## The light baryon resonance spectrum in a coupled-channel approach

## Recent results from the Jülich-Bonn model - NSTAR 2022

October 18, 2022 | Deborah Rönchen | Institute for Advanced Simulation, Forschungszentrum Jülich

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Supported by DFG, NSFC and MKW NRW
HPC support by Jülich Supercomputing Centre

## The excited baryon spectrum:

## Connection between experiment and QCD in the non-perturbative regime

Experimental study of hadronic reactions

source: ELSA; data: ELSA, JLab, MAMI

Theoretical predictions of excited hadrons e.g. from relativistic quark models:


Löring et al. EPJ A 10, 395 (2001), experimental spectrum: PDG 2000

Major source of information:
In recent years: photoproduction reactions

- enlarged data base with high quality (double) polarization observables, towards a complete experiment Reviews: Prog.Part.Nucl.Phys. 125, 103949 (2022), Prog.Part.Nucl.Phys. 111 (2020) 103752

In the future: electroproduction reactions

- $10^{5}$ data points for $\pi N, \eta N, K Y, \pi \pi N$ already available
- access the $Q^{2}$ dependence of the amplitude Reviews: Prog.Part.Nucl.Phys. 67 (2012)


## The excited baryon spectrum:

## Connection between experiment and QCD in the non-perturbative regime

Experimental study of hadronic reactions

source: ELSA; data: ELSA, JLab, MAMI
$\Rightarrow$ Partial wave decomposition: decompose data with respect to a conserved quantum number:
total angular momentum and parity $J^{P}$

Theoretical predictions of excited hadrons e.g. from relativistic quark models:


Löring et al. EPJ A 10, 395 (2001), experimental spectrum: PDG 2000
$\Rightarrow$ search for resonances/excited states in those partial waves: poles on the $2^{\text {nd }}$ Riemann sheet
(Breit-Wigner problematic in baryon spectroscopy)

## The Jülich-Bonn DCC approach for $N^{*}$ and $\Delta^{*}$

 pion-induced reactionsDynamical coupled-channels (DCC): simultaneous analysis of different reactions

## The scattering equation in partial-wave basis

$$
\begin{aligned}
&\left\langle L^{\prime} S^{\prime} p^{\prime}\right| T_{\mu \nu}^{\prime \prime}|L S p\rangle=\left\langle L^{\prime} S^{\prime} p^{\prime}\right| V_{\mu \nu}^{\prime \prime}|L S p\rangle+ \\
& \sum_{\gamma, L^{\prime \prime} S^{\prime \prime}} \int_{0}^{\infty} d q q^{2} \quad\left\langle L^{\prime} S^{\prime} p^{\prime}\right| V_{\mu \gamma}^{\prime \prime}\left|L^{\prime \prime} S^{\prime \prime} q\right\rangle \frac{1}{E-E_{\gamma}(q)+i \epsilon}\left\langle L^{\prime \prime} S^{\prime \prime} q\right| T_{\gamma \nu}^{\prime \prime}|L S p\rangle
\end{aligned}
$$

- channels $\nu, \mu, \gamma$ :



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\end{aligned}
$$



- potentials $V$ constructed from effective $\mathcal{L}$
- s-channel diagrams: $T^{P}$
genuine resonance states
- $t$ - and $u$-channel: $T^{N P}$ dynamical generation of poles partial waves strongly correlated
- contact terms


## Thresholds of inelastic channels

- (2 body) unitarity and analyticity respected (no on-shell factorization, dispersive parts included)
- opening of inelastic channels $\Rightarrow$ branch point and new Riemann sheet


## 3-body $\pi \pi N$ channel:

- parameterized effectively as $\pi \Delta, \sigma N, \rho N$
- $\pi N / \pi \pi$ subsystems fit the respective phase shifts
$\square$ branch points move into complex plane



Example: $\rho N$ branch point at

$$
M_{N}+m_{\text {rho }}=1700 \pm i 75 \mathrm{MeV}
$$

Inclusion of branch points important to avoid false resonance signal!

