

# Measuring the muon magnetic anomaly $a_{\mu}$ with the Muon g-2 experiment at Fermilab 

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## What is " $g-2$ " ?

$$
\begin{gathered}
\vec{\mu}_{p}=-g_{p} \frac{e}{2 m_{p}} \vec{S} \\
a_{p}=\frac{g_{p}-2}{2}
\end{gathered}
$$

- $\boldsymbol{g}_{p}$ : proportionality constant between spin and magnetic moment for particle $P$
- $\boldsymbol{a}_{\boldsymbol{p}}$ : magnetic anomaly
- $a_{p}=0$ at tree level (purely Dirac particle)
- Using modern language, the term (g-2)/2 reflects the magnitude of the Feynmann diagrams beyond leading order



## Standard Model determination of $a_{\mu}$



|  | VALUE $\left(\times 10^{-10}\right)$ UNITS |
| :--- | ---: |
| QED $(\gamma+\ell)$ | $11658471.8951 \pm 0.0009 \pm 0.0019 \pm 0.0007 \pm 0.0077_{\alpha}$ |
| HVP $(\mathrm{lo})$ Davier17 | $692.6 \pm 3.33$ |
| HLbL Glasgow | $10.5 \pm 2.6$ |
| EW | $15.4 \pm 0.1$ |
| Total SM Davier17 | $11659181.7 \pm 4.2$ |

Theory Initiative White Paper (arXiv 2006:08443) $a_{\mu}=(116591810 \pm 43) \times 10^{-11} \rightarrow 370 \mathrm{ppb}$

## A rich history of g-2 Theory and Experiment

History of muon anomaly measurements and predictions


Tension between theory and experiment

## The Fundamental Experimental Principle

- Difference between spin precession and cyclotron revolution for a muon (charged particle with spin) in a magnetic field*:

$$
\omega_{a}=\omega_{s}-\omega_{c}=g \frac{e}{2 m} B-\frac{e}{m} B=\frac{g-2}{2} \frac{e}{m} B=a_{\mu} \frac{e}{m} B
$$

*s and $\mathbf{p}$ are assumed to be in a plane perpendicular to $\mathbf{B}$

- simple classical calculation
- the relativistic approach provides the same result



## How do we measure the spin direction?

- Use V-A structure of weak decays to build a polarized beam...

- ... and to measure the muon polarization looking for energetic positrons



## Measuring the spin precession

- The number of observed positrons above a threshold energy oscillates with the $\omega_{a} / 2 \pi$ frequency due to spin precession

- exponential decay modulated by spin precession
- note that the $x$-axis "wraps up" every 100 $\mu$ sec for a total of $\sim 700 \mu \mathrm{~s} \rightarrow \sim 10$ muon lifetimes


## Extracting $a_{\mu}$ (simplified)

$$
\omega_{a}=a_{\mu}(e / m) B \rightarrow a_{\mu}=\omega_{a} / B(m / e)
$$

by expressing $B$ in terms of the (shielded) proton precession frequency:
( $B=\hbar \omega_{p}^{\prime} / 2 \mu_{p}^{\prime}$ ):

$$
a_{\mu}=\frac{\omega_{a}}{\widetilde{\omega}_{p}^{\prime}} \cdot \frac{\mu_{p}^{\prime}}{\mu_{e}} \frac{m_{\mu}}{m_{e}} \frac{g_{e}}{2}=R_{\mu}^{\prime} \cdot \frac{\mu_{p}^{\prime}}{\mu_{e}} \frac{m_{\mu}}{m_{e}} \frac{g_{e}}{2}
$$ measure

$$
R_{\mu}^{\prime}=\frac{\omega_{a}}{\widetilde{\omega}_{p}^{\prime}}
$$

ratio of muon to proton precessions in the same magnetic field
$\widetilde{\omega}^{\prime} p=($ shielded) Proton angular velocity weighted for the muon distribution

## The key ingredients

- Neglecting the corrections, the three key ingredients to measure the muon magnetic anomaly are:

$$
\begin{aligned}
& \omega_{\mathrm{a}}=\text { muon spin precession respect to } \\
& \text { momentum (in B field) } \\
& R_{\mu}^{\prime}=\frac{\omega_{a}}{\widetilde{\omega}_{p}^{\prime}} \sim \\
& \widetilde{\boldsymbol{\omega}}_{\boldsymbol{p}}^{\prime}=\omega_{p}^{\prime}(\boldsymbol{x}, \boldsymbol{y}, \varphi) \cdot \boldsymbol{M}(\boldsymbol{x}, \boldsymbol{y}, \varphi) \\
& \omega_{\mathrm{p}}^{\prime}=\text { proton precession frequency } \\
& \text { M=muon spatial distribution }
\end{aligned}
$$



