



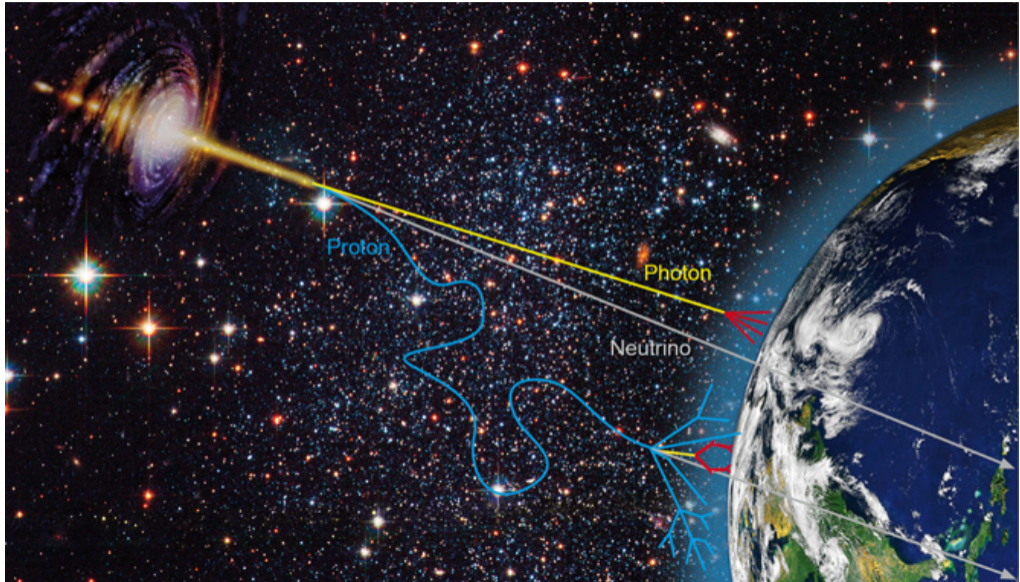
Multi-messenger Neutrino Astrophysics with the Super-Kamiokande detector (*and beyond*)

Fellini seminar

Mathieu Lamoureux
INFN Padova
Secondment in APC (Uni.Paris)

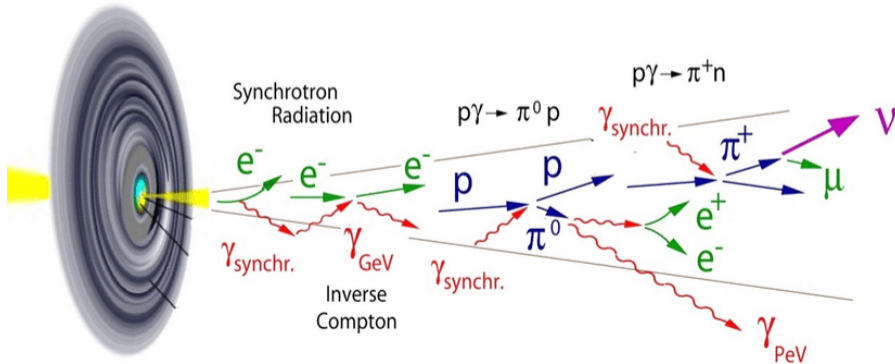
- 1 Neutrino astrophysics
- 2 Super-Kamiokande
- 3 Follow-up of GWTC-2 events
- 4 Outlook

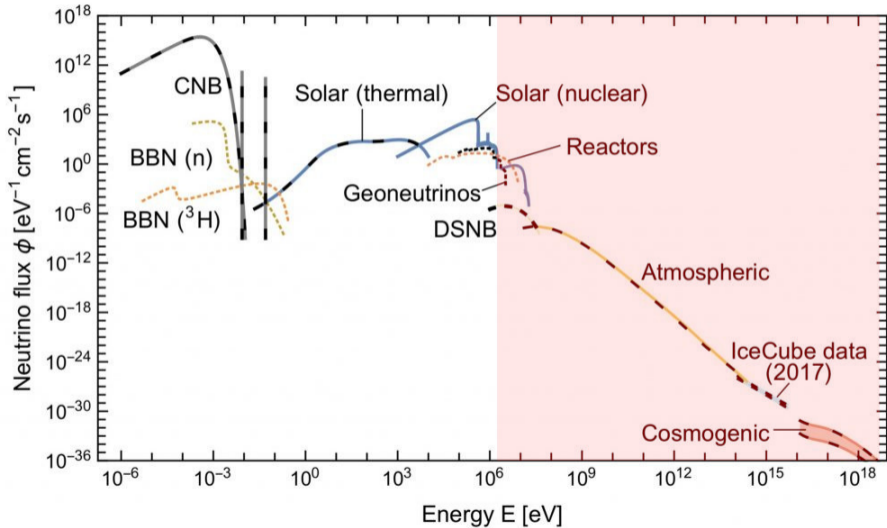
Neutrino astrophysics



Cosmic ray interactions at acceleration sites or during their propagation are producing pions:

- Charged pion decays are emitting neutrinos
- Neutral pion decays are emitting high-energy gammas.

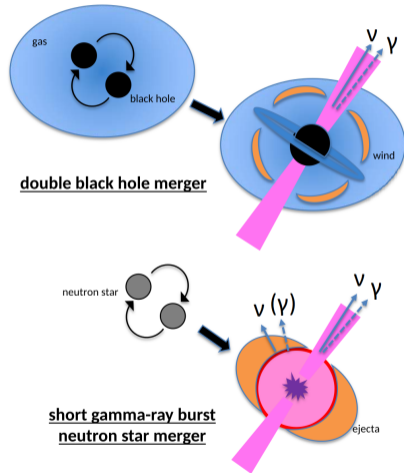




Since 2015, LIGO/Virgo interferometers are detecting gravitational waves from mergers of binary objects.

- Binary Neutrino Star (BNS): may produce short Gamma-Ray Bursts (GRB) with neutrino production
- Binary Black Hole (BBH): neutrino production in the accretion disks of the black holes
- Neutron Star - Black Hole (NSBH)

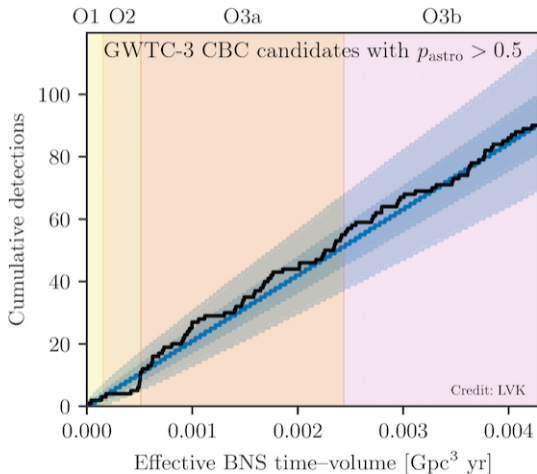
Detecting neutrinos from these objects would allow a better understanding of the mechanisms behind them.



- First detection in 2015
- **GWTC-1:** 11 events from 2015 to 2017 (O1 and O2)
- **GWTC-2:** 39 events from the first half of the third observing run (O3a)
- **GWTC-2.1:** update of GWTC-2, containing events with higher FAR
- **GWTC-3:** 35 events from the second half of the third observing run (O3b)

~ 100 events in total

We expect > 100 events/year starting from O4.



