

Dark matter indirect detection limits including complete annihilation patterns

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DARK MATTER

85% of the total matter of our Universe

Relic density observed experimentally by Planck:

$$\Omega_{\chi} h^2 \simeq 0.1200 \pm 0.0012$$

Ref: Ade et al. 2016, *Astrophys. J.* 594, A13

Its identification would **reveal new Physics**

Proving its existence and nature would **improve our understanding** of the Universe

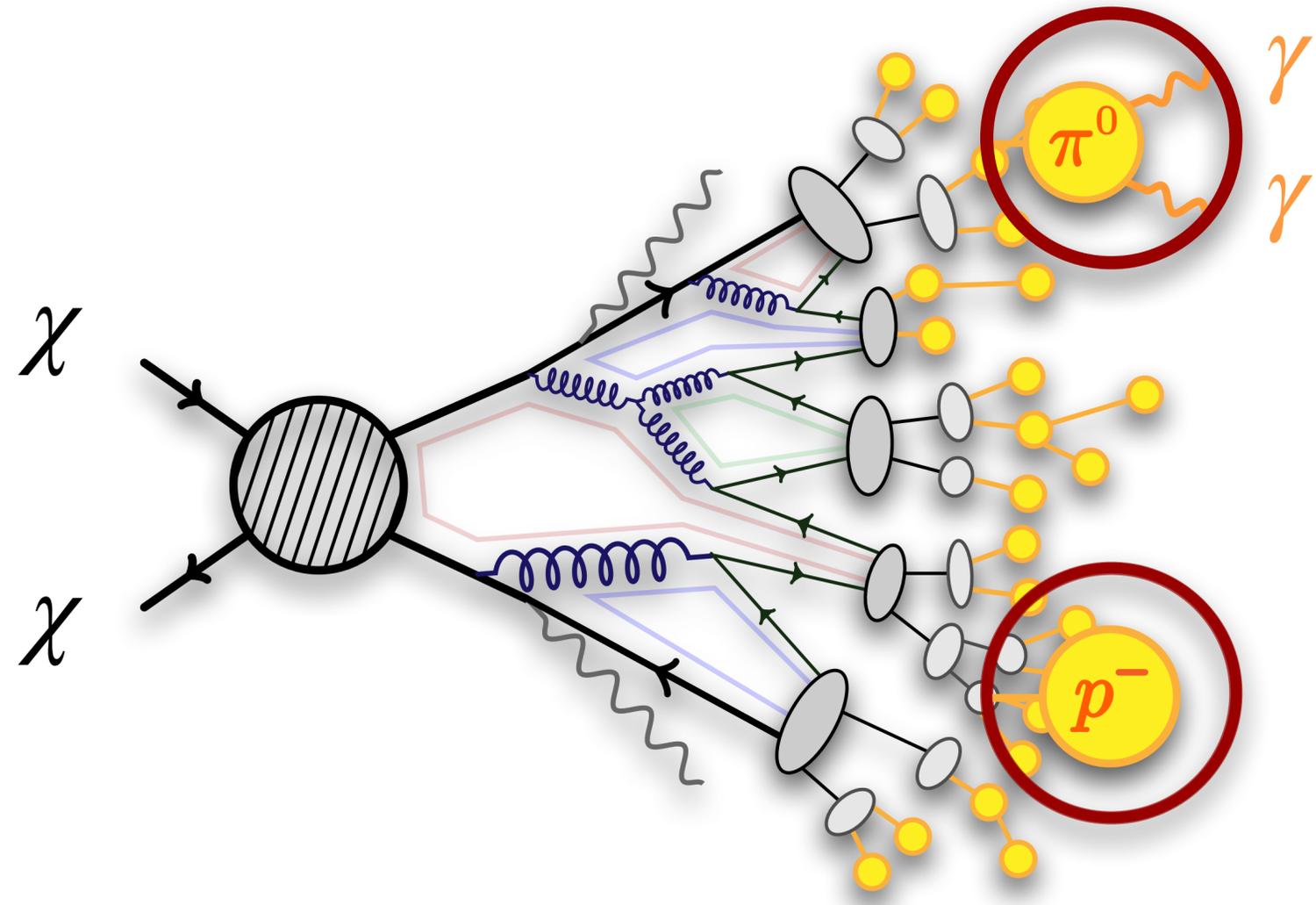
GOALS

Study of the impact of a more complete particle model

New prediction of DM upper limits with CTA mockdata of Sculptor

- **Previously:** use of individual annihilation channels
- **This work:** Collaboration with a theoretician to include a more complex and more complete model

INDIRECT SEARCHES



Dark Matter (DM)
annihilation



Standard Model particles
(bosons, quarks, leptons)



Final state products
such as γ rays

INDIRECT SEARCHES

Expected γ -ray flux from DM annihilation

Astrophysical
J factor

$$\frac{d\Phi(\langle\sigma v\rangle, J)}{dE} = \frac{1}{4\pi} \frac{\langle\sigma v\rangle}{2m_\chi^2} \sum_f \text{BR}_f \frac{dN_f}{dE} \times \int_{\Delta\Omega} \int_{\text{los}} \rho_{\text{DM}}^2 ds d\Omega$$

Particle Physics
factor

where

$\langle\sigma v\rangle$ = annihilation cross-section

m_χ = DM particle mass

BR_f = branching ratio

dN_f/dE = differential spectrum

ρ_{DM} = DM density

STATISTICAL ANALYSIS

LOG-LIKELIHOOD RATIO TEST STATISTICS

$$\Lambda = -2 \ln \frac{\mathcal{L}_{H_0}}{\mathcal{L}_{H_1}} = -2 \ln \frac{\mathcal{L}(\langle \sigma v \rangle_0 | \hat{N}_B, \hat{J})}{\mathcal{L}(\langle \hat{\sigma} v \rangle, \hat{N}_B, \hat{J})}$$

