Nucleon Electromagnetic Form Factors at Large Momentum from Lattice QCD

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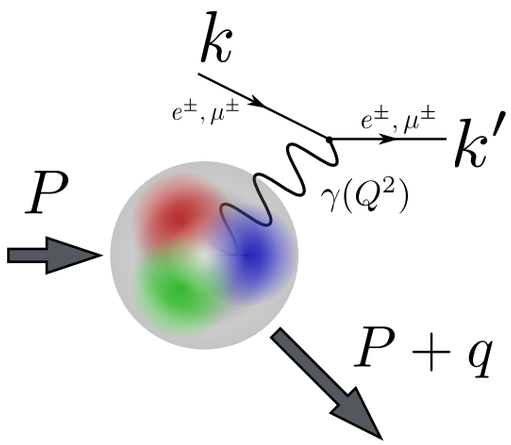
Outline

- Nucleon vector form factors : introduction and motivation
- Challenges for high-Q^2 hadron structure on Lattice
- Preliminary results and comparison to experiment & phenomenology
- Examining systematic effects
- Summary and Outlook

Nucleon Elastic E&M Form Factors



$$Q^2 = -(k-k')^2 > 0$$

Elastic e^-p amplitude

$$\langle P+q|\,\bar{q}\gamma^\mu q\,|P\rangle = \bar{U}_{P+q}\Big[\underbrace{F_1(Q^2)}_{\text{(Dirac)}}\gamma^\mu + \underbrace{F_2(Q^2)}_{\text{(Pauli)}}\frac{i\sigma^{\mu\nu}q_\nu}{2M_N}\Big]U_P$$

Sachs Electric $\quad G_E(Q^2) = F_1(Q^2) - \dfrac{Q^2}{4M^2}F_2(Q^2)$

Magnetic $\quad G_M(Q^2) = F_1(Q^2) + F_2(Q^2)$

Elastic e^-p cross-section

- $G_{E,M}$ from ϵ-dep. at fixed $\tau(Q^2)$ ("***Rosenbluth separation***")
- dominated by G_M at large Q^2
- 2γ corrections at $Q^2 \gtrsim 1$ GeV^2

$$\frac{d\sigma}{d\Omega} = \frac{\sigma_{\rm Mott}}{1+\tau}\left[G_E^2 + \frac{\tau}{\epsilon}G_M^2\right]$$

$$\tau = \frac{Q^2}{4M_N^2} \qquad \epsilon = \left[1 + 2(1+\tau)\tan^2\frac{\theta}{2}\right]^{-1}$$

Polarization transfer: polarized e^-beam
+ detect polarization of recoil nucleon
(alt.: transverse asymmetry on pol. target)

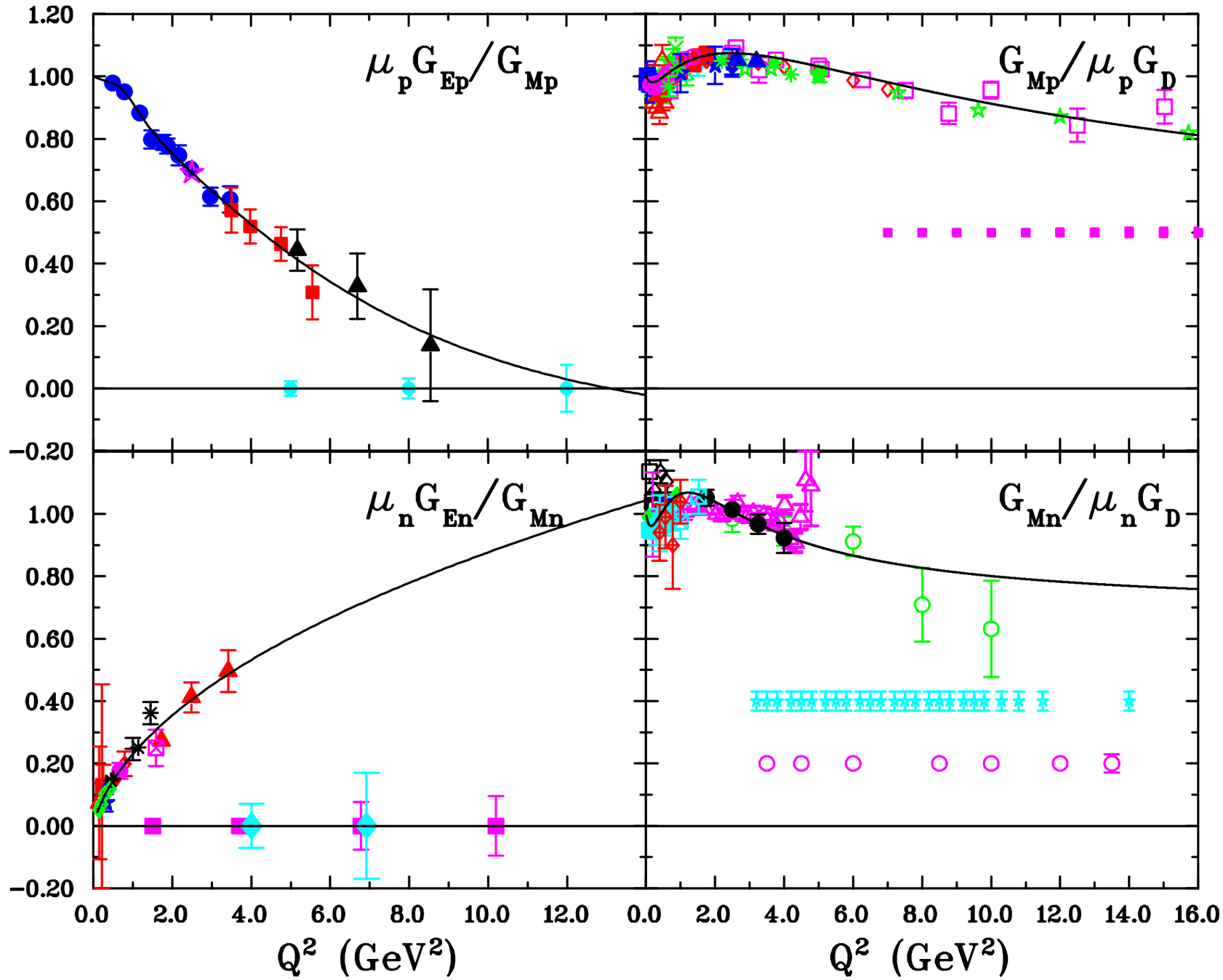
- G_E/G_M *ratio* (only small radiative corrections)

$$P_t/P_l \propto G_E/G_M$$

Recent/Ongoing Experiments

Experiments at JLab@12GeV

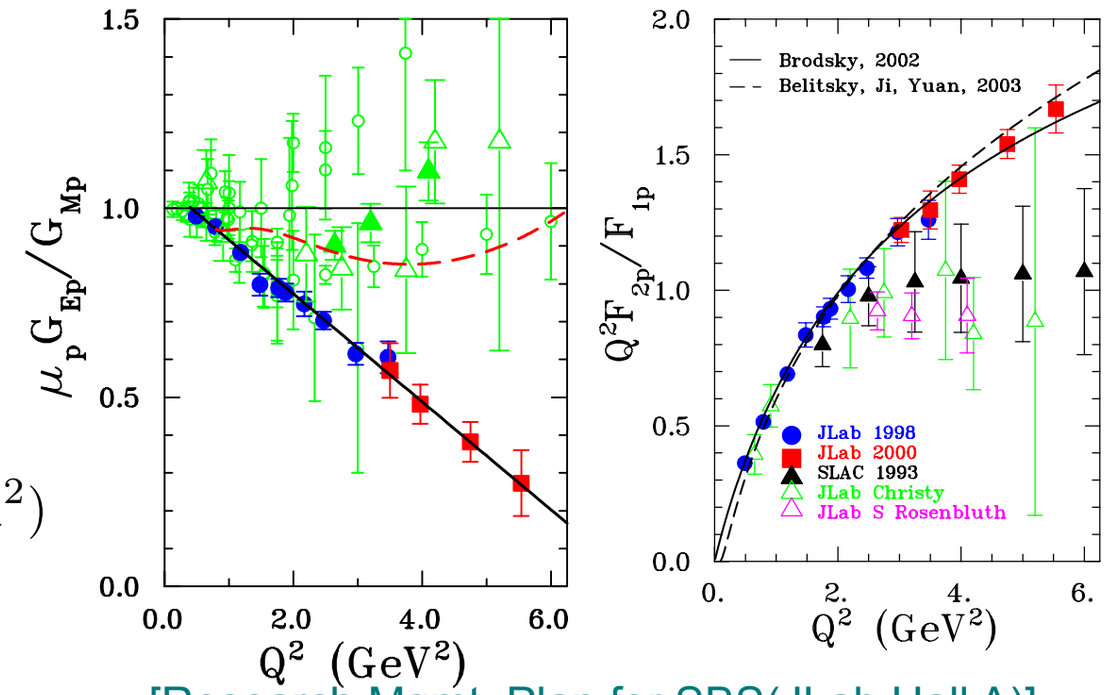
- Hall A (HRS, SBS):
  G_{Mp} @ $Q^2 \lesssim 17.5$ GeV^2
  G_{Ep}/G_{Mp} @ $Q^2 \lesssim 15$ GeV^2 ;
  G_{Mn} @ $Q^2 \lesssim 18$ GeV^2
  G_{En}/G_{Mn} @ $Q^2 \lesssim 10.2$ GeV^2 ;
- Hall B (CLAS12):
  G_{Mn} @ $Q^2 \lesssim 14$ GeV^2
- Hall C :
  G_{En}/G_{Mn} @ $Q^2 \lesssim 6.9$ GeV^2



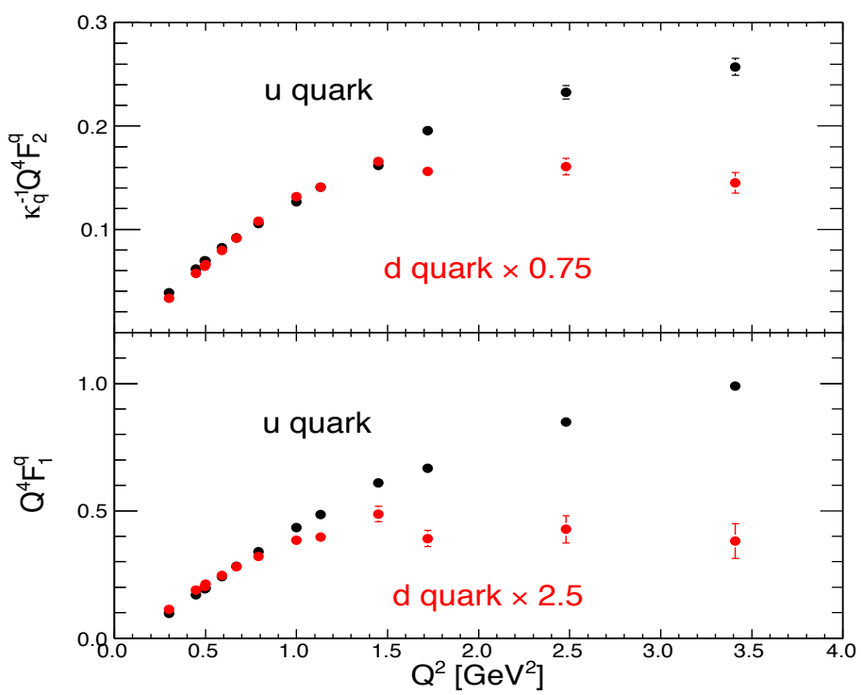
Projected new precision on proton & neutron form factors
[V. Punjabi et al, EPJ A51: 79 (2015); arXiv: 1503.01452]

