NEN 500P2022

The 2nd INFN School on Underground Physics: Theory & Experiments

New Photodetectors

Ettore Segreto Universidade Estadual de Campinas - UNICAMP (Brazil)

SoUP 2022 – Laboratori Nazional del Gran Sasso – 24th June 2022





Outline

- Introduction
- Vacuum Photosensor *YPMTs, MaPMTs, LAPPD, ...*
- Solid State detectors *` APD, SiPM*
- Hybrid Photosensors
- Estimating the Light Yield of a detector

Light detection in particle physics

- Light detection played a central role in particle physics in the last century
- Light production is **one of the mechanisms through which the energy deposited by a charged particle is transformed spontaneously into a detectable signal** (as ionization, phonons, acoustic, heat,...)
- Light signal can be produced through **different processes**: scintillation, Cherenkov radiation, transition radiation, ...
- Light signals carry a lot of information about the incoming particle: time of arrival, type, direction, energy loss (related to dE/dx and LET)
- Enough reasons to justify the enormous effort that experimental particle physicists have done to detect them with high efficiency to extract as much as information as possible







Determined by the characteristics of the incoming particle:

- Depostied energy (or better LET linear energy transfer): the number of emitted photons is usually proportional to the energy deposted in the medium (unless quenching process are in act)
- Particle type: in many scintillators the emission of light happens throught the *transition of an excited molecule from the first excited state to ground state*. Depending on the total spin of the excited state (triplet or singlet) *the deexcitation can be fast or slow*. The abundance of fast and slow components in the light signals depends on the particle type
- Direction: Especially in Cherenkov detectors ==> The shape of the Cherenkov rings give the direction of the particle



This is the transport/propagation of the procuced photons to the active photo-sensors and comprises:

- Attenuation of the signal due to absortion of the radiator or of impurities present on the radiator
- Rayleygh elastic scattering
- *Wavelght shifting of the photons* to mach the sensitivity curve of the photosensors
 - Can be done diluting optically active compound in liquid or solid scintillators or depositing shifting compounds on the photo-sensors and/or on the inner surface of the detector



LIGHT PRODUCTION

The conversion of photons into an electrical signal that we measure:

- It does not exist the perfect photodetector but the most suited for our application
- The most suited photon detector is chosen on the basis of the *information we want to extract from the light signal, from the characteristics of the light signal (wavelenght, shape, intensity,..) itself, geometrical constraints for the photosensor, cost,*





Spinthariscope

- Invented by William Crookes in 1903!
- Allows to observe the scintillation light emitted from a thin foil of zinc sulfide with a lens in a dark little box
- Crookes named his device from Ancient Greek: σπινθήρ (spinthḗr) "spark"
- Photon detector is human eye...



Beautiful story of this device on wikipedia: https://en.wikipedia.org/wiki/Spinthariscope

Photomultipliers tube

- The *first demonstration of a photomultiplier* (PMT) dates back to 1934. The floor for the invention of PMTs had been set up by the discovery (H. Hertz, 1887) and theoretical explanation (A. Einstein, 1905) of the **photoelectric effect** and by the discovery of the **secondary electron emission** (Villard, 1902)
- PMTs were a **sort of spin-of**f of the technology developed to build the **first practical television cameras** (iconoscope, orthicon...)
- H. lams and B. Salzberg realized the first prototype of PMT integrating a photocathode and a single stage of multiplication. The tube had a gain of 8 and operated at frequencies above 10 kHz!



