

Mounting evidence for a 95 GeV Higgs boson

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in collaboration with Sven Heinemeyer and Georg Weiglein

[2203.13180,2204.05975]

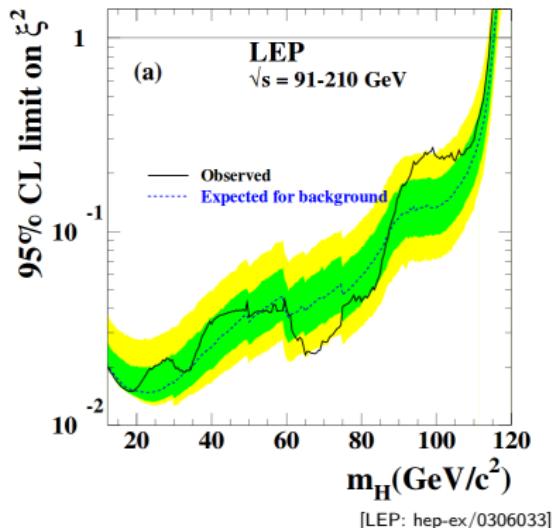
ICHEP 2022 in Bologna

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CLUSTER OF EXCELLENCE
QUANTUM UNIVERSE

“The 95GeV excesses”



~ 2σ local excess at 95 - 98GeV

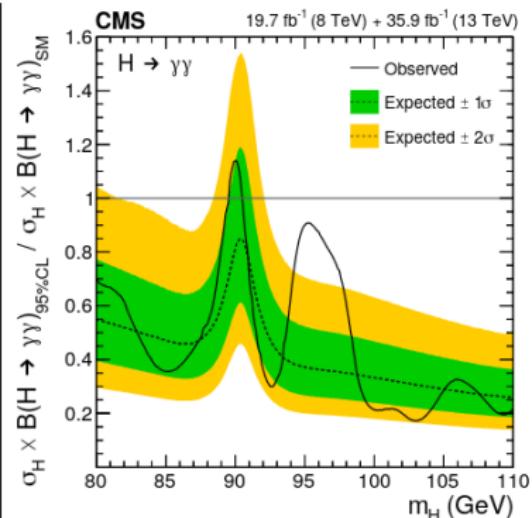
Extracted signal strength:

$$\mu_{bb} (e^+ e^- \rightarrow Zh \rightarrow Zb\bar{b}) = 0.117 \pm 0.057$$

[1612.08522]

Many model interpretations with common origin of both excesses, including N2HDM and NMSSM

see [T.B, M. Chakraborti, S. Heinemeyer: 2003.05422] for a list models



[CMS: 1811.08459]

Run I/II data: Local excess of ~ 3σ

Extracted signal strength:

$$\mu_{\gamma\gamma} (gg \rightarrow h \rightarrow \gamma\gamma) = 0.6 \pm 0.2$$

The Next-to 2 Higgs Doublet Model: N2HDM

$$\begin{aligned}\text{N2HDM} &= \text{SM}(\phi_1) + \text{Second Higgs Doublet}(\phi_2) + \text{Real Scalar Singlet}(\phi_s) \\ &= \text{2HDM}(\phi_1, \phi_2) + \text{Real Scalar Singlet}(\phi_s)\end{aligned}$$

Scalar tree-level potential

$$\begin{aligned}V &= m_{11}^2 |\Phi_1|^2 + m_{22}^2 |\Phi_2|^2 - m_{12}^2 (\Phi_1^\dagger \Phi_2 + \text{h.c.}) + \frac{\lambda_1}{2} (\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^\dagger \Phi_2)^2 \\ &\quad + \lambda_3 (\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1) + \frac{\lambda_5}{2} [(\Phi_1^\dagger \Phi_2)^2 + \text{h.c.}] \\ &\quad + \frac{1}{2} m_S^2 \Phi_S^2 + \frac{\lambda_6}{8} \Phi_S^4 + \frac{\lambda_7}{2} (\Phi_1^\dagger \Phi_1) \Phi_S^2 + \frac{\lambda_8}{2} (\Phi_2^\dagger \Phi_2) \Phi_S^2\end{aligned}$$

Symmetries: Z_2 : $\phi_1 \rightarrow \phi_1$, $\phi_2 \rightarrow -\phi_2$ and $\phi_s \rightarrow \phi_s$, only softly broken by m_{12}^2
 Z'_2 : $\phi_1 \rightarrow \phi_1$, $\phi_2 \rightarrow \phi_2$ and $\phi_s \rightarrow -\phi_s$, spontaneously broken by v_s

Extension of Z_2 to Yukawa sector \Rightarrow 4 types of the (N)2HDM

EW vacuum:

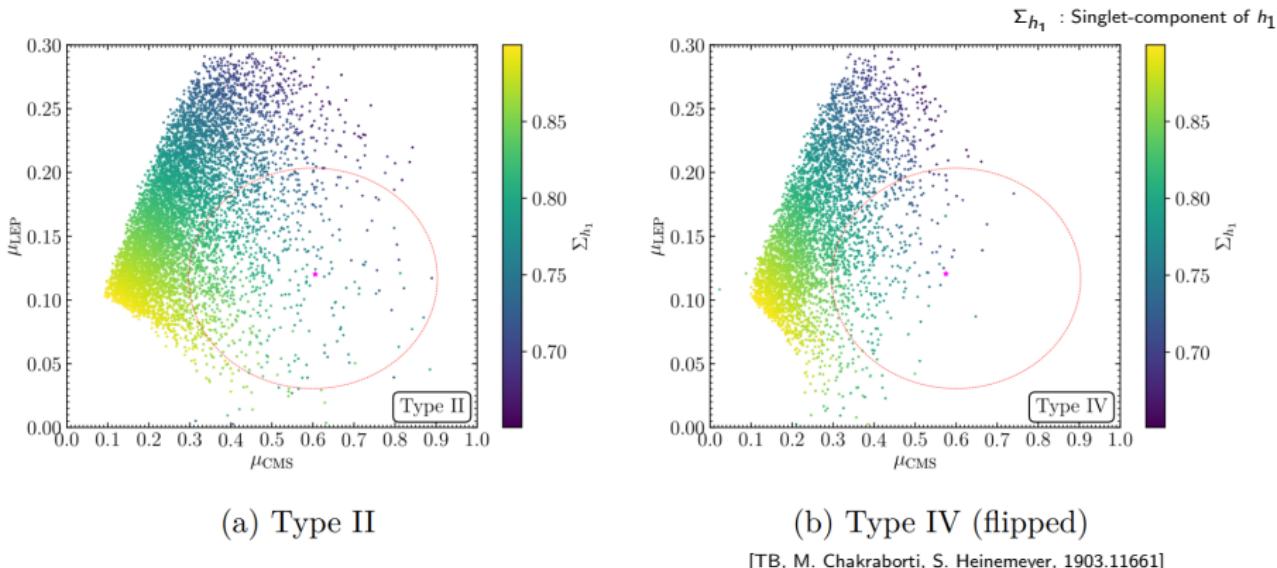
$$\langle \phi_1 \rangle = \begin{pmatrix} 0 \\ v_1/\sqrt{2} \end{pmatrix}, \quad \langle \phi_2 \rangle = \begin{pmatrix} 0 \\ v_2/\sqrt{2} \end{pmatrix}, \quad \langle \phi_S \rangle = v_S/\sqrt{2} \in \mathbb{R}$$

BSM Particles:

CP-even scalars h_1, h_2, h_3 , CP-odd scalar A , charged scalars H^\pm

N2HDM interpretation

$$\begin{aligned} \text{N2HDM} &= \text{SM}(\phi_1) + \text{Second Higgs Doublet}(\phi_2) + \text{Real Scalar Singlet}(\phi_s) \\ &= \text{2HDM}(\phi_1, \phi_2) + \text{Real Scalar Singlet}(\phi_s) \end{aligned}$$

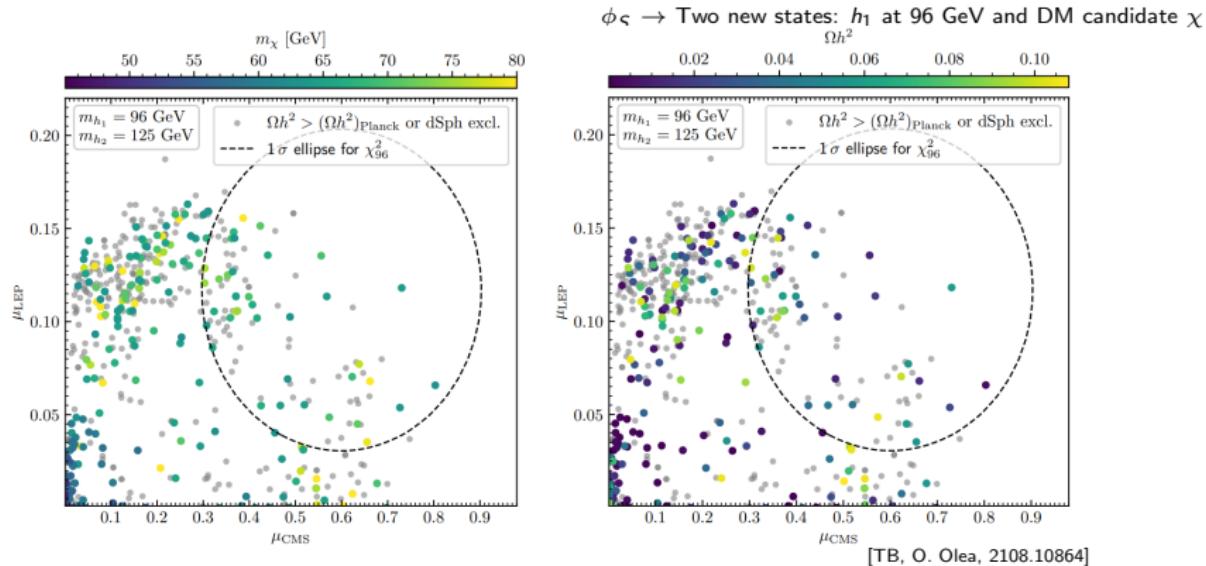


A singlet-like scalar (h_1) at 95GeV that mixes with h_{125} can accommodate the excesses

