

Quantum sensor applications to (low and high)-energy physics

M. Doser, CERN

“Sala TV” but no popcorn...

What are the challenges?

① Some words on the landscape

Clarification of terms

Quantum sensors for low energy particle physics

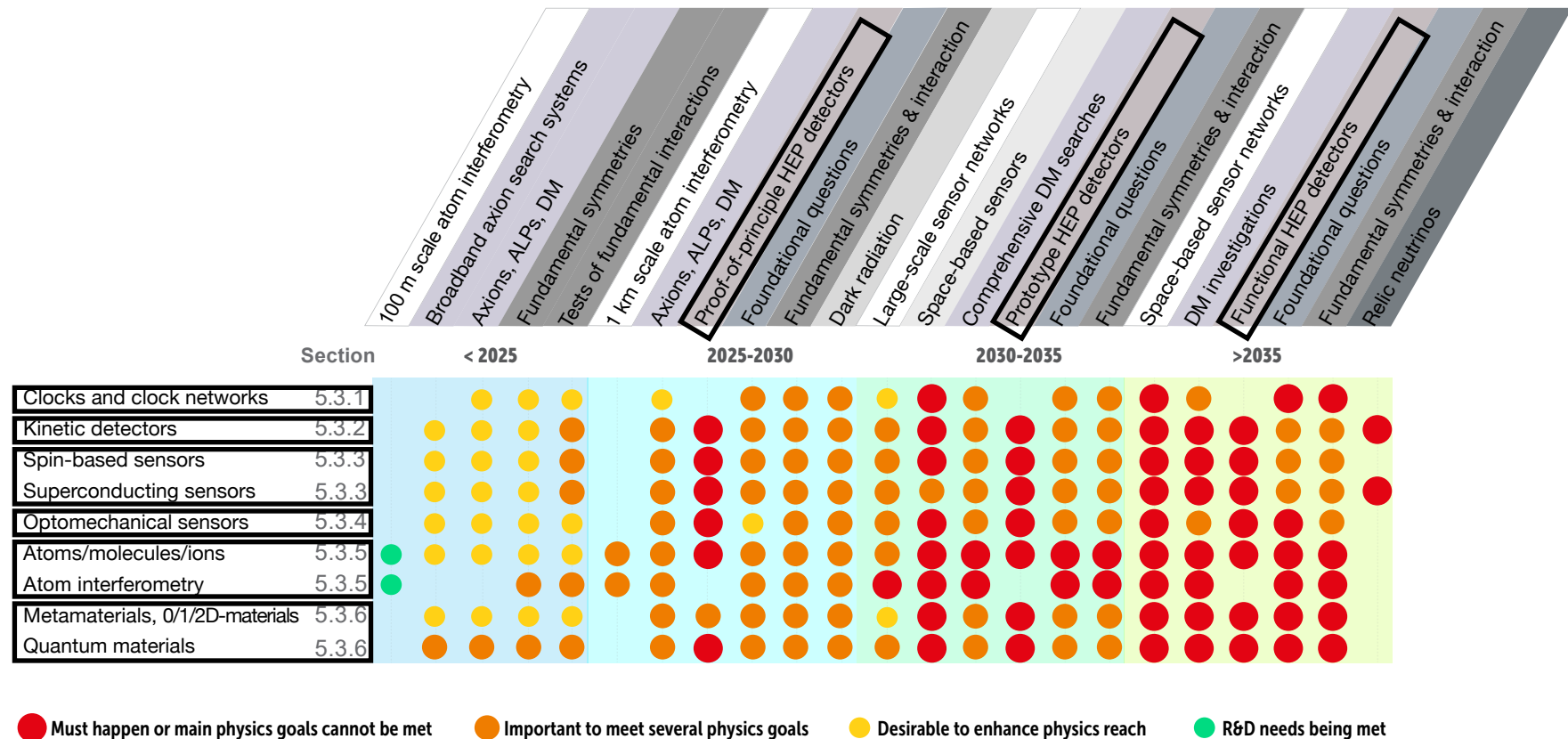
② Quantum sensors for new particle physics experiments

③ Quantum detectors for high energy particle physics

RECFA Detector R&D roadmap 2021

<https://cds.cern.ch/record/2784893>

Chapter 5: Quantum and Emerging Technologies Detectors



Chapter 4: Particle Identification and Photon Detectors

It is recommended that several “blue-sky” R&D activities be pursued. The development of solid state photon detectors from novel materials is an important future line of research, as is the development of **cryogenic superconducting photosensors** for accelerator-based experiments. Regarding advances in PID techniques, gaseous photon detectors for visible light should be advanced. **Meta-materials** such as **photonic crystals** should be developed, giving **tune-able refractive indices** for PID at high momentum. Finally, for TRD imaging detectors, the detection of transition radiation with silicon sensors is an important line of future research.

ECFA Detector R&D Roadmap Symposium of Task Force 5 Quantum and Emerging Technologies

Symposium: April 12, 2022

<https://indico.cern.ch/event/999818/>

14 presentations

first block covering
physics landscape

following blocks
focusing on
technologies

discussion of three
important points

ECFA Detector R&D Roadmap Symposium of Task Force 5: Quantum and emerging technologies

Monday 12 Apr 2021, 09:00 → 18:30 Europe/Zurich

09:00 → 09:15 **Introduction**

09:15 → 11:00 science targets – Overview and Landscape

9:15 EDM searches & tests of fundamental symmetries **Peter Fierlinger / TU Munich**

9:45 **Tests of QM** [wavefunction collapse, size effects, temporal separation, decoherence]

10:15 Multimessenger detection [including atom interferometer or magnetometer networks] **Giovanni Barontoni / Birmingham**

10:45 Axion and other DM (as well as non-DM Ultra-light) particle searches **Mina Arvanitaki / Perimeter Institute**

11:15 → 11:30 **Coffee break**

11:30 → 12:30 Experimental methods and techniques - Overview and Landscape

11:30 Precision spectroscopy and clocks, networks of sensors and of entangled systems [optical atomic clocks] **David Hume / NIST**

12:00 Novel ionic, atomic and molecular systems [RaF, multiatomic molecules, exotic atoms] **Marianna Safranova / U. Delaware**

12:30 → 13:30 **Lunch break**

13:30 → 16:00 Experimental and technological challenges, New Developments

13:30 **Superconducting platforms** [detectors: TES, SNSPD, Haloscopes, including single photon detection]

14:00 High sensitivity superconducting cryogenic electronics, low noise amplifiers **Stafford Withington / Cambridge**

14:30 Broadband axion detection **Kent Irwin / Stanford**

15:00 Mechanical / optomechanical detectors **Andrew Geraci / Northwestern**

15:30 Spin-based techniques, NV-diamonds, Magnetometry **Dima Budker / Mainz**

16:00 → 16:15 **Coffee break**

16:15 → 18:30 Experimental and technological challenges, New Developments

16:15 Calorimetric techniques for neutrinos and axions **potential speaker identified**

16:35 Quantum techniques for scintillators **potential speaker identified**

16:55 Atom interferometry at large scales (ground based, space based) **Jason Hogan / Stanford**

17:25 → 18:15 **Discussion session** : discussion points

- Scaling up from table-top systems
- Networking – identifying commonalities with neighboring communities
- Applying quantum technologies to high energy detectors

18:15 → 18:30 **Wrap-up**

CERN quantum initiative

<https://quantum.web.cern.ch/>



- Assess the **areas of potential quantum advantage** in HEP applications (QML, classification, anomaly detection, tracking)
- Develop **common libraries of algorithms, methods, tools**; benchmark as technology evolves
- Collaborate to the development of shared, **hybrid classic-quantum infrastructures**
-

Computing & Algorithms



- Identify and develop techniques for **quantum simulation** in collider physics, QCD, cosmology within and beyond the SM
- Co-develop quantum computing and sensing approaches by providing **theoretical foundations** to the identifications of the areas of interest
-

Simulation & Theory



- Develop and promote **expertise in quantum sensing** in low- and high-energy physics applications
- Develop quantum sensing approaches with emphasis on **low-energy particle physics measurements**
- Assess **novel technologies and materials** for HEP applications
-

Sensing, Metrology & Materials

currently: 2.5 PhD's



- **Co-develop CERN technologies relevant to quantum infrastructures** (time synch, frequency distribution, lasers)
- Contribute to the **deployment and validation of quantum infrastructures**
- Assess requirements and **impact of quantum communication on computing applications** (security, privacy)

Communications & Networks

<https://quantum.web.cern.ch/>

quantum sensors register a change of quantum state caused by the interaction with an external system:

- transition between superconducting and normal-conducting
- transition of an atom from one state to another
- change of resonant frequency of a system (quantized)

Then, a "quantum sensor" is a device, the measurement (sensing) capabilities of which are enabled by our ability to manipulate and read out its quantum states.

and because the commensurate energies are very low, unsurprisingly, quantum sensors are ideally matched to low energy (particle) physics;

→ focus on CERN activities both in low energy and high energy particle physics

(I will not however be talking about entanglement and its potential applications)

particle physics: what are we talking about?

Background

- Light DM candidate have large mode volume occupation number -> can be treated as **classical fields**
- QCD Axions and ALPs $\mathcal{L}_{axion} \supset \sum_f \frac{c_f}{\Lambda} \partial_\mu a \bar{f} \gamma^\mu \gamma^5 f \rightarrow H \propto \sum_f \nabla a \cdot \mathbf{S}_f$
 - ∇a acts as a pseudo magnetic field ➔ can be detected by **atomic magnetometers**
- Scalar fields $\mathcal{L}_{scalar} \supset \frac{\phi^n}{\Lambda_\gamma^n} F_{\mu\nu} F^{\mu\nu} - \sum_f \frac{\phi^n}{\Lambda_f^n} m_f \bar{f} f$
 - Λ_γ^n alter the fine structure constant α , Λ_f^n the fermionic masses -> manifest as **variations of fundamental constants**
- symmetry violations probed via **precision measurements** (~~CP~~, ~~CPT~~, ~~Lorenz~~, ~~WEP~~, ...) ➔ SME

D. Colladay, V.A. Kostelecký, "CPT violation and the standard model".
Physical Review D. 55 (11) (1997) 6760–6774. [arXiv:hep-ph/9703464](https://arxiv.org/abs/hep-ph/9703464)

quantum sensors & particle physics: what are we talking about?

domains of physics

search for NP / BSM

Axions, ALP's, DM & non-DM
UL-particle searches

tests of QM

wavefunction collapse,
decoherence

EDM searches & tests of
fundamental symmetries

quantum technologies

① superconducting devices (TES, SNSPD, ...) / cryo-electronics

② spin-based, NV-diamonds

③ optical clocks

④ ionic / atomic / molecular

⑤ optomechanical sensors

⑥ metamaterials, 0/1/2-D materials

1 superconducting devices (TES, SNSPD, ...) / cryo-electronics

superconducting bolometers in the form of a transition edge sensor (TES) : 1990

invention of kinetic inductance detectors (KID): 2000's

Photons incident on a strip of superconducting material break Cooper pairs and create excess quasiparticles. The kinetic inductance of the superconducting strip is inversely proportional to the density of Cooper pairs, and thus the kinetic inductance increases upon photon absorption. This inductance is combined with a capacitor to form a microwave resonator whose resonant frequency changes with the absorption of photons

introduction of the travelling wave parametric amplifier (TWPA) for quantum noise limited coherent amplification of large bandwidth: 2010

Superconducting qubits for quantum computing, superconductor/spin-system quantum memory elements: 2020

Superconducting device physics is a whole technology, based on thin film deposition techniques and lithographic patterning and allowing large scale integration with high degree of functionality

Superconducting nanowire single photon detector

Superconducting quantum interference device

Josephson junction parametric amplifier

	Microwave	Submillimetre	Far infrared	Optical	High energy
	10 – 100 GHz 3 cm- 3 mm	100 GHz – 1 THz 3 mm – 300 µm	1 – 10 THz 300 – 30 µm	2 µm – 300 nm	UV, Yray and Xray
SIS mixers		•			
HEB			•		
CEB		•			
TES	•	•	•	•	•
KID	•	•	•	•	
SNSPD			•	•	
SQUID	•				
JJPA	•				
TWPA	•	•			

1 superconducting devices (TES, SNSPD, ...) / cryo-electronics

superconducting bolometers in the form of a transition edge sensor (TES) : x-ray imaging

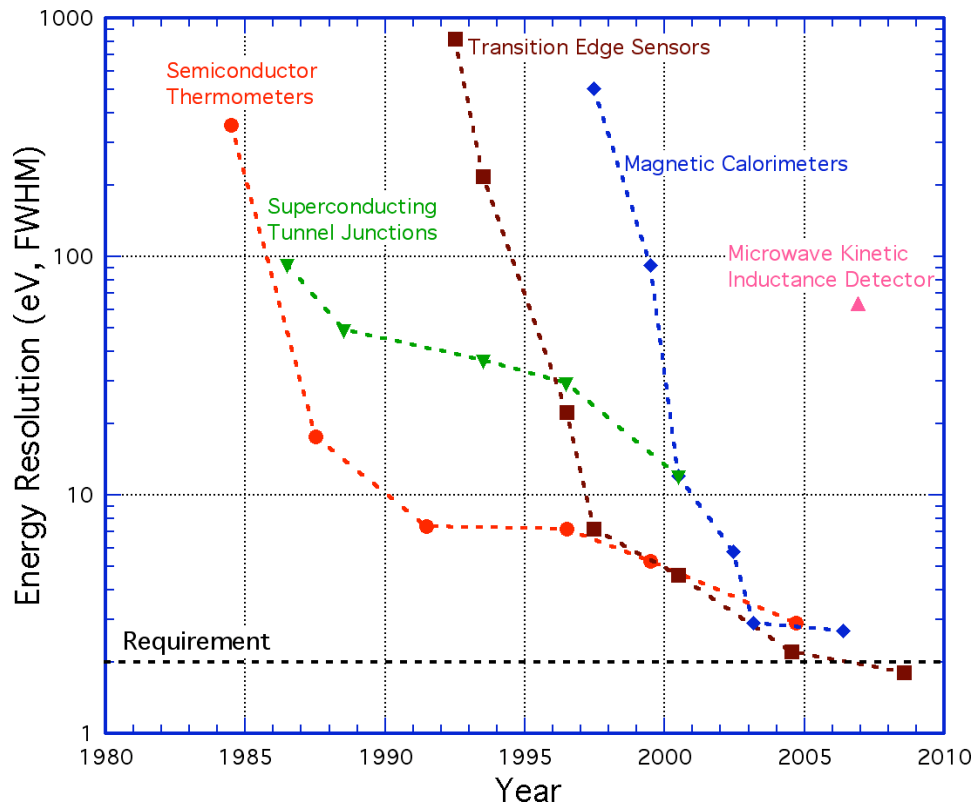
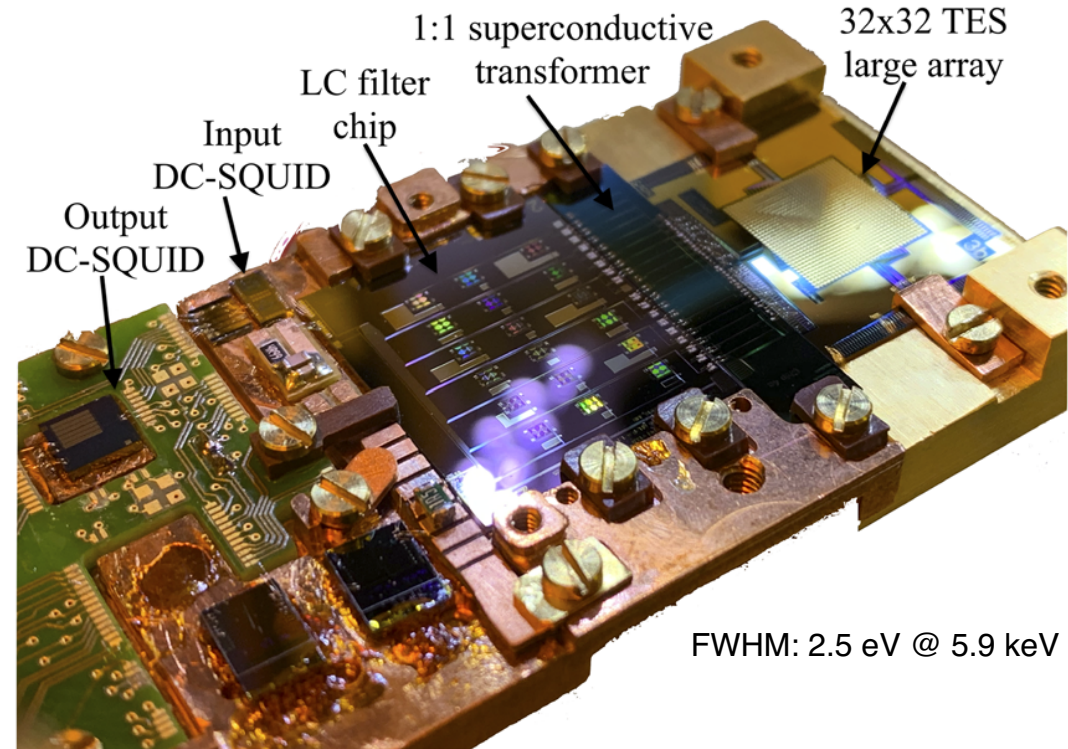


Fig. 1. Time evolution of energy resolution at 6 keV as a function of time for three microcalorimeter technologies.

Development of Low-Temperature Detectors For Generation-X and Other Missions Requiring High-Resolution, Large-Format, X-ray Detector Arrays
A White Paper submitted to *Electromagnetic Observations from Space (EOS)* Discipline Program, Simon Bandler *et al.*



Picture of the setup used to characterize the kilo-pixels array hung on the mixing chamber of a dilution refrigerator and held at a base temperature of 50 mK. Main parts are highlighted by arrows.

Performance and uniformity of a kilo- pixel array of Ti/Au transition-edge sensor microcalorimeters
E. Taralli, M. D'Andrea, L. Gottardi, et al., *Rev. Sci. Instrum.* 92, 023101 (2021);
<https://doi.org/10.1063/5.0027750>

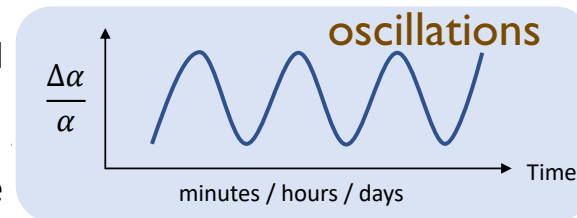
2 spin-based, NV-diamonds, ...

search for NP / BSM

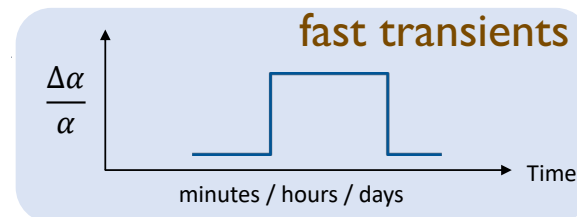
3 optical clocks

Signal characteristics

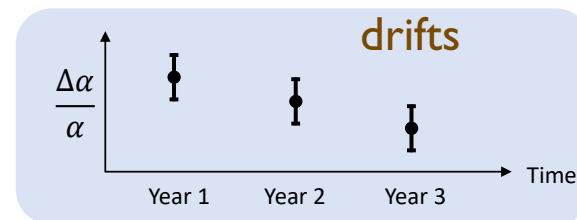
- The only possibility of detecting **transient** events such as topological defects, solitons, Q balls and dark stars
- **Oscillations** of dark matter fields at different locations as long as the distance is below the coherence length (100 km: mass $\sim 10^{-9}$ eV)
- Sensors with **similar sensitivities and different systematics** are necessary to confirm any measurements and reject false positives
- Using multiple sensors increases the detection confidence and sensitivity
- **Multimessenger** detection, discriminating between different couplings



Very light DM



DM- topological defects

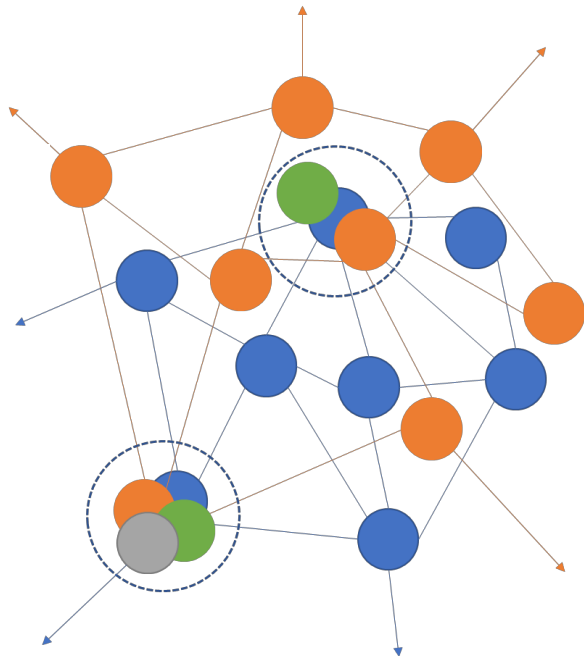


New physics

particle physics: what are we talking about?

search for NP / BSM

networks of sensors



magnetometers

2

Afach et al, arXiv:2102.13379v2

atomic clocks

nuclear, HCl,
molecules

3

Wcislo et al, Sci.Adv. 4, 4869 (2018)

optical fiber networks

3

Roberts et al, New J. Phys. 22, 093010 (2020)

Investigate very light scalar and pseudo-scalar DM candidates
over ~ 10 orders of magnitude in mass and different couplings

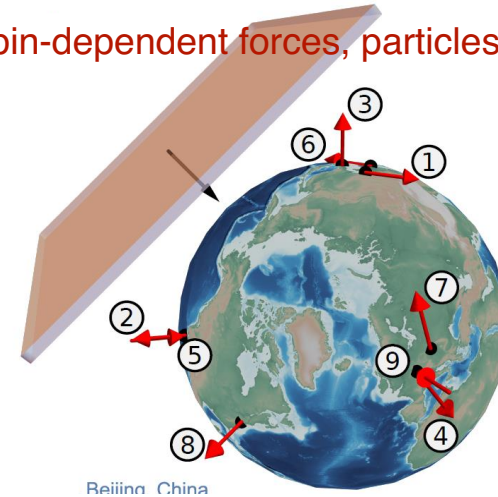
2 spin-based, NV-diamonds, ...



GNOME: arXiv:2102.13379 (2021)

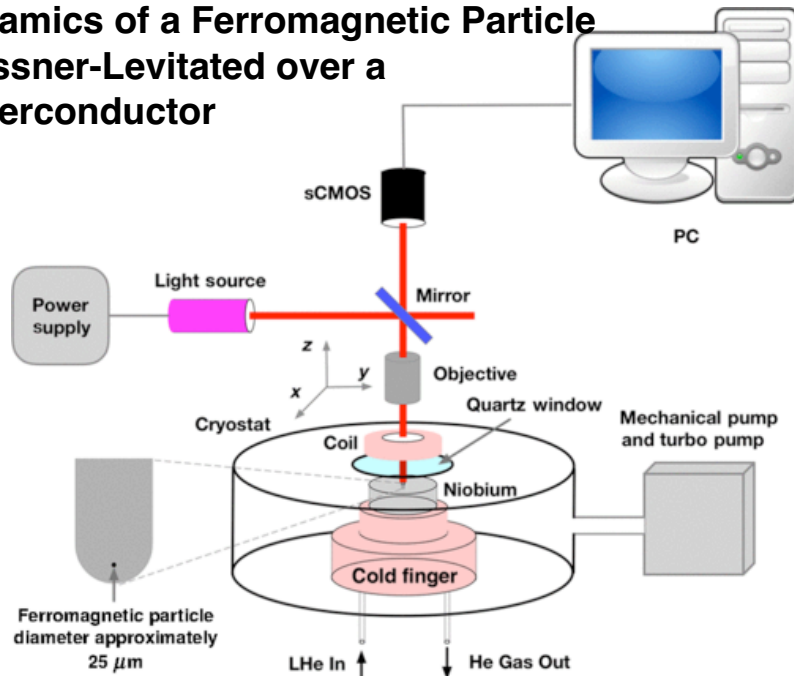
Global Network of **Optical (atomic)**
Magnetometers for Exotic searches

→ Spin-dependent forces, particles coupling to spin

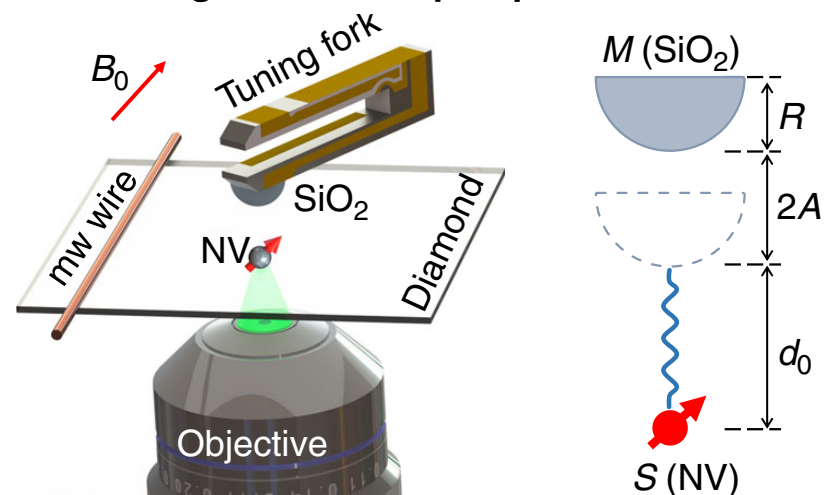


ALP's can interact with atomic spins; passage of an
ALP domain wall → simultaneous / correlated signals
~ **transient magnetic field pulse** (down to O(pT))

Dynamics of a Ferromagnetic Particle Meissner-Levitated over a Superconductor



Searching for an exotic spin-dependent interaction with a single electron-spin quantum sensor



Rong, X., Wang, M., Geng, J. *et al.*, *Nat Commun* **9**, 739 (2018).
<https://doi.org/10.1038/s41467-018-03152-9>

3 optical clocks

- **Atomic clocks** measure with extreme precision atomic and molecular spectra

- Spectroscopy lends itself to measure variations of:

$$\alpha = \frac{1}{4\pi\epsilon_0} \frac{e^2}{\hbar c}$$

$$\mu = \frac{m_p}{m_e}$$

- Different clocks have different sensitivities to variations of α and μ

$$\frac{\delta\omega}{\omega} = K_\alpha \frac{\delta\alpha}{\alpha} + K_\mu \frac{\delta\mu}{\mu}$$

