

# Topology in Lattice QCD

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INFN Firenze



**EuroHPC**  
Joint Undertaking

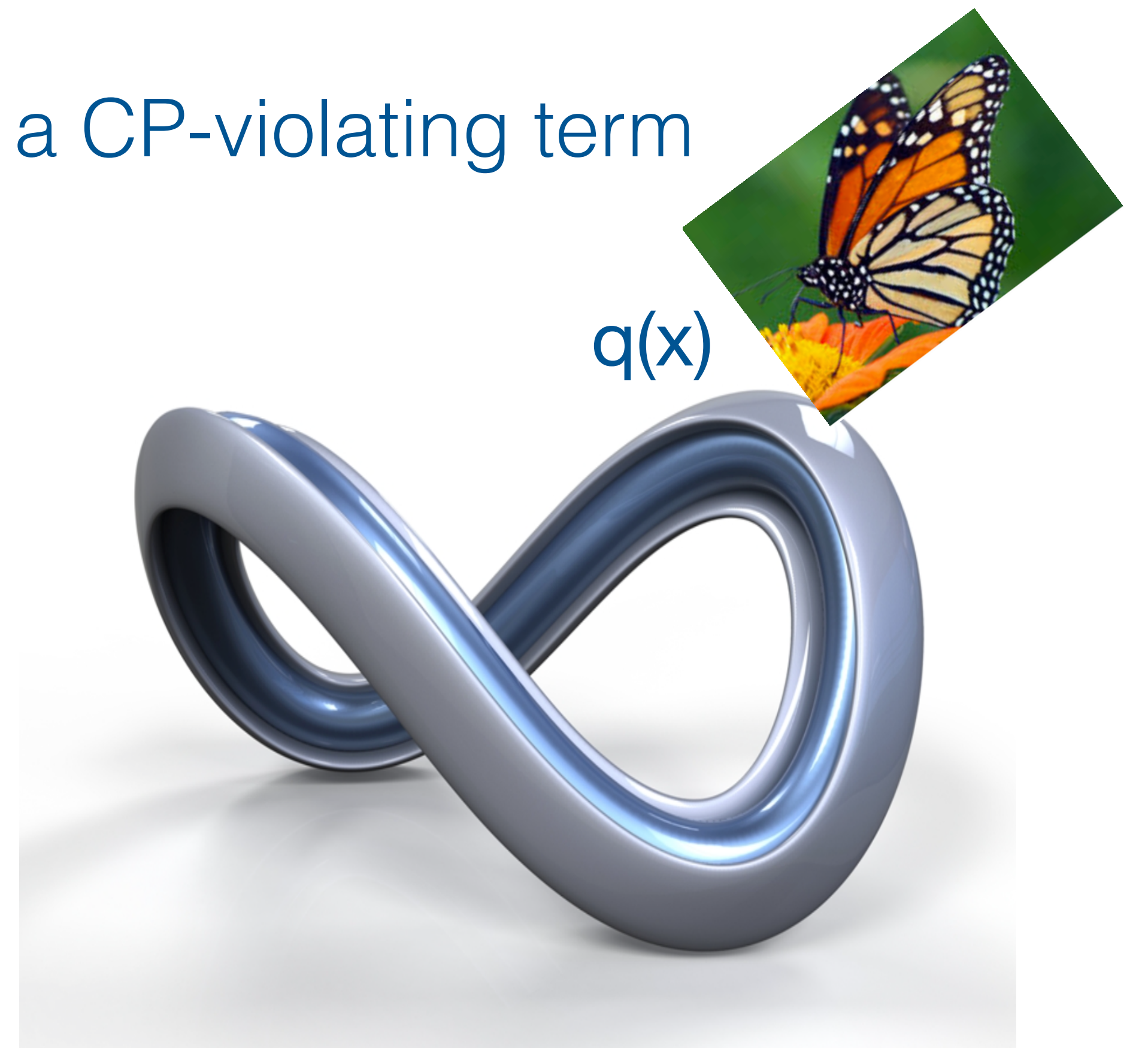
It is possible to add to the QCD Lagrangian a CP-violating term

$$\mathcal{L} = \mathcal{L}_{QCD} + \theta \frac{g^2}{32\pi^2} F_{\mu\nu}^a \tilde{F}_a^{\mu\nu}$$

$$\frac{g^2}{32\pi^2} F_{\mu\nu}^a \tilde{F}_a^{\mu\nu} = q(x)$$

$$Q = \sum q(x)$$

Defines a topological sector



Q — topological charge

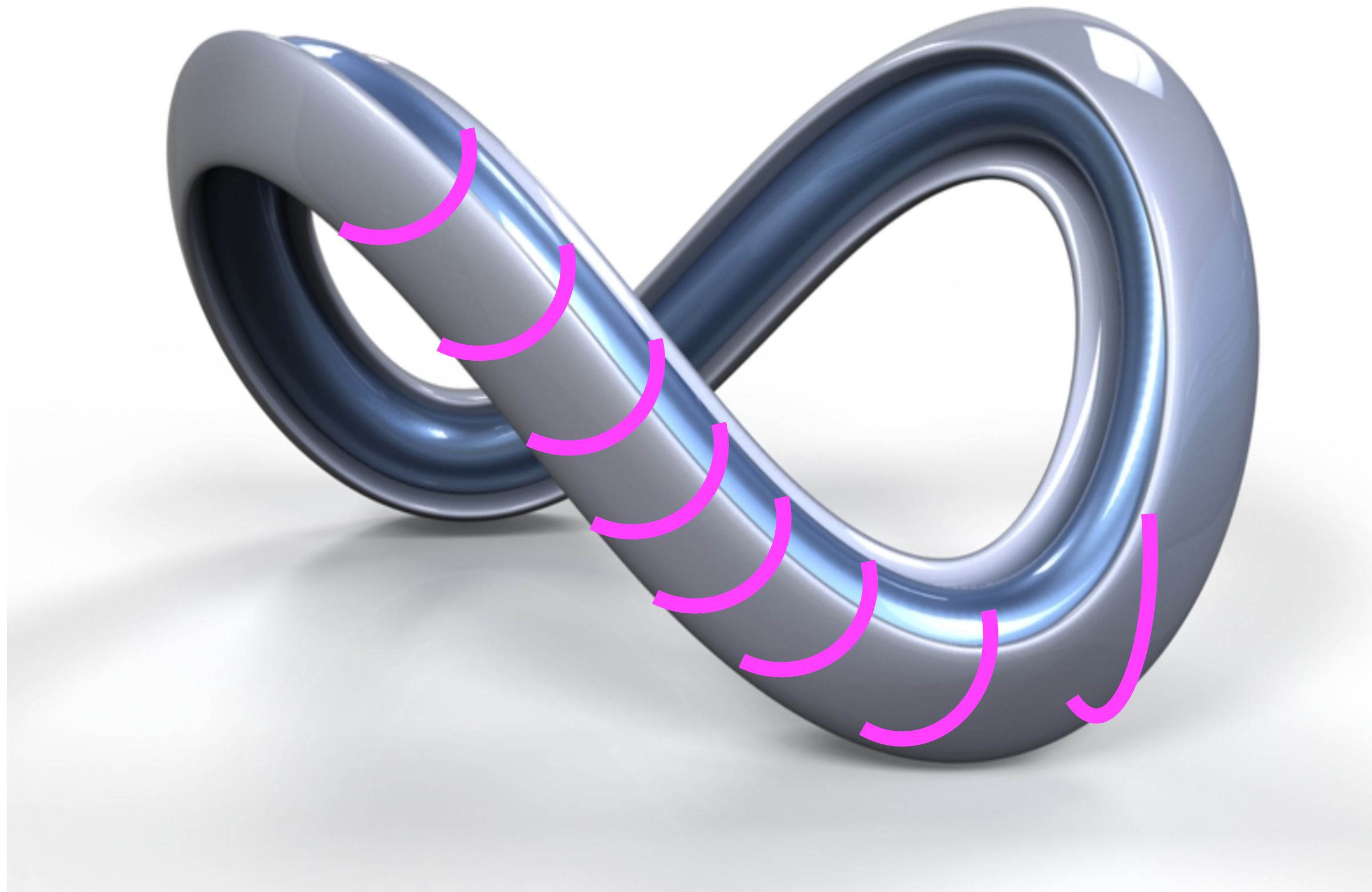
CP-violating term



Lattice discretization — see talk by Jana Guenther —

is particularly delicate..

..the very concept of ‘topology’ seems lost



The talks by:

Claudio Bonanno  
Francesco D'Angelo

discuss technical  
aspects and strategies

# Outline

## Topology in QCD

1. From low to high temperatures — a threshold in the plasma?
2. From low to high densities - experimental results and first theoretical analysis
3. Topology and the QCD axion

# TOPOLOGY, SYMMETRIES AND SPECTRUM

$$\mathcal{L} = \sum_{a=1}^n \bar{q}_{La} \not{\partial} q_{La} + \bar{q}_{Ra} \not{\partial} q_{Ra} - m(\bar{q}_{La} q_{La} + \bar{q}_{Ra} q_{Ra}) + \theta \frac{g^2}{32\pi^2} F_{\mu\nu}^a \tilde{F}_a^{\mu\nu} + \mathcal{L}_{gauge}$$

With  $m = 0$ , invariant under

$$q_L \rightarrow V_L q_L, q_R \rightarrow V_R q_R, \text{ with } V \in U(n)$$

Global symmetry:

$$U(n)_L \times U(n)_R \cong SU(n) \times SU(n) \times U(1)_V \times U(1)_A$$

.

Spontaneously Broken

baryon  
number

Explicitly broken

# TOPOLOGY, SYMMETRIES AND SPECTRUM

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baryon  
number

Spontaneously Broken

Explicitly broken

Experimental evidence

$(n^2 - 1)$  pseudoGB

Heavy  $\eta'$

Topology,  $\eta'$  and the  $U_A(1)$  problem:

It can be proven that

$$\frac{1}{32\pi^2} \int d^4x F \tilde{F} = Q \quad \text{Gluonic definition}$$

and

$$Q = n_+ - n_- \quad \text{Fermionic definition}$$

The  $\eta'$  mass may now be computed from the decay of the correlation

$$\langle \partial_\mu j_5^\mu(x) \partial_\mu j_5^\mu(y) \rangle \propto \frac{1}{N^2} \langle F(x) \tilde{F}(x) F(y) \tilde{F}(y) \rangle$$

which at leading order gives the Witten-Veneziano formula

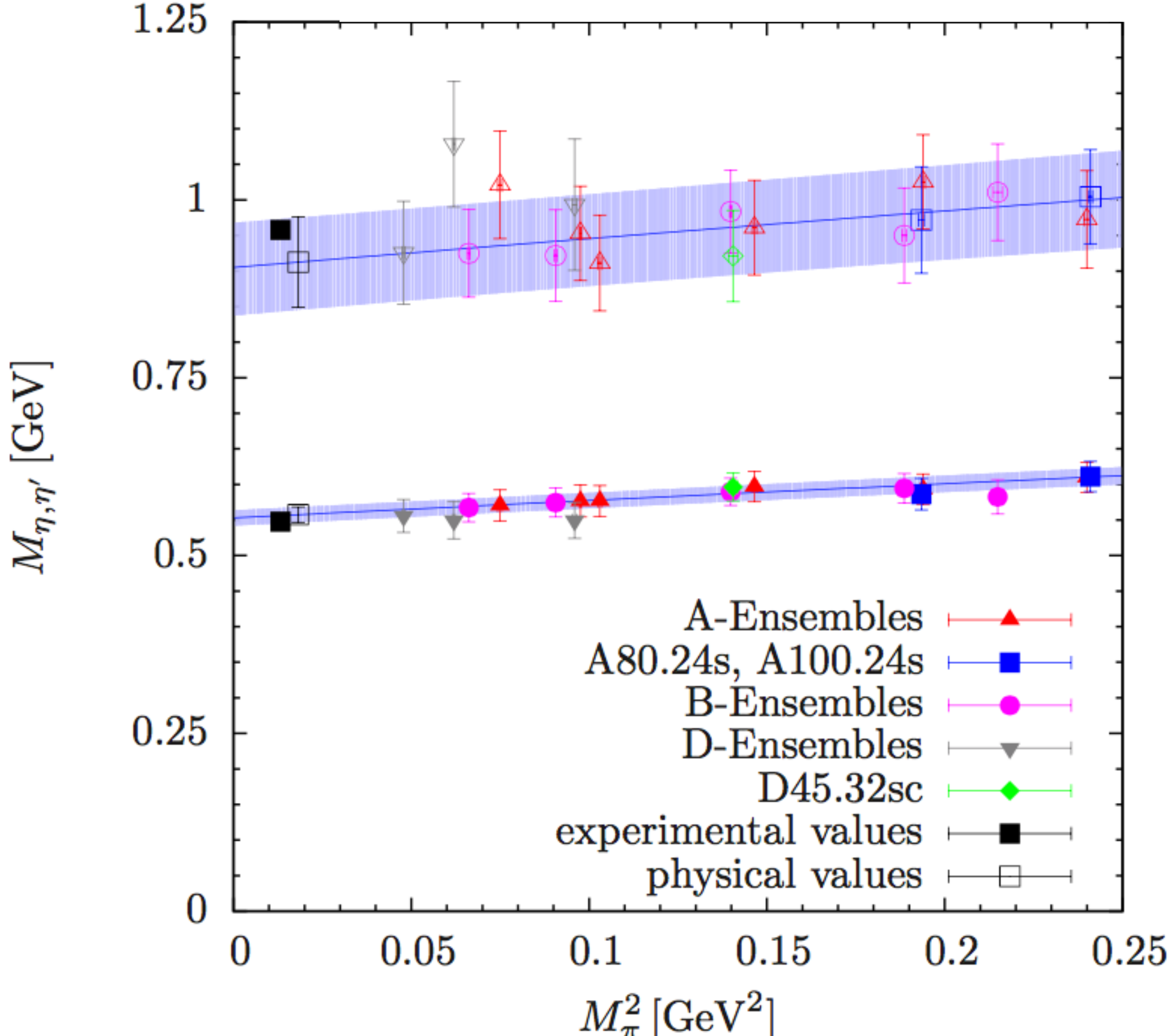
$$m_{\eta'}^2 = \frac{2N_f}{F_\pi^2} \chi_t^{\text{qu}}$$

Successful  
at T=0



Topology observable effects:

EVIDENCES OF THE EXPLICIT  $U(1)_A$  BREAKING.