## Additional corrections

- The ratio $R_{\mu}^{\prime}$ requires additional corrections related to beam dynamics and to magnetic transient fields:

$$
R_{\mu}^{\prime}=\frac{\omega_{a}}{\widetilde{\omega}_{p}^{\prime}} \cdot \frac{1+C_{e}+C_{p}+C_{m l}+C_{p a}}{1+\underbrace{B_{k}+B_{q}}_{\text {Corrections due to beam dynamics }}} \cdot f_{c l o c k}
$$

- $f_{\text {clock }}=$ blinding frequency


## Measuring $\omega_{a}: 5$ parameters fit function

- Fit with simple positron oscillation:

$$
N_{e}(t)=N_{0} \exp \left(-t / \tau_{\mu}\right)\left[1+A \cos \left(\omega_{a} t+\varphi\right)\right]
$$

- This simple fit is clearly not sufficient and well defined resonances are observed in the residuals



## The complete 22 parameters fit function

$\omega_{y,} \omega_{V W}$ vertical oscillations
$\omega_{C B O}, \omega_{2 C B O}$ radial oscillations

$$
\begin{aligned}
& \begin{array}{l}
N_{0} e^{-\frac{t}{\gamma \tau}}\left(1+A \cdot A_{B O}(t) \cos \left(\omega_{a} t+\phi \cdot \phi_{B O}(t)\right)\right) \cdot N_{\mathrm{CBO}}(t) \cdot N_{\mathrm{VW}}(t) \cdot N_{y}(t) \cdot N_{2 \mathrm{CBO}}(t) \cdot J(t) \\
A_{\mathrm{BO}}(t)=1+A_{A} \cos \left(\omega_{\mathrm{CBO}}(t)+\phi_{A}\right) e^{-\frac{t}{\tau_{\mathrm{CBO}}}} \\
\phi_{\mathrm{BO}}(t)=1+A_{\phi} \cos \left(\omega_{\mathrm{CBO}}(t)+\phi_{\phi}\right) e^{-\frac{t}{\tau_{\mathrm{CBO}}}} \\
N_{\mathrm{CBO}}(t)=1+A_{\mathrm{CBO}} \cos \left(\omega_{\mathrm{CBO}}(t)+\phi_{\mathrm{CBO}}\right) e^{-\frac{t}{\tau_{\mathrm{CBO}}}} \\
N_{2 \mathrm{CBO}}(t)=1+A_{2 \mathrm{CBO}} \cos \left(2 \omega_{\mathrm{CBO}}(t)+\phi_{2 \mathrm{CBO}}\right) e^{-\frac{t}{2 \tau \mathrm{CBO}}} \\
N_{\mathrm{VW}}(t)=1+A_{\mathrm{VW}} \cos \left(\omega_{\mathrm{VW}}(t) t+\phi_{\mathrm{VW}}\right) e^{-\frac{t}{\tau_{\mathrm{VW}}}} \\
\qquad N_{y}(t)=1+A_{y} \cos \left(\omega_{y}(t) t+\phi_{y}\right) e^{-\frac{t}{\tau_{y}}} \\
\text { Red = free parameters } \\
\text { Blue }=\text { fixed parameters } \\
\qquad J(t)=1-k_{L M} \int_{t_{0}}^{t} \Lambda(t) d t \text { Lost muons }(\mu \text { hitting } \\
\operatorname{collimators}) \\
\omega_{\mathrm{CBO}}(t)=\omega_{0} t+A e^{-\frac{t}{\tau_{A}}}+B e^{-\frac{t}{\tau_{B}}} \\
\omega_{y}(t)=F \omega_{\mathrm{CBO}(t)} \sqrt{2 \omega_{c} / F \omega_{\mathrm{CBO}}(t)-1} \\
\omega_{\mathrm{VW}}(t)=\omega_{c}-2 \omega_{y}(t)
\end{array}
\end{aligned}
$$

## Final fit to get $\omega_{a}$



## Magnetic field $\vec{B}$ determination

The "trolley"
The "trolley" inside the beam "pipe"
The map


- A Cylinder with 17 NMR probes ("trolley") runs inside the ring every 2-3 days to map the field experienced by muons
- A set of 378 fixed probes, located in 72 azimuthal positions, continuosly measures the field
- Absolute probes for calibration tested at Argonne (ANL) magnet
$\omega_{p}^{\prime} \rightarrow \widetilde{\omega}_{p}^{\prime}:$ muon distribution inside the Ring
- Two tracker stations, made of straw tube modules, placed at $\phi \sim 180^{\circ}$ and $\phi \sim 270^{\circ}$, are used to trace back positrons and get the muon distribution
- Use Beam Dynamics models to extrapolate the distribution all around the ring
- Systematic uncertainties mostly due to Beam Dynamics models used for extrapolation and to tracker alignment


