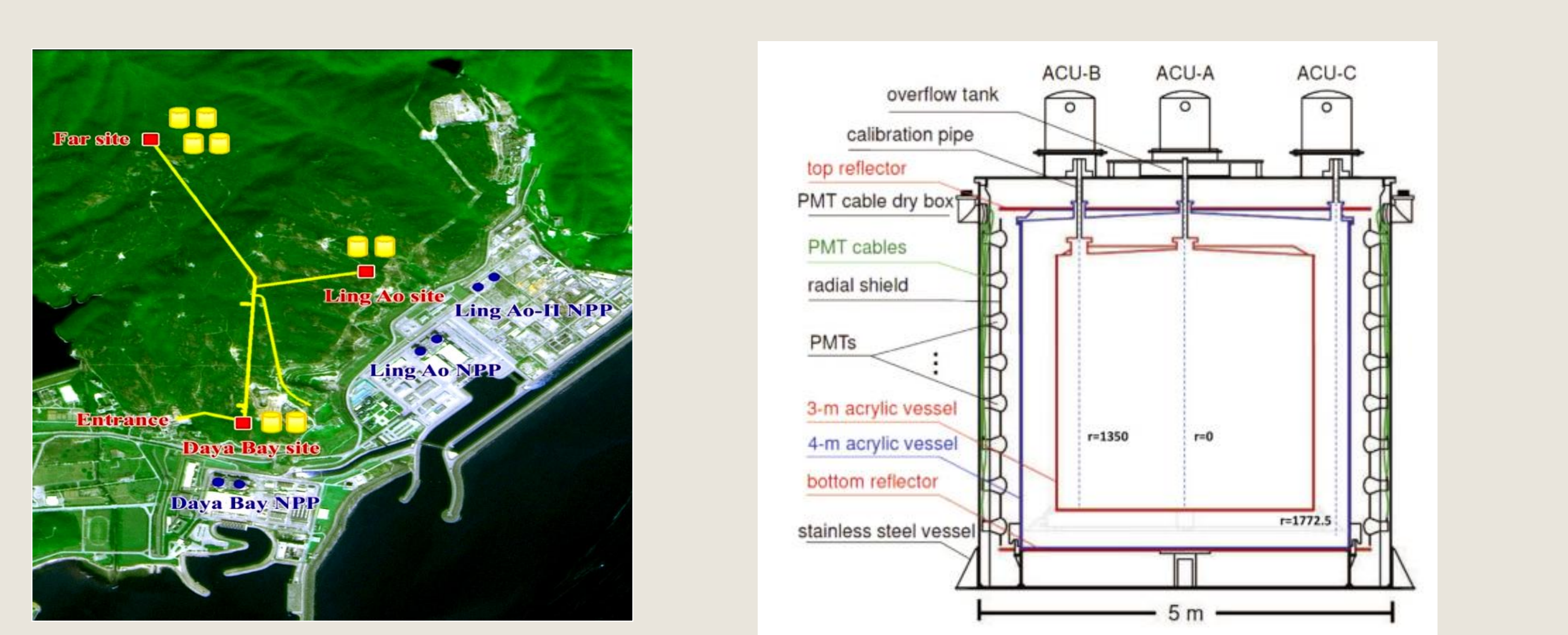


1. Daya Bay Experiment [1]

The Daya Bay Experiment provides precise measurement of reactor anti-neutrino disappearance via Inverse Beta Decays (IBDs), and the IBDs are tagged by neutron capture on gadolinium (nGd) or on hydrogen (nH).



- 6 low enriched uranium (LEU) commercial reactors, each with 2.9 GW thermal power.
- These reactors provide abundant electron-antineutrinos.

2. Gravitational-Wave Events

- The direct observation of gravitational waves (GWs) provides an important probe for investigating the dynamical origin of high-energy cosmic transients.
- Providing a possible connection between neutrino emission and gravitational-wave (GW) bursts is important to our understanding of the physical processes that occur when black holes or neutron stars merge.
- A list of GWs of interest for this study is listed [2-8].

