# Majorana neutrinos and rare decays: where we are

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#### Work developed in collaboration with:

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- J. Menéndez, University of Barcelona
- F. Vissani, Laboratori Nazionali del Gran Sasso

A lot of content and plots from <u>arXiv:2202.01787</u>

## Historical landscape



~1950	Geochemical and direct experiments
~1970	First searches with Ge(Li) and scintillators
~1990	HPGe, TPCs, tracking detectors
~2000	O(10) kg HPGe sensitivity; KK claim 😲
~2010	O(100) kg HPGe, TPCs, liquid scintillators
now	moving towards the ton-scale!

# Which are the fundamental symmetries and conserved quantities?

- B,  $L_e$ ,  $L_\mu$ ,  $L_\tau$  and  $L=L_e+L_\mu+L_\tau$  are accidental symmetries that emerge without being *a priori* required
- Actually, these are anomalous symmetries spoiled by the full quantum theory
- Exact *global* symmetries of the standard model:

 $L_e^-L_u^- \rightarrow Violation$  seen in appearance mode by T2K

- $L_{\mu}$ - $L_{\tau} \rightarrow$  Violation seen in appearance mode by OPERA
- $\textbf{B-L} \quad \rightarrow \text{ We don't know yet}$



# What generated the matter-antimatter asymmetry in the Universe?

- Non-perturbative effect violate B+L, but can't explain *quantitatively* the asymmetry
   → We need a *dynamical* explanation yielding the asymmetry at some point in history: *baryogenesis*
- Possible solution: baryogenesis through *leptogenesis* 
  - $\rightarrow$  Observing the violation of B-L in the lepton sector would give a strong qualitative indication for the correctness of the baryogenesis hypothesis

# From $0\nu\beta\beta$ decay to the neutrino mass

Why are neutrino masses so small? Maybe, are they Majorana masses?

• Connection between  $0\nu\beta\beta$  decay rate and new physics terms:

$$\frac{\Gamma_{0\nu}}{\ln 2} = \frac{1}{T_{1/2}^{0\nu}} = \sum_{i} G_i g_i^4 M_i^2 f_i(\Lambda) + \text{interference terms},$$

• In the case of D=5 (light neutrino exchange):

$$\frac{1}{T_{1/2}^{0\nu}} = G_{01} g_A^2 \left( M_{light}^{0\nu} \right)^2 \frac{m_{\beta\beta}^2}{m_e^2}$$
$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}^2| e^{i\varphi_i} m_i \right|$$



## From $0\nu\beta\beta$ decay to the neutrino mass

$$m_{\beta\beta} \equiv |m_1 c_{12}^2 c_{13}^2 + m_2 s_{12}^2 c_{13}^2 e^{i\alpha_{21}} + m_3 s_{13}^2 e^{i(\alpha_{31} - \delta)}|$$



#### Majorana phases

$$m_{\beta\beta} \equiv |m_1 c_{12}^2 c_{13}^2 + m_2 s_{12}^2 c_{13}^2 e^{i\alpha_{21}} + m_3 s_{13}^2 e^{i(\alpha_{31} - \delta)}|$$

"Traditional" approach

- Assume worst- and best-case scenario and merge them
- Compute allowed space, no probability distribution!



#### Frequentist approach

- Maximum likelihood analysis
- Compute allowed space also for Majorana phases

20 meV

 $\rho_{\text{true}} = 0^{\circ}$ 

 $\rho_{\rm true} = 90^{\circ}$ 

 $\rho_{\rm true} = 180^\circ$ 

360°

270°

90°

0

270° Q 180°

90°

270°

م 180° 90° 0

Q 180°

10 meV

5 10

Exposure  $\xi$  (ton·yr)

 $|M_{0y}| = 3$ 

#### **Bayesian approach**

- Assume priors on phases
- Compute probability distribution on all parameters
- Strong dependence on priors!