Nucleon Form Factors: Open Questions

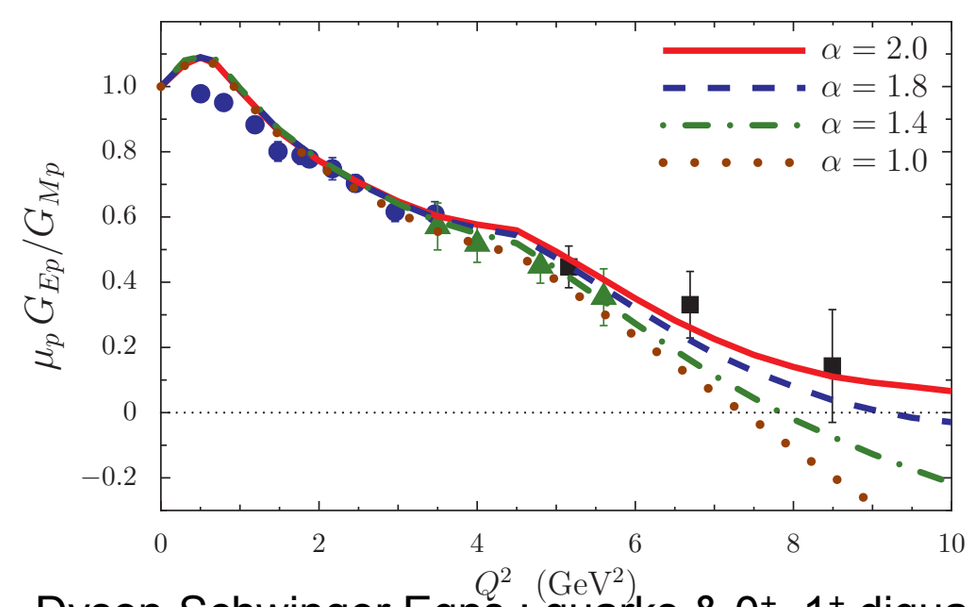
- *Are model descriptions of the nucleon viable ?*
  Nucleon models disagree where data unavailable
- *Role of diquark correlations in elastic scattering ?*
  Neutron & proton G_E/G_M at/above $Q^2 = 8$ GeV^2
- *Scale of transition to perturbative QCD ?*
  (F_2/F_1) scaling at large Q^2 : $Q^2 F_{2p}/F_{1p} \overset{?}{\propto} \log^2(Q^2/\Lambda^2)$
- *What are contributions from u and d flavors?*
  Proton <u>and</u> neutron data needed in wide Q^2 range



[Research Mgmt. Plan for SBS(JLab Hall A)]



[G.D.Cates, C.W.de Jager, S.Riordan, B.Wojtsekhovski, PRL106:252003, arXiv:1103.1808]



Dyson-Schwinger Eqns : quarks & 0^+, 1^+ diquarks
($\alpha \approx$ rate of transition CQ $\rightarrow$ pQCD with Q^2)
[Cloet, Roberts, Prog.Part.Nucl.Phys 77:1 (2014)]

Nucleon Structure from QCD on a Lattice

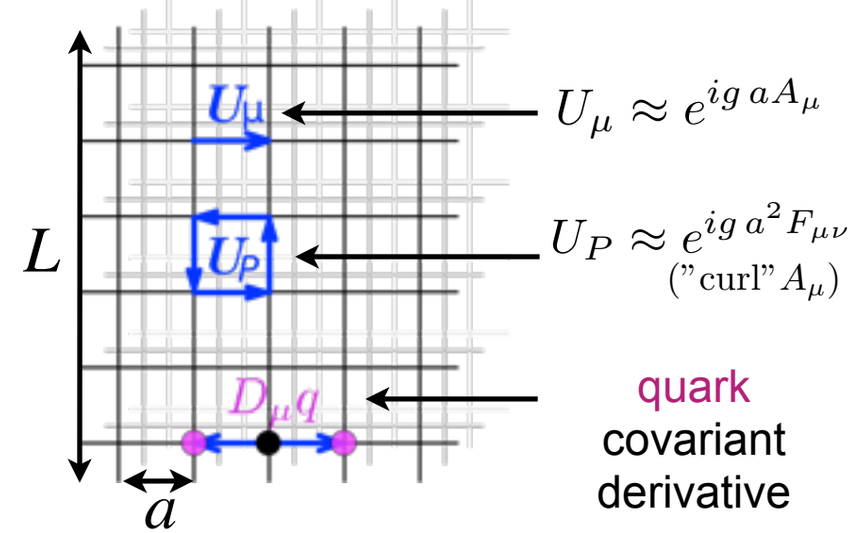
Lattice Field Theory ⇔ Numerical evaluation of the Path Integral

$$\langle q_x \bar{q}_y \ldots \rangle = \int \mathcal{D}\left(Glue\right) \int \mathcal{D}\left(Quarks\right) \; e^{-S_{Glue} - \bar{q}\left(\not{D}+m\right)q} \; \left[q_x \bar{q}_y \ldots\right]$$

Grassmann integration

$$= \int \underbrace{\mathcal{D}\left(Glue\right) e^{-S_{Glue}} \, \mathrm{Det}\left(\not{D}+m\right)}_{\text{Hybrid Monte Carlo sampling of gluon background}} \; \underbrace{\left[\left(\not{D}+m\right)^{-1}_{x,y} \ldots\right]}_{\text{Quark Eqns. of Motion}}$$

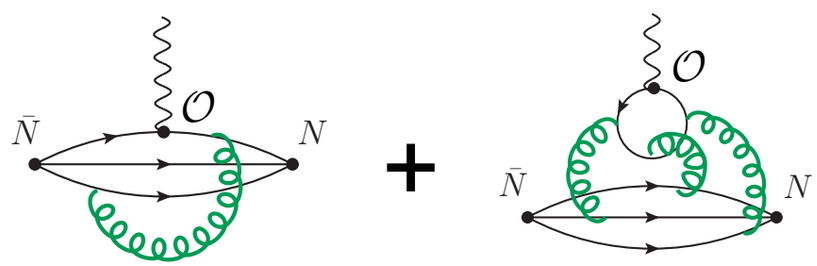
QCD on Euclidean 4D lattice



Hadron Matrix Elements:

$$C^{\mathcal{O}}_{3\mathrm{pt}}(T) = \langle N(T)\mathcal{O}(\tau)\bar{N}(0)\rangle =$$

$$[N \propto (uC\gamma_5 d)u]$$



quark-connected quark-disconnected

$$\langle N(T)\mathcal{O}(\tau)N(0)\rangle = \sum_{n,m} Z_m e^{-E_n(T-\tau)} \langle n|\mathcal{O}|m\rangle e^{-E_m \tau} Z^*_n$$

$$\xrightarrow[T\to\infty]{} Z_{00} e^{-M_N T}\left[\langle P'|\mathcal{O}|P\rangle + \mathcal{O}\Big(\underbrace{e^{-\Delta E_{10} T}, e^{-\Delta E_{10}\tau}, e^{-\Delta E_{10}(T-\tau)}}_{\text{excited states}}\Big)\right]$$

Ground state form factors

Fit and discard

Systematic uncertainties

- discretization effects
- unphys.-heavy pion mass
- finite volume
- excited state contributions
- (costly) disconnected diagrams

Challenges at Large Q^2

- Multiscale problem: $a \ll p_N^{-1}, Q^{-1}\,; \quad 4m_\pi^{-1} \lesssim L$

- Discretization effects:
  O(a) Correction to current operator

$$(V_\mu)_I = [\bar{q}\gamma_\mu q] + c_V\, a\, \underbrace{\partial_\nu[\bar{q} i\sigma_{\mu\nu} q]}_{\propto Q}$$

- Stochastic noise grows faster with *T* [Lepage'89]:

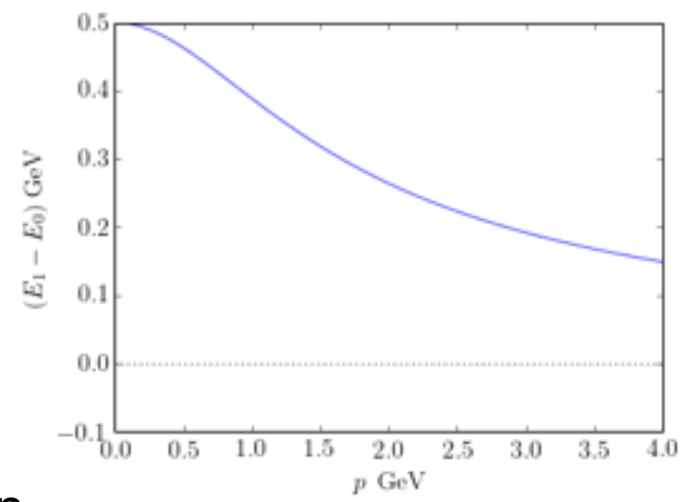
$$\begin{aligned}
&\text{Signal} \quad \langle N(T)\bar{N}(0)\rangle && \sim e^{-E_N T}\\
&\text{Noise} \quad \langle |N(T)\bar{N}(0)|^2\rangle - |\langle N(T)\bar{N}(0)\rangle|^2 && \sim e^{-3m_\pi T}\\
&\text{Signal/Noise} && \sim e^{-(E_N - \frac{3}{2}m_\pi)T}
\end{aligned}$$

SNR reduction
at 1 fm/c ***~ O(10^{-6})***
(phys.quarks, $Q^2 \approx 8$ GeV^2)

