

# Impact of radiative corrections on decays of the charged and CP-odd Higgs bosons

Based on Nucl. Phys. B973 (2021) 115518 [hep-ph: 2108.11868]  
hep-ph: 2207.01032

Masashi Aiko (KEK)

Collaborators : Shinya Kanemura (Osaka University), Kodai Sakurai (Tohoku University)



# Introduction

## Standard Model

We have problems that cannot be explained within the SM.

Baryon asymmetry of the universe, Dark matter, Neutrino's tiny mass, etc.

SM must be extended to solve these problems.

## Extended Higgs model

- One  $SU(2)_L$  doublet is a theoretical assumption in the SM.
- The above problems can be solved.

The structure of the Higgs sector is still a mystery.

# Two Higgs doublet model (2HDM)

The model with two scalar doublet  $\Phi_1$  and  $\Phi_2$  with  $Y = 1/2$ .

$$V(\Phi_1, \Phi_2) = m_1^2 |\Phi_1|^2 + m_2^2 |\Phi_2|^2 - m_{12}^2 (\Phi_1^\dagger \Phi_2 + h.c.) \\ + \frac{1}{2} \lambda_1 |\Phi_1|^4 + \frac{1}{2} \lambda_2 |\Phi_2|^4 + \lambda_3 |\Phi_1|^2 |\Phi_2|^2 + \lambda_4 |\Phi_1^\dagger \Phi_2|^2 + \frac{1}{2} \lambda_5 [(\Phi_1^\dagger \Phi_2)^2 + h.c.], \quad \Phi_i = \begin{pmatrix} \omega_i^\pm \\ \frac{1}{\sqrt{2}}(v_i + h_i + iz_i) \end{pmatrix}$$

Softly-broken  $Z_2$  symmetry suppresses flavor-changing neutral current. Glashow, Weinberg, PRD15 (1977)  
Paschos, PRD15 (1966)

- 2HDM is classified into Type-I, II, X and Y.

Barger et al. PRD41 (1990), Aoki et al. PRD80 (2009)

**Scalar particles**  $h$  (SM-like Higgs boson),  $H$ ,  $A$ ,  $H^\pm$

**Parameters**  $v (=246 \text{ GeV})$ ,  $m_h (=125 \text{ GeV})$ ,  $m_H$ ,  $m_A$ ,  $m_{H^\pm}$ ,  $M^2 = m_{12}^2/(s_\beta c_\beta)$ ,  $\tan \beta$ ,  $s_{\beta-\alpha}$

**Higgs couplings**  $g_{hVV} = s_{\beta-\alpha} g_{hVV}^{\text{SM}}$ ,  $g_{hff} = (s_{\beta-\alpha} - c_{\beta-\alpha} \zeta_f) g_{hff}^{\text{SM}}$  ( $\zeta_f = -\tan \beta$  or  $\cot \beta$ )

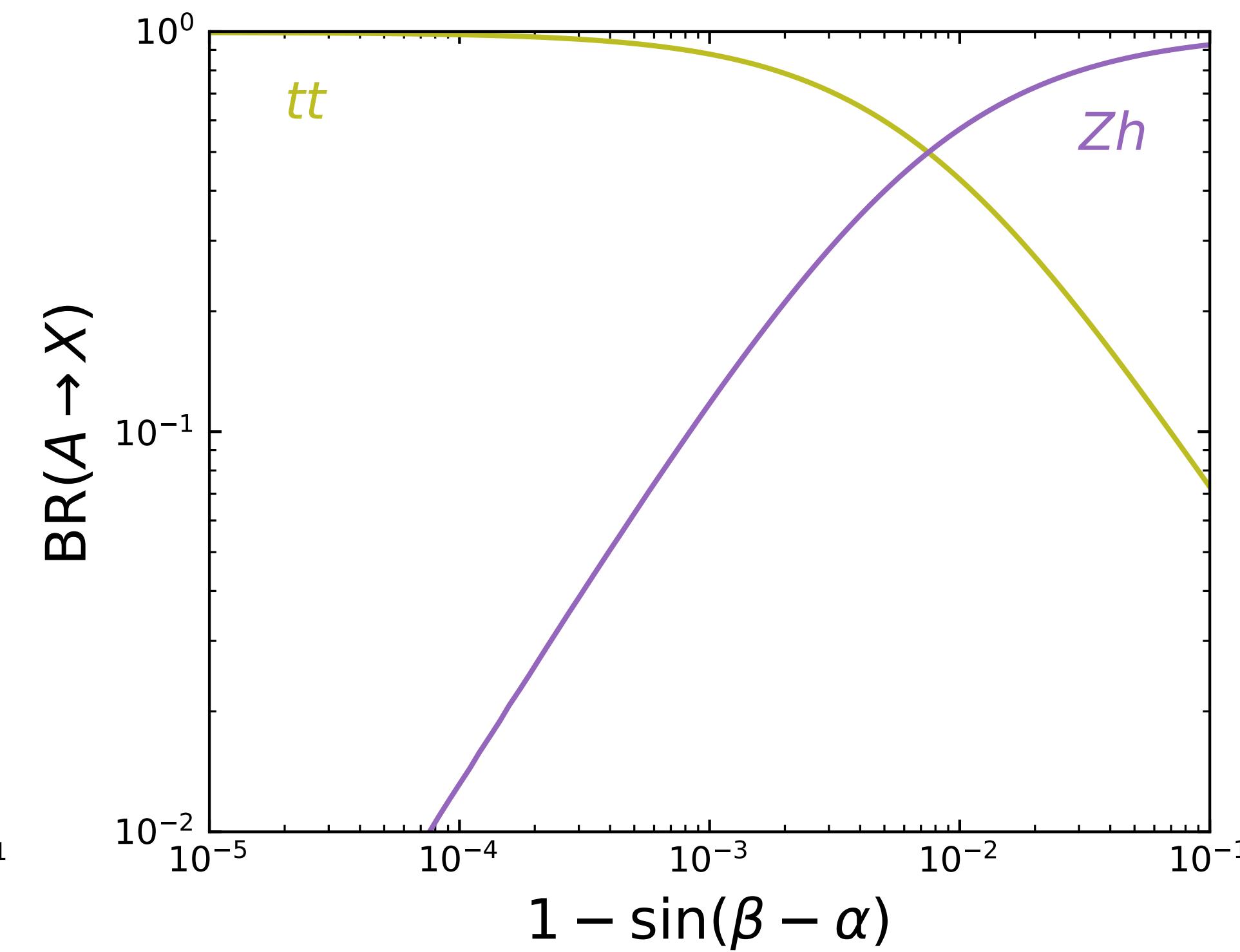
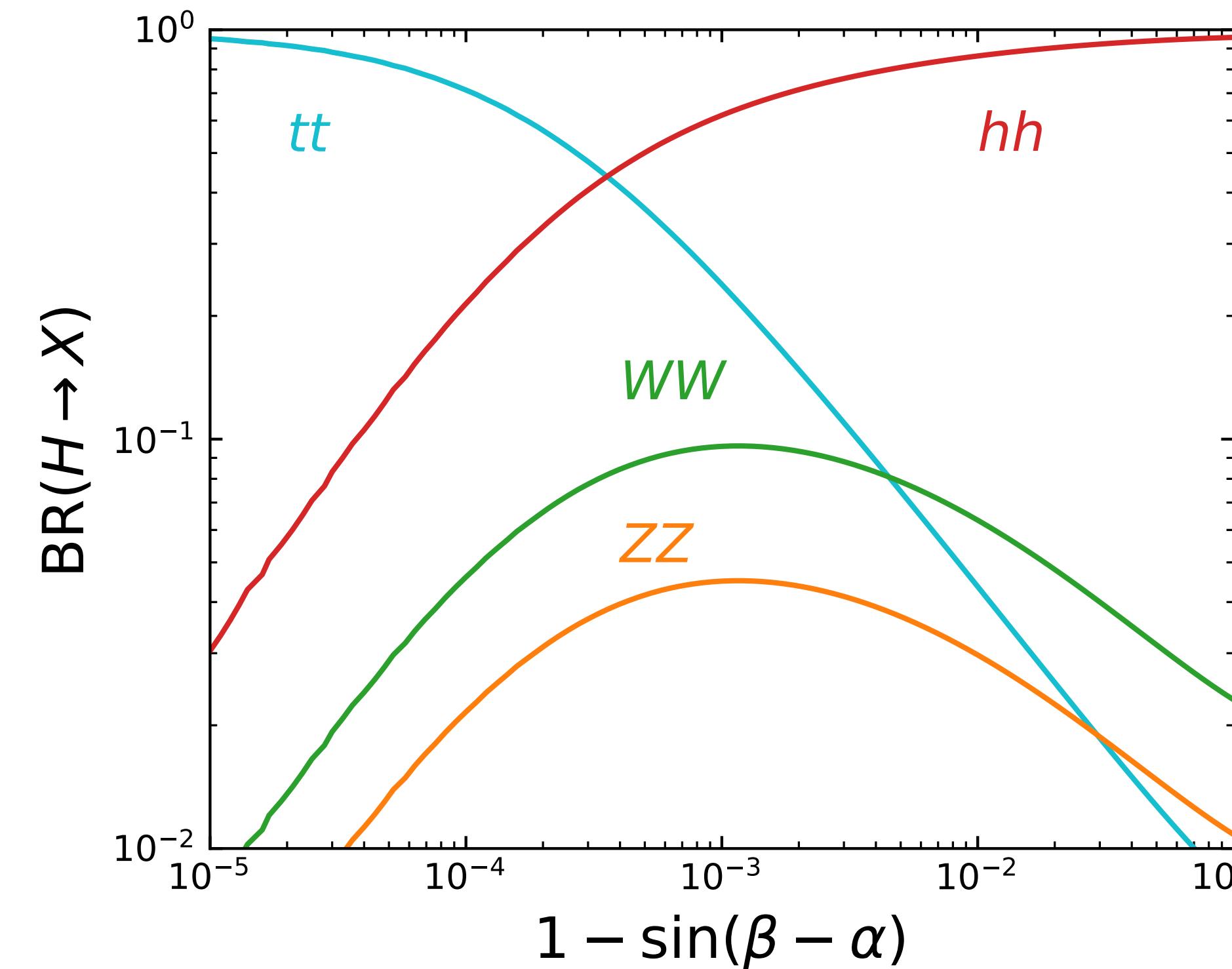
- **Alignment limit** :  $s_{\beta-\alpha} \rightarrow 1$  (tree-level Higgs couplings take SM-values.)
- LHC data indicate  $s_{\beta-\alpha} \simeq 1$ . Aad et al. PRD101 (2020)