GW events	Type of merged bodies	Detection time (UTC)	Distance D_{LIGO} (Mpc)
GW150914	Black holes	2015.09.14 09:50:45	410^{+160}_{-180}
GW151012	Black holes	2015.10.12 09:54:43	1100^{+500}_{-500}
GW151226	Black holes	2015.12.26 03:38:53	440^{+180}_{-190}
GW170104	Black holes	2017.01.04 10:11:58	880^{+450}_{-390}
GW170608	Black holes	2017.06.08 02:01:16	340^{+140}_{-140}
GW170814	Black holes	2017.08.14 10:30:43	540^{+130}_{-210}
GW170817	Neutron stars	2017.08.17 12:41:04	40^{+8}_{-14}

3. Neutrino Fluence Measurement Method

- We limited our search for $\bar{\nu}_e$ with energies below 100 MeV.
- To accommodate the uncertainties, we adopted three time windows to search for neutrino bursts associated with the GW events, ± 10 s, ± 500 s, ± 1000 s.
- We first measured the electron-antineutrino fluence, Φ_{FD} , with a normalized, pinched Fermi-Dirac spectrum, with zero chemical potential and a pinching factor of $\eta = 0$, as applied in the KamLAND experiment [9]. Using the number of electron-antineutrino candidates N_ν within the searching window, the electron-antineutrino fluence is calculated as

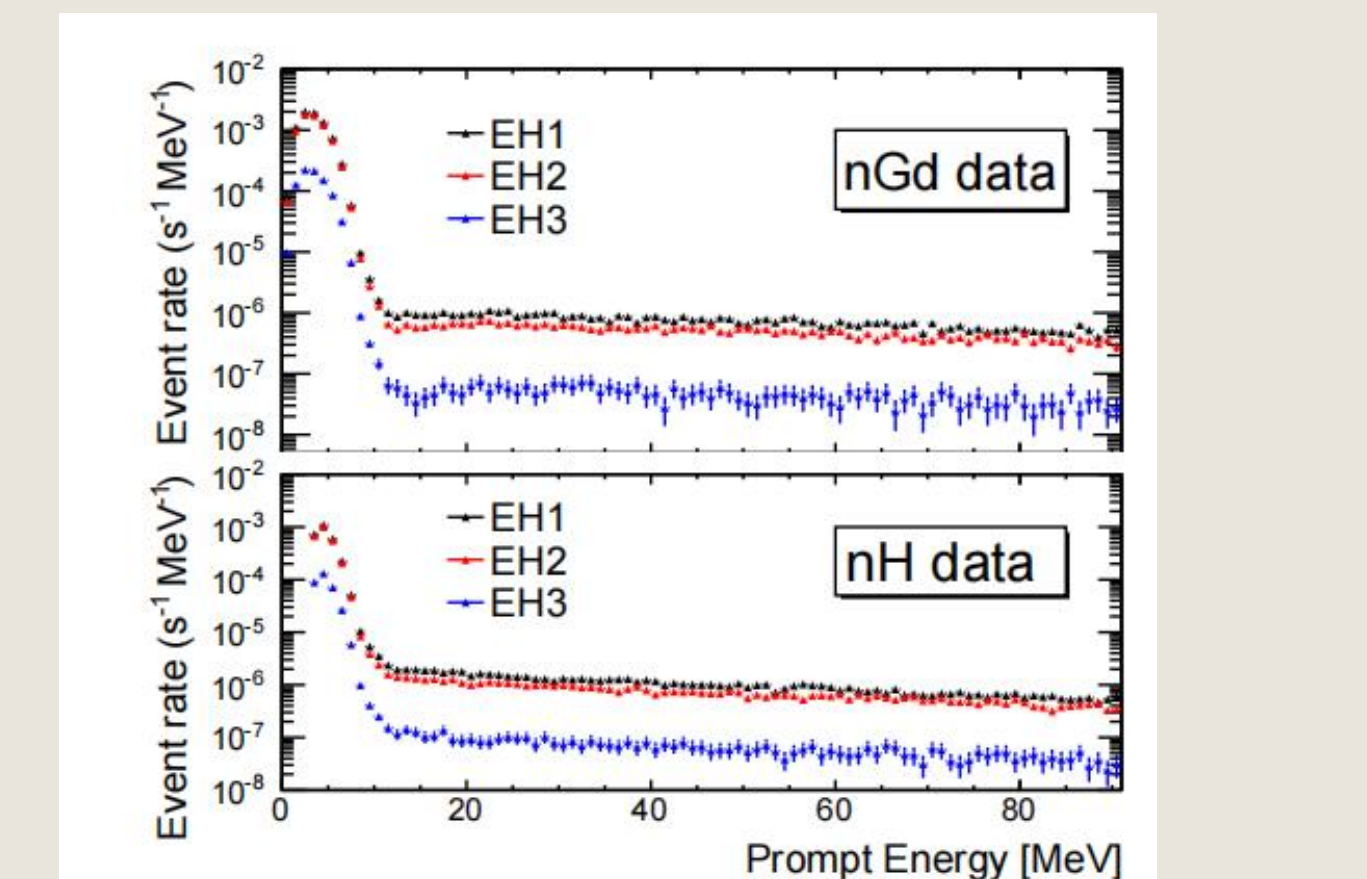
$$\Phi_{FD} = \frac{N_\nu}{N_p \int \sigma(E_\nu) \epsilon(E_\nu) \phi(E_\nu) dE_\nu}, \quad \phi(E_\nu) = \frac{1}{T^3 F_2(\eta)} \frac{E_\nu^2}{e^{E_\nu/T - \eta} + 1},$$

