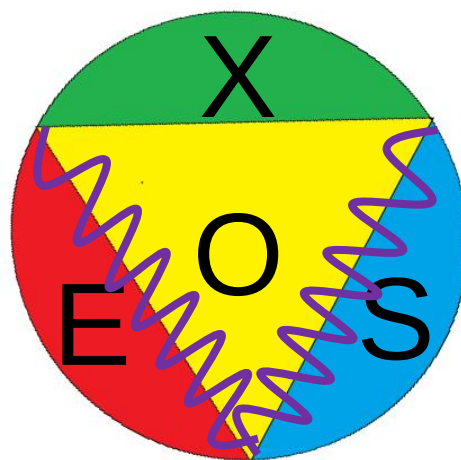


HAPG mosaic crystal Von Hamos spectrometer for high precision exotic atoms spectroscopy



Alessandro Scordo

Laboratori Nazionali di Frascati, INFN

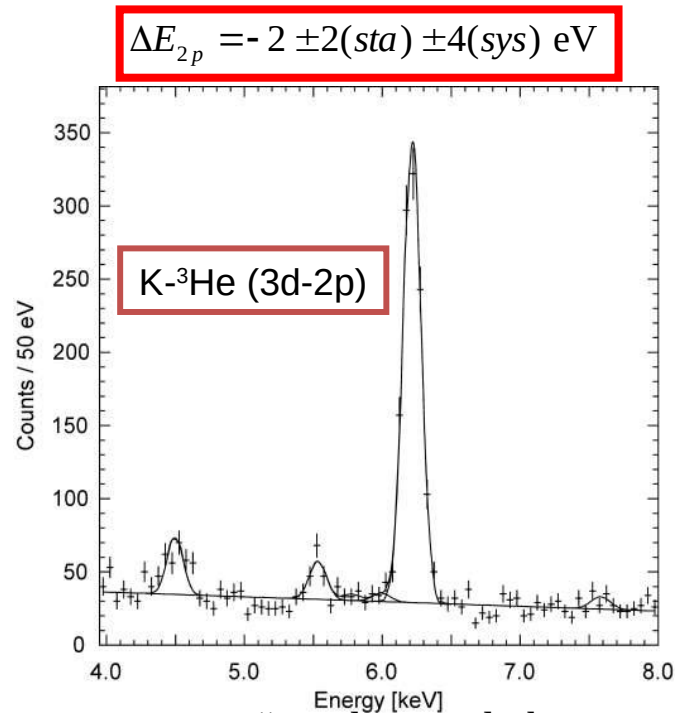
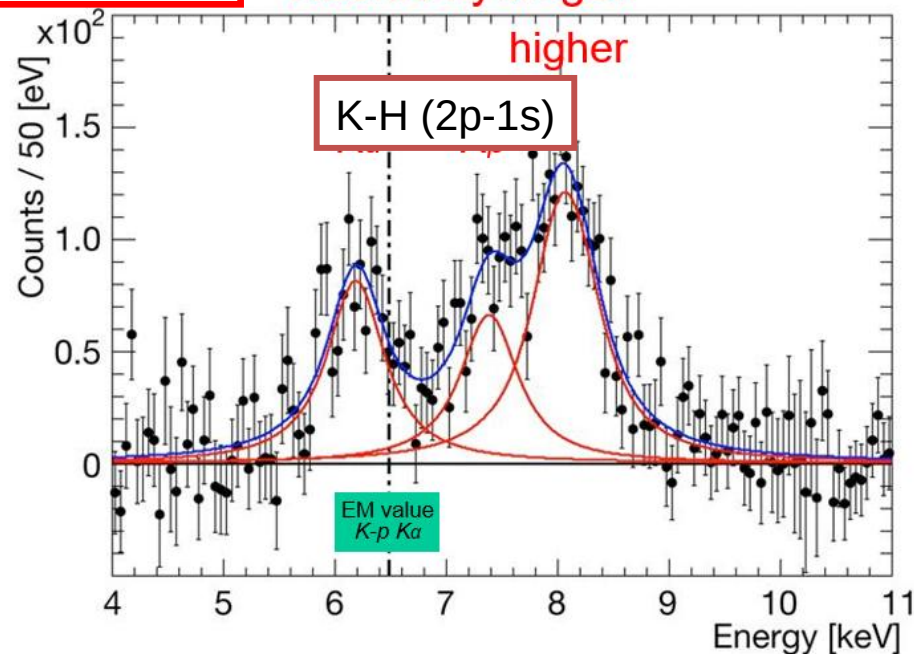
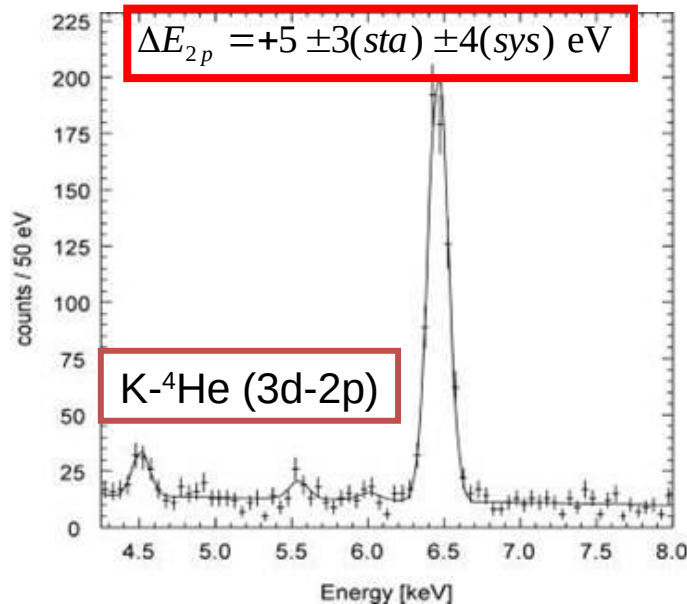
Latest results on kaonic atoms @ DAΦNE



$$\epsilon_{1s} = -283 \pm 36(\text{stat}) \pm 6(\text{syst}) \text{ eV}$$

$$\Gamma_{1s} = 541 \pm 89(\text{stat}) \pm 22(\text{syst}) \text{ eV}$$

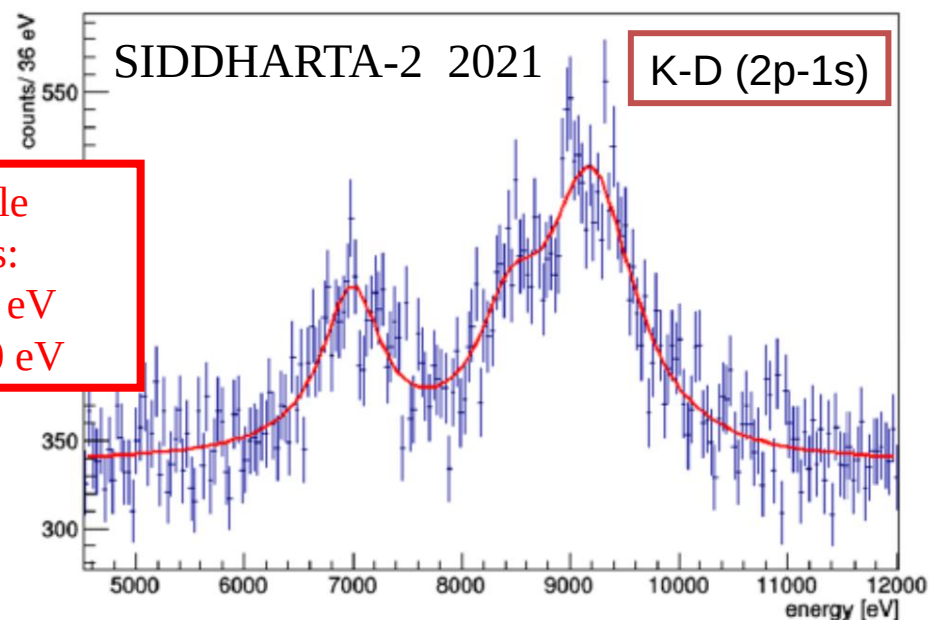
Kaonic hydrogen



Achievable
precisions:

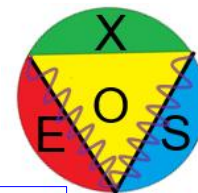
$$\Delta\epsilon(1s) = 30 \text{ eV}$$

$$\Delta\Gamma(1s) = 70 \text{ eV}$$



“Fundamental physics with exotic atoms and radiation detectors – LNF – 25-26/09/2021”

Needs for higher precisions



FWHM obtained in these measurements are already at the Fano limit for solid state detectors

Precisions of $1 \sim 50$ eV, depending on the statistics, can be reached with this FWHM

These values of FWHM and σE are not enough for many other measurements:

Example: KHe widths measured by SIDDHARTA

$$\Gamma_{2p}({}^3\text{He}) = 6 \pm 6 \text{ (stat.)} \pm 7 \text{ (syst.) eV}$$

$$\Gamma_{2p}({}^4\text{He}) = 14 \pm 8 \text{ (stat.)} \pm 5 \text{ (syst.) eV}$$

Example: Upper level measurements with very small Γ

An advantage of "upper levels"*

SŁAWOMIR WYCECH

In analogy to antiprotons the scenario under the $\bar{K}N$ threshold is determined by a resonant state $\Lambda(1405)$ with a pole close to E_{cm} 1410 MeV that is in the ${}^3\text{He}$ region. On the other side one has $\Sigma(1385)$ state which exerts maximum repulsive effect in the ${}^4\text{He}$ region. Apparently these two main agents yield attractive shift in ${}^3\text{He}$ and repulsive in ${}^4\text{He}$. Now, in order to go above the errors one has to magnify the shifts and enhance the atomic-nuclear overlaps. The proper targets would be ${}^8\text{Be}$ and ${}^6\text{Li}$. These offer similar values of E_{cm} as ${}^4\text{He}$ and ${}^3\text{He}$. A simple re-scaling of overlaps generates the level shifts of about 100 eV. One should perhaps consider also studies of 3D levels in these atoms. One interesting outcome might be the estimate where the isospin 0 $\text{Re } T(\bar{K}N \rightarrow \bar{K}N)$ amplitude crosses zero. That will help to settle the controversy as to where is the $\Lambda(1405)$ pole in the complex plane located.

Example: Kaon mass measurement

Charged Kaon Mass

Claude Amsler¹ and Simon Eidelman^{2,3,4}

¹Stefan Meyer Institute, Vienna, Austria

²Budker Institute of Nuclear Physics, Novosibirsk, Russia

³Novosibirsk State University, Novosibirsk, Russia

⁴Lebedev Physical Institute, Moscow, Russia

January 10, 2021

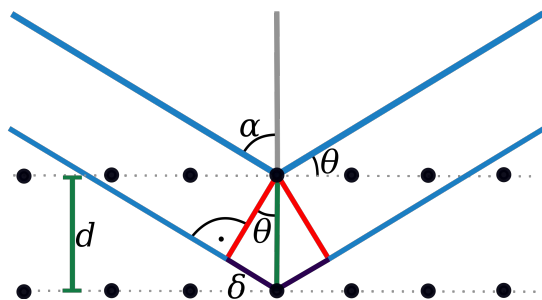
2003 [6]) was the first enigmatic state whose properties cannot be fully understood in the framework of the quark model. Despite very extensive efforts (the discovery paper with 1880 citations is one of the most cited experimental publications), there is no consensus today about its internal structure. The most popular explanation is that it is a mixture of a regular $q\bar{q}$ state and a $D^0\bar{D}^{*0}$ molecule. To test the validity of the molecular hypothesis it is of vital importance to know precisely how far the $\chi_{c1}(3872)$ state lies from the $D^0\bar{D}^{*0}$ threshold. Recently LHCb performed a study of $\chi_{c1}(3872)$ produced in decays of B^\pm mesons and other b hadrons [7, 8]. Using the world-largest sample of almost 20k $\chi_{c1}(3872) \rightarrow J/\psi\pi^+\pi^-$ decays, LHCb performed the most precise measurement of the $\chi_{c1}(3872)$ mass and of the energy difference $\delta E = m(D^0) + m(D^{*0}) - m(\chi_{c1}(3872)) = 0.07 \pm 0.12$ MeV. Again, the precision is limited by that of the charged kaon mass.

The precision on the D^0 mass also affects the mixing parameters in the $D^0\text{-}\bar{D}^0$ system [4], and in the long run, a more accurate kaon mass may become interesting for first-principle calculations on the lattice [9].

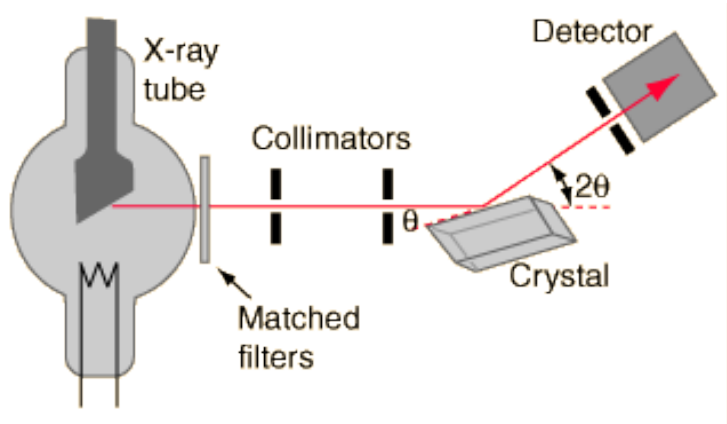
Example: Fine splitting of kaonic atoms
levels for cascade processes

"Fundamental physics with exotic atoms and radiation detectors – LNF – 25-26/09/2021"

Bragg spectroscopy



$$n\lambda = 2d\sin\theta_B$$



FWHM $\sim 1-10$ eV can be achieved depending on the quality of the crystal and the dimensions of the detectors

Natural background reduction from geometry

But....

- Small solid angles can be covered
- Typical efficiencies : $10^{-5} - 10^{-8}$
- Typical d (Si) ≈ 5.5 Å (good for $E < 6$ keV)
- Typical source size $10-100$ μm

The x-ray source, which was used for the measurements, is a low power microfocus x-ray tube (IfG) with a source diameter of about $50 \mu\text{m}$. Measurements were performed with the Cu K_α emission of a Cu anode at 8 keV. The spec-

III. SPECTROMETER SETUP

The spectrometer consists of three principal components: the X-ray source, the HAPG optic, and the position sensitive detector. As source a watercooled 100 W micro focus X-ray tube with a tungsten anode and a focus size of $50 \mu\text{m}$ is used. The emitted radiation is focused onto the sample by a polycapillary full lense with a spot size of $35 \mu\text{m}$. The HAPG

Laser-produced plasmas were created using the “Phoenix” Nd glass laser (the Lebedev Physical Institute) operated at a wavelength of $0.53 \mu\text{m}$ with pulse energy up to 10 J and 2 ns pulse duration. The laser beam was focused onto massive Mg, Al, Ti, or Fe targets (see Fig. 2). The focal spot diameter was about $\sim 15 \mu\text{m}$.



Mosaic crystal consist in a large number of nearly perfect small crystallites.