# Decay of the additional Higgs bosons (LO)

4

## Decay patterns

Type-I 2HDM:  $m_\Phi = m_H = m_A = m_{H^\pm} = 400$  GeV,  $\tan\beta = 10$



Higgs-to-Higgs decays can be dominant if  $s_{\beta-\alpha} \neq 1$

# Synergy between direct and indirect searches 5

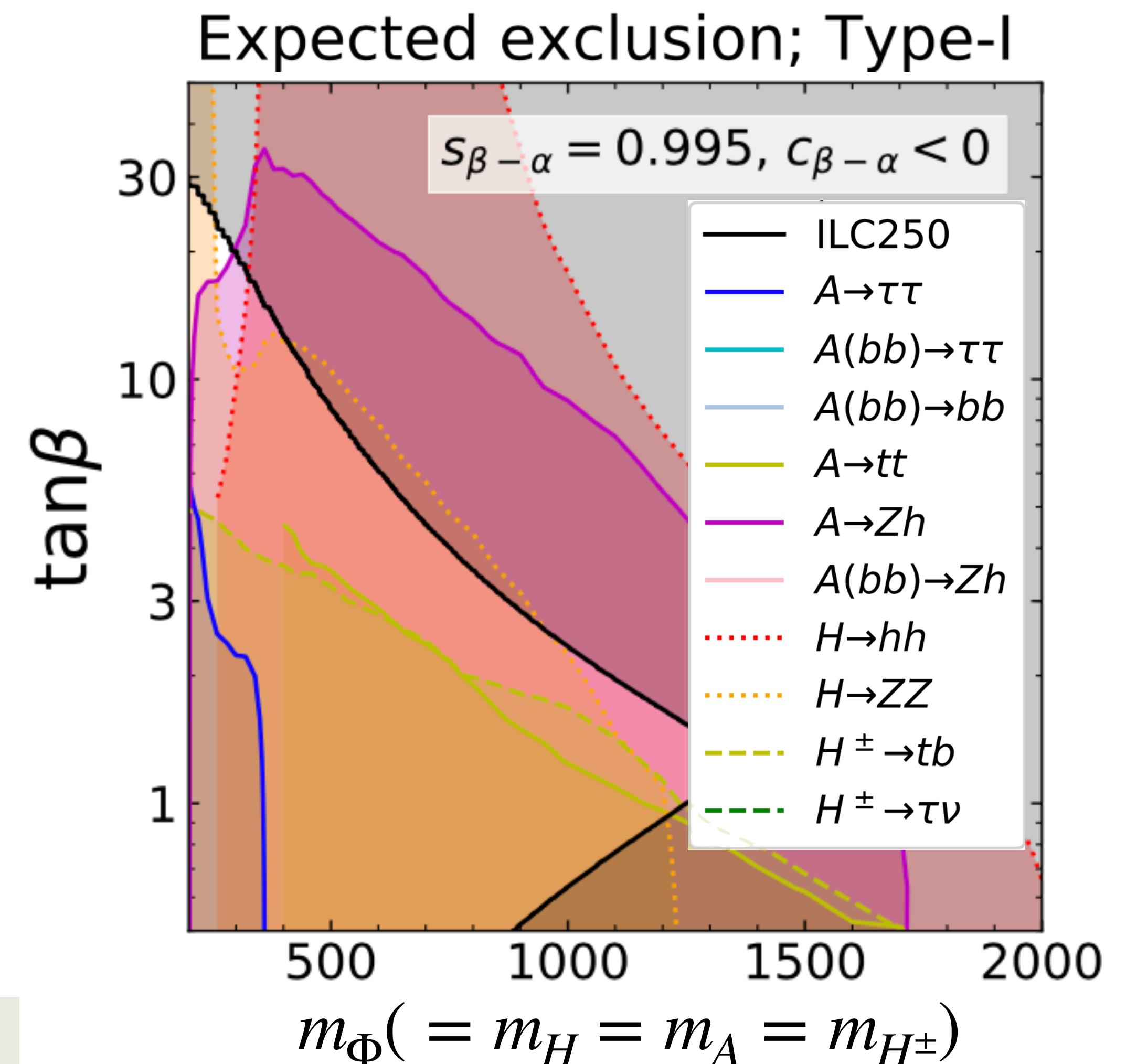
## Direct search

- HL-LHC gives lower bounds on  $m_\Phi$ .
- $H \rightarrow hh, A \rightarrow Zh$  decays exclude wide parameter region

## Indirect search

- If the SM-like Higgs boson couplings deviate from those in the SM, an upper bound on  $m_\Phi$  can be deduced.

