Phenomenological aspects of A_5' modular symmetry on linear seesaw with leptogenesis

Mitesh Kumar Behera^a, Rukmani Mohanta^b





 ${
m miteshbehera 1304@gmail.com^a, rukmani 98@gmail.com^b} \ {
m University of Hyderabad}^{a,b}$

Abstract

- The uniqueness of this work is that it employs the Linear Seesaw mechanism in **SUSY** to create a model that explains neutrino oscillation data under A'_5 modular symmetry.
- The Standard Model (SM) is supplemented by three RH neutrinos (N_{R_i}) and three Sterile neutrinos (S_{L_i}) , as well as one weighton that eliminates unnecessary components in the superpotential and maintains the model consistent.

Introduction

• A_5' Modular Symmetry: The utilisation of flavon fields is minimised with $\Gamma_5' \approx A_5'$ modular symmetry, and the modular group is the group of linear fractional transformations occurring on a complex variable. τ , used in expressing Dedekind eta-function $\eta(\tau)$.

$$\tau \longrightarrow \gamma \tau = \frac{a\tau + b}{c\tau + d}$$
, where $a, b, c, d \in \mathbb{Z}$

and
$$ad - bc = 1$$
, $Im[\tau] > 0$,

- $\eta(\tau) = e^{\frac{i\pi\tau}{12}} \prod_{n=1}^{\infty} (1 e^{2in\tau\pi}).$ (1)
- The dimension of Γ_5' is 5k+1 hence their are six basis vectors used is expressing lowest weight Yukawa coupling $Y_6^{(1)}$

$$\hat{e}_1 = 1 + 3q + 4q^2 + 2q^3 + q^4 + 3q^5 + 6q^6 + \cdots,$$

$$\hat{e}_2 = q^{1/5} \left(1 + 2q + 2q^2 + q^3 + 2q^4 + 2q^5 + \cdots \right),$$

$$\hat{e}_3 = q^{2/5} \left(1 + q + q^2 + q^3 + 2q^4 + q^6 + \cdots \right) ,$$

$$\hat{e}_4 = q^{3/5} \left(1 + q^2 + q^3 + q^4 - q^5 + \cdots \right) ,$$

$$\widehat{e}_5 = q^{4/5} \left(1 - q + 2q^2 + 2q^6 - 2q^7 + \cdots \right) ,$$

$$\widehat{e}_{6} = q \left(1 - 2q + 4q^{2} - 3q^{3} + q^{4} + 2q^{5} + \cdots \right) .$$
 (2)

• Linear Seesaw: The light neutrino mass matrix under the linear seesaw in the flavor basis of (ν_L, N_R, S_L) is expressed as

$$\mathbb{M} = \begin{pmatrix} 0 & M_D & M_{LS} \\ M_D^T & 0 & M_{RS} \\ M_{LS}^T & M_{RS}^T & 0 \end{pmatrix},$$

• The resulting mass formula

$$m_{\nu} = M_D M_{RS}^{-1} M_{LS}^T + \text{transpose.}$$
 (3)

Conclusion

- We investigated the linear seesaw mechanism under modular A'_5 symmetry and the local $U(1)_{B-L}$, which aids in the prohibition of needless terms and the constraint structures of crucial Yukawa interactions.
- In addition, our model is good enough to handle Leptogenesis.

References

- F. Feruglio, doi:10.1142/9789813238053_0012 [arXiv:1706.08749 [hep-ph]].
- X. G. Liu and G. J. Ding, JHEP **08**, 134 (2019) doi:10.1007/JHEP08(2019)134 [arXiv:1907.01488 [hep-ph]].
- M. K. Behera and R. Mohanta, Front. in Phys. 10, 854595 (2022) doi:10.3389/fphy.2022.854595 [arXiv:2201.10429 [hep-ph]].

Acknowledgements

- I want to acknowledge DST-INSPIRE for its financial support.
- I want to acknowledge IOE-UOH and SERB, IN-DIA for financial support.
- I want to acknowledge ICHEP-2022 for its financial support.