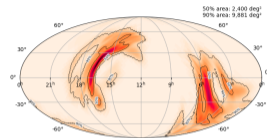
For each GW, we have:

- time of the event
- sky localisation
- estimated distance
- estimated masses

Classification:

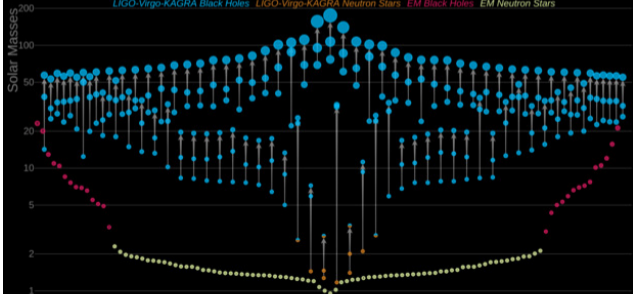
Events can be classified based on object masses:

- $m < 3 M_{\odot} = \text{NS}$
- $m > 3 M_{\odot} = \text{BH}$



Masses in the Stellar Graveyard

LIGO-Virgo-KAGRA Black Holes LIGO-Virgo-KAGRA Neutron Stars EM Black Holes EM Neutron Stars



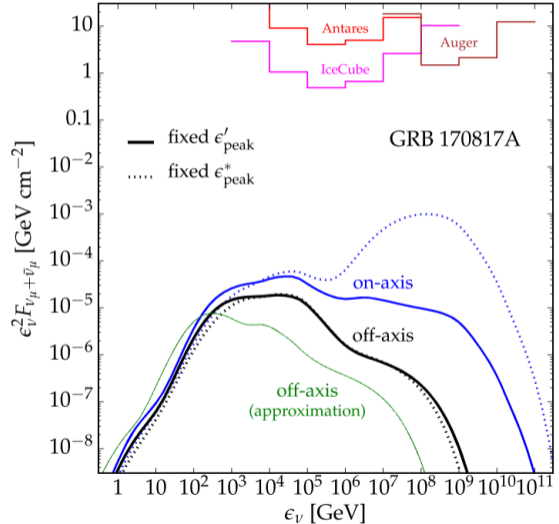
In total:

- 2 BNS
- 5 NSBH
- the rest are BBH

- Very few models available on the market.
- What is the jetting of the neutrino emission?
 - isotropic: easy to constrain but unrealistic, especially for high-energy emission
 - structured jet: in absence of any (EM) counterpart, it is difficult to constrain
- What is the timescale of the neutrino emission? Prompt, precursor, delayed?

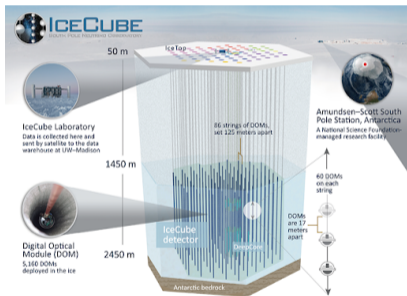
Usual time window

Neutrino signal is usually searched in a 1000 s time window centered on GW.

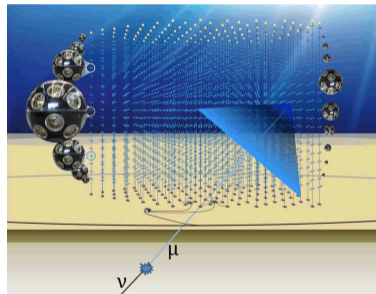


Need very large detectors to detect such low fluxes...

Two big players



IceCube



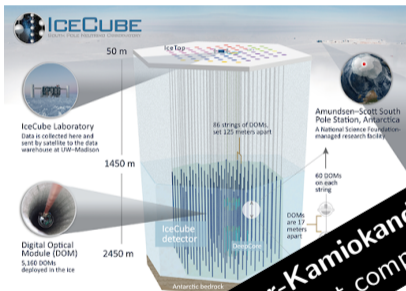
KM3NeT

One kilometer-cube of Ice at the South Pole, instrumented with $\sim 5k$ optical modules.

Kilometer-cube scale infrastructure in the Mediterranean Sea. Currently under deployment.

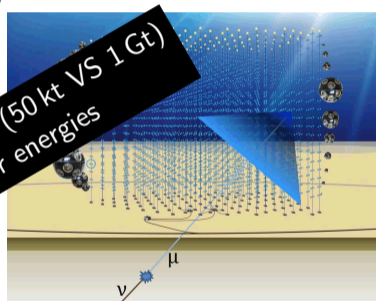
Need very large detectors to detect such low fluxes...

Two big players



Ice

Super-Kamiokande is much smaller (50 kt VS 1 Gt) but complementary for lower energies



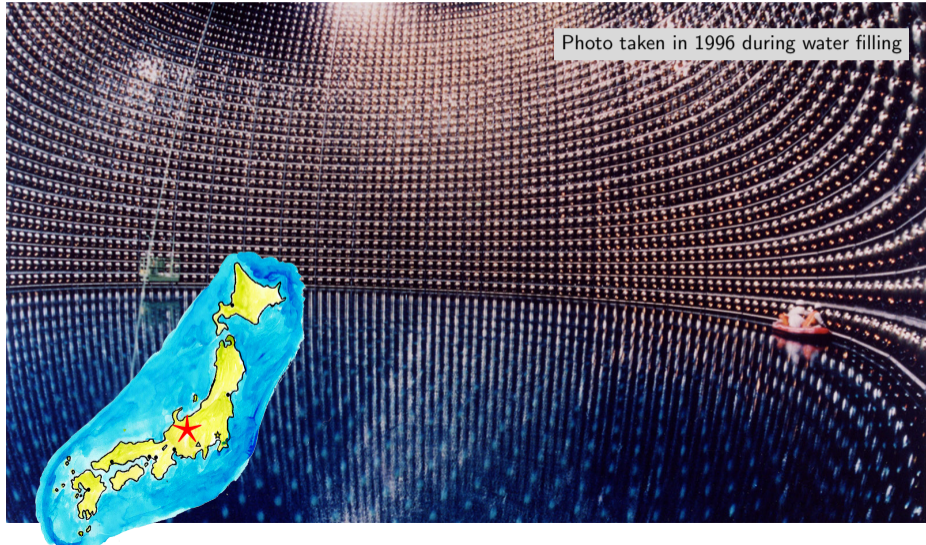
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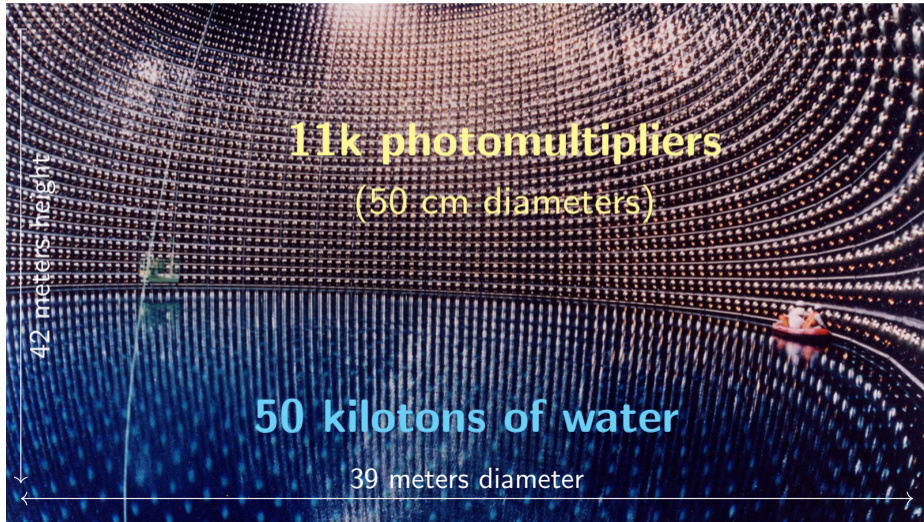
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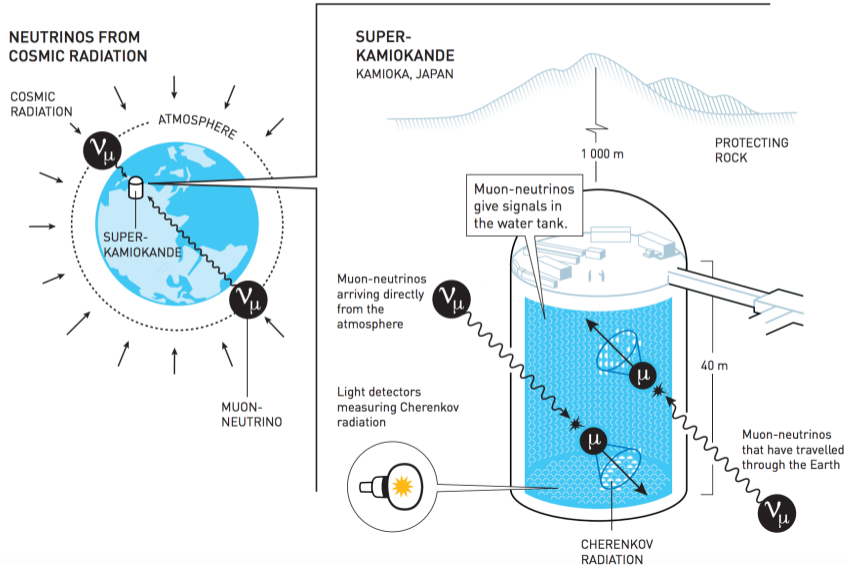
Super-Kamiokande

