# Dark sector freeze-out due to a non-Boltzmann suppression 

## Sougata Ganguly

Indian Association for the Cultivation of Science

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## Thermal dark matter: Idea of WIMP

## Assumptions:

- DM is in thermal equilibrium with the SM particles in the early Universe.
, DM interacts with the SM particles with weak scale couplings.


WIMP annihilation cross section considering $\mathrm{m}_{\mathrm{DM}}=1 \mathrm{TeV}$
$\langle\sigma v\rangle_{\mathrm{DMDM} \rightarrow \mathrm{SMSM}} \simeq \frac{g_{2}^{4}}{32 \pi m_{\mathrm{DM}}^{2}}=1.46 \times 10^{-9} \mathrm{GeV}^{-2}$

DM of mass 1 TeV satisfies the relic density constraint if it couples with the SM particles with weak scale coupling

## WIMP Searches

, Spin-independent DM-nucleon cross-section put stringent constraint on the DM mass and its couplings with the SM particles.

- In colliders, produce WIMP in association with the SM particles and calculate the missing energy to predict the mass of the WIMP.


Need to go beyond WIMP

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Classifications of particle dark matter: A schematic picture


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## Secluded Dark Sector



- Constraints from direct searches can be evaded easily
- Determines the thermal properties of the dark sector
- $\Phi$ should decay before BBN
$\rightarrow$ Can be multi-component.
$\rightarrow$ Can be degenerate.
$\rightarrow$ Can be thermally decoupled.


## Degenerate Dark Sector

, The dark sector contains two particles A and $B$ and they are degenerate i.e. $m_{A}=$ $m_{B}=m$
, B can decay out-of-equilibrium into the SM particles.


Dark sector becomes nonrelativistic

B starts to decay

- Number density of A starts to decrease when $B$ starts to decay into the SM particles.

A freezes-out

Co-Decaying Dark Matter

Final abundance of $\mathbf{A}$ depends on $\langle\sigma v\rangle_{A A \rightarrow B B}$ as well as $\Gamma_{B}$

## Dynamics of a multicomponent degenerate dark sector

- We need a DM candidate.
- We need a massive mediator particle, which can decay into SM particles.

The dark sector contains

- Two Majorana fermions $\chi_{1}$ and $\chi_{2}$

Two mediator particles $Z^{\prime}$ and $h_{d}$

A natural choice $\mathrm{U}(1)_{\mathrm{x}}$ extension of SM .

- Contains $\xi_{L}$ and $\xi_{2 L}$ with $U(1)_{x}$ charges +1 and -1 respectively $\longrightarrow$ Anomaly free
- A gauge boson $Z^{\prime}$ corresponding to the $U(1)_{X}$ gauge symmetry and it kinetically mixes with SM Z boson.
- A complex scalar $\eta$ to break $U(1)_{x}$ gauge symmetry $\longrightarrow$ Massive $Z^{\prime}$
- Two component dark matter
- $\gamma$ and $\nu$ signals from DM annihilation via one step cascade processes


## Lagrangian

$$
\mathcal{L}=\mathcal{L}_{\mathrm{SM}}+\mathcal{L}_{\text {gauge }}+\mathcal{L}_{\mathrm{DM}-\text { gauge }}+\mathcal{L}_{\mathrm{DM}-\text { Yukawa }}+\mathcal{L}_{\text {scalar }},
$$

where

$$
\begin{aligned}
\mathcal{L}_{\text {gauge }} & =-\frac{1}{4} X_{\mu \nu} X^{\mu \nu}-\frac{\epsilon}{2} B_{\mu \nu} X^{\mu \nu}, \\
\mathcal{L}_{\text {DM-gauge }} & =i \overline{\xi_{1 L}} D \xi_{1 L}+i \overline{\xi_{2 L}} D \xi_{2 L} \\
\mathcal{L}_{\text {DM-Yukawa }} & =-\left(\frac{y}{2} \overline{\left.\xi_{1 L}^{c} \xi_{1 L} \eta+\frac{y}{2} \overline{\xi_{2 L}^{c}} \xi_{2 L} \eta^{\dagger}+h . c\right)}\right. \\
\mathcal{L}_{\text {scalar }} & =\left(D_{\mu} \eta^{\dagger}\right)\left(D^{\mu} \eta\right)+\mu_{X}^{2}\left(\eta^{\dagger} \eta\right)-\lambda_{X}\left(\eta^{\dagger} \eta\right)^{2}-\lambda^{\prime}\left(\eta^{\dagger} \eta\right)\left(\Phi^{\dagger} \Phi\right)
\end{aligned}
$$

## DM annihilation channels



## Results




## Constraints on the mediator particles




## m-gx plane



## Summary

The out-of-equilibrium decay of long-lived mediator particle such as $\mathbf{Z}$ ' and $h_{d}$ leading to delayed freeze-out of Dark Matter and this is known as "Co-Decaying Dark Matter".

The dark sector enters into the "Cannibal" phase due to the presence of $3 \rightarrow 2$ processes and during this phase the temperature evolution changes significantly.

We have investigated the allowed region of parameter space from $\nu$ and $\gamma$ ray signals from DM annihilation via one step cascade process.

The bounds from direct, indirect, laboratory and astrophysical searches can be easily evaded in case of degenerate dark sector. However for non-degenerate dark sector, a certain region of the parameter space is significantly constrained from the measurement of diffuse $\gamma$ ray flux by INTEGRAL, CMB anisotropy, and positron flux by AMS-02 experiment.

Thank You

## Back-up slides