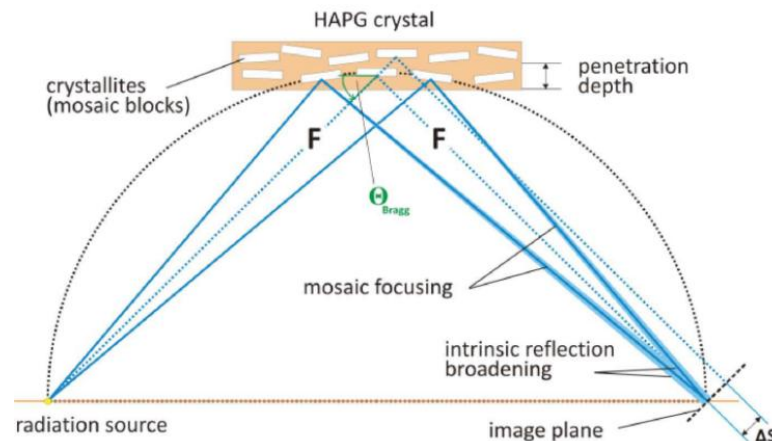
Mosaicity makes it possible that even for a fixed incidence angle on the crystal surface, an energetic distribution of photons can be reflected

Increase of efficiency
(focusing) ~ 50

Loss in resolution

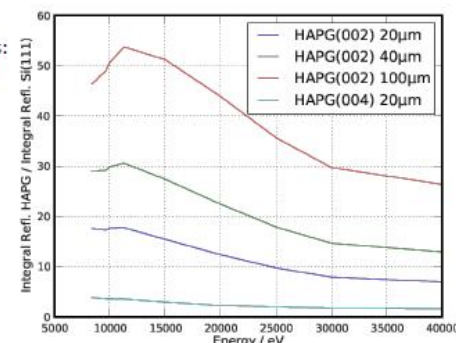
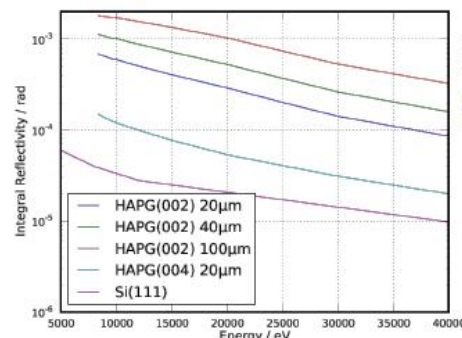
Pyrolytic Graphite mosaic crystals ($d = 3.354 \text{ \AA}$):

- Bending does not influence resolution and intensity
- Mosaic spread down to 0.05 degree
- Integral reflectivity $\sim 10^2$ higher than for other crystals
- Variable thickness (efficiency)
- Excellent thermal and radiation stability



Integral reflectivity

- Measured integral reflectivities (synchrotron measurements)



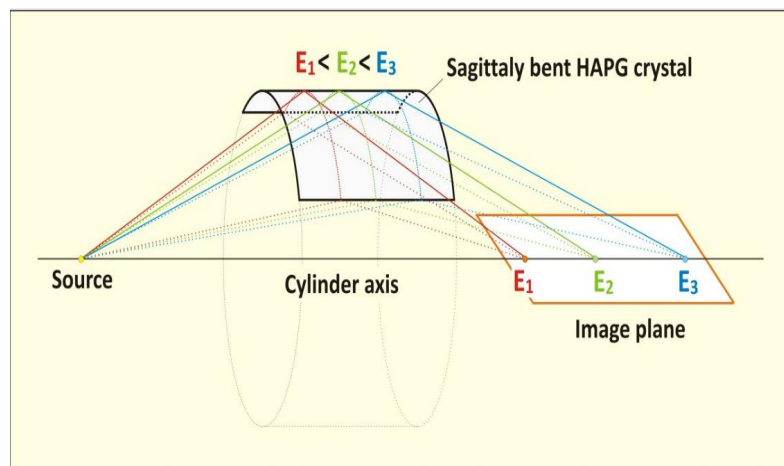
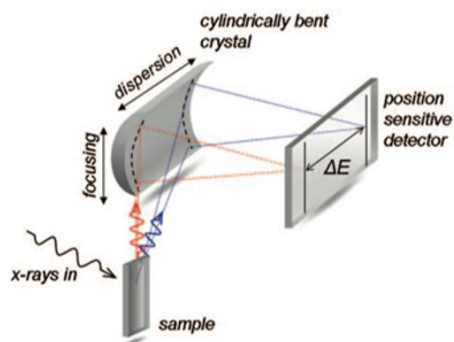
- The integral reflectivity can be more than 50 times higher compared to Si(111) reflection.
- The use of the von Hamos geometry can increase the overall efficiency even more.

Characterization of HAPG mosaic crystals using synchrotron radiation

Martin Gerlach,^a Lars Anklamm,^b Alexander Antonov,^c Inna Grigorieva,^c Ina Holfelder,^a Birgit Kanngießer,^b Herbert Legall,^c Wolfgang Malzer,^b Christopher Schlesiger^b and Burkhard Beckhoff^{a*}

J. Appl. Cryst. (2015). 48

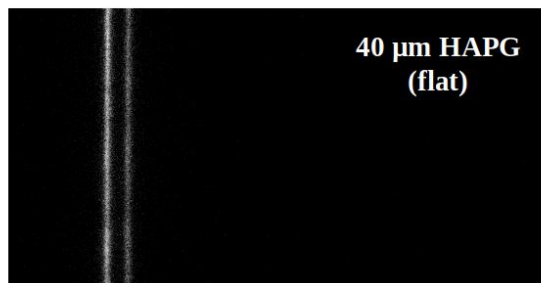
Von Hamos configuration: improving solid angle



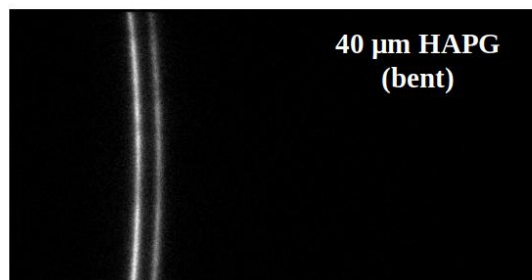
VH configuration can further improve the signal collection efficiency.

In this configuration, also the vertical dimension of the X-ray source can be exploited

distance: $F = 400$ mm in (004)-reflexion @ 8 keV (Cu K_α)

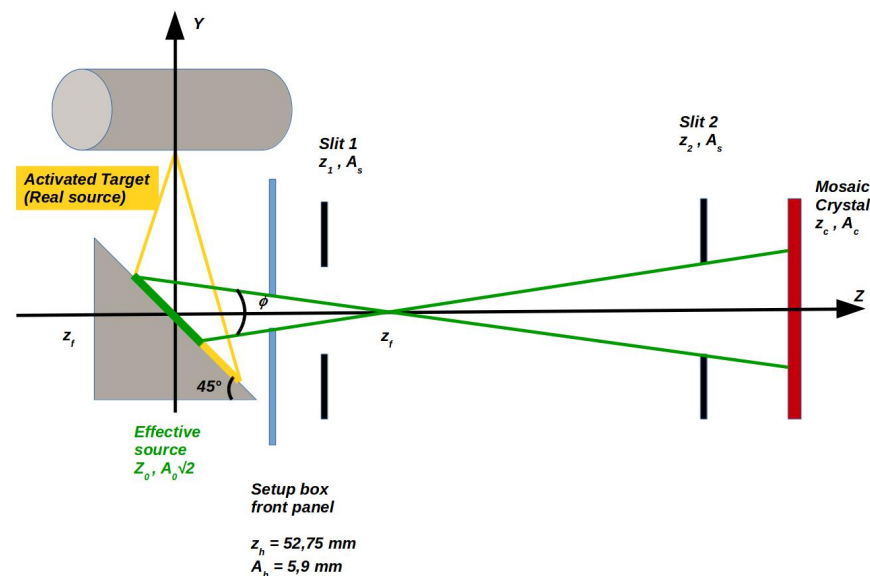


40 μ m HAPG
(flat)

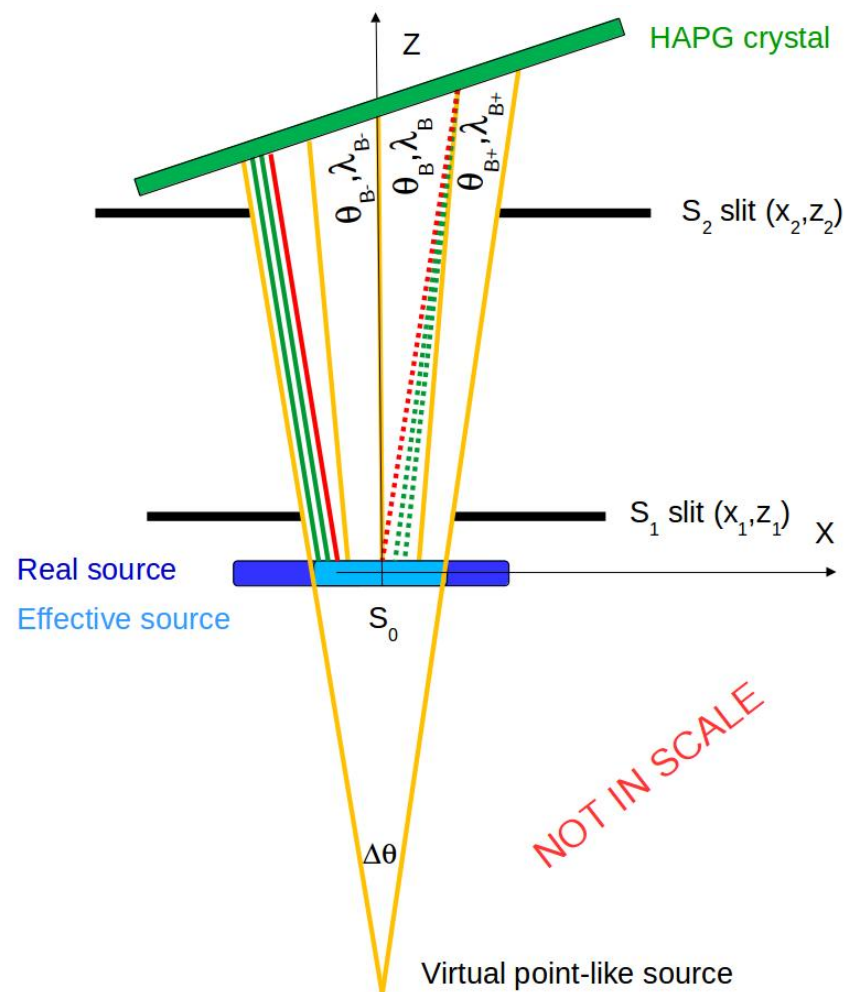
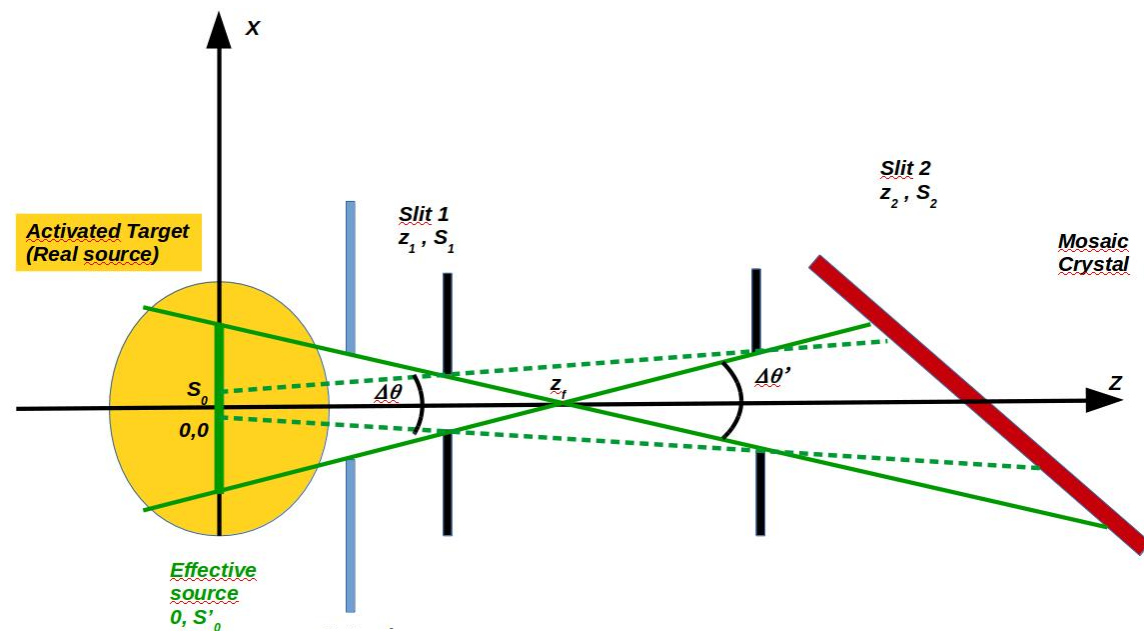


40 μ m HAPG
(bent)

Spectral resolution of bent HAPG/HOPG crystal is comparable to the flat one !



VOXES: enlarging the source size

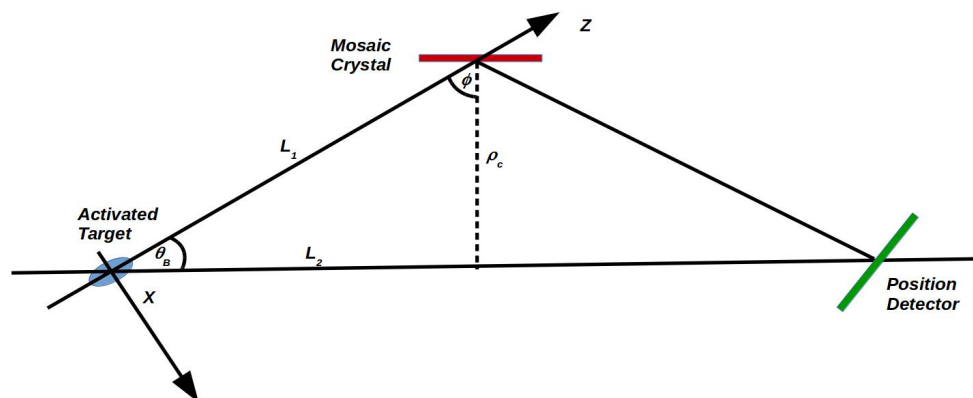


The shape of the “signal” beam can be optimized to increase the effective source size

On the other hand, this may lead to a worsening of the resolution

How big can a source be keeping FWHM < 10 eV?