We can explore the non-alignment scenario.



# Need of higher-order calculations

Higher-order corrections would sizably modify tree-level analysis especially in  $s_{\beta-\alpha} \approx 1$ .

## SM-like Higgs boson

Kanemura, Kikuchi, Mawatari, Sakurai, Yagyu  
NPB949 (2019), CPC257 (2020)

c.f. other public tools  
2HDECAY; Krause, Mühlleitner, Spira, CPC246 (2020)  
Prophecy4f; Denner, Dittmaier, Mück, CPC254 (2020)

- Expected accuracies in future experiments are  $\mathcal{O}(1)\%$ 
  - **H-COUP ver.2** : Two and Three-body decays of  $h(125)$  with NLO EW and QCD

## Additional Higgs bosons

## H-COUP ver.3

- Higgs-to-Higgs decays ( $H \rightarrow hh, A \rightarrow Zh, H^\pm \rightarrow W^\pm h$ ) are quite sensitive to  $c_{\beta-\alpha}$ .
  - Decays of additional CP-even Higgs boson  $H$  with NLO EW.

Krause, Mühlleitner JHEP04 (2020)  
Kanemura, Kikuchi, Yagyu 2203.08337

We discuss the impact of radiative corrections on **decays of  $A$  and  $H^\pm$** .

# Higher-order calculation

## Decay modes of $A$

MA, Kanemura, Sakurai, 2207.01032

$A \rightarrow f\bar{f}$  ( $f = q, \ell$ ),  $A \rightarrow Zh/H$   $W^\pm H^\mp$ ,  $A \rightarrow W^+W^-$ ,  $ZZ$ ,  $Z\gamma$ ,  $\gamma\gamma$  and  $gg$  (loop-induced)

## Decay modes of $H^\pm$

MA, Kanemura, Sakurai, NPB973 (2021)

$H^+ \rightarrow t\bar{b}$ ,  $c\bar{s}$ ,  $\bar{\tau}\nu$     $H^+ \rightarrow W^+h/H/A$ ,    $H^+ \rightarrow W^+Z$  and  $W^+Z\gamma$  (loop-induced)

## Improved on-shell scheme

Kanemura, Kikuchi, Sakurai, Yagyu, PRD96 (2017)  
See also Krause, Lorenz, Muhlleitner, Santos, Ziesche, JHEP09 (2016)

- UV divergences are renormalized in the on-shell scheme.
- Gauge dependencies are removed by the pinch technique. [Papavassiliou, PRD50, 5958](#)
- Infrared divergences are removed by adding real photon emission.
- For  $A \rightarrow t\bar{t}$  and  $H^\pm \rightarrow tb$  decays, NLO on-shell and NNLO  $\overline{\text{MS}}$  QCD corrections are complemented. [Djouadi et al. CPC108 \(1998\)](#)

We focus on the EW corrections to  $A \rightarrow Zh$  and  $H^\pm \rightarrow W^\pm h$ .

# Impact of NLO corrections

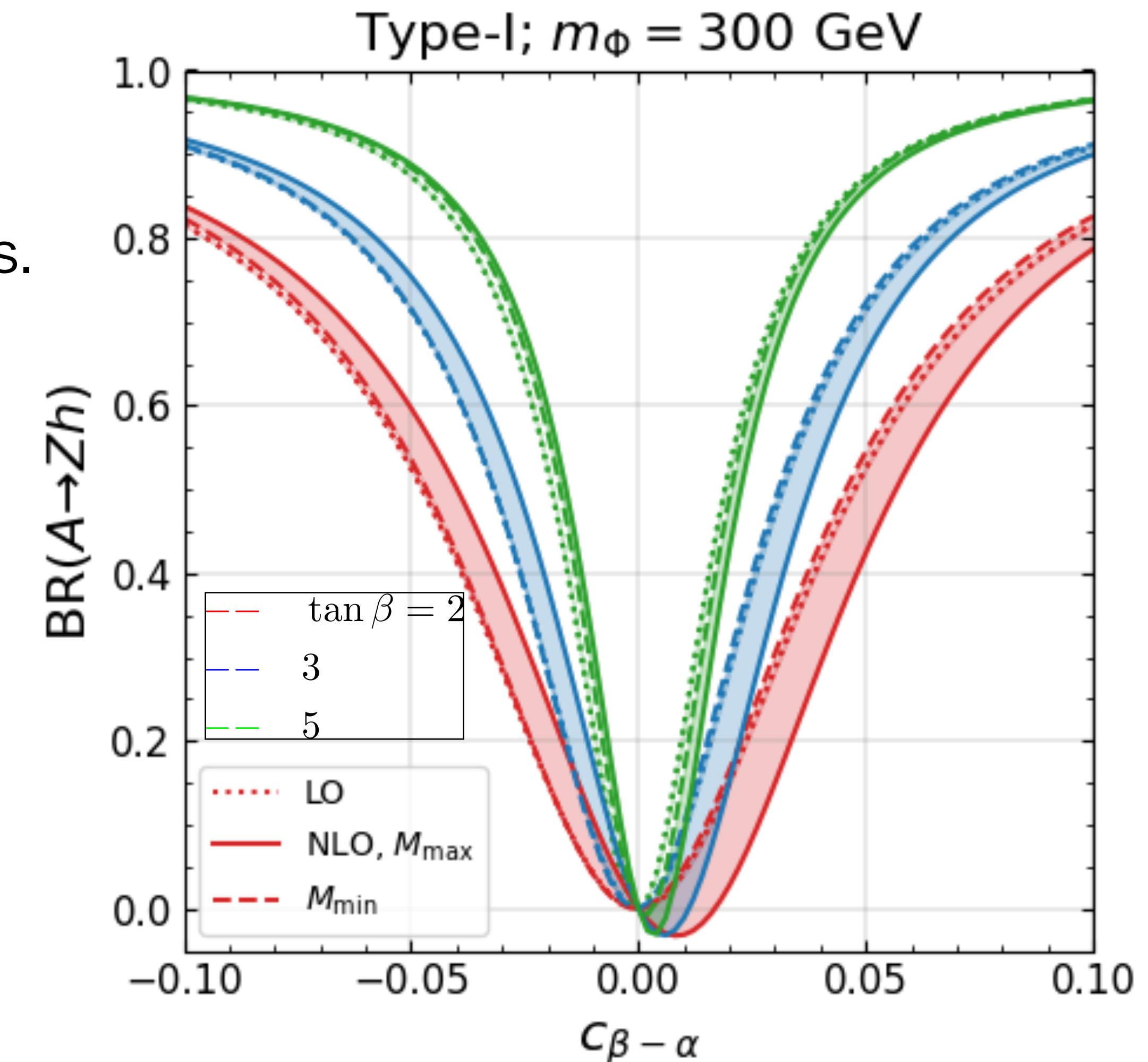
8

$A \rightarrow Zh$

- When  $\tan \beta$  is small, we have  $\mathcal{O}(10)\%$  corrections.
- When  $c_{\beta-\alpha} \approx 0$ , NLO corrections become larger than LO corrections  $\rightarrow \text{Br}(A \rightarrow Zh) < 0$
- NNLO corrections restore  $\text{Br}(A \rightarrow Zh) \geq 0$

$$|\mathcal{M}|^2 \simeq |\mathcal{M}_{\text{LO}}|^2 + 2 \operatorname{Re}(\mathcal{M}_{\text{LO}} \mathcal{M}_{\text{NLO}}^*)$$

$$|\mathcal{M}|^2 \simeq |\mathcal{M}_{\text{LO}} + \mathcal{M}_{\text{NLO}}|^2$$