- Clocks are “naturally” networked, need to compare at least 2

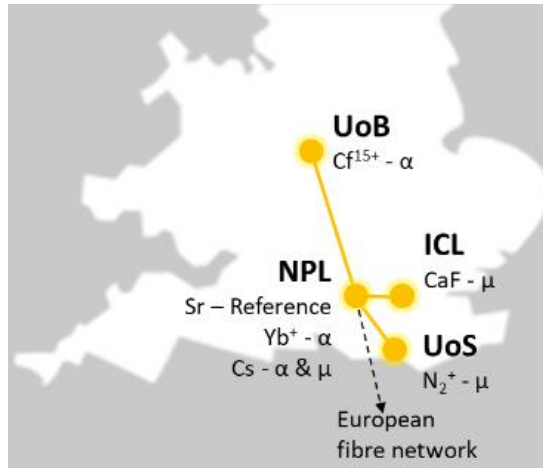
A common, stable, and insensitive frequency reference (a Sr clock), against which all the clocks of the network can measure variations

Clock		K_α	K_μ
Higly-charged ion clock	Cf ¹⁵⁺ (775 nm)	59	0
Atomic clock	Yb ⁺ (467 nm)	-5.95	0
Molecular ion clock	N ₂ ⁺ (2.31 μm)	0	0.5
Molecular clock	CaF (17 μm)	0	0.5
Atomic clock	Sr (698 nm)	0.06	0
	Cs (32.6 mm)	2.83	1

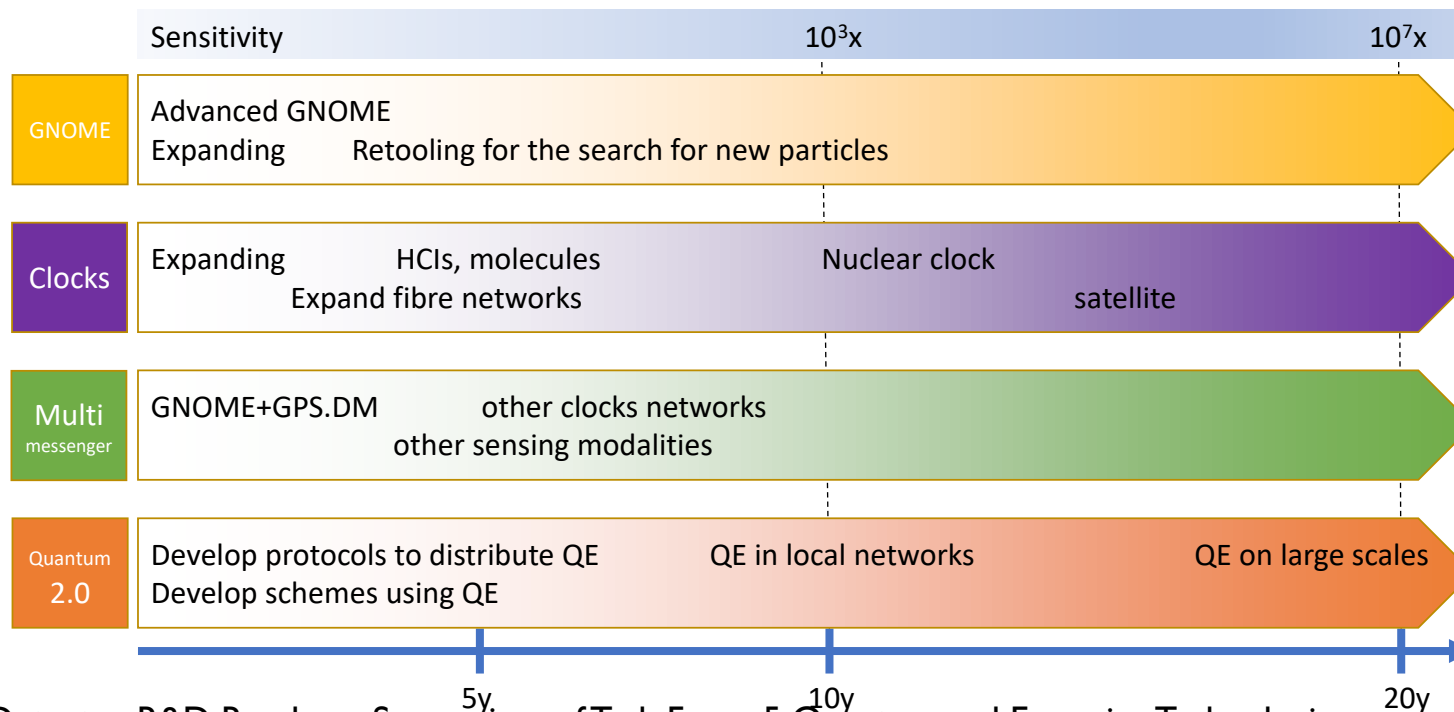
Quantum 2.0

- Entanglement between sensors can give an advantage when measuring multiple non-commuting parameters of dealing with “nuisance” parameters [PRA 95, 012326 (2017)]
- Creating a super-stable global network of clocks synchronized with entanglement [Nat. Phys. 10, 582 (2014)]
- Need more measurement schemes

3 optical fiber networks



QSNET



ECFA Detector R&D Roadmap Symposium of Task Force 5 Quantum and Emerging Technologies

<https://indico.cern.ch/event/999818/>

Giovanni Barontini (Birmingham)

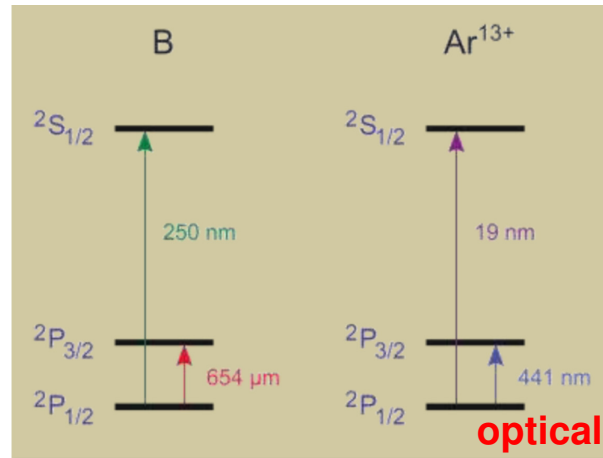
4 ionic / atomic / molecular

novel systems with greatly enhanced sensitivity wrt standard atomic systems

HCI's

e.g. $^{18}\text{Ar} \rightarrow \text{Ar}^{13+}$
(5 e⁻ remain)

Lorentz violation searches
variation of fund. constants
DM searches
tests of QED
5th force searches



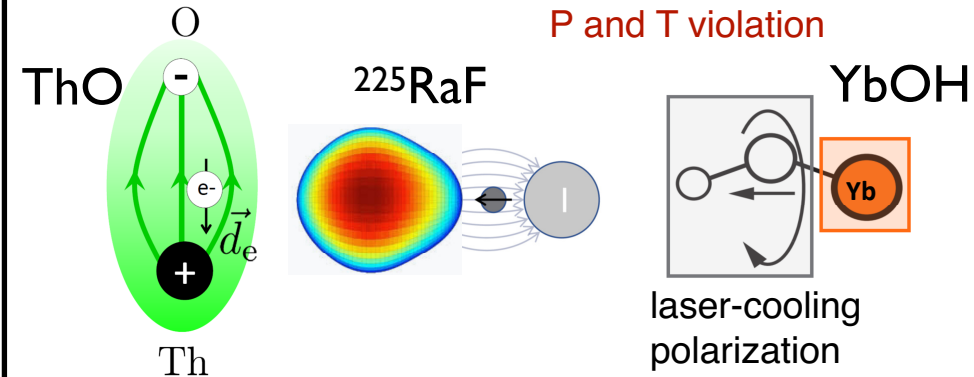
Quantum Sensors for New-Physics Discoveries
<https://iopscience.iop.org/journal/2058-9565/page/Focus-on-Quantum-Sensors-for-New-Physics-Discoveries>

Search for new physics with atoms and molecules
Rev. Mod. Phys. 90, 025008 (2018)

Barontini et al. EPJ Quantum Technology (2022) 9:12
<https://doi.org/10.1140/epjqt/s40507-022-00130-5>

eEDM's in molecules

P and T violation



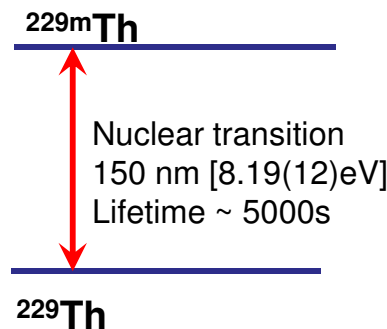
nuclear clock (^{229}Th)

typical nuclear energy levels are in MeV. Six orders of magnitude from ~few eV we can access by lasers!

only ONE exception:

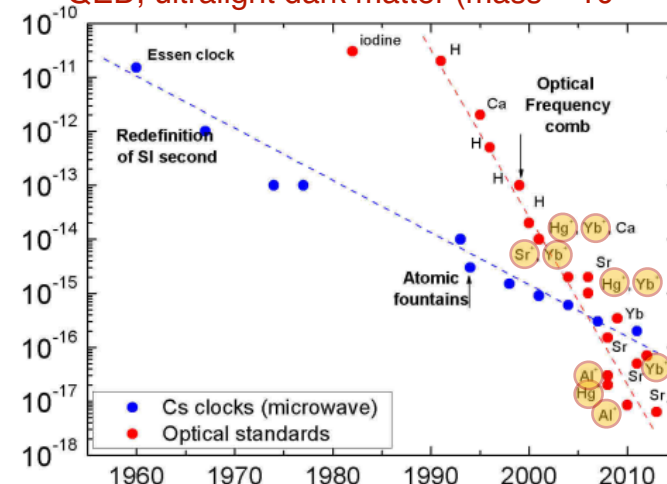
sensitivity to dark matter, variation of α
~ 8.19 eV / MeV ~ 10⁵ wrt current clocks

Review: E. Peik, et al., arXiv:2012.09304, in press, Quantum Science and Technology (2021)



molecular / ion clocks

QED, ultralight dark matter (mass ~ 10⁻²² – 10⁻¹⁵ eV)



Hg⁺
Al⁺
Yb⁺
CaF
N₂⁺

N₂⁺: 10⁻¹⁸ <https://doi.org/10.1103/PhysRevA.89.032509>

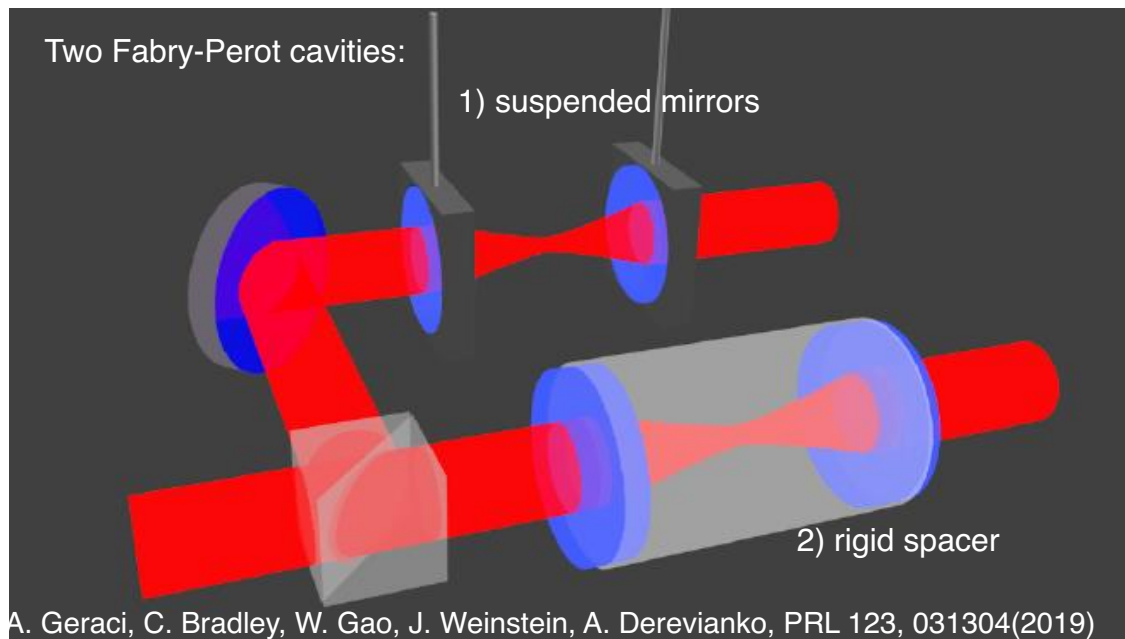
ECFA Detector R&D Roadmap Symposium of Task Force 5 Quantum and Emerging Technologies

<https://indico.cern.ch/event/999818/>

Marianna Safronova (University of Delaware) David Hume (NIST)

5 optomechanical sensors

- Wavelike DM
 - Scalar-like: Optical-cavity-based detectors
 - Vector-like: Accelerometer detectors
- Particle/extended-object DM
 - Gravitational Wave detectors for DM (also Axions)
 - Levitated microspheres
 - Windchime experiment



Isotropic strain in material objects due to variation of atomic size from interaction with ultra-light DM [Manley et al. PRL 124, 151301 \(2020\)](#)

$$\phi(t, \mathbf{r}) \approx \frac{\hbar}{m_\phi c} \sqrt{2\rho_{\text{DM}}} \cos[2\pi f_\phi t - \mathbf{k}_\phi \cdot \mathbf{r} + \dots]$$

$$\mathbf{k}_\phi = m_\phi \mathbf{v} / \hbar$$

m_ϕ is the mass of the DM field,
 \mathbf{v} is the relative speed of the DM

$$\frac{\delta m_e(t, \mathbf{r})}{m_{e,0}} = d_{m_e} \sqrt{4\pi\hbar c} E_P^{-1} \phi(t, \mathbf{r})$$

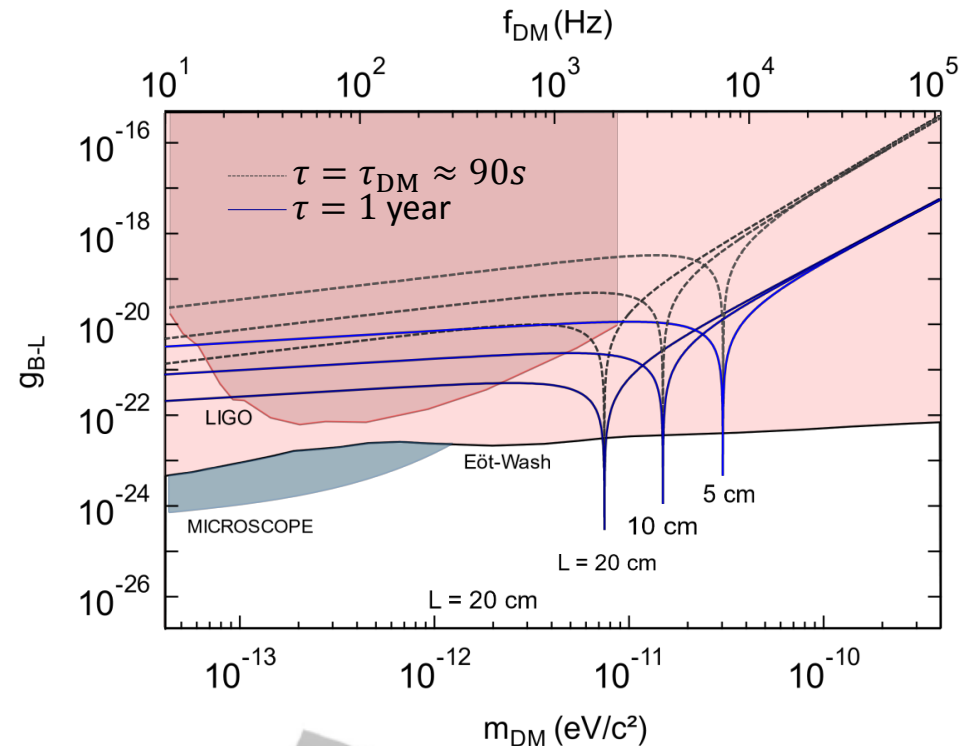
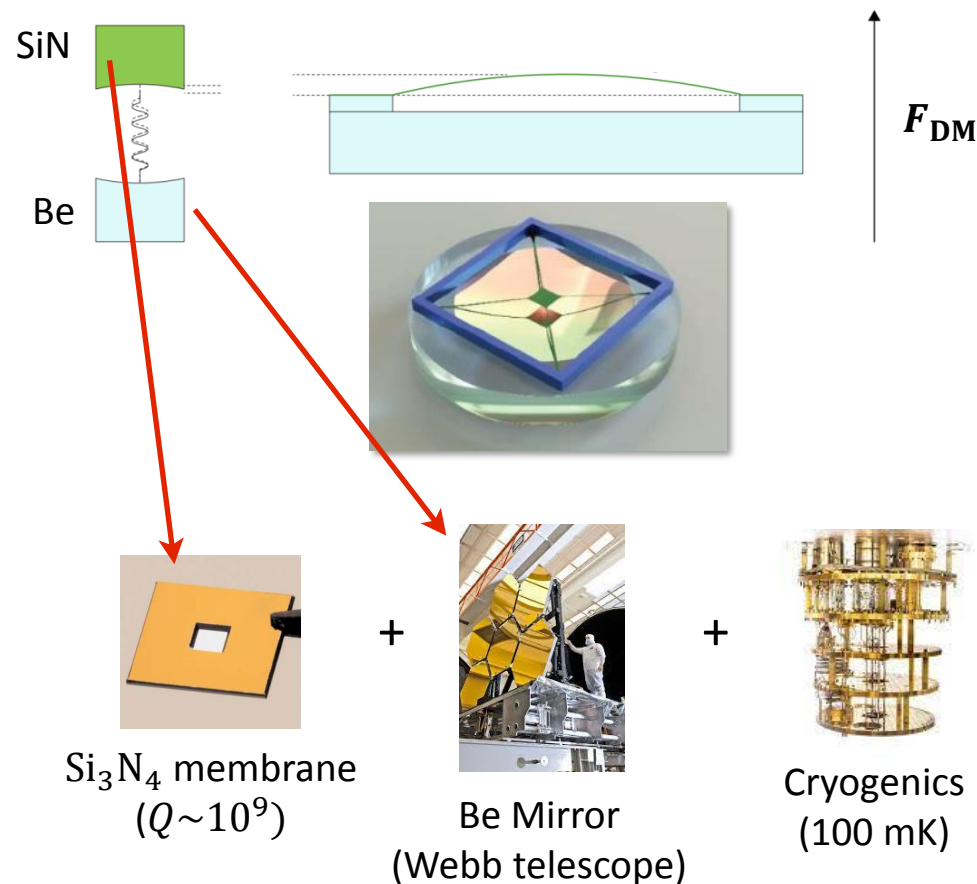
$$\frac{\delta \alpha(t, \mathbf{r})}{\alpha_0} = d_e \sqrt{4\pi\hbar c} E_P^{-1} \phi(t, \mathbf{r})$$

strain :

$$h = -\frac{\delta \alpha}{\alpha_0} - \frac{\delta m_e}{m_{e,0}},$$

accelerometers

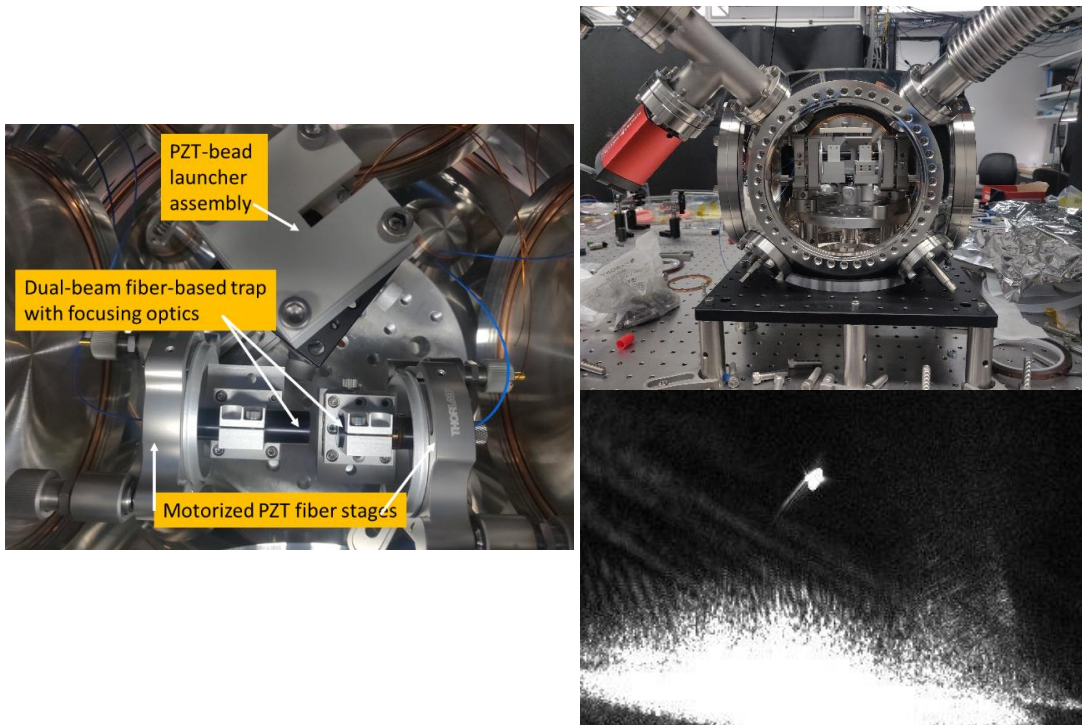
differential acceleration between materials of different composition (probe of B-L couplings)



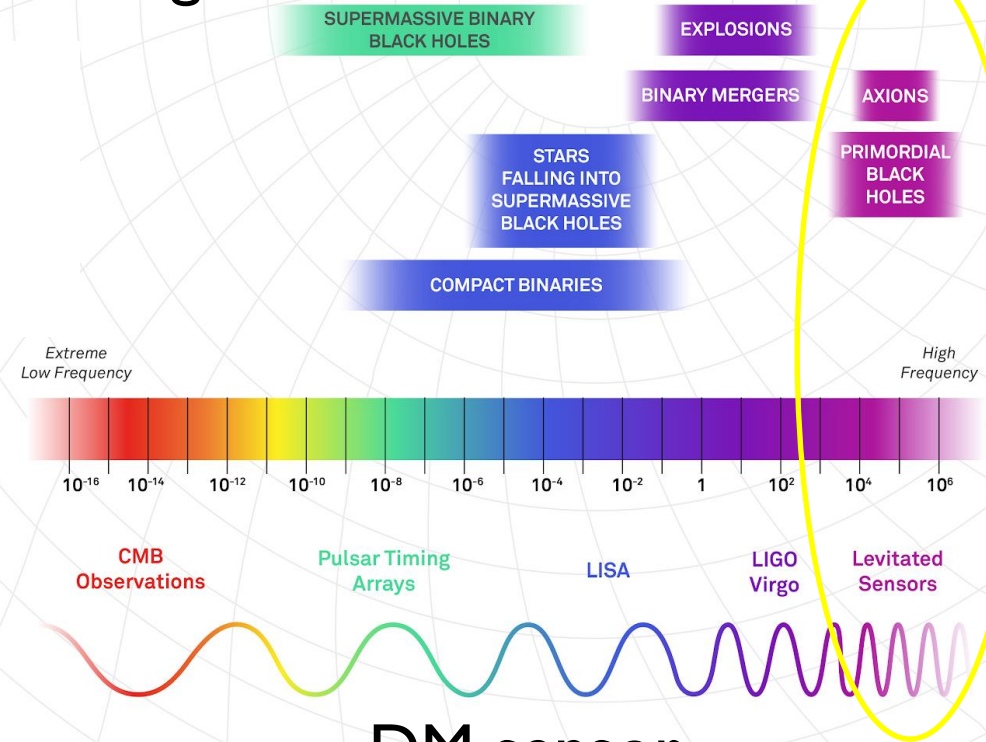
Manley et. al. PRL 126 061301 (2021)

5 levitated microspheres

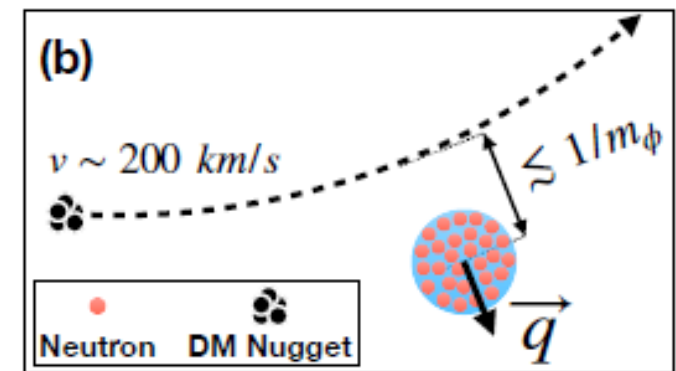
nanoparticle standing-wave trap (optical lattice)



gravitational wave sensor



DM sensor



Search for Composite Dark Matter with Optically Levitated Sensors

Fernando Monteiro, Gadi Afek, Daniel Carney, Gordan Krnjaic, Jiaxiang Wang, and David C. Moore
Phys. Rev. Lett. **125**, 181102 – Published 28 October 2020

= zeptonewton sensing

ECFA Detector R&D Roadmap Symposium of Task Force 5 Quantum and Emerging Technologies

<https://indico.cern.ch/event/999818/>

Andrew Geraci (Northwestern)

5

“Wind chime”

Planck-scale DM: measure the gravitational effect of flying-by DM on an array of accelerometers

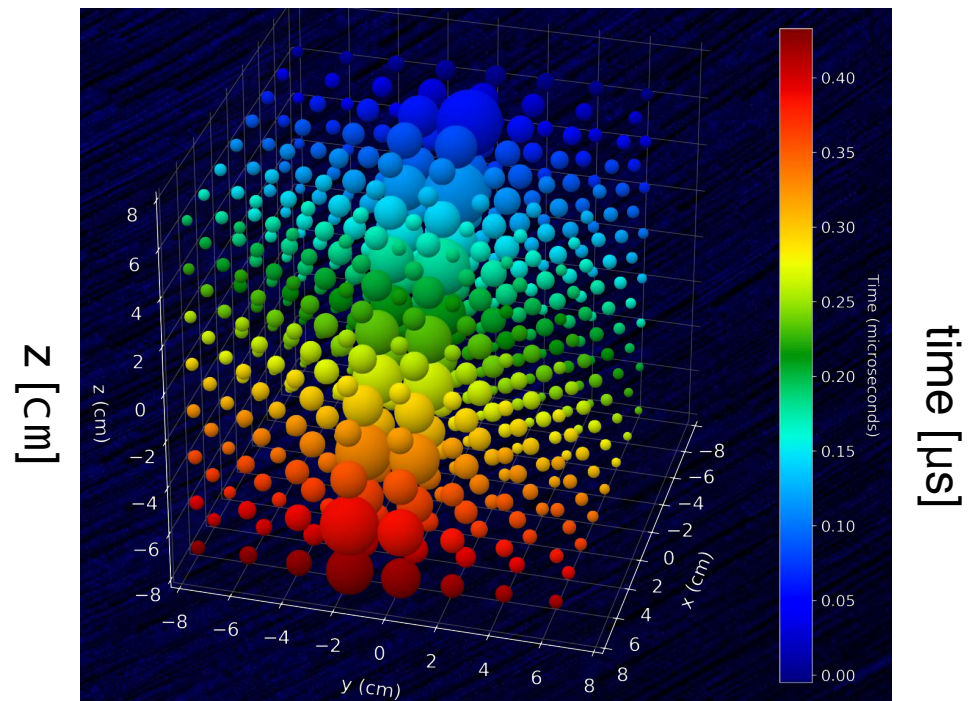
vision: “Our proposed strategy is to build a **three-dimensional array of force sensors**. A heavy DM particle passing through the array will exert a **small but correlated force** on the sensors nearest its trajectory. Much **like tracking a particle in a bubble chamber**, we can then pick out this correlated force signal along the DM “track” through the array.”

