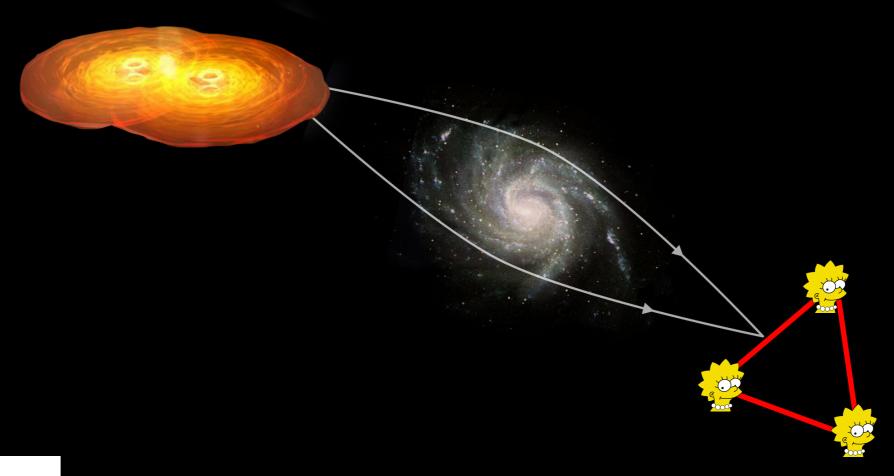
Observability of lensing of gravitational waves from massive black hole binaries with LISA

