

QCD phase diagram in strong magnetic background

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Introduction

Magnetic
field on
the lattice

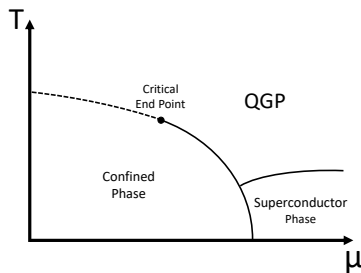
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Phase
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Summary



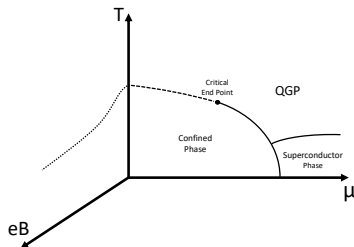
Quantum Chromodynamics (QCD) is the quantum field theory describing strong interactions. The fundamental degrees of freedom of the theory are quark and gluons, which carry color charge.

The strong interacting matter described by the theory can exist in different phases.

At ordinary temperatures and baryonic chemical potential it is in confined phase: quark and gluons cannot freely propagate.

At high temperatures and chemical potentials, the model undergoes a phase transitions to a deconfined phase called Quark Gluon Plasma (QGP).

The phase diagram in this plane is known to have a critical line and an end point.



When a magnetic field is present, it is able to interact with the matter fields which carry electric charge (i.e. quarks).

Vacuum properties are affected by the external field. Its effects are relevant even in the pure gauge sector of the theory, because of vacuum polarization effects.

The magnetic field, as well as the chemical potential, has effects on the thermodynamics of the theory. A drop in the critical temperature is the first signature of its presence.

Different predictions suggest that at extremely strong magnetic field intensities, the thermal crossover is expected to turn in a discontinuous phase transition.

Strong interacting matter can be found under such extreme conditions in the first stages of the Early Universe, in Magnetars and in heavy ion collisions.

Continuum Dirac Operator in the presence of Abelian and non-Abelian gauge fields:

$$\bar{\psi}^f(x) D_\mu^f(x) \psi^f(x) = \bar{\psi}^f(x) \left(\partial_\mu + ig G_\mu^a(x) T^a + iq^f A_\mu(x) + m^f \right) \psi^f(x)$$

The discretization can be performed through the following substitution in the usual LQCD quark action

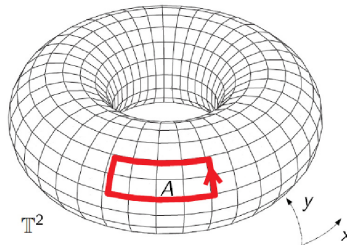
$$U_{i;\mu} \rightarrow u_{i;\mu}^f U_{i;\mu}$$

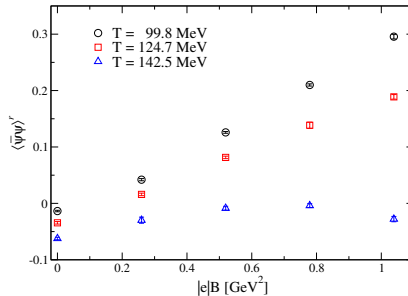
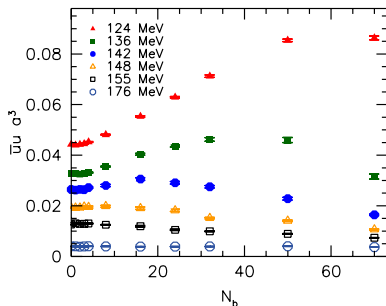
Where the Abelian phases $u_{i;\mu}^f$ are defined as follows

$$u_{i;y}^f = e^{ia^2 q_f B_z i_x}, \quad u_{i;x}^f|_{i_x=L_x} = e^{-ia^2 q_f L_x B_z i_y}$$

And are quantized in order to avoid ambiguities in the definition:

$$eB = \frac{6\pi b_z}{a^2 L_x L_y}, \quad \text{with } b_z \in \mathbb{Z} \quad \text{and} \quad b_z < \frac{L_x L_y}{6}.$$





The magnetic catalysis is known to revert when the temperature approaches the pseudo critical region in which the crossover transition to the QGP phase takes place.

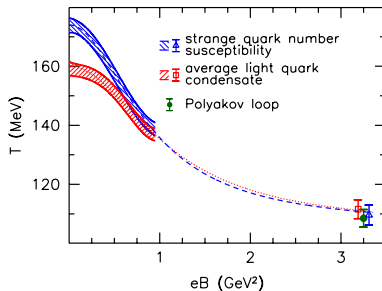
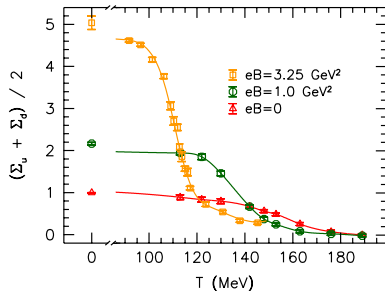
The stronger the magnetic field is, the more relevant the “*reverse chiral catalysis*” is. This effect combined with the chiral catalysis observed at low temperatures, results in a “strengthening” of the transition with the magnetic field.



