# Thermalization with a multibath: an investigation in simple models 

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December 20th, 2022

## SM\&FT2022

Preprint: 2212.00527 (with F. Corberi and G. Gonnella)

## Spins on a lattice, out-of-equilibrium



- Ising Hamiltonian

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\mathcal{H}=-\sum_{i j} J_{i j} S_{i} S_{j}, \quad S_{i}= \pm 1
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$S_{i}$ in contact with a thermal bath

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- Add a dynamics (ex. Glauber, 1963) to study the evolution towards equilibrium (if present)
- for $J_{i j}$ without fixed signs $\rightarrow$ long relaxation time


## Aging

1D Ising chain, quenching at $T=0$


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- Idea introduced in the context of Langevin dynamics (Cugliandolo, Kurchan \& Peliti, 1996, Cugliandolo \& Kurchan, 1998, 1999)
- Naturally out of equilibrium, but rather simple models
- Lack of literature for lattice models (Piscitelli et al., 2008, 2009, Borchers et al., 2012, Contucci et al., 2020)


## Nested partition function

Contucci et al. 2019, 2021

- Hamiltonian $\mathcal{H}\left(\mathbf{x}_{1}, \ldots, \mathbf{x}_{r}\right)$, each group of variables in contact with a different thermal bath, temperatures $T_{a}=\beta_{a}^{-1}$, relaxation times $\tau_{r} \ll \tau_{r-1} \ll \ldots \ll \tau_{1}, \zeta_{a}=\frac{\beta_{a}}{\beta_{r}}$


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- Nested partition function:

$$
Z=\left\{\int d \mathbf{x}_{1}\left[\int d \mathbf{x}_{2} \ldots\left[\int d \mathbf{x}_{r} e^{-\beta_{r} H\left(\mathbf{x}_{1}, \mathbf{x}_{2}, \mathbf{x}_{3}, \ldots, \mathbf{x}_{r}\right)}\right]^{\zeta_{r-1} / \zeta_{r}} \ldots\right]^{\zeta_{1} / \zeta_{2}} \ldots\right\}^{1 / \zeta_{1}}
$$

## Multibath: relation with disordered systems theory

- Bibath: $\vec{\sigma}$ in contact at $T_{1}, \vec{S}$ in contact at $T_{2}, \tau_{2} \ll \tau_{1}$, $\zeta_{1}=\frac{\beta_{1}}{\beta_{2}}, \zeta_{2}=\frac{\beta_{2}}{\beta_{2}}=1$

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F=-\frac{1}{\beta_{2}} \log Z=-\frac{1}{\beta_{2} \zeta_{1}} \log \left[\mathbb{E}_{\sigma}\left[\operatorname{tr}_{S} e^{-\beta_{2} H}\right]^{\zeta_{1}}\right]
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- if $T_{1} \rightarrow \infty: \zeta_{1} \rightarrow 0$,

$$
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- Fluctuation-Dissipation Theorem: at equilibrium

$$
T \chi\left(t, t^{\prime}\right)=1-C\left(t, t^{\prime}\right)
$$

## Effective temperature



Effective temperature (Cugliandolo, Kurchan \& Parisi, 1994):

$$
T_{\mathrm{eff}}=-\left(\frac{d(T \chi)}{d C}\right)^{-1}
$$

## Problematic?

- 1D Ising chain, quenching at $T=0$ (Lippiello \& Zannetti, 2000): $T \chi\left(t, t^{\prime}\right)=\frac{\sqrt{2}}{\pi} \arctan \left(\sqrt{2} \cot \left(\frac{\pi}{2} C\left(t, t^{\prime}\right)\right)\right)$



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- This plot depends on the observables in the correlator (Sollich, Mayer \& Fielding, 2002)


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- $\sigma_{i}$ evolve slower, with relative timescale $\tau>1$.


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advantage: interacting and highly non trivial phenomenology; disadvantage: not analytically solvable;
3. $\mathcal{H}=-\sum_{i}\left(S_{i} \sigma_{i}+\sigma_{i} S_{i+1}\right)$ (alternated).

## First result: stationarization

Energy per spin $B=\infty, \beta=0$


## Paramagnet:

$$
E(t)=-\frac{m+\frac{\mu}{\tau}}{1+\frac{1}{\tau}}\left(1-e^{-t\left(1+\frac{1}{\tau}\right)}\right), \quad m=\tanh B, \quad \mu=\tanh \beta
$$

## Second result: effective temperatures

Paramagnet (analytical curves, $\tau=100, B=1$ ):


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Ferromagnet (numerical simulation, $\tau=20, B=\infty$ ):


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2.2 effective temperatures:
2.2.1 paramagnet: standard picture;
2.2.2 ferromagnet: complicated picture.