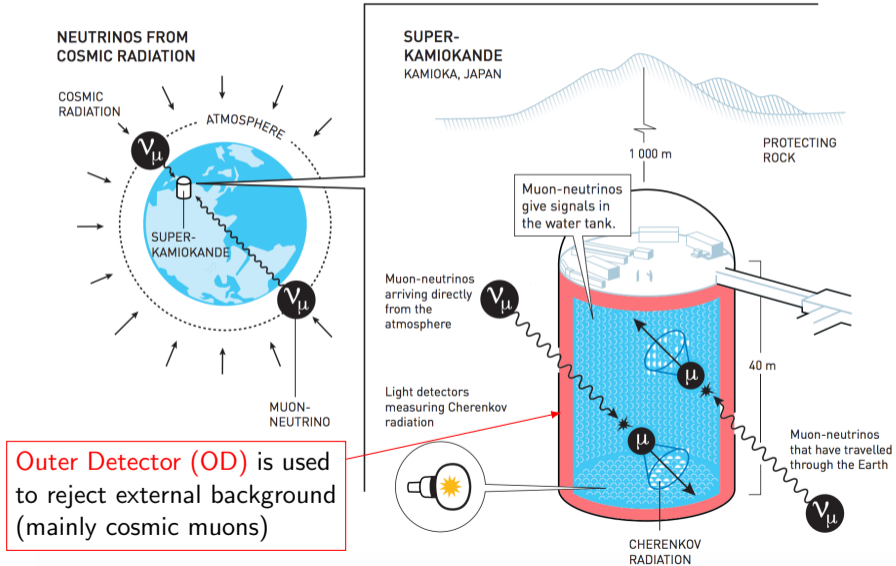
Experiment running since 1998, located in the Mozumi mine in Japan.



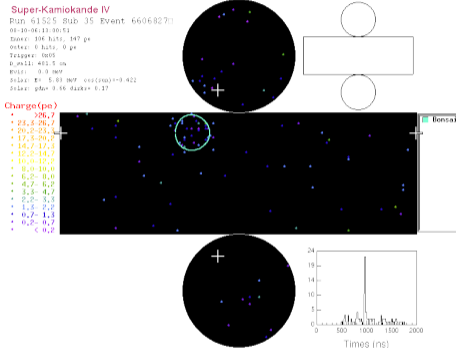
Experiment running since 1998, located in the Mozumi mine in Japan.







Low-energy events



$$E_\nu = 3.5 - 100 \text{ MeV}$$

sparse PMT hits, faint ring



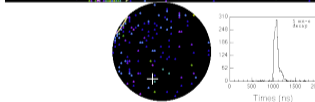
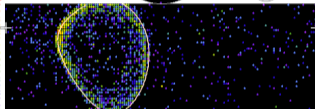
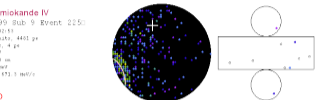
Fully-Contained (FC) events

LOW-ENERGY EVENTS (LE)

$E_\nu \sim 3.5 - 100 \text{ MeV}$

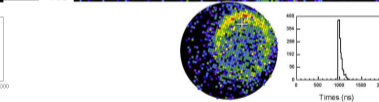
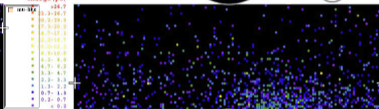
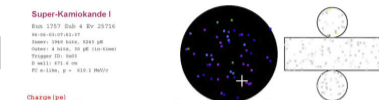
Super-Kamiokande IV
Run: 999999 Sub: 9 Event: 2255
ID: 18-01-07-19:02:53
Date: 1705 bits, 4401 pe
offset: 4 bits, 4 pe
Trigger: 0037
D.wall: 712.3 cm
D.top: 150.7 mtd
wt-Like: p = 473.3 mtd/c

Charge (pe)
• 226.7
• 23.3-28.7
• 20.2-25.3
• 17.3-20.0
• 14.7-17.1
• 12.7-14.7
• 10.6-12.6
• 8.6-10.6
• 6.7-8.6
• 4.7-6.7
• 3.3-4.7
• 2.2-3.3
• 1.5-2.2
• 0.7-1.5
• 0.2-0.7
• 0.0-0.2



Super-Kamiokande I
Run: 1757 Sub: 4 Ev: 25716
ID: 06-06-07-03:197
Date: 2948 bits, 5243 pe
offset: 4 bits, 10 pe (10x-Kami)
Trigger ID: 0000
D.wall: 471.4 cm
D.C = 1146, p = 410.1 Mtd/c

Charge (pe)
• 226.7
• 23.3-28.7
• 20.2-25.3
• 17.3-20.0
• 14.7-17.1
• 12.7-14.7
• 10.6-12.6
• 8.6-10.6
• 6.7-8.6
• 4.7-6.7
• 3.3-4.7
• 2.2-3.3
• 1.5-2.2
• 0.7-1.5
• 0.2-0.7
• 0.0-0.2



$$E_\nu = 0.1 - 10 \text{ GeV}$$

no activity in OD

μ : sharp ring

e: fuzzy ring

INCREASING ENERGY

LOW-ENERGY EVENTS (LE)
 $E_\nu \sim 3.5 - 100 \text{ MeV}$

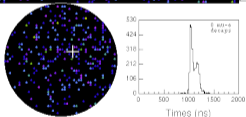
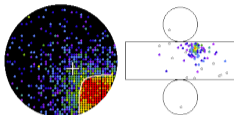
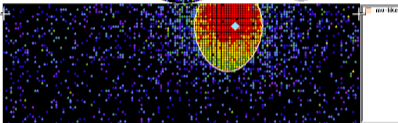
FULLY-CONTAINED EVENTS (FC)
 $E_\nu \sim 0.1 - 10 \text{ GeV}$

Partially-Contained (PC) events

Super-Kamiokande IV
 Run 999999 Sub 6 Event 1170
 18-01-07:13:03:53
 Dtime: 2046 hits, 10482 ps
 otime: 71 hits, 153 ps
 Trigger: mbf
 E_pull: 322.8 cm
 muir: 3.6 sev
 mu-Like, p = 9989.3 MW/c

Charge (pe)

- 206.7
- 23.5-26.4
- 20.2-23.3
- 17.5-20.2
- 14.7-17.5
- 12.2-14.7
- 10.0-12.2
- 8.0-10.0
- 6.2- 8.0
- 4.7- 6.2
- 3.5- 4.7
- 2.2- 3.5
- 1.5- 2.2
- 0.7- 1.5
- 0.2- 0.7
- < 0.2



$E_\nu = 0.1 - 100 \text{ GeV}$
 exit activity in OD

INCREASING ENERGY

LOW-ENERGY EVENTS (LE)

$E_\nu \sim 3.5 - 100 \text{ MeV}$

FULLY-CONTAINED EVENTS (FC)

$E_\nu \sim 0.1 - 10 \text{ GeV}$

PARTIALLY-CONTAINED EVENTS (PC)

$E_\nu \sim 0.1 - 100 \text{ GeV}$

INCREASING ENERGY

Upward-going muons (UPMU)

