A last chance for kinetic mixing: explaining $(g-2)_\mu$ with semi-visible dark photons

A. Abdullahi, M. Hostert, **D. Massaro**, S. Pascoli

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ALMA MATER STUDIORUM
UNIVERSITÀ DI BOLOGNA

UCLouvain

FNAL measurement

FNAL confirmed $(g\text{-}2)_\mu$ discrepancy with SM

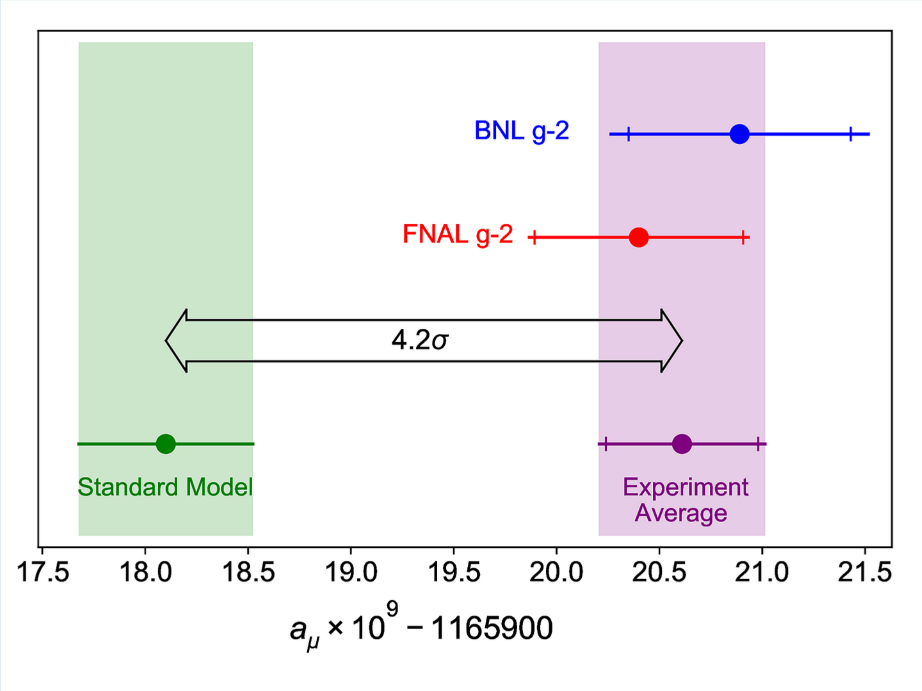
$$\Delta a_\mu = a_\mu^{\mathrm{EXP}} - a_\mu^{\mathrm{SM}} = (251 \pm 59) \times 10^{-11}$$

4.2 σ

Hints for new physics?

Not too far out of reach

$$\Delta a_\mu \propto \frac{g^2}{16\pi^2}\frac{m_\mu^2}{\Lambda^2} \Rightarrow \frac{\Lambda}{g} \approx \mathcal{O}(10^2)\ \mathrm{GeV}$$



[B. Abi et al. Phys. Rev. Lett. 126.14 (2021), p. 141801]

*Recent comparison of Lattice Results for Leading Order Hadron Vacuum Polarization would reduce the significance to 2.9 σ. [Cè et al. (2022), arXiv:2206.06582]

Dark photon as a solution

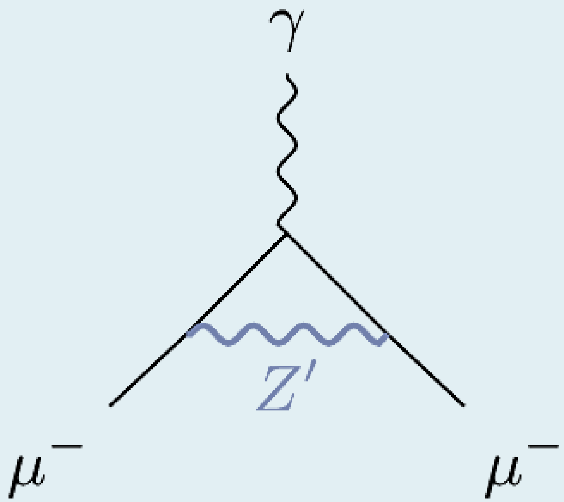
$$\mathcal{L} \supset \frac{\varepsilon}{2} B^{\mu\nu} X_{\mu\nu}$$

New U(1)' gauge group with vector boson **Z'**, kinetically mixed to $U(1)_Y$ gauge boson

Positive contribution to a_μ

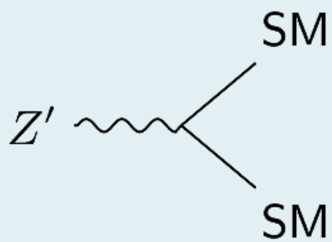
$$\Delta a_\mu = \frac{\alpha}{2\pi}\varepsilon^2 \int_0^1 \mathrm{d}z \frac{2m_\mu^2 z(1-z)^2}{m_\mu^2(1-z)^2 + m_{Z'}^2 z}$$

[M. Pospelov. Phys. Rev. D 80 (2009), p. 095002]