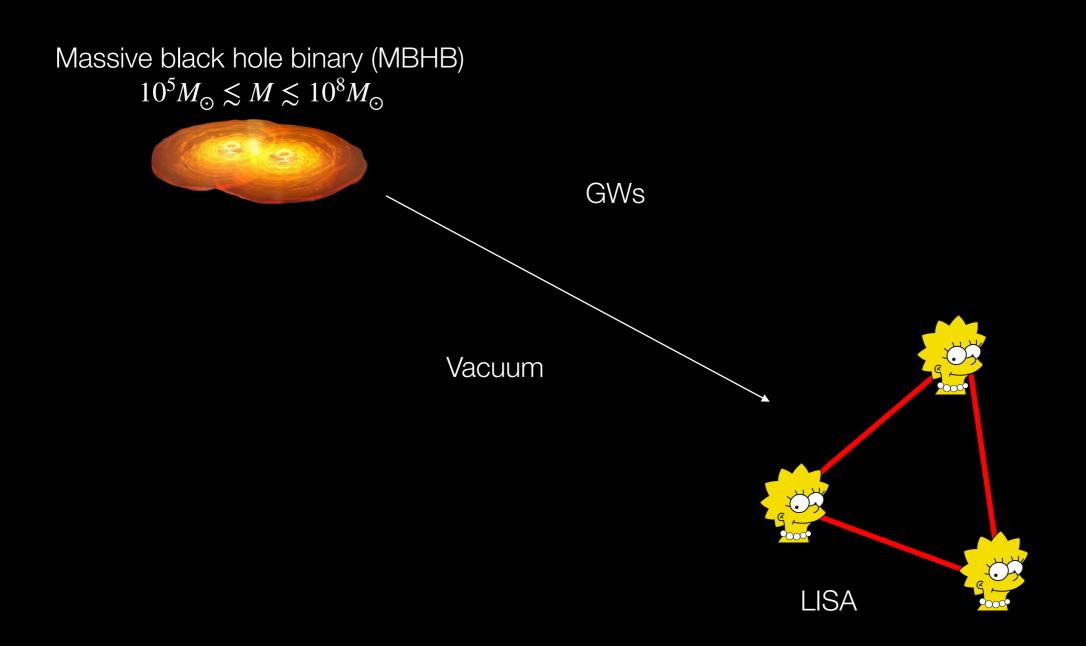


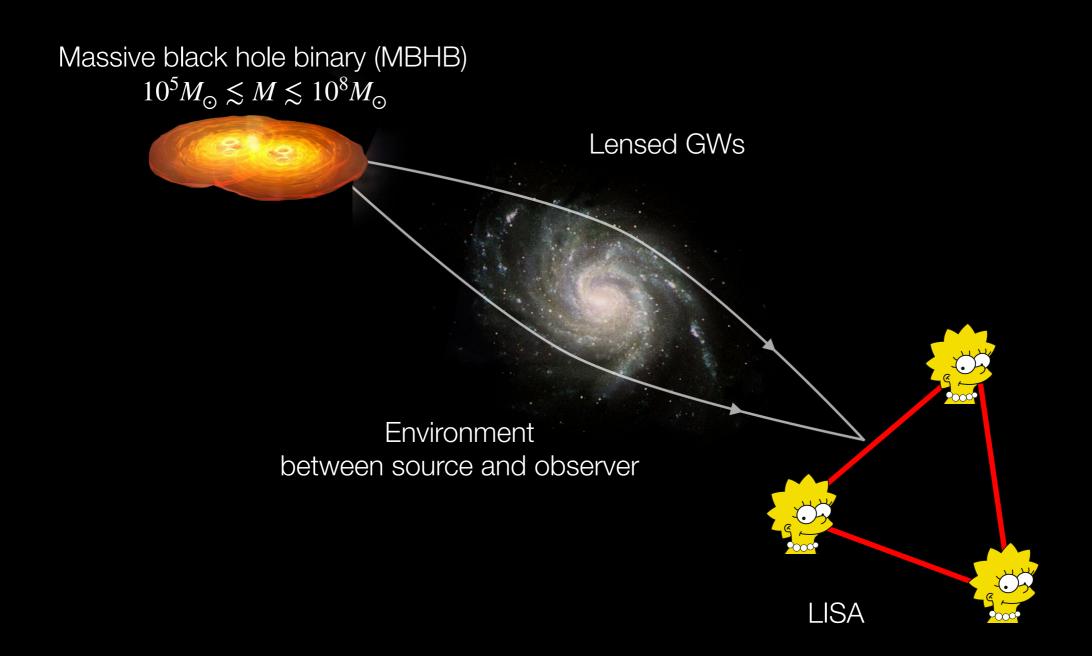


Roberto Cotesta



See the paper on arXiv:2206.02803 by Mesut Çalışkan, Lingyuan Ji, Roberto Cotesta, Emanuele Berti, Mark Kamionkowski and Sylvain Marsat





$$\tilde{h}^{L}(f; \theta^{S}, \theta^{L}) = F(f, \mathbf{y}) \, \tilde{h}(f; \theta^{S})$$

$$\theta^{\rm S} \equiv {\rm parameters} \ {\rm of} \ {\rm the} \ {\rm source}$$

$$\theta^{\mathrm{L}} = \{M_{\mathrm{L}}, \mathbf{y}\}$$

 $M_{\rm L} \equiv {\rm mass} \ {\rm of} \ {\rm the \ lens}$

 $y \equiv$ coordinates of the source in source plane

$$F(f, \mathbf{y}) \equiv \text{diffraction integral}$$

$$F(f, \mathbf{y}) \equiv \frac{D_{S}(1 + z_{L})\xi_{0}^{2}}{D_{L}D_{LS}} \frac{f}{i} \int d^{2}\mathbf{x} \exp[2\pi i f t_{d}(\mathbf{x}, \mathbf{y})]$$

