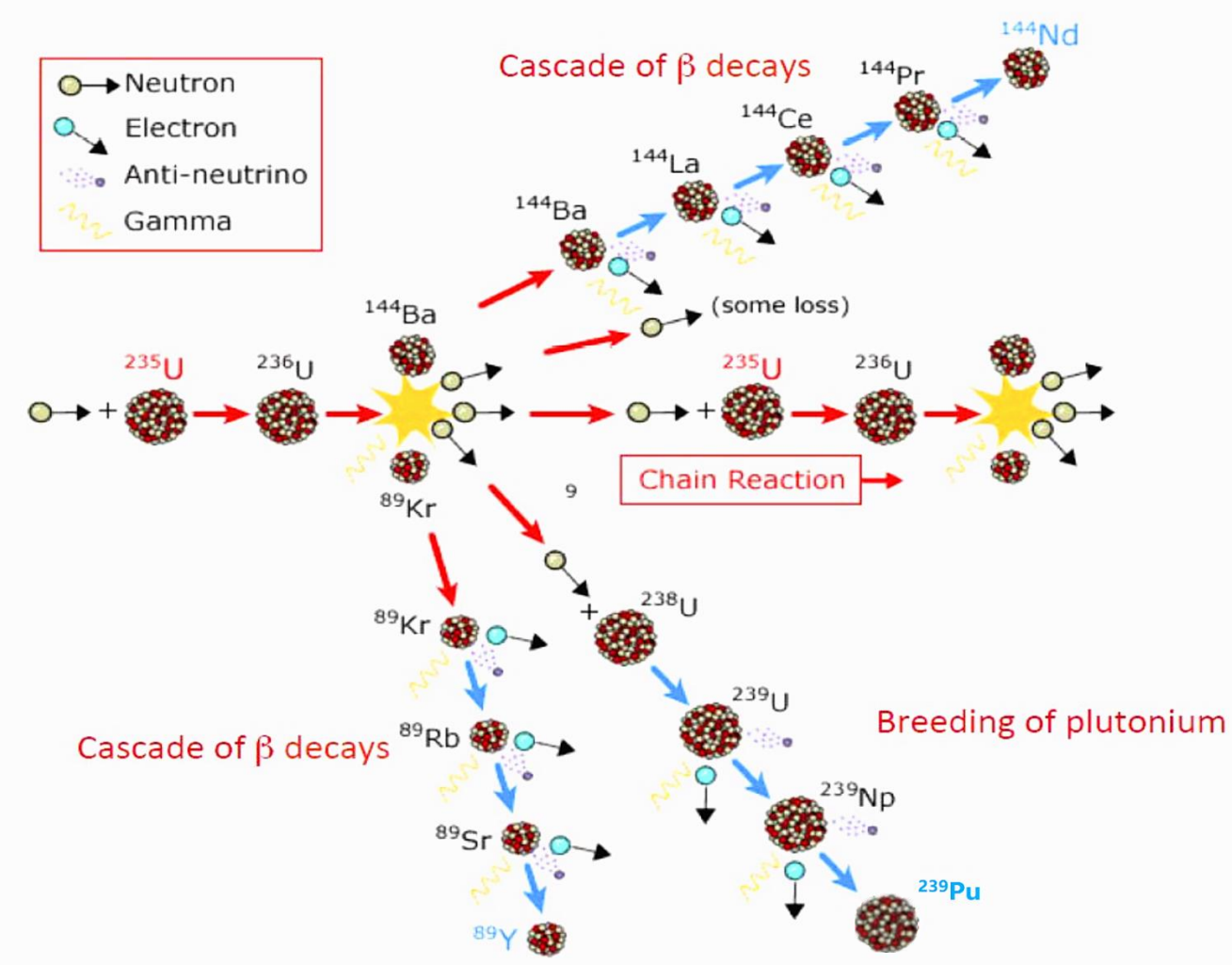