PMT scheme



PMT basic principles: window



WAVELENGTH (nm)

PMT basic principles: photocatode

Photoelectric emission and cathode Quantum efficiency

$$QE(\lambda) = (1-R)P_V \frac{1}{1+\lambda_a/\lambda_e}P_{escape}$$

ALKALI PHOTOCATHODE





PMT ba

Photoelectric emissi Quantum efficiency

 $QE(\lambda) = (1-R)I$

ALKALI PHOTOCATHODE



Curve code	Photocathode material	Window material			Spectral response			
			sensitivit (Typ.)	Spectral range	Peak wavelength			
					Radiant sensitivity		Quantum efficienc	
			(µA/Im	(nm)	(mA/W)	(nm)	(%)	(nm)
100M	Cs-I	MgF2	—	115 to 200	14	140	13	130
200S	Cs-Te	Silica	—	160 to 320	29	240	16	210
200M	Cs-Te	MgF2	—	115 to 320	29	240	17	200
201S	Cs-Te	Silica	—	160 to 320	31	240	17	210
400K	Bialkali	Borosilicate	95	300 to 650	88	420	27	390
400U	Bialkali	UV	95	185 to 650	88	420	27	390
400S	Bialkali	Silica	95	160 to 650	88	420	27	390
401K	High temp. bialkali	Borosilicate	40	300 to 650	51	375	17	375
402K	Low noise bialkali	Borosilicate	40	300 to 650	54	375	18	375
500K	Multialkali	Borosilicate	150	300 to 850	64	420	20	375
500U	Multialkali	UV	150	185 to 850	64	420	25	280
500S	Multialkali	Silica	150	160 to 850	64	420	25	280
501K	Extended red multialkali	Borosilicate	200	300 to 900	40	600	8	580
502K	Multialkali	Borosilicate (prism)	230	300 to 900	69	420	20	390
600K	GaAsP(Cs)	Borosilicate	700	280 to 720	180	550 to 650	40	480 to 53
601K	Extended red GaAsP(Cs)	Borosilicate	750	280 to 820	160	550 to 650	36	480 to 53
602K	GaAs(Cs)	Borosilicate	700	370 to 920	85	750 to 850	12	600 to 75
700K	Ag-O-Cs	Borosilicate	20	400 to 1200	2.2	800	0.36	740
900S	InP/InGaAsP(Cs)	Silica	—	950 to 1200	18	1100	2	1000 to 110
901S	InP/InGaAs(Cs)	Silica	—	950 to 1700	24	1500	2	1000 to 155
440K	Super bialkali	Borosilicate	105	300 to 650	110	400	35	350
441K	Ultra bialkali	Borosilicate	135	300 to 650	130	400	43	350
442K	Super bialkali	Borosilicate	105	230 to 700	110	400	35	350
443K	Ultra bialkali	Borosilicate	135	230 to 700	130	400	43	350
444K	Extended green bialkali	Borosilicate	160	300 to 700	127	420	40	380

Transmission mode photocathodes





PMT basic principles: dynodes



- Photoelectrons from the cathode are focused on the first multiplication stage (dynode)
 If multiplication for the stars of an a dynamic stars of a star
- If multiplication factor of one dynode is $\delta =>$ overall gain is δ^n , where n is the number of dynodes



PMT basic principles: gain



•
$$G = \delta^n = (aE^k)^n$$

PMT basic principles: timing





PMT basic principles: timing

Unit: ns

Dynode type	Rise time	Fall time	Pulse width (FWHM)	Electron transit time	T.T.S.
Linear-focused	0.7 to 3	1 to 10	1.3 to 5	16 to 50	0.37 to 1.1
Circular-cage	3.4	10	7	31	3.6
Box-and-grid	to 7	25	13 to 20	57 to 70	Less than 10
Venetian blind	to 7	25	25	60	Less than 10
Mesh	2.5 to 2.7	4 to 6	5	15	Less than 0.45
Metal channel	0.65 to 1.5	1 to 3	1.5 to 3	4.7 to 8.8	0.4

T.T.S.: Transit Time Spread

PMT dark current and afterpulses



- **Region a** dominated by *leakage current*
- **Region b** dominated by *thermoionic emisison from the cathode*
- **Region c** dominated by field emission and scintillation of the materials
- Region b in general offers the best S/N
- Afterpulsing is a correlated noise, which manifests itself as delayed pulses
- Produced by the ionization of the residual gas



Photomultipliers: pros and cons



- Well known and mastered technology
- Good detection efficiency (10 – 30%)
- Needs high voltage (1000 3000 Volt)
- Typically large devices which occupies a big volume
- Not so easy to make them radiopure