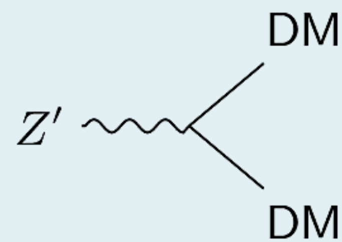


Detection: decay modes in minimal models

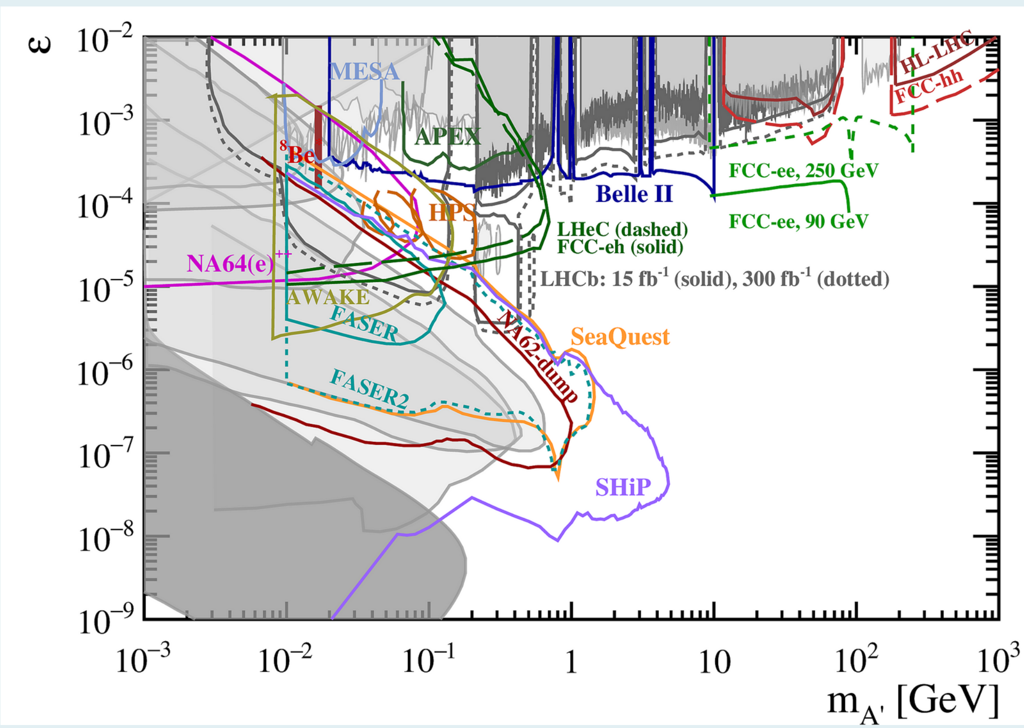
Dark photon decay modes



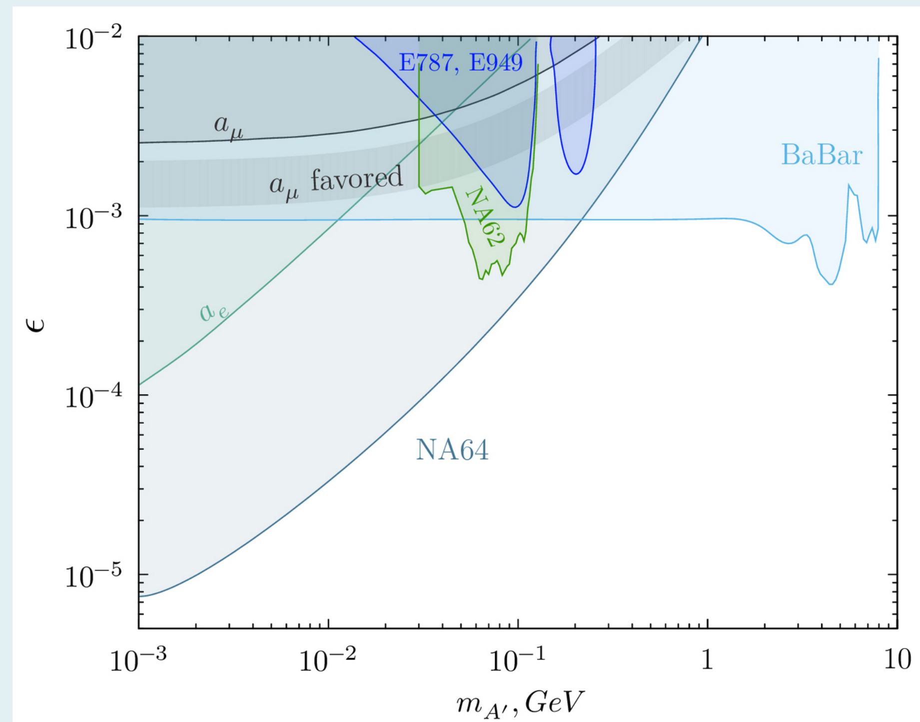
Visible



Invisible



[M. Fabbrichesi et al. (2020). arXiv: 2005.01515]

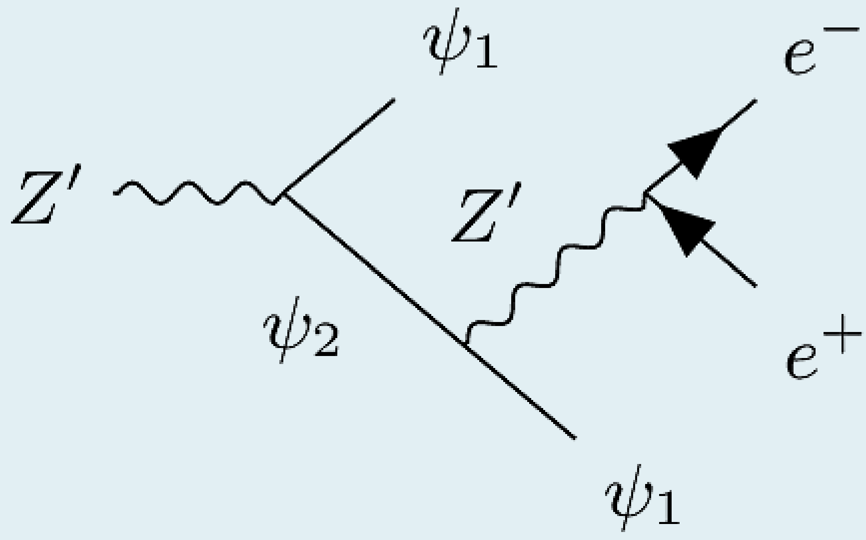


[D. Banerjee et al. Phys. Rev. Lett. 123.12 (2019), p. 121801]

Visible decays in SM final states are **excluded** by colliders in resonance searches

Invisible decays **fully excluded** by BaBar and NA64

Semi-visible to the rescue



[E. Izaguirre et al. Phys. Rev. D 96.5 (2017), p. 055007]
[G. Mohlabeng. Phys. Rev. D 99.11 (2019), p. 115001]

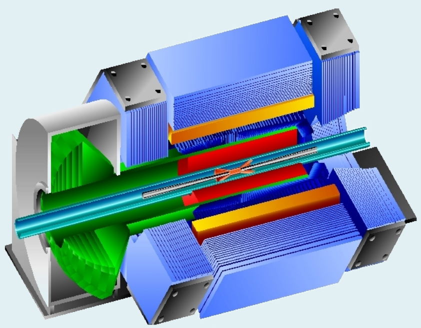
Potential solution

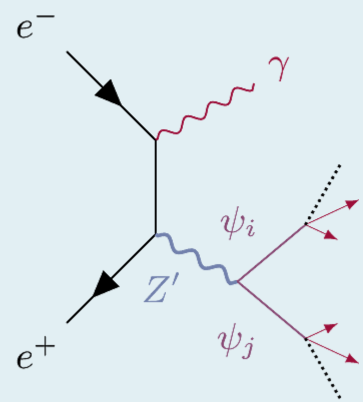
- Evade visible constraints with **missing energy**
- Evade invisible constraints with **visible final states**

Main constraints

BaBar

monophoton



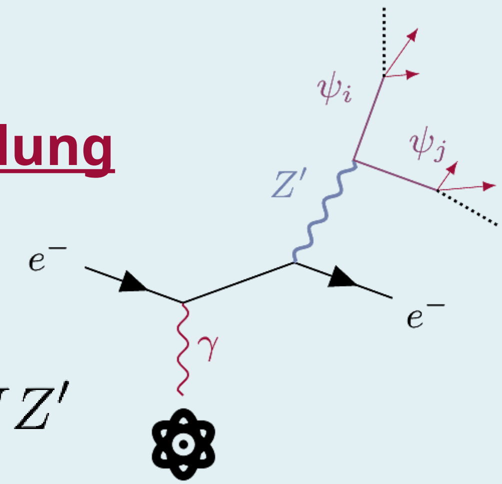


$$e^+e^- \to \gamma Z'$$

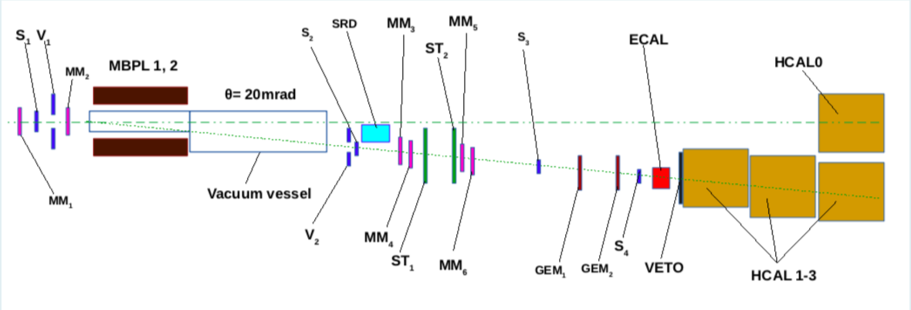
Monophoton + missing energy

NA64

Z’ bremsstrahlung



$$e^- N \to e^- N Z'$$



Missing energy events with dark photon produced in electron bremsstrahlung

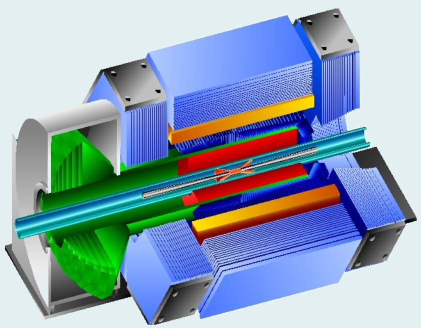
Are semi-visible Z’ viable?