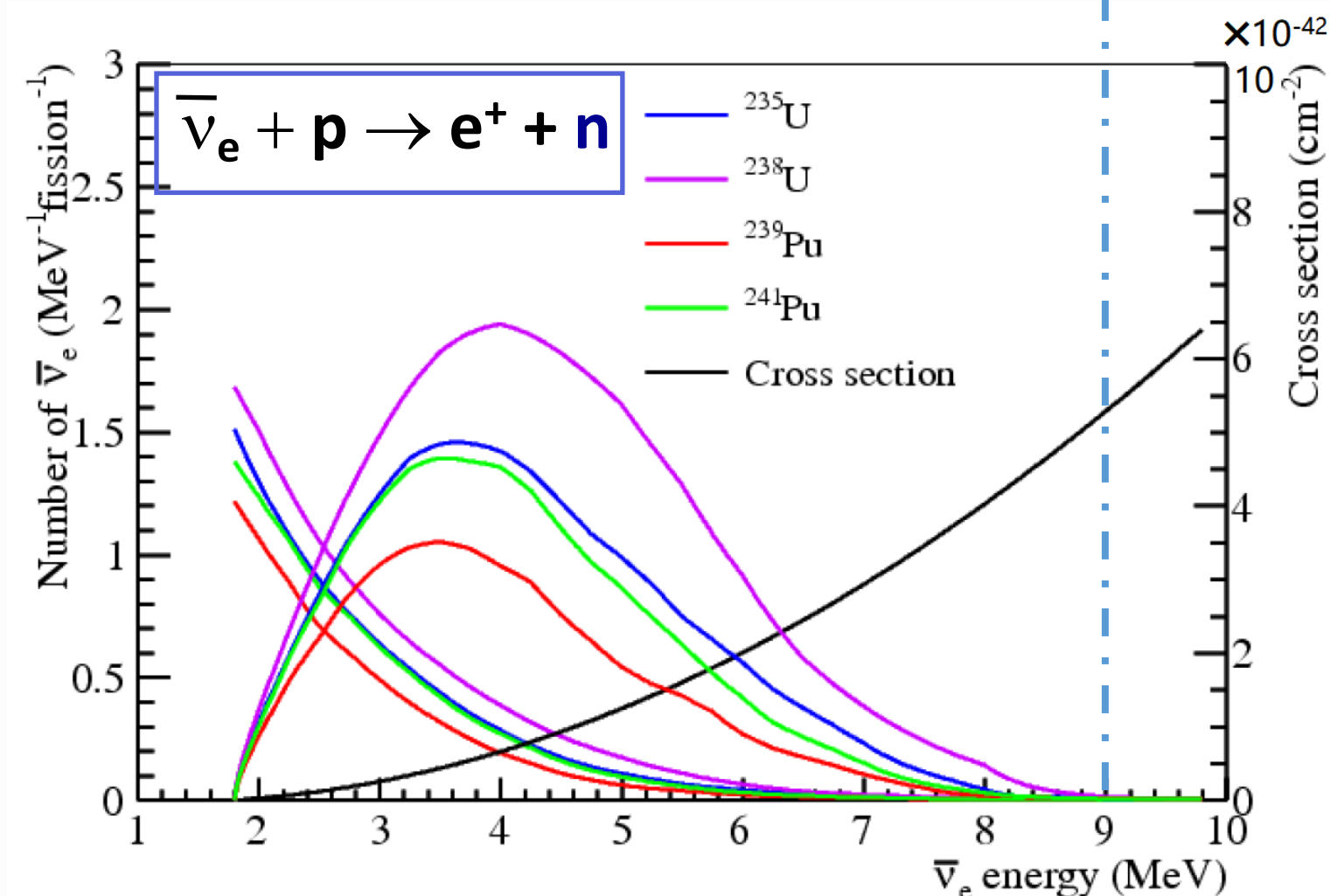