PMTs in Astroparticle: LVD



- Observaotry for Supernova neutrino Burst operating since 1992
- **840 scintillator counters** of 1,5 m³ in three independent towers
- Each counter viewed by three 15 cm PMTS => FEU49b or FEU125
- Detects (anti) neutrino through inverse beta decay => two correlated signals: a prompt one from emission and annihilation of the positron and a delayed neutron absorption





PMTs in Astroparticle: LVD



the Inverse Beta Decay:

$$\overline{v_e} + p \rightarrow n + e^+$$

The visible energy for this interaction:

$$E_{vis} = E_{\overline{v_e}} - 0.789 MeV$$



PMTs in Astroparticle: Borexino



- Detects *solar neutrinos* through elastic neutrino scattering on electrons.
- Graded shielding form outside to inside
- Inner scintillator mass (pseudocumene+PPO) is ~300t

```
2200 ETL9351 => 8" PMTs
High gain => 10<sup>7</sup>
Short TTS=> 1 ns
```





PMTs in Astroparticle: DarkSide50



- Experiment for the *direct detection of Dark Matter*
- Dual phase depleted Ar (from ³⁹AR) TPC
- 38 PMTs in the inner LAr detector=> Hamamatsu 3" R11065 crygenic
- High quantum efficiency > 35%

ETL 9351

- Very low radioactive background
- 110 Hamamatsu R5912 (8") in the liquid scintillator of the veto
- 80 ETL 9351 8" PMTs in the water Cherenkov muon veto.



PMTs in Astroparticle: XenonNT



- Dual phase Xe experiment for direct Dark Matter detection
- XenonNT uses 8 ton of Lxe readoout by 494 PMTs
- Hamamatsu 3" R11410-21
- Cryogenic
- High QE ~ 35% -UV sensitive
- Ultra low background



KM3NeT multi-PMT optical module



Undrsea experiment for atmospheric and high energy neutrinos: two sites in the Mediterranean Sea

Optical module is made of a **glass sphere containing 31 3**" **PMTs**, *calibration and and readout electronics*

- A Photocathode area of about 1300 cm² three times the area of a single 10" PMT,
- Uniform angular coverage
 - Sensitivity to the incoming direction of detected photons
- Good photon counting performance;
- Possibility to define local triggers

 (implemented onshore) based on the
 pattern of PMT signals => background
 soppression
 30

ICECUBE mDOM





The *ICECUBE upgrade* foresee the installation of ~400 new multiMpmt digital optical modules

- amera 24 small (3") Hamamatsu R15458-02 PMTs arranged in a spherical geometry
 - Large optical comerage
- Homogeneus solid angle coverge
 - Allows to access directional information
 - Implement multiplicity triggern inside oone module

Multi Anode PMT PHOTOCATHODE Electron multiplication in narrow, independet channel FOCUSING MESH • The read-out is performed with a **segmented anode** Position sensitive Tolerance to modest B field METAL CHANNEL DYNODES Exist in many anode segmentation. Dimensions ~5 cm x 5 cm High gain, high QE • Limited cross talk <2% MULTIANODE *MicroPMT*: very small, flat single channel devices Micro PMT internal structure SECONDARY ELECTRONS CONNECTION TERMINA FOCUSING ELECTRODE VACUUN PHOTOCATHODE ANODE ANODE DIRECTION ELECTRON MULTIPLIER (DYNODES LAST DYNODE PHOTOFI ECTRO GLASS ELECTRON MULTIPLIER (DYNODE **TOP VIEW** [SECTIONAL VIEW]

Multi Anode PMTs: Applications

AMS Rich detector

10,880

photosensors

- ----

- *RICH detectors* (HERA-B first with MaPMTS)
- Positron Emission Tomography

Particle

Intensity $\Rightarrow Z^2$

Z = 13 (AI)

= 9.148 TeV/c

Header Level1 Level2

Aerogel



LAPPD Large Area Picosecond Photodetector



- QE ~20%
- Time resolution < 100 ps for single photon events
- **Space resolution < 1 mm** in both directions
- Gain 10⁷ @ 1 kV
- Dark noise ~20 Cts/s/cm²