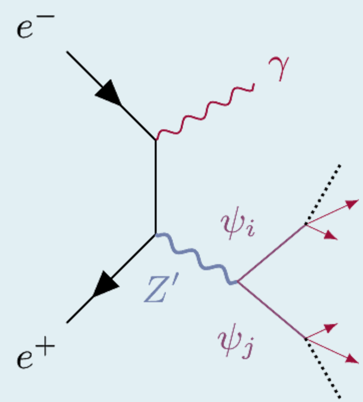
- ☑ Simulate production of semi-visible Z’ in BaBar and NA64 detector
- ☑ Apply selection criteria
- ☑ Recast the bound

Main constraints

BaBar

monophoton





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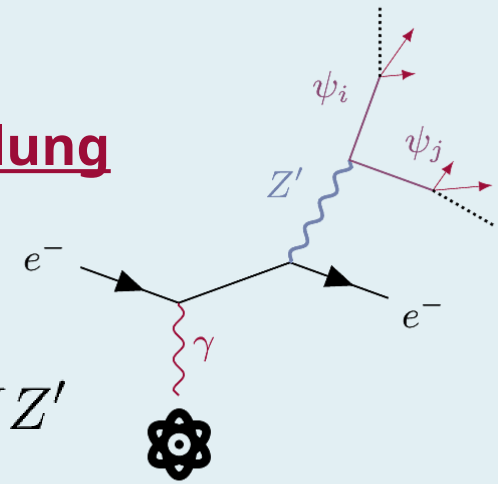
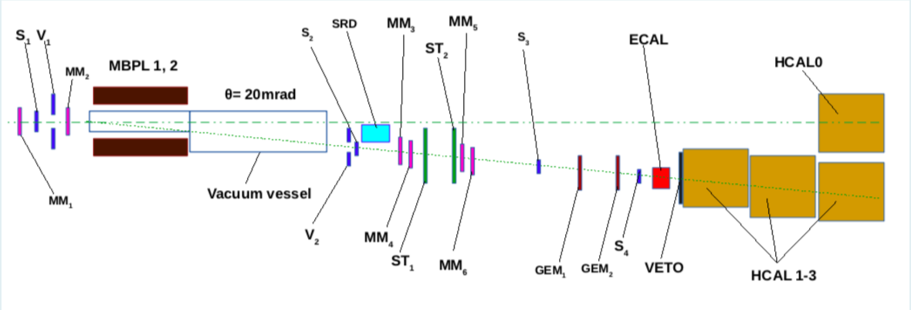


$$e^-N \to e^-NZ'$$



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- ☑ Apply selection criteria
- ☑ Recast the bound

$$\varepsilon_{\text{bound}} = \varepsilon_{\text{new}} \sqrt{P^{\text{inv}}(\varepsilon_{\text{new}}, m_{Z'})}$$

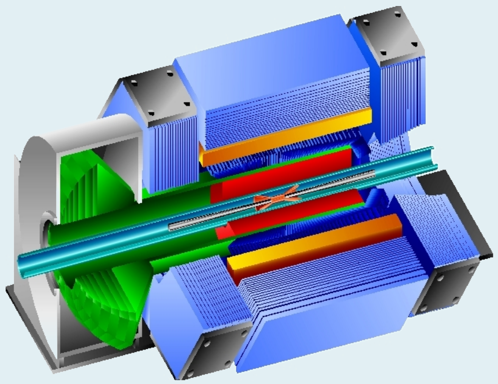
The BaBar simulation: generation

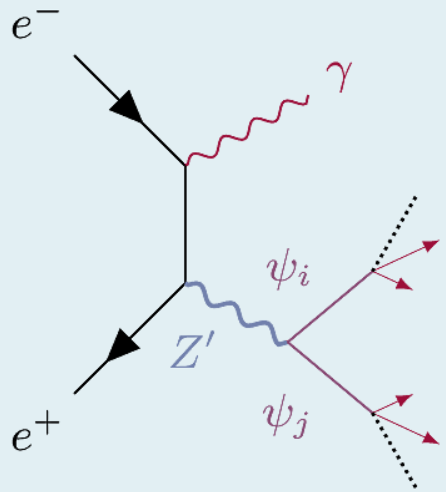
Process: $e^+e^- \to \gamma Z'$

- ☑ Generate $O(10^6)$ events for given $m_{Z'}$ and ε
- ☑ Apply primary selection cut on monophoton CM angle
- ☑ Apply veto criteria and compute new P^{inv}

Z' decays instantly to HNL, producing pairs of leptons final states:

$$Z' \to (\psi_i \to (\psi_k \to \ldots)e^+e^-)\,(\psi_j \to (\psi_l \to \ldots)e^+e^-)$$





Events containing visible final states are vetoed

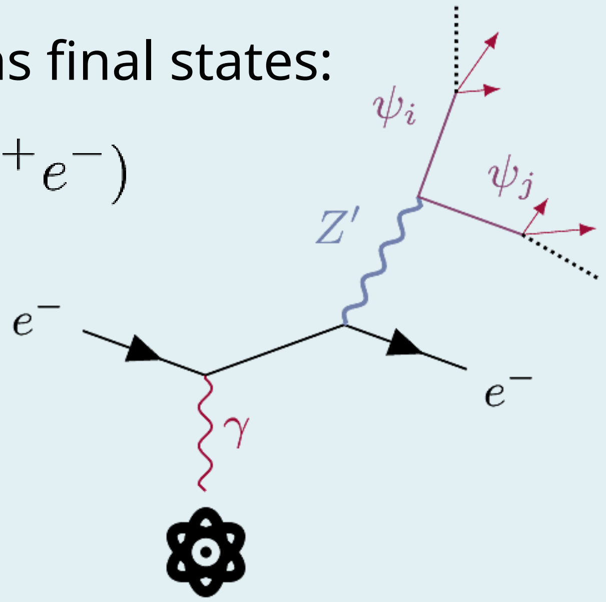
The NA64 simulation: generation

Process: $e^- N \to e^- N Z'$

- ☑ Generate O(10^5) events for given $m_{Z'}$ and ε
- ☑ Apply veto criteria and compute new P^{inv}

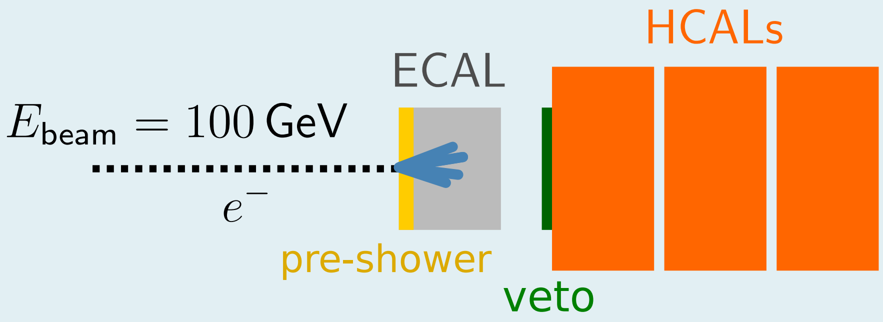
Z' decays instantly to HNL, producing pairs of leptons final states:

$$Z' \to (\psi_i \to \psi_k e^+ e^-)\,(\psi_j \to \psi_l e^+ e^-)$$



Recast: solve $\varepsilon_{\mathrm{NA64}} = \varepsilon \sqrt{P^{\mathrm{inv}}(\varepsilon, m_{Z'})}$