Reactor neutrinos

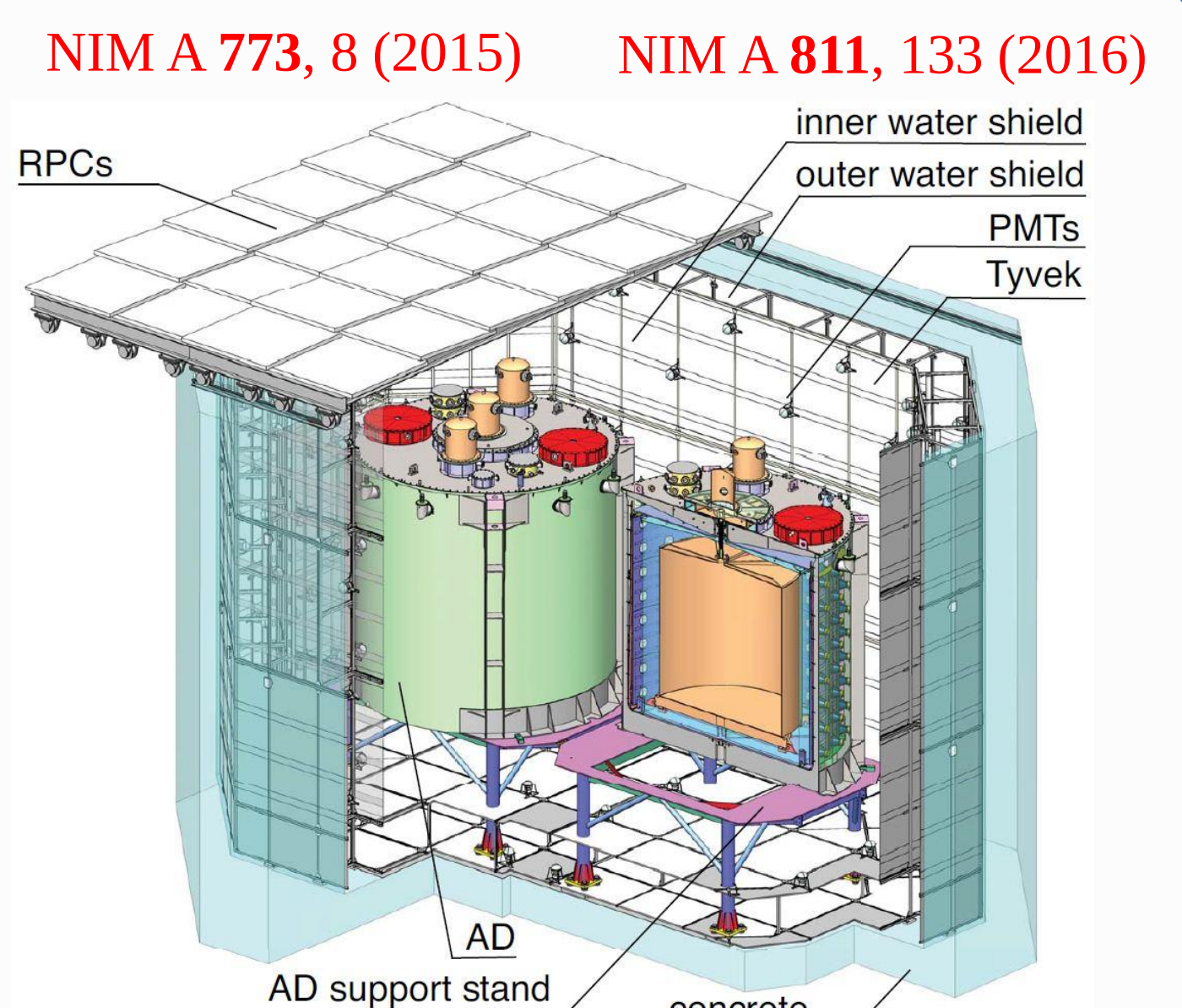