Super-Kamiokande IV

Point 76304 Sub 27 Event 261998870

17-09-04 14:36:26

Diags: 6559 hits, 32043 ps

outer: 174 hits, 681 ps

Trigger: on1000000

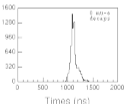
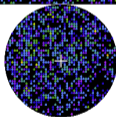
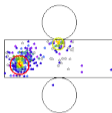
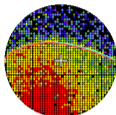
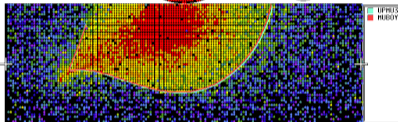
0.9911: 1630.0 cm

Rate: 0.0 sec

Throughgoing sign

Charge (pe)

- 326.7
- 23.5-26.7
- 20.0-23.3
- 17.5-20.8
- 14.7-17.3
- 12.0-14.7
- 10.0-12.0
- 8.0-10.0
- 6.0-8.0
- 4.7-6.0
- 3.5-4.7
- 2.0-3.3
- 1.5-2.0
- 0.7-1.3
- 0.20-0.27
- C 0.2



$E_\nu = 1.6 \text{ GeV} - 10 \text{ TeV}$

entering activity in OD

stopping ($p_\mu > 1.6 \text{ GeV}$) or through-going ($L > 7 \text{ m}$)

LOW-ENERGY EVENTS (LE)
 $E_\nu \sim 3.5 - 100 \text{ MeV}$

FULLY-CONTAINED EVENTS (FC)
 $E_\nu \sim 0.1 - 10 \text{ GeV}$

PARTIALLY-CONTAINED EVENTS (PC)
 $E_\nu \sim 0.1 - 100 \text{ GeV}$

UPGOING MUONS (UPMU)
 $E_\nu \sim 1.6 - 10^5 \text{ GeV}$

INCREASING ENERGY

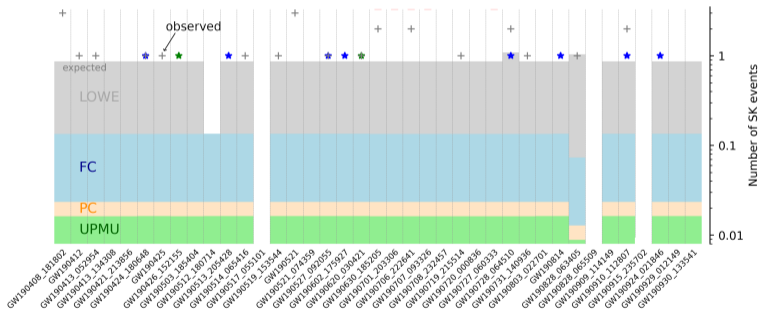
- Four samples
- Energies ranging from MeV to TeV
 - * solar neutrinos
 - * supernova neutrinos
 - * accelerator neutrinos (T2K)
 - * atmospheric neutrinos
 - * astrophysical neutrinos?
- Relatively small size w.r.t. IceCube but wide sensitivity in the lower energy range

Follow-up of GWTC-2 events

- Focus on GWTC-2 catalogue
 - Define a ± 500 s centered on GW time
 - Search for events within this time window, in the four SK samples
 - Compare observation with expected background and extract neutrino flux upper limits
- and compute eventual signal significance by comparing neutrino directions and GW localisation (only for high-energy SK samples)

Low-energy sample	High-energy samples		
	FC	PC	UPMU
Standard solar/SRN selection + 7 MeV energy threshold to ensure stable bkg rate	Standard atmospheric selection		
expected background in 1000 seconds = 0.729	0.112	0.007	0.016

Performed the analysis for the 39 GW in GWTC-2. Three of them were associated to SK downtime (due to calibration) (one less for low-energy due to HV issues).

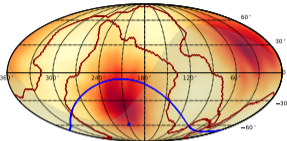


In total:

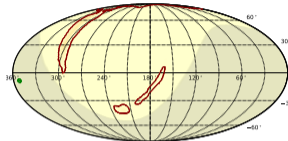
<i>Sample</i>	N_{obs}	N_{exp}
LOWE+	24	24.97
FC*	8	3.95
PC*	0	0.26
UPMU*	2	0.58

No significant excess was observed in the follow-up analysis.

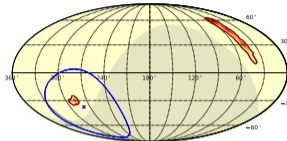
GW190424_180648



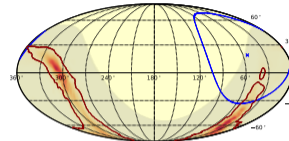
GW190426_152155



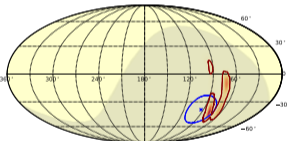
GW190513_205428



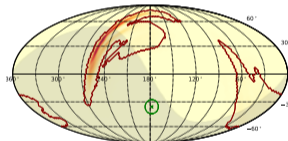
GW190527_092055



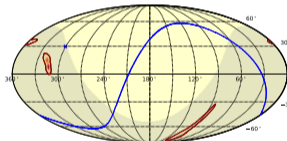
GW190602_175927



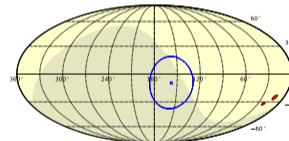
GW190620_030421



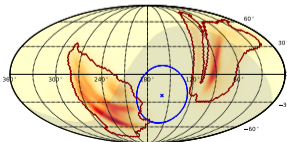
GW190728_064510



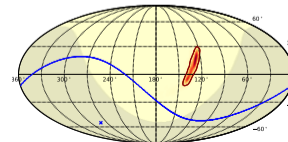
GW190814



GW190910_112807



GW190924_021846



Skymaps in equatorial coordinates

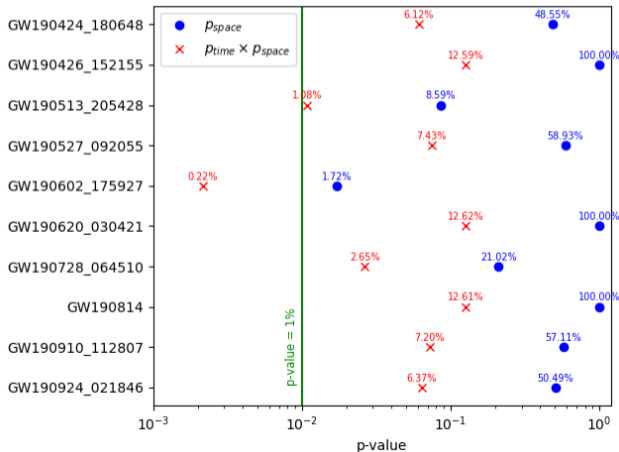
Red: GW localisation and 90% contour

Blue: SK FC events with 1σ angular uncertainty

Green: SK UPMU events.

Shaded area: SK upgoing sky.

Test statistic (TS) has been built to separate signal (point-source) from background (full-sky). It is used to compute p-values (compared observed TS to background distribution).



The most significant GW+ ν coincidence is for GW190602_175927:

$$p = 0.22\%$$

Considering the number of trials ($N = 36$ follow-ups), we get a **post-trial** p-value:

$$P = 7.8\%$$

The neutrino flux is assumed as $\frac{dn}{dE_\nu} = \phi_0 E_\nu^{-2}$ and
 $N_{\text{expected signal}} = \int_{E_{\text{min}}}^{E_{\text{max}}} dE_\nu A_{\text{eff}}^{s,f}(E_\nu, \theta) \times \frac{dn}{dE_\nu}$.