- Large area (20 cm X 20 cm) vacuum device
- Fused silica window with MultiAlkali photochatode on the inside
- Two strage of electron multiplication through microchannel plates
- **Strip anode** =>spatial reconstruction of the
 - hit



LAPPD Large Area Picosecond Photodetector

Micro Plates Channel



New generation large-area high performance MCPs - enabled by two technological breakthroughs

Produce *large blocks of low cost, hollow, glass capillary arrays with micron-sized* pores (Incom Inc). No need of chemical etching

Atomic layer deposition coating methods to impart the necessary resistive and secondary emission properties for high electronic gain and robust performance

LAPPD: ANNIE detector

- Accelerator Neutrino Neutron Interaction Experiment
- 26-ton Gd-doped water Cherenkov Front Muon Veto detector installed in the BNB at (FMV) Fermilab
- Study the neutron multiplicity in CC neutrino-nucleus interaction in water
- Demonstrate the use of LAPPD

Identifying and counting final state neutrons is important to understand and reduce the systematic uncertainties of the neutrino energy reconstruction in oscillation experiments


LAPPD: ANNIE detector

- Fundamental measuring with precision the interaction vertex of the neutino to fiducialize the active volume => maximize the capture prob. of neutron by Gd after thermalization
- Using just a small coverage of LAPPD, complementing standard PMT (5 over 128) significantly improves the localization of the vertex (factor 3)





Silicon devices: PN junction



KAPDC0071EB

Members of a larger family with PIN diode and APD (Avalanche Photo Diode). They differ for the amplification mechanism

Hamamatsu handbook: https://www.hamamatsu.com/resources/pdf/ssd/mppc_kapd9005e.pdf³⁸

APD

Gain



- **Depletion layer can be deepened** biasing the PN junction **in reverse**
- Increases the electric force on charge carriers => more kinetic energy between collisions
- If E is high enough => mean energy of carriers can exceed the silicon band gap energy between collisions => ionize lattice atoms upon impact and release at least another electron-hole pair into conduction and valence bands per impact

$$G=2+2P^2+2P^3+...=\sum P^i=\frac{1}{1-P}$$

Geiger Mode APD (GAPD)



Silicon Photomultipliers equivalent circuit



single cell

Silicon Photomultipliers (SiPM)



Gain and Photon Detection Efficiency



KAPDB0141EB

Cross talk and afterpulses



- During avalanche process, kinetic energy of avalanching carriers can be lost through photon ewmission
- Photons can travel to the neighboring pixels and initiate avalanches

- During avalanche process a small fraction of avalanching carriers can be *trapped in impurityenergy levels*
- Released after short delays

Cross talk and afterpulses





Silicon Photomultipliers – pros and cons

- Low bias voltage O(50 Volt)
- Digital detectors
- High gain → Single Photo electron reconstruction capability
- Small amount of material => Good for low background experiments
- But small active area O(1 cm²) each
- High Dark count rate at room temperature
- Need of ganging schemes to reach the coverage of a standard PMT
- and/or coupling with passive photon collectors

Example: DUNE photon detection system



- 1. A high-power, wide-band neutrino beam (~ GeV energy range).
- 2. A ≈ 40 kt liquid-argon Far Detector in South Dakota, located 1478 m underground in a former gold mine.
- 3. A Near Detector located approximately 575 m from the neutrino source at Fermilab close to Chicago.