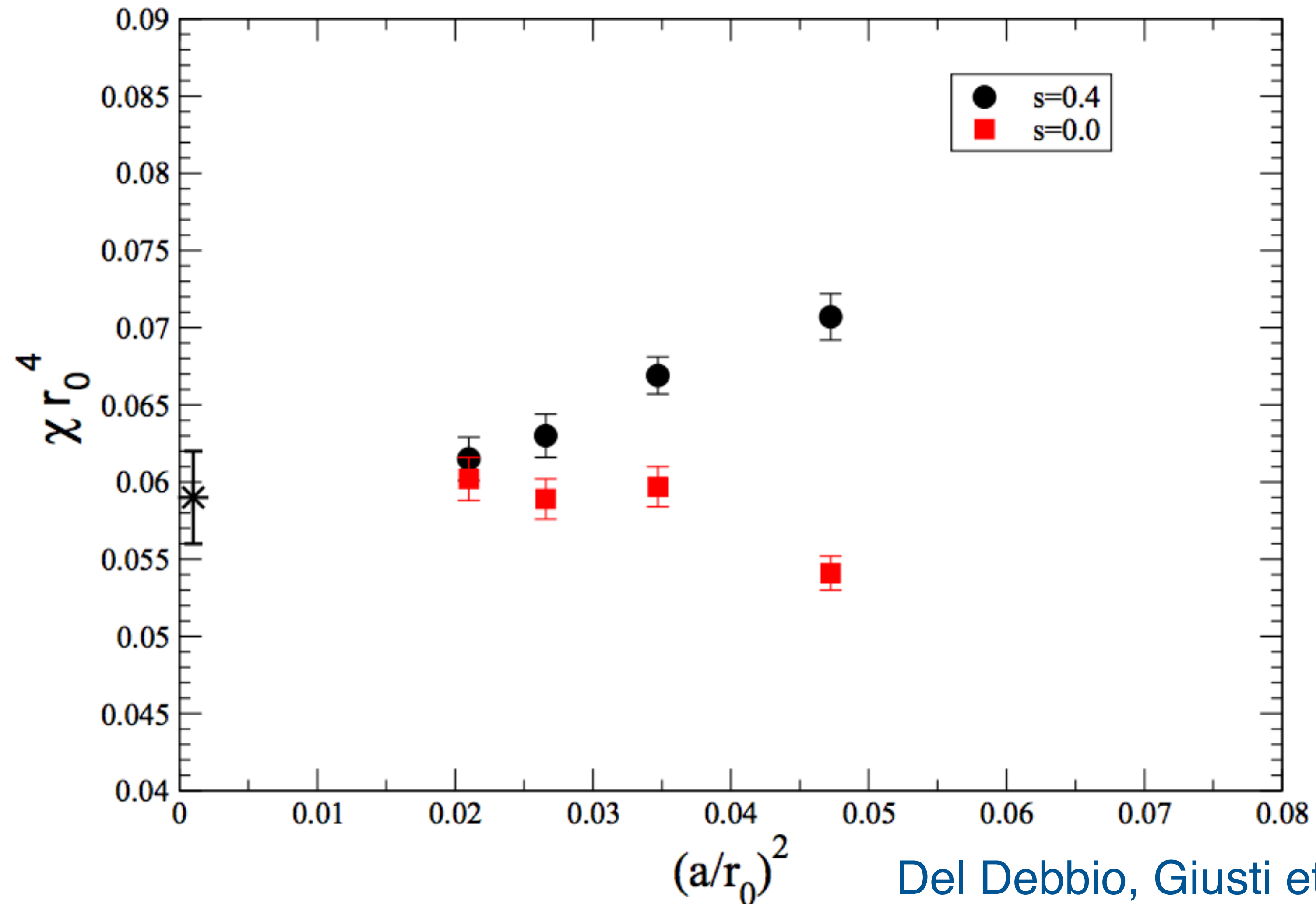




# Topology observable effects:

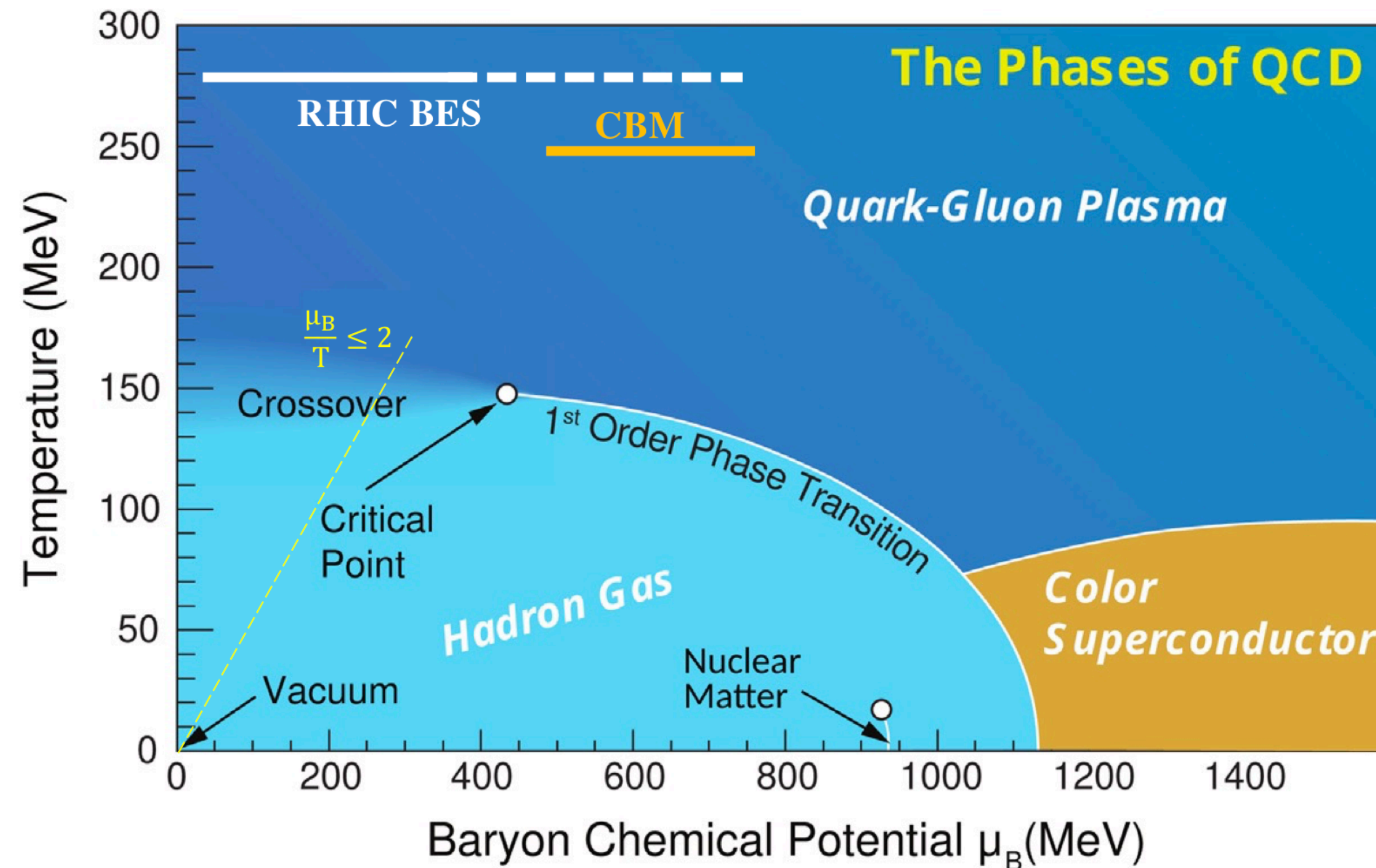
EVIDENCES OF THE EXPLICIT  $U(1)_A$  BREAKING.

$\chi = (191 \pm 5 \text{ MeV})^4$ , Yang-Mills Topological Susceptibility



Del Debbio, Giusti et al. (2015)

# 1. From low to high temperatures - A threshold in the plasma ?

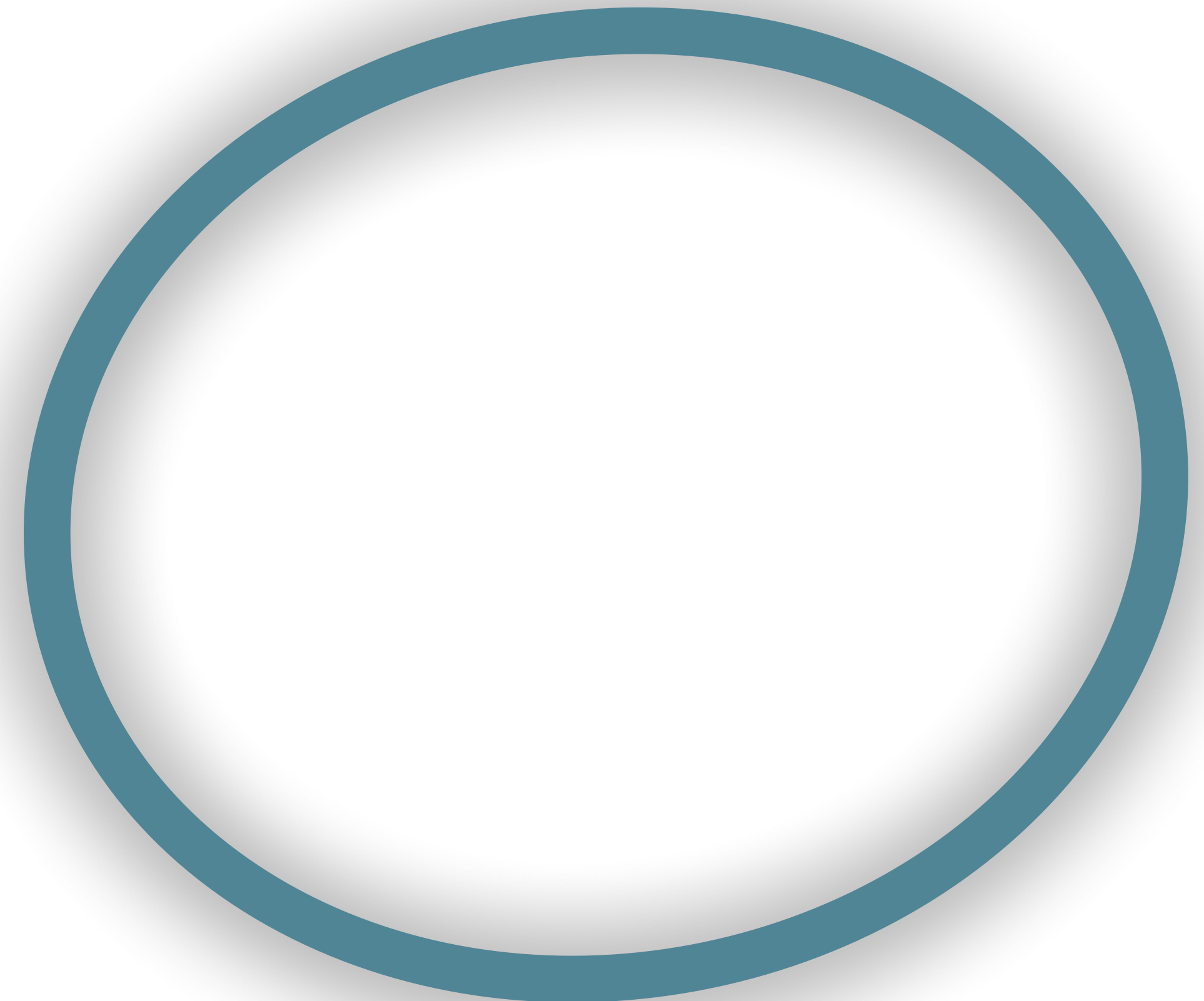


Almanool et al. 2022



# Topology from low to high Temperature

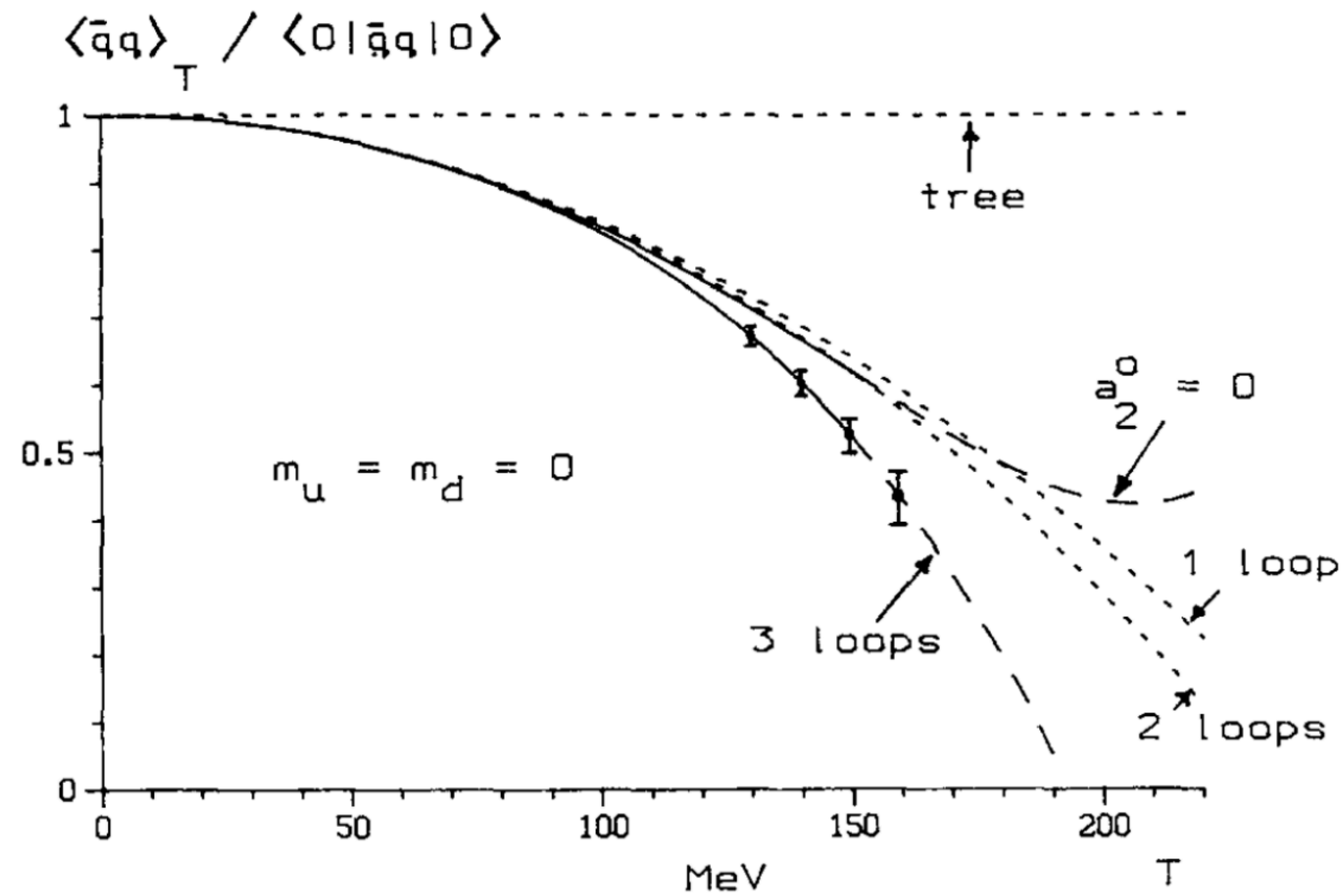
What happens to topology in the Quark Gluon Plasma?



# Finite temperature

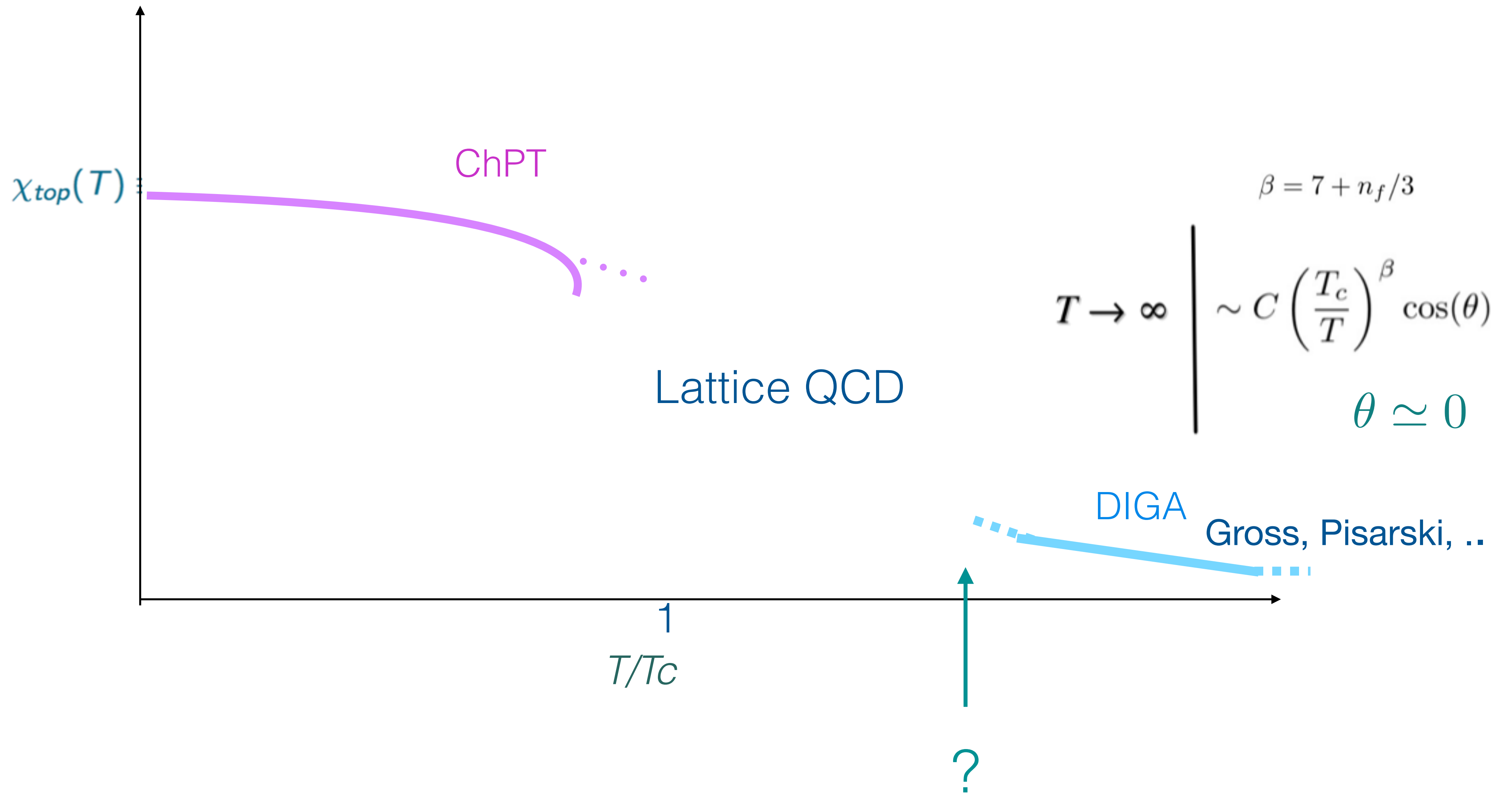
$$\frac{\chi_{top}(T)}{\chi_{top}} \stackrel{\text{NLO}}{=} \frac{m_\pi^2(T) f_\pi^2(T)}{m_\pi^2 f_\pi^2} = \frac{\langle \bar{q}q \rangle_T}{\langle \bar{q}q \rangle}$$

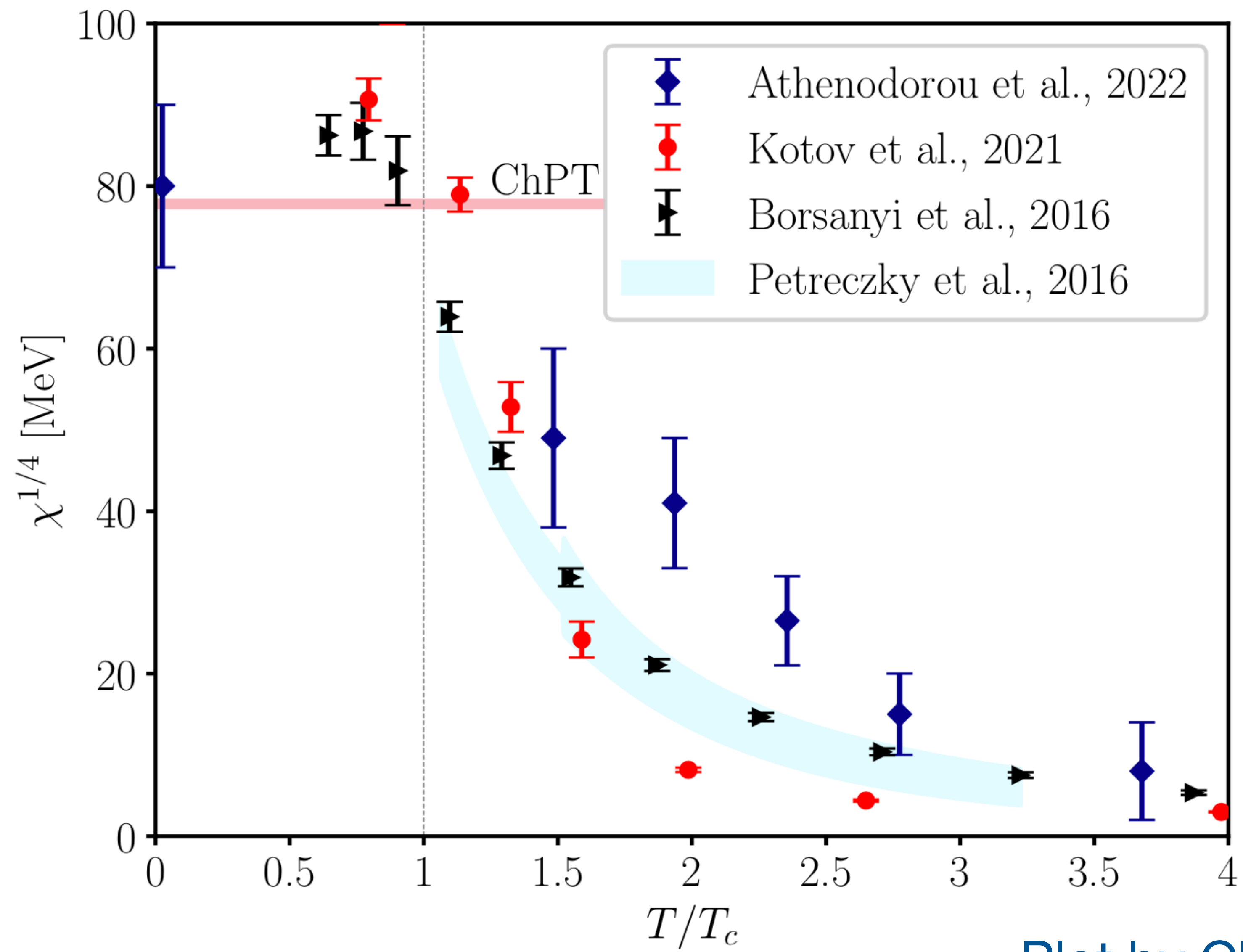
$$\langle \bar{q}q \rangle \stackrel{m \rightarrow 0}{=} \langle 0 | \bar{q}q | 0 \rangle \left( 1 - \frac{T^2}{8F^2} - \frac{T^4}{384F^4} - \frac{T^6}{288F^6} \ln \frac{\Lambda_q}{T} + \mathcal{O}(T^8) \right)$$





What do we know about  $\chi_{top}(T) \equiv \left. \frac{\partial^2 F(\theta, T)}{\partial \theta^2} \right|_{\theta=0}$



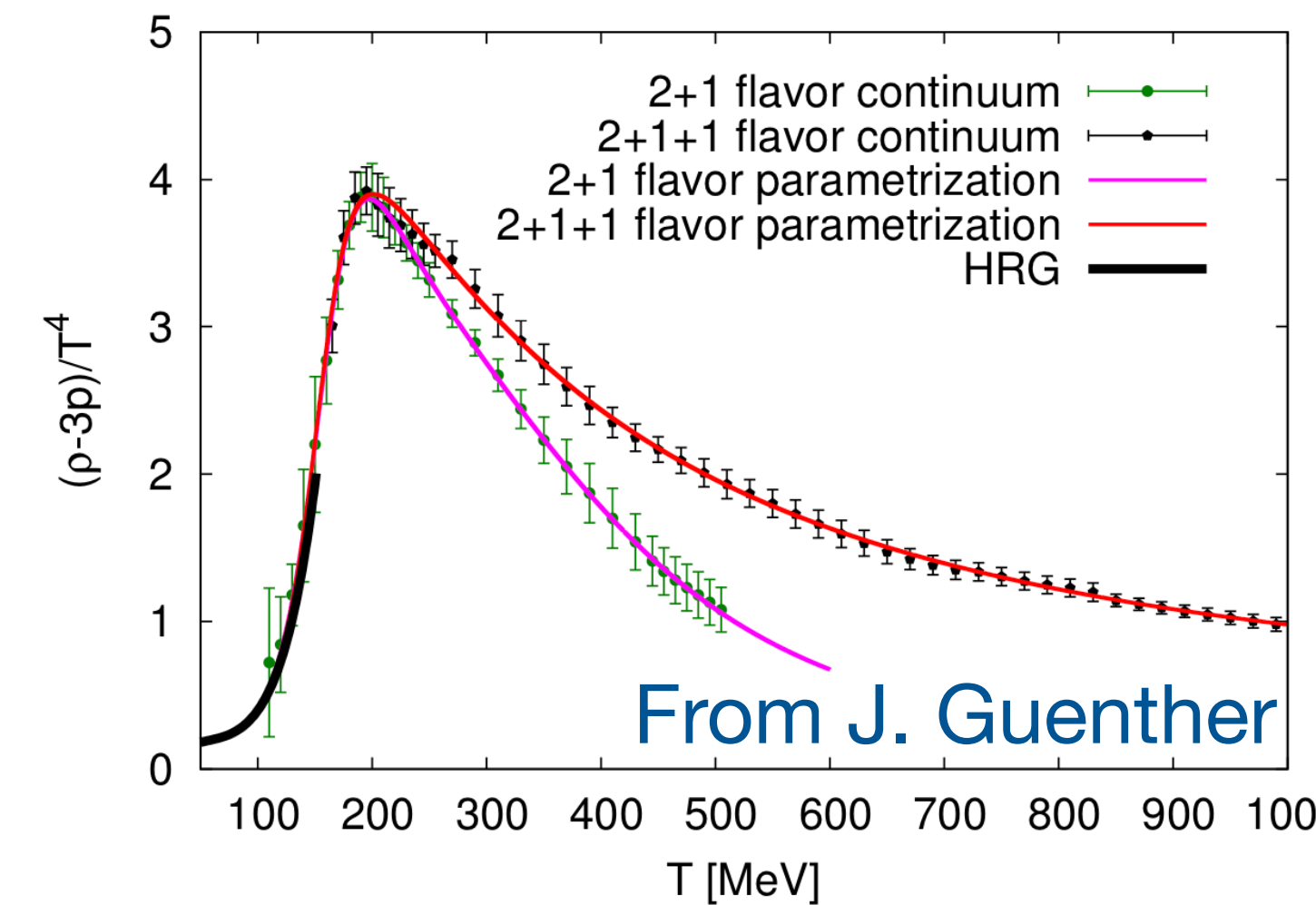
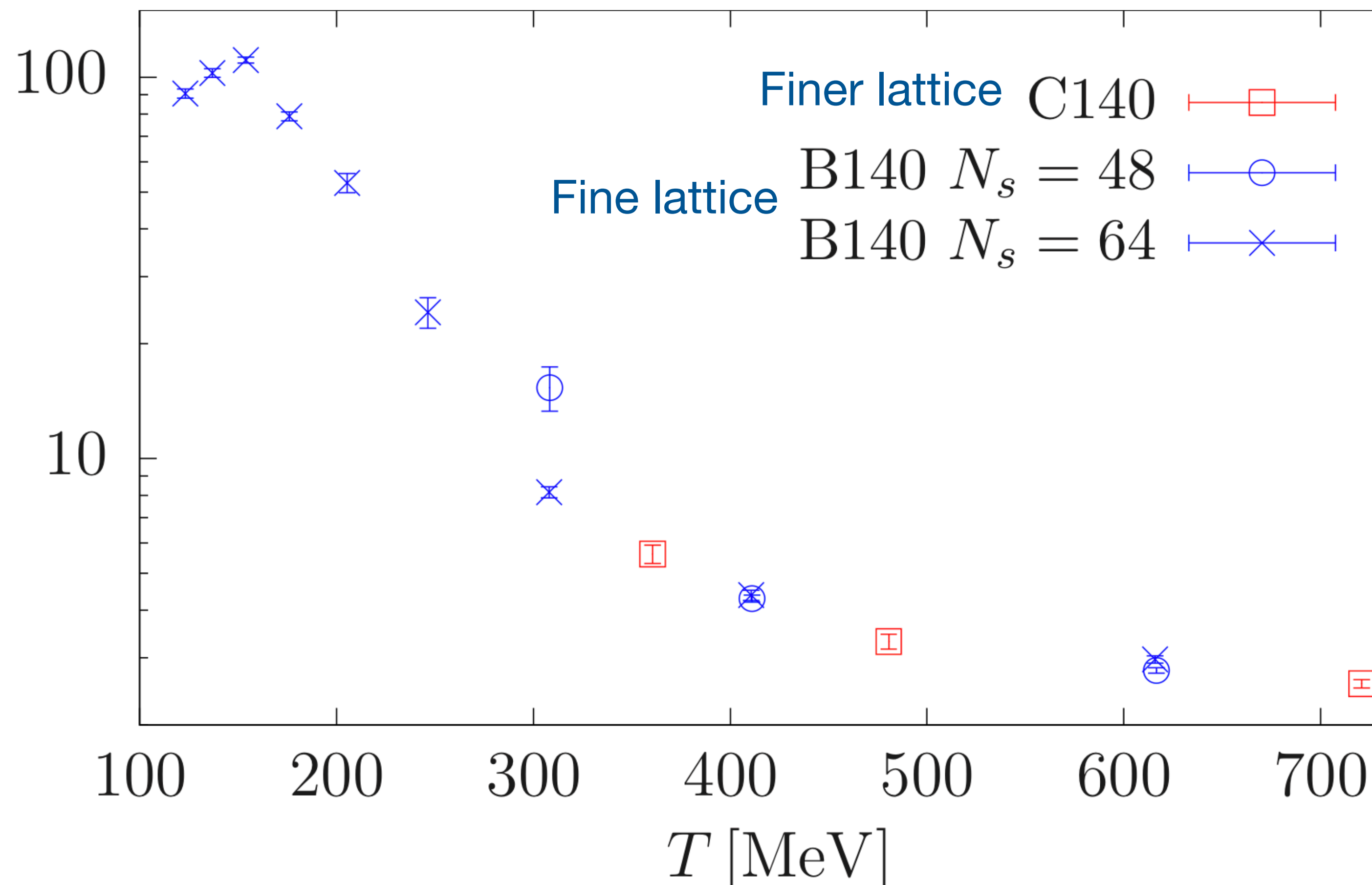


Plot by Claudio Bonanno

# Systematics from twisted mass Wilson fermions

2+1+1 flavours

$\chi_{top}^{1/4}$  MeV



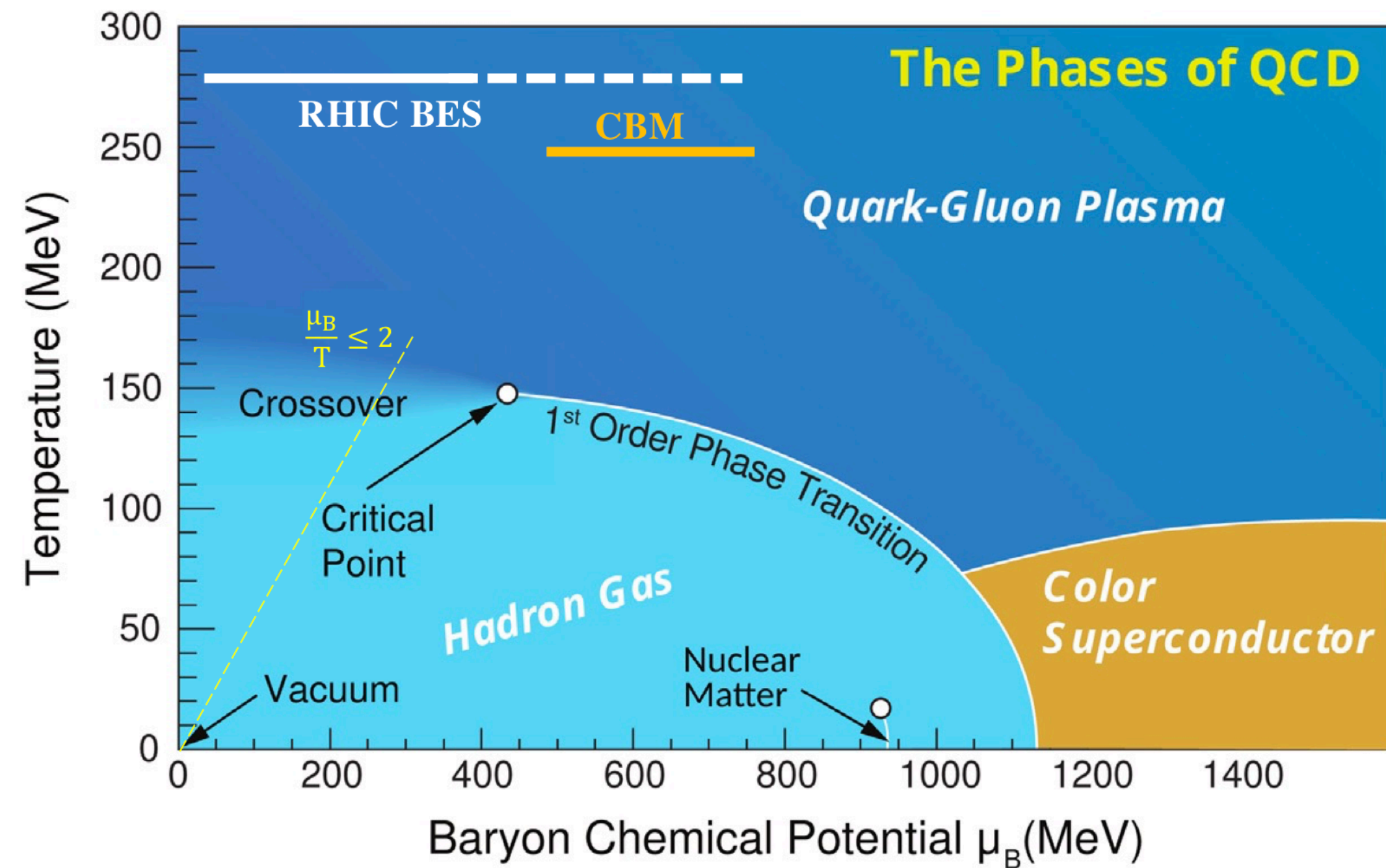
Kotov MpL Trunin (in progress)

# A threshold in the plasma ?

Alexandru and Horvath (2019-2021);

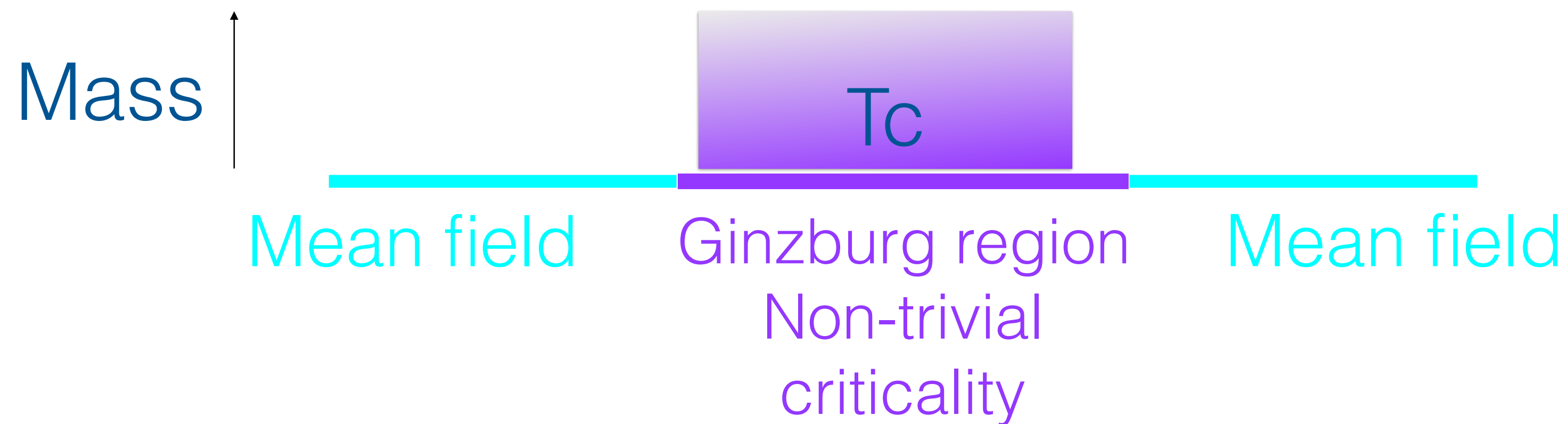
Glozman et al; Glozman, Philipsen, Pisarski (2016-2022);

Burger, Kotov, MpL, Trunin (2018-2022..)





## Generic features of a critical region



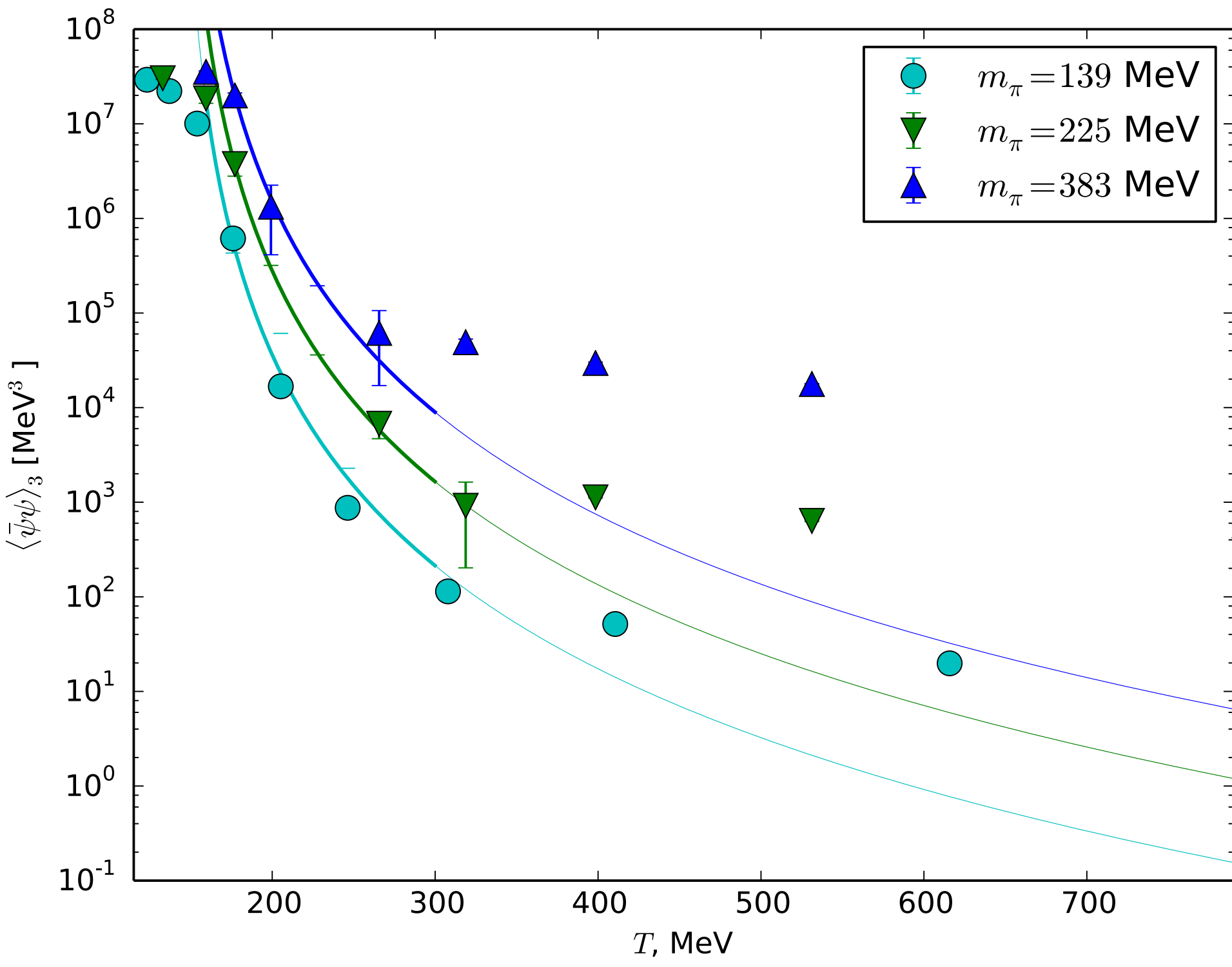
- . Is the 'threshold' related with the crossover from the Ginzburg region to mean field ?
- . Is the threshold related with topology?