## Photoproduction in a semi-phenomenological approach

## Multipole amplitude

$$
M_{\mu \gamma}^{\prime \prime}=V_{\mu \gamma}^{\prime \prime}+\sum_{\kappa} T_{\mu \kappa}^{\prime \prime} G_{\kappa} V_{\kappa \gamma}^{\prime \prime}
$$

(partial wave basis)

$T_{\mu \kappa}$ : full hadronic $T$-matrix as in pion-induced reactions
Photoproduction potential: approximated by energy-dependent polynomials (field-theoretical description numerically too expensive )


$$
=\frac{\tilde{\gamma}_{\mu}^{a}(q)}{m_{N}} P_{\mu}^{\mathrm{NP}}(E)+\sum_{i} \frac{\gamma_{\mu ; i}^{a}(q) P_{i}^{P}(E)}{E-m_{i}^{b}}
$$

## Simultaneous fit of pion- \& photon-induced reactions

## Free parameters

$\pi N \rightarrow \pi N, \eta N, K Y:$ s-channel: resonances $\left(T^{P}\right)$


- $\gamma p \rightarrow \pi N, \eta N, K Y$ : couplings of the polynomials and $s$-channel parameters

- couplings in contact terms: one per PW, couplings to $\pi N, \eta N$, ( $\pi \Delta$,) $K \Lambda, K \Sigma$
- $t$ - \& $u$-channel parameters: cut-offs, mostly fixed to values of previous JüBo studies (couplings fixed from $\mathrm{SU}(3)$ )
$\Rightarrow \quad>900$ fit parameters in total, $\sim 72,000$ data points
$\bigsqcup$ calculations on a supercomputer [JURECA, Julich Supercomputing Centre, Journal of large-scale research facilities, 2, A62 (2016)]
- large number of fit parameters, many from polynomials
- can be regarded as advantage: prevents the inclusion of superfluous $s$-channel states to improve fit


## Extension to $K \Sigma$ photoproduction on the proton

JüBo2022 arXiv:2208.00089 [nucl-th], accepted at EPJ A

Unique opportunities in $\gamma p \rightarrow K^{+} \Sigma^{0}, K^{0} \Sigma^{+}$:

- coupling of $N^{*}$ 's, $\Delta^{*}$ 's to strangeness channels: missing resonances not seen in $\pi N$ scattering?
- mixed isospin $\rightarrow$ more information on $\Delta$ states
- self-analyzing decay of Y's: recoil polarization from angular distribution of decay products
(important for "complete experiment")
- better data quality than in $\pi N \rightarrow K Y$

Selected fit results


## Extension to $K \Sigma$ photoproduction on the proton

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Simultaneous analysis of $\pi N \rightarrow \pi N, \eta N, K \Lambda, K \Sigma$ and $\gamma p \rightarrow \pi N, \eta N, K \Lambda, K \Sigma$

- almost 72,000 data points in total, $W_{\max }=2.4 \mathrm{GeV}$

$$
\begin{aligned}
& \gamma p \rightarrow K^{+} \Sigma^{0}: d \sigma / d \Omega, P, \Sigma, T, C_{x^{\prime}, z^{\prime}}, O_{x, z}=5,652 \\
& \gamma p \rightarrow K^{0} \Sigma^{+}: d \sigma / d \Omega, P=448
\end{aligned}
$$

- polarizations scaled by new $\Lambda$ decay constant $\alpha$ - (Ireland PRL 123 (2019), 182301), if applicable
- $\chi^{2}$ minimization with MINUIT on JURECA [Jülich

Supercomputing Centre, JURECA: JLSRF 2, A62 (2016)]

Resonance analysis:

- all 4 -star $N$ and $\Delta$ states up to $J=9 / 2$ are seen (exception: $N(1895) 1 / 2^{-}$) + some states rated less than 4 stars
- no additional s-channel diagram, but indications for new dyn. gen. poles


## Resonance contributions to $K \Sigma$ photoproduction

$$
\gamma p \rightarrow K^{+} \Sigma^{0}
$$


(Data not included in fit)


JüBo2O22 arXiv:2208.00089 [nucl-th]