- Where N_p is the number of target protons, $\sigma(E_\nu)$ is the IBD cross-section and $\epsilon(E_\nu)$ is the detector efficiency.
- For the Monochromatic Spectra,

$$\Phi_D(E_\nu) = \frac{N_\nu}{N_p \sigma(E_\nu) \epsilon(E_\nu)}.$$

4. Data Analysis

- IBD candidate selection [10]
 - Background
 - Candidates and background comparison(result for GW150914 as an example)
 - Detection efficiency
- In all possible cases, we used the realtime estimation, i.e. the result in ± 5 days around a GW, for background estimation. But when the background is very low, we used the average of the entire data set.
- Average background rate (per second per antineutrino detector).



	EH1	EH2	EH3
nGd			
Low E	$(7.65 \pm 0.01) \times 10^{-3}$	$(6.82 \pm 0.01) \times 10^{-3}$	$(8.45 \pm 0.01) \times 10^{-4}$
High E	$(6.35 \pm 0.04) \times 10^{-5}$	$(4.32 \pm 0.04) \times 10^{-5}$	$(3.83 \pm 0.08) \times 10^{-6}$
nH			
Low E	$(28.75 \pm 0.04) \times 10^{-4}$	$(25.76 \pm 0.03) \times 10^{-4}$	$(3.25 \pm 0.01) \times 10^{-4}$
High E	$(9.20 \pm 0.05) \times 10^{-5}$	$(6.30 \pm 0.05) \times 10^{-5}$	$(5.65 \pm 0.10) \times 10^{-6}$

- Low E : $E < 10\text{MeV}$.
- High E : $E > 10\text{MeV}$.

	nGd Low E	nGd High E	nH Low E	nH High E
Candidate	4	0	4	0
BKG. (± 5 days)	6.96 ± 0.08	0.060 ± 0.008	2.32 ± 0.06	0.080 ± 0.009
BKG. (Averaged)	7.65 ± 0.01	0.064 ± 0.001	2.88 ± 0.01	0.092 ± 0.001
Candidate	5	0	1	0
BKG. (± 5 days)	6.95 ± 0.08	0.054 ± 0.007	2.54 ± 0.05	0.072 ± 0.008
BKG. (Averaged)	7.65 ± 0.01	0.064 ± 0.001	2.88 ± 0.01	0.092 ± 0.001
Candidate	4	0	2	0
BKG. (± 5 days)	6.62 ± 0.08	0.037 ± 0.006	2.37 ± 0.05	0.041 ± 0.006
BKG. (Averaged)	6.82 ± 0.01	0.043 ± 0.001	2.58 ± 0.01	0.063 ± 0.001
Candidate	8	0	1	0
BKG. (± 5 days)	6.46 ± 0.08	0.027 ± 0.005	2.35 ± 0.05	0.056 ± 0.006
BKG. (Averaged)	6.82 ± 0.01	0.043 ± 0.001	2.58 ± 0.01	0.063 ± 0.001
Candidate	0	0	0	0
BKG. (± 5 days)	0.97 ± 0.03	0.004 ± 0.002	0.37 ± 0.02	0.008 ± 0.003
BKG. (Averaged)	0.850 ± 0.001	0.0038 ± 0.0001	0.330 ± 0.001	0.0056 ± 0.0001
Candidate	0	0	0	0
BKG. (± 5 days)	1.00 ± 0.03	0.003 ± 0.002	0.36 ± 0.02	0.007 ± 0.003
BKG. (Averaged)	0.850 ± 0.001	0.0038 ± 0.0001	0.330 ± 0.001	0.0056 ± 0.0001
Candidate	0	0	0	0
BKG. (± 5 days)	0.97 ± 0.03	0.004 ± 0.002	0.37 ± 0.02	0.008 ± 0.003
BKG. (Averaged)	0.850 ± 0.001	0.0038 ± 0.0001	0.330 ± 0.001	0.0056 ± 0.0001
Candidate	1	0	0	0
BKG. (± 5 days)	0.97 ± 0.03	0.007 ± 0.002	0.36 ± 0.02	0.005 ± 0.002
BKG. (Averaged)	0.850 ± 0.001	0.0038 ± 0.0001	0.330 ± 0.001	0.0056 ± 0.0001