# Searching for the scaling window in temperature

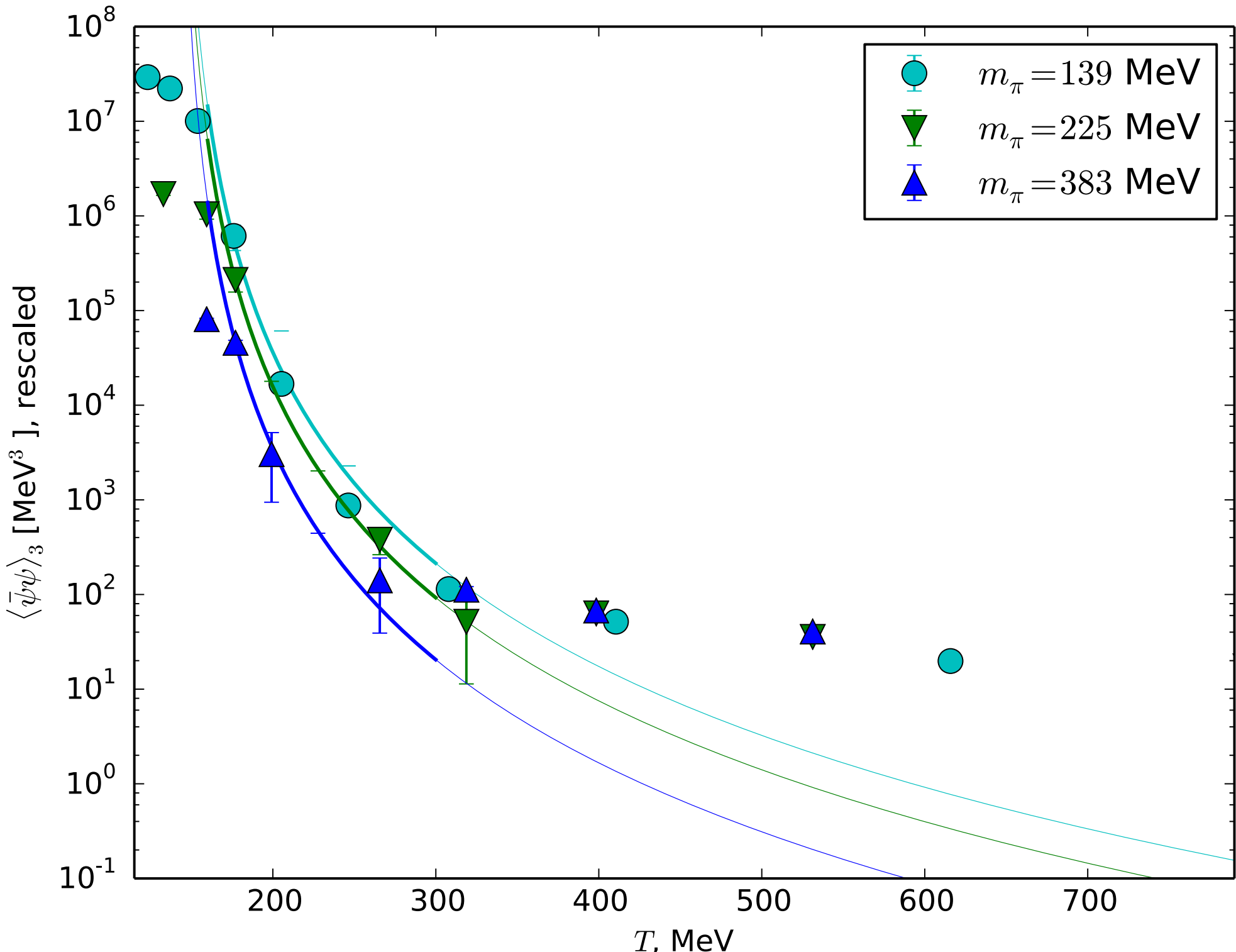
Kotov, MpL, Trunin 2021

Ginzburg region  $T < 300$  MeV

Simple analytic behaviour  $T > 300$  MeV



$$\Delta_3 \propto t^{-\gamma-2\beta\delta}$$



$$\Delta_3 \propto m_\pi^6$$

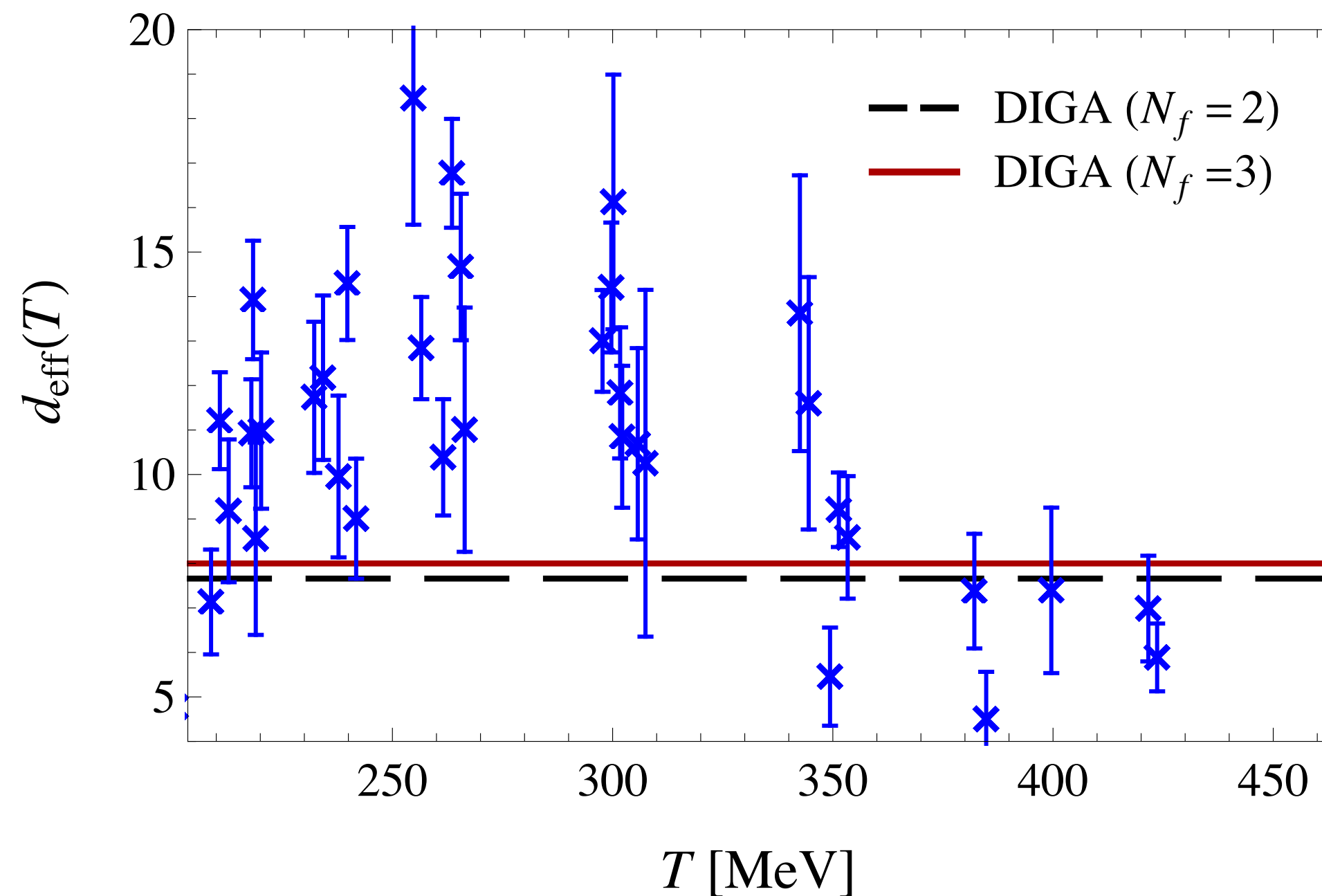
# Power-law decay for $T > 300$ MeV

For instanton gas

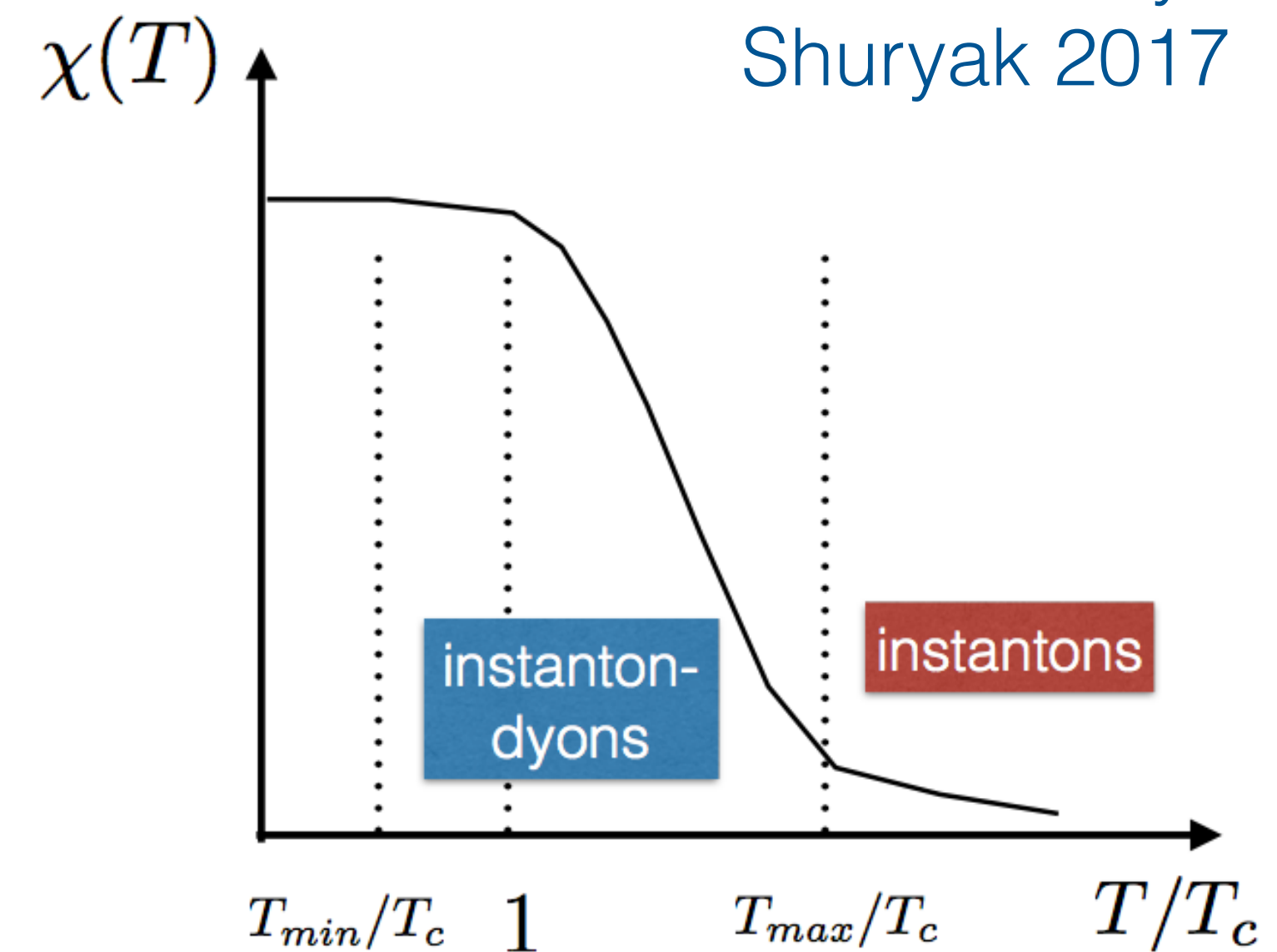
$$\chi^{0.25}(T) = aT^{-d(T)}$$

$$d(T) \equiv \text{const} \simeq \left(7 + \frac{N_f}{3}\right)$$

$$d(T) = -T \frac{d}{dT} \ln \chi^{0.25}(T)$$



Possibly consistent with instant-dyon?  
Shuryak 2017

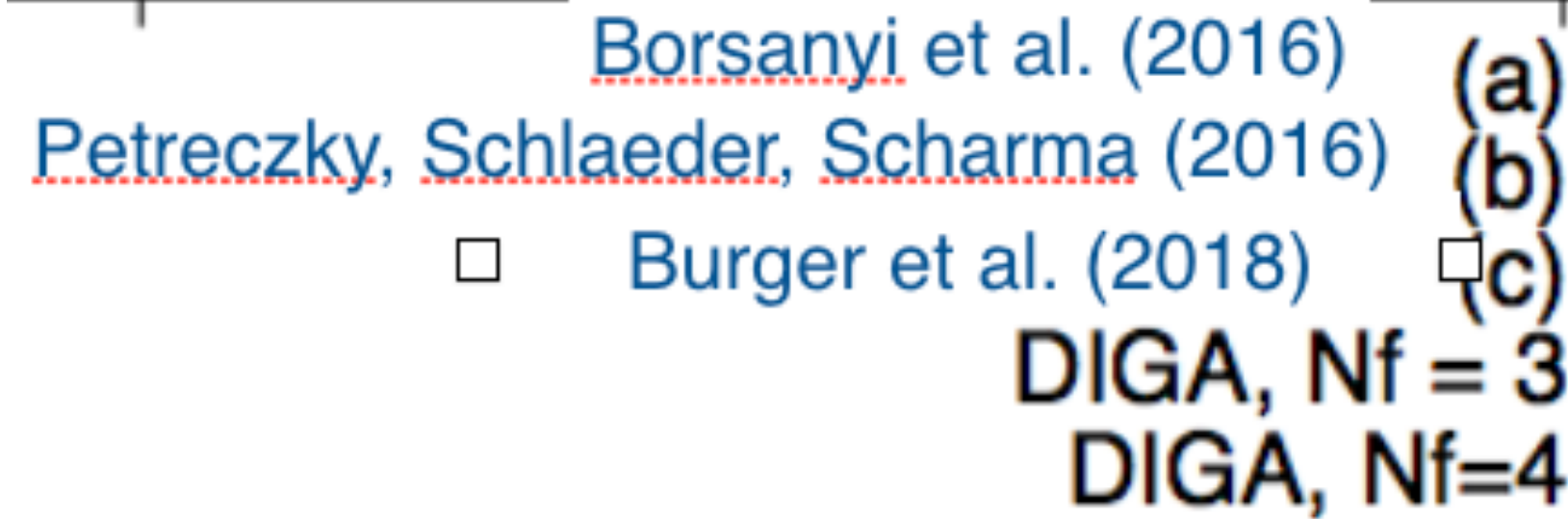


Faster decrease before DIGA sets in

QCD - Summary of b parameter

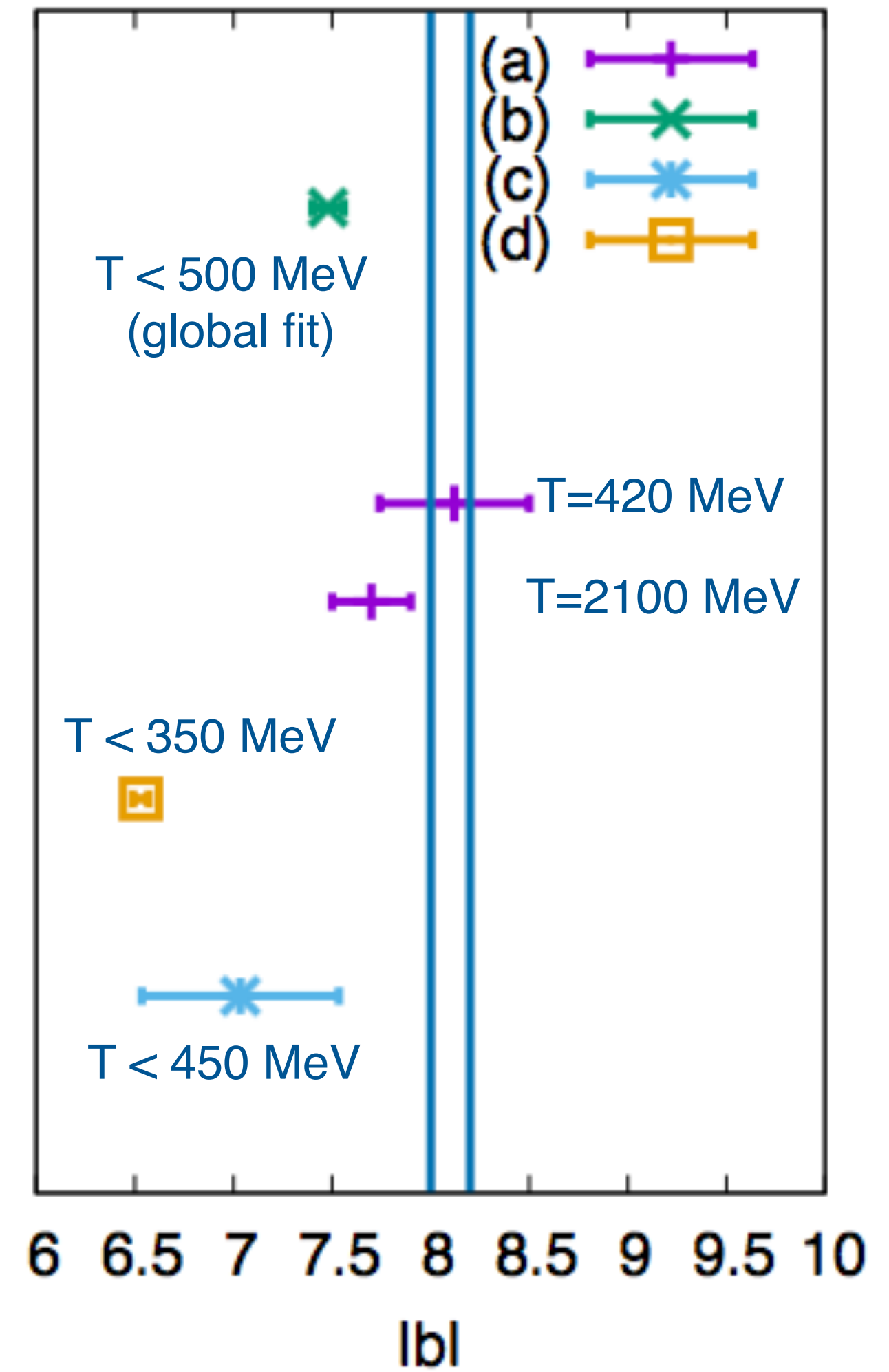
$$\chi(T) = A T^b$$

Y. Taniguchi, K. Kanaya, H. Suzuki and T. Umeda (2017) (d),



For  $T > 300$  MeV the DIGA exp is approached from below

$T_c < T < 250 - 300$  MeV ??





