

Plasma Photocathode

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Acknowledgements



BELLA (Berkeley Lab Laser Accelerator) Center:



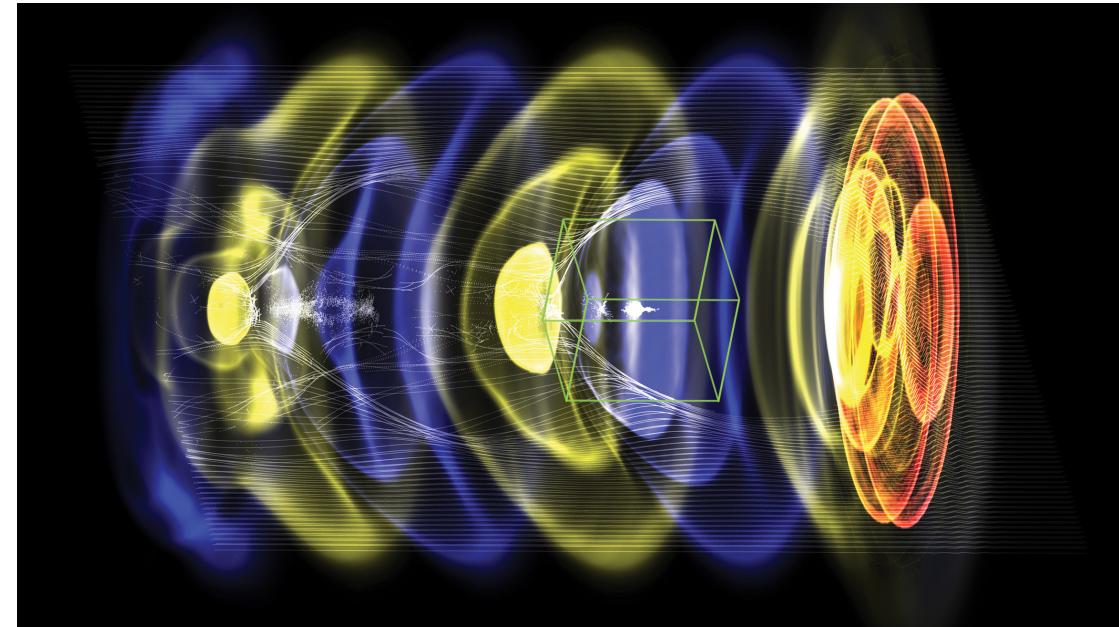
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Outline

- Plasma-based accelerators (PBAs):
 - Multi-GeV electron beams
 - Femtosecond bunch lengths
 - Sub-micron emittance
 - High-charge bunches
 - PBAs enable high brightness e-beams
- Beam quality depends on injection methods:
 - Self-trapping
 - Plasma density gradient injection
 - Ionization injection
- Path to higher brightness electron beams with the “Plasma photocathode” (relevant for colliders or radiation sources, e.g. x-ray FEL)



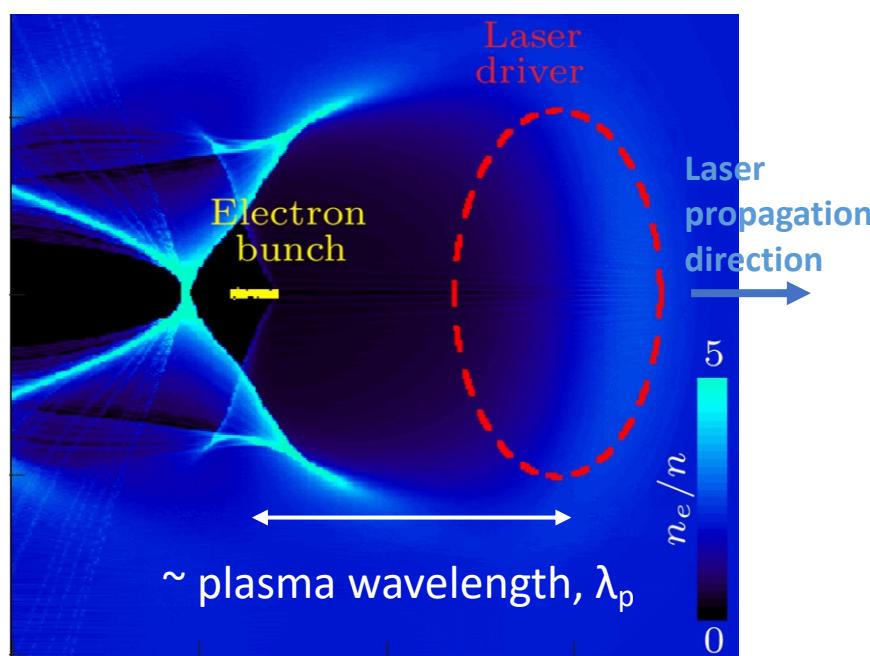
Intense laser (red) drives a density wave (blue to yellow) in plasma, electrons (white) are trapped and accelerated with fields of order 10 GV/m.

Plasma-based acceleration of electron bunches

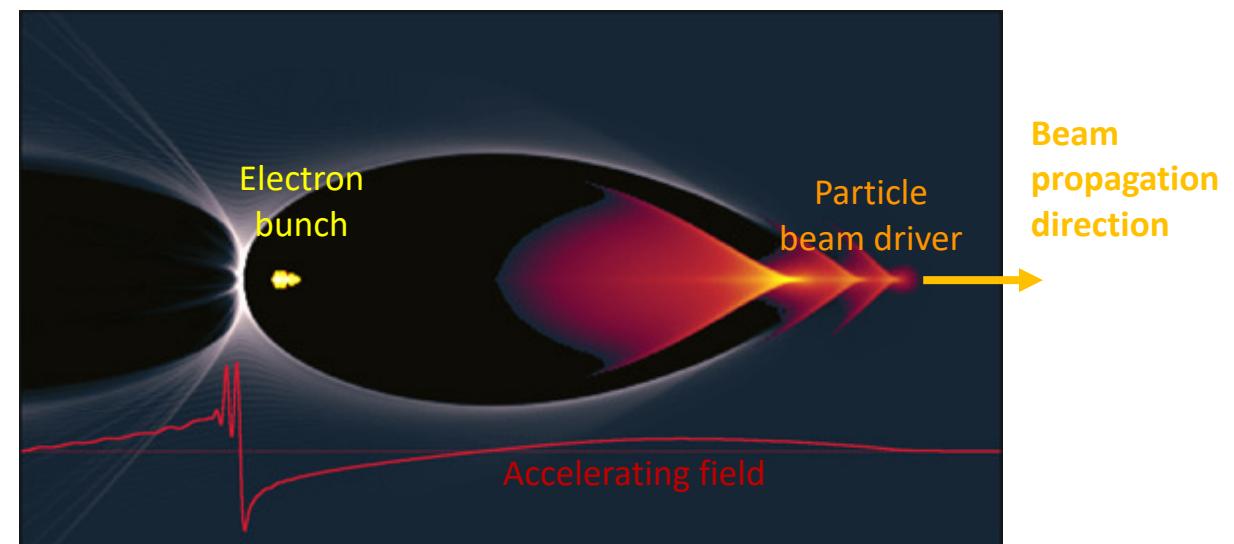
- Plasma wave (wakefield) driven by ponderomotive force of intense laser [or space-charge field of particle beam]
- Electron displaced by driver (ions remain stationary) → EM fields from charge separation → plasma oscillations
- Focusing and acceleration provided by plasma wave
- Characteristic size of the plasma wave: $\lambda_p = 100 - 10 \mu\text{m}$ for plasma density $10^{16} - 10^{19} \text{ cm}^{-3}$
- Characteristic strength of the wakefields: $10 - 100 \text{ GV/m}$ for plasma density $10^{16} - 10^{19} \text{ cm}^{-3}$

$$E \sim E_0 = m_e c \omega_p / e \simeq 96 \sqrt{n_0 [\text{cm}^{-3}]}$$

Laser-driven plasma wakefield accelerator (LWFA):