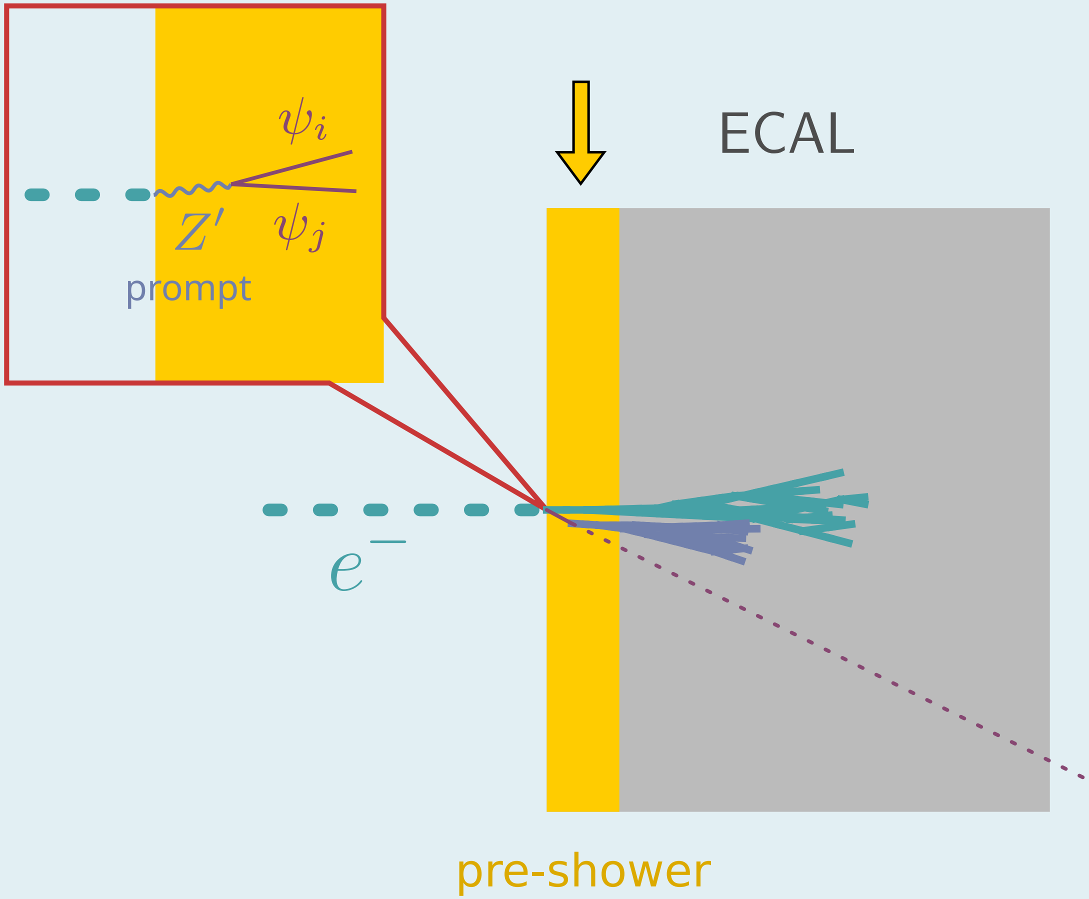
where $P^{\mathrm{inv}}(\varepsilon, m_{Z'}) = 1 - \dfrac{N_{\mathrm{veto}}(\varepsilon, m_{Z'})}{N(\varepsilon, m_{Z'})}$



e^+e^- pairs deposit energy in the calorimeters or in the veto, so new events are vetoed according cut conditions.

Recast of the "invisible" bound

[Phys. Rev. Lett. 123.12 (2019), p. 121801]



✔ Pre-shower ECAL

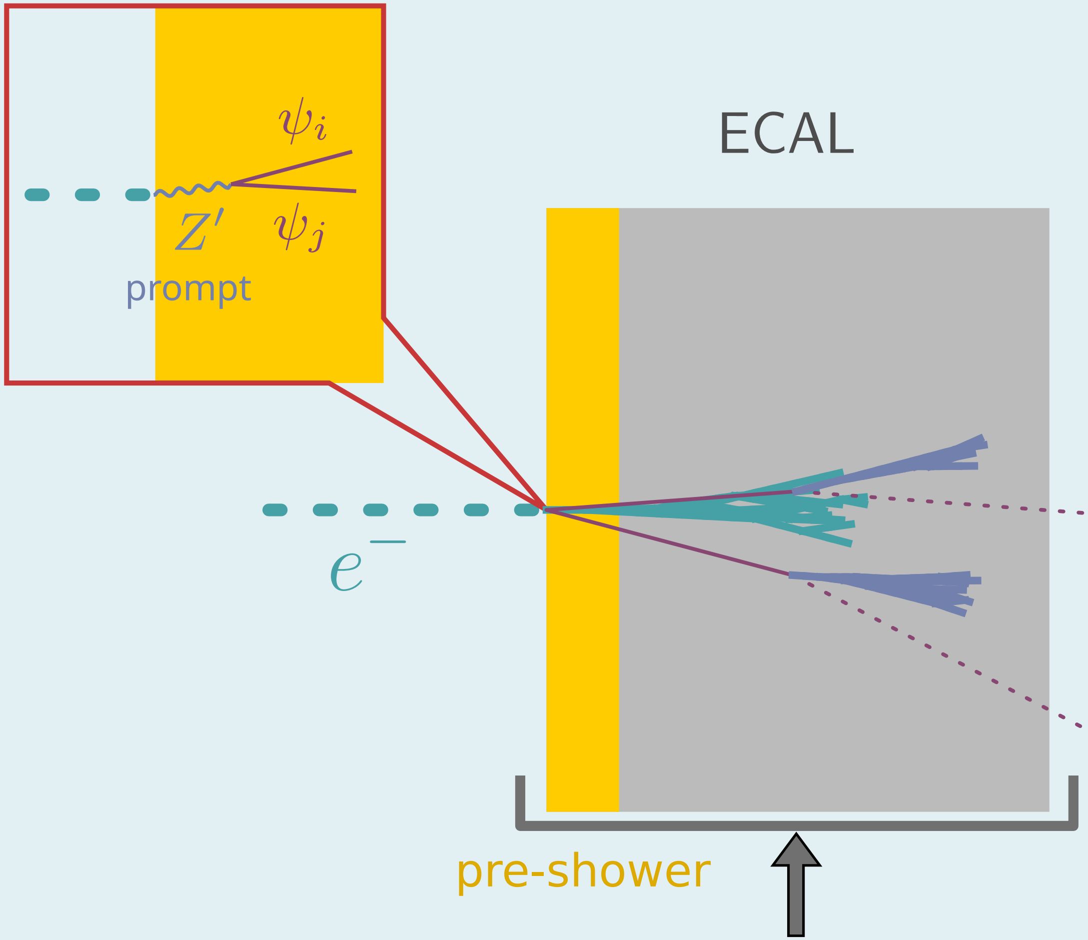
$E_{\text{tot}} \notin [0.5,\, 10]$ GeV

Assumptions:

- e^+e^- sufficiently overlapped: sum energies;
- Everything is forward: neglect *leakage* of particles from detector.