Constrained
minimization

Global
minimization

Ref: Cowan et al, 2010
Eur.Phys.J.C71:1554,2011

$\langle \sigma v \rangle$

Parameter of interest

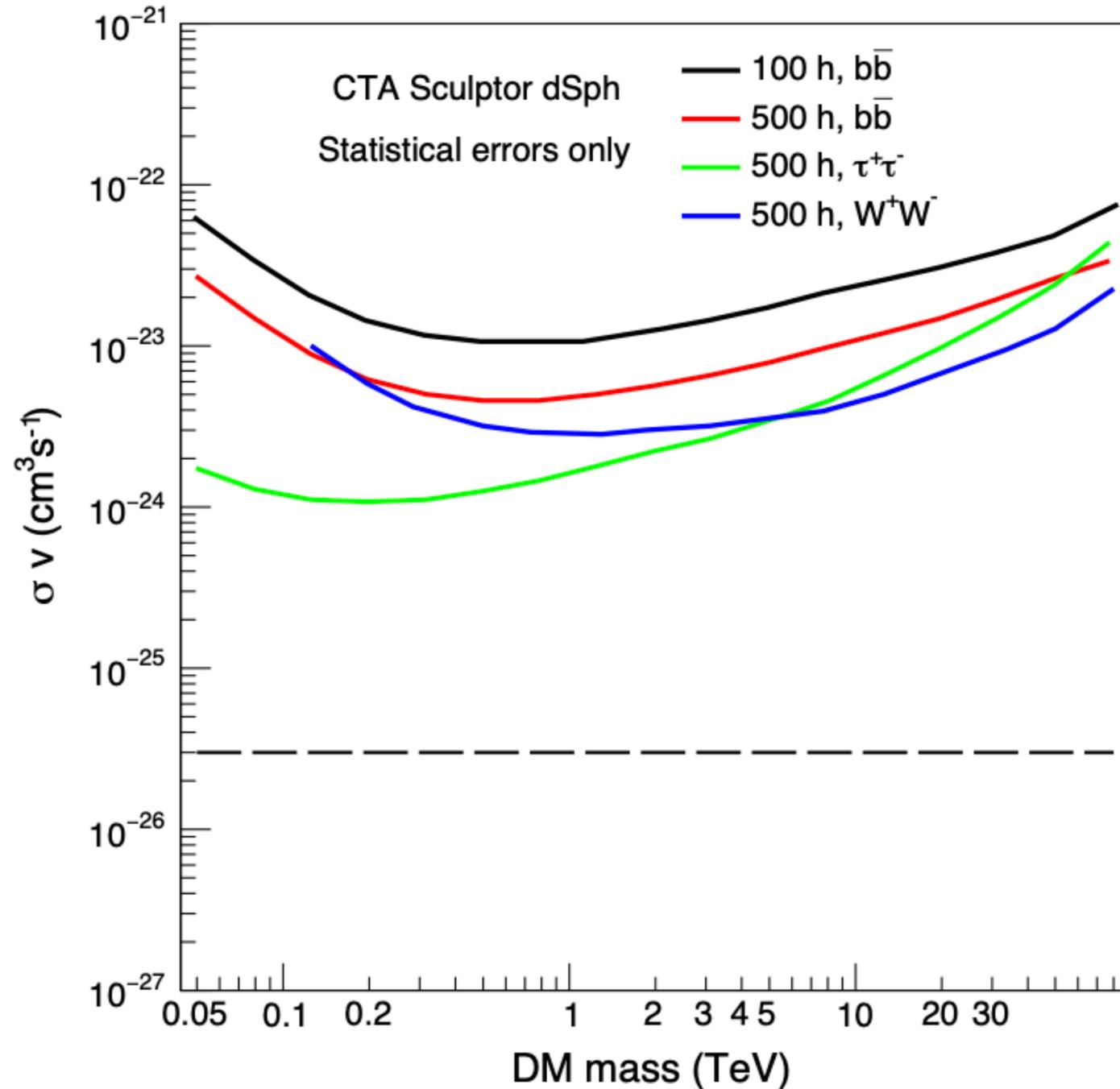
N_B, J

Nuisance parameters

Λ

2.71 at 95% Confidence Level

UPPER LIMITS



- Each annihilation channel **treated independently**
- Corresponding to a **branching ratio of 100%**
- **Simplest** model possible where all DM particles annihilate through the same channel

WHAT IF

We change the particle physics model?

SINGLET SCALAR DARK MATTER

Standard model extended by an additional scalar field (DM)

$$V_{\text{scalar}} \supset 2\lambda_H v^2 h^2 + \frac{1}{2}\mu_S^2 S^2 + \frac{1}{4}\lambda_{SH} v^2 S^2 + \frac{1}{4}\lambda_{SH} v S^2 h + \lambda_{SH} S^2 h^2$$

DM mass

$$m_S^2 = \mu_S^2 + \frac{1}{2}\lambda_{SH} v^2$$

DM – Higgs interaction

("Higgs portal")

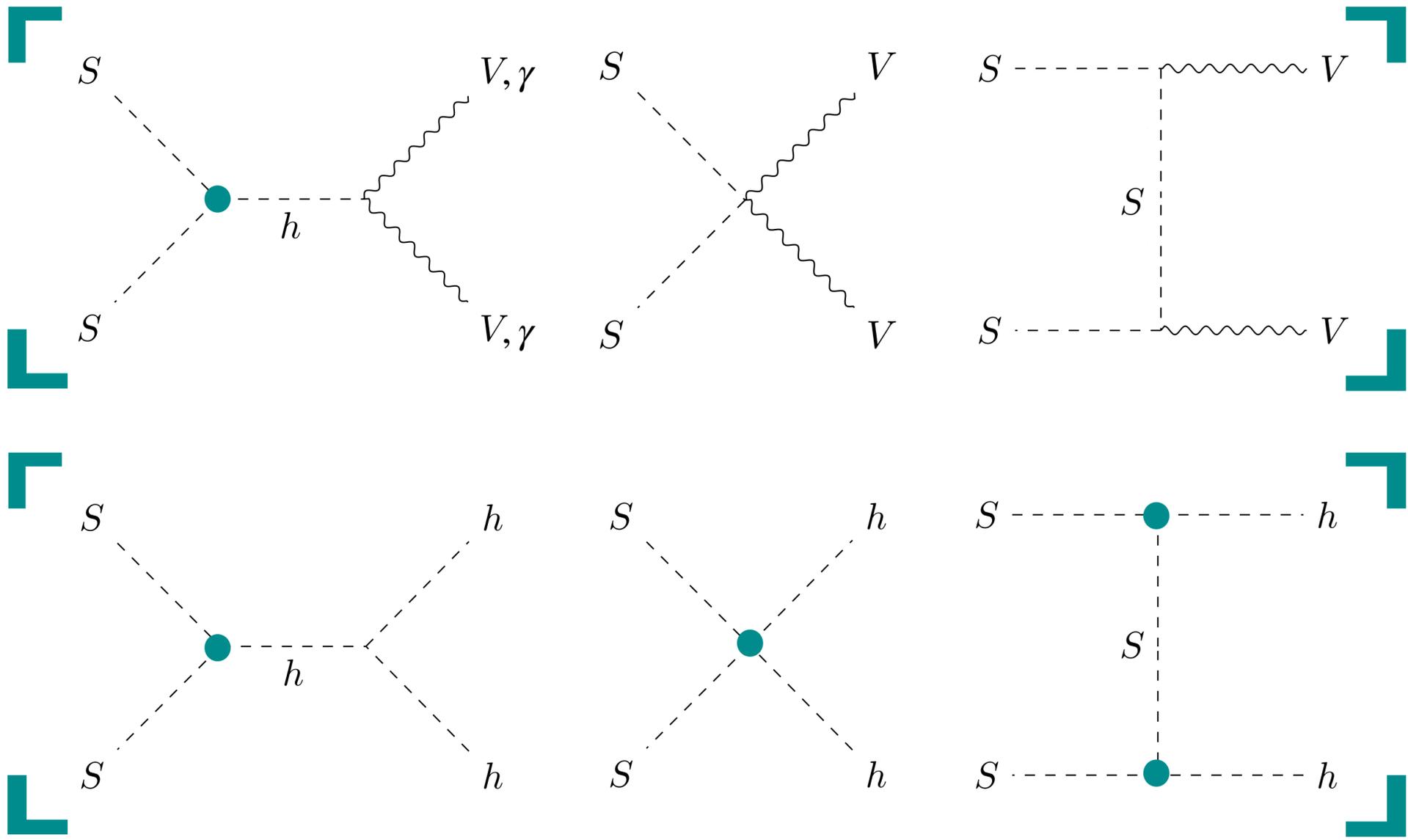
Phenomenology governed by

m_S (DM mass)

λ_{SH} (DM coupling)

SINGLET SCALAR DARK MATTER

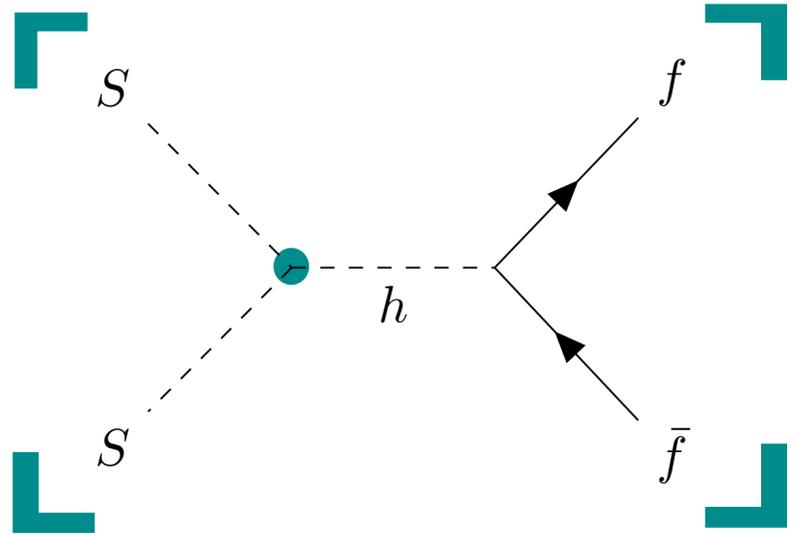
Possible dark matter annihilation channels (DM relic density + indirect detection)