LAr read-out technologies

Module #1

-3.6 m horizontal drift -vertical anode wire planes -vertical resistive cathode photon detectors $\rightarrow XARAPUCAs$

Module #2

-6.5 m vertical drift
-horizontal PCB anode readout (CRP)
-horizontal grid cathode
photon detectors → XARAPUCAs



48 6x6 mm² ganged together in a single electronic channel

Ganging scheme:

- SiPM passively ganged in parallel in groups of 6
- 8 groups of 6 SiPMs ganged actively together throug an OpAmp
- Passive ganging of SiPM reduces the amplitude of single photon signals given the increased overall capacitance
- Active ganging *does not affect the amplitude* but adds electronic noise

Total number of SiPM ~ 300,000



Example: DarkSide 20k

- 50t of Ultrapure depleted Arg
- Primary (S1) and secndary (S2) scintillation and electroluminescence signals detected by *two planes of SiPM arrays*
- SiPM allow for a denser coverage of the surfaces, higher PDE and improved radiopurity



Example: DarkSide 20k

- Basic photo-detection module is an array of 24 SiPM => PDM
- SiPM passively ganged in groups of 6 mixed series/parallel and actively ganged in 4 groups of 6
- 25 PDMs arrangend in a Motherboard Salar: 12.48 4.1 2.19 888: 24.1 Single SiPM • 8,200 PDMs => 196,800 cryogenic SiPMs **本 本 本 本 本** 本 * * * * * * Prototype Tile TIA TIA TIA TIA cryo summing amp





Hybrids: HPD, HAPD, QUPID, VSIPM, ABALONE

- Hybrid detectors are **cobinations of vacuum and silicon devices**
- They have *large photocathode* which ensure a large photoelectron conversion area
- Photolectrons are accelerated by a strong high field (10-20 kV) and focused on a silicon device
- Photoelectrons loose energy producing a number of electron-hole pairs proportional to the electron kinetic energy, eventually multiplied by the intrinsic gain of the silicon device

 $\sigma = \gamma_{i}$

52

- G = e $G_{si}(\Delta V V_{th})/W_{si}$ W_{si} = 3.6 eV for silicon
- Very low fluctuation in the Gain => Fano factor of Si ~0.12



LHCb HPD

- Used in the RICH detector of LHCb
- Cross focused HPD HV = 20kV, G = 5,000
- 8192 pixels of size 500µm x 62.5µm
- QE ~30% @ 350nm
- Vacuum degradation causes ion feedback and aging of the device









out

Hybrids: ABALONE

- Ultra high vacuum device
- Photocathide coupled to a scintillator read-out by a SiPM
- Great optical aperture $\sim 2\pi$
- QE determined by the photocatode
- Excellent gain ~10⁸
- Proposed as photosensor for the *future dark matter experiment DARWIN*



V. D'Andrea et al 2022 JINST 17 C01038



Hybrids: ABALONE

- Operated at 25 kV
- LYSO scintillator
- Gain of the Scintillator ~100
- Overall gain ~10⁸
- Non negligible fraction of electrons escapes from the crystal and does not deposit entire energy (non-returning)









ARAPUCA concept



ARAPUCA concept



- ARAPUCA in the language of *native Brazilian* means *trap* for birds
- Trap photons inside a box with highly reflective internal surfaces => Increase detection efficiency even with a limited active coverage of its internal surface
- Detection efficiency can be tuned by varying the number of SiPMs (ratio between acceptance window and SiPM areas)..
- Baseline DUNE design of Photon Detection System of 2 far detector modules

Dichroic filter

- The core of the device is a **dichroic filter**. It is a dielectric interference **film** deposited on a fused silica substrate.
- It has the property of being highly transparent for wavelength below a cutoff and highly reflective above it.



Foton VUV 128nm \rightarrow p-Therphenyl \rightarrow 350nm \rightarrow Filtro cutoff 400nm \rightarrow TPB \rightarrow 430nm

ARAPUCA Operating Principle

Two shifts of the light:

- One on the external side of filter (Para-Therphenyl 127nm => 350nm)
- The other on the internal side or on the internal walls (Tetraphenyl- γ Butadiene 350nm=>430nm) λ≈127nm



ARAPUCA for neutrinos!







https://www.symmetrymagazine.org/article/arapuca-letthere-be-light-traps?language=pt-br



Illustration by Sandbox Studio, Chicago with Pedro Rivas

ARAPUCA: Façam-se armadilhas para a luz

10/24/19 | By Lauren Biron

As instituições latino-americanas são imprescindíveis para a criação dos detetores de fotões usados na Experiência de Neutrinos em Grande Profundidade.

