

3d $N=2$ dualities for SQCD with D-type superpotential

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Based on 2204.06961 w/ Simone Rota
(see also 1809.03796, 1509.02199 & 1409.8623)

Motivations

SUSY dualities: why are they still important/useful?

- Links with mathematics
- Non perturbative dynamics
- Conformal window and AdS/CFT
- 3d bosonization
- Connections among dualities in different dimensions
- SUSY and global symmetry enhancements
- Relating quantum aspects in different dimensions (e.g. role of the monopoles)

Plan

- 3d reduction of 4d dualities
- Basic dualities and basic web
- The web for the A-type superpotential
- D-type
 - 4d aspects
 - The web
 - Three sphere partition function
- Conclusions

4d/3d reduction

This work fits in the program of reduction of 4d Seiberg-like dualities to 3d (ARSW, [Aharony-Razamat-Seiberg-Willett `13](#)).

Original case: SU(N) SQCD with F flavors and $F > N+1$ ([Seiberg `94](#))

Then extension to SO/USp SQCD ([Intriligator-Seiberg `95](#) and [Intriligator-Pouliot `95](#)) and to cases w/ one two-index tensor (Adj, Symm. or Anti-Symm) and A-type superpotential ([Kutasov-Schwimmer-Seiberg `95](#), [Intriligator-Leigh-Strassler `95](#)).

The ARSW procedure has been (or can be) applied to all these cases. Often the reduction can be studied with the aim of localization.

D-type superpotential? Curious, uncommon and more delicate dualities obtained by [Brodie `96](#) and [Brodie-Strassler `96](#).

$$\begin{array}{l}
 \widehat{O} \\
 \widehat{A} \\
 \widehat{D} \\
 \widehat{E} \\
 A_k \\
 D_{k+2} \\
 E_6 \\
 E_7 \\
 E_8
 \end{array}
 \qquad
 \begin{array}{l}
 W_{\widehat{O}} = 0 \\
 W_{\widehat{A}} = \text{Tr} Y^2 \\
 W_{\widehat{D}} = \text{Tr} X Y^2 \\
 W_{\widehat{E}} = \text{Tr} Y^3 \\
 W_{A_k} = \text{Tr}(X^{k+1} + Y^2) \\
 W_{D_{k+2}} = \text{Tr}(X^{k+1} + X Y^2) \\
 W_{E_6} = \text{Tr}(Y^3 + X^4) \\
 W_{E_7} = \text{Tr}(Y^3 + Y X^3) \\
 W_{E_8} = \text{Tr}(Y^3 + X^5).
 \end{array}$$

Arnold's ADE classification of singularities, which precisely coincides with the possible relevant deformation superpotentials, listed as A_k , D_{k+2} and $E_{6,7,8}$

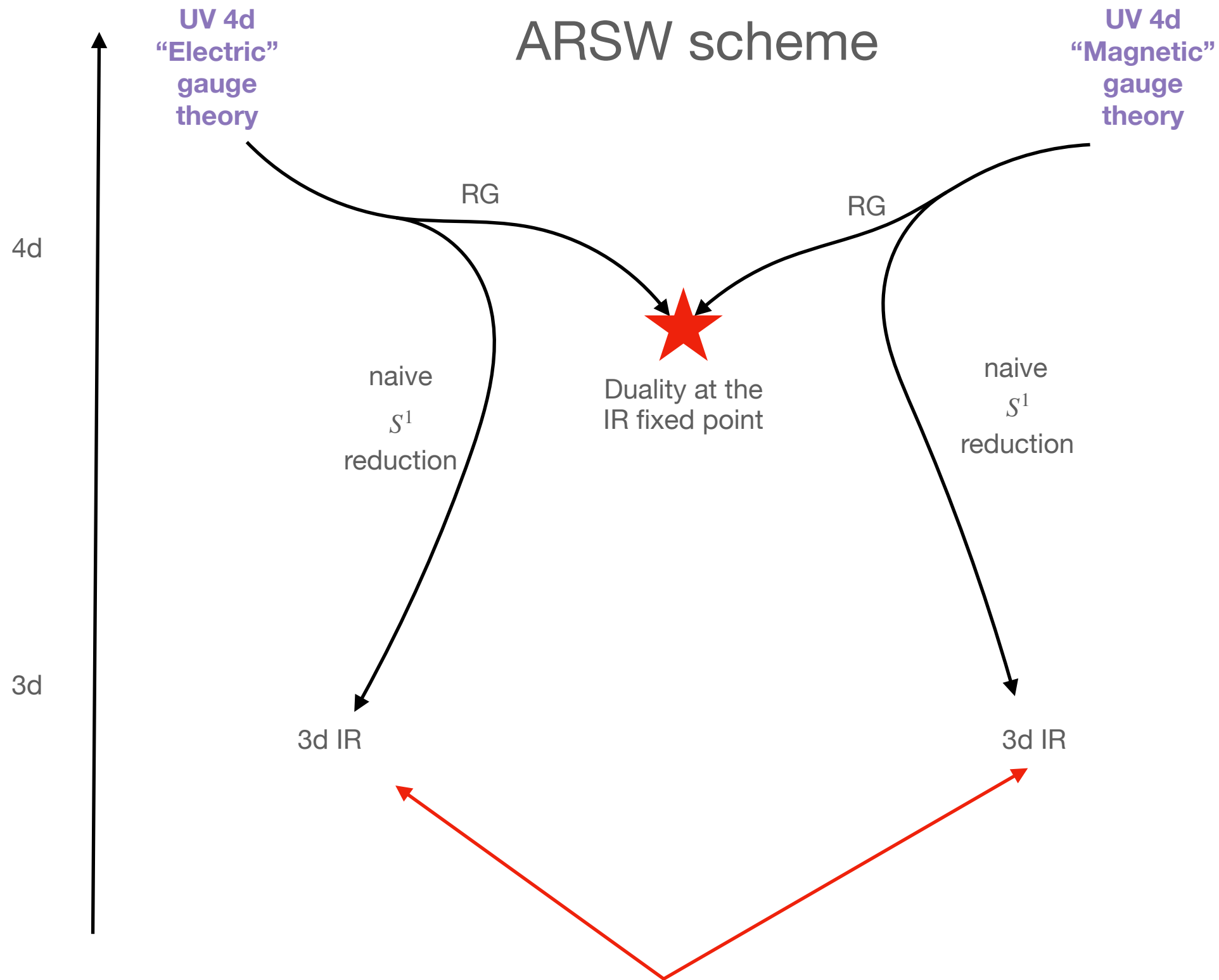
4d and 3d SUSY with four supercharges: a brief overview