Gauge boson final states

$$V = Z^0, W^\pm$$

Higgs boson final states



Quark or lepton final states

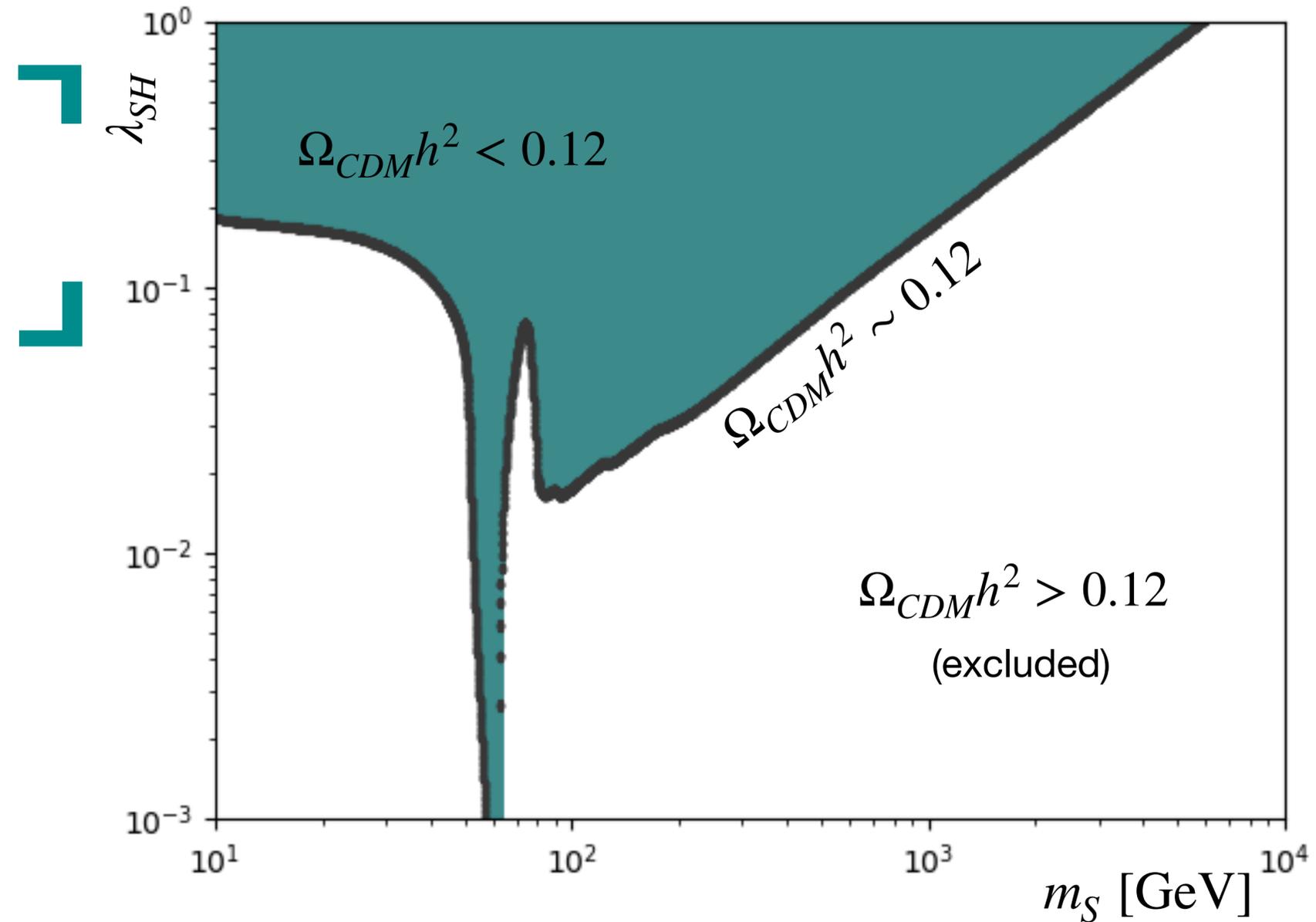
$$f = u, d, c, s, b, t, e, \mu, \tau$$

SINGLET SCALAR DARK MATTER

DM coupling vs DM mass

Relic density and branching ratio grid
computed using micrOMEGAs

Ref: Bélanger, Pukhov et al. 2003 - 2022



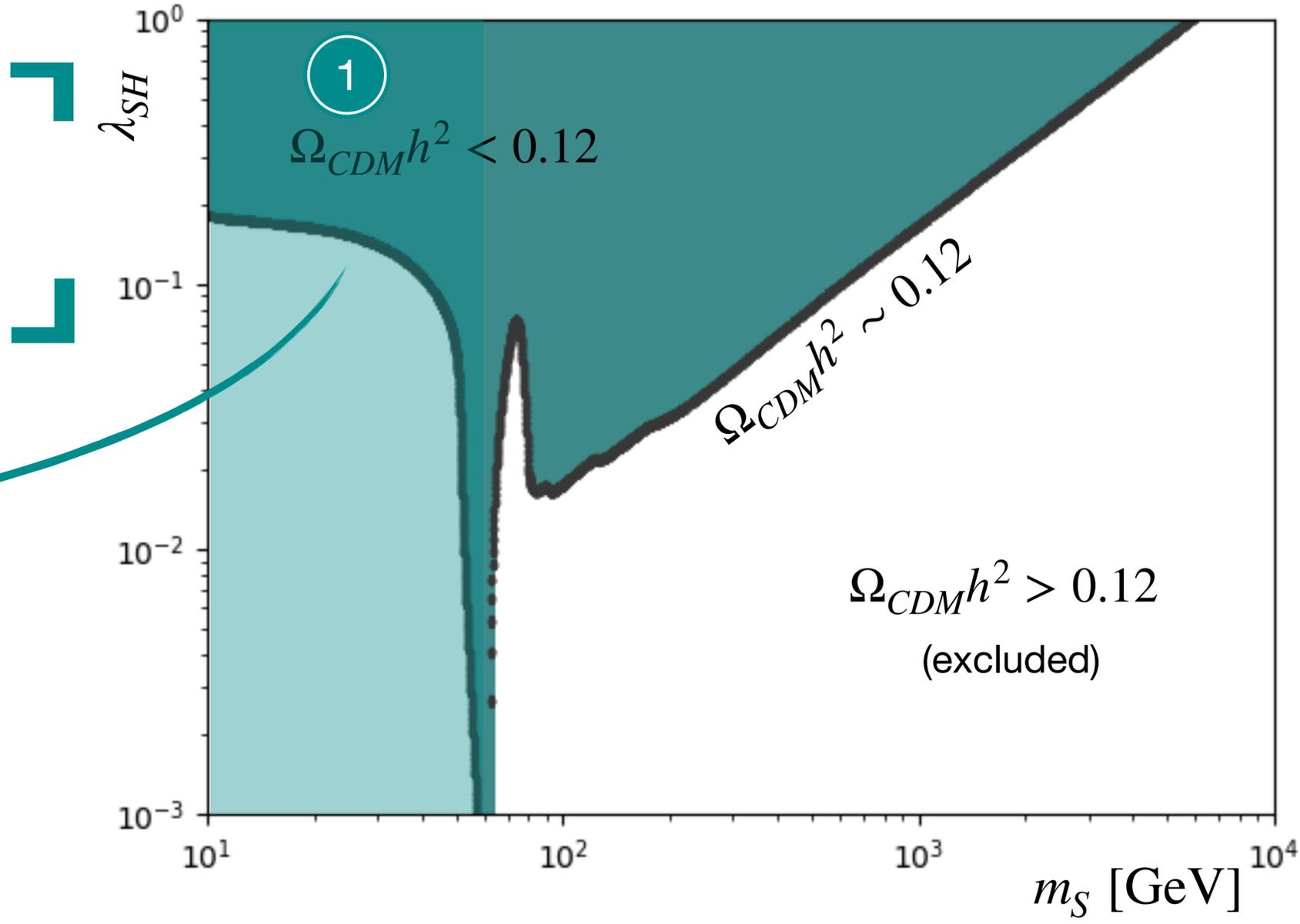
SINGLET SCALAR DARK MATTER

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$SS \rightarrow b\bar{b} \sim 90\%$
 $SS \rightarrow \tau^+\tau^- \sim 10\%$