- dominant partial waves: $I=3 / 2$

Exception: $P_{13}$ partial wave ( $/=1 / 2$ ):

| $N(1720) 3 / 2^{+}$ | $\operatorname{Re} E_{0}$ | $-2 \operatorname{lm} E_{0}$ | $\frac{\Gamma_{\pi N}^{1 / 2} \Gamma_{K \Sigma}^{1 / 2}}{\Gamma_{\text {tot }}}$ | $\theta_{\pi N \rightarrow K \Sigma}$ |
| :--- | :--- | :--- | :--- | :--- |
| $* * * *$ | $[\mathrm{MeV}]$ | $[\mathrm{MeV}]$ | $[\%]$ | $[\mathrm{deg}]$ |
| 2022 | $1726(8)$ | $185(12)$ | $5.9(1)$ | $82(6)$ |
| 2017 | $1689(4)$ | $191(3)$ | $0.6(0.4)$ | $26(58)$ |
| PDG 2021 | $1675 \pm 15$ | $250_{-100}^{+150}$ | - | - |


| $N(1900) 3 / 2^{+}$ <br> $* * * *$ | $\operatorname{Re} E_{0}$ | $-2 \operatorname{lm} E_{0}$ | $\frac{\Gamma_{\pi N}^{1 / 2} \Gamma_{K \Sigma}^{1 / 2}}{\Gamma_{\text {tot }}}$ | $\theta_{\pi N \rightarrow K \Sigma}$ |
| :--- | :--- | :--- | :--- | :--- |
| $[\mathrm{MeV}]$ | $[\mathrm{MeV}]$ | $[\%]$ | $[\mathrm{deg}]$ |  |
| 2022 | $1905(3)$ | $93(4)$ | $1.3(0.3)$ | $-40(18)$ |
| 2017 | $1923(2)$ | $217(23)$ | $10(7)$ | $-34(74)$ |
| PDG 2021 | $1920 \pm 20$ | $150 \pm 50$ | $4 \pm 2$ | $110 \pm 30$ |

- drop in cross section due to $N(1900) 3 / 2^{+}$
- "cusp-like structure" only qualitatively explained


## Resonance contributions to $K \Sigma$ photoproduction

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$$
\gamma p \rightarrow K^{+} \Sigma^{0}
$$



Data: Jude et al. (BGOOD) PLB 820 (2021)

## Selected results $\gamma p \rightarrow K^{0} \Sigma^{+}$



W [MeV]
Selected fit results:


- much less data than for $K^{+} \Sigma^{0}$ (448 vs 5,652 data points)
- in parts inconsistent data
$\rightarrow$ difficult to achieve a good fit result
- cusp in $\sigma_{\text {tot }}$ at $\sim 2 \mathrm{GeV}$ not reproduced (data not included in fit)

Data: open squares: SPAHIR 1999, cyan: SAPHIR 2005, orange:
CBELSA/TAPS 2007, black squares: CBELSA/TAPS 2011, open circles: A2 2018, open triangles: A2 2013, black triangles: Hall B 2003, black circles: CLAS 2013


## New data for $\gamma p \rightarrow \eta p$ from CBELSA/TAPS

included in JüBo2O22

- $T, P, H, G, E$ Müller PLB 803, 135323 (2020): very first data on $H, G$ (and $P$ ) in this channel

- $\sum_{\text {Afzal PRL } 125,152002 \text { (2020): }}$ Backward peak in data
$\rightarrow$ Observation of $\eta^{\prime} N$ cusp + importance of $N(1895) 1 / 2^{-}$(BnGa)


| $N(1535) 1 / 2^{-}$ | $\operatorname{Re} E_{0}$ | $-2 \operatorname{lm} E_{0}$ | $\frac{\Gamma_{\pi N}^{1 / 2} \Gamma_{\eta N}^{1 / 2}}{\Gamma_{\text {tot }}}$ | $\theta_{\pi N \rightarrow K \Sigma}$ |
| :--- | :--- | :--- | :--- | :--- |
| $* * * *$ | $[\mathrm{MeV}]$ | $[\mathrm{MeV}]$ | $[\%]$ | $[\mathrm{deg}]$ |
| 2022 | $1504(0)$ | $74(1)$ | $50(3)$ | $118(3)$ |
| 2017 | $1495(2)$ | $112(1)$ | $51(1)$ | $105(3)$ |
| PDG 2022 | $1510 \pm 10$ | $130 \pm 20$ | $43 \pm 3$ | $-76 \pm 5$ |