ARAPUCA in protoDUNE Run I (2018-2020)





X-ARAPUCA



X-ARAPUCA detection efficiency



In units of 8" PMT: 1 X-ARAPUCA bar is equivalent to 2.5 8" PMTs

LAr Read-out Technologies of DUNE far detectors

Module #1

-3.6 m horizontal drift -vertical anode wire planes -vertical resistive cathode photon detectors $\rightarrow XARAPUCAs$

Module #2

- -6.5 m vertical drift
- -horizontal PCB anode readout (CRP)

-horizontal grid cathode

photon detectors → XARAPUCAs



X-ARAPUCA in ProtoDUNE Run II (2022-2023)



X-ARAPUCA in ProtoDUNE Run II (2022-2023)



- X-ARAPUCA is the baseline design for the two DUNE far detector modules
- Design and construction managed by a Consortium of more than 50 Institutions in 3 continents
- 1,500 bar shaped modules in module 1 (Horizontal Drift) => 300,000 SiPMs ~10 m² area
- 640 square X-ARAPUCA modules for module 2 (Vertical Drift) => 100,000 siPMs 3 m² area





Dichroicon: sorting photons



Dichroicon: sorting photons

- Charged particle traveling through liquid scintillator creates both scintillation (~10,000 photons/MeV) and Cherenkov light (~100 photons/MeV)
- Challenge is to detect the Cherenkov light, which provides the direction of the traveling particle. Typically use timing and directionality.
- High light yield from scintillator provides excellent energy and position resolution and low energy thresholds



Dichroicon: combining two technologies

Winston Cones





https://arxiv.org/pdf/physics/0310076.pd



Dichroic Filters





Tanner Kaptanoglu -






THEIA Experiment (proposed)



- Water based Liquid Scintillator
- Read-out system possibly based on LAPPD, PMTs and Dichroicon
- Large-scale detector that can *discriminate between Cherenkov and scintillation signals.*
- Reconstruct *particle direction and species using Cherenkov light* while having *excellent energy resolution and low threshold of a scintillator detector*
 - *low- and high-energy solar neutrinos*
 - Neutrino mass ordering
 - measurement of the neutrino CP-violating phase
 - observations of diffuse supernova neutrinos and neutrinos from a supernova burst
 - nucleon decay

Light Yield of a scintillation detector

LY can be factorized as:

$LY = N_y \, \varepsilon_{PSD} \, \varepsilon_{OPT}$

Where:

- ✓ N_{γ} is the photon yield of the scintillator => number of photons produced per unit of deposited energy by a certain radiation
- ✓ ε_{PSD} is the conversion efficiency of the PSDs => the efficiency of the PSD system in converting photons into signal (photo-electrons)
- ✓ ε_{opt} is the optical efficiency => the fraction of the originally produced photons that manages to cross the windows of the PSDs N_γ and ε_{PSD} in the majority of the cases are precisely known. On the other side ε_{opt} is typically unknown and needs to be estimated.

Recursivity of light propagation (think to a sphere...)

The propagation of photons inside a scintillation detector is an intrinsically recursive process.

Reflectivity = **R**

Optical window – photcatodic coverage = **f**

A photon produced in a **random point inside the sphere** with a random direction when reaches the boundary surface has an *average probability* **f** to be detected

And a probability **R(1-f)** to be sent back into the chamber

Reflected photons has again a detection probability equal to f and a probability to be reflected equal to R(1 - f) when hit the boundary surface for the second time

The same situation will repeat again *identical to itself after any reflection*.