- Excited states: boosting "shrinks" the energy gap

$$E_1 - E_0 = \sqrt{M_1^2 + \vec{p}^2} - \sqrt{M_2^2 + \vec{p}^2} < M_1 - M_0$$

  - N(~1500): pN→1.5 GeV ⇒ ΔE = 500 → 300 MeV

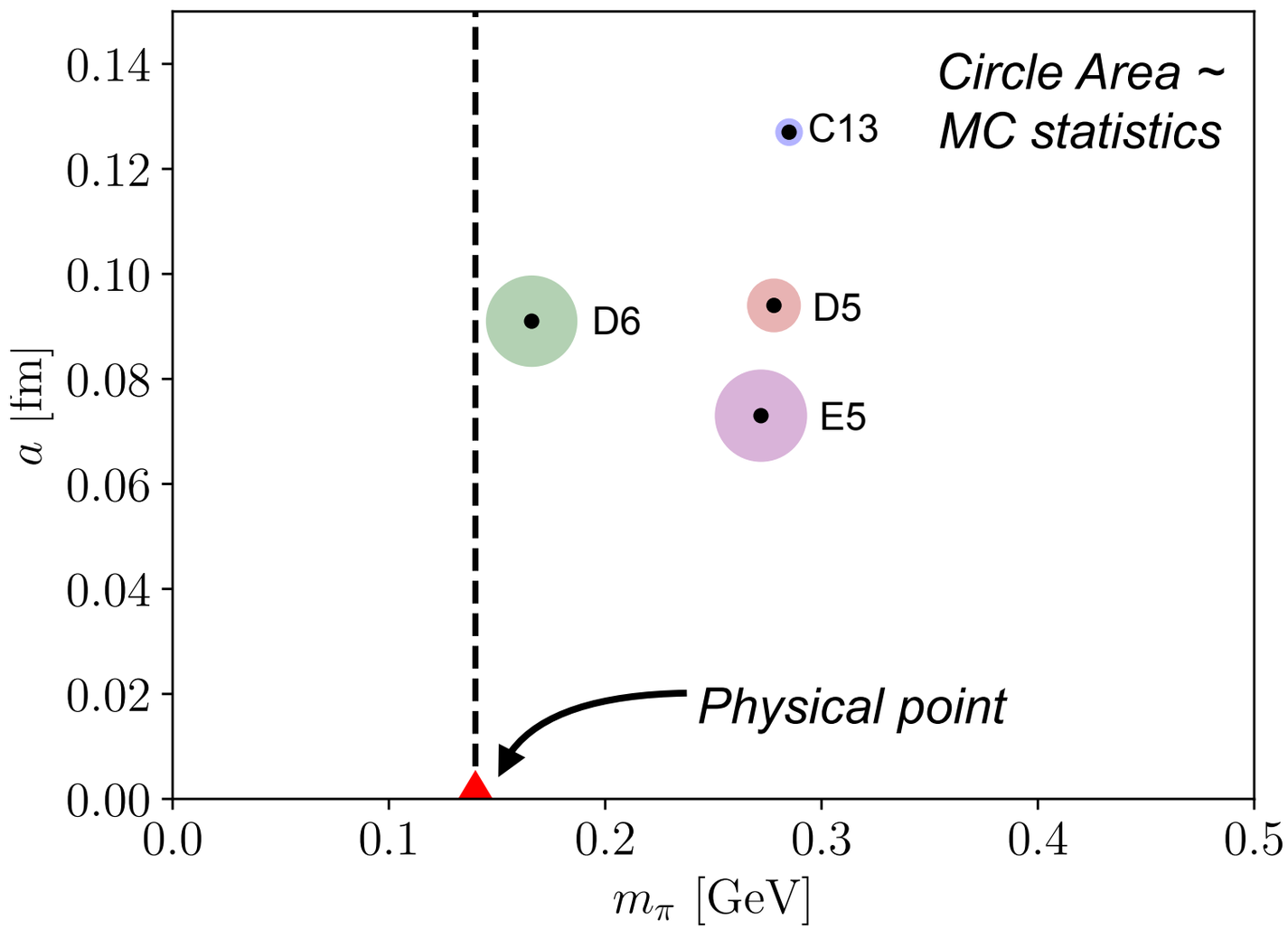


- Large p_N: no reliable EFT/ChPT for mπ-, lattice size-extrapolation

Large statistics required to suppress MC noise in lattice correlators

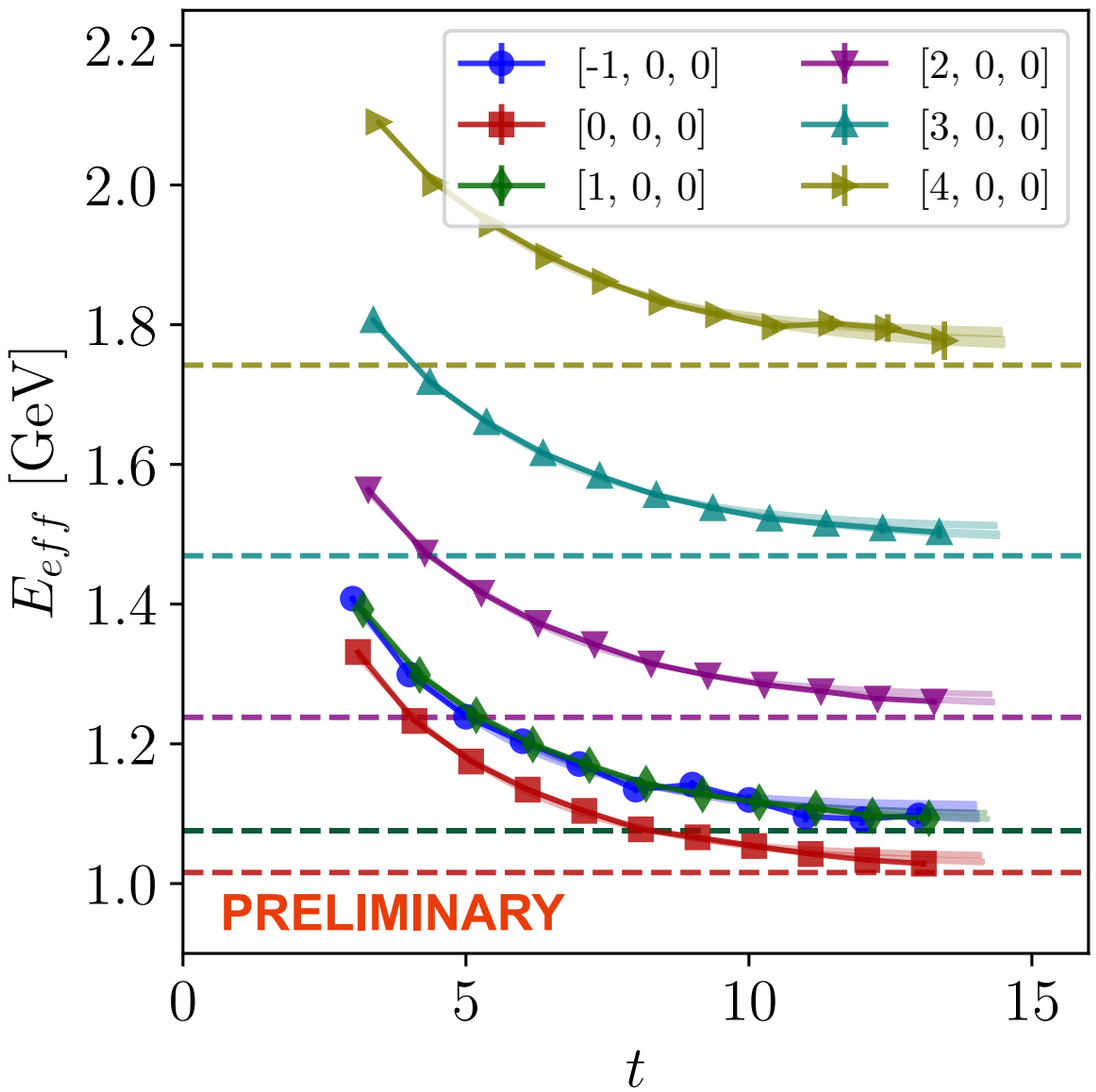
Present QCD Simulation Parameters

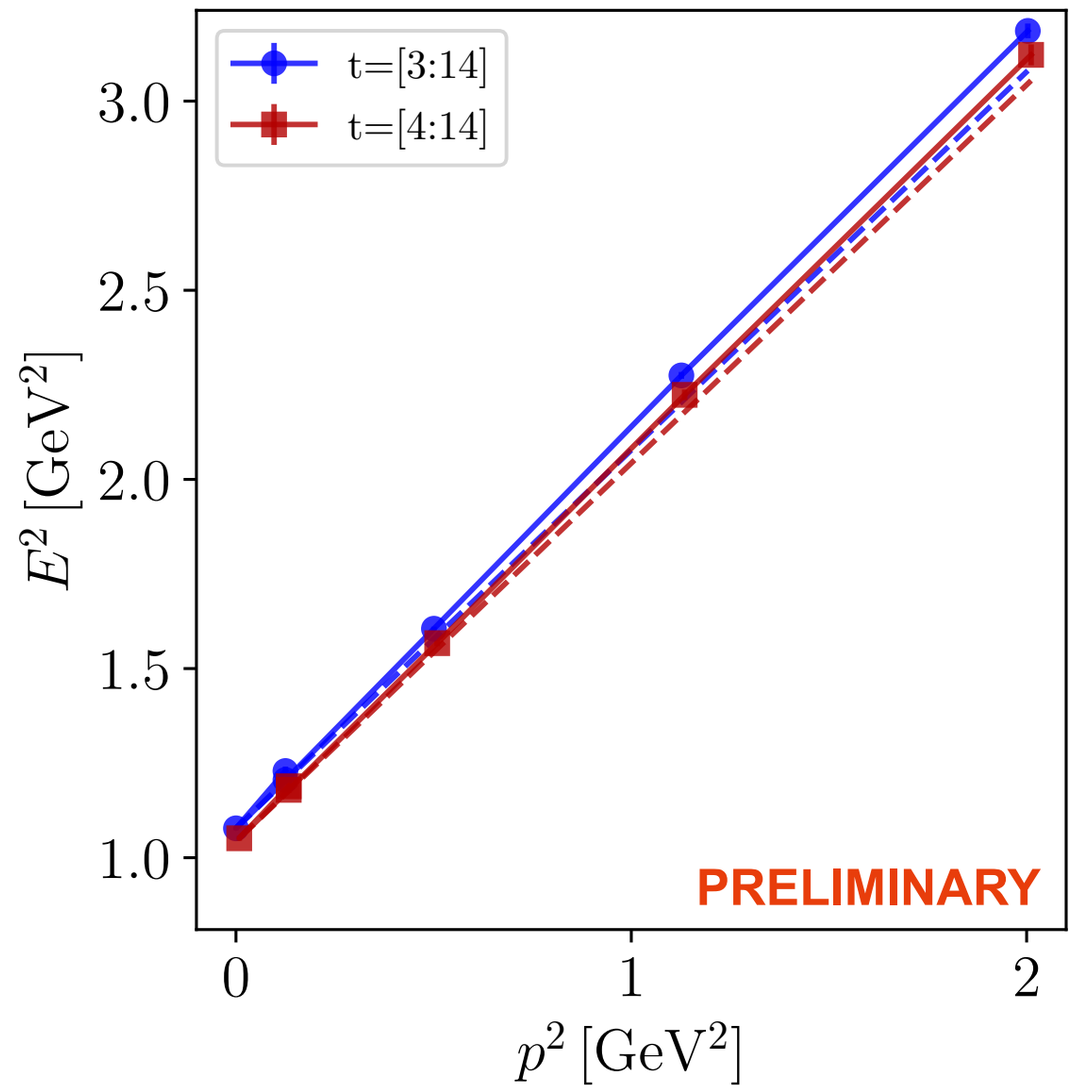
- $N_F = 2+1$ clover-improved Wilson fermion ensembles (JLab / W&M / LANL / MIT)
- lattice spacing $a \approx 0.07 \div 0.12$ fm
- unphysical light quark masses : $m\pi = 170 \div 280$ MeV
- large physical volume $L \gtrsim 3.7\ (m\pi)^{-1}$



Lattice Nucleon Energy & Dispersion Relation (E5)

