



University  
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on High Energy Physics  
Bologna (Italy)

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13 07 2022



# Latest Hermes results on azimuthal single- and double-spin asymmetries in semi-inclusive deep-inelastic scattering by transversely polarized protons

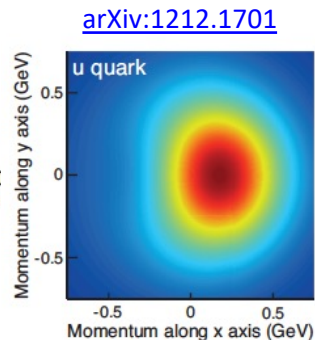
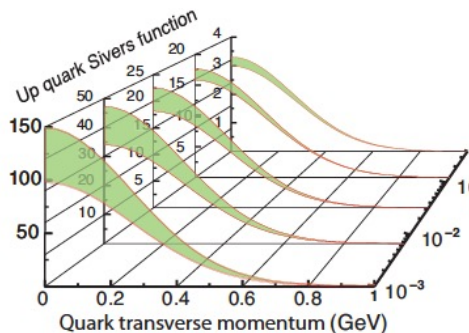
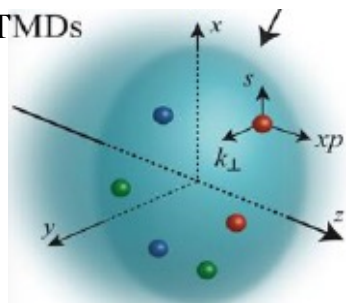
Luciano L. Pappalardo (for the HERMES Collaboration)

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# TMDs in SIDIS

TMDs

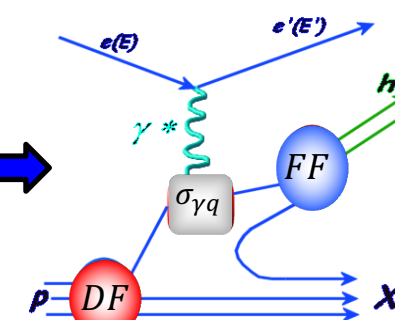
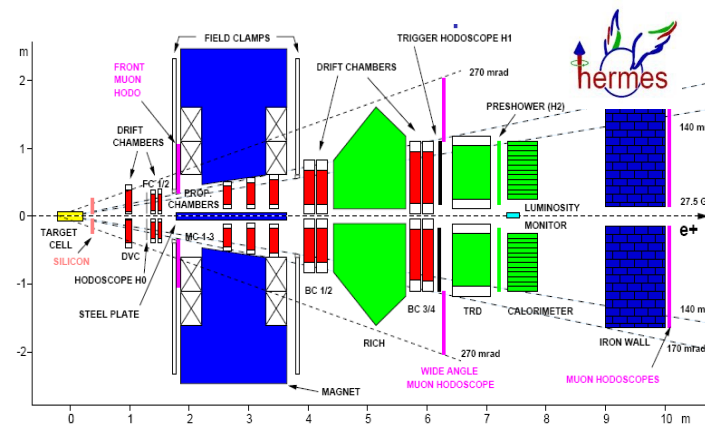


- encode flavour-dependent correlations between  $p_T$  and the spin orientation of the parent hadron or of the quark itself
- 3D description of nucleon structure in momentum space ( $\rightarrow$  nucleon tomography)

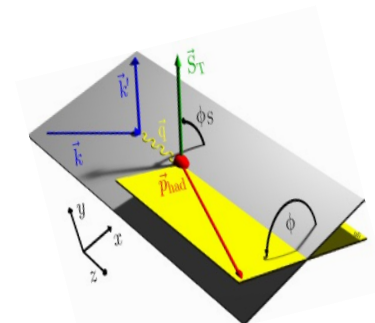
quark polarisation

nucleon polarisation	quark polarisation		
	U	L	T
	<b>U</b> $f_1$ number density PRD 87 (2013) 074029		$h_1^\perp$ Boer-Mulders PRD 87 (2013) 012010
		<b>L</b> $g_1$ helicity PRD 75 (2007) 012007	$h_{1L}^\perp$ worm-gear PLB 562 (2003) 182 PRL 84 (2000) 4047
T	$f_{1T}^\perp$ Sivers PRL 94 (2005) 012002 PRL 103 (2009) 152002	$g_{1T}$ worm-gear released	$h_{1T}^\perp$ transversity PRL 94 (2005) 012002 PLB 693 (2010) 11 released

## Semi-inclusive DIS processes (SIDIS)



$$\begin{aligned}
 & \frac{d\sigma}{dx_B dy d\psi dz_h d\phi_h dP_{h\perp}^2} \\
 &= \frac{\alpha^2}{x_B y Q^2} \frac{y^2}{2(1-\varepsilon)} \left( 1 + \frac{\gamma^2}{2x_B} \right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\
 &+ \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \\
 &+ S_{\parallel} \left[ \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \\
 &+ S_{\parallel} \lambda_e \left[ \sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\
 &+ |S_{\perp}| \left[ \sin(\phi_h - \phi_S) \left( F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
 &+ \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\
 &+ \left. \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right] \\
 &+ |S_{\perp}| \lambda_e \left[ \sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right. \\
 &+ \left. \left. \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\}, \quad \text{Bacchetta et al JHEP 08}
 \end{aligned}$$



# The "Hermes TMDs Bible"

PREPARED FOR SUBMISSION TO JHEP  
DESY REPORT 20-119

Azimuthal single- and double-spin asymmetries in semi-inclusive deep-inelastic lepton scattering by transversely polarized protons

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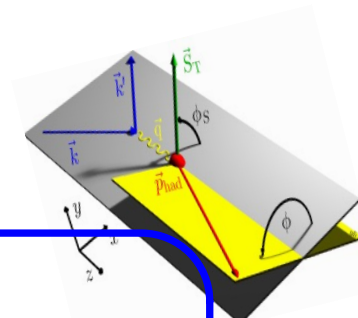
<sup>13</sup>II. Physikalisches Institut, Justus-Liebig Universität Gießen, 35392 Gießen, Germany

<sup>a</sup>Deceased.

[JHEP 12 (2020) 010] [arXiv:2007.07755v1](https://arxiv.org/abs/2007.07755v1)

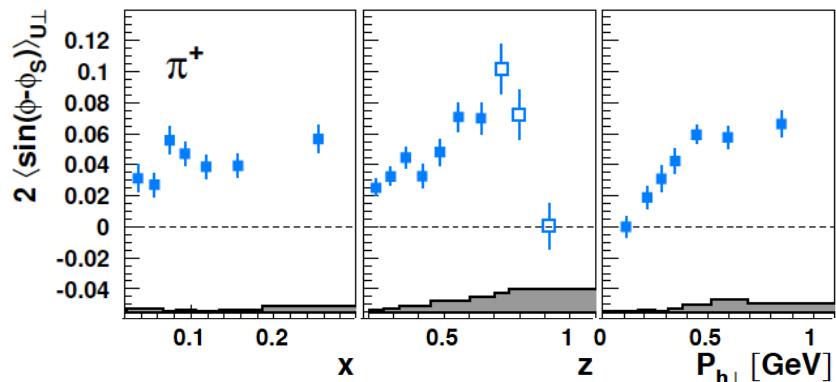
- Compendium of HERMES TMDs results obtained with transv. Pol. H target (84 pages!)
- 10 azimuthal modulations ( $6 A_{U\perp} + 4 A_{L\perp}$ ) for 7 hadron types
- Many advances w.r.t previously published analyses
  - 3D binning in  $x, z, P_{h\perp}$  (before only 1D)
  - $p/\bar{p}$  asymmetries (in addition to  $\pi^\pm, \pi^0, K^\pm$ )
  - 1D binning optimized and extended to the high- $z$  ("semi-exclusive") region ( $0.7 < z < 1.2$ )
  - The  $x$  range is extended up to 0.6 (before was up to 0.4)
  - ...and more

$$\begin{aligned} & \frac{d\sigma}{dx_B dy d\psi dz_h d\phi_h dP_{h\perp}^2} \\ &= \frac{\alpha^2}{x_B y Q^2} \frac{y^2}{2(1-\varepsilon)} \left( 1 + \frac{\gamma^2}{2x_B} \right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\ &+ \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \\ &+ S_{\parallel} \left[ \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \\ &+ S_{\parallel} \lambda_e \left[ \sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\ &+ |S_{\perp}| \left[ \sin(\phi_h - \phi_S) \left( F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\ &+ \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\ &+ \left. \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right] \\ &+ |S_{\perp}| \lambda_e \left[ \sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right. \\ &+ \left. \left. \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\}, \quad \text{Bacchetta et al JHEP 08} \end{aligned}$$

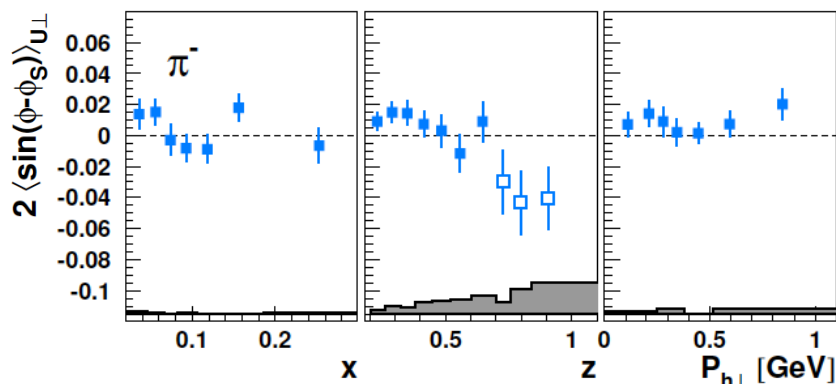


# Selected results

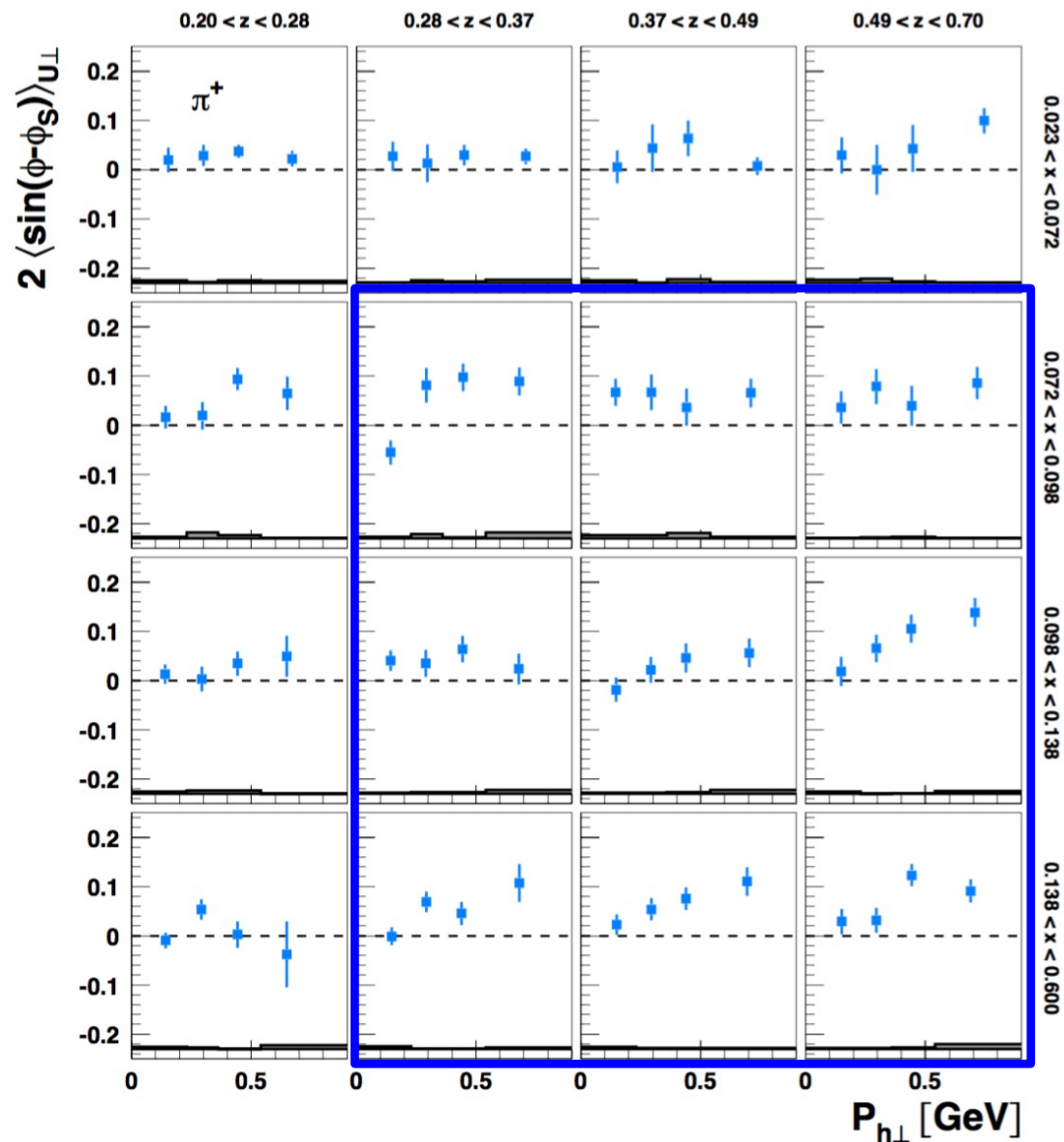
# Sivers amplitudes: pions results



- large positive amplitude  $\rightarrow$  clear evidence of non-zero  $f_{1T}^{\perp,u}$
- signal rises with  $x$ ,  $z$  and  $P_{h\perp}$  in SIDIS region ( $0.2 < z < 0.7$ )
- More informative 3D projections confirm and further detail the rise of the amplitude at large  $x$ ,  $z$  and  $P_{h\perp}$

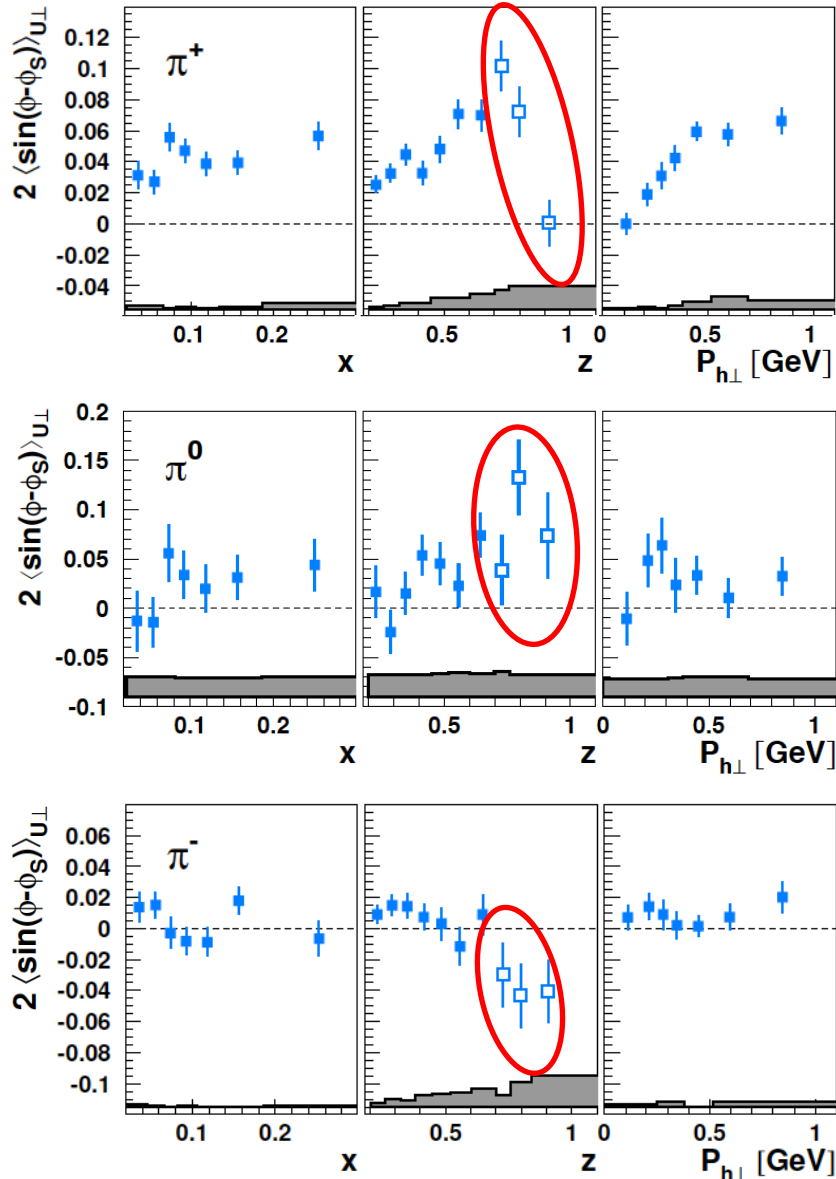


Vanishing due to the cancellation of the opposite Sivers effect for  $u$  and  $d$  quarks

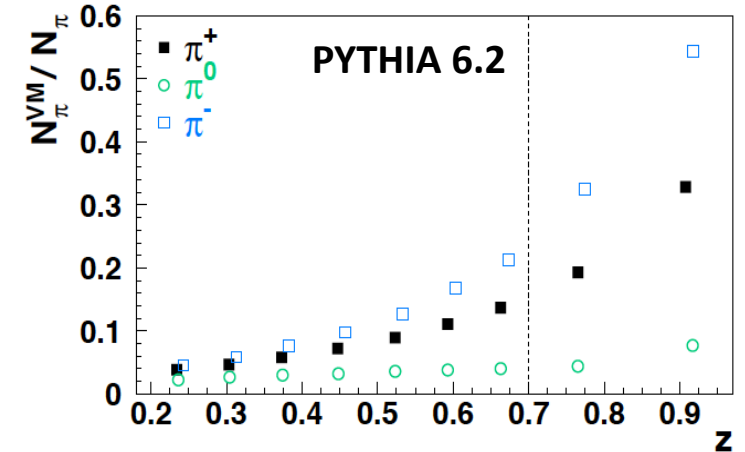




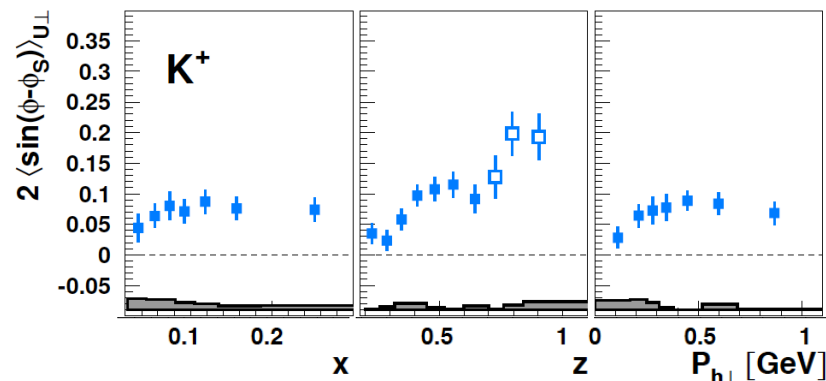
# Sivers amplitudes: pions results



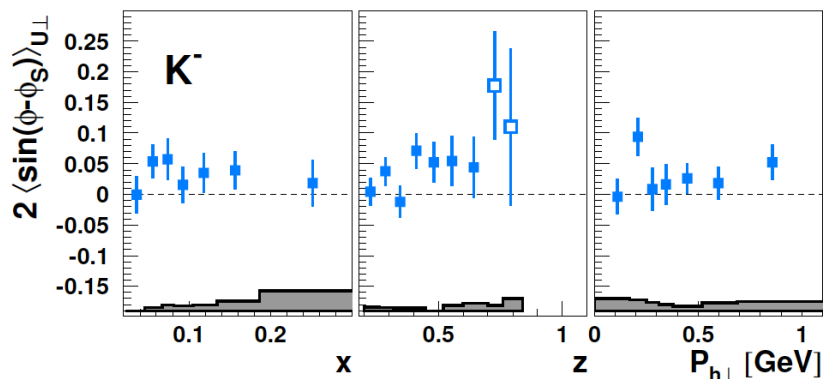
- Sudden drop at large- $z$  ( $> 0.7$ ) reveals a change of mechanism in this **semi-exclusive region**
- Contributions from decays of exclusively produced  $\rho^0$  into  $\pi^+\pi^-$  are large in this region!
- intermediate size between those of  $\pi^+$  and  $\pi^-$  reflects isospin symmetry at the amplitude level
- $\pi^0$  amplitude is much less susceptible to VM decays and no sudden change is observed at large  $z \rightarrow$  observed positive signal cannot be attributed solely to contributions from VM
- An alternative (concurrent?) explanation: at large  $z$ , favored fragmentation ( $d \rightarrow \pi^-$ ) prevails over the disfavored one ( $u \rightarrow \pi^-$ )  $\rightarrow$  no cancellation and a non-zero amplitude opposite to that of  $\pi^+$  is observed.