4d

- Global Anomalies
- Absence of axial symmetry
- U(1) gauge groups are IR free
- Exact R by a-maximization
- SYM action
- $V = (A, \lambda, D)$
- $\Lambda_{holo} = \text{Exp}(1/g_4^2)$
- Instanton: $W_{eff} = N\Lambda_{holo}^3$

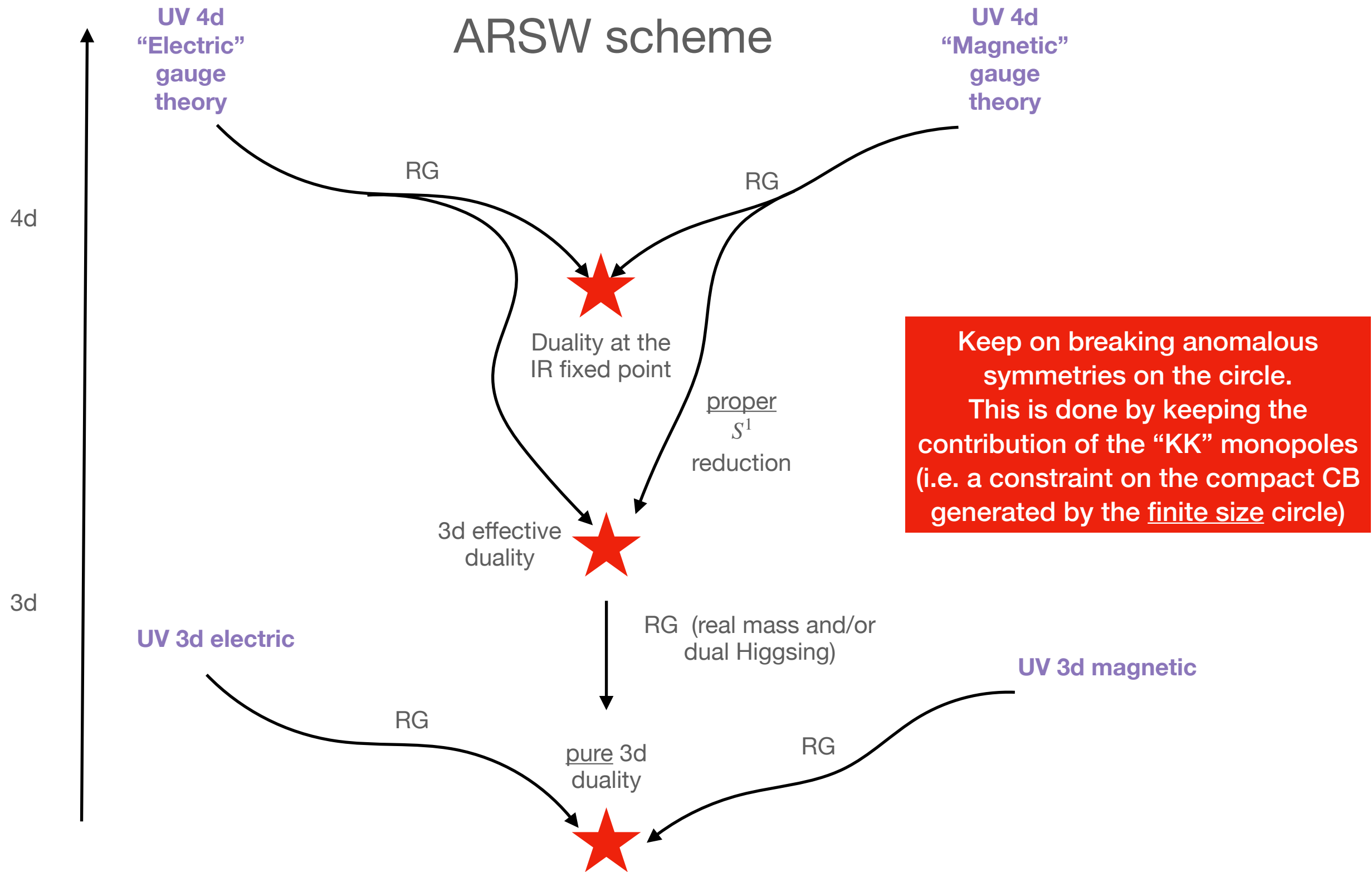
3d

- Parity Anomalies
- Possible axial symmetry (broken by monopoles)
- Interacting U(1) gauge groups in the IR
- Exact R by F-maximization
- SYM and CS action
- $V = (A, \lambda, D, \sigma)$ w/ $\sigma \propto A_4$ & $\varphi = d^* F$
- CB : $Y \propto \text{Exp}(\sigma/g_3^2 + i\varphi)$
- BPS Monopoles: $W_{eff} \propto \sum 1/Y_i$



4d Duality not preserved (in general) in 3d by a "naive" dimensional reduction
Reason: extra symmetries anomalous in 4d (e.g. axial in SQCD) can spoil the duality
(e.g. different IR mixing with the R-symmetry in the two phases - check w/ F-max)

ARSW scheme



Keep on breaking anomalous symmetries on the circle. This is done by keeping the contribution of the "KK" monopoles (i.e. a constraint on the compact CB generated by the finite size circle)

A bit naive: 3d RG does not necessarily eliminates all the finite size effects. There can be more sophisticated structures involving the CB.