VOXES: setup



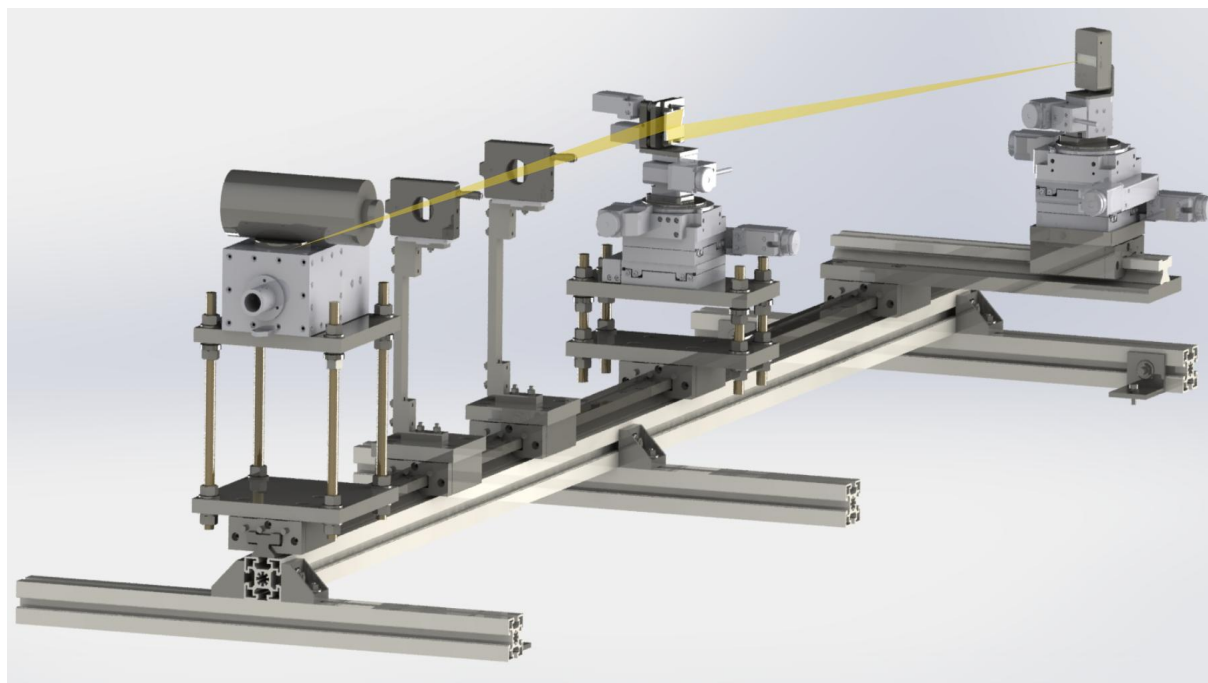
For a given X-ray energy the Bragg angle (θ_B) and the curvature radius of the crystal (ρ_c) completely determine the position of the source, the crystal and the position detector

$$L_1 = \frac{\rho_c}{\sin \theta_B}$$

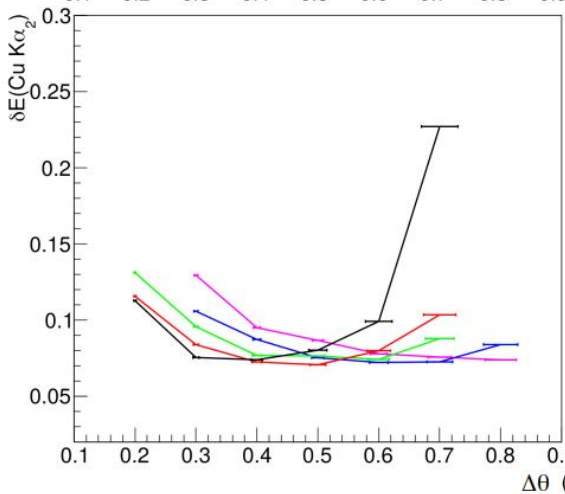
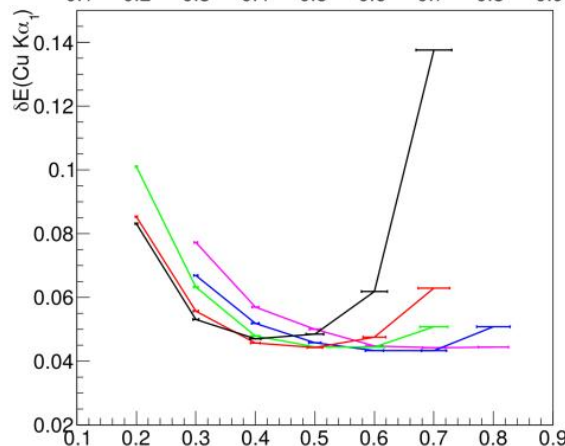
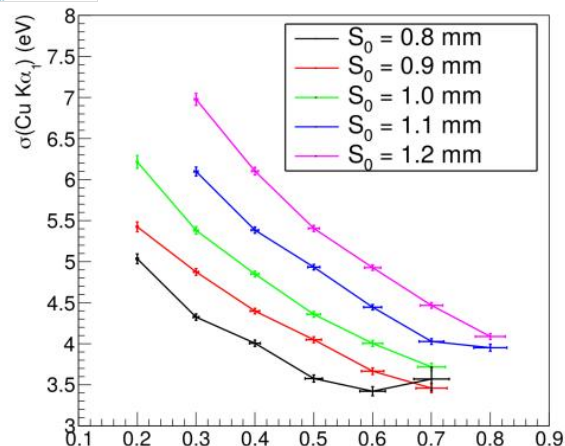
$$L_2 = L_1 \sin \phi$$

Table 1 List of the X-ray lines measured in this work and the corresponding Bragg angles θ_B

Line	E (eV)	θ_B (°)
Fe($K_{\alpha 1}$)	6403,84	16,77
Fe($K_{\alpha 2}$)	6390,84	16,81
Cu($K_{\alpha 1}$)	8047,78	13,28
Cu($K_{\alpha 2}$)	8027,83	13,31
Ni(K_{β})	8264,66	12,92
Zn($K_{\alpha 1}$)	8638,86	12,35
Zn($K_{\alpha 2}$)	8615,78	12,39
Mo($K_{\alpha 1}$)	17479,34	6,07
Mo($K_{\alpha 2}$)	17374,30	6,11
Nb(K_{β})	18622,50	5,70



VOXES: results



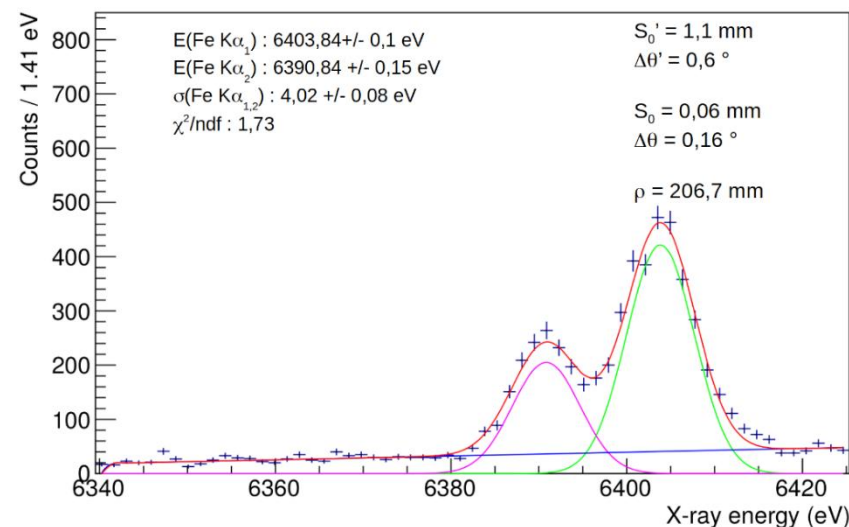
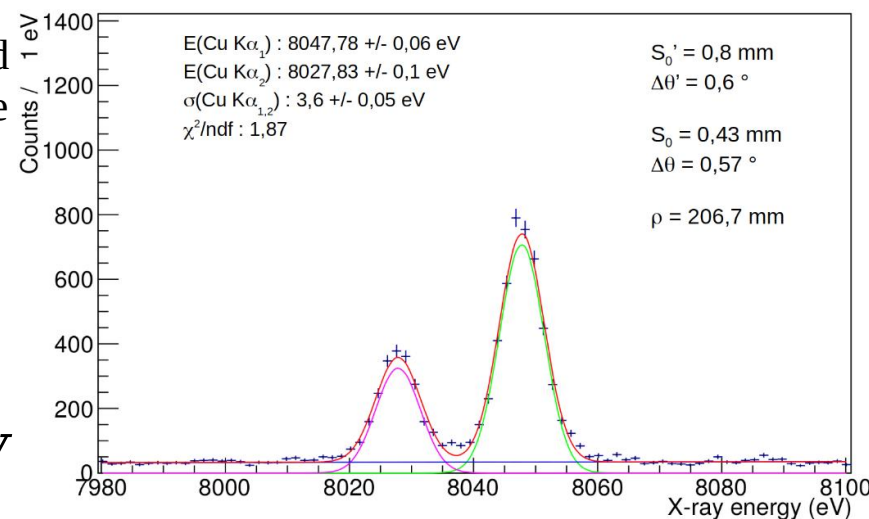
In the limit of a background free pure gaussian peak, the precision is related to the resolution via:

$$\delta E = \frac{\sigma E}{\sqrt{N}}$$

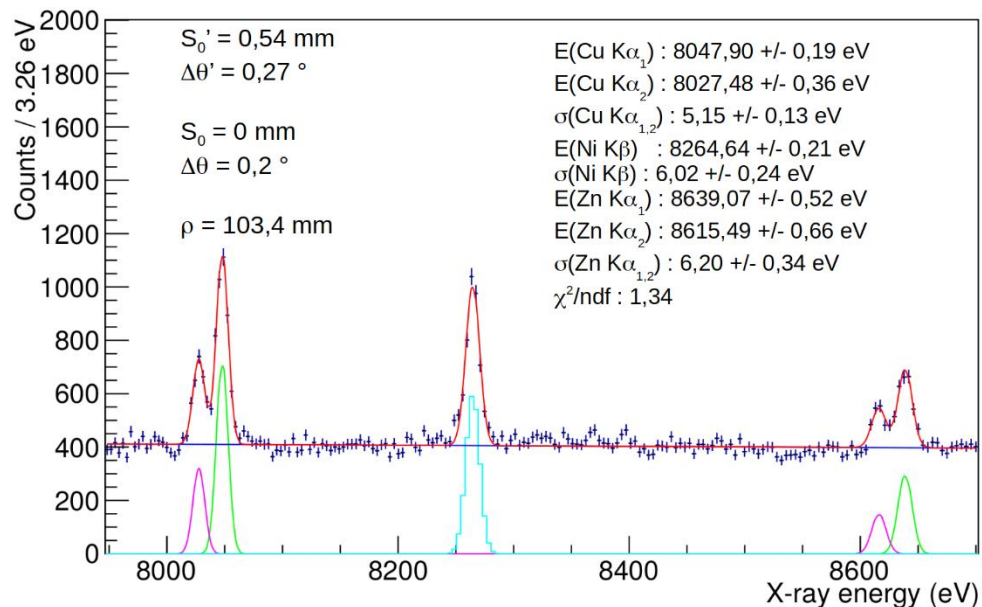
$$\frac{3,6}{\sqrt{4323}} = 0,0547 \text{ eV}$$

Given the energy and ρ_c it is always possible to find the optimal configuration to obtain the best peak position precision

Valid for all energies (tested for 6-20 keV)



VOXES: results



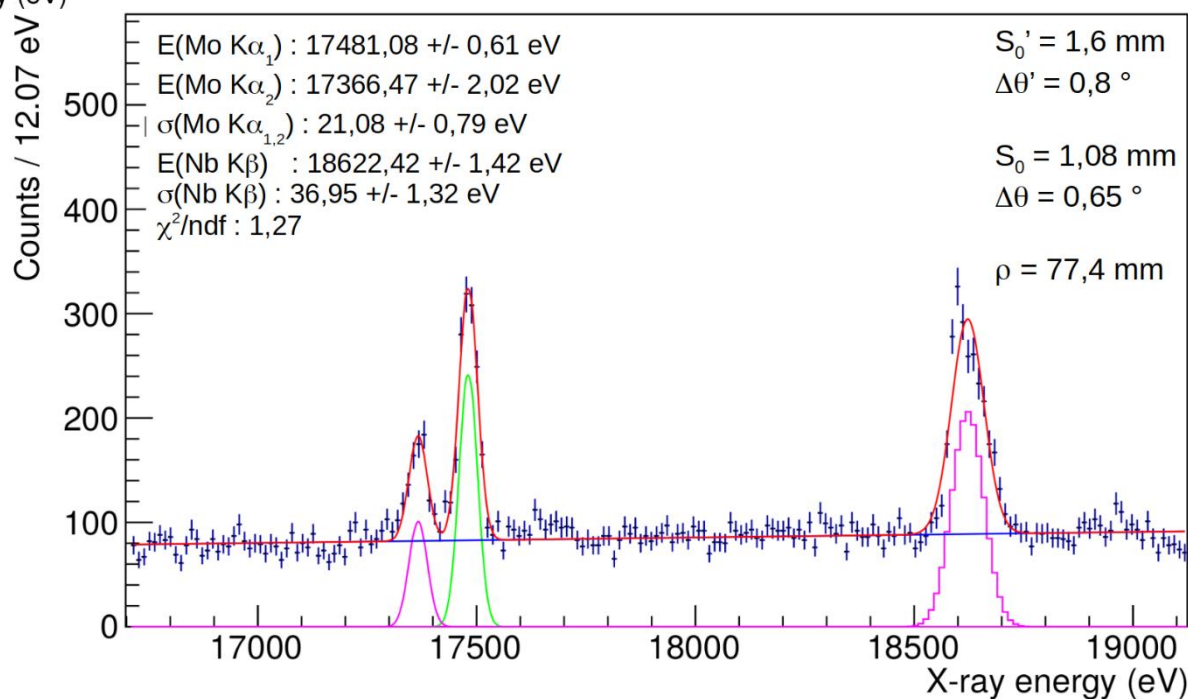
700 eV dynamic range (one shot)

Resolution still at 6 eV level (σ)