<https://arxiv.org/abs/1903.00492>

flux: $1/\text{m}^2/\text{year}$

“With a **billion** detectors at the gram scale, Planck-scale gravitational DM detection is achievable.”

sensors: “cryogenic opto-mechanical devices”, e.g.: atoms in a lattice that are continuously optically probed



6 metamaterials, 0/1/2-D materia

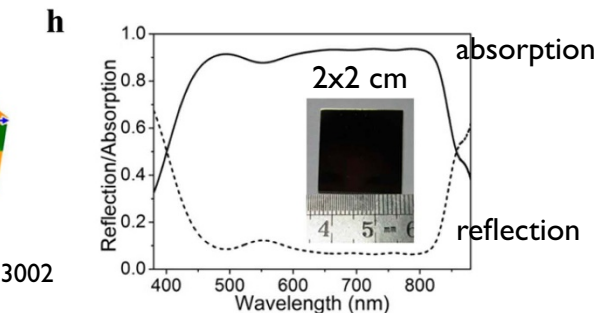
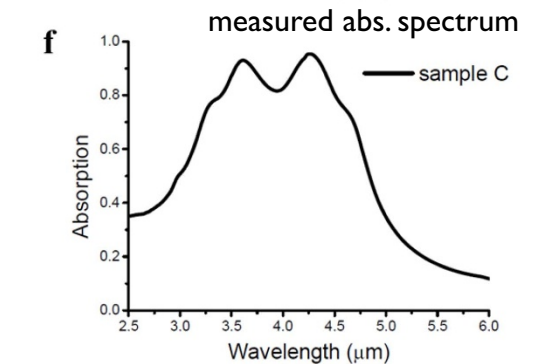
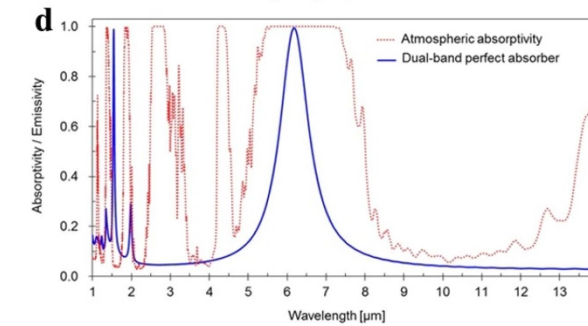
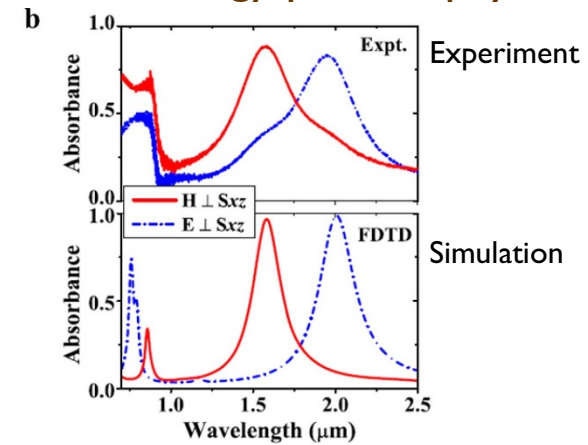
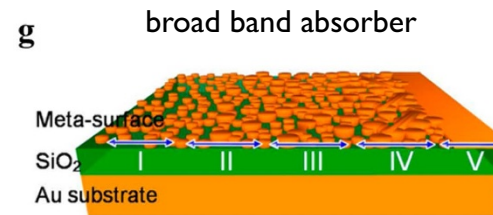
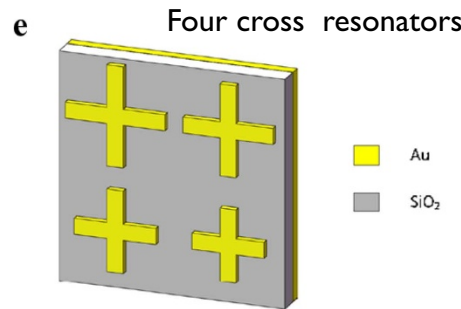
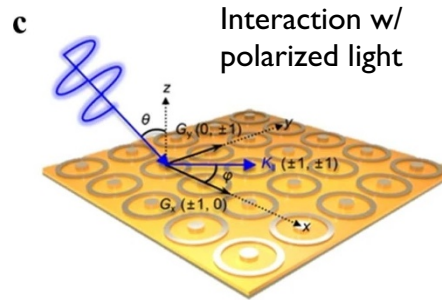
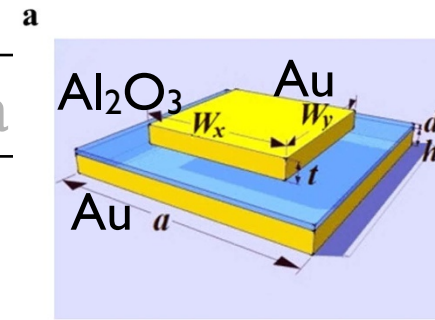
metamaterial : materials that obtain their properties from their *structure* rather than the material of which they are *composed*

These are *engineered composite materials* mainly consisting of artificially designed periodic sub-wavelength structures.

One particular application revolves around the absorption / reflectance characteristics in the IR in very compact devices.

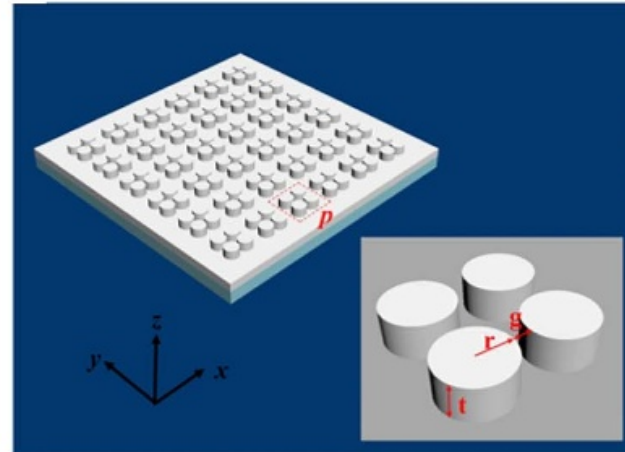
(relatively) *broadband* absorption:

Importantly: the properties can be designed!

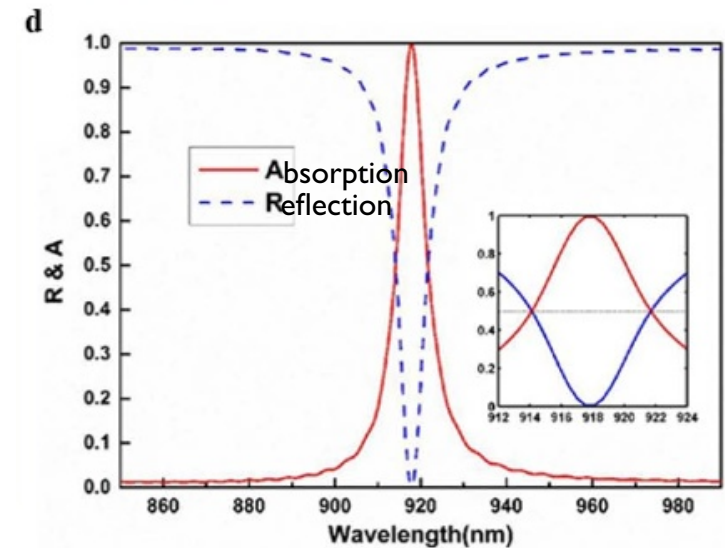


6 metamaterials, 0/1/2-D materials

Schematic of the MMA
 $r = 80\text{nm}$
 $t = 100\text{ nm}$
 $g = 20\text{ nm}$



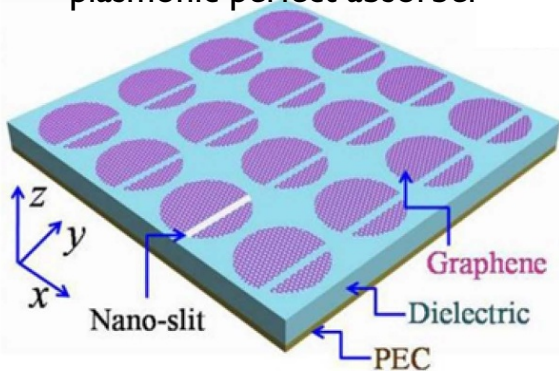
(relatively) *narrowband* absorption:



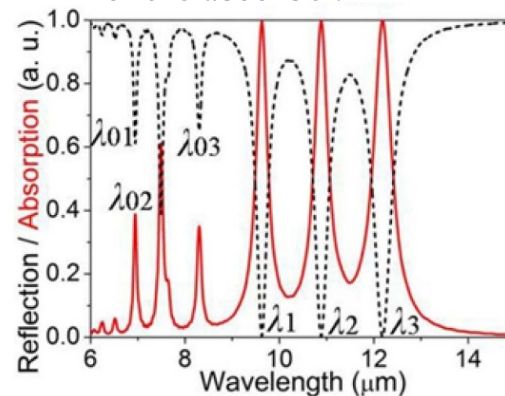
Y.Yao et al., J. Phys. D: Appl. Phys. 54 (2021) 113002 <https://doi.org/10.1088/1361-6463/abccf0>

2D materials based MMAs: strong tunability by tuning gate voltage

Schematic of the graphene plasmonic perfect absorber



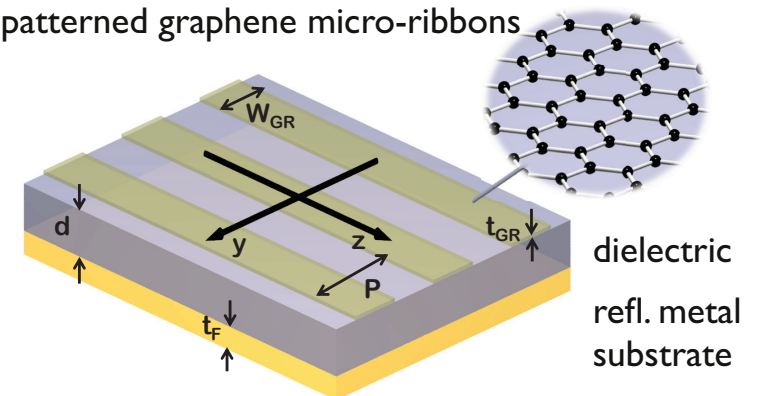
Absorption and reflection of the absorber.



Y.Yao et al., J. Phys. D: Appl. Phys. 54 (2021) 113002 <https://doi.org/10.1088/1361-6463/abccf0>

plasmonic metamaterials: perfect absorption of light

patterned graphene micro-ribbons



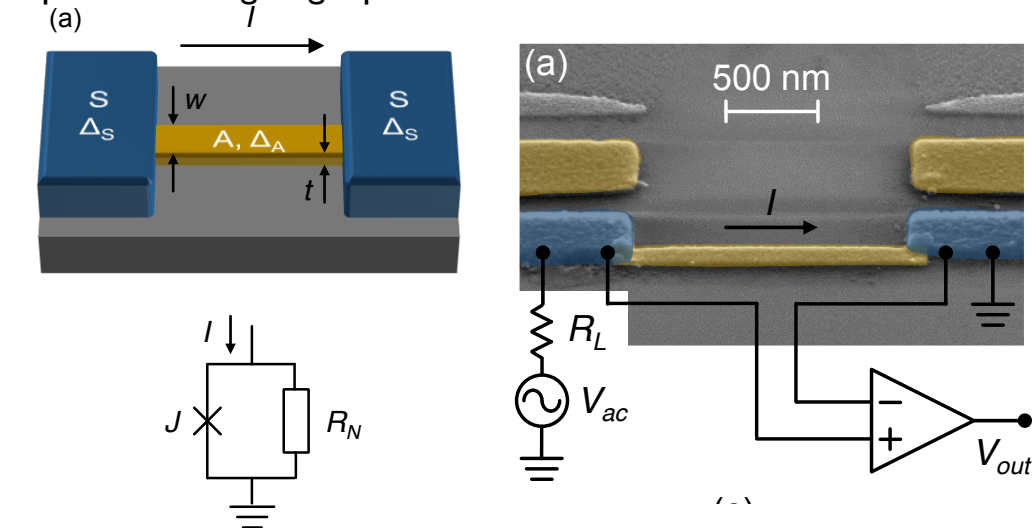
R.Alaei et al., Opt. Express. 20 (2012) 28017 <https://doi.org/10.1364/OE.20.028017>

6 metamaterials, 0/1/2-D materials

I-D junction: JES: Josephson escape sensor,
nanoscale transition edge sensor (nano-TES)

well suited to detecting dark matter / axion interactions with electrons

superconducting single-photon detection arXiv:2101.08558v2

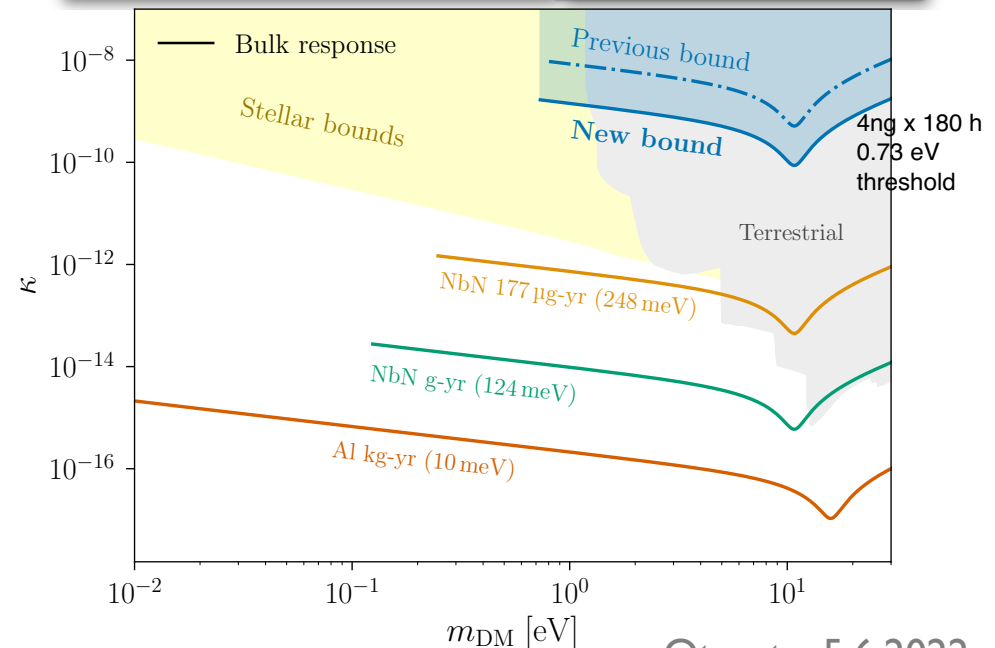
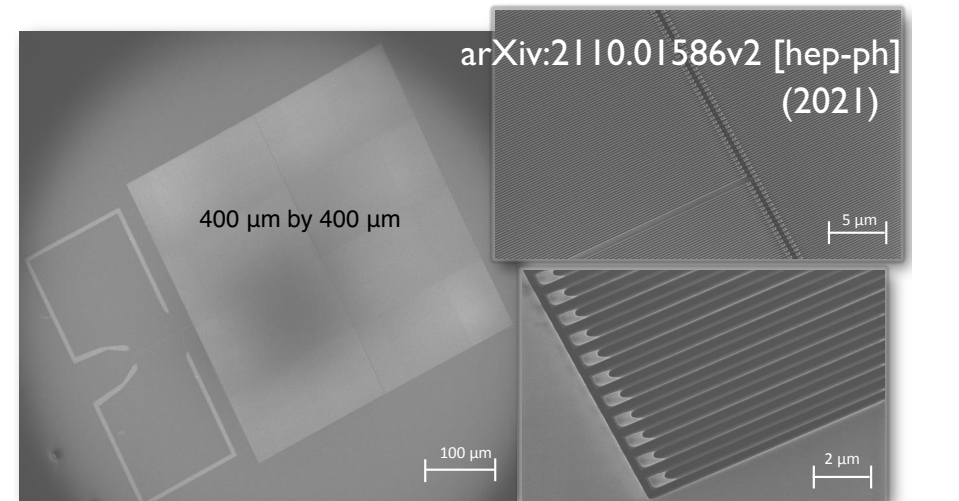


directional detection of light dark matter

dark matter scattering produces initial excitations with an
anisotropic distribution which is preserved during de-excitation

<https://arxiv.org/abs/2109.04473>

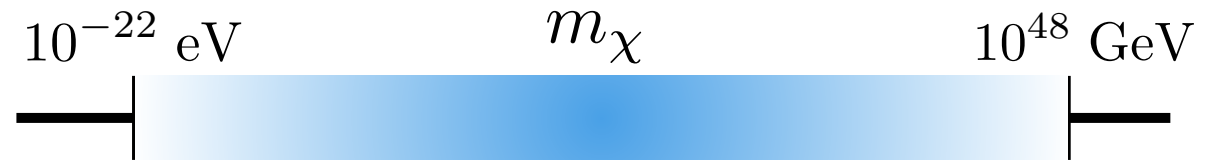
Superconducting nanowire single-photon detectors
(SNSPDs): sub-eV energy deposition thresholds



Quantum sensors for new particle physics experiments

quantum sensors & particle physics: what are we talking about?

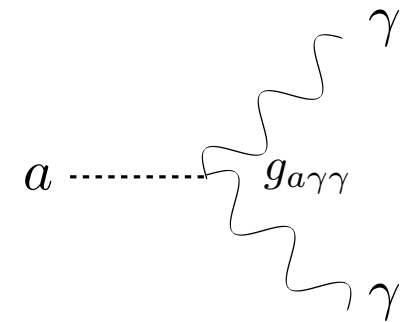
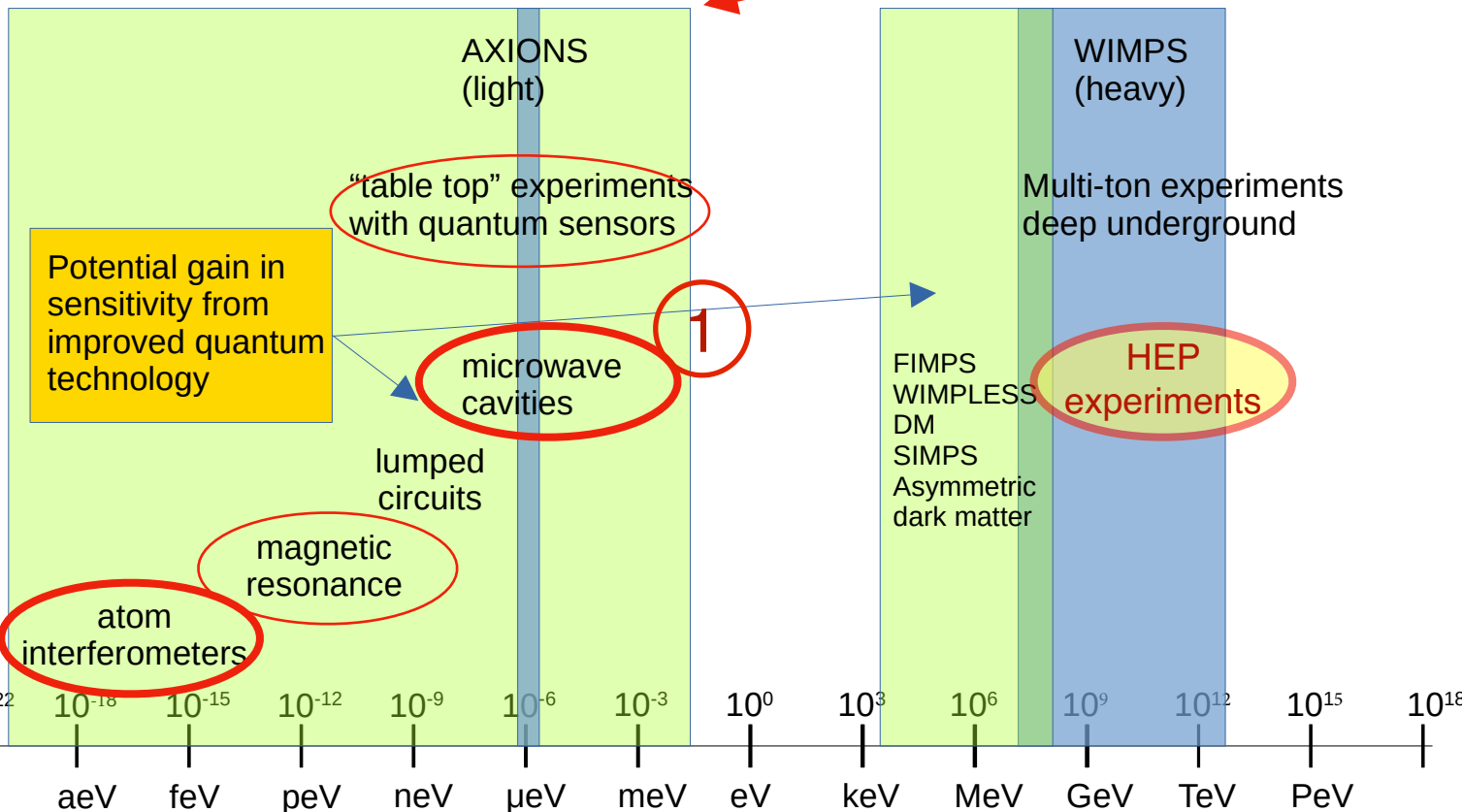
Axions, ALP's, DM & non-DM UL-particle searches



cavity size = axion size
axion mass = unknown

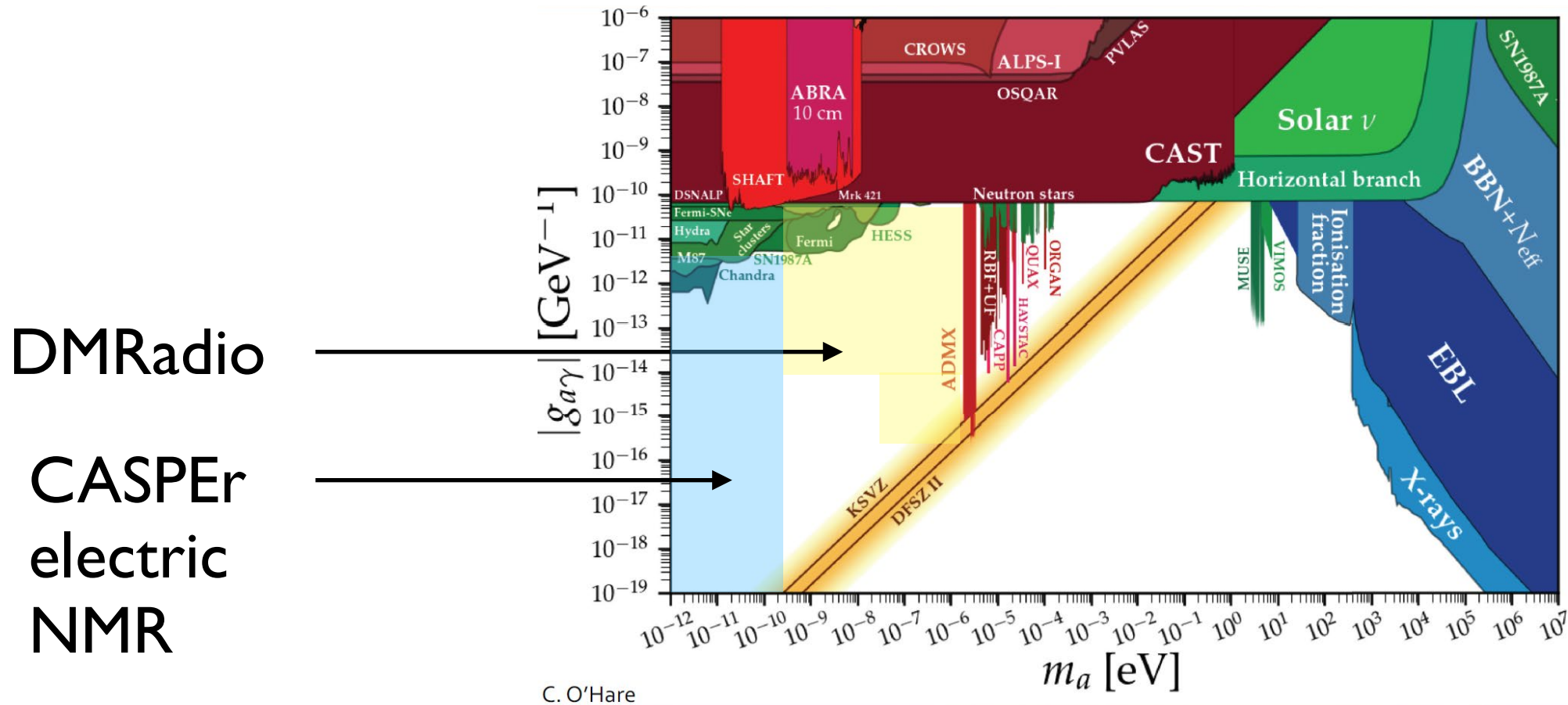
$$F \sim g_{a\gamma}^2 m_A^2 B^4 V^2 T_{sys}^{-2} G^4 Q,$$

system noise temperature
cryo-amplifiers JJPA



(but not only...)

particle physics: what are we talking about?