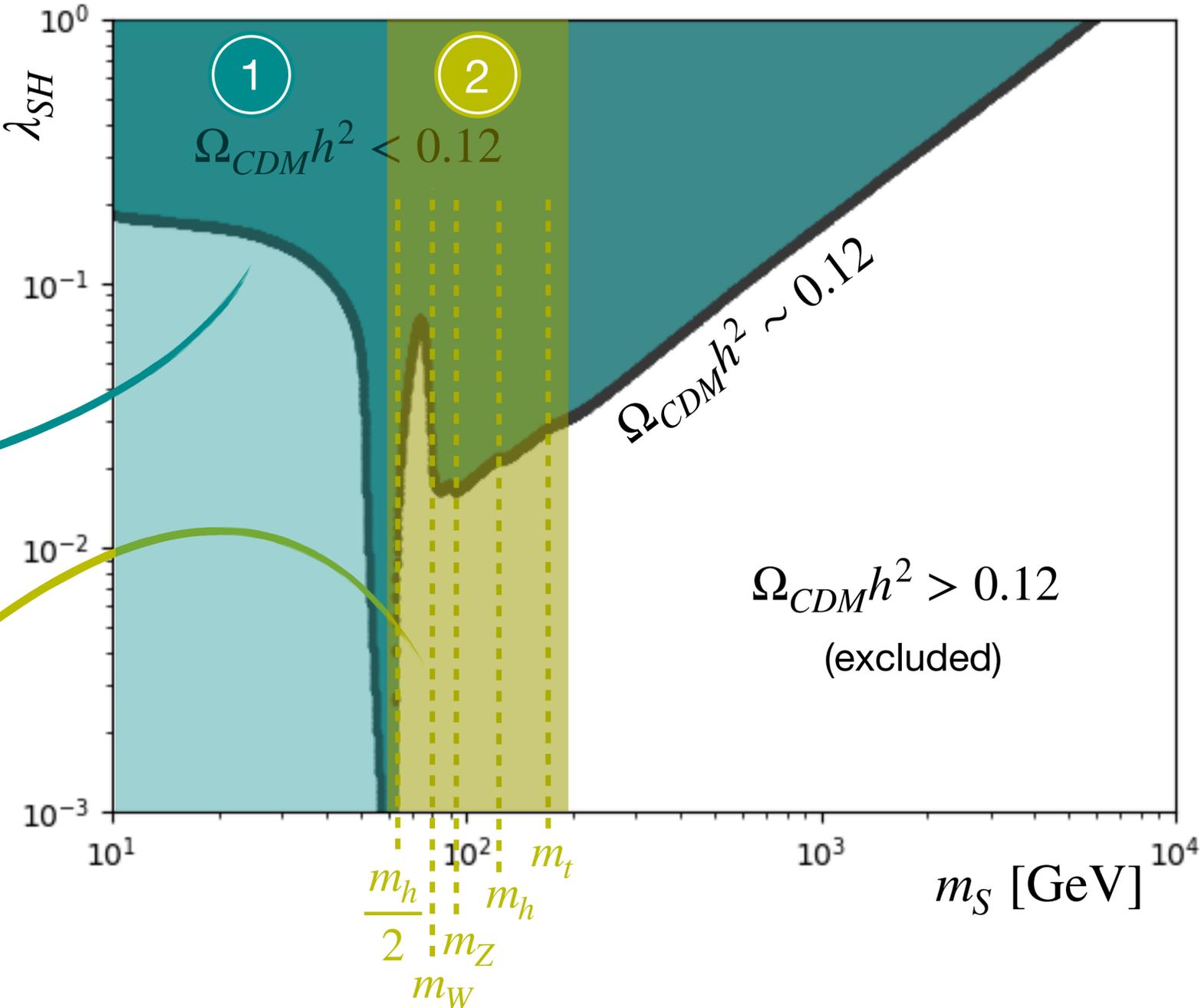


SINGLET SCALAR DARK MATTER

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Region of frequent change of the dominant annihilation channel

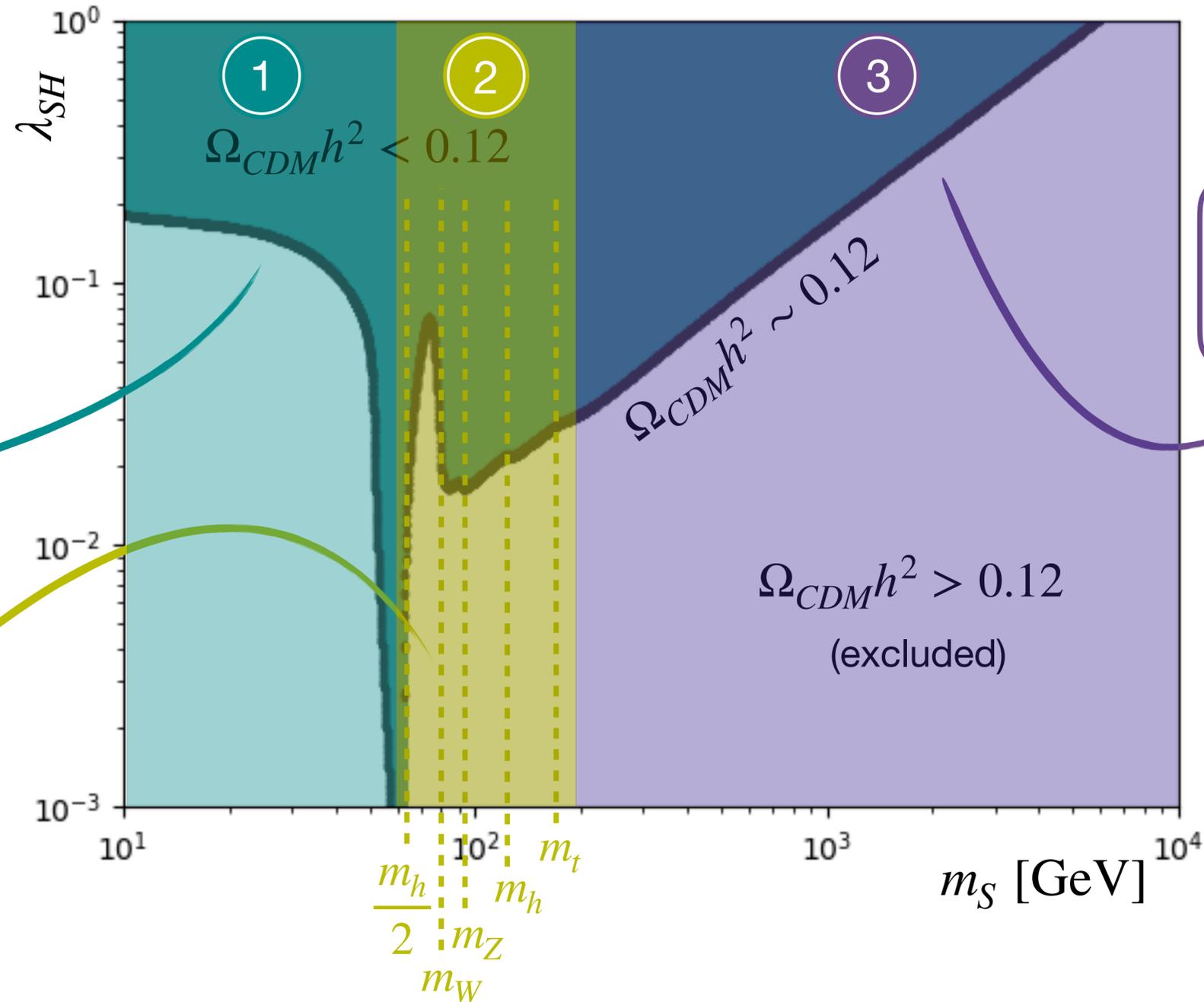
New annihilation channels open up, Higgs resonance at $m_h/2$