G. S. Bali, F. Bruckmann, G. Endrodi, Z. Fodor, S. D. Katz, S. Krieg, A. Schafer and K. K. Szabo, JHEP **1202**, 044 (2012) [arXiv:1111.4956 [hep-lat]].



C. Bonati, M. D’Elia, M. Mariti, M. Mesiti, F. Negro, A. Rucci and F. Sanfilippo Phys. Rev. D **94**, 094007 (2016) [arXiv:1607.08160 [hep-lat]].

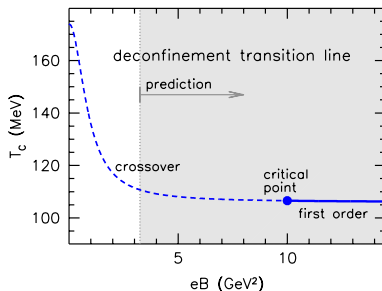


In this picture the strengthening of the transition is even clearer and another effect is highlighted: a strong background magnetic field causes a drop in the (pseudo) critical temperature.

The critical temperatures looking at different order parameters are closer, meaning that the crossover transition is slowly going to become a real and discontinuous phase transition.



G. Endrodi, JHEP **1507**, 173 (2015) [arXiv:1504.08280 [hep-lat]].



The same author provided a speculative proposal for the position of the critical point in the QCD phase diagram in the (eB, T) plane.

The existence of a first order transition in a strong field regime is deduced from a model of anisotropic pure gauge theory, which is expected to approximate the $eB \gg \Lambda_{QCD}^2$ limit of QCD.

The localization of the critical point is based on an extrapolation to zero of the width of the susceptibility peak at the transition. No direct observations were performed.



G. Endrodi, JHEP **1507**, 173 (2015) [arXiv:1504.08280 [hep-lat]].

We study a $N_f = 2 + 1$ QCD, with physical quark masses, using a tree-level improved Symanzik gauge action and stout improved staggered fermions.

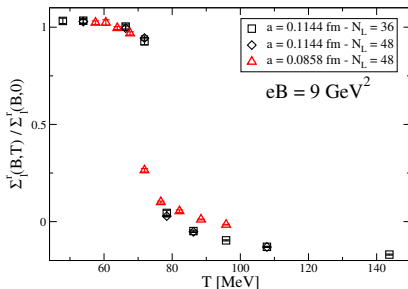
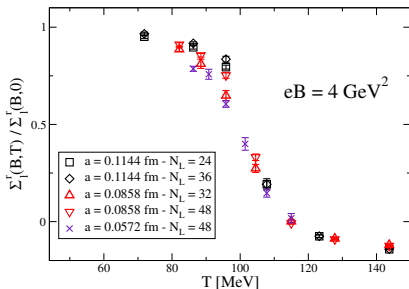
We explored theory at different lattice spacings $a = 0.114, 0.086$ and 0.057 fm. And for two different values of the magnetic field: $eB = 4$ and 9 GeV².

We performed simulations at finite temperature by imposing periodic boundary conditions on the Euclidean time direction of the lattice.

The temperature is linked to the physical extent of the time axis according to

$$T(\beta, m_l, m_s, N_t) = \frac{1}{a(\beta, m_l, m_s)N_t} \quad (\text{with } K_B = 1).$$

So, different temperatures can be probed both changing the lattice spacing a (suitably tuning the parameters) and changing the lattice sites number in the temporal axis N_t . We chose the second approach, and we refer to it as “fixed scale prescription”.



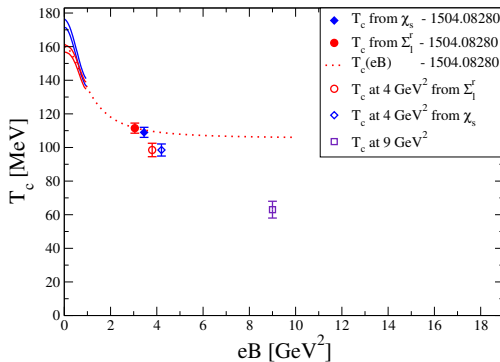
We brought up to $eB = 9 \text{ GeV}^2$ the limit of magnetic background field simulated.

The critical temperature for the transition is still dropping down to lower values as the magnetic field is increased.