 $D_{\rm S} \equiv$ distance from the source

 $D_{\rm L} \equiv$ distance from the lens

 $z_{\rm L} \equiv {\rm redshift}$ of the lens

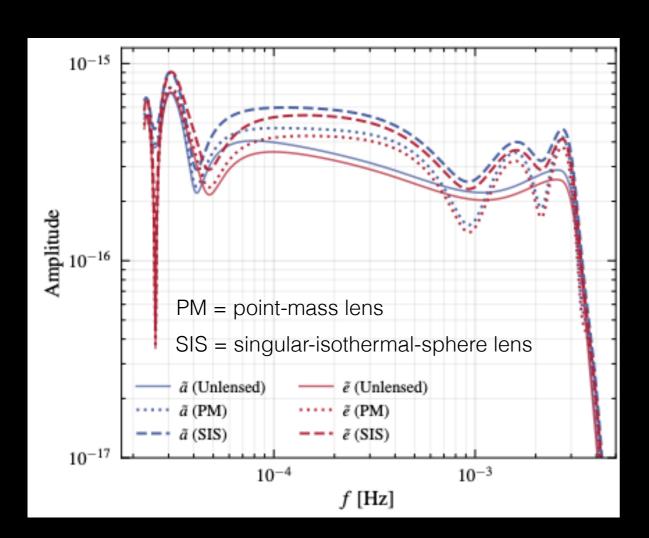
 $D_{\rm LS} \equiv$ distance lens-source

 $t_{\rm d}({\bf x},{\bf y}) \equiv {\rm time\ delay\ function}$

 $x \equiv$ coordinates of the image in lens plane

$$\tilde{h}^{L}(f; \theta^{S}, \theta^{L}) = F(f, \mathbf{y}) \, \tilde{h}(f; \theta^{S})$$

 $heta^{\mathrm{S}} \equiv \mathrm{parameters}$ of the source $heta^{\mathrm{L}} = \{M_{\mathrm{L}}, \mathbf{y}\}$ $M_{\mathrm{L}} \equiv \mathrm{mass}$ of the lens $\mathbf{y} \equiv \mathrm{coordinates}$ of the source in source plane $F(f, \mathbf{y}) \equiv \mathrm{diffraction}$ integral



$$F(f, \mathbf{y}) \equiv \frac{D_{\mathrm{S}}(1 + z_{\mathrm{L}})\xi_0^2}{D_{\mathrm{L}}D_{\mathrm{LS}}} \frac{f}{i} \int \mathrm{d}^2\mathbf{x} \, \exp[2\pi i f t_{\mathrm{d}}(\mathbf{x}, \mathbf{y})]$$

 $D_{\rm S} \equiv$ distance from the source $D_{\rm L} \equiv$ distance from the lens $z_{\rm L} \equiv$ redshift of the lens $D_{\rm LS} \equiv$ distance lens-source $t_{\rm d}({\bf x},{\bf y}) \equiv$ time delay function ${\bf x} \equiv$ coordinates of the image in lens plane

$$\tilde{h}^{L}(f; \theta^{S}, \theta^{L}) = F(f, \mathbf{y}) \, \tilde{h}(f; \theta^{S})$$

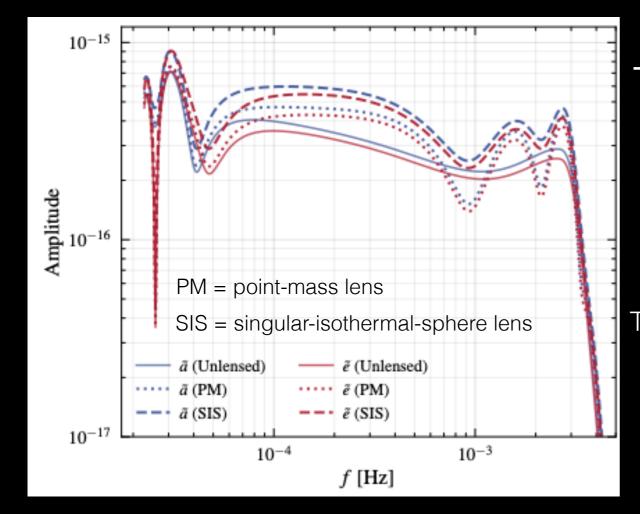
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 $D_{\rm LS} \equiv$ distance lens-source

 $t_{\rm d}({\bf x},{\bf y}) \equiv {\rm time\ delay\ function}$

 $x \equiv$ coordinates of the image in lens plane

Lensing effects depend on frequency!