- "E5" ensemble: $m\pi = 272$ MeV , spacing $a = 0.073$ fm , 266k MC samples





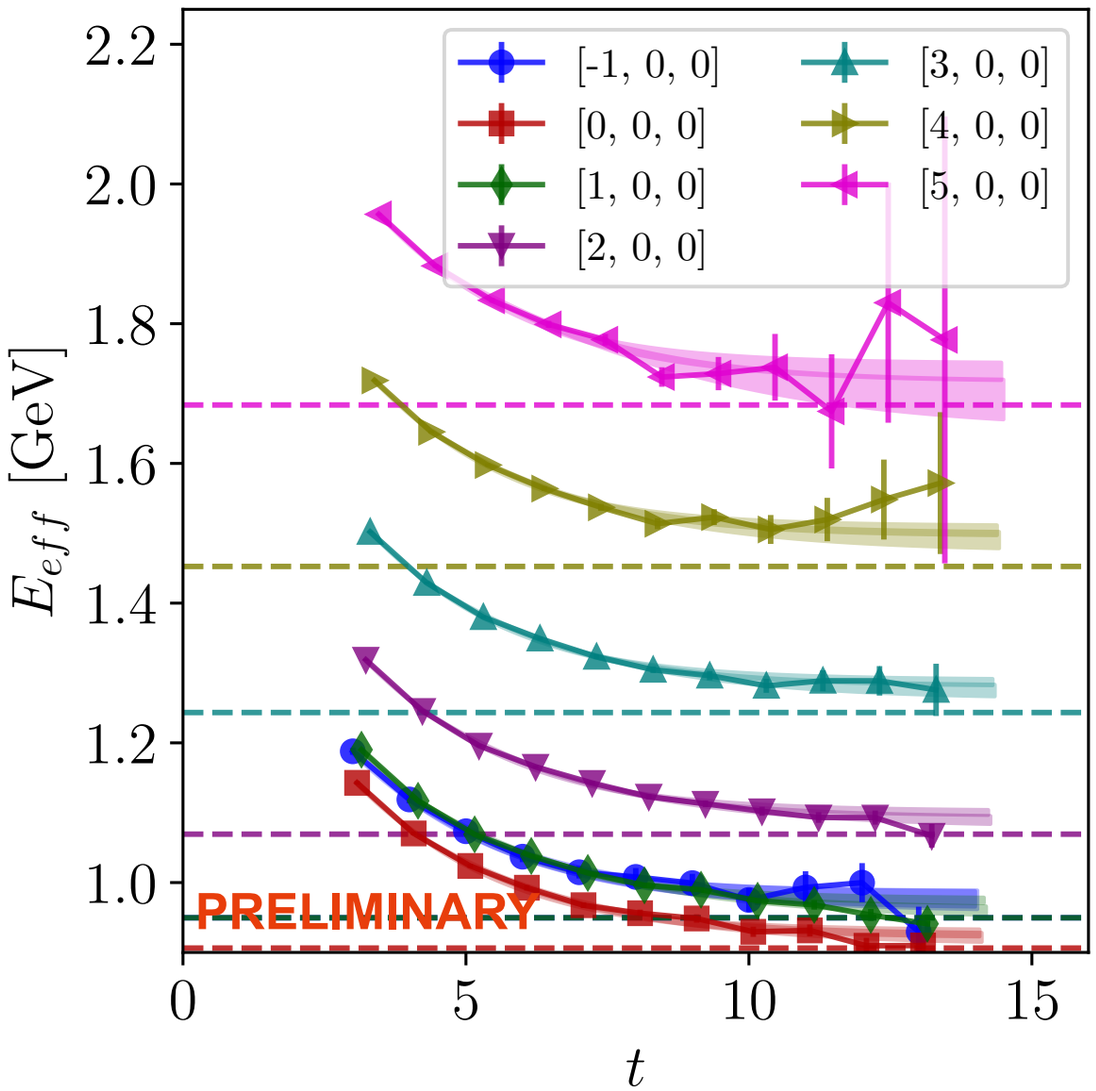
- Effective energy and 2-state fits

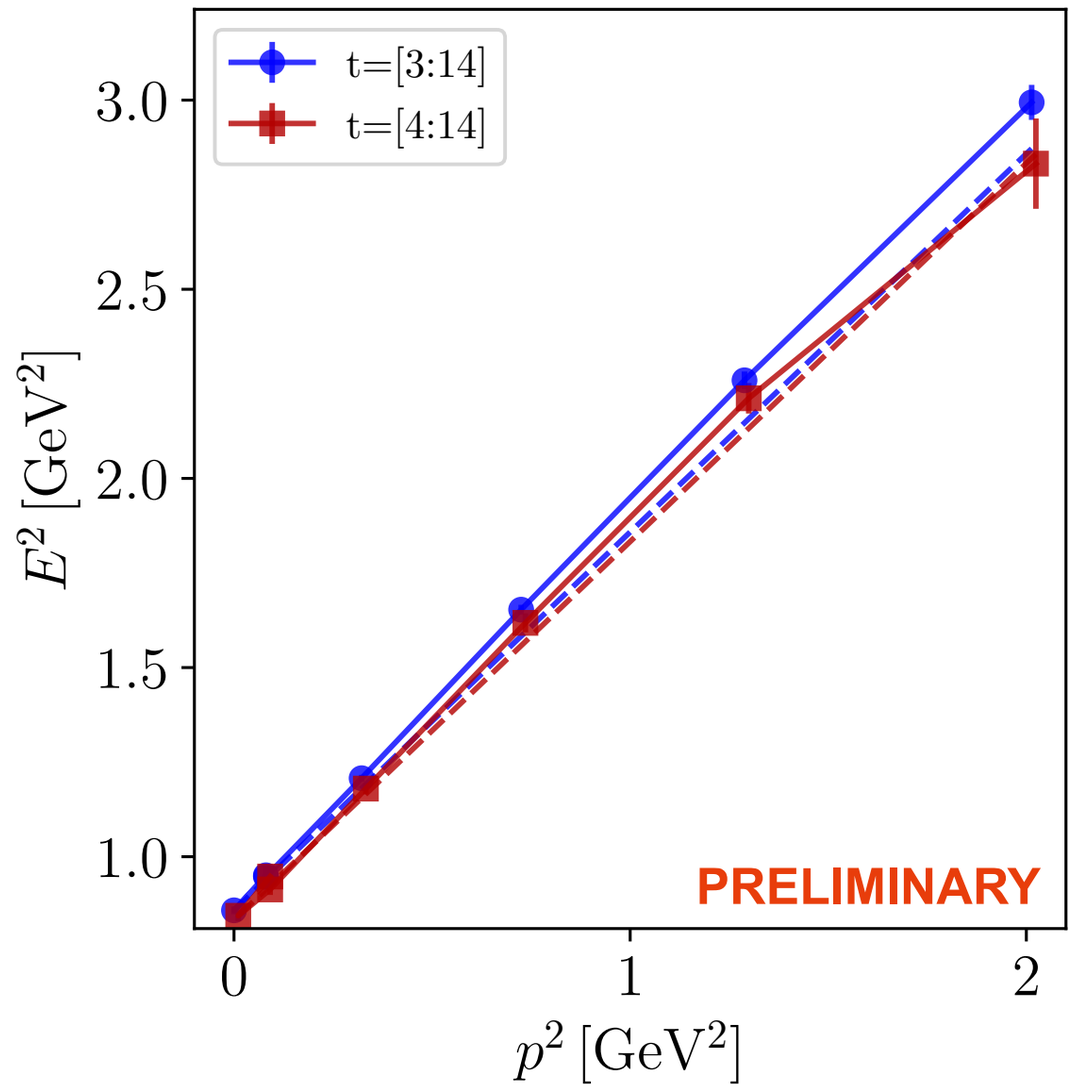
$$E_{eff} = \frac{1}{a} \log \frac{C_{N\bar{N}}(t)}{C_{N\bar{N}}(t+a)}$$

- Dispersion relation
  Dashed lines: cont. $E^2(p) = E^2(0) + p^2$

Lattice Nucleon Energy & Dispersion Relation (D6)

- "D6" ensemble: $m\pi = 166$ MeV , spacing $a = 0.091$ fm , 261k MC samples



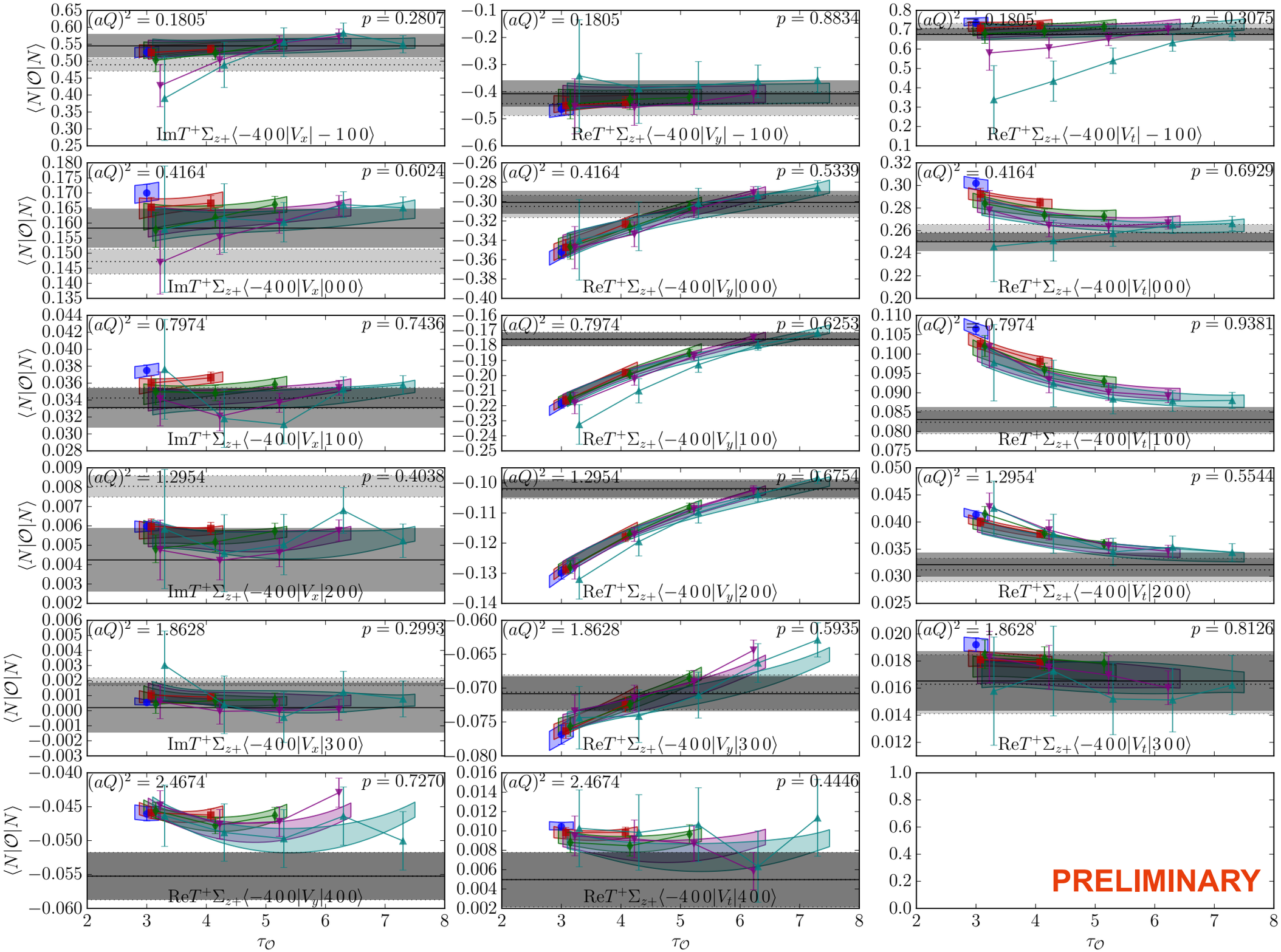


- Effective energy and 2-state fits

$$E_{eff} = \frac{1}{a} \log \frac{C_{N\bar{N}}(t)}{C_{N\bar{N}}(t+a)}$$

- Dispersion relation
  Dashed lines: cont. $E^2(p) = E^2(0) + p^2$

Nucleon Matrix Element & Form Factor Fits (D5)