Recast of the “invisible” bound

[Phys. Rev. Lett. 123.12 (2019), p. 121801]



✔ Pre-shower ECAL
$E_{\text{tot}} \notin [0.5,\ 10]$ GeV

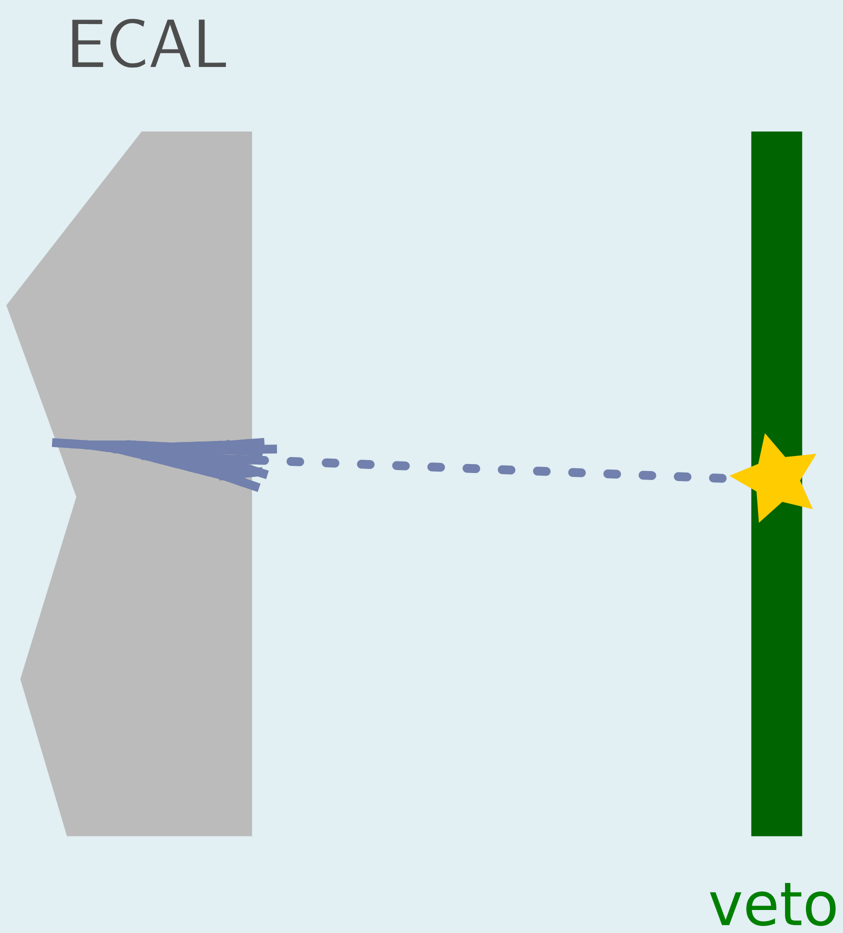
✔ ECAL
$E_{\text{tot}} > 50\,\text{GeV}$

Assumptions:

- e^+e^- sufficiently overlapped: sum energies;
- Everything is forward: neglect *leakage* of particles from detector.

Recast of the “invisible” bound

[Phys. Rev. Lett. 123.12 (2019), p. 121801]



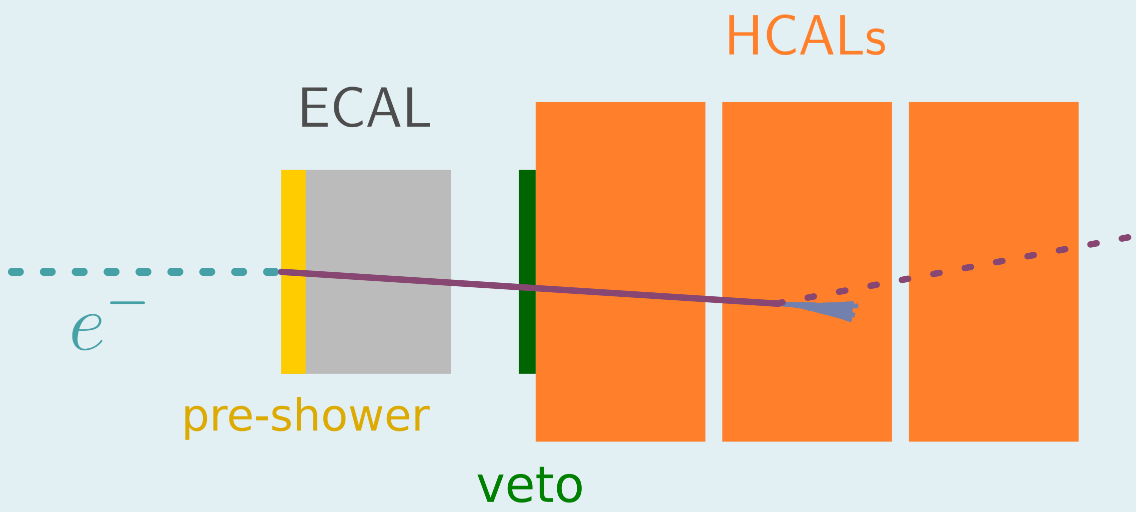
- ✔ Pre-shower ECAL
  $E_{\rm tot} \notin [0.5,\ 10]\ {\rm GeV}$
- ✔ ECAL
  $E_{\rm tot} > 50\,{\rm GeV}$
- ✔ Veto
  $E_{\pm}(x_{\rm veto}) > 10\,{\rm MeV}$

Assumptions:

- e^+e^- sufficiently overlapped: sum energies;
- Everything is forward: neglect *leakage* of particles from detector.

Recast of the "invisible" bound

[Phys. Rev. Lett. 123.12 (2019), p. 121801]



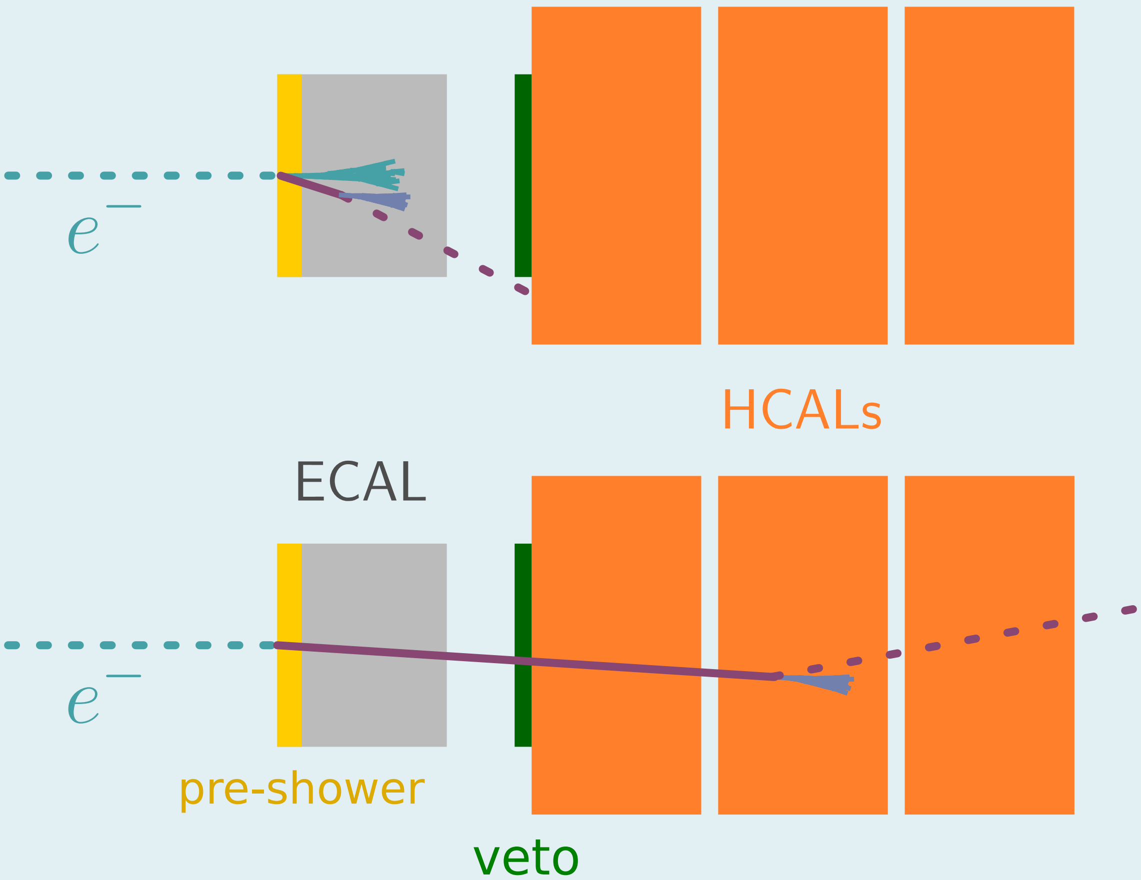
- ✔ Pre-shower ECAL
  $E_{\text{tot}} \notin [0.5,\, 10]$ GeV
- ✔ ECAL
  $E_{\text{tot}} > 50\,\text{GeV}$
- ✔ Veto
  $E_{\pm}(x_{\text{veto}}) > 10\,\text{MeV}$
- ✔ HCAL
  $E_{\pm} > 1\,\text{GeV}$