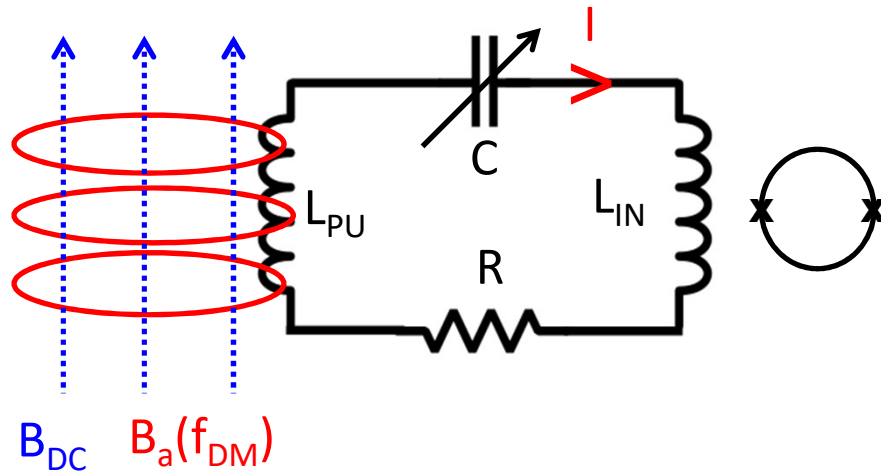


spin-based, NMR (2)	microwave cavities (1)	bolometers, TES (1)
electromagn. resonators		

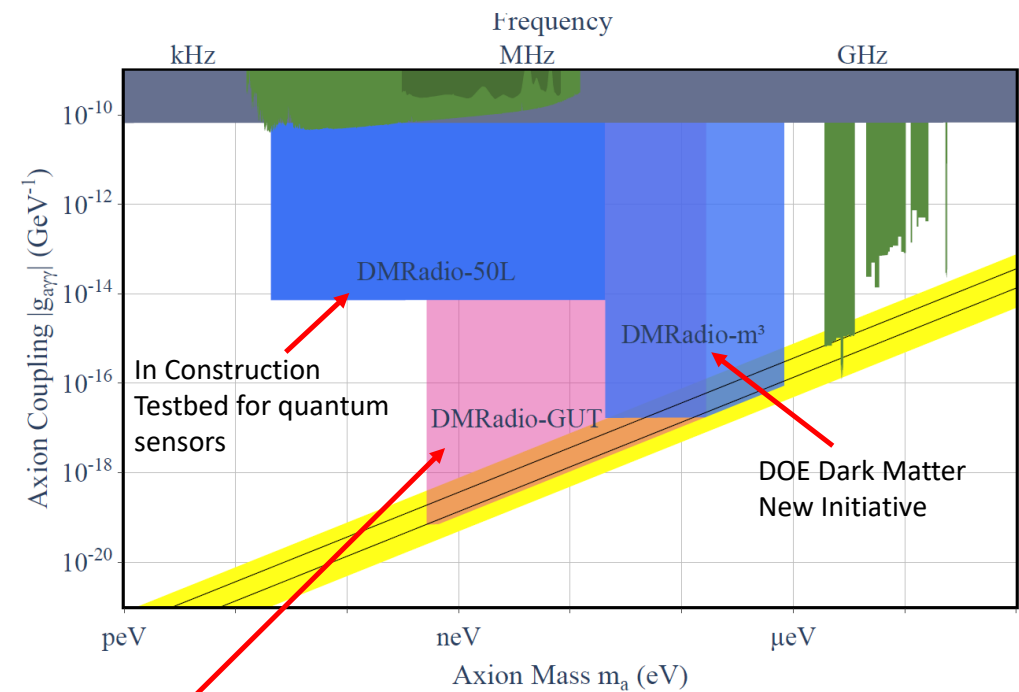
DMRadio

Focus on electromagnetic interaction: axions and photons mix in the presence of a strong magnetic field

Focus on detecting not a particle (a photon), but a field



- Axion field converts to oscillating EM signal in background DC magnetic field
- Detect using tunable resonator
- Signal enhancement when resonance frequency matches rest-mass frequency $\nu_{DM}=mc^2/h$
- SQUID's, RF Quantum upconverters, cryoamplifiers (e.g. JJPA)



CASPER electric NMR

Focus on different interactions: the **electric dipole moment (EDM) interaction** and the **gradient interaction** with **nuclear spin** I . The EDM interaction arises from the coupling of the axion to the gluon field.

→ spin σ to axion coupling:

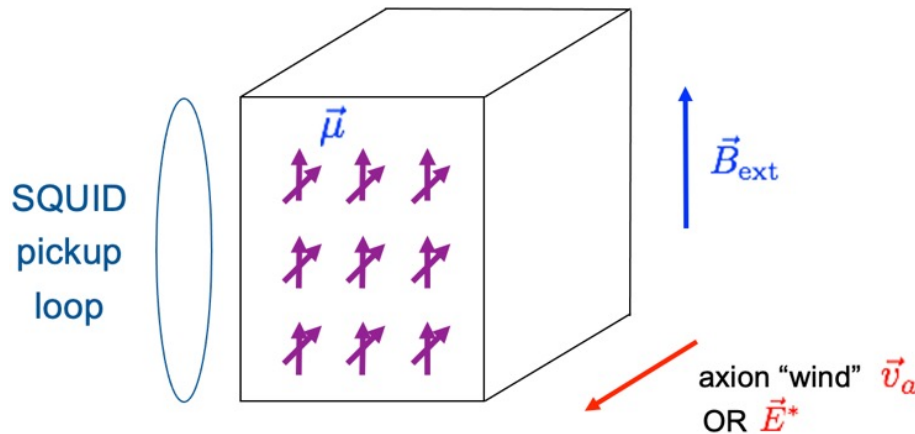
$$H_e \propto a \sigma \cdot E^*$$

CASPER-electric

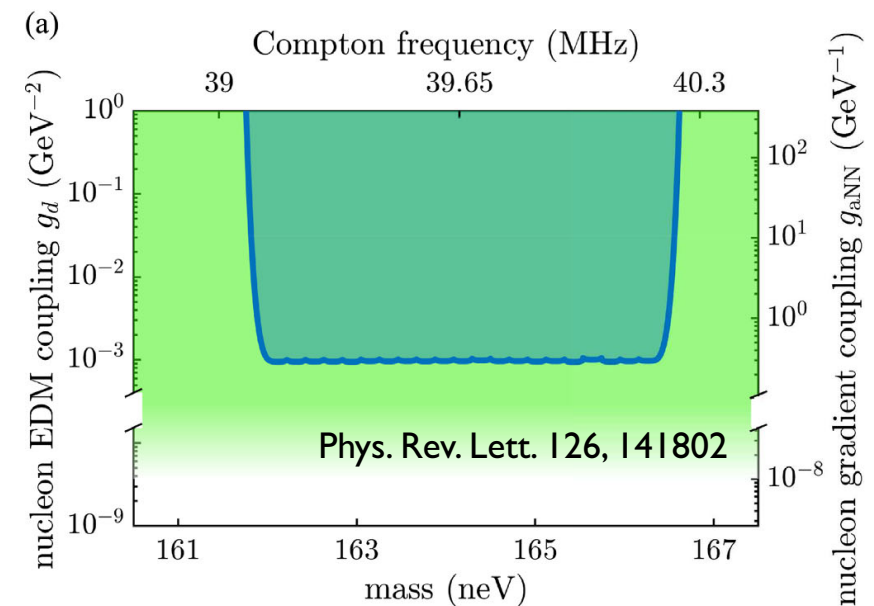
→ spin σ to axion gradient coupling:

$$H_g \propto \sigma \cdot \nabla a$$

CASPER-gradient



Cosmic Axion Spin Precession Experiment is based on a precision measurement of ^{207}Pb solid-state nuclear magnetic resonance in a polarized ferroelectric crystal. Axion-like dark matter can exert an oscillating torque on ^{207}Pb nuclear spins via the electric dipole moment coupling g_d or via the gradient coupling g_{aNN} .



numerous improvements possible → many orders of magnitude in mass and sensitivity range

Initial experiments with quantum sensors world-wide

→ rapid investigation of new phase space

→ scaling up to larger systems, improved devices
→ expanding explored phase space

→ particles, atoms, ions, nuclei: tests of QED, symmetries

→ atomic interferometers: DM searches

→ RF cavities: axion searches

@ CERN: PBC, large low energy physics community...

<https://indico.cern.ch/event/1002356/> PBC technology annual workshop 2021 (focus on quantum sensing)
<https://indico.cern.ch/event/1057715/> PBC technology mini workshop: superconducting RF (Sep. 2021)

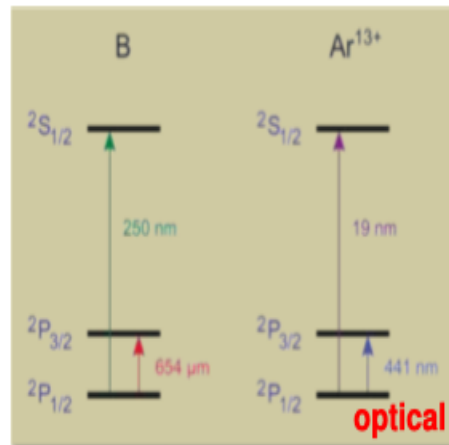
particle physics @ CERN : what are we talking about?

tests of QED, T-violation, P, Lorentz-violation, DM searches

HCI's in Penning traps

Scaling with a nuclear charge Z

Binding energy $\sim Z^2$
 Hyperfine splitting $\sim Z^3$
 QED effects $\sim Z^4$
 Stark shifts $\sim Z^{-6}$

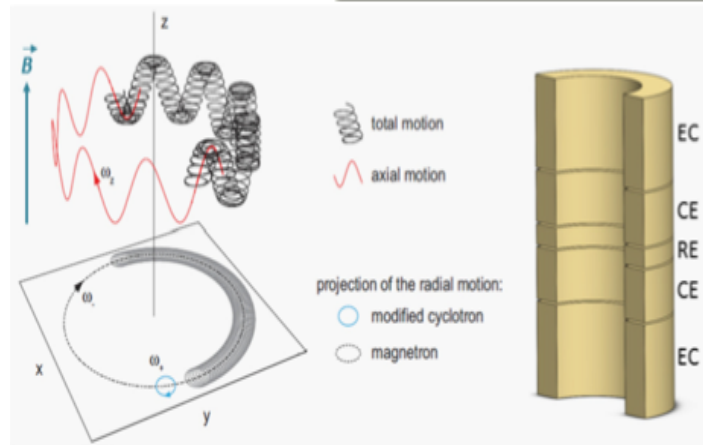


@ CERN:

eEDM's in molecules

nuclear clock (^{229}Th)

molecular / ion clocks



K. Blaum et al., Quantum Sci. Technol. 6 014002 (2021)

Quantum Sensors for New-Physics Discoveries

<https://iopscience.iop.org/journal/2058-9565/page/Focus-on-Quantum-Sensors-for-New-Physics-Discoveries>

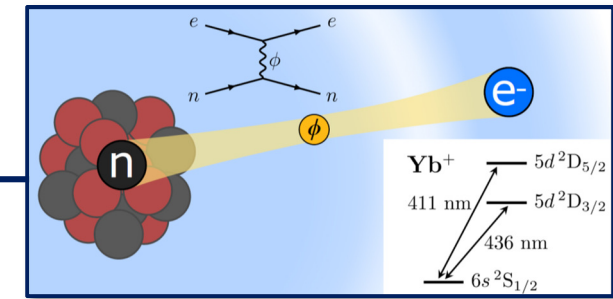
Present status: 5×10^{-12} relative mass precision

R. X. Schüssler et al., Nature 581, 42 (2020)

HCLs: **much larger** sensitivity to variation of α and dark matter searches than current clocks

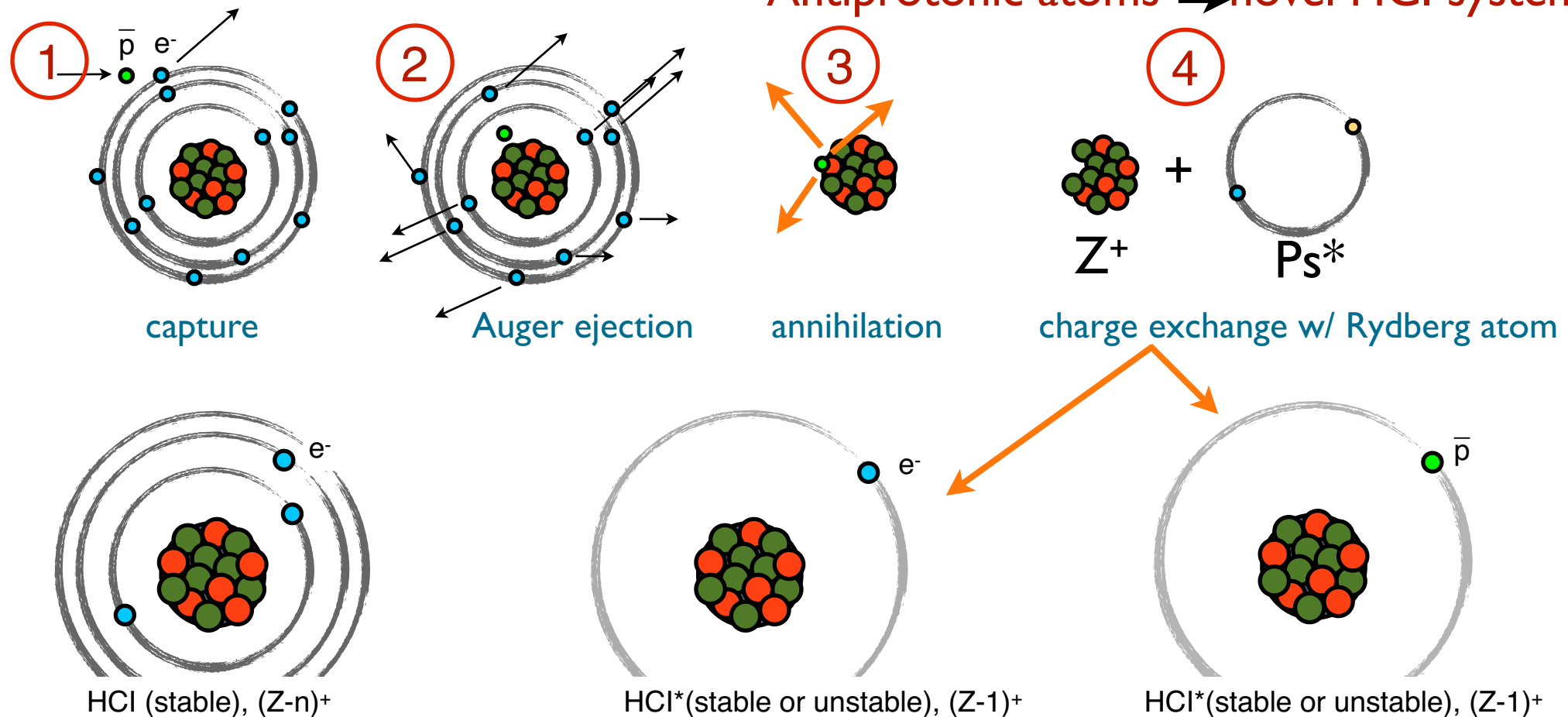
- Searches for the variation of fundamental constants
- Tests of QED: precision spectroscopy
- Fifth force searches: precision measurements of isotope shifts with HCLs to study non-linearity of the King plot

linear relation in isotope shifts between two transitions caused by e.g. new force mediated by weakly interacting boson
Mikami, K. Et al.. Probing new intra-atomic force with isotope shifts. The European Physical Journal C. 77. 10.1140/epjc/s10052-017-5467-4.

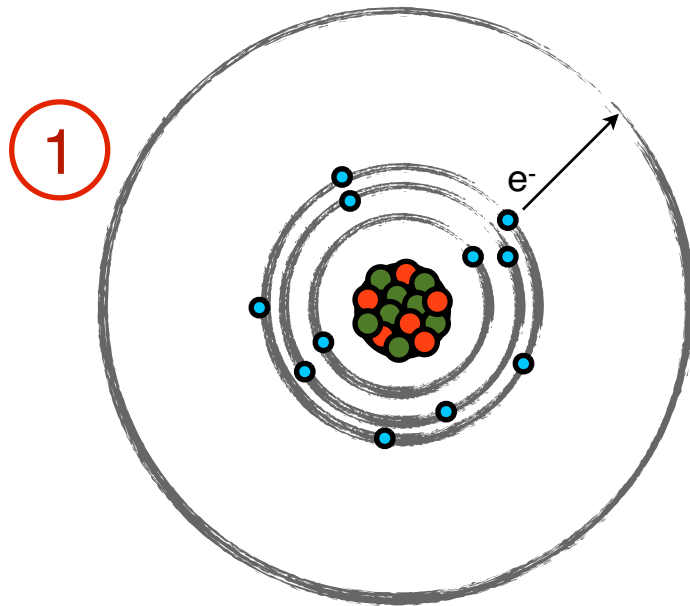


Review on HCLs for optical clocks: Kozlov *et al.*, Rev. Mod. Phys. **90**, 045005 (2018)

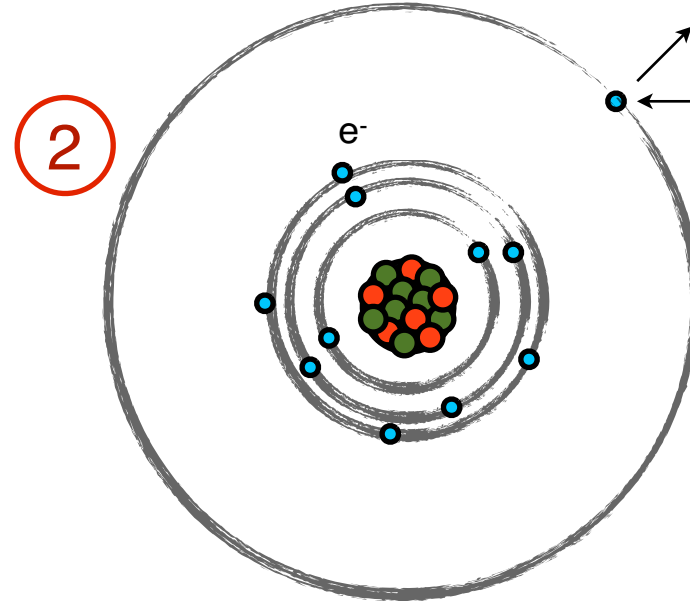
Antiprotonic atoms \rightarrow novel HCL systems



Antiprotonic Rydberg atom: exotic couplings, similar approach as spectroscopy of muonic atoms, CPT tests



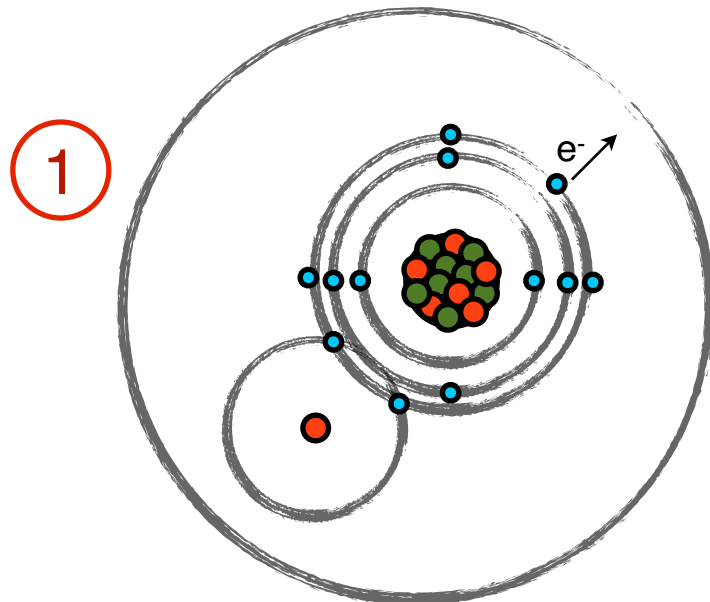
Rydberg excitation



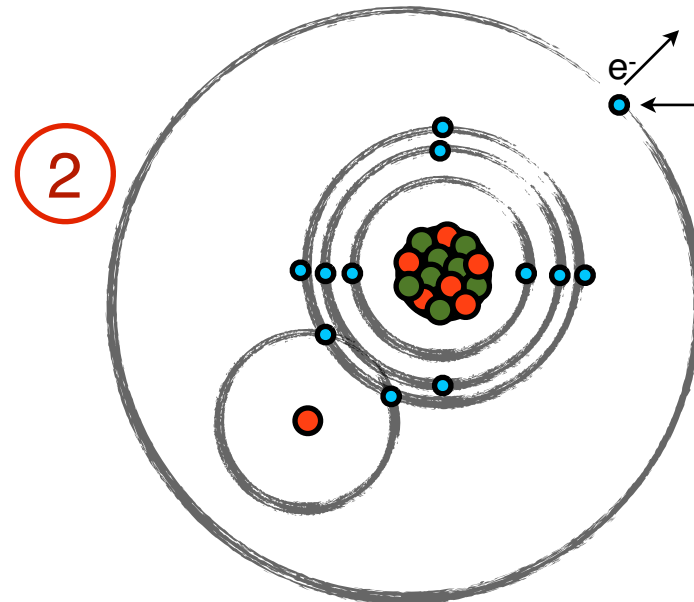
charge exchange

at end of cascade, p is very close to nucleus... investigate long-range behavior of strong interaction?

Antiprotonic Rydberg molecule: \bar{p} EDM?



Rydberg excitation



charge exchange

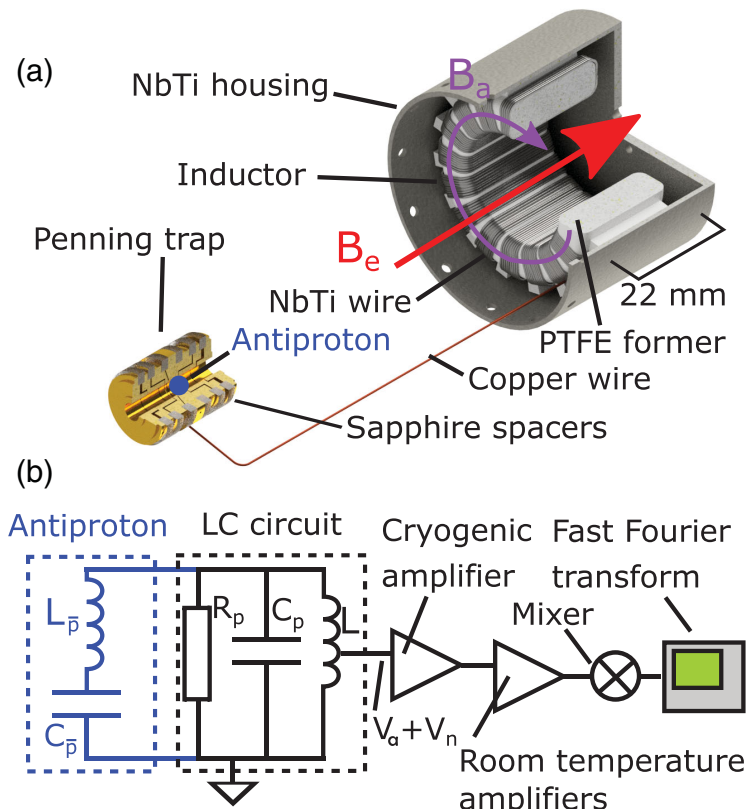
similar approach as eEDM in molecules

search the noise spectrum of fixed-frequency resonant circuit for peaks caused by dark matter ALPs converting into photons in the strong magnetic field of the Penning-trap magnet

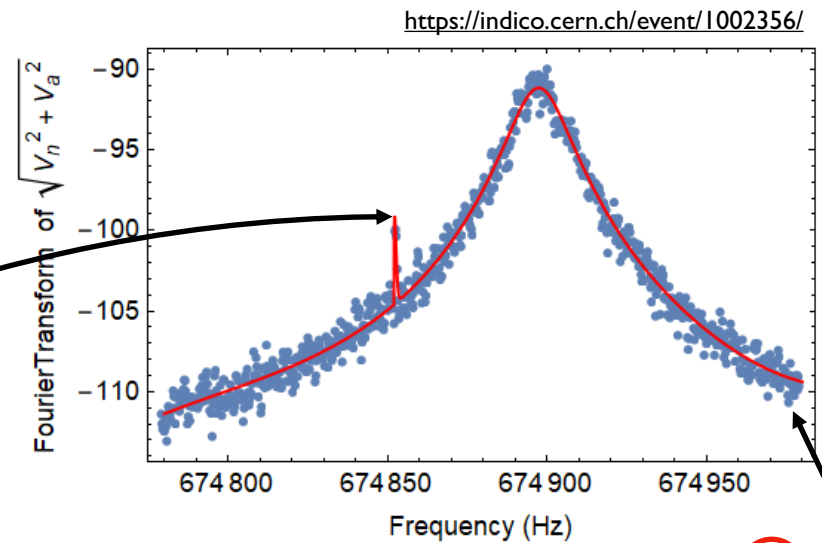
Resolving single antiproton spin flips requires the highest Q and lowest temperature LC resonant detectors ever built: **BASE**-CERN is the state of the art

Constraints on the Coupling between Axion-like Dark Matter and Photons Using an Antiproton Superconducting Tuned Detection Circuit in a Cryogenic Penning Trap

J. Devlin et al., **BASE** collaboration, Physical Review Letters 126, 041301 (2021)



H. Nagahama et al., Rev. Sci. Instrum. 87, 113305 (2016)



resonator background $\propto \sqrt{T_z}$
from antiproton spin-flip

The axion signal

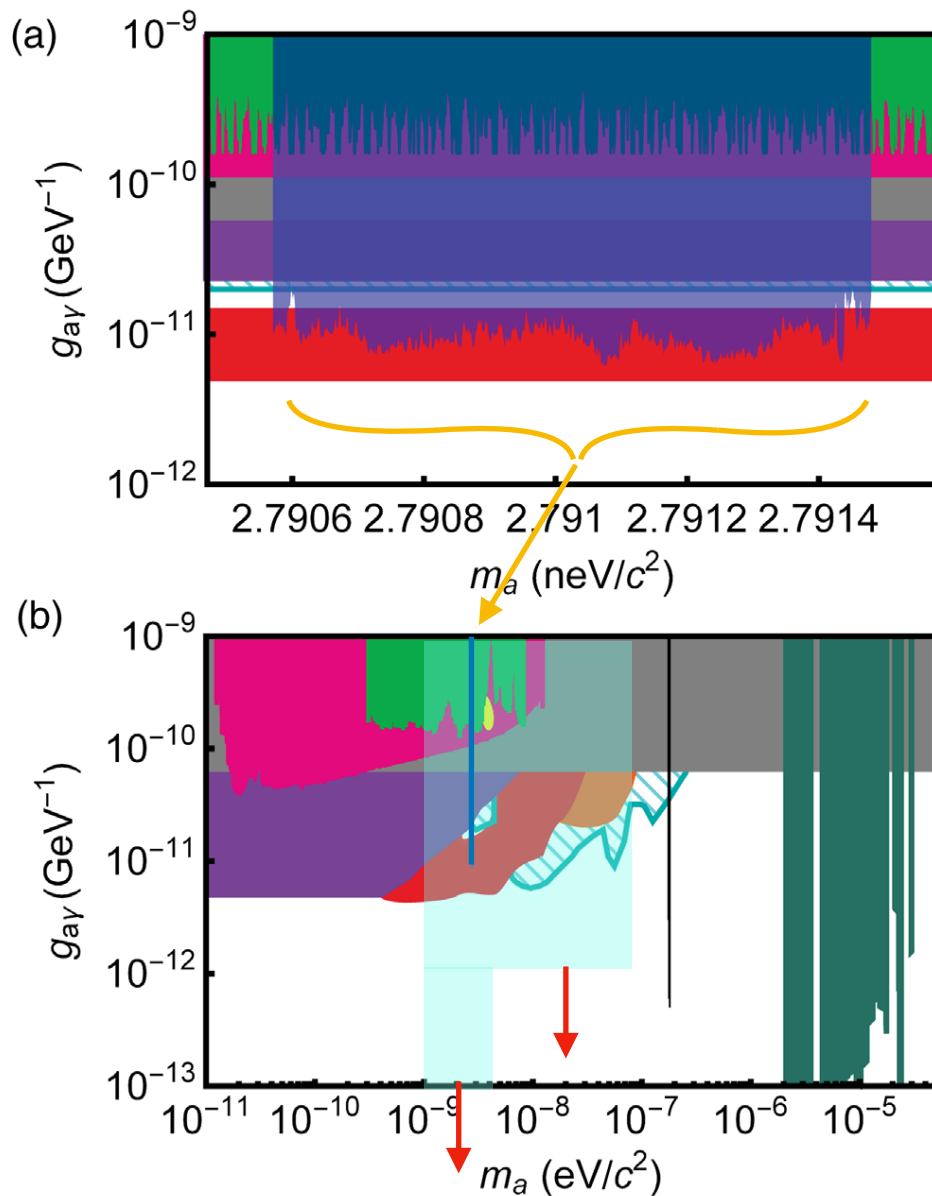
$$V_a = \frac{\pi}{2} Q \sqrt{f(\nu, Q, \mathbf{q})} \kappa \nu_a l N_T (r_2^2 - r_1^2) g_{a\gamma} |\mathbf{B}_e| \sqrt{\rho_a \hbar c}.$$

$f(\nu, Q, \mathbf{q})$ is a lorentzian line-shape function proportional to $\text{Re}[Z]$
 e_n is the equivalent input noise of the amplifier
 κ is the coupling constant
 Q is the resonator Q-factor
 N_T is the number of turns
 l is the length of the toroid along the magnet B field

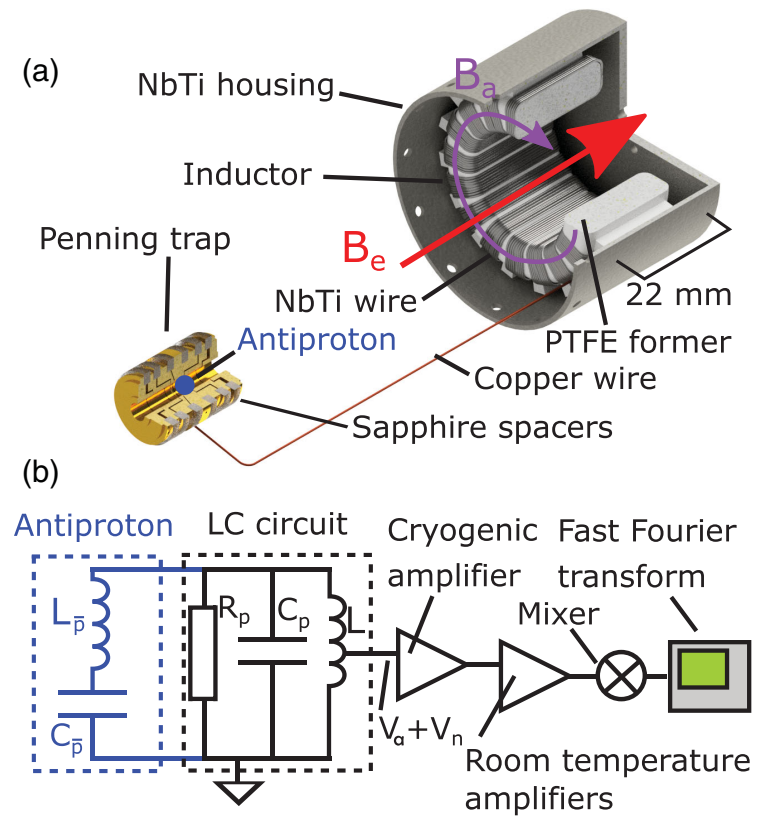
r_1 is the inner radius of the toroid
 r_2 is the outer radius
 $g_{a\gamma}$ is the coupling constant
 B is the static magnetic field
 ρ_a is the dark matter density

Tunability!