See also [2109.01128] for interpretation in combination with a 400 GeV pseudoscalar
that accommodates a local 3.5σ excess in $pp \rightarrow t\bar{t}$

S2HDM interpretation

$$\begin{aligned} \text{S2HDM} &= \text{SM}(\phi_1) + \text{Second Higgs Doublet}(\phi_2) + \text{Complex Scalar Singlet}(\phi_s) + \text{U(1)} \\ &= \text{2HDM}(\phi_1, \phi_2) + \text{Complex Scalar Singlet}(\phi_s) + \text{U(1)} \end{aligned}$$

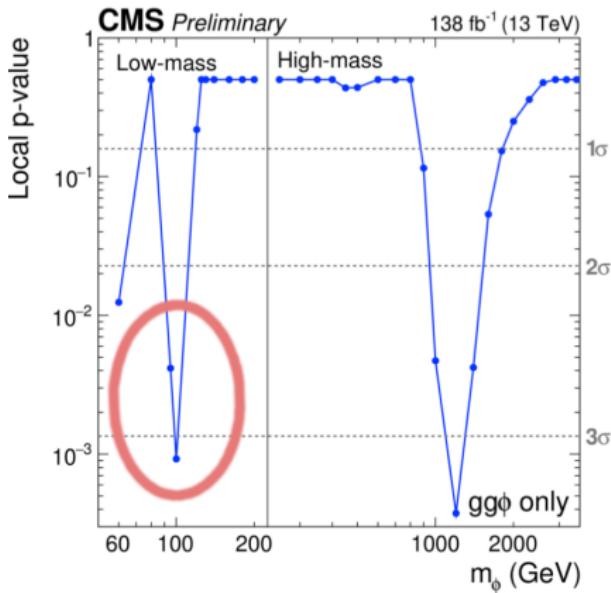
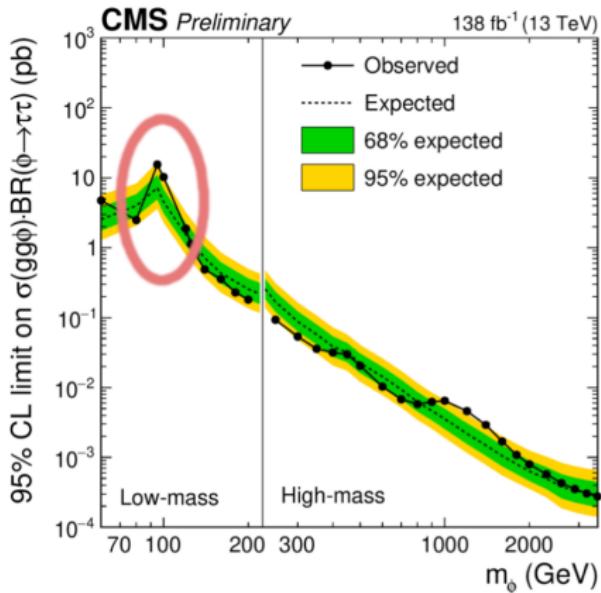


✓ Relic abundance ✓ Galactic-center excess ✓ AMS \bar{p} excess ✓ CMS excess ✓ LEP excess

A softly broken global U(1) makes χ a pseudo-Nambu-Goldstone DM candidate
 → vanishing DM direct-detection scattering cross sections at tree-level

Recent news from CMS

Searches for additional Higgs bosons in $\tau^+\tau^-$ final states



[CMS-HIG-21-001]

Local $3.1(2.6)\sigma$ excess at 100(95) GeV

$$\mu_{\tau\tau}^{\text{exp}} = \frac{\sigma^{\text{exp}}(\text{gg} \rightarrow \phi \rightarrow \tau^+\tau^-)}{\sigma^{\text{SM}}(\text{gg} \rightarrow H \rightarrow \tau^+\tau^-)} = 1.2 \pm 0.5$$

N2HDM: Can we accommodate also the new excess in $\tau^+\tau^-$?

1. $\tau^+\tau^-$ excess in combination with $\gamma\gamma$ excess
2. $\tau^+\tau^-$ excess in combination with $\gamma\gamma$ and $b\bar{b}$ excesses
3. And what about M_W ?

Strategie: Random scan to minimize χ^2 in order to find valid points that accommodate the excesses