Model Framework

• The complete superpotential is given by

$$W = A_{M_l} \left[(L_L l_R^c)_{\mathbf{3}} Y_{\mathbf{3}}^{k_Y} \right] H_d + \mu H_u H_d + G_d \left[(L_L N_R^c)_{\mathbf{5}} Y_{\mathbf{5}}^{(2)} \right] H_u + G_{ls} \left[(L_L S_L)_{\mathbf{4}} H_u \sum_{i=1}^2 Y_{\mathbf{4},i}^{(6)} \right] \frac{\zeta}{\Lambda} + B_{rs} \left[(S_L N_R^c)_{\mathbf{5}} \sum_{i=1}^2 Y_{\mathbf{5},i}^{(6)} \right] \zeta'.$$

$$(4)$$

Fields	c	,, <u>c</u>	$ au^c$	T. T	N^c	S_{τ}	H ,	<i>\</i>	<i>>1</i>
	e_R	μ_R	$^{\prime\prime}R$	DL	IV R	$\mathcal{O}L$	$H_{u,d}$	ς	<u> </u>
$SU(2)_L$	1	1	1	2	1	$\mid 1 \mid$	2	1	1
$U(1)_Y$	1	1	1	$-\frac{1}{2}$	0	0	$\frac{1}{2}, -\frac{1}{2}$	0	0
$U(1)_{B-L}$	1	1	1	-1	1	0	0	1	-1
A_5'	1	1	1	3	3'	3'	1	1	1
k_I	1	3	5	1	1	4	0	1	1

Table 1: The particle spectrum and their charges under the symmetry groups $SU(2)_L \times U(1)_Y \times U(1)_{B-L} \times A_5'$ while k_I represents the modular weight.

Mass matrices expression

$$M_{D} = \frac{v_{u}G_{d}}{\sqrt{3}0} \begin{pmatrix} \sqrt{3}(Y_{\mathbf{5}}^{(2)})_{1} & (Y_{\mathbf{5}}^{(2)})_{4} & (Y_{\mathbf{5}}^{(2)})_{3} \\ (Y_{\mathbf{5}}^{(2)})_{5} & -\sqrt{2}(Y_{\mathbf{5}}^{(2)})_{3} - \sqrt{2}(Y_{\mathbf{5}}^{(2)})_{2} \\ (Y_{\mathbf{5}}^{(2)})_{2} & -\sqrt{2}(Y_{\mathbf{5}}^{(2)})_{5} - \sqrt{2}(Y_{\mathbf{5}}^{(2)})_{4} \end{pmatrix}, M_{RS} = \frac{v_{\zeta}B_{rs}}{\sqrt{6}0} \begin{pmatrix} 2\left(\sum\limits_{i=1}^{2}Y_{\mathbf{5},i}^{(6)}\right)_{1} & -\sqrt{3}\left(\sum\limits_{i=1}^{2}Y_{\mathbf{5},i}^{(6)}\right)_{4} - \sqrt{3}\left(\sum\limits_{i=1}^{2}Y_{\mathbf{5},i}^{(6)}\right)_{4} - \sqrt{3}\left(\sum\limits_{1}^{2}Y_{\mathbf{5},i}^{(6)}\right)_{4} - \sqrt{3}\left(\sum\limits_{i=1}^{2}Y_{\mathbf{5},i}^{(6)}$$

$$M_{LS} = \frac{v_u}{2\sqrt{6}} \left(\frac{v_\zeta}{\sqrt{2}\Lambda}\right) G_{ls} \begin{pmatrix} 0 & -\sqrt{2} \left(\sum_{i=1}^{2} Y_{\mathbf{4},\mathbf{i}}^{(6)}\right)_3 & -\sqrt{2} \left(\sum_{i=1}^{2} Y_{\mathbf{4},\mathbf{i}}^{(6)}\right)_2 \\ \sqrt{2} \left(\sum_{i=1}^{2} Y_{\mathbf{4},\mathbf{i}}^{(6)}\right)_4 & -\left(\sum_{i=1}^{2} Y_{\mathbf{4},\mathbf{i}}^{(6)}\right)_2 & \left(\sum_{i=1}^{2} Y_{\mathbf{4},\mathbf{i}}^{(6)}\right)_1 \\ \left(\sum_{i=1}^{2} Y_{\mathbf{4},\mathbf{i}}^{(6)}\right)_1 & \left(\sum_{i=1}^{2} Y_{\mathbf{4},\mathbf{i}}^{(6)}\right)_4 & -\left(\sum_{i=1}^{2} Y_{\mathbf{4},\mathbf{i}}^{(6)}\right)_3 \end{pmatrix}.$$
 (5)

