# Curing Sterile Neutrino Dark Matter with a Dark Force

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Based on Torsten Bringmann, Paul Frederik Depta, Marco Hufnagel, JK, Joshua T. Ruderman, Kai Schmidt-Hoberg, arXiv:2206.10630

# Introduction

- WIMP Dark Matter theoretically appealing but not observed
- Alternative: sterile neutrino  $\nu_s$  with mass  $m_s \sim \text{keV}$
- Mixing angle  $\theta$  with Standard Model neutrino  $\nu_a$
- Dodelson-Widrow production by oscillations in early Universe Dodelson & Widrow, PRL 72 (1994)

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- Decay  $\nu_s \rightarrow \nu_a \gamma$  at 1-loop level  $\Gamma \propto \sin^2 2\theta m_s^{-5}$ 
  - $\rightsquigarrow$  Bounds from searches for X-ray line

Abazajian, Fuller, Tucker, APJ 562 (2001)

# Dodelson-Widrow Mechanism Cannot Produce All DM



# **Alternative Production Mechanisms**

 Resonant production in presence of lepton asymmetry Shi & Fuller, PRL 82 (1999)

### Scalar decay

Shaposhnikov & Tkachev, PLB **639** (2006) Kusenko, PRL **97** (2006); Petraki & Kusenko, PRD **77** (2008) Roland, Shakya, Wells, PRD **92** (2015) Merle & Totzauer, JCAP **06** (2015); König, Merle, Totzauer, JCAP **11** (2016)

### New interactions of Standard Model neutrinos

De Gouvêa, Sen, Tangarife, Zhang, PRL **124** (2020) Kelly, Sen, Tangarife, Zhang, PRD **101** (2020) Chichiri, Gelmini, Lu, Takhistov, arXiv:2111.04087 Benso, Rodejohann, Sen, Ramachandran, PRD **105** (2022)

### Extended gauge sector

Bezrukov, Hettmansperger, Lindner, PRD **81** (2010) Kusenko, Takahashi, Yanagida, PLB **693** (2010) Dror, Dunsky, Hall, Harigaya, JHEP **07** (2020)

• (Insert your favorite model here)

#### • New interactions of sterile neutrinos

Hansen & Vogl, PRL **119** (2017) Johns & Fuller, PRD **100** (2019) Bringmann, Depta, Hufnagel, JK, Ruderman, Schmidt-Hoberg, arXiv:2206.10630

- Standard Model neutrino  $\nu_a$
- Sterile neutrino  $\nu_s$ , mass  $m_s \sim \text{keV}$ , right-handed spinor
- Scalar singlet  $\phi$ , mass  $m_{\phi} \gtrsim 2 m_s$
- Yukawa coupling y between scalar and gauge eigenstate  $\nu'_{s}$

$$\mathcal{L}_{\phi}^{\text{int}} = rac{\mathbf{y}}{2} \overline{\nu_{s}^{\prime c}} \phi \, \nu_{s}^{\prime} + \text{h.c.}$$

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Rotation to mass eigenstate basis by mixing angle θ
 → Neutrino-scalar interactions

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• Assume scalar self-interactions and Higgs portal to be irrelevant

### • Hansen & Vogl, PRL **119** (2017)

- $m_\phi \sim 10^2 \, {
  m keV}$
- Number-changing interactions (e.g., in scalar potential)
- Viable  $\nu_s$  Dark Matter for  $m_s \simeq 4 \text{ keV}$  or lepton asymmetry

• Johns & Fuller, PRD 100 (2019)

- *m*<sub>\phi</sub> > 1 GeV
- Interactions affect in-medium neutrino potential
- Either no impact or runaway production for  $T \sim 10^2 \text{ MeV}$
- Bringmann, Depta, Hufnagel, JK, Ruderman, Schmidt-Hoberg
  - $m_{\phi} \gtrsim 2 m_s$
  - Viable  $\nu_s$  Dark Matter due to exponential growth of abundance

Conversion of heat bath particles  $\psi$  to DM particles  $\chi$   $\rightsquigarrow$  phase of exponential growth