# Impact of NLO corrections

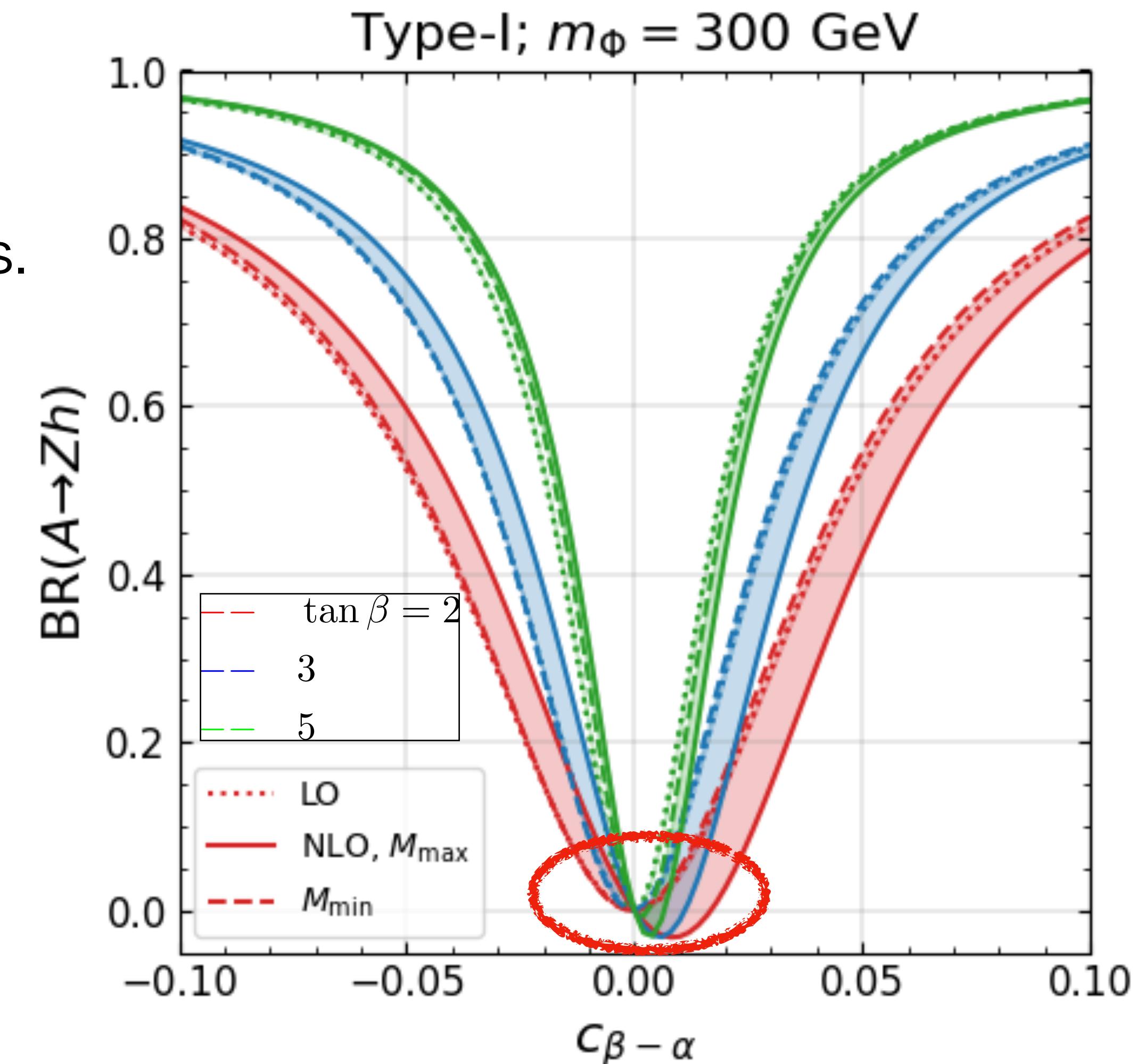
8

$A \rightarrow Zh$

- When  $\tan \beta$  is small, we have  $\mathcal{O}(10)\%$  corrections.
- When  $c_{\beta-\alpha} \simeq 0$ , NLO corrections become larger than LO corrections  $\rightarrow \text{Br}(A \rightarrow Zh) < 0$
- NNLO corrections restore  $\text{Br}(A \rightarrow Zh) \geq 0$

$$|\mathcal{M}|^2 \simeq |\mathcal{M}_{\text{LO}}|^2 + 2 \operatorname{Re}(\mathcal{M}_{\text{LO}} \mathcal{M}_{\text{NLO}}^*)$$

$$|\mathcal{M}|^2 \simeq |\mathcal{M}_{\text{LO}} + \mathcal{M}_{\text{NLO}}|^2$$