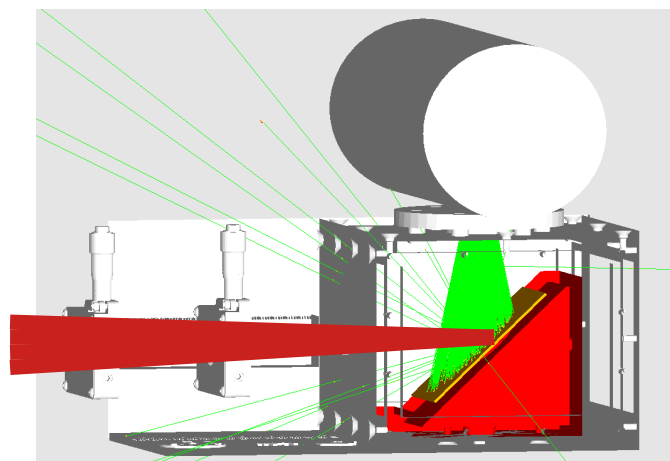
>1 keV dynamic range

Resolution still at 1,2 ‰

Almost 2 mm source size



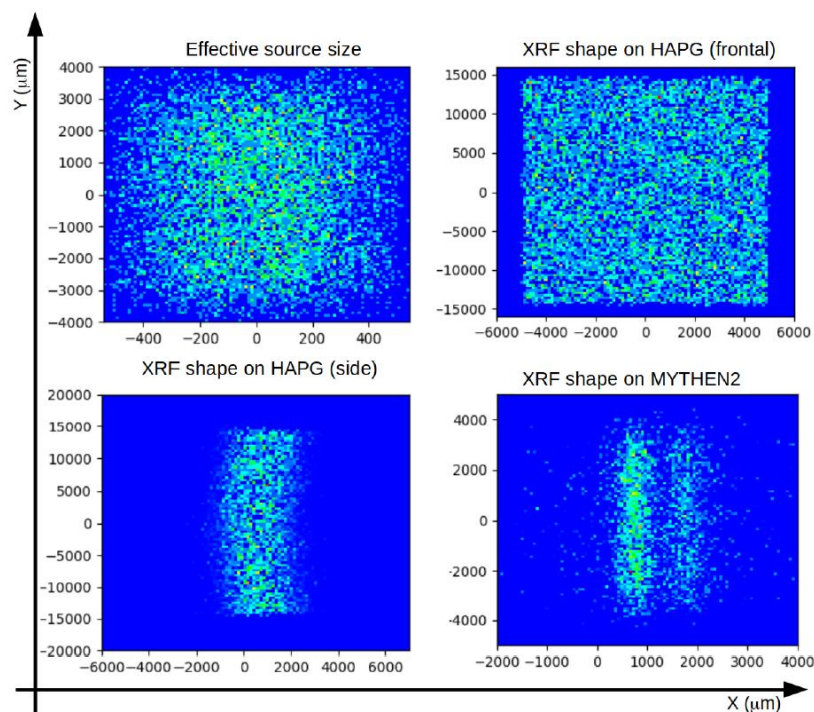
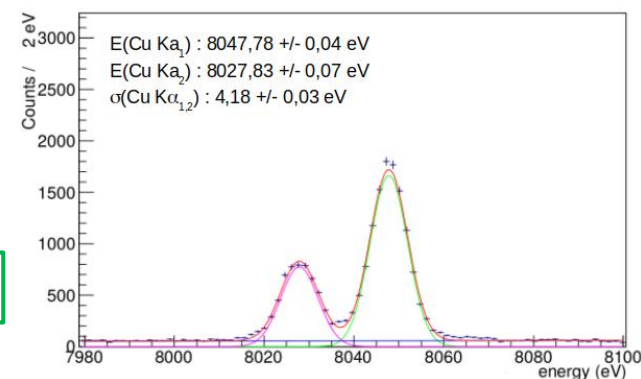
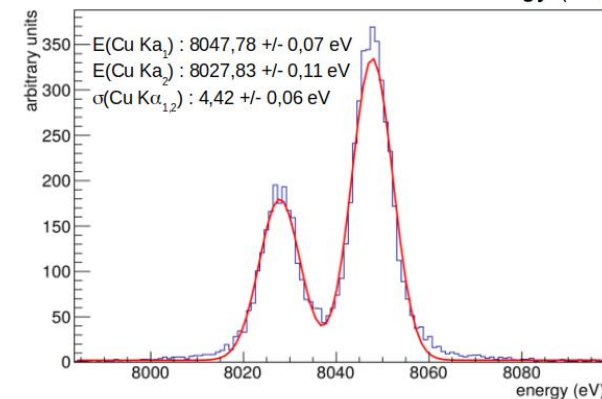
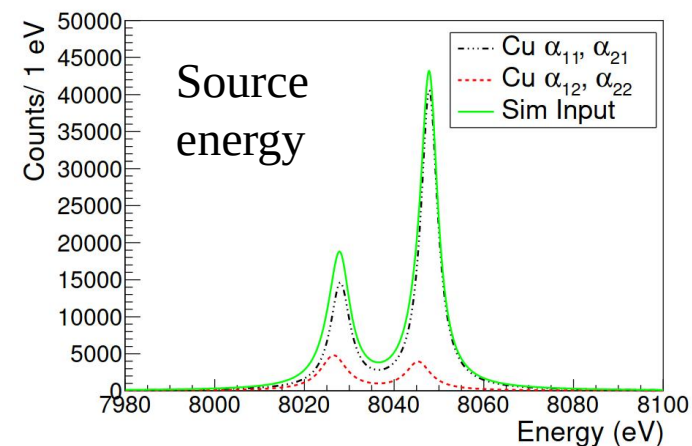
VOXES: ray tracing simulations



Software:
XOP → HAPG and ray tracing
SHADOW → Visualization

Integrated in OASYS tool

22 G. Hölzer et al., *Phys. Rev. A*, 1997, **56**, 4554–4568.

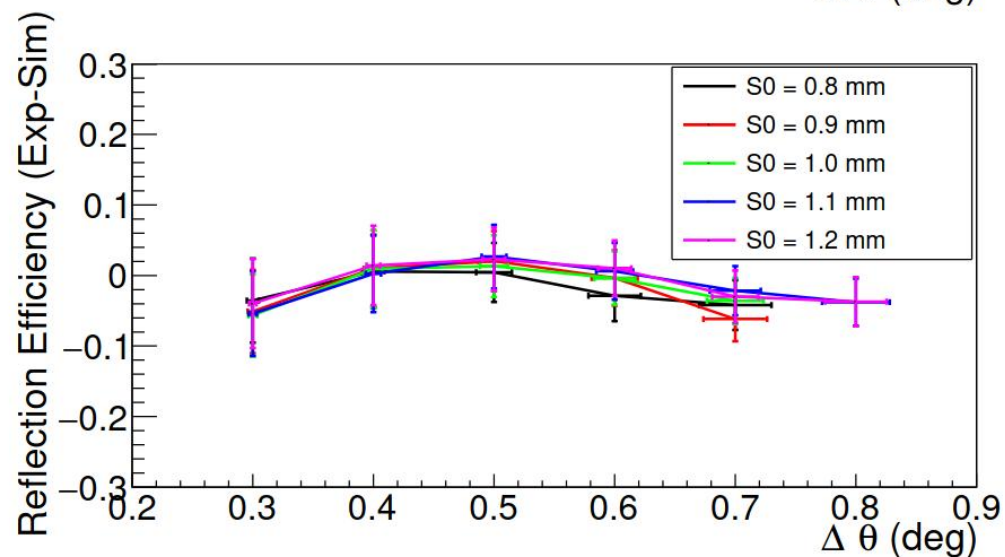
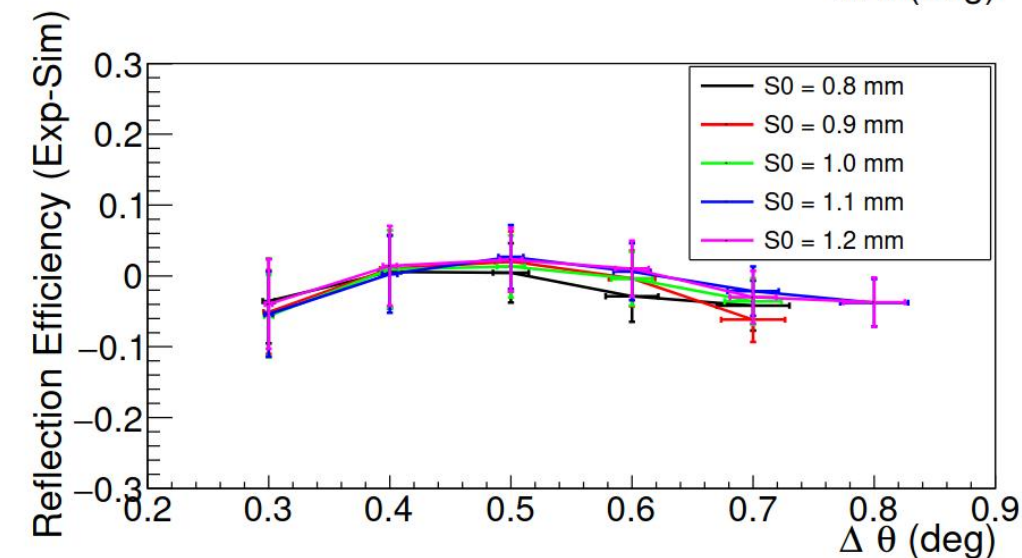
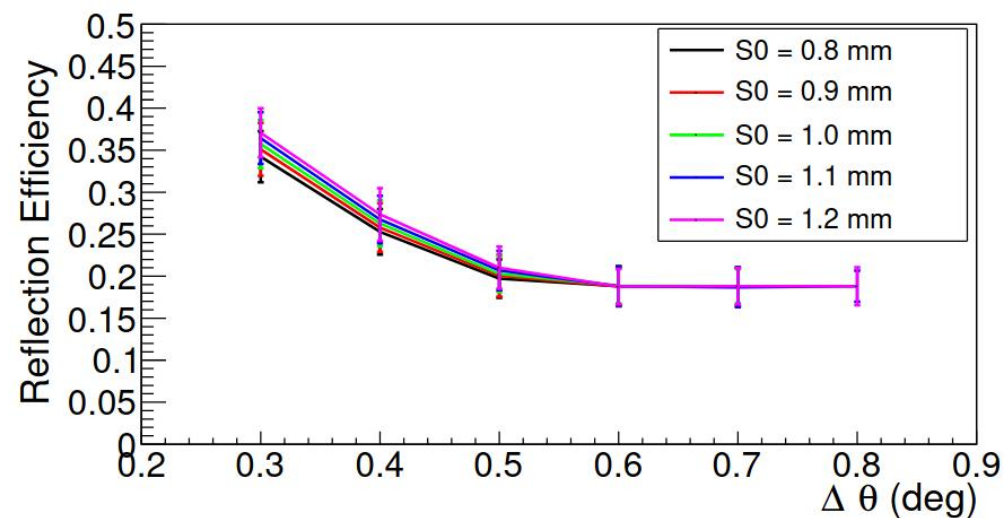
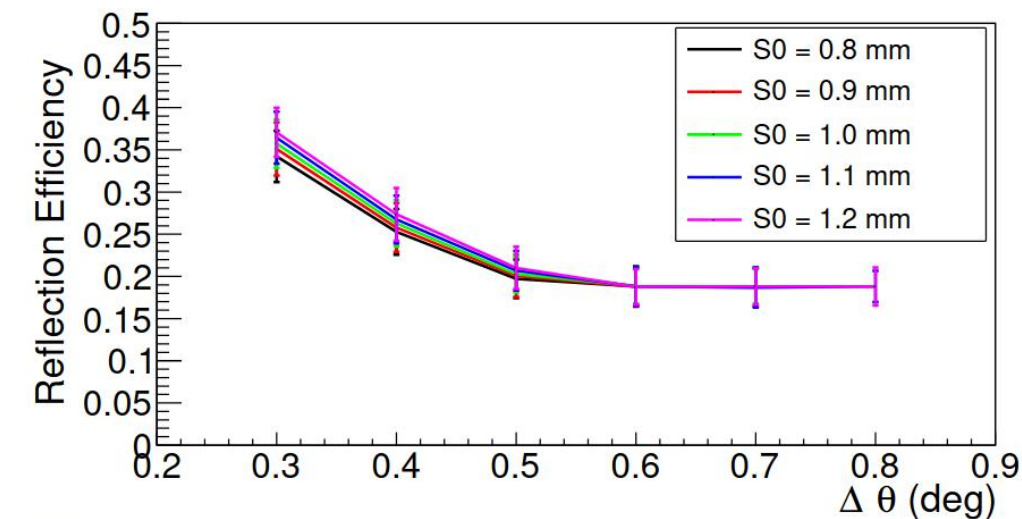


The physical source size is large (2,5x 2,5 cm)

We want to check the effective source size dimensions as coming from ray tracing simulations

Peak position and σ well reproduced

VOXES: ray tracing simulations



Reflection efficiencies are also well reproduced and under control

Possible kaonic transitions to be measured with HAPG crystal spectrometer:

$K^3He(3 \rightarrow 2) : 6.2 \text{ keV}$
 $K^3He(4 \rightarrow 2) : 8.4 \text{ keV}$
 $K^3He(5 \rightarrow 2) : 9.4 \text{ keV}$
 $K^3He(6 \rightarrow 2) : 9.9 \text{ keV}$
 $K^3He(7 \rightarrow 2) : 10.2 \text{ keV}$

$K^4He(3 \rightarrow 2) : 6.4 \text{ keV}$
 $K^4He(4 \rightarrow 2) : 8.7 \text{ keV}$
 $K^4He(5 \rightarrow 2) : 9.7 \text{ keV}$
 $K^4He(6 \rightarrow 2) : 10.3 \text{ keV}$
 $K^4He(7 \rightarrow 2) : 10.7 \text{ keV}$

$KN(6 \rightarrow 5) : 7.6 \text{ keV}$
 $KN(7 \rightarrow 5) : 12.1 \text{ keV}$
 $KN(8 \rightarrow 5) : 15.1 \text{ keV}$
 $KN(7 \rightarrow 6) : 4.6 \text{ keV}$
 $KN(8 \rightarrow 6) : 7.5 \text{ keV}$
 $KN(9 \rightarrow 6) : 9.6 \text{ keV}$
 $KN(10 \rightarrow 6) : 11 \text{ keV}$
 $KN(11 \rightarrow 6) : 12.1 \text{ keV}$
 $KN(10 \rightarrow 7) : 6.5 \text{ keV}$
 $KN(11 \rightarrow 7) : 7.5 \text{ keV}$
 $KN(12 \rightarrow 7) : 8.3 \text{ keV}$

Expected Impact:

- - Kaon mass measurements from different lines in parallel
- Cascade processes
- Impact on dark matter search driven experiments using exotic atoms in space (accurate cascade models calculations)
 - Upper level measurements with very small Γ
 - Proton radius puzzle (???)

Manifestation of interest from international institution and research centers (PSI, ...)