## $\mathrm{a}_{\mu}$ : Unblinding and result

- 462 ppb total uncertainty
- The combined $a_{\mu}$ value shows a $4.2 \sigma$ tension with the standard model 2020 prediction in the Theory Initiative
Group White $\begin{array}{lllllllll}17.5 & 18.0 & 18.5 & 19.0 & 19.5 & 20.0 & 20.5 & 21.0 & 21.5\end{array}$ Paper



## The Theory Initiative

- ~130 physicists collaborated for 3 years [June 2017 - June 2020] in seven workshops to produce a reference number for am to be used by FNAL g-2 experiment as a benchmark

Muon g-2 Theory Initiative defines SM benchmark value that our collaboration will use for comparison. We don't "pick and choose" other individual results.

Group photo from the Seattle workshop in September 2019, https://indico.fnal.gov/event/21626/


## Organizers:

Aida El-Khadra Martin Hoferichter DWH

| Contribution | Section | Equation | Value $\times 10^{11}$ | References |
| :--- | :--- | :--- | ---: | :--- |
| Experiment (E821) |  | Eq. $(8.13)$ | $116592089(63)$ | Ref. [1] |
| HVP LO $\left(e^{+} e^{-}\right)$ | Sec. 2.3 .7 | Eq. $(2.33)$ | $6931(40)$ | Refs. [2-7] |
| HVP NLO $\left(e^{+} e^{-}\right)$ | Sec. 2.3 .8 | Eq. $(2.34)$ | $-98.3(7)$ | Ref. [7] |
| HVP NNLO $\left(e^{+} e^{-}\right)$ | Sec. 2.3 .8 | Eq. $(2.35)$ | $12.4(1)$ | Ref. [8] |
| HVP LO (lattice, $u d s c)$ | Sec. 3.5 .1 | Eq. $(3.49)$ | $7116(184)$ | Refs. [9-17] |
| HLbL (phenomenology) | Sec. 4.9 .4 | Eq. $(4.92)$ | $92(19)$ | Refs. [18-30] |
| HLbL NLO (phenomenology) | Sec. 4.8 | Eq. $(4.91)$ | $2(1)$ | Ref. [31] |
| HLbL (lattice, $u d s)$ | Sec. 5.7 | Eq. $(5.49)$ | $79(35)$ | Ref. [32] |
| HLbL (phenomenology + lattice) | Sec. 8 | Eq. $(8.10)$ | $90(17)$ | Refs. [18-30,32] |
| QED | Sec. 6.5 | Eq. $(6.30)$ | $116584718.931(104)$ | Refs. [33,34] |
| Electroweak | Sec. 7.4 | Eq. $(7.16)$ | $153.6(1.0)$ | Refs. [35,36] |
| HVP $\left(e^{+} e^{-}\right.$, LO + NLO + NNLO) | Sec. 8 | Eq. $(8.5)$ | $6845(40)$ | Refs. [2-8] |
| HLbL (phenomenology + lattice + NLO) | Sec. 8 | Eq. $(8.11)$ | $92(18)$ | Refs. [18-32] |
| Total SM Value | Sec. 8 | Eq. $(8.12)$ | $116591810(43)$ | Refs. [2-8, 18-24,31-36] |
| Difference: $\Delta a_{\mu}:=a_{\mu}^{\text {exp }}-a_{\mu}^{\text {SM }}$ | Sec. 8 | Eq. $(8.14)$ | $279(76)$ |  |

## The lattice evaluation of $a_{\mu}$

g-2

- The BMW (Budapest, Marseille, Wuppertal) collaboration published the result of their lattice calculation in april 2021 (preprint in 2020): closer to the experimental value than with the Theory Initiative one

- The Theory Initiative evaluated the (leading order) hadronic contribution to $a_{\mu}$ using the dispersion integral:

$$
a_{\mu}^{H L O}=\frac{1}{4 \pi^{3}} \int_{m_{\pi}^{2}}^{\infty} K(s) \sigma^{h a d}(s) d s \text {; with } K(s) \propto \frac{1}{s}
$$

- which is based on the experimental measurement of the cross section $\sigma^{\text {had }}(s) \Rightarrow \sigma\left(e^{-} e^{+} \rightarrow\right.$ hadrons $)$


## News on lattice calculations (ICHEP 2022)

- Use the
"Intermediate region", more under control
- This region corresponds to ~235/700=33\% of the full hadronic contribution

2022 news: lattice \& R-ratio results in time/energy windows
Covering whole region at high precision is challenging for lattice calculations
$\rightarrow$ but individual time/energy regions are more accessible
lattice: time windows