# Impact of NLO corrections

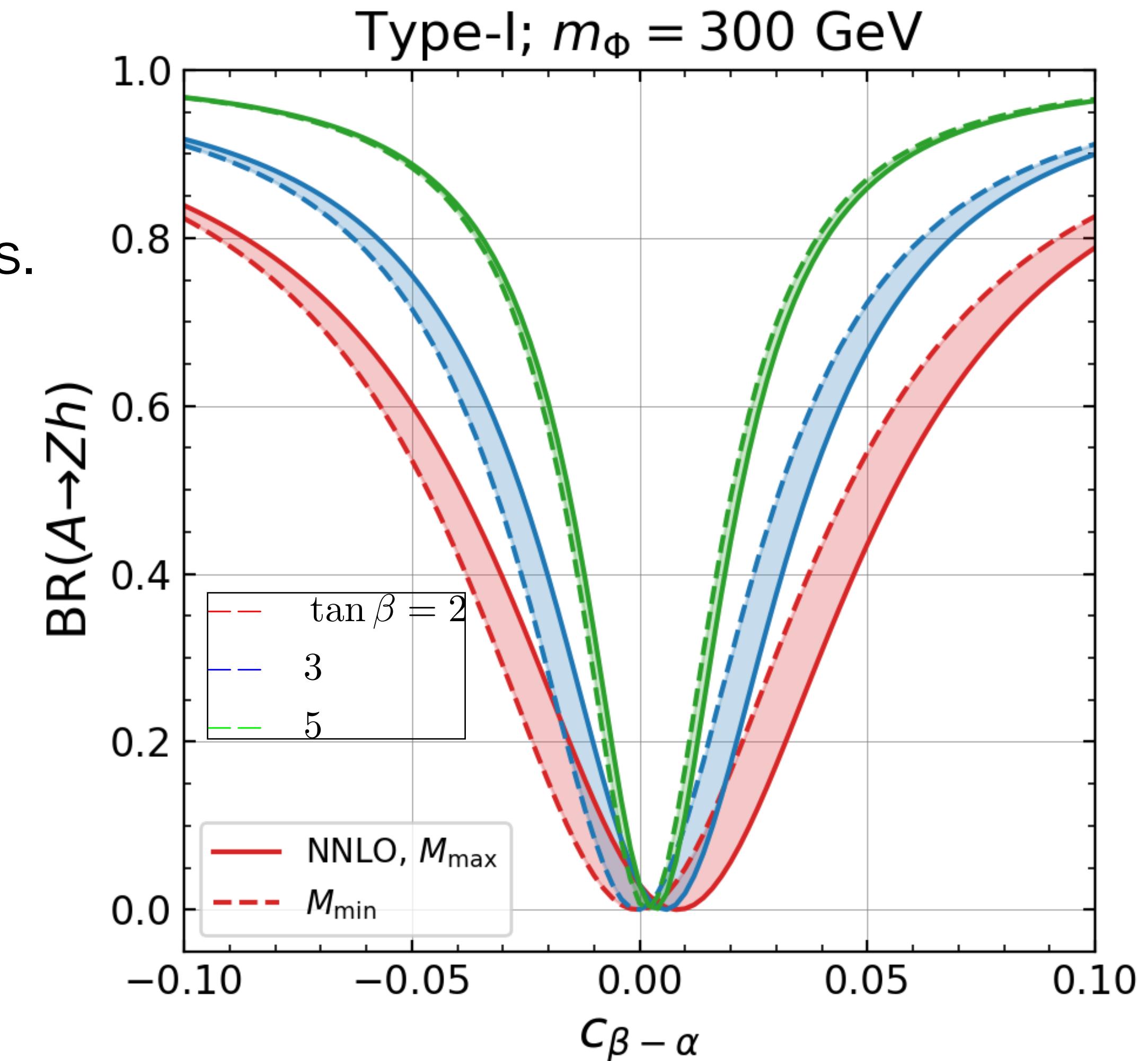
8

$A \rightarrow Zh$

- When  $\tan \beta$  is small, we have  $\mathcal{O}(10)\%$  corrections.
- When  $c_{\beta-\alpha} \simeq 0$ , NLO corrections become larger than LO corrections  $\rightarrow \text{Br}(A \rightarrow Zh) < 0$
- NNLO corrections restore  $\text{Br}(A \rightarrow Zh) \geq 0$

$$|\mathcal{M}|^2 \simeq |\mathcal{M}_{\text{LO}}|^2 + 2 \operatorname{Re}(\mathcal{M}_{\text{LO}} \mathcal{M}_{\text{NLO}}^*)$$

$$|\mathcal{M}|^2 \simeq |\mathcal{M}_{\text{LO}} + \mathcal{M}_{\text{NLO}}|^2$$

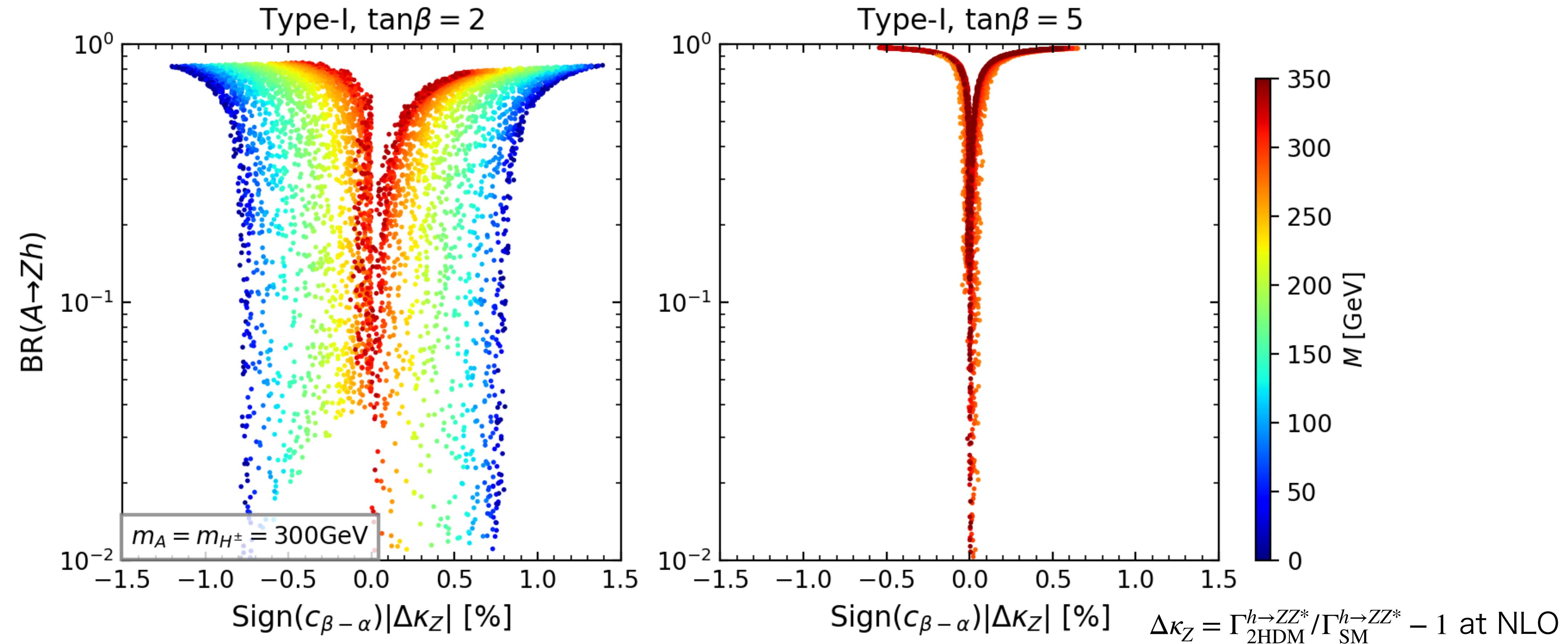


# $\text{BR}(A \rightarrow Zh) \text{ vs } \Delta\kappa_Z \text{ at NLO}$

9

$A \rightarrow Zh$

$\Gamma(A \rightarrow Zh) \propto c_{\beta-\alpha}^2, \Delta\kappa_Z \propto c_{\beta-\alpha}^2$  at LO,  $\lambda_{\phi\phi'\phi''} \propto (m_\Phi^2 - M^2)/v^2$  MA, Kanemura, Sakurai, 2207.01032