Quantum sensors for new particle physics experiments: Penning traps



Limits		Hints
SN-1987A	CAST	ADMX-SLIC
H.E.S.S.	BASE	ABRACADABRA
Cavities	SHAFT	FERMI-LAT
		Excess
		γ rays
		Pulsars



currently developing **superconducting tunable capacitors**
& **laser-cooled resonators**

7 T magnet + broader FFT span: one month \longrightarrow
2 and 5 neV to an upper limit of $1.5 \times 10^{-11} \text{ GeV}^{-1}$

Axion heterodyne detection

problem: cavity resonance generally fixed

Conceptual Theory Level Proposal:

A. Berlin, Raffaele Tito D'Agnolo, S. Ellis, C. Nantista, J. Nielson, P. Schuster, S. Tantawi, N. Toro, K. Zhou, *JHEP* 07 (2020) 07, 088

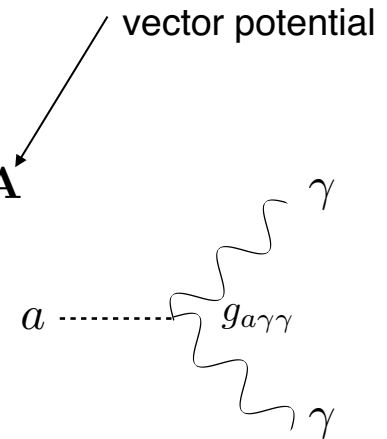
A. Berlin, Raffaele Tito D'Agnolo, S. Ellis, K. Zhou, arXiv:2007.15656

Axion DM coupling to electromagnetism through

$$-\mathcal{L} \supset \frac{1}{4} g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu} \supset \frac{1}{2} \mathbf{J}_{\text{eff}} \cdot \mathbf{A}$$

In the presence of \mathbf{B} , axion \rightarrow effective current density

$$\mathbf{J}_{\text{eff}} \simeq g_{a\gamma\gamma} \partial_t a \mathbf{B}.$$



Static $\mathbf{B} \rightarrow \mathbf{J}_{\text{eff}}$ oscillates with the same frequency as the axion field

Resonant cavities possible down to μeV ; below that, need huge volume

\rightarrow frequency conversion: driving “**pump mode**” at $\omega_0 \sim \text{GHz}$ allows axion to resonantly drive power into “**signal mode**” at $\omega_1 \sim \omega_0 \pm m_a$

\rightarrow scan over axion masses $m_a =$ **slight perturbation of cavity geometry**, which modulates the frequency splitting $\omega_0 - \omega_1$

\rightarrow **superconducting RF cavities**

Tunability!

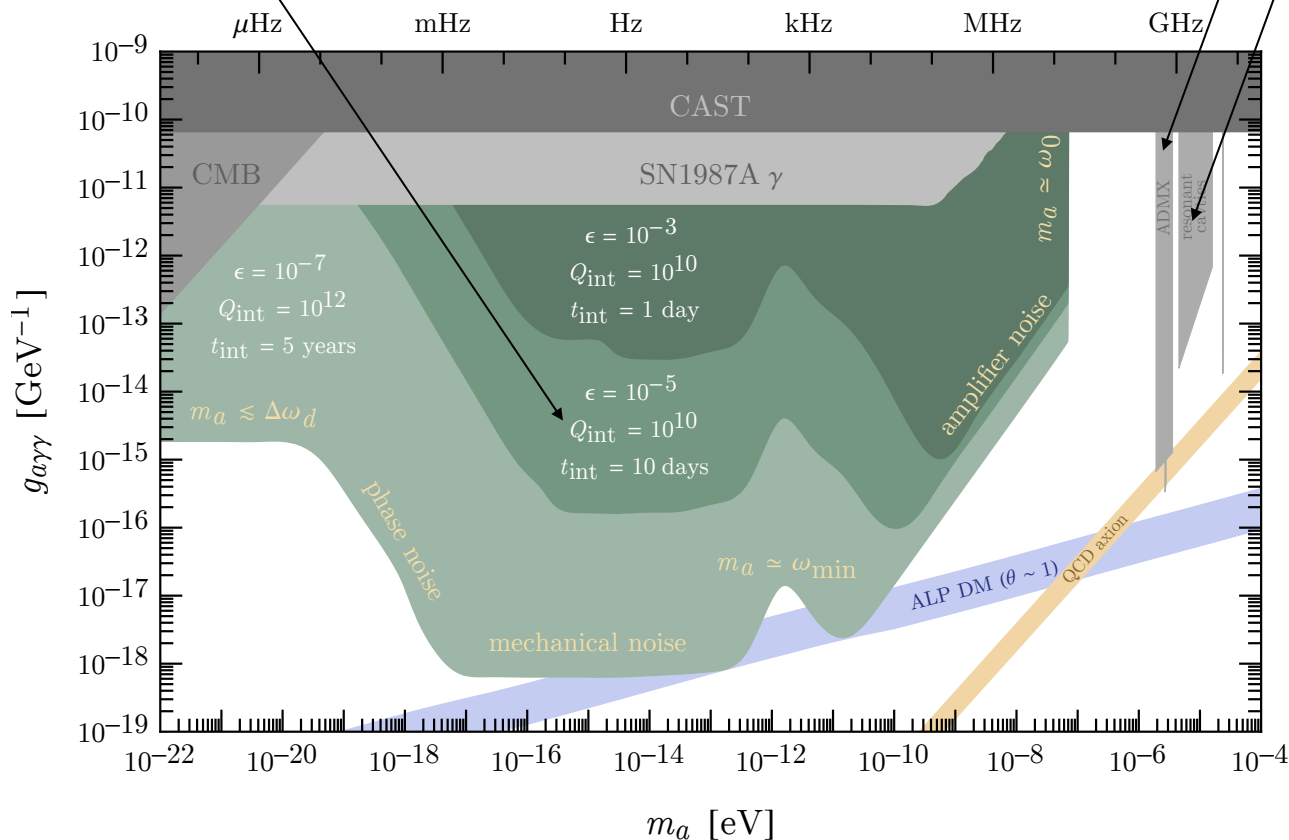
Quantum sensors for new particle physics experiments: tunable RF cavities

$Q_{\text{int}} \gtrsim 10^{10}$ achieved by DarkSRF collaboration
(sub-nm cavity wall displacements)

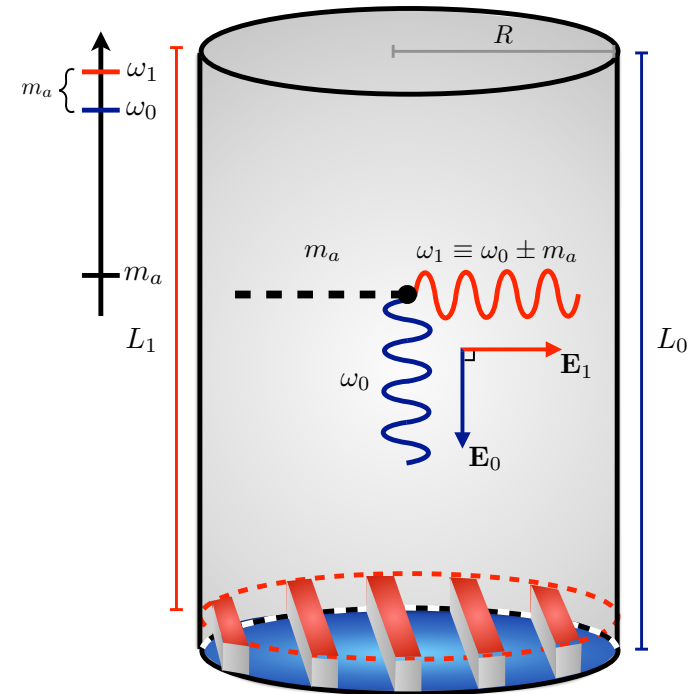
A. Grassellino, "SRF-based dark matter search: Experiment," 2019. <https://indico.fnal.gov/event/19433/session/2/contribution/2/material/slides/0.pdf>

$$\text{frequency} = m_a/2\pi$$

ADMX experiment
resonant cavities



A. Berlin, Raffaele Tito D'Agnolo, S. Ellis, C. Nantista, J. Nielson, P. Schuster, S. Tantawi, N. Toro, K. Zhou, *JHEP* 07 (2020) 07, 088



(a) Cartoon of cavity setup.

Asher Berlin, Raffaele Tito D'Agnolo, Sebastian A. R. Ellis, Christopher Nantista, Jeffrey Neilson, Philip Schuster, Sami Tantawi, Natalia Toro, Kevin Zhou, <https://arxiv.org/abs/1912.11048>

"The cavity is designed to have two nearly degenerate resonant modes at ω_0 and $\omega_1 = \omega_0 + m_a$. One possibility is to split the frequencies of the two polarizations of a hybrid HE_{11p} mode in a corrugated cylindrical cavity. These two polarizations effectively see distinct cavity lengths, L_0 and L_1 , allowing ω_0 and ω_1 to be tuned independently."

Axion searches with resonant haloscopes: resonant cavity immersed in a high and static magnetic field

Relic Axion Detector Exploratory Setup (RADES) searches for axion dark matter with $m_a > 30 \mu\text{eV}$

$$F \sim g_{a\gamma}^2 m_A^2 B^4 V^2 T_{sys}^{-2} G^4 Q,$$

figure of merit magnetic field quality factor

Cavity coatings: type II superconductor with a critical magnetic field B_c well above 11 T at 4.2 K

Thin Film (High Temperature) Superconducting Radiofrequency Cavities for the Search of Axion Dark Matter

S. Golm, ..., Sergio Calatroni, ... et al. <https://ieeexplore.ieee.org/document/9699394> DOI: 10.1109/TASC.2022.3147741

—————> developments of HTS for coatings is essential in improving the sensitivity of resonant haloscopes

Multiple cavities: optimal coupling with external B field, very selective (high Q), centered on resonant ν

Universe **2022**, 8(1), 5; <https://doi.org/10.3390/universe8010005>

other frequencies: e.g. solenoidal magnet in dilution cryostat at 10 mK (Canfranc Underground Lab.)

—————> to exploit the ultra-low temperatures and go beyond the standard quantum limit:
Josephson parametric amplifiers (JPA), superconducting qubit-based single photon counters,
(or for higher frequencies, kinetic inductor devices (KID))

AION: atom interferometer (start small, ultimately → space)

L. Badurina et al., AION: An Atom Interferometer Observatory and Network, *JCAP* 05 (2020) 011, [arXiv:1911.11755].

Topological Dark Matter (TDM)

TDM can be expressed as a scalar field that couples to fundamental constants, thus producing variations in the transition frequencies of atomic clocks at its passage.

Ultralight Dark Matter

spatial variation of the fundamental constants associated with a change in the gravitational potential

Local Lorentz Invariance (LLI)

independence of any local test experiment from the velocity of the freely-falling apparatus.

searching for daily variations of the relative frequency difference of e.g. Sr optical lattice clocks or Yb⁺ clocks confined in two traps with quantization axis aligned along non-parallel directions

Local Position Invariance (LPI)

independence of any local test experiment from when and where it is performed in the Universe

comparing atomic clocks based on different transitions can be used to constrain the time variation of fundamental constants and their couplings, comparison of two ¹⁷¹Yb⁺ clocks and two Cs clocks → limits on the time variation of the fine structure constant and of the electron-to-proton mass ratio

Gravitational wave detector

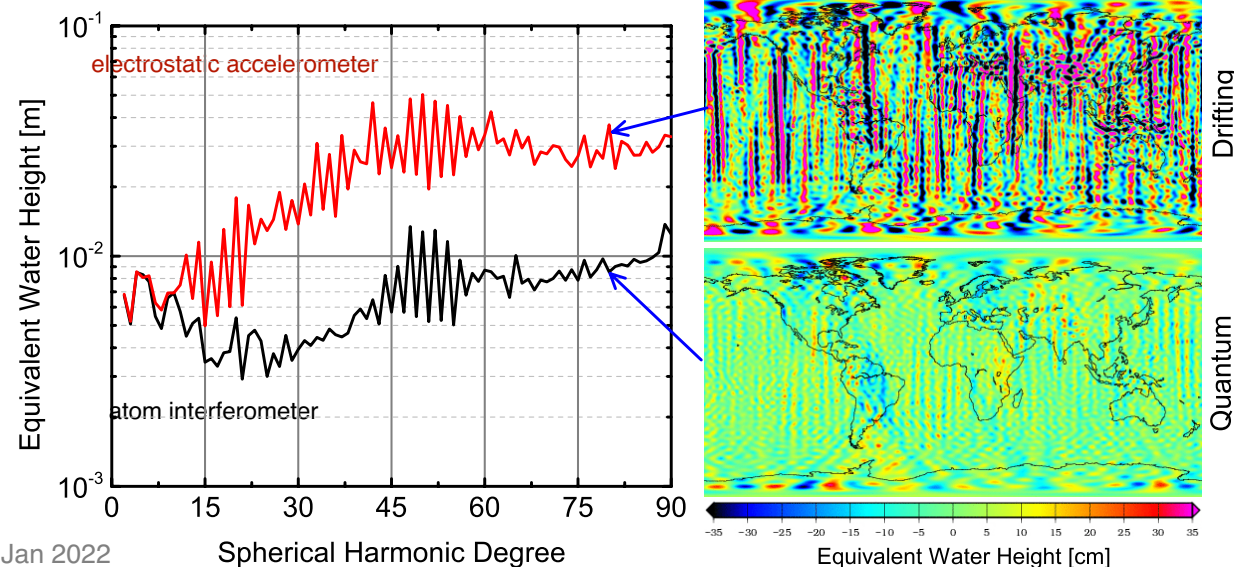
clocks act as narrowband detectors of the Doppler shift on the laser frequency due to the relative velocity between the satellites induced by the incoming gravitational wave

R & D needed:

Optical lattice clocks at up to 1×10^{-18} relative accuracy

& expanded optical fibre network (operated between a number of European metrology institutes)

& develop cold atom technology for robust, long-term operation



Ultralight Dark Matter

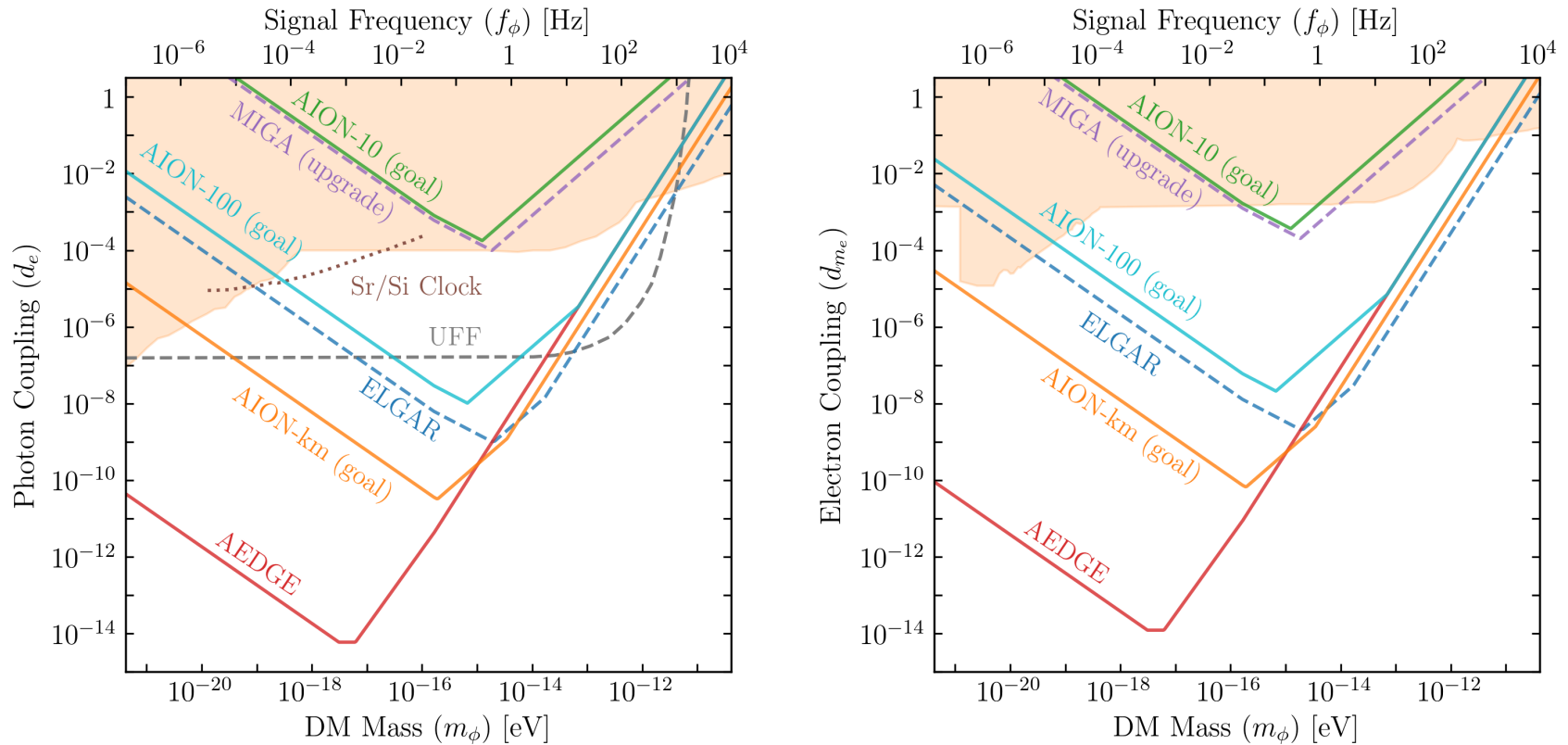


Figure 10. Sensitivities to the possible couplings of scalar ULDM to photons (left panel) and to electrons (right panel) of the terrestrial AION [181], MIGA [182] and ELGAR [180] experiments, and of the AEDGE space-borne concept [5]. Also shown is a combination of the current constraints from MICROSCOPE [189] and terrestrial torsion balance and atomic clock experiments and (in the left panel) the sensitivity of a prospective measurement of the universality of free fall (UFF) with a precision of 10^{-17} [6].

arXiv:2201.07789v1 [astro-ph.IM] 19 Jan 2022

Where does this fit in? Go after $10^{-20} \text{ eV} < m_a < 10^{-12} \text{ eV}$

atom interferometry at macroscopic scales:

arXiv:2201.07789v1 [astro-ph.IM] 19 Jan 2022

MIGA^{France}

AION^{UK}

ZAIGA^{China}



shafts (100~500 m ideal testing ground),
cryogenics, vacuum, complexity...

MAGIS^{Fermilab}

M. Abe, P. Adamson, M. Borcean, D. Bortoletto, K. Bridges, S.
P. Carman et al., *Matter-wave Atomic Gradiometer
Interferometric Sensor (MAGIS-100)*, [arXiv:2104.02835v1](https://arxiv.org/abs/2104.02835v1).

MAGIS collaboration, Graham PW, Hogan JM, Kasevich MA,
Rajendran S, Romani RW. *Mid-band gravitational wave
detection with precision atomic sensors*. [arXiv:1711.02225](https://arxiv.org/abs/1711.02225)

satellite missions:

ACES (Atomic Clock Ensemble in Space): 2024-2025

ESA mission for ISS

two on-board clocks rely on atomic transitions in the microwave domain

probe time variations of fundamental constants, and to perform tests of the Lorentz-Violating Standard Model Extension (SME). Possibly topological dark matter

pathfinder / technology development missions:

I-SOC: key optical clock technology (laser cooling, trapping, optical resonators) for space; Sr optical lattice clock / Sr ion clock;
microwave and optical link technology;

FOCOS (Fundamental physics with an Optical Clock Orbiting in Space): Yb optical lattice clock with 1×10^{-18} stability

~2030

AEDGE: ~2045

satellite mission

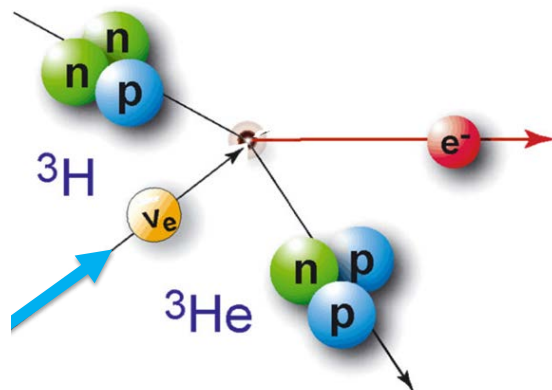
El-Neaj, Y.A., Alpigiani, C., Amairi-Pyka, S. *et al.* *AEDGE: Atomic
Experiment for Dark Matter and Gravity Exploration in Space*. *EPJ Quantum
Technol.* **7**, 6 (2020). <https://doi.org/10.1140/epjqt/s40507-020-0080-0>

Challenge: Detection of primordial neutrino flux: $E_\nu \sim \mu\text{eV}$

neutrino capture on β -decay nuclei as a detection method for the Cosmic Neutrino Background (CNB) was laid out in the original paper by Steven Weinberg in 1962

S. Weinberg, Universal Neutrino Degeneracy, Phys. Rev. 128, 1457 (1962)

detection!
directionality!
energy!

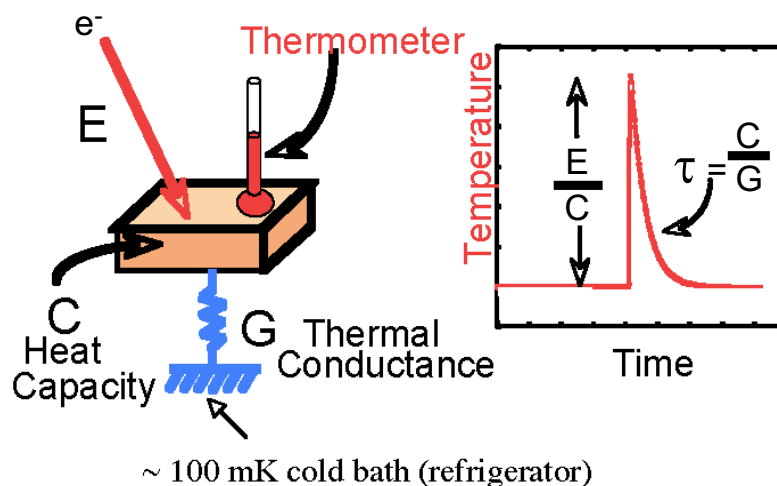
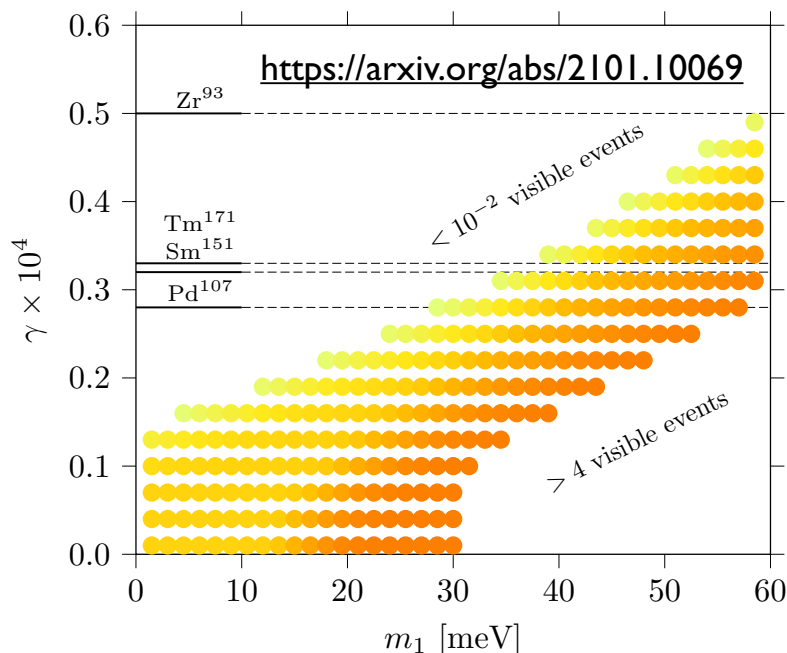


PTOLEMY experimental quest

relic neutrino capture
RF tracker (EM radiation in B field)
microcalorimeter (TES: sub-eV for $\sim\text{eV}$)

<https://indico.cern.ch/event/1090250/>

event rate low!
Better use
heavier
isotopes than
tritium as
targets!