# Sivers amplitudes: Kaons results

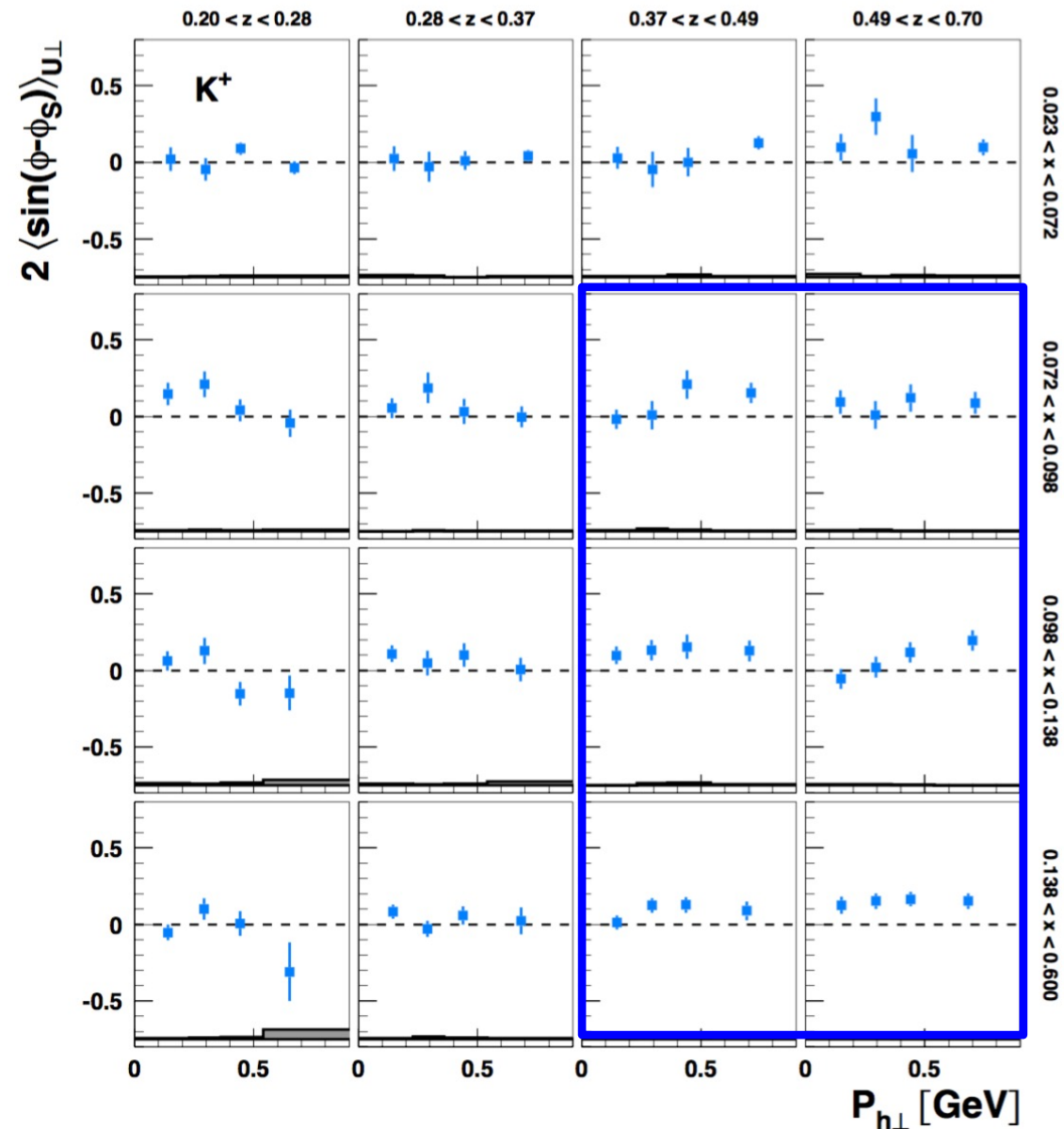


Large positive amplitude, similar kinematic dep. of  $\pi^+$

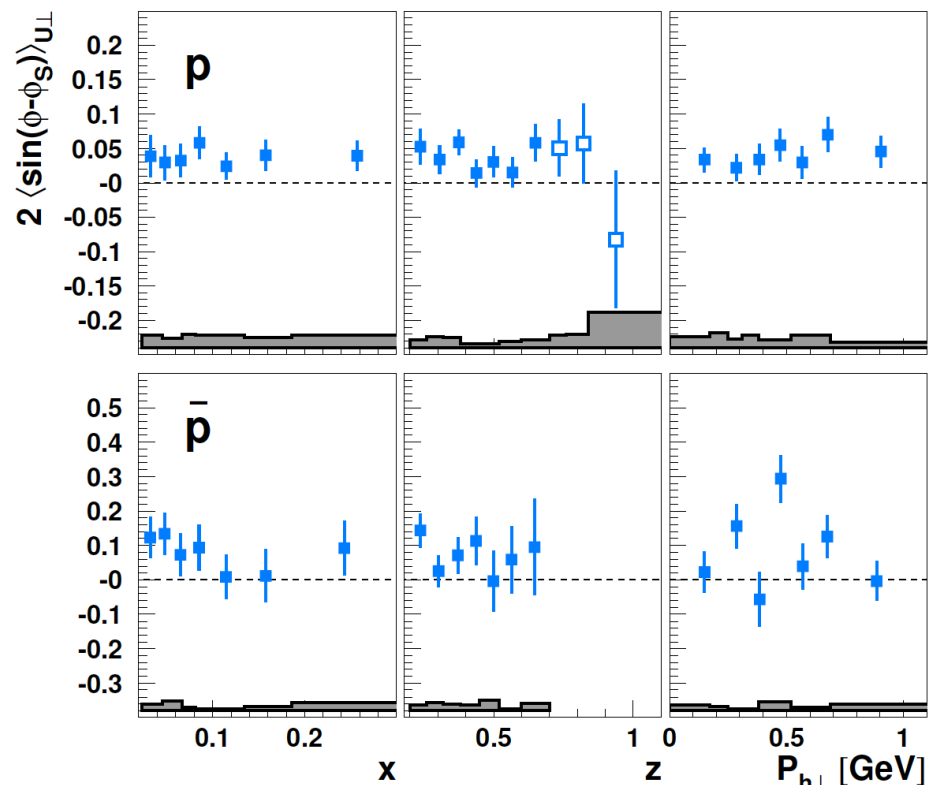


Positive amplitude, different than  $\pi^-$

$K^-$  is a pure sea object with no valence quarks in common with target proton



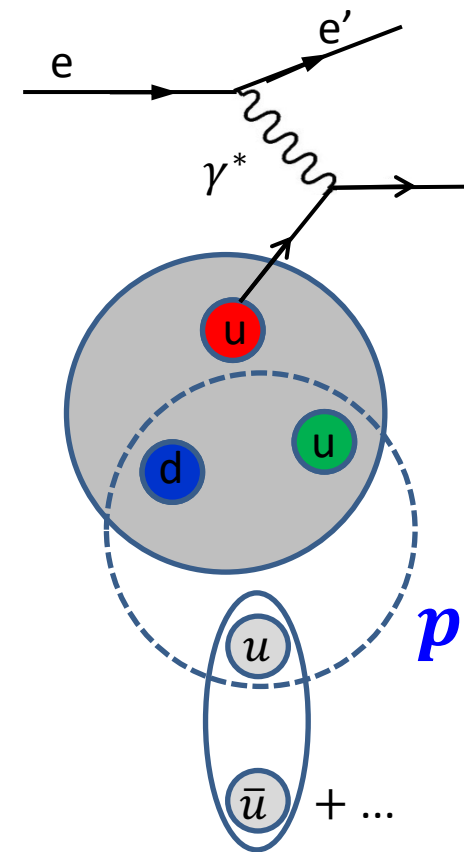
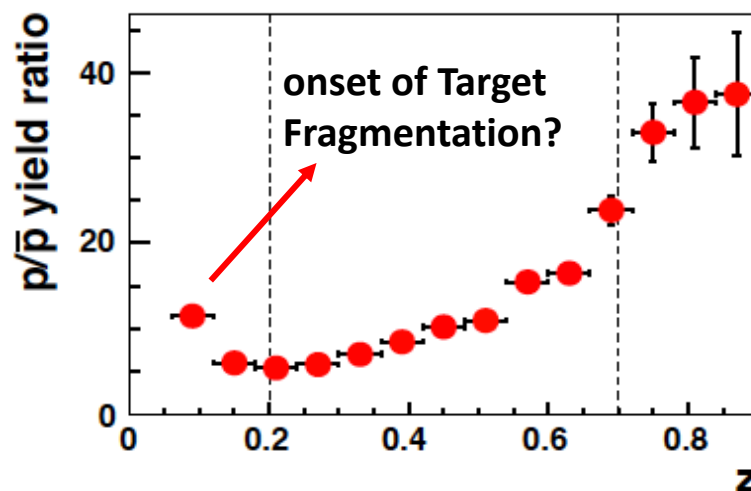
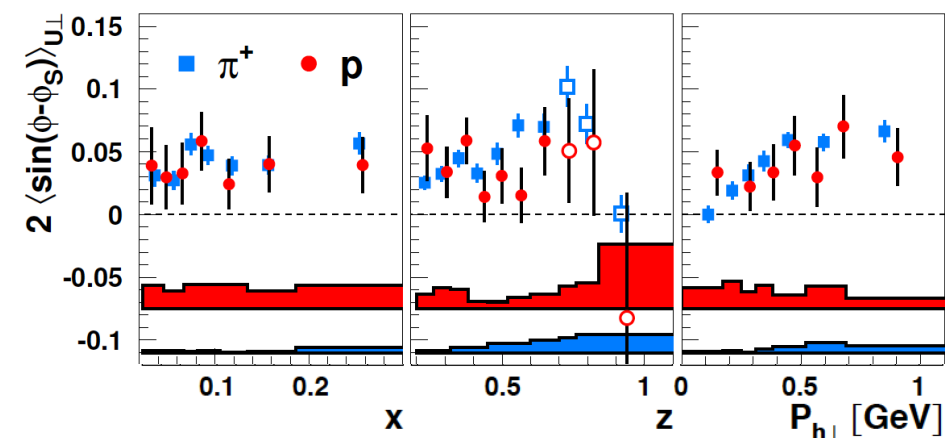
# Sivers amplitudes: protons results



First measurement of Sivers asymmetries for  $p, \bar{p}$  in SIDIS

Both amplitudes are non-zero and positive

Proton production is particularly susceptible to receive contributions from **Target Fragmentation**

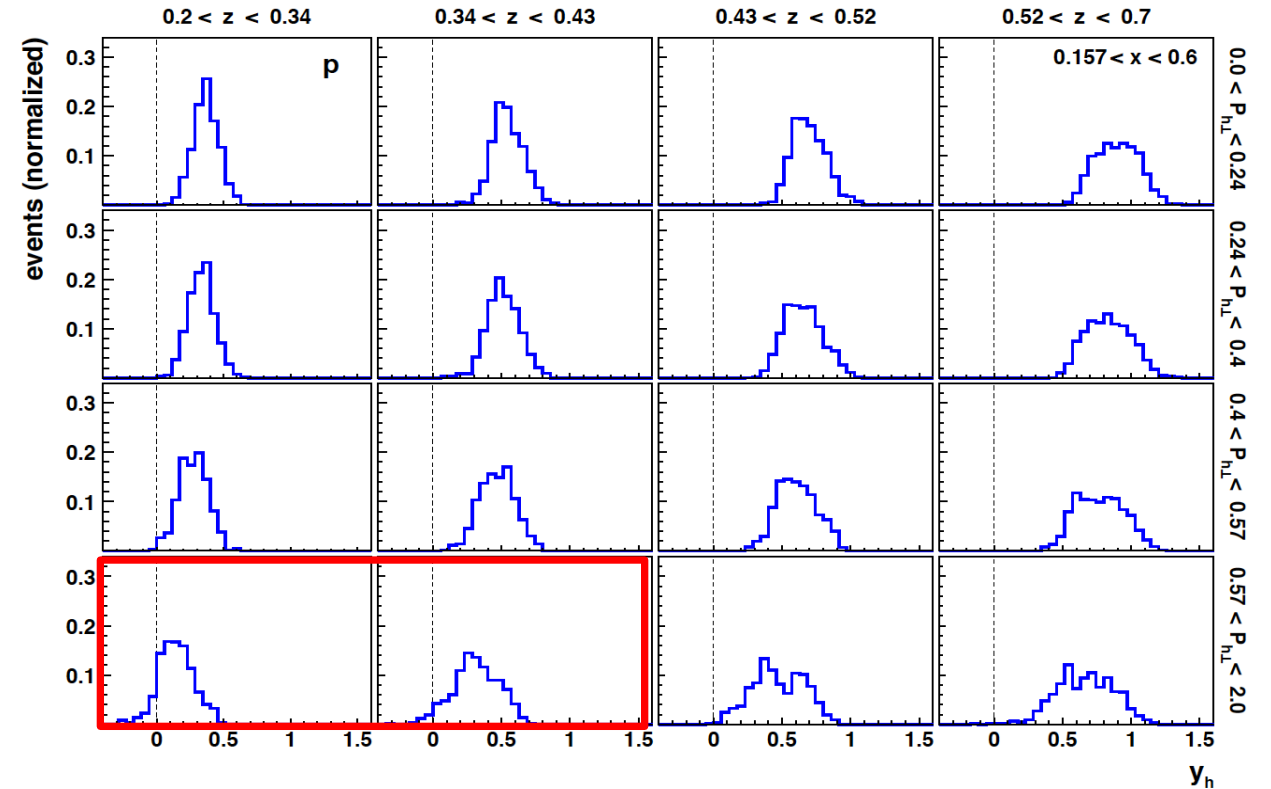
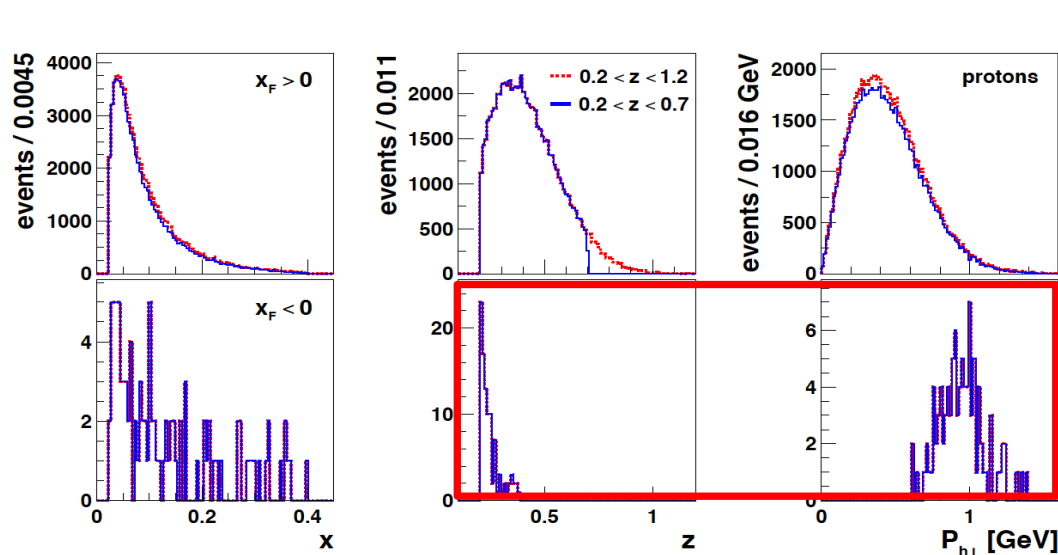


(low  $z$ , high  $P_{h\perp}$ )



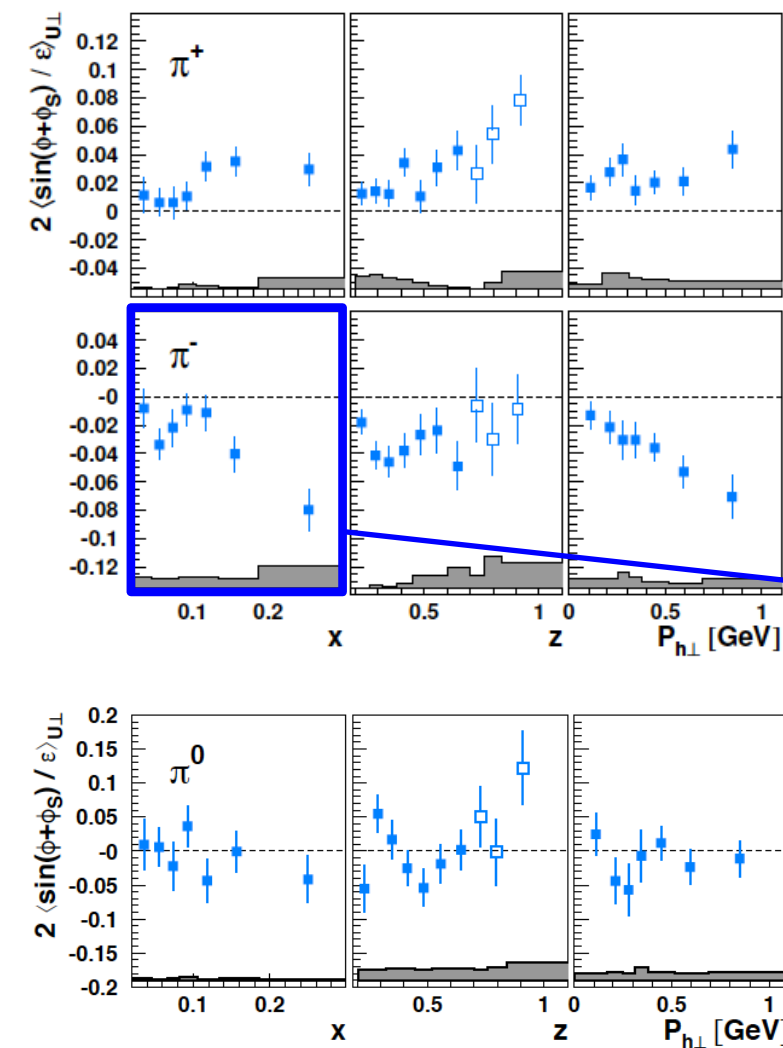
# Sivers amplitudes: protons results (CFR vs. TFR)

- No generally-accepted recipe exists
- positive values of  $x_F$  and rapidity ( $y_h$ ) are typically associated with hadrons produced from the struck quark (CFR)
- negative values point at target fragmentation (TFR)



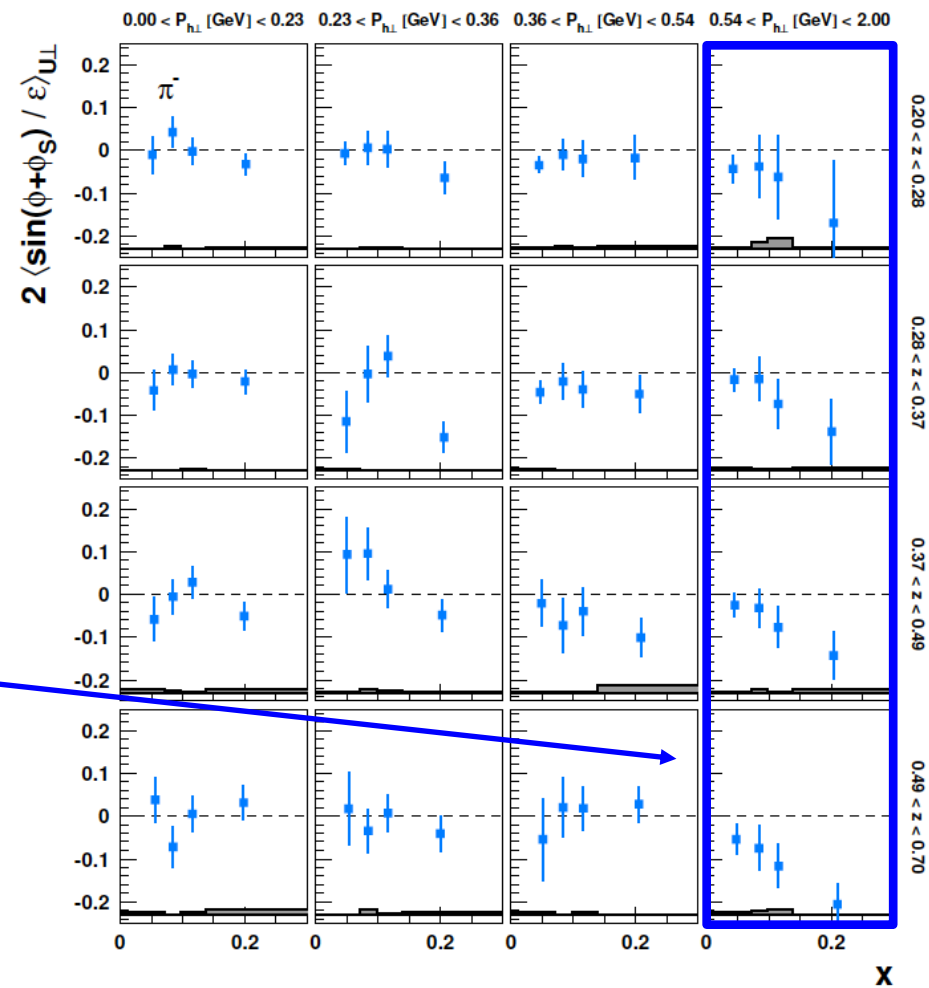
At the selected kinematics the vast majority of protons are compatible with being produced in CFR

# Collins amplitudes: SFA pion results



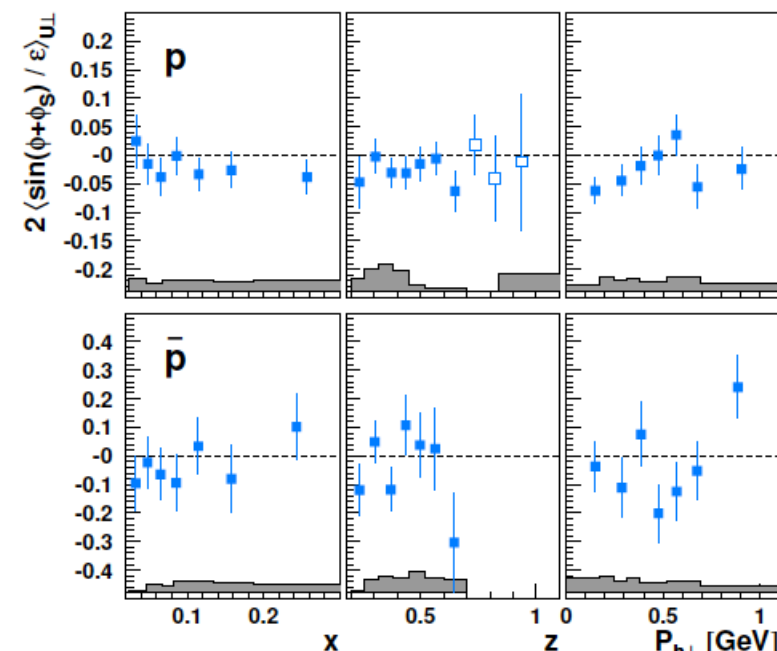
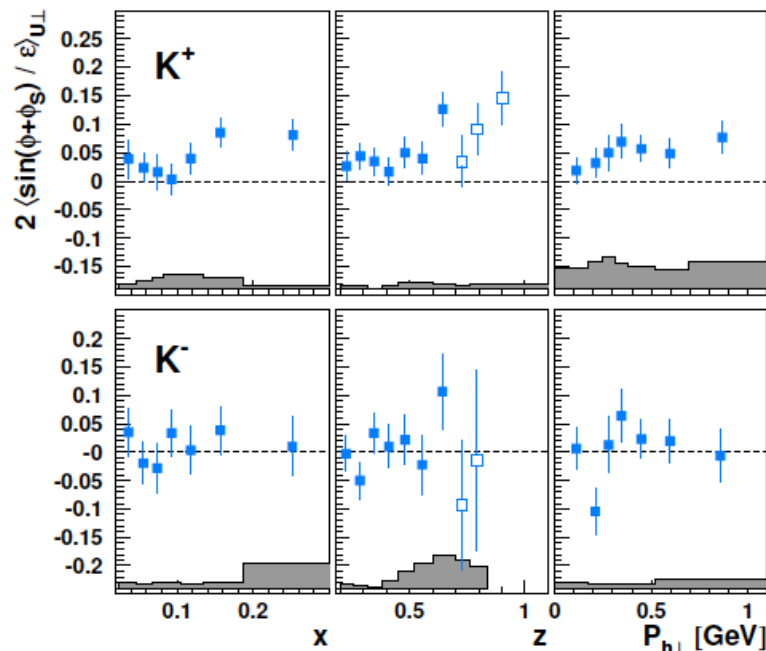
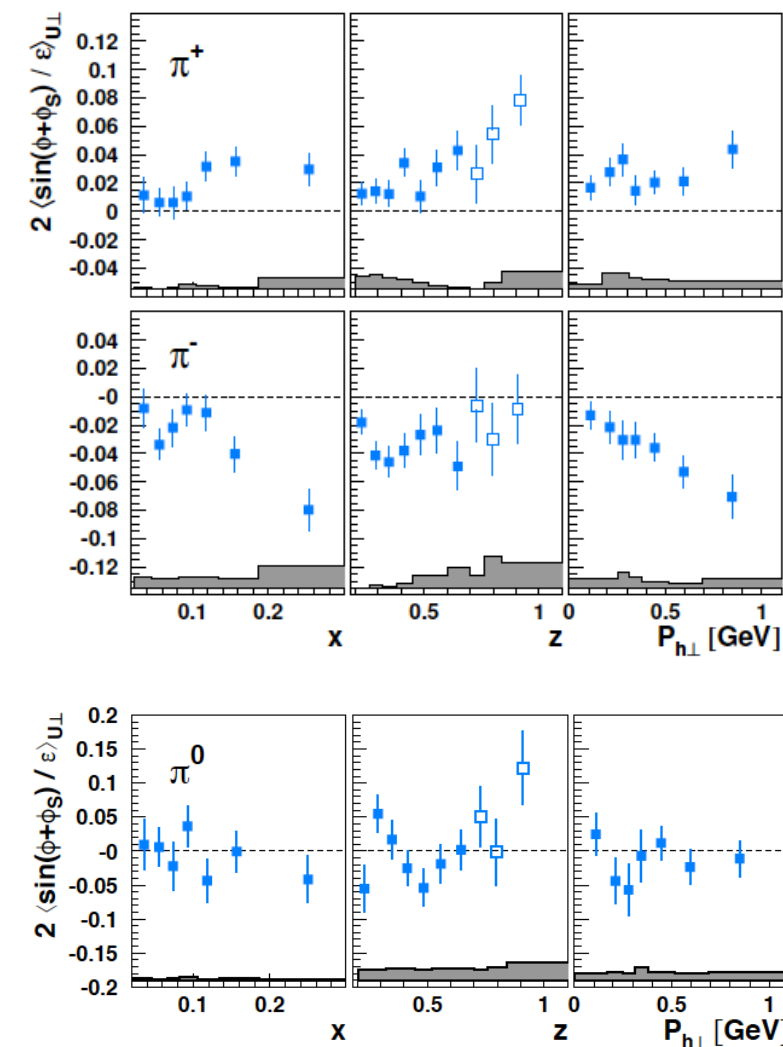
- Large and opposite amplitudes
- Clear evidence of non-zero transversity
- Negative  $\pi^-$  amplitude points to large disfavoured ( $u \rightarrow \pi^-$ ) Collins FF opposite to the favoured one ( $d \rightarrow \pi^-$ )