SINGLET SCALAR DARK MATTER

DM coupling vs DM mass

Relic density and branching ratio grid computed using micrOMEGAs

Ref: Bélanger, Pukhov et al. 2002 - 2022



$SS \rightarrow b\bar{b} \sim 90\%$
 $SS \rightarrow \tau^+\tau^- \sim 10\%$

$SS \rightarrow W^+W^- \sim 50\%$
 $SS \rightarrow Z^0Z^0 \sim 25\%$
 $SS \rightarrow h^0h^0 \sim 25\%$

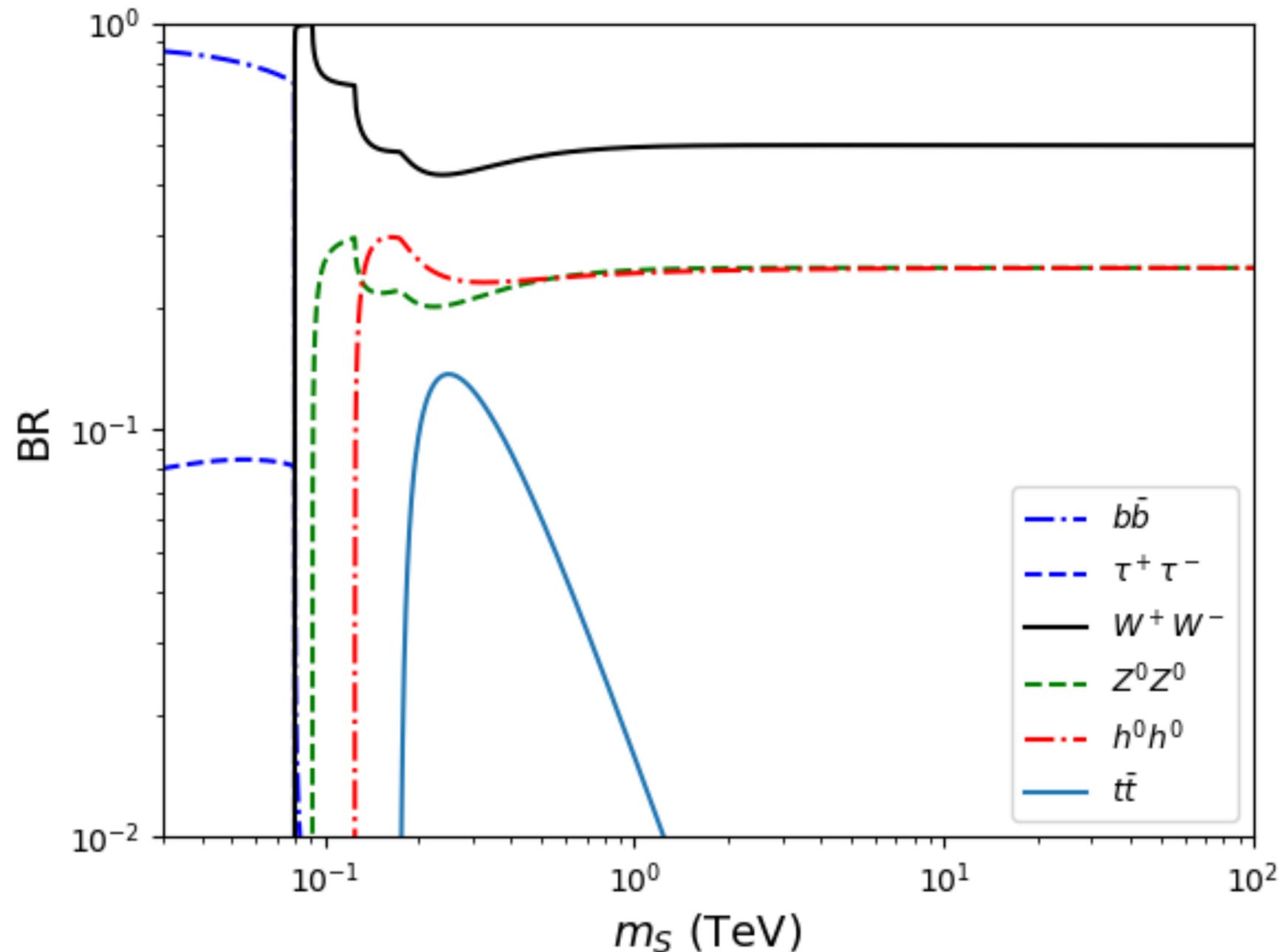
Region of frequent change of the dominant annihilation channel

New annihilation channels open, Higgs resonance at $m_h/2$

All annihilation channels **treated all together** whose **branching ratio varies** with respect to the DM mass

SINGLET SCALAR DARK MATTER

Branching Ratio according to the relic density constraint



None of the annihilation channels are at 100% branching ratio over the full mass range

For the remaining part, we focus on the case where the relic density constraint is satisfied (black line in previous figure):

$$\Omega_\chi h^2 \simeq 0.1200 \pm 0.0012$$

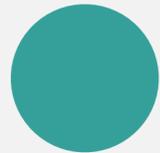
Even in such a simple setup, the
“100% hypothesis” is not justified...

More complex models invoke an even
richer phenomenology...