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# Interplay between $0\nu\beta\beta$ decay and cosmology



- Robust constraints on Σ from cosmological observations
  - → Single results are model dependent, but limit fairly independent from selected dataset/assumption
- Possibility to evaluate discovery probability for m<sub>ββ</sub> with minimal assumptions

 0vββ decay triggered by BSM physics, but takes place in a highly correlated many-body system

 $\rightarrow$  Nuclear model(s) required to connect experimental sensitivity/results on  $0\nu\beta\beta$  decay rate to underlying BSM physics

- Connection to any BSM mechanism possible with just two approximations:
  - $\circ$  Light: m<sub>v</sub> << p ~ 200 MeV
  - Heavy:  $m_v >> p \sim 200 \text{ MeV}$



# Effective field theory

- BSM models defined at higher energy than electroweak scale (250 GeV), hadron scale (1 GeV) and nuclei (200 MeV)
  - $\rightarrow$  EFT suitable to study and organize all  $0\nu\beta\beta$  contributions
- Master formula for  $0\nu\beta\beta$  decay rate:

$$T_{1/2}^{-1} = G_{01} g_A^4 \left( M_{\text{long}}^{0\nu} + M_{\text{short}}^{0\nu} \right)^2 \frac{m_{\beta\beta}^2}{m_e^2} + \frac{m_N^2}{m_e^2} \tilde{G} \tilde{g}^4 \tilde{M}^2 \left( \frac{v}{\tilde{\Lambda}} \right)^6 + \frac{m_N^4}{m_e^2 v^2} \tilde{G'} \tilde{g'}^4 \tilde{M'}^2 \left( \frac{v}{\tilde{\Lambda'}} \right)^{10} + \cdots$$

- $\rightarrow$  Compute contribution due to *any* odd-dimensional operator
- $\rightarrow$  Phase spaces well known (computed)
- $\rightarrow$  NMEs computed by nuclear theory
- $\rightarrow$  Can be used to place limits on new physics scale
- $\rightarrow$  Short-range exchange of high-energy light neutrinos might not be negligible

# Many-body methods

Name	Method	Performance	
Nuclear Shell Model	Mix nuclear 1- and 2-shell configurations within given space	Good nuclear spectroscopy from O to Pb; effective Hamiltonian tuned on each nucleus	
Quasiparticle Random-Phase Approximation	Large configuration spaces with several shells; few nuclear correlations	Overestimated NMEs for non-spherical nuclei $\rightarrow$ fixed recently	
Interacting Boson Model	Models nucleus as set of bosons and map bosonic operators to nucleon DOF	No p-p correlations $\rightarrow$ Overestimates NMEs	
Energy-Density Functional	Mean-field description + correlations	Very good spectroscopy; Expensive inclusion of correlation → Overestimates NMEs	

- Systematic trend: for all nuclei, NSM predicts the smallest NME, EDF the largest
- Factor ~3 difference between the methods
- Most predicted NME for single-β decay overpredicts the experiment
  - $\rightarrow$  Caused by missing nuclear correlations or suboptimal transition operator
  - $\rightarrow$  Same effect expected for  $\beta\beta$  decay
  - $\rightarrow$  Commonly fixed by **ad-hoc** *quenching of*  $g_A$

## Ab-initio methods

- Treat explicitly all nucleons in the nucleus interacting with realistic forces
  - $\rightarrow$  Include nucleon-nucleon and three nucleon forces
  - $\rightarrow$  Include 1-body operators and 2-body correlations
- NMEs similar or smaller than from NSM

Name	Method	Performance
Quantum MC	Time evolution of trial nuclear state (via Hamiltonian) towards lowest-energy configuration	Most accurate for A≲12 Exponential scaling vs A
No-Core Shell Model	Extension of NSM: lowest-energy nucleons treated explicitly + high-energy orbitals	Limited to A≲22 Exponential scaling vs A
In-Medium Similarity Renormalization Group	Based on unitary transformation that simplify the solution of the many-body problem + correlations to reference state	Polinomial scaling vs A
In-Medium Generator Coordinate Method	Combination of reference state + GCM to include correlations and proton-proton pairing	Polinomial scaling vs A; smaller M <sub>long</sub> but enhanced M <sub>short</sub>
Coupled-Cluster	Up to triple correlations to reference state	Polinomial scaling, limited to spherical nuclei near magic ones

