

The background of the slide is a vibrant cosmic scene. It features a large, bright galaxy with a yellow and orange core, surrounded by blue and purple nebulae. The sky is filled with numerous stars and smaller galaxies, creating a rich, multi-colored space environment.

Strong interactions in Cosmology:

*strong CP problem, neutrino mass,
Dark Matter, Dark Energy and
gravitational radiowaves*

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ICHEP 2022 Bologna

Vacuum in Quantum Physics vs in Cosmology

Vacuum energy

in Quantum Physics

*“...the worst theoretical prediction in the history of physics“
(Hobson 2006)*

in Cosmology

$$\epsilon_{vac} \sim 10^{-2} \text{GeV}^4$$

Topological QCD vacuum
unique strongly-coupled subsystem!

$$\Lambda_{\text{cosm}} \sim 10^{-47} \text{GeV}^4$$

$$\sim 10^8 \text{GeV}^4$$

Higgs condensate

“Old” CC problem: Why such small and positive?
“New” CC problem: Why non-zero and exists at all?

Vacuum in Quantum Physics has incredibly wrong energy scale!

Quantum-topological (chromomagnetic) vacuum in QCD

$$\begin{aligned} \epsilon_{vac(top)} &= -\frac{9}{32} \langle 0 | : \frac{\alpha_s}{\pi} F_{ik}^a(x) F_a^{ik}(x) : | 0 \rangle + \frac{1}{4} \left(\langle 0 | : m_u \bar{u}u : | 0 \rangle + \langle 0 | : m_d \bar{d}d : | 0 \rangle + \langle 0 | : m_s \bar{s}s : | 0 \rangle \right) \\ &\simeq -(5 \pm 1) \times 10^9 \text{ MeV}^4. \end{aligned}$$

Two possible approaches to this problem:

- Let's forget about the “bare” vacuum (DE: “phantom”, “quintessence”, “ghost”... etc)
Zero vacuum density in the Minkowski limit, by (Casimir-like) definition, then (Zhitnitsky et al)

$$\Lambda_{\text{cosm}} \equiv \epsilon_{\text{FLRW}} - \epsilon_{\text{Mink}} \quad \text{simply imposing a cancellation of the “bare” vacuum by hands!!}$$

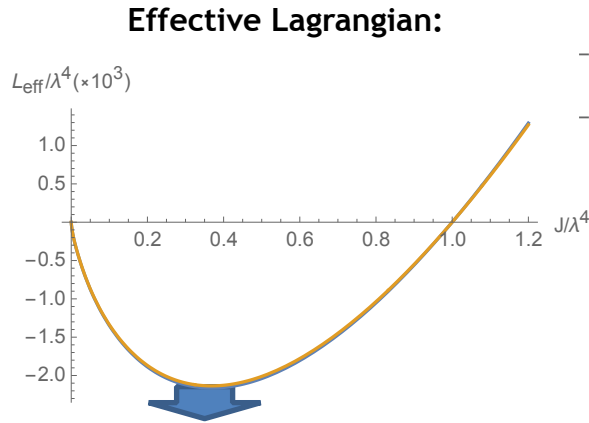
- Let's look closer at the vacuum state — why/how does it become “invisible” to gravity?

Effective YM action and Savvidy vacuum

At least, for SU(2) gauge symmetry,
the all-loop and one-loop effective Lagrangians
are practically indistinguishable (by FRG approach)

P. Dona, A. Marciano, Y. Zhang and C. Antolini, Phys. Rev. D **93** (2016) no.4, 043012.

A. Eichhorn, H. Gies and J. M. Pawłowski, Phys. Rev. D **83** (2011) 045014 [Phys. Rev. D **83** (2011) 069903].



gluon condensate (Savvidy vacuum)

Discovery of chromomagnetic condensate:

G. K. Savvidy, Phys. Lett. **71B**, 133 (1977)

G. Savvidy, Eur. Phys. J. C **80** (2020) 165

NOTE: the RG equation

$$\frac{d \ln |\bar{g}^2|}{d \ln |\mathcal{J}|/\mu_0^4} = \frac{\beta(\bar{g}^2)}{2}$$

H. Pagels and E. Tomboulis, Nucl. Phys. B **143**, 485 (1978).

Classical YM Lagrangian:

$$\mathcal{L}_{\text{cl}} = -\frac{1}{4} F_{\mu\nu}^a F_a^{\mu\nu}$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g_{\text{YM}} f^{abc} A_\mu^b A_\nu^c$$

$$A_\mu^a \equiv g_{\text{YM}} A_\mu^a$$

Effective YM Lagrangian: $\mathcal{F}_{\mu\nu}^a \equiv g_{\text{YM}} F_{\mu\nu}^a$

$$\mathcal{L}_{\text{eff}} = \frac{\mathcal{J}}{4\bar{g}^2(\mathcal{J})}, \quad \mathcal{J} = -\mathcal{F}_{\mu\nu}^a \mathcal{F}_a^{\mu\nu}$$