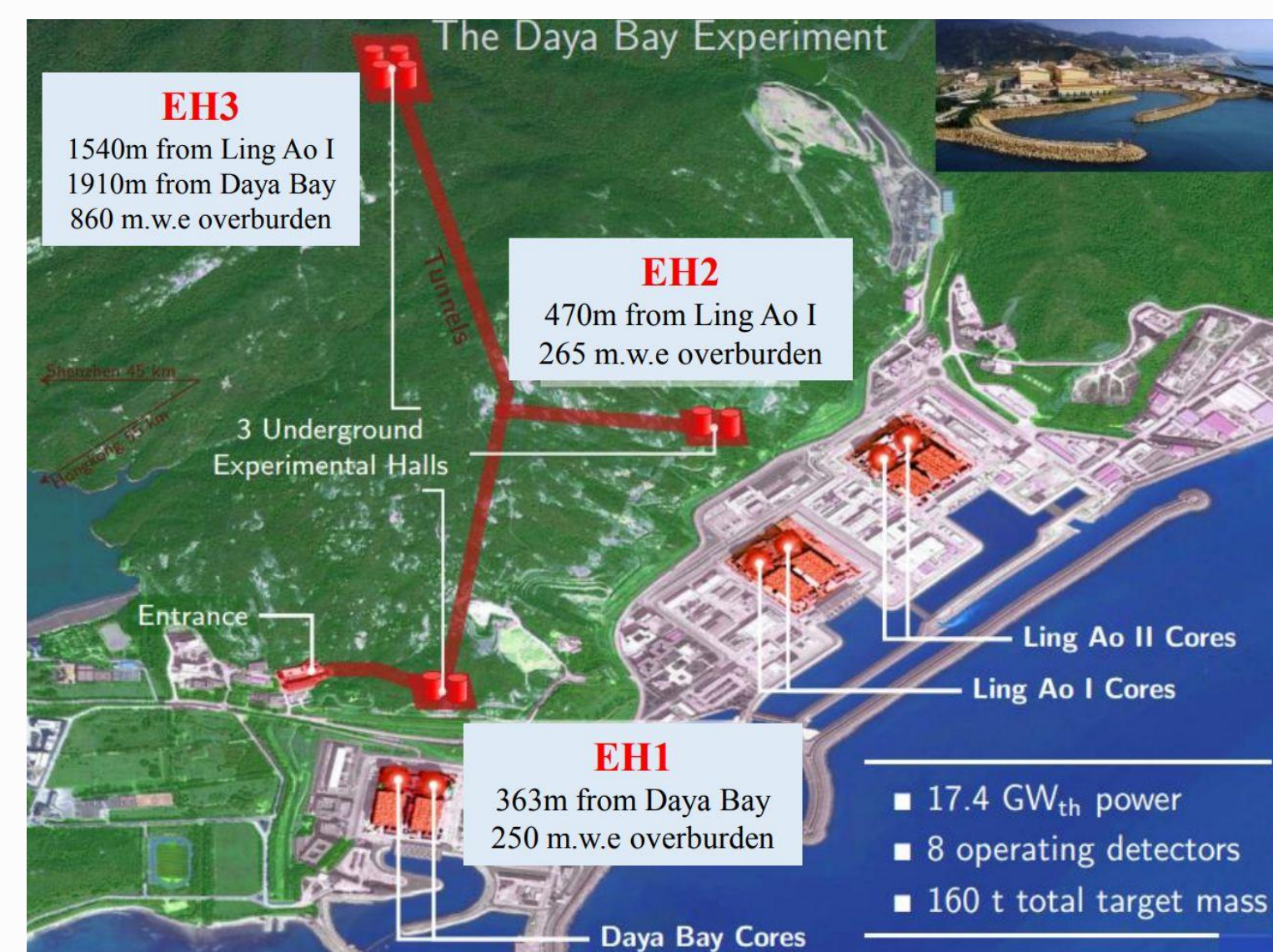