- QMC and NCSM can reproduce experimental NMEs of single-β decay with few percent precision
- Reliable NME calculation for  $0\nu\beta\beta$  decay not available yet  $\rightarrow$  Nuclear spectroscopy not matching data yet
- Seems to confirm that "the quenching of  $g_{A}$  is not a theory" (cit. F. Vissani)



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#### Heavy neutrino exchange



- Better general agreement, except for QRPA
- Large uncertainty due to ignored short-range correlations

## Experimental aspects: $0\nu\beta\beta$ decay signature



•  $0\nu\beta\beta$  decay can be searched only with isotopes for which the regular  $\beta$  decay is forbidden

- $0\nu\beta\beta$  decay signature is an excess at the Q-value ( $Q_{\beta\beta}$ ) in the sum electron spectrum
- 2vββ decay yields a continuum up to Q<sub>ββ</sub>

#### Isotope selection



- High Q-value  $(Q_{\beta\beta})$  is highly preferable:
  - Larger phase space
  - Lower background from natural radioactivity
- High isotopic abundance preferable  $\rightarrow$  Easier enrichment

Isotope	Daughter	$Q_{etaeta}{}^{\mathbf{a}}$	${f_{\mathrm{nat}}}^{\mathbf{b}}$	$f_{\rm enr}{}^{\rm c}$	$T_{1/2}^{2 uetaeta\mathrm{d}}$	$T_{1/2}^{0\nu\beta\beta e}$
		$[\mathrm{keV}]$	[%]	[%]	$[\mathbf{yr}]$	[yr]
$^{48}Ca$	$^{48}\mathrm{Ti}$	4267.98(32)	0.187(21)	16	$(6.4^{+0.7}_{-0.6}(\text{stat})^{+1.2}_{-0.9}(\text{syst})) \cdot 10^{19}$	$> 5.8 \cdot 10^{22}$
$^{76}\mathrm{Ge}$	$^{76}$ Se	2039.061(7)	7.75(12)	92	$(1.926 \pm 94) \cdot 10^{21}$	$> 1.8 \cdot 10^{26}$
$^{82}$ Se	$^{82}$ Kr	2997.9(3)	8.82(15)	96.3	$(8.60 \pm 0.03(\text{stat})^{+0.19}_{-0.13}(\text{syst})) \cdot 10^{19}$	$> 3.5 \cdot 10^{24}$
$^{96}$ Zr	$^{96}Mo$	3356.097(86)	2.80(2)	86	$(2.35 \pm 0.14(\text{stat}) \pm 0.16(\text{syst})) \cdot 10^{19}$	$> 9.2 \cdot 10^{21}$
$^{100}\mathrm{Mo}$	$^{100}$ Ru	3034.40(17)	9.744(65)	99.5	$(7.12^{+0.18}_{-0.14}(\text{stat}) \pm 0.10(\text{syst})) \cdot 10^{18}$	$> 1.5 \cdot 10^{24}$
$^{116}\mathrm{Cd}$	$^{116}$ Sn	2813.50(13)	7.512(54)	82	$(2.63^{+0.11}_{-0.12}) \cdot 10^{19}$	$> 2.2 \cdot 10^{23}$
$^{130}\mathrm{Te}$	$^{130}$ Xe	2527.518(13)	34.08(62)	92	$(7.71^{+0.08}_{-0.06}(\text{stat})^{+0.12}_{0.15}(\text{syst})) \cdot 10^{20}$	$> 2.2 \cdot 10^{25}$
$^{136}$ Xe	$^{136}$ Ba	2457.83(37)	8.857(72)	90	$(2.165 \pm 0.016(\text{stat}) \pm 0.059(\text{syst})) \cdot 10^{21}$	$> 1.1 \cdot 10^{26}$
$^{150}\mathrm{Nd}$	$^{150}\mathrm{Sm}$	3371.38(20)	5.638(28)	91	$(9.34 \pm 0.22(\text{stat})^{+0.62}_{-0.60}(\text{syst})) \cdot 10^{18}$	$> 2.0 \cdot 10^{22}$