The energy-momentum tensor:

$$T_\mu^\nu = \frac{1}{\bar{g}^2} \left[\frac{\beta(\bar{g}^2)}{2} - 1 \right] \left(\mathcal{F}_{\mu\lambda}^a \mathcal{F}_a^{\nu\lambda} + \frac{1}{4} \delta_\mu^\nu \mathcal{J} \right) - \delta_\mu^\nu \frac{\beta(\bar{g}^2)}{8\bar{g}^2} \mathcal{J}$$

Equations of motion:

$$\vec{\mathcal{D}}_\nu^{ab} \left[\frac{\mathcal{F}_b^{\mu\nu}}{\bar{g}^2} \left(1 - \frac{\beta(\bar{g}^2)}{2} \right) \right] = 0,$$

$$\vec{\mathcal{D}}_\nu^{ab} \equiv \left(\delta^{ab} \vec{\partial}_\nu - f^{abc} A_\nu^c \right),$$

trace anomaly:

$$T_\mu^\mu = -\frac{\beta(\bar{g}^2)}{2\bar{g}^2} \mathcal{J}$$

appears to be
invariant under

$$\mathcal{J} \longleftrightarrow -\mathcal{J}$$

$$\bar{g}^2 = \bar{g}^2(|\mathcal{J}|)$$

Real-time evolution of the gluon condensate

FLRW metric in conformal time:

$$\mathcal{J} = \frac{2}{\sqrt{-g}} \sum_a (\mathbf{E}_a \cdot \mathbf{E}_a - \mathbf{B}_a \cdot \mathbf{B}_a) \equiv \frac{2}{\sqrt{-g}} (\mathbf{E}^2 - \mathbf{B}^2)$$

$$g \equiv \det(g_{\mu\nu}), \quad g_{\mu\nu} = a(\eta)^2 \text{diag}(1, -1, -1, -1)$$

$$\sqrt{-g} = a^4(\eta), \quad t = \int a(\eta) d\eta$$

- **Basic qualitative features on the non-perturbative YM action are noticed already at one loop**

Einstein-YM equations of motion for the effective YM theory:

$$\frac{1}{\varkappa} \left(R^\nu_\mu - \frac{1}{2} \delta^\nu_\mu R \right) = \bar{\epsilon} \delta^\nu_\mu + \frac{b}{32\pi^2} \frac{1}{\sqrt{-g}} \left[\left(-\mathcal{F}_{\mu\lambda}^a \mathcal{F}_a^{\nu\lambda} + \frac{1}{4} \delta^\nu_\mu \mathcal{F}_{\sigma\lambda}^a \mathcal{F}_a^{\sigma\lambda} \right) \ln \frac{e |\mathcal{F}_{\alpha\beta}^a \mathcal{F}_a^{\alpha\beta}|}{\sqrt{-g} \lambda^4} - \frac{1}{4} \delta^\nu_\mu \mathcal{F}_{\sigma\lambda}^a \mathcal{F}_a^{\sigma\lambda} \right],$$

$$\left(\frac{\delta^{ab}}{\sqrt{-g}} \vec{\partial}_\nu \sqrt{-g} - f^{abc} A_\nu^c \right) \left(\frac{\mathcal{F}_b^{\mu\nu}}{\sqrt{-g}} \ln \frac{e |\mathcal{F}_{\alpha\beta}^a \mathcal{F}_a^{\alpha\beta}|}{\sqrt{-g} \lambda^4} \right) = 0$$

**temporal (Hamilton)
gauge**

$$A_0^a = 0$$

$$e_i^a A_k^a \equiv A_{ik}$$

$$e_i^a e_k^a = \delta_{ik}$$

$$e_i^a e_j^b = \delta_{ab}$$

due to local SU(2) ~ SO(3) isomorphism

$$A_{ik}(t, \vec{x}) = \delta_{ik} U(t) + \tilde{A}_{ik}(t, \vec{x})$$

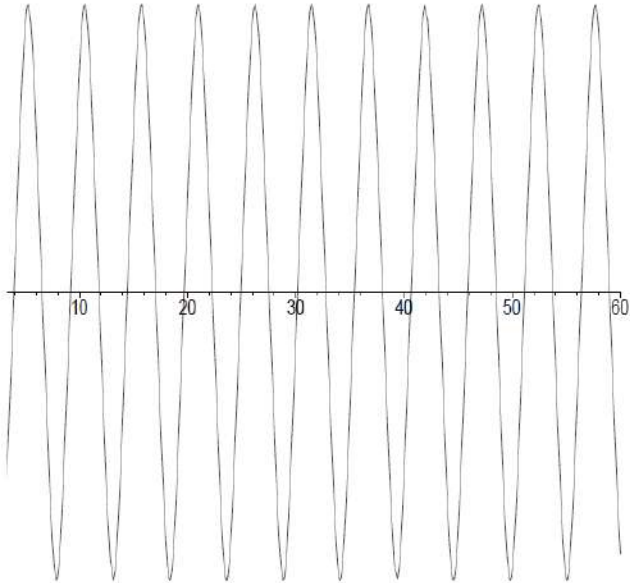
The resulting equations:

$$\frac{6}{\varkappa} \frac{a''}{a^3} = 4\bar{\epsilon} + T_\mu^{\mu,U}, \quad T_\mu^{\mu,U} = \frac{3b}{16\pi^2 a^4} \left[(U')^2 - \frac{1}{4} U^4 \right],$$

$$\frac{\partial}{\partial \eta} \left(U' \ln \frac{6e |(U')^2 - \frac{1}{4} U^4|}{a^4 \lambda^4} \right) + \frac{1}{2} U^3 \ln \frac{6e |(U')^2 - \frac{1}{4} U^4|}{a^4 \lambda^4} = 0$$

Gluon condensate on non-stationary (FLRW) background

Classical YM condensate



“Radiation” medium

$$\epsilon_{\text{YM}} \propto 1/a^4$$

Unstable solution!

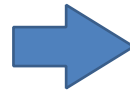
$$Q \equiv \frac{32}{11} \pi^2 e (\xi \Lambda_{\text{QCD}})^{-4} T_{\mu}^{\mu}[U]$$

$$= 6e \left[(U')^2 - \frac{1}{4} U^4 \right] a^{-4} (\xi \Lambda_{\text{QCD}})^{-4}$$

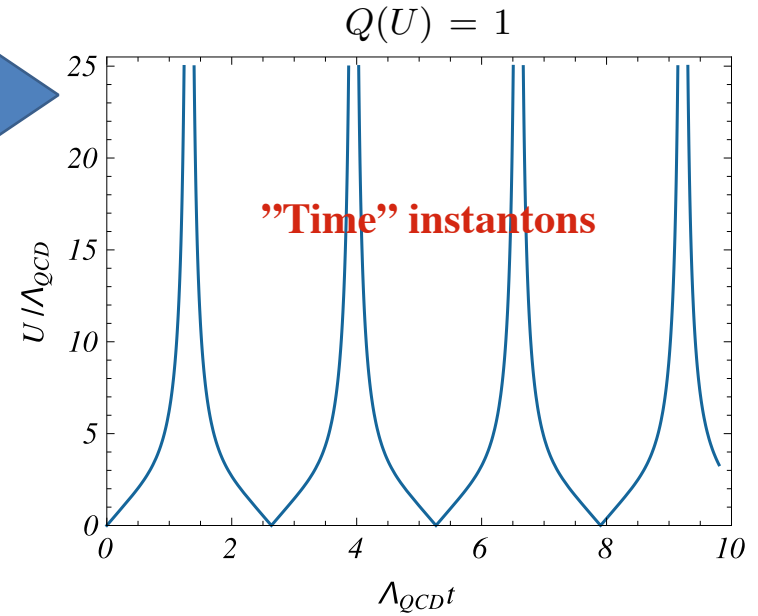
Exact partial solution:

$$|Q| = 1$$

Quantum corrections



Quantum YM vacuum



QCD vacuum:

**a ferromagnetic undergoing
spontaneous magnetisation
(Pagels&Tomboulis)**

Asymptotic (attractor) solution

$$\epsilon_{\text{CE}} \rightarrow +\text{const} \quad t \rightarrow \infty$$

Stable solution!

- In fact, both chromoelectric and chromomagnetic condensates are stable on non-stationary (FLRW) background of expanding Universe**

“Mirror” symmetry of the ground state

In a vicinity of the ground state, the effective Lagrangian

$$\mathcal{L}_{\text{eff}} = \frac{\mathcal{J}}{4\bar{g}^2} \quad \mathcal{J} \simeq \mathcal{J}^*$$

is invariant under

$$\mathbb{Z}_2: \quad \mathcal{J}^* \longleftrightarrow -\mathcal{J}^*, \quad \bar{g}^2(\mathcal{J}^*) \longleftrightarrow -\bar{g}^2(\mathcal{J}^*), \quad \beta(\bar{g}_*^2) \longleftrightarrow -\beta(\bar{g}_*^2)$$