- $\approx 0$ : intermediate between  $\pi^+$  and  $\pi^-$



- 3D projections confirm and further detail the rise of the amplitude at large  $x$  and  $P_{h\perp}$

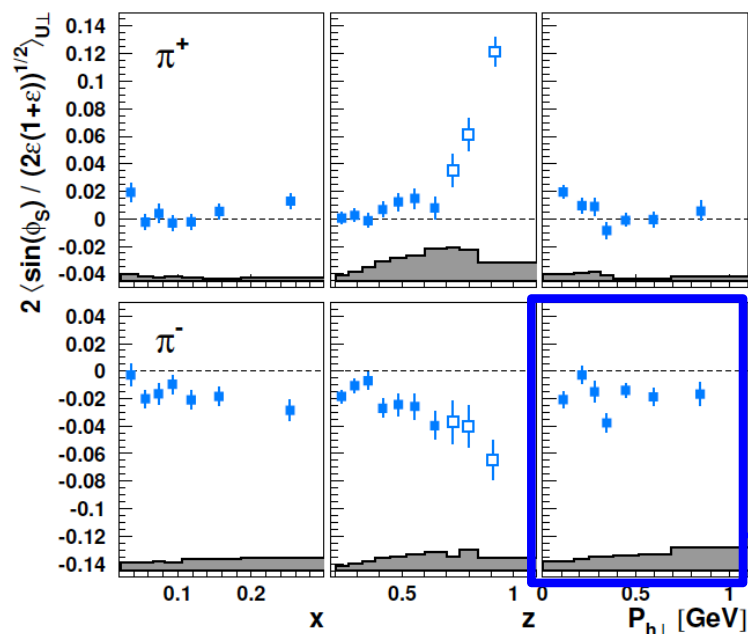
# Collins amplitudes: all SFA 1D results



- $K^+$  exhibits a very similar kinematic dependence as  $\pi^+$ , but amplitude is twice as large!
- $K^- \approx 0$ : only disfavored and opposite ( $u \rightarrow K^-$ ,  $d \rightarrow K^-$ ) fragmentation mechanisms can contribute

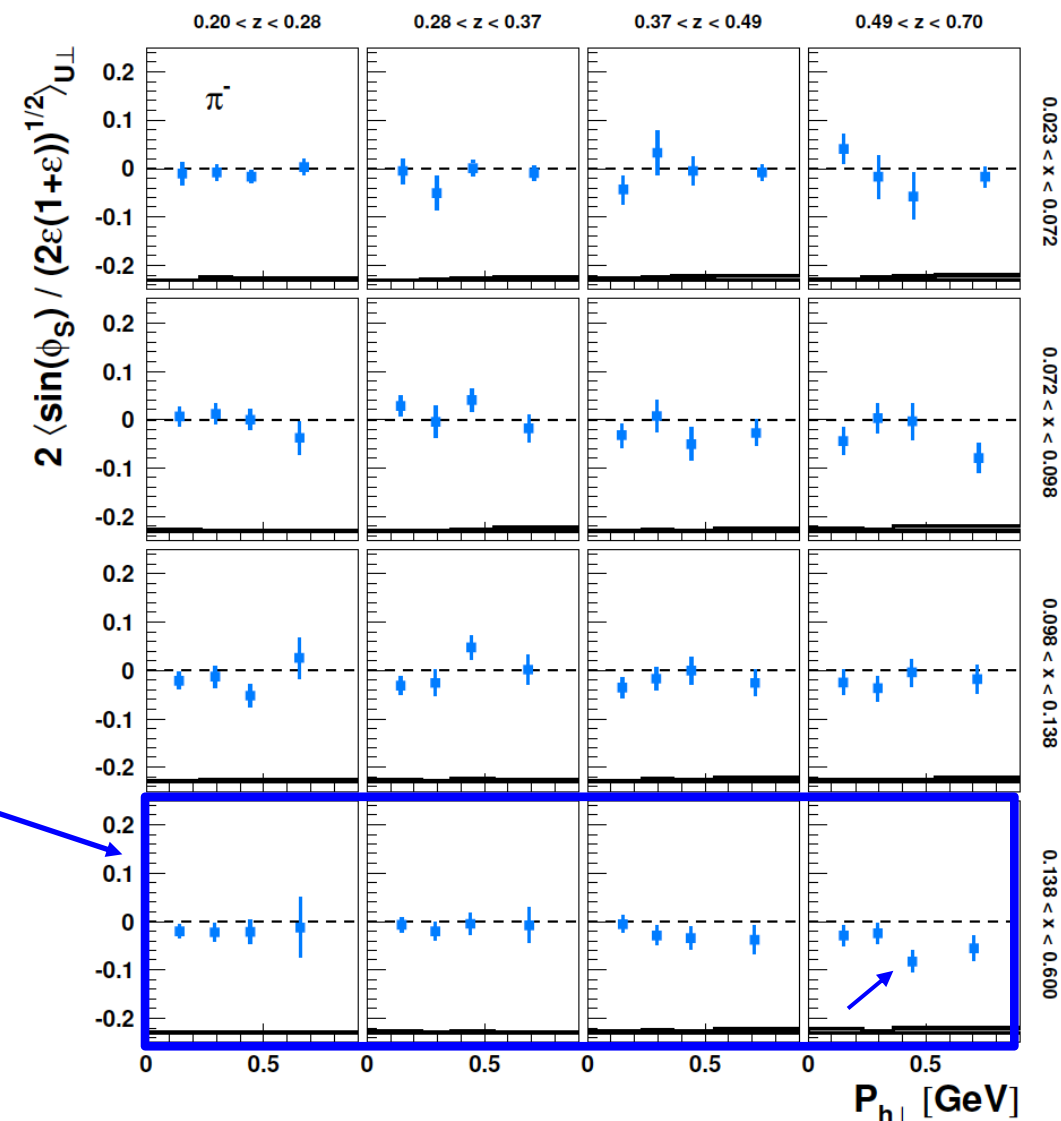
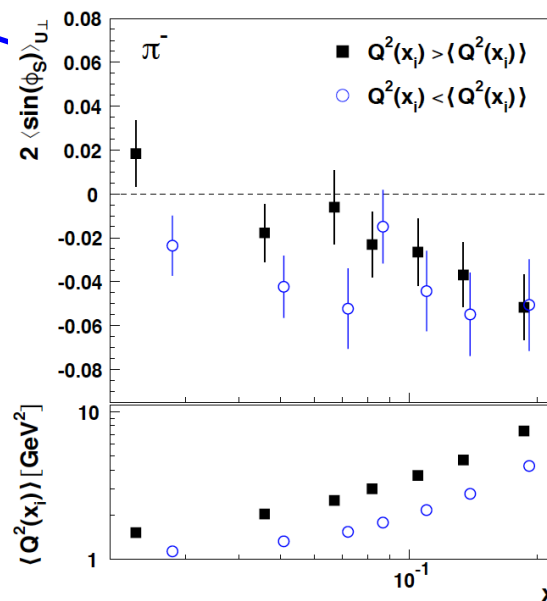
- **First measurement of Collins asymm. for protons/antiprotons!**
- **proton amplitude is non zero (negative)**
- antiproton amplitude  $\approx 0$
- Collins effect is a fragmentation process, but too little is known about this effect for spin- $\frac{1}{2}$  hadron production

# The sub-leading twist $\sin \phi_S$ term: pions SFA results

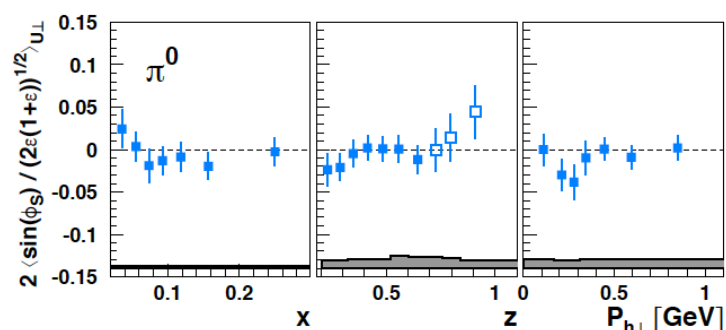
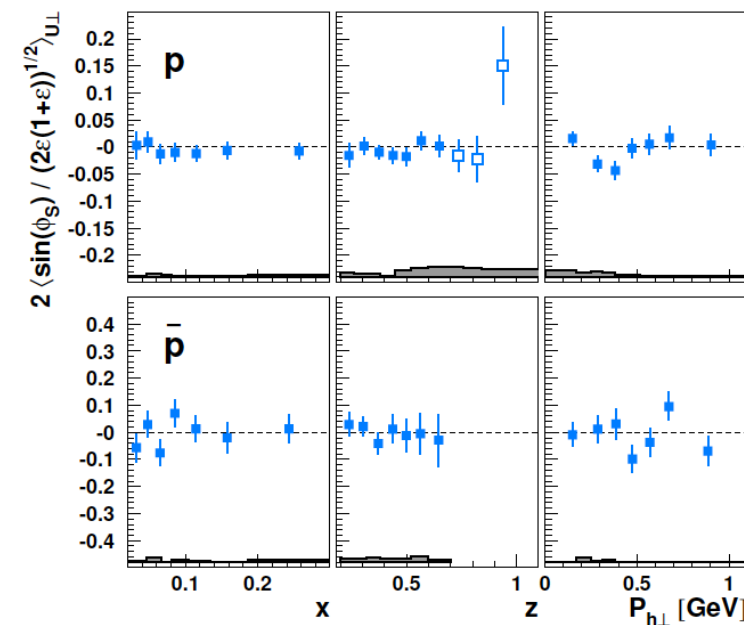
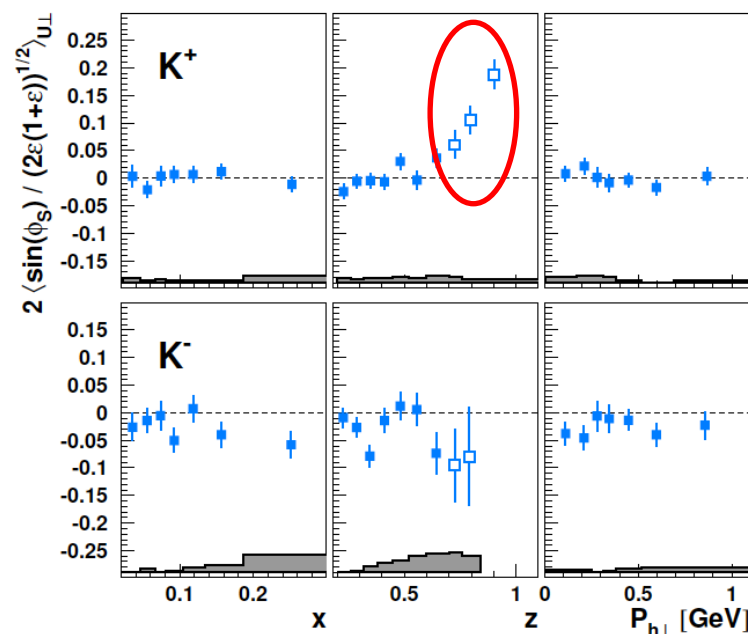
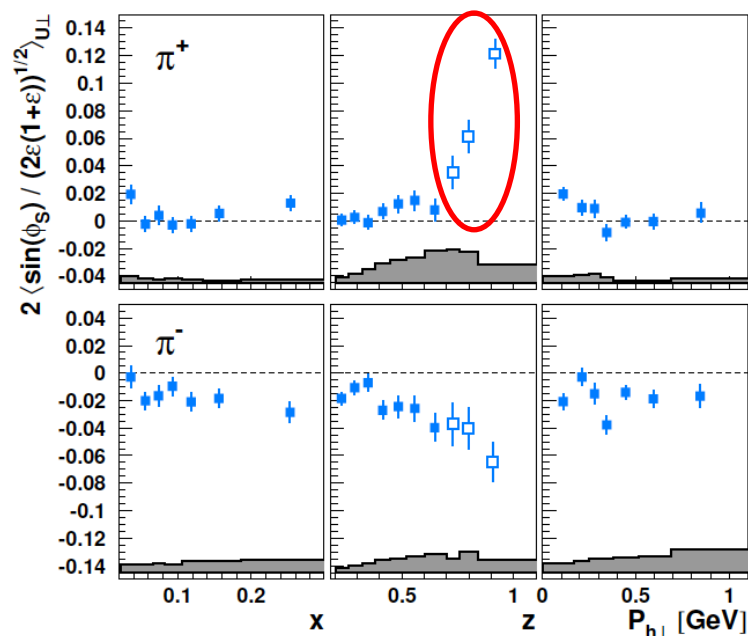


- Charged pions amplitudes non-zero and opposite
- Negative  $\pi^-$  amplitude increases with  $x$  and  $z$
- Overall similar behaviour of Collins asymmetries!

- Subleading-twist term: interesting to study the  $Q^2$  dependence
- Split each  $x$ -bin in two  $Q^2$  regions
- Hint of suppression at higher  $Q^2$



# The sub-leading twist $\sin \phi_S$ term: all SFA 1D results



- $\pi^+$  and  $K^+$  amplitudes in SIDIS region ( $0.2 < z < 0.7$ ) are similar: small and positive
- $K^-$  negative and similar to  $\pi^-$
- $\pi^0, p, \bar{p}$  results vanishing
- Striking  $z$ -dependence in “semi-exclusive region” for  $\pi^+ / K^+$  consistent with large  $\sin(\phi_S)$  amplitude observed in exclusive  $\pi^+$  electroproduction [Phys. Lett. B 682 (2010)]

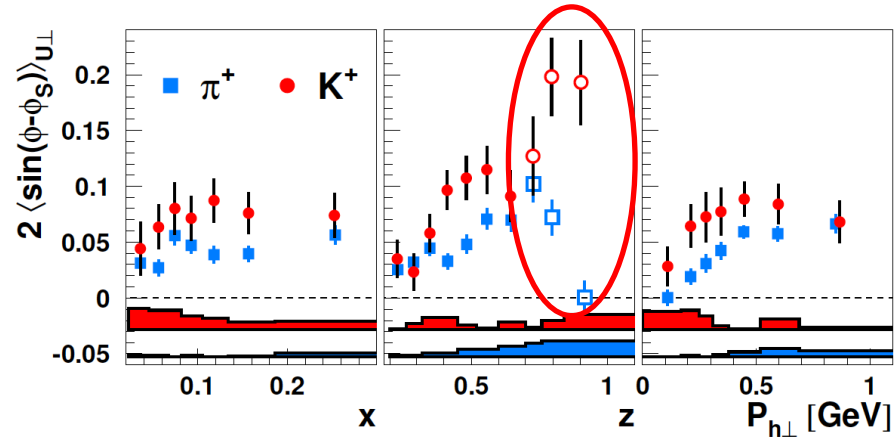


# Conclusions

- The full collection of leading- and subleading-twist SSAs and DSAs with a transversely polarized H target has recently been published, based on an improved analysis including proton/antiproton results, as well as results in a 3D binning and extended to the large- $z$  ("semi-exclusive" ) region.
- A **rich phenomenology** and surprising effects arise when intrinsic transverse degrees of freedom (spin, momentum) are not integrated out!
- **Flavor sensitivity** ensured by the excellent hadron ID of the HERMES experiment reveals interesting and unexpected facets of data (e.g.  $\pi \leftrightarrow K$ , see backup)
- The **3D imaging of the nucleon** is a fascinating and fast evolving research field. HERMES has been a pioneer experiment in this field and continues to play a key role in these studies.

Backup

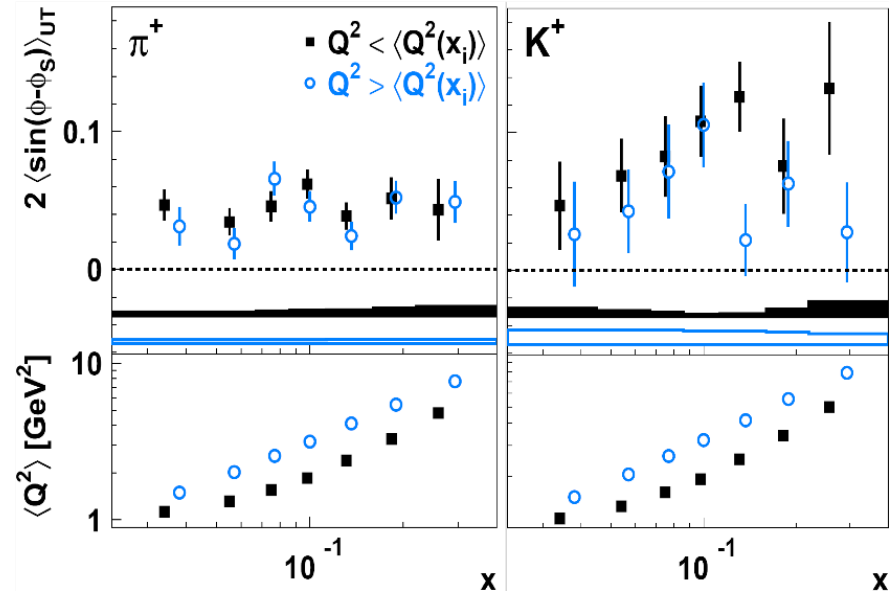
# Sivers amplitudes: the $K^+$ vs. $\pi^+$ issue



Similar kinematic dependence in SIDIS region but  $K^+$  is substantially larger!

- both expected to be mainly produced from scattering off  $u$ -quarks
- but **different sea-quark content**
- there could be a different  $k_T$  dependence of the fragmentation functions for different (sea) quarks flavours (entering the convolution integral)?
- different impact of higher-twist effects ?
- $K^+$  amplitude keeps rising with  $z$  in semi-exclusive region (no sudden change)  $\rightarrow$  Contribution from exclusive VM decays much less pronounced for Kaons than for pions.

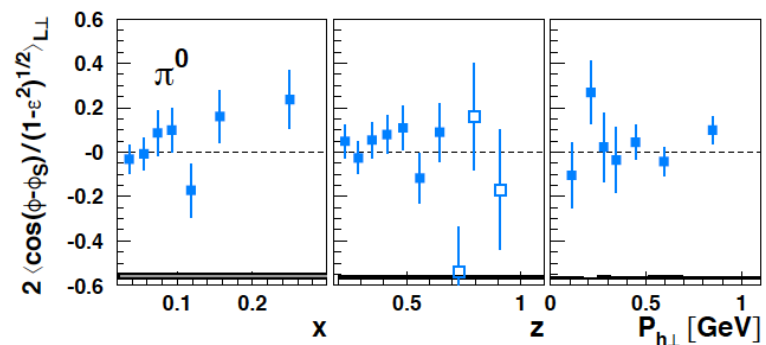
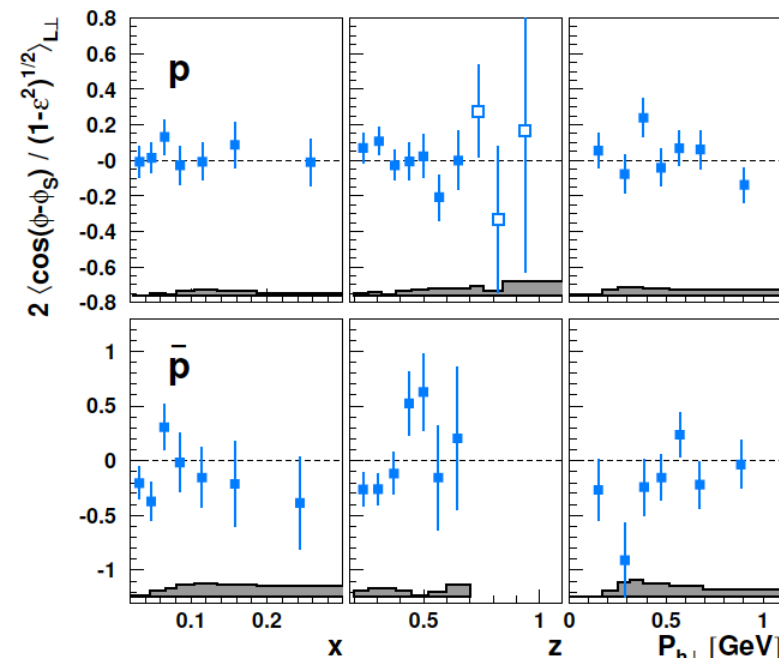
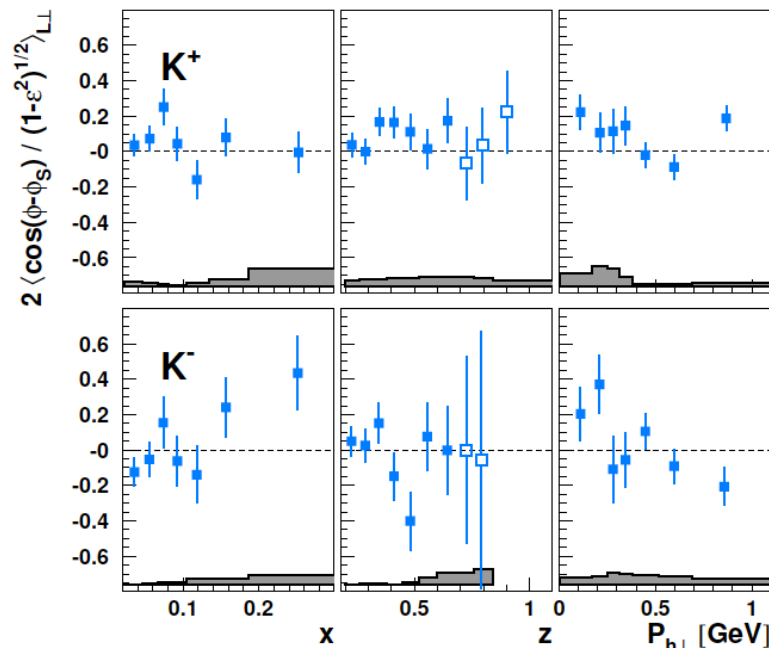
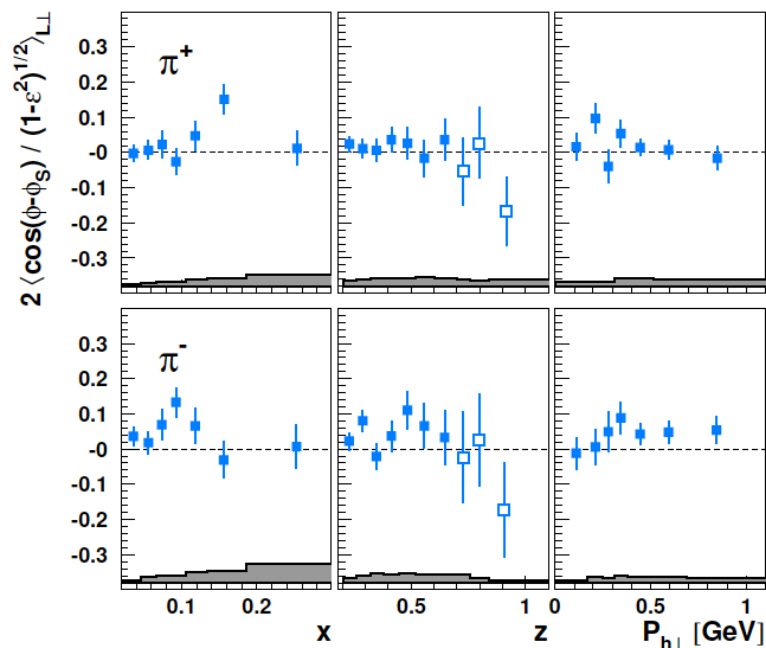
Phys. Rev. Lett. 103 (2009) 152002



- $x$ - $Q^2$  strongly correlated  $\rightarrow$  split each  $x$  bin in two  $Q^2$  regions:  $\leq \langle Q^2 \rangle$  of each  $x$  bin
- no effect for pions, but hint of suppression at larger  $Q^2$  for kaons

The other SFA results...