Sample-by-sample flux limits

For each sample and flavour ($\nu_e, \bar{\nu}_e, \nu_\mu, \bar{\nu}_\mu$), we define the flux likelihood:

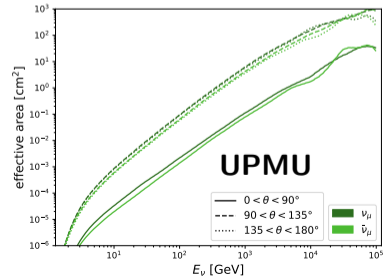
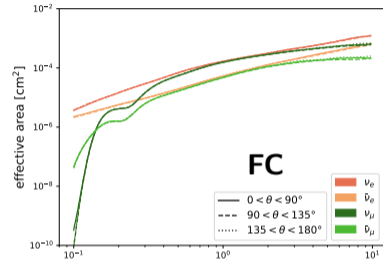
$$\mathcal{L}(\phi_0; n_B, N) = \int \frac{(c(\Omega)\phi_0 + n_B)^N}{N!} e^{-(c(\Omega)\phi_0 + n_B)} \mathcal{P}_{\text{GW}}(\Omega) d\Omega$$

with $c(\Omega) = \int_{E_{\text{min}}}^{E_{\text{max}}} dE_\nu A_{\text{eff}}(E_\nu, \theta) E_\nu^{-2}$ and the 90% U.L on the flux ϕ^{up} is obtained by solving $\int_0^{\phi^{\text{up}}} \mathcal{L}(\phi) d\phi = 0.9$

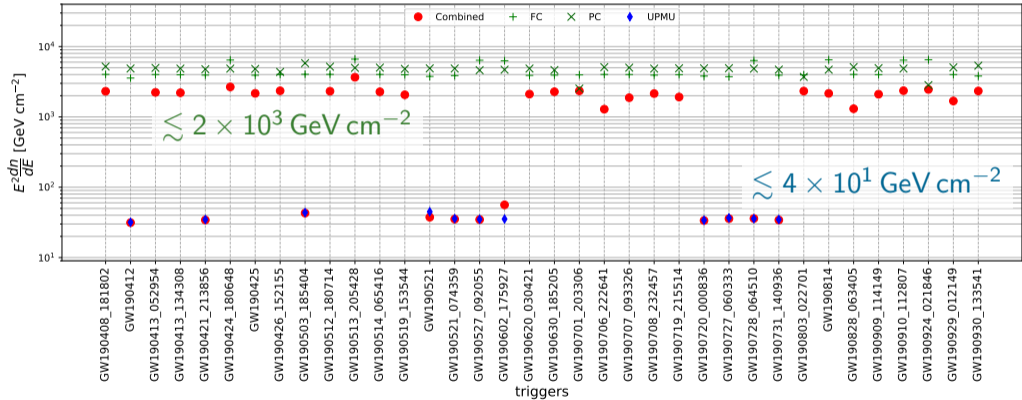
Combined flux limits

Limits combining FC, PC and UPMU are obtained by using the combined TS defined before.

Effective area A_{eff}

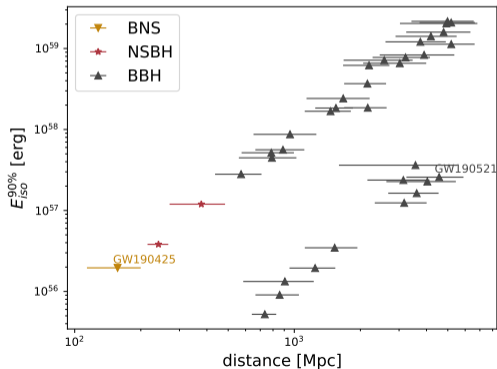
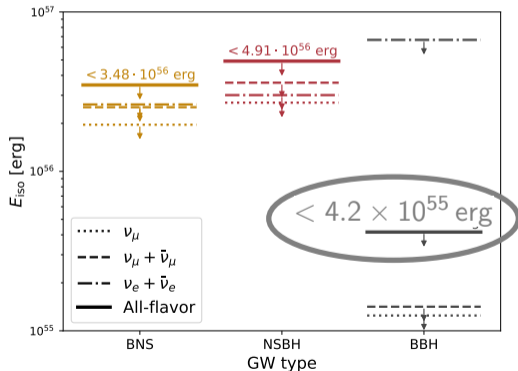


Example of limits for ν_μ flavour:



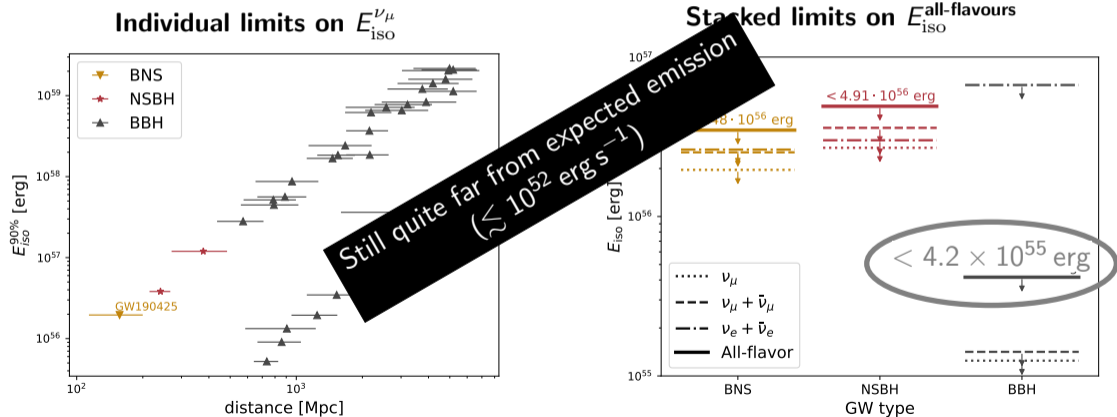
Better limits with the UPMU sample when the GW is below the local horizon. Combined limits are close to the best individual one.

- The total energy in ν from the source (assuming **isotropy**) is $E_{\text{ISO}} = 4\pi d^2 \int \frac{dn}{dE} \times E dE$
 $\Rightarrow E_{\text{ISO}}$ limits obtained by using the 3D localisation skymap from the LVC data release¹.
- We can stack events by nature, assuming same emission (or $E_{\text{ISO}} \propto M_{\text{source}}$).

Individual limits on $E_{\text{ISO}}^{\nu\mu}$ Stacked limits on $E_{\text{ISO}}^{\text{all-flavours}}$ 

¹This is done assuming the flux at Earth is equally distributed between the flavours ($\nu_e : \nu_\mu : \nu_\tau = 1 : 1 : 1$)

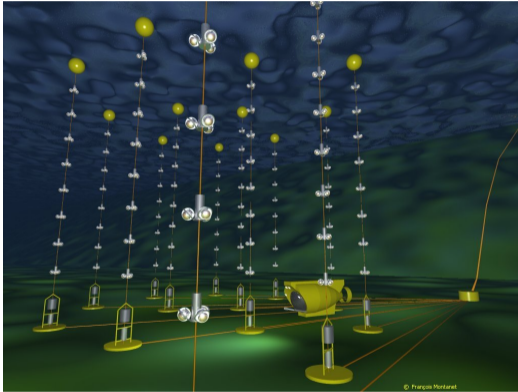
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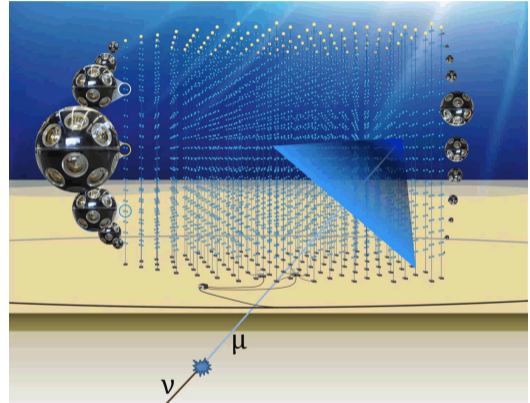
Outlook

ANTARES



- 12 lines instrumenting 10 Mt
- Took data continuously from 2007 to 02/2022

KM3NeT

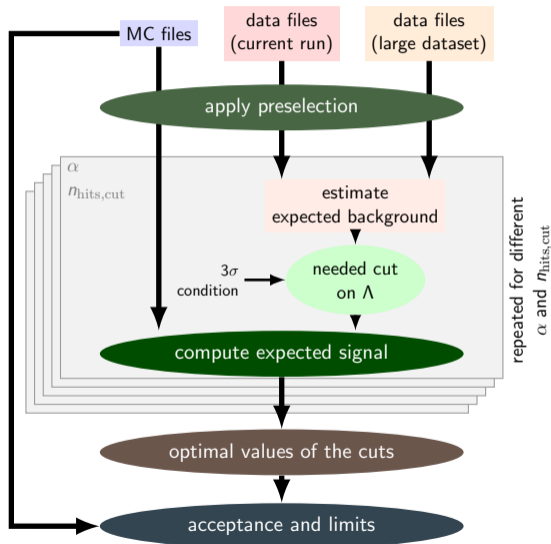


- 3×115 lines instrumenting 7 Mt + 2×0.5 Mt
- Under construction but already taking data with a limited number of lines

Detector is located in a less stable environment with respect to Super-Kamiokande.