Bringmann, Depta, Hufnagel, Ruderman, Schmidt-Hoberg, PRL 127 (2021)

# **Exponentially Growing Sterile Neutrinos**

- Heat bath particle:  $\psi = \nu_a$
- Dark Matter particle:  $\chi = \nu_s$
- Conversion via scalar  $\phi$  (almost on-shell  $\rightsquigarrow$  *s*-channel dominates)



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- Small initial v<sub>s</sub> abundance produced by Dodelson-Widrow
- Freeze-in production  $\nu_a \nu_a \rightarrow \nu_s \nu_s$  negligible due to tiny mixing

# Calculation of Dark Sector Densities

- Initial condition: ν<sub>s</sub> abundance from Dodelson-Widrow mechanism Asaka, Laine, Shaposhnikov, JHEP 01 (2007)
- $\nu_{s}\nu_{s} \leftrightarrow \phi \rightsquigarrow \text{dark sector thermalizes}$
- Boltzmann equations for number and energy densities

$$\dot{n}_{s} + 3Hn_{s} = C_{n_{s}}$$
  
 $\dot{n}_{\phi} + 3Hn_{\phi} = C_{n_{\phi}}$   
 $\dot{\rho} + 3H(\rho + P) = C_{\rho}$ 







#### $m_s = 12 \text{ keV}, m_{\phi} = 36 \text{ keV}, \sin^2 2\theta = 2.5 \cdot 10^{-13}, y = 1.905 \cdot 10^{-4}$



 $T_{\nu} \ll m_{\phi} \rightsquigarrow$  transmission inefficient  $\rightsquigarrow$  exponential growth ends





#### $m_s = 20 \text{ keV}, m_{\phi} = 60 \text{ keV}, \sin^2 2\theta = 3.0 \cdot 10^{-15}, y = 1.602 \cdot 10^{-3}$



 $\theta \downarrow \rightsquigarrow y \uparrow \rightsquigarrow \nu_s \nu_s \rightarrow \phi \phi (\rightarrow 4 \nu_s)$  becomes relevant

- X-ray searches  $\rightsquigarrow$  upper limit on  $\theta$  as function of  $m_s$
- Lyman- $\alpha$  forest  $\rightsquigarrow$  lower limit on  $m_s$  as function of  $\theta$ 
  - $\nu_s \nu_s$  scatterings before kinetic decoupling
  - $\nu_{s}$  free-streaming after kinetic decoupling

• DM self-interactions  $\rightsquigarrow$  upper limit on  $y \rightsquigarrow$  lower limit on  $\theta$ 

# Viable Parameter Space



Yukawa coupling y chosen to yield observed Dark Matter density

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### Conclusions

- Sterile neutrinos: well-motivated, straightforward extension of SM
- $\bullet \ \text{Mass} \sim \text{keV} \rightsquigarrow \text{excellent } \textbf{Dark Matter candidate}$
- Automatically produced in early Universe
- Testable by X-ray searches ~> simplest scenario ruled out
- Minimal extension: addition of scalar singlet
   viable model with qualitatively new production mechanism

### Conclusions

- Sterile neutrinos: well-motivated, straightforward extension of SM
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- Automatically produced in early Universe
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- Minimal extension: addition of scalar singlet
   viable model with qualitatively new production mechanism
- Future tests
  - More sensitive X-ray searches
  - Better understanding of Ly- $\alpha$  forest
  - Stronger lower limit on m<sub>s</sub> from 21 cm observations
  - Stronger constraints on Dark Matter self-interactions
  - Suppression of astrophysical neutrino flux via  $\nu_a \nu_s \rightarrow \phi$
  - Impact on supernovae? Johns & Fuller, PRD 100 (2019)

# **Chemical Potentials**

DW production  $\rightsquigarrow p_s \sim p_a$  but  $n_s \ll n_a$  $\rightsquigarrow$  negative chemical potential  $\mu_s \ll -T_d$ 



 $BP2: \nu_{s}\nu_{s} \rightarrow \phi\phi \rightarrow 4\nu_{s} \rightsquigarrow n_{s} \uparrow, T_{d} \downarrow \rightsquigarrow \mu_{s} \uparrow$