Ending inflation with a bang: Higgs vacuum metastability in $R + R^2$ gravity

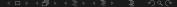
Andreas Mantziris

with A. Rajantie and T. Markkanen (arxiv:2011.03763 and 2207.00696)

Imperial College London

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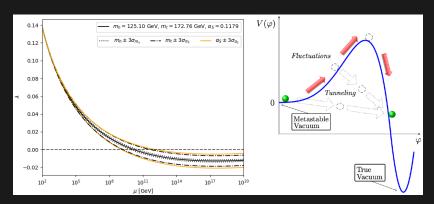


The EW vacuum metastability

Experimental values of SM particle masses m_h, m_t indicate that:

ullet currently in metastable EW vacuum o constrain fundamental physics.

$$V_{\rm H}(h,\mu,R) = \frac{\xi(\mu)}{2}Rh^2 + \frac{\lambda(\mu)}{4}h^4$$



Bubble nucleation from vacuum decay

 \bullet Decay expands at c with singularity within \to true vacuum bubbles:

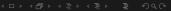
$$d\langle \mathcal{N} \rangle = \mathbf{\Gamma} d\mathcal{V} \Rightarrow \langle \mathcal{N} \rangle = \int_{\text{past}} d^4x \sqrt{-g} \mathbf{\Gamma}(\mathbf{x})$$

ullet Universe still in metastable vacuum o no bubbles in past light-cone:

$$\langle \mathcal{N} \rangle \lesssim 1$$

Vacuum bubbles expectation value (during inflation)

$$\langle \mathcal{N} \rangle = \frac{4\pi}{3} \int_{0}^{N_{\text{start}}} dN \left(\frac{a_{\text{inf}} \left(\eta_{0} - \eta \left(N \right) \right)}{e^{N}} \right)^{3} \frac{\Gamma(N)}{H(N)} \leq 1$$



Overview of computation

- Calculate $\Delta V_{
 m H}$ and plug it in $\Gamma pprox \left(rac{R}{12}
 ight)^2 e^{-rac{384\pi^2\Delta V_{
 m H}}{R^2}}$.
- 2 Cosmological quantities according to the inflationary model $V_{\rm I}(\tilde{\phi})$; for Starobinsky/ R^2 inflation $V_{\rm I}(\phi)=\frac{3M^2M_P^4}{4}\left(1-e^{-\sqrt{\frac{2}{3}}\frac{\phi}{M_P}}\right)^2$.
- 3 Complete calculation of $\langle \mathcal{N} \rangle$ imposing the condition $\langle \mathcal{N} \rangle \leq 1$.

$$\left\langle \mathcal{N} \right\rangle = rac{4\pi}{3} \int_{0}^{N_{\mathrm{start}}} dN \left(rac{a_{\mathrm{inf}} \left(\eta_{0} - \eta \left(N
ight)
ight)}{e^{N}}
ight)^{3} rac{\Gamma(N)}{H(N)} \leq 1$$

4 Result: constraints on $\xi \geq \xi_{\langle \mathcal{N} \rangle = 1}$ and cosmological implications from the time of predominant bubble nucleation.

Higgs potential in curved spacetime

Minkowski terms to 3-loops, curvature corrections in dS at 1-loop:

$$V_{\rm H}(h,\mu,R) = \frac{\xi}{2}Rh^2 + \frac{\lambda}{4}h^4 + \frac{\alpha}{144}R^2 + \Delta V_{\rm loops},$$

where the loop contribution can be parametrized as

$$\Delta V_{\text{loops}} = \frac{1}{64\pi^2} \sum_{i=1}^{31} \left\{ n_i \mathcal{M}_i^4 \left[\log \left(\frac{|\mathcal{M}_i^2|}{\mu^2} \right) - d_i \right] + \frac{n_i' R^2}{144} \log \left(\frac{|\mathcal{M}_i^2|}{\mu^2} \right) \right\}$$

RGI: choose $\mu = \mu_*(h,R)$ such that $\Delta V_{\text{loops}}(h,\mu_*,R) = 0 \rightarrow$