The step in the chiral condensate across the transition looks steeper as the field grows, signaling that the phase transition could actually be a first order.



M. D'Elia, LM, F. Sanfilippo and A. Stanzione, Phys. Rev. D **105** (2022) no.3, 034511 [arXiv:2111.11237 [hep-lat]].

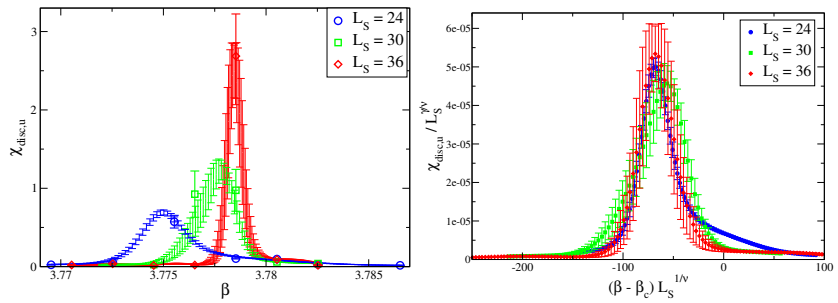


The drop in T_c that we observe is much higher than the predicted one.

An extrapolation on these data to $T_c = 0$ would return the deconfining magnetic field intensity $eB \simeq 18 \text{ GeV}^2$.

There is no a $\sim 4 \text{ GeV}$ energy scale in QCD which would explain such a value for a deconfining magnetic field.





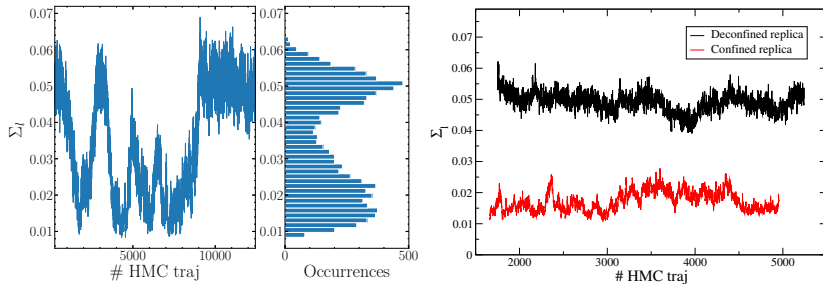
Because of the fixed scale prescription, a study of the nature of the transition at $eB = 9 \text{ GeV}^2$ on the same constant physics line was not affordable with the available computational resources.

Since the presence of a first order transition is stable under small variations of the parameters, we performed a fine tuning in order to cross the critical line.

Using as starting point a simulation performed at a temperature and magnetic field close to the transition, we slightly changed β until the transition occurred.

Using this prescription, a finite size scaling was now affordable without lack in reliability, and it clearly shows a first order transition.





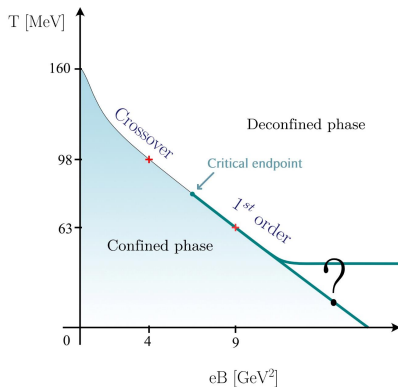
Other smoking guns of a first order transition are found on these samples.

A simulation running on the pseudo critical value of β shows indeed a bi-stable history of the chiral condensate.

In a larger volume, the tunneling probability becomes small and simulations with the same parameters, but different starting points, show parallel histories.



M. D'Elia, LM, F. Sanfilippo and A. Stanzione, Phys. Rev. D **105** (2022) no.3, 034511 [arXiv:2111.11237 [hep-lat]].



We propose a new hint for the phase diagram of the QCD in the (eB, T) plane, based upon our results.

Since the nature of the transition appears to be a crossover at $eB = 4 \text{ GeV}^2$ and a first order transition at $eB = 9 \text{ GeV}^2$, the critical endpoint should be located somewhere in that interval.



Summary

- The drop in T_c for $|e|B = 9 \text{ GeV}^2$ is larger than expected;
- We found an analytical crossover transition at $eB = 4 \text{ GeV}^2$, and a first order transition at $eB = 9 \text{ GeV}^2$;
- A critical endpoint is expected to lie between this two magnetic fields at a temperature in the range $T_{CEP} \in (63, 98) \text{ MeV}$.

Outlooks

- Increase precision on the phase point location;
- Characterize the two phases, studying other relevant quantities (i.e., latent heat, pressure, transport properties, and so on);
- Including further external conditions, in order to study the phase diagram in a more general and multidimensional framework.

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Thanks for your attention