Maybe you are more familiar with...

$$F(f, \mathbf{y}) \approx \sum_{j} \sqrt{|\mu_{j}|} \exp[2\pi i f t_{d}(\mathbf{x}_{j}, \mathbf{y}_{j}) - i \pi n_{j}]$$

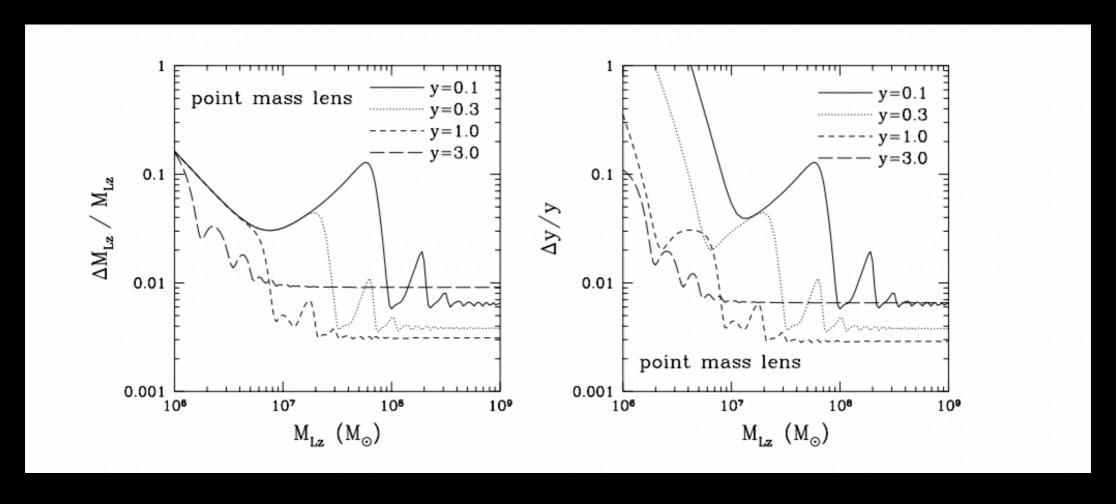
The **geometric optics** limit is **NOT** generically **valid** for **LISA**!

Diffraction effects **important** when

$$M_{\rm L} \lesssim 10^5 M_{\odot}/(f/{\rm Hz})$$

Can we measure lensing from LISA observations?

YES!



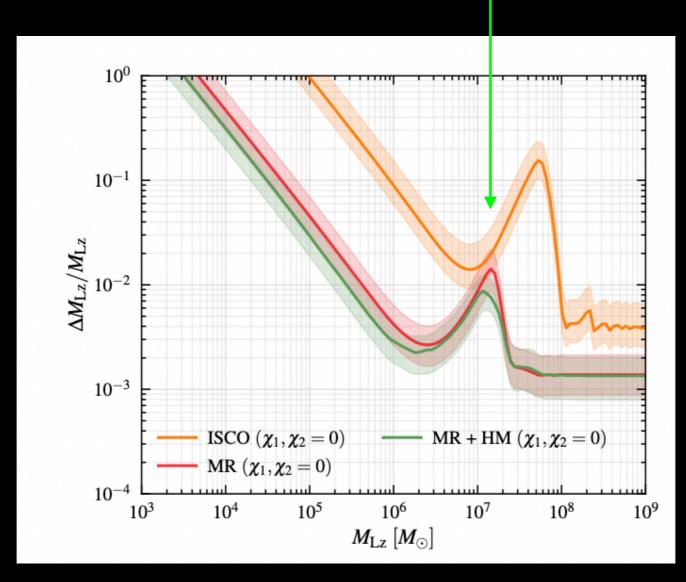
Takahashi and Nakamura (2003)

Fisher matrix analysis for non-spinning $10^6 + 10^6 M_{\odot}$ @ $z_s = 1$

Lensing is measurable when $M_{\rm L} \gtrsim 10^6 M_{\odot}, y \le 10^8 M_{\odot}$