Detector Key Points:

- Tunable energy range from 2-20 keV
- Extremely high resolutions of few eV
- Very low background after shielding

Feasibility:

- Working principle tested in laboratory
- Dependence from HAPG parameters well investigated and published (thickness, mosaicity, ...)
 - Consistent Ray Tracing simulations available
- Few eV resolutions confirmed for solid sources with millimetric dimensions

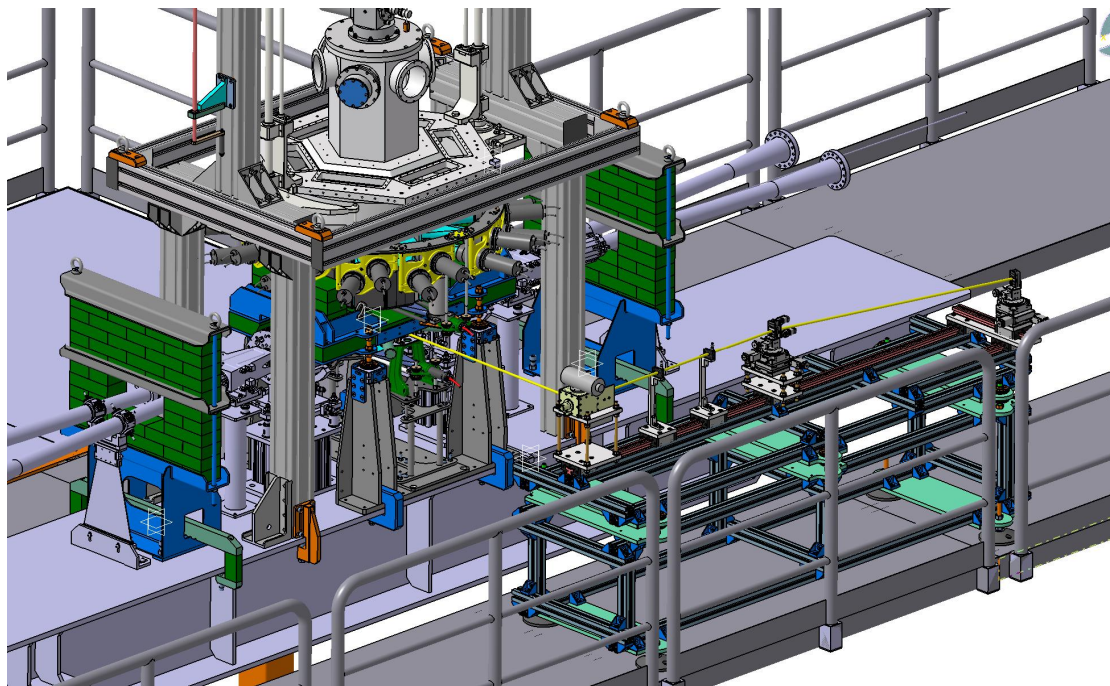
VOXES: a possible preliminary run



First run with KC for a feasibility test and background evaluation

Available:

- 1) Multi - Crystal support structure
- 2) Target (Solid or Liquid/Gas)
- 3) Optics
- 4) Alignment support
- 5) Target box
- 6) Detector
- 7) DAQ (integ. KM)

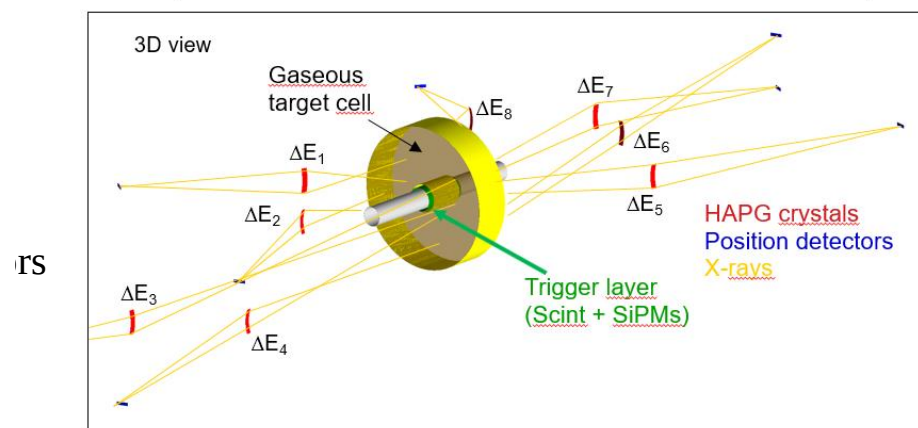
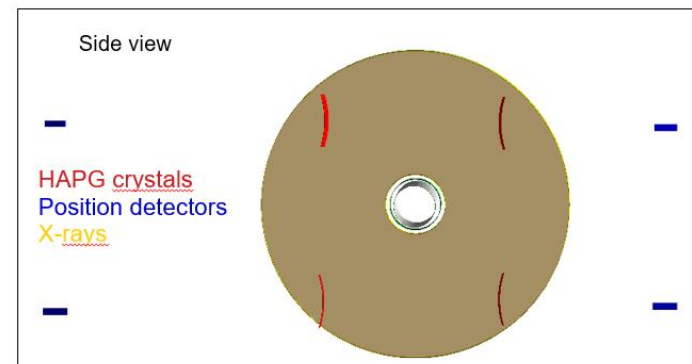
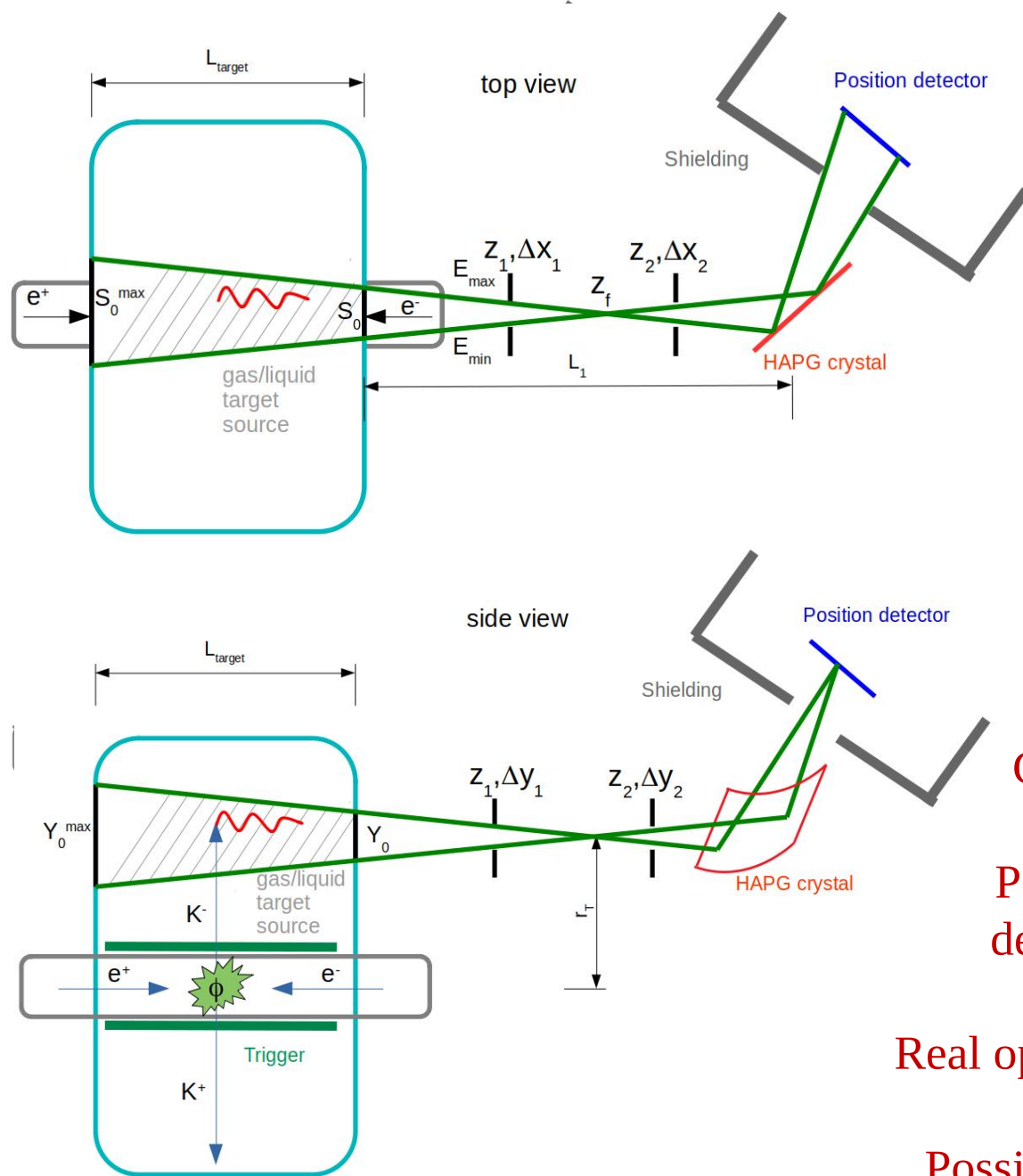


Future implementations:

- Shielding around Detector
- Solid support structure

Possible run in parallel with SIDDHARTA-2 @ LNF
in spring 2022

VOXES: future scenarios on DAΦNE



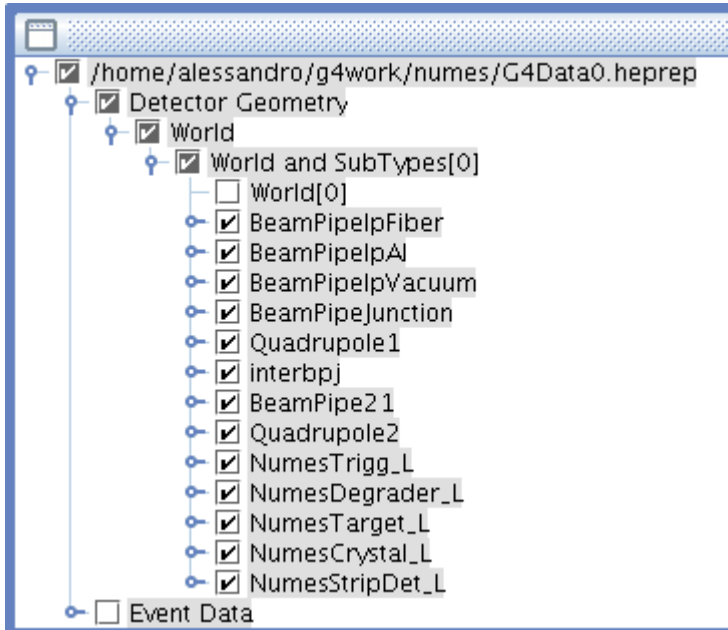
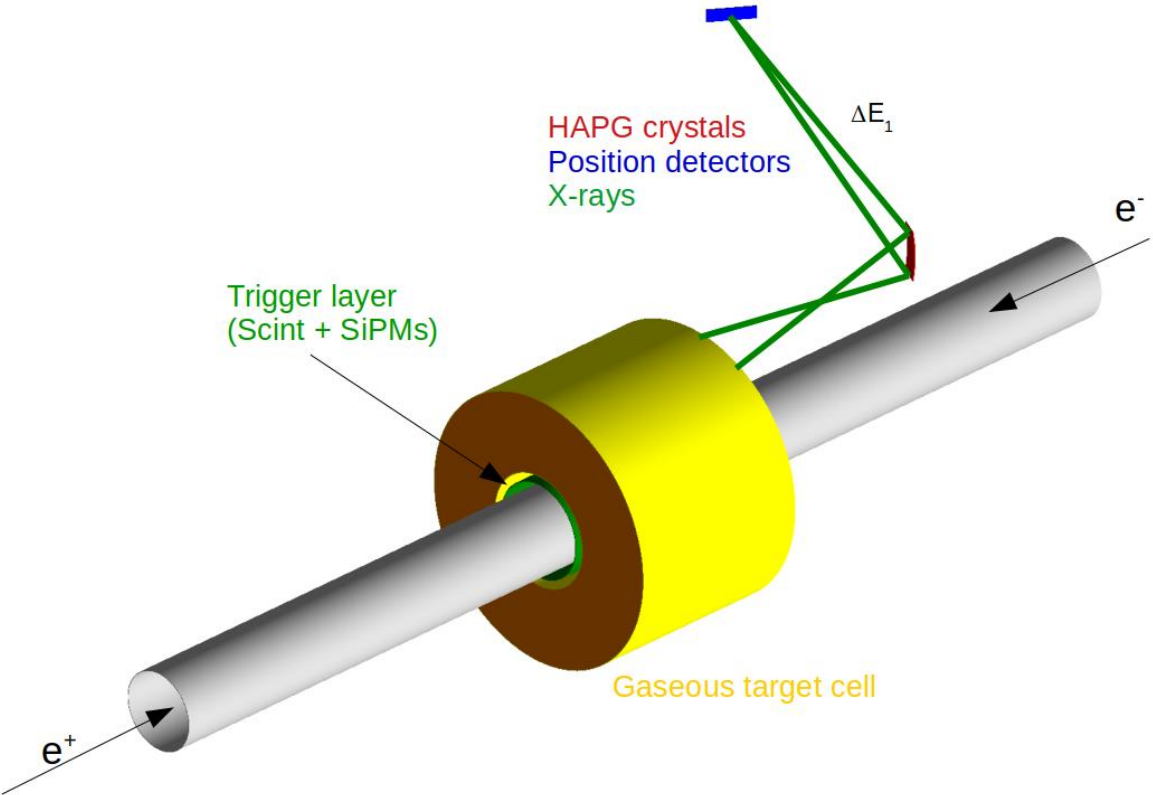
Completely new experiment / setup

Position Detector and HAPG crystals development with R&D opportunities

Real opportunity to apply for external fundings

Possibility to attract new interested institutes

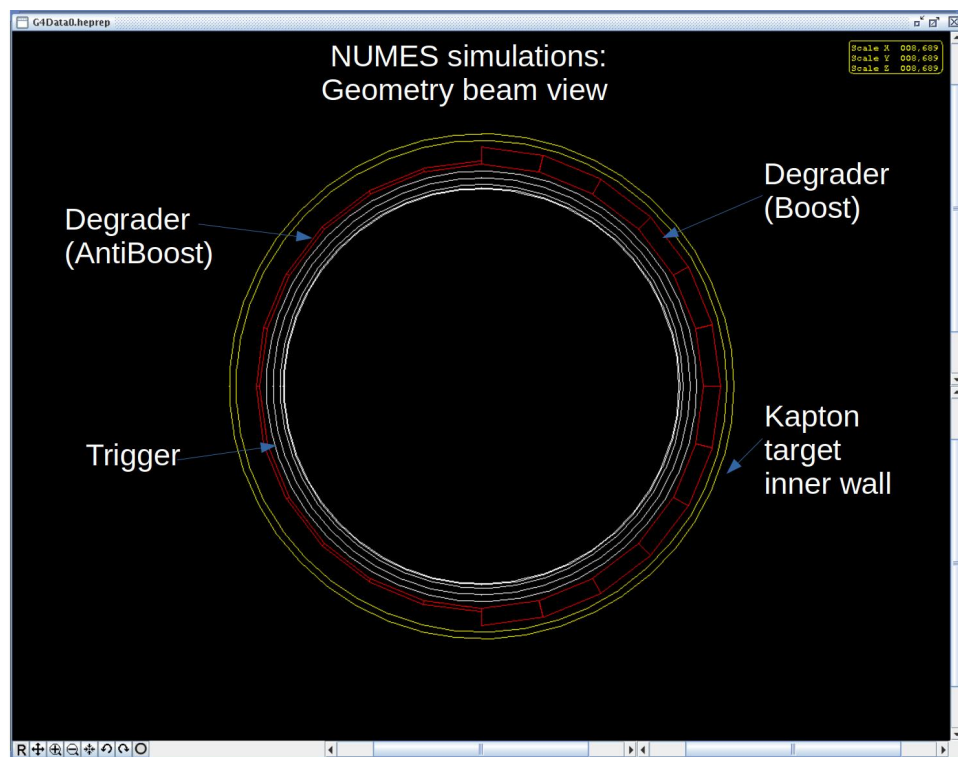
NUMES : **N**ovel X-ray detector system for **U**ltra-high precision **M**easurements of **E**xotic atoms from macroscopic **S**ources



Geant4:
Starting point is SIDDHARTA-2 MC

- Geometry:
- BeamPipe & Quadrupoles from SIDDHARTA-2
 - Other parts are removed
 - Trigger: 1mm BC420 cylinder
 - Degrader: variable thickness Mylar cylinder
 - Cylindrical Target: $\Delta R \geq 10$ cm, $L = 40$ cm