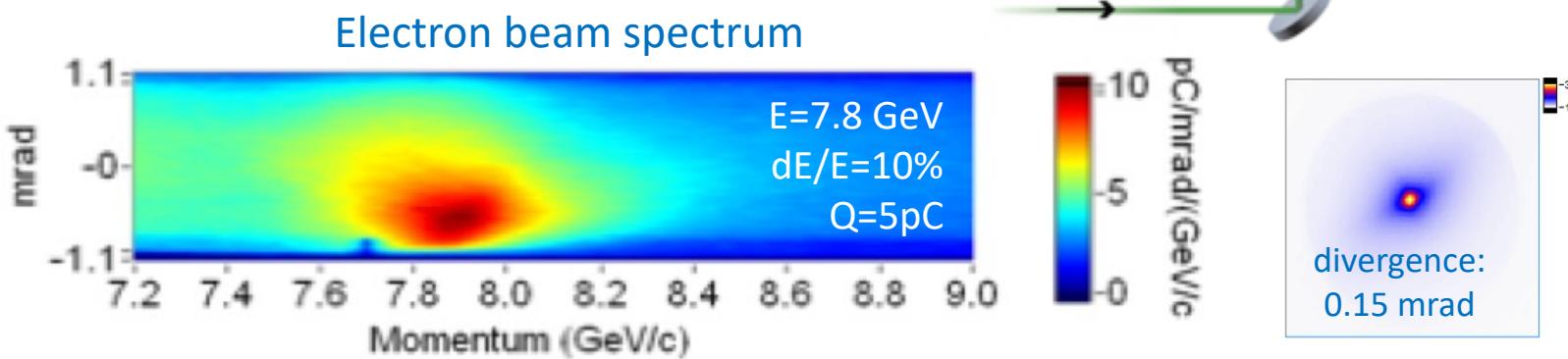
Beam-driven plasma wakefield accelerator (PWFA):



Laser-driven plasma accelerator state-of-the-art: multi-GeV accelerator (~8 GeV in 20 cm)

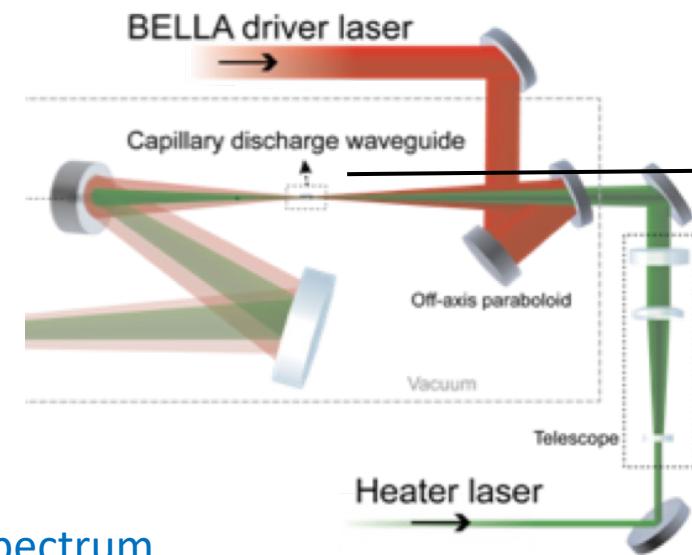


PW (40 J, 30 fs, 1Hz)

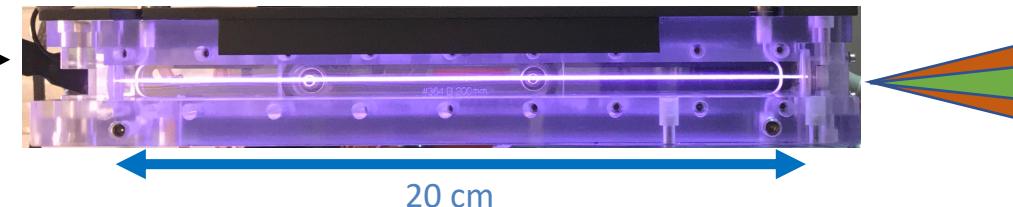


- Key technology: plasma channel creation for guiding at low density ($3 \times 10^{17} \text{ cm}^{-3}$)
 - discharge capillary + laser heater
- New technologies emerging: hydrodynamic optical-field-ionized plasma channels

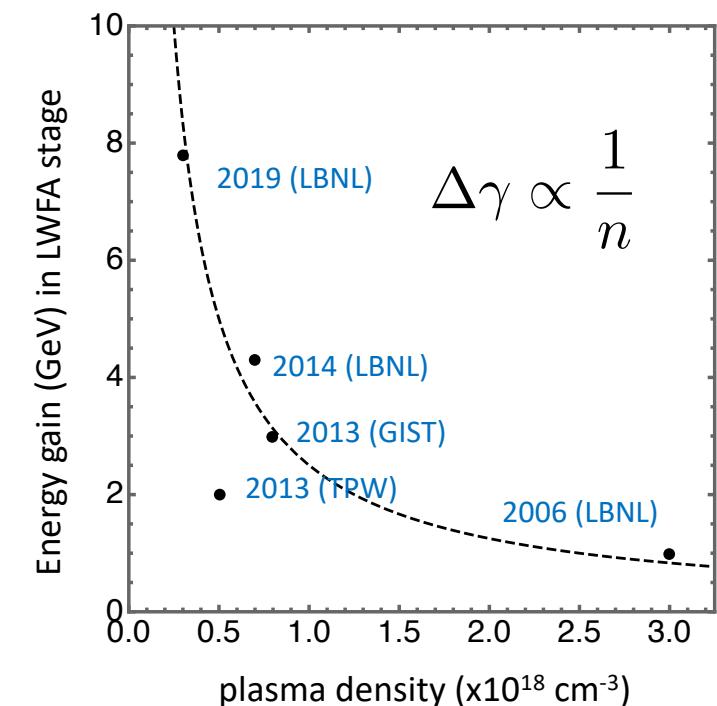
Single-stage acceleration: 8 GeV in 20 cm (40 GV/m)



Laser-heater-assisted discharge-capillary plasma waveguide



Gonsalves et al. PRL (2019)



Plasma-based accelerators enable high-brightness beams

- Plasma accelerators have the potential to yield high-brightness beams:
 - Rapid acceleration ($>10 \text{ GV/m}$) + shielding from background ions reduce detrimental space-charge effects
 - Ultra-short bunches: $\sigma_z \ll \lambda_p$ (e.g., $\lambda_p \sim 10 \mu\text{m}$ for 10^{19} cm^{-3}) → **1-10 fs** (exp. measured) Buck et al., Nature Phys (2011)
Lundh et al., Nature Phys (2011)
 - Low normalized emittances → **0.1 – 1 μm** (exp. measured) Weingartner et al., Phys Rev ST AB (2012)
Plateau et al., PRL (2012)
 - High charge in a monochromatic peak (< a few %) → **10 -100 pC** (exp. measured) Kirchen et al., PRL (2021)
Götzfried et al., PRX (2020)
Wang et al., PRL (2016)

- Typical brightness:

$$B_{5D} = \frac{2I}{\epsilon_n^2} \sim 10^{18} \text{ A}/(\text{m-rad})^2$$

↑ $\sim (10-100 \text{ pC})/(1-10 \text{ fs}) \sim 10 \text{ kA}$
↑ $\sim 0.1 - 1 \mu\text{m}$

