Inclusive Photoproduction of Quarkonium-pair at the US EIC

Sangem Rajesh

INFN Perugia

In collaboration with

F. A. Ceccopieri, J. P Lansberg, M. Rinaldi and H. S. Shao





July 6-13th, 2022, Bologna, Italy

Outline

- **Quarkonium production**
- Double parton scattering
- $\hfill\square$ Di-J/ ψ SPS and DPS photoproduction at EIC
- □ Summary

Three models for describing the formation of Quarkonium and they are successful at different regions

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Quarkonium-pair measurment

• Only after 8 years of J/ψ discovery

For the first time, di- J/ψ was measured by NA3 collaboration

Phys.Lett.B 114 (1982) 457-460 Phys. Lett. B 158 (1985) 85-91

EVIDENCE FOR $\psi\psi$ PRODUCTION IN π^- INTERACTIONS AT 150 AND 280 GeV/c

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J. BADIER^d, J. BOUCROT^e, J. BOUROTTE^d, G. BURGUN^a, O. CALLOT^e, Ph. CHARPENTIER^a, M. CROZON^c, D. DECAMP^e, P. DELPIERRE^c, B. GANDOIS^a, R. HAGELBERG^b, M. HANSROUL^b, Y. KARYOTAKIS^d, W. KIENZLE^b, P. Le DÛ^a, J. LEFRANÇOIS^c, Th. LERAY^c, J. MAILLARD^c, A. MICHELINI^b, Ph. MINÉ^d, G. RAHAL^a, O. RUNOLFSSON^b, P. SIEGRIST^a, A. TILQUIN^c, J. VALENTIN^c, R. VANDERHAGHEN^d and S. WEISZ^d ^a *CEN*, Saclay, France ^b CERN, Geneva, Switzerland ^c College de France, Paris, France ^d Ecole Polytechnique, Palaiseau, France ^e Laboratoire de l'Accélérateur Linéaire, Orsay, France Received 27 May 1982

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Later at LHC di- J/ψ

- LHCb : PLB 707 (2012) 52–59, JHEP 06 (2017) 047
- D0 : PRD 90 (2014) 11, 111101
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- ATLAS : EPJC 77 (2) (2017) 76
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di- Υ at LHC

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CSM LO theoretical computation

V.G. Kartvelishvili and Sh.M. Esakiya Yad.Fiz. 38 (1983) 722-726

- B. Humpert and P. Mery Z. Phys. C 20 (1983) 83
- C.F. Qiao, Phys. Rev. D 66 (2002) 057504

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Puzzle at large Δy (M)



JP Lansberg, HS Shao, N Yamanaka, and YJ Zhang, EPJC 79 (2019) 12, 1006, PLB 807 (2020), 135559

JP Lansberg, HS Shao, NPB 900 (2015) 273–294, PLB 751 (2015) 479–486

H. S. Shao and Y. J. Zhang, PRL 117 (2016) 6, 062001 J. P. Lansberg, Phys. Rept. 889 (2020), 1-106

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• A simple solution to this excess is the independent production of di- J/ψ

(Double Parton Scattering)

DPS and the LHC data



JP Lansberg, HS Shao, NPB 900 (2015) 273-294

CMS coll. JHEP 09 (2014) 094

• DPS is the simplest explanation for the gap between SPS prediction and the data at large Δy and $M_{\psi\psi}$

DPS and the LHC data



JP Lansberg, HS Shao, NPB 900 (2015) 273-294

CMS coll. JHEP 09 (2014) 094

- DPS is the simplest explanation for the gap between SPS prediction and the data at large Δy and $M_{\psi\psi}$
- Same observation with the ATLAS data





7





However, differential cross section of DPS could be of the same order as that of SPS

7

• Phenomenologically, pocket formula is a tool for DPS cross section

$$\sigma_{\psi\psi}^{DPS} = \frac{1}{2} \frac{\sigma_{\psi}\sigma_{\psi}}{\sigma_{eff}}$$

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• D0 $\sigma_{eff} = 4.8 \pm 2.5 \text{ mb}$

Do coll. Phys.Rev.D 90 (2014) 11, 111101

• CMS $\sigma_{eff} = 8.2 \pm 2.0 \pm 2.9 \text{ mb}$

JP Lansberg, H.-S.Shao PLB 751 (2015) 479, CMS coll. JHEP 09 (2014) 094

• ATLAS $\sigma_{eff} = 6.3 \pm 1.6 \pm 1.0 \text{ mb}$

ATLAS coll. EPJC (2017) 77:76

DPS in Photoproduction(HERA)



S. Chekanov et al., Nucl.Phys.B 792 (2008) 1-47

• 3, 4-jet photoproduction cross sections are better described when MPI are included (Indication of Double Parton Scattering even in photoproduction)



Illustration of DPS for $\gamma + p \rightarrow J/\psi + J/\psi + X$

We consider the possibility of **resolved** photon to estimate the DPS cross section in quarkonium-pair photoproduction at the EIC 10