Ray tracing simulations:
Starting point are VOXES simulations

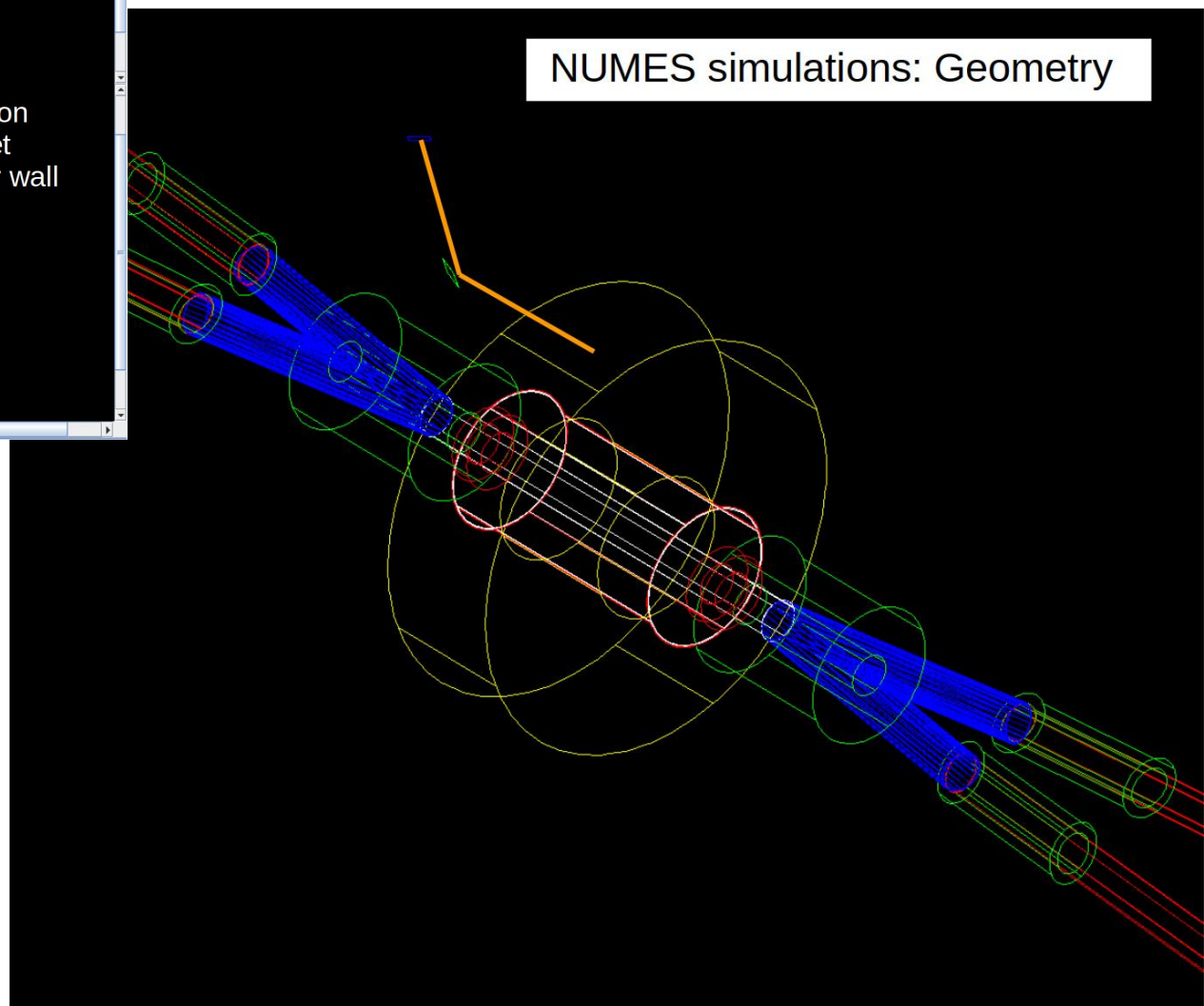


Kaons stopping number and position is strongly dependent on the geometry and on the degrader / trigger materials and thicknesses

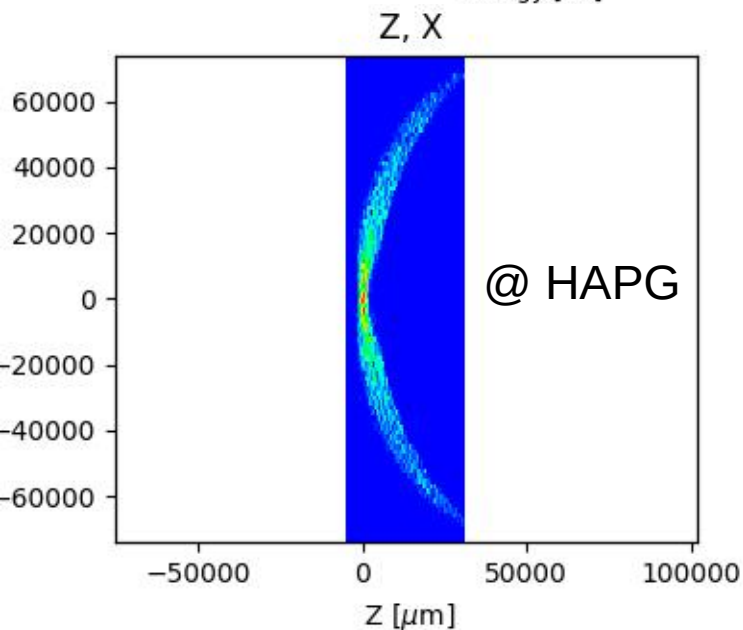
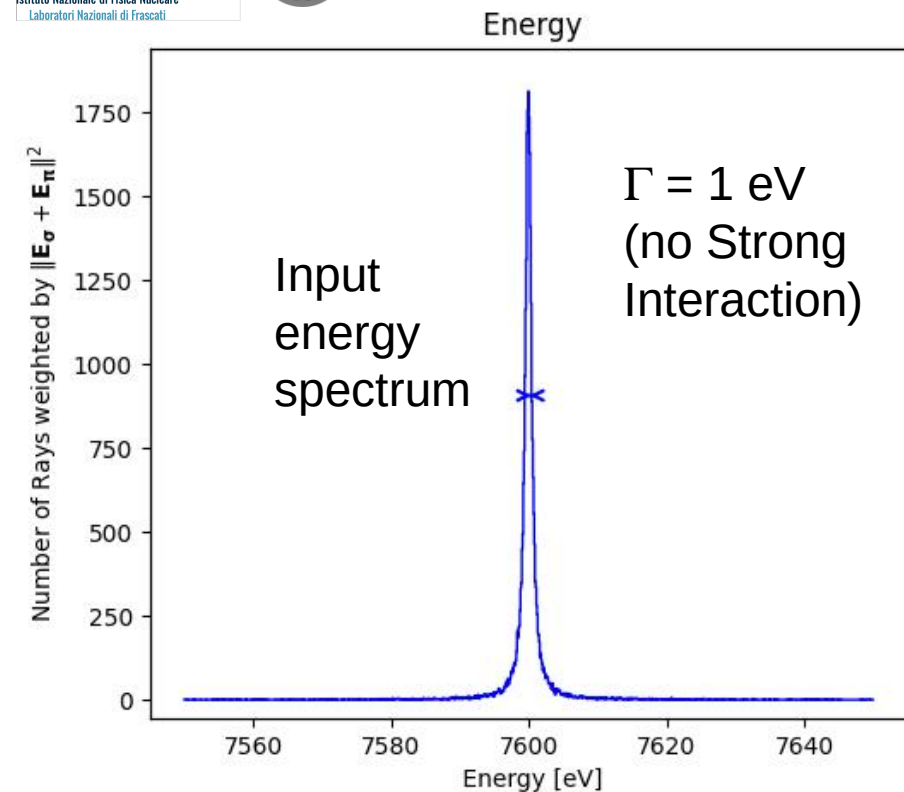
→ A proper optimization requires a big effort

2-steps simulation:

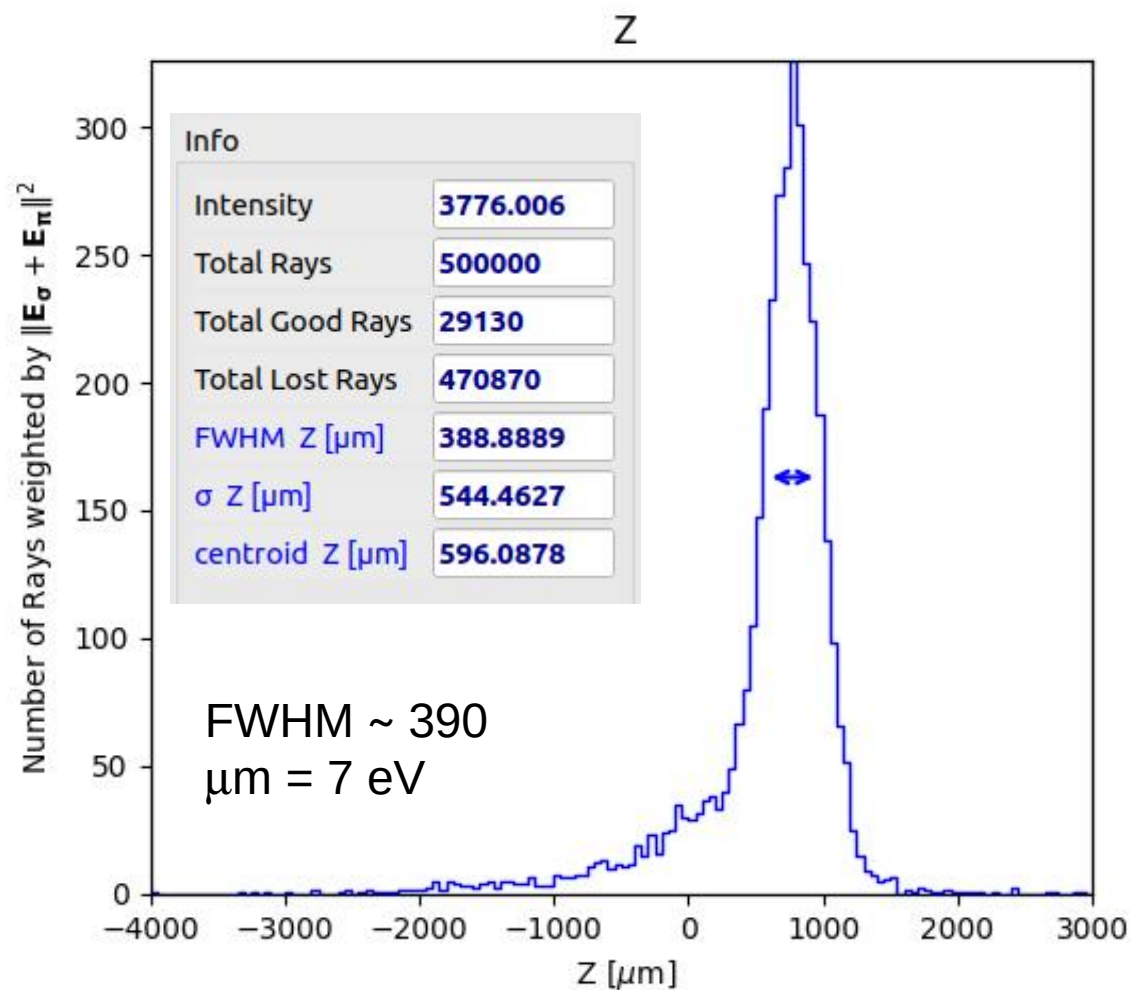
The geometrical distribution of the K-stopped is used as input source for ray tracing simulations



KN $6 \rightarrow 5$ @ 7.6 keV

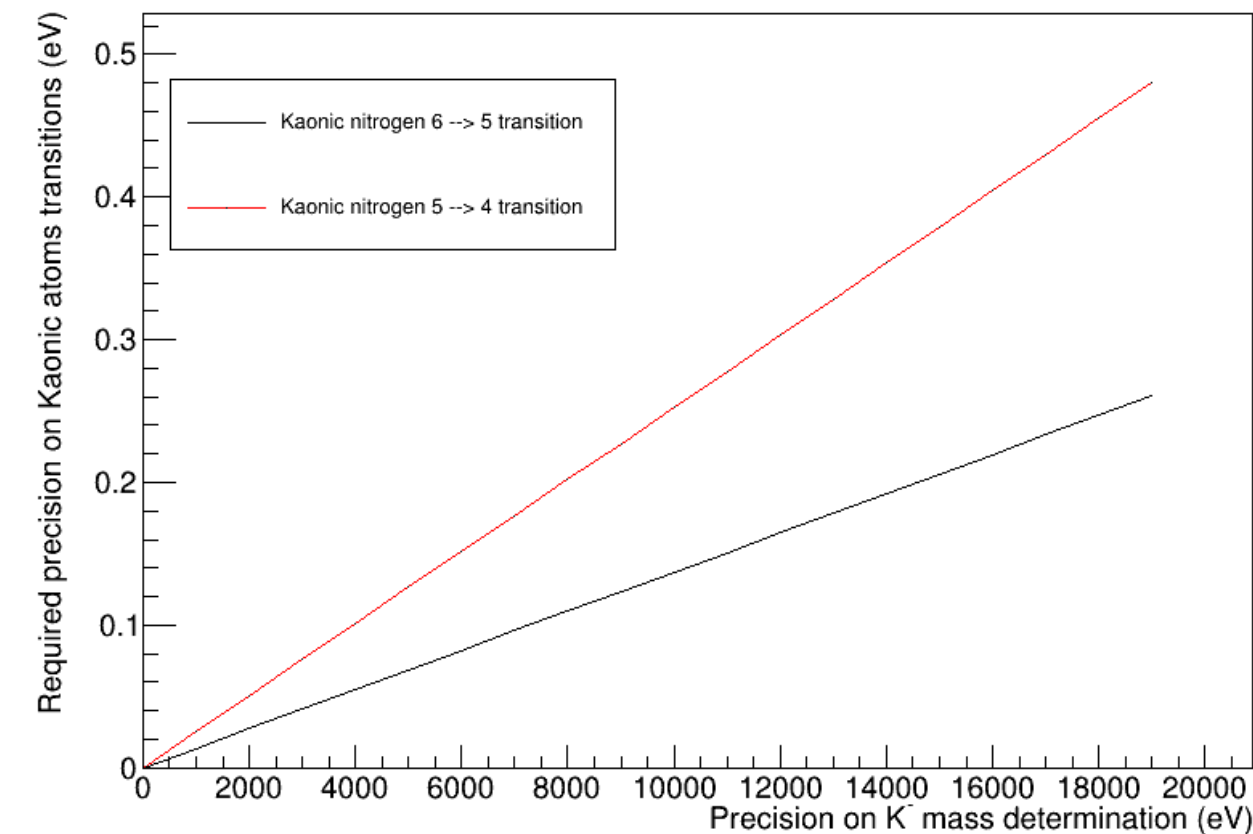


Simulated Bragg spectrum
(1 bin = 1 50 μm strip)