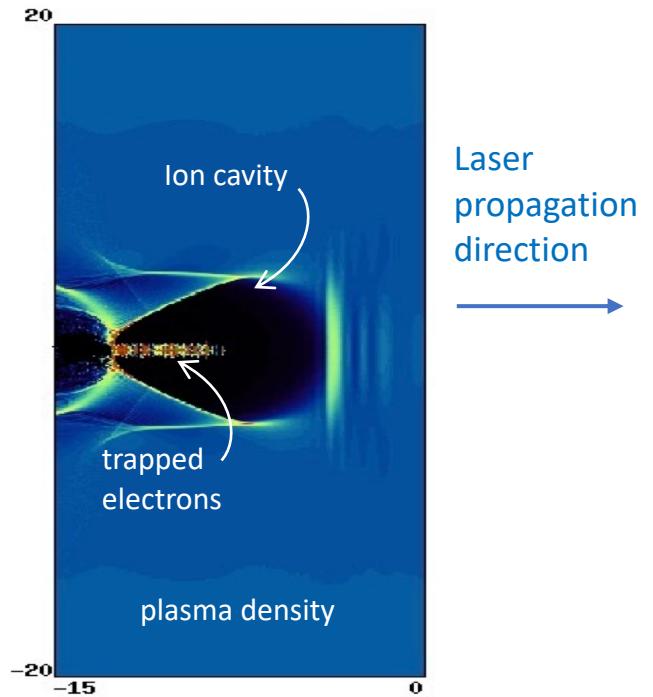
→ Higher quality beam generation (10s nm emittance) possible but requires “controlled / triggered” injection into plasma wave

Injection methods determine electron beam properties and quality

Injection (from background plasma) methods:

- Self-injection
- Triggered injection:
 - Plasma density gradient injection
 - Laser triggered injection
 - Colliding pulse injection
 - Laser ionization injection
 - Single laser pulse
 - Driver + laser injector
("plasma photocathode")

Self-injection



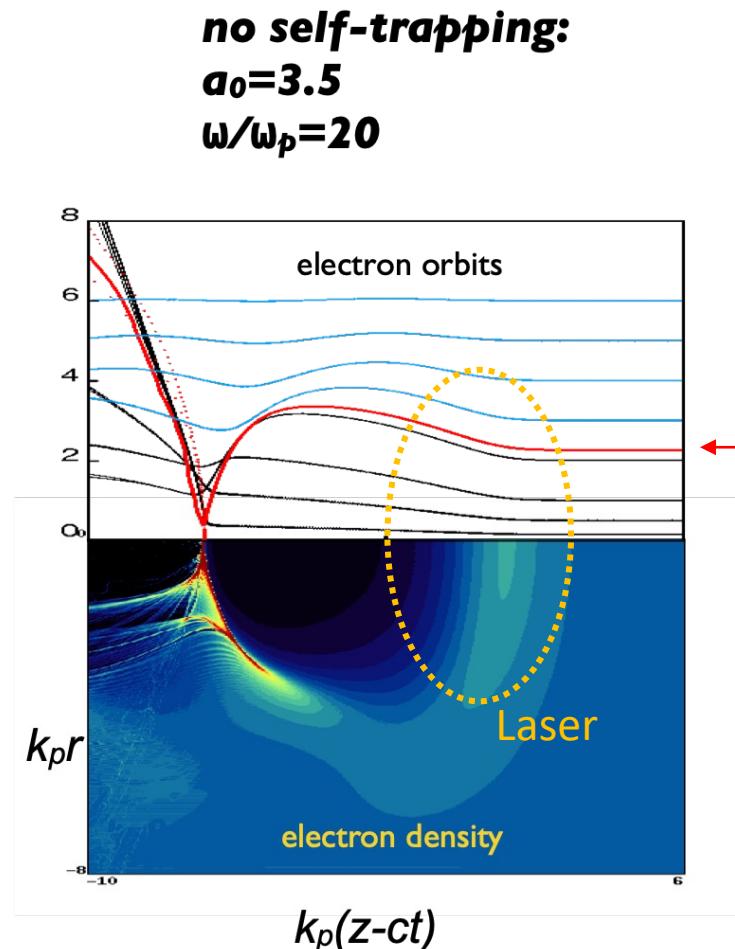
- Requires large amplitude wakefield → high laser intensity

Laser strength: $a_0 = 0.85\lambda[\mu\text{m}] \left(I[10^{18}\text{W/cm}^2] \right)^{1/2}$

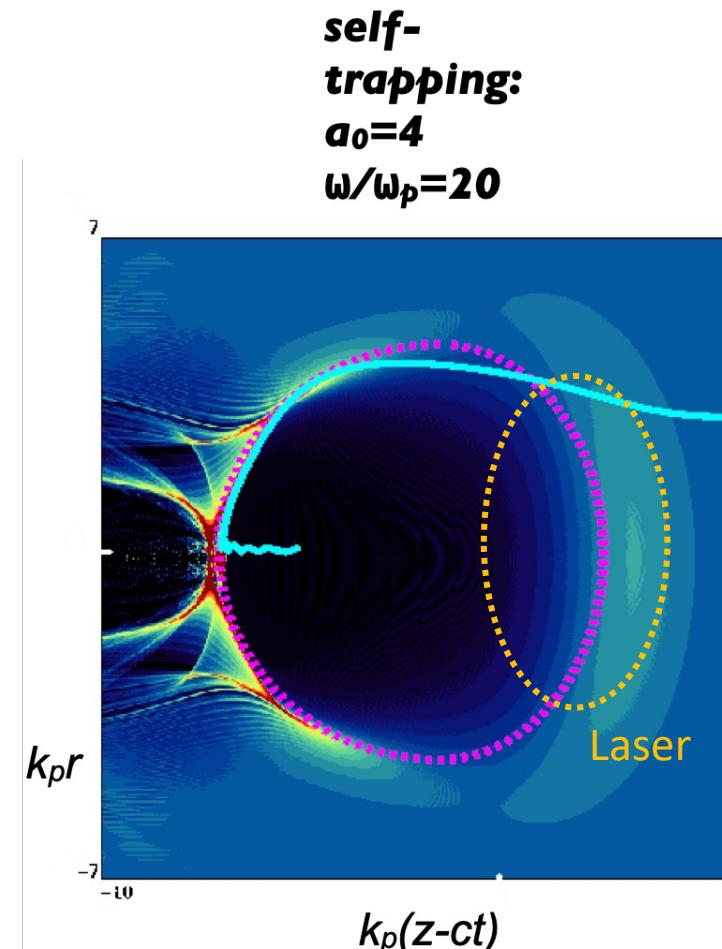
- Requires slow plasma wave velocity → high plasma density

Wake velocity: $\gamma_p \sim \frac{\omega_L}{\omega_p} \propto n^{-1/2}$

Electron orbits in a laser-produced wakefield



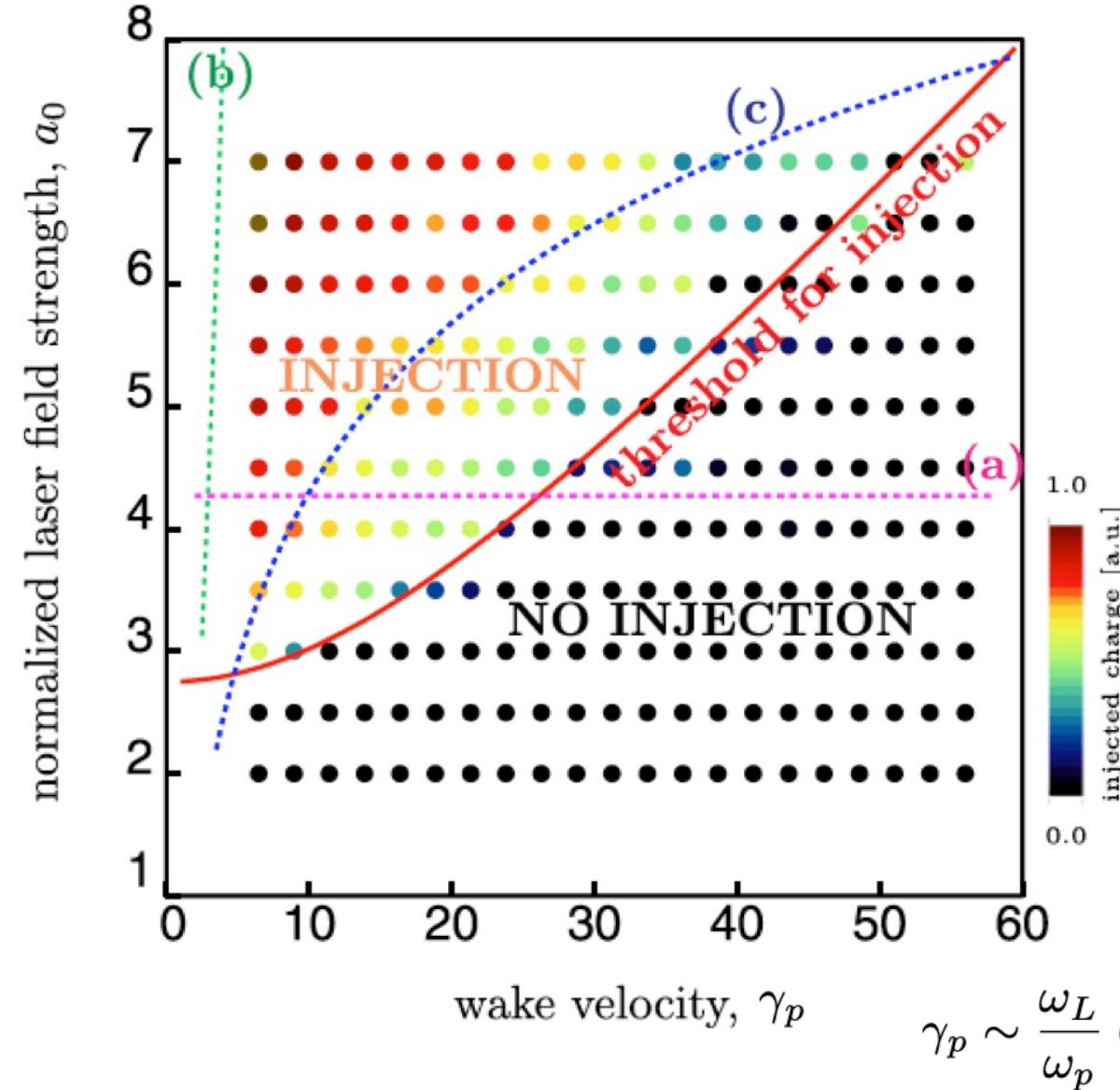
electrons orbits along bubble “border” can reach back of bubble and be trapped