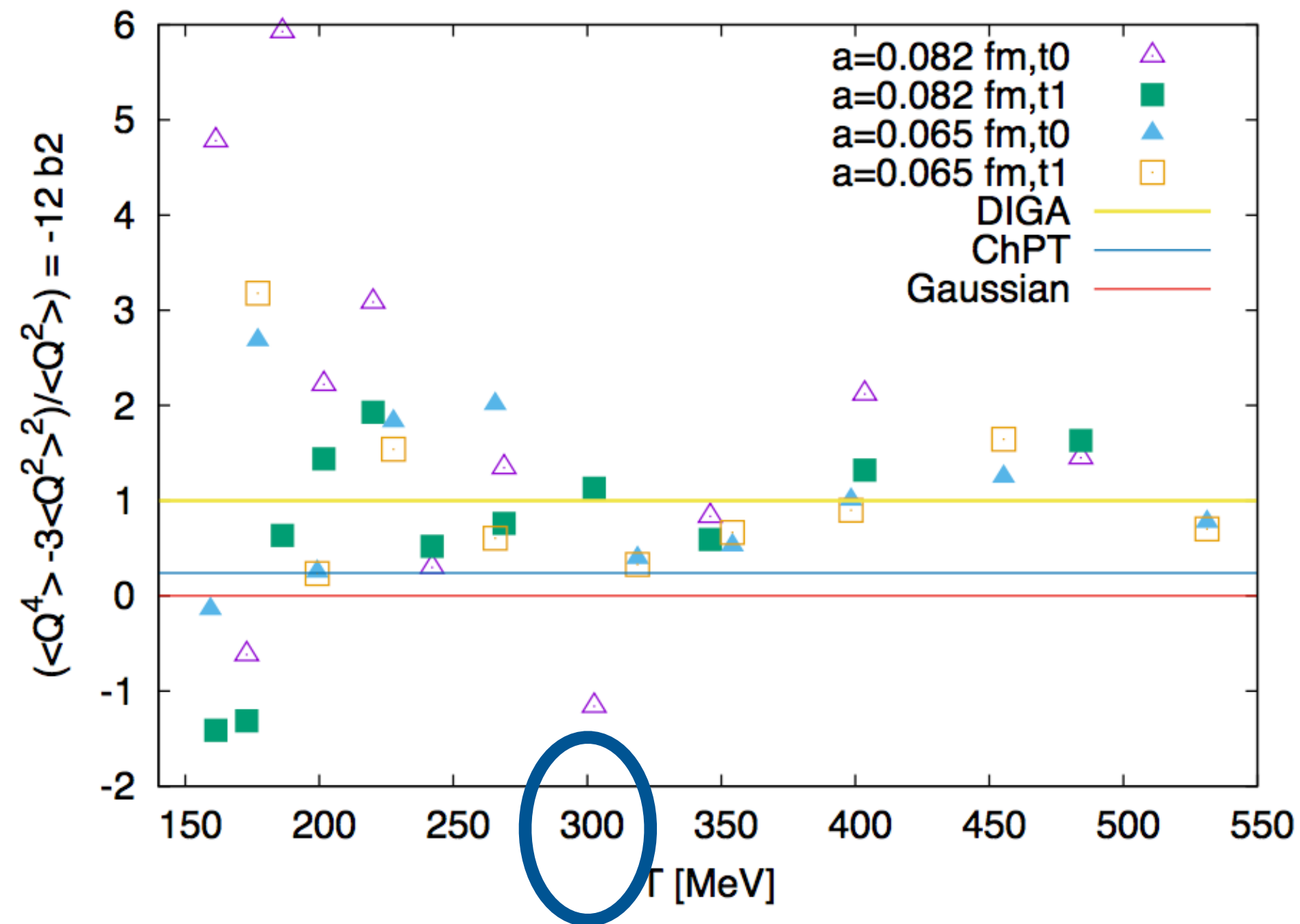
# Further evidence of a fast crossover in topology

Beyond topological susceptibility

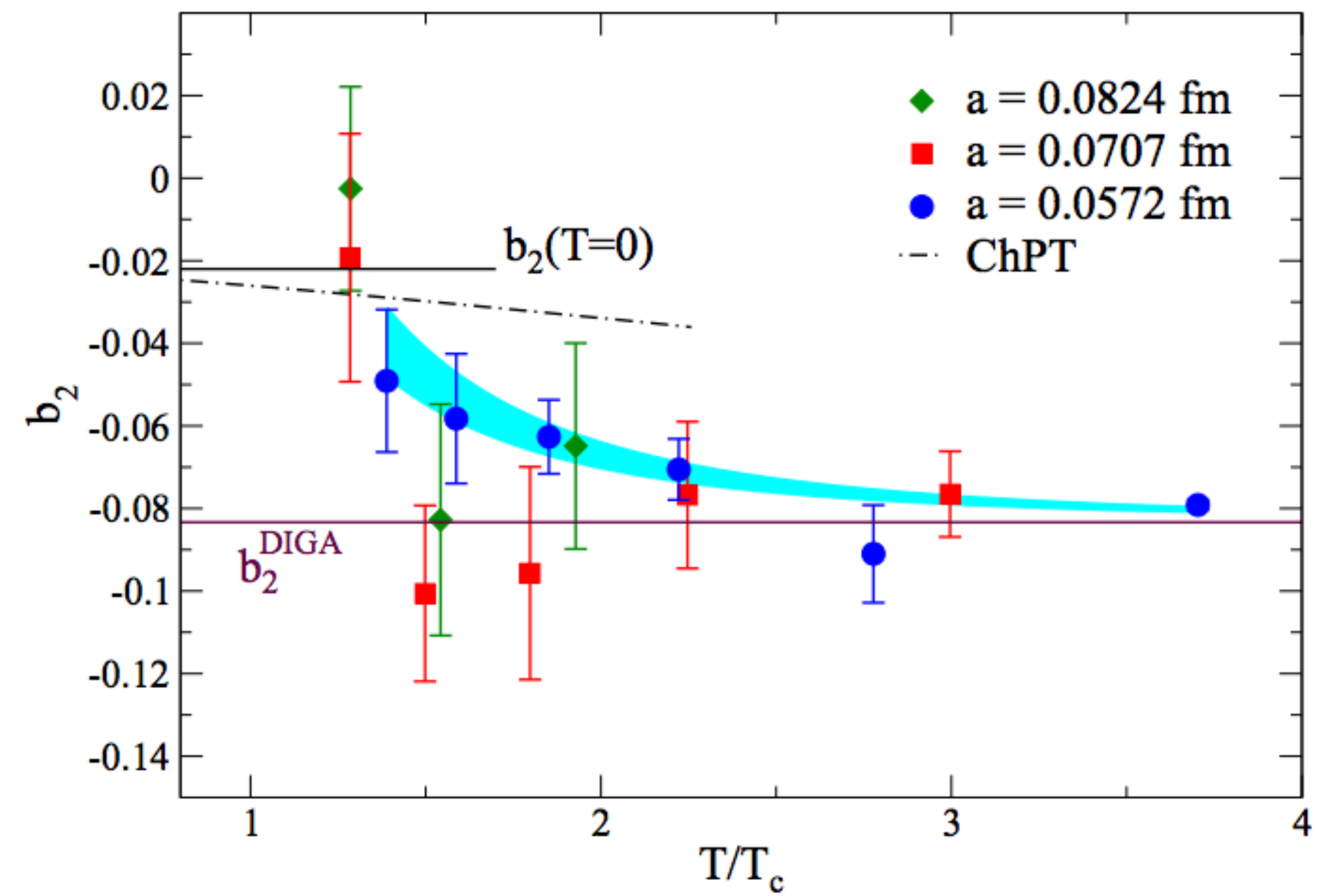
$T > 250\text{-}300\text{ MeV}$

$$C_n = (-1)^{n+1} \frac{d^{2n}}{d\theta^{2n}} F(\theta, T) \Big|_{\theta=0} = \langle Q^{2n} \rangle_{conn.}$$

d'Elia, Vicari (2013)



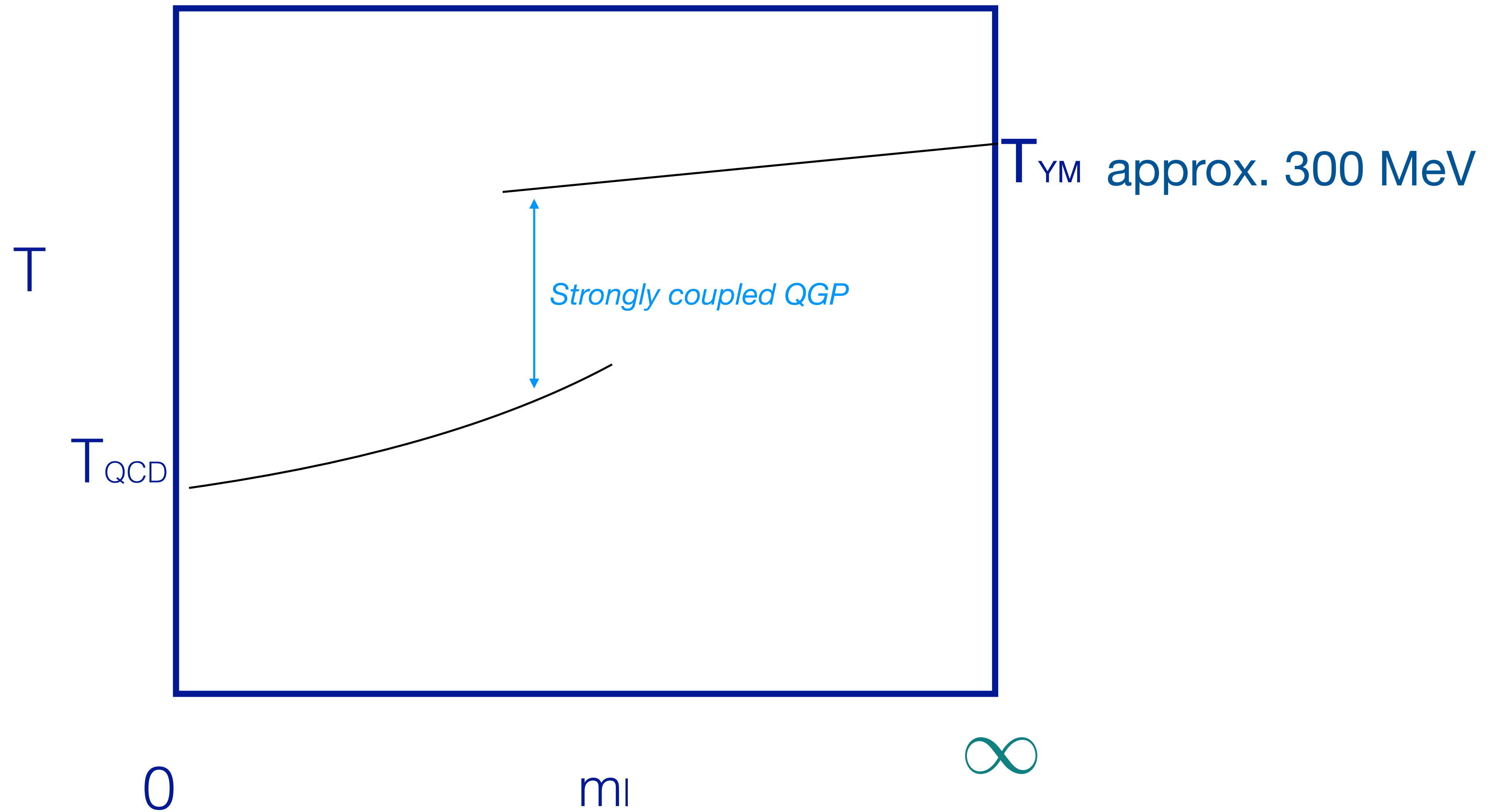
Trunin et al (2018)



Bonati et al. (2016)

.. a speculation...

Is a resilient YM topology producing a threshold in full QCD?



# Topology from low to high Temperature - effects on the spectrum

In the hadronic phase topology solves the puzzle by explicit breaking  $U(1)_A$   $\eta'$

What happens to  $\eta'$  in the Quark Gluon Plasma?

PHYSICAL REVIEW D

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1 MAY 1996

## Return of the prodigal Goldstone boson

J. Kapusta

*School of Physics and Astronomy, University of Minnesota, Minneapolis, Minnesota 55455*

D. Kharzeev

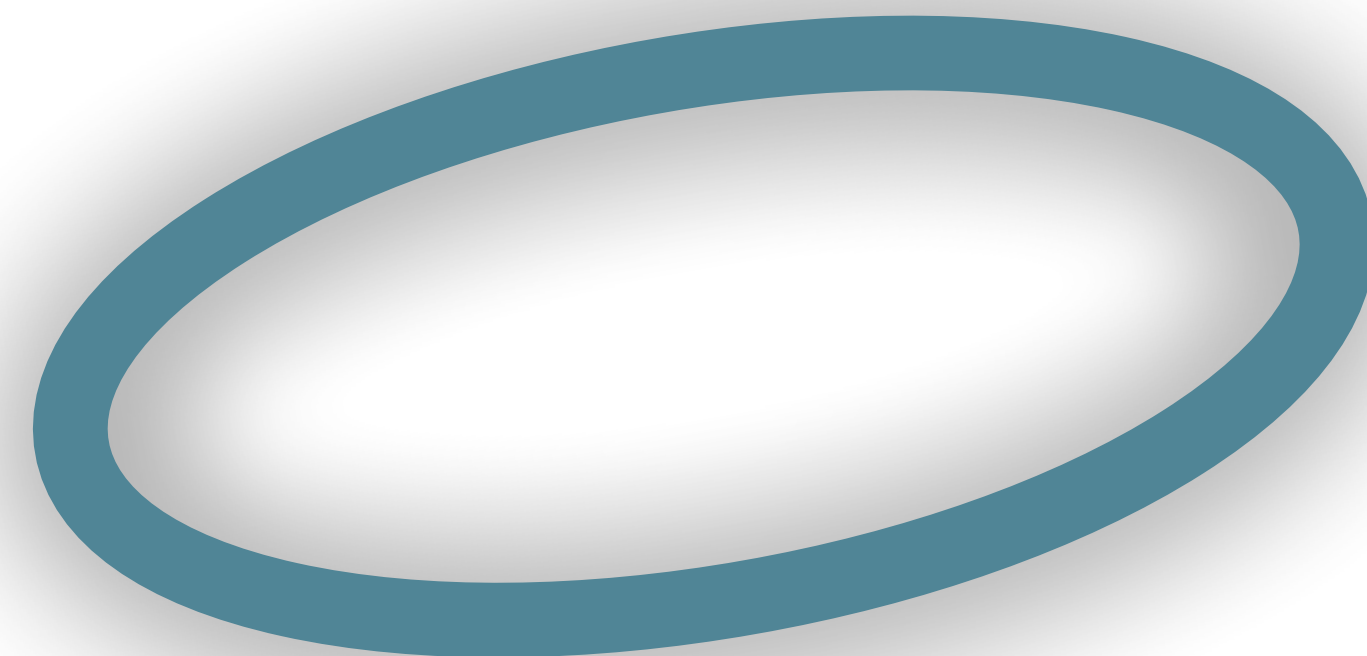
*Theory Division, CERN, Geneva, Switzerland  
and Fakultät für Physik, Universität Bielefeld, Bielefeld, Germany*

L. McLerran

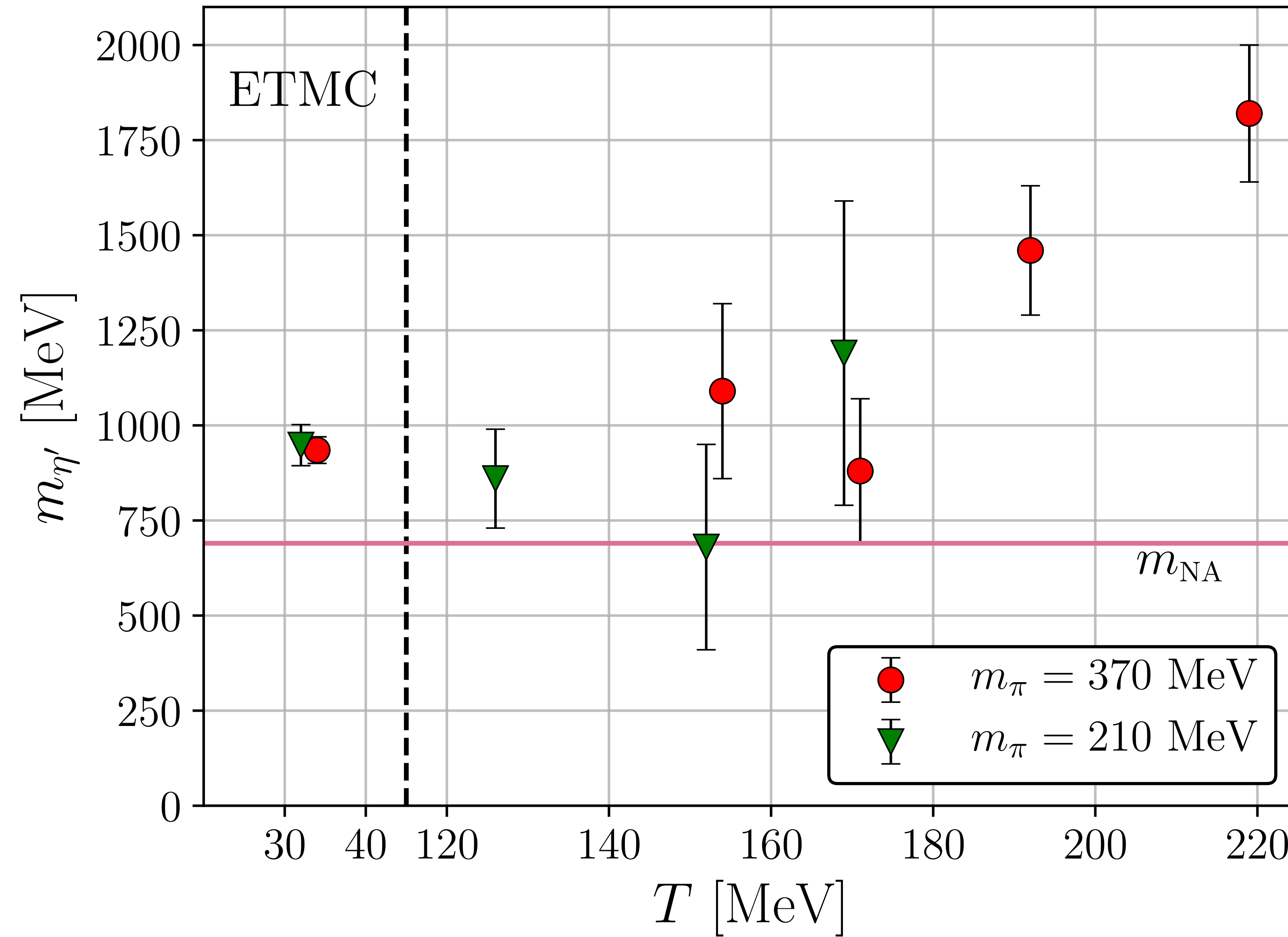
*School of Physics and Astronomy, University of Minnesota, Minneapolis, Minnesota 55455*

(Received 14 July 1995)

We propose that the mass of the  $\eta'$  meson is a particularly sensitive probe of the properties of finite energy density hadronic matter and quark-gluon plasma. We argue that the mass of the  $\eta'$  excitation in hot and dense matter should be small, and, therefore, that the  $\eta'$  production cross section should be much increased relative to that for  $pp$  collisions. This may have observable consequences in dilepton and diphoton experiments.