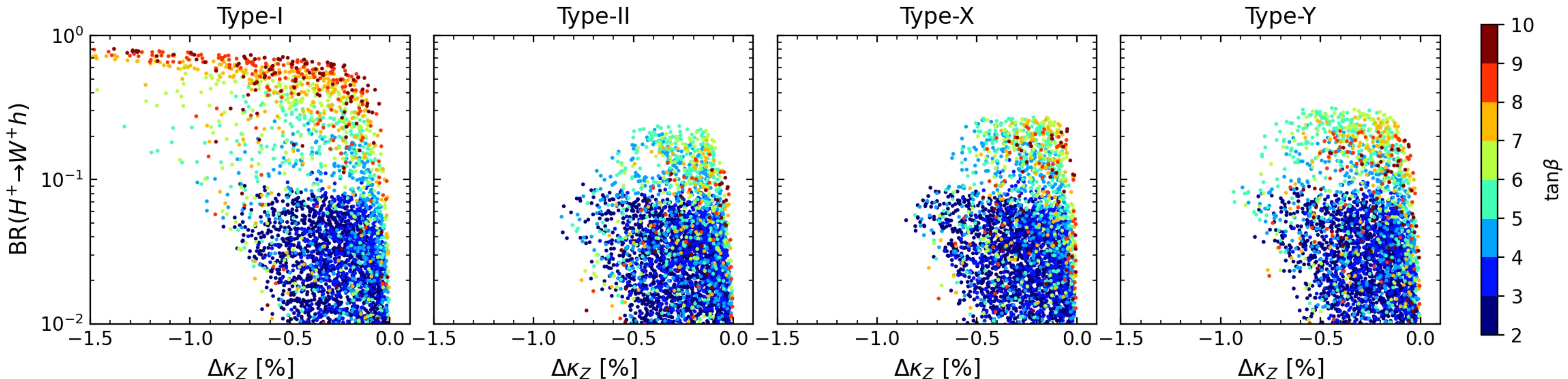
$\text{BR}(A \rightarrow Zh)$  can be small while  $\Delta\kappa_Z$  is sizable to detect.

# $\text{BR}(H^+ \rightarrow W^+ h)$ vs $\Delta\kappa_Z$ at NLO

10

- $B \rightarrow X_s\gamma$  excludes  $m_{H^\pm} \lesssim 800$  GeV in Type-II and Type-Y Misiak, Steinhauser, JHEP06 (2020)  
 → The additional Higgs bosons almost decouple ( $s_{\beta-\alpha} \simeq 1$ )

$m_{H^\pm} = m_A = 1000$  GeV Other parameters are scanned



MA, Kanemura, Sakurai, NPB973 (2021)

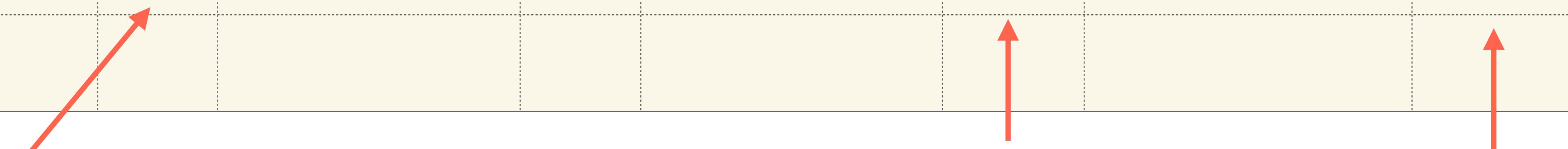
- Even if the additional Higgs bosons are heavy,  $H^\pm \rightarrow W^\pm h$  can be a dominant decay mode.

# H-COUP ver.3

11

✓ : H-COUP ver.2    ✓ : Our works    ✓ : Kanemura, Kikuchi, Yagyu 2203.08337

125GeV Higgs	CP-even	CP-odd	Charged
$h \rightarrow ff$	✓	$H \rightarrow ff$	✓
$h \rightarrow VV^*$	✓	$H \rightarrow VV$	✓
$h \rightarrow \gamma\gamma/Z\gamma/gg$	✓	$H \rightarrow hh$	✓
$e^+e^- \rightarrow hZ$	✓	$A \rightarrow VV$	✓



# Summary

## Motivation

- The phenomenology of the additional Higgs bosons is drastically changed whether  $s_{\beta-\alpha} = 1$  or not.
- NLO EW corrections would play an important role, especially if  $s_{\beta-\alpha} \simeq 1$ .

## New points

- Decays of  $A$  and  $H^\pm$  are comprehensively analyzed.
- Correlation between the decay branching ratios and  $\Delta\kappa_Z$  are exhibited.

## What we found

- Branching ratios of  $A \rightarrow Zh$  and  $H^\pm \rightarrow W^\pm h$  receive  $\mathcal{O}(10)$  % corrections if  $\tan\beta \simeq 2$ .
- Higgs-to-Higgs decay can be dominant decay modes even if  $\Delta\kappa_Z$  is small.

NLO corrections are important for direct searches of the additional Higgs bosons.

# Backup

# Synergy between direct and indirect searches

## Direct search

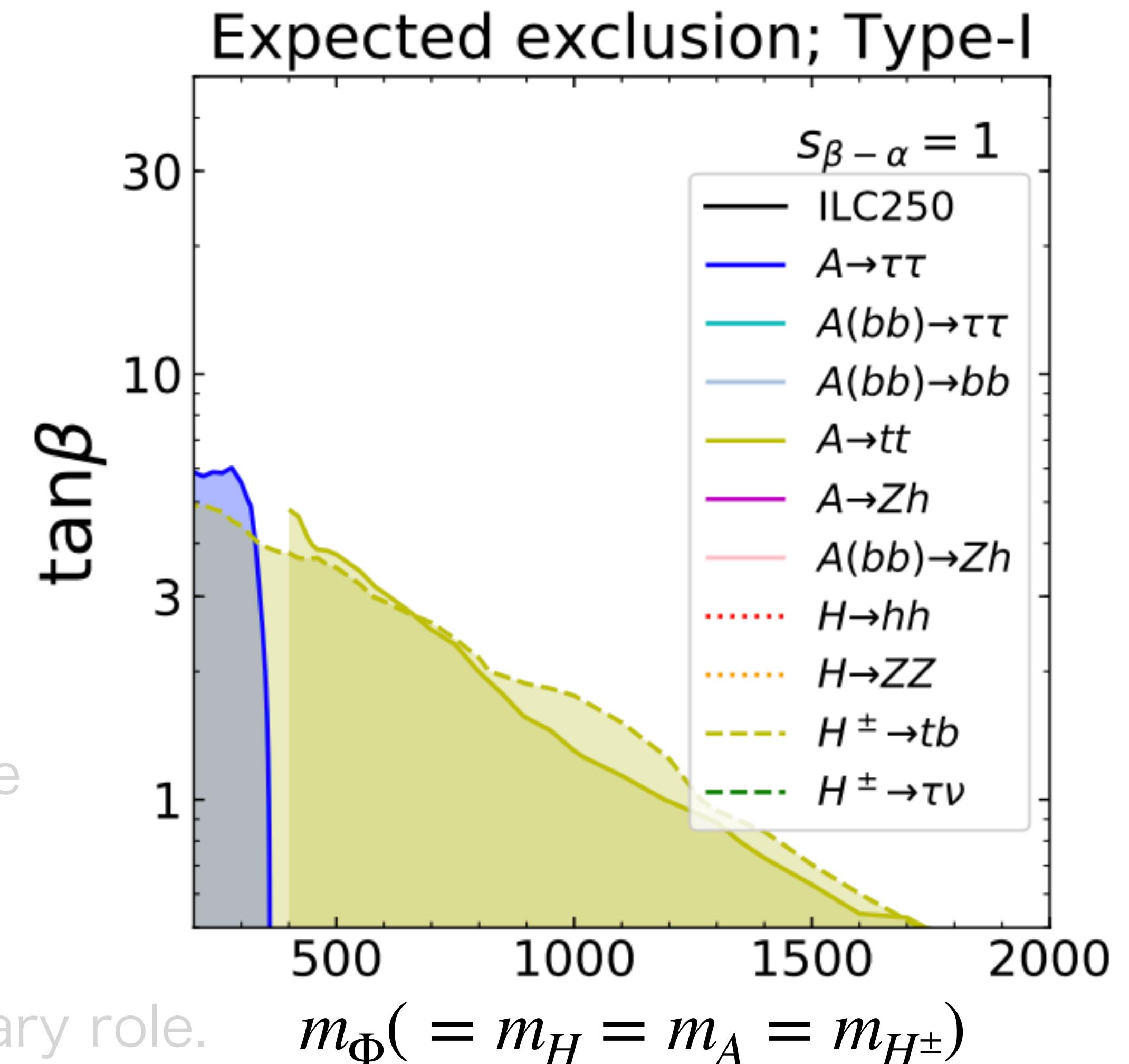
- HL-LHC gives lower bounds on  $m_\Phi$ .
- $H \rightarrow hh, A \rightarrow Zh$  decays exclude wide parameter region