MUCH R&D needed!

Quantum detectors for high energy particle physics

typically not obvious, given that most detectors rely on detecting the product of many interactions between a particle and the detector (ionization, scintillation, Cerenkov photons, ...)

handful of ideas that rely on quantum devices, or are inspired by them, but do not necessarily use them as quantum detectors per se, but rather their properties to enhance / permit measurements that are more difficult to achieve otherwise

these are not developed concepts, but rather the kind of approaches one might contemplate working towards



very speculative!

main focus on tracking / calorimetry / timing

→ closely related: nanostructured materials
Frontiers of Physics, M. Doser et al., 2022

what are the challenges?

- tracking: hit positions, material budget, vertexing
- timing: TOF \sim sub-ns (ideally ps) for PID
- calorimetry: shower shape, timing, granularity, particle flow
- redundancy / independent modes of measurement
- novel observables: helicity / polarization
- power budget, event rate, PU (timing)

Metamaterials, 0 / 1 / 2-dimensional materials (quantum dots, nanolayers)

ultra-fast scintillators based on perovskites

chromatic calorimetry (QDs)

active scintillators (QCL, QWs, QDs)

5.3.6 *

GEMs (graphene)

Atoms, molecules, ions

Rydberg TPC's

5.3.5 *

Spin-based sensors

helicity detectors

5.3.3 *

* <https://cds.cern.ch/record/2784893>

Metamaterials, 0 / 1 / 2-dimensional materials (quantum dots, nanolayers)

ultra-fast scintillators based on perovskites

chromatic calorimetry (QDs)

active scintillators (QCL, QWs, QDs)

5.3.6

GEMs (graphene)

Atoms, molecules, ions

Rydberg TPC's

5.3.5

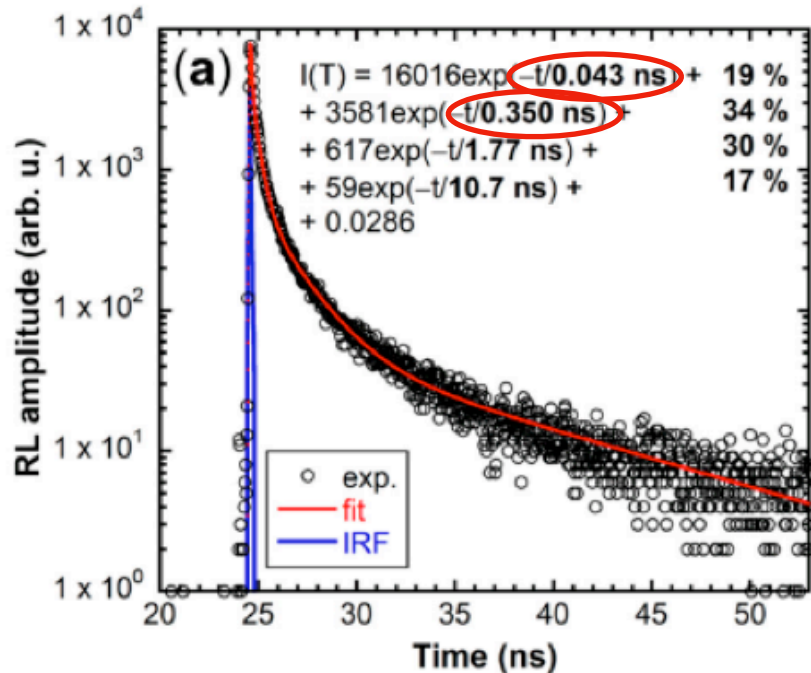
Spin-based sensors

helicity detectors

5.3.3

Quantum dots: timing

Etiennette Auffray-Hillemans / CERN



Scintillation decay time spectra from CsPbBr₃ nanocrystal deposited on glass

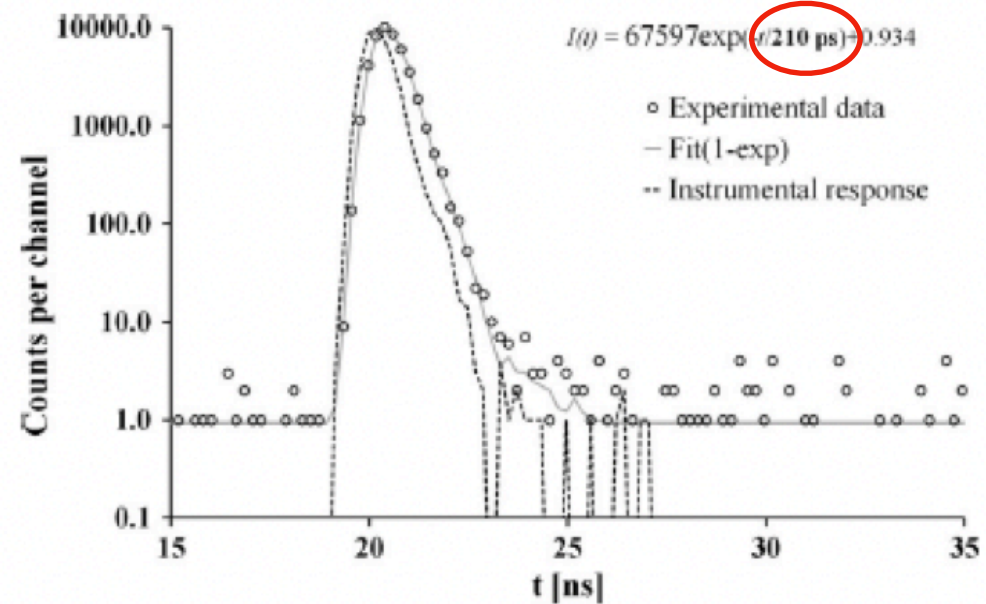


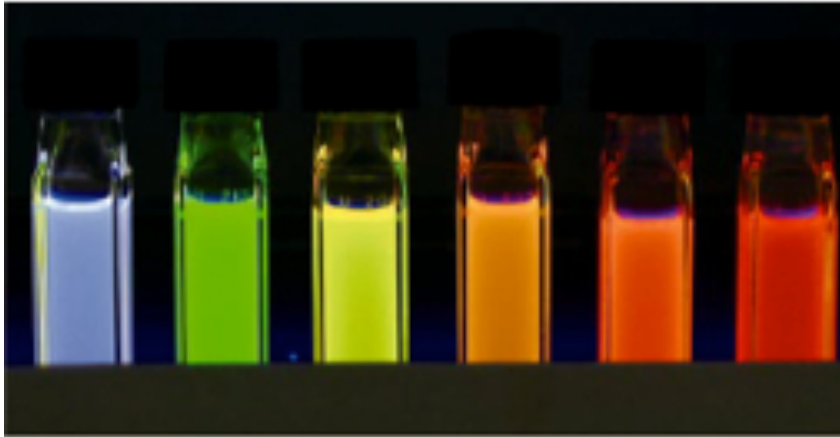
Fig. 9. Photoluminescence decay of ZnO:Ga sample at room temperature. Excitation nanoLED 339 nm, emission wavelength set at 390 nm. Decay curve is approximated by the convolution of instrumental response (also in figure) and single exponential function $I(t)$ provided in the figure.

Lenka Prochazkova et al., Optical Materials 47 (2015) 67–71

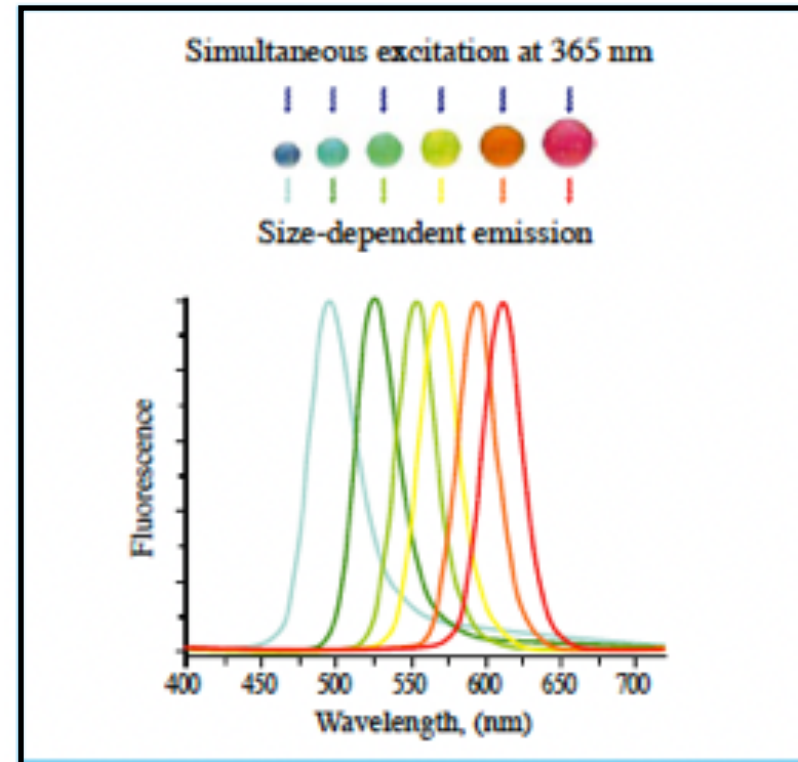
Concerns: integrated light yield (need many photons to benefit from rapid rise time)

Quantum dots: timing

Etiennette Auffray-Hillemans / CERN



Hideki Ooba, "Synthesis of Unique High Quality Fluorescence Quantum Dots for the Biochemical Measurements," AIST TODAY Vol.6 , No.6 (2006) p.26- 27

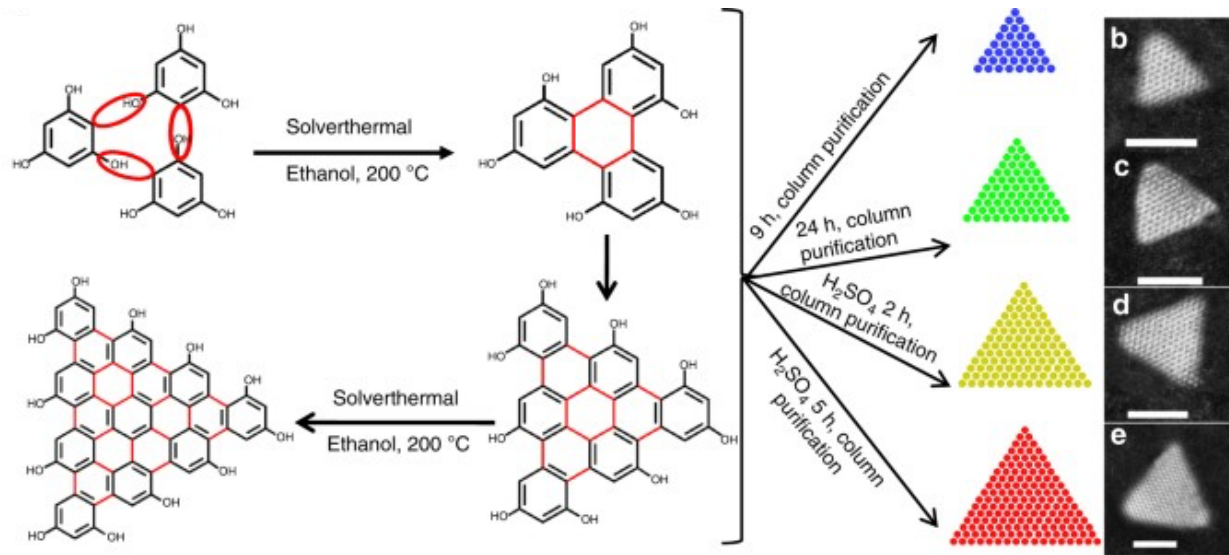


chromatic tunability → optimize for quantum efficiency of PD (fast, optimizable WLS)

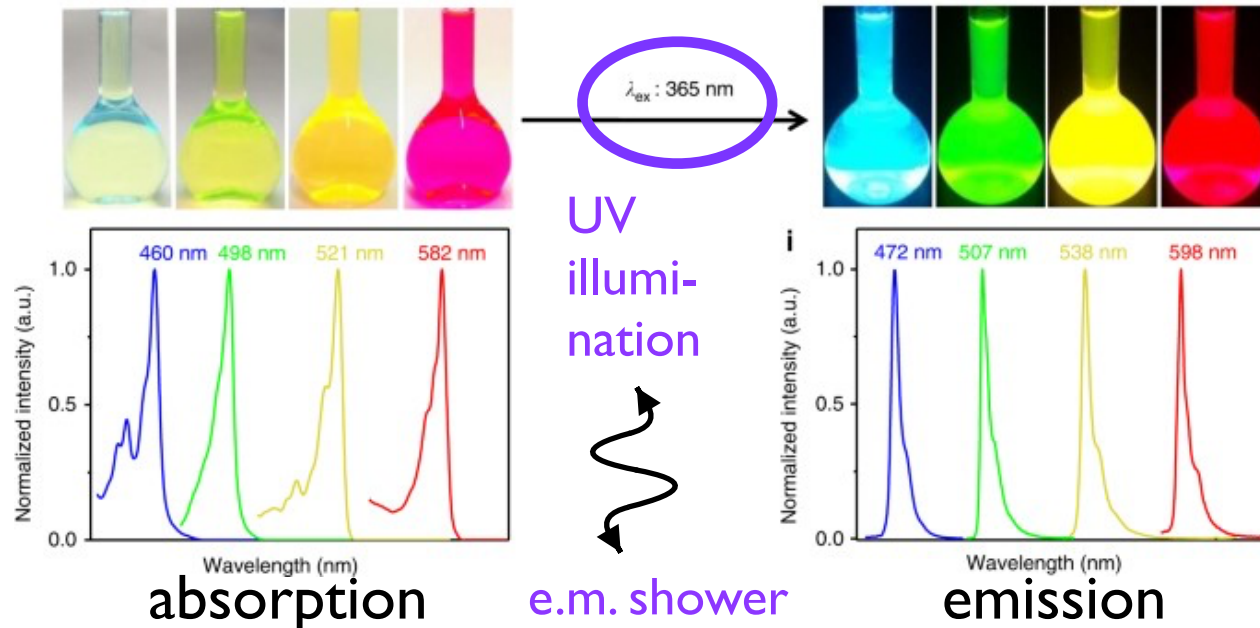
deposit on surface of high-Z material → thin layers of UV → VIS WLS

embed in high-Z material ? two-species (nanodots + microcrystals) embedded in polymer matrix?
 → quasi continuous VIS-light emitter (but what about re-absorbtion?)

Quantum dots: chromatic calorimetry



idea: seed different parts of a “crystal” with nanodots emitting at different wavelengths, such that the wavelength of a stimulated fluorescence photon is uniquely assignable to a specific nanodot position



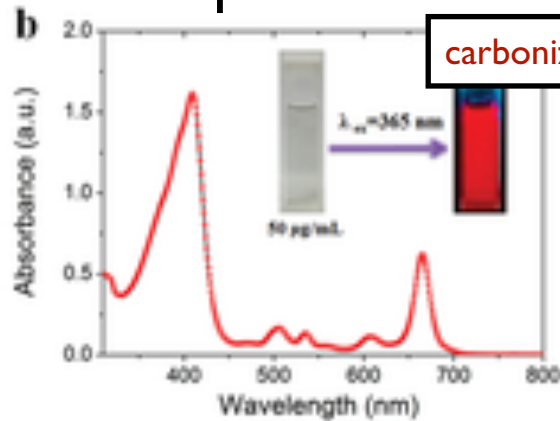
requires:

- narrowband emission (~20nm)
- only absorption at longer wavelengths
- short rise / decay times

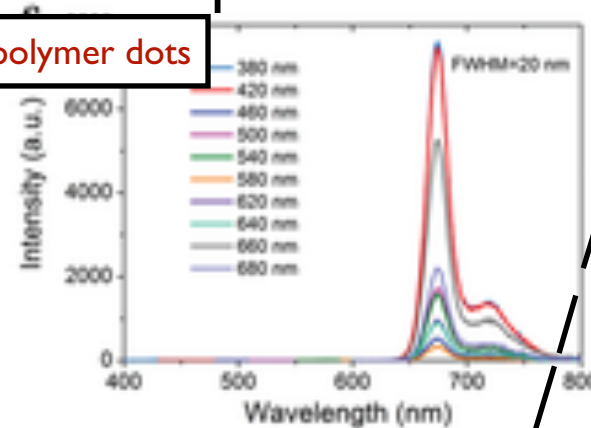
select appropriate nanodots

e.g. **triangular carbon nanodots**

absorption spectrum



emission spectrum



leftmost nanodots:

absorb wavelengths < 650 nm
emit at > 680 nm

next band:

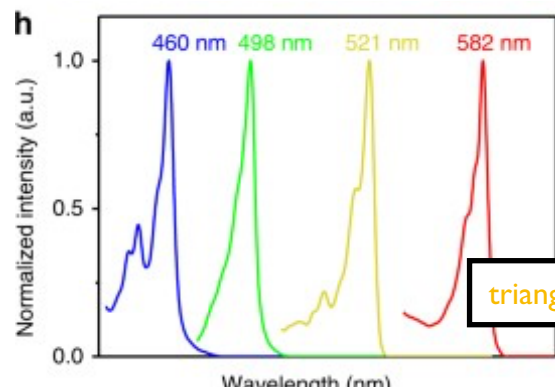
absorb wavelengths < 590 nm
emit at > 590 nm

...

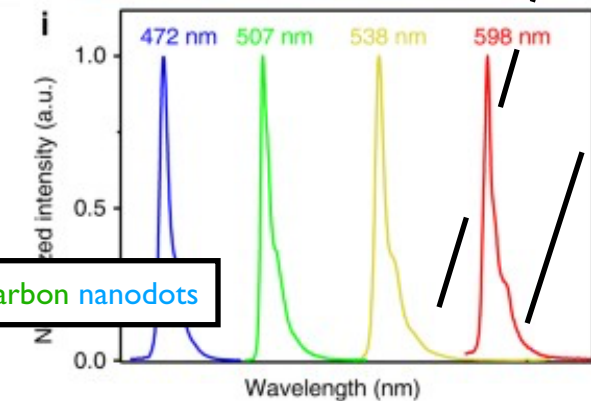
rightmost nanodots:

absorb wavelengths < 410 nm
emit at > 420 nm

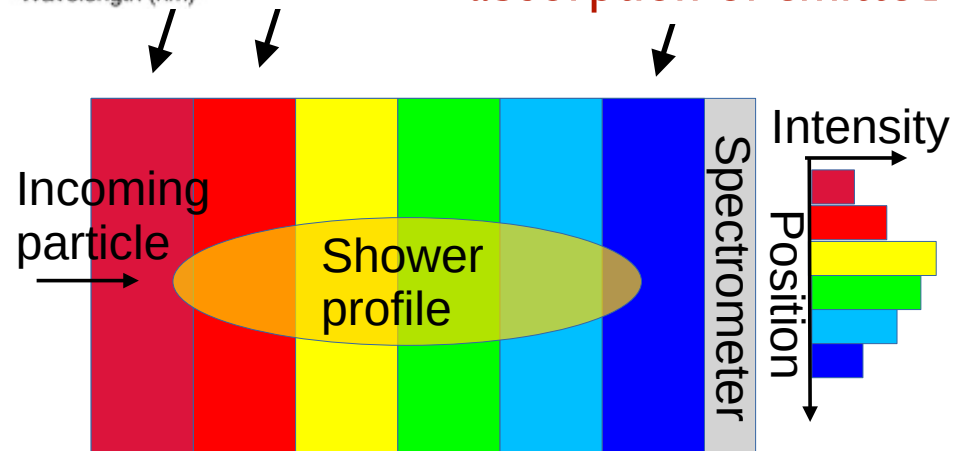
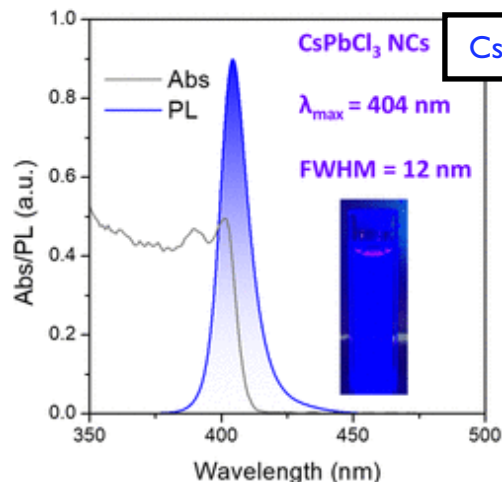
if high-Z substrate transparent
in 400-700 nm, then no re-
absorption of emitted light



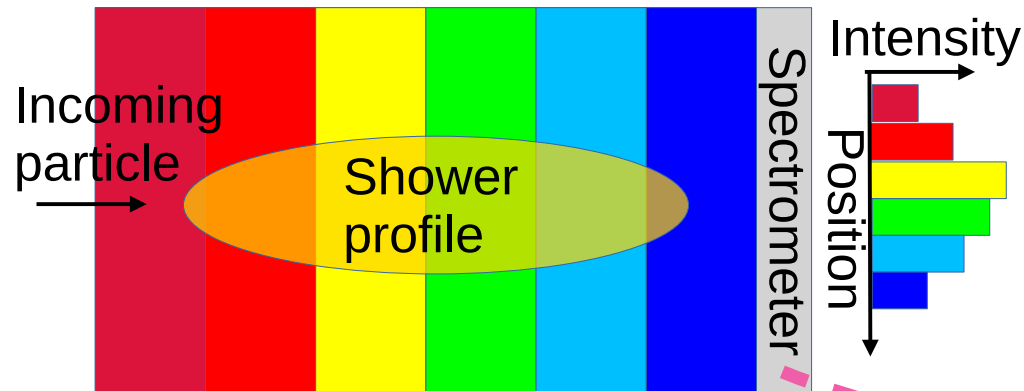
triangular carbon nanodots



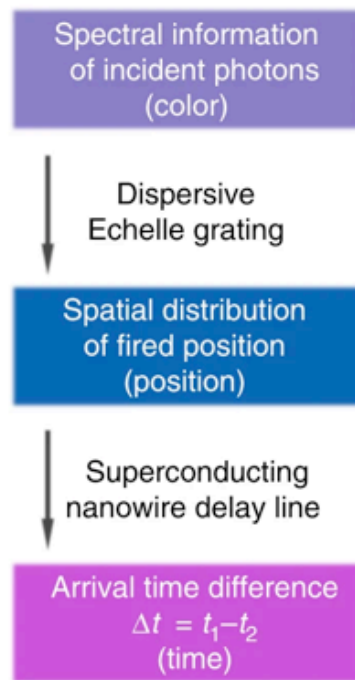
CsPbCl₃ nanocrystals



Quantum dots: chromatic calorimetry (shower profile via spectrometry)



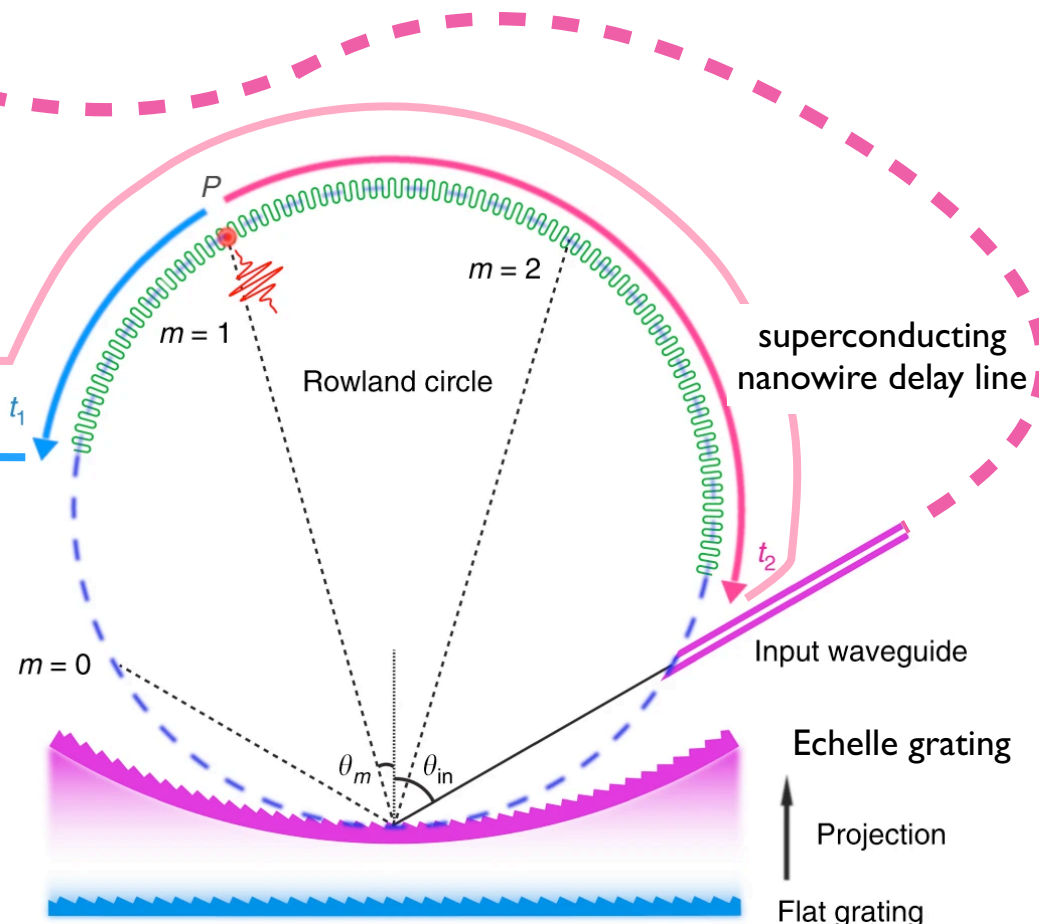
- different options for spectrometry:
- monochromators + PD
 - light guiding fiber / each layer
 - light guiding fiber to spectrometer



cryogenic amplifier

DC current

cryogenic amplifier



R. Cheng, H. X. Tang, et al., Broadband on-chip single-photon spectrometer, Nat Commun 10 (2019) 4104; <https://www.nature.com/articles/s41467-019-12149-x>

Active scintillators (QWs, QDs, QWDs, QCLs)

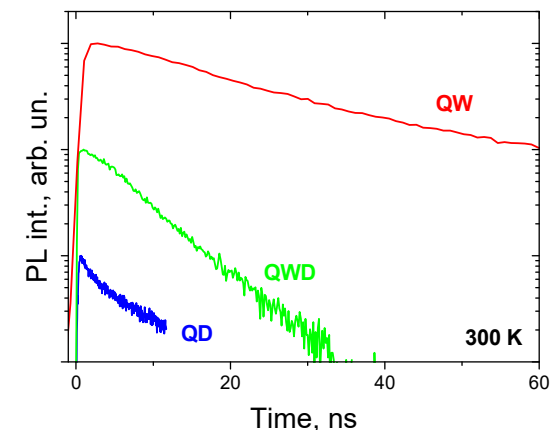
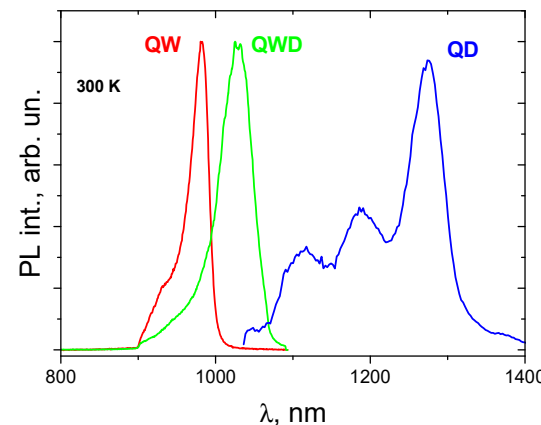
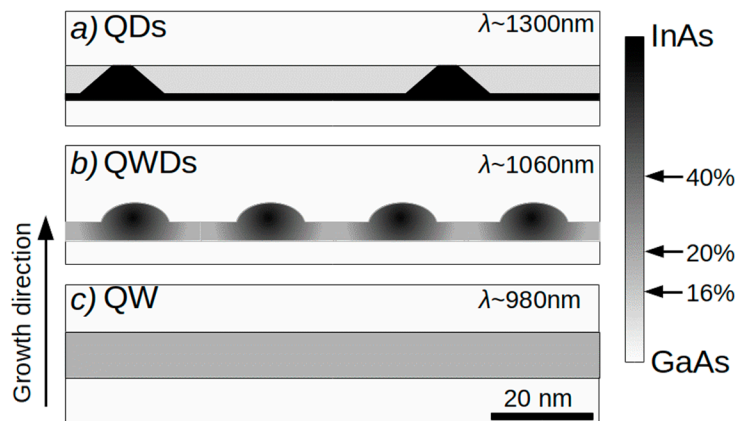
standard scintillating materials are **passive**

- can not be amplified
- can not be turned on/off
- can not be modified once they are in place

is it possible to produce **active** scintillating materials?