# Evidence for topology suppression in the spectrum



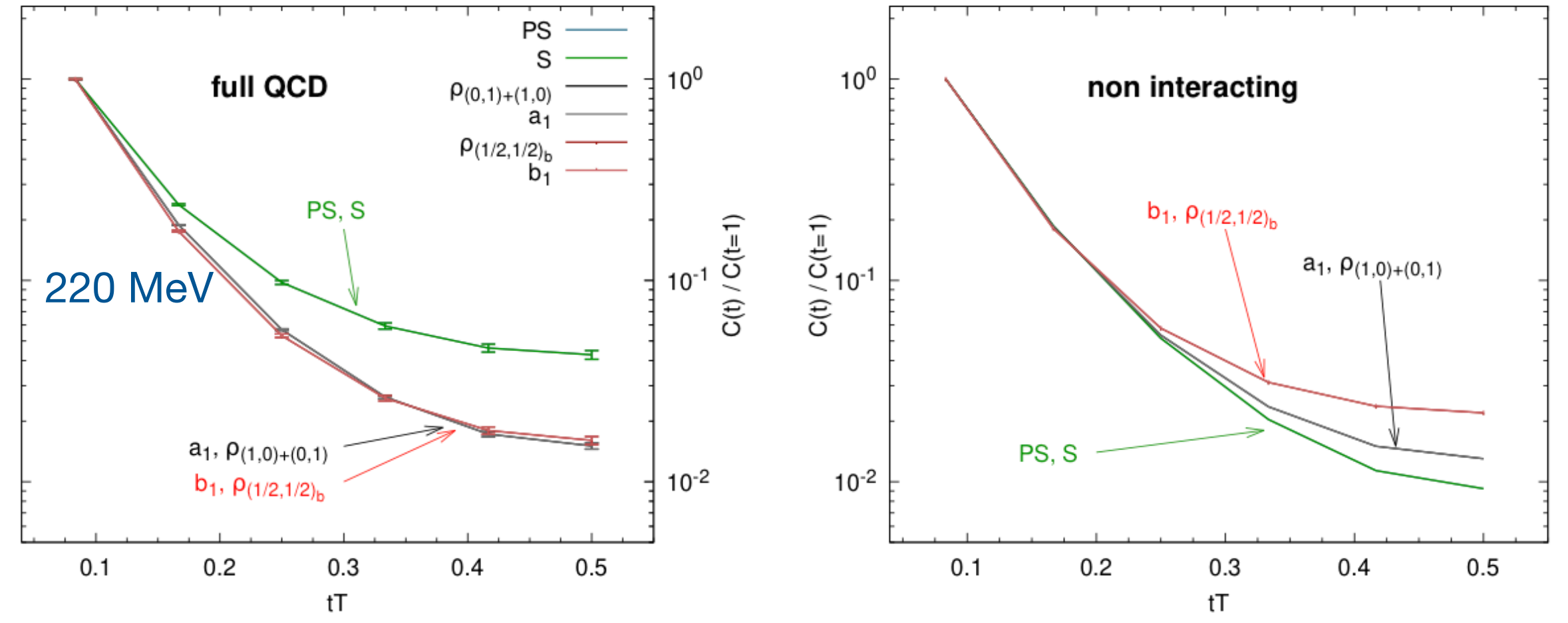
Non anomalous:

$$\eta' \simeq 700 \text{ MeV}$$

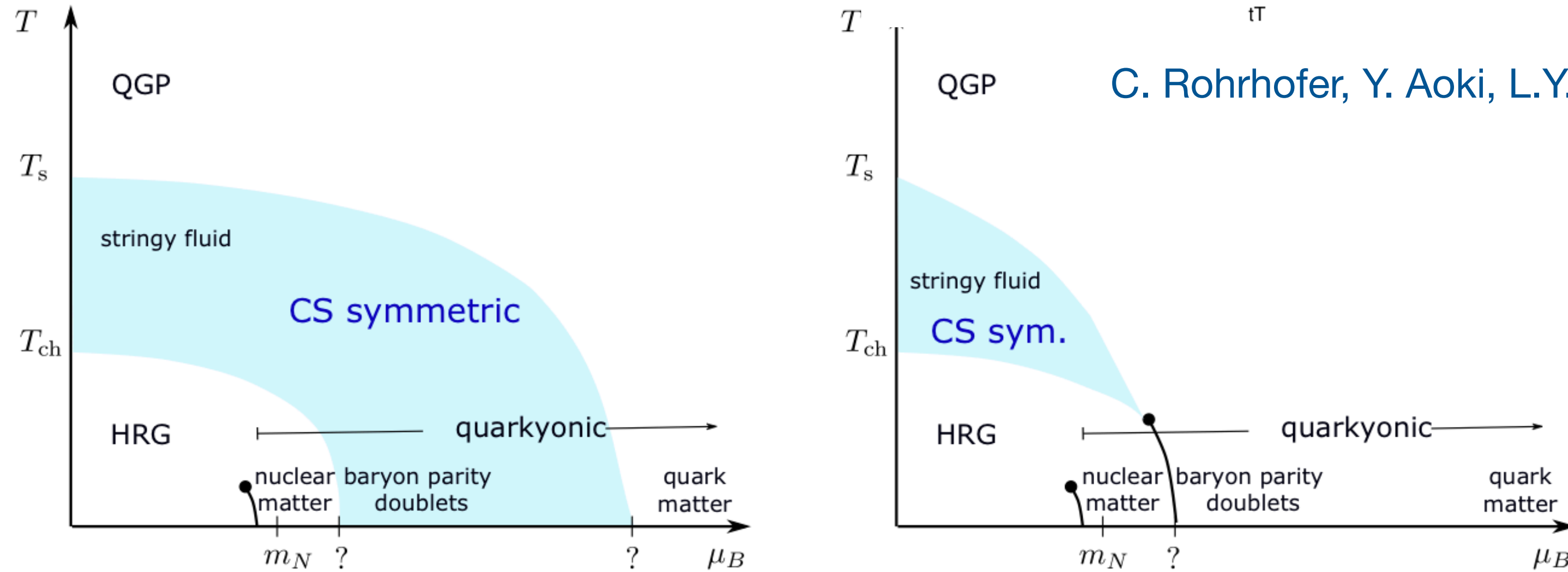
(strange only)



# Spectrum and the anomalous threshold



C. Rohrhofer, Y. Aoki, L.Y. Glozman and S. Hashimoto, 2021



**Figure 8:** Possibilities for the QCD phase diagram with a chiral spin and  $SU(4)$ -symmetric band.

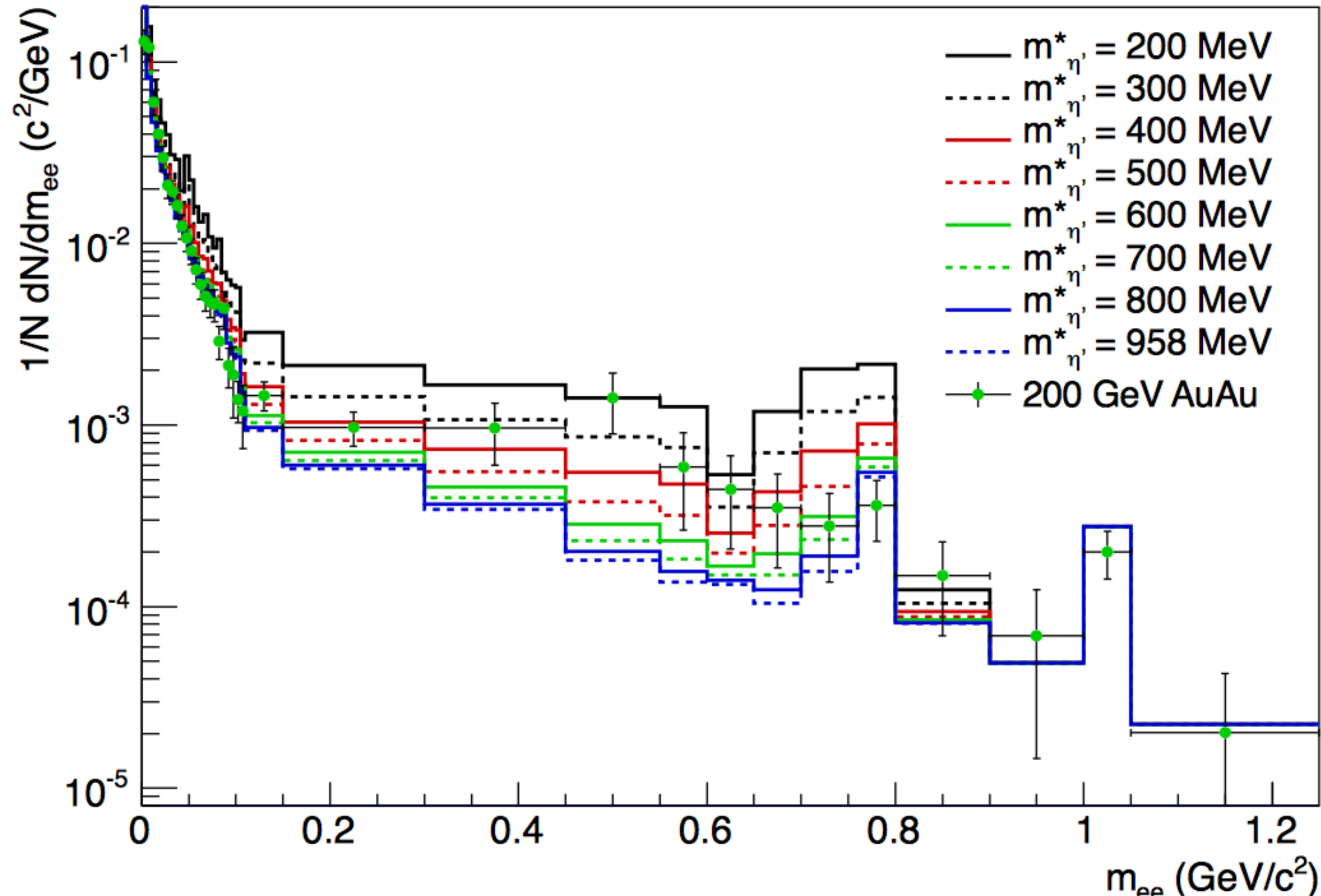
Philipsen et al. 2022

From low to high densities -  
Experimental results and  
first theoretical analysis

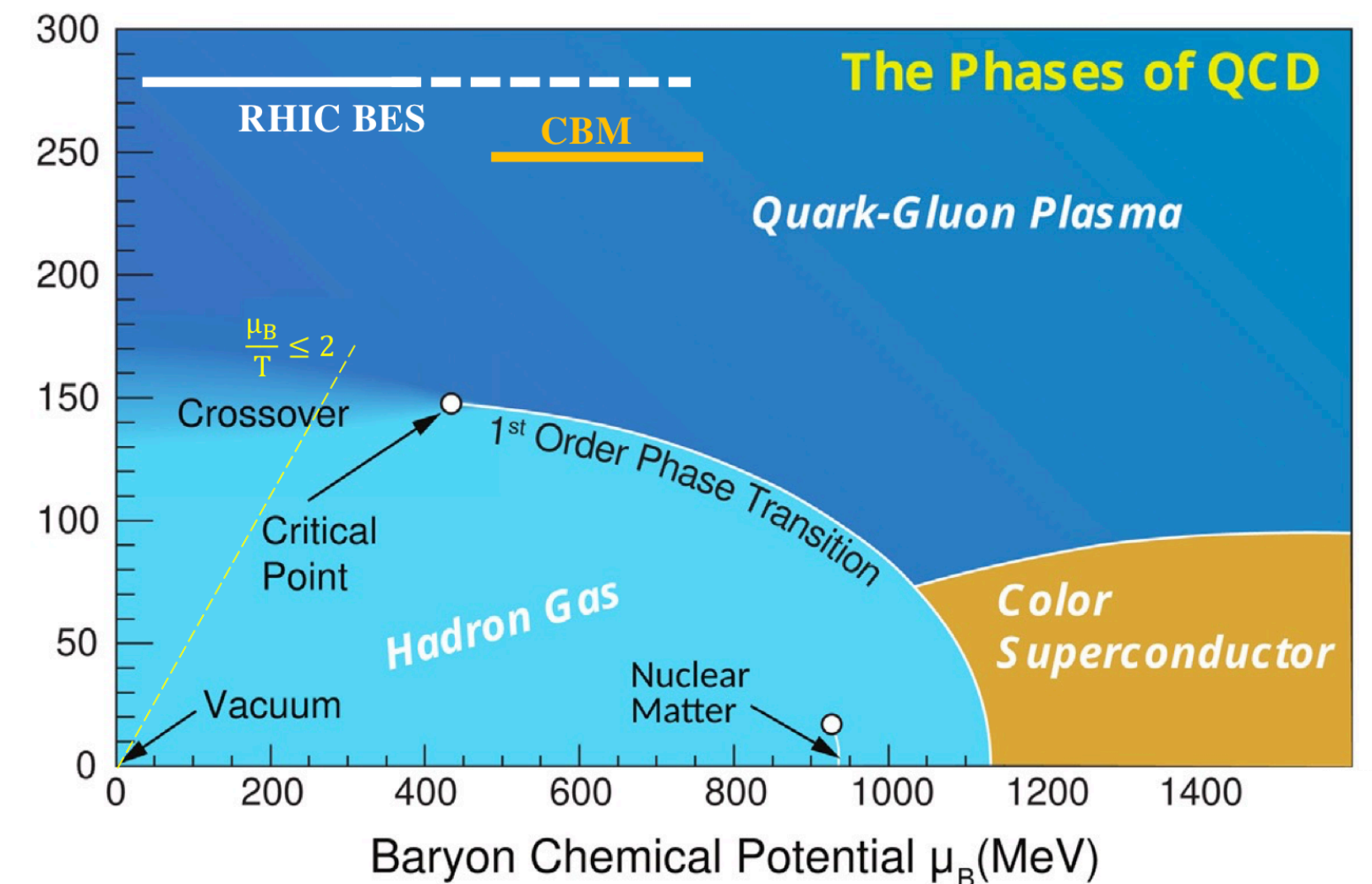
# Indication of topology suppression in PHENIX@RHIC

Effects of chain decays, radial flow and  $U_A(1)$  restoration on the low-mass dilepton enhancement in  $\sqrt{s_{NN}}=200$  GeV Au+Au reactions

Mátyás Veszteg<sup>a,b,1</sup> Tamás Csörgő<sup>b,2</sup> Dénes Váti<sup>a,b,c,3</sup>

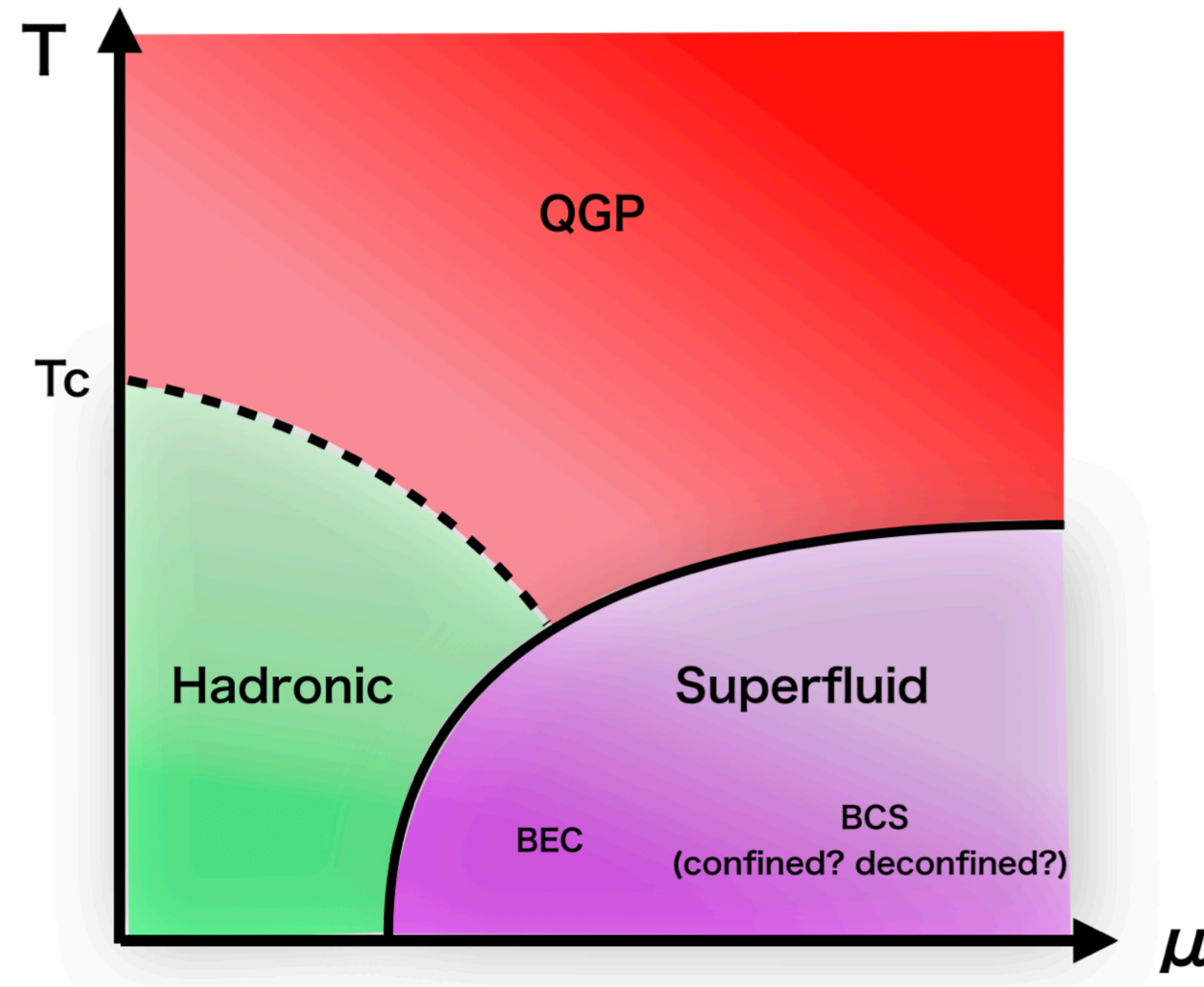


This is at finite density!