## Indirect search

- If the Higgs boson couplings deviate from those in the SM, an upper bound on  $m_\Phi$  can be deduced.

Direct and indirect searches play a complementary role.

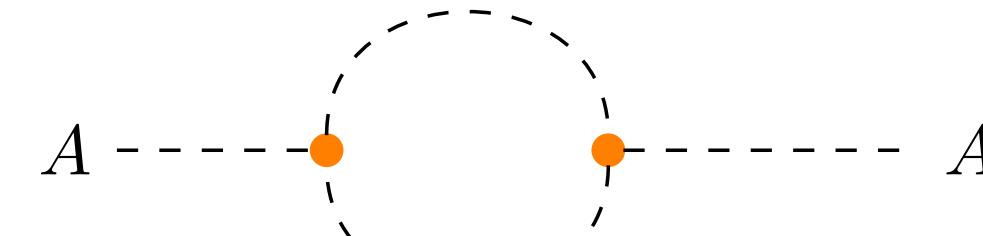
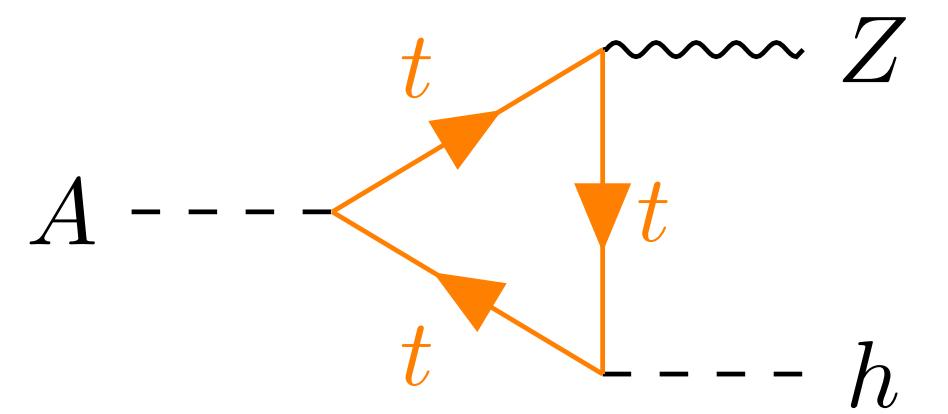
Non-alignment scenarios can be explored in detail



# Behavior of NLO Electroweak correction

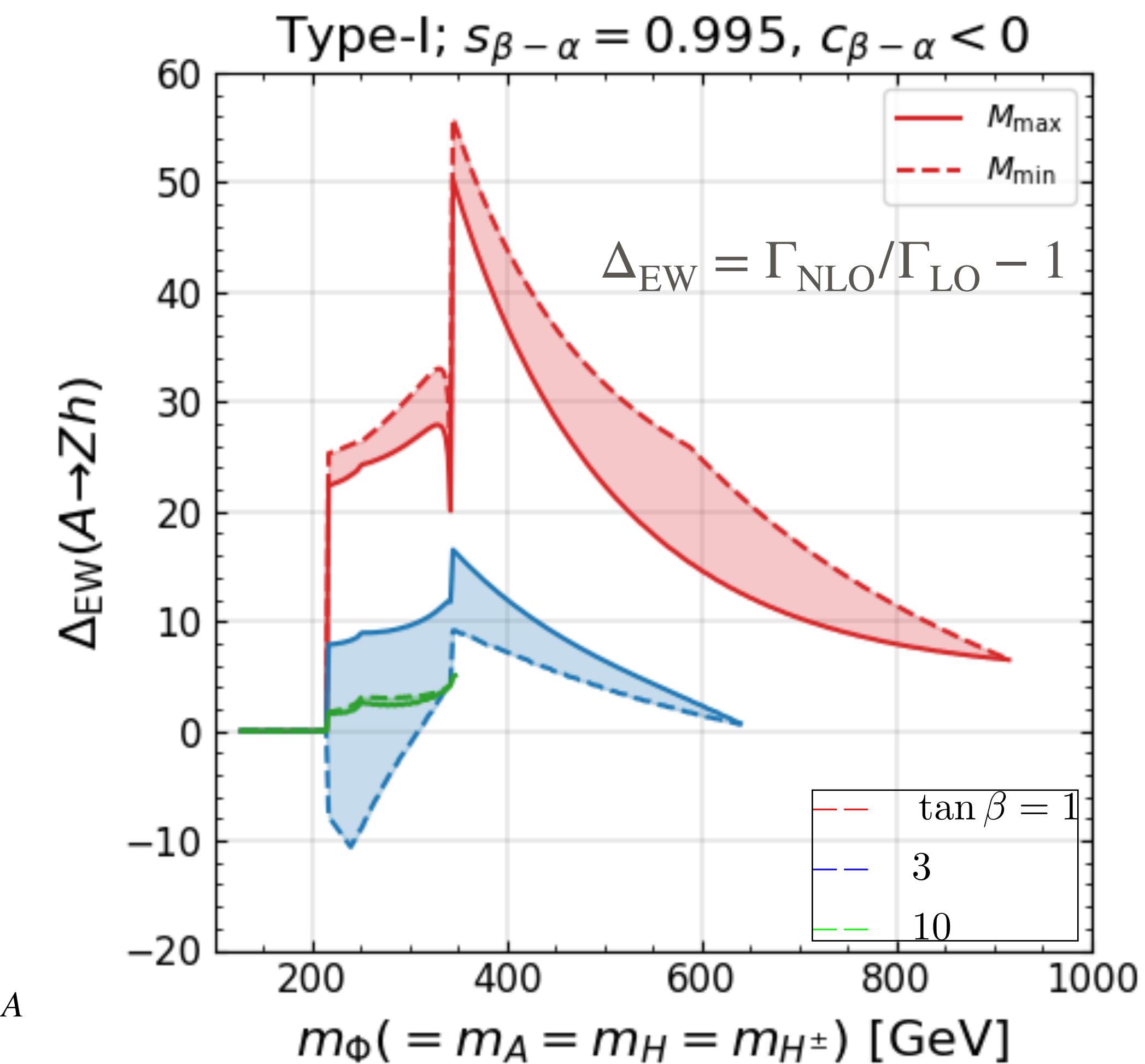
$A \rightarrow Zh$

- LO partial decay width is proportional to  $c_{\beta-\alpha}^2$ .
- The top-quark triangle diagram gives large threshold correction at  $m_A \simeq 2m_t$
- When  $m_\Phi$  is large,  $\delta Z_A$  and  $\delta\beta$  gives  $\mathcal{O}(\lambda_{SS'S''}^2)$  corrections, and they give dominant effects.