- Nuclear reactors produce pure $\bar{\nu}_e$ from beta decays of fission daughters.
- Reactor neutrinos have been an important tool for both discovery and precision measurement in the history of neutrino studies [1].
- Previous measurements of reactor neutrinos only focus on the prompt energy (E_p) less than 9 MeV.
 - Daya Bay's high statistic samples enable new studies for high energy reactor neutrinos > 9 MeV [2].

How about high-energy reactor neutrinos?



Daya Bay experiment



- 6 low enriched uranium (LEU) commercial reactors, each with 2.9 GW thermal power.
- Eight identically designed underground detectors deployed at different baselines [3].
- The largest dataset (~6 million) of reactor neutrinos collected from December 2011 to December 2022.

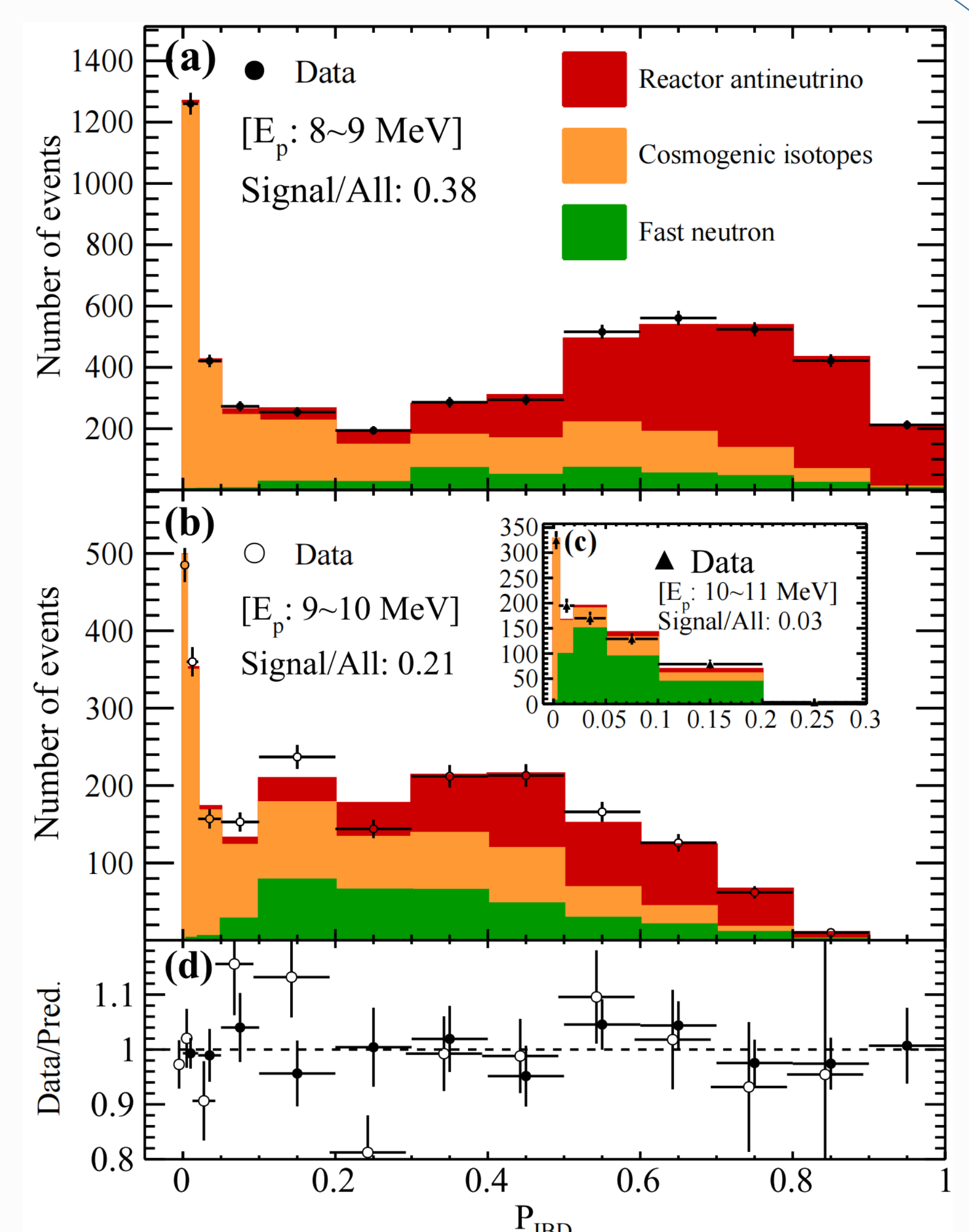
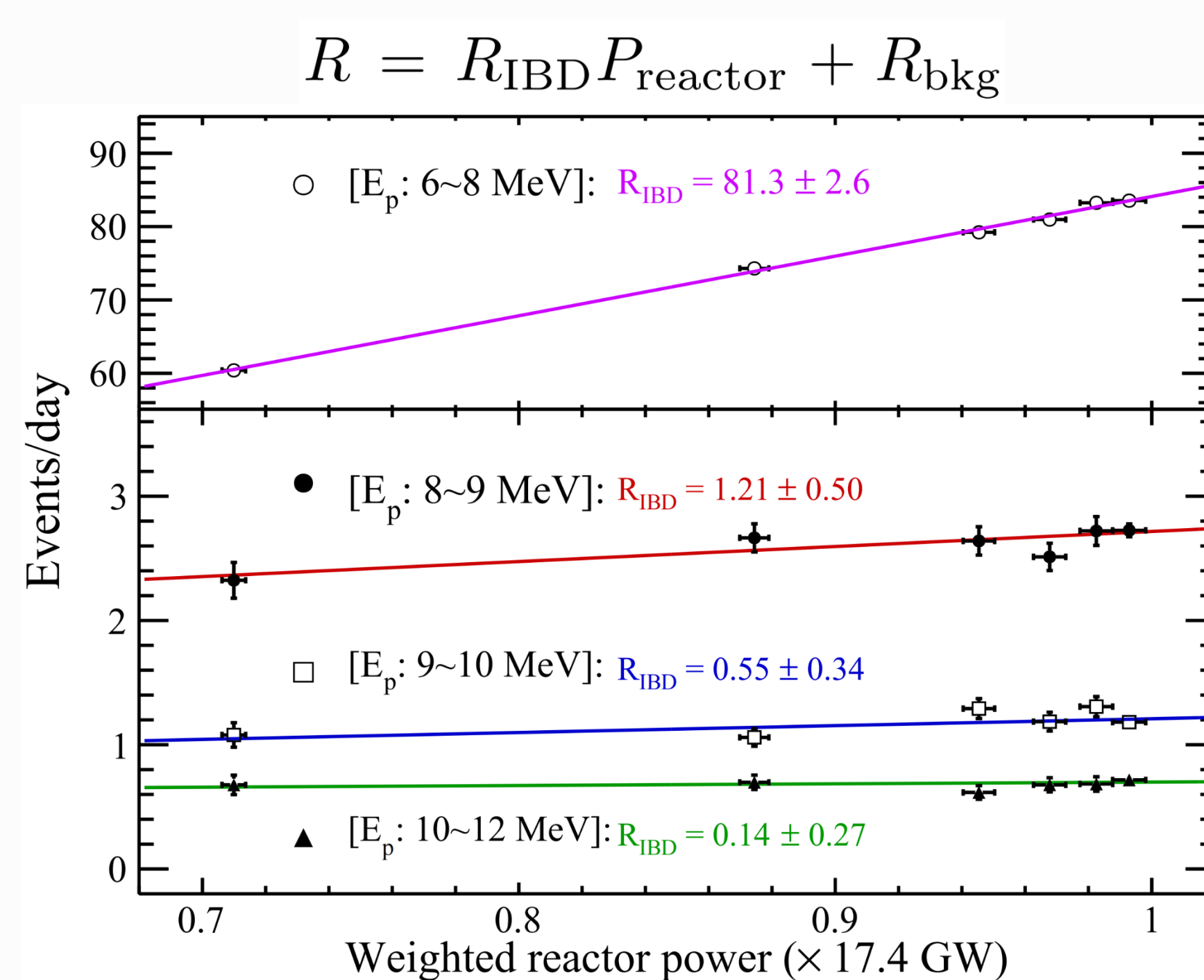
Data selection