For pure gluodynamics at **one-loop**: $\beta_{(1)} = -\frac{bN}{48\pi^2} \bar{g}_{(1)}^2 \quad b = 11$

$$\alpha_s = \frac{\bar{g}^2}{4\pi} \quad \alpha_s(\mu^2) = \frac{\alpha_s(\mu_0^2)}{1 + \beta_0 \alpha_s(\mu_0^2) \ln(\mu^2/\mu_0^2)} \quad \mu^2 \equiv \sqrt{|\mathcal{J}|}$$

Choosing the ground state value of the condensate $\mu_0^2 \equiv \sqrt{|\mathcal{J}^*|}$ as the physical scale

we observe that **the mirror symmetry**, indeed, holds provided

$$\mathcal{J} \simeq \mathcal{J}^* \quad \alpha_s(\mu_0^2) \longleftrightarrow -\alpha_s(\mu_0^2)$$

i.e. in the ground state only!

Heterogenous quantum ground state: two-scale vacuum

The running coupling at one-loop

$$\bar{g}_1^2(\mathcal{J}) = \frac{\bar{g}_1^2(\mu_0^4)}{1 + \frac{bN}{96\pi^2} \bar{g}_1^2(\mu_0^4) \ln(|\mathcal{J}|/\mu_0^4)} = \frac{96\pi^2}{bN \ln(|\mathcal{J}|/\lambda_{\pm}^4)}$$

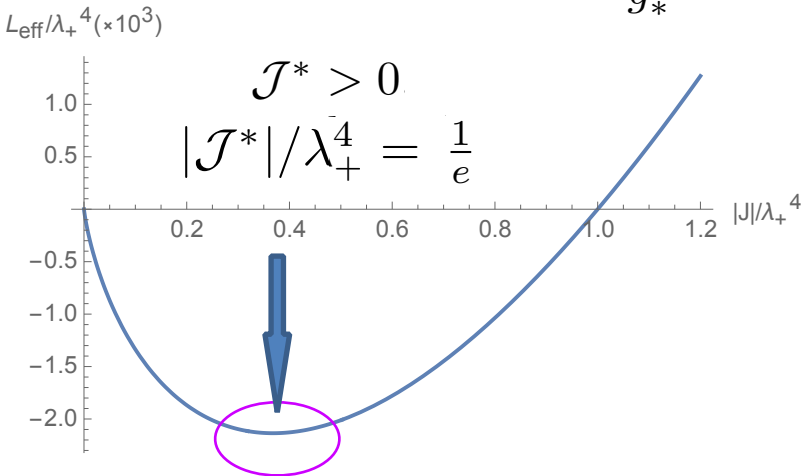
$$\mathcal{L}_{\text{eff}}^{(1)} = \frac{bN}{384\pi^2} \mathcal{J} \ln\left(\frac{|\mathcal{J}|}{\lambda_{\pm}^4}\right) \quad \text{with two energy scales}$$

$$\lambda_{\pm}^4 \equiv |\mathcal{J}^*| \exp\left[\mp \frac{96\pi^2}{bN |\bar{g}_1^2(\mathcal{J}^*)|}\right] \quad |\mathcal{J}^*| = \lambda_+^2 \lambda_-^2$$

CE vacuum: $\beta(\bar{g}_*^2) = 2$

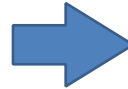
e.o.m. is automatically satisfied!

Trace anomaly: $T_{\mu, \text{CE}}^{\mu} = -\frac{1}{\bar{g}_*^2} \mathcal{J}^*$



Cosmological CE attractor

Mirror symmetry



CM vacuum: $\beta(\bar{g}_*^2) = -2$

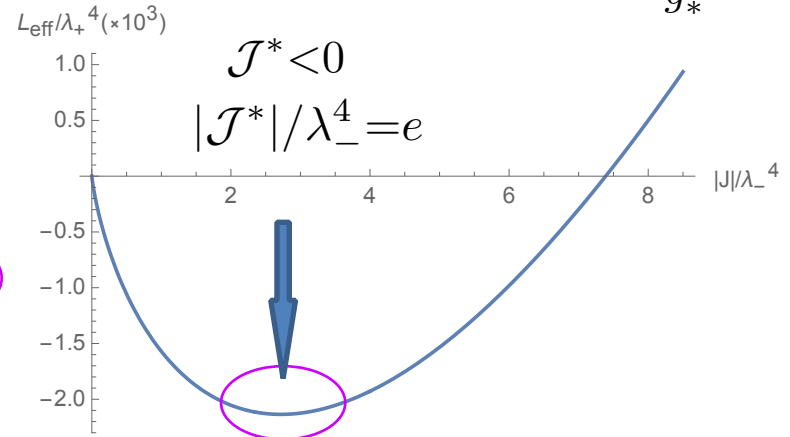
Reduces to the standard YM e.o.m. discussed in e.g. in instanton theory

$$\vec{D}_{\nu}^{ab} \left[\frac{\mathcal{F}_b^{\mu\nu}}{\bar{g}^2} \right] = 0, \quad \bar{g}^2 \simeq \bar{g}_*^2$$