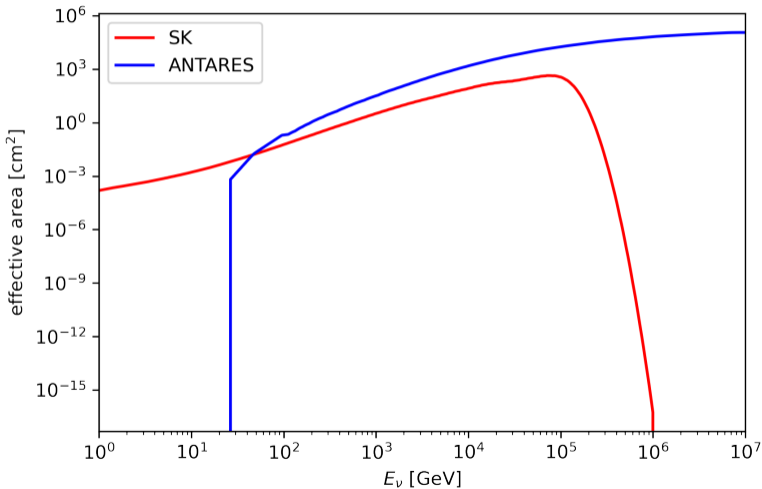
⇒ need **event-by-event treatment**

- Impact of background may fluctuate depending on sea conditions
- Selection cuts need to be optimized to ensure low background
- Detector acceptance will then vary
- Limits are obtained and a stacking analysis can be done as for SK



Complementarity between experiments (1)

Let's consider ANTARES and Super-Kamiokande, and their published effective areas.



What is the expected gain by considering both experiments simultaneously to compute upper limits on $F = \int \frac{dn}{dE} dE$ with $\frac{dn}{dE} \propto E^{-\gamma} e^{-E/E_{\text{cut}}}$?

Simple test with Poisson likelihood (one per experiment and a combined one):

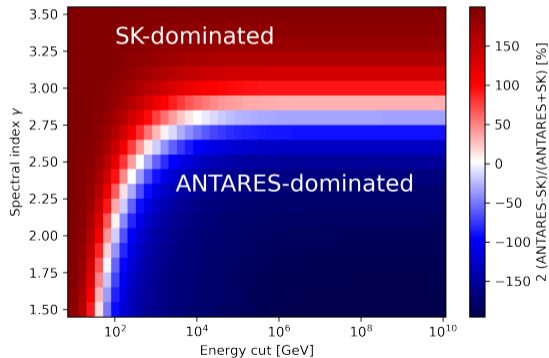


Fig: Relative diff. between ANTARES and SK limits.

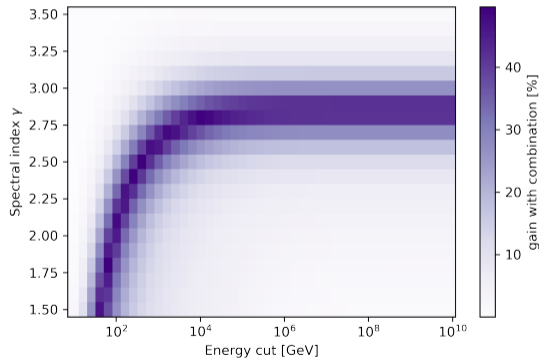
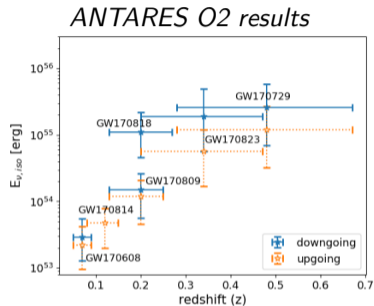


Fig: Relative gain with the combination.

- Currently finalizing follow-up of O3 events with ANTARES (similar strategy to SK's)
- Developing a Python package to handle various detector and allow combinations
→ can be used to investigate further SK+ANTARES
- At the end of Fellini, will start to work within KM3NeT and IceCube, with the aim to expand their low energy selection and make population studies
- Next GW observation run (O4) expected to start in December 2022 / early 2023.



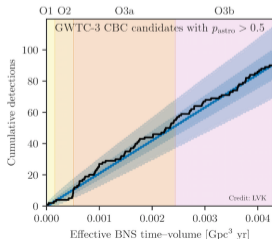
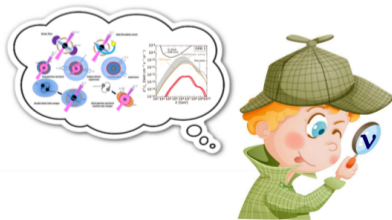
PyJANG package (under development)

JANG

New contract in Louvain (Belgium)

UCLouvain

- 1 Common sources of neutrinos and gravitational waves are expected (in particular from astrophysical objects involving neutron stars), but have not yet been detected.
- 2 SK is sensitive to such neutrinos with energies range from MeV to TeV .
- 3 Search was performed using latest GW catalog, but *no significant excess* has been found.
- 4 There are interesting prospects in existing or under construction experiments:
 - **LOWERING** the energy threshold
 - **ANALYSING** together the different detectors
 - **COMBINING** all sources in catalogs for population studies

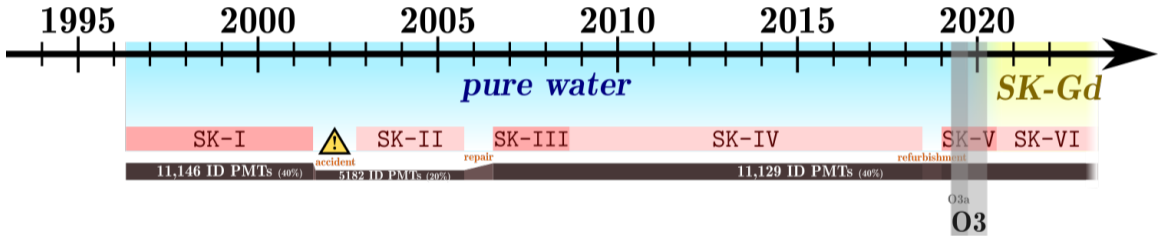


THANK YOU



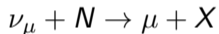
SK collaboration meeting, Toyama, 15/11/2019

Backups



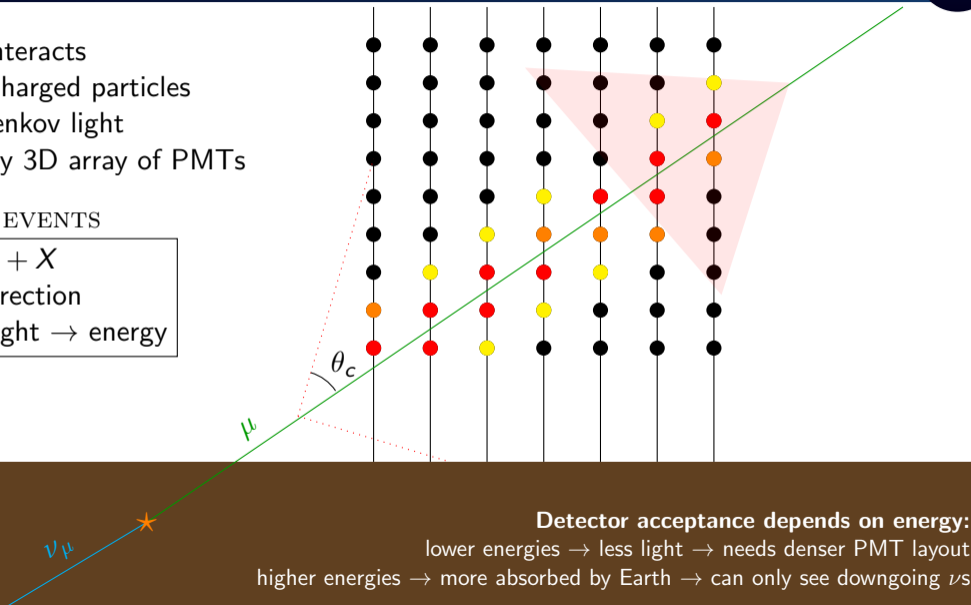
1. Neutrino interacts
2. Produces charged particles
3. Emit Cherenkov light
4. Detected by 3D array of PMTs