$(aQ)^2 = 0.18$
$N_{dof} = 1$

$(aQ)^2 = 0.42$
$N_{dof} = 1$

$(aQ)^2 = 0.80$
$N_{dof} = 1$

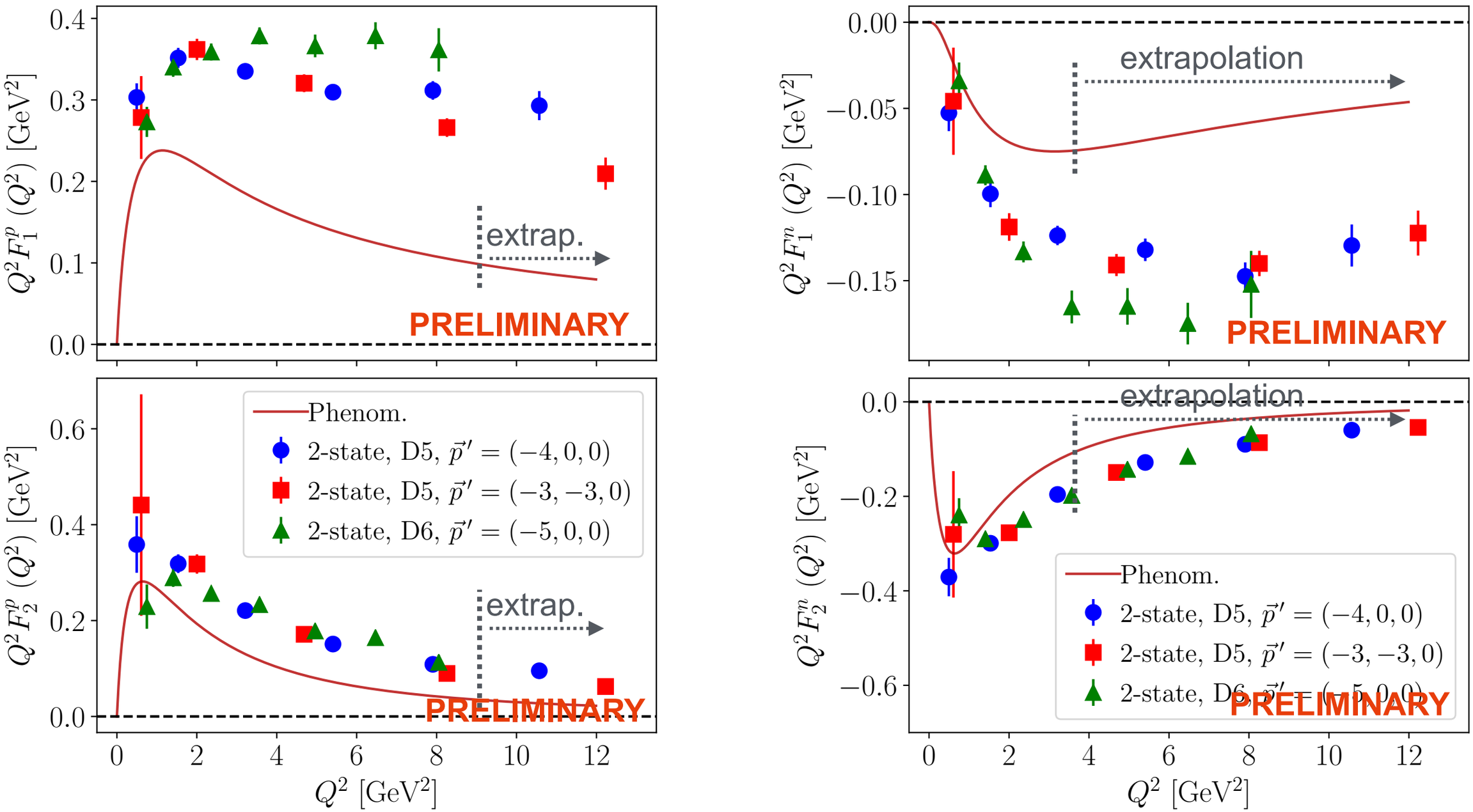
$(aQ)^2 = 1.30$
$N_{dof} = 1$

$(aQ)^2 = 1.86$
$N_{dof} = 1$

$(aQ)^2 = 2.47$
$N_{dof} = 0$

Results : Nucleon Form Factors

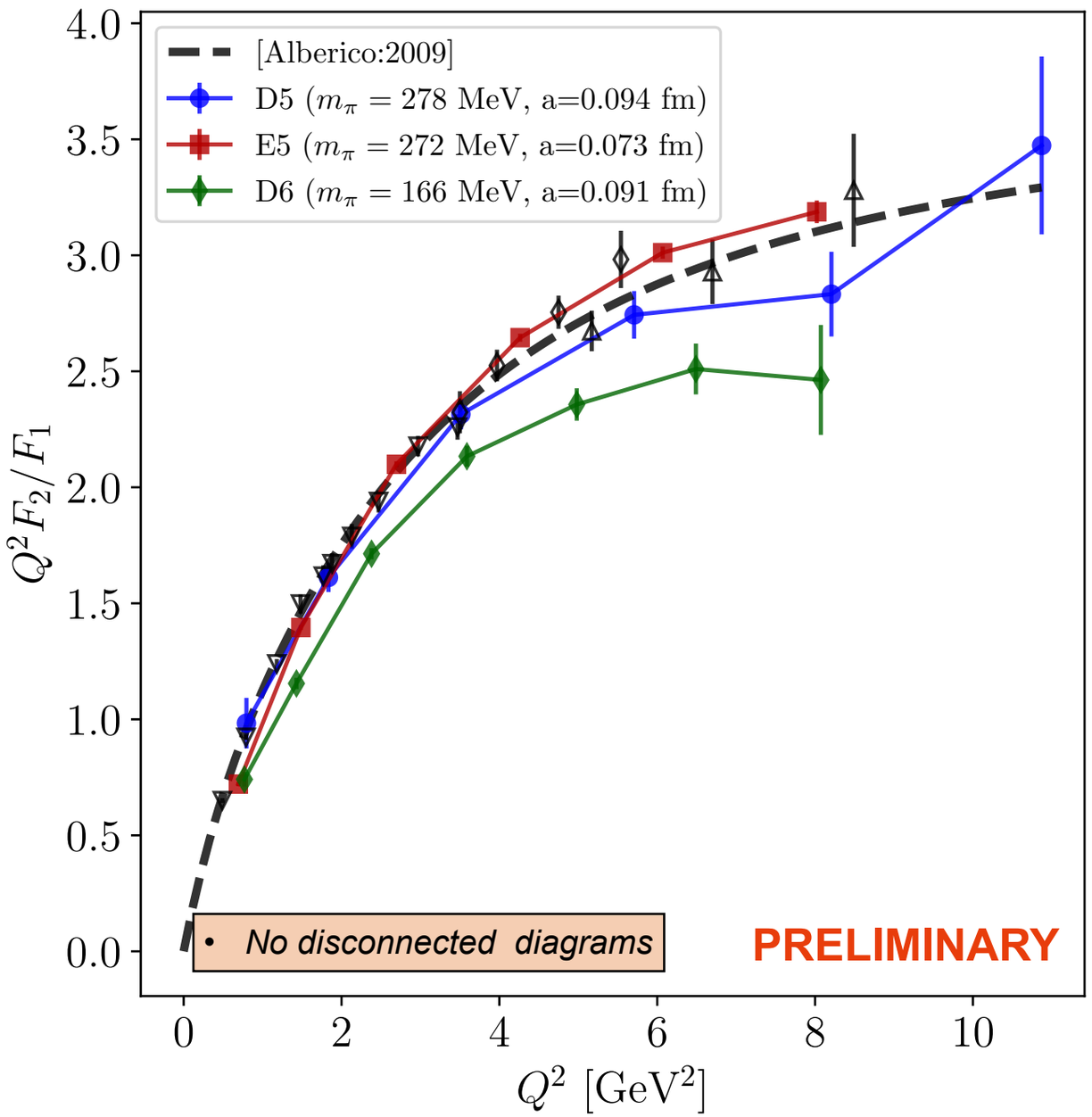
- Phenomenology curves : [Alberico et al, PRC79:065204 (2008)]

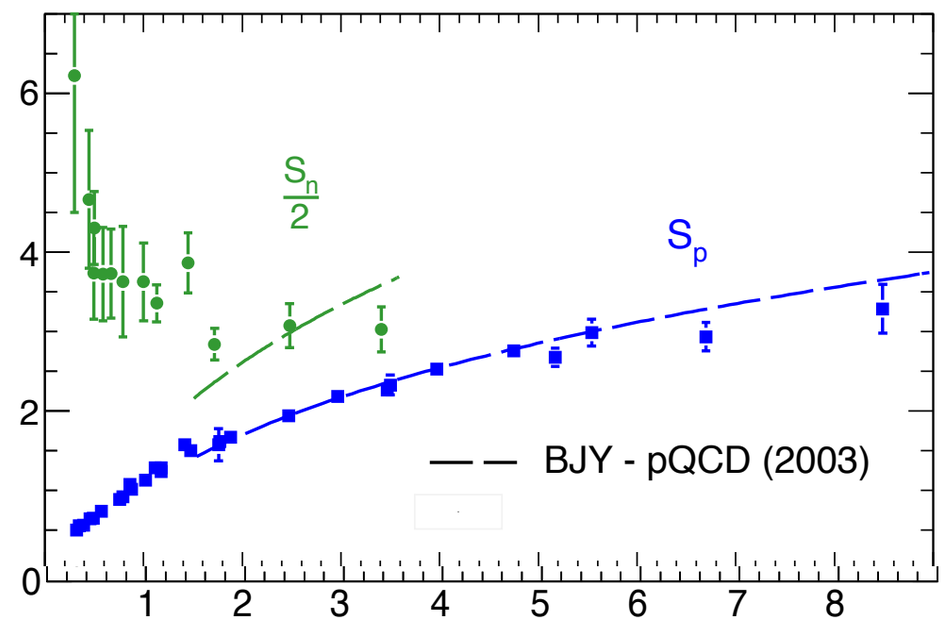


- *No disconnected diagrams*

F_{2p}/F_{1p} Ratio, Proton

- Lattice data: 2-state fits
- Phenomenology curves : [Alberico et al, PRC79:065204 (2008)]
- Comparison to experiments (black points)

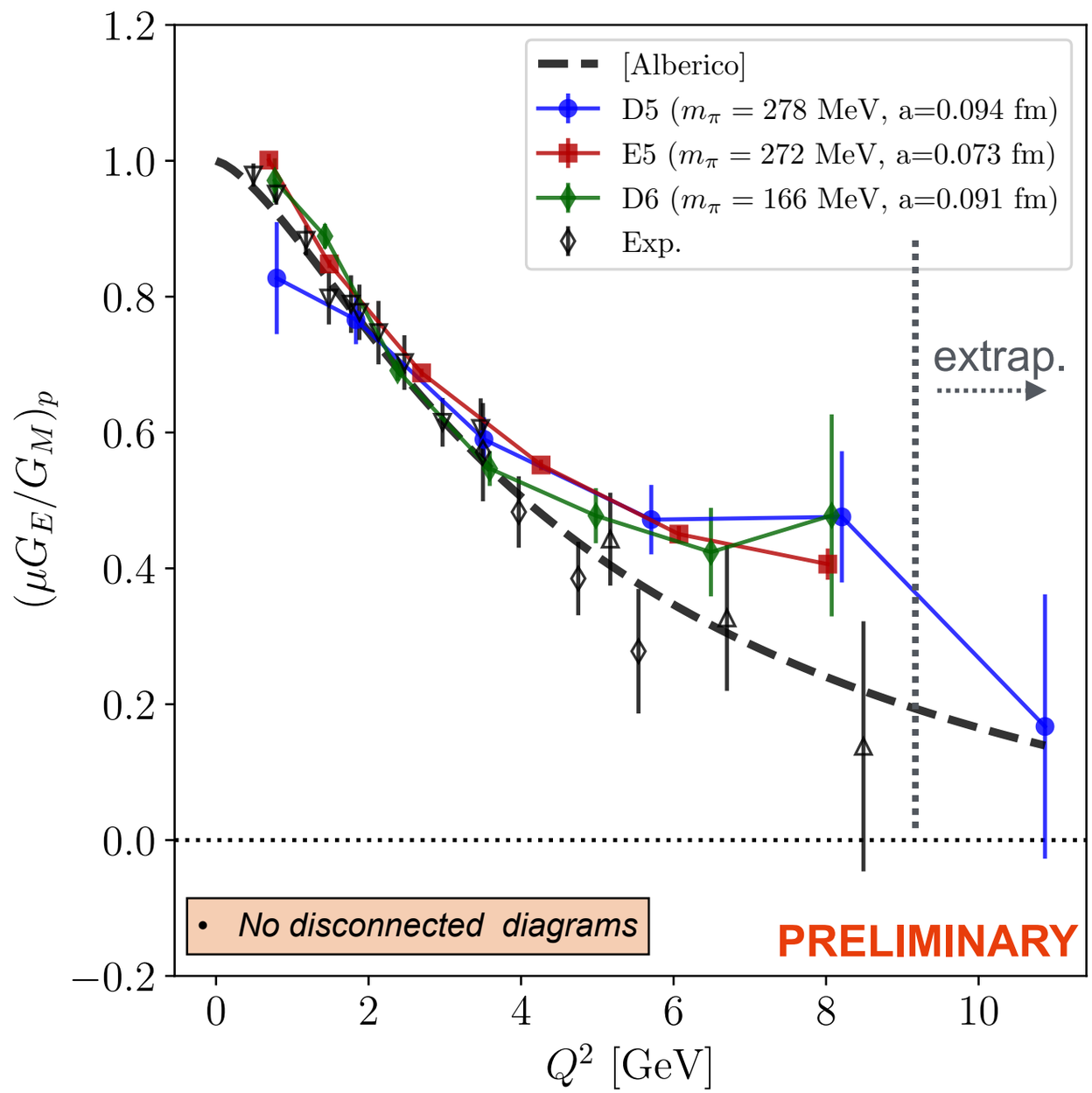


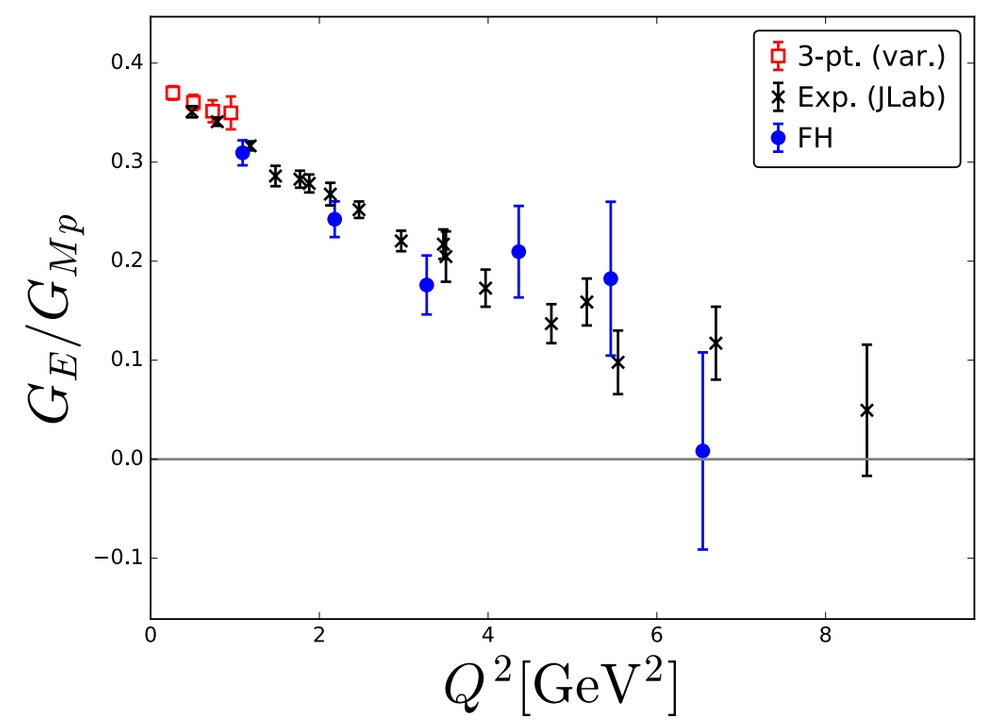


[G.D.Cates, et al, PRL106:252003 (2011)]

G_{Ep}/G_{Mp} Form Factor Ratio (Proton)

- Lattice data: 2-state fits
- Phenomenology curves : [Alberico et al, PRC79:065204 (2008)]
- Comparison to experiments (black points)





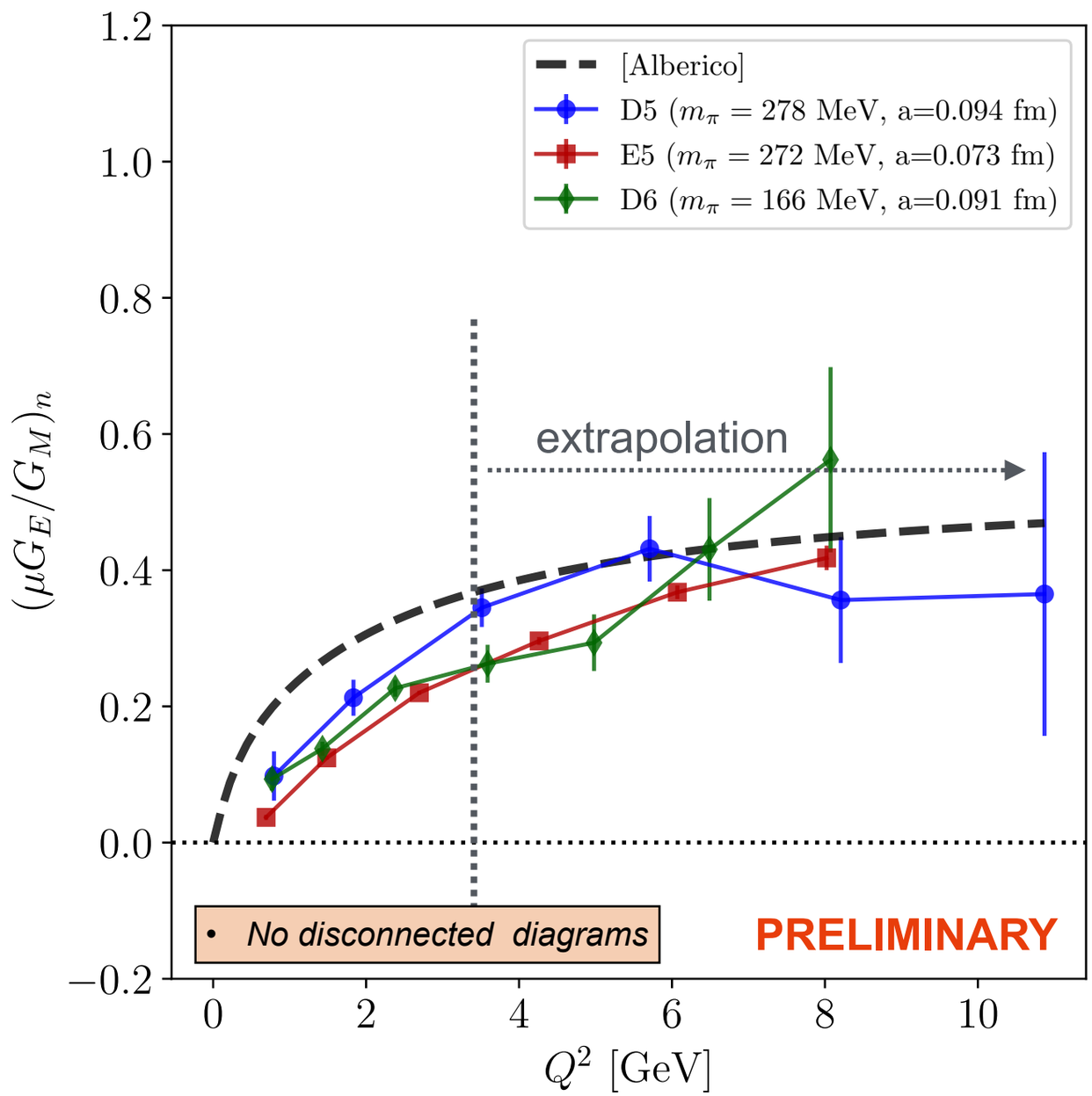
Earler calculation: (a=0.074 fm, mπ=470MeV)