Impact of merger-ringdown and higher harmonics

Takahashi and Nakamura used inspiral-only waveforms without higher harmonics

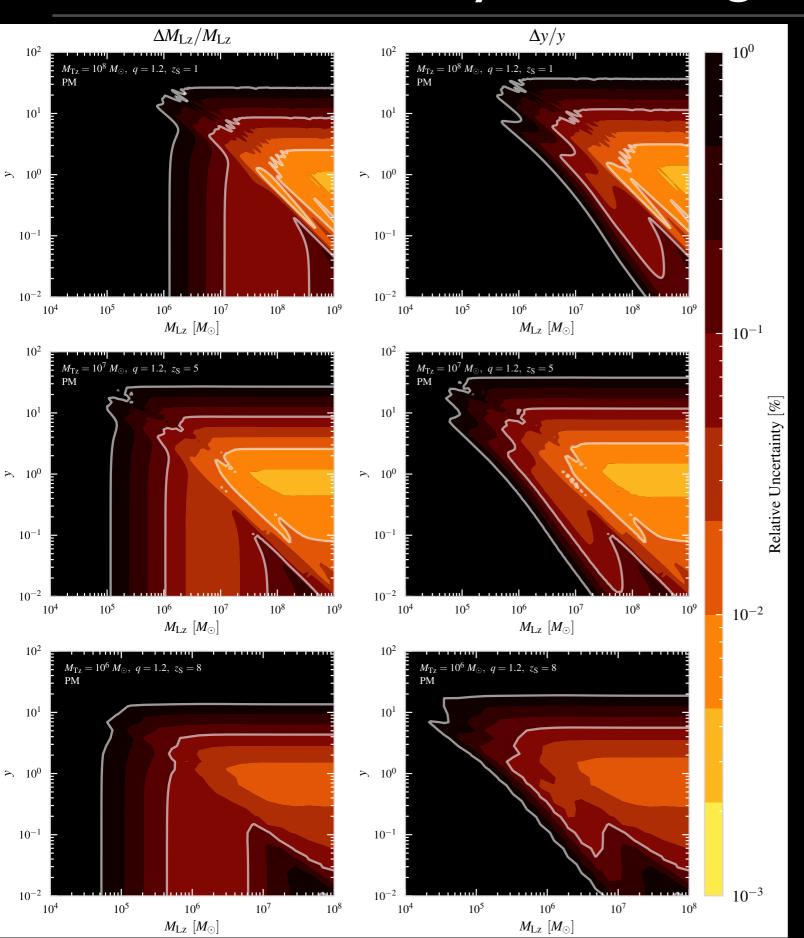


Çalışkan,...RC et al.

Fisher matrix analysis for non-spinning $10^6 + 10^6 M_{\odot}$ @ $z_s = 1, y = 0.1$ with analytic expression of diffraction integral and its derivatives for Fisher matrix we use lisabeta Marsat+(2020)

Including merger-ringdown in the waveform improves precision on lens mass measurement by an order of magnitude, higher harmonics by a factor of 2

Measurability of lensing effects in MBHBs



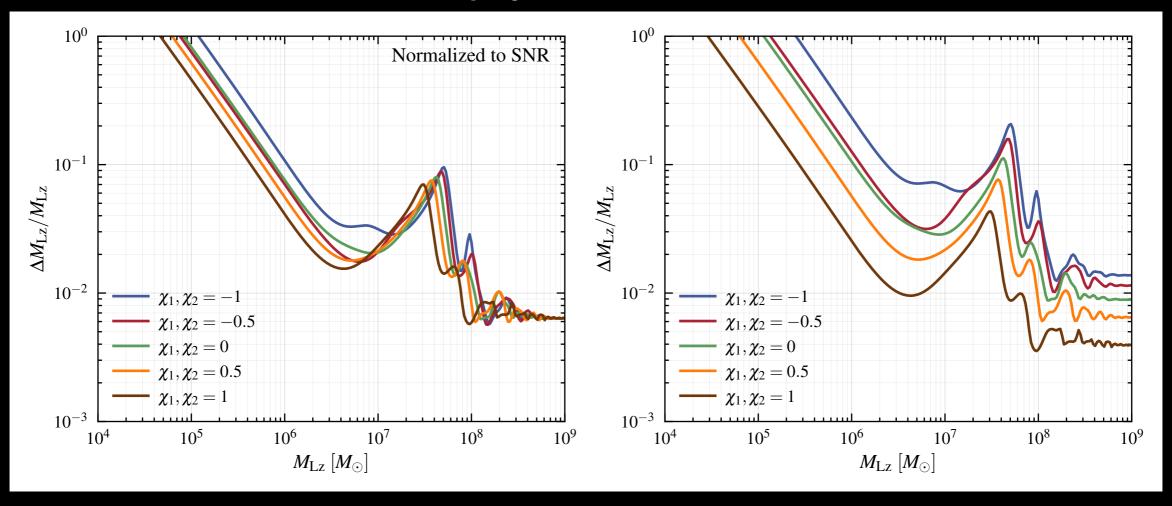
Çalışkan,...RC et al.