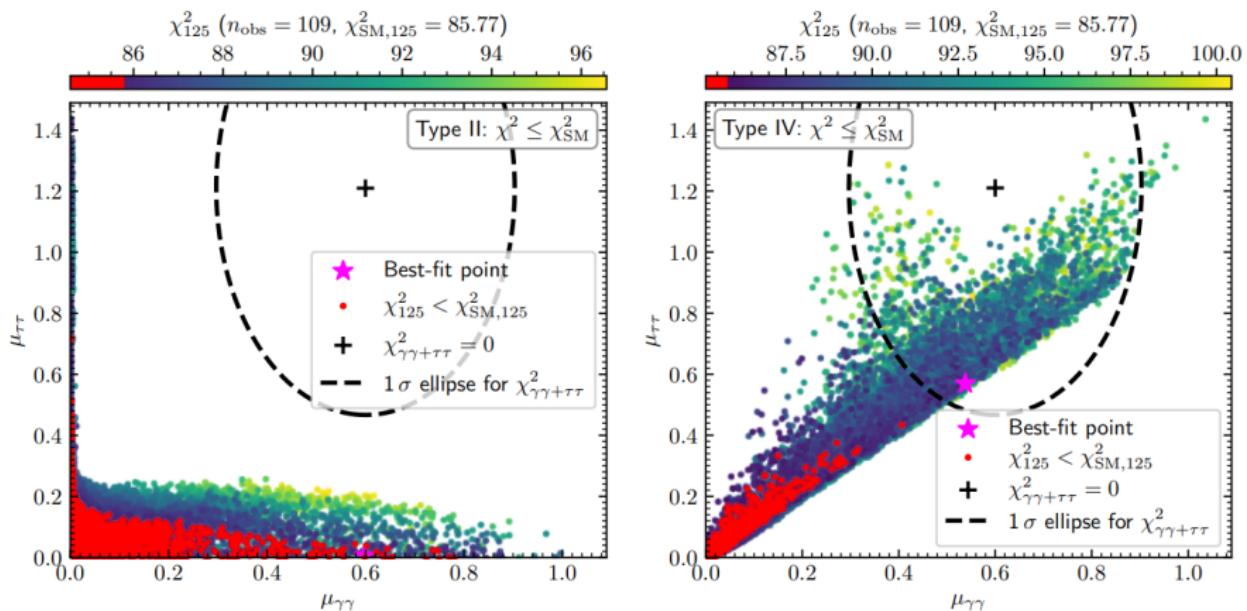
$$\chi^2 = \chi_{\gamma\gamma}^2 + \chi_{\tau\tau}^2 (+\chi_{bb}^2) + \chi_{125}^2 , \quad \chi_{\gamma\gamma, \tau\tau, bb}^2 = \frac{(\mu_{\gamma\gamma, \tau\tau, bb} - \mu_{\gamma\gamma, \tau\tau, bb}^{\text{exp}})^2}{(\Delta \mu_{\gamma\gamma, \tau\tau, bb}^{\text{exp}})^2}$$

Constraints: Collider searches, measurements of $h_{125} \rightarrow \chi_{125}^2$, flavour physics, vacuum stability, tree-level perturbative unitarity, electroweak precision observables

For M_W -Tevatron:

$$\chi_{M_W^{\text{CDF-new}}}^2 = \frac{(M_W^{\text{N2HDM}} - M_W^{\text{CDF-new}})^2}{(\Delta M_W^{\text{CDF-new}})^2} \leq 4 , \quad M_W^{\text{N2HDM}} = M_W^{\text{N2HDM}}(S, T, U)$$

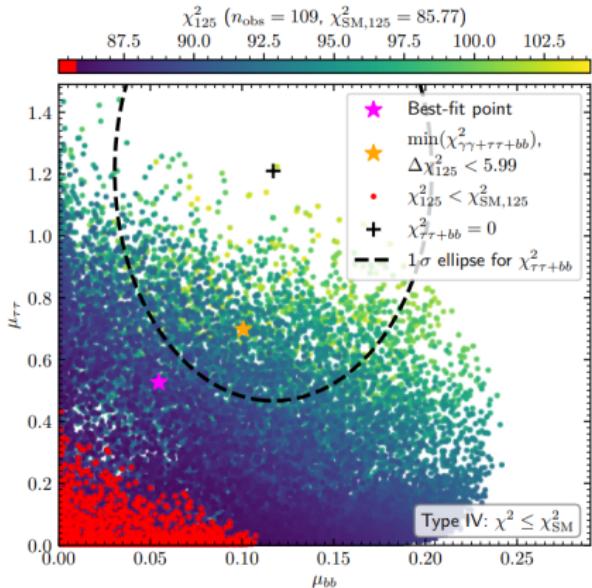
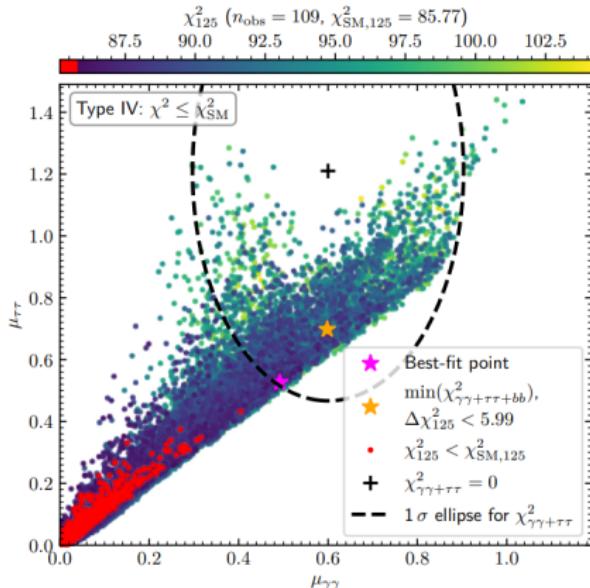
N2HDM interpretation: $\tau^+\tau^-$ and $\gamma\gamma$



[TB, S. Heinemeyer, G. Weiglein, 2203.13180]

Type II does not work, but type IV does!