...let's generalize a little bit (I)

- Consider a **general scintillation detector** and assume that the **process is recursive**
- It can be divided into a series of subsequent and indistinguishable steps and it is possible to define two quantities:
 - \checkmark **a** is the average probability per step that a photon is detected
 - *β* is the average probability per step that a photon is regenerated, that is the probability that it is not lost (detected or absorbed) and that some physical process randomizes again its direction (reflection for instance)
 - \checkmark α and β are constant for all the steps (recursivity of the process)
 - $\checkmark \alpha \le 1$ and $\beta < 1$
- Detection and regeneration probabilities after n steps are easily calculated:

	detection probability	regeneration probability
step 0	α	β
step 1	αβ	β^2
step 2	$\alpha \beta^2$	β^3
step n	$\alpha \beta^n$	β^n

78

...let's generalize a little bit (II)

$$\varepsilon_{opt} = \sum \alpha \beta^n = \frac{\alpha}{1 - \beta}$$

For the *spherical scintillator* (and one can safely extend **to scintillators of regular shapes**) this means that:

$$\varepsilon_{opt} = \frac{f}{1 - R(1 - f)} \qquad \frac{\alpha = f}{\beta = R(1 - f)}$$

If the optical window has transmissivity T_w and reflectivity R_{w} :

$$\varepsilon_{opt} = \frac{T_w f}{1 - R(1 - f) - R_w f} \qquad \frac{\alpha = T_w f}{\beta = R(1 - f) + R_w f}$$

Including Rayleigh scattering and absorption

Define:

 $\lambda_{R} \rightarrow Rayleigh scattering length$ $<math>\lambda_{A} \rightarrow absorption length$

and:

probability that a photon reaches the end of the step, as defined in absence of scattering/absorption. without interactions

It is possible to define:

 $\overline{\lambda} = \overline{\lambda}$

Photon regenerated by reflections

Photon regenerated by scattering

with α_0 and β_0 detection and regeneration probabilities in absence of scattering and absorption

 $\begin{aligned} \alpha &= U_{\scriptscriptstyle RA} \alpha_0 \\ \beta &= U_{\scriptscriptstyle RA} \beta_0 + (1 - U_{\scriptscriptstyle RA}) \end{aligned}$

Including Rayleigh scattering and absorption (II)

After little algebra one finds:

$$\mathcal{E}_{opt} = \frac{\alpha_{0}}{Q - \beta_{0}} \begin{bmatrix} \alpha_{0} & \alpha$$

For our simple detector, under reasonable hypotheses, one obtains:

$$U_{RA} = \frac{\tilde{\lambda}}{\tilde{L}} (1 - e^{-\tilde{L}/\tilde{\lambda}}) \quad \text{with:} \quad \tilde{L} = 6\frac{V}{S}$$

It is the characteristic linear

dimension of the detector



Monte Carlo tests of the model (II)



Monte Carlo tests of the model (scattering and absorption)



1 PSD Specular reflectivity R = 0.95 Rayleigh scattering = 10 cm

1 PSD PSD radius = 4 cm Specular reflectivity R = 0.95 Absorption length = 50 cm

$$U_{RA} = \frac{\lambda}{\tilde{L}}(1 - e^{-\tilde{L}/\tilde{\lambda}})$$

Monte Carlo tests of the model: Parallelepiped



Parallelepiped 10 cm × 10 cm × 30 cm (l × w × h).
 2 PSDs on the opposite faces
 Windows' reflectivity is set at 0.3 and transmissivity at 0.5 (as for the cubic scintillator)

Parallelepiped (cont.)



RED LINES :

$$\varepsilon_{opt} = \frac{T_w f}{1 - R(1 - f) - R_w f}$$

Extremely asymmetric detector

- Surprisingly good agreement between MC simulation and model outcomes
- For specular reflectivity, above 0.9, discrepancies at the level of few percent are found
- In all other cases they are of the order of 10%.

LY estimation of a real detector



LAr scintillation chamber - PTFE cell (h = 9.0 cm and $\varphi = 8.4 \text{ cm}$) - observed by a single 3" photomultiplier (Hamamatsu R11065) Internal surface of the cell completely covered with a reflective foil deposited with Tetra Phenyl Butadiene.

87

WARP Collaboration, Demonstration and comparison of photomultiplier tubes at liquid Argon temperature, 2012 JINST 7 P01016 [arXiv:1108.5584].