- Inverse beta decay (IBD) events: the interaction candidates are selected following the criteria similar to "Selection A" described in Ref. [4]
 - $E_p \approx E_{\bar{\nu}} - 0.8 \text{ MeV}$
- Main background: cosmogenic fast neutrons and cosmogenic isotopes.
- Challenge: large background to signal ratio in the 8-12 MeV prompt energy region.

Multivariate analysis

- The probability distribution function (PDF) $F(\mathbf{r}; \Delta\mathbf{t}, z, w)$ is constructed in each prompt-energy bin for the multivariate event-by-event fitter with pull terms $g(\epsilon)$:

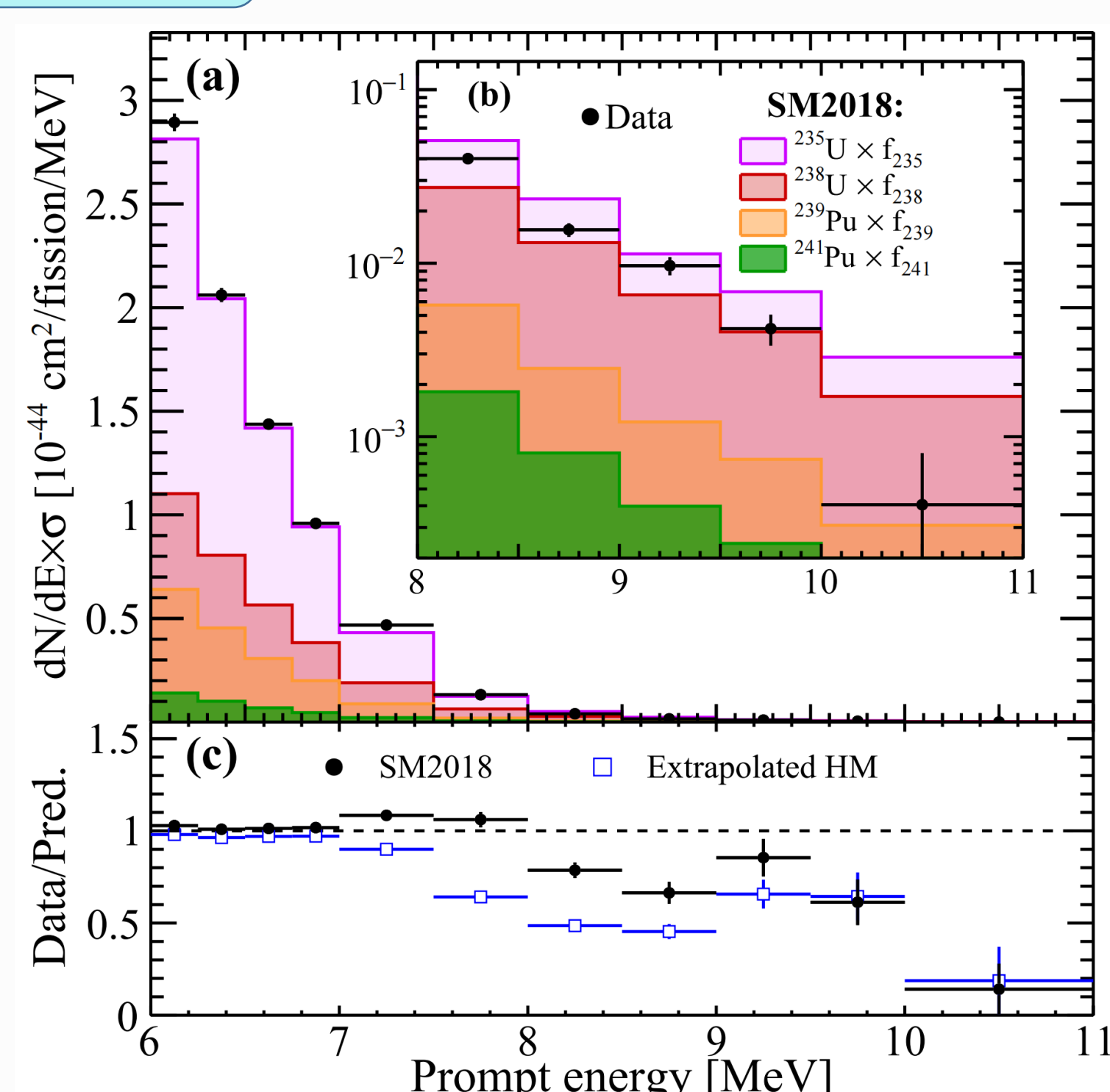
$$F(\mathbf{r}; \Delta\mathbf{t}, z, w) = \sum_p r_p f_p(\Delta\mathbf{t}) h_p(z) k_p(w) \quad \chi^2(\mathbf{r}) = -2 \sum [\log F(\mathbf{r}; \Delta\mathbf{t}, z, w)] + g(\epsilon)$$
 - p : event types (IBD, cosmogenic isotopes, or fast neutron);
 - $\mathbf{r}$: vector of r_p , the ratio of evens in type p over the total number of events in each E_p bin;
 - $\Delta\mathbf{t}$: vector for time difference between event and its preceding muon events;
 - z : vertical position of prompt signal;
 - w : weighted reactor power at the occurrence of the candidate;
- With the best-fit values of $\mathbf{r}$, the probability of being an IBD signal can be calculated for each event:
 - $f_{IBD}(\Delta\mathbf{t})$: time to last muon PDF; $h_{IBD}(z)$: vertical position PDF; $k_{IBD}(w)$: reactor power PDF;
- Separate 2500 signal events from about 9000 IBD candidates above 8 MeV. (using the dataset taken in 1958 calendar days)
- The P_{IBD} distributions from data and predictions are consistent within statistical uncertainty.



$$P_{IBD} = \frac{r_{IBD} f_{IBD}(\Delta\mathbf{t}) h_{IBD}(z) k_{IBD}(w)}{F(\mathbf{r}; \Delta\mathbf{t}, z, w)}$$