4d U(N) **Seiberg '95**
duality

U(N) SQCD w/F+1 flavors
and $W=0$

U(F+1-N) dual SQCD w/F+1 flavors,
 F^2 singlets (meson M) and $W = Mq\tilde{q}$

Circle
reduction

3d U(N) **ARSW '13** effective
duality w/ KK monopoles

U(N) SQCD w/F+1 flavors
and $W = \eta Y_+ Y_-$

U(F+1-N) SQCD w/F+1 flavors
and $W = \tilde{\eta} \tilde{Y}_+ \tilde{Y}_- + Mq\tilde{q}$

Real masses
and
dual Higgsing

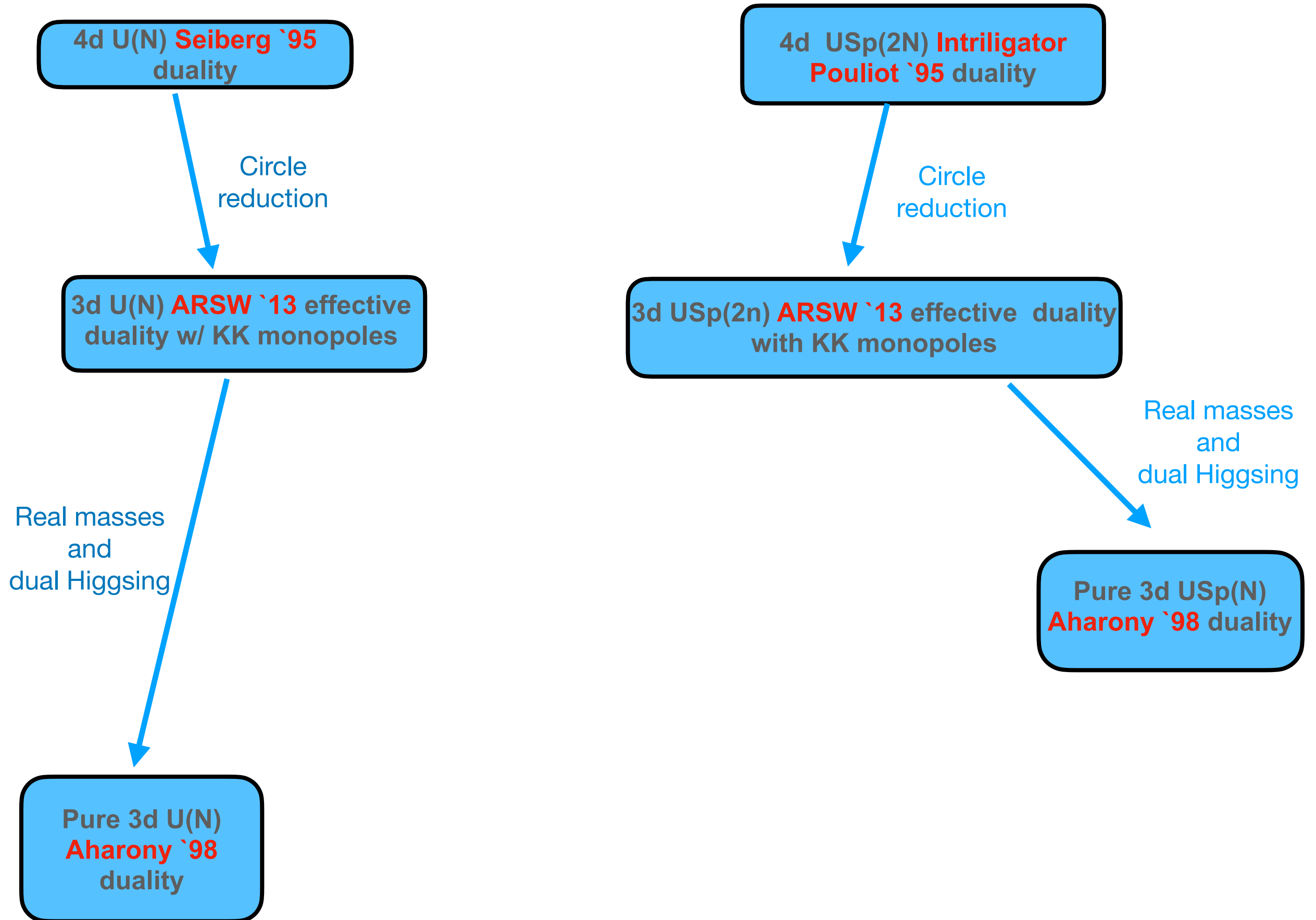
Pure 3d U(N)
Aharony '98
duality

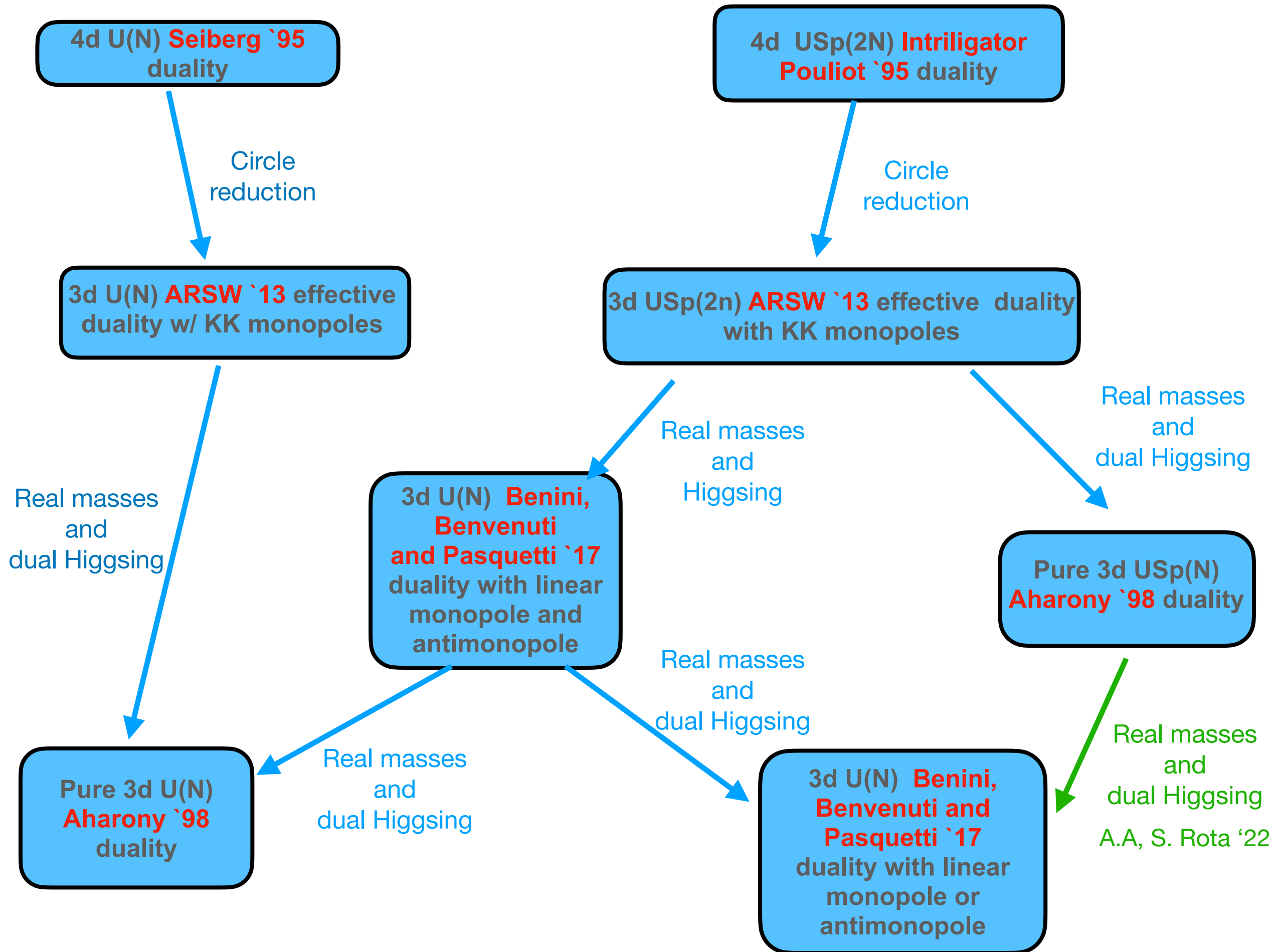
U(N) SQCD w/F flavors
and $W = 0$

U(F-N) w/F coupled with U(1) w/1

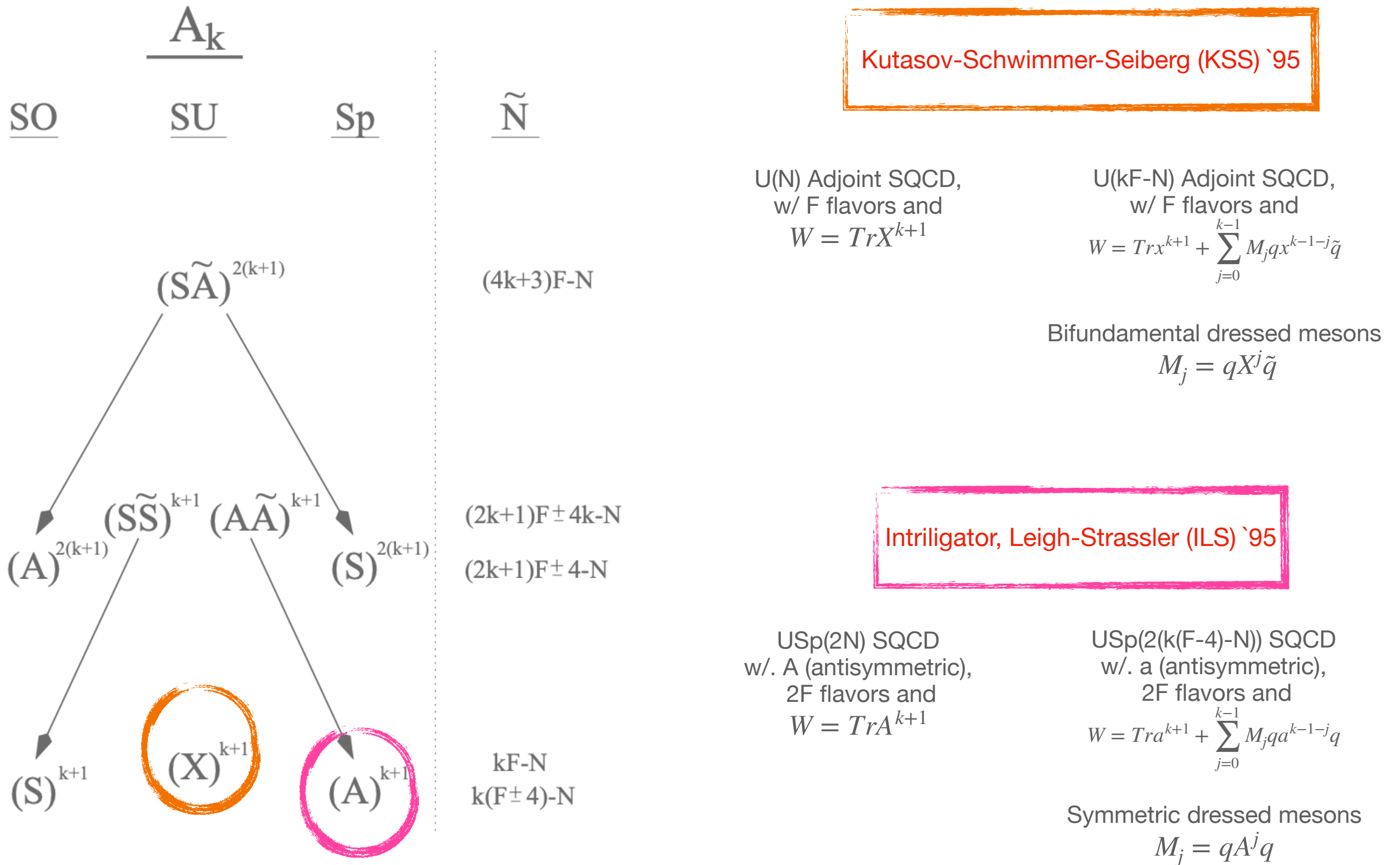