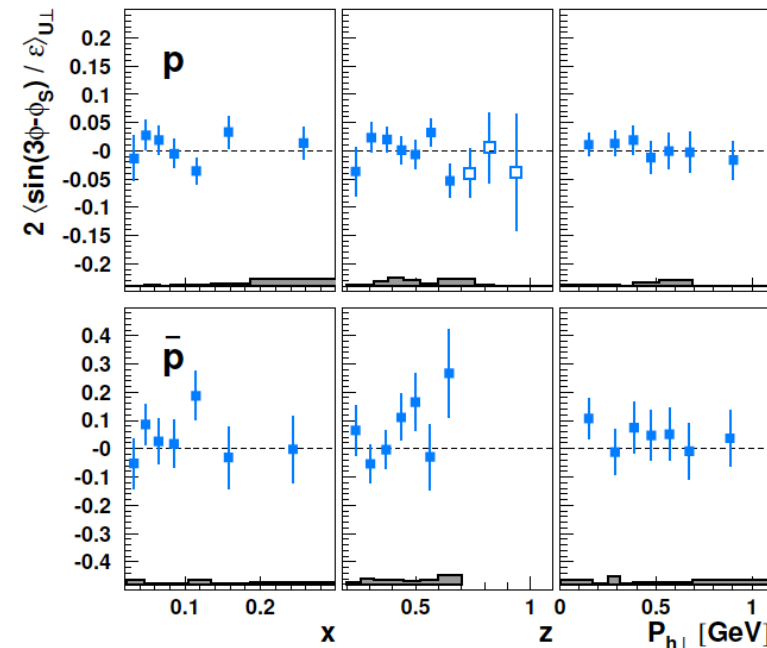
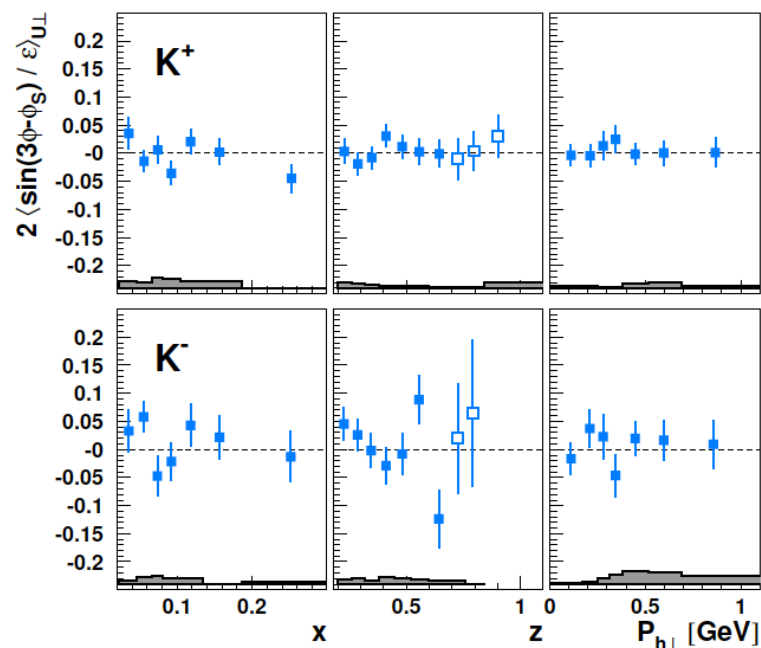
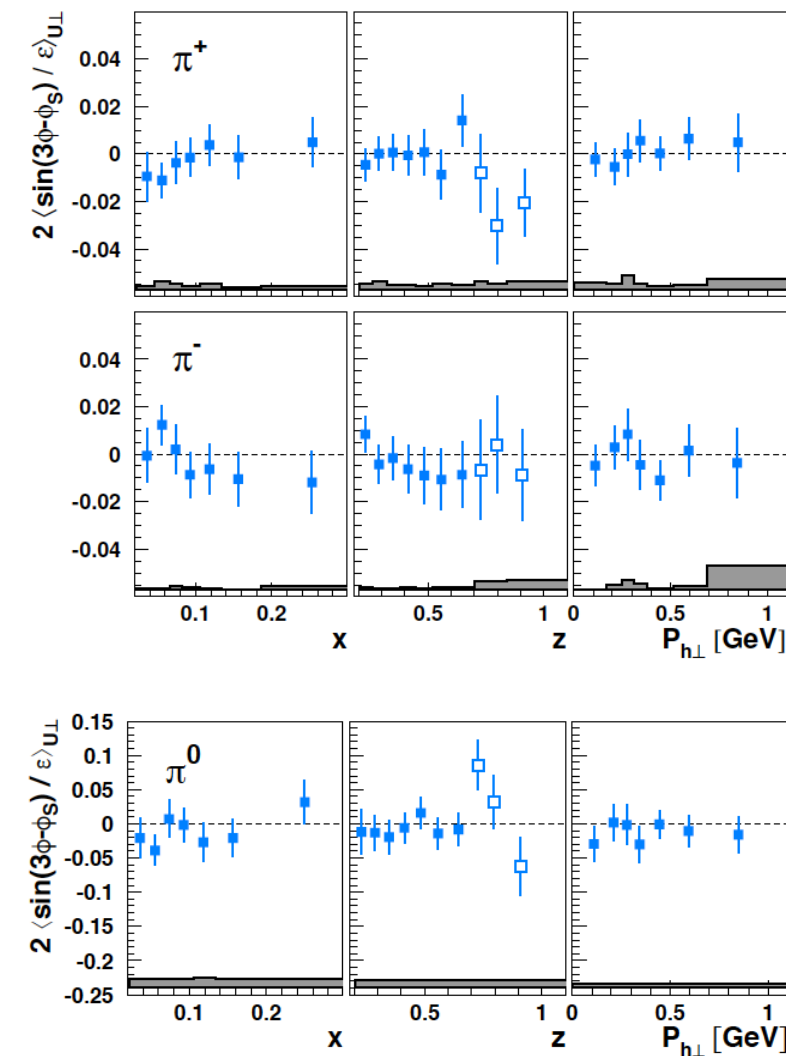
# The $\cos(\phi - \phi_S)$ DSA: all SFA 1D results



- $\pi^+$ ,  $\pi^-$  and  $K^+$  amplitudes are non-zero in SIDIS region ( $0.2 < z < 0.7$ )
- indication of a non-zero worm-gear function  $g_{1T}$
- amplitudes consistent with zero for all other hadron species
- Larger stat. errors (compared to SSAs) due to low beam polarization



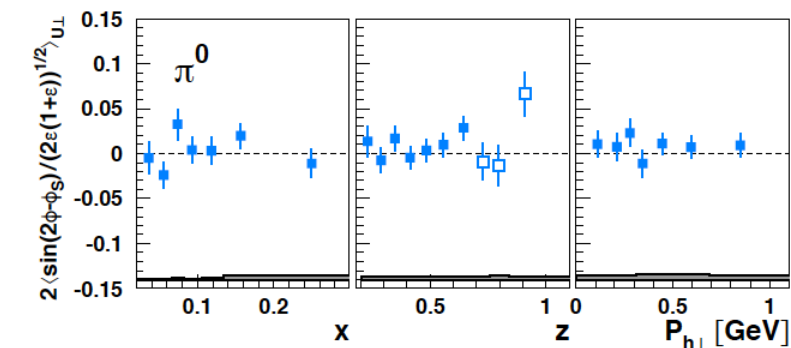
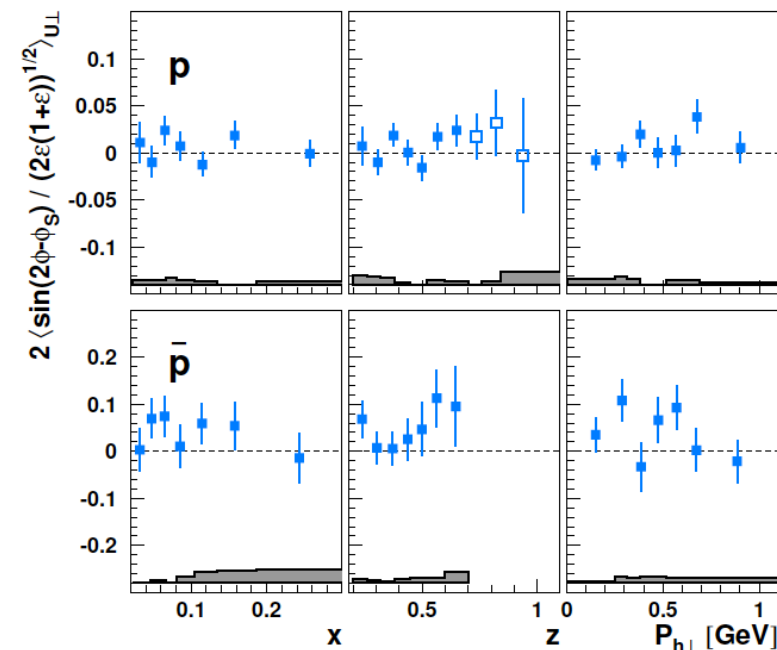
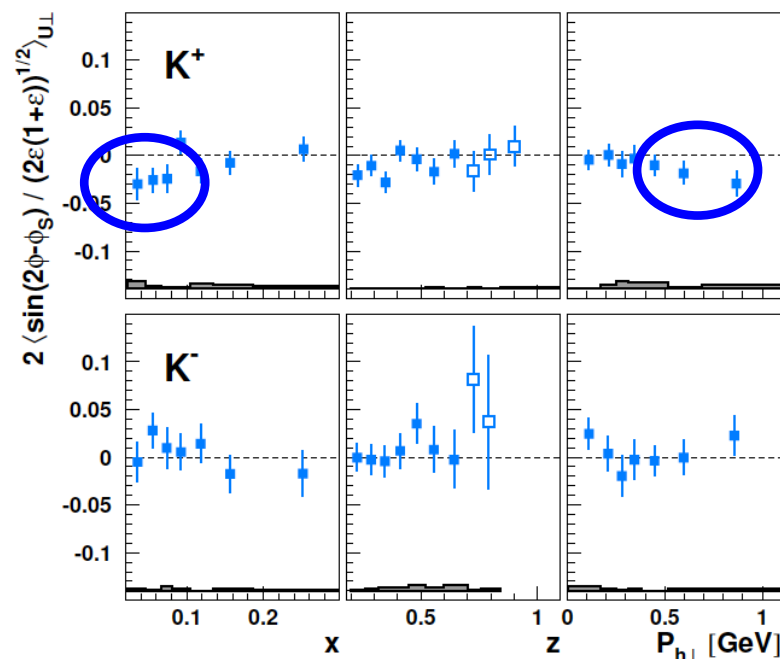
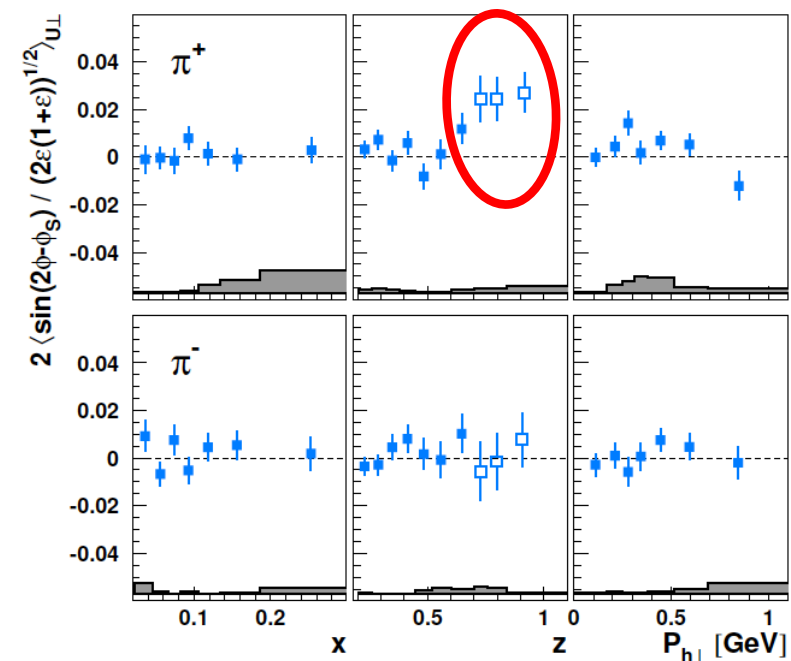
# $\langle \sin(3\phi - \phi_S) / \varepsilon \rangle_{U\perp}$ (Pretzelosity): all 1D results



Suppressed by two powers of  $P_{h\perp}$  w.r.t. Collins and Sivers amplitudes

All results are consistent with zero

$$\left\langle \sin(2\phi - \phi_S) / \sqrt{2\varepsilon(1 + \varepsilon)} \right\rangle_{U\perp} : \text{all 1D results}$$



**Semi-Inclusive region ( $0.2 < z < 0.7$ ):**

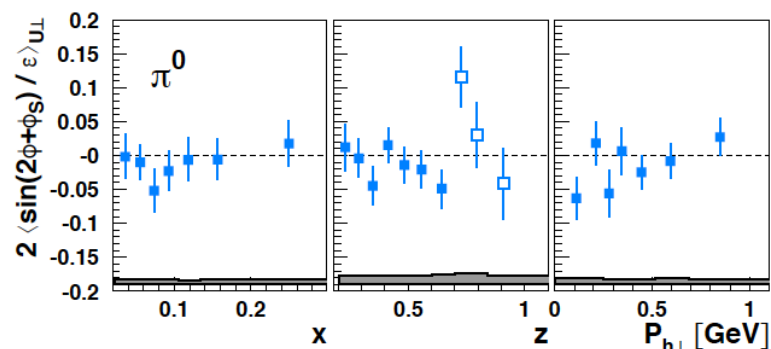
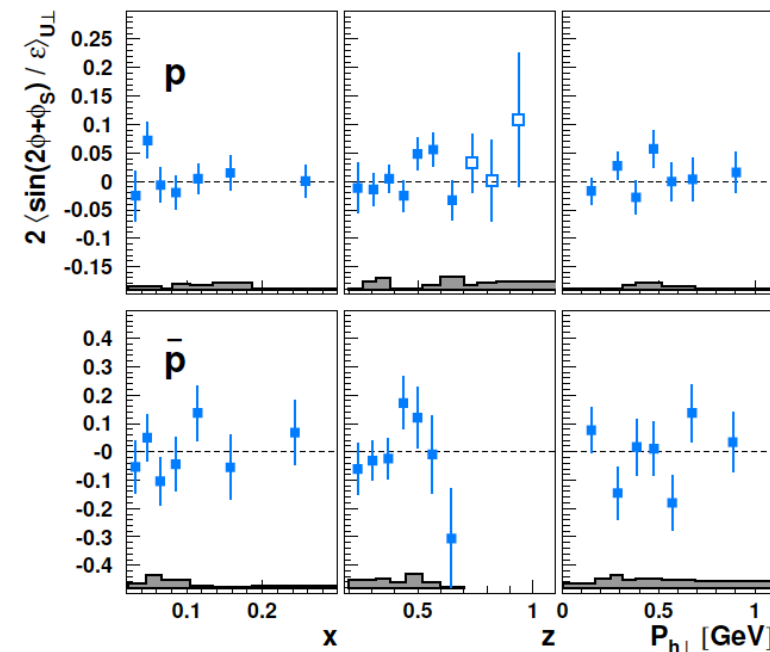
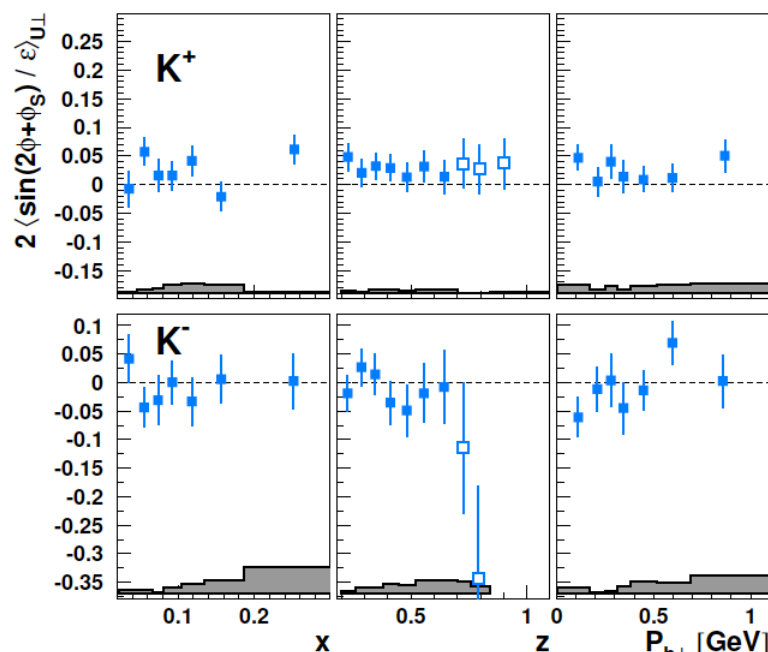
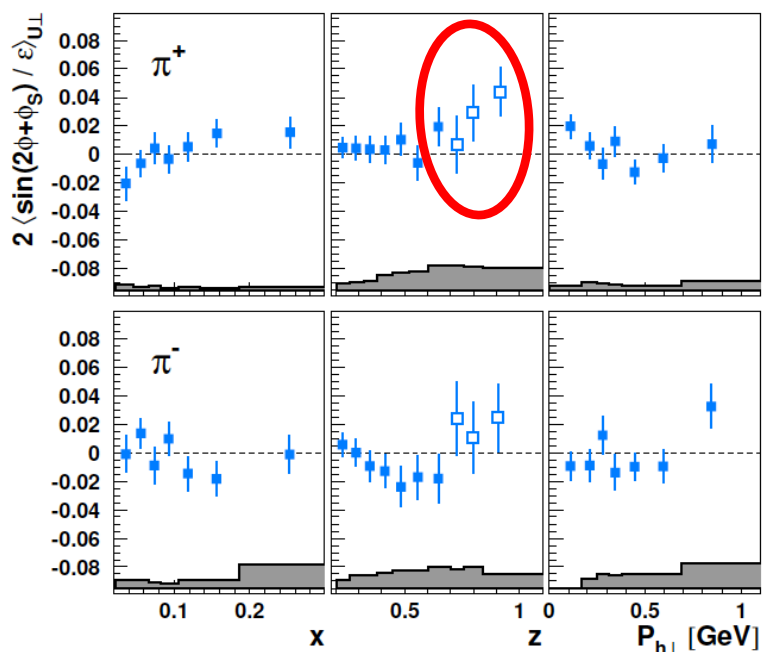
$K^+$ : hint of non-zero signal at small  $x$  and large  $P_{h\perp}$

$\bar{p}$ : hint of positive amplitude rising with  $z$

**Semi-Exclusive region ( $z > 0.7$ ):**

$\pi^+$ : positive amplitude ( $\sim 2\%$ )  $\rightarrow$  consistent with positive  $\sin(2\phi - \phi_S)$   
amplitude observed for exclusive  $\pi^+$  electroproduction [Phys. Lett. B 682 (2010)]

# $\langle \sin(2\phi + \phi_S) / \varepsilon \rangle_{U\perp}$ : all 1D results



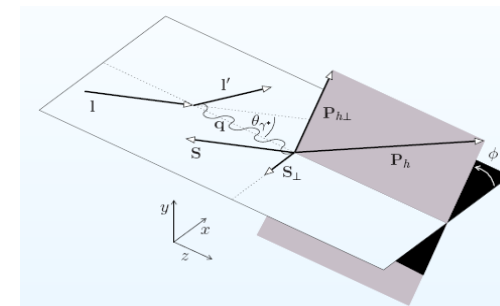
Arises solely from the small longitudinal target polarization component

**Semi-Inclusive region ( $0.2 < z < 0.7$ ):**

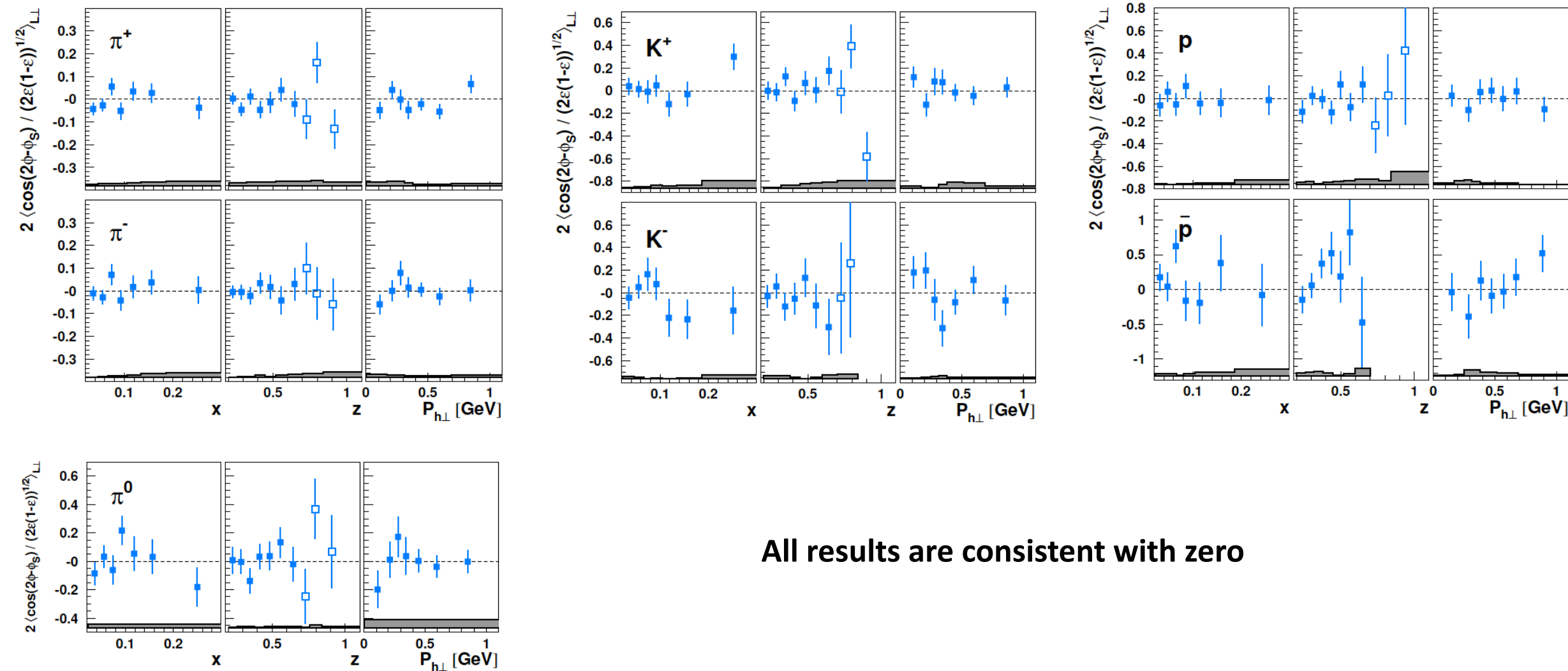
$K^+$ : positive amplitude over full  $z$  range

**Semi-Exclusive region ( $z > 0.7$ ):**

$\pi^+$ : positive amplitude rising with  $z \rightarrow$  consistent with positive  $\sin(2\phi + \phi_S)$  amplitude observed for exclusive  $\pi^+$  electroproduction [Phys. Lett. B 682 (2010)]