			$\Delta M_{ m Lz}/M_{ m Lz}$		$\Delta y/y$
$M_{ m Tz} \; [M_{\odot}]$	q	$z_{ m S}$	$M_{ m Lz}^{ m crit} \; [M_{\odot}]$	$y^{ m crit}$	$y^{ m crit}$
10^{8}	1.2	1	$\gtrsim 10^6$	$\lesssim 30$	$\lesssim 40$
10^{7}	1.2	5	$\gtrsim 10^5$	$\lesssim 30$	$\lesssim 40$
10^{6}	1.2	8	$\gtrsim 5 \times 10^4$	$\lesssim 15$	$\lesssim 20$

We have results also for SIS lens in partial agreement with Gao+ (2022)

Impact of BH spins

Çalışkan,...RC et al.



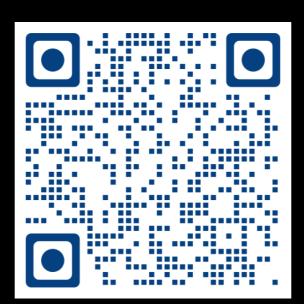
Fisher matrix analysis for $5.5 \times 10^6 + 4.5 \times 10^6 M_{\odot}$ @ $z_s = 5$, y = 0.1

Large spins aligned with the angular momentum improve measurability by an order of magnitude (a factor of 3 when normalizing by the SNR), when compared to large anti-aligned spins

Conclusions and future work

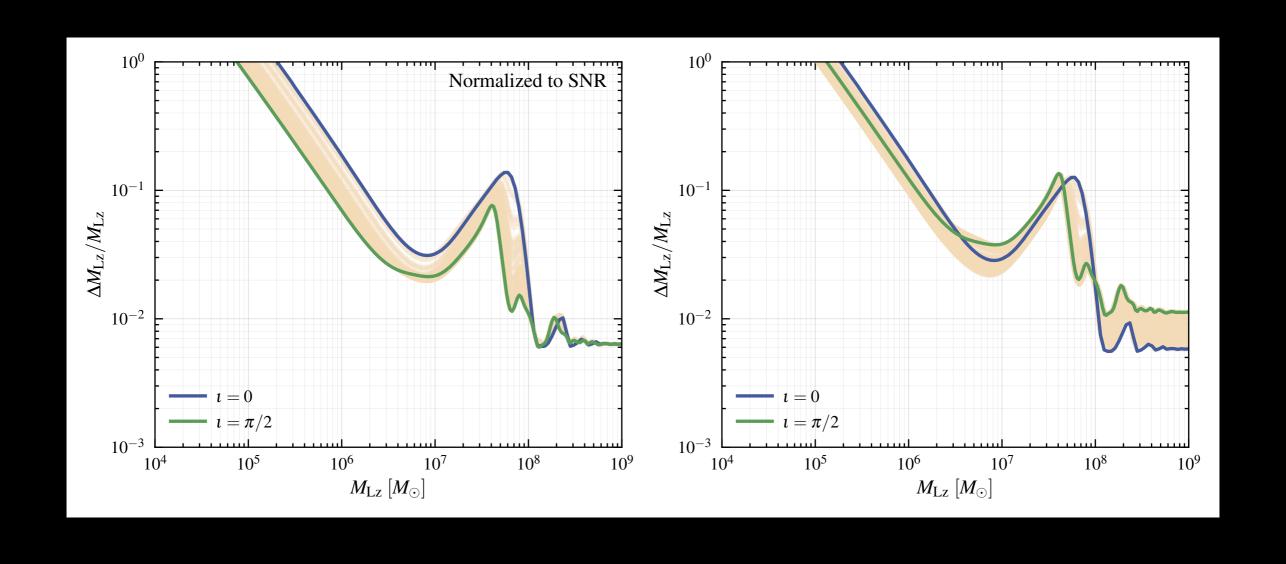
- For MBHBs in LISA we cannot always use the geometric optics approximation, as diffraction effects are important
- Lensing effects are detectable for certain MBHBs in LISA
- Previous studies underestimated by an order of magnitude the precision on the measurement
 of the lens mass and by a factor of 2 that on the impact parameter y, because of neglecting the
 merger-ringdown part of the waveform
- Using Fisher matrix and analytical solution for the diffraction integral and its derivatives, we have analyzed the
 measurability of lensing for "typical" MBHBs in LISA using both PM and SIS lens
- We investigated the effect on all binary parameters (BH spins, source-observer inclination, sky position...) on the measurability of lensing effects
- Use our Fisher matrix analysis on a population of MBHBs to estimate the number of GW signals with detectable lensing effects given a population of lenses

Thanks!

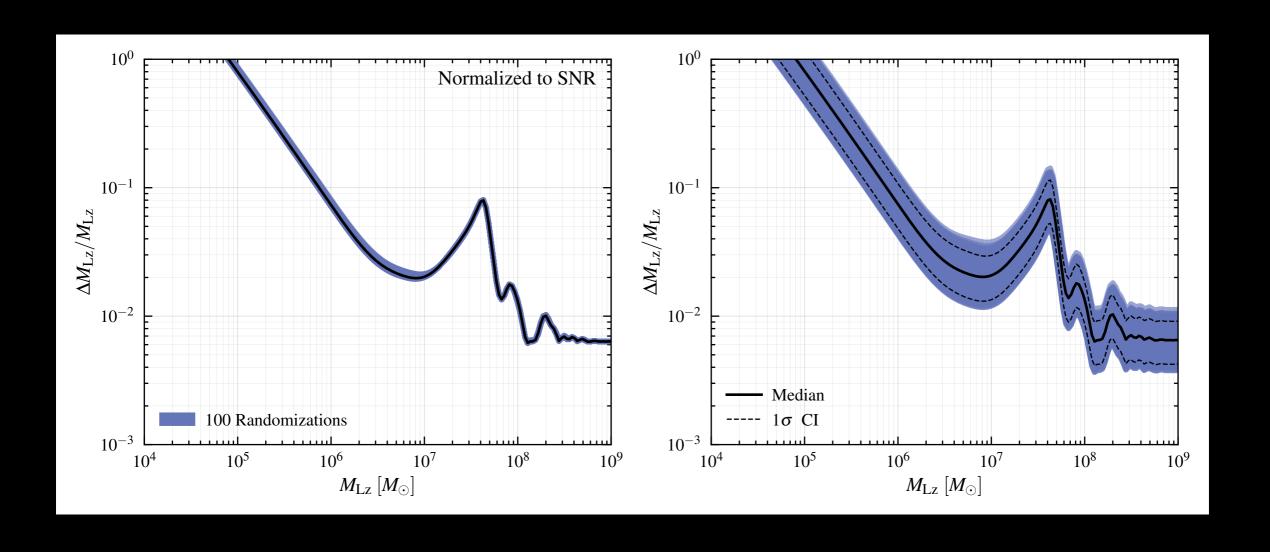


Read this QR code for the paper! arXiv:2206.02803

Impact of source-observer angle



Impact of sky-localization and other angles



SIS results