“Local”
Mirror
SQED - XYZ

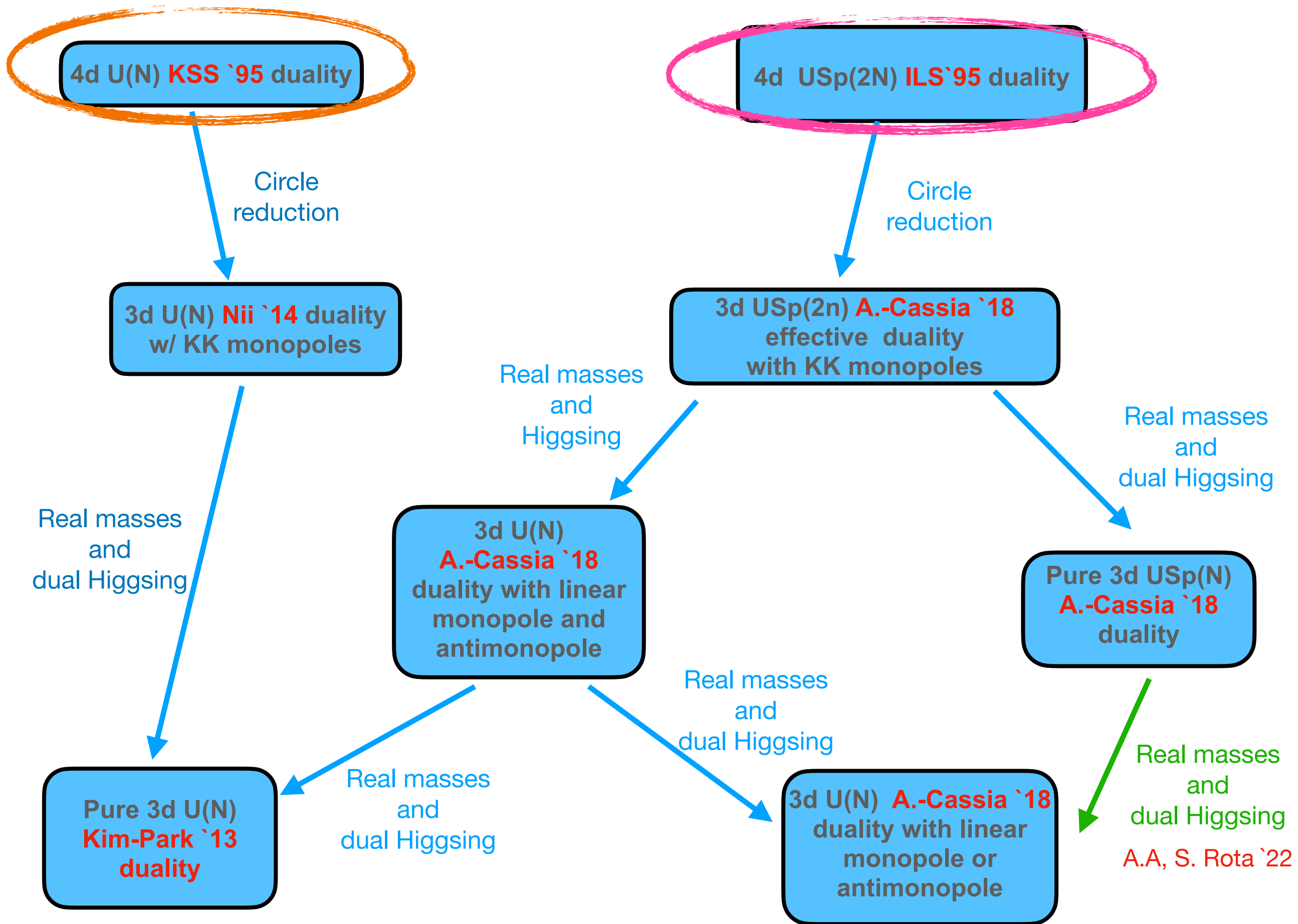
U(F-N) w/F flavors
 $W = Mq\tilde{q} + V_+ v_- + V_- v_+$





Let's borrow this picture from Brodie-Strassler '96, the orange and pink circles generalize Seiberg '94 and Intriligator-Pouliot '95 for the A_k type superpotential