- Trapping determined by the strength of the accelerating wakefield (laser strength) + wake velocity (plasma density)

Self-injection trapping threshold

Benedetti et al., Phys Plasmas (2013)



- Trapping threshold a function of amplitude of the accelerating wake field (laser strength) and wake phase velocity (plasma density)
- Sub-micron emittance observed from self-trapping

Density gradient injection (downramp injection)

- Reduce phase velocity by exciting plasma wave on a density gradient

S. V. Bulanov, PRE (1998)

C. G. R. Geddes, PRL (2008)

X. L. Xu, PRAB (2017)

$$E \sim \cos[k_p(z)(z - \beta_g t)] = \cos \psi$$

Phase velocity:

$$\beta_p = \frac{\omega}{k} = \frac{-\partial_t \psi}{\partial_z \psi}$$

$$\beta_p \simeq \beta_g [1 - |\zeta| k_p^{-1} dk_p/dz]^{-1}$$

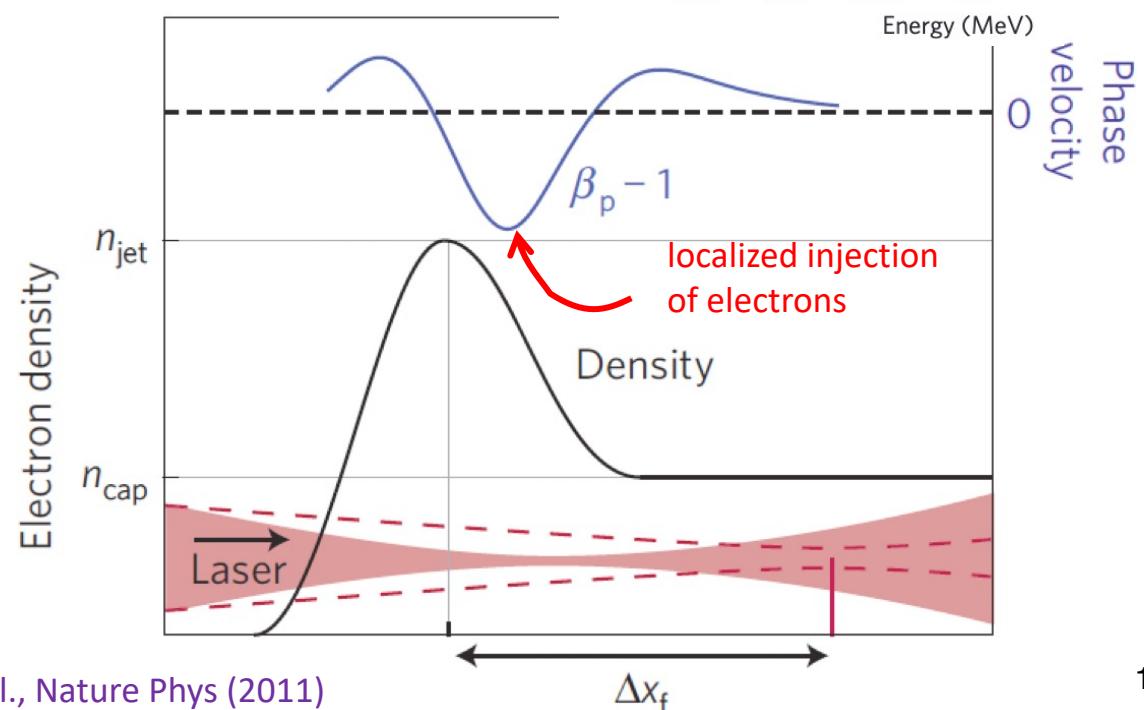
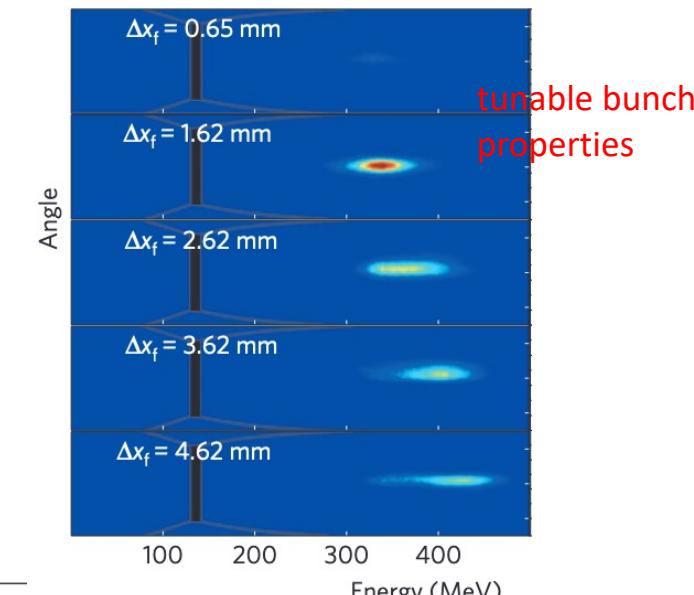
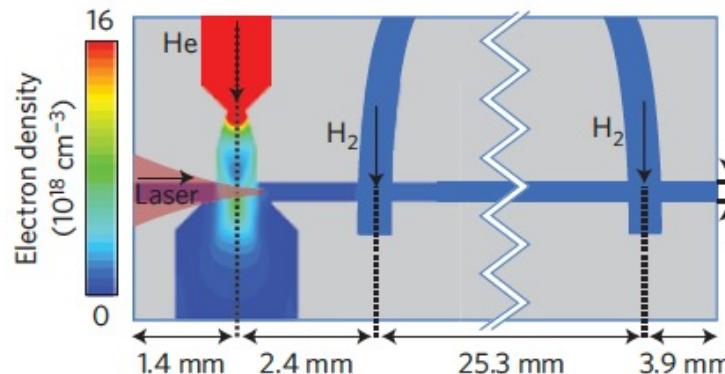
Laser velocity



Control phase velocity using density variation and laser evolution:

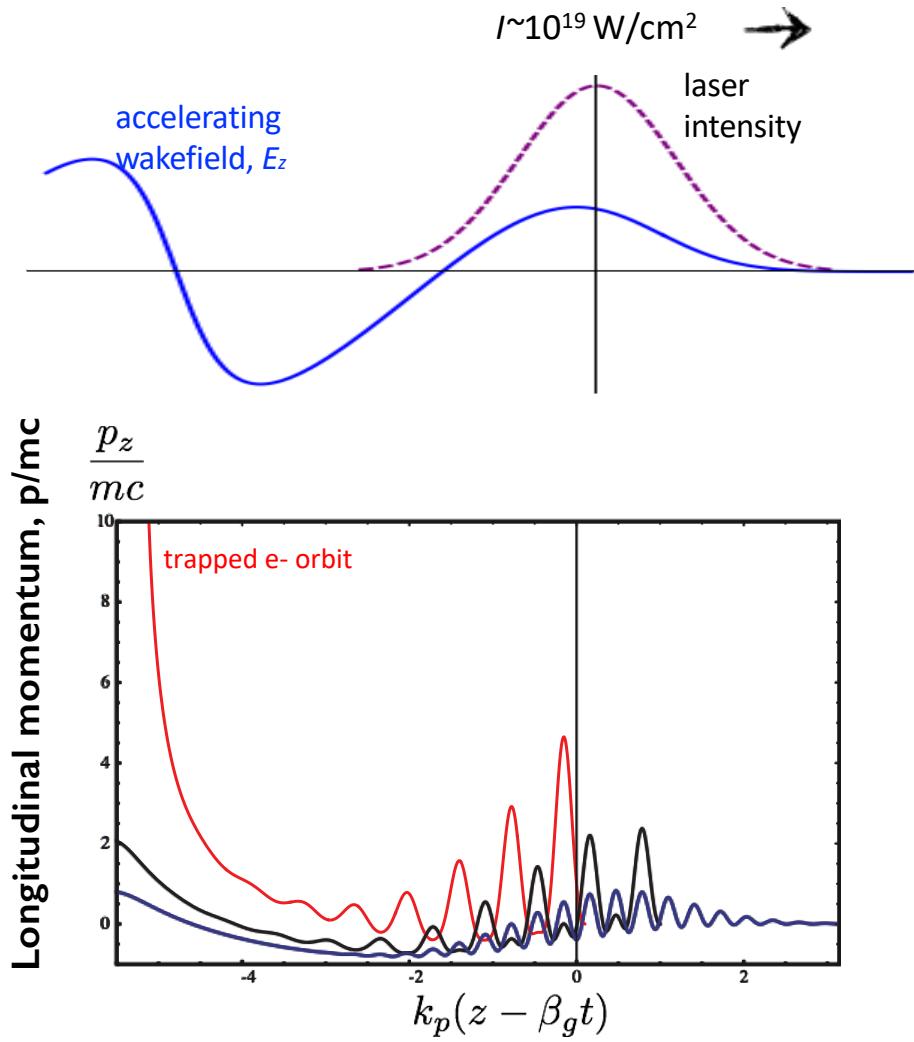
$$\frac{1}{k_p} \frac{dk_p}{dz} = \frac{1}{2n} \frac{dn}{dz} - \frac{1}{\lambda} \frac{d\lambda}{da} \frac{da}{dz}$$

Density gradient Laser evolution



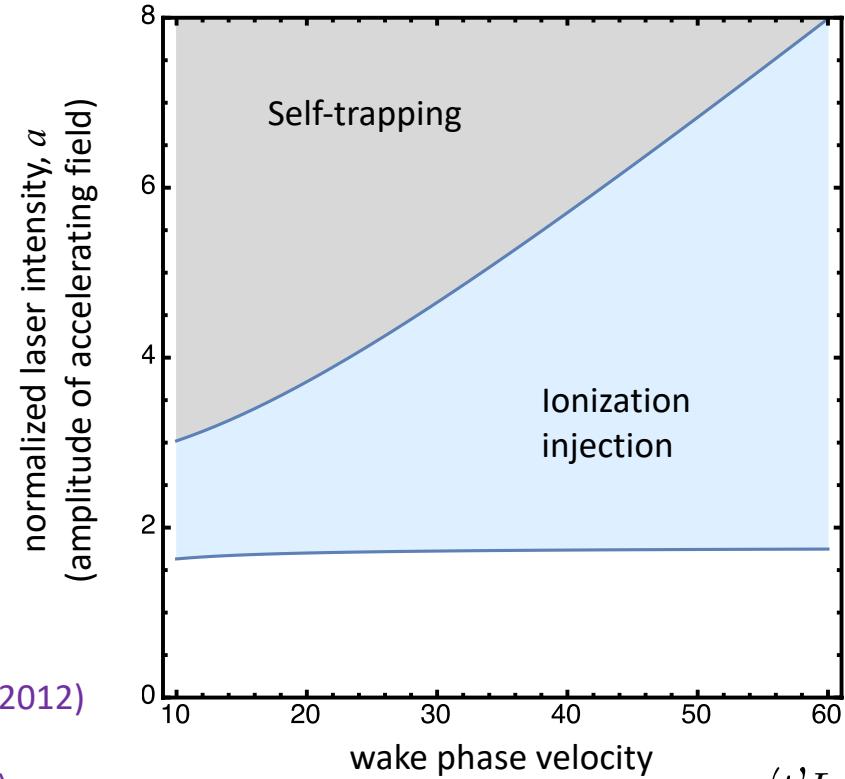
Ionization injection

Ionization injection: ionize bound electron (at rest) in trapped orbit of plasma wave using high-Z gas (e.g., Hydrogen doped with Nitrogen)



Chen et al., Phys. Plasmas (2012)
 Pak et al., PRL(2010)
 McGuffey et al., PRL (2010)

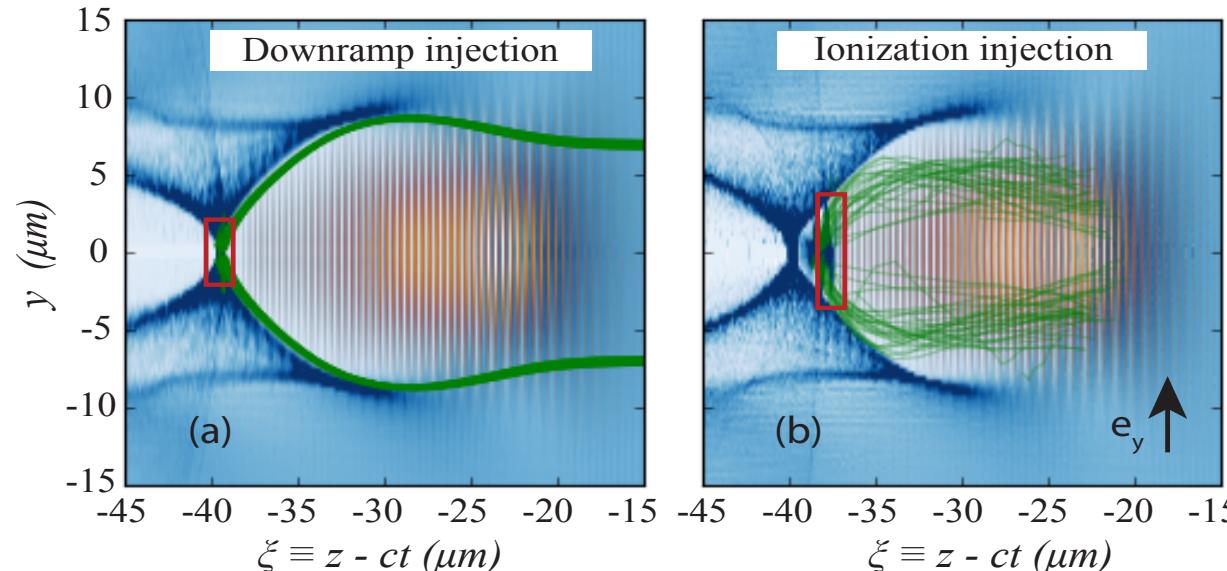
- Weak dependence on density (wake velocity)
- Amplitude for trapping reduced: $a > 1.7$
 → enables trapping at lower laser intensity and lower plasma densities
- Large energy gains achieved operating at low density



$$\gamma_p \sim \frac{\omega_L}{\omega_p} \propto n^{-1/2}$$

Emittance dependence on injection method

- Characteristics of trapped e-beam determined by initial distribution of injected electrons (injection method)