Trace anomaly: $T_{\mu, \text{CM}}^{\mu} = +\frac{1}{\bar{g}_*^2} \mathcal{J}^*$

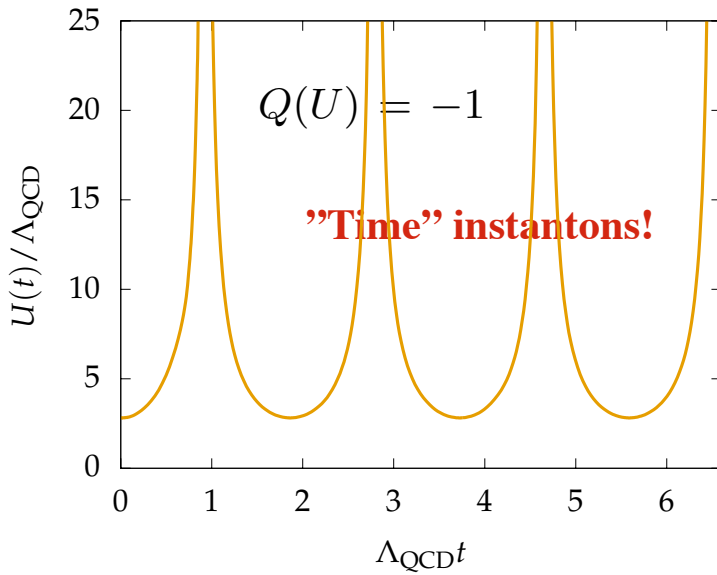
One-loop:

$$\lambda_+^2 / \lambda_-^2 = e$$



Cosmological CM attractor

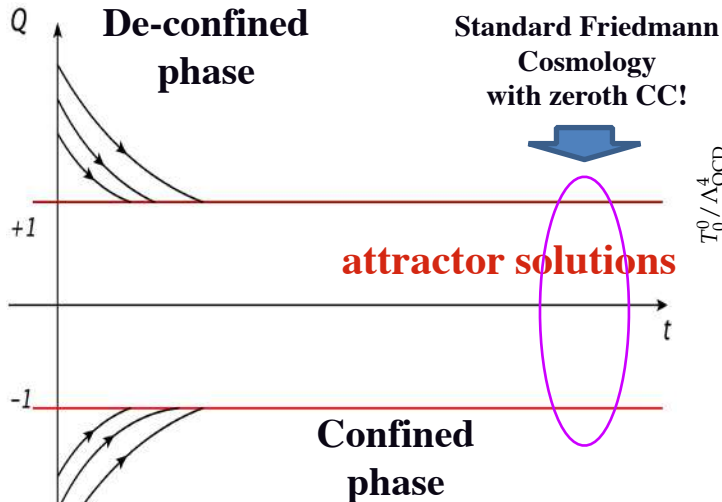
Infrared restoration of conformal invariance



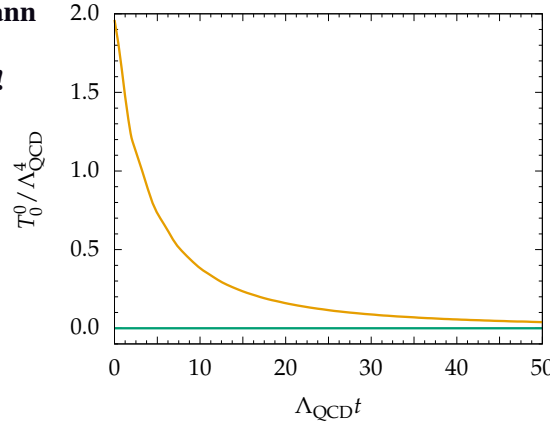
$$\epsilon_{\text{vac}} \equiv \frac{1}{4} \langle T_{\mu}^{\mu} \rangle_{\text{vac}} = \mp \mathcal{L}_{\text{eff}}(\mathcal{J}^*)$$

$$\epsilon_{\text{vac}}^{\text{CE}} |_{\mathcal{J}^* > 0} + \epsilon_{\text{vac}}^{\text{CM}} |_{\mathcal{J}^* < 0} \equiv 0$$

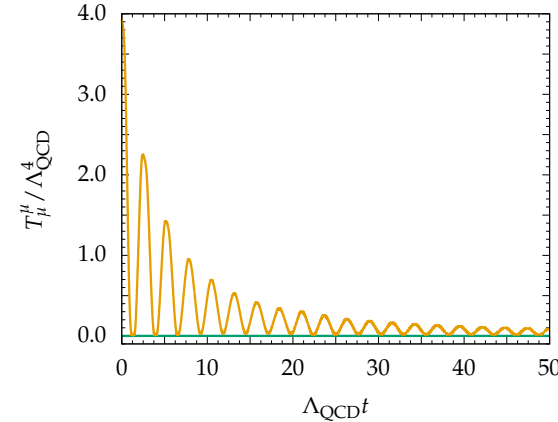
Exact compensation of CM and CE vacua as soon as the cosmological attractor is achieved!