D_{k+2}

SO

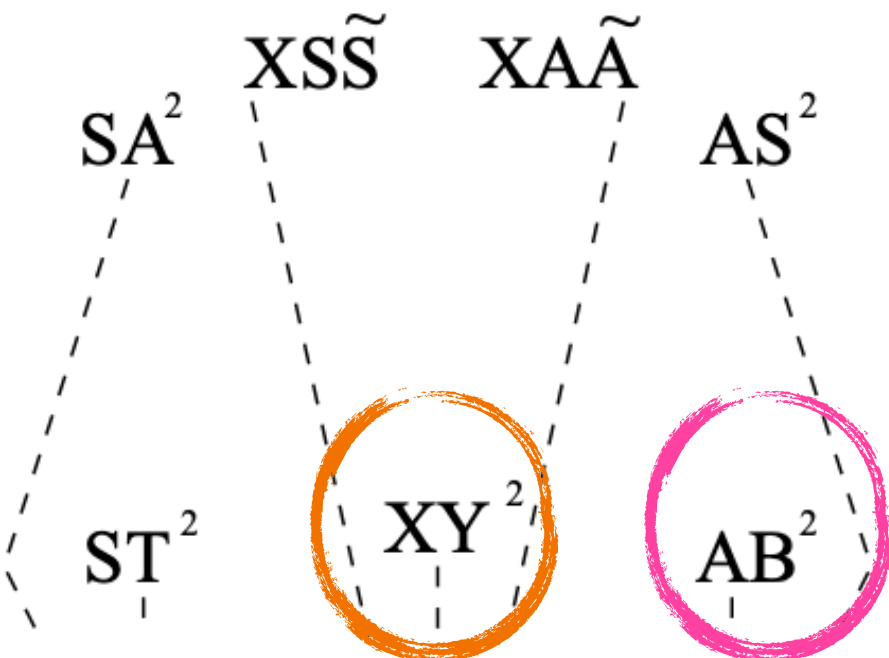
SU

Sp

$S^{k+1} +$

$X^{k+1} +$

$A^{k+1} +$



4d $SU(n)$ SQCD with F flavors Q and \tilde{Q} with two adjoints X and Y interacting through the superpotential

$$W = \text{Tr } XY^2 + \text{Tr } X^{k+1}$$

with k odd.

4d $SU(\tilde{n} = 3kF - n)$ SQCD with F dual flavors q and \tilde{q} with two adjoints x and y interacting through the superpotential

$$W = \text{Tr } xy^2 + \text{Tr } x^{k+1} + \sum_{j=0}^{k-1} \sum_{\ell=0}^2 \text{Tr } \mathcal{M}_{j,\ell} q x^{k-1-j} y^\ell \tilde{q}$$

where the singlets $\mathcal{M}_{j,\ell}$ correspond under the duality map to the dressed mesons $QX^jY^\ell\tilde{Q}$ for $j = 0, \dots, k-1$ and $\ell = 0, 1, 2$ of the electric phase.

$Usp(2n)$ SQCD with $2f$ fundamentals Q and two rank-two antisymmetric tensors A and B interacting through a superpotential

$$W = \text{Tr } A^{k+1} + \text{Tr } AB^2$$

$USp(2\tilde{n} = 2(3kf - n - 4k - 2))$ SQCD with $2f$ fundamentals q , two rank-two antisymmetric tensors a and b and the dressed mesons $M_{rs}^{(j,\ell)} = Q_r A^j B^\ell Q_s$ symmetric (antisymmetric) for $j\ell$ odd (even) with $j = 0, \dots, k-1$ and $\ell = 0, 1, 2$. The superpotential of the dual phase is

$$W = \text{Tr } a^{k+1} + \text{Tr } ab^2 + \sum_{j=0}^{k-1} \sum_{\ell=0}^2 M_{j\ell} q a^{k-j-1} b^{2-\ell} \tilde{q}$$

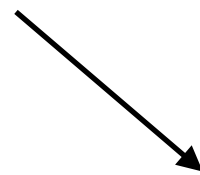
where the singlets $\mathcal{M}_{j,\ell}$ correspond under the duality map to the dressed mesons $QX^jY^\ell Q$ for $j = 0, \dots, k-1$ and $\ell = 0, 1, 2$ of the electric phase. Furthermore k here is required to be odd, otherwise the global anomalies do not match.

	G_{ele}	G_{mag}
SQCD	SU(N)	SU(F-N)
A_k	SU(N)	SU(kF-N)
D_{k+2}	SU(N)	SU(3kF-N)

	G_{ele}	G_{mag}
SQCD	USp(2N)	USp(F-N-4)
A_k	USp(N)	SU(k(F-4)-N)
D_{k+2}	USp(N)	SU(3k F-4k-2-N)

Chiral ring truncation: classical for odd k , quantum (conjectured) in the case of even k

Duality only for even k (mismatch of 't Hooft anomalies for odd k)



$$F_Y = XY + YX = 0$$

$$F_X = Y^2 + X^k = 0$$

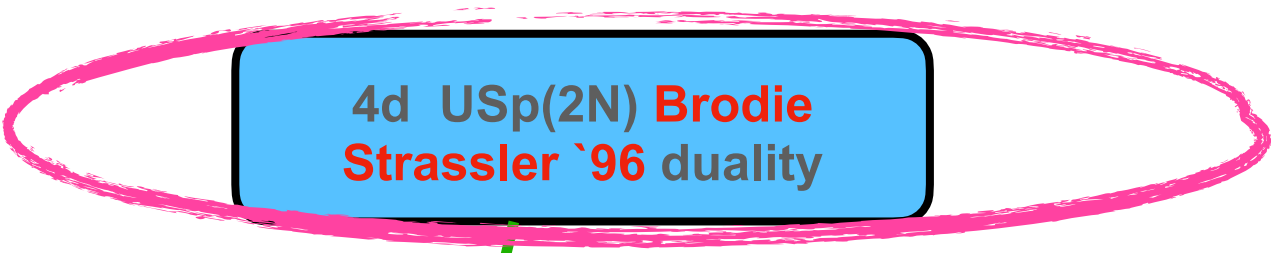
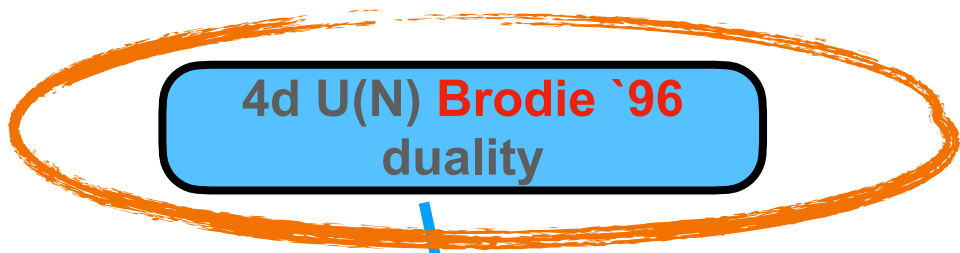
$$YF_X + F_XY = 0 \rightarrow YX^k + X^kY =$$