RGI effective Higgs potential

$$V_{\rm H}^{\rm RGI}(h,R) = \frac{\xi(\mu_*(h,R))}{2}Rh^2 + \frac{\lambda(\mu_*(h,R))}{4}h^4 + \frac{\alpha(\mu_*(h,R))}{144}R^2$$

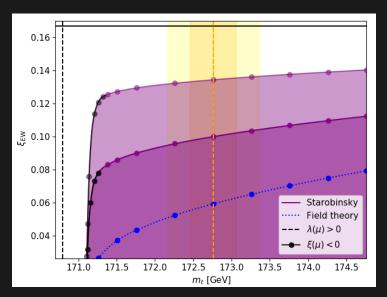
Markkanen et al, "The 1-loop effective potential for the Standard Model in curved spacetime", 2018.

RGI effective Higgs potential in $R+R^2$ gravity

$$\begin{split} S &= \int d^4x \sqrt{-g_J} \left[\frac{M_P^2}{2} \left(1 - \frac{\xi h^2}{M_P^2} \right) R_J + \frac{1}{12M^2} R_J^2 + \frac{1}{2} g_J^{\mu\nu} \partial_\mu h \partial_\nu h - \frac{\lambda}{4} h^4 \right] \\ &\Rightarrow \ldots \Rightarrow \mathcal{L} \approx \frac{M_P^2}{2} R + \frac{1}{2} \partial_\mu \tilde{\phi} \partial^\mu \tilde{\phi} + \frac{1}{2} \partial_\mu \rho \partial^\mu \rho - \tilde{U}(\tilde{\phi}, \rho) \,, \\ \tilde{U}(\tilde{\phi}, \rho) &= V_I(\tilde{\phi}) + m_{\text{eff}}^2(\tilde{\phi}, \mu_*) \frac{\rho^2}{2} + \lambda_{\text{eff}}(\tilde{\phi}, \mu_*) \frac{\rho^4}{4} + \frac{\alpha(\mu_*)}{144} R^2(\tilde{\phi}) + \mathcal{O}(\frac{\rho^6}{M_P^2}), \\ \text{where } \Xi(\mu_*) &= \xi(\mu_*) - \frac{1}{6} \text{ and} \end{split}$$

$$\begin{split} V_{\rm I}(\tilde{\phi}) &= \frac{3M^2 M_P^4}{4} \left(1 - e^{-\sqrt{\frac{2}{3}} \frac{\tilde{\phi}}{M_P}} \right)^2 \,, \\ m_{\rm eff}^2 &= \xi R + 3M^2 M_P^2 \Xi \left(1 - e^{-\sqrt{\frac{2}{3}} \frac{\tilde{\phi}}{M_P}} \right) e^{-\sqrt{\frac{2}{3}} \frac{\tilde{\phi}}{M_P}} + \frac{\Xi}{M_P^2} \partial_\mu \tilde{\phi} \partial^\mu \tilde{\phi} \,, \\ \lambda_{\rm eff} &= \lambda + 3M^2 \Xi^2 e^{-2\sqrt{\frac{2}{3}} \frac{\tilde{\phi}}{M_P}} + \frac{4 \left[\xi R + \Delta m_1^2 \right] \Xi^2}{M_P^2} + \frac{4\Xi^3}{M_P^4} \partial_\mu \tilde{\phi} \partial^\mu \tilde{\phi} \,. \end{split}$$

Results: Lower ξ -bounds for varying top quark mass



Conclusions

• Minimal model of the early universe: SM + Starobinsky inflation (observationally favoured) from modification of gravity $R+R^2$.

• Vacuum decay constraints on the Higgs-curvature coupling, with state-of-the-art $V_{\rm eff}^{\rm RGI}$ (3-loop couplings, 1-loop dS corrections):

$$\xi_{\rm EW} \gtrsim 0.1 > 0.06$$
,

give stricter ξ -bounds from extra negative terms in $V_{
m H}^{
m RGI}$.

- Bubble nucleation in the last moments of inflation: breakdown of dS approximations and necessity to consider the dynamics of reheating.
- Possibly hints against eternal inflation (again)