CE energy density



CE EMT trace

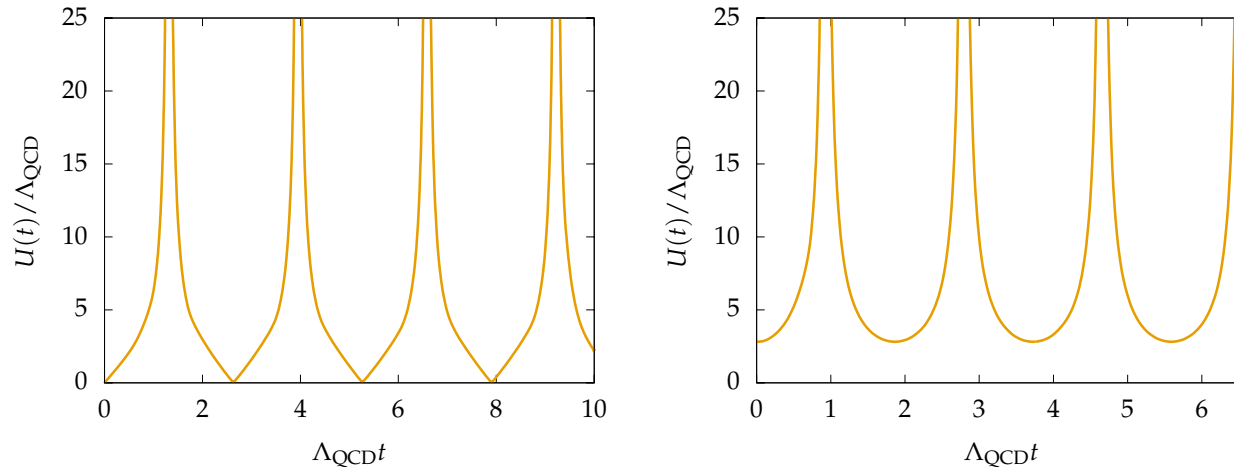


System with very unusual dynamical properties!

Addazi, A.; Marcianò, A.; Pasechnik, R.; Prokhorov, G. Mirror Symmetry of quantum Yang-Mills vacua and cosmological implications. *Eur. Phys. J. C* **2019**, *79*, 251, [[arXiv:hep-th/1804.09826](https://arxiv.org/abs/1804.09826)].

QCD “time crystal”

- The emergence of spikes localised in time at a characteristic QCD time lapse $\Delta t \simeq \Lambda_{\text{QCD}}^{-1}$ and extended in 3-space dimensions reveals the presence of an order state of **space-like soliton/domain wall solutions (chronons)**



- A time-ordered classical solution spontaneously breaking time translational invariance down to a **discrete time shift symmetry** $T_n : t \rightarrow t + n\Lambda_{\text{QCD}}^{-1}$ is known as the **“time crystal”** first discovered by Wilczek in the context of superconductors and superfluids in F. Wilczek, Phys. Rev. Lett. **109**, 160401 (2012)

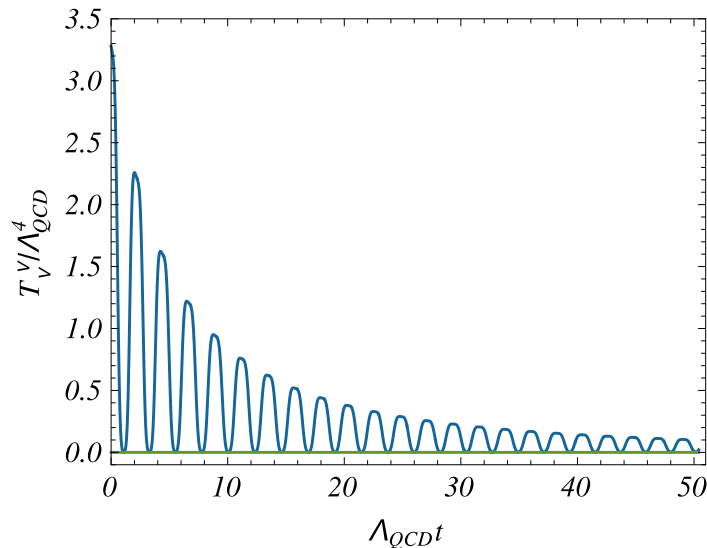
- The kink (anti-kink) profile localised in time corresponds to a space-like domain wall

$$U(\eta) \simeq \frac{v}{\sqrt{2}} \tanh\left[\frac{v}{\sqrt{2}}(\eta - \eta_0)\right] \quad v \simeq \Lambda_{\text{QCD}}$$