$$= ((-1)^k + 1)X^kY = -2Y^3$$

Use F_Y

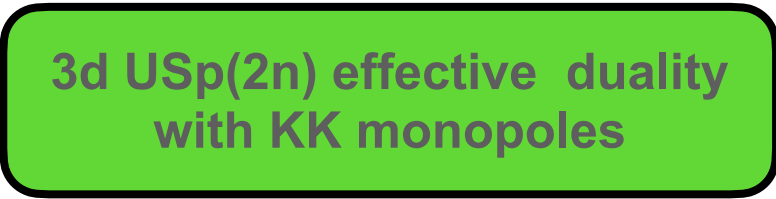
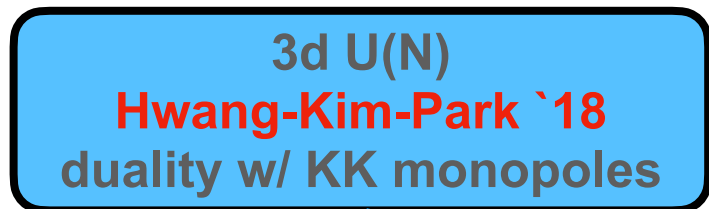


NOTE: absence of a type IIA description of the D-type dualities in the Hanany-Witten setup



Circle reduction

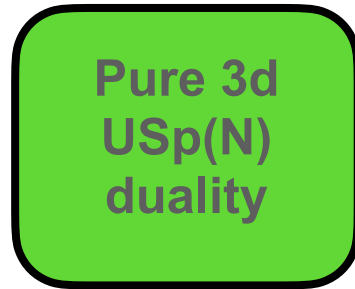
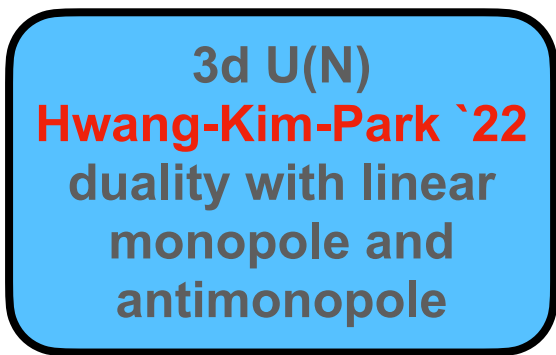
Circle reduction



Real masses and Higgsing

Real masses and dual Higgsing

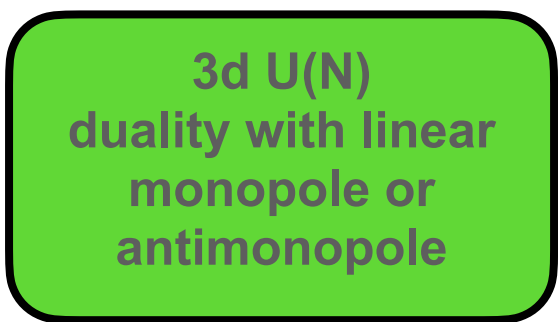
Real masses and dual Higgsing



Real masses and dual Higgsing



Real masses and dual Higgsing



Real masses and dual Higgsing

Hwang-Kim-Park '18:

S^1 reduction of U(n) Brodie duality by breaking the gauge symmetry in SQCD sectors
Aharony like dual W involving also flux-two monopoles $|\pm 1, \pm 1, 0, \dots, 0\rangle$

Hwang-Kim-Park '22:

New U(n) dualities with linear monopole and anti-monopole W in both phases
4d origin not discussed

A.A.-S.Rota '22:

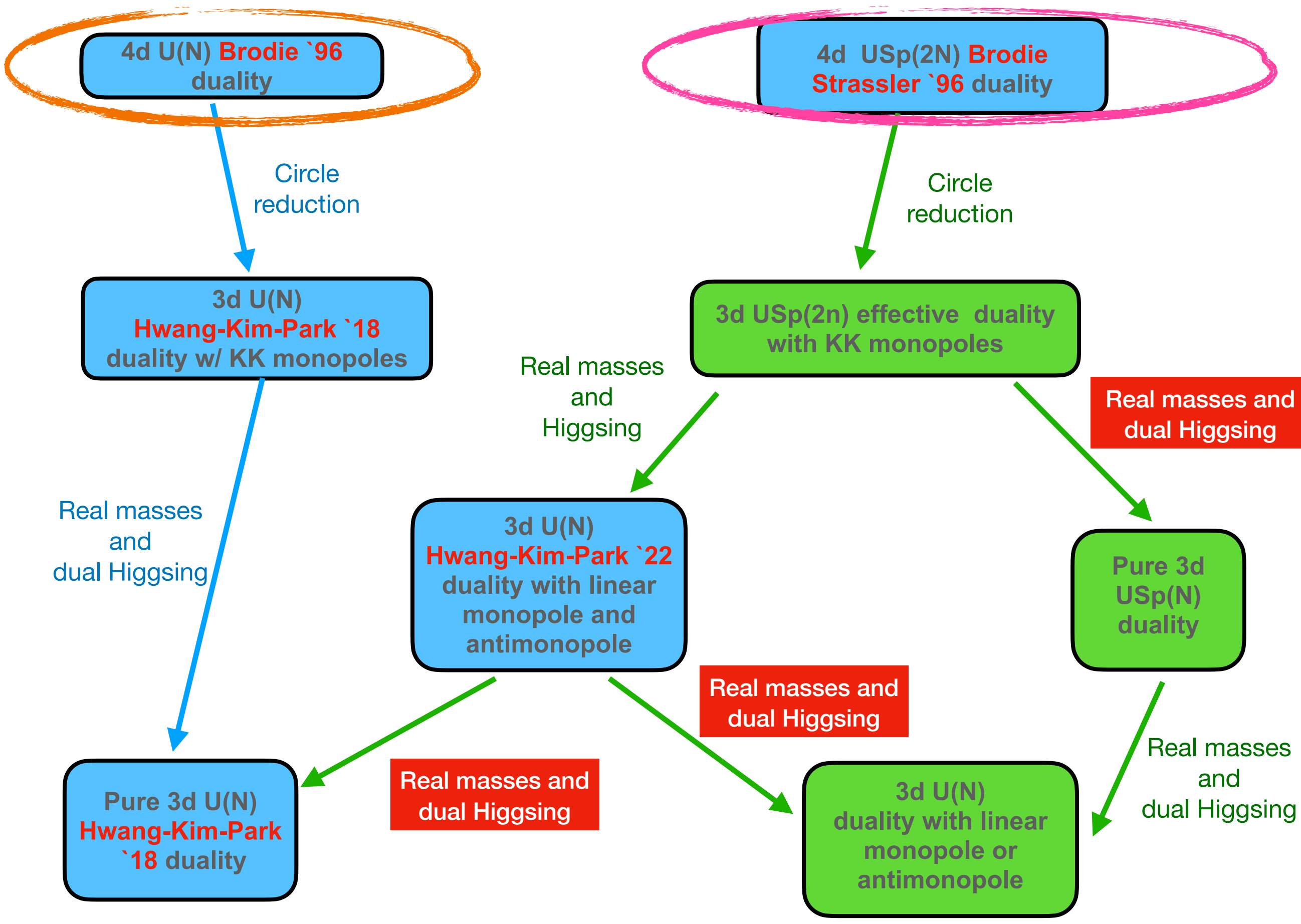
Unified picture, embedding of the various models in the general web