R -ratio: energy windows


- It (approximately) corresponds to the energy window 1-3 GeV
- Note: the largest contribution to $a_{\mu}^{\text {had }}$ still comes from the $\rho$ resonance region, i.e below 1 GeV


## New results from ICHEP22

- Recent lattice results are in agreement with BMW calculation in the intermediate region
- strong tension with the value obtained from $e^{+} e^{-}$data
intermediate window: 2 new lattice results (ETMC-22 \& CLS/Mainz-22)


$$
\begin{aligned}
& \text { - strong tension with } a_{\mu}^{W} \text { (HVP-LO) results driven by experimental } e^{+} e^{-} \text {data: } \\
& \text { at } \sim 4.2 \sigma_{\text {combined }} \text { if WP-proc.('22) (2205.12963, Colangelo et al.), see light-red band, is used } \\
& \text { at } \sim 5.8 \sigma_{\text {combined }} \text { if KNT('19-'22) }(1911.00367+\text { private comm. ), see dashed lines, is used }
\end{aligned}
$$

ICHEP 2022

## Hadronic cross section

- Is it possible a "mistake" in $\sigma^{\text {had }}(s)$ ?
- In principle yes (experimental cuts) but
- many different experiments should have same bias
- an upward shift in $\sigma^{\text {had }}(s)$ induces an increase of $\Delta \alpha^{\text {had }}\left(M_{Z}\right)$

$$
\begin{gathered}
a_{\mu}^{H L O}=\frac{1}{4 \pi^{3}} \int_{m_{\pi}^{2}}^{\infty} K(s) \sigma^{h a d}(s) d s ; \text { with } K(s) \propto \frac{1}{s} \\
\Delta \alpha^{h a d}=\int_{m_{\pi}^{2}}^{\infty} g(s) \sigma^{h a d}(s) d s ; \text { with } g(s) \propto \frac{M_{Z}^{2}}{M_{Z}^{2}-s}
\end{gathered}
$$

- Similar dispersion integral with a different kernel function


## Hadronic cross section and fine structure constant

- An upward shift in the hadronic cross section to move the theoretical prediction towards the lattice one, or towards the exp. value, also "pushes down" the preferred Higgs Mass value in the global electroweak fit (green band)



# Clearly there is something going on there. 

## Our task: reduce as much as possible the experimental uncertainty on $g$-2!

## FNAL g-2: next publications

Collected data: 18.8
BNL equivalent
Publications:

- Run1

7/apr/2021:
$\pm 462$ ppb

- Run2,3
spring 2023:
$\pm 220 \mathrm{ppb}$
- Run4,5,6
beginning 2025:

Last update: 2022-06-27 23:28; Total = 18.80 (xBNL)

$\pm 140 \mathrm{ppb}$ (maybe
$\pm 120 \mathrm{ppb})$

## Main improvements wrt Run1



- Thermalization of experimental Hall
- Kicker HV to design value



## Main improvements wrt Run1

- Thermalization of experimental Hall
- Kicker strength to design value
- New Radio Frequency System mounted on quadrupoles which reduces Beam Betatron oscillations




## Perspectives on final uncertainty Run 1-5

Total Uncertainties

|  | Run-1 | Run-2/3 | Run-4/5 | Run-1/3 | Run-1/5 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\omega_{a}$ (stat) | 434 | 202 | 119 | 183 | 102 |
| $\omega_{a}$ (syst) | 109 | 70 | 60 |  |  |
| $\omega_{p}^{\prime}$ (tot) | 114 | 65 | 62 |  |  |
| $R_{\mu}$ (tot syst) | 157 | 95 | 86 | 106 | 92 |
| Total | 462 | 223 | 147 | 212 | 136 |

- The final TDR goal of 140ppb is at reach!
- And we can do more with Run 6 (next slide)..


## Run 6

- In August 2022, the new Fermilab Director Lia Merminga confirmed a Run 6 with MU+ in Nov22-Jul23
- The priority, however, will be given to machine development for MU2E experiment
- Aim to reduce systematics, rather than to increase statistics

$$
\sigma^{t o t}=136 p p b \quad \underline{R_{\mu} \text { uncertainty vs. n. of Run } 6}
$$



## Conclusions

- There is clearly something going on in the muon anomalous magnetic moment
- Lots of theoretical and experimental activity to pin down the hadronic contribution to $a_{\mu}$
- The "Muon g-2" experiment has collected data to reduce by a factor of 4 the published uncertainty on $a_{\mu}$, thus reaching the TDR goal, and it has the potentiality to even gain some 1020\% additional reduction in Run6
- This will be the more precise measurement of the muon anomaly for several decades, thus every possible reduction of the finale uncertainty should (and will) be pursued in the next 2-3 years