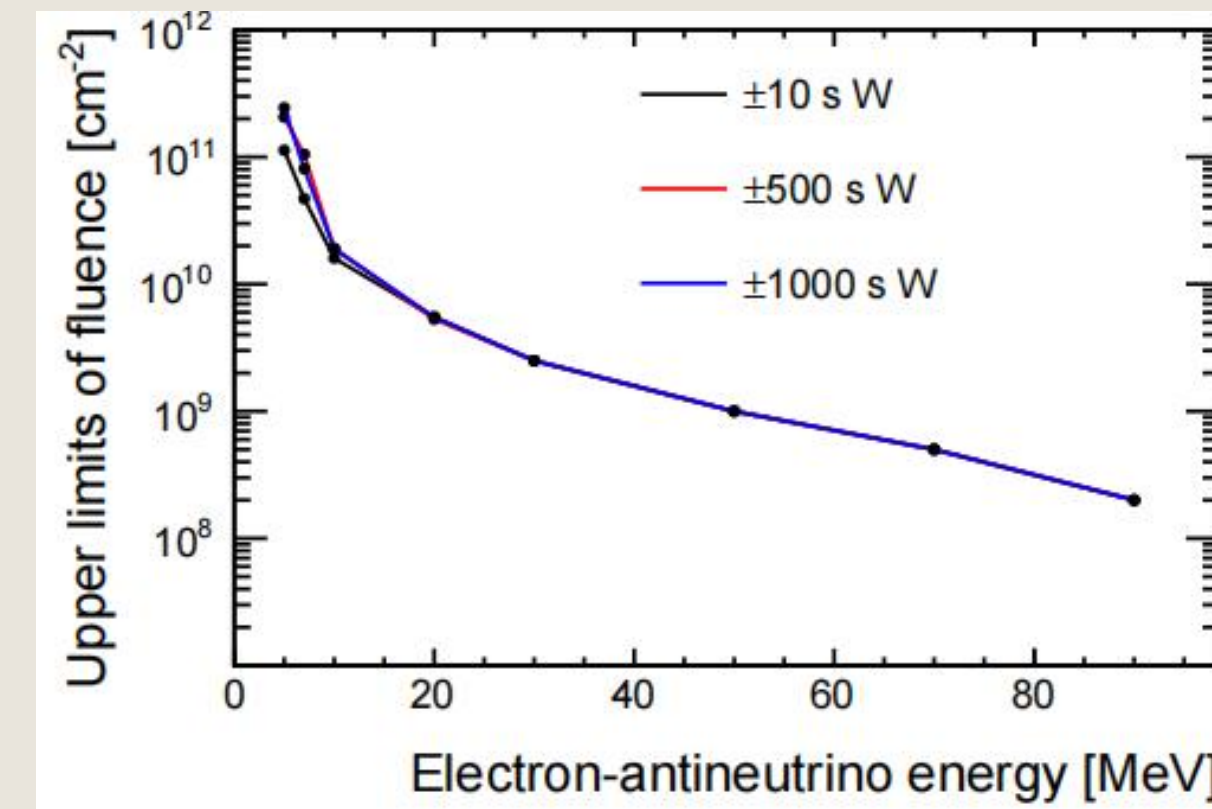
- In all cases, the number of candidates is consistent with the background number.

5. Upper Limits on $\bar{\nu}_e$ Fluence

- Under two different neutrino spectrum assumptions, upper limits (90% C.L.) of fluence for three search time windows are listed.

Fluence ($\times 10^{10} \text{ cm}^{-2}$)	Monochromatic Spectra								Fermi-Dirac Spectrum
E_ν	5 MeV	7 MeV	10 MeV	20 MeV	30 MeV	50 MeV	70 MeV	90 MeV	(1.8, 100) MeV
± 10 s	11.3	4.7	1.6	0.55	0.25	0.10	0.05	0.02	0.70
± 500 s	20.6	10.5	1.9	0.53	0.25	0.10	0.05	0.02	0.63
± 1000 s	24.4	8.1	1.9	0.55	0.25	0.10	0.05	0.02	0.54

- For the Monochromatic Spectra, we plot the results in the table above below.



6. References

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