KN $6 \rightarrow 5$ @ 7.6 keV

N to have 0,15 eV precisions	Lumi (pb-1) / 8 arm needed	N to have 0,2 eV precisions	Lumi (pb-1) / 8 arm needed	N to have 0,25 eV precisions	Lumi (pb-1) / 8 arm needed
394,35	878,66	221,82	494,25	141,96	316,32

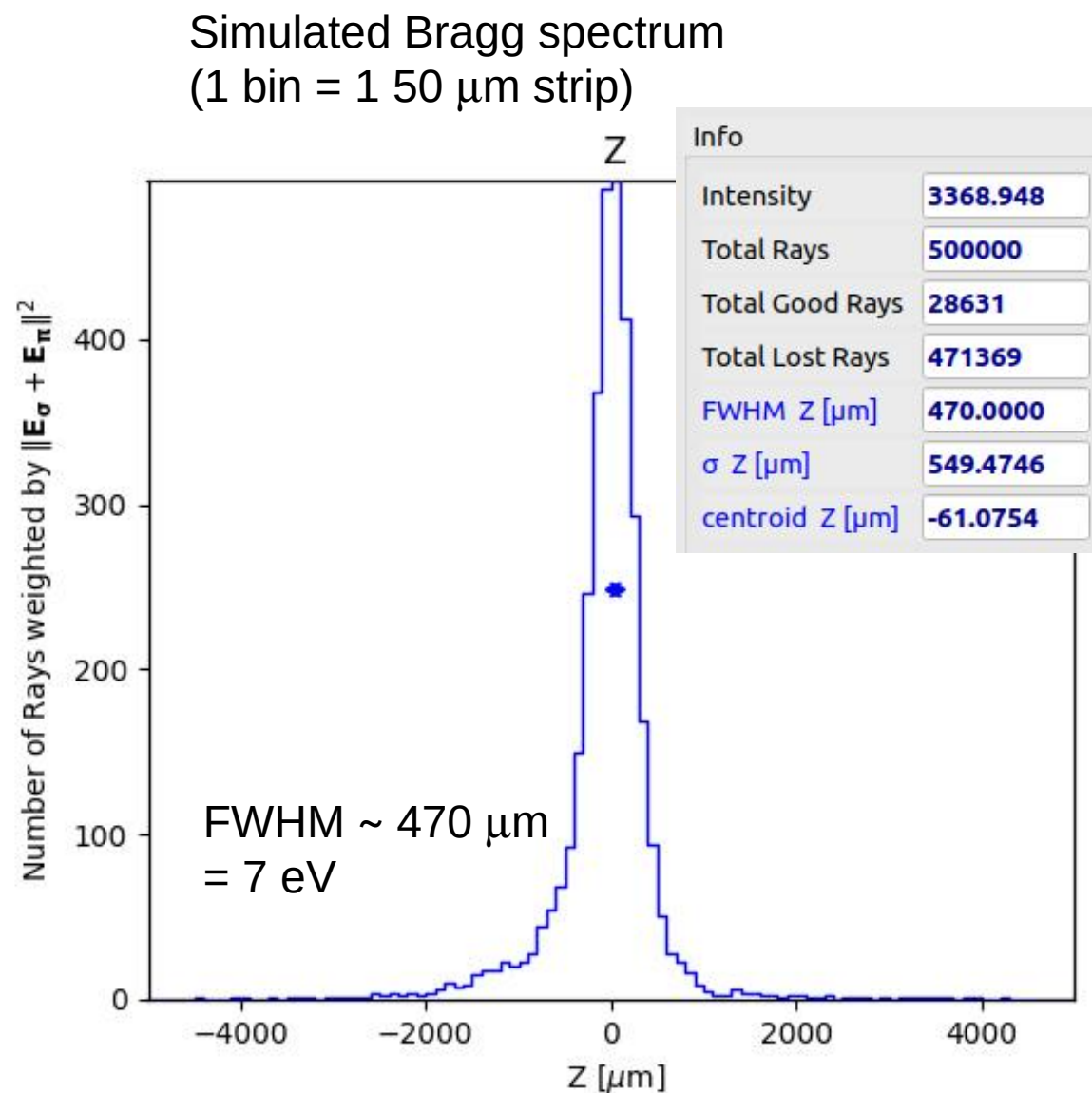
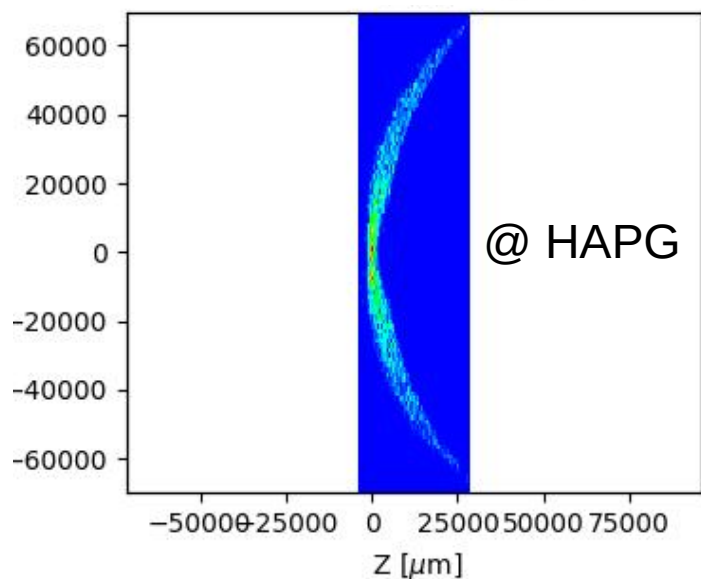
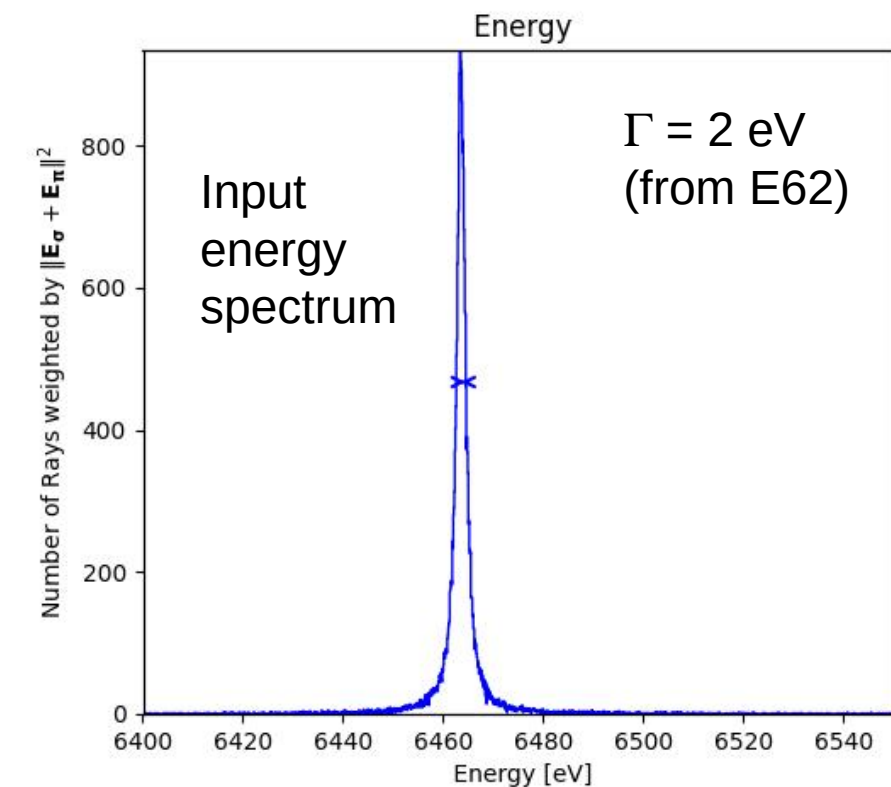


Depending on true DAFNE luminosity (in Siddharta it was in average $2-3 \times 10^{32} \sim 20 \text{ pb}^{-1} / \text{day}$) we have the possibility to perform a measurement of the K- mass well below 7 keV

Using 2 crystals for each arm, there is the possibility to focus also KN(5 \rightarrow 4) line @ 14 keV on the same position detector.

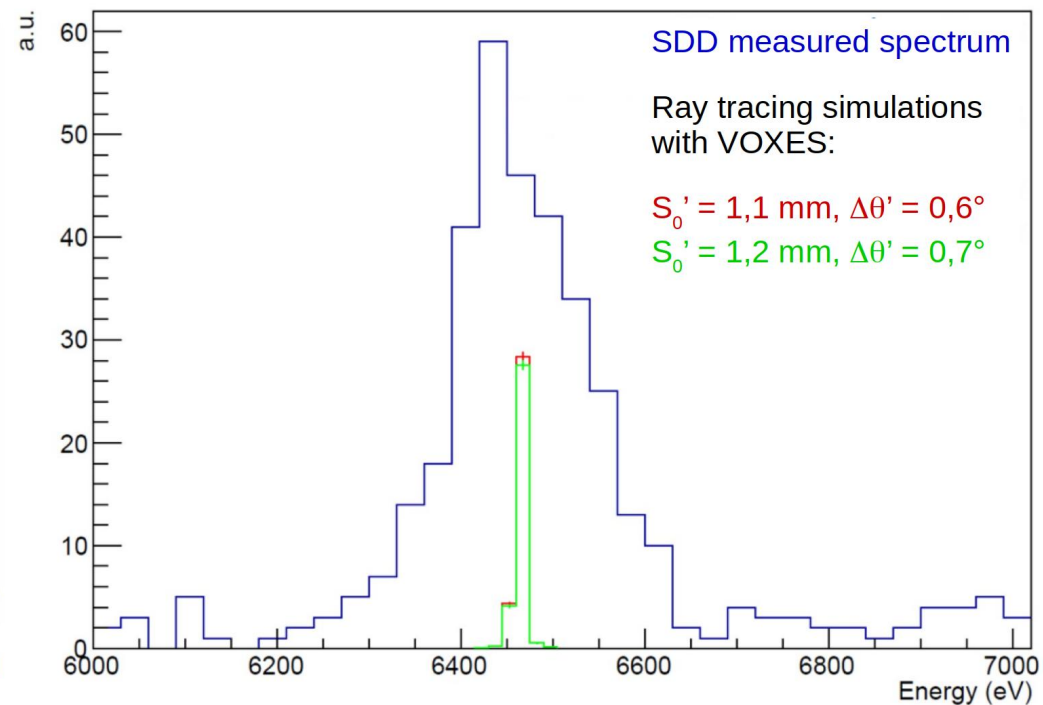
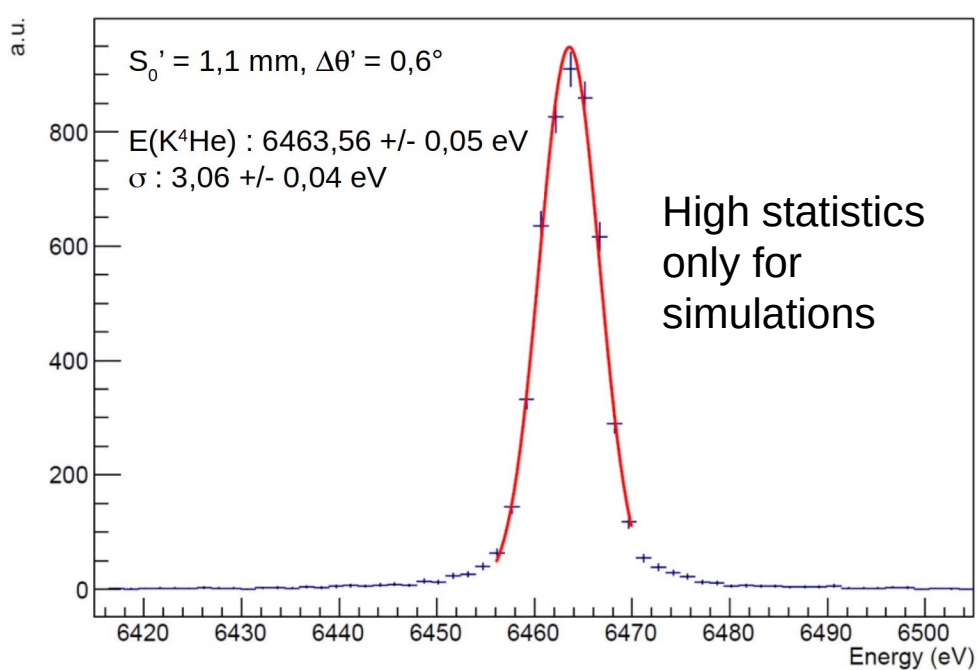
In this way, more statistics could be achieved with low costs

KHe @ 6.4 keV



KHe @ 6.4 keV

N to have 0,1 eV precisions	Lumi (pb-1) / 8 arm needed	N to have 0,15 eV precisions	Lumi (pb-1) / 8 arm needed	N to have 0,2 eV precisions	Lumi (pb-1) / 8 arm needed
862,11	1558,92	383,16	692,85	215,53	389,73



Conclusions



- HAPG based Bragg spectrometers represents a concrete possibility for future sub-eV precision kaonic atoms measurements
- VOXES collaboration developed in Frascati a version of such a spectrometer, to be used also with sources up to mm/cm dimensions
- Detailed investigation and optimization of crystal parameters, calibration procedure and peak shape description has been carried on
- The obtained results are very promising, showing precisions and resolution (well) below 1 eV and 10 eV, respectively
- MC ray tracing simulations have been also performed, which proved to be solid and to perfectly match the data. All these ingredients represent a fundamental starting point for future application
- With a first preliminary test and an expanded ad-hoc setup, VOXES spectrometer has been included in the proposal for future experiments to be carried on at DAΦNE after SIDDHARTA-2:
- *“Fundamental physics at the strangeness frontier at DAΦNE. Outline of a proposal for future measurements”*, arXiv:2104.06076v2 [nucl-ex]