Leptogenesis

• To account for leptogenesis, a higher dimensional mass term for the Majorana fermion (N_R) is introduced as in eqn.(6), resulting in a tiny mass splitting between the heavy fermions, where G_R is the coupling.

$$\mathcal{W}_{M_R} = -G_R \left[\sum_{i=1}^2 Y_{5,i}^{(4)} N_R^c N_R^c \right] \frac{\zeta'^2}{\Lambda}. \tag{6}$$

Boltzmann equations

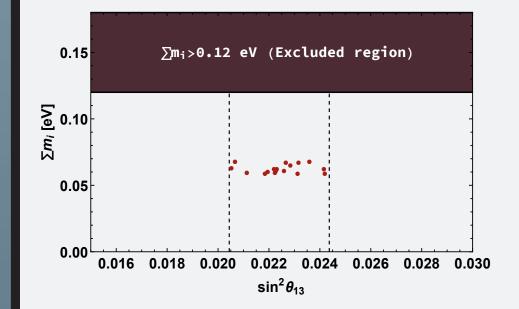
$$\frac{dY_{N^{-}}}{dz} = -\frac{z}{sH(M_{1}^{-})} \left[\left(\frac{Y_{N^{-}}}{Y_{N^{-}}^{eq}} - 1 \right) \gamma_{D} + \left(\left(\frac{Y_{N^{-}}}{Y_{N^{-}}^{eq}} \right)^{2} - 1 \right) \gamma_{S} \right],$$

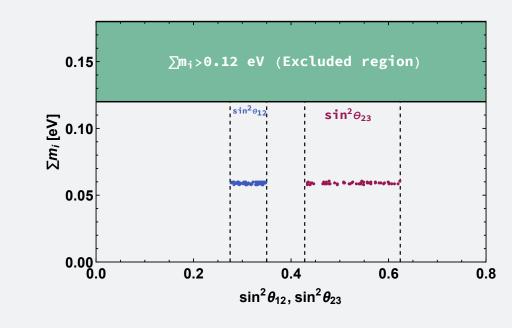
$$\frac{dY_{B-L}}{dz} = -\frac{z}{sH(M_{-}^{-})} \left[\epsilon_{N^{-}} \left(\frac{Y_{N^{-}}}{Y_{eq}^{eq}} - 1 \right) \gamma_{D} - \frac{Y_{B-L}}{Y_{eq}^{eq}} \frac{\gamma_{D}}{2} \right].$$
(7)

Here, s denotes the entropy density, $z = M_1^-/T$ and the equilibrium number densities are given by

$$Y_{N^{-}}^{\text{eq}} = \frac{135g_{N^{-}}}{16\pi^4 g_{\star}} z^2 K_2(z), \quad Y_{\ell}^{\text{eq}} = \frac{135\zeta(3)g_{\ell}}{8\pi^4 g_{\star}}.$$
 (8)

Results





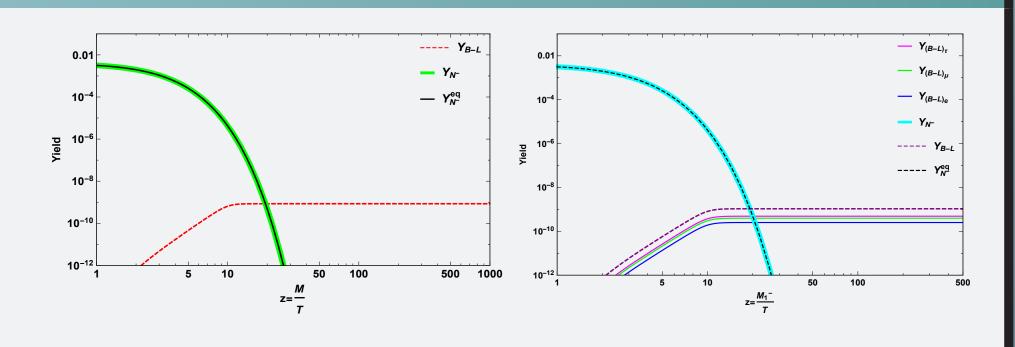


Figure 1: Extreme left plot represents $\sin^2 \theta_{13}$, $(\sin^2 \theta_{12} \text{ and } \sin^2 \theta_{23})$ (middle) versus $\sum m_i[\text{eV}]$, right middle and extreme right plot shows evolution of Y_{B-L} (dashed) as a function of $z = M_1^-/T$ for one flavor approximation and flavoured case respectively.