two **soft** particles emitted at **large angle** get easily recombined into a single jet: counterintuitive and perturbatively troublesome

e+e- kt (Durham) algorithm

[Catani, Dokshitzer, Olsson, Turnock, Webber '91]

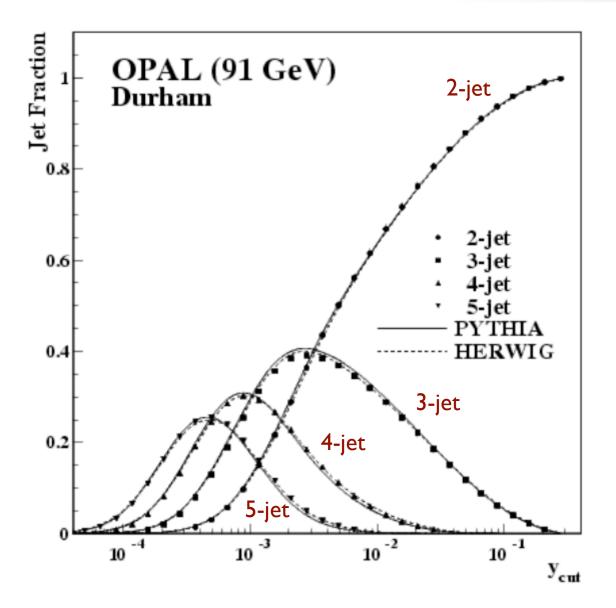
Distance:

$$y_{ij} = \frac{2\min(E_i^2, E_j^2)(1 - \cos\theta_{ij})}{Q^2}$$

In the collinear limit, the numerator reduces to the **relative transverse momentum** (squared) of the two particles, hence the name of the algorithm

- Find the minimum y_{min} of all y_{ij}
- If y_{min} is below some jet resolution threshold y_{cut}, recombine i and j into a single new particle ('pseudojet'), and repeat
- If no $y_{min} < y_{cut}$ are left, all remaining particles are jets

e+e- kt (Durham) algorithm in action



Characterise events in terms of number of jets (as a function of y_{cut})

Resummed calculations for distributions of y_{cut} doable with the k_t algorithm

e⁺e⁻ k_t (Durham) algorithm v. QCD

kt is a sequential recombination type algorithm

One key feature of the k_t algorithm is its relation to the structure of QCD divergences:

$$\frac{dP_{k\to ij}}{dE_i d\theta_{ij}} \sim \frac{\alpha_s}{\min(E_i, E_j)\theta_{ij}}$$

The yij distance is the inverse of the emission probability

- ▶ The k_t algorithm roughly inverts the QCD branching sequence (the pair which is recombined first is the one with the largest probability to have branched)
- ▶ The history of successive clusterings has physical meaning

kt algorithm in hadron collisions

(Inclusive and longitudinally invariant version)

$$d_{ij} = \min(p_{ti}^2, p_{tj}^2) \frac{\Delta R_{ij}^2}{R^2} \qquad d_{iB} = p_{ti}^2$$

- ▶ Calculate the distances between the particles: d_{ij}
- ► Calculate the beam distances: diB
- Combine particles with smallest distance d_{ij} or, if d_{iB} is smallest, call it a jet
- Find again smallest distance and repeat procedure until no particles are left (this stopping criterion leads to the inclusive version of the kt algorithm)
- Only use jets with $p_t > p_{t,min}$

The kt algorithm and its siblings

$$d_{ij} = \min(p_{ti}^{2p}, p_{tj}^{2p}) \frac{\Delta y^2 + \Delta \phi^2}{R^2} \qquad d_{iB} = p_{ti}^{2p}$$

p = I k_t algorithm

S. Catani, Y. Dokshitzer, M. Seymour and B. Webber, Nucl. Phys. B406 (1993) 187 S.D. Ellis and D.E. Soper, Phys. Rev. D48 (1993) 3160

p = 0 Cambridge/Aachen algorithm

Y. Dokshitzer, G. Leder, S.Moretti and B. Webber, JHEP 08 (1997) 001 M.Wobisch and T.Wengler, hep-ph/9907280

p = -I anti- k_t algorithm

MC, G. Salam and G. Soyez, arXiv:0802.1189

NB: in anti-kt pairs with a **hard** particle will cluster first: if no other hard particles are close by, the algorithm will give **perfect cones**

Quite ironically, a sequential recombination algorithm is the 'perfect' cone algorithm

IRC safety of generalised-kt algorithms

$$d_{ij} = \min(p_{ti}^{2p}, p_{tj}^{2p}) \frac{\Delta y^2 + \Delta \phi^2}{R^2} \qquad d_{iB} = p_{ti}^{2p}$$

p > 0

New **soft** particle $(p_t \to 0)$ means that $d \to 0 \Rightarrow$ clustered first, no effect on jets New **collinear** particle $(\Delta y^2 + \Delta \Phi^2 \to 0)$ means that $d \to 0 \Rightarrow$ clustered first, no effect on jets

p = 0

New **soft** particle $(p_t \to 0)$ can be new jet of zero momentum \Rightarrow no effect on hard jets New **collinear** particle $(\Delta y^2 + \Delta \Phi^2 \to 0)$ means that $d \to 0 \Rightarrow$ clustered first, no effect on jets

p < 0

New **soft** particle $(p_t \to 0)$ means $d \to \infty \Rightarrow$ clustered last or new zero-jet, no effect on hard jets New **collinear** particle $(\Delta y^2 + \Delta \Phi^2 \to 0)$ means that $d \to 0 \Rightarrow$ clustered first, no effect on jets

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k _t	SR $d_{ij} = \min(p_{ti}^{2}, p_{tj}^{2}) \Delta R_{ij}^{2}/R^{2}$ hierarchical in rel P_{t}	Catani et al '91 Ellis, Soper '93	NInN
Cambridge/ Aachen	SR $d_{ij} = \Delta R_{ij}^2/R^2$ hierarchical in angle	Dokshitzer et al '97 Wengler, Wobish '98	NInN
anti-k _t	SR $d_{ij} = min(p_{ti}^{-2}, p_{tj}^{-2}) \Delta R_{ij}^{2}/R^{2}$ gives perfectly conical hard jets	MC, Salam, Soyez '08 (Delsart, Loch)	N 3/2
SISCone	Seedless iterative cone with split-merge gives 'economical' jets	Salam, Soyez '07	N ² InN

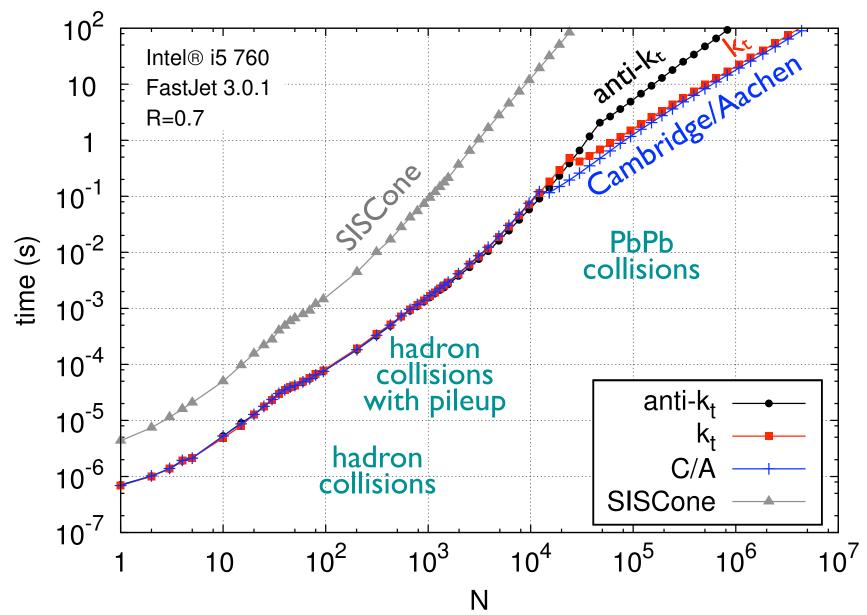
'second-generation' algorithms

All are available in FastJet, http://fastjet.fr

(As well as many IRC unsafe ones)

FastJet speed

Time needed to cluster an event with N particles



Anti-kt in action

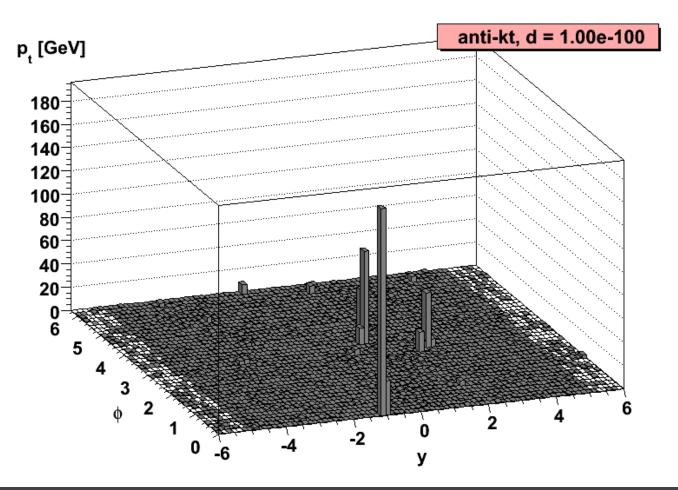
Clustering grows around hard cores

$$d_{ij} = \frac{1}{\max(p_{ti}^2, p_{tj}^2)} \frac{\Delta R_{ij}^2}{R^2}, \quad d_{iB} = \frac{1}{p_{ti}^2}$$