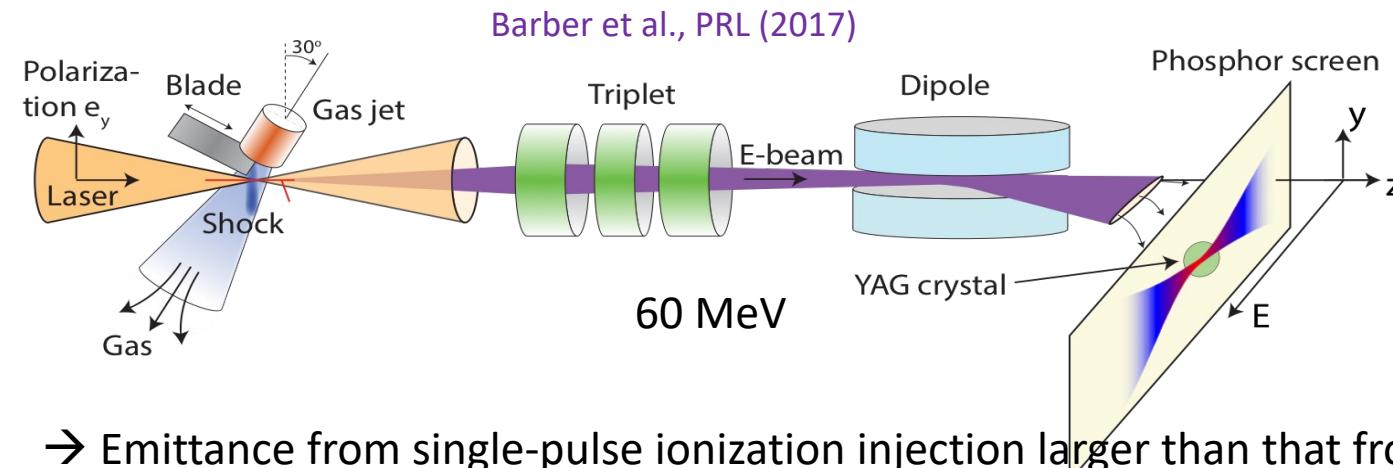
- transverse size \sim laser waist

$$\sigma_x \propto w_0 \sim 2k_p^{-1} \sqrt{a_0}$$

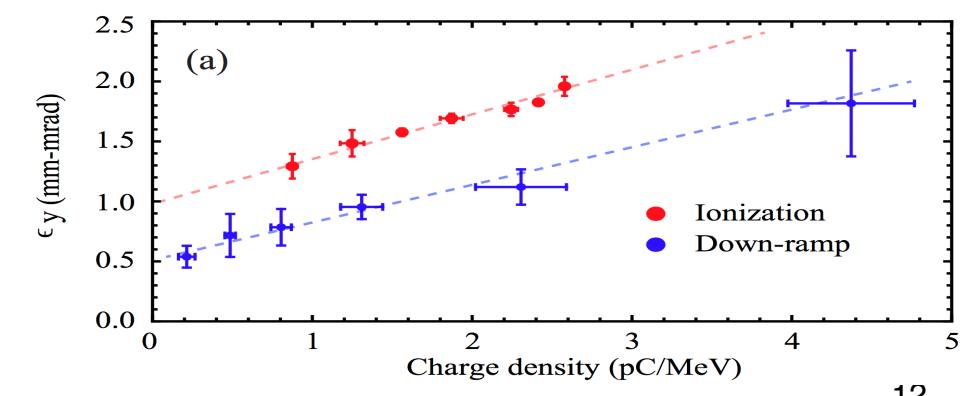
- transverse momentum: $\sigma_p = [\sigma_{\text{quiver}}^2 + \sigma_{\text{PMF}}^2]^{1/2}$

$$\sigma_{\text{quiver}} \propto a_0$$

$$\sigma_{\text{PMF}} \propto \frac{a_0^2}{\sqrt{1 + a_0^2/2}}$$

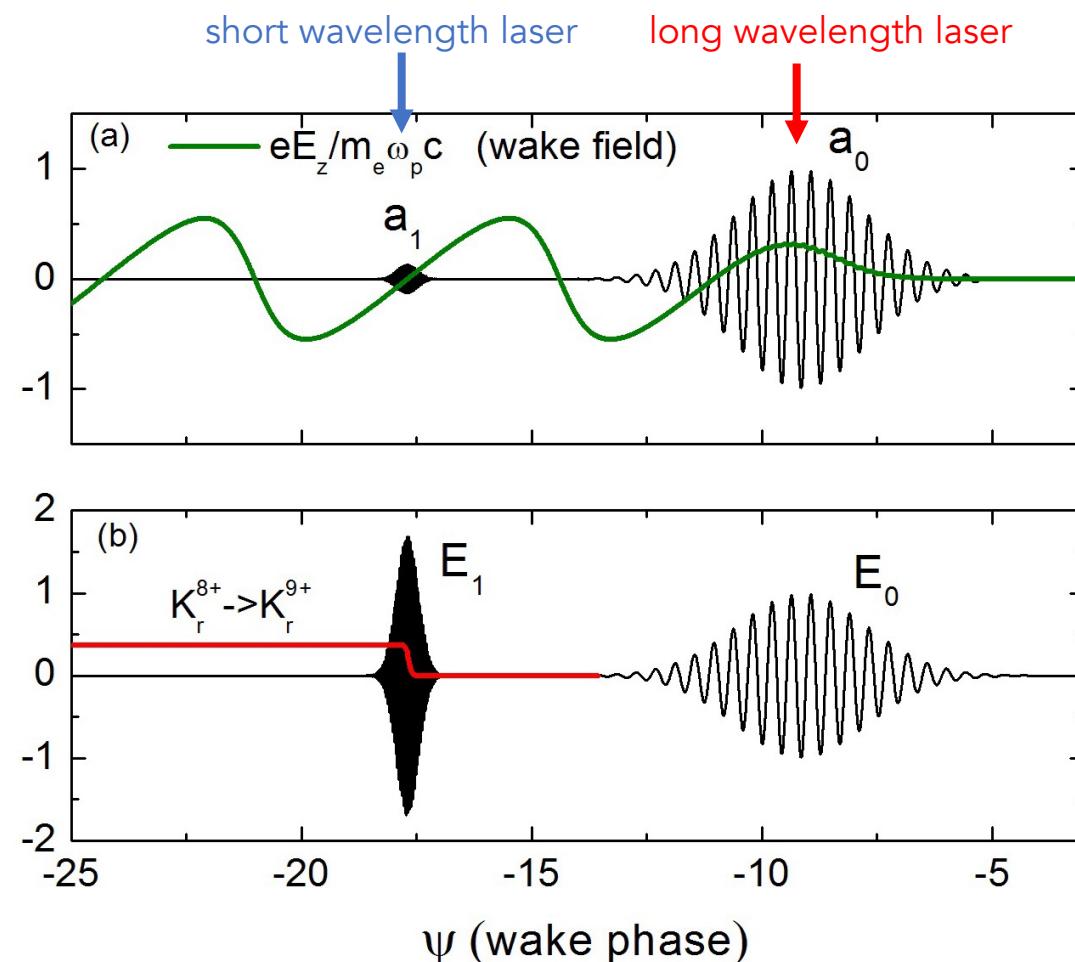


→ Emittance from single-pulse ionization injection larger than that from downramp injection due large spread in initial transverse momenta



Plasma photocathode: higher brightness e-beams from two-color laser ionization injection