TRACK EVENTS



- fit line \rightarrow direction

- amount of light \rightarrow energy



Detector acceptance depends on energy:

lower energies \rightarrow less light \rightarrow needs denser PMT layout

higher energies \rightarrow more absorbed by Earth \rightarrow can only see downgoing ν_s

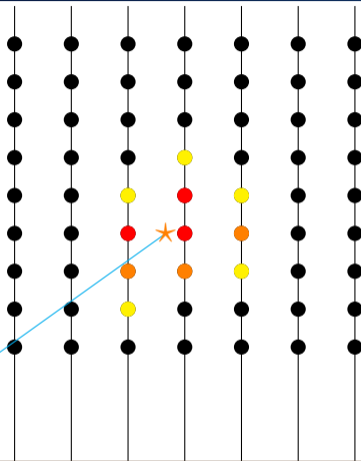
1. Neutrino interacts
2. Produces charged particles
3. Emit Cherenkov light
4. Detected by 3D array of PMTs

SHOWER EVENTS

$\nu_e + \nu_\tau$ charged current interactions

$\nu_e + \nu_\mu + \nu_\tau$ neutral current interactions

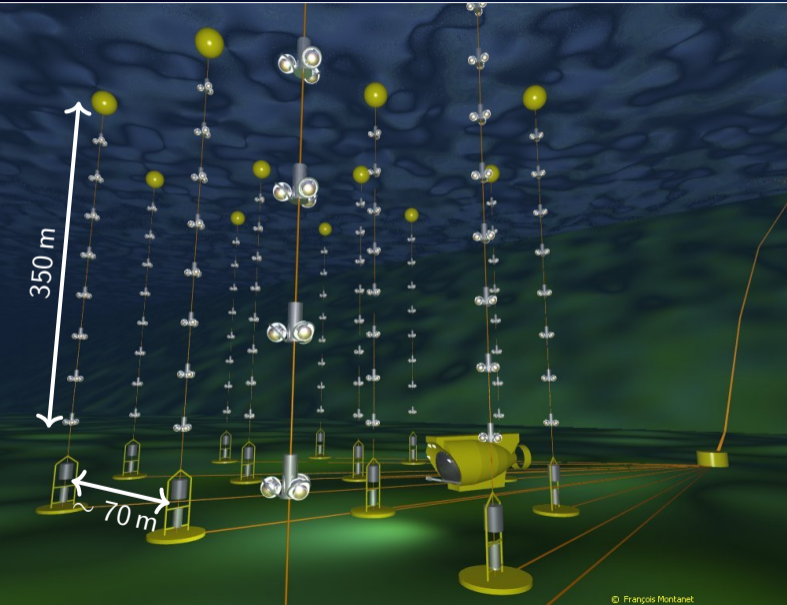
ν_e



Detector acceptance depends on energy:

lower energies \rightarrow less light \rightarrow needs denser PMT layout

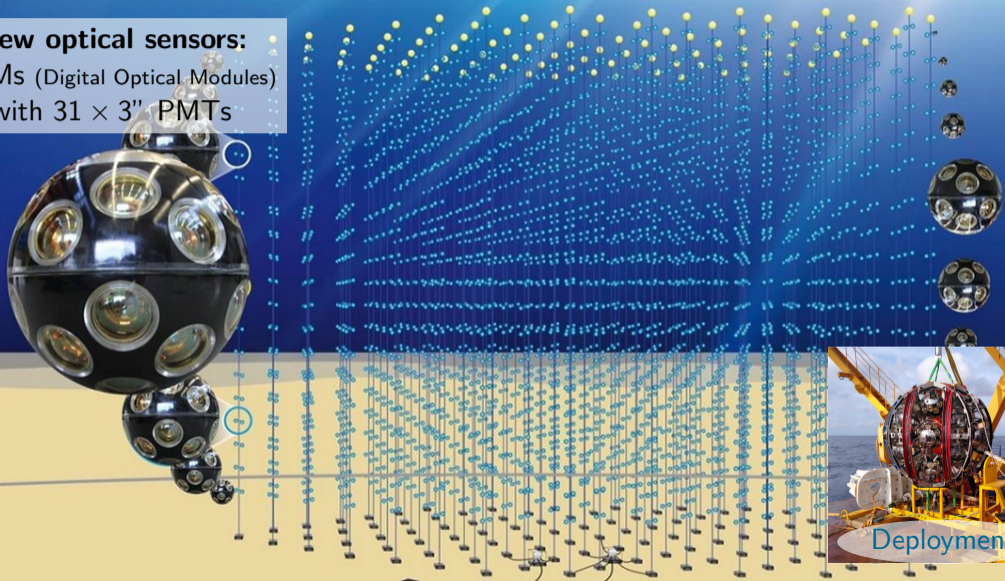
higher energies \rightarrow more absorbed by Earth \rightarrow can only see downgoing ν_s



- In operation since 2006 (completed in 2008)
- Off the coast of Toulon
- 12 lines
- 25 storeys/line
- 3 PMTs / storey

*Total instrumented
volume:
10 Mt*

New optical sensors:
DOMs (Digital Optical Modules)
with $31 \times 3''$ PMTs

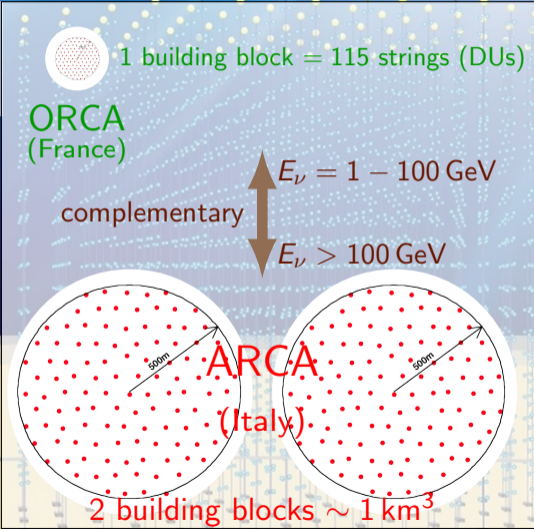


18 DOMs / string



Deployment

New optical sensors:
DOMs (Digital Optical Modules)
with $31 \times 3''$ PMTs



18 DOMs / string



Deployment

How likely the SK observation is associated to background, given time+space correlations?

The p-value can be dissociated in $p = p_{\text{time}} \times p_{\text{space}}$, with:

- $p_{\text{time}} = \text{Prob}(N \geq 1) = 1 - e^{-n_B} \sim 12.6\%$ for $n_B = \text{total background}_{(FC+PC+UPMU)} = 0.13$
- p_{space} is obtained by comparing neutrino direction and GW localisation²
 - For each sample ($k = \text{FC, PC or UPMU}$), define the point-source likelihood $\mathcal{L}_\nu^{(k)}(n_S^{(k)}, \gamma; \Omega_S)$ that separates background from signal ($dn/dE \propto E^{-\gamma}$, direction Ω_S).
 - Compute the maximum log-likelihood ratio Λ (GW localisation \mathcal{P}_{GW} used as prior) and find the source direction Ω_S that maximises it:

$$\Lambda(\Omega_S) = 2 \sum_k \ln \left[\frac{\mathcal{L}_\nu(\widehat{n}_S^{(k)}, \widehat{\gamma}^{(k)}; \Omega_S)}{\mathcal{L}_\nu(n_S^{(k)} = 0; \Omega_S)} \right] + 2 \ln \mathcal{P}_{\text{GW}}(\Omega_S) \text{ and } \boxed{\text{TS} = \max_{\Omega} [\Lambda(\Omega)]}$$

- Compare TS_{data} with the expected background distribution (with $N \geq 1$) to obtain p_{space} .