N2HDM interpretation: $\tau^+\tau^-$ and $\gamma\gamma$ and $b\bar{b}$

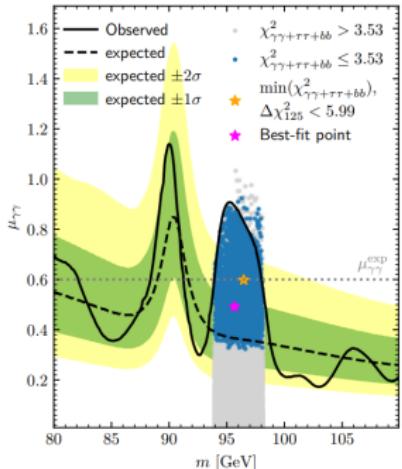


[TB, S. Heinemeyer, G. Weiglein, 2203.13180]

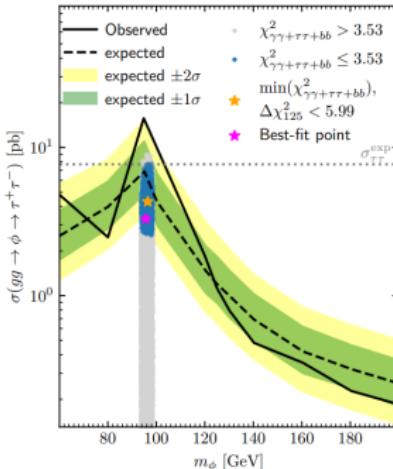
N2HDM Type IV can accommodate the excesses in all three decay modes

N2HDM interpretation: $\tau^+\tau^-$ and $\gamma\gamma$ and $b\bar{b}$

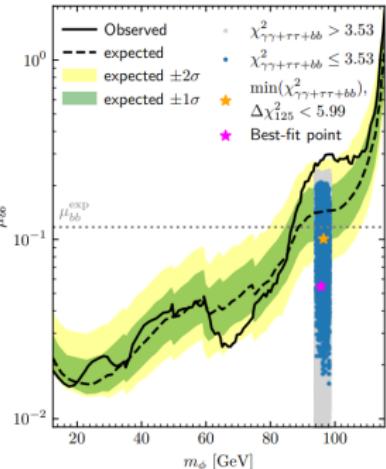
$gg \rightarrow h_{95} \rightarrow \gamma\gamma$



$gg \rightarrow h_{95} \rightarrow \tau^+\tau^-$



$e^+e^- \rightarrow Z(h_{95}) \rightarrow (b\bar{b})$

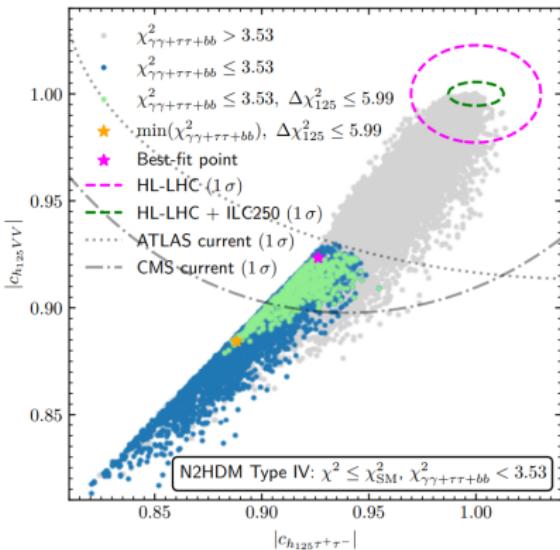
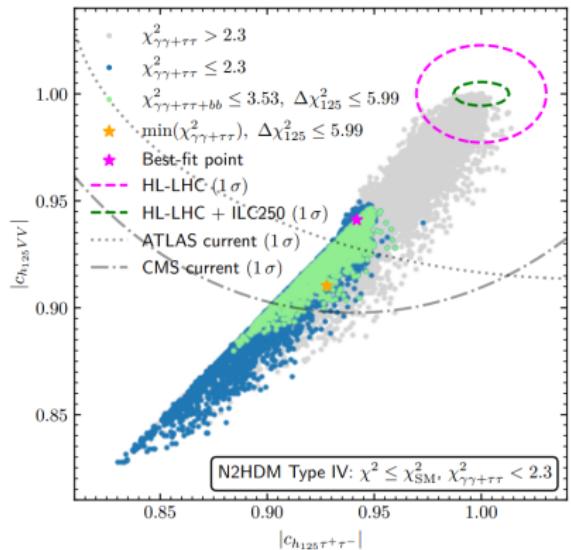


[TB, S. Heinemeyer, G. Weiglein, 2203.13180]

Blue points = Describe the excesses within 1σ confidence level: $\chi^2_{\gamma\gamma + \tau\tau + b\bar{b}} \leq 3.53$