The sign problem hampers simulations in dense matter at low temperature

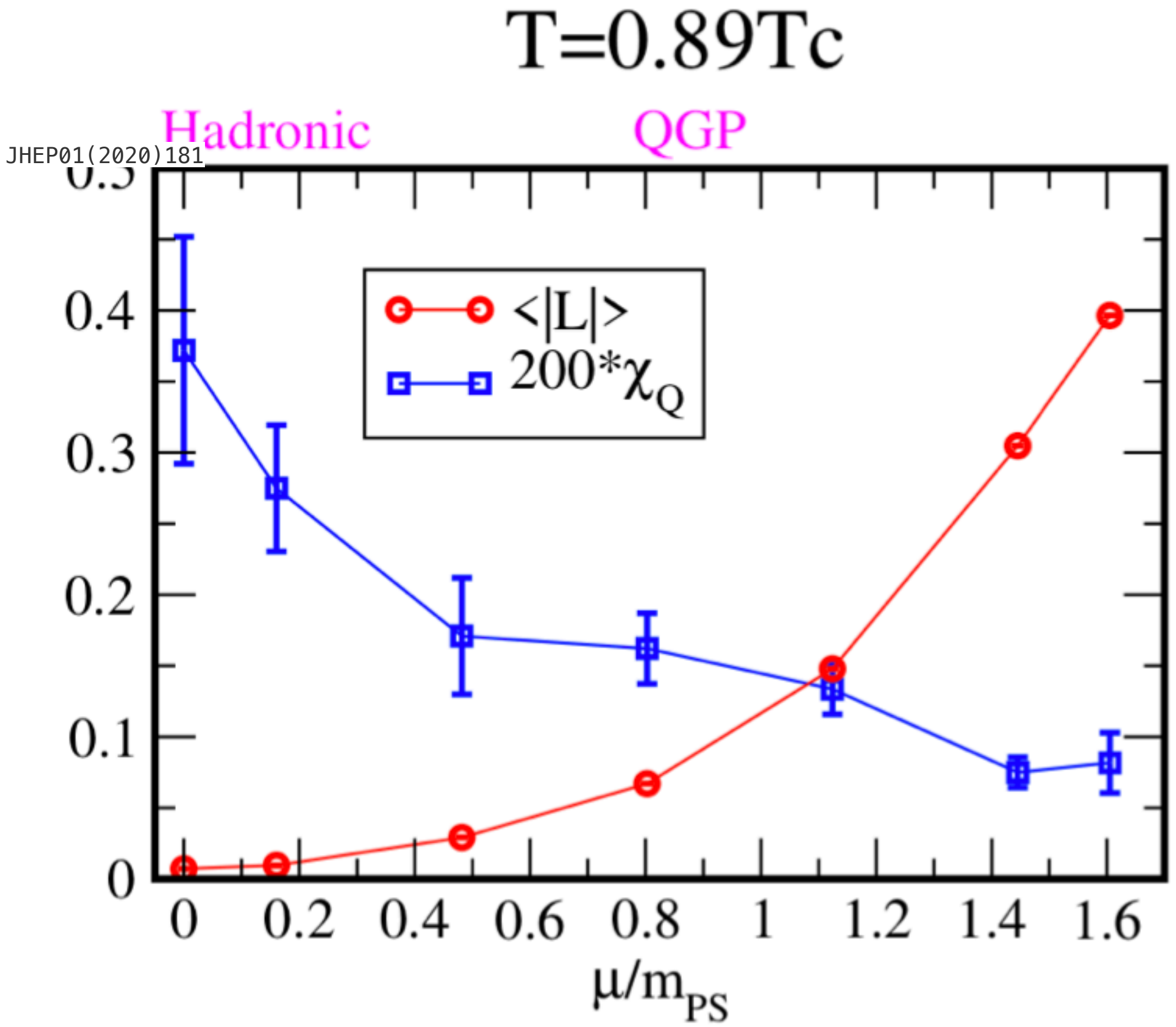
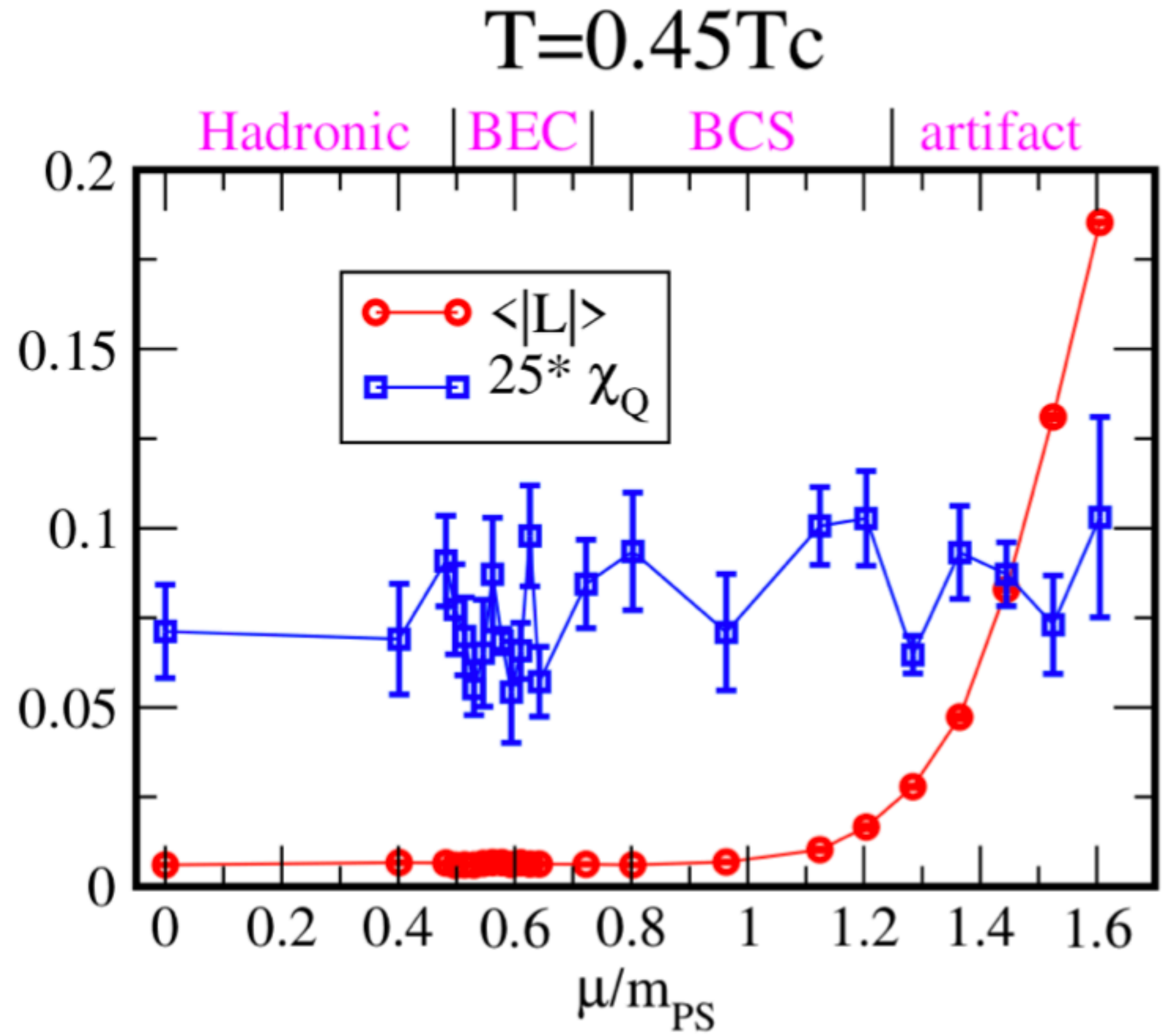
Some lessons may be learned from two colors QCD:





# Topology in two-color QCD: susceptibility ‘survives’ beyond the transition!

Role of instanton molecules?  
 Trivial in QCD2?  
 Need to go beyond susceptibility?

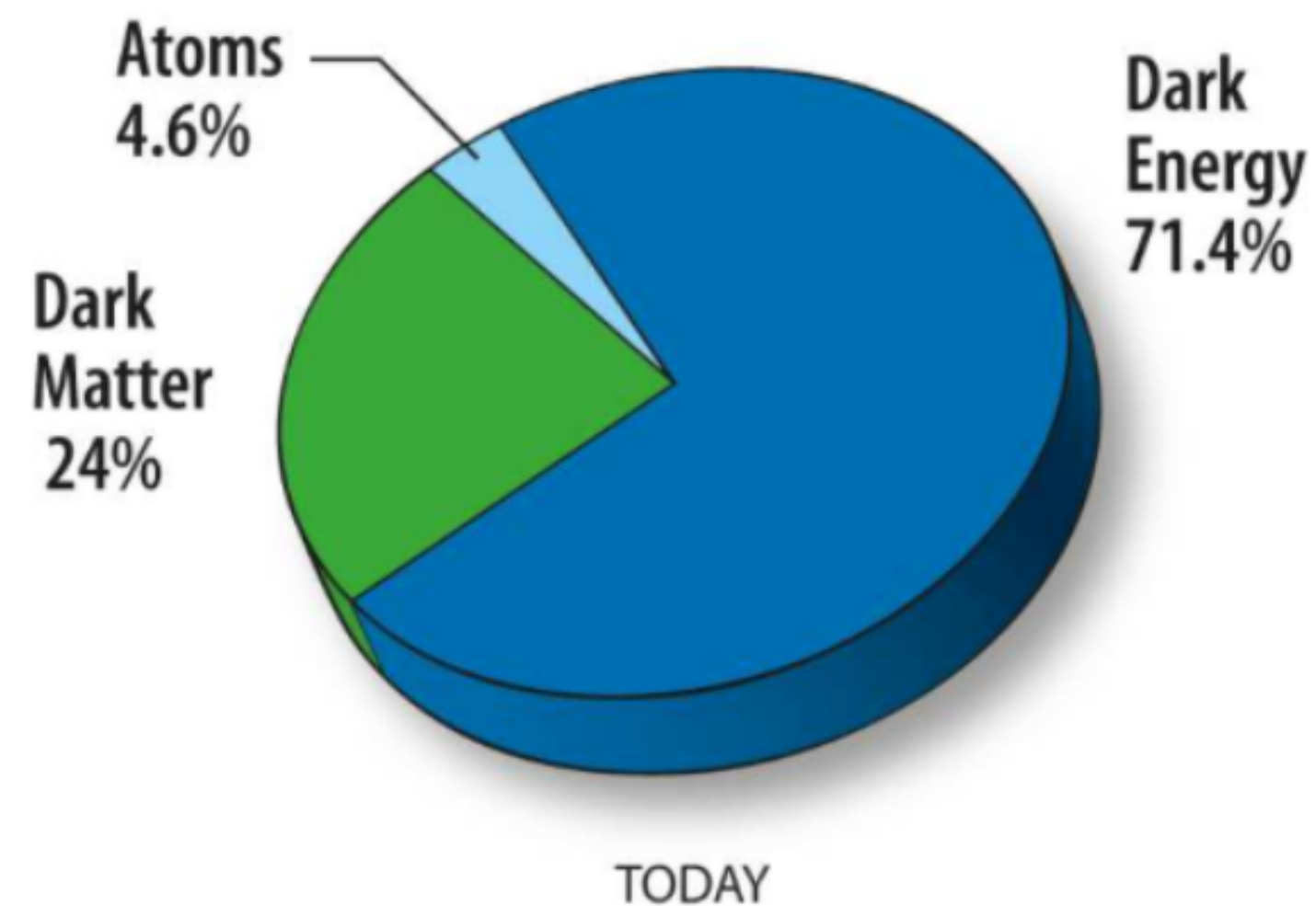


# 3. Topology and axioms

# the strong CP problem in QCD

..can be solved by introducing the AXION

a new particle which is a viable **dark matter candidate**



Crucial ingredient:  $\langle Q^2(T) \rangle$

# THE AXION MASS

At leading order in  $1/f_A$  – well justified as  $f_A \gtrsim 4 \times 10^8$  GeV

$$m_A^2(T) f_A^2 = \left. \frac{\partial^2 F(\theta, T)}{\partial \theta^2} \right|_{\theta=0} \equiv \chi_{top}(T)$$

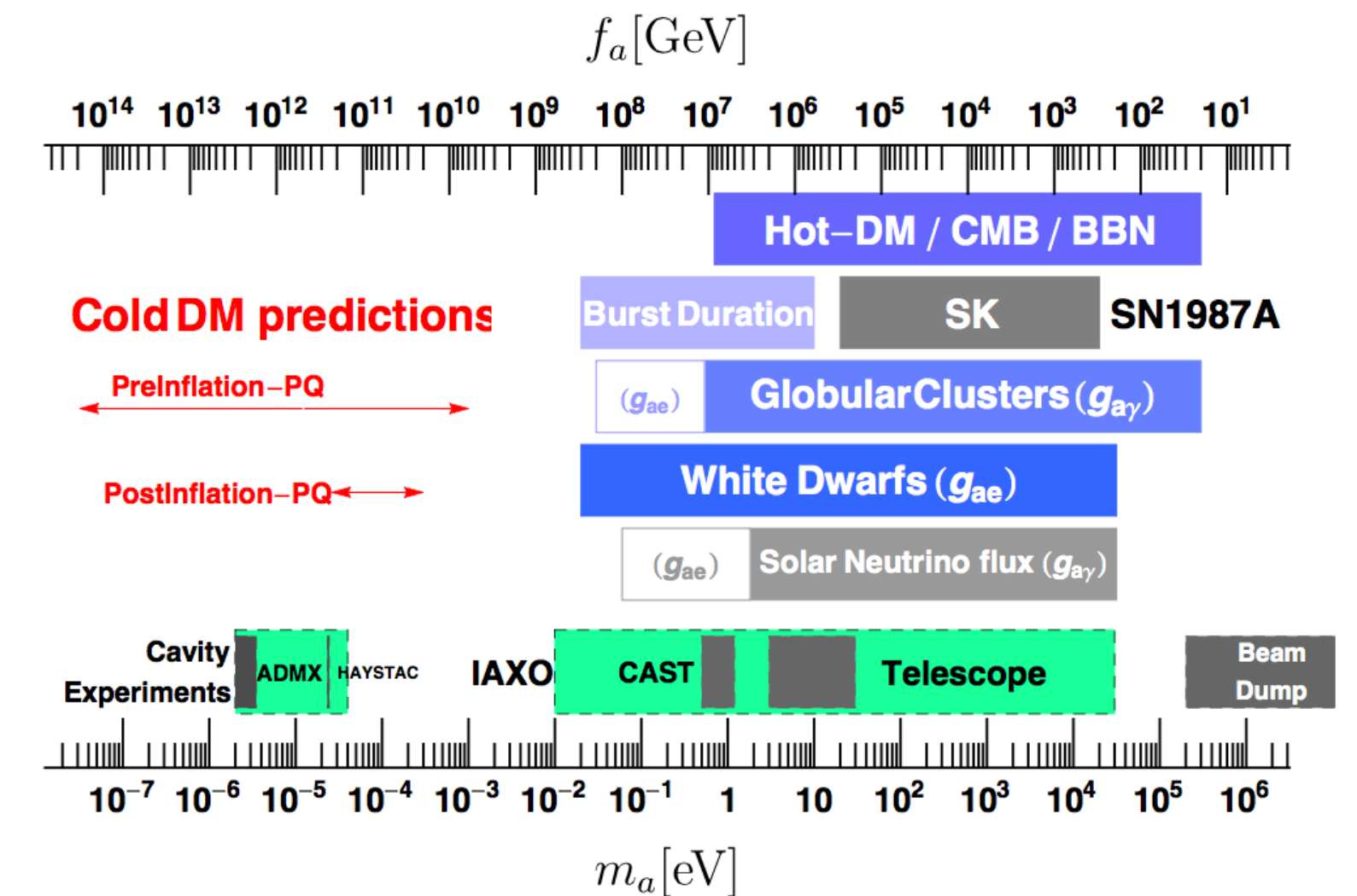
At low temperature ChPT gives:

$$m_A^2 = \frac{m_u m_d}{(m_u + m_d)^2} \frac{m_\pi^2 f_\pi^2}{f_A^2},$$

In general

$$m_A^2 f_A^2 = \chi_{top},$$

valid for any temperature.



$$m_A = 5.70(6)(4) \mu\text{eV} \left( \frac{10^{12} \text{ GeV}}{f_A} \right),$$



Time from Big Bang



Axions's freezout

$$3H(T) = m_a(T)$$

**Axions' mass  
and density  
today**

After freezout  $\frac{n_a}{s}$  constant

$$\rho_{a,0} = \frac{n_a}{s} m_a s_0$$

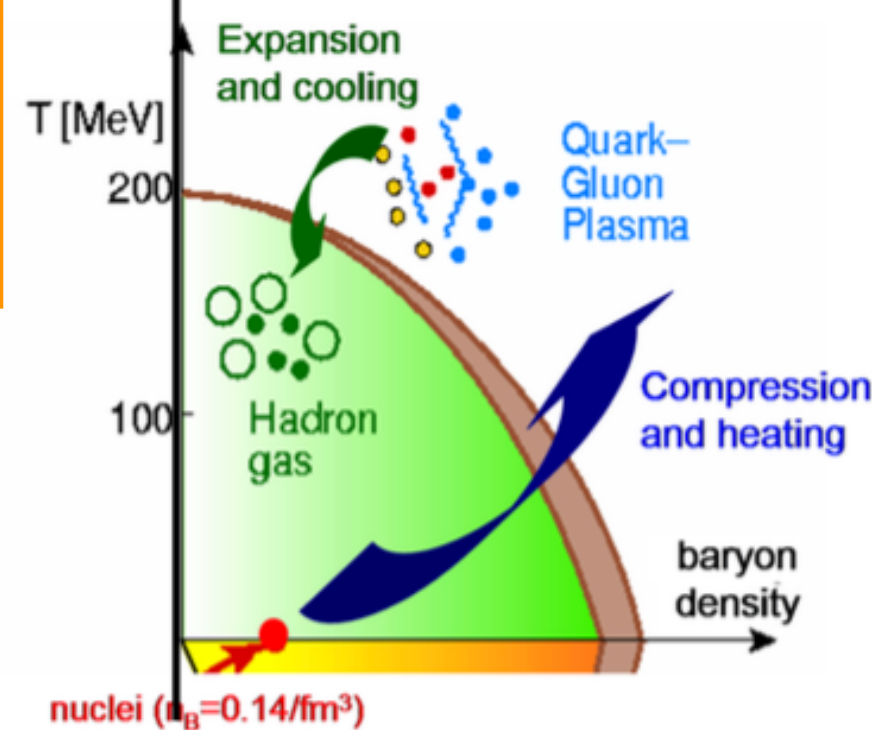
Wantz, Shellard 2010

Temperature

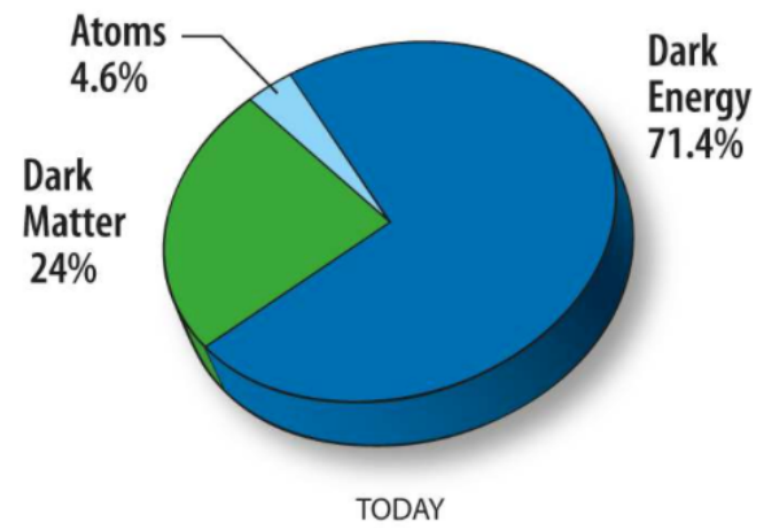
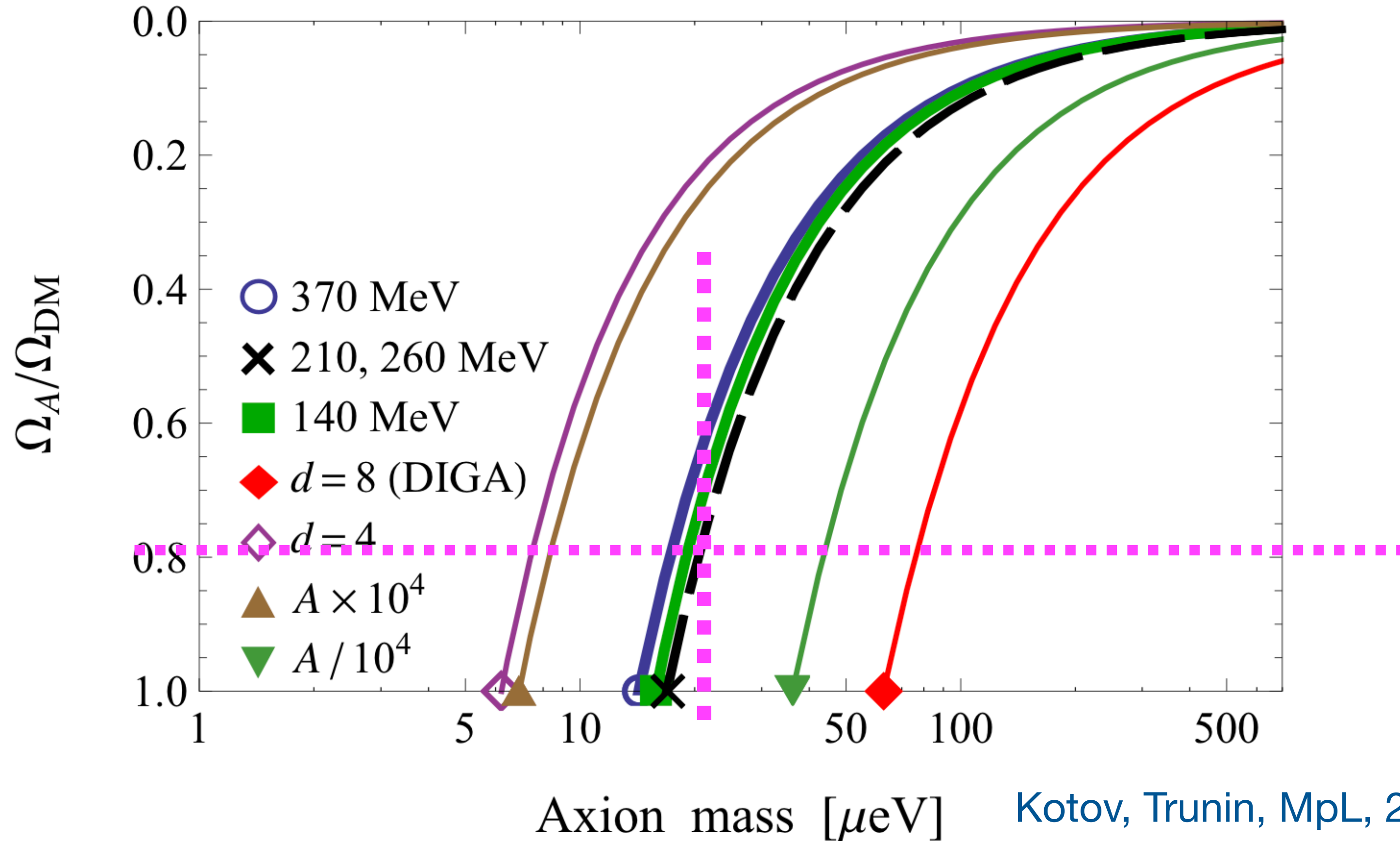
Hubble parameter  
 $H(T) \simeq T^2 / M_P$