²IceCube collaboration. IceCube Search for Neutrinos Coincident with Compact Binary Mergers from LIGO-Virgo's First Gravitational-wave Transient Catalog. [Astrophys.J.Lett. 898 \(2020\) 1, L10](#)

For each sample k , we define the likelihood:

$$\mathcal{L}_\nu^{(k)}(n_S^{(k)}, \gamma; \Omega_S) = \frac{e^{-(n_S^{(k)} + n_B^{(k)})} (n_S^{(k)} + n_B^{(k)})^{N^{(k)}}}{N^{(k)}!} \prod_{i=1}^{N^{(k)}} \frac{n_S^{(k)} \mathcal{S}^{(k)}(\vec{x}_i, E_i; \Omega_S, \gamma) + n_B^{(k)} \mathcal{B}^{(k)}(\vec{x}_i, E_i)}{n_S^{(k)} + n_B^{(k)}}$$

where $\mathcal{S}^{(k)}$ and $\mathcal{B}^{(k)}$ are the signal/background p.d.f. (characterizing detector response).

Then, we compute the log-likelihood ratio:

$$\Lambda(\Omega_S) = 2 \sum_k \ln \left[\frac{\mathcal{L}_\nu(n_S^{(k)}, \widehat{\gamma}^{(k)}; \Omega_S)}{\mathcal{L}_\nu(n_S^{(k)} = 0; \Omega_S)} \right] + 2 \ln \mathcal{P}_{GW}(\Omega_S)$$

The final test statistic and p-value are:

$$TS = \max_{\Omega} [\Lambda(\Omega)] \text{ and } p_{\text{space}} = \int_{TS_{\text{data}}}^{\infty} \mathcal{P}_{\text{bkg}}(TS) dTS$$

where $\mathcal{P}_{\text{bkg}}(TS)$ is the expected background distribution.

- **Flux:** We define the following likelihood by using the TS defined before:

$$\mathcal{L}(\phi_0; TS_{\text{data}}, \mathcal{P}_{GW}) = \int \sum_{k=0}^2 \left[\frac{(c(\Omega)\phi_0)^k}{k!} e^{-c(\Omega)\phi_0} \times \mathcal{P}_k(TS_{\text{data}}) \right] \times \mathcal{P}_{GW}(\Omega) d\Omega$$

where $P_i(TS)$ is the distribution of the test statistic assuming the signal consists in i events, assuming E^{-2} spectrum ($dn/dE = \phi_0 E^{-2}$). The 90% upper limit is obtained as above ($\int_0^{\phi_0^{\text{up}}} \mathcal{L}(\phi_0) d\phi_0 = 0.90$).

- **Total energy:** Same for E_{iso} limits:

$$\mathcal{L}(E_{\text{iso}}; TS_{\text{data}}^{(i)}, \mathcal{V}_{GW}^{(i)}) = \int \sum_{k=0}^2 \left[\frac{(c'(r, \Omega)E_{\text{iso}})^k}{k!} e^{-c'(r, \Omega)E_{\text{iso}}} \times \mathcal{P}_k^{(i)}(TS_{\text{data}}^{(i)}) \right] \times \mathcal{V}_{GW}^{(i)}(r, \Omega) d\Omega$$

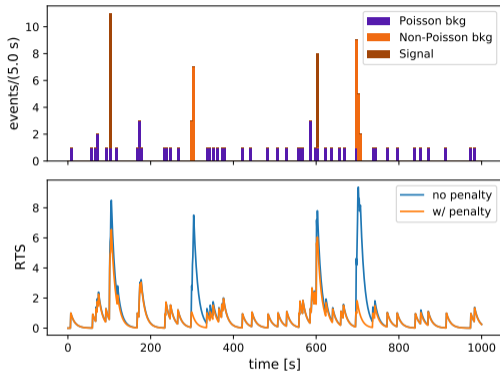
A galactic supernova would create a burst of low-energy neutrinos in any significantly big detector. In the search for such bursts, the usual technique is to:

- define a fixed-width time window (e.g. $w = 20$ s)
- compute multiplicity m of selected “cluster”
- send alert if $m >$ threshold (+ additional checks)

But this does not take benefit of the special time distribution of signal (events very concentrated in time) with respect to background (uniformly distributed in time) and additional checks are needed to remove spatially-correlated background.

⇒ This is critical for small detectors and/or far SNe, where the expected number of signal events is low.

We propose a new approach where the clustering is replaced by a continuous variable that characterises the “signalness” of the data.



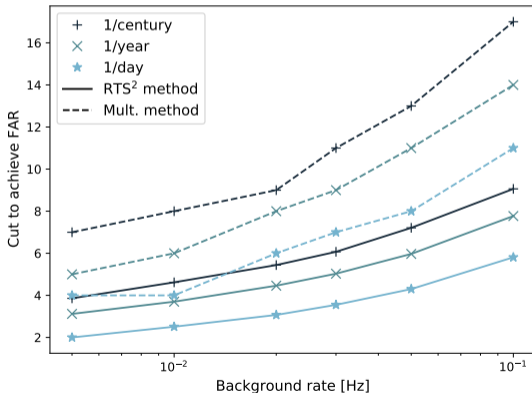
Generated positions:



New RTS² method

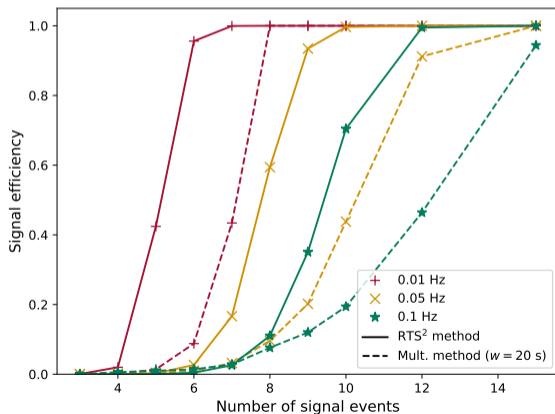
We define $\mathcal{F}(t) = \sum_i f(t; t_i)$ with $f(t; t_i) = e^{-\frac{t-t_i}{T_c}} \times H(t - t_i) \times \prod_{j < i} (1 - e^{-\frac{t_i-t_j}{T_{sp}}} e^{-\frac{\|\vec{x}_i - \vec{x}_j\|^2}{2\sigma_{sp}^2}})$
 where T_c is a chosen time characteristic (one naive choice would be $T_c = 1/\text{rate}$).

Which cut should be applied on \mathcal{F} to achieve a false alarm rate (FAR) of e.g. one per century or one per day?



We can compare the limits on \mathcal{F} to achieve a given FAR (\sim minimum number of signal events) to the multiplicity cut from the standard approach.

In the following, signal is simulation by assuming a simple double-exponential shape: $p(t) = (1 - e^{-t/\tau_1}) \times e^{-t/\tau_2}$ (in the following, $\tau_1 = 10$ ms, $\tau_2 = 1$ s).



- Simulated background is Poisson distributed with a fixed rate r .
- $N \gg 1$ centuries are simulated in order to find the cut \mathcal{C} to apply to get a given false alarm rate.
- Inject signal with fixed multiplicity (on top of background).
- The signal efficiency is the ratio of events passing \mathcal{C} .

- The developed tool may be used for the detection of far SNe (in the local cluster) in Hyper-Kamiokande or other large detectors.
- Similar technique may be exploited for the identification of GeV astrophysical neutrinos in KM3NeT and IceCube:
 - Background: time independent and localised (within a detector)
 - Signal: burst