Anti-kt in action

Clustering grows around hard cores

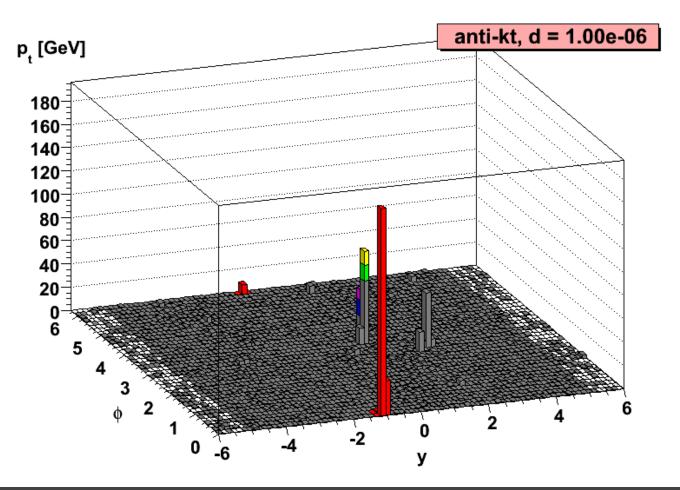
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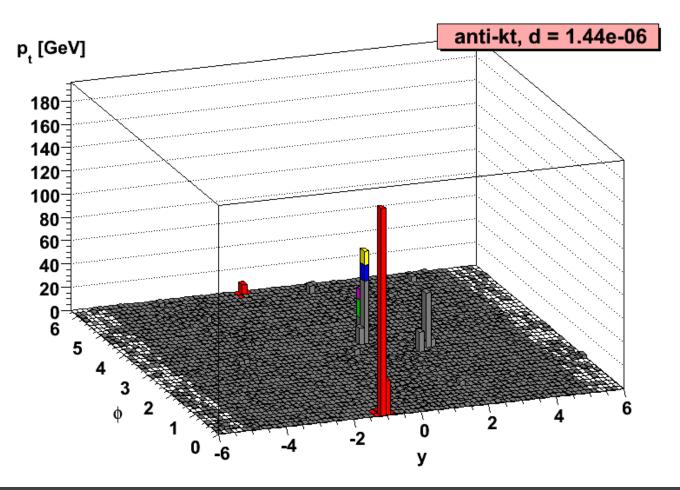
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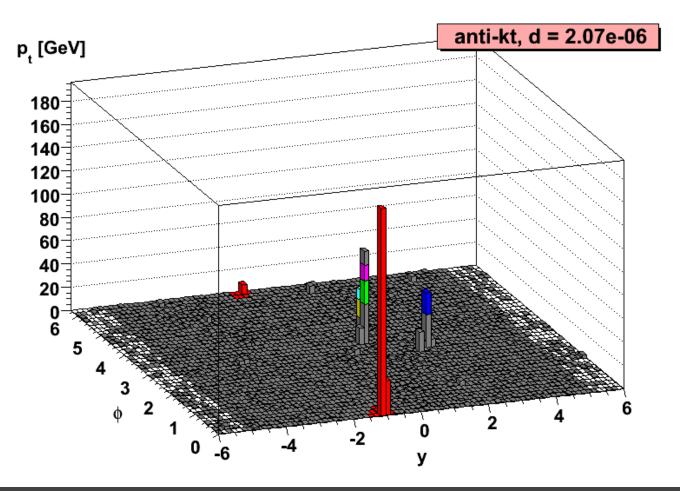
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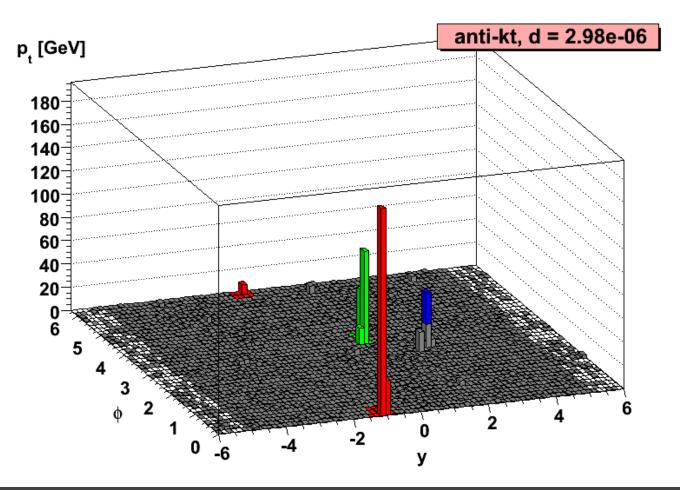
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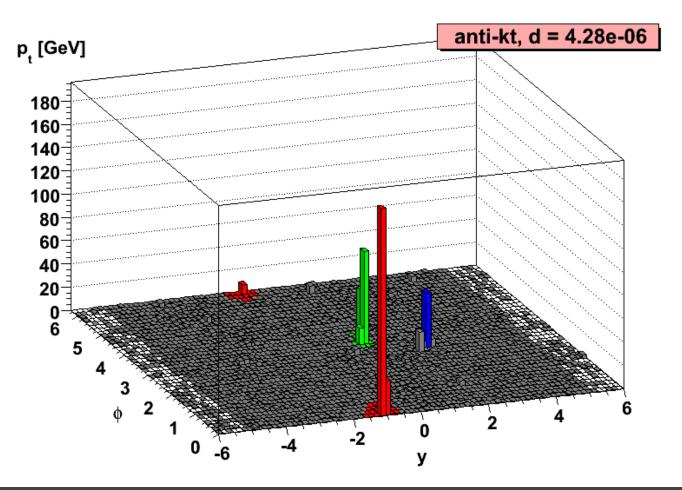
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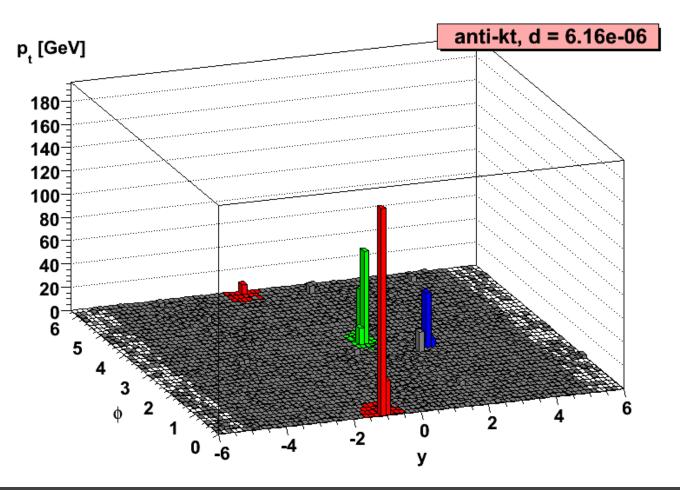
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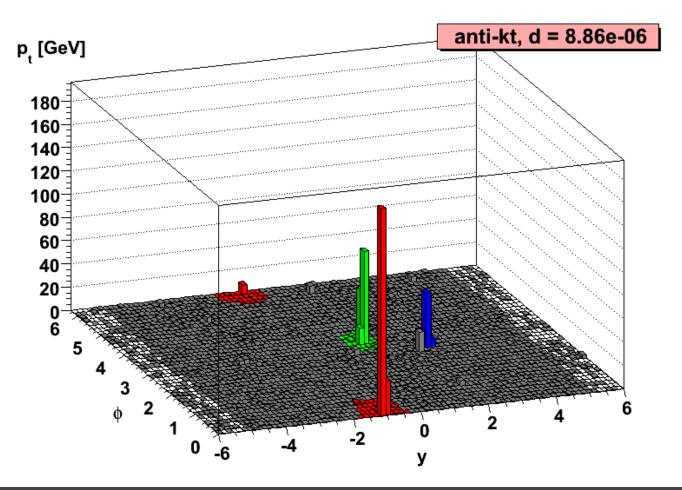
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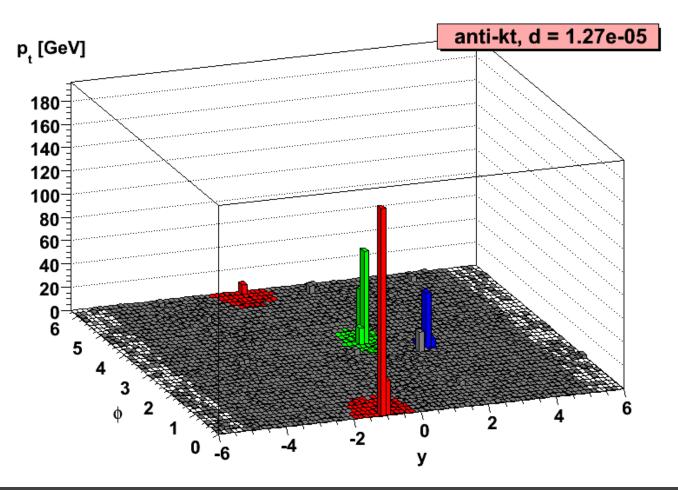
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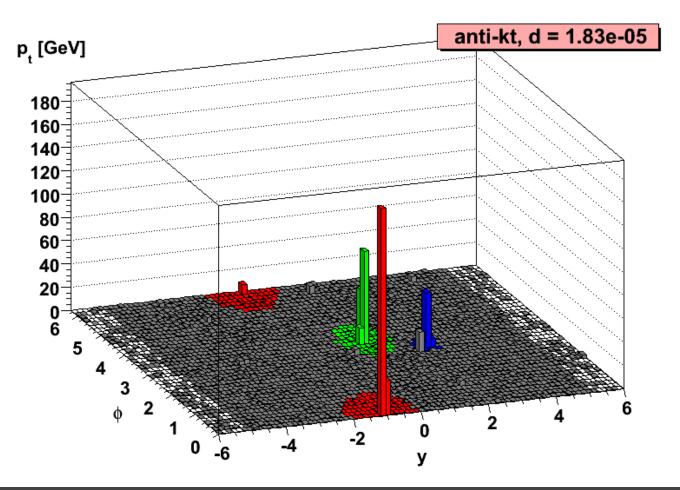
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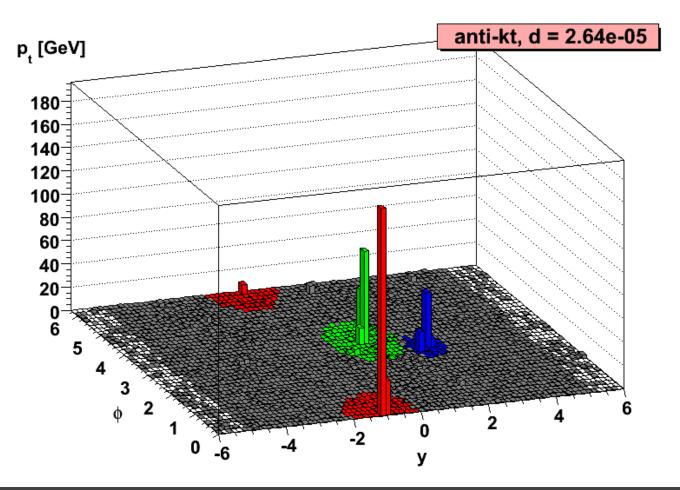
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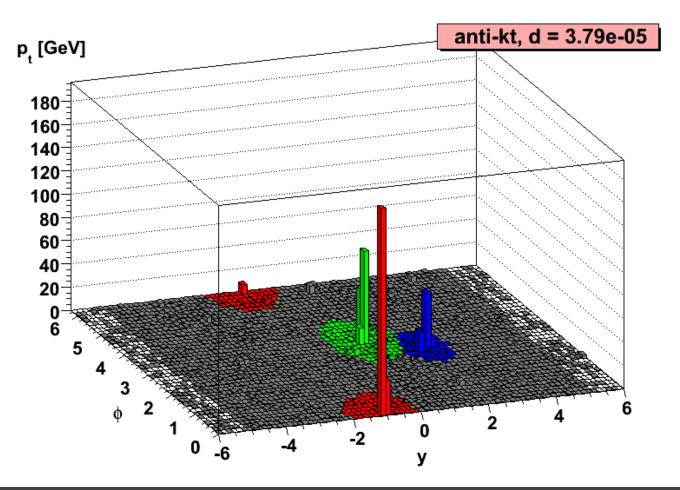
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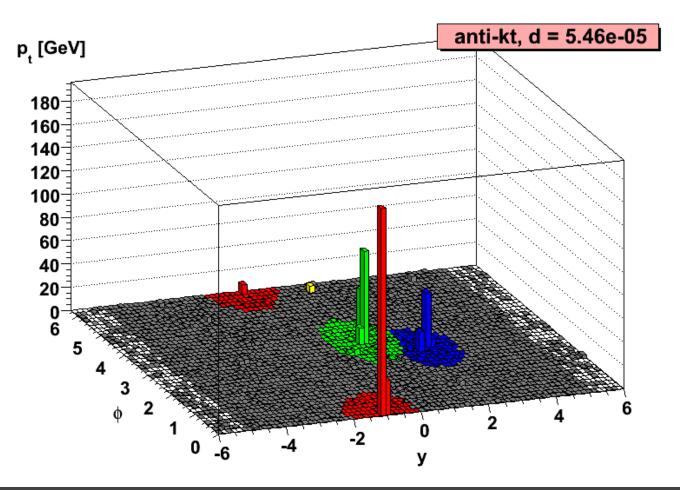
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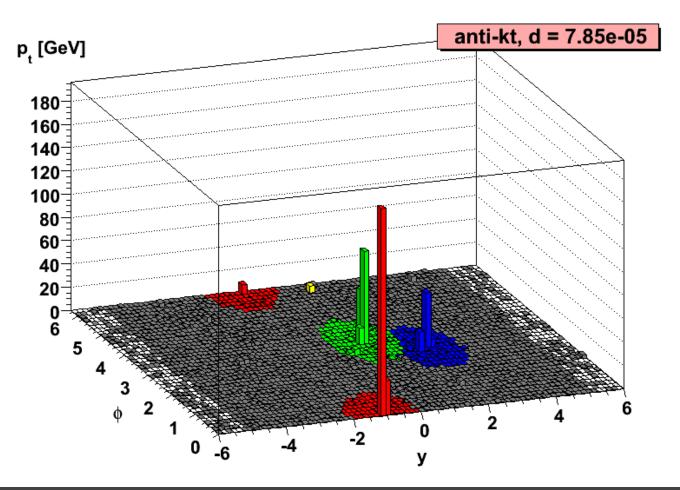
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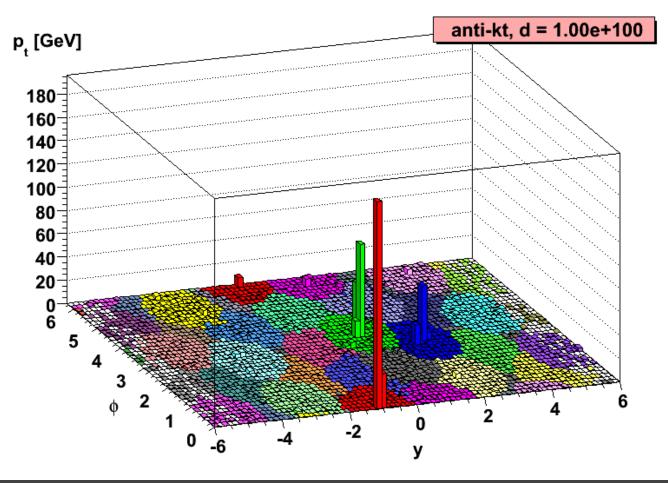


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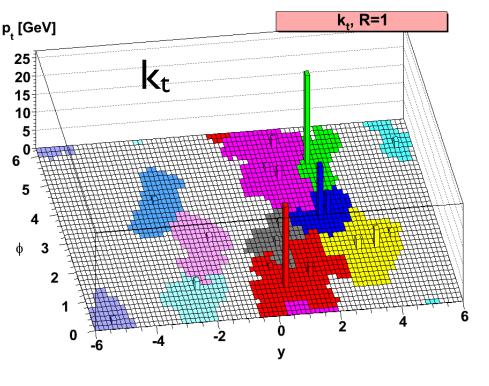


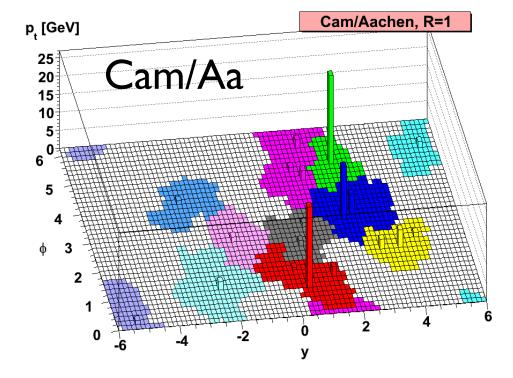
Clustering grows around hard cores

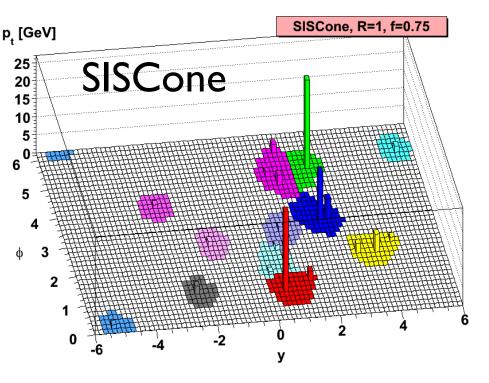
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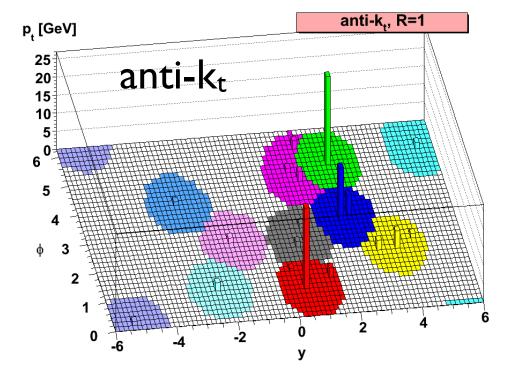


Anti-k_t gives circular jets ("cone-like") in a way that's infrared safe

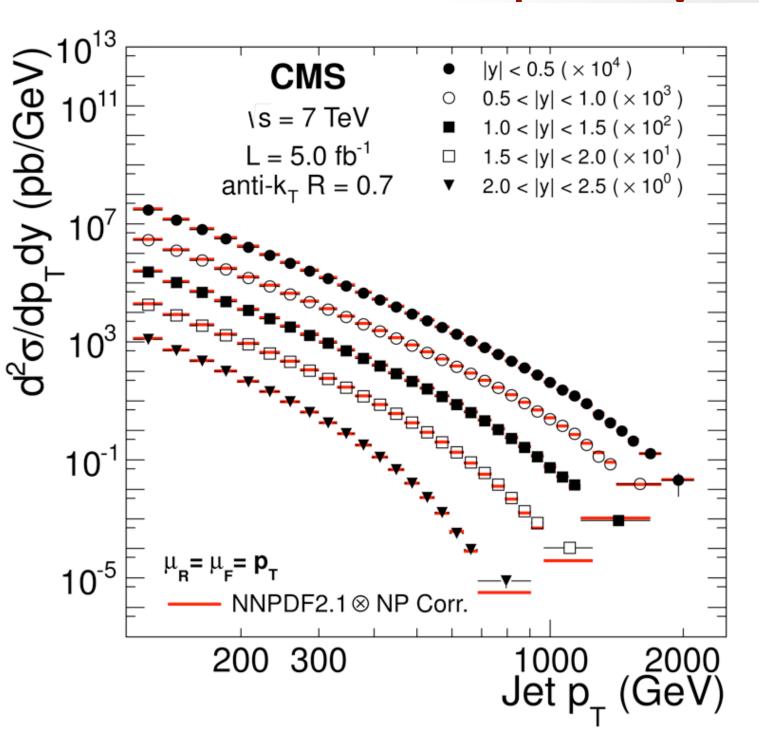








Example of jet observable



Inclusive jet cross section

Excellent
theory-data
agreement over
many orders of
magnitude

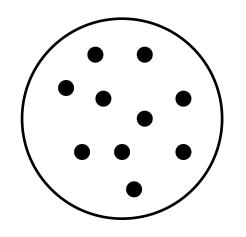
Take home points

- A vast zoology of jet algorithms has been reduced in the past few years to 4 infrared and collinear safe algorithms
 - ▶ All are implemented in an efficient and fast way
 - ▶ Of these, **anti-k**_t is used by all the LHC collaborations as their main algorithm for "finding" jets and measuring inclusive cross sections
- ▶ The four algorithms have quite different characteristics, which makes them non easily swappable when specific properties are needed for specific tasks. On the other hand, chances are that one can chose the algorithm which is most appropriate for a specific job

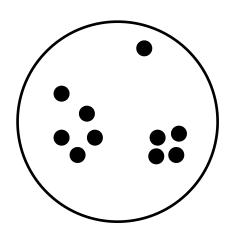
At the end of a jet finding (i.e. clustering) procedure, a jet is a **collection of constituents** to which we assign a 4-momentum

(related to the sum of the 4-momenta of the constituents)

What is the **arrangement** of the constituents inside the jet?



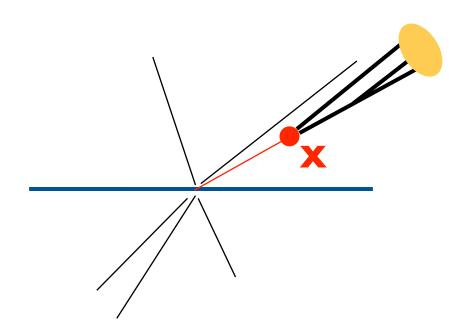


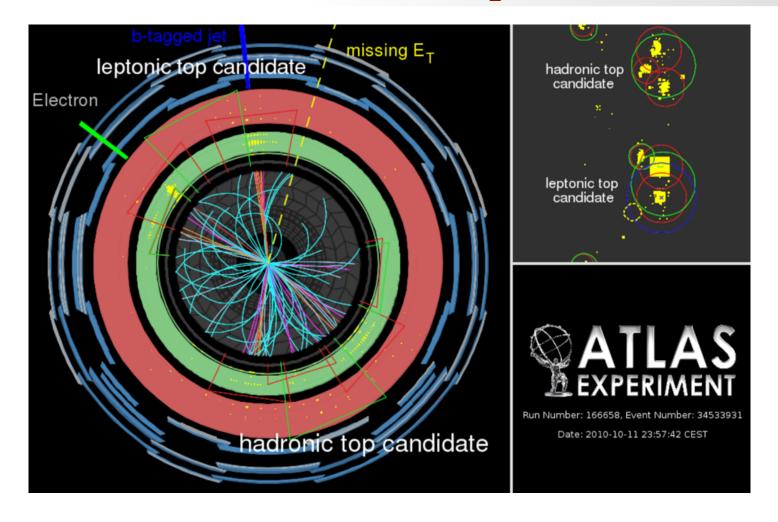


First studied by Mike Seymour in the early '90s to distinguish W jets from QCD jets

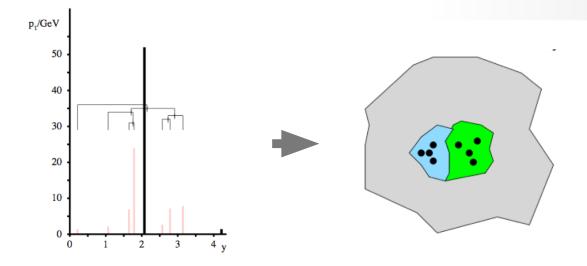
Topic revived about fifteen years ago in order to study boosted objects

[Butterworth, Davison, Rubin, Salam, 0802.2470]





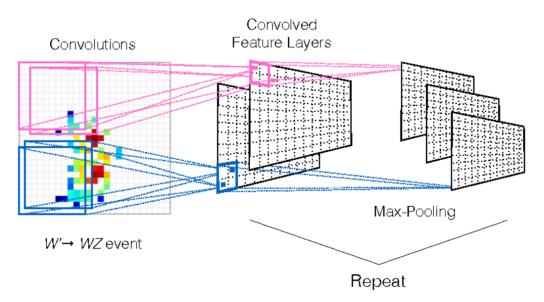
The past fifteen years have seen en explosion in jet substructure studies, i.e. how radiation is arranged within jets, and what it can tell us



Jet declustering

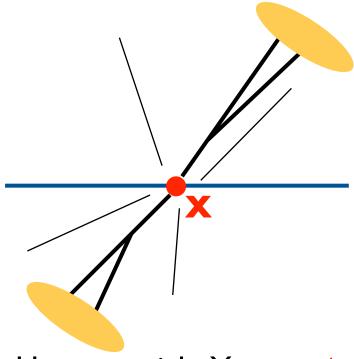
Jet shapes

(calculate a function from radiation distribution



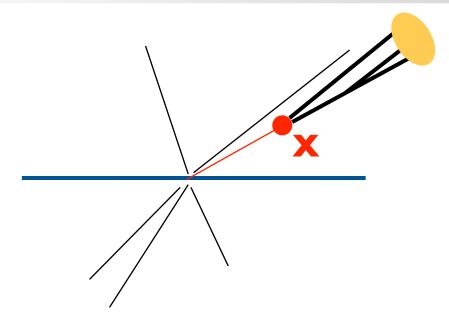
Machine learning

Why boosted objects



Heavy particle X at **rest**

Easy to resolve jets and calculate invariant mass, but signal very likely swamped by background (eg H→bb v. tt →WbWb)

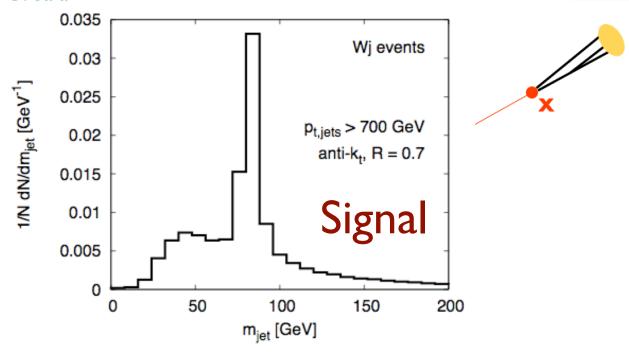


Boosted heavy particle X

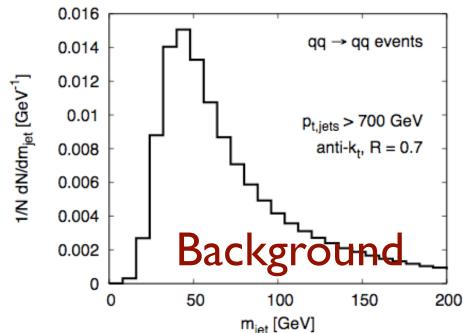
Cross section very much reduced, but acceptance better and some backgrounds smaller/ reducible

Mass of a single jet

G. Salam



A heavy object decaying into a single jet naturally gives it a mass...

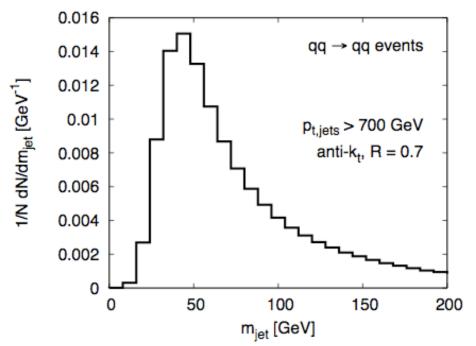


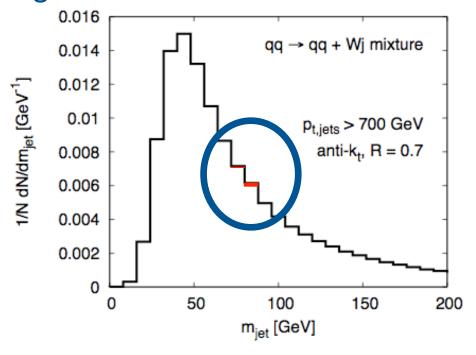
... but pure QCD jets can be massive too:

$$rac{dN}{d\ln m} \sim lpha_{
m s} \ln rac{p_t R}{m} imes {
m Sudakov}$$

Mass of a single jet

Summing 'signal' and 'background' (with appropriate cross sections) shows how much the background dominates





Background only

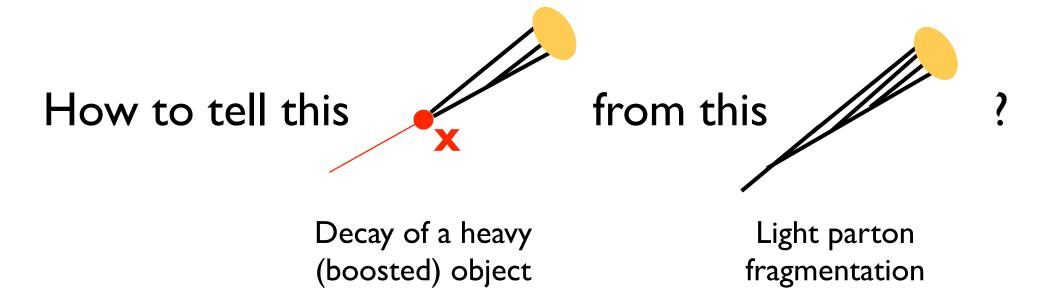
Signal + background

Practically identical

This means that one can't rely on the invariant mass only.

An appropriate strategy must be found to reduce the background and enhance the signal

Tagging



Tagging and Grooming

- The substructure of a jet can be exploited to
 - ▶ **tag** a particular structure inside the jet, i.e. a massive particle
 - ▶ First examples: Higgs (2-prong decay), top (3-prong decay)
 - remove background contamination from the jet or its components, while keeping the bulk of the perturbative radiation (often generically denoted as **grooming**)
 - ▶ First examples: filtering, trimming, pruning

Nomenclature

Groomer

procedure that always returns an output jet (i.e. it only subtracts uncorrelated 'UE/pileup' radiation from it. This is used to "clean" the jets from radiation largely unrelated to the fragmentation of the particle of interest)

Tagger

procedure that might not return an output jet (i.e. it either tags a heavy particle originating the jet or returns zero. This is used to identify a specific particle originating the jet.)

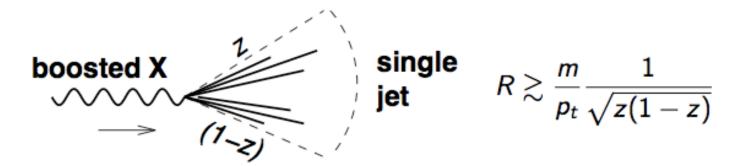
In practice, this classification is not always followed.

In some cases it also denoted a 'tagger' a procedure that rejects background jets more often than signal jets

Why substructure

Scales: $m \sim 100 \text{ GeV}$, $p_t \sim 500 \text{ GeV}$

(e.g. electroweak particle from decay of ~ ITeV BSM particle)



- ▶ need small R (< $2m/p_t \sim 0.4$) to resolve two prongs
- ▶ need large R (> \sim 3m/p_t \sim 0.6) to cluster into a single jet

Possible strategies

- ▶ Use large R, get a single jet : background large
- ▶ Use small R, resolve the jets : what is the right scale?
 - Also: small jets lead to huge combinatorial issues

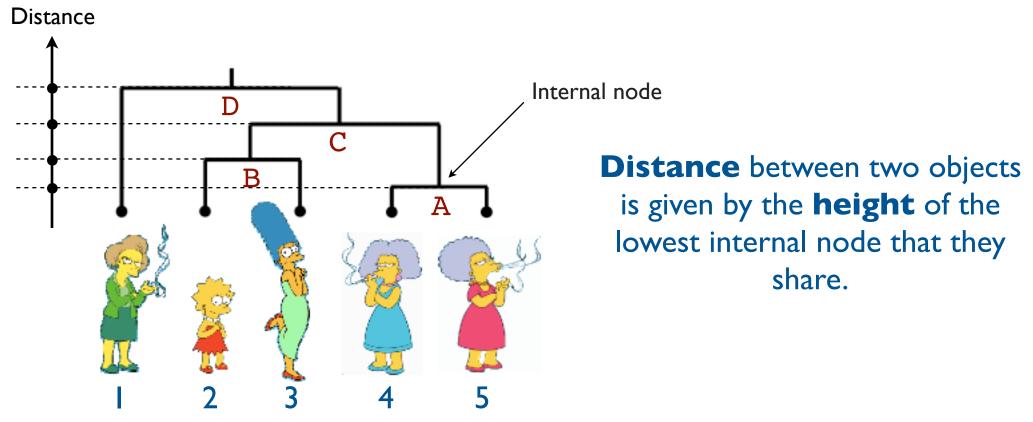
Let an algorithm find the 'right' substructure

What jets to use for substructure?

Different jet algorithms will give different 'pictures' of what's inside a jet

Dendrogram

Used to represent graphically the sequence of clustering steps in a sequential recombination algorithm



Order of clustering here is A, B, C, D

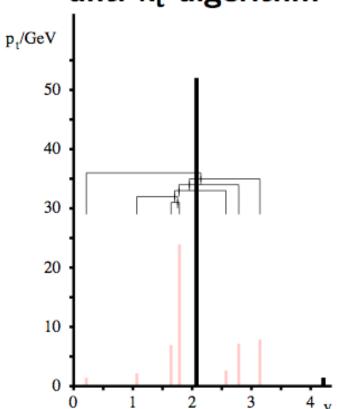
The clustering sequence is 4-5 (A), 2-3 (B), 23-45 (C), 1-2345 (D)

First try

anti-kt

Hierarchical substructure

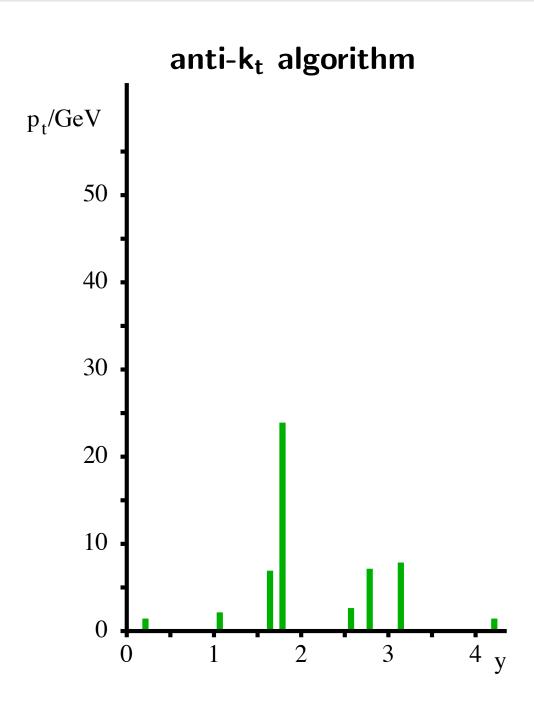




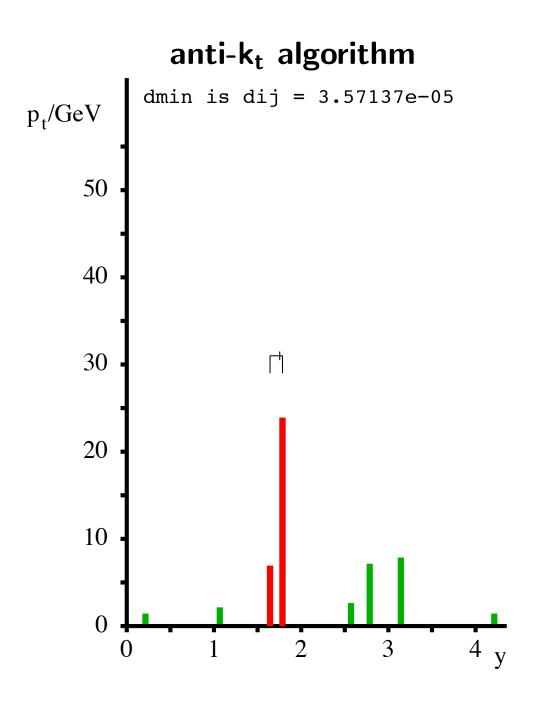
Anti-kt distance measure

$$d_{ij} = \min\left(\frac{1}{p_{ti}^2}, \frac{1}{p_{tj}^2}\right) \frac{\Delta y^2 + \Delta \phi^2}{R^2}$$

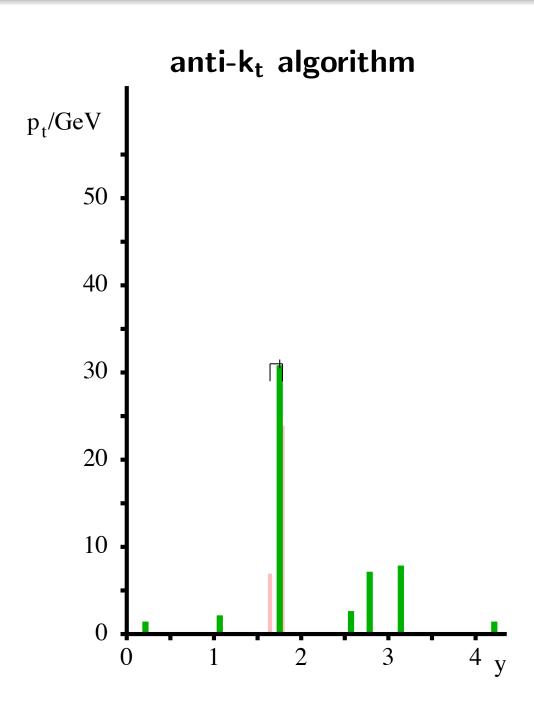
Cluster by merging to the **hardest/closest** particle



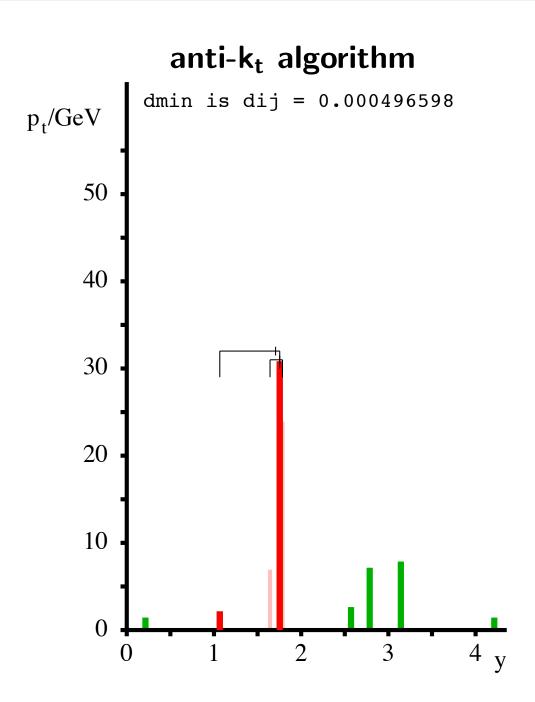
How well can an algorithm identify the "blobs" of energy inside a jet that come from different partons?



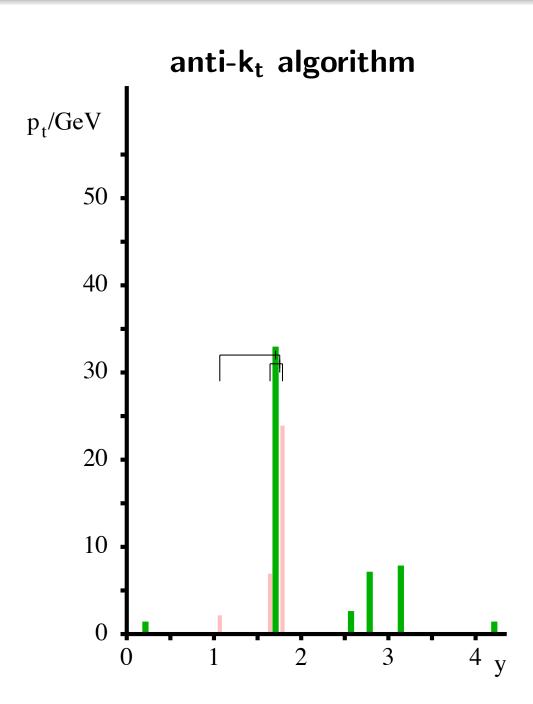
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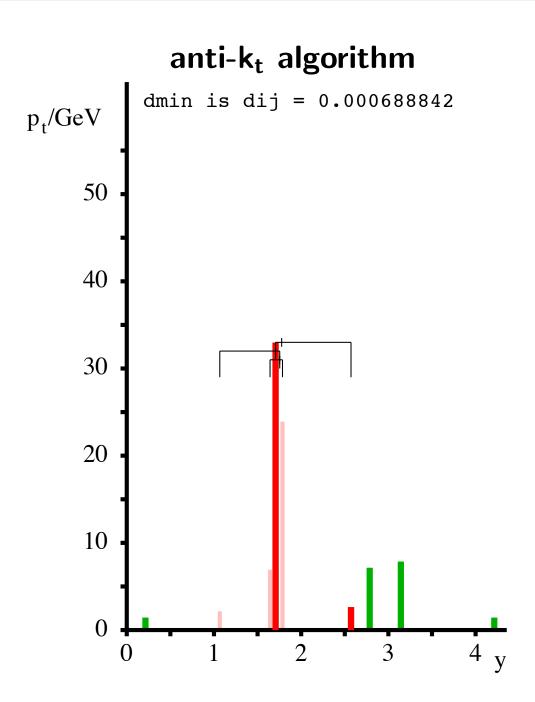
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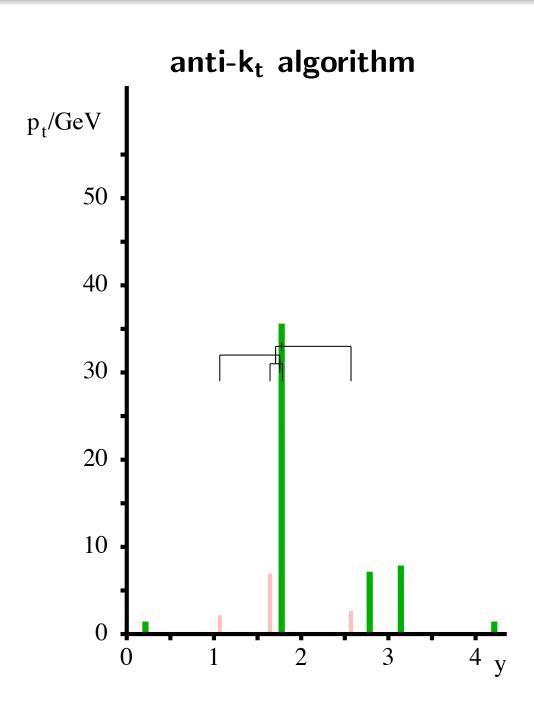
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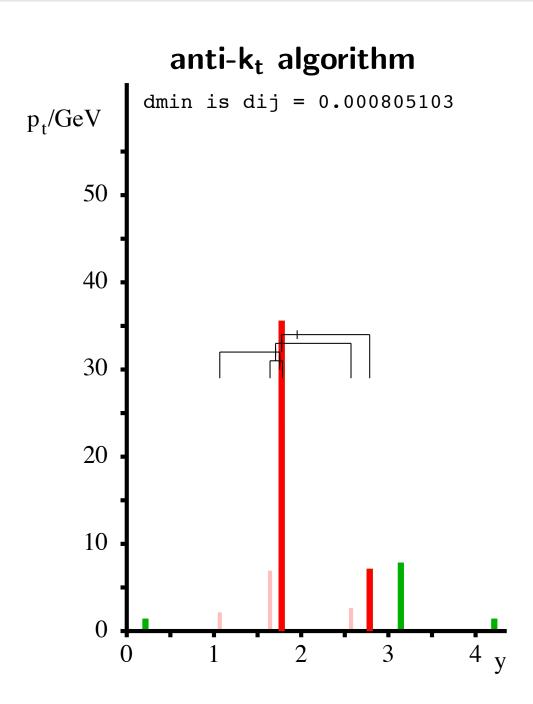
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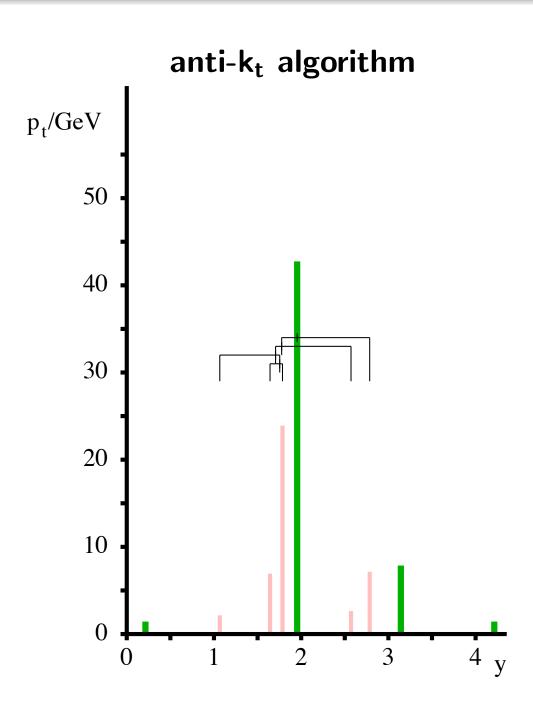


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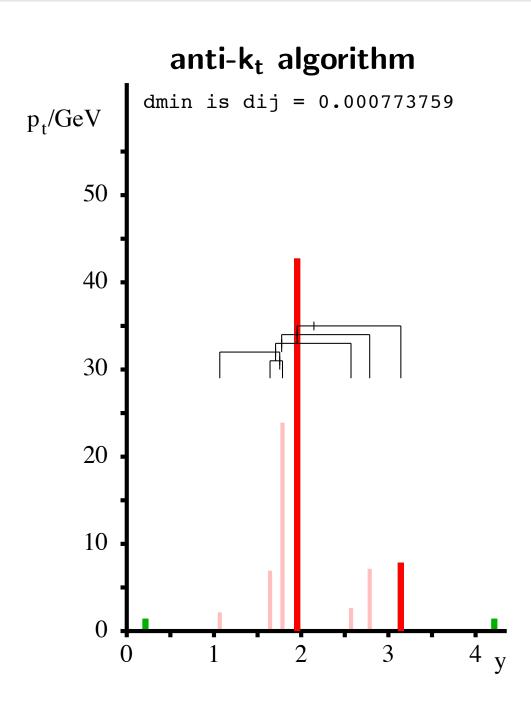
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This is crucial for identifying the kinematic variables of the partons in the jet (e.g. z).



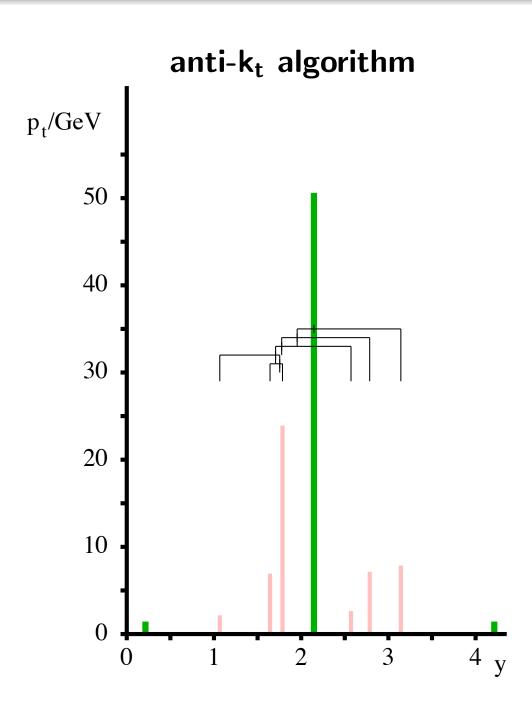
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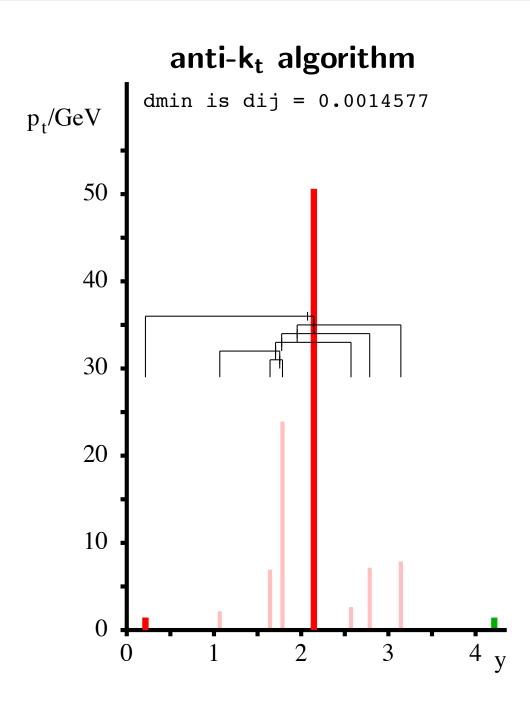
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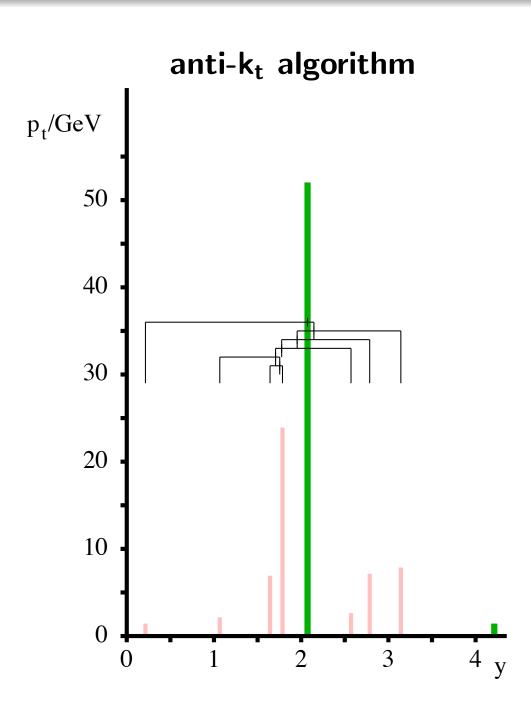
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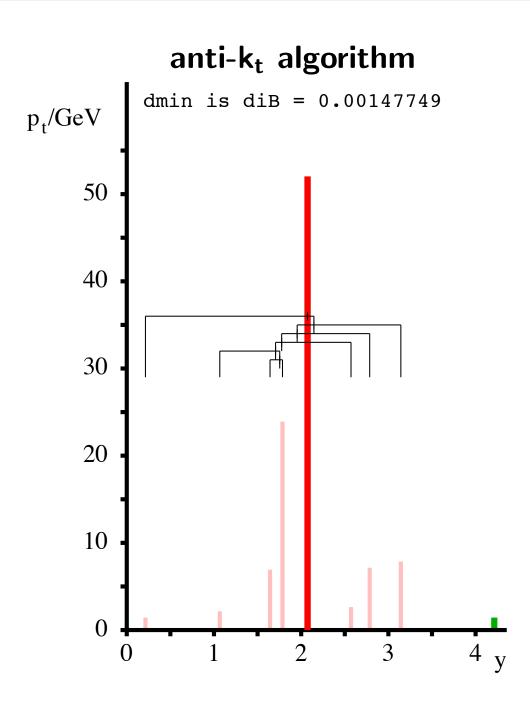
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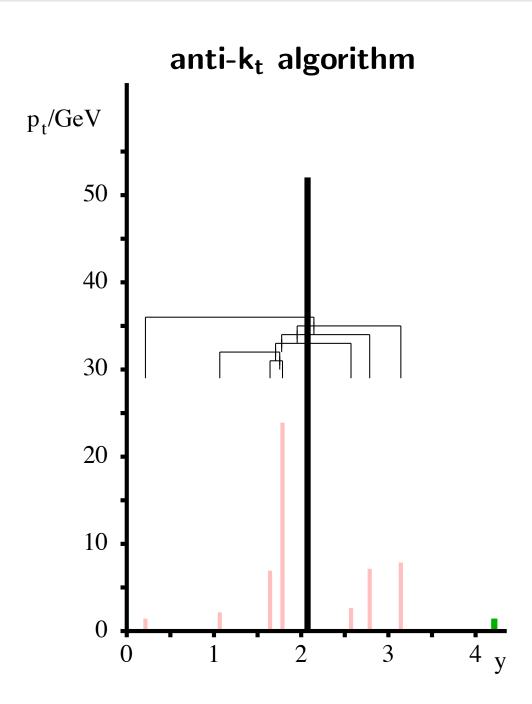
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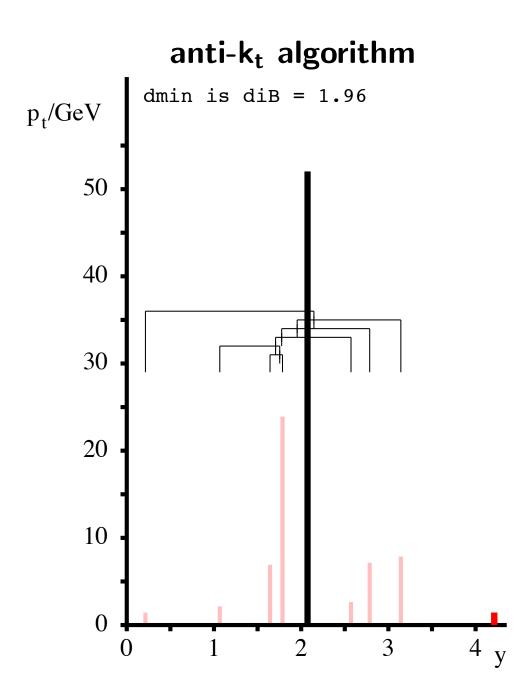
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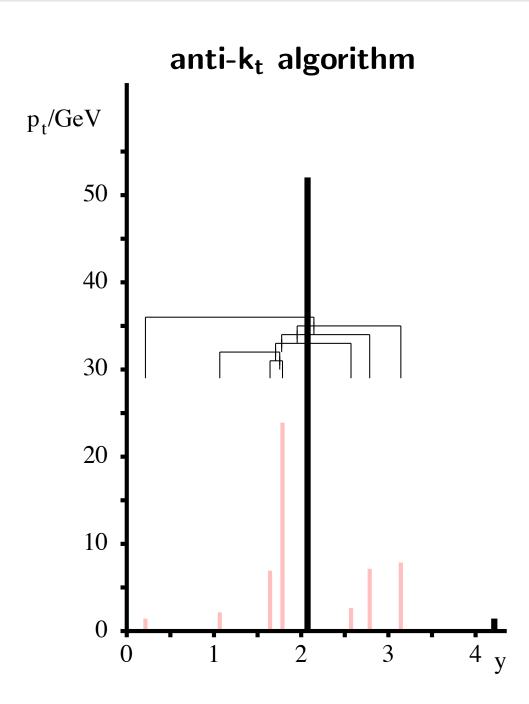
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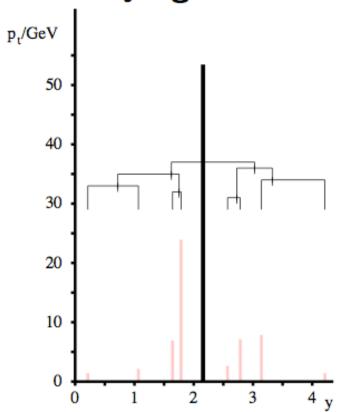
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Second try

kt

Hierarchical substructure

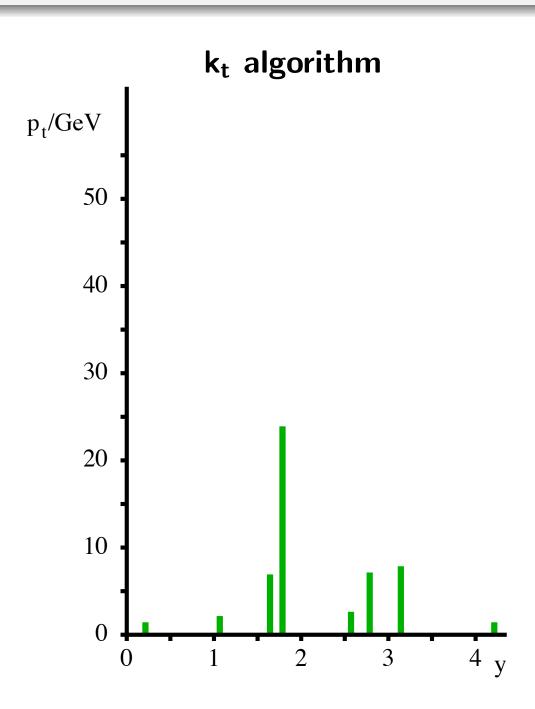




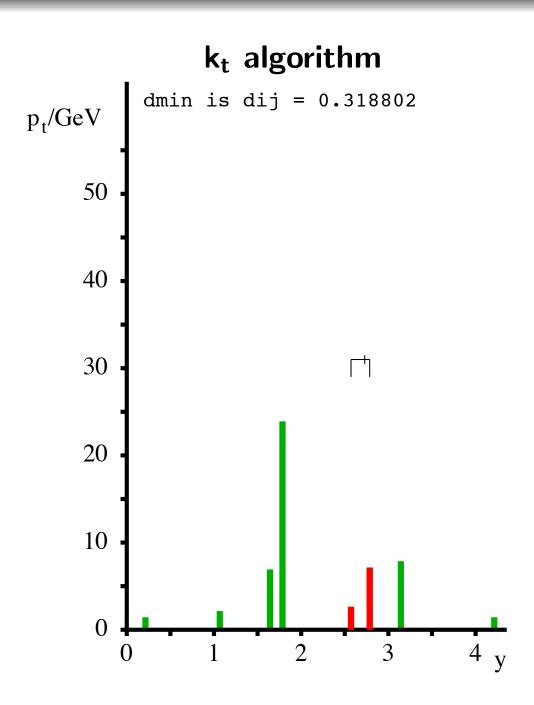
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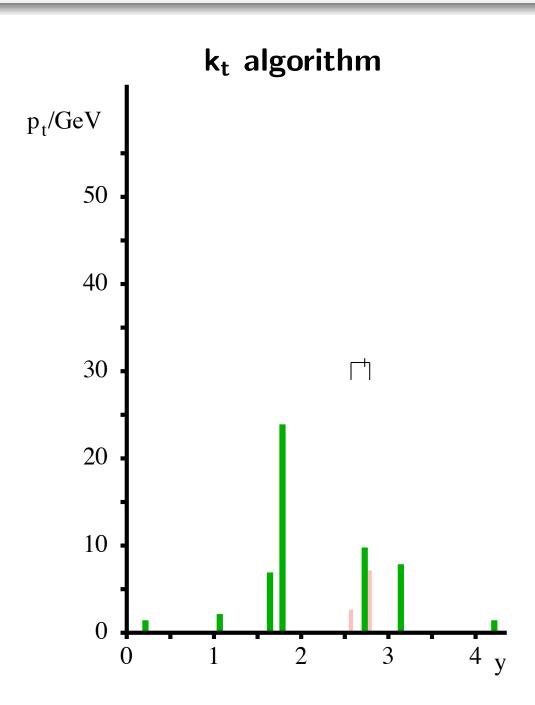
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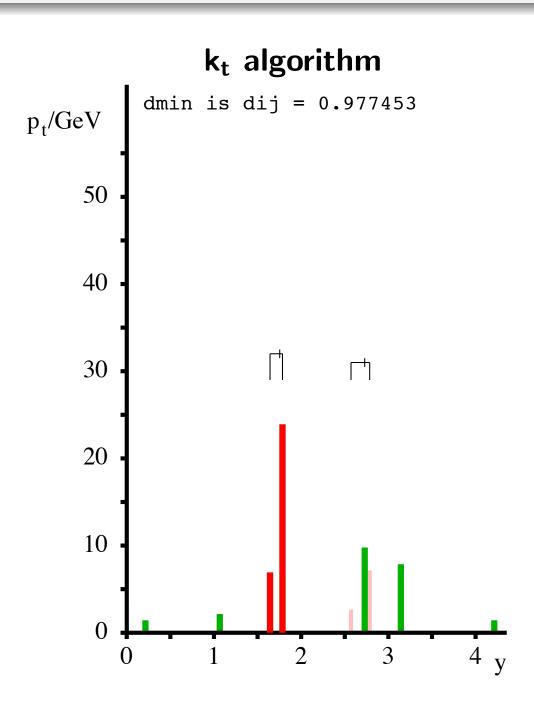
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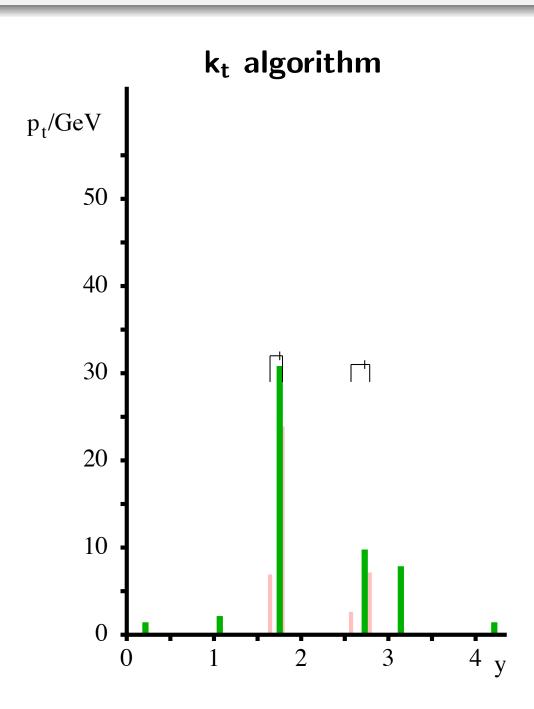
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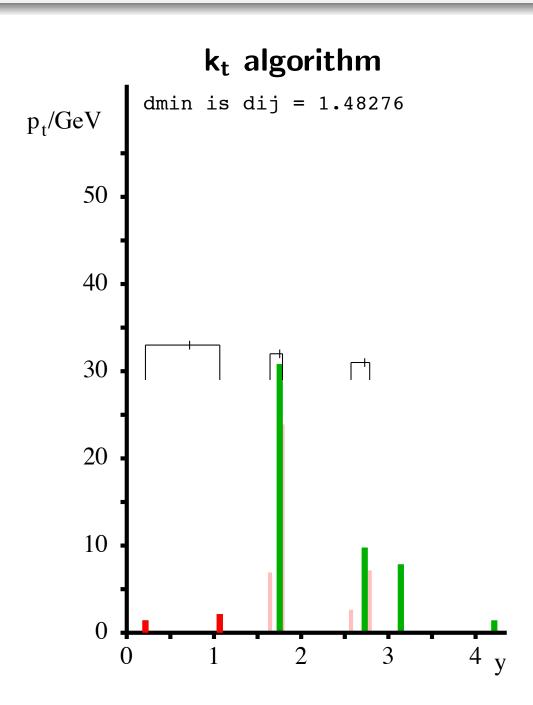
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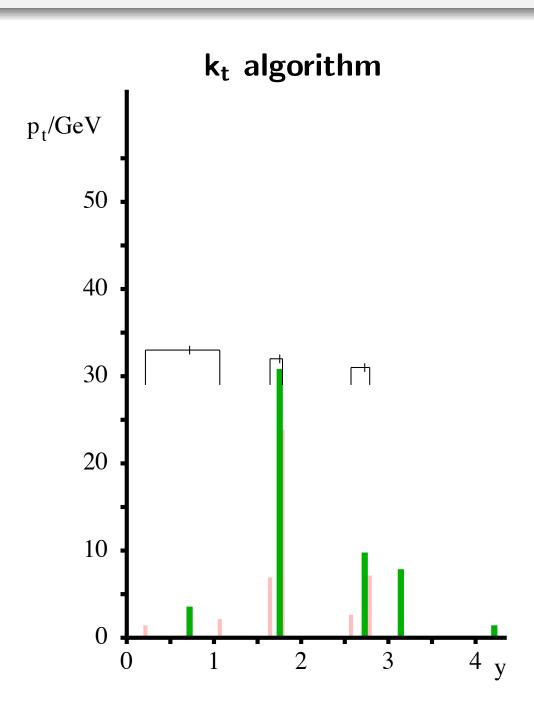


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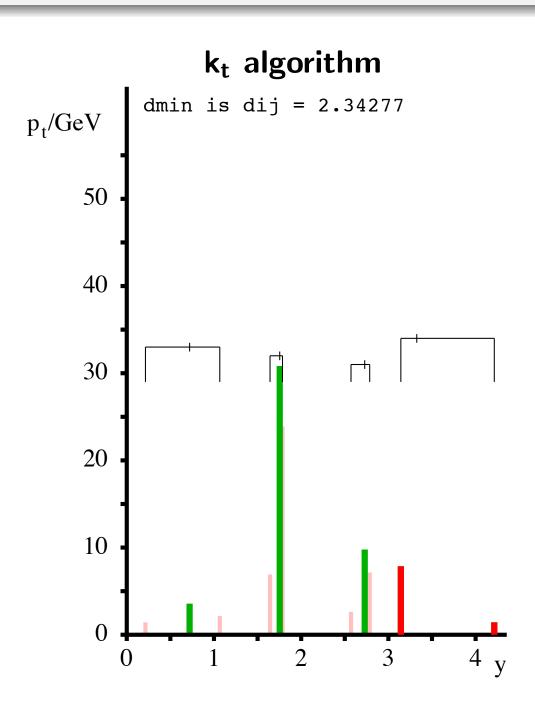
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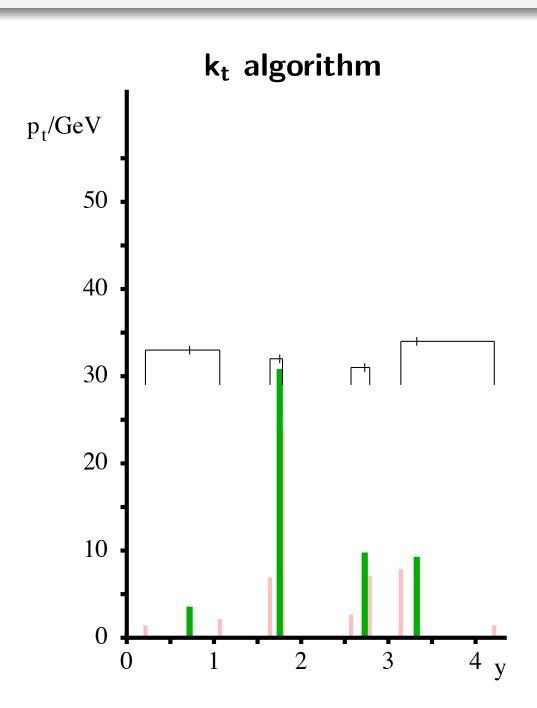
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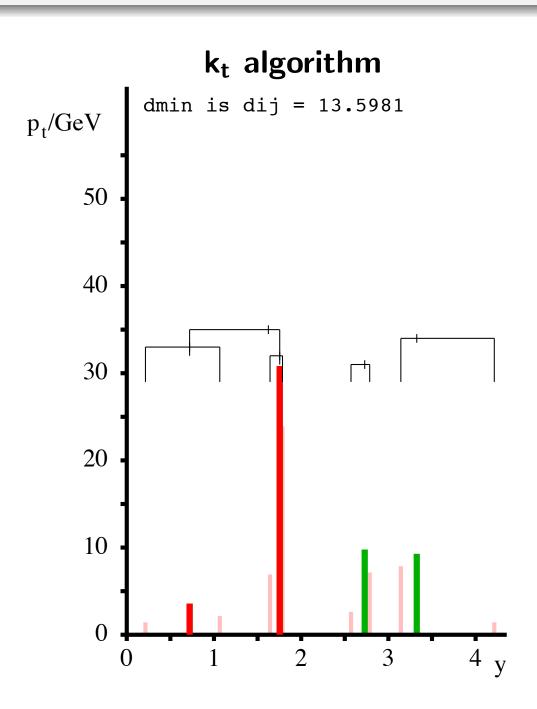
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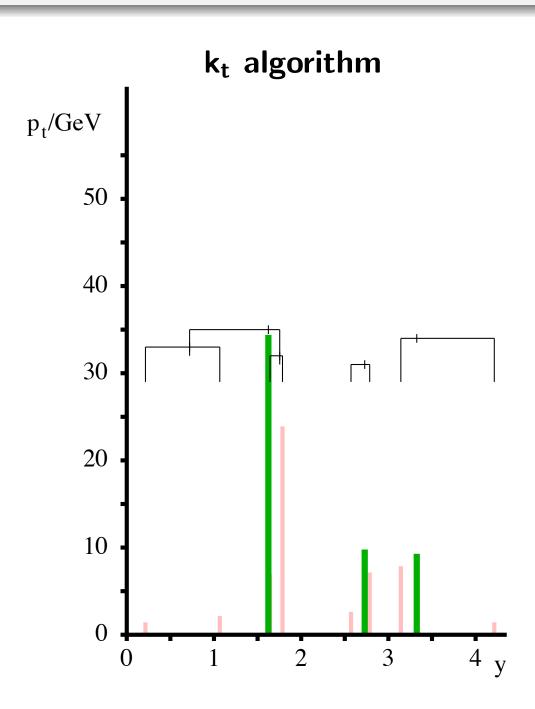
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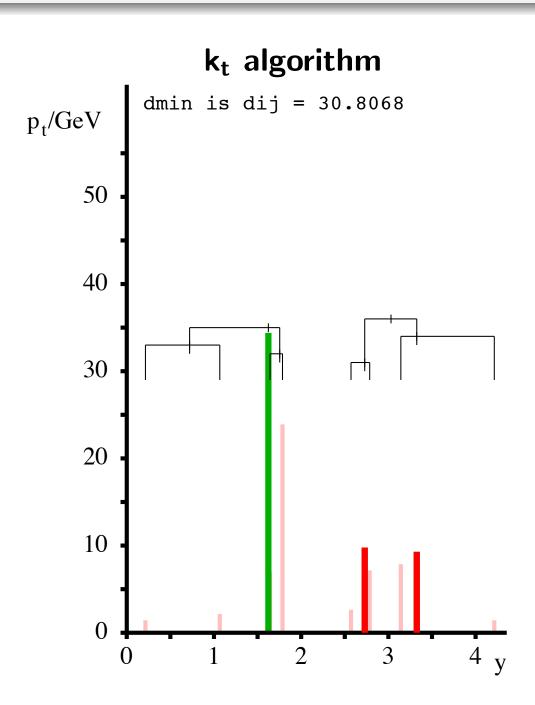
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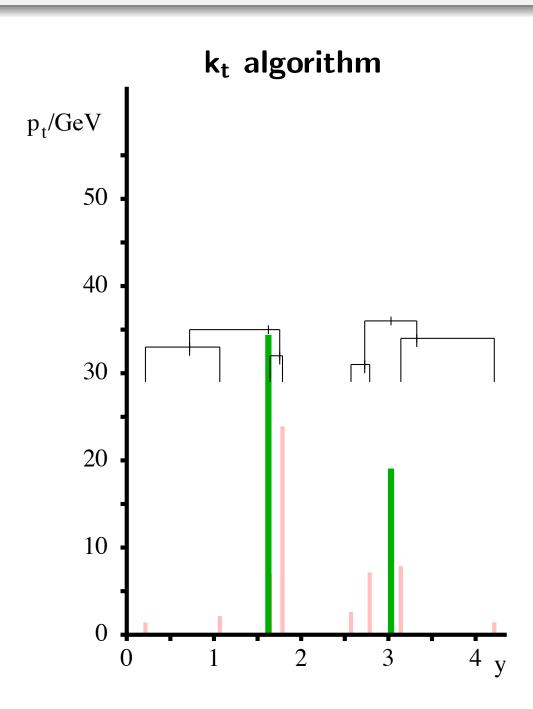
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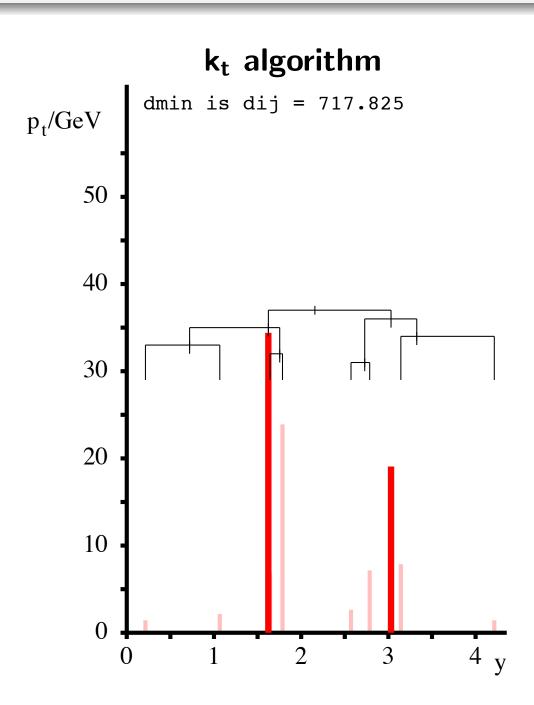
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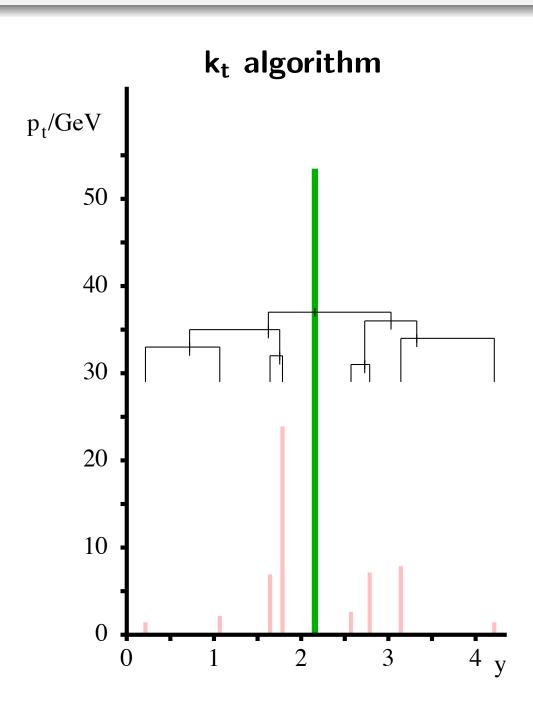


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 k_t clusters soft "junk" early on in the clustering

Its last step is to merge two hard pieces. Easily undone to identify underlying kinematics

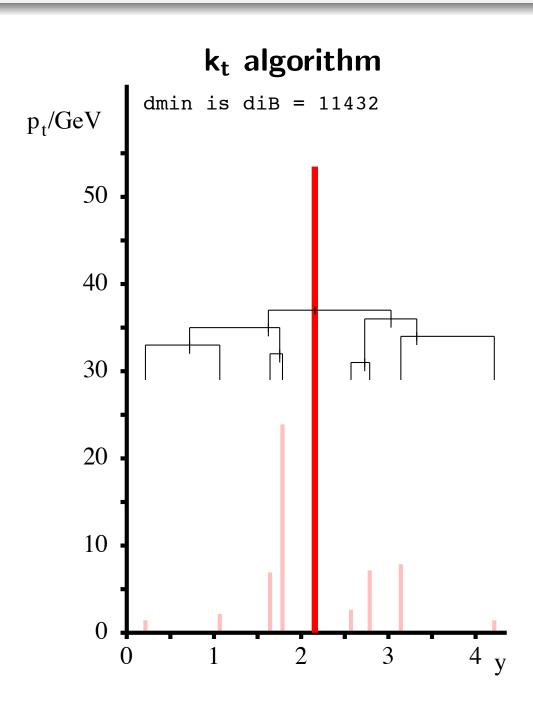


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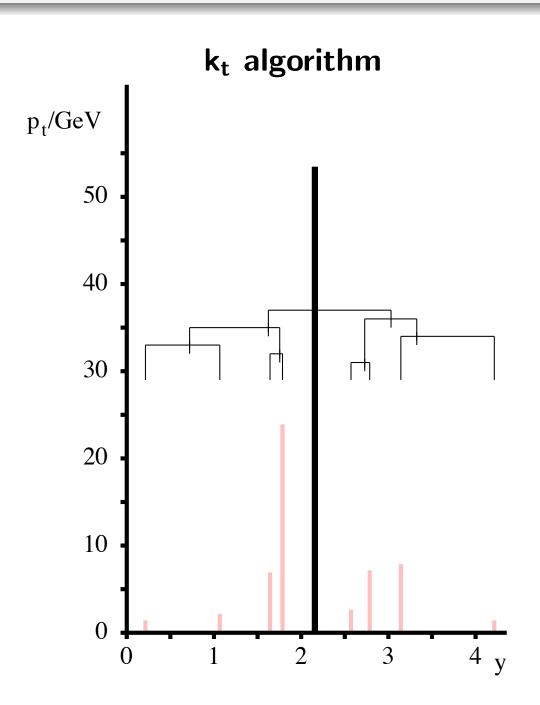


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This meant it was the first algorithm to be used for jet substructure.

Seymour '93

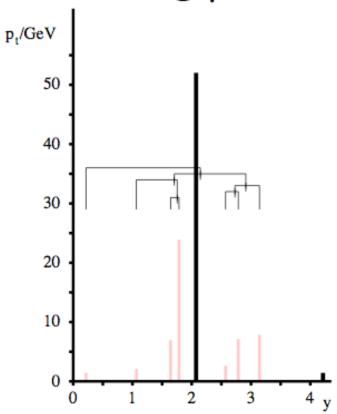
Butterworth, Cox & Forshaw '02

Third try

Cambridge/Aachen

Hierarchical substructure

Cambridge/Aachen



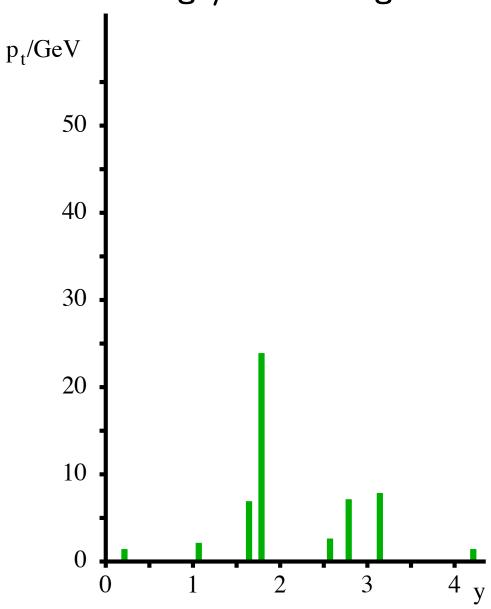
C/A distance measure

$$d_{ij} = \frac{\Delta y^2 + \Delta \phi^2}{R^2}$$

Cluster by merging the **closest** particles

Identifying jet substructure: Cam/Aachen

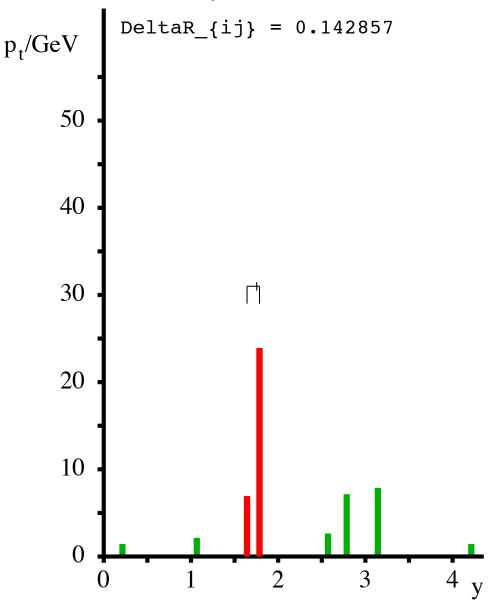
Cambridge/Aachen algorithm



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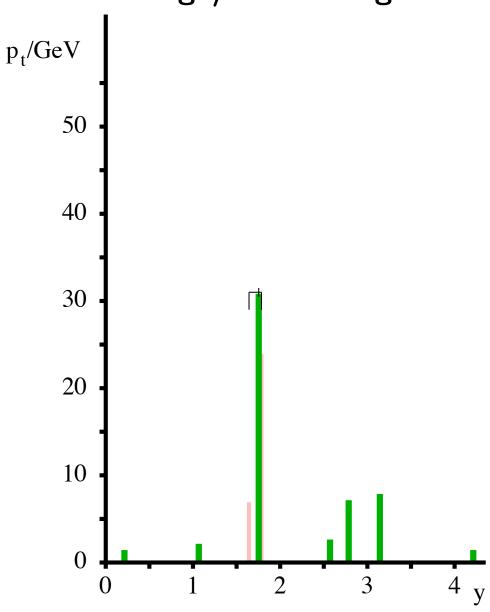
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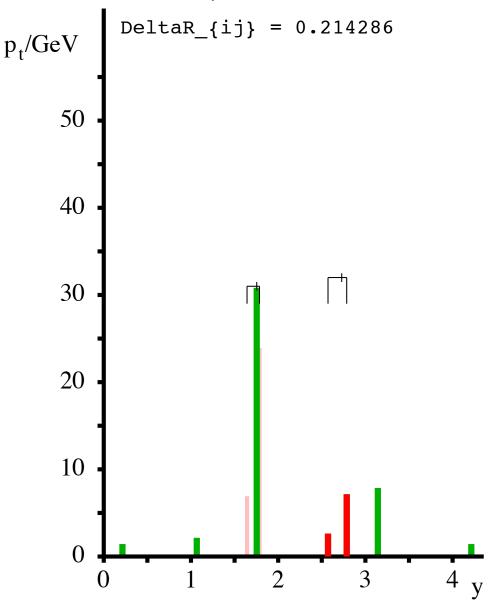
Identifying jet substructure: Cam/Aachen

Cambridge/Aachen algorithm

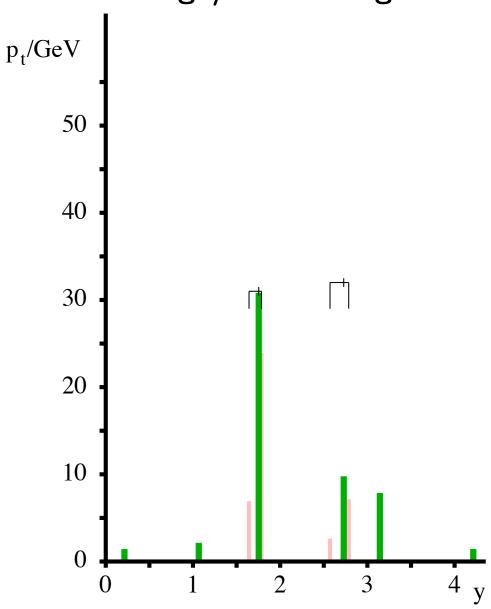


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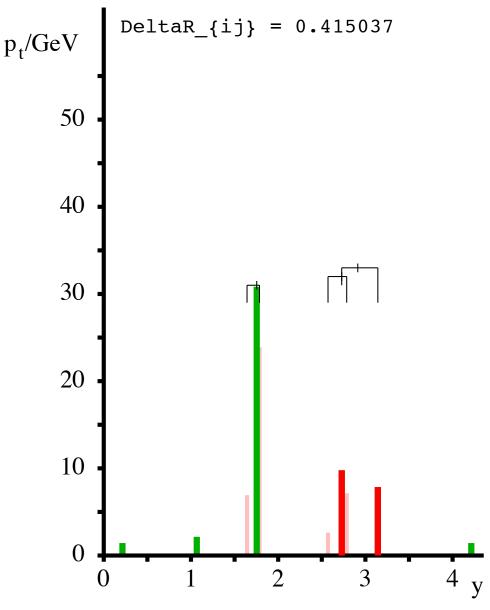
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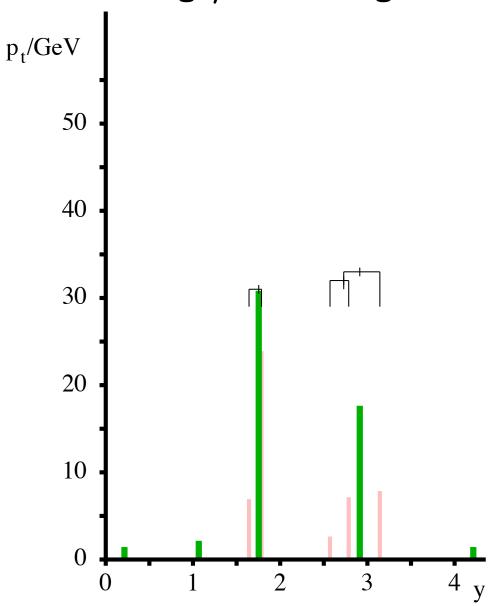
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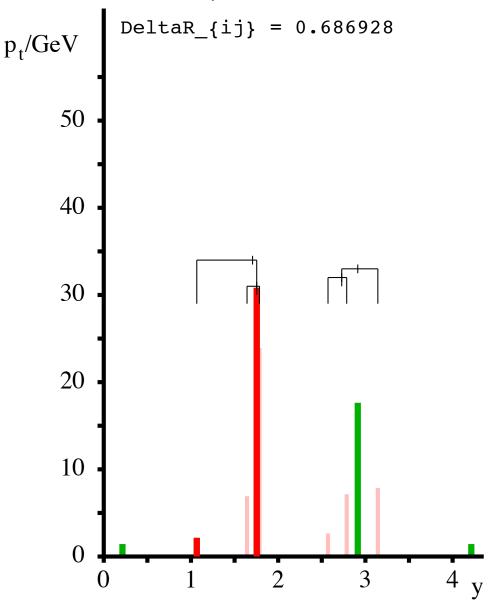
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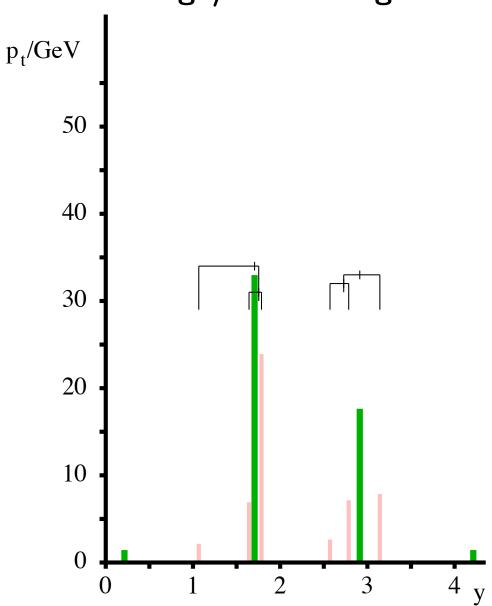
Cambridge/Aachen algorithm



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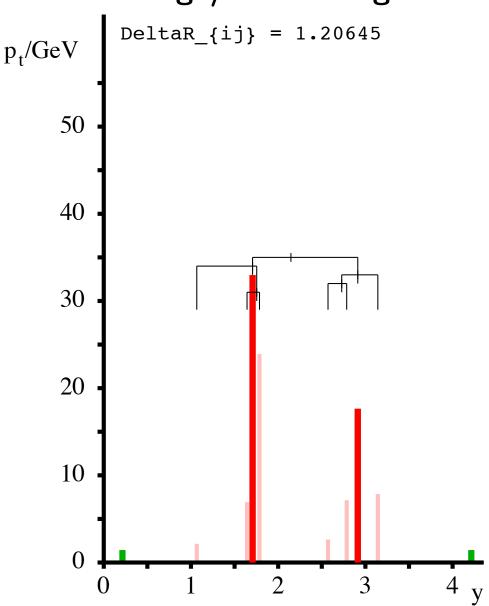
Cambridge/Aachen algorithm



How well can an algorithm identify the "blobs" of energy inside a jet that come from different partons?