Assumptions:

- e^+e^- sufficiently overlapped: sum energies;
- Everything is forward: neglect *leakage* of particles from detector.

Semi-visible optimized cuts

[Eur. Phy. J. C 81.10 (2021), p. 959]



✔ Any e^+e^- pair triggering an activity before HCAL 2

✔ Max 1 e^+e^- pair in HCAL

Model dependent: inelastic dark matter model: 1 dark photon and 2 HNFs

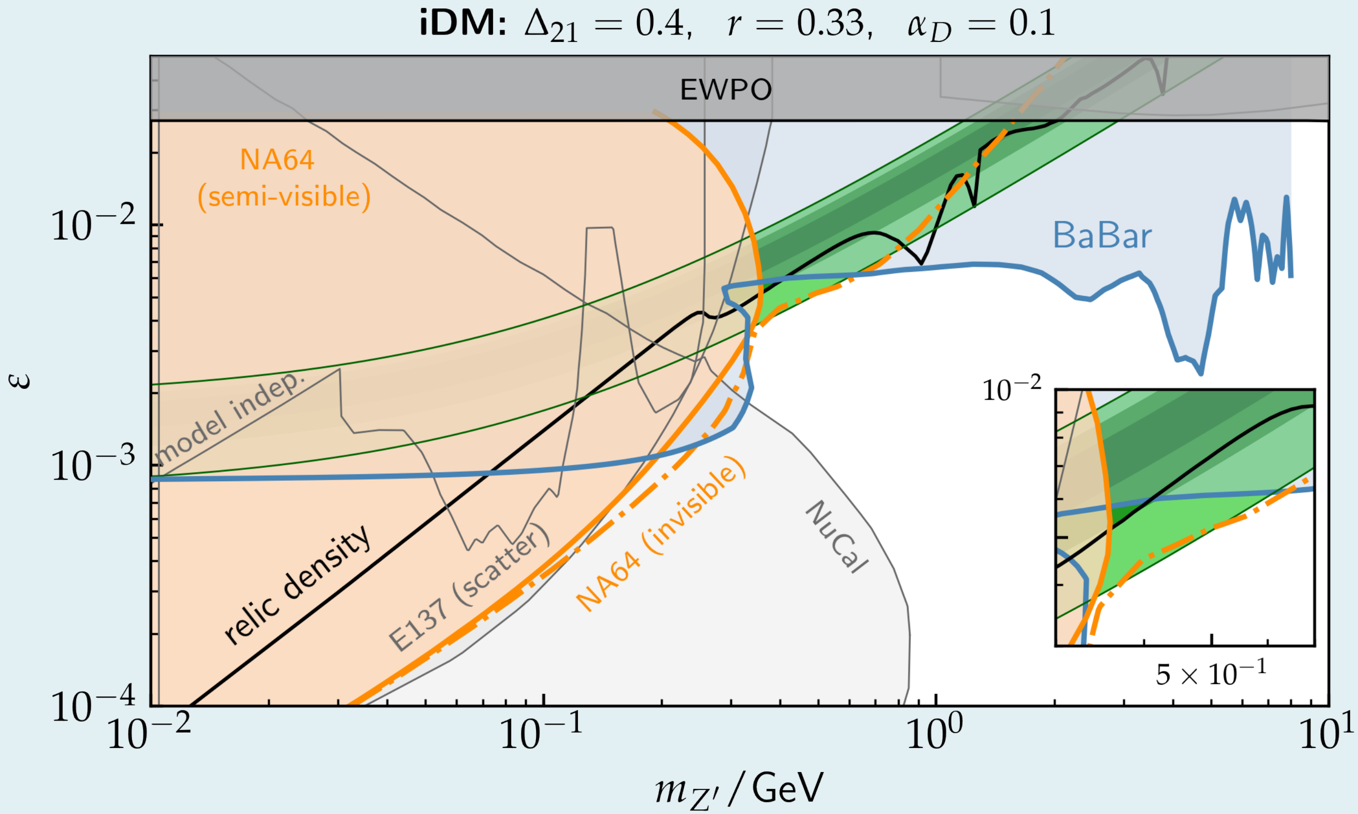
Analysis with optimized cuts.

Model 1: inelastic dark matter (iDM) with one pseudo-Dirac pair

$$\mathcal{L} \supset g_{\mathrm{D}} \overline{\psi}_1 \not{Z}' \psi_2$$

- After diagonalisation, lightest state ψ_1 is DM, with ψ_2 decaying semi-visibly
- DM relic abundance from co-annihilations