(Unresolved/direct)
$$\sigma_{SPS}^{(J/\psi, J/\psi)} \propto \sum_{a=g,q} \int dx_{p_a} f_{a/p}(x_{p_a}, \mu) \, d\hat{\sigma}^{\gamma a \to J/\psi + J/\psi + a}$$

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Proton PDF

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Photon PDF Proton PDF

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Single J/ψ SPS resolved (namely same partonic cross section as hadroproduction)





- GRV photon PDF is used PRD 46, 1973 (1992), while CT18NLO PDF for proton T.J. Hou et al., PRD 103, 014013 (2021)
- HELAC-Onia latest version is used for generating matrix elements HS Shao, CPC 184, 2562 (2013), 198, 238 (2016)
- CO LDMEs are taken from M. Butenschoen and B. A. Kniehl, PRD 84, 051501 (2011)
- We expect at least 600 four-muon events with 100 fb⁻¹ luminosity



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Range of cross sections in CSM
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 $\sigma_{SPS}^{(J/\psi, J/\psi)} \times Br^2 = 4 - 30 \text{ fb}$ (Resolved) $\sigma_{eff}^{\gamma p} = 10 \text{ mb for DPS}$
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$\Delta y \text{ and } M_{\psi\psi} \text{ at EIC}$

Kinematical distributions useful to separate SPS and DPS



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LDME set (HU): M. Butenschoen and B. A. Kniehl, PRD 84, 051501 (2011)

z distribution



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z distribution



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- For z < 0.1 resolved SPS dominates over unresolved/direct
- Unique opportunity to study the photon structure
- At larger z one can test quarkonium production mechanism via direct photoproduction
- Resolved case: gluon channel dominates in the low *z* region, and quark channel at high *z*
- CS and CO states are considered: CO states contribution is only significant (for some LDMEs) in unresolved but not in the resolved case

Summary

- Quarkonium production is a rich channel to probe the parton correlations through DPS
- We have estimated SPS and DPS cross sections for quarkonium-pair photoproduction at the EIC using the NRQCD framework
- > DPS total cross section is small compared to the SPS but could be measured if σ_{eff} small
- > Differential spectra in $|\Delta y(\psi, \psi)|$ and $M_{\psi\psi}$ are useful to separate DPS and SPS
- \triangleright Differential spectrum in z is useful to separate resolved and unresolved
- \succ Bulk of the yield at low *z* where resolved SPS is dominant
- Quarkonium-pair photoproduction is a promising channel to probe the gluonic content of the photon structure

Thank you for the attention

Backup slides

Effective X-section

JP Lansberg's slide MPI 19



z-spectrum (Resolved)









z-spectrum (glu-quark)







19

Pocket formula details

 $f_{i/\gamma}(x;\mu^2) = K \alpha f_{i/\pi}(x;\mu^2)$ Gluck et al, PRD 46, 1973 (1992) How to build up the photon dPDFs from the PDF? 091 If we assume that evolved dPDFs can be written as product of PDFs: $\tilde{F}_{ij/\gamma}(x_1, x_2, k_{\perp}; \mu^2) = N f_{i/\gamma}(x_1; \mu^2) f_{j/\gamma}(x_2; \mu^2) \tilde{F}_{\gamma}(k_{\perp}) \theta(1 - x_1 - x_2)$ Normalization Effective form factor constant M. Rinaldi and F. A. Ceccopieri, PRD 97 (2018) 7, 071501 Then by using sum rules: $\sum_{i} \int_{0}^{1-x_{1}} dx_{2} x_{2} \tilde{F}_{ji/\gamma}(x_{1}, x_{2}, k_{\perp} = 0; \mu^{2}) = (1-x_{1})f_{i/\gamma}(x_{1}; \mu^{2})$ J. R. Gaunt et al, JHEP 03, 005 (2010) Therefore, from our ansatz we get: $\sum_{\alpha} \int_{0}^{1-x_{1}} dx_{2} x_{2} \tilde{F}_{ji/\gamma}(x_{1}, x_{2}, k_{\perp} = 0; \mu^{2}) = NK\alpha f_{i/\gamma}(x_{1}; \mu^{2})(1-x_{1})$ Thus the normalization constant is $N = \frac{1}{K\alpha}$ The new pocket formula would be: $\sigma_{DPS}^{A+B} = \frac{m}{2\alpha K} \frac{\sigma_{SPS}^{A} \sigma_{SPS}^{D}}{\sigma_{cc}^{\gamma p}}$ Where the effective cross-section keeps the $\sigma_{eff}^{\gamma p} = \frac{1}{\int \frac{d^2 k_{\perp}}{(2\pi^2)} \tilde{F}_{\gamma}(\vec{k}_{\perp}) \tilde{F}_{p}(-\vec{k}_{\perp})}$ 20 usual geometrical interpretation, i.e.: