



<u>An absolute *v*-mass measurement</u> at the DUNE detector

Federica Pompa

In collaboration with: Dr. Olga Mena Dr. Michel Sorel Dr. Francesco Capozzi



07 July 2022







The most elusive of the known fermions

 ${\cal U}$

Evidence that "new physics" is required



IN THE STANDARD MODEL

 u_{lpha} , $lpha=e,\mu, au$

 $m_{\nu} = 0$







THE STANDARD MC

 $2, \mu, \tau$ ν_{α} ,

 $m_{\nu}=0$

ν FLAVOUR OSCILLATIONS

$$\nu_{\alpha} = \sum_{i=1}^{3} U_{\alpha i} \nu_{i}$$
$$m_{\nu} \neq 0$$







<u>arXiv:2106.15267</u>

From $0\nu\beta\beta$ measurements: KamLAND-Zen

 $m_{\beta\beta} \in [36, 156] \text{ meV (90\% CL)}$

<u>arXiv:2203.02139</u>



ν FLAVOUR OSCILLATIONS $\nu_{\alpha} = \sum_{i=1}^{} U_{\alpha i} \, \nu_{i}$

 $m_{\nu} \neq 0$

From kinematic measurements: KATRIN $m_{eta} < 0.8 \; {\rm eV} \; (90\% \; {\rm CL})$ <u>arXiv:2105.08533</u>

Time-of-flight constraints: Kamiokande-II (SN1987A)

 $m_{\nu} < 5.7 \text{ eV}$ (95% CL)

<u>arXiv:1002.3349</u>





Copiously produced during the Supernova explosion

99% of the released energy emitted through ν and $\bar{\nu}$ of all flavors with mean energies of $\mathcal{O}(10)$ MeV



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Aready detected! **SN1987A** $M \approx 20 M_{\odot}$ $@ \sim 50 \text{ kpc}$ Kamiokande-II: Phys. Rev. Lett. 58, 1490





Neutronization burst first 50 ms of the explosion

 $p + e^- \rightarrow n + \nu_e$





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Liquid Argon Time Projection Chamber

Charged-current electron neutrino interactions on Ar nuclei

 $\nu_e + {}^{40}Ar \rightarrow e^- + {}^{40}K^*$

Why DUNE?





Neutronization burst first 50 ms of the explosion

 $p + e^- \rightarrow n + \nu_e$



Effects of non-zero neutrino mass:

$$t_{i} = \delta t_{i} - \Delta t_{i} + t_{off}$$
$$\Delta t_{i} = \frac{D}{2c} \left(\frac{m_{\nu}}{E_{i}}\right)^{2}$$



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$R(t, E) = N_{target} \sigma_{CC}(E) \Phi_{\nu_e}(t, E)$

DUNE far-detector fiducial mass: 40 kton of LAr





Non oscillated and oscillated fluxes (Large Mixing Angle MSW)

 $\Phi_{\nu_{\rho}} = p \ \Phi^{0}_{\nu_{\rho}} + (1-p) \ \Phi^{0}_{\nu_{r}}$

Normal mass Ordering : $p = |U_{e3}|^2$ Inverted mass Ordering : $p = |U_{e2}|^2$

Likelihood analysis

- $(\delta t_i, E_i)$ generation by fixing D • $L(t, m_{\nu}) = \prod_{i=1}^{R} \int R(t_i, E_i) G_i(E, 0.1E) dE$ • $\chi^2(t_i, m_\nu) = -2\log(L)$
- $\Delta \chi^2(m_\nu) = \chi^2(m_\nu) \chi^2_{min}(m_\nu)$



 $R(t, E) = N_{target} \sigma_{CC}(E) \Phi_{\nu_e}(t, E)$







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Results: D = 10 kpc









UPPER BOUNDS ON

$$m_{\nu} = \sqrt{\sum_{i=1}^{3} |U_{ei}|}$$

 $m_{\nu} \leq 0.51^{+0.20}_{-0.19} \text{ eV}$ $m_{\nu} \le 0.91^{+0.30}_{-0.33} \text{ eV}$ $m_{\nu} \leq 2.01^{+0.69}_{-0.55} \text{ eV}$