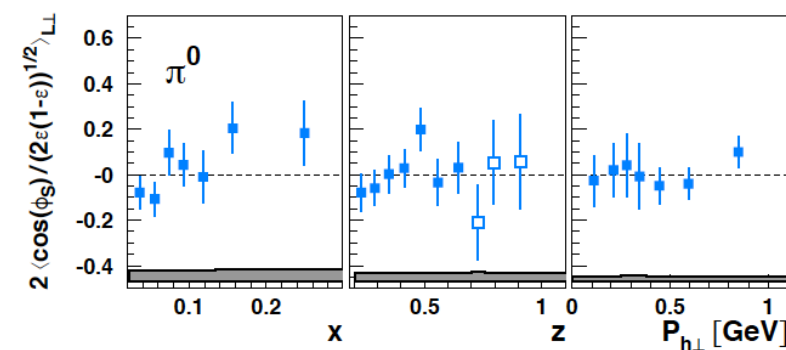
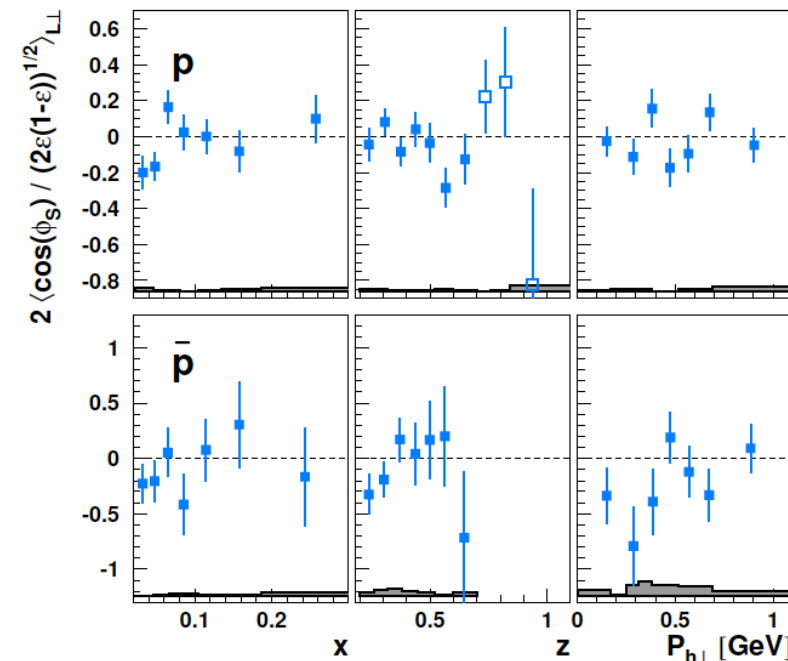
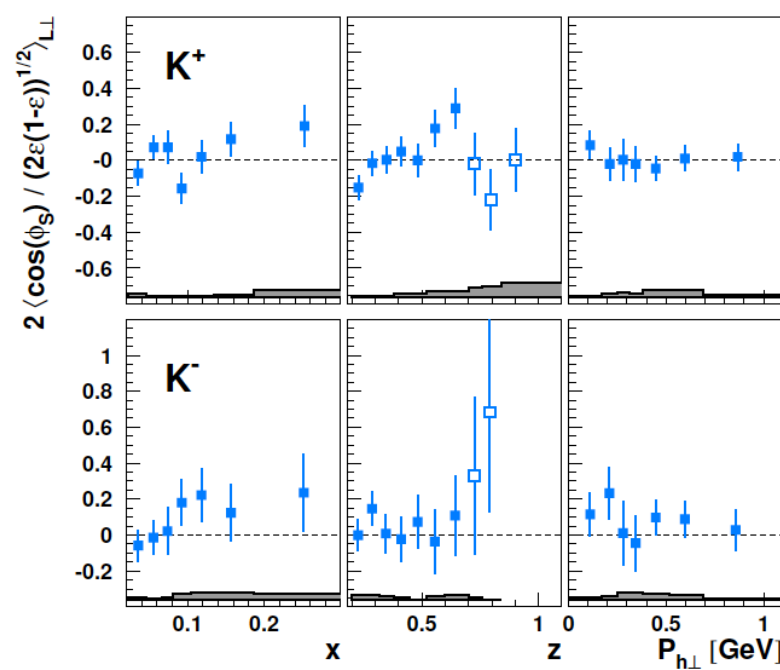
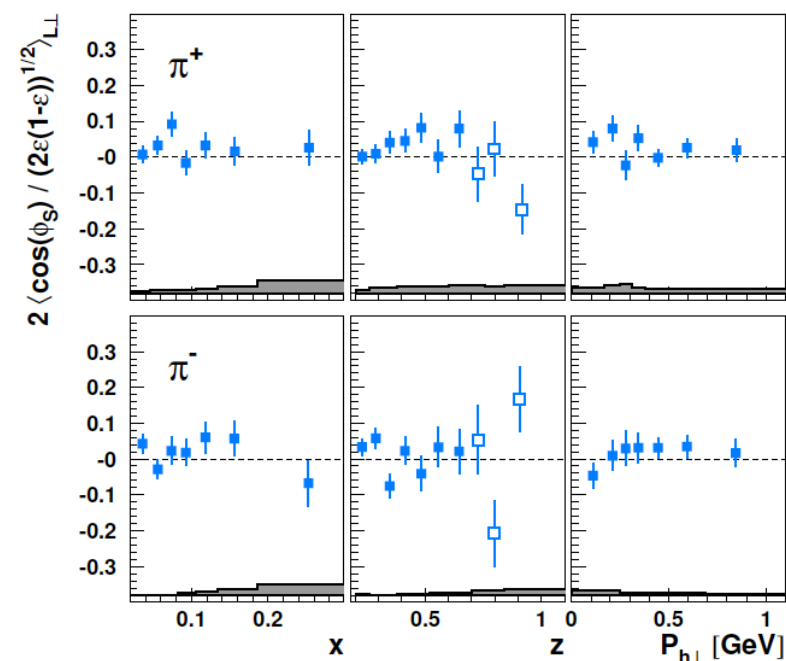


$$\left\langle \cos(2\phi - \phi_S) / \sqrt{2\varepsilon(1-\varepsilon)} \right\rangle_{L\perp} : \text{all 1D results}$$



All results are consistent with zero

$\left\langle \cos(\phi_S) / \sqrt{2\varepsilon(1-\varepsilon)} \right\rangle_{L\perp}$  : all 1D results

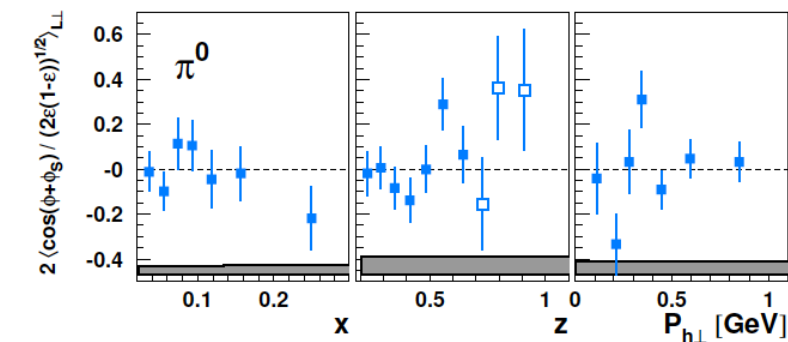
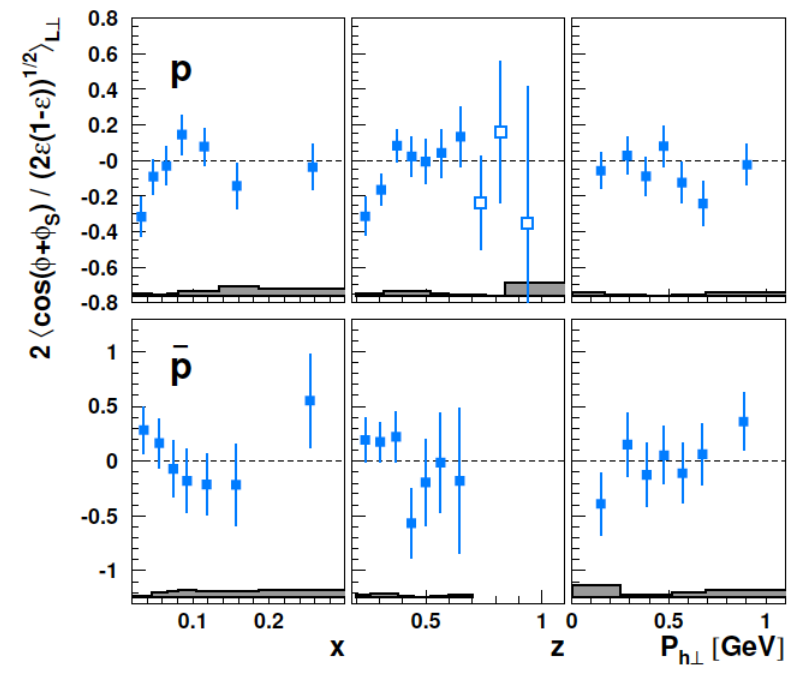
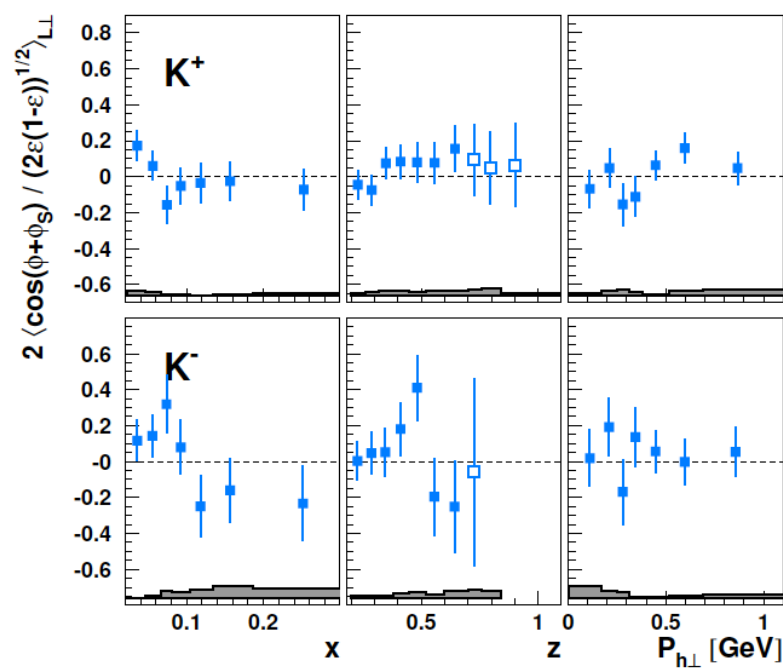
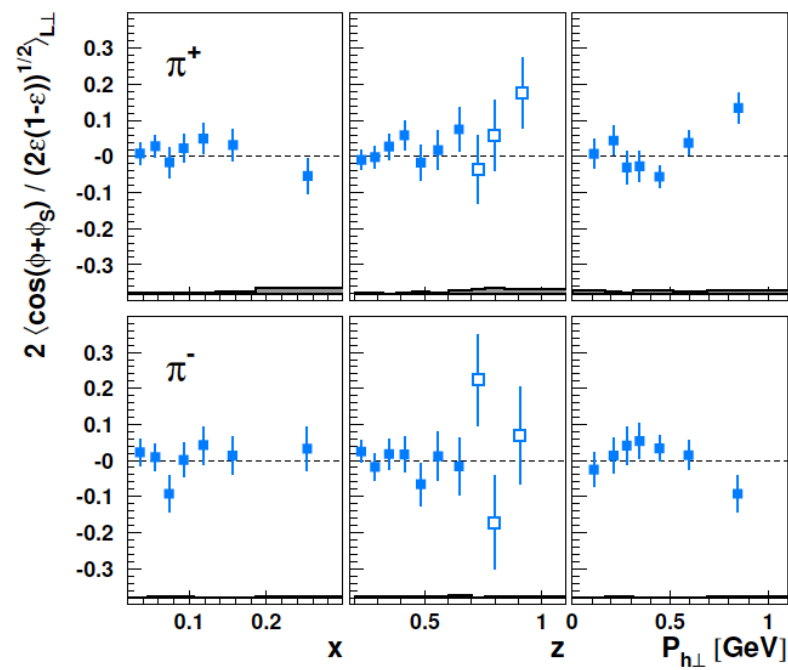


Can receive contributions from the longitudinal target polarization component

$K^-$  : small positive amplitude

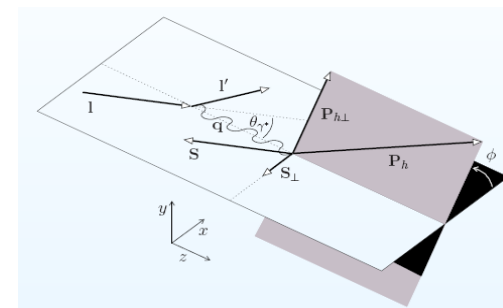


$$\left\langle \cos(\phi + \phi_S) / \sqrt{2\varepsilon(1 - \varepsilon)} \right\rangle_{L\perp} : \text{all 1D results}$$



Arises solely from the small longit.  
target polarization component

All results consistent with zero



# Miscellania

# The CSA amplitudes

The probability-density function used for the **CSA decomposition** of the cross section

$$\begin{aligned}
 & \mathbb{P}\left(x, z, P_{h\perp}, \phi, \phi_S, P_l, S_\perp : 2 \langle \sin(\phi - \phi_S) \rangle_{U\perp}^h, \dots, 2 \langle \cos(\phi + \phi_S) \rangle_{L\perp}^h \right) \\
 &= \left[ 1 + S_\perp \left( 2 \langle \sin(\phi - \phi_S) \rangle_{U\perp}^h \sin(\phi - \phi_S) + 2 \langle \sin(\phi + \phi_S) \rangle_{U\perp}^h \sin(\phi + \phi_S) + \right. \right. \\
 &\quad \left. \left. 2 \langle \sin(3\phi - \phi_S) \rangle_{U\perp}^h \sin(3\phi - \phi_S) + 2 \langle \sin(\phi_S) \rangle_{U\perp}^h \sin(\phi_S) + \right. \right. \\
 &\quad \left. \left. 2 \langle \sin(2\phi - \phi_S) \rangle_{U\perp}^h \sin(2\phi - \phi_S) + 2 \langle \sin(2\phi + \phi_S) \rangle_{U\perp}^h \sin(2\phi + \phi_S) \right) \right. \\
 &\quad \left. + P_l S_\perp \left( 2 \langle \cos(\phi - \phi_S) \rangle_{L\perp}^h \cos(\phi - \phi_S) + 2 \langle \cos(\phi_S) \rangle_{L\perp}^h \cos(\phi_S) + \right. \right. \\
 &\quad \left. \left. 2 \langle \cos(2\phi - \phi_S) \rangle_{L\perp}^h \cos(2\phi - \phi_S) + 2 \langle \cos(\phi + \phi_S) \rangle_{L\perp}^h \cos(\phi + \phi_S) \right) \right]^w
 \end{aligned}$$

$\left. \begin{array}{l} \text{Blue bracket} \\ \text{Red bracket} \end{array} \right\} \begin{array}{l} A_{U\perp} \text{ SSAs} \\ A_{L\perp} \text{ DSAs} \end{array}$

## 10 Fourier components:

- 6  $A_{U\perp}$  SSAs (4 leading-twist + 2 subleading twist)
- 4  $A_{L\perp}$  DSAs (2 leading-twist + 2 subleading twist)
- $\sin(2\phi + \phi_S)$  and  $\cos(\phi + \phi_S)$  terms arise purely from the small but non-vanishing longitudinal target-polarization component (target polarization states are referred to the lepton beam direction)
- **The CSA amplitudes include in their definition the  $\varepsilon$ -dependent kinematic prefactors that enter the various cross section terms**

# The SFA amplitudes (NEW!)

The probability-density function used for the **SFA decomposition** of the cross section

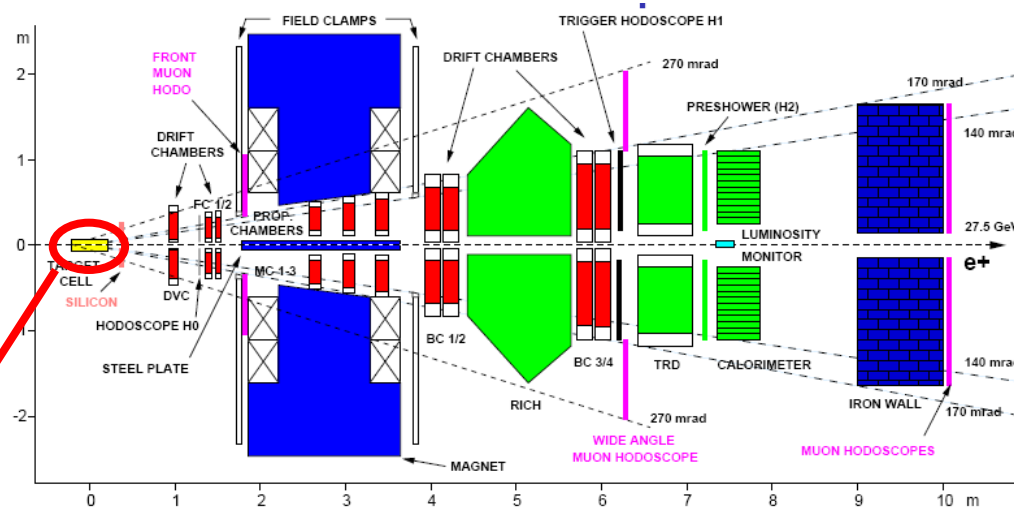
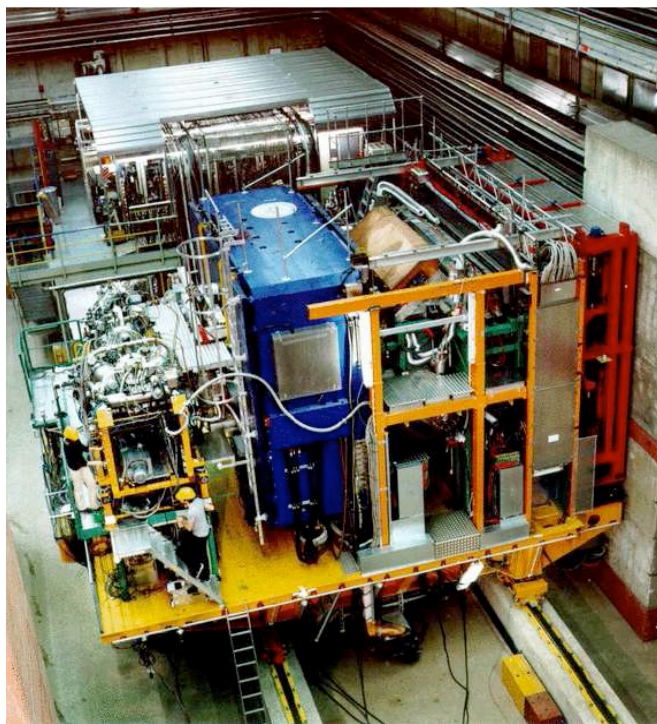
$$\begin{aligned}
 & \mathbb{P}\left(x, z, \epsilon, P_{h\perp}, \phi, \phi_S, P_l, S_\perp : 2\langle \sin(\phi - \phi_S) \rangle_{U\perp}^h, \dots, 2\langle \cos(\phi + \phi_S) / \sqrt{2\epsilon(1-\epsilon)} \rangle_{L\perp}^h \right) \\
 &= \left[ 1 + S_\perp \left( 2\langle \sin(\phi - \phi_S) \rangle_{U\perp}^h \sin(\phi - \phi_S) + \epsilon 2\langle \sin(\phi + \phi_S) / \epsilon \rangle_{U\perp}^h \sin(\phi + \phi_S) + \right. \right. \\
 & \quad \epsilon 2\langle \sin(3\phi - \phi_S) / \epsilon \rangle_{U\perp}^h \sin(3\phi - \phi_S) + \sqrt{2\epsilon(1+\epsilon)} 2\langle \sin(\phi_S) / \sqrt{2\epsilon(1+\epsilon)} \rangle_{U\perp}^h \sin(\phi_S) + \\
 & \quad \left. \sqrt{2\epsilon(1+\epsilon)} 2\langle \sin(2\phi - \phi_S) / \sqrt{2\epsilon(1+\epsilon)} \rangle_{U\perp}^h \sin(2\phi - \phi_S) + \epsilon 2\langle \sin(2\phi + \phi_S) / \epsilon \rangle_{U\perp}^h \sin(2\phi + \phi_S) \right) \\
 & \quad + P_l S_\perp \left( \sqrt{1-\epsilon^2} 2\langle \cos(\phi - \phi_S) / \sqrt{1-\epsilon^2} \rangle_{L\perp}^h \cos(\phi - \phi_S) + \sqrt{2\epsilon(1-\epsilon)} 2\langle \cos(\phi_S) / \sqrt{2\epsilon(1-\epsilon)} \rangle_{L\perp}^h \cos(\phi_S) + \right. \\
 & \quad \left. \sqrt{2\epsilon(1-\epsilon)} 2\langle \cos(2\phi - \phi_S) / \sqrt{2\epsilon(1-\epsilon)} \rangle_{L\perp}^h \cos(2\phi - \phi_S) + \sqrt{2\epsilon(1-\epsilon)} 2\langle \cos(\phi + \phi_S) / \sqrt{2\epsilon(1-\epsilon)} \rangle_{L\perp}^h \cos(\phi + \phi_S) \right) \Big]^w
 \end{aligned}$$

}  $A_{U\perp}$  SSAs  
}  $A_{L\perp}$  DSAs

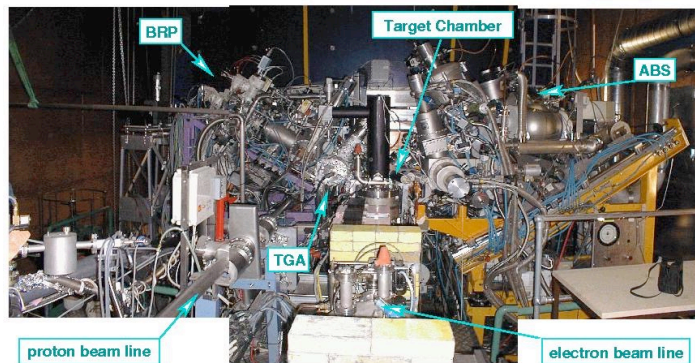
## 10 Fourier components:

- 6  $A_{U\perp}$  SSAs (4 leading-twist + 2 subleading twist)
- 4  $A_{L\perp}$  DSAs (2 leading-twist + 2 subleading twist)
- $\sin(2\phi + \phi_S)$  and  $\cos(\phi + \phi_S)$  terms arise purely from the small but non-vanishing longitudinal target-polarization component (target polarization states are referred to the lepton beam direction)
- **The SFA amplitudes do not include the  $\epsilon$ -dependent kinematic prefactors of the various cross section terms.**
- They are obtained by including explicitly the  $\epsilon$ -dependent kinematic prefactors in the probability-density function separated from the fit parameters.