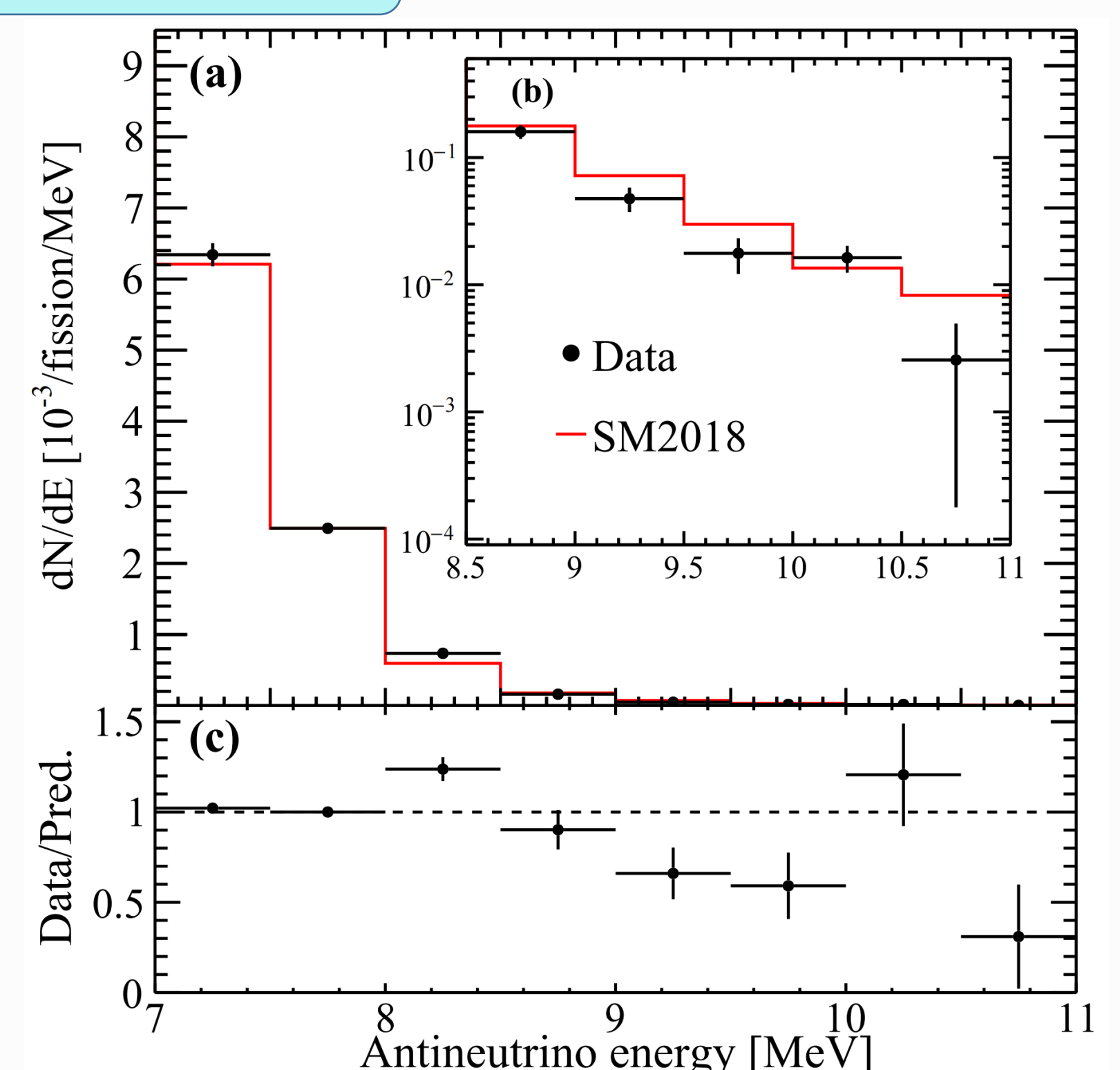
Prompt-energy Spectrum

- Measured IBD yield compared with a recent summation model (SM2018 [5]):
 - 3% larger for 6-8 MeV;
 - 29% smaller for 8-11 MeV;
- Above 10 MeV, a non-zero signal rate of greater than 5σ is predicted by the SM2018 model while the significance is only 1σ for the data.



Unfolded Antineutrino Spectrum

- We provide the unfolded antineutrino spectrum above 7 MeV as a data-based reference for other experiments.
- Significance in rejecting the hypothesis of no reactor antineutrinos above 10 MeV is determined to be 6.2σ .
- For the first time, this work extends the energy region of reactor antineutrinos above 10 MeV by direct measurement.



References:

- L. J. Wen et al. Ann.Rev.Nucl.Part.Sci. 67 (2017) 183-211
- F. P. An et al. (Daya Bay Collaboration), arXiv:2203.06686 [hep-ex]
- F. P. An et al. (Daya Bay Collaboration), Phys. Rev. D 95, 072006 (2017)
- F. P. An et al. (Daya Bay), Phys. Rev. D 95, 072006 (2017)
- M. Estienne et al., Phys. Rev. Lett. 123, 022502 (2019)