 $m_{\nu} \leq 0.56^{+0.20}_{-0.21} \text{ eV}$ $m_{\nu} \leq 0.85^{+0.30}_{-0.25} \text{ eV}$ $m_{\rm e} < 1.65^{+0.54} {\rm eV}$ -0.40











Results: including Earth matter effects

Affect only the Inverted mass Ordering in the neutrino channel

<u>arXiv:9702343</u> <u>arXiv:1205.5254</u>



 $\Phi_{\nu_e}^{\oplus NO}(t, E) = \Phi_{\nu_e}^{NO}(t, E)$ $\Phi_{\nu_e}^{\oplus IO}(t,E) = P_{2e} \Phi_{\nu_e}^0(t,E) + (1 - P_{2e}) \Phi_{\nu_x}^0(t,E)$



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$\mathcal{T}_{\alpha\beta} = \mathcal{T}(\overline{P_{det}P_1}) \,\mathcal{T}(\overline{P_1P_2}) \cdots \mathcal{T}(\overline{P_MP_{prod}}) \quad \rightarrow \quad P_{2e}(E, \cos\theta) = \mathcal{T}_{e\beta} \cdot U_{PMNS, 2}$

<u>arXiv:9702343</u> <u>arXiv:1205.5254</u>

$$\Phi_{\nu_e}^{\bigoplus NO}(t, E) = \Phi_{\nu_e}^{NO}(t, E)$$

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D = 10 kpc $M = 8.8 M_{\odot}$

EFFECTS ON RATE

Mild variation respect to the case of absence of Earth matter







This project has received funding and support from the European Union's Horizon 2020 research and innovation programme under the Marie Skłodowska-Curie grant agreement No 860881-HIDDeN

The neutrino signal coming from the Supernova neutronization burst, visible only in the ν_{ρ} spectrum, constitutes an important tool to constrain the absolute value of the neutrino mass and it can give a complementary (and independent) measurement to β -decays and cosmology.



Supernova parameters uncertainties: luminosity





The neutronization burst results to be a robust, **model independent** prediction of the Supernova models.

Very slight variations as a function of progenitor mass (left panel), microphysics of neutrino interactions (middle panel) and equation of state (right panel).



Supernova parameters uncertainties: mean energy





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Dependence on the Supernova distance from the detector

How the 95% CL upper bounds on m_{ν} shift with the Supernova distance D from the detector



Supernova neutrinos emission: details

$$\Phi^0_{\nu_{\beta}}(E,t) = \frac{L_{\nu_{\beta}}(t)}{4\pi D^2} \frac{\varphi_{\nu_{\beta}}(E,t)}{\langle E_{\nu_{\beta}}(t) \rangle} \qquad \Phi^0_{\nu_{\mu}}, \Phi^0_{\nu_{\mu}}$$

$$\varphi_{\nu_{\beta}}(E,t) = \xi_{\beta}(t) \left(\frac{E}{\langle E_{\nu_{\beta}}(t) \rangle}\right)^{\alpha_{\beta}(t)} e^{\left\{\frac{-[\alpha_{\beta}(t)+1]}{\langle E_{\nu_{\beta}}(t) \rangle}\right\}}$$

$$\alpha_{\beta}(t) = \frac{2\langle E_{\nu_{\beta}}(t) \rangle^{2} - \langle E_{\nu_{\beta}}^{2}(t) \rangle}{\langle E_{\nu_{\beta}}^{2}(t) \rangle - \langle E_{\nu_{\beta}}(t) \rangle^{2}}$$

 $\equiv \Phi^0_{\nu_x}$

 $\left| \frac{E}{E} \right\rangle$

Evolution operator definition

 $\mathcal{T}(\overline{P_{j-1}P_j}) = \exp\{-i(H_0 - V_{matter,j}) \cdot l_j\}$ $H_0 = \frac{U_{PMNS} M_{mass} U_{PMNS}^{\dagger}}{2E}$ $V_{matter,j} = \operatorname{diag}\left(\sqrt{2} \ G_F \ \overline{N_j}(x)\right)$ $\overline{N_j}(x) = \frac{1}{l_j} \int_{x_{i-1}}^{x_j} N_j(x) \, dx$ $N_i(x) = \alpha_i + \beta_i x^2 + \gamma_i x^4$



Mikheyev-Smirnov-Wolfenstein effect

Adiabatic or partially adiabatic neutrino flavor conversion in medium with varying density