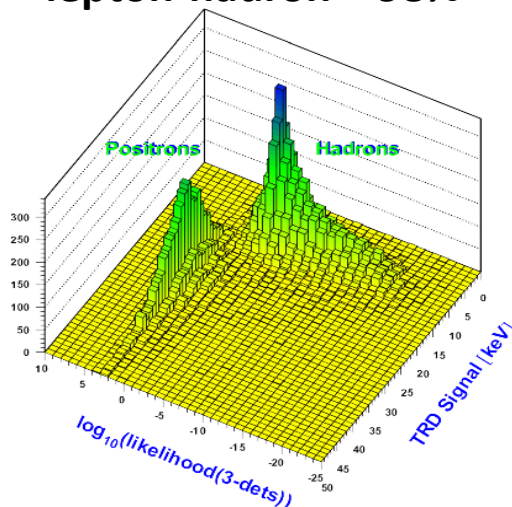
# The HERMES experiment at HERA (1995-2007)



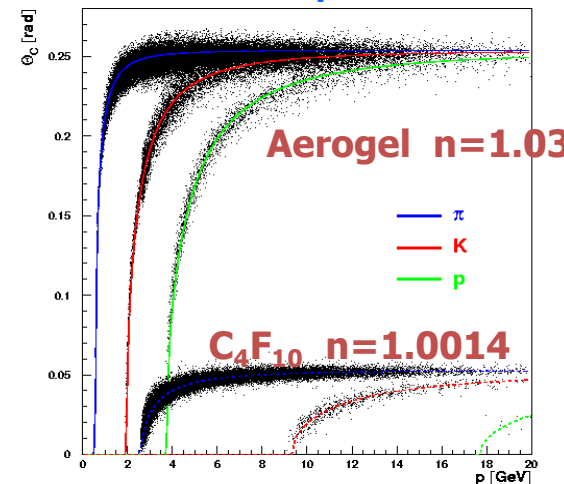
The polarized gas target



TRD, Calorimeter,  
preshower, RICH:  
lepton-hadron > 98%



hadron separation

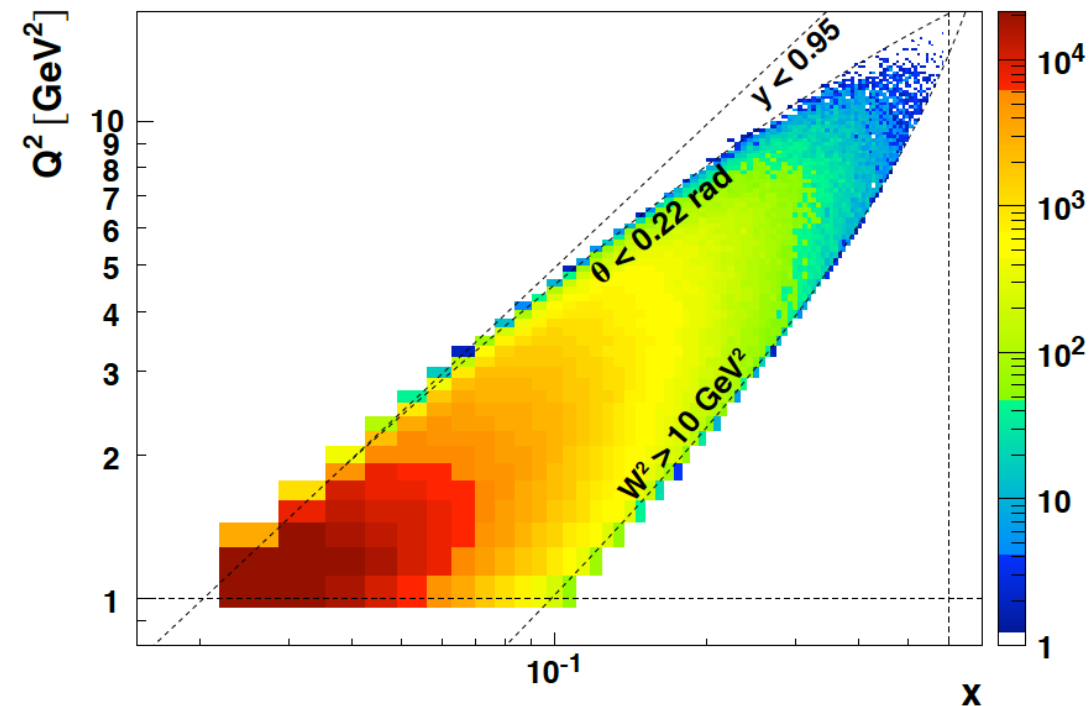
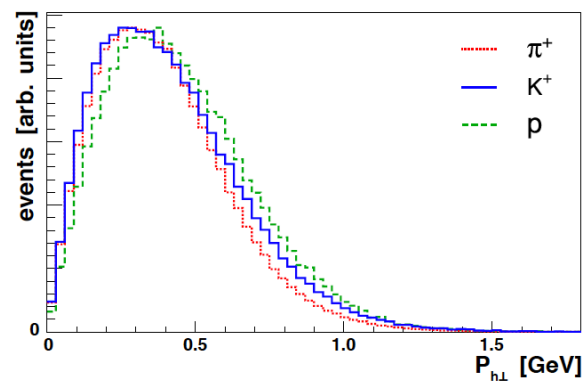
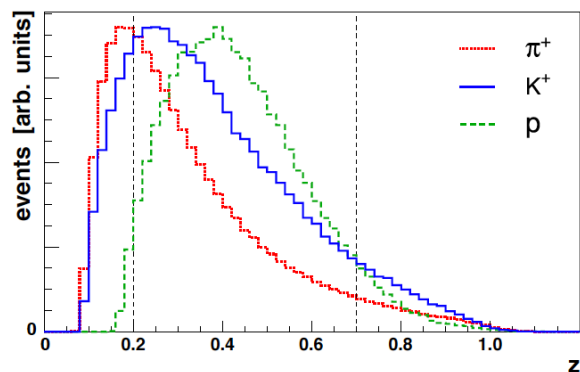


$\pi \sim 98\%$ ,  $K \sim 88\%$ ,  $P \sim 85\%$

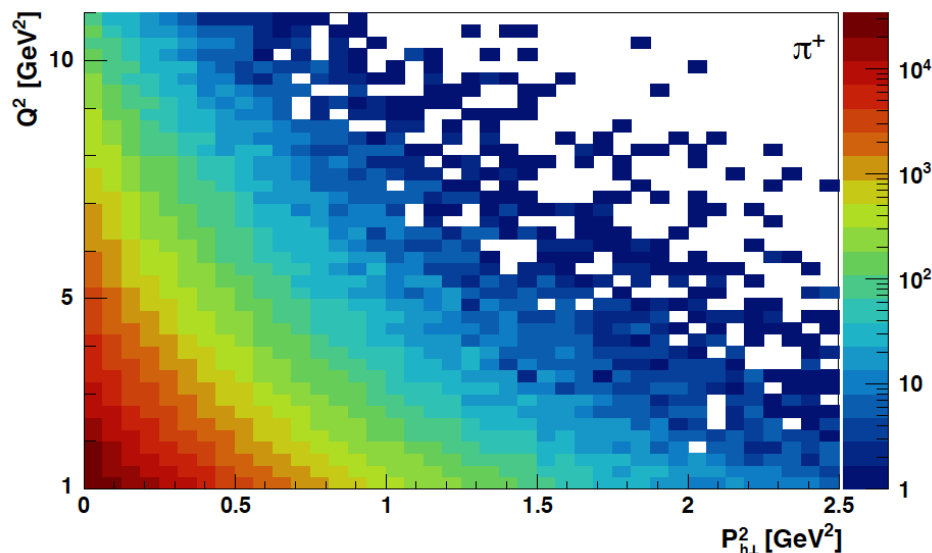


# Kinematic coverage

Scattered lepton:	$Q^2$	$> 1 \text{ GeV}^2$
	$W^2$	$> 10 \text{ GeV}^2$
	$0.023 < x$	$< 0.6$
	$0.1 < y$	$< 0.95$
Detected hadrons:	$2 \text{ GeV} <  \mathbf{P}_h $	$< 15 \text{ GeV}$ charged mesons
	$4 \text{ GeV} <  \mathbf{P}_h $	$< 15 \text{ GeV}$ (anti)protons
	$ \mathbf{P}_h $	$> 2 \text{ GeV}$ neutral pions
	$P_{h\perp}$	$< 2 \text{ GeV}$
	$0.2 < z$	$< 0.7$ (1.2 for the “semi-exclusive” region)

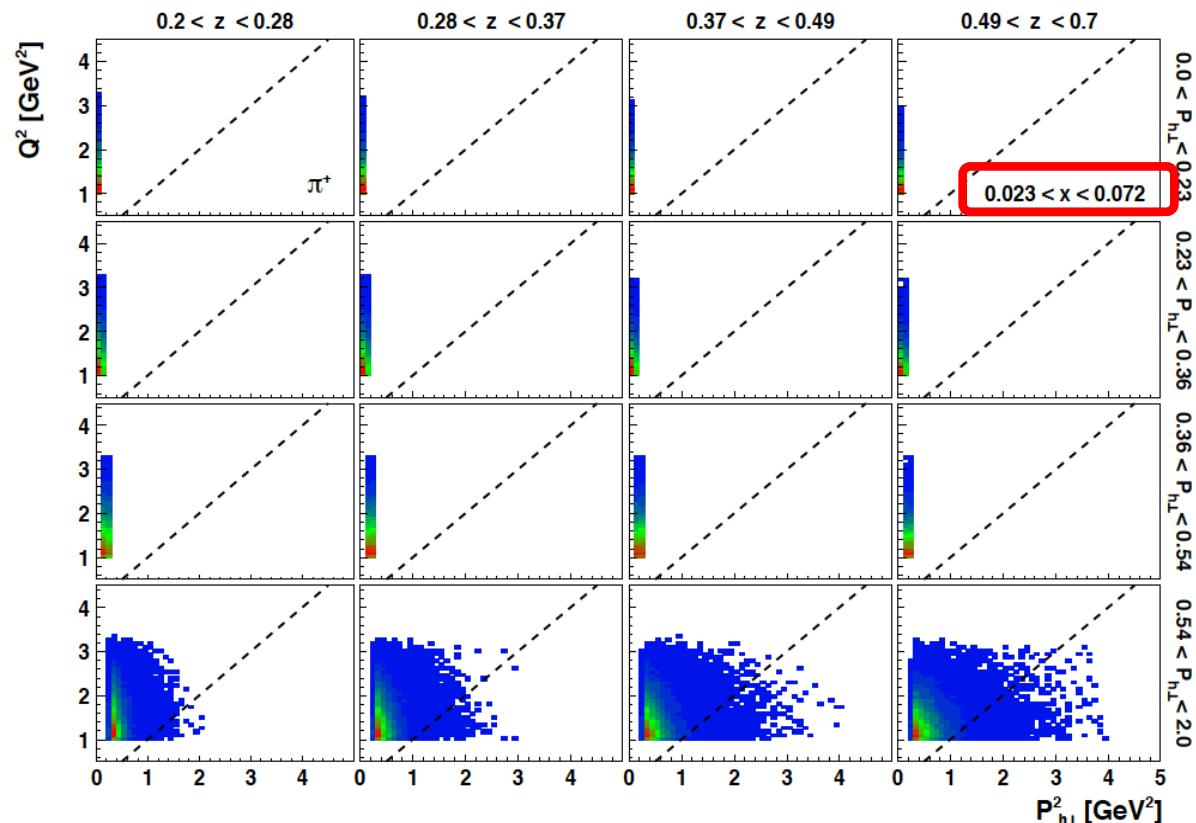


# Kinematic coverage and factorization requirements



- Factorization requirement  $P_{h\perp}^2 \ll Q^2$  fulfilled for most of the selected DIS events
- the stricter constraint  $P_{h\perp}^2 \ll z^2 Q^2$  is violated at large  $P_{h\perp}$  in the region of small  $x$  and small  $z$
- detailed studies in appendix B of the paper (and next slides)

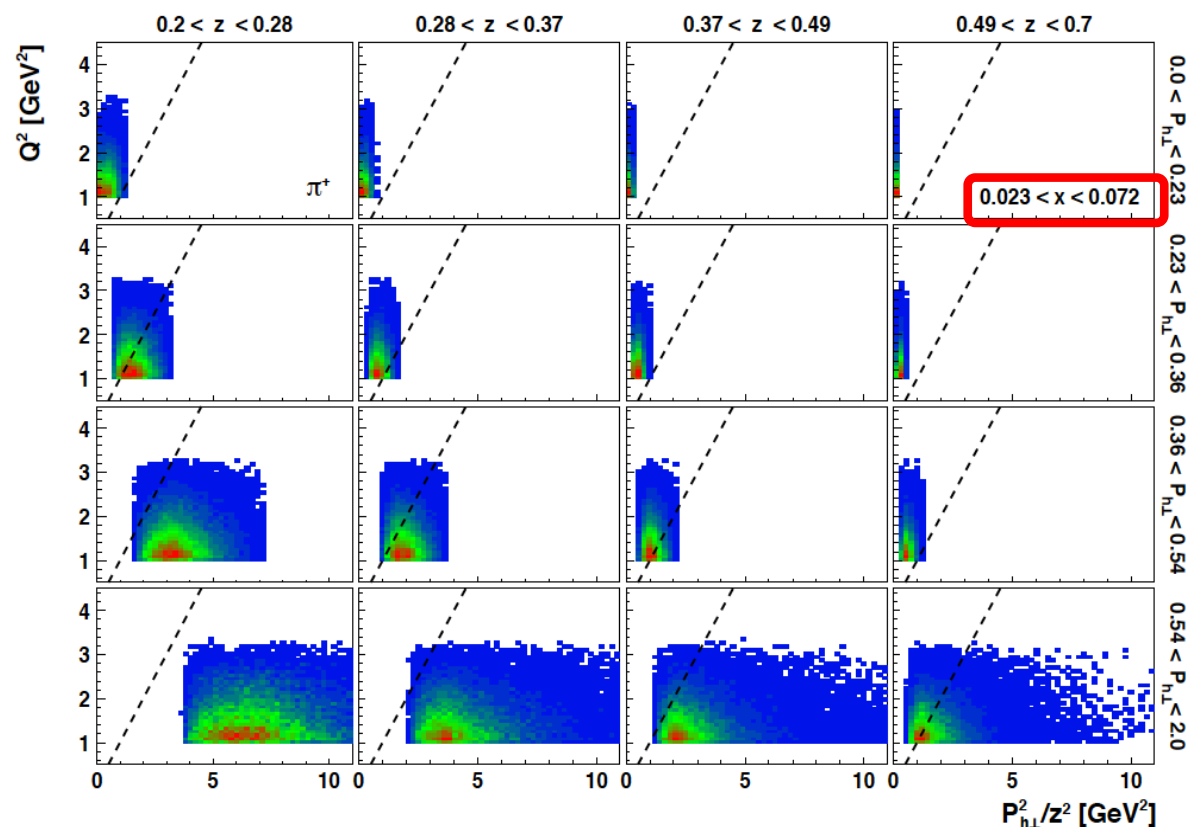
Due to  $x$ - $Q^2$  correlation, the **first  $x$  bin** corresponds to the **small  $Q^2$  region**, where the TMD-factorization requirement  $P_{h\perp}^2 \ll Q^2$  is less favourable.



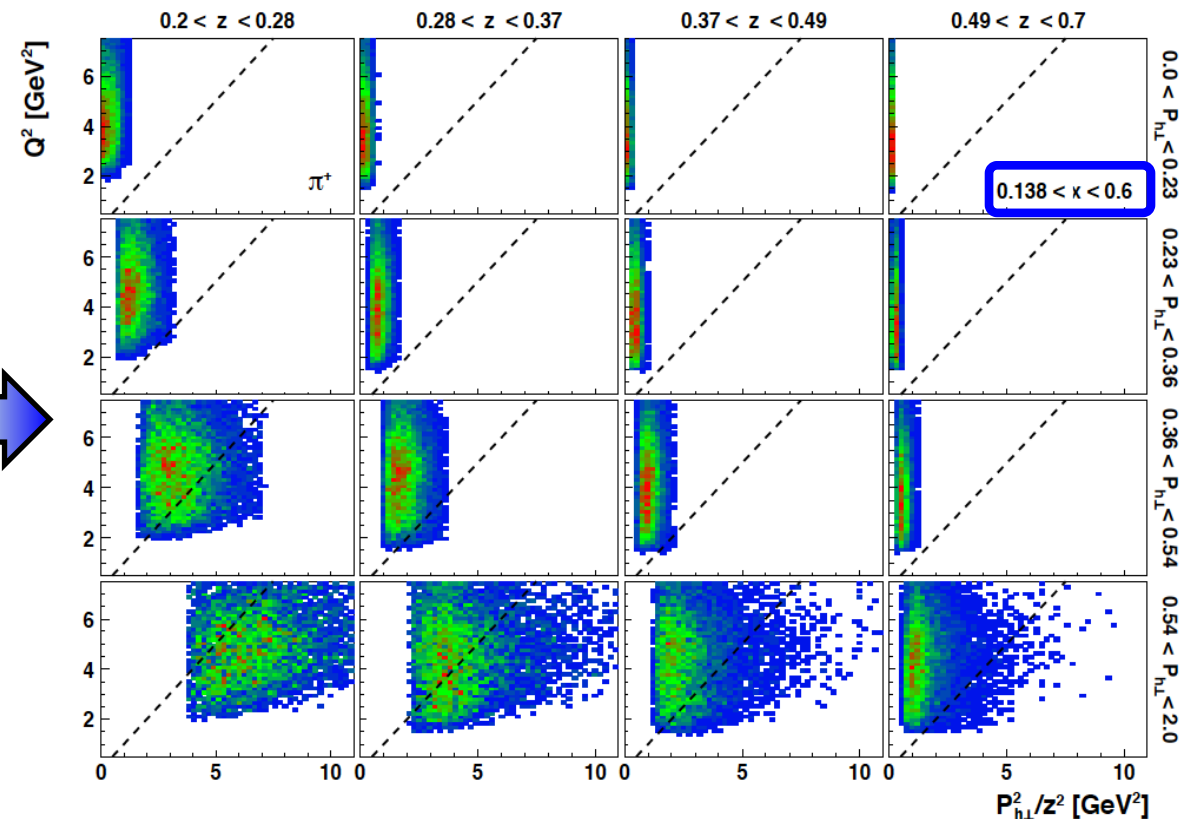
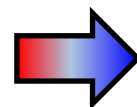
**TMD-factorization requirement  $P_{h\perp}^2 \ll Q^2$   
fulfilled for most of the selected DIS events!**

# Factorization requirements

Due to the  $1/z^2$  factor, which becomes large at small  $z$ , the **stricter condition**  $P_{h\perp}^2/z^2 \ll Q^2$  is unfulfilled for the majority of the events:

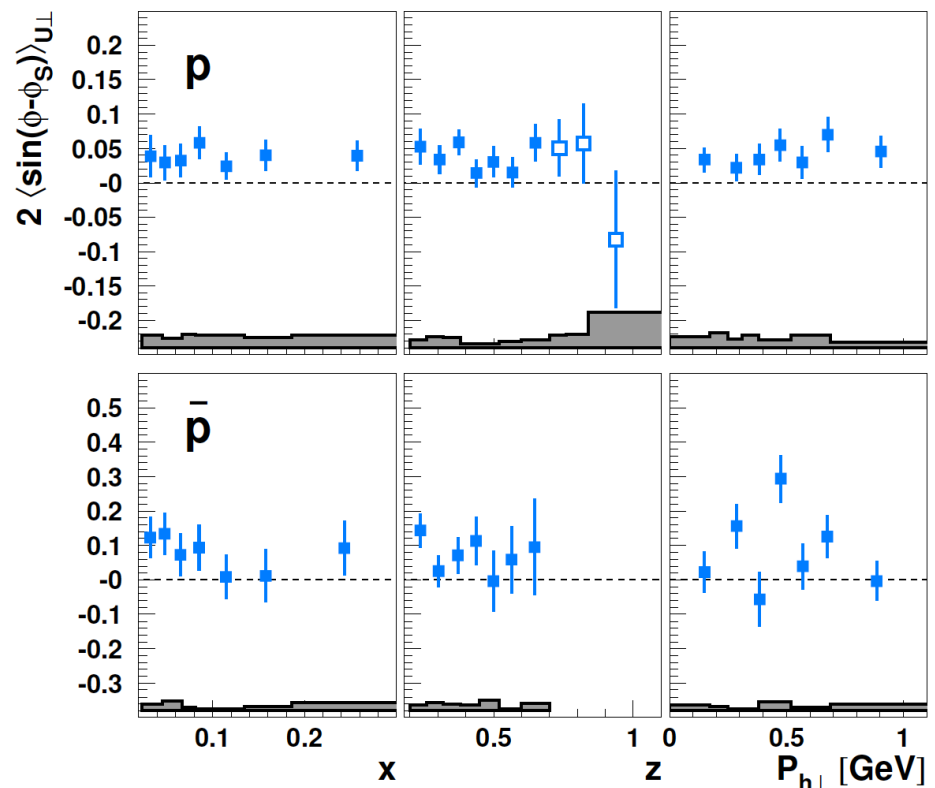


First  $x$ -bin  $\rightarrow$  smaller  $Q^2$



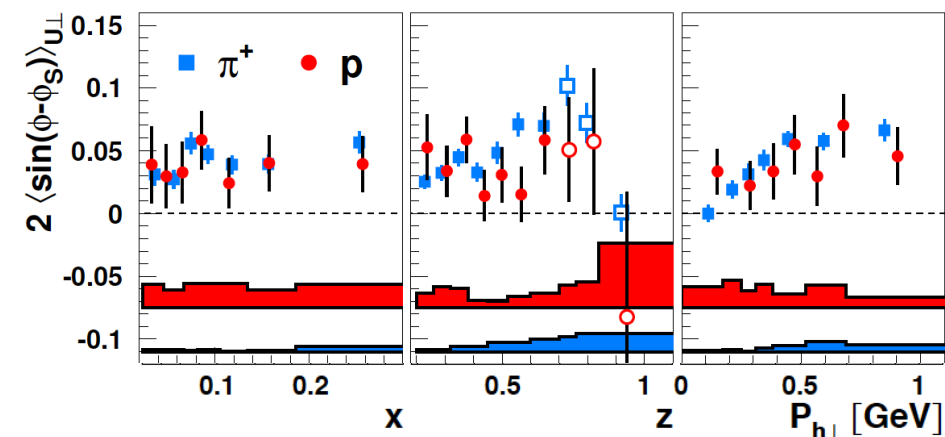
Last  $x$ -bin  $\rightarrow$  larger  $Q^2$

# Sivers amplitudes: protons results



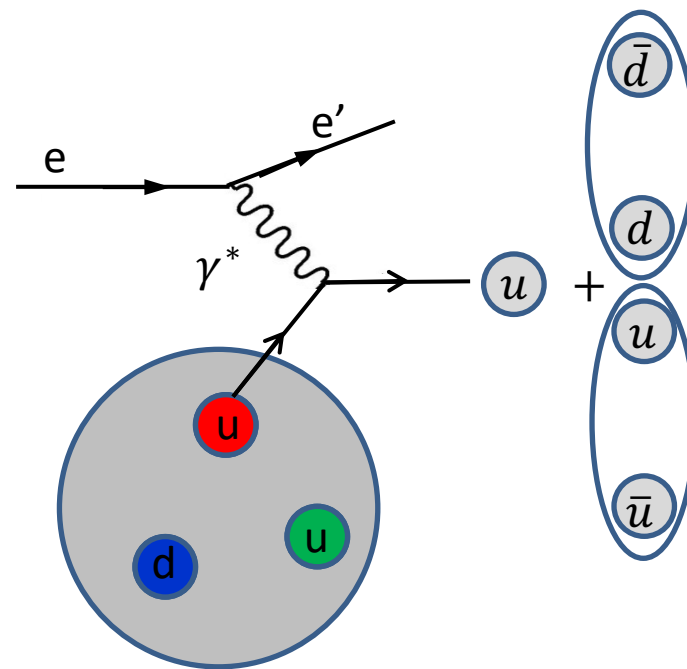
First measurement of Sivers asymmetries for  $p, \bar{p}$  in SIDIS