N2HDM interpretation: $\tau^+\tau^-$ and $\gamma\gamma$ and $b\bar{b}$

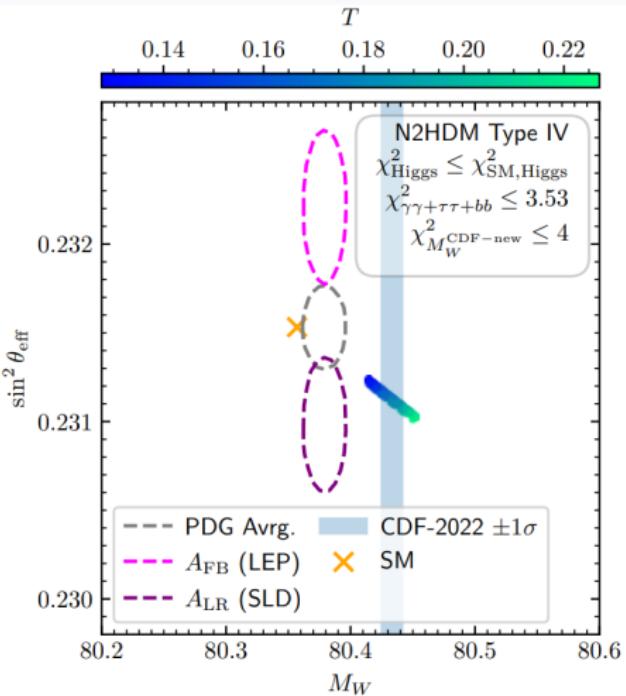
Can be probed indirectly via h_{125} :



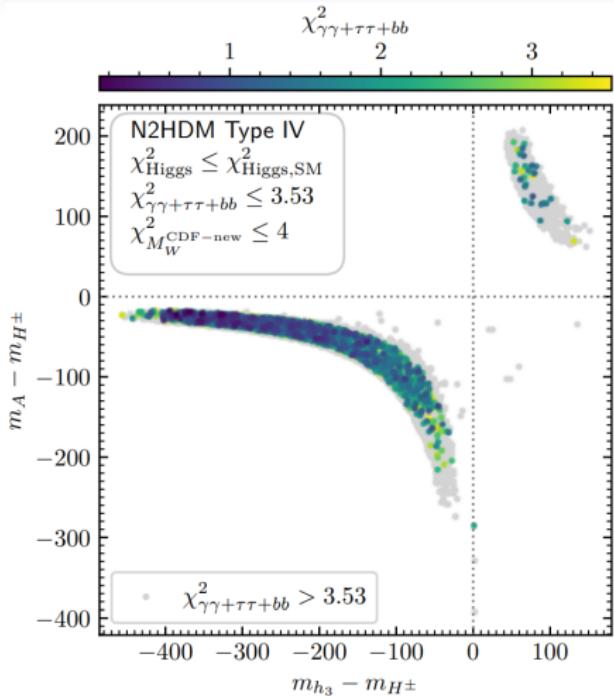
[TB, S. Heinemeyer, G. Weiglein, 2203.13180]

HL-LHC precision sufficient to exclude the scenario

Tevatron and M_W in the N2HDM



Preference for value of effective weak mixing angle as extracted from A_{LR} measured at SLD



Preferred region with mass hierarchy for heavy Higgs bosons: $m_{h_3} < m_A \approx m_{H^\pm}$

[TB, S. Heinemeyer, G. Weiglein, 2204.05975]

Conclusions

Three local excesses ($3\sigma + 3\sigma + 2\sigma$) at the same mass

→ two different production modes (ggF and Higgsstrahlung)

→ three different decay modes ($\gamma\gamma$, $\tau\tau$, bb)

Can be described economically in the N2HDM

→ valid DM candidate if N2HDM → S2HDM (galactic-center excess)

→ Can simultaneously accommodate M_W Tevatron measurement

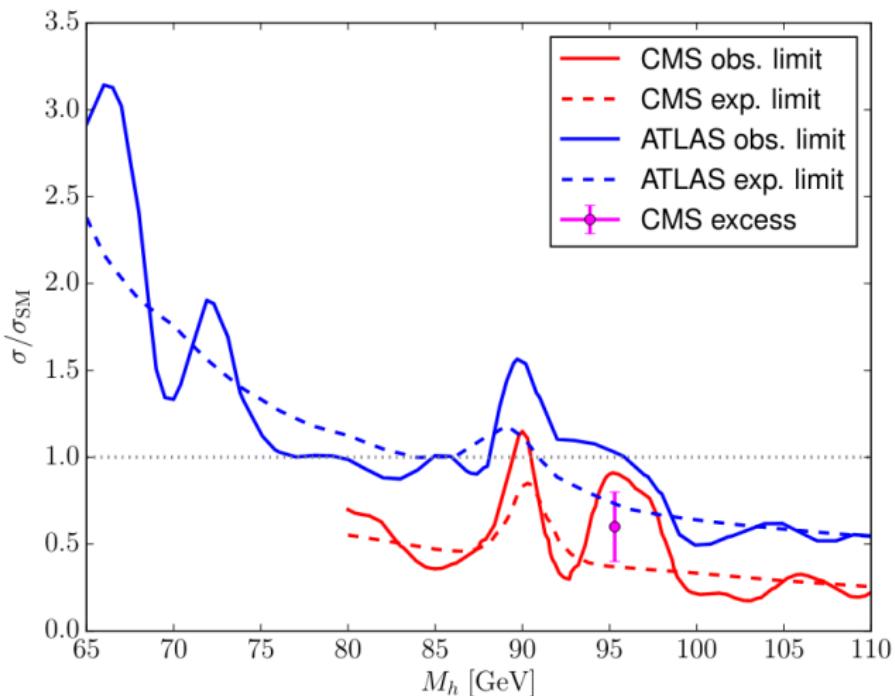
HL-LHC h_{125} measurements have potential to exclude all of this

CMS/ATLAS $\gamma\gamma$ low-mass searches with Full Run II eagerly awaited

Thanks!

