Thermalization with a multibath: an investigation in simple models

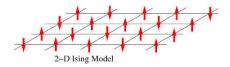
> Giovanni Battista Carollo Università degli Studi di Bari

> > December 20th, 2022

## **SM&FT**2022

Preprint: 2212.00527 (with F. Corberi and G. Gonnella)

## Spins on a lattice, out-of-equilibrium

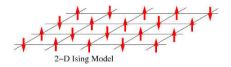


Ising Hamiltonian

$$\mathcal{H} = -\sum_{ij} J_{ij} S_i S_j , \quad S_i = \pm 1;$$

 $S_i$  in contact with a thermal bath

## Spins on a lattice, out-of-equilibrium



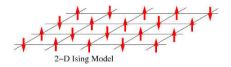
Ising Hamiltonian

$$\mathcal{H} = -\sum_{ij} J_{ij} S_i S_j , \quad S_i = \pm 1;$$

 $S_i$  in contact with a thermal bath

 Add a dynamics (ex. Glauber, 1963) to study the evolution towards equilibrium (if present)

## Spins on a lattice, out-of-equilibrium



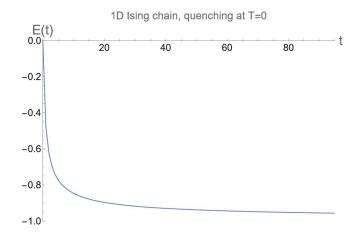
Ising Hamiltonian

$$\mathcal{H} = -\sum_{ij} J_{ij} S_i S_j , \quad S_i = \pm 1;$$

 $S_i$  in contact with a thermal bath

- Add a dynamics (ex. Glauber, 1963) to study the evolution towards equilibrium (if present)
- for  $J_{ij}$  without fixed signs  $\rightarrow$  long relaxation time

Aging



# Different groups of variables attached to different thermal baths?

- Different groups of variables attached to different thermal baths?
- Idea introduced in the context of Langevin dynamics (Cugliandolo, Kurchan & Peliti, 1996, Cugliandolo & Kurchan, 1998, 1999)

- Different groups of variables attached to different thermal baths?
- Idea introduced in the context of Langevin dynamics (Cugliandolo, Kurchan & Peliti, 1996, Cugliandolo & Kurchan, 1998, 1999)
- Naturally out of equilibrium, but rather simple models

- Different groups of variables attached to different thermal baths?
- Idea introduced in the context of Langevin dynamics (Cugliandolo, Kurchan & Peliti, 1996, Cugliandolo & Kurchan, 1998, 1999)
- Naturally out of equilibrium, but rather simple models
- Lack of literature for lattice models (Piscitelli et al., 2008, 2009, Borchers et al., 2012, Contucci et al., 2020)

#### Contucci et al. 2019, 2021

► Hamiltonian  $\mathcal{H}(\mathbf{x}_1, \ldots, \mathbf{x}_r)$ , each group of variables in contact with a different thermal bath, temperatures  $T_a = \beta_a^{-1}$ , relaxation times  $\tau_r \ll \tau_{r-1} \ll \ldots \ll \tau_1$ ,  $\zeta_a = \frac{\beta_a}{\beta_r}$ 

#### Contucci et al. 2019, 2021

- ► Hamiltonian  $\mathcal{H}(\mathbf{x}_1, \ldots, \mathbf{x}_r)$ , each group of variables in contact with a different thermal bath, temperatures  $T_a = \beta_a^{-1}$ , relaxation times  $\tau_r \ll \tau_{r-1} \ll \ldots \ll \tau_1$ ,  $\zeta_a = \frac{\beta_a}{\beta_r}$
- Nested partition function:

$$Z = \left\{ \int d\mathbf{x}_1 \left[ \int d\mathbf{x}_2 \dots \left[ \int d\mathbf{x}_r e^{-\beta_r H(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3, \dots, \mathbf{x}_r)} \right]^{\zeta_{r-1}/\zeta_r} \dots \right]^{\zeta_1/\zeta_2} \dots \right\}^{1/\zeta_1}$$

## Multibath: relation with disordered systems theory

▶ Bibath:  $\vec{\sigma}$  in contact at  $T_1$ ,  $\vec{S}$  in contact at  $T_2$ ,  $\tau_2 \ll \tau_1$ ,  $\zeta_1 = \frac{\beta_1}{\beta_2}$ ,  $\zeta_2 = \frac{\beta_2}{\beta_2} = 1$ 

$$F = -\frac{1}{\beta_2} \log Z = -\frac{1}{\beta_2 \zeta_1} \log \left[ \mathbb{E}_{\sigma} \left[ \operatorname{tr}_S e^{-\beta_2 H} \right]^{\zeta_1} \right]$$

## Multibath: relation with disordered systems theory

Bibath: \$\vec{\sigma}\$ in contact at \$T\_1\$, \$\vec{S}\$ in contact at \$T\_2\$, \$\tau\_2\$ < \$\vec{\approx}\$\_1\$, \$\zeta\_1\$ = \$\frac{\beta\_1}{\beta\_2}\$, \$\zeta\_2\$ = \$\vec{\beta\_2}{\beta\_2}\$ = \$1\$</li>
\$F = \$-\frac{1}{\beta\_2}\$ log \$Z = \$-\frac{1}{\beta\_2\zeta\_1}\$ log \$\left[ \mathbb{E}\_{\sigma}\$ \left[ \text{tr}\_S e^{-\beta\_2 H} \right]^{\zeta\_1}\$ \right]\$
\$if \$T\_1 = \$T\_2\$: \$\zeta\_1\$ = \$1\$, \$F = \$-\frac{1}{\beta\_2}\$ log \$\left[ \mathbb{E}\_{\sigma}\$ [\text{tr}\_S e^{-\beta\_2 H} \right]^{\zeta\_1}\$]\$ (annealed)

## Multibath: relation with disordered systems theory

**b** Bibath:  $\vec{\sigma}$  in contact at  $T_1$ ,  $\vec{S}$  in contact at  $T_2$ ,  $\tau_2 \ll \tau_1$ ,  $\zeta_1 = \frac{\beta_1}{\beta_2}$ ,  $\zeta_2 = \frac{\beta_2}{\beta_2} = 1$   $F = -\frac{1}{\beta_2} \log Z = -\frac{1}{\beta_2 \zeta_1} \log \left[ \mathbb{E}_{\sigma} \left[ \operatorname{tr}_S e^{-\beta_2 H} \right]^{\zeta_1} \right]$  **b** if  $T_1 = T_2$ :  $\zeta_1 = 1$ ,  $F = -\frac{1}{\beta_2} \log \left[ \mathbb{E}_{\sigma} \left[ \operatorname{tr}_S e^{-\beta_2 H} \right] \right]$  (annealed) **b** if  $T_1 \to \infty$ :  $\zeta_1 \to 0$ ,  $F = -\lim_{\zeta_1 \to 0} \frac{1}{\beta_2 \zeta_1} \log \left[ \mathbb{E}_{\sigma} \left[ \operatorname{tr}_S e^{-\beta_2 H} \right]^{\zeta_1} \right]$  (quenched)

• Autocorrelation function:  $C(t, t') = \langle S_i(t)S_i(t') \rangle$  (t > t')

- ▶ Autocorrelation function:  $C(t,t') = \langle S_i(t)S_i(t') \rangle$  (t > t')
- Integrated response function (h linear perturbation of  $\mathcal{H}$ ):

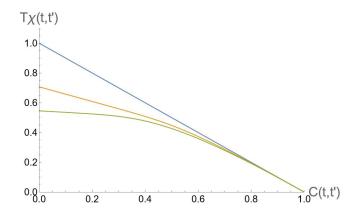
$$\chi(t,t') = \int_{t'}^t R(t,t'')dt'' \quad \text{where} \quad R(t,t') = \frac{\delta \langle S_i \rangle}{\delta h_i} \bigg|_{h=0}$$

- ▶ Autocorrelation function:  $C(t,t') = \langle S_i(t)S_i(t') \rangle$  (t > t')
- Integrated response function (h linear perturbation of  $\mathcal{H}$ ):

$$\chi(t,t') = \int_{t'}^t R(t,t'') dt'' \quad \text{where} \quad R(t,t') = \left. \frac{\delta \langle S_i \rangle}{\delta h_i} \right|_{h=0}$$

Fluctuation-Dissipation Theorem: at equilibrium  $T\chi(t,t') = 1 - C(t,t')$ 

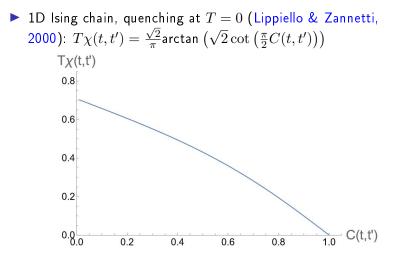
## Effective temperature



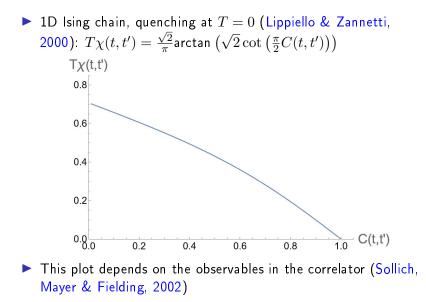
Effective temperature (Cugliandolo, Kurchan & Parisi, 1994):

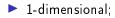
$$T_{\rm eff} = -\left(\frac{d(T\chi)}{dC}\right)^{-1}$$

## Problematic?



## Problematic?





• Binary Variables  $(\pm 1)$ :

• Binary Variables  $(\pm 1)$ :

•  $S_i$  coupled to a bath with reciprocal temperature B;

- Binary Variables  $(\pm 1)$ :
  - $S_i$  coupled to a bath with reciprocal temperature B;
  - $\sigma_i$  coupled to a bath with reciprocal temperature  $\beta$ .

- Binary Variables  $(\pm 1)$ :
  - $S_i$  coupled to a bath with reciprocal temperature B;
  - $\sigma_i$  coupled to a bath with reciprocal temperature  $\beta$ .
- $S_i$  colder  $(B > \beta)$ .

- Binary Variables  $(\pm 1)$ :
  - $S_i$  coupled to a bath with reciprocal temperature B;
  - $\sigma_i$  coupled to a bath with reciprocal temperature  $\beta$ .
- $S_i$  colder  $(B > \beta)$ .
- Dynamical evolution:

- Binary Variables  $(\pm 1)$ :
  - $S_i$  coupled to a bath with reciprocal temperature B;
  - $\sigma_i$  coupled to a bath with reciprocal temperature  $\beta$ .
- $S_i$  colder  $(B > \beta)$ .
- Dynamical evolution:
  - Glauber dynamics (simple);

- ► Binary Variables (±1):
  - $\triangleright$   $S_i$  coupled to a bath with reciprocal temperature B;
  - $\sigma_i$  coupled to a bath with reciprocal temperature  $\beta$ .
- $S_i$  colder  $(B > \beta)$ .
- Dynamical evolution:
  - Glauber dynamics (simple);
  - $\sigma_i$  evolve slower, with relative timescale  $\tau > 1$ .

Basic models (with PBC):

1.  $\mathcal{H} = -\sum_{i} S_i \sigma_i$  (paramagnet); advantage: fully analitically solvable; disadvantage: not interacting (trivial); Basic models (with PBC):

1.  $\mathcal{H} = -\sum_{i} S_i \sigma_i$  (paramagnet); advantage: fully analitically solvable; disadvantage: not interacting (trivial);

2. 
$$\mathcal{H} = -\sum_{i} S_i \sigma_i S_{i+1}$$
 (ferromagnet);

advantage: interacting and highly non trivial phenomenology; disadvantage: not analytically solvable;

Basic models (with PBC):

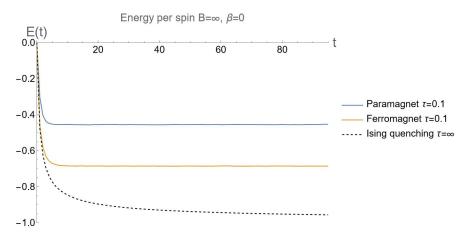
1.  $\mathcal{H} = -\sum_{i} S_i \sigma_i$  (paramagnet); advantage: fully analitically solvable; disadvantage: not interacting (trivial);

2. 
$$\mathcal{H} = -\sum_{i} S_i \sigma_i S_{i+1}$$
 (ferromagnet);

advantage: interacting and highly non trivial phenomenology; disadvantage: not analytically solvable;

3.  $\mathcal{H} = -\sum_i (S_i \sigma_i + \sigma_i S_{i+1})$  (alternated).

## First result: stationarization

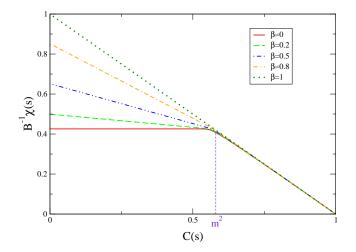


#### Paramagnet:

$$E(t) = -\frac{m + \frac{\mu}{\tau}}{1 + \frac{1}{\tau}} \left( 1 - e^{-t\left(1 + \frac{1}{\tau}\right)} \right) \ , \quad m = \tanh B \ , \quad \mu = \tanh \beta$$
12/15

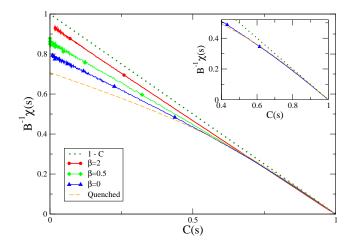
## Second result: effective temperatures

Paramagnet (analytical curves,  $\tau = 100, B = 1$ ):



## Second result: effective temperatures

Ferromagnet (numerical simulation,  $\tau = 20$ ,  $B = \infty$ ):



1. In lattice models one can compute averages by coupling the variables to a multibath and studying the dynamics.

- 1. In lattice models one can compute averages by coupling the variables to a multibath and studying the dynamics.
- 2. We considered the simplest models in this contest and we found that:

- 1. In lattice models one can compute averages by coupling the variables to a multibath and studying the dynamics.
- 2. We considered the simplest models in this contest and we found that:
  - 2.1 the multibath leads the system to a nonequilibrium stationary state, destroying aging;

- 1. In lattice models one can compute averages by coupling the variables to a multibath and studying the dynamics.
- 2. We considered the simplest models in this contest and we found that:
  - 2.1 the multibath leads the system to a nonequilibrium stationary state, destroying aging;
  - 2.2 effective temperatures:

- 1. In lattice models one can compute averages by coupling the variables to a multibath and studying the dynamics.
- 2. We considered the simplest models in this contest and we found that:
  - 2.1 the multibath leads the system to a nonequilibrium stationary state, destroying aging;
  - 2.2 effective temperatures:
    - 2.2.1 paramagnet: standard picture;

- 1. In lattice models one can compute averages by coupling the variables to a multibath and studying the dynamics.
- 2. We considered the simplest models in this contest and we found that:
  - 2.1 the multibath leads the system to a nonequilibrium stationary state, destroying aging;
  - 2.2 effective temperatures:
    - 2.2.1 paramagnet: standard picture;
    - 2.2.2 ferromagnet: complicated picture.