TARGET SOURCE



Dwarf galaxy selected for the
CTA dark matter program



South Hemisphere
 $l = 287.62^\circ, b = -83.16^\circ$



J factor
 $\text{Log}_{10} J_{0.1} = 18.3 \pm 0.3$



Ref: Bonnivard et al, 2015 ApJ 808 L3

Sculptor

Mock data prepared
with Gammapy 0.18.2

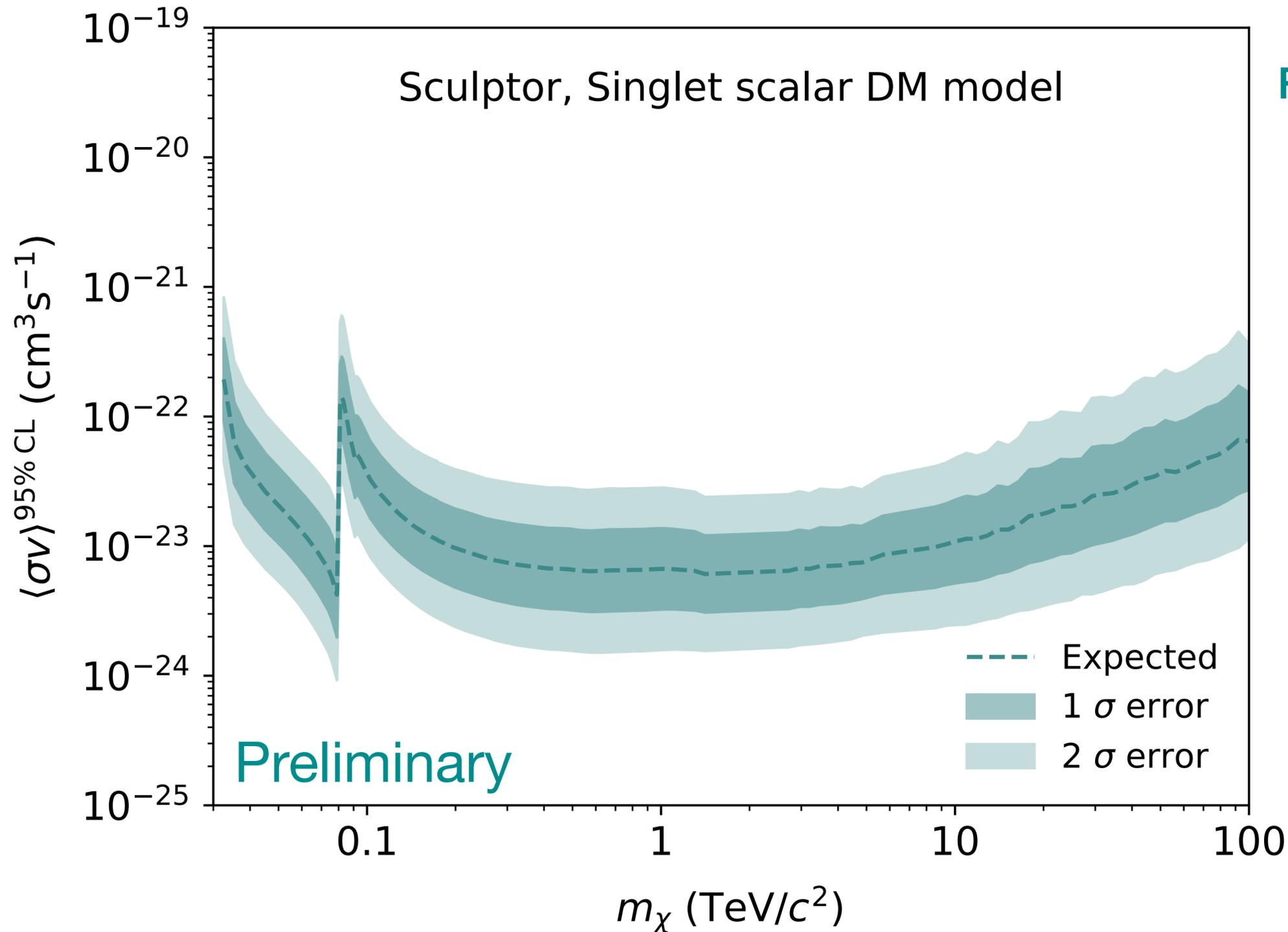
Simulated events for 500h of
observation at 20° zenith angle

NEW UPPER LIMITS

Computation of the predicted DM cross section
VS
DM particle mass

- 1 **Expected** limits - Sample of **300 Poisson realizations** of the simulated background events
- 2 **Mean expected limits**
Mean of the derived $\langle\sigma v\rangle$ distribution
- Statistical uncertainty bands**
Standard deviation at 1 and 2σ

RESULTS



Predicted upper limit and uncertainties

Assuming a singlet scalar DM model
 γ -ray spectra taken from Cirelli et al.
JCAP 03 (2011) 051