Novelty: non-trivial “topological” sectors, more complicated than in the other cases

Caveat: one **conjecture** necessary (w/o proof, but consistent w/ the whole web)

CONJECTURE: U(k-1) with $W = \text{Tr}XY^2 + \text{Tr}X^{k+1} + V_{\pm}$ is confining

(See Klein '98 for 4d confining limits of Brodie-Strassler '96 dualities)



4d U(N) **Brodie '96** duality

4d USp(2N) **Brodie Strassler '96** duality

Circle reduction

Circle reduction

3d U(N) **Hwang-Kim-Park '18** duality w/ KK monopoles

3d USp(2n) effective duality with KK monopoles

Real masses and Higgsing

Real masses and dual Higgsing

Real masses and dual Higgsing

3d U(N) **Hwang-Kim-Park '22** duality with linear monopole and antimonopole

Pure 3d USp(N) duality

Real masses and dual Higgsing

Pure 3d U(N) **Hwang-Kim-Park '18** duality

Real masses and dual Higgsing

3d U(N) duality with linear monopole or antimonopole

Real masses and dual Higgsing

Results

	G_{electric}	G_{magnetic}	
W_{KK} on S^1 A.A., S. Rota '22	$Usp(2n)$	$USp(2(3kf - n - 4k - 2))$	
	$U(n)$	$U(3kf - n - 4k - 2)$	$W_{mon^+} + W_{mon^-}$ HKP '22
W_{mon^\pm} A.A., S. Rota '22	$U(n)$	$U(3kf - n - 2k - 1)$	
	$U(n)$	$U(3kf - n)$	W_{KK} on S^1 HKP '18
Aharony-like W A.A., S. Rota '22	$Usp(2n)$	$USp(2(3kf - n - 2k - 1))$	
	$U(n)$	$U(3kf - n)$	Aharony-like W HKP '18

A further check: the three sphere partition function

$$Z_{G;k}(\lambda; \vec{\mu}) = \frac{1}{|W|} \int \prod_{i=1}^G \frac{d\sigma_i}{\sqrt{-\omega_1\omega_2}} e^{\frac{k\pi i\sigma_i^2}{\omega_1\omega_2} + \frac{2\pi i\lambda\sigma_i}{\omega_1\omega_2}} \frac{\prod_I \Gamma_h(\omega\Delta_I + \rho_I(\sigma) + \tilde{\rho}_I(\mu))}{\prod_{\alpha \in G_+} \Gamma_h(\pm\alpha(\sigma))},$$

$$\Gamma_h(x; \omega_1, \omega_2) \equiv \Gamma_h(x) \equiv e^{\frac{\pi i}{2\omega_1\omega_2} \left((x-\omega)^2 - \frac{\omega_1^2 + \omega_2^2}{12} \right)} \prod_{j=0}^{\infty} \frac{1 - e^{\frac{2\pi i}{\omega_1}(\omega_2 - x)} e^{\frac{2\pi i\omega_2 j}{\omega_1}}}{1 - e^{-\frac{2\pi i}{\omega_2}x} e^{-\frac{2\pi i\omega_1 j}{\omega_2}}}$$

- Reduction of 4d SCI (or $Z_{S^3 \times S^1}$) to Z_{S^3}
- 4d integral identities (not yet proven!!) become 3d identities
- Interpretation: effective 3d dualities on S^1
- Real mass flows & dual Higgsing
 - Cancellation of divergent contributions in the new (IR) identities
 - “Topological” vacua related to new gauge sectors
 - Eliminated by conjecturing an integral identity for $U(k-1)$ w/ 1 flavor & $W = \text{Tr}XY^2 + \text{Tr}X^{k+1} + V_{\pm}$

In [A.A.-S.Rota '22](#): we checked the consistency of the whole web (i.e. flows and dualities) using this strategy

Conclusions

- Prove the confining duality (using field theory methods and/or localization).
Mirror symmetry?
- Dualities with linear superpotential for the flux-two monopole(s)
- Generalizations
 - Chern-Simons
 - Chiral matter
 - $SU(n)$
 - Other rank-two tensors
 - Orthogonal case
- F-maximization
 - Relevancy of the monopole deformations
 - Conformal window
- Brane picture
- Case of even k