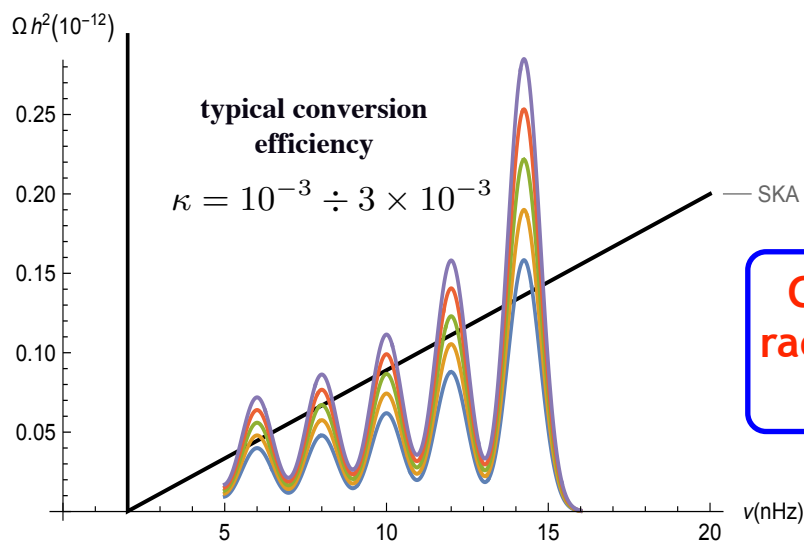
- As the T-invariance is broken, a massless moduli field $\eta_0(x, y, z)$ localised on the domain wall world sheet x, y, z arises and corresponds to a Nambu-Goldstone boson

Gravitational radio-waves from QCD relaxation

The pressure kinks get efficiently transmitted to the primordial plasma inducing shock sound waves and turbulence in it



A. Addazi, A. Marcianò, RP,
CPC 43 (2019) 6, 065101
arXiv: 1812.07376



GW signal lies at the radio-astronomy pulsar timing scale

$10^{-9} \div 10^{-8}$ Hz

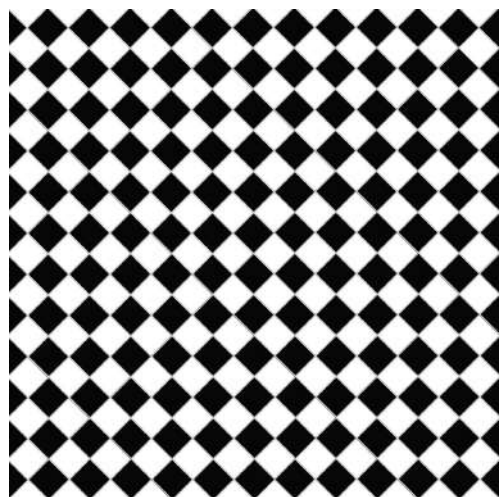
SKA should be able to probe QCD relaxation through detection of primordial GW radio waves

Breaking of Mirror symmetry and Cosmological Constant

Exact mirror symmetry
of the YM ground state



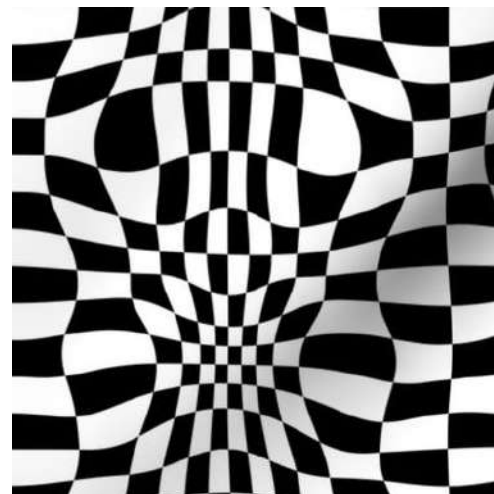
Exact conformal invariance
at macroscopic scales



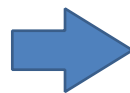
Quantum Gravity in the quasi classical
approximation



Mirror symmetry and conformal invariance
breakdown at cosmological scales



Gravity



Pasechnik, R.; Beylin, V.; Vereshkov, G. Dark Energy from graviton-mediated interactions in the QCD vacuum. *JCAP* 2013, 06, 011, [[arXiv:gr-qc/1302.6456](https://arxiv.org/abs/1302.6456)].