| $N(1650) 1 / 2^{-}$ | $\operatorname{Re} E_{0}$ | $-2 \operatorname{lm} E_{0}$ | $\frac{\Gamma_{\pi N}^{1 / 2} \Gamma_{\eta N}^{1 / 2}}{\Gamma_{\text {tot }}}$ | $\theta_{\pi N \rightarrow K \Sigma}$ |
| :--- | :--- | :--- | :--- | :--- |
| $\quad * * *$ | $[\mathrm{MeV}]$ | $[\mathrm{MeV}]$ | $[\%]$ | $[\mathrm{deg}]$ |
| 2022 | $1678(3)$ | $127(3)$ | $34(12)$ | $71(45)$ |
| 2017 | $1674(3)$ | $130(9)$ | $18(3)$ | $28(5)$ |
| PDG 2022 | $1655 \pm 15$ | $135 \pm 35$ | $29 \pm 3$ | $134 \pm 10$ |

$\rightarrow \eta N$ residue $N(1650) 1 / 2^{-}$much larger (similarly observed by BnGa)

JüBo2022:

- no $\eta^{\prime} N$ channel (or cusp), to be included in the future
- no $N(1895) 1 / 2^{-}$(not needed)
- backward peak from $N(1720) \& N(1900) 3 / 2^{+}$


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$\rightarrow$ Observation of $\eta^{\prime} N$ cusp + importance of $N(1895) 1 / 2^{-}$(BnGa)


| $\begin{gathered} N(1535) 1 / 2^{-} \\ * * * * \end{gathered}$ | $\operatorname{Re} E_{0}$ <br> [MeV] | $-2 \operatorname{Im} E_{0}$ <br> [MeV] | $\frac{\Gamma_{\pi N}^{1 / 2} \Gamma_{\eta N}^{1 / 2}}{\Gamma_{\text {tot }}}$ <br> [\%] | $\begin{aligned} & \theta_{\pi N \rightarrow K \Sigma} \\ & {[\mathrm{deg}]} \\ & \hline \end{aligned}$ |
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- no $N(1895) 1 / 2^{-}$(not needed)
- backward peak from $N(1720) \& N(1900) 3 / 2^{+}$ (turquoise lines: both states off) J JULICH


## Inclusion of the $\omega N$ channnel: $\pi N \rightarrow \omega N$ channel

## Wang et al. 2208.03061 [nucl-th]

- Preparation of the study of $\gamma N \rightarrow \omega N$ (abundant high quality data)
- importance of $\omega$ in nuclear matter [H. Shen et al. 1998 NPA]
- Scattering length $a_{\omega N} \rightarrow$ whether or not there are in-medium bound states

Selected fit results: Total cross section, backward/forward differential cross section




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## Scattering length:

- $\operatorname{Re} \bar{a}>0 \rightarrow$ in-medium bound states
- Result of fit $A(B)$
$\bar{a}=-0.24(-0.21)+0.05 i(0.05 i)$

§) JüLICH


## Summary

Jülich-Bonn dynamical coupled-channel analysis:

- Extraction of the $N^{*}$ and $\Delta^{*}$ spectrum in a simultaneous analysis of pion- and photon-induced reactions:
$-\pi N \rightarrow \pi N, \eta N, K \Lambda$ and $K \Sigma$
lagrangian based description, unitarity \& analyticity respected
- $\gamma N \rightarrow \pi N, \eta N, K \Lambda$ and $K \Sigma$ in a semi-phenomenological approach hadronic final state interaction: JüBo DCC analysis
$\rightarrow$ analysis of almost 72,000 data points
- $\pi N \rightarrow \omega N$ channel included, prerequisite for $\omega$ photoproduction
- Electroproduction: Jülich-Bonn-Washington approach Mai et al. PRC 103 (2021), PRC 106 (2022)
- JüBo photoproduction amplitude as input at $Q^{2}=0$
- New interactive web interface: https://jbw.phys.gwu.edu (multipoles, observables, data)
$\rightarrow$ Talk by Maxim Mai on Wednesday
Thank you for your attention!


[^0]:    Data: Kraemer et al. 1964 PR, Danburg et al. 1970 PRD, Binnie et al. 1973 PRD, Keyne et al. 1976 PRD, Karami et al. 1979 NPB