C/A identifies two hard blobs with limited soft contamination

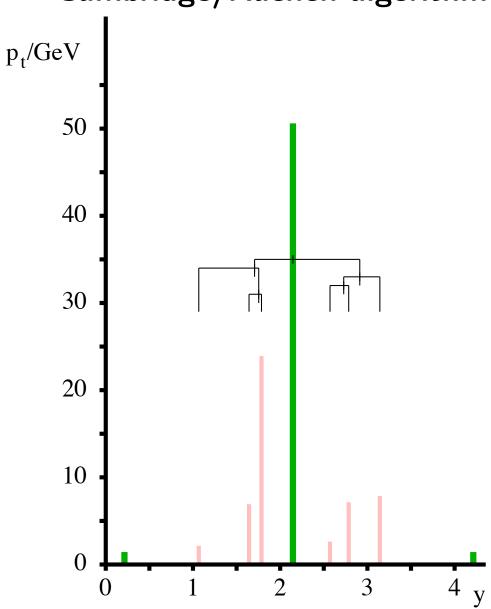
Cambridge/Aachen algorithm



How well can an algorithm identify the "blobs" of energy inside a jet that come from different partons?

C/A identifies two hard blobs with limited soft contamination, joins them

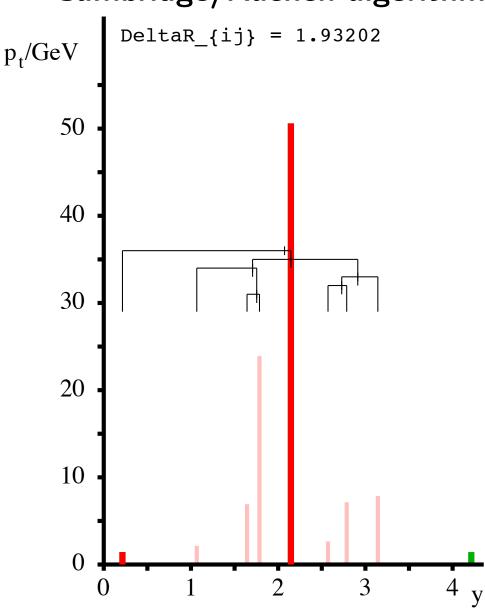
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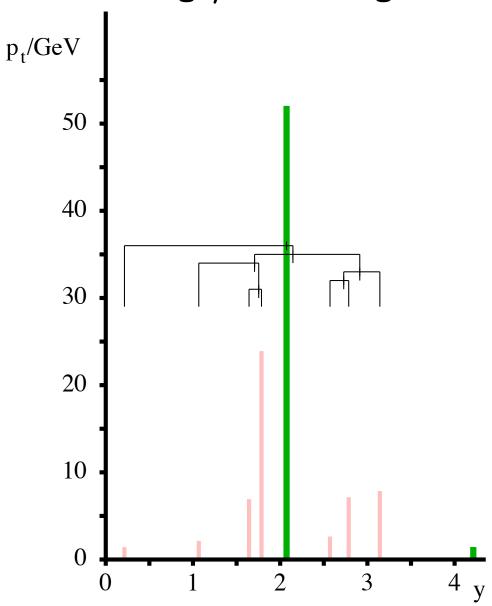
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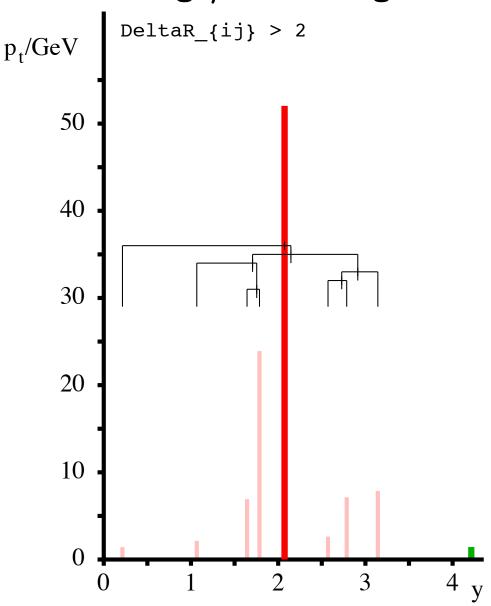
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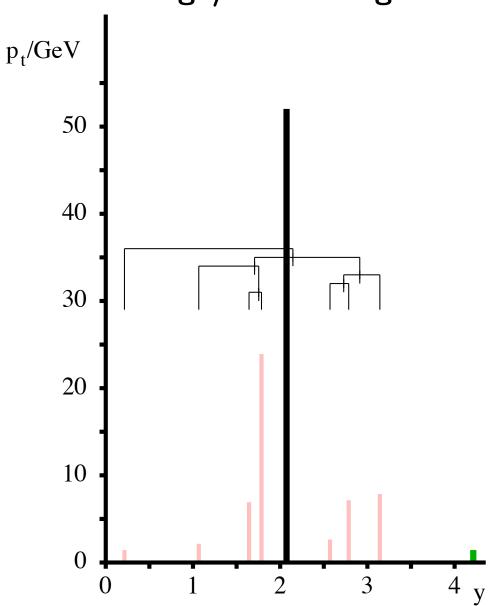
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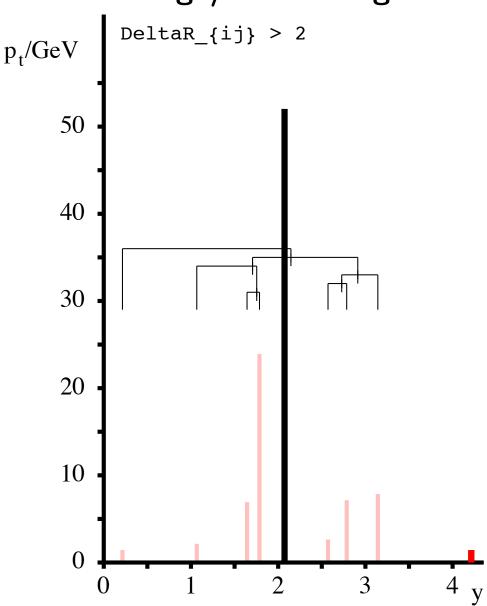
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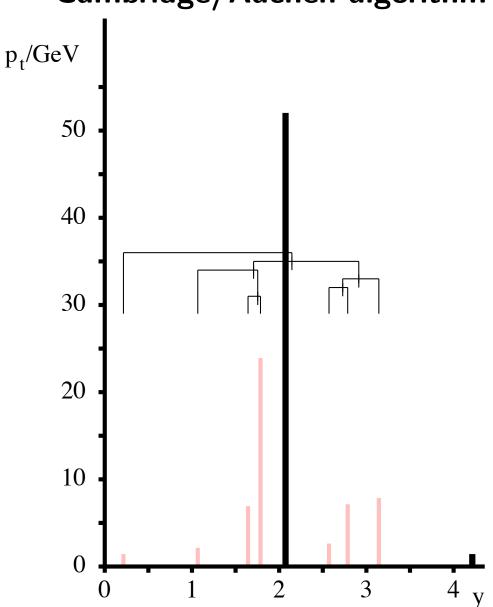
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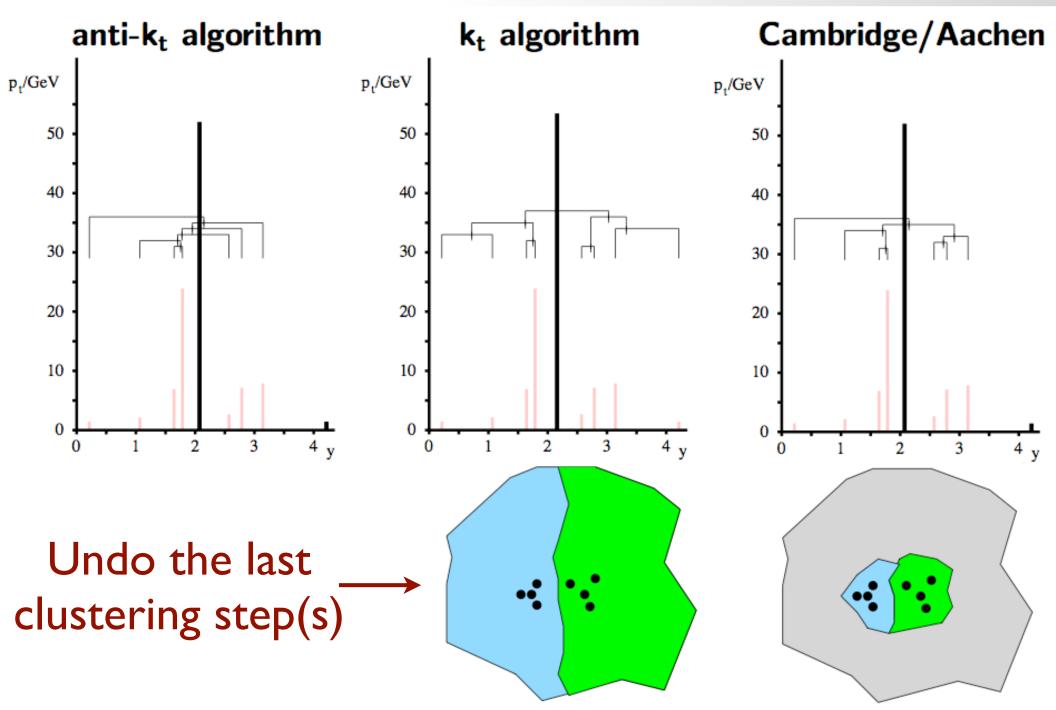
How well can an algorithm identify the "blobs" of energy inside a jet that come from different partons?

C/A identifies two hard blobs with limited soft contamination, joins them, and then adds in remaining soft junk

The interesting substructure is buried inside the clustering sequence — it's less contamined by soft junk, but needs to be pulled out with special techniques

Butterworth, Davison, Rubin & GPS '08 Kaplan, Schwartz, Reherman & Tweedie '08 Butterworth, Ellis, Rubin & GPS '09 Ellis, Vermilion & Walsh '09

Hierarchical substructure



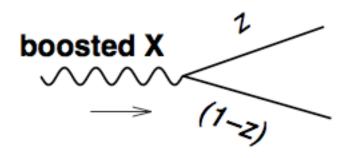
The IRC safe algorithms

	Speed	Regularity	UE contamination	Backreaction	Hierarchical substructure
k _t	© © ©		T	**	9000
Cambridge /Aachen	⊕ ⊕ ⊕			**	◎ ◎ ◎
anti-k _t	© © ©	<u> </u>	♣ /☺	❷⊙⊙®	×
SISCone	⊚	•	9000	•	×

Array of tools with different characteristics. Pick the right one for the job

QCD v. heavy decay

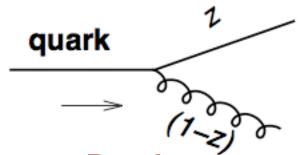
A possible approach for reducing the QCD background is to identify the two prongs of the heavy particle decay, and put a cut on their momentum fraction



Signal:

$$P(z) \sim 1$$

Will split mainly symmetrically



$$P(z) \sim \frac{1+z^2}{1-z}$$

$$P(z) \sim \frac{1+z^2}{1-z}$$
 $P(z) \sim \frac{1+(1-z)^2}{z}$

Will split mainly asymmetrically

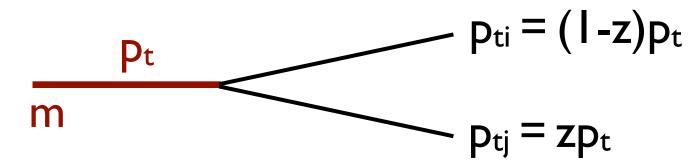
Potential tagger: asymmetric splitting

Possibly implemented via a cut on

$$y = min(p_{ti}^2, p_{tj}^2) \frac{\Delta R_{ij}^2}{m^2} \simeq \frac{min(p_{ti}, p_{tj})}{max(p_{ti}, p_{tj})}$$

Splittings and distances

Quasi-collinear splitting $(p_{tj} < p_{ti})$



Invariant mass:

$$m^2 \simeq p_{ti} p_{tj} \Delta R_{ij}^2 = (1-z) z p_t^2 \Delta R_{ij}^2$$

k_t distance:

$$d_{ij} \stackrel{\text{(ptj < pti)}}{=} z^2 p_t^2 \Delta R_{ij}^2 \simeq \frac{z}{1-z} m^2$$

For a given mass, the **background** will have smaller distance d_{ij} than the signal, i.e. it will tend to **cluster earlier** in the k_t algorithm

Potential tagger: last clustering in kt algorithm

This is where the hierarchy of the k_t algorithm becomes relevant. QCD radiation is clustered first, and only at the end the symmetric, large-angle splittings due to decays are reclustered

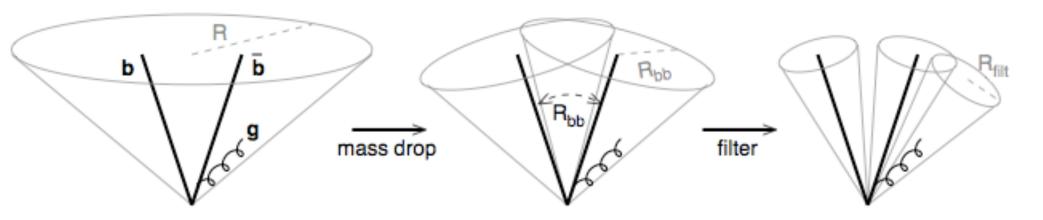
Alternative algorithms

- \blacktriangleright Suppose that for some reasons (which will become clearer later) one does not with to use the k_t algorithm
 - ▶ One must then find a way to determine what the **relevant splitting** (i.e. the one due to the decay, not to QCD radiation) is.