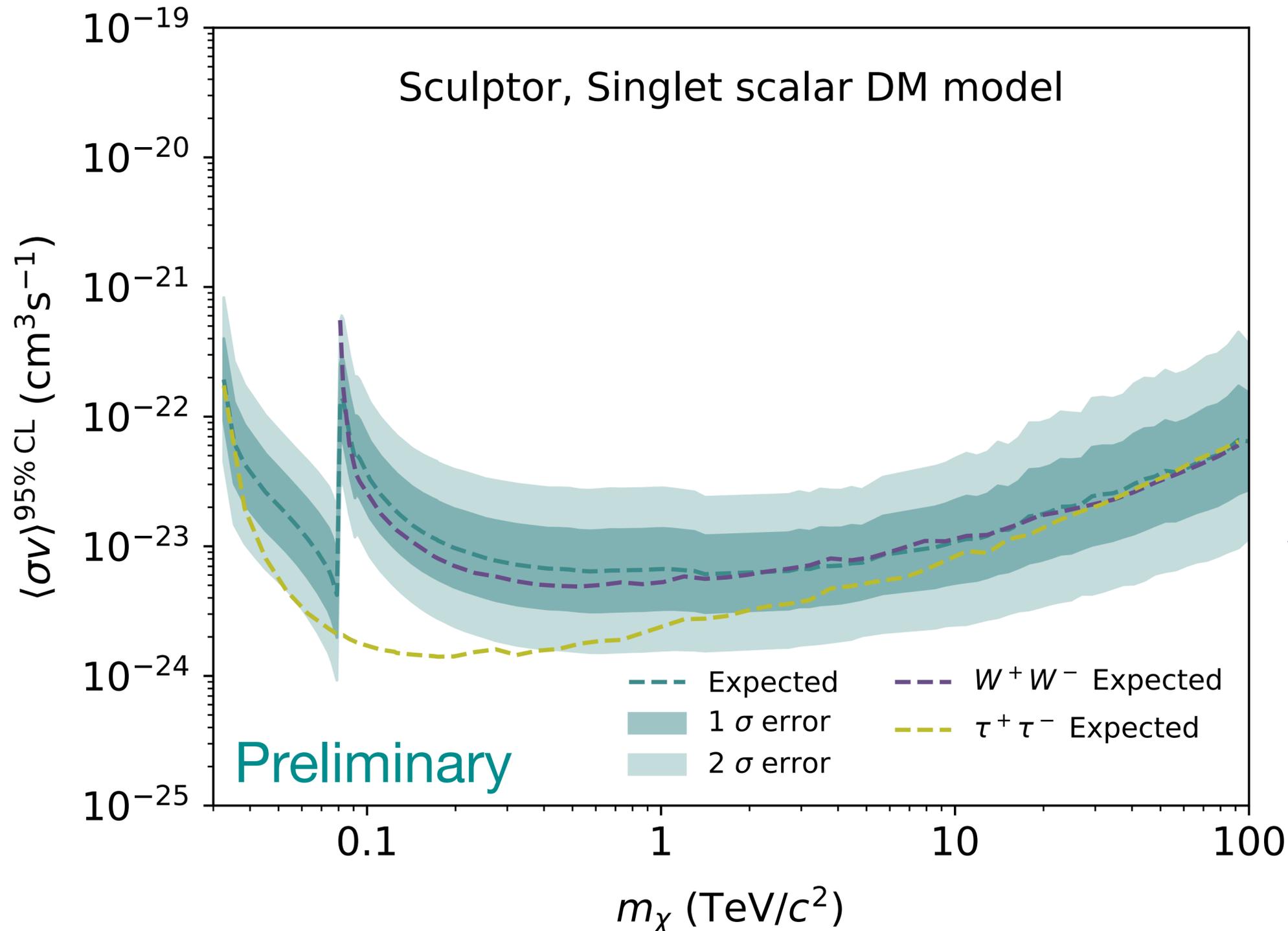
Inflection point

Due to the Higgs resonance

Sudden increase

Due to the opening of the WW channel

COMPARISONS



SINGLET SCALAR MODEL VS 100% W^+W^-

More conservative limit with the singlet scalar DM model

Below the W mass

No upper limit for 100% WW since the WW channel **does not exist**

$\sim 0.1-1$ TeV

Slight difference between 100% WW and singlet scalar DM model

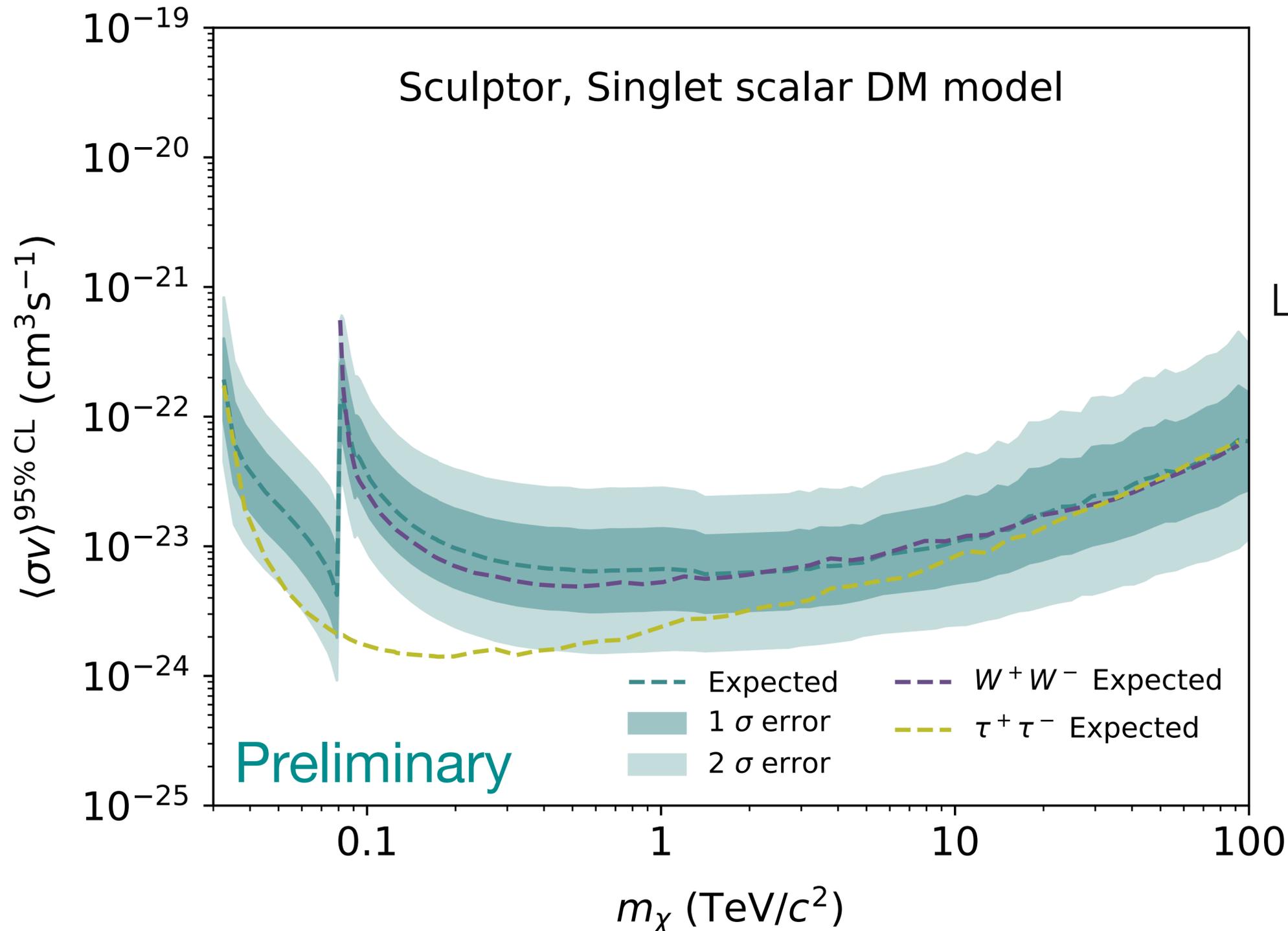
Additional contributions: ZZ, hh, tt

Above 1 TeV

Stability with $\sim 50\%$ WW - 25% ZZ - 25% hh
Limits similar to the 100% WW case since WW, ZZ, hh lead to similar γ -ray spectra

COMPARISONS

SINGLET SCALAR MODEL
VS
100% $\tau^+\tau^-$



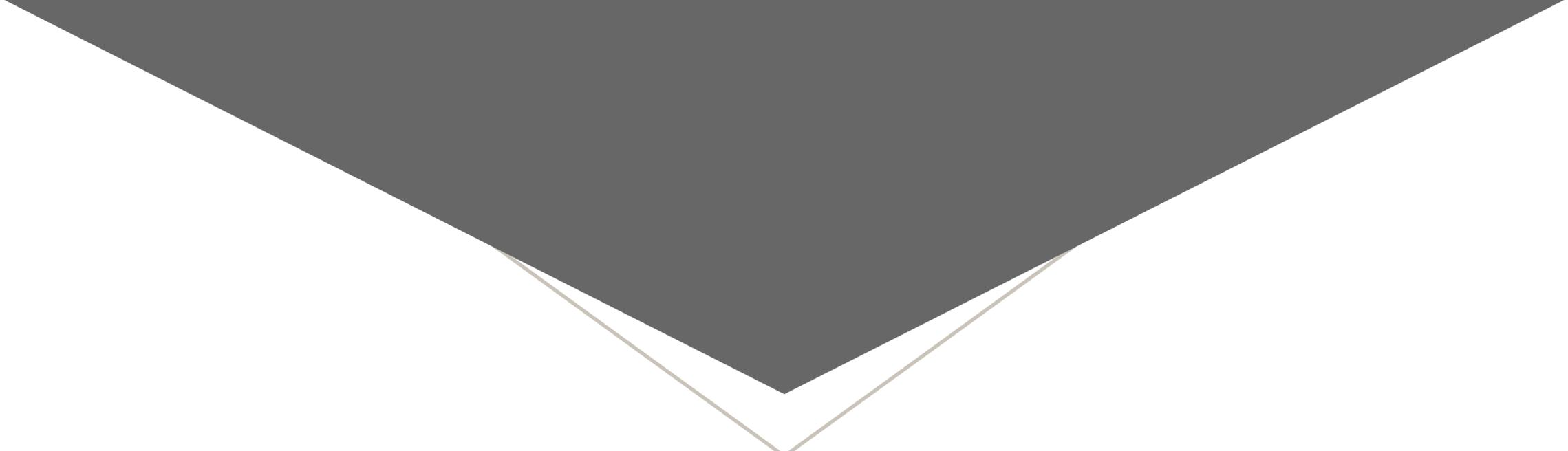
100% $\tau^+\tau^-$ produces more γ rays
Leads to **more constraining** upper limits

However, in the singlet scalar model,
this $\tau^+\tau^-$ channel is never dominant

100% $\tau^+\tau^-$ = over estimation of the contribution

CONCLUSION & PERSPECTIVES

- Use of a **more complex and more complete** particle physics model
- Takes into account the **full phenomenology** with all annihilation channels at once
- **Change of dominant annihilation channel(s)** along with the DM particle mass
- **Affects** the predicted upper limits
- Feature can be **expected in any particle physics model**
- Derivation of a **predicted upper limit and its 1σ and 2σ uncertainty bands** over the energy range of CTA
- Particle physics model could be used as well **on the future data of CTA**
- Paper in preparation



Thanks for your attention



STATISTICAL ANALYSIS

Total likelihood

$$\mathcal{L}(\langle \sigma v \rangle, N_B, J) = \prod_{i=1} \mathcal{L}_{P_i}(\langle \sigma v \rangle, N_{B_i}, J | N_{\text{ON}_i}, N_{\text{OFF}_i}, \alpha) \mathcal{L}^J(J | \bar{J}, \sigma_J)$$

Poisson likelihood

Log-normal likelihood

Poisson likelihood for each energy bin

$$\mathcal{L}_i^P = \frac{(N_{S_i} + N_{B_i})^{N_{\text{ON}_i}}}{N_{\text{ON}_i}!} e^{-(N_{S_i} + N_{B_i})} \cdot \frac{(\alpha N_{B_i})^{N_{\text{OFF}_i}}}{N_{\text{OFF}_i}!} e^{-\alpha N_{B_i}}$$