- electronically amplified / modulable
- pulsed / primed
- gain adapted in situ

existing QD's, QWD's are elements of optoelectronic devices, typically running at 10 GHz, quite insensitive to temperature



Light Emitting Devices Based on Quantum Well-Dots, Appl. Sci. 2020, 10, 1038; doi:10.3390/app10031038

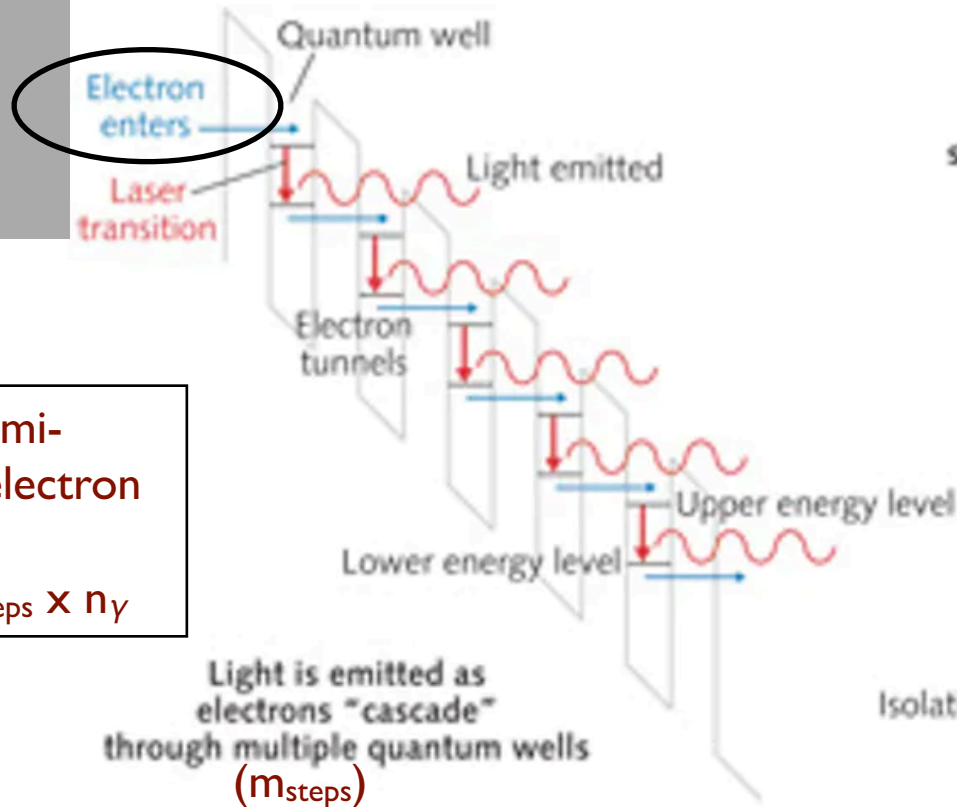
Emission in **IR**! Silicon is transparent at these wavelengths...
Can this IR light be transported *through* a tracker to outside PDs?

QD's are radiation resistant

R. Leon et al., "Effects of proton irradiation on luminescence emission and carrier dynamics of self-assembled III-V quantum dots," in IEEE Transactions on Nuclear Science, 49, 6, 2844-2851 (2002), doi: 10.1109/TNS.2002.806018.

Active scintillators (QCLs, QWs, QDs, QWDs)

<https://www.laserfocusworld.com/test-measurement/spectroscopy/article/16556856/quantumcascade-lasers-qcls-enable-applications-in-ir-spectroscopy>

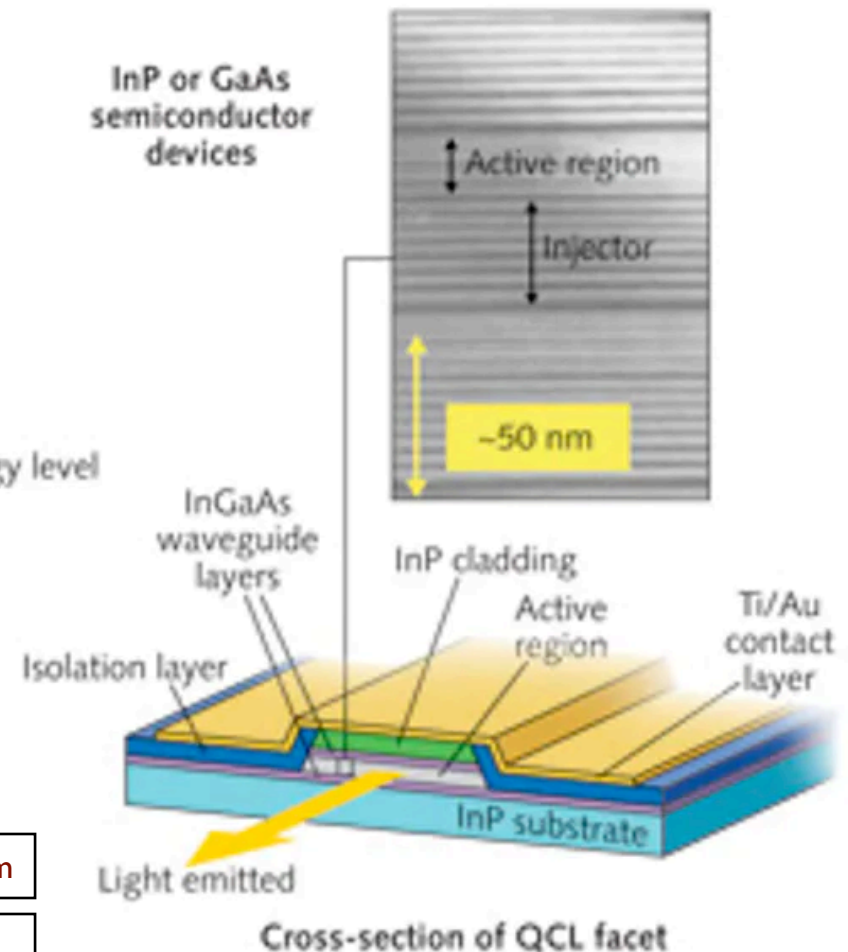


Couple bulk semiconductor to electron injection layer:

$$n_e \longrightarrow m_{\text{steps}} \times n_y$$

Emitted light is IR~THz, normally mono-chromatic but tunable from $3 \mu\text{m} \sim 12 \mu\text{m}$

Radiation resistant ([Radiation Physics and Chemistry 174](#), 2020, 108983)



Quantum dots and wells: <https://arxiv.org/abs/2202.11828>

submicron pixels

DoTPiX

= single n-channel MOS transistor, in which a buried quantum well gate performs two functions:

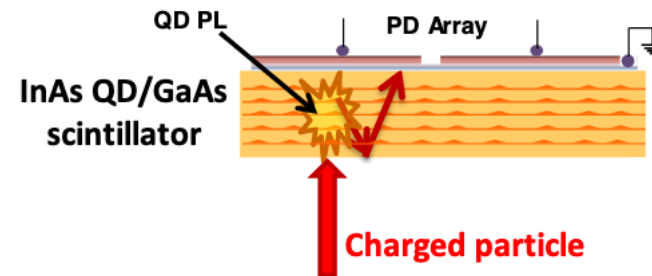
- as a hole-collecting electrode and
- as a channel current modulation gate

Novel Sensors for Particle Tracking: a Contribution to the Snowmass Community Planning Exercise of 2021

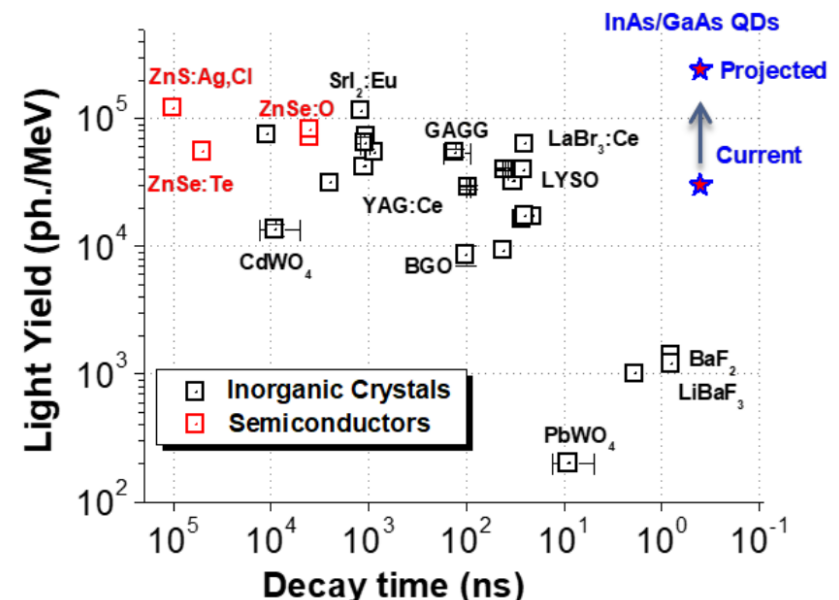
[M.R. Hoferkamp](#), [S. Seidel](#), [S. Kim](#), [J. Metcalfe](#), [A. Sumant](#), [H. Kagan](#), [W. Trischuk](#), [M. Boscardin](#), [G.-F. Dalla Betta](#), [D.M.S. Sultan](#), [N.T. Fourches](#), [C. Renard](#), [A. Barbier](#), [T. Mahajan](#), [A. Minns](#), [V. Tokranov](#), [M. Yakimov](#), [S. Oktyabrsky](#), [C. Gingu](#), [P. Murat](#), [M.T. Hedges](#)

<https://arxiv.org/abs/2202.11828>

scintillating (chromatic) tracker



IR emission from InAs QD's
integrated PD's (1-2 μm thick)



2-D materials for MPGDs

Florian Brunbauer / CERN

State-of-the-art MPGDs:

- high spatial resolution
- good energy resolution
- timing resolution $< 25\text{ps}$ (PICOSEC Micromegas)

use of 2-D materials to improve:

- tailor the primary charge production process,
- protect sensitive photocathodes in harsh environments
- improve the performance of the amplification stage

tunable work function

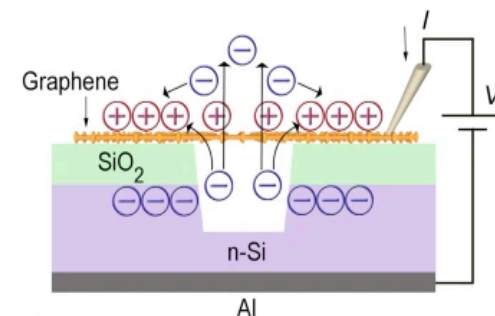
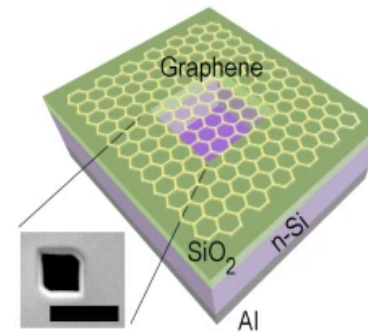
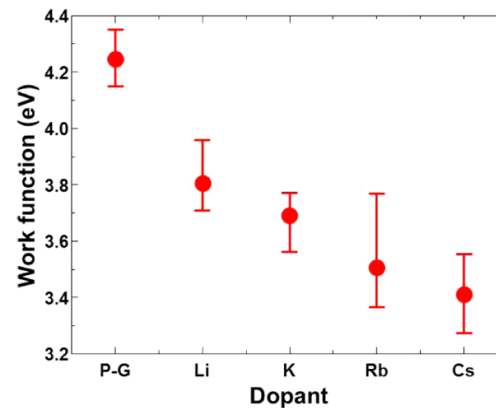
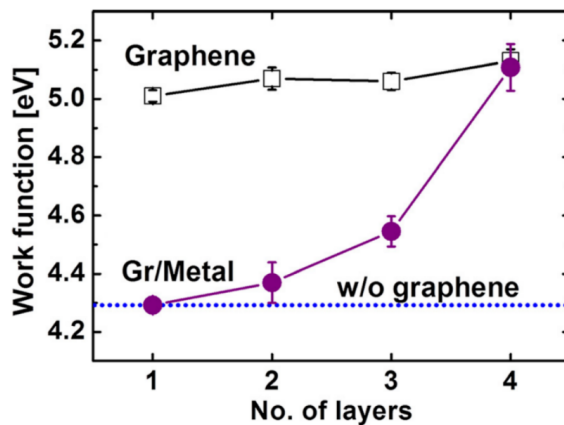
efficiency of the photocathode \rightarrow timing resolution; QE
tune via resonant processes in low dimensional coating structures

(additionally, encapsulation of semiconductive as well as metallic (i.e. Cu) photocathodes increases operational lifetime)

amplification

back flow of positive ions created during charge amplification to the drift region can lead to significant distortions of electric fields

Graphene has been proposed as selective filter to **suppress** ion back flow while **permitting** electrons to pass:
Good transparency (up to $\sim 99.9\%$) to very low energy ($< 3\text{ eV}$) electrons (?)



Tuning the work function of graphene toward application as anode and cathode, Samira Naghdi, Gonzalo Sanchez-Arriaga, Kyong Yop Rhee, <https://arxiv.org/abs/1905.06594>

Space charge neutralization by electron-transparent suspended graphene, Siwapon Srisophonpan, Myungji Kim & Hong Koo Kim, [Scientific Reports](#) 4, 3764 (2014)

Metamaterials, 0 / 1 / 2-dimensional materials (quantum dots, nanolayers)

ultra-fast scintillators based on perovskites

chromatic calorimetry (QDs)

active scintillators (QCL, QWs, QDs)

5.3.6

GEMs (graphene)

Atoms, molecules, ions

Rydberg TPC's

5.3.5

Spin-based sensors

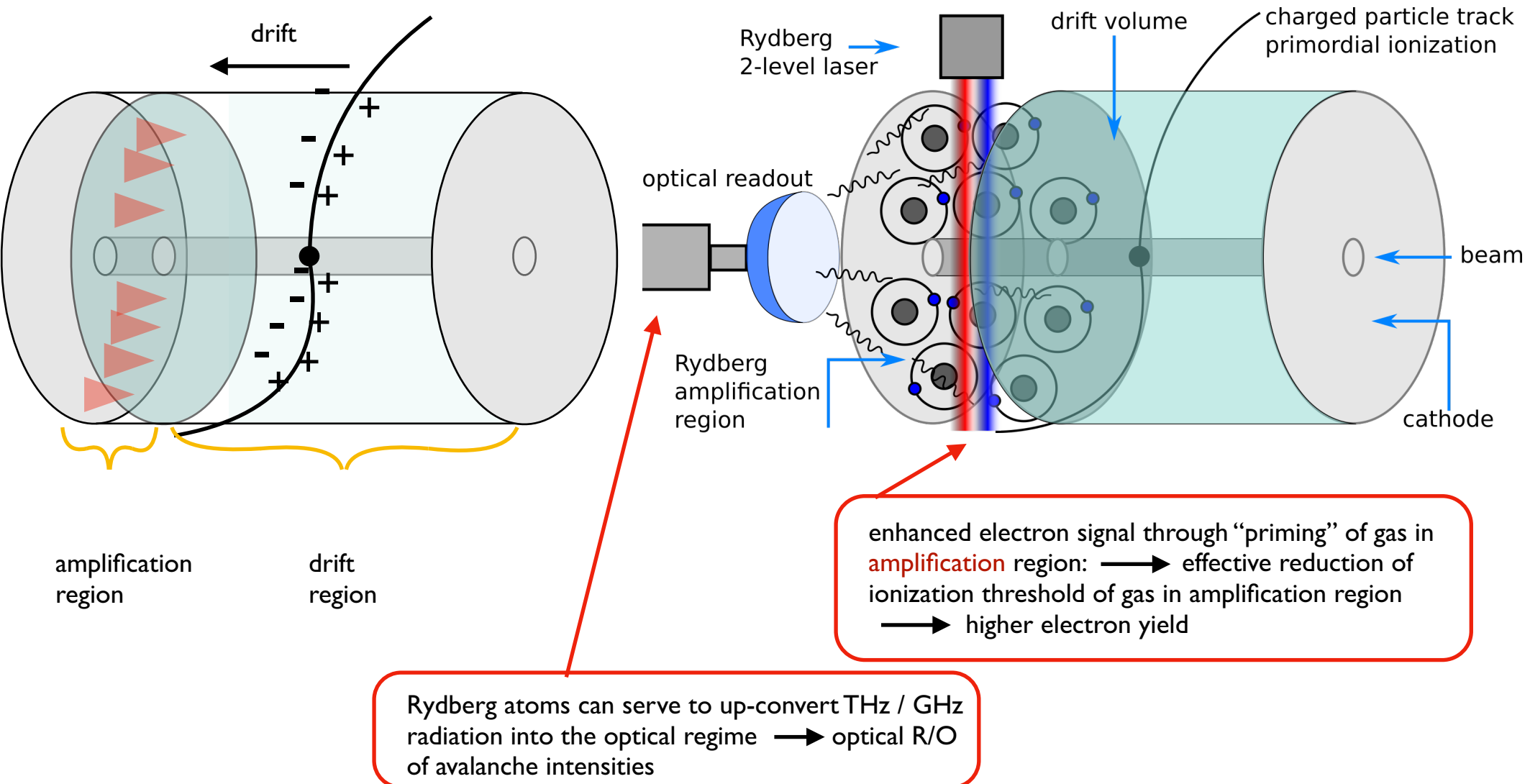
helicity detectors

5.3.3

Rydberg atom TPC's

Georgy Kornakov / WUT

Act on the amplification region



Rydberg atom TPC's

Georgy Kornakov / WUT

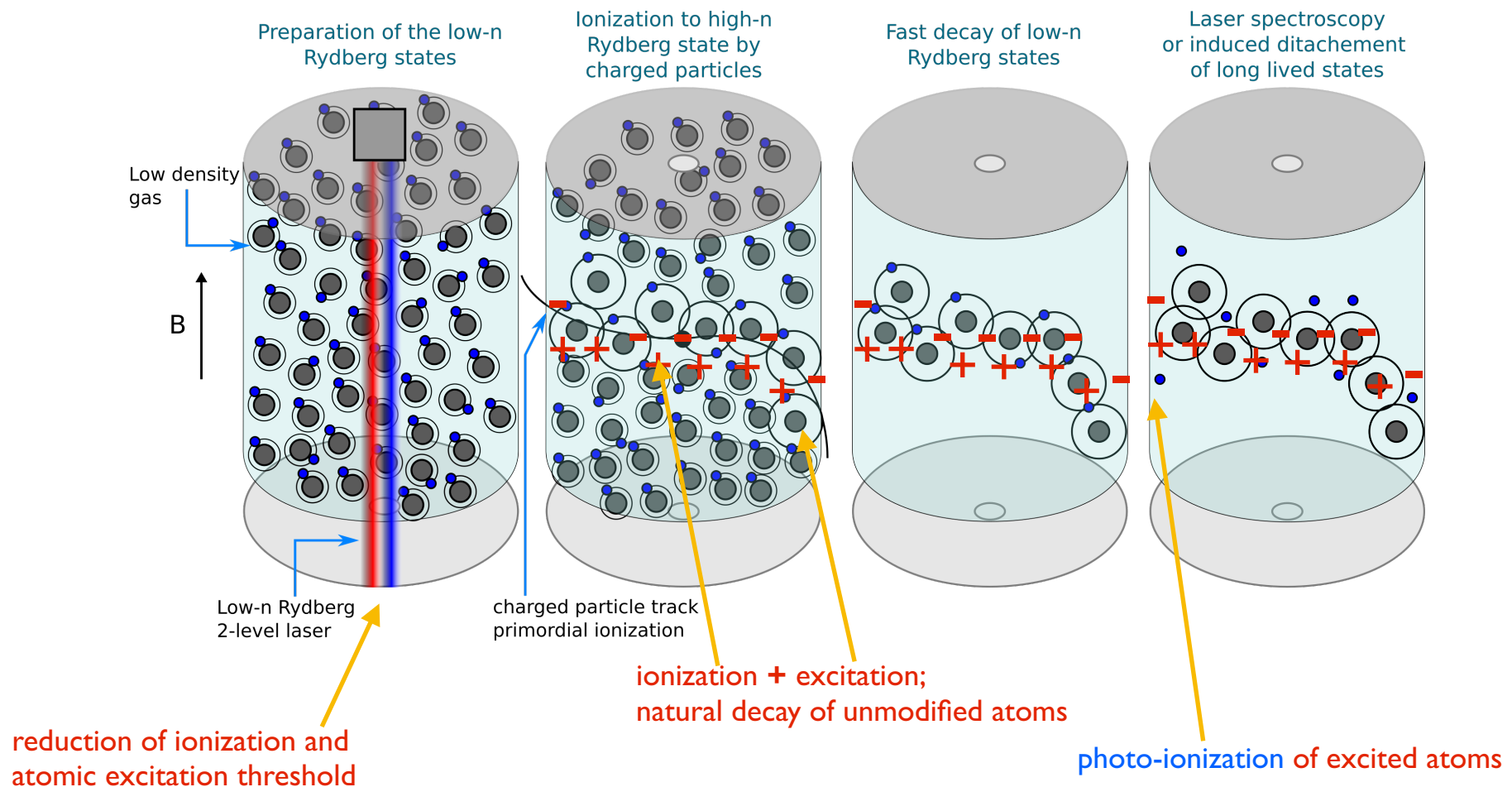
Act on the drift region

principle carries over to drift region:

enhanced electron signal through “priming” of gas in **drift** region:

effective reduction of ionization threshold of gas in drift region

increased dE/dx through standard **primary** ionization + **photo-ionization** of atoms excited by mip's

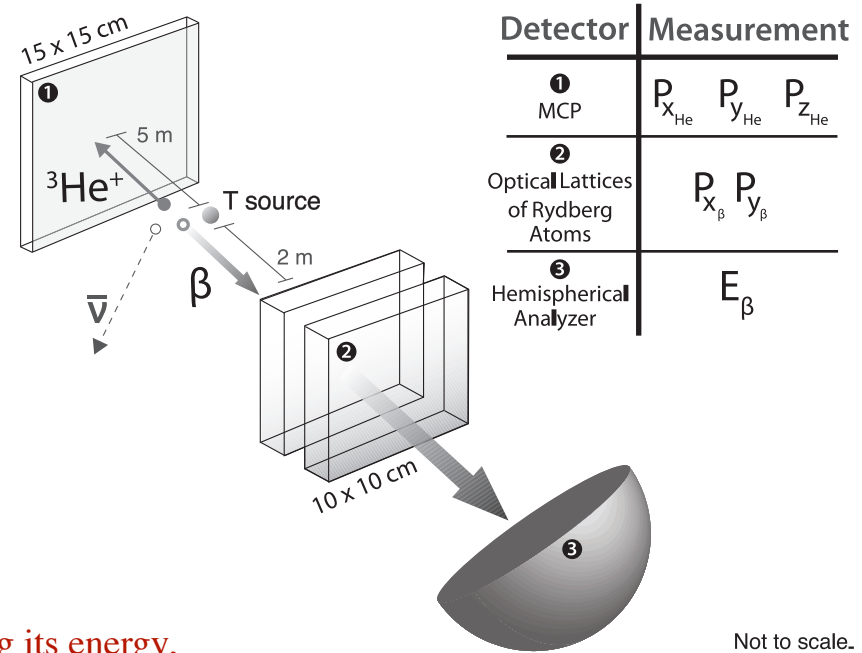
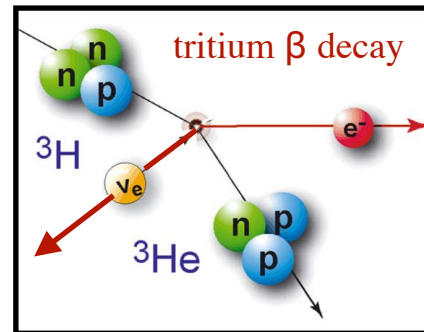


Rydberg atom TPC's

Using cold atoms to measure neutrino mass

M Jerkins, J R Klein, J H Majors, F Robicheaux and M G Raizen
2010, *New J. Phys.* **12** 043022

<https://iopscience.iop.org/article/10.1088/1367-2630/12/4/043022/meta>



“In order to measure the momentum of the β without significantly altering its energy, we propose exploiting the effect of a passing electron on Rydberg atoms [22, 23]. In the β 's flight path before it reaches the spectrometer, we create an optical lattice filled with rubidium atoms in the ground state [24, 25]. Using laser excitation, we can excite the atoms to a high Rydberg state [26, 27], such as 53s. When the β passes one of these atoms, it can excite the atom from a 53s state to a 53p state, and the atom will remain trapped in its optical lattice position. We propose slowing the electrons with a controlled voltage soon after they leave the source so that by the time they reach the optical lattice, they have a maximum energy of 900 eV, which increases their cross section for exciting a Rydberg atom to $0.36 \times 10^{-9} \text{ cm}^2$. When a β signal is detected downstream in the spectrometer, the 53 s atoms are optically de-excited using stimulated Raman adiabatic passage (STIRAP) [27], and an electric field of 100 V cm^{-1} is ramped within $\sim 130 \text{ ns}$ to ionize any Rydberg atoms in a 53p state. Once the atoms are ionized, they will be detected by a multi-hit position-sensitive MCP. Based on realistic density limits, the β will excite several Rydberg atoms as it passes through the optical lattice, so we will be able to obtain the projection of a track from the passing β .”