Minimal iDM is almost excluded



Model 2: iDM with several pseudo-Dirac pair

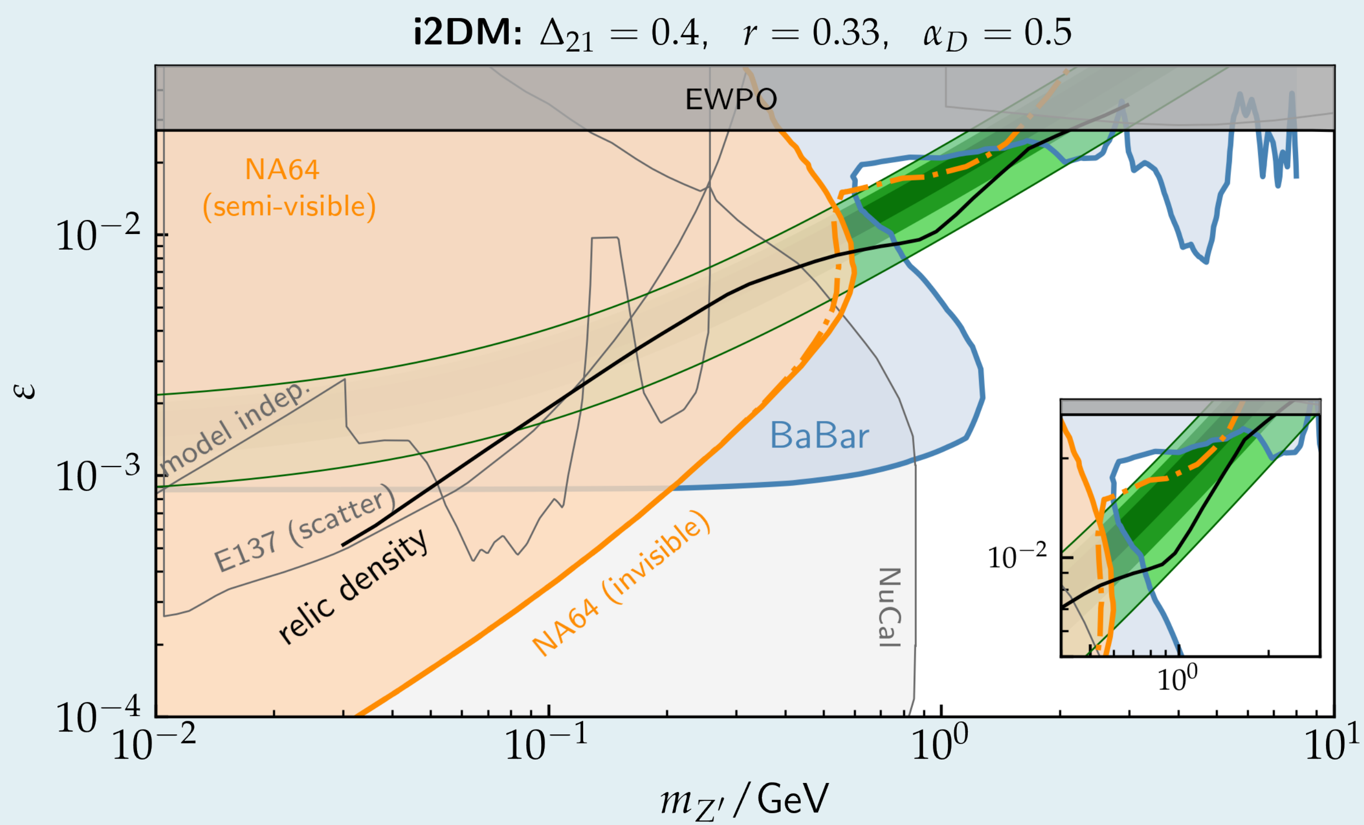
$$\mathcal{L} \supset g_D \theta^2 \overline{\psi_1} \not{Z}' \psi_1 + g_D \theta \overline{\psi_2} \not{Z}' \psi_1 + g_D \overline{\psi_2} \not{Z}' \psi_2$$

- After diagonalisation: light neutral state ψ_1 and heavier state ψ_2 more strongly coupled to dark force with small mixing ϑ

Tip: inject more visible final states

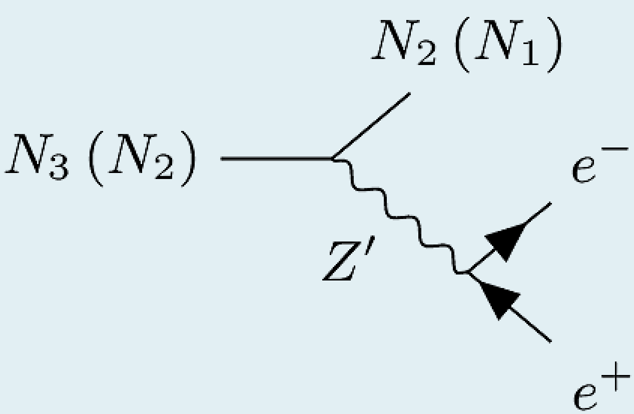
- $Z' \to \psi_1 \psi_1$
- $Z' \to \psi_1 \psi_2$
- $Z' \to \psi_2 \psi_2$

Modified iDM allowed

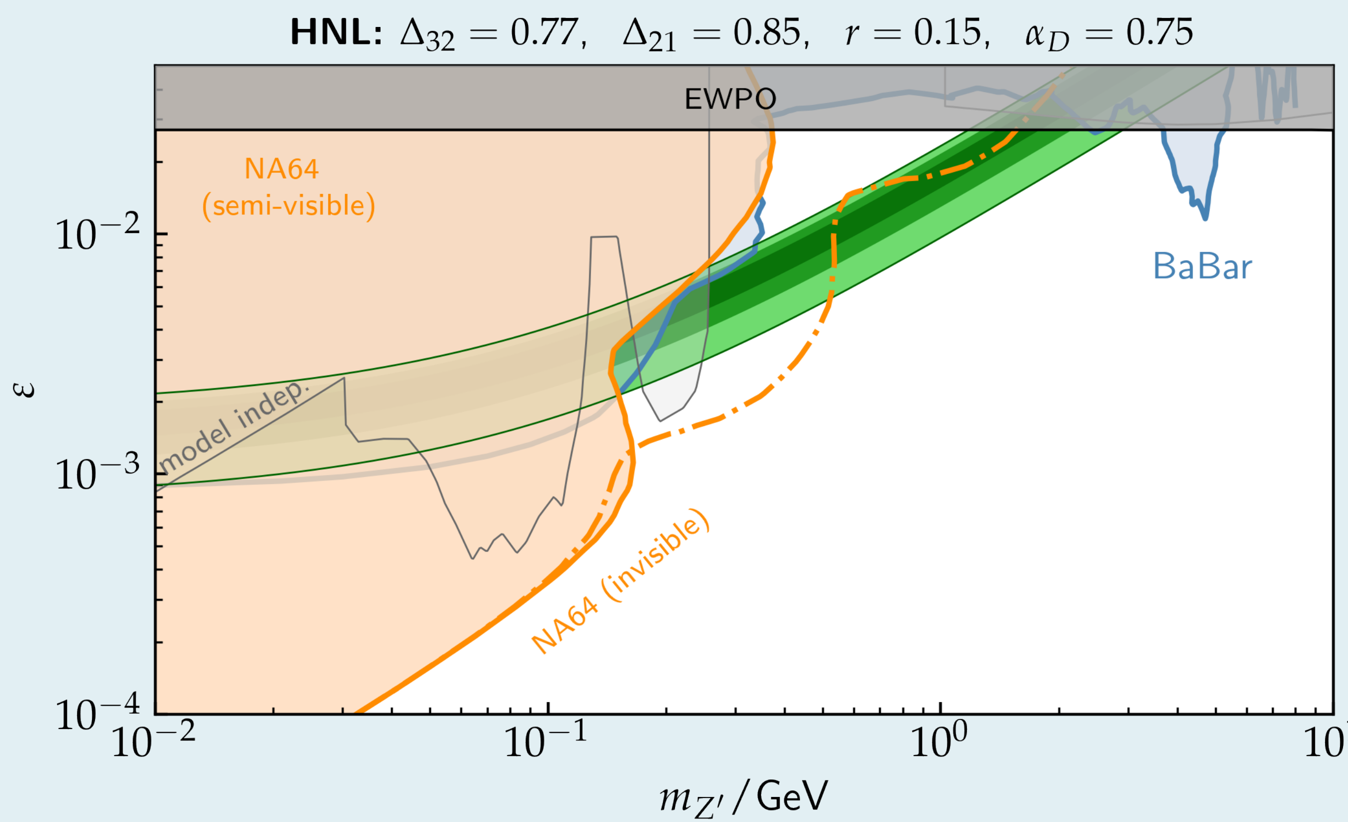


$$\mathcal{L} \supset g_{\mathrm{D}} V_{ij} \overline{N_i} \not{Z}' N_j + \text{h.c.}$$