Detector type	Description	Typical isotope mass	Containment efficiency	Pros and cons
Solid state granular	Crystals containing ββ isotope	0.1-1 kg (single crystal)	70-95%	<ul><li>+ Scalable</li><li>+ High resolution</li><li>- Many crystals required</li></ul>
Monolithic liquid or gaseous	Material = isotope or isotope dissolved in liquid	100-1000 kg	100%	<ul><li>+ Single detector</li><li>+ Self shielding</li><li>- Poor resolution</li></ul>
Composite	Foil with isotope in low-pressure gas tracker	10 kg	~50%	+ Ultra-low background + Multiple isotopes - Very hard to scale

Readout channel	Energy resolution	Particle identification	Sensitivity to position	Applicable to multiple isotopes
Ionization	0.1-1%	Only in gas	Yes	Not really
Phonons	~0.2%	Nope	Nope	Yes
Scintillation	Few %	a vs β	In liquids and gases	Yes
Cherenkov	Forget it!	Visible only for $\beta$ 's	Maybe	Yes

# $0\nu\beta\beta$ decay experimental fauna



#### Germanium experiments



- Low Q-value: 2039 keV
- Highest energy resolution: ~0.1% FWHM
- **Discrimination of \beta\beta vs a vs \gamma + possibility of active \gamma veto** 
  - $\rightarrow$  Extremely low bkg: ~5·10<sup>-4</sup> counts/keV/kg/yr reached by GERDA
  - $\rightarrow$  Operating next to linear sensitivity regime
- MAJORANA and GERDA completed, joining for next-gen experiment LEGEND

#### $\rightarrow$ GERDA talk by T. Comellato, LEGEND talk by R. Brugnera



#### Majorana demonstrator



## Gas-Xe TPCs: NEXT

- High-pressure ~10 bar with double readout (ionization + scintillation)
- Energy resolution ~1% FWHM demonstrated
- Particle tracking
  - $\rightarrow$  Discrimination of  $\beta\beta$  from a and single-  $\beta$  or  $\gamma$  events
- 2vββ decay measured by NEXT-White using 3.5 kg of Xe only!
- NEXT-100 under construction @ Canfranc
  - $\rightarrow$  Expected background 5·10<sup>-4</sup> counts/keV/kg/yr
  - $\rightarrow$  Expected resolution 0.5-0.7% FWHM





#### Liquid scintillator experiments

#### KamLAND-Zen



- Very large volume
  - $\rightarrow$  Isotope in central part
  - $\rightarrow$  Highly effective **self**

#### shielding

- Isotope (<sup>130</sup>Te or <sup>136</sup>Xe) dissolved in liquid scintillator
  - $\rightarrow \text{Easily } \textbf{scalable}$
  - $\rightarrow$  Enrichment not strictly required
- High intrinsic backgrounds
- Readout of scintillation only
   → Particle identification
  - possible
  - $\rightarrow$  Energy resolution of few %

Good for limit setting, not for discovery

 $\rightarrow$  KamLAND-Zen talk by K. Ichimura

SNO+



# Liquid scintillator experiments: SNO+

- Repurposed SNO experiment with 0.5% Te-loaded scintillator
- Te-loading not as easy as Xe-loading
- Separate background and signal+background measurements
  - 2017-2019 as pure-water Cherenkov detector  $\rightarrow$  Characterization of detector response
  - 2019-now with unloaded LAB scintillator
    - $\rightarrow$  Study scintillator response
    - $\rightarrow$  Study scintillator-related backgrounds
  - Late 2022-20?? with Te-loaded scintillator
    - $\rightarrow$  Looking forward to hear good news!





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#### **Cryogenic calorimeters**



# Tracking calorimeters: SuperNEMO

- Measure both energy and momentum
  - $\rightarrow$  Background suppression
  - $\rightarrow$  Single electrons resolved
  - $\rightarrow$  Possible to study  $0\nu\beta\beta$  decay mechanism
- Source ≠ detector
  - $\rightarrow$  Limited isotope mass
  - $\rightarrow$  Any isotope is usable
- Perfect technology for precision measurement of  $0\nu\beta\beta$  and  $2\nu\beta\beta$  decay
- SuperNEMO demonstrator under commissioning









# **THANK YOU!**