Metamaterials, 0 / 1 / 2-dimensional materials (quantum dots, nanolayers)

ultra-fast scintillators based on perovskites

chromatic calorimetry (QDs)

active scintillators (QCL, QWs, QDs)

5.3.6

GEMs (graphene)

Atoms, molecules, ions

Rydberg TPC's

5.3.5

Spin-based sensors

helicity detectors

5.3.3

optically polarizable elements: Nitrogen-vacancy diamonds (NVD)

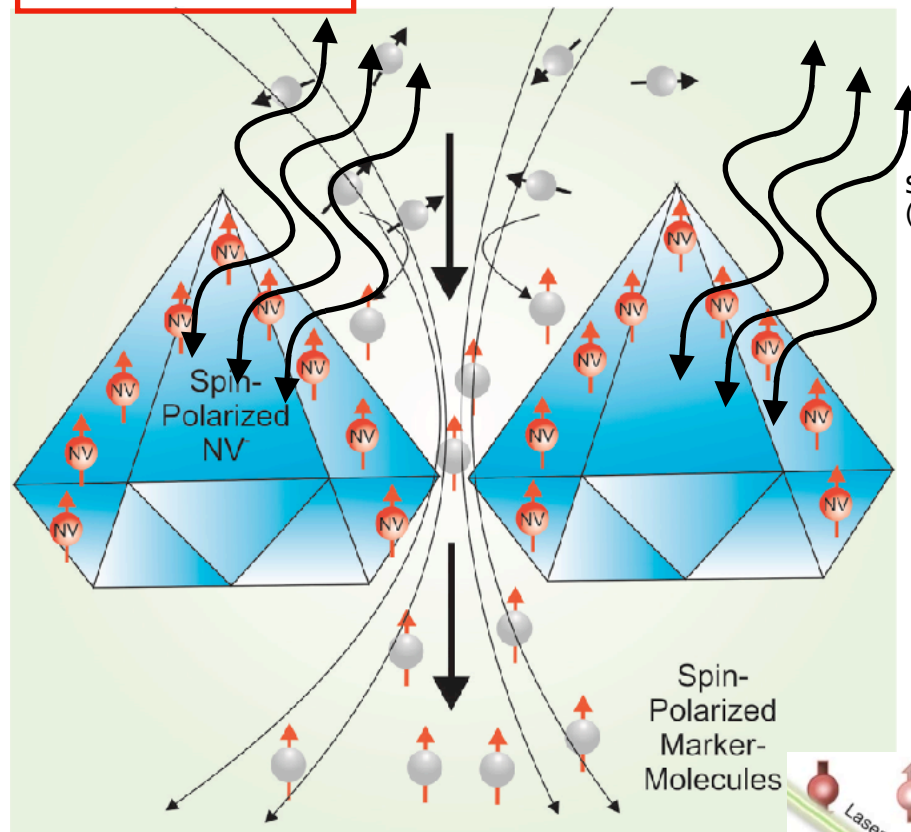
Georgy Kornakov / WUT

spin-spin scattering for helicity determination: usually with polarized beams and/or polarized targets

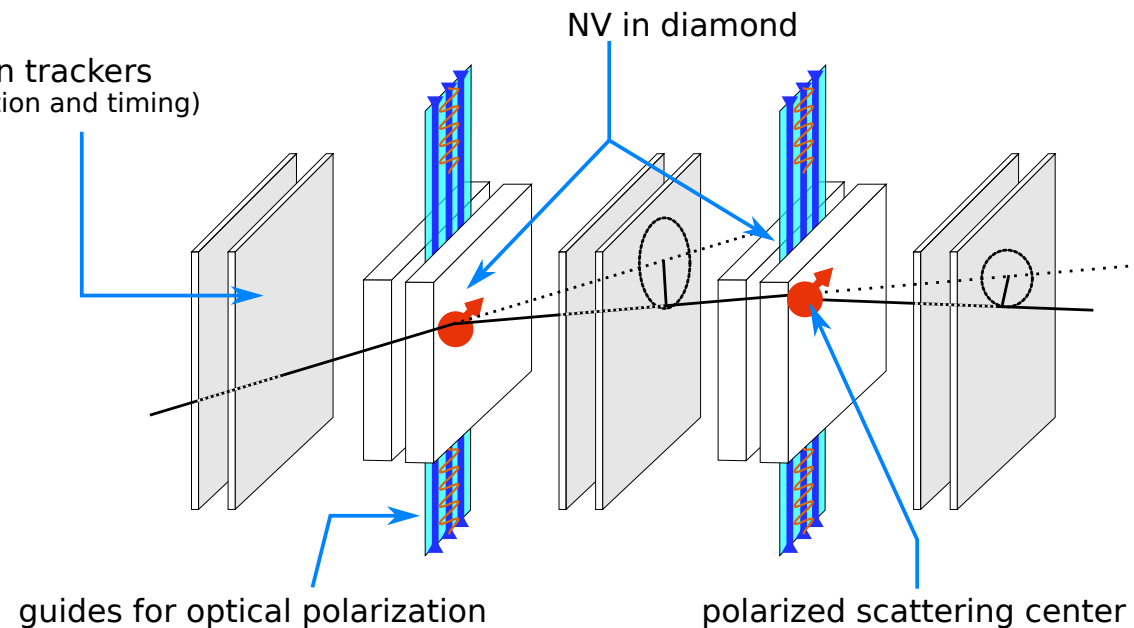
introduce polarized scattering planes to extract track-by-track particle helicity

$10^{16} \sim 10^{18} / \text{cm}^3$

use NV-diamonds to scatter mips (spin-dependent scattering)



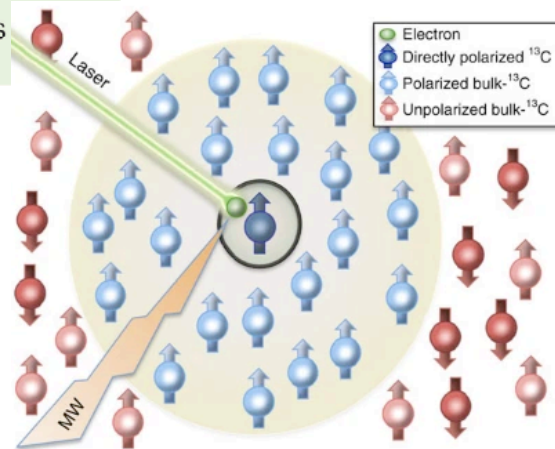
silicon trackers
(direction and timing)



© Dr. Christoph Nebel, Fraunhofer IAF

https://www.metaboliqs.eu/en/news-events/MetaboliQs_PM_first_year.html

Diamond plates of up to $8 \times 8 \text{ mm}^2$ in size, fabricated by Element Six



Local and bulk ^{13}C hyperpolarization in nitrogen-vacancy centred diamonds at variable fields and orientations, G. Alvarez et al., *Nature Communications* **6**, 8456 (2015)

<https://www.nature.com/articles/ncomms9456>

$\times 10^2$

optically polarizable elements: Nitrogen-vacancy diamonds (NVD)

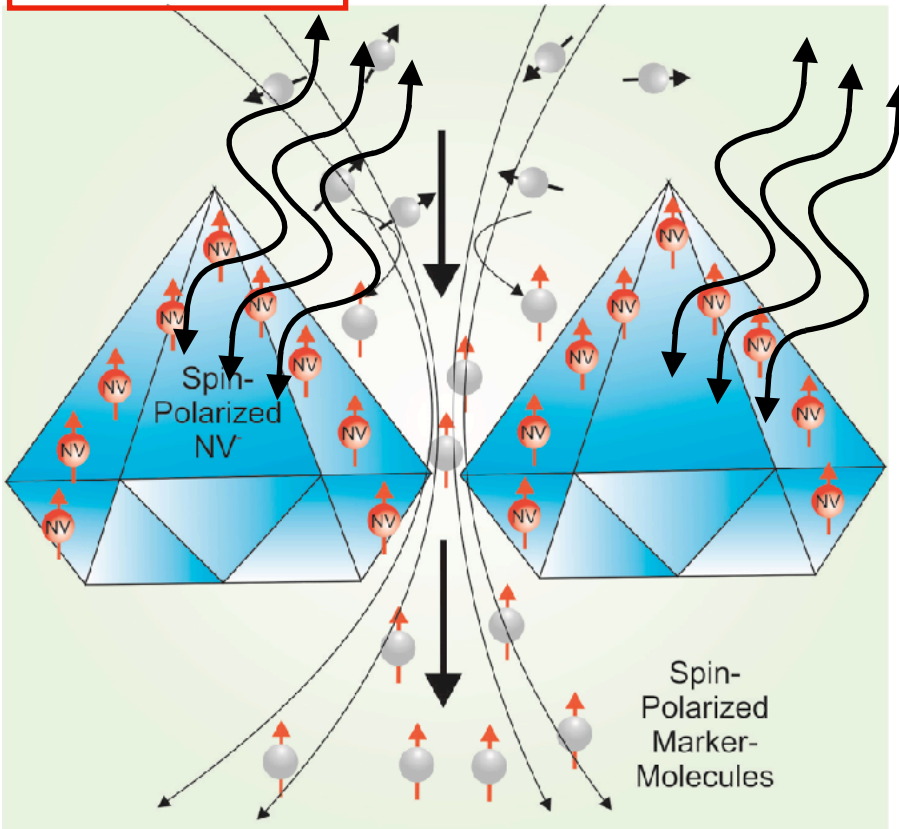
Georgy Kornakov / WUT

spin-spin scattering for helicity determination: usually with polarized beams and/or polarized targets

polarimetry of target fragments (active target)

$10^{16} \sim 10^{18} / \text{cm}^3$

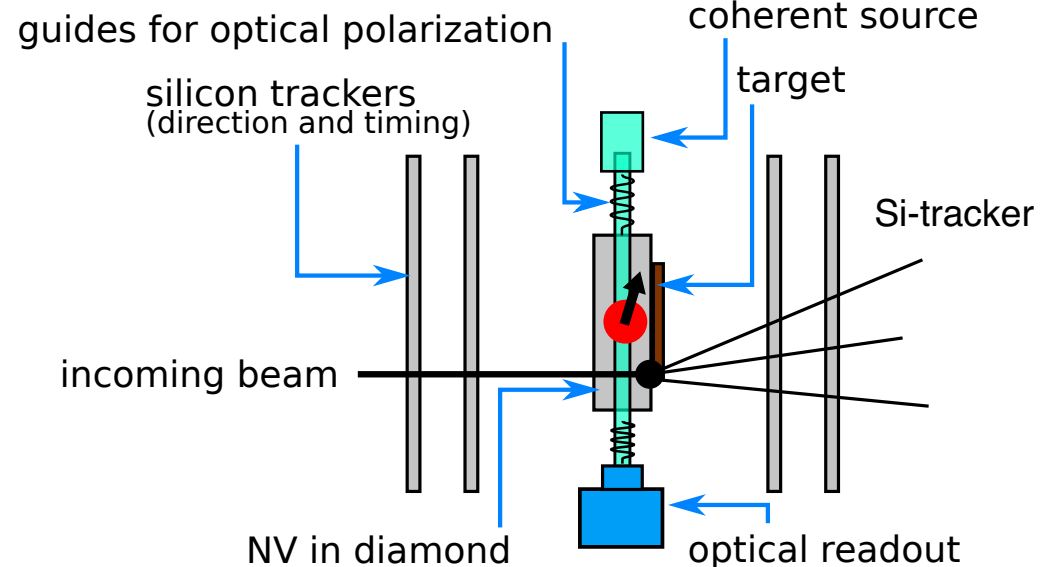
use NV-diamonds to sense spin of heavy nuclear spectator fragment



© Dr. Christoph Nebel, Fraunhofer IAF

https://www.metaboliqs.eu/en/news-events/MetaboliQs_PM_first_year.html

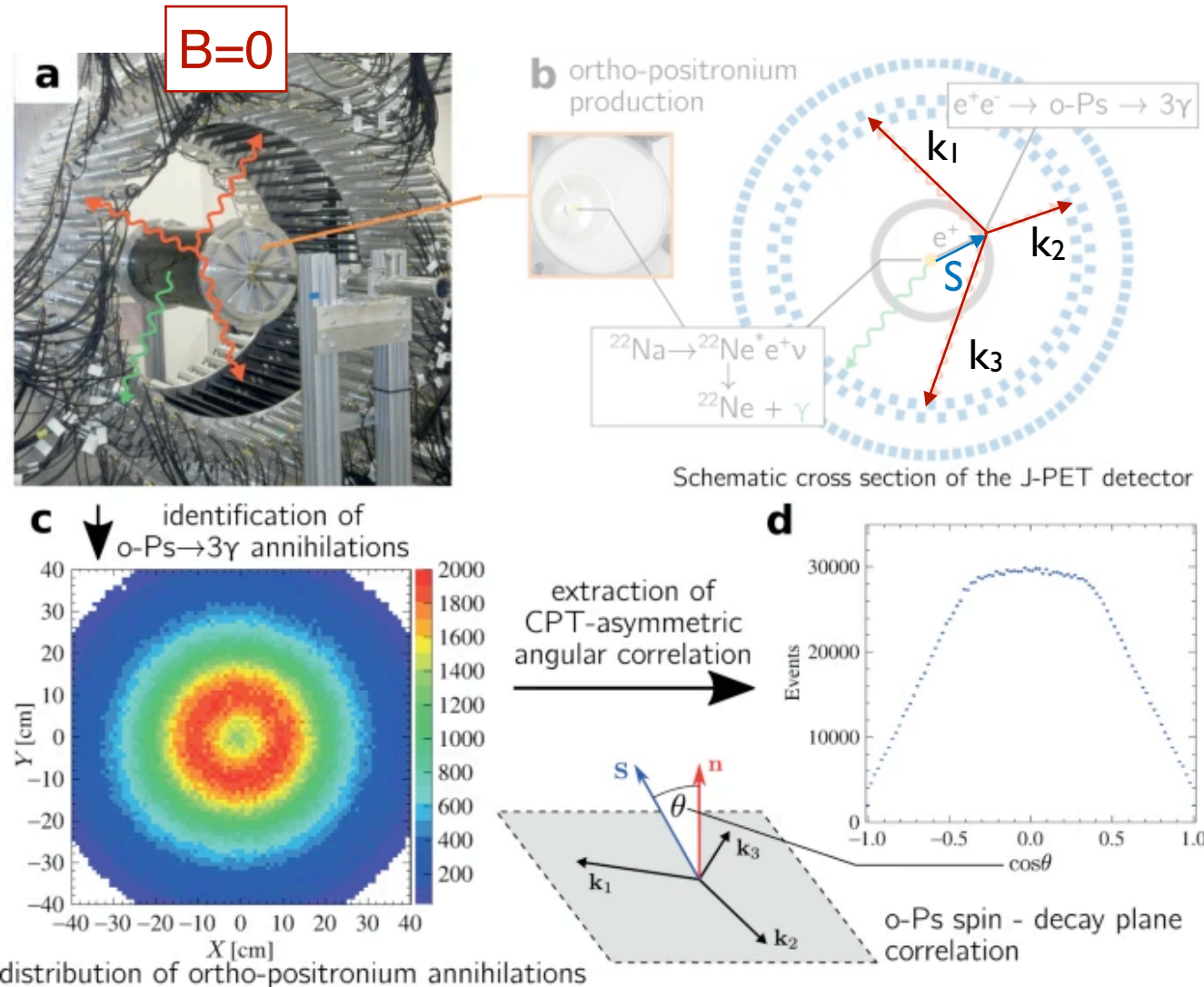
Diamond plates of up to $8 \times 8 \text{ mm}^2$ in size, fabricated by Element Six



measurement of photon polarizations

Pawel Moskal / Krakow

J-PET:3-photon correlations / entanglement



discrete symmetry operators

Operator	C	P	T	CP	CPT
$S \cdot k_1$	+	-	+	-	-
$S \cdot (k_1 \times k_2)$	+	+	-	+	-
$(S \cdot k_1) \cdot (S \cdot (k_1 \times k_2))$	+	-	-	-	+

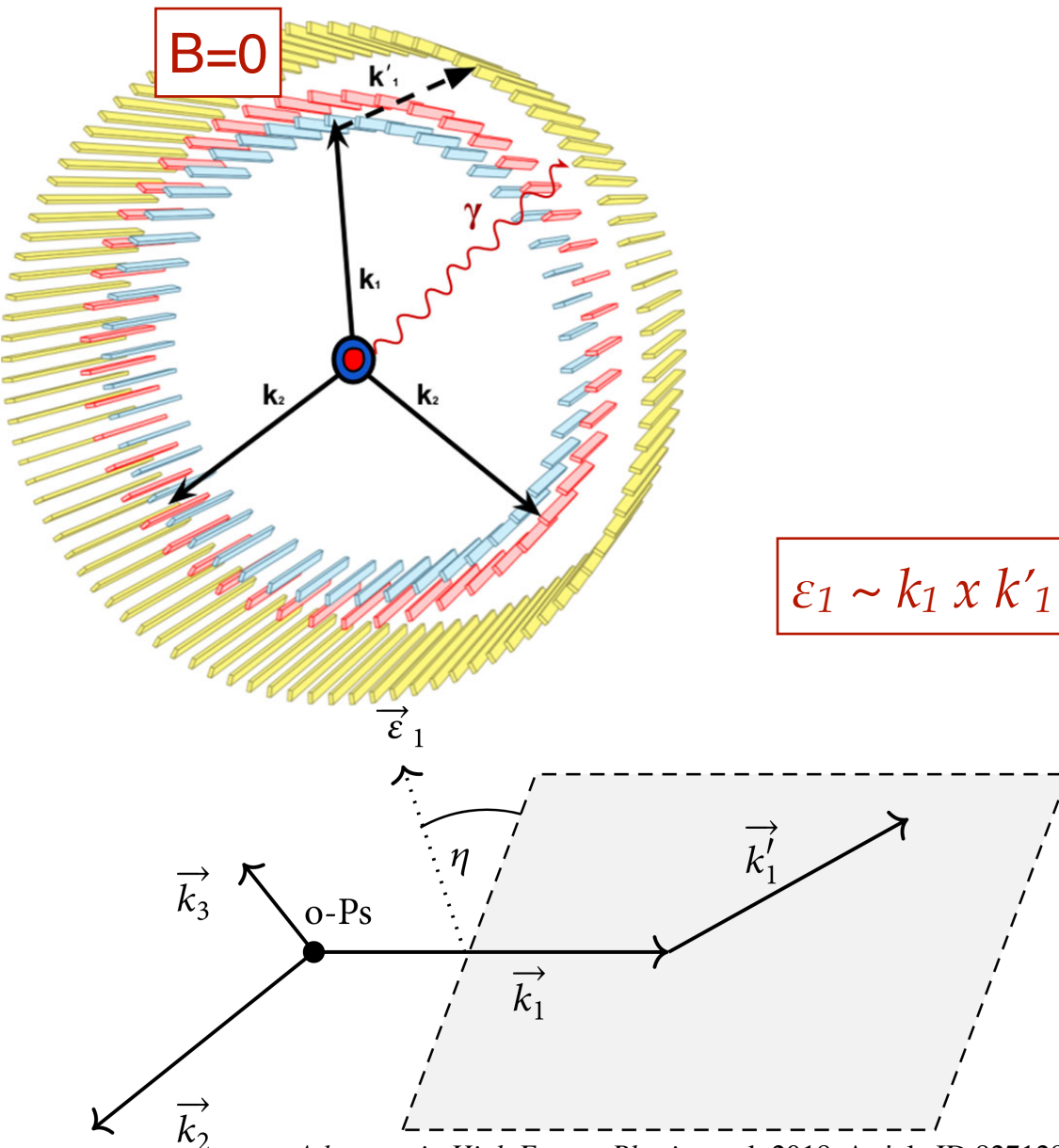
Operator	C	P	T	CP	CPT
$\epsilon_1 \cdot k_2$	+	-	-	-	+
$S \cdot \epsilon_1$	+	+	-	+	-
$S \cdot (k_2 \times \epsilon_2)$	+	-	+	-	-

limit on CPT violation in purely leptonic system

Hyperfine Interact (2018) 239: 56 <https://doi.org/10.1007/s10751-018-1527-x>

measurement of photon polarizations

Pawel Moskal / Krakow

Compton scattering \rightarrow polarization

$$\epsilon_1 \sim k_1 \times k'_1$$

discrete symmetry operators

Operator	C	P	T	CP	CPT
$S \cdot k_1$	+	-	+	-	-
$S \cdot (k_1 \times k_2)$	+	+	-	+	-
$(S \cdot k_1) \cdot (S \cdot (k_1 \times k_2))$	+	-	-	-	+

Operator	C	P	T	CP	CPT
$\epsilon_1 \cdot k_2$	+	-	-	-	+
$S \cdot \epsilon_1$	+	+	-	+	-
$S \cdot (k_2 \times \epsilon_2)$	+	-	+	-	-

novel symmetry operators

challenge: is it possible
 \rightarrow to identify Compton
 scattered HE γ 's?

What's next?

These potential applications of quantum sensors also in HEP require dedicated R&D to evaluate their potential and feasibility.

In line with the RECFA R&D roadmap, it makes sense to consider a quantum-sensing R&D program that brings together the following strands:

Clocks and clock networks 5.3.1

Kinetic detectors 5.3.2

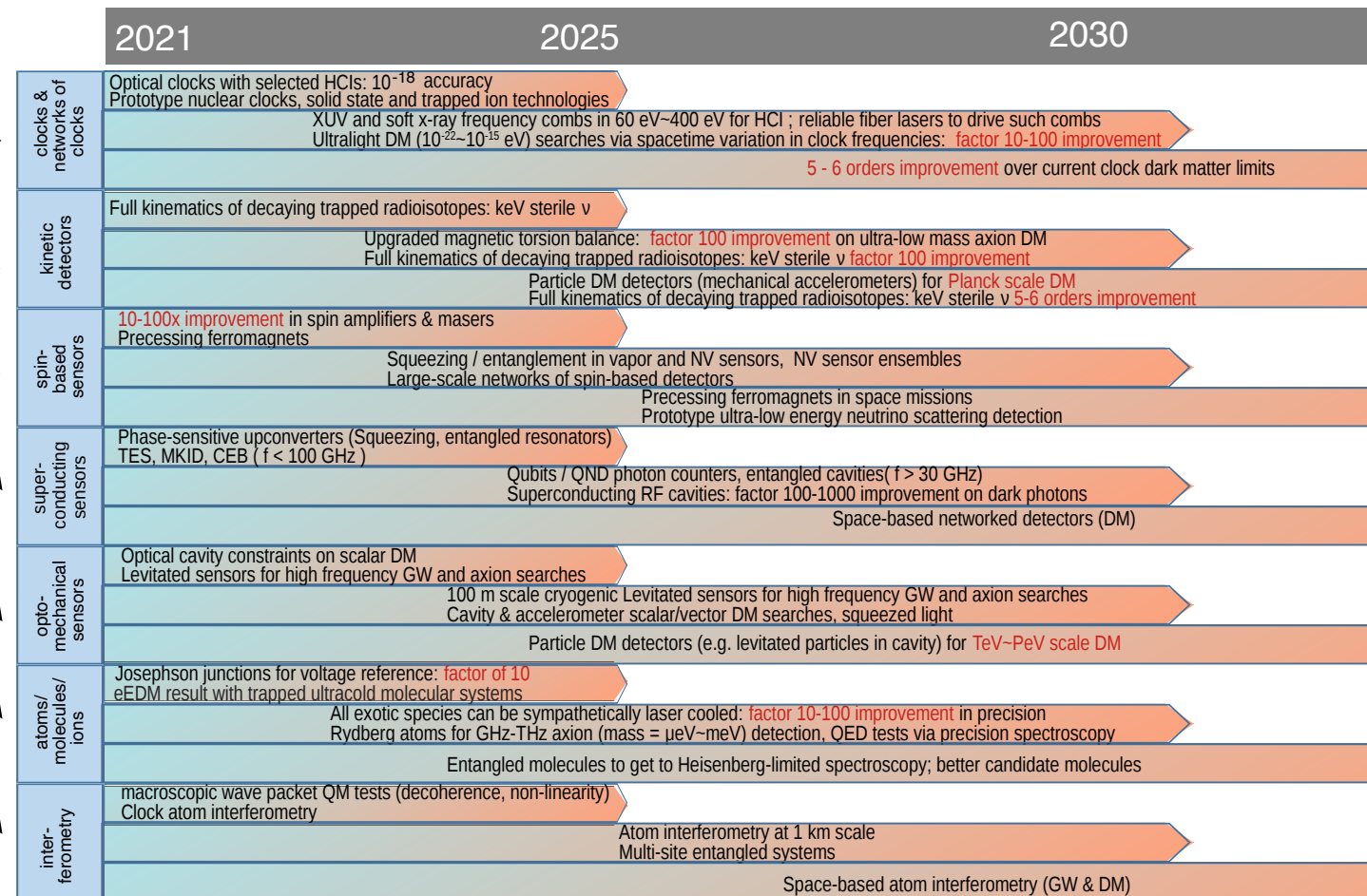
Spin-based sensors 5.3.3
Superconducting sensors 5.3.3

Optomechanical sensors 5.3.4

Atoms/molecules/ions 5.3.5
Atom interferometry 5.3.5

Metamaterials, 0/1/2D-materials
Quantum materials 5.3.6

also for HEP!



Status of TF5

Anna Grassellino, Marcel Demarteau,
Michael Doser, Caterina Braggio,
Stafford Withington, Peter Graham,
John March-Russel, Andrew Geraci

“Recommendations”

- many fascinating opportunities in nascent fields
- encourage exploratory approaches
- adapt funding profiles to both exploratory as well as consolidation approaches:
 - exploratory: funding cycle of 3 years, lightweight grant application, “fail early / fail often / proof-of-principle” mindset
 - consolidation: funding cycle of 10 years, after initial proof of principle, proposal
- importance of interdisciplinarity
 - training not only of early stage researchers but also of established researchers
 - opportunistic (awareness of developments elsewhere - physics or industry)

thank you!