A possible approach is to use a Mass-Drop requirement: the clustering is **progressively undone**, and a splitting is the relevant one if both subjects are much less massive than their combination

The BDRS tagger/groomer

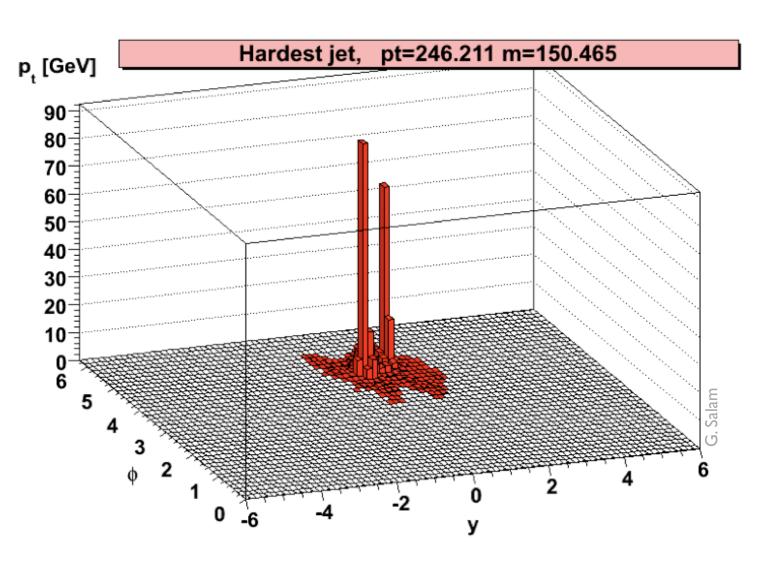
Butterworth, Davison, Rubin, Salam, 2008



- A two-prong tagger/groomer for boosted Higgs, which
 - ▶ Uses the Cambridge/Aachen algorithm (because it's 'physical')
 - ▶ Employs a Mass-Drop condition, as well as an asymmetry cut to find the relevant splitting (i.e. 'tag' the heavy particle)
 - Includes a post-processing step, using 'filtering' (introduced in the same paper) to clean as much as possible the resulting jets of UE contamination ('grooming')

BDRS: tagging

$PP \rightarrow ZH \rightarrow V\bar{V}b\bar{b}$



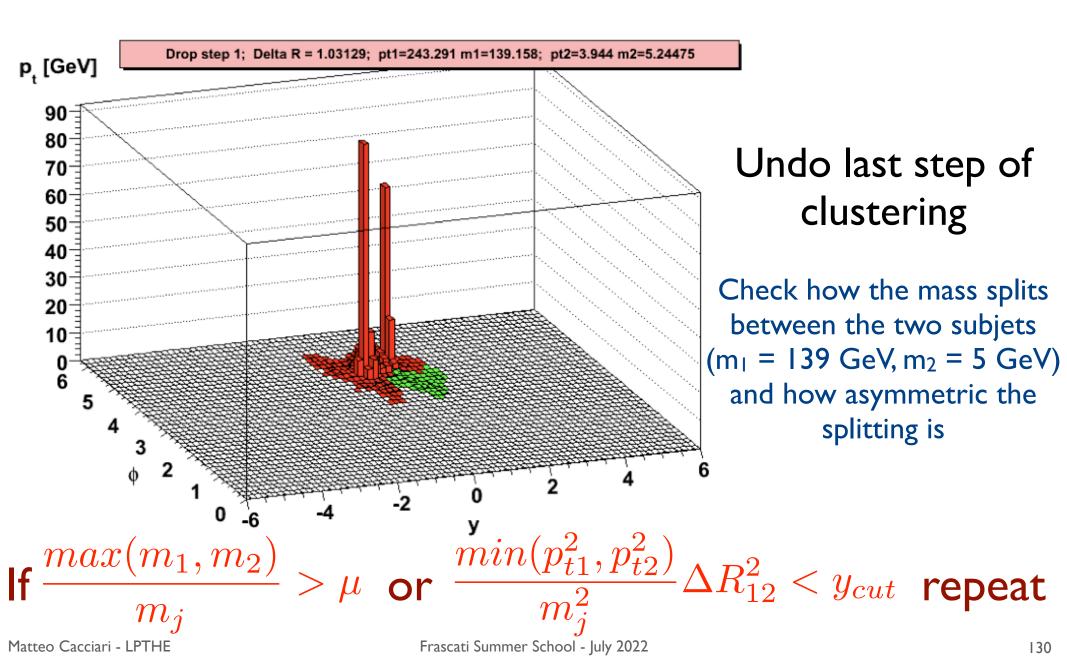
Start with the hardest jet

Use C/A with large R=1.2

 $m_j = 150 \text{ GeV}$

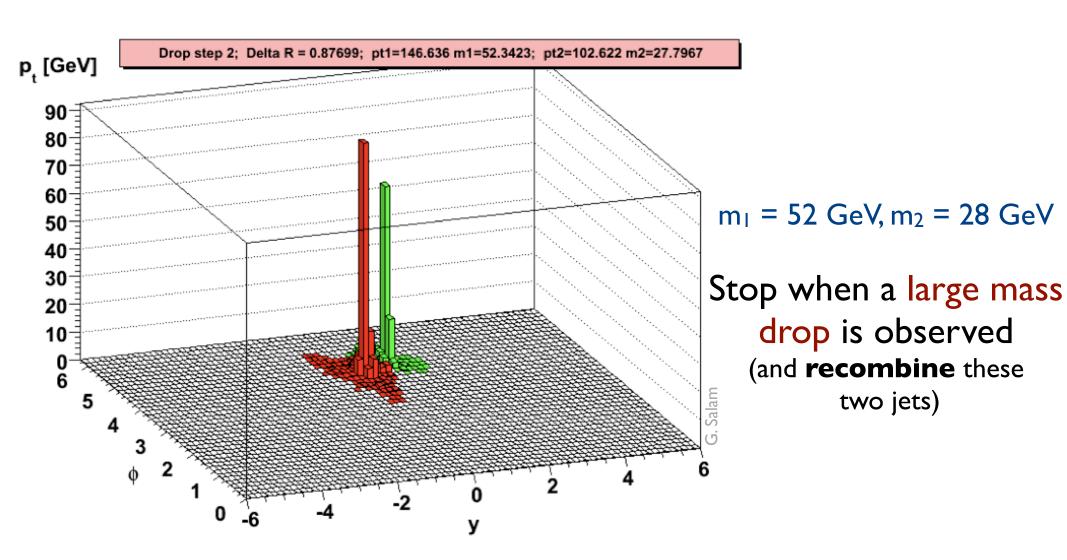
BDRS: tagging

$pp \rightarrow ZH \rightarrow vvbb$



BDRS: tagging

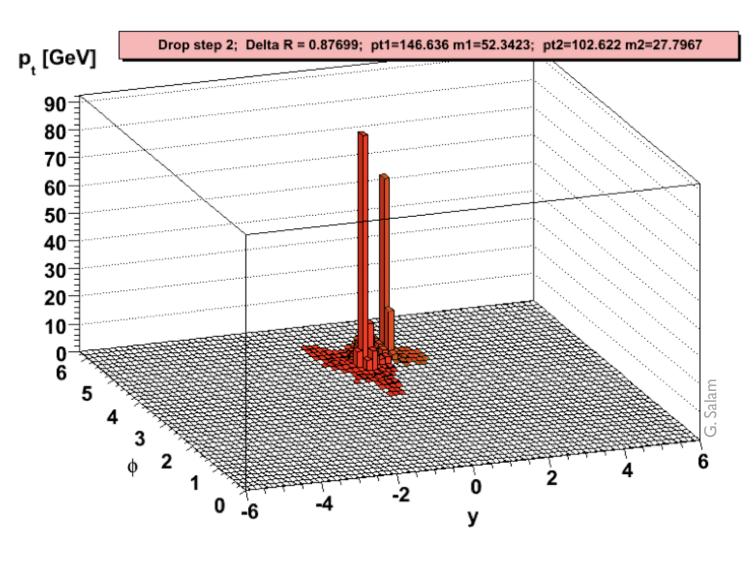
$pp \rightarrow ZH \rightarrow vvbb$



[NB. Parameters used $\mu = 0.67$ and $y_{cut} = 0.09$]

BDRS: filtering

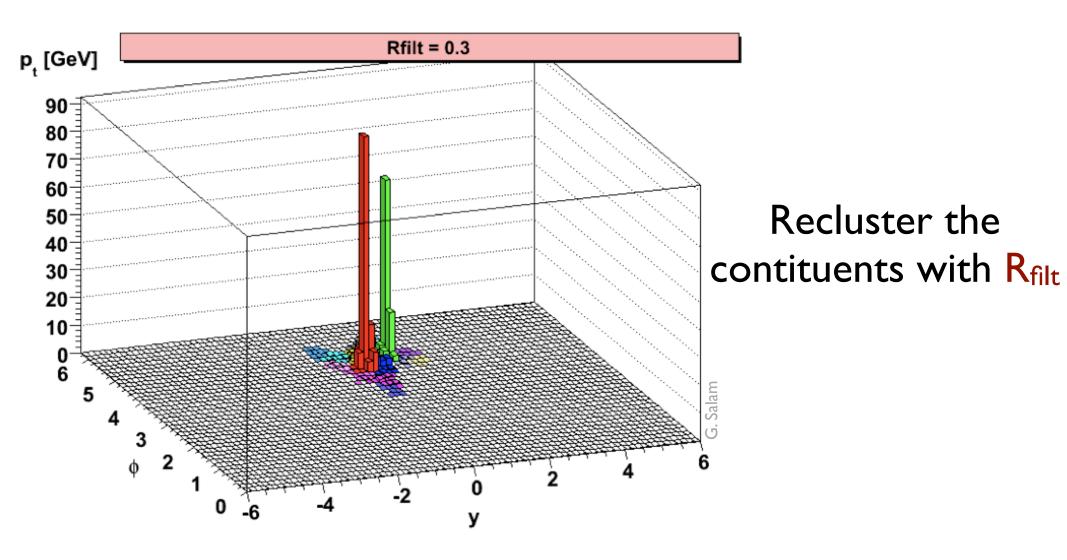
$pp \rightarrow ZH \rightarrow vvbb$



Start with the recombined jet

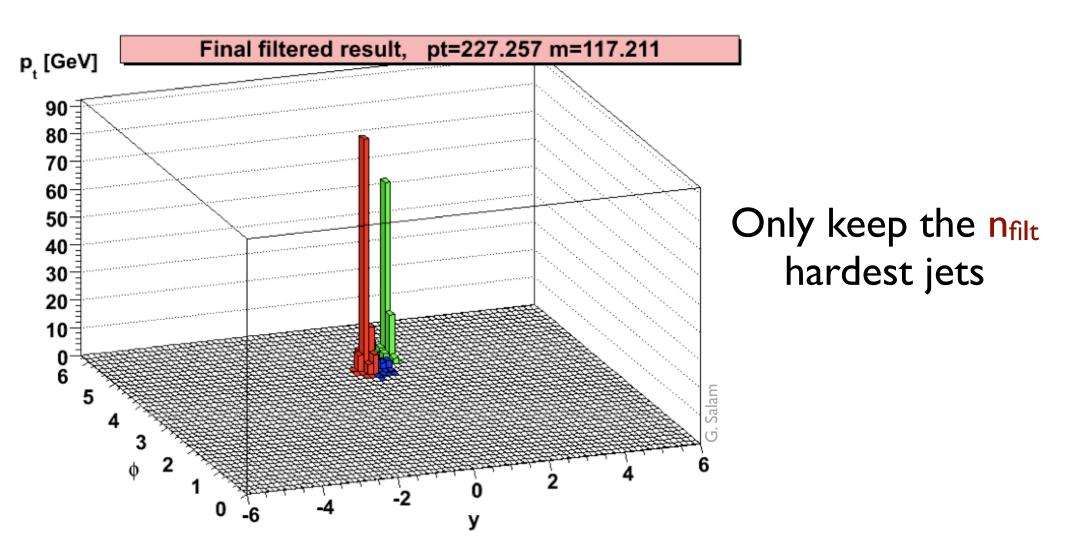
BDRS: filtering

$pp \rightarrow ZH \rightarrow vvbb$



BDRS: filtering

$pp \rightarrow ZH \rightarrow vvbb$

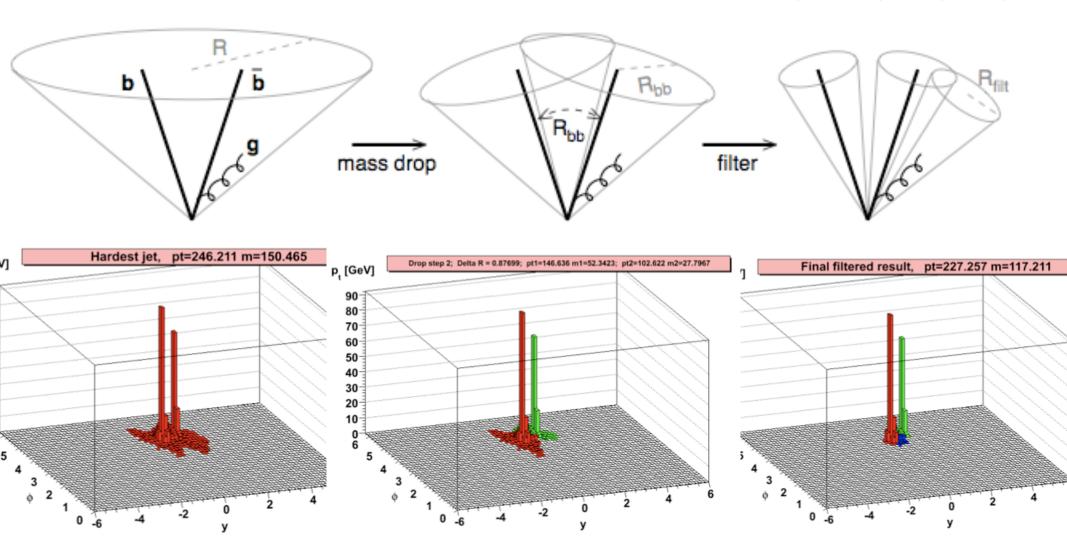


The low-momentum stuff surrounding the hard particles has been removed

$pp \rightarrow ZH \rightarrow v\bar{v}b\bar{b}$

Visualisation of BDRS

Butterworth, Davison, Rubin, Salam, 2008



Cluster with a large R

Undo the clustering into subjets, until a large asymmetry/mass drop is observed: tagging step

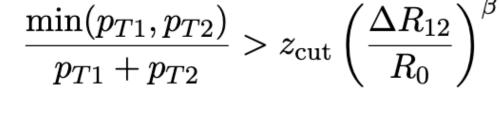
Re-cluster with smaller R, and keep only 3 hardest jets: grooming step

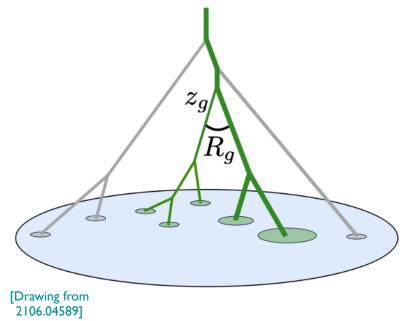
Soft Drop declustering

Larkoski, Marzani, Soyez, Thaler, 2014

A generalisation of the (modified) Mass-Drop tagger. Progressively decluster and drop constituent unless

Soft Drop Condition:
$$\frac{mn}{r}$$





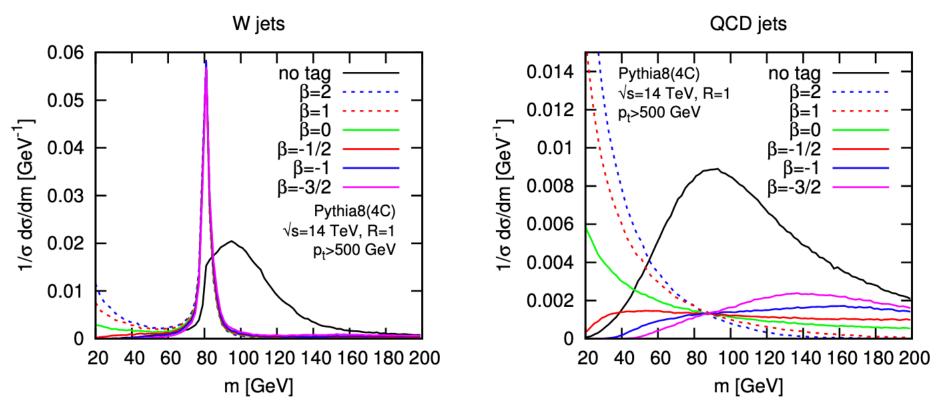
i.e. remove large-angle soft radiation from a jet

Soft Drop declustering

Larkoski, Marzani, Soyez, Thaler, 2014

The paper contains

- √ analytical calculations and comparisons to Monte Carlos
- √ study of effect of non-perturbative corrections
- ✓ performance studies



Example of SoftDrop performance when used as a boosted W tagger

Take home points

The big news of the past fifteen years has been the development of robust taggers and groomers using properties of jet substructure, through

- declustering
- jet shapes
- direct analysis of images (machine learning)

These techniques have been commissioned by experimental collaborations and have proven their worth in 'Standard Model' analyses. They are now being implemented in BSM searches