- After diagonalisation: we obtain 3 HNLs, interacting by V_{ij}
- BPs and model: [A. Abdullahi et al. Phys. Lett. B 820 (2021), p. 136531]



HNLs model allowed



Backup slides

The BaBar simulation: veto criteria

Veto criteria: if decays happen in instrumented region of detector then veto event if both

(1) Energy of lepton tracks exceeds BaBar detection threshold

(1) $\theta_{\rm sep} > 10^\circ: \;\; E_\pm > 100\,{\rm MeV}$ well separated pairs

(2) $\theta_{\rm sep} < 10^\circ: \;\; E_+ + E_- > 100\,{\rm MeV}$ overlapping pairs

(2) Wide polar angle to avoid electrons escaping along the beam pipeline

$$17^\circ < \theta_{\rm pol} < 142^\circ$$

Recast: solve $\varepsilon_{\rm BaBar} = \varepsilon\sqrt{P^{\rm inv}(\varepsilon, m_{Z'})}$

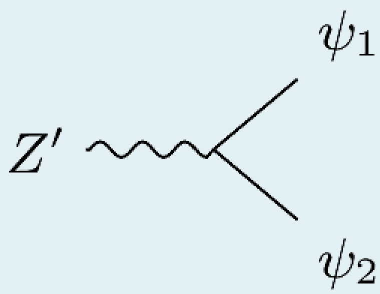
where $P^{\rm inv}(\varepsilon, m_{Z'}) = 1 - \dfrac{N_{\rm veto}(\varepsilon, m_{Z'})}{N(\varepsilon, m_{Z'})}$

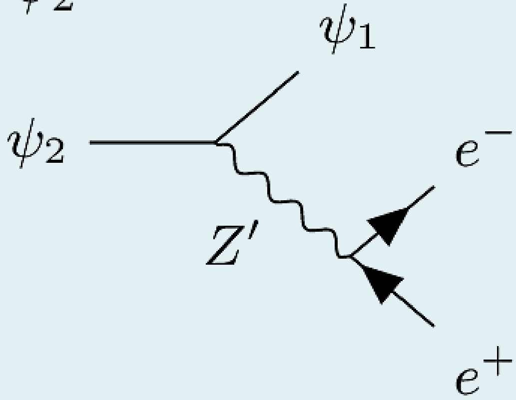
Model 1: inelastic dark matter (iDM) with one pseudo-Dirac pair

$$\mathcal{L} \supset \frac{1}{2} \begin{pmatrix} \overline{\psi_{\mathrm{L}}} & \overline{\psi^{\mathrm{c}}_{\mathrm{R}}} \end{pmatrix} \begin{pmatrix} \mu_{\mathrm{L}} & m_{\mathrm{D}} \\ m_{\mathrm{D}} & \mu_{\mathrm{R}} \end{pmatrix} \begin{pmatrix} \psi^{\mathrm{c}}_{\mathrm{L}} \\ \psi_{\mathrm{R}} \end{pmatrix} + \text{h.c.}$$

- $\mu_{\mathrm{L}} \approx \mu_{\mathrm{R}} \ll m_{\mathrm{D}}$: on-diagonal couplings suppressed
- After diagonalisation, lightest state ψ_1 is DM, with ψ_2 decaying visibly
- DM relic abundance from co-annihilations

$$\mathcal{L} \supset g_{\mathrm{D}} \overline{\psi}_1 \not{Z}' \psi_2$$





- Proposed in [G. Mohlabeng. Phys. Rev. D 99.11 (2019), p. 115001], where the study assumed BaBar energy threshold of 60 MeV (optimistic, we used 100 MeV, which is the same used in BaBar selection criteria)

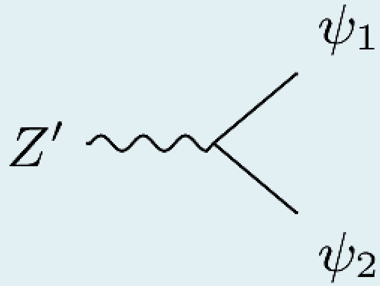
Model 2: iDM with several pseudo-Dirac pair

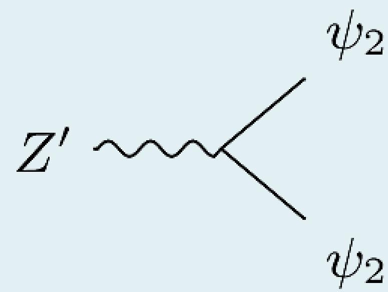
$$\mathcal{L} \supset \begin{pmatrix} \eta_{\rm L} & \chi_{\rm L} \end{pmatrix} \begin{pmatrix} M_1 & M_{\rm L} \\ M_{\rm R} & M_2 \end{pmatrix} \begin{pmatrix} \eta_{\rm R} \\ \chi_{\rm R} \end{pmatrix} + \text{h.c.}$$

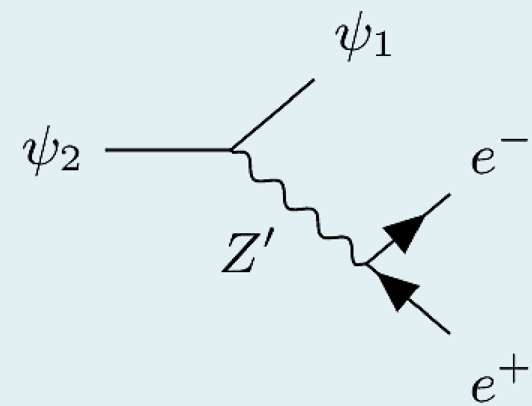
- neutral (η) and U(1)’ charged (χ) pseudo-Dirac pairs
- After diagonalisation: light neutral state ψ_1 and heavier state ψ_2 more strongly coupled to dark force with small mixing ϑ

$$\mathcal{L} \supset g_{\rm D}\theta^2\, \overline{\psi_1} \not{Z}' \psi_1 + g_{\rm D}\theta\, \overline{\psi_2} \not{Z}' \psi_1 + g_{\rm D} \overline{\psi_2} \not{Z}' \psi_2$$

- $Z' \to \psi_1 \psi_1$
- $Z' \to \psi_1 \psi_2$
- $Z' \to \psi_2 \psi_2$







- ψ_1 can not be dark matter, unless secluded annihilations
- If not DM, then must decay (e.g. mixing w/ SM ν).

Model 3: HNLs

[A. Abdullahi et al. Phys. Lett. B 820 (2021), p. 136531]

$$\mathcal{L} \supset \left(\overline{N} \quad \overline{\psi}_{\mathrm{L}} \quad \overline{\psi}^{\mathsf{c}}_{\mathrm{R}}\right) \begin{pmatrix} M_N & \Lambda_{\mathrm{L}} & \Lambda_{\mathrm{R}} \\ \Lambda_{\mathrm{L}}^{\mathsf{T}} & 0 & M_X \\ \Lambda_{\mathrm{R}}^{\mathsf{T}} & M_X^{\mathsf{T}} & 0 \end{pmatrix} \begin{pmatrix} N^{\mathsf{c}} \\ \psi_{\mathrm{L}}^{\mathsf{c}} \\ \psi_{\mathrm{R}} \end{pmatrix} + \text{h.c.}$$

- With neutral state N and U(1)'-charged fermions
- After diagonalisation we obtain 3 HNLs, interacting by V_{ij} (in principle they can mix also with SM ν).

$$\mathcal{L} \supset g_{\mathrm{D}} V_{ij} \overline{N_i} \not{Z}' N_j + \text{h.c.}$$