Both amplitudes are non-zero and positive



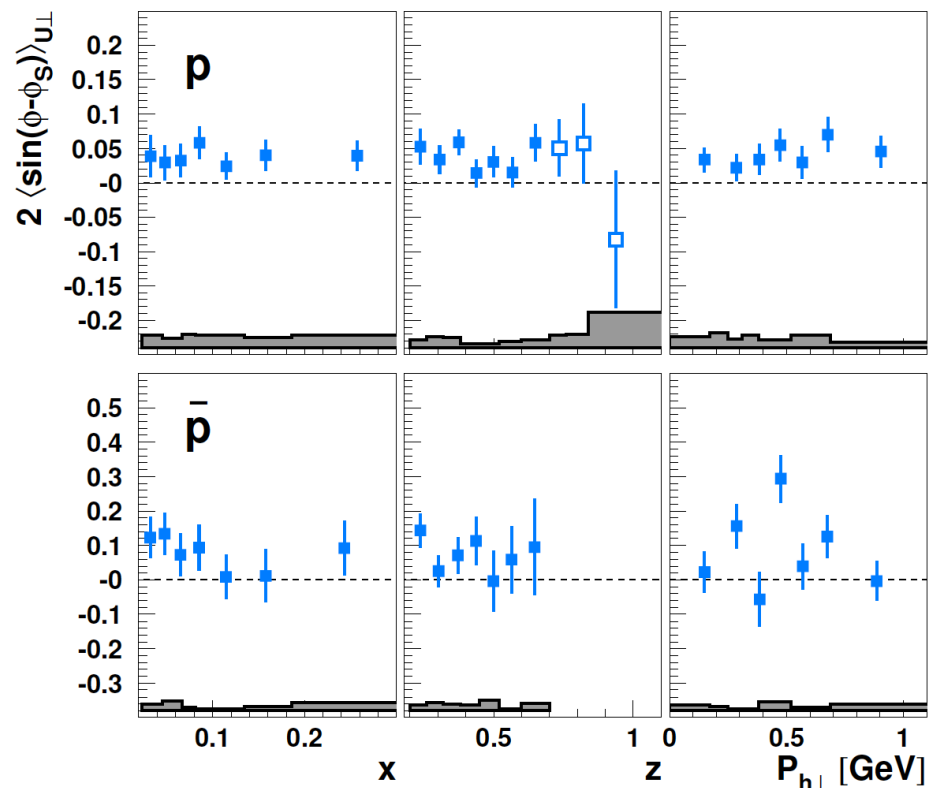
Similar agreement between  $\bar{p}$  and  $\pi^+$  (but with larger statistical errors)

A naive fragmentation process that can lead to  $p/\bar{p}$ :



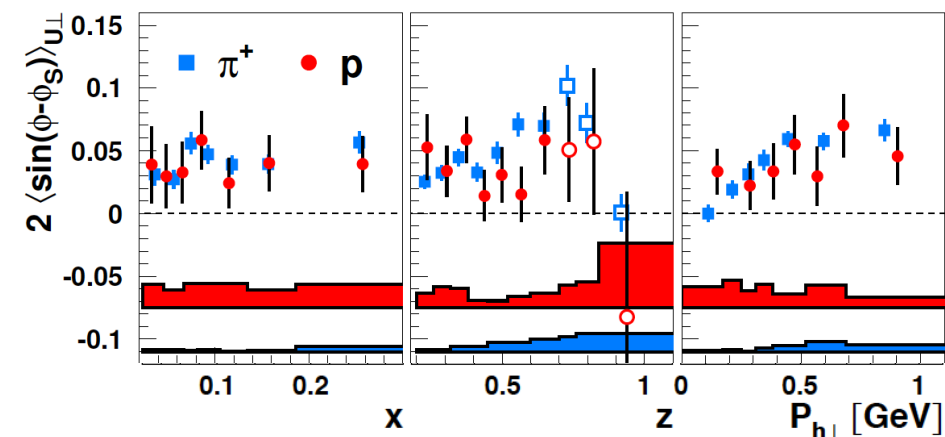
Let's assume scattering off the up quark (dominance of  $u$ -quarks in  $p/\bar{p}$  production supported by global fits of FF [[Phys.Rev.D76:074033,2007](#)])

# Sivers amplitudes: protons results



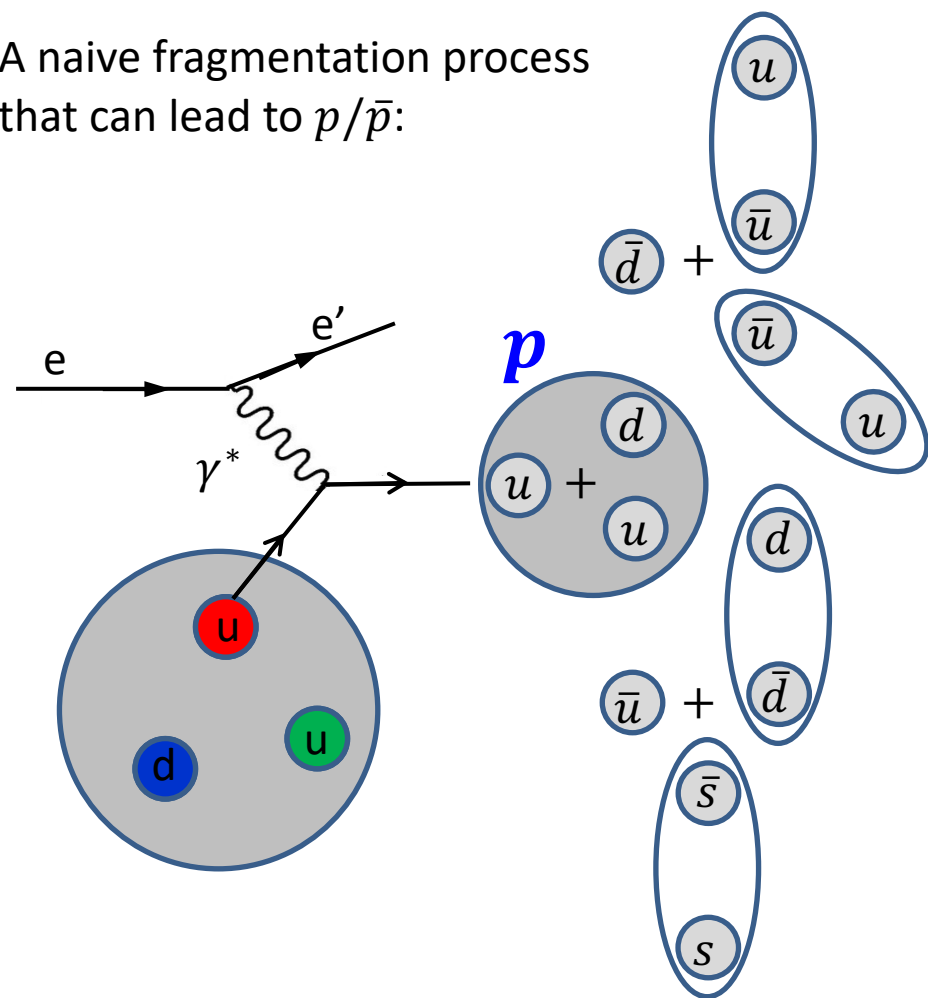
First measurement of Sivers asymmetries for  $p, \bar{p}$  in SIDIS

Both amplitudes are non-zero and positive



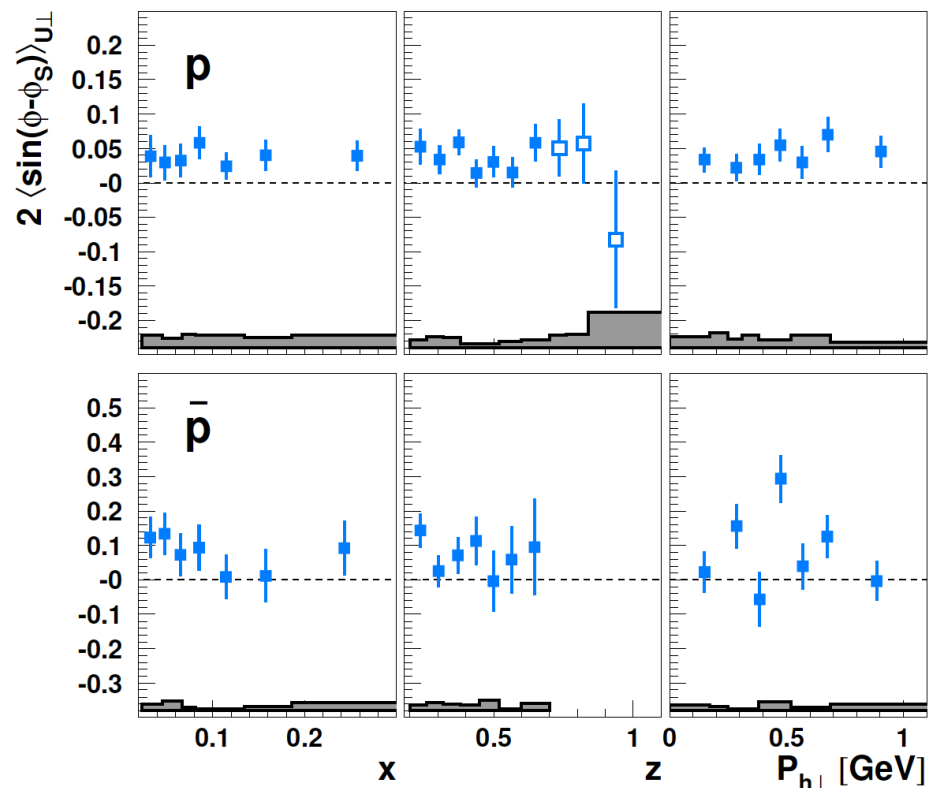
Similar agreement between  $\bar{p}$  and  $\pi^+$  (but with larger statistical errors)

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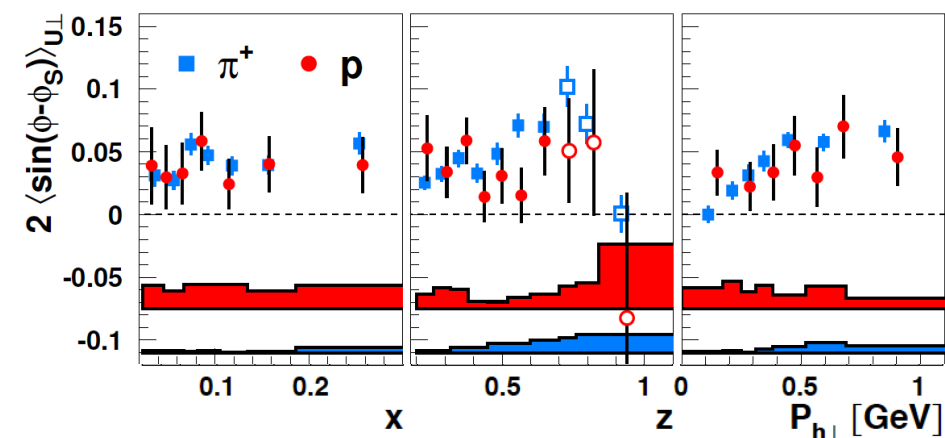
Let's assume scattering off the up quark (dominance of  $u$ -quarks in  $p/\bar{p}$  production supported by global fits of FF [[Phys.Rev.D76:074033,2007](#)])

# Sivers amplitudes: protons results



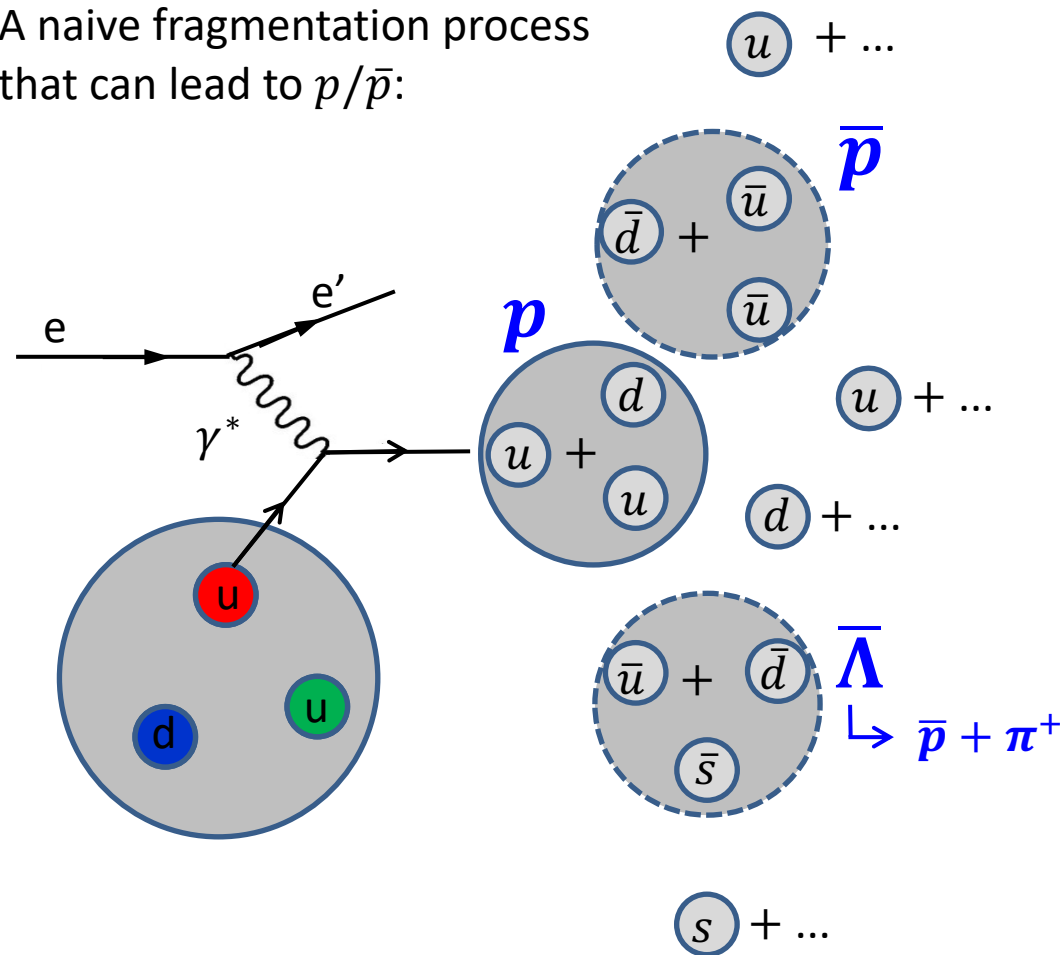
First measurement of Sivers asymmetries for  $p, \bar{p}$  in SIDIS

Both amplitudes are non-zero and positive



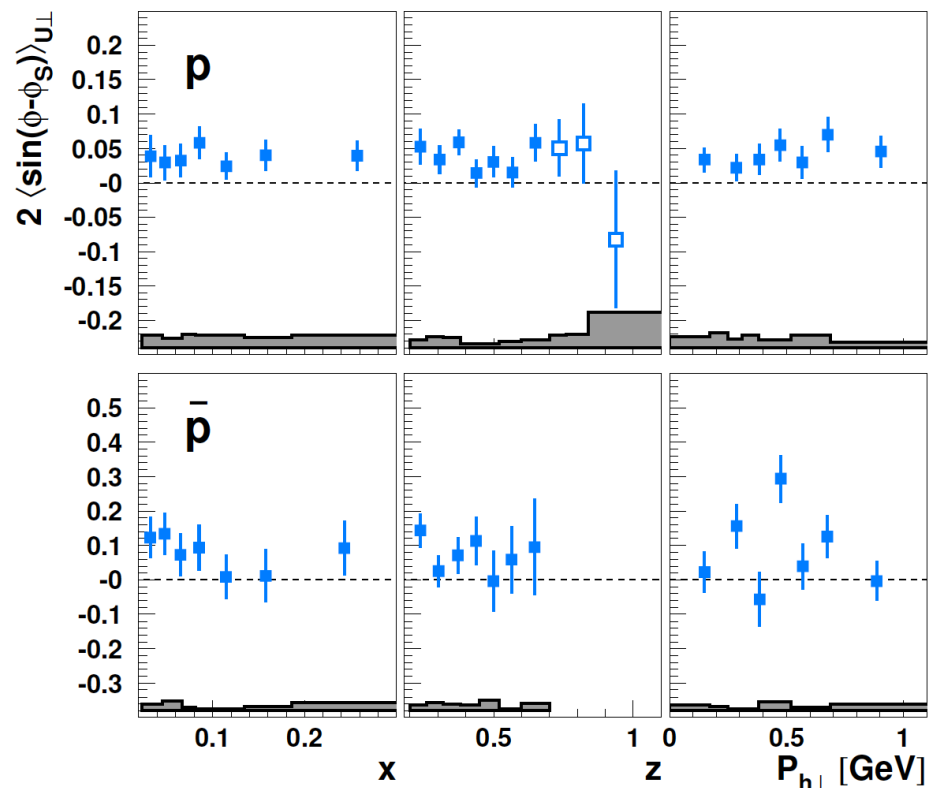
Similar agreement between  $\bar{p}$  and  $\pi^+$  (but with larger statistical errors)

A naive fragmentation process that can lead to  $p/\bar{p}$ :



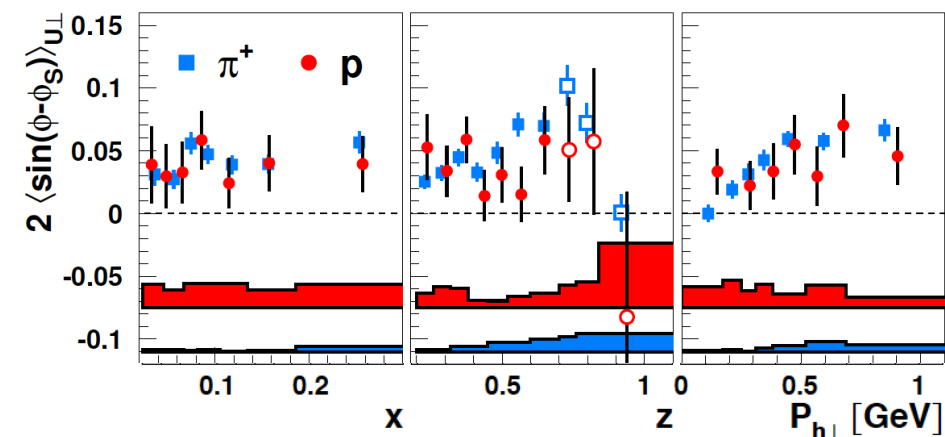
Let's assume scattering off the up quark (dominance of  $u$ -quarks in  $p/\bar{p}$  production supported by global fits of FF [[Phys.Rev.D76:074033,2007](#)])

# Sivers amplitudes: protons results



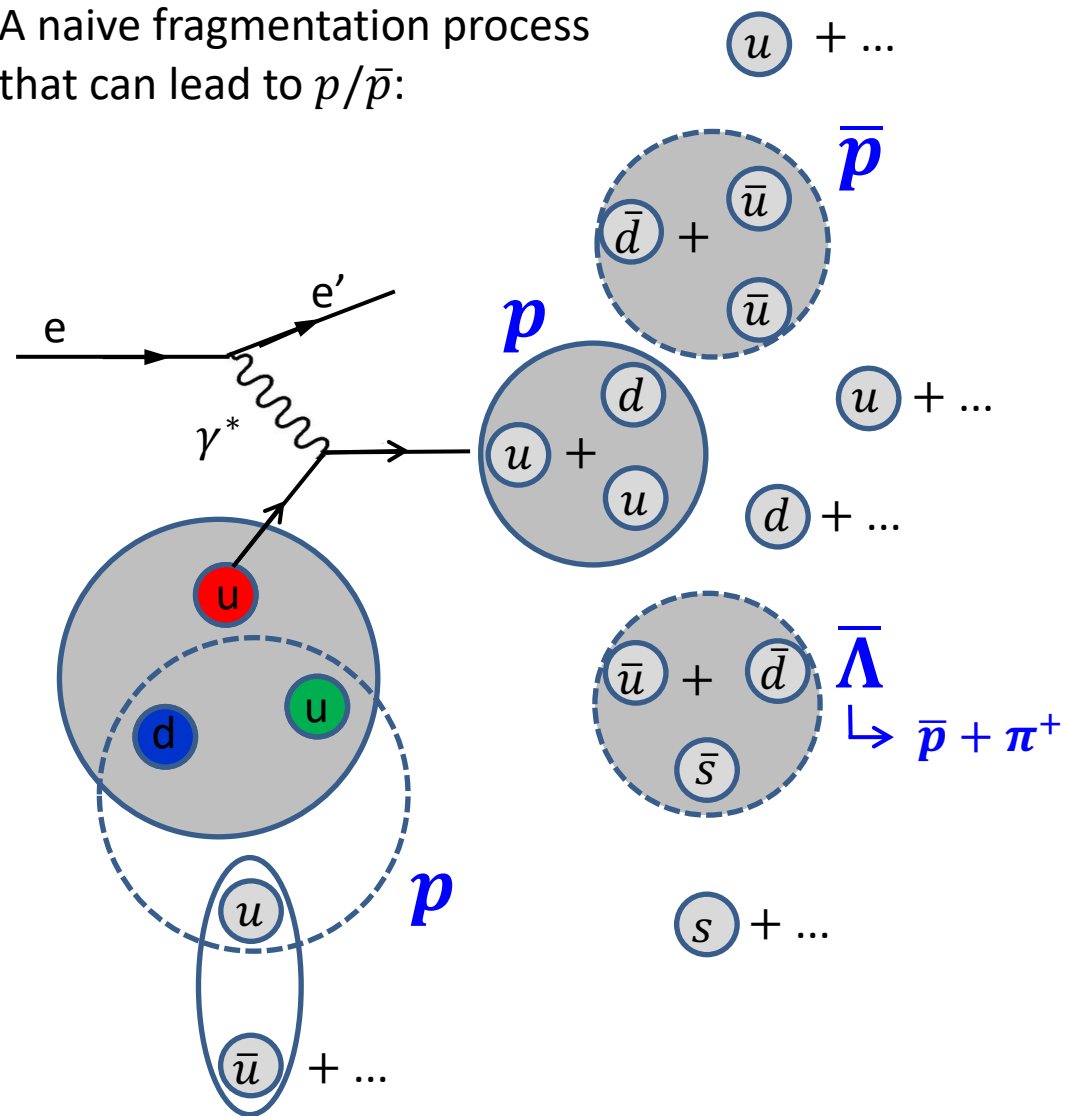
First measurement of Sivers asymmetries for  $p, \bar{p}$  in SIDIS

Both amplitudes are non-zero and positive



Similar agreement between  $\bar{p}$  and  $\pi^+$  (but with larger statistical errors)

A naive fragmentation process that can lead to  $p/\bar{p}$ :