Feynman-Hellman method

[Chambers et al (CSSM), PRD96: 114509]

G_{En}/G_{Mn} Form Factor Ratio (Neutron)

- Lattice data: 2-state fits
- Phenomenology curves : [Alberico et al, PRC79:065204 (2008)]
- Comparison to experiments (black points)



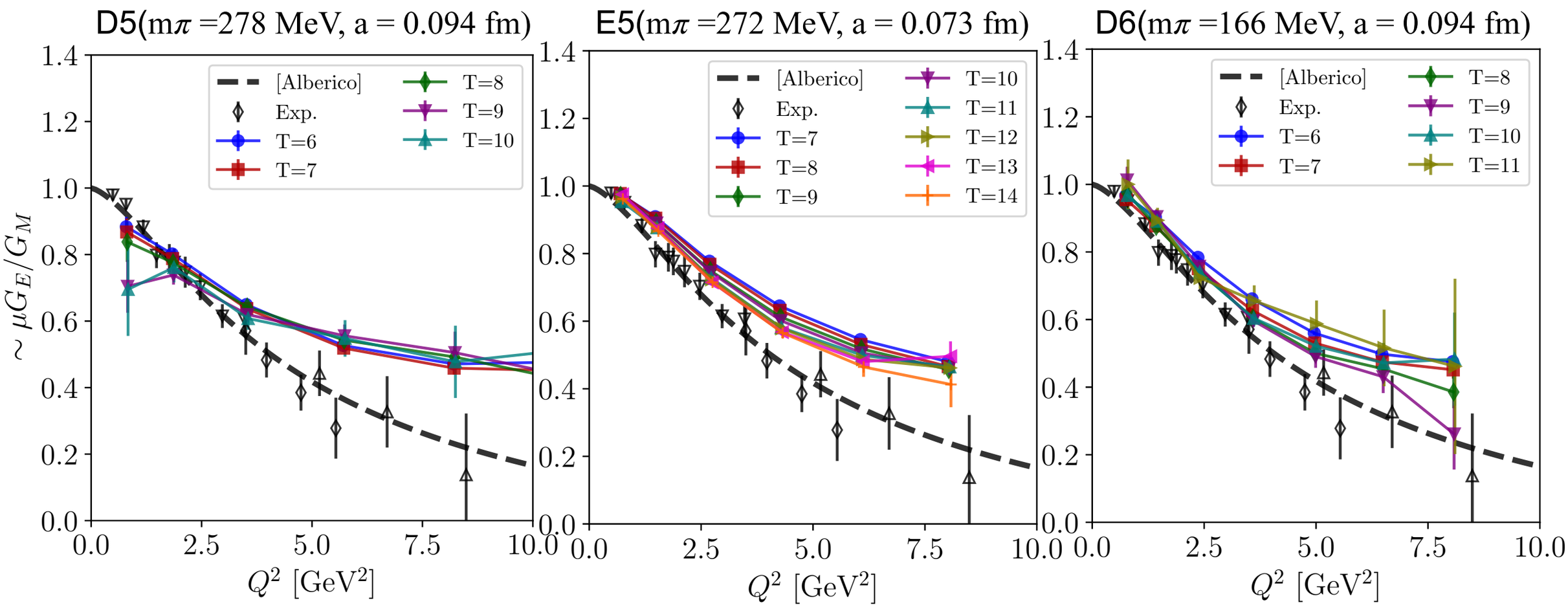
Examine Exc.States in G_{Ep}/G_{Mp}

- Robust estimator from nucleon-current correlators:
  avoid lattice correlators fits to ~Σ exp(–Et)

$$\mathrm{Re}\,\langle p'\hat{x}|J_t|p\hat{x}\rangle \;\propto\; \cosh\frac{\lambda'+\lambda}{2}G_E$$

$$\mathrm{Re}\,\langle p'\hat{x}|J_y|p\hat{x}\rangle \;\propto\; \sinh\frac{\lambda'-\lambda}{2}G_M$$

where $\left(\begin{array}{ll} p^{(\prime)} & = m_N \sinh\lambda^{(\prime)} \\ E^{(\prime)} & = m_N \cosh\lambda^{(\prime)} \end{array}\right)$



$$\left(\frac{\sinh\frac{\lambda'-\lambda}{2}}{\cosh\frac{\lambda'+\lambda}{2}}\right)\frac{\mathrm{Re}\langle N_\uparrow(p'_x,T)\,J_t(T/2)\,\bar{N}_\uparrow(p_x,0)\rangle}{\mathrm{Re}\langle N_\uparrow(p'_x,T)\,J_y(T/2)\,\bar{N}_\uparrow(p_x,0)\rangle} \overset{T\to\infty}{=} G_E/G_M$$

Summary

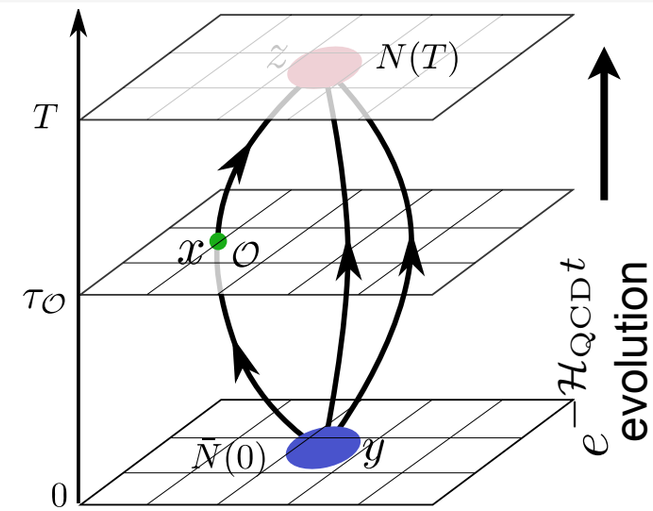
- Preliminary results for high MC-statistics high-momentum form factors
  up to $Q^2 \lesssim 12$ GeV^2, two lattice spacings $a \gtrsim 0.07$ fm, two pion masses $m\pi \gtrsim 170$ MeV
  (No quark-disconnected contributions yet)

- Form factor results overshoot experimental data x(2 ... 2.5) ;
  G_E/G_M ratios in qualitative agreement
  ~~*Discretization?*~~
  ~~*Excited states?*~~
  Non-physical quark masses?
  Quark-Disconnected contributions?

- Comparison to experiment important to validate lattice methods
  for computing relativistic nucleon matrix elements
  Impact on lattice methodology for TMDs, PDFs, DAs calculation

BACKUP

Nucleon Form Factors: Lattice Recipe



Quark lines = $(\not{D}+m)^{-1}\cdot\psi$

- Generate lattice ensemble

$$P(A_\mu) \propto e^{-S[A_\mu]} \prod_q \det(\not{D}+m_q)$$

- Compute nucleon correlation functions

2-point $$\langle N(\vec{p}\,', t_{\rm sep})\, \bar{N}(0)\rangle = \sum_y e^{-i\vec{p}\,'\vec{x}} \langle N_{t_{\rm sep},\vec{y}}\, \bar{N}(0,\vec{0})\rangle$$

3-point $$\langle N(\vec{p}\,', t_{\rm sep})\, \mathcal{O}(\vec{q},\tau)\, \bar{N}(0)\rangle = \sum_{y,z} e^{-i\vec{p}\,'\vec{x}+i\vec{q}\vec{z}} \langle N_{t_{\rm sep},\vec{y}}\, \mathcal{O}_{z,\tau}\, \bar{N}(0,\vec{0})\rangle$$

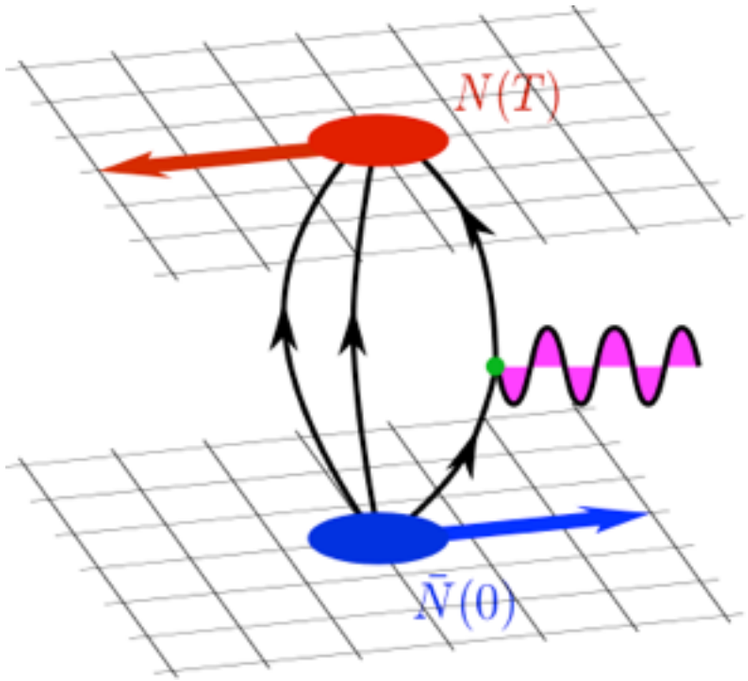
- Extract ground-state matrix elements (2-state fits typical)

$$\begin{aligned}\langle N(p', t_{\rm sep})\, \mathcal{O}(\tau)\, \bar{N}(p,0)\rangle \sim e^{-E_0'(t_{\rm tsep}-\tau)-E_0\tau} \Big[& \langle 0(p')|\mathcal{O}|0(p)\rangle \\ & + \langle 0(p')|\mathcal{O}|1(p)\rangle e^{-\Delta E_{01}\tau} + \langle 1(p')|\mathcal{O}|0(p)\rangle e^{-\Delta E_{01}'(t_{\rm sep}-\tau)} \\ & + \langle 1(p')|\mathcal{O}|1(p)\rangle e^{-\Delta E_{01}\tau - \Delta E_{01}'(t_{\rm sep}-\tau)} \Big]\end{aligned}$$

- Reduce m.e.'s to form factors: fit over momentum combinations, polarization, etc

$$\langle p', \sigma'|J^\mu|p,\sigma\rangle = [\bar{u}'\gamma^\mu u]F_1 + [\bar{u}'\frac{i\sigma^{\mu\nu}q_\nu}{2m_N}u]F_2$$

Accessing Large Q^2 : Breit Frame on a Lattice



Minimize $E_{in,out}$ for target Q^2 :

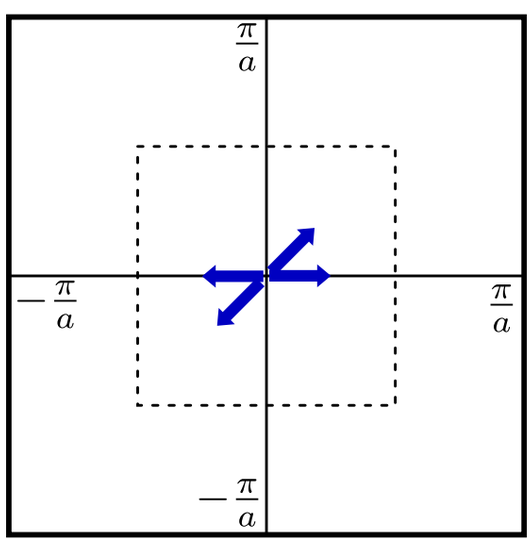
$$Q^2 = (\vec{p}_{in} - \vec{p}_{out})^2 - (E_{in} - E_{out})^2$$