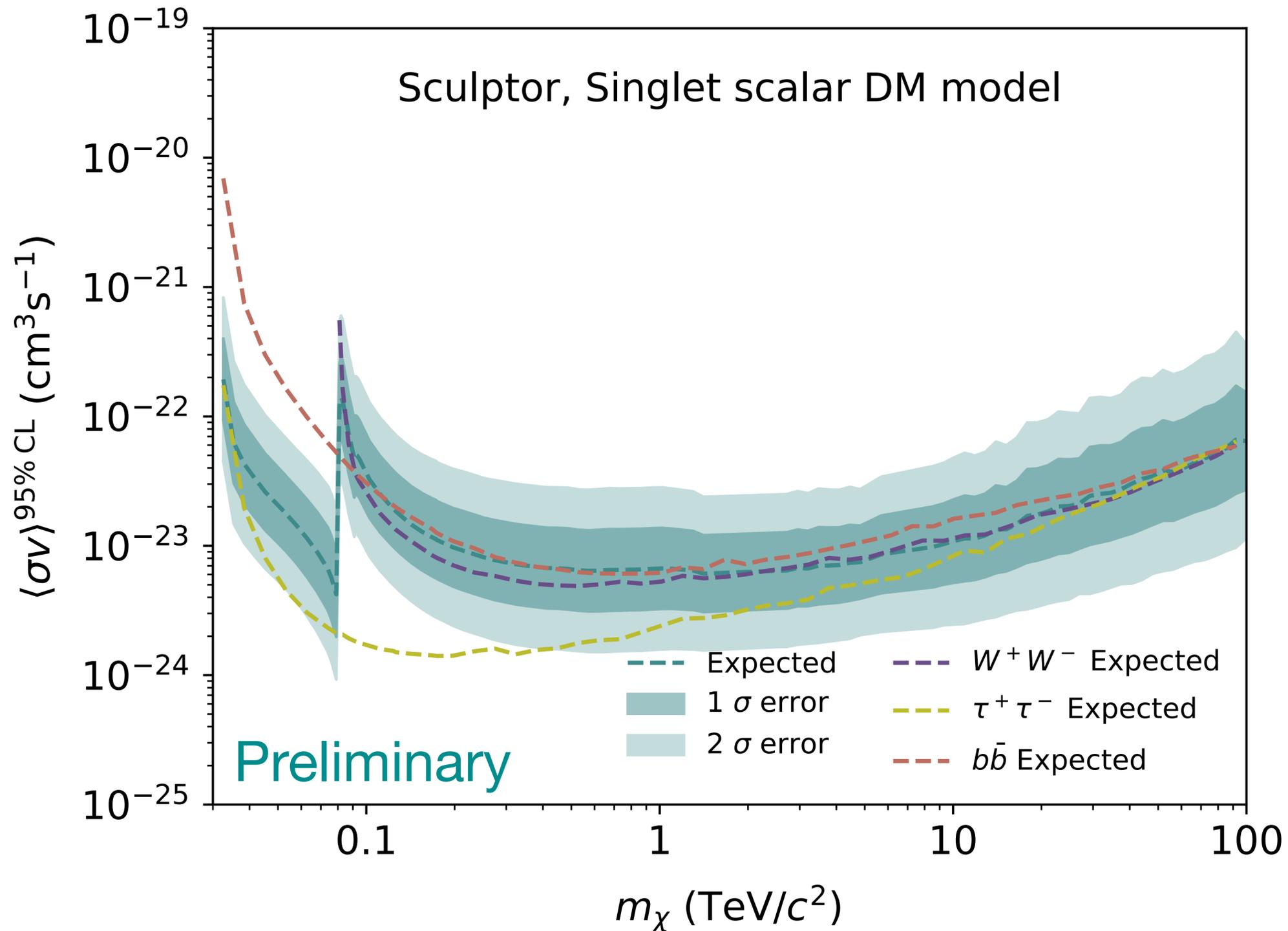
ON REGION

OFF REGION

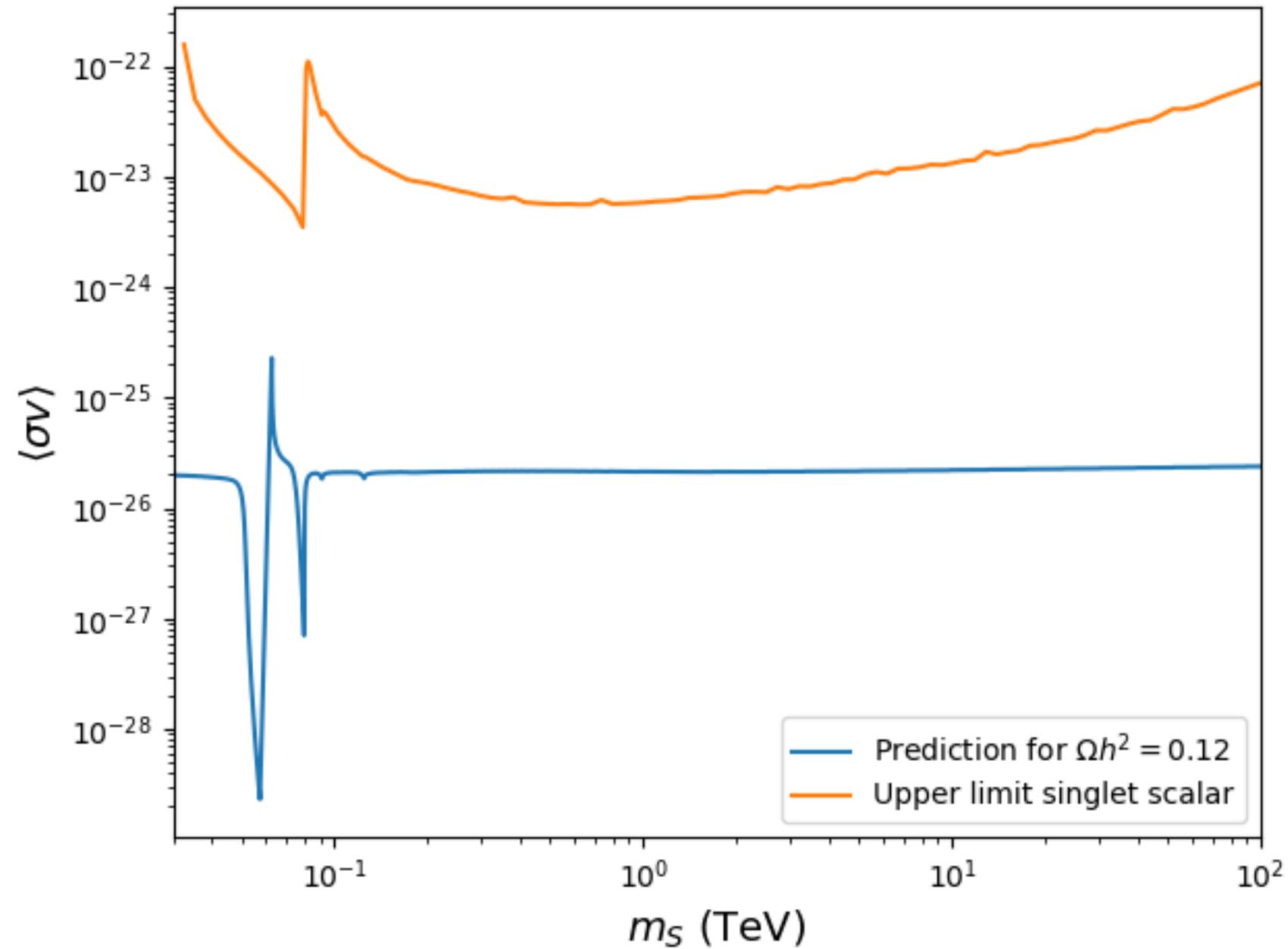
Log-normal likelihood to model the uncertainties of the J factor

$$\mathcal{L}^J = \frac{1}{\ln(10)\sqrt{2\pi}\sigma_J J} \exp - \frac{(\log_{10} J - \log_{10} \bar{J})^2}{2\sigma_J^2}$$

COMPARISONS



COMPARISONS



COMPARISONS

Ref: Cirelli et al. *JCAP* 03 (2011) 051