ATLAS searches below 125 GeV (80fb^{-1})

[ATLAS-CONF-2018-025]



[S. Heinemeyer, T. Stefaniak, 1812.05864]

The broken phase

Assuming CP conservation: 3 CP-even, 1 CP-odd, 1 charged Higgs bosons

Electroweak symmetry breaking

$$\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \frac{1}{\sqrt{2}}(v_1 + \rho_1 + i\eta_1) \end{pmatrix}, \quad \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}}(v_2 + \rho_2 + i\eta_2) \end{pmatrix}, \quad \Phi_S = v_S + \rho_S$$

CP-even scalar sector: Mixing from gauge eigenstate to mass eigenstate basis

$$\begin{pmatrix} h_1 \\ h_2 \\ h_3 \end{pmatrix} = R \begin{pmatrix} \rho_1 \\ \rho_2 \\ \rho_S \end{pmatrix}, \quad R = \begin{pmatrix} c_{\alpha_1} c_{\alpha_2} & s_{\alpha_1} c_{\alpha_2} & s_{\alpha_2} \\ -(c_{\alpha_1} s_{\alpha_2} s_{\alpha_3} + s_{\alpha_1} c_{\alpha_3}) & c_{\alpha_1} c_{\alpha_3} - s_{\alpha_1} s_{\alpha_2} s_{\alpha_3} & c_{\alpha_2} s_{\alpha_3} \\ -c_{\alpha_1} s_{\alpha_2} c_{\alpha_3} + s_{\alpha_1} s_{\alpha_3} & -(c_{\alpha_1} s_{\alpha_3} + s_{\alpha_1} s_{\alpha_2} c_{\alpha_3}) & c_{\alpha_2} c_{\alpha_3} \end{pmatrix}$$

Convention: $m_{h_1} < m_{h_2} < m_{h_3}$, $-\frac{\pi}{2} \leq \alpha_i \leq \frac{\pi}{2}$



Coupling to SM fermions:

$$-\mathcal{L}_{\text{Yuk}} = \sum_{i=1}^2 \frac{\sqrt{2}m_f}{v} \cancel{c_{h_i f \bar{f}}} \bar{\psi}_f \psi_f h_i$$

$$t_\beta = \frac{v_2}{v_1}$$

	$c_{h_i b \bar{b}}$	$c_{h_i t \bar{t}}$	$c_{h_i \tau \bar{\tau}}$
type I	R_{i2}/s_β	R_{i2}/s_β	R_{i2}/s_β
type II	R_{i1}/c_β	R_{i2}/s_β	R_{i1}/c_β
lepton-specific	R_{i2}/s_β	R_{i2}/s_β	R_{i1}/c_β
flipped	R_{i1}/c_β	R_{i2}/s_β	R_{i2}/s_β

S2HDM: Singlet-extended 2 Higgs doublet model

ϕ_1, ϕ_2 : SU(2) doublets, ϕ_S : SM singlet, charged under global U(1)

Scalar potential:

$$\begin{aligned} V = & \mu_{11}^2 (\phi_1^\dagger \phi_1) + \mu_{22}^2 (\phi_2^\dagger \phi_2) - \mu_{12}^2 ((\phi_1^\dagger \phi_2) + (\phi_2^\dagger \phi_1)) + \frac{1}{2} \mu_S^2 |\phi_S|^2 - \frac{1}{4} \mu_\chi^2 (\phi_S^2 + (\phi_S^*)^2) \\ & + \frac{1}{2} \lambda_1 (\phi_1^\dagger \phi_1)^2 + \frac{1}{2} \lambda_2 (\phi_2^\dagger \phi_2)^2 + \lambda_3 (\phi_1^\dagger \phi_1) (\phi_2^\dagger \phi_2) + \lambda_4 (\phi_1^\dagger \phi_2) (\phi_2^\dagger \phi_1) \\ & + \frac{1}{2} \lambda_5 ((\phi_1^\dagger \phi_2)^2 + (\phi_2^\dagger \phi_1)^2) + \frac{1}{2} \lambda_6 (|\phi_S|^2)^2 + \lambda_7 (\phi_1^\dagger \phi_1) |\phi_S|^2 + \lambda_8 (\phi_2^\dagger \phi_2) |\phi_S|^2 \end{aligned}$$

EW vacuum: $\langle \phi_1 \rangle = \begin{pmatrix} 0 \\ v_1/\sqrt{2} \end{pmatrix}, \quad \langle \phi_2 \rangle = \begin{pmatrix} 0 \\ v_2/\sqrt{2} \end{pmatrix}, \quad \langle \phi_S \rangle = v_S/\sqrt{2} \in \mathbb{R}$

BSM particles:

$h_{1,2,3}$: CP-even Higgs bosons

H^\pm : Charged Higgs bosons

A : CP-odd Higgs boson

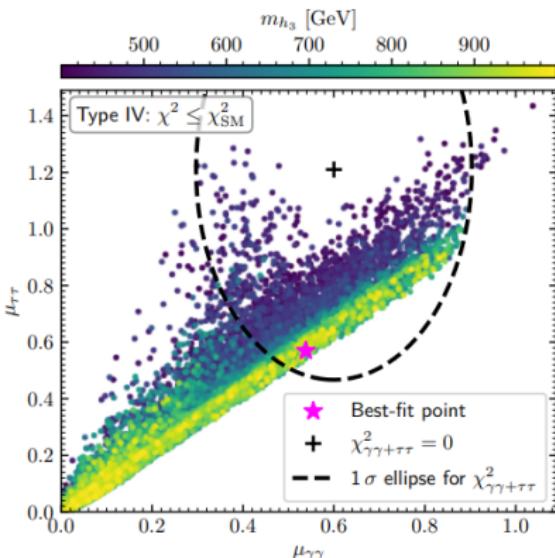
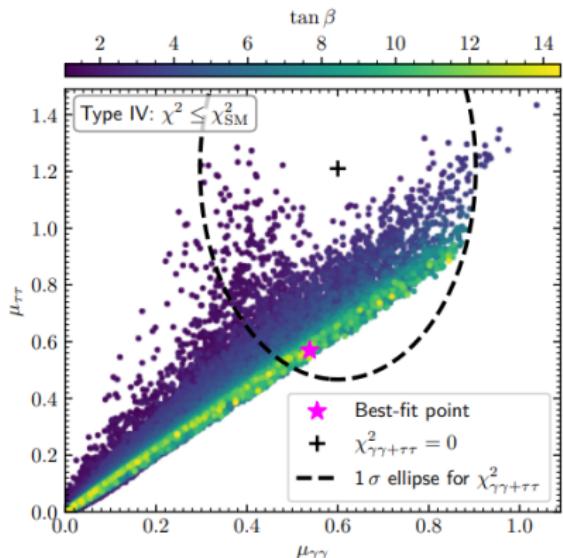
χ : pNG DM

$$\begin{pmatrix} h_1 \\ h_2 \\ h_3 \end{pmatrix} = R(\alpha_1, \alpha_2, \alpha_3) \cdot \begin{pmatrix} \text{Re}(\phi_1^0) \\ \text{Re}(\phi_2^0) \\ \text{Re}(\phi_S^0) \end{pmatrix}$$

Free parameters (Yukawa type II):

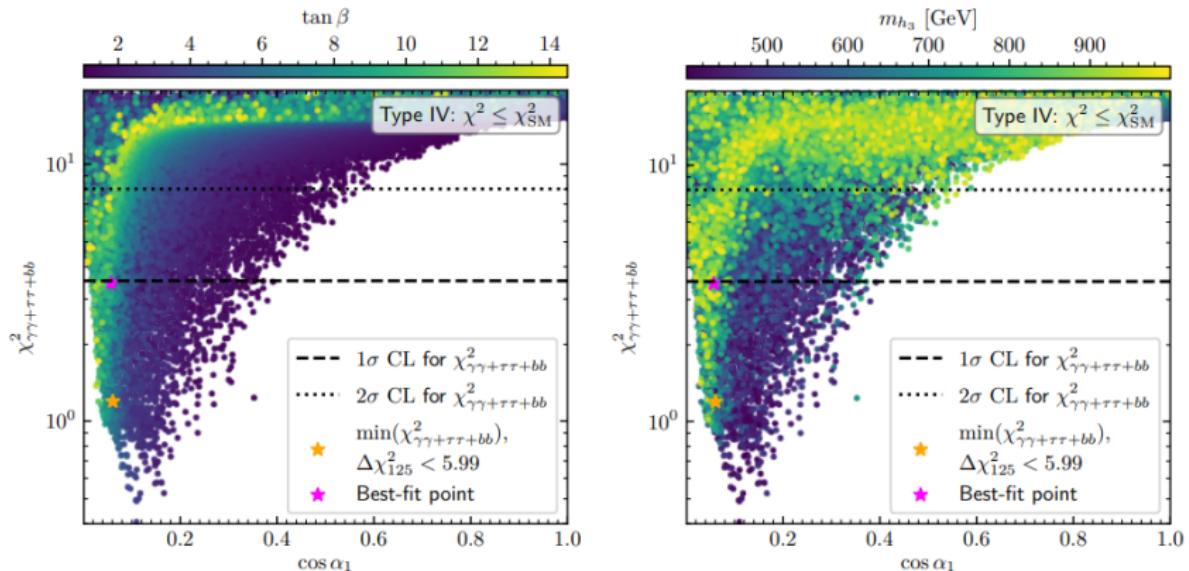
$$m_{h_{1,2,3}}, \quad m_A, \quad m_{H^\pm}, \quad m_\chi, \quad \alpha_{1,2,3}, \quad \tan \beta, \quad M = \sqrt{\mu_{12}^2 / (s_\beta c_\beta)}, \quad v_S$$

Favoured parameter space



[TB, S. Heinemeyer, G. Weiglein, 2203.13180]

Favoured parameter space



[TB, S. Heinemeyer, G. Weiglein, 2203.13180]

Correlations of $\mu_{\gamma\gamma}$, $\mu_{\tau\tau}$, μ_{bb} and $c_{h_{125}VV}^2$