Ya. Zeldovich (1967):

$$\Lambda \sim Gm^6$$

A. Sakharov (1967):

extra terms describing an effect of graviton exchanges between *identical particles* (bosons occupying the same quantum state) should appear in the *right hand side of Einstein equations* (averaged over quantum ensemble)



Graviton

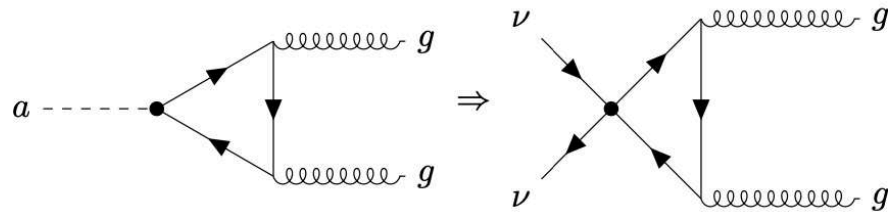
$$\epsilon_{\Lambda} \sim G\Lambda_{\text{QCD}}^6$$

Composite QCD axion as a source of Dark Matter

A. Addazi, A. Marcianò, RP, K. Zeng, Phys.Dark Univ. 36 (2022) 101007; arXiv: 2106.03549

- **Ginzburg, Zharkov 1967**: neutrinos can be non-relativistically produced in a superfluid phase, related to the neutrino mass gap and Dark Energy
- We **postulate** that neutrinos have a Peccei-Quinn (PQ) scale suppressed **portal with the strongly-coupled QCD sector** through quantum anomalies:

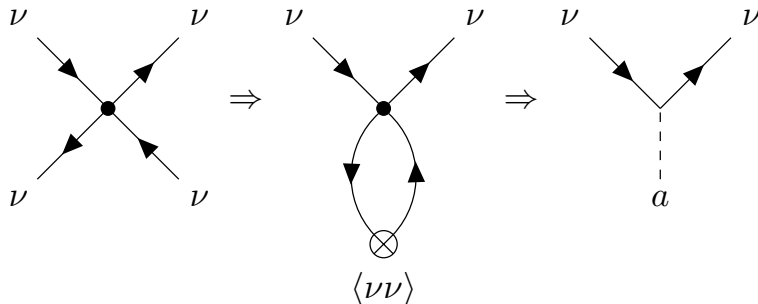
$$\frac{1}{f_{\text{PQ}} f_a^2} (\nu^T \mathcal{C}^{-1} \gamma_5 \nu) G \tilde{G}$$



- The **QCD axion** is then identified with **the composite pair of neutrinos**:

$$\mathcal{L} = \partial\Phi^\dagger \partial\Phi - V(\Phi), \quad V(\Phi) = (|\Phi|^2 - f_{\text{PQ}}^2)^2, \quad \Phi = \rho e^{i \frac{a}{f_{\text{PQ}}}}, \quad a = \frac{1}{f_a^2} \nu^T \mathcal{C}^{-1} \gamma_5 \nu$$

- A four-fermion interaction responsible **for condensation of neutrino pairs and for generation of the neutrino mass scale in the SM**:



Neutrinos can be produced in pairs as cold during the QCD phase transition and condense as non-relativistic Cooper pairs via the misalignment mechanism

- **Neutrino condensate** can naturally account for **the right amount of Dark Matter**

Concluding remarks

- **Local loss of continuous time-translational invariance** leads to “time crystal”-type configurations in the QCD vacuum
- Even starting from an initial non-zero energy-density, the evolution of localised 3-space “pockets” of the CE and CM condensates trigger **a mutual screening**, flowing towards **a zero-energy density attractor and accompanying by a formation of the domain walls** corresponding to an asymptotic restoration of the Z_2 (Mirror) symmetry and effectively protecting the “false” CE vacua pockets from further decay
- The vacua cancellation mechanism seems to **naturally marry the existing confinement pictures** related to a formation of a network of t’Hooft monopoles or chromovortices. In this approach, **the scalar kink profile may correspond the J-invariant** whose change may be related to the presence of monopole or vortex solutions localise inside the space-time domain walls. This implies **the existence of space-time solitonic objects of a new type**.
- **Breaking of the Mirror symmetry by gravitational interactions** induces non-vanishing leading order contribution to the QCD ground state energy **compatible with the observed cosmological constant value**
- **Pressure oscillations** during the QCD relaxation epoch trigger **multi-peaked primordial gravitational wave spectrum in the radio-frequency range** that can be potentially **probed by the SKA telescope**
- Cold neutrino pairs can be **produced during the QCD transition** and condense into axions through a possible **four-fermion neutrino interaction and a coupling to the QCD anomaly enabling neutrino mass gap and Dark Matter generation**