Back-to-back

$$Q^2 = 4\vec{p}^2$$

For $(Q^2)_{max}$ = *10 GeV²*

$$|\vec{p}| = \frac{1}{2}\sqrt{Q^2_{\max}} \approx 1.6 \text{ GeV} \qquad (E_N \approx 1.9 \text{ GeV})$$



lattice kinematics for $Q^2 \approx 10$ GeV²

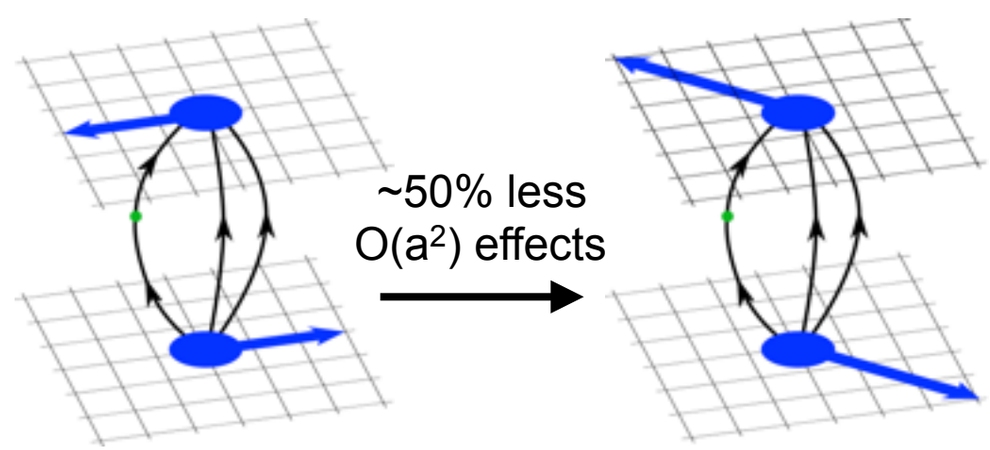
Nucleon momentum ~ Brillouin zone
⇒ unknown distortion of effective nucleon Dirac eqn.