$\delta Z_A$ : Wave-function renormalization constant of  $A$

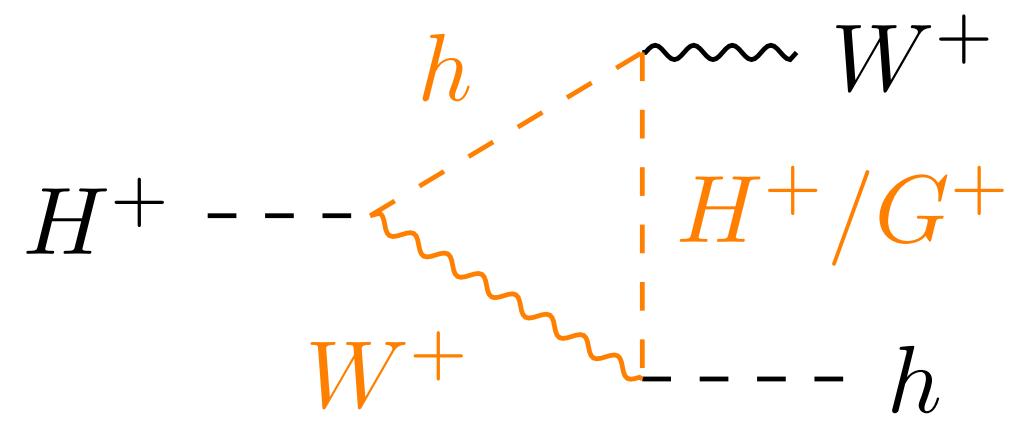
$\delta\beta$ : Counter-term of the mixing angle



# Behavior of NLO Electroweak correction

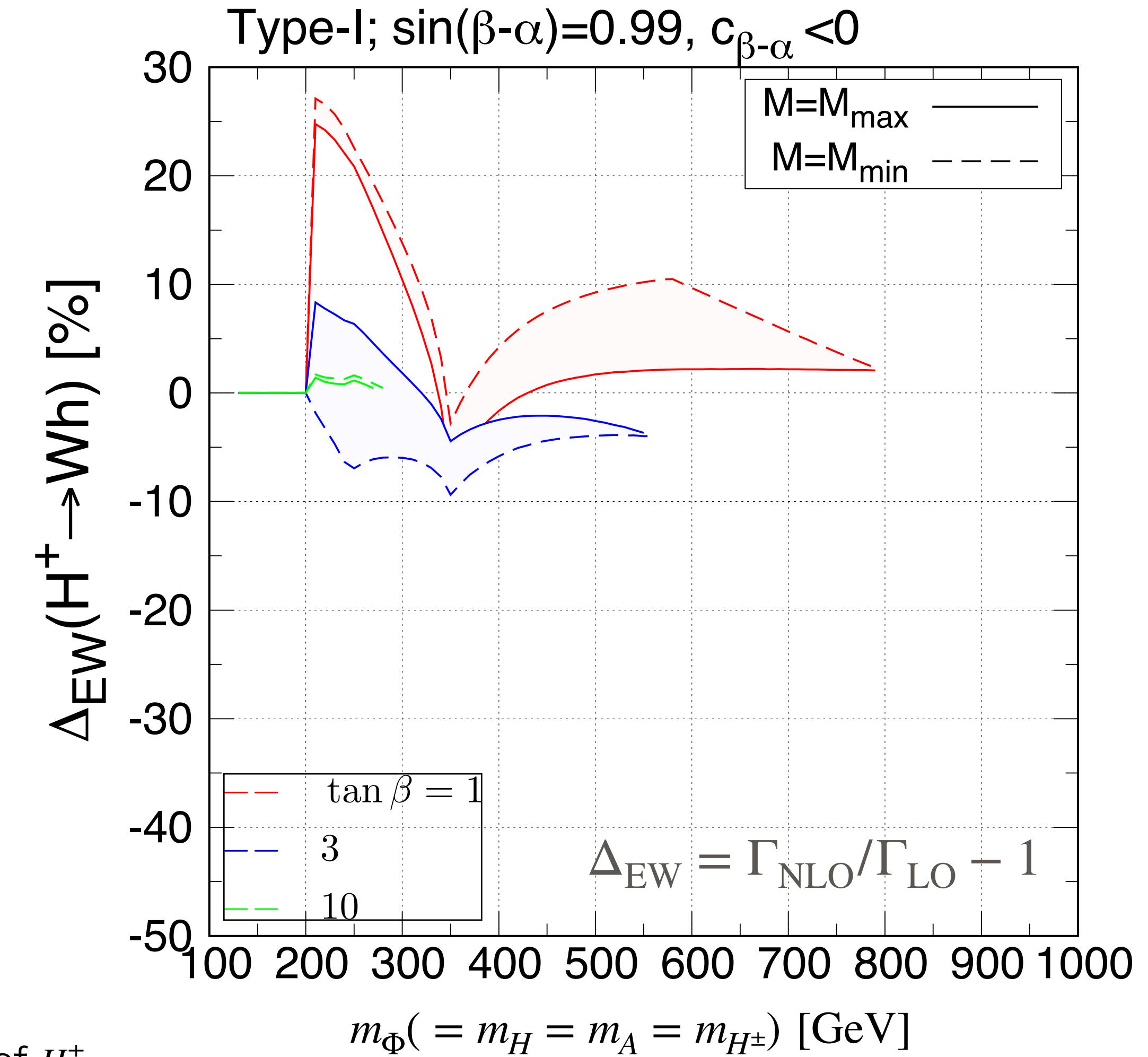
$H^+ \rightarrow W^+ h$

- LO partial decay width is proportional to  $c_{\beta-\alpha}^2$ .
- Triangle diagrams give large threshold corrections at  $m_{H^\pm} \simeq m_h + m_W$  ( $\sim 30\%$ ).
- When  $m_\Phi$  is large,  $\delta Z_{H^\pm}$  and  $\delta\beta$  gives  $\mathcal{O}(\lambda_{SS'S''}^2)$  corrections, and they give dominant effects.



$\delta Z_{H^\pm}$ : Wave-function renormalization constant of  $H^\pm$

$\delta\beta$ : Counter-term of the mixing angle



MA, Kanemura, Sakurai, NPB973 (2021) 115518

# Behavior of NLO Electroweak correction

$H^+ \rightarrow W^+ h$

- LO amplitude is proportional to  $c_{\beta-\alpha}$ , and this process becomes loop-induced if  $s_{\beta-\alpha} = 1$ .
- In the alignment limit, we have no infrared divergence.
- Decay branching ratio reaches 0.1 % if  $M^2$  is relatively small.