- **Two-color, two-pulse laser-ionization injection:** use two lasers of different wavelengths to separate plasma wave excitation (**long wavelength**) and ionization injection (**short wavelength**)



$$\text{Laser electric field: } E = (2\pi m_e c^2 / e) \frac{a}{\lambda}$$

→ Long wavelength laser for large ponderomotive force and plasma wave excitation

$$\text{ponderomotive force: } F_{\text{PMF}} = m_e c^2 \nabla a^2 / 2$$

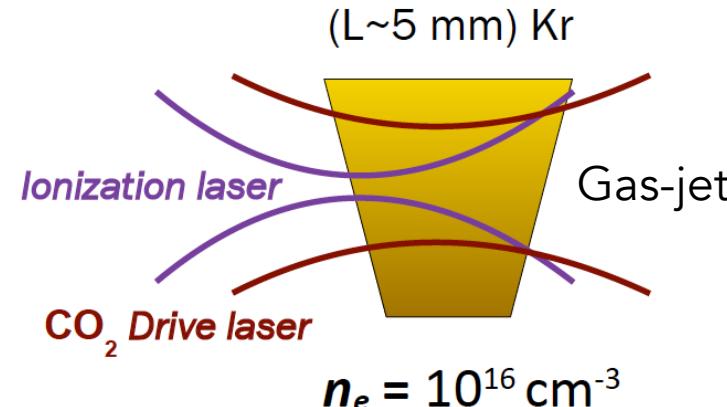
→ Short wavelength laser for large peak electric field and ionization

$$\text{Ionization rate: } w \propto \frac{E_a}{E} \exp \left[-\frac{2}{3} \frac{E_a}{E} \left(\frac{U_{\text{ion}}}{U_H} \right)^{3/2} \right]$$

► Requires: $a_i > \left(\frac{\lambda_i}{\lambda_0} \right) a_0$

Two-color ionization injection: charge vs emittance

Two-color ionization injection using CO₂ drive laser can achieve 10s nm emittance with 10s pc charge.

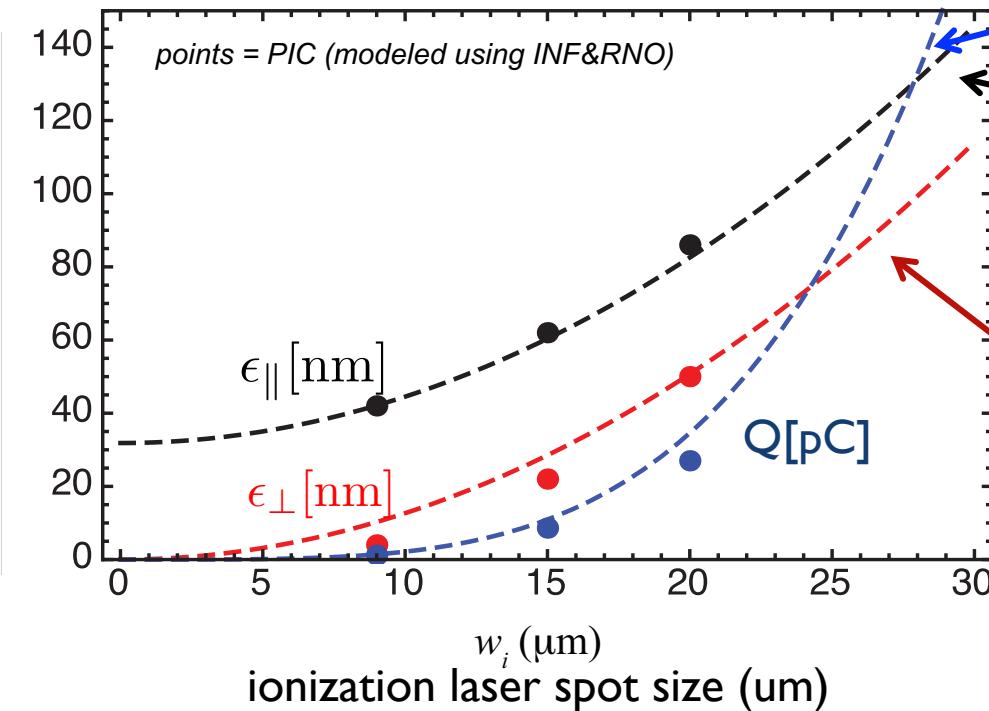
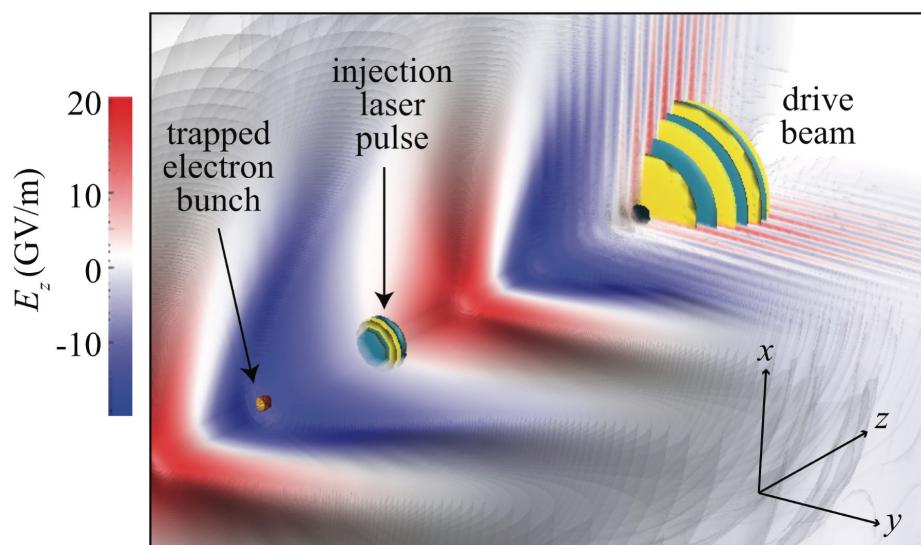


CO₂ pump laser pulse:
 $a=2$
 10 um wavelength
 470 fs (FWHM intensity)
 155 um spot ($ZR = 7.5$ mm)
 10 J

Frequency-doubled Ti:Sa injection pulse:
 $a=0.13$
 0.4 um wavelength
 50 fs
 20 um spot ($ZR = 3$ mm)
 120 mJ
 \rightarrow ionizes $\text{Kr}^{+8} \rightarrow \text{Kr}^{+9}$

charge:

$$Q = eN_t = en_g(\pi w_i^2 \Delta^2) Z_{Ri} \propto w_i^4$$



normalized emittance (in laser pol. plane):

$$\epsilon_{\parallel} = \left[1 + \left(\frac{2a_i}{k_p w_i} \right)^2 \right] \frac{k_p w_i^2}{4\sqrt{2}} \Delta^2$$

normalized emittance (out of pol. plane):

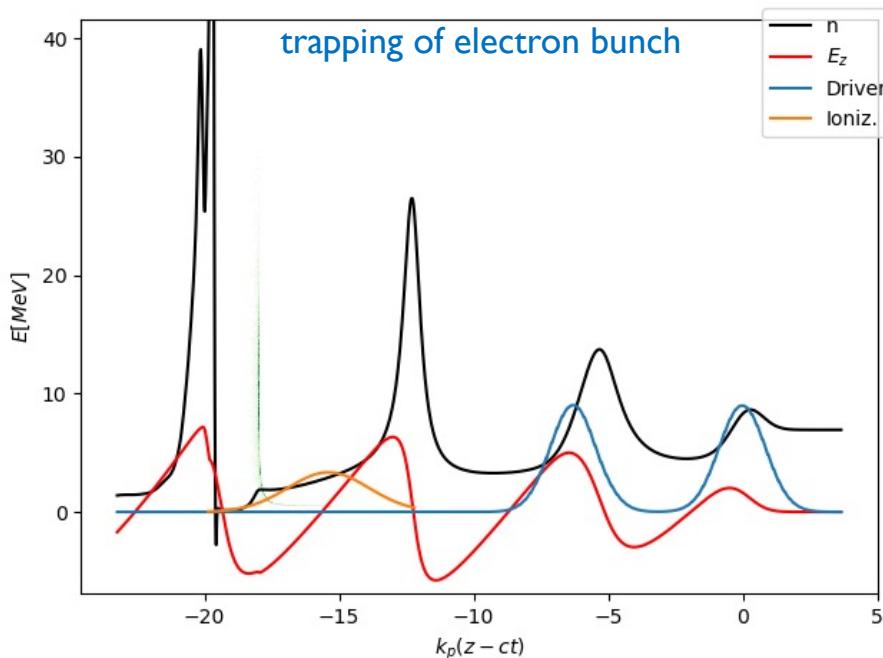
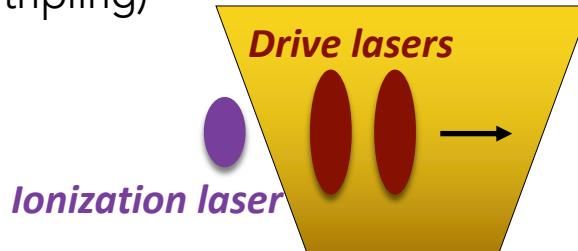
$$\epsilon_{\perp} = \frac{k_p w_i^2}{4\sqrt{2}} \Delta^2 \propto w_i^2$$