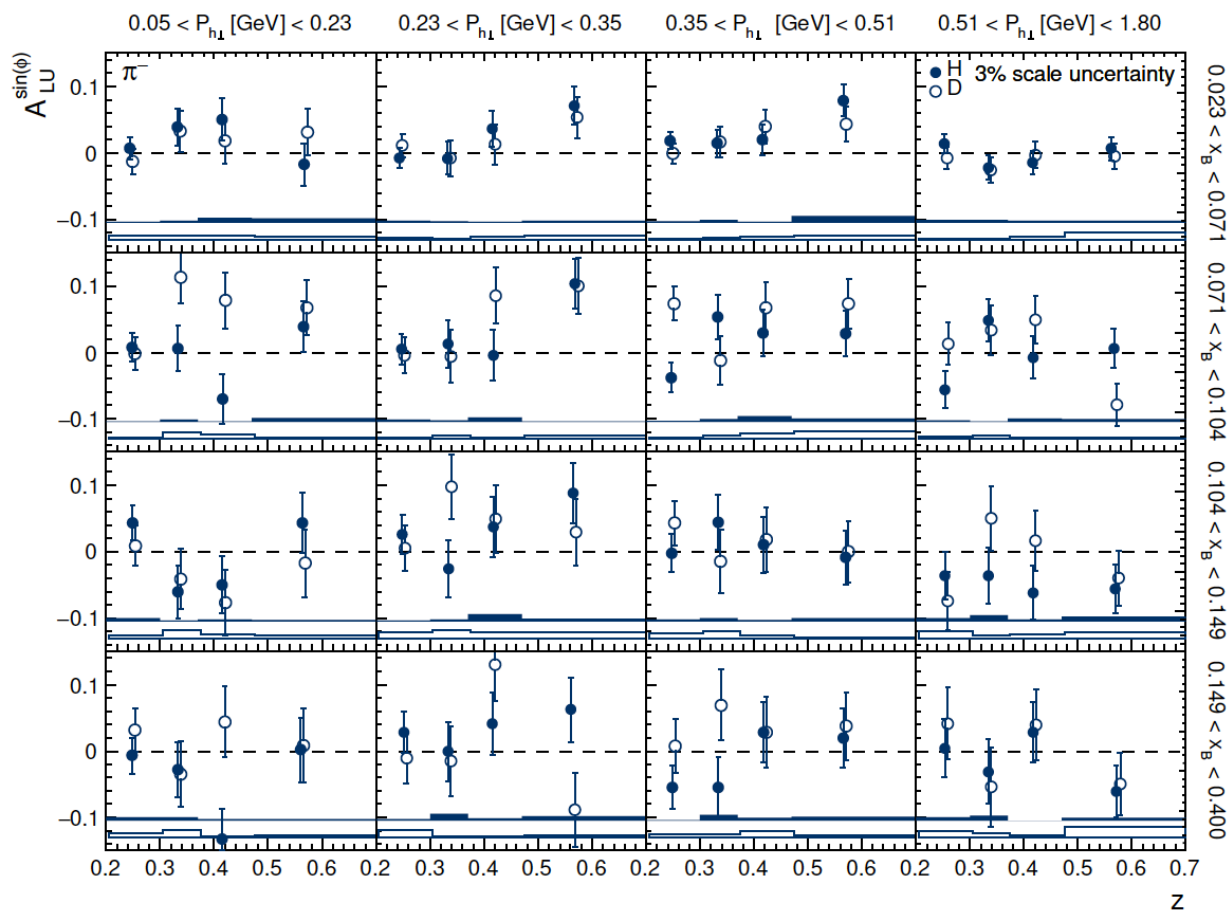
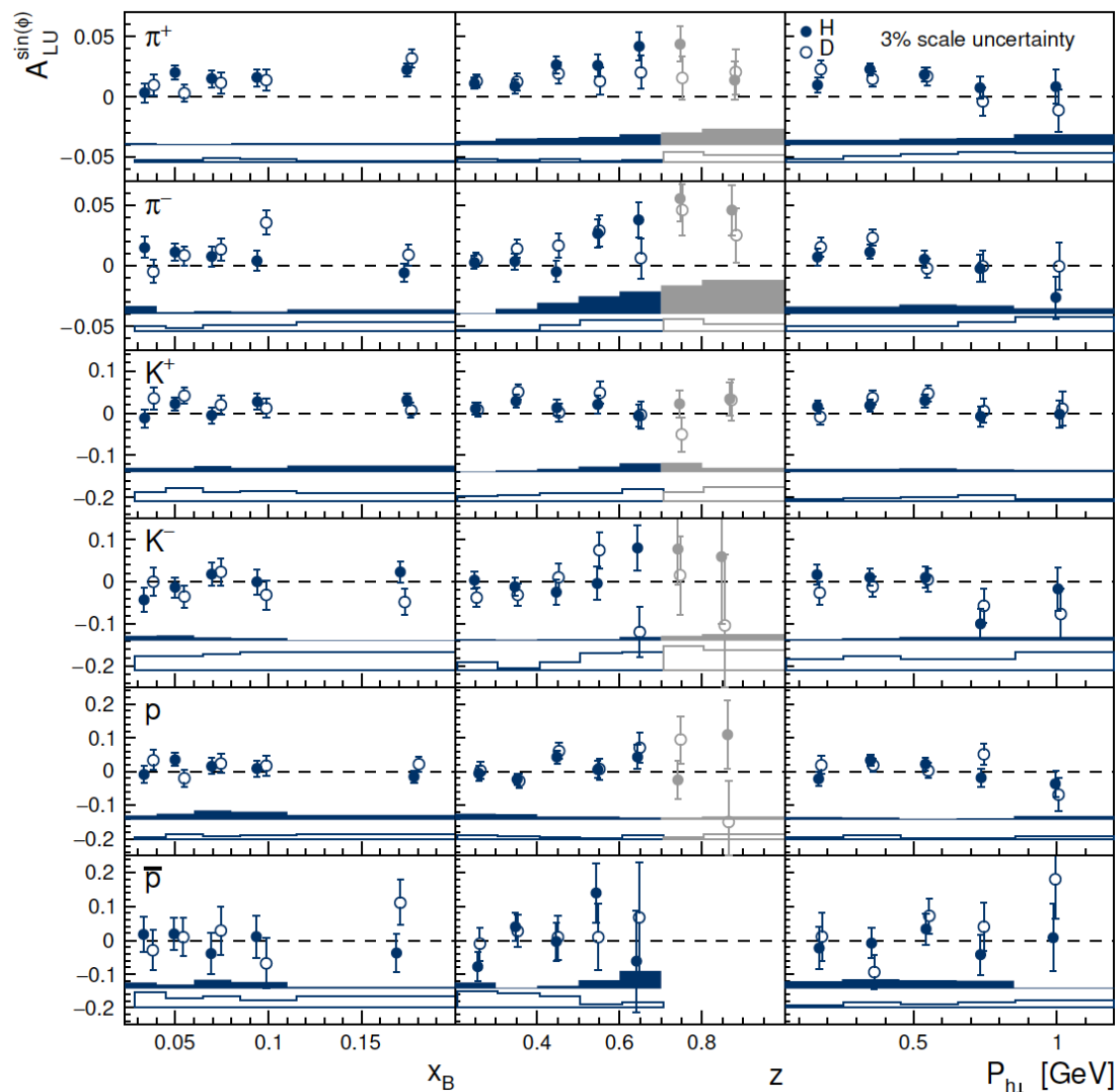
...also from TFR (low  $z$ , high  $P_{h\perp}$ )

# Other HERMES results



# Sub-leading twist $\sin(\phi)$ BSA

*Phys. Lett. B* 797 (2019) 134886

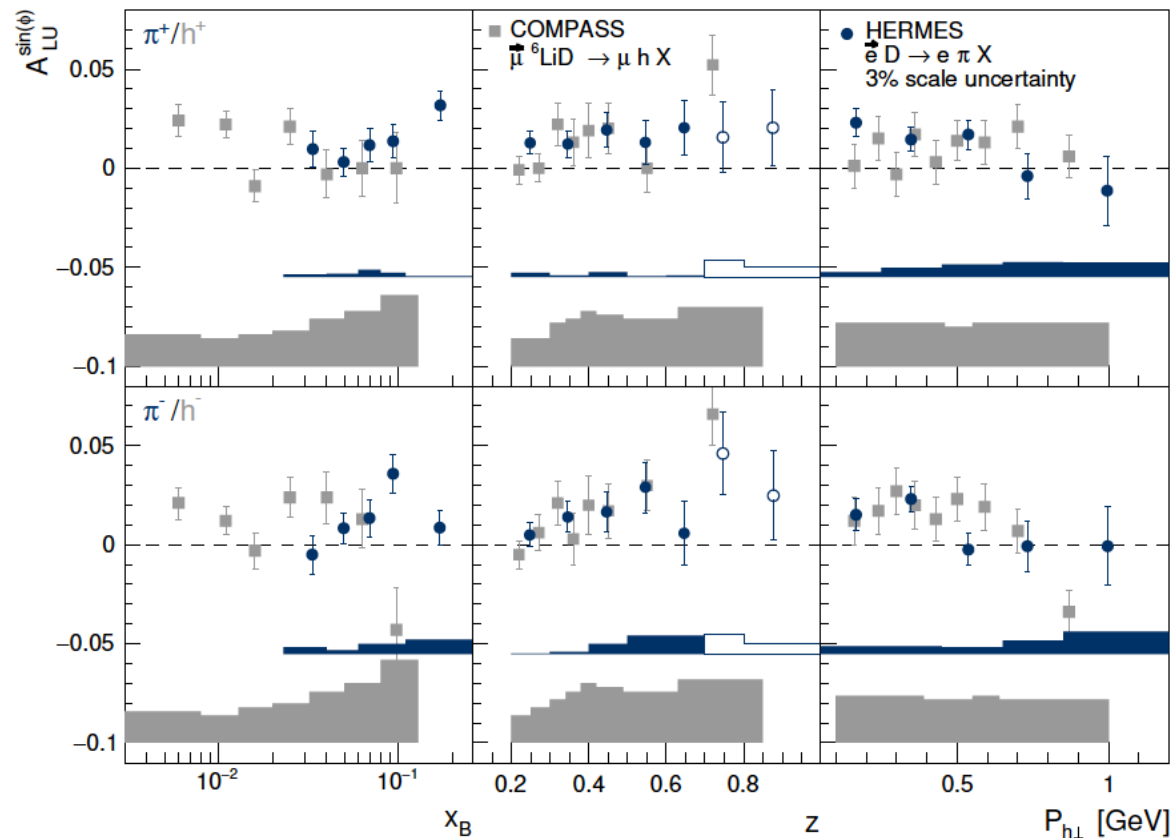


- Large positive amplitudes rising with  $z$  for  $\pi^+$  and  $\pi^-$
- Small positive amplitude with mild kinematic dep. for  $K^+$
- Results compatible with zero for  $K^-$ ,  $p$  and  $\bar{p}$

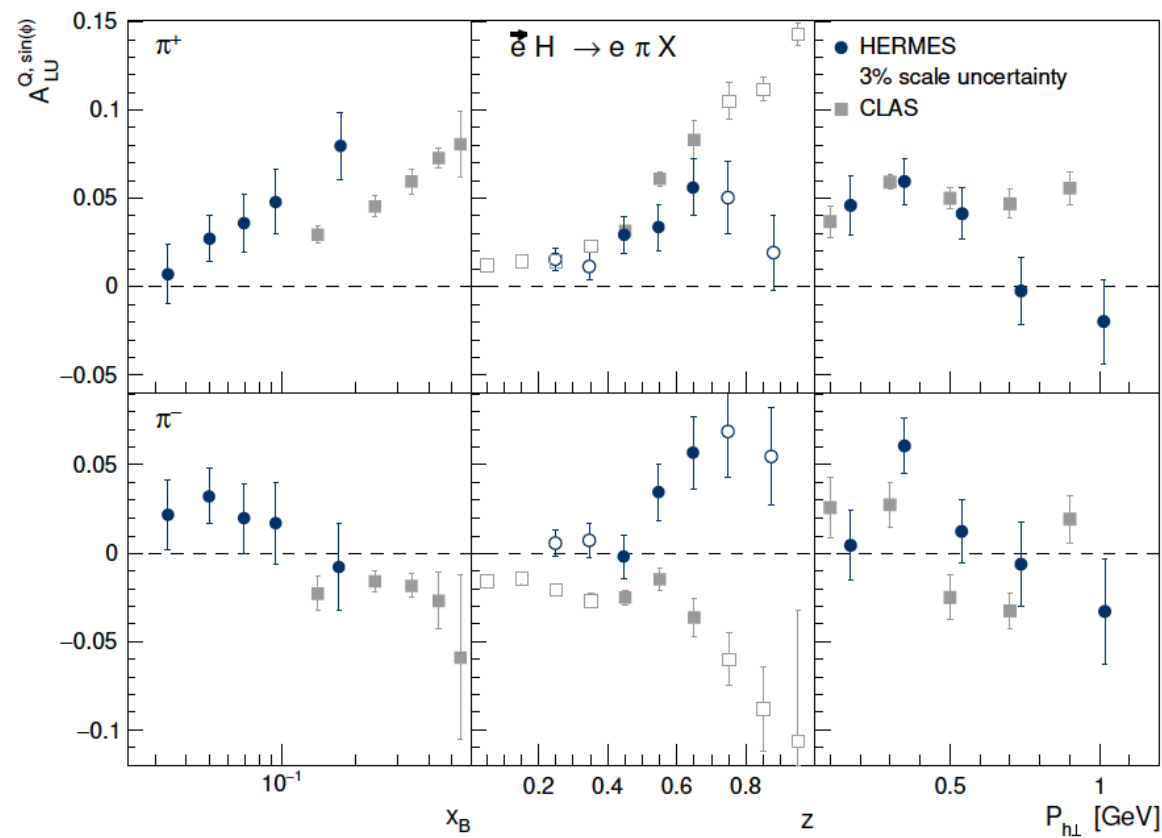
# Sub-leading twist $\sin(\phi)$ BSA

Phys. Lett. B 797 (2019) 134886

HERMES vs. COMPASS

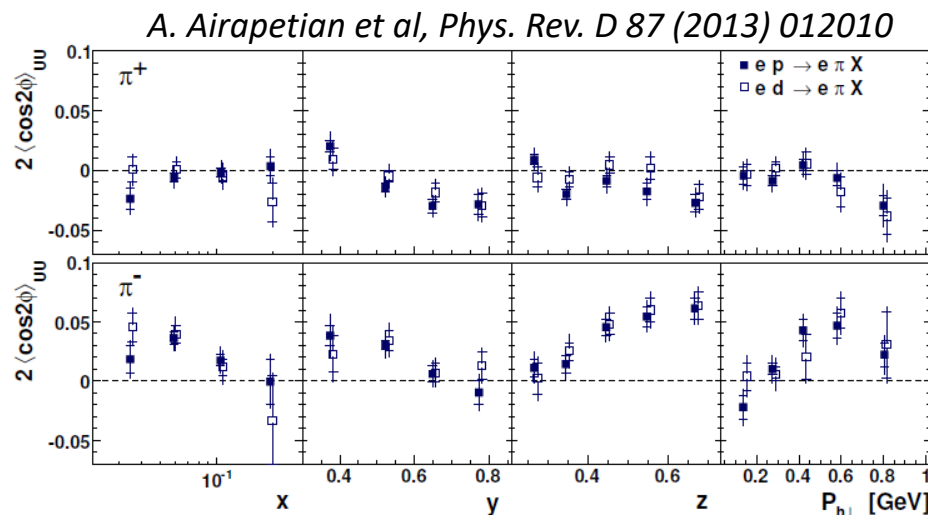


HERMES vs. CLAS



Sign change with increasing  $x$ ?

# The $\cos 2\phi$ amplitudes $\propto h_1^\perp(x, p_T^2) \otimes H_1^\perp(z, k_T^2)$



negative

positive

- Amplitudes are significant

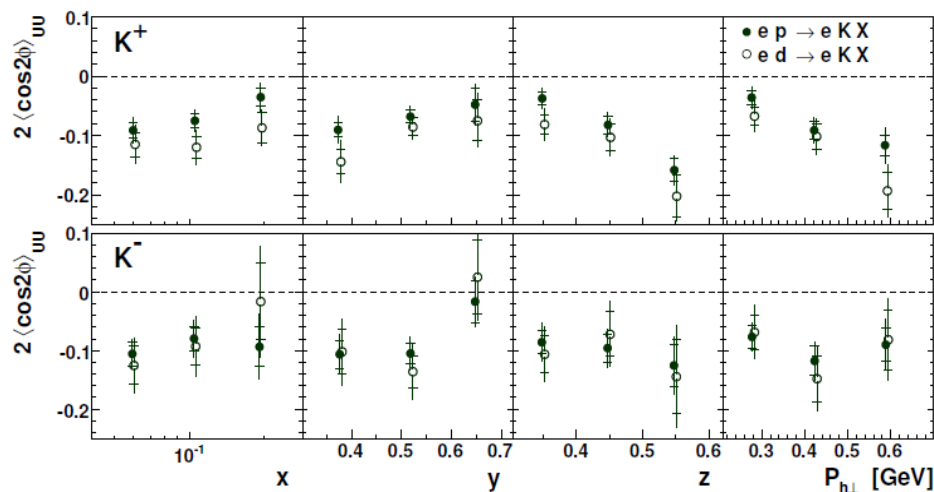
→ evidence of BM effect

- similar results for H & D

→  $h_1^{\perp,u} \approx h_1^{\perp,d}$

- Opposite sign for  $\pi^+/\pi^-$

→ opposite signs of fav/unfav  
Collins FF



Large  
and  
negative

Large  
and  
negative

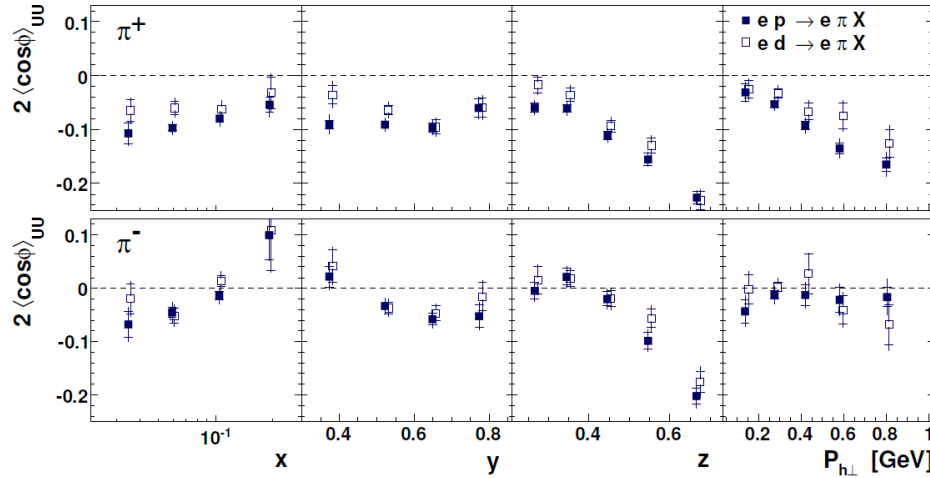
-  $K^+/K^-$  amplitudes larger  
than for pions, have different  
kinematic dependencies than  
pions and are both negative

→ different role of Collins FF  
for pions and kaons?

→ significant contribution from  
scattering off strange quarks?

# The $\cos\phi$ amplitudes $\propto +\frac{1}{Q} [h_1^\perp \otimes H_1^\perp + f_1 \otimes D_1 \dots]$

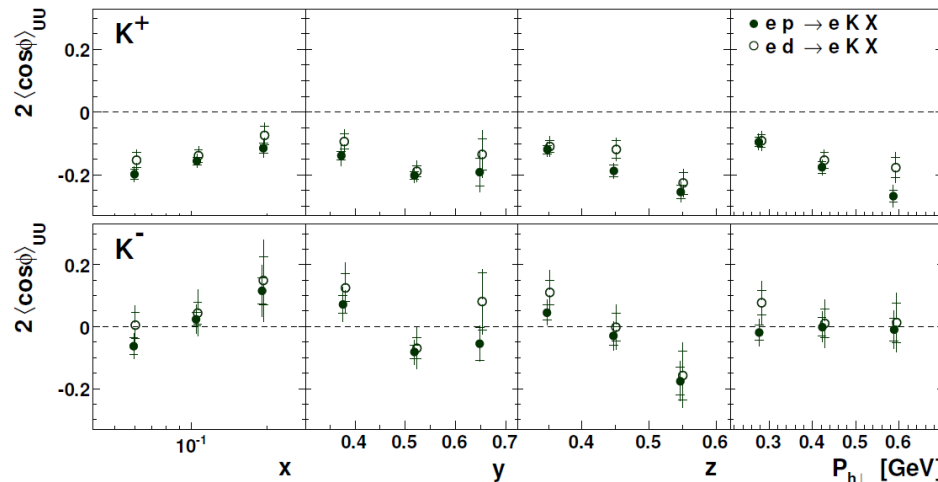
A. Airapetian et al, Phys. Rev. D 87 (2013) 012010



negative

negative

- Significant and of same sign → Chan effect weakly flavor dependent?
- Clear rise with  $z$  for  $\pi^+$  &  $\pi^-$  and  $P_{h\perp}$  for  $\pi^+$
- Different  $P_{h\perp}$  dependence → contrib. of flavor dependent effects (e.g. BM) for  $\pi^-$ ?



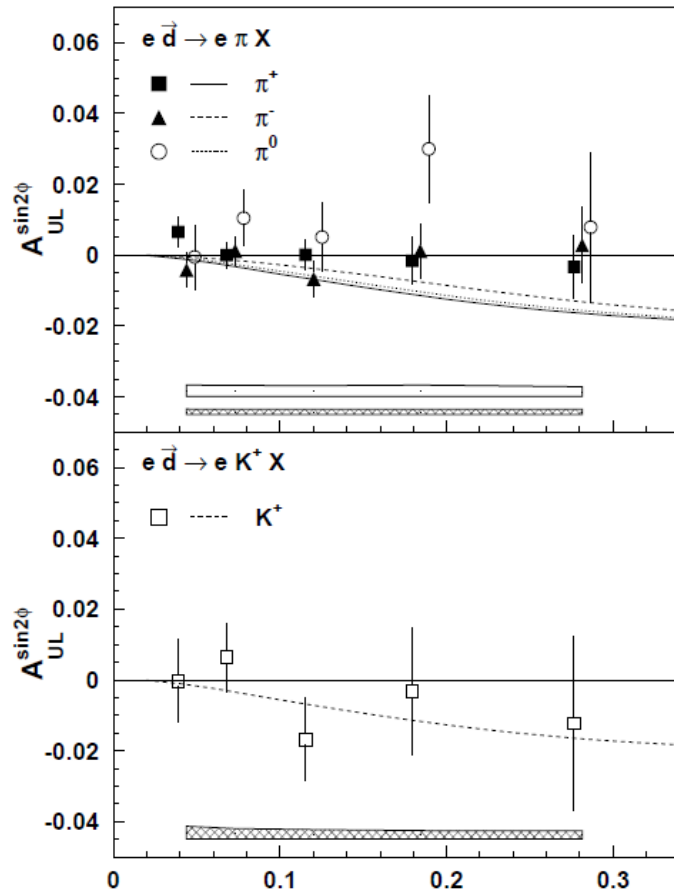
Large and negative

Consist. with 0

- $K^+$  amplitudes larger than  $\pi^+$  → different Collins FF for  $\pi$  &  $K$
- $K^- \approx 0$  different than  $K^+$  (in contrast to  $\cos 2\phi$ )
- Significant contrib from interaction dependent terms?

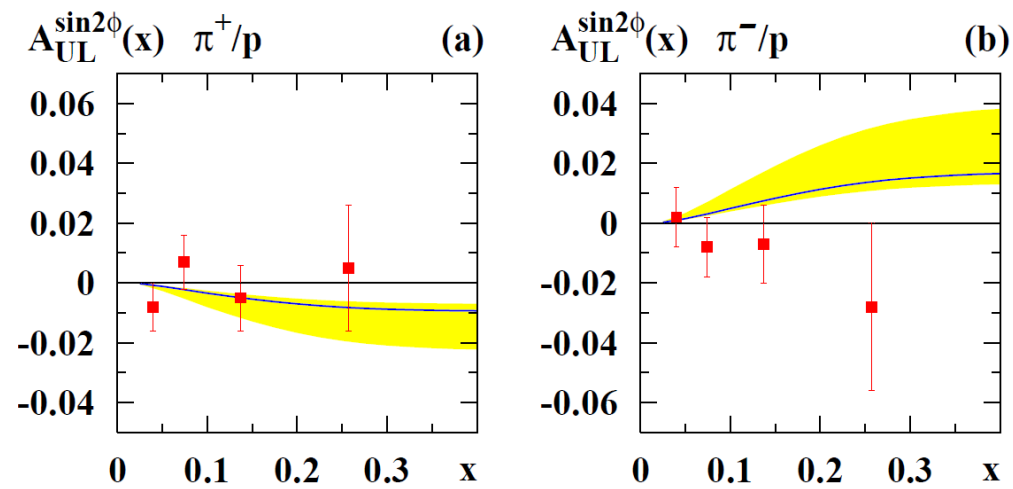
# The $\sin(2\phi)$ amplitude $\propto h_{1L}^\perp(x, p_T^2) \otimes H_1^\perp(z, k_T^2)$

Deuterium target



A. Airapetian et al, *Phys. Lett. B* 562 (2003)

Hydrogen target



A. Airapetian et al, *Phys. Rev. Lett.* 84 (2000)

Amplitudes consistent with zero for all mesons and for both H and D targets