LY estimation of a real detector (I)

Parameters used for LY estimation

photon yield	$N_{\gamma} = 40$ photons/keV [11]
photocathodic coverage	f = 13%
transmissivity of PMT window	$T_w = 0.94$ [12]
reflectivity of PMT window	$R_w = 0$
conversion efficiency of PMT	$\varepsilon_{PSD} = 28\%$
no absorption of VUV photons	$Q_{VUV} = 1$
no absorption of visible photons	$Q_{vis} = 1$
conversion efficiency of passive surface	$\varepsilon_{WLS} = 1$ [13]
conversion efficiency of PMT window	
(no shifter)	$\varepsilon_{wls} = 0.$
reflectivity of passive surface (reflector+TPB)	R = 0.95 [14]

LAr is assumed to be pure => no absorption (Q = 1) Photons converted on the walls $LY_{estimated} = N_{\gamma} \varepsilon_{PSD} \varepsilon_{WLS} (1-f) \frac{T_{w}f}{1-R(1-f)} = 6.9 \frac{phel}{keV}$ $LY_{measured} = 7 \frac{phel}{keV} \pm 5\%$

BACK-UP

Hamamatsu R5912



Journal of Instrumentation 13(10):P10030-P10030 ICARUS Collaboration

Hamamatsu R11065





PHOTOMULTIPLIER TUBE

R11065

For Low Temperature Operation down to -186 deg. C Special Bialkali Photocathode (Bialkali LT), Low Radioactivity 76 mm (3 Inch) Diameter, 12-stage, Head-on Type, Synthetic Silica

General

Parameter		Description / Value	Unit
Spectral response		160 to 650	nm
Wavelength of Maximum Response		420	nm
Window material		Synthetic silica	
Thetesatheda	Material	Bialkali	1 H
Photocathode	Minimum Effective Area	64	mm dia_
D 1	Structure	Box & Linear-focused	
Dynode	Number of Stages	12	
Operating Ambient Temperature		-186 to +50	deg. C
Storage Temperature		-186 to +50	deg. C

Electronics – Cold active ganging

Read out electronic is divided into *two stages*: Cold active ganging board and digitizing board

• Cold ganging (summing) board:

- SiPMs are small devices. They need to be ganged in order to contain the number of readout channels. There is a limitation on the number of channels per ARAPUCA bar due to the space available to route the cables inside APA (see D. Warner talk). There will be 4 readout channels per module → one channel per X-ARAPUCA supercell.
- **48 SiPMs will be ganged together**. The ganging is active, that is using active components (Operational Amplifier)
- The active ganging board is installed on the X-ARAPUCA module and *operates at LAr temperature*
- Two active ganging circuits developed by the Consortium. With different degrees of maturity at this moment. Both demonstrated LAr operations in the consortium of the consortium. With different degrees of maturity at



X-ARAPUCA concept

- There are *two main mechanisms* through which a photon can be detected by the X-ARAPUCA:
 - Standard ARAPUCA mechanism. The photon, after entering the X-ARAPUCA box, is converted by the WLS of the inner slab, but is not captured by total internal reflection. In this case the photon bounces a few times on the inner surfaces of the box until when it is or detected or absorbed;
 - Total internal reflection. The photon, converted by the filter and the slab, gets trapped by total internal reflection. It will be guided towards one end of the slab where it will be eventually detected. This represents an improvement with respect to a conventional ARAPUCA, which contributes to reduce the effective number of reflections on the internal surfaces. The sides of the slab where there are not active photo-sensors will be coated with a reflective layer which will allow to keep the photon trapped by total internal reflection.

Operating principle



- The simplest geometry is a flattened box with highly reflective internal surfaces (Teflon, VIKUITI, VM2000) with an open side.
- The open side hosts the **dichroic filter** that is the acceptance window of the device
- The filter is deposited with **TWO WAVELENGTH SHIFTERS (WLS) – one on each side**
- The shifter on the external side, S1, converts LAr scintillation light (128 nm) to a wavelength L1, with L1 < cutoff
- The shifter on the internal side, S2, converts S1 shifted photons to a wavelength L2, with L2 > cutoff
- The internal surface of the ARAPUCA is observed by one or more SiPM
 94