Two-color ionization injection: proof-of-principle demo planned at LBNL



Ionization laser pulse - frequency tripled Ti:Sa

- $E = 1.9 \text{ J}$ (before frequency tripling)
- Tripling efficiency = 15%
- $w_0 = 7 \text{ um}$
- $T_{FWHM} = 70 \text{ fs}$
- $a_0 = 0.51$



PIC simulation using INF&RNO (D.Terzani)

Electron beam after 1cm

- emittance: $0.08 \mu\text{m} / 0.35 \mu\text{m}$
- 15 pC
- 400 MeV ($dE/E < 3\%$)

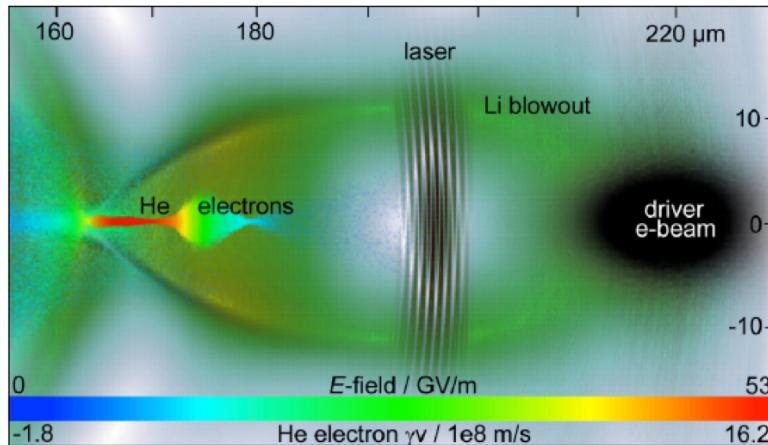
Laser wake drivers:

- $E = 13 \text{ J}$ (total for both pulses)
- $w_0 = 52 \text{ um}$
- $a_0 = 0.92$
- $T_{FWHM} = 40 \text{ fs}$
- $P = 153 \text{ TW}$ (single pulse)

Plasma

- $n_0 = 3e17 \text{ cm}^{-3}$
- $P_{crit} = 100 \text{ TW}$
- $L \sim 1\text{cm}$ (driver Rayleigh range)
- Nitrogen

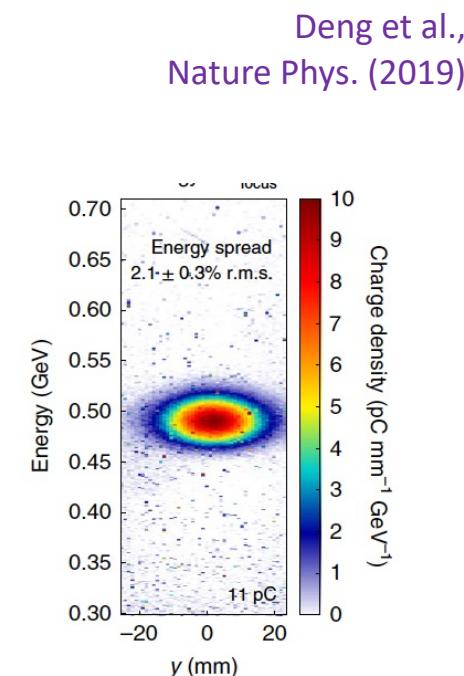
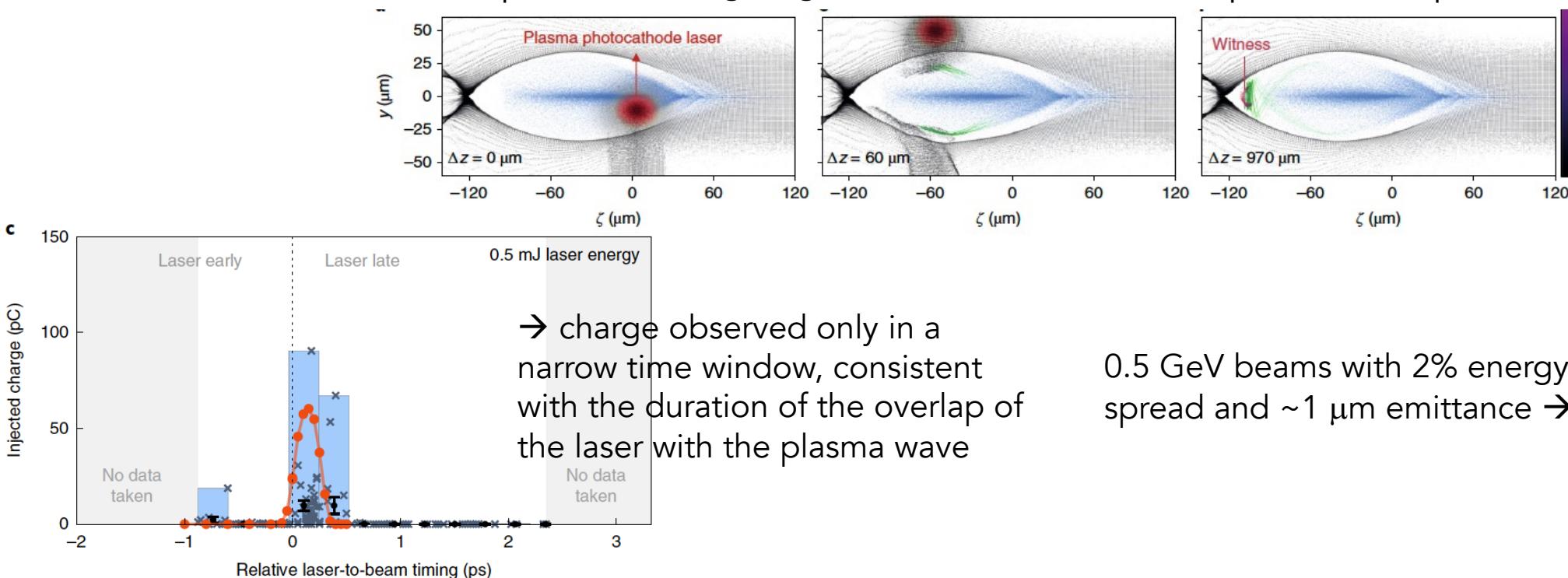
Plasma photocathode in a PWFA: experiments ongoing at SLAC



B. Hidding et al., PRL (2012)

- Use beam diver to excite plasma wave + ionization laser behind driver at low laser intensities resulting in lower beam transverse emittance ("Trojan horse" scheme)
- 10s of nm emittance observed in PIC simulations

Experiments ongoing at SLAC – FACET (completed) & II (planned)



Conclusions

- Plasma-based acceleration of electron beams:
 - Compact (>10 GV/m) source of 100 MeV to multi-GeV electron beams
 - Bunch length: ~few fs
 - Sub-micron emittance
 - Brightness: $B_{5D} \sim 10^{18}$ A/(m-rad) 2
- Injection method enables control of initial beam phase space
- Path to higher brightness – plasma accelerators using laser-triggered injection (plasma photocathode):
 $10s$ pC charge + $10s$ nm emittance and potential to reach $B_{5D} \sim 10^{20}$ A/(m-rad) 2

Thank you for the attention!