$$\langle N\bar{N}\rangle^{-1}(p) \stackrel{?}{=} -i\not{p}^{\text{lat}} + m_N$$

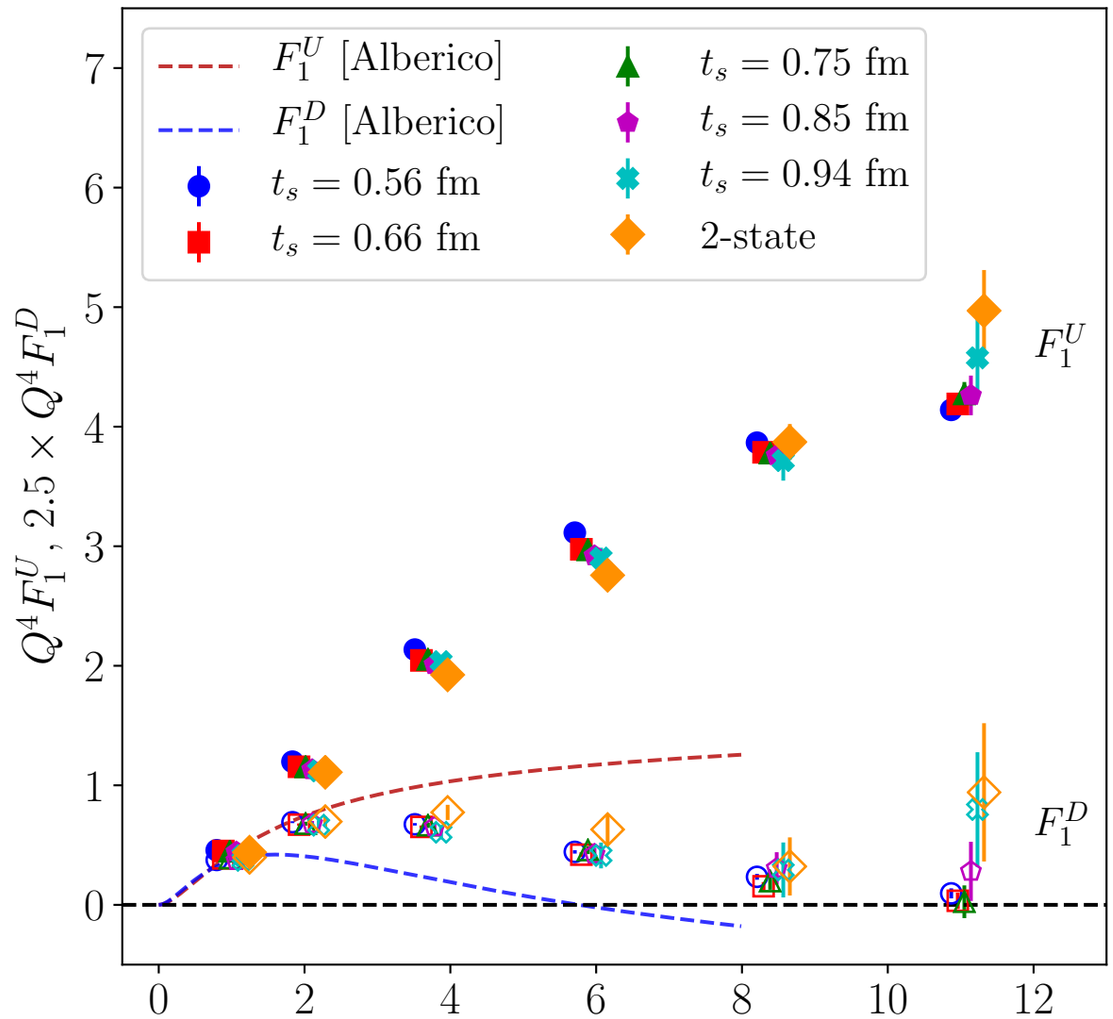
$$p_\mu^{\text{lat}} = k_\mu + O(k^3)$$

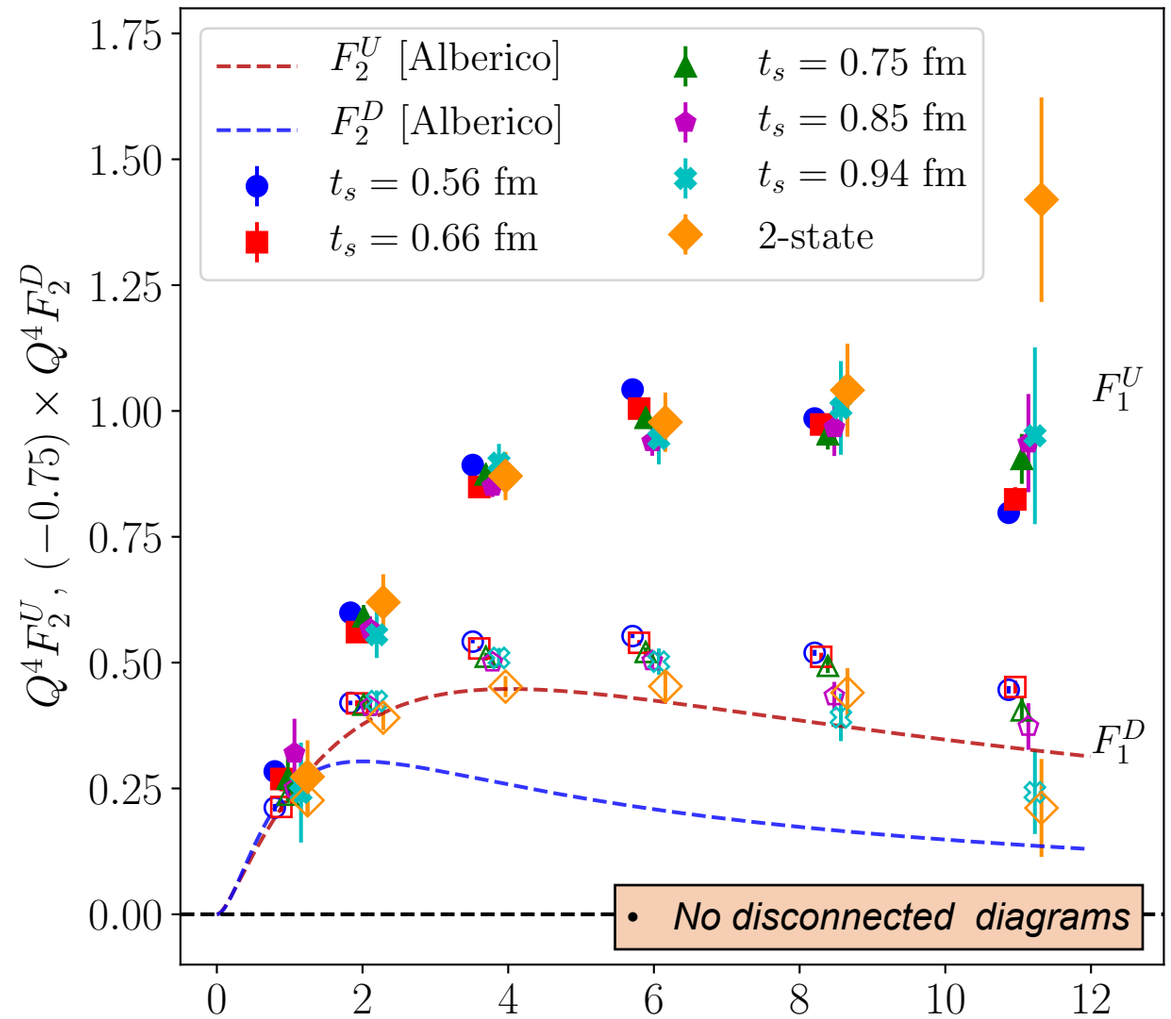
⇒ **addl**. O(a^2) corrections

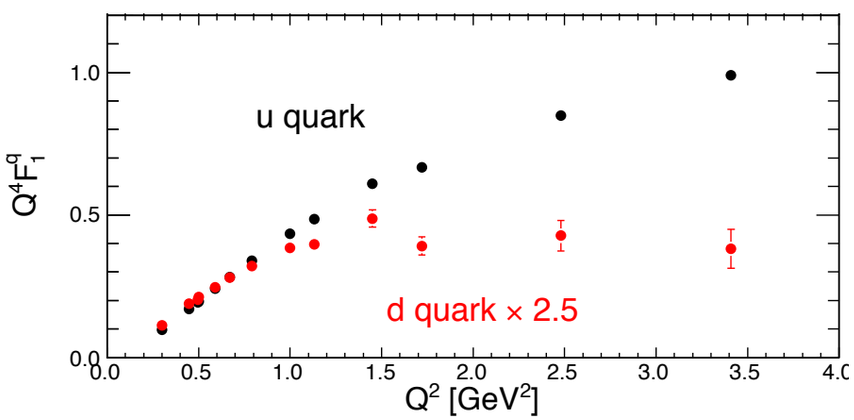
from lattice nucleon polarization

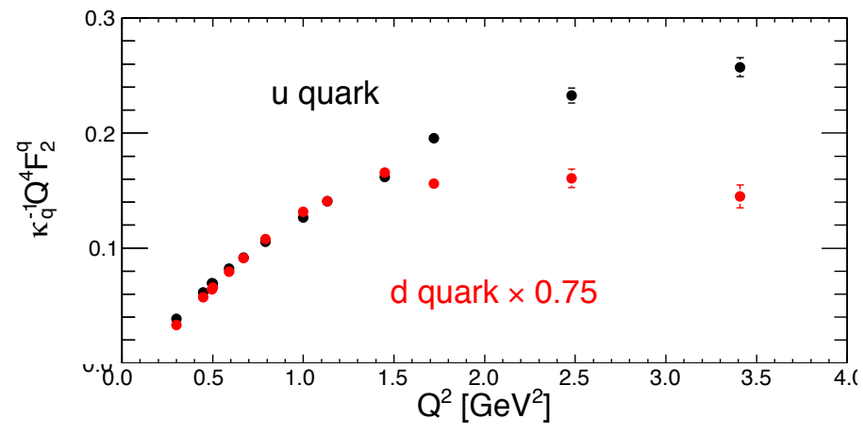


Results* : Flavor Decomposition



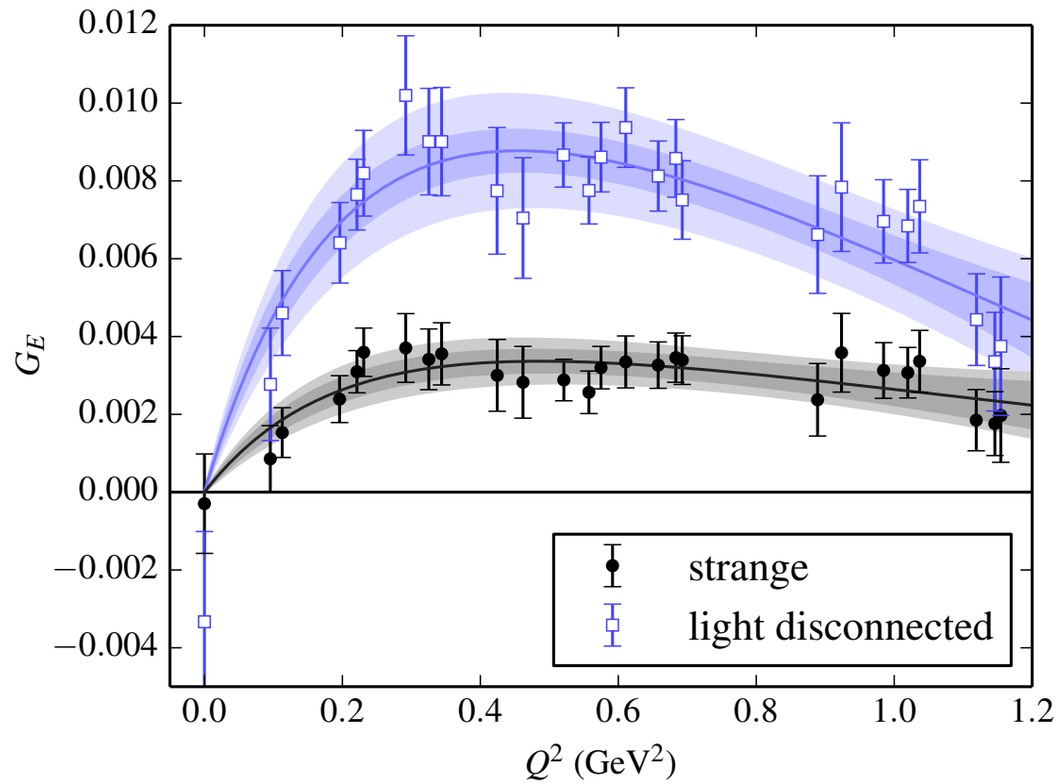






- Reproduce qual.features of flavor dependence [G.D.Cates, et al, PRL106:252003(2011)]
- Larger form factors: nucleon (+ exc.states?) on a lattice is more "compact"

Prior work: Disc.Light & Strange Quark F.F's



N_f=2+1 dynamical fermions, $m_\pi \approx$ 320 MeV
(the "coarse" JLab Clover ensemble)

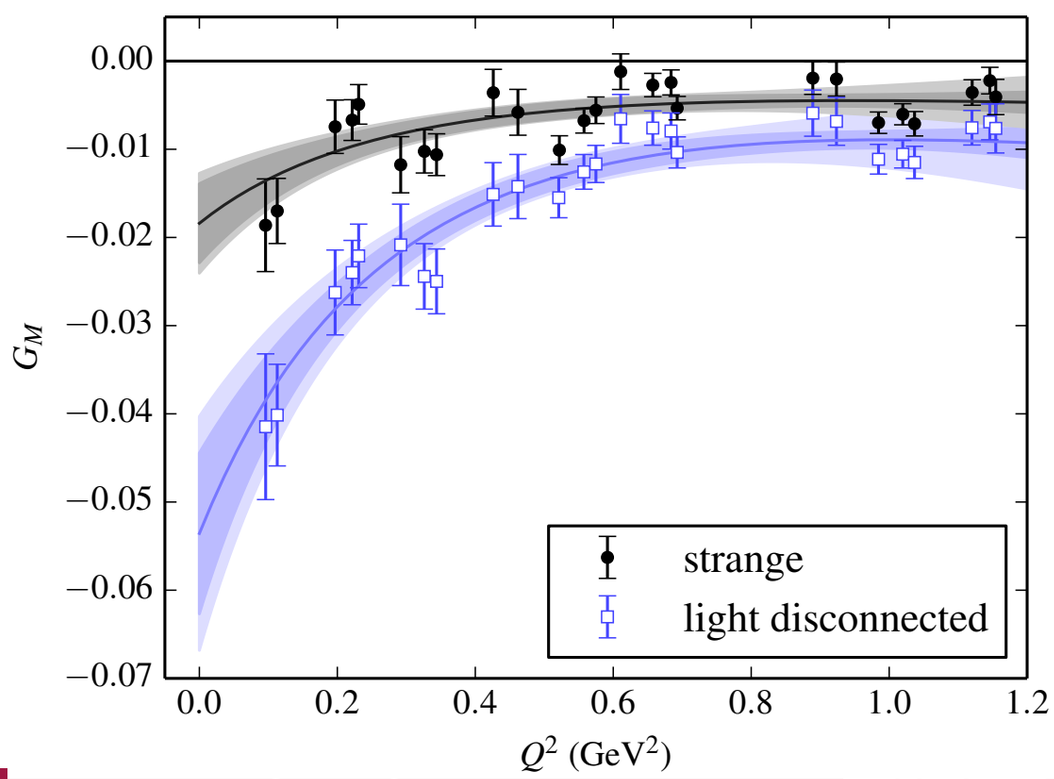
$$|(G_E^{u/d})_{\rm disc}| \lesssim 0.010 \text{ of } |(G_E^{u-d})_{\rm conn}|$$

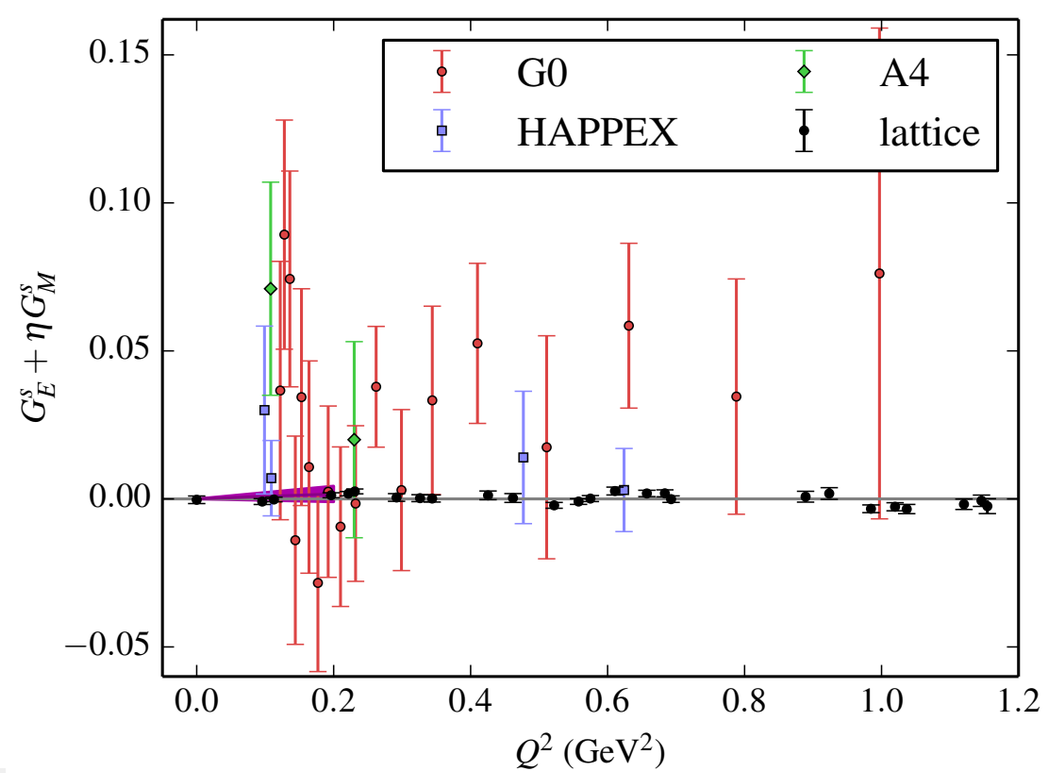
$$|(G_E^{s})_{\rm disc}| \lesssim 0.005 \text{ of } |(G_E^{u-d})_{\rm conn}|$$

$$|(G_M^{u/d})_{\rm disc}| \lesssim 0.015 \text{ of } |(G_M^{u-d})_{\rm conn}|$$

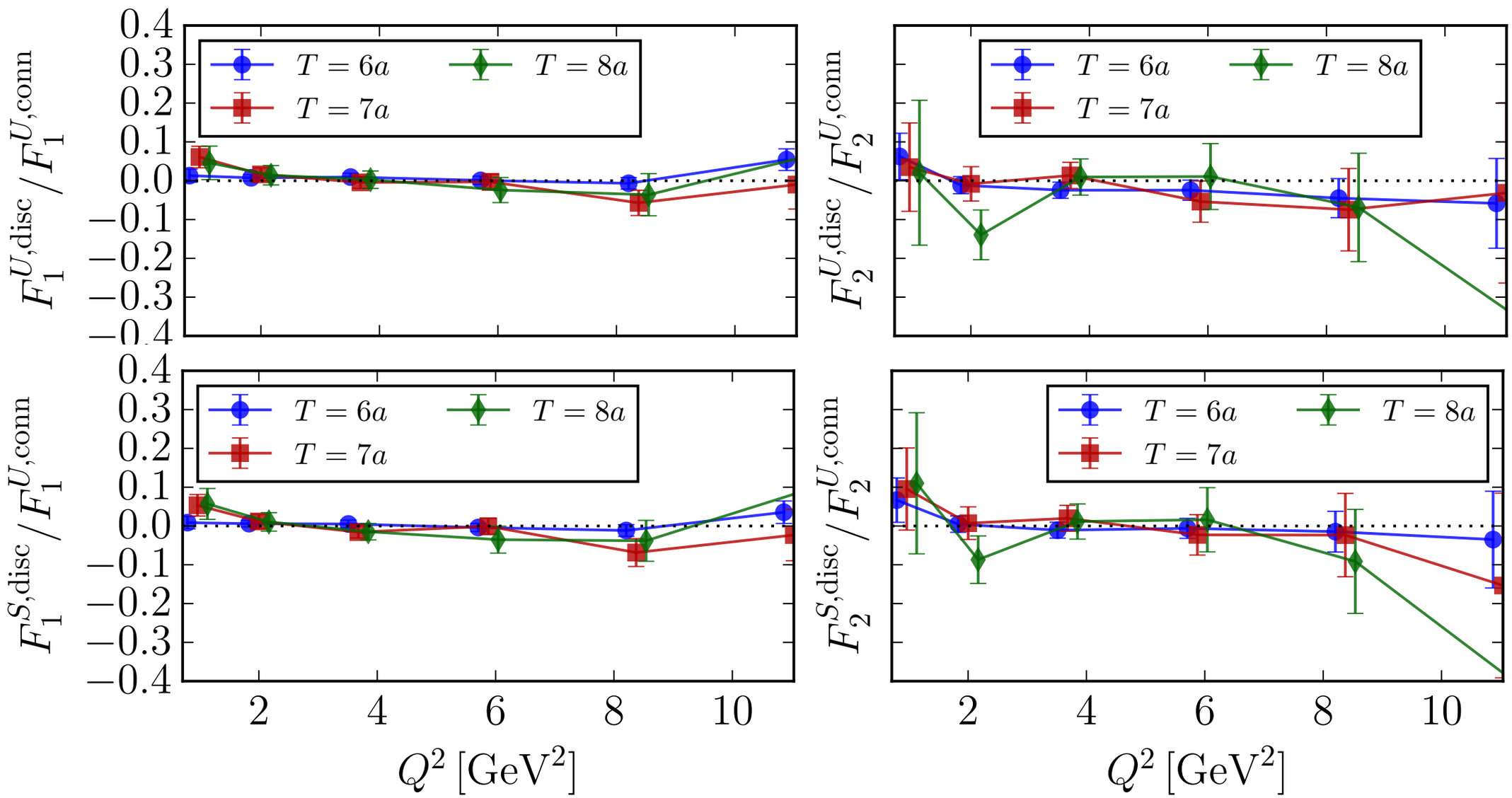
$$|(G_M^{s})_{\rm disc}| \lesssim 0.005 \text{ of } |(G_M^{u-d})_{\rm conn}|$$

[J. Green, S. Meinel, S.S. et al; PRD92:031501 (2015)]





Disconnected Light, Strange vs. Connected



- D5 ensemble(mπ=280 MeV, a=0.094 fm)
- s-, disconnected *u,d*- contributions are still small at high Q^2 up to ≲ 10 GeV^2

$$|F_1^s| \lesssim |\left(F_1^{u/d}\right)_{\mathrm{disc}}| \lesssim \ 10\% \text{ of } |F_2^{u,d}|$$

$$|F_2^s| \lesssim |\left(F_2^{u/d}\right)_{\mathrm{disc}}| \lesssim \ 20\% \text{ of } |F_2^{u,d}|$$