$$m_a(T) = \sqrt{\chi(T)} / f_a$$

**Quark Gluon Plasma:  
Topology**



$$\Omega_A = F(A, d, \dots) m_A^{-\frac{3.053+d/2}{2.027+d/2}}$$



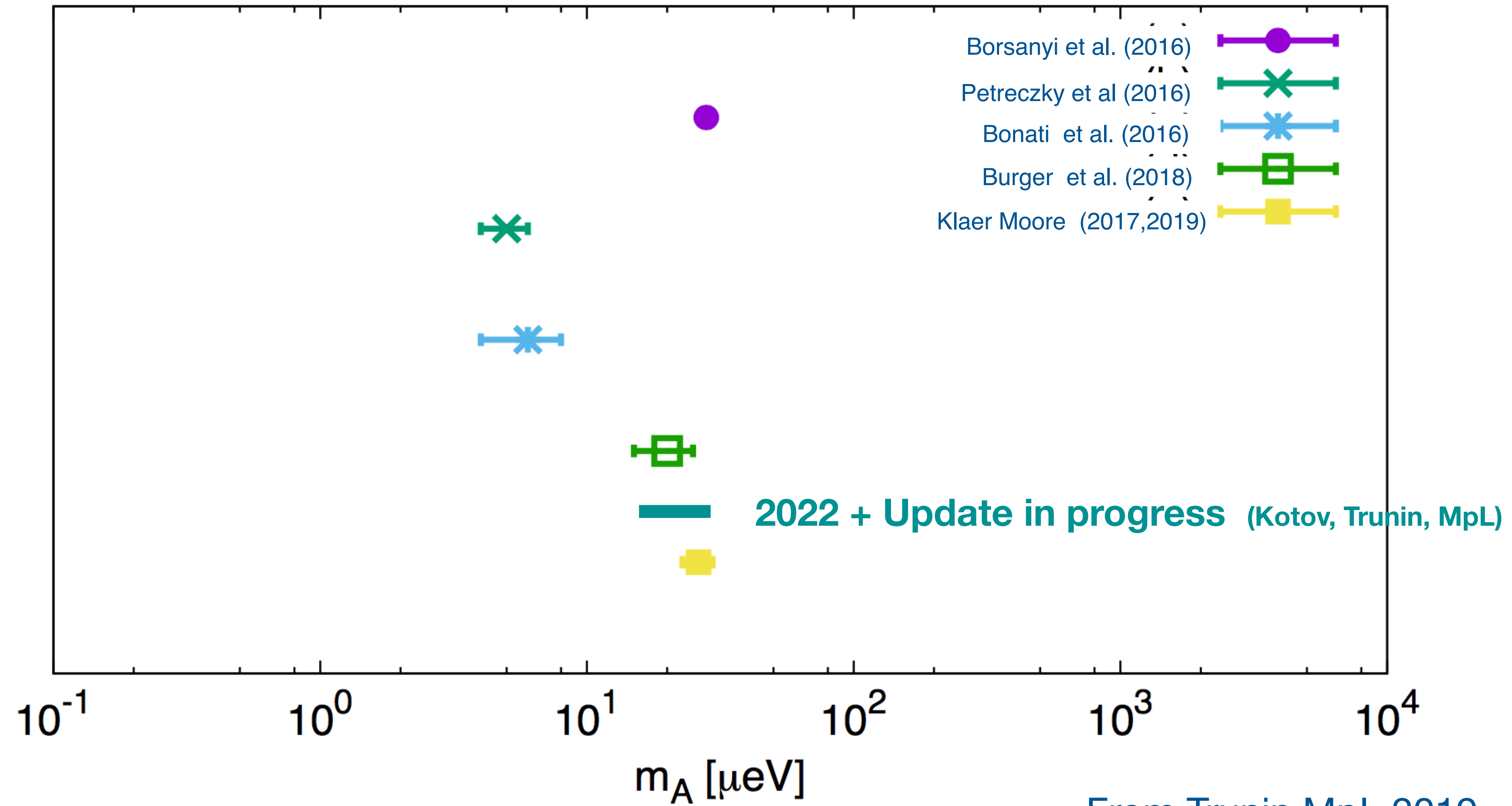
$$\Omega_a = \frac{\rho_{a,0}}{\rho_c};$$

Example: if axions constitute 80% DM,  
our results give a lower bound for the  
axion mass of  $\simeq 30 \mu\text{eV}$

Kotov, Trunin, MpL, 2022

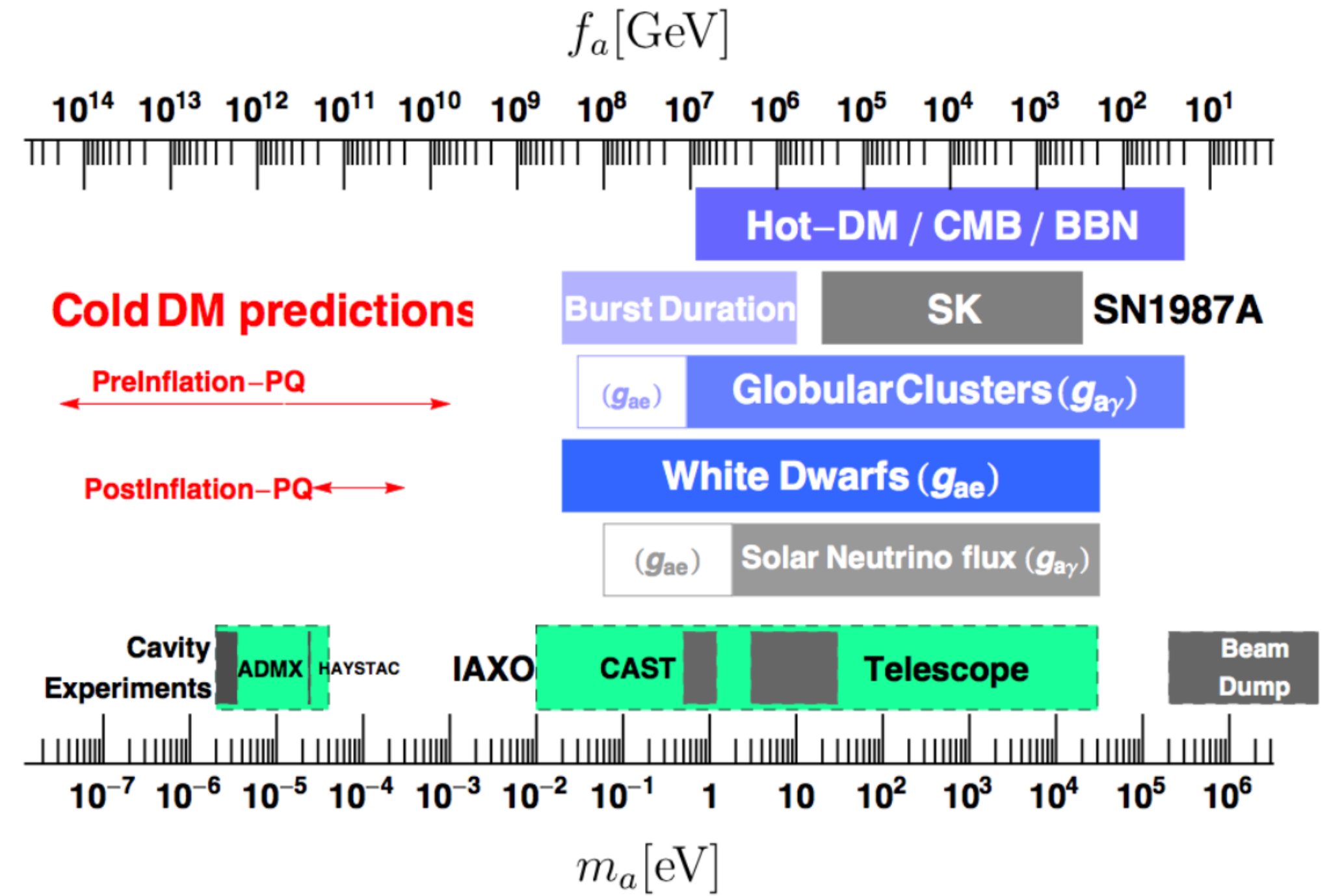
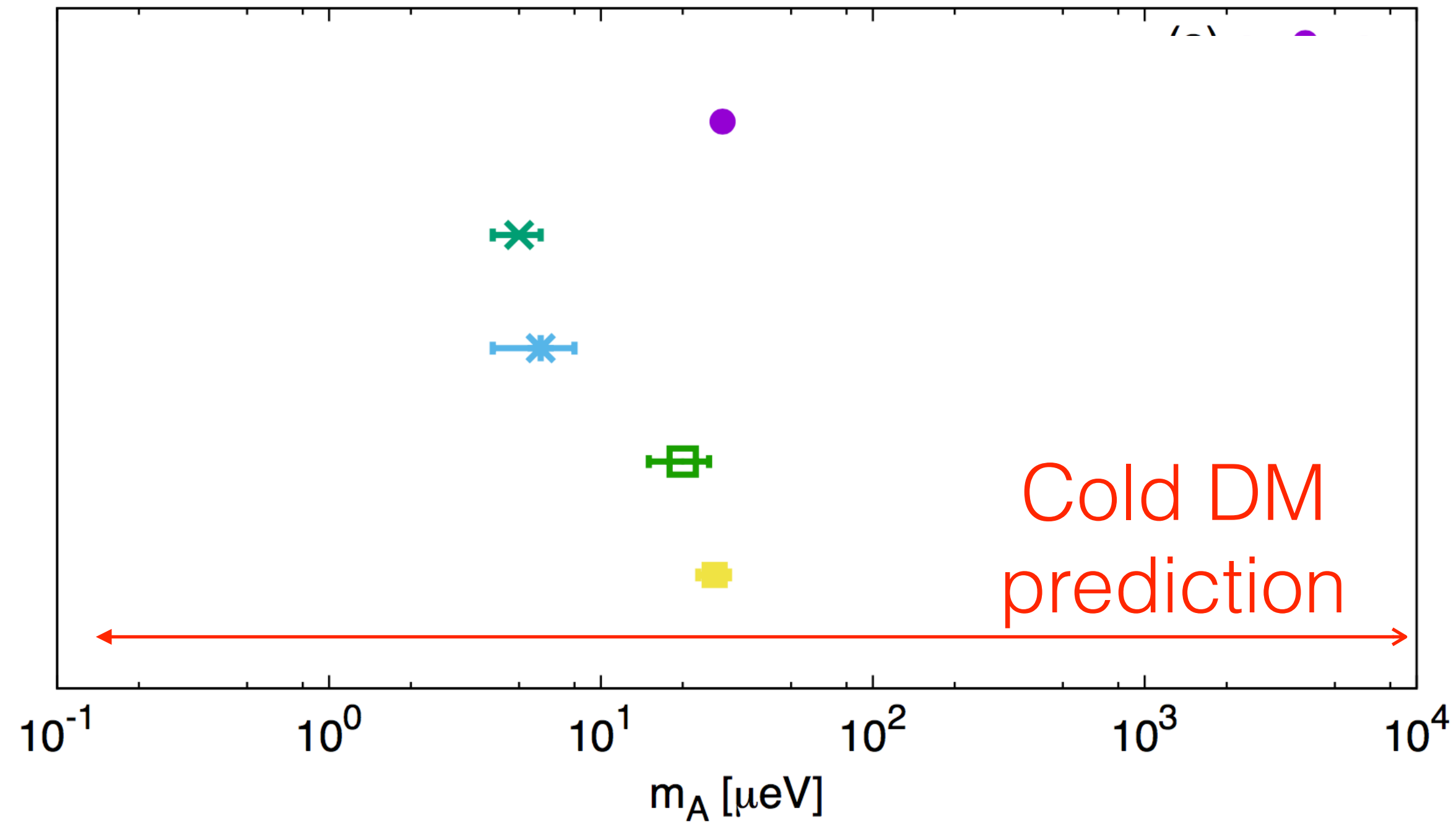
Burger, Trunin, Ilgenfritz, Mueller-Preussker, MpL 2019

# Lower limits on post-infl. axion mass from lattice QCD



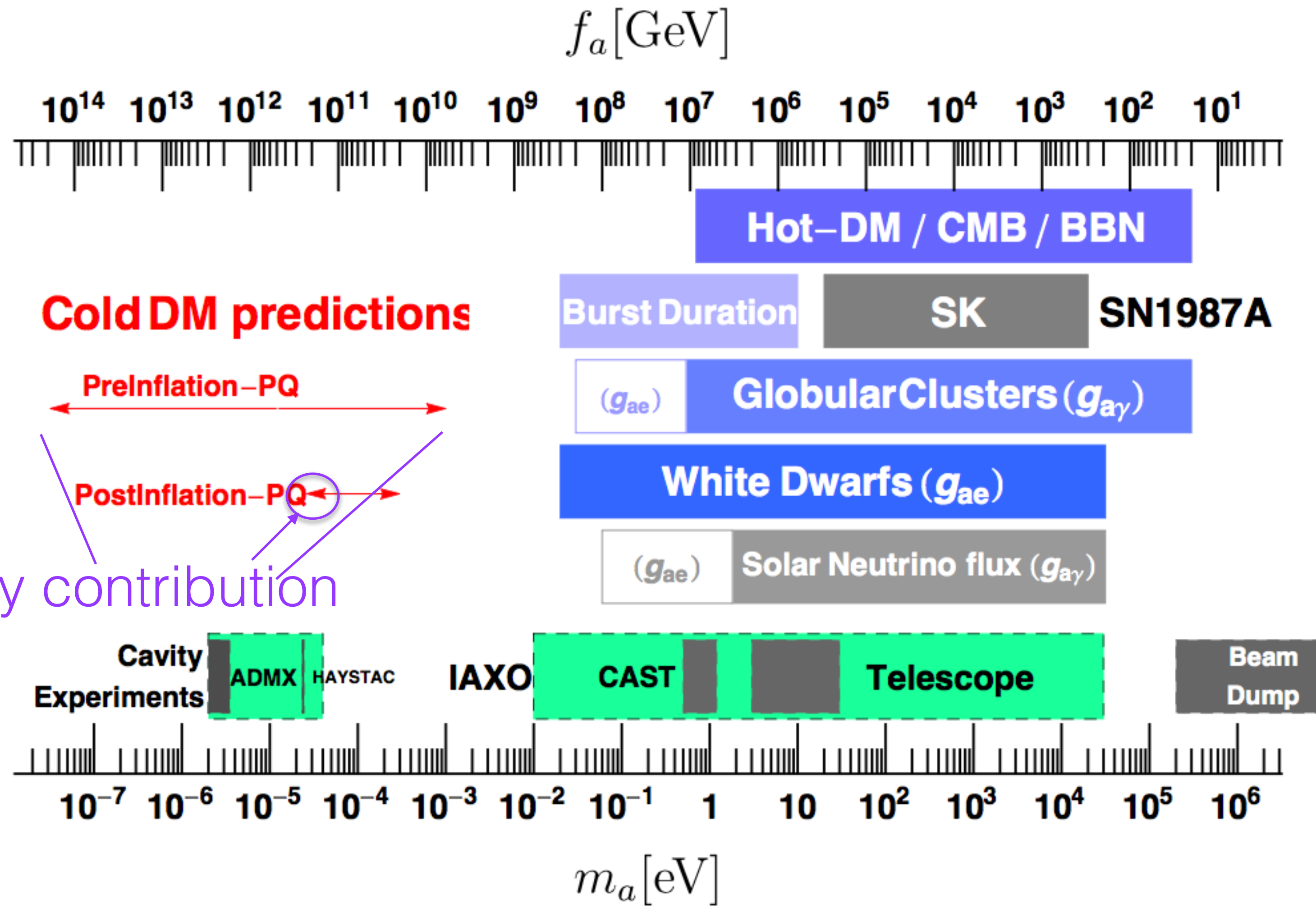
From Trunin MpL 2019

# Lower limits on post-inflationary axion mass





# Limits on the axion mass



## Summary

### Topology in QCD

1. From low to high temperatures — a threshold in the plasma?  
Common temperature for the crossover for topology and the limit of the scaling window
2. From low to high densities - experimental results and first theoretical analysis  
Experimental observation of topology suppression  
Lattice results confirm this to some extent, but with puzzling features at low temperatures in two colors QCD
3. Topology and the QCD axion  
Limits on the post-inflationary axion mass; results not entirely settled ; extrapolation to high temperature regime may be subtle and needs further studies; Axion potential?