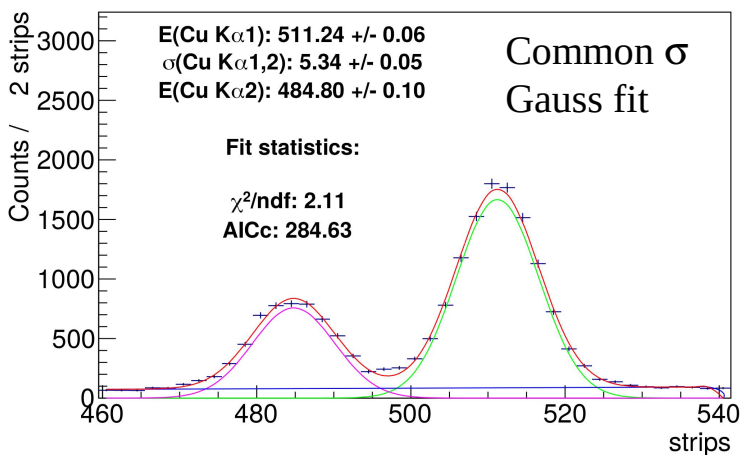
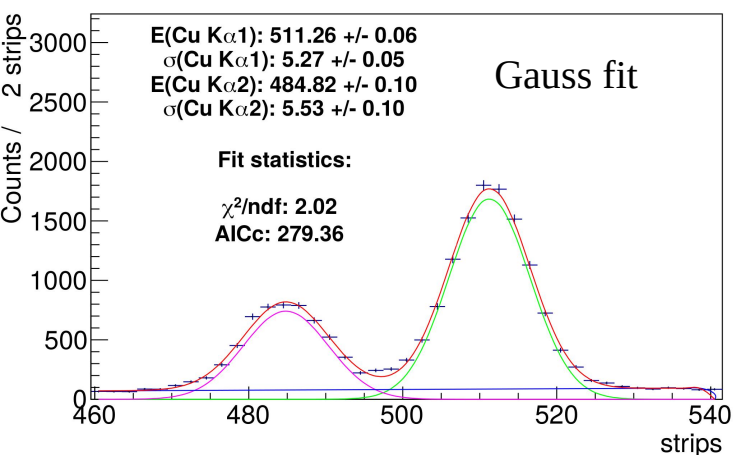
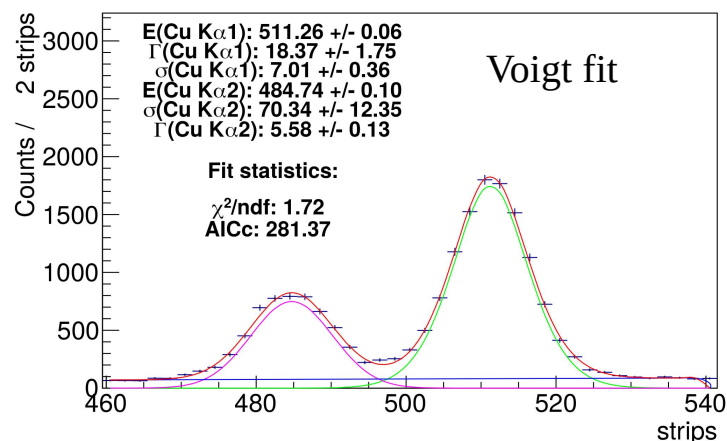
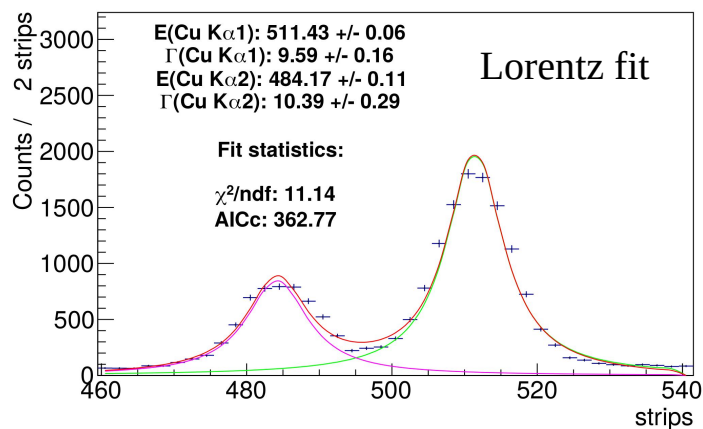
References:

- A. Scordo et al., J. Anal. At. Spectrom., 2020, 35, 155–168.
- A. Scordo et al., Condensed Matter, 2019, 4, 59.
- A. Scordo et al., JINST, 2018, 13, C04002.
- A. Scordo et al., Acta Phys. Polon., 2017, B48, 1715.

Spares



Which is the correct shape to be used for peak fitting?
(Natural linewidths are Lorentzian but....)



$$V(X) = \frac{A}{2\pi} \frac{\Gamma}{(x-x_0)^2 + \frac{\Gamma^2}{4}} \frac{e^{-\frac{(x-x_0)^2}{2\sigma^2}}}{\sigma\sqrt{2\pi}}$$

Is Voigt really better?

Akaike Information Criteria:

$$AIC = 2p + N \cdot \ln\left(\frac{R}{N}\right)$$

N = num of fitted points

p = num of fit parameters

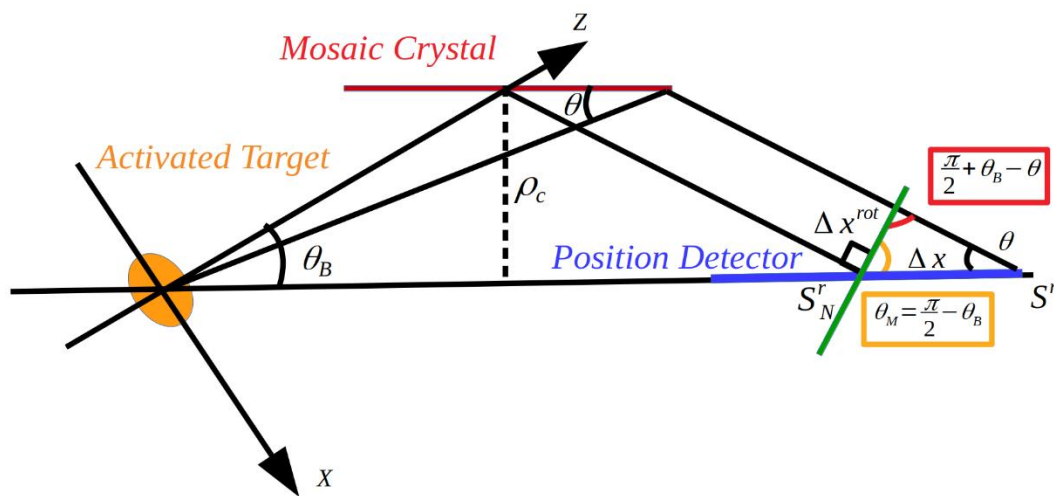
$$AICc = AIC + \frac{2 \cdot p \cdot (p+1)}{N - p - 1}$$

(for N/p < 40)

Not much information
loss using gaussian
shape

For each model i the quantity $e^{-0.5(AIC_{\text{cmin}} - AIC_i)}$ is proportional to the probability of the i-th model to minimize the (estimated) information loss as good as the minimum AICc one.

VOXES: energy calibration



Standard VH calibration: $\Delta x = 2\rho_c \cdot (\cot \theta_B - \cot \theta)$

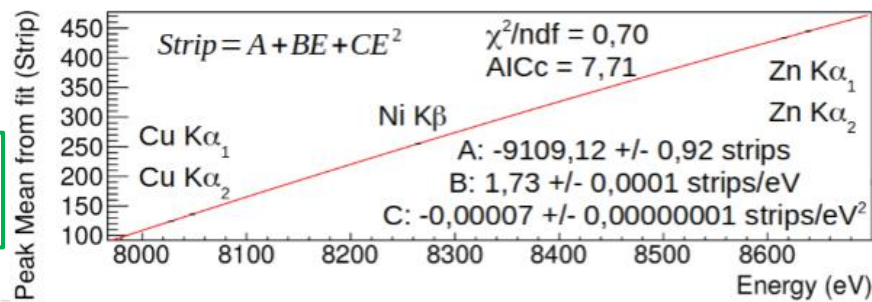
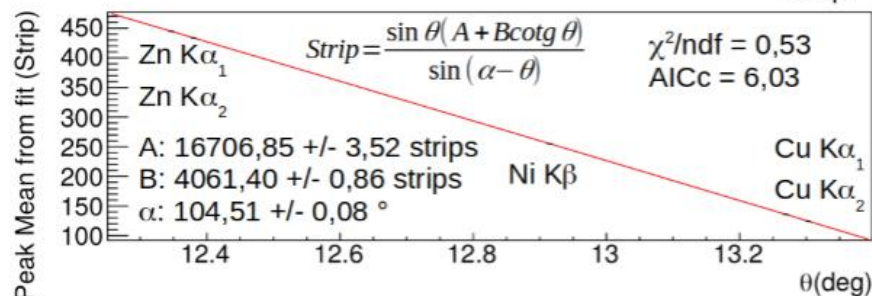
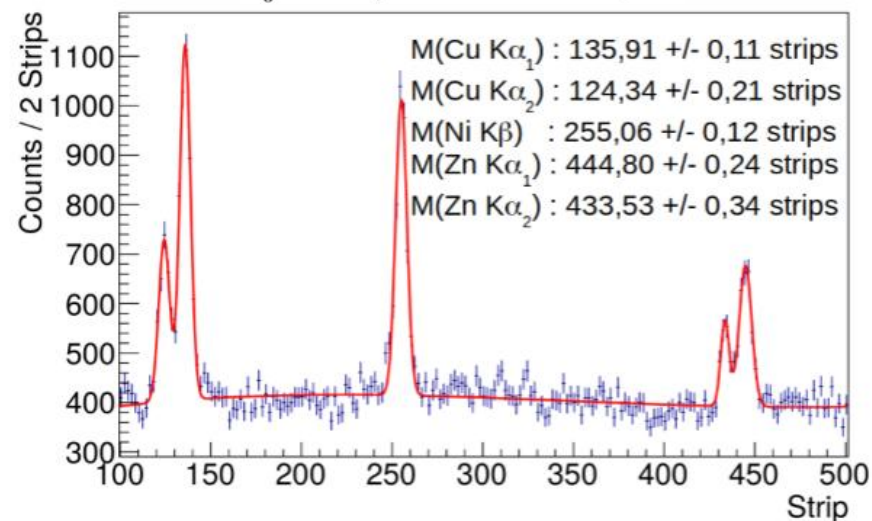
Semi VH calibration: $\Delta x^{rot} = \frac{2\rho_c \cdot (\cot \theta_B - \cot \theta) \cdot \sin \theta}{\sin(\frac{\pi}{2} - \theta_M - \theta)}$

Parametric form: $\Delta x^{rot} = \frac{(A + B \cdot \cot \theta) \cdot \sin \theta}{\sin(\alpha - \theta)}$

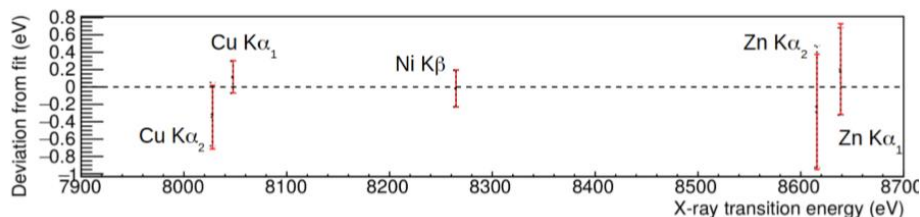
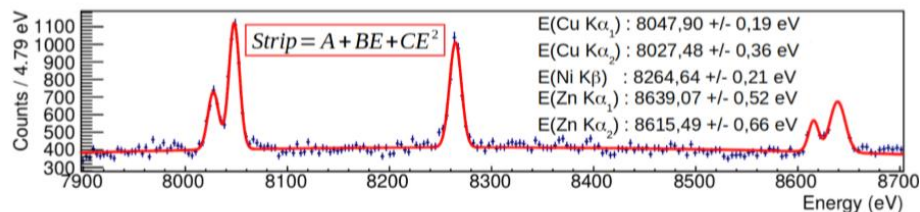
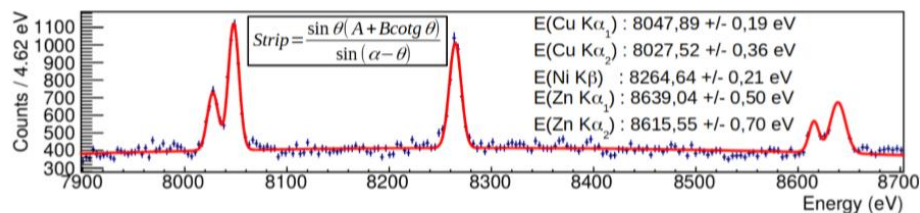
Given the small θ values ($\sin \theta \approx \theta$),
is it also possible to calibrate with a polynomial?

Information loss???

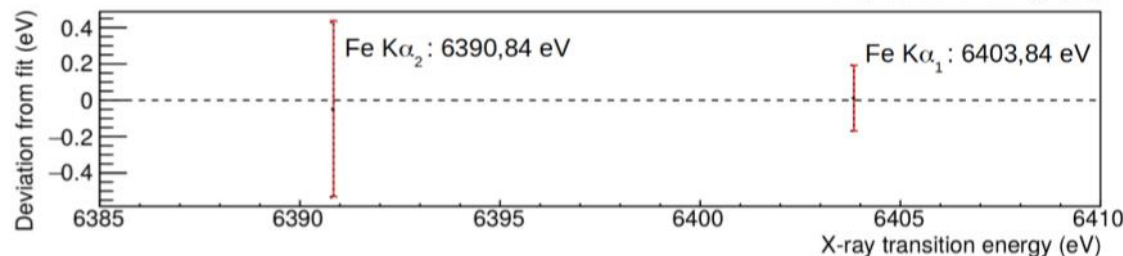
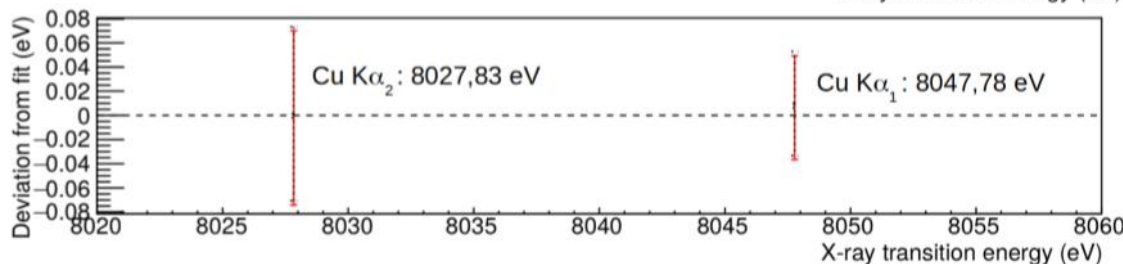
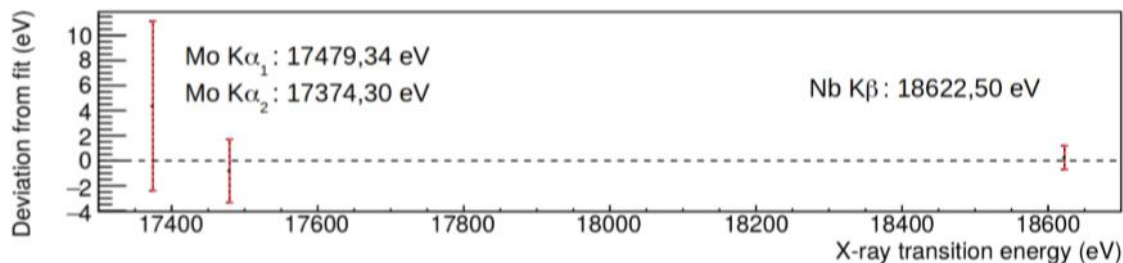
$$S'_0 = 540 \mu m \text{ and } \Delta \theta' = 0,27^\circ$$



VOXES: energy calibration



The (small) information loss is not influencing the peak positions



Also valid for higher and wider energy ranges (